A Solution Manual For

## Ordinary Differential Equations, Robert H. Martin, 1983



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## 1 Problem 1.1-2, page 6

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## 1.1 problem 1.1-2 (a)

1.1.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 3
1.1.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 4

Internal problem ID [2447]
Internal file name [OUTPUT/1939_Sunday_June_05_2022_02_40_08_AM_92199210/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-2, page 6
Problem number: 1.1-2 (a).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}=t^{2}+3
$$

### 1.1.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
y & =\int t^{2}+3 \mathrm{~d} t \\
& =\frac{1}{3} t^{3}+3 t+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{1}{3} t^{3}+3 t+c_{1} \tag{1}
\end{equation*}
$$



Figure 1: Slope field plot

## Verification of solutions

$$
y=\frac{1}{3} t^{3}+3 t+c_{1}
$$

Verified OK.

### 1.1.2 Maple step by step solution

Let's solve

$$
y^{\prime}=t^{2}+3
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Integrate both sides with respect to $t$
$\int y^{\prime} d t=\int\left(t^{2}+3\right) d t+c_{1}$
- Evaluate integral
$y=\frac{1}{3} t^{3}+3 t+c_{1}$
- $\quad$ Solve for $y$

$$
y=\frac{1}{3} t^{3}+3 t+c_{1}
$$

Maple trace

- Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
<- quadrature successful`
$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 14

```
dsolve(diff(y(t),t)=t^2+3,y(t), singsol=all)
```

$$
y(t)=\frac{1}{3} t^{3}+3 t+c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.003 (sec). Leaf size: 18

```
DSolve[y'[t]==t^2+3,y[t],t,IncludeSingularSolutions -> True]
```

$$
y(t) \rightarrow \frac{t^{3}}{3}+3 t+c_{1}
$$

## 1.2 problem 1.1-2 (b)

1.2.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 6
1.2.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 7

Internal problem ID [2448]
Internal file name [OUTPUT/1940_Sunday_June_05_2022_02_40_11_AM_83674871/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-2, page 6
Problem number: 1.1-2 (b).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}=\mathrm{e}^{2 t} t
$$

### 1.2.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
y & =\int \mathrm{e}^{2 t} t \mathrm{~d} t \\
& =\frac{(2 t-1) \mathrm{e}^{2 t}}{4}+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{(2 t-1) \mathrm{e}^{2 t}}{4}+c_{1} \tag{1}
\end{equation*}
$$



Figure 2: Slope field plot

Verification of solutions

$$
y=\frac{(2 t-1) \mathrm{e}^{2 t}}{4}+c_{1}
$$

Verified OK.

### 1.2.2 Maple step by step solution

Let's solve

$$
y^{\prime}=\mathrm{e}^{2 t} t
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Integrate both sides with respect to $t$

$$
\int y^{\prime} d t=\int \mathrm{e}^{2 t} t d t+c_{1}
$$

- Evaluate integral
$y=\frac{(2 t-1) \mathrm{e}^{2 t}}{4}+c_{1}$
- $\quad$ Solve for $y$

$$
y=\frac{\mathrm{e}^{2 t} t}{2}-\frac{\mathrm{e}^{2 t}}{4}+c_{1}
$$

Maple trace

- Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
<- quadrature successful`
$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 17

```
dsolve(diff(y(t),t)=t*exp(2*t),y(t), singsol=all)
```

$$
y(t)=\frac{(2 t-1) \mathrm{e}^{2 t}}{4}+c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.004 (sec). Leaf size: 22
DSolve[y' t$]==\mathrm{t} * \operatorname{Exp}[2 * \mathrm{t}], \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow \frac{1}{4} e^{2 t}(2 t-1)+c_{1}
$$

## 1.3 problem 1.1-2 (c)

1.3.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 9
1.3.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 10

Internal problem ID [2449]
Internal file name [OUTPUT/1941_Sunday_June_05_2022_02_40_12_AM_45605752/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-2, page 6
Problem number: 1.1-2 (c).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}=\sin (3 t)
$$

### 1.3.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
y & =\int \sin (3 t) \mathrm{d} t \\
& =-\frac{\cos (3 t)}{3}+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\cos (3 t)}{3}+c_{1} \tag{1}
\end{equation*}
$$



Figure 3: Slope field plot
Verification of solutions

$$
y=-\frac{\cos (3 t)}{3}+c_{1}
$$

Verified OK.

### 1.3.2 Maple step by step solution

Let's solve
$y^{\prime}=\sin (3 t)$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Integrate both sides with respect to $t$
$\int y^{\prime} d t=\int \sin (3 t) d t+c_{1}$
- Evaluate integral
$y=-\frac{\cos (3 t)}{3}+c_{1}$
- $\quad$ Solve for $y$

$$
y=-\frac{\cos (3 t)}{3}+c_{1}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
<- quadrature successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 12

```
dsolve(diff(y(t),t)=\operatorname{sin}(3*t),y(t), singsol=all)
```

$$
y(t)=-\frac{\cos (3 t)}{3}+c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.005 (sec). Leaf size: 16
DSolve[y'[t]==Sin[3*t], y[t],t,IncludeSingularSolutions $->$ True]

$$
y(t) \rightarrow-\frac{1}{3} \cos (3 t)+c_{1}
$$

## 1.4 problem 1.1-2 (d)

1.4.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 12
1.4.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 13

Internal problem ID [2450]
Internal file name [OUTPUT/1942_Sunday_June_05_2022_02_40_14_AM_83224243/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-2, page 6
Problem number: 1.1-2 (d).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}=\sin (t)^{2}
$$

### 1.4.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
y & =\int \sin (t)^{2} \mathrm{~d} t \\
& =-\frac{\sin (t) \cos (t)}{2}+\frac{t}{2}+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\sin (t) \cos (t)}{2}+\frac{t}{2}+c_{1} \tag{1}
\end{equation*}
$$



Figure 4: Slope field plot

Verification of solutions

$$
y=-\frac{\sin (t) \cos (t)}{2}+\frac{t}{2}+c_{1}
$$

Verified OK.

### 1.4.2 Maple step by step solution

Let's solve

$$
y^{\prime}=\sin (t)^{2}
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Integrate both sides with respect to $t$
$\int y^{\prime} d t=\int \sin (t)^{2} d t+c_{1}$
- Evaluate integral

$$
y=-\frac{\sin (t) \cos (t)}{2}+\frac{t}{2}+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=-\frac{\sin (t) \cos (t)}{2}+\frac{t}{2}+c_{1}
$$

Maple trace

```
`Methods for first order ODEs:
    Trying classification methods ---
trying a quadrature
<- quadrature successful`
```

$\checkmark$ Solution by Maple
Time used: 0.015 (sec). Leaf size: 15

```
dsolve(diff(y(t),t)=sin(t)^2,y(t), singsol=all)
```

$$
y(t)=\frac{t}{2}+c_{1}-\frac{\sin (2 t)}{4}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.012 (sec). Leaf size: 21
DSolve[y'[t]==Sin[t]~2,y[t],t,IncludeSingularSolutions $->$ True]

$$
y(t) \rightarrow \frac{t}{2}-\frac{1}{4} \sin (2 t)+c_{1}
$$

## 1.5 problem 1.1-2 (e)

1.5.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 15
1.5.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 16

Internal problem ID [2451]
Internal file name [OUTPUT/1943_Sunday_June_05_2022_02_40_15_AM_39058871/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-2, page 6
Problem number: 1.1-2 (e).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}=\frac{t}{t^{2}+4}
$$

### 1.5.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
y & =\int \frac{t}{t^{2}+4} \mathrm{~d} t \\
& =\frac{\ln \left(t^{2}+4\right)}{2}+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\ln \left(t^{2}+4\right)}{2}+c_{1} \tag{1}
\end{equation*}
$$



Figure 5: Slope field plot

Verification of solutions

$$
y=\frac{\ln \left(t^{2}+4\right)}{2}+c_{1}
$$

Verified OK.

### 1.5.2 Maple step by step solution

Let's solve

$$
y^{\prime}=\frac{t}{t^{2}+4}
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Integrate both sides with respect to $t$
$\int y^{\prime} d t=\int \frac{t}{t^{2}+4} d t+c_{1}$
- Evaluate integral

$$
y=\frac{\ln \left(t^{2}+4\right)}{2}+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\frac{\ln \left(t^{2}+4\right)}{2}+c_{1}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
<- quadrature successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 14

```
dsolve(diff(y(t),t)=t/(t^2+4),y(t), singsol=all)
```

$$
y(t)=\frac{\ln \left(t^{2}+4\right)}{2}+c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.006 (sec). Leaf size: 18
DSolve[y'[t]==t/(t^2+4),y[t],t,IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow \frac{1}{2} \log \left(t^{2}+4\right)+c_{1}
$$

## 1.6 problem 1.1-2 (f)

1.6.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 18
1.6.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 19

Internal problem ID [2452]
Internal file name [OUTPUT/1944_Sunday_June_05_2022_02_40_17_AM_22616209/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-2, page 6
Problem number: 1.1-2 (f).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}=\ln (t)
$$

### 1.6.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
y & =\int \ln (t) \mathrm{d} t \\
& =t \ln (t)-t+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=t \ln (t)-t+c_{1} \tag{1}
\end{equation*}
$$



Figure 6: Slope field plot

Verification of solutions

$$
y=t \ln (t)-t+c_{1}
$$

Verified OK.

### 1.6.2 Maple step by step solution

Let's solve

$$
y^{\prime}=\ln (t)
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Integrate both sides with respect to $t$

$$
\int y^{\prime} d t=\int \ln (t) d t+c_{1}
$$

- Evaluate integral

$$
y=t \ln (t)-t+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=t \ln (t)-t+c_{1}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
<- quadrature successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 13

```
dsolve(diff(y(t),t)=ln(t),y(t), singsol=all)
```

$$
y(t)=t \ln (t)-t+c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.004 (sec). Leaf size: 15

```
DSolve[y'[t]==Log[t],y[t],t,IncludeSingularSolutions -> True]
```

$$
y(t) \rightarrow-t+t \log (t)+c_{1}
$$

## 1.7 problem 1.1-2 (g)

1.7.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 21
1.7.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 22

Internal problem ID [2453]
Internal file name [OUTPUT/1945_Sunday_June_05_2022_02_40_20_AM_75628296/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-2, page 6
Problem number: 1.1-2 (g).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}=\frac{t}{\sqrt{t}+1}
$$

1.7.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
y & =\int \frac{t}{\sqrt{t}+1} \mathrm{~d} t \\
& =\frac{2 t^{\frac{3}{2}}}{3}-t+2 \sqrt{t}-2 \ln (\sqrt{t}+1)+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{2 t^{\frac{3}{2}}}{3}-t+2 \sqrt{t}-2 \ln (\sqrt{t}+1)+c_{1} \tag{1}
\end{equation*}
$$



Figure 7: Slope field plot

Verification of solutions

$$
y=\frac{2 t^{\frac{3}{2}}}{3}-t+2 \sqrt{t}-2 \ln (\sqrt{t}+1)+c_{1}
$$

Verified OK.

### 1.7.2 Maple step by step solution

Let's solve

$$
y^{\prime}=\frac{t}{\sqrt{t}+1}
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Integrate both sides with respect to $t$
$\int y^{\prime} d t=\int \frac{t}{\sqrt{t}+1} d t+c_{1}$
- Evaluate integral

$$
y=\frac{2 t^{\frac{3}{2}}}{3}-t+2 \sqrt{t}-2 \ln (\sqrt{t}+1)+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\frac{2 t^{\frac{3}{2}}}{3}-t+2 \sqrt{t}-2 \ln (\sqrt{t}+1)+c_{1}
$$

Maple trace

- Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
<- quadrature successful`
$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 27
dsolve(diff( $y(t), t)=t /(\operatorname{sqrt}(t)+1), y(t), \quad$ singsol=all)

$$
y(t)=\frac{2 t^{\frac{3}{2}}}{3}-t+2 \sqrt{t}-2 \ln (\sqrt{t}+1)+c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.012 (sec). Leaf size: 25
DSolve[y' $[t]==1 /(1+$ Sqrt [ t$])$, $\mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $->$ True]

$$
y(t) \rightarrow 2 \sqrt{t}-2 \log (\sqrt{t}+1)+c_{1}
$$

## 2 Problem 1.1-3, page 6

2.1 problem 1.1-3 (a) ..... 25
2.2 problem 1.1-3 (b) ..... 29
2.3 problem 1.1-3 (c) ..... 34
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2.5 problem 1.1-3 (e) ..... 51
2.6 problem 1.1-3 (f) ..... 55

## 2.1 problem 1.1-3 (a)

2.1.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 25
2.1.2 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 26
2.1.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 27

Internal problem ID [2454]
Internal file name [OUTPUT/1946_Sunday_June_05_2022_02_40_22_AM_70399770/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-3, page 6
Problem number: 1.1-3 (a).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}-2 y=-4
$$

With initial conditions

$$
[y(0)=5]
$$

### 2.1.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =-2 \\
q(t) & =-4
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-2 y=-4
$$

The domain of $p(t)=-2$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=0$ is inside this domain. The domain of $q(t)=-4$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=0$ is also inside this domain. Hence solution exists and is unique.

### 2.1.2 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{2 y-4} d y & =\int d t \\
\frac{\ln (y-2)}{2} & =t+c_{1}
\end{aligned}
$$

Raising both side to exponential gives

$$
\sqrt{y-2}=\mathrm{e}^{t+c_{1}}
$$

Which simplifies to

$$
\sqrt{y-2}=c_{2} \mathrm{e}^{t}
$$

Initial conditions are used to solve for $c_{2}$. Substituting $t=0$ and $y=5$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 5=c_{2}^{2}+2 \\
& c_{2}=-\sqrt{3}
\end{aligned}
$$

Substituting $c_{2}$ found above in the general solution gives

$$
y=3 \mathrm{e}^{2 t}+2
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=3 \mathrm{e}^{2 t}+2 \tag{1}
\end{equation*}
$$



Verification of solutions

$$
y=3 \mathrm{e}^{2 t}+2
$$

Verified OK.

### 2.1.3 Maple step by step solution

Let's solve
$\left[y^{\prime}-2 y=-4, y(0)=5\right]$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Separate variables
$\frac{y^{\prime}}{2 y-4}=1$
- Integrate both sides with respect to $t$
$\int \frac{y^{\prime}}{2 y-4} d t=\int 1 d t+c_{1}$
- Evaluate integral
$\frac{\ln (y-2)}{2}=t+c_{1}$
- $\quad$ Solve for $y$

$$
y=\mathrm{e}^{2 t+2 c_{1}}+2
$$

- Use initial condition $y(0)=5$

$$
5=\mathrm{e}^{2 c_{1}}+2
$$

- $\quad$ Solve for $c_{1}$
$c_{1}=\frac{\ln (3)}{2}$
- Substitute $c_{1}=\frac{\ln (3)}{2}$ into general solution and simplify

$$
y=3 \mathrm{e}^{2 t}+2
$$

- $\quad$ Solution to the IVP

$$
y=3 \mathrm{e}^{2 t}+2
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 12

```
dsolve([diff(y(t),t)=2*y(t)-4,y(0) = 5],y(t), singsol=all)
```

$$
y(t)=2+3 \mathrm{e}^{2 t}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.025 (sec). Leaf size: 14

```
DSolve[{y'[t]==2*y[t]-4,y[0]==5},y[t],t,IncludeSingularSolutions -> True]
```

$$
y(t) \rightarrow 3 e^{2 t}+2
$$

## 2.2 problem 1.1-3 (b)

2.2.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 29
2.2.2 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 30
2.2.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 31

Internal problem ID [2455]
Internal file name [OUTPUT/1947_Sunday_June_05_2022_02_40_25_AM_30918361/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-3, page 6
Problem number: 1.1-3 (b).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}+y^{3}=0
$$

With initial conditions

$$
[y(1)=3]
$$

### 2.2.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$
\begin{aligned}
y^{\prime} & =f(t, y) \\
& =-y^{3}
\end{aligned}
$$

The $y$ domain of $f(t, y)$ when $t=1$ is

$$
\{-\infty<y<\infty\}
$$

And the point $y_{0}=3$ is inside this domain. Now we will look at the continuity of

$$
\begin{aligned}
\frac{\partial f}{\partial y} & =\frac{\partial}{\partial y}\left(-y^{3}\right) \\
& =-3 y^{2}
\end{aligned}
$$

The $y$ domain of $\frac{\partial f}{\partial y}$ when $t=1$ is

$$
\{-\infty<y<\infty\}
$$

And the point $y_{0}=3$ is inside this domain. Therefore solution exists and is unique.

### 2.2.2 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
\int-\frac{1}{y^{3}} d y & =t+c_{1} \\
\frac{1}{2 y^{2}} & =t+c_{1}
\end{aligned}
$$

Solving for $y$ gives these solutions

$$
\begin{aligned}
& y_{1}=\frac{1}{\sqrt{2 t+2 c_{1}}} \\
& y_{2}=-\frac{1}{\sqrt{2 t+2 c_{1}}}
\end{aligned}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=1$ and $y=3$ in the above solution gives an equation to solve for the constant of integration.

$$
3=-\frac{1}{\sqrt{2+2 c_{1}}}
$$

Warning: Unable to solve for $c_{1}$. No particular solution can be found using given initial conditions for this solution. removing this solution as not valid. Initial conditions are used to solve for $c_{1}$. Substituting $t=1$ and $y=3$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
3=\frac{1}{\sqrt{2+2 c_{1}}} \\
c_{1}=-\frac{17}{18}
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{3}{\sqrt{18 t-17}}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{3}{\sqrt{18 t-17}} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

Verification of solutions

$$
y=\frac{3}{\sqrt{18 t-17}}
$$

Verified OK.

### 2.2.3 Maple step by step solution

Let's solve
$\left[y^{\prime}+y^{3}=0, y(1)=3\right]$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- $\quad$ Separate variables
$\frac{y^{\prime}}{y^{3}}=-1$
- Integrate both sides with respect to $t$

$$
\int \frac{y^{\prime}}{y^{3}} d t=\int(-1) d t+c_{1}
$$

- Evaluate integral
$-\frac{1}{2 y^{2}}=-t+c_{1}$
- $\quad$ Solve for $y$
$\left\{y=\frac{1}{\sqrt{-2 c_{1}+2 t}}, y=-\frac{1}{\sqrt{-2 c_{1}+2 t}}\right\}$
- Use initial condition $y(1)=3$

$$
3=\frac{1}{\sqrt{-2 c_{1}+2}}
$$

- $\quad$ Solve for $c_{1}$
$c_{1}=\frac{17}{18}$
- Substitute $c_{1}=\frac{17}{18}$ into general solution and simplify
$y=\frac{3}{\sqrt{18 t-17}}$
- Use initial condition $y(1)=3$
$3=-\frac{1}{\sqrt{-2 c_{1}+2}}$
- Solution does not satisfy initial condition
- $\quad$ Solution to the IVP
$y=\frac{3}{\sqrt{18 t-17}}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

$\checkmark$ Solution by Maple
Time used: 0.047 (sec). Leaf size: 13
dsolve([diff $(y(t), t)=-y(t) \wedge 3, y(1)=3], y(t)$, singsol=all)

$$
y(t)=\frac{3}{\sqrt{18 t-17}}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.005 (sec). Leaf size: 16
DSolve[\{y' $[\mathrm{t}]==-\mathrm{y}[\mathrm{t}] \wedge 3, \mathrm{y}[1]==3\}, \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow \frac{3}{\sqrt{18 t-17}}
$$

## 2.3 problem 1.1-3 (c)

2.3.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 34
2.3.2 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 35
2.3.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 37
2.3.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 42
2.3.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 45

Internal problem ID [2456]
Internal file name [OUTPUT/1948_Sunday_June_05_2022_02_40_29_AM_86989023/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-3, page 6
Problem number: 1.1-3 (c).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "separable", "first__order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
y^{\prime}-\frac{\mathrm{e}^{t}}{y}=0
$$

With initial conditions

$$
[y(\ln (2))=-8]
$$

### 2.3.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$
\begin{aligned}
y^{\prime} & =f(t, y) \\
& =\frac{\mathrm{e}^{t}}{y}
\end{aligned}
$$

The $t$ domain of $f(t, y)$ when $y=-8$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=\ln (2)$ is inside this domain. The $y$ domain of $f(t, y)$ when $t=\ln (2)$ is

$$
\{y<0 \vee 0<y\}
$$

And the point $y_{0}=-8$ is inside this domain. Now we will look at the continuity of

$$
\begin{aligned}
\frac{\partial f}{\partial y} & =\frac{\partial}{\partial y}\left(\frac{\mathrm{e}^{t}}{y}\right) \\
& =-\frac{\mathrm{e}^{t}}{y^{2}}
\end{aligned}
$$

The $t$ domain of $\frac{\partial f}{\partial y}$ when $y=-8$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=\ln (2)$ is inside this domain. The $y$ domain of $\frac{\partial f}{\partial y}$ when $t=\ln (2)$ is

$$
\{y<0 \vee 0<y\}
$$

And the point $y_{0}=-8$ is inside this domain. Therefore solution exists and is unique.

### 2.3.2 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(t, y) \\
& =f(t) g(y) \\
& =\frac{\mathrm{e}^{t}}{y}
\end{aligned}
$$

Where $f(t)=\mathrm{e}^{t}$ and $g(y)=\frac{1}{y}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{1}{y}} d y & =\mathrm{e}^{t} d t \\
\int \frac{1}{\frac{1}{y}} d y & =\int \mathrm{e}^{t} d t \\
\frac{y^{2}}{2} & =\mathrm{e}^{t}+c_{1}
\end{aligned}
$$

Which results in

$$
\begin{aligned}
& y=\sqrt{2 \mathrm{e}^{t}+2 c_{1}} \\
& y=-\sqrt{2 \mathrm{e}^{t}+2 c_{1}}
\end{aligned}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\ln (2)$ and $y=-8$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
-8=-\sqrt{4+2 c_{1}} \\
c_{1}=30
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=-\sqrt{2 \mathrm{e}^{t}+60}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\ln (2)$ and $y=-8$ in the above solution gives an equation to solve for the constant of integration.

$$
-8=\sqrt{4+2 c_{1}}
$$

Warning: Unable to solve for constant of integration. $\frac{\text { Summary }}{\text { The solution(s) found are the following }}$

$$
y=-\sqrt{2 \mathrm{e}^{t}+60}
$$



## Verification of solutions

$$
y=-\sqrt{2 \mathrm{e}^{t}+60}
$$

Verified OK.

### 2.3.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{\mathrm{e}^{t}}{y} \\
y^{\prime} & =\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the $\operatorname{PDE}(\mathrm{A})$, and can just use the lookup table shown below to find $\xi, \eta$

Table 10: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=\mathrm{e}^{-t} \\
& \eta(t, y)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=y
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d t \\
& =\int \frac{1}{\mathrm{e}^{-t}} d t
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{t}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=\frac{\mathrm{e}^{t}}{y}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =0 \\
R_{y} & =1 \\
S_{t} & =\mathrm{e}^{t} \\
S_{y} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=y \tag{2A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=R
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{R^{2}}{2}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\mathrm{e}^{t}=\frac{y^{2}}{2}+c_{1}
$$

Which simplifies to

$$
\mathrm{e}^{t}=\frac{y^{2}}{2}+c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=\frac{\mathrm{e}^{t}}{y}$ |  | $\frac{d S}{d R}=R$ |
|  |  |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{\rightarrow \rightarrow \rightarrow}$ |  | 吅 |
|  |  |  |
|  |  |  |
| $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ ¢ $\uparrow \uparrow \uparrow \uparrow$ |  |  |
|  | $R=y$ |  |
|  | $S=\mathrm{e}^{t}$ |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\infty]{ }$ |  |  |
| $\rightarrow \rightarrow \rightarrow$ |  |  |
|  |  |  |
|  |  |  |
| $\rightarrow-\cdots 1$ |  |  |

Initial conditions are used to solve for $c_{1}$. Substituting $t=\ln (2)$ and $y=-8$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
2=32+c_{1} \\
c_{1}=-30
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
\mathrm{e}^{t}=\frac{y^{2}}{2}-30
$$

Solving for $y$ from the above gives

$$
y=-\sqrt{2 \mathrm{e}^{t}+60}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\sqrt{2 \mathrm{e}^{t}+60} \tag{1}
\end{equation*}
$$



(a) Solution plot
(b) Slope field plot

Verification of solutions

$$
y=-\sqrt{2 \mathrm{e}^{t}+60}
$$

Verified OK.

### 2.3.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(y) \mathrm{d} y & =\left(\mathrm{e}^{t}\right) \mathrm{d} t \\
\left(-\mathrm{e}^{t}\right) \mathrm{d} t+(y) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =-\mathrm{e}^{t} \\
N(t, y) & =y
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\mathrm{e}^{t}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(y) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int M \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\mathrm{e}^{t} \mathrm{~d} t \\
\phi & =-\mathrm{e}^{t}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=y$. Therefore equation (4) becomes

$$
\begin{equation*}
y=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=y
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int(y) \mathrm{d} y \\
f(y) & =\frac{y^{2}}{2}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\mathrm{e}^{t}+\frac{y^{2}}{2}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\mathrm{e}^{t}+\frac{y^{2}}{2}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\ln (2)$ and $y=-8$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 30=c_{1} \\
& c_{1}=30
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
-\mathrm{e}^{t}+\frac{y^{2}}{2}=30
$$

Solving for $y$ from the above gives

$$
y=-\sqrt{2 \mathrm{e}^{t}+60}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\sqrt{2 \mathrm{e}^{t}+60} \tag{1}
\end{equation*}
$$



## Verification of solutions

$$
y=-\sqrt{2 \mathrm{e}^{t}+60}
$$

Verified OK.

### 2.3.5 Maple step by step solution

Let's solve

$$
\left[y^{\prime}-\frac{\mathrm{e}^{t}}{y}=0, y(\ln (2))=-8\right]
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables
$y^{\prime} y=\mathrm{e}^{t}$
- Integrate both sides with respect to $t$
$\int y^{\prime} y d t=\int \mathrm{e}^{t} d t+c_{1}$
- Evaluate integral
$\frac{y^{2}}{2}=\mathrm{e}^{t}+c_{1}$
- $\quad$ Solve for $y$

$$
\left\{y=\sqrt{2 \mathrm{e}^{t}+2 c_{1}}, y=-\sqrt{2 \mathrm{e}^{t}+2 c_{1}}\right\}
$$

- Use initial condition $y(\ln (2))=-8$

$$
-8=\sqrt{4+2 c_{1}}
$$

- Solution does not satisfy initial condition
- Use initial condition $y(\ln (2))=-8$
$-8=-\sqrt{4+2 c_{1}}$
- $\quad$ Solve for $c_{1}$
$c_{1}=30$
- Substitute $c_{1}=30$ into general solution and simplify $y=-\sqrt{2 \mathrm{e}^{t}+60}$
- Solution to the IVP

$$
y=-\sqrt{2 \mathrm{e}^{t}+60}
$$

Maple trace

- Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
$\checkmark$ Solution by Maple
Time used: 0.078 (sec). Leaf size: 14

```
dsolve([diff(y(t),t)=exp(t)/y(t),y(ln(2)) = -8],y(t), singsol=all)
```

$$
y(t)=-\sqrt{2 \mathrm{e}^{t}+60}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.594 (sec). Leaf size: 21

```
DSolve[{y'[t]==Exp[t]/y[t],y[Log[2]]==-8},y[t],t,IncludeSingularSolutions -> True]
```

$$
y(t) \rightarrow-\sqrt{2} \sqrt{e^{t}+30}
$$

## 2.4 problem 1.1-3 (d)

2.4.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 47
2.4.2 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 48
2.4.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 49

Internal problem ID [2457]
Internal file name [OUTPUT/1949_Sunday_June_05_2022_02_40_31_AM_54207246/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-3, page 6
Problem number: 1.1-3 (d).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}=\mathrm{e}^{2 t} t
$$

With initial conditions

$$
[y(1)=5]
$$

### 2.4.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =0 \\
q(t) & =\mathrm{e}^{2 t} t
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}=\mathrm{e}^{2 t} t
$$

The domain of $p(t)=0$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=1$ is inside this domain. The domain of $q(t)=\mathrm{e}^{2 t} t$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=1$ is also inside this domain. Hence solution exists and is unique.

### 2.4.2 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
y & =\int \mathrm{e}^{2 t} t \mathrm{~d} t \\
& =\frac{(2 t-1) \mathrm{e}^{2 t}}{4}+c_{1}
\end{aligned}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=1$ and $y=5$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
5=\frac{\mathrm{e}^{2}}{4}+c_{1} \\
c_{1}=-\frac{\mathrm{e}^{2}}{4}+5
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{\mathrm{e}^{2 t} t}{2}-\frac{\mathrm{e}^{2 t}}{4}+5-\frac{\mathrm{e}^{2}}{4}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\mathrm{e}^{2 t} t}{2}-\frac{\mathrm{e}^{2 t}}{4}+5-\frac{\mathrm{e}^{2}}{4} \tag{1}
\end{equation*}
$$



## Verification of solutions

$$
y=\frac{\mathrm{e}^{2 t} t}{2}-\frac{\mathrm{e}^{2 t}}{4}+5-\frac{\mathrm{e}^{2}}{4}
$$

Verified OK.

### 2.4.3 Maple step by step solution

Let's solve

$$
\left[y^{\prime}=\mathrm{e}^{2 t} t, y(1)=5\right]
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Integrate both sides with respect to $t$

$$
\int y^{\prime} d t=\int \mathrm{e}^{2 t} t d t+c_{1}
$$

- Evaluate integral
$y=\frac{(2 t-1) \mathrm{e}^{2 t}}{4}+c_{1}$
- $\quad$ Solve for $y$

$$
y=\frac{\mathrm{e}^{2 t} t}{2}-\frac{\mathrm{e}^{2 t}}{4}+c_{1}
$$

- Use initial condition $y(1)=5$
$5=\frac{\mathrm{e}^{2}}{4}+c_{1}$
- $\quad$ Solve for $c_{1}$
$c_{1}=-\frac{\mathrm{e}^{2}}{4}+5$
- Substitute $c_{1}=-\frac{\mathrm{e}^{2}}{4}+5$ into general solution and simplify

$$
y=\frac{(2 t-1) \mathrm{e}^{2 t}}{4}-\frac{\mathrm{e}^{2}}{4}+5
$$

- Solution to the IVP
$y=\frac{(2 t-1) \mathrm{e}^{2 t}}{4}-\frac{\mathrm{e}^{2}}{4}+5$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
<- quadrature successful`
```

$\checkmark$ Solution by Maple
Time used: 0.047 (sec). Leaf size: 21

```
dsolve([diff(y(t),t)=t*exp(2*t),y(1) = 5],y(t), singsol=all)
```

$$
y(t)=\frac{(2 t-1) \mathrm{e}^{2 t}}{4}+5-\frac{\mathrm{e}^{2}}{4}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.007 (sec). Leaf size: 27

```
DSolve[{y'[t]==t*Exp[2*t],y[1]==5},y[t],t,IncludeSingularSolutions -> True]
```

$$
y(t) \rightarrow \frac{1}{4}\left(e^{2 t}(2 t-1)-e^{2}+20\right)
$$

## 2.5 problem 1.1-3 (e)

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Internal problem ID [2458]
Internal file name [OUTPUT/1950_Sunday_June_05_2022_02_40_34_AM_63720430/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-3, page 6
Problem number: 1.1-3 (e).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}=\sin (t)^{2}
$$

With initial conditions

$$
\left[y\left(\frac{\pi}{6}\right)=3\right]
$$

### 2.5.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =0 \\
q(t) & =\sin (t)^{2}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}=\sin (t)^{2}
$$

The domain of $p(t)=0$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=\frac{\pi}{6}$ is inside this domain. The domain of $q(t)=\sin (t)^{2}$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=\frac{\pi}{6}$ is also inside this domain. Hence solution exists and is unique.

### 2.5.2 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
y & =\int \sin (t)^{2} \mathrm{~d} t \\
& =-\frac{\sin (t) \cos (t)}{2}+\frac{t}{2}+c_{1}
\end{aligned}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\frac{\pi}{6}$ and $y=3$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 3=-\frac{\sqrt{3}}{8}+\frac{\pi}{12}+c_{1} \\
& c_{1}=\frac{\sqrt{3}}{8}-\frac{\pi}{12}+3
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=-\frac{\sin (2 t)}{4}+\frac{t}{2}+\frac{\sqrt{3}}{8}-\frac{\pi}{12}+3
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\sin (2 t)}{4}+\frac{t}{2}+\frac{\sqrt{3}}{8}-\frac{\pi}{12}+3 \tag{1}
\end{equation*}
$$



## Verification of solutions

$$
y=-\frac{\sin (2 t)}{4}+\frac{t}{2}+\frac{\sqrt{3}}{8}-\frac{\pi}{12}+3
$$

Verified OK.

### 2.5.3 Maple step by step solution

Let's solve

$$
\left[y^{\prime}=\sin (t)^{2}, y\left(\frac{\pi}{6}\right)=3\right]
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Integrate both sides with respect to $t$
$\int y^{\prime} d t=\int \sin (t)^{2} d t+c_{1}$
- Evaluate integral
$y=-\frac{\sin (t) \cos (t)}{2}+\frac{t}{2}+c_{1}$
- $\quad$ Solve for $y$
$y=-\frac{\sin (t) \cos (t)}{2}+\frac{t}{2}+c_{1}$
- Use initial condition $y\left(\frac{\pi}{6}\right)=3$

$$
3=-\frac{\sqrt{3}}{8}+\frac{\pi}{12}+c_{1}
$$

- $\quad$ Solve for $c_{1}$

$$
c_{1}=\frac{\sqrt{3}}{8}-\frac{\pi}{12}+3
$$

- $\quad$ Substitute $c_{1}=\frac{\sqrt{3}}{8}-\frac{\pi}{12}+3$ into general solution and simplify

$$
y=-\frac{\sin (2 t)}{4}+\frac{t}{2}+\frac{\sqrt{3}}{8}-\frac{\pi}{12}+3
$$

- $\quad$ Solution to the IVP

$$
y=-\frac{\sin (2 t)}{4}+\frac{t}{2}+\frac{\sqrt{3}}{8}-\frac{\pi}{12}+3
$$

Maple trace
-Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature <- quadrature successful`
$\checkmark$ Solution by Maple
Time used: 0.032 (sec). Leaf size: 23

```
dsolve([diff(y(t),t)=sin(t)^2,y(1/6*Pi) = 3],y(t), singsol=all)
```

$$
y(t)=\frac{t}{2}+3-\frac{\pi}{12}+\frac{\sqrt{3}}{8}-\frac{\sin (2 t)}{4}
$$

$\sqrt{ }$ Solution by Mathematica
Time used: 0.008 (sec). Leaf size: 31

```
DSolve[{y'[t]==Sin[t]^2,y[Pi/6]==3},y[t],t,IncludeSingularSolutions -> True]
```

$$
y(t) \rightarrow \frac{1}{24}(3(4 t+\sqrt{3}+24)-6 \sin (2 t)-2 \pi)
$$

## 2.6 problem 1.1-3 (f)

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2.6.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 57

Internal problem ID [2459]
Internal file name [OUTPUT/1951_Sunday_June_05_2022_02_40_37_AM_90808783/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-3, page 6
Problem number: 1.1-3 (f).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}=8 \mathrm{e}^{4 t}+t
$$

With initial conditions

$$
[y(0)=12]
$$

### 2.6.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =0 \\
q(t) & =8 \mathrm{e}^{4 t}+t
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}=8 \mathrm{e}^{4 t}+t
$$

The domain of $p(t)=0$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=0$ is inside this domain. The domain of $q(t)=8 \mathrm{e}^{4 t}+t$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=0$ is also inside this domain. Hence solution exists and is unique.

### 2.6.2 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
y & =\int 8 \mathrm{e}^{4 t}+t \mathrm{~d} t \\
& =\frac{t^{2}}{2}+2 \mathrm{e}^{4 t}+c_{1}
\end{aligned}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=0$ and $y=12$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
12=2+c_{1} \\
c_{1}=10
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{t^{2}}{2}+2 \mathrm{e}^{4 t}+10
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t^{2}}{2}+2 \mathrm{e}^{4 t}+10 \tag{1}
\end{equation*}
$$


(a) Solution plot

## Verification of solutions

$$
y=\frac{t^{2}}{2}+2 \mathrm{e}^{4 t}+10
$$

Verified OK.

### 2.6.3 Maple step by step solution

Let's solve

$$
\left[y^{\prime}=8 \mathrm{e}^{4 t}+t, y(0)=12\right]
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Integrate both sides with respect to $t$

$$
\int y^{\prime} d t=\int\left(8 \mathrm{e}^{4 t}+t\right) d t+c_{1}
$$

- Evaluate integral

$$
y=\frac{t^{2}}{2}+2 \mathrm{e}^{4 t}+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\frac{t^{2}}{2}+2 \mathrm{e}^{4 t}+c_{1}
$$

- Use initial condition $y(0)=12$

$$
12=2+c_{1}
$$

- $\quad$ Solve for $c_{1}$

$$
c_{1}=10
$$

- Substitute $c_{1}=10$ into general solution and simplify

$$
y=\frac{t^{2}}{2}+2 \mathrm{e}^{4 t}+10
$$

- Solution to the IVP

$$
y=\frac{t^{2}}{2}+2 \mathrm{e}^{4 t}+10
$$

Maple trace

- Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
<- quadrature successful`
$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 17

```
dsolve([diff(y(t),t)=8*exp(4*t)+t,y(0) = 12],y(t), singsol=all)
```

$$
y(t)=\frac{t^{2}}{2}+2 \mathrm{e}^{4 t}+10
$$

$\checkmark$ Solution by Mathematica
Time used: 0.011 (sec). Leaf size: 21
DSolve[\{y' $[\mathrm{t}]==8 * \operatorname{Exp}[4 * \mathrm{t}]+\mathrm{t}, \mathrm{y}[0]==12\}, \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow \frac{1}{2}\left(t^{2}+4 e^{4 t}+20\right)
$$

3 Problem 1.1-4, page 7
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## 3.1 problem 1.1-4 (a)

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60
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Internal problem ID [2460]
Internal file name [OUTPUT/1952_Sunday_June_05_2022_02_40_40_AM_57457162/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-4, page 7
Problem number: 1.1-4 (a).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"
Maple gives the following as the ode type
[_separable]

$$
y^{\prime}-\frac{y}{t}=0
$$

### 3.1.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(t, y) \\
& =f(t) g(y) \\
& =\frac{y}{t}
\end{aligned}
$$

Where $f(t)=\frac{1}{t}$ and $g(y)=y$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{y} d y & =\frac{1}{t} d t \\
\int \frac{1}{y} d y & =\int \frac{1}{t} d t \\
\ln (y) & =\ln (t)+c_{1} \\
y & =\mathrm{e}^{\ln (t)+c_{1}} \\
& =c_{1} t
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} t \tag{1}
\end{equation*}
$$



Figure 16: Slope field plot

Verification of solutions

$$
y=c_{1} t
$$

Verified OK.

### 3.1.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =-\frac{1}{t} \\
q(t) & =0
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{y}{t}=0
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{1}{t} d t} \\
& =\frac{1}{t}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t} \mu y & =0 \\
\frac{\mathrm{~d}}{\mathrm{~d} t}\left(\frac{y}{t}\right) & =0
\end{aligned}
$$

Integrating gives

$$
\frac{y}{t}=c_{1}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{t}$ results in

$$
y=c_{1} t
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} t \tag{1}
\end{equation*}
$$



Figure 17: Slope field plot

Verification of solutions

$$
y=c_{1} t
$$

Verified OK.

### 3.1.3 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(t) t$ on the above ode results in new ode in $u(t)$

$$
u^{\prime}(t) t=0
$$

Integrating both sides gives

$$
\begin{aligned}
u(t) & =\int 0 \mathrm{~d} t \\
& =c_{2}
\end{aligned}
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =t u \\
& =t c_{2}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=t c_{2} \tag{1}
\end{equation*}
$$



Figure 18: Slope field plot
Verification of solutions

$$
y=t c_{2}
$$

Verified OK.

### 3.1.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{y}{t} \\
& y^{\prime}=\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 16: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=t \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{t} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{t}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=\frac{y}{t}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =-\frac{y}{t^{2}} \\
S_{y} & =\frac{1}{t}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\frac{y}{t}=c_{1}
$$

Which simplifies to

$$
\frac{y}{t}=c_{1}
$$

Which gives

$$
y=c_{1} t
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=\frac{y}{t}$ |  | $\frac{d S}{d R}=0$ |
|  |  | $\xrightarrow[\rightarrow \rightarrow+\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{ }$ |
| $\therefore$ - |  | $\xrightarrow[\rightarrow \rightarrow+\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \longrightarrow \rightarrow+]{ }$ |
|  |  | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \text { S }}$ |
|  |  | $\xrightarrow{+\rightarrow \rightarrow \rightarrow \rightarrow \longrightarrow \rightarrow}$ |
| $\xrightarrow{\rightarrow \rightarrow-4 \rightarrow \rightarrow \rightarrow- \pm \rightarrow+\infty} \mid$ |  |  |
|  |  | $\xrightarrow{\longrightarrow \rightarrow \longrightarrow \rightarrow \longrightarrow}$ |
|  |  | $\xrightarrow{-2, \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow}$ |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow+\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \longrightarrow]{ }$ |
|  |  | $\rightarrow$ |

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} t \tag{1}
\end{equation*}
$$



Figure 19: Slope field plot
Verification of solutions

$$
y=c_{1} t
$$

Verified OK.

### 3.1.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{y}\right) \mathrm{d} y & =\left(\frac{1}{t}\right) \mathrm{d} t \\
\left(-\frac{1}{t}\right) \mathrm{d} t+\left(\frac{1}{y}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(t, y)=-\frac{1}{t} \\
& N(t, y)=\frac{1}{y}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{t}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}\left(\frac{1}{y}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int M \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\frac{1}{t} \mathrm{~d} t \\
\phi & =-\ln (t)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{y}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{y}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{y}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{y}\right) \mathrm{d} y \\
f(y) & =\ln (y)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\ln (t)+\ln (y)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\ln (t)+\ln (y)
$$

The solution becomes

$$
y=\mathrm{e}^{c_{1}} t
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{c_{1}} t \tag{1}
\end{equation*}
$$



Figure 20: Slope field plot
Verification of solutions

$$
y=\mathrm{e}^{c_{1}} t
$$

Verified OK.

### 3.1.6 Maple step by step solution

Let's solve

$$
y^{\prime}-\frac{y}{t}=0
$$

- Highest derivative means the order of the ODE is 1

```
y
```

- Separate variables

$$
\frac{y^{\prime}}{y}=\frac{1}{t}
$$

- Integrate both sides with respect to $t$

$$
\int \frac{y^{\prime}}{y} d t=\int \frac{1}{t} d t+c_{1}
$$

- Evaluate integral
$\ln (y)=\ln (t)+c_{1}$
- $\quad$ Solve for $y$

$$
y=\mathrm{e}^{c_{1}} t
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 7

```
dsolve(diff(y(t),t)=y(t)/t,y(t), singsol=all)
```

$$
y(t)=c_{1} t
$$

$\checkmark$ Solution by Mathematica
Time used: 0.022 (sec). Leaf size: 14
DSolve[y' $[t]==y[t] / t, y[t], t$, IncludeSingularSolutions $->$ True]

$$
\begin{aligned}
& y(t) \rightarrow c_{1} t \\
& y(t) \rightarrow 0
\end{aligned}
$$

## 3.2 problem 1.1-4 (b)

3.2.1 Solving as separable ode ..... 74
3.2.2 Solving as homogeneousTypeD2 ode ..... 76
3.2.3 Solving as differentialType ode ..... 78
3.2.4 Solving as first order ode lie symmetry lookup ode ..... 79
3.2.5 Solving as exact ode ..... 83
3.2.6 Maple step by step solution ..... 87

Internal problem ID [2461]
Internal file name [OUTPUT/1953_Sunday_June_05_2022_02_40_42_AM_12687518/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-4, page 7
Problem number: 1.1-4 (b).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "separable", "differentialType", "homogeneousTypeD2", "first_order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
y^{\prime}+\frac{t}{y}=0
$$

### 3.2.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(t, y) \\
& =f(t) g(y) \\
& =-\frac{t}{y}
\end{aligned}
$$

Where $f(t)=-t$ and $g(y)=\frac{1}{y}$. Integrating both sides gives

$$
\frac{1}{\frac{1}{y}} d y=-t d t
$$

$$
\begin{aligned}
\int \frac{1}{\frac{1}{y}} d y & =\int-t d t \\
\frac{y^{2}}{2} & =-\frac{t^{2}}{2}+c_{1}
\end{aligned}
$$

Which results in

$$
\begin{aligned}
& y=\sqrt{-t^{2}+2 c_{1}} \\
& y=-\sqrt{-t^{2}+2 c_{1}}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\sqrt{-t^{2}+2 c_{1}}  \tag{1}\\
& y=-\sqrt{-t^{2}+2 c_{1}} \tag{2}
\end{align*}
$$



Figure 21: Slope field plot

## Verification of solutions

$$
y=\sqrt{-t^{2}+2 c_{1}}
$$

Verified OK.

$$
y=-\sqrt{-t^{2}+2 c_{1}}
$$

Verified OK.

### 3.2.2 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(t) t$ on the above ode results in new ode in $u(t)$

$$
u^{\prime}(t) t+u(t)+\frac{1}{u(t)}=0
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(t, u) \\
& =f(t) g(u) \\
& =-\frac{u^{2}+1}{t u}
\end{aligned}
$$

Where $f(t)=-\frac{1}{t}$ and $g(u)=\frac{u^{2}+1}{u}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{u^{2}+1}{u}} d u & =-\frac{1}{t} d t \\
\int \frac{1}{\frac{u^{2}+1}{u}} d u & =\int-\frac{1}{t} d t \\
\frac{\ln \left(u^{2}+1\right)}{2} & =-\ln (t)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\sqrt{u^{2}+1}=\mathrm{e}^{-\ln (t)+c_{2}}
$$

Which simplifies to

$$
\sqrt{u^{2}+1}=\frac{c_{3}}{t}
$$

Which simplifies to

$$
\sqrt{u(t)^{2}+1}=\frac{c_{3} \mathrm{e}^{c_{2}}}{t}
$$

The solution is

$$
\sqrt{u(t)^{2}+1}=\frac{c_{3} \mathrm{e}^{c_{2}}}{t}
$$

Replacing $u(t)$ in the above solution by $\frac{y}{t}$ results in the solution for $y$ in implicit form

$$
\begin{aligned}
\sqrt{\frac{y^{2}}{t^{2}}+1} & =\frac{c_{3} \mathrm{e}^{c_{2}}}{t} \\
\sqrt{\frac{y^{2}+t^{2}}{t^{2}}} & =\frac{c_{3} \mathrm{e}^{c_{2}}}{t}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\sqrt{\frac{y^{2}+t^{2}}{t^{2}}}=\frac{c_{3} \mathrm{c}_{2}}{t} \tag{1}
\end{equation*}
$$



Figure 22: Slope field plot

Verification of solutions

$$
\sqrt{\frac{y^{2}+t^{2}}{t^{2}}}=\frac{c_{3} \mathrm{e}^{c_{2}}}{t}
$$

Verified OK.

### 3.2.3 Solving as differentialType ode

Writing the ode as

$$
\begin{equation*}
y^{\prime}=-\frac{t}{y} \tag{1}
\end{equation*}
$$

Which becomes

$$
\begin{equation*}
(y) d y=(-t) d t \tag{2}
\end{equation*}
$$

But the RHS is complete differential because

$$
(-t) d t=d\left(-\frac{t^{2}}{2}\right)
$$

Hence (2) becomes

$$
(y) d y=d\left(-\frac{t^{2}}{2}\right)
$$

Integrating both sides gives gives these solutions

$$
\begin{aligned}
& y=\sqrt{-t^{2}+2 c_{1}}+c_{1} \\
& y=-\sqrt{-t^{2}+2 c_{1}}+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\sqrt{-t^{2}+2 c_{1}}+c_{1}  \tag{1}\\
& y=-\sqrt{-t^{2}+2 c_{1}}+c_{1} \tag{2}
\end{align*}
$$



Figure 23: Slope field plot
Verification of solutions

$$
y=\sqrt{-t^{2}+2 c_{1}}+c_{1}
$$

Verified OK.

$$
y=-\sqrt{-t^{2}+2 c_{1}}+c_{1}
$$

Verified OK.

### 3.2.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{t}{y} \\
y^{\prime} & =\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the $\operatorname{PDE}(A)$, and can just use the lookup table shown below to find $\xi, \eta$

Table 19: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=-\frac{1}{t} \\
& \eta(t, y)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=y
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d t \\
& =\int \frac{1}{-\frac{1}{t}} d t
\end{aligned}
$$

Which results in

$$
S=-\frac{t^{2}}{2}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=-\frac{t}{y}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =0 \\
R_{y} & =1 \\
S_{t} & =-t \\
S_{y} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=y \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=R
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{R^{2}}{2}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
-\frac{t^{2}}{2}=\frac{y^{2}}{2}+c_{1}
$$

Which simplifies to

$$
-\frac{t^{2}}{2}=\frac{y^{2}}{2}+c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical <br> coordinates <br> transformation | ODE in canonical coordinates <br> $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=-\frac{t}{y}$ |  | $\frac{d S}{d R}=R$ |
| 多 |  |  |
| 多 |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
-\frac{t^{2}}{2}=\frac{y^{2}}{2}+c_{1} \tag{1}
\end{equation*}
$$



Figure 24: Slope field plot
Verification of solutions

$$
-\frac{t^{2}}{2}=\frac{y^{2}}{2}+c_{1}
$$

Verified OK.

### 3.2.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(-y) \mathrm{d} y & =(t) \mathrm{d} t \\
(-t) \mathrm{d} t+(-y) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =-t \\
N(t, y) & =-y
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(-t) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(-y) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int M \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-t \mathrm{~d} t \\
\phi & =-\frac{t^{2}}{2}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=-y$. Therefore equation (4) becomes

$$
\begin{equation*}
-y=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=-y
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int(-y) \mathrm{d} y \\
f(y) & =-\frac{y^{2}}{2}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\frac{t^{2}}{2}-\frac{y^{2}}{2}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\frac{t^{2}}{2}-\frac{y^{2}}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-\frac{t^{2}}{2}-\frac{y^{2}}{2}=c_{1} \tag{1}
\end{equation*}
$$



Figure 25: Slope field plot
Verification of solutions

$$
-\frac{t^{2}}{2}-\frac{y^{2}}{2}=c_{1}
$$

Verified OK.

### 3.2.6 Maple step by step solution

Let's solve

$$
y^{\prime}+\frac{t}{y}=0
$$

- Highest derivative means the order of the ODE is 1

```
y'
```

- Separate variables

$$
y^{\prime} y=-t
$$

- Integrate both sides with respect to $t$

$$
\int y^{\prime} y d t=\int-t d t+c_{1}
$$

- Evaluate integral
$\frac{y^{2}}{2}=-\frac{t^{2}}{2}+c_{1}$
- $\quad$ Solve for $y$

$$
\left\{y=\sqrt{-t^{2}+2 c_{1}}, y=-\sqrt{-t^{2}+2 c_{1}}\right\}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 27

```
dsolve(diff(y(t),t)=-t/y(t),y(t), singsol=all)
```

$$
\begin{aligned}
& y(t)=\sqrt{-t^{2}+c_{1}} \\
& y(t)=-\sqrt{-t^{2}+c_{1}}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.105 (sec). Leaf size: 39
DSolve[y'[t]==-t/y[t],y[t],t,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(t) \rightarrow-\sqrt{-t^{2}+2 c_{1}} \\
& y(t) \rightarrow \sqrt{-t^{2}+2 c_{1}}
\end{aligned}
$$

## 4 Problem 1.1-5, page 7

4.1 problem 1.1-5 . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 90

## 4.1 problem 1.1-5

4.1.1 Solving as quadrature ode90
4.1.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 91

Internal problem ID [2462]
Internal file name [OUTPUT/1954_Sunday_June_05_2022_02_40_45_AM_34903495/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-5, page 7
Problem number: 1.1-5.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}-y^{2}+y=0
$$

### 4.1.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{y^{2}-y} d y & =\int d t \\
\ln (y-1)-\ln (y) & =t+c_{1}
\end{aligned}
$$

Raising both side to exponential gives

$$
\mathrm{e}^{\ln (y-1)-\ln (y)}=\mathrm{e}^{t+c_{1}}
$$

Which simplifies to

$$
\frac{y-1}{y}=c_{2} \mathrm{e}^{t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{1}{-1+c_{2} \mathrm{e}^{t}} \tag{1}
\end{equation*}
$$



Figure 26: Slope field plot
Verification of solutions

$$
y=-\frac{1}{-1+c_{2} \mathrm{e}^{t}}
$$

Verified OK.

### 4.1.2 Maple step by step solution

Let's solve
$y^{\prime}-y^{2}+y=0$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables
$\frac{y^{\prime}}{y^{2}-y}=1$
- Integrate both sides with respect to $t$
$\int \frac{y^{\prime}}{y^{2}-y} d t=\int 1 d t+c_{1}$
- Evaluate integral

$$
\ln (y-1)-\ln (y)=t+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=-\frac{1}{\mathrm{e}^{t+c_{1}-1}}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 12

```
dsolve(diff(y(t),t)=y(t)^2-y(t),y(t), singsol=all)
```

$$
y(t)=\frac{1}{1+\mathrm{e}^{t} c_{1}}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.242 (sec). Leaf size: 25
DSolve[y' $[t]==y[t] \wedge 2-y[t], y[t], t$, IncludeSingularSolutions $->$ True]

$$
\begin{aligned}
& y(t) \rightarrow \frac{1}{1+e^{t+c_{1}}} \\
& y(t) \rightarrow 0 \\
& y(t) \rightarrow 1
\end{aligned}
$$

5 Problem 1.1-6, page 7
5.1 problem 1.1-6 (a) ..... 94
5.2 problem 1.1-6 (b) ..... 97
5.3 problem 1.1-6 (c) ..... 100
5.4 problem 1.1-6 (d) ..... 103

## 5.1 problem 1.1-6 (a)

> 5.1.1 Solving as quadrature ode
5.1.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 95

Internal problem ID [2463]
Internal file name [OUTPUT/1955_Sunday_June_05_2022_02_40_48_AM_33782510/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-6, page 7
Problem number: 1.1-6 (a).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}-y=-1
$$

### 5.1.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{y-1} d y & =\int d t \\
\ln (y-1) & =t+c_{1}
\end{aligned}
$$

Raising both side to exponential gives

$$
y-1=\mathrm{e}^{t+c_{1}}
$$

Which simplifies to

$$
y-1=c_{2} \mathrm{e}^{t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{2} \mathrm{e}^{t}+1 \tag{1}
\end{equation*}
$$



Figure 27: Slope field plot
Verification of solutions

$$
y=c_{2} \mathrm{e}^{t}+1
$$

## Verified OK.

### 5.1.2 Maple step by step solution

> Let's solve

$$
y^{\prime}-y=-1
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- $\quad$ Separate variables
$\frac{y^{\prime}}{y-1}=1$
- Integrate both sides with respect to $t$
$\int \frac{y^{\prime}}{y-1} d t=\int 1 d t+c_{1}$
- Evaluate integral

$$
\ln (y-1)=t+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\mathrm{e}^{t+c_{1}}+1
$$

Maple trace

- Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 10
dsolve(diff $(y(t), t)=y(t)-1, y(t)$, singsol=all)

$$
y(t)=1+\mathrm{e}^{t} c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.023 (sec). Leaf size: 18
DSolve [y' $[\mathrm{t}]==\mathrm{y}[\mathrm{t}]-1, \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $->$ True]

$$
\begin{aligned}
& y(t) \rightarrow 1+c_{1} e^{t} \\
& y(t) \rightarrow 1
\end{aligned}
$$

## 5.2 problem 1.1-6 (b)

$$
\text { 5.2.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . } 97
$$

5.2.2 Maple step by step solution ..... 98

Internal problem ID [2464]
Internal file name [OUTPUT/1956_Sunday_June_05_2022_02_40_50_AM_17490583/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-6, page 7
Problem number: 1.1-6 (b).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}+y=1
$$

### 5.2.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{1-y} d y & =\int d t \\
-\ln (1-y) & =t+c_{1}
\end{aligned}
$$

Raising both side to exponential gives

$$
\frac{1}{1-y}=\mathrm{e}^{t+c_{1}}
$$

Which simplifies to

$$
\frac{1}{1-y}=c_{2} \mathrm{e}^{t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\mathrm{e}^{-t}}{c_{2}}+1 \tag{1}
\end{equation*}
$$



Figure 28: Slope field plot

Verification of solutions

$$
y=-\frac{\mathrm{e}^{-t}}{c_{2}}+1
$$

Verified OK.

### 5.2.2 Maple step by step solution

Let's solve

$$
y^{\prime}+y=1
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables
$\frac{y^{\prime}}{1-y}=1$
- Integrate both sides with respect to $t$

$$
\int \frac{y^{\prime}}{1-y} d t=\int 1 d t+c_{1}
$$

- Evaluate integral

$$
-\ln (1-y)=t+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=-\mathrm{e}^{-t-c_{1}}+1
$$

Maple trace

```
Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 12

```
dsolve(diff(y(t),t)=1-y(t),y(t), singsol=all)
```

$$
y(t)=1+\mathrm{e}^{-t} c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.023 (sec). Leaf size: 20
DSolve[y'[t]==1-y[t],y[t],t,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(t) \rightarrow 1+c_{1} e^{-t} \\
& y(t) \rightarrow 1
\end{aligned}
$$

## 5.3 problem 1.1-6 (c)

5.3.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 100
5.3.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 101

Internal problem ID [2465]
Internal file name [OUTPUT/1957_Sunday_June_05_2022_02_40_52_AM_89102964/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-6, page 7
Problem number: 1.1-6 (c).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}-y^{3}+y^{2}=0
$$

### 5.3.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{y^{3}-y^{2}} d y & =\int d t \\
\int^{y} \frac{1}{-^{3}-\_a^{2}} d \_a & =t+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\int^{y} \frac{1}{-a^{3}-\_a^{2}} d \_a=t+c_{1} \tag{1}
\end{equation*}
$$



Figure 29: Slope field plot

Verification of solutions

$$
\int^{y} \frac{1}{-a^{3}-\_a^{2}} d \_a=t+c_{1}
$$

Verified OK.

### 5.3.2 Maple step by step solution

Let's solve

$$
y^{\prime}-y^{3}+y^{2}=0
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- $\quad$ Separate variables

$$
\frac{y^{\prime}}{y^{3}-y^{2}}=1
$$

- Integrate both sides with respect to $t$

$$
\int \frac{y^{\prime}}{y^{3}-y^{2}} d t=\int 1 d t+c_{1}
$$

- Evaluate integral

$$
\ln (y-1)+\frac{1}{y}-\ln (y)=t+c_{1}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

$\checkmark$ Solution by Maple
Time used: 0.109 (sec). Leaf size: 16

```
dsolve(diff(y(t),t)=y(t)^3-y(t)^2,y(t), singsol=all)
```

$$
y(t)=\frac{1}{\text { LambertW }\left(-c_{1} \mathrm{e}^{t-1}\right)+1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.227 (sec). Leaf size: 38

```
DSolve[y'[t]==y[t]^3-y[t]^2,y[t],t,IncludeSingularSolutions -> True]
```

$$
\begin{aligned}
& y(t) \rightarrow \text { InverseFunction }\left[\frac{1}{\# 1}+\log (1-\# 1)-\log (\# 1) \&\right]\left[t+c_{1}\right] \\
& y(t) \rightarrow 0 \\
& y(t) \rightarrow 1
\end{aligned}
$$

## 5.4 problem 1.1-6 (d)

5.4.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 103
5.4.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 104

Internal problem ID [2466]
Internal file name [OUTPUT/1958_Sunday_June_05_2022_02_40_58_AM_15291467/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.1-6, page 7
Problem number: 1.1-6 (d).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}+y^{2}=1
$$

### 5.4.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{-y^{2}+1} d y & =t+c_{1} \\
\operatorname{arctanh}(y) & =t+c_{1}
\end{aligned}
$$

Solving for $y$ gives these solutions

$$
y_{1}=\tanh \left(t+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\tanh \left(t+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 30: Slope field plot
Verification of solutions

$$
y=\tanh \left(t+c_{1}\right)
$$

Verified OK.

### 5.4.2 Maple step by step solution

Let's solve

$$
y^{\prime}+y^{2}=1
$$

- Highest derivative means the order of the ODE is 1
- $\quad$ Separate variables

$$
\frac{y^{\prime}}{1-y^{2}}=1
$$

- Integrate both sides with respect to $t$
$\int \frac{y^{\prime}}{1-y^{2}} d t=\int 1 d t+c_{1}$
- Evaluate integral

$$
\operatorname{arctanh}(y)=t+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\tanh \left(t+c_{1}\right)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 8

```
dsolve(diff(y(t),t)=1-y(t)^2,y(t), singsol=all)
```

$$
y(t)=\tanh \left(t+c_{1}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.713 (sec). Leaf size: 44
DSolve[y'[t]==1-y[t]^2,y[t],t,IncludeSingularSolutions -> True]

$$
\begin{aligned}
& y(t) \rightarrow \frac{e^{2 t}-e^{2 c_{1}}}{e^{2 t}+e^{2 c_{1}}} \\
& y(t) \rightarrow-1 \\
& y(t) \rightarrow 1
\end{aligned}
$$

6 Problem 1.2-1, page 12
6.1 problem 1.2-1 (a) ..... 107
6.2 problem 1.2-1 (b) ..... 122
6.3 problem 1.2-1 (c) ..... 125
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6.5 problem 1.2-1 (e) ..... 151
6.6 problem 1.2-1 (f) ..... 164
6.7 problem 1.2-1 (g) ..... 177
6.8 problem 1.2-1 (h) ..... 190
6.9 problem 1.2-1 (i) ..... 203

## 6.1 problem 1.2-1 (a)

6.1.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 107
6.1.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 109
6.1.3 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 110
6.1.4 Solving as first order ode lie symmetry lookup ode . . . . . . . 112
6.1.5 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 116
6.1.6 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 120

Internal problem ID [2467]
Internal file name [OUTPUT/1959_Sunday_June_05_2022_02_41_00_AM_8212680/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-1, page 12
Problem number: 1.2-1 (a).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
y^{\prime}-\left(t^{2}+1\right) y=0
$$

### 6.1.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(t, y) \\
& =f(t) g(y) \\
& =\left(t^{2}+1\right) y
\end{aligned}
$$

Where $f(t)=t^{2}+1$ and $g(y)=y$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{y} d y & =t^{2}+1 d t \\
\int \frac{1}{y} d y & =\int t^{2}+1 d t \\
\ln (y) & =\frac{1}{3} t^{3}+t+c_{1} \\
y & =\mathrm{e}^{\frac{1}{3} t^{3}+t+c_{1}} \\
& =c_{1} \mathrm{e}^{\frac{1}{t^{3}}+t}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} \mathrm{e}^{\frac{1}{3} t^{3}+t} \tag{1}
\end{equation*}
$$



Figure 31: Slope field plot
Verification of solutions

$$
y=c_{1} \mathrm{e}^{\frac{1}{3} t^{3}+t}
$$

Verified OK.

### 6.1.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
& p(t)=-t^{2}-1 \\
& q(t)=0
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+\left(-t^{2}-1\right) y=0
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int\left(-t^{2}-1\right) d t} \\
& =\mathrm{e}^{-\frac{1}{3} t^{3}-t}
\end{aligned}
$$

Which simplifies to

$$
\mu=\mathrm{e}^{-\frac{t\left(t^{2}+3\right)}{3}}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t} \mu y & =0 \\
\frac{\mathrm{~d}}{\mathrm{~d} t}\left(\mathrm{e}^{-\frac{t\left(t^{2}+3\right)}{3}} y\right) & =0
\end{aligned}
$$

Integrating gives

$$
\mathrm{e}^{-\frac{t\left(t^{2}+3\right)}{3}} y=c_{1}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{-\frac{t\left(t^{2}+3\right)}{3}}$ results in

$$
y=c_{1} \mathrm{e}^{\frac{t\left(t^{2}+3\right)}{3}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} \mathrm{e}^{\frac{t\left(t^{2}+3\right)}{3}} \tag{1}
\end{equation*}
$$



Figure 32: Slope field plot
Verification of solutions

$$
y=c_{1} \mathrm{e}^{\frac{t\left(t^{2}+3\right)}{3}}
$$

Verified OK.

### 6.1.3 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(t) t$ on the above ode results in new ode in $u(t)$

$$
u^{\prime}(t) t+u(t)-\left(t^{2}+1\right) u(t) t=0
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(t, u) \\
& =f(t) g(u) \\
& =\frac{u\left(t^{3}+t-1\right)}{t}
\end{aligned}
$$

Where $f(t)=\frac{t^{3}+t-1}{t}$ and $g(u)=u$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{u} d u & =\frac{t^{3}+t-1}{t} d t \\
\int \frac{1}{u} d u & =\int \frac{t^{3}+t-1}{t} d t \\
\ln (u) & =\frac{t^{3}}{3}+t-\ln (t)+c_{2} \\
u & =\mathrm{e}^{\frac{t^{3}}{3}+t-\ln (t)+c_{2}} \\
& =c_{2} \mathrm{e}^{\frac{t^{3}}{3}+t-\ln (t)}
\end{aligned}
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =u t \\
& =t c_{2} \mathrm{e}^{t^{3}}+t-\ln (t)
\end{aligned}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=t c_{2} \mathrm{e}^{t^{3}+t-\ln (t)} \tag{1}
\end{equation*}
$$



Figure 33: Slope field plot

Verification of solutions

$$
y=t c_{2} \mathrm{e}^{t^{3}+t-\ln (t)}
$$

Verified OK.

### 6.1.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\left(t^{2}+1\right) y \\
& y^{\prime}=\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 27: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{2}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=\mathrm{e}^{\frac{1}{3} t^{3}+t} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{\frac{1}{3} t^{3}+t}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{-\frac{1}{3} t^{3}-t} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=\left(t^{2}+1\right) y
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =-\left(t^{2}+1\right) \mathrm{e}^{-\frac{t\left(t^{2}+3\right)}{3}} y \\
S_{y} & =\mathrm{e}^{-\frac{t\left(t^{2}+3\right)}{3}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\mathrm{e}^{-\frac{t\left(t^{2}+3\right)}{3}} y=c_{1}
$$

Which simplifies to

$$
\mathrm{e}^{-\frac{t\left(t^{2}+3\right)}{3}} y=c_{1}
$$

Which gives

$$
y=c_{1} \mathrm{e}^{\frac{t\left(t^{2}+3\right)}{3}}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=\left(t^{2}+1\right) y$ |  | $\frac{d S}{d R}=0$ |
|  |  | $\xrightarrow{\text { a }} \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ 为 |
|  |  | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow}$ |
|  |  |  |
|  |  | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow 2}$ 为 |
|  | $R=t$ | $\rightarrow$ |
|  | $t\left(t^{2}+3\right)$ | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow}$ |
|  | $S=\mathrm{e}^{-\frac{(t 2+3)}{3}} y$ | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+R^{+}}$ |
| Lexty |  | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+2}$ |
|  |  | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow}$ |
| . |  | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-4 \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \longrightarrow}$ |

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} \mathrm{e}^{\frac{t\left(t^{2}+3\right)}{3}} \tag{1}
\end{equation*}
$$



Figure 34: Slope field plot

## Verification of solutions

$$
y=c_{1} \mathrm{e}^{\frac{t\left(t^{2}+3\right)}{3}}
$$

Verified OK.

### 6.1.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{y}\right) \mathrm{d} y & =\left(t^{2}+1\right) \mathrm{d} t \\
\left(-t^{2}-1\right) \mathrm{d} t+\left(\frac{1}{y}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =-t^{2}-1 \\
N(t, y) & =\frac{1}{y}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-t^{2}-1\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}\left(\frac{1}{y}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int M \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-t^{2}-1 \mathrm{~d} t \\
\phi & =-\frac{1}{3} t^{3}-t+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{y}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{y}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{y}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{y}\right) \mathrm{d} y \\
f(y) & =\ln (y)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\frac{t^{3}}{3}-t+\ln (y)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\frac{t^{3}}{3}-t+\ln (y)
$$

The solution becomes

$$
y=\mathrm{e}^{\frac{1}{3} t^{3}+t+c_{1}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{\frac{1}{3} t^{3}+t+c_{1}} \tag{1}
\end{equation*}
$$



Figure 35: Slope field plot

Verification of solutions

$$
y=\mathrm{e}^{\frac{1}{3} t^{3}+t+c_{1}}
$$

Verified OK.

### 6.1.6 Maple step by step solution

Let's solve
$y^{\prime}-\left(t^{2}+1\right) y=0$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables
$\frac{y^{\prime}}{y}=t^{2}+1$
- Integrate both sides with respect to $t$
$\int \frac{y^{\prime}}{y} d t=\int\left(t^{2}+1\right) d t+c_{1}$
- Evaluate integral
$\ln (y)=\frac{1}{3} t^{3}+t+c_{1}$
- $\quad$ Solve for $y$
$y=\mathrm{e}^{\frac{1}{3} t^{3}+t+c_{1}}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 15

```
dsolve(diff(y(t),t)=(t^2+1)*y(t),y(t), singsol=all)
```

$$
y(t)=c_{1} \mathrm{e}^{\frac{t\left(t^{2}+3\right)}{3}}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.025 (sec). Leaf size: 24
DSolve[y'[ t$]==(\mathrm{t} \sim 2+1) * \mathrm{y}[\mathrm{t}], \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(t) \rightarrow c_{1} e^{t^{3}}+t \\
& y(t) \rightarrow 0
\end{aligned}
$$

## 6.2 problem 1.2-1 (b)

6.2.1 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 122
6.2.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 123

Internal problem ID [2468]
Internal file name [OUTPUT/1960_Sunday_June_05_2022_02_41_02_AM_48108646/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-1, page 12
Problem number: 1.2-1 (b).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}+y=0
$$

### 6.2.1 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
\int-\frac{1}{y} d y & =\int d t \\
-\ln (y) & =t+c_{1}
\end{aligned}
$$

Raising both side to exponential gives

$$
\frac{1}{y}=\mathrm{e}^{t+c_{1}}
$$

Which simplifies to

$$
\frac{1}{y}=c_{2} \mathrm{e}^{t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\mathrm{e}^{-t}}{c_{2}} \tag{1}
\end{equation*}
$$



Figure 36: Slope field plot
Verification of solutions

$$
y=\frac{\mathrm{e}^{-t}}{c_{2}}
$$

Verified OK.

### 6.2.2 Maple step by step solution

Let's solve

$$
y^{\prime}+y=0
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables

$$
\frac{y^{\prime}}{y}=-1
$$

- Integrate both sides with respect to $t$

$$
\int \frac{y^{\prime}}{y} d t=\int(-1) d t+c_{1}
$$

- Evaluate integral

$$
\ln (y)=-t+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\mathrm{e}^{-t+c_{1}}
$$

Maple trace

```
Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 10

```
dsolve(diff(y(t),t)=-y(t),y(t), singsol=all)
```

$$
y(t)=\mathrm{e}^{-t} c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.021 (sec). Leaf size: 18
DSolve[y'[t]==-y[t],y[t],t,IncludeSingularSolutions -> True]

$$
\begin{aligned}
& y(t) \rightarrow c_{1} e^{-t} \\
& y(t) \rightarrow 0
\end{aligned}
$$

## 6.3 problem 1.2-1 (c)

6.3.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 125
6.3.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 127
6.3.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 131
6.3.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 135

Internal problem ID [2469]
Internal file name [OUTPUT/1961_Sunday_June_05_2022_02_41_05_AM_14012892/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-1, page 12
Problem number: 1.2-1 (c).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[[_linear, `class A`]]

$$
y^{\prime}-2 y=\mathrm{e}^{-3 t}
$$

### 6.3.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =-2 \\
q(t) & =\mathrm{e}^{-3 t}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-2 y=\mathrm{e}^{-3 t}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int(-2) d t} \\
& =\mathrm{e}^{-2 t}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)\left(\mathrm{e}^{-3 t}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} t}\left(\mathrm{e}^{-2 t} y\right) & =\left(\mathrm{e}^{-2 t}\right)\left(\mathrm{e}^{-3 t}\right) \\
\mathrm{d}\left(\mathrm{e}^{-2 t} y\right) & =\mathrm{e}^{-5 t} \mathrm{~d} t
\end{aligned}
$$

## Integrating gives

$$
\begin{aligned}
& \mathrm{e}^{-2 t} y=\int \mathrm{e}^{-5 t} \mathrm{~d} t \\
& \mathrm{e}^{-2 t} y=-\frac{\mathrm{e}^{-5 t}}{5}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{-2 t}$ results in

$$
y=-\frac{\mathrm{e}^{2 t} \mathrm{e}^{-5 t}}{5}+c_{1} \mathrm{e}^{2 t}
$$

which simplifies to

$$
y=\frac{\left(5 c_{1} \mathrm{e}^{5 t}-1\right) \mathrm{e}^{-3 t}}{5}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\left(5 c_{1} \mathrm{e}^{5 t}-1\right) \mathrm{e}^{-3 t}}{5} \tag{1}
\end{equation*}
$$



Figure 37: Slope field plot
Verification of solutions

$$
y=\frac{\left(5 c_{1} \mathrm{e}^{5 t}-1\right) \mathrm{e}^{-3 t}}{5}
$$

Verified OK.

### 6.3.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =2 y+\mathrm{e}^{-3 t} \\
y^{\prime} & =\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the $\operatorname{PDE}(\mathrm{A})$, and can just use the lookup table shown below to find $\xi, \eta$

Table 31: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=\mathrm{e}^{2 t} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{2 t}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{-2 t} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=2 y+\mathrm{e}^{-3 t}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =-2 \mathrm{e}^{-2 t} y \\
S_{y} & =\mathrm{e}^{-2 t}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\mathrm{e}^{-5 t} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\mathrm{e}^{-5 R}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\frac{\mathrm{e}^{-5 R}}{5}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\mathrm{e}^{-2 t} y=-\frac{\mathrm{e}^{-5 t}}{5}+c_{1}
$$

Which simplifies to

$$
\mathrm{e}^{-2 t} y=-\frac{\mathrm{e}^{-5 t}}{5}+c_{1}
$$

Which gives

$$
y=-\frac{\left(\mathrm{e}^{-5 t}-5 c_{1}\right) \mathrm{e}^{2 t}}{5}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=2 y+\mathrm{e}^{-3 t}$ |  | $\frac{d S}{d R}=\mathrm{e}^{-5 R}$ |
|  |  | 11111111 to |
|  |  |  |
|  |  |  |
| - 1 Patatapatat |  | $1+$ |
| 1 + dogeseses | $R=t$ | $\xrightarrow{\text { cti }} \rightarrow$ |
|  | $S=\mathrm{e}^{-2 t} y$ |  |
|  |  | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \longrightarrow}$ |
|  |  | , $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ + |
| ! |  | $\xrightarrow{\text { L }} \xrightarrow{\text { L }}$ |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\left(\mathrm{e}^{-5 t}-5 c_{1}\right) \mathrm{e}^{2 t}}{5} \tag{1}
\end{equation*}
$$



Figure 38: Slope field plot

## Verification of solutions

$$
y=-\frac{\left(\mathrm{e}^{-5 t}-5 c_{1}\right) \mathrm{e}^{2 t}}{5}
$$

Verified OK.

### 6.3.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(2 y+\mathrm{e}^{-3 t}\right) \mathrm{d} t \\
\left(-2 y-\mathrm{e}^{-3 t}\right) \mathrm{d} t+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =-2 y-\mathrm{e}^{-3 t} \\
N(t, y) & =1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-2 y-\mathrm{e}^{-3 t}\right) \\
& =-2
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =1((-2)-(0)) \\
& =-2
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int-2 \mathrm{~d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-2 t} \\
& =\mathrm{e}^{-2 t}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\mathrm{e}^{-2 t}\left(-2 y-\mathrm{e}^{-3 t}\right) \\
& =-\mathrm{e}^{-2 t}\left(2 y+\mathrm{e}^{-3 t}\right)
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\mathrm{e}^{-2 t}(1) \\
& =\mathrm{e}^{-2 t}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
\left(-\mathrm{e}^{-2 t}\left(2 y+\mathrm{e}^{-3 t}\right)\right)+\left(\mathrm{e}^{-2 t}\right) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\mathrm{e}^{-2 t}\left(2 y+\mathrm{e}^{-3 t}\right) \mathrm{d} t \\
\phi & =\frac{\mathrm{e}^{-5 t}}{5}+\mathrm{e}^{-2 t} y+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\mathrm{e}^{-2 t}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\mathrm{e}^{-2 t}$. Therefore equation (4) becomes

$$
\begin{equation*}
\mathrm{e}^{-2 t}=\mathrm{e}^{-2 t}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{\mathrm{e}^{-5 t}}{5}+\mathrm{e}^{-2 t} y+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{\mathrm{e}^{-5 t}}{5}+\mathrm{e}^{-2 t} y
$$

The solution becomes

$$
y=-\frac{\left(\mathrm{e}^{-5 t}-5 c_{1}\right) \mathrm{e}^{2 t}}{5}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\left(\mathrm{e}^{-5 t}-5 c_{1}\right) \mathrm{e}^{2 t}}{5} \tag{1}
\end{equation*}
$$



Figure 39: Slope field plot
Verification of solutions

$$
y=-\frac{\left(\mathrm{e}^{-5 t}-5 c_{1}\right) \mathrm{e}^{2 t}}{5}
$$

Verified OK.

### 6.3.4 Maple step by step solution

Let's solve
$y^{\prime}-2 y=\mathrm{e}^{-3 t}$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Isolate the derivative
$y^{\prime}=2 y+\mathrm{e}^{-3 t}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-2 y=\mathrm{e}^{-3 t}$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}-2 y\right)=\mu(t) \mathrm{e}^{-3 t}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}-2 y\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=-2 \mu(t)$
- Solve to find the integrating factor
$\mu(t)=\mathrm{e}^{-2 t}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \mu(t) \mathrm{e}^{-3 t} d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \mu(t) \mathrm{e}^{-3 t} d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \mu(t) \mathrm{e}^{-3 t} d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\mathrm{e}^{-2 t}$
$y=\frac{\int \mathrm{e}^{-3 t} \mathrm{e}^{-2 t} d t+c_{1}}{\mathrm{e}^{-2 t}}$
- Evaluate the integrals on the rhs
$y=\frac{-\frac{\mathrm{e}^{-5 t}}{5}+c_{1}}{\mathrm{e}^{-2 t}}$
- Simplify
$y=\frac{\left(5 c \mathrm{e}^{5 t}-1\right) \mathrm{e}^{-3 t}}{5}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 19

```
dsolve(diff(y(t),t)=2*y(t)+exp(-3*t),y(t), singsol=all)
```

$$
y(t)=\frac{\left(5 c_{1} \mathrm{e}^{5 t}-1\right) \mathrm{e}^{-3 t}}{5}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.065 (sec). Leaf size: 23
DSolve [y' $[\mathrm{t}]==2 * \mathrm{y}[\mathrm{t}]+\operatorname{Exp}[-3 * \mathrm{t}], \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $->$ True]

$$
y(t) \rightarrow-\frac{e^{-3 t}}{5}+c_{1} e^{2 t}
$$

## 6.4 problem 1.2-1 (d)

6.4.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 138
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Internal problem ID [2470]
Internal file name [OUTPUT/1962_Sunday_June_05_2022_02_41_08_AM_79796114/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-1, page 12
Problem number: 1.2-1 (d).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[[_linear, `class A`]]

$$
y^{\prime}-2 y=\mathrm{e}^{2 t}
$$

### 6.4.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
& p(t)=-2 \\
& q(t)=\mathrm{e}^{2 t}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-2 y=\mathrm{e}^{2 t}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int(-2) d t} \\
& =\mathrm{e}^{-2 t}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)\left(\mathrm{e}^{2 t}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} t}\left(\mathrm{e}^{-2 t} y\right) & =\left(\mathrm{e}^{-2 t}\right)\left(\mathrm{e}^{2 t}\right) \\
\mathrm{d}\left(\mathrm{e}^{-2 t} y\right) & =\mathrm{d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \mathrm{e}^{-2 t} y=\int \mathrm{d} t \\
& \mathrm{e}^{-2 t} y=t+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{-2 t}$ results in

$$
y=\mathrm{e}^{2 t} t+c_{1} \mathrm{e}^{2 t}
$$

which simplifies to

$$
y=\mathrm{e}^{2 t}\left(t+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{2 t}\left(t+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 40: Slope field plot

Verification of solutions

$$
y=\mathrm{e}^{2 t}\left(t+c_{1}\right)
$$

Verified OK.

### 6.4.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=2 y+\mathrm{e}^{2 t} \\
& y^{\prime}=\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 34: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=\mathrm{e}^{2 t} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{2 t}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{-2 t} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=2 y+\mathrm{e}^{2 t}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =-2 \mathrm{e}^{-2 t} y \\
S_{y} & =\mathrm{e}^{-2 t}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=1 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=1
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=R+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\mathrm{e}^{-2 t} y=t+c_{1}
$$

Which simplifies to

$$
\mathrm{e}^{-2 t} y=t+c_{1}
$$

Which gives

$$
y=\mathrm{e}^{2 t}\left(t+c_{1}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=2 y+\mathrm{e}^{2 t}$ |  | $\frac{d S}{d R}=1$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=t$ |  |
|  |  |  |
|  | $S=\mathrm{e}^{-2 t} y$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{2 t}\left(t+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 41: Slope field plot

## Verification of solutions

$$
y=\mathrm{e}^{2 t}\left(t+c_{1}\right)
$$

Verified OK.

### 6.4.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(2 y+\mathrm{e}^{2 t}\right) \mathrm{d} t \\
\left(-2 y-\mathrm{e}^{2 t}\right) \mathrm{d} t+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =-2 y-\mathrm{e}^{2 t} \\
N(t, y) & =1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-2 y-\mathrm{e}^{2 t}\right) \\
& =-2
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =1((-2)-(0)) \\
& =-2
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int-2 \mathrm{~d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-2 t} \\
& =\mathrm{e}^{-2 t}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\mathrm{e}^{-2 t}\left(-2 y-\mathrm{e}^{2 t}\right) \\
& =-2 \mathrm{e}^{-2 t} y-1
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\mathrm{e}^{-2 t}(1) \\
& =\mathrm{e}^{-2 t}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
\left(-2 \mathrm{e}^{-2 t} y-1\right)+\left(\mathrm{e}^{-2 t}\right) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-2 \mathrm{e}^{-2 t} y-1 \mathrm{~d} t \\
\phi & =-t+\mathrm{e}^{-2 t} y+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\mathrm{e}^{-2 t}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\mathrm{e}^{-2 t}$. Therefore equation (4) becomes

$$
\begin{equation*}
\mathrm{e}^{-2 t}=\mathrm{e}^{-2 t}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-t+\mathrm{e}^{-2 t} y+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-t+\mathrm{e}^{-2 t} y
$$

The solution becomes

$$
y=\mathrm{e}^{2 t}\left(t+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{2 t}\left(t+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 42: Slope field plot

## Verification of solutions

$$
y=\mathrm{e}^{2 t}\left(t+c_{1}\right)
$$

Verified OK.

### 6.4.4 Maple step by step solution

Let's solve
$y^{\prime}-2 y=\mathrm{e}^{2 t}$

- Highest derivative means the order of the ODE is 1

```
y
```

- Isolate the derivative
$y^{\prime}=2 y+\mathrm{e}^{2 t}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-2 y=\mathrm{e}^{2 t}$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}-2 y\right)=\mu(t) \mathrm{e}^{2 t}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}-2 y\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=-2 \mu(t)$
- $\quad$ Solve to find the integrating factor
$\mu(t)=\mathrm{e}^{-2 t}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \mu(t) \mathrm{e}^{2 t} d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \mu(t) \mathrm{e}^{2 t} d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \mu(t) e^{2 t} d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\mathrm{e}^{-2 t}$
$y=\frac{\int \mathrm{e}^{2 t} \mathrm{e}^{-2 t} d t+c_{1}}{\mathrm{e}^{-2 t}}$
- Evaluate the integrals on the rhs
$y=\frac{t+c_{1}}{\mathrm{e}^{-2 t}}$
- Simplify
$y=\mathrm{e}^{2 t}\left(t+c_{1}\right)$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 12
dsolve( $\operatorname{diff}(y(t), t)=2 * y(t)+\exp (2 * t), y(t), \quad$ singsol=all)

$$
y(t)=\left(t+c_{1}\right) \mathrm{e}^{2 t}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.041 (sec). Leaf size: 15
DSolve [y' $[\mathrm{t}]==2 * \mathrm{y}[\mathrm{t}]+\operatorname{Exp}[2 * \mathrm{t}], \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow e^{2 t}\left(t+c_{1}\right)
$$

## 6.5 problem 1.2-1 (e)

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Internal problem ID [2471]

Internal file name [OUTPUT/1963_Sunday_June_05_2022_02_41_10_AM_20563059/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-1, page 12
Problem number: 1.2-1 (e).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"
Maple gives the following as the ode type
[[_linear, `class A`]]

$$
y^{\prime}+y=t
$$

### 6.5.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =1 \\
q(t) & =t
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+y=t
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int 1 d t} \\
& =\mathrm{e}^{t}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)(t) \\
\frac{\mathrm{d}}{\mathrm{~d} t}\left(\mathrm{e}^{t} y\right) & =\left(\mathrm{e}^{t}\right)(t) \\
\mathrm{d}\left(\mathrm{e}^{t} y\right) & =\left(\mathrm{e}^{t} t\right) \mathrm{d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \mathrm{e}^{t} y=\int \mathrm{e}^{t} t \mathrm{~d} t \\
& \mathrm{e}^{t} y=(t-1) \mathrm{e}^{t}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{t}$ results in

$$
y=\mathrm{e}^{-t}(t-1) \mathrm{e}^{t}+c_{1} \mathrm{e}^{-t}
$$

which simplifies to

$$
y=t-1+c_{1} \mathrm{e}^{-t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=t-1+c_{1} \mathrm{e}^{-t} \tag{1}
\end{equation*}
$$



Figure 43: Slope field plot

Verification of solutions

$$
y=t-1+c_{1} \mathrm{e}^{-t}
$$

Verified OK.

### 6.5.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=-y+t \\
& y^{\prime}=\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 37: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=\mathrm{e}^{-t} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{-t}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{t} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=-y+t
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =\mathrm{e}^{t} y \\
S_{y} & =\mathrm{e}^{t}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\mathrm{e}^{t} t \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\mathrm{e}^{R} R
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=(R-1) \mathrm{e}^{R}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\mathrm{e}^{t} y=(t-1) \mathrm{e}^{t}+c_{1}
$$

Which simplifies to

$$
\mathrm{e}^{t} y=(t-1) \mathrm{e}^{t}+c_{1}
$$

Which gives

$$
y=\left(\mathrm{e}^{t} t-\mathrm{e}^{t}+c_{1}\right) \mathrm{e}^{-t}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=-y+t$ |  | $\frac{d S}{d R}=\mathrm{e}^{R} R$ |
|  |  |  |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-\infty \rightarrow+\infty]{\rightarrow \rightarrow \rightarrow-\infty}$ |
|  |  | $\rightarrow \rightarrow$ S ${ }_{\text {S }} \rightarrow$ NT |
|  |  |  |
|  | $R=t$ |  |
|  |  |  |
|  | $S=\mathrm{e}^{t} y$ | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow-\infty]{ }$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  | - $\boldsymbol{\rightarrow} \uparrow \uparrow \uparrow \uparrow \uparrow \uparrow \uparrow$ |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\left(\mathrm{e}^{t} t-\mathrm{e}^{t}+c_{1}\right) \mathrm{e}^{-t} \tag{1}
\end{equation*}
$$



Figure 44: Slope field plot

## Verification of solutions

$$
y=\left(\mathrm{e}^{t} t-\mathrm{e}^{t}+c_{1}\right) \mathrm{e}^{-t}
$$

Verified OK.

### 6.5.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =(-y+t) \mathrm{d} t \\
(y-t) \mathrm{d} t+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =y-t \\
N(t, y) & =1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(y-t) \\
& =1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =1((1)-(0)) \\
& =1
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int 1 \mathrm{~d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{t} \\
& =\mathrm{e}^{t}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\mathrm{e}^{t}(y-t) \\
& =-\mathrm{e}^{t}(-y+t)
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\mathrm{e}^{t}(1) \\
& =\mathrm{e}^{t}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
\left(-\mathrm{e}^{t}(-y+t)\right)+\left(\mathrm{e}^{t}\right) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\mathrm{e}^{t}(-y+t) \mathrm{d} t \\
\phi & =-(t-y-1) \mathrm{e}^{t}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\mathrm{e}^{t}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\mathrm{e}^{t}$. Therefore equation (4) becomes

$$
\begin{equation*}
\mathrm{e}^{t}=\mathrm{e}^{t}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-(t-y-1) \mathrm{e}^{t}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-(t-y-1) \mathrm{e}^{t}
$$

The solution becomes

$$
y=\left(\mathrm{e}^{t} t-\mathrm{e}^{t}+c_{1}\right) \mathrm{e}^{-t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\left(\mathrm{e}^{t} t-\mathrm{e}^{t}+c_{1}\right) \mathrm{e}^{-t} \tag{1}
\end{equation*}
$$



Figure 45: Slope field plot
Verification of solutions

$$
y=\left(\mathrm{e}^{t} t-\mathrm{e}^{t}+c_{1}\right) \mathrm{e}^{-t}
$$

Verified OK.

### 6.5.4 Maple step by step solution

Let's solve
$y^{\prime}+y=t$

- Highest derivative means the order of the ODE is 1

```
y'
```

- Isolate the derivative
$y^{\prime}=-y+t$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}+y=t$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}+y\right)=\mu(t) t$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}+y\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=\mu(t)$
- Solve to find the integrating factor
$\mu(t)=\mathrm{e}^{t}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \mu(t) t d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \mu(t) t d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \mu(t) t d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\mathrm{e}^{t}$
$y=\frac{\int \mathrm{e}^{t} t d t+c_{1}}{\mathrm{e}^{t}}$
- Evaluate the integrals on the rhs
$y=\frac{(t-1) \mathrm{e}^{t}+c_{1}}{\mathrm{e}^{t}}$
- Simplify

$$
y=t-1+c_{1} \mathrm{e}^{-t}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 13
dsolve(diff( $y(t), t)=-y(t)+t, y(t)$, singsol=all)

$$
y(t)=t-1+\mathrm{e}^{-t} c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.026 (sec). Leaf size: 16
DSolve[y' [ t$]==-\mathrm{y}[\mathrm{t}]+\mathrm{t}, \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow t+c_{1} e^{-t}-1
$$

## 6.6 problem 1.2-1 (f)

6.6.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 164
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Internal problem ID [2472]
Internal file name [OUTPUT/1964_Sunday_June_05_2022_02_41_14_AM_71309339/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-1, page 12
Problem number: 1.2-1 (f).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime} t+2 y=\sin (t)
$$

### 6.6.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =\frac{2}{t} \\
q(t) & =\frac{\sin (t)}{t}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+\frac{2 y}{t}=\frac{\sin (t)}{t}
$$

The integrating factor $\mu$ is

$$
\begin{gathered}
\mu=\mathrm{e}^{\int \frac{2}{t} d t} \\
=t^{2}
\end{gathered}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)\left(\frac{\sin (t)}{t}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} t}\left(t^{2} y\right) & =\left(t^{2}\right)\left(\frac{\sin (t)}{t}\right) \\
\mathrm{d}\left(t^{2} y\right) & =(t \sin (t)) \mathrm{d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& t^{2} y=\int t \sin (t) \mathrm{d} t \\
& t^{2} y=-t \cos (t)+\sin (t)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=t^{2}$ results in

$$
y=\frac{-t \cos (t)+\sin (t)}{t^{2}}+\frac{c_{1}}{t^{2}}
$$

which simplifies to

$$
y=\frac{-t \cos (t)+\sin (t)+c_{1}}{t^{2}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{-t \cos (t)+\sin (t)+c_{1}}{t^{2}} \tag{1}
\end{equation*}
$$



Figure 46: Slope field plot

## Verification of solutions

$$
y=\frac{-t \cos (t)+\sin (t)+c_{1}}{t^{2}}
$$

Verified OK.

### 6.6.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{-2 y+\sin (t)}{t} \\
& y^{\prime}=\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 40: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=\frac{1}{t^{2}} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{1}{t^{2}}} d y
\end{aligned}
$$

Which results in

$$
S=t^{2} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=\frac{-2 y+\sin (t)}{t}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =2 y t \\
S_{y} & =t^{2}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=t \sin (t) \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=R \sin (R)
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\sin (R)-R \cos (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
y t^{2}=-t \cos (t)+\sin (t)+c_{1}
$$

Which simplifies to

$$
y t^{2}=-t \cos (t)+\sin (t)+c_{1}
$$

Which gives

$$
y=\frac{-t \cos (t)+\sin (t)+c_{1}}{t^{2}}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=\frac{-2 y+\sin (t)}{t}$ |  | $\frac{d S}{d R}=R \sin (R)$ |
|  |  |  |
|  |  | ! ! |
|  |  |  |
|  |  | ! |
|  |  |  |
|  |  |  |
|  |  |  |
| $\triangle$ 边 | $S=t^{2} y$ | 速 |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{-t \cos (t)+\sin (t)+c_{1}}{t^{2}} \tag{1}
\end{equation*}
$$



Figure 47: Slope field plot

## Verification of solutions

$$
y=\frac{-t \cos (t)+\sin (t)+c_{1}}{t^{2}}
$$

Verified OK.

### 6.6.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(t) \mathrm{d} y & =(-2 y+\sin (t)) \mathrm{d} t \\
(2 y-\sin (t)) \mathrm{d} t+(t) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =2 y-\sin (t) \\
N(t, y) & =t
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(2 y-\sin (t)) \\
& =2
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(t) \\
& =1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =\frac{1}{t}((2)-(1)) \\
& =\frac{1}{t}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int \frac{1}{t} \mathrm{~d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{\ln (t)} \\
& =t
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =t(2 y-\sin (t)) \\
& =(2 y-\sin (t)) t
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =t(t) \\
& =t^{2}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
((2 y-\sin (t)) t)+\left(t^{2}\right) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int(2 y-\sin (t)) t \mathrm{~d} t \\
\phi & =t^{2} y-\sin (t)+t \cos (t)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=t^{2}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=t^{2}$. Therefore equation (4) becomes

$$
\begin{equation*}
t^{2}=t^{2}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=t^{2} y-\sin (t)+t \cos (t)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=t^{2} y-\sin (t)+t \cos (t)
$$

The solution becomes

$$
y=\frac{-t \cos (t)+\sin (t)+c_{1}}{t^{2}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{-t \cos (t)+\sin (t)+c_{1}}{t^{2}} \tag{1}
\end{equation*}
$$



Figure 48: Slope field plot

Verification of solutions

$$
y=\frac{-t \cos (t)+\sin (t)+c_{1}}{t^{2}}
$$

Verified OK.

### 6.6.4 Maple step by step solution

Let's solve
$y^{\prime} t+2 y=\sin (t)$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=-\frac{2 y}{t}+\frac{\sin (t)}{t}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}+\frac{2 y}{t}=\frac{\sin (t)}{t}$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}+\frac{2 y}{t}\right)=\frac{\mu(t) \sin (t)}{t}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}+\frac{2 y}{t}\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=\frac{2 \mu(t)}{t}$
- $\quad$ Solve to find the integrating factor
$\mu(t)=t^{2}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \frac{\mu(t) \sin (t)}{t} d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \frac{\mu(t) \sin (t)}{t} d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \frac{\mu(t) \sin (t)}{t} d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=t^{2}$
$y=\frac{\int t \sin (t) d t+c_{1}}{t^{2}}$
- Evaluate the integrals on the rhs
$y=\frac{-t \cos (t)+\sin (t)+c_{1}}{t^{2}}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 17

```
dsolve(t*diff(y(t),t)+2*y(t)=sin(t),y(t), singsol=all)
```

$$
y(t)=\frac{-\cos (t) t+\sin (t)+c_{1}}{t^{2}}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.034 (sec). Leaf size: 19
DSolve[t*y' $[\mathrm{t}]+2 * \mathrm{y}[\mathrm{t}]==\operatorname{Sin}[\mathrm{t}], \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow \frac{\sin (t)-t \cos (t)+c_{1}}{t^{2}}
$$

## 6.7 problem 1.2-1 (g)

6.7.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 177
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6.7.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 183
6.7.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 187

Internal problem ID [2473]
Internal file name [OUTPUT/1965_Sunday_June_05_2022_02_41_16_AM_55837158/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-1, page 12
Problem number: 1.2-1 (g).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}-y \tan (t)=\sec (t)
$$

### 6.7.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =-\tan (t) \\
q(t) & =\sec (t)
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-y \tan (t)=\sec (t)
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\tan (t) d t} \\
& =\cos (t)
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)(\sec (t)) \\
\frac{\mathrm{d}}{\mathrm{~d} t}(\cos (t) y) & =(\cos (t))(\sec (t)) \\
\mathrm{d}(\cos (t) y) & =\mathrm{d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \cos (t) y=\int \mathrm{d} t \\
& \cos (t) y=t+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\cos (t)$ results in

$$
y=\sec (t) t+c_{1} \sec (t)
$$

which simplifies to

$$
y=\sec (t)\left(t+c_{1}\right)
$$

Summary
The solution(s) found are the following


Figure 49: Slope field plot

Verification of solutions

$$
y=\sec (t)\left(t+c_{1}\right)
$$

Verified OK.

### 6.7.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =y \tan (t)+\sec (t) \\
y^{\prime} & =\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 43: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
\xi(t, y) & =0 \\
\eta(t, y) & =\frac{1}{\cos (t)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{1}{\cos (t)}} d y
\end{aligned}
$$

Which results in

$$
S=\cos (t) y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=y \tan (t)+\sec (t)
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =-\sin (t) y \\
S_{y} & =\cos (t)
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=1 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=1
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=R+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
y \cos (t)=t+c_{1}
$$

Which simplifies to

$$
y \cos (t)=t+c_{1}
$$

Which gives

$$
y=\frac{t+c_{1}}{\cos (t)}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=y \tan (t)+\sec (t)$ |  | $\frac{d S}{d R}=1$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=t$ |  |
|  | $S=\cos (t) y$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t+c_{1}}{\cos (t)} \tag{1}
\end{equation*}
$$



Figure 50: Slope field plot

## Verification of solutions

$$
y=\frac{t+c_{1}}{\cos (t)}
$$

Verified OK.

### 6.7.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =(y \tan (t)+\sec (t)) \mathrm{d} t \\
(-y \tan (t)-\sec (t)) \mathrm{d} t+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =-y \tan (t)-\sec (t) \\
N(t, y) & =1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(-y \tan (t)-\sec (t)) \\
& =-\tan (t)
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =1((-\tan (t))-(0)) \\
& =-\tan (t)
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int-\tan (t) \mathrm{d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{\ln (\cos (t))} \\
& =\cos (t)
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\cos (t)(-y \tan (t)-\sec (t)) \\
& =-\sin (t) y-1
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\cos (t)(1) \\
& =\cos (t)
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
(-\sin (t) y-1)+(\cos (t)) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\sin (t) y-1 \mathrm{~d} t \\
\phi & =-t+\cos (t) y+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\cos (t)+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\cos (t)$. Therefore equation (4) becomes

$$
\begin{equation*}
\cos (t)=\cos (t)+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-t+\cos (t) y+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-t+\cos (t) y
$$

The solution becomes

$$
y=\frac{t+c_{1}}{\cos (t)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t+c_{1}}{\cos (t)} \tag{1}
\end{equation*}
$$



Figure 51: Slope field plot
Verification of solutions

$$
y=\frac{t+c_{1}}{\cos (t)}
$$

Verified OK.

### 6.7.4 Maple step by step solution

Let's solve

$$
y^{\prime}-y \tan (t)=\sec (t)
$$

- Highest derivative means the order of the ODE is 1

$$
y^{\prime}
$$

- Isolate the derivative
$y^{\prime}=y \tan (t)+\sec (t)$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-y \tan (t)=\sec (t)$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}-y \tan (t)\right)=\mu(t) \sec (t)$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}-y \tan (t)\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=-\mu(t) \tan (t)$
- $\quad$ Solve to find the integrating factor
$\mu(t)=\cos (t)$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \mu(t) \sec (t) d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \mu(t) \sec (t) d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \mu(t) \sec (t) d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\cos (t)$
$y=\frac{\int \sec (t) \cos (t) d t+c_{1}}{\cos (t)}$
- Evaluate the integrals on the rhs
$y=\frac{t+c_{1}}{\cos (t)}$
- Simplify
$y=\sec (t)\left(t+c_{1}\right)$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 10

```
dsolve(diff(y(t),t)=y(t)*\operatorname{tan}(t)+sec(t),y(t), singsol=all)
```

$$
y(t)=\sec (t)\left(t+c_{1}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.04 (sec). Leaf size: 12

```
DSolve[y'[t]==y[t]*Tan[t]+Sec[t],y[t],t,IncludeSingularSolutions -> True]
```

$$
y(t) \rightarrow\left(t+c_{1}\right) \sec (t)
$$

## 6.8 problem 1.2-1 (h)

6.8.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 190
6.8.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 192
6.8.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 196
6.8.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 201

Internal problem ID [2474]
Internal file name [OUTPUT/1966_Sunday_June_05_2022_02_41_19_AM_18769874/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-1, page 12
Problem number: 1.2-1 (h).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}-\frac{2 t y}{t^{2}+1}=t+1
$$

### 6.8.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
& p(t)=-\frac{2 t}{t^{2}+1} \\
& q(t)=t+1
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{2 t y}{t^{2}+1}=t+1
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2 t}{t^{2}+1} d t} \\
& =\frac{1}{t^{2}+1}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)(t+1) \\
\frac{\mathrm{d}}{\mathrm{~d} t}\left(\frac{y}{t^{2}+1}\right) & =\left(\frac{1}{t^{2}+1}\right)(t+1) \\
\mathrm{d}\left(\frac{y}{t^{2}+1}\right) & =\left(\frac{t+1}{t^{2}+1}\right) \mathrm{d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{y}{t^{2}+1}=\int \frac{t+1}{t^{2}+1} \mathrm{~d} t \\
& \frac{y}{t^{2}+1}=\frac{\ln \left(t^{2}+1\right)}{2}+\arctan (t)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{t^{2}+1}$ results in

$$
y=\left(t^{2}+1\right)\left(\frac{\ln \left(t^{2}+1\right)}{2}+\arctan (t)\right)+c_{1}\left(t^{2}+1\right)
$$

which simplifies to

$$
y=\left(t^{2}+1\right)\left(\frac{\ln \left(t^{2}+1\right)}{2}+\arctan (t)+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\left(t^{2}+1\right)\left(\frac{\ln \left(t^{2}+1\right)}{2}+\arctan (t)+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 52: Slope field plot

## Verification of solutions

$$
y=\left(t^{2}+1\right)\left(\frac{\ln \left(t^{2}+1\right)}{2}+\arctan (t)+c_{1}\right)
$$

Verified OK.

### 6.8.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{t^{3}+t^{2}+2 y t+t+1}{t^{2}+1} \\
& y^{\prime}=\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 46: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=t^{2}+1 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{t^{2}+1} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{t^{2}+1}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=\frac{t^{3}+t^{2}+2 y t+t+1}{t^{2}+1}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =-\frac{2 y t}{\left(t^{2}+1\right)^{2}} \\
S_{y} & =\frac{1}{t^{2}+1}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{t+1}{t^{2}+1} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{R+1}{R^{2}+1}
$$

The above is a quadrature ode．This is the whole point of Lie symmetry method． It converts an ode，no matter how complicated it is，to one that can be solved by integration when the ode is in the canonical coordiates $R, S$ ．Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{\ln \left(R^{2}+1\right)}{2}+\arctan (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution，we just need to transform（4）back to $t, y$ coordinates．This results in

$$
\frac{y}{t^{2}+1}=\frac{\ln \left(t^{2}+1\right)}{2}+\arctan (t)+c_{1}
$$

Which simplifies to

$$
\frac{y}{t^{2}+1}=\frac{\ln \left(t^{2}+1\right)}{2}+\arctan (t)+c_{1}
$$

Which gives

$$
y=\frac{\left(t^{2}+1\right)\left(\ln \left(t^{2}+1\right)+2 \arctan (t)+2 c_{1}\right)}{2}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown．

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=\frac{t^{3}+t^{2}+2 y t+t+1}{t^{2}+1}$ |  | $\frac{d S}{d R}=\frac{R+1}{R^{2}+1}$ |
|  |  |  |
|  |  |  |
|  |  | $\rightarrow \rightarrow \rightarrow$ 乐多 |
|  |  |  |
|  |  |  |
|  | $R=t$ |  |
|  |  |  |
| $x^{-4}$ | $S=\frac{y}{t^{2}+1}$ |  |
|  | $t^{2}+1$ |  |
|  |  | \％ |
|  |  | $\rightarrow \rightarrow \rightarrow \rightarrow$ 或 |
|  |  |  |
|  |  | $\rightarrow \rightarrow 0$ |

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\left(t^{2}+1\right)\left(\ln \left(t^{2}+1\right)+2 \arctan (t)+2 c_{1}\right)}{2} \tag{1}
\end{equation*}
$$



Figure 53: Slope field plot
Verification of solutions

$$
y=\frac{\left(t^{2}+1\right)\left(\ln \left(t^{2}+1\right)+2 \arctan (t)+2 c_{1}\right)}{2}
$$

Verified OK.

### 6.8.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(\frac{2 t y}{t^{2}+1}+t+1\right) \mathrm{d} t \\
\left(-\frac{2 t y}{t^{2}+1}-t-1\right) \mathrm{d} t+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(t, y)=-\frac{2 t y}{t^{2}+1}-t-1 \\
& N(t, y)=1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{2 t y}{t^{2}+1}-t-1\right) \\
& =-\frac{2 t}{t^{2}+1}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =1\left(\left(-\frac{2 t}{t^{2}+1}\right)-(0)\right) \\
& =-\frac{2 t}{t^{2}+1}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int-\frac{2 t}{t^{2}+1} \mathrm{~d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-\ln \left(t^{2}+1\right)} \\
& =\frac{1}{t^{2}+1}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{t^{2}+1}\left(-\frac{2 t y}{t^{2}+1}-t-1\right) \\
& =\frac{-1-t^{3}-t^{2}+(-2 y-1) t}{\left(t^{2}+1\right)^{2}}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{t^{2}+1}(1) \\
& =\frac{1}{t^{2}+1}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
\left(\frac{-1-t^{3}-t^{2}+(-2 y-1) t}{\left(t^{2}+1\right)^{2}}\right)+\left(\frac{1}{t^{2}+1}\right) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \frac{-1-t^{3}-t^{2}+(-2 y-1) t}{\left(t^{2}+1\right)^{2}} \mathrm{~d} t \\
\phi & =\frac{y}{t^{2}+1}-\frac{\ln \left(t^{2}+1\right)}{2}-\arctan (t)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\frac{1}{t^{2}+1}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{t^{2}+1}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{t^{2}+1}=\frac{1}{t^{2}+1}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{y}{t^{2}+1}-\frac{\ln \left(t^{2}+1\right)}{2}-\arctan (t)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{y}{t^{2}+1}-\frac{\ln \left(t^{2}+1\right)}{2}-\arctan (t)
$$

## Summary

The solution(s) found are the following


Figure 54: Slope field plot

## Verification of solutions

$$
\frac{y}{t^{2}+1}-\frac{\ln \left(t^{2}+1\right)}{2}-\arctan (t)=c_{1}
$$

Verified OK.

### 6.8.4 Maple step by step solution

Let's solve
$y^{\prime}-\frac{2 t y}{t^{2}+1}=t+1$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=\frac{2 t y}{t^{2}+1}+t+1$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-\frac{2 t y}{t^{2}+1}=t+1$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}-\frac{2 t y}{t^{2}+1}\right)=\mu(t)(t+1)$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}-\frac{2 t y}{t^{2}+1}\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=-\frac{2 \mu(t) t}{t^{2}+1}$
- Solve to find the integrating factor
$\mu(t)=\frac{1}{t^{2}+1}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \mu(t)(t+1) d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \mu(t)(t+1) d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \mu(t)(t+1) d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\frac{1}{t^{2}+1}$
$y=\left(t^{2}+1\right)\left(\int \frac{t+1}{t^{2}+1} d t+c_{1}\right)$
- Evaluate the integrals on the rhs
$y=\left(t^{2}+1\right)\left(\frac{\ln \left(t^{2}+1\right)}{2}+\arctan (t)+c_{1}\right)$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 22

```
dsolve(diff(y(t),t)=2*t/(t^2+1)*y(t)+t+1,y(t), singsol=all)
```

$$
y(t)=\left(\frac{\ln \left(t^{2}+1\right)}{2}+\arctan (t)+c_{1}\right)\left(t^{2}+1\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.035 (sec). Leaf size: 26
DSolve[y'[t]==2*t/(t^2+1)*y[t]+t+1,y[t],t,IncludeSingularSolutions $->$ True]

$$
y(t) \rightarrow\left(t^{2}+1\right)\left(\arctan (t)+\frac{1}{2} \log \left(t^{2}+1\right)+c_{1}\right)
$$

## 6.9 problem 1.2-1 (i)

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Internal problem ID [2475]
Internal file name [OUTPUT/1967_Sunday_June_05_2022_02_41_21_AM_12422856/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-1, page 12
Problem number: 1.2-1 (i).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}-y \tan (t)=\sec (t)^{3}
$$

### 6.9.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =-\tan (t) \\
q(t) & =\sec (t)^{3}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-y \tan (t)=\sec (t)^{3}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\tan (t) d t} \\
& =\cos (t)
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)\left(\sec (t)^{3}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} t}(\cos (t) y) & =(\cos (t))\left(\sec (t)^{3}\right) \\
\mathrm{d}(\cos (t) y) & =\sec (t)^{2} \mathrm{~d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \cos (t) y=\int \sec (t)^{2} \mathrm{~d} t \\
& \cos (t) y=\tan (t)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\cos (t)$ results in

$$
y=\sec (t) \tan (t)+c_{1} \sec (t)
$$

which simplifies to

$$
y=\sec (t)\left(\tan (t)+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\sec (t)\left(\tan (t)+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 55: Slope field plot

Verification of solutions

$$
y=\sec (t)\left(\tan (t)+c_{1}\right)
$$

Verified OK.

### 6.9.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =y \tan (t)+\sec (t)^{3} \\
y^{\prime} & =\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 49: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
\xi(t, y) & =0 \\
\eta(t, y) & =\frac{1}{\cos (t)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{1}{\cos (t)}} d y
\end{aligned}
$$

Which results in

$$
S=\cos (t) y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=y \tan (t)+\sec (t)^{3}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =-\sin (t) y \\
S_{y} & =\cos (t)
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\sec (t)^{2} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\sec (R)^{2}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\tan (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
y \cos (t)=\tan (t)+c_{1}
$$

Which simplifies to

$$
y \cos (t)=\tan (t)+c_{1}
$$

Which gives

$$
y=\frac{\tan (t)+c_{1}}{\cos (t)}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=y \tan (t)+\sec (t)^{3}$ |  | $\frac{d S}{d R}=\sec (R)^{2}$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=t$ |  |
|  |  |  |
|  | $S=\cos (t) y$ |  |
| 14, |  |  |
| 19 |  |  |
| ${ }^{+1}$ |  |  |
|  |  |  |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\tan (t)+c_{1}}{\cos (t)} \tag{1}
\end{equation*}
$$



Figure 56: Slope field plot

## Verification of solutions

$$
y=\frac{\tan (t)+c_{1}}{\cos (t)}
$$

Verified OK.

### 6.9.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(y \tan (t)+\sec (t)^{3}\right) \mathrm{d} t \\
\left(-y \tan (t)-\sec (t)^{3}\right) \mathrm{d} t+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =-y \tan (t)-\sec (t)^{3} \\
N(t, y) & =1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-y \tan (t)-\sec (t)^{3}\right) \\
& =-\tan (t)
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =1((-\tan (t))-(0)) \\
& =-\tan (t)
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int-\tan (t) \mathrm{d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{\ln (\cos (t))} \\
& =\cos (t)
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\cos (t)\left(-y \tan (t)-\sec (t)^{3}\right) \\
& =-\sin (t) y-\sec (t)^{2}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\cos (t)(1) \\
& =\cos (t)
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t}=0 \\
\left(-\sin (t) y-\sec (t)^{2}\right)+(\cos (t)) \frac{\mathrm{d} y}{\mathrm{~d} t}=0
\end{array}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\sin (t) y-\sec (t)^{2} \mathrm{~d} t \\
\phi & =-\tan (t)+\cos (t) y+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\cos (t)+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\cos (t)$. Therefore equation (4) becomes

$$
\begin{equation*}
\cos (t)=\cos (t)+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\tan (t)+\cos (t) y+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\tan (t)+\cos (t) y
$$

The solution becomes

$$
y=\frac{\tan (t)+c_{1}}{\cos (t)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\tan (t)+c_{1}}{\cos (t)} \tag{1}
\end{equation*}
$$



Figure 57: Slope field plot

## Verification of solutions

$$
y=\frac{\tan (t)+c_{1}}{\cos (t)}
$$

Verified OK.

### 6.9.4 Maple step by step solution

Let's solve
$y^{\prime}-y \tan (t)=\sec (t)^{3}$

- Highest derivative means the order of the ODE is 1

```
y'
```

- Isolate the derivative
$y^{\prime}=y \tan (t)+\sec (t)^{3}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-y \tan (t)=\sec (t)^{3}$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}-y \tan (t)\right)=\mu(t) \sec (t)^{3}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}-y \tan (t)\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=-\mu(t) \tan (t)$
- $\quad$ Solve to find the integrating factor
$\mu(t)=\cos (t)$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \mu(t) \sec (t)^{3} d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \mu(t) \sec (t)^{3} d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \mu(t) \sec ()^{3} d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\cos (t)$
$y=\frac{\int \sec (t)^{3} \cos (t) d t+c_{1}}{\cos (t)}$
- Evaluate the integrals on the rhs
$y=\frac{\tan (t)+c_{1}}{\cos (t)}$
- Simplify
$y=\sec (t)\left(\tan (t)+c_{1}\right)$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 11

```
dsolve(diff(y(t),t)=y(t)*\operatorname{tan}(t)+sec(t)^3,y(t), singsol=all)
```

$$
y(t)=\sec (t)\left(\tan (t)+c_{1}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.057 (sec). Leaf size: 13
DSolve[y' $[t]==y[t] * \operatorname{Tan}[t]+\operatorname{Sec}[t] \wedge 3, y[t], t$, IncludeSingularSolutions $->$ True $]$

$$
y(t) \rightarrow \sec (t)\left(\tan (t)+c_{1}\right)
$$

## 7 Problem 1.2-2, page 12

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## 7.1 problem 1.2-2 (a)

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Internal problem ID [2476]
Internal file name [OUTPUT/1968_Sunday_June_05_2022_02_41_23_AM_85041395/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-2, page 12
Problem number: 1.2-2 (a).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}-y=0
$$

With initial conditions

$$
[y(0)=2]
$$

### 7.1.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =-1 \\
q(t) & =0
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-y=0
$$

The domain of $p(t)=-1$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=0$ is inside this domain. Hence solution exists and is unique.

### 7.1.2 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{y} d y & =t+c_{1} \\
\ln (y) & =t+c_{1} \\
y & =\mathrm{e}^{t+c_{1}} \\
y & =c_{1} \mathrm{e}^{t}
\end{aligned}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=0$ and $y=2$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 2=c_{1} \\
& c_{1}=2
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=2 \mathrm{e}^{t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=2 \mathrm{e}^{t} \tag{1}
\end{equation*}
$$


(a) Solution plot (b) Slope field plot


Verification of solutions

$$
y=2 \mathrm{e}^{t}
$$

Verified OK.

### 7.1.3 Maple step by step solution

Let's solve
$\left[y^{\prime}-y=0, y(0)=2\right]$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables
$\frac{y^{\prime}}{y}=1$
- Integrate both sides with respect to $t$
$\int \frac{y^{\prime}}{y} d t=\int 1 d t+c_{1}$
- Evaluate integral
$\ln (y)=t+c_{1}$
- $\quad$ Solve for $y$
$y=\mathrm{e}^{t+c_{1}}$
- Use initial condition $y(0)=2$
$2=\mathrm{e}^{c_{1}}$
- $\quad$ Solve for $c_{1}$
$c_{1}=\ln (2)$
- $\quad$ Substitute $c_{1}=\ln (2)$ into general solution and simplify $y=2 \mathrm{e}^{t}$
- Solution to the IVP
$y=2 \mathrm{e}^{t}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.015 (sec). Leaf size: 8

```
dsolve([diff(y(t),t)=y(t),y(0) = 2],y(t), singsol=all)
```

$$
y(t)=2 \mathrm{e}^{t}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.023 (sec). Leaf size: 10

```
DSolve[{y'[t]==y[t],y[0]==2},y[t],t,IncludeSingularSolutions -> True]
```

$$
y(t) \rightarrow 2 e^{t}
$$

## 7.2 problem 1.2-2 (b)

7.2.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 221
7.2.2 Solving as quadrature ode . . . . . . . . . . . . . . . . . . . . . 222
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Internal problem ID [2477]
Internal file name [OUTPUT/1969_Sunday_June_05_2022_02_41_26_AM_44077002/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-2, page 12
Problem number: 1.2-2 (b).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y^{\prime}-2 y=0
$$

With initial conditions

$$
[y(\ln (3))=3]
$$

### 7.2.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =-2 \\
q(t) & =0
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-2 y=0
$$

The domain of $p(t)=-2$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=\ln (3)$ is inside this domain. Hence solution exists and is unique.

### 7.2.2 Solving as quadrature ode

Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{2 y} d y & =\int d t \\
\frac{\ln (y)}{2} & =t+c_{1}
\end{aligned}
$$

Raising both side to exponential gives

$$
\sqrt{y}=\mathrm{e}^{t+c_{1}}
$$

Which simplifies to

$$
\sqrt{y}=c_{2} \mathrm{e}^{t}
$$

Initial conditions are used to solve for $c_{2}$. Substituting $t=\ln (3)$ and $y=3$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
3=9 c_{2}^{2} \\
c_{2}=-\frac{\sqrt{3}}{3}
\end{gathered}
$$

Substituting $c_{2}$ found above in the general solution gives

$$
y=\frac{\mathrm{e}^{2 t}}{3}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\mathrm{e}^{2 t}}{3} \tag{1}
\end{equation*}
$$



Verification of solutions

$$
y=\frac{\mathrm{e}^{2 t}}{3}
$$

Verified OK.

### 7.2.3 Maple step by step solution

Let's solve

$$
\left[y^{\prime}-2 y=0, y(\ln (3))=3\right]
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Separate variables
$\frac{y^{\prime}}{y}=2$
- Integrate both sides with respect to $t$
$\int \frac{y^{\prime}}{y} d t=\int 2 d t+c_{1}$
- Evaluate integral
$\ln (y)=2 t+c_{1}$
- $\quad$ Solve for $y$
$y=\mathrm{e}^{2 t+c_{1}}$
- Use initial condition $y(\ln (3))=3$
$3=\mathrm{e}^{2 \ln (3)+c_{1}}$
- $\quad$ Solve for $c_{1}$

$$
c_{1}=-\ln (3)
$$

- $\quad$ Substitute $c_{1}=-\ln (3)$ into general solution and simplify $y=\frac{\mathrm{e}^{2 t}}{3}$
- Solution to the IVP
$y=\frac{\mathrm{e}^{2 t}}{3}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 10

```
dsolve([diff(y(t),t)=2*y(t),y(ln(3)) = 3],y(t), singsol=all)
```

$$
y(t)=\frac{\mathrm{e}^{2 t}}{3}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.022 (sec). Leaf size: 14
DSolve[\{y' $[t]==2 * y[t], y[\log [3]]==3\}, y[t], t$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow \frac{e^{2 t}}{3}
$$

## 7.3 problem 1.2-2 (c)

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Internal problem ID [2478]
Internal file name [OUTPUT/1970_Sunday_June_05_2022_02_41_28_AM_89034612/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-2, page 12
Problem number: 1.2-2 (c).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "homogeneousTypeD2", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry__lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime} t-y=t^{3}
$$

With initial conditions

$$
[y(1)=-2]
$$

### 7.3.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
& p(t)=-\frac{1}{t} \\
& q(t)=t^{2}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{y}{t}=t^{2}
$$

The domain of $p(t)=-\frac{1}{t}$ is

$$
\{t<0 \vee 0<t\}
$$

And the point $t_{0}=1$ is inside this domain. The domain of $q(t)=t^{2}$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=1$ is also inside this domain. Hence solution exists and is unique.

### 7.3.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{1}{t} d t} \\
& =\frac{1}{t}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)\left(t^{2}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} t}\left(\frac{y}{t}\right) & =\left(\frac{1}{t}\right)\left(t^{2}\right) \\
\mathrm{d}\left(\frac{y}{t}\right) & =t \mathrm{~d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{y}{t}=\int t \mathrm{~d} t \\
& \frac{y}{t}=\frac{t^{2}}{2}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{t}$ results in

$$
y=\frac{1}{2} t^{3}+c_{1} t
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=1$ and $y=-2$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
-2 & =\frac{1}{2}+c_{1} \\
c_{1} & =-\frac{5}{2}
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{1}{2} t^{3}-\frac{5}{2} t
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{1}{2} t^{3}-\frac{5}{2} t \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

## Verification of solutions

$$
y=\frac{1}{2} t^{3}-\frac{5}{2} t
$$

Verified OK.

### 7.3.3 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(t) t$ on the above ode results in new ode in $u(t)$

$$
\left(u^{\prime}(t) t+u(t)\right) t-u(t) t=t^{3}
$$

Integrating both sides gives

$$
\begin{aligned}
u(t) & =\int t \mathrm{~d} t \\
& =\frac{t^{2}}{2}+c_{2}
\end{aligned}
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =u t \\
& =t\left(\frac{t^{2}}{2}+c_{2}\right)
\end{aligned}
$$

Initial conditions are used to solve for $c_{2}$. Substituting $t=1$ and $y=-2$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
-2=\frac{1}{2}+c_{2} \\
c_{2}=-\frac{5}{2}
\end{gathered}
$$

Substituting $c_{2}$ found above in the general solution gives

$$
y=\frac{t\left(t^{2}-5\right)}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t\left(t^{2}-5\right)}{2} \tag{1}
\end{equation*}
$$



## Verification of solutions

$$
y=\frac{t\left(t^{2}-5\right)}{2}
$$

Verified OK.

### 7.3.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{t^{3}+y}{t} \\
y^{\prime} & =\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 54: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=t \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{t} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{t}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=\frac{t^{3}+y}{t}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =-\frac{y}{t^{2}} \\
S_{y} & =\frac{1}{t}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=t \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=R
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{R^{2}}{2}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\frac{y}{t}=\frac{t^{2}}{2}+c_{1}
$$

Which simplifies to

$$
\frac{y}{t}=\frac{t^{2}}{2}+c_{1}
$$

Which gives

$$
y=\frac{t\left(t^{2}+2 c_{1}\right)}{2}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=\frac{t^{3}+y}{t}$ |  | $\frac{d S}{d R}=R$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=t$ |  |
|  | $c \_y$ |  |
|  | $S=\frac{y}{t}$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |

Initial conditions are used to solve for $c_{1}$. Substituting $t=1$ and $y=-2$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
-2 & =\frac{1}{2}+c_{1} \\
c_{1} & =-\frac{5}{2}
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{t\left(t^{2}-5\right)}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t\left(t^{2}-5\right)}{2} \tag{1}
\end{equation*}
$$



(a) Solution plot

Verification of solutions

$$
y=\frac{t\left(t^{2}-5\right)}{2}
$$

## Verified OK.

### 7.3.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(t) \mathrm{d} y & =\left(t^{3}+y\right) \mathrm{d} t \\
\left(-t^{3}-y\right) \mathrm{d} t+(t) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(t, y)=-t^{3}-y \\
& N(t, y)=t
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-t^{3}-y\right) \\
& =-1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(t) \\
& =1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =\frac{1}{t}((-1)-(1)) \\
& =-\frac{2}{t}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int-\frac{2}{t} \mathrm{~d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-2 \ln (t)} \\
& =\frac{1}{t^{2}}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{t^{2}}\left(-t^{3}-y\right) \\
& =\frac{-t^{3}-y}{t^{2}}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{t^{2}}(t) \\
& =\frac{1}{t}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
\left(\frac{-t^{3}-y}{t^{2}}\right)+\left(\frac{1}{t}\right) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \frac{-t^{3}-y}{t^{2}} \mathrm{~d} t \\
\phi & =\frac{-t^{3}+2 y}{2 t}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\frac{1}{t}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{t}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{t}=\frac{1}{t}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{-t^{3}+2 y}{2 t}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{-t^{3}+2 y}{2 t}
$$

The solution becomes

$$
y=\frac{t\left(t^{2}+2 c_{1}\right)}{2}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=1$ and $y=-2$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
-2=\frac{1}{2}+c_{1} \\
c_{1}=-\frac{5}{2}
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{t\left(t^{2}-5\right)}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t\left(t^{2}-5\right)}{2} \tag{1}
\end{equation*}
$$



## Verification of solutions

$$
y=\frac{t\left(t^{2}-5\right)}{2}
$$

Verified OK.

### 7.3.6 Maple step by step solution

Let's solve

$$
\left[y^{\prime} t-y=t^{3}, y(1)=-2\right]
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative

$$
y^{\prime}=\frac{y}{t}+t^{2}
$$

- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-\frac{y}{t}=t^{2}$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}-\frac{y}{t}\right)=\mu(t) t^{2}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}-\frac{y}{t}\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=-\frac{\mu(t)}{t}$
- Solve to find the integrating factor
$\mu(t)=\frac{1}{t}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \mu(t) t^{2} d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \mu(t) t^{2} d t+c_{1}$
- Solve for $y$
$y=\frac{\int \mu(t) t^{2} d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\frac{1}{t}$
$y=t\left(\int t d t+c_{1}\right)$
- Evaluate the integrals on the rhs
$y=t\left(\frac{t^{2}}{2}+c_{1}\right)$
- Simplify
$y=\frac{t\left(t^{2}+2 c_{1}\right)}{2}$
- Use initial condition $y(1)=-2$
$-2=\frac{1}{2}+c_{1}$
- Solve for $c_{1}$
$c_{1}=-\frac{5}{2}$
- Substitute $c_{1}=-\frac{5}{2}$ into general solution and simplify
$y=\frac{t\left(t^{2}-5\right)}{2}$
- Solution to the IVP
$y=\frac{t\left(t^{2}-5\right)}{2}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.015 (sec). Leaf size: 12

```
dsolve([t*diff(y(t),t)=y(t)+t^3,y(1) = - 2],y(t), singsol=all)
```

$$
y(t)=\frac{\left(t^{2}-5\right) t}{2}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.048 (sec). Leaf size: 27
DSolve[\{y' $\left.[t]==y[t]+t^{\wedge} 3, y[1]==-2\right\}, y[t], t$, IncludeSingularSolutions $->$ True]

$$
y(t) \rightarrow-t^{3}-3 t^{2}-6 t+14 e^{t-1}-6
$$

## 7.4 problem 1.2-2 (d)

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Internal problem ID [2479]
Internal file name [OUTPUT/1971_Sunday_June_05_2022_02_41_30_AM_86709491/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-2, page 12
Problem number: 1.2-2 (d).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}+y \tan (t)=\sec (t)
$$

With initial conditions

$$
[y(0)=0]
$$

### 7.4.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =\tan (t) \\
q(t) & =\sec (t)
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+y \tan (t)=\sec (t)
$$

The domain of $p(t)=\tan (t)$ is

$$
\left\{t<\frac{1}{2} \pi+\pi \_Z 83 \vee \frac{1}{2} \pi+\pi \_Z 83<t\right\}
$$

And the point $t_{0}=0$ is inside this domain. The domain of $q(t)=\sec (t)$ is

$$
\left\{t<\frac{1}{2} \pi+\pi_{-} Z 83 \vee \frac{1}{2} \pi+\pi_{-} Z 83<t\right\}
$$

And the point $t_{0}=0$ is also inside this domain. Hence solution exists and is unique.

### 7.4.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \tan (t) d t} \\
& =\frac{1}{\cos (t)}
\end{aligned}
$$

Which simplifies to

$$
\mu=\sec (t)
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)(\sec (t)) \\
\frac{\mathrm{d}}{\mathrm{~d} t}(\sec (t) y) & =(\sec (t))(\sec (t)) \\
\mathrm{d}(\sec (t) y) & =\sec (t)^{2} \mathrm{~d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \sec (t) y=\int \sec (t)^{2} \mathrm{~d} t \\
& \sec (t) y=\tan (t)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\sec (t)$ results in

$$
y=\cos (t) \tan (t)+c_{1} \cos (t)
$$

which simplifies to

$$
y=c_{1} \cos (t)+\sin (t)
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=0$ and $y=0$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 0=c_{1} \\
& c_{1}=0
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\sin (t)
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\sin (t) \tag{1}
\end{equation*}
$$


(a) Solution plot

(b) Slope field plot

## Verification of solutions

$$
y=\sin (t)
$$

Verified OK.

### 7.4.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-y \tan (t)+\sec (t) \\
y^{\prime} & =\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 57: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | special | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |  |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}$ |
| $a_{1} b_{2}-a_{2} b_{1}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |  |  |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=\cos (t) \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\cos (t)} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{\cos (t)}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=-y \tan (t)+\sec (t)
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =\sec (t) \tan (t) y \\
S_{y} & =\sec (t)
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\sec (t)^{2} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\sec (R)^{2}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\tan (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\sec (t) y=\tan (t)+c_{1}
$$

Which simplifies to

$$
\sec (t) y=\tan (t)+c_{1}
$$

Which gives

$$
y=\frac{\tan (t)+c_{1}}{\sec (t)}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates |  | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=-y \tan (t)+\sec (t)$ |  | $\frac{d S}{d R}=\sec (R)^{2}$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=t$ |  |
|  | $S=\sec (t) y$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |

Initial conditions are used to solve for $c_{1}$. Substituting $t=0$ and $y=0$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 0=c_{1} \\
& c_{1}=0
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\sin (t)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\sin (t) \tag{1}
\end{equation*}
$$



## Verification of solutions

$$
y=\sin (t)
$$

Verified OK.

### 7.4.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =(-y \tan (t)+\sec (t)) \mathrm{d} t \\
(y \tan (t)-\sec (t)) \mathrm{d} t+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =y \tan (t)-\sec (t) \\
N(t, y) & =1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(y \tan (t)-\sec (t)) \\
& =\tan (t)
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =1((\tan (t))-(0)) \\
& =\tan (t)
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int \tan (t) \mathrm{d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-\ln (\cos (t))} \\
& =\sec (t)
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\sec (t)(y \tan (t)-\sec (t)) \\
& =\sec (t)^{2}(\sin (t) y-1)
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\sec (t)(1) \\
& =\sec (t)
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
\left(\sec (t)^{2}(\sin (t) y-1)\right)+(\sec (t)) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \sec (t)^{2}(\sin (t) y-1) \mathrm{d} t \\
\phi & =\sec (t) y-\tan (t)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\sec (t)+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\sec (t)$. Therefore equation (4) becomes

$$
\begin{equation*}
\sec (t)=\sec (t)+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\sec (t) y-\tan (t)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\sec (t) y-\tan (t)
$$

The solution becomes

$$
y=\frac{\tan (t)+c_{1}}{\sec (t)}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=0$ and $y=0$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 0=c_{1} \\
& c_{1}=0
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\sin (t)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\sin (t) \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

## Verification of solutions

$$
y=\sin (t)
$$

Verified OK.

### 7.4.5 Maple step by step solution

Let's solve

$$
\left[y^{\prime}+y \tan (t)=\sec (t), y(0)=0\right]
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=-y \tan (t)+\sec (t)$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}+y \tan (t)=\sec (t)$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}+y \tan (t)\right)=\mu(t) \sec (t)$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}+y \tan (t)\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=\mu(t) \tan (t)$
- $\quad$ Solve to find the integrating factor
$\mu(t)=\frac{1}{\cos (t)}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \mu(t) \sec (t) d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \mu(t) \sec (t) d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \mu(t) \sec (t) d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\frac{1}{\cos (t)}$
$y=\cos (t)\left(\int \frac{\sec (t)}{\cos (t)} d t+c_{1}\right)$
- Evaluate the integrals on the rhs
$y=\cos (t)\left(\tan (t)+c_{1}\right)$
- Simplify
$y=c_{1} \cos (t)+\sin (t)$
- Use initial condition $y(0)=0$
$0=c_{1}$
- $\quad$ Solve for $c_{1}$
$c_{1}=0$
- $\quad$ Substitute $c_{1}=0$ into general solution and simplify
$y=\sin (t)$
- Solution to the IVP
$y=\sin (t)$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 6

```
dsolve([diff(y(t),t)=-tan(t)*y(t)+\operatorname{sec}(t),y(0)=0],y(t), singsol=all)
```

$$
y(t)=\sin (t)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.04 (sec). Leaf size: 7
DSolve[\{y' [ t$]==-\operatorname{Tan}[\mathrm{t}] * \mathrm{y}[\mathrm{t}]+\mathrm{Sec}[\mathrm{t}], \mathrm{y}[0]==0\}, \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow \sin (t)
$$

## 7.5 problem 1.2-2 (e)

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Internal problem ID [2480]
Internal file name [OUTPUT/1972_Sunday_June_05_2022_02_41_34_AM_36487731/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-2, page 12
Problem number: 1.2-2 (e).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "homogeneousTypeMapleC", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
y^{\prime}-\frac{2 y}{t+1}=0
$$

With initial conditions

$$
[y(0)=6]
$$

### 7.5.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
& p(t)=-\frac{2}{t+1} \\
& q(t)=0
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{2 y}{t+1}=0
$$

The domain of $p(t)=-\frac{2}{t+1}$ is

$$
\{t<-1 \vee-1<t\}
$$

And the point $t_{0}=0$ is inside this domain. Hence solution exists and is unique.

### 7.5.2 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(t, y) \\
& =f(t) g(y) \\
& =\frac{2 y}{t+1}
\end{aligned}
$$

Where $f(t)=\frac{2}{t+1}$ and $g(y)=y$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{y} d y & =\frac{2}{t+1} d t \\
\int \frac{1}{y} d y & =\int \frac{2}{t+1} d t \\
\ln (y) & =2 \ln (t+1)+c_{1} \\
y & =\mathrm{e}^{2 \ln (t+1)+c_{1}} \\
& =c_{1}(t+1)^{2}
\end{aligned}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=0$ and $y=6$ in the above solution gives an equation to solve for the constant of integration.

$$
6=c_{1}
$$

$$
c_{1}=6
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=6(t+1)^{2}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=6(t+1)^{2} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

Verification of solutions

$$
y=6(t+1)^{2}
$$

Verified OK.

### 7.5.3 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2}{t+1} d t} \\
& =\frac{1}{(t+1)^{2}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t} \mu y & =0 \\
\frac{\mathrm{~d}}{\mathrm{~d} t}\left(\frac{y}{(t+1)^{2}}\right) & =0
\end{aligned}
$$

Integrating gives

$$
\frac{y}{(t+1)^{2}}=c_{1}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{(t+1)^{2}}$ results in

$$
y=c_{1}(t+1)^{2}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=0$ and $y=6$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 6=c_{1} \\
& c_{1}=6
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=6(t+1)^{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=6(t+1)^{2} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

## Verification of solutions

$$
y=6(t+1)^{2}
$$

Verified OK.

### 7.5.4 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(t) t$ on the above ode results in new ode in $u(t)$

$$
u^{\prime}(t) t+u(t)-\frac{2 u(t) t}{t+1}=0
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(t, u) \\
& =f(t) g(u) \\
& =\frac{u(t-1)}{t(t+1)}
\end{aligned}
$$

Where $f(t)=\frac{t-1}{t(t+1)}$ and $g(u)=u$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{u} d u & =\frac{t-1}{t(t+1)} d t \\
\int \frac{1}{u} d u & =\int \frac{t-1}{t(t+1)} d t \\
\ln (u) & =2 \ln (t+1)-\ln (t)+c_{2} \\
u & =\mathrm{e}^{2 \ln (t+1)-\ln (t)+c_{2}} \\
& =c_{2} \mathrm{e}^{2 \ln (t+1)-\ln (t)}
\end{aligned}
$$

Which simplifies to

$$
u(t)=c_{2}\left(t+2+\frac{1}{t}\right)
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =u t \\
& =t c_{2}\left(t+2+\frac{1}{t}\right)
\end{aligned}
$$

Initial conditions are used to solve for $c_{2}$. Substituting $t=0$ and $y=6$ in the above solution gives an equation to solve for the constant of integration.

$$
6=c_{2}
$$

$$
c_{2}=6
$$

Substituting $c_{2}$ found above in the general solution gives

$$
y=6 t^{2}+12 t+6
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=6 t^{2}+12 t+6 \tag{1}
\end{equation*}
$$



(b) Slope field plot

## Verification of solutions

$$
y=6 t^{2}+12 t+6
$$

## Verified OK.

### 7.5.5 Solving as homogeneousTypeMapleC ode

Let $Y=y+y_{0}$ and $X=t+x_{0}$ then the above is transformed to new ode in $Y(X)$

$$
\frac{d}{d X} Y(X)=\frac{2 Y(X)+2 y_{0}}{X+x_{0}+1}
$$

Solving for possible values of $x_{0}$ and $y_{0}$ which makes the above ode a homogeneous ode results in

$$
\begin{aligned}
& x_{0}=-1 \\
& y_{0}=0
\end{aligned}
$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$
\frac{d}{d X} Y(X)=\frac{2 Y(X)}{X}
$$

In canonical form, the ODE is

$$
\begin{align*}
Y^{\prime} & =F(X, Y) \\
& =\frac{2 Y}{X} \tag{1}
\end{align*}
$$

An ode of the form $Y^{\prime}=\frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order $n$ if

$$
f\left(t^{n} X, t^{n} Y\right)=t^{n} f(X, Y)
$$

In this case, it can be seen that both $M=2 Y$ and $N=X$ are both homogeneous and of the same order $n=1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u=\frac{Y}{X}$, or $Y=u X$. Hence

$$
\frac{\mathrm{d} Y}{\mathrm{~d} X}=\frac{\mathrm{d} u}{\mathrm{~d} X} X+u
$$

Applying the transformation $Y=u X$ to the above ODE in (1) gives

$$
\begin{aligned}
\frac{\mathrm{d} u}{\mathrm{~d} X} X+u & =2 u \\
\frac{\mathrm{~d} u}{\mathrm{~d} X} & =\frac{u(X)}{X}
\end{aligned}
$$

Or

$$
\frac{d}{d X} u(X)-\frac{u(X)}{X}=0
$$

Or

$$
\left(\frac{d}{d X} u(X)\right) X-u(X)=0
$$

Which is now solved as separable in $u(X)$. Which is now solved in $u(X)$. In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(X, u) \\
& =f(X) g(u) \\
& =\frac{u}{X}
\end{aligned}
$$

Where $f(X)=\frac{1}{X}$ and $g(u)=u$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{u} d u & =\frac{1}{X} d X \\
\int \frac{1}{u} d u & =\int \frac{1}{X} d X \\
\ln (u) & =\ln (X)+c_{2} \\
u & =\mathrm{e}^{\ln (X)+c_{2}} \\
& =c_{2} X
\end{aligned}
$$

Now $u$ in the above solution is replaced back by $Y$ using $u=\frac{Y}{X}$ which results in the solution

$$
Y(X)=X^{2} c_{2}
$$

Using the solution for $Y(X)$

$$
Y(X)=X^{2} c_{2}
$$

And replacing back terms in the above solution using

$$
\begin{aligned}
& Y=y+y_{0} \\
& X=t+x_{0}
\end{aligned}
$$

Or

$$
\begin{aligned}
& Y=y \\
& X=t-1
\end{aligned}
$$

Then the solution in $y$ becomes

$$
y=(t+1)^{2} c_{2}
$$

Initial conditions are used to solve for $c_{2}$. Substituting $t=0$ and $y=6$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 6=c_{2} \\
& c_{2}=6
\end{aligned}
$$

Substituting $c_{2}$ found above in the general solution gives

$$
y=6(t+1)^{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=6(t+1)^{2} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

## Verification of solutions

$$
y=6(t+1)^{2}
$$

Verified OK.

### 7.5.6 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{2 y}{t+1} \\
y^{\prime} & =\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 60: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=(t+1)^{2} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{(t+1)^{2}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{(t+1)^{2}}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=\frac{2 y}{t+1}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =-\frac{2 y}{(t+1)^{3}} \\
S_{y} & =\frac{1}{(t+1)^{2}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\frac{y}{(t+1)^{2}}=c_{1}
$$

Which simplifies to

$$
\frac{y}{(t+1)^{2}}=c_{1}
$$

Which gives

$$
y=c_{1}(t+1)^{2}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=\frac{2 y}{t+1}$ |  | $\frac{d S}{d R}=0$ |
|  |  | $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ 边 |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow 40]{ }$ |
| -6. |  | $S(R)$ |
|  |  | S |
|  | $R=t$ | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow}$ |
| $\xrightarrow{\rightarrow \rightarrow \rightarrow-4 \times 1}$ | $y$ | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow}$ |
|  | $S=\frac{y}{(t+1)^{2}}$ |  |
|  | (t+1) | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-{ }_{-2}}$ |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{\text { a }}$ |
|  |  |  |
|  |  |  |

Initial conditions are used to solve for $c_{1}$. Substituting $t=0$ and $y=6$ in the above solution gives an equation to solve for the constant of integration.

$$
6=c_{1}
$$

$$
c_{1}=6
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=6(t+1)^{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=6(t+1)^{2} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

## Verification of solutions

$$
y=6(t+1)^{2}
$$

## Verified OK.

### 7.5.7 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{2 y}\right) \mathrm{d} y & =\left(\frac{1}{t+1}\right) \mathrm{d} t \\
\left(-\frac{1}{t+1}\right) \mathrm{d} t+\left(\frac{1}{2 y}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =-\frac{1}{t+1} \\
N(t, y) & =\frac{1}{2 y}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{t+1}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}\left(\frac{1}{2 y}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int M \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\frac{1}{t+1} \mathrm{~d} t \\
\phi & =-\ln (t+1)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{2 y}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{2 y}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{2 y}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{2 y}\right) \mathrm{d} y \\
f(y) & =\frac{\ln (y)}{2}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\ln (t+1)+\frac{\ln (y)}{2}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\ln (t+1)+\frac{\ln (y)}{2}
$$

The solution becomes

$$
y=\mathrm{e}^{2 c_{1}}(t+1)^{2}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=0$ and $y=6$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
6=\mathrm{e}^{2 c_{1}} \\
c_{1}=\frac{\ln (6)}{2}
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=6 t^{2}+12 t+6
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=6 t^{2}+12 t+6 \tag{1}
\end{equation*}
$$



## Verification of solutions

$$
y=6 t^{2}+12 t+6
$$

Verified OK.

### 7.5.8 Maple step by step solution

Let's solve
$\left[y^{\prime}-\frac{2 y}{t+1}=0, y(0)=6\right]$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- $\quad$ Separate variables
$\frac{y^{\prime}}{y}=\frac{2}{t+1}$
- Integrate both sides with respect to $t$
$\int \frac{y^{\prime}}{y} d t=\int \frac{2}{t+1} d t+c_{1}$
- Evaluate integral
$\ln (y)=2 \ln (t+1)+c_{1}$
- $\quad$ Solve for $y$
$y=\mathrm{e}^{c_{1}}(t+1)^{2}$
- Use initial condition $y(0)=6$
$6=e^{c_{1}}$
- $\quad$ Solve for $c_{1}$
$c_{1}=\ln (6)$
- $\quad$ Substitute $c_{1}=\ln (6)$ into general solution and simplify $y=6(t+1)^{2}$
- $\quad$ Solution to the IVP
$y=6(t+1)^{2}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 11

```
dsolve([diff(y(t),t)=2/(1+t)*y(t),y(0) = 6],y(t), singsol=all)
```

$$
y(t)=6(t+1)^{2}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.027 (sec). Leaf size: 12
DSolve [\{y' $[\mathrm{t}]==2 /(1+\mathrm{t}) * \mathrm{y}[\mathrm{t}], \mathrm{y}[0]==6\}, \mathrm{y}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow 6(t+1)^{2}
$$

## 7.6 problem 1.2-2 (f)

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Internal problem ID [2481]
Internal file name [OUTPUT/1973_Sunday_June_05_2022_02_41_37_AM_55133636/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-2, page 12
Problem number: 1.2-2 (f).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "differentialType", "first_order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime} t+y=t^{3}
$$

With initial conditions

$$
[y(1)=2]
$$

### 7.6.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
& p(t)=\frac{1}{t} \\
& q(t)=t^{2}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+\frac{y}{t}=t^{2}
$$

The domain of $p(t)=\frac{1}{t}$ is

$$
\{t<0 \vee 0<t\}
$$

And the point $t_{0}=1$ is inside this domain. The domain of $q(t)=t^{2}$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=1$ is also inside this domain. Hence solution exists and is unique.

### 7.6.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor $\mu$ is

$$
\begin{aligned}
& \mu=\mathrm{e}^{\int \frac{1}{t} d t} \\
& =t
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)\left(t^{2}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} t}(y t) & =(t)\left(t^{2}\right) \\
\mathrm{d}(y t) & =t^{3} \mathrm{~d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& y t=\int t^{3} \mathrm{~d} t \\
& y t=\frac{t^{4}}{4}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=t$ results in

$$
y=\frac{t^{3}}{4}+\frac{c_{1}}{t}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=1$ and $y=2$ in the above solution gives an equation to solve for the constant of integration.

$$
2=c_{1}+\frac{1}{4}
$$

$$
c_{1}=\frac{7}{4}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{t^{4}+7}{4 t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t^{4}+7}{4 t} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot


Verification of solutions

$$
y=\frac{t^{4}+7}{4 t}
$$

Verified OK.

### 7.6.3 Solving as differentialType ode

Writing the ode as

$$
\begin{equation*}
y^{\prime}=\frac{-y+t^{3}}{t} \tag{1}
\end{equation*}
$$

Which becomes

$$
\begin{equation*}
0=(-t) d y+\left(t^{3}-y\right) d t \tag{2}
\end{equation*}
$$

But the RHS is complete differential because

$$
(-t) d y+\left(t^{3}-y\right) d t=d\left(\frac{1}{4} t^{4}-y t\right)
$$

Hence (2) becomes

$$
0=d\left(\frac{1}{4} t^{4}-y t\right)
$$

Integrating both sides gives gives these solutions

$$
y=\frac{t^{4}+4 c_{1}}{4 t}+c_{1}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=1$ and $y=2$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
2=2 c_{1}+\frac{1}{4} \\
c_{1}=\frac{7}{8}
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{2 t^{4}+7 t+7}{8 t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{2 t^{4}+7 t+7}{8 t} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

## Verification of solutions

$$
y=\frac{2 t^{4}+7 t+7}{8 t}
$$

Warning, solution could not be verified

### 7.6.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{-t^{3}+y}{t} \\
y^{\prime} & =\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 63: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=\frac{1}{t} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{1}{t}} d y
\end{aligned}
$$

Which results in

$$
S=y t
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=-\frac{-t^{3}+y}{t}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =y \\
S_{y} & =t
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=t^{3} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=R^{3}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{R^{4}}{4}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
y t=\frac{t^{4}}{4}+c_{1}
$$

Which simplifies to

$$
y t=\frac{t^{4}}{4}+c_{1}
$$

Which gives

$$
y=\frac{t^{4}+4 c_{1}}{4 t}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=-\frac{-t^{3}+y}{t}$ |  | $\frac{d S}{d R}=R^{3}$ |
|  |  |  |
|  |  | $\cdots$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=t$ |  |
|  | $S=y t$ |  |
|  |  |  |
|  |  | $\rightarrow \rightarrow 14+$ |
|  |  |  |
|  |  |  |
| ¢ ¢ ¢ ¢ ¢ ¢ - ¢ blat ¢ ¢ ¢ ¢ ¢ ¢ ¢ ¢ |  |  |

Initial conditions are used to solve for $c_{1}$. Substituting $t=1$ and $y=2$ in the above solution gives an equation to solve for the constant of integration.

$$
2=c_{1}+\frac{1}{4}
$$

$$
c_{1}=\frac{7}{4}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{t^{4}+7}{4 t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t^{4}+7}{4 t} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot


Verification of solutions

$$
y=\frac{t^{4}+7}{4 t}
$$

Verified OK.

### 7.6.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(t) \mathrm{d} y & =\left(t^{3}-y\right) \mathrm{d} t \\
\left(-t^{3}+y\right) \mathrm{d} t+(t) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =-t^{3}+y \\
N(t, y) & =t
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-t^{3}+y\right) \\
& =1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(t) \\
& =1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int M \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-t^{3}+y \mathrm{~d} t \\
\phi & =-\frac{1}{4} t^{4}+y t+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=t+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=t$. Therefore equation (4) becomes

$$
\begin{equation*}
t=t+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\frac{1}{4} t^{4}+y t+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\frac{1}{4} t^{4}+y t
$$

The solution becomes

$$
y=\frac{t^{4}+4 c_{1}}{4 t}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=1$ and $y=2$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
2=c_{1}+\frac{1}{4} \\
c_{1}=\frac{7}{4}
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{t^{4}+7}{4 t}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t^{4}+7}{4 t} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

## Verification of solutions

$$
y=\frac{t^{4}+7}{4 t}
$$

Verified OK.

### 7.6.6 Maple step by step solution

Let's solve

$$
\left[y^{\prime} t+y=t^{3}, y(1)=2\right]
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=-\frac{y}{t}+t^{2}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE
$y^{\prime}+\frac{y}{t}=t^{2}$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}+\frac{y}{t}\right)=\mu(t) t^{2}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}+\frac{y}{t}\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- $\quad$ Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=\frac{\mu(t)}{t}$
- Solve to find the integrating factor
$\mu(t)=t$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \mu(t) t^{2} d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \mu(t) t^{2} d t+c_{1}$
- Solve for $y$
$y=\frac{\int \mu(t) t^{2} d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=t$
$y=\frac{\int t^{3} d t+c_{1}}{t}$
- Evaluate the integrals on the rhs
$y=\frac{\frac{t^{4}}{4}+c_{1}}{t}$
- Simplify
$y=\frac{t^{4}+4 c_{1}}{4 t}$
- Use initial condition $y(1)=2$
$2=c_{1}+\frac{1}{4}$
- $\quad$ Solve for $c_{1}$
$c_{1}=\frac{7}{4}$
- Substitute $c_{1}=\frac{7}{4}$ into general solution and simplify
$y=\frac{t^{4}+7}{4 t}$
- Solution to the IVP
$y=\frac{t^{4}+7}{4 t}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 14

```
dsolve([t*diff(y(t),t)=-y(t)+t^3,y(1) = 2],y(t), singsol=all)
```

$$
y(t)=\frac{t^{4}+7}{4 t}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.048 (sec). Leaf size: 27
DSolve[\{y' $[t]==-y[t]+t \backsim 3, y[1]==2\}, y[t], t$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(t) \rightarrow t^{3}-3 t^{2}+6 t+4 e^{1-t}-6
$$

## 8 Problem 1.2-3, page 12

8.1 problem 1.2-3 (a) ..... 289
8.2 problem 1.2-3 (b) ..... 303
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## 8.1 problem 1.2-3 (a)

8.1.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 289
8.1.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 290
8.1.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 292
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8.1.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 301

Internal problem ID [2482]
Internal file name [OUTPUT/1974_Sunday_June_05_2022_02_41_41_AM_94026061/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-3, page 12
Problem number: 1.2-3 (a).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
y^{\prime}+4 \tan (2 t) y=\tan (2 t)
$$

With initial conditions

$$
\left[y\left(\frac{\pi}{8}\right)=2\right]
$$

### 8.1.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =4 \tan (2 t) \\
q(t) & =\tan (2 t)
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+4 \tan (2 t) y=\tan (2 t)
$$

The domain of $p(t)=4 \tan (2 t)$ is

$$
\left\{t<\frac{1}{4} \pi+\frac{1}{2} \pi \_Z 116 \vee \frac{1}{4} \pi+\frac{1}{2} \pi \_Z 116<t\right\}
$$

And the point $t_{0}=\frac{\pi}{8}$ is inside this domain. The domain of $q(t)=\tan (2 t)$ is

$$
\left\{t<\frac{1}{4} \pi+\frac{1}{2} \pi \_Z 116 \vee \frac{1}{4} \pi+\frac{1}{2} \pi \_Z 116<t\right\}
$$

And the point $t_{0}=\frac{\pi}{8}$ is also inside this domain. Hence solution exists and is unique.

### 8.1.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int 4 \tan (2 t) d t} \\
& =\frac{1}{\cos (2 t)^{2}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)(\tan (2 t)) \\
\frac{\mathrm{d}}{\mathrm{~d} t}\left(\frac{y}{\cos (2 t)^{2}}\right) & =\left(\frac{1}{\cos (2 t)^{2}}\right)(\tan (2 t)) \\
\mathrm{d}\left(\frac{y}{\cos (2 t)^{2}}\right) & =\left(\tan (2 t) \sec (2 t)^{2}\right) \mathrm{d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{y}{\cos (2 t)^{2}}=\int \tan (2 t) \sec (2 t)^{2} \mathrm{~d} t \\
& \frac{y}{\cos (2 t)^{2}}=\frac{\sec (2 t)^{2}}{4}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{\cos (2 t)^{2}}$ results in

$$
y=\frac{\cos (2 t)^{2} \sec (2 t)^{2}}{4}+c_{1} \cos (2 t)^{2}
$$

which simplifies to

$$
y=c_{1} \cos (2 t)^{2}+\frac{1}{4}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\frac{\pi}{8}$ and $y=2$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
2=\frac{c_{1}}{2}+\frac{1}{4} \\
c_{1}=\frac{7}{2}
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{1}{4}+\frac{7 \cos (2 t)^{2}}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{1}{4}+\frac{7 \cos (2 t)^{2}}{2} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

Verification of solutions

$$
y=\frac{1}{4}+\frac{7 \cos (2 t)^{2}}{2}
$$

Verified OK.

### 8.1.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=-4 \tan (2 t) y+\tan (2 t) \\
& y^{\prime}=\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 66: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=\frac{1}{1+\tan (2 t)^{2}} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{1}{1+\tan (2 t)^{2}}} d y
\end{aligned}
$$

Which results in

$$
S=\left(1+\tan (2 t)^{2}\right) y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=-4 \tan (2 t) y+\tan (2 t)
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =4 \sec (2 t)^{2} y \tan (2 t) \\
S_{y} & =\sec (2 t)^{2}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\tan (2 t) \sec (2 t)^{2} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\tan (2 R) \sec (2 R)^{2}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$ ．Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{\sec (2 R)^{2}}{4}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution，we just need to transform（4）back to $t, y$ coordinates．This results in

$$
\sec (2 t)^{2} y=\frac{\sec (2 t)^{2}}{4}+c_{1}
$$

Which simplifies to

$$
\sec (2 t)^{2} y=\frac{\sec (2 t)^{2}}{4}+c_{1}
$$

Which gives

$$
y=\frac{\sec (2 t)^{2}+4 c_{1}}{4 \sec (2 t)^{2}}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown．

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=-4 \tan (2 t) y+\tan (2 t)$ |  | $\frac{d S}{d R}=\tan (2 R) \sec (2 R)^{2}$ |
|  |  | －ッさ入け入入 |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=t$ | $x \rightarrow+x+y$ |
|  | $S=\sec (2 t)^{2} y$ |  |
|  | $S=\sec (2 t)^{2} y$ |  |
|  |  |  |
| $1+1$ |  | －可吅 |
|  |  |  |
|  |  |  |

Initial conditions are used to solve for $c_{1}$ ．Substituting $t=\frac{\pi}{8}$ and $y=2$ in the above solution gives an equation to solve for the constant of integration．

$$
2=\frac{c_{1}}{2}+\frac{1}{4}
$$

$$
c_{1}=\frac{7}{2}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{1}{4}+\frac{7 \cos (2 t)^{2}}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{1}{4}+\frac{7 \cos (2 t)^{2}}{2} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot

## Verification of solutions

$$
y=\frac{1}{4}+\frac{7 \cos (2 t)^{2}}{2}
$$

Verified OK.

### 8.1.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =(-4 \tan (2 t) y+\tan (2 t)) \mathrm{d} t \\
(4 \tan (2 t) y-\tan (2 t)) \mathrm{d} t+\mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =4 \tan (2 t) y-\tan (2 t) \\
N(t, y) & =1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(4 \tan (2 t) y-\tan (2 t)) \\
& =4 \tan (2 t)
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =1((4 \tan (2 t))-(0)) \\
& =4 \tan (2 t)
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int 4 \tan (2 t) \mathrm{d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{\ln \left(1+\tan (2 t)^{2}\right)} \\
& =\sec (2 t)^{2}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\sec (2 t)^{2}(4 \tan (2 t) y-\tan (2 t)) \\
& =\tan (2 t)(4 y-1) \sec (2 t)^{2}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\sec (2 t)^{2}(1) \\
& =\sec (2 t)^{2}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t}=0 \\
\left(\tan (2 t)(4 y-1) \sec (2 t)^{2}\right)+\left(\sec (2 t)^{2}\right) \frac{\mathrm{d} y}{\mathrm{~d} t}
\end{array}=0
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \tan (2 t)(4 y-1) \sec (2 t)^{2} \mathrm{~d} t \\
\phi & =\frac{\sec (2 t)^{2}(4 y-1)}{4}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\sec (2 t)^{2}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\sec (2 t)^{2}$. Therefore equation (4) becomes

$$
\begin{equation*}
\sec (2 t)^{2}=\sec (2 t)^{2}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{\sec (2 t)^{2}(4 y-1)}{4}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{\sec (2 t)^{2}(4 y-1)}{4}
$$

The solution becomes

$$
y=\frac{\sec (2 t)^{2}+4 c_{1}}{4 \sec (2 t)^{2}}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\frac{\pi}{8}$ and $y=2$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
2=\frac{c_{1}}{2}+\frac{1}{4} \\
c_{1}=\frac{7}{2}
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{1}{4}+\frac{7 \cos (2 t)^{2}}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{1}{4}+\frac{7 \cos (2 t)^{2}}{2} \tag{1}
\end{equation*}
$$


(a) Solution plot

(b) Slope field plot

Verification of solutions

$$
y=\frac{1}{4}+\frac{7 \cos (2 t)^{2}}{2}
$$

Verified OK.

### 8.1.5 Maple step by step solution

Let's solve

$$
\left[y^{\prime}+4 \tan (2 t) y=\tan (2 t), y\left(\frac{\pi}{8}\right)=2\right]
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- $\quad$ Separate variables
$\frac{y^{\prime}}{4 y-1}=-\tan (2 t)$
- Integrate both sides with respect to $t$
$\int \frac{y^{\prime}}{4 y-1} d t=\int-\tan (2 t) d t+c_{1}$
- Evaluate integral

$$
\frac{\ln (4 y-1)}{4}=-\frac{\ln \left(1+\tan (2 t)^{2}\right)}{4}+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\frac{\mathrm{e}^{4 c_{1}} \cos (2 t)^{2}}{4}+\frac{1}{4}
$$

- Use initial condition $y\left(\frac{\pi}{8}\right)=2$

$$
2=\frac{\mathrm{e}^{4 c_{1}}}{8}+\frac{1}{4}
$$

- $\quad$ Solve for $c_{1}$
$c_{1}=\frac{\ln (14)}{4}$
- Substitute $c_{1}=\frac{\ln (14)}{4}$ into general solution and simplify

$$
y=2+\frac{7 \cos (4 t)}{4}
$$

- Solution to the IVP
$y=2+\frac{7 \cos (4 t)}{4}$

Maple trace
-Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

Solution by Maple
Time used: 0.046 (sec). Leaf size: 12

```
dsolve([diff(y(t),t)+4*\operatorname{tan}(2*t)*y(t)=tan(2*t),y(1/8*Pi) = 2],y(t), singsol=all)
```

$$
y(t)=2+\frac{7 \cos (4 t)}{4}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.098 (sec). Leaf size: 15
DSolve $\left[\left\{y^{\prime}[t]+4 * \operatorname{Tan}[2 * t] * y[t]==\operatorname{Tan}[2 * t], y[\mathrm{Pi} / 8]==2\right\}, y[t], t\right.$, IncludeSingularSolutions $->$ True $]$

$$
y(t) \rightarrow \frac{7}{4} \cos (4 t)+2
$$

## 8.2 problem 1.2-3 (b)

8.2.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 303
8.2.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 304
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Internal problem ID [2483]
Internal file name [OUTPUT/1975_Sunday_June_05_2022_02_41_44_AM_15660872/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-3, page 12
Problem number: 1.2-3 (b).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
t \ln (t) y^{\prime}+y=t \ln (t)
$$

With initial conditions

$$
[y(\mathrm{e})=1]
$$

### 8.2.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
& p(t)=\frac{1}{t \ln (t)} \\
& q(t)=1
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+\frac{y}{t \ln (t)}=1
$$

The domain of $p(t)=\frac{1}{t \ln (t)}$ is

$$
\{0<t<1,1<t \leq \infty\}
$$

And the point $t_{0}=\mathrm{e}$ is inside this domain. The domain of $q(t)=1$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=\mathrm{e}$ is also inside this domain. Hence solution exists and is unique.

### 8.2.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor $\mu$ is

$$
\begin{gathered}
\mu=\mathrm{e}^{\int \frac{1}{t \ln (t)} d t} \\
=\ln (t)
\end{gathered}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =\mu \\
\frac{\mathrm{d}}{\mathrm{~d} t}(\ln (t) y) & =\ln (t) \\
\mathrm{d}(\ln (t) y) & =\ln (t) \mathrm{d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \ln (t) y=\int \ln (t) \mathrm{d} t \\
& \ln (t) y=t \ln (t)-t+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\ln (t)$ results in

$$
y=\frac{t \ln (t)-t}{\ln (t)}+\frac{c_{1}}{\ln (t)}
$$

which simplifies to

$$
y=\frac{t \ln (t)+c_{1}-t}{\ln (t)}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\mathrm{e}$ and $y=1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 1=c_{1} \\
& c_{1}=1
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{t \ln (t)+1-t}{\ln (t)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t \ln (t)+1-t}{\ln (t)} \tag{1}
\end{equation*}
$$



(a) Solution plot
(b) Slope field plot

## Verification of solutions

$$
y=\frac{t \ln (t)+1-t}{\ln (t)}
$$

Verified OK.

### 8.2.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{t \ln (t)-y}{t \ln (t)} \\
y^{\prime} & =\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 69: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $y$ |  |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | special | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |  |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $\underline{a}_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}$ |
| $a_{1} b_{2}-a_{2} b_{1}$ |  |  |  |$\underline{\underline{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}} \frac{a_{1} b_{2}-a_{2} b_{1}}{}}$| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 |
| :--- | :--- | :--- |
| $-\int(n-1) f(x) d x y^{n}$ |  |  |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=\frac{1}{\ln (t)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{1}{\ln (t)}} d y
\end{aligned}
$$

Which results in

$$
S=\ln (t) y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=\frac{t \ln (t)-y}{t \ln (t)}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =\frac{y}{t} \\
S_{y} & =\ln (t)
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\ln (t) \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\ln (R)
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=R \ln (R)-R+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\ln (t) y=t \ln (t)+c_{1}-t
$$

Which simplifies to

$$
\ln (t) y=t \ln (t)+c_{1}-t
$$

Which gives

$$
y=\frac{t \ln (t)+c_{1}-t}{\ln (t)}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
|  | $\begin{aligned} R & =t \\ S & =\ln (t) y \end{aligned}$ |  |

Initial conditions are used to solve for $c_{1}$. Substituting $t=\mathrm{e}$ and $y=1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 1=c_{1} \\
& c_{1}=1
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{t \ln (t)+1-t}{\ln (t)}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t \ln (t)+1-t}{\ln (t)} \tag{1}
\end{equation*}
$$


(b) Slope field plot


## Verification of solutions

$$
y=\frac{t \ln (t)+1-t}{\ln (t)}
$$

Verified OK.

### 8.2.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(t \ln (t)) \mathrm{d} y & =(t \ln (t)-y) \mathrm{d} t \\
(-t \ln (t)+y) \mathrm{d} t+(t \ln (t)) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =-t \ln (t)+y \\
N(t, y) & =t \ln (t)
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(-t \ln (t)+y) \\
& =1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(t \ln (t)) \\
& =\ln (t)+1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =\frac{1}{t \ln (t)}((1)-(\ln (t)+1)) \\
& =-\frac{1}{t}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int-\frac{1}{t} \mathrm{~d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-\ln (t)} \\
& =\frac{1}{t}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{t}(-t \ln (t)+y) \\
& =\frac{-t \ln (t)+y}{t}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{t}(t \ln (t)) \\
& =\ln (t)
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
\left(\frac{-t \ln (t)+y}{t}\right)+(\ln (t)) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \frac{-t \ln (t)+y}{t} \mathrm{~d} t \\
\phi & =(y-t) \ln (t)+t+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\ln (t)+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\ln (t)$. Therefore equation (4) becomes

$$
\begin{equation*}
\ln (t)=\ln (t)+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=(y-t) \ln (t)+t+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=(y-t) \ln (t)+t
$$

The solution becomes

$$
y=\frac{t \ln (t)+c_{1}-t}{\ln (t)}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\mathrm{e}$ and $y=1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 1=c_{1} \\
& c_{1}=1
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=\frac{t \ln (t)+1-t}{\ln (t)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{t \ln (t)+1-t}{\ln (t)} \tag{1}
\end{equation*}
$$



(a) Solution plot
(b) Slope field plot

## Verification of solutions

$$
y=\frac{t \ln (t)+1-t}{\ln (t)}
$$

Verified OK.

### 8.2.5 Maple step by step solution

Let's solve
$\left[t \ln (t) y^{\prime}+y=t \ln (t), y(\mathrm{e})=1\right]$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=1-\frac{y}{t \ln (t)}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE
$y^{\prime}+\frac{y}{t \ln (t)}=1$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}+\frac{y}{t \ln (t)}\right)=\mu(t)$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}+\frac{y}{t \ln (t)}\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=\frac{\mu(t)}{t \ln (t)}$
- Solve to find the integrating factor
$\mu(t)=\ln (t)$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int \mu(t) d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int \mu(t) d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \mu(t) d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\ln (t)$
$y=\frac{\int \ln (t) d t+c_{1}}{\ln (t)}$
- Evaluate the integrals on the rhs
$y=\frac{t \ln (t)+c_{1}-t}{\ln (t)}$
- Use initial condition $y(\mathrm{e})=1$
$1=c_{1}$
- $\quad$ Solve for $c_{1}$
$c_{1}=1$
- $\quad$ Substitute $c_{1}=1$ into general solution and simplify
$y=\frac{t \ln (t)+1-t}{\ln (t)}$
- $\quad$ Solution to the IVP
$y=\frac{t \ln (t)+1-t}{\ln (t)}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

Solution by Maple
Time used: 0.016 (sec). Leaf size: 18

```
dsolve([t* ln (t)*diff(y(t),t)=t*\operatorname{ln}(t)-y(t),y(exp(1)) = 1],y(t), singsol=all)
```

$$
y(t)=\frac{t \ln (t)-t+1}{\ln (t)}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.036 (sec). Leaf size: 19
DSolve $[\{t * \log [t] * y$ ' $[t]==t * \log [t]-y[t], y[\operatorname{Exp}[1]]==1\}, y[t], t$, IncludeSingularSolutions $->$ True $]$

$$
y(t) \rightarrow \frac{-t+t \log (t)+1}{\log (t)}
$$

## 8.3 problem 1.2-3 (c)

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8.3.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 318
8.3.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 320
8.3.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 325
8.3.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 330

Internal problem ID [2484]
Internal file name [OUTPUT/1976_Sunday_June_05_2022_02_41_48_AM_14092017/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-3, page 12
Problem number: 1.2-3 (c).
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}-\frac{2 y}{-t^{2}+1}=3
$$

With initial conditions

$$
\left[y\left(\frac{1}{2}\right)=1\right]
$$

### 8.3.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
& p(t)=\frac{2}{t^{2}-1} \\
& q(t)=3
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+\frac{2 y}{t^{2}-1}=3
$$

The domain of $p(t)=\frac{2}{t^{2}-1}$ is

$$
\{-\infty \leq t<-1,-1<t<1,1<t \leq \infty\}
$$

And the point $t_{0}=\frac{1}{2}$ is inside this domain. The domain of $q(t)=3$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=\frac{1}{2}$ is also inside this domain. Hence solution exists and is unique.

### 8.3.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \frac{2}{t^{2}-1} d t} \\
& =\frac{-t^{2}+1}{(t+1)^{2}}
\end{aligned}
$$

Which simplifies to

$$
\mu=\frac{-t+1}{t+1}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)(3) \\
\frac{\mathrm{d}}{\mathrm{~d} t}\left(\frac{(-t+1) y}{t+1}\right) & =\left(\frac{-t+1}{t+1}\right)(3) \\
\mathrm{d}\left(\frac{(-t+1) y}{t+1}\right) & =\left(\frac{-3 t+3}{t+1}\right) \mathrm{d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{(-t+1) y}{t+1}=\int \frac{-3 t+3}{t+1} \mathrm{~d} t \\
& \frac{(-t+1) y}{t+1}=-3 t+6 \ln (t+1)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{-t+1}{t+1}$ results in

$$
y=\frac{(-t-1)(-3 t+6 \ln (t+1))}{t-1}+\frac{c_{1}(-t-1)}{t-1}
$$

which simplifies to

$$
y=\frac{(t+1)\left(3 t-6 \ln (t+1)-c_{1}\right)}{t-1}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\frac{1}{2}$ and $y=1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
1=-\frac{9}{2}+18 \ln (3)-18 \ln (2)+3 c_{1} \\
c_{1}=\frac{11}{6}-6 \ln (3)+6 \ln (2)
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives
$y=\frac{-36 \ln (t+1) t+36 \ln (3) t-36 \ln (2) t+18 t^{2}-36 \ln (t+1)+36 \ln (3)-36 \ln (2)+7 t-11}{6 t-6}$

## Summary

The solution(s) found are the following

## $y$

$=\frac{-36 \ln (t+1) t+36 \ln (3) t-36 \ln (2) t+18 t^{2}-36 \ln (t+1)+36 \ln (3)-36 \ln (2)+7 t-11}{6 t-6}$



## Verification of solutions

$y$
$=\frac{-36 \ln (t+1) t+36 \ln (3) t-36 \ln (2) t+18 t^{2}-36 \ln (t+1)+36 \ln (3)-36 \ln (2)+7 t-11}{6 t-6}$
Verified OK.

### 8.3.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=-\frac{-3 t^{2}+2 y+3}{t^{2}-1} \\
& y^{\prime}=\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 72: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, y)=0 \\
& \eta(t, y)=\frac{(t+1)^{2}}{-t^{2}+1} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{(t+1)^{2}}{-t^{2}+1}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{\left(-t^{2}+1\right) y}{(t+1)^{2}}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=-\frac{-3 t^{2}+2 y+3}{t^{2}-1}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =-\frac{2 y}{(t+1)^{2}} \\
S_{y} & =\frac{-t+1}{t+1}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{-3 t+3}{t+1} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{-3 R+3}{R+1}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-3 R+6 \ln (R+1)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
-\frac{(t-1) y}{t+1}=-3 t+6 \ln (t+1)+c_{1}
$$

Which simplifies to

$$
-\frac{(t-1) y}{t+1}=-3 t+6 \ln (t+1)+c_{1}
$$

Which gives

$$
y=-\frac{(t+1)\left(-3 t+6 \ln (t+1)+c_{1}\right)}{t-1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=-\frac{-3 t^{2}+2 y+3}{t^{2}-1}$ |  | $\frac{d S}{d R}=\frac{-3 R+3}{R+1}$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=t$ |  |
|  | $(t-1) y$ |  |
|  | $S=-\frac{(t-1) y}{t+1}$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |

Initial conditions are used to solve for $c_{1}$. Substituting $t=\frac{1}{2}$ and $y=1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
1=-\frac{9}{2}+18 \ln (3)-18 \ln (2)+3 c_{1} \\
c_{1}=\frac{11}{6}-6 \ln (3)+6 \ln (2)
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives
$y=\frac{-36 \ln (t+1) t+36 \ln (3) t-36 \ln (2) t+18 t^{2}-36 \ln (t+1)+36 \ln (3)-36 \ln (2)+7 t-11}{6 t-6}$

## Summary

The solution(s) found are the following
$y$
$=\frac{-36 \ln (t+1) t+36 \ln (3) t-36 \ln (2) t+18 t^{2}-36 \ln (t+1)+36 \ln (3)-36 \ln (2)+7 t-11}{6 t-6}$

(b) Slope field plot

## Verification of solutions

$y$
$=\frac{-36 \ln (t+1) t+36 \ln (3) t-36 \ln (2) t+18 t^{2}-36 \ln (t+1)+36 \ln (3)-36 \ln (2)+7 t-11}{6 t-6}$
Verified OK.

### 8.3.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(\frac{2 y}{-t^{2}+1}+3\right) \mathrm{d} t \\
\left(-\frac{2 y}{-t^{2}+1}-3\right) \mathrm{d} t+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(t, y)=-\frac{2 y}{-t^{2}+1}-3 \\
& N(t, y)=1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{2 y}{-t^{2}+1}-3\right) \\
& =\frac{2}{t^{2}-1}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =1\left(\left(-\frac{2}{-t^{2}+1}\right)-(0)\right) \\
& =\frac{2}{t^{2}-1}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int \frac{2}{t^{2}-1} \mathrm{~d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-2 \operatorname{arctanh}(t)} \\
& =\frac{-t+1}{t+1}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{-t+1}{t+1}\left(-\frac{2 y}{-t^{2}+1}-3\right) \\
& =\frac{3 t^{2}-2 y-3}{(t+1)^{2}}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{-t+1}{t+1}(1) \\
& =\frac{-t+1}{t+1}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
\left(\frac{3 t^{2}-2 y-3}{(t+1)^{2}}\right)+\left(\frac{-t+1}{t+1}\right) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \frac{3 t^{2}-2 y-3}{(t+1)^{2}} \mathrm{~d} t \\
\phi & =3 t-6 \ln (t+1)+\frac{2 y}{t+1}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\frac{2}{t+1}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{-t+1}{t+1}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{-t+1}{t+1}=\frac{2}{t+1}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=-1
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int(-1) \mathrm{d} y \\
f(y) & =-y+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=3 t-6 \ln (t+1)+\frac{2 y}{t+1}-y+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=3 t-6 \ln (t+1)+\frac{2 y}{t+1}-y
$$

The solution becomes

$$
y=-\frac{(t+1)\left(-3 t+6 \ln (t+1)+c_{1}\right)}{t-1}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\frac{1}{2}$ and $y=1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
1=-\frac{9}{2}+18 \ln (3)-18 \ln (2)+3 c_{1} \\
c_{1}=\frac{11}{6}-6 \ln (3)+6 \ln (2)
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives
$y=\frac{-36 \ln (t+1) t+36 \ln (3) t-36 \ln (2) t+18 t^{2}-36 \ln (t+1)+36 \ln (3)-36 \ln (2)+7 t-11}{6 t-6}$

## Summary

The solution(s) found are the following
$y$
$=\frac{-36 \ln (t+1) t+36 \ln (3) t-36 \ln (2) t+18 t^{2}-36 \ln (t+1)+36 \ln (3)-36 \ln (2)+7 t-11}{6 t-6}$


(a) Solution plot
(b) Slope field plot

## Verification of solutions

$y$
$=\frac{-36 \ln (t+1) t+36 \ln (3) t-36 \ln (2) t+18 t^{2}-36 \ln (t+1)+36 \ln (3)-36 \ln (2)+7 t-11}{6 t-6}$
Verified OK.

### 8.3.5 Maple step by step solution

Let's solve

$$
\left[y^{\prime}-\frac{2 y}{-t^{2}+1}=3, y\left(\frac{1}{2}\right)=1\right]
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=3-\frac{2 y}{t^{2}-1}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE
$y^{\prime}+\frac{2 y}{t^{2}-1}=3$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}+\frac{2 y}{t^{2}-1}\right)=3 \mu(t)$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}+\frac{2 y}{t^{2}-1}\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=\frac{2 \mu(t)}{t^{2}-1}$
- Solve to find the integrating factor
$\mu(t)=\frac{t-1}{t+1}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int 3 \mu(t) d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int 3 \mu(t) d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int 3 \mu(t) d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\frac{t-1}{t+1}$
$y=\frac{(t+1)\left(\int \frac{3(t-1)}{t+1} d t+c_{1}\right)}{t-1}$
- Evaluate the integrals on the rhs
$y=\frac{(t+1)\left(3 t-6 \ln (t+1)+c_{1}\right)}{t-1}$
- Use initial condition $y\left(\frac{1}{2}\right)=1$

$$
1=-\frac{9}{2}+18 \ln \left(\frac{3}{2}\right)-3 c_{1}
$$

- $\quad$ Solve for $c_{1}$

$$
c_{1}=-\frac{11}{6}+6 \ln \left(\frac{3}{2}\right)
$$

- Substitute $c_{1}=-\frac{11}{6}+6 \ln \left(\frac{3}{2}\right)$ into general solution and simplify $y=\frac{(18 t-36 \ln (t+1)-11+36 \ln (3)-36 \ln (2))(t+1)}{6 t-6}$
- Solution to the IVP
$y=\frac{(18 t-36 \ln (t+1)-11+36 \ln (3)-36 \ln (2))(t+1)}{6 t-6}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.032 (sec). Leaf size: 34

```
dsolve([diff(y(t),t)=2/(1-t^2)*y(t)+3,y(1/2) = 1],y(t), singsol=all)
```

$$
y(t)=\frac{(t+1)(18 t-36 \ln (t+1)-11+36 \ln (3)-36 \ln (2))}{6 t-6}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.041 (sec). Leaf size: 34
DSolve[\{y' $\left.[t]==2 /\left(1-t^{\wedge} 2\right) * y[t]+3, y[1 / 2]==1\right\}, y[t], t$, IncludeSingularSolutions $->$ True]

$$
y(t) \rightarrow \frac{(t+1)\left(18 t-36 \log (t+1)-11+36 \log \left(\frac{3}{2}\right)\right)}{6(t-1)}
$$

## 8.4 problem 1.2-3 (d)

8.4.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 332
8.4.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 333
8.4.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 335
8.4.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 339
8.4.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 343

Internal problem ID [2485]
Internal file name [OUTPUT/1977_Sunday_June_05_2022_02_41_52_AM_30503089/index.tex]
Book: Ordinary Differential Equations, Robert H. Martin, 1983
Section: Problem 1.2-3, page 12
Problem number: 1.2-3 (d).
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}+\cot (t) y=6 \cos (t)^{2}
$$

With initial conditions

$$
\left[y\left(\frac{\pi}{4}\right)=3\right]
$$

### 8.4.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime}+p(t) y=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =\cot (t) \\
q(t) & =6 \cos (t)^{2}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+\cot (t) y=6 \cos (t)^{2}
$$

The domain of $p(t)=\cot (t)$ is

$$
\left\{t<\pi \_Z 133 \vee \pi \_Z 133<t\right\}
$$

And the point $t_{0}=\frac{\pi}{4}$ is inside this domain. The domain of $q(t)=6 \cos (t)^{2}$ is

$$
\{-\infty<t<\infty\}
$$

And the point $t_{0}=\frac{\pi}{4}$ is also inside this domain. Hence solution exists and is unique.

### 8.4.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \cot (t) d t} \\
& =\sin (t)
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu y) & =(\mu)\left(6 \cos (t)^{2}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} t}(\sin (t) y) & =(\sin (t))\left(6 \cos (t)^{2}\right) \\
\mathrm{d}(\sin (t) y) & =\left(6 \sin (t) \cos (t)^{2}\right) \mathrm{d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \sin (t) y=\int 6 \sin (t) \cos (t)^{2} \mathrm{~d} t \\
& \sin (t) y=-2 \cos (t)^{3}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\sin (t)$ results in

$$
y=-2 \csc (t) \cos (t)^{3}+c_{1} \csc (t)
$$

which simplifies to

$$
y=\csc (t)\left(-2 \cos (t)^{3}+c_{1}\right)
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\frac{\pi}{4}$ and $y=3$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
3=-1+\sqrt{2} c_{1} \\
c_{1}=2 \sqrt{2}
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=-2 \csc (t) \cos (t)^{3}+2 \csc (t) \sqrt{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-2 \csc (t) \cos (t)^{3}+2 \csc (t) \sqrt{2} \tag{1}
\end{equation*}
$$


(a) Solution plot (b) Slope field plot

Verification of solutions

$$
y=-2 \csc (t) \cos (t)^{3}+2 \csc (t) \sqrt{2}
$$

Verified OK.

### 8.4.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=-\cot (t) y+6 \cos (t)^{2} \\
& y^{\prime}=\omega(t, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{y}-\xi_{t}\right)-\omega^{2} \xi_{y}-\omega_{t} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 75: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | special | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |  |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}$ |
| $a_{1} b_{2}-a_{2} b_{1}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |  |  |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |

The above table shows that

$$
\begin{align*}
\xi(t, y) & =0 \\
\eta(t, y) & =\frac{1}{\sin (t)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(t, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial y}\right) S(t, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{1}{\sin (t)}} d y
\end{aligned}
$$

Which results in

$$
S=\sin (t) y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, y) S_{y}}{R_{t}+\omega(t, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{y}, S_{t}, S_{y}$ are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$
\omega(t, y)=-\cot (t) y+6 \cos (t)^{2}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{y} & =0 \\
S_{t} & =\cos (t) y \\
S_{y} & =\sin (t)
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=6 \sin (t) \cos (t)^{2} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=6 \sin (R) \cos (R)^{2}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-2 \cos (R)^{3}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, y$ coordinates. This results in

$$
\sin (t) y=-2 \cos (t)^{3}+c_{1}
$$

Which simplifies to

$$
\sin (t) y=-2 \cos (t)^{3}+c_{1}
$$

Which gives

$$
y=-\frac{2 \cos (t)^{3}-c_{1}}{\sin (t)}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d t}=-\cot (t) y+6 \cos (t)^{2}$ |  | $\frac{d S}{d R}=6 \sin (R) \cos (R)^{2}$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=t$ |  |
|  |  |  |
|  | $S=\sin (t) y$ | $\rightarrow \rightarrow 40 \times \pm$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |

Initial conditions are used to solve for $c_{1}$. Substituting $t=\frac{\pi}{4}$ and $y=3$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
3=-1+\sqrt{2} c_{1} \\
c_{1}=2 \sqrt{2}
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=-2 \csc (t) \cos (t)^{3}+2 \csc (t) \sqrt{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-2 \csc (t) \cos (t)^{3}+2 \csc (t) \sqrt{2} \tag{1}
\end{equation*}
$$



## Verification of solutions

$$
y=-2 \csc (t) \cos (t)^{3}+2 \csc (t) \sqrt{2}
$$

Verified OK.

### 8.4.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, y) \mathrm{d} t+N(t, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(-\cot (t) y+6 \cos (t)^{2}\right) \mathrm{d} t \\
\left(\cot (t) y-6 \cos (t)^{2}\right) \mathrm{d} t+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, y) & =\cot (t) y-6 \cos (t)^{2} \\
N(t, y) & =1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\cot (t) y-6 \cos (t)^{2}\right) \\
& =\cot (t)
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial t}\right) \\
& =1((\cot (t))-(0)) \\
& =\cot (t)
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int \cot (t) \mathrm{d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{\ln (\sin (t))} \\
& =\sin (t)
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\sin (t)\left(\cot (t) y-6 \cos (t)^{2}\right) \\
& =\cos (t)(-6 \sin (t) \cos (t)+y)
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\sin (t)(1) \\
& =\sin (t)
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} t} & =0 \\
(\cos (t)(-6 \sin (t) \cos (t)+y))+(\sin (t)) \frac{\mathrm{d} y}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \cos (t)(-6 \sin (t) \cos (t)+y) \mathrm{d} t \\
\phi & =\sin (t) y+2 \cos (t)^{3}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\sin (t)+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\sin (t)$. Therefore equation (4) becomes

$$
\begin{equation*}
\sin (t)=\sin (t)+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\sin (t) y+2 \cos (t)^{3}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\sin (t) y+2 \cos (t)^{3}
$$

The solution becomes

$$
y=-\frac{2 \cos (t)^{3}-c_{1}}{\sin (t)}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $t=\frac{\pi}{4}$ and $y=3$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
3=-1+\sqrt{2} c_{1} \\
c_{1}=2 \sqrt{2}
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
y=-2 \csc (t) \cos (t)^{3}+2 \csc (t) \sqrt{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-2 \csc (t) \cos (t)^{3}+2 \csc (t) \sqrt{2} \tag{1}
\end{equation*}
$$


(a) Solution plot
(b) Slope field plot


## Verification of solutions

$$
y=-2 \csc (t) \cos (t)^{3}+2 \csc (t) \sqrt{2}
$$

Verified OK.

### 8.4.5 Maple step by step solution

Let's solve

$$
\left[y^{\prime}+\cot (t) y=6 \cos (t)^{2}, y\left(\frac{\pi}{4}\right)=3\right]
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=-\cot (t) y+6 \cos (t)^{2}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}+\cot (t) y=6 \cos (t)^{2}$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(y^{\prime}+\cot (t) y\right)=6 \mu(t) \cos (t)^{2}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) y)$
$\mu(t)\left(y^{\prime}+\cot (t) y\right)=\mu^{\prime}(t) y+\mu(t) y^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=\mu(t) \cot (t)$
- $\quad$ Solve to find the integrating factor
$\mu(t)=\sin (t)$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) y)\right) d t=\int 6 \mu(t) \cos (t)^{2} d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) y=\int 6 \mu(t) \cos (t)^{2} d t+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int 6 \mu(t) \cos (t)^{2} d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\sin (t)$
$y=\frac{\int 6 \sin (t) \cos (t)^{2} d t+c_{1}}{\sin (t)}$
- Evaluate the integrals on the rhs
$y=\frac{-2 \cos (t)^{3}+c_{1}}{\sin (t)}$
- Simplify
$y=\csc (t)\left(-2 \cos (t)^{3}+c_{1}\right)$
- Use initial condition $y\left(\frac{\pi}{4}\right)=3$
$3=\sqrt{2}\left(-\frac{\sqrt{2}}{2}+c_{1}\right)$
- $\quad$ Solve for $c_{1}$
$c_{1}=2 \sqrt{2}$
- $\quad$ Substitute $c_{1}=2 \sqrt{2}$ into general solution and simplify
$y=-2\left(\cos (t)^{3}-\sqrt{2}\right) \csc (t)$
- Solution to the IVP
$y=-2\left(\cos (t)^{3}-\sqrt{2}\right) \csc (t)$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.031 (sec). Leaf size: 18

```
dsolve([diff (y (t),t)=-cot(t)*y(t)+6*\operatorname{cos}(t)~2,y(1/4*Pi) = 3],y(t), singsol=all)
```

$$
y(t)=-2 \csc (t)\left(\cos (t)^{3}-\sqrt{2}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.06 (sec). Leaf size: 23

```
DSolve[{y'[t]==-Cot[t]*y[t]+6*Cos[t]^2,y[Pi/4]==3},y[t],t,IncludeSingularSolutions -> True]
```

$$
y(t) \rightarrow 2 \sqrt{2} \csc (t)-2 \cos ^{2}(t) \cot (t)
$$

