

Solving the Van Der Pol nonlinear differential equation using first order approximation perturbation method

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Abstract

The Van Der Pol differential equation

$$x''(t) - \alpha(1 - x^2(t))x'(t) + x(t) = 0$$

Was solved using perturbation with first order approximation. Two different solutions were obtained. The first solution restricted the initial conditions to be $x(0)^2 + x'(0)^2 = 4$ which resulted in forcing function that caused resonance to be eliminated. This gave a stable solution but with initial conditions restricted to be near the origin of the phase plane space.

The second solution allowed arbitrary initial conditions any where in the phase plane but the the resulting forcing function caused resonance resulting in a solution which became unstable after some time.

Phase plane plots are used to compare the two solutions.

1 First solution: Restriction on initial condition. No resonance

The Van Der Pol equation is

$$\ddot{x} - \alpha(1 - x^2)\dot{x} + x = 0$$

Initial conditions are $x(0) = \varphi$ and $\dot{x}(0) = \xi$.

If $\alpha \lll 1$, then the equation becomes $\ddot{x}_0 + x_0 = 0$ which has an exact solution. We perturb the exact solution for $\ddot{x}_0 + x_0 = 0$ to approximate the solution of the nonlinear equation $\ddot{x} - \alpha(1 - x^2)\dot{x} + x = 0$ in terms of perturbation parameter α Let the solution of the Van Der Pol equation be $x(t)$ and the solution to the exact equation $\ddot{x}_0 + x_0 = 0$ be $x_0(t)$, then

$$x(t) = x_0(t) + \alpha x_1(t) + \alpha^2 x_2(t) + \dots$$

First order approximation results in

$$x(t) = x_0(t) + \alpha x_1(t)$$

Now we need to determine $x_0(t)$ and $x_1(t)$. Substituting (2) into (1) gives

$$\ddot{x}_0 + \alpha \ddot{x}_1 - \alpha \dot{x}_0 + \alpha x_0^2 \dot{x}_0 + \alpha^3 x_1^2 \dot{x}_0 + 2\alpha^2 x_1 x_0 \dot{x}_0 - \alpha^2 \dot{x}_1 + \alpha^2 x_0^2 \dot{x}_1 + \alpha^4 x_1^2 \dot{x}_1 + 2\alpha^3 x_1 x_0 \dot{x}_1 + x_0 + \alpha x_1 = 0$$

Collecting terms with same power of α gives

$$\alpha^0(\ddot{x}_0 + x_0) + \alpha(\ddot{x}_1 - \dot{x}_0 + x_0^2 \dot{x}_0 + x_1) + \alpha^2(x_0^2 \dot{x}_1 + 2x_1 x_0 \dot{x}_0 - \dot{x}_1) + \alpha^3(x_1^2 \dot{x}_0 + 2x_1 x_0 \dot{x}_1) + \alpha^4 x_1^2 \dot{x}_1 = 0$$

Setting terms which multiply by higher power α to zero and since it is assumed that α is very small therefore

$$(\ddot{x}_0 + x_0) + \alpha(\ddot{x}_1 - \dot{x}_0 + x_0^2 \dot{x}_0 + x_1) = 0$$

For the LHS to be zero implies that

$$\ddot{x}_0 + x_0 = 0$$

And

$$\ddot{x}_1 - \dot{x}_0 + x_0^2 \dot{x}_0 + x_1 = 0$$

Equation (3) is solved for x_0 and (4) is solved for x_1 in order to find the solution $x(t)$. Solution to (3) is

$$x_0(t) = A_0 \cos t + B_0 \sin t$$

Assuming initial conditions for $x_1(t)$ are zero, then initial conditions for $x_0(t)$ can be taken to be those given for $x(t)$. Solving for A_0, B_0 gives

$$x_0(t) = \varphi \cos t + \xi \sin t$$

Substituting the solution just found (and its derivative) into (4) gives

$$\begin{aligned} \ddot{x}_1 + x_1 = & -\varphi \sin t + \xi \cos t - \xi^3 \cos t \sin^2 t - 2\xi^2 \varphi \cos^2 t \sin t + \xi^2 \varphi \sin^3 t \\ & - \xi \varphi^2 \cos^3 t + 2\xi \varphi^2 \cos t \sin^2 t + \varphi^3 \cos^2 t \sin t \end{aligned}$$

Using $\sin t \cos^2 t = \frac{1}{4}(\sin t + \sin 3t)$ and $\cos t \sin^2 t = \frac{1}{4}(\cos t - \cos 3t)$ the above can be simplified to

$$\begin{aligned} \ddot{x}_1 + x_1 = & \left(-\varphi - \frac{\xi^2 \varphi}{2} + \frac{\varphi^3}{4}\right) \sin t + \left(\xi - \frac{\xi^3}{4} + \frac{\xi \varphi^2}{2}\right) \cos t + \left(\frac{\xi^3}{4} - \frac{\xi \varphi^2}{2}\right) \cos 3t \\ & + \left(-\frac{\xi^2 \varphi}{2} + \frac{\varphi^3}{4}\right) \sin 3t + \xi \varphi (\xi \sin^3 t - \varphi \cos^3 t) \end{aligned}$$

Using $\xi \sin^3 t - \varphi \cos^3 t = \frac{1}{4}(-3\varphi \cos t - \varphi \cos 3t + 3\xi \sin t - \xi \sin 3t)$ the above can be simplified further to

$$\begin{aligned} \ddot{x}_1 + x_1 = & \left(-\varphi - \frac{\xi^2 \varphi}{2} + \frac{\varphi^3}{4} + \frac{3\xi^2 \varphi}{4}\right) \sin t + \left(\xi - \frac{\xi^3}{4} + \frac{\xi \varphi^2}{2} - \frac{3}{4}\xi \varphi^2\right) \cos t \\ & + \left(\frac{\xi^3}{4} - \frac{\xi \varphi^2}{2} - \frac{\xi \varphi^2}{4}\right) \cos 3t + \left(-\frac{\xi^2 \varphi}{2} + \frac{\varphi^3}{4} - \frac{\xi^2 \varphi}{4}\right) \sin 3t \end{aligned} \quad (5)$$

The above is the equation we will now solve for $x_1(t)$, which has four forcing functions, hence four particular solutions. two of these particular solutions will cause a complete solution blows up as time increases (the first two in the RHS shown above). If we however restrict the initial conditions such that

$$\left(-\varphi - \frac{\xi^2 \varphi}{2} + \frac{\varphi^3}{4} + \frac{3\xi^2 \varphi}{4}\right) = 0$$

And

$$\left(\xi - \frac{\xi^3}{4} + \frac{\xi \varphi^2}{2} - \frac{3}{4}\xi \varphi^2\right) = 0$$

Then those forcing function will vanish. The above two equation result in the same restriction, which can be found to be

$$4 = \xi^2 + \varphi^2$$

By restricting $\xi^2 + \varphi^2$ to be exactly 4, then the solution we obtain for $x_1(t)$ from (5) will not blow up. Hence (4) can now be rewritten without the two forcing functions which caused resonance resulting in

$$\ddot{x}_1 + x_1 = \xi \left(\frac{\xi^2 - 3\varphi^2}{4} \right) \cos 3t + \varphi \left(\frac{\varphi^2 - 3\xi^2}{4} \right) \sin 3t$$

But $\varphi^2 = 4 - \xi^2$ and $\xi^2 = 4 - \varphi^2$, hence the above becomes after some simplification

$$\ddot{x}_1 + x_1 = \xi(\xi^2 - 3) \cos 3t + \varphi(\varphi^2 - 3) \sin 3t \quad (6)$$

The homogeneous solution to the above is

$$x_{1,h} = c_1 \cos t + c_2 \sin t$$

And we guess two particular solutions $x_{1,p1} = a_1 \cos 3t + a_2 \sin 3t$ and $x_{1,p2} = d_1 \cos 3t + d_2 \sin 3t$. By substituting each of these particular solutions into (6) one at a time, and comparing coefficients, we can determine a_1, a_2, d_1, d_2 . This results in

$$x_{1,p1} = \frac{\xi(3 - \xi^2)}{8} \cos 3t$$

Similarly,

$$x_{1,p2} = \frac{\varphi(3 - \varphi^2)}{8} \sin 3t$$

Hence the solution to (6) becomes

$$\begin{aligned} x_1(t) &= x_{1,h} + x_{1,p1} + x_{1,p2} \\ &= (c_1 \cos t + c_2 \sin t) + \frac{\xi(3 - \xi^2)}{8} \cos 3t + \frac{\varphi(3 - \varphi^2)}{8} \sin 3t \end{aligned} \quad (7)$$

Now applying initial conditions, which are $x_1(0) = 0$ and $\dot{x}_1(0) = 0$, to (7) and determining c_1 and c_2 results in

$$x_1(t) = \frac{\xi(\xi^2 - 3)}{8} \cos t + \frac{3}{8}\varphi(\varphi^2 - 3) \sin t + \frac{\xi(3 - \xi^2)}{8} \cos 3t + \frac{\varphi(3 - \varphi^2)}{8} \sin 3t$$

Therefore we now have the final perturbation solution, which is

$$\begin{aligned} x(t) &= x_0(t) + \alpha x_1(t) \\ &= \varphi \cos t + \xi \sin t + \alpha \left(\frac{\xi(\xi^2 - 3)}{8} \cos t + \frac{3}{8}\varphi(\varphi^2 - 3) \sin t + \frac{\xi(3 - \xi^2)}{8} \cos 3t + \frac{\varphi(3 - \varphi^2)}{8} \sin 3t \right) \end{aligned}$$

Where $x(0) = \varphi$ and $\dot{x}(0) = \xi$ and the above solution is valid under the restriction that

$$x^2(0) + \dot{x}^2(0) = 4$$

We notice that the above restriction implies that both $x(0)$ and $\dot{x}(0)$ have to be less than 4 to avoid getting a quantity which is complex. This implies that this solution is valid near the center of the phase plane only. In addition, it is valid only for small α . To illustrate this solution, We show the phase plot which results from the above solution, and also plot the solution $x(t)$. Selecting $x(0) = 1.5$ and $\dot{x}(0) = \sqrt{4 - 1.5^2} = 1.3229$

```

In[271]:= initialPosition = 1.5;
initialSpeed = Sqrt[4 - initialPosition^2];
x = x0 Cos[t] + v0 Sin[t] +
  α ( (v0 (v0^2 - 3) / 8) Cos[t] + (3 / 8) x0 (x0^2 - 3) Sin[t] + (v0 (3 - v0^2) / 8) Cos[3 t] +
    (x0 (3 - x0^2) / 8) Sin[3 t] );

xα = x /. {α -> 0.01, x0 -> initialPosition, v0 -> initialSpeed};
vα = D[x, t] /. {α -> 0.01, x0 -> initialPosition, v0 -> initialSpeed};
data = Table[{xα, vα}, {t, 0, 60, .01}];
p = ListPlot[data, AxesLabel -> {"x(t)", "v(t)"},
  PlotLabel -> "Phase plane for I.C. x(0)=" <> ToString[initialPosition] <>
    " and v(0)=" <> ToString[initialSpeed] <> " α=0.01"];
p2 = Graphics[{PointSize[.03], Point[{initialPosition, initialSpeed}]}];
Show[{p, p2}, AspectRatio -> 1]

```

Phase plane for I.C. x(0)=1.5 and v(0)=1.32288 α=0.01

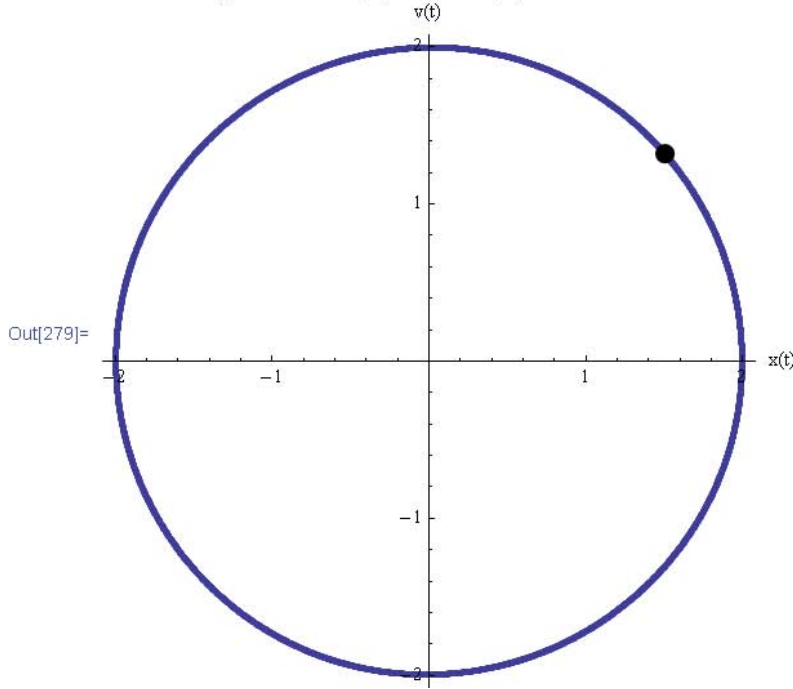


Figure 1: No resonance

2 Second solution. No restriction on initial conditions. Resonance present in the solution

The same initial steps are repeated as in the first solution, up the equation to solve for $x_1(t)$

$$\begin{aligned}
 \ddot{x}_1 + x_1 = & \varphi \left(\frac{\varphi^2 + \xi^2}{4} - 1 \right) \sin t + \xi \left(1 - \frac{\xi^2 - \varphi^2}{4} \right) \cos t \\
 & + \xi \left(\frac{\xi^2 - 3\varphi^2}{4} \right) \cos 3t + \varphi \left(\frac{\varphi^2 - 3\xi^2}{4} \right) \sin 3t
 \end{aligned} \tag{8}$$

There exists four particular solutions relating to the above four forcing functions

$$\begin{aligned}
 f_1(t) &= \varphi \left(\frac{\varphi^2 + \xi^2}{4} - 1 \right) \sin t \\
 f_2(t) &= \xi \left(1 - \frac{\xi^2 - \varphi^2}{4} \right) \cos t \\
 f_3(t) &= \xi \left(\frac{\xi^2 - 3\varphi^2}{4} \right) \cos 3t \\
 f_4(t) &= \varphi \left(\frac{\varphi^2 - 3\xi^2}{4} \right) \sin 3t
 \end{aligned}$$

For each of the above, we guess a particular solution, and determine the solution parameters. We guess $x_{1,p1} = t(c_1 \sin t + c_2 \cos t)$, $x_{1,p2} = t(c_1 \sin t + c_2 \cos t)$, $x_{1,p3} = c_1 \sin 3t + c_2 \cos 3t$,

and $x_{1,p4} = c_1 \sin 3t + c_2 \cos 3t$. By substituting each of the above into (8) one at a time and comparing coefficients, we arrive at the following solutions

$$\begin{aligned} x_{1,p1} &= t \left(\frac{1}{2} \varphi \left(1 - \frac{\varphi^2 + \xi^2}{4} \right) \cos t \right) \\ x_{1,p2} &= t \left(\frac{1}{2} \xi \left(1 - \frac{\xi^2 - \varphi^2}{4} \right) \sin t \right) \\ x_{1,p3} &= \frac{1}{8} \xi \left(\frac{\xi^2 - \varphi^2}{4} - 1 \right) \cos 3t \\ x_{1,p4} &= \frac{1}{8} \varphi \left(\frac{3\xi^2 - \varphi^2}{4} \right) \sin 3t \end{aligned}$$

$$\begin{aligned} x_1(t) &= x_{1,h} + x_{1,p4} + x_{2,p4} + x_{3,p4} + x_{4,p4} \\ &= (c_1 \cos t + c_2 \sin t) + t \left(\frac{1}{2} \varphi \left(1 - \frac{\varphi^2 + \xi^2}{4} \right) \cos t \right) + t \left(\frac{1}{2} \xi \left(1 - \frac{\xi^2 - \varphi^2}{4} \right) \sin t \right) \\ &\quad + \frac{1}{8} \xi \left(\frac{\xi^2 - \varphi^2}{4} - 1 \right) \cos 3t + \frac{1}{8} \varphi \left(\frac{3\xi^2 - \varphi^2}{4} \right) \sin 3t \end{aligned}$$

Now applying initial conditions which are $x_1(0) = 0$ and $\dot{x}_1(0) = 0$ and determining c_1 and c_2 results in

$$\begin{aligned} x_1(t) &= \left(-\frac{1}{32} \xi^3 + \frac{1}{32} \xi \varphi^2 + \frac{1}{8} \xi \right) \cos t - \frac{1}{32} \varphi (5\xi^2 - 7\varphi^2 + 16) \sin t \\ &\quad + t \left(\frac{1}{2} \varphi \left(1 - \frac{\varphi^2 + \xi^2}{4} \right) \cos t \right) + t \left(\frac{1}{2} \xi \left(1 - \frac{\xi^2 - \varphi^2}{4} \right) \sin t \right) \\ &\quad + \frac{1}{8} \xi \left(\frac{\xi^2 - \varphi^2}{4} - 1 \right) \cos 3t + \frac{1}{8} \varphi \left(\frac{3\xi^2 - \varphi^2}{4} \right) \sin 3t \end{aligned}$$

We can now write the final solution for $x(t)$ as

$$\begin{aligned} x(t) &= x_0(t) + \alpha x_1(t) \\ &= \varphi \cos t + \xi \sin t + \\ &\quad \alpha \left(-\frac{1}{32} \xi^3 + \frac{1}{32} \xi \varphi^2 + \frac{1}{8} \xi \right) \cos t - \alpha \frac{1}{32} \varphi (5\xi^2 - 7\varphi^2 + 16) \sin t \\ &\quad + \alpha t \left(\frac{1}{2} \varphi \left(1 - \frac{\varphi^2 + \xi^2}{4} \right) \cos t \right) + \alpha t \left(\frac{1}{2} \xi \left(1 - \frac{\xi^2 - \varphi^2}{4} \right) \sin t \right) \\ &\quad + \alpha \frac{1}{8} \xi \left(\frac{\xi^2 - \varphi^2}{4} - 1 \right) \cos 3t + \alpha \frac{1}{8} \varphi \left(\frac{3\xi^2 - \varphi^2}{4} \right) \sin 3t \end{aligned}$$

To illustrate this solution, we show the phase plot which results from the above solution, and also plot the solution $x(t)$. We select $x(0) = 1.5$ and $\dot{x}(0) = 1.3229$. We note that the same initial conditions are used as in the earlier solution to compare both solutions.

```

In[262]:= initialPosition = 1.5;
initialSpeed = Sqrt[4 - initialPosition^2];

x = x0 Cos[t] + v0 Sin[t] +  $\alpha \left( -\frac{v0^3}{32} + \frac{v0 x0^2}{32} + \frac{1}{8} v0 \right) \text{Cos}[t] -$ 
 $\frac{\alpha}{32} x0 (5 v0^2 - 7 x0^2 + 16) \text{Sin}[t] + \alpha t \left( \frac{1}{2} x0 \left( 1 - \frac{x0^2 + v0^2}{4} \right) \text{Cos}[t] \right) +$ 
 $\alpha t \left( \frac{1}{2} v0 \left( 1 - \frac{v0^2 - x0^2}{4} \right) \text{Sin}[t] \right) + \frac{\alpha}{8} v0 \left( \frac{v0^2 - x0^2}{4} - 1 \right) \text{Cos}[3 t] +$ 
 $\frac{\alpha}{8} x0 \left( \frac{3 v0^2 - x0^2}{4} \right) \text{Sin}[3 t];$ 

x $\alpha$  = x /. { $\alpha$  -> 0.01, x0 -> initialPosition, v0 -> initialSpeed};
v $\alpha$  = D[x, t] /. { $\alpha$  -> 0.01, x0 -> initialPosition, v0 -> initialSpeed};
data = Table[{x $\alpha$ , v $\alpha$ }, {t, 0, 60, .01}];
p = ListPlot[data, AxesLabel -> {"x(t)", "v(t)"},
  PlotLabel -> "Phase plane for I.C. x(0)=" <> ToString[initialPosition] <>
    " and v(0)=" <> ToString[initialSpeed] <> "  $\alpha=0.01$ " ];
p2 = Graphics[{PointSize[.03], Point[{initialPosition, initialSpeed}]}];
Show[{p, p2}, AspectRatio -> 1]

```

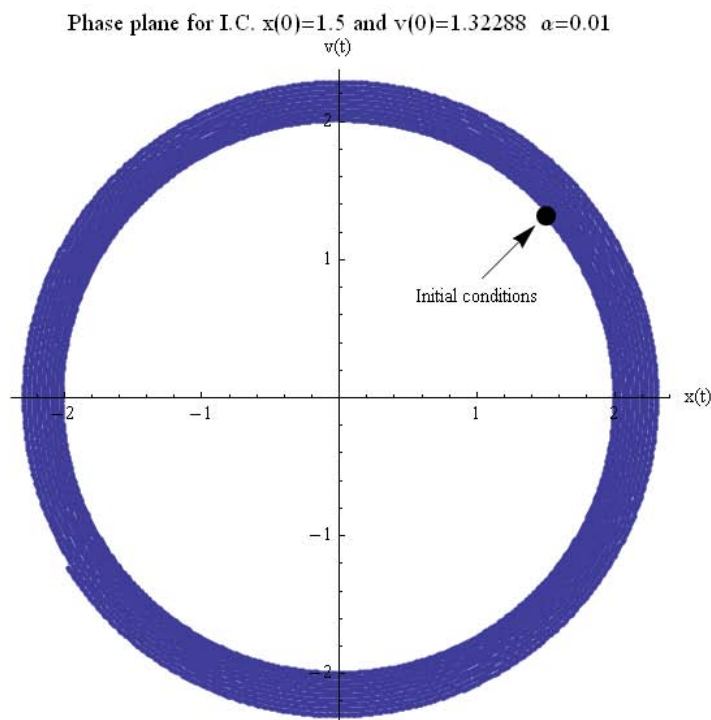


Figure 2: Rresonance

We see clearly the effect of including resonance particular solutions. The limit cycle boundary is increasing with time.

The effect becomes more clear if we plot the solution $x(t)$ itself, and compare it to the solution earlier with no resonance. Below, We show $x(t)$ from both solutions. We clearly see the effect of including the resonance terms. This is the solution with no resonance terms

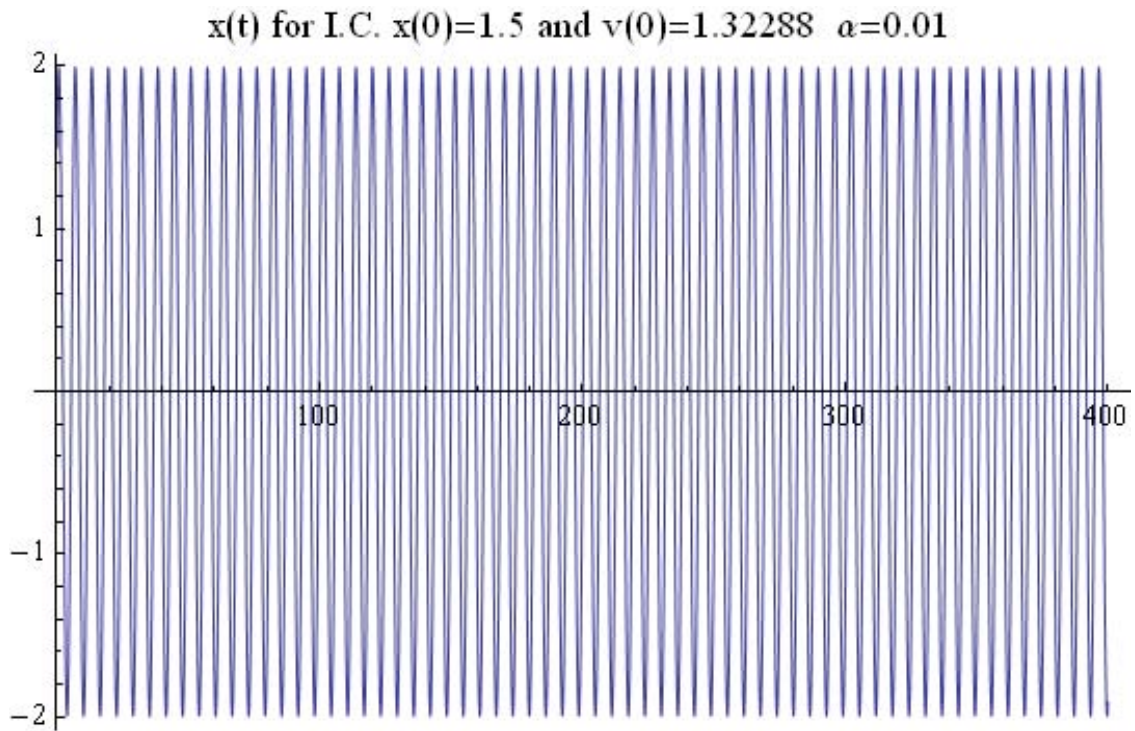


Figure 3: No rresonance terms

This is the solution with resonance terms

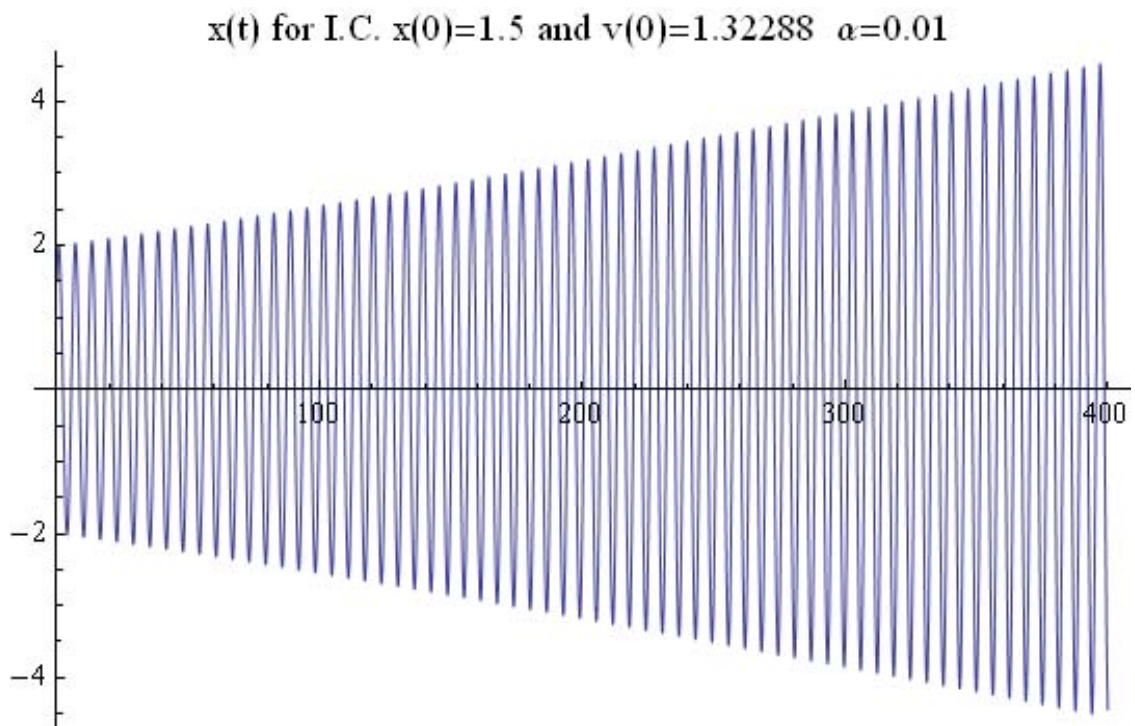


Figure 4: Rresonance terms

In the second solution (with resonance), there is no restriction on initial conditions. Hence we can for example look at the solution starting from any $x(0)$ and $\dot{x}(0)$. This is the phase plot for $x(0) = 5$ and $\dot{x}(0) = 2$. This would not be possible using the first solution, due to the restriction on $x(0)^2 + \dot{x}(0)^2$ having to be equal to 4

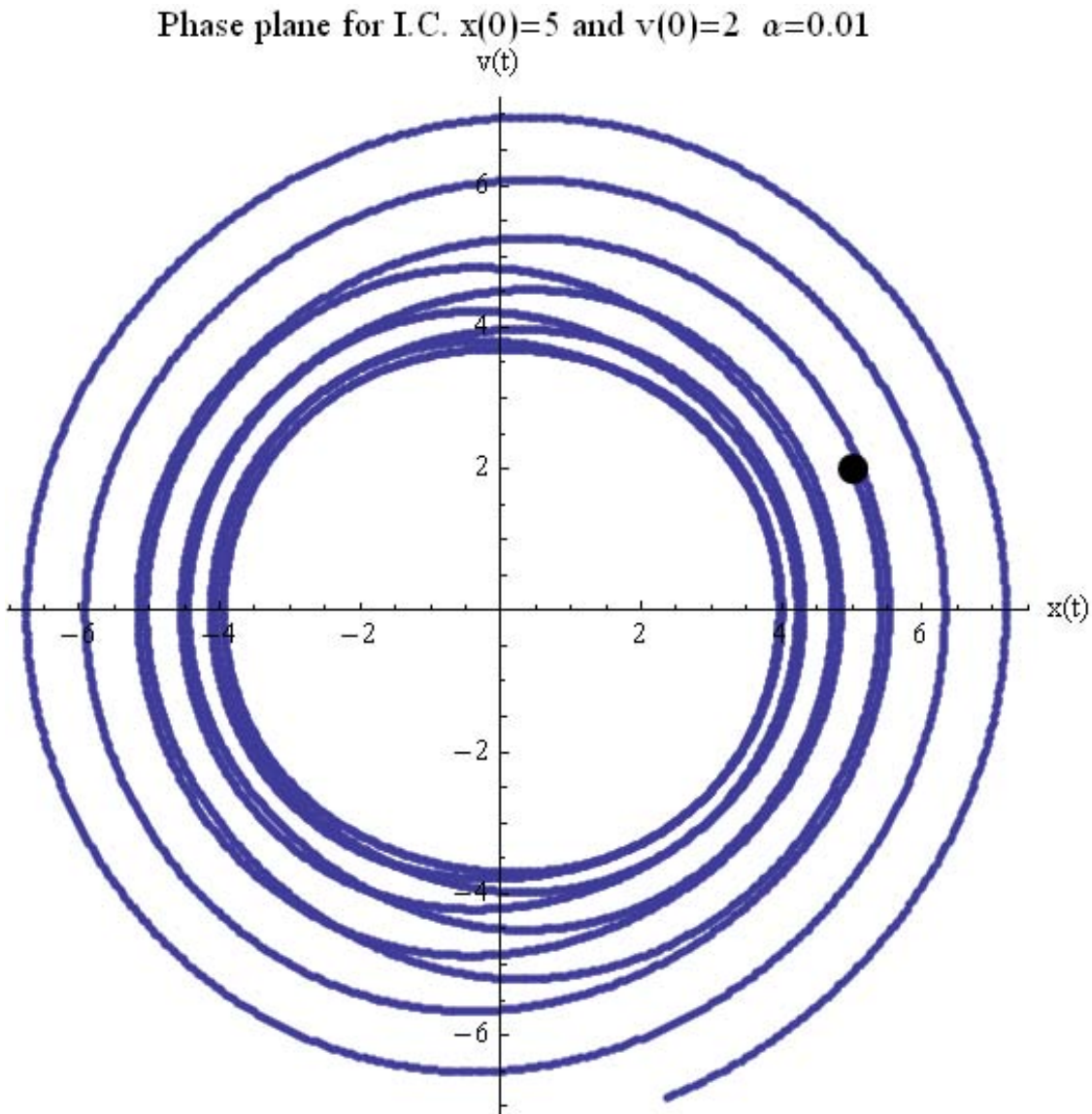


Figure 5: phase plot

3 Conclusion

Van Der Pol was solved using perturbation for first order approximation. It was solved using two methods. In the first method, the solution was restricted to not include forcing functions which leads to resonance.

In the second method, no such restriction was made. In both methods, this problem was solved for general initial conditions (i.e. both $x(0)$ and $\dot{x}(0)$ can both be nonzero).

In the first method, it was found that we must restrict the initial conditions to be such that $x(0)^2 + \dot{x}(0)^2 = 4$. This is the condition which resulted in the resonance terms vanishing from the solution. This leads to a solution which generated a limit cycle which did not blow up as the solution is run for longer time. In other words, once the solution enters a limit cycle, it remains in the limit cycle.

In the second method, no restriction on the initial conditions was made. This allowed the solution to start from any state. However, the limit cycle would grow with time, and the solution will suffer from fluttering due to the presence of the resonance terms. However, even though the solution was not stable in the long term, this second approach allowed one to examine the solution for a shorter time but with the flexibility of choosing any initial conditions.

It was also observed that increasing the value of the perturbation parameter α gradually resulted in an inaccurate solution as would be expected, as the solutions derived here all assumed a very small¹ α was used in generating the solutions and plots.

¹ $\alpha = 0.01$ value was used

For future research, it would be interesting to consider the effect of higher order approximation on the solutions and compare with accurate numerical solutions.