## A Solution Manual For

## DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.



Nasser M. Abbasi

May 15, 2024

## Contents

1 Chapter 8. Differential equations. Exercises page 595
1 Chapter 8. Differential equations. Exercises page 595
1.1 problem 1 ..... 6
1.2 problem 2 ..... 19
1.3 problem 3 ..... 23
1.4 problem 4 ..... 29
1.5 problem 5 ..... 43
1.6 problem 6 ..... 50
1.7 problem 7 ..... 61
1.8 problem 8 ..... 76
1.9 problem 9 ..... 90
1.10 problem 10 ..... 104
1.11 problem 11 ..... 116
1.12 problem 12 ..... 134
1.13 problem 13 ..... 146
1.14 problem 14 ..... 156
1.15 problem 15 ..... 167
1.16 problem 16 ..... 181
1.17 problem 17 ..... 195
1.18 problem 21 ..... 210
1.19 problem 22 ..... 222
1.20 problem 23 ..... 234
1.21 problem 24 ..... 246
1.22 problem 39 ..... 262
1.23 problem 40 ..... 278
1.24 problem 41 ..... 293
1.25 problem 42 ..... 307
1.26 problem 43 ..... 315
1.27 problem 44 ..... 324
1.28 problem 45 ..... 332
1.29 problem 46 ..... 346
1.30 problem 47 ..... 362
1.31 problem 48 ..... 377
1.32 problem 49 ..... 388
1.33 problem 50 ..... 396
1.34 problem 52 ..... 412
1.35 problem 53 ..... 425
1.36 problem 55 ..... 437
1.37 problem 56 ..... 449
1.38 problem 57 ..... 457
1.39 problem 58 ..... 470
1.40 problem 59 ..... 481
1.41 problem 60 ..... 492
1.42 problem 61 ..... 505
1.43 problem 62 ..... 518
1.44 problem 63 ..... 529
1.45 problem 64 ..... 540
1.46 problem 65 ..... 553
1.47 problem 66 ..... 566
1.48 problem 67 ..... 575
1.49 problem 68 ..... 589
1.50 problem 69 ..... 601
1.51 problem 70 ..... 603
1.52 problem 71 ..... 614
1.53 problem 72 ..... 625
1.54 problem 73 ..... 633
1.55 problem 74 ..... 641
1.56 problem 75 ..... 655
1.57 problem 76 ..... 662
1.58 problem 77 ..... 670
1.59 problem 78 ..... 684
1.60 problem 79 ..... 699
1.61 problem 80 ..... 714
1.62 problem 89 ..... 727
1.63 problem 90 ..... 733
1.64 problem 91 ..... 738
1.65 problem 92 ..... 743
1.66 problem 94 ..... 749
1.67 problem 95 ..... 752
1.68 problem 96 ..... 756
1.69 problem 97 ..... 774
1.70 problem 98 ..... 778
1.71 problem 110 ..... 783
1.72 problem 116 ..... 798
1.73 problem 117 ..... 803
1.74 problem 118 ..... 813
1.75 problem 120 ..... 818
1.76 problem 121 ..... 827
1.77 problem 122 ..... 833
1.78 problem 123 ..... 860
1.79 problem 124 ..... 866
1.80 problem 125 ..... 873
1.81 problem 126 ..... 887
1.82 problem 127 ..... 897
1.83 problem 128 ..... 899
1.84 problem 129 ..... 901
1.85 problem 130 ..... 911
1.86 problem 131 ..... 921
1.87 problem 132 ..... 931
1.88 problem 133 ..... 939
1.89 problem 134 ..... 948
1.90 problem 135 ..... 956
1.91 problem 136 ..... 964
1.92 problem 137 ..... 973
1.93 problem 140 ..... 982
1.94 problem 141 ..... 988
1.95 problem 142 ..... 993
1.96 problem 143 ..... 995
1.97 problem 144 ..... 1002
1.98 problem 145 ..... 1009
1.99 problem 146 ..... 1016
1.100problem 147 ..... 1023
1.101problem 148 ..... 1025
1.102problem 149 ..... 1035
1.103problem 150 ..... 1045
1.104problem 151 ..... 1056
1.105problem 152 ..... 1067
1.106problem 153 ..... 1078
1.107problem 154 ..... 1089
1.108problem 155 ..... 1100
1.109problem 156 ..... 1120
1.110problem 157 ..... 1131
1.111problem 158 ..... 1142
1.112problem 159 ..... 1150
1.113problem 160 ..... 1154
1.114problem 162 ..... 1160
1.115problem 163 ..... 1170
1.116problem 167 ..... 1183
1.117problem 168 ..... 1194
1.118problem 169 ..... 1207
1.119problem 170 ..... 1221
1.120problem 171 ..... 1231
1.121 problem 172 ..... 1238
1.122problem 181 ..... 1250
1.123problem 182 ..... 1257
1.124problem 183 ..... 1272
1.125problem 184 ..... 1277
1.126problem 185 ..... 1290
1.127problem 186 ..... 1303
1.128problem 187 ..... 1318
1.129problem 188 ..... 1329
1.130problem 189 ..... 1346
1.131problem 190 ..... 1354
1.132problem 191 ..... 1366
1.133problem 192 ..... 1375
1.134problem 193 ..... 1384
1.135problem 194 ..... 1393
1.136problem 195 ..... 1398
1.137problem 196 ..... 1411

## 1.1 problem 1

1.1.1 Solving as linear ode
1.1.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 8
1.1.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 12
1.1.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 17

Internal problem ID [12418]
Internal file name [OUTPUT/11070_Monday_October_16_2023_09_41_18_PM_64086885/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 1.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}+y \cos (x)=\frac{\sin (2 x)}{2}
$$

### 1.1.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=\cos (x) \\
& q(x)=\frac{\sin (2 x)}{2}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+y \cos (x)=\frac{\sin (2 x)}{2}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \cos (x) d x} \\
& =\mathrm{e}^{\sin (x)}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\frac{\sin (2 x)}{2}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\mathrm{e}^{\sin (x)} y\right) & =\left(\mathrm{e}^{\sin (x)}\right)\left(\frac{\sin (2 x)}{2}\right) \\
\mathrm{d}\left(\mathrm{e}^{\sin (x)} y\right) & =\left(\frac{\sin (2 x) \mathrm{e}^{\sin (x)}}{2}\right) \mathrm{d} x
\end{aligned}
$$

## Integrating gives

$$
\begin{aligned}
& \mathrm{e}^{\sin (x)} y=\int \frac{\sin (2 x) \mathrm{e}^{\sin (x)}}{2} \mathrm{~d} x \\
& \mathrm{e}^{\sin (x)} y=\sin (x) \mathrm{e}^{\sin (x)}-\mathrm{e}^{\sin (x)}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{\sin (x)}$ results in

$$
y=\mathrm{e}^{-\sin (x)}\left(\sin (x) \mathrm{e}^{\sin (x)}-\mathrm{e}^{\sin (x)}\right)+c_{1} \mathrm{e}^{-\sin (x)}
$$

which simplifies to

$$
y=\sin (x)-1+c_{1} \mathrm{e}^{-\sin (x)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\sin (x)-1+c_{1} \mathrm{e}^{-\sin (x)} \tag{1}
\end{equation*}
$$



Figure 1: Slope field plot
Verification of solutions

$$
y=\sin (x)-1+c_{1} \mathrm{e}^{-\sin (x)}
$$

Verified OK.

### 1.1.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=-\cos (x) y+\frac{\sin (2 x)}{2} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 1: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{2}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\mathrm{e}^{-\sin (x)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{-\sin (x)}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{\sin (x)} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\cos (x) y+\frac{\sin (2 x)}{2}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\cos (x) \mathrm{e}^{\sin (x)} y \\
S_{y} & =\mathrm{e}^{\sin (x)}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{\sin (2 x) \mathrm{e}^{\sin (x)}}{2} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{\sin (2 R) \mathrm{e}^{\sin (R)}}{2}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1}+\mathrm{e}^{\sin (R)}(-1+\sin (R)) \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\mathrm{e}^{\sin (x)} y=\mathrm{e}^{\sin (x)}(-1+\sin (x))+c_{1}
$$

Which simplifies to

$$
\mathrm{e}^{\sin (x)} y=\mathrm{e}^{\sin (x)}(-1+\sin (x))+c_{1}
$$

Which gives

$$
y=\mathrm{e}^{-\sin (x)}\left(\sin (x) \mathrm{e}^{\sin (x)}-\mathrm{e}^{\sin (x)}+c_{1}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical <br> coordinates <br> transformation | ODE in canonical coordinates <br> $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=-\cos (x) y+\frac{\sin (2 x)}{2}$ |  | $\frac{d S}{d R}=\frac{\sin (2 R) e^{\sin (R)}}{2}$ |
| 为 |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{-\sin (x)}\left(\sin (x) \mathrm{e}^{\sin (x)}-\mathrm{e}^{\sin (x)}+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 2: Slope field plot

## Verification of solutions

$$
y=\mathrm{e}^{-\sin (x)}\left(\sin (x) \mathrm{e}^{\sin (x)}-\mathrm{e}^{\sin (x)}+c_{1}\right)
$$

Verified OK.

### 1.1.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(-\cos (x) y+\frac{\sin (2 x)}{2}\right) \mathrm{d} x \\
\left(\cos (x) y-\frac{\sin (2 x)}{2}\right) \mathrm{d} x+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=\cos (x) y-\frac{\sin (2 x)}{2} \\
& N(x, y)=1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\cos (x) y-\frac{\sin (2 x)}{2}\right) \\
& =\cos (x)
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =1((\cos (x))-(0)) \\
& =\cos (x)
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int \cos (x) \mathrm{d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{\sin (x)} \\
& =\mathrm{e}^{\sin (x)}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\mathrm{e}^{\sin (x)}\left(\cos (x) y-\frac{\sin (2 x)}{2}\right) \\
& =\cos (x)(-\sin (x)+y) \mathrm{e}^{\sin (x)}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\mathrm{e}^{\sin (x)}(1) \\
& =\mathrm{e}^{\sin (x)}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x}=0 \\
\left(\cos (x)(-\sin (x)+y) \mathrm{e}^{\sin (x)}\right)+\left(\mathrm{e}^{\sin (x)}\right) \frac{\mathrm{d} y}{\mathrm{~d} x}=0
\end{array}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \cos (x)(-\sin (x)+y) \mathrm{e}^{\sin (x)} \mathrm{d} x \\
\phi & =(y-\sin (x)+1) \mathrm{e}^{\sin (x)}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\mathrm{e}^{\sin (x)}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\mathrm{e}^{\sin (x)}$. Therefore equation (4) becomes

$$
\begin{equation*}
\mathrm{e}^{\sin (x)}=\mathrm{e}^{\sin (x)}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=(y-\sin (x)+1) \mathrm{e}^{\sin (x)}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=(y-\sin (x)+1) \mathrm{e}^{\sin (x)}
$$

The solution becomes

$$
y=\mathrm{e}^{-\sin (x)}\left(\sin (x) \mathrm{e}^{\sin (x)}-\mathrm{e}^{\sin (x)}+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{-\sin (x)}\left(\sin (x) \mathrm{e}^{\sin (x)}-\mathrm{e}^{\sin (x)}+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 3: Slope field plot

Verification of solutions

$$
y=\mathrm{e}^{-\sin (x)}\left(\sin (x) \mathrm{e}^{\sin (x)}-\mathrm{e}^{\sin (x)}+c_{1}\right)
$$

## Verified OK.

### 1.1.4 Maple step by step solution

Let's solve
$y^{\prime}+y \cos (x)=\frac{\sin (2 x)}{2}$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=-y \cos (x)+\frac{\sin (2 x)}{2}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE
$y^{\prime}+y \cos (x)=\frac{\sin (2 x)}{2}$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}+y \cos (x)\right)=\frac{\mu(x) \sin (2 x)}{2}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$
$\mu(x)\left(y^{\prime}+y \cos (x)\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=\mu(x) \cos (x)$
- Solve to find the integrating factor
$\mu(x)=\mathrm{e}^{\sin (x)}$
- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int \frac{\mu(x) \sin (2 x)}{2} d x+c_{1}$
- Evaluate the integral on the lhs
$\mu(x) y=\int \frac{\mu(x) \sin (2 x)}{2} d x+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \frac{\mu(x) \sin (2 x)}{2} d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\mathrm{e}^{\sin (x)}$
$y=\frac{\int \frac{\sin (2 x)^{\sin (x)}}{\mathrm{e}^{\sin (x)}} d x+c_{1}}{}$
- Evaluate the integrals on the rhs
$y=\frac{\sin (x) \mathrm{e}^{\sin (x)}-\mathrm{e}^{\sin (x)}+c_{1}}{\mathrm{e}^{\sin (x)}}$
- Simplify

$$
y=\sin (x)-1+c_{1} \mathrm{e}^{-\sin (x)}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 15
dsolve $(\operatorname{diff}(y(x), x)+y(x) * \cos (x)=1 / 2 * \sin (2 * x), y(x), \quad$ singsol=all)

$$
y(x)=\sin (x)-1+\mathrm{e}^{-\sin (x)} c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.082 (sec). Leaf size: 18
DSolve [y' $[x]+y[x] * \operatorname{Cos}[x]==1 / 2 * \operatorname{Sin}[2 * x], y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow \sin (x)+c_{1} e^{-\sin (x)}-1
$$

## 1.2 problem 2

$$
\text { 1.2.1 Solving as clairaut ode . . . . . . . . . . . . . . . . . . . . . . . } 19
$$

Internal problem ID [12419]
Internal file name [OUTPUT/11071_Monday_October_16_2023_09_41_20_PM_35752472/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 2.
ODE order: 1.
ODE degree: 2 .

The type(s) of ODE detected by this program : "clairaut"
Maple gives the following as the ode type

```
[[_1st_order, _with_linear_symmetries], _Clairaut]
```

$$
y^{\prime 2}-y^{\prime}-y^{\prime} x+y=0
$$

### 1.2.1 Solving as clairaut ode

This is Clairaut ODE. It has the form

$$
y=y^{\prime} x+g\left(y^{\prime}\right)
$$

Where $g$ is function of $y^{\prime}(x)$. Let $p=y^{\prime}$ the ode becomes

$$
p^{2}-p x-p+y=0
$$

Solving for $y$ from the above results in

$$
\begin{equation*}
y=-p^{2}+p x+p \tag{1~A}
\end{equation*}
$$

The above ode is a Clairaut ode which is now solved. We start by replacing $y^{\prime}$ by $p$ which gives

$$
\begin{aligned}
y & =-p^{2}+p x+p \\
& =-p^{2}+p x+p
\end{aligned}
$$

Writing the ode as

$$
y=p x+g(p)
$$

We now write $g \equiv g(p)$ to make notation simpler but we should always remember that $g$ is function of $p$ which in turn is function of $x$. Hence the above becomes

$$
\begin{equation*}
y=p x+g \tag{1}
\end{equation*}
$$

Then we see that

$$
g=-p^{2}+p
$$

Taking derivative of (1) w.r.t. $x$ gives

$$
\begin{aligned}
& p=\frac{d}{d x}(x p+g) \\
& p=\left(p+x \frac{d p}{d x}\right)+\left(g^{\prime} \frac{d p}{d x}\right) \\
& p=p+\left(x+g^{\prime}\right) \frac{d p}{d x} \\
& 0=\left(x+g^{\prime}\right) \frac{d p}{d x}
\end{aligned}
$$

Where $g^{\prime}$ is derivative of $g(p)$ w.r.t. $p$. The general solution is given by

$$
\begin{aligned}
\frac{d p}{d x} & =0 \\
p & =c_{1}
\end{aligned}
$$

Substituting this in (1) gives the general solution as

$$
y=-c_{1}^{2}+c_{1} x+c_{1}
$$

The singular solution is found from solving for $p$ from

$$
x+g^{\prime}(p)=0
$$

And substituting the result back in (1). Since we found above that $g=-p^{2}+p$, then the above equation becomes

$$
\begin{aligned}
x+g^{\prime}(p) & =x-2 p+1 \\
& =0
\end{aligned}
$$

Solving the above for $p$ results in

$$
p_{1}=\frac{x}{2}+\frac{1}{2}
$$

Substituting the above back in (1) results in

$$
y_{1}=\frac{(x+1)^{2}}{4}
$$

## Summary

The solution(s) found are the following

$$
\begin{align*}
& y=-c_{1}^{2}+c_{1} x+c_{1}  \tag{1}\\
& y=\frac{(x+1)^{2}}{4} \tag{2}
\end{align*}
$$

Verification of solutions

$$
y=-c_{1}^{2}+c_{1} x+c_{1}
$$

Verified OK.

$$
y=\frac{(x+1)^{2}}{4}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
    *** Sublevel 2 ***
    Methods for first order ODEs:
    -> Solving 1st order ODE of high degree, 1st attempt
    trying 1st order WeierstrassP solution for high degree ODE
    trying 1st order WeierstrassPPrime solution for high degree ODE
    trying 1st order JacobiSN solution for high degree ODE
    trying 1st order ODE linearizable_by_differentiation
    trying differential order: 1; missing variables
    trying dAlembert
    <- dAlembert successful`
```

$\checkmark$ Solution by Maple
Time used: 0.078 (sec). Leaf size: 22
dsolve(diff $(y(x), x) \sim 2-\operatorname{diff}(y(x), x)-x * \operatorname{diff}(y(x), x)+y(x)=0, y(x)$, singsol=all)

$$
\begin{aligned}
& y(x)=\frac{(1+x)^{2}}{4} \\
& y(x)=c_{1}\left(-c_{1}+x+1\right)
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.011 (sec). Leaf size: 28
DSolve[(y'[x])~2-y'[x]-x*y'[x]+y[x]==0,y[x],x,IncludeSingularSolutions -> True]

$$
\begin{aligned}
& y(x) \rightarrow c_{1}\left(x+1-c_{1}\right) \\
& y(x) \rightarrow \frac{1}{4}(x+1)^{2}
\end{aligned}
$$

## 1.3 problem 3

> 1.3.1 Solving as dAlembert ode

Internal problem ID [12420]
Internal file name [OUTPUT/11072_Monday_October_16_2023_09_41_20_PM_39372808/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 3.
ODE order: 1.
ODE degree: 2 .

The type(s) of ODE detected by this program : "dAlembert"
Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _dAlembert]
```

$$
y y^{\prime 2}+2 y^{\prime} x-y=0
$$

### 1.3.1 Solving as dAlembert ode

Let $p=y^{\prime}$ the ode becomes

$$
y p^{2}+2 p x-y=0
$$

Solving for $y$ from the above results in

$$
\begin{equation*}
y=-\frac{2 p x}{p^{2}-1} \tag{1~A}
\end{equation*}
$$

This has the form

$$
\begin{equation*}
y=x f(p)+g(p) \tag{}
\end{equation*}
$$

Where $f, g$ are functions of $p=y^{\prime}(x)$. The above ode is dAlembert ode which is now solved. Taking derivative of $\left({ }^{*}\right)$ w.r.t. $x$ gives

$$
\begin{align*}
p & =f+\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \\
p-f & =\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \tag{2}
\end{align*}
$$

Comparing the form $y=x f+g$ to (1A) shows that

$$
\begin{aligned}
& f=-\frac{2 p}{p^{2}-1} \\
& g=0
\end{aligned}
$$

Hence (2) becomes

$$
\begin{equation*}
p+\frac{2 p}{p^{2}-1}=x\left(-\frac{2}{p^{2}-1}+\frac{4 p^{2}}{\left(p^{2}-1\right)^{2}}\right) p^{\prime}(x) \tag{2~A}
\end{equation*}
$$

The singular solution is found by setting $\frac{d p}{d x}=0$ in the above which gives

$$
p+\frac{2 p}{p^{2}-1}=0
$$

Solving for $p$ from the above gives

$$
\begin{aligned}
& p=0 \\
& p=i \\
& p=-i
\end{aligned}
$$

Substituting these in (1A) gives

$$
\begin{aligned}
& y=0 \\
& y=-i x \\
& y=i x
\end{aligned}
$$

The general solution is found when $\frac{\mathrm{d} p}{\mathrm{~d} x} \neq 0$. From eq. (2A). This results in

$$
\begin{equation*}
p^{\prime}(x)=\frac{p(x)+\frac{2 p(x)}{p(x)^{2}-1}}{x\left(-\frac{2}{p(x)^{2}-1}+\frac{4 p(x)^{2}}{\left(p(x)^{2}-1\right)^{2}}\right)} \tag{3}
\end{equation*}
$$

This ODE is now solved for $p(x)$.
Inverting the above ode gives

$$
\begin{equation*}
\frac{d}{d p} x(p)=\frac{x(p)\left(-\frac{2}{p^{2}-1}+\frac{4 p^{2}}{\left(p^{2}-1\right)^{2}}\right)}{p+\frac{2 p}{p^{2}-1}} \tag{4}
\end{equation*}
$$

This ODE is now solved for $x(p)$.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
\frac{d}{d p} x(p)+p(p) x(p)=q(p)
$$

Where here

$$
\begin{aligned}
& p(p)=-\frac{2}{p^{3}-p} \\
& q(p)=0
\end{aligned}
$$

Hence the ode is

$$
\frac{d}{d p} x(p)-\frac{2 x(p)}{p^{3}-p}=0
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2}{p^{3}-p} d p} \\
& =\mathrm{e}^{-\ln (p-1)-\ln (p+1)+2 \ln (p)}
\end{aligned}
$$

Which simplifies to

$$
\mu=\frac{p^{2}}{p^{2}-1}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} p} \mu x & =0 \\
\frac{\mathrm{~d}}{\mathrm{~d} p}\left(\frac{p^{2} x}{p^{2}-1}\right) & =0
\end{aligned}
$$

Integrating gives

$$
\frac{p^{2} x}{p^{2}-1}=c_{3}
$$

Dividing both sides by the integrating factor $\mu=\frac{p^{2}}{p^{2}-1}$ results in

$$
x(p)=\frac{c_{3}\left(p^{2}-1\right)}{p^{2}}
$$

Now we need to eliminate $p$ between the above and (1A). One way to do this is by solving (1) for $p$. This results in

$$
\begin{aligned}
& p=\frac{\sqrt{x^{2}+y^{2}}-x}{y} \\
& p=-\frac{x+\sqrt{x^{2}+y^{2}}}{y}
\end{aligned}
$$

Substituting the above in the solution for $x$ found above gives

$$
\begin{aligned}
& x=-\frac{2 c_{3} x}{\sqrt{x^{2}+y^{2}}-x} \\
& x=\frac{2 c_{3} x}{x+\sqrt{x^{2}+y^{2}}}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=0  \tag{1}\\
& y=-i x  \tag{2}\\
& y=i x  \tag{3}\\
& x=-\frac{2 c_{3} x}{\sqrt{x^{2}+y^{2}}-x}  \tag{4}\\
& x=\frac{2 c_{3} x}{x+\sqrt{x^{2}+y^{2}}} \tag{5}
\end{align*}
$$

Verification of solutions

$$
y=0
$$

Verified OK.

$$
y=-i x
$$

Verified OK.

$$
y=i x
$$

Verified OK.

$$
x=-\frac{2 c_{3} x}{\sqrt{x^{2}+y^{2}}-x}
$$

Verified OK.

$$
x=\frac{2 c_{3} x}{x+\sqrt{x^{2}+y^{2}}}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
    *** Sublevel 2 ***
    Methods for first order ODEs:
    -> Solving 1st order ODE of high degree, 1st attempt
    trying 1st order WeierstrassP solution for high degree ODE
    trying 1st order WeierstrassPPrime solution for high degree ODE
    trying 1st order JacobiSN solution for high degree ODE
    trying 1st order ODE linearizable_by_differentiation
    trying differential order: 1; missing variables
    trying simple symmetries for implicit equations
    <- symmetries for implicit equations successful`
```

$\checkmark$ Solution by Maple
Time used: 0.672 (sec). Leaf size: 71
dsolve( $y(x) * \operatorname{diff}(y(x), x)^{\sim} 2+2 * x * \operatorname{diff}(y(x), x)-y(x)=0, y(x)$, singsol=all)

$$
\begin{aligned}
& y(x)=-i x \\
& y(x)=i x \\
& y(x)=0 \\
& y(x)=\sqrt{c_{1}\left(-2 x+c_{1}\right)} \\
& y(x)=\sqrt{c_{1}\left(c_{1}+2 x\right)} \\
& y(x)=-\sqrt{c_{1}\left(-2 x+c_{1}\right)} \\
& y(x)=-\sqrt{c_{1}\left(c_{1}+2 x\right)}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.759 (sec). Leaf size: 126
DSolve $\left[y[x] * y '[x] \sim 2+2 * x * y^{\prime}[x]-y[x]==0, y[x], x\right.$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow-e^{\frac{c_{1}}{2}} \sqrt{-2 x+e^{c_{1}}} \\
& y(x) \rightarrow e^{\frac{c_{1}}{2}} \sqrt{-2 x+e^{c_{1}}} \\
& y(x) \rightarrow-e^{\frac{c_{1}}{2}} \sqrt{2 x+e^{c_{1}}} \\
& y(x) \rightarrow e^{\frac{c_{1}}{2}} \sqrt{2 x+e^{c_{1}}} \\
& y(x) \rightarrow 0 \\
& y(x) \rightarrow-i x \\
& y(x) \rightarrow i x
\end{aligned}
$$

## 1.4 problem 4

Internal problem ID [12421]
Internal file name [OUTPUT/11073_Monday_October_16_2023_09_41_23_PM_40250331/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 4.
ODE order: 1.
ODE degree: 2.

The type(s) of ODE detected by this program : "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type
[_rational]

$$
x y\left(-y^{\prime 2}+1\right)-\left(x^{2}-y^{2}-a^{2}\right) y^{\prime}=0
$$

Solving the given ode for $y^{\prime}$ results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$
\begin{align*}
& y^{\prime}=\frac{y^{2}+a^{2}-x^{2}+\sqrt{y^{4}+2 y^{2} a^{2}+2 x^{2} y^{2}+a^{4}-2 a^{2} x^{2}+x^{4}}}{2 y x}  \tag{1}\\
& y^{\prime}=\frac{y^{2}+a^{2}-x^{2}-\sqrt{y^{4}+2 y^{2} a^{2}+2 x^{2} y^{2}+a^{4}-2 a^{2} x^{2}+x^{4}}}{2 y x} \tag{2}
\end{align*}
$$

Now each one of the above ODE is solved.
Solving equation (1)
Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{y^{2}+a^{2}-x^{2}+\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}{2 y x} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 3 to use as anstaz gives

$$
\begin{align*}
& \xi=x^{3} a_{7}+x^{2} y a_{8}+x y^{2} a_{9}+y^{3} a_{10}+x^{2} a_{4}+x y a_{5}+y^{2} a_{6}+x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x^{3} b_{7}+x^{2} y b_{8}+x y^{2} b_{9}+y^{3} b_{10}+x^{2} b_{4}+x y b_{5}+y^{2} b_{6}+x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, a_{4}, a_{5}, a_{6}, a_{7}, a_{8}, a_{9}, a_{10}, b_{1}, b_{2}, b_{3}, b_{4}, b_{5}, b_{6}, b_{7}, b_{8}, b_{9}, b_{10}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
& 3 x^{2} b_{7}+2 x y b_{8}+y^{2} b_{9}+2 x b_{4}+y b_{5}+b_{2}  \tag{5E}\\
& +\frac{\left(y^{2}+a^{2}-x^{2}+\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}\right)\left(-3 x^{2} a_{7}+x^{2} b_{8}-2 x y a_{8}+2 x y b_{9}-y^{2} a_{9}+3 y^{2}\right.}{2 y x} \\
& -\frac{\left(y^{2}+a^{2}-x^{2}+\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}\right)^{2}\left(x^{2} a_{8}+2 x y a_{9}+3 y^{2} a_{10}+x a_{5}+2 y a_{6}+a_{3}\right)}{4 y^{2} x^{2}} \\
& -\left(-\frac{y^{2}+a^{2}-x^{2}+\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}{2 y x^{2}}\right. \\
& \left.+\frac{-2 x+\frac{-4 a^{2} x+4 x^{3}+4 x y^{2}}{2 \sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}}{2 y x}\right)\left(x^{3} a_{7}+x^{2} y a_{8}\right. \\
& \left.+x y^{2} a_{9}+y^{3} a_{10}+x^{2} a_{4}+x y a_{5}+y^{2} a_{6}+x a_{2}+y a_{3}+a_{1}\right) \\
& -\left(-\frac{y^{2}+a^{2}-x^{2}+\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}{2 y^{2} x}\right. \\
& \left.+\frac{2 y+\frac{4 a^{2} y+4 x^{2} y+4 y^{3}}{2 \sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}}{2 y x}\right)\left(x^{3} b_{7}+x^{2} y b_{8}\right. \\
& \left.+x y^{2} b_{9}+y^{3} b_{10}+x^{2} b_{4}+x y b_{5}+y^{2} b_{6}+x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
$$

Putting the above in normal form gives
Expression too large to display

Setting the numerator to zero gives
Expression too large to display

Simplifying the above gives
Expression too large to display

Since the PDE has radicals, simplifying gives

## Expression too large to display

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\left\{x, y, \sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}\right\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}, \sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}=v_{3}\right\}
$$

The above PDE (6E) now becomes

> Expression too large to display

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}, v_{3}\right\}
$$

Equation (7E) now becomes

> Expression too large to display

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
& -2 a_{6}=0 \\
& -4 a_{10}=0 \\
& -14 a^{2} a_{10}=0 \\
& -10 a^{2} a_{10}=0 \\
& -2 a^{4} a_{3}=0 \\
& -2 a^{6} a_{3}=0 \\
& -2 a_{4}-4 a_{6}=0 \\
& -2 a_{5}-2 b_{4}=0 \\
& -2 a_{5}+2 b_{6}=0 \\
& -2 a_{5}+6 b_{6}=0 \\
& 2 a_{5}+2 b_{4}=0 \\
& -2 a_{8}-2 b_{7}=0 \\
& -4 a_{9}+4 b_{10}=0 \\
& 2 b_{7}+2 a_{8}=0 \\
& -8 a_{4}+2 a_{6}+4 b_{5}=0 \\
& -6 a_{4}+4 a_{6}+4 b_{5}=0 \\
& -2 a_{4}+2 a_{6}+4 b_{5}=0 \\
& 6 a_{4}-4 a_{6}-4 b_{5}=0 \\
& -6 a_{5}-2 b_{4}+8 b_{6}=0 \\
& 4 a_{5}+6 b_{4}-6 b_{6}=0 \\
& -4 a_{8}-6 a_{10}+2 b_{9}=0 \\
& -8 a_{9}+12 b_{10}-4 a_{7}=0 \\
& -4 a_{9}+8 a_{7}-4 b_{8}=0 \\
& 4 a_{9}-8 a_{7}+4 b_{8}=0 \\
& 2 a_{10}+6 b_{9}-4 a_{8}=0 \\
& 6 b_{9}+6 a_{10}-4 a_{8}=0 \\
& 8 b_{10}-12 a_{7}+4 b_{8}=0 \\
& -10 a_{8}+4 a_{10}-2 b_{7}+8 b_{9}=0 \\
& -6 a_{10}+10 b_{7}+6 a_{8}-6 b_{9}=0 \\
& 8 b_{8}-4 a_{7}+4 a_{9}-8 b_{10}=0 \\
& -8 a^{2} a_{6}+2 a_{1}=0 \\
& -6 a^{2} a_{6}+2 a_{1}=0 \\
& -10 a^{4} a_{6}+4 a^{2} a_{1}=0 \\
& -4 a^{4} a_{6}+2 a^{2} a_{1}=0 \\
& -4 a^{6} a_{6}+2 a^{4} a_{1}=0 \\
& -16 a^{4} a_{10}-2 a^{2} a_{3}=0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =-a^{2} b_{10} \\
a_{3} & =0 \\
a_{4} & =0 \\
a_{5} & =0 \\
a_{6} & =0 \\
a_{7} & =b_{10} \\
a_{8} & =0 \\
a_{9} & =b_{10} \\
a_{10} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
b_{3} & =a^{2} b_{10} \\
b_{4} & =0 \\
b_{5} & =0 \\
b_{6} & =0 \\
b_{7} & =0 \\
b_{8} & =b_{10} \\
b_{9} & =0 \\
b_{10} & =b_{10}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E, 2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=-a^{2} x+x^{3}+x y^{2} \\
& \eta=a^{2} y+x^{2} y+y^{3}
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Therefore

$$
\begin{aligned}
\frac{d y}{d x} & =\frac{\eta}{\xi} \\
& =\frac{a^{2} y+x^{2} y+y^{3}}{-a^{2} x+x^{3}+x y^{2}} \\
& =-\frac{y\left(a^{2}+x^{2}+y^{2}\right)}{x\left(a^{2}-x^{2}-y^{2}\right)}
\end{aligned}
$$

This is easily solved to give

$$
\frac{1}{\frac{1}{y^{2}}-\frac{1}{a^{2}-x^{2}}}=-\frac{\sqrt{x-a} \sqrt{x+a} x}{\sqrt{c_{1}+8 a^{2}\left(\frac{1}{4 a(x-a)}-\frac{1}{4 a(x+a)}\right)}}-\frac{(a-x)(x+a)}{2}
$$

Where now the coordinate $R$ is taken as the constant of integration. Hence

$$
R=\frac{4 a^{4}-8 a^{2} x^{2}+8 y^{2} a^{2}+4 x^{4}+8 x^{2} y^{2}+4 y^{4}}{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}-2 x^{2} y^{2}+y^{4}}
$$

Since $\xi$ depends on $y$ and $\eta$ depends on $x$ then we can use either one to find $S$. Let us use

$$
\begin{aligned}
d S & =\frac{d x}{\xi} \\
& =\frac{d x}{-a^{2} x+x^{3}+x y^{2}}
\end{aligned}
$$

But we have now to replace $y$ in $\xi$ from its value from the solution of $\frac{d y}{d x}=\frac{\eta}{\xi}$ found above. This results in

$$
\xi=-a^{2} x+x^{3}+x\left(-\frac{\sqrt{x-a} \sqrt{x+a} x}{\sqrt{c_{1}+8 a^{2}\left(\frac{1}{4 a(x-a)}-\frac{1}{4 a(x+a)}\right)}}-\frac{(a-x)(x+a)}{2}\right)^{2}
$$

Integrating gives

$$
S=\frac{d x}{-a^{2} x+x^{3}+x\left(-\frac{\sqrt{x-a} \sqrt{x+a} x}{\sqrt{c_{1}+8 a^{2}\left(\frac{1}{4 a(x-a)}-\frac{1}{4 a(x+a)}\right)}}-\frac{(a-x)(x+a)}{2}\right)^{2}}
$$

$=$ Expression too large to display

Where the constant of integration is set to zero as we just need one solution. Replacing back $c_{1}=\frac{4 a^{4}-8 a^{2} x^{2}+8 y^{2} a^{2}+4 x^{4}+8 x^{2} y^{2}+4 y^{4}}{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}-2 x^{2} y^{2}+y^{4}}$ then the above becomes

$$
S=\text { Expression too large to display }
$$

Unable to determine ODE type.
Solving equation (2)
Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{y^{2}+a^{2}-x^{2}-\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}{2 y x} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 3 to use as anstaz gives

$$
\begin{align*}
& \xi=x^{3} a_{7}+x^{2} y a_{8}+x y^{2} a_{9}+y^{3} a_{10}+x^{2} a_{4}+x y a_{5}+y^{2} a_{6}+x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x^{3} b_{7}+x^{2} y b_{8}+x y^{2} b_{9}+y^{3} b_{10}+x^{2} b_{4}+x y b_{5}+y^{2} b_{6}+x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, a_{4}, a_{5}, a_{6}, a_{7}, a_{8}, a_{9}, a_{10}, b_{1}, b_{2}, b_{3}, b_{4}, b_{5}, b_{6}, b_{7}, b_{8}, b_{9}, b_{10}\right\}
$$

Substituting equations (1E, 2 E ) and $\omega$ into (A) gives

$$
\begin{aligned}
& 3 x^{2} b_{7}+2 x y b_{8}+y^{2} b_{9}+2 x b_{4}+y b_{5}+b_{2} \\
& +\frac{\left(y^{2}+a^{2}-x^{2}-\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}\right)\left(-3 x^{2} a_{7}+x^{2} b_{8}-2 x y a_{8}+2 x y b_{9}-y^{2} a_{9}+3 y\right.}{2 y x} \\
& -\frac{\left(y^{2}+a^{2}-x^{2}-\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}\right)^{2}\left(x^{2} a_{8}+2 x y a_{9}+3 y^{2} a_{10}+x a_{5}+2 y a_{6}+a_{3}\right)}{4 y^{2} x^{2}} \\
& -\left(-\frac{y^{2}+a^{2}-x^{2}-\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}{2 y x^{2}}\right. \\
& \left.+\frac{-2 x-\frac{-4 a^{2} x+4 x^{3}+4 x y^{2}}{2 \sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}}{2 y x}\right)\left(x^{3} a_{7}+x^{2} y a_{8}\right. \\
& \left.+x y^{2} a_{9}+y^{3} a_{10}+x^{2} a_{4}+x y a_{5}+y^{2} a_{6}+x a_{2}+y a_{3}+a_{1}\right) \\
& -\left(-\frac{y^{2}+a^{2}-x^{2}-\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}{2 y^{2} x}\right. \\
& \left.+\frac{2 y-\frac{4 a^{2} y+4 x^{2} y+4 y^{3}}{2 \sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}}{2 y x}\right)\left(x^{3} b_{7}+x^{2} y b_{8}\right. \\
& \left.+x y^{2} b_{9}+y^{3} b_{10}+x^{2} b_{4}+x y b_{5}+y^{2} b_{6}+x b_{2}+y b_{3}+b_{1}\right)=0
\end{aligned}
$$

Putting the above in normal form gives

Expression too large to display

Setting the numerator to zero gives

> Expression too large to display

Simplifying the above gives

Expression too large to display

Since the PDE has radicals, simplifying gives
Expression too large to display

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\left\{x, y, \sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}\right\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}, \sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}=v_{3}\right\}
$$

The above PDE (6E) now becomes

> Expression too large to display

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}, v_{3}\right\}
$$

Equation (7E) now becomes

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
& -2 a_{6}=0 \\
& 2 a_{6}=0 \\
& -4 a_{10}=0 \\
& 4 a_{10}=0 \\
& -10 a^{2} a_{10}=0 \\
& 14 a^{2} a_{10}=0 \\
& -2 a^{4} a_{3}=0 \\
& 2 a^{6} a_{3}=0 \\
& 2 a_{4}+4 a_{6}=0 \\
& -2 a_{5}-2 b_{4}=0 \\
& -2 a_{5}+2 b_{6}=0 \\
& 2 a_{5}-6 b_{6}=0 \\
& 2 a_{5}-2 b_{6}=0 \\
& -2 a_{8}-2 b_{7}=0 \\
& -4 a_{9}+4 b_{10}=0 \\
& 4 a_{9}-4 b_{10}=0 \\
& -2 a_{4}+2 a_{6}+4 b_{5}=0 \\
& 6 a_{4}-4 a_{6}-4 b_{5}=0 \\
& 8 a_{4}-2 a_{6}-4 b_{5}=0 \\
& 4 a_{5}+6 b_{4}-6 b_{6}=0 \\
& 6 a_{5}+2 b_{4}-8 b_{6}=0 \\
& 4 a_{7}+8 a_{9}-12 b_{10}=0 \\
& -4 a_{8}+2 a_{10}+6 b_{9}=0 \\
& 4 a_{8}+6 a_{10}-2 b_{9}=0 \\
& -4 a_{9}+8 a_{7}-4 b_{8}=0 \\
& -6 b_{9}-6 a_{10}+4 a_{8}=0 \\
& -8 b_{10}+12 a_{7}-4 b_{8}=0 \\
& 10 a_{8}-4 a_{10}+2 b_{7}-8 b_{9}=0 \\
& -6 a_{10}+10 b_{7}+6 a_{8}-6 b_{9}=0 \\
& 8 b_{8}-4 a_{7}+4 a_{9}-8 b_{10}=0 \\
& -6 a^{2} a_{6}+2 a_{1}=0 \\
& 8 a^{2} a_{6}-2 a_{1}=0 \\
& -4 a^{4} a_{6}+2 a^{2} a_{1}=0 \\
& 10 a^{4} a_{6}-4 a^{2} a_{1}=0 \\
& 4 a^{6} a_{6}-2 a^{4} a_{1}=0 \\
& -6 a^{4} a_{10}-2 a^{2} a_{3}=0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =-a^{2} b_{10} \\
a_{3} & =0 \\
a_{4} & =0 \\
a_{5} & =0 \\
a_{6} & =0 \\
a_{7} & =b_{10} \\
a_{8} & =0 \\
a_{9} & =b_{10} \\
a_{10} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
b_{3} & =a^{2} b_{10} \\
b_{4} & =0 \\
b_{5} & =0 \\
b_{6} & =0 \\
b_{7} & =0 \\
b_{8} & =b_{10} \\
b_{9} & =0 \\
b_{10} & =b_{10}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E, 2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=-a^{2} x+x^{3}+x y^{2} \\
& \eta=a^{2} y+x^{2} y+y^{3}
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =a^{2} y+x^{2} y+y^{3}-\left(\frac{y^{2}+a^{2}-x^{2}-\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}{2 y x}\right)\left(-a^{2} x+x^{3}+x y^{2}\right) \\
& =\frac{a^{4} x-2 x^{3} a^{2}+2 a^{2} x y^{2}+x^{5}+2 y^{2} x^{3}+x y^{4}-\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}} a^{2} x+x^{3} \sqrt{a^{4}-}}{2 x y} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{a^{4} x-2 x^{3} a^{2}+2 a^{2} x y^{2}+x^{5}+2 y^{2} x^{3}+x y^{4}-\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}} a^{2} x+x^{3} \sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}+x y^{2} \sqrt{a^{4}-}}{2 x y}}
\end{aligned}
$$

Which results in
$S=-\frac{\ln \left(\frac{2 a^{4}-4 a^{2} x^{2}+2 x^{4}+\left(2 a^{2}+2 x^{2}\right) y^{2}+2 \sqrt{\left(-a^{2}+x^{2}\right)^{2}} \sqrt{y^{4}+\left(2 a^{2}+2 x^{2}\right) y^{2}+a^{4}-2 a^{2} x^{2}+x^{4}}}{y^{2}}\right)}{4 \sqrt{\left(-a^{2}+x^{2}\right)^{2}}}+\frac{\ln \left(a^{2}-2 a x+x^{2}+y^{2}\right)}{16 a x}-$
Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{y^{2}+a^{2}-x^{2}-\sqrt{a^{4}-2 a^{2} x^{2}+2 y^{2} a^{2}+x^{4}+2 x^{2} y^{2}+y^{4}}}{2 y x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
& R_{x}=1 \\
& R_{y}=0 \\
& S_{x}=\frac{2 x}{\sqrt{\left(a^{2}-2 a x+x^{2}+y^{2}\right)\left(a^{2}+2 a x+x^{2}+y^{2}\right)}\left(a^{2}+x^{2}+y^{2}+\sqrt{\left(a^{2}-2 a x+x^{2}+y^{2}\right)\left(a^{2}+2 a x+x\right.}\right.} \\
& S_{y}=\frac{\left(x^{4}+\left(-2 a^{2}+y^{2}\right) x^{2}+a^{4}\right) \sqrt{\left(a^{2}-2 a x+x^{2}+y^{2}\right)\left(a^{2}+2 a x+x^{2}+y^{2}\right)}-x^{6}+\left(3 a^{2}-2 y^{2}\right) x^{4}}{\sqrt{\left(a^{2}-2 a x+x^{2}+y^{2}\right)\left(a^{2}+2 a x+x^{2}+y^{2}\right)} y a^{2}\left(\left(a^{2}-x^{2}\right) \sqrt{\left(a^{2}-2 a x+x^{2}+y^{2}\right)\left(a^{2}+2 a x+x^{2}\right.}\right.}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{-\ln (2)-\ln \left(\left(a^{2}-x^{2}\right) \sqrt{\left(y^{2}+(x+a)^{2}\right)\left(y^{2}+(a-x)^{2}\right)}+a^{4}+\left(y^{2}-2 x^{2}\right) a^{2}+x^{4}+x^{2} y^{2}\right)+4 \ln (y)-}{4 a^{2}}
$$

Which simplifies to

$$
\frac{-\ln (2)-\ln \left(\left(a^{2}-x^{2}\right) \sqrt{\left(y^{2}+(x+a)^{2}\right)\left(y^{2}+(a-x)^{2}\right)}+a^{4}+\left(y^{2}-2 x^{2}\right) a^{2}+x^{4}+x^{2} y^{2}\right)+4 \ln (y)-}{4 a^{2}}
$$

## Summary

The solution(s) found are the following

$$
\begin{aligned}
& -\ln (2)-\ln \left(\left(a^{2}-x^{2}\right) \sqrt{\left(y^{2}+(x+a)^{2}\right)\left(y^{2}+(a-x)^{2}\right)}+a^{4}+\left(y^{2}-2 x^{2}\right) a^{2}+x^{4}+x^{2} y^{2}\right)+4 \ln (y)- \\
& =c_{1}
\end{aligned}
$$

## Verification of solutions

$-\ln (2)-\ln \left(\left(a^{2}-x^{2}\right) \sqrt{\left(y^{2}+(x+a)^{2}\right)\left(y^{2}+(a-x)^{2}\right)}+a^{4}+\left(y^{2}-2 x^{2}\right) a^{2}+x^{4}+x^{2} y^{2}\right)+4 \ln (y)-$
$=c_{1}$
Verified OK.

X Solution by Maple
dsolve $(x * y(x) *(1-\operatorname{diff}(y(x), x) \wedge 2)=(x \wedge 2-y(x) \wedge 2-a \wedge 2) * \operatorname{diff}(y(x), x), y(x)$, singsol=all)

No solution found
$\checkmark$ Solution by Mathematica
Time used: 0.612 (sec). Leaf size: 75
DSolve $\left[x * y[x] *\left(1-y^{\prime}[x] \wedge 2\right)==\left(x^{\wedge} 2-y[x] \wedge 2-a^{\wedge} 2\right) * y^{\prime}[x], y[x], x\right.$, IncludeSingularSolutions $\rightarrow$ True

$$
\begin{aligned}
& y(x) \rightarrow \sqrt{c_{1}\left(x^{2}-\frac{a^{2}}{1+c_{1}}\right)} \\
& y(x) \rightarrow-i(a-x) \\
& y(x) \rightarrow i(a-x) \\
& y(x) \rightarrow-i(a+x) \\
& y(x) \rightarrow i(a+x)
\end{aligned}
$$

## 1.5 problem 5

1.5.1 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 45

Internal problem ID [12422]
Internal file name [OUTPUT/11074_Monday_October_16_2023_09_47_11_PM_10630605/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 5.
ODE order: 3.
ODE degree: 1.

The type(s) of ODE detected by this program : "higher_order_missing_y"
Maple gives the following as the ode type
[[_3rd_order, _missing_y]]

$$
y^{\prime \prime \prime}+\frac{3 y^{\prime \prime}}{x}=0
$$

Since $y$ is missing from the ode then we can use the substitution $y^{\prime}=v(x)$ to reduce the order by one. The ODE becomes

$$
v^{\prime \prime}(x) x+3 v^{\prime}(x)=0
$$

Integrating both sides of the ODE w.r.t $x$ gives

$$
\begin{gathered}
\int\left(v^{\prime \prime}(x) x+3 v^{\prime}(x)\right) d x=0 \\
v^{\prime}(x) x+2 v(x)=c_{1}
\end{gathered}
$$

Which is now solved for $v(x)$. In canonical form the ODE is

$$
\begin{aligned}
v^{\prime} & =F(x, v) \\
& =f(x) g(v) \\
& =\frac{-2 v+c_{1}}{x}
\end{aligned}
$$

Where $f(x)=\frac{1}{x}$ and $g(v)=-2 v+c_{1}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{-2 v+c_{1}} d v & =\frac{1}{x} d x \\
\int \frac{1}{-2 v+c_{1}} d v & =\int \frac{1}{x} d x \\
-\frac{\ln \left(-2 v+c_{1}\right)}{2} & =\ln (x)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\frac{1}{\sqrt{-2 v+c_{1}}}=\mathrm{e}^{\ln (x)+c_{2}}
$$

Which simplifies to

$$
\frac{1}{\sqrt{-2 v+c_{1}}}=c_{3} x
$$

But since $y^{\prime}=v(x)$ then we now need to solve the ode $y^{\prime}=\frac{\left(c_{1} c_{3}^{2} \mathrm{e}^{2 c_{2}} x^{2}-1\right) \mathrm{e}^{-2 c_{2}}}{2 c_{3}^{2} x^{2}}$. Integrating both sides gives

$$
\begin{aligned}
y & =\int \frac{\left(c_{1} c_{3}^{2} \mathrm{e}^{2 c_{2}} x^{2}-1\right) \mathrm{e}^{-2 c_{2}}}{2 c_{3}^{2} x^{2}} \mathrm{~d} x \\
& =\frac{\mathrm{e}^{-2 c_{2}}\left(\mathrm{e}^{2 c_{2}} c_{1} c_{3}^{2} x+\frac{1}{x}\right)}{2 c_{3}^{2}}+c_{4}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\mathrm{e}^{-2 c_{2}}\left(\mathrm{e}^{2 c_{2}} c_{1} c_{3}^{2} x+\frac{1}{x}\right)}{2 c_{3}^{2}}+c_{4} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{\mathrm{e}^{-2 c_{2}}\left(\mathrm{e}^{2 c_{2}} c_{1} c_{3}^{2} x+\frac{1}{x}\right)}{2 c_{3}^{2}}+c_{4}
$$

Verified OK.

### 1.5.1 Maple step by step solution

Let's solve

$$
y^{\prime \prime \prime} x+3 y^{\prime \prime}=0
$$

- Highest derivative means the order of the ODE is 3
$y^{\prime \prime \prime}$
- Isolate 3rd derivative
$y^{\prime \prime \prime}=-\frac{3 y^{\prime \prime}}{x}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear
$y^{\prime \prime \prime}+\frac{3 y^{\prime \prime}}{x}=0$
- Multiply by denominators of the ODE
$y^{\prime \prime \prime} x+3 y^{\prime \prime}=0$
- Make a change of variables
$t=\ln (x)$
Substitute the change of variables back into the ODE
- Calculate the 1st derivative of y with respect to x , using the chain rule $y^{\prime}=\left(\frac{d}{d t} y(t)\right) t^{\prime}(x)$
- Compute derivative
$y^{\prime}=\frac{\frac{d}{d} y(t)}{x}$
- Calculate the 2nd derivative of y with respect to x , using the chain rule $y^{\prime \prime}=\left(\frac{d^{2}}{d t^{2}} y(t)\right) t^{\prime}(x)^{2}+t^{\prime \prime}(x)\left(\frac{d}{d t} y(t)\right)$
- Compute derivative
$y^{\prime \prime}=\frac{\frac{d^{2}}{d t^{2}} y(t)}{x^{2}}-\frac{\frac{d}{d t} y(t)}{x^{2}}$
- Calculate the 3 rd derivative of y with respect to x , using the chain rule

$$
y^{\prime \prime \prime}=\left(\frac{d^{3}}{d t^{3}} y(t)\right) t^{\prime}(x)^{3}+3 t^{\prime}(x) t^{\prime \prime}(x)\left(\frac{d^{2}}{d t^{2}} y(t)\right)+t^{\prime \prime \prime}(x)\left(\frac{d}{d t} y(t)\right)
$$

- Compute derivative
$y^{\prime \prime \prime}=\frac{d^{3}}{d t^{3}} y(t)-\frac{3\left(\frac{d^{2}}{d t^{2}} y(t)\right)}{x^{3}}+\frac{2\left(\frac{d}{d t} y(t)\right)}{x^{3}}$
Substitute the change of variables back into the ODE

$$
\left(\frac{\frac{d^{3}}{d t^{3}} y(t)}{x^{3}}-\frac{3\left(\frac{d^{2}}{d t^{2}} y(t)\right)}{x^{3}}+\frac{2\left(\frac{d}{d t} y(t)\right)}{x^{3}}\right) x+\frac{3\left(\frac{d^{2}}{d t^{2}} y(t)\right)}{x^{2}}-\frac{3\left(\frac{d}{d t} y(t)\right)}{x^{2}}=0
$$

- Simplify

$$
\frac{\frac{d^{3}}{d t^{3}} y(t)-\frac{d}{d t} y(t)}{x^{2}}=0
$$

- Isolate 3rd derivative

$$
\frac{d^{3}}{d t^{3}} y(t)=\frac{d}{d t} y(t)
$$

- Group terms with $y(t)$ on the lhs of the ODE and the rest on the rhs of the ODE; ODE is lin $\frac{d^{3}}{d t^{3}} y(t)-\frac{d}{d t} y(t)=0$

Convert linear ODE into a system of first order ODEs

- Define new variable $y_{1}(t)$

$$
y_{1}(t)=y(t)
$$

- Define new variable $y_{2}(t)$

$$
y_{2}(t)=\frac{d}{d t} y(t)
$$

- Define new variable $y_{3}(t)$

$$
y_{3}(t)=\frac{d^{2}}{d t^{2}} y(t)
$$

- Isolate for $\frac{d}{d t} y_{3}(t)$ using original ODE

$$
\frac{d}{d t} y_{3}(t)=y_{2}(t)
$$

Convert linear ODE into a system of first order ODEs

$$
\left[y_{2}(t)=\frac{d}{d t} y_{1}(t), y_{3}(t)=\frac{d}{d t} y_{2}(t), \frac{d}{d t} y_{3}(t)=y_{2}(t)\right]
$$

- Define vector

$$
\vec{y}(t)=\left[\begin{array}{l}
y_{1}(t) \\
y_{2}(t) \\
y_{3}(t)
\end{array}\right]
$$

- System to solve

$$
\frac{d}{d t} \vec{y}(t)=\left[\begin{array}{lll}
0 & 1 & 0 \\
0 & 0 & 1 \\
0 & 1 & 0
\end{array}\right] \cdot \vec{y}(t)
$$

- Define the coefficient matrix

$$
A=\left[\begin{array}{lll}
0 & 1 & 0 \\
0 & 0 & 1 \\
0 & 1 & 0
\end{array}\right]
$$

- Rewrite the system as
$\frac{d}{d t} \vec{y}(t)=A \cdot \vec{y}(t)$
- To solve the system, find the eigenvalues and eigenvectors of $A$
- $\quad$ Eigenpairs of $A$

$$
\left[\left[-1,\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]\right],\left[0,\left[\begin{array}{l}
1 \\
0 \\
0
\end{array}\right]\right],\left[1,\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]\right]\right]
$$

- Consider eigenpair

$$
\left[-1,\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]\right]
$$

- Solution to homogeneous system from eigenpair

$$
\vec{y}_{1}=\mathrm{e}^{-t} \cdot\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]
$$

- Consider eigenpair

$$
\left[0,\left[\begin{array}{l}
1 \\
0 \\
0
\end{array}\right]\right]
$$

- Solution to homogeneous system from eigenpair

$$
\vec{y}_{2}=\left[\begin{array}{l}
1 \\
0 \\
0
\end{array}\right]
$$

- Consider eigenpair

$$
\left[1,\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]\right]
$$

- $\quad$ Solution to homogeneous system from eigenpair

$$
\vec{y}_{3}=\mathrm{e}^{t} \cdot\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]
$$

- General solution to the system of ODEs

$$
\vec{y}=c_{1} \vec{y}_{1}+c_{2} \vec{y}_{2}+c_{3} \vec{y}_{3}
$$

- $\quad$ Substitute solutions into the general solution

$$
\vec{y}=c_{1} \mathrm{e}^{-t} \cdot\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]+c_{3} \mathrm{e}^{t} \cdot\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]+\left[\begin{array}{c}
c_{2} \\
0 \\
0
\end{array}\right]
$$

- First component of the vector is the solution to the ODE
$y(t)=c_{1} \mathrm{e}^{-t}+c_{3} \mathrm{e}^{t}+c_{2}$
- Change variables back using $t=\ln (x)$

$$
y=\frac{c_{1}}{x}+c_{3} x+c_{2}
$$

Maple trace

```
`Methods for third order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 14
dsolve(diff( $y(x), x \$ 3)+3 / x * \operatorname{diff}(y(x), x \$ 2)=0, y(x)$, singsol=all)

$$
y(x)=c_{1}+\frac{c_{2}}{x}+c_{3} x
$$

$\checkmark$ Solution by Mathematica
Time used: 0.042 (sec). Leaf size: 21
DSolve[y'''[x]+3/x*y''[x]==0,y[x],x,IncludeSingularSolutions -> True]

$$
y(x) \rightarrow \frac{c_{1}}{2 x}+c_{3} x+c_{2}
$$

## 1.6 problem 6

1.6.1 Solving as second order linear constant coeff ode . . . . . . . . 50
$\begin{array}{ll}\text { 1.6.2 } & \text { Solving as linear second order ode solved by an integrating factor } \\ & \text { ode . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 53 }\end{array}$
1.6.3 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 54
1.6.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 58

Internal problem ID [12423]
Internal file name [OUTPUT/11075_Monday_October_16_2023_09_47_11_PM_92580685/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 6.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by__an_integrating_factor"

Maple gives the following as the ode type
[[_2nd_order, _with_linear_symmetries]]

$$
y^{\prime \prime}-2 k y^{\prime}+y k^{2}=\mathrm{e}^{x}
$$

### 1.6.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
$$

Where $A=1, B=-2 k, C=k^{2}, f(x)=\mathrm{e}^{x}$. Let the solution be

$$
y=y_{h}+y_{p}
$$

Where $y_{h}$ is the solution to the homogeneous ODE $A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0$, and $y_{p}$ is a particular solution to the non-homogeneous ODE $A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)$. $y_{h}$ is the solution to

$$
y^{\prime \prime}-2 k y^{\prime}+y k^{2}=0
$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
$$

Where in the above $A=1, B=-2 k, C=k^{2}$. Let the solution be $y=e^{\lambda x}$. Substituting this into the ODE gives

$$
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-2 k \lambda \mathrm{e}^{\lambda x}+k^{2} \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
$$

Since exponential function is never zero, then dividing $\operatorname{Eq}(2)$ throughout by $e^{\lambda x}$ gives

$$
\begin{equation*}
k^{2}-2 k \lambda+\lambda^{2}=0 \tag{2}
\end{equation*}
$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula

$$
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
$$

Substituting $A=1, B=-2 k, C=k^{2}$ into the above gives

$$
\begin{aligned}
\lambda_{1,2} & =\frac{2 k}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(-2 k)^{2}-(4)(1)\left(k^{2}\right)} \\
& =k
\end{aligned}
$$

Hence this is the case of a double root $\lambda_{1,2}=-k$. Therefore the solution is

$$
\begin{equation*}
y=c_{1} \mathrm{e}^{k x}+c_{2} x \mathrm{e}^{k x} \tag{1}
\end{equation*}
$$

Therefore the homogeneous solution $y_{h}$ is

$$
y_{h}=c_{1} \mathrm{e}^{k x}+c_{2} x \mathrm{e}^{k x}
$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$
\mathrm{e}^{x}
$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$
\left[\left\{\mathrm{e}^{x}\right\}\right]
$$

While the set of the basis functions for the homogeneous solution found earlier is

$$
\left\{x \mathrm{e}^{k x}, \mathrm{e}^{k x}\right\}
$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$
y_{p}=A_{1} \mathrm{e}^{x}
$$

The unknowns $\left\{A_{1}\right\}$ are found by substituting the above trial solution $y_{p}$ into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$
A_{1} \mathrm{e}^{x}-2 k A_{1} \mathrm{e}^{x}+A_{1} \mathrm{e}^{x} k^{2}=\mathrm{e}^{x}
$$

Solving for the unknowns by comparing coefficients results in

$$
\left[A_{1}=\frac{1}{(k-1)^{2}}\right]
$$

Substituting the above back in the above trial solution $y_{p}$, gives the particular solution

$$
y_{p}=\frac{\mathrm{e}^{x}}{(k-1)^{2}}
$$

Therefore the general solution is

$$
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{k x}+c_{2} x \mathrm{e}^{k x}\right)+\left(\frac{\mathrm{e}^{x}}{(k-1)^{2}}\right)
\end{aligned}
$$

Which simplifies to

$$
y=\mathrm{e}^{k x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(k-1)^{2}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{k x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(k-1)^{2}} \tag{1}
\end{equation*}
$$

## Verification of solutions

$$
y=\mathrm{e}^{k x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(k-1)^{2}}
$$

Verified OK.

### 1.6.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$
y^{\prime \prime}+p(x) y^{\prime}+\frac{\left(p(x)^{2}+p^{\prime}(x)\right) y}{2}=f(x)
$$

Where $p(x)=-2 k$. Therefore, there is an integrating factor given by

$$
\begin{aligned}
M(x) & =e^{\frac{1}{2} \int p d x} \\
& =e^{\int-2 k d x} \\
& =\mathrm{e}^{-k x}
\end{aligned}
$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$
\begin{aligned}
(M(x) y)^{\prime \prime} & =\mathrm{e}^{-k x} \mathrm{e}^{x} \\
\left(\mathrm{e}^{-k x} y\right)^{\prime \prime} & =\mathrm{e}^{-k x} \mathrm{e}^{x}
\end{aligned}
$$

Integrating once gives

$$
\left(\mathrm{e}^{-k x} y\right)^{\prime}=-\frac{\mathrm{e}^{-(k-1) x}}{k-1}+c_{1}
$$

Integrating again gives

$$
\left(\mathrm{e}^{-k x} y\right)=\frac{\mathrm{e}^{-(k-1) x}+x\left(k^{2}-2 k+1\right) c_{1}}{(k-1)^{2}}+c_{2}
$$

Hence the solution is

$$
y=\frac{\frac{\mathrm{e}^{-(k-1) x}+x\left(k^{2}-2 k+1\right) c_{1}}{(k-1)^{2}}+c_{2}}{\mathrm{e}^{-k x}}
$$

Or

$$
y=c_{2} \mathrm{e}^{k x}+\left(\frac{k^{2} x \mathrm{e}^{k x}}{(k-1)^{2}}-\frac{2 k x \mathrm{e}^{k x}}{(k-1)^{2}}+\frac{x \mathrm{e}^{k x}}{(k-1)^{2}}\right) c_{1}+\frac{\mathrm{e}^{x}}{(k-1)^{2}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{2} \mathrm{e}^{k x}+\left(\frac{k^{2} x \mathrm{e}^{k x}}{(k-1)^{2}}-\frac{2 k x \mathrm{e}^{k x}}{(k-1)^{2}}+\frac{x \mathrm{e}^{k x}}{(k-1)^{2}}\right) c_{1}+\frac{\mathrm{e}^{x}}{(k-1)^{2}} \tag{1}
\end{equation*}
$$

## Verification of solutions

$$
y=c_{2} \mathrm{e}^{k x}+\left(\frac{k^{2} x \mathrm{e}^{k x}}{(k-1)^{2}}-\frac{2 k x \mathrm{e}^{k x}}{(k-1)^{2}}+\frac{x \mathrm{e}^{k x}}{(k-1)^{2}}\right) c_{1}+\frac{\mathrm{e}^{x}}{(k-1)^{2}}
$$

Verified OK.

### 1.6.3 Solving using Kovacic algorithm

Writing the ode as

$$
\begin{align*}
y^{\prime \prime}-2 k y^{\prime}+y k^{2} & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
$$

Comparing (1) and (2) shows that

$$
\begin{align*}
& A=1 \\
& B=-2 k  \tag{3}\\
& C=k^{2}
\end{align*}
$$

Applying the Liouville transformation on the dependent variable gives

$$
z(x)=y e^{\int \frac{B}{2 A} d x}
$$

Then (2) becomes

$$
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
$$

Where $r$ is given by

$$
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
$$

Substituting the values of $A, B, C$ from (3) in the above and simplifying gives

$$
\begin{equation*}
r=\frac{0}{1} \tag{6}
\end{equation*}
$$

Comparing the above to (5) shows that

$$
\begin{aligned}
& s=0 \\
& t=1
\end{aligned}
$$

Therefore eq. (4) becomes

$$
\begin{equation*}
z^{\prime \prime}(x)=0 \tag{7}
\end{equation*}
$$

Equation (7) is now solved. After finding $z(x)$ then $y$ is found using the inverse transformation

$$
y=z(x) e^{-\int \frac{B}{2 A} d x}
$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of $r$ and the order of $r$ at $\infty$. The following table summarizes these cases.

| Case | Allowed pole order for $r$ | Allowed value for $\mathcal{O}(\infty)$ |
| :--- | :--- | :--- |
| 1 | $\{0,1,2,4,6,8, \cdots\}$ | $\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}$ |
| 2 | Need to have at least one pole that <br> is either order 2 or odd order greater <br> than 2. Any other pole order is <br> allowed as long as the above condi- | no condition |
| tion is satisfied. Hence the following |  |  |
| set of pole orders are all allowed. |  |  |
| $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}$. |  |  |
| 3 | $\{1,2\}$ | $\{2,3,4,5,6,7, \cdots\}$ |

Table 5: Necessary conditions for each Kovacic case

The order of $r$ at $\infty$ is the degree of $t$ minus the degree of $s$. Therefore

$$
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0--\infty \\
& =\infty
\end{aligned}
$$

There are no poles in $r$. Therefore the set of poles $\Gamma$ is empty. Since there is no odd order pole larger than 2 and the order at $\infty$ is infinity then the necessary conditions for case one are met. Therefore

$$
L=[1]
$$

Since $r=0$ is not a function of $x$, then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z^{\prime \prime}=r z$ as one solution is

$$
z_{1}(x)=1
$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in $y$ is found from

$$
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{2 k}{1} d x} \\
& =z_{1} e^{k x} \\
& =z_{1}\left(\mathrm{e}^{k x}\right)
\end{aligned}
$$

Which simplifies to

$$
y_{1}=\mathrm{e}^{k x}
$$

The second solution $y_{2}$ to the original ode is found using reduction of order

$$
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
$$

Substituting gives

$$
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-2 k}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{2 k x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}(x)
\end{aligned}
$$

Therefore the solution is

$$
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{k x}\right)+c_{2}\left(\mathrm{e}^{k x}(x)\right)
\end{aligned}
$$

This is second order nonhomogeneous ODE. Let the solution be

$$
y=y_{h}+y_{p}
$$

Where $y_{h}$ is the solution to the homogeneous ODE $A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0$, and $y_{p}$ is a particular solution to the nonhomogeneous ODE $A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)$. $y_{h}$ is the solution to

$$
y^{\prime \prime}-2 k y^{\prime}+y k^{2}=0
$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$
y_{h}=c_{1} \mathrm{e}^{k x}+c_{2} x \mathrm{e}^{k x}
$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$
\mathrm{e}^{x}
$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$
\left[\left\{\mathrm{e}^{x}\right\}\right]
$$

While the set of the basis functions for the homogeneous solution found earlier is

$$
\left\{x \mathrm{e}^{k x}, \mathrm{e}^{k x}\right\}
$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$
y_{p}=A_{1} \mathrm{e}^{x}
$$

The unknowns $\left\{A_{1}\right\}$ are found by substituting the above trial solution $y_{p}$ into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$
A_{1} \mathrm{e}^{x}-2 k A_{1} \mathrm{e}^{x}+A_{1} \mathrm{e}^{x} k^{2}=\mathrm{e}^{x}
$$

Solving for the unknowns by comparing coefficients results in

$$
\left[A_{1}=\frac{1}{(k-1)^{2}}\right]
$$

Substituting the above back in the above trial solution $y_{p}$, gives the particular solution

$$
y_{p}=\frac{\mathrm{e}^{x}}{(k-1)^{2}}
$$

Therefore the general solution is

$$
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{k x}+c_{2} x \mathrm{e}^{k x}\right)+\left(\frac{\mathrm{e}^{x}}{(k-1)^{2}}\right)
\end{aligned}
$$

Which simplifies to

$$
y=\mathrm{e}^{k x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(k-1)^{2}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{k x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(k-1)^{2}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\mathrm{e}^{k x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(k-1)^{2}}
$$

Verified OK.

### 1.6.4 Maple step by step solution

Let's solve
$y^{\prime \prime}-2 k y^{\prime}+y k^{2}=\mathrm{e}^{x}$

- Highest derivative means the order of the ODE is 2
$y^{\prime \prime}$
- Characteristic polynomial of homogeneous ODE
$k^{2}-2 k r+r^{2}=0$
- Factor the characteristic polynomial
$(k-r)^{2}=0$
- Root of the characteristic polynomial
$r=k$
- $\quad 1$ st solution of the homogeneous ODE
$y_{1}(x)=\mathrm{e}^{k x}$
- Repeated root, multiply $y_{1}(x)$ by $x$ to ensure linear independence $y_{2}(x)=x \mathrm{e}^{k x}$
- General solution of the ODE
$y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)$
- Substitute in solutions of the homogeneous ODE

$$
y=c_{1} \mathrm{e}^{k x}+c_{2} x \mathrm{e}^{k x}+y_{p}(x)
$$

Find a particular solution $y_{p}(x)$ of the ODE

- Use variation of parameters to find $y_{p}$ here $f(x)$ is the forcing function

$$
\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=\mathrm{e}^{x}\right]
$$

- Wronskian of solutions of the homogeneous equation

$$
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
\mathrm{e}^{k x} & x \mathrm{e}^{k x} \\
k \mathrm{e}^{k x} & \mathrm{e}^{k x}+k x \mathrm{e}^{k x}
\end{array}\right]
$$

- Compute Wronskian

$$
W\left(y_{1}(x), y_{2}(x)\right)=\mathrm{e}^{2 k x}
$$

- Substitute functions into equation for $y_{p}(x)$

$$
y_{p}(x)=\mathrm{e}^{k x}\left(-\left(\int x \mathrm{e}^{-(k-1) x} d x\right)+\left(\int \mathrm{e}^{-(k-1) x} d x\right) x\right)
$$

- Compute integrals

$$
y_{p}(x)=\frac{\mathrm{e}^{x}}{(k-1)^{2}}
$$

- Substitute particular solution into general solution to ODE

$$
y=c_{1} \mathrm{e}^{k x}+c_{2} x \mathrm{e}^{k x}+\frac{\mathrm{e}^{x}}{(k-1)^{2}}
$$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 28
dsolve(diff $(y(x), x \$ 2)-2 * k * \operatorname{diff}(y(x), x)+k^{\wedge} 2 * y(x)=\exp (x), y(x)$, singsol=all)

$$
y(x)=\frac{(-1+k)^{2}\left(c_{1} x+c_{2}\right) \mathrm{e}^{k x}+\mathrm{e}^{x}}{(-1+k)^{2}}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.221 (sec). Leaf size: 28
DSolve [y' ' $[\mathrm{x}]-2 * \mathrm{k} * \mathrm{y}$ ' $[\mathrm{x}]+\mathrm{k}^{\wedge} 2 * \mathrm{y}[\mathrm{x}]==\operatorname{Exp}[\mathrm{x}], \mathrm{y}[\mathrm{x}], \mathrm{x}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow \frac{e^{x}}{(k-1)^{2}}+\left(c_{2} x+c_{1}\right) e^{k x}
$$

## 1.7 problem 7

1.7.1 Solving as second order change of variable on $x$ method 2 ode . 61
1.7.2 Solving as second order change of variable on $x$ method 1 ode . 64
1.7.3 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 66
1.7.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 72

Internal problem ID [12424]
Internal file name [OUTPUT/11076_Monday_October_16_2023_09_47_14_PM_53640396/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 7 .
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_change__of_variable_on_x_method_1", "second_order_change_of__variable_on_x_method_2"

Maple gives the following as the ode type

```
[_Gegenbauer, [_2nd_order, _linear, ` _with_symmetry_[0,F(x)]`]]
```

$$
\left(-x^{2}+1\right) y^{\prime \prime}-y^{\prime} x-y a^{2}=0
$$

### 1.7.1 Solving as second order change of variable on $x$ method 2 ode

In normal form the ode

$$
\begin{equation*}
\left(-x^{2}+1\right) y^{\prime \prime}-y^{\prime} x-y a^{2}=0 \tag{1}
\end{equation*}
$$

Becomes

$$
\begin{equation*}
y^{\prime \prime}+p(x) y^{\prime}+q(x) y=0 \tag{2}
\end{equation*}
$$

Where

$$
\begin{aligned}
p(x) & =\frac{x}{x^{2}-1} \\
q(x) & =\frac{a^{2}}{x^{2}-1}
\end{aligned}
$$

Applying change of variables $\tau=g(x)$ to (2) gives

$$
\begin{equation*}
\frac{d^{2}}{d \tau^{2}} y(\tau)+p_{1}\left(\frac{d}{d \tau} y(\tau)\right)+q_{1} y(\tau)=0 \tag{3}
\end{equation*}
$$

Where $\tau$ is the new independent variable, and

$$
\begin{align*}
& p_{1}(\tau)=\frac{\tau^{\prime \prime}(x)+p(x) \tau^{\prime}(x)}{\tau^{\prime}(x)^{2}}  \tag{4}\\
& q_{1}(\tau)=\frac{q(x)}{\tau^{\prime}(x)^{2}} \tag{5}
\end{align*}
$$

Let $p_{1}=0 . \mathrm{Eq}(4)$ simplifies to

$$
\tau^{\prime \prime}(x)+p(x) \tau^{\prime}(x)=0
$$

This ode is solved resulting in

$$
\begin{align*}
\tau & =\int \mathrm{e}^{-\left(\int p(x) d x\right)} d x \\
& =\int \mathrm{e}^{-\left(\int \frac{x}{x^{2}-1} d x\right)} d x \\
& =\int e^{-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}} d x \\
& =\int \frac{1}{\sqrt{x-1} \sqrt{x+1}} d x \\
& =\frac{\sqrt{(x-1)(x+1)} \ln \left(x+\sqrt{x^{2}-1}\right)}{\sqrt{x-1} \sqrt{x+1}} \tag{6}
\end{align*}
$$

Using (6) to evaluate $q_{1}$ from (5) gives

$$
\begin{align*}
q_{1}(\tau) & =\frac{q(x)}{\tau^{\prime}(x)^{2}} \\
& =\frac{\frac{a^{2}}{x^{2}-1}}{(x-1)(x+1)} \\
& =a^{2} \tag{7}
\end{align*}
$$

Substituting the above in (3) and noting that now $p_{1}=0$ results in

$$
\begin{gathered}
\frac{d^{2}}{d \tau^{2}} y(\tau)+q_{1} y(\tau)=0 \\
\frac{d^{2}}{d \tau^{2}} y(\tau)+a^{2} y(\tau)=0
\end{gathered}
$$

The above ode is now solved for $y(\tau)$.This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$
A y^{\prime \prime}(\tau)+B y^{\prime}(\tau)+C y(\tau)=0
$$

Where in the above $A=1, B=0, C=a^{2}$. Let the solution be $y(\tau)=e^{\lambda \tau}$. Substituting this into the ODE gives

$$
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda \tau}+a^{2} \mathrm{e}^{\lambda \tau}=0 \tag{1}
\end{equation*}
$$

Since exponential function is never zero, then dividing $\operatorname{Eq}(2)$ throughout by $e^{\lambda \tau}$ gives

$$
\begin{equation*}
a^{2}+\lambda^{2}=0 \tag{2}
\end{equation*}
$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula

$$
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
$$

Substituting $A=1, B=0, C=a^{2}$ into the above gives

$$
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)\left(a^{2}\right)} \\
& = \pm \sqrt{-a^{2}}
\end{aligned}
$$

Hence

$$
\begin{aligned}
& \lambda_{1}=+\sqrt{-a^{2}} \\
& \lambda_{2}=-\sqrt{-a^{2}}
\end{aligned}
$$

Which simplifies to

$$
\begin{aligned}
& \lambda_{1}=\sqrt{-a^{2}} \\
& \lambda_{2}=-\sqrt{-a^{2}}
\end{aligned}
$$

Since roots are real and distinct, then the solution is

$$
\begin{aligned}
& y(\tau)=c_{1} e^{\lambda_{1} \tau}+c_{2} e^{\lambda_{2} \tau} \\
& y(\tau)=c_{1} e^{\left(\sqrt{-a^{2}}\right) \tau}+c_{2} e^{\left(-\sqrt{-a^{2}}\right) \tau}
\end{aligned}
$$

Or

$$
y(\tau)=c_{1} \mathrm{e}^{\sqrt{-a^{2}} \tau}+c_{2} \mathrm{e}^{-\sqrt{-a^{2}} \tau}
$$

The above solution is now transformed back to $y$ using (6) which results in

$$
y=c_{1}\left(x+\sqrt{x^{2}-1}\right)^{\frac{\sqrt{-a^{2}} \sqrt{x^{2}-1}}{\sqrt{x-1} \sqrt{x+1}}}+c_{2}\left(x+\sqrt{x^{2}-1}\right)^{-\frac{\sqrt{-a^{2}} \sqrt{x^{2}-1}}{\sqrt{x-1} \sqrt{x+1}}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1}\left(x+\sqrt{x^{2}-1}\right)^{\frac{\sqrt{-a^{2}} \sqrt{x^{2}-1}}{\sqrt{x-1} \sqrt{x+1}}}+c_{2}\left(x+\sqrt{x^{2}-1}\right)^{-\frac{\sqrt{-a^{2}} \sqrt{x^{2}-1}}{\sqrt{x-1} \sqrt{x+1}}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=c_{1}\left(x+\sqrt{x^{2}-1}\right)^{\frac{\sqrt{-a^{2}} \sqrt{x^{2}-1}}{\sqrt{x}-1} \sqrt{x+1}}+c_{2}\left(x+\sqrt{x^{2}-1}\right)^{-\frac{\sqrt{-a^{2}} \sqrt{x^{2}-1}}{\sqrt{x-1} \sqrt{x+1}}}
$$

Verified OK.

### 1.7.2 Solving as second order change of variable on $x$ method 1 ode

In normal form the ode

$$
\begin{equation*}
\left(-x^{2}+1\right) y^{\prime \prime}-y^{\prime} x-y a^{2}=0 \tag{1}
\end{equation*}
$$

Becomes

$$
\begin{equation*}
y^{\prime \prime}+p(x) y^{\prime}+q(x) y=0 \tag{2}
\end{equation*}
$$

Where

$$
\begin{aligned}
p(x) & =\frac{x}{x^{2}-1} \\
q(x) & =\frac{a^{2}}{x^{2}-1}
\end{aligned}
$$

Applying change of variables $\tau=g(x)$ to (2) results

$$
\begin{equation*}
\frac{d^{2}}{d \tau^{2}} y(\tau)+p_{1}\left(\frac{d}{d \tau} y(\tau)\right)+q_{1} y(\tau)=0 \tag{3}
\end{equation*}
$$

Where $\tau$ is the new independent variable, and

$$
\begin{align*}
& p_{1}(\tau)=\frac{\tau^{\prime \prime}(x)+p(x) \tau^{\prime}(x)}{\tau^{\prime}(x)^{2}}  \tag{4}\\
& q_{1}(\tau)=\frac{q(x)}{\tau^{\prime}(x)^{2}} \tag{5}
\end{align*}
$$

Let $q_{1}=c^{2}$ where $c$ is some constant. Therefore from (5)

$$
\begin{align*}
\tau^{\prime} & =\frac{1}{c} \sqrt{q} \\
& =\frac{\sqrt{\frac{a^{2}}{x^{2}-1}}}{c}  \tag{6}\\
\tau^{\prime \prime} & =-\frac{a^{2} x}{c \sqrt{\frac{a^{2}}{x^{2}-1}}\left(x^{2}-1\right)^{2}}
\end{align*}
$$

Substituting the above into (4) results in

$$
\begin{aligned}
p_{1}(\tau) & =\frac{\tau^{\prime \prime}(x)+p(x) \tau^{\prime}(x)}{\tau^{\prime}(x)^{2}} \\
& =\frac{-\frac{a^{2} x}{c \sqrt{\frac{a^{2}}{x^{2}-1}}\left(x^{2}-1\right)^{2}}+\frac{x}{x^{2}-1} \frac{\sqrt{\frac{a^{2}}{x^{2}-1}}}{c}}{\left(\frac{\sqrt{\frac{a^{2}}{x^{2}-1}}}{c}\right)^{2}} \\
& =0
\end{aligned}
$$

Therefore ode (3) now becomes

$$
\begin{align*}
y(\tau)^{\prime \prime}+p_{1} y(\tau)^{\prime}+q_{1} y(\tau) & =0 \\
\frac{d^{2}}{d \tau^{2}} y(\tau)+c^{2} y(\tau) & =0 \tag{7}
\end{align*}
$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$
y(\tau)=c_{1} \cos (c \tau)+c_{2} \sin (c \tau)
$$

Now from (6)

$$
\begin{aligned}
\tau & =\int \frac{1}{c} \sqrt{q} d x \\
& =\frac{\int \sqrt{\frac{a^{2}}{x^{2}-1}} d x}{c} \\
& =\frac{\sqrt{\frac{a^{2}}{x^{2}-1}} \sqrt{x^{2}-1} \ln \left(x+\sqrt{x^{2}-1}\right)}{c}
\end{aligned}
$$

Substituting the above into the solution obtained gives

$$
\begin{aligned}
y= & c_{1} \cos \left(a \sqrt{\frac{1}{x^{2}-1}} \sqrt{x^{2}-1} \ln \left(x+\sqrt{x^{2}-1}\right)\right) \\
& +c_{2} \sin \left(a \sqrt{\frac{1}{x^{2}-1}} \sqrt{x^{2}-1} \ln \left(x+\sqrt{x^{2}-1}\right)\right)
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
y= & c_{1} \cos \left(a \sqrt{\frac{1}{x^{2}-1}} \sqrt{x^{2}-1} \ln \left(x+\sqrt{x^{2}-1}\right)\right)  \tag{1}\\
& +c_{2} \sin \left(a \sqrt{\frac{1}{x^{2}-1}} \sqrt{x^{2}-1} \ln \left(x+\sqrt{x^{2}-1}\right)\right)
\end{align*}
$$

Verification of solutions

$$
\begin{aligned}
y= & c_{1} \cos \left(a \sqrt{\frac{1}{x^{2}-1}} \sqrt{x^{2}-1} \ln \left(x+\sqrt{x^{2}-1}\right)\right) \\
& +c_{2} \sin \left(a \sqrt{\frac{1}{x^{2}-1}} \sqrt{x^{2}-1} \ln \left(x+\sqrt{x^{2}-1}\right)\right)
\end{aligned}
$$

Verified OK.

### 1.7.3 Solving using Kovacic algorithm

Writing the ode as

$$
\begin{align*}
\left(-x^{2}+1\right) y^{\prime \prime}-y^{\prime} x-y a^{2} & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
$$

Comparing (1) and (2) shows that

$$
\begin{align*}
& A=-x^{2}+1 \\
& B=-x  \tag{3}\\
& C=-a^{2}
\end{align*}
$$

Applying the Liouville transformation on the dependent variable gives

$$
z(x)=y e^{\int \frac{B}{2 A} d x}
$$

Then (2) becomes

$$
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
$$

Where $r$ is given by

$$
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
$$

Substituting the values of $A, B, C$ from (3) in the above and simplifying gives

$$
\begin{equation*}
r=\frac{-4 a^{2} x^{2}+4 a^{2}-x^{2}-2}{4\left(x^{2}-1\right)^{2}} \tag{6}
\end{equation*}
$$

Comparing the above to (5) shows that

$$
\begin{aligned}
& s=-4 a^{2} x^{2}+4 a^{2}-x^{2}-2 \\
& t=4\left(x^{2}-1\right)^{2}
\end{aligned}
$$

Therefore eq. (4) becomes

$$
\begin{equation*}
z^{\prime \prime}(x)=\left(\frac{-4 a^{2} x^{2}+4 a^{2}-x^{2}-2}{4\left(x^{2}-1\right)^{2}}\right) z(x) \tag{7}
\end{equation*}
$$

Equation (7) is now solved. After finding $z(x)$ then $y$ is found using the inverse transformation

$$
y=z(x) e^{-\int \frac{B}{2 A} d x}
$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of $r$ and the order of $r$ at $\infty$. The following table summarizes these cases.

| Case | Allowed pole order for $r$ | Allowed value for $\mathcal{O}(\infty)$ |
| :--- | :--- | :--- |
| 1 | $\{0,1,2,4,6,8, \cdots\}$ | $\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}$ |
| 2 | Need to have at least one pole that <br> is either order 2 or odd order greater <br> than 2. Any other pole order is <br> allowed as long as the above condi- <br> tion is satisfied. Hence the following <br> set of pole orders are all allowed. <br> $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}$. |  |
| 3 | $\{1,2\}$ | $\{2,3,4,5,6,7, \cdots\}$ |

Table 7: Necessary conditions for each Kovacic case

The order of $r$ at $\infty$ is the degree of $t$ minus the degree of $s$. Therefore

$$
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =4-2 \\
& =2
\end{aligned}
$$

The poles of $r$ in eq. (7) and the order of each pole are determined by solving for the roots of $t=4\left(x^{2}-1\right)^{2}$. There is a pole at $x=1$ of order 2 . There is a pole at $x=-1$ of order 2 . Since there is no odd order pole larger than 2 and the order at $\infty$ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at $\infty$ is 2 then the necessary conditions for case three are met. Therefore

$$
L=[1,2,4,6,12]
$$

Attempting to find a solution using case $n=1$.
Unable to find solution using case one
Attempting to find a solution using case $n=2$.
Looking at poles of order 2. The partial fractions decomposition of $r$ is

$$
r=-\frac{3}{16(x-1)^{2}}+\frac{\frac{1}{16}-\frac{a^{2}}{2}}{x-1}-\frac{3}{16(x+1)^{2}}+\frac{-\frac{1}{16}+\frac{a^{2}}{2}}{x+1}
$$

For the pole at $x=1$ let $b$ be the coefficient of $\frac{1}{(x-1)^{2}}$ in the partial fractions decomposition of $r$ given above. Therefore $b=-\frac{3}{16}$. Hence

$$
\begin{aligned}
E_{c} & =\{2,2+2 \sqrt{1+4 b}, 2-2 \sqrt{1+4 b}\} \\
& =\{1,2,3\}
\end{aligned}
$$

For the pole at $x=-1$ let $b$ be the coefficient of $\frac{1}{(x+1)^{2}}$ in the partial fractions decomposition of $r$ given above. Therefore $b=-\frac{3}{16}$. Hence

$$
\begin{aligned}
E_{c} & =\{2,2+2 \sqrt{1+4 b}, 2-2 \sqrt{1+4 b}\} \\
& =\{1,2,3\}
\end{aligned}
$$

Since the order of $r$ at $\infty$ is 2 then let $b$ be the coefficient of $\frac{1}{x^{2}}$ in the Laurent series expansion of $r$ at $\infty$. which can be found by dividing the leading coefficient of $s$ by the leading coefficient of $t$ from

$$
r=\frac{s}{t}=\frac{-4 a^{2} x^{2}+4 a^{2}-x^{2}-2}{4\left(x^{2}-1\right)^{2}}
$$

Since the $\operatorname{gcd}(s, t)=1$. This gives $b=-1$. Hence

$$
\begin{aligned}
E_{\infty} & =\{2,2+2 \sqrt{1+4 b}, 2-2 \sqrt{1+4 b}\} \\
& =\{2\}
\end{aligned}
$$

The following table summarizes the findings so far for poles and for the order of $r$ at $\infty$ for case 2 of Kovacic algorithm.

| pole $c$ location | pole order | $E_{c}$ |
| :---: | :---: | :---: |
| 1 | 2 | $\{1,2,3\}$ |
| -1 | 2 | $\{1,2,3\}$ |


| Order of $r$ at $\infty$ | $E_{\infty}$ |
| :---: | :---: |
| 2 | $\{2\}$ |

Using the family $\left\{e_{1}, e_{2}, \ldots, e_{\infty}\right\}$ given by

$$
e_{1}=1, e_{2}=1, e_{\infty}=2
$$

Gives a non negative integer $d$ (the degree of the polynomial $p(x)$ ), which is generated using

$$
\begin{aligned}
d & =\frac{1}{2}\left(e_{\infty}-\sum_{c \in \Gamma} e_{c}\right) \\
& =\frac{1}{2}(2-(1+(1))) \\
& =0
\end{aligned}
$$

We now form the following rational function

$$
\begin{aligned}
\theta & =\frac{1}{2} \sum_{c \in \Gamma} \frac{e_{c}}{x-c} \\
& =\frac{1}{2}\left(\frac{1}{(x-(1))}+\frac{1}{(x-(-1))}\right) \\
& =\frac{1}{2 x-2}+\frac{1}{2 x+2}
\end{aligned}
$$

Now we search for a monic polynomial $p(x)$ of degree $d=0$ such that

$$
\begin{equation*}
p^{\prime \prime \prime}+3 \theta p^{\prime \prime}+\left(3 \theta^{2}+3 \theta^{\prime}-4 r\right) p^{\prime}+\left(\theta^{\prime \prime}+3 \theta \theta^{\prime}+\theta^{3}-4 r \theta-2 r^{\prime}\right) p=0 \tag{1A}
\end{equation*}
$$

Since $d=0$, then letting

$$
\begin{equation*}
p=1 \tag{2A}
\end{equation*}
$$

Substituting $p$ and $\theta$ into Eq. (1A) gives

$$
0=0
$$

And solving for $p$ gives

$$
p=1
$$

Now that $p(x)$ is found let

$$
\begin{aligned}
\phi & =\theta+\frac{p^{\prime}}{p} \\
& =\frac{1}{2 x-2}+\frac{1}{2 x+2}
\end{aligned}
$$

Let $\omega$ be the solution of

$$
\omega^{2}-\phi \omega+\left(\frac{1}{2} \phi^{\prime}+\frac{1}{2} \phi^{2}-r\right)=0
$$

Substituting the values for $\phi$ and $r$ into the above equation gives

$$
w^{2}-\left(\frac{1}{2 x-2}+\frac{1}{2 x+2}\right) w+\frac{4 a^{2} x^{2}-4 a^{2}+x^{2}}{4\left(x^{2}-1\right)^{2}}=0
$$

Solving for $\omega$ gives

$$
\omega=\frac{x+2 a \sqrt{-x^{2}+1}}{2(x-1)(x+1)}
$$

Therefore the first solution to the ode $z^{\prime \prime}=r z$ is

$$
\begin{aligned}
z_{1}(x) & =e^{\int \omega d x} \\
& =\mathrm{e}^{\int \frac{x+2 a \sqrt{-x^{2}+1}}{2(x-1)(x+1)} d x} \\
& =\left(x^{2}-1\right)^{\frac{1}{4}} \mathrm{e}^{-a \arcsin (x)}
\end{aligned}
$$

The first solution to the original ode in $y$ is found from

$$
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{-x}{-x^{2}+1} d x} \\
& =z_{1} e^{-\frac{\ln (x-1)}{4}-\frac{\ln (x+1)}{4}} \\
& =z_{1}\left(\frac{1}{(x-1)^{\frac{1}{4}}(x+1)^{\frac{1}{4}}}\right)
\end{aligned}
$$

Which simplifies to

$$
y_{1}=\frac{\left(x^{2}-1\right)^{\frac{1}{4}} \mathrm{e}^{-a \arcsin (x)}}{(x-1)^{\frac{1}{4}}(x+1)^{\frac{1}{4}}}
$$

The second solution $y_{2}$ to the original ode is found using reduction of order

$$
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
$$

Substituting gives

$$
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-x}{-x^{2}+1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\int \frac{\mathrm{e}^{2 a \arcsin (x)}}{\sqrt{x^{2}-1}} d x\right)
\end{aligned}
$$

Therefore the solution is

$$
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\frac{\left(x^{2}-1\right)^{\frac{1}{4}} \mathrm{e}^{-a \arcsin (x)}}{(x-1)^{\frac{1}{4}}(x+1)^{\frac{1}{4}}}\right)+c_{2}\left(\frac{\left(x^{2}-1\right)^{\frac{1}{4}} \mathrm{e}^{-a \arcsin (x)}}{(x-1)^{\frac{1}{4}}(x+1)^{\frac{1}{4}}}\left(\int \frac{\mathrm{e}^{2 a \arcsin (x)}}{\sqrt{x^{2}-1}} d x\right)\right)
\end{aligned}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{c_{1}\left(x^{2}-1\right)^{\frac{1}{4}} \mathrm{e}^{-a \arcsin (x)}}{(x-1)^{\frac{1}{4}}(x+1)^{\frac{1}{4}}}+\frac{c_{2}\left(x^{2}-1\right)^{\frac{1}{4}} \mathrm{e}^{-a \arcsin (x)}\left(\int \frac{\mathrm{e}^{2 a \arcsin (x)}}{\sqrt{x^{2}-1}} d x\right)}{(x-1)^{\frac{1}{4}}(x+1)^{\frac{1}{4}}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{c_{1}\left(x^{2}-1\right)^{\frac{1}{4}} \mathrm{e}^{-a \arcsin (x)}}{(x-1)^{\frac{1}{4}}(x+1)^{\frac{1}{4}}}+\frac{c_{2}\left(x^{2}-1\right)^{\frac{1}{4}} \mathrm{e}^{-a \arcsin (x)}\left(\int \frac{\mathrm{e}^{2 a \arcsin (x)}}{\sqrt{x^{2}-1}} d x\right)}{(x-1)^{\frac{1}{4}}(x+1)^{\frac{1}{4}}}
$$

Verified OK.

### 1.7.4 Maple step by step solution

Let's solve

$$
\left(-x^{2}+1\right) y^{\prime \prime}-y^{\prime} x-y a^{2}=0
$$

- Highest derivative means the order of the ODE is 2
$y^{\prime \prime}$
- Isolate 2nd derivative

$$
y^{\prime \prime}=-\frac{x y^{\prime}}{x^{2}-1}-\frac{a^{2} y}{x^{2}-1}
$$

- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$
y^{\prime \prime}+\frac{x y^{\prime}}{x^{2}-1}+\frac{a^{2} y}{x^{2}-1}=0
$$

Check to see if $x_{0}$ is a regular singular point

- Define functions
$\left[P_{2}(x)=\frac{x}{x^{2}-1}, P_{3}(x)=\frac{a^{2}}{x^{2}-1}\right]$
- $(x+1) \cdot P_{2}(x)$ is analytic at $x=-1$
$\left.\left((x+1) \cdot P_{2}(x)\right)\right|_{x=-1}=\frac{1}{2}$
- $(x+1)^{2} \cdot P_{3}(x)$ is analytic at $x=-1$
$\left.\left((x+1)^{2} \cdot P_{3}(x)\right)\right|_{x=-1}=0$
- $x=-1$ is a regular singular point

Check to see if $x_{0}$ is a regular singular point
$x_{0}=-1$

- Multiply by denominators
$y^{\prime \prime}\left(x^{2}-1\right)+y^{\prime} x+y a^{2}=0$
- $\quad$ Change variables using $x=u-1$ so that the regular singular point is at $u=0$
$\left(u^{2}-2 u\right)\left(\frac{d^{2}}{d u^{2}} y(u)\right)+(u-1)\left(\frac{d}{d u} y(u)\right)+a^{2} y(u)=0$
- Assume series solution for $y(u)$
$y(u)=\sum_{k=0}^{\infty} a_{k} u^{k+r}$
Rewrite ODE with series expansions
- Convert $u^{m} \cdot\left(\frac{d}{d u} y(u)\right)$ to series expansion for $m=0 . .1$

$$
u^{m} \cdot\left(\frac{d}{d u} y(u)\right)=\sum_{k=0}^{\infty} a_{k}(k+r) u^{k+r-1+m}
$$

- Shift index using $k->k+1-m$

$$
u^{m} \cdot\left(\frac{d}{d u} y(u)\right)=\sum_{k=-1+m}^{\infty} a_{k+1-m}(k+1-m+r) u^{k+r}
$$

- Convert $u^{m} \cdot\left(\frac{d^{2}}{d u^{2}} y(u)\right)$ to series expansion for $m=1 . .2$

$$
u^{m} \cdot\left(\frac{d^{2}}{d u^{2}} y(u)\right)=\sum_{k=0}^{\infty} a_{k}(k+r)(k+r-1) u^{k+r-2+m}
$$

- Shift index using $k->k+2-m$
$u^{m} \cdot\left(\frac{d^{2}}{d u^{2}} y(u)\right)=\sum_{k=-2+m}^{\infty} a_{k+2-m}(k+2-m+r)(k+1-m+r) u^{k+r}$
Rewrite ODE with series expansions
$-a_{0} r(-1+2 r) u^{-1+r}+\left(\sum_{k=0}^{\infty}\left(-a_{k+1}(k+1+r)(2 k+1+2 r)+a_{k}\left(a^{2}+k^{2}+2 k r+r^{2}\right)\right) u^{k+r}\right)$
- $\quad a_{0}$ cannot be 0 by assumption, giving the indicial equation
$-r(-1+2 r)=0$
- Values of $r$ that satisfy the indicial equation
$r \in\left\{0, \frac{1}{2}\right\}$
- Each term in the series must be 0 , giving the recursion relation

$$
-2(k+1+r)\left(k+\frac{1}{2}+r\right) a_{k+1}+a_{k}\left(a^{2}+k^{2}+2 k r+r^{2}\right)=0
$$

- Recursion relation that defines series solution to ODE
$a_{k+1}=\frac{a_{k}\left(a^{2}+k^{2}+2 k r+r^{2}\right)}{(k+1+r)(2 k+1+2 r)}$
- Recursion relation for $r=0$
$a_{k+1}=\frac{a_{k}\left(a^{2}+k^{2}\right)}{(k+1)(2 k+1)}$
- $\quad$ Solution for $r=0$

$$
\left[y(u)=\sum_{k=0}^{\infty} a_{k} u^{k}, a_{k+1}=\frac{a_{k}\left(a^{2}+k^{2}\right)}{(k+1)(2 k+1)}\right]
$$

- $\quad$ Revert the change of variables $u=x+1$

$$
\left[y=\sum_{k=0}^{\infty} a_{k}(x+1)^{k}, a_{k+1}=\frac{a_{k}\left(a^{2}+k^{2}\right)}{(k+1)(2 k+1)}\right]
$$

- Recursion relation for $r=\frac{1}{2}$

$$
a_{k+1}=\frac{a_{k}\left(a^{2}+k^{2}+k+\frac{1}{4}\right)}{\left(k+\frac{3}{2}\right)(2 k+2)}
$$

- $\quad$ Solution for $r=\frac{1}{2}$

$$
\left[y(u)=\sum_{k=0}^{\infty} a_{k} u^{k+\frac{1}{2}}, a_{k+1}=\frac{a_{k}\left(a^{2}+k^{2}+k+\frac{1}{4}\right)}{\left(k+\frac{3}{2}\right)(2 k+2)}\right]
$$

- $\quad$ Revert the change of variables $u=x+1$

$$
\left[y=\sum_{k=0}^{\infty} a_{k}(x+1)^{k+\frac{1}{2}}, a_{k+1}=\frac{a_{k}\left(a^{2}+k^{2}+k+\frac{1}{4}\right)}{\left(k+\frac{3}{2}\right)(2 k+2)}\right]
$$

- Combine solutions and rename parameters

$$
\left[y=\left(\sum_{k=0}^{\infty} b_{k}(x+1)^{k}\right)+\left(\sum_{k=0}^{\infty} c_{k}(x+1)^{k+\frac{1}{2}}\right), b_{1+k}=\frac{b_{k}\left(a^{2}+k^{2}\right)}{(1+k)(2 k+1)}, c_{1+k}=\frac{c_{k}\left(a^{2}+k^{2}+k+\frac{1}{4}\right)}{\left(k+\frac{3}{2}\right)(2 k+2)}\right]
$$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
<- linear_1 successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 37
dsolve ( $\left(1-x^{\wedge} 2\right) * \operatorname{diff}(y(x), x \$ 2)-x * \operatorname{diff}(y(x), x)-a^{\wedge} 2 * y(x)=0, y(x), \quad$ singsol=all)

$$
y(x)=c_{1}\left(x+\sqrt{x^{2}-1}\right)^{i a}+c_{2}\left(x+\sqrt{x^{2}-1}\right)^{-i a}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.151 (sec). Leaf size: 89
DSolve[(1- $\left.x^{\wedge} 2\right) * y '$ '[x]-x*y'[x]-a^2*y[x]==0,y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
y(x) \rightarrow & c_{1} \cos \left(\frac{1}{2} a\left(\log \left(1-\frac{x}{\sqrt{x^{2}-1}}\right)-\log \left(\frac{x}{\sqrt{x^{2}-1}}+1\right)\right)\right) \\
& -c_{2} \sin \left(\frac{1}{2} a\left(\log \left(1-\frac{x}{\sqrt{x^{2}-1}}\right)-\log \left(\frac{x}{\sqrt{x^{2}-1}}+1\right)\right)\right)
\end{aligned}
$$

## 1.8 problem 8

1.8.1 Solving as second order integrable as is ode ..... [77
1.8.2 Solving as second order ode missing y ode ..... 78
1.8.3 Solving as second order ode non constant coeff transformation on B ode ..... 79
1.8.4 Solving as type second_order_integrable_as_is (not using ABC version) ..... 81
1.8.5 Solving using Kovacic algorithm ..... 82
1.8.6 Solving as exact linear second order ode ode ..... 85
1.8.7 Maple step by step solution ..... 87

Internal problem ID [12425]
Internal file name [OUTPUT/11077_Monday_October_16_2023_09_47_15_PM_7666333/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 8.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "exact linear second order ode", "second_order_integrable_as_is", "second_order_ode_missing_y", "second_order_ode__non_constant_coeff_transformation__on_B"

Maple gives the following as the ode type
[[_2nd_order, _missing_y]]

$$
y^{\prime \prime}+\frac{2 y^{\prime}}{x}=0
$$

### 1.8.1 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t $x$ gives

$$
\begin{gathered}
\int\left(y^{\prime \prime} x+2 y^{\prime}\right) d x=0 \\
y^{\prime} x+y=c_{1}
\end{gathered}
$$

Which is now solved for $y$. In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{-y+c_{1}}{x}
\end{aligned}
$$

Where $f(x)=\frac{1}{x}$ and $g(y)=-y+c_{1}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{-y+c_{1}} d y & =\frac{1}{x} d x \\
\int \frac{1}{-y+c_{1}} d y & =\int \frac{1}{x} d x \\
-\ln \left(-y+c_{1}\right) & =\ln (x)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\frac{1}{-y+c_{1}}=\mathrm{e}^{\ln (x)+c_{2}}
$$

Which simplifies to

$$
\frac{1}{-y+c_{1}}=c_{3} x
$$

Which simplifies to

$$
y=\frac{\left(c_{3} \mathrm{e}^{c_{2}} x c_{1}-1\right) \mathrm{e}^{-c_{2}}}{c_{3} x}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\left(c_{3} \mathrm{e}^{c_{2}} x c_{1}-1\right) \mathrm{e}^{-c_{2}}}{c_{3} x} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{\left(c_{3} \mathrm{e}^{c_{2}} x c_{1}-1\right) \mathrm{e}^{-c_{2}}}{c_{3} x}
$$

Verified OK.

### 1.8.2 Solving as second order ode missing y ode

This is second order ode with missing dependent variable $y$. Let

$$
p(x)=y^{\prime}
$$

Then

$$
p^{\prime}(x)=y^{\prime \prime}
$$

Hence the ode becomes

$$
p^{\prime}(x) x+2 p(x)=0
$$

Which is now solve for $p(x)$ as first order ode. In canonical form the ODE is

$$
\begin{aligned}
p^{\prime} & =F(x, p) \\
& =f(x) g(p) \\
& =-\frac{2 p}{x}
\end{aligned}
$$

Where $f(x)=-\frac{2}{x}$ and $g(p)=p$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{p} d p & =-\frac{2}{x} d x \\
\int \frac{1}{p} d p & =\int-\frac{2}{x} d x \\
\ln (p) & =-2 \ln (x)+c_{1} \\
p & =\mathrm{e}^{-2 \ln (x)+c_{1}} \\
& =\frac{c_{1}}{x^{2}}
\end{aligned}
$$

Since $p=y^{\prime}$ then the new first order ode to solve is

$$
y^{\prime}=\frac{c_{1}}{x^{2}}
$$

Integrating both sides gives

$$
\begin{aligned}
y & =\int \frac{c_{1}}{x^{2}} \mathrm{~d} x \\
& =-\frac{c_{1}}{x}+c_{2}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{c_{1}}{x}+c_{2} \tag{1}
\end{equation*}
$$

## Verification of solutions

$$
y=-\frac{c_{1}}{x}+c_{2}
$$

Verified OK.

### 1.8.3 Solving as second order ode non constant coeff transformation on B ode

Given an ode of the form

$$
A y^{\prime \prime}+B y^{\prime}+C y=F(x)
$$

This method reduces the order ode the ODE by one by applying the transformation

$$
y=B v
$$

This results in

$$
\begin{aligned}
y^{\prime} & =B^{\prime} v+v^{\prime} B \\
y^{\prime \prime} & =B^{\prime \prime} v+B^{\prime} v^{\prime}+v^{\prime \prime} B+v^{\prime} B^{\prime} \\
& =v^{\prime \prime} B+2 v^{\prime}+B^{\prime}+B^{\prime \prime} v
\end{aligned}
$$

And now the original ode becomes

$$
\begin{align*}
A\left(v^{\prime \prime} B+2 v^{\prime} B^{\prime}+B^{\prime \prime} v\right)+B\left(B^{\prime} v+v^{\prime} B\right)+C B v & =0 \\
A B v^{\prime \prime}+\left(2 A B^{\prime}+B^{2}\right) v^{\prime}+\left(A B^{\prime \prime}+B B^{\prime}+C B\right) v & =0 \tag{1}
\end{align*}
$$

If the term $A B^{\prime \prime}+B B^{\prime}+C B$ is zero, then this method works and can be used to solve

$$
A B v^{\prime \prime}+\left(2 A B^{\prime}+B^{2}\right) v^{\prime}=0
$$

By Using $u=v^{\prime}$ which reduces the order of the above ode to one. The new ode is

$$
A B u^{\prime}+\left(2 A B^{\prime}+B^{2}\right) u=0
$$

The above ode is first order ode which is solved for $u$. Now a new ode $v^{\prime}=u$ is solved for $v$ as first order ode. Then the final solution is obtain from $y=B v$.

This method works only if the term $A B^{\prime \prime}+B B^{\prime}+C B$ is zero. The given ODE shows that

$$
\begin{aligned}
& A=x \\
& B=2 \\
& C=0 \\
& F=0
\end{aligned}
$$

The above shows that for this ode

$$
\begin{aligned}
A B^{\prime \prime}+B B^{\prime}+C B & =(x)(0)+(2)(0)+(0)(2) \\
& =0
\end{aligned}
$$

Hence the ode in $v$ given in (1) now simplifies to

$$
2 x v^{\prime \prime}+(4) v^{\prime}=0
$$

Now by applying $v^{\prime}=u$ the above becomes

$$
2 x u^{\prime}(x)+4 u(x)=0
$$

Which is now solved for $u$. In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =-\frac{2 u}{x}
\end{aligned}
$$

Where $f(x)=-\frac{2}{x}$ and $g(u)=u$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{u} d u & =-\frac{2}{x} d x \\
\int \frac{1}{u} d u & =\int-\frac{2}{x} d x \\
\ln (u) & =-2 \ln (x)+c_{1} \\
u & =\mathrm{e}^{-2 \ln (x)+c_{1}} \\
& =\frac{c_{1}}{x^{2}}
\end{aligned}
$$

The ode for $v$ now becomes

$$
\begin{aligned}
v^{\prime} & =u \\
& =\frac{c_{1}}{x^{2}}
\end{aligned}
$$

Which is now solved for $v$. Integrating both sides gives

$$
\begin{aligned}
v(x) & =\int \frac{c_{1}}{x^{2}} \mathrm{~d} x \\
& =-\frac{c_{1}}{x}+c_{2}
\end{aligned}
$$

Therefore the solution is

$$
\begin{aligned}
y(x) & =B v \\
& =(2)\left(-\frac{c_{1}}{x}+c_{2}\right) \\
& =-\frac{2 c_{1}}{x}+2 c_{2}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{2 c_{1}}{x}+2 c_{2} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=-\frac{2 c_{1}}{x}+2 c_{2}
$$

Verified OK.

### 1.8.4 Solving as type second_order_integrable_as_is (not using ABC version)

Writing the ode as

$$
y^{\prime \prime} x+2 y^{\prime}=0
$$

Integrating both sides of the ODE w.r.t $x$ gives

$$
\begin{gathered}
\int\left(y^{\prime \prime} x+2 y^{\prime}\right) d x=0 \\
y^{\prime} x+y=c_{1}
\end{gathered}
$$

Which is now solved for $y$. In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{-y+c_{1}}{x}
\end{aligned}
$$

Where $f(x)=\frac{1}{x}$ and $g(y)=-y+c_{1}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{-y+c_{1}} d y & =\frac{1}{x} d x \\
\int \frac{1}{-y+c_{1}} d y & =\int \frac{1}{x} d x \\
-\ln \left(-y+c_{1}\right) & =\ln (x)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\frac{1}{-y+c_{1}}=\mathrm{e}^{\ln (x)+c_{2}}
$$

Which simplifies to

$$
\frac{1}{-y+c_{1}}=c_{3} x
$$

Which simplifies to

$$
y=\frac{\left(c_{3} \mathrm{e}^{c_{2}} x c_{1}-1\right) \mathrm{e}^{-c_{2}}}{c_{3} x}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\left(c_{3} \mathrm{e}^{c_{2}} x c_{1}-1\right) \mathrm{e}^{-c_{2}}}{c_{3} x} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{\left(c_{3} \mathrm{e}^{c_{2}} x c_{1}-1\right) \mathrm{e}^{-c_{2}}}{c_{3} x}
$$

Verified OK.

### 1.8.5 Solving using Kovacic algorithm

Writing the ode as

$$
\begin{array}{r}
y^{\prime \prime} x+2 y^{\prime}=0 \\
A y^{\prime \prime}+B y^{\prime}+C y=0 \tag{2}
\end{array}
$$

Comparing (1) and (2) shows that

$$
\begin{align*}
A & =x \\
B & =2  \tag{3}\\
C & =0
\end{align*}
$$

Applying the Liouville transformation on the dependent variable gives

$$
z(x)=y e^{\int \frac{B}{2 A} d x}
$$

Then (2) becomes

$$
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
$$

Where $r$ is given by

$$
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
$$

Substituting the values of $A, B, C$ from (3) in the above and simplifying gives

$$
\begin{equation*}
r=\frac{0}{1} \tag{6}
\end{equation*}
$$

Comparing the above to (5) shows that

$$
\begin{aligned}
& s=0 \\
& t=1
\end{aligned}
$$

Therefore eq. (4) becomes

$$
\begin{equation*}
z^{\prime \prime}(x)=0 \tag{7}
\end{equation*}
$$

Equation (7) is now solved. After finding $z(x)$ then $y$ is found using the inverse transformation

$$
y=z(x) e^{-\int \frac{B}{2 A} d x}
$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of $r$ and the order of $r$ at $\infty$. The following table summarizes these cases.

| Case | Allowed pole order for $r$ | Allowed value for $\mathcal{O}(\infty)$ |
| :--- | :--- | :--- |
| 1 | $\{0,1,2,4,6,8, \cdots\}$ | $\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}$ |
| 2 | Need to have at least one pole that <br> is either order 2 or odd order greater <br> than 2. Any other pole order is <br> allowed as long as the above condi- | no condition |
| tion is satisfied. Hence the following |  |  |
| set of pole orders are all allowed. |  |  |
| $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}$. |  |  |
| 3 | $\{1,2\}$ | $\{2,3,4,5,6,7, \cdots\}$ |

Table 9: Necessary conditions for each Kovacic case

The order of $r$ at $\infty$ is the degree of $t$ minus the degree of $s$. Therefore

$$
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0--\infty \\
& =\infty
\end{aligned}
$$

There are no poles in $r$. Therefore the set of poles $\Gamma$ is empty. Since there is no odd order pole larger than 2 and the order at $\infty$ is infinity then the necessary conditions for case one are met. Therefore

$$
L=[1]
$$

Since $r=0$ is not a function of $x$, then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z^{\prime \prime}=r z$ as one solution is

$$
z_{1}(x)=1
$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in $y$ is found from

$$
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{2}{x} d x} \\
& =z_{1} e^{-\ln (x)} \\
& =z_{1}\left(\frac{1}{x}\right)
\end{aligned}
$$

Which simplifies to

$$
y_{1}=\frac{1}{x}
$$

The second solution $y_{2}$ to the original ode is found using reduction of order

$$
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
$$

Substituting gives

$$
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{2}{x} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{-2 \ln (x)}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}(x)
\end{aligned}
$$

Therefore the solution is

$$
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\frac{1}{x}\right)+c_{2}\left(\frac{1}{x}(x)\right)
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{c_{1}}{x}+c_{2} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{c_{1}}{x}+c_{2}
$$

Verified OK.

### 1.8.6 Solving as exact linear second order ode ode

An ode of the form

$$
p(x) y^{\prime \prime}+q(x) y^{\prime}+r(x) y=s(x)
$$

is exact if

$$
\begin{equation*}
p^{\prime \prime}(x)-q^{\prime}(x)+r(x)=0 \tag{1}
\end{equation*}
$$

For the given ode we have

$$
\begin{aligned}
p(x) & =x \\
q(x) & =2 \\
r(x) & =0 \\
s(x) & =0
\end{aligned}
$$

Hence

$$
\begin{aligned}
p^{\prime \prime}(x) & =0 \\
q^{\prime}(x) & =0
\end{aligned}
$$

Therefore (1) becomes

$$
0-(0)+(0)=0
$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$
\left(p(x) y^{\prime}+\left(q(x)-p^{\prime}(x)\right) y\right)^{\prime}=s(x)
$$

Integrating gives

$$
p(x) y^{\prime}+\left(q(x)-p^{\prime}(x)\right) y=\int s(x) d x
$$

Substituting the above values for $p, q, r, s$ gives

$$
y^{\prime} x+y=c_{1}
$$

We now have a first order ode to solve which is

$$
y^{\prime} x+y=c_{1}
$$

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{-y+c_{1}}{x}
\end{aligned}
$$

Where $f(x)=\frac{1}{x}$ and $g(y)=-y+c_{1}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{-y+c_{1}} d y & =\frac{1}{x} d x \\
\int \frac{1}{-y+c_{1}} d y & =\int \frac{1}{x} d x \\
-\ln \left(-y+c_{1}\right) & =\ln (x)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\frac{1}{-y+c_{1}}=\mathrm{e}^{\ln (x)+c_{2}}
$$

Which simplifies to

$$
\frac{1}{-y+c_{1}}=c_{3} x
$$

Which simplifies to

$$
y=\frac{\left(c_{3} \mathrm{e}^{c_{2}} x c_{1}-1\right) \mathrm{e}^{-c_{2}}}{c_{3} x}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\left(c_{3} \mathrm{e}^{c_{2}} x c_{1}-1\right) \mathrm{e}^{-c_{2}}}{c_{3} x} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{\left(c_{3} \mathrm{e}^{c_{2}} x c_{1}-1\right) \mathrm{e}^{-c_{2}}}{c_{3} x}
$$

Verified OK.

### 1.8.7 Maple step by step solution

Let's solve

$$
y^{\prime \prime} x+2 y^{\prime}=0
$$

- Highest derivative means the order of the ODE is 2

$$
y^{\prime \prime}
$$

- Isolate 2nd derivative

$$
y^{\prime \prime}=-\frac{2 y^{\prime}}{x}
$$

- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$
y^{\prime \prime}+\frac{2 y^{\prime}}{x}=0
$$

- Multiply by denominators of the ODE

$$
y^{\prime \prime} x+2 y^{\prime}=0
$$

- Make a change of variables
$t=\ln (x)$
Substitute the change of variables back into the ODE
- Calculate the 1st derivative of y with respect to x , using the chain rule

$$
y^{\prime}=\left(\frac{d}{d t} y(t)\right) t^{\prime}(x)
$$

- Compute derivative

$$
y^{\prime}=\frac{\frac{d}{d t} y(t)}{x}
$$

- Calculate the 2nd derivative of y with respect to x , using the chain rule $y^{\prime \prime}=\left(\frac{d^{2}}{d t^{2}} y(t)\right) t^{\prime}(x)^{2}+t^{\prime \prime}(x)\left(\frac{d}{d t} y(t)\right)$
- Compute derivative
$y^{\prime \prime}=\frac{\frac{d^{2}}{d t^{2}} y(t)}{x^{2}}-\frac{\frac{d}{d t} y(t)}{x^{2}}$
Substitute the change of variables back into the ODE
$\left(\frac{\frac{d^{2}}{d t^{2}} y(t)}{x^{2}}-\frac{d}{d t} y(t) x^{2}\right) x+\frac{2\left(\frac{d}{d t} y(t)\right)}{x}=0$
- $\quad$ Simplify

$$
\frac{\frac{d^{2}}{d t^{2}} y(t)+\frac{d}{d t} y(t)}{x}=0
$$

- Isolate 2 nd derivative
$\frac{d^{2}}{d t^{2}} y(t)=-\frac{d}{d t} y(t)$
- Group terms with $y(t)$ on the lhs of the ODE and the rest on the rhs of the ODE; ODE is lin $\frac{d^{2}}{d t^{2}} y(t)+\frac{d}{d t} y(t)=0$
- Characteristic polynomial of ODE
$r^{2}+r=0$
- Factor the characteristic polynomial
$r(r+1)=0$
- Roots of the characteristic polynomial
$r=(-1,0)$
- $\quad$ 1st solution of the ODE
$y_{1}(t)=\mathrm{e}^{-t}$
- $\quad 2$ nd solution of the ODE
$y_{2}(t)=1$
- General solution of the ODE
$y(t)=c_{1} y_{1}(t)+c_{2} y_{2}(t)$
- $\quad$ Substitute in solutions
$y(t)=c_{1} \mathrm{e}^{-t}+c_{2}$
- Change variables back using $t=\ln (x)$
$y=\frac{c_{1}}{x}+c_{2}$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 11

```
dsolve(diff(y(x),x$2)+2/x*\operatorname{diff}(\textrm{y}(\textrm{x}),\textrm{x})=0,y(x), singsol=all)
```

$$
y(x)=c_{1}+\frac{c_{2}}{x}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.018 (sec). Leaf size: 15
DSolve[y''[x]+2/x*y'[x]==0,y[x],x,IncludeSingularSolutions -> True]

$$
y(x) \rightarrow c_{2}-\frac{c_{1}}{x}
$$

## 1.9 problem 9

1.9.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 90
1.9.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 92
1.9.3 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 93
1.9.4 Solving as first order ode lie symmetry lookup ode . . . . . . . 94
1.9.5 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 98
1.9.6 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 102

Internal problem ID [12426]
Internal file name [OUTPUT/11078_Monday_October_16_2023_09_47_16_PM_30437443/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 9.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
y-y^{\prime} x=0
$$

### 1.9.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{y}{x}
\end{aligned}
$$

Where $f(x)=\frac{1}{x}$ and $g(y)=y$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{y} d y & =\frac{1}{x} d x \\
\int \frac{1}{y} d y & =\int \frac{1}{x} d x \\
\ln (y) & =\ln (x)+c_{1} \\
y & =\mathrm{e}^{\ln (x)+c_{1}} \\
& =c_{1} x
\end{aligned}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} x \tag{1}
\end{equation*}
$$



Figure 4: Slope field plot

Verification of solutions

$$
y=c_{1} x
$$

Verified OK.

### 1.9.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{1}{x} \\
& q(x)=0
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{y}{x}=0
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{1}{x} d x} \\
& =\frac{1}{x}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x} \mu y & =0 \\
\frac{\mathrm{~d}}{\mathrm{~d} x}\left(\frac{y}{x}\right) & =0
\end{aligned}
$$

Integrating gives

$$
\frac{y}{x}=c_{1}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x}$ results in

$$
y=c_{1} x
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} x \tag{1}
\end{equation*}
$$



Figure 5: Slope field plot
Verification of solutions

$$
y=c_{1} x
$$

Verified OK.

### 1.9.3 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
u(x) x-\left(u^{\prime}(x) x+u(x)\right) x=0
$$

Integrating both sides gives

$$
\begin{aligned}
u(x) & =\int 0 \mathrm{~d} x \\
& =c_{2}
\end{aligned}
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =u x \\
& =c_{2} x
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{2} x \tag{1}
\end{equation*}
$$



Figure 6: Slope field plot

Verification of solutions

$$
y=c_{2} x
$$

Verified OK.

### 1.9.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{y}{x} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 11: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=x \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{x} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{x}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{y}{x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{y}{x^{2}} \\
S_{y} & =\frac{1}{x}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode．This is the whole point of Lie symmetry method． It converts an ode，no matter how complicated it is，to one that can be solved by integration when the ode is in the canonical coordiates $R, S$ ．Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution，we just need to transform（4）back to $x, y$ coordinates．This results in

$$
\frac{y}{x}=c_{1}
$$

Which simplifies to

$$
\frac{y}{x}=c_{1}
$$

Which gives

$$
y=c_{1} x
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown．

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{y}{x}$ |  | $\frac{d S}{d R}=0$ |
|  |  | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \longrightarrow \rightarrow \rightarrow \rightarrow \longrightarrow \rightarrow \longrightarrow}$ |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\rightarrow \longrightarrow \rightarrow \longrightarrow \rightarrow \longrightarrow]{ }$ |
|  |  |  |
| $\cdots$ |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \longrightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \longrightarrow]{ }$ |
|  | $R=x$ S |  |
|  | $=\frac{y}{x}$ | $\xrightarrow{\sim \rightarrow \rightarrow \rightarrow \rightarrow-R_{0 \rightarrow \rightarrow}}$ |
| 多多多夝早新： |  | $\xrightarrow{-2 \rightarrow \longrightarrow \rightarrow \longrightarrow \rightarrow \longrightarrow \longrightarrow}$ |
|  |  | $\xrightarrow{\sim \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow}$ |
|  |  | $\xrightarrow{+}$ |

Summary
The solution（s）found are the following

$$
\begin{equation*}
y=c_{1} x \tag{1}
\end{equation*}
$$



Figure 7: Slope field plot
Verification of solutions

$$
y=c_{1} x
$$

Verified OK.

### 1.9.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{y}\right) \mathrm{d} y & =\left(\frac{1}{x}\right) \mathrm{d} x \\
\left(-\frac{1}{x}\right) \mathrm{d} x+\left(\frac{1}{y}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=-\frac{1}{x} \\
& N(x, y)=\frac{1}{y}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{x}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{y}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{1}{x} \mathrm{~d} x \\
\phi & =-\ln (x)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{y}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{y}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{y}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{y}\right) \mathrm{d} y \\
f(y) & =\ln (y)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\ln (x)+\ln (y)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\ln (x)+\ln (y)
$$

The solution becomes

$$
y=\mathrm{e}^{c_{1}} x
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{c_{1}} x \tag{1}
\end{equation*}
$$



Figure 8: Slope field plot
Verification of solutions

$$
y=\mathrm{e}^{c_{1}} x
$$

Verified OK.

### 1.9.6 Maple step by step solution

Let's solve

$$
y-y^{\prime} x=0
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables

$$
\frac{y^{\prime}}{y}=\frac{1}{x}
$$

- Integrate both sides with respect to $x$

$$
\int \frac{y^{\prime}}{y} d x=\int \frac{1}{x} d x+c_{1}
$$

- Evaluate integral

$$
\ln (y)=\ln (x)+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\mathrm{e}^{c_{1}} x
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\sqrt{ }$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 7

```
dsolve(y(x)-x*diff(y(x),x)=0,y(x), singsol=all)
```

$$
y(x)=c_{1} x
$$

$\checkmark$ Solution by Mathematica
Time used: 0.037 (sec). Leaf size: 14
DSolve[y[x]-x*y'[x]==0,y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow c_{1} x \\
& y(x) \rightarrow 0
\end{aligned}
$$

### 1.10 problem 10

$$
\text { 1.10.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . } 104
$$

1.10.2 Solving as first order ode lie symmetry lookup ode ..... 106
1.10.3 Solving as exact ode ..... 110
1.10.4 Maple step by step solution ..... 114

Internal problem ID [12427]
Internal file name [OUTPUT/11079_Monday_October_16_2023_09_47_17_PM_6484944/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 10.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first__order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type

```
[_separable]
```

$$
(u+1) v+(1-v) u v^{\prime}=0
$$

### 1.10.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
v^{\prime} & =F(u, v) \\
& =f(u) g(v) \\
& =\frac{v(u+1)}{(-1+v) u}
\end{aligned}
$$

Where $f(u)=\frac{u+1}{u}$ and $g(v)=\frac{v}{-1+v}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{v}{-1+v}} d v & =\frac{u+1}{u} d u \\
\int \frac{1}{\frac{v}{-1+v}} d v & =\int \frac{u+1}{u} d u \\
v-\ln (v) & =u+\ln (u)+c_{1}
\end{aligned}
$$

Which results in

$$
v=-\operatorname{LambertW}\left(-\frac{\mathrm{e}^{-u-c_{1}}}{u}\right)
$$

Since $c_{1}$ is constant, then exponential powers of this constant are constants also, and these can be simplified to just $c_{1}$ in the above solution. Which simplifies to

$$
v=-\operatorname{LambertW}\left(-\frac{\mathrm{e}^{-u-c_{1}}}{u}\right)
$$

gives

$$
v=-\operatorname{LambertW}\left(-\frac{\mathrm{e}^{-u}}{c_{1} u}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
v=- \text { LambertW }\left(-\frac{\mathrm{e}^{-u}}{c_{1} u}\right) \tag{1}
\end{equation*}
$$



Figure 9: Slope field plot

## Verification of solutions

$$
v=-\operatorname{LambertW}\left(-\frac{\mathrm{e}^{-u}}{c_{1} u}\right)
$$

Verified OK.

### 1.10.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
v^{\prime} & =\frac{v(u+1)}{(-1+v) u} \\
v^{\prime} & =\omega(u, v)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{u}+\omega\left(\eta_{v}-\xi_{u}\right)-\omega^{2} \xi_{v}-\omega_{u} \xi-\omega_{v} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the $\operatorname{PDE}(\mathrm{A})$, and can just use the lookup table shown below to find $\xi, \eta$

Table 14: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(u, v)=\frac{u}{u+1} \\
& \eta(u, v)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(u, v) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d u}{\xi}=\frac{d v}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial u}+\eta \frac{\partial}{\partial v}\right) S(u, v)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=v
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d u \\
& =\int \frac{1}{\frac{u}{u+1}} d u
\end{aligned}
$$

Which results in

$$
S=u+\ln (u)
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{u}+\omega(u, v) S_{v}}{R_{u}+\omega(u, v) R_{v}} \tag{2}
\end{equation*}
$$

Where in the above $R_{u}, R_{v}, S_{u}, S_{v}$ are all partial derivatives and $\omega(u, v)$ is the right hand side of the original ode given by

$$
\omega(u, v)=\frac{v(u+1)}{(-1+v) u}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{u} & =0 \\
R_{v} & =1 \\
S_{u} & =1+\frac{1}{u} \\
S_{v} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{-1+v}{v} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $u, v$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{-1+R}{R}
$$

The above is a quadrature ode．This is the whole point of Lie symmetry method． It converts an ode，no matter how complicated it is，to one that can be solved by integration when the ode is in the canonical coordiates $R, S$ ．Integrating the above gives

$$
\begin{equation*}
S(R)=R-\ln (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution，we just need to transform（4）back to $u, v$ coordinates．This results in

$$
u+\ln (u)=v-\ln (v)+c_{1}
$$

Which simplifies to

$$
u+\ln (u)=v-\ln (v)+c_{1}
$$

Which gives

$$
v=-\operatorname{LambertW}\left(-\frac{\mathrm{e}^{-u+c_{1}}}{u}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown．

| Original ode in $u, v$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d v}{d u}=\frac{v(u+1)}{(-1+v) u}$ |  | $\frac{d S}{d R}=\frac{-1+R}{R}$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
| なしなくなくすかなatata |  |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow 0]{ }$ | $R=v$ | $\rightarrow$ 为 |
| $\rightarrow \rightarrow-4 \rightarrow- \pm \rightarrow \infty$ | $S=u+\ln (u)$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  | 加加加加加期 |

## Summary

The solution（s）found are the following

$$
\begin{equation*}
v=-\operatorname{LambertW}\left(-\frac{\mathrm{e}^{-u+c_{1}}}{u}\right) \tag{1}
\end{equation*}
$$



Figure 10: Slope field plot

## Verification of solutions

$$
v=-\operatorname{LambertW}\left(-\frac{\mathrm{e}^{-u+c_{1}}}{u}\right)
$$

Verified OK.

### 1.10.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(u, v) \mathrm{d} u+N(u, v) \mathrm{d} v=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{-1+v}{v}\right) \mathrm{d} v & =\left(\frac{u+1}{u}\right) \mathrm{d} u \\
\left(-\frac{u+1}{u}\right) \mathrm{d} u+\left(\frac{-1+v}{v}\right) \mathrm{d} v & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(u, v)=-\frac{u+1}{u} \\
& N(u, v)=\frac{-1+v}{v}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial v}=\frac{\partial N}{\partial u}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial v} & =\frac{\partial}{\partial v}\left(-\frac{u+1}{u}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial u} & =\frac{\partial}{\partial u}\left(\frac{-1+v}{v}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial v}=\frac{\partial N}{\partial u}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(u, v)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial u}=M  \tag{1}\\
& \frac{\partial \phi}{\partial v}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $u$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial u} \mathrm{~d} u & =\int M \mathrm{~d} u \\
\int \frac{\partial \phi}{\partial u} \mathrm{~d} u & =\int-\frac{u+1}{u} \mathrm{~d} u \\
\phi & =-u-\ln (u)+f(v) \tag{3}
\end{align*}
$$

Where $f(v)$ is used for the constant of integration since $\phi$ is a function of both $u$ and $v$. Taking derivative of equation (3) w.r.t $v$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial v}=0+f^{\prime}(v) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial v}=\frac{-1+v}{v}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{-1+v}{v}=0+f^{\prime}(v) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(v)$ gives

$$
f^{\prime}(v)=\frac{-1+v}{v}
$$

Integrating the above w.r.t $v$ gives

$$
\begin{aligned}
\int f^{\prime}(v) \mathrm{d} v & =\int\left(\frac{-1+v}{v}\right) \mathrm{d} v \\
f(v) & =v-\ln (v)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(v)$ into equation (3) gives $\phi$

$$
\phi=-u-\ln (u)+v-\ln (v)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-u-\ln (u)+v-\ln (v)
$$

The solution becomes

$$
v=- \text { LambertW }\left(-\frac{\mathrm{e}^{-u-c_{1}}}{u}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
v=-\operatorname{LambertW}\left(-\frac{\mathrm{e}^{-u-c_{1}}}{u}\right) \tag{1}
\end{equation*}
$$



Figure 11: Slope field plot

Verification of solutions

$$
v=-\operatorname{LambertW}\left(-\frac{\mathrm{e}^{-u-c_{1}}}{u}\right)
$$

Verified OK.

### 1.10.4 Maple step by step solution

Let's solve

$$
(u+1) v+(1-v) u v^{\prime}=0
$$

- Highest derivative means the order of the ODE is 1 $v^{\prime}$
- Separate variables
$\frac{v^{\prime}(1-v)}{v}=-\frac{u+1}{u}$
- Integrate both sides with respect to $u$
$\int \frac{v^{\prime}(1-v)}{v} d u=\int-\frac{u+1}{u} d u+c_{1}$
- Evaluate integral
$-v+\ln (v)=-u-\ln (u)+c_{1}$
- $\quad$ Solve for $v$
$v=-$ Lambert $W\left(-\frac{\mathrm{e}^{-u+c_{1}}}{u}\right)$

Maple trace

```
Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 19
dsolve ((1+u) $\mathrm{v}(\mathrm{u})+(1-\mathrm{v}(\mathrm{u})) * u * \operatorname{diff}(\mathrm{v}(\mathrm{u}), \mathrm{u})=0, \mathrm{v}(\mathrm{u})$, singsol=all)

$$
v(u)=- \text { LambertW }\left(-\frac{\mathrm{e}^{-u}}{c_{1} u}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 4.942 (sec). Leaf size: 28
DSolve[(1+u)*v[u]+(1-v[u])*u*v'[u]==0,v[u],u,IncludeSingularSolutions -> True]

$$
\begin{aligned}
& v(u) \rightarrow-W\left(-\frac{e^{-u-c_{1}}}{u}\right) \\
& v(u) \rightarrow 0
\end{aligned}
$$

### 1.11 problem 11

1.11.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 116
1.11.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 118
1.11.3 Solving as differentialType ode . . . . . . . . . . . . . . . . . . 120
1.11.4 Solving as homogeneousTypeMapleC ode . . . . . . . . . . . . . 121
1.11.5 Solving as first order ode lie symmetry lookup ode . . . . . . . 124
1.11.6 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 128
1.11.7 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 132

Internal problem ID [12428]
Internal file name [OUTPUT/11080_Monday_October_16_2023_09_47_17_PM_8291719/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 11.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "separable", "differentialType", "homogeneousTypeMapleC", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
y-(1-x) y^{\prime}=-1
$$

### 1.11.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{-y-1}{x-1}
\end{aligned}
$$

Where $f(x)=\frac{1}{x-1}$ and $g(y)=-y-1$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{-y-1} d y & =\frac{1}{x-1} d x \\
\int \frac{1}{-y-1} d y & =\int \frac{1}{x-1} d x \\
-\ln (y+1) & =\ln (x-1)+c_{1}
\end{aligned}
$$

Raising both side to exponential gives

$$
\frac{1}{y+1}=\mathrm{e}^{\ln (x-1)+c_{1}}
$$

Which simplifies to

$$
\frac{1}{y+1}=c_{2}(x-1)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\left(c_{2} \mathrm{e}^{\ln (x-1)+c_{1}}-1\right) \mathrm{e}^{-c_{1}}}{c_{2}(x-1)} \tag{1}
\end{equation*}
$$



Figure 12: Slope field plot

## Verification of solutions

$$
y=-\frac{\left(c_{2} \mathrm{e}^{\ln (x-1)+c_{1}}-1\right) \mathrm{e}^{-c_{1}}}{c_{2}(x-1)}
$$

Verified OK.

### 1.11.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=\frac{1}{x-1} \\
& q(x)=-\frac{1}{x-1}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+\frac{y}{x-1}=-\frac{1}{x-1}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \frac{1}{x-1} d x} \\
& =x-1
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(-\frac{1}{x-1}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}((x-1) y) & =(x-1)\left(-\frac{1}{x-1}\right) \\
\mathrm{d}((x-1) y) & =-1 \mathrm{~d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& (x-1) y=\int-1 \mathrm{~d} x \\
& (x-1) y=-x+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=x-1$ results in

$$
y=-\frac{x}{x-1}+\frac{c_{1}}{x-1}
$$

which simplifies to

$$
y=\frac{-x+c_{1}}{x-1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{-x+c_{1}}{x-1} \tag{1}
\end{equation*}
$$



Figure 13: Slope field plot

Verification of solutions

$$
y=\frac{-x+c_{1}}{x-1}
$$

Verified OK.

### 1.11.3 Solving as differentialType ode

Writing the ode as

$$
\begin{equation*}
y^{\prime}=\frac{-y-1}{x-1} \tag{1}
\end{equation*}
$$

Which becomes

$$
\begin{equation*}
0=(1-x) d y+(-y-1) d x \tag{2}
\end{equation*}
$$

But the RHS is complete differential because

$$
(1-x) d y+(-y-1) d x=d(-(y+1) x+y)
$$

Hence (2) becomes

$$
0=d(-(y+1) x+y)
$$

Integrating both sides gives gives these solutions

$$
y=\frac{-x+c_{1}}{x-1}+c_{1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{-x+c_{1}}{x-1}+c_{1} \tag{1}
\end{equation*}
$$



Figure 14: Slope field plot

## Verification of solutions

$$
y=\frac{-x+c_{1}}{x-1}+c_{1}
$$

Verified OK.

### 1.11.4 Solving as homogeneousTypeMapleC ode

Let $Y=y+y_{0}$ and $X=x+x_{0}$ then the above is transformed to new ode in $Y(X)$

$$
\frac{d}{d X} Y(X)=-\frac{Y(X)+y_{0}+1}{X+x_{0}-1}
$$

Solving for possible values of $x_{0}$ and $y_{0}$ which makes the above ode a homogeneous ode results in

$$
\begin{aligned}
& x_{0}=1 \\
& y_{0}=-1
\end{aligned}
$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$
\frac{d}{d X} Y(X)=-\frac{Y(X)}{X}
$$

In canonical form, the ODE is

$$
\begin{align*}
Y^{\prime} & =F(X, Y) \\
& =-\frac{Y}{X} \tag{1}
\end{align*}
$$

An ode of the form $Y^{\prime}=\frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order $n$ if

$$
f\left(t^{n} X, t^{n} Y\right)=t^{n} f(X, Y)
$$

In this case, it can be seen that both $M=-Y$ and $N=X$ are both homogeneous and of the same order $n=1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u=\frac{Y}{X}$, or $Y=u X$. Hence

$$
\frac{\mathrm{d} Y}{\mathrm{~d} X}=\frac{\mathrm{d} u}{\mathrm{~d} X} X+u
$$

Applying the transformation $Y=u X$ to the above ODE in (1) gives

$$
\begin{aligned}
\frac{\mathrm{d} u}{\mathrm{~d} X} X+u & =-u \\
\frac{\mathrm{~d} u}{\mathrm{~d} X} & =-\frac{2 u(X)}{X}
\end{aligned}
$$

Or

$$
\frac{d}{d X} u(X)+\frac{2 u(X)}{X}=0
$$

Or

$$
\left(\frac{d}{d X} u(X)\right) X+2 u(X)=0
$$

Which is now solved as separable in $u(X)$. Which is now solved in $u(X)$. In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(X, u) \\
& =f(X) g(u) \\
& =-\frac{2 u}{X}
\end{aligned}
$$

Where $f(X)=-\frac{2}{X}$ and $g(u)=u$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{u} d u & =-\frac{2}{X} d X \\
\int \frac{1}{u} d u & =\int-\frac{2}{X} d X \\
\ln (u) & =-2 \ln (X)+c_{2} \\
u & =\mathrm{e}^{-2 \ln (X)+c_{2}} \\
& =\frac{c_{2}}{X^{2}}
\end{aligned}
$$

Now $u$ in the above solution is replaced back by $Y$ using $u=\frac{Y}{X}$ which results in the solution

$$
Y(X)=\frac{c_{2}}{X}
$$

Using the solution for $Y(X)$

$$
Y(X)=\frac{c_{2}}{X}
$$

And replacing back terms in the above solution using

$$
\begin{aligned}
Y & =y+y_{0} \\
X & =x+x_{0}
\end{aligned}
$$

Or

$$
\begin{aligned}
& Y=y-1 \\
& X=x+1
\end{aligned}
$$

Then the solution in $y$ becomes

$$
y+1=\frac{c_{2}}{x-1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y+1=\frac{c_{2}}{x-1} \tag{1}
\end{equation*}
$$



Figure 15: Slope field plot

## Verification of solutions

$$
y+1=\frac{c_{2}}{x-1}
$$

Verified OK.

### 1.11.5 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{y+1}{x-1} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 17: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{2}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\underline{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}} \frac{a_{1} b_{2}-a_{2} b_{1}}{}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\frac{1}{x-1} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{1}{x-1}} d y
\end{aligned}
$$

Which results in

$$
S=(x-1) y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{y+1}{x-1}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =y \\
S_{y} & =x-1
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-1 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-1
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-R+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
y(x-1)=-x+c_{1}
$$

Which simplifies to

$$
y(x-1)=-x+c_{1}
$$

Which gives

$$
y=\frac{-x+c_{1}}{x-1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical <br> coordinates <br> transformation | ODE in canonical coordinates <br> $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=-\frac{y+1}{x-1}$ |  | $\frac{d S}{d R}=-1$ |
| 为 |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{-x+c_{1}}{x-1} \tag{1}
\end{equation*}
$$



Figure 16: Slope field plot

## Verification of solutions

$$
y=\frac{-x+c_{1}}{x-1}
$$

Verified OK.

### 1.11.6 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{-y-1}\right) \mathrm{d} y & =\left(\frac{1}{x-1}\right) \mathrm{d} x \\
\left(-\frac{1}{x-1}\right) \mathrm{d} x+\left(\frac{1}{-y-1}\right) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=-\frac{1}{x-1} \\
& N(x, y)=\frac{1}{-y-1}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{x-1}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{-y-1}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{1}{x-1} \mathrm{~d} x \\
\phi & =-\ln (x-1)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{-y-1}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{-y-1}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=-\frac{1}{y+1}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(-\frac{1}{y+1}\right) \mathrm{d} y \\
f(y) & =-\ln (y+1)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\ln (x-1)-\ln (y+1)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\ln (x-1)-\ln (y+1)
$$

The solution becomes

$$
y=-\frac{\left(\mathrm{e}^{c_{1}} x-\mathrm{e}^{c_{1}}-1\right) \mathrm{e}^{-c_{1}}}{x-1}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\left(\mathrm{e}^{c_{1}} x-\mathrm{e}^{c_{1}}-1\right) \mathrm{e}^{-c_{1}}}{x-1} \tag{1}
\end{equation*}
$$



Figure 17: Slope field plot

Verification of solutions

$$
y=-\frac{\left(\mathrm{e}^{c_{1}} x-\mathrm{e}^{c_{1}}-1\right) \mathrm{e}^{-c_{1}}}{x-1}
$$

Verified OK.

### 1.11.7 Maple step by step solution

Let's solve

$$
y-(1-x) y^{\prime}=-1
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Integrate both sides with respect to $x$

$$
\int\left(y-(1-x) y^{\prime}\right) d x=\int(-1) d x+c_{1}
$$

- Evaluate integral

$$
y(x-1)=-x+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\frac{-x+c_{1}}{x-1}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

Solution by Maple
Time used: 0.0 (sec). Leaf size: 15

```
dsolve((1+y(x))-(1-x)*diff (y(x),x)=0,y(x), singsol=all)
```

$$
y(x)=\frac{c_{1}-x}{-1+x}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.044 (sec). Leaf size: 22
DSolve[(1+y[x])-(1-x)*y'[x]==0,y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow \frac{x+c_{1}}{1-x} \\
& y(x) \rightarrow-1
\end{aligned}
$$

### 1.12 problem 12

1.12.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 134
1.12.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 136
1.12.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 140
1.12.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 144

Internal problem ID [12429]
Internal file name [OUTPUT/11081_Monday_October_16_2023_09_47_18_PM_89448285/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 12.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "separable", "first__order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
\left(t^{2}+x t^{2}\right) x^{\prime}+x^{2}+t x^{2}=0
$$

### 1.12.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
x^{\prime} & =F(t, x) \\
& =f(t) g(x) \\
& =-\frac{x^{2}(1+t)}{t^{2}(x+1)}
\end{aligned}
$$

Where $f(t)=-\frac{1+t}{t^{2}}$ and $g(x)=\frac{x^{2}}{x+1}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{x^{2}}{x+1}} d x & =-\frac{1+t}{t^{2}} d t \\
\int \frac{1}{\frac{x^{2}}{x+1}} d x & =\int-\frac{1+t}{t^{2}} d t
\end{aligned}
$$

$$
-\frac{1}{x}+\ln (x)=\frac{1}{t}-\ln (t)+c_{1}
$$

Which results in

$$
x=\frac{1}{\operatorname{LambertW}\left(\mathrm{e}^{\frac{t \ln (t)-c_{1} t-1}{t}}\right)}
$$

Since $c_{1}$ is constant, then exponential powers of this constant are constants also, and these can be simplified to just $c_{1}$ in the above solution. Which simplifies to

$$
x=\frac{1}{\operatorname{LambertW}\left(\mathrm{e}^{\frac{t \ln (t)-c_{1} t-1}{t}}\right)}
$$

gives

$$
x=\frac{1}{\operatorname{LambertW}\left(\frac{t \mathrm{e}^{-\frac{1}{t}}}{c_{1}}\right)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
x=\frac{1}{\text { LambertW }\left(\frac{t \mathrm{e}^{-\frac{1}{t}}}{c_{1}}\right)} \tag{1}
\end{equation*}
$$



Figure 18: Slope field plot

## Verification of solutions

$$
x=\frac{1}{\text { LambertW }\left(\frac{t \mathrm{e}^{-\frac{1}{t}}}{c_{1}}\right)}
$$

Verified OK.

### 1.12.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& x^{\prime}=-\frac{x^{2}(1+t)}{t^{2}(x+1)} \\
& x^{\prime}=\omega(t, x)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{x}-\xi_{t}\right)-\omega^{2} \xi_{x}-\omega_{t} \xi-\omega_{x} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 20: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, x)=-\frac{t^{2}}{1+t} \\
& \eta(t, x)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(t, x) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.
The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d x}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial x}\right) S(t, x)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d t \\
& =\int \frac{1}{-\frac{t^{2}}{1+t}} d t
\end{aligned}
$$

Which results in

$$
S=\frac{1}{t}-\ln (t)
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, x) S_{x}}{R_{t}+\omega(t, x) R_{x}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{x}, S_{t}, S_{x}$ are all partial derivatives and $\omega(t, x)$ is the right hand side of the original ode given by

$$
\omega(t, x)=-\frac{x^{2}(1+t)}{t^{2}(x+1)}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =0 \\
R_{x} & =1 \\
S_{t} & =\frac{-1-t}{t^{2}} \\
S_{x} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{x+1}{x^{2}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, x$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{R+1}{R^{2}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\frac{1}{R}+\ln (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, x$ coordinates. This results in

$$
\frac{-t \ln (t)+1}{t}=-\frac{1}{x}+\ln (x)+c_{1}
$$

Which simplifies to

$$
\frac{-t \ln (t)+1}{t}=-\frac{1}{x}+\ln (x)+c_{1}
$$

Which gives

$$
x=\frac{1}{\operatorname{LambertW}\left(\mathrm{e}^{\frac{t \ln (t)+c_{1} t-1}{t}}\right)}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, x$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d x}{d t}=-\frac{x^{2}(1+t)}{t^{2}(x+1)}$ |  | $\frac{d S}{d R}=\frac{R+1}{R^{2}}$ |
|  |  | $\rightarrow \rightarrow \rightarrow$ - |
|  |  | $\xrightarrow{+\infty}$ |
|  |  |  |
|  |  | $\rightarrow \rightarrow(R) \rightarrow+$ + $+\rightarrow \rightarrow \rightarrow \infty$ |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \infty]{ }$ |  | $\rightarrow \rightarrow \rightarrow \rightarrow+$ + |
| $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\infty}$ | $R=x$ | $\rightarrow \rightarrow \rightarrow \infty$ |
|  | $S=-t \ln (t)+1$ | $\xrightarrow{\rightarrow \rightarrow-4 \rightarrow \rightarrow \rightarrow-\infty}$ |
|  | $S=\frac{t}{}$ |  |
|  |  | 䢖 |
|  |  |  |
|  |  | $\rightarrow \rightarrow$ - |
|  |  | 为 |

Summary
The solution(s) found are the following

$$
\begin{equation*}
x=\frac{1}{\operatorname{LambertW}\left(\mathrm{e}^{\frac{t \ln (t)+c_{1} t-1}{t}}\right)} \tag{1}
\end{equation*}
$$



Figure 19: Slope field plot

Verification of solutions

$$
x=\frac{1}{\operatorname{LambertW}\left(\mathrm{e}^{\frac{t \ln (t)+c_{1} t-1}{t}}\right)}
$$

Verified OK.

### 1.12.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the
ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, x) \mathrm{d} t+N(t, x) \mathrm{d} x=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(-\frac{x+1}{x^{2}}\right) \mathrm{d} x & =\left(\frac{1+t}{t^{2}}\right) \mathrm{d} t \\
\left(-\frac{1+t}{t^{2}}\right) \mathrm{d} t+\left(-\frac{x+1}{x^{2}}\right) \mathrm{d} x & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, x) & =-\frac{1+t}{t^{2}} \\
N(t, x) & =-\frac{x+1}{x^{2}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial x}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial x} & =\frac{\partial}{\partial x}\left(-\frac{1+t}{t^{2}}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}\left(-\frac{x+1}{x^{2}}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial x}=\frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, x)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=M  \tag{1}\\
& \frac{\partial \phi}{\partial x}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int M \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\frac{1+t}{t^{2}} \mathrm{~d} t \\
\phi & =\frac{1}{t}-\ln (t)+f(x) \tag{3}
\end{align*}
$$

Where $f(x)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $x$. Taking derivative of equation (3) w.r.t $x$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}=0+f^{\prime}(x) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial x}=-\frac{x+1}{x^{2}}$. Therefore equation (4) becomes

$$
\begin{equation*}
-\frac{x+1}{x^{2}}=0+f^{\prime}(x) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(x)$ gives

$$
f^{\prime}(x)=-\frac{x+1}{x^{2}}
$$

Integrating the above w.r.t $x$ gives

$$
\begin{aligned}
\int f^{\prime}(x) \mathrm{d} x & =\int\left(\frac{-x-1}{x^{2}}\right) \mathrm{d} x \\
f(x) & =\frac{1}{x}-\ln (x)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(x)$ into equation (3) gives $\phi$

$$
\phi=\frac{1}{t}-\ln (t)+\frac{1}{x}-\ln (x)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{1}{t}-\ln (t)+\frac{1}{x}-\ln (x)
$$

The solution becomes

$$
x=\frac{1}{\text { LambertW }\left(\mathrm{e}^{\frac{t \ln (t)+c_{1} t-1}{t}}\right)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
x=\frac{1}{\text { LambertW }\left(\mathrm{e}^{\frac{t \ln (t)+c_{1} t-1}{t}}\right)} \tag{1}
\end{equation*}
$$



Figure 20: Slope field plot

Verification of solutions

$$
x=\frac{1}{\operatorname{LambertW}\left(\mathrm{e}^{\frac{t \ln (t)+c_{1} t-1}{t}}\right)}
$$

Verified OK.

### 1.12.4 Maple step by step solution

Let's solve

$$
\left(t^{2}+x t^{2}\right) x^{\prime}+x^{2}+t x^{2}=0
$$

- Highest derivative means the order of the ODE is 1

$$
x^{\prime}
$$

- $\quad$ Separate variables

$$
\frac{x^{\prime}(x+1)}{x^{2}}=-\frac{1+t}{t^{2}}
$$

- Integrate both sides with respect to $t$

$$
\int \frac{x^{\prime}(x+1)}{x^{2}} d t=\int-\frac{1+t}{t^{2}} d t+c_{1}
$$

- Evaluate integral

$$
-\frac{1}{x}+\ln (x)=\frac{1}{t}-\ln (t)+c_{1}
$$

- $\quad$ Solve for $x$

$$
x=\frac{1}{\operatorname{Lambert} W\left(\mathrm{e}^{\frac{t \ln (t)-c_{1} t-1}{t}}\right)}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

Solution by Maple
Time used: 0.032 (sec). Leaf size: 16

```
dsolve((t^2+x(t)*t^2)*diff(x(t),t)+x(t)^2+t*x(t)^2=0,x(t), singsol=all)
```

$$
x(t)=\frac{1}{\text { LambertW }\left(c_{1} t \mathrm{e}^{-\frac{1}{t}}\right)}
$$

$\checkmark$ Solution by Mathematica
Time used: 5.02 (sec). Leaf size: 27
DSolve[(t^2+x[t]*t^2)*x'[t]+x[t]^2+t*x[t]^2=0,x[t],t,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& x(t) \rightarrow \frac{1}{W\left(t e^{-\frac{1}{t}-c_{1}}\right)} \\
& x(t) \rightarrow 0
\end{aligned}
$$

### 1.13 problem 13

1.13.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 146
1.13.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 147
1.13.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 148
1.13.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 151
1.13.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 154

Internal problem ID [12430]
Internal file name [OUTPUT/11082_Monday_October_16_2023_09_47_20_PM_81844345/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 13.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "separable", "first_order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
x^{2} y^{\prime}+y=a
$$

### 1.13.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{a-y}{x^{2}}
\end{aligned}
$$

Where $f(x)=\frac{1}{x^{2}}$ and $g(y)=a-y$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{a-y} d y & =\frac{1}{x^{2}} d x \\
\int \frac{1}{a-y} d y & =\int \frac{1}{x^{2}} d x
\end{aligned}
$$

$$
-\ln (a-y)=-\frac{1}{x}+c_{1}
$$

Raising both side to exponential gives

$$
\frac{1}{a-y}=\mathrm{e}^{-\frac{1}{x}+c_{1}}
$$

Which simplifies to

$$
\frac{1}{a-y}=\mathrm{e}^{-\frac{1}{x}} c_{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\left(c_{2} \mathrm{e}^{\frac{c_{1} x-1}{x}} a-1\right) \mathrm{e}^{-\frac{c_{1} x-1}{x}}}{c_{2}} \tag{1}
\end{equation*}
$$

## Verification of solutions

$$
y=\frac{\left(c_{2} \mathrm{e}^{\frac{c_{1} x-1}{x}} a-1\right) \mathrm{e}^{-\frac{c_{1} x-1}{x}}}{c_{2}}
$$

Verified OK.

### 1.13.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=\frac{1}{x^{2}} \\
& q(x)=\frac{a}{x^{2}}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+\frac{y}{x^{2}}=\frac{a}{x^{2}}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \frac{1}{x^{2}} d x} \\
& =\mathrm{e}^{-\frac{1}{x}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\frac{a}{x^{2}}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\mathrm{e}^{-\frac{1}{x}} y\right) & =\left(\mathrm{e}^{-\frac{1}{x}}\right)\left(\frac{a}{x^{2}}\right) \\
\mathrm{d}\left(\mathrm{e}^{-\frac{1}{x}} y\right) & =\left(\frac{a \mathrm{e}^{-\frac{1}{x}}}{x^{2}}\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \mathrm{e}^{-\frac{1}{x}} y=\int \frac{a \mathrm{e}^{-\frac{1}{x}}}{x^{2}} \mathrm{~d} x \\
& \mathrm{e}^{-\frac{1}{x}} y=a \mathrm{e}^{-\frac{1}{x}}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{-\frac{1}{x}}$ results in

$$
y=\mathrm{e}^{\frac{1}{x}} a \mathrm{e}^{-\frac{1}{x}}+c_{1} \mathrm{e}^{\frac{1}{x}}
$$

which simplifies to

$$
y=a+c_{1} \mathrm{e}^{\frac{1}{x}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=a+c_{1} \mathrm{e}^{\frac{1}{x}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=a+c_{1} \mathrm{e}^{\frac{1}{x}}
$$

Verified OK.

### 1.13.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=-\frac{-a+y}{x^{2}} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 23: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\mathrm{e}^{\frac{1}{x}} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{\frac{1}{x}}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{-\frac{1}{x}} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{-a+y}{x^{2}}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{\mathrm{e}^{-\frac{1}{x}} y}{x^{2}} \\
S_{y} & =\mathrm{e}^{-\frac{1}{x}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{a \mathrm{e}^{-\frac{1}{x}}}{x^{2}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{a \mathrm{e}^{-\frac{1}{R}}}{R^{2}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=a \mathrm{e}^{-\frac{1}{R}}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\mathrm{e}^{-\frac{1}{x}} y=a \mathrm{e}^{-\frac{1}{x}}+c_{1}
$$

Which simplifies to

$$
\mathrm{e}^{-\frac{1}{x}}(y-a)-c_{1}=0
$$

Which gives

$$
y=\left(a \mathrm{e}^{-\frac{1}{x}}+c_{1}\right) \mathrm{e}^{\frac{1}{x}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\left(a \mathrm{e}^{-\frac{1}{x}}+c_{1}\right) \mathrm{e}^{\frac{1}{x}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\left(a \mathrm{e}^{-\frac{1}{x}}+c_{1}\right) \mathrm{e}^{\frac{1}{x}}
$$

Verified OK.

### 1.13.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{a-y}\right) \mathrm{d} y & =\left(\frac{1}{x^{2}}\right) \mathrm{d} x \\
\left(-\frac{1}{x^{2}}\right) \mathrm{d} x+\left(\frac{1}{a-y}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =-\frac{1}{x^{2}} \\
N(x, y) & =\frac{1}{a-y}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{x^{2}}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{a-y}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{1}{x^{2}} \mathrm{~d} x \\
\phi & =\frac{1}{x}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{a-y}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{a-y}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{a-y}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{a-y}\right) \mathrm{d} y \\
f(y) & =-\ln (a-y)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{1}{x}-\ln (a-y)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{1}{x}-\ln (a-y)
$$

The solution becomes

$$
y=-\mathrm{e}^{-\frac{c_{1} x-1}{x}}+a
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\mathrm{e}^{-\frac{c_{1} x-1}{x}}+a \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=-\mathrm{e}^{-\frac{c_{1} x-1}{x}}+a
$$

Verified OK.

### 1.13.5 Maple step by step solution

Let's solve

$$
x^{2} y^{\prime}+y=a
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables

$$
\frac{y^{\prime}}{a-y}=\frac{1}{x^{2}}
$$

- Integrate both sides with respect to $x$
$\int \frac{y^{\prime}}{a-y} d x=\int \frac{1}{x^{2}} d x+c_{1}$
- Evaluate integral

$$
-\ln (a-y)=-\frac{1}{x}+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=-\mathrm{e}^{-\frac{c_{1} x-1}{x}}+a
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 12

```
dsolve((y(x)-a)+x^2*diff(y(x),x)=0,y(x), singsol=all)
```

$$
y(x)=a+\mathrm{e}^{\frac{1}{x}} c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.048 (sec). Leaf size: 20
DSolve[(y[x]-a)+x^2*y'[x]==0,y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow a+c_{1} e^{\frac{1}{x}} \\
& y(x) \rightarrow a
\end{aligned}
$$

### 1.14 problem 14

1.14.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 156
1.14.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 157
1.14.3 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 158
1.14.4 Solving as first order ode lie symmetry lookup ode . . . . . . . 159
1.14.5 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 162
1.14.6 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 165

Internal problem ID [12431]
Internal file name [OUTPUT/11083_Monday_October_16_2023_09_47_21_PM_48193375/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 14.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
z-\left(-a^{2}+t^{2}\right) z^{\prime}=0
$$

### 1.14.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
z^{\prime} & =F(t, z) \\
& =f(t) g(z) \\
& =-\frac{z}{a^{2}-t^{2}}
\end{aligned}
$$

Where $f(t)=-\frac{1}{a^{2}-t^{2}}$ and $g(z)=z$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{z} d z & =-\frac{1}{a^{2}-t^{2}} d t \\
\int \frac{1}{z} d z & =\int-\frac{1}{a^{2}-t^{2}} d t \\
\ln (z) & =\frac{\ln (-a+t)}{2 a}-\frac{\ln (a+t)}{2 a}+c_{1} \\
z & =\mathrm{e}^{\frac{\ln (-a+t)}{2 a}-\frac{\ln (a+t)}{2 a}+c_{1}} \\
& =c_{1} \mathrm{e}^{\frac{\ln (-a+t)}{2 a}-\frac{\ln (a+t)}{2 a}}
\end{aligned}
$$

Which simplifies to

$$
z=c_{1}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
z=c_{1}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
z=c_{1}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}}
$$

Verified OK.

### 1.14.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
z^{\prime}+p(t) z=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =\frac{1}{a^{2}-t^{2}} \\
q(t) & =0
\end{aligned}
$$

Hence the ode is

$$
z^{\prime}+\frac{z}{a^{2}-t^{2}}=0
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \frac{1}{a^{2}-t^{2}} d t} \\
& =\mathrm{e}^{-\frac{\ln (-a+t)}{2 a}+\frac{\ln (a+t)}{2 a}}
\end{aligned}
$$

Which simplifies to

$$
\mu=(-a+t)^{-\frac{1}{2 a}}(a+t)^{\frac{1}{2 a}}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t} \mu z & =0 \\
\frac{\mathrm{~d}}{\mathrm{~d} t}\left((-a+t)^{-\frac{1}{2 a}}(a+t)^{\frac{1}{2 a}} z\right) & =0
\end{aligned}
$$

Integrating gives

$$
(-a+t)^{-\frac{1}{2 a}}(a+t)^{\frac{1}{2 a}} z=c_{1}
$$

Dividing both sides by the integrating factor $\mu=(-a+t)^{-\frac{1}{2 a}}(a+t)^{\frac{1}{2 a}}$ results in

$$
z=c_{1}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
z=c_{1}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
z=c_{1}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}}
$$

Verified OK.

### 1.14.3 Solving as homogeneousTypeD2 ode

Using the change of variables $z=u(t) t$ on the above ode results in new ode in $u(t)$

$$
u(t) t-\left(-a^{2}+t^{2}\right)\left(u^{\prime}(t) t+u(t)\right)=0
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(t, u) \\
& =f(t) g(u) \\
& =-\frac{u\left(a^{2}-t^{2}+t\right)}{\left(a^{2}-t^{2}\right) t}
\end{aligned}
$$

Where $f(t)=-\frac{a^{2}-t^{2}+t}{\left(a^{2}-t^{2}\right) t}$ and $g(u)=u$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{u} d u & =-\frac{a^{2}-t^{2}+t}{\left(a^{2}-t^{2}\right) t} d t \\
\int \frac{1}{u} d u & =\int-\frac{a^{2}-t^{2}+t}{\left(a^{2}-t^{2}\right) t} d t \\
\ln (u) & =\frac{\ln (-a+t)}{2 a}-\frac{\ln (a+t)}{2 a}-\ln (t)+c_{2} \\
u & =\mathrm{e}^{\frac{\ln (-a+t)}{2 a}-\frac{\ln (a+t)}{2 a}-\ln (t)+c_{2}} \\
& =c_{2} \mathrm{e}^{\frac{\ln (-a+t)}{2 a}-\frac{\ln (a+t)}{2 a}-\ln (t)}
\end{aligned}
$$

Which simplifies to

$$
u(t)=\frac{c_{2}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}}}{t}
$$

Therefore the solution $z$ is

$$
\begin{aligned}
z & =u t \\
& =c_{2}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
z=c_{2}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
z=c_{2}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}}
$$

Verified OK.

### 1.14.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
z^{\prime} & =-\frac{z}{a^{2}-t^{2}} \\
z^{\prime} & =\omega(t, z)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{z}-\xi_{t}\right)-\omega^{2} \xi_{z}-\omega_{t} \xi-\omega_{z} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 26: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, z)=0 \\
& \eta(t, z)=\mathrm{e}^{\frac{\ln (-a+t)}{2 a}-\frac{\ln (a+t)}{2 a}} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, z) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d z}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial z}\right) S(t, z)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{\frac{\ln (-a+t)}{2 a}-\frac{\ln (a+t)}{2 a}}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{\frac{-\ln (-a+t)}{2}+\frac{\ln (a+t)}{2}} z
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, z) S_{z}}{R_{t}+\omega(t, z) R_{z}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{z}, S_{t}, S_{z}$ are all partial derivatives and $\omega(t, z)$ is the right hand side of the original ode given by

$$
\omega(t, z)=-\frac{z}{a^{2}-t^{2}}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
& R_{t}=1 \\
& R_{z}=0 \\
& S_{t}=\frac{z\left(-(-a+t)^{-\frac{2 a+1}{2 a}}(a+t)^{\frac{1}{2 a}}+(a+t)^{-\frac{-1+2 a}{2 a}}(-a+t)^{-\frac{1}{2 a}}\right)}{2 a} \\
& S_{z}=(-a+t)^{-\frac{1}{2 a}}(a+t)^{\frac{1}{2 a}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.
$\frac{d S}{d R}=\frac{z\left(\left(-a^{2}+t^{2}\right)(a+t)^{\frac{1}{2 a}}(-a+t)^{-\frac{2 a+1}{2 a}}+\left((a+t)^{-\frac{-1+2 a}{2 a}}\left(a^{2}-t^{2}\right)-2(a+t)^{\frac{1}{2 a}} a\right)(-a+t)^{-\frac{1}{2 a}}\right)}{2 a^{3}-2 t^{2} a}$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, z$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{S(R)(a+R)^{\frac{1}{2 a}}(-a+R)^{-\frac{1}{2 a}}\left(\left(-R^{2}+a^{2}\right)(a+R)^{-\frac{1}{2 a}}(-a+R)^{-\frac{-1+2 a}{2 a}}+(-a+R)^{\frac{1}{2 a}}\left(\left(R^{2}-a^{2}\right)(c\right.\right.}{2 R^{2} a-2 a^{3}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, z$ coordinates. This results in

$$
(-a+t)^{-\frac{1}{2 a}}(a+t)^{\frac{1}{2 a}} z=c_{1}
$$

Which simplifies to

$$
(-a+t)^{-\frac{1}{2 a}}(a+t)^{\frac{1}{2 a}} z=c_{1}
$$

Which gives

$$
z=c_{1}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
z=c_{1}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
z=c_{1}(-a+t)^{\frac{1}{2 a}}(a+t)^{-\frac{1}{2 a}}
$$

Verified OK.

### 1.14.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, z) \mathrm{d} t+N(t, z) \mathrm{d} z=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(-\frac{1}{z}\right) \mathrm{d} z & =\left(\frac{1}{a^{2}-t^{2}}\right) \mathrm{d} t \\
\left(-\frac{1}{a^{2}-t^{2}}\right) \mathrm{d} t+\left(-\frac{1}{z}\right) \mathrm{d} z & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, z) & =-\frac{1}{a^{2}-t^{2}} \\
N(t, z) & =-\frac{1}{z}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial z}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial z} & =\frac{\partial}{\partial z}\left(-\frac{1}{a^{2}-t^{2}}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}\left(-\frac{1}{z}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial z}=\frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, z)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=M  \tag{1}\\
& \frac{\partial \phi}{\partial z}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int M \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\frac{1}{a^{2}-t^{2}} \mathrm{~d} t \\
\phi & =\frac{\ln (-a+t)-\ln (a+t)}{2 a}+f(z) \tag{3}
\end{align*}
$$

Where $f(z)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $z$. Taking derivative of equation (3) w.r.t $z$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial z}=0+f^{\prime}(z) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial z}=-\frac{1}{z}$. Therefore equation (4) becomes

$$
\begin{equation*}
-\frac{1}{z}=0+f^{\prime}(z) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(z)$ gives

$$
f^{\prime}(z)=-\frac{1}{z}
$$

Integrating the above w.r.t $z$ gives

$$
\begin{aligned}
\int f^{\prime}(z) \mathrm{d} z & =\int\left(-\frac{1}{z}\right) \mathrm{d} z \\
f(z) & =-\ln (z)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(z)$ into equation (3) gives $\phi$

$$
\phi=\frac{\ln (-a+t)-\ln (a+t)}{2 a}-\ln (z)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{\ln (-a+t)-\ln (a+t)}{2 a}-\ln (z)
$$

The solution becomes

$$
z=\mathrm{e}^{-\frac{2 c_{1} a+\ln (a+t)-\ln (-a+t)}{2 a}}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
z=\mathrm{e}^{-\frac{2 c_{1} a+\ln (a+t)-\ln (-a+t)}{2 a}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
z=\mathrm{e}^{-\frac{2 c_{1} a+\ln (a+t)-\ln (-a+t)}{2 a}}
$$

Verified OK.

### 1.14.6 Maple step by step solution

Let's solve

$$
z-\left(-a^{2}+t^{2}\right) z^{\prime}=0
$$

- Highest derivative means the order of the ODE is 1 $z^{\prime}$
- Separate variables

$$
\frac{z^{\prime}}{z}=\frac{1}{-a^{2}+t^{2}}
$$

- Integrate both sides with respect to $t$

$$
\int \frac{z^{\prime}}{z} d t=\int \frac{1}{-a^{2}+t^{2}} d t+c_{1}
$$

- Evaluate integral

$$
\ln (z)=\frac{\ln (-a+t)}{2 a}-\frac{\ln (a+t)}{2 a}+c_{1}
$$

- $\quad$ Solve for $z$

$$
z=\mathrm{e}^{-\frac{-2 c_{1} a+\ln (a+t)-\ln (-a+t)}{2 a}}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 26

```
dsolve(z(t)-(t~2-a^2)*diff(z(t),t)=0,z(t), singsol=all)
```

$$
z(t)=c_{1}(a+t)^{-\frac{1}{2 a}}(t-a)^{\frac{1}{2 a}}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.059 (sec). Leaf size: 26
DSolve[z[t]-(t^2-a^2)*z'[t]==0,z[t],t,IncludeSingularSolutions -> True]

$$
\begin{aligned}
& z(t) \rightarrow c_{1} e^{-\frac{\operatorname{arctanh}\left(\frac{t}{a}\right)}{a}} \\
& z(t) \rightarrow 0
\end{aligned}
$$

### 1.15 problem 15

1.15.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 167
1.15.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 169
1.15.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 173
1.15.4 Solving as riccati ode . . . . . . . . . . . . . . . . . . . . . . . . 177
1.15.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 179

Internal problem ID [12432]
Internal file name [OUTPUT/11084_Monday_October_16_2023_09_47_21_PM_35154629/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 15.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "first_order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
y^{\prime}-\frac{1+y^{2}}{x^{2}+1}=0
$$

### 1.15.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{y^{2}+1}{x^{2}+1}
\end{aligned}
$$

Where $f(x)=\frac{1}{x^{2}+1}$ and $g(y)=y^{2}+1$. Integrating both sides gives

$$
\frac{1}{y^{2}+1} d y=\frac{1}{x^{2}+1} d x
$$

$$
\begin{aligned}
\int \frac{1}{y^{2}+1} d y & =\int \frac{1}{x^{2}+1} d x \\
\arctan (y) & =\arctan (x)+c_{1}
\end{aligned}
$$

Which results in

$$
y=\tan \left(\arctan (x)+c_{1}\right)
$$

## Summary

The solution(s) found are the following


Figure 21: Slope field plot

Verification of solutions

$$
y=\tan \left(\arctan (x)+c_{1}\right)
$$

Verified OK.

### 1.15.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{y^{2}+1}{x^{2}+1} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the $\operatorname{PDE}$ (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 29: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | special | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |  |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}$ |
| $a_{1} b_{2}-a_{2} b_{1}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |  |  |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=x^{2}+1 \\
& \eta(x, y)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=y
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d x \\
& =\int \frac{1}{x^{2}+1} d x
\end{aligned}
$$

Which results in

$$
S=\arctan (x)
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{y^{2}+1}{x^{2}+1}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =0 \\
R_{y} & =1 \\
S_{x} & =\frac{1}{x^{2}+1} \\
S_{y} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{1}{y^{2}+1} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{1}{R^{2}+1}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\arctan (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\arctan (x)=\arctan (y)+c_{1}
$$

Which simplifies to

$$
\arctan (x)=\arctan (y)+c_{1}
$$

Which gives

$$
y=-\tan \left(-\arctan (x)+c_{1}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown．

| Original ode in $x, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{y^{2}+1}{x^{2}+1}$ |  | $\frac{d S}{d R}=\frac{1}{R^{2}+1}$ |
|  |  | $\rightarrow \rightarrow \rightarrow \rightarrow$－ $\mid$ 他 |
|  |  |  |
|  |  | $\rightarrow \rightarrow \rightarrow- \pm$ 何 |
| $\rightarrow \rightarrow-\infty$－ |  | $\rightarrow \rightarrow$ 为 $⿻ 上 丨 匕 力$ |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \infty]{ }$ |  | $\rightarrow \rightarrow+\infty$ |
| $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-2}$ |  |  |
|  | $S=\arctan (x)$ | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow-\infty]{ }$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  | 为 |

## Summary

The solution（s）found are the following

$$
\begin{equation*}
y=-\tan \left(-\arctan (x)+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 22: Slope field plot

## Verification of solutions

$$
y=-\tan \left(-\arctan (x)+c_{1}\right)
$$

Verified OK.

### 1.15.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{y^{2}+1}\right) \mathrm{d} y & =\left(\frac{1}{x^{2}+1}\right) \mathrm{d} x \\
\left(-\frac{1}{x^{2}+1}\right) \mathrm{d} x+\left(\frac{1}{y^{2}+1}\right) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =-\frac{1}{x^{2}+1} \\
N(x, y) & =\frac{1}{y^{2}+1}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{x^{2}+1}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{y^{2}+1}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{1}{x^{2}+1} \mathrm{~d} x \\
\phi & =-\arctan (x)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{y^{2}+1}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{y^{2}+1}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{y^{2}+1}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{y^{2}+1}\right) \mathrm{d} y \\
f(y) & =\arctan (y)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\arctan (x)+\arctan (y)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\arctan (x)+\arctan (y)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-\arctan (x)+\arctan (y)=c_{1} \tag{1}
\end{equation*}
$$



Figure 23: Slope field plot

Verification of solutions

$$
-\arctan (x)+\arctan (y)=c_{1}
$$

## Verified OK.

### 1.15.4 Solving as riccati ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{y^{2}+1}{x^{2}+1}
\end{aligned}
$$

This is a Riccati ODE. Comparing the ODE to solve

$$
y^{\prime}=\frac{y^{2}}{x^{2}+1}+\frac{1}{x^{2}+1}
$$

With Riccati ODE standard form

$$
y^{\prime}=f_{0}(x)+f_{1}(x) y+f_{2}(x) y^{2}
$$

Shows that $f_{0}(x)=\frac{1}{x^{2}+1}, f_{1}(x)=0$ and $f_{2}(x)=\frac{1}{x^{2}+1}$. Let

$$
\begin{align*}
y & =\frac{-u^{\prime}}{f_{2} u} \\
& =\frac{-u^{\prime}}{\frac{u}{x^{2}+1}} \tag{1}
\end{align*}
$$

Using the above substitution in the given ODE results (after some simplification)in a second order ODE to solve for $u(x)$ which is

$$
\begin{equation*}
f_{2} u^{\prime \prime}(x)-\left(f_{2}^{\prime}+f_{1} f_{2}\right) u^{\prime}(x)+f_{2}^{2} f_{0} u(x)=0 \tag{2}
\end{equation*}
$$

But

$$
\begin{aligned}
f_{2}^{\prime} & =-\frac{2 x}{\left(x^{2}+1\right)^{2}} \\
f_{1} f_{2} & =0 \\
f_{2}^{2} f_{0} & =\frac{1}{\left(x^{2}+1\right)^{3}}
\end{aligned}
$$

Substituting the above terms back in equation (2) gives

$$
\frac{u^{\prime \prime}(x)}{x^{2}+1}+\frac{2 x u^{\prime}(x)}{\left(x^{2}+1\right)^{2}}+\frac{u(x)}{\left(x^{2}+1\right)^{3}}=0
$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$
u(x)=\frac{c_{1} x+c_{2}}{\sqrt{x^{2}+1}}
$$

The above shows that

$$
u^{\prime}(x)=\frac{-c_{2} x+c_{1}}{\left(x^{2}+1\right)^{\frac{3}{2}}}
$$

Using the above in (1) gives the solution

$$
y=-\frac{-c_{2} x+c_{1}}{c_{1} x+c_{2}}
$$

Dividing both numerator and denominator by $c_{1}$ gives, after renaming the constant $\frac{c_{2}}{c_{1}}=c_{3}$ the following solution

$$
y=\frac{-c_{3}+x}{c_{3} x+1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{-c_{3}+x}{c_{3} x+1} \tag{1}
\end{equation*}
$$



Figure 24: Slope field plot

Verification of solutions

$$
y=\frac{-c_{3}+x}{c_{3} x+1}
$$

Verified OK.

### 1.15.5 Maple step by step solution

Let's solve
$y^{\prime}-\frac{1+y^{2}}{x^{2}+1}=0$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables
$\frac{y^{\prime}}{1+y^{2}}=\frac{1}{x^{2}+1}$
- Integrate both sides with respect to $x$
$\int \frac{y^{\prime}}{1+y^{2}} d x=\int \frac{1}{x^{2}+1} d x+c_{1}$
- Evaluate integral
$\arctan (y)=\arctan (x)+c_{1}$
- $\quad$ Solve for $y$

$$
y=\tan \left(\arctan (x)+c_{1}\right)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 9
dsolve(diff $(y(x), x)=(1+y(x) \wedge 2) /\left(1+x^{\wedge} 2\right), y(x)$, singsol=all)

$$
y(x)=\tan \left(\arctan (x)+c_{1}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.408 (sec). Leaf size: 25
DSolve $\left[y^{\prime}[x]==(1+y[x] \sim 2) /\left(1+x^{\wedge} 2\right), y[x], x\right.$, IncludeSingularSolutions $->$ True]

$$
\begin{aligned}
& y(x) \rightarrow \tan \left(\arctan (x)+c_{1}\right) \\
& y(x) \rightarrow-i \\
& y(x) \rightarrow i
\end{aligned}
$$

### 1.16 problem 16

$$
\text { 1.16.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . } 181
$$

1.16.2 Solving as first order ode lie symmetry lookup ode ..... 183
1.16.3 Solving as exact ode ..... 187
1.16.4 Solving as riccati ode ..... 191
1.16.5 Maple step by step solution ..... 193

Internal problem ID [12433]
Internal file name [OUTPUT/11085_Monday_October_16_2023_09_47_23_PM_8027738/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 16.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "first_order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
s^{2}-\sqrt{t} s^{\prime}=-1
$$

### 1.16.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
s^{\prime} & =F(t, s) \\
& =f(t) g(s) \\
& =\frac{s^{2}+1}{\sqrt{t}}
\end{aligned}
$$

Where $f(t)=\frac{1}{\sqrt{ } t}$ and $g(s)=s^{2}+1$. Integrating both sides gives

$$
\frac{1}{s^{2}+1} d s=\frac{1}{\sqrt{t}} d t
$$

$$
\begin{aligned}
\int \frac{1}{s^{2}+1} d s & =\int \frac{1}{\sqrt{t}} d t \\
\arctan (s) & =2 \sqrt{t}+c_{1}
\end{aligned}
$$

Which results in

$$
s=\tan \left(2 \sqrt{t}+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
s=\tan \left(2 \sqrt{t}+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 25: Slope field plot

Verification of solutions

$$
s=\tan \left(2 \sqrt{t}+c_{1}\right)
$$

Verified OK.

### 1.16.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& s^{\prime}=\frac{s^{2}+1}{\sqrt{t}} \\
& s^{\prime}=\omega(t, s)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{s}-\xi_{t}\right)-\omega^{2} \xi_{s}-\omega_{t} \xi-\omega_{s} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 32: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, s)=\sqrt{t} \\
& \eta(t, s)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, s) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d s}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial s}\right) S(t, s)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=s
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d t \\
& =\int \frac{1}{\sqrt{t}} d t
\end{aligned}
$$

Which results in

$$
S=2 \sqrt{t}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, s) S_{s}}{R_{t}+\omega(t, s) R_{s}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{s}, S_{t}, S_{s}$ are all partial derivatives and $\omega(t, s)$ is the right hand side of the original ode given by

$$
\omega(t, s)=\frac{s^{2}+1}{\sqrt{t}}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =0 \\
R_{s} & =1 \\
S_{t} & =\frac{1}{\sqrt{t}} \\
S_{s} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{1}{s^{2}+1} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, s$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{1}{R^{2}+1}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\arctan (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, s$ coordinates. This results in

$$
2 \sqrt{t}=\arctan (s)+c_{1}
$$

Which simplifies to

$$
2 \sqrt{t}=\arctan (s)+c_{1}
$$

Which gives

$$
s=-\tan \left(-2 \sqrt{t}+c_{1}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, s$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d s}{d t}=\frac{s^{2}+1}{\sqrt{t}}$ |  | $\frac{d S}{d R}=\frac{1}{R^{2}+1}$ |
| + $\uparrow \uparrow \uparrow \uparrow$ |  | $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ Nオオ |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow 0]{ }$ |
| s(t) |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+]{ }+$ |
|  |  |  |
|  | $R=s$ | $\rightarrow \rightarrow \rightarrow+\infty$ |
|  | $S=2 \sqrt{t}$ |  |
|  |  | $\rightarrow \rightarrow+$ |
|  |  | $\rightarrow$ |
|  |  | $\rightarrow \rightarrow$ |
| ¢ ¢ ¢ ¢ ¢ ¢ ¢ $\uparrow$ |  | $\rightarrow \rightarrow \rightarrow \rightarrow \infty$ |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
s=-\tan \left(-2 \sqrt{t}+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 26: Slope field plot

Verification of solutions

$$
s=-\tan \left(-2 \sqrt{t}+c_{1}\right)
$$

Verified OK.

### 1.16.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, s) \mathrm{d} t+N(t, s) \mathrm{d} s=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{s^{2}+1}\right) \mathrm{d} s & =\left(\frac{1}{\sqrt{t}}\right) \mathrm{d} t \\
\left(-\frac{1}{\sqrt{t}}\right) \mathrm{d} t+\left(\frac{1}{s^{2}+1}\right) \mathrm{d} s & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(t, s)=-\frac{1}{\sqrt{t}} \\
& N(t, s)=\frac{1}{s^{2}+1}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial s}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial s} & =\frac{\partial}{\partial s}\left(-\frac{1}{\sqrt{t}}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}\left(\frac{1}{s^{2}+1}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial s}=\frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, s)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=M  \tag{1}\\
& \frac{\partial \phi}{\partial s}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int M \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\frac{1}{\sqrt{t}} \mathrm{~d} t \\
\phi & =-2 \sqrt{t}+f(s) \tag{3}
\end{align*}
$$

Where $f(s)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $s$. Taking derivative of equation (3) w.r.t $s$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial s}=0+f^{\prime}(s) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial s}=\frac{1}{s^{2}+1}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{s^{2}+1}=0+f^{\prime}(s) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(s)$ gives

$$
f^{\prime}(s)=\frac{1}{s^{2}+1}
$$

Integrating the above w.r.t $s$ gives

$$
\begin{aligned}
\int f^{\prime}(s) \mathrm{d} s & =\int\left(\frac{1}{s^{2}+1}\right) \mathrm{d} s \\
f(s) & =\arctan (s)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(s)$ into equation (3) gives $\phi$

$$
\phi=-2 \sqrt{t}+\arctan (s)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-2 \sqrt{t}+\arctan (s)
$$

The solution becomes

$$
s=\tan \left(2 \sqrt{t}+c_{1}\right)
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
s=\tan \left(2 \sqrt{t}+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 27: Slope field plot

## Verification of solutions

$$
s=\tan \left(2 \sqrt{t}+c_{1}\right)
$$

Verified OK.

### 1.16.4 Solving as riccati ode

In canonical form the ODE is

$$
\begin{aligned}
s^{\prime} & =F(t, s) \\
& =\frac{s^{2}+1}{\sqrt{t}}
\end{aligned}
$$

This is a Riccati ODE. Comparing the ODE to solve

$$
s^{\prime}=\frac{s^{2}}{\sqrt{t}}+\frac{1}{\sqrt{t}}
$$

With Riccati ODE standard form

$$
s^{\prime}=f_{0}(t)+f_{1}(t) s+f_{2}(t) s^{2}
$$

Shows that $f_{0}(t)=\frac{1}{\sqrt{t}}, f_{1}(t)=0$ and $f_{2}(t)=\frac{1}{\sqrt{t}}$. Let

$$
\begin{align*}
s & =\frac{-u^{\prime}}{f_{2} u} \\
& =\frac{-u^{\prime}}{\frac{u}{\sqrt{t}}} \tag{1}
\end{align*}
$$

Using the above substitution in the given ODE results (after some simplification)in a second order ODE to solve for $u(x)$ which is

$$
\begin{equation*}
f_{2} u^{\prime \prime}(t)-\left(f_{2}^{\prime}+f_{1} f_{2}\right) u^{\prime}(t)+f_{2}^{2} f_{0} u(t)=0 \tag{2}
\end{equation*}
$$

But

$$
\begin{aligned}
f_{2}^{\prime} & =-\frac{1}{2 t^{\frac{3}{2}}} \\
f_{1} f_{2} & =0 \\
f_{2}^{2} f_{0} & =\frac{1}{t^{\frac{3}{2}}}
\end{aligned}
$$

Substituting the above terms back in equation (2) gives

$$
\frac{u^{\prime \prime}(t)}{\sqrt{t}}+\frac{u^{\prime}(t)}{2 t^{\frac{3}{2}}}+\frac{u(t)}{t^{\frac{3}{2}}}=0
$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$
u(t)=c_{1} \sin (2 \sqrt{t})+c_{2} \cos (2 \sqrt{t})
$$

The above shows that

$$
u^{\prime}(t)=\frac{c_{1} \cos (2 \sqrt{t})-c_{2} \sin (2 \sqrt{t})}{\sqrt{t}}
$$

Using the above in (1) gives the solution

$$
s=-\frac{c_{1} \cos (2 \sqrt{t})-c_{2} \sin (2 \sqrt{t})}{c_{1} \sin (2 \sqrt{t})+c_{2} \cos (2 \sqrt{t})}
$$

Dividing both numerator and denominator by $c_{1}$ gives, after renaming the constant $\frac{c_{2}}{c_{1}}=c_{3}$ the following solution

$$
s=\frac{-c_{3} \cos (2 \sqrt{t})+\sin (2 \sqrt{t})}{c_{3} \sin (2 \sqrt{t})+\cos (2 \sqrt{t})}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
s=\frac{-c_{3} \cos (2 \sqrt{t})+\sin (2 \sqrt{t})}{c_{3} \sin (2 \sqrt{t})+\cos (2 \sqrt{t})} \tag{1}
\end{equation*}
$$



Figure 28: Slope field plot

Verification of solutions

$$
s=\frac{-c_{3} \cos (2 \sqrt{t})+\sin (2 \sqrt{t})}{c_{3} \sin (2 \sqrt{t})+\cos (2 \sqrt{t})}
$$

Verified OK.

### 1.16.5 Maple step by step solution

Let's solve

$$
s^{2}-\sqrt{t} s^{\prime}=-1
$$

- Highest derivative means the order of the ODE is 1 $s^{\prime}$
- Separate variables

$$
\frac{s^{\prime}}{-1-s^{2}}=-\frac{1}{\sqrt{t}}
$$

- Integrate both sides with respect to $t$

$$
\int \frac{s^{\prime}}{-1-s^{2}} d t=\int-\frac{1}{\sqrt{t}} d t+c_{1}
$$

- Evaluate integral

$$
-\arctan (s)=-2 \sqrt{t}+c_{1}
$$

- $\quad$ Solve for $s$

$$
s=-\tan \left(-2 \sqrt{t}+c_{1}\right)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 12

```
dsolve((1+s(t)~2)-sqrt(t)*diff(s(t),t)=0,s(t), singsol=all)
```

$$
s(t)=\tan \left(2 \sqrt{t}+c_{1}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.332 (sec). Leaf size: 30
DSolve[(1+s[t]~2)-Sqrt[t]*s'[t]==0,s[t],t,IncludeSingularSolutions -> True]

$$
\begin{aligned}
& s(t) \rightarrow \tan \left(2 \sqrt{t}+c_{1}\right) \\
& s(t) \rightarrow-i \\
& s(t) \rightarrow i
\end{aligned}
$$

### 1.17 problem 17

1.17.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 195
1.17.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 197
1.17.3 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 198
1.17.4 Solving as first order ode lie symmetry lookup ode . . . . . . . 200
1.17.5 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 204
1.17.6 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 208

Internal problem ID [12434]
Internal file name [OUTPUT/11086_Monday_October_16_2023_09_47_23_PM_19040029/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 17.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
r^{\prime}+r \tan (t)=0
$$

### 1.17.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
r^{\prime} & =F(t, r) \\
& =f(t) g(r) \\
& =-r \tan (t)
\end{aligned}
$$

Where $f(t)=-\tan (t)$ and $g(r)=r$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{r} d r & =-\tan (t) d t \\
\int \frac{1}{r} d r & =\int-\tan (t) d t \\
\ln (r) & =\ln (\cos (t))+c_{1} \\
r & =\mathrm{e}^{\ln (\cos (t))+c_{1}} \\
& =c_{1} \cos (t)
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
r=c_{1} \cos (t) \tag{1}
\end{equation*}
$$



Figure 29: Slope field plot

Verification of solutions

$$
r=c_{1} \cos (t)
$$

Verified OK.

### 1.17.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
r^{\prime}+p(t) r=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =\tan (t) \\
q(t) & =0
\end{aligned}
$$

Hence the ode is

$$
r^{\prime}+r \tan (t)=0
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \tan (t) d t} \\
& =\frac{1}{\cos (t)}
\end{aligned}
$$

Which simplifies to

$$
\mu=\sec (t)
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t} \mu r & =0 \\
\frac{\mathrm{~d}}{\mathrm{~d} t}(\sec (t) r) & =0
\end{aligned}
$$

Integrating gives

$$
\sec (t) r=c_{1}
$$

Dividing both sides by the integrating factor $\mu=\sec (t)$ results in

$$
r=c_{1} \cos (t)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
r=c_{1} \cos (t) \tag{1}
\end{equation*}
$$



Figure 30: Slope field plot

Verification of solutions

$$
r=c_{1} \cos (t)
$$

Verified OK.

### 1.17.3 Solving as homogeneousTypeD2 ode

Using the change of variables $r=u(t) t$ on the above ode results in new ode in $u(t)$

$$
u^{\prime}(t) t+u(t)+u(t) t \tan (t)=0
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(t, u) \\
& =f(t) g(u) \\
& =-\frac{u(\tan (t) t+1)}{t}
\end{aligned}
$$

Where $f(t)=-\frac{\tan (t) t+1}{t}$ and $g(u)=u$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{u} d u & =-\frac{\tan (t) t+1}{t} d t \\
\int \frac{1}{u} d u & =\int-\frac{\tan (t) t+1}{t} d t \\
\ln (u) & =\ln (\cos (t))-\ln (t)+c_{2} \\
u & =\mathrm{e}^{\ln (\cos (t))-\ln (t)+c_{2}} \\
& =c_{2} \mathrm{e}^{\ln (\cos (t))-\ln (t)}
\end{aligned}
$$

Which simplifies to

$$
u(t)=\frac{c_{2} \cos (t)}{t}
$$

Therefore the solution $r$ is

$$
\begin{aligned}
r & =t u \\
& =c_{2} \cos (t)
\end{aligned}
$$

## Summary

The solution(s) found are the following


Figure 31: Slope field plot

Verification of solutions

$$
r=c_{2} \cos (t)
$$

Verified OK.

### 1.17.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
r^{\prime} & =-r \tan (t) \\
r^{\prime} & =\omega(t, r)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{r}-\xi_{t}\right)-\omega^{2} \xi_{r}-\omega_{t} \xi-\omega_{r} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 35: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, r)=0 \\
& \eta(t, r)=\cos (t) \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, r) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d r}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial r}\right) S(t, r)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\cos (t)} d y
\end{aligned}
$$

Which results in

$$
S=\frac{r}{\cos (t)}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, r) S_{r}}{R_{t}+\omega(t, r) R_{r}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{r}, S_{t}, S_{r}$ are all partial derivatives and $\omega(t, r)$ is the right hand side of the original ode given by

$$
\omega(t, r)=-r \tan (t)
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{r} & =0 \\
S_{t} & =\sec (t) \tan (t) r \\
S_{r} & =\sec (t)
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, r$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, r$ coordinates. This results in

$$
\sec (t) r=c_{1}
$$

Which simplifies to

$$
\sec (t) r=c_{1}
$$

Which gives

$$
r=\frac{c_{1}}{\sec (t)}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, r$ coordinates | Canonical <br> coordinates <br> transformation | ODE in canonical coordinates <br> $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d r}{d t}=-r \tan (t)$ |  | $\frac{d S}{d R}=0$ |
| 性 1 分 |  |  |

Summary
The solution(s) found are the following

$$
\begin{equation*}
r=\frac{c_{1}}{\sec (t)} \tag{1}
\end{equation*}
$$



Figure 32: Slope field plot

## Verification of solutions

$$
r=\frac{c_{1}}{\sec (t)}
$$

Verified OK.

### 1.17.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, r) \mathrm{d} t+N(t, r) \mathrm{d} r=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(-\frac{1}{r}\right) \mathrm{d} r & =(\tan (t)) \mathrm{d} t \\
(-\tan (t)) \mathrm{d} t+\left(-\frac{1}{r}\right) \mathrm{d} r & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, r) & =-\tan (t) \\
N(t, r) & =-\frac{1}{r}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial r}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial r} & =\frac{\partial}{\partial r}(-\tan (t)) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}\left(-\frac{1}{r}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial r}=\frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, r)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=M  \tag{1}\\
& \frac{\partial \phi}{\partial r}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int M \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int-\tan (t) \mathrm{d} t \\
\phi & =\ln (\cos (t))+f(r) \tag{3}
\end{align*}
$$

Where $f(r)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $r$. Taking derivative of equation (3) w.r.t $r$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial r}=0+f^{\prime}(r) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial r}=-\frac{1}{r}$. Therefore equation (4) becomes

$$
\begin{equation*}
-\frac{1}{r}=0+f^{\prime}(r) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(r)$ gives

$$
f^{\prime}(r)=-\frac{1}{r}
$$

Integrating the above w.r.t $r$ gives

$$
\begin{aligned}
\int f^{\prime}(r) \mathrm{d} r & =\int\left(-\frac{1}{r}\right) \mathrm{d} r \\
f(r) & =-\ln (r)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(r)$ into equation (3) gives $\phi$

$$
\phi=\ln (\cos (t))-\ln (r)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\ln (\cos (t))-\ln (r)
$$

The solution becomes

$$
r=\mathrm{e}^{-c_{1}} \cos (t)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
r=\mathrm{e}^{-c_{1}} \cos (t) \tag{1}
\end{equation*}
$$



Figure 33: Slope field plot
Verification of solutions

$$
r=\mathrm{e}^{-c_{1}} \cos (t)
$$

Verified OK.

### 1.17.6 Maple step by step solution

Let's solve

$$
r^{\prime}+r \tan (t)=0
$$

- Highest derivative means the order of the ODE is 1

$$
r^{\prime}
$$

- Separate variables

$$
\frac{r^{\prime}}{r}=-\tan (t)
$$

- Integrate both sides with respect to $t$

$$
\int \frac{r^{\prime}}{r} d t=\int-\tan (t) d t+c_{1}
$$

- Evaluate integral

$$
\ln (r)=\ln (\cos (t))+c_{1}
$$

- $\quad$ Solve for $r$

$$
r=\mathrm{e}^{c_{1}} \cos (t)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 8

```
dsolve(diff(r(t),t)+r(t)*tan(t)=0,r(t), singsol=all)
```

$$
r(t)=\cos (t) c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.068 (sec). Leaf size: 15
DSolve[r' $[\mathrm{t}]+\mathrm{r}[\mathrm{t}] * \operatorname{Tan}[\mathrm{t}]==0, \mathrm{r}[\mathrm{t}], \mathrm{t}$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& r(t) \rightarrow c_{1} \cos (t) \\
& r(t) \rightarrow 0
\end{aligned}
$$

### 1.18 problem 21

1.18.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 210
1.18.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 212
1.18.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 216
1.18.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 220

Internal problem ID [12435]
Internal file name [OUTPUT/11087_Monday_October_16_2023_09_47_24_PM_16412134/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 21.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first__order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
\left(x^{2}+1\right) y^{\prime}-\sqrt{1-y^{2}}=0
$$

### 1.18.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{\sqrt{-y^{2}+1}}{x^{2}+1}
\end{aligned}
$$

Where $f(x)=\frac{1}{x^{2}+1}$ and $g(y)=\sqrt{-y^{2}+1}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\sqrt{-y^{2}+1}} d y & =\frac{1}{x^{2}+1} d x \\
\int \frac{1}{\sqrt{-y^{2}+1}} d y & =\int \frac{1}{x^{2}+1} d x \\
\arcsin (y) & =\arctan (x)+c_{1}
\end{aligned}
$$

Which results in

$$
y=\sin \left(\arctan (x)+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\sin \left(\arctan (x)+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 34: Slope field plot

Verification of solutions

$$
y=\sin \left(\arctan (x)+c_{1}\right)
$$

Verified OK.

### 1.18.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{\sqrt{-y^{2}+1}}{x^{2}+1} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the $\operatorname{PDE}(\mathrm{A})$, and can just use the lookup table shown below to find $\xi, \eta$

Table 38: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | special | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |  |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}$ |
| $a_{1} b_{2}-a_{2} b_{1}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |  |  |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=x^{2}+1 \\
& \eta(x, y)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=y
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d x \\
& =\int \frac{1}{x^{2}+1} d x
\end{aligned}
$$

Which results in

$$
S=\arctan (x)
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{\sqrt{-y^{2}+1}}{x^{2}+1}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =0 \\
R_{y} & =1 \\
S_{x} & =\frac{1}{x^{2}+1} \\
S_{y} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{1}{\sqrt{-y^{2}+1}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{1}{\sqrt{-R^{2}+1}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\arcsin (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\arctan (x)=\arcsin (y)+c_{1}
$$

Which simplifies to

$$
\arctan (x)=\arcsin (y)+c_{1}
$$

Which gives

$$
y=-\sin \left(-\arctan (x)+c_{1}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
|  | $\begin{aligned} R & =y \\ S & =\arctan (x) \end{aligned}$ |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=-\sin \left(-\arctan (x)+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 35: Slope field plot
Verification of solutions

$$
y=-\sin \left(-\arctan (x)+c_{1}\right)
$$

Verified OK.

### 1.18.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{\sqrt{-y^{2}+1}}\right) \mathrm{d} y & =\left(\frac{1}{x^{2}+1}\right) \mathrm{d} x \\
\left(-\frac{1}{x^{2}+1}\right) \mathrm{d} x+\left(\frac{1}{\sqrt{-y^{2}+1}}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=-\frac{1}{x^{2}+1} \\
& N(x, y)=\frac{1}{\sqrt{-y^{2}+1}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{x^{2}+1}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{\sqrt{-y^{2}+1}}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{1}{x^{2}+1} \mathrm{~d} x \\
\phi & =-\arctan (x)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{\sqrt{-y^{2}+1}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{\sqrt{-y^{2}+1}}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{\sqrt{-y^{2}+1}}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{\sqrt{-y^{2}+1}}\right) \mathrm{d} y \\
f(y) & =\arcsin (y)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\arctan (x)+\arcsin (y)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\arctan (x)+\arcsin (y)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-\arctan (x)+\arcsin (y)=c_{1} \tag{1}
\end{equation*}
$$



Figure 36: Slope field plot

Verification of solutions

$$
-\arctan (x)+\arcsin (y)=c_{1}
$$

Verified OK.

### 1.18.4 Maple step by step solution

Let's solve

$$
\left(x^{2}+1\right) y^{\prime}-\sqrt{1-y^{2}}=0
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables

$$
\frac{y^{\prime}}{\sqrt{1-y^{2}}}=\frac{1}{x^{2}+1}
$$

- Integrate both sides with respect to $x$
$\int \frac{y^{\prime}}{\sqrt{1-y^{2}}} d x=\int \frac{1}{x^{2}+1} d x+c_{1}$
- Evaluate integral
$\arcsin (y)=\arctan (x)+c_{1}$
- $\quad$ Solve for $y$
$y=\sin \left(\arctan (x)+c_{1}\right)$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 9
dsolve $\left(\left(1+x^{\wedge} 2\right) * \operatorname{diff}(y(x), x)-\operatorname{sqrt}\left(1-y(x)^{\wedge} 2\right)=0, y(x)\right.$, singsol=all)

$$
y(x)=\sin \left(\arctan (x)+c_{1}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.479 (sec). Leaf size: 29
DSolve[(1+x~2)*y'[x]-Sqrt[1-y[x]~2]==0,y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow \cos \left(\arctan (x)+c_{1}\right) \\
& y(x) \rightarrow-1 \\
& y(x) \rightarrow 1 \\
& y(x) \rightarrow \text { Interval }[\{-1,1\}]
\end{aligned}
$$

### 1.19 problem 22

$$
\text { 1.19.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . } 222
$$

1.19.2 Solving as first order ode lie symmetry lookup ode ..... 224
1.19.3 Solving as exact ode ..... 228
1.19.4 Maple step by step solution ..... 232

Internal problem ID [12436]
Internal file name [OUTPUT/11088_Monday_October_16_2023_09_47_25_PM_47365630/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 22.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first__order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
\sqrt{-x^{2}+1} y^{\prime}-\sqrt{1-y^{2}}=0
$$

### 1.19.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{\sqrt{-y^{2}+1}}{\sqrt{-x^{2}+1}}
\end{aligned}
$$

Where $f(x)=\frac{1}{\sqrt{-x^{2}+1}}$ and $g(y)=\sqrt{-y^{2}+1}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\sqrt{-y^{2}+1}} d y & =\frac{1}{\sqrt{-x^{2}+1}} d x \\
\int \frac{1}{\sqrt{-y^{2}+1}} d y & =\int \frac{1}{\sqrt{-x^{2}+1}} d x
\end{aligned}
$$

$$
\arcsin (y)=-\frac{\sqrt{-(x-1)^{2}-2 x+2}}{2}+\arcsin (x)+\frac{\sqrt{-(x+1)^{2}+2 x+2}}{2}+c_{1}
$$

Which results in

$$
y=\sin \left(\arcsin (x)+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\sin \left(\arcsin (x)+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 37: Slope field plot
Verification of solutions

$$
y=\sin \left(\arcsin (x)+c_{1}\right)
$$

## Verified OK.

### 1.19.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{\sqrt{-y^{2}+1}}{\sqrt{-x^{2}+1}} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the $\operatorname{PDE}(\mathrm{A})$, and can just use the lookup table shown below to find $\xi, \eta$

Table 41: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=\sqrt{-x^{2}+1} \\
& \eta(x, y)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=y
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d x \\
& =\int \frac{1}{\sqrt{-x^{2}+1}} d x
\end{aligned}
$$

Which results in

$$
S=\arcsin (x)
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{\sqrt{-y^{2}+1}}{\sqrt{-x^{2}+1}}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =0 \\
R_{y} & =1 \\
S_{x} & =\frac{1}{\sqrt{-x^{2}+1}} \\
S_{y} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{1}{\sqrt{-y^{2}+1}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{1}{\sqrt{-R^{2}+1}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\arcsin (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\arcsin (x)=\arcsin (y)+c_{1}
$$

Which simplifies to

$$
\arcsin (x)=\arcsin (y)+c_{1}
$$

Which gives

$$
y=-\sin \left(-\arcsin (x)+c_{1}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.


## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=-\sin \left(-\arcsin (x)+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 38: Slope field plot

Verification of solutions

$$
y=-\sin \left(-\arcsin (x)+c_{1}\right)
$$

Verified OK.

### 1.19.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{\sqrt{-y^{2}+1}}\right) \mathrm{d} y & =\left(\frac{1}{\sqrt{-x^{2}+1}}\right) \mathrm{d} x \\
\left(-\frac{1}{\sqrt{-x^{2}+1}}\right) \mathrm{d} x+\left(\frac{1}{\sqrt{-y^{2}+1}}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=-\frac{1}{\sqrt{-x^{2}+1}} \\
& N(x, y)=\frac{1}{\sqrt{-y^{2}+1}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{\sqrt{-x^{2}+1}}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{\sqrt{-y^{2}+1}}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{1}{\sqrt{-x^{2}+1}} \mathrm{~d} x \\
\phi & =-\arcsin (x)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{\sqrt{-y^{2}+1}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{\sqrt{-y^{2}+1}}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{\sqrt{-y^{2}+1}}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{\sqrt{-y^{2}+1}}\right) \mathrm{d} y \\
f(y) & =\arcsin (y)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\arcsin (x)+\arcsin (y)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\arcsin (x)+\arcsin (y)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-\arcsin (x)+\arcsin (y)=c_{1} \tag{1}
\end{equation*}
$$



Figure 39: Slope field plot

Verification of solutions

$$
-\arcsin (x)+\arcsin (y)=c_{1}
$$

Verified OK.

### 1.19.4 Maple step by step solution

Let's solve
$\sqrt{-x^{2}+1} y^{\prime}-\sqrt{1-y^{2}}=0$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables

$$
\frac{y^{\prime}}{\sqrt{1-y^{2}}}=\frac{1}{\sqrt{-x^{2}+1}}
$$

- Integrate both sides with respect to $x$
$\int \frac{y^{\prime}}{\sqrt{1-y^{2}}} d x=\int \frac{1}{\sqrt{-x^{2}+1}} d x+c_{1}$
- Evaluate integral

$$
\arcsin (y)=\arcsin (x)+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\sin \left(\arcsin (x)+c_{1}\right)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 9

```
dsolve(sqrt(1-x^2)*diff(y(x),x)-sqrt(1-y(x)^2)=0,y(x), singsol=all)
```

$$
y(x)=\sin \left(\arcsin (x)+c_{1}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.488 (sec). Leaf size: 49
DSolve[Sqrt [1- $\left.\mathrm{x}^{\wedge} 2\right] * \mathrm{y}^{\prime}[\mathrm{x}]$-Sqrt[1-y[x]~2]=$=0, \mathrm{y}[\mathrm{x}], \mathrm{x}$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow \cos \left(2 \arctan \left(\frac{\sqrt{1-x^{2}}}{x+1}\right)-c_{1}\right) \\
& y(x) \rightarrow-1 \\
& y(x) \rightarrow 1 \\
& y(x) \rightarrow \text { Interval }[\{-1,1\}]
\end{aligned}
$$

### 1.20 problem 23

$$
\text { 1.20.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . } 234
$$

1.20.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 236
1.20.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 240
1.20.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 244

Internal problem ID [12437]
Internal file name [OUTPUT/11089_Monday_October_16_2023_09_47_27_PM_925777/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 23.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first__order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
3 \mathrm{e}^{x} \tan (y)+\left(1-\mathrm{e}^{x}\right) \sec (y)^{2} y^{\prime}=0
$$

### 1.20.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{3 \mathrm{e}^{x} \sin (2 y)}{2\left(-1+\mathrm{e}^{x}\right)}
\end{aligned}
$$

Where $f(x)=\frac{3 \mathrm{e}^{x}}{-1+\mathrm{e}^{x}}$ and $g(y)=\frac{\sin (2 y)}{2}$. Integrating both sides gives

$$
\begin{gathered}
\frac{1}{\frac{\sin (2 y)}{2}} d y=\frac{3 \mathrm{e}^{x}}{-1+\mathrm{e}^{x}} d x \\
\int \frac{1}{\frac{\sin (2 y)}{2}} d y=\int \frac{3 \mathrm{e}^{x}}{-1+\mathrm{e}^{x}} d x
\end{gathered}
$$

$$
\ln (\csc (2 y)-\cot (2 y))=3 \ln \left(-1+\mathrm{e}^{x}\right)+c_{1}
$$

Raising both side to exponential gives

$$
\csc (2 y)-\cot (2 y)=\mathrm{e}^{3 \ln \left(-1+\mathrm{e}^{x}\right)+c_{1}}
$$

Which simplifies to

$$
\csc (2 y)-\cot (2 y)=c_{2}\left(-1+\mathrm{e}^{x}\right)^{3}
$$

## Summary

The solution(s) found are the following

## $y$

$=\frac{\arctan \left(\frac{2 \mathrm{e}^{c_{1}} c_{2}\left(\mathrm{e}^{3 x}-3 \mathrm{e}^{2 x}+3 \mathrm{e}^{x}-1\right)}{\mathrm{e}^{6 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-20 \mathrm{e}^{3 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{2 x} \mathrm{e}^{2 c_{1} c_{1}^{2}-6 \mathrm{e}^{x}} \mathrm{e}^{2 c_{1} c_{1}^{2}+c_{2}^{2} \mathrm{e}^{2 c_{1}}+1}},-\frac{\mathrm{e}^{6 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c} c_{1} c}{\mathrm{e}^{6 x} \mathrm{e}^{2 c} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c 1} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c_{1} c}}\right.}{2}$


Figure 40: Slope field plot

## Verification of solutions

$y$
$=\frac{\arctan \left(\frac{2 \mathrm{e}^{c_{1}} c_{2}\left(\mathrm{e}^{3 x}-3 \mathrm{e}^{2 x}+3 \mathrm{e}^{x}-1\right)}{\mathrm{e}^{6 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-20 \mathrm{e}^{3 x} \mathrm{e}^{c_{1}} c_{2}^{2}+15 \mathrm{e}^{2 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-6 \mathrm{e}^{x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+c_{2}^{2} \mathrm{e}^{2 c_{1}}+1},-\frac{\mathrm{e}^{6 x} \mathrm{e}^{2 c_{1} c_{2}^{2}-6} \mathrm{e}^{5 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c_{1}} \mathrm{c}}{\mathrm{e}^{6 x} \mathrm{e}^{c_{1}} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c_{1}} \mathrm{c}}\right.}{2}$
Verified OK.

### 1.20.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{3 \mathrm{e}^{x} \tan (y)}{\left(-1+\mathrm{e}^{x}\right) \sec (y)^{2}} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the $\operatorname{PDE}$ (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 44: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=\frac{\mathrm{e}^{-x}\left(-1+\mathrm{e}^{x}\right)}{3} \\
& \eta(x, y)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=y
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d x \\
& =\int \frac{1}{\frac{\mathrm{e}^{-x}\left(-1+\mathrm{e}^{x}\right)}{3}} d x
\end{aligned}
$$

Which results in

$$
S=3 \ln \left(-1+\mathrm{e}^{x}\right)
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{3 \mathrm{e}^{x} \tan (y)}{\left(-1+\mathrm{e}^{x}\right) \sec (y)^{2}}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =0 \\
R_{y} & =1 \\
S_{x} & =\frac{3 \mathrm{e}^{x}}{-1+\mathrm{e}^{x}} \\
S_{y} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\sec (y) \csc (y) \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\sec (R) \csc (R)
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\ln (\tan (R))+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
3 \ln \left(-1+\mathrm{e}^{x}\right)=\ln (\tan (y))+c_{1}
$$

Which simplifies to

$$
3 \ln \left(-1+\mathrm{e}^{x}\right)=\ln (\tan (y))+c_{1}
$$

Which gives

$$
y=\arctan \left(\mathrm{e}^{-c_{1}}\left(-1+\mathrm{e}^{x}\right)^{3}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{3 \mathrm{e}^{x} \tan (y)}{\left(-1+\mathrm{e}^{x}\right) \sec (y)^{2}}$ |  | $\frac{d S}{d R}=\sec (R) \csc (R)$ |
|  |  |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-\infty]{\rightarrow \rightarrow \rightarrow+\infty}$ |  | 星: |
|  |  |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \infty]{ }$ |  |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\infty]{ }$ |  |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\infty]{ }$ | $R=y$ | +iditativatidiapti |
|  | $S=3 \ln \left(-1+\mathrm{e}^{x}\right)$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-\infty]{ }$ |  |  |

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\arctan \left(\mathrm{e}^{-c_{1}}\left(-1+\mathrm{e}^{x}\right)^{3}\right) \tag{1}
\end{equation*}
$$



Figure 41: Slope field plot
Verification of solutions

$$
y=\arctan \left(\mathrm{e}^{-c_{1}}\left(-1+\mathrm{e}^{x}\right)^{3}\right)
$$

Verified OK.

### 1.20.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{\sec (y)^{2}}{3 \tan (y)}\right) \mathrm{d} y & =\left(\frac{\mathrm{e}^{x}}{-1+\mathrm{e}^{x}}\right) \mathrm{d} x \\
\left(-\frac{\mathrm{e}^{x}}{-1+\mathrm{e}^{x}}\right) \mathrm{d} x+\left(\frac{\sec (y)^{2}}{3 \tan (y)}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=-\frac{\mathrm{e}^{x}}{-1+\mathrm{e}^{x}} \\
& N(x, y)=\frac{\sec (y)^{2}}{3 \tan (y)}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{\mathrm{e}^{x}}{-1+\mathrm{e}^{x}}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{\sec (y)^{2}}{3 \tan (y)}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{\mathrm{e}^{x}}{-1+\mathrm{e}^{x}} \mathrm{~d} x \\
\phi & =-\ln \left(-1+\mathrm{e}^{x}\right)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{\sec (y)^{2}}{3 \tan (y)}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{\sec (y)^{2}}{3 \tan (y)}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
\begin{aligned}
f^{\prime}(y) & =\frac{\sec (y)^{2}}{3 \tan (y)} \\
& =\frac{\sec (y) \csc (y)}{3}
\end{aligned}
$$

Integrating the above w.r.t $y$ results in

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{\sec (y) \csc (y)}{3}\right) \mathrm{d} y \\
f(y) & =\frac{\ln (\tan (y))}{3}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\ln \left(-1+\mathrm{e}^{x}\right)+\frac{\ln (\tan (y))}{3}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\ln \left(-1+\mathrm{e}^{x}\right)+\frac{\ln (\tan (y))}{3}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-\ln \left(-1+\mathrm{e}^{x}\right)+\frac{\ln (\tan (y))}{3}=c_{1} \tag{1}
\end{equation*}
$$



Figure 42: Slope field plot

Verification of solutions

$$
-\ln \left(-1+\mathrm{e}^{x}\right)+\frac{\ln (\tan (y))}{3}=c_{1}
$$

Verified OK.

### 1.20.4 Maple step by step solution

Let's solve
$3 \mathrm{e}^{x} \tan (y)+\left(1-\mathrm{e}^{x}\right) \sec (y)^{2} y^{\prime}=0$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Separate variables
$\frac{y^{\prime} \sec (y)^{2}}{\tan (y)}=-\frac{3 \mathrm{e}^{x}}{1-\mathrm{e}^{x}}$
- Integrate both sides with respect to $x$
$\int \frac{y^{\prime} \sec (y)^{2}}{\tan (y)} d x=\int-\frac{3 \mathrm{e}^{x}}{1-\mathrm{e}^{x}} d x+c_{1}$
- Evaluate integral

$$
\ln (\tan (y))=3 \ln \left(1-\mathrm{e}^{x}\right)+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=-\arctan \left(\mathrm{e}^{c_{1}}\left(-1+\mathrm{e}^{x}\right)^{3}\right)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

$\checkmark$ Solution by Maple
Time used: 0.031 (sec). Leaf size: 205

```
dsolve(3*exp(x)*\operatorname{tan}(y(x))+(1-exp(x))*sec(y(x))^2*diff (y(x),x)=0,y(x), singsol=all)
```

$y(x)$
$=\frac{\arctan \left(-\frac{2 c_{1}\left(\mathrm{e}^{3 x}-3 \mathrm{e}^{2 x}+3 \mathrm{e}^{x}-1\right)}{-c_{1}^{2} \mathrm{e}^{6 x}+6 c_{1}^{2} 5^{5 x}-15 c_{1}^{2} \mathrm{e}^{4 x}+20 c_{1}^{2} \mathrm{e}^{3 x}-15 c_{1}^{2} \mathrm{e}^{2 x}+6 c_{1}^{2} \mathrm{e}^{x}-c_{1}^{2}-1}, \frac{c_{1}^{2} e^{6 x}-6 c_{1}^{2} 5^{5 x}+15 c_{1}^{2} \mathrm{e}^{4 x}-20 c_{1}^{2} \mathrm{e}^{3 x}+15 c_{1}^{2} \mathrm{e}^{2 x}-6 c_{1}^{2} \mathrm{e}^{x}+c_{1}^{2}-1}{-c_{1}^{2} \mathrm{e}^{6 x}+6 c_{1}^{2} \mathrm{e}^{5 x}-15 c_{1}^{2} \mathrm{e}^{4 x}+20 c_{1}^{2} \mathrm{e}^{3 x}-15 c_{1}^{2} \mathrm{e}^{2 x}+6 c_{1}^{2} \mathrm{e}^{x}-c_{1}^{2}-1}\right)}{2}$
$\sqrt{ }$ Solution by Mathematica
Time used: 1.829 (sec). Leaf size: 74
DSolve $\left[3 * \operatorname{Exp}[x] * \operatorname{Tan}[y[x]]+(1-\operatorname{Exp}[x]) * \operatorname{Sec}[y[x]]^{\wedge} 2 * y^{\prime}[x]==0, y[x], x\right.$, IncludeSingularSolutions $\rightarrow$

$$
\begin{aligned}
& y(x) \rightarrow-\frac{1}{2} \arccos \left(-\tanh \left(3 \log \left(e^{x}-1\right)+2 c_{1}\right)\right) \\
& y(x) \rightarrow \frac{1}{2} \arccos \left(-\tanh \left(3 \log \left(e^{x}-1\right)+2 c_{1}\right)\right) \\
& y(x) \rightarrow 0 \\
& y(x) \rightarrow-\frac{\pi}{2} \\
& y(x) \rightarrow \frac{\pi}{2}
\end{aligned}
$$

### 1.21 problem 24

$$
\text { 1.21.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . } 246
$$

1.21.2 Solving as first order ode lie symmetry lookup ode ..... 248
1.21.3 Solving as bernoulli ode ..... 252
1.21.4 Solving as exact ode ..... 256
1.21.5 Maple step by step solution ..... 260

Internal problem ID [12438]
Internal file name [OUTPUT/11090_Monday_October_16_2023_09_47_29_PM_63877805/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 24.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "bernoulli", "separable", "first_order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
-y^{2} x+\left(y-x^{2} y\right) y^{\prime}=-x
$$

### 1.21.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =-\frac{x\left(y^{2}-1\right)}{y\left(x^{2}-1\right)}
\end{aligned}
$$

Where $f(x)=-\frac{x}{x^{2}-1}$ and $g(y)=\frac{y^{2}-1}{y}$. Integrating both sides gives

$$
\frac{1}{\frac{y^{2}-1}{y}} d y=-\frac{x}{x^{2}-1} d x
$$

$$
\begin{aligned}
\int \frac{1}{\frac{y^{2}-1}{y}} d y & =\int-\frac{x}{x^{2}-1} d x \\
\frac{\ln (y-1)}{2}+\frac{\ln (y+1)}{2} & =-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}+c_{1}
\end{aligned}
$$

The above can be written as

$$
\begin{aligned}
\left(\frac{1}{2}\right)(\ln (y-1)+\ln (y+1)) & =-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}+2 c_{1} \\
\ln (y-1)+\ln (y+1) & =(2)\left(-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}+2 c_{1}\right) \\
& =-\ln (x-1)-\ln (x+1)+4 c_{1}
\end{aligned}
$$

Raising both side to exponential gives

$$
\mathrm{e}^{\ln (y-1)+\ln (y+1)}=\mathrm{e}^{-\ln (x-1)-\ln (x+1)+2 c_{1}}
$$

Which simplifies to

$$
\begin{aligned}
y^{2}-1 & =2 c_{1} \mathrm{e}^{-\ln (x-1)-\ln (x+1)} \\
& =c_{2} \mathrm{e}^{-\ln (x-1)-\ln (x+1)}
\end{aligned}
$$

Which simplifies to

$$
y^{2}-1=\frac{c_{2}}{(x-1)(x+1)}
$$

The solution is

$$
y^{2}-1=\frac{c_{2}}{(x-1)(x+1)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y^{2}-1=\frac{c_{2}}{(x-1)(x+1)} \tag{1}
\end{equation*}
$$



Figure 43: Slope field plot
Verification of solutions

$$
y^{2}-1=\frac{c_{2}}{(x-1)(x+1)}
$$

Verified OK.

### 1.21.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{x\left(y^{2}-1\right)}{y\left(x^{2}-1\right)} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 47: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=-\frac{x^{2}-1}{x} \\
& \eta(x, y)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=y
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d x \\
& =\int \frac{1}{-\frac{x^{2}-1}{x}} d x
\end{aligned}
$$

Which results in

$$
S=-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{x\left(y^{2}-1\right)}{y\left(x^{2}-1\right)}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =0 \\
R_{y} & =1 \\
S_{x} & =-\frac{x}{x^{2}-1} \\
S_{y} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{y}{y^{2}-1} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{R}{R^{2}-1}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{\ln (R-1)}{2}+\frac{\ln (R+1)}{2}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}=\frac{\ln (y-1)}{2}+\frac{\ln (y+1)}{2}+c_{1}
$$

Which simplifies to

$$
-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}=\frac{\ln (y-1)}{2}+\frac{\ln (y+1)}{2}+c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=-\frac{x\left(y^{2}-1\right)}{y\left(x^{2}-1\right)}$ | $\begin{aligned} & R=y \\ & S=-\frac{\ln (x-1)}{2}- \end{aligned}$ | $\frac{d S}{d R}=\frac{R}{R^{2}-1}$ |
|  |  |  |
|  |  | $\rightarrow \rightarrow+\infty$ |
|  |  | ¢ |
|  |  | (R) $1 \rightarrow+1$ |
|  |  |  |
| - +1 |  | 19, - 介 |
|  |  | 2 $19.0+1$ |
|  |  | , |
|  |  | $1{ }^{1} 1$ |
|  |  | 19 |
|  |  | 19 |
|  |  | 㭗 |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}=\frac{\ln (y-1)}{2}+\frac{\ln (y+1)}{2}+c_{1} \tag{1}
\end{equation*}
$$



Figure 44: Slope field plot

## Verification of solutions

$$
-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}=\frac{\ln (y-1)}{2}+\frac{\ln (y+1)}{2}+c_{1}
$$

Verified OK.

### 1.21.3 Solving as bernoulli ode

In canonical form, the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =-\frac{x\left(y^{2}-1\right)}{y\left(x^{2}-1\right)}
\end{aligned}
$$

This is a Bernoulli ODE.

$$
\begin{equation*}
y^{\prime}=-\frac{x}{x^{2}-1} y+\frac{x}{x^{2}-1} \frac{1}{y} \tag{1}
\end{equation*}
$$

The standard Bernoulli ODE has the form

$$
\begin{equation*}
y^{\prime}=f_{0}(x) y+f_{1}(x) y^{n} \tag{2}
\end{equation*}
$$

The first step is to divide the above equation by $y^{n}$ which gives

$$
\begin{equation*}
\frac{y^{\prime}}{y^{n}}=f_{0}(x) y^{1-n}+f_{1}(x) \tag{3}
\end{equation*}
$$

The next step is use the substitution $w=y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$
\begin{aligned}
f_{0}(x) & =-\frac{x}{x^{2}-1} \\
f_{1}(x) & =\frac{x}{x^{2}-1} \\
n & =-1
\end{aligned}
$$

Dividing both sides of ODE (1) by $y^{n}=\frac{1}{y}$ gives

$$
\begin{equation*}
y^{\prime} y=-\frac{x y^{2}}{x^{2}-1}+\frac{x}{x^{2}-1} \tag{4}
\end{equation*}
$$

Let

$$
\begin{align*}
w & =y^{1-n} \\
& =y^{2} \tag{5}
\end{align*}
$$

Taking derivative of equation (5) w.r.t $x$ gives

$$
\begin{equation*}
w^{\prime}=2 y y^{\prime} \tag{6}
\end{equation*}
$$

Substituting equations (5) and (6) into equation (4) gives

$$
\begin{align*}
\frac{w^{\prime}(x)}{2} & =-\frac{x w(x)}{x^{2}-1}+\frac{x}{x^{2}-1} \\
w^{\prime} & =-\frac{2 x w}{x^{2}-1}+\frac{2 x}{x^{2}-1} \tag{7}
\end{align*}
$$

The above now is a linear ODE in $w(x)$ which is now solved.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
w^{\prime}(x)+p(x) w(x)=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=\frac{2 x}{x^{2}-1} \\
& q(x)=\frac{2 x}{x^{2}-1}
\end{aligned}
$$

Hence the ode is

$$
w^{\prime}(x)+\frac{2 x w(x)}{x^{2}-1}=\frac{2 x}{x^{2}-1}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \frac{2 x}{x^{2}-1} d x} \\
& =\mathrm{e}^{\ln (x-1)+\ln (x+1)}
\end{aligned}
$$

Which simplifies to

$$
\mu=x^{2}-1
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu w) & =(\mu)\left(\frac{2 x}{x^{2}-1}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\left(x^{2}-1\right) w\right) & =\left(x^{2}-1\right)\left(\frac{2 x}{x^{2}-1}\right) \\
\mathrm{d}\left(\left(x^{2}-1\right) w\right) & =(2 x) \mathrm{d} x
\end{aligned}
$$

## Integrating gives

$$
\begin{aligned}
& \left(x^{2}-1\right) w=\int 2 x \mathrm{~d} x \\
& \left(x^{2}-1\right) w=x^{2}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=x^{2}-1$ results in

$$
w(x)=\frac{x^{2}}{x^{2}-1}+\frac{c_{1}}{x^{2}-1}
$$

which simplifies to

$$
w(x)=\frac{x^{2}+c_{1}}{x^{2}-1}
$$

Replacing $w$ in the above by $y^{2}$ using equation (5) gives the final solution.

$$
y^{2}=\frac{x^{2}+c_{1}}{x^{2}-1}
$$

Solving for $y$ gives

$$
\begin{aligned}
& y(x)=\frac{\sqrt{\left(x^{2}-1\right)\left(x^{2}+c_{1}\right)}}{x^{2}-1} \\
& y(x)=-\frac{\sqrt{\left(x^{2}-1\right)\left(x^{2}+c_{1}\right)}}{x^{2}-1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\frac{\sqrt{\left(x^{2}-1\right)\left(x^{2}+c_{1}\right)}}{x^{2}-1}  \tag{1}\\
& y=-\frac{\sqrt{\left(x^{2}-1\right)\left(x^{2}+c_{1}\right)}}{x^{2}-1} \tag{2}
\end{align*}
$$



Figure 45: Slope field plot

## Verification of solutions

$$
y=\frac{\sqrt{\left(x^{2}-1\right)\left(x^{2}+c_{1}\right)}}{x^{2}-1}
$$

Verified OK.

$$
y=-\frac{\sqrt{\left(x^{2}-1\right)\left(x^{2}+c_{1}\right)}}{x^{2}-1}
$$

Verified OK.

### 1.21.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might
or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(-\frac{y}{y^{2}-1}\right) \mathrm{d} y & =\left(\frac{x}{x^{2}-1}\right) \mathrm{d} x \\
\left(-\frac{x}{x^{2}-1}\right) \mathrm{d} x+\left(-\frac{y}{y^{2}-1}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =-\frac{x}{x^{2}-1} \\
N(x, y) & =-\frac{y}{y^{2}-1}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{x}{x^{2}-1}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(-\frac{y}{y^{2}-1}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{x}{x^{2}-1} \mathrm{~d} x \\
\phi & =-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=-\frac{y}{y^{2}-1}$. Therefore equation (4) becomes

$$
\begin{equation*}
-\frac{y}{y^{2}-1}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=-\frac{y}{y^{2}-1}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(-\frac{y}{y^{2}-1}\right) \mathrm{d} y \\
f(y) & =-\frac{\ln (y-1)}{2}-\frac{\ln (y+1)}{2}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}-\frac{\ln (y-1)}{2}-\frac{\ln (y+1)}{2}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}-\frac{\ln (y-1)}{2}-\frac{\ln (y+1)}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}-\frac{\ln (y-1)}{2}-\frac{\ln (y+1)}{2}=c_{1} \tag{1}
\end{equation*}
$$



Figure 46: Slope field plot

Verification of solutions

$$
-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}-\frac{\ln (y-1)}{2}-\frac{\ln (y+1)}{2}=c_{1}
$$

Verified OK.

### 1.21.5 Maple step by step solution

Let's solve

$$
-y^{2} x+\left(y-x^{2} y\right) y^{\prime}=-x
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Integrate both sides with respect to $x$

$$
\int\left(-y^{2} x+\left(y-x^{2} y\right) y^{\prime}\right) d x=\int-x d x+c_{1}
$$

- Evaluate integral

$$
-\frac{y^{2}(x-1)(x+1)}{2}=-\frac{x^{2}}{2}+c_{1}
$$

- $\quad$ Solve for $y$

$$
\left\{y=\frac{\sqrt{-\left(x^{2}-1\right)\left(-x^{2}+2 c_{1}\right)}}{x^{2}-1}, y=-\frac{\sqrt{-\left(x^{2}-1\right)\left(-x^{2}+2 c_{1}\right)}}{x^{2}-1}\right\}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 50

```
dsolve((x-y(x)^2*x)+(y(x)-x^2*y(x))*diff(y(x),x)=0,y(x), singsol=all)
```

$$
\begin{aligned}
& y(x)=\frac{\sqrt{\left(x^{2}-1\right)\left(x^{2}+c_{1}\right)}}{x^{2}-1} \\
& y(x)=-\frac{\sqrt{\left(x^{2}-1\right)\left(x^{2}+c_{1}\right)}}{x^{2}-1}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.652 (sec). Leaf size: 74
DSolve[( $x-y[x] \sim 2 * x)+\left(y[x]-x^{\wedge} 2 * y[x]\right) * y{ }^{\prime}[x]==0, y[x], x$, IncludeSingularSolutions $\rightarrow$ True $]$

$$
\begin{aligned}
& y(x) \rightarrow-\frac{\sqrt{x^{2}-1-e^{2 c_{1}}}}{\sqrt{x^{2}-1}} \\
& y(x) \rightarrow \frac{\sqrt{x^{2}-1-e^{2 c_{1}}}}{\sqrt{x^{2}-1}} \\
& y(x) \rightarrow-1 \\
& y(x) \rightarrow 1
\end{aligned}
$$

### 1.22 problem 39

1.22.1 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 262
1.22.2 Solving as differentialType ode . . . . . . . . . . . . . . . . . . 264
1.22.3 Solving as first order ode lie symmetry calculated ode . . . . . . 266
1.22.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 271
1.22.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 275

Internal problem ID [12439]
Internal file name [OUTPUT/11091_Monday_October_16_2023_09_47_31_PM_43528772/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 39.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "differentialType", "homogeneousTypeD2", "first__order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type
[[_homogeneous, `class A`], _exact, _rational, [_Abel, `2nd type`, `class A`]

$$
y+(y+x) y^{\prime}=x
$$

### 1.22.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
u(x) x+(u(x) x+x)\left(u^{\prime}(x) x+u(x)\right)=x
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =-\frac{u^{2}+2 u-1}{x(u+1)}
\end{aligned}
$$

Where $f(x)=-\frac{1}{x}$ and $g(u)=\frac{u^{2}+2 u-1}{u+1}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{u^{2}+2 u-1}{u+1}} d u & =-\frac{1}{x} d x \\
\int \frac{1}{\frac{u^{2}+2 u-1}{u+1}} d u & =\int-\frac{1}{x} d x \\
\frac{\ln \left(u^{2}+2 u-1\right)}{2} & =-\ln (x)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\sqrt{u^{2}+2 u-1}=\mathrm{e}^{-\ln (x)+c_{2}}
$$

Which simplifies to

$$
\sqrt{u^{2}+2 u-1}=\frac{c_{3}}{x}
$$

Which simplifies to

$$
\sqrt{u(x)^{2}+2 u(x)-1}=\frac{c_{3} \mathrm{e}^{c_{2}}}{x}
$$

The solution is

$$
\sqrt{u(x)^{2}+2 u(x)-1}=\frac{c_{3} \mathrm{e}^{c_{2}}}{x}
$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for $y$ in implicit form

$$
\begin{aligned}
\sqrt{\frac{y^{2}}{x^{2}}+\frac{2 y}{x}-1} & =\frac{c_{3} \mathrm{e}^{c_{2}}}{x} \\
\sqrt{\frac{y^{2}+2 y x-x^{2}}{x^{2}}} & =\frac{c_{3} \mathrm{e}^{c_{2}}}{x}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\sqrt{\frac{y^{2}+2 y x-x^{2}}{x^{2}}}=\frac{c_{3} \mathrm{e}^{c_{2}}}{x} \tag{1}
\end{equation*}
$$



Figure 47: Slope field plot

Verification of solutions

$$
\sqrt{\frac{y^{2}+2 y x-x^{2}}{x^{2}}}=\frac{c_{3} \mathrm{e}^{c_{2}}}{x}
$$

Verified OK.

### 1.22.2 Solving as differentialType ode

Writing the ode as

$$
\begin{equation*}
y^{\prime}=\frac{-y+x}{y+x} \tag{1}
\end{equation*}
$$

Which becomes

$$
\begin{equation*}
(y) d y=(-x) d y+(-y+x) d x \tag{2}
\end{equation*}
$$

But the RHS is complete differential because

$$
(-x) d y+(-y+x) d x=d\left(\frac{1}{2} x^{2}-x y\right)
$$

Hence (2) becomes

$$
\text { (y) } d y=d\left(\frac{1}{2} x^{2}-x y\right)
$$

Integrating both sides gives gives these solutions

$$
\begin{aligned}
& y=-x+\sqrt{2 x^{2}+2 c_{1}}+c_{1} \\
& y=-x-\sqrt{2 x^{2}+2 c_{1}}+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=-x+\sqrt{2 x^{2}+2 c_{1}}+c_{1}  \tag{1}\\
& y=-x-\sqrt{2 x^{2}+2 c_{1}}+c_{1} \tag{2}
\end{align*}
$$



Figure 48: Slope field plot
Verification of solutions

$$
y=-x+\sqrt{2 x^{2}+2 c_{1}}+c_{1}
$$

Verified OK.

$$
y=-x-\sqrt{2 x^{2}+2 c_{1}}+c_{1}
$$

Verified OK.

### 1.22.3 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{y-x}{y+x} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & -\frac{(y-x)\left(b_{3}-a_{2}\right)}{y+x}-\frac{(y-x)^{2} a_{3}}{(y+x)^{2}}-\left(\frac{1}{y+x}+\frac{y-x}{(y+x)^{2}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\left(-\frac{1}{y+x}+\frac{y-x}{(y+x)^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
$$

Putting the above in normal form gives

$$
\begin{aligned}
& -\frac{x^{2} a_{2}+x^{2} a_{3}-3 x^{2} b_{2}-x^{2} b_{3}+2 x y a_{2}-2 x y a_{3}-2 x y b_{2}-2 x y b_{3}-y^{2} a_{2}+3 y^{2} a_{3}-y^{2} b_{2}+y^{2} b_{3}-2 x b_{1}+2}{(y+x)^{2}} \\
& =0
\end{aligned}
$$

Setting the numerator to zero gives

$$
\begin{align*}
& -x^{2} a_{2}-x^{2} a_{3}+3 x^{2} b_{2}+x^{2} b_{3}-2 x y a_{2}+2 x y a_{3}+2 x y b_{2}  \tag{6E}\\
& \quad+2 x y b_{3}+y^{2} a_{2}-3 y^{2} a_{3}+y^{2} b_{2}-y^{2} b_{3}+2 x b_{1}-2 y a_{1}=0
\end{align*}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\{x, y\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& -a_{2} v_{1}^{2}-2 a_{2} v_{1} v_{2}+a_{2} v_{2}^{2}-a_{3} v_{1}^{2}+2 a_{3} v_{1} v_{2}-3 a_{3} v_{2}^{2}+3 b_{2} v_{1}^{2}  \tag{7E}\\
& +2 b_{2} v_{1} v_{2}+b_{2} v_{2}^{2}+b_{3} v_{1}^{2}+2 b_{3} v_{1} v_{2}-b_{3} v_{2}^{2}-2 a_{1} v_{2}+2 b_{1} v_{1}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(-a_{2}-a_{3}+3 b_{2}+b_{3}\right) v_{1}^{2}+\left(-2 a_{2}+2 a_{3}+2 b_{2}+2 b_{3}\right) v_{1} v_{2}  \tag{8E}\\
& \quad+2 b_{1} v_{1}+\left(a_{2}-3 a_{3}+b_{2}-b_{3}\right) v_{2}^{2}-2 a_{1} v_{2}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
-2 a_{1} & =0 \\
2 b_{1} & =0 \\
-2 a_{2}+2 a_{3}+2 b_{2}+2 b_{3} & =0 \\
-a_{2}-a_{3}+3 b_{2}+b_{3} & =0 \\
a_{2}-3 a_{3}+b_{2}-b_{3} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =2 b_{2}+b_{3} \\
a_{3} & =b_{2} \\
b_{1} & =0 \\
b_{2} & =b_{2} \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E, 2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
\xi & =x \\
\eta & =y
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =y-\left(-\frac{y-x}{y+x}\right)(x) \\
& =\frac{-x^{2}+2 x y+y^{2}}{y+x} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{-x^{2}+2 x y+y^{2}}{y+x}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{\ln \left(-x^{2}+2 x y+y^{2}\right)}{2}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{y-x}{y+x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{-y+x}{x^{2}-2 x y-y^{2}} \\
S_{y} & =\frac{-y-x}{x^{2}-2 x y-y^{2}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{\ln \left(y^{2}+2 y x-x^{2}\right)}{2}=c_{1}
$$

Which simplifies to

$$
\frac{\ln \left(y^{2}+2 y x-x^{2}\right)}{2}=c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=-\frac{y-x}{y+x}$ |  | $\frac{d S}{d R}=0$ |
|  |  |  |
| 促 |  |  |
|  |  |  |
|  |  |  |
|  | $R=x$ | $\rightarrow$ |
|  | $S=\underline{\ln \left(-x^{2}+2 x y+y^{2}\right)}$ |  |
|  | $S=\frac{2}{2}$ |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \infty]{ }$ |  |  |
|  |  |  |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \longrightarrow]{ }$ |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
\frac{\ln \left(y^{2}+2 y x-x^{2}\right)}{2}=c_{1} \tag{1}
\end{equation*}
$$



Figure 49: Slope field plot
Verification of solutions

$$
\frac{\ln \left(y^{2}+2 y x-x^{2}\right)}{2}=c_{1}
$$

Verified OK.

### 1.22.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(y+x) \mathrm{d} y & =(-y+x) \mathrm{d} x \\
(y-x) \mathrm{d} x+(y+x) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =y-x \\
N(x, y) & =y+x
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(y-x) \\
& =1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(y+x) \\
& =1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int y-x \mathrm{~d} x \\
\phi & =-\frac{x(x-2 y)}{2}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=x+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=y+x$. Therefore equation (4) becomes

$$
\begin{equation*}
y+x=x+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=y
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int(y) \mathrm{d} y \\
f(y) & =\frac{y^{2}}{2}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\frac{x(x-2 y)}{2}+\frac{y^{2}}{2}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\frac{x(x-2 y)}{2}+\frac{y^{2}}{2}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
-\frac{x(x-2 y)}{2}+\frac{y^{2}}{2}=c_{1} \tag{1}
\end{equation*}
$$



Figure 50: Slope field plot

Verification of solutions

$$
-\frac{x(x-2 y)}{2}+\frac{y^{2}}{2}=c_{1}
$$

Verified OK.

### 1.22.5 Maple step by step solution

Let's solve
$y+(y+x) y^{\prime}=x$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
Check if ODE is exact
- ODE is exact if the lhs is the total derivative of a $C^{2}$ function
$F^{\prime}(x, y)=0$
- Compute derivative of lhs
$F^{\prime}(x, y)+\left(\frac{\partial}{\partial y} F(x, y)\right) y^{\prime}=0$
- Evaluate derivatives
$1=1$
- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$
\left[F(x, y)=c_{1}, M(x, y)=F^{\prime}(x, y), N(x, y)=\frac{\partial}{\partial y} F(x, y)\right]
$$

- $\quad$ Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to $x$
$F(x, y)=\int(y-x) d x+f_{1}(y)$
- Evaluate integral
$F(x, y)=x y-\frac{x^{2}}{2}+f_{1}(y)$
- $\quad$ Take derivative of $F(x, y)$ with respect to $y$
$N(x, y)=\frac{\partial}{\partial y} F(x, y)$
- Compute derivative
$y+x=x+\frac{d}{d y} f_{1}(y)$
- $\quad$ Isolate for $\frac{d}{d y} f_{1}(y)$
$\frac{d}{d y} f_{1}(y)=y$
- $\quad$ Solve for $f_{1}(y)$
$f_{1}(y)=\frac{y^{2}}{2}$
- $\quad$ Substitute $f_{1}(y)$ into equation for $F(x, y)$

$$
F(x, y)=x y-\frac{1}{2} x^{2}+\frac{1}{2} y^{2}
$$

- $\quad$ Substitute $F(x, y)$ into the solution of the ODE

$$
x y-\frac{1}{2} x^{2}+\frac{1}{2} y^{2}=c_{1}
$$

- $\quad$ Solve for $y$

$$
\left\{y=-x-\sqrt{2 x^{2}+2 c_{1}}, y=-x+\sqrt{2 x^{2}+2 c_{1}}\right\}
$$

Maple trace

```
Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.063 (sec). Leaf size: 51

```
dsolve((y(x)-x)+(y(x)+x)*diff (y(x),x)=0,y(x), singsol=all)
```

$$
\begin{aligned}
& y(x)=\frac{-c_{1} x-\sqrt{2 c_{1}^{2} x^{2}+1}}{c_{1}} \\
& y(x)=\frac{-c_{1} x+\sqrt{2 c_{1}^{2} x^{2}+1}}{c_{1}}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.763 (sec). Leaf size: 94
DSolve $[(y[x]-x)+(y[x]+x) * y$ ' $[x]==0, y[x], x$, IncludeSingularSolutions $->$ True $]$

$$
\begin{aligned}
& y(x) \rightarrow-x-\sqrt{2 x^{2}+e^{2 c_{1}}} \\
& y(x) \rightarrow-x+\sqrt{2 x^{2}+e^{2 c_{1}}} \\
& y(x) \rightarrow-\sqrt{2} \sqrt{x^{2}}-x \\
& y(x) \rightarrow \sqrt{2} \sqrt{x^{2}}-x
\end{aligned}
$$

### 1.23 problem 40

$$
\text { 1.23.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . } 278
$$

1.23.2 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 280
1.23.3 Solving as differentialType ode . . . . . . . . . . . . . . . . . . 282
1.23.4 Solving as first order ode lie symmetry lookup ode . . . . . . . 283
1.23.5 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 287
1.23.6 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 291

Internal problem ID [12440]
Internal file name [OUTPUT/11092_Monday_October_16_2023_09_47_34_PM_96130630/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 40.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "linear", "differentialType", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"
Maple gives the following as the ode type
[_linear]

$$
y^{\prime} x+y=-x
$$

### 1.23.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=\frac{1}{x} \\
& q(x)=-1
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+\frac{y}{x}=-1
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
& \mu=\mathrm{e}^{\int \frac{1}{x} d x} \\
& =x
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)(-1) \\
\frac{\mathrm{d}}{\mathrm{~d} x}(x y) & =(x)(-1) \\
\mathrm{d}(x y) & =(-x) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& x y=\int-x \mathrm{~d} x \\
& x y=-\frac{x^{2}}{2}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=x$ results in

$$
y=-\frac{x}{2}+\frac{c_{1}}{x}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{x}{2}+\frac{c_{1}}{x} \tag{1}
\end{equation*}
$$



Figure 51: Slope field plot
Verification of solutions

$$
y=-\frac{x}{2}+\frac{c_{1}}{x}
$$

Verified OK.

### 1.23.2 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
\left(u^{\prime}(x) x+u(x)\right) x+u(x) x=-x
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =\frac{-2 u-1}{x}
\end{aligned}
$$

Where $f(x)=\frac{1}{x}$ and $g(u)=-2 u-1$. Integrating both sides gives

$$
\frac{1}{-2 u-1} d u=\frac{1}{x} d x
$$

$$
\begin{aligned}
\int \frac{1}{-2 u-1} d u & =\int \frac{1}{x} d x \\
-\frac{\ln (-2 u-1)}{2} & =\ln (x)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\frac{1}{\sqrt{-2 u-1}}=\mathrm{e}^{\ln (x)+c_{2}}
$$

Which simplifies to

$$
\frac{1}{\sqrt{-2 u-1}}=c_{3} x
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =u x \\
& =-\frac{\left(c_{3}^{2} \mathrm{e}^{2 c_{2}} x^{2}+1\right) \mathrm{e}^{-2 c_{2}}}{2 x c_{3}^{2}}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\left(c_{3}^{2} \mathrm{e}^{2 c_{2}} x^{2}+1\right) \mathrm{e}^{-2 c_{2}}}{2 x c_{3}^{2}} \tag{1}
\end{equation*}
$$



Figure 52: Slope field plot

## Verification of solutions

$$
y=-\frac{\left(c_{3}^{2} \mathrm{e}^{2 c_{2}} x^{2}+1\right) \mathrm{e}^{-2 c_{2}}}{2 x c_{3}^{2}}
$$

Verified OK.

### 1.23.3 Solving as differentialType ode

Writing the ode as

$$
\begin{equation*}
y^{\prime}=\frac{-y-x}{x} \tag{1}
\end{equation*}
$$

Which becomes

$$
\begin{equation*}
0=(-x) d y+(-y-x) d x \tag{2}
\end{equation*}
$$

But the RHS is complete differential because

$$
(-x) d y+(-y-x) d x=d\left(-\frac{1}{2} x^{2}-x y\right)
$$

Hence (2) becomes

$$
0=d\left(-\frac{1}{2} x^{2}-x y\right)
$$

Integrating both sides gives gives these solutions

$$
y=\frac{-x^{2}+2 c_{1}}{2 x}+c_{1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{-x^{2}+2 c_{1}}{2 x}+c_{1} \tag{1}
\end{equation*}
$$



Figure 53: Slope field plot
Verification of solutions

$$
y=\frac{-x^{2}+2 c_{1}}{2 x}+c_{1}
$$

Verified OK.

### 1.23.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{y+x}{x} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 51: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\frac{1}{x} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{1}{x}} d y
\end{aligned}
$$

Which results in

$$
S=x y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{y+x}{x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =y \\
S_{y} & =x
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-x \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-R
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\frac{R^{2}}{2}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
y x=-\frac{x^{2}}{2}+c_{1}
$$

Which simplifies to

$$
y x=-\frac{x^{2}}{2}+c_{1}
$$

Which gives

$$
y=\frac{-x^{2}+2 c_{1}}{2 x}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=-\frac{y+x}{x}$ |  | $\frac{d S}{d R}=-R$ |
|  |  |  |
|  |  |  |
| $\xrightarrow{\rightarrow \rightarrow \rightarrow-)^{1}}$ |  |  |
|  |  |  |
| $x_{0}$ |  |  |
|  | $R=x$ |  |
| $x^{4}+2 x^{2}$ |  |  |
| $x^{2}$ |  |  |
|  |  |  |
| $L_{1}$ |  |  |
|  |  |  |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{-x^{2}+2 c_{1}}{2 x} \tag{1}
\end{equation*}
$$



Figure 54: Slope field plot
Verification of solutions

$$
y=\frac{-x^{2}+2 c_{1}}{2 x}
$$

Verified OK.

### 1.23.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(x) \mathrm{d} y & =(-y-x) \mathrm{d} x \\
(y+x) \mathrm{d} x+(x) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =y+x \\
N(x, y) & =x
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(y+x) \\
& =1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(x) \\
& =1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int y+x \mathrm{~d} x \\
\phi & =\frac{x(x+2 y)}{2}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=x+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=x$. Therefore equation (4) becomes

$$
\begin{equation*}
x=x+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{x(x+2 y)}{2}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{x(x+2 y)}{2}
$$

The solution becomes

$$
y=\frac{-x^{2}+2 c_{1}}{2 x}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{-x^{2}+2 c_{1}}{2 x} \tag{1}
\end{equation*}
$$



Figure 55: Slope field plot

Verification of solutions

$$
y=\frac{-x^{2}+2 c_{1}}{2 x}
$$

Verified OK.

### 1.23.6 Maple step by step solution

Let's solve
$y^{\prime} x+y=-x$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=-1-\frac{y}{x}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}+\frac{y}{x}=-1$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}+\frac{y}{x}\right)=-\mu(x)$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$
$\mu(x)\left(y^{\prime}+\frac{y}{x}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=\frac{\mu(x)}{x}$
- Solve to find the integrating factor
$\mu(x)=x$
- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int-\mu(x) d x+c_{1}$
- Evaluate the integral on the lhs
$\mu(x) y=\int-\mu(x) d x+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int-\mu(x) d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=x$
$y=\frac{\int-x d x+c_{1}}{x}$
- Evaluate the integrals on the rhs
$y=\frac{-\frac{x^{2}}{2}+c_{1}}{x}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 13

```
dsolve((x+y(x))+x*diff(y(x),x)=0,y(x), singsol=all)
```

$$
y(x)=-\frac{x}{2}+\frac{c_{1}}{x}
$$

$\sqrt{ }$ Solution by Mathematica
Time used: 0.038 (sec). Leaf size: 17
DSolve $[(x+y[x])+x * y$ ' $[x]==0, y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow-\frac{x}{2}+\frac{c_{1}}{x}
$$

### 1.24 problem 41

1.24.1 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 293
1.24.2 Solving as first order ode lie symmetry calculated ode . . . . . . 295
1.24.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 300

Internal problem ID [12441]
Internal file name [OUTPUT/11093_Monday_October_16_2023_09_47_35_PM_17453700/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 41.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exactByInspection", "homogeneousTypeD2", "first_order_ode__lie_symmetry__calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`,
    class A`]]
```

$$
y+(y-x) y^{\prime}=-x
$$

### 1.24.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
u(x) x+(u(x) x-x)\left(u^{\prime}(x) x+u(x)\right)=-x
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =-\frac{u^{2}+1}{x(u-1)}
\end{aligned}
$$

Where $f(x)=-\frac{1}{x}$ and $g(u)=\frac{u^{2}+1}{u-1}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{u^{2}+1}{u-1}} d u & =-\frac{1}{x} d x \\
\int \frac{1}{\frac{u^{2}+1}{u-1}} d u & =\int-\frac{1}{x} d x \\
\frac{\ln \left(u^{2}+1\right)}{2}-\arctan (u) & =-\ln (x)+c_{2}
\end{aligned}
$$

The solution is

$$
\frac{\ln \left(u(x)^{2}+1\right)}{2}-\arctan (u(x))+\ln (x)-c_{2}=0
$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for $y$ in implicit form

$$
\begin{aligned}
& \frac{\ln \left(\frac{y^{2}}{x^{2}}+1\right)}{2}-\arctan \left(\frac{y}{x}\right)+\ln (x)-c_{2}=0 \\
& \frac{\ln \left(\frac{y^{2}}{x^{2}}+1\right)}{2}-\arctan \left(\frac{y}{x}\right)+\ln (x)-c_{2}=0
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{\ln \left(\frac{y^{2}}{x^{2}}+1\right)}{2}-\arctan \left(\frac{y}{x}\right)+\ln (x)-c_{2}=0 \tag{1}
\end{equation*}
$$



Figure 56: Slope field plot

## Verification of solutions

$$
\frac{\ln \left(\frac{y^{2}}{x^{2}}+1\right)}{2}-\arctan \left(\frac{y}{x}\right)+\ln (x)-c_{2}=0
$$

Verified OK.

### 1.24.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{y+x}{y-x} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E, 2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & -\frac{(y+x)\left(b_{3}-a_{2}\right)}{y-x}-\frac{(y+x)^{2} a_{3}}{(y-x)^{2}}-\left(-\frac{1}{y-x}-\frac{y+x}{(y-x)^{2}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\left(-\frac{1}{y-x}+\frac{y+x}{(y-x)^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
$$

Putting the above in normal form gives

$$
\begin{aligned}
& -\frac{x^{2} a_{2}+x^{2} a_{3}+x^{2} b_{2}-x^{2} b_{3}-2 x y a_{2}+2 x y a_{3}+2 x y b_{2}+2 x y b_{3}-y^{2} a_{2}-y^{2} a_{3}-y^{2} b_{2}+y^{2} b_{3}+2 x b_{1}-2 y a}{(-y+x)^{2}} \\
& =0
\end{aligned}
$$

Setting the numerator to zero gives

$$
\begin{align*}
& -x^{2} a_{2}-x^{2} a_{3}-x^{2} b_{2}+x^{2} b_{3}+2 x y a_{2}-2 x y a_{3}-2 x y b_{2}  \tag{6E}\\
& \quad-2 x y b_{3}+y^{2} a_{2}+y^{2} a_{3}+y^{2} b_{2}-y^{2} b_{3}-2 x b_{1}+2 y a_{1}=0
\end{align*}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\{x, y\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& -a_{2} v_{1}^{2}+2 a_{2} v_{1} v_{2}+a_{2} v_{2}^{2}-a_{3} v_{1}^{2}-2 a_{3} v_{1} v_{2}+a_{3} v_{2}^{2}-b_{2} v_{1}^{2}  \tag{7E}\\
& \quad-2 b_{2} v_{1} v_{2}+b_{2} v_{2}^{2}+b_{3} v_{1}^{2}-2 b_{3} v_{1} v_{2}-b_{3} v_{2}^{2}+2 a_{1} v_{2}-2 b_{1} v_{1}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(-a_{2}-a_{3}-b_{2}+b_{3}\right) v_{1}^{2}+\left(2 a_{2}-2 a_{3}-2 b_{2}-2 b_{3}\right) v_{1} v_{2}  \tag{8E}\\
& \quad-2 b_{1} v_{1}+\left(a_{2}+a_{3}+b_{2}-b_{3}\right) v_{2}^{2}+2 a_{1} v_{2}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
2 a_{1} & =0 \\
-2 b_{1} & =0 \\
-a_{2}-a_{3}-b_{2}+b_{3} & =0 \\
a_{2}+a_{3}+b_{2}-b_{3} & =0 \\
2 a_{2}-2 a_{3}-2 b_{2}-2 b_{3} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =b_{3} \\
a_{3} & =-b_{2} \\
b_{1} & =0 \\
b_{2} & =b_{2} \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=x \\
& \eta=y
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =y-\left(-\frac{y+x}{y-x}\right)(x) \\
& =\frac{-x^{2}-y^{2}}{-y+x} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{-x^{2}-y^{2}}{-y+x}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{\ln \left(x^{2}+y^{2}\right)}{2}-\arctan \left(\frac{y}{x}\right)
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{y+x}{y-x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{y+x}{x^{2}+y^{2}} \\
S_{y} & =\frac{y-x}{x^{2}+y^{2}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{\ln \left(x^{2}+y^{2}\right)}{2}-\arctan \left(\frac{y}{x}\right)=c_{1}
$$

Which simplifies to

$$
\frac{\ln \left(x^{2}+y^{2}\right)}{2}-\arctan \left(\frac{y}{x}\right)=c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.


## Summary

The solution(s) found are the following

$$
\begin{equation*}
\frac{\ln \left(x^{2}+y^{2}\right)}{2}-\arctan \left(\frac{y}{x}\right)=c_{1} \tag{1}
\end{equation*}
$$



Figure 57: Slope field plot

## Verification of solutions

$$
\frac{\ln \left(x^{2}+y^{2}\right)}{2}-\arctan \left(\frac{y}{x}\right)=c_{1}
$$

Verified OK.

### 1.24.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(y-x) \mathrm{d} y & =(-y-x) \mathrm{d} x \\
(y+x) \mathrm{d} x+(y-x) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=y+x \\
& N(x, y)=y-x
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(y+x) \\
& =1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(y-x) \\
& =-1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. By inspection $\frac{1}{x^{2}+y^{2}}$ is an integrating factor. Therefore by multiplying $M=y+x$ and $N=y-x$ by this integrating factor the ode becomes exact. The new $M, N$ are

$$
\begin{aligned}
M & =\frac{y+x}{x^{2}+y^{2}} \\
N & =\frac{y-x}{x^{2}+y^{2}}
\end{aligned}
$$

To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might
or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{y-x}{x^{2}+y^{2}}\right) \mathrm{d} y & =\left(-\frac{y+x}{x^{2}+y^{2}}\right) \mathrm{d} x \\
\left(\frac{y+x}{x^{2}+y^{2}}\right) \mathrm{d} x+\left(\frac{y-x}{x^{2}+y^{2}}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =\frac{y+x}{x^{2}+y^{2}} \\
N(x, y) & =\frac{y-x}{x^{2}+y^{2}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\frac{y+x}{x^{2}+y^{2}}\right) \\
& =\frac{x^{2}-2 x y-y^{2}}{\left(x^{2}+y^{2}\right)^{2}}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{y-x}{x^{2}+y^{2}}\right) \\
& =\frac{x^{2}-2 x y-y^{2}}{\left(x^{2}+y^{2}\right)^{2}}
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{y+x}{x^{2}+y^{2}} \mathrm{~d} x \\
\phi & =\frac{\ln \left(x^{2}+y^{2}\right)}{2}+\arctan \left(\frac{x}{y}\right)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{align*}
\frac{\partial \phi}{\partial y} & =\frac{y}{x^{2}+y^{2}}-\frac{x}{y^{2}\left(\frac{x^{2}}{y^{2}}+1\right)}+f^{\prime}(y)  \tag{4}\\
& =\frac{y-x}{x^{2}+y^{2}}+f^{\prime}(y)
\end{align*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{y-x}{x^{2}+y^{2}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{y-x}{x^{2}+y^{2}}=\frac{y-x}{x^{2}+y^{2}}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{\ln \left(x^{2}+y^{2}\right)}{2}+\arctan \left(\frac{x}{y}\right)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{\ln \left(x^{2}+y^{2}\right)}{2}+\arctan \left(\frac{x}{y}\right)
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
\frac{\ln \left(x^{2}+y^{2}\right)}{2}+\arctan \left(\frac{x}{y}\right)=c_{1} \tag{1}
\end{equation*}
$$



Figure 58: Slope field plot

Verification of solutions

$$
\frac{\ln \left(x^{2}+y^{2}\right)}{2}+\arctan \left(\frac{x}{y}\right)=c_{1}
$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 ( sec ). Leaf size: 24

```
dsolve((x+y(x))+(y(x)-x)*diff(y(x),x)=0,y(x), singsol=all)
```

$$
y(x)=\tan \left(\text { RootOf }\left(-2 \_Z+\ln \left(\sec \left(\_Z\right)^{2}\right)+2 \ln (x)+2 c_{1}\right)\right) x
$$

Solution by Mathematica
Time used: 0.054 (sec). Leaf size: 36
DSolve $[(x+y[x])+(y[x]-x) * y$ ' $[x]==0, y[x], x$, IncludeSingularSolutions $->$ True $]$

$$
\text { Solve }\left[\frac{1}{2} \log \left(\frac{y(x)^{2}}{x^{2}}+1\right)-\arctan \left(\frac{y(x)}{x}\right)=-\log (x)+c_{1}, y(x)\right]
$$

### 1.25 problem 42

1.25.1 Solving as first order ode lie symmetry calculated ode

Internal problem ID [12442]
Internal file name [OUTPUT/11094_Monday_October_16_2023_09_47_38_PM_50688309/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 42.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "first_order_ode__lie__symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _dAlembert]
```

$$
y^{\prime} x-y-\sqrt{x^{2}+y^{2}}=0
$$

### 1.25.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{y+\sqrt{x^{2}+y^{2}}}{x} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & +\frac{\left(y+\sqrt{x^{2}+y^{2}}\right)\left(b_{3}-a_{2}\right)}{x}-\frac{\left(y+\sqrt{x^{2}+y^{2}}\right)^{2} a_{3}}{x^{2}} \\
& -\left(\frac{1}{\sqrt{x^{2}+y^{2}}}-\frac{y+\sqrt{x^{2}+y^{2}}}{x^{2}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\frac{\left(1+\frac{y}{\sqrt{x^{2}+y^{2}}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)}{x}=0
\end{align*}
$$

Putting the above in normal form gives

$$
\begin{aligned}
& -\frac{\left(x^{2}+y^{2}\right)^{\frac{3}{2}} a_{3}+x^{3} a_{2}-x^{3} b_{3}+2 x^{2} y a_{3}+x^{2} y b_{2}+y^{3} a_{3}+\sqrt{x^{2}+y^{2}} x b_{1}-\sqrt{x^{2}+y^{2}} y a_{1}+x y b_{1}-y^{2} a_{1}}{\sqrt{x^{2}+y^{2}} x^{2}} \\
& =0
\end{aligned}
$$

Setting the numerator to zero gives

$$
\begin{gather*}
-\left(x^{2}+y^{2}\right)^{\frac{3}{2}} a_{3}-x^{3} a_{2}+x^{3} b_{3}-2 x^{2} y a_{3}-x^{2} y b_{2}-y^{3} a_{3}  \tag{6E}\\
-\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}-x y b_{1}+y^{2} a_{1}=0
\end{gather*}
$$

Simplifying the above gives

$$
\begin{align*}
& -\left(x^{2}+y^{2}\right)^{\frac{3}{2}} a_{3}+\left(x^{2}+y^{2}\right) x b_{3}-\left(x^{2}+y^{2}\right) y a_{3}-x^{3} a_{2}-x^{2} y a_{3}-x^{2} y b_{2}  \tag{6E}\\
& \quad-x y^{2} b_{3}+\left(x^{2}+y^{2}\right) a_{1}-\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}-x^{2} a_{1}-x y b_{1}=0
\end{align*}
$$

Since the PDE has radicals, simplifying gives

$$
\begin{aligned}
& -x^{3} a_{2}+x^{3} b_{3}-x^{2} \sqrt{x^{2}+y^{2}} a_{3}-2 x^{2} y a_{3}-x^{2} y b_{2}-\sqrt{x^{2}+y^{2}} y^{2} a_{3} \\
& \quad-y^{3} a_{3}-\sqrt{x^{2}+y^{2}} x b_{1}-x y b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}+y^{2} a_{1}=0
\end{aligned}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\left\{x, y, \sqrt{x^{2}+y^{2}}\right\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}, \sqrt{x^{2}+y^{2}}=v_{3}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& -v_{1}^{3} a_{2}-2 v_{1}^{2} v_{2} a_{3}-v_{1}^{2} v_{3} a_{3}-v_{2}^{3} a_{3}-v_{3} v_{2}^{2} a_{3}-v_{1}^{2} v_{2} b_{2}  \tag{7E}\\
& +v_{1}^{3} b_{3}+v_{2}^{2} a_{1}+v_{3} v_{2} a_{1}-v_{1} v_{2} b_{1}-v_{3} v_{1} b_{1}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}, v_{3}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(b_{3}-a_{2}\right) v_{1}^{3}+\left(-2 a_{3}-b_{2}\right) v_{1}^{2} v_{2}-v_{1}^{2} v_{3} a_{3}-v_{1} v_{2} b_{1}  \tag{8E}\\
& \quad-v_{3} v_{1} b_{1}-v_{2}^{3} a_{3}-v_{3} v_{2}^{2} a_{3}+v_{2}^{2} a_{1}+v_{3} v_{2} a_{1}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
a_{1} & =0 \\
-a_{3} & =0 \\
-b_{1} & =0 \\
-2 a_{3}-b_{2} & =0 \\
b_{3}-a_{2} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =b_{3} \\
a_{3} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E, 2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
\xi & =x \\
\eta & =y
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =y-\left(\frac{y+\sqrt{x^{2}+y^{2}}}{x}\right)(x) \\
& =-\sqrt{x^{2}+y^{2}} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{-\sqrt{x^{2}+y^{2}}} d y
\end{aligned}
$$

Which results in

$$
S=-\ln \left(y+\sqrt{x^{2}+y^{2}}\right)
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{y+\sqrt{x^{2}+y^{2}}}{x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
& R_{x}=1 \\
& R_{y}=0 \\
& S_{x}=-\frac{x}{\sqrt{x^{2}+y^{2}}\left(y+\sqrt{x^{2}+y^{2}}\right)} \\
& S_{y}=-\frac{1}{\sqrt{x^{2}+y^{2}}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-\frac{2\left(\sqrt{x^{2}+y^{2}} y+x^{2}+y^{2}\right)}{x \sqrt{x^{2}+y^{2}}\left(y+\sqrt{x^{2}+y^{2}}\right)} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-\frac{2}{R}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-2 \ln (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
-\ln \left(y+\sqrt{x^{2}+y^{2}}\right)=-2 \ln (x)+c_{1}
$$

Which simplifies to

$$
-\ln \left(y+\sqrt{x^{2}+y^{2}}\right)=-2 \ln (x)+c_{1}
$$

Which gives

$$
y=-\frac{\mathrm{e}^{-c_{1}}\left(\mathrm{e}^{2 c_{1}}-x^{2}\right)}{2}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{y+\sqrt{x^{2}+y^{2}}}{x}$ |  | $\frac{d S}{d R}=-\frac{2}{R}$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=x$ |  |
|  |  |  |
|  | $S=-\ln (y+\sqrt{x}$ |  |
| $\xrightarrow{2}$ |  |  |
| - - - - - - |  |  |
| $\rightarrow \rightarrow \rightarrow-\infty$ |  |  |
| 为 |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\mathrm{e}^{-c_{1}}\left(\mathrm{e}^{2 c_{1}}-x^{2}\right)}{2} \tag{1}
\end{equation*}
$$



Figure 59: Slope field plot

Verification of solutions

$$
y=-\frac{\mathrm{e}^{-c_{1}}\left(\mathrm{e}^{2 c_{1}}-x^{2}\right)}{2}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying homogeneous types:
trying homogeneous G
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.047 (sec). Leaf size: 26
dsolve( $x * \operatorname{diff}(y(x), x)-y(x)=\operatorname{sqrt}\left(x^{\wedge} 2+y(x)^{\wedge} 2\right), y(x)$, singsol=all)

$$
\frac{-c_{1} x^{2}+y(x)+\sqrt{y(x)^{2}+x^{2}}}{x^{2}}=0
$$

$\checkmark$ Solution by Mathematica
Time used: 0.575 (sec). Leaf size: 27
DSolve[x*y'[x]-y[x]==Sqrt[x^2+y[x]~2],y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow \frac{1}{2} e^{-c_{1}}\left(-1+e^{2 c_{1}} x^{2}\right)
$$

### 1.26 problem 43

1.26.1 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 315
1.26.2 Solving as first order ode lie symmetry calculated ode . . . . . . 317

Internal problem ID [12443]
Internal file name [OUTPUT/11095_Monday_October_16_2023_09_47_41_PM_89202063/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 43.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "homogeneousTypeD2", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `
    class A`]]
```

$$
8 y+(5 y+7 x) y^{\prime}=-10 x
$$

### 1.26.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
8 u(x) x+(5 u(x) x+7 x)\left(u^{\prime}(x) x+u(x)\right)=-10 x
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =-\frac{5\left(u^{2}+3 u+2\right)}{x(5 u+7)}
\end{aligned}
$$

Where $f(x)=-\frac{5}{x}$ and $g(u)=\frac{u^{2}+3 u+2}{5 u+7}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{u^{2}+3 u+2}{5 u+7}} d u & =-\frac{5}{x} d x \\
\int \frac{1}{\frac{u^{2}+3 u+2}{5 u+7}} d u & =\int-\frac{5}{x} d x \\
2 \ln (u+1)+3 \ln (u+2) & =-5 \ln (x)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\mathrm{e}^{2 \ln (u+1)+3 \ln (u+2)}=\mathrm{e}^{-5 \ln (x)+c_{2}}
$$

Which simplifies to

$$
(u+1)^{2}(u+2)^{3}=\frac{c_{3}}{x^{5}}
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =x u \\
& =\operatorname{RootOf}\left(\_Z^{5}+8 x \_Z^{4}+25 \_Z^{3} x^{2}+38 x^{3} \_Z^{2}+28 \_Z x^{4}+8 x^{5}-c_{3}\right)
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\operatorname{RootOf}\left(\_Z^{5}+8 x \_Z^{4}+25 \_Z^{3} x^{2}+38 x^{3} \_Z^{2}+28 \_Z x^{4}+8 x^{5}-c_{3}\right) \tag{1}
\end{equation*}
$$



Figure 60: Slope field plot
Verification of solutions

$$
y=\operatorname{RootOf}\left(\_Z^{5}+8 x \_Z^{4}+25 \_Z^{3} x^{2}+38 x^{3} \_Z^{2}+28 \_Z x^{4}+8 x^{5}-c_{3}\right)
$$

Verified OK.

### 1.26.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=-\frac{2(5 x+4 y)}{5 y+7 x} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & -\frac{2(5 x+4 y)\left(b_{3}-a_{2}\right)}{5 y+7 x}-\frac{4(5 x+4 y)^{2} a_{3}}{(5 y+7 x)^{2}} \\
& -\left(-\frac{10}{5 y+7 x}+\frac{70 x+56 y}{(5 y+7 x)^{2}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\left(-\frac{8}{5 y+7 x}+\frac{50 x+40 y}{(5 y+7 x)^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
$$

Putting the above in normal form gives

$$
\begin{aligned}
& \frac{70 x^{2} a_{2}-100 x^{2} a_{3}+55 x^{2} b_{2}-70 x^{2} b_{3}+100 x y a_{2}-160 x y a_{3}+70 x y b_{2}-100 x y b_{3}+40 y^{2} a_{2}-70 y^{2} a_{3}+25 y}{(5 y+7 x)^{2}} \\
& =0
\end{aligned}
$$

Setting the numerator to zero gives

$$
\begin{array}{r}
70 x^{2} a_{2}-100 x^{2} a_{3}+55 x^{2} b_{2}-70 x^{2} b_{3}+100 x y a_{2}-160 x y a_{3}+70 x y b_{2}  \tag{6E}\\
-100 x y b_{3}+40 y^{2} a_{2}-70 y^{2} a_{3}+25 y^{2} b_{2}-40 y^{2} b_{3}+6 x b_{1}-6 y a_{1}=0
\end{array}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\{x, y\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& 70 a_{2} v_{1}^{2}+100 a_{2} v_{1} v_{2}+40 a_{2} v_{2}^{2}-100 a_{3} v_{1}^{2}-160 a_{3} v_{1} v_{2}-70 a_{3} v_{2}^{2}+55 b_{2} v_{1}^{2}  \tag{7E}\\
& +70 b_{2} v_{1} v_{2}+25 b_{2} v_{2}^{2}-70 b_{3} v_{1}^{2}-100 b_{3} v_{1} v_{2}-40 b_{3} v_{2}^{2}-6 a_{1} v_{2}+6 b_{1} v_{1}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(70 a_{2}-100 a_{3}+55 b_{2}-70 b_{3}\right) v_{1}^{2}+\left(100 a_{2}-160 a_{3}+70 b_{2}-100 b_{3}\right) v_{1} v_{2}  \tag{8E}\\
& +6 b_{1} v_{1}+\left(40 a_{2}-70 a_{3}+25 b_{2}-40 b_{3}\right) v_{2}^{2}-6 a_{1} v_{2}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
-6 a_{1} & =0 \\
6 b_{1} & =0 \\
40 a_{2}-70 a_{3}+25 b_{2}-40 b_{3} & =0 \\
70 a_{2}-100 a_{3}+55 b_{2}-70 b_{3} & =0 \\
100 a_{2}-160 a_{3}+70 b_{2}-100 b_{3} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
& a_{1}=0 \\
& a_{2}=3 a_{3}+b_{3} \\
& a_{3}=a_{3} \\
& b_{1}=0 \\
& b_{2}=-2 a_{3} \\
& b_{3}=b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E, 2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=x \\
& \eta=y
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =y-\left(-\frac{2(5 x+4 y)}{5 y+7 x}\right)(x) \\
& =\frac{10 x^{2}+15 x y+5 y^{2}}{5 y+7 x} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{10 x^{2}+15 x y+5 y^{2}}{5 y+7 x}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{2 \ln (y+x)}{5}+\frac{3 \ln (2 x+y)}{5}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{2(5 x+4 y)}{5 y+7 x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{2 x+\frac{8 y}{5}}{(y+x)(2 x+y)} \\
S_{y} & =\frac{5 y+7 x}{5(y+x)(2 x+y)}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{2 \ln (y+x)}{5}+\frac{3 \ln (2 x+y)}{5}=c_{1}
$$

Which simplifies to

$$
\frac{2 \ln (y+x)}{5}+\frac{3 \ln (2 x+y)}{5}=c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.


Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{2 \ln (y+x)}{5}+\frac{3 \ln (2 x+y)}{5}=c_{1} \tag{1}
\end{equation*}
$$



Figure 61: Slope field plot

Verification of solutions

$$
\frac{2 \ln (y+x)}{5}+\frac{3 \ln (2 x+y)}{5}=c_{1}
$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.734 (sec). Leaf size: 38

```
dsolve((8*y(x)+10*x)+(5*y(x)+7*x)*diff (y(x),x)=0,y(x), singsol=all)
```

$$
y(x)=x\left(-2+\operatorname{RootOf}\left(\_Z^{25} c_{1} x^{5}-2 \_Z^{20} c_{1} x^{5}+\_Z^{15} c_{1} x^{5}-1\right)^{5}\right)
$$

Solution by Mathematica
Time used: 3.57 (sec). Leaf size: 276
DSolve $\left[(8 * y[x]+10 * x)+(5 * y[x]+7 * x) * y{ }^{\prime}[x]==0, y[x], x\right.$, IncludeSingularSolutions $->$ True]

$$
\begin{aligned}
& y(x) \rightarrow \operatorname{Root}\left[\# 1^{5}+8 \# 1^{4} x+25 \# 1^{3} x^{2}+38 \# 1^{2} x^{3}+28 \# 1 x^{4}+8 x^{5}-e^{c_{1}} \&, 1\right] \\
& y(x) \rightarrow \operatorname{Root}\left[\# 1^{5}+8 \# 1^{4} x+25 \# 1^{3} x^{2}+38 \# 1^{2} x^{3}+28 \# 1 x^{4}+8 x^{5}-e^{c_{1}} \&, 2\right] \\
& y(x) \rightarrow \operatorname{Root}\left[\# 1^{5}+8 \# 1^{4} x+25 \# 1^{3} x^{2}+38 \# 1^{2} x^{3}+28 \# 1 x^{4}+8 x^{5}-e^{c_{1}} \&, 3\right] \\
& y(x) \rightarrow \operatorname{Root}\left[\# 1^{5}+8 \# 1^{4} x+25 \# 1^{3} x^{2}+38 \# 1^{2} x^{3}+28 \# 1 x^{4}+8 x^{5}-e^{c_{1}} \&, 4\right] \\
& y(x) \rightarrow \operatorname{Root}\left[\# 1^{5}+8 \# 1^{4} x+25 \# 1^{3} x^{2}+38 \# 1^{2} x^{3}+28 \# 1 x^{4}+8 x^{5}-e^{c_{1}} \&, 5\right]
\end{aligned}
$$

### 1.27 problem 44

1.27.1 Solving as first order ode lie symmetry calculated ode . . . . . . 324

Internal problem ID [12444]
Internal file name [OUTPUT/11096_Monday_October_16_2023_09_47_47_PM_82011611/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 44.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type
[[_homogeneous, `class A`], _dAlembert]

$$
2 \sqrt{s t}-s+t s^{\prime}=0
$$

### 1.27.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
& s^{\prime}=-\frac{2 \sqrt{s t}-s}{t} \\
& s^{\prime}=\omega(t, s)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{s}-\xi_{t}\right)-\omega^{2} \xi_{s}-\omega_{t} \xi-\omega_{s} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=s a_{3}+t a_{2}+a_{1}  \tag{1E}\\
& \eta=s b_{3}+t b_{2}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & -\frac{(2 \sqrt{s t}-s)\left(b_{3}-a_{2}\right)}{t}-\frac{(2 \sqrt{s t}-s)^{2} a_{3}}{t^{2}}  \tag{5E}\\
& -\left(-\frac{s}{\sqrt{s t} t}+\frac{2 \sqrt{s t}-s}{t^{2}}\right)\left(s a_{3}+t a_{2}+a_{1}\right)+\frac{\left(\frac{t}{\sqrt{s t}}-1\right)\left(s b_{3}+t b_{2}+b_{1}\right)}{t}=0
\end{align*}
$$

Putting the above in normal form gives

$$
-\frac{4(s t)^{\frac{3}{2}} a_{3}-3 s^{2} t a_{3}+s t^{2} b_{3}-s t^{2} a_{2}-t^{3} b_{2}+s t a_{1}-\sqrt{s t} s a_{1}+\sqrt{s t} t b_{1}-t^{2} b_{1}}{\sqrt{s t} t^{2}}=0
$$

Setting the numerator to zero gives

$$
\begin{equation*}
-4(s t)^{\frac{3}{2}} a_{3}+3 s^{2} t a_{3}+s t^{2} a_{2}-s t^{2} b_{3}+t^{3} b_{2}+\sqrt{s t} s a_{1}-\sqrt{s t} t b_{1}-s t a_{1}+t^{2} b_{1}=0 \tag{6E}
\end{equation*}
$$

Since the PDE has radicals, simplifying gives

$$
3 s^{2} t a_{3}+s t^{2} a_{2}-s t^{2} b_{3}-4 s t \sqrt{s t} a_{3}+t^{3} b_{2}-s t a_{1}+\sqrt{s t} s a_{1}+t^{2} b_{1}-\sqrt{s t} t b_{1}=0
$$

Looking at the above PDE shows the following are all the terms with $\{s, t\}$ in them.

$$
\{s, t, \sqrt{s t}\}
$$

The following substitution is now made to be able to collect on all terms with $\{s, t\}$ in them

$$
\left\{s=v_{1}, t=v_{2}, \sqrt{s t}=v_{3}\right\}
$$

The above PDE (6E) now becomes
$v_{1} v_{2}^{2} a_{2}+3 v_{1}^{2} v_{2} a_{3}-4 v_{1} v_{2} v_{3} a_{3}+v_{2}^{3} b_{2}-v_{1} v_{2}^{2} b_{3}-v_{1} v_{2} a_{1}+v_{3} v_{1} a_{1}+v_{2}^{2} b_{1}-v_{3} v_{2} b_{1}=0$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}, v_{3}\right\}
$$

Equation (7E) now becomes

$$
\begin{equation*}
3 v_{1}^{2} v_{2} a_{3}+\left(-b_{3}+a_{2}\right) v_{1} v_{2}^{2}-4 v_{1} v_{2} v_{3} a_{3}-v_{1} v_{2} a_{1}+v_{3} v_{1} a_{1}+v_{2}^{3} b_{2}+v_{2}^{2} b_{1}-v_{3} v_{2} b_{1}=0 \tag{8E}
\end{equation*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{array}{r}
a_{1}=0 \\
b_{1}=0 \\
b_{2}=0 \\
-a_{1}=0 \\
-4 a_{3}=0 \\
3 a_{3}=0 \\
-b_{1}=0 \\
-b_{3}+a_{2}=0
\end{array}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =b_{3} \\
a_{3} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=t \\
& \eta=s
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(t, s) \xi \\
& =s-\left(-\frac{2 \sqrt{s t}-s}{t}\right)(t) \\
& =2 \sqrt{s t} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(t, s) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d s}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial s}\right) S(t, s)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{2 \sqrt{s t}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{s}{\sqrt{s t}}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, s) S_{s}}{R_{t}+\omega(t, s) R_{s}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{s}, S_{t}, S_{s}$ are all partial derivatives and $\omega(t, s)$ is the right hand side of the original ode given by

$$
\omega(t, s)=-\frac{2 \sqrt{s t}-s}{t}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{s} & =0 \\
S_{t} & =-\frac{\sqrt{s}}{2 t^{\frac{3}{2}}} \\
S_{s} & =\frac{1}{2 \sqrt{s} \sqrt{t}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-\frac{\sqrt{s t}}{\sqrt{s} t^{\frac{3}{2}}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, s$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-\frac{1}{R}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\ln (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, s$ coordinates. This results in

$$
\frac{\sqrt{s}}{\sqrt{t}}=-\ln (t)+c_{1}
$$

Which simplifies to

$$
\frac{\sqrt{s}}{\sqrt{t}}=-\ln (t)+c_{1}
$$

Which gives

$$
s=t \ln (t)^{2}-2 t \ln (t) c_{1}+t c_{1}^{2}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.


## Summary

The solution(s) found are the following

$$
\begin{equation*}
s=t \ln (t)^{2}-2 t \ln (t) c_{1}+t c_{1}^{2} \tag{1}
\end{equation*}
$$



Figure 62: Slope field plot

Verification of solutions

$$
s=t \ln (t)^{2}-2 t \ln (t) c_{1}+t c_{1}^{2}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying homogeneous types:
trying homogeneous G
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.015 (sec). Leaf size: 18
dsolve( $(2 * \operatorname{sqrt}(s(t) * t)-s(t))+t * \operatorname{diff}(s(t), t)=0, s(t), \quad$ singsol=all)

$$
\frac{s(t)}{\sqrt{s(t) t}}+\ln (t)-c_{1}=0
$$

$\checkmark$ Solution by Mathematica
Time used: 0.315 (sec). Leaf size: 19
DSolve[(2*Sqrt[s[t]*t]-s[t])+t*s'[t]==0,s[t],t,IncludeSingularSolutions $\rightarrow$ True]

$$
s(t) \rightarrow \frac{1}{4} t\left(-2 \log (t)+c_{1}\right)^{2}
$$

### 1.28 problem 45

1.28.1 Solving as linear ode332
1.28.2 Solving as homogeneousTypeD2 ode ..... 334
1.28.3 Solving as first order ode lie symmetry lookup ode ..... 335
1.28.4 Solving as exact ode ..... 339
1.28.5 Maple step by step solution ..... 344

Internal problem ID [12445]
Internal file name [OUTPUT/11097_Monday_October_16_2023_09_47_52_PM_93722533/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 45.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "homogeneousTypeD2", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry__lookup"

Maple gives the following as the ode type
[_linear]

$$
-s+t s^{\prime}=-t
$$

### 1.28.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
s^{\prime}+p(t) s=q(t)
$$

Where here

$$
\begin{aligned}
& p(t)=-\frac{1}{t} \\
& q(t)=-1
\end{aligned}
$$

Hence the ode is

$$
s^{\prime}-\frac{s}{t}=-1
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{1}{t} d t} \\
& =\frac{1}{t}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu s) & =(\mu)(-1) \\
\frac{\mathrm{d}}{\mathrm{~d} t}\left(\frac{s}{t}\right) & =\left(\frac{1}{t}\right)(-1) \\
\mathrm{d}\left(\frac{s}{t}\right) & =\left(-\frac{1}{t}\right) \mathrm{d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{s}{t}=\int-\frac{1}{t} \mathrm{~d} t \\
& \frac{s}{t}=-\ln (t)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{t}$ results in

$$
s=-t \ln (t)+c_{1} t
$$

which simplifies to

$$
s=t\left(-\ln (t)+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
s=t\left(-\ln (t)+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 63: Slope field plot
Verification of solutions

$$
s=t\left(-\ln (t)+c_{1}\right)
$$

Verified OK.

### 1.28.2 Solving as homogeneousTypeD2 ode

Using the change of variables $s=u(t) t$ on the above ode results in new ode in $u(t)$

$$
-u(t) t+t\left(u^{\prime}(t) t+u(t)\right)=-t
$$

Integrating both sides gives

$$
\begin{aligned}
u(t) & =\int-\frac{1}{t} \mathrm{~d} t \\
& =-\ln (t)+c_{2}
\end{aligned}
$$

Therefore the solution $s$ is

$$
\begin{aligned}
s & =u t \\
& =t\left(-\ln (t)+c_{2}\right)
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
s=t\left(-\ln (t)+c_{2}\right) \tag{1}
\end{equation*}
$$



Figure 64: Slope field plot
Verification of solutions

$$
s=t\left(-\ln (t)+c_{2}\right)
$$

Verified OK.

### 1.28.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& s^{\prime}=\frac{-t+s}{t} \\
& s^{\prime}=\omega(t, s)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{s}-\xi_{t}\right)-\omega^{2} \xi_{s}-\omega_{t} \xi-\omega_{s} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 54: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, s)=0 \\
& \eta(t, s)=t \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, s) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d s}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial s}\right) S(t, s)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{t} d y
\end{aligned}
$$

Which results in

$$
S=\frac{s}{t}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, s) S_{s}}{R_{t}+\omega(t, s) R_{s}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{s}, S_{t}, S_{s}$ are all partial derivatives and $\omega(t, s)$ is the right hand side of the original ode given by

$$
\omega(t, s)=\frac{-t+s}{t}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{s} & =0 \\
S_{t} & =-\frac{s}{t^{2}} \\
S_{s} & =\frac{1}{t}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-\frac{1}{t} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, s$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-\frac{1}{R}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\ln (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, s$ coordinates. This results in

$$
\frac{s}{t}=-\ln (t)+c_{1}
$$

Which simplifies to

$$
\frac{s}{t}=-\ln (t)+c_{1}
$$

Which gives

$$
s=-t\left(\ln (t)-c_{1}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, s$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d s}{d t}=\frac{-t+s}{t}$ |  | $\frac{d S}{d R}=-\frac{1}{R}$ |
|  |  |  |
|  |  |  |
|  |  | $\xrightarrow[\rightarrow \infty \rightarrow \infty]{ }$ |
|  |  |  |
| -1. | $R=t$ | $\rightarrow$ - |
| $\frac{1}{x}$ | $S=\frac{s}{t}$ |  |
| - |  | $\xrightarrow{\sim}$ |
|  |  |  |
|  |  | $\mathrm{m}_{\rightarrow \pm \pm \pm}$ |
|  |  |  |

Summary
The solution(s) found are the following

$$
\begin{equation*}
s=-t\left(\ln (t)-c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 65: Slope field plot

## Verification of solutions

$$
s=-t\left(\ln (t)-c_{1}\right)
$$

Verified OK.

### 1.28.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, s) \mathrm{d} t+N(t, s) \mathrm{d} s=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(t) \mathrm{d} s & =(-t+s) \mathrm{d} t \\
(t-s) \mathrm{d} t+(t) \mathrm{d} s & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, s) & =t-s \\
N(t, s) & =t
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial s}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial s} & =\frac{\partial}{\partial s}(t-s) \\
& =-1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(t) \\
& =1
\end{aligned}
$$

Since $\frac{\partial M}{\partial s} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial s}-\frac{\partial N}{\partial t}\right) \\
& =\frac{1}{t}((-1)-(1)) \\
& =-\frac{2}{t}
\end{aligned}
$$

Since $A$ does not depend on $s$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int-\frac{2}{t} \mathrm{~d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-2 \ln (t)} \\
& =\frac{1}{t^{2}}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{t^{2}}(t-s) \\
& =\frac{t-s}{t^{2}}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{t^{2}}(t) \\
& =\frac{1}{t}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} s}{\mathrm{~d} t} & =0 \\
\left(\frac{t-s}{t^{2}}\right)+\left(\frac{1}{t}\right) \frac{\mathrm{d} s}{\mathrm{~d} t} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(t, s)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial s}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \frac{t-s}{t^{2}} \mathrm{~d} t \\
\phi & =\frac{s}{t}+\ln (t)+f(s) \tag{3}
\end{align*}
$$

Where $f(s)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $s$. Taking derivative of equation (3) w.r.t $s$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial s}=\frac{1}{t}+f^{\prime}(s) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial s}=\frac{1}{t}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{t}=\frac{1}{t}+f^{\prime}(s) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(s)$ gives

$$
f^{\prime}(s)=0
$$

Therefore

$$
f(s)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(s)$ into equation (3) gives $\phi$

$$
\phi=\frac{s}{t}+\ln (t)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{s}{t}+\ln (t)
$$

The solution becomes

$$
s=-t\left(\ln (t)-c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
s=-t\left(\ln (t)-c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 66: Slope field plot

Verification of solutions

$$
s=-t\left(\ln (t)-c_{1}\right)
$$

Verified OK.

### 1.28.5 Maple step by step solution

Let's solve
$-s+t s^{\prime}=-t$

- Highest derivative means the order of the ODE is 1
$s^{\prime}$
- Isolate the derivative
$s^{\prime}=-1+\frac{s}{t}$
- Group terms with $s$ on the lhs of the ODE and the rest on the rhs of the ODE $s^{\prime}-\frac{s}{t}=-1$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(s^{\prime}-\frac{s}{t}\right)=-\mu(t)$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) s)$
$\mu(t)\left(s^{\prime}-\frac{s}{t}\right)=\mu^{\prime}(t) s+\mu(t) s^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=-\frac{\mu(t)}{t}$
- $\quad$ Solve to find the integrating factor
$\mu(t)=\frac{1}{t}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) s)\right) d t=\int-\mu(t) d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) s=\int-\mu(t) d t+c_{1}$
- $\quad$ Solve for $s$
$s=\frac{\int-\mu(t) d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\frac{1}{t}$
$s=t\left(\int-\frac{1}{t} d t+c_{1}\right)$
- Evaluate the integrals on the rhs
$s=t\left(-\ln (t)+c_{1}\right)$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 12

```
dsolve((t-s(t))+t*diff(s(t),t)=0,s(t), singsol=all)
```

$$
s(t)=\left(c_{1}-\ln (t)\right) t
$$

$\checkmark$ Solution by Mathematica
Time used: 0.041 (sec). Leaf size: 14
DSolve[(t-s[t])+t*s'[t]==0,s[t],t,IncludeSingularSolutions $\rightarrow$ True]

$$
s(t) \rightarrow t\left(-\log (t)+c_{1}\right)
$$

### 1.29 problem 46

$$
\text { 1.29.1 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . } 346
$$

1.29.2 Solving as first order ode lie symmetry lookup ode ..... 348
1.29.3 Solving as bernoulli ode ..... 352
1.29.4 Solving as exact ode ..... 356

Internal problem ID [12446]

Internal file name [OUTPUT/11098_Monday_October_16_2023_09_47_53_PM_23089025/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 46.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "bernoulli", "homogeneousTypeD2", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry__lookup"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _Bernoulli]
```

$$
x y^{2} y^{\prime}-y^{3}=x^{3}
$$

### 1.29.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
x^{3} u(x)^{2}\left(u^{\prime}(x) x+u(x)\right)-u(x)^{3} x^{3}=x^{3}
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =\frac{1}{u^{2} x}
\end{aligned}
$$

Where $f(x)=\frac{1}{x}$ and $g(u)=\frac{1}{u^{2}}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{1}{u^{2}}} d u & =\frac{1}{x} d x \\
\int \frac{1}{\frac{1}{u^{2}}} d u & =\int \frac{1}{x} d x \\
\frac{u^{3}}{3} & =\ln (x)+c_{2}
\end{aligned}
$$

The solution is

$$
\frac{u(x)^{3}}{3}-\ln (x)-c_{2}=0
$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for $y$ in implicit form

$$
\begin{aligned}
& \frac{y^{3}}{3 x^{3}}-\ln (x)-c_{2}=0 \\
& \frac{y^{3}}{3 x^{3}}-\ln (x)-c_{2}=0
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{y^{3}}{3 x^{3}}-\ln (x)-c_{2}=0 \tag{1}
\end{equation*}
$$



Figure 67: Slope field plot

Verification of solutions

$$
\frac{y^{3}}{3 x^{3}}-\ln (x)-c_{2}=0
$$

Verified OK.

### 1.29.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{x^{3}+y^{3}}{x y^{2}} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type Bernoulli. Therefore we do not need to solve the $\operatorname{PDE}(\mathrm{A})$, and can just use the lookup table shown below to find $\xi, \eta$

Table 57: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
\xi(x, y) & =0 \\
\eta(x, y) & =\frac{x^{3}}{y^{2}} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{x^{3}}{y^{2}}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y^{3}}{3 x^{3}}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{x^{3}+y^{3}}{x y^{2}}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{y^{3}}{x^{4}} \\
S_{y} & =\frac{y^{2}}{x^{3}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{1}{x} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{1}{R}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\ln (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{y^{3}}{3 x^{3}}=\ln (x)+c_{1}
$$

Which simplifies to

$$
\frac{y^{3}}{3 x^{3}}=\ln (x)+c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{x^{3}+y^{3}}{x y^{2}}$ |  | $\frac{d S}{d R}=\frac{1}{R}$ |
|  |  |  |
|  |  | $\rightarrow \rightarrow \rightarrow \Delta x+1$ |
|  |  |  |
|  |  | STRU $+1+0 \rightarrow 0 \rightarrow 0$ |
|  |  |  |
|  | $R=x$ |  |
|  |  | $\cdots$ |
|  | $S=\frac{y}{3 x^{3}}$ | $\rightarrow \rightarrow$ ard |
|  |  | - $\sim_{4} \uparrow$ |
|  |  | - 4 |
|  |  | + 4 |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
\frac{y^{3}}{3 x^{3}}=\ln (x)+c_{1} \tag{1}
\end{equation*}
$$



Figure 68: Slope field plot
Verification of solutions

$$
\frac{y^{3}}{3 x^{3}}=\ln (x)+c_{1}
$$

Verified OK.

### 1.29.3 Solving as bernoulli ode

In canonical form, the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{x^{3}+y^{3}}{x y^{2}}
\end{aligned}
$$

This is a Bernoulli ODE.

$$
\begin{equation*}
y^{\prime}=\frac{1}{x} y+x^{2} \frac{1}{y^{2}} \tag{1}
\end{equation*}
$$

The standard Bernoulli ODE has the form

$$
\begin{equation*}
y^{\prime}=f_{0}(x) y+f_{1}(x) y^{n} \tag{2}
\end{equation*}
$$

The first step is to divide the above equation by $y^{n}$ which gives

$$
\begin{equation*}
\frac{y^{\prime}}{y^{n}}=f_{0}(x) y^{1-n}+f_{1}(x) \tag{3}
\end{equation*}
$$

The next step is use the substitution $w=y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$
\begin{aligned}
f_{0}(x) & =\frac{1}{x} \\
f_{1}(x) & =x^{2} \\
n & =-2
\end{aligned}
$$

Dividing both sides of ODE (1) by $y^{n}=\frac{1}{y^{2}}$ gives

$$
\begin{equation*}
y^{\prime} y^{2}=\frac{y^{3}}{x}+x^{2} \tag{4}
\end{equation*}
$$

Let

$$
\begin{align*}
w & =y^{1-n} \\
& =y^{3} \tag{5}
\end{align*}
$$

Taking derivative of equation (5) w.r.t $x$ gives

$$
\begin{equation*}
w^{\prime}=3 y^{2} y^{\prime} \tag{6}
\end{equation*}
$$

Substituting equations (5) and (6) into equation (4) gives

$$
\begin{align*}
\frac{w^{\prime}(x)}{3} & =\frac{w(x)}{x}+x^{2} \\
w^{\prime} & =\frac{3 w}{x}+3 x^{2} \tag{7}
\end{align*}
$$

The above now is a linear ODE in $w(x)$ which is now solved.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
w^{\prime}(x)+p(x) w(x)=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{3}{x} \\
& q(x)=3 x^{2}
\end{aligned}
$$

Hence the ode is

$$
w^{\prime}(x)-\frac{3 w(x)}{x}=3 x^{2}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{3}{x} d x} \\
& =\frac{1}{x^{3}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu w) & =(\mu)\left(3 x^{2}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{w}{x^{3}}\right) & =\left(\frac{1}{x^{3}}\right)\left(3 x^{2}\right) \\
\mathrm{d}\left(\frac{w}{x^{3}}\right) & =\left(\frac{3}{x}\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
\frac{w}{x^{3}} & =\int \frac{3}{x} \mathrm{~d} x \\
\frac{w}{x^{3}} & =3 \ln (x)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x^{3}}$ results in

$$
w(x)=3 x^{3} \ln (x)+c_{1} x^{3}
$$

which simplifies to

$$
w(x)=x^{3}\left(3 \ln (x)+c_{1}\right)
$$

Replacing $w$ in the above by $y^{3}$ using equation (5) gives the final solution.

$$
y^{3}=x^{3}\left(3 \ln (x)+c_{1}\right)
$$

Solving for $y$ gives

$$
\begin{aligned}
& y(x)=\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}} x \\
& y(x)=\frac{\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}}(i \sqrt{3}-1) x}{2} \\
& y(x)=-\frac{\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}}(1+i \sqrt{3}) x}{2}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}} x  \tag{1}\\
& y=\frac{\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}}(i \sqrt{3}-1) x}{2}  \tag{2}\\
& y=-\frac{\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}}(1+i \sqrt{3}) x}{2} \tag{3}
\end{align*}
$$



Figure 69: Slope field plot

## Verification of solutions

$$
y=\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}} x
$$

Verified OK.

$$
y=\frac{\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}}(i \sqrt{3}-1) x}{2}
$$

Verified OK.

$$
y=-\frac{\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}}(1+i \sqrt{3}) x}{2}
$$

Verified OK.

### 1.29.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition
$\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(x y^{2}\right) \mathrm{d} y & =\left(x^{3}+y^{3}\right) \mathrm{d} x \\
\left(-x^{3}-y^{3}\right) \mathrm{d} x+\left(x y^{2}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =-x^{3}-y^{3} \\
N(x, y) & =x y^{2}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-x^{3}-y^{3}\right) \\
& =-3 y^{2}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(x y^{2}\right) \\
& =y^{2}
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =\frac{1}{x y^{2}}\left(\left(-3 y^{2}\right)-\left(y^{2}\right)\right) \\
& =-\frac{4}{x}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-\frac{4}{x} \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-4 \ln (x)} \\
& =\frac{1}{x^{4}}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{x^{4}}\left(-x^{3}-y^{3}\right) \\
& =\frac{-x^{3}-y^{3}}{x^{4}}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{x^{4}}\left(x y^{2}\right) \\
& =\frac{y^{2}}{x^{3}}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x} & =0 \\
\left(\frac{-x^{3}-y^{3}}{x^{4}}\right)+\left(\frac{y^{2}}{x^{3}}\right) \frac{\mathrm{d} y}{\mathrm{~d} x} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{-x^{3}-y^{3}}{x^{4}} \mathrm{~d} x \\
\phi & =\frac{y^{3}}{3 x^{3}}-\ln (x)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\frac{y^{2}}{x^{3}}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{y^{2}}{x^{3}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{y^{2}}{x^{3}}=\frac{y^{2}}{x^{3}}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{y^{3}}{3 x^{3}}-\ln (x)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{y^{3}}{3 x^{3}}-\ln (x)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{y^{3}}{3 x^{3}}-\ln (x)=c_{1} \tag{1}
\end{equation*}
$$



Figure 70: Slope field plot

Verification of solutions

$$
\frac{y^{3}}{3 x^{3}}-\ln (x)=c_{1}
$$

Verified OK.
Maple trace

- Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful-
$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 58
dsolve $\left(x * y(x) \wedge 2 * \operatorname{diff}(y(x), x)=\left(x^{\wedge} 3+y(x) \wedge 3\right), y(x)\right.$, singsol=all)

$$
\begin{aligned}
& y(x)=\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}} x \\
& y(x)=-\frac{\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}}(1+i \sqrt{3}) x}{2} \\
& y(x)=\frac{\left(3 \ln (x)+c_{1}\right)^{\frac{1}{3}}(i \sqrt{3}-1) x}{2}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.331 (sec). Leaf size: 63
DSolve $\left[x * y[x] \sim 2 * y\right.$ ' $[x]==\left(x^{\wedge} 3+y[x] \wedge 3\right), y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow x \sqrt[3]{3 \log (x)+c_{1}} \\
& y(x) \rightarrow-\sqrt[3]{-1} x \sqrt[3]{3 \log (x)+c_{1}} \\
& y(x) \rightarrow(-1)^{2 / 3} x \sqrt[3]{3 \log (x)+c_{1}}
\end{aligned}
$$

### 1.30 problem 47

1.30.1 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 362
1.30.2 Solving as first order ode lie symmetry calculated ode . . . . . . 364
1.30.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 371

Internal problem ID [12447]
Internal file name [OUTPUT/11099_Monday_October_16_2023_09_48_01_PM_34561763/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 47.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "homogeneousTypeD2", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type
[[_homogeneous, `class A`], _dAlembert]

$$
x \cos \left(\frac{y}{x}\right)\left(y^{\prime} x+y\right)-y \sin \left(\frac{y}{x}\right)\left(y^{\prime} x-y\right)=0
$$

### 1.30.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$ $x \cos (u(x))\left(\left(u^{\prime}(x) x+u(x)\right) x+u(x) x\right)-u(x) x \sin (u(x))\left(\left(u^{\prime}(x) x+u(x)\right) x-u(x) x\right)=0$ In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =\frac{2 \cos (u) u}{x(u \sin (u)-\cos (u))}
\end{aligned}
$$

Where $f(x)=\frac{2}{x}$ and $g(u)=\frac{\cos (u) u}{u \sin (u)-\cos (u)}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{\cos (u) u}{u \sin (u)-\cos (u)}} d u & =\frac{2}{x} d x \\
\int \frac{1}{\frac{\cos (u) u}{u \sin (u)-\cos (u)}} d u & =\int \frac{2}{x} d x \\
-\ln (\cos (u))-\ln (u) & =2 \ln (x)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\mathrm{e}^{-\ln (\cos (u))-\ln (u)}=\mathrm{e}^{2 \ln (x)+c_{2}}
$$

Which simplifies to

$$
\frac{1}{\cos (u) u}=c_{3} x^{2}
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =x u \\
& =x \operatorname{RootOf}\left(\_Z c_{3} x^{2} \cos \left(\_Z\right)-1\right)
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x \operatorname{RootOf}\left(\_Z c_{3} x^{2} \cos \left(\_Z\right)-1\right) \tag{1}
\end{equation*}
$$



Figure 71: Slope field plot
Verification of solutions

$$
y=x \operatorname{RootOf}\left(\_Z c_{3} x^{2} \cos \left(\_Z\right)-1\right)
$$

Verified OK.

### 1.30.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=-\frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)}{x\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & -\frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)\left(b_{3}-a_{2}\right)}{x\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)}-\frac{y^{2}\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)^{2} a_{3}}{x^{2}\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)^{2}} \\
& -\left(-\frac{y\left(-\frac{y^{2} \cos \left(\frac{y}{x}\right)}{x^{2}}+\frac{y \sin \left(\frac{y}{x}\right)}{x}+\cos \left(\frac{y}{x}\right)\right)}{x\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)}+\frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)}{x^{2}\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)}\right. \\
& \left.+\frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)\left(\frac{y \sin \left(\frac{y}{x}\right)}{x}+\cos \left(\frac{y}{x}\right)+\frac{y^{2} \cos \left(\frac{y}{x}\right)}{x^{2}}\right)}{x\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)^{2}}\right)\left(x a_{2}+y a_{3}\right.  \tag{5E}\\
& \left.+a_{1}\right)-\left(-\frac{\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x}{x\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)}-\frac{y^{2} \cos \left(\frac{y}{x}\right)}{x^{2}\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)}\right. \\
& \left.+\frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)\left(-2 \sin \left(\frac{y}{x}\right)-\frac{\cos \left(\frac{y}{x}\right) y}{x}\right)}{x\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
$$

Putting the above in normal form gives
$2 \cos \left(\frac{y}{x}\right)^{2} x^{4} b_{2}-2 \cos \left(\frac{y}{x}\right)^{2} x^{2} y^{2} a_{3}+2 \cos \left(\frac{y}{x}\right)^{2} x^{2} y^{2} b_{2}-2 \cos \left(\frac{y}{x}\right)^{2} x y^{3} a_{2}+2 \cos \left(\frac{y}{x}\right)^{2} x y^{3} b_{3}-2 \cos \left(\frac{y}{x}\right)^{2} y$ $=0$

Setting the numerator to zero gives

$$
\begin{align*}
& 2 \cos \left(\frac{y}{x}\right)^{2} x^{4} b_{2}-2 \cos \left(\frac{y}{x}\right)^{2} x^{2} y^{2} a_{3}+2 \cos \left(\frac{y}{x}\right)^{2} x^{2} y^{2} b_{2} \\
& -2 \cos \left(\frac{y}{x}\right)^{2} x y^{3} a_{2}+2 \cos \left(\frac{y}{x}\right)^{2} x y^{3} b_{3}-2 \cos \left(\frac{y}{x}\right)^{2} y^{4} a_{3} \\
& -2 \cos \left(\frac{y}{x}\right) \sin \left(\frac{y}{x}\right) x^{2} y^{2} a_{2}+2 \cos \left(\frac{y}{x}\right) \sin \left(\frac{y}{x}\right) x^{2} y^{2} b_{3} \\
& -4 \cos \left(\frac{y}{x}\right) \sin \left(\frac{y}{x}\right) x y^{3} a_{3}+2 \sin \left(\frac{y}{x}\right)^{2} x^{2} y^{2} b_{2}-2 \sin \left(\frac{y}{x}\right)^{2} x y^{3} a_{2}  \tag{6E}\\
& +2 \sin \left(\frac{y}{x}\right)^{2} x y^{3} b_{3}-2 \sin \left(\frac{y}{x}\right)^{2} y^{4} a_{3}+\cos \left(\frac{y}{x}\right)^{2} x^{3} b_{1}-\cos \left(\frac{y}{x}\right)^{2} x^{2} y a_{1} \\
& +2 \cos \left(\frac{y}{x}\right)^{2} x y^{2} b_{1}-2 \cos \left(\frac{y}{x}\right)^{2} y^{3} a_{1}+2 \cos \left(\frac{y}{x}\right) \sin \left(\frac{y}{x}\right) x^{2} y b_{1} \\
& -2 \cos \left(\frac{y}{x}\right) \sin \left(\frac{y}{x}\right) x y^{2} a_{1}+\sin \left(\frac{y}{x}\right)^{2} x y^{2} b_{1}-\sin \left(\frac{y}{x}\right)^{2} y^{3} a_{1}=0
\end{align*}
$$

Simplifying the above gives

$$
\begin{aligned}
& x\left(-2 x^{2} y^{2} a_{2} \sin \left(\frac{2 y}{x}\right)+2 x^{2} y^{2} b_{3} \sin \left(\frac{2 y}{x}\right)-4 x y^{3} a_{3} \sin \left(\frac{2 y}{x}\right)+2 x^{4} b_{2} \cos \left(\frac{2 y}{x}\right)-2 x^{2} y^{2} a_{3} \cos \left(\frac{2 y}{x}\right)+2 x^{2} y b_{1} \operatorname{sir}\right. \\
& =0
\end{aligned}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\left\{x, y, \cos \left(\frac{2 y}{x}\right), \sin \left(\frac{2 y}{x}\right)\right\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}, \cos \left(\frac{2 y}{x}\right)=v_{3}, \sin \left(\frac{2 y}{x}\right)=v_{4}\right\}
$$

The above PDE (6E) now becomes
$\underline{v_{1}\left(-2 v_{1}^{2} v_{2}^{2} a_{2} v_{4}-2 v_{1}^{2} v_{2}^{2} a_{3} v_{3}-4 v_{1} v_{2}^{3} a_{3} v_{4}+2 v_{1}^{4} b_{2} v_{3}+2 v_{1}^{2} v_{2}^{2} b_{3} v_{4}-v_{1}^{2} v_{2} a_{1} v_{3}-2 v_{1} v_{2}^{2} a_{\left(7 \mathrm{H}_{4}\right)}-v_{2}^{3} a_{1} v_{3}-4 v_{1} v_{2}^{3} c\right.}$ $=0$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}, v_{3}, v_{4}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& b_{2} v_{3} v_{1}^{5}+b_{2} v_{1}^{5}+\frac{b_{1} v_{1}^{4}}{2}+\frac{b_{1} v_{3} v_{1}^{4}}{2}+\left(b_{3}-a_{2}\right) v_{4} v_{2}^{2} v_{1}^{3}-a_{3} v_{3} v_{2}^{2} v_{1}^{3}-\frac{a_{1} v_{3} v_{2} v_{1}^{3}}{2} \\
& +b_{1} v_{4} v_{2} v_{1}^{3}+\left(-a_{3}+2 b_{2}\right) v_{2}^{2} v_{1}^{3}-\frac{a_{1} v_{2} v_{1}^{3}}{2}-a_{1} v_{4} v_{2}^{2} v_{1}^{2}+\frac{b_{1} v_{3} v_{2}^{2} v_{1}^{2}}{2}-2 a_{3} v_{4} v_{2}^{3} v_{1}^{2}  \tag{8E}\\
& \quad+\frac{3 b_{1} v_{2}^{2} v_{1}^{2}}{2}+\left(-2 a_{2}+2 b_{3}\right) v_{2}^{3} v_{1}^{2}-2 a_{3} v_{2}^{4} v_{1}-\frac{a_{1} v_{2}^{3} v_{3} v_{1}}{2}-\frac{3 a_{1} v_{2}^{3} v_{1}}{2}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
b_{1} & =0 \\
b_{2} & =0 \\
-a_{1} & =0 \\
-\frac{3 a_{1}}{2} & =0 \\
-\frac{a_{1}}{2} & =0 \\
-2 a_{3} & =0 \\
-a_{3} & =0 \\
\frac{b_{1}}{2} & =0 \\
\frac{3 b_{1}}{2} & =0 \\
-2 a_{2}+2 b_{3} & =0 \\
-a_{3}+2 b_{2} & =0 \\
b_{3}-a_{2} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =b_{3} \\
a_{3} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=x \\
& \eta=y
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =y-\left(-\frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)}{x\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)}\right)(x) \\
& =\frac{2 y \cos \left(\frac{y}{x}\right) x}{\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{2 y \cos \left(\frac{y}{x}\right) x}{\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{\ln \left(\cos \left(\frac{y}{x}\right)\right)}{2}+\frac{\ln \left(\frac{y}{x}\right)}{2}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)}{x\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{-x+\tan \left(\frac{y}{x}\right) y}{2 x^{2}} \\
S_{y} & =-\frac{\tan \left(\frac{y}{x}\right)}{2 x}+\frac{1}{2 y}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-\frac{1}{x} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-\frac{1}{R}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\ln (R)+c_{1} \tag{4}
\end{equation*}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown．

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=-\frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)}{x\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)}$ |  | $\frac{d S}{d R}=-\frac{1}{R}$ |
|  |  |  |
|  |  | $\rightarrow \rightarrow \infty$ |
|  |  |  |
|  |  |  |
| $\rightarrow \rightarrow \rightarrow \infty \times$ 乐 9 | $R=x$ | $0 \rightarrow 49$ |
|  | $\ln \left(\cos \left(\frac{y}{x}\right)\right)-\ln (x)$ |  |
|  | $S=\frac{2}{2}-\frac{1}{2}$ | $0 \rightarrow$ 分 |
|  |  | 勿分幸 |
|  |  |  |
|  |  | － A $^{4}$ |
|  |  | $\rightarrow \rightarrow \rightarrow$ 为 |

## Summary

The solution（s）found are the following

$$
\begin{equation*}
\frac{\ln \left(\cos \left(\frac{y}{x}\right)\right)}{2}-\frac{\ln (x)}{2}+\frac{\ln (y)}{2}=-\ln (x)+c_{1} \tag{1}
\end{equation*}
$$



Figure 72: Slope field plot

## Verification of solutions

$$
\frac{\ln \left(\cos \left(\frac{y}{x}\right)\right)}{2}-\frac{\ln (x)}{2}+\frac{\ln (y)}{2}=-\ln (x)+c_{1}
$$

Verified OK.

### 1.30.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(-y \sin \left(\frac{y}{x}\right) x+x^{2} \cos \left(\frac{y}{x}\right)\right) \mathrm{d} y & =\left(-\sin \left(\frac{y}{x}\right) y^{2}-y \cos \left(\frac{y}{x}\right) x\right) \mathrm{d} x \\
\left(\sin \left(\frac{y}{x}\right) y^{2}+y \cos \left(\frac{y}{x}\right) x\right) \mathrm{d} x+\left(-y \sin \left(\frac{y}{x}\right) x+x^{2} \cos \left(\frac{y}{x}\right)\right) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=\sin \left(\frac{y}{x}\right) y^{2}+y \cos \left(\frac{y}{x}\right) x \\
& N(x, y)=-y \sin \left(\frac{y}{x}\right) x+x^{2} \cos \left(\frac{y}{x}\right)
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\sin \left(\frac{y}{x}\right) y^{2}+y \cos \left(\frac{y}{x}\right) x\right) \\
& =\frac{\left(x^{2}+y^{2}\right) \cos \left(\frac{y}{x}\right)+y \sin \left(\frac{y}{x}\right) x}{x}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(-y \sin \left(\frac{y}{x}\right) x+x^{2} \cos \left(\frac{y}{x}\right)\right) \\
& =\cos \left(\frac{y}{x}\right)\left(\frac{y^{2}}{x}+2 x\right)
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =\frac{1}{x\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right)}\left(\left(\frac{\cos \left(\frac{y}{x}\right) y^{2}}{x}+\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)-\left(\frac{\cos \left(\frac{y}{x}\right) y^{2}}{x}+2 \cos \left(\frac{y}{x}\right) x\right)\right) \\
& =-\frac{1}{x}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-\frac{1}{x} \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-\ln (x)} \\
& =\frac{1}{x}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{x}\left(\sin \left(\frac{y}{x}\right) y^{2}+y \cos \left(\frac{y}{x}\right) x\right) \\
& =\frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)}{x}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{x}\left(-y \sin \left(\frac{y}{x}\right) x+x^{2} \cos \left(\frac{y}{x}\right)\right) \\
& =\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x} & =0 \\
\left(\frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)}{x}\right)+\left(\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y\right) \frac{\mathrm{d} y}{\mathrm{~d} x} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{y\left(\sin \left(\frac{y}{x}\right) y+\cos \left(\frac{y}{x}\right) x\right)}{x} \mathrm{~d} x \\
\phi & =y \cos \left(\frac{y}{x}\right) x+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y$. Therefore equation (4) becomes

$$
\begin{equation*}
\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y=\cos \left(\frac{y}{x}\right) x-\sin \left(\frac{y}{x}\right) y+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=y \cos \left(\frac{y}{x}\right) x+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=y \cos \left(\frac{y}{x}\right) x
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y \cos \left(\frac{y}{x}\right) x=c_{1} \tag{1}
\end{equation*}
$$



Figure 73: Slope field plot
Verification of solutions

$$
y \cos \left(\frac{y}{x}\right) x=c_{1}
$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.265 (sec). Leaf size: 18
$\operatorname{dsolve}(x * \cos (y(x) / x) *(y(x)+x * \operatorname{diff}(y(x), x))=y(x) * \sin (y(x) / x) *(x * \operatorname{diff}(y(x), x)-y(x)), y(x), \quad \operatorname{sing}$

$$
y(x)=x \operatorname{RootOf}\left(\_Z x^{2} \cos \left(\_Z\right)-c_{1}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.569 (sec). Leaf size: 31
DSolve $\left[x * \operatorname{Cos}[y[x] / x] *\left(y[x]+x * y{ }^{\prime}[x]\right)==y[x] * \operatorname{Sin}[y[x] / x] *(x * y '[x]-y[x]), y[x], x\right.$, IncludeSingulars

Solve $\left[-\log \left(\frac{y(x)}{x}\right)-\log \left(\cos \left(\frac{y(x)}{x}\right)\right)=2 \log (x)+c_{1}, y(x)\right]$

### 1.31 problem 48

$$
\text { 1.31.1 Solving as homogeneousTypeMapleC ode . . . . . . . . . . . . . } 377
$$

1.31.2 Solving as first order ode lie symmetry calculated ode . . . . . . 381

Internal problem ID [12448]
Internal file name [OUTPUT/11100_Monday_October_16_2023_09_48_05_PM_65448582/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 48.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "homogeneousTypeMapleC", "first_order_ode_lie_symmetry__calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class C`], _rational, [_Abel, `2nd type`,`
    class A`]]
```

$$
3 y-(3 x-7 y-3) y^{\prime}=7 x-7
$$

### 1.31.1 Solving as homogeneousTypeMapleC ode

Let $Y=y+y_{0}$ and $X=x+x_{0}$ then the above is transformed to new ode in $Y(X)$

$$
\frac{d}{d X} Y(X)=-\frac{3 Y(X)+3 y_{0}-7 X-7 x_{0}+7}{-3 X-3 x_{0}+7 Y(X)+7 y_{0}+3}
$$

Solving for possible values of $x_{0}$ and $y_{0}$ which makes the above ode a homogeneous ode results in

$$
\begin{aligned}
& x_{0}=1 \\
& y_{0}=0
\end{aligned}
$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$
\frac{d}{d X} Y(X)=-\frac{3 Y(X)-7 X}{-3 X+7 Y(X)}
$$

In canonical form, the ODE is

$$
\begin{align*}
Y^{\prime} & =F(X, Y) \\
& =-\frac{3 Y-7 X}{-3 X+7 Y} \tag{1}
\end{align*}
$$

An ode of the form $Y^{\prime}=\frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order $n$ if

$$
f\left(t^{n} X, t^{n} Y\right)=t^{n} f(X, Y)
$$

In this case, it can be seen that both $M=3 Y-7 X$ and $N=3 X-7 Y$ are both homogeneous and of the same order $n=1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u=\frac{Y}{X}$, or $Y=u X$. Hence

$$
\frac{\mathrm{d} Y}{\mathrm{~d} X}=\frac{\mathrm{d} u}{\mathrm{~d} X} X+u
$$

Applying the transformation $Y=u X$ to the above ODE in (1) gives

$$
\begin{aligned}
\frac{\mathrm{d} u}{\mathrm{~d} X} X+u & =\frac{-3 u+7}{7 u-3} \\
\frac{\mathrm{~d} u}{\mathrm{~d} X} & =\frac{\frac{-3 u(X)+7}{7 u(X)-3}-u(X)}{X}
\end{aligned}
$$

Or

$$
\frac{d}{d X} u(X)-\frac{\frac{-3 u(X)+7}{7 u(X)-3}-u(X)}{X}=0
$$

Or

$$
7\left(\frac{d}{d X} u(X)\right) X u(X)-3\left(\frac{d}{d X} u(X)\right) X+7 u(X)^{2}-7=0
$$

Or

$$
-7+X(7 u(X)-3)\left(\frac{d}{d X} u(X)\right)+7 u(X)^{2}=0
$$

Which is now solved as separable in $u(X)$. Which is now solved in $u(X)$. In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(X, u) \\
& =f(X) g(u) \\
& =-\frac{7\left(u^{2}-1\right)}{X(7 u-3)}
\end{aligned}
$$

Where $f(X)=-\frac{7}{X}$ and $g(u)=\frac{u^{2}-1}{7 u-3}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{u^{2}-1}{7 u-3}} d u & =-\frac{7}{X} d X \\
\int \frac{1}{\frac{u^{2}-1}{7 u-3}} d u & =\int-\frac{7}{X} d X \\
2 \ln (u-1)+5 \ln (u+1) & =-7 \ln (X)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\mathrm{e}^{2 \ln (u-1)+5 \ln (u+1)}=\mathrm{e}^{-7 \ln (X)+c_{2}}
$$

Which simplifies to

$$
(u-1)^{2}(u+1)^{5}=\frac{c_{3}}{X^{7}}
$$

Now $u$ in the above solution is replaced back by $Y$ using $u=\frac{Y}{X}$ which results in the solution
$Y(X)=$ RootOf $\left(X^{7}+3 X^{6} \_Z+X^{5} \_Z^{2}-5 X^{4} \_Z^{3}-5 X^{3} \_^{4}+X^{2} \_^{5}+3 X \_Z^{6}+\_Z^{7}-c_{3}\right)$
Using the solution for $Y(X)$
$Y(X)=$ RootOf $\left(X^{7}+3 X^{6} \_Z+X^{5} \_Z^{2}-5 X^{4} \_Z^{3}-5 X^{3} \_Z^{4}+X^{2} \_^{5}+3 X \_Z^{6}+\_Z^{7}-c_{3}\right)$
And replacing back terms in the above solution using

$$
\begin{aligned}
& Y=y+y_{0} \\
& X=x+x_{0}
\end{aligned}
$$

Or

$$
\begin{aligned}
& Y=y \\
& X=x+1
\end{aligned}
$$

Then the solution in $y$ becomes
$y=\operatorname{RootOf}\left(\_Z^{7}+(-3+3 x) \_Z^{6}+\left(x^{2}-2 x+1\right) \_Z^{5}+\left(-5 x^{3}+15 x^{2}-15 x+5\right) \_Z^{4}+\left(-5 x^{4}+20 x\right.\right.$

## Summary

The solution(s) found are the following

$$
\begin{array}{r}
y=\operatorname{RootOf}\left(\_Z^{7}+(-3+3 x) \_Z^{6}+\left(x^{2}-2 x+1\right) \_Z^{5}+\left(-5 x^{3}+15 x^{2}-15 x+5\right) \_Z^{4}\right. \\
+\left(-5 x^{4}+20 x^{3}-30 x^{2}+20 x-5\right) \_Z^{3}+\left(x^{5}-5 x^{4}+10 x^{3}-10 x^{2}+5 x-1\right) \_Z^{2} \\
+\left(3 x^{6}-18 x^{5}+45 x^{4}-60 x^{3}+45 x^{2}-18 x+3\right) \_Z+x^{7}-7 x^{6}+21 x^{5}-35 x^{4} \\
\left.+35 x^{3}-21 x^{2}-c_{3}+7 x-1\right) \tag{1}
\end{array}
$$



Figure 74: Slope field plot

Verification of solutions

$$
\begin{array}{r}
y=\operatorname{RootOf}\left(\_Z^{7}+(-3+3 x) \_Z^{6}+\left(x^{2}-2 x+1\right) \_Z^{5}+\left(-5 x^{3}+15 x^{2}-15 x+5\right) \_Z^{4}\right. \\
+\left(-5 x^{4}+20 x^{3}-30 x^{2}+20 x-5\right) \_Z^{3}+\left(x^{5}-5 x^{4}+10 x^{3}-10 x^{2}+5 x-1\right) \_Z^{2} \\
+\left(3 x^{6}-18 x^{5}+45 x^{4}-60 x^{3}+45 x^{2}-18 x+3\right) \_Z+x^{7}-7 x^{6}+21 x^{5}-35 x^{4} \\
\left.+35 x^{3}-21 x^{2}-c_{3}+7 x-1\right)
\end{array}
$$

Verified OK.

### 1.31.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{3 y-7 x+7}{-3 x+7 y+3} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & -\frac{(3 y-7 x+7)\left(b_{3}-a_{2}\right)}{-3 x+7 y+3}-\frac{(3 y-7 x+7)^{2} a_{3}}{(-3 x+7 y+3)^{2}} \\
& -\left(\frac{7}{-3 x+7 y+3}-\frac{3(3 y-7 x+7)}{(-3 x+7 y+3)^{2}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\left(-\frac{3}{-3 x+7 y+3}+\frac{21 y-49 x+49}{(-3 x+7 y+3)^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
$$

Putting the above in normal form gives
$\underline{21 x^{2} a_{2}-49 x^{2} a_{3}+49 x^{2} b_{2}-21 x^{2} b_{3}-98 x y a_{2}+42 x y a_{3}-42 x y b_{2}+98 x y b_{3}+21 y^{2} a_{2}-49 y^{2} a_{3}+49 y^{2} b_{2}-1 . . ~}$ $=0$

Setting the numerator to zero gives

$$
\begin{align*}
& 21 x^{2} a_{2}-49 x^{2} a_{3}+49 x^{2} b_{2}-21 x^{2} b_{3}-98 x y a_{2}+42 x y a_{3}-42 x y b_{2}+98 x y b_{3}  \tag{6E}\\
& +21 y^{2} a_{2}-49 y^{2} a_{3}+49 y^{2} b_{2}-21 y^{2} b_{3}-42 x a_{2}+98 x a_{3}+40 x b_{1}-58 x b_{2}+42 x b_{3} \\
& \quad-40 y a_{1}+58 y a_{2}-42 y a_{3}+42 y b_{2}-98 y b_{3}+21 a_{2}-49 a_{3}-40 b_{1}+9 b_{2}-21 b_{3}=0
\end{align*}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\{x, y\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& 21 a_{2} v_{1}^{2}-98 a_{2} v_{1} v_{2}+21 a_{2} v_{2}^{2}-49 a_{3} v_{1}^{2}+42 a_{3} v_{1} v_{2}-49 a_{3} v_{2}^{2}+49 b_{2} v_{1}^{2} \\
& \quad-42 b_{2} v_{1} v_{2}+49 b_{2} v_{2}^{2}-21 b_{3} v_{1}^{2}+98 b_{3} v_{1} v_{2}-21 b_{3} v_{2}^{2}-40 a_{1} v_{2}  \tag{7E}\\
& \quad-42 a_{2} v_{1}+58 a_{2} v_{2}+98 a_{3} v_{1}-42 a_{3} v_{2}+40 b_{1} v_{1}-58 b_{2} v_{1}+42 b_{2} v_{2} \\
& \quad+42 b_{3} v_{1}-98 b_{3} v_{2}+21 a_{2}-49 a_{3}-40 b_{1}+9 b_{2}-21 b_{3}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(21 a_{2}-49 a_{3}+49 b_{2}-21 b_{3}\right) v_{1}^{2}+\left(-98 a_{2}+42 a_{3}-42 b_{2}+98 b_{3}\right) v_{1} v_{2}  \tag{8E}\\
& \quad+\left(-42 a_{2}+98 a_{3}+40 b_{1}-58 b_{2}+42 b_{3}\right) v_{1}+\left(21 a_{2}-49 a_{3}+49 b_{2}-21 b_{3}\right) v_{2}^{2} \\
& \quad+\left(-40 a_{1}+58 a_{2}-42 a_{3}+42 b_{2}-98 b_{3}\right) v_{2}+21 a_{2}-49 a_{3}-40 b_{1}+9 b_{2}-21 b_{3}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{array}{r}
-98 a_{2}+42 a_{3}-42 b_{2}+98 b_{3}=0 \\
21 a_{2}-49 a_{3}+49 b_{2}-21 b_{3}=0 \\
-40 a_{1}+58 a_{2}-42 a_{3}+42 b_{2}-98 b_{3}=0 \\
-42 a_{2}+98 a_{3}+40 b_{1}-58 b_{2}+42 b_{3}=0 \\
21 a_{2}-49 a_{3}-40 b_{1}+9 b_{2}-21 b_{3}=0
\end{array}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =-b_{3} \\
a_{2} & =b_{3} \\
a_{3} & =b_{2} \\
b_{1} & =-b_{2} \\
b_{2} & =b_{2} \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E, 2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=y \\
& \eta=x-1
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =x-1-\left(-\frac{3 y-7 x+7}{-3 x+7 y+3}\right)(y) \\
& =\frac{3 x^{2}-3 y^{2}-6 x+3}{3 x-7 y-3} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{3 x^{2}-3 y^{2}-6 x+3}{3 x-7 y-3}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{2 \ln (1-x+y)}{3}+\frac{5 \ln (x-1+y)}{3}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{3 y-7 x+7}{-3 x+7 y+3}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{2}{3 x-3-3 y}+\frac{5}{3 x-3+3 y} \\
S_{y} & =-\frac{2}{3 x-3-3 y}+\frac{5}{3 x-3+3 y}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{2 \ln (1-x+y)}{3}+\frac{5 \ln (y-1+x)}{3}=c_{1}
$$

Which simplifies to

$$
\frac{2 \ln (1-x+y)}{3}+\frac{5 \ln (y-1+x)}{3}=c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown．

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=-\frac{3 y-7 x+7}{-3 x+7 y+3}$ |  | $\frac{d S}{d R}=0$ |
|  |  | $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ |
| divivivivivic） |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{ }$ |
|  |  |  |
|  |  |  |
|  | $R=x$ | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{ }$ |
|  | $S=\frac{2 \ln (1-x+y)}{3}$ |  |
| $\underline{L}$ | $S=\frac{3}{3}+$ |  |
| $1 x^{1+4}$ |  | $\xrightarrow{ }$ |
|  |  |  |
| 多多办他 $\rightarrow \rightarrow \rightarrow-\infty \rightarrow+\infty$ |  |  |

## Summary

The solution（s）found are the following

$$
\begin{equation*}
\frac{2 \ln (1-x+y)}{3}+\frac{5 \ln (y-1+x)}{3}=c_{1} \tag{1}
\end{equation*}
$$



Figure 75: Slope field plot
Verification of solutions

$$
\frac{2 \ln (1-x+y)}{3}+\frac{5 \ln (y-1+x)}{3}=c_{1}
$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous C
trying homogeneous types:
trying homogeneous D
<- homogeneous successful
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 1.0 (sec). Leaf size: 1814

```
dsolve((3*y(x)-7*x+7)-(3*x-7*y(x)-3)*diff (y(x),x)=0,y(x), singsol=all)
```

Expression too large to display
$\sqrt{ }$ Solution by Mathematica
Time used: 61.254 (sec). Leaf size: 7785
DSolve[(3*y[x]-7*x+7)-(3*x-7*y[x]-3)*y'[x]==0,y[x],x,IncludeSingularSolutions $\rightarrow$ True]

Too large to display

### 1.32 problem 49

1.32.1 Solving as first order ode lie symmetry calculated ode

Internal problem ID [12449]
Internal file name [OUTPUT/11101_Monday_October_16_2023_09_48_10_PM_32949649/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 49.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "first__order_ode__lie_symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class C`], _rational, [_Abel, `2nd type`, `
    class A`]]
```

$$
2 y-(4 y+2 x+3) y^{\prime}=-x-1
$$

### 1.32.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{x+2 y+1}{4 y+2 x+3} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E, 2 E ) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & +\frac{(x+2 y+1)\left(b_{3}-a_{2}\right)}{4 y+2 x+3}-\frac{(x+2 y+1)^{2} a_{3}}{(4 y+2 x+3)^{2}} \\
& -\left(\frac{1}{4 y+2 x+3}-\frac{2(x+2 y+1)}{(4 y+2 x+3)^{2}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\left(\frac{2}{4 y+2 x+3}-\frac{4(x+2 y+1)}{(4 y+2 x+3)^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
$$

Putting the above in normal form gives

$$
\begin{aligned}
& -\frac{2 x^{2} a_{2}+x^{2} a_{3}-4 x^{2} b_{2}-2 x^{2} b_{3}+8 x y a_{2}+4 x y a_{3}-16 x y b_{2}-8 x y b_{3}+8 y^{2} a_{2}+4 y^{2} a_{3}-16 y^{2} b_{2}-8 y^{2} b_{3}+}{(4 y+2 x+} \\
& =0
\end{aligned}
$$

Setting the numerator to zero gives

$$
\begin{align*}
& -2 x^{2} a_{2}-x^{2} a_{3}+4 x^{2} b_{2}+2 x^{2} b_{3}-8 x y a_{2}-4 x y a_{3}+16 x y b_{2}+8 x y b_{3}  \tag{6E}\\
& \quad-8 y^{2} a_{2}-4 y^{2} a_{3}+16 y^{2} b_{2}+8 y^{2} b_{3}-6 x a_{2}-2 x a_{3}+10 x b_{2}+5 x b_{3} \\
& \quad-10 y a_{2}-5 y a_{3}+24 y b_{2}+8 y b_{3}-a_{1}-3 a_{2}-a_{3}-2 b_{1}+9 b_{2}+3 b_{3}=0
\end{align*}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\{x, y\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& -2 a_{2} v_{1}^{2}-8 a_{2} v_{1} v_{2}-8 a_{2} v_{2}^{2}-a_{3} v_{1}^{2}-4 a_{3} v_{1} v_{2}-4 a_{3} v_{2}^{2}+4 b_{2} v_{1}^{2}+16 b_{2} v_{1} v_{2}  \tag{7E}\\
& +16 b_{2} v_{2}^{2}+2 b_{3} v_{1}^{2}+8 b_{3} v_{1} v_{2}+8 b_{3} v_{2}^{2}-6 a_{2} v_{1}-10 a_{2} v_{2}-2 a_{3} v_{1}-5 a_{3} v_{2} \\
& +10 b_{2} v_{1}+24 b_{2} v_{2}+5 b_{3} v_{1}+8 b_{3} v_{2}-a_{1}-3 a_{2}-a_{3}-2 b_{1}+9 b_{2}+3 b_{3}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(-2 a_{2}-a_{3}+4 b_{2}+2 b_{3}\right) v_{1}^{2}+\left(-8 a_{2}-4 a_{3}+16 b_{2}+8 b_{3}\right) v_{1} v_{2}  \tag{8E}\\
& \quad+\left(-6 a_{2}-2 a_{3}+10 b_{2}+5 b_{3}\right) v_{1}+\left(-8 a_{2}-4 a_{3}+16 b_{2}+8 b_{3}\right) v_{2}^{2} \\
& \quad+\left(-10 a_{2}-5 a_{3}+24 b_{2}+8 b_{3}\right) v_{2}-a_{1}-3 a_{2}-a_{3}-2 b_{1}+9 b_{2}+3 b_{3}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
-10 a_{2}-5 a_{3}+24 b_{2}+8 b_{3} & =0 \\
-8 a_{2}-4 a_{3}+16 b_{2}+8 b_{3} & =0 \\
-6 a_{2}-2 a_{3}+10 b_{2}+5 b_{3} & =0 \\
-2 a_{2}-a_{3}+4 b_{2}+2 b_{3} & =0 \\
-a_{1}-3 a_{2}-a_{3}-2 b_{1}+9 b_{2}+3 b_{3} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
& a_{1}=5 b_{2}-2 b_{1} \\
& a_{2}=2 b_{2} \\
& a_{3}=4 b_{2} \\
& b_{1}=b_{1} \\
& b_{2}=b_{2} \\
& b_{3}=2 b_{2}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E, 2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=-2 \\
& \eta=1
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =1-\left(\frac{x+2 y+1}{4 y+2 x+3}\right)(-2) \\
& =\frac{4 x+8 y+5}{4 y+2 x+3} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{4 x+8 y+5}{4 y+2 x+3}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{2}+\frac{\ln (4 x+8 y+5)}{16}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{x+2 y+1}{4 y+2 x+3}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{1}{16 x+32 y+20} \\
S_{y} & =\frac{4 y+2 x+3}{4 x+8 y+5}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{1}{4} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{1}{4}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{R}{4}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{y}{2}+\frac{\ln (4 x+8 y+5)}{16}=\frac{x}{4}+c_{1}
$$

Which simplifies to

$$
\frac{y}{2}+\frac{\ln (4 x+8 y+5)}{16}=\frac{x}{4}+c_{1}
$$

Which gives

$$
y=\frac{\text { LambertW }\left(\mathrm{e}^{8 x+5+16 c_{1}}\right)}{8}-\frac{x}{2}-\frac{5}{8}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.


## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\text { LambertW }\left(\mathrm{e}^{8 x+5+16 c_{1}}\right)}{8}-\frac{x}{2}-\frac{5}{8} \tag{1}
\end{equation*}
$$



Figure 76: Slope field plot

## Verification of solutions

$$
y=\frac{\text { LambertW }\left(\mathrm{e}^{8 x+5+16 c_{1}}\right)}{8}-\frac{x}{2}-\frac{5}{8}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous C
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.047 (sec). Leaf size: 20
dsolve $((x+2 * y(x)+1)-(2 * x+4 * y(x)+3) * \operatorname{diff}(y(x), x)=0, y(x)$, singsol=all)

$$
y(x)=-\frac{x}{2}+\frac{\text { LambertW }\left(c_{1} \mathrm{e}^{5+8 x}\right)}{8}-\frac{5}{8}
$$

$\sqrt{ }$ Solution by Mathematica
Time used: 6.228 (sec). Leaf size: 39
DSolve $\left[(x+2 * y[x]+1)-(2 * x+4 * y[x]+3) * y{ }^{\prime}[x]==0, y[x], x\right.$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow \frac{1}{8}\left(W\left(-e^{8 x-1+c_{1}}\right)-4 x-5\right) \\
& y(x) \rightarrow \frac{1}{8}(-4 x-5)
\end{aligned}
$$

### 1.33 problem 50

1.33.1 Solving as linear ode ..... 396
1.33.2 Solving as homogeneousTypeMapleC ode ..... 398
1.33.3 Solving as first order ode lie symmetry lookup ode ..... 401
1.33.4 Solving as exact ode ..... 405
1.33.5 Maple step by step solution ..... 410

Internal problem ID [12450]
Internal file name [OUTPUT/11102_Monday_October_16_2023_09_48_13_PM_50550288/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 50.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "homogeneousTypeMapleC", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"
Maple gives the following as the ode type
[_linear]

$$
2 y-(2 x-3) y^{\prime}=-x-1
$$

### 1.33.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{2}{2 x-3} \\
& q(x)=\frac{x+1}{2 x-3}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{2 y}{2 x-3}=\frac{x+1}{2 x-3}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2}{2 x-3} d x} \\
& =\frac{1}{2 x-3}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\frac{x+1}{2 x-3}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{y}{2 x-3}\right) & =\left(\frac{1}{2 x-3}\right)\left(\frac{x+1}{2 x-3}\right) \\
\mathrm{d}\left(\frac{y}{2 x-3}\right) & =\left(\frac{x+1}{(2 x-3)^{2}}\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{y}{2 x-3}=\int \frac{x+1}{(2 x-3)^{2}} \mathrm{~d} x \\
& \frac{y}{2 x-3}=\frac{\ln (2 x-3)}{4}-\frac{5}{4(2 x-3)}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{2 x-3}$ results in

$$
y=(2 x-3)\left(\frac{\ln (2 x-3)}{4}-\frac{5}{4(2 x-3)}\right)+c_{1}(2 x-3)
$$

which simplifies to

$$
y=-\frac{5}{4}+\frac{(2 x-3) \ln (2 x-3)}{4}+c_{1}(2 x-3)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{5}{4}+\frac{(2 x-3) \ln (2 x-3)}{4}+c_{1}(2 x-3) \tag{1}
\end{equation*}
$$



Figure 77: Slope field plot
Verification of solutions

$$
y=-\frac{5}{4}+\frac{(2 x-3) \ln (2 x-3)}{4}+c_{1}(2 x-3)
$$

Verified OK.

### 1.33.2 Solving as homogeneousTypeMapleC ode

Let $Y=y+y_{0}$ and $X=x+x_{0}$ then the above is transformed to new ode in $Y(X)$

$$
\frac{d}{d X} Y(X)=\frac{2 Y(X)+2 y_{0}+X+x_{0}+1}{2 X+2 x_{0}-3}
$$

Solving for possible values of $x_{0}$ and $y_{0}$ which makes the above ode a homogeneous ode results in

$$
\begin{aligned}
& x_{0}=\frac{3}{2} \\
& y_{0}=-\frac{5}{4}
\end{aligned}
$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$
\frac{d}{d X} Y(X)=\frac{2 Y(X)+X}{2 X}
$$

In canonical form, the ODE is

$$
\begin{align*}
Y^{\prime} & =F(X, Y) \\
& =\frac{2 Y+X}{2 X} \tag{1}
\end{align*}
$$

An ode of the form $Y^{\prime}=\frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order $n$ if

$$
f\left(t^{n} X, t^{n} Y\right)=t^{n} f(X, Y)
$$

In this case, it can be seen that both $M=2 Y+X$ and $N=2 X$ are both homogeneous and of the same order $n=1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u=\frac{Y}{X}$, or $Y=u X$. Hence

$$
\frac{\mathrm{d} Y}{\mathrm{~d} X}=\frac{\mathrm{d} u}{\mathrm{~d} X} X+u
$$

Applying the transformation $Y=u X$ to the above ODE in (1) gives

$$
\begin{aligned}
\frac{\mathrm{d} u}{\mathrm{~d} X} X+u & =u+\frac{1}{2} \\
\frac{\mathrm{~d} u}{\mathrm{~d} X} & =\frac{1}{2 X}
\end{aligned}
$$

Or

$$
\frac{d}{d X} u(X)-\frac{1}{2 X}=0
$$

Or

$$
2\left(\frac{d}{d X} u(X)\right) X-1=0
$$

Which is now solved as separable in $u(X)$. Which is now solved in $u(X)$. Integrating both sides gives

$$
\begin{aligned}
u(X) & =\int \frac{1}{2 X} \mathrm{~d} X \\
& =\frac{\ln (X)}{2}+c_{2}
\end{aligned}
$$

Now $u$ in the above solution is replaced back by $Y$ using $u=\frac{Y}{X}$ which results in the solution

$$
Y(X)=X\left(\frac{\ln (X)}{2}+c_{2}\right)
$$

Using the solution for $Y(X)$

$$
Y(X)=X\left(\frac{\ln (X)}{2}+c_{2}\right)
$$

And replacing back terms in the above solution using

$$
\begin{aligned}
Y & =y+y_{0} \\
X & =x+x_{0}
\end{aligned}
$$

Or

$$
\begin{aligned}
& Y=y-\frac{5}{4} \\
& X=x+\frac{3}{2}
\end{aligned}
$$

Then the solution in $y$ becomes

$$
y+\frac{5}{4}=\left(x-\frac{3}{2}\right)\left(\frac{\ln \left(x-\frac{3}{2}\right)}{2}+c_{2}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y+\frac{5}{4}=\left(x-\frac{3}{2}\right)\left(\frac{\ln \left(x-\frac{3}{2}\right)}{2}+c_{2}\right) \tag{1}
\end{equation*}
$$



Figure 78: Slope field plot

Verification of solutions

$$
y+\frac{5}{4}=\left(x-\frac{3}{2}\right)\left(\frac{\ln \left(x-\frac{3}{2}\right)}{2}+c_{2}\right)
$$

Verified OK.

### 1.33.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{x+2 y+1}{2 x-3} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 59: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{2}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=2 x-3 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{2 x-3} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{2 x-3}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{x+2 y+1}{2 x-3}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{2 y}{(2 x-3)^{2}} \\
S_{y} & =\frac{1}{2 x-3}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{x+1}{(2 x-3)^{2}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{R+1}{(2 R-3)^{2}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{\ln (2 R-3)}{4}-\frac{5}{4(2 R-3)}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{y}{2 x-3}=\frac{\ln (2 x-3)}{4}-\frac{5}{4(2 x-3)}+c_{1}
$$

Which simplifies to

$$
\frac{y}{2 x-3}=\frac{\ln (2 x-3)}{4}-\frac{5}{4(2 x-3)}+c_{1}
$$

Which gives

$$
y=\frac{\ln (2 x-3) x}{2}+2 c_{1} x-\frac{3 \ln (2 x-3)}{4}-3 c_{1}-\frac{5}{4}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{x+2 y+1}{2 x-3}$ |  | $\frac{d S}{d R}=\frac{R+1}{(2 R-3)^{2}}$ |
|  |  | $\rightarrow \rightarrow \rightarrow$ ( |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\rightarrow \text { ¢ }]{ }+$ |
|  |  | $\rightarrow \rightarrow$ |
|  |  |  |
|  |  |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-\infty]{ }$ | $R=x$ | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{ }+\wedge^{+}$ |
|  | $S=\quad y$ | $\xrightarrow{\rightarrow \rightarrow-4 \rightarrow \rightarrow-2 \rightarrow \rightarrow-0 \rightarrow 74}$ |
|  | $S=\frac{}{2 x-3}$ |  |
|  |  |  |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow 0]{ }{ }^{\text {¢ }}$ ¢ |
|  |  |  |
|  |  |  |

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\ln (2 x-3) x}{2}+2 c_{1} x-\frac{3 \ln (2 x-3)}{4}-3 c_{1}-\frac{5}{4} \tag{1}
\end{equation*}
$$



Figure 79: Slope field plot

## Verification of solutions

$$
y=\frac{\ln (2 x-3) x}{2}+2 c_{1} x-\frac{3 \ln (2 x-3)}{4}-3 c_{1}-\frac{5}{4}
$$

Verified OK.

### 1.33.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(-2 x+3) \mathrm{d} y & =(-x-2 y-1) \mathrm{d} x \\
(x+2 y+1) \mathrm{d} x+(-2 x+3) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =x+2 y+1 \\
N(x, y) & =-2 x+3
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(x+2 y+1) \\
& =2
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(-2 x+3) \\
& =-2
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =\frac{1}{-2 x+3}((2)-(-2)) \\
& =-\frac{4}{2 x-3}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-\frac{4}{2 x-3} \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-2 \ln (2 x-3)} \\
& =\frac{1}{(2 x-3)^{2}}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{(2 x-3)^{2}}(x+2 y+1) \\
& =\frac{x+2 y+1}{(2 x-3)^{2}}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{(2 x-3)^{2}}(-2 x+3) \\
& =-\frac{1}{2 x-3}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x} & =0 \\
\left(\frac{x+2 y+1}{(2 x-3)^{2}}\right)+\left(-\frac{1}{2 x-3}\right) \frac{\mathrm{d} y}{\mathrm{~d} x} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{x+2 y+1}{(2 x-3)^{2}} \mathrm{~d} x \\
\phi & =\frac{-5-4 y+(2 x-3) \ln (2 x-3)}{8 x-12}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=-\frac{4}{8 x-12}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=-\frac{1}{2 x-3}$. Therefore equation (4) becomes

$$
\begin{equation*}
-\frac{1}{2 x-3}=-\frac{1}{2 x-3}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{-5-4 y+(2 x-3) \ln (2 x-3)}{8 x-12}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{-5-4 y+(2 x-3) \ln (2 x-3)}{8 x-12}
$$

The solution becomes

$$
y=\frac{\ln (2 x-3) x}{2}-2 c_{1} x-\frac{3 \ln (2 x-3)}{4}+3 c_{1}-\frac{5}{4}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\ln (2 x-3) x}{2}-2 c_{1} x-\frac{3 \ln (2 x-3)}{4}+3 c_{1}-\frac{5}{4} \tag{1}
\end{equation*}
$$



Figure 80: Slope field plot

## Verification of solutions

$$
y=\frac{\ln (2 x-3) x}{2}-2 c_{1} x-\frac{3 \ln (2 x-3)}{4}+3 c_{1}-\frac{5}{4}
$$

Verified OK.

### 1.33.5 Maple step by step solution

Let's solve
$2 y-(2 x-3) y^{\prime}=-x-1$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=\frac{2 y}{2 x-3}+\frac{x+1}{2 x-3}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE
$y^{\prime}-\frac{2 y}{2 x-3}=\frac{x+1}{2 x-3}$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}-\frac{2 y}{2 x-3}\right)=\frac{\mu(x)(x+1)}{2 x-3}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$
$\mu(x)\left(y^{\prime}-\frac{2 y}{2 x-3}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=-\frac{2 \mu(x)}{2 x-3}$
- Solve to find the integrating factor
$\mu(x)=\frac{1}{2 x-3}$
- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int \frac{\mu(x)(x+1)}{2 x-3} d x+c_{1}$
- Evaluate the integral on the lhs
$\mu(x) y=\int \frac{\mu(x)(x+1)}{2 x-3} d x+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \frac{\mu(x)(x+1)}{2 x-3} d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\frac{1}{2 x-3}$

$$
y=(2 x-3)\left(\int \frac{x+1}{(2 x-3)^{2}} d x+c_{1}\right)
$$

- Evaluate the integrals on the rhs

$$
y=(2 x-3)\left(\frac{\ln (2 x-3)}{4}-\frac{5}{4(2 x-3)}+c_{1}\right)
$$

- $\quad$ Simplify

$$
y=-\frac{5}{4}+\frac{(2 x-3) \ln (2 x-3)}{4}+c_{1}(2 x-3)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 26

```
dsolve((x+2*y(x)+1)-(2*x-3)*diff(y(x),x)=0,y(x), singsol=all)
```

$$
y(x)=-\frac{5}{4}+\frac{(2 x-3) \ln (2 x-3)}{4}+(2 x-3) c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.064 (sec). Leaf size: 32
DSolve $[(x+2 * y[x]+1)-(2 * x-3) * y '[x]==0, y[x], x$, IncludeSingularSolutions $->$ True $]$

$$
y(x) \rightarrow \frac{1}{4}\left((2 x-3) \log (3-2 x)+4 c_{1}(2 x-3)-5\right)
$$

### 1.34 problem 52

1.34.1 Solving as first order ode lie symmetry calculated ode . . . . . . 412
1.34.2 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 417

Internal problem ID [12451]
Internal file name [OUTPUT/11103_Monday_October_16_2023_09_48_13_PM_84134013/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 52.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type
[[_homogeneous, `class A`], _dAlembert]

$$
\frac{y-y^{\prime} x}{\sqrt{x^{2}+y^{2}}}=m
$$

### 1.34.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{-m \sqrt{x^{2}+y^{2}}+y}{x} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{gather*}
\xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
\eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{gather*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
& b_{2}+\frac{\left(-m \sqrt{x^{2}+y^{2}}+y\right)\left(b_{3}-a_{2}\right)}{x}-\frac{\left(-m \sqrt{x^{2}+y^{2}}+y\right)^{2} a_{3}}{x^{2}} \\
& -\left(-\frac{m}{\sqrt{x^{2}+y^{2}}}-\frac{-m \sqrt{x^{2}+y^{2}}+y}{x^{2}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\frac{\left(-\frac{m y}{\sqrt{x^{2}+y^{2}}}+1\right)\left(x b_{2}+y b_{3}+b_{1}\right)}{x}=0
\end{align*}
$$

Putting the above in normal form gives
$-\frac{\left(x^{2}+y^{2}\right)^{\frac{3}{2}} m^{2} a_{3}-m x^{3} a_{2}+m x^{3} b_{3}-2 m x^{2} y a_{3}-m x^{2} y b_{2}-m y^{3} a_{3}-m x y b_{1}+m y^{2} a_{1}+\sqrt{x^{2}+y^{2}} x b_{1}}{\sqrt{x^{2}+y^{2}} x^{2}}$
$=0 \quad$ $=0$

Setting the numerator to zero gives

$$
\begin{align*}
& -\left(x^{2}+y^{2}\right)^{\frac{3}{2}} m^{2} a_{3}+m x^{3} a_{2}-m x^{3} b_{3}+2 m x^{2} y a_{3}+m x^{2} y b_{2}  \tag{6E}\\
& \quad+m y^{3} a_{3}+m x y b_{1}-m y^{2} a_{1}-\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}=0
\end{align*}
$$

Simplifying the above gives

$$
\begin{align*}
& -\left(x^{2}+y^{2}\right)^{\frac{3}{2}} m^{2} a_{3}-\left(x^{2}+y^{2}\right) m x b_{3}+\left(x^{2}+y^{2}\right) m y a_{3}  \tag{6E}\\
& +m x^{3} a_{2}+m x^{2} y a_{3}+m x^{2} y b_{2}+m x y^{2} b_{3}-\left(x^{2}+y^{2}\right) m a_{1} \\
& +m x^{2} a_{1}+m x y b_{1}-\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}=0
\end{align*}
$$

Since the PDE has radicals, simplifying gives

$$
\begin{aligned}
& -m^{2} x^{2} \sqrt{x^{2}+y^{2}} a_{3}-m^{2} \sqrt{x^{2}+y^{2}} y^{2} a_{3}+m x^{3} a_{2}-m x^{3} b_{3}+2 m x^{2} y a_{3} \\
& \quad+m x^{2} y b_{2}+m y^{3} a_{3}+m x y b_{1}-m y^{2} a_{1}-\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}=0
\end{aligned}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\left\{x, y, \sqrt{x^{2}+y^{2}}\right\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}, \sqrt{x^{2}+y^{2}}=v_{3}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& -m^{2} v_{1}^{2} v_{3} a_{3}-m^{2} v_{3} v_{2}^{2} a_{3}+m v_{1}^{3} a_{2}+2 m v_{1}^{2} v_{2} a_{3}+m v_{2}^{3} a_{3}  \tag{7E}\\
& \quad+m v_{1}^{2} v_{2} b_{2}-m v_{1}^{3} b_{3}-m v_{2}^{2} a_{1}+m v_{1} v_{2} b_{1}+v_{3} v_{2} a_{1}-v_{3} v_{1} b_{1}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}, v_{3}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(m a_{2}-m b_{3}\right) v_{1}^{3}+\left(2 m a_{3}+m b_{2}\right) v_{1}^{2} v_{2}-m^{2} v_{1}^{2} v_{3} a_{3}+m v_{1} v_{2} b_{1}  \tag{8E}\\
& -v_{3} v_{1} b_{1}+m v_{2}^{3} a_{3}-m^{2} v_{3} v_{2}^{2} a_{3}-m v_{2}^{2} a_{1}+v_{3} v_{2} a_{1}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
a_{1} & =0 \\
m a_{3} & =0 \\
m b_{1} & =0 \\
-b_{1} & =0 \\
-m a_{1} & =0 \\
-m^{2} a_{3} & =0 \\
m a_{2}-m b_{3} & =0 \\
2 m a_{3}+m b_{2} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =b_{3} \\
a_{3} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E, 2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=x \\
& \eta=y
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =y-\left(\frac{-m \sqrt{x^{2}+y^{2}}+y}{x}\right)(x) \\
& =m \sqrt{x^{2}+y^{2}} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{m \sqrt{x^{2}+y^{2}}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{\ln \left(y+\sqrt{x^{2}+y^{2}}\right)}{m}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{-m \sqrt{x^{2}+y^{2}}+y}{x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{x}{m \sqrt{x^{2}+y^{2}}\left(y+\sqrt{x^{2}+y^{2}}\right)} \\
S_{y} & =\frac{1}{m \sqrt{x^{2}+y^{2}}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-\frac{\left(\sqrt{x^{2}+y^{2}} y+x^{2}+y^{2}\right)(m-1)}{\sqrt{x^{2}+y^{2}} m\left(y+\sqrt{x^{2}+y^{2}}\right) x} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{-m+1}{R m}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\ln (R)+\frac{\ln (R)}{m}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{\ln \left(y+\sqrt{x^{2}+y^{2}}\right)}{m}=-\ln (x)+\frac{\ln (x)}{m}+c_{1}
$$

Which simplifies to

$$
\frac{\ln \left(y+\sqrt{x^{2}+y^{2}}\right)}{m}=-\ln (x)+\frac{\ln (x)}{m}+c_{1}
$$

Which gives

$$
y=\frac{x\left(\mathrm{e}^{-2 m\left(\ln (x)-c_{1}\right)}-1\right) \mathrm{e}^{m\left(\ln (x)-c_{1}\right)}}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{x\left(\mathrm{e}^{-2 m\left(\ln (x)-c_{1}\right)}-1\right) \mathrm{e}^{m\left(\ln (x)-c_{1}\right)}}{2} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{x\left(\mathrm{e}^{-2 m\left(\ln (x)-c_{1}\right)}-1\right) \mathrm{e}^{m\left(\ln (x)-c_{1}\right)}}{2}
$$

Verified OK.

### 1.34.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(-\frac{x}{\sqrt{x^{2}+y^{2}}}\right) \mathrm{d} y & =\left(-\frac{y}{\sqrt{x^{2}+y^{2}}}+m\right) \mathrm{d} x \\
\left(\frac{y}{\sqrt{x^{2}+y^{2}}}-m\right) \mathrm{d} x+\left(-\frac{x}{\sqrt{x^{2}+y^{2}}}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=\frac{y}{\sqrt{x^{2}+y^{2}}}-m \\
& N(x, y)=-\frac{x}{\sqrt{x^{2}+y^{2}}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\frac{y}{\sqrt{x^{2}+y^{2}}}-m\right) \\
& =\frac{x^{2}}{\left(x^{2}+y^{2}\right)^{\frac{3}{2}}}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(-\frac{x}{\sqrt{x^{2}+y^{2}}}\right) \\
& =-\frac{y^{2}}{\left(x^{2}+y^{2}\right)^{\frac{3}{2}}}
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =-\frac{\sqrt{x^{2}+y^{2}}}{x}\left(\left(\frac{1}{\sqrt{x^{2}+y^{2}}}-\frac{y^{2}}{\left(x^{2}+y^{2}\right)^{\frac{3}{2}}}\right)-\left(-\frac{1}{\sqrt{x^{2}+y^{2}}}+\frac{x^{2}}{\left(x^{2}+y^{2}\right)^{\frac{3}{2}}}\right)\right) \\
& =-\frac{1}{x}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-\frac{1}{x} \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-\ln (x)} \\
& =\frac{1}{x}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{x}\left(\frac{y}{\sqrt{x^{2}+y^{2}}}-m\right) \\
& =\frac{\frac{y}{\sqrt{x^{2}+y^{2}}}-m}{x}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{x}\left(-\frac{x}{\sqrt{x^{2}+y^{2}}}\right) \\
& =-\frac{1}{\sqrt{x^{2}+y^{2}}}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x} & =0 \\
\left(\frac{\frac{y}{\sqrt{x^{2}+y^{2}}}-m}{x}\right)+\left(-\frac{1}{\sqrt{x^{2}+y^{2}}}\right) \frac{\mathrm{d} y}{\mathrm{~d} x} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives
$\int \frac{\partial \phi}{\partial x} \mathrm{~d} x=\int \bar{M} \mathrm{~d} x$
$\int \frac{\partial \phi}{\partial x} \mathrm{~d} x=\int \frac{\frac{y}{\sqrt{x^{2}+y^{2}}}-m}{x} \mathrm{~d} x$

$$
\phi=-\operatorname{csgn}(y) \ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right)-\operatorname{csgn}(y) \ln (2)-m \ln (x)+f((3 y))
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{align*}
\frac{\partial \phi}{\partial y}= & -\operatorname{csgn}(1, y) \ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right) \\
& \operatorname{csgn}(y)\left(\frac{\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y}{x}+\frac{y\left(\frac{\operatorname{csgn}(y) y}{\sqrt{x^{2}+y^{2}}+\sqrt{x^{2}+y^{2}} \operatorname{csgn}(1, y)+1}\right)}{x}\right) x  \tag{4}\\
- & \frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{} \\
- & \operatorname{csgn}(1, y) \ln (2)+f^{\prime}(y)
\end{align*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=-\frac{1}{\sqrt{x^{2}+y^{2}}}$. Therefore equation (4) becomes

$$
\begin{aligned}
- & \frac{1}{\sqrt{x^{2}+y^{2}}}= \\
& -\frac{\left(\left(\ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right)+\ln (2)+1\right)\left(x^{2}+y^{2}\right) \operatorname{csgn}(y)+\left(\ln (2)+\ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right)\right) y\right.}{\sqrt{x^{2}+y^{2}} y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)} \\
& +f^{\prime}(y)
\end{aligned}
$$

Solving equation (5) for $f^{\prime}(y)$ gives
$f^{\prime}(y)$
$=\xlongequal{\operatorname{csgn}(y) \operatorname{csgn}(1, y) \ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right) x^{2} y+\operatorname{csgn}(y) \operatorname{csgn}(1, y) \ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right) y^{3}+\operatorname{csgn}}$

$$
=\frac{\left(\left(\ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right)+\ln (2)+1\right)\left(x^{2}+y^{2}\right) \operatorname{csgn}(y)+\left(\ln (2)+\ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right)\right)\right.}{\sqrt{x^{2}+y^{2}} y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}
$$

Integrating the above w.r.t $y$ results in

$$
\begin{aligned}
& \int f^{\prime}(y) \mathrm{d} y \\
& =\int\left(\frac{\left(\left(\ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right)+\ln (2)+1\right)\left(x^{2}+y^{2}\right) \operatorname{csgn}(y)+\left(\ln (2)+\ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right)\right)\right.}{\sqrt{x^{2}+y^{2}} y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}\right. \\
& \quad f(y) \\
& \quad=\int_{0}^{y} \frac{\left(\left(\ln \left(\frac{-a\left(\sqrt{-^{2}+x^{2}} \operatorname{csgn}(\ldots a)+\ldots a\right)}{x}\right)+\ln (2)+1\right)\left(\_a^{2}+x^{2}\right) \operatorname{csgn}\left(\_a\right)+(\ln (2)+\ln (-a( \right.}{} \quad+c_{1}
\end{aligned}
$$

Assuming $0<\_a$ then
$f(y)=\int_{0}^{y} \frac{\left(\_a \sqrt{\square^{2}+x^{2}}+\operatorname{csgn}\left(\_a\right)\left(\_a^{2}+x^{2}\right)\right) \operatorname{csgn}\left(1, \_a\right) \_a \ln \left(\frac{-a\left(\sqrt{\left.-^{a^{2}+x^{2}} \operatorname{csgn}\left(\_a\right)+\_a\right)}\right.}{x}\right)+\_a}{\sqrt{\_^{2}+x^{2}} \_a(1)}$

Assuming $0<1$ then

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\begin{aligned}
\phi & =-\operatorname{csgn}(y) \ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right)-\operatorname{csgn}(y) \ln (2)-m \ln (x) \\
& +\int_{0}^{y} \frac{\left(-a \sqrt{a^{2}+x^{2}}+\operatorname{csgn}\left(\_a\right)\left(\_^{2}+x^{2}\right)\right) \operatorname{csgn}\left(1, \_a\right) \_a \ln \left(\frac{-a\left(\sqrt{-a^{2}+x^{2}} \operatorname{csgn}\left(\_a\right)+\_a\right)}{x}\right)+\_a(-}{\sqrt{-^{2}+x^{2}}-a(\sqrt{-}}
\end{aligned}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
\begin{aligned}
c_{1} & =-\operatorname{csgn}(y) \ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right)-\operatorname{csgn}(y) \ln (2)-m \ln (x) \\
& +\int_{0}^{y} \frac{\left(-a \sqrt{a^{2}+x^{2}}+\operatorname{csgn}\left(\_a\right)\left(\_a^{2}+x^{2}\right)\right) \operatorname{csgn}\left(1, \_a\right) \_a \ln \left(\frac{a\left(\sqrt{-a^{2}+x^{2}} \operatorname{csgn}\left(\_a\right)+\_a\right)}{x}\right)+\_a(-}{\sqrt{-^{2}+x^{2}} \_a(\sqrt{-}}
\end{aligned}
$$

Simplifying the solution $-\operatorname{csgn}(y) \ln \left(\frac{y\left(\sqrt{x^{2}+y^{2}} \operatorname{csgn}(y)+y\right)}{x}\right)-\operatorname{csgn}(y) \ln (2)-m \ln (x)+$ $\int_{0}^{y} \frac{\left(\_a \sqrt{\left.-^{a^{2}+x^{2}}+\operatorname{csgn}\left(\_a\right)\left(\_a^{2}+x^{2}\right)\right) \operatorname{csgn}\left(1, \_a\right) \_a \ln \left(\frac{-a\left(\sqrt{-^{2}+x^{2}} \operatorname{csgn}\left(\_a\right)+\_a\right)}{x}\right)+\_a\left(\_a \ln (2) \operatorname{csgn}\left(1, \_a\right)+\operatorname{csgn}\left(\_a\right)\right.}\right.}{\sqrt{-^{a^{2}+x^{2}}} \_a\left(\sqrt{\left.-^{a^{2}+x^{2}} \operatorname{csgn}\left(\_a\right)+\_a\right)}\right.}$
$c_{1}$ to $-\ln \left(\frac{y\left(y+\sqrt{x^{2}+y^{2}}\right)}{x}\right)-\ln (2)-m \ln (x)+\int_{0}^{y} \frac{\left(-a \sqrt{-^{a^{2}+x^{2}}+a^{2}+x^{2}}\right) \_a \ln \left(\frac{-a\left(\sqrt{-^{a^{2}+x^{2}}+\_a}\right)}{x}\right)+\_a \_a \ln (2}{\sqrt{-^{a^{2}+x^{2}}} \_^{a\left(\sqrt{-^{a^{2}+x^{2}}}+\ldots\right.}}$

## Summary

The solution(s) found are the following

$$
\begin{aligned}
& c_{1}-\ln \left(\frac{y\left(y+\sqrt{x^{2}+y^{2}}\right)}{x}\right)-\ln (2)-m \ln (x) \\
& +\int_{0}^{y} \frac{\left(\_a \sqrt{\_^{2}+x^{2}}+\_a^{2}+x^{2}\right) \_a \ln \left(\frac{-a\left(\sqrt{-a^{2}+x^{2}}+\ldots a\right)}{x}\right)+\_a\left(\_a \ln (2)+1\right) \sqrt{a^{2}+x^{2}}+(1+}{\sqrt{a^{2}+x^{2}} \_a\left(\sqrt{a^{2}+x^{2}}+\ldots a\right)} \\
& \quad=c_{1}
\end{aligned}
$$

Verification of solutions

$$
\begin{aligned}
& -\ln \left(\frac{y\left(y+\sqrt{x^{2}+y^{2}}\right)}{x}\right)-\ln (2)-m \ln (x) \\
& +\int_{0}^{y} \frac{\left(\_a \sqrt{\_^{2}+x^{2}}+\_a^{2}+x^{2}\right) \_a \ln \left(\frac{-a\left(\sqrt{\left.-^{a^{2}+x^{2}}+\_a\right)}\right.}{x}\right)+\_a\left(\_a \ln (2)+1\right) \sqrt{\_^{2}+x^{2}}+\left(1+\_a( \right.}{\sqrt{\_^{2}+x^{2}}} \_a\left(\sqrt{a^{2}+x^{2}}+\_a\right) \\
& =c_{1}
\end{aligned}
$$

Verified OK. \{1::positive, _a::positive\}
Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying homogeneous types:
trying homogeneous G
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 32

```
dsolve((y(x)-x*diff(y(x),x))/sqrt(x^2+y(x)^2)=m,y(x), singsol=all)
```

$$
\frac{x^{m} y(x)+x^{m} \sqrt{y(x)^{2}+x^{2}}-c_{1} x}{x}=0
$$

$\checkmark$ Solution by Mathematica
Time used: 0.442 (sec). Leaf size: 36
DSolve[(y[x]-x*y'[x])/Sqrt[x^2+y[x] 2$]==m, y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow \frac{1}{2} e^{-c_{1}} x^{1-m}\left(-x^{2 m}+e^{2 c_{1}}\right)
$$

### 1.35 problem 53

1.35.1 Solving as first order ode lie symmetry calculated ode . . . . . . 425
1.35.2 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 431
1.35.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 434

Internal problem ID [12452]
Internal file name [OUTPUT/11104_Monday_October_16_2023_09_49_11_PM_95249135/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 53.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type
[[_homogeneous, `class A`], _exact, _dAlembert]

$$
\frac{x+y y^{\prime}}{\sqrt{x^{2}+y^{2}}}=m
$$

### 1.35.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{m \sqrt{x^{2}+y^{2}}-x}{y} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & +\frac{\left(m \sqrt{x^{2}+y^{2}}-x\right)\left(b_{3}-a_{2}\right)}{y}-\frac{\left(m \sqrt{x^{2}+y^{2}}-x\right)^{2} a_{3}}{y^{2}} \\
& -\frac{\left(\frac{m x}{\sqrt{x^{2}+y^{2}}}-1\right)\left(x a_{2}+y a_{3}+a_{1}\right)}{y}  \tag{5E}\\
& -\left(\frac{m}{\sqrt{x^{2}+y^{2}}}-\frac{m \sqrt{x^{2}+y^{2}}-x}{y^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
$$

Putting the above in normal form gives

$$
\begin{aligned}
& -\underline{\left(x^{2}+y^{2}\right)^{\frac{3}{2}} m^{2} a_{3}-2 m x^{3} a_{3}-m x^{3} b_{2}+2 m x^{2} y a_{2}-2 m x^{2} y b_{3}-m x y^{2} a_{3}+m y^{3} a_{2}-m y^{3} b_{3}+\sqrt{x^{2}+y^{2}}} \\
& =0
\end{aligned}
$$

Setting the numerator to zero gives

$$
\begin{align*}
& -\left(x^{2}+y^{2}\right)^{\frac{3}{2}} m^{2} a_{3}+2 m x^{3} a_{3}+m x^{3} b_{2}-2 m x^{2} y a_{2}+2 m x^{2} y b_{3} \\
& +m x y^{2} a_{3}-m y^{3} a_{2}+m y^{3} b_{3}-\sqrt{x^{2}+y^{2}} x^{2} a_{3}-\sqrt{x^{2}+y^{2}} x^{2} b_{2}  \tag{6E}\\
& +2 \sqrt{x^{2}+y^{2}} x y a_{2}-2 \sqrt{x^{2}+y^{2}} x y b_{3}+\sqrt{x^{2}+y^{2}} y^{2} a_{3}+b_{2} \sqrt{x^{2}+y^{2}} y^{2} \\
& +m x^{2} b_{1}-m x y a_{1}-\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}=0
\end{align*}
$$

Simplifying the above gives

$$
\begin{align*}
& -\left(x^{2}+y^{2}\right)^{\frac{3}{2}} m^{2} a_{3}+2\left(x^{2}+y^{2}\right) m x a_{3}+\left(x^{2}+y^{2}\right) m x b_{2} \\
& \quad-\left(x^{2}+y^{2}\right) m y a_{2}+2\left(x^{2}+y^{2}\right) m y b_{3}-m x^{2} y a_{2}-m x y^{2} a_{3}-m x y^{2} b_{2}  \tag{6E}\\
& \quad-m y^{3} b_{3}+\left(x^{2}+y^{2}\right) m b_{1}-\sqrt{x^{2}+y^{2}} x^{2} a_{3}-\sqrt{x^{2}+y^{2}} x^{2} b_{2} \\
& +2 \sqrt{x^{2}+y^{2}} x y a_{2}-2 \sqrt{x^{2}+y^{2}} x y b_{3}+\sqrt{x^{2}+y^{2}} y^{2} a_{3}+b_{2} \sqrt{x^{2}+y^{2}} y^{2} \\
& -m x y a_{1}-m y^{2} b_{1}-\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}=0
\end{align*}
$$

Since the PDE has radicals, simplifying gives

$$
\begin{aligned}
& -m^{2} x^{2} \sqrt{x^{2}+y^{2}} a_{3}-m^{2} \sqrt{x^{2}+y^{2}} y^{2} a_{3}+2 m x^{3} a_{3}+m x^{3} b_{2}-2 m x^{2} y a_{2} \\
& +2 m x^{2} y b_{3}+m x y^{2} a_{3}-m y^{3} a_{2}+m y^{3} b_{3}+m x^{2} b_{1}-m x y a_{1} \\
& -\sqrt{x^{2}+y^{2}} x^{2} a_{3}-\sqrt{x^{2}+y^{2}} x^{2} b_{2}+2 \sqrt{x^{2}+y^{2}} x y a_{2}-2 \sqrt{x^{2}+y^{2}} x y b_{3} \\
& +\sqrt{x^{2}+y^{2}} y^{2} a_{3}+b_{2} \sqrt{x^{2}+y^{2}} y^{2}-\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}=0
\end{aligned}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\left\{x, y, \sqrt{x^{2}+y^{2}}\right\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}, \sqrt{x^{2}+y^{2}}=v_{3}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& -m^{2} v_{1}^{2} v_{3} a_{3}-m^{2} v_{3} v_{2}^{2} a_{3}-2 m v_{1}^{2} v_{2} a_{2}-m v_{2}^{3} a_{2}+2 m v_{1}^{3} a_{3}+m v_{1} v_{2}^{2} a_{3}  \tag{7E}\\
& \quad+m v_{1}^{3} b_{2}+2 m v_{1}^{2} v_{2} b_{3}+m v_{2}^{3} b_{3}-m v_{1} v_{2} a_{1}+m v_{1}^{2} b_{1}+2 v_{3} v_{1} v_{2} a_{2} \\
& \quad-v_{3} v_{1}^{2} a_{3}+v_{3} v_{2}^{2} a_{3}-v_{3} v_{1}^{2} b_{2}+b_{2} v_{3} v_{2}^{2}-2 v_{3} v_{1} v_{2} b_{3}+v_{3} v_{2} a_{1}-v_{3} v_{1} b_{1}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}, v_{3}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(2 m a_{3}+m b_{2}\right) v_{1}^{3}+\left(-2 m a_{2}+2 m b_{3}\right) v_{1}^{2} v_{2}+\left(-m^{2} a_{3}-a_{3}-b_{2}\right) v_{1}^{2} v_{3}  \tag{8E}\\
& +m v_{1}^{2} b_{1}+m v_{1} v_{2}^{2} a_{3}+\left(2 a_{2}-2 b_{3}\right) v_{1} v_{2} v_{3}-m v_{1} v_{2} a_{1}-v_{3} v_{1} b_{1} \\
& +\left(-m a_{2}+m b_{3}\right) v_{2}^{3}+\left(-m^{2} a_{3}+a_{3}+b_{2}\right) v_{2}^{2} v_{3}+v_{3} v_{2} a_{1}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
a_{1} & =0 \\
m a_{3} & =0 \\
m b_{1} & =0 \\
-b_{1} & =0 \\
-m a_{1} & =0 \\
2 a_{2}-2 b_{3} & =0 \\
-2 m a_{2}+2 m b_{3} & =0 \\
-m a_{2}+m b_{3} & =0 \\
2 m a_{3}+m b_{2} & =0 \\
-m^{2} a_{3}-a_{3}-b_{2} & =0 \\
-m^{2} a_{3}+a_{3}+b_{2} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =b_{3} \\
a_{3} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=x \\
& \eta=y
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Therefore

$$
\begin{aligned}
\frac{d y}{d x} & =\frac{\eta}{\xi} \\
& =\frac{y}{x} \\
& =\frac{y}{x}
\end{aligned}
$$

This is easily solved to give

$$
y=c_{1} x
$$

Where now the coordinate $R$ is taken as the constant of integration. Hence

$$
R=\frac{y}{x}
$$

And $S$ is found from

$$
\begin{aligned}
d S & =\frac{d x}{\xi} \\
& =\frac{d x}{x}
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
S & =\int \frac{d x}{T} \\
& =\ln (x)
\end{aligned}
$$

Where the constant of integration is set to zero as we just need one solution. Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{m \sqrt{x^{2}+y^{2}}-x}{y}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =-\frac{y}{x^{2}} \\
R_{y} & =\frac{1}{x} \\
S_{x} & =\frac{1}{x} \\
S_{y} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{x y}{\sqrt{x^{2}+y^{2}} m x-x^{2}-y^{2}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-\frac{R}{R^{2}-\sqrt{R^{2}+1} m+1}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\int \frac{R}{\sqrt{R^{2}+1} m-R^{2}-1} d R+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\ln (x)=\int^{\frac{y}{x}} \frac{\_a}{\sqrt{\_^{2}+1} m-\_a^{2}-1} d \_a+c_{1}
$$

Which simplifies to

$$
\ln (x)=\int^{\frac{y}{x}} \frac{\_a}{\sqrt{\_^{2}+1} m-\_a^{2}-1} d \_a+c_{1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\ln (x)=\int^{\frac{y}{x}} \frac{\_a}{\sqrt{\_^{2}+1} m-\_a^{2}-1} d \_a+c_{1} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
\ln (x)=\int^{\frac{y}{x}} \frac{\_a}{\sqrt{a^{2}+1} m-\_a^{2}-1} d \_a+c_{1}
$$

Verified OK.

### 1.35.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{y}{\sqrt{x^{2}+y^{2}}}\right) \mathrm{d} y & =\left(-\frac{x}{\sqrt{x^{2}+y^{2}}}+m\right) \mathrm{d} x \\
\left(\frac{x}{\sqrt{x^{2}+y^{2}}}-m\right) \mathrm{d} x+\left(\frac{y}{\sqrt{x^{2}+y^{2}}}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=\frac{x}{\sqrt{x^{2}+y^{2}}}-m \\
& N(x, y)=\frac{y}{\sqrt{x^{2}+y^{2}}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\frac{x}{\sqrt{x^{2}+y^{2}}}-m\right) \\
& =-\frac{x y}{\left(x^{2}+y^{2}\right)^{\frac{3}{2}}}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{y}{\sqrt{x^{2}+y^{2}}}\right) \\
& =-\frac{x y}{\left(x^{2}+y^{2}\right)^{\frac{3}{2}}}
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{x}{\sqrt{x^{2}+y^{2}}}-m \mathrm{~d} x \\
\phi & =-m x+\sqrt{x^{2}+y^{2}}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\frac{y}{\sqrt{x^{2}+y^{2}}}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{y}{\sqrt{x^{2}+y^{2}}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{y}{\sqrt{x^{2}+y^{2}}}=\frac{y}{\sqrt{x^{2}+y^{2}}}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-m x+\sqrt{x^{2}+y^{2}}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-m x+\sqrt{x^{2}+y^{2}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-m x+\sqrt{x^{2}+y^{2}}=c_{1} \tag{1}
\end{equation*}
$$

## Verification of solutions

$$
-m x+\sqrt{x^{2}+y^{2}}=c_{1}
$$

Verified OK.

### 1.35.3 Maple step by step solution

Let's solve
$\frac{x+y y^{\prime}}{\sqrt{x^{2}+y^{2}}}=m$

- Highest derivative means the order of the ODE is 1


## $y^{\prime}$

Check if ODE is exact

- ODE is exact if the lhs is the total derivative of a $C^{2}$ function $F^{\prime}(x, y)=0$
- Compute derivative of lhs

$$
F^{\prime}(x, y)+\left(\frac{\partial}{\partial y} F(x, y)\right) y^{\prime}=0
$$

- Evaluate derivatives
$-\frac{x y}{\left(x^{2}+y^{2}\right)^{\frac{3}{2}}}=-\frac{x y}{\left(x^{2}+y^{2}\right)^{\frac{3}{2}}}$
- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$
\left[F(x, y)=c_{1}, M(x, y)=F^{\prime}(x, y), N(x, y)=\frac{\partial}{\partial y} F(x, y)\right]
$$

- $\quad$ Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to $x$
$F(x, y)=\int\left(\frac{x}{\sqrt{x^{2}+y^{2}}}-m\right) d x+f_{1}(y)$
- Evaluate integral
$F(x, y)=-m x+\sqrt{x^{2}+y^{2}}+f_{1}(y)$
- $\quad$ Take derivative of $F(x, y)$ with respect to $y$
$N(x, y)=\frac{\partial}{\partial y} F(x, y)$
- Compute derivative
$\frac{y}{\sqrt{x^{2}+y^{2}}}=\frac{y}{\sqrt{x^{2}+y^{2}}}+\frac{d}{d y} f_{1}(y)$
- Isolate for $\frac{d}{d y} f_{1}(y)$
$\frac{d}{d y} f_{1}(y)=0$
- $\quad$ Solve for $f_{1}(y)$
$f_{1}(y)=0$
- $\quad$ Substitute $f_{1}(y)$ into equation for $F(x, y)$
$F(x, y)=-m x+\sqrt{x^{2}+y^{2}}$
- $\quad$ Substitute $F(x, y)$ into the solution of the ODE

$$
-m x+\sqrt{x^{2}+y^{2}}=c_{1}
$$

- $\quad$ Solve for $y$

$$
\left\{y=\sqrt{m^{2} x^{2}+2 c_{1} m x+c_{1}^{2}-x^{2}}, y=-\sqrt{m^{2} x^{2}+2 c_{1} m x+c_{1}^{2}-x^{2}}\right\}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying homogeneous types:
trying homogeneous G
trying an integrating factor from the invariance group
<- integrating factor successful
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.078 (sec). Leaf size: 178

```
dsolve((x+y(x)*diff(y(x),x))/sqrt(x^2+y(x)^2)=m,y(x), singsol=all)
```

$\int_{-b}^{x} \frac{m \sqrt{a^{2}+y(x)^{2}}-\_a}{-m \sqrt{a^{2}+y(x)^{2}} \_a+y(x)^{2}+\ldots a^{2}} d \_a$

$\checkmark$ Solution by Mathematica
Time used: 2.496 (sec). Leaf size: 103
DSolve[( $\mathrm{x}+\mathrm{y}[\mathrm{x}] * \mathrm{y}$ ' $[\mathrm{x}]) /$ Sqrt $\left[\mathrm{x}^{\wedge} 2+\mathrm{y}[\mathrm{x}] \sim 2\right]==\mathrm{m}, \mathrm{y}[\mathrm{x}], \mathrm{x}$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow-\sqrt{\left(m^{2}-1\right) x^{2}-2 e^{c_{1}} m x+e^{2 c_{1}}} \\
& y(x) \rightarrow \sqrt{\left(m^{2}-1\right) x^{2}-2 e^{c_{1}} m x+e^{2 c_{1}}} \\
& y(x) \rightarrow-\sqrt{\left(m^{2}-1\right) x^{2}} \\
& y(x) \rightarrow \sqrt{\left(m^{2}-1\right) x^{2}}
\end{aligned}
$$

### 1.36 problem 55

1.36.1 Solving as first order ode lie symmetry calculated ode . . . . . . 437
1.36.2 Solving as dAlembert ode . . . . . . . . . . . . . . . . . . . . . 444

Internal problem ID [12453]
Internal file name [OUTPUT/11105_Monday_October_16_2023_09_49_21_PM_82205635/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR
PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 55 .
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "dAlembert", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _dAlembert]
```

$$
y+\frac{x}{y^{\prime}}-\sqrt{x^{2}+y^{2}}=0
$$

### 1.36.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{x}{y-\sqrt{x^{2}+y^{2}}} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & -\frac{x\left(b_{3}-a_{2}\right)}{y-\sqrt{x^{2}+y^{2}}}-\frac{x^{2} a_{3}}{\left(y-\sqrt{x^{2}+y^{2}}\right)^{2}} \\
& -\left(-\frac{1}{y-\sqrt{x^{2}+y^{2}}}-\frac{x^{2}}{\left(y-\sqrt{x^{2}+y^{2}}\right)^{2} \sqrt{x^{2}+y^{2}}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\frac{x\left(1-\frac{y}{\sqrt{x^{2}+y^{2}}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)}{\left(y-\sqrt{x^{2}+y^{2}}\right)^{2}}=0
\end{align*}
$$

Putting the above in normal form gives

$$
\begin{aligned}
& \frac{\left(x^{2}+y^{2}\right)^{\frac{3}{2}} b_{2}-x^{2} a_{3} \sqrt{x^{2}+y^{2}}-\sqrt{x^{2}+y^{2}} x^{2} b_{2}+2 \sqrt{x^{2}+y^{2}} x y a_{2}-2 \sqrt{x^{2}+y^{2}} x y b_{3}+\sqrt{x^{2}+y^{2}} y^{2} a_{3}+\sqrt{ }}{\left(\sqrt{x^{2}+}\right.} \\
& =0
\end{aligned}
$$

Setting the numerator to zero gives

$$
\begin{align*}
& \left(x^{2}+y^{2}\right)^{\frac{3}{2}} b_{2}-x^{2} a_{3} \sqrt{x^{2}+y^{2}}-\sqrt{x^{2}+y^{2}} x^{2} b_{2}+2 \sqrt{x^{2}+y^{2}} x y a_{2} \\
& \quad-2 \sqrt{x^{2}+y^{2}} x y b_{3}+\sqrt{x^{2}+y^{2}} y^{2} a_{3}+\sqrt{x^{2}+y^{2}} y^{2} b_{2}  \tag{6E}\\
& -x^{3} a_{2}+x^{3} b_{3}-x^{2} y b_{2}-2 x y^{2} a_{2}+2 x y^{2} b_{3}-y^{3} a_{3}-2 y^{3} b_{2} \\
& -\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}+x y b_{1}-y^{2} a_{1}=0
\end{align*}
$$

Simplifying the above gives

$$
\begin{align*}
& \left(x^{2}+y^{2}\right)^{\frac{3}{2}} b_{2}-2\left(x^{2}+y^{2}\right) x a_{2}+\left(x^{2}+y^{2}\right) x b_{3}-\left(x^{2}+y^{2}\right) y a_{3} \\
& \quad-2\left(x^{2}+y^{2}\right) y b_{2}-x^{2} a_{3} \sqrt{x^{2}+y^{2}}-\sqrt{x^{2}+y^{2}} x^{2} b_{2}+2 \sqrt{x^{2}+y^{2}} x y a_{2}  \tag{6E}\\
& -2 \sqrt{x^{2}+y^{2}} x y b_{3}+\sqrt{x^{2}+y^{2}} y^{2} a_{3}+\sqrt{x^{2}+y^{2}} y^{2} b_{2}+x^{3} a_{2}+x^{2} y a_{3}+x^{2} y b_{2} \\
& +x y^{2} b_{3}-\left(x^{2}+y^{2}\right) a_{1}-\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}+x^{2} a_{1}+x y b_{1}=0
\end{align*}
$$

Since the PDE has radicals, simplifying gives

$$
\begin{aligned}
& -x^{3} a_{2}+x^{3} b_{3}-x^{2} a_{3} \sqrt{x^{2}+y^{2}}-x^{2} y b_{2}+2 \sqrt{x^{2}+y^{2}} x y a_{2}-2 \sqrt{x^{2}+y^{2}} x y b_{3} \\
& \quad-2 x y^{2} a_{2}+2 x y^{2} b_{3}+\sqrt{x^{2}+y^{2}} y^{2} a_{3}+2 \sqrt{x^{2}+y^{2}} y^{2} b_{2}-y^{3} a_{3} \\
& -2 y^{3} b_{2}-\sqrt{x^{2}+y^{2}} x b_{1}+x y b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}-y^{2} a_{1}=0
\end{aligned}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\left\{x, y, \sqrt{x^{2}+y^{2}}\right\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}, \sqrt{x^{2}+y^{2}}=v_{3}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& -v_{1}^{3} a_{2}-2 v_{1} v_{2}^{2} a_{2}+2 v_{3} v_{1} v_{2} a_{2}-v_{1}^{2} a_{3} v_{3}-v_{2}^{3} a_{3}+v_{3} v_{2}^{2} a_{3}-v_{1}^{2} v_{2} b_{2}-2 v_{2}^{3} b_{2}  \tag{7E}\\
& +2 v_{3} v_{2}^{2} b_{2}+v_{1}^{3} b_{3}+2 v_{1} v_{2}^{2} b_{3}-2 v_{3} v_{1} v_{2} b_{3}-v_{2}^{2} a_{1}+v_{3} v_{2} a_{1}+v_{1} v_{2} b_{1}-v_{3} v_{1} b_{1}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}, v_{3}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(b_{3}-a_{2}\right) v_{1}^{3}-v_{1}^{2} v_{2} b_{2}-v_{1}^{2} a_{3} v_{3}+\left(-2 a_{2}+2 b_{3}\right) v_{1} v_{2}^{2}+\left(2 a_{2}-2 b_{3}\right) v_{1} v_{2} v_{3}  \tag{8E}\\
& \quad+v_{1} v_{2} b_{1}-v_{3} v_{1} b_{1}+\left(-a_{3}-2 b_{2}\right) v_{2}^{3}+\left(a_{3}+2 b_{2}\right) v_{2}^{2} v_{3}-v_{2}^{2} a_{1}+v_{3} v_{2} a_{1}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
a_{1} & =0 \\
b_{1} & =0 \\
-a_{1} & =0 \\
-a_{3} & =0 \\
-b_{1} & =0 \\
-b_{2} & =0 \\
-2 a_{2}+2 b_{3} & =0 \\
2 a_{2}-2 b_{3} & =0 \\
-a_{3}-2 b_{2} & =0 \\
a_{3}+2 b_{2} & =0 \\
b_{3}-a_{2} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =b_{3} \\
a_{3} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=x \\
& \eta=y
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =y-\left(-\frac{x}{y-\sqrt{x^{2}+y^{2}}}\right)(x) \\
& =\frac{-x^{2}-y^{2}+\sqrt{x^{2}+y^{2}} y}{\sqrt{x^{2}+y^{2}}-y} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{-x^{2}-y^{2}+\sqrt{x^{2}+y^{2}} y}{\sqrt{x^{2}+y^{2}-y}} d y}
\end{aligned}
$$

Which results in

$$
S=-\ln \left(y+\sqrt{x^{2}+y^{2}}\right)
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{x}{y-\sqrt{x^{2}+y^{2}}}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{x}{\sqrt{x^{2}+y^{2}}\left(y+\sqrt{x^{2}+y^{2}}\right)} \\
S_{y} & =-\frac{1}{\sqrt{x^{2}+y^{2}}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-\frac{2}{x} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-\frac{2}{R}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-2 \ln (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
-\ln \left(y+\sqrt{x^{2}+y^{2}}\right)=-2 \ln (x)+c_{1}
$$

Which simplifies to

$$
-\ln \left(y+\sqrt{x^{2}+y^{2}}\right)=-2 \ln (x)+c_{1}
$$

Which gives

$$
y=-\frac{\mathrm{e}^{-c_{1}}\left(\mathrm{e}^{2 c_{1}}-x^{2}\right)}{2}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=-\frac{x}{y-\sqrt{x^{2}+y^{2}}}$ |  | $\frac{d S}{d R}=-\frac{2}{R}$ |
| + ${ }^{\text {¢ }}$ ¢ |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=x$ |  |
|  |  |  |
| 禺 | $S=-\ln \left(y+\sqrt{x^{2}}\right.$ |  |
|  |  |  |
| - - - - - |  |  |
| ${ }_{\rightarrow \rightarrow \rightarrow \infty} \rightarrow+\infty \rightarrow \infty$ |  |  |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\mathrm{e}^{-c_{1}}\left(\mathrm{e}^{2 c_{1}}-x^{2}\right)}{2} \tag{1}
\end{equation*}
$$



Figure 81: Slope field plot
Verification of solutions

$$
y=-\frac{\mathrm{e}^{-c_{1}}\left(\mathrm{e}^{2 c_{1}}-x^{2}\right)}{2}
$$

Verified OK.

### 1.36.2 Solving as dAlembert ode

Let $p=y^{\prime}$ the ode becomes

$$
y+\frac{x}{p}-\sqrt{x^{2}+y^{2}}=0
$$

Solving for $y$ from the above results in

$$
\begin{equation*}
y=\frac{x\left(p^{2}-1\right)}{2 p} \tag{1A}
\end{equation*}
$$

This has the form

$$
\begin{equation*}
y=x f(p)+g(p) \tag{}
\end{equation*}
$$

Where $f, g$ are functions of $p=y^{\prime}(x)$. The above ode is dAlembert ode which is now solved. Taking derivative of $\left({ }^{*}\right)$ w.r.t. $x$ gives

$$
\begin{align*}
p & =f+\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \\
p-f & =\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \tag{2}
\end{align*}
$$

Comparing the form $y=x f+g$ to (1A) shows that

$$
\begin{aligned}
& f=\frac{p^{2}-1}{2 p} \\
& g=0
\end{aligned}
$$

Hence (2) becomes

$$
\begin{equation*}
p-\frac{p^{2}-1}{2 p}=x\left(1-\frac{p^{2}-1}{2 p^{2}}\right) p^{\prime}(x) \tag{2~A}
\end{equation*}
$$

The singular solution is found by setting $\frac{d p}{d x}=0$ in the above which gives

$$
p-\frac{p^{2}-1}{2 p}=0
$$

Solving for $p$ from the above gives

$$
\begin{aligned}
& p=i \\
& p=-i
\end{aligned}
$$

Substituting these in (1A) gives

$$
\begin{aligned}
& y=-i x \\
& y=i x
\end{aligned}
$$

The general solution is found when $\frac{\mathrm{d} p}{\mathrm{~d} x} \neq 0$. From eq. (2A). This results in

$$
\begin{equation*}
p^{\prime}(x)=\frac{p(x)-\frac{p(x)^{2}-1}{2 p(x)}}{x\left(1-\frac{p(x)^{2}-1}{2 p(x)^{2}}\right)} \tag{3}
\end{equation*}
$$

This ODE is now solved for $p(x)$.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
p^{\prime}(x)+p(x) p(x)=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{1}{x} \\
& q(x)=0
\end{aligned}
$$

Hence the ode is

$$
p^{\prime}(x)-\frac{p(x)}{x}=0
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{1}{x} d x} \\
& =\frac{1}{x}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x} \mu p & =0 \\
\frac{\mathrm{~d}}{\mathrm{~d} x}\left(\frac{p}{x}\right) & =0
\end{aligned}
$$

Integrating gives

$$
\frac{p}{x}=c_{1}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x}$ results in

$$
p(x)=c_{1} x
$$

Substituing the above solution for $p$ in (2A) gives

$$
y=\frac{c_{1}^{2} x^{2}-1}{2 c_{1}}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=-i x  \tag{1}\\
& y=i x  \tag{2}\\
& y=\frac{c_{1}^{2} x^{2}-1}{2 c_{1}} \tag{3}
\end{align*}
$$



Figure 82: Slope field plot

Verification of solutions

$$
y=-i x
$$

Verified OK.

$$
y=i x
$$

Verified OK.

$$
y=\frac{c_{1}^{2} x^{2}-1}{2 c_{1}}
$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying homogeneous types:
trying homogeneous C
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 26

```
dsolve(y(x)+x/diff(y(x),x)=sqrt(x^2+y(x)^2),y(x), singsol=all)
```

$$
\frac{-c_{1} x^{2}+y(x)+\sqrt{y(x)^{2}+x^{2}}}{x^{2}}=0
$$

Solution by Mathematica
Time used: 0.527 (sec). Leaf size: 27
DSolve[y[x]+x/y'[x]==Sqrt[x^2+y[x]^2],y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow \frac{1}{2} e^{-c_{1}}\left(x^{2}-e^{2 c_{1}}\right)
$$

### 1.37 problem 56

1.37.1 Solving as first order ode lie symmetry calculated ode

Internal problem ID [12454]
Internal file name [OUTPUT/11106_Monday_October_16_2023_09_49_25_PM_38914913/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 56.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "first__order_ode__lie_symmetry_calculated"

Maple gives the following as the ode type
[[_homogeneous, `class A`], _rational, _dAlembert]

$$
y y^{\prime}-\sqrt{x^{2}+y^{2}}=-x
$$

### 1.37.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{\sqrt{x^{2}+y^{2}}-x}{y} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{gather*}
\xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
\eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{gather*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & +\frac{\left(\sqrt{x^{2}+y^{2}}-x\right)\left(b_{3}-a_{2}\right)}{y}-\frac{\left(\sqrt{x^{2}+y^{2}}-x\right)^{2} a_{3}}{y^{2}} \\
& -\frac{\left(\frac{x}{\sqrt{x^{2}+y^{2}}}-1\right)\left(x a_{2}+y a_{3}+a_{1}\right)}{y}  \tag{5E}\\
& -\left(\frac{1}{\sqrt{x^{2}+y^{2}}}-\frac{\sqrt{x^{2}+y^{2}}-x}{y^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
$$

Putting the above in normal form gives

$=0$

Setting the numerator to zero gives

$$
\begin{align*}
& -\left(x^{2}+y^{2}\right)^{\frac{3}{2}} a_{3}-\sqrt{x^{2}+y^{2}} x^{2} a_{3}-\sqrt{x^{2}+y^{2}} x^{2} b_{2}+2 \sqrt{x^{2}+y^{2}} x y a_{2} \\
& -2 \sqrt{x^{2}+y^{2}} x y b_{3}+\sqrt{x^{2}+y^{2}} y^{2} a_{3}+b_{2} \sqrt{x^{2}+y^{2}} y^{2}  \tag{6E}\\
& +2 x^{3} a_{3}+x^{3} b_{2}-2 x^{2} y a_{2}+2 x^{2} y b_{3}+x y^{2} a_{3}-y^{3} a_{2}+y^{3} b_{3} \\
& -\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}+x^{2} b_{1}-x y a_{1}=0
\end{align*}
$$

Simplifying the above gives

$$
\begin{align*}
& -\left(x^{2}+y^{2}\right)^{\frac{3}{2}} a_{3}+2\left(x^{2}+y^{2}\right) x a_{3}+\left(x^{2}+y^{2}\right) x b_{2}-\left(x^{2}+y^{2}\right) y a_{2} \\
& +2\left(x^{2}+y^{2}\right) y b_{3}-\sqrt{x^{2}+y^{2}} x^{2} a_{3}-\sqrt{x^{2}+y^{2}} x^{2} b_{2}  \tag{6E}\\
& +2 \sqrt{x^{2}+y^{2}} x y a_{2}-2 \sqrt{x^{2}+y^{2}} x y b_{3}+\sqrt{x^{2}+y^{2}} y^{2} a_{3} \\
& +b_{2} \sqrt{x^{2}+y^{2}} y^{2}-x^{2} y a_{2}-x y^{2} a_{3}-x y^{2} b_{2}-y^{3} b_{3}+\left(x^{2}+y^{2}\right) b_{1} \\
& -\sqrt{x^{2}+y^{2}} x b_{1}+\sqrt{x^{2}+y^{2}} y a_{1}-x y a_{1}-y^{2} b_{1}=0
\end{align*}
$$

Since the PDE has radicals, simplifying gives

$$
\begin{aligned}
& 2 x^{3} a_{3}+x^{3} b_{2}-2 \sqrt{x^{2}+y^{2}} x^{2} a_{3}-\sqrt{x^{2}+y^{2}} x^{2} b_{2}-2 x^{2} y a_{2}+2 x^{2} y b_{3} \\
& +2 \sqrt{x^{2}+y^{2}} x y a_{2}-2 \sqrt{x^{2}+y^{2}} x y b_{3}+x y^{2} a_{3}+b_{2} \sqrt{x^{2}+y^{2}} y^{2} \\
& \quad-y^{3} a_{2}+y^{3} b_{3}+x^{2} b_{1}-\sqrt{x^{2}+y^{2}} x b_{1}-x y a_{1}+\sqrt{x^{2}+y^{2}} y a_{1}=0
\end{aligned}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\left\{x, y, \sqrt{x^{2}+y^{2}}\right\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}, \sqrt{x^{2}+y^{2}}=v_{3}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& -2 v_{1}^{2} v_{2} a_{2}+2 v_{3} v_{1} v_{2} a_{2}-v_{2}^{3} a_{2}+2 v_{1}^{3} a_{3}-2 v_{3} v_{1}^{2} a_{3}+v_{1} v_{2}^{2} a_{3}+v_{1}^{3} b_{2}-v_{3} v_{1}^{2} b_{2}  \tag{7E}\\
& \quad+b_{2} v_{3} v_{2}^{2}+2 v_{1}^{2} v_{2} b_{3}-2 v_{3} v_{1} v_{2} b_{3}+v_{2}^{3} b_{3}-v_{1} v_{2} a_{1}+v_{3} v_{2} a_{1}+v_{1}^{2} b_{1}-v_{3} v_{1} b_{1}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}, v_{3}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(2 a_{3}+b_{2}\right) v_{1}^{3}+\left(-2 a_{2}+2 b_{3}\right) v_{1}^{2} v_{2}+\left(-2 a_{3}-b_{2}\right) v_{1}^{2} v_{3}+v_{1}^{2} b_{1}+v_{1} v_{2}^{2} a_{3}  \tag{8E}\\
& \quad+\left(2 a_{2}-2 b_{3}\right) v_{1} v_{2} v_{3}-v_{1} v_{2} a_{1}-v_{3} v_{1} b_{1}+\left(b_{3}-a_{2}\right) v_{2}^{3}+b_{2} v_{3} v_{2}^{2}+v_{3} v_{2} a_{1}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
a_{1} & =0 \\
a_{3} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
-a_{1} & =0 \\
-b_{1} & =0 \\
-2 a_{2}+2 b_{3} & =0 \\
2 a_{2}-2 b_{3} & =0 \\
-2 a_{3}-b_{2} & =0 \\
2 a_{3}+b_{2} & =0 \\
b_{3}-a_{2} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =b_{3} \\
a_{3} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=x \\
& \eta=y
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =y-\left(\frac{\sqrt{x^{2}+y^{2}}-x}{y}\right)(x) \\
& =\frac{x^{2}-\sqrt{x^{2}+y^{2}} x+y^{2}}{y} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{x^{2}-\sqrt{x^{2}+y^{2}} x+y^{2}}{y}} d y
\end{aligned}
$$

Which results in

$$
S=\ln (y)-\frac{x \ln \left(\frac{2 x^{2}+2 \sqrt{x^{2}} \sqrt{x^{2}+y^{2}}}{y}\right)}{\sqrt{x^{2}}}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{\sqrt{x^{2}+y^{2}}-x}{y}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{\sqrt{x^{2}+y^{2}}+x}{x \sqrt{x^{2}+y^{2}}} \\
S_{y} & =\frac{2 x^{2}+y^{2}+2 \sqrt{x^{2}+y^{2}} x}{y \sqrt{x^{2}+y^{2}}\left(\sqrt{x^{2}+y^{2}}+x\right)}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-\frac{\sqrt{x^{2}+y^{2}} x+x^{2}+y^{2}}{x \sqrt{x^{2}+y^{2}}\left(\sqrt{x^{2}+y^{2}}+x\right)} \tag{2A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-\frac{1}{R}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\ln (R)+c_{1} \tag{4}
\end{equation*}
$$

Which gives

$$
y=\mathrm{e}^{\frac{\ln (2)}{2}+\frac{\ln \left(2 \mathrm{e}^{c_{1}}+2 x\right)}{2}+\frac{c_{1}}{2}}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.


## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{\frac{\ln (2)}{2}+\frac{\ln \left(2 \mathrm{e}^{c_{1}}+2 x\right)}{2}+\frac{c_{1}}{2}} \tag{1}
\end{equation*}
$$



Figure 83: Slope field plot

Verification of solutions

$$
y=\mathrm{e}^{\frac{\ln (2)}{2}+\frac{\ln \left(2 \mathrm{e}^{c_{1}}+2 x\right)}{2}+\frac{c_{1}}{2}}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying homogeneous types:
trying homogeneous C
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.109 (sec). Leaf size: 27
dsolve $\left(y(x) * \operatorname{diff}(y(x), x)=-x+\operatorname{sqrt}\left(x^{\wedge} 2+y(x) \wedge 2\right), y(x), \quad\right.$ singsol=all)

$$
\frac{-c_{1} y(x)^{2}+\sqrt{y(x)^{2}+x^{2}}+x}{y(x)^{2}}=0
$$

$\checkmark$ Solution by Mathematica
Time used: 0.656 (sec). Leaf size: 57
DSolve[y[x]*y'[x]==-x+Sqrt[x^2+y[x]~2],y[x],x,IncludeSingularSolutions $->$ True]

$$
\begin{aligned}
& y(x) \rightarrow-e^{\frac{c_{1}}{2}} \sqrt{2 x+e^{c_{1}}} \\
& y(x) \rightarrow e^{\frac{c_{1}}{2}} \sqrt{2 x+e^{c_{1}}} \\
& y(x) \rightarrow 0
\end{aligned}
$$

### 1.38 problem 57

1.38.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 457
1.38.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 459
1.38.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 463
1.38.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 468

Internal problem ID [12455]
Internal file name [OUTPUT/11107_Monday_October_16_2023_09_49_28_PM_33336505/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 57.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}-\frac{2 y}{x+1}=(x+1)^{3}
$$

### 1.38.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{2}{x+1} \\
& q(x)=(x+1)^{3}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{2 y}{x+1}=(x+1)^{3}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2}{x+1} d x} \\
& =\frac{1}{(x+1)^{2}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left((x+1)^{3}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{y}{(x+1)^{2}}\right) & =\left(\frac{1}{(x+1)^{2}}\right)\left((x+1)^{3}\right) \\
\mathrm{d}\left(\frac{y}{(x+1)^{2}}\right) & =(x+1) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{y}{(x+1)^{2}}=\int x+1 \mathrm{~d} x \\
& \frac{y}{(x+1)^{2}}=\frac{1}{2} x^{2}+x+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{(x+1)^{2}}$ results in

$$
y=(x+1)^{2}\left(\frac{1}{2} x^{2}+x\right)+c_{1}(x+1)^{2}
$$

which simplifies to

$$
y=\frac{(x+1)^{2}\left(x^{2}+2 c_{1}+2 x\right)}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{(x+1)^{2}\left(x^{2}+2 c_{1}+2 x\right)}{2} \tag{1}
\end{equation*}
$$



Figure 84: Slope field plot
Verification of solutions

$$
y=\frac{(x+1)^{2}\left(x^{2}+2 c_{1}+2 x\right)}{2}
$$

Verified OK.

### 1.38.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{x^{4}+4 x^{3}+6 x^{2}+4 x+2 y+1}{x+1} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 63: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=(x+1)^{2} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{(x+1)^{2}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{(x+1)^{2}}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{x^{4}+4 x^{3}+6 x^{2}+4 x+2 y+1}{x+1}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{2 y}{(x+1)^{3}} \\
S_{y} & =\frac{1}{(x+1)^{2}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=x+1 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=R+1
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{1}{2} R^{2}+R+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{y}{(x+1)^{2}}=\frac{1}{2} x^{2}+x+c_{1}
$$

Which simplifies to

$$
\frac{y}{(x+1)^{2}}=\frac{1}{2} x^{2}+x+c_{1}
$$

Which gives

$$
y=\frac{(x+1)^{2}\left(x^{2}+2 c_{1}+2 x\right)}{2}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{x^{4}+4 x^{3}+6 x^{2}+4 x+2 y+1}{x+1}$ |  | $\frac{d S}{d R}=R+1$ |
|  |  |  |
|  |  |  |
| bldetapatalatal |  |  |
|  |  | : |
| ¢ |  |  |
|  |  |  |
|  | $S=\quad y$ |  |
| $\therefore \uparrow+L^{+1}+1+x_{1}$ | $S=\overline{(x+1)^{2}}$ |  |
| $14.2{ }^{4}+$ |  |  |
| 4 4.1 |  | Lf |
| - $\square^{4} \boldsymbol{y}$ |  |  |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{(x+1)^{2}\left(x^{2}+2 c_{1}+2 x\right)}{2} \tag{1}
\end{equation*}
$$



Figure 85: Slope field plot
Verification of solutions

$$
y=\frac{(x+1)^{2}\left(x^{2}+2 c_{1}+2 x\right)}{2}
$$

Verified OK.

### 1.38.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(\frac{2 y}{x+1}+(x+1)^{3}\right) \mathrm{d} x \\
\left(-\frac{2 y}{x+1}-(x+1)^{3}\right) \mathrm{d} x+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =-\frac{2 y}{x+1}-(x+1)^{3} \\
N(x, y) & =1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{2 y}{x+1}-(x+1)^{3}\right) \\
& =-\frac{2}{x+1}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =1\left(\left(-\frac{2}{x+1}\right)-(0)\right) \\
& =-\frac{2}{x+1}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-\frac{2}{x+1} \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-2 \ln (x+1)} \\
& =\frac{1}{(x+1)^{2}}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{(x+1)^{2}}\left(-\frac{2 y}{x+1}-(x+1)^{3}\right) \\
& =\frac{-x^{4}-4 x^{3}-6 x^{2}-4 x-2 y-1}{(x+1)^{3}}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{(x+1)^{2}}(1) \\
& =\frac{1}{(x+1)^{2}}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x} & =0 \\
\left(\frac{-x^{4}-4 x^{3}-6 x^{2}-4 x-2 y-1}{(x+1)^{3}}\right)+\left(\frac{1}{(x+1)^{2}}\right) \frac{\mathrm{d} y}{\mathrm{~d} x} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{-x^{4}-4 x^{3}-6 x^{2}-4 x-2 y-1}{(x+1)^{3}} \mathrm{~d} x \\
\phi & =-\frac{x^{2}}{2}-x+\frac{y}{(x+1)^{2}}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\frac{1}{(x+1)^{2}}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{(x+1)^{2}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{(x+1)^{2}}=\frac{1}{(x+1)^{2}}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\frac{x^{2}}{2}-x+\frac{y}{(x+1)^{2}}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\frac{x^{2}}{2}-x+\frac{y}{(x+1)^{2}}
$$

The solution becomes

$$
y=\frac{(x+1)^{2}\left(x^{2}+2 c_{1}+2 x\right)}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{(x+1)^{2}\left(x^{2}+2 c_{1}+2 x\right)}{2} \tag{1}
\end{equation*}
$$



Figure 86: Slope field plot

## Verification of solutions

$$
y=\frac{(x+1)^{2}\left(x^{2}+2 c_{1}+2 x\right)}{2}
$$

Verified OK.

### 1.38.4 Maple step by step solution

Let's solve
$y^{\prime}-\frac{2 y}{x+1}=(x+1)^{3}$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Isolate the derivative
$y^{\prime}=\frac{2 y}{x+1}+(x+1)^{3}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-\frac{2 y}{x+1}=(x+1)^{3}$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}-\frac{2 y}{x+1}\right)=\mu(x)(x+1)^{3}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$
$\mu(x)\left(y^{\prime}-\frac{2 y}{x+1}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=-\frac{2 \mu(x)}{x+1}$
- Solve to find the integrating factor
$\mu(x)=\frac{1}{(x+1)^{2}}$
- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int \mu(x)(x+1)^{3} d x+c_{1}$
- Evaluate the integral on the lhs
$\mu(x) y=\int \mu(x)(x+1)^{3} d x+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \mu(x)(x+1)^{3} d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\frac{1}{(x+1)^{2}}$

$$
y=(x+1)^{2}\left(\int(x+1) d x+c_{1}\right)
$$

- Evaluate the integrals on the rhs

$$
y=(x+1)^{2}\left(\frac{1}{2} x^{2}+x+c_{1}\right)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 18

```
dsolve(diff(y(x),x)-2*y(x)/(x+1)=(x+1)^3,y(x), singsol=all)
```

$$
y(x)=\left(x+\frac{1}{2} x^{2}+c_{1}\right)(1+x)^{2}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.052 (sec). Leaf size: 22
DSolve[y'[x]-2*y[x]/(x+1)==(x+1) $3, y[x], x$, IncludeSingularSolutions $->$ True]

$$
y(x) \rightarrow(x+1)^{2}\left(\frac{x^{2}}{2}+x+c_{1}\right)
$$

### 1.39 problem 58

$$
\text { 1.39.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . } 470
$$

1.39.2 Solving as first order ode lie symmetry lookup ode ..... 472
1.39.3 Solving as exact ode ..... 475
1.39.4 Maple step by step solution ..... 479

Internal problem ID [12456]
Internal file name [OUTPUT/11108_Monday_October_16_2023_09_49_29_PM_86812676/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 58.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}-\frac{a y}{x}=\frac{x+1}{x}
$$

### 1.39.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
p(x) & =-\frac{a}{x} \\
q(x) & =\frac{x+1}{x}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{a y}{x}=\frac{x+1}{x}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{a}{x} d x} \\
& =\mathrm{e}^{-a \ln (x)}
\end{aligned}
$$

Which simplifies to

$$
\mu=x^{-a}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\frac{x+1}{x}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(x^{-a} y\right) & =\left(x^{-a}\right)\left(\frac{x+1}{x}\right) \\
\mathrm{d}\left(x^{-a} y\right) & =\left(x^{-a-1}(x+1)\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& x^{-a} y=\int x^{-a-1}(x+1) \mathrm{d} x \\
& x^{-a} y=-\frac{x^{-a}(a x+a-1)}{a(a-1)}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=x^{-a}$ results in

$$
y=-\frac{x^{a} x^{-a}(a x+a-1)}{a(a-1)}+c_{1} x^{a}
$$

which simplifies to

$$
y=\frac{1+(-x-1) a}{a(a-1)}+c_{1} x^{a}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{1+(-x-1) a}{a(a-1)}+c_{1} x^{a} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{1+(-x-1) a}{a(a-1)}+c_{1} x^{a}
$$

Verified OK.

### 1.39.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{a y+x+1}{x} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 66: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\mathrm{e}^{a \ln (x)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{a \ln (x)}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{-a \ln (x)} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{a y+x+1}{x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-a y x^{-a-1} \\
S_{y} & =x^{-a}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=x^{-a-1}(x+1) \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=R^{-a-1}(R+1)
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\frac{R^{-a}(a R+a-1)}{a(a-1)}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
x^{-a} y=-\frac{x^{-a}(a x+a-1)}{a(a-1)}+c_{1}
$$

Which simplifies to

$$
x^{-a} y=-\frac{x^{-a}(a x+a-1)}{a(a-1)}+c_{1}
$$

Which gives

$$
y=-\frac{\left(x^{-a} a x-c_{1} a^{2}+x^{-a} a+c_{1} a-x^{-a}\right) x^{a}}{a(a-1)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\left(x^{-a} a x-c_{1} a^{2}+x^{-a} a+c_{1} a-x^{-a}\right) x^{a}}{a(a-1)} \tag{1}
\end{equation*}
$$

$\underline{\text { Verification of solutions }}$

$$
y=-\frac{\left(x^{-a} a x-c_{1} a^{2}+x^{-a} a+c_{1} a-x^{-a}\right) x^{a}}{a(a-1)}
$$

Verified OK.

### 1.39.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(\frac{a y}{x}+\frac{x+1}{x}\right) \mathrm{d} x \\
\left(-\frac{a y}{x}-\frac{x+1}{x}\right) \mathrm{d} x+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=-\frac{a y}{x}-\frac{x+1}{x} \\
& N(x, y)=1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{a y}{x}-\frac{x+1}{x}\right) \\
& =-\frac{a}{x}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =1\left(\left(-\frac{a}{x}\right)-(0)\right) \\
& =-\frac{a}{x}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-\frac{a}{x} \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-a \ln (x)} \\
& =x^{-a}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =x^{-a}\left(-\frac{a y}{x}-\frac{x+1}{x}\right) \\
& =-x^{-a-1}(a y+x+1)
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =x^{-a}(1) \\
& =x^{-a}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x} & =0 \\
\left(-x^{-a-1}(a y+x+1)\right)+\left(x^{-a}\right) \frac{\mathrm{d} y}{\mathrm{~d} x} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-x^{-a-1}(a y+x+1) \mathrm{d} x \\
\phi & =\frac{x^{-a}\left(-1+a^{2} y+(-y+1+x) a\right)}{(a-1) a}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{align*}
\frac{\partial \phi}{\partial y} & =\frac{x^{-a}\left(a^{2}-a\right)}{(a-1) a}+f^{\prime}(y)  \tag{4}\\
& =x^{-a}+f^{\prime}(y)
\end{align*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=x^{-a}$. Therefore equation (4) becomes

$$
\begin{equation*}
x^{-a}=x^{-a}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{x^{-a}\left(-1+a^{2} y+(-y+1+x) a\right)}{(a-1) a}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{x^{-a}\left(-1+a^{2} y+(-y+1+x) a\right)}{(a-1) a}
$$

The solution becomes

$$
y=-\frac{\left(x^{-a} a x-c_{1} a^{2}+x^{-a} a+c_{1} a-x^{-a}\right) x^{a}}{a(a-1)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\left(x^{-a} a x-c_{1} a^{2}+x^{-a} a+c_{1} a-x^{-a}\right) x^{a}}{a(a-1)} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=-\frac{\left(x^{-a} a x-c_{1} a^{2}+x^{-a} a+c_{1} a-x^{-a}\right) x^{a}}{a(a-1)}
$$

Verified OK.

### 1.39.4 Maple step by step solution

Let's solve
$y^{\prime}-\frac{a y}{x}=\frac{x+1}{x}$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=\frac{a y}{x}+\frac{x+1}{x}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-\frac{a y}{x}=\frac{x+1}{x}$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}-\frac{a y}{x}\right)=\frac{\mu(x)(x+1)}{x}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$
$\mu(x)\left(y^{\prime}-\frac{a y}{x}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=-\frac{\mu(x) a}{x}$
- Solve to find the integrating factor
$\mu(x)=\frac{1}{x^{a}}$
- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int \frac{\mu(x)(x+1)}{x} d x+c_{1}$
- Evaluate the integral on the lhs
$\mu(x) y=\int \frac{\mu(x)(x+1)}{x} d x+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \frac{\mu(x)(x+1)}{x} d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\frac{1}{x^{a}}$
$y=x^{a}\left(\int \frac{x+1}{x^{a} x} d x+c_{1}\right)$
- Evaluate the integrals on the rhs
$y=x^{a}\left(-\frac{a x+a-1}{(a-1) a x^{a}}+c_{1}\right)$
- Simplify

$$
y=\frac{c_{1} x^{a}(a-1) a+1+(-x-1) a}{(a-1) a}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 31

```
dsolve(diff(y(x),x)-a*y(x)/x=(x+1)/x,y(x), singsol=all)
```

$$
y(x)=\left(-\frac{x^{-a}(a x+a-1)}{a(a-1)}+c_{1}\right) x^{a}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.136 (sec). Leaf size: 28
DSolve[y' $[\mathrm{x}]-\mathrm{a} * \mathrm{y}[\mathrm{x}] / \mathrm{x}==(\mathrm{x}+1) / \mathrm{x}, \mathrm{y}[\mathrm{x}], \mathrm{x}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow-\frac{a x+a-1}{(a-1) a}+c_{1} x^{a}
$$

### 1.40 problem 59

1.40.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 481
1.40.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 483
1.40.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 486
1.40.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 490

Internal problem ID [12457]
Internal file name [OUTPUT/11109_Monday_October_16_2023_09_49_29_PM_82463922/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 59.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
\left(-x^{2}+x\right) y^{\prime}+\left(2 x^{2}-1\right) y=a x^{3}
$$

### 1.40.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{2 x^{2}-1}{x(x-1)} \\
& q(x)=-\frac{x^{2} a}{x-1}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{\left(2 x^{2}-1\right) y}{x(x-1)}=-\frac{x^{2} a}{x-1}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2 x^{2}-1}{x(x-1)} d x} \\
& =\mathrm{e}^{-2 x-\ln (x-1)-\ln (x)}
\end{aligned}
$$

Which simplifies to

$$
\mu=\frac{\mathrm{e}^{-2 x}}{x(x-1)}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(-\frac{x^{2} a}{x-1}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{\mathrm{e}^{-2 x} y}{x(x-1)}\right) & =\left(\frac{\mathrm{e}^{-2 x}}{x(x-1)}\right)\left(-\frac{x^{2} a}{x-1}\right) \\
\mathrm{d}\left(\frac{\mathrm{e}^{-2 x} y}{x(x-1)}\right) & =\left(-\frac{x a \mathrm{e}^{-2 x}}{(x-1)^{2}}\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{\mathrm{e}^{-2 x} y}{x(x-1)}=\int-\frac{x a \mathrm{e}^{-2 x}}{(x-1)^{2}} \mathrm{~d} x \\
& \frac{\mathrm{e}^{-2 x} y}{x(x-1)}=-a\left(\mathrm{e}^{-2} \exp \text { Integral }_{1}(2 x-2)+\frac{2 \mathrm{e}^{-2 x}}{-2 x+2}\right)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{\mathrm{e}^{-2 x}}{x(x-1)}$ results in

$$
y=-x(x-1) \mathrm{e}^{2 x} a\left(\mathrm{e}^{-2} \exp \operatorname{Integral}_{1}(2 x-2)+\frac{2 \mathrm{e}^{-2 x}}{-2 x+2}\right)+c_{1} x(x-1) \mathrm{e}^{2 x}
$$

which simplifies to

$$
y=-\left(a \mathrm{e}^{2 x-2}(x-1) \exp \operatorname{Integral}_{1}(2 x-2)-c_{1}(x-1) \mathrm{e}^{2 x}-a\right) x
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\left(a \mathrm{e}^{2 x-2}(x-1) \exp \text { Integral }_{1}(2 x-2)-c_{1}(x-1) \mathrm{e}^{2 x}-a\right) x \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=-\left(a \mathrm{e}^{2 x-2}(x-1) \exp \operatorname{Integral}_{1}(2 x-2)-c_{1}(x-1) \mathrm{e}^{2 x}-a\right) x
$$

Verified OK.

### 1.40.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{-a x^{3}+2 x^{2} y-y}{x(x-1)} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 69: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | special | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |  |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |$\frac{\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}}{}$| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 |
| :--- | :--- | :--- |
| $-\int(n-1) f(x) d x y^{n}$ |  |  |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\mathrm{e}^{2 x+\ln (x-1)+\ln (x)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{2 x+\ln (x-1)+\ln (x)}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{\mathrm{e}^{-2 x} y}{x(x-1)}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{-a x^{3}+2 x^{2} y-y}{x(x-1)}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{y\left(-2 x^{2}+1\right) \mathrm{e}^{-2 x}}{x^{2}(x-1)^{2}} \\
S_{y} & =\frac{\mathrm{e}^{-2 x}}{x(x-1)}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-\frac{x a \mathrm{e}^{-2 x}}{(x-1)^{2}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-\frac{R a \mathrm{e}^{-2 R}}{(R-1)^{2}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\mathrm{e}^{-2} \exp \operatorname{Integral}_{1}(2 R-2) a-\frac{2 \mathrm{e}^{-2 R} a}{2-2 R}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{\mathrm{e}^{-2 x} y}{x(x-1)}=-\mathrm{e}^{-2} \operatorname{expIntegral} l_{1}(2 x-2) a-\frac{2 \mathrm{e}^{-2 x} a}{-2 x+2}+c_{1}
$$

Which simplifies to

$$
\frac{\mathrm{e}^{-2 x} y}{x(x-1)}=-\mathrm{e}^{-2} \operatorname{expIntegral} l_{1}(2 x-2) a-\frac{2 \mathrm{e}^{-2 x} a}{-2 x+2}+c_{1}
$$

Which gives
$y=-\left(\mathrm{e}^{-2} \operatorname{expIntegral}{ }_{1}(2 x-2) a x-\mathrm{e}^{-2} \exp\right.$ Integral $\left._{1}(2 x-2) a-a \mathrm{e}^{-2 x}-c_{1} x+c_{1}\right) \mathrm{e}^{2 x} x$
Summary
The solution(s) found are the following

$$
\begin{array}{r}
y=-\left(\mathrm{e}^{-2} \exp \text { Integral }_{1}(2 x-2) a x-\mathrm{e}^{-2} \exp \text { Integral }_{1}(2 x-2) a-a \mathrm{e}^{-2 x}-c_{1} x\right. \\
\left.+c_{1}\right) \mathrm{e}^{2 x} x
\end{array}
$$

## Verification of solutions

$y=-\left(\mathrm{e}^{-2} \operatorname{expIntegral}{ }_{1}(2 x-2) a x-\mathrm{e}^{-2} \exp \operatorname{Integral}_{1}(2 x-2) a-a \mathrm{e}^{-2 x}-c_{1} x+c_{1}\right) \mathrm{e}^{2 x} x$
Verified OK.

### 1.40.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(-x^{2}+x\right) \mathrm{d} y & =\left(-\left(2 x^{2}-1\right) y+a x^{3}\right) \mathrm{d} x \\
\left(\left(2 x^{2}-1\right) y-a x^{3}\right) \mathrm{d} x+\left(-x^{2}+x\right) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =\left(2 x^{2}-1\right) y-a x^{3} \\
N(x, y) & =-x^{2}+x
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\left(2 x^{2}-1\right) y-a x^{3}\right) \\
& =2 x^{2}-1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(-x^{2}+x\right) \\
& =-2 x+1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =\frac{1}{-x^{2}+x}\left(\left(2 x^{2}-1\right)-(-2 x+1)\right) \\
& =\frac{-2 x^{2}-2 x+2}{x(x-1)}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int \frac{-2 x^{2}-2 x+2}{x(x-1)} \mathrm{d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-2 x-2 \ln (x-1)-2 \ln (x)} \\
& =\frac{\mathrm{e}^{-2 x}}{x^{2}(x-1)^{2}}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{\mathrm{e}^{-2 x}}{x^{2}(x-1)^{2}}\left(\left(2 x^{2}-1\right) y-a x^{3}\right) \\
& =-\frac{\mathrm{e}^{-2 x}\left(a x^{3}-2 x^{2} y+y\right)}{x^{2}(x-1)^{2}}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{\mathrm{e}^{-2 x}}{x^{2}(x-1)^{2}}\left(-x^{2}+x\right) \\
& =-\frac{\mathrm{e}^{-2 x}}{x(x-1)}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x}=0 \\
\left(-\frac{\mathrm{e}^{-2 x}\left(a x^{3}-2 x^{2} y+y\right)}{x^{2}(x-1)^{2}}\right)+\left(-\frac{\mathrm{e}^{-2 x}}{x(x-1)}\right) \frac{\mathrm{d} y}{\mathrm{~d} x}=0
\end{array}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{\mathrm{e}^{-2 x}\left(a x^{3}-2 x^{2} y+y\right)}{x^{2}(x-1)^{2}} \mathrm{~d} x \\
\phi & =\frac{-a x \mathrm{e}^{-2}(x-1) \exp \operatorname{Integral}_{1}(2 x-2)+\mathrm{e}^{-2 x}(a x-y)}{(x-1) x}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=-\frac{\mathrm{e}^{-2 x}}{x(x-1)}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=-\frac{\mathrm{e}^{-2 x}}{x(x-1)}$. Therefore equation (4) becomes

$$
\begin{equation*}
-\frac{\mathrm{e}^{-2 x}}{x(x-1)}=-\frac{\mathrm{e}^{-2 x}}{x(x-1)}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{-a x \mathrm{e}^{-2}(x-1) \exp \operatorname{Integral}_{1}(2 x-2)+\mathrm{e}^{-2 x}(a x-y)}{(x-1) x}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
\left.c_{1}=\frac{-a x \mathrm{e}^{-2}(x-1) \operatorname{expIntegral}}{1}(2 x-2)+\mathrm{e}^{-2 x}(a x-y)\right)
$$

The solution becomes
$y=-x\left(\mathrm{e}^{-2} \operatorname{expIntegral}{ }_{1}(2 x-2) a x-\mathrm{e}^{-2} \exp \operatorname{Integral}_{1}(2 x-2) a-a \mathrm{e}^{-2 x}+c_{1} x-c_{1}\right) \mathrm{e}^{2 x}$ Summary
The solution(s) found are the following

$$
\begin{aligned}
y=-x\left(\mathrm{e}^{-2} \exp \operatorname{Integral}_{1}(2 x-2) a x-\mathrm{e}^{-2} \exp \text { Integral }_{1}(2 x-2) a-a \mathrm{e}^{-2 x}\right. & +c_{1} x \\
& \left.-c_{1}\right) \mathrm{e}^{2 x}
\end{aligned}
$$

## Verification of solutions

$y=-x\left(\mathrm{e}^{-2} \exp\right.$ Integral $\left._{1}(2 x-2) a x-\mathrm{e}^{-2} \operatorname{expIntegral}{ }_{1}(2 x-2) a-a \mathrm{e}^{-2 x}+c_{1} x-c_{1}\right) \mathrm{e}^{2 x}$ Verified OK.

### 1.40.4 Maple step by step solution

Let's solve
$\left(-x^{2}+x\right) y^{\prime}+\left(2 x^{2}-1\right) y=a x^{3}$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=\frac{\left(2 x^{2}-1\right) y}{x(x-1)}-\frac{x^{2} a}{x-1}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-\frac{\left(2 x^{2}-1\right) y}{x(x-1)}=-\frac{x^{2} a}{x-1}$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}-\frac{\left(2 x^{2}-1\right) y}{x(x-1)}\right)=-\frac{\mu(x) x^{2} a}{x-1}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$
$\mu(x)\left(y^{\prime}-\frac{\left(2 x^{2}-1\right) y}{x(x-1)}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=-\frac{\mu(x)\left(2 x^{2}-1\right)}{x(x-1)}$
- Solve to find the integrating factor
$\mu(x)=\frac{\mathrm{e}^{-2 x}}{x(x-1)}$
- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int-\frac{\mu(x) x^{2} a}{x-1} d x+c_{1}$
- Evaluate the integral on the lhs
$\mu(x) y=\int-\frac{\mu(x) x^{2} a}{x-1} d x+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int-\frac{\mu(x) x^{2} a}{x-1} d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\frac{\mathrm{e}^{-2 x}}{x(x-1)}$
$y=\frac{x(x-1)\left(\int-\frac{x a e^{-2 x}}{(x-1)^{2}} d x+c_{1}\right)}{\mathrm{e}^{-2 x}}$
- Evaluate the integrals on the rhs

$$
y=\frac{x(x-1)\left(-a\left(\mathrm{e}^{-2} \mathrm{Ei}_{1}(2 x-2)+\frac{2 \mathrm{e}^{-2 x}}{-2 x+2}\right)+c_{1}\right)}{\mathrm{e}^{-2 x}}
$$

- Simplify

$$
y=-\left(a \mathrm{e}^{2 x-2}(x-1) \operatorname{Ei}_{1}(2 x-2)-c_{1}(x-1) \mathrm{e}^{2 x}-a\right) x
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

Solution by Maple
Time used: 0.0 (sec). Leaf size: 39

```
dsolve((x-x^2)*diff (y(x),x)+(2*x^2-1)*y(x)-a*x^3=0,y(x), singsol=all)
```

$$
y(x)=-\left(a \mathrm{e}^{2 x-2}(-1+x) \exp \text { Integral }_{1}(2 x-2)-c_{1}(-1+x) \mathrm{e}^{2 x}-a\right) x
$$

$\checkmark$ Solution by Mathematica
Time used: 0.473 (sec). Leaf size: 39
DSolve[( $\left.x-x^{\wedge} 2\right) * y$ ' $[x]+\left(2 * x^{\wedge} 2-1\right) * y[x]-a * x^{\wedge} 3==0, y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow x\left(a e^{2 x-2}(x-1) \text { ExpIntegralEi }(2-2 x)+a-c_{1} e^{2 x}(x-1)\right)
$$

### 1.41 problem 60

> 1.41.1 Solving as linear ode
1.41.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 494
1.41.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 498
1.41.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 502

Internal problem ID [12458]
Internal file name [OUTPUT/11110_Monday_October_16_2023_09_49_30_PM_76967086/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 60.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
s^{\prime} \cos (t)+s \sin (t)=1
$$

### 1.41.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
s^{\prime}+p(t) s=q(t)
$$

Where here

$$
\begin{aligned}
p(t) & =\tan (t) \\
q(t) & =\sec (t)
\end{aligned}
$$

Hence the ode is

$$
s^{\prime}+\tan (t) s=\sec (t)
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \tan (t) d t} \\
& =\frac{1}{\cos (t)}
\end{aligned}
$$

Which simplifies to

$$
\mu=\sec (t)
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu s) & =(\mu)(\sec (t)) \\
\frac{\mathrm{d}}{\mathrm{~d} t}(\sec (t) s) & =(\sec (t))(\sec (t)) \\
\mathrm{d}(\sec (t) s) & =\sec (t)^{2} \mathrm{~d} t
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \sec (t) s=\int \sec (t)^{2} \mathrm{~d} t \\
& \sec (t) s=\tan (t)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\sec (t)$ results in

$$
s=\cos (t) \tan (t)+c_{1} \cos (t)
$$

which simplifies to

$$
s=c_{1} \cos (t)+\sin (t)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
s=c_{1} \cos (t)+\sin (t) \tag{1}
\end{equation*}
$$



Figure 87: Slope field plot
Verification of solutions

$$
s=c_{1} \cos (t)+\sin (t)
$$

Verified OK.

### 1.41.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& s^{\prime}=-\frac{s \sin (t)-1}{\cos (t)} \\
& s^{\prime}=\omega(t, s)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{s}-\xi_{t}\right)-\omega^{2} \xi_{s}-\omega_{t} \xi-\omega_{s} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 72: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(t, s)=0 \\
& \eta(t, s)=\cos (t) \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, s) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d s}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial s}\right) S(t, s)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\cos (t)} d y
\end{aligned}
$$

Which results in

$$
S=\frac{s}{\cos (t)}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, s) S_{s}}{R_{t}+\omega(t, s) R_{s}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{s}, S_{t}, S_{s}$ are all partial derivatives and $\omega(t, s)$ is the right hand side of the original ode given by

$$
\omega(t, s)=-\frac{s \sin (t)-1}{\cos (t)}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{s} & =0 \\
S_{t} & =\sec (t) \tan (t) s \\
S_{s} & =\sec (t)
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\sec (t)^{2} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, s$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\sec (R)^{2}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\tan (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, s$ coordinates. This results in

$$
\sec (t) s=\tan (t)+c_{1}
$$

Which simplifies to

$$
\sec (t) s=\tan (t)+c_{1}
$$

Which gives

$$
s=\frac{\tan (t)+c_{1}}{\sec (t)}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, s$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d s}{d t}=-\frac{s \sin (t)-1}{\cos (t)}$ |  | $\frac{d S}{d R}=\sec (R)^{2}$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=t$ |  |
|  |  |  |
|  | $S=\sec (t) s$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
s=\frac{\tan (t)+c_{1}}{\sec (t)} \tag{1}
\end{equation*}
$$



Figure 88: Slope field plot

## Verification of solutions

$$
s=\frac{\tan (t)+c_{1}}{\sec (t)}
$$

Verified OK.

### 1.41.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, s) \mathrm{d} t+N(t, s) \mathrm{d} s=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(\cos (t)) \mathrm{d} s & =(-s \sin (t)+1) \mathrm{d} t \\
(s \sin (t)-1) \mathrm{d} t+(\cos (t)) \mathrm{d} s & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(t, s) & =s \sin (t)-1 \\
N(t, s) & =\cos (t)
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial s}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial s} & =\frac{\partial}{\partial s}(s \sin (t)-1) \\
& =\sin (t)
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(\cos (t)) \\
& =-\sin (t)
\end{aligned}
$$

Since $\frac{\partial M}{\partial s} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial s}-\frac{\partial N}{\partial t}\right) \\
& =\sec (t)((\sin (t))-(-\sin (t))) \\
& =2 \tan (t)
\end{aligned}
$$

Since $A$ does not depend on $s$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int 2 \tan (t) \mathrm{d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-2 \ln (\cos (t))} \\
& =\sec (t)^{2}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\sec (t)^{2}(s \sin (t)-1) \\
& =(s \sin (t)-1) \sec (t)^{2}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\sec (t)^{2}(\cos (t)) \\
& =\sec (t)
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} s}{\mathrm{~d} t}=0 \\
\left((s \sin (t)-1) \sec (t)^{2}\right)+(\sec (t)) \frac{\mathrm{d} s}{\mathrm{~d} t}=0
\end{array}
$$

The following equations are now set up to solve for the function $\phi(t, s)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial s}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int(s \sin (t)-1) \sec (t)^{2} \mathrm{~d} t \\
\phi & =\sec (t) s-\tan (t)+f(s) \tag{3}
\end{align*}
$$

Where $f(s)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $s$. Taking derivative of equation (3) w.r.t $s$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial s}=\sec (t)+f^{\prime}(s) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial s}=\sec (t)$. Therefore equation (4) becomes

$$
\begin{equation*}
\sec (t)=\sec (t)+f^{\prime}(s) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(s)$ gives

$$
f^{\prime}(s)=0
$$

Therefore

$$
f(s)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(s)$ into equation (3) gives $\phi$

$$
\phi=\sec (t) s-\tan (t)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\sec (t) s-\tan (t)
$$

The solution becomes

$$
s=\frac{\tan (t)+c_{1}}{\sec (t)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
s=\frac{\tan (t)+c_{1}}{\sec (t)} \tag{1}
\end{equation*}
$$



Figure 89: Slope field plot

## Verification of solutions

$$
s=\frac{\tan (t)+c_{1}}{\sec (t)}
$$

Verified OK.

### 1.41.4 Maple step by step solution

Let's solve
$s^{\prime} \cos (t)+s \sin (t)=1$

- Highest derivative means the order of the ODE is 1
$s^{\prime}$
- Isolate the derivative
$s^{\prime}=-\frac{\sin (t) s}{\cos (t)}+\frac{1}{\cos (t)}$
- Group terms with $s$ on the lhs of the ODE and the rest on the rhs of the ODE $s^{\prime}+\frac{\sin (t) s}{\cos (t)}=\frac{1}{\cos (t)}$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(s^{\prime}+\frac{\sin (t) s}{\cos (t)}\right)=\frac{\mu(t)}{\cos (t)}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) s)$
$\mu(t)\left(s^{\prime}+\frac{\sin (t) s}{\cos (t)}\right)=\mu^{\prime}(t) s+\mu(t) s^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=\frac{\mu(t) \sin (t)}{\cos (t)}$
- Solve to find the integrating factor
$\mu(t)=\frac{1}{\cos (t)}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) s)\right) d t=\int \frac{\mu(t)}{\cos (t)} d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) s=\int \frac{\mu(t)}{\cos (t)} d t+c_{1}$
- $\quad$ Solve for $s$
$s=\frac{\int \frac{\mu(t)}{\cos (t)} d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\frac{1}{\cos (t)}$
$s=\cos (t)\left(\int \frac{1}{\cos (t)^{2}} d t+c_{1}\right)$
- Evaluate the integrals on the rhs
$s=\cos (t)\left(\tan (t)+c_{1}\right)$
- Simplify
$s=c_{1} \cos (t)+\sin (t)$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 11

```
dsolve(diff(s(t),t)*\operatorname{cos}(t)+s(t)*sin(t)=1,s(t), singsol=all)
```

$$
s(t)=\cos (t) c_{1}+\sin (t)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.068 (sec). Leaf size: 13
DSolve[s'[t]*Cos[t]+s[t]*Sin[t]==1,s[t],t,IncludeSingularSolutions $->$ True]

$$
s(t) \rightarrow \sin (t)+c_{1} \cos (t)
$$

### 1.42 problem 61

1.42.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 505
1.42.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 507
1.42.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 511
1.42.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 516

Internal problem ID [12459]
Internal file name [OUTPUT/11111_Monday_October_16_2023_09_49_32_PM_86818481/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 61.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
s^{\prime}+s \cos (t)=\frac{\sin (2 t)}{2}
$$

### 1.42.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
s^{\prime}+p(t) s=q(t)
$$

Where here

$$
\begin{aligned}
& p(t)=\cos (t) \\
& q(t)=\frac{\sin (2 t)}{2}
\end{aligned}
$$

Hence the ode is

$$
s^{\prime}+s \cos (t)=\frac{\sin (2 t)}{2}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \cos (t) d t} \\
& =\mathrm{e}^{\sin (t)}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t}(\mu s) & =(\mu)\left(\frac{\sin (2 t)}{2}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} t}\left(\mathrm{e}^{\sin (t)} s\right) & =\left(\mathrm{e}^{\sin (t)}\right)\left(\frac{\sin (2 t)}{2}\right) \\
\mathrm{d}\left(\mathrm{e}^{\sin (t)} s\right) & =\left(\frac{\sin (2 t) \mathrm{e}^{\sin (t)}}{2}\right) \mathrm{d} t
\end{aligned}
$$

## Integrating gives

$$
\begin{aligned}
& \mathrm{e}^{\sin (t)} s=\int \frac{\sin (2 t) \mathrm{e}^{\sin (t)}}{2} \mathrm{~d} t \\
& \mathrm{e}^{\sin (t)} s=\sin (t) \mathrm{e}^{\sin (t)}-\mathrm{e}^{\sin (t)}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{\sin (t)}$ results in

$$
s=\mathrm{e}^{-\sin (t)}\left(\sin (t) \mathrm{e}^{\sin (t)}-\mathrm{e}^{\sin (t)}\right)+c_{1} \mathrm{e}^{-\sin (t)}
$$

which simplifies to

$$
s=\sin (t)-1+c_{1} \mathrm{e}^{-\sin (t)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
s=\sin (t)-1+c_{1} \mathrm{e}^{-\sin (t)} \tag{1}
\end{equation*}
$$



Figure 90: Slope field plot
Verification of solutions

$$
s=\sin (t)-1+c_{1} \mathrm{e}^{-\sin (t)}
$$

Verified OK.

### 1.42.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& s^{\prime}=-s \cos (t)+\frac{\sin (2 t)}{2} \\
& s^{\prime}=\omega(t, s)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{t}+\omega\left(\eta_{s}-\xi_{t}\right)-\omega^{2} \xi_{s}-\omega_{t} \xi-\omega_{s} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 75: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{2}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
\xi(t, s) & =0 \\
\eta(t, s) & =\mathrm{e}^{-\sin (t)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(t, s) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d t}{\xi}=\frac{d s}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t}+\eta \frac{\partial}{\partial s}\right) S(t, s)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=t
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{-\sin (t)}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{\sin (t)} s
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{t}+\omega(t, s) S_{s}}{R_{t}+\omega(t, s) R_{s}} \tag{2}
\end{equation*}
$$

Where in the above $R_{t}, R_{s}, S_{t}, S_{s}$ are all partial derivatives and $\omega(t, s)$ is the right hand side of the original ode given by

$$
\omega(t, s)=-s \cos (t)+\frac{\sin (2 t)}{2}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{t} & =1 \\
R_{s} & =0 \\
S_{t} & =\cos (t) \mathrm{e}^{\sin (t)} s \\
S_{s} & =\mathrm{e}^{\sin (t)}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{\sin (2 t) \mathrm{e}^{\sin (t)}}{2} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $t, s$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{\sin (2 R) \mathrm{e}^{\sin (R)}}{2}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1}+\mathrm{e}^{\sin (R)}(-1+\sin (R)) \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $t, s$ coordinates. This results in

$$
\mathrm{e}^{\sin (t)} s=\mathrm{e}^{\sin (t)}(-1+\sin (t))+c_{1}
$$

Which simplifies to

$$
\mathrm{e}^{\sin (t)} s=\mathrm{e}^{\sin (t)}(-1+\sin (t))+c_{1}
$$

Which gives

$$
s=\mathrm{e}^{-\sin (t)}\left(\sin (t) \mathrm{e}^{\sin (t)}-\mathrm{e}^{\sin (t)}+c_{1}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $t, s$ coordinates | Canonical <br> coordinates <br> transformation | ODE in canonical coordinates <br> $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d s}{d t}=-s \cos (t)+\frac{\sin (2 t)}{2}$ |  | $\frac{d S}{d R}=\frac{\sin (2 R) e^{\sin (R)}}{2}$ |
| 为 |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
s=\mathrm{e}^{-\sin (t)}\left(\sin (t) \mathrm{e}^{\sin (t)}-\mathrm{e}^{\sin (t)}+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 91: Slope field plot

## Verification of solutions

$$
s=\mathrm{e}^{-\sin (t)}\left(\sin (t) \mathrm{e}^{\sin (t)}-\mathrm{e}^{\sin (t)}+c_{1}\right)
$$

Verified OK.

### 1.42.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(t, s) \mathrm{d} t+N(t, s) \mathrm{d} s=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} s & =\left(-s \cos (t)+\frac{\sin (2 t)}{2}\right) \mathrm{d} t \\
\left(s \cos (t)-\frac{\sin (2 t)}{2}\right) \mathrm{d} t+\mathrm{d} s & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(t, s)=s \cos (t)-\frac{\sin (2 t)}{2} \\
& N(t, s)=1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial s}=\frac{\partial N}{\partial t}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial s} & =\frac{\partial}{\partial s}\left(s \cos (t)-\frac{\sin (2 t)}{2}\right) \\
& =\cos (t)
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial t} & =\frac{\partial}{\partial t}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial s} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial s}-\frac{\partial N}{\partial t}\right) \\
& =1((\cos (t))-(0)) \\
& =\cos (t)
\end{aligned}
$$

Since $A$ does not depend on $s$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} t} \\
& =e^{\int \cos (t) \mathrm{d} t}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{\sin (t)} \\
& =\mathrm{e}^{\sin (t)}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\mathrm{e}^{\sin (t)}\left(s \cos (t)-\frac{\sin (2 t)}{2}\right) \\
& =\cos (t)(s-\sin (t)) \mathrm{e}^{\sin (t)}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\mathrm{e}^{\sin (t)}(1) \\
& =\mathrm{e}^{\sin (t)}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} s}{\mathrm{~d} t}=0 \\
\left(\cos (t)(s-\sin (t)) \mathrm{e}^{\sin (t)}\right)+\left(\mathrm{e}^{\sin (t)}\right) \frac{\mathrm{d} s}{\mathrm{~d} t}=0
\end{array}
$$

The following equations are now set up to solve for the function $\phi(t, s)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial t}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial s}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $t$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \bar{M} \mathrm{~d} t \\
\int \frac{\partial \phi}{\partial t} \mathrm{~d} t & =\int \cos (t)(s-\sin (t)) \mathrm{e}^{\sin (t)} \mathrm{d} t \\
\phi & =(s-\sin (t)+1) \mathrm{e}^{\sin (t)}+f(s) \tag{3}
\end{align*}
$$

Where $f(s)$ is used for the constant of integration since $\phi$ is a function of both $t$ and $s$. Taking derivative of equation (3) w.r.t $s$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial s}=\mathrm{e}^{\sin (t)}+f^{\prime}(s) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial s}=\mathrm{e}^{\sin (t)}$. Therefore equation (4) becomes

$$
\begin{equation*}
\mathrm{e}^{\sin (t)}=\mathrm{e}^{\sin (t)}+f^{\prime}(s) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(s)$ gives

$$
f^{\prime}(s)=0
$$

Therefore

$$
f(s)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(s)$ into equation (3) gives $\phi$

$$
\phi=(s-\sin (t)+1) \mathrm{e}^{\sin (t)}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=(s-\sin (t)+1) \mathrm{e}^{\sin (t)}
$$

The solution becomes

$$
s=\mathrm{e}^{-\sin (t)}\left(\sin (t) \mathrm{e}^{\sin (t)}-\mathrm{e}^{\sin (t)}+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
s=\mathrm{e}^{-\sin (t)}\left(\sin (t) \mathrm{e}^{\sin (t)}-\mathrm{e}^{\sin (t)}+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 92: Slope field plot

Verification of solutions

$$
s=\mathrm{e}^{-\sin (t)}\left(\sin (t) \mathrm{e}^{\sin (t)}-\mathrm{e}^{\sin (t)}+c_{1}\right)
$$

## Verified OK.

### 1.42.4 Maple step by step solution

Let's solve
$s^{\prime}+s \cos (t)=\frac{\sin (2 t)}{2}$

- Highest derivative means the order of the ODE is 1
$s^{\prime}$
- Isolate the derivative
$s^{\prime}=-s \cos (t)+\frac{\sin (2 t)}{2}$
- Group terms with $s$ on the lhs of the ODE and the rest on the rhs of the ODE
$s^{\prime}+s \cos (t)=\frac{\sin (2 t)}{2}$
- The ODE is linear; multiply by an integrating factor $\mu(t)$
$\mu(t)\left(s^{\prime}+s \cos (t)\right)=\frac{\mu(t) \sin (2 t)}{2}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d t}(\mu(t) s)$
$\mu(t)\left(s^{\prime}+s \cos (t)\right)=\mu^{\prime}(t) s+\mu(t) s^{\prime}$
- Isolate $\mu^{\prime}(t)$
$\mu^{\prime}(t)=\mu(t) \cos (t)$
- Solve to find the integrating factor
$\mu(t)=\mathrm{e}^{\sin (t)}$
- Integrate both sides with respect to $t$
$\int\left(\frac{d}{d t}(\mu(t) s)\right) d t=\int \frac{\mu(t) \sin (2 t)}{2} d t+c_{1}$
- Evaluate the integral on the lhs
$\mu(t) s=\int \frac{\mu(t) \sin (2 t)}{2} d t+c_{1}$
- $\quad$ Solve for $s$
$s=\frac{\int \frac{\mu(t) \sin (2 t)}{2} d t+c_{1}}{\mu(t)}$
- $\quad$ Substitute $\mu(t)=\mathrm{e}^{\sin (t)}$
$s=\frac{\int \frac{\sin (2 t) \mathrm{e}^{\sin (t)}}{2} d t+c_{1}}{\mathrm{e}^{\sin (t)}}$
- Evaluate the integrals on the rhs
$s=\frac{\sin (t) \mathrm{e}^{\sin (t)}-\mathrm{e}^{\sin (t)}+c_{1}}{\mathrm{e}^{\sin (t)}}$
- Simplify

$$
s=\sin (t)-1+c_{1} \mathrm{e}^{-\sin (t)}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 15

```
dsolve(diff(s(t),t)+s(t)*\operatorname{cos}(t)=1/2*\operatorname{sin}(2*t),s(t), singsol=all)
```

$$
s(t)=\sin (t)-1+\mathrm{e}^{-\sin (t)} c_{1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.087 (sec). Leaf size: 18
DSolve[s'[t]+s[t]*Cos[t]==1/2*Sin[2*t],s[t],t,IncludeSingularSolutions -> True]

$$
s(t) \rightarrow \sin (t)+c_{1} e^{-\sin (t)}-1
$$

### 1.43 problem 62

$$
\text { 1.43.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . } 518
$$

1.43.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 520
1.43.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 523
1.43.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 526

Internal problem ID [12460]
Internal file name [OUTPUT/11112_Monday_October_16_2023_09_49_33_PM_86601182/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 62.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}-\frac{n y}{x}=\mathrm{e}^{x} x^{n}
$$

### 1.43.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{n}{x} \\
& q(x)=\mathrm{e}^{x} x^{n}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{n y}{x}=\mathrm{e}^{x} x^{n}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{n}{x} d x} \\
& =\mathrm{e}^{-n \ln (x)}
\end{aligned}
$$

Which simplifies to

$$
\mu=x^{-n}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\mathrm{e}^{x} x^{n}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(x^{-n} y\right) & =\left(x^{-n}\right)\left(\mathrm{e}^{x} x^{n}\right) \\
\mathrm{d}\left(x^{-n} y\right) & =\mathrm{e}^{x} \mathrm{~d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& x^{-n} y=\int \mathrm{e}^{x} \mathrm{~d} x \\
& x^{-n} y=\mathrm{e}^{x}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=x^{-n}$ results in

$$
y=\mathrm{e}^{x} x^{n}+c_{1} x^{n}
$$

which simplifies to

$$
y=\left(\mathrm{e}^{x}+c_{1}\right) x^{n}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\left(\mathrm{e}^{x}+c_{1}\right) x^{n} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\left(\mathrm{e}^{x}+c_{1}\right) x^{n}
$$

Verified OK.

### 1.43.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{\mathrm{e}^{x} x^{n} x+y n}{x} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 78: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | special | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |  |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $\underline{a}_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}$ |
| $a_{1} b_{2}-a_{2} b_{1}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |  |  |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\mathrm{e}^{n \ln (x)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{n \ln (x)}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{-n \ln (x)} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{\mathrm{e}^{x} x^{n} x+y n}{x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-n y x^{-n-1} \\
S_{y} & =x^{-n}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\mathrm{e}^{x} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\mathrm{e}^{R}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\mathrm{e}^{R}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
y x^{-n}=\mathrm{e}^{x}+c_{1}
$$

Which simplifies to

$$
y x^{-n}=\mathrm{e}^{x}+c_{1}
$$

Which gives

$$
y=\left(\mathrm{e}^{x}+c_{1}\right) x^{n}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\left(\mathrm{e}^{x}+c_{1}\right) x^{n} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\left(\mathrm{e}^{x}+c_{1}\right) x^{n}
$$

Verified OK.

### 1.43.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(\frac{n y}{x}+\mathrm{e}^{x} x^{n}\right) \mathrm{d} x \\
\left(-\frac{n y}{x}-\mathrm{e}^{x} x^{n}\right) \mathrm{d} x+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=-\frac{n y}{x}-\mathrm{e}^{x} x^{n} \\
& N(x, y)=1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{n y}{x}-\mathrm{e}^{x} x^{n}\right) \\
& =-\frac{n}{x}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =1\left(\left(-\frac{n}{x}\right)-(0)\right) \\
& =-\frac{n}{x}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-\frac{n}{x} \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-n \ln (x)} \\
& =x^{-n}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =x^{-n}\left(-\frac{n y}{x}-\mathrm{e}^{x} x^{n}\right) \\
& =-n y x^{-n-1}-\mathrm{e}^{x}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =x^{-n}(1) \\
& =x^{-n}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x}=0 \\
\left(-n y x^{-n-1}-\mathrm{e}^{x}\right)+\left(x^{-n}\right) \frac{\mathrm{d} y}{\mathrm{~d} x}=0
\end{array}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-n y x^{-n-1}-\mathrm{e}^{x} \mathrm{~d} x \\
\phi & =x^{-n} y-\mathrm{e}^{x}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=x^{-n}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=x^{-n}$. Therefore equation (4) becomes

$$
\begin{equation*}
x^{-n}=x^{-n}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=x^{-n} y-\mathrm{e}^{x}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=x^{-n} y-\mathrm{e}^{x}
$$

The solution becomes

$$
y=\left(\mathrm{e}^{x}+c_{1}\right) x^{n}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\left(\mathrm{e}^{x}+c_{1}\right) x^{n} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\left(\mathrm{e}^{x}+c_{1}\right) x^{n}
$$

Verified OK.

### 1.43.4 Maple step by step solution

Let's solve
$y^{\prime}-\frac{n y}{x}=\mathrm{e}^{x} x^{n}$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=\frac{n y}{x}+\mathrm{e}^{x} x^{n}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-\frac{n y}{x}=\mathrm{e}^{x} x^{n}$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}-\frac{n y}{x}\right)=\mu(x) \mathrm{e}^{x} x^{n}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$

$$
\mu(x)\left(y^{\prime}-\frac{n y}{x}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}
$$

- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=-\frac{\mu(x) n}{x}$
- Solve to find the integrating factor

$$
\mu(x)=\frac{1}{x^{n}}
$$

- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int \mu(x) \mathrm{e}^{x} x^{n} d x+c_{1}$
- Evaluate the integral on the lhs

$$
\mu(x) y=\int \mu(x) \mathrm{e}^{x} x^{n} d x+c_{1}
$$

- $\quad$ Solve for $y$
$y=\frac{\int \mu(x) \mathrm{e}^{x} x^{n} d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\frac{1}{x^{n}}$
$y=x^{n}\left(\int \mathrm{e}^{x} d x+c_{1}\right)$
- Evaluate the integrals on the rhs

$$
y=\left(\mathrm{e}^{x}+c_{1}\right) x^{n}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 12

```
dsolve(diff(y(x),x)-n/x*y(x)=exp(x)*x^n,y(x), singsol=all)
```

$$
y(x)=\left(\mathrm{e}^{x}+c_{1}\right) x^{n}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.095 (sec). Leaf size: 15
DSolve[y'[x]-n/x*y[x]==Exp[x]*x^n,y[x],x,IncludeSingularSolutions -> True]

$$
y(x) \rightarrow\left(e^{x}+c_{1}\right) x^{n}
$$

### 1.44 problem 63

1.44.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 529
1.44.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 531
1.44.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 534
1.44.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 537

Internal problem ID [12461]
Internal file name [OUTPUT/11113_Monday_October_16_2023_09_49_33_PM_89991684/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 63.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}+\frac{n y}{x}=x^{-n} a
$$

### 1.44.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=\frac{n}{x} \\
& q(x)=x^{-n} a
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+\frac{n y}{x}=x^{-n} a
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \frac{n}{x} d x} \\
& =\mathrm{e}^{n \ln (x)}
\end{aligned}
$$

Which simplifies to

$$
\mu=x^{n}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(x^{-n} a\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(y x^{n}\right) & =\left(x^{n}\right)\left(x^{-n} a\right) \\
\mathrm{d}\left(y x^{n}\right) & =a \mathrm{~d} x
\end{aligned}
$$

## Integrating gives

$$
\begin{aligned}
& y x^{n}=\int a \mathrm{~d} x \\
& y x^{n}=a x+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=x^{n}$ results in

$$
y=a x x^{-n}+c_{1} x^{-n}
$$

which simplifies to

$$
y=x^{-n}\left(a x+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x^{-n}\left(a x+c_{1}\right) \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=x^{-n}\left(a x+c_{1}\right)
$$

Verified OK.

### 1.44.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{\left(y n x^{n}-a x\right) x^{-n}}{x} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 81: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | special | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |  |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |$\frac{\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}}{}$| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 |
| :--- | :--- | :--- |
| $e^{-\int(n-1) f(x) d x} y^{n}$ |  |  |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\mathrm{e}^{-n \ln (x)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{-n \ln (x)}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{n \ln (x)} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{\left(y n x^{n}-a x\right) x^{-n}}{x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =n y x^{n-1} \\
S_{y} & =x^{n}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=a \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=a
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=a R+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
y x^{n}=a x+c_{1}
$$

Which simplifies to

$$
y x^{n}=a x+c_{1}
$$

Which gives

$$
y=x^{-n}\left(a x+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x^{-n}\left(a x+c_{1}\right) \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=x^{-n}\left(a x+c_{1}\right)
$$

Verified OK.

### 1.44.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(-\frac{n y}{x}+x^{-n} a\right) \mathrm{d} x \\
\left(\frac{n y}{x}-x^{-n} a\right) \mathrm{d} x+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=\frac{n y}{x}-x^{-n} a \\
& N(x, y)=1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\frac{n y}{x}-x^{-n} a\right) \\
& =\frac{n}{x}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =1\left(\left(\frac{n}{x}\right)-(0)\right) \\
& =\frac{n}{x}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int \frac{n}{x} \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{n \ln (x)} \\
& =x^{n}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =x^{n}\left(\frac{n y}{x}-x^{-n} a\right) \\
& =n y x^{n-1}-a
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =x^{n}(1) \\
& =x^{n}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x}=0 \\
\left(n y x^{n-1}-a\right)+\left(x^{n}\right) \frac{\mathrm{d} y}{\mathrm{~d} x}=0
\end{array}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int n y x^{n-1}-a \mathrm{~d} x \\
\phi & =y x^{n}-a x+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=x^{n}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=x^{n}$. Therefore equation (4) becomes

$$
\begin{equation*}
x^{n}=x^{n}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=y x^{n}-a x+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=y x^{n}-a x
$$

The solution becomes

$$
y=x^{-n}\left(a x+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x^{-n}\left(a x+c_{1}\right) \tag{1}
\end{equation*}
$$

## Verification of solutions

$$
y=x^{-n}\left(a x+c_{1}\right)
$$

Verified OK.

### 1.44.4 Maple step by step solution

Let's solve
$y^{\prime}+\frac{n y}{x}=\frac{a}{x^{n}}$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=-\frac{n y}{x}+\frac{a}{x^{n}}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE
$y^{\prime}+\frac{n y}{x}=\frac{a}{x^{n}}$
- $\quad$ The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}+\frac{n y}{x}\right)=\frac{\mu(x) a}{x^{n}}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$
$\mu(x)\left(y^{\prime}+\frac{n y}{x}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=\frac{\mu(x) n}{x}$
- $\quad$ Solve to find the integrating factor

$$
\mu(x)=x^{n}
$$

- Integrate both sides with respect to $x$

$$
\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int \frac{\mu(x) a}{x^{n}} d x+c_{1}
$$

- Evaluate the integral on the lhs
$\mu(x) y=\int \frac{\mu(x) a}{x^{n}} d x+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \frac{\mu(x) a}{x^{n}} d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=x^{n}$
$y=\frac{\int a d x+c_{1}}{x^{n}}$
- Evaluate the integrals on the rhs

$$
y=\frac{a x+c_{1}}{x^{n}}
$$

- Simplify
$y=x^{-n}\left(a x+c_{1}\right)$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 15
dsolve(diff $(y(x), x)+n / x * y(x)=a / x^{\wedge} n, y(x)$, singsol=all)

$$
y(x)=\left(a x+c_{1}\right) x^{-n}
$$

$\sqrt{ }$ Solution by Mathematica
Time used: 0.098 (sec). Leaf size: 17
DSolve[y' $[\mathrm{x}]+\mathrm{n} / \mathrm{x} * \mathrm{y}[\mathrm{x}]==\mathrm{a} / \mathrm{x}^{\wedge} \mathrm{n}, \mathrm{y}[\mathrm{x}], \mathrm{x}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow x^{-n}\left(a x+c_{1}\right)
$$

### 1.45 problem 64

1.45.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 540
1.45.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 542
1.45.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 546
1.45.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 550

Internal problem ID [12462]
Internal file name [OUTPUT/11114_Monday_October_16_2023_09_49_35_PM_74416962/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 64.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[[_linear, `class A`]]

$$
y^{\prime}+y=\mathrm{e}^{-x}
$$

### 1.45.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
p(x) & =1 \\
q(x) & =\mathrm{e}^{-x}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}+y=\mathrm{e}^{-x}
$$

The integrating factor $\mu$ is

$$
\begin{gathered}
\mu=\mathrm{e}^{\int 1 d x} \\
=\mathrm{e}^{x}
\end{gathered}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\mathrm{e}^{-x}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\mathrm{e}^{x} y\right) & =\left(\mathrm{e}^{x}\right)\left(\mathrm{e}^{-x}\right) \\
\mathrm{d}\left(\mathrm{e}^{x} y\right) & =\mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \mathrm{e}^{x} y=\int \mathrm{d} x \\
& \mathrm{e}^{x} y=x+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{x}$ results in

$$
y=x \mathrm{e}^{-x}+c_{1} \mathrm{e}^{-x}
$$

which simplifies to

$$
y=\mathrm{e}^{-x}\left(x+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{-x}\left(x+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 93: Slope field plot

Verification of solutions

$$
y=\mathrm{e}^{-x}\left(x+c_{1}\right)
$$

Verified OK.

### 1.45.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-y+\mathrm{e}^{-x} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 84: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\mathrm{e}^{-x} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{-x}} d y
\end{aligned}
$$

Which results in

$$
S=\mathrm{e}^{x} y
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-y+\mathrm{e}^{-x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\mathrm{e}^{x} y \\
S_{y} & =\mathrm{e}^{x}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=1 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=1
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by
integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=R+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
y \mathrm{e}^{x}=x+c_{1}
$$

Which simplifies to

$$
y \mathrm{e}^{x}=x+c_{1}
$$

Which gives

$$
y=\mathrm{e}^{-x}\left(x+c_{1}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=-y+\mathrm{e}^{-x}$ |  | $\frac{d S}{d R}=1$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=x$ |  |
|  |  |  |
|  | $S=\mathrm{e}^{x} y$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{-x}\left(x+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 94: Slope field plot

Verification of solutions

$$
y=\mathrm{e}^{-x}\left(x+c_{1}\right)
$$

Verified OK.

### 1.45.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(-y+\mathrm{e}^{-x}\right) \mathrm{d} x \\
\left(y-\mathrm{e}^{-x}\right) \mathrm{d} x+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=y-\mathrm{e}^{-x} \\
& N(x, y)=1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(y-\mathrm{e}^{-x}\right) \\
& =1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =1((1)-(0)) \\
& =1
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int 1 \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{x} \\
& =\mathrm{e}^{x}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\mathrm{e}^{x}\left(y-\mathrm{e}^{-x}\right) \\
& =\mathrm{e}^{x} y-1
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\mathrm{e}^{x}(1) \\
& =\mathrm{e}^{x}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x}=0 \\
\left(\mathrm{e}^{x} y-1\right)+\left(\mathrm{e}^{x}\right) \frac{\mathrm{d} y}{\mathrm{~d} x}=0
\end{array}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \mathrm{e}^{x} y-1 \mathrm{~d} x \\
\phi & =-x+\mathrm{e}^{x} y+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\mathrm{e}^{x}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\mathrm{e}^{x}$. Therefore equation (4) becomes

$$
\begin{equation*}
\mathrm{e}^{x}=\mathrm{e}^{x}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-x+\mathrm{e}^{x} y+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-x+\mathrm{e}^{x} y
$$

The solution becomes

$$
y=\mathrm{e}^{-x}\left(x+c_{1}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{-x}\left(x+c_{1}\right) \tag{1}
\end{equation*}
$$



Figure 95: Slope field plot
Verification of solutions

$$
y=\mathrm{e}^{-x}\left(x+c_{1}\right)
$$

Verified OK.

### 1.45.4 Maple step by step solution

Let's solve
$y^{\prime}+y=\mathrm{e}^{-x}$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=-y+\mathrm{e}^{-x}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}+y=\mathrm{e}^{-x}$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}+y\right)=\mu(x) \mathrm{e}^{-x}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$
$\mu(x)\left(y^{\prime}+y\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=\mu(x)$
- Solve to find the integrating factor
$\mu(x)=\mathrm{e}^{x}$
- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int \mu(x) \mathrm{e}^{-x} d x+c_{1}$
- Evaluate the integral on the lhs
$\mu(x) y=\int \mu(x) \mathrm{e}^{-x} d x+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int \mu(x) \mathrm{e}^{-x} d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\mathrm{e}^{x}$
$y=\frac{\int \mathrm{e}^{-x} \mathrm{e}^{x} d x+c_{1}}{\mathrm{e}^{x}}$
- Evaluate the integrals on the rhs
$y=\frac{x+c_{1}}{\mathrm{e}^{x}}$
- Simplify

$$
y=\mathrm{e}^{-x}\left(x+c_{1}\right)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 12
dsolve(diff $(y(x), x)+y(x)=\exp (-x), y(x)$, singsol=all)

$$
y(x)=\left(c_{1}+x\right) \mathrm{e}^{-x}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.079 (sec). Leaf size: 15
DSolve $[y$ ' $[x]+y[x]==\operatorname{Exp}[-x], y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow e^{-x}\left(x+c_{1}\right)
$$

### 1.46 problem 65

1.46.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 553
1.46.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 555
1.46.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 559
1.46.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 564

Internal problem ID [12463]
Internal file name [OUTPUT/11115_Monday_October_16_2023_09_49_36_PM_66917123/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 65 .
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}+\frac{(-2 x+1) y}{x^{2}}=1
$$

### 1.46.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{2 x-1}{x^{2}} \\
& q(x)=1
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{(2 x-1) y}{x^{2}}=1
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2 x-1}{x^{2}} d x} \\
& =\mathrm{e}^{-\frac{1}{x}-2 \ln (x)}
\end{aligned}
$$

Which simplifies to

$$
\mu=\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =\mu \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{\mathrm{e}^{-\frac{1}{x}} y}{x^{2}}\right) & =\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}} \\
\mathrm{~d}\left(\frac{\mathrm{e}^{-\frac{1}{x}} y}{x^{2}}\right) & =\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}} \mathrm{~d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{\mathrm{e}^{-\frac{1}{x}} y}{x^{2}}=\int \frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}} \mathrm{~d} x \\
& \frac{\mathrm{e}^{-\frac{1}{x}} y}{x^{2}}=\mathrm{e}^{-\frac{1}{x}}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}$ results in

$$
y=x^{2} \mathrm{e}^{\frac{1}{x}} \mathrm{e}^{-\frac{1}{x}}+c_{1} x^{2} \mathrm{e}^{\frac{1}{x}}
$$

which simplifies to

$$
y=x^{2}\left(1+c_{1} \mathrm{e}^{\frac{1}{x}}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x^{2}\left(1+c_{1} \mathrm{e}^{\frac{1}{x}}\right) \tag{1}
\end{equation*}
$$



Figure 96: Slope field plot
Verification of solutions

$$
y=x^{2}\left(1+c_{1} \mathrm{e}^{\frac{1}{x}}\right)
$$

Verified OK.

### 1.46.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{x^{2}+2 x y-y}{x^{2}} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 87: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\mathrm{e}^{2 \ln (x)+\frac{1}{x}} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\mathrm{e}^{2 \ln (x)+\frac{1}{x}}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{\mathrm{e}^{-\frac{1}{x}} y}{x^{2}}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{x^{2}+2 x y-y}{x^{2}}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
& R_{x}=1 \\
& R_{y}=0 \\
& S_{x}=\frac{y(-2 x+1) \mathrm{e}^{-\frac{1}{x}}}{x^{4}} \\
& S_{y}=\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{\mathrm{e}^{-\frac{1}{R}}}{R^{2}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\mathrm{e}^{-\frac{1}{R}}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{\mathrm{e}^{-\frac{1}{x}} y}{x^{2}}=\mathrm{e}^{-\frac{1}{x}}+c_{1}
$$

Which simplifies to

$$
\frac{\mathrm{e}^{-\frac{1}{x}} y}{x^{2}}=\mathrm{e}^{-\frac{1}{x}}+c_{1}
$$

Which gives

$$
y=x^{2}\left(\mathrm{e}^{-\frac{1}{x}}+c_{1}\right) \mathrm{e}^{\frac{1}{x}}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{x^{2}+2 x y-y}{x^{2}}$ |  | $\frac{d S}{d R}=\frac{\mathrm{e}^{-\frac{1}{R}}}{R^{2}}$ |
|  |  | $\rightarrow \rightarrow \rightarrow$ ¢ |
|  |  | 他 $\rightarrow \rightarrow \rightarrow \rightarrow \infty$ |
|  |  | 他 $\rightarrow \rightarrow \rightarrow \infty$ |
| $\rightarrow \rightarrow \rightarrow x^{\rightarrow \rightarrow-\infty}$ |  |  |
|  |  | , |
|  | $R=x$ | 01 |
|  | $\mathrm{e}^{-\frac{1}{x}} y$ |  |
|  | $S=\frac{\mathrm{e} x y}{x^{2}}$ | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \infty]{ }$ |
|  | $x^{2}$ | , |
| $\xrightarrow{+1}$ |  | 1 |
| 4t |  | - |
|  |  | $\rightarrow \ggg+14$ |

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x^{2}\left(\mathrm{e}^{-\frac{1}{x}}+c_{1}\right) \mathrm{e}^{\frac{1}{x}} \tag{1}
\end{equation*}
$$



Figure 97: Slope field plot

Verification of solutions

$$
y=x^{2}\left(\mathrm{e}^{-\frac{1}{x}}+c_{1}\right) \mathrm{e}^{\frac{1}{x}}
$$

Verified OK.

### 1.46.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(-\frac{(-2 x+1) y}{x^{2}}+1\right) \mathrm{d} x \\
\left(\frac{(-2 x+1) y}{x^{2}}-1\right) \mathrm{d} x+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=\frac{(-2 x+1) y}{x^{2}}-1 \\
& N(x, y)=1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\frac{(-2 x+1) y}{x^{2}}-1\right) \\
& =\frac{-2 x+1}{x^{2}}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =1\left(\left(\frac{-2 x+1}{x^{2}}\right)-(0)\right) \\
& =\frac{-2 x+1}{x^{2}}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int \frac{-2 x+1}{x^{2}} \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-\frac{1}{x}-2 \ln (x)} \\
& =\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}\left(\frac{(-2 x+1) y}{x^{2}}-1\right) \\
& =-\frac{\mathrm{e}^{-\frac{1}{x}}\left(x^{2}+2 x y-y\right)}{x^{4}}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}(1) \\
& =\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x}=0 \\
\left(-\frac{\mathrm{e}^{-\frac{1}{x}}\left(x^{2}+2 x y-y\right)}{x^{4}}\right)+\left(\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}\right) \frac{\mathrm{d} y}{\mathrm{~d} x}=0
\end{array}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{\mathrm{e}^{-\frac{1}{x}}\left(x^{2}+2 x y-y\right)}{x^{4}} \mathrm{~d} x \\
\phi & =-\frac{\left(x^{2}-y\right) \mathrm{e}^{-\frac{1}{x}}}{x^{2}}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}=\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\frac{\left(x^{2}-y\right) \mathrm{e}^{-\frac{1}{x}}}{x^{2}}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\frac{\left(x^{2}-y\right) \mathrm{e}^{-\frac{1}{x}}}{x^{2}}
$$

The solution becomes

$$
y=x^{2}\left(\mathrm{e}^{-\frac{1}{x}}+c_{1}\right) \mathrm{e}^{\frac{1}{x}}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=x^{2}\left(\mathrm{e}^{-\frac{1}{x}}+c_{1}\right) \mathrm{e}^{\frac{1}{x}} \tag{1}
\end{equation*}
$$



Figure 98: Slope field plot

## Verification of solutions

$$
y=x^{2}\left(\mathrm{e}^{-\frac{1}{x}}+c_{1}\right) \mathrm{e}^{\frac{1}{x}}
$$

Verified OK.

### 1.46.4 Maple step by step solution

Let's solve
$y^{\prime}+\frac{(-2 x+1) y}{x^{2}}=1$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=1+\frac{(2 x-1) y}{x^{2}}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-\frac{(2 x-1) y}{x^{2}}=1$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}-\frac{(2 x-1) y}{x^{2}}\right)=\mu(x)$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$
$\mu(x)\left(y^{\prime}-\frac{(2 x-1) y}{x^{2}}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=-\frac{\mu(x)(2 x-1)}{x^{2}}$
- Solve to find the integrating factor
$\mu(x)=\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}$
- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int \mu(x) d x+c_{1}$
- Evaluate the integral on the lhs

$$
\mu(x) y=\int \mu(x) d x+c_{1}
$$

- $\quad$ Solve for $y$
$y=\frac{\int \mu(x) d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}}$

$$
y=\frac{x^{2}\left(\int \frac{\mathrm{e}^{-\frac{1}{x}}}{x^{2}} d x+c_{1}\right)}{\mathrm{e}^{-\frac{1}{x}}}
$$

- Evaluate the integrals on the rhs

$$
y=\frac{x^{2}\left(\mathrm{e}^{-\frac{1}{x}}+c_{1}\right)}{\mathrm{e}^{-\frac{1}{x}}}
$$

- Simplify

$$
y=x^{2}\left(1+c_{1} \mathrm{e}^{\frac{1}{x}}\right)
$$

## Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 16

```
dsolve(diff(y(x),x)+(1-2*x)/x^2*y(x)-1=0,y(x), singsol=all)
```

$$
y(x)=x^{2}\left(1+\mathrm{e}^{\frac{1}{x}} c_{1}\right)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.057 (sec). Leaf size: 19
DSolve[y'[x]+(1-2*x)/x^2*y[x]-1==0,y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow x^{2}\left(1+c_{1} e^{\frac{1}{x}}\right)
$$

### 1.47 problem 66

1.47.1 Solving as first order ode lie symmetry lookup ode . . . . . . . 566
1.47.2 Solving as bernoulli ode . . . . . . . . . . . . . . . . . . . . . . 570

Internal problem ID [12464]
Internal file name [OUTPUT/11116_Monday_October_16_2023_09_49_36_PM_98065178/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR
PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 66.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "bernoulli", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_Bernoulli]

$$
y^{\prime}+y x-y^{3} x^{3}=0
$$

### 1.47.1 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=x^{3} y^{3}-x y \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type Bernoulli. Therefore we do not need to solve the $\operatorname{PDE}(A)$, and can just use the lookup table shown below to find $\xi, \eta$

Table 90: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{2}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=y^{3} \mathrm{e}^{x^{2}} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{y^{3} \mathrm{e}^{x^{2}}} d y
\end{aligned}
$$

Which results in

$$
S=-\frac{\mathrm{e}^{-x^{2}}}{2 y^{2}}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=x^{3} y^{3}-x y
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{x \mathrm{e}^{-x^{2}}}{y^{2}} \\
S_{y} & =\frac{\mathrm{e}^{-x^{2}}}{y^{3}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=x^{3} \mathrm{e}^{-x^{2}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=R^{3} \mathrm{e}^{-R^{2}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\frac{\left(R^{2}+1\right) \mathrm{e}^{-R^{2}}}{2}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
-\frac{\mathrm{e}^{-x^{2}}}{2 y^{2}}=-\frac{\left(x^{2}+1\right) \mathrm{e}^{-x^{2}}}{2}+c_{1}
$$

Which simplifies to

$$
-\frac{\mathrm{e}^{-x^{2}}}{2 y^{2}}=-\frac{\left(x^{2}+1\right) \mathrm{e}^{-x^{2}}}{2}+c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=x^{3} y^{3}-x y$ |  | $\frac{d S}{d R}=R^{3} \mathrm{e}^{-R^{2}}$ |
|  |  | $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ 辺 |
|  |  |  |
| dre $x^{1} \mathrm{~V}^{-1}+14$ |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow S(R)]{ }$ |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{ }$ |
|  | $R=x$ | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{ }$ |
|  | $\mathrm{e}^{-x^{2}}$ | $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-4}$ |
|  | $S=-\frac{}{2 y^{2}}$ | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\infty]{ }$ |
| $\rightarrow 1$ |  | ${ }_{9}^{2}$ |
| ! |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow+]{ }$ |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow+]{ }$ |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
-\frac{\mathrm{e}^{-x^{2}}}{2 y^{2}}=-\frac{\left(x^{2}+1\right) \mathrm{e}^{-x^{2}}}{2}+c_{1} \tag{1}
\end{equation*}
$$



Figure 99: Slope field plot
Verification of solutions

$$
-\frac{\mathrm{e}^{-x^{2}}}{2 y^{2}}=-\frac{\left(x^{2}+1\right) \mathrm{e}^{-x^{2}}}{2}+c_{1}
$$

Verified OK.

### 1.47.2 Solving as bernoulli ode

In canonical form, the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =x^{3} y^{3}-x y
\end{aligned}
$$

This is a Bernoulli ODE.

$$
\begin{equation*}
y^{\prime}=-x y+x^{3} y^{3} \tag{1}
\end{equation*}
$$

The standard Bernoulli ODE has the form

$$
\begin{equation*}
y^{\prime}=f_{0}(x) y+f_{1}(x) y^{n} \tag{2}
\end{equation*}
$$

The first step is to divide the above equation by $y^{n}$ which gives

$$
\begin{equation*}
\frac{y^{\prime}}{y^{n}}=f_{0}(x) y^{1-n}+f_{1}(x) \tag{3}
\end{equation*}
$$

The next step is use the substitution $w=y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$
\begin{aligned}
f_{0}(x) & =-x \\
f_{1}(x) & =x^{3} \\
n & =3
\end{aligned}
$$

Dividing both sides of ODE (1) by $y^{n}=y^{3}$ gives

$$
\begin{equation*}
y^{\prime} \frac{1}{y^{3}}=-\frac{x}{y^{2}}+x^{3} \tag{4}
\end{equation*}
$$

Let

$$
\begin{align*}
w & =y^{1-n} \\
& =\frac{1}{y^{2}} \tag{5}
\end{align*}
$$

Taking derivative of equation (5) w.r.t $x$ gives

$$
\begin{equation*}
w^{\prime}=-\frac{2}{y^{3}} y^{\prime} \tag{6}
\end{equation*}
$$

Substituting equations (5) and (6) into equation (4) gives

$$
\begin{align*}
-\frac{w^{\prime}(x)}{2} & =-w(x) x+x^{3} \\
w^{\prime} & =-2 x^{3}+2 x w \tag{7}
\end{align*}
$$

The above now is a linear ODE in $w(x)$ which is now solved.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
w^{\prime}(x)+p(x) w(x)=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-2 x \\
& q(x)=-2 x^{3}
\end{aligned}
$$

Hence the ode is

$$
w^{\prime}(x)-2 w(x) x=-2 x^{3}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-2 x d x} \\
& =\mathrm{e}^{-x^{2}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu w) & =(\mu)\left(-2 x^{3}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\mathrm{e}^{-x^{2}} w\right) & =\left(\mathrm{e}^{-x^{2}}\right)\left(-2 x^{3}\right) \\
\mathrm{d}\left(\mathrm{e}^{-x^{2}} w\right) & =\left(-2 x^{3} \mathrm{e}^{-x^{2}}\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \mathrm{e}^{-x^{2}} w=\int-2 x^{3} \mathrm{e}^{-x^{2}} \mathrm{~d} x \\
& \mathrm{e}^{-x^{2}} w=\left(x^{2}+1\right) \mathrm{e}^{-x^{2}}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{-x^{2}}$ results in

$$
w(x)=\mathrm{e}^{x^{2}}\left(x^{2}+1\right) \mathrm{e}^{-x^{2}}+c_{1} \mathrm{e}^{x^{2}}
$$

which simplifies to

$$
w(x)=x^{2}+1+c_{1} \mathrm{e}^{x^{2}}
$$

Replacing $w$ in the above by $\frac{1}{y^{2}}$ using equation (5) gives the final solution.

$$
\frac{1}{y^{2}}=x^{2}+1+c_{1} \mathrm{e}^{x^{2}}
$$

Solving for $y$ gives

$$
\begin{aligned}
& y(x)=\frac{1}{\sqrt{x^{2}+1+c_{1} \mathrm{e}^{x^{2}}}} \\
& y(x)=-\frac{1}{\sqrt{x^{2}+1+c_{1} \mathrm{e}^{x^{2}}}}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\frac{1}{\sqrt{x^{2}+1+c_{1} \mathrm{e}^{x^{2}}}}  \tag{1}\\
& y=-\frac{1}{\sqrt{x^{2}+1+c_{1} \mathrm{e}^{x^{2}}}} \tag{2}
\end{align*}
$$



Figure 100: Slope field plot

## Verification of solutions

$$
y=\frac{1}{\sqrt{x^{2}+1+c_{1} \mathrm{e}^{x^{2}}}}
$$

Verified OK.

$$
y=-\frac{1}{\sqrt{x^{2}+1+c_{1} \mathrm{e}^{x^{2}}}}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 35
dsolve(diff $(y(x), x)+x * y(x)=x^{\wedge} 3 * y(x) \wedge 3, y(x)$, singsol=all)

$$
\begin{aligned}
& y(x)=\frac{1}{\sqrt{\mathrm{e}^{x^{2}} c_{1}+x^{2}+1}} \\
& y(x)=-\frac{1}{\sqrt{\mathrm{e}^{x^{2}} c_{1}+x^{2}+1}}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 7.379 (sec). Leaf size: 50
DSolve $\left[y\right.$ ' $[x]+x * y[x]==x^{\wedge} 3 * y[x] \sim 3, y[x], x$, IncludeSingularSolutions $->$ True]

$$
\begin{aligned}
& y(x) \rightarrow-\frac{1}{\sqrt{x^{2}+c_{1} e^{x^{2}}+1}} \\
& y(x) \rightarrow \frac{1}{\sqrt{x^{2}+c_{1} e^{x^{2}}+1}} \\
& y(x) \rightarrow 0
\end{aligned}
$$

### 1.48 problem 67

1.48.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 575
1.48.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 576
1.48.3 Solving as bernoulli ode . . . . . . . . . . . . . . . . . . . . . . 579
1.48.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 582
1.48.5 Solving as riccati ode . . . . . . . . . . . . . . . . . . . . . . . . 585
1.48.6 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 587

Internal problem ID [12465]
Internal file name [OUTPUT/11117_Monday_October_16_2023_09_49_37_PM_15169171/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 67.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "bernoulli", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
\left(-x^{2}+1\right) y^{\prime}-y x+y^{2} a x=0
$$

### 1.48.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{x y(a y-1)}{x^{2}-1}
\end{aligned}
$$

Where $f(x)=\frac{x}{x^{2}-1}$ and $g(y)=y(a y-1)$. Integrating both sides gives

$$
\frac{1}{y(a y-1)} d y=\frac{x}{x^{2}-1} d x
$$

$$
\begin{aligned}
\int \frac{1}{y(a y-1)} d y & =\int \frac{x}{x^{2}-1} d x \\
\ln (a y-1)-\ln (y) & =\frac{\ln (x-1)}{2}+\frac{\ln (x+1)}{2}+c_{1}
\end{aligned}
$$

Raising both side to exponential gives

$$
\mathrm{e}^{\ln (a y-1)-\ln (y)}=\mathrm{e}^{\frac{\ln (x-1)}{2}+\frac{\ln (x+1)}{2}+c_{1}}
$$

Which simplifies to

$$
\frac{a y-1}{y}=c_{2} \mathrm{e}^{\frac{\ln (x-1)}{2}+\frac{\ln (x+1)}{2}}
$$

Which simplifies to

$$
y=-\frac{1}{c_{2} \sqrt{x-1} \sqrt{x+1}-a}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{1}{c_{2} \sqrt{x-1} \sqrt{x+1}-a} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=-\frac{1}{c_{2} \sqrt{x-1} \sqrt{x+1}-a}
$$

Verified OK.

### 1.48.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{x y(a y-1)}{x^{2}-1} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 92: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=\frac{x^{2}-1}{x} \\
& \eta(x, y)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=y
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d x \\
& =\int \frac{1}{\frac{x^{2}-1}{x}} d x
\end{aligned}
$$

Which results in

$$
S=\frac{\ln (x-1)}{2}+\frac{\ln (x+1)}{2}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{x y(a y-1)}{x^{2}-1}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =0 \\
R_{y} & =1 \\
S_{x} & =\frac{x}{x^{2}-1} \\
S_{y} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{1}{y(a y-1)} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{1}{R(R a-1)}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\ln (R a-1)-\ln (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{\ln (x-1)}{2}+\frac{\ln (x+1)}{2}=\ln (a y-1)-\ln (y)+c_{1}
$$

Which simplifies to

$$
\frac{\ln (x-1)}{2}+\frac{\ln (x+1)}{2}=\ln (a y-1)-\ln (y)+c_{1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{\ln (x-1)}{2}+\frac{\ln (x+1)}{2}=\ln (a y-1)-\ln (y)+c_{1} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
\frac{\ln (x-1)}{2}+\frac{\ln (x+1)}{2}=\ln (a y-1)-\ln (y)+c_{1}
$$

Verified OK.

### 1.48.3 Solving as bernoulli ode

In canonical form, the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{x y(a y-1)}{x^{2}-1}
\end{aligned}
$$

This is a Bernoulli ODE.

$$
\begin{equation*}
y^{\prime}=-\frac{x}{x^{2}-1} y+\frac{a x}{x^{2}-1} y^{2} \tag{1}
\end{equation*}
$$

The standard Bernoulli ODE has the form

$$
\begin{equation*}
y^{\prime}=f_{0}(x) y+f_{1}(x) y^{n} \tag{2}
\end{equation*}
$$

The first step is to divide the above equation by $y^{n}$ which gives

$$
\begin{equation*}
\frac{y^{\prime}}{y^{n}}=f_{0}(x) y^{1-n}+f_{1}(x) \tag{3}
\end{equation*}
$$

The next step is use the substitution $w=y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$
\begin{aligned}
f_{0}(x) & =-\frac{x}{x^{2}-1} \\
f_{1}(x) & =\frac{a x}{x^{2}-1} \\
n & =2
\end{aligned}
$$

Dividing both sides of ODE (1) by $y^{n}=y^{2}$ gives

$$
\begin{equation*}
y^{\prime} \frac{1}{y^{2}}=-\frac{x}{\left(x^{2}-1\right) y}+\frac{a x}{x^{2}-1} \tag{4}
\end{equation*}
$$

Let

$$
\begin{align*}
w & =y^{1-n} \\
& =\frac{1}{y} \tag{5}
\end{align*}
$$

Taking derivative of equation (5) w.r.t $x$ gives

$$
\begin{equation*}
w^{\prime}=-\frac{1}{y^{2}} y^{\prime} \tag{6}
\end{equation*}
$$

Substituting equations (5) and (6) into equation (4) gives

$$
\begin{align*}
-w^{\prime}(x) & =-\frac{x w(x)}{x^{2}-1}+\frac{a x}{x^{2}-1} \\
w^{\prime} & =\frac{x w}{x^{2}-1}-\frac{a x}{x^{2}-1} \tag{7}
\end{align*}
$$

The above now is a linear ODE in $w(x)$ which is now solved.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
w^{\prime}(x)+p(x) w(x)=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{x}{x^{2}-1} \\
& q(x)=-\frac{a x}{x^{2}-1}
\end{aligned}
$$

Hence the ode is

$$
w^{\prime}(x)-\frac{x w(x)}{x^{2}-1}=-\frac{a x}{x^{2}-1}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{x}{x^{2}-1} d x} \\
& =\mathrm{e}^{-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}}
\end{aligned}
$$

Which simplifies to

$$
\mu=\frac{1}{\sqrt{x-1} \sqrt{x+1}}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu w) & =(\mu)\left(-\frac{a x}{x^{2}-1}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{w}{\sqrt{x-1} \sqrt{x+1}}\right) & =\left(\frac{1}{\sqrt{x-1} \sqrt{x+1}}\right)\left(-\frac{a x}{x^{2}-1}\right) \\
\mathrm{d}\left(\frac{w}{\sqrt{x-1} \sqrt{x+1}}\right) & =\left(-\frac{a x}{\left(x^{2}-1\right) \sqrt{x-1} \sqrt{x+1}}\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{w}{\sqrt{x-1} \sqrt{x+1}}=\int-\frac{a x}{\left(x^{2}-1\right) \sqrt{x-1} \sqrt{x+1}} \mathrm{~d} x \\
& \frac{w}{\sqrt{x-1} \sqrt{x+1}}=\frac{\sqrt{x-1} \sqrt{x+1} a}{x^{2}-1}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{\sqrt{x-1} \sqrt{x+1}}$ results in

$$
w(x)=\frac{(x-1)(x+1) a}{x^{2}-1}+c_{1} \sqrt{x-1} \sqrt{x+1}
$$

which simplifies to

$$
w(x)=a+c_{1} \sqrt{x-1} \sqrt{x+1}
$$

Replacing $w$ in the above by $\frac{1}{y}$ using equation (5) gives the final solution.

$$
\frac{1}{y}=a+c_{1} \sqrt{x-1} \sqrt{x+1}
$$

Or

$$
y=\frac{1}{a+c_{1} \sqrt{x-1} \sqrt{x+1}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{1}{a+c_{1} \sqrt{x-1} \sqrt{x+1}} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{1}{a+c_{1} \sqrt{x-1} \sqrt{x+1}}
$$

Verified OK.

### 1.48.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition
$\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{y(a y-1)}\right) \mathrm{d} y & =\left(\frac{x}{x^{2}-1}\right) \mathrm{d} x \\
\left(-\frac{x}{x^{2}-1}\right) \mathrm{d} x+\left(\frac{1}{y(a y-1)}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =-\frac{x}{x^{2}-1} \\
N(x, y) & =\frac{1}{y(a y-1)}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{x}{x^{2}-1}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{y(a y-1)}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{x}{x^{2}-1} \mathrm{~d} x \\
\phi & =-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{y(a y-1)}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{y(a y-1)}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{y(a y-1)}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{y(a y-1)}\right) \mathrm{d} y \\
f(y) & =\ln (a y-1)-\ln (y)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}+\ln (a y-1)-\ln (y)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}+\ln (a y-1)-\ln (y)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}+\ln (a y-1)-\ln (y)=c_{1} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
-\frac{\ln (x-1)}{2}-\frac{\ln (x+1)}{2}+\ln (a y-1)-\ln (y)=c_{1}
$$

Verified OK.

### 1.48.5 Solving as riccati ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{x y(a y-1)}{x^{2}-1}
\end{aligned}
$$

This is a Riccati ODE. Comparing the ODE to solve

$$
y^{\prime}=\frac{x y^{2} a}{x^{2}-1}-\frac{x y}{x^{2}-1}
$$

With Riccati ODE standard form

$$
y^{\prime}=f_{0}(x)+f_{1}(x) y+f_{2}(x) y^{2}
$$

Shows that $f_{0}(x)=0, f_{1}(x)=-\frac{x}{x^{2}-1}$ and $f_{2}(x)=\frac{a x}{x^{2}-1}$. Let

$$
\begin{align*}
y & =\frac{-u^{\prime}}{f_{2} u} \\
& =\frac{-u^{\prime}}{\frac{a x u}{x^{2}-1}} \tag{1}
\end{align*}
$$

Using the above substitution in the given ODE results (after some simplification)in a second order ODE to solve for $u(x)$ which is

$$
\begin{equation*}
f_{2} u^{\prime \prime}(x)-\left(f_{2}^{\prime}+f_{1} f_{2}\right) u^{\prime}(x)+f_{2}^{2} f_{0} u(x)=0 \tag{2}
\end{equation*}
$$

But

$$
\begin{aligned}
f_{2}^{\prime} & =\frac{a}{x^{2}-1}-\frac{2 a x^{2}}{\left(x^{2}-1\right)^{2}} \\
f_{1} f_{2} & =-\frac{a x^{2}}{\left(x^{2}-1\right)^{2}} \\
f_{2}^{2} f_{0} & =0
\end{aligned}
$$

Substituting the above terms back in equation (2) gives

$$
\frac{a x u^{\prime \prime}(x)}{x^{2}-1}-\left(\frac{a}{x^{2}-1}-\frac{3 a x^{2}}{\left(x^{2}-1\right)^{2}}\right) u^{\prime}(x)=0
$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$
u(x)=c_{1}+\frac{c_{2}}{\sqrt{x^{2}-1}}
$$

The above shows that

$$
u^{\prime}(x)=-\frac{c_{2} x}{\left(x^{2}-1\right)^{\frac{3}{2}}}
$$

Using the above in (1) gives the solution

$$
y=\frac{c_{2}}{\sqrt{x^{2}-1} a\left(c_{1}+\frac{c_{2}}{\sqrt{x^{2}-1}}\right)}
$$

Dividing both numerator and denominator by $c_{1}$ gives, after renaming the constant $\frac{c_{2}}{c_{1}}=c_{3}$ the following solution

$$
y=\frac{1}{a\left(c_{3} \sqrt{x^{2}-1}+1\right)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{1}{a\left(c_{3} \sqrt{x^{2}-1}+1\right)} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{1}{a\left(c_{3} \sqrt{x^{2}-1}+1\right)}
$$

Verified OK.

### 1.48.6 Maple step by step solution

Let's solve

$$
\left(-x^{2}+1\right) y^{\prime}-y x+y^{2} a x=0
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables

$$
\frac{y^{\prime}}{y(a y-1)}=\frac{x}{(x-1)(x+1)}
$$

- Integrate both sides with respect to $x$

$$
\int \frac{y^{\prime}}{y(a y-1)} d x=\int \frac{x}{(x-1)(x+1)} d x+c_{1}
$$

- Evaluate integral

$$
\ln (a y-1)-\ln (y)=\frac{\ln ((x-1)(x+1))}{2}+c_{1}
$$

- $\quad$ Solve for $y$

$$
\left\{y=\frac{-a+\sqrt{x^{2} \mathrm{e}^{2 c_{1}}-\mathrm{e}^{2 c_{1}}}}{x^{2} \mathrm{e}^{\mathrm{e}^{2} c_{1}-a^{2}-\mathrm{e}^{2 c_{1}}}}, y=-\frac{a+\sqrt{x^{2} \mathrm{e}^{2 c_{1}}-\mathrm{e}^{2 c_{1}}}}{x^{2} \mathrm{e}^{2 c_{1}}-a^{2}-\mathrm{e}^{c_{1}}}\right\}
$$

## Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 20

```
dsolve((1-x^2)*diff (y(x),x)-x*y(x)+a*x*y(x)^2=0,y(x), singsol=all)
```

$$
y(x)=\frac{1}{\sqrt{-1+x} \sqrt{1+x} c_{1}+a}
$$

$\checkmark$ Solution by Mathematica
Time used: 4.118 (sec). Leaf size: 35
DSolve[(1- $\left.x^{\wedge} 2\right) * y '[x]-x * y[x]+a * x * y[x] \sim 2==0, y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow \frac{1}{a+e^{c_{1}} \sqrt{x^{2}-1}} \\
& y(x) \rightarrow 0 \\
& y(x) \rightarrow \frac{1}{a}
\end{aligned}
$$

### 1.49 problem 68

1.49.1 Solving as first order ode lie symmetry lookup ode . . . . . . . 589
1.49.2 Solving as bernoulli ode . . . . . . . . . . . . . . . . . . . . . . 592
1.49.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 595

Internal problem ID [12466]
Internal file name [OUTPUT/11118_Monday_October_16_2023_09_49_39_PM_45817509/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 68.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "bernoulli", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

```
[_rational, _Bernoulli]
```

$$
3 y^{\prime} y^{2}-a y^{3}=x+1
$$

### 1.49.1 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{a y^{3}+x+1}{3 y^{2}} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type Bernoulli. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 95: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\frac{\mathrm{e}^{a x}}{y^{2}} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{a^{a x}}{y^{2}}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y^{3} \mathrm{e}^{-a x}}{3}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{a y^{3}+x+1}{3 y^{2}}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{y^{3} a \mathrm{e}^{-a x}}{3} \\
S_{y} & =y^{2} \mathrm{e}^{-a x}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{\mathrm{e}^{-a x}(x+1)}{3} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{\mathrm{e}^{-a R}(R+1)}{3}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\frac{(a R+a+1) \mathrm{e}^{-a R}}{3 a^{2}}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{y^{3} \mathrm{e}^{-a x}}{3}=-\frac{(a x+a+1) \mathrm{e}^{-a x}}{3 a^{2}}+c_{1}
$$

Which simplifies to

$$
\frac{y^{3} \mathrm{e}^{-a x}}{3}=-\frac{(a x+a+1) \mathrm{e}^{-a x}}{3 a^{2}}+c_{1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{y^{3} \mathrm{e}^{-a x}}{3}=-\frac{(a x+a+1) \mathrm{e}^{-a x}}{3 a^{2}}+c_{1} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
\frac{y^{3} \mathrm{e}^{-a x}}{3}=-\frac{(a x+a+1) \mathrm{e}^{-a x}}{3 a^{2}}+c_{1}
$$

Verified OK.

### 1.49.2 Solving as bernoulli ode

In canonical form, the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{a y^{3}+x+1}{3 y^{2}}
\end{aligned}
$$

This is a Bernoulli ODE.

$$
\begin{equation*}
y^{\prime}=\frac{a}{3} y+\frac{x}{3}+\frac{1}{3} \frac{1}{y^{2}} \tag{1}
\end{equation*}
$$

The standard Bernoulli ODE has the form

$$
\begin{equation*}
y^{\prime}=f_{0}(x) y+f_{1}(x) y^{n} \tag{2}
\end{equation*}
$$

The first step is to divide the above equation by $y^{n}$ which gives

$$
\begin{equation*}
\frac{y^{\prime}}{y^{n}}=f_{0}(x) y^{1-n}+f_{1}(x) \tag{3}
\end{equation*}
$$

The next step is use the substitution $w=y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$
\begin{aligned}
f_{0}(x) & =\frac{a}{3} \\
f_{1}(x) & =\frac{x}{3}+\frac{1}{3} \\
n & =-2
\end{aligned}
$$

Dividing both sides of ODE (1) by $y^{n}=\frac{1}{y^{2}}$ gives

$$
\begin{equation*}
y^{\prime} y^{2}=\frac{a y^{3}}{3}+\frac{x}{3}+\frac{1}{3} \tag{4}
\end{equation*}
$$

Let

$$
\begin{align*}
w & =y^{1-n} \\
& =y^{3} \tag{5}
\end{align*}
$$

Taking derivative of equation (5) w.r.t $x$ gives

$$
\begin{equation*}
w^{\prime}=3 y^{2} y^{\prime} \tag{6}
\end{equation*}
$$

Substituting equations (5) and (6) into equation (4) gives

$$
\begin{align*}
\frac{w^{\prime}(x)}{3} & =\frac{a w(x)}{3}+\frac{x}{3}+\frac{1}{3} \\
w^{\prime} & =a w+x+1 \tag{7}
\end{align*}
$$

The above now is a linear ODE in $w(x)$ which is now solved.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
w^{\prime}(x)+p(x) w(x)=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-a \\
& q(x)=x+1
\end{aligned}
$$

Hence the ode is

$$
w^{\prime}(x)-a w(x)=x+1
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-a d x} \\
& =\mathrm{e}^{-a x}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu w) & =(\mu)(x+1) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\mathrm{e}^{-a x} w\right) & =\left(\mathrm{e}^{-a x}\right)(x+1) \\
\mathrm{d}\left(\mathrm{e}^{-a x} w\right) & =\left(\mathrm{e}^{-a x}(x+1)\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \mathrm{e}^{-a x} w=\int \mathrm{e}^{-a x}(x+1) \mathrm{d} x \\
& \mathrm{e}^{-a x} w=-\frac{(a x+a+1) \mathrm{e}^{-a x}}{a^{2}}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{-a x}$ results in

$$
w(x)=-\frac{\mathrm{e}^{a x}(a x+a+1) \mathrm{e}^{-a x}}{a^{2}}+c_{1} \mathrm{e}^{a x}
$$

which simplifies to

$$
w(x)=\frac{c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a}{a^{2}}
$$

Replacing $w$ in the above by $y^{3}$ using equation (5) gives the final solution.

$$
y^{3}=\frac{c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a}{a^{2}}
$$

Solving for $y$ gives

$$
\begin{aligned}
& y(x)=\frac{\left(\left(c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}}{a} \\
& y(x)=\frac{\left(\left(c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}(i \sqrt{3}-1)}{2 a} \\
& y(x)=-\frac{\left(\left(c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}(1+i \sqrt{3})}{2 a}
\end{aligned}
$$

## Summary

The solution(s) found are the following

$$
\begin{align*}
& y=\frac{\left(\left(c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}}{a}  \tag{1}\\
& y=\frac{\left(\left(c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}(i \sqrt{3}-1)}{2 a}  \tag{2}\\
& y=-\frac{\left(\left(c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}(1+i \sqrt{3})}{2 a} \tag{3}
\end{align*}
$$

Verification of solutions

$$
y=\frac{\left(\left(c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}}{a}
$$

Verified OK.

$$
y=\frac{\left(\left(c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}(i \sqrt{3}-1)}{2 a}
$$

Verified OK.

$$
y=-\frac{\left(\left(c_{1} \mathrm{e}^{a x} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}(1+i \sqrt{3})}{2 a}
$$

Verified OK.

### 1.49.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(3 y^{2}\right) \mathrm{d} y & =\left(a y^{3}+x+1\right) \mathrm{d} x \\
\left(-a y^{3}-x-1\right) \mathrm{d} x+\left(3 y^{2}\right) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =-a y^{3}-x-1 \\
N(x, y) & =3 y^{2}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-a y^{3}-x-1\right) \\
& =-3 a y^{2}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(3 y^{2}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =\frac{1}{3 y^{2}}\left(\left(-3 a y^{2}\right)-(0)\right) \\
& =-a
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-a \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-a x} \\
& =\mathrm{e}^{-a x}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\mathrm{e}^{-a x}\left(-a y^{3}-x-1\right) \\
& =-\mathrm{e}^{-a x}\left(a y^{3}+x+1\right)
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\mathrm{e}^{-a x}\left(3 y^{2}\right) \\
& =3 y^{2} \mathrm{e}^{-a x}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{array}{r}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x}=0 \\
\left(-\mathrm{e}^{-a x}\left(a y^{3}+x+1\right)\right)+\left(3 y^{2} \mathrm{e}^{-a x}\right) \frac{\mathrm{d} y}{\mathrm{~d} x}=0
\end{array}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\mathrm{e}^{-a x}\left(a y^{3}+x+1\right) \mathrm{d} x \\
\phi & =\frac{\left(y^{3} a^{2}+a x+a+1\right) \mathrm{e}^{-a x}}{a^{2}}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=3 y^{2} \mathrm{e}^{-a x}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=3 y^{2} \mathrm{e}^{-a x}$. Therefore equation (4) becomes

$$
\begin{equation*}
3 y^{2} \mathrm{e}^{-a x}=3 y^{2} \mathrm{e}^{-a x}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{\left(y^{3} a^{2}+a x+a+1\right) \mathrm{e}^{-a x}}{a^{2}}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{\left(y^{3} a^{2}+a x+a+1\right) \mathrm{e}^{-a x}}{a^{2}}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
\frac{\left(y^{3} a^{2}+a x+a+1\right) \mathrm{e}^{-a x}}{a^{2}}=c_{1} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
\frac{\left(y^{3} a^{2}+a x+a+1\right) \mathrm{e}^{-a x}}{a^{2}}=c_{1}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 106

```
dsolve(3*y(x)^ 2*diff(y(x),x)-a*y(x)^3-x-1=0,y(x), singsol=all)
```

$$
\begin{aligned}
& y(x)=\frac{\left(\left(\mathrm{e}^{a x} c_{1} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}}{a} \\
& y(x)=-\frac{\left(\left(\mathrm{e}^{a x} c_{1} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}(1+i \sqrt{3})}{2 a} \\
& y(x)=\frac{\left(\left(\mathrm{e}^{a x} c_{1} a^{2}-1+(-x-1) a\right) a\right)^{\frac{1}{3}}(i \sqrt{3}-1)}{2 a}
\end{aligned}
$$

$\sqrt{ }$ Solution by Mathematica
Time used: 17.994 (sec). Leaf size: 111
DSolve[3*y $[x] \sim 2 * y$ ' $[x]-a * y[x] \sim 3-x-1==0, y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow \frac{\sqrt[3]{a^{2} c_{1} e^{a x}-a(x+1)-1}}{a^{2 / 3}} \\
& y(x) \rightarrow-\frac{\sqrt[3]{-1} \sqrt[3]{a^{2} c_{1} e^{a x}-a(x+1)-1}}{a^{2 / 3}} \\
& y(x) \rightarrow \frac{(-1)^{2 / 3} \sqrt[3]{a^{2} c_{1} e^{a x}-a(x+1)-1}}{a^{2 / 3}}
\end{aligned}
$$

### 1.50 problem 69

Internal problem ID [12467]
Internal file name [OUTPUT/11119_Monday_October_16_2023_09_49_42_PM_22236623/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 69.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "unknown"
Maple gives the following as the ode type

```
[[_1st_order, `_with_symmetry_[F(x)*G(y),0]`]]
```

Unable to solve or complete the solution.

$$
y^{\prime}\left(y^{3} x^{2}+y x\right)=1
$$

Unable to determine ODE type.
Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
Looking for potential symmetries
trying inverse_Riccati
<- Bernoulli successful
<- inverse_Riccati successful`
```

$\checkmark$ Solution by Maple
Time used: 0.031 (sec). Leaf size: 78
dsolve(diff $(y(x), x) *\left(x^{\wedge} 2 * y(x) \wedge 3+x * y(x)\right)=1, y(x)$, singsol=all)

$$
\begin{aligned}
& y(x)=\frac{\sqrt{2 x^{2} \text { LambertW }\left(\frac{c_{1} \mathrm{e}^{-\frac{-1+2 x}{2 x}}}{2}\right)+2 x^{2}-x}}{x} \\
& y(x)=-\frac{\sqrt{2 x^{2} \text { LambertW }\left(\frac{c_{1} \mathrm{e}^{-\frac{-1+2 x}{2 x}}}{2}\right)+2 x^{2}-x}}{x}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.18 (sec). Leaf size: 76

DSolve[y'[x]*(x^2*y[x]^3+x*y[x])==1,y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow-\frac{\sqrt{2 x W\left(c_{1} e^{\frac{1}{2 x}-1}\right)+2 x-1}}{\sqrt{x}} \\
& y(x) \rightarrow \frac{\sqrt{2 x W\left(c_{1} e^{\frac{1}{2 x}-1}\right)+2 x-1}}{\sqrt{x}}
\end{aligned}
$$

### 1.51 problem 70

1.51.1 Solving as first order ode lie symmetry lookup ode . . . . . . . 603
1.51.2 Solving as bernoulli ode . . . . . . . . . . . . . . . . . . . . . . 607]
1.51.3 Solving as riccati ode . . . . . . . . . . . . . . . . . . . . . . . . 611

Internal problem ID [12468]
Internal file name [OUTPUT/11120_Monday_October_16_2023_09_49_43_PM_72894314/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 70.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "riccati", "bernoulli", "first_order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_Bernoulli]

$$
y^{\prime} x-(y \ln (x)-2) y=0
$$

### 1.51.1 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{(y \ln (x)-2) y}{x} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type Bernoulli. Therefore we do not need to solve the $\operatorname{PDE}(A)$, and can just use the lookup table shown below to find $\xi, \eta$

Table 97: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=x^{2} y^{2} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{x^{2} y^{2}} d y
\end{aligned}
$$

Which results in

$$
S=-\frac{1}{x^{2} y}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{(y \ln (x)-2) y}{x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{2}{x^{3} y} \\
S_{y} & =\frac{1}{x^{2} y^{2}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{\ln (x)}{x^{3}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{\ln (R)}{R^{3}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\frac{\ln (R)}{2 R^{2}}-\frac{1}{4 R^{2}}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
-\frac{1}{x^{2} y}=-\frac{\ln (x)}{2 x^{2}}-\frac{1}{4 x^{2}}+c_{1}
$$

Which simplifies to

$$
-\frac{1}{x^{2} y}=-\frac{\ln (x)}{2 x^{2}}-\frac{1}{4 x^{2}}+c_{1}
$$

Which gives

$$
y=\frac{4}{-4 c_{1} x^{2}+2 \ln (x)+1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.


## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{4}{-4 c_{1} x^{2}+2 \ln (x)+1} \tag{1}
\end{equation*}
$$



Figure 101: Slope field plot

Verification of solutions

$$
y=\frac{4}{-4 c_{1} x^{2}+2 \ln (x)+1}
$$

Verified OK.

### 1.51.2 Solving as bernoulli ode

In canonical form, the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{(y \ln (x)-2) y}{x}
\end{aligned}
$$

This is a Bernoulli ODE.

$$
\begin{equation*}
y^{\prime}=-\frac{2}{x} y+\frac{\ln (x)}{x} y^{2} \tag{1}
\end{equation*}
$$

The standard Bernoulli ODE has the form

$$
\begin{equation*}
y^{\prime}=f_{0}(x) y+f_{1}(x) y^{n} \tag{2}
\end{equation*}
$$

The first step is to divide the above equation by $y^{n}$ which gives

$$
\begin{equation*}
\frac{y^{\prime}}{y^{n}}=f_{0}(x) y^{1-n}+f_{1}(x) \tag{3}
\end{equation*}
$$

The next step is use the substitution $w=y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.
This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$
\begin{aligned}
f_{0}(x) & =-\frac{2}{x} \\
f_{1}(x) & =\frac{\ln (x)}{x} \\
n & =2
\end{aligned}
$$

Dividing both sides of ODE (1) by $y^{n}=y^{2}$ gives

$$
\begin{equation*}
y^{\prime} \frac{1}{y^{2}}=-\frac{2}{x y}+\frac{\ln (x)}{x} \tag{4}
\end{equation*}
$$

Let

$$
\begin{align*}
w & =y^{1-n} \\
& =\frac{1}{y} \tag{5}
\end{align*}
$$

Taking derivative of equation (5) w.r.t $x$ gives

$$
\begin{equation*}
w^{\prime}=-\frac{1}{y^{2}} y^{\prime} \tag{6}
\end{equation*}
$$

Substituting equations (5) and (6) into equation (4) gives

$$
\begin{align*}
-w^{\prime}(x) & =-\frac{2 w(x)}{x}+\frac{\ln (x)}{x} \\
w^{\prime} & =\frac{2 w}{x}-\frac{\ln (x)}{x} \tag{7}
\end{align*}
$$

The above now is a linear ODE in $w(x)$ which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
w^{\prime}(x)+p(x) w(x)=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{2}{x} \\
& q(x)=-\frac{\ln (x)}{x}
\end{aligned}
$$

Hence the ode is

$$
w^{\prime}(x)-\frac{2 w(x)}{x}=-\frac{\ln (x)}{x}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2}{x} d x} \\
& =\frac{1}{x^{2}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu w) & =(\mu)\left(-\frac{\ln (x)}{x}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{w}{x^{2}}\right) & =\left(\frac{1}{x^{2}}\right)\left(-\frac{\ln (x)}{x}\right) \\
\mathrm{d}\left(\frac{w}{x^{2}}\right) & =\left(-\frac{\ln (x)}{x^{3}}\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
\frac{w}{x^{2}} & =\int-\frac{\ln (x)}{x^{3}} \mathrm{~d} x \\
\frac{w}{x^{2}} & =\frac{\ln (x)}{2 x^{2}}+\frac{1}{4 x^{2}}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x^{2}}$ results in

$$
w(x)=x^{2}\left(\frac{\ln (x)}{2 x^{2}}+\frac{1}{4 x^{2}}\right)+c_{1} x^{2}
$$

which simplifies to

$$
w(x)=\frac{1}{4}+\frac{\ln (x)}{2}+c_{1} x^{2}
$$

Replacing $w$ in the above by $\frac{1}{y}$ using equation (5) gives the final solution.

$$
\frac{1}{y}=\frac{1}{4}+\frac{\ln (x)}{2}+c_{1} x^{2}
$$

Or

$$
y=\frac{1}{\frac{1}{4}+\frac{\ln (x)}{2}+c_{1} x^{2}}
$$

Which is simplified to

$$
y=\frac{4}{4 c_{1} x^{2}+2 \ln (x)+1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{4}{4 c_{1} x^{2}+2 \ln (x)+1} \tag{1}
\end{equation*}
$$



Figure 102: Slope field plot
Verification of solutions

$$
y=\frac{4}{4 c_{1} x^{2}+2 \ln (x)+1}
$$

Verified OK.

### 1.51.3 Solving as riccati ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{(y \ln (x)-2) y}{x}
\end{aligned}
$$

This is a Riccati ODE. Comparing the ODE to solve

$$
y^{\prime}=\frac{y^{2} \ln (x)}{x}-\frac{2 y}{x}
$$

With Riccati ODE standard form

$$
y^{\prime}=f_{0}(x)+f_{1}(x) y+f_{2}(x) y^{2}
$$

Shows that $f_{0}(x)=0, f_{1}(x)=-\frac{2}{x}$ and $f_{2}(x)=\frac{\ln (x)}{x}$. Let

$$
\begin{align*}
y & =\frac{-u^{\prime}}{f_{2} u} \\
& =\frac{-u^{\prime}}{\frac{\ln (x) u}{x}} \tag{1}
\end{align*}
$$

Using the above substitution in the given ODE results (after some simplification)in a second order ODE to solve for $u(x)$ which is

$$
\begin{equation*}
f_{2} u^{\prime \prime}(x)-\left(f_{2}^{\prime}+f_{1} f_{2}\right) u^{\prime}(x)+f_{2}^{2} f_{0} u(x)=0 \tag{2}
\end{equation*}
$$

But

$$
\begin{aligned}
f_{2}^{\prime} & =-\frac{\ln (x)}{x^{2}}+\frac{1}{x^{2}} \\
f_{1} f_{2} & =-\frac{2 \ln (x)}{x^{2}} \\
f_{2}^{2} f_{0} & =0
\end{aligned}
$$

Substituting the above terms back in equation (2) gives

$$
\frac{\ln (x) u^{\prime \prime}(x)}{x}-\left(-\frac{3 \ln (x)}{x^{2}}+\frac{1}{x^{2}}\right) u^{\prime}(x)=0
$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$
u(x)=c_{1}+\frac{(-2 \ln (x)-1) c_{2}}{4 x^{2}}
$$

The above shows that

$$
u^{\prime}(x)=\frac{c_{2} \ln (x)}{x^{3}}
$$

Using the above in (1) gives the solution

$$
y=-\frac{c_{2}}{x^{2}\left(c_{1}+\frac{(-2 \ln (x)-1) c_{2}}{4 x^{2}}\right)}
$$

Dividing both numerator and denominator by $c_{1}$ gives, after renaming the constant $\frac{c_{2}}{c_{1}}=c_{3}$ the following solution

$$
y=\frac{4}{-4 c_{3} x^{2}+2 \ln (x)+1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{4}{-4 c_{3} x^{2}+2 \ln (x)+1} \tag{1}
\end{equation*}
$$



Figure 103: Slope field plot

Verification of solutions

$$
y=\frac{4}{-4 c_{3} x^{2}+2 \ln (x)+1}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 20

```
dsolve(x*diff(y(x),x)=(y(x)*\operatorname{ln}(\textrm{x})-2)*y(x),y(x), singsol=all)
```

$$
y(x)=\frac{4}{1+4 c_{1} x^{2}+2 \ln (x)}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.254 (sec). Leaf size: 27
DSolve $\left[x * y y^{\prime}[x]==(y[x] * \log [x]-2) * y[x], y[x], x\right.$, IncludeSingularSolutions $->$ True]

$$
\begin{aligned}
& y(x) \rightarrow \frac{4}{4 c_{1} x^{2}+2 \log (x)+1} \\
& y(x) \rightarrow 0
\end{aligned}
$$

### 1.52 problem 71

1.52.1 Solving as first order ode lie symmetry lookup ode . . . . . . . 614
1.52.2 Solving as bernoulli ode . . . . . . . . . . . . . . . . . . . . . . 618
1.52.3 Solving as riccati ode . . . . . . . . . . . . . . . . . . . . . . . . 621

Internal problem ID [12469]
Internal file name [OUTPUT/11121_Monday_October_16_2023_09_49_43_PM_16125643/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 71.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "riccati", "bernoulli", "first_order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_Bernoulli]

$$
y-\cos (x) y^{\prime}-y^{2} \cos (x)(-\sin (x)+1)=0
$$

### 1.52.1 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{y(\cos (x) \sin (x) y-\cos (x) y+1)}{\cos (x)} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type Bernoulli. Therefore we do not need to solve the $\operatorname{PDE}(\mathrm{A})$, and can just use the lookup table shown below to find $\xi, \eta$

Table 99: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\frac{y^{2}}{\sec (x)+\tan (x)} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{y^{2}}{\sec (x)+\tan (x)}} d y
\end{aligned}
$$

Which results in

$$
S=-\frac{\sec (x)+\tan (x)}{y}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{y(\cos (x) \sin (x) y-\cos (x) y+1)}{\cos (x)}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{1}{(-1+\sin (x)) y} \\
S_{y} & =\frac{\sec (x)+\tan (x)}{y^{2}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-\cos (x) \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-\cos (R)
$$

The above is a quadrature ode．This is the whole point of Lie symmetry method． It converts an ode，no matter how complicated it is，to one that can be solved by integration when the ode is in the canonical coordiates $R, S$ ．Integrating the above gives

$$
\begin{equation*}
S(R)=-\sin (R)+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution，we just need to transform（4）back to $x, y$ coordinates．This results in

$$
\frac{-\sec (x)-\tan (x)}{y}=-\sin (x)+c_{1}
$$

Which simplifies to

$$
\frac{-\sec (x)-\tan (x)}{y}=-\sin (x)+c_{1}
$$

Which gives

$$
y=\frac{\sec (x)+\tan (x)}{\sin (x)-c_{1}}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown．

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{y(\cos (x) \sin (x) y-\cos (x) y+1)}{\cos (x)}$ |  | $\frac{d S}{d R}=-\cos (R)$ |
|  |  |  |
|  |  | $\rightarrow 0$ |
|  |  | $\rightarrow \gg 刀 S T R x^{2}$ |
|  |  |  |
|  | $R=x$ |  |
|  | $S-\sec (x)-\tan (x)$ | $\xrightarrow{x}$ |
|  | $S=\frac{y}{y}$ |  |
|  | $y$ | ア0ッ |
|  |  | ス |
| 人日＊＊ |  |  |
|  |  | タップ |

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{\sec (x)+\tan (x)}{\sin (x)-c_{1}} \tag{1}
\end{equation*}
$$



Figure 104: Slope field plot

## Verification of solutions

$$
y=\frac{\sec (x)+\tan (x)}{\sin (x)-c_{1}}
$$

Verified OK.

### 1.52.2 Solving as bernoulli ode

In canonical form, the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{y(\cos (x) \sin (x) y-\cos (x) y+1)}{\cos (x)}
\end{aligned}
$$

This is a Bernoulli ODE.

$$
\begin{equation*}
y^{\prime}=\frac{1}{\cos (x)} y+\frac{\cos (x) \sin (x)-\cos (x)}{\cos (x)} y^{2} \tag{1}
\end{equation*}
$$

The standard Bernoulli ODE has the form

$$
\begin{equation*}
y^{\prime}=f_{0}(x) y+f_{1}(x) y^{n} \tag{2}
\end{equation*}
$$

The first step is to divide the above equation by $y^{n}$ which gives

$$
\begin{equation*}
\frac{y^{\prime}}{y^{n}}=f_{0}(x) y^{1-n}+f_{1}(x) \tag{3}
\end{equation*}
$$

The next step is use the substitution $w=y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$
\begin{aligned}
f_{0}(x) & =\frac{1}{\cos (x)} \\
f_{1}(x) & =\frac{\cos (x) \sin (x)-\cos (x)}{\cos (x)} \\
n & =2
\end{aligned}
$$

Dividing both sides of ODE (1) by $y^{n}=y^{2}$ gives

$$
\begin{equation*}
y^{\prime} \frac{1}{y^{2}}=\frac{1}{\cos (x) y}+\frac{\cos (x) \sin (x)-\cos (x)}{\cos (x)} \tag{4}
\end{equation*}
$$

Let

$$
\begin{align*}
w & =y^{1-n} \\
& =\frac{1}{y} \tag{5}
\end{align*}
$$

Taking derivative of equation (5) w.r.t $x$ gives

$$
\begin{equation*}
w^{\prime}=-\frac{1}{y^{2}} y^{\prime} \tag{6}
\end{equation*}
$$

Substituting equations (5) and (6) into equation (4) gives

$$
\begin{align*}
-w^{\prime}(x) & =\frac{w(x)}{\cos (x)}+\frac{\cos (x) \sin (x)-\cos (x)}{\cos (x)} \\
w^{\prime} & =-\frac{w}{\cos (x)}-\frac{\cos (x) \sin (x)-\cos (x)}{\cos (x)} \tag{7}
\end{align*}
$$

The above now is a linear ODE in $w(x)$ which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
w^{\prime}(x)+p(x) w(x)=q(x)
$$

Where here

$$
\begin{aligned}
p(x) & =\sec (x) \\
q(x) & =-\sin (x)+1
\end{aligned}
$$

Hence the ode is

$$
w^{\prime}(x)+\sec (x) w(x)=-\sin (x)+1
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \sec (x) d x} \\
& =\sec (x)+\tan (x)
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu w) & =(\mu)(-\sin (x)+1) \\
\frac{\mathrm{d}}{\mathrm{~d} x}((\sec (x)+\tan (x)) w) & =(\sec (x)+\tan (x))(-\sin (x)+1) \\
\mathrm{d}((\sec (x)+\tan (x)) w) & =\cos (x) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& (\sec (x)+\tan (x)) w=\int \cos (x) \mathrm{d} x \\
& (\sec (x)+\tan (x)) w=\sin (x)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\sec (x)+\tan (x)$ results in

$$
w(x)=\frac{\sin (x)}{\sec (x)+\tan (x)}+\frac{c_{1}}{\sec (x)+\tan (x)}
$$

Replacing $w$ in the above by $\frac{1}{y}$ using equation (5) gives the final solution.

$$
\frac{1}{y}=\frac{\sin (x)}{\sec (x)+\tan (x)}+\frac{c_{1}}{\sec (x)+\tan (x)}
$$

Or

$$
y=\frac{1}{\frac{\sin (x)}{\sec (x)+\tan (x)}+\frac{c_{1}}{\sec (x)+\tan (x)}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{1}{\frac{\sin (x)}{\sec (x)+\tan (x)}+\frac{c_{1}}{\sec (x)+\tan (x)}} \tag{1}
\end{equation*}
$$



Figure 105: Slope field plot
Verification of solutions

$$
y=\frac{1}{\frac{\sin (x)}{\sec (x)+\tan (x)}+\frac{c_{1}}{\sec (x)+\tan (x)}}
$$

Verified OK.

### 1.52.3 Solving as riccati ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{y(\cos (x) \sin (x) y-\cos (x) y+1)}{\cos (x)}
\end{aligned}
$$

This is a Riccati ODE. Comparing the ODE to solve

$$
y^{\prime}=y^{2} \sin (x)-y^{2}+\frac{y}{\cos (x)}
$$

With Riccati ODE standard form

$$
y^{\prime}=f_{0}(x)+f_{1}(x) y+f_{2}(x) y^{2}
$$

Shows that $f_{0}(x)=0, f_{1}(x)=\frac{1}{\cos (x)}$ and $f_{2}(x)=\frac{\cos (x) \sin (x)-\cos (x)}{\cos (x)}$. Let

$$
\begin{align*}
y & =\frac{-u^{\prime}}{f_{2} u} \\
& =\frac{-u^{\prime}}{\frac{(\cos (x) \sin (x)-\cos (x)) u}{\cos (x)}} \tag{1}
\end{align*}
$$

Using the above substitution in the given ODE results (after some simplification)in a second order ODE to solve for $u(x)$ which is

$$
\begin{equation*}
f_{2} u^{\prime \prime}(x)-\left(f_{2}^{\prime}+f_{1} f_{2}\right) u^{\prime}(x)+f_{2}^{2} f_{0} u(x)=0 \tag{2}
\end{equation*}
$$

But

$$
\begin{aligned}
f_{2}^{\prime} & =\frac{-\sin (x)^{2}+\cos (x)^{2}+\sin (x)}{\cos (x)}+\frac{(\cos (x) \sin (x)-\cos (x)) \sin (x)}{\cos (x)^{2}} \\
f_{1} f_{2} & =\frac{\cos (x) \sin (x)-\cos (x)}{\cos (x)^{2}} \\
f_{2}^{2} f_{0} & =0
\end{aligned}
$$

Substituting the above terms back in equation (2) gives

$$
\frac{(\cos (x) \sin (x)-\cos (x)) u^{\prime \prime}(x)}{\cos (x)}-\left(\frac{-\sin (x)^{2}+\cos (x)^{2}+\sin (x)}{\cos (x)}+\frac{(\cos (x) \sin (x)-\cos (x)) \sin (x)}{\cos (x)^{2}}+\frac{\mathrm{c}}{}\right.
$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$
u(x)=c_{1}+c_{2} \sin (x)
$$

The above shows that

$$
u^{\prime}(x)=c_{2} \cos (x)
$$

Using the above in (1) gives the solution

$$
y=-\frac{c_{2} \cos (x)^{2}}{(\cos (x) \sin (x)-\cos (x))\left(c_{1}+c_{2} \sin (x)\right)}
$$

Dividing both numerator and denominator by $c_{1}$ gives, after renaming the constant $\frac{c_{2}}{c_{1}}=c_{3}$ the following solution

$$
y=-\frac{\cos (x)}{(-1+\sin (x))\left(c_{3}+\sin (x)\right)}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{\cos (x)}{(-1+\sin (x))\left(c_{3}+\sin (x)\right)} \tag{1}
\end{equation*}
$$



Figure 106: Slope field plot

Verification of solutions

$$
y=-\frac{\cos (x)}{(-1+\sin (x))\left(c_{3}+\sin (x)\right)}
$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 27

```
dsolve(y(x)-diff (y(x),x)*\operatorname{cos}(x)=y(x)^2*\operatorname{cos}(x)*(1-\operatorname{sin}(\textrm{x})),y(x), singsol=all)
```

$$
y(x)=\frac{\cos (x)+\sin (x)+1}{\left(\sin (x)+c_{1}\right)(-\sin (x)+\cos (x)+1)}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.778 (sec). Leaf size: 41

```
DSolve[y[x]-y'[x]*Cos[x]==y[x] 2*Cos[x]*(1-Sin[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$
\begin{aligned}
& y(x) \rightarrow \frac{e^{2 \operatorname{arctanh}\left(\tan \left(\frac{x}{2}\right)\right)}}{\cos (x) e^{2 \operatorname{arctanh}\left(\tan \left(\frac{x}{2}\right)\right)}+c_{1}} \\
& y(x) \rightarrow 0
\end{aligned}
$$

### 1.53 problem 72

1.53.1 Solving as differentialType ode625
1.53.2 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 627
1.53.3 Maple step by step solution 630

Internal problem ID [12470]
Internal file name [OUTPUT/11122_Monday_October_16_2023_09_49_47_PM_88426782/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 72.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "differentialType"
Maple gives the following as the ode type
[_exact, _rational, [_1st_order, ` _with_symmetry_[F(x),G(x)]`], [_Abel, `2nd type`, `class A`]]

$$
y+(x-2 y) y^{\prime}=-x^{2}
$$

### 1.53.1 Solving as differentialType ode

Writing the ode as

$$
\begin{equation*}
y^{\prime}=\frac{-x^{2}-y}{x-2 y} \tag{1}
\end{equation*}
$$

Which becomes

$$
\begin{equation*}
(-2 y) d y=(-x) d y+\left(-x^{2}-y\right) d x \tag{2}
\end{equation*}
$$

But the RHS is complete differential because

$$
(-x) d y+\left(-x^{2}-y\right) d x=d\left(-\frac{1}{3} x^{3}-x y\right)
$$

Hence (2) becomes

$$
(-2 y) d y=d\left(-\frac{1}{3} x^{3}-x y\right)
$$

Integrating both sides gives gives these solutions

$$
\begin{aligned}
& y=\frac{x}{2}+\frac{\sqrt{12 x^{3}+9 x^{2}-36 c_{1}}}{6}+c_{1} \\
& y=\frac{x}{2}-\frac{\sqrt{12 x^{3}+9 x^{2}-36 c_{1}}}{6}+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\frac{x}{2}+\frac{\sqrt{12 x^{3}+9 x^{2}-36 c_{1}}}{6}+c_{1}  \tag{1}\\
& y=\frac{x}{2}-\frac{\sqrt{12 x^{3}+9 x^{2}-36 c_{1}}}{6}+c_{1} \tag{2}
\end{align*}
$$

Figure 107: Slope field plot

## Verification of solutions

$$
y=\frac{x}{2}+\frac{\sqrt{12 x^{3}+9 x^{2}-36 c_{1}}}{6}+c_{1}
$$

Verified OK.

$$
y=\frac{x}{2}-\frac{\sqrt{12 x^{3}+9 x^{2}-36 c_{1}}}{6}+c_{1}
$$

Verified OK.

### 1.53.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(x-2 y) \mathrm{d} y & =\left(-x^{2}-y\right) \mathrm{d} x \\
\left(x^{2}+y\right) \mathrm{d} x+(x-2 y) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =x^{2}+y \\
N(x, y) & =x-2 y
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(x^{2}+y\right) \\
& =1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(x-2 y) \\
& =1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int x^{2}+y \mathrm{~d} x \\
\phi & =\frac{1}{3} x^{3}+x y+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=x+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=x-2 y$. Therefore equation (4) becomes

$$
\begin{equation*}
x-2 y=x+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=-2 y
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int(-2 y) \mathrm{d} y \\
f(y) & =-y^{2}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{1}{3} x^{3}+x y-y^{2}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{1}{3} x^{3}+x y-y^{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{x^{3}}{3}+y x-y^{2}=c_{1} \tag{1}
\end{equation*}
$$



Figure 108: Slope field plot
Verification of solutions

$$
\frac{x^{3}}{3}+y x-y^{2}=c_{1}
$$

Verified OK.

### 1.53.3 Maple step by step solution

Let's solve

$$
y+(x-2 y) y^{\prime}=-x^{2}
$$

- Highest derivative means the order of the ODE is 1

$$
y^{\prime}
$$

Check if ODE is exact

- ODE is exact if the lhs is the total derivative of a $C^{2}$ function

$$
F^{\prime}(x, y)=0
$$

- Compute derivative of lhs

$$
F^{\prime}(x, y)+\left(\frac{\partial}{\partial y} F(x, y)\right) y^{\prime}=0
$$

- Evaluate derivatives

$$
1=1
$$

- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$
\left[F(x, y)=c_{1}, M(x, y)=F^{\prime}(x, y), N(x, y)=\frac{\partial}{\partial y} F(x, y)\right]
$$

- $\quad$ Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to $x$

$$
F(x, y)=\int\left(x^{2}+y\right) d x+f_{1}(y)
$$

- Evaluate integral

$$
F(x, y)=\frac{x^{3}}{3}+x y+f_{1}(y)
$$

- $\quad$ Take derivative of $F(x, y)$ with respect to $y$

$$
N(x, y)=\frac{\partial}{\partial y} F(x, y)
$$

- Compute derivative
$x-2 y=x+\frac{d}{d y} f_{1}(y)$
- $\quad$ Isolate for $\frac{d}{d y} f_{1}(y)$
$\frac{d}{d y} f_{1}(y)=-2 y$
- $\quad$ Solve for $f_{1}(y)$
$f_{1}(y)=-y^{2}$
- $\quad$ Substitute $f_{1}(y)$ into equation for $F(x, y)$
$F(x, y)=\frac{1}{3} x^{3}+x y-y^{2}$
- $\quad$ Substitute $F(x, y)$ into the solution of the ODE $\frac{1}{3} x^{3}+x y-y^{2}=c_{1}$
- $\quad$ Solve for $y$

$$
\left\{y=\frac{x}{2}-\frac{\sqrt{12 x^{3}+9 x^{2}-36 c_{1}}}{6}, y=\frac{x}{2}+\frac{\sqrt{12 x^{3}+9 x^{2}-36 c_{1}}}{6}\right\}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 51

```
dsolve((x^2+y(x))+(x-2*y(x))*diff(y(x),x)=0,y(x), singsol=all)
```

$$
\begin{aligned}
& y(x)=\frac{x}{2}-\frac{\sqrt{12 x^{3}+9 x^{2}+36 c_{1}}}{6} \\
& y(x)=\frac{x}{2}+\frac{\sqrt{12 x^{3}+9 x^{2}+36 c_{1}}}{6}
\end{aligned}
$$

$\sqrt{\checkmark}$ Solution by Mathematica
Time used: 0.252 (sec). Leaf size: 81
DSolve $\left[\left(x^{\wedge} 2+y[x]\right)+(x-2 * y[x]) * y\right.$ ' $[x]==0, y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow \frac{1}{6}\left(3 x-i \sqrt{3} \sqrt{-4 x^{3}-3 x^{2}-12 c_{1}}\right) \\
& y(x) \rightarrow \frac{1}{6}\left(3 x+i \sqrt{3} \sqrt{-4 x^{3}-3 x^{2}-12 c_{1}}\right)
\end{aligned}
$$

### 1.54 problem 73

1.54.1 Solving as differentialType ode . . . . . . . . . . . . . . . . . . 633
1.54.2 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 635
1.54.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 638

Internal problem ID [12471]
Internal file name [OUTPUT/11123_Monday_October_16_2023_09_49_50_PM_64466181/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 73.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "differentialType"
Maple gives the following as the ode type

```
[_exact, _rational, [_1st_order, ` _with_symmetry_[F(x),G(x)]`],
    [_Abel, `2nd type`, `class A`]]
```

$$
y-(4 y-x) y^{\prime}=3 x^{2}
$$

### 1.54.1 Solving as differentialType ode

Writing the ode as

$$
\begin{equation*}
y^{\prime}=\frac{-y+3 x^{2}}{-4 y+x} \tag{1}
\end{equation*}
$$

Which becomes

$$
\begin{equation*}
(-4 y) d y=(-x) d y+\left(3 x^{2}-y\right) d x \tag{2}
\end{equation*}
$$

But the RHS is complete differential because

$$
(-x) d y+\left(3 x^{2}-y\right) d x=d\left(x^{3}-x y\right)
$$

Hence (2) becomes

$$
(-4 y) d y=d\left(x^{3}-x y\right)
$$

Integrating both sides gives gives these solutions

$$
\begin{aligned}
& y=\frac{x}{4}+\frac{\sqrt{-8 x^{3}+x^{2}-8 c_{1}}}{4}+c_{1} \\
& y=\frac{x}{4}-\frac{\sqrt{-8 x^{3}+x^{2}-8 c_{1}}}{4}+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{aligned}
& y=\frac{x}{4}+\frac{\sqrt{-8 x^{3}+x^{2}-8 c_{1}}}{4}+c_{1} \\
& y=\frac{x}{4}-\frac{\sqrt{-8 x^{3}+x^{2}-8 c_{1}}}{4}+c_{1}
\end{aligned}
$$

Figure 109: Slope field plot

## Verification of solutions

$$
y=\frac{x}{4}+\frac{\sqrt{-8 x^{3}+x^{2}-8 c_{1}}}{4}+c_{1}
$$

Verified OK.

$$
y=\frac{x}{4}-\frac{\sqrt{-8 x^{3}+x^{2}-8 c_{1}}}{4}+c_{1}
$$

Verified OK.

### 1.54.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
(-4 y+x) \mathrm{d} y & =\left(3 x^{2}-y\right) \mathrm{d} x \\
\left(-3 x^{2}+y\right) \mathrm{d} x+(-4 y+x) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =-3 x^{2}+y \\
N(x, y) & =-4 y+x
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-3 x^{2}+y\right) \\
& =1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(-4 y+x) \\
& =1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-3 x^{2}+y \mathrm{~d} x \\
\phi & =-x^{3}+x y+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=x+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=-4 y+x$. Therefore equation (4) becomes

$$
\begin{equation*}
-4 y+x=x+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=-4 y
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int(-4 y) \mathrm{d} y \\
f(y) & =-2 y^{2}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-x^{3}+x y-2 y^{2}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-x^{3}+x y-2 y^{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-x^{3}+y x-2 y^{2}=c_{1} \tag{1}
\end{equation*}
$$



Figure 110: Slope field plot
Verification of solutions

$$
-x^{3}+y x-2 y^{2}=c_{1}
$$

Verified OK.

### 1.54.3 Maple step by step solution

Let's solve

$$
y-(4 y-x) y^{\prime}=3 x^{2}
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
Check if ODE is exact
- ODE is exact if the lhs is the total derivative of a $C^{2}$ function

$$
F^{\prime}(x, y)=0
$$

- Compute derivative of lhs

$$
F^{\prime}(x, y)+\left(\frac{\partial}{\partial y} F(x, y)\right) y^{\prime}=0
$$

- Evaluate derivatives

$$
1=1
$$

- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$
\left[F(x, y)=c_{1}, M(x, y)=F^{\prime}(x, y), N(x, y)=\frac{\partial}{\partial y} F(x, y)\right]
$$

- $\quad$ Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to $x$

$$
F(x, y)=\int\left(-3 x^{2}+y\right) d x+f_{1}(y)
$$

- Evaluate integral

$$
F(x, y)=-x^{3}+x y+f_{1}(y)
$$

- $\quad$ Take derivative of $F(x, y)$ with respect to $y$

$$
N(x, y)=\frac{\partial}{\partial y} F(x, y)
$$

- Compute derivative

$$
-4 y+x=x+\frac{d}{d y} f_{1}(y)
$$

- Isolate for $\frac{d}{d y} f_{1}(y)$

$$
\frac{d}{d y} f_{1}(y)=-4 y
$$

- $\quad$ Solve for $f_{1}(y)$
$f_{1}(y)=-2 y^{2}$
- $\quad$ Substitute $f_{1}(y)$ into equation for $F(x, y)$
$F(x, y)=-x^{3}+x y-2 y^{2}$
- $\quad$ Substitute $F(x, y)$ into the solution of the ODE

$$
-x^{3}+x y-2 y^{2}=c_{1}
$$

- $\quad$ Solve for $y$

$$
\left\{y=\frac{x}{4}-\frac{\sqrt{-8 x^{3}+x^{2}-8 c_{1}}}{4}, y=\frac{x}{4}+\frac{\sqrt{-8 x^{3}+x^{2}-8 c_{1}}}{4}\right\}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 47

```
dsolve((y(x)-3*x^2)-(4*y(x)-x)*diff(y(x),x)=0,y(x), singsol=all)
```

$$
\begin{aligned}
& y(x)=\frac{x}{4}-\frac{\sqrt{-8 x^{3}+x^{2}+8 c_{1}}}{4} \\
& y(x)=\frac{x}{4}+\frac{\sqrt{-8 x^{3}+x^{2}+8 c_{1}}}{4}
\end{aligned}
$$

$\sqrt{ }$ Solution by Mathematica
Time used: 0.229 (sec). Leaf size: 67
DSolve[( $\left.\mathrm{y}[\mathrm{x}]-3 * \mathrm{x}^{\wedge} 2\right)-(4 * y[\mathrm{x}]-\mathrm{x}) * \mathrm{y}^{\prime}[\mathrm{x}]==0, \mathrm{y}[\mathrm{x}], \mathrm{x}$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow \frac{1}{4}\left(x-i \sqrt{8 x^{3}-x^{2}-16 c_{1}}\right) \\
& y(x) \rightarrow \frac{1}{4}\left(x+i \sqrt{8 x^{3}-x^{2}-16 c_{1}}\right)
\end{aligned}
$$

### 1.55 problem 74

1.55.1 Solving as differentialType ode . . . . . . . . . . . . . . . . . . 641
1.55.2 Solving as first order ode lie symmetry calculated ode . . . . . . 643
1.55.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 648
1.55.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 652

Internal problem ID [12472]
Internal file name [OUTPUT/11124_Monday_October_16_2023_09_49_52_PM_75845069/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 74.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "differentialType", "first__order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class G`], _exact, _rational]
```

$$
\left(y^{3}-x\right) y^{\prime}-y=0
$$

### 1.55.1 Solving as differentialType ode

Writing the ode as

$$
\begin{equation*}
y^{\prime}=\frac{y}{y^{3}-x} \tag{1}
\end{equation*}
$$

Which becomes

$$
\begin{equation*}
\left(-y^{3}\right) d y=(-x) d y+(-y) d x \tag{2}
\end{equation*}
$$

But the RHS is complete differential because

$$
(-x) d y+(-y) d x=d(-x y)
$$

Hence (2) becomes

$$
\left(-y^{3}\right) d y=d(-x y)
$$

Integrating both sides gives gives the solution as

$$
-\frac{y^{4}}{4}=-y x+c_{1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-\frac{y^{4}}{4}=-y x+c_{1} \tag{1}
\end{equation*}
$$



Figure 111: Slope field plot

Verification of solutions

$$
-\frac{y^{4}}{4}=-y x+c_{1}
$$

Verified OK.

### 1.55.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{y}{y^{3}-x} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{gather*}
b_{2}+\frac{y\left(b_{3}-a_{2}\right)}{y^{3}-x}-\frac{y^{2} a_{3}}{\left(y^{3}-x\right)^{2}}-\frac{y\left(x a_{2}+y a_{3}+a_{1}\right)}{\left(y^{3}-x\right)^{2}}  \tag{5E}\\
\quad-\left(\frac{1}{y^{3}-x}-\frac{3 y^{3}}{\left(y^{3}-x\right)^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{gather*}
$$

Putting the above in normal form gives

$$
\frac{y^{6} b_{2}-y^{4} a_{2}+3 y^{4} b_{3}+2 y^{3} b_{1}+2 x^{2} b_{2}-2 y^{2} a_{3}+x b_{1}-y a_{1}}{\left(-y^{3}+x\right)^{2}}=0
$$

Setting the numerator to zero gives

$$
\begin{equation*}
y^{6} b_{2}-y^{4} a_{2}+3 y^{4} b_{3}+2 y^{3} b_{1}+2 x^{2} b_{2}-2 y^{2} a_{3}+x b_{1}-y a_{1}=0 \tag{6E}
\end{equation*}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\{x, y\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{equation*}
b_{2} v_{2}^{6}-a_{2} v_{2}^{4}+3 b_{3} v_{2}^{4}+2 b_{1} v_{2}^{3}-2 a_{3} v_{2}^{2}+2 b_{2} v_{1}^{2}-a_{1} v_{2}+b_{1} v_{1}=0 \tag{7E}
\end{equation*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}\right\}
$$

Equation (7E) now becomes

$$
\begin{equation*}
2 b_{2} v_{1}^{2}+b_{1} v_{1}+b_{2} v_{2}^{6}+\left(-a_{2}+3 b_{3}\right) v_{2}^{4}+2 b_{1} v_{2}^{3}-2 a_{3} v_{2}^{2}-a_{1} v_{2}=0 \tag{8E}
\end{equation*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
b_{1} & =0 \\
b_{2} & =0 \\
-a_{1} & =0 \\
-2 a_{3} & =0 \\
2 b_{1} & =0 \\
2 b_{2} & =0 \\
-a_{2}+3 b_{3} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
a_{1} & =0 \\
a_{2} & =3 b_{3} \\
a_{3} & =0 \\
b_{1} & =0 \\
b_{2} & =0 \\
b_{3} & =b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=3 x \\
& \eta=y
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =y-\left(\frac{y}{y^{3}-x}\right)(3 x) \\
& =\frac{-y^{4}+4 x y}{-y^{3}+x} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{-y^{4}+4 x y}{-y^{3}+x}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{\ln \left(y\left(y^{3}-4 x\right)\right)}{4}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{y}{y^{3}-x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{1}{-y^{3}+4 x} \\
S_{y} & =\frac{-y^{3}+x}{y\left(-y^{3}+4 x\right)}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{\ln (y)}{4}+\frac{\ln \left(y^{3}-4 x\right)}{4}=c_{1}
$$

Which simplifies to

$$
\frac{\ln (y)}{4}+\frac{\ln \left(y^{3}-4 x\right)}{4}=c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown．

| Original ode in $x, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{y}{y^{3}-x}$ |  | $\frac{d S}{d R}=0$ |
| $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ 遇 |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{ }$ |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{ }$ |  | $\rightarrow \rightarrow \rightarrow$ |
| $\rightarrow \rightarrow \rightarrow \rightarrow$ 岛 $x]_{0}$ |  | S（R）$\longrightarrow \rightarrow \rightarrow \rightarrow \longrightarrow \longrightarrow$ |
| $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \infty$ |  |  |
| $\xrightarrow{\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow}$ | $\ln (u) \quad \ln \left(u^{3}-4 x\right)$ | ． |
| 边 | $S=\frac{\ln (y)}{4}+\frac{\ln \left(y^{3}-4 x\right)}{4}$ |  |
| $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-\infty \rightarrow \rightarrow \rightarrow \rightarrow \infty \rightarrow \infty$ | － 4 |  |
| $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{ }$ |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow- \pm+\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow]{ }$ |  | $\rightarrow \rightarrow \rightarrow \rightarrow-4_{4} \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ |

## Summary

The solution（s）found are the following

$$
\begin{equation*}
\frac{\ln (y)}{4}+\frac{\ln \left(y^{3}-4 x\right)}{4}=c_{1} \tag{1}
\end{equation*}
$$



Figure 112: Slope field plot
Verification of solutions

$$
\frac{\ln (y)}{4}+\frac{\ln \left(y^{3}-4 x\right)}{4}=c_{1}
$$

Verified OK.

### 1.55.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(y^{3}-x\right) \mathrm{d} y & =(y) \mathrm{d} x \\
(-y) \mathrm{d} x+\left(y^{3}-x\right) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =-y \\
N(x, y) & =y^{3}-x
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(-y) \\
& =-1
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(y^{3}-x\right) \\
& =-1
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-y \mathrm{~d} x \\
\phi & =-x y+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=-x+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=y^{3}-x$. Therefore equation (4) becomes

$$
\begin{equation*}
y^{3}-x=-x+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=y^{3}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(y^{3}\right) \mathrm{d} y \\
f(y) & =\frac{y^{4}}{4}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-x y+\frac{1}{4} y^{4}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-x y+\frac{1}{4} y^{4}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{y^{4}}{4}-y x=c_{1} \tag{1}
\end{equation*}
$$



Figure 113: Slope field plot

Verification of solutions

$$
\frac{y^{4}}{4}-y x=c_{1}
$$

Verified OK.

### 1.55.4 Maple step by step solution

Let's solve

$$
\left(y^{3}-x\right) y^{\prime}-y=0
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
Check if ODE is exact
- ODE is exact if the lhs is the total derivative of a $C^{2}$ function
$F^{\prime}(x, y)=0$
- Compute derivative of lhs
$F^{\prime}(x, y)+\left(\frac{\partial}{\partial y} F(x, y)\right) y^{\prime}=0$
- Evaluate derivatives

$$
-1=-1
$$

- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$
\left[F(x, y)=c_{1}, M(x, y)=F^{\prime}(x, y), N(x, y)=\frac{\partial}{\partial y} F(x, y)\right]
$$

- $\quad$ Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to $x$ $F(x, y)=\int-y d x+f_{1}(y)$
- Evaluate integral

$$
F(x, y)=-x y+f_{1}(y)
$$

- $\quad$ Take derivative of $F(x, y)$ with respect to $y$

$$
N(x, y)=\frac{\partial}{\partial y} F(x, y)
$$

- Compute derivative

$$
y^{3}-x=-x+\frac{d}{d y} f_{1}(y)
$$

- $\quad$ Isolate for $\frac{d}{d y} f_{1}(y)$
$\frac{d}{d y} f_{1}(y)=y^{3}$
- $\quad$ Solve for $f_{1}(y)$
$f_{1}(y)=\frac{y^{4}}{4}$
- $\quad$ Substitute $f_{1}(y)$ into equation for $F(x, y)$

$$
F(x, y)=-x y+\frac{1}{4} y^{4}
$$

- $\quad$ Substitute $F(x, y)$ into the solution of the ODE

$$
-x y+\frac{1}{4} y^{4}=c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\operatorname{RootOf}\left(\_Z^{4}-4 \_Z x-4 c_{1}\right)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
<- 1st order linear successful
<- inverse linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 18

```
dsolve((y(x)^3-x)*diff(y(x),x)=y(x),y(x), singsol=all)
```

$$
-\frac{c_{1}}{y(x)}+x-\frac{y(x)^{3}}{4}=0
$$

## Solution by Mathematic

Time used: 57.499 (sec). Leaf size: 996

```
DSolve[(y[x]^3-x)*y'[x]==y[x],y[x],x,IncludeSingularSolutions -> True]
```

$$
\begin{aligned}
&-\frac{\sqrt{\frac{\left(9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}\right)^{2 / 3}+4 \sqrt[3]{3} c_{1}}{\sqrt[3]{9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}}}}}{\sqrt{2}(x)} \\
& \sqrt{-\frac{\sqrt[3]{3}}{\sqrt{\frac{\left(9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}\right)^{2 / 3}+4 \sqrt[3]{3} c_{1}}{\sqrt[3]{9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}}}}}-\frac{2 \sqrt[3]{9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}}}{3^{2 / 3}}-\frac{8 c_{1}}{\sqrt[3]{27 x^{2}-3 \sqrt{81 x^{4}-192 c_{1}^{3}}}}}
\end{aligned}
$$

$$
y(x)
$$

$$
\rightarrow \frac{1}{2}\left(\sqrt{-\frac{4 \sqrt{2} \sqrt[3]{3} x}{\sqrt{\frac{\left(9 x^{2}-\sqrt{81 x^{4}-19 c_{1}^{3}}\right)^{2 / 3}+4 \sqrt[3]{3} c_{1}}{\sqrt[3]{9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}}}}-\frac{2 \sqrt[3]{9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}}}{3^{2 / 3}}-\frac{8 c_{1}}{\sqrt[3]{27 x^{2}-3 \sqrt{81 x^{4}-192 c_{1}^{3}}}}}} \begin{array}{c}
\sqrt{2} \sqrt{\frac{\left(9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}\right)^{2 / 3}+4 \sqrt[3]{3} c_{1}}{\sqrt[3]{9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}}}} \\
-\frac{\sqrt[3]{3}}{\sqrt{3}}
\end{array}\right)
$$

$$
y(x) \rightarrow \frac{\sqrt{\frac{\left(9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}\right)^{2 / 3}+4 \sqrt[3]{3} c_{1}}{\sqrt[3]{9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}}}}}{\sqrt{2} \sqrt[3]{3}}
$$

$$
-\frac{1}{2} \sqrt{\frac{4 \sqrt{2} \sqrt[3]{3} x}{\sqrt{\frac{\left(9 x^{2}-\sqrt{81 x^{4}-192 c_{1}{ }^{2}}\right)^{2 / 3}+4 \sqrt[3]{3} c_{1}}{}}}-\frac{2 \sqrt[3]{9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}}}{3^{2 / 3}}}-\frac{8 c_{1}}{\sqrt[3]{27 x^{2}-3 \sqrt{81 x^{4}-192 c_{1}^{3}}}}
$$

$$
\sqrt{ } \sqrt{\frac{\left(9 x^{2}-\sqrt{ } 81 x^{4}-192 c_{1}{ }^{3}\right)^{2 / 5}+4 \vee J c_{1}}{\sqrt[3]{9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}}}}
$$

$$
y(x) \rightarrow \frac{1}{2}\left(\frac{\sqrt{2} \sqrt{\frac{\left(9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}\right)^{2 / 3}+4 \sqrt[3]{3} c_{1}}{\sqrt[3]{9 x^{2}-\sqrt{81 x^{4}-192 c_{1}^{3}}}}}}{\sqrt[3]{3}}\right.
$$

### 1.56 problem 75

1.56.1 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 655
1.56.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 659

Internal problem ID [12473]
Internal file name [OUTPUT/11125_Monday_October_16_2023_09_49_53_PM_18849345/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 75 .
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact"
Maple gives the following as the ode type
[_exact, _rational]

$$
\frac{y^{2}}{(-y+x)^{2}}+\left(\frac{1}{y}-\frac{x^{2}}{(-y+x)^{2}}\right) y^{\prime}=\frac{1}{x}
$$

### 1.56.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{y}-\frac{x^{2}}{(-y+x)^{2}}\right) \mathrm{d} y & =\left(-\frac{y^{2}}{(-y+x)^{2}}+\frac{1}{x}\right) \mathrm{d} x \\
\left(\frac{y^{2}}{(-y+x)^{2}}-\frac{1}{x}\right) \mathrm{d} x+\left(\frac{1}{y}-\frac{x^{2}}{(-y+x)^{2}}\right) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=\frac{y^{2}}{(-y+x)^{2}}-\frac{1}{x} \\
& N(x, y)=\frac{1}{y}-\frac{x^{2}}{(-y+x)^{2}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\frac{y^{2}}{(-y+x)^{2}}-\frac{1}{x}\right) \\
& =\frac{2 x y}{(-y+x)^{3}}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{y}-\frac{x^{2}}{(-y+x)^{2}}\right) \\
& =\frac{2 x y}{(-y+x)^{3}}
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{y^{2}}{(-y+x)^{2}}-\frac{1}{x} \mathrm{~d} x \\
\phi & =-\frac{y^{2}}{-y+x}-\ln (x)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{align*}
\frac{\partial \phi}{\partial y} & =-\frac{2 y}{-y+x}-\frac{y^{2}}{(-y+x)^{2}}+f^{\prime}(y)  \tag{4}\\
& =\frac{-2 x y+y^{2}}{(-y+x)^{2}}+f^{\prime}(y)
\end{align*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{y}-\frac{x^{2}}{(-y+x)^{2}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{y}-\frac{x^{2}}{(-y+x)^{2}}=\frac{-2 x y+y^{2}}{(-y+x)^{2}}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=-\frac{y-1}{y}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1-y}{y}\right) \mathrm{d} y \\
f(y) & =-y+\ln (y)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\frac{y^{2}}{-y+x}-\ln (x)-y+\ln (y)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\frac{y^{2}}{-y+x}-\ln (x)-y+\ln (y)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
-\frac{y^{2}}{-y+x}-\ln (x)-y+\ln (y)=c_{1} \tag{1}
\end{equation*}
$$



Figure 114: Slope field plot

Verification of solutions

$$
-\frac{y^{2}}{-y+x}-\ln (x)-y+\ln (y)=c_{1}
$$

Verified OK.

### 1.56.2 Maple step by step solution

Let's solve

$$
\frac{y^{2}}{(-y+x)^{2}}+\left(\frac{1}{y}-\frac{x^{2}}{(-y+x)^{2}}\right) y^{\prime}=\frac{1}{x}
$$

- Highest derivative means the order of the ODE is 1

$$
y^{\prime}
$$

$\square \quad$ Check if ODE is exact

- ODE is exact if the lhs is the total derivative of a $C^{2}$ function $F^{\prime}(x, y)=0$
- Compute derivative of lhs

$$
F^{\prime}(x, y)+\left(\frac{\partial}{\partial y} F(x, y)\right) y^{\prime}=0
$$

- Evaluate derivatives
$\frac{2 y}{(-y+x)^{2}}+\frac{2 y^{2}}{(-y+x)^{3}}=-\frac{2 x}{(-y+x)^{2}}+\frac{2 x^{2}}{(-y+x)^{3}}$
- Simplify

$$
\frac{2 x y}{(-y+x)^{3}}=\frac{2 x y}{(-y+x)^{3}}
$$

- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$
\left[F(x, y)=c_{1}, M(x, y)=F^{\prime}(x, y), N(x, y)=\frac{\partial}{\partial y} F(x, y)\right]
$$

- $\quad$ Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to $x$

$$
F(x, y)=\int\left(\frac{y^{2}}{(-y+x)^{2}}-\frac{1}{x}\right) d x+f_{1}(y)
$$

- Evaluate integral
$F(x, y)=-\frac{y^{2}}{-y+x}-\ln (x)+f_{1}(y)$
- $\quad$ Take derivative of $F(x, y)$ with respect to $y$
$N(x, y)=\frac{\partial}{\partial y} F(x, y)$
- Compute derivative
$\frac{1}{y}-\frac{x^{2}}{(-y+x)^{2}}=-\frac{2 y}{-y+x}-\frac{y^{2}}{(-y+x)^{2}}+\frac{d}{d y} f_{1}(y)$
- $\quad$ Isolate for $\frac{d}{d y} f_{1}(y)$

$$
\frac{d}{d y} f_{1}(y)=\frac{1}{y}-\frac{x^{2}}{(-y+x)^{2}}+\frac{2 y}{-y+x}+\frac{y^{2}}{(-y+x)^{2}}
$$

- $\quad$ Solve for $f_{1}(y)$
$f_{1}(y)=\ln (y)-y-\frac{x^{2}}{y-x}-\frac{x^{2}}{-y+x}$
- $\quad$ Substitute $f_{1}(y)$ into equation for $F(x, y)$

$$
F(x, y)=-\frac{y^{2}}{-y+x}-\ln (x)+\ln (y)-y-\frac{x^{2}}{y-x}-\frac{x^{2}}{-y+x}
$$

- $\quad$ Substitute $F(x, y)$ into the solution of the ODE

$$
-\frac{y^{2}}{-y+x}-\ln (x)+\ln (y)-y-\frac{x^{2}}{y-x}-\frac{x^{2}}{-y+x}=c_{1}
$$

- $\quad$ Solve for $y$
$y=\mathrm{e}^{\operatorname{RootOf}\left(\mathrm{e}^{Z} \ln (x)+c_{1} \mathrm{e}^{Z}-\mathrm{e}^{Z}-Z-x \mathrm{e}^{Z}-x \ln (x)-c_{1} x+\_Z x\right)}$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`
```

$\checkmark$ Solution by Maple
Time used: 0.094 (sec). Leaf size: 37
dsolve( $\left(y(x) \wedge 2 /(x-y(x))^{\wedge} 2-1 / x\right)+\left(1 / y(x)-x^{\wedge} 2 /(x-y(x))^{\wedge} 2\right) * \operatorname{diff}(y(x), x)=0, y(x)$, singsol=all)

$$
y(x)=\mathrm{e}^{\operatorname{RootOf}\left(-\ln (x) \mathrm{e}^{Z}+c_{1} \mathrm{e}^{Z}+\_Z \mathrm{e}^{Z}+x \mathrm{e}^{Z}+\ln (x) x-c_{1} x-\_Z x\right)}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.721 (sec). Leaf size: 29
DSolve $\left[(y[x] \sim 2 /(x-y[x]) \wedge 2-1 / x)+\left(1 / y[x]-x^{\wedge} 2 /(x-y[x]) \wedge 2\right) * y '[x]==0, y[x], x\right.$, IncludeSingularSolut

Solve $\left[\frac{y(x)^{2}}{x-y(x)}+y(x)-\log (y(x))+\log (x)=c_{1}, y(x)\right]$

### 1.57 problem 76

1.57.1 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 662
1.57.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 666

Internal problem ID [12474]
Internal file name [OUTPUT/11126_Monday_October_16_2023_09_49_56_PM_1537815/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 76.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact"
Maple gives the following as the ode type

```
[_exact, _rational]
```

$$
6 y^{2} x+3\left(2 x^{2} y+y^{2}\right) y^{\prime}=-4 x^{3}
$$

### 1.57.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(6 x^{2} y+3 y^{2}\right) \mathrm{d} y & =\left(-4 x^{3}-6 x y^{2}\right) \mathrm{d} x \\
\left(4 x^{3}+6 x y^{2}\right) \mathrm{d} x+\left(6 x^{2} y+3 y^{2}\right) \mathrm{d} y & =0 \tag{2A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
M(x, y) & =4 x^{3}+6 x y^{2} \\
N(x, y) & =6 x^{2} y+3 y^{2}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(4 x^{3}+6 x y^{2}\right) \\
& =12 x y
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(6 x^{2} y+3 y^{2}\right) \\
& =12 x y
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int 4 x^{3}+6 x y^{2} \mathrm{~d} x \\
\phi & =\frac{\left(2 x^{2}+3 y^{2}\right)^{2}}{4}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{align*}
\frac{\partial \phi}{\partial y} & =3\left(2 x^{2}+3 y^{2}\right) y+f^{\prime}(y)  \tag{4}\\
& =6 x^{2} y+9 y^{3}+f^{\prime}(y)
\end{align*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=6 x^{2} y+3 y^{2}$. Therefore equation (4) becomes

$$
\begin{equation*}
6 x^{2} y+3 y^{2}=6 x^{2} y+9 y^{3}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=-9 y^{3}+3 y^{2}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(-9 y^{3}+3 y^{2}\right) \mathrm{d} y \\
f(y) & =-\frac{9}{4} y^{4}+y^{3}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{\left(2 x^{2}+3 y^{2}\right)^{2}}{4}-\frac{9 y^{4}}{4}+y^{3}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{\left(2 x^{2}+3 y^{2}\right)^{2}}{4}-\frac{9 y^{4}}{4}+y^{3}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{\left(2 x^{2}+3 y^{2}\right)^{2}}{4}-\frac{9 y^{4}}{4}+y^{3}=c_{1} \tag{1}
\end{equation*}
$$



Figure 115: Slope field plot

Verification of solutions

$$
\frac{\left(2 x^{2}+3 y^{2}\right)^{2}}{4}-\frac{9 y^{4}}{4}+y^{3}=c_{1}
$$

Verified OK.

### 1.57.2 Maple step by step solution

Let's solve

$$
6 y^{2} x+3\left(2 x^{2} y+y^{2}\right) y^{\prime}=-4 x^{3}
$$

- Highest derivative means the order of the ODE is 1

$$
y^{\prime}
$$

## Check if ODE is exact

- ODE is exact if the lhs is the total derivative of a $C^{2}$ function $F^{\prime}(x, y)=0$
- Compute derivative of lhs

$$
F^{\prime}(x, y)+\left(\frac{\partial}{\partial y} F(x, y)\right) y^{\prime}=0
$$

- Evaluate derivatives

$$
12 x y=12 x y
$$

- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$
\left[F(x, y)=c_{1}, M(x, y)=F^{\prime}(x, y), N(x, y)=\frac{\partial}{\partial y} F(x, y)\right]
$$

- $\quad$ Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to $x$

$$
F(x, y)=\int\left(4 x^{3}+6 x y^{2}\right) d x+f_{1}(y)
$$

- Evaluate integral

$$
F(x, y)=\frac{\left(2 x^{2}+3 y^{2}\right)^{2}}{4}+f_{1}(y)
$$

- $\quad$ Take derivative of $F(x, y)$ with respect to $y$

$$
N(x, y)=\frac{\partial}{\partial y} F(x, y)
$$

- Compute derivative

$$
6 x^{2} y+3 y^{2}=3\left(2 x^{2}+3 y^{2}\right) y+\frac{d}{d y} f_{1}(y)
$$

- Isolate for $\frac{d}{d y} f_{1}(y)$

$$
\frac{d}{d y} f_{1}(y)=6 x^{2} y+3 y^{2}-3\left(2 x^{2}+3 y^{2}\right) y
$$

- $\quad$ Solve for $f_{1}(y)$

$$
f_{1}(y)=-\frac{9}{4} y^{4}+y^{3}
$$

- $\quad$ Substitute $f_{1}(y)$ into equation for $F(x, y)$

$$
F(x, y)=\frac{\left(2 x^{2}+3 y^{2}\right)^{2}}{4}-\frac{9 y^{4}}{4}+y^{3}
$$

- $\quad$ Substitute $F(x, y)$ into the solution of the ODE

$$
\frac{\left(2 x^{2}+3 y^{2}\right)^{2}}{4}-\frac{9 y^{4}}{4}+y^{3}=c_{1}
$$

- $\quad$ Solve for $y$

$$
\left\{y=\frac{\left(-4 x^{4}+4 c_{1}-8 x^{6}+4 \sqrt{4 x^{10}+x^{8}-4 c_{1} x^{6}-2 c_{1} x^{4}+c_{1}^{2}}\right)^{\frac{1}{3}}}{2}+\frac{2 x^{4}}{\left(-4 x^{4}+4 c_{1}-8 x^{6}+4 \sqrt{4 x^{10}+x^{8}-4 c_{1} x^{6}-2 c_{1} x^{4}+c_{1}^{2}}\right)^{\frac{1}{3}}}-x^{2},\right.
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 422
dsolve ( $2 *\left(3 * x * y(x) \wedge 2+2 * x^{\wedge} 3\right)+3 *\left(2 * x^{\wedge} 2 * y(x)+y(x) \wedge 2\right) * \operatorname{diff}(y(x), x)=0, y(x)$, singsol=all)

$$
\begin{aligned}
& y(x)= \frac{\left(-4 x^{4}-4 c_{1}-8 x^{6}+4 \sqrt{4 x^{10}+x^{8}+4 c_{1} x^{6}+2 c_{1} x^{4}+c_{1}^{2}}\right)^{\frac{1}{3}}}{2} \\
&+\frac{2 x^{4}}{\left(-4 x^{4}-4 c_{1}-8 x^{6}+4 \sqrt{4 x^{10}+x^{8}+4 c_{1} x^{6}+2 c_{1} x^{4}+c_{1}^{2}}\right)^{\frac{1}{3}}}-x^{2} \\
& y(x)
\end{aligned}
$$

$$
\begin{aligned}
& =\frac{4 i \sqrt{3} x^{4}-i \sqrt{3}\left(-4 x^{4}-4 c_{1}-8 x^{6}+4 \sqrt{\left(4 x^{6}+x^{4}+c_{1}\right)\left(x^{4}+c_{1}\right)}\right)^{\frac{2}{3}}-4 x^{4}-4 x^{2}\left(-4 x^{4}-4 c_{1}-8 x^{6}+\right.}{4\left(-4 x^{4}-4 c_{1}-8 x^{6}+4 \sqrt{\left(4 x^{6}+x\right.}\right.} \\
& y(x)=\frac{(i \sqrt{3}-1)\left(-4 x^{4}-4 c_{1}-8 x^{6}+4 \sqrt{\left(4 x^{6}+x^{4}+c_{1}\right)\left(x^{4}+c_{1}\right)}\right)^{\frac{1}{3}}}{4} \\
& \quad-\frac{\left(i \sqrt{3} x^{2}+x^{2}+\left(-4 x^{4}-4 c_{1}-8 x^{6}+4 \sqrt{\left(4 x^{6}+x^{4}+c_{1}\right)\left(x^{4}+c_{1}\right)}\right)^{\frac{1}{3}}\right) x^{2}}{\left(-4 x^{4}-4 c_{1}-8 x^{6}+4 \sqrt{\left(4 x^{6}+x^{4}+c_{1}\right)\left(x^{4}+c_{1}\right)}\right)^{\frac{1}{3}}}
\end{aligned}
$$

## Solution by Mathematica

Time used: 25.227 (sec). Leaf size: 419
DSolve $\left[2 *\left(3 * x * y[x] \wedge 2+2 * x^{\wedge} 3\right)+3 *(2 * x \wedge 2 * y[x]+y[x] \sim 2) * y \cdot[x]==0, y[x], x\right.$, IncludeSingularSolutions

$$
\begin{aligned}
y(x) \rightarrow & -x^{2}+\frac{\sqrt[3]{2} x^{4}}{\sqrt[3]{-2 x^{6}-x^{4}+\sqrt{4 x^{10}+x^{8}-4 c_{1} x^{6}-2 c_{1} x^{4}+c_{1}^{2}}+c_{1}}} \\
+ & \frac{\sqrt[3]{-2 x^{6}-x^{4}+\sqrt{4 x^{10}+x^{8}-4 c_{1} x^{6}-2 c_{1} x^{4}+c_{1}^{2}}+c_{1}}}{\sqrt[3]{2}} \\
y(x) \rightarrow & \frac{1}{4}\left(-4 x^{2}-\frac{2 \sqrt[3]{2}(1+i \sqrt{3}) x^{4}}{\sqrt[3]{-2 x^{6}-x^{4}+\sqrt{4 x^{10}+x^{8}-4 c_{1} x^{6}-2 c_{1} x^{4}+c_{1}^{2}}+c_{1}}}\right. \\
& \left.+i 2^{2 / 3}(\sqrt{3}+i) \sqrt[3]{-2 x^{6}-x^{4}+\sqrt{4 x^{10}+x^{8}-4 c_{1} x^{6}-2 c_{1} x^{4}+c_{1}^{2}}+c_{1}}\right) \\
y(x) \rightarrow & \frac{1}{4}\left(-4 x^{2}+\frac{2 i \sqrt[3]{2}(\sqrt{3}+i) x^{4}}{\sqrt[3]{-2 x^{6}-x^{4}+\sqrt{4 x^{10}+x^{8}-4 c_{1} x^{6}-2 c_{1} x^{4}+c_{1}^{2}}+c_{1}}}\right. \\
& \left.+2^{2 / 3}(-1-i \sqrt{3}) \sqrt[3]{-2 x^{6}-x^{4}+\sqrt{4 x^{10}+x^{8}-4 c_{1} x^{6}-2 c_{1} x^{4}+c_{1}^{2}}+c_{1}}\right)
\end{aligned}
$$

### 1.58 problem 77

1.58.1 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 670
1.58.2 Solving as first order ode lie symmetry calculated ode . . . . . . 672
1.58.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 677
1.58.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 681

Internal problem ID [12475]
Internal file name [OUTPUT/11127_Monday_October_16_2023_09_49_58_PM_34374306/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 77.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "homogeneousTypeD2", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _exact, _rational, [_Abel, `2nd type`, `class C`], _dAlembert]
```

$$
\frac{x}{(y+x)^{2}}+\frac{(2 x+y) y^{\prime}}{(y+x)^{2}}=0
$$

### 1.58.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
\frac{x}{(u(x) x+x)^{2}}+\frac{(2 x+u(x) x)\left(u^{\prime}(x) x+u(x)\right)}{(u(x) x+x)^{2}}=0
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =-\frac{(u+1)^{2}}{x(u+2)}
\end{aligned}
$$

Where $f(x)=-\frac{1}{x}$ and $g(u)=\frac{(u+1)^{2}}{u+2}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{(u+1)^{2}}{u+2}} d u & =-\frac{1}{x} d x \\
\int \frac{1}{\frac{(u+1)^{2}}{u+2}} d u & =\int-\frac{1}{x} d x \\
\ln (u+1)-\frac{1}{u+1} & =-\ln (x)+c_{2}
\end{aligned}
$$

The solution is

$$
\ln (u(x)+1)-\frac{1}{u(x)+1}+\ln (x)-c_{2}=0
$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for $y$ in implicit form

$$
\begin{aligned}
& \ln \left(\frac{y}{x}+1\right)-\frac{1}{\frac{y}{x}+1}+\ln (x)-c_{2}=0 \\
& \ln \left(\frac{y+x}{x}\right)-\frac{x}{y+x}+\ln (x)-c_{2}=0
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\ln \left(\frac{y+x}{x}\right)-\frac{x}{y+x}+\ln (x)-c_{2}=0 \tag{1}
\end{equation*}
$$



Figure 116: Slope field plot
Verification of solutions

$$
\ln \left(\frac{y+x}{x}\right)-\frac{x}{y+x}+\ln (x)-c_{2}=0
$$

Verified OK.

### 1.58.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{x}{2 x+y} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & -\frac{x\left(b_{3}-a_{2}\right)}{2 x+y}-\frac{x^{2} a_{3}}{(2 x+y)^{2}}-\left(-\frac{1}{2 x+y}+\frac{2 x}{(2 x+y)^{2}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\frac{x\left(x b_{2}+y b_{3}+b_{1}\right)}{(2 x+y)^{2}}=0
\end{align*}
$$

Putting the above in normal form gives

$$
\frac{2 x^{2} a_{2}-x^{2} a_{3}+3 x^{2} b_{2}-2 x^{2} b_{3}+2 x y a_{2}+4 x y b_{2}-2 x y b_{3}+y^{2} a_{3}+y^{2} b_{2}-x b_{1}+y a_{1}}{(2 x+y)^{2}}=0
$$

Setting the numerator to zero gives
$2 x^{2} a_{2}-x^{2} a_{3}+3 x^{2} b_{2}-2 x^{2} b_{3}+2 x y a_{2}+4 x y b_{2}-2 x y b_{3}+y^{2} a_{3}+y^{2} b_{2}-x b_{1}+y a_{1}=0$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\{x, y\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& 2 a_{2} v_{1}^{2}+2 a_{2} v_{1} v_{2}-a_{3} v_{1}^{2}+a_{3} v_{2}^{2}+3 b_{2} v_{1}^{2}+4 b_{2} v_{1} v_{2}  \tag{7E}\\
& +b_{2} v_{2}^{2}-2 b_{3} v_{1}^{2}-2 b_{3} v_{1} v_{2}+a_{1} v_{2}-b_{1} v_{1}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}\right\}
$$

Equation (7E) now becomes
$\left(2 a_{2}-a_{3}+3 b_{2}-2 b_{3}\right) v_{1}^{2}+\left(2 a_{2}+4 b_{2}-2 b_{3}\right) v_{1} v_{2}-b_{1} v_{1}+\left(a_{3}+b_{2}\right) v_{2}^{2}+a_{1} v_{2}=0$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
a_{1} & =0 \\
-b_{1} & =0 \\
a_{3}+b_{2} & =0 \\
2 a_{2}+4 b_{2}-2 b_{3} & =0 \\
2 a_{2}-a_{3}+3 b_{2}-2 b_{3} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
& a_{1}=0 \\
& a_{2}=-2 b_{2}+b_{3} \\
& a_{3}=-b_{2} \\
& b_{1}=0 \\
& b_{2}=b_{2} \\
& b_{3}=b_{3}
\end{aligned}
$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
& \xi=x \\
& \eta=y
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =y-\left(-\frac{x}{2 x+y}\right)(x) \\
& =\frac{x^{2}+2 x y+y^{2}}{2 x+y} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{x^{2}+2 x y+y^{2}}{2 x+y}} d y
\end{aligned}
$$

Which results in

$$
S=\ln (y+x)-\frac{x}{y+x}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{x}{2 x+y}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{x}{(y+x)^{2}} \\
S_{y} & =\frac{2 x+y}{(y+x)^{2}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{(y+x) \ln (y+x)-x}{y+x}=c_{1}
$$

Which simplifies to

$$
\frac{(y+x) \ln (y+x)-x}{y+x}=c_{1}
$$

Which gives

$$
y=\mathrm{e}^{\mathrm{LambertW}\left(x \mathrm{e}^{-c_{1}}\right)+c_{1}}-x
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.


Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{\mathrm{LambertW}\left(x \mathrm{e}^{-c_{1}}\right)+c_{1}}-x \tag{1}
\end{equation*}
$$



Figure 117: Slope field plot

Verification of solutions

$$
y=\mathrm{e}^{\mathrm{LambertW}\left(x \mathrm{e}^{-c_{1}}\right)+c_{1}}-x
$$

Verified OK.

### 1.58.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{2 x+y}{(y+x)^{2}}\right) \mathrm{d} y & =\left(-\frac{x}{(y+x)^{2}}\right) \mathrm{d} x \\
\left(\frac{x}{(y+x)^{2}}\right) \mathrm{d} x+\left(\frac{2 x+y}{(y+x)^{2}}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=\frac{x}{(y+x)^{2}} \\
& N(x, y)=\frac{2 x+y}{(y+x)^{2}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\frac{x}{(y+x)^{2}}\right) \\
& =-\frac{2 x}{(y+x)^{3}}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{2 x+y}{(y+x)^{2}}\right) \\
& =-\frac{2 x}{(y+x)^{3}}
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{x}{(y+x)^{2}} \mathrm{~d} x \\
\phi & =\frac{y}{y+x}+\ln (y+x)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{align*}
\frac{\partial \phi}{\partial y} & =\frac{2}{y+x}-\frac{y}{(y+x)^{2}}+f^{\prime}(y)  \tag{4}\\
& =\frac{2 x+y}{(y+x)^{2}}+f^{\prime}(y)
\end{align*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{2 x+y}{(y+x)^{2}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{2 x+y}{(y+x)^{2}}=\frac{2 x+y}{(y+x)^{2}}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{y}{y+x}+\ln (y+x)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{y}{y+x}+\ln (y+x)
$$

The solution becomes

$$
y=\mathrm{e}^{\mathrm{Lambert} \mathrm{~W}\left(x \mathrm{e}^{1-c_{1}}\right)-1+c_{1}}-x
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{\mathrm{LambertW}\left(x \mathrm{e}^{1-c_{1}}\right)-1+c_{1}}-x \tag{1}
\end{equation*}
$$



Figure 118: Slope field plot

## Verification of solutions

$$
y=\mathrm{e}^{\mathrm{LambertW}\left(x \mathrm{e}^{1-c_{1}}\right)-1+c_{1}}-x
$$

Verified OK.

### 1.58.4 Maple step by step solution

Let's solve
$\frac{x}{(y+x)^{2}}+\frac{(2 x+y) y^{\prime}}{(y+x)^{2}}=0$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
Check if ODE is exact
- ODE is exact if the lhs is the total derivative of a $C^{2}$ function $F^{\prime}(x, y)=0$
- Compute derivative of lhs
$F^{\prime}(x, y)+\left(\frac{\partial}{\partial y} F(x, y)\right) y^{\prime}=0$
- Evaluate derivatives
$-\frac{2 x}{(y+x)^{3}}=\frac{2}{(y+x)^{2}}-\frac{2(2 x+y)}{(y+x)^{3}}$
- Simplify
$-\frac{2 x}{(y+x)^{3}}=-\frac{2 x}{(y+x)^{3}}$
- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$
\left[F(x, y)=c_{1}, M(x, y)=F^{\prime}(x, y), N(x, y)=\frac{\partial}{\partial y} F(x, y)\right]
$$

- $\quad$ Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to $x$ $F(x, y)=\int \frac{x}{(y+x)^{2}} d x+f_{1}(y)$
- Evaluate integral

$$
F(x, y)=\frac{y}{y+x}+\ln (y+x)+f_{1}(y)
$$

- $\quad$ Take derivative of $F(x, y)$ with respect to $y$
$N(x, y)=\frac{\partial}{\partial y} F(x, y)$
- Compute derivative

$$
\frac{2 x+y}{(y+x)^{2}}=\frac{2}{y+x}-\frac{y}{(y+x)^{2}}+\frac{d}{d y} f_{1}(y)
$$

- Isolate for $\frac{d}{d y} f_{1}(y)$

$$
\frac{d}{d y} f_{1}(y)=\frac{2 x+y}{(y+x)^{2}}-\frac{2}{y+x}+\frac{y}{(y+x)^{2}}
$$

- $\quad$ Solve for $f_{1}(y)$

$$
f_{1}(y)=0
$$

- $\quad$ Substitute $f_{1}(y)$ into equation for $F(x, y)$
$F(x, y)=\frac{y}{y+x}+\ln (y+x)$
- $\quad$ Substitute $F(x, y)$ into the solution of the ODE
$\frac{y}{y+x}+\ln (y+x)=c_{1}$
- $\quad$ Solve for $y$
$y=\mathrm{e}^{\text {Lambert } W\left(x \mathrm{e}^{1-c_{1}}\right)-1+c_{1}}-x$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`
```

$\checkmark$ Solution by Maple
Time used: 0.031 (sec). Leaf size: 19

```
dsolve(x/(x+y(x))^2+(2*x+y(x))/(x+y(x))^2*diff (y(x), x)=0,y(x), singsol=all)
```

$$
y(x)=-\frac{x\left(\operatorname{LambertW}\left(c_{1} x\right)-1\right)}{\operatorname{LambertW}\left(c_{1} x\right)}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.192 (sec). Leaf size: 33
DSolve $[x /(x+y[x]) \sim 2+(2 * x+y[x]) /(x+y[x]) \sim 2 * y '[x]==0, y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
\text { Solve }\left[\log \left(\frac{y(x)}{x}+1\right)-\frac{1}{\frac{y(x)}{x}+1}=-\log (x)+c_{1}, y(x)\right]
$$

### 1.59 problem 78

> 1.59.1 Solving as homogeneousTypeD2 ode
1.59.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 686
1.59.3 Solving as bernoulli ode . . . . . . . . . . . . . . . . . . . . . . 690
1.59.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 693
1.59.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 697

Internal problem ID [12476]
Internal file name [OUTPUT/11128_Monday_October_16_2023_09_50_00_PM_62518707/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 78.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "bernoulli", "homogeneousTypeD2", "first__order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[[_homogeneous, `class A`], _exact, _rational, _Bernoulli]

$$
\frac{3 y^{2}}{x^{4}}-\frac{2 y y^{\prime}}{x^{3}}=-\frac{1}{x^{2}}
$$

### 1.59.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
\frac{3 u(x)^{2}}{x^{2}}-\frac{2 u(x)\left(u^{\prime}(x) x+u(x)\right)}{x^{2}}=-\frac{1}{x^{2}}
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =\frac{u^{2}+1}{2 u x}
\end{aligned}
$$

Where $f(x)=\frac{1}{2 x}$ and $g(u)=\frac{u^{2}+1}{u}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{u^{2}+1}{u}} d u & =\frac{1}{2 x} d x \\
\int \frac{1}{\frac{u^{2}+1}{u}} d u & =\int \frac{1}{2 x} d x \\
\frac{\ln \left(u^{2}+1\right)}{2} & =\frac{\ln (x)}{2}+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\sqrt{u^{2}+1}=\mathrm{e}^{\frac{\ln (x)}{2}+c_{2}}
$$

Which simplifies to

$$
\sqrt{u^{2}+1}=c_{3} \sqrt{x}
$$

Which simplifies to

$$
\sqrt{u(x)^{2}+1}=c_{3} \sqrt{x} \mathrm{e}^{c_{2}}
$$

The solution is

$$
\sqrt{u(x)^{2}+1}=c_{3} \sqrt{x} \mathrm{e}^{c_{2}}
$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for $y$ in implicit form

$$
\begin{aligned}
\sqrt{\frac{y^{2}}{x^{2}}+1} & =c_{3} \sqrt{x} \mathrm{e}^{c_{2}} \\
\sqrt{\frac{x^{2}+y^{2}}{x^{2}}} & =c_{3} \sqrt{x} \mathrm{e}^{c_{2}}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\sqrt{\frac{x^{2}+y^{2}}{x^{2}}}=c_{3} \sqrt{x} \mathrm{e}^{c_{2}} \tag{1}
\end{equation*}
$$



Figure 119: Slope field plot
Verification of solutions

$$
\sqrt{\frac{x^{2}+y^{2}}{x^{2}}}=c_{3} \sqrt{x} \mathrm{e}^{c_{2}}
$$

Verified OK.

### 1.59.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{x^{2}+3 y^{2}}{2 x y} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type Bernoulli. Therefore we do not need to solve the $\operatorname{PDE}(\mathrm{A})$, and can just use the lookup table shown below to find $\xi, \eta$

Table 107: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | special | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |  |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}$ |
| $a_{1} b_{2}-a_{2} b_{1}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |  |  |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |

The above table shows that

$$
\begin{align*}
\xi(x, y) & =0 \\
\eta(x, y) & =\frac{x^{3}}{y} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{x^{3}}{y}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y^{2}}{2 x^{3}}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{x^{2}+3 y^{2}}{2 x y}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{3 y^{2}}{2 x^{4}} \\
S_{y} & =\frac{y}{x^{3}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{1}{2 x^{2}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{1}{2 R^{2}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\frac{1}{2 R}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{y^{2}}{2 x^{3}}=-\frac{1}{2 x}+c_{1}
$$

Which simplifies to

$$
\frac{y^{2}}{2 x^{3}}=-\frac{1}{2 x}+c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{x^{2}+3 y^{2}}{2 x y}$ |  | $\frac{d S}{d R}=\frac{1}{2 R^{2}}$ |
|  |  |  |
|  |  |  |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+]{ }+\uparrow \xrightarrow{+}$ |
|  |  |  |
|  |  | $\rightarrow$, |
|  | $R=x$ | - |
|  |  |  |
|  | $S=\frac{y^{2}}{2 x^{3}}$ |  |
|  |  | $\rightarrow$ 号新 |
|  |  | $\rightarrow \rightarrow \uparrow$ |
|  |  | $\rightarrow$ - $4+$ |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
\frac{y^{2}}{2 x^{3}}=-\frac{1}{2 x}+c_{1} \tag{1}
\end{equation*}
$$



Figure 120: Slope field plot
Verification of solutions

$$
\frac{y^{2}}{2 x^{3}}=-\frac{1}{2 x}+c_{1}
$$

Verified OK.

### 1.59.3 Solving as bernoulli ode

In canonical form, the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{x^{2}+3 y^{2}}{2 x y}
\end{aligned}
$$

This is a Bernoulli ODE.

$$
\begin{equation*}
y^{\prime}=\frac{3}{2 x} y+\frac{x}{2} \frac{1}{y} \tag{1}
\end{equation*}
$$

The standard Bernoulli ODE has the form

$$
\begin{equation*}
y^{\prime}=f_{0}(x) y+f_{1}(x) y^{n} \tag{2}
\end{equation*}
$$

The first step is to divide the above equation by $y^{n}$ which gives

$$
\begin{equation*}
\frac{y^{\prime}}{y^{n}}=f_{0}(x) y^{1-n}+f_{1}(x) \tag{3}
\end{equation*}
$$

The next step is use the substitution $w=y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$
\begin{aligned}
f_{0}(x) & =\frac{3}{2 x} \\
f_{1}(x) & =\frac{x}{2} \\
n & =-1
\end{aligned}
$$

Dividing both sides of ODE (1) by $y^{n}=\frac{1}{y}$ gives

$$
\begin{equation*}
y^{\prime} y=\frac{3 y^{2}}{2 x}+\frac{x}{2} \tag{4}
\end{equation*}
$$

Let

$$
\begin{align*}
w & =y^{1-n} \\
& =y^{2} \tag{5}
\end{align*}
$$

Taking derivative of equation (5) w.r.t $x$ gives

$$
\begin{equation*}
w^{\prime}=2 y y^{\prime} \tag{6}
\end{equation*}
$$

Substituting equations (5) and (6) into equation (4) gives

$$
\begin{align*}
\frac{w^{\prime}(x)}{2} & =\frac{3 w(x)}{2 x}+\frac{x}{2} \\
w^{\prime} & =\frac{3 w}{x}+x \tag{7}
\end{align*}
$$

The above now is a linear ODE in $w(x)$ which is now solved.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
w^{\prime}(x)+p(x) w(x)=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{3}{x} \\
& q(x)=x
\end{aligned}
$$

Hence the ode is

$$
w^{\prime}(x)-\frac{3 w(x)}{x}=x
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{3}{x} d x} \\
& =\frac{1}{x^{3}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu w) & =(\mu)(x) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{w}{x^{3}}\right) & =\left(\frac{1}{x^{3}}\right)(x) \\
\mathrm{d}\left(\frac{w}{x^{3}}\right) & =\frac{1}{x^{2}} \mathrm{~d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
\frac{w}{x^{3}} & =\int \frac{1}{x^{2}} \mathrm{~d} x \\
\frac{w}{x^{3}} & =-\frac{1}{x}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x^{3}}$ results in

$$
w(x)=c_{1} x^{3}-x^{2}
$$

Replacing $w$ in the above by $y^{2}$ using equation (5) gives the final solution.

$$
y^{2}=c_{1} x^{3}-x^{2}
$$

Solving for $y$ gives

$$
\begin{aligned}
& y(x)=\sqrt{c_{1} x-1} x \\
& y(x)=-\sqrt{c_{1} x-1} x
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\sqrt{c_{1} x-1} x  \tag{1}\\
& y=-\sqrt{c_{1} x-1} x \tag{2}
\end{align*}
$$



Figure 121: Slope field plot
Verification of solutions

$$
y=\sqrt{c_{1} x-1} x
$$

Verified OK.

$$
y=-\sqrt{c_{1} x-1} x
$$

Verified OK.

### 1.59.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(-\frac{2 y}{x^{3}}\right) \mathrm{d} y & =\left(-\frac{1}{x^{2}}-\frac{3 y^{2}}{x^{4}}\right) \mathrm{d} x \\
\left(\frac{1}{x^{2}}+\frac{3 y^{2}}{x^{4}}\right) \mathrm{d} x+\left(-\frac{2 y}{x^{3}}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=\frac{1}{x^{2}}+\frac{3 y^{2}}{x^{4}} \\
& N(x, y)=-\frac{2 y}{x^{3}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(\frac{1}{x^{2}}+\frac{3 y^{2}}{x^{4}}\right) \\
& =\frac{6 y}{x^{4}}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(-\frac{2 y}{x^{3}}\right) \\
& =\frac{6 y}{x^{4}}
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{1}{x^{2}}+\frac{3 y^{2}}{x^{4}} \mathrm{~d} x \\
\phi & =\frac{-x^{2}-y^{2}}{x^{3}}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=-\frac{2 y}{x^{3}}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=-\frac{2 y}{x^{3}}$. Therefore equation (4) becomes

$$
\begin{equation*}
-\frac{2 y}{x^{3}}=-\frac{2 y}{x^{3}}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{-x^{2}-y^{2}}{x^{3}}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{-x^{2}-y^{2}}{x^{3}}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{-x^{2}-y^{2}}{x^{3}}=c_{1} \tag{1}
\end{equation*}
$$



Figure 122: Slope field plot

Verification of solutions

$$
\frac{-x^{2}-y^{2}}{x^{3}}=c_{1}
$$

Verified OK.

### 1.59.5 Maple step by step solution

Let's solve
$\frac{3 y^{2}}{x^{4}}-\frac{2 y y^{\prime}}{x^{3}}=-\frac{1}{x^{2}}$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
Check if ODE is exact
- ODE is exact if the lhs is the total derivative of a $C^{2}$ function
$F^{\prime}(x, y)=0$
- Compute derivative of lhs
$F^{\prime}(x, y)+\left(\frac{\partial}{\partial y} F(x, y)\right) y^{\prime}=0$
- Evaluate derivatives
$\frac{6 y}{x^{4}}=\frac{6 y}{x^{4}}$
- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$
\left[F(x, y)=c_{1}, M(x, y)=F^{\prime}(x, y), N(x, y)=\frac{\partial}{\partial y} F(x, y)\right]
$$

- $\quad$ Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to $x$

$$
F(x, y)=\int\left(\frac{1}{x^{2}}+\frac{3 y^{2}}{x^{4}}\right) d x+f_{1}(y)
$$

- $\quad$ Evaluate integral
$F(x, y)=-\frac{1}{x}-\frac{y^{2}}{x^{3}}+f_{1}(y)$
- $\quad$ Take derivative of $F(x, y)$ with respect to $y$
$N(x, y)=\frac{\partial}{\partial y} F(x, y)$
- Compute derivative
$-\frac{2 y}{x^{3}}=-\frac{2 y}{x^{3}}+\frac{d}{d y} f_{1}(y)$
- $\quad$ Isolate for $\frac{d}{d y} f_{1}(y)$
$\frac{d}{d y} f_{1}(y)=0$
- $\quad$ Solve for $f_{1}(y)$

$$
f_{1}(y)=0
$$

- $\quad$ Substitute $f_{1}(y)$ into equation for $F(x, y)$

$$
F(x, y)=-\frac{1}{x}-\frac{y^{2}}{x^{3}}
$$

- $\quad$ Substitute $F(x, y)$ into the solution of the ODE

$$
-\frac{1}{x}-\frac{y^{2}}{x^{3}}=c_{1}
$$

- $\quad$ Solve for $y$

$$
\left\{y=\sqrt{-c_{1} x-1} x, y=-\sqrt{-c_{1} x-1} x\right\}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 26

```
dsolve(1/x^2+ 3*y(x)^2/x^4=2*y(x)/x^3*diff(y(x),x),y(x), singsol=all)
```

$$
\begin{aligned}
& y(x)=\sqrt{c_{1} x-1} x \\
& y(x)=-\sqrt{c_{1} x-1} x
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.465 (sec). Leaf size: 34
DSolve $\left[1 / x^{\wedge} 2+3 * y[x] \wedge 2 / x^{\wedge} 4==2 * y[x] / x^{\wedge}-3 * y \cdot[x], y[x], x\right.$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow-x \sqrt{-1+c_{1} x} \\
& y(x) \rightarrow x \sqrt{-1+c_{1} x}
\end{aligned}
$$

### 1.60 problem 79

1.60.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 699
1.60.2 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 701
1.60.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 702
1.60.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 706
1.60.5 Solving as riccati ode . . . . . . . . . . . . . . . . . . . . . . . . 710
1.60.6 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 712

Internal problem ID [12477]
Internal file name [OUTPUT/11129_Monday_October_16_2023_09_50_03_PM_13121631/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 79.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]

$$
\frac{x^{2} y^{\prime}}{(-y+x)^{2}}-\frac{y^{2}}{(-y+x)^{2}}=0
$$

### 1.60.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{y^{2}}{x^{2}}
\end{aligned}
$$

Where $f(x)=\frac{1}{x^{2}}$ and $g(y)=y^{2}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{y^{2}} d y & =\frac{1}{x^{2}} d x \\
\int \frac{1}{y^{2}} d y & =\int \frac{1}{x^{2}} d x \\
-\frac{1}{y} & =-\frac{1}{x}+c_{1}
\end{aligned}
$$

Which results in

$$
y=-\frac{x}{c_{1} x-1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{x}{c_{1} x-1} \tag{1}
\end{equation*}
$$



Figure 123: Slope field plot

## Verification of solutions

$$
y=-\frac{x}{c_{1} x-1}
$$

Verified OK.

### 1.60.2 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
\frac{x^{2}\left(u^{\prime}(x) x+u(x)\right)}{(-u(x) x+x)^{2}}-\frac{u(x)^{2} x^{2}}{(-u(x) x+x)^{2}}=0
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =\frac{u(u-1)}{x}
\end{aligned}
$$

Where $f(x)=\frac{1}{x}$ and $g(u)=u(u-1)$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{u(u-1)} d u & =\frac{1}{x} d x \\
\int \frac{1}{u(u-1)} d u & =\int \frac{1}{x} d x \\
\ln (u-1)-\ln (u) & =\ln (x)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\mathrm{e}^{\ln (u-1)-\ln (u)}=\mathrm{e}^{\ln (x)+c_{2}}
$$

Which simplifies to

$$
\frac{u-1}{u}=c_{3} x
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =x u \\
& =-\frac{x}{c_{3} x-1}
\end{aligned}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{x}{c_{3} x-1} \tag{1}
\end{equation*}
$$



Figure 124: Slope field plot

## Verification of solutions

$$
y=-\frac{x}{c_{3} x-1}
$$

Verified OK.

### 1.60.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{y^{2}}{x^{2}} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type separable. Therefore we do not need to solve the $\operatorname{PDE}(\mathrm{A})$, and can just use the lookup table shown below to find $\xi, \eta$

Table 110: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :---: | :---: | :---: | :---: |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order special form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=x^{2} \\
& \eta(x, y)=0 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\eta=0$ then in this special case

$$
R=y
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\xi} d x \\
& =\int \frac{1}{x^{2}} d x
\end{aligned}
$$

Which results in

$$
S=-\frac{1}{x}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{y^{2}}{x^{2}}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =0 \\
R_{y} & =1 \\
S_{x} & =\frac{1}{x^{2}} \\
S_{y} & =0
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=\frac{1}{y^{2}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=\frac{1}{R^{2}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\frac{1}{R}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
-\frac{1}{x}=-\frac{1}{y}+c_{1}
$$

Which simplifies to

$$
-\frac{1}{x}=-\frac{1}{y}+c_{1}
$$

Which gives

$$
y=\frac{x}{c_{1} x+1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{y^{2}}{x^{2}}$ |  | $\frac{d S}{d R}=\frac{1}{R^{2}}$ |
|  |  | $\rightarrow \rightarrow \rightarrow \rightarrow$ 性 $\uparrow \rightarrow \rightarrow$ |
|  |  | $\rightarrow \rightarrow \infty 14$ |
|  |  | $\rightarrow{ }^{\prime \prime} 4^{4} \uparrow$ |
|  |  |  |
|  |  |  |
| $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-\infty]{ }+\uparrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ | $R=y$ | $\rightarrow>+\uparrow+\infty$ |
| $\xrightarrow{\rightarrow \rightarrow \rightarrow-\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow 0}$, |  |  |
|  | $S=-\frac{1}{x}$ | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-\infty]{ }$ |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow-\infty]{ }$ |
|  |  | $\xrightarrow[\rightarrow \rightarrow \rightarrow-\infty]{ } \uparrow+$ |
|  |  |  |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{x}{c_{1} x+1} \tag{1}
\end{equation*}
$$



Figure 125: Slope field plot
Verification of solutions

$$
y=\frac{x}{c_{1} x+1}
$$

Verified OK.

### 1.60.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{y^{2}}\right) \mathrm{d} y & =\left(\frac{1}{x^{2}}\right) \mathrm{d} x \\
\left(-\frac{1}{x^{2}}\right) \mathrm{d} x+\left(\frac{1}{y^{2}}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=-\frac{1}{x^{2}} \\
& N(x, y)=\frac{1}{y^{2}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{x^{2}}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{y^{2}}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{1}{x^{2}} \mathrm{~d} x \\
\phi & =\frac{1}{x}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{y^{2}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{y^{2}}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{y^{2}}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{y^{2}}\right) \mathrm{d} y \\
f(y) & =-\frac{1}{y}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{1}{x}-\frac{1}{y}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{1}{x}-\frac{1}{y}
$$

The solution becomes

$$
y=-\frac{x}{c_{1} x-1}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\frac{x}{c_{1} x-1} \tag{1}
\end{equation*}
$$



Figure 126: Slope field plot

## Verification of solutions

$$
y=-\frac{x}{c_{1} x-1}
$$

Verified OK.

### 1.60.5 Solving as riccati ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{y^{2}}{x^{2}}
\end{aligned}
$$

This is a Riccati ODE. Comparing the ODE to solve

$$
y^{\prime}=\frac{y^{2}}{x^{2}}
$$

With Riccati ODE standard form

$$
y^{\prime}=f_{0}(x)+f_{1}(x) y+f_{2}(x) y^{2}
$$

Shows that $f_{0}(x)=0, f_{1}(x)=0$ and $f_{2}(x)=\frac{1}{x^{2}}$. Let

$$
\begin{align*}
y & =\frac{-u^{\prime}}{f_{2} u} \\
& =\frac{-u^{\prime}}{\frac{u}{x^{2}}} \tag{1}
\end{align*}
$$

Using the above substitution in the given ODE results (after some simplification)in a second order ODE to solve for $u(x)$ which is

$$
\begin{equation*}
f_{2} u^{\prime \prime}(x)-\left(f_{2}^{\prime}+f_{1} f_{2}\right) u^{\prime}(x)+f_{2}^{2} f_{0} u(x)=0 \tag{2}
\end{equation*}
$$

But

$$
\begin{aligned}
f_{2}^{\prime} & =-\frac{2}{x^{3}} \\
f_{1} f_{2} & =0 \\
f_{2}^{2} f_{0} & =0
\end{aligned}
$$

Substituting the above terms back in equation (2) gives

$$
\frac{u^{\prime \prime}(x)}{x^{2}}+\frac{2 u^{\prime}(x)}{x^{3}}=0
$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$
u(x)=c_{1}+\frac{c_{2}}{x}
$$

The above shows that

$$
u^{\prime}(x)=-\frac{c_{2}}{x^{2}}
$$

Using the above in (1) gives the solution

$$
y=\frac{c_{2}}{c_{1}+\frac{c_{2}}{x}}
$$

Dividing both numerator and denominator by $c_{1}$ gives, after renaming the constant $\frac{c_{2}}{c_{1}}=c_{3}$ the following solution

$$
y=\frac{1}{c_{3}+\frac{1}{x}}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\frac{1}{c_{3}+\frac{1}{x}} \tag{1}
\end{equation*}
$$



Figure 127: Slope field plot

Verification of solutions

$$
y=\frac{1}{c_{3}+\frac{1}{x}}
$$

Verified OK.

### 1.60.6 Maple step by step solution

Let's solve

$$
\frac{x^{2} y^{\prime}}{(-y+x)^{2}}-\frac{y^{2}}{(-y+x)^{2}}=0
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Integrate both sides with respect to $x$

$$
\int\left(\frac{x^{2} y^{\prime}}{(-y+x)^{2}}-\frac{y^{2}}{(-y+x)^{2}}\right) d x=\int 0 d x+c_{1}
$$

- Evaluate integral

$$
\frac{x^{2}}{-y+x}-x=c_{1}
$$

- $\quad$ Solve for $y$

$$
y=\frac{c_{1} x}{x+c_{1}}
$$

Maple trace

- Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 13

```
dsolve(x^2/(x-y(x))^2*diff(y(x),x)- y(x)^2/(x-y(x))^2=0,y(x), singsol=all)
```

$$
y(x)=\frac{x}{c_{1} x+1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.193 (sec). Leaf size: 21
DSolve $\left[x^{\wedge} 2 /(x-y[x])^{\wedge} 2 * y^{\prime}[x]-y[x] \sim 2 /(x-y[x]) \wedge 2==0, y[x], x\right.$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow \frac{x}{1-c_{1} x} \\
& y(x) \rightarrow 0
\end{aligned}
$$

### 1.61 problem 80

1.61.1 Solving as first order ode lie symmetry calculated ode . . . . . . 714
1.61.2 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 720
1.61.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 724

Internal problem ID [12478]
Internal file name [OUTPUT/11130_Monday_October_16_2023_09_50_05_PM_63728957/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 80.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type
[[_1st_order, _with_linear_symmetries], _exact, _rational]

$$
y y^{\prime}-\frac{y}{x^{2}+y^{2}}+\frac{x y^{\prime}}{x^{2}+y^{2}}=-x
$$

### 1.61.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =-\frac{x^{3}+x y^{2}-y}{x^{2} y+y^{3}+x} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is not in the lookup table. To determine $\xi, \eta$ then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
$$

Where the unknown coefficients are

$$
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
$$

Substituting equations (1E,2E) and $\omega$ into (A) gives

$$
\begin{align*}
b_{2} & -\frac{\left(x^{3}+x y^{2}-y\right)\left(b_{3}-a_{2}\right)}{x^{2} y+y^{3}+x}-\frac{\left(x^{3}+x y^{2}-y\right)^{2} a_{3}}{\left(x^{2} y+y^{3}+x\right)^{2}} \\
& -\left(-\frac{3 x^{2}+y^{2}}{x^{2} y+y^{3}+x}+\frac{\left(x^{3}+x y^{2}-y\right)(2 x y+1)}{\left(x^{2} y+y^{3}+x\right)^{2}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\left(-\frac{2 x y-1}{x^{2} y+y^{3}+x}+\frac{\left(x^{3}+x y^{2}-y\right)\left(x^{2}+3 y^{2}\right)}{\left(x^{2} y+y^{3}+x\right)^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
$$

Putting the above in normal form gives

$$
\begin{aligned}
& -\frac{x^{6} a_{3}+x^{6} b_{2}-2 x^{5} y a_{2}+2 x^{5} y b_{3}+x^{4} y^{2} a_{3}+x^{4} y^{2} b_{2}-4 x^{3} y^{3} a_{2}+4 x^{3} y^{3} b_{3}-x^{2} y^{4} a_{3}-x^{2} y^{4} b_{2}-2 x y^{5} a_{2}+}{}=0
\end{aligned}
$$

Setting the numerator to zero gives

$$
\begin{align*}
& -x^{6} a_{3}-x^{6} b_{2}+2 x^{5} y a_{2}-2 x^{5} y b_{3}-x^{4} y^{2} a_{3}-x^{4} y^{2} b_{2}+4 x^{3} y^{3} a_{2} \\
& \quad-4 x^{3} y^{3} b_{3}+x^{2} y^{4} a_{3}+x^{2} y^{4} b_{2}+2 x y^{5} a_{2}-2 x y^{5} b_{3}+y^{6} a_{3}+y^{6} b_{2}  \tag{6E}\\
& -x^{5} b_{1}+x^{4} y a_{1}-2 x^{3} y^{2} b_{1}+2 x^{2} y^{3} a_{1}-x y^{4} b_{1}+y^{5} a_{1}+3 x^{4} a_{2} \\
& -x^{4} b_{3}+4 x^{3} y a_{3}+4 x^{3} y b_{2}+2 x^{2} y^{2} a_{2}+2 x^{2} y^{2} b_{3}+4 x y^{3} a_{3}+4 x y^{3} b_{2} \\
& -y^{4} a_{2}+3 y^{4} b_{3}+2 x^{3} a_{1}+2 x^{2} y b_{1}+2 x y^{2} a_{1}+2 y^{3} b_{1}-x b_{1}+y a_{1}=0
\end{align*}
$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$
\{x, y\}
$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$
\left\{x=v_{1}, y=v_{2}\right\}
$$

The above PDE (6E) now becomes

$$
\begin{align*}
& 2 a_{2} v_{1}^{5} v_{2}+4 a_{2} v_{1}^{3} v_{2}^{3}+2 a_{2} v_{1} v_{2}^{5}-a_{3} v_{1}^{6}-a_{3} v_{1}^{4} v_{2}^{2}+a_{3} v_{1}^{2} v_{2}^{4}+a_{3} v_{2}^{6}-b_{2} v_{1}^{6} \\
& \quad-b_{2} v_{1}^{4} v_{2}^{2}+b_{2} v_{1}^{2} v_{2}^{4}+b_{2} v_{2}^{6}-2 b_{3} v_{1}^{5} v_{2}-4 b_{3} v_{1}^{3} v_{2}^{3}-2 b_{3} v_{1} v_{2}^{5}+a_{1} v_{1}^{4} v_{2}  \tag{7E}\\
& \quad+2 a_{1} v_{1}^{2} v_{2}^{3}+a_{1} v_{2}^{5}-b_{1} v_{1}^{5}-2 b_{1} v_{1}^{3} v_{2}^{2}-b_{1} v_{1} v_{2}^{4}+3 a_{2} v_{1}^{4}+2 a_{2} v_{1}^{2} v_{2}^{2} \\
& \quad-a_{2} v_{2}^{4}+4 a_{3} v_{1}^{3} v_{2}+4 a_{3} v_{1} v_{2}^{3}+4 b_{2} v_{1}^{3} v_{2}+4 b_{2} v_{1} v_{2}^{3}-b_{3} v_{1}^{4}+2 b_{3} v_{1}^{2} v_{2}^{2} \\
& \quad+3 b_{3} v_{2}^{4}+2 a_{1} v_{1}^{3}+2 a_{1} v_{1} v_{2}^{2}+2 b_{1} v_{1}^{2} v_{2}+2 b_{1} v_{2}^{3}+a_{1} v_{2}-b_{1} v_{1}=0
\end{align*}
$$

Collecting the above on the terms $v_{i}$ introduced, and these are

$$
\left\{v_{1}, v_{2}\right\}
$$

Equation (7E) now becomes

$$
\begin{align*}
& \left(-a_{3}-b_{2}\right) v_{1}^{6}+\left(2 a_{2}-2 b_{3}\right) v_{1}^{5} v_{2}-b_{1} v_{1}^{5}+\left(-a_{3}-b_{2}\right) v_{1}^{4} v_{2}^{2}+a_{1} v_{1}^{4} v_{2} \\
& \quad+\left(3 a_{2}-b_{3}\right) v_{1}^{4}+\left(4 a_{2}-4 b_{3}\right) v_{1}^{3} v_{2}^{3}-2 b_{1} v_{1}^{3} v_{2}^{2}+\left(4 a_{3}+4 b_{2}\right) v_{1}^{3} v_{2}  \tag{8E}\\
& \quad+2 a_{1} v_{1}^{3}+\left(a_{3}+b_{2}\right) v_{1}^{2} v_{2}^{4}+2 a_{1} v_{1}^{2} v_{2}^{3}+\left(2 a_{2}+2 b_{3}\right) v_{1}^{2} v_{2}^{2}+2 b_{1} v_{1}^{2} v_{2} \\
& \quad+\left(2 a_{2}-2 b_{3}\right) v_{1} v_{2}^{5}-b_{1} v_{1} v_{2}^{4}+\left(4 a_{3}+4 b_{2}\right) v_{1} v_{2}^{3}+2 a_{1} v_{1} v_{2}^{2}-b_{1} v_{1} \\
& \quad+\left(a_{3}+b_{2}\right) v_{2}^{6}+a_{1} v_{2}^{5}+\left(-a_{2}+3 b_{3}\right) v_{2}^{4}+2 b_{1} v_{2}^{3}+a_{1} v_{2}=0
\end{align*}
$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$
\begin{aligned}
a_{1} & =0 \\
2 a_{1} & =0 \\
-2 b_{1} & =0 \\
-b_{1} & =0 \\
2 b_{1} & =0 \\
-a_{2}+3 b_{3} & =0 \\
2 a_{2}-2 b_{3} & =0 \\
2 a_{2}+2 b_{3} & =0 \\
3 a_{2}-b_{3} & =0 \\
4 a_{2}-4 b_{3} & =0 \\
-a_{3}-b_{2} & =0 \\
a_{3}+b_{2} & =0 \\
4 a_{3}+4 b_{2} & =0
\end{aligned}
$$

Solving the above equations for the unknowns gives

$$
\begin{aligned}
& a_{1}=0 \\
& a_{2}=0 \\
& a_{3}=-b_{2} \\
& b_{1}=0 \\
& b_{2}=b_{2} \\
& b_{3}=0
\end{aligned}
$$

Substituting the above solution in the anstaz (1E, 2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$
\begin{aligned}
\xi & =-y \\
\eta & =x
\end{aligned}
$$

Shifting is now applied to make $\xi=0$ in order to simplify the rest of the computation

$$
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =x-\left(-\frac{x^{3}+x y^{2}-y}{x^{2} y+y^{3}+x}\right)(-y) \\
& =\frac{x^{2}+y^{2}}{x^{2} y+y^{3}+x} \\
\xi & =0
\end{aligned}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates $\operatorname{map}(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{x^{2}+y^{2}}{x^{2} y+y^{3}+x}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y^{2}}{2}+\arctan \left(\frac{y}{x}\right)
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=-\frac{x^{3}+x y^{2}-y}{x^{2} y+y^{3}+x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{y}{x^{2}+y^{2}} \\
S_{y} & =y+\frac{x}{x^{2}+y^{2}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-x \tag{2A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-R
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=-\frac{R^{2}}{2}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{y^{2}}{2}+\arctan \left(\frac{y}{x}\right)=-\frac{x^{2}}{2}+c_{1}
$$

Which simplifies to

$$
\frac{y^{2}}{2}+\arctan \left(\frac{y}{x}\right)=-\frac{x^{2}}{2}+c_{1}
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=-\frac{x^{3}+x y^{2}-y}{x^{2} y+y^{3}+x}$ |  | $\frac{d S}{d R}=-R$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  | $R=x$ |  |
|  |  |  |
|  | $S=\frac{y}{2}+\arctan \left(\frac{y}{x}\right)$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |
| $\rightarrow \rightarrow \rightarrow \rightarrow \infty$ |  |  |

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{y^{2}}{2}+\arctan \left(\frac{y}{x}\right)=-\frac{x^{2}}{2}+c_{1} \tag{1}
\end{equation*}
$$



Figure 128: Slope field plot

Verification of solutions

$$
\frac{y^{2}}{2}+\arctan \left(\frac{y}{x}\right)=-\frac{x^{2}}{2}+c_{1}
$$

Verified OK.

### 1.61.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(y+\frac{x}{x^{2}+y^{2}}\right) \mathrm{d} y & =\left(-x+\frac{y}{x^{2}+y^{2}}\right) \mathrm{d} x \\
\left(x-\frac{y}{x^{2}+y^{2}}\right) \mathrm{d} x+\left(y+\frac{x}{x^{2}+y^{2}}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=x-\frac{y}{x^{2}+y^{2}} \\
& N(x, y)=y+\frac{x}{x^{2}+y^{2}}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(x-\frac{y}{x^{2}+y^{2}}\right) \\
& =\frac{-x^{2}+y^{2}}{\left(x^{2}+y^{2}\right)^{2}}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(y+\frac{x}{x^{2}+y^{2}}\right) \\
& =\frac{-x^{2}+y^{2}}{\left(x^{2}+y^{2}\right)^{2}}
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int x-\frac{y}{x^{2}+y^{2}} \mathrm{~d} x \\
\phi & =\frac{x^{2}}{2}-\arctan \left(\frac{x}{y}\right)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{align*}
\frac{\partial \phi}{\partial y} & =\frac{x}{y^{2}\left(\frac{x^{2}}{y^{2}}+1\right)}+f^{\prime}(y)  \tag{4}\\
& =\frac{x}{x^{2}+y^{2}}+f^{\prime}(y)
\end{align*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=y+\frac{x}{x^{2}+y^{2}}$. Therefore equation (4) becomes

$$
\begin{equation*}
y+\frac{x}{x^{2}+y^{2}}=\frac{x}{x^{2}+y^{2}}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=y
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int(y) \mathrm{d} y \\
f(y) & =\frac{y^{2}}{2}+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{x^{2}}{2}-\arctan \left(\frac{x}{y}\right)+\frac{y^{2}}{2}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{x^{2}}{2}-\arctan \left(\frac{x}{y}\right)+\frac{y^{2}}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
\frac{x^{2}}{2}-\arctan \left(\frac{x}{y}\right)+\frac{y^{2}}{2}=c_{1} \tag{1}
\end{equation*}
$$



Figure 129: Slope field plot

Verification of solutions

$$
\frac{x^{2}}{2}-\arctan \left(\frac{x}{y}\right)+\frac{y^{2}}{2}=c_{1}
$$

Verified OK.

### 1.61.3 Maple step by step solution

Let's solve

$$
y y^{\prime}-\frac{y}{x^{2}+y^{2}}+\frac{x y^{\prime}}{x^{2}+y^{2}}=-x
$$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
Check if ODE is exact
- ODE is exact if the lhs is the total derivative of a $C^{2}$ function

$$
F^{\prime}(x, y)=0
$$

- Compute derivative of lhs

$$
F^{\prime}(x, y)+\left(\frac{\partial}{\partial y} F(x, y)\right) y^{\prime}=0
$$

- Evaluate derivatives

$$
\frac{2 y^{2}}{\left(x^{2}+y^{2}\right)^{2}}-\frac{1}{x^{2}+y^{2}}=\frac{1}{x^{2}+y^{2}}-\frac{2 x^{2}}{\left(x^{2}+y^{2}\right)^{2}}
$$

- Simplify

$$
\frac{-x^{2}+y^{2}}{\left(x^{2}+y^{2}\right)^{2}}=\frac{-x^{2}+y^{2}}{\left(x^{2}+y^{2}\right)^{2}}
$$

- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$
\left[F(x, y)=c_{1}, M(x, y)=F^{\prime}(x, y), N(x, y)=\frac{\partial}{\partial y} F(x, y)\right]
$$

- $\quad$ Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to $x$

$$
F(x, y)=\int\left(x-\frac{y}{x^{2}+y^{2}}\right) d x+f_{1}(y)
$$

- Evaluate integral

$$
F(x, y)=\frac{x^{2}}{2}-\arctan \left(\frac{x}{y}\right)+f_{1}(y)
$$

- $\quad$ Take derivative of $F(x, y)$ with respect to $y$

$$
N(x, y)=\frac{\partial}{\partial y} F(x, y)
$$

- Compute derivative

$$
y+\frac{x}{x^{2}+y^{2}}=\frac{x}{y^{2}\left(\frac{x^{2}}{y^{2}}+1\right)}+\frac{d}{d y} f_{1}(y)
$$

- $\quad$ Isolate for $\frac{d}{d y} f_{1}(y)$
$\frac{d}{d y} f_{1}(y)=y+\frac{x}{x^{2}+y^{2}}-\frac{x}{y^{2}\left(\frac{x^{2}}{y^{2}}+1\right)}$
- $\quad$ Solve for $f_{1}(y)$
$f_{1}(y)=\frac{y^{2}}{2}$
- $\quad$ Substitute $f_{1}(y)$ into equation for $F(x, y)$

$$
F(x, y)=\frac{x^{2}}{2}-\arctan \left(\frac{x}{y}\right)+\frac{y^{2}}{2}
$$

- $\quad$ Substitute $F(x, y)$ into the solution of the ODE

$$
\frac{x^{2}}{2}-\arctan \left(\frac{x}{y}\right)+\frac{y^{2}}{2}=c_{1}
$$

- $\quad$ Solve for $y$

$$
y=-\frac{x}{\tan \left(\operatorname{RootOf}\left(-x^{2} \tan \left(\_Z\right)^{2}+2 c_{1} \tan \left(\_Z\right)^{2}-2 \_Z \tan \left(\_Z\right)^{2}-x^{2}\right)\right)}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
differential order: 1; found: 1 linear symmetries. Trying reduction of order
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful`
```

$\checkmark$ Solution by Maple
Time used: 0.172 (sec). Leaf size: 26

```
dsolve(x+y(x)*\operatorname{diff}(y(x),x)= y(x)/(x^2+y(x)^2)- x/( (x^2+y(x)^2)*diff (y(x),x),y(x), singsol=all
```

$$
y(x)=\cot \left(\operatorname{RootOf}\left(2 \sin \left(\_Z\right)^{2} c_{1}-2 \_Z \sin \left(\_Z\right)^{2}+x^{2}\right)\right) x
$$

Solution by Mathematica
Time used: 0.18 (sec). Leaf size: 31
DSolve $\left[x+y[x] * y{ }^{\prime}[x]==y[x] /\left(x^{\wedge} 2+y[x] \sim 2\right)-x /\left(x^{\wedge} 2+y[x] \sim 2\right) * y '[x], y[x], x\right.$, IncludeSingularSolution

$$
\text { Solve }\left[-\arctan \left(\frac{x}{y(x)}\right)+\frac{x^{2}}{2}+\frac{y(x)^{2}}{2}=c_{1}, y(x)\right]
$$

### 1.62 problem 89

$$
\text { 1.62.1 Solving as dAlembert ode . . . . . . . . . . . . . . . . . . . . . } 727
$$

Internal problem ID [12479]
Internal file name [OUTPUT/11131_Monday_October_16_2023_09_50_06_PM_60401402/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 89.
ODE order: 1.
ODE degree: 2.

The type(s) of ODE detected by this program : "dAlembert"
Maple gives the following as the ode type

```
[[_1st_order, _with_linear_symmetries], _dAlembert]
```

$$
y-2 y^{\prime} x-y^{\prime 2}=0
$$

### 1.62.1 Solving as dAlembert ode

Let $p=y^{\prime}$ the ode becomes

$$
-p^{2}-2 p x+y=0
$$

Solving for $y$ from the above results in

$$
\begin{equation*}
y=p^{2}+2 p x \tag{1A}
\end{equation*}
$$

This has the form

$$
\begin{equation*}
y=x f(p)+g(p) \tag{}
\end{equation*}
$$

Where $f, g$ are functions of $p=y^{\prime}(x)$. The above ode is dAlembert ode which is now solved. Taking derivative of $\left({ }^{*}\right)$ w.r.t. $x$ gives

$$
\begin{align*}
p & =f+\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \\
p-f & =\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \tag{2}
\end{align*}
$$

Comparing the form $y=x f+g$ to (1A) shows that

$$
\begin{aligned}
& f=2 p \\
& g=p^{2}
\end{aligned}
$$

Hence (2) becomes

$$
\begin{equation*}
-p=(2 x+2 p) p^{\prime}(x) \tag{2~A}
\end{equation*}
$$

The singular solution is found by setting $\frac{d p}{d x}=0$ in the above which gives

$$
-p=0
$$

Solving for $p$ from the above gives

$$
p=0
$$

Substituting these in (1A) gives

$$
y=0
$$

The general solution is found when $\frac{\mathrm{d} p}{\mathrm{~d} x} \neq 0$. From eq. (2A). This results in

$$
\begin{equation*}
p^{\prime}(x)=-\frac{p(x)}{2 x+2 p(x)} \tag{3}
\end{equation*}
$$

This ODE is now solved for $p(x)$.
Inverting the above ode gives

$$
\begin{equation*}
\frac{d}{d p} x(p)=-\frac{2 x(p)+2 p}{p} \tag{4}
\end{equation*}
$$

This ODE is now solved for $x(p)$.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
\frac{d}{d p} x(p)+p(p) x(p)=q(p)
$$

Where here

$$
\begin{aligned}
& p(p)=\frac{2}{p} \\
& q(p)=-2
\end{aligned}
$$

Hence the ode is

$$
\frac{d}{d p} x(p)+\frac{2 x(p)}{p}=-2
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \frac{2}{p} d p} \\
& =p^{2}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} p}(\mu x) & =(\mu)(-2) \\
\frac{\mathrm{d}}{\mathrm{~d} p}\left(p^{2} x\right) & =\left(p^{2}\right)(-2) \\
\mathrm{d}\left(p^{2} x\right) & =\left(-2 p^{2}\right) \mathrm{d} p
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
p^{2} x & =\int-2 p^{2} \mathrm{~d} p \\
p^{2} x & =-\frac{2 p^{3}}{3}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=p^{2}$ results in

$$
x(p)=-\frac{2 p}{3}+\frac{c_{1}}{p^{2}}
$$

Now we need to eliminate $p$ between the above and (1A). One way to do this is by solving (1) for $p$. This results in

$$
\begin{aligned}
& p=-x+\sqrt{x^{2}+y} \\
& p=-x-\sqrt{x^{2}+y}
\end{aligned}
$$

Substituting the above in the solution for $x$ found above gives

$$
\begin{aligned}
& x=\frac{\left(-8 x^{2}-2 y\right) \sqrt{x^{2}+y}+8 x^{3}+6 y x+3 c_{1}}{3\left(x-\sqrt{x^{2}+y}\right)^{2}} \\
& x=\frac{\left(8 x^{2}+2 y\right) \sqrt{x^{2}+y}+8 x^{3}+6 y x+3 c_{1}}{3\left(x+\sqrt{x^{2}+y}\right)^{2}}
\end{aligned}
$$

## Summary

The solution(s) found are the following

$$
\begin{align*}
& y=0  \tag{1}\\
& x=\frac{\left(-8 x^{2}-2 y\right) \sqrt{x^{2}+y}+8 x^{3}+6 y x+3 c_{1}}{3\left(x-\sqrt{x^{2}+y}\right)^{2}}  \tag{2}\\
& x=\frac{\left(8 x^{2}+2 y\right) \sqrt{x^{2}+y}+8 x^{3}+6 y x+3 c_{1}}{3\left(x+\sqrt{x^{2}+y}\right)^{2}} \tag{3}
\end{align*}
$$

Verification of solutions

$$
y=0
$$

Verified OK.

$$
x=\frac{\left(-8 x^{2}-2 y\right) \sqrt{x^{2}+y}+8 x^{3}+6 y x+3 c_{1}}{3\left(x-\sqrt{x^{2}+y}\right)^{2}}
$$

Verified OK.

$$
x=\frac{\left(8 x^{2}+2 y\right) \sqrt{x^{2}+y}+8 x^{3}+6 y x+3 c_{1}}{3\left(x+\sqrt{x^{2}+y}\right)^{2}}
$$

Verified OK.
Maple trace

```
Methods for first order ODEs:
-> Solving 1st order ODE of high degree, 1st attempt
trying 1st order WeierstrassP solution for high degree ODE
trying 1st order WeierstrassPPrime solution for high degree ODE
trying 1st order JacobiSN solution for high degree ODE
trying 1st order ODE linearizable_by_differentiation
trying differential order: 1; missing variables
trying dAlembert
<- dAlembert successful`
```

$\checkmark$ Solution by Maple
Time used: 0.078 (sec). Leaf size: 650

```
dsolve(y(x)=2*x*diff(y(x),x)+diff(y(x),x)^2,y(x), singsol=all)
```

$$
=\frac{\left(x^{2}-x\left(-x^{3}+2 \sqrt{3} \sqrt{-c_{1}\left(x^{3}-3 c_{1}\right)}+6 c_{1}\right)^{\frac{1}{3}}+\left(-x^{3}+2 \sqrt{3} \sqrt{-c_{1}\left(x^{3}-3 c_{1}\right)}+6 c_{1}\right)^{\frac{2}{3}}\right)\left(x^{2}+3 x(-\right.}{4\left(-x^{3}+2 \sqrt{3} \sqrt{-c_{1}\left(x^{3}-3 c_{1}\right)}\right.}
$$

$y(x)$
$=\underline{\left(i \sqrt{3}\left(-x^{3}+2 \sqrt{3} \sqrt{-c_{1}\left(x^{3}-3 c_{1}\right)}+6 c_{1}\right)^{\frac{2}{3}}-i \sqrt{3} x^{2}+\left(-x^{3}+2 \sqrt{3} \sqrt{-c_{1}\left(x^{3}-3 c_{1}\right)}+6 c_{1}\right)^{\frac{2}{3}}+2 x( \right.}$
$y(x)$
$=\underline{\left(i \sqrt{3} x^{2}-i \sqrt{3}\left(-x^{3}+2 \sqrt{3} \sqrt{-c_{1}\left(x^{3}-3 c_{1}\right)}+6 c_{1}\right)^{\frac{2}{3}}+x^{2}+2 x\left(-x^{3}+2 \sqrt{3} \sqrt{-c_{1}\left(x^{3}-3 c_{1}\right)}+6 c_{1}\right)\right.}$

## Solution by Mathematica

Time used: 60.162 (sec). Leaf size: 931
DSolve[y[x]==2*x*y'[x]+(y'[x])~2,y[x],x,IncludeSingularSolutions -> True]

$$
y(x) \rightarrow \frac{1}{72}\left(-18 x^{2}+\frac{9(1+i \sqrt{3}) x\left(-x^{3}+8 e^{3 c_{1}}\right)}{\sqrt[3]{-x^{6}-20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}}\right.
$$

$$
\left.+9 i(\sqrt{3}+i) \sqrt[3]{-x^{6}-20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}\right)
$$

$$
y(x) \rightarrow \frac{1}{72}\left(-18 x^{2}+\frac{9 i(\sqrt{3}+i) x\left(x^{3}-8 e^{3 c_{1}}\right)}{\sqrt[3]{-x^{6}-20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}}\left(x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}\right.
$$

$$
\left.-9(1+i \sqrt{3}) \sqrt[3]{-x^{6}-20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}\right)
$$

$$
\begin{aligned}
& y(x) \rightarrow \frac{1}{4}\left(-x^{2}+\frac{x\left(x^{3}+8 e^{3 c_{1}}\right)}{\sqrt[3]{-x^{6}+20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(-x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}}\right. \\
& \left.+\sqrt[3]{-x^{6}+20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(-x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}\right) \\
& y(x) \rightarrow \frac{1}{72}\left(-18 x^{2}-\frac{9 i(\sqrt{3}-i) x\left(x^{3}+8 e^{3 c_{1}}\right)}{\sqrt[3]{-x^{6}+20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(-x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}}\right. \\
& \left.+9 i(\sqrt{3}+i) \sqrt[3]{-x^{6}+20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(-x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}\right) \\
& y(x) \rightarrow \frac{1}{72}\left(-18 x^{2}+\frac{9 i(\sqrt{3}+i) x\left(x^{3}+8 e^{3 c_{1}}\right)}{\sqrt[3]{-x^{6}+20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(-x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}}\right. \\
& \left.-9(1+i \sqrt{3}) \sqrt[3]{-x^{6}+20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(-x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}\right) \\
& y(x) \rightarrow \frac{1}{4}\left(-x^{2}+\frac{x\left(x^{3}-8 e^{3 c_{1}}\right)}{\sqrt[3]{-x^{6}-20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}}\right. \\
& \left.+\sqrt[3]{-x^{6}-20 e^{3 c_{1}} x^{3}+8 \sqrt{e^{3 c_{1}}\left(x^{3}+e^{3 c_{1}}\right)^{3}}+8 e^{6 c_{1}}}\right)
\end{aligned}
$$

### 1.63 problem 90

> 1.63.1 Solving as dAlembert ode

Internal problem ID [12480]
Internal file name [OUTPUT/11132_Monday_October_16_2023_09_51_36_PM_33290913/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 90.
ODE order: 1.
ODE degree: 2 .

The type(s) of ODE detected by this program : "dAlembert"
Maple gives the following as the ode type

```
[[_homogeneous, `class C`], _rational, _dAlembert]
```

$$
y-x y^{\prime 2}-y^{\prime 2}=0
$$

### 1.63.1 Solving as dAlembert ode

Let $p=y^{\prime}$ the ode becomes

$$
-x p^{2}-p^{2}+y=0
$$

Solving for $y$ from the above results in

$$
\begin{equation*}
y=x p^{2}+p^{2} \tag{1A}
\end{equation*}
$$

This has the form

$$
\begin{equation*}
y=x f(p)+g(p) \tag{}
\end{equation*}
$$

Where $f, g$ are functions of $p=y^{\prime}(x)$. The above ode is dAlembert ode which is now solved. Taking derivative of $\left({ }^{*}\right)$ w.r.t. $x$ gives

$$
\begin{align*}
p & =f+\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \\
p-f & =\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \tag{2}
\end{align*}
$$

Comparing the form $y=x f+g$ to (1A) shows that

$$
\begin{aligned}
& f=p^{2} \\
& g=p^{2}
\end{aligned}
$$

Hence (2) becomes

$$
\begin{equation*}
-p^{2}+p=(2 x p+2 p) p^{\prime}(x) \tag{2~A}
\end{equation*}
$$

The singular solution is found by setting $\frac{d p}{d x}=0$ in the above which gives

$$
-p^{2}+p=0
$$

Solving for $p$ from the above gives

$$
\begin{aligned}
& p=0 \\
& p=1
\end{aligned}
$$

Substituting these in (1A) gives

$$
\begin{aligned}
& y=0 \\
& y=x+1
\end{aligned}
$$

The general solution is found when $\frac{\mathrm{d} p}{\mathrm{~d} x} \neq 0$. From eq. (2A). This results in

$$
\begin{equation*}
p^{\prime}(x)=\frac{-p(x)^{2}+p(x)}{2 p(x) x+2 p(x)} \tag{3}
\end{equation*}
$$

This ODE is now solved for $p(x)$.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
p^{\prime}(x)+p(x) p(x)=q(x)
$$

Where here

$$
\begin{aligned}
p(x) & =\frac{1}{2 x+2} \\
q(x) & =\frac{1}{2 x+2}
\end{aligned}
$$

Hence the ode is

$$
p^{\prime}(x)+\frac{p(x)}{2 x+2}=\frac{1}{2 x+2}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \frac{1}{2 x+2} d x} \\
& =\sqrt{x+1}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu p) & =(\mu)\left(\frac{1}{2 x+2}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}(\sqrt{x+1} p) & =(\sqrt{x+1})\left(\frac{1}{2 x+2}\right) \\
\mathrm{d}(\sqrt{x+1} p) & =\left(\frac{1}{2 \sqrt{x+1}}\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \sqrt{x+1} p=\int \frac{1}{2 \sqrt{x+1}} \mathrm{~d} x \\
& \sqrt{x+1} p=\sqrt{x+1}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\sqrt{x+1}$ results in

$$
p(x)=1+\frac{c_{1}}{\sqrt{x+1}}
$$

Substituing the above solution for $p$ in (2A) gives

$$
y=x\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2}+\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=0  \tag{1}\\
& y=x+1  \tag{2}\\
& y=x\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2}+\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2} \tag{3}
\end{align*}
$$

## Verification of solutions

$$
y=0
$$

Verified OK.

$$
y=x+1
$$

Verified OK.

$$
y=x\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2}+\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2}
$$

Verified OK.
Maple trace

- Methods for first order ODEs:
*** Sublevel 2 ***
Methods for first order ODEs:
-> Solving 1st order ODE of high degree, 1st attempt
trying 1st order WeierstrassP solution for high degree ODE
trying 1st order WeierstrassPPrime solution for high degree ODE
trying 1st order JacobiSN solution for high degree ODE
trying 1st order ODE linearizable_by_differentiation
trying differential order: 1; missing variables
trying dAlembert
<- dAlembert successful`
$\checkmark$ Solution by Maple
Time used: 0.078 (sec). Leaf size: 53

```
dsolve(y(x)=x*diff(y(x),x)^2+diff(y(x),x)^2,y(x), singsol=all)
```

$$
\begin{aligned}
& y(x)=0 \\
& y(x)=\frac{\left(x+1+\sqrt{(1+x)\left(c_{1}+1\right)}\right)^{2}}{1+x} \\
& y(x)=\frac{\left(-x-1+\sqrt{(1+x)\left(c_{1}+1\right)}\right)^{2}}{1+x}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.111 (sec). Leaf size: 57
DSolve[y[x]==x*(y'[x]) $2+(y \prime[x]) \sim 2, y[x], x$, IncludeSingularSolutions $->$ True]

$$
\begin{aligned}
& y(x) \rightarrow x-c_{1} \sqrt{x+1}+1+\frac{c_{1}^{2}}{4} \\
& y(x) \rightarrow x+c_{1} \sqrt{x+1}+1+\frac{c_{1}^{2}}{4} \\
& y(x) \rightarrow 0
\end{aligned}
$$

### 1.64 problem 91

1.64.1 Solving as dAlembert ode

Internal problem ID [12481]
Internal file name [OUTPUT/11133_Monday_October_16_2023_09_51_37_PM_66995158/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 91.
ODE order: 1.
ODE degree: 2 .

The type(s) of ODE detected by this program : "dAlembert"
Maple gives the following as the ode type

```
[[_1st_order, _with_linear_symmetries], _dAlembert]
```

$$
y-x\left(1+y^{\prime}\right)-y^{\prime 2}=0
$$

### 1.64.1 Solving as dAlembert ode

Let $p=y^{\prime}$ the ode becomes

$$
y-x(1+p)-p^{2}=0
$$

Solving for $y$ from the above results in

$$
\begin{equation*}
y=x(1+p)+p^{2} \tag{1~A}
\end{equation*}
$$

This has the form

$$
\begin{equation*}
y=x f(p)+g(p) \tag{*}
\end{equation*}
$$

Where $f, g$ are functions of $p=y^{\prime}(x)$. The above ode is dAlembert ode which is now solved. Taking derivative of $\left({ }^{*}\right)$ w.r.t. $x$ gives

$$
\begin{align*}
p & =f+\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \\
p-f & =\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \tag{2}
\end{align*}
$$

Comparing the form $y=x f+g$ to (1A) shows that

$$
\begin{aligned}
& f=1+p \\
& g=p^{2}
\end{aligned}
$$

Hence (2) becomes

$$
\begin{equation*}
-1=(x+2 p) p^{\prime}(x) \tag{2~A}
\end{equation*}
$$

The singular solution is found by setting $\frac{d p}{d x}=0$ in the above which gives

$$
-1=0
$$

No singular solution are found
The general solution is found when $\frac{\mathrm{d} p}{\mathrm{~d} x} \neq 0$. From eq. (2A). This results in

$$
\begin{equation*}
p^{\prime}(x)=-\frac{1}{x+2 p(x)} \tag{3}
\end{equation*}
$$

This ODE is now solved for $p(x)$.
Inverting the above ode gives

$$
\begin{equation*}
\frac{d}{d p} x(p)=-x(p)-2 p \tag{4}
\end{equation*}
$$

This ODE is now solved for $x(p)$.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
\frac{d}{d p} x(p)+p(p) x(p)=q(p)
$$

Where here

$$
\begin{aligned}
p(p) & =1 \\
q(p) & =-2 p
\end{aligned}
$$

Hence the ode is

$$
\frac{d}{d p} x(p)+x(p)=-2 p
$$

The integrating factor $\mu$ is

$$
\begin{gathered}
\mu=\mathrm{e}^{\int 1 d p} \\
=\mathrm{e}^{p}
\end{gathered}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} p}(\mu x) & =(\mu)(-2 p) \\
\frac{\mathrm{d}}{\mathrm{~d} p}\left(\mathrm{e}^{p} x\right) & =\left(\mathrm{e}^{p}\right)(-2 p) \\
\mathrm{d}\left(\mathrm{e}^{p} x\right) & =\left(-2 p \mathrm{e}^{p}\right) \mathrm{d} p
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \mathrm{e}^{p} x=\int-2 p \mathrm{e}^{p} \mathrm{~d} p \\
& \mathrm{e}^{p} x=-2(p-1) \mathrm{e}^{p}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\mathrm{e}^{p}$ results in

$$
x(p)=-2 \mathrm{e}^{-p}(p-1) \mathrm{e}^{p}+c_{1} \mathrm{e}^{-p}
$$

which simplifies to

$$
x(p)=-2 p+2+c_{1} \mathrm{e}^{-p}
$$

Now we need to eliminate $p$ between the above and (1A). One way to do this is by solving (1) for $p$. This results in

$$
\begin{aligned}
& p=-\frac{x}{2}+\frac{\sqrt{x^{2}+4 y-4 x}}{2} \\
& p=-\frac{x}{2}-\frac{\sqrt{x^{2}+4 y-4 x}}{2}
\end{aligned}
$$

Substituting the above in the solution for $x$ found above gives

$$
\begin{aligned}
& x=x-\sqrt{x^{2}+4 y-4 x}+2+c_{1} \mathrm{e}^{\frac{x}{2}-\frac{\sqrt{x^{2}+4 y-4 x}}{2}} \\
& x=x+\sqrt{x^{2}+4 y-4 x}+2+c_{1} \mathrm{e}^{\frac{x}{2}+\frac{\sqrt{x^{2}+4 y-4 x}}{2}}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& x=x-\sqrt{x^{2}+4 y-4 x}+2+c_{1} \mathrm{e}^{\frac{x}{2}-\frac{\sqrt{x^{2}+4 y-4 x}}{2}}  \tag{1}\\
& x=x+\sqrt{x^{2}+4 y-4 x}+2+c_{1} \mathrm{e}^{\frac{x}{2}+\frac{\sqrt{x^{2}+4 y-4 x}}{2}} \tag{2}
\end{align*}
$$

Verification of solutions

$$
x=x-\sqrt{x^{2}+4 y-4 x}+2+c_{1} \mathrm{e}^{\frac{x}{2}-\frac{\sqrt{x^{2}+4 y-4 x}}{2}}
$$

Verified OK.

$$
x=x+\sqrt{x^{2}+4 y-4 x}+2+c_{1} \mathrm{e}^{\frac{x}{2}+\frac{\sqrt{x^{2}+4 y-4 x}}{2}}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
-> Solving 1st order ODE of high degree, 1st attempt
trying 1st order WeierstrassP solution for high degree ODE
trying 1st order WeierstrassPPrime solution for high degree ODE
trying 1st order JacobiSN solution for high degree ODE
trying 1st order ODE linearizable_by_differentiation
trying differential order: 1; missing variables
trying dAlembert
<- dAlembert successful`
```

Solution by Maple
Time used: 0.063 (sec). Leaf size: 36

```
dsolve(y(x)=x*(1+\operatorname{diff}(y(x),x))+\operatorname{diff}(y(x),x)^2,y(x), singsol=all)
```

$$
y(x)=x-\frac{x^{2}}{4}+\text { LambertW }\left(\frac{c_{1} \mathrm{e}^{\frac{x}{2}-1}}{2}\right)^{2}+2 \operatorname{LambertW}\left(\frac{c_{1} \mathrm{e}^{\frac{x}{2}-1}}{2}\right)+1
$$

$\checkmark$ Solution by Mathematica
Time used: 3.313 (sec). Leaf size: 177
DSolve[y[x]==x*(1+y'[x])+(y'[x])~2,y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& \text { Solve }\left[-\sqrt{x^{2}+4 y(x)-4 x}+2 \log \left(\sqrt{x^{2}+4 y(x)-4 x}-x+2\right)\right. \\
& \left.\quad-2 \log \left(-x \sqrt{x^{2}+4 y(x)-4 x}+x^{2}+4 y(x)-2 x-4\right)+x=c_{1}, y(x)\right] \\
& \text { Solve }\left[-4 \operatorname{arctanh}\left(\frac{(x-5) \sqrt{x^{2}+4 y(x)-4 x}-x^{2}-4 y(x)+7 x-6}{(x-3) \sqrt{x^{2}+4 y(x)-4 x}-x^{2}-4 y(x)+5 x-2}\right)\right. \\
& \left.+\sqrt{x^{2}+4 y(x)-4 x}+x=c_{1}, y(x)\right]
\end{aligned}
$$

### 1.65 problem 92

1.65.1 Solving as dAlembert ode

Internal problem ID [12482]
Internal file name [OUTPUT/11134_Monday_October_16_2023_09_51_37_PM_96570058/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 92.
ODE order: 1.
ODE degree: 2 .

The type(s) of ODE detected by this program : "dAlembert"
Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _dAlembert]
```

$$
y-y y^{\prime 2}-2 y^{\prime} x=0
$$

### 1.65.1 Solving as dAlembert ode

Let $p=y^{\prime}$ the ode becomes

$$
-y p^{2}-2 p x+y=0
$$

Solving for $y$ from the above results in

$$
\begin{equation*}
y=-\frac{2 p x}{p^{2}-1} \tag{1A}
\end{equation*}
$$

This has the form

$$
\begin{equation*}
y=x f(p)+g(p) \tag{}
\end{equation*}
$$

Where $f, g$ are functions of $p=y^{\prime}(x)$. The above ode is dAlembert ode which is now solved. Taking derivative of $\left({ }^{*}\right)$ w.r.t. $x$ gives

$$
\begin{align*}
p & =f+\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \\
p-f & =\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \tag{2}
\end{align*}
$$

Comparing the form $y=x f+g$ to (1A) shows that

$$
\begin{aligned}
& f=-\frac{2 p}{p^{2}-1} \\
& g=0
\end{aligned}
$$

Hence (2) becomes

$$
\begin{equation*}
p+\frac{2 p}{p^{2}-1}=x\left(-\frac{2}{p^{2}-1}+\frac{4 p^{2}}{\left(p^{2}-1\right)^{2}}\right) p^{\prime}(x) \tag{2~A}
\end{equation*}
$$

The singular solution is found by setting $\frac{d p}{d x}=0$ in the above which gives

$$
p+\frac{2 p}{p^{2}-1}=0
$$

Solving for $p$ from the above gives

$$
\begin{aligned}
& p=0 \\
& p=i \\
& p=-i
\end{aligned}
$$

Substituting these in (1A) gives

$$
\begin{aligned}
& y=0 \\
& y=-i x \\
& y=i x
\end{aligned}
$$

The general solution is found when $\frac{\mathrm{d} p}{\mathrm{~d} x} \neq 0$. From eq. (2A). This results in

$$
\begin{equation*}
p^{\prime}(x)=\frac{p(x)+\frac{2 p(x)}{p(x)^{2}-1}}{x\left(-\frac{2}{p(x)^{2}-1}+\frac{4 p(x)^{2}}{\left(p(x)^{2}-1\right)^{2}}\right)} \tag{3}
\end{equation*}
$$

This ODE is now solved for $p(x)$.
Inverting the above ode gives

$$
\begin{equation*}
\frac{d}{d p} x(p)=\frac{x(p)\left(-\frac{2}{p^{2}-1}+\frac{4 p^{2}}{\left(p^{2}-1\right)^{2}}\right)}{p+\frac{2 p}{p^{2}-1}} \tag{4}
\end{equation*}
$$

This ODE is now solved for $x(p)$.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
\frac{d}{d p} x(p)+p(p) x(p)=q(p)
$$

Where here

$$
\begin{aligned}
& p(p)=-\frac{2}{p^{3}-p} \\
& q(p)=0
\end{aligned}
$$

Hence the ode is

$$
\frac{d}{d p} x(p)-\frac{2 x(p)}{p^{3}-p}=0
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2}{p^{3}-p} d p} \\
& =\mathrm{e}^{-\ln (p-1)-\ln (p+1)+2 \ln (p)}
\end{aligned}
$$

Which simplifies to

$$
\mu=\frac{p^{2}}{p^{2}-1}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} p} \mu x & =0 \\
\frac{\mathrm{~d}}{\mathrm{~d} p}\left(\frac{p^{2} x}{p^{2}-1}\right) & =0
\end{aligned}
$$

Integrating gives

$$
\frac{p^{2} x}{p^{2}-1}=c_{3}
$$

Dividing both sides by the integrating factor $\mu=\frac{p^{2}}{p^{2}-1}$ results in

$$
x(p)=\frac{c_{3}\left(p^{2}-1\right)}{p^{2}}
$$

Now we need to eliminate $p$ between the above and (1A). One way to do this is by solving (1) for $p$. This results in

$$
\begin{aligned}
& p=\frac{\sqrt{x^{2}+y^{2}}-x}{y} \\
& p=-\frac{x+\sqrt{x^{2}+y^{2}}}{y}
\end{aligned}
$$

Substituting the above in the solution for $x$ found above gives

$$
\begin{aligned}
& x=-\frac{2 c_{3} x}{\sqrt{x^{2}+y^{2}}-x} \\
& x=\frac{2 c_{3} x}{x+\sqrt{x^{2}+y^{2}}}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=0  \tag{1}\\
& y=-i x  \tag{2}\\
& y=i x  \tag{3}\\
& x=-\frac{2 c_{3} x}{\sqrt{x^{2}+y^{2}}-x}  \tag{4}\\
& x=\frac{2 c_{3} x}{x+\sqrt{x^{2}+y^{2}}} \tag{5}
\end{align*}
$$

Verification of solutions

$$
y=0
$$

Verified OK.

$$
y=-i x
$$

Verified OK.

$$
y=i x
$$

Verified OK.

$$
x=-\frac{2 c_{3} x}{\sqrt{x^{2}+y^{2}}-x}
$$

Verified OK.

$$
x=\frac{2 c_{3} x}{x+\sqrt{x^{2}+y^{2}}}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
    *** Sublevel 2 ***
    Methods for first order ODEs:
    -> Solving 1st order ODE of high degree, 1st attempt
    trying 1st order WeierstrassP solution for high degree ODE
    trying 1st order WeierstrassPPrime solution for high degree ODE
    trying 1st order JacobiSN solution for high degree ODE
    trying 1st order ODE linearizable_by_differentiation
    trying differential order: 1; missing variables
    trying simple symmetries for implicit equations
    <- symmetries for implicit equations successful`
```

$\checkmark$ Solution by Maple
Time used: 0.109 (sec). Leaf size: 71
dsolve( $y(x)=y(x) * \operatorname{diff}(y(x), x)^{\wedge} 2+2 * x * \operatorname{diff}(y(x), x), y(x)$, singsol=all)

$$
\begin{aligned}
& y(x)=-i x \\
& y(x)=i x \\
& y(x)=0 \\
& y(x)=\sqrt{c_{1}\left(-2 x+c_{1}\right)} \\
& y(x)=\sqrt{c_{1}\left(c_{1}+2 x\right)} \\
& y(x)=-\sqrt{c_{1}\left(-2 x+c_{1}\right)} \\
& y(x)=-\sqrt{c_{1}\left(c_{1}+2 x\right)}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.788 (sec). Leaf size: 126
DSolve $\left[y[x]==y[x] *\left(y y^{\prime}[x]\right)^{\sim} 2+2 * x * y\right.$ ' $[x], y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow-e^{\frac{c_{1}}{2}} \sqrt{-2 x+e^{c_{1}}} \\
& y(x) \rightarrow e^{\frac{c_{1}}{2}} \sqrt{-2 x+e^{c_{1}}} \\
& y(x) \rightarrow-e^{\frac{c_{1}}{2}} \sqrt{2 x+e^{c_{1}}} \\
& y(x) \rightarrow e^{\frac{c_{1}}{2}} \sqrt{2 x+e^{c_{1}}} \\
& y(x) \rightarrow 0 \\
& y(x) \rightarrow-i x \\
& y(x) \rightarrow i x
\end{aligned}
$$

### 1.66 problem 94

1.66.1 Maple step by step solution

750
Internal problem ID [12483]
Internal file name [OUTPUT/11135_Monday_October_16_2023_09_51_43_PM_60606324/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 94.
ODE order: 1.
ODE degree: 2 .

The type(s) of ODE detected by this program : "quadrature"
Maple gives the following as the ode type
[_quadrature]

$$
y-y y^{\prime}-y^{\prime}+y^{\prime 2}=0
$$

Solving the given ode for $y^{\prime}$ results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$
\begin{align*}
& y^{\prime}=1  \tag{1}\\
& y^{\prime}=y \tag{2}
\end{align*}
$$

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives

$$
\begin{aligned}
y & =\int 1 \mathrm{~d} x \\
& =x+c_{1}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x+c_{1} \tag{1}
\end{equation*}
$$

## Verification of solutions

$$
y=x+c_{1}
$$

Verified OK.
Solving equation (2)
Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{y} d y & =x+c_{2} \\
\ln (y) & =x+c_{2} \\
y & =\mathrm{e}^{x+c_{2}} \\
y & =c_{2} \mathrm{e}^{x}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{2} \mathrm{e}^{x} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=c_{2} \mathrm{e}^{x}
$$

Verified OK.

### 1.66.1 Maple step by step solution

Let's solve

$$
y-y y^{\prime}-y^{\prime}+y^{\prime 2}=0
$$

- Highest derivative means the order of the ODE is 1


## $y^{\prime}$

- Integrate both sides with respect to $x$

$$
\int\left(y-y y^{\prime}-y^{\prime}+y^{\prime 2}\right) d x=\int 0 d x+c_{1}
$$

- Cannot compute integral

$$
\int\left(y-y y^{\prime}-y^{\prime}+y^{\prime 2}\right) d x=c_{1}
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful
Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 14

```
dsolve(y(x)=y(x)*diff(y(x),x)+diff(y(x),x)-diff(y(x),x)^2,y(x), singsol=all)
```

$$
\begin{aligned}
& y(x)=c_{1}+x \\
& y(x)=c_{1} \mathrm{e}^{x}
\end{aligned}
$$

$\sqrt{ }$ Solution by Mathematica
Time used: 0.019 (sec). Leaf size: 19
DSolve[y[x]==y[x]*y'[x]+y'[x]-(y'[x])~2,y[x],x,IncludeSingularSolutions -> True]

$$
\begin{aligned}
& y(x) \rightarrow c_{1} e^{x} \\
& y(x) \rightarrow x+c_{1}
\end{aligned}
$$

### 1.67 problem 95

1.67.1 Solving as clairaut ode

Internal problem ID [12484]
Internal file name [OUTPUT/11136_Monday_October_16_2023_09_51_44_PM_5596819/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 95.
ODE order: 1.
ODE degree: 2.

The type(s) of ODE detected by this program : "clairaut"
Maple gives the following as the ode type

```
[[_1st_order, _with_linear_symmetries], _rational, _Clairaut]
```

$$
y-y^{\prime} x-\sqrt{-y^{\prime 2}+1}=0
$$

### 1.67.1 Solving as clairaut ode

This is Clairaut ODE. It has the form

$$
y=y^{\prime} x+g\left(y^{\prime}\right)
$$

Where $g$ is function of $y^{\prime}(x)$. Let $p=y^{\prime}$ the ode becomes

$$
y-p x-\sqrt{-p^{2}+1}=0
$$

Solving for $y$ from the above results in

$$
\begin{equation*}
y=p x+\sqrt{-p^{2}+1} \tag{1~A}
\end{equation*}
$$

The above ode is a Clairaut ode which is now solved. We start by replacing $y^{\prime}$ by $p$ which gives

$$
\begin{aligned}
y & =p x+\sqrt{-p^{2}+1} \\
& =p x+\sqrt{-p^{2}+1}
\end{aligned}
$$

Writing the ode as

$$
y=p x+g(p)
$$

We now write $g \equiv g(p)$ to make notation simpler but we should always remember that $g$ is function of $p$ which in turn is function of $x$. Hence the above becomes

$$
\begin{equation*}
y=p x+g \tag{1}
\end{equation*}
$$

Then we see that

$$
g=\sqrt{-p^{2}+1}
$$

Taking derivative of (1) w.r.t. $x$ gives

$$
\begin{aligned}
& p=\frac{d}{d x}(x p+g) \\
& p=\left(p+x \frac{d p}{d x}\right)+\left(g^{\prime} \frac{d p}{d x}\right) \\
& p=p+\left(x+g^{\prime}\right) \frac{d p}{d x} \\
& 0=\left(x+g^{\prime}\right) \frac{d p}{d x}
\end{aligned}
$$

Where $g^{\prime}$ is derivative of $g(p)$ w.r.t. $p$. The general solution is given by

$$
\begin{aligned}
\frac{d p}{d x} & =0 \\
p & =c_{1}
\end{aligned}
$$

Substituting this in (1) gives the general solution as

$$
y=c_{1} x+\sqrt{-c_{1}^{2}+1}
$$

The singular solution is found from solving for $p$ from

$$
x+g^{\prime}(p)=0
$$

And substituting the result back in (1). Since we found above that $g=\sqrt{-p^{2}+1}$, then the above equation becomes

$$
\begin{aligned}
x+g^{\prime}(p) & =x-\frac{p}{\sqrt{-p^{2}+1}} \\
& =0
\end{aligned}
$$

Solving the above for $p$ results in

$$
p_{1}=\frac{x}{\sqrt{x^{2}+1}}
$$

Substituting the above back in (1) results in

$$
y_{1}=\sqrt{x^{2}+1}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=c_{1} x+\sqrt{-c_{1}^{2}+1}  \tag{1}\\
& y=\sqrt{x^{2}+1} \tag{2}
\end{align*}
$$

Verification of solutions

$$
y=c_{1} x+\sqrt{-c_{1}^{2}+1}
$$

Verified OK.

$$
y=\sqrt{x^{2}+1}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
-> Solving 1st order ODE of high degree, 1st attempt
trying 1st order WeierstrassP solution for high degree ODE
trying 1st order WeierstrassPPrime solution for high degree ODE
trying 1st order JacobiSN solution for high degree ODE
trying 1st order ODE linearizable_by_differentiation
trying differential order: 1; missing variables
trying dAlembert
<- dAlembert successful`
```

$\checkmark$ Solution by Maple
Time used: 0.344 (sec). Leaf size: 17

```
dsolve(y(x)=x*diff(y(x),x)+sqrt(1-diff(y(x),x)^2),y(x), singsol=all)
```

$$
y(x)=c_{1} x+\sqrt{-c_{1}^{2}+1}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.08 (sec). Leaf size: 27
DSolve $[y[x]==x * y$ ' $[x]+$ Sqrt $[1-y$ ' $[x] \sim 2], y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow c_{1} x+\sqrt{1-c_{1}^{2}} \\
& y(x) \rightarrow 1
\end{aligned}
$$

### 1.68 problem 96

1.68.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 756
1.68.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 758
1.68.3 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 759
1.68.4 Solving as homogeneousTypeMapleC ode . . . . . . . . . . . . . 761
1.68.5 Solving as first order ode lie symmetry lookup ode . . . . . . . 764
1.68.6 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 768
1.68.7 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 772

Internal problem ID [12485]
Internal file name [OUTPUT/11137_Monday_October_16_2023_09_51_47_PM_19113061/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 96.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "homogeneousTypeMapleC", "first_order_ode_lie_symmetry__lookup"

Maple gives the following as the ode type
[_separable]

$$
y-y^{\prime} x-y^{\prime}=0
$$

### 1.68.1 Solving as separable ode

In canonical form the ODE is

$$
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{y}{x+1}
\end{aligned}
$$

Where $f(x)=\frac{1}{x+1}$ and $g(y)=y$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{y} d y & =\frac{1}{x+1} d x \\
\int \frac{1}{y} d y & =\int \frac{1}{x+1} d x \\
\ln (y) & =\ln (x+1)+c_{1} \\
y & =\mathrm{e}^{\ln (x+1)+c_{1}} \\
& =c_{1}(x+1)
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1}(x+1) \tag{1}
\end{equation*}
$$



Figure 130: Slope field plot

Verification of solutions

$$
y=c_{1}(x+1)
$$

Verified OK.

### 1.68.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{1}{x+1} \\
& q(x)=0
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{y}{x+1}=0
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{1}{x+1} d x} \\
& =\frac{1}{x+1}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x} \mu y & =0 \\
\frac{\mathrm{~d}}{\mathrm{~d} x}\left(\frac{y}{x+1}\right) & =0
\end{aligned}
$$

Integrating gives

$$
\frac{y}{x+1}=c_{1}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x+1}$ results in

$$
y=c_{1}(x+1)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1}(x+1) \tag{1}
\end{equation*}
$$



Figure 131: Slope field plot
Verification of solutions

$$
y=c_{1}(x+1)
$$

Verified OK.

### 1.68.3 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
u(x) x-\left(u^{\prime}(x) x+u(x)\right) x-u^{\prime}(x) x-u(x)=0
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =-\frac{u}{x(x+1)}
\end{aligned}
$$

Where $f(x)=-\frac{1}{x(x+1)}$ and $g(u)=u$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{u} d u & =-\frac{1}{x(x+1)} d x \\
\int \frac{1}{u} d u & =\int-\frac{1}{x(x+1)} d x \\
\ln (u) & =\ln (x+1)-\ln (x)+c_{2} \\
u & =\mathrm{e}^{\ln (x+1)-\ln (x)+c_{2}} \\
& =c_{2} \mathrm{e}^{\ln (x+1)-\ln (x)}
\end{aligned}
$$

Which simplifies to

$$
u(x)=c_{2}\left(1+\frac{1}{x}\right)
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =u x \\
& =x c_{2}\left(1+\frac{1}{x}\right)
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x c_{2}\left(1+\frac{1}{x}\right) \tag{1}
\end{equation*}
$$



Figure 132: Slope field plot
Verification of solutions

$$
y=x c_{2}\left(1+\frac{1}{x}\right)
$$

Verified OK.

### 1.68.4 Solving as homogeneousTypeMapleC ode

Let $Y=y+y_{0}$ and $X=x+x_{0}$ then the above is transformed to new ode in $Y(X)$

$$
\frac{d}{d X} Y(X)=\frac{Y(X)+y_{0}}{X+x_{0}+1}
$$

Solving for possible values of $x_{0}$ and $y_{0}$ which makes the above ode a homogeneous ode results in

$$
\begin{aligned}
x_{0} & =-1 \\
y_{0} & =0
\end{aligned}
$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$
\frac{d}{d X} Y(X)=\frac{Y(X)}{X}
$$

In canonical form, the ODE is

$$
\begin{align*}
Y^{\prime} & =F(X, Y) \\
& =\frac{Y}{X} \tag{1}
\end{align*}
$$

An ode of the form $Y^{\prime}=\frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order $n$ if

$$
f\left(t^{n} X, t^{n} Y\right)=t^{n} f(X, Y)
$$

In this case, it can be seen that both $M=Y$ and $N=X$ are both homogeneous and of the same order $n=1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u=\frac{Y}{X}$, or $Y=u X$. Hence

$$
\frac{\mathrm{d} Y}{\mathrm{~d} X}=\frac{\mathrm{d} u}{\mathrm{~d} X} X+u
$$

Applying the transformation $Y=u X$ to the above ODE in (1) gives

$$
\begin{aligned}
\frac{\mathrm{d} u}{\mathrm{~d} X} X+u & =u \\
\frac{\mathrm{~d} u}{\mathrm{~d} X} & =0
\end{aligned}
$$

Or

$$
\frac{d}{d X} u(X)=0
$$

Which is now solved as separable in $u(X)$. Which is now solved in $u(X)$. Integrating both sides gives

$$
\begin{aligned}
u(X) & =\int 0 \mathrm{~d} X \\
& =c_{2}
\end{aligned}
$$

Now $u$ in the above solution is replaced back by $Y$ using $u=\frac{Y}{X}$ which results in the solution

$$
Y(X)=X c_{2}
$$

Using the solution for $Y(X)$

$$
Y(X)=X c_{2}
$$

And replacing back terms in the above solution using

$$
\begin{aligned}
& Y=y+y_{0} \\
& X=x+x_{0}
\end{aligned}
$$

Or

$$
\begin{aligned}
& Y=y \\
& X=x-1
\end{aligned}
$$

Then the solution in $y$ becomes

$$
y=(x+1) c_{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=(x+1) c_{2} \tag{1}
\end{equation*}
$$



Figure 133: Slope field plot
Verification of solutions

$$
y=(x+1) c_{2}
$$

Verified OK.

### 1.68.5 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
y^{\prime} & =\frac{y}{x+1} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 115: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{\prime}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=x+1 \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{x+1} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{x+1}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{y}{x+1}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{y}{(x+1)^{2}} \\
S_{y} & =\frac{1}{x+1}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=0 \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=0
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{y}{x+1}=c_{1}
$$

Which simplifies to

$$
\frac{y}{x+1}=c_{1}
$$

Which gives

$$
y=c_{1}(x+1)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | $\begin{gathered} \text { Canonical } \\ \text { coordinates } \\ \text { transformation } \end{gathered}$ | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{y}{x+1}$ |  | $\frac{d S}{d R}=0$ |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  | $\xrightarrow{\text { a }}$ 他 $\rightarrow \rightarrow \rightarrow \rightarrow 22 \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ |
| $\rightarrow \rightarrow+1$ | $R=x$ | $\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow$ 他 |
| $\xrightarrow[\rightarrow \rightarrow-9 \rightarrow 0]{\rightarrow \rightarrow-\infty}$ | $S=\underline{y}$ |  |
| $\rightarrow \rightarrow \rightarrow 0$ | $\overline{x+1}$ |  |
|  |  |  |
|  |  | $\rightarrow$ |
|  |  |  |

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=c_{1}(x+1) \tag{1}
\end{equation*}
$$



Figure 134: Slope field plot
Verification of solutions

$$
y=c_{1}(x+1)
$$

Verified OK.

### 1.68.6 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\left(\frac{1}{y}\right) \mathrm{d} y & =\left(\frac{1}{x+1}\right) \mathrm{d} x \\
\left(-\frac{1}{x+1}\right) \mathrm{d} x+\left(\frac{1}{y}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=-\frac{1}{x+1} \\
& N(x, y)=\frac{1}{y}
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{x+1}\right) \\
& =0
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{y}\right) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{1}{x+1} \mathrm{~d} x \\
\phi & =-\ln (x+1)+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{y}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{y}=0+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=\frac{1}{y}
$$

Integrating the above w.r.t $y$ gives

$$
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{y}\right) \mathrm{d} y \\
f(y) & =\ln (y)+c_{1}
\end{aligned}
$$

Where $c_{1}$ is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives $\phi$

$$
\phi=-\ln (x+1)+\ln (y)+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=-\ln (x+1)+\ln (y)
$$

The solution becomes

$$
y=\mathrm{e}^{c_{1}}(x+1)
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{c_{1}}(x+1) \tag{1}
\end{equation*}
$$



Figure 135: Slope field plot
Verification of solutions

$$
y=\mathrm{e}^{c_{1}}(x+1)
$$

Verified OK.

### 1.68.7 Maple step by step solution

Let's solve

$$
y-y^{\prime} x-y^{\prime}=0
$$

- Highest derivative means the order of the ODE is 1 $y^{\prime}$
- Separate variables

$$
\frac{y^{\prime}}{y}=-\frac{1}{-x-1}
$$

- Integrate both sides with respect to $x$

$$
\int \frac{y^{\prime}}{y} d x=\int-\frac{1}{-x-1} d x+c_{1}
$$

- Evaluate integral

$$
\ln (y)=\ln (-x-1)+c_{1}
$$

- $\quad$ Solve for $y$

$$
y=-\mathrm{e}^{c_{1}}(x+1)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\sqrt{ }$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 9

```
dsolve(y(x)=x*diff(y(x),x)+diff(y(x),x),y(x), singsol=all)
```

$$
y(x)=c_{1}(1+x)
$$

$\checkmark$ Solution by Mathematica
Time used: 0.042 (sec). Leaf size: 16
DSolve[y[x]==x*y'[x]+y'[x],y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow c_{1}(x+1) \\
& y(x) \rightarrow 0
\end{aligned}
$$

### 1.69 problem 97

$$
\text { 1.69.1 Solving as clairaut ode . . . . . . . . . . . . . . . . . . . . . . . } 774
$$

Internal problem ID [12486]
Internal file name [OUTPUT/11138_Monday_October_16_2023_09_51_47_PM_68887557/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 97.
ODE order: 1.
ODE degree: 2 .

The type(s) of ODE detected by this program : "clairaut"
Maple gives the following as the ode type

```
[[_homogeneous, `class G`], _rational, _Clairaut]
```

$$
y-y^{\prime} x-\frac{1}{y^{\prime}}=0
$$

### 1.69.1 Solving as clairaut ode

This is Clairaut ODE. It has the form

$$
y=y^{\prime} x+g\left(y^{\prime}\right)
$$

Where $g$ is function of $y^{\prime}(x)$. Let $p=y^{\prime}$ the ode becomes

$$
y-p x-\frac{1}{p}=0
$$

Solving for $y$ from the above results in

$$
\begin{equation*}
y=\frac{p^{2} x+1}{p} \tag{1~A}
\end{equation*}
$$

The above ode is a Clairaut ode which is now solved. We start by replacing $y^{\prime}$ by $p$ which gives

$$
\begin{aligned}
y & =p x+\frac{1}{p} \\
& =p x+\frac{1}{p}
\end{aligned}
$$

Writing the ode as

$$
y=p x+g(p)
$$

We now write $g \equiv g(p)$ to make notation simpler but we should always remember that $g$ is function of $p$ which in turn is function of $x$. Hence the above becomes

$$
\begin{equation*}
y=p x+g \tag{1}
\end{equation*}
$$

Then we see that

$$
g=\frac{1}{p}
$$

Taking derivative of (1) w.r.t. $x$ gives

$$
\begin{aligned}
& p=\frac{d}{d x}(x p+g) \\
& p=\left(p+x \frac{d p}{d x}\right)+\left(g^{\prime} \frac{d p}{d x}\right) \\
& p=p+\left(x+g^{\prime}\right) \frac{d p}{d x} \\
& 0=\left(x+g^{\prime}\right) \frac{d p}{d x}
\end{aligned}
$$

Where $g^{\prime}$ is derivative of $g(p)$ w.r.t. $p$. The general solution is given by

$$
\begin{aligned}
\frac{d p}{d x} & =0 \\
p & =c_{1}
\end{aligned}
$$

Substituting this in (1) gives the general solution as

$$
y=c_{1} x+\frac{1}{c_{1}}
$$

The singular solution is found from solving for $p$ from

$$
x+g^{\prime}(p)=0
$$

And substituting the result back in (1). Since we found above that $g=\frac{1}{p}$, then the above equation becomes

$$
\begin{aligned}
x+g^{\prime}(p) & =x-\frac{1}{p^{2}} \\
& =0
\end{aligned}
$$

Solving the above for $p$ results in

$$
\begin{aligned}
& p_{1}=\frac{1}{\sqrt{x}} \\
& p_{2}=-\frac{1}{\sqrt{x}}
\end{aligned}
$$

Substituting the above back in (1) results in

$$
\begin{aligned}
& y_{1}=2 \sqrt{x} \\
& y_{2}=-2 \sqrt{x}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=c_{1} x+\frac{1}{c_{1}}  \tag{1}\\
& y=2 \sqrt{x}  \tag{2}\\
& y=-2 \sqrt{x} \tag{3}
\end{align*}
$$

Verification of solutions

$$
y=c_{1} x+\frac{1}{c_{1}}
$$

Verified OK.

$$
y=2 \sqrt{x}
$$

Verified OK.

$$
y=-2 \sqrt{x}
$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:
    *** Sublevel 2 ***
    Methods for first order ODEs:
    -> Solving 1st order ODE of high degree, 1st attempt
    trying 1st order WeierstrassP solution for high degree ODE
    trying 1st order WeierstrassPPrime solution for high degree ODE
    trying 1st order JacobiSN solution for high degree ODE
    trying 1st order ODE linearizable_by_differentiation
    trying differential order: 1; missing variables
    trying dAlembert
    <- dAlembert successful
```

Solution by Maple
Time used: 0.078 (sec). Leaf size: 27

```
dsolve(y(x)=x*diff(y(x),x)+1/diff (y(x),x),y(x), singsol=all)
```

$$
\begin{aligned}
& y(x)=-2 \sqrt{x} \\
& y(x)=2 \sqrt{x} \\
& y(x)=c_{1} x+\frac{1}{c_{1}}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.017 (sec). Leaf size: 41
DSolve[y[x]==x*y'[x]+1/y'[x],y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow c_{1} x+\frac{1}{c_{1}} \\
& y(x) \rightarrow \text { Indeterminate } \\
& y(x) \rightarrow-2 \sqrt{x} \\
& y(x) \rightarrow 2 \sqrt{x}
\end{aligned}
$$

### 1.70 problem 98

$$
\text { 1.70.1 Solving as clairaut ode . . . . . . . . . . . . . . . . . . . . . . . } 778
$$

Internal problem ID [12487]
Internal file name [OUTPUT/11139_Monday_October_16_2023_09_51_48_PM_44382532/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 98.
ODE order: 1.
ODE degree: 3 .

The type(s) of ODE detected by this program : "clairaut"
Maple gives the following as the ode type

```
[[_1st_order, _with_linear_symmetries], _Clairaut]
```

$$
y-y^{\prime} x+\frac{1}{y^{\prime 2}}=0
$$

### 1.70.1 Solving as clairaut ode

This is Clairaut ODE. It has the form

$$
y=y^{\prime} x+g\left(y^{\prime}\right)
$$

Where $g$ is function of $y^{\prime}(x)$. Let $p=y^{\prime}$ the ode becomes

$$
y-p x+\frac{1}{p^{2}}=0
$$

Solving for $y$ from the above results in

$$
\begin{equation*}
y=\frac{p^{3} x-1}{p^{2}} \tag{1~A}
\end{equation*}
$$

The above ode is a Clairaut ode which is now solved. We start by replacing $y^{\prime}$ by $p$ which gives

$$
\begin{aligned}
y & =p x-\frac{1}{p^{2}} \\
& =p x-\frac{1}{p^{2}}
\end{aligned}
$$

Writing the ode as

$$
y=p x+g(p)
$$

We now write $g \equiv g(p)$ to make notation simpler but we should always remember that $g$ is function of $p$ which in turn is function of $x$. Hence the above becomes

$$
\begin{equation*}
y=p x+g \tag{1}
\end{equation*}
$$

Then we see that

$$
g=-\frac{1}{p^{2}}
$$

Taking derivative of (1) w.r.t. $x$ gives

$$
\begin{aligned}
& p=\frac{d}{d x}(x p+g) \\
& p=\left(p+x \frac{d p}{d x}\right)+\left(g^{\prime} \frac{d p}{d x}\right) \\
& p=p+\left(x+g^{\prime}\right) \frac{d p}{d x} \\
& 0=\left(x+g^{\prime}\right) \frac{d p}{d x}
\end{aligned}
$$

Where $g^{\prime}$ is derivative of $g(p)$ w.r.t. $p$. The general solution is given by

$$
\begin{aligned}
\frac{d p}{d x} & =0 \\
p & =c_{1}
\end{aligned}
$$

Substituting this in (1) gives the general solution as

$$
y=c_{1} x-\frac{1}{c_{1}^{2}}
$$

The singular solution is found from solving for $p$ from

$$
x+g^{\prime}(p)=0
$$

And substituting the result back in (1). Since we found above that $g=-\frac{1}{p^{2}}$, then the above equation becomes

$$
\begin{aligned}
x+g^{\prime}(p) & =x+\frac{2}{p^{3}} \\
& =0
\end{aligned}
$$

Solving the above for $p$ results in

$$
\begin{aligned}
& p_{1}=\frac{\left(-2 x^{2}\right)^{\frac{1}{3}}}{x} \\
& p_{2}=-\frac{\left(-2 x^{2}\right)^{\frac{1}{3}}}{2 x}-\frac{i \sqrt{3}\left(-2 x^{2}\right)^{\frac{1}{3}}}{2 x} \\
& p_{3}=-\frac{\left(-2 x^{2}\right)^{\frac{1}{3}}}{2 x}+\frac{i \sqrt{3}\left(-2 x^{2}\right)^{\frac{1}{3}}}{2 x}
\end{aligned}
$$

Substituting the above back in (1) results in

$$
\begin{aligned}
& y_{1}=-\frac{3 x^{2} 2^{\frac{1}{3}}}{2\left(-x^{2}\right)^{\frac{2}{3}}} \\
& y_{2}=-\frac{3 x^{2} 2^{\frac{1}{3}}}{\left(-x^{2}\right)^{\frac{2}{3}}(i \sqrt{3}-1)} \\
& y_{3}=\frac{3 x^{2} 2^{\frac{1}{3}}}{\left(-x^{2}\right)^{\frac{2}{3}}(1+i \sqrt{3})}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=c_{1} x-\frac{1}{c_{1}^{2}}  \tag{1}\\
& y=-\frac{3 x^{2} 2^{\frac{1}{3}}}{2\left(-x^{2}\right)^{\frac{2}{3}}}  \tag{2}\\
& y=-\frac{3 x^{2} 2^{\frac{1}{3}}}{\left(-x^{2}\right)^{\frac{2}{3}}(i \sqrt{3}-1)}  \tag{3}\\
& y=\frac{3 x^{2} 2^{\frac{1}{3}}}{\left(-x^{2}\right)^{\frac{2}{3}}(1+i \sqrt{3})} \tag{4}
\end{align*}
$$

## Verification of solutions

$$
y=c_{1} x-\frac{1}{c_{1}^{2}}
$$

Verified OK.

$$
y=-\frac{3 x^{2} 2^{\frac{1}{3}}}{2\left(-x^{2}\right)^{\frac{2}{3}}}
$$

Verified OK.

$$
y=-\frac{3 x^{2} 2^{\frac{1}{3}}}{\left(-x^{2}\right)^{\frac{2}{3}}(i \sqrt{3}-1)}
$$

Verified OK.

$$
y=\frac{3 x^{2} 2^{\frac{1}{3}}}{\left(-x^{2}\right)^{\frac{2}{3}}(1+i \sqrt{3})}
$$

Verified OK.
Maple trace

```
`Methods for first order ODEs:
    *** Sublevel 2 ***
    Methods for first order ODEs:
    -> Solving 1st order ODE of high degree, 1st attempt
    trying 1st order WeierstrassP solution for high degree ODE
    trying 1st order WeierstrassPPrime solution for high degree ODE
    trying 1st order JacobiSN solution for high degree ODE
    trying 1st order ODE linearizable_by_differentiation
    trying differential order: 1; missing variables
    trying dAlembert
    <- dAlembert successful
```

$\checkmark$ Solution by Maple
Time used: 0.078 (sec). Leaf size: 74
dsolve( $y(x)=x * \operatorname{diff}(y(x), x)-1 / \operatorname{diff}(y(x), x) \sim 2, y(x)$, singsol=all)

$$
\begin{aligned}
& y(x)=\frac{32^{\frac{1}{3}}\left(-x^{2}\right)^{\frac{1}{3}}}{2} \\
& y(x)=-\frac{32^{\frac{1}{3}}\left(-x^{2}\right)^{\frac{1}{3}}(1+i \sqrt{3})}{4} \\
& y(x)=\frac{32^{\frac{1}{3}}\left(-x^{2}\right)^{\frac{1}{3}}(i \sqrt{3}-1)}{4} \\
& y(x)=c_{1} x-\frac{1}{c_{1}^{2}}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.019 (sec). Leaf size: 71
DSolve[y[x]==x*y'[x]-1/(y'[x])~2,y[x],x,IncludeSingularSolutions -> True]

$$
\begin{aligned}
& y(x) \rightarrow c_{1} x-\frac{1}{c_{1}^{2}} \\
& y(x) \rightarrow-3\left(-\frac{1}{2}\right)^{2 / 3} x^{2 / 3} \\
& y(x) \rightarrow-\frac{3 x^{2 / 3}}{2^{2 / 3}} \\
& y(x) \rightarrow \frac{3 \sqrt[3]{-1} x^{2 / 3}}{2^{2 / 3}}
\end{aligned}
$$

### 1.71 problem 110

1.71.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 783
1.71.2 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 785
1.71.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 787
1.71.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 791
1.71.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 796

Internal problem ID [12488]
Internal file name [OUTPUT/11140_Monday_October_16_2023_09_51_51_PM_40834021/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 110.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "homogeneousTypeD2", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry__lookup"

Maple gives the following as the ode type
[_linear]

$$
y^{\prime}-\frac{2 y}{x}=-\sqrt{3}
$$

### 1.71.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{2}{x} \\
& q(x)=-\sqrt{3}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{2 y}{x}=-\sqrt{3}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2}{x} d x} \\
& =\frac{1}{x^{2}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)(-\sqrt{3}) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{y}{x^{2}}\right) & =\left(\frac{1}{x^{2}}\right)(-\sqrt{3}) \\
\mathrm{d}\left(\frac{y}{x^{2}}\right) & =\left(-\frac{\sqrt{3}}{x^{2}}\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
\frac{y}{x^{2}} & =\int-\frac{\sqrt{3}}{x^{2}} \mathrm{~d} x \\
\frac{y}{x^{2}} & =\frac{\sqrt{3}}{x}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x^{2}}$ results in

$$
y=\sqrt{3} x+c_{1} x^{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\sqrt{3} x+c_{1} x^{2} \tag{1}
\end{equation*}
$$



Figure 136: Slope field plot
Verification of solutions

$$
y=\sqrt{3} x+c_{1} x^{2}
$$

Verified OK.

### 1.71.2 Solving as homogeneousTypeD2 ode

Using the change of variables $y=u(x) x$ on the above ode results in new ode in $u(x)$

$$
u^{\prime}(x) x-u(x)=-\sqrt{3}
$$

In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =\frac{u-\sqrt{3}}{x}
\end{aligned}
$$

Where $f(x)=\frac{1}{x}$ and $g(u)=u-\sqrt{3}$. Integrating both sides gives

$$
\frac{1}{u-\sqrt{3}} d u=\frac{1}{x} d x
$$

$$
\begin{aligned}
\int \frac{1}{u-\sqrt{3}} d u & =\int \frac{1}{x} d x \\
\ln (u-\sqrt{3}) & =\ln (x)+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
u-\sqrt{3}=\mathrm{e}^{\ln (x)+c_{2}}
$$

Which simplifies to

$$
u-\sqrt{3}=c_{3} x
$$

Which simplifies to

$$
u(x)=c_{3} \mathrm{e}^{c_{2}} x+\sqrt{3}
$$

Therefore the solution $y$ is

$$
\begin{aligned}
y & =x u \\
& =x\left(c_{3} \mathrm{e}^{c_{2}} x+\sqrt{3}\right)
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x\left(c_{3} \mathrm{e}^{c_{2}} x+\sqrt{3}\right) \tag{1}
\end{equation*}
$$



Figure 137: Slope field plot
Verification of solutions

$$
y=x\left(c_{3} \mathrm{e}^{c_{2}} x+\sqrt{3}\right)
$$

Verified OK.

### 1.71.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$
\begin{aligned}
& y^{\prime}=\frac{-\sqrt{3} x+2 y}{x} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
$$

The condition of Lie symmetry is the linearized PDE given by

$$
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
$$

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find $\xi, \eta$

Table 118: Lie symmetry infinitesimal lookup table for known first order ODE's

| ODE class | Form | $\xi$ | $\eta$ |
| :--- | :--- | :--- | :--- |
| linear ode | $y^{\prime}=f(x) y(x)+g(x)$ | 0 | $e^{\int f d x}$ |
| separable ode | $y^{\prime}=f(x) g(y)$ | $\frac{1}{f}$ | 0 |
| quadrature ode | $y^{\prime}=f(x)$ | 0 | 1 |
| quadrature ode | $y^{\prime}=g(y)$ | 1 | 0 |
| homogeneous ODEs of <br> Class A | $y^{\prime}=f\left(\frac{y}{x}\right)$ | $x$ | $y$ |
| homogeneous ODEs of <br> Class C | $y^{\prime}=(a+b x+c y)^{\frac{n}{m}}$ | 1 | $-\frac{b}{c}$ |
| homogeneous class D | $y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)$ | $x^{2}$ | $x y$ |
| First order <br> form ID 1 | $y^{2}=g(x) e^{h(x)+b y}+f(x)$ | $\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}$ | $\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}$ |
| polynomial type ode | $y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}$ | $\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ | $\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}$ |
| Bernoulli ode | $y^{\prime}=f(x) y+g(x) y^{n}$ | 0 | $e^{-\int(n-1) f(x) d x} y^{n}$ |
| Reduced Riccati | $y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}$ | 0 | $e^{-\int f_{1} d x}$ |

The above table shows that

$$
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=x^{2} \tag{A1}
\end{align*}
$$

The next step is to determine the canonical coordinates $R, S$. The canonical coordinates map $(x, y) \rightarrow(R, S)$ where $(R, S)$ are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable $R$ in the
canonical coordinates, where $S(R)$. Since $\xi=0$ then in this special case

$$
R=x
$$

$S$ is found from

$$
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{x^{2}} d y
\end{aligned}
$$

Which results in

$$
S=\frac{y}{x^{2}}
$$

Now that $R, S$ are found, we need to setup the ode in these coordinates. This is done by evaluating

$$
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
$$

Where in the above $R_{x}, R_{y}, S_{x}, S_{y}$ are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$
\omega(x, y)=\frac{-\sqrt{3} x+2 y}{x}
$$

Evaluating all the partial derivatives gives

$$
\begin{aligned}
& R_{x}=1 \\
& R_{y}=0 \\
& S_{x}=-\frac{2 y}{x^{3}} \\
& S_{y}=\frac{1}{x^{2}}
\end{aligned}
$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$
\begin{equation*}
\frac{d S}{d R}=-\frac{\sqrt{3}}{x^{2}} \tag{2~A}
\end{equation*}
$$

We now need to express the RHS as function of $R$ only. This is done by solving for $x, y$ in terms of $R, S$ from the result obtained earlier and simplifying. This gives

$$
\frac{d S}{d R}=-\frac{\sqrt{3}}{R^{2}}
$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates $R, S$. Integrating the above gives

$$
\begin{equation*}
S(R)=\frac{\sqrt{3}}{R}+c_{1} \tag{4}
\end{equation*}
$$

To complete the solution, we just need to transform (4) back to $x, y$ coordinates. This results in

$$
\frac{y}{x^{2}}=\frac{\sqrt{3}}{x}+c_{1}
$$

Which simplifies to

$$
\frac{y}{x^{2}}=\frac{\sqrt{3}}{x}+c_{1}
$$

Which gives

$$
y=x\left(c_{1} x+\sqrt{3}\right)
$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

| Original ode in $x, y$ coordinates | Canonical coordinates transformation | ODE in canonical coordinates $(R, S)$ |
| :---: | :---: | :---: |
| $\frac{d y}{d x}=\frac{-\sqrt{3} x+2 y}{x}$ |  | $\frac{d S}{d R}=-\frac{\sqrt{3}}{R^{2}}$ |
|  |  |  |
|  |  | $\xrightarrow[\rightarrow \rightarrow-\infty]{ } \rightarrow+\infty$ |
|  |  | $\rightarrow \rightarrow$ Pr $\rightarrow$, ${ }^{\text {a }}$ |
|  |  | 1 |
| $\frac{10 y y}{}$ |  |  |
|  | $S=\frac{y}{x^{2}}$ | $\xrightarrow[\rightarrow \rightarrow-\infty]{ }+\infty$ |
|  |  | , |
| tarners |  |  |
|  |  | C. |

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x\left(c_{1} x+\sqrt{3}\right) \tag{1}
\end{equation*}
$$



Figure 138: Slope field plot

Verification of solutions

$$
y=x\left(c_{1} x+\sqrt{3}\right)
$$

Verified OK.

### 1.71.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form

$$
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
$$

We assume there exists a function $\phi(x, y)=c$ where $c$ is constant, that satisfies the ode. Taking derivative of $\phi$ w.r.t. $x$ gives

$$
\frac{d}{d x} \phi(x, y)=0
$$

Hence

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
$$

Comparing ( $\mathrm{A}, \mathrm{B}$ ) shows that

$$
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
$$

But since $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1~A}
\end{equation*}
$$

Therefore

$$
\begin{align*}
\mathrm{d} y & =\left(\frac{2 y}{x}-\sqrt{3}\right) \mathrm{d} x \\
\left(-\frac{2 y}{x}+\sqrt{3}\right) \mathrm{d} x+\mathrm{d} y & =0 \tag{2~A}
\end{align*}
$$

Comparing (1A) and (2A) shows that

$$
\begin{aligned}
& M(x, y)=-\frac{2 y}{x}+\sqrt{3} \\
& N(x, y)=1
\end{aligned}
$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
$$

Using result found above gives

$$
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{2 y}{x}+\sqrt{3}\right) \\
& =-\frac{2}{x}
\end{aligned}
$$

And

$$
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(1) \\
& =0
\end{aligned}
$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =1\left(\left(-\frac{2}{x}\right)-(0)\right) \\
& =-\frac{2}{x}
\end{aligned}
$$

Since $A$ does not depend on $y$, then it can be used to find an integrating factor. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-\frac{2}{x} \mathrm{~d} x}
\end{aligned}
$$

The result of integrating gives

$$
\begin{aligned}
\mu & =e^{-2 \ln (x)} \\
& =\frac{1}{x^{2}}
\end{aligned}
$$

$M$ and $N$ are multiplied by this integrating factor, giving new $M$ and new $N$ which are called $\bar{M}$ and $\bar{N}$ for now so not to confuse them with the original $M$ and $N$.

$$
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{x^{2}}\left(-\frac{2 y}{x}+\sqrt{3}\right) \\
& =\frac{\sqrt{3} x-2 y}{x^{3}}
\end{aligned}
$$

And

$$
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{x^{2}}(1) \\
& =\frac{1}{x^{2}}
\end{aligned}
$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x} & =0 \\
\left(\frac{\sqrt{3} x-2 y}{x^{3}}\right)+\left(\frac{1}{x^{2}}\right) \frac{\mathrm{d} y}{\mathrm{~d} x} & =0
\end{aligned}
$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
$$

Integrating (1) w.r.t. $x$ gives

$$
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{\sqrt{3} x-2 y}{x^{3}} \mathrm{~d} x \\
\phi & =\frac{-\sqrt{3} x+y}{x^{2}}+f(y) \tag{3}
\end{align*}
$$

Where $f(y)$ is used for the constant of integration since $\phi$ is a function of both $x$ and $y$. Taking derivative of equation (3) w.r.t $y$ gives

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}=\frac{1}{x^{2}}+f^{\prime}(y) \tag{4}
\end{equation*}
$$

But equation (2) says that $\frac{\partial \phi}{\partial y}=\frac{1}{x^{2}}$. Therefore equation (4) becomes

$$
\begin{equation*}
\frac{1}{x^{2}}=\frac{1}{x^{2}}+f^{\prime}(y) \tag{5}
\end{equation*}
$$

Solving equation (5) for $f^{\prime}(y)$ gives

$$
f^{\prime}(y)=0
$$

Therefore

$$
f(y)=c_{1}
$$

Where $c_{1}$ is constant of integration. Substituting this result for $f(y)$ into equation (3) gives $\phi$

$$
\phi=\frac{-\sqrt{3} x+y}{x^{2}}+c_{1}
$$

But since $\phi$ itself is a constant function, then let $\phi=c_{2}$ where $c_{2}$ is new constant and combining $c_{1}$ and $c_{2}$ constants into new constant $c_{1}$ gives the solution as

$$
c_{1}=\frac{-\sqrt{3} x+y}{x^{2}}
$$

The solution becomes

$$
y=x\left(c_{1} x+\sqrt{3}\right)
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=x\left(c_{1} x+\sqrt{3}\right) \tag{1}
\end{equation*}
$$



Figure 139: Slope field plot

## Verification of solutions

$$
y=x\left(c_{1} x+\sqrt{3}\right)
$$

Verified OK.

### 1.71.5 Maple step by step solution

Let's solve
$y^{\prime}-\frac{2 y}{x}=-\sqrt{3}$

- Highest derivative means the order of the ODE is 1
$y^{\prime}$
- Isolate the derivative
$y^{\prime}=\frac{2 y}{x}-\sqrt{3}$
- Group terms with $y$ on the lhs of the ODE and the rest on the rhs of the ODE $y^{\prime}-\frac{2 y}{x}=-\sqrt{3}$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(y^{\prime}-\frac{2 y}{x}\right)=-\mu(x) \sqrt{3}$
- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) y)$
$\mu(x)\left(y^{\prime}-\frac{2 y}{x}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=-\frac{2 \mu(x)}{x}$
- Solve to find the integrating factor
$\mu(x)=\frac{1}{x^{2}}$
- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int-\mu(x) \sqrt{3} d x+c_{1}$
- Evaluate the integral on the lhs
$\mu(x) y=\int-\mu(x) \sqrt{3} d x+c_{1}$
- $\quad$ Solve for $y$
$y=\frac{\int-\mu(x) \sqrt{3} d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\frac{1}{x^{2}}$

$$
y=x^{2}\left(\int-\frac{\sqrt{3}}{x^{2}} d x+c_{1}\right)
$$

- Evaluate the integrals on the rhs

$$
y=x^{2}\left(\frac{\sqrt{3}}{x}+c_{1}\right)
$$

- Simplify

$$
y=x\left(c_{1} x+\sqrt{3}\right)
$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

$\checkmark$ Solution by Maple
Time used: 0.015 (sec). Leaf size: 13

```
dsolve(diff(y(x),x)=2*y(x)/x-sqrt(3),y(x), singsol=all)
```

$$
y(x)=\left(c_{1} x+\sqrt{3}\right) x
$$

$\sqrt{ }$ Solution by Mathematica
Time used: 0.066 (sec). Leaf size: 17
DSolve[y' $[x]==2 * y[x] / x-S q r t[3], y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow x\left(\sqrt{3}+c_{1} x\right)
$$

### 1.72 problem 116

1.72.1 Maple step by step solution

799
Internal problem ID [12489]
Internal file name [OUTPUT/11141_Monday_October_16_2023_09_51_52_PM_66535871/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 116.
ODE order: 3.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_linear_constant__coefficients_ODE"

Maple gives the following as the ode type
[[_3rd_order, _missing_x]]

$$
y^{\prime \prime \prime}-2 y^{\prime \prime}-y^{\prime}+2 y=0
$$

The characteristic equation is

$$
\lambda^{3}-2 \lambda^{2}-\lambda+2=0
$$

The roots of the above equation are

$$
\begin{aligned}
& \lambda_{1}=1 \\
& \lambda_{2}=2 \\
& \lambda_{3}=-1
\end{aligned}
$$

Therefore the homogeneous solution is

$$
y_{h}(x)=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{3}
$$

The fundamental set of solutions for the homogeneous solution are the following

$$
\begin{aligned}
& y_{1}=\mathrm{e}^{-x} \\
& y_{2}=\mathrm{e}^{x} \\
& y_{3}=\mathrm{e}^{2 x}
\end{aligned}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{3} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{3}
$$

Verified OK.

### 1.72.1 Maple step by step solution

Let's solve

$$
y^{\prime \prime \prime}-2 y^{\prime \prime}-y^{\prime}+2 y=0
$$

- Highest derivative means the order of the ODE is 3

$$
y^{\prime \prime \prime}
$$

Convert linear ODE into a system of first order ODEs

- Define new variable $y_{1}(x)$

$$
y_{1}(x)=y
$$

- Define new variable $y_{2}(x)$

$$
y_{2}(x)=y^{\prime}
$$

- Define new variable $y_{3}(x)$

$$
y_{3}(x)=y^{\prime \prime}
$$

- Isolate for $y_{3}^{\prime}(x)$ using original ODE

$$
y_{3}^{\prime}(x)=2 y_{3}(x)+y_{2}(x)-2 y_{1}(x)
$$

Convert linear ODE into a system of first order ODEs

$$
\left[y_{2}(x)=y_{1}^{\prime}(x), y_{3}(x)=y_{2}^{\prime}(x), y_{3}^{\prime}(x)=2 y_{3}(x)+y_{2}(x)-2 y_{1}(x)\right]
$$

- Define vector

$$
\vec{y}(x)=\left[\begin{array}{l}
y_{1}(x) \\
y_{2}(x) \\
y_{3}(x)
\end{array}\right]
$$

- System to solve

$$
\vec{y}^{\prime}(x)=\left[\begin{array}{ccc}
0 & 1 & 0 \\
0 & 0 & 1 \\
-2 & 1 & 2
\end{array}\right] \cdot \vec{y}(x)
$$

- Define the coefficient matrix

$$
A=\left[\begin{array}{ccc}
0 & 1 & 0 \\
0 & 0 & 1 \\
-2 & 1 & 2
\end{array}\right]
$$

- Rewrite the system as

$$
\vec{y}^{\prime}(x)=A \cdot \vec{y}(x)
$$

- To solve the system, find the eigenvalues and eigenvectors of $A$
- $\quad$ Eigenpairs of $A$

$$
\left[\left[-1,\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]\right],\left[1,\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]\right],\left[2,\left[\begin{array}{c}
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]\right]\right]
$$

- Consider eigenpair

$$
\left[-1,\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]\right]
$$

- $\quad$ Solution to homogeneous system from eigenpair

$$
\vec{y}_{1}=\mathrm{e}^{-x} \cdot\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]
$$

- Consider eigenpair
$\left[1,\left[\begin{array}{l}1 \\ 1 \\ 1\end{array}\right]\right]$
- Solution to homogeneous system from eigenpair

$$
\vec{y}_{2}=\mathrm{e}^{x} \cdot\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]
$$

- Consider eigenpair

$$
\left[2,\left[\begin{array}{c}
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]\right]
$$

- Solution to homogeneous system from eigenpair

$$
\vec{y}_{3}=\mathrm{e}^{2 x} \cdot\left[\begin{array}{c}
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]
$$

- General solution to the system of ODEs

$$
\vec{y}=c_{1} \vec{y}_{1}+c_{2} \vec{y}_{2}+c_{3} \vec{y}_{3}
$$

- $\quad$ Substitute solutions into the general solution

$$
\vec{y}=c_{1} \mathrm{e}^{-x} \cdot\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]+c_{2} \mathrm{e}^{x} \cdot\left[\begin{array}{c}
1 \\
1 \\
1
\end{array}\right]+\mathrm{e}^{2 x} c_{3} \cdot\left[\begin{array}{c}
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]
$$

- First component of the vector is the solution to the ODE $y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+\frac{\mathrm{e}^{2 x} c_{3}}{4}$

Maple trace

```
`Methods for third order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 21
dsolve(diff $(y(x), x \$ 3)-2 * \operatorname{diff}(y(x), x \$ 2)-\operatorname{diff}(y(x), x)+2 * y(x)=0, y(x)$, singsol=all)

$$
y(x)=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+c_{3} \mathrm{e}^{2 x}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.006 (sec). Leaf size: 28
DSolve[y' ' ' $[x]-2 * y$ '' $[x]-y$ ' $[x]+2 * y[x]==0, y[x], x$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow c_{1} e^{-x}+c_{2} e^{x}+c_{3} e^{2 x}
$$

### 1.73 problem 117

1.73.1 Solving as second order integrable as is ode . . . . . . . . . . . 803
1.73.2 Solving as second order ode missing y ode . . . . . . . . . . . . 805
1.73.3 Solving as second order ode missing $x$ ode . . . . . . . . . . . . 806
1.73.4 Solving as exact nonlinear second order ode ode . . . . . . . . . 808
1.73.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 810

Internal problem ID [12490]
Internal file name [OUTPUT/11142_Monday_October_16_2023_09_51_52_PM_94770677/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 117.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "second_order_integrable_as_is", "second_order_ode_missing_x", "second_order_ode_missing_y", "exact nonlinear second order ode"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x], [_2nd_order, _reducible, _mu_poly_yn ]]
```

$$
y^{\prime \prime}-\frac{1}{2 y^{\prime}}=0
$$

### 1.73.1 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t $x$ gives

$$
\begin{aligned}
& \int_{y^{\prime 2}=x+c_{1}} 2 y^{\prime} y^{\prime \prime} d x=\int 1 d x \\
& y^{\prime 2}
\end{aligned}
$$

Which is now solved for $y$. Solving the given ode for $y^{\prime}$ results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$
\begin{align*}
& y^{\prime}=\sqrt{x+c_{1}}  \tag{1}\\
& y^{\prime}=-\sqrt{x+c_{1}} \tag{2}
\end{align*}
$$

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives

$$
\begin{aligned}
y & =\int \sqrt{x+c_{1}} \mathrm{~d} x \\
& =\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\end{aligned}
$$

Solving equation (2)
Integrating both sides gives

$$
\begin{aligned}
y & =\int-\sqrt{x+c_{1}} \mathrm{~d} x \\
& =-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}  \tag{1}\\
& y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3} \tag{2}
\end{align*}
$$

Verification of solutions

$$
y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
$$

Verified OK.

$$
y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
$$

Verified OK.

### 1.73.2 Solving as second order ode missing y ode

This is second order ode with missing dependent variable $y$. Let

$$
p(x)=y^{\prime}
$$

Then

$$
p^{\prime}(x)=y^{\prime \prime}
$$

Hence the ode becomes

$$
2 p(x) p^{\prime}(x)-1=0
$$

Which is now solve for $p(x)$ as first order ode. Integrating both sides gives

$$
\begin{aligned}
\int 2 p d p & =x+c_{1} \\
p^{2} & =x+c_{1}
\end{aligned}
$$

Solving for $p$ gives these solutions

$$
\begin{aligned}
& p_{1}=\sqrt{x+c_{1}} \\
& p_{2}=-\sqrt{x+c_{1}}
\end{aligned}
$$

For solution (1) found earlier, since $p=y^{\prime}$ then the new first order ode to solve is

$$
y^{\prime}=\sqrt{x+c_{1}}
$$

Integrating both sides gives

$$
\begin{aligned}
y & =\int \sqrt{x+c_{1}} \mathrm{~d} x \\
& =\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\end{aligned}
$$

Since $p=y^{\prime}$ then the new first order ode to solve is

$$
y^{\prime}=-\sqrt{x+c_{1}}
$$

Integrating both sides gives

$$
\begin{aligned}
y & =\int-\sqrt{x+c_{1}} \mathrm{~d} x \\
& =-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}  \tag{1}\\
& y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3} \tag{2}
\end{align*}
$$

Verification of solutions

$$
y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
$$

Verified OK.

$$
y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
$$

Verified OK.

### 1.73.3 Solving as second order ode missing $x$ ode

This is missing independent variable second order ode. Solved by reduction of order by using substitution which makes the dependent variable $y$ an independent variable. Using

$$
y^{\prime}=p(y)
$$

Then

$$
\begin{aligned}
y^{\prime \prime} & =\frac{d p}{d x} \\
& =\frac{d y}{d x} \frac{d p}{d y} \\
& =p \frac{d p}{d y}
\end{aligned}
$$

Hence the ode becomes

$$
2 p(y)^{2}\left(\frac{d}{d y} p(y)\right)=1
$$

Which is now solved as first order ode for $p(y)$. Integrating both sides gives

$$
\begin{aligned}
\int 2 p^{2} d p & =y+c_{1} \\
\frac{2 p^{3}}{3} & =y+c_{1}
\end{aligned}
$$

Solving for $p$ gives these solutions

$$
\begin{aligned}
& p_{1}=\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{2} \\
& p_{2}=-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}-\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4} \\
& p_{3}=-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}+\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}
\end{aligned}
$$

For solution (1) found earlier, since $p=y^{\prime}$ then we now have a new first order ode to solve which is

$$
y^{\prime}=\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{2}
$$

Integrating both sides gives

$$
\begin{aligned}
\int \frac{2}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}} d y & =\int d x \\
\frac{3 y+3 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}} & =x+c_{2}
\end{aligned}
$$

For solution (2) found earlier, since $p=y^{\prime}$ then we now have a new first order ode to solve which is

$$
y^{\prime}=-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}-\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}
$$

Integrating both sides gives

$$
\begin{aligned}
& \int \frac{1}{-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}-\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}} d y \\
&=\int d x \\
&-\frac{6\left(y+c_{1}\right)}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(1+i \sqrt{3})}=c_{3}+x
\end{aligned}
$$

For solution (3) found earlier, since $p=y^{\prime}$ then we now have a new first order ode to solve which is

$$
y^{\prime}=-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}+\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}
$$

Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}+\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}} d y & =\int d x \\
\frac{6 y+6 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(i \sqrt{3}-1)} & =x+c_{4}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
\frac{3 y+3 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}} & =x+c_{2}  \tag{1}\\
-\frac{6\left(y+c_{1}\right)}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(1+i \sqrt{3})} & =c_{3}+x  \tag{2}\\
\frac{6 y+6 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(i \sqrt{3}-1)} & =x+c_{4} \tag{3}
\end{align*}
$$

Verification of solutions

$$
\frac{3 y+3 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}=x+c_{2}
$$

Verified OK.

$$
-\frac{6\left(y+c_{1}\right)}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(1+i \sqrt{3})}=c_{3}+x
$$

Verified OK.

$$
\frac{6 y+6 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(i \sqrt{3}-1)}=x+c_{4}
$$

Verified OK.

### 1.73.4 Solving as exact nonlinear second order ode ode

An exact non-linear second order ode has the form

$$
a_{2}\left(x, y, y^{\prime}\right) y^{\prime \prime}+a_{1}\left(x, y, y^{\prime}\right) y^{\prime}+a_{0}\left(x, y, y^{\prime}\right)=0
$$

Where the following conditions are satisfied

$$
\begin{aligned}
\frac{\partial a_{2}}{\partial y} & =\frac{\partial a_{1}}{\partial y^{\prime}} \\
\frac{\partial a_{2}}{\partial x} & =\frac{\partial a_{0}}{\partial y^{\prime}} \\
\frac{\partial a_{1}}{\partial x} & =\frac{\partial a_{0}}{\partial y}
\end{aligned}
$$

Looking at the the ode given we see that

$$
\begin{aligned}
& a_{2}=2 y^{\prime} \\
& a_{1}=0 \\
& a_{0}=-1
\end{aligned}
$$

Applying the conditions to the above shows this is a nonlinear exact second order ode. Therefore it can be reduced to first order ode given by

$$
\begin{aligned}
& \int a_{2} d y^{\prime}+\int a_{1} d y+\int a_{0} d x=c_{1} \\
& \int 2 y^{\prime} d y^{\prime}+\int 0 d y+\int-1 d x=c_{1}
\end{aligned}
$$

Which results in

$$
y^{\prime 2}-x=c_{1}
$$

Which is now solved Solving the given ode for $y^{\prime}$ results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$
\begin{align*}
y^{\prime} & =\sqrt{x+c_{1}}  \tag{1}\\
y^{\prime} & =-\sqrt{x+c_{1}} \tag{2}
\end{align*}
$$

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives

$$
\begin{aligned}
y & =\int \sqrt{x+c_{1}} \mathrm{~d} x \\
& =\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\end{aligned}
$$

Solving equation (2)

Integrating both sides gives

$$
\begin{aligned}
y & =\int-\sqrt{x+c_{1}} \mathrm{~d} x \\
& =-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}  \tag{1}\\
& y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3} \tag{2}
\end{align*}
$$

Verification of solutions

$$
y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
$$

Verified OK.

$$
y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
$$

Verified OK.

### 1.73.5 Maple step by step solution

Let's solve

$$
2 y^{\prime} y^{\prime \prime}=1
$$

- Highest derivative means the order of the ODE is 2

$$
y^{\prime \prime}
$$

- Make substitution $u=y^{\prime}$ to reduce order of ODE

$$
2 u(x) u^{\prime}(x)=1
$$

- Integrate both sides with respect to $x$
$\int 2 u(x) u^{\prime}(x) d x=\int 1 d x+c_{1}$
- Evaluate integral
$u(x)^{2}=x+c_{1}$
- $\quad$ Solve for $u(x)$

$$
\left\{u(x)=\sqrt{x+c_{1}}, u(x)=-\sqrt{x+c_{1}}\right\}
$$

- $\quad$ Solve 1st ODE for $u(x)$

$$
u(x)=\sqrt{x+c_{1}}
$$

- Make substitution $u=y^{\prime}$

$$
y^{\prime}=\sqrt{x+c_{1}}
$$

- Integrate both sides to solve for $y$

$$
\int y^{\prime} d x=\int \sqrt{x+c_{1}} d x+c_{2}
$$

- Compute integrals

$$
y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
$$

- $\quad$ Solve 2nd ODE for $u(x)$

$$
u(x)=-\sqrt{x+c_{1}}
$$

- Make substitution $u=y^{\prime}$

$$
y^{\prime}=-\sqrt{x+c_{1}}
$$

- Integrate both sides to solve for $y$

$$
\int y^{\prime} d x=\int-\sqrt{x+c_{1}} d x+c_{2}
$$

- Compute integrals

$$
y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
$$

## Maple trace

-Methods for second order ODEs:
--- Trying classification methods ---
trying 2nd order Liouville
trying 2nd order WeierstrassP
trying 2nd order JacobiSN
differential order: 2; trying a linearization to 3rd order
trying 2nd order ODE linearizable_by_differentiation
trying 2nd order, 2 integrating factors of the form $m u(x, y)$
trying differential order: 2; missing variables
, --> Computing symmetries using: way $=3$
$\rightarrow$ Calling odsolve with the ODE`, \(\operatorname{diff}\left(\_b\left(\_a\right), \quad a\right)=(1 / 2) / \_b\left(\_a\right), \quad b\left(\_a\right), H I N T=[[1,0]\), symmetry methods on request , `1st order, trying reduction of order with given symmetries:`[1, 0], [_a, 1/2*_b]
$\checkmark$ Solution by Maple
Time used: 0.047 (sec). Leaf size: 39

```
dsolve(diff(y(x),x$2)=1/(2*\operatorname{diff}(y(x),x)),y(x), singsol=all)
```

$$
\begin{aligned}
& y(x)=\frac{\left(2 x+2 c_{1}\right) \sqrt{c_{1}+x}}{3}+c_{2} \\
& y(x)=\frac{\left(-2 x-2 c_{1}\right) \sqrt{c_{1}+x}}{3}+c_{2}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.023 (sec). Leaf size: 43
DSolve[y''[x]==1/(2*y'[x]),y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
\begin{aligned}
& y(x) \rightarrow c_{2}-\frac{2}{3}\left(x+2 c_{1}\right)^{3 / 2} \\
& y(x) \rightarrow \frac{2}{3}\left(x+2 c_{1}\right)^{3 / 2}+c_{2}
\end{aligned}
$$

### 1.74 problem 118

Internal problem ID [12491]
Internal file name [OUTPUT/11143_Monday_October_16_2023_09_51_55_PM_10421388/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 118.
ODE order: 3.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_missing_y"
Maple gives the following as the ode type
[[_3rd_order, _quadrature]]

$$
x y^{\prime \prime \prime}=2
$$

Since $y$ is missing from the ode then we can use the substitution $y^{\prime}=v(x)$ to reduce the order by one. The ODE becomes

$$
v^{\prime \prime}(x) x=0
$$

Integrating twice gives the solution

$$
v(x)=c_{1} x+c_{2}
$$

But since $y^{\prime}=v(x)$ then we now need to solve the ode $y^{\prime}=c_{1} x+c_{2}$. Integrating both sides gives

$$
\begin{aligned}
y & =\int c_{1} x+c_{2} \mathrm{~d} x \\
& =\frac{1}{2} c_{1} x^{2}+c_{2} x+c_{3}
\end{aligned}
$$

This is higher order nonhomogeneous ODE. Let the solution be

$$
y=y_{h}+y_{p}
$$

Where $y_{h}$ is the solution to the homogeneous ODE And $y_{p}$ is a particular solution to the nonhomogeneous ODE. $y_{h}$ is the solution to

$$
x y^{\prime \prime \prime}=0
$$

Let the particular solution be

$$
y_{p}=U_{1} y_{1}+U_{2} y_{2}+U_{3} y_{3}
$$

Where $y_{i}$ are the basis solutions found above for the homogeneous solution $y_{h}$ and $U_{i}(x)$ are functions to be determined as follows

$$
U_{i}=(-1)^{n-i} \int \frac{F(x) W_{i}(x)}{a W(x)} d x
$$

Where $W(x)$ is the Wronskian and $W_{i}(x)$ is the Wronskian that results after deleting the last row and the $i$-th column of the determinant and $n$ is the order of the ODE or equivalently, the number of basis solutions, and $a$ is the coefficient of the leading derivative in the ODE, and $F(x)$ is the RHS of the ODE. Therefore, the first step is to find the Wronskian $W(x)$. This is given by

$$
W(x)=\left|\begin{array}{lll}
y_{1} & y_{2} & y_{3} \\
y_{1}^{\prime} & y_{2}^{\prime} & y_{3}^{\prime} \\
y_{1}^{\prime \prime} & y_{2}^{\prime \prime} & y_{3}^{\prime \prime}
\end{array}\right|
$$

Substituting the fundamental set of solutions $y_{i}$ found above in the Wronskian gives

$$
\begin{aligned}
& W=\left[\begin{array}{ccc}
1 & x & x^{2} \\
0 & 1 & 2 x \\
0 & 0 & 2
\end{array}\right] \\
&|W|=2
\end{aligned}
$$

The determinant simplifies to

$$
|W|=2
$$

Now we determine $W_{i}$ for each $U_{i}$.

$$
\begin{aligned}
W_{1}(x) & =\operatorname{det}\left[\begin{array}{ll}
x & x^{2} \\
1 & 2 x
\end{array}\right] \\
& =x^{2} \\
W_{2}(x) & =\operatorname{det}\left[\begin{array}{ll}
1 & x^{2} \\
0 & 2 x
\end{array}\right] \\
& =2 x
\end{aligned}
$$

$$
\begin{aligned}
W_{3}(x) & =\operatorname{det}\left[\begin{array}{ll}
1 & x \\
0 & 1
\end{array}\right] \\
& =1
\end{aligned}
$$

Now we are ready to evaluate each $U_{i}(x)$.

$$
\begin{aligned}
U_{1} & =(-1)^{3-1} \int \frac{F(x) W_{1}(x)}{a W(x)} d x \\
& =(-1)^{2} \int \frac{(2)\left(x^{2}\right)}{(x)(2)} d x \\
& =\int \frac{2 x^{2}}{2 x} d x \\
& =\int(x) d x \\
& =\frac{x^{2}}{2}
\end{aligned}
$$

$$
U_{2}=(-1)^{3-2} \int \frac{F(x) W_{2}(x)}{a W(x)} d x
$$

$$
=(-1)^{1} \int \frac{(2)(2 x)}{(x)(2)} d x
$$

$$
=-\int \frac{4 x}{2 x} d x
$$

$$
=-\int(2) d x
$$

$$
=-2 x
$$

$$
U_{3}=(-1)^{3-3} \int \frac{F(x) W_{3}(x)}{a W(x)} d x
$$

$$
=(-1)^{0} \int \frac{(2)(1)}{(x)(2)} d x
$$

$$
=\int \frac{2}{2 x} d x
$$

$$
=\int\left(\frac{1}{x}\right) d x
$$

$$
=\ln (x)
$$

Now that all the $U_{i}$ functions have been determined, the particular solution is found from

$$
y_{p}=U_{1} y_{1}+U_{2} y_{2}+U_{3} y_{3}
$$

Hence

$$
\begin{aligned}
y_{p} & =\left(\frac{x^{2}}{2}\right) \\
& +(-2 x)(x) \\
& +(\ln (x))\left(x^{2}\right)
\end{aligned}
$$

Therefore the particular solution is

$$
y_{p}=x^{2}\left(-\frac{3}{2}+\ln (x)\right)
$$

Therefore the general solution is

$$
\begin{aligned}
y & =y_{h}+y_{p} \\
& =(y \\
& \left.=\frac{1}{2} c_{1} x^{2}+c_{2} x+c_{3}\right)+\left(x^{2}\left(-\frac{3}{2}+\ln (x)\right)\right)
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\frac{c_{1} x^{2}}{2}+c_{2} x+c_{3}+x^{2}\left(-\frac{3}{2}+\ln (x)\right) \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\frac{c_{1} x^{2}}{2}+c_{2} x+c_{3}+x^{2}\left(-\frac{3}{2}+\ln (x)\right)
$$

Verified OK.
Maple trace

```
`Methods for third order ODEs:
--- Trying classification methods ---
trying a quadrature
<- quadrature successful`
```

$\checkmark$ Solution by Maple
Time used: 0.016 (sec). Leaf size: 23
dsolve( $x * \operatorname{diff}(y(x), x \$ 3)=2, y(x)$, singsol=all)

$$
y(x)=\ln (x) x^{2}+\frac{\left(c_{1}-3\right) x^{2}}{2}+c_{2} x+c_{3}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.005 (sec). Leaf size: 28
DSolve[x*y' ' ' $[\mathrm{x}]==2, \mathrm{y}[\mathrm{x}], \mathrm{x}$, IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow x^{2} \log (x)+\left(-\frac{3}{2}+c_{3}\right) x^{2}+c_{2} x+c_{1}
$$

### 1.75 problem 120

1.75.1 Solving as second order linear constant coeff ode . . . . . . . . 818
1.75.2 Solving as second order ode can be made integrable ode . . . . 820
1.75.3 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 821
1.75.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 825

Internal problem ID [12492]
Internal file name [OUTPUT/11144_Monday_October_16_2023_09_51_56_PM_19853961/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 120.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second__order_linear_constant_coeff", "second__order_ode_can_bbe_made_integrable"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$
y^{\prime \prime}-y a^{2}=0
$$

### 1.75.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
$$

Where in the above $A=1, B=0, C=-a^{2}$. Let the solution be $y=e^{\lambda x}$. Substituting this into the ODE gives

$$
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-a^{2} \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
$$

Since exponential function is never zero, then dividing $\mathrm{Eq}(2)$ throughout by $e^{\lambda x}$ gives

$$
\begin{equation*}
-a^{2}+\lambda^{2}=0 \tag{2}
\end{equation*}
$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula

$$
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
$$

Substituting $A=1, B=0, C=-a^{2}$ into the above gives

$$
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)\left(-a^{2}\right)} \\
& = \pm \sqrt{a^{2}}
\end{aligned}
$$

Hence

$$
\begin{aligned}
& \lambda_{1}=+\sqrt{a^{2}} \\
& \lambda_{2}=-\sqrt{a^{2}}
\end{aligned}
$$

Which simplifies to

$$
\begin{aligned}
& \lambda_{1}=\sqrt{a^{2}} \\
& \lambda_{2}=-\sqrt{a^{2}}
\end{aligned}
$$

Since roots are real and distinct, then the solution is

$$
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{\left(\sqrt{a^{2}}\right) x}+c_{2} e^{\left(-\sqrt{a^{2}}\right) x}
\end{aligned}
$$

Or

$$
y=c_{1} \mathrm{e}^{\sqrt{a^{2}} x}+c_{2} \mathrm{e}^{-\sqrt{a^{2}} x}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} \mathrm{e}^{\sqrt{a^{2}} x}+c_{2} \mathrm{e}^{-\sqrt{a^{2}} x} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=c_{1} \mathrm{e}^{\sqrt{a^{2}} x}+c_{2} \mathrm{e}^{-\sqrt{a^{2}} x}
$$

Verified OK.

### 1.75.2 Solving as second order ode can be made integrable ode

Multiplying the ode by $y^{\prime}$ gives

$$
y^{\prime} y^{\prime \prime}-y^{\prime} a^{2} y=0
$$

Integrating the above w.r.t $x$ gives

$$
\begin{gathered}
\int\left(y^{\prime} y^{\prime \prime}-y^{\prime} a^{2} y\right) d x=0 \\
\frac{y^{\prime 2}}{2}-\frac{y^{2} a^{2}}{2}=c_{2}
\end{gathered}
$$

Which is now solved for $y$. Solving the given ode for $y^{\prime}$ results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$
\begin{align*}
& y^{\prime}=\sqrt{y^{2} a^{2}+2 c_{1}}  \tag{1}\\
& y^{\prime}=-\sqrt{y^{2} a^{2}+2 c_{1}} \tag{2}
\end{align*}
$$

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{\sqrt{y^{2} a^{2}+2 c_{1}}} d y & =\int d x \\
\frac{\ln \left(\frac{a^{2} y}{\sqrt{a^{2}}}+\sqrt{y^{2} a^{2}+2 c_{1}}\right)}{\sqrt{a^{2}}} & =x+c_{2}
\end{aligned}
$$

Raising both side to exponential gives

$$
\mathrm{e}^{\frac{\ln \left(\frac{a^{2} y}{\sqrt{a^{2}}}+\sqrt{y^{2} a^{2}+2 c_{1}}\right)}{\sqrt{a^{2}}}}=\mathrm{e}^{x+c_{2}}
$$

Which simplifies to

$$
\left(a y \operatorname{csgn}(a)+\sqrt{y^{2} a^{2}+2 c_{1}}\right)^{\frac{1}{\sqrt{a^{2}}}}=c_{3} \mathrm{e}^{x}
$$

Simplifying the solution $y=\frac{\operatorname{csgn}(a)\left(\left(c_{3} \mathrm{e}^{x}\right)^{\operatorname{csgn}(a) a}-2\left(c_{3} \mathrm{e}^{x}\right)^{-\operatorname{csgn}(a) a} c_{1}\right)}{2 a}$ to $y=\frac{\left(c_{3} \mathrm{e}^{x}\right)^{a}-2\left(c_{3} \mathrm{e}^{x}\right)^{-a} c_{1}}{2 a}$ Solving equation (2)

Integrating both sides gives

$$
\begin{aligned}
\int-\frac{1}{\sqrt{y^{2} a^{2}+2 c_{1}}} d y & =\int d x \\
-\frac{\ln \left(\frac{a^{2} y}{\sqrt{a^{2}}}+\sqrt{y^{2} a^{2}+2 c_{1}}\right)}{\sqrt{a^{2}}} & =x+c_{4}
\end{aligned}
$$

Raising both side to exponential gives

$$
\mathrm{e}^{-\frac{\ln \left(\frac{a^{2} y}{\sqrt{a^{2}}+\sqrt{y^{2} a^{2}+2 c_{1}}}\right)}{\sqrt{a^{2}}}}=\mathrm{e}^{x+c_{4}}
$$

Which simplifies to

$$
\left(a y \operatorname{csgn}(a)+\sqrt{y^{2} a^{2}+2 c_{1}}\right)^{-\frac{\operatorname{csgn}(a)}{a}}=c_{5} \mathrm{e}^{x}
$$

Simplifying the solution $y=-\frac{\operatorname{csgn}(a)\left(2\left(c_{5} \mathrm{e}^{x}\right)^{\operatorname{cssn}(a) a} c_{1}-\left(c_{5} \mathrm{e}^{x}\right)^{-\operatorname{csgn}(a) a}\right)}{2 a}$ to $y=-\frac{2\left(c_{5} \mathrm{e}^{x}\right)^{a} c_{1}-\left(c_{5} \mathrm{e}^{x}\right)^{-a}}{2 a}$ Summary
The solution(s) found are the following

$$
\begin{align*}
& y=\frac{\left(c_{3} \mathrm{e}^{x}\right)^{a}-2\left(c_{3} \mathrm{e}^{x}\right)^{-a} c_{1}}{2 a}  \tag{1}\\
& y=-\frac{2\left(c_{5} \mathrm{e}^{x}\right)^{a} c_{1}-\left(c_{5} \mathrm{e}^{x}\right)^{-a}}{2 a} \tag{2}
\end{align*}
$$

Verification of solutions

$$
y=\frac{\left(c_{3} \mathrm{e}^{x}\right)^{a}-2\left(c_{3} \mathrm{e}^{x}\right)^{-a} c_{1}}{2 a}
$$

Verified OK.

$$
y=-\frac{2\left(c_{5} \mathrm{e}^{x}\right)^{a} c_{1}-\left(c_{5} \mathrm{e}^{x}\right)^{-a}}{2 a}
$$

Verified OK.

### 1.75.3 Solving using Kovacic algorithm

Writing the ode as

$$
\begin{array}{r}
y^{\prime \prime}-y a^{2}=0 \\
A y^{\prime \prime}+B y^{\prime}+C y=0 \tag{2}
\end{array}
$$

Comparing (1) and (2) shows that

$$
\begin{align*}
A & =1 \\
B & =0  \tag{3}\\
C & =-a^{2}
\end{align*}
$$

Applying the Liouville transformation on the dependent variable gives

$$
z(x)=y e^{\int \frac{B}{2 A} d x}
$$

Then (2) becomes

$$
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
$$

Where $r$ is given by

$$
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
$$

Substituting the values of $A, B, C$ from (3) in the above and simplifying gives

$$
\begin{equation*}
r=\frac{a^{2}}{1} \tag{6}
\end{equation*}
$$

Comparing the above to (5) shows that

$$
\begin{aligned}
& s=a^{2} \\
& t=1
\end{aligned}
$$

Therefore eq. (4) becomes

$$
\begin{equation*}
z^{\prime \prime}(x)=\left(a^{2}\right) z(x) \tag{7}
\end{equation*}
$$

Equation (7) is now solved. After finding $z(x)$ then $y$ is found using the inverse transformation

$$
y=z(x) e^{-\int \frac{B}{2 A} d x}
$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of $r$ and the order of $r$ at $\infty$. The following table summarizes these cases.

| Case | Allowed pole order for $r$ | Allowed value for $\mathcal{O}(\infty)$ |
| :--- | :--- | :--- |
| 1 | $\{0,1,2,4,6,8, \cdots\}$ | $\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}$ |
| 2 | Need to have at least one pole that <br> is either order 2 or odd order greater <br> than 2. Any other pole order is <br> allowed as long as the above condi- | no condition |
| tion is satisfied. Hence the following |  |  |
| set of pole orders are all allowed. |  |  |
| $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}$. |  |  |
| 3 | $\{1,2\}$ | $\{2,3,4,5,6,7, \cdots\}$ |

Table 123: Necessary conditions for each Kovacic case

The order of $r$ at $\infty$ is the degree of $t$ minus the degree of $s$. Therefore

$$
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
$$

There are no poles in $r$. Therefore the set of poles $\Gamma$ is empty. Since there is no odd order pole larger than 2 and the order at $\infty$ is 0 then the necessary conditions for case one are met. Therefore

$$
L=[1]
$$

Since $r=a^{2}$ is not a function of $x$, then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z^{\prime \prime}=r z$ as one solution is

$$
z_{1}(x)=\mathrm{e}^{\sqrt{a^{2}} x}
$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in $y$ is found from

$$
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
$$

Since $B=0$ then the above reduces to

$$
\begin{aligned}
y_{1} & =z_{1} \\
& =\mathrm{e}^{\sqrt{a^{2}} x}
\end{aligned}
$$

Which simplifies to

$$
y_{1}=\mathrm{e}^{\sqrt{a^{2}} x}
$$

The second solution $y_{2}$ to the original ode is found using reduction of order

$$
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
$$

Since $B=0$ then the above becomes

$$
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\mathrm{e}^{\sqrt{a^{2}} x} \int \frac{1}{\mathrm{e}^{2 \sqrt{a^{2}} x}} d x \\
& =\mathrm{e}^{\sqrt{a^{2}} x}\left(-\frac{\operatorname{csgn}(a) \mathrm{e}^{-2 \operatorname{csgn}(a) a x}}{2 a}\right)
\end{aligned}
$$

Therefore the solution is

$$
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{\sqrt{a^{2}} x}\right)+c_{2}\left(\mathrm{e}^{\sqrt{a^{2}} x}\left(-\frac{\operatorname{csgn}(a) \mathrm{e}^{-2 \operatorname{csgn}(a) a x}}{2 a}\right)\right)
\end{aligned}
$$

Simplifying the solution $y=c_{1} \mathrm{e}^{\sqrt{a^{2}} x}-\frac{c_{2} \operatorname{csgn}(a) \mathrm{e}^{-\operatorname{csgn}(a) a x}}{2 a}$ to $y=c_{1} \mathrm{e}^{\sqrt{a^{2}} x}-\frac{c_{2} \mathrm{e}^{-a x}}{2 a}$ Summary
The solution(s) found are the following

$$
\begin{equation*}
y=c_{1} \mathrm{e}^{\sqrt{a^{2}} x}-\frac{c_{2} \mathrm{e}^{-a x}}{2 a} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=c_{1} \mathrm{e}^{\sqrt{a^{2}} x}-\frac{c_{2} \mathrm{e}^{-a x}}{2 a}
$$

Verified OK.

### 1.75.4 Maple step by step solution

Let's solve
$y^{\prime \prime}-y a^{2}=0$

- Highest derivative means the order of the ODE is 2

$$
y^{\prime \prime}
$$

- Characteristic polynomial of ODE
$-a^{2}+r^{2}=0$
- Factor the characteristic polynomial
$-(a-r)(a+r)=0$
- Roots of the characteristic polynomial
$r=(a,-a)$
- $\quad$ 1st solution of the ODE
$y_{1}(x)=\mathrm{e}^{a x}$
- 2nd solution of the ODE
$y_{2}(x)=\mathrm{e}^{-a x}$
- General solution of the ODE
$y=c_{1} y_{1}(x)+c_{2} y_{2}(x)$
- Substitute in solutions
$y=c_{1} \mathrm{e}^{a x}+c_{2} \mathrm{e}^{-a x}$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`
```

$\checkmark$ Solution by Maple
Time used: 0.0 (sec). Leaf size: 18
dsolve(diff $(y(x), x \$ 2)=a^{\wedge} 2 * y(x), y(x)$, singsol=all)

$$
y(x)=\mathrm{e}^{a x} c_{1}+c_{2} \mathrm{e}^{-a x}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.022 (sec). Leaf size: 23
DSolve[y''[x]==a^2*y[x],y[x],x,IncludeSingularSolutions $\rightarrow$ True]

$$
y(x) \rightarrow c_{1} e^{a x}+c_{2} e^{-a x}
$$

### 1.76 problem 121

1.76.1 Solving as second order ode missing x ode . . . . . . . . . . . . 827
1.76.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 829

Internal problem ID [12493]
Internal file name [OUTPUT/11145_Monday_October_16_2023_09_51_56_PM_61541068/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 121.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "second__order_ode_missing_x"
Maple gives the following as the ode type
[[_2nd_order, _missing_x], [_2nd_order, _reducible, _mu_x_y1]]

$$
y^{\prime \prime}-\frac{a}{y^{3}}=0
$$

### 1.76.1 Solving as second order ode missing $x$ ode

This is missing independent variable second order ode. Solved by reduction of order by using substitution which makes the dependent variable $y$ an independent variable. Using

$$
y^{\prime}=p(y)
$$

Then

$$
\begin{aligned}
y^{\prime \prime} & =\frac{d p}{d x} \\
& =\frac{d y}{d x} \frac{d p}{d y} \\
& =p \frac{d p}{d y}
\end{aligned}
$$

Hence the ode becomes

$$
p(y)\left(\frac{d}{d y} p(y)\right) y^{3}=a
$$

Which is now solved as first order ode for $p(y)$. In canonical form the ODE is

$$
\begin{aligned}
p^{\prime} & =F(y, p) \\
& =f(y) g(p) \\
& =\frac{a}{p y^{3}}
\end{aligned}
$$

Where $f(y)=\frac{a}{y^{3}}$ and $g(p)=\frac{1}{p}$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{\frac{1}{p}} d p & =\frac{a}{y^{3}} d y \\
\int \frac{1}{\frac{1}{p}} d p & =\int \frac{a}{y^{3}} d y \\
\frac{p^{2}}{2} & =-\frac{a}{2 y^{2}}+c_{1}
\end{aligned}
$$

The solution is

$$
\frac{p(y)^{2}}{2}+\frac{a}{2 y^{2}}-c_{1}=0
$$

For solution (1) found earlier, since $p=y^{\prime}$ then we now have a new first order ode to solve which is

$$
\frac{y^{\prime 2}}{2}+\frac{a}{2 y^{2}}-c_{1}=0
$$

Solving the given ode for $y^{\prime}$ results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$
\begin{align*}
& y^{\prime}=\frac{\sqrt{2 c_{1} y^{2}-a}}{y}  \tag{1}\\
& y^{\prime}=-\frac{\sqrt{2 c_{1} y^{2}-a}}{y} \tag{2}
\end{align*}
$$

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives

$$
\begin{aligned}
\int \frac{y}{\sqrt{2 c_{1} y^{2}-a}} d y & =\int d x \\
\frac{\sqrt{2 c_{1} y^{2}-a}}{2 c_{1}} & =x+c_{2}
\end{aligned}
$$

Solving equation (2)
Integrating both sides gives

$$
\begin{aligned}
\int-\frac{y}{\sqrt{2 c_{1} y^{2}-a}} d y & =\int d x \\
-\frac{\sqrt{2 c_{1} y^{2}-a}}{2 c_{1}} & =c_{3}+x
\end{aligned}
$$

Summary
The solution(s) found are the following

$$
\begin{align*}
\frac{\sqrt{2 c_{1} y^{2}-a}}{2 c_{1}} & =x+c_{2}  \tag{1}\\
-\frac{\sqrt{2 c_{1} y^{2}-a}}{2 c_{1}} & =c_{3}+x \tag{2}
\end{align*}
$$

Verification of solutions

$$
\frac{\sqrt{2 c_{1} y^{2}-a}}{2 c_{1}}=x+c_{2}
$$

Verified OK.

$$
-\frac{\sqrt{2 c_{1} y^{2}-a}}{2 c_{1}}=c_{3}+x
$$

Verified OK.

### 1.76.2 Maple step by step solution

Let's solve
$y^{\prime \prime} y^{3}=a$

- Highest derivative means the order of the ODE is 2
$y^{\prime \prime}$
- $\quad$ Define new dependent variable $u$
$u(x)=y^{\prime}$
- Compute $y^{\prime \prime}$
$u^{\prime}(x)=y^{\prime \prime}$
- Use chain rule on the lhs
$y^{\prime}\left(\frac{d}{d y} u(y)\right)=y^{\prime \prime}$
- $\quad$ Substitute in the definition of $u$
$u(y)\left(\frac{d}{d y} u(y)\right)=y^{\prime \prime}$
- Make substitutions $y^{\prime}=u(y), y^{\prime \prime}=u(y)\left(\frac{d}{d y} u(y)\right)$ to reduce order of ODE $u(y)\left(\frac{d}{d y} u(y)\right) y^{3}=a$
- $\quad$ Separate variables
$u(y)\left(\frac{d}{d y} u(y)\right)=\frac{a}{y^{3}}$
- Integrate both sides with respect to $y$
$\int u(y)\left(\frac{d}{d y} u(y)\right) d y=\int \frac{a}{y^{3}} d y+c_{1}$
- Evaluate integral
$\frac{u(y)^{2}}{2}=-\frac{a}{2 y^{2}}+c_{1}$
- $\quad$ Solve for $u(y)$
$\left\{u(y)=\frac{\sqrt{2 c_{1} y^{2}-a}}{y}, u(y)=-\frac{\sqrt{2 c_{1} y^{2}-a}}{y}\right\}$
- $\quad$ Solve 1st ODE for $u(y)$
$u(y)=\frac{\sqrt{2 c_{1} y^{2}-a}}{y}$
- Revert to original variables with substitution $u(y)=y^{\prime}, y=y$
$y^{\prime}=\frac{\sqrt{2 c_{1} y^{2}-a}}{y}$
- $\quad$ Separate variables

$$
\frac{y y^{\prime}}{\sqrt{2 c_{1} y^{2}-a}}=1
$$

- Integrate both sides with respect to $x$
$\int \frac{y y^{\prime}}{\sqrt{2 c_{1} y^{2}-a}} d x=\int 1 d x+c_{2}$
- Evaluate integral
$\frac{\sqrt{2 c_{1} y^{2}-a}}{2 c_{1}}=x+c_{2}$
- $\quad$ Solve for $y$
$\left\{y=-\frac{\sqrt{2} \sqrt{c_{1}\left(4 c_{1}^{2} c_{2}^{2}+8 c_{1}^{2} c_{2} x+4 c_{1}^{2} x^{2}+a\right)}}{2 c_{1}}, y=\frac{\sqrt{2} \sqrt{c_{1}\left(4 c_{1}^{2} c_{2}^{2}+8 c_{1}^{2} c_{2} x+4 c_{1}^{2} x^{2}+a\right)}}{2 c_{1}}\right\}$
- $\quad$ Solve 2nd ODE for $u(y)$

$$
u(y)=-\frac{\sqrt{2 c_{1} y^{2}-a}}{y}
$$

- Revert to original variables with substitution $u(y)=y^{\prime}, y=y$

$$
y^{\prime}=-\frac{\sqrt{2 c_{1} y^{2}-a}}{y}
$$

- Separate variables

$$
\frac{y y^{\prime}}{\sqrt{2 c_{1} y^{2}-a}}=-1
$$

- Integrate both sides with respect to $x$

$$
\int \frac{y y^{\prime}}{\sqrt{2 c_{1} y^{2}-a}} d x=\int(-1) d x+c_{2}
$$

- Evaluate integral

$$
\frac{\sqrt{2 c_{1} y^{2}-a}}{2 c_{1}}=-x+c_{2}
$$

- $\quad$ Solve for $y$

$$
\left\{y=-\frac{\sqrt{2} \sqrt{c_{1}\left(4 c_{1}^{2} c_{2}^{2}-8 c_{1}^{2} c_{2} x+4 c_{1}^{2} x^{2}+a\right)}}{2 c_{1}}, y=\frac{\sqrt{2} \sqrt{c_{1}\left(4 c_{1}^{2} c_{2}^{2}-8 c_{1}^{2} c_{2} x+4 c_{1}^{2} x^{2}+a\right)}}{2 c_{1}}\right\}
$$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying 2nd order Liouville
trying 2nd order WeierstrassP
trying 2nd order JacobiSN
differential order: 2; trying a linearization to 3rd order
trying 2nd order ODE linearizable_by_differentiation
trying 2nd order, 2 integrating factors of the form mu(x,y)
trying differential order: 2; missing variables
`, `-> Computing symmetries using: way = 3
<- differential order: 2; canonical coordinates successful
<- differential order 2; missing variables successful`
```

$\checkmark$ Solution by Maple
Time used: 0.047 (sec). Leaf size: 46
dsolve(diff $(y(x), x \$ 2)=a / y(x) \sim 3, y(x)$, singsol=all)

$$
\begin{aligned}
& y(x)=\frac{\sqrt{c_{1}\left(\left(c_{2}+x\right)^{2} c_{1}^{2}+a\right)}}{c_{1}} \\
& y(x)=-\frac{\sqrt{c_{1}\left(\left(c_{2}+x\right)^{2} c_{1}^{2}+a\right)}}{c_{1}}
\end{aligned}
$$

$\checkmark$ Solution by Mathematica
Time used: 4.493 (sec). Leaf size: 63
DSolve $\left[y^{\prime \prime}[x]==a / y[x] \sim 3, y[x], x\right.$, IncludeSingularSolutions $->$ True]

$$
\begin{aligned}
& y(x) \rightarrow-\frac{\sqrt{a+c_{1}^{2}\left(x+c_{2}\right)^{2}}}{\sqrt{c_{1}}} \\
& y(x) \rightarrow \frac{\sqrt{a+c_{1}^{2}\left(x+c_{2}\right)^{2}}}{\sqrt{c_{1}}} \\
& y(x) \rightarrow \text { Indeterminate }
\end{aligned}
$$

### 1.77 problem 122

1.77.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 834
1.77.2 Solving as second order integrable as is ode . . . . . . . . . . . 834
1.77.3 Solving as second order ode missing y ode . . . . . . . . . . . . 836


1.77.6 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 846
1.77.7 Solving as exact linear second order ode ode . . . . . . . . . . . 854
1.77.8 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 857

Internal problem ID [12494]
Internal file name [OUTPUT/11146_Monday_October_16_2023_09_51_59_PM_15915211/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 122.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "exact linear second order ode", "second__order_integrable_as_is", "second_order_ode_missing_y", "second__order_ode__non_constant_coeff_transformation_on_B"

Maple gives the following as the ode type
[[_2nd_order, _missing_y]]

$$
x y^{\prime \prime}-y^{\prime}=x^{2} \mathrm{e}^{x}
$$

With initial conditions

$$
\left[y(0)=-1, y^{\prime}(0)=0\right]
$$

### 1.77.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime \prime}+p(x) y^{\prime}+q(x) y=F
$$

Where here

$$
\begin{aligned}
p(x) & =-\frac{1}{x} \\
q(x) & =0 \\
F & =x \mathrm{e}^{x}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime \prime}-\frac{y^{\prime}}{x}=x \mathrm{e}^{x}
$$

The domain of $p(x)=-\frac{1}{x}$ is

$$
\{x<0 \vee 0<x\}
$$

But the point $x_{0}=0$ is not inside this domain. Hence existence and uniqueness theorem does not apply. There could be infinite number of solutions, or one solution or no solution at all.

### 1.77.2 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t $x$ gives

$$
\begin{array}{r}
\int\left(x y^{\prime \prime}-y^{\prime}\right) d x=\int x^{2} \mathrm{e}^{x} d x \\
y^{\prime} x-2 y=\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}
\end{array}
$$

Which is now solved for $y$.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{2}{x} \\
& q(x)=\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{2 y}{x}=\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2}{x} d x} \\
& =\frac{1}{x^{2}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{y}{x^{2}}\right) & =\left(\frac{1}{x^{2}}\right)\left(\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}\right) \\
\mathrm{d}\left(\frac{y}{x^{2}}\right) & =\left(\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x^{3}}\right) \mathrm{d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
\frac{y}{x^{2}} & =\int \frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x^{3}} \mathrm{~d} x \\
\frac{y}{x^{2}} & =-\frac{c_{1}}{2 x^{2}}-\frac{\mathrm{e}^{x}}{x^{2}}+\frac{\mathrm{e}^{x}}{x}+c_{2}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x^{2}}$ results in

$$
y=x^{2}\left(-\frac{c_{1}}{2 x^{2}}-\frac{\mathrm{e}^{x}}{x^{2}}+\frac{\mathrm{e}^{x}}{x}\right)+c_{2} x^{2}
$$

which simplifies to

$$
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}-\frac{c_{1}}{2}
$$

Initial conditions are used to solve for the constants of integration.
Looking at the above solution

$$
\begin{equation*}
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}-\frac{c_{1}}{2} \tag{1}
\end{equation*}
$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y=-1$ and $x=0$ in the above gives

$$
\begin{equation*}
-1=-1-\frac{c_{1}}{2} \tag{1A}
\end{equation*}
$$

Taking derivative of the solution gives

$$
y^{\prime}=\mathrm{e}^{x}(x-1)+\mathrm{e}^{x}+2 c_{2} x
$$

substituting $y^{\prime}=0$ and $x=0$ in the above gives

$$
\begin{equation*}
0=0 \tag{2~A}
\end{equation*}
$$

Equations $\{1 \mathrm{~A}, 2 \mathrm{~A}\}$ are now solved for $\left\{c_{1}, c_{2}\right\}$. Solving for the constants gives

$$
c_{1}=0
$$

Substituting these values back in above solution results in

$$
y=x \mathrm{e}^{x}-\mathrm{e}^{x}+c_{2} x^{2}
$$

Which simplifies to

$$
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}
$$

Verified OK.

### 1.77.3 Solving as second order ode missing y ode

This is second order ode with missing dependent variable $y$. Let

$$
p(x)=y^{\prime}
$$

Then

$$
p^{\prime}(x)=y^{\prime \prime}
$$

Hence the ode becomes

$$
p^{\prime}(x) x-p(x)-x^{2} \mathrm{e}^{x}=0
$$

Which is now solve for $p(x)$ as first order ode.

Entering Linear first order ODE solver. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{1}{x} d x} \\
& =\frac{1}{x}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu p) & =(\mu)\left(x \mathrm{e}^{x}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{p}{x}\right) & =\left(\frac{1}{x}\right)\left(x \mathrm{e}^{x}\right) \\
\mathrm{d}\left(\frac{p}{x}\right) & =\mathrm{e}^{x} \mathrm{~d} x
\end{aligned}
$$

Integrating gives

$$
\begin{aligned}
& \frac{p}{x}=\int \mathrm{e}^{x} \mathrm{~d} x \\
& \frac{p}{x}=\mathrm{e}^{x}+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x}$ results in

$$
p(x)=x \mathrm{e}^{x}+c_{1} x
$$

which simplifies to

$$
p(x)=\left(\mathrm{e}^{x}+c_{1}\right) x
$$

Initial conditions are used to solve for $c_{1}$. Substituting $x=0$ and $p=0$ in the above solution gives an equation to solve for the constant of integration.

$$
0=0
$$

This solution is valid for any $c_{1}$. Hence there are infinite number of solutions.
Since $p=y^{\prime}$ then the new first order ode to solve is

$$
y^{\prime}=\left(\mathrm{e}^{x}+c_{1}\right) x
$$

Integrating both sides gives

$$
\begin{aligned}
y & =\int\left(\mathrm{e}^{x}+c_{1}\right) x \mathrm{~d} x \\
& =x \mathrm{e}^{x}-\mathrm{e}^{x}+\frac{c_{1} x^{2}}{2}+c_{2}
\end{aligned}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration.

$$
-1=-1+c_{2}
$$

This solution is valid for any $c_{1}$. Hence there are infinite number of solutions.
Initial conditions are used to solve for the constants of integration.
Looking at the above solution

$$
\begin{equation*}
y=x \mathrm{e}^{x}-\mathrm{e}^{x}+\frac{c_{1} x^{2}}{2}+c_{2} \tag{1}
\end{equation*}
$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y=-1$ and $x=0$ in the above gives

$$
\begin{equation*}
-1=-1+c_{2} \tag{1A}
\end{equation*}
$$

Taking derivative of the solution gives

$$
y^{\prime}=x \mathrm{e}^{x}+c_{1} x
$$

substituting $y^{\prime}=0$ and $x=0$ in the above gives

$$
\begin{equation*}
0=0 \tag{2A}
\end{equation*}
$$

Equations $\{1 \mathrm{~A}, 2 \mathrm{~A}\}$ are now solved for $\left\{c_{1}, c_{2}\right\}$. Solving for the constants gives

$$
c_{2}=0
$$

Substituting these values back in above solution results in

$$
y=x \mathrm{e}^{x}-\mathrm{e}^{x}+\frac{c_{1} x^{2}}{2}
$$

Which simplifies to

$$
y=\mathrm{e}^{x}(x-1)+\frac{c_{1} x^{2}}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{x}(x-1)+\frac{c_{1} x^{2}}{2} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\mathrm{e}^{x}(x-1)+\frac{c_{1} x^{2}}{2}
$$

Verified OK.

### 1.77.4 Solving as second order ode non constant coeff transformation on B ode

Given an ode of the form

$$
A y^{\prime \prime}+B y^{\prime}+C y=F(x)
$$

This method reduces the order ode the ODE by one by applying the transformation

$$
y=B v
$$

This results in

$$
\begin{aligned}
y^{\prime} & =B^{\prime} v+v^{\prime} B \\
y^{\prime \prime} & =B^{\prime \prime} v+B^{\prime} v^{\prime}+v^{\prime \prime} B+v^{\prime} B^{\prime} \\
& =v^{\prime \prime} B+2 v^{\prime}+B^{\prime}+B^{\prime \prime} v
\end{aligned}
$$

And now the original ode becomes

$$
\begin{array}{r}
A\left(v^{\prime \prime} B+2 v^{\prime} B^{\prime}+B^{\prime \prime} v\right)+B\left(B^{\prime} v+v^{\prime} B\right)+C B v=0 \\
A B v^{\prime \prime}+\left(2 A B^{\prime}+B^{2}\right) v^{\prime}+\left(A B^{\prime \prime}+B B^{\prime}+C B\right) v=0 \tag{1}
\end{array}
$$

If the term $A B^{\prime \prime}+B B^{\prime}+C B$ is zero, then this method works and can be used to solve

$$
A B v^{\prime \prime}+\left(2 A B^{\prime}+B^{2}\right) v^{\prime}=0
$$

By Using $u=v^{\prime}$ which reduces the order of the above ode to one. The new ode is

$$
A B u^{\prime}+\left(2 A B^{\prime}+B^{2}\right) u=0
$$

The above ode is first order ode which is solved for $u$. Now a new ode $v^{\prime}=u$ is solved for $v$ as first order ode. Then the final solution is obtain from $y=B v$.

This method works only if the term $A B^{\prime \prime}+B B^{\prime}+C B$ is zero. The given ODE shows that

$$
\begin{aligned}
& A=x \\
& B=-1 \\
& C=0 \\
& F=x^{2} \mathrm{e}^{x}
\end{aligned}
$$

The above shows that for this ode

$$
\begin{aligned}
A B^{\prime \prime}+B B^{\prime}+C B & =(x)(0)+(-1)(0)+(0)(-1) \\
& =0
\end{aligned}
$$

Hence the ode in $v$ given in (1) now simplifies to

$$
-x v^{\prime \prime}+(1) v^{\prime}=0
$$

Now by applying $v^{\prime}=u$ the above becomes

$$
-x u^{\prime}(x)+u(x)=0
$$

Which is now solved for $u$. In canonical form the ODE is

$$
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =\frac{u}{x}
\end{aligned}
$$

Where $f(x)=\frac{1}{x}$ and $g(u)=u$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{u} d u & =\frac{1}{x} d x \\
\int \frac{1}{u} d u & =\int \frac{1}{x} d x \\
\ln (u) & =\ln (x)+c_{1} \\
u & =\mathrm{e}^{\ln (x)+c_{1}} \\
& =c_{1} x
\end{aligned}
$$

The ode for $v$ now becomes

$$
\begin{aligned}
v^{\prime} & =u \\
& =c_{1} x
\end{aligned}
$$

Which is now solved for $v$. Integrating both sides gives

$$
\begin{aligned}
v(x) & =\int c_{1} x \mathrm{~d} x \\
& =\frac{c_{1} x^{2}}{2}+c_{2}
\end{aligned}
$$

Therefore the homogeneous solution is

$$
\begin{aligned}
y_{h}(x) & =B v \\
& =(-1)\left(\frac{c_{1} x^{2}}{2}+c_{2}\right) \\
& =-\frac{c_{1} x^{2}}{2}-c_{2}
\end{aligned}
$$

And now the particular solution $y_{p}(x)$ will be found. The particular solution $y_{p}$ can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on $x$ as well. Let

$$
\begin{equation*}
y_{p}(x)=u_{1} y_{1}+u_{2} y_{2} \tag{1}
\end{equation*}
$$

Where $u_{1}, u_{2}$ to be determined, and $y_{1}, y_{2}$ are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$
\begin{aligned}
& y_{1}=-1 \\
& y_{2}=x^{2}
\end{aligned}
$$

In the Variation of parameters $u_{1}, u_{2}$ are found using

$$
\begin{align*}
& u_{1}=-\int \frac{y_{2} f(x)}{a W(x)}  \tag{2}\\
& u_{2}=\int \frac{y_{1} f(x)}{a W(x)} \tag{3}
\end{align*}
$$

Where $W(x)$ is the Wronskian and $a$ is the coefficient in front of $y^{\prime \prime}$ in the given ODE.
The Wronskian is given by $W=\left|\begin{array}{ll}y_{1} & y_{2} \\ y_{1}^{\prime} & y_{2}^{\prime}\end{array}\right|$. Hence

$$
W=\left|\begin{array}{cc}
-1 & x^{2} \\
\frac{d}{d x}(-1) & \frac{d}{d x}\left(x^{2}\right)
\end{array}\right|
$$

Which gives

$$
W=\left|\begin{array}{cc}
-1 & x^{2} \\
0 & 2 x
\end{array}\right|
$$

Therefore

$$
W=(-1)(2 x)-\left(x^{2}\right)(0)
$$

Which simplifies to

$$
W=-2 x
$$

Which simplifies to

$$
W=-2 x
$$

Therefore Eq. (2) becomes

$$
u_{1}=-\int \frac{x^{4} \mathrm{e}^{x}}{-2 x^{2}} d x
$$

Which simplifies to

$$
u_{1}=-\int-\frac{x^{2} \mathrm{e}^{x}}{2} d x
$$

Hence

$$
u_{1}=\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}}{2}
$$

And Eq. (3) becomes

$$
u_{2}=\int \frac{-x^{2} \mathrm{e}^{x}}{-2 x^{2}} d x
$$

Which simplifies to

$$
u_{2}=\int \frac{\mathrm{e}^{x}}{2} d x
$$

Hence

$$
u_{2}=\frac{\mathrm{e}^{x}}{2}
$$

Therefore the particular solution, from equation (1) is

$$
y_{p}(x)=-\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}}{2}+\frac{x^{2} \mathrm{e}^{x}}{2}
$$

Which simplifies to

$$
y_{p}(x)=\mathrm{e}^{x}(x-1)
$$

Hence the complete solution is

$$
\begin{aligned}
y(x) & =y_{h}+y_{p} \\
& =\left(-\frac{c_{1} x^{2}}{2}-c_{2}\right)+\left(\mathrm{e}^{x}(x-1)\right) \\
& =-\frac{c_{1} x^{2}}{2}-c_{2}+\mathrm{e}^{x}(x-1)
\end{aligned}
$$

Initial conditions are used to solve for the constants of integration.
Looking at the above solution

$$
\begin{equation*}
y=-\frac{c_{1} x^{2}}{2}-c_{2}+\mathrm{e}^{x}(x-1) \tag{1}
\end{equation*}
$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y=-1$ and $x=0$ in the above gives

$$
\begin{equation*}
-1=-1-c_{2} \tag{1~A}
\end{equation*}
$$

Taking derivative of the solution gives

$$
y^{\prime}=-c_{1} x+\mathrm{e}^{x}(x-1)+\mathrm{e}^{x}
$$

substituting $y^{\prime}=0$ and $x=0$ in the above gives

$$
\begin{equation*}
0=0 \tag{2A}
\end{equation*}
$$

Equations $\{1 \mathrm{~A}, 2 \mathrm{~A}\}$ are now solved for $\left\{c_{1}, c_{2}\right\}$. Solving for the constants gives

$$
c_{2}=0
$$

Substituting these values back in above solution results in

$$
y=-\frac{c_{1} x^{2}}{2}+x \mathrm{e}^{x}-\mathrm{e}^{x}
$$

Which simplifies to

$$
y=\mathrm{e}^{x}(x-1)-\frac{c_{1} x^{2}}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{x}(x-1)-\frac{c_{1} x^{2}}{2} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\mathrm{e}^{x}(x-1)-\frac{c_{1} x^{2}}{2}
$$

Verified OK.

### 1.77.5 Solving as type second_order_integrable_as_is (not using ABC version)

Writing the ode as

$$
x y^{\prime \prime}-y^{\prime}=x^{2} \mathrm{e}^{x}
$$

Integrating both sides of the ODE w.r.t $x$ gives

$$
\begin{array}{r}
\int\left(x y^{\prime \prime}-y^{\prime}\right) d x=\int x^{2} \mathrm{e}^{x} d x \\
y^{\prime} x-2 y=\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}
\end{array}
$$

Which is now solved for $y$.
Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{2}{x} \\
& q(x)=\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{2 y}{x}=\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2}{x} d x} \\
& =\frac{1}{x^{2}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{y}{x^{2}}\right) & =\left(\frac{1}{x^{2}}\right)\left(\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}\right) \\
\mathrm{d}\left(\frac{y}{x^{2}}\right) & =\left(\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x^{3}}\right) \mathrm{d} x
\end{aligned}
$$

## Integrating gives

$$
\begin{aligned}
\frac{y}{x^{2}} & =\int \frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x^{3}} \mathrm{~d} x \\
\frac{y}{x^{2}} & =-\frac{c_{1}}{2 x^{2}}-\frac{\mathrm{e}^{x}}{x^{2}}+\frac{\mathrm{e}^{x}}{x}+c_{2}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x^{2}}$ results in

$$
y=x^{2}\left(-\frac{c_{1}}{2 x^{2}}-\frac{\mathrm{e}^{x}}{x^{2}}+\frac{\mathrm{e}^{x}}{x}\right)+c_{2} x^{2}
$$

which simplifies to

$$
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}-\frac{c_{1}}{2}
$$

Initial conditions are used to solve for the constants of integration.
Looking at the above solution

$$
\begin{equation*}
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}-\frac{c_{1}}{2} \tag{1}
\end{equation*}
$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y=-1$ and $x=0$ in the above gives

$$
\begin{equation*}
-1=-1-\frac{c_{1}}{2} \tag{1~A}
\end{equation*}
$$

Taking derivative of the solution gives

$$
y^{\prime}=\mathrm{e}^{x}(x-1)+\mathrm{e}^{x}+2 c_{2} x
$$

substituting $y^{\prime}=0$ and $x=0$ in the above gives

$$
\begin{equation*}
0=0 \tag{2~A}
\end{equation*}
$$

Equations $\{1 \mathrm{~A}, 2 \mathrm{~A}\}$ are now solved for $\left\{c_{1}, c_{2}\right\}$. Solving for the constants gives

$$
c_{1}=0
$$

Substituting these values back in above solution results in

$$
y=x \mathrm{e}^{x}-\mathrm{e}^{x}+c_{2} x^{2}
$$

Which simplifies to

$$
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}
$$

Verified OK.

### 1.77.6 Solving using Kovacic algorithm

Writing the ode as

$$
\begin{array}{r}
x y^{\prime \prime}-y^{\prime}=0 \\
A y^{\prime \prime}+B y^{\prime}+C y=0 \tag{2}
\end{array}
$$

Comparing (1) and (2) shows that

$$
\begin{align*}
& A=x \\
& B=-1  \tag{3}\\
& C=0
\end{align*}
$$

Applying the Liouville transformation on the dependent variable gives

$$
z(x)=y e^{\int \frac{B}{2 A} d x}
$$

Then (2) becomes

$$
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
$$

Where $r$ is given by

$$
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
$$

Substituting the values of $A, B, C$ from (3) in the above and simplifying gives

$$
\begin{equation*}
r=\frac{3}{4 x^{2}} \tag{6}
\end{equation*}
$$

Comparing the above to (5) shows that

$$
\begin{aligned}
& s=3 \\
& t=4 x^{2}
\end{aligned}
$$

Therefore eq. (4) becomes

$$
\begin{equation*}
z^{\prime \prime}(x)=\left(\frac{3}{4 x^{2}}\right) z(x) \tag{7}
\end{equation*}
$$

Equation (7) is now solved. After finding $z(x)$ then $y$ is found using the inverse transformation

$$
y=z(x) e^{-\int \frac{B}{2 A} d x}
$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of $r$ and the order of $r$ at $\infty$. The following table summarizes these cases.

| Case | Allowed pole order for $r$ | Allowed value for $\mathcal{O}(\infty)$ |
| :--- | :--- | :--- |
| 1 | $\{0,1,2,4,6,8, \cdots\}$ | $\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}$ |
| 2 | Need to have at least one pole that <br> is either order 2 or odd order greater <br> than 2. Any other pole order is <br> allowed as long as the above condi- <br> tion is satisfied. Hence the following <br> set of pole orders are all allowed. <br> $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}$. |  |
| 3 | $\{1,2\}$ | $\{2,3,4,5,6,7, \cdots\}$ |

Table 126: Necessary conditions for each Kovacic case

The order of $r$ at $\infty$ is the degree of $t$ minus the degree of $s$. Therefore

$$
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =2-0 \\
& =2
\end{aligned}
$$

The poles of $r$ in eq. (7) and the order of each pole are determined by solving for the roots of $t=4 x^{2}$. There is a pole at $x=0$ of order 2 . Since there is no odd order pole larger than 2 and the order at $\infty$ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at $\infty$ is 2 then the necessary conditions for case three are met. Therefore

$$
L=[1,2,4,6,12]
$$

Attempting to find a solution using case $n=1$.
Looking at poles of order 2. The partial fractions decomposition of $r$ is

$$
r=\frac{3}{4 x^{2}}
$$

For the pole at $x=0$ let $b$ be the coefficient of $\frac{1}{x^{2}}$ in the partial fractions decomposition of $r$ given above. Therefore $b=\frac{3}{4}$. Hence

$$
\begin{aligned}
{[\sqrt{r}]_{c} } & =0 \\
\alpha_{c}^{+} & =\frac{1}{2}+\sqrt{1+4 b}=\frac{3}{2} \\
\alpha_{c}^{-} & =\frac{1}{2}-\sqrt{1+4 b}=-\frac{1}{2}
\end{aligned}
$$

Since the order of $r$ at $\infty$ is 2 then $[\sqrt{r}]_{\infty}=0$. Let $b$ be the coefficient of $\frac{1}{x^{2}}$ in the Laurent series expansion of $r$ at $\infty$. which can be found by dividing the leading coefficient of $s$ by the leading coefficient of $t$ from

$$
r=\frac{s}{t}=\frac{3}{4 x^{2}}
$$

Since the $\operatorname{gcd}(s, t)=1$. This gives $b=\frac{3}{4}$. Hence

$$
\begin{aligned}
{[\sqrt{r}]_{\infty} } & =0 \\
\alpha_{\infty}^{+} & =\frac{1}{2}+\sqrt{1+4 b}=\frac{3}{2} \\
\alpha_{\infty}^{-} & =\frac{1}{2}-\sqrt{1+4 b}=-\frac{1}{2}
\end{aligned}
$$

The following table summarizes the findings so far for poles and for the order of $r$ at $\infty$ where $r$ is

$$
r=\frac{3}{4 x^{2}}
$$

| pole $c$ location | pole order | $[\sqrt{r}]_{c}$ | $\alpha_{c}^{+}$ | $\alpha_{c}^{-}$ |
| :---: | :---: | :---: | :---: | :---: |
| 0 | 2 | 0 | $\frac{3}{2}$ | $-\frac{1}{2}$ |


| Order of $r$ at $\infty$ | $[\sqrt{r}]_{\infty}$ | $\alpha_{\infty}^{+}$ | $\alpha_{\infty}^{-}$ |
| :---: | :---: | :---: | :---: |
| 2 | 0 | $\frac{3}{2}$ | $-\frac{1}{2}$ |

Now that the all $[\sqrt{r}]_{c}$ and its associated $\alpha_{c}^{ \pm}$have been determined for all the poles in the set $\Gamma$ and $[\sqrt{r}]_{\infty}$ and its associated $\alpha_{\infty}^{ \pm}$have also been found, the next step is to
determine possible non negative integer $d$ from these using

$$
d=\alpha_{\infty}^{s(\infty)}-\sum_{c \in \Gamma} \alpha_{c}^{s(c)}
$$

Where $s(c)$ is either + or - and $s(\infty)$ is the sign of $\alpha_{\infty}^{ \pm}$. This is done by trial over all set of families $s=(s(c))_{c \in \Gamma \cup \infty}$ until such $d$ is found to work in finding candidate $\omega$. Trying $\alpha_{\infty}^{-}=-\frac{1}{2}$ then

$$
\begin{aligned}
d & =\alpha_{\infty}^{-}-\left(\alpha_{c_{1}}^{-}\right) \\
& =-\frac{1}{2}-\left(-\frac{1}{2}\right) \\
& =0
\end{aligned}
$$

Since $d$ an integer and $d \geq 0$ then it can be used to find $\omega$ using

$$
\omega=\sum_{c \in \Gamma}\left(s(c)[\sqrt{r}]_{c}+\frac{\alpha_{c}^{s(c)}}{x-c}\right)+s(\infty)[\sqrt{r}]_{\infty}
$$

The above gives

$$
\begin{aligned}
\omega & =\left((-)[\sqrt{r}]_{c_{1}}+\frac{\alpha_{c_{1}}^{-}}{x-c_{1}}\right)+(-)[\sqrt{r}]_{\infty} \\
& =-\frac{1}{2 x}+(-)(0) \\
& =-\frac{1}{2 x} \\
& =-\frac{1}{2 x}
\end{aligned}
$$

Now that $\omega$ is determined, the next step is find a corresponding minimal polynomial $p(x)$ of degree $d=0$ to solve the ode. The polynomial $p(x)$ needs to satisfy the equation

$$
\begin{equation*}
p^{\prime \prime}+2 \omega p^{\prime}+\left(\omega^{\prime}+\omega^{2}-r\right) p=0 \tag{1~A}
\end{equation*}
$$

Let

$$
\begin{equation*}
p(x)=1 \tag{2~A}
\end{equation*}
$$

Substituting the above in eq. (1A) gives

$$
\begin{array}{r}
(0)+2\left(-\frac{1}{2 x}\right)(0)+\left(\left(\frac{1}{2 x^{2}}\right)+\left(-\frac{1}{2 x}\right)^{2}-\left(\frac{3}{4 x^{2}}\right)\right)=0 \\
0=0
\end{array}
$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z^{\prime \prime}=r z$ is

$$
\begin{aligned}
z_{1}(x) & =p e^{\int \omega d x} \\
& =\mathrm{e}^{\int-\frac{1}{2 x} d x} \\
& =\frac{1}{\sqrt{x}}
\end{aligned}
$$

The first solution to the original ode in $y$ is found from

$$
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{1}{x} d x} \\
& =z_{1} e^{\frac{\ln (x)}{2}} \\
& =z_{1}(\sqrt{x})
\end{aligned}
$$

Which simplifies to

$$
y_{1}=1
$$

The second solution $y_{2}$ to the original ode is found using reduction of order

$$
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
$$

Substituting gives

$$
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-1}{x} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{\ln (x)}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\frac{x^{2}}{2}\right)
\end{aligned}
$$

Therefore the solution is

$$
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}(1)+c_{2}\left(1\left(\frac{x^{2}}{2}\right)\right)
\end{aligned}
$$

This is second order nonhomogeneous ODE. Let the solution be

$$
y=y_{h}+y_{p}
$$

Where $y_{h}$ is the solution to the homogeneous ODE $A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0$, and $y_{p}$ is a particular solution to the nonhomogeneous ODE $A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)$. $y_{h}$ is the solution to

$$
x y^{\prime \prime}-y^{\prime}=0
$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$
y_{h}=c_{1}+\frac{c_{2} x^{2}}{2}
$$

The particular solution $y_{p}$ can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on $x$ as well. Let

$$
\begin{equation*}
y_{p}(x)=u_{1} y_{1}+u_{2} y_{2} \tag{1}
\end{equation*}
$$

Where $u_{1}, u_{2}$ to be determined, and $y_{1}, y_{2}$ are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$
\begin{aligned}
& y_{1}=1 \\
& y_{2}=\frac{x^{2}}{2}
\end{aligned}
$$

In the Variation of parameters $u_{1}, u_{2}$ are found using

$$
\begin{align*}
& u_{1}=-\int \frac{y_{2} f(x)}{a W(x)}  \tag{2}\\
& u_{2}=\int \frac{y_{1} f(x)}{a W(x)} \tag{3}
\end{align*}
$$

Where $W(x)$ is the Wronskian and $a$ is the coefficient in front of $y^{\prime \prime}$ in the given ODE.
The Wronskian is given by $W=\left|\begin{array}{ll}y_{1} & y_{2} \\ y_{1}^{\prime} & y_{2}^{\prime}\end{array}\right|$. Hence

$$
W=\left|\begin{array}{cc}
1 & \frac{x^{2}}{2} \\
\frac{d}{d x}(1) & \frac{d}{d x}\left(\frac{x^{2}}{2}\right)
\end{array}\right|
$$

Which gives

$$
W=\left|\begin{array}{cc}
1 & \frac{x^{2}}{2} \\
0 & x
\end{array}\right|
$$

Therefore

$$
W=(1)(x)-\left(\frac{x^{2}}{2}\right)(0)
$$

Which simplifies to

$$
W=x
$$

Which simplifies to

$$
W=x
$$

Therefore Eq. (2) becomes

$$
u_{1}=-\int \frac{\frac{x^{4} \mathrm{e}^{x}}{2}}{x^{2}} d x
$$

Which simplifies to

$$
u_{1}=-\int \frac{x^{2} \mathrm{e}^{x}}{2} d x
$$

Hence

$$
u_{1}=-\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}}{2}
$$

And Eq. (3) becomes

$$
u_{2}=\int \frac{x^{2} \mathrm{e}^{x}}{x^{2}} d x
$$

Which simplifies to

$$
u_{2}=\int \mathrm{e}^{x} d x
$$

Hence

$$
u_{2}=\mathrm{e}^{x}
$$

Therefore the particular solution, from equation (1) is

$$
y_{p}(x)=-\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}}{2}+\frac{x^{2} \mathrm{e}^{x}}{2}
$$

Which simplifies to

$$
y_{p}(x)=\mathrm{e}^{x}(x-1)
$$

Therefore the general solution is

$$
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1}+\frac{c_{2} x^{2}}{2}\right)+\left(\mathrm{e}^{x}(x-1)\right)
\end{aligned}
$$

Initial conditions are used to solve for the constants of integration.
Looking at the above solution

$$
\begin{equation*}
y=c_{1}+\frac{c_{2} x^{2}}{2}+\mathrm{e}^{x}(x-1) \tag{1}
\end{equation*}
$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y=-1$ and $x=0$ in the above gives

$$
\begin{equation*}
-1=-1+c_{1} \tag{1~A}
\end{equation*}
$$

Taking derivative of the solution gives

$$
y^{\prime}=c_{2} x+\mathrm{e}^{x}(x-1)+\mathrm{e}^{x}
$$

substituting $y^{\prime}=0$ and $x=0$ in the above gives

$$
\begin{equation*}
0=0 \tag{2~A}
\end{equation*}
$$

Equations $\{1 \mathrm{~A}, 2 \mathrm{~A}\}$ are now solved for $\left\{c_{1}, c_{2}\right\}$. Solving for the constants gives

$$
c_{1}=0
$$

Substituting these values back in above solution results in

$$
y=\frac{c_{2} x^{2}}{2}+x \mathrm{e}^{x}-\mathrm{e}^{x}
$$

Which simplifies to

$$
y=\mathrm{e}^{x}(x-1)+\frac{c_{2} x^{2}}{2}
$$

Summary
The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{x}(x-1)+\frac{c_{2} x^{2}}{2} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\mathrm{e}^{x}(x-1)+\frac{c_{2} x^{2}}{2}
$$

Verified OK.

### 1.77.7 Solving as exact linear second order ode ode

An ode of the form

$$
p(x) y^{\prime \prime}+q(x) y^{\prime}+r(x) y=s(x)
$$

is exact if

$$
\begin{equation*}
p^{\prime \prime}(x)-q^{\prime}(x)+r(x)=0 \tag{1}
\end{equation*}
$$

For the given ode we have

$$
\begin{aligned}
p(x) & =x \\
q(x) & =-1 \\
r(x) & =0 \\
s(x) & =x^{2} \mathrm{e}^{x}
\end{aligned}
$$

Hence

$$
\begin{aligned}
p^{\prime \prime}(x) & =0 \\
q^{\prime}(x) & =0
\end{aligned}
$$

Therefore (1) becomes

$$
0-(0)+(0)=0
$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$
\left(p(x) y^{\prime}+\left(q(x)-p^{\prime}(x)\right) y\right)^{\prime}=s(x)
$$

## Integrating gives

$$
p(x) y^{\prime}+\left(q(x)-p^{\prime}(x)\right) y=\int s(x) d x
$$

Substituting the above values for $p, q, r, s$ gives

$$
y^{\prime} x-2 y=\int x^{2} \mathrm{e}^{x} d x
$$

We now have a first order ode to solve which is

$$
y^{\prime} x-2 y=\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}
$$

Entering Linear first order ODE solver. In canonical form a linear first order is

$$
y^{\prime}+p(x) y=q(x)
$$

Where here

$$
\begin{aligned}
& p(x)=-\frac{2}{x} \\
& q(x)=\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}
\end{aligned}
$$

Hence the ode is

$$
y^{\prime}-\frac{2 y}{x}=\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}
$$

The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{2}{x} d x} \\
& =\frac{1}{x^{2}}
\end{aligned}
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{y}{x^{2}}\right) & =\left(\frac{1}{x^{2}}\right)\left(\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x}\right) \\
\mathrm{d}\left(\frac{y}{x^{2}}\right) & =\left(\frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x^{3}}\right) \mathrm{d} x
\end{aligned}
$$

## Integrating gives

$$
\begin{aligned}
\frac{y}{x^{2}} & =\int \frac{\left(x^{2}-2 x+2\right) \mathrm{e}^{x}+c_{1}}{x^{3}} \mathrm{~d} x \\
\frac{y}{x^{2}} & =-\frac{c_{1}}{2 x^{2}}-\frac{\mathrm{e}^{x}}{x^{2}}+\frac{\mathrm{e}^{x}}{x}+c_{2}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\frac{1}{x^{2}}$ results in

$$
y=x^{2}\left(-\frac{c_{1}}{2 x^{2}}-\frac{\mathrm{e}^{x}}{x^{2}}+\frac{\mathrm{e}^{x}}{x}\right)+c_{2} x^{2}
$$

which simplifies to

$$
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}-\frac{c_{1}}{2}
$$

Initial conditions are used to solve for the constants of integration.
Looking at the above solution

$$
\begin{equation*}
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}-\frac{c_{1}}{2} \tag{1}
\end{equation*}
$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y=-1$ and $x=0$ in the above gives

$$
\begin{equation*}
-1=-1-\frac{c_{1}}{2} \tag{1~A}
\end{equation*}
$$

Taking derivative of the solution gives

$$
y^{\prime}=\mathrm{e}^{x}(x-1)+\mathrm{e}^{x}+2 c_{2} x
$$

substituting $y^{\prime}=0$ and $x=0$ in the above gives

$$
\begin{equation*}
0=0 \tag{2~A}
\end{equation*}
$$

Equations $\{1 \mathrm{~A}, 2 \mathrm{~A}\}$ are now solved for $\left\{c_{1}, c_{2}\right\}$. Solving for the constants gives

$$
c_{1}=0
$$

Substituting these values back in above solution results in

$$
y=x \mathrm{e}^{x}-\mathrm{e}^{x}+c_{2} x^{2}
$$

Which simplifies to

$$
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}
$$

## Summary

The solution(s) found are the following

$$
\begin{equation*}
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2} \tag{1}
\end{equation*}
$$

Verification of solutions

$$
y=\mathrm{e}^{x}(x-1)+c_{2} x^{2}
$$

Verified OK.

### 1.77.8 Maple step by step solution

Let's solve

$$
\left[y^{\prime \prime} x-y^{\prime}=x^{2} \mathrm{e}^{x}, y(0)=-1,\left.y^{\prime}\right|_{\{x=0\}}=0\right]
$$

- Highest derivative means the order of the ODE is 2

$$
y^{\prime \prime}
$$

- Make substitution $u=y^{\prime}$ to reduce order of ODE

$$
u^{\prime}(x) x-u(x)=x^{2} \mathrm{e}^{x}
$$

- Isolate the derivative

$$
u^{\prime}(x)=\frac{u(x)}{x}+x \mathrm{e}^{x}
$$

- Group terms with $u(x)$ on the lhs of the ODE and the rest on the rhs of the ODE $u^{\prime}(x)-\frac{u(x)}{x}=x \mathrm{e}^{x}$
- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$
\mu(x)\left(u^{\prime}(x)-\frac{u(x)}{x}\right)=\mu(x) x \mathrm{e}^{x}
$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) u(x))$

$$
\mu(x)\left(u^{\prime}(x)-\frac{u(x)}{x}\right)=\mu^{\prime}(x) u(x)+\mu(x) u^{\prime}(x)
$$

- $\quad$ Isolate $\mu^{\prime}(x)$

$$
\mu^{\prime}(x)=-\frac{\mu(x)}{x}
$$

- $\quad$ Solve to find the integrating factor

$$
\mu(x)=\frac{1}{x}
$$

- Integrate both sides with respect to $x$

$$
\int\left(\frac{d}{d x}(\mu(x) u(x))\right) d x=\int \mu(x) x \mathrm{e}^{x} d x+c_{1}
$$

- Evaluate the integral on the lhs
$\mu(x) u(x)=\int \mu(x) x \mathrm{e}^{x} d x+c_{1}$
- $\quad$ Solve for $u(x)$
$u(x)=\frac{\int \mu(x) x e^{x} d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\frac{1}{x}$
$u(x)=x\left(\int \mathrm{e}^{x} d x+c_{1}\right)$
- Evaluate the integrals on the rhs
$u(x)=\left(\mathrm{e}^{x}+c_{1}\right) x$
- $\quad$ Solve 1st ODE for $u(x)$
$u(x)=\left(\mathrm{e}^{x}+c_{1}\right) x$
- $\quad$ Make substitution $u=y^{\prime}$
$y^{\prime}=\left(\mathrm{e}^{x}+c_{1}\right) x$
- Integrate both sides to solve for $y$
$\int y^{\prime} d x=\int\left(\mathrm{e}^{x}+c_{1}\right) x d x+c_{2}$
- Compute integrals
$y=x \mathrm{e}^{x}-\mathrm{e}^{x}+\frac{c_{1} x^{2}}{2}+c_{2}$
Check validity of solution $y=x \mathrm{e}^{x}-\mathrm{e}^{x}+\frac{c_{1} x^{2}}{2}+c_{2}$
- Use initial condition $y(0)=-1$

$$
-1=-1+c_{2}
$$

- Compute derivative of the solution
$y^{\prime}=x \mathrm{e}^{x}+c_{1} x$
- Use the initial condition $\left.y^{\prime}\right|_{\{x=0\}}=0$
$0=0$
- Solve for $c_{1}$ and $c_{2}$
$\left\{c_{1}=c_{1}, c_{2}=0\right\}$
- Substitute constant values into general solution and simplify

$$
y=\mathrm{e}^{x}(x-1)+\frac{c_{1} x^{2}}{2}
$$

- $\quad$ Solution to the IVP

$$
y=\mathrm{e}^{x}(x-1)+\frac{c_{1} x^{2}}{2}
$$

Maple trace

```
-Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
-> Calling odsolve with the ODE`, diff(_b(_a), _a) = (_a^2*exp(_a)+_b(_a))/_a, _b(_a)`
    Methods for first order ODEs:
    --- Trying classification methods ---
    trying a quadrature
    trying 1st order linear
    <- 1st order linear successful
<- high order exact linear fully integrable successful`
```

Solution by Maple
Time used: 0.046 (sec). Leaf size: 17

```
dsolve([x*diff(y(x),x$2)-diff (y(x),x)=\mp@subsup{x}{}{~}2*\operatorname{exp}(x),y(0)=-1, D(y)(0)=0],y(x), singsol=all)
```

$$
y(x)=(-1+x) \mathrm{e}^{x}+\frac{c_{1} x^{2}}{2}
$$

Solution by Mathematica
Time used: 0.114 (sec). Leaf size: 22
DSolve $\left[\left\{x * y\right.\right.$ '' $[x]-y$ ' $\left.[x]==x^{\wedge} 2 * \operatorname{Exp}[x],\left\{y[0]==-1, y^{\prime}[0]==0\right\}\right\}, y[x], x$, IncludeSingularSolutions $->~ T$

$$
y(x) \rightarrow e^{x}(x-1)+\frac{c_{1} x^{2}}{2}
$$

### 1.78 problem 123

1.78.1 Solving as second order ode missing $x$ ode . . . . . . . . . . . . 860
1.78.2 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 863

Internal problem ID [12495]
Internal file name [OUTPUT/11147_Monday_October_16_2023_09_52_00_PM_4354937/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 123.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "second__order_ode_missing_x"
Maple gives the following as the ode type

```
[[_2nd_order, _missing_x], [_2nd_order, _reducible, _mu_x_y1],
    [_2nd_order, _reducible, _mu_y_y1]]
```

$$
y y^{\prime \prime}-y^{\prime 2}+y^{\prime 3}=0
$$

With initial conditions

$$
\left[y(0)=-1, y^{\prime}(0)=0\right]
$$

### 1.78.1 Solving as second order ode missing $x$ ode

This is missing independent variable second order ode. Solved by reduction of order by using substitution which makes the dependent variable $y$ an independent variable. Using

$$
y^{\prime}=p(y)
$$

Then

$$
\begin{aligned}
y^{\prime \prime} & =\frac{d p}{d x} \\
& =\frac{d y}{d x} \frac{d p}{d y} \\
& =p \frac{d p}{d y}
\end{aligned}
$$

Hence the ode becomes

$$
y p(y)\left(\frac{d}{d y} p(y)\right)+\left(-p(y)+p(y)^{2}\right) p(y)=0
$$

Which is now solved as first order ode for $p(y)$. In canonical form the ODE is

$$
\begin{aligned}
p^{\prime} & =F(y, p) \\
& =f(y) g(p) \\
& =-\frac{p(p-1)}{y}
\end{aligned}
$$

Where $f(y)=-\frac{1}{y}$ and $g(p)=p(p-1)$. Integrating both sides gives

$$
\begin{aligned}
\frac{1}{p(p-1)} d p & =-\frac{1}{y} d y \\
\int \frac{1}{p(p-1)} d p & =\int-\frac{1}{y} d y \\
\ln (p-1)-\ln (p) & =-\ln (y)+c_{1}
\end{aligned}
$$

Raising both side to exponential gives

$$
\mathrm{e}^{\ln (p-1)-\ln (p)}=\mathrm{e}^{-\ln (y)+c_{1}}
$$

Which simplifies to

$$
\frac{p-1}{p}=\frac{c_{2}}{y}
$$

Initial conditions are used to solve for $c_{2}$. Substituting $y=-1$ and $p=0$ in the above solution gives an equation to solve for the constant of integration.

$$
0=\frac{1}{1+c_{2}}
$$

Unable to solve for constant of integration. Since $\lim _{c_{1} \rightarrow \infty}$ gives $p=-\frac{y}{c_{2}-y}=p=0$ and this result satisfies the given initial condition. For solution (1) found earlier, since $p=y^{\prime}$ then we now have a new first order ode to solve which is

$$
y^{\prime}=0
$$

Integrating both sides gives

$$
\begin{aligned}
y & =\int 0 \mathrm{~d} x \\
& =c_{3}
\end{aligned}
$$

Initial conditions are used to solve for $c_{3}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& -1=c_{3} \\
& c_{3}=-1
\end{aligned}
$$

Substituting $c_{3}$ found above in the general solution gives

$$
y=-1
$$

Initial conditions are used to solve for the constants of integration.
Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-1 \tag{1}
\end{equation*}
$$



Figure 140: Solution plot

Verification of solutions

$$
y=-1
$$

Verified OK.

### 1.78.2 Maple step by step solution

Let's solve

$$
\left[y y^{\prime \prime}+\left(-y^{\prime}+y^{\prime 2}\right) y^{\prime}=0, y(0)=-1,\left.y^{\prime}\right|_{\{x=0\}}=0\right]
$$

- Highest derivative means the order of the ODE is 2

$$
y^{\prime \prime}
$$

- $\quad$ Define new dependent variable $u$

$$
u(x)=y^{\prime}
$$

- Compute $y^{\prime \prime}$

$$
u^{\prime}(x)=y^{\prime \prime}
$$

- Use chain rule on the lhs

$$
y^{\prime}\left(\frac{d}{d y} u(y)\right)=y^{\prime \prime}
$$

- $\quad$ Substitute in the definition of $u$
$u(y)\left(\frac{d}{d y} u(y)\right)=y^{\prime \prime}$
- Make substitutions $y^{\prime}=u(y), y^{\prime \prime}=u(y)\left(\frac{d}{d y} u(y)\right)$ to reduce order of ODE $y u(y)\left(\frac{d}{d y} u(y)\right)+\left(-u(y)+u(y)^{2}\right) u(y)=0$
- $\quad$ Separate variables

$$
\frac{\frac{d}{d y} u(y)}{-u(y)+u(y)^{2}}=-\frac{1}{y}
$$

- Integrate both sides with respect to $y$
$\int \frac{\frac{d}{d y} u(y)}{-u(y)+u(y)^{2}} d y=\int-\frac{1}{y} d y+c_{1}$
- $\quad$ Evaluate integral
$\ln (u(y)-1)-\ln (u(y))=-\ln (y)+c_{1}$
- $\quad$ Solve for $u(y)$

$$
u(y)=-\frac{y}{\mathrm{e}^{c_{1}}-y}
$$

- $\quad$ Solve 1st ODE for $u(y)$

$$
u(y)=-\frac{y}{\mathrm{e}^{c_{1}}-y}
$$

- Revert to original variables with substitution $u(y)=y^{\prime}, y=y$

$$
y^{\prime}=-\frac{y}{\mathrm{e}^{c_{1}}-y}
$$

- $\quad$ Separate variables

$$
\frac{y^{\prime}\left(\mathrm{e}^{c_{1}}-y\right)}{y}=-1
$$

- Integrate both sides with respect to $x$
$\int \frac{y^{\prime}\left(\mathrm{e}^{c_{1}}-y\right)}{y} d x=\int(-1) d x+c_{2}$
- $\quad$ Evaluate integral
$-y+\mathrm{e}^{c_{1}} \ln (y)=-x+c_{2}$
- $\quad$ Solve for $y$
$y=\mathrm{e}^{-\frac{\text { Lambert } W}{}\left(-\mathrm{e}^{-\frac{c_{1} \mathrm{e}^{c_{1}-c_{2}+x}}{\mathrm{e}^{c_{1}}}}\right) \mathrm{e}^{c_{1}-c_{2}+x}} \mathrm{e}^{c_{1}}$
Check validity of solution $y=\mathrm{e}^{-\frac{\text { LambertW }\left(-\mathrm{e}^{-\frac{c_{1} \mathrm{e}^{c_{1}}-c_{2}+x}{\mathrm{e}^{c_{1}}}}\right)}{\mathrm{e}^{c_{1}-c_{2}+x}}} \mathrm{e}^{c_{1}}$
- Use initial condition $y(0)=-1$
$-1=\mathrm{e}^{-\frac{\text { Lambert } W\left(-\mathrm{e}^{-\frac{c_{1} \mathrm{e}^{c_{1}}-c_{2}}{\mathrm{e}^{c_{1}}}}\right) \mathrm{e}^{c_{1}-c_{2}}}{\mathrm{e}^{c_{1}}}}$
- Compute derivative of the solution
$y^{\prime}=-\frac{\left(-\frac{\operatorname{Lambert} W\left(-\mathrm{e}^{\left.-\frac{c_{1} \mathrm{e}^{c_{1}-c_{2}+x}}{\mathrm{e}^{c_{1}}}\right)}\right.}{1+\operatorname{Lambert} W\left(-\mathrm{e}^{\left.-\frac{c_{1} \mathrm{e}^{c_{1}-c_{2}+x}}{\mathrm{e}^{c_{1}}}\right)}+1\right.}\right) \mathrm{e}^{\left.-\frac{\operatorname{LambertW}\left(-\mathrm{e}^{-\frac{c_{1} \mathrm{e}^{c_{1}-c_{2}+x}}{\mathrm{e}^{c} 1}}\right) \mathrm{e}^{c_{1}}-c_{2}+x}{\mathrm{e}^{c_{1}}}\right)}}{\mathrm{e}^{c_{1}}}$
- Use the initial condition $\left.y^{\prime}\right|_{\{x=0\}}=0$
$0=-\frac{\left(-\frac{\operatorname{Lambert} W\left(-\frac{c_{1} \mathrm{e}^{c_{1}-c_{2}}}{\mathrm{e}^{c_{1}}}\right)}{1+\operatorname{Lambert} W\left(-\mathrm{e}^{\left.-\frac{c_{1} \mathrm{e}^{c_{1}-c_{2}}}{\mathrm{e}^{c_{1}}}\right)}+1\right.}\right) \mathrm{e}^{-\frac{\operatorname{LambertW}\left(-\mathrm{e}^{-\frac{c_{1} \mathrm{e}^{c_{1}-c_{2}}}{\mathrm{e}^{c_{1}}}}\right) \mathrm{e}^{c_{1}}-c_{2}}{\mathrm{e}^{c_{1}}}} \mathrm{e}^{c_{1}}}{}$
- Solve for $c_{1}$ and $c_{2}$
- The solution does not satisfy the initial conditions

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying 2nd order Liouville
trying 2nd order WeierstrassP
trying 2nd order JacobiSN
differential order: 2; trying a linearization to 3rd order
--- trying a change of variables {x -> y(x), y(x) -> x}
differential order: 2; trying a linearization to 3rd order
trying 2nd order ODE linearizable_by_differentiation
trying 2nd order, 2 integrating factors of the form mu(x,y)
trying differential order: 2; missing variables
-> Calling odsolve with the ODE`, (diff(_b(_a), _a))*_b(_a)+_b(_a)^2*(_b(_a)-1)/_a = 0, _b(_
    Methods for first order ODEs:
    --- Trying classification methods ---
    trying a quadrature
    trying 1st order linear
    trying Bernoulli
    <- Bernoulli successful
<- differential order: 2; canonical coordinates successful
<- differential order 2; missing variables successful`
```

$\checkmark$ Solution by Maple
Time used: 0.063 (sec). Leaf size: 5

```
dsolve([y(x)*diff(y(x),x$2)-diff(y(x),x)^2+\operatorname{diff}(y(x),x)^3=0,y(0) = -1, D(y)(0) = 0],y(x), si
```

$$
y(x)=-1
$$

$X$ Solution by Mathematica
Time used: 0.0 (sec). Leaf size: 0
DSolve $\left[\left\{y[x] * y{ }^{\prime}{ }^{\prime}[x]-(y '[x]) \sim 2+(y '[x]) \wedge 3==0,\{y[0]==-1, y\right.\right.$ ' $\left.[0]==0\}\right\}, y[x], x$, IncludeSingularSoluti

### 1.79 problem 124

1.79.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 866
1.79.2 Solving as second order ode missing y ode . . . . . . . . . . . . 867
1.79.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 869

Internal problem ID [12496]
Internal file name [OUTPUT/11148_Monday_October_16_2023_09_52_00_PM_71604979/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 124.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "second_order_ode_missing_y"
Maple gives the following as the ode type
[[_2nd_order, _missing_y]]

$$
y^{\prime \prime}+y^{\prime} \tan (x)=\sin (2 x)
$$

With initial conditions

$$
\left[y(0)=-1, y^{\prime}(0)=0\right]
$$

### 1.79.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$
y^{\prime \prime}+p(x) y^{\prime}+q(x) y=F
$$

Where here

$$
\begin{aligned}
p(x) & =\tan (x) \\
q(x) & =0 \\
F & =\sin (2 x)
\end{aligned}
$$

Hence the ode is

$$
y^{\prime \prime}+y^{\prime} \tan (x)=\sin (2 x)
$$

The domain of $p(x)=\tan (x)$ is

$$
\left\{x<\frac{1}{2} \pi+\pi \_Z 117 \vee \frac{1}{2} \pi+\pi \_Z 117<x\right\}
$$

And the point $x_{0}=0$ is inside this domain. The domain of $F=\sin (2 x)$ is

$$
\{-\infty<x<\infty\}
$$

And the point $x_{0}=0$ is also inside this domain. Hence solution exists and is unique.

### 1.79.2 Solving as second order ode missing y ode

This is second order ode with missing dependent variable $y$. Let

$$
p(x)=y^{\prime}
$$

Then

$$
p^{\prime}(x)=y^{\prime \prime}
$$

Hence the ode becomes

$$
p^{\prime}(x)+p(x) \tan (x)-\sin (2 x)=0
$$

Which is now solve for $p(x)$ as first order ode.
Entering Linear first order ODE solver. The integrating factor $\mu$ is

$$
\begin{aligned}
\mu & =\mathrm{e}^{\int \tan (x) d x} \\
& =\frac{1}{\cos (x)}
\end{aligned}
$$

Which simplifies to

$$
\mu=\sec (x)
$$

The ode becomes

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu p) & =(\mu)(\sin (2 x)) \\
\frac{\mathrm{d}}{\mathrm{~d} x}(\sec (x) p) & =(\sec (x))(\sin (2 x)) \\
\mathrm{d}(\sec (x) p) & =(2 \sin (x)) \mathrm{d} x
\end{aligned}
$$

## Integrating gives

$$
\begin{aligned}
& \sec (x) p=\int 2 \sin (x) \mathrm{d} x \\
& \sec (x) p=-2 \cos (x)+c_{1}
\end{aligned}
$$

Dividing both sides by the integrating factor $\mu=\sec (x)$ results in

$$
p(x)=-2 \cos (x)^{2}+c_{1} \cos (x)
$$

which simplifies to

$$
p(x)=\cos (x)\left(-2 \cos (x)+c_{1}\right)
$$

Initial conditions are used to solve for $c_{1}$. Substituting $x=0$ and $p=0$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
0=-2+c_{1} \\
c_{1}=2
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
p(x)=-2 \cos (x)(\cos (x)-1)
$$

Since $p=y^{\prime}$ then the new first order ode to solve is

$$
y^{\prime}=-2 \cos (x)(\cos (x)-1)
$$

Integrating both sides gives

$$
\begin{aligned}
y & =\int-2 \cos (x)(\cos (x)-1) \mathrm{d} x \\
& =-\cos (x) \sin (x)-x+2 \sin (x)+c_{2}
\end{aligned}
$$

Initial conditions are used to solve for $c_{2}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
-1=c_{2} \\
c_{2}=-1
\end{gathered}
$$

Substituting $c_{2}$ found above in the general solution gives

$$
y=-\cos (x) \sin (x)-x+2 \sin (x)-1
$$

Initial conditions are used to solve for the constants of integration.
Summary
The solution(s) found are the following

$$
\begin{equation*}
y=-\cos (x) \sin (x)-x+2 \sin (x)-1 \tag{1}
\end{equation*}
$$



Figure 141: Solution plot

Verification of solutions

$$
y=-\cos (x) \sin (x)-x+2 \sin (x)-1
$$

Verified OK.
1.79.3 Maple step by step solution

Let's solve

$$
\left[y^{\prime \prime}+y^{\prime} \tan (x)=\sin (2 x), y(0)=-1,\left.y^{\prime}\right|_{\{x=0\}}=0\right]
$$

- Highest derivative means the order of the ODE is 2
$\square$
- Make substitution $u=y^{\prime}$ to reduce order of ODE
$u^{\prime}(x)+u(x) \tan (x)=\sin (2 x)$
- Isolate the derivative
$u^{\prime}(x)=-u(x) \tan (x)+\sin (2 x)$
- Group terms with $u(x)$ on the lhs of the ODE and the rest on the rhs of the ODE
$u^{\prime}(x)+u(x) \tan (x)=\sin (2 x)$
- The ODE is linear; multiply by an integrating factor $\mu(x)$
$\mu(x)\left(u^{\prime}(x)+u(x) \tan (x)\right)=\mu(x) \sin (2 x)$
- $\quad$ Assume the lhs of the ODE is the total derivative $\frac{d}{d x}(\mu(x) u(x))$
$\mu(x)\left(u^{\prime}(x)+u(x) \tan (x)\right)=\mu^{\prime}(x) u(x)+\mu(x) u^{\prime}(x)$
- Isolate $\mu^{\prime}(x)$
$\mu^{\prime}(x)=\mu(x) \tan (x)$
- $\quad$ Solve to find the integrating factor
$\mu(x)=\frac{1}{\cos (x)}$
- Integrate both sides with respect to $x$
$\int\left(\frac{d}{d x}(\mu(x) u(x))\right) d x=\int \mu(x) \sin (2 x) d x+c_{1}$
- Evaluate the integral on the lhs

$$
\mu(x) u(x)=\int \mu(x) \sin (2 x) d x+c_{1}
$$

- $\quad$ Solve for $u(x)$
$u(x)=\frac{\int \mu(x) \sin (2 x) d x+c_{1}}{\mu(x)}$
- $\quad$ Substitute $\mu(x)=\frac{1}{\cos (x)}$
$u(x)=\cos (x)\left(\int \frac{\sin (2 x)}{\cos (x)} d x+c_{1}\right)$
- Evaluate the integrals on the rhs
$u(x)=\cos (x)\left(-2 \cos (x)+c_{1}\right)$
- $\quad$ Solve 1st ODE for $u(x)$

$$
u(x)=\cos (x)\left(-2 \cos (x)+c_{1}\right)
$$

- $\quad$ Make substitution $u=y^{\prime}$

$$
y^{\prime}=\cos (x)\left(-2 \cos (x)+c_{1}\right)
$$

- Integrate both sides to solve for $y$
$\int y^{\prime} d x=\int \cos (x)\left(-2 \cos (x)+c_{1}\right) d x+c_{2}$
- Compute integrals
$y=-\cos (x) \sin (x)-x+\sin (x) c_{1}+c_{2}$

Check validity of solution $y=-\cos (x) \sin (x)-x+\sin (x) c_{1}+c_{2}$

- Use initial condition $y(0)=-1$

$$
-1=c_{2}
$$

- Compute derivative of the solution

$$
y^{\prime}=\sin (x)^{2}-\cos (x)^{2}-1+c_{1} \cos (x)
$$

- Use the initial condition $\left.y^{\prime}\right|_{\{x=0\}}=0$

$$
0=-2+c_{1}
$$

- Solve for $c_{1}$ and $c_{2}$

$$
\left\{c_{1}=2, c_{2}=-1\right\}
$$

- Substitute constant values into general solution and simplify

$$
y=\sin (x)(2-\cos (x))-x-1
$$

- $\quad$ Solution to the IVP

$$
y=\sin (x)(2-\cos (x))-x-1
$$

Maple trace

```
-Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
-> Calling odsolve with the ODE`, diff(_b(_a), _a) = -tan(_a)*_b(_a)+sin(2*_a), _b(_a)`
    Methods for first order ODEs:
    --- Trying classification methods ---
    trying a quadrature
    trying 1st order linear
    <- 1st order linear successful
<- high order exact linear fully integrable successful`
```

$\checkmark$ Solution by Maple
Time used: 0.047 (sec). Leaf size: 19
dsolve $([\operatorname{diff}(y(x), x \$ 2)+\tan (x) * \operatorname{diff}(y(x), x)=\sin (2 * x), y(0)=-1, D(y)(0)=0], y(x)$, singsol=al

$$
y(x)=-x-1+2 \sin (x)-\frac{\sin (2 x)}{2}
$$

$\checkmark$ Solution by Mathematica
Time used: 0.154 (sec). Leaf size: 18
DSolve[\{y' ' $[x]+\operatorname{Tan}[x] * y$ ' $\left.[x]==\operatorname{Sin}[2 * x],\left\{y[0]==-1, y^{\prime}[0]==0\right\}\right\}, y[x], x$, IncludeSingularSolutions

$$
y(x) \rightarrow-x-\sin (x)(\cos (x)-2)-1
$$

### 1.80 problem 125

1.80.1 Solving as second order ode missing y ode . . . . . . . . . . . . 873
1.80.2 Solving as second order ode missing x ode . . . . . . . . . . . . 876
1.80.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 883

Internal problem ID [12497]
Internal file name [OUTPUT/11149_Monday_October_16_2023_09_52_02_PM_17799216/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 125.
ODE order: 2.
ODE degree: 2.

The type(s) of ODE detected by this program : "second_order_ode_missing_x", "second__order__ode_missing_y"

Maple gives the following as the ode type
[[_2nd_order, _missing_x]]

$$
y^{\prime \prime 2}+y^{\prime 2}=a^{2}
$$

With initial conditions

$$
\left[y(0)=-1, y^{\prime}(0)=0\right]
$$

### 1.80.1 Solving as second order ode missing y ode

This is second order ode with missing dependent variable $y$. Let

$$
p(x)=y^{\prime}
$$

Then

$$
p^{\prime}(x)=y^{\prime \prime}
$$

Hence the ode becomes

$$
p^{\prime}(x)^{2}+p(x)^{2}-a^{2}=0
$$

Which is now solve for $p(x)$ as first order ode. Solving the given ode for $p^{\prime}(x)$ results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$
\begin{align*}
& p^{\prime}(x)=\sqrt{-p(x)^{2}+a^{2}}  \tag{1}\\
& p^{\prime}(x)=-\sqrt{-p(x)^{2}+a^{2}} \tag{2}
\end{align*}
$$

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{\sqrt{a^{2}-p^{2}}} d p & =\int d x \\
\arctan \left(\frac{p(x)}{\sqrt{-p(x)^{2}+a^{2}}}\right) & =x+c_{1}
\end{aligned}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $x=0$ and $p=0$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& 0=c_{1} \\
& c_{1}=0
\end{aligned}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
\arctan \left(\frac{p}{\sqrt{a^{2}-p^{2}}}\right)=x
$$

Solving equation (2)
Integrating both sides gives

$$
\begin{aligned}
\int-\frac{1}{\sqrt{a^{2}-p^{2}}} d p & =\int d x \\
-\arctan \left(\frac{p(x)}{\sqrt{-p(x)^{2}+a^{2}}}\right) & =x+c_{2}
\end{aligned}
$$

Initial conditions are used to solve for $c_{2}$. Substituting $x=0$ and $p=0$ in the above solution gives an equation to solve for the constant of integration.

$$
0=c_{2}
$$

$$
c_{2}=0
$$

Substituting $c_{2}$ found above in the general solution gives

$$
-\arctan \left(\frac{p}{\sqrt{a^{2}-p^{2}}}\right)=x
$$

For solution (1) found earlier, since $p=y^{\prime}$ then we now have a new first order ode to solve which is

$$
\arctan \left(\frac{y^{\prime}}{\sqrt{-y^{\prime 2}+a^{2}}}\right)=x
$$

Integrating both sides gives

$$
\begin{aligned}
y & =\int \tan (x) \sqrt{\frac{a^{2}}{\tan (x)^{2}+1}} \mathrm{~d} x \\
& =-\sqrt{\frac{a^{2}}{\tan (x)^{2}+1}}+c_{3}
\end{aligned}
$$

Initial conditions are used to solve for $c_{3}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
-1=-\operatorname{csgn}(a) a+c_{3} \\
c_{3}=a-1
\end{gathered}
$$

Substituting $c_{3}$ found above in the general solution gives

$$
y=-\sqrt{\cos (x)^{2} a^{2}}+a-1
$$

For solution (2) found earlier, since $p=y^{\prime}$ then we now have a new first order ode to solve which is

$$
-\arctan \left(\frac{y^{\prime}}{\sqrt{-y^{\prime 2}+a^{2}}}\right)=x
$$

Integrating both sides gives

$$
\begin{aligned}
y & =\int-\tan (x) \sqrt{\frac{a^{2}}{\tan (x)^{2}+1}} \mathrm{~d} x \\
& =\sqrt{\frac{a^{2}}{\tan (x)^{2}+1}}+c_{4}
\end{aligned}
$$

Initial conditions are used to solve for $c_{4}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
-1=\operatorname{csgn}(a) a+c_{4} \\
c_{4}=-a-1
\end{gathered}
$$

Substituting $c_{4}$ found above in the general solution gives

$$
y=\sqrt{\cos (x)^{2} a^{2}}-a-1
$$

Initial conditions are used to solve for the constants of integration.
Summary
The solution(s) found are the following

$$
\begin{align*}
& y=-\sqrt{\cos (x)^{2} a^{2}}+a-1  \tag{1}\\
& y=\sqrt{\cos (x)^{2} a^{2}}-a-1 \tag{2}
\end{align*}
$$

Verification of solutions

$$
y=-\sqrt{\cos (x)^{2} a^{2}}+a-1
$$

Verified OK.

$$
y=\sqrt{\cos (x)^{2} a^{2}}-a-1
$$

Verified OK.

### 1.80.2 Solving as second order ode missing x ode

This is missing independent variable second order ode. Solved by reduction of order by using substitution which makes the dependent variable $y$ an independent variable. Using

$$
y^{\prime}=p(y)
$$

Then

$$
\begin{aligned}
y^{\prime \prime} & =\frac{d p}{d x} \\
& =\frac{d y}{d x} \frac{d p}{d y} \\
& =p \frac{d p}{d y}
\end{aligned}
$$

Hence the ode becomes

$$
p(y)^{2}\left(\frac{d}{d y} p(y)\right)^{2}+p(y)^{2}=a^{2}
$$

Which is now solved as first order ode for $p(y)$. Solving the given ode for $\frac{d}{d y} p(y)$ results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$
\begin{align*}
\frac{d}{d y} p(y) & =\frac{\sqrt{-p(y)^{2}+a^{2}}}{p(y)}  \tag{1}\\
\frac{d}{d y} p(y) & =-\frac{\sqrt{-p(y)^{2}+a^{2}}}{p(y)} \tag{2}
\end{align*}
$$

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives

$$
\begin{aligned}
\int \frac{p}{\sqrt{a^{2}-p^{2}}} d p & =\int d y \\
-\frac{(a-p(y))(p(y)+a)}{\sqrt{-p(y)^{2}+a^{2}}} & =y+c_{1}
\end{aligned}
$$

Initial conditions are used to solve for $c_{1}$. Substituting $y=-1$ and $p=0$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
-\operatorname{csgn}(a) a=-1+c_{1} \\
c_{1}=-a+1
\end{gathered}
$$

Substituting $c_{1}$ found above in the general solution gives

$$
-\frac{(a-p)(a+p)}{\sqrt{a^{2}-p^{2}}}=y-a+1
$$

The above simplifies to

$$
a \sqrt{a^{2}-p^{2}}-y \sqrt{a^{2}-p^{2}}-a^{2}+p^{2}-\sqrt{a^{2}-p^{2}}=0
$$

Solving equation (2)

Integrating both sides gives

$$
\begin{aligned}
\int-\frac{p}{\sqrt{a^{2}-p^{2}}} d p & =\int d y \\
\frac{(a-p(y))(p(y)+a)}{\sqrt{-p(y)^{2}+a^{2}}} & =y+c_{2}
\end{aligned}
$$

Initial conditions are used to solve for $c_{2}$. Substituting $y=-1$ and $p=0$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{gathered}
\operatorname{csgn}(a) a=-1+c_{2} \\
c_{2}=a+1
\end{gathered}
$$

Substituting $c_{2}$ found above in the general solution gives

$$
\frac{(a-p)(a+p)}{\sqrt{a^{2}-p^{2}}}=y+a+1
$$

The above simplifies to

$$
-a \sqrt{a^{2}-p^{2}}-y \sqrt{a^{2}-p^{2}}+a^{2}-p^{2}-\sqrt{a^{2}-p^{2}}=0
$$

For solution (1) found earlier, since $p=y^{\prime}$ then we now have a new first order ode to solve which is

$$
(a-y-1) \sqrt{-y^{\prime 2}+a^{2}}-a^{2}+y^{\prime 2}=0
$$

Solving the given ode for $y^{\prime}$ results in 4 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$
\begin{align*}
& y^{\prime}=a  \tag{1}\\
& y^{\prime}=-a  \tag{2}\\
& y^{\prime}=\sqrt{-1+2 a y-y^{2}+2 a-2 y}  \tag{3}\\
& y^{\prime}=-\sqrt{-1+2 a y-y^{2}+2 a-2 y} \tag{4}
\end{align*}
$$

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives

$$
\begin{aligned}
y & =\int a \mathrm{~d} x \\
& =a x+c_{3}
\end{aligned}
$$

Initial conditions are used to solve for $c_{3}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& -1=c_{3} \\
& c_{3}=-1
\end{aligned}
$$

Substituting $c_{3}$ found above in the general solution gives

$$
y=a x-1
$$

Solving equation (2)
Integrating both sides gives

$$
\begin{aligned}
y & =\int-a \mathrm{~d} x \\
& =-a x+c_{4}
\end{aligned}
$$

Initial conditions are used to solve for $c_{4}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& -1=c_{4} \\
& c_{4}=-1
\end{aligned}
$$

Substituting $c_{4}$ found above in the general solution gives

$$
y=-a x-1
$$

Solving equation (3)
Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{\sqrt{2 a y-y^{2}+2 a-2 y-1}} d y & =\int d x \\
\arctan \left(\frac{-a+y+1}{\sqrt{-y^{2}+(-2+2 a) y-1+2 a}}\right) & =x+c_{5}
\end{aligned}
$$

Initial conditions are used to solve for $c_{5}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration. Unable to solve for
constant of integration. Warning: Unable to solve for constant of integration. Unable to determine ODE type.

Solving equation (4)
Integrating both sides gives

$$
\begin{aligned}
\int-\frac{1}{\sqrt{2 a y-y^{2}+2 a-2 y-1}} d y & =\int d x \\
-\arctan \left(\frac{-a+y+1}{\sqrt{-y^{2}+(-2+2 a) y-1+2 a}}\right) & =x+c_{6}
\end{aligned}
$$

Initial conditions are used to solve for $c_{6}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration. Unable to solve for constant of integration. Warning: Unable to solve for constant of integration. Unable to determine ODE type.
For solution (2) found earlier, since $p=y^{\prime}$ then we now have a new first order ode to solve which is

$$
(-a-y-1) \sqrt{-y^{\prime 2}+a^{2}}+a^{2}-y^{\prime 2}=0
$$

Solving the given ode for $y^{\prime}$ results in 4 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$
\begin{align*}
& y^{\prime}=a  \tag{1}\\
& y^{\prime}=-a  \tag{2}\\
& y^{\prime}=\sqrt{-1-2 a y-y^{2}-2 a-2 y}  \tag{3}\\
& y^{\prime}=-\sqrt{-1-2 a y-y^{2}-2 a-2 y} \tag{4}
\end{align*}
$$

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives

$$
\begin{aligned}
y & =\int a \mathrm{~d} x \\
& =a x+c_{7}
\end{aligned}
$$

Initial conditions are used to solve for $c_{7}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration.

$$
-1=c_{7}
$$

$$
c_{7}=-1
$$

Substituting $c_{7}$ found above in the general solution gives

$$
y=a x-1
$$

Solving equation (2)
Integrating both sides gives

$$
\begin{aligned}
y & =\int-a \mathrm{~d} x \\
& =-a x+c_{8}
\end{aligned}
$$

Initial conditions are used to solve for $c_{8}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration.

$$
\begin{aligned}
& -1=c_{8} \\
& c_{8}=-1
\end{aligned}
$$

Substituting $c_{8}$ found above in the general solution gives

$$
y=-a x-1
$$

Solving equation (3)
Integrating both sides gives

$$
\begin{aligned}
\int \frac{1}{\sqrt{-2 a y-y^{2}-2 a-2 y-1}} d y & =\int d x \\
\arctan \left(\frac{a+y+1}{\sqrt{-y^{2}+(-2-2 a) y-1-2 a}}\right) & =x+c_{9}
\end{aligned}
$$

Initial conditions are used to solve for $c_{9}$. Substituting $x=0$ and $y=-1$ in the above solution gives an equation to solve for the constant of integration. Unable to solve for constant of integration. Warning: Unable to solve for constant of integration. Unable to determine ODE type.

Solving equation (4)

Integrating both sides gives

$$
\begin{aligned}
\int-\frac{1}{\sqrt{-2 a y-y^{2}-2 a-2 y-1}} d y & =\int d x \\
-\arctan \left(\frac{a+y+1}{\sqrt{-y^{2}+(-2-2 a) y-1-2 a}}\right) & =x+\_C 10
\end{aligned}
$$

Initial conditions are used to solve for the constants of integration.
Looking at the Fifth solution

$$
\begin{equation*}
y=-\tan (x+\ldots C 10) \sqrt{\frac{a^{2}}{\tan (x+\ldots C 10)^{2}+1}}-a-1 \tag{5}
\end{equation*}
$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y=-1$ and $x=0$ in the above gives

$$
\begin{equation*}
-1=-\tan \left(\_C 10\right) \sqrt{\cos \left(\_C 10\right)^{2} a^{2}}-a-1 \tag{1~A}
\end{equation*}
$$

Taking derivative of the solution gives

$$
y^{\prime}=-\left(\tan \left(x+\_C 10\right)^{2}+1\right) \sqrt{\frac{a^{2}}{\tan (x+\ldots C 10)^{2}+1}}+\frac{\tan (x+\ldots C 10)^{2} a^{2}}{\sqrt{\frac{a^{2}}{\tan (x+\ldots C 10)^{2}+1}}\left(\tan (x+\ldots C 10)^{2}+1\right)}
$$

substituting $y^{\prime}=0$ and $x=0$ in the above gives

$$
\begin{equation*}
0=-\frac{\cos \left(\_C 10\right)^{2} a^{2}}{\sqrt{\cos \left(\_C 10\right)^{2} a^{2}}} \tag{2~A}
\end{equation*}
$$

Equations $\{1 \mathrm{~A}, 2 \mathrm{~A}\}$ are now solved for $\left\{\_C 10\right\}$. There is no solution for the constants of integrations. This solution is removed.
Summary
The solution(s) found are the following

$$
\begin{align*}
& y=a x-1  \tag{1}\\
& y=-a x-1  \tag{2}\\
& y=a x-1  \tag{3}\\
& y=-a x-1 \tag{4}
\end{align*}
$$

Verification of solutions

$$
y=a x-1
$$

Verified OK.

$$
y=-a x-1
$$

Verified OK.

$$
y=a x-1
$$

Verified OK.

$$
y=-a x-1
$$

Verified OK.

### 1.80.3 Maple step by step solution

Let's solve

$$
\left[y^{\prime \prime 2}+y^{\prime 2}=a^{2}, y(0)=-1,\left.y^{\prime}\right|_{\{x=0\}}=0\right]
$$

- Highest derivative means the order of the ODE is 2

```
y'
```

- Make substitution $u=y^{\prime}$ to reduce order of ODE

$$
u^{\prime}(x)^{2}+u(x)^{2}=a^{2}
$$

- $\quad$ Separate variables

$$
\frac{u^{\prime}(x)}{\sqrt{-u(x)^{2}+a^{2}}}=1
$$

- Integrate both sides with respect to $x$
$\int \frac{u^{\prime}(x)}{\sqrt{-u(x)^{2}+a^{2}}} d x=\int 1 d x+c_{1}$
- Evaluate integral

$$
\arctan \left(\frac{u(x)}{\sqrt{-u(x)^{2}+a^{2}}}\right)=x+c_{1}
$$

- $\quad$ Solve for $u(x)$

$$
u(x)=\tan \left(x+c_{1}\right) \sqrt{\frac{a^{2}}{\tan \left(x+c_{1}\right)^{2}+1}}
$$

- $\quad$ Solve 1st ODE for $u(x)$

$$
u(x)=\tan \left(x+c_{1}\right) \sqrt{\frac{a^{2}}{\tan \left(x+c_{1}\right)^{2}+1}}
$$

- $\quad$ Make substitution $u=y^{\prime}$

$$
y^{\prime}=\tan \left(x+c_{1}\right) \sqrt{\frac{a^{2}}{\tan \left(x+c_{1}\right)^{2}+1}}
$$

- Integrate both sides to solve for $y$
$\int y^{\prime} d x=\int \tan \left(x+c_{1}\right) \sqrt{\frac{a^{2}}{\tan \left(x+c_{1}\right)^{2}+1}} d x+c_{2}$
- Compute integrals

$$
y=-\sqrt{\frac{a^{2}}{\tan \left(x+c_{1}\right)^{2}+1}}+c_{2}
$$

$$
\text { Check validity of solution } y=-\sqrt{\frac{a^{2}}{\tan \left(x+c_{1}\right)^{2}+1}}+c_{2}
$$

- Use initial condition $y(0)=-1$

$$
-1=-\sqrt{\frac{a^{2}}{\tan \left(c_{1}\right)^{2}+1}}+c_{2}
$$

- Compute derivative of the solution

$$
y^{\prime}=\frac{a^{2} \tan \left(x+c_{1}\right)}{\sqrt{\frac{a^{2}}{\tan \left(x+c_{1}\right)^{2}+1}}\left(\tan \left(x+c_{1}\right)^{2}+1\right)}
$$

- Use the initial condition $\left.y^{\prime}\right|_{\{x=0\}}=0$

$$
0=\frac{a^{2} \tan \left(c_{1}\right)}{\sqrt{\frac{a^{2}}{\tan \left(c_{1}\right)^{2}+1}}\left(\tan \left(c_{1}\right)^{2}+1\right)}
$$

- Solve for $c_{1}$ and $c_{2}$

$$
\left\{c_{1}=0, c_{2}=-1+\sqrt{a^{2}}\right\}
$$

- Substitute constant values into general solution and simplify

$$
y=-\sqrt{\cos (x)^{2} a^{2}}-1+\operatorname{csgn}(a) a
$$

- $\quad$ Solution to the IVP

$$
y=-\sqrt{\cos (x)^{2} a^{2}}-1+\operatorname{csgn}(a) a
$$

Maple trace

```
`Methods for second order ODEs:
    *** Sublevel 2 ***
    Methods for second order ODEs:
    Successful isolation of d^2y/dx^2: 2 solutions were found. Trying to solve each resulting
        *** Sublevel 3 ***
        Methods for second order ODEs:
        --- Trying classification methods ---
        trying 2nd order Liouville
        trying 2nd order WeierstrassP
        trying 2nd order JacobiSN
        differential order: 2; trying a linearization to 3rd order
        trying 2nd order ODE linearizable_by_differentiation
        -> Calling odsolve with the ODE`, diff(diff(diff(y(x), x), x), x)+diff(y(x), x), y(x)
            Methods for third order ODEs:
            --- Trying classification methods ---
            trying a quadrature
            checking if the LODE has constant coefficients
            <- constant coefficients successful
            <- 2nd order ODE linearizable_by_differentiation successful
    * Tackling next ODE.
        *** Sublevel 3 ***
        Methods for second order ODEs:
        --- Trying classification methods ---
        trying 2nd order Liouville
        trying 2nd order WeierstrassP
        trying 2nd order JacobiSN
        differential order: 2; trying a linearization to 3rd order
        trying 2nd order ODE linearizable_by_differentiation
        <- 2nd order ODE linearizable_by_differentiation successful
-> Calling odsolve with the ODE`, diff(y(x), x) = -a, y(x), singsol = none`
    Methods for first order ODEs:
    --- Trying classification methods ---
    trying a quadrature
    trying 1st order linear
    <- 1st order linear successful
-> Calling odsolve with the ODE`, diff(y(x), x) = a, y(x), singsol = none` *** Sublevel 2
    Methods for first order ODEs:
    --- Trying classification methods ---
    trying a quadrature
    trying 1st order linear
    <- 1st order linear successful`
\(\checkmark\) Solution by Maple
Time used: 0.657 (sec). Leaf size: 24
dsolve([diff \((y(x), x \$ 2) \wedge 2+\operatorname{diff}(y(x), x) \wedge 2=a \wedge 2, y(0)=-1, D(y)(0)=0], y(x)\), singsol=all)
\[
\begin{aligned}
& y(x)=-a-1+a \cos (x) \\
& y(x)=a-1-a \cos (x)
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 15.637 (sec). Leaf size: 37
DSolve \(\left[\left\{\left(y^{\prime \prime}[x]\right) \wedge 2+(y \prime[x]) \wedge 2==a \wedge 2,\left\{y[0]==-1, y^{\prime}[0]==0\right\}\right\}, y[x], x\right.\), IncludeSingularSolutions \(\rightarrow \mathrm{Tr}\)
\[
\begin{aligned}
& y(x) \rightarrow a\left(\frac{1}{\sqrt{\sec ^{2}(x)}}-1\right)-1 \\
& y(x) \rightarrow-\frac{a}{\sqrt{\sec ^{2}(x)}}+a-1
\end{aligned}
\]

\subsection*{1.81 problem 126}
\[
\text { 1.81.1 Solving as second order integrable as is ode . . . . . . . . . . . } 887
\]
1.81.2 Solving as second order ode missing y ode ..... 889
1.81.3 Solving as second order ode missing x ode ..... 890
1.81.4 Solving as exact nonlinear second order ode ode ..... 892
1.81.5 Maple step by step solution ..... 894

Internal problem ID [12498]
Internal file name [OUTPUT/11150_Monday_October_16_2023_09_54_03_PM_52250683/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 126.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "second_order_integrable_as_is", "second_order_ode_missing_x", "second_order_ode_missing_y", "exact nonlinear second order ode"

Maple gives the following as the ode type
```

[[_2nd_order, _missing_x], [_2nd_order, _reducible, _mu_poly_yn ]]

```
\[
y^{\prime \prime}-\frac{1}{2 y^{\prime}}=0
\]

\subsection*{1.81.1 Solving as second order integrable as is ode}

Integrating both sides of the ODE w.r.t \(x\) gives
\[
\begin{aligned}
& \int_{y^{\prime 2}=x+c_{1}} 2 y^{\prime} y^{\prime \prime} d x=\int 1 d x \\
&
\end{aligned}
\]

Which is now solved for \(y\). Solving the given ode for \(y^{\prime}\) results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are
\[
\begin{align*}
y^{\prime} & =\sqrt{x+c_{1}}  \tag{1}\\
y^{\prime} & =-\sqrt{x+c_{1}} \tag{2}
\end{align*}
\]

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives
\[
\begin{aligned}
y & =\int \sqrt{x+c_{1}} \mathrm{~d} x \\
& =\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\end{aligned}
\]

Solving equation (2)
Integrating both sides gives
\[
\begin{aligned}
y & =\int-\sqrt{x+c_{1}} \mathrm{~d} x \\
& =-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{align*}
& y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}  \tag{1}\\
& y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3} \tag{2}
\end{align*}
\]

Verification of solutions
\[
y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\]

Verified OK.
\[
y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
\]

Verified OK.

\subsection*{1.81.2 Solving as second order ode missing y ode}

This is second order ode with missing dependent variable \(y\). Let
\[
p(x)=y^{\prime}
\]

Then
\[
p^{\prime}(x)=y^{\prime \prime}
\]

Hence the ode becomes
\[
2 p(x) p^{\prime}(x)-1=0
\]

Which is now solve for \(p(x)\) as first order ode. Integrating both sides gives
\[
\begin{aligned}
\int 2 p d p & =x+c_{1} \\
p^{2} & =x+c_{1}
\end{aligned}
\]

Solving for \(p\) gives these solutions
\[
\begin{aligned}
& p_{1}=\sqrt{x+c_{1}} \\
& p_{2}=-\sqrt{x+c_{1}}
\end{aligned}
\]

For solution (1) found earlier, since \(p=y^{\prime}\) then the new first order ode to solve is
\[
y^{\prime}=\sqrt{x+c_{1}}
\]

Integrating both sides gives
\[
\begin{aligned}
y & =\int \sqrt{x+c_{1}} \mathrm{~d} x \\
& =\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\end{aligned}
\]

Since \(p=y^{\prime}\) then the new first order ode to solve is
\[
y^{\prime}=-\sqrt{x+c_{1}}
\]

Integrating both sides gives
\[
\begin{aligned}
y & =\int-\sqrt{x+c_{1}} \mathrm{~d} x \\
& =-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{align*}
& y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}  \tag{1}\\
& y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3} \tag{2}
\end{align*}
\]

Verification of solutions
\[
y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\]

Verified OK.
\[
y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
\]

Verified OK.

\subsection*{1.81.3 Solving as second order ode missing \(x\) ode}

This is missing independent variable second order ode. Solved by reduction of order by using substitution which makes the dependent variable \(y\) an independent variable. Using
\[
y^{\prime}=p(y)
\]

Then
\[
\begin{aligned}
y^{\prime \prime} & =\frac{d p}{d x} \\
& =\frac{d y}{d x} \frac{d p}{d y} \\
& =p \frac{d p}{d y}
\end{aligned}
\]

Hence the ode becomes
\[
2 p(y)^{2}\left(\frac{d}{d y} p(y)\right)=1
\]

Which is now solved as first order ode for \(p(y)\). Integrating both sides gives
\[
\begin{aligned}
\int 2 p^{2} d p & =y+c_{1} \\
\frac{2 p^{3}}{3} & =y+c_{1}
\end{aligned}
\]

Solving for \(p\) gives these solutions
\[
\begin{aligned}
& p_{1}=\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{2} \\
& p_{2}=-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}-\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4} \\
& p_{3}=-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}+\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}
\end{aligned}
\]

For solution (1) found earlier, since \(p=y^{\prime}\) then we now have a new first order ode to solve which is
\[
y^{\prime}=\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{2}
\]

Integrating both sides gives
\[
\begin{aligned}
\int \frac{2}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}} d y & =\int d x \\
\frac{3 y+3 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}} & =x+c_{2}
\end{aligned}
\]

For solution (2) found earlier, since \(p=y^{\prime}\) then we now have a new first order ode to solve which is
\[
y^{\prime}=-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}-\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}
\]

Integrating both sides gives
\[
\begin{aligned}
& \int \frac{1}{-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}-\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}} d y \\
&=\int d x \\
&-\frac{6\left(y+c_{1}\right)}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(1+i \sqrt{3})}=c_{3}+x
\end{aligned}
\]

For solution (3) found earlier, since \(p=y^{\prime}\) then we now have a new first order ode to solve which is
\[
y^{\prime}=-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}+\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}
\]

Integrating both sides gives
\[
\begin{aligned}
\int \frac{1}{-\frac{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}+\frac{i \sqrt{3}\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}{4}} d y & =\int d x \\
\frac{6 y+6 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(i \sqrt{3}-1)} & =x+c_{4}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{align*}
\frac{3 y+3 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}} & =x+c_{2}  \tag{1}\\
-\frac{6\left(y+c_{1}\right)}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(1+i \sqrt{3})} & =c_{3}+x  \tag{2}\\
\frac{6 y+6 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(i \sqrt{3}-1)} & =x+c_{4} \tag{3}
\end{align*}
\]

Verification of solutions
\[
\frac{3 y+3 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}}=x+c_{2}
\]

Verified OK.
\[
-\frac{6\left(y+c_{1}\right)}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(1+i \sqrt{3})}=c_{3}+x
\]

Verified OK.
\[
\frac{6 y+6 c_{1}}{\left(12 y+12 c_{1}\right)^{\frac{1}{3}}(i \sqrt{3}-1)}=x+c_{4}
\]

Verified OK.

\subsection*{1.81.4 Solving as exact nonlinear second order ode ode}

An exact non-linear second order ode has the form
\[
a_{2}\left(x, y, y^{\prime}\right) y^{\prime \prime}+a_{1}\left(x, y, y^{\prime}\right) y^{\prime}+a_{0}\left(x, y, y^{\prime}\right)=0
\]

Where the following conditions are satisfied
\[
\begin{aligned}
\frac{\partial a_{2}}{\partial y} & =\frac{\partial a_{1}}{\partial y^{\prime}} \\
\frac{\partial a_{2}}{\partial x} & =\frac{\partial a_{0}}{\partial y^{\prime}} \\
\frac{\partial a_{1}}{\partial x} & =\frac{\partial a_{0}}{\partial y}
\end{aligned}
\]

Looking at the the ode given we see that
\[
\begin{aligned}
& a_{2}=2 y^{\prime} \\
& a_{1}=0 \\
& a_{0}=-1
\end{aligned}
\]

Applying the conditions to the above shows this is a nonlinear exact second order ode. Therefore it can be reduced to first order ode given by
\[
\begin{aligned}
& \int a_{2} d y^{\prime}+\int a_{1} d y+\int a_{0} d x=c_{1} \\
& \int 2 y^{\prime} d y^{\prime}+\int 0 d y+\int-1 d x=c_{1}
\end{aligned}
\]

Which results in
\[
y^{\prime 2}-x=c_{1}
\]

Which is now solved Solving the given ode for \(y^{\prime}\) results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are
\[
\begin{align*}
y^{\prime} & =\sqrt{x+c_{1}}  \tag{1}\\
y^{\prime} & =-\sqrt{x+c_{1}} \tag{2}
\end{align*}
\]

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives
\[
\begin{aligned}
y & =\int \sqrt{x+c_{1}} \mathrm{~d} x \\
& =\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\end{aligned}
\]

Solving equation (2)

Integrating both sides gives
\[
\begin{aligned}
y & =\int-\sqrt{x+c_{1}} \mathrm{~d} x \\
& =-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{align*}
& y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}  \tag{1}\\
& y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3} \tag{2}
\end{align*}
\]

Verification of solutions
\[
y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\]

Verified OK.
\[
y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{3}
\]

Verified OK.

\subsection*{1.81.5 Maple step by step solution}

Let's solve
\[
2 y^{\prime} y^{\prime \prime}=1
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Make substitution \(u=y^{\prime}\) to reduce order of ODE
\[
2 u(x) u^{\prime}(x)=1
\]
- Integrate both sides with respect to \(x\)
\(\int 2 u(x) u^{\prime}(x) d x=\int 1 d x+c_{1}\)
- Evaluate integral
\(u(x)^{2}=x+c_{1}\)
- \(\quad\) Solve for \(u(x)\)
\[
\left\{u(x)=\sqrt{x+c_{1}}, u(x)=-\sqrt{x+c_{1}}\right\}
\]
- \(\quad\) Solve 1st ODE for \(u(x)\)
\[
u(x)=\sqrt{x+c_{1}}
\]
- Make substitution \(u=y^{\prime}\)
\[
y^{\prime}=\sqrt{x+c_{1}}
\]
- Integrate both sides to solve for \(y\)
\[
\int y^{\prime} d x=\int \sqrt{x+c_{1}} d x+c_{2}
\]
- Compute integrals
\[
y=\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\]
- \(\quad\) Solve 2nd ODE for \(u(x)\)
\[
u(x)=-\sqrt{x+c_{1}}
\]
- Make substitution \(u=y^{\prime}\)
\[
y^{\prime}=-\sqrt{x+c_{1}}
\]
- Integrate both sides to solve for \(y\)
\(\int y^{\prime} d x=\int-\sqrt{x+c_{1}} d x+c_{2}\)
- Compute integrals
\[
y=-\frac{2\left(x+c_{1}\right)^{\frac{3}{2}}}{3}+c_{2}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying 2nd order Liouville trying 2nd order WeierstrassP trying 2nd order JacobiSN differential order: 2; trying a linearization to 3rd order trying 2nd order ODE linearizable_by_differentiation trying 2nd order, 2 integrating factors of the form mu(x,y) trying differential order: 2; missing variables `, `-> Computing symmetries using: way = 3 <- differential order: 2; canonical coordinates successful <- differential order 2; missing variables successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.032 (sec). Leaf size: 39
```

dsolve(diff(y(x),x\$2)=1/(2*\operatorname{diff}(y(x),x)),y(x), singsol=all)

```
\[
\begin{aligned}
& y(x)=\frac{\left(2 x+2 c_{1}\right) \sqrt{c_{1}+x}}{3}+c_{2} \\
& y(x)=\frac{\left(-2 x-2 c_{1}\right) \sqrt{c_{1}+x}}{3}+c_{2}
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.024 (sec). Leaf size: 43
DSolve[y''[x]==1/(2*y'[x]),y[x],x,IncludeSingularSolutions \(\rightarrow\) True]
\[
\begin{aligned}
& y(x) \rightarrow c_{2}-\frac{2}{3}\left(x+2 c_{1}\right)^{3 / 2} \\
& y(x) \rightarrow \frac{2}{3}\left(x+2 c_{1}\right)^{3 / 2}+c_{2}
\end{aligned}
\]

\subsection*{1.82 problem 127}

Internal problem ID [12499]
Internal file name [OUTPUT/11151_Monday_October_16_2023_09_54_06_PM_81038991/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.

Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 127.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "unknown"
Maple gives the following as the ode type
```

[[_3rd_order, _missing_x], [_3rd_order, _missing_y], [
_3rd_order, _with_linear_symmetries], [_3rd_order, _reducible
, _mu_y2], [_3rd_order, _reducible, _mu_poly_yn]]

```

Unable to solve or complete the solution.
Unable to parse ODE.
Maple trace
```

Methods for third order ODEs:
--- Trying classification methods ---
trying 3rd order ODE linearizable_by_differentiation
differential order: 3; trying a linearization to 4th order
trying differential order: 3; missing variables
`, `-> Computing symmetries using: way = 3
-> Calling odsolve with the ODE`, diff(_b(_a), _a) = _b(_a)~2, _b(_a), HINT = [[1, 0], [_a,     symmetry methods on request `, `1st order, trying reduction of order with given symmetries:` [1, 0], [_a, __b]

```
\(\checkmark\) Solution by Maple
Time used: 0.078 (sec). Leaf size: 24
```

dsolve(diff(y(x),x\$3)=diff(y(x),x\$2)^2,y(x), singsol=all)

```
\[
y(x)=\left(-c_{1}-x\right) \ln \left(c_{1}+x\right)+\left(c_{2}+1\right) x+c_{3}+c_{1}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.607 (sec). Leaf size: 24
DSolve[y'''[x]==(y''[x])~2,y[x],x,IncludeSingularSolutions \(\rightarrow\) True]
\[
y(x) \rightarrow x+c_{3} x-\left(x+c_{1}\right) \log \left(x+c_{1}\right)+c_{2}
\]

\subsection*{1.83 problem 128}

Internal problem ID [12500]
Internal file name [OUTPUT/11152_Monday_October_16_2023_09_54_06_PM_48316860/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 128.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "unknown"
Maple gives the following as the ode type
```

[[_3rd_order, _missing_x], [_3rd_order, _missing_y], [
_3rd_order, _with_linear_symmetries], [_3rd_order, _reducible
, _mu_y2], [_3rd_order, _reducible, _mu_poly_yn]]

```

Unable to solve or complete the solution.
Unable to parse ODE.
Maple trace
```

`Methods for third order ODEs: --- Trying classification methods --- trying 3rd order ODE linearizable_by_differentiation differential order: 3; trying a linearization to 4th order trying differential order: 3; missing variables `, `-> Computing symmetries using: way = 3 -> Calling odsolve with the ODE`, (diff(diff(_b(_a), _a), _a))*_b(_a)^2-2*(diff(_b(_a), _a))
symmetry methods on request
, `2nd order, trying reduction of order with given symmetries:`[_a, _b], [1, 0], [0, _b^2],

```
\(\checkmark\) Solution by Maple
Time used: 0.156 (sec). Leaf size: 59
dsolve(diff \((y(x), x) * \operatorname{diff}(y(x), x \$ 3)-3 * \operatorname{diff}(y(x), x \$ 2) \wedge 2=0, y(x), \quad\) singsol=all)
\[
\begin{aligned}
& y(x)=c_{1} \\
& y(x)=\frac{-c_{1} c_{2}+\sqrt{-2\left(-\frac{c_{1} c_{2}^{2}}{2}+x+c_{3}\right) c_{1}}}{c_{1}} \\
& y(x)=\frac{-c_{1} c_{2}-\sqrt{-2\left(-\frac{c_{1} c_{2}^{2}}{2}+x+c_{3}\right) c_{1}}}{c_{1}}
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.307 (sec). Leaf size: 21
DSolve[y'[x]*y'''[x]-3*(y''[x])~2==0,y[x],x,IncludeSingularSolutions -> True]
\[
y(x) \rightarrow c_{2} \sqrt{2 x+c_{1}}+c_{3}
\]

\subsection*{1.84 problem 129}
1.84.1 Solving as second order linear constant coeff ode . . . . . . . . 901
1.84.2 Solving as second order ode can be made integrable ode . . . . 903
1.84.3 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 905
1.84.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 909

Internal problem ID [12501]
Internal file name [OUTPUT/11153_Monday_October_16_2023_09_54_07_PM_79758852/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 129.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "second__order_ode_can_bbe_made_integrable"

Maple gives the following as the ode type
```

[[_2nd_order, _missing_x]]

```
\[
y^{\prime \prime}-9 y=0
\]

\subsection*{1.84.1 Solving as second order linear constant coeff ode}

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=-9\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-9 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\mathrm{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}-9=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=-9\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)(-9)} \\
& = \pm 3
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=+3 \\
& \lambda_{2}=-3
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=3 \\
& \lambda_{2}=-3
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{(3) x}+c_{2} e^{(-3) x}
\end{aligned}
\]

Or
\[
y=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{-3 x}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{-3 x} \tag{1}
\end{equation*}
\]


Figure 142: Slope field plot

Verification of solutions
\[
y=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{-3 x}
\]

Verified OK.

\subsection*{1.84.2 Solving as second order ode can be made integrable ode}

Multiplying the ode by \(y^{\prime}\) gives
\[
y^{\prime} y^{\prime \prime}-9 y y^{\prime}=0
\]

Integrating the above w.r.t \(x\) gives
\[
\begin{gathered}
\int\left(y^{\prime} y^{\prime \prime}-9 y y^{\prime}\right) d x=0 \\
\frac{y^{\prime 2}}{2}-\frac{9 y^{2}}{2}=c_{2}
\end{gathered}
\]

Which is now solved for \(y\). Solving the given ode for \(y^{\prime}\) results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are
\[
\begin{align*}
y^{\prime} & =\sqrt{9 y^{2}+2 c_{1}}  \tag{1}\\
y^{\prime} & =-\sqrt{9 y^{2}+2 c_{1}} \tag{2}
\end{align*}
\]

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives
\[
\begin{aligned}
\int \frac{1}{\sqrt{9 y^{2}+2 c_{1}}} d y & =\int d x \\
\frac{\ln \left(y \sqrt{9}+\sqrt{9 y^{2}+2 c_{1}}\right) \sqrt{9}}{9} & =x+c_{2}
\end{aligned}
\]

Raising both side to exponential gives
\[
\mathrm{e}^{\frac{\ln \left(y \sqrt{9}+\sqrt{9 y^{2}+2 c_{1}}\right) \sqrt{9}}{9}}=\mathrm{e}^{x+c_{2}}
\]

Which simplifies to
\[
\left(3 y+\sqrt{9 y^{2}+2 c_{1}}\right)^{\frac{1}{3}}=c_{3} \mathrm{e}^{x}
\]

Solving equation (2)
Integrating both sides gives
\[
\begin{aligned}
\int-\frac{1}{\sqrt{9 y^{2}+2 c_{1}}} d y & =\int d x \\
-\frac{\ln \left(y \sqrt{9}+\sqrt{9 y^{2}+2 c_{1}}\right) \sqrt{9}}{9} & =x+c_{4}
\end{aligned}
\]

Raising both side to exponential gives
\[
\mathrm{e}^{-\frac{\ln \left(y \sqrt{9}+\sqrt{9 y^{2}+2 c_{1}}\right) \sqrt{9}}{9}}=\mathrm{e}^{x+c_{4}}
\]

Which simplifies to
\[
\frac{1}{\left(3 y+\sqrt{9 y^{2}+2 c_{1}}\right)^{\frac{1}{3}}}=c_{5} \mathrm{e}^{x}
\]

Summary
The solution(s) found are the following
\[
\begin{align*}
& y=\frac{\left(\mathrm{e}^{6 x} c_{3}^{6}-2 c_{1}\right) \mathrm{e}^{-3 x}}{6 c_{3}^{3}}  \tag{1}\\
& y=-\frac{\left(2 c_{1} c_{5}^{6} \mathrm{e}^{6 x}-1\right) \mathrm{e}^{-3 x}}{6 c_{5}^{3}} \tag{2}
\end{align*}
\]


Figure 143: Slope field plot

Verification of solutions
\[
y=\frac{\left(\mathrm{e}^{6 x} c_{3}^{6}-2 c_{1}\right) \mathrm{e}^{-3 x}}{6 c_{3}^{3}}
\]

Verified OK.
\[
y=-\frac{\left(2 c_{1} c_{5}^{6} \mathrm{e}^{6 x}-1\right) \mathrm{e}^{-3 x}}{6 c_{5}^{3}}
\]

Verified OK.

\subsection*{1.84.3 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}-9 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=0  \tag{3}\\
& C=-9
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{9}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=9 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=9 z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 132: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=9\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-3 x}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\mathrm{e}^{-3 x}
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{-3 x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\mathrm{e}^{-3 x} \int \frac{1}{\mathrm{e}^{-6 x}} d x \\
& =\mathrm{e}^{-3 x}\left(\frac{\mathrm{e}^{6 x}}{6}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{-3 x}\right)+c_{2}\left(\mathrm{e}^{-3 x}\left(\frac{\mathrm{e}^{6 x}}{6}\right)\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{-3 x}+\frac{\mathrm{e}^{3 x} c_{2}}{6} \tag{1}
\end{equation*}
\]


Figure 144: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \mathrm{e}^{-3 x}+\frac{\mathrm{e}^{3 x} c_{2}}{6}
\]

Verified OK.

\subsection*{1.84.4 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}-9 y=0
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of ODE
\[
r^{2}-9=0
\]
- Factor the characteristic polynomial
\[
(r-3)(r+3)=0
\]
- Roots of the characteristic polynomial
\[
r=(-3,3)
\]
- \(\quad 1\) st solution of the ODE
\[
y_{1}(x)=\mathrm{e}^{-3 x}
\]
- \(\quad 2 n d\) solution of the ODE
\[
y_{2}(x)=\mathrm{e}^{3 x}
\]
- General solution of the ODE
\[
y=c_{1} y_{1}(x)+c_{2} y_{2}(x)
\]
- Substitute in solutions
\[
y=c_{1} \mathrm{e}^{-3 x}+\mathrm{e}^{3 x} c_{2}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 17
```

dsolve(diff(y(x),x\$2)=9*y(x),y(x), singsol=all)

```
\[
y(x)=c_{1} \mathrm{e}^{3 x}+\mathrm{e}^{-3 x} c_{2}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.023 (sec). Leaf size: 22
```

DSolve[y''[x]==9*y[x],y[x],x,IncludeSingularSolutions -> True]

```
\[
y(x) \rightarrow e^{-3 x}\left(c_{1} e^{6 x}+c_{2}\right)
\]

\subsection*{1.85 problem 130}
1.85.1 Solving as second order linear constant coeff ode . . . . . . . . 911
1.85.2 Solving as second order ode can be made integrable ode . . . . 913
1.85.3 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 915
1.85.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 919

Internal problem ID [12502]
Internal file name [OUTPUT/11154_Monday_October_16_2023_09_54_09_PM_70184344/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 130.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second__order_linear_constant_coeff", "second__order_ode_can_bbe_made_integrable"

Maple gives the following as the ode type
```

[[_2nd_order, _missing_x]]

```
\[
y^{\prime \prime}+y=0
\]

\subsection*{1.85.1 Solving as second order linear constant coeff ode}

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=1\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+\mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\mathrm{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+1=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=1\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)(1)} \\
& = \pm i
\end{aligned}
\]

Hence
\[
\begin{gathered}
\lambda_{1}=+i \\
\lambda_{2}=-i
\end{gathered}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=i \\
& \lambda_{2}=-i
\end{aligned}
\]

Since roots are complex conjugate of each others, then let the roots be
\[
\lambda_{1,2}=\alpha \pm i \beta
\]

Where \(\alpha=0\) and \(\beta=1\). Therefore the final solution, when using Euler relation, can be written as
\[
y=e^{\alpha x}\left(c_{1} \cos (\beta x)+c_{2} \sin (\beta x)\right)
\]

Which becomes
\[
y=e^{0}\left(c_{1} \cos (x)+c_{2} \sin (x)\right)
\]

Or
\[
y=c_{1} \cos (x)+c_{2} \sin (x)
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (x)+c_{2} \sin (x) \tag{1}
\end{equation*}
\]


Figure 145: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \cos (x)+c_{2} \sin (x)
\]

Verified OK.

\subsection*{1.85.2 Solving as second order ode can be made integrable ode}

Multiplying the ode by \(y^{\prime}\) gives
\[
y^{\prime} y^{\prime \prime}+y y^{\prime}=0
\]

Integrating the above w.r.t \(x\) gives
\[
\begin{gathered}
\int\left(y^{\prime} y^{\prime \prime}+y y^{\prime}\right) d x=0 \\
\frac{y^{\prime 2}}{2}+\frac{y^{2}}{2}=c_{2}
\end{gathered}
\]

Which is now solved for \(y\). Solving the given ode for \(y^{\prime}\) results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are
\[
\begin{align*}
& y^{\prime}=\sqrt{-y^{2}+2 c_{1}}  \tag{1}\\
& y^{\prime}=-\sqrt{-y^{2}+2 c_{1}} \tag{2}
\end{align*}
\]

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives
\[
\begin{aligned}
\int \frac{1}{\sqrt{-y^{2}+2 c_{1}}} d y & =\int d x \\
\arctan \left(\frac{y}{\sqrt{-y^{2}+2 c_{1}}}\right) & =x+c_{2}
\end{aligned}
\]

Solving equation (2)
Integrating both sides gives
\[
\begin{aligned}
\int-\frac{1}{\sqrt{-y^{2}+2 c_{1}}} d y & =\int d x \\
-\arctan \left(\frac{y}{\sqrt{-y^{2}+2 c_{1}}}\right) & =c_{3}+x
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{align*}
\arctan \left(\frac{y}{\sqrt{-y^{2}+2 c_{1}}}\right) & =x+c_{2}  \tag{1}\\
-\arctan \left(\frac{y}{\sqrt{-y^{2}+2 c_{1}}}\right) & =c_{3}+x \tag{2}
\end{align*}
\]


Figure 146: Slope field plot

Verification of solutions
\[
\arctan \left(\frac{y}{\sqrt{-y^{2}+2 c_{1}}}\right)=x+c_{2}
\]

Verified OK.
\[
-\arctan \left(\frac{y}{\sqrt{-y^{2}+2 c_{1}}}\right)=c_{3}+x
\]

Verified OK.

\subsection*{1.85.3 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
A & =1 \\
B & =0  \tag{3}\\
C & =1
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{-1}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=-1 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=-z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi-
\end{tabular} & no condition \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\). & \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 134: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=-1\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\cos (x)
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\cos (x)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\cos (x)
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\cos (x) \int \frac{1}{\cos (x)^{2}} d x \\
& =\cos (x)(\tan (x))
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}(\cos (x))+c_{2}(\cos (x)(\tan (x)))
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (x)+c_{2} \sin (x) \tag{1}
\end{equation*}
\]


Figure 147: Slope field plot

Verification of solutions
\[
y=c_{1} \cos (x)+c_{2} \sin (x)
\]

Verified OK.

\subsection*{1.85.4 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}+y=0
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of ODE
\[
r^{2}+1=0
\]
- Use quadratic formula to solve for \(r\)
\[
r=\frac{0 \pm(\sqrt{-4})}{2}
\]
- Roots of the characteristic polynomial
\[
r=(-\mathrm{I}, \mathrm{I})
\]
- 1st solution of the ODE
\[
y_{1}(x)=\cos (x)
\]
- \(\quad 2 n d\) solution of the ODE
\[
y_{2}(x)=\sin (x)
\]
- General solution of the ODE
\[
y=c_{1} y_{1}(x)+c_{2} y_{2}(x)
\]
- Substitute in solutions
\[
y=c_{1} \cos (x)+c_{2} \sin (x)
\]

Maple trace
```

Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 13
```

dsolve(diff(y(x),x\$2)+y(x)=0,y(x), singsol=all)

```
\[
y(x)=c_{1} \sin (x)+c_{2} \cos (x)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.017 (sec). Leaf size: 16
```

DSolve[y''[x]+y[x]==0,y[x],x,IncludeSingularSolutions -> True]

```
\[
y(x) \rightarrow c_{1} \cos (x)+c_{2} \sin (x)
\]

\subsection*{1.86 problem 131}
1.86.1 Solving as second order linear constant coeff ode . . . . . . . . 921
1.86.2 Solving as second order ode can be made integrable ode . . . . 923
1.86.3 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 925
1.86.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 929

Internal problem ID [12503]
Internal file name [OUTPUT/11155_Monday_October_16_2023_09_54_10_PM_6500354/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 131.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second__order_linear_constant_coeff", "second__order_ode_can__be_made_integrable"

Maple gives the following as the ode type
```

[[_2nd_order, _missing_x]]

```
\[
y^{\prime \prime}-y=0
\]

\subsection*{1.86.1 Solving as second order linear constant coeff ode}

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=-1\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-\mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\mathrm{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}-1=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=-1\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)(-1)} \\
& = \pm 1
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=+1 \\
& \lambda_{2}=-1
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=1 \\
& \lambda_{2}=-1
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{(1) x}+c_{2} e^{(-1) x}
\end{aligned}
\]

Or
\[
y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-x}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-x} \tag{1}
\end{equation*}
\]


Figure 148: Slope field plot

\section*{Verification of solutions}
\[
y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-x}
\]

Verified OK.

\subsection*{1.86.2 Solving as second order ode can be made integrable ode}

Multiplying the ode by \(y^{\prime}\) gives
\[
y^{\prime} y^{\prime \prime}-y y^{\prime}=0
\]

Integrating the above w.r.t \(x\) gives
\[
\begin{gathered}
\int\left(y^{\prime} y^{\prime \prime}-y y^{\prime}\right) d x=0 \\
\frac{y^{\prime 2}}{2}-\frac{y^{2}}{2}=c_{2}
\end{gathered}
\]

Which is now solved for \(y\). Solving the given ode for \(y^{\prime}\) results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are
\[
\begin{align*}
& y^{\prime}=\sqrt{y^{2}+2 c_{1}}  \tag{1}\\
& y^{\prime}=-\sqrt{y^{2}+2 c_{1}} \tag{2}
\end{align*}
\]

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives
\[
\begin{aligned}
\int \frac{1}{\sqrt{y^{2}+2 c_{1}}} d y & =\int d x \\
\ln \left(y+\sqrt{y^{2}+2 c_{1}}\right) & =x+c_{2}
\end{aligned}
\]

Raising both side to exponential gives
\[
y+\sqrt{y^{2}+2 c_{1}}=\mathrm{e}^{x+c_{2}}
\]

Which simplifies to
\[
y+\sqrt{y^{2}+2 c_{1}}=c_{3} \mathrm{e}^{x}
\]

Solving equation (2)
Integrating both sides gives
\[
\begin{aligned}
\int-\frac{1}{\sqrt{y^{2}+2 c_{1}}} d y & =\int d x \\
-\ln \left(y+\sqrt{y^{2}+2 c_{1}}\right) & =x+c_{4}
\end{aligned}
\]

Raising both side to exponential gives
\[
\frac{1}{y+\sqrt{y^{2}+2 c_{1}}}=\mathrm{e}^{x+c_{4}}
\]

Which simplifies to
\[
\frac{1}{y+\sqrt{y^{2}+2 c_{1}}}=c_{5} \mathrm{e}^{x}
\]

Summary
The solution(s) found are the following
\[
\begin{align*}
& y=\frac{\left(\mathrm{e}^{2 x} c_{3}^{2}-2 c_{1}\right) \mathrm{e}^{-x}}{2 c_{3}}  \tag{1}\\
& y=-\frac{\left(2 c_{1} c_{5}^{2} \mathrm{e}^{2 x}-1\right) \mathrm{e}^{-x}}{2 c_{5}} \tag{2}
\end{align*}
\]


Figure 149: Slope field plot

\section*{Verification of solutions}
\[
y=\frac{\left(\mathrm{e}^{2 x} c_{3}^{2}-2 c_{1}\right) \mathrm{e}^{-x}}{2 c_{3}}
\]

Verified OK.
\[
y=-\frac{\left(2 c_{1} c_{5}^{2} \mathrm{e}^{2 x}-1\right) \mathrm{e}^{-x}}{2 c_{5}}
\]

Verified OK.

\subsection*{1.86.3 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{array}{r}
y^{\prime \prime}-y=0 \\
A y^{\prime \prime}+B y^{\prime}+C y=0 \tag{2}
\end{array}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=0  \tag{3}\\
& C=-1
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{1}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=1 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 136: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=1\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-x}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\mathrm{e}^{-x}
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{-x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\mathrm{e}^{-x} \int \frac{1}{\mathrm{e}^{-2 x}} d x \\
& =\mathrm{e}^{-x}\left(\frac{\mathrm{e}^{2 x}}{2}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{-x}\right)+c_{2}\left(\mathrm{e}^{-x}\left(\frac{\mathrm{e}^{2 x}}{2}\right)\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{c_{2} \mathrm{e}^{x}}{2}+c_{1} \mathrm{e}^{-x} \tag{1}
\end{equation*}
\]


Figure 150: Slope field plot

Verification of solutions
\[
y=\frac{c_{2} \mathrm{e}^{x}}{2}+c_{1} \mathrm{e}^{-x}
\]

Verified OK.

\subsection*{1.86.4 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}-y=0
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of ODE
\[
r^{2}-1=0
\]
- Factor the characteristic polynomial
\[
(r-1)(r+1)=0
\]
- Roots of the characteristic polynomial
\[
r=(-1,1)
\]
- 1st solution of the ODE
\[
y_{1}(x)=\mathrm{e}^{-x}
\]
- 2 nd solution of the ODE
\[
y_{2}(x)=\mathrm{e}^{x}
\]
- General solution of the ODE
\[
y=c_{1} y_{1}(x)+c_{2} y_{2}(x)
\]
- Substitute in solutions
\[
y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 15
```

dsolve(diff(y(x),x\$2)-y(x)=0,y(x), singsol=all)

```
\[
y(x)=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.02 (sec). Leaf size: 20
```

DSolve[y''[x]-y[x]==0,y[x],x,IncludeSingularSolutions -> True]

```
\[
y(x) \rightarrow c_{1} e^{x}+c_{2} e^{-x}
\]

\subsection*{1.87 problem 132}
1.87.1 Solving as second order linear constant coeff ode . . . . . . . . 931
1.87.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 933
1.87.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 937

Internal problem ID [12504]
Internal file name [OUTPUT/11156_Monday_October_16_2023_09_54_13_PM_8617977/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 132.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _missing_x]]
\[
y^{\prime \prime}-7 y^{\prime}+12 y=0
\]

\subsection*{1.87.1 Solving as second order linear constant coeff ode}

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=-7, C=12\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-7 \lambda \mathrm{e}^{\lambda x}+12 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}-7 \lambda+12=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=-7, C=12\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{7}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-7^{2}-(4)(1)(12)} \\
& =\frac{7}{2} \pm \frac{1}{2}
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=\frac{7}{2}+\frac{1}{2} \\
& \lambda_{2}=\frac{7}{2}-\frac{1}{2}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=4 \\
& \lambda_{2}=3
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{(4) x}+c_{2} e^{(3) x}
\end{aligned}
\]

Or
\[
y=c_{1} \mathrm{e}^{4 x}+\mathrm{e}^{3 x} c_{2}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{4 x}+\mathrm{e}^{3 x} c_{2} \tag{1}
\end{equation*}
\]


Figure 151: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \mathrm{e}^{4 x}+\mathrm{e}^{3 x} c_{2}
\]

Verified OK.

\subsection*{1.87.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}-7 y^{\prime}+12 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=-7  \tag{3}\\
& C=12
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{1}{4} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=1 \\
& t=4
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=\frac{z(x)}{4} \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi-
\end{tabular} & no condition \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\). & \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 138: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=\frac{1}{4}\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-\frac{x}{2}}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{7}{1} d x} \\
& =z_{1} e^{\frac{7 x}{2}} \\
& =z_{1}\left(\mathrm{e}^{\frac{7 x}{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{3 x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-7}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{7 x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\mathrm{e}^{x}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{3 x}\right)+c_{2}\left(\mathrm{e}^{3 x}\left(\mathrm{e}^{x}\right)\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{4 x} \tag{1}
\end{equation*}
\]


Figure 152: Slope field plot

Verification of solutions
\[
y=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{4 x}
\]

Verified OK.

\subsection*{1.87.3 Maple step by step solution}

Let's solve
\(y^{\prime \prime}-7 y^{\prime}+12 y=0\)
- Highest derivative means the order of the ODE is 2
\(y^{\prime \prime}\)
- Characteristic polynomial of ODE
\(r^{2}-7 r+12=0\)
- Factor the characteristic polynomial
\((r-3)(r-4)=0\)
- Roots of the characteristic polynomial
\(r=(3,4)\)
- \(\quad 1\) st solution of the ODE
\(y_{1}(x)=\mathrm{e}^{3 x}\)
- \(\quad 2 n d\) solution of the ODE
\(y_{2}(x)=\mathrm{e}^{4 x}\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)\)
- Substitute in solutions
\(y=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{4 x}\)

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 17
dsolve(diff \((y(x), x \$ 2)+12 * y(x)=7 * \operatorname{diff}(y(x), x), y(x)\), singsol=all)
\[
y(x)=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{4 x}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.023 (sec). Leaf size: 20
DSolve[y''[x]+12*y[x]==7*y'[x],y[x],x,IncludeSingularSolutions -> True]
\[
y(x) \rightarrow e^{3 x}\left(c_{2} e^{x}+c_{1}\right)
\]

\subsection*{1.88 problem 133}
\[
\text { 1.88.1 Solving as second order linear constant coeff ode . . . . . . . . } 939
\]
1.88.2 Solving as linear second order ode solved by an integrating factor ode ..... 941
1.88.3 Solving using Kovacic algorithm ..... 942
1.88.4 Maple step by step solution ..... 946

Internal problem ID [12505]
Internal file name [OUTPUT/11157_Monday_October_16_2023_09_54_15_PM_411349/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 133.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear__second_order_ode_solved_by__an_integrating_factor"

Maple gives the following as the ode type
[[_2nd_order, _missing_x]]
\[
y^{\prime \prime}-4 y^{\prime}+4 y=0
\]

\subsection*{1.88.1 Solving as second order linear constant coeff ode}

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=-4, C=4\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-4 \lambda \mathrm{e}^{\lambda x}+4 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}-4 \lambda+4=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=-4, C=4\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{4}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(-4)^{2}-(4)(1)(4)} \\
& =2
\end{aligned}
\]

Hence this is the case of a double root \(\lambda_{1,2}=-2\). Therefore the solution is
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{2 x} x \tag{1}
\end{equation*}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{2 x} x \tag{1}
\end{equation*}
\]


Figure 153: Slope field plot
Verification of solutions
\[
y=c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{2 x} x
\]

Verified OK.

\subsection*{1.88.2 Solving as linear second order ode solved by an integrating factor ode}

The ode satisfies this form
\[
y^{\prime \prime}+p(x) y^{\prime}+\frac{\left(p(x)^{2}+p^{\prime}(x)\right) y}{2}=f(x)
\]

Where \(p(x)=-4\). Therefore, there is an integrating factor given by
\[
\begin{aligned}
M(x) & =e^{\frac{1}{2} \int p d x} \\
& =e^{\int-4 d x} \\
& =\mathrm{e}^{-2 x}
\end{aligned}
\]

Multiplying both sides of the ODE by the integrating factor \(M(x)\) makes the left side of the ODE a complete differential
\[
\begin{aligned}
(M(x) y)^{\prime \prime} & =0 \\
\left(\mathrm{e}^{-2 x} y\right)^{\prime \prime} & =0
\end{aligned}
\]

Integrating once gives
\[
\left(\mathrm{e}^{-2 x} y\right)^{\prime}=c_{1}
\]

Integrating again gives
\[
\left(\mathrm{e}^{-2 x} y\right)=c_{1} x+c_{2}
\]

Hence the solution is
\[
y=\frac{c_{1} x+c_{2}}{\mathrm{e}^{-2 x}}
\]

Or
\[
y=\mathrm{e}^{2 x} c_{1} x+c_{2} \mathrm{e}^{2 x}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{2 x} c_{1} x+c_{2} \mathrm{e}^{2 x} \tag{1}
\end{equation*}
\]


Figure 154: Slope field plot

\section*{Verification of solutions}
\[
y=\mathrm{e}^{2 x} c_{1} x+c_{2} \mathrm{e}^{2 x}
\]

Verified OK.

\subsection*{1.88.3 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}-4 y^{\prime}+4 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=-4  \tag{3}\\
& C=4
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{0}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=0 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=0 \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 140: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0--\infty \\
& =\infty
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is infinity then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=0\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=1
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{-4}{1} d x} \\
& =z_{1} e^{2 x} \\
& =z_{1}\left(\mathrm{e}^{2 x}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{2 x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-4}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{4 x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}(x)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{2 x}\right)+c_{2}\left(\mathrm{e}^{2 x}(x)\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{2 x} x \tag{1}
\end{equation*}
\]


Figure 155: Slope field plot

Verification of solutions
\[
y=c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{2 x} x
\]

Verified OK.

\subsection*{1.88.4 Maple step by step solution}

Let's solve
\(y^{\prime \prime}-4 y^{\prime}+4 y=0\)
- Highest derivative means the order of the ODE is 2
\(y^{\prime \prime}\)
- Characteristic polynomial of ODE
\(r^{2}-4 r+4=0\)
- Factor the characteristic polynomial
\((r-2)^{2}=0\)
- Root of the characteristic polynomial
\[
r=2
\]
- \(\quad 1\) st solution of the ODE
\[
y_{1}(x)=\mathrm{e}^{2 x}
\]
- Repeated root, multiply \(y_{1}(x)\) by \(x\) to ensure linear independence \(y_{2}(x)=\mathrm{e}^{2 x} x\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)\)
- Substitute in solutions
\[
y=c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{2 x} x
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 14
dsolve(diff \((y(x), x \$ 2)-4 * \operatorname{diff}(y(x), x)+4 * y(x)=0, y(x)\), singsol=all)
\[
y(x)=\mathrm{e}^{2 x}\left(c_{2} x+c_{1}\right)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.025 (sec). Leaf size: 18
DSolve[y'' \([x]-4 * y\) ' \([x]+4 * y[x]==0, y[x], x\), IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow e^{2 x}\left(c_{2} x+c_{1}\right)
\]

\subsection*{1.89 problem 134}

\subsection*{1.89.1 Solving as second order linear constant coeff ode 948}
1.89.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 950
1.89.3 Maple step by step solution 954

Internal problem ID [12506]
Internal file name [OUTPUT/11158_Monday_October_16_2023_09_54_16_PM_15038119/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 134.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _missing_x]]
\[
y^{\prime \prime}+2 y^{\prime}+10 y=0
\]

\subsection*{1.89.1 Solving as second order linear constant coeff ode}

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=2, C=10\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+2 \lambda \mathrm{e}^{\lambda x}+10 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\mathrm{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+2 \lambda+10=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=2, C=10\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{-2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{2^{2}-(4)(1)(10)} \\
& =-1 \pm 3 i
\end{aligned}
\]

Hence
\[
\begin{gathered}
\lambda_{1}=-1+3 i \\
\lambda_{2}=-1-3 i
\end{gathered}
\]

Which simplifies to
\[
\begin{gathered}
\lambda_{1}=-1+3 i \\
\lambda_{2}=-1-3 i
\end{gathered}
\]

Since roots are complex conjugate of each others, then let the roots be
\[
\lambda_{1,2}=\alpha \pm i \beta
\]

Where \(\alpha=-1\) and \(\beta=3\). Therefore the final solution, when using Euler relation, can be written as
\[
y=e^{\alpha x}\left(c_{1} \cos (\beta x)+c_{2} \sin (\beta x)\right)
\]

Which becomes
\[
y=e^{-x}\left(c_{1} \cos (3 x)+c_{2} \sin (3 x)\right)
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{-x}\left(c_{1} \cos (3 x)+c_{2} \sin (3 x)\right) \tag{1}
\end{equation*}
\]


Figure 156: Slope field plot

\section*{Verification of solutions}
\[
y=\mathrm{e}^{-x}\left(c_{1} \cos (3 x)+c_{2} \sin (3 x)\right)
\]

Verified OK.

\subsection*{1.89.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{array}{r}
y^{\prime \prime}+2 y^{\prime}+10 y=0 \\
A y^{\prime \prime}+B y^{\prime}+C y=0 \tag{2}
\end{array}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=2  \tag{3}\\
& C=10
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{-9}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
s & =-9 \\
t & =1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=-9 z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 142: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=-9\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\cos (3 x)
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{2}{1} d x} \\
& =z_{1} e^{-x} \\
& =z_{1}\left(\mathrm{e}^{-x}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{-x} \cos (3 x)
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{2}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{-2 x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\frac{\tan (3 x)}{3}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{-x} \cos (3 x)\right)+c_{2}\left(\mathrm{e}^{-x} \cos (3 x)\left(\frac{\tan (3 x)}{3}\right)\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{-x} \cos (3 x) c_{1}+\frac{\mathrm{e}^{-x} \sin (3 x) c_{2}}{3} \tag{1}
\end{equation*}
\]


Figure 157: Slope field plot

Verification of solutions
\[
y=\mathrm{e}^{-x} \cos (3 x) c_{1}+\frac{\mathrm{e}^{-x} \sin (3 x) c_{2}}{3}
\]

Verified OK.

\subsection*{1.89.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}+2 y^{\prime}+10 y=0
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of ODE
\[
r^{2}+2 r+10=0
\]
- Use quadratic formula to solve for \(r\)
\(r=\frac{(-2) \pm(\sqrt{-36})}{2}\)
- Roots of the characteristic polynomial
\[
r=(-1-3 \mathrm{I},-1+3 \mathrm{I})
\]
- \(\quad 1\) st solution of the ODE
\(y_{1}(x)=\mathrm{e}^{-x} \cos (3 x)\)
- \(\quad 2 \mathrm{nd}\) solution of the ODE
\(y_{2}(x)=\mathrm{e}^{-x} \sin (3 x)\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)\)
- Substitute in solutions
\[
y=\mathrm{e}^{-x} \cos (3 x) c_{1}+\mathrm{e}^{-x} \sin (3 x) c_{2}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 22
dsolve( \(\operatorname{diff}(y(x), x \$ 2)+2 * \operatorname{diff}(y(x), x)+10 * y(x)=0, y(x)\), singsol=all)
\[
y(x)=\mathrm{e}^{-x}\left(c_{1} \sin (3 x)+c_{2} \cos (3 x)\right)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.033 (sec). Leaf size: 26
DSolve[y'' \([x]+2 * y\) ' \([x]+10 * y[x]==0, y[x], x\), IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow e^{-x}\left(c_{2} \cos (3 x)+c_{1} \sin (3 x)\right)
\]

\subsection*{1.90 problem 135}
1.90.1 Solving as second order linear constant coeff ode
1.90.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 958
1.90.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 962

Internal problem ID [12507]
Internal file name [OUTPUT/11159_Monday_October_16_2023_09_54_18_PM_28113913/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 135.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _missing_x]]
\[
y^{\prime \prime}+3 y^{\prime}-2 y=0
\]

\subsection*{1.90.1 Solving as second order linear constant coeff ode}

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=3, C=-2\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+3 \lambda \mathrm{e}^{\lambda x}-2 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+3 \lambda-2=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=3, C=-2\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{-3}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{3^{2}-(4)(1)(-2)} \\
& =-\frac{3}{2} \pm \frac{\sqrt{17}}{2}
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=-\frac{3}{2}+\frac{\sqrt{17}}{2} \\
& \lambda_{2}=-\frac{3}{2}-\frac{\sqrt{17}}{2}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=-\frac{3}{2}+\frac{\sqrt{17}}{2} \\
& \lambda_{2}=-\frac{3}{2}-\frac{\sqrt{17}}{2}
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{\left(-\frac{3}{2}+\frac{\sqrt{17}}{2}\right) x}+c_{2} e^{\left(-\frac{3}{2}-\frac{\sqrt{17}}{2}\right) x}
\end{aligned}
\]

Or
\[
y=c_{1} \mathrm{e}^{\left(-\frac{3}{2}+\frac{\sqrt{17}}{2}\right) x}+c_{2} \mathrm{e}^{\left(-\frac{3}{2}-\frac{\sqrt{17}}{2}\right) x}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{\left(-\frac{3}{2}+\frac{\sqrt{17}}{2}\right) x}+c_{2} \mathrm{e}^{\left(-\frac{3}{2}-\frac{\sqrt{17}}{2}\right) x} \tag{1}
\end{equation*}
\]


Figure 158: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \mathrm{e}^{\left(-\frac{3}{2}+\frac{\sqrt{17}}{2}\right) x}+c_{2} \mathrm{e}^{\left(-\frac{3}{2}-\frac{\sqrt{17}}{2}\right) x}
\]

Verified OK.

\subsection*{1.90.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+3 y^{\prime}-2 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=3  \tag{3}\\
& C=-2
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{17}{4} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=17 \\
& t=4
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=\frac{17 z(x)}{4} \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi-
\end{tabular} & no condition \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\). & \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 144: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=\frac{17}{4}\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-\frac{x \sqrt{17}}{2}}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{3}{1} d x} \\
& =z_{1} e^{-\frac{3 x}{2}} \\
& =z_{1}\left(\mathrm{e}^{-\frac{3 x}{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{-\frac{(3+\sqrt{17}) x}{2}}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{3}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{-3 x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\frac{\sqrt{17} \mathrm{e}^{x \sqrt{17}}}{17}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{-\frac{(3+\sqrt{17}) x}{2}}\right)+c_{2}\left(\mathrm{e}^{-\frac{(3+\sqrt{17}) x}{2}}\left(\frac{\sqrt{17} \mathrm{e}^{x \sqrt{17}}}{17}\right)\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{-\frac{(3+\sqrt{17}) x}{2}}+\frac{c_{2} \sqrt{17} \mathrm{e}^{\frac{(-3+\sqrt{17}) x}{2}}}{17} \tag{1}
\end{equation*}
\]


Figure 159: Slope field plot

Verification of solutions
\[
y=c_{1} \mathrm{e}^{-\frac{(3+\sqrt{17}) x}{2}}+\frac{c_{2} \sqrt{17} \mathrm{e}^{\frac{(-3+\sqrt{17}) x}{2}}}{17}
\]

Verified OK.

\subsection*{1.90.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}+3 y^{\prime}-2 y=0
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of ODE
\(r^{2}+3 r-2=0\)
- Use quadratic formula to solve for \(r\)
\(r=\frac{(-3) \pm(\sqrt{17})}{2}\)
- Roots of the characteristic polynomial
\(r=\left(-\frac{3}{2}-\frac{\sqrt{17}}{2},-\frac{3}{2}+\frac{\sqrt{17}}{2}\right)\)
- \(\quad 1\) st solution of the ODE
\(y_{1}(x)=\mathrm{e}^{\left(-\frac{3}{2}-\frac{\sqrt{17}}{2}\right) x}\)
- \(\quad 2\) nd solution of the ODE
\(y_{2}(x)=\mathrm{e}^{\left(-\frac{3}{2}+\frac{\sqrt{17}}{2}\right) x}\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)\)
- Substitute in solutions
\(y=c_{1} \mathrm{e}^{\left(-\frac{3}{2}-\frac{\sqrt{17}}{2}\right) x}+c_{2} \mathrm{e}^{\left(-\frac{3}{2}+\frac{\sqrt{17}}{2}\right) x}\)

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 27
dsolve(diff \((y(x), x \$ 2)+3 * \operatorname{diff}(y(x), x)-2 * y(x)=0, y(x)\), singsol=all)
\[
y(x)=c_{1} \mathrm{e}^{\frac{(-3+\sqrt{17}) x}{2}}+c_{2} \mathrm{e}^{-\frac{(3+\sqrt{17}) x}{2}}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.041 (sec). Leaf size: 35
DSolve[y'' \([x]+3 * y\) ' \([x]-2 * y[x]==0, y[x], x\), IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow e^{-\frac{1}{2}(3+\sqrt{17}) x}\left(c_{2} e^{\sqrt{17} x}+c_{1}\right)
\]

\subsection*{1.91 problem 136}
\[
\text { 1.91.1 Solving as second order linear constant coeff ode . . . . . . . . } 964
\]
1.91.2 Solving as linear second order ode solved by an integrating factor
ode . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 966
1.91.3 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 967
1.91.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 971

Internal problem ID [12508]
Internal file name [OUTPUT/11160_Monday_October_16_2023_09_54_19_PM_35388541/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 136.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear__second_order_ode_solved_by__an_integrating_factor"

Maple gives the following as the ode type
[[_2nd_order, _missing_x]]
\[
4 y^{\prime \prime}-12 y^{\prime}+9 y=0
\]

\subsection*{1.91.1 Solving as second order linear constant coeff ode}

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=4, B=-12, C=9\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
4 \lambda^{2} \mathrm{e}^{\lambda x}-12 \lambda \mathrm{e}^{\lambda x}+9 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
4 \lambda^{2}-12 \lambda+9=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=4, B=-12, C=9\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{12}{(2)(4)} \pm \frac{1}{(2)(4)} \sqrt{(-12)^{2}-(4)(4)(9)} \\
& =\frac{3}{2}
\end{aligned}
\]

Hence this is the case of a double root \(\lambda_{1,2}=-\frac{3}{2}\). Therefore the solution is
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{\frac{3 x}{2}}+c_{2} x \mathrm{e}^{\frac{3 x}{2}} \tag{1}
\end{equation*}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{\frac{3 x}{2}}+c_{2} x \mathrm{e}^{\frac{3 x}{2}} \tag{1}
\end{equation*}
\]


Figure 160: Slope field plot
Verification of solutions
\[
y=c_{1} \mathrm{e}^{\frac{3 x}{2}}+c_{2} x \mathrm{e}^{\frac{3 x}{2}}
\]

Verified OK.

\subsection*{1.91.2 Solving as linear second order ode solved by an integrating factor ode}

The ode satisfies this form
\[
y^{\prime \prime}+p(x) y^{\prime}+\frac{\left(p(x)^{2}+p^{\prime}(x)\right) y}{2}=f(x)
\]

Where \(p(x)=-3\). Therefore, there is an integrating factor given by
\[
\begin{aligned}
M(x) & =e^{\frac{1}{2} \int p d x} \\
& =e^{\int-3 d x} \\
& =\mathrm{e}^{-\frac{3 x}{2}}
\end{aligned}
\]

Multiplying both sides of the ODE by the integrating factor \(M(x)\) makes the left side of the ODE a complete differential
\[
\begin{aligned}
& (M(x) y)^{\prime \prime}=0 \\
& \left(\mathrm{e}^{-\frac{3 x}{2}} y\right)^{\prime \prime}=0
\end{aligned}
\]

Integrating once gives
\[
\left(\mathrm{e}^{-\frac{3 x}{2}} y\right)^{\prime}=c_{1}
\]

Integrating again gives
\[
\left(\mathrm{e}^{-\frac{3 x}{2}} y\right)=c_{1} x+c_{2}
\]

Hence the solution is
\[
y=\frac{c_{1} x+c_{2}}{\mathrm{e}^{-\frac{3 x}{2}}}
\]

Or
\[
y=c_{1} x \mathrm{e}^{\frac{3 x}{2}}+\mathrm{e}^{\frac{3 x}{2}} c_{2}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} x \mathrm{e}^{\frac{3 x}{2}}+\mathrm{e}^{\frac{3 x}{2}} c_{2} \tag{1}
\end{equation*}
\]


Figure 161: Slope field plot

Verification of solutions
\[
y=c_{1} x \mathrm{e}^{\frac{3 x}{2}}+\mathrm{e}^{\frac{3 x}{2}} c_{2}
\]

Verified OK.

\subsection*{1.91.3 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
4 y^{\prime \prime}-12 y^{\prime}+9 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=4 \\
& B=-12  \tag{3}\\
& C=9
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{0}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=0 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=0 \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 146: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0--\infty \\
& =\infty
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is infinity then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=0\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=1
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{-12}{4} d x} \\
& =z_{1} e^{\frac{3 x}{2}} \\
& =z_{1}\left(\mathrm{e}^{\frac{3 x}{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{\frac{3 x}{2}}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-12}{4} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{3 x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}(x)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{\frac{3 x}{2}}\right)+c_{2}\left(\mathrm{e}^{\frac{3 x}{2}}(x)\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{\frac{3 x}{2}}+c_{2} x \mathrm{e}^{\frac{3 x}{2}} \tag{1}
\end{equation*}
\]


Figure 162: Slope field plot

Verification of solutions
\[
y=c_{1} \mathrm{e}^{\frac{3 x}{2}}+c_{2} x \mathrm{e}^{\frac{3 x}{2}}
\]

\section*{Verified OK.}

\subsection*{1.91.4 Maple step by step solution}

Let's solve
\[
4 y^{\prime \prime}-12 y^{\prime}+9 y=0
\]
- Highest derivative means the order of the ODE is 2
\(y^{\prime \prime}\)
- Isolate 2nd derivative
\(y^{\prime \prime}=3 y^{\prime}-\frac{9 y}{4}\)
- Group terms with \(y\) on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear \(y^{\prime \prime}-3 y^{\prime}+\frac{9 y}{4}=0\)
- Characteristic polynomial of ODE
\[
r^{2}-3 r+\frac{9}{4}=0
\]
- Factor the characteristic polynomial
\(\frac{(2 r-3)^{2}}{4}=0\)
- Root of the characteristic polynomial
\(r=\frac{3}{2}\)
- 1st solution of the ODE
\(y_{1}(x)=\mathrm{e}^{\frac{3 x}{2}}\)
- Repeated root, multiply \(y_{1}(x)\) by \(x\) to ensure linear independence \(y_{2}(x)=x \mathrm{e}^{\frac{3 x}{2}}\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)\)
- Substitute in solutions
\(y=c_{1} \mathrm{e}^{\frac{3 x}{2}}+c_{2} x \mathrm{e}^{\frac{3 x}{2}}\)

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.016 (sec). Leaf size: 14
```

dsolve(4*diff(y(x),x\$2)-12*diff(y(x),x)+9*y(x)=0,y(x), singsol=all)

```
\[
y(x)=\mathrm{e}^{\frac{3 x}{2}}\left(c_{2} x+c_{1}\right)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.024 (sec). Leaf size: 20
DSolve[4*y''[x]-12*y'[x]+9*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
\[
y(x) \rightarrow e^{3 x / 2}\left(c_{2} x+c_{1}\right)
\]

\subsection*{1.92 problem 137}
1.92.1 Solving as second order linear constant coeff ode . . . . . . . . 973
1.92.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 975
1.92.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 979

Internal problem ID [12509]
Internal file name [OUTPUT/11161_Monday_October_16_2023_09_54_21_PM_74855947/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 137.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _missing_x]]
\[
y^{\prime \prime}+y^{\prime}+y=0
\]

\subsection*{1.92.1 Solving as second order linear constant coeff ode}

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=1, C=1\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+\lambda \mathrm{e}^{\lambda x}+\mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\mathrm{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+\lambda+1=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=1, C=1\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{-1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{1^{2}-(4)(1)(1)} \\
& =-\frac{1}{2} \pm \frac{i \sqrt{3}}{2}
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=-\frac{1}{2}+\frac{i \sqrt{3}}{2} \\
& \lambda_{2}=-\frac{1}{2}-\frac{i \sqrt{3}}{2}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=-\frac{1}{2}+\frac{i \sqrt{3}}{2} \\
& \lambda_{2}=-\frac{1}{2}-\frac{i \sqrt{3}}{2}
\end{aligned}
\]

Since roots are complex conjugate of each others, then let the roots be
\[
\lambda_{1,2}=\alpha \pm i \beta
\]

Where \(\alpha=-\frac{1}{2}\) and \(\beta=\frac{\sqrt{3}}{2}\). Therefore the final solution, when using Euler relation, can be written as
\[
y=e^{\alpha x}\left(c_{1} \cos (\beta x)+c_{2} \sin (\beta x)\right)
\]

Which becomes
\[
y=e^{-\frac{x}{2}}\left(c_{1} \cos \left(\frac{\sqrt{3} x}{2}\right)+c_{2} \sin \left(\frac{\sqrt{3} x}{2}\right)\right)
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{-\frac{x}{2}}\left(c_{1} \cos \left(\frac{\sqrt{3} x}{2}\right)+c_{2} \sin \left(\frac{\sqrt{3} x}{2}\right)\right) \tag{1}
\end{equation*}
\]


Figure 163: Slope field plot

Verification of solutions
\[
y=\mathrm{e}^{-\frac{x}{2}}\left(c_{1} \cos \left(\frac{\sqrt{3} x}{2}\right)+c_{2} \sin \left(\frac{\sqrt{3} x}{2}\right)\right)
\]

Verified OK.

\subsection*{1.92.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+y^{\prime}+y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
A & =1 \\
B & =1  \tag{3}\\
C & =1
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{-3}{4} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=-3 \\
& t=4
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=-\frac{3 z(x)}{4} \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi-
\end{tabular} & no condition \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\). & \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 148: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=-\frac{3}{4}\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\cos \left(\frac{\sqrt{3} x}{2}\right)
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{1}{1} d x} \\
& =z_{1} e^{-\frac{x}{2}} \\
& =z_{1}\left(\mathrm{e}^{-\frac{x}{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{-\frac{x}{2}} \cos \left(\frac{\sqrt{3} x}{2}\right)
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{1}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{-x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\frac{2 \sqrt{3} \tan \left(\frac{\sqrt{3} x}{2}\right)}{3}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{-\frac{x}{2}} \cos \left(\frac{\sqrt{3} x}{2}\right)\right)+c_{2}\left(\mathrm{e}^{-\frac{x}{2}} \cos \left(\frac{\sqrt{3} x}{2}\right)\left(\frac{2 \sqrt{3} \tan \left(\frac{\sqrt{3} x}{2}\right)}{3}\right)\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{-\frac{x}{2}} \cos \left(\frac{\sqrt{3} x}{2}\right)+\frac{2 \mathrm{e}^{-\frac{x}{2}} \sqrt{3} \sin \left(\frac{\sqrt{3} x}{2}\right) c_{2}}{3} \tag{1}
\end{equation*}
\]


Figure 164: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \mathrm{e}^{-\frac{x}{2}} \cos \left(\frac{\sqrt{3} x}{2}\right)+\frac{2 \mathrm{e}^{-\frac{x}{2}} \sqrt{3} \sin \left(\frac{\sqrt{3} x}{2}\right) c_{2}}{3}
\]

Verified OK.

\subsection*{1.92.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}+y^{\prime}+y=0
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of ODE
\[
r^{2}+r+1=0
\]
- Use quadratic formula to solve for \(r\)
\[
r=\frac{(-1) \pm(\sqrt{-3})}{2}
\]
- Roots of the characteristic polynomial
\[
r=\left(-\frac{1}{2}-\frac{\mathrm{I} \sqrt{3}}{2},-\frac{1}{2}+\frac{\mathrm{I} \sqrt{3}}{2}\right)
\]
- 1st solution of the ODE
\[
y_{1}(x)=\mathrm{e}^{-\frac{x}{2}} \cos \left(\frac{\sqrt{3} x}{2}\right)
\]
- 2 nd solution of the ODE
\(y_{2}(x)=\mathrm{e}^{-\frac{x}{2}} \sin \left(\frac{\sqrt{3} x}{2}\right)\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)\)
- \(\quad\) Substitute in solutions
\[
y=c_{1} \mathrm{e}^{-\frac{x}{2}} \cos \left(\frac{\sqrt{3} x}{2}\right)+c_{2} \mathrm{e}^{-\frac{x}{2}} \sin \left(\frac{\sqrt{3} x}{2}\right)
\]

\section*{Maple trace}
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.015 (sec). Leaf size: 28
```

dsolve(diff(y(x),x\$2)+diff(y(x),x)+y(x)=0,y(x), singsol=all)

```
\[
y(x)=\mathrm{e}^{-\frac{x}{2}}\left(c_{1} \sin \left(\frac{\sqrt{3} x}{2}\right)+c_{2} \cos \left(\frac{\sqrt{3} x}{2}\right)\right)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.041 (sec). Leaf size: 42
DSolve[y'' \([x]+y\) ' \([x]+y[x]==0, y[x], x\), IncludeSingularSolutions \(\rightarrow\) True]
\[
y(x) \rightarrow e^{-x / 2}\left(c_{2} \cos \left(\frac{\sqrt{3} x}{2}\right)+c_{1} \sin \left(\frac{\sqrt{3} x}{2}\right)\right)
\]

\subsection*{1.93 problem 140}
1.93.1 Maple step by step solution

983
Internal problem ID [12510]
Internal file name [OUTPUT/11162_Monday_October_16_2023_09_54_22_PM_55984425/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 140.
ODE order: 4.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_linear_constant__coefficients_ODE"

Maple gives the following as the ode type
[[_high_order, _missing_x]]
\[
y^{\prime \prime \prime \prime}-5 y^{\prime \prime}+4 y=0
\]

The characteristic equation is
\[
\lambda^{4}-5 \lambda^{2}+4=0
\]

The roots of the above equation are
\[
\begin{aligned}
\lambda_{1} & =2 \\
\lambda_{2} & =-2 \\
\lambda_{3} & =1 \\
\lambda_{4} & =-1
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{-2 x}+c_{3} \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{4}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=\mathrm{e}^{-x} \\
& y_{2}=\mathrm{e}^{-2 x} \\
& y_{3}=\mathrm{e}^{x} \\
& y_{4}=\mathrm{e}^{2 x}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{-2 x}+c_{3} \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{4} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{-2 x}+c_{3} \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{4}
\]

Verified OK.

\subsection*{1.93.1 Maple step by step solution}

Let's solve
\[
y^{\prime \prime \prime \prime}-5 y^{\prime \prime}+4 y=0
\]
- Highest derivative means the order of the ODE is 4 \(y^{\prime \prime \prime \prime}\)
Convert linear ODE into a system of first order ODEs
- Define new variable \(y_{1}(x)\)
\(y_{1}(x)=y\)
- Define new variable \(y_{2}(x)\)
\(y_{2}(x)=y^{\prime}\)
- Define new variable \(y_{3}(x)\)
\[
y_{3}(x)=y^{\prime \prime}
\]
- Define new variable \(y_{4}(x)\)
\(y_{4}(x)=y^{\prime \prime \prime}\)
- Isolate for \(y_{4}^{\prime}(x)\) using original ODE
\(y_{4}^{\prime}(x)=5 y_{3}(x)-4 y_{1}(x)\)
Convert linear ODE into a system of first order ODEs
\[
\left[y_{2}(x)=y_{1}^{\prime}(x), y_{3}(x)=y_{2}^{\prime}(x), y_{4}(x)=y_{3}^{\prime}(x), y_{4}^{\prime}(x)=5 y_{3}(x)-4 y_{1}(x)\right]
\]
- Define vector
\[
\vec{y}(x)=\left[\begin{array}{l}
y_{1}(x) \\
y_{2}(x) \\
y_{3}(x) \\
y_{4}(x)
\end{array}\right]
\]
- System to solve
\[
\vec{y}^{\prime}(x)=\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
-4 & 0 & 5 & 0
\end{array}\right] \cdot \vec{y}(x)
\]
- Define the coefficient matrix
\[
A=\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
-4 & 0 & 5 & 0
\end{array}\right]
\]
- Rewrite the system as
\(\vec{y}^{\prime}(x)=A \cdot \vec{y}(x)\)
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- \(\quad\) Eigenpairs of \(A\)
\(\left[\left[\left[-2,\left[\begin{array}{c}-\frac{1}{8} \\ \frac{1}{4} \\ -\frac{1}{2} \\ 1\end{array}\right]\right],\left[-1,\left[\begin{array}{c}-1 \\ 1 \\ -1 \\ 1\end{array}\right]\right],\left[1,\left[\begin{array}{l}1 \\ 1 \\ 1 \\ 1\end{array}\right]\right],\left[2,\left[\begin{array}{c}\frac{1}{8} \\ \frac{1}{4} \\ \frac{1}{2} \\ 1\end{array}\right]\right]\right]\right.\)
- Consider eigenpair
\[
\left[-2,\left[\begin{array}{c}
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]\right]
\]
- Solution to homogeneous system from eigenpair
\[
\vec{y}_{1}=\mathrm{e}^{-2 x} \cdot\left[\begin{array}{c}
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]
\]
- Consider eigenpair
\[
\left[-1,\left[\begin{array}{c}
-1 \\
1 \\
-1 \\
1
\end{array}\right]\right]
\]
- Solution to homogeneous system from eigenpair
\[
\vec{y}_{2}=\mathrm{e}^{-x} \cdot\left[\begin{array}{c}
-1 \\
1 \\
-1 \\
1
\end{array}\right]
\]
- Consider eigenpair
\(\left[1,\left[\begin{array}{l}1 \\ 1 \\ 1 \\ 1\end{array}\right]\right]\)
- \(\quad\) Solution to homogeneous system from eigenpair
\[
\vec{y}_{3}=\mathrm{e}^{x} \cdot\left[\begin{array}{l}
1 \\
1 \\
1 \\
1
\end{array}\right]
\]
- Consider eigenpair
\[
\left[2,\left[\begin{array}{c}
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]\right]
\]
- \(\quad\) Solution to homogeneous system from eigenpair
\[
\vec{y}_{4}=\mathrm{e}^{2 x} \cdot\left[\begin{array}{c}
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]
\]
- General solution to the system of ODEs
\[
\vec{y}=c_{1} \vec{y}_{1}+c_{2} \vec{y}_{2}+c_{3} \vec{y}_{3}+c_{4} \vec{y}_{4}
\]
- \(\quad\) Substitute solutions into the general solution
\[
\vec{y}=c_{1} \mathrm{e}^{-2 x} \cdot\left[\begin{array}{c}
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]+c_{2} \mathrm{e}^{-x} \cdot\left[\begin{array}{c}
-1 \\
1 \\
-1 \\
1
\end{array}\right]+c_{3} \mathrm{e}^{x} \cdot\left[\begin{array}{c}
1 \\
1 \\
1 \\
1
\end{array}\right]+\mathrm{e}^{2 x} c_{4} \cdot\left[\begin{array}{c}
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]
\]
- First component of the vector is the solution to the ODE
\(y=-\frac{\left(-\mathrm{e}^{4 x} c_{4}-8 c_{3} \mathrm{e}^{3 x}+8 c_{2} \mathrm{e}^{x}+c_{1}\right) \mathrm{e}^{-2 x}}{8}\)

Maple trace
```

`Methods for high order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 27
dsolve(diff \((y(x), x \$ 4)-5 * \operatorname{diff}(y(x), x \$ 2)+4 * y(x)=0, y(x)\), singsol=all)
\[
y(x)=\left(c_{3} \mathrm{e}^{4 x}+c_{2} \mathrm{e}^{3 x}+c_{1} \mathrm{e}^{x}+c_{4}\right) \mathrm{e}^{-2 x}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.005 (sec). Leaf size: 35
DSolve[y'' '' \([\mathrm{x}]-5 * \mathrm{y}\) '' \([\mathrm{x}]+4 * \mathrm{y}[\mathrm{x}]==0, \mathrm{y}[\mathrm{x}], \mathrm{x}\), IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow e^{-2 x}\left(c_{2} e^{x}+e^{3 x}\left(c_{4} e^{x}+c_{3}\right)+c_{1}\right)
\]

\subsection*{1.94 problem 141}
1.94.1 Maple step by step solution

989
Internal problem ID [12511]
Internal file name [OUTPUT/11163_Monday_October_16_2023_09_54_22_PM_74100822/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 141.
ODE order: 3.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_linear_constant_coefficients_ODE"

Maple gives the following as the ode type
[[_3rd_order, _missing_x]]
\[
y^{\prime \prime \prime}-2 y^{\prime \prime}-y^{\prime}+2 y=0
\]

The characteristic equation is
\[
\lambda^{3}-2 \lambda^{2}-\lambda+2=0
\]

The roots of the above equation are
\[
\begin{aligned}
& \lambda_{1}=1 \\
& \lambda_{2}=2 \\
& \lambda_{3}=-1
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{3}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=\mathrm{e}^{-x} \\
& y_{2}=\mathrm{e}^{x} \\
& y_{3}=\mathrm{e}^{2 x}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{3} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{3}
\]

Verified OK.

\subsection*{1.94.1 Maple step by step solution}

Let's solve
\[
y^{\prime \prime \prime}-2 y^{\prime \prime}-y^{\prime}+2 y=0
\]
- Highest derivative means the order of the ODE is 3
\[
y^{\prime \prime \prime}
\]

Convert linear ODE into a system of first order ODEs
- Define new variable \(y_{1}(x)\)
\[
y_{1}(x)=y
\]
- Define new variable \(y_{2}(x)\)
\[
y_{2}(x)=y^{\prime}
\]
- Define new variable \(y_{3}(x)\)
\[
y_{3}(x)=y^{\prime \prime}
\]
- Isolate for \(y_{3}^{\prime}(x)\) using original ODE
\[
y_{3}^{\prime}(x)=2 y_{3}(x)+y_{2}(x)-2 y_{1}(x)
\]

Convert linear ODE into a system of first order ODEs
\[
\left[y_{2}(x)=y_{1}^{\prime}(x), y_{3}(x)=y_{2}^{\prime}(x), y_{3}^{\prime}(x)=2 y_{3}(x)+y_{2}(x)-2 y_{1}(x)\right]
\]
- Define vector
\[
\vec{y}(x)=\left[\begin{array}{l}
y_{1}(x) \\
y_{2}(x) \\
y_{3}(x)
\end{array}\right]
\]
- System to solve
\[
\vec{y}^{\prime}(x)=\left[\begin{array}{ccc}
0 & 1 & 0 \\
0 & 0 & 1 \\
-2 & 1 & 2
\end{array}\right] \cdot \vec{y}(x)
\]
- Define the coefficient matrix
\[
A=\left[\begin{array}{ccc}
0 & 1 & 0 \\
0 & 0 & 1 \\
-2 & 1 & 2
\end{array}\right]
\]
- Rewrite the system as
\[
\vec{y}^{\prime}(x)=A \cdot \vec{y}(x)
\]
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- \(\quad\) Eigenpairs of \(A\)
\[
\left[\left[-1,\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]\right],\left[1,\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]\right],\left[2,\left[\begin{array}{c}
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]\right]\right]
\]
- Consider eigenpair
\[
\left[-1,\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]\right]
\]
- \(\quad\) Solution to homogeneous system from eigenpair
\[
\vec{y}_{1}=\mathrm{e}^{-x} \cdot\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]
\]
- Consider eigenpair
\[
\left[1,\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]\right]
\]
- Solution to homogeneous system from eigenpair
\[
\vec{y}_{2}=\mathrm{e}^{x} \cdot\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]
\]
- Consider eigenpair
\[
\left[2,\left[\begin{array}{c}
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]\right]
\]
- Solution to homogeneous system from eigenpair
\[
\vec{y}_{3}=\mathrm{e}^{2 x} \cdot\left[\begin{array}{c}
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]
\]
- General solution to the system of ODEs
\[
\vec{y}=c_{1} \vec{y}_{1}+c_{2} \vec{y}_{2}+c_{3} \vec{y}_{3}
\]
- \(\quad\) Substitute solutions into the general solution
\[
\vec{y}=c_{1} \mathrm{e}^{-x} \cdot\left[\begin{array}{c}
1 \\
-1 \\
1
\end{array}\right]+c_{2} \mathrm{e}^{x} \cdot\left[\begin{array}{c}
1 \\
1 \\
1
\end{array}\right]+\mathrm{e}^{2 x} c_{3} \cdot\left[\begin{array}{c}
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]
\]
- First component of the vector is the solution to the ODE \(y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+\frac{\mathrm{e}^{2 x} c_{3}}{4}\)

Maple trace
```

`Methods for third order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 21
dsolve(diff \((y(x), x \$ 3)-2 * \operatorname{diff}(y(x), x \$ 2)-\operatorname{diff}(y(x), x)+2 * y(x)=0, y(x)\), singsol=all)
\[
y(x)=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+c_{3} \mathrm{e}^{2 x}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.005 (sec). Leaf size: 28
DSolve[y'''[x]-2*y''[x]-y'[x]+2*y[x]==0,y[x],x,IncludeSingularSolutions \(\rightarrow\) True]
\[
y(x) \rightarrow c_{1} e^{-x}+c_{2} e^{x}+c_{3} e^{2 x}
\]

\subsection*{1.95 problem 142}

Internal problem ID [12512]
Internal file name [OUTPUT/11164_Monday_October_16_2023_09_54_23_PM_48092678/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR
PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 142.
ODE order: 3.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_linear_constant_coefficients_ODE"

Maple gives the following as the ode type
[[_3rd_order, _missing_x]]
\[
y^{\prime \prime \prime}-3 a y^{\prime \prime}+3 a^{2} y^{\prime}-a^{3} y=0
\]

The characteristic equation is
\[
-a^{3}+3 a^{2} \lambda-3 a \lambda^{2}+\lambda^{3}=0
\]

The roots of the above equation are
\[
\begin{aligned}
\lambda_{1} & =a \\
\lambda_{2} & =a \\
\lambda_{3} & =a
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=c_{1} \mathrm{e}^{a x}+x \mathrm{e}^{a x} c_{2}+x^{2} \mathrm{e}^{a x} c_{3}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=\mathrm{e}^{a x} \\
& y_{2}=\mathrm{e}^{a x} x \\
& y_{3}=x^{2} \mathrm{e}^{a x}
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{a x}+x \mathrm{e}^{a x} c_{2}+x^{2} \mathrm{e}^{a x} c_{3} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=c_{1} \mathrm{e}^{a x}+x \mathrm{e}^{a x} c_{2}+x^{2} \mathrm{e}^{a x} c_{3}
\]

Verified OK.
Maple trace
- Methods for third order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 19
dsolve(diff( \(y(x), x \$ 3)-3 * a * \operatorname{diff}(y(x), x \$ 2)+3 * a^{\wedge} 2 * \operatorname{diff}(y(x), x)-a^{\wedge} 3 * y(x)=0, y(x)\), singsol=all)
\[
y(x)=\mathrm{e}^{a x}\left(c_{3} x^{2}+c_{2} x+c_{1}\right)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.005 (sec). Leaf size: 23
DSolve[y'''[x]-3*a*y''[x]+3*a^2*y'[x]-a^3*y[x]==0,y[x],x,IncludeSingularSolutions \(\rightarrow\) True]
\[
y(x) \rightarrow e^{a x}\left(x\left(c_{3} x+c_{2}\right)+c_{1}\right)
\]

\subsection*{1.96 problem 143}
1.96.1 Maple step by step solution

Internal problem ID [12513]
Internal file name [OUTPUT/11165_Monday_October_16_2023_09_54_23_PM_97915133/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 143.
ODE order: 5.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_linear_constant_coefficients_ODE"

Maple gives the following as the ode type
[[_high_order, _missing_x]]
\[
y^{(5)}-4 y^{\prime \prime \prime}=0
\]

The characteristic equation is
\[
\lambda^{5}-4 \lambda^{3}=0
\]

The roots of the above equation are
\[
\begin{aligned}
& \lambda_{1}=0 \\
& \lambda_{2}=0 \\
& \lambda_{3}=0 \\
& \lambda_{4}=2 \\
& \lambda_{5}=-2
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=c_{3} x^{2}+c_{2} x+c_{1}+\mathrm{e}^{-2 x} c_{4}+\mathrm{e}^{2 x} c_{5}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=1 \\
& y_{2}=x \\
& y_{3}=x^{2} \\
& y_{4}=\mathrm{e}^{-2 x} \\
& y_{5}=\mathrm{e}^{2 x}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{3} x^{2}+c_{2} x+c_{1}+\mathrm{e}^{-2 x} c_{4}+\mathrm{e}^{2 x} c_{5} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=c_{3} x^{2}+c_{2} x+c_{1}+\mathrm{e}^{-2 x} c_{4}+\mathrm{e}^{2 x} c_{5}
\]

Verified OK.

\subsection*{1.96.1 Maple step by step solution}

Let's solve
\(y^{(5)}-4 y^{\prime \prime \prime}=0\)
- Highest derivative means the order of the ODE is 5
\(y^{(5)}\)

\section*{Convert linear ODE into a system of first order ODEs}
- Define new variable \(y_{1}(x)\)
\[
y_{1}(x)=y
\]
- Define new variable \(y_{2}(x)\) \(y_{2}(x)=y^{\prime}\)
- Define new variable \(y_{3}(x)\)
\[
y_{3}(x)=y^{\prime \prime}
\]
- Define new variable \(y_{4}(x)\)
\[
y_{4}(x)=y^{\prime \prime \prime}
\]
- Define new variable \(y_{5}(x)\)
\[
y_{5}(x)=y^{\prime \prime \prime \prime}
\]
- Isolate for \(y_{5}^{\prime}(x)\) using original ODE
\[
y_{5}^{\prime}(x)=4 y_{4}(x)
\]

Convert linear ODE into a system of first order ODEs
\[
\left[y_{2}(x)=y_{1}^{\prime}(x), y_{3}(x)=y_{2}^{\prime}(x), y_{4}(x)=y_{3}^{\prime}(x), y_{5}(x)=y_{4}^{\prime}(x), y_{5}^{\prime}(x)=4 y_{4}(x)\right]
\]
- Define vector
\[
\vec{y}(x)=\left[\begin{array}{l}
y_{1}(x) \\
y_{2}(x) \\
y_{3}(x) \\
y_{4}(x) \\
y_{5}(x)
\end{array}\right]
\]
- System to solve
\[
\vec{y}^{\prime}(x)=\left[\begin{array}{ccccc}
0 & 1 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 & 1 \\
0 & 0 & 0 & 4 & 0
\end{array}\right] \cdot \vec{y}(x)
\]
- Define the coefficient matrix
\[
A=\left[\begin{array}{lllll}
0 & 1 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 & 1 \\
0 & 0 & 0 & 4 & 0
\end{array}\right]
\]
- Rewrite the system as
\[
\vec{y}^{\prime}(x)=A \cdot \vec{y}(x)
\]
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- Eigenpairs of \(A\)
\[
\left[\left[-2,\left[\begin{array}{c}
\frac{1}{16} \\
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]\right],\left[0,\left[\begin{array}{l}
1 \\
0 \\
0 \\
0 \\
0
\end{array}\right]\right],\left[0,\left[\begin{array}{l}
0 \\
0 \\
0 \\
0 \\
0
\end{array}\right]\right],\left[0,\left[\begin{array}{l}
0 \\
0 \\
0 \\
0 \\
0
\end{array}\right]\right],\left[2,\left[\begin{array}{c}
\frac{1}{16} \\
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]\right]\right]
\]
- Consider eigenpair
\(\left[-2,\left[\begin{array}{c}\frac{1}{16} \\ -\frac{1}{8} \\ \frac{1}{4} \\ -\frac{1}{2} \\ 1\end{array}\right]\right]\)
- Solution to homogeneous system from eigenpair
\[
\vec{y}_{1}=\mathrm{e}^{-2 x} \cdot\left[\begin{array}{c}
\frac{1}{16} \\
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]
\]
- Consider eigenpair
\(\left[0,\left[\begin{array}{l}1 \\ 0 \\ 0 \\ 0 \\ 0\end{array}\right]\right]\)
- \(\quad\) Solution to homogeneous system from eigenpair
\[
\vec{y}_{2}=\left[\begin{array}{l}
1 \\
0 \\
0 \\
0 \\
0
\end{array}\right]
\]
- Consider eigenpair
\(\left[0,\left[\begin{array}{l}0 \\ 0 \\ 0 \\ 0 \\ 0\end{array}\right]\right]\)
- \(\quad\) Solution to homogeneous system from eigenpair
\(\vec{y}_{3}=\left[\begin{array}{l}0 \\ 0 \\ 0 \\ 0 \\ 0\end{array}\right]\)
- Consider eigenpair
\(\left[0,\left[\begin{array}{l}0 \\ 0 \\ 0 \\ 0 \\ 0\end{array}\right]\right]\)
- Solution to homogeneous system from eigenpair
\[
\vec{y}_{4}=\left[\begin{array}{l}
0 \\
0 \\
0 \\
0 \\
0
\end{array}\right]
\]
- Consider eigenpair
\(\left[2,\left[\begin{array}{c}\frac{1}{16} \\ \frac{1}{8} \\ \frac{1}{4} \\ \frac{1}{2} \\ 1\end{array}\right]\right]\)
- \(\quad\) Solution to homogeneous system from eigenpair
\[
\vec{y}_{5}=\mathrm{e}^{2 x} \cdot\left[\begin{array}{c}
\frac{1}{16} \\
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]
\]
- General solution to the system of ODEs
\[
\vec{y}=c_{1} \vec{y}_{1}+c_{2} \vec{y}_{2}+c_{3} \vec{y}_{3}+c_{4} \vec{y}_{4}+c_{5} \vec{y}_{5}
\]
- \(\quad\) Substitute solutions into the general solution
\[
\vec{y}=c_{1} \mathrm{e}^{-2 x} \cdot\left[\begin{array}{c}
\frac{1}{16} \\
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]+\mathrm{e}^{2 x} c_{5} \cdot\left[\begin{array}{c}
\frac{1}{16} \\
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]+\left[\begin{array}{c}
c_{2} \\
0 \\
0 \\
0 \\
0
\end{array}\right]
\]
- First component of the vector is the solution to the ODE \(y=\frac{c_{1} \mathrm{e}^{-2 x}}{16}+\frac{\mathrm{e}^{2 x} c_{5}}{16}+c_{2}\)

Maple trace
```

`Methods for high order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 26
```

dsolve(diff(y(x),x\$5)-4*\operatorname{diff}(y(x),x\$3)=0,y(x), singsol=all)

```
\[
y(x)=c_{1}+c_{2} x+c_{3} x^{2}+c_{4} \mathrm{e}^{2 x}+c_{5} \mathrm{e}^{-2 x}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.073 (sec). Leaf size: 39
DSolve[y'''''[x]-4*y'''[x]==0,y[x],x, IncludeSingularSolutions -> True]
\[
y(x) \rightarrow \frac{1}{8} c_{1} e^{2 x}-\frac{1}{8} c_{2} e^{-2 x}+x\left(c_{5} x+c_{4}\right)+c_{3}
\]

\subsection*{1.97 problem 144}
1.97.1 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1003

Internal problem ID [12514]
Internal file name [OUTPUT/11166_Monday_October_16_2023_09_54_23_PM_42698315/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 144.
ODE order: 4.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_linear_constant_coefficients_ODE"

Maple gives the following as the ode type
[[_high_order, _missing_x]]
\[
y^{\prime \prime \prime \prime}+2 y^{\prime \prime}+9 y=0
\]

The characteristic equation is
\[
\lambda^{4}+2 \lambda^{2}+9=0
\]

The roots of the above equation are
\[
\begin{aligned}
& \lambda_{1}=i \sqrt{2}-1 \\
& \lambda_{2}=-1-i \sqrt{2} \\
& \lambda_{3}=1-i \sqrt{2} \\
& \lambda_{4}=1+i \sqrt{2}
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=\mathrm{e}^{(1+i \sqrt{2}) x} c_{1}+\mathrm{e}^{(1-i \sqrt{2}) x} c_{2}+\mathrm{e}^{(i \sqrt{2}-1) x} c_{3}+\mathrm{e}^{(-1-i \sqrt{2}) x} c_{4}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=\mathrm{e}^{(1+i \sqrt{2}) x} \\
& y_{2}=\mathrm{e}^{(1-i \sqrt{2}) x} \\
& y_{3}=\mathrm{e}^{(i \sqrt{2}-1) x} \\
& y_{4}=\mathrm{e}^{(-1-i \sqrt{2}) x}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{(1+i \sqrt{2}) x} c_{1}+\mathrm{e}^{(1-i \sqrt{2}) x} c_{2}+\mathrm{e}^{(i \sqrt{2}-1) x} c_{3}+\mathrm{e}^{(-1-i \sqrt{2}) x} c_{4} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=\mathrm{e}^{(1+i \sqrt{2}) x} c_{1}+\mathrm{e}^{(1-i \sqrt{2}) x} c_{2}+\mathrm{e}^{(i \sqrt{2}-1) x} c_{3}+\mathrm{e}^{(-1-i \sqrt{2}) x} c_{4}
\]

Verified OK.

\subsection*{1.97.1 Maple step by step solution}

Let's solve
\(y^{\prime \prime \prime \prime}+2 y^{\prime \prime}+9 y=0\)
- Highest derivative means the order of the ODE is 4 \(y^{\prime \prime \prime \prime}\)
Convert linear ODE into a system of first order ODEs
- Define new variable \(y_{1}(x)\)
\[
y_{1}(x)=y
\]
- Define new variable \(y_{2}(x)\)
\[
y_{2}(x)=y^{\prime}
\]
- Define new variable \(y_{3}(x)\)
\[
y_{3}(x)=y^{\prime \prime}
\]
- Define new variable \(y_{4}(x)\)
\[
y_{4}(x)=y^{\prime \prime \prime}
\]
- Isolate for \(y_{4}^{\prime}(x)\) using original ODE
\(y_{4}^{\prime}(x)=-2 y_{3}(x)-9 y_{1}(x)\)

Convert linear ODE into a system of first order ODEs
\(\left[y_{2}(x)=y_{1}^{\prime}(x), y_{3}(x)=y_{2}^{\prime}(x), y_{4}(x)=y_{3}^{\prime}(x), y_{4}^{\prime}(x)=-2 y_{3}(x)-9 y_{1}(x)\right]\)
- Define vector
\(\vec{y}(x)=\left[\begin{array}{l}y_{1}(x) \\ y_{2}(x) \\ y_{3}(x) \\ y_{4}(x)\end{array}\right]\)
- \(\quad\) System to solve
\[
\vec{y}^{\prime}(x)=\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
-9 & 0 & -2 & 0
\end{array}\right] \cdot \vec{y}(x)
\]
- Define the coefficient matrix
\(A=\left[\begin{array}{cccc}0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -9 & 0 & -2 & 0\end{array}\right]\)
- Rewrite the system as
\(\vec{y}^{\prime}(x)=A \cdot \vec{y}(x)\)
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- \(\quad\) Eigenpairs of \(A\)

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored
\[
\left[-1-\mathrm{I} \sqrt{2},\left[\begin{array}{c}
\frac{1}{(-1-\mathrm{I} \sqrt{2})^{3}} \\
\frac{1}{(-1-\mathrm{I} \sqrt{2})^{2}} \\
\frac{1}{-1-\mathrm{I} \sqrt{2}} \\
1
\end{array}\right]\right]
\]
- Solution from eigenpair
\[
\mathrm{e}^{(-1-\mathrm{I} \sqrt{2}) x} \cdot\left[\begin{array}{c}
\frac{1}{(-1-\mathrm{I} \sqrt{2})^{3}} \\
\frac{1}{(-1-\mathrm{I} \sqrt{2})^{2}} \\
\frac{1}{-1-\mathrm{I} \sqrt{2}} \\
1
\end{array}\right]
\]
- Use Euler identity to write solution in terms of sin and cos
\[
\mathrm{e}^{-x} \cdot(\cos (\sqrt{2} x)-\mathrm{I} \sin (\sqrt{2} x)) \cdot\left[\begin{array}{c}
\frac{1}{(-1-\mathrm{I} \sqrt{2})^{3}} \\
\frac{1}{(-1-\mathrm{I} \sqrt{2})^{2}} \\
\frac{1}{-1-\mathrm{I} \sqrt{2}} \\
1
\end{array}\right]
\]
- Simplify expression
\[
\mathrm{e}^{-x} \cdot\left[\begin{array}{c}
\frac{\cos (\sqrt{2} x)-\mathrm{I} \sin (\sqrt{2} x)}{(-1-\mathrm{I} \sqrt{2})^{3}} \\
\frac{\cos (\sqrt{2} x)-\mathrm{I} \sin (\sqrt{2} x)}{(-1-\mathrm{I} \sqrt{2})^{2}} \\
\frac{\cos (\sqrt{2} x)-\mathrm{I} \sin (\sqrt{2} x)}{-1-\mathrm{I} \sqrt{2}} \\
\cos (\sqrt{2} x)-\mathrm{I} \sin (\sqrt{2} x)
\end{array}\right]
\]
- Both real and imaginary parts are solutions to the homogeneous system
\[
\left[\vec{y}_{1}(x)=\mathrm{e}^{-x} \cdot\left[\begin{array}{c}
\frac{5 \cos (\sqrt{2} x)}{27}+\frac{\sin (\sqrt{2} x) \sqrt{2}}{27} \\
-\frac{\cos (\sqrt{2} x)}{9}-\frac{2 \sin (\sqrt{2} x) \sqrt{2}}{9} \\
-\frac{\cos (\sqrt{2} x)}{3}+\frac{\sin (\sqrt{2} x) \sqrt{2}}{3} \\
\cos (\sqrt{2} x)
\end{array}\right], \vec{y}_{2}(x)=\mathrm{e}^{-x} \cdot\left[\begin{array}{c}
\frac{\cos (\sqrt{2} x) \sqrt{2}}{27}-\frac{5 \sin (\sqrt{2} x)}{27} \\
-\frac{2 \cos (\sqrt{2} x) \sqrt{2}}{9}+\frac{\sin (\sqrt{2} x)}{9} \\
\frac{\cos (\sqrt{2} x) \sqrt{2}}{3}+\frac{\sin (\sqrt{2} x)}{3} \\
-\sin (\sqrt{2} x)
\end{array}\right]\right]
\]
- Consider complex eigenpair, complex conjugate eigenvalue can be ignored
\[
\left[1-\mathrm{I} \sqrt{2},\left[\begin{array}{c}
\frac{1}{(1-\mathrm{I} \sqrt{2})^{3}} \\
\frac{1}{(1-\mathrm{I} \sqrt{2})^{2}} \\
\frac{1}{1-\mathrm{I} \sqrt{2}} \\
1
\end{array}\right]\right]
\]
- Solution from eigenpair
\[
\mathrm{e}^{(1-\mathrm{I} \sqrt{2}) x} \cdot\left[\begin{array}{c}
\frac{1}{(1-\mathrm{I} \sqrt{2})^{3}} \\
\frac{1}{(1-\mathrm{I} \sqrt{2})^{2}} \\
\frac{1}{1-\mathrm{I} \sqrt{2}} \\
1
\end{array}\right]
\]
- Use Euler identity to write solution in terms of sin and cos
\[
\mathrm{e}^{x} \cdot(\cos (\sqrt{2} x)-\mathrm{I} \sin (\sqrt{2} x)) \cdot\left[\begin{array}{c}
\frac{1}{(1-\mathrm{I} \sqrt{2})^{3}} \\
\frac{1}{(1-\mathrm{I} \sqrt{2})^{2}} \\
\frac{1}{1-\mathrm{I} \sqrt{2}} \\
1
\end{array}\right]
\]
- Simplify expression
\[
\mathrm{e}^{x} \cdot\left[\begin{array}{c}
\frac{\cos (\sqrt{2} x)-\mathrm{I} \sin (\sqrt{2} x)}{(1-\mathrm{I} \sqrt{2})^{3}} \\
\frac{\cos (\sqrt{2} x)-\mathrm{I} \sin (\sqrt{2} x)}{(1-\mathrm{I} \sqrt{2})^{2}} \\
\frac{\cos (\sqrt{2} x)-\mathrm{I} \sin (\sqrt{2} x)}{1-\mathrm{I} \sqrt{2}} \\
\cos (\sqrt{2} x)-\mathrm{I} \sin (\sqrt{2} x)
\end{array}\right]
\]
- Both real and imaginary parts are solutions to the homogeneous system
\[
\left[\vec{y}_{3}(x)=\mathrm{e}^{x} \cdot\left[\begin{array}{c}
-\frac{5 \cos (\sqrt{2} x)}{27}+\frac{\sin (\sqrt{2} x) \sqrt{2}}{27} \\
-\frac{\cos (\sqrt{2} x)}{9}+\frac{2 \sin (\sqrt{2} x) \sqrt{2}}{9} \\
\frac{\cos (\sqrt{2} x)}{3}+\frac{\sin (\sqrt{2} x) \sqrt{2}}{3} \\
\cos (\sqrt{2} x)
\end{array}\right], \vec{y}_{4}(x)=\mathrm{e}^{x} \cdot\left[\begin{array}{c}
\frac{\cos (\sqrt{2} x) \sqrt{2}}{27}+\frac{5 \sin (\sqrt{2} x)}{27} \\
\frac{2 \cos (\sqrt{2} x) \sqrt{2}}{9}+\frac{\sin (\sqrt{2} x)}{9} \\
\frac{\cos (\sqrt{2} x) \sqrt{2}}{3}-\frac{\sin (\sqrt{2} x)}{3} \\
-\sin (\sqrt{2} x)
\end{array}\right]\right]
\]
- General solution to the system of ODEs
\[
\vec{y}=c_{1} \vec{y}_{1}(x)+c_{2} \vec{y}_{2}(x)+c_{3} \vec{y}_{3}(x)+c_{4} \vec{y}_{4}(x)
\]
- \(\quad\) Substitute solutions into the general solution
- First component of the vector is the solution to the ODE
\[
y=\frac{\left(\left(\sqrt{2} c_{2}+5 c_{1}\right) \mathrm{e}^{-x}-5\left(-\frac{c_{4} \sqrt{2}}{5}+c_{3}\right) \mathrm{e}^{x}\right) \cos (\sqrt{2} x)}{27}+\frac{\sin (\sqrt{2} x)\left(\left(\sqrt{2} c_{1}-5 c_{2}\right) \mathrm{e}^{-x}+\mathrm{e}^{x}\left(\sqrt{2} c_{3}+5 c_{4}\right)\right)}{27}
\]

Maple trace
```

`Methods for high order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 41
```

dsolve(diff(y(x),x\$4)+2*diff (y(x),x\$2)+9*y(x)=0,y(x), singsol=all)

```
\[
y(x)=\left(c_{2} \mathrm{e}^{x}+c_{4} \mathrm{e}^{-x}\right) \cos (x \sqrt{2})+\sin (x \sqrt{2})\left(c_{1} \mathrm{e}^{x}+c_{3} \mathrm{e}^{-x}\right)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.007 (sec). Leaf size: 52
DSolve[y''''[x]+2*y''[x]+9*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
\[
y(x) \rightarrow e^{-x}\left(\left(c_{4} e^{2 x}+c_{2}\right) \cos (\sqrt{2} x)+\left(c_{3} e^{2 x}+c_{1}\right) \sin (\sqrt{2} x)\right)
\]

\subsection*{1.98 problem 145}
1.98.1 Maple step by step solution

1010
Internal problem ID [12515]
Internal file name [OUTPUT/11167_Monday_October_16_2023_09_54_23_PM_15705254/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 145.
ODE order: 4.
ODE degree: 1.

The type(s) of ODE detected by this program : "higher_order_linear_constant_coefficients_ODE"

Maple gives the following as the ode type
[[_high_order, _missing_x]]
\[
y^{\prime \prime \prime \prime}-8 y^{\prime \prime}+16 y=0
\]

The characteristic equation is
\[
\lambda^{4}-8 \lambda^{2}+16=0
\]

The roots of the above equation are
\[
\begin{aligned}
\lambda_{1} & =2 \\
\lambda_{2} & =2 \\
\lambda_{3} & =-2 \\
\lambda_{4} & =-2
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=c_{1} \mathrm{e}^{-2 x}+x \mathrm{e}^{-2 x} c_{2}+\mathrm{e}^{2 x} c_{3}+x \mathrm{e}^{2 x} c_{4}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=\mathrm{e}^{-2 x} \\
& y_{2}=x \mathrm{e}^{-2 x} \\
& y_{3}=\mathrm{e}^{2 x} \\
& y_{4}=\mathrm{e}^{2 x} x
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{-2 x}+x \mathrm{e}^{-2 x} c_{2}+\mathrm{e}^{2 x} c_{3}+x \mathrm{e}^{2 x} c_{4} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=c_{1} \mathrm{e}^{-2 x}+x \mathrm{e}^{-2 x} c_{2}+\mathrm{e}^{2 x} c_{3}+x \mathrm{e}^{2 x} c_{4}
\]

Verified OK.

\subsection*{1.98.1 Maple step by step solution}

Let's solve
\[
y^{\prime \prime \prime \prime}-8 y^{\prime \prime}+16 y=0
\]
- Highest derivative means the order of the ODE is 4 \(y^{\prime \prime \prime \prime}\)

Convert linear ODE into a system of first order ODEs
- Define new variable \(y_{1}(x)\)
\[
y_{1}(x)=y
\]
- Define new variable \(y_{2}(x)\)
\[
y_{2}(x)=y^{\prime}
\]
- Define new variable \(y_{3}(x)\)
\[
y_{3}(x)=y^{\prime \prime}
\]
- Define new variable \(y_{4}(x)\)
\[
y_{4}(x)=y^{\prime \prime \prime}
\]
- Isolate for \(y_{4}^{\prime}(x)\) using original ODE
\(y_{4}^{\prime}(x)=8 y_{3}(x)-16 y_{1}(x)\)
Convert linear ODE into a system of first order ODEs
\(\left[y_{2}(x)=y_{1}^{\prime}(x), y_{3}(x)=y_{2}^{\prime}(x), y_{4}(x)=y_{3}^{\prime}(x), y_{4}^{\prime}(x)=8 y_{3}(x)-16 y_{1}(x)\right]\)
- Define vector
\[
\vec{y}(x)=\left[\begin{array}{l}
y_{1}(x) \\
y_{2}(x) \\
y_{3}(x) \\
y_{4}(x)
\end{array}\right]
\]
- System to solve
\[
\vec{y}^{\prime}(x)=\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
-16 & 0 & 8 & 0
\end{array}\right] \cdot \vec{y}(x)
\]
- Define the coefficient matrix
\[
A=\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
-16 & 0 & 8 & 0
\end{array}\right]
\]
- Rewrite the system as
\[
\vec{y}^{\prime}(x)=A \cdot \vec{y}(x)
\]
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- \(\quad\) Eigenpairs of \(A\)
\begin{tabular}{|c|c|c|c|}
\hline \(\left[-2,\left[\begin{array}{c}-\frac{1}{8} \\ \frac{1}{4} \\ -\frac{1}{2} \\ 1\end{array}\right]\right.\) & ,\(\left[-2,\left[\begin{array}{l}0 \\ 0 \\ 0 \\ 0\end{array}\right]\right.\) & , \(\left[2,\left[\begin{array}{c}\frac{1}{8} \\ \frac{1}{4} \\ \frac{1}{2} \\ 1\end{array}\right]\right.\) & ,\(\left[2,\left[\begin{array}{l}0 \\ 0 \\ 0 \\ 0\end{array}\right]\right]\) \\
\hline
\end{tabular}
- Consider eigenpair, with eigenvalue of algebraic multiplicity 2
\[
\left[-2,\left[\begin{array}{c}
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]\right]
\]
- \(\quad\) First solution from eigenvalue -2
\[
\vec{y}_{1}(x)=\mathrm{e}^{-2 x} \cdot\left[\begin{array}{c}
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]
\]
- Form of the 2nd homogeneous solution where \(\vec{p}\) is to be solved for, \(\lambda=-2\) is the eigenvalue, a
\[
\vec{y}_{2}(x)=\mathrm{e}^{\lambda x}(x \vec{v}+\vec{p})
\]
- \(\quad\) Note that the \(x\) multiplying \(\vec{v}\) makes this solution linearly independent to the 1 st solution obt
- \(\quad\) Substitute \(\vec{y}_{2}(x)\) into the homogeneous system
\(\lambda \mathrm{e}^{\lambda x}(x \vec{v}+\vec{p})+\mathrm{e}^{\lambda x} \vec{v}=\left(\mathrm{e}^{\lambda x} A\right) \cdot(x \vec{v}+\vec{p})\)
- Use the fact that \(\vec{v}\) is an eigenvector of \(A\)
\[
\lambda \mathrm{e}^{\lambda x}(x \vec{v}+\vec{p})+\mathrm{e}^{\lambda x} \vec{v}=\mathrm{e}^{\lambda x}(\lambda x \vec{v}+A \cdot \vec{p})
\]
- Simplify equation
\(\lambda \vec{p}+\vec{v}=A \cdot \vec{p}\)
- \(\quad\) Make use of the identity matrix I
\[
(\lambda \cdot I) \cdot \vec{p}+\vec{v}=A \cdot \vec{p}
\]
- Condition \(\vec{p}\) must meet for \(\vec{y}_{2}(x)\) to be a solution to the homogeneous system
\[
(A-\lambda \cdot I) \cdot \vec{p}=\vec{v}
\]
- Choose \(\vec{p}\) to use in the second solution to the homogeneous system from eigenvalue -2
\[
\left(\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
-16 & 0 & 8 & 0
\end{array}\right]-(-2) \cdot\left[\begin{array}{cccc}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1
\end{array}\right]\right) \cdot \vec{p}=\left[\begin{array}{c}
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]
\]
- \(\quad\) Choice of \(\vec{p}\)
\(\vec{p}=\left[\begin{array}{c}-\frac{1}{16} \\ 0 \\ 0 \\ 0\end{array}\right]\)
- \(\quad\) Second solution from eigenvalue -2
\[
\vec{y}_{2}(x)=\mathrm{e}^{-2 x} \cdot\left(x \cdot\left[\begin{array}{c}
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]+\left[\begin{array}{c}
-\frac{1}{16} \\
0 \\
0 \\
0
\end{array}\right]\right)
\]
- Consider eigenpair, with eigenvalue of algebraic multiplicity 2
\[
\left[2,\left[\begin{array}{c}
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]\right]
\]
- First solution from eigenvalue 2
\[
\vec{y}_{3}(x)=\mathrm{e}^{2 x} \cdot\left[\begin{array}{c}
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]
\]
- Form of the 2nd homogeneous solution where \(\vec{p}\) is to be solved for, \(\lambda=2\) is the eigenvalue, an \(\vec{y}_{4}(x)=\mathrm{e}^{\lambda x}(x \vec{v}+\vec{p})\)
- \(\quad\) Note that the \(x\) multiplying \(\vec{v}\) makes this solution linearly independent to the 1 st solution obt
- \(\quad\) Substitute \(\vec{y}_{4}(x)\) into the homogeneous system
\(\lambda \mathrm{e}^{\lambda x}(x \vec{v}+\vec{p})+\mathrm{e}^{\lambda x} \vec{v}=\left(\mathrm{e}^{\lambda x} A\right) \cdot(x \vec{v}+\vec{p})\)
- Use the fact that \(\vec{v}\) is an eigenvector of \(A\)
\(\lambda \mathrm{e}^{\lambda x}(x \vec{v}+\vec{p})+\mathrm{e}^{\lambda x} \vec{v}=\mathrm{e}^{\lambda x}(\lambda x \vec{v}+A \cdot \vec{p})\)
- \(\quad\) Simplify equation
\(\lambda \vec{p}+\vec{v}=A \cdot \vec{p}\)
- Make use of the identity matrix I
\((\lambda \cdot I) \cdot \vec{p}+\vec{v}=A \cdot \vec{p}\)
- Condition \(\vec{p}\) must meet for \(\vec{y}_{4}(x)\) to be a solution to the homogeneous system
\[
(A-\lambda \cdot I) \cdot \vec{p}=\vec{v}
\]
- Choose \(\vec{p}\) to use in the second solution to the homogeneous system from eigenvalue 2
\[
\left(\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
-16 & 0 & 8 & 0
\end{array}\right]-2 \cdot\left[\begin{array}{cccc}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1
\end{array}\right]\right) \cdot \vec{p}=\left[\begin{array}{c}
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]
\]
- \(\quad\) Choice of \(\vec{p}\)
\(\vec{p}=\left[\begin{array}{c}-\frac{1}{16} \\ 0 \\ 0 \\ 0\end{array}\right]\)
- \(\quad\) Second solution from eigenvalue 2
\(\vec{y}_{4}(x)=\mathrm{e}^{2 x} \cdot\left(x \cdot\left[\begin{array}{c}\frac{1}{8} \\ \frac{1}{4} \\ \frac{1}{2} \\ 1\end{array}\right]+\left[\begin{array}{c}-\frac{1}{16} \\ 0 \\ 0 \\ 0\end{array}\right]\right)\)
- General solution to the system of ODEs
\[
\vec{y}=c_{1} \vec{y}_{1}(x)+c_{2} \vec{y}_{2}(x)+c_{3} \vec{y}_{3}(x)+c_{4} \vec{y}_{4}(x)
\]
- \(\quad\) Substitute solutions into the general solution
\[
\vec{y}=c_{1} \mathrm{e}^{-2 x} \cdot\left[\begin{array}{c}
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]+c_{2} \mathrm{e}^{-2 x} \cdot\left(x \cdot\left[\begin{array}{c}
-\frac{1}{8} \\
\frac{1}{4} \\
-\frac{1}{2} \\
1
\end{array}\right]+\left[\begin{array}{c}
-\frac{1}{16} \\
0 \\
0 \\
0
\end{array}\right]\right)+\mathrm{e}^{2 x} c_{3} \cdot\left[\begin{array}{c}
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]+\mathrm{e}^{2 x} c_{4} \cdot\left[x \cdot\left[\begin{array}{c}
\frac{1}{8} \\
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]\right.
\]
- First component of the vector is the solution to the ODE
\(y=\frac{\left((-2 x-1) c_{2}-2 c_{1}\right) \mathrm{e}^{-2 x}}{16}+\frac{\left(\left(x-\frac{1}{2}\right) c_{4}+c_{3}\right) \mathrm{e}^{2 x}}{8}\)

Maple trace
```

`Methods for high order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.016 (sec). Leaf size: 25
```

dsolve(diff (y (x),x\$4)-8*diff (y(x),x\$2)+16*y(x)=0,y(x), singsol=all)

```
\[
y(x)=\left(c_{4} x+c_{3}\right) \mathrm{e}^{-2 x}+\mathrm{e}^{2 x}\left(c_{2} x+c_{1}\right)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.004 (sec). Leaf size: 35
```

DSolve[y''''[x]-8*y''[x]+16*y[x]==0,y[x],x,IncludeSingularSolutions -> True]

```
\[
y(x) \rightarrow e^{-2 x}\left(c_{3} e^{4 x}+x\left(c_{4} e^{4 x}+c_{2}\right)+c_{1}\right)
\]

\subsection*{1.99 problem 146}
1.99.1 Maple step by step solution

1017
Internal problem ID [12516]
Internal file name [OUTPUT/11168_Monday_October_16_2023_09_54_23_PM_61106002/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 146.
ODE order: 4.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_linear_constant_coefficients_ODE"

Maple gives the following as the ode type
[[_high_order, _missing_x]]
\[
y^{\prime \prime \prime \prime}+y=0
\]

The characteristic equation is
\[
\lambda^{4}+1=0
\]

The roots of the above equation are
\[
\begin{aligned}
& \lambda_{1}=\frac{\sqrt{2}}{2}+\frac{i \sqrt{2}}{2} \\
& \lambda_{2}=-\frac{\sqrt{2}}{2}+\frac{i \sqrt{2}}{2} \\
& \lambda_{3}=-\frac{\sqrt{2}}{2}-\frac{i \sqrt{2}}{2} \\
& \lambda_{4}=\frac{\sqrt{2}}{2}-\frac{i \sqrt{2}}{2}
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=\mathrm{e}^{\left(\frac{\sqrt{2}}{2}+\frac{i \sqrt{2}}{2}\right) x} c_{1}+\mathrm{e}^{\left(-\frac{\sqrt{2}}{2}-\frac{i \sqrt{2}}{2}\right) x} c_{2}+\mathrm{e}^{\left(\frac{\sqrt{2}}{2}-\frac{i \sqrt{2}}{2}\right) x} c_{3}+\mathrm{e}^{\left(-\frac{\sqrt{2}}{2}+\frac{i \sqrt{2}}{2}\right) x} c_{4}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=\mathrm{e}^{\left(\frac{\sqrt{2}}{2}+\frac{i \sqrt{2}}{2}\right) x} \\
& y_{2}=\mathrm{e}^{\left(-\frac{\sqrt{2}}{2}-\frac{i \sqrt{2}}{2}\right) x} \\
& y_{3}=\mathrm{e}^{\left(\frac{\sqrt{2}}{2}-\frac{i \sqrt{2}}{2}\right) x} \\
& y_{4}=\mathrm{e}^{\left(-\frac{\sqrt{2}}{2}+\frac{i \sqrt{2}}{2}\right) x}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{\left(\frac{\sqrt{2}}{2}+\frac{i \sqrt{2}}{2}\right) x} c_{1}+\mathrm{e}^{\left(-\frac{\sqrt{2}}{2}-\frac{i \sqrt{2}}{2}\right) x} c_{2}+\mathrm{e}^{\left(\frac{\sqrt{2}}{2}-\frac{i \sqrt{2}}{2}\right) x} c_{3}+\mathrm{e}^{\left(-\frac{\sqrt{2}}{2}+\frac{i \sqrt{2}}{2}\right) x} c_{4} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=\mathrm{e}^{\left(\frac{\sqrt{2}}{2}+\frac{i \sqrt{2}}{2}\right) x} c_{1}+\mathrm{e}^{\left(-\frac{\sqrt{2}}{2}-\frac{i \sqrt{2}}{2}\right) x} c_{2}+\mathrm{e}^{\left(\frac{\sqrt{2}}{2}-\frac{i \sqrt{2}}{2}\right) x} c_{3}+\mathrm{e}^{\left(-\frac{\sqrt{2}}{2}+\frac{i \sqrt{2}}{2}\right) x} c_{4}
\]

Verified OK.

\subsection*{1.99.1 Maple step by step solution}

Let's solve
\(y^{\prime \prime \prime \prime}+y=0\)
- Highest derivative means the order of the ODE is 4 \(y^{\prime \prime \prime \prime}\)

Convert linear ODE into a system of first order ODEs
- Define new variable \(y_{1}(x)\)
\[
y_{1}(x)=y
\]
- Define new variable \(y_{2}(x)\)
\(y_{2}(x)=y^{\prime}\)
- Define new variable \(y_{3}(x)\)
\(y_{3}(x)=y^{\prime \prime}\)
- Define new variable \(y_{4}(x)\)
\[
y_{4}(x)=y^{\prime \prime \prime}
\]
- Isolate for \(y_{4}^{\prime}(x)\) using original ODE
\(y_{4}^{\prime}(x)=-y_{1}(x)\)

Convert linear ODE into a system of first order ODEs
\(\left[y_{2}(x)=y_{1}^{\prime}(x), y_{3}(x)=y_{2}^{\prime}(x), y_{4}(x)=y_{3}^{\prime}(x), y_{4}^{\prime}(x)=-y_{1}(x)\right]\)
- Define vector
\(\vec{y}(x)=\left[\begin{array}{l}y_{1}(x) \\ y_{2}(x) \\ y_{3}(x) \\ y_{4}(x)\end{array}\right]\)
- \(\quad\) System to solve
\[
\vec{y}^{\prime}(x)=\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
-1 & 0 & 0 & 0
\end{array}\right] \cdot \vec{y}(x)
\]
- Define the coefficient matrix
\[
A=\left[\begin{array}{cccc}
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
-1 & 0 & 0 & 0
\end{array}\right]
\]
- Rewrite the system as
\[
\vec{y}^{\prime}(x)=A \cdot \vec{y}(x)
\]
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- \(\quad\) Eigenpairs of \(A\)

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored
\[
\left[-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2},\left[\begin{array}{c}
\frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{3}} \\
\frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{2}} \\
\frac{1}{-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}} \\
1
\end{array}\right]\right]
\]
- Solution from eigenpair
\[
\mathrm{e}^{\left(-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right) x} \cdot\left[\begin{array}{c}
\frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{3}} \\
\frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{2}} \\
\frac{1}{-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}} \\
1
\end{array}\right]
\]
- Use Euler identity to write solution in terms of sin and cos
\[
\mathrm{e}^{-\frac{\sqrt{2}}{2} x} \cdot\left(\cos \left(\frac{\sqrt{2} x}{2}\right)-\mathrm{I} \sin \left(\frac{\sqrt{2} x}{2}\right)\right) \cdot\left[\begin{array}{c}
\frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{3}} \\
\frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{2}} \\
\frac{1}{-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}} \\
1
\end{array}\right]
\]
- Simplify expression
\[
\mathrm{e}^{-\frac{\sqrt{2} x}{2}} \cdot\left[\begin{array}{c}
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right)-\mathrm{I} \sin \left(\frac{\sqrt{2} x}{2}\right)}{\left(-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{3}} \\
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right)-\mathrm{I} \sin \left(\frac{\sqrt{2} x}{2}\right)}{\left(-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{2}} \\
\frac{\cos \left(\frac{\sqrt{2} x}{2} x\right)-\mathrm{I} \sin \left(\frac{\sqrt{2} x}{2}\right)}{-\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}} \\
\cos \left(\frac{\sqrt{2} x}{2}\right)-\mathrm{I} \sin \left(\frac{\sqrt{2} x}{2}\right)
\end{array}\right]
\]
- Both real and imaginary parts are solutions to the homogeneous system
\[
\left[\vec{y}_{1}(x)=\mathrm{e}^{-\frac{\sqrt{2} x}{2}} \cdot\left[\begin{array}{c}
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}+\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
-\sin \left(\frac{\sqrt{2} x}{2}\right) \\
-\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}+\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
\cos \left(\frac{\sqrt{2} x}{2}\right)
\end{array}\right], \vec{y}_{2}(x)=\mathrm{e}^{-\frac{\sqrt{2} x}{2}} \cdot\left[\begin{array}{c}
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}-\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
-\cos \left(\frac{\sqrt{2} x}{2}\right) \\
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}+\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
-\sin \left(\frac{\sqrt{2} x}{2}\right)
\end{array}\right]\right.
\]
- Consider complex eigenpair, complex conjugate eigenvalue can be ignored
- Solution from eigenpair
\[
\mathrm{e}^{\left(\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right) x} \cdot\left[\begin{array}{c}
\frac{1}{\left(\frac{\sqrt{2}}{2}-\frac{1 \sqrt{2}}{2}\right)^{3}} \\
\frac{1}{\left(\frac{\sqrt{2}}{2}-\frac{1 \sqrt{2}}{2}\right)^{2}} \\
\frac{1}{\frac{\sqrt{2}}{2}-\frac{1 \sqrt{2}}{2}} \\
1
\end{array}\right]
\]
- Use Euler identity to write solution in terms of sin and cos
\[
\mathrm{e}^{\frac{\sqrt{2}}{2} x} \cdot\left(\cos \left(\frac{\sqrt{2} x}{2}\right)-\mathrm{I} \sin \left(\frac{\sqrt{2} x}{2}\right)\right) \cdot\left[\begin{array}{c}
\frac{1}{\left(\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{3}} \\
\frac{1}{\left(\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{2}} \\
\frac{1}{\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}} \\
1
\end{array}\right]
\]
- Simplify expression
\[
\mathrm{e}^{\frac{\sqrt{2} x}{2} x} \cdot\left[\begin{array}{c}
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right)-\mathrm{I} \sin \left(\frac{\sqrt{2} x}{2}\right)}{\left(\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{3}} \\
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right)-\mathrm{I} \sin \left(\frac{\sqrt{2} x}{2}\right)}{\left(\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}\right)^{2}} \\
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right)-\mathrm{I} \sin \left(\frac{\sqrt{2} x}{2}\right)}{\frac{\sqrt{2}}{2}-\frac{\mathrm{I} \sqrt{2}}{2}} \\
\cos \left(\frac{\sqrt{2} x}{2}\right)-\mathrm{I} \sin \left(\frac{\sqrt{2} x}{2}\right)
\end{array}\right]
\]
- Both real and imaginary parts are solutions to the homogeneous system
\[
\left[\vec{y}_{3}(x)=\mathrm{e}^{\frac{\sqrt{2} x}{2}} \cdot\left[\begin{array}{c}
-\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}+\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
\sin \left(\frac{\sqrt{2} x}{2}\right) \\
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}+\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
\cos \left(\frac{\sqrt{2} x}{2}\right)
\end{array}\right], \vec{y}_{4}(x)=\mathrm{e}^{\frac{\sqrt{2} x}{2}} \cdot\left[\begin{array}{c}
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}+\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
\cos \left(\frac{\sqrt{2} x}{2}\right) \\
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}-\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
-\sin \left(\frac{\sqrt{2} x}{2}\right)
\end{array}\right]\right]
\]
- General solution to the system of ODEs
\[
\vec{y}=c_{1} \vec{y}_{1}(x)+c_{2} \vec{y}_{2}(x)+c_{3} \vec{y}_{3}(x)+c_{4} \vec{y}_{4}(x)
\]
- \(\quad\) Substitute solutions into the general solution
\[
\vec{y}=c_{1} \mathrm{e}^{-\frac{\sqrt{2} x}{2}} \cdot\left[\begin{array}{c}
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}+\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
-\sin \left(\frac{\sqrt{2} x}{2}\right) \\
-\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}+\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
\cos \left(\frac{\sqrt{2} x}{2}\right)
\end{array}\right]+c_{2} \mathrm{e}^{-\frac{\sqrt{2} x}{2}} \cdot\left[\begin{array}{c}
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}-\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
-\cos \left(\frac{\sqrt{2} x}{2}\right) \\
\frac{\cos \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2}+\frac{\sin \left(\frac{\sqrt{2} x}{2}\right) \sqrt{2}}{2} \\
-\sin \left(\frac{\sqrt{2} x}{2}\right)
\end{array}\right]+c_{3} \mathrm{e}^{\frac{\sqrt{2} x}{2}}\left[\begin{array}{c} 
\\
-
\end{array}\right]
\]
- First component of the vector is the solution to the ODE
\(y=\frac{\sqrt{2}\left(\left(\left(c_{1}+c_{2}\right) \mathrm{e}^{-\frac{\sqrt{2} x}{2}}-\mathrm{e}^{\frac{\sqrt{2} x}{2} x}\left(c_{3}-c_{4}\right)\right) \cos \left(\frac{\sqrt{2} x}{2}\right)+\sin \left(\frac{\sqrt{2} x}{2}\right)\left(\left(c_{1}-c_{2}\right) \mathrm{e}^{-\frac{\sqrt{2} x}{2} x}+\mathrm{e}^{\frac{\sqrt{2} x}{2}}\left(c_{3}+c_{4}\right)\right)\right)}{2}\)

Maple trace
- Methods for high order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 61
```

dsolve(diff(y(x),x\$4)+y(x)=0,y(x), singsol=all)

```
\[
y(x)=\left(-c_{1} \mathrm{e}^{-\frac{x \sqrt{2}}{2}}-c_{2} \mathrm{e}^{\frac{x \sqrt{2}}{2}}\right) \sin \left(\frac{x \sqrt{2}}{2}\right)+\left(c_{3} \mathrm{e}^{-\frac{x \sqrt{2}}{2}}+c_{4} \mathrm{e}^{\frac{x \sqrt{2}}{2}}\right) \cos \left(\frac{x \sqrt{2}}{2}\right)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.006 (sec). Leaf size: 65
DSolve[y'''' \([x]+y[x]==0, y[x], x\), IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow e^{-\frac{x}{\sqrt{2}}}\left(\left(c_{1} e^{\sqrt{2} x}+c_{2}\right) \cos \left(\frac{x}{\sqrt{2}}\right)+\left(c_{4} e^{\sqrt{2} x}+c_{3}\right) \sin \left(\frac{x}{\sqrt{2}}\right)\right)
\]

\subsection*{1.100 problem 147}

Internal problem ID [12517]
Internal file name [OUTPUT/11169_Monday_October_16_2023_09_54_23_PM_29243652/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 147.
ODE order: 4.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_linear_constant_coefficients_ODE"

Maple gives the following as the ode type
[[_high_order, _missing_x]]
\[
y^{\prime \prime \prime \prime}-a^{4} y=0
\]

The characteristic equation is
\[
-a^{4}+\lambda^{4}=0
\]

The roots of the above equation are
\[
\begin{aligned}
& \lambda_{1}=a \\
& \lambda_{2}=-a \\
& \lambda_{3}=i a \\
& \lambda_{4}=-i a
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=\mathrm{e}^{i a x} c_{1}+\mathrm{e}^{-i a x} c_{2}+\mathrm{e}^{-a x} c_{3}+\mathrm{e}^{a x} c_{4}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=\mathrm{e}^{i a x} \\
& y_{2}=\mathrm{e}^{-i a x} \\
& y_{3}=\mathrm{e}^{-a x} \\
& y_{4}=\mathrm{e}^{a x}
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{i a x} c_{1}+\mathrm{e}^{-i a x} c_{2}+\mathrm{e}^{-a x} c_{3}+\mathrm{e}^{a x} c_{4} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=\mathrm{e}^{i a x} c_{1}+\mathrm{e}^{-i a x} c_{2}+\mathrm{e}^{-a x} c_{3}+\mathrm{e}^{a x} c_{4}
\]

Verified OK.
Maple trace
- Methods for high order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 30
dsolve(diff \((y(x), x \$ 4)-a^{\wedge} 4 * y(x)=0, y(x)\), singsol=all)
\[
y(x)=\mathrm{e}^{a x} c_{1}+c_{2} \mathrm{e}^{-a x}+c_{3} \sin (a x)+c_{4} \cos (a x)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.004 (sec). Leaf size: 37
DSolve[y''''[x]-a^4*y[x]==0,y[x],x,IncludeSingularSolutions \(\rightarrow\) True]
\[
y(x) \rightarrow c_{2} e^{-a x}+c_{4} e^{a x}+c_{1} \cos (a x)+c_{3} \sin (a x)
\]

\subsection*{1.101 problem 148}
1.101.1 Solving as second order linear constant coeff ode . . . . . . . . 1025
1.101.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1028
1.101.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1033

Internal problem ID [12518]
Internal file name [OUTPUT/11170_Monday_October_16_2023_09_54_23_PM_27570738/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 148.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _with_linear_symmetries]]
\[
y^{\prime \prime}-7 y^{\prime}+12 y=x
\]

\subsection*{1.101.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=-7, C=12, f(x)=x\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-7 y^{\prime}+12 y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=-7, C=12\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-7 \lambda \mathrm{e}^{\lambda x}+12 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}-7 \lambda+12=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=-7, C=12\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{7}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-7^{2}-(4)(1)(12)} \\
& =\frac{7}{2} \pm \frac{1}{2}
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=\frac{7}{2}+\frac{1}{2} \\
& \lambda_{2}=\frac{7}{2}-\frac{1}{2}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=4 \\
& \lambda_{2}=3
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{(4) x}+c_{2} e^{(3) x}
\end{aligned}
\]

Or
\[
y=c_{1} \mathrm{e}^{4 x}+\mathrm{e}^{3 x} c_{2}
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \mathrm{e}^{4 x}+\mathrm{e}^{3 x} c_{2}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

\section*{\(x\)}

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{1, x\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{3 x}, \mathrm{e}^{4 x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{2} x+A_{1}
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
12 A_{2} x+12 A_{1}-7 A_{2}=x
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{7}{144}, A_{2}=\frac{1}{12}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{x}{12}+\frac{7}{144}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{4 x}+\mathrm{e}^{3 x} c_{2}\right)+\left(\frac{x}{12}+\frac{7}{144}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{4 x}+\mathrm{e}^{3 x} c_{2}+\frac{x}{12}+\frac{7}{144} \tag{1}
\end{equation*}
\]


Figure 165: Slope field plot

Verification of solutions
\[
y=c_{1} \mathrm{e}^{4 x}+\mathrm{e}^{3 x} c_{2}+\frac{x}{12}+\frac{7}{144}
\]

Verified OK.

\subsection*{1.101.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}-7 y^{\prime}+12 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=-7  \tag{3}\\
& C=12
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{1}{4} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=1 \\
& t=4
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=\frac{z(x)}{4} \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi-
\end{tabular} & no condition \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\). & \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 156: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=\frac{1}{4}\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-\frac{x}{2}}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{7}{1} d x} \\
& =z_{1} e^{\frac{7 x}{2}} \\
& =z_{1}\left(\mathrm{e}^{\frac{7 x}{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{3 x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-7}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{7 x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\mathrm{e}^{x}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{3 x}\right)+c_{2}\left(\mathrm{e}^{3 x}\left(\mathrm{e}^{x}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous \(\operatorname{ODE} A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-7 y^{\prime}+12 y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{4 x}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
x
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{1, x\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{3 x}, \mathrm{e}^{4 x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{2} x+A_{1}
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
12 A_{2} x+12 A_{1}-7 A_{2}=x
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{7}{144}, A_{2}=\frac{1}{12}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{x}{12}+\frac{7}{144}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{4 x}\right)+\left(\frac{x}{12}+\frac{7}{144}\right)
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{4 x}+\frac{x}{12}+\frac{7}{144} \tag{1}
\end{equation*}
\]


Figure 166: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{4 x}+\frac{x}{12}+\frac{7}{144}
\]

Verified OK.

\subsection*{1.101.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}-7 y^{\prime}+12 y=x
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\[
r^{2}-7 r+12=0
\]
- Factor the characteristic polynomial
\[
(r-3)(r-4)=0
\]
- Roots of the characteristic polynomial
\[
r=(3,4)
\]
- \(\quad\) 1st solution of the homogeneous ODE
\(y_{1}(x)=\mathrm{e}^{3 x}\)
- \(\quad 2 n d\) solution of the homogeneous ODE
\[
y_{2}(x)=\mathrm{e}^{4 x}
\]
- General solution of the ODE
\[
y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)
\]
- \(\quad\) Substitute in solutions of the homogeneous ODE
\[
y=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{4 x}+y_{p}(x)
\]

Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function \(\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=x\right]\)
- Wronskian of solutions of the homogeneous equation
\[
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
\mathrm{e}^{3 x} & \mathrm{e}^{4 x} \\
3 \mathrm{e}^{3 x} & 4 \mathrm{e}^{4 x}
\end{array}\right]
\]
- Compute Wronskian
\[
W\left(y_{1}(x), y_{2}(x)\right)=\mathrm{e}^{7 x}
\]
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=-\mathrm{e}^{3 x}\left(\int x \mathrm{e}^{-3 x} d x\right)+\mathrm{e}^{4 x}\left(\int \mathrm{e}^{-4 x} x d x\right)
\]
- Compute integrals
\[
y_{p}(x)=\frac{x}{12}+\frac{7}{144}
\]
- Substitute particular solution into general solution to ODE
\[
y=c_{1} \mathrm{e}^{3 x}+c_{2} \mathrm{e}^{4 x}+\frac{x}{12}+\frac{7}{144}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 21
```

dsolve(diff(y(x),x\$2)-7*diff(y(x),x)+12*y(x)=x,y(x), singsol=all)

```
\[
y(x)=c_{2} \mathrm{e}^{3 x}+c_{1} \mathrm{e}^{4 x}+\frac{x}{12}+\frac{7}{144}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.023 (sec). Leaf size: 30
```

DSolve[y''[x]-7*y'[x]+12*y[x]==x,y[x],x,IncludeSingularSolutions -> True]

```
\[
y(x) \rightarrow \frac{x}{12}+c_{1} e^{3 x}+c_{2} e^{4 x}+\frac{7}{144}
\]

\subsection*{1.102 problem 149}
1.102.1 Solving as second order linear constant coeff ode
1.102.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1038
1.102.3 Maple step by step solution 1042

Internal problem ID [12519]
Internal file name [OUTPUT/11171_Monday_October_16_2023_09_54_26_PM_71095273/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 149.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _with_linear_symmetries]]
\[
s^{\prime \prime}-a^{2} s=1+t
\]

\subsection*{1.102.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A s^{\prime \prime}(t)+B s^{\prime}(t)+C s(t)=f(t)
\]

Where \(A=1, B=0, C=-a^{2}, f(t)=1+t\). Let the solution be
\[
s=s_{h}+s_{p}
\]

Where \(s_{h}\) is the solution to the homogeneous ODE \(A s^{\prime \prime}(t)+B s^{\prime}(t)+C s(t)=0\), and \(s_{p}\) is a particular solution to the non-homogeneous ODE \(A s^{\prime \prime}(t)+B s^{\prime}(t)+C s(t)=f(t)\). \(s_{h}\) is the solution to
\[
s^{\prime \prime}-a^{2} s=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A s^{\prime \prime}(t)+B s^{\prime}(t)+C s(t)=0
\]

Where in the above \(A=1, B=0, C=-a^{2}\). Let the solution be \(s=e^{\lambda t}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda t}-a^{2} \mathrm{e}^{\lambda t}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\mathrm{Eq}(2)\) throughout by \(e^{\lambda t}\) gives
\[
\begin{equation*}
-a^{2}+\lambda^{2}=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=-a^{2}\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)\left(-a^{2}\right)} \\
& = \pm \sqrt{a^{2}}
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=+\sqrt{a^{2}} \\
& \lambda_{2}=-\sqrt{a^{2}}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=\sqrt{a^{2}} \\
& \lambda_{2}=-\sqrt{a^{2}}
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& s=c_{1} e^{\lambda_{1} t}+c_{2} e^{\lambda_{2} t} \\
& s=c_{1} e^{\left(\sqrt{a^{2}}\right) t}+c_{2} e^{\left(-\sqrt{a^{2}}\right) t}
\end{aligned}
\]

Or
\[
s=c_{1} \mathrm{e}^{\sqrt{a^{2}} t}+c_{2} \mathrm{e}^{-\sqrt{a^{2}} t}
\]

Therefore the homogeneous solution \(s_{h}\) is
\[
s_{h}=c_{1} \mathrm{e}^{\sqrt{a^{2}} t}+c_{2} \mathrm{e}^{-\sqrt{a^{2}} t}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
1+t
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{1, t\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{\sqrt{a^{2}} t}, \mathrm{e}^{-\sqrt{a^{2}} t}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
s_{p}=A_{2} t+A_{1}
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(s_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-a^{2}\left(A_{2} t+A_{1}\right)=1+t
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-\frac{1}{a^{2}}, A_{2}=-\frac{1}{a^{2}}\right]
\]

Substituting the above back in the above trial solution \(s_{p}\), gives the particular solution
\[
s_{p}=-\frac{t}{a^{2}}-\frac{1}{a^{2}}
\]

Therefore the general solution is
\[
\begin{aligned}
s & =s_{h}+s_{p} \\
& =\left(c_{1} \mathrm{e}^{\sqrt{a^{2}} t}+c_{2} \mathrm{e}^{-\sqrt{a^{2}} t}\right)+\left(-\frac{t}{a^{2}}-\frac{1}{a^{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
s=c_{1} \mathrm{e}^{\operatorname{csgn}(a) a t}+c_{2} \mathrm{e}^{-\operatorname{csgn}(a) a t}-\frac{t}{a^{2}}-\frac{1}{a^{2}}
\]

Simplifying the solution \(s=c_{1} \mathrm{e}^{\operatorname{csgn}(a) a t}+c_{2} \mathrm{e}^{-\operatorname{csgn}(a) a t}-\frac{t}{a^{2}}-\frac{1}{a^{2}}\) to \(s=\mathrm{e}^{t a} c_{1}+c_{2} \mathrm{e}^{-t a}-\frac{t}{a^{2}}-\frac{1}{a^{2}}\) Summary
The solution(s) found are the following
\[
\begin{equation*}
s=\mathrm{e}^{t a} c_{1}+c_{2} \mathrm{e}^{-t a}-\frac{t}{a^{2}}-\frac{1}{a^{2}} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
s=\mathrm{e}^{t a} c_{1}+c_{2} \mathrm{e}^{-t a}-\frac{t}{a^{2}}-\frac{1}{a^{2}}
\]

Verified OK.

\subsection*{1.102.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
s^{\prime \prime}-a^{2} s & =0  \tag{1}\\
A s^{\prime \prime}+B s^{\prime}+C s & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
A & =1 \\
B & =0  \tag{3}\\
C & =-a^{2}
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(t)=s e^{\int \frac{B}{2 A} d t}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(t)=r z(t) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{a^{2}}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=a^{2} \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(t)=\left(a^{2}\right) z(t) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(t)\) then \(s\) is found using the inverse transformation
\[
s=z(t) e^{-\int \frac{B}{2 A} d t}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 158: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=a^{2}\) is not a function of \(t\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(t)=\mathrm{e}^{\sqrt{a^{2}} t}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(s\) is found from
\[
s_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d t}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
s_{1} & =z_{1} \\
& =\mathrm{e}^{\sqrt{a^{2}} t}
\end{aligned}
\]

Which simplifies to
\[
s_{1}=\mathrm{e}^{\sqrt{a^{2}} t}
\]

The second solution \(s_{2}\) to the original ode is found using reduction of order
\[
s_{2}=s_{1} \int \frac{e^{\int-\frac{B}{A} d t}}{s_{1}^{2}} d t
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
s_{2} & =s_{1} \int \frac{1}{s_{1}^{2}} d t \\
& =\mathrm{e}^{\sqrt{a^{2}}} t \int \frac{1}{\mathrm{e}^{2 \sqrt{a^{2}} t} d t} \\
& =\mathrm{e}^{\sqrt{a^{2}} t}\left(-\frac{\operatorname{csgn}(a) \mathrm{e}^{-2 \operatorname{csgn}(a) a t}}{2 a}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
& s=c_{1} s_{1}+c_{2} s_{2} \\
& =c_{1}\left(\mathrm{e}^{\sqrt{a^{2}} t}\right)+c_{2}\left(\mathrm{e}^{\sqrt{a^{2}} t}\left(-\frac{\operatorname{csgn}(a) \mathrm{e}^{-2 \operatorname{csgn}(a) a t}}{2 a}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
s=s_{h}+s_{p}
\]

Where \(s_{h}\) is the solution to the homogeneous ODE \(A s^{\prime \prime}(t)+B s^{\prime}(t)+C s(t)=0\), and \(s_{p}\) is a particular solution to the nonhomogeneous ODE \(A s^{\prime \prime}(t)+B s^{\prime}(t)+C s(t)=f(t)\). \(s_{h}\) is the solution to
\[
s^{\prime \prime}-a^{2} s=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
s_{h}=c_{1} \mathrm{e}^{\sqrt{a^{2}} t}-\frac{c_{2} \operatorname{csgn}(a) \mathrm{e}^{-\operatorname{csgn}(a) a t}}{2 a}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
1+t
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{1, t\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{-\frac{\operatorname{csgn}(a) \mathrm{e}^{-\operatorname{csgn}(a) a t}}{2 a}, \mathrm{e}^{\sqrt{a^{2}} t}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
s_{p}=A_{2} t+A_{1}
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(s_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-a^{2}\left(A_{2} t+A_{1}\right)=1+t
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-\frac{1}{a^{2}}, A_{2}=-\frac{1}{a^{2}}\right]
\]

Substituting the above back in the above trial solution \(s_{p}\), gives the particular solution
\[
s_{p}=-\frac{t}{a^{2}}-\frac{1}{a^{2}}
\]

Therefore the general solution is
\[
\begin{aligned}
& s=s_{h}+s_{p} \\
& =\left(c_{1} \mathrm{e}^{\sqrt{a^{2}} t}-\frac{c_{2} \operatorname{csgn}(a) \mathrm{e}^{-\operatorname{csgn}(a) a t}}{2 a}\right)+\left(-\frac{t}{a^{2}}-\frac{1}{a^{2}}\right)
\end{aligned}
\]

Simplifying the solution \(s=c_{1} \mathrm{e}^{\sqrt{a^{2}} t}-\frac{c_{2} \operatorname{csgn}(a) \mathrm{e}^{-\operatorname{cscn}(a) a t}}{2 a}-\frac{t}{a^{2}}-\frac{1}{a^{2}}\) to \(s=c_{1} \mathrm{e}^{\sqrt{a^{2}}} t-\frac{c_{2} \mathrm{e}^{-t a}}{2 a}-\) Summary
The solution(s) found are the following
\(\frac{t}{a^{2}}-\frac{1}{a^{2}}\)
\[
\begin{equation*}
s=c_{1} \mathrm{e}^{\sqrt{a^{2}} t}-\frac{c_{2} \mathrm{e}^{-t a}}{2 a}-\frac{t}{a^{2}}-\frac{1}{a^{2}} \tag{1}
\end{equation*}
\]

\section*{Verification of solutions}
\[
s=c_{1} \mathrm{e}^{\sqrt{a^{2}} t}-\frac{c_{2} \mathrm{e}^{-t a}}{2 a}-\frac{t}{a^{2}}-\frac{1}{a^{2}}
\]

Verified OK.

\subsection*{1.102.3 Maple step by step solution}

Let's solve
\[
s^{\prime \prime}-a^{2} s=1+t
\]
- Highest derivative means the order of the ODE is 2
\[
s^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\[
-a^{2}+r^{2}=0
\]
- Factor the characteristic polynomial
\(-(a-r)(a+r)=0\)
- Roots of the characteristic polynomial
\[
r=(a,-a)
\]
- \(\quad 1\) st solution of the homogeneous ODE
\[
s_{1}(t)=\mathrm{e}^{t a}
\]
- \(\quad 2\) nd solution of the homogeneous ODE
\(s_{2}(t)=\mathrm{e}^{-t a}\)
- General solution of the ODE
\(s=c_{1} s_{1}(t)+c_{2} s_{2}(t)+s_{p}(t)\)
- Substitute in solutions of the homogeneous ODE
\(s=\mathrm{e}^{t a} c_{1}+c_{2} \mathrm{e}^{-t a}+s_{p}(t)\)
Find a particular solution \(s_{p}(t)\) of the ODE
- Use variation of parameters to find \(s_{p}\) here \(f(t)\) is the forcing function
\(\left[s_{p}(t)=-s_{1}(t)\left(\int \frac{s_{2}(t) f(t)}{W\left(s_{1}(t), s_{2}(t)\right)} d t\right)+s_{2}(t)\left(\int \frac{s_{1}(t) f(t)}{W\left(s_{1}(t), s_{2}(t)\right)} d t\right), f(t)=1+t\right]\)
- Wronskian of solutions of the homogeneous equation
\(W\left(s_{1}(t), s_{2}(t)\right)=\left[\begin{array}{cc}\mathrm{e}^{t a} & \mathrm{e}^{-t a} \\ a \mathrm{e}^{t a} & -a \mathrm{e}^{-t a}\end{array}\right]\)
- Compute Wronskian
\(W\left(s_{1}(t), s_{2}(t)\right)=-2 a\)
- Substitute functions into equation for \(s_{p}(t)\)
\(s_{p}(t)=\frac{\mathrm{e}^{t a}\left(\int \mathrm{e}^{-t a}(1+t) d t\right)-\mathrm{e}^{-t a}\left(\int \mathrm{e}^{t a}(1+t) d t\right)}{2 a}\)
- Compute integrals
\(s_{p}(t)=\frac{-1-t}{a^{2}}\)
- Substitute particular solution into general solution to ODE
\(s=\mathrm{e}^{t a} c_{1}+c_{2} \mathrm{e}^{-t a}+\frac{-1-t}{a^{2}}\)

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 32
```

dsolve(diff(s(t),t\$2)-a^2*s(t)=t+1,s(t), singsol=all)

```
\[
s(t)=\frac{\mathrm{e}^{a t} c_{2} a^{2}+\mathrm{e}^{-a t} c_{1} a^{2}-t-1}{a^{2}}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.029 (sec). Leaf size: 31
```

DSolve[s''[t]-a^2*s[t]==1+t,s[t],t,IncludeSingularSolutions -> True]

```
\[
s(t) \rightarrow-\frac{t+1}{a^{2}}+c_{1} e^{a t}+c_{2} e^{-a t}
\]

\subsection*{1.103 problem 150}
1.103.1 Solving as second order linear constant coeff ode
1.103.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1048
1.103.3 Maple step by step solution 1053

Internal problem ID [12520]
Internal file name [OUTPUT/11172_Monday_October_16_2023_09_54_26_PM_6197742/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 150.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _linear, _nonhomogeneous]]
\[
y^{\prime \prime}+y^{\prime}-2 y=8 \sin (2 x)
\]

\subsection*{1.103.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=1, C=-2, f(x)=8 \sin (2 x)\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+y^{\prime}-2 y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=1, C=-2\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+\lambda \mathrm{e}^{\lambda x}-2 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+\lambda-2=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=1, C=-2\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{-1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{1^{2}-(4)(1)(-2)} \\
& =-\frac{1}{2} \pm \frac{3}{2}
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=-\frac{1}{2}+\frac{3}{2} \\
& \lambda_{2}=-\frac{1}{2}-\frac{3}{2}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
\lambda_{1} & =1 \\
\lambda_{2} & =-2
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{(1) x}+c_{2} e^{(-2) x}
\end{aligned}
\]

Or
\[
y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-2 x}
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-2 x}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
8 \sin (2 x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{\cos (2 x), \sin (2 x)\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{x}, \mathrm{e}^{-2 x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \cos (2 x)+A_{2} \sin (2 x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-6 A_{1} \cos (2 x)-6 A_{2} \sin (2 x)-2 A_{1} \sin (2 x)+2 A_{2} \cos (2 x)=8 \sin (2 x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-\frac{2}{5}, A_{2}=-\frac{6}{5}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=-\frac{2 \cos (2 x)}{5}-\frac{6 \sin (2 x)}{5}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-2 x}\right)+\left(-\frac{2 \cos (2 x)}{5}-\frac{6 \sin (2 x)}{5}\right)
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-2 x}-\frac{2 \cos (2 x)}{5}-\frac{6 \sin (2 x)}{5} \tag{1}
\end{equation*}
\]


Figure 167: Slope field plot

Verification of solutions
\[
y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-2 x}-\frac{2 \cos (2 x)}{5}-\frac{6 \sin (2 x)}{5}
\]

Verified OK.

\subsection*{1.103.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+y^{\prime}-2 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=1  \tag{3}\\
& C=-2
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{9}{4} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=9 \\
& t=4
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=\frac{9 z(x)}{4} \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi-
\end{tabular} & no condition \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\). & \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 160: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=\frac{9}{4}\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-\frac{3 x}{2}}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{1}{1} d x} \\
& =z_{1} e^{-\frac{x}{2}} \\
& =z_{1}\left(\mathrm{e}^{-\frac{x}{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{-2 x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{1}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{-x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\frac{\mathrm{e}^{3 x}}{3}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{-2 x}\right)+c_{2}\left(\mathrm{e}^{-2 x}\left(\frac{\mathrm{e}^{3 x}}{3}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+y^{\prime}-2 y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \mathrm{e}^{-2 x}+\frac{c_{2} \mathrm{e}^{x}}{3}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
8 \sin (2 x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{\cos (2 x), \sin (2 x)\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\frac{\mathrm{e}^{x}}{3}, \mathrm{e}^{-2 x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \cos (2 x)+A_{2} \sin (2 x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-6 A_{1} \cos (2 x)-6 A_{2} \sin (2 x)-2 A_{1} \sin (2 x)+2 A_{2} \cos (2 x)=8 \sin (2 x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-\frac{2}{5}, A_{2}=-\frac{6}{5}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=-\frac{2 \cos (2 x)}{5}-\frac{6 \sin (2 x)}{5}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{-2 x}+\frac{c_{2} \mathrm{e}^{x}}{3}\right)+\left(-\frac{2 \cos (2 x)}{5}-\frac{6 \sin (2 x)}{5}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{-2 x}+\frac{c_{2} \mathrm{e}^{x}}{3}-\frac{2 \cos (2 x)}{5}-\frac{6 \sin (2 x)}{5} \tag{1}
\end{equation*}
\]


Figure 168: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \mathrm{e}^{-2 x}+\frac{c_{2} \mathrm{e}^{x}}{3}-\frac{2 \cos (2 x)}{5}-\frac{6 \sin (2 x)}{5}
\]

Verified OK.

\subsection*{1.103.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}+y^{\prime}-2 y=8 \sin (2 x)
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\[
r^{2}+r-2=0
\]
- Factor the characteristic polynomial
\((r+2)(r-1)=0\)
- Roots of the characteristic polynomial
\(r=(-2,1)\)
- \(\quad 1\) st solution of the homogeneous ODE
\(y_{1}(x)=\mathrm{e}^{-2 x}\)
- \(\quad\) 2nd solution of the homogeneous ODE
\(y_{2}(x)=\mathrm{e}^{x}\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)\)
- Substitute in solutions of the homogeneous ODE
\(y=c_{1} \mathrm{e}^{-2 x}+c_{2} \mathrm{e}^{x}+y_{p}(x)\)Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\(\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=8 \sin (2 x)\right]\)
- Wronskian of solutions of the homogeneous equation
\(W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}\mathrm{e}^{-2 x} & \mathrm{e}^{x} \\ -2 \mathrm{e}^{-2 x} & \mathrm{e}^{x}\end{array}\right]\)
- Compute Wronskian
\(W\left(y_{1}(x), y_{2}(x)\right)=3 \mathrm{e}^{-x}\)
- Substitute functions into equation for \(y_{p}(x)\)
\(y_{p}(x)=\frac{8\left(\mathrm{e}^{3 x}\left(\int \mathrm{e}^{-x} \sin (2 x) d x\right)-\left(\int \mathrm{e}^{2 x} \sin (2 x) d x\right)\right) \mathrm{e}^{-2 x}}{3}\)
- Compute integrals
\(y_{p}(x)=-\frac{2 \cos (2 x)}{5}-\frac{6 \sin (2 x)}{5}\)
- Substitute particular solution into general solution to ODE
\(y=c_{1} \mathrm{e}^{-2 x}+c_{2} \mathrm{e}^{x}-\frac{2 \cos (2 x)}{5}-\frac{6 \sin (2 x)}{5}\)

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 36
```

dsolve(diff(y(x),x\$2)+diff(y(x),x)-2*y(x)=8*\operatorname{sin}(2*x),y(x), singsol=all)

```
\[
y(x)=\left(\frac{2(-\cos (2 x)-3 \sin (2 x)) \mathrm{e}^{2 x}}{5}+c_{2} \mathrm{e}^{3 x}+c_{1}\right) \mathrm{e}^{-2 x}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.035 (sec). Leaf size: 35
DSolve[y''[x]+y'[x]-2*y[x]==8*Sin[2*x],y[x],x,IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow c_{1} e^{-2 x}+c_{2} e^{x}-\frac{2}{5}(3 \sin (2 x)+\cos (2 x))
\]

\subsection*{1.104 problem 151}
1.104.1 Solving as second order linear constant coeff ode
1.104.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1059
1.104.3 Maple step by step solution 1064

Internal problem ID [12521]
Internal file name [OUTPUT/11173_Tuesday_October_17_2023_07_20_24_AM_64086885/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 151.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _with_linear_symmetries]]
\[
y^{\prime \prime}-y=2+5 x
\]

\subsection*{1.104.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=0, C=-1, f(x)=2+5 x\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous \(\operatorname{ODE} A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=-1\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-\mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}-1=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=-1\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)(-1)} \\
& = \pm 1
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=+1 \\
& \lambda_{2}=-1
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
\lambda_{1} & =1 \\
\lambda_{2} & =-1
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{(1) x}+c_{2} e^{(-1) x}
\end{aligned}
\]

Or
\[
y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-x}
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-x}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
x+1
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{1, x\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{x}, \mathrm{e}^{-x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{2} x+A_{1}
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-A_{2} x-A_{1}=2+5 x
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-2, A_{2}=-5\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=-2-5 x
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-x}\right)+(-2-5 x)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-x}-2-5 x \tag{1}
\end{equation*}
\]


Figure 169: Slope field plot

\section*{Verification of solutions}
\[
y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{-x}-2-5 x
\]

Verified OK.

\subsection*{1.104.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}-y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=0  \tag{3}\\
& C=-1
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{1}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=1 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 162: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=1\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-x}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\mathrm{e}^{-x}
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{-x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\mathrm{e}^{-x} \int \frac{1}{\mathrm{e}^{-2 x}} d x \\
& =\mathrm{e}^{-x}\left(\frac{\mathrm{e}^{2 x}}{2}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{-x}\right)+c_{2}\left(\mathrm{e}^{-x}\left(\frac{\mathrm{e}^{2 x}}{2}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=\frac{c_{2} \mathrm{e}^{x}}{2}+c_{1} \mathrm{e}^{-x}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
x+1
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{1, x\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\frac{\mathrm{e}^{x}}{2}, \mathrm{e}^{-x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{2} x+A_{1}
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-A_{2} x-A_{1}=2+5 x
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-2, A_{2}=-5\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=-2-5 x
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(\frac{c_{2} \mathrm{e}^{x}}{2}+c_{1} \mathrm{e}^{-x}\right)+(-2-5 x)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{c_{2} \mathrm{e}^{x}}{2}+c_{1} \mathrm{e}^{-x}-2-5 x \tag{1}
\end{equation*}
\]


Figure 170: Slope field plot

\section*{Verification of solutions}
\[
y=\frac{c_{2} \mathrm{e}^{x}}{2}+c_{1} \mathrm{e}^{-x}-2-5 x
\]

Verified OK.

\subsection*{1.104.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}-y=2+5 x
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\(r^{2}-1=0\)
- Factor the characteristic polynomial
\[
(r-1)(r+1)=0
\]
- Roots of the characteristic polynomial
\(r=(-1,1)\)
- \(\quad 1\) st solution of the homogeneous ODE
\(y_{1}(x)=\mathrm{e}^{-x}\)
- \(\quad\) 2nd solution of the homogeneous ODE
\(y_{2}(x)=\mathrm{e}^{x}\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)\)
- Substitute in solutions of the homogeneous ODE
\(y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}+y_{p}(x)\)
Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\(\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=2+5 x\right]\)
- Wronskian of solutions of the homogeneous equation
\(W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}\mathrm{e}^{-x} & \mathrm{e}^{x} \\ -\mathrm{e}^{-x} & \mathrm{e}^{x}\end{array}\right]\)
- Compute Wronskian
\(W\left(y_{1}(x), y_{2}(x)\right)=2\)
- Substitute functions into equation for \(y_{p}(x)\)
\(y_{p}(x)=-\frac{\mathrm{e}^{-x}\left(\int \mathrm{e}^{x}(2+5 x) d x\right)}{2}+\frac{\mathrm{e}^{x}\left(\int \mathrm{e}^{-x}(2+5 x) d x\right)}{2}\)
- Compute integrals
\(y_{p}(x)=-2-5 x\)
- Substitute particular solution into general solution to ODE
\(y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{x}-2-5 x\)

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 19
```

dsolve(diff(y(x),x\$2)-y(x)=5*x+2,y(x), singsol=all)

```
\[
y(x)=c_{2} \mathrm{e}^{-x}+c_{1} \mathrm{e}^{x}-2-5 x
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.021 (sec). Leaf size: 24
```

DSolve[y''[x]-y[x]==5*x+2,y[x],x,IncludeSingularSolutions -> True]

```
\[
y(x) \rightarrow-5 x+c_{1} e^{x}+c_{2} e^{-x}-2
\]

\subsection*{1.105 problem 152}
1.105.1 Solving as second order linear constant coeff ode 1067
1.105.2 Solving as linear second order ode solved by an integrating factor
ode . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 1070
1.105.3 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1071
1.105.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1075

Internal problem ID [12522]
Internal file name [OUTPUT/11174_Tuesday_October_17_2023_07_20_27_AM_14524618/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 152.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear__second_order_ode_solved_by__an_integrating_factor"

Maple gives the following as the ode type
[[_2nd_order, _with_linear_symmetries]]
\[
y^{\prime \prime}-2 a y^{\prime}+y a^{2}=\mathrm{e}^{x}
\]

\subsection*{1.105.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=-2 a, C=a^{2}, f(x)=\mathrm{e}^{x}\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-2 a y^{\prime}+y a^{2}=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=-2 a, C=a^{2}\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-2 a \lambda \mathrm{e}^{\lambda x}+a^{2} \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
a^{2}-2 a \lambda+\lambda^{2}=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=-2 a, C=a^{2}\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{2 a}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(-2 a)^{2}-(4)(1)\left(a^{2}\right)} \\
& =a
\end{aligned}
\]

Hence this is the case of a double root \(\lambda_{1,2}=-a\). Therefore the solution is
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{a x}+c_{2} \mathrm{e}^{a x} x \tag{1}
\end{equation*}
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \mathrm{e}^{a x}+x \mathrm{e}^{a x} c_{2}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
\mathrm{e}^{x}
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{x}\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{a x} x, \mathrm{e}^{a x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{x}
\]

The unknowns \(\left\{A_{1}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
A_{1} \mathrm{e}^{x}-2 a A_{1} \mathrm{e}^{x}+A_{1} \mathrm{e}^{x} a^{2}=\mathrm{e}^{x}
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{1}{(a-1)^{2}}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{\mathrm{e}^{x}}{(a-1)^{2}}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{a x}+x \mathrm{e}^{a x} c_{2}\right)+\left(\frac{\mathrm{e}^{x}}{(a-1)^{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
y=\mathrm{e}^{a x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(a-1)^{2}}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{a x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(a-1)^{2}} \tag{1}
\end{equation*}
\]

\section*{Verification of solutions}
\[
y=\mathrm{e}^{a x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(a-1)^{2}}
\]

Verified OK.

\subsection*{1.105.2 Solving as linear second order ode solved by an integrating factor ode}

The ode satisfies this form
\[
y^{\prime \prime}+p(x) y^{\prime}+\frac{\left(p(x)^{2}+p^{\prime}(x)\right) y}{2}=f(x)
\]

Where \(p(x)=-2 a\). Therefore, there is an integrating factor given by
\[
\begin{aligned}
M(x) & =e^{\frac{1}{2} \int p d x} \\
& =e^{\int-2 a d x} \\
& =\mathrm{e}^{-a x}
\end{aligned}
\]

Multiplying both sides of the ODE by the integrating factor \(M(x)\) makes the left side of the ODE a complete differential
\[
\begin{aligned}
(M(x) y)^{\prime \prime} & =\mathrm{e}^{-a x} \mathrm{e}^{x} \\
\left(\mathrm{e}^{-a x} y\right)^{\prime \prime} & =\mathrm{e}^{-a x} \mathrm{e}^{x}
\end{aligned}
\]

Integrating once gives
\[
\left(\mathrm{e}^{-a x} y\right)^{\prime}=-\frac{\mathrm{e}^{-x(a-1)}}{a-1}+c_{1}
\]

Integrating again gives
\[
\left(\mathrm{e}^{-a x} y\right)=\frac{\mathrm{e}^{-x(a-1)}+x\left(a^{2}-2 a+1\right) c_{1}}{(a-1)^{2}}+c_{2}
\]

Hence the solution is
\[
y=\frac{\frac{\mathrm{e}^{-x(a-1)}+x\left(a^{2}-2 a+1\right) c_{1}}{(a-1)^{2}}+c_{2}}{\mathrm{e}^{-a x}}
\]

Or
\[
y=c_{2} \mathrm{e}^{a x}+\left(\frac{a^{2} x \mathrm{e}^{a x}}{(a-1)^{2}}-\frac{2 a x \mathrm{e}^{a x}}{(a-1)^{2}}+\frac{x \mathrm{e}^{a x}}{(a-1)^{2}}\right) c_{1}+\frac{\mathrm{e}^{x}}{(a-1)^{2}}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{2} \mathrm{e}^{a x}+\left(\frac{a^{2} x \mathrm{e}^{a x}}{(a-1)^{2}}-\frac{2 a x \mathrm{e}^{a x}}{(a-1)^{2}}+\frac{x \mathrm{e}^{a x}}{(a-1)^{2}}\right) c_{1}+\frac{\mathrm{e}^{x}}{(a-1)^{2}} \tag{1}
\end{equation*}
\]

\section*{Verification of solutions}
\[
y=c_{2} \mathrm{e}^{a x}+\left(\frac{a^{2} x \mathrm{e}^{a x}}{(a-1)^{2}}-\frac{2 a x \mathrm{e}^{a x}}{(a-1)^{2}}+\frac{x \mathrm{e}^{a x}}{(a-1)^{2}}\right) c_{1}+\frac{\mathrm{e}^{x}}{(a-1)^{2}}
\]

Verified OK.

\subsection*{1.105.3 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}-2 a y^{\prime}+y a^{2} & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=-2 a  \tag{3}\\
& C=a^{2}
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{0}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=0 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=0 \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 164: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0--\infty \\
& =\infty
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is infinity then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=0\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=1
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2}-\frac{2 a}{1} d x} \\
& =z_{1} e^{a x} \\
& =z_{1}\left(\mathrm{e}^{a x}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{a x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-2 a}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{2 a x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}(x)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{a x}\right)+c_{2}\left(\mathrm{e}^{a x}(x)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-2 a y^{\prime}+y a^{2}=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \mathrm{e}^{a x}+x \mathrm{e}^{a x} c_{2}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
\mathrm{e}^{x}
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{x}\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{a x} x, \mathrm{e}^{a x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{x}
\]

The unknowns \(\left\{A_{1}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
A_{1} \mathrm{e}^{x}-2 a A_{1} \mathrm{e}^{x}+A_{1} \mathrm{e}^{x} a^{2}=\mathrm{e}^{x}
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{1}{(a-1)^{2}}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{\mathrm{e}^{x}}{(a-1)^{2}}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{a x}+x \mathrm{e}^{a x} c_{2}\right)+\left(\frac{\mathrm{e}^{x}}{(a-1)^{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
y=\mathrm{e}^{a x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(a-1)^{2}}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{a x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(a-1)^{2}} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=\mathrm{e}^{a x}\left(c_{2} x+c_{1}\right)+\frac{\mathrm{e}^{x}}{(a-1)^{2}}
\]

Verified OK.

\subsection*{1.105.4 Maple step by step solution}

Let's solve
\(y^{\prime \prime}-2 a y^{\prime}+y a^{2}=\mathrm{e}^{x}\)
- Highest derivative means the order of the ODE is 2
\(y^{\prime \prime}\)
- Characteristic polynomial of homogeneous ODE
\(a^{2}-2 a r+r^{2}=0\)
- Factor the characteristic polynomial
\((a-r)^{2}=0\)
- Root of the characteristic polynomial
\(r=a\)
- \(\quad 1\) st solution of the homogeneous ODE
\(y_{1}(x)=\mathrm{e}^{a x}\)
- Repeated root, multiply \(y_{1}(x)\) by \(x\) to ensure linear independence \(y_{2}(x)=\mathrm{e}^{a x} x\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)\)
- Substitute in solutions of the homogeneous ODE
\[
y=c_{1} \mathrm{e}^{a x}+x \mathrm{e}^{a x} c_{2}+y_{p}(x)
\]

Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\[
\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=\mathrm{e}^{x}\right]
\]
- Wronskian of solutions of the homogeneous equation
\[
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
\mathrm{e}^{a x} & \mathrm{e}^{a x} x \\
a \mathrm{e}^{a x} & a x \mathrm{e}^{a x}+\mathrm{e}^{a x}
\end{array}\right]
\]
- Compute Wronskian
\[
W\left(y_{1}(x), y_{2}(x)\right)=\mathrm{e}^{2 a x}
\]
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=\mathrm{e}^{a x}\left(-\left(\int x \mathrm{e}^{-x(a-1)} d x\right)+\left(\int \mathrm{e}^{-x(a-1)} d x\right) x\right)
\]
- Compute integrals
\[
y_{p}(x)=\frac{\mathrm{e}^{x}}{(a-1)^{2}}
\]
- Substitute particular solution into general solution to ODE
\[
y=c_{1} \mathrm{e}^{a x}+x \mathrm{e}^{a x} c_{2}+\frac{\mathrm{e}^{x}}{(a-1)^{2}}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.016 (sec). Leaf size: 28
dsolve(diff \((y(x), x \$ 2)-2 * a * \operatorname{diff}(y(x), x)+a^{\wedge} 2 * y(x)=\exp (x), y(x)\), singsol=all)
\[
y(x)=\frac{(a-1)^{2}\left(c_{1} x+c_{2}\right) \mathrm{e}^{a x}+\mathrm{e}^{x}}{(a-1)^{2}}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.224 (sec). Leaf size: 28
DSolve[y''[x]-2*a*y'[x]+a^2*y[x]==Exp[x],y[x],x,IncludeSingularSolutions \(\rightarrow\) True]
\[
y(x) \rightarrow \frac{e^{x}}{(a-1)^{2}}+e^{a x}\left(c_{2} x+c_{1}\right)
\]

\subsection*{1.106 problem 153}
1.106.1 Solving as second order linear constant coeff ode . . . . . . . . 1078
1.106.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1081
1.106.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1086

Internal problem ID [12523]
Internal file name [OUTPUT/11175_Tuesday_October_17_2023_07_20_28_AM_3851012/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 153.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _with_linear_symmetries]]
\[
y^{\prime \prime}+6 y^{\prime}+5 y=\mathrm{e}^{2 x}
\]

\subsection*{1.106.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=6, C=5, f(x)=\mathrm{e}^{2 x}\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+6 y^{\prime}+5 y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=6, C=5\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+6 \lambda \mathrm{e}^{\lambda x}+5 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+6 \lambda+5=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=6, C=5\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{-6}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{6^{2}-(4)(1)(5)} \\
& =-3 \pm 2
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=-3+2 \\
& \lambda_{2}=-3-2
\end{aligned}
\]

Which simplifies to
\[
\begin{gathered}
\lambda_{1}=-1 \\
\lambda_{2}=-5
\end{gathered}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{(-1) x}+c_{2} e^{(-5) x}
\end{aligned}
\]

Or
\[
y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{-5 x}
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{-5 x}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
\mathrm{e}^{2 x}
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{2 x}\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{-5 x}, \mathrm{e}^{-x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{2 x}
\]

The unknowns \(\left\{A_{1}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
21 A_{1} \mathrm{e}^{2 x}=\mathrm{e}^{2 x}
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{1}{21}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{\mathrm{e}^{2 x}}{21}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{-5 x}\right)+\left(\frac{\mathrm{e}^{2 x}}{21}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{-5 x}+\frac{\mathrm{e}^{2 x}}{21} \tag{1}
\end{equation*}
\]


Figure 171: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \mathrm{e}^{-x}+c_{2} \mathrm{e}^{-5 x}+\frac{\mathrm{e}^{2 x}}{21}
\]

Verified OK.

\subsection*{1.106.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+6 y^{\prime}+5 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=6  \tag{3}\\
& C=5
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{4}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=4 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=4 z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 166: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=4\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-2 x}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{6}{1} d x} \\
& =z_{1} e^{-3 x} \\
& =z_{1}\left(\mathrm{e}^{-3 x}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{-5 x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{6}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{-6 x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\frac{\mathrm{e}^{4 x}}{4}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{-5 x}\right)+c_{2}\left(\mathrm{e}^{-5 x}\left(\frac{\mathrm{e}^{4 x}}{4}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+6 y^{\prime}+5 y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \mathrm{e}^{-5 x}+\frac{c_{2} \mathrm{e}^{-x}}{4}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
\mathrm{e}^{2 x}
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{2 x}\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\frac{\mathrm{e}^{-x}}{4}, \mathrm{e}^{-5 x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{2 x}
\]

The unknowns \(\left\{A_{1}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
21 A_{1} \mathrm{e}^{2 x}=\mathrm{e}^{2 x}
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{1}{21}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{\mathrm{e}^{2 x}}{21}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{-5 x}+\frac{c_{2} \mathrm{e}^{-x}}{4}\right)+\left(\frac{\mathrm{e}^{2 x}}{21}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{-5 x}+\frac{c_{2} \mathrm{e}^{-x}}{4}+\frac{\mathrm{e}^{2 x}}{21} \tag{1}
\end{equation*}
\]


Figure 172: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \mathrm{e}^{-5 x}+\frac{c_{2} \mathrm{e}^{-x}}{4}+\frac{\mathrm{e}^{2 x}}{21}
\]

Verified OK.

\subsection*{1.106.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}+6 y^{\prime}+5 y=\mathrm{e}^{2 x}
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE \(r^{2}+6 r+5=0\)
- Factor the characteristic polynomial
\[
(r+5)(r+1)=0
\]
- Roots of the characteristic polynomial
\[
r=(-5,-1)
\]
- \(\quad\) 1st solution of the homogeneous ODE
\[
y_{1}(x)=\mathrm{e}^{-5 x}
\]
- \(\quad 2\) nd solution of the homogeneous ODE
\[
y_{2}(x)=\mathrm{e}^{-x}
\]
- General solution of the ODE
\[
y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)
\]
- \(\quad\) Substitute in solutions of the homogeneous ODE
\(y=c_{1} \mathrm{e}^{-5 x}+c_{2} \mathrm{e}^{-x}+y_{p}(x)\)Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function \(\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=\mathrm{e}^{2 x}\right]\)
- Wronskian of solutions of the homogeneous equation
\[
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
\mathrm{e}^{-5 x} & \mathrm{e}^{-x} \\
-5 \mathrm{e}^{-5 x} & -\mathrm{e}^{-x}
\end{array}\right]
\]
- Compute Wronskian
\(W\left(y_{1}(x), y_{2}(x)\right)=4 \mathrm{e}^{-6 x}\)
- Substitute functions into equation for \(y_{p}(x)\)
\(y_{p}(x)=-\frac{\mathrm{e}^{-5 x}\left(\int \mathrm{e}^{7 x} d x\right)}{4}+\frac{\mathrm{e}^{-x}\left(\int \mathrm{e}^{3 x} d x\right)}{4}\)
- Compute integrals
\[
y_{p}(x)=\frac{\mathrm{e}^{2 x}}{21}
\]
- Substitute particular solution into general solution to ODE
\(y=c_{1} \mathrm{e}^{-5 x}+c_{2} \mathrm{e}^{-x}+\frac{\mathrm{e}^{2 x}}{21}\)

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 25
```

dsolve(diff(y(x),x\$2)+6*diff(y(x),x)+5*y(x)=exp(2*x),y(x), singsol=all)

```
\[
y(x)=\frac{\left(\mathrm{e}^{7 x}+21 c_{2} \mathrm{e}^{4 x}+21 c_{1}\right) \mathrm{e}^{-5 x}}{21}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.069 (sec). Leaf size: 31
DSolve[y''[x]+6*y'[x]+5*y[x]==Exp[2*x],y[x],x,IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow \frac{e^{2 x}}{21}+c_{1} e^{-5 x}+c_{2} e^{-x}
\]

\subsection*{1.107 problem 154}
1.107.1 Solving as second order linear constant coeff ode
1.107.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1092
1.107.3 Maple step by step solution 1097

Internal problem ID [12524]
Internal file name [OUTPUT/11176_Tuesday_October_17_2023_07_20_29_AM_66644115/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 154.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _with_linear_symmetries]]
\[
y^{\prime \prime}+9 y=6 \mathrm{e}^{3 x}
\]

\subsection*{1.107.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=0, C=9, f(x)=6 \mathrm{e}^{3 x}\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+9 y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=9\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+9 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+9=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=9\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)(9)} \\
& = \pm 3 i
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=+3 i \\
& \lambda_{2}=-3 i
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=3 i \\
& \lambda_{2}=-3 i
\end{aligned}
\]

Since roots are complex conjugate of each others, then let the roots be
\[
\lambda_{1,2}=\alpha \pm i \beta
\]

Where \(\alpha=0\) and \(\beta=3\). Therefore the final solution, when using Euler relation, can be written as
\[
y=e^{\alpha x}\left(c_{1} \cos (\beta x)+c_{2} \sin (\beta x)\right)
\]

Which becomes
\[
y=e^{0}\left(c_{1} \cos (3 x)+c_{2} \sin (3 x)\right)
\]

Or
\[
y=c_{1} \cos (3 x)+c_{2} \sin (3 x)
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \cos (3 x)+c_{2} \sin (3 x)
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
6 \mathrm{e}^{3 x}
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{3 x}\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\{\cos (3 x), \sin (3 x)\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{3 x}
\]

The unknowns \(\left\{A_{1}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
18 A_{1} \mathrm{e}^{3 x}=6 \mathrm{e}^{3 x}
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{1}{3}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{\mathrm{e}^{3 x}}{3}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \cos (3 x)+c_{2} \sin (3 x)\right)+\left(\frac{\mathrm{e}^{3 x}}{3}\right)
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (3 x)+c_{2} \sin (3 x)+\frac{\mathrm{e}^{3 x}}{3} \tag{1}
\end{equation*}
\]


Figure 173: Slope field plot

Verification of solutions
\[
y=c_{1} \cos (3 x)+c_{2} \sin (3 x)+\frac{\mathrm{e}^{3 x}}{3}
\]

Verified OK.

\subsection*{1.107.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+9 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
A & =1 \\
B & =0  \tag{3}\\
C & =9
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{-9}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
s & =-9 \\
t & =1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=-9 z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi-
\end{tabular} & no condition \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\). & \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 168: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=-9\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\cos (3 x)
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\cos (3 x)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\cos (3 x)
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\cos (3 x) \int \frac{1}{\cos (3 x)^{2}} d x \\
& =\cos (3 x)\left(\frac{\tan (3 x)}{3}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}(\cos (3 x))+c_{2}\left(\cos (3 x)\left(\frac{\tan (3 x)}{3}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+9 y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \cos (3 x)+\frac{c_{2} \sin (3 x)}{3}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
6 \mathrm{e}^{3 x}
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{3 x}\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\frac{\sin (3 x)}{3}, \cos (3 x)\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{3 x}
\]

The unknowns \(\left\{A_{1}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
18 A_{1} \mathrm{e}^{3 x}=6 \mathrm{e}^{3 x}
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{1}{3}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{\mathrm{e}^{3 x}}{3}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \cos (3 x)+\frac{c_{2} \sin (3 x)}{3}\right)+\left(\frac{\mathrm{e}^{3 x}}{3}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (3 x)+\frac{c_{2} \sin (3 x)}{3}+\frac{\mathrm{e}^{3 x}}{3} \tag{1}
\end{equation*}
\]


Figure 174: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \cos (3 x)+\frac{c_{2} \sin (3 x)}{3}+\frac{\mathrm{e}^{3 x}}{3}
\]

Verified OK.

\subsection*{1.107.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}+9 y=6 \mathrm{e}^{3 x}
\]
- Highest derivative means the order of the ODE is 2
\(y^{\prime \prime}\)
- Characteristic polynomial of homogeneous ODE \(r^{2}+9=0\)
- Use quadratic formula to solve for \(r\)
\(r=\frac{0 \pm(\sqrt{-36})}{2}\)
- Roots of the characteristic polynomial
\[
r=(-3 \mathrm{I}, 3 \mathrm{I})
\]
- \(\quad 1\) st solution of the homogeneous ODE
\[
y_{1}(x)=\cos (3 x)
\]
- \(\quad\) 2nd solution of the homogeneous ODE
\[
y_{2}(x)=\sin (3 x)
\]
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)\)
- \(\quad\) Substitute in solutions of the homogeneous ODE
\[
y=c_{1} \cos (3 x)+c_{2} \sin (3 x)+y_{p}(x)
\]

Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\[
\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=6 \mathrm{e}^{3 x}\right]
\]
- Wronskian of solutions of the homogeneous equation
\[
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
\cos (3 x) & \sin (3 x) \\
-3 \sin (3 x) & 3 \cos (3 x)
\end{array}\right]
\]
- Compute Wronskian
\(W\left(y_{1}(x), y_{2}(x)\right)=3\)
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=-2 \cos (3 x)\left(\int \mathrm{e}^{3 x} \sin (3 x) d x\right)+2 \sin (3 x)\left(\int \mathrm{e}^{3 x} \cos (3 x) d x\right)
\]
- Compute integrals
\[
y_{p}(x)=\frac{\mathrm{e}^{3 x}}{3}
\]
- Substitute particular solution into general solution to ODE
\[
y=c_{1} \cos (3 x)+c_{2} \sin (3 x)+\frac{\mathrm{e}^{3 x}}{3}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 23
```

dsolve(diff(y(x),x\$2)+9*y(x)=6*exp(3*x),y(x), singsol=all)

```
\[
y(x)=c_{2} \sin (3 x)+c_{1} \cos (3 x)+\frac{\mathrm{e}^{3 x}}{3}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.029 (sec). Leaf size: 29
DSolve[y'' \([x]+9 * y[x]==6 * \operatorname{Exp}[3 * x], y[x], x\), IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow \frac{e^{3 x}}{3}+c_{1} \cos (3 x)+c_{2} \sin (3 x)
\]

\subsection*{1.108 problem 155}
1.108.1 Solving as second order linear constant coeff ode . . . . . . . . 1100
1.108.2 Solving as second order integrable as is ode . . . . . . . . . . . 1104
1.108.3 Solving as second order ode missing y ode . . . . . . . . . . . . 1106
1.108.4 Solving as type second_order_integrable_as_is (not using ABC
version) . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 1108
1.108.5 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1110
1.108.6 Solving as exact linear second order ode ode . . . . . . . . . . . 1114
1.108.7 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1117

Internal problem ID [12525]
Internal file name [OUTPUT/11177_Tuesday_October_17_2023_07_20_30_AM_62813000/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 155.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "exact linear second order ode", "second_order_integrable_as_is", "second_order_ode_missing_y", "second__order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _missing_y]]
\[
y^{\prime \prime}-3 y^{\prime}=-6 x+2
\]

\subsection*{1.108.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=-3, C=0, f(x)=-6 x+2\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-3 y^{\prime}=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=-3, C=0\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-3 \lambda \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}-3 \lambda=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=-3, C=0\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{3}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-3^{2}-(4)(1)(0)} \\
& =\frac{3}{2} \pm \frac{3}{2}
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=\frac{3}{2}+\frac{3}{2} \\
& \lambda_{2}=\frac{3}{2}-\frac{3}{2}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=3 \\
& \lambda_{2}=0
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{(3) x}+c_{2} e^{(0) x}
\end{aligned}
\]

Or
\[
y=c_{1} \mathrm{e}^{3 x}+c_{2}
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \mathrm{e}^{3 x}+c_{2}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
x+1
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{1, x\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{1, \mathrm{e}^{3 x}\right\}
\]

Since 1 is duplicated in the UC_set, then this basis is multiplied by extra \(x\). The UC_set becomes
\[
\left[\left\{x, x^{2}\right\}\right]
\]

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.
\[
y_{p}=A_{2} x^{2}+A_{1} x
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-6 x A_{2}-3 A_{1}+2 A_{2}=-6 x+2
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=0, A_{2}=1\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=x^{2}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{3 x}+c_{2}\right)+\left(x^{2}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{aligned}
& y=c_{1} \mathrm{e}^{3 x}+c_{2}+x^{2}
\end{aligned}
\]

Figure 175: Slope field plot

Verification of solutions
\[
y=c_{1} \mathrm{e}^{3 x}+c_{2}+x^{2}
\]

Verified OK.

\subsection*{1.108.2 Solving as second order integrable as is ode}

Integrating both sides of the ODE w.r.t \(x\) gives
\[
\begin{aligned}
& \int\left(y^{\prime \prime}-3 y^{\prime}\right) d x=\int(-6 x+2) d x \\
-3 y+y^{\prime}= & -3 x^{2}+2 x+c_{1}
\end{aligned}
\]

Which is now solved for \(y\).
Entering Linear first order ODE solver. In canonical form a linear first order is
\[
y^{\prime}+p(x) y=q(x)
\]

Where here
\[
\begin{aligned}
& p(x)=-3 \\
& q(x)=-3 x^{2}+c_{1}+2 x
\end{aligned}
\]

Hence the ode is
\[
-3 y+y^{\prime}=-3 x^{2}+c_{1}+2 x
\]

The integrating factor \(\mu\) is
\[
\begin{aligned}
\mu & =\mathrm{e}^{\int(-3) d x} \\
& =\mathrm{e}^{-3 x}
\end{aligned}
\]

The ode becomes
\[
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(-3 x^{2}+c_{1}+2 x\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\mathrm{e}^{-3 x} y\right) & =\left(\mathrm{e}^{-3 x}\right)\left(-3 x^{2}+c_{1}+2 x\right) \\
\mathrm{d}\left(\mathrm{e}^{-3 x} y\right) & =\left(\left(-3 x^{2}+c_{1}+2 x\right) \mathrm{e}^{-3 x}\right) \mathrm{d} x
\end{aligned}
\]

Integrating gives
\[
\begin{aligned}
& \mathrm{e}^{-3 x} y=\int\left(-3 x^{2}+c_{1}+2 x\right) \mathrm{e}^{-3 x} \mathrm{~d} x \\
& \mathrm{e}^{-3 x} y=\frac{\mathrm{e}^{-3 x}\left(3 x^{2}-c_{1}\right)}{3}+c_{2}
\end{aligned}
\]

Dividing both sides by the integrating factor \(\mu=\mathrm{e}^{-3 x}\) results in
\[
y=\frac{\mathrm{e}^{3 x} \mathrm{e}^{-3 x}\left(3 x^{2}-c_{1}\right)}{3}+\mathrm{e}^{3 x} c_{2}
\]
which simplifies to
\[
y=x^{2}-\frac{c_{1}}{3}+\mathrm{e}^{3 x} c_{2}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=x^{2}-\frac{c_{1}}{3}+\mathrm{e}^{3 x} c_{2} \tag{1}
\end{equation*}
\]


Figure 176: Slope field plot

Verification of solutions
\[
y=x^{2}-\frac{c_{1}}{3}+\mathrm{e}^{3 x} c_{2}
\]

Verified OK.

\subsection*{1.108.3 Solving as second order ode missing y ode}

This is second order ode with missing dependent variable \(y\). Let
\[
p(x)=y^{\prime}
\]

Then
\[
p^{\prime}(x)=y^{\prime \prime}
\]

Hence the ode becomes
\[
p^{\prime}(x)-3 p(x)+6 x-2=0
\]

Which is now solve for \(p(x)\) as first order ode.
Entering Linear first order ODE solver. In canonical form a linear first order is
\[
p^{\prime}(x)+p(x) p(x)=q(x)
\]

Where here
\[
\begin{aligned}
p(x) & =-3 \\
q(x) & =-6 x+2
\end{aligned}
\]

Hence the ode is
\[
p^{\prime}(x)-3 p(x)=-6 x+2
\]

The integrating factor \(\mu\) is
\[
\begin{aligned}
\mu & =\mathrm{e}^{\int(-3) d x} \\
& =\mathrm{e}^{-3 x}
\end{aligned}
\]

The ode becomes
\[
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu p) & =(\mu)(-6 x+2) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\mathrm{e}^{-3 x} p\right) & =\left(\mathrm{e}^{-3 x}\right)(-6 x+2) \\
\mathrm{d}\left(\mathrm{e}^{-3 x} p\right) & =\left((-6 x+2) \mathrm{e}^{-3 x}\right) \mathrm{d} x
\end{aligned}
\]

Integrating gives
\[
\begin{aligned}
& \mathrm{e}^{-3 x} p=\int(-6 x+2) \mathrm{e}^{-3 x} \mathrm{~d} x \\
& \mathrm{e}^{-3 x} p=2 x \mathrm{e}^{-3 x}+c_{1}
\end{aligned}
\]

Dividing both sides by the integrating factor \(\mu=\mathrm{e}^{-3 x}\) results in
\[
p(x)=2 \mathrm{e}^{3 x} x \mathrm{e}^{-3 x}+c_{1} \mathrm{e}^{3 x}
\]
which simplifies to
\[
p(x)=2 x+c_{1} \mathrm{e}^{3 x}
\]

Since \(p=y^{\prime}\) then the new first order ode to solve is
\[
y^{\prime}=2 x+c_{1} \mathrm{e}^{3 x}
\]

Integrating both sides gives
\[
\begin{aligned}
y & =\int 2 x+c_{1} \mathrm{e}^{3 x} \mathrm{~d} x \\
& =\frac{c_{1} \mathrm{e}^{3 x}}{3}+x^{2}+c_{2}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{c_{1} \mathrm{e}^{3 x}}{3}+x^{2}+c_{2} \tag{1}
\end{equation*}
\]


Figure 177: Slope field plot

\section*{Verification of solutions}
\[
y=\frac{c_{1} \mathrm{e}^{3 x}}{3}+x^{2}+c_{2}
\]

Verified OK.

\subsection*{1.108.4 Solving as type second_order_integrable_as_is (not using ABC version)}

Writing the ode as
\[
y^{\prime \prime}-3 y^{\prime}=-6 x+2
\]

Integrating both sides of the ODE w.r.t \(x\) gives
\[
\begin{aligned}
& \int\left(y^{\prime \prime}-3 y^{\prime}\right) d x=\int(-6 x+2) d x \\
-3 y+y^{\prime}= & -3 x^{2}+2 x+c_{1}
\end{aligned}
\]

Which is now solved for \(y\).
Entering Linear first order ODE solver. In canonical form a linear first order is
\[
y^{\prime}+p(x) y=q(x)
\]

Where here
\[
\begin{aligned}
p(x) & =-3 \\
q(x) & =-3 x^{2}+c_{1}+2 x
\end{aligned}
\]

Hence the ode is
\[
-3 y+y^{\prime}=-3 x^{2}+c_{1}+2 x
\]

The integrating factor \(\mu\) is
\[
\begin{aligned}
\mu & =\mathrm{e}^{\int(-3) d x} \\
& =\mathrm{e}^{-3 x}
\end{aligned}
\]

The ode becomes
\[
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(-3 x^{2}+c_{1}+2 x\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\mathrm{e}^{-3 x} y\right) & =\left(\mathrm{e}^{-3 x}\right)\left(-3 x^{2}+c_{1}+2 x\right) \\
\mathrm{d}\left(\mathrm{e}^{-3 x} y\right) & =\left(\left(-3 x^{2}+c_{1}+2 x\right) \mathrm{e}^{-3 x}\right) \mathrm{d} x
\end{aligned}
\]

Integrating gives
\[
\begin{aligned}
& \mathrm{e}^{-3 x} y=\int\left(-3 x^{2}+c_{1}+2 x\right) \mathrm{e}^{-3 x} \mathrm{~d} x \\
& \mathrm{e}^{-3 x} y=\frac{\mathrm{e}^{-3 x}\left(3 x^{2}-c_{1}\right)}{3}+c_{2}
\end{aligned}
\]

Dividing both sides by the integrating factor \(\mu=\mathrm{e}^{-3 x}\) results in
\[
y=\frac{\mathrm{e}^{3 x} \mathrm{e}^{-3 x}\left(3 x^{2}-c_{1}\right)}{3}+\mathrm{e}^{3 x} c_{2}
\]
which simplifies to
\[
y=x^{2}-\frac{c_{1}}{3}+\mathrm{e}^{3 x} c_{2}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=x^{2}-\frac{c_{1}}{3}+\mathrm{e}^{3 x} c_{2} \tag{1}
\end{equation*}
\]


Figure 178: Slope field plot
\(\underline{\text { Verification of solutions }}\)
\[
y=x^{2}-\frac{c_{1}}{3}+\mathrm{e}^{3 x} c_{2}
\]

Verified OK.

\subsection*{1.108.5 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}-3 y^{\prime} & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=-3  \tag{3}\\
& C=0
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{9}{4} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
s & =9 \\
t & =4
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=\frac{9 z(x)}{4} \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 170: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=\frac{9}{4}\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-\frac{3 x}{2}}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{3}{1} d x}
\end{aligned}
\]
\[
\begin{aligned}
& =z_{1} e^{\frac{3 x}{2}} \\
& =z_{1}\left(\mathrm{e}^{\frac{3 x}{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=1
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-3}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{3 x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\frac{\mathrm{e}^{3 x}}{3}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}(1)+c_{2}\left(1\left(\frac{\mathrm{e}^{3 x}}{3}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-3 y^{\prime}=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1}+\frac{\mathrm{e}^{3 x} c_{2}}{3}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
x+1
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{1, x\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{1, \frac{\mathrm{e}^{3 x}}{3}\right\}
\]

Since 1 is duplicated in the UC_set, then this basis is multiplied by extra \(x\). The UC_set becomes
\[
\left[\left\{x, x^{2}\right\}\right]
\]

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.
\[
y_{p}=A_{2} x^{2}+A_{1} x
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-6 x A_{2}-3 A_{1}+2 A_{2}=-6 x+2
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=0, A_{2}=1\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=x^{2}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1}+\frac{\mathrm{e}^{3 x} c_{2}}{3}\right)+\left(x^{2}\right)
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=c_{1}+\frac{\mathrm{e}^{3 x} c_{2}}{3}+x^{2} \tag{1}
\end{equation*}
\]


Figure 179: Slope field plot

Verification of solutions
\[
y=c_{1}+\frac{\mathrm{e}^{3 x} c_{2}}{3}+x^{2}
\]

Verified OK.

\subsection*{1.108.6 Solving as exact linear second order ode ode}

An ode of the form
\[
p(x) y^{\prime \prime}+q(x) y^{\prime}+r(x) y=s(x)
\]
is exact if
\[
\begin{equation*}
p^{\prime \prime}(x)-q^{\prime}(x)+r(x)=0 \tag{1}
\end{equation*}
\]

For the given ode we have
\[
\begin{aligned}
p(x) & =1 \\
q(x) & =-3 \\
r(x) & =0 \\
s(x) & =-6 x+2
\end{aligned}
\]

Hence
\[
\begin{aligned}
p^{\prime \prime}(x) & =0 \\
q^{\prime}(x) & =0
\end{aligned}
\]

Therefore (1) becomes
\[
0-(0)+(0)=0
\]

Hence the ode is exact. Since we now know the ode is exact, it can be written as
\[
\left(p(x) y^{\prime}+\left(q(x)-p^{\prime}(x)\right) y\right)^{\prime}=s(x)
\]

Integrating gives
\[
p(x) y^{\prime}+\left(q(x)-p^{\prime}(x)\right) y=\int s(x) d x
\]

Substituting the above values for \(p, q, r, s\) gives
\[
-3 y+y^{\prime}=\int-6 x+2 d x
\]

We now have a first order ode to solve which is
\[
-3 y+y^{\prime}=-3 x^{2}+c_{1}+2 x
\]

Entering Linear first order ODE solver. In canonical form a linear first order is
\[
y^{\prime}+p(x) y=q(x)
\]

Where here
\[
\begin{aligned}
p(x) & =-3 \\
q(x) & =-3 x^{2}+c_{1}+2 x
\end{aligned}
\]

Hence the ode is
\[
-3 y+y^{\prime}=-3 x^{2}+c_{1}+2 x
\]

The integrating factor \(\mu\) is
\[
\begin{aligned}
\mu & =\mathrm{e}^{\int(-3) d x} \\
& =\mathrm{e}^{-3 x}
\end{aligned}
\]

The ode becomes
\[
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(-3 x^{2}+c_{1}+2 x\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\mathrm{e}^{-3 x} y\right) & =\left(\mathrm{e}^{-3 x}\right)\left(-3 x^{2}+c_{1}+2 x\right) \\
\mathrm{d}\left(\mathrm{e}^{-3 x} y\right) & =\left(\left(-3 x^{2}+c_{1}+2 x\right) \mathrm{e}^{-3 x}\right) \mathrm{d} x
\end{aligned}
\]

Integrating gives
\[
\begin{aligned}
& \mathrm{e}^{-3 x} y=\int\left(-3 x^{2}+c_{1}+2 x\right) \mathrm{e}^{-3 x} \mathrm{~d} x \\
& \mathrm{e}^{-3 x} y=\frac{\mathrm{e}^{-3 x}\left(3 x^{2}-c_{1}\right)}{3}+c_{2}
\end{aligned}
\]

Dividing both sides by the integrating factor \(\mu=\mathrm{e}^{-3 x}\) results in
\[
y=\frac{\mathrm{e}^{3 x} \mathrm{e}^{-3 x}\left(3 x^{2}-c_{1}\right)}{3}+\mathrm{e}^{3 x} c_{2}
\]
which simplifies to
\[
y=x^{2}-\frac{c_{1}}{3}+\mathrm{e}^{3 x} c_{2}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=x^{2}-\frac{c_{1}}{3}+\mathrm{e}^{3 x} c_{2} \tag{1}
\end{equation*}
\]


Figure 180: Slope field plot

\section*{Verification of solutions}
\[
y=x^{2}-\frac{c_{1}}{3}+\mathrm{e}^{3 x} c_{2}
\]

Verified OK.

\subsection*{1.108.7 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}-3 y^{\prime}=-6 x+2
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\[
r^{2}-3 r=0
\]
- Factor the characteristic polynomial
\[
r(r-3)=0
\]
- Roots of the characteristic polynomial
\[
r=(0,3)
\]
- \(\quad 1\) st solution of the homogeneous ODE
\[
y_{1}(x)=1
\]
- \(\quad 2 n d\) solution of the homogeneous ODE
\[
y_{2}(x)=\mathrm{e}^{3 x}
\]
- General solution of the ODE
\[
y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)
\]
- Substitute in solutions of the homogeneous ODE
\[
y=c_{1}+\mathrm{e}^{3 x} c_{2}+y_{p}(x)
\]Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\[
\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=-6 x+2\right]
\]
- Wronskian of solutions of the homogeneous equation
\[
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
1 & \mathrm{e}^{3 x} \\
0 & 3 \mathrm{e}^{3 x}
\end{array}\right]
\]
- Compute Wronskian
\[
W\left(y_{1}(x), y_{2}(x)\right)=3 \mathrm{e}^{3 x}
\]
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=\frac{2\left(\int(3 x-1) d x\right)}{3}-\frac{2 \mathrm{e}^{3 x}\left(\int(3 x-1) \mathrm{e}^{-3 x} d x\right)}{3}
\]
- Compute integrals
\[
y_{p}(x)=x^{2}
\]
- Substitute particular solution into general solution to ODE
\[
y=c_{1}+\mathrm{e}^{3 x} c_{2}+x^{2}
\]

Maple trace
```

-Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
-> Calling odsolve with the ODE`, diff(_b(_a), _a) = 3*_b(_a)-6*_a+2, _b(_a)`
Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful
<- high order exact linear fully integrable successful`

```

Solution by Maple
Time used: 0.0 (sec). Leaf size: 16
```

dsolve(diff (y (x),x\$2)-3*diff (y (x),x)=2-6*x,y(x), singsol=all)

```
\[
y(x)=\frac{c_{1} \mathrm{e}^{3 x}}{3}+x^{2}+c_{2}
\]

Solution by Mathematica
Time used: 0.121 (sec). Leaf size: 22
DSolve[y'' \([\mathrm{x}]-3 * y\) ' \([\mathrm{x}]==2-6 * \mathrm{x}, \mathrm{y}[\mathrm{x}], \mathrm{x}\), IncludeSingularSolutions \(\rightarrow\) True]
\[
y(x) \rightarrow x^{2}+\frac{1}{3} c_{1} e^{3 x}+c_{2}
\]

\subsection*{1.109 problem 156}
1.109.1 Solving as second order linear constant coeff ode 1120
1.109.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1123
1.109.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1128

Internal problem ID [12526]
Internal file name [OUTPUT/11178_Tuesday_October_17_2023_07_20_31_AM_79597876/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 156.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _linear, _nonhomogeneous]]
\[
y^{\prime \prime}-2 y^{\prime}+3 y=\mathrm{e}^{-x} \cos (x)
\]

\subsection*{1.109.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=-2, C=3, f(x)=\mathrm{e}^{-x} \cos (x)\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-2 y^{\prime}+3 y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=-2, C=3\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-2 \lambda \mathrm{e}^{\lambda x}+3 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}-2 \lambda+3=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=-2, C=3\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-2^{2}-(4)(1)(3)} \\
& =1 \pm i \sqrt{2}
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=1+i \sqrt{2} \\
& \lambda_{2}=1-i \sqrt{2}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=1+i \sqrt{2} \\
& \lambda_{2}=1-i \sqrt{2}
\end{aligned}
\]

Since roots are complex conjugate of each others, then let the roots be
\[
\lambda_{1,2}=\alpha \pm i \beta
\]

Where \(\alpha=1\) and \(\beta=\sqrt{2}\). Therefore the final solution, when using Euler relation, can be written as
\[
y=e^{\alpha x}\left(c_{1} \cos (\beta x)+c_{2} \sin (\beta x)\right)
\]

Which becomes
\[
y=e^{x}\left(c_{1} \cos (\sqrt{2} x)+c_{2} \sin (\sqrt{2} x)\right)
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=\mathrm{e}^{x}\left(c_{1} \cos (\sqrt{2} x)+c_{2} \sin (\sqrt{2} x)\right)
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
\mathrm{e}^{-x} \cos (x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{-x} \cos (x), \mathrm{e}^{-x} \sin (x)\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{x} \cos (\sqrt{2} x), \mathrm{e}^{x} \sin (\sqrt{2} x)\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{-x} \cos (x)+A_{2} \mathrm{e}^{-x} \sin (x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
4 A_{1} \mathrm{e}^{-x} \sin (x)-4 A_{2} \mathrm{e}^{-x} \cos (x)+5 A_{1} \mathrm{e}^{-x} \cos (x)+5 A_{2} \mathrm{e}^{-x} \sin (x)=\mathrm{e}^{-x} \cos (x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{5}{41}, A_{2}=-\frac{4}{41}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{5 \mathrm{e}^{-x} \cos (x)}{41}-\frac{4 \mathrm{e}^{-x} \sin (x)}{41}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(\mathrm{e}^{x}\left(c_{1} \cos (\sqrt{2} x)+c_{2} \sin (\sqrt{2} x)\right)\right)+\left(\frac{5 \mathrm{e}^{-x} \cos (x)}{41}-\frac{4 \mathrm{e}^{-x} \sin (x)}{41}\right)
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{x}\left(c_{1} \cos (\sqrt{2} x)+c_{2} \sin (\sqrt{2} x)\right)+\frac{5 \mathrm{e}^{-x} \cos (x)}{41}-\frac{4 \mathrm{e}^{-x} \sin (x)}{41} \tag{1}
\end{equation*}
\]


Figure 181: Slope field plot

\section*{Verification of solutions}
\[
y=\mathrm{e}^{x}\left(c_{1} \cos (\sqrt{2} x)+c_{2} \sin (\sqrt{2} x)\right)+\frac{5 \mathrm{e}^{-x} \cos (x)}{41}-\frac{4 \mathrm{e}^{-x} \sin (x)}{41}
\]

Verified OK.

\subsection*{1.109.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}-2 y^{\prime}+3 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=-2  \tag{3}\\
& C=3
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{-2}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=-2 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=-2 z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 172: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=-2\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\cos (\sqrt{2} x)
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{2}{1} d x} \\
& =z_{1} e^{x} \\
& =z_{1}\left(\mathrm{e}^{x}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{x} \cos (\sqrt{2} x)
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-2}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{2 x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\frac{\sqrt{2} \tan (\sqrt{2} x)}{2}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{x} \cos (\sqrt{2} x)\right)+c_{2}\left(\mathrm{e}^{x} \cos (\sqrt{2} x)\left(\frac{\sqrt{2} \tan (\sqrt{2} x)}{2}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-2 y^{\prime}+3 y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=\mathrm{e}^{x} \cos (\sqrt{2} x) c_{1}+\frac{c_{2} \sin (\sqrt{2} x) \mathrm{e}^{x} \sqrt{2}}{2}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
\mathrm{e}^{-x} \cos (x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{-x} \cos (x), \mathrm{e}^{-x} \sin (x)\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{x} \cos (\sqrt{2} x), \frac{\sin (\sqrt{2} x) \mathrm{e}^{x} \sqrt{2}}{2}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{-x} \cos (x)+A_{2} \mathrm{e}^{-x} \sin (x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
4 A_{1} \mathrm{e}^{-x} \sin (x)-4 A_{2} \mathrm{e}^{-x} \cos (x)+5 A_{1} \mathrm{e}^{-x} \cos (x)+5 A_{2} \mathrm{e}^{-x} \sin (x)=\mathrm{e}^{-x} \cos (x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{5}{41}, A_{2}=-\frac{4}{41}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{5 \mathrm{e}^{-x} \cos (x)}{41}-\frac{4 \mathrm{e}^{-x} \sin (x)}{41}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(\mathrm{e}^{x} \cos (\sqrt{2} x) c_{1}+\frac{c_{2} \sin (\sqrt{2} x) \mathrm{e}^{x} \sqrt{2}}{2}\right)+\left(\frac{5 \mathrm{e}^{-x} \cos (x)}{41}-\frac{4 \mathrm{e}^{-x} \sin (x)}{41}\right)
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{x} \cos (\sqrt{2} x) c_{1}+\frac{c_{2} \sin (\sqrt{2} x) \mathrm{e}^{x} \sqrt{2}}{2}+\frac{5 \mathrm{e}^{-x} \cos (x)}{41}-\frac{4 \mathrm{e}^{-x} \sin (x)}{41} \tag{1}
\end{equation*}
\]


Figure 182: Slope field plot

Verification of solutions
\[
y=\mathrm{e}^{x} \cos (\sqrt{2} x) c_{1}+\frac{c_{2} \sin (\sqrt{2} x) \mathrm{e}^{x} \sqrt{2}}{2}+\frac{5 \mathrm{e}^{-x} \cos (x)}{41}-\frac{4 \mathrm{e}^{-x} \sin (x)}{41}
\]

Verified OK.

\subsection*{1.109.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}-2 y^{\prime}+3 y=\mathrm{e}^{-x} \cos (x)
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\[
r^{2}-2 r+3=0
\]
- Use quadratic formula to solve for \(r\)
\(r=\frac{2 \pm(\sqrt{-8})}{2}\)
- Roots of the characteristic polynomial
\(r=(1-\mathrm{I} \sqrt{2}, 1+\mathrm{I} \sqrt{2})\)
- \(\quad 1\) st solution of the homogeneous ODE
\(y_{1}(x)=\mathrm{e}^{x} \cos (\sqrt{2} x)\)
- \(\quad 2\) nd solution of the homogeneous ODE
\(y_{2}(x)=\mathrm{e}^{x} \sin (\sqrt{2} x)\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)\)
- \(\quad\) Substitute in solutions of the homogeneous ODE
\(y=\mathrm{e}^{x} \cos (\sqrt{2} x) c_{1}+\mathrm{e}^{x} \sin (\sqrt{2} x) c_{2}+y_{p}(x)\)
Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\[
\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=\mathrm{e}^{-x} \cos (x)\right]
\]
- Wronskian of solutions of the homogeneous equation
\[
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
\mathrm{e}^{x} \cos (\sqrt{2} x) & \mathrm{e}^{x} \sin (\sqrt{2} x) \\
\mathrm{e}^{x} \cos (\sqrt{2} x)-\sin (\sqrt{2} x) \mathrm{e}^{x} \sqrt{2} & \mathrm{e}^{x} \sin (\sqrt{2} x)+\mathrm{e}^{x} \cos (\sqrt{2} x) \sqrt{2}
\end{array}\right]
\]
- Compute Wronskian
\[
W\left(y_{1}(x), y_{2}(x)\right)=\sqrt{2} \mathrm{e}^{2 x}
\]
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=-\frac{\mathrm{e}^{x} \sqrt{2}\left(\cos (\sqrt{2} x)\left(\int \cos (x) \mathrm{e}^{-2 x} \sin (\sqrt{2} x) d x\right)-\sin (\sqrt{2} x)\left(\int \cos (x) \mathrm{e}^{-2 x} \cos (\sqrt{2} x) d x\right)\right)}{2}
\]
- Compute integrals
\(y_{p}(x)=\frac{\mathrm{e}^{-x}(5 \cos (x)-4 \sin (x))}{41}\)
- Substitute particular solution into general solution to ODE
\(y=\mathrm{e}^{x} \cos (\sqrt{2} x) c_{1}+\mathrm{e}^{x} \sin (\sqrt{2} x) c_{2}+\frac{\mathrm{e}^{-x}(5 \cos (x)-4 \sin (x))}{41}\)

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.016 (sec). Leaf size: 38
dsolve(diff \((y(x), x \$ 2)-2 * \operatorname{diff}(y(x), x)+3 * y(x)=\exp (-x) * \cos (x), y(x), \quad\) singsol=all)
\[
y(x)=\mathrm{e}^{x} \cos (x \sqrt{2}) c_{1}+\mathrm{e}^{x} \sin (x \sqrt{2}) c_{2}+\frac{5 \mathrm{e}^{-x}\left(\cos (x)-\frac{4 \sin (x)}{5}\right)}{41}
\]

Solution by Mathematica
Time used: 1.089 (sec). Leaf size: 56
DSolve \([y\) '' \([x]-2 * y\) ' \([x]+3 * y[x]==\operatorname{Exp}[-x] * \operatorname{Cos}[x], y[x], x\), IncludeSingularSolutions \(\rightarrow\) True \(]\)
\[
y(x) \rightarrow-\frac{4}{41} e^{-x} \sin (x)+\frac{5}{41} e^{-x} \cos (x)+c_{2} e^{x} \cos (\sqrt{2} x)+c_{1} e^{x} \sin (\sqrt{2} x)
\]

\subsection*{1.110 problem 157}
1.110.1 Solving as second order linear constant coeff ode . . . . . . . . 1131
1.110.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1135
1.110.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1140

Internal problem ID [12527]
Internal file name [OUTPUT/11179_Tuesday_October_17_2023_07_20_32_AM_84307621/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 157.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _linear, _nonhomogeneous]]
\[
y^{\prime \prime}+4 y=2 \sin (2 x)
\]

\subsection*{1.110.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=0, C=4, f(x)=2 \sin (2 x)\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+4 y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=4\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+4 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+4=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=4\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)(4)} \\
& = \pm 2 i
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=+2 i \\
& \lambda_{2}=-2 i
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=2 i \\
& \lambda_{2}=-2 i
\end{aligned}
\]

Since roots are complex conjugate of each others, then let the roots be
\[
\lambda_{1,2}=\alpha \pm i \beta
\]

Where \(\alpha=0\) and \(\beta=2\). Therefore the final solution, when using Euler relation, can be written as
\[
y=e^{\alpha x}\left(c_{1} \cos (\beta x)+c_{2} \sin (\beta x)\right)
\]

Which becomes
\[
y=e^{0}\left(c_{1} \cos (2 x)+c_{2} \sin (2 x)\right)
\]

Or
\[
y=c_{1} \cos (2 x)+c_{2} \sin (2 x)
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \cos (2 x)+c_{2} \sin (2 x)
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
2 \sin (2 x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{\cos (2 x), \sin (2 x)\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\{\cos (2 x), \sin (2 x)\}
\]

Since \(\cos (2 x)\) is duplicated in the UC_set, then this basis is multiplied by extra \(x\). The UC_set becomes
\[
[\{x \cos (2 x), \sin (2 x) x\}]
\]

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.
\[
y_{p}=A_{1} x \cos (2 x)+A_{2} \sin (2 x) x
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-4 A_{1} \sin (2 x)+4 A_{2} \cos (2 x)=2 \sin (2 x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-\frac{1}{2}, A_{2}=0\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=-\frac{x \cos (2 x)}{2}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \cos (2 x)+c_{2} \sin (2 x)\right)+\left(-\frac{x \cos (2 x)}{2}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (2 x)+c_{2} \sin (2 x)-\frac{x \cos (2 x)}{2} \tag{1}
\end{equation*}
\]


Figure 183: Slope field plot

Verification of solutions
\[
y=c_{1} \cos (2 x)+c_{2} \sin (2 x)-\frac{x \cos (2 x)}{2}
\]

Verified OK.

\subsection*{1.110.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+4 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=0  \tag{3}\\
& C=4
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{-4}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=-4 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=-4 z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 174: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=-4\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\cos (2 x)
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\cos (2 x)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\cos (2 x)
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\cos (2 x) \int \frac{1}{\cos (2 x)^{2}} d x \\
& =\cos (2 x)\left(\frac{\tan (2 x)}{2}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}(\cos (2 x))+c_{2}\left(\cos (2 x)\left(\frac{\tan (2 x)}{2}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+4 y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \cos (2 x)+\frac{c_{2} \sin (2 x)}{2}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
2 \sin (2 x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{\cos (2 x), \sin (2 x)\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\frac{\sin (2 x)}{2}, \cos (2 x)\right\}
\]

Since \(\cos (2 x)\) is duplicated in the UC_set, then this basis is multiplied by extra \(x\). The UC_set becomes
\[
[\{x \cos (2 x), \sin (2 x) x\}]
\]

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.
\[
y_{p}=A_{1} x \cos (2 x)+A_{2} \sin (2 x) x
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-4 A_{1} \sin (2 x)+4 A_{2} \cos (2 x)=2 \sin (2 x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-\frac{1}{2}, A_{2}=0\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=-\frac{x \cos (2 x)}{2}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \cos (2 x)+\frac{c_{2} \sin (2 x)}{2}\right)+\left(-\frac{x \cos (2 x)}{2}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (2 x)+\frac{c_{2} \sin (2 x)}{2}-\frac{x \cos (2 x)}{2} \tag{1}
\end{equation*}
\]


Figure 184: Slope field plot

Verification of solutions
\[
y=c_{1} \cos (2 x)+\frac{c_{2} \sin (2 x)}{2}-\frac{x \cos (2 x)}{2}
\]

Verified OK.

\subsection*{1.110.3 Maple step by step solution}

Let's solve
\(y^{\prime \prime}+4 y=2 \sin (2 x)\)
- Highest derivative means the order of the ODE is 2
\(y^{\prime \prime}\)
- Characteristic polynomial of homogeneous ODE
\(r^{2}+4=0\)
- Use quadratic formula to solve for \(r\)
\(r=\frac{0 \pm(\sqrt{-16})}{2}\)
- Roots of the characteristic polynomial
\[
r=(-2 \mathrm{I}, 2 \mathrm{I})
\]
- \(\quad 1\) st solution of the homogeneous ODE
\(y_{1}(x)=\cos (2 x)\)
- \(\quad 2\) nd solution of the homogeneous ODE
\(y_{2}(x)=\sin (2 x)\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)\)
- Substitute in solutions of the homogeneous ODE
\(y=c_{1} \cos (2 x)+c_{2} \sin (2 x)+y_{p}(x)\)
Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\[
\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=2 \sin (2 x)\right]
\]
- Wronskian of solutions of the homogeneous equation
\[
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
\cos (2 x) & \sin (2 x) \\
-2 \sin (2 x) & 2 \cos (2 x)
\end{array}\right]
\]
- Compute Wronskian
\(W\left(y_{1}(x), y_{2}(x)\right)=2\)
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=-\cos (2 x)\left(\int \sin (2 x)^{2} d x\right)+\frac{\sin (2 x)\left(\int \sin (4 x) d x\right)}{2}
\]
- Compute integrals
\[
y_{p}(x)=\frac{\sin (2 x)}{8}-\frac{x \cos (2 x)}{2}
\]
- Substitute particular solution into general solution to ODE
\[
y=c_{1} \cos (2 x)+c_{2} \sin (2 x)+\frac{\sin (2 x)}{8}-\frac{x \cos (2 x)}{2}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 24
```

dsolve(diff(y(x),x\$2)+4*y(x)=2*\operatorname{sin}(2*x),y(x), singsol=all)

```
\[
y(x)=\frac{\left(-x+2 c_{1}\right) \cos (2 x)}{2}+c_{2} \sin (2 x)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.17 (sec). Leaf size: 33
```

DSolve[y''[x]+4*y[x]==2*Sin[2*x],y[x],x,IncludeSingularSolutions -> True]

```
\[
y(x) \rightarrow \frac{1}{8}\left(\left(1+8 c_{2}\right) \sin (2 x)-4\left(x-2 c_{1}\right) \cos (2 x)\right)
\]

\subsection*{1.111 problem 158}
1.111.1 Maple step by step solution

Internal problem ID [12528]
Internal file name [OUTPUT/11180_Tuesday_October_17_2023_07_20_33_AM_91064150/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 158.
ODE order: 3.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_linear_constant_coefficients_ODE"

Maple gives the following as the ode type
[[_3rd_order, _with_linear_symmetries]]
\[
y^{\prime \prime \prime}-4 y^{\prime \prime}+5 y^{\prime}-2 y=3+2 x
\]

This is higher order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE And \(y_{p}\) is a particular solution to the nonhomogeneous ODE. \(y_{h}\) is the solution to
\[
y^{\prime \prime \prime}-4 y^{\prime \prime}+5 y^{\prime}-2 y=0
\]

The characteristic equation is
\[
\lambda^{3}-4 \lambda^{2}+5 \lambda-2=0
\]

The roots of the above equation are
\[
\begin{aligned}
& \lambda_{1}=2 \\
& \lambda_{2}=1 \\
& \lambda_{3}=1
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=\mathrm{e}^{x} c_{1}+c_{2} x \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{3}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=\mathrm{e}^{x} \\
& y_{2}=x \mathrm{e}^{x} \\
& y_{3}=\mathrm{e}^{2 x}
\end{aligned}
\]

Now the particular solution to the given ODE is found
\[
y^{\prime \prime \prime}-4 y^{\prime \prime}+5 y^{\prime}-2 y=3+2 x
\]

The particular solution is found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
x+1
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{1, x\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{x \mathrm{e}^{x}, \mathrm{e}^{x}, \mathrm{e}^{2 x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{2} x+A_{1}
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-2 A_{2} x-2 A_{1}+5 A_{2}=3+2 x
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-4, A_{2}=-1\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=-x-4
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(\mathrm{e}^{x} c_{1}+c_{2} x \mathrm{e}^{x}+\mathrm{e}^{2 x} c_{3}\right)+(-x-4)
\end{aligned}
\]

Which simplifies to
\[
y=\mathrm{e}^{2 x} c_{3}+\mathrm{e}^{x}\left(c_{2} x+c_{1}\right)-x-4
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{2 x} c_{3}+\mathrm{e}^{x}\left(c_{2} x+c_{1}\right)-x-4 \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=\mathrm{e}^{2 x} c_{3}+\mathrm{e}^{x}\left(c_{2} x+c_{1}\right)-x-4
\]

Verified OK.

\subsection*{1.111.1 Maple step by step solution}

Let's solve
\[
y^{\prime \prime \prime}-4 y^{\prime \prime}+5 y^{\prime}-2 y=3+2 x
\]
- Highest derivative means the order of the ODE is 3
\(y^{\prime \prime \prime}\)
Convert linear ODE into a system of first order ODEs
- Define new variable \(y_{1}(x)\)
\(y_{1}(x)=y\)
- Define new variable \(y_{2}(x)\)
\[
y_{2}(x)=y^{\prime}
\]
- Define new variable \(y_{3}(x)\)
\[
y_{3}(x)=y^{\prime \prime}
\]
- Isolate for \(y_{3}^{\prime}(x)\) using original ODE
\(y_{3}^{\prime}(x)=3+2 x+4 y_{3}(x)-5 y_{2}(x)+2 y_{1}(x)\)
Convert linear ODE into a system of first order ODEs
\[
\left[y_{2}(x)=y_{1}^{\prime}(x), y_{3}(x)=y_{2}^{\prime}(x), y_{3}^{\prime}(x)=3+2 x+4 y_{3}(x)-5 y_{2}(x)+2 y_{1}(x)\right]
\]
- Define vector
\[
\vec{y}(x)=\left[\begin{array}{l}
y_{1}(x) \\
y_{2}(x) \\
y_{3}(x)
\end{array}\right]
\]
- System to solve
\[
\vec{y}^{\prime}(x)=\left[\begin{array}{ccc}
0 & 1 & 0 \\
0 & 0 & 1 \\
2 & -5 & 4
\end{array}\right] \cdot \vec{y}(x)+\left[\begin{array}{c}
0 \\
0 \\
3+2 x
\end{array}\right]
\]
- Define the forcing function
\[
\vec{f}(x)=\left[\begin{array}{c}
0 \\
0 \\
3+2 x
\end{array}\right]
\]
- Define the coefficient matrix
\[
A=\left[\begin{array}{ccc}
0 & 1 & 0 \\
0 & 0 & 1 \\
2 & -5 & 4
\end{array}\right]
\]
- Rewrite the system as
\[
\vec{y}^{\prime}(x)=A \cdot \vec{y}(x)+\vec{f}
\]
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- \(\quad\) Eigenpairs of \(A\)
\[
\left[\left[1,\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]\right],\left[1,\left[\begin{array}{l}
0 \\
0 \\
0
\end{array}\right]\right],\left[2,\left[\begin{array}{c}
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]\right]\right]
\]
- Consider eigenpair, with eigenvalue of algebraic multiplicity 2
\[
\left[1,\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]\right]
\]
- \(\quad\) First solution from eigenvalue 1
\[
\vec{y}_{1}(x)=\mathrm{e}^{x} \cdot\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]
\]
- Form of the 2 nd homogeneous solution where \(\vec{p}\) is to be solved for, \(\lambda=1\) is the eigenvalue, an \(\vec{y}_{2}(x)=\mathrm{e}^{\lambda x}(x \vec{v}+\vec{p})\)
- \(\quad\) Note that the \(x\) multiplying \(\vec{v}\) makes this solution linearly independent to the 1 st solution obt
- \(\quad\) Substitute \(\vec{y}_{2}(x)\) into the homogeneous system
\(\lambda \mathrm{e}^{\lambda x}(x \vec{v}+\vec{p})+\mathrm{e}^{\lambda x} \vec{v}=\left(\mathrm{e}^{\lambda x} A\right) \cdot(x \vec{v}+\vec{p})\)
- Use the fact that \(\vec{v}\) is an eigenvector of \(A\)
\(\lambda \mathrm{e}^{\lambda x}(x \vec{v}+\vec{p})+\mathrm{e}^{\lambda x} \vec{v}=\mathrm{e}^{\lambda x}(\lambda x \vec{v}+A \cdot \vec{p})\)
- Simplify equation
\(\lambda \vec{p}+\vec{v}=A \cdot \vec{p}\)
- Make use of the identity matrix I
\((\lambda \cdot I) \cdot \vec{p}+\vec{v}=A \cdot \vec{p}\)
- Condition \(\vec{p}\) must meet for \(\vec{y}_{2}(x)\) to be a solution to the homogeneous system \((A-\lambda \cdot I) \cdot \vec{p}=\vec{v}\)
- Choose \(\vec{p}\) to use in the second solution to the homogeneous system from eigenvalue 1
\[
\left(\left[\begin{array}{ccc}
0 & 1 & 0 \\
0 & 0 & 1 \\
2 & -5 & 4
\end{array}\right]-1 \cdot\left[\begin{array}{ccc}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right]\right) \cdot \vec{p}=\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]
\]
- Choice of \(\vec{p}\)
\[
\vec{p}=\left[\begin{array}{c}
-1 \\
0 \\
0
\end{array}\right]
\]
- \(\quad\) Second solution from eigenvalue 1
\[
\vec{y}_{2}(x)=\mathrm{e}^{x} \cdot\left(x \cdot\left[\begin{array}{l}
1 \\
1 \\
1
\end{array}\right]+\left[\begin{array}{c}
-1 \\
0 \\
0
\end{array}\right]\right)
\]
- Consider eigenpair
\(\left[2,\left[\begin{array}{c}\frac{1}{4} \\ \frac{1}{2} \\ 1\end{array}\right]\right]\)
- Solution to homogeneous system from eigenpair
\[
\vec{y}_{3}=\mathrm{e}^{2 x} \cdot\left[\begin{array}{c}
\frac{1}{4} \\
\frac{1}{2} \\
1
\end{array}\right]
\]
- General solution of the system of ODEs can be written in terms of the particular solution \(\vec{y}_{p}\) \(\vec{y}(x)=c_{1} \vec{y}_{1}(x)+c_{2} \vec{y}_{2}(x)+c_{3} \vec{y}_{3}+\vec{y}_{p}(x)\)
Fundamental matrix
- Let \(\phi(x)\) be the matrix whose columns are the independent solutions of the homogeneous syst \(\phi(x)=\left[\begin{array}{ccc}\mathrm{e}^{x} & \mathrm{e}^{x}(x-1) & \frac{\mathrm{e}^{2 x}}{4} \\ \mathrm{e}^{x} & x \mathrm{e}^{x} & \frac{\mathrm{e}^{2 x}}{2} \\ \mathrm{e}^{x} & x \mathrm{e}^{x} & \mathrm{e}^{2 x}\end{array}\right]\)
- The fundamental matrix, \(\Phi(x)\) is a normalized version of \(\phi(x)\) satisfying \(\Phi(0)=I\) where \(I\) is t \(\Phi(x)=\phi(x) \cdot \frac{1}{\phi(0)}\)
- Substitute the value of \(\phi(x)\) and \(\phi(0)\)
\[
\Phi(x)=\left[\begin{array}{ccc}
\mathrm{e}^{x} & \mathrm{e}^{x}(x-1) & \frac{\mathrm{e}^{2 x}}{4} \\
\mathrm{e}^{x} & x \mathrm{e}^{x} & \frac{\mathrm{e}^{2 x}}{2} \\
\mathrm{e}^{x} & x \mathrm{e}^{x} & \mathrm{e}^{2 x}
\end{array}\right] \cdot \frac{1}{\left[\begin{array}{ccc}
1 & -1 & \frac{1}{4} \\
1 & 0 & \frac{1}{2} \\
1 & 0 & 1
\end{array}\right]}
\]
- Evaluate and simplify to get the fundamental matrix
\[
\Phi(x)=\left[\begin{array}{ccc}
-\mathrm{e}^{x}(x-1) & -\frac{\mathrm{e}^{2 x}}{2}+\frac{\mathrm{e}^{x}(3 x+1)}{2} & \frac{\mathrm{e}^{2 x}}{2}+\frac{\mathrm{e}^{x}(-x-1)}{2} \\
-x \mathrm{e}^{x} & 2 \mathrm{e}^{x}+\frac{3 x \mathrm{e}^{x}}{2}-\mathrm{e}^{2 x} & -\mathrm{e}^{x}-\frac{x \mathrm{e}^{x}}{2}+\mathrm{e}^{2 x} \\
-x \mathrm{e}^{x} & 2 \mathrm{e}^{x}+\frac{3 x \mathrm{e}^{x}}{2}-2 \mathrm{e}^{2 x} & -\mathrm{e}^{x}-\frac{x \mathrm{e}^{x}}{2}+2 \mathrm{e}^{2 x}
\end{array}\right]
\]

Find a particular solution of the system of ODEs using variation of parameters
- Let the particular solution be the fundamental matrix multiplied by \(\vec{v}(x)\) and solve for \(\vec{v}(x)\) \(\vec{y}_{p}(x)=\Phi(x) \cdot \vec{v}(x)\)
- Take the derivative of the particular solution
\(\vec{y}_{p}^{\prime}(x)=\Phi^{\prime}(x) \cdot \vec{v}(x)+\Phi(x) \cdot \vec{v}^{\prime}(x)\)
- Substitute particular solution and its derivative into the system of ODEs
\(\Phi^{\prime}(x) \cdot \vec{v}(x)+\Phi(x) \cdot \vec{v}^{\prime}(x)=A \cdot \Phi(x) \cdot \vec{v}(x)+\vec{f}(x)\)
- The fundamental matrix has columns that are solutions to the homogeneous system so its der \(A \cdot \Phi(x) \cdot \vec{v}(x)+\Phi(x) \cdot \vec{v}^{\prime}(x)=A \cdot \Phi(x) \cdot \vec{v}(x)+\vec{f}(x)\)
- Cancel like terms
\[
\Phi(x) \cdot \vec{v}^{\prime}(x)=\vec{f}(x)
\]
- Multiply by the inverse of the fundamental matrix
\(\vec{v}^{\prime}(x)=\frac{1}{\Phi(x)} \cdot \vec{f}(x)\)
- Integrate to solve for \(\vec{v}(x)\)
\(\vec{v}(x)=\int_{0}^{x} \frac{1}{\Phi(s)} \cdot \vec{f}(s) d s\)
- Plug \(\vec{v}(x)\) into the equation for the particular solution \(\vec{y}_{p}(x)=\Phi(x) \cdot\left(\int_{0}^{x} \frac{1}{\Phi(s)} \cdot \vec{f}(s) d s\right)\)
- Plug in the fundamental matrix and the forcing function and compute
\[
\vec{y}_{p}(x)=\left[\begin{array}{c}
-\frac{5 x \mathrm{e}^{x}}{2}+\mathrm{e}^{2 x}-2+\mathrm{e}^{x}-\frac{x}{2} \\
2 \mathrm{e}^{2 x}-\frac{1}{2}+\frac{(-5 x-3) \mathrm{e}^{x}}{2} \\
4 \mathrm{e}^{2 x}+\frac{(-5 x-3) \mathrm{e}^{x}}{2}-x-\frac{5}{2}
\end{array}\right]
\]
- Plug particular solution back into general solution
\[
\vec{y}(x)=c_{1} \vec{y}_{1}(x)+c_{2} \vec{y}_{2}(x)+c_{3} \vec{y}_{3}+\left[\begin{array}{c}
-\frac{5 x \mathrm{e}^{x}}{2}+\mathrm{e}^{2 x}-2+\mathrm{e}^{x}-\frac{x}{2} \\
2 \mathrm{e}^{2 x}-\frac{1}{2}+\frac{(-5 x-3) \mathrm{e}^{x}}{2} \\
4 \mathrm{e}^{2 x}+\frac{(-5 x-3) \mathrm{e}^{x}}{2}-x-\frac{5}{2}
\end{array}\right]
\]
- First component of the vector is the solution to the ODE
\[
y=-2+\frac{\left(c_{3}+4\right) \mathrm{e}^{2 x}}{4}+\frac{\left(2+\left(2 c_{2}-5\right) x+2 c_{1}-2 c_{2}\right) \mathrm{e}^{x}}{2}-\frac{x}{2}
\]

\section*{Maple trace}
```

`Methods for third order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 3; linear nonhomogeneous with symmetry [0,1] trying high order linear exact nonhomogeneous trying differential order: 3; missing the dependent variable checking if the LODE has constant coefficients <- constant coefficients successful`

```

Solution by Maple
Time used: 0.0 (sec). Leaf size: 23
```

dsolve(diff(y(x),x\$3)-4*diff(y(x),x\$2)+5*diff(y(x),x)-2*y(x)=2*x+3,y(x), singsol=all)

```
\[
y(x)=c_{2} \mathrm{e}^{2 x}+\left(c_{3} x+c_{1}\right) \mathrm{e}^{x}-x-4
\]

Solution by Mathematica
Time used: 0.006 (sec). Leaf size: 31
```

DSolve[y'''[x]-4*y''[x]+5*y'[x]-2*y[x]==2*x+3,y[x],x,IncludeSingularSolutions -> True]

```
\[
y(x) \rightarrow c_{1} e^{x}+x\left(-1+c_{2} e^{x}\right)+c_{3} e^{2 x}-4
\]

\subsection*{1.112 problem 159}

Internal problem ID [12529]
Internal file name [OUTPUT/11181_Tuesday_October_17_2023_07_20_33_AM_47053148/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 159.
ODE order: 4.
ODE degree: 1 .

The type(s) of ODE detected by this program : "higher_order_linear_constant_coefficients_ODE"

Maple gives the following as the ode type
```

[[_high_order, _linear, _nonhomogeneous]]

```
\[
y^{\prime \prime \prime \prime}-a^{4} y=5 a^{4} \mathrm{e}^{a x} \sin (a x)
\]

This is higher order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE And \(y_{p}\) is a particular solution to the nonhomogeneous ODE. \(y_{h}\) is the solution to
\[
y^{\prime \prime \prime \prime}-a^{4} y=0
\]

The characteristic equation is
\[
-a^{4}+\lambda^{4}=0
\]

The roots of the above equation are
\[
\begin{aligned}
\lambda_{1} & =a \\
\lambda_{2} & =-a \\
\lambda_{3} & =i a \\
\lambda_{4} & =-i a
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=\mathrm{e}^{i a x} c_{1}+\mathrm{e}^{-i a x} c_{2}+\mathrm{e}^{-a x} c_{3}+\mathrm{e}^{a x} c_{4}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=\mathrm{e}^{i a x} \\
& y_{2}=\mathrm{e}^{-i a x} \\
& y_{3}=\mathrm{e}^{-a x} \\
& y_{4}=\mathrm{e}^{a x}
\end{aligned}
\]

Now the particular solution to the given ODE is found
\[
y^{\prime \prime \prime \prime}-a^{4} y=5 a^{4} \mathrm{e}^{a x} \sin (a x)
\]

The particular solution is found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
5 a^{4} \mathrm{e}^{a x} \sin (a x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{a x} \cos (a x), \mathrm{e}^{a x} \sin (a x)\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{a x}, \mathrm{e}^{i a x}, \mathrm{e}^{-a x}, \mathrm{e}^{-i a x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{a x} \cos (a x)+A_{2} \mathrm{e}^{a x} \sin (a x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
\begin{aligned}
& -4 A_{1} a^{4} \mathrm{e}^{a x} \cos (a x)-4 A_{2} a^{4} \mathrm{e}^{a x} \sin (a x)-a^{4}\left(A_{1} \mathrm{e}^{a x} \cos (a x)+A_{2} \mathrm{e}^{a x} \sin (a x)\right) \\
& =5 a^{4} \mathrm{e}^{a x} \sin (a x)
\end{aligned}
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=0, A_{2}=-1\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=-\mathrm{e}^{a x} \sin (a x)
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(\mathrm{e}^{i a x} c_{1}+\mathrm{e}^{-i a x} c_{2}+\mathrm{e}^{-a x} c_{3}+\mathrm{e}^{a x} c_{4}\right)+\left(-\mathrm{e}^{a x} \sin (a x)\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{i a x} c_{1}+\mathrm{e}^{-i a x} c_{2}+\mathrm{e}^{-a x} c_{3}+\mathrm{e}^{a x} c_{4}-\mathrm{e}^{a x} \sin (a x) \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=\mathrm{e}^{i a x} c_{1}+\mathrm{e}^{-i a x} c_{2}+\mathrm{e}^{-a x} c_{3}+\mathrm{e}^{a x} c_{4}-\mathrm{e}^{a x} \sin (a x)
\]

Verified OK.
Maple trace
```

`Methods for high order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 4; linear nonhomogeneous with symmetry [0,1] trying high order linear exact nonhomogeneous trying differential order: 4; missing the dependent variable checking if the LODE has constant coefficients <- constant coefficients successful`

```

Solution by Maple
Time used: 0.0 (sec). Leaf size: 37
```

dsolve(diff(y(x),x\$4)-a^4*y(x)=5*a^4*exp(a*x)*\operatorname{sin}(a*x),y(x), singsol=all)

```
\[
y(x)=c_{4} \mathrm{e}^{-a x}+\left(c_{2}-\sin (a x)\right) \mathrm{e}^{a x}+c_{1} \cos (a x)+c_{3} \sin (a x)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.02 (sec). Leaf size: 45
DSolve[y''''[x]-a~4*y[x]==5*a^4*Exp[a*x]*Sin[a*x],y[x],x,IncludeSingularSolutions \(\rightarrow\) True]
\[
y(x) \rightarrow c_{2} e^{-a x}+c_{4} e^{a x}+c_{1} \cos (a x)+\left(-e^{a x}+c_{3}\right) \sin (a x)
\]

\subsection*{1.113 problem 160}

Internal problem ID [12530]
Internal file name [OUTPUT/11182_Tuesday_October_17_2023_07_20_33_AM_31298264/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 160.
ODE order: 4.
ODE degree: 1.

The type(s) of ODE detected by this program : "higher_order_linear_constant_coefficients_ODE"

Maple gives the following as the ode type
[[_high_order, _linear, _nonhomogeneous]]
\[
y^{\prime \prime \prime \prime}+2 a^{2} y^{\prime \prime}+a^{4} y=8 \cos (a x)
\]

This is higher order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE And \(y_{p}\) is a particular solution to the nonhomogeneous ODE. \(y_{h}\) is the solution to
\[
y^{\prime \prime \prime \prime}+2 a^{2} y^{\prime \prime}+a^{4} y=0
\]

The characteristic equation is
\[
a^{4}+2 a^{2} \lambda^{2}+\lambda^{4}=0
\]

The roots of the above equation are
\[
\begin{aligned}
\lambda_{1} & =i a \\
\lambda_{2} & =-i a \\
\lambda_{3} & =i a \\
\lambda_{4} & =-i a
\end{aligned}
\]

Therefore the homogeneous solution is
\[
y_{h}(x)=\mathrm{e}^{i a x} c_{1}+x \mathrm{e}^{i a x} c_{2}+\mathrm{e}^{-i a x} c_{3}+x \mathrm{e}^{-i a x} c_{4}
\]

The fundamental set of solutions for the homogeneous solution are the following
\[
\begin{aligned}
& y_{1}=\mathrm{e}^{i a x} \\
& y_{2}=x \mathrm{e}^{i a x} \\
& y_{3}=\mathrm{e}^{-i a x} \\
& y_{4}=x \mathrm{e}^{-i a x}
\end{aligned}
\]

Now the particular solution to the given ODE is found
\[
y^{\prime \prime \prime \prime}+2 a^{2} y^{\prime \prime}+a^{4} y=8 \cos (a x)
\]

Let the particular solution be
\[
y_{p}=U_{1} y_{1}+U_{2} y_{2}+U_{3} y_{3}+U_{4} y_{4}
\]

Where \(y_{i}\) are the basis solutions found above for the homogeneous solution \(y_{h}\) and \(U_{i}(x)\) are functions to be determined as follows
\[
U_{i}=(-1)^{n-i} \int \frac{F(x) W_{i}(x)}{a W(x)} d x
\]

Where \(W(x)\) is the Wronskian and \(W_{i}(x)\) is the Wronskian that results after deleting the last row and the \(i\)-th column of the determinant and \(n\) is the order of the ODE or equivalently, the number of basis solutions, and \(a\) is the coefficient of the leading derivative in the ODE, and \(F(x)\) is the RHS of the ODE. Therefore, the first step is to find the Wronskian \(W(x)\). This is given by
\[
W(x)=\left|\begin{array}{cccc}
y_{1} & y_{2} & y_{3} & y_{4} \\
y_{1}^{\prime} & y_{2}^{\prime} & y_{3}^{\prime} & y_{4}^{\prime} \\
y_{1}^{\prime \prime} & y_{2}^{\prime \prime} & y_{3}^{\prime \prime} & y_{4}^{\prime \prime} \\
y_{1}^{\prime \prime \prime} & y_{2}^{\prime \prime \prime} & y_{3}^{\prime \prime \prime} & y_{4}^{\prime \prime \prime}
\end{array}\right|
\]

Substituting the fundamental set of solutions \(y_{i}\) found above in the Wronskian gives
\[
\begin{aligned}
& W=\left[\begin{array}{cccc}
\mathrm{e}^{i a x} & x \mathrm{e}^{i a x} & \mathrm{e}^{-i a x} & x \mathrm{e}^{-i a x} \\
i a \mathrm{e}^{i a x} & \mathrm{e}^{i a x}(i a x+1) & -i a \mathrm{e}^{-i a x} & \mathrm{e}^{-i a x}(-i a x+1) \\
-a^{2} \mathrm{e}^{i a x} & \mathrm{e}^{i a x} a(-a x+2 i) & -a^{2} \mathrm{e}^{-i a x} & -2\left(\frac{a x}{2}+i\right) \mathrm{e}^{-i a x} a \\
-i a^{3} \mathrm{e}^{i a x} & -\mathrm{e}^{i a x} a^{2}(i a x+3) & i a^{3} \mathrm{e}^{-i a x} & \mathrm{e}^{-i a x} a^{2}(i a x-3)
\end{array}\right] \\
&|W|=16 \mathrm{e}^{2 i a x} \mathrm{e}^{-2 i a x} a^{4}
\end{aligned}
\]

The determinant simplifies to
\[
|W|=16 a^{4}
\]

Now we determine \(W_{i}\) for each \(U_{i}\).
\[
\begin{aligned}
& W_{1}(x)=\operatorname{det}\left[\begin{array}{ccc}
x \mathrm{e}^{i a x} & \mathrm{e}^{-i a x} & x \mathrm{e}^{-i a x} \\
\mathrm{e}^{i a x}(i a x+1) & -i a \mathrm{e}^{-i a x} & \mathrm{e}^{-i a x}(-i a x+1) \\
\mathrm{e}^{i a x} a(-a x+2 i) & -a^{2} \mathrm{e}^{-i a x} & -2\left(\frac{a x}{2}+i\right) \mathrm{e}^{-i a x} a
\end{array}\right] \\
& =-4 \mathrm{e}^{-i a x} a(a x-i) \\
& \begin{aligned}
& W_{2}(x)=\operatorname{det}\left[\begin{array}{ccc}
\mathrm{e}^{i a x} & \mathrm{e}^{-i a x} & x \mathrm{e}^{-i a x} \\
i a \mathrm{e}^{i a x} & -i a \mathrm{e}^{-i a x} & \mathrm{e}^{-i a x}(-i a x+1) \\
-a^{2} \mathrm{e}^{i a x} & -a^{2} \mathrm{e}^{-i a x} & -2\left(\frac{a x}{2}+i\right) \mathrm{e}^{-i a x} a
\end{array}\right] \\
&=-4 a^{2} \mathrm{e}^{-i a x} \\
&\left.\left.\begin{array}{rl}
W_{3}(x) & =\operatorname{det}\left[\begin{array}{ccc}
\mathrm{e}^{i a x} & x \mathrm{e}^{i a x} & x \mathrm{e}^{-i a x} \\
i a \mathrm{e}^{i a x} & \mathrm{e}^{i a x}(i a x+1) & \mathrm{e}^{-i a x}(-i a x+1) \\
-a^{2} \mathrm{e}^{i a x} & \mathrm{e}^{i a x} a(-a x+2 i) & -2\left(\frac{a x}{2}+i\right) \mathrm{e}^{-i a x} a
\end{array}\right] \\
=-4 \mathrm{e}^{i a x} a(a x+i) & x \mathrm{e}^{i a x}
\end{array}\right] \begin{array}{c}
\mathrm{e}^{-i a x} \\
W_{4}(x)
\end{array}\right] \\
&=\operatorname{det}\left[\begin{array}{ccc}
\mathrm{e}^{i a x} & \mathrm{e}^{i a x}(i a x+1) & -i a \mathrm{e}^{-i a x} \\
i a \mathrm{e}^{i a x} \\
-a^{2} \mathrm{e}^{i a x} & \mathrm{e}^{i a x} a(-a x+2 i) & -a^{2} \mathrm{e}^{-i a x}
\end{array}\right]
\end{aligned}
\end{aligned}
\]

Now we are ready to evaluate each \(U_{i}(x)\).
\[
\begin{aligned}
U_{1} & =(-1)^{4-1} \int \frac{F(x) W_{1}(x)}{a W(x)} d x \\
& =(-1)^{3} \int \frac{(8 \cos (a x))\left(-4 \mathrm{e}^{-i a x} a(a x-i)\right)}{(1)\left(16 a^{4}\right)} d x \\
& =-\int \frac{-32 \cos (a x) \mathrm{e}^{-i a x} a(a x-i)}{16 a^{4}} d x \\
& =-\int\left(-\frac{2 \cos (a x) \mathrm{e}^{-i a x}(a x-i)}{a^{3}}\right) d x \\
& =-\left(\int-\frac{2 \cos (a x) \mathrm{e}^{-i a x}(a x-i)}{a^{3}} d x\right)
\end{aligned}
\]
\[
\begin{aligned}
& U_{2}=(-1)^{4-2} \int \frac{F(x) W_{2}(x)}{a W(x)} d x \\
& =(-1)^{2} \int \frac{(8 \cos (a x))\left(-4 a^{2} \mathrm{e}^{-i a x}\right)}{(1)\left(16 a^{4}\right)} d x \\
& =\int \frac{-32 \cos (a x) a^{2} \mathrm{e}^{-i a x}}{16 a^{4}} d x \\
& =\int\left(-\frac{2 \cos (a x) \mathrm{e}^{-i a x}}{a^{2}}\right) d x \\
& =\frac{\frac{x \mathrm{e}^{-i a x} \tan \left(\frac{a x}{2}\right)^{2}}{a}-\frac{x \mathrm{e}^{-i a x}}{a}-\frac{2 i x \mathrm{e}^{-i a x} \tan \left(\frac{a x}{2}\right)}{a}-\frac{2 \mathrm{e}^{-i a x} \tan \left(\frac{a x}{2}\right)}{a^{2}}}{a\left(1+\tan \left(\frac{a x}{2}\right)^{2}\right)} \\
& U_{3}=(-1)^{4-3} \int \frac{F(x) W_{3}(x)}{a W(x)} d x \\
& =(-1)^{1} \int \frac{(8 \cos (a x))\left(-4 \mathrm{e}^{i a x} a(a x+i)\right)}{(1)\left(16 a^{4}\right)} d x \\
& =-\int \frac{-32 \cos (a x) \mathrm{e}^{i a x} a(a x+i)}{16 a^{4}} d x \\
& =-\int\left(-\frac{2 \cos (a x) \mathrm{e}^{i a x}(a x+i)}{a^{3}}\right) d x \\
& =-\left(\int-\frac{2 \cos (a x) \mathrm{e}^{i a x}(a x+i)}{a^{3}} d x\right) \\
& U_{4}=(-1)^{4-4} \int \frac{F(x) W_{4}(x)}{a W(x)} d x \\
& =(-1)^{0} \int \frac{(8 \cos (a x))\left(-4 a^{2} \mathrm{e}^{i a x}\right)}{(1)\left(16 a^{4}\right)} d x \\
& =\int \frac{-32 \cos (a x) a^{2} \mathrm{e}^{i a x}}{16 a^{4}} d x \\
& =\int\left(-\frac{2 \cos (a x) \mathrm{e}^{i a x}}{a^{2}}\right) d x \\
& =\frac{\frac{x \mathrm{e}^{i a x} \tan \left(\frac{a x}{2}\right)^{2}}{a}-\frac{x \mathrm{e}^{i a x}}{a}+\frac{2 i x \mathrm{e}^{i a x} \tan \left(\frac{a x}{2}\right)}{a}-\frac{2 \mathrm{e}^{i a x} \tan \left(\frac{a x}{2}\right)}{a^{2}}}{a\left(1+\tan \left(\frac{a x}{2}\right)^{2}\right)}
\end{aligned}
\]

Now that all the \(U_{i}\) functions have been determined, the particular solution is found from
\[
y_{p}=U_{1} y_{1}+U_{2} y_{2}+U_{3} y_{3}+U_{4} y_{4}
\]

Hence
\[
\begin{aligned}
y_{p} & =\left(-\left(\int-\frac{2 \cos (a x) \mathrm{e}^{-i a x}(a x-i)}{a^{3}} d x\right)\right)\left(\mathrm{e}^{i a x}\right) \\
& +\left(\frac{\frac{x \mathrm{e}^{-i a x} \tan \left(\frac{a x}{2}\right)^{2}}{a}-\frac{x \mathrm{e}^{-i a x}}{a}-\frac{2 i x \mathrm{e}^{-i a x} \tan \left(\frac{a x}{2}\right)}{a}-\frac{2 \mathrm{e}^{-i a x} \tan \left(\frac{a x}{2}\right)}{a^{2}}}{a\left(1+\tan \left(\frac{a x}{2}\right)^{2}\right)}\right)\left(x \mathrm{e}^{i a x}\right) \\
& +\left(-\left(\int-\frac{2 \cos (a x) \mathrm{e}^{i a x}(a x+i)}{a^{3}} d x\right)\right)\left(\mathrm{e}^{-i a x}\right) \\
& +\left(\frac{\frac{x \mathrm{e}^{i a x} \tan \left(\frac{a x}{2}\right)^{2}}{a}-\frac{x \mathrm{e}^{i a x}}{a}+\frac{2 i x \mathrm{e}^{i a x} \tan \left(\frac{a x}{2}\right)}{a}-\frac{2 \mathrm{e}^{i a x} \tan \left(\frac{a x}{2}\right)}{a^{2}}}{a\left(1+\tan \left(\frac{a x}{2}\right)^{2}\right)}\right)\left(x \mathrm{e}^{-i a x}\right)
\end{aligned}
\]

Therefore the particular solution is
\[
y_{p}=\frac{-2 a x^{2} \cos (a x)-2\left(\int \cos (a x) \mathrm{e}^{-i a x}(-a x+i) d x\right) \mathrm{e}^{i a x}-2 x \sin (a x)+2\left(\int \cos (a x) \mathrm{e}^{i a x}(a x+i) d x\right) \mathrm{e}}{a^{3}}
\]

Which simplifies to
\[
y_{p}=\frac{-2 a x^{2} \cos (a x)+4\left(\int \cos (a x)(\cos (a x) a x-\sin (a x)) d x\right) \cos (a x)+4\left(\int \cos (a x)(\sin (a x) a x+\cos \right.}{a^{3}}
\]

Therefore the general solution is
\[
\begin{aligned}
y= & y_{h}+y_{p} \\
= & \left(\mathrm{e}^{i a x} c_{1}+x \mathrm{e}^{i a x} c_{2}+\mathrm{e}^{-i a x} c_{3}+x \mathrm{e}^{-i a x} c_{4}\right) \\
& +\left(\frac{-2 a x^{2} \cos (a x)+4\left(\int \cos (a x)(\cos (a x) a x-\sin (a x)) d x\right) \cos (a x)+4\left(\int \cos (a x)(\sin (a x) a x+\mathrm{c}\right.}{a^{3}}\right.
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
y & =\left(c_{4} x+c_{3}\right) \mathrm{e}^{-i a x}+\left(c_{2} x+c_{1}\right) \mathrm{e}^{i a x} \\
& +\frac{-2 a x^{2} \cos (a x)+4\left(\int \cos (a x)(\cos (a x) a x-\sin (a x)) d x\right) \cos (a x)+4\left(\int \cos (a x)(\sin (a x) a x+\cos ( \right.}{a^{3}}
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{aligned}
y & =\left(c_{4} x+c_{3}\right) \mathrm{e}^{-i a x}+\left(c_{2} x+c_{1}\right) \mathrm{e}^{i a x} \\
& +\frac{-2 a x^{2} \cos (a x)+4\left(\int \cos (a x)(\cos (a x) a x-\sin (a x)) d x\right) \cos (a x)+4\left(\int \cos (a x)(\sin (a x) a x+\cos ( \right.}{a^{3}}
\end{aligned}
\]

\section*{Verification of solutions}
\[
\begin{aligned}
y & =\left(c_{4} x+c_{3}\right) \mathrm{e}^{-i a x}+\left(c_{2} x+c_{1}\right) \mathrm{e}^{i a x} \\
& +\frac{-2 a x^{2} \cos (a x)+4\left(\int \cos (a x)(\cos (a x) a x-\sin (a x)) d x\right) \cos (a x)+4\left(\int \cos (a x)(\sin (a x) a x+\cos \right.}{a^{3}}
\end{aligned}
\]

Verified OK.
Maple trace
```

`Methods for high order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 4; linear nonhomogeneous with symmetry [0,1] trying high order linear exact nonhomogeneous trying differential order: 4; missing the dependent variable checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.016 (sec). Leaf size: 52
```

dsolve(diff(y(x),x\$4)+2*a^2*diff(y(x),x\$2)+a^4*y(x)=8*\operatorname{cos(a*x),y(x), singsol=all)}

```
\[
y(x)=\frac{\left(2+\left(c_{3} x+c_{1}\right) a^{4}-a^{2} x^{2}\right) \cos (a x)+\sin (a x)\left(\left(c_{4} x+c_{2}\right) a^{3}+3 x\right) a}{a^{4}}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.187 (sec). Leaf size: 64
DSolve[y''''[x]+2*a^2*y''[x]+a^4*y[x]==8*Cos[a*x],y[x],x,IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow \frac{2 a\left(x\left(2+a^{3} c_{4}\right)+a^{3} c_{3}\right) \sin (a x)+\left(2 a^{4}\left(c_{2} x+c_{1}\right)-2 a^{2} x^{2}+5\right) \cos (a x)}{2 a^{4}}
\]

\subsection*{1.114 problem 162}
1.114.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 1160
1.114.2 Solving as second order linear constant coeff ode . . . . . . . . 1161
1.114.3 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1163
1.114.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1167

Internal problem ID [12531]
Internal file name [OUTPUT/11183_Tuesday_October_17_2023_07_20_34_AM_90181001/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 162.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second__order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _missing_x]]
\[
y^{\prime \prime}+2 h y^{\prime}+y n^{2}=0
\]

With initial conditions
\[
\left[y(0)=a, y^{\prime}(0)=c\right]
\]

\subsection*{1.114.1 Existence and uniqueness analysis}

This is a linear ODE. In canonical form it is written as
\[
y^{\prime \prime}+p(x) y^{\prime}+q(x) y=F
\]

Where here
\[
\begin{aligned}
p(x) & =2 h \\
q(x) & =n^{2} \\
F & =0
\end{aligned}
\]

Hence the ode is
\[
y^{\prime \prime}+2 h y^{\prime}+y n^{2}=0
\]

The domain of \(p(x)=2 h\) is
\[
\{-\infty<x<\infty\}
\]

And the point \(x_{0}=0\) is inside this domain. The domain of \(q(x)=n^{2}\) is
\[
\{-\infty<x<\infty\}
\]

And the point \(x_{0}=0\) is also inside this domain. Hence solution exists and is unique.

\subsection*{1.114.2 Solving as second order linear constant coeff ode}

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=2 h, C=n^{2}\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+2 h \lambda \mathrm{e}^{\lambda x}+n^{2} \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
2 h \lambda+\lambda^{2}+n^{2}=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=2 h, C=n^{2}\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{-2 h}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{2 h^{2}-(4)(1)\left(n^{2}\right)} \\
& =-h \pm \sqrt{h^{2}-n^{2}}
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=-h+\sqrt{h^{2}-n^{2}} \\
& \lambda_{2}=-h-\sqrt{h^{2}-n^{2}}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=-h+\sqrt{h^{2}-n^{2}} \\
& \lambda_{2}=-h-\sqrt{h^{2}-n^{2}}
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}+c_{2} e^{\left(-h-\sqrt{h^{2}-n^{2}}\right) x}
\end{aligned}
\]

Or
\[
y=c_{1} \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}+c_{2} \mathrm{e}^{\left(-h-\sqrt{h^{2}-n^{2}}\right) x}
\]

Initial conditions are used to solve for the constants of integration.
Looking at the above solution
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}+c_{2} \mathrm{e}^{\left(-h-\sqrt{h^{2}-n^{2}}\right) x} \tag{1}
\end{equation*}
\]

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting \(y=a\) and \(x=0\) in the above gives
\[
\begin{equation*}
a=c_{1}+c_{2} \tag{1~A}
\end{equation*}
\]

Taking derivative of the solution gives
\[
y^{\prime}=c_{1}\left(-h+\sqrt{h^{2}-n^{2}}\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}+c_{2}\left(-h-\sqrt{h^{2}-n^{2}}\right) \mathrm{e}^{\left(-h-\sqrt{h^{2}-n^{2}}\right) x}
\]
substituting \(y^{\prime}=c\) and \(x=0\) in the above gives
\[
\begin{equation*}
c=\left(c_{1}-c_{2}\right) \sqrt{h^{2}-n^{2}}-h\left(c_{1}+c_{2}\right) \tag{2~A}
\end{equation*}
\]

Equations \(\{1 \mathrm{~A}, 2 \mathrm{~A}\}\) are now solved for \(\left\{c_{1}, c_{2}\right\}\). Solving for the constants gives
\[
\begin{aligned}
& c_{1}=\frac{\sqrt{h^{2}-n^{2}} a+h a+c}{2 \sqrt{h^{2}-n^{2}}} \\
& c_{2}=\frac{\sqrt{h^{2}-n^{2}} a-h a-c}{2 \sqrt{h^{2}-n^{2}}}
\end{aligned}
\]

Substituting these values back in above solution results in
\(y=\frac{\sqrt{h^{2}-n^{2}} \mathrm{e}^{-\left(h-\sqrt{h^{2}-n^{2}}\right) x} a+\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x} \sqrt{h^{2}-n^{2}} a+\mathrm{e}^{-\left(h-\sqrt{h^{2}-n^{2}}\right) x} a h-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x} h a+\mathrm{e}^{-(h-}}{2 \sqrt{h^{2}-n^{2}}}\)
Which simplifies to
\[
y=\frac{\left(\sqrt{h^{2}-n^{2}} a+h a+c\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}\left(-\sqrt{h^{2}-n^{2}} a+h a+c\right)}{2 \sqrt{h^{2}-n^{2}}}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=\frac{\left(\sqrt{h^{2}-n^{2}} a+h a+c\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}\left(-\sqrt{h^{2}-n^{2}} a+h a+c\right)}{2 \sqrt{h^{2}-n^{2}}} \tag{1}
\end{equation*}
\]

\section*{Verification of solutions}
\[
y=\frac{\left(\sqrt{h^{2}-n^{2}} a+h a+c\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}\left(-\sqrt{h^{2}-n^{2}} a+h a+c\right)}{2 \sqrt{h^{2}-n^{2}}}
\]

\section*{Verified OK.}

\subsection*{1.114.3 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+2 h y^{\prime}+y n^{2} & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=2 h  \tag{3}\\
& C=n^{2}
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{h^{2}-n^{2}}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=h^{2}-n^{2} \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=\left(h^{2}-n^{2}\right) z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 177: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=h^{2}-n^{2}\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{x \sqrt{h^{2}-n^{2}}}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{12 h}{2} \frac{1}{1} d x} \\
& =z_{1} e^{-x h} \\
& =z_{1}\left(\mathrm{e}^{-x h}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{2 h}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{-2 x h}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(-\frac{\mathrm{e}^{-2 x \sqrt{h^{2}-n^{2}}}}{2 \sqrt{h^{2}-n^{2}}}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}\right)+c_{2}\left(\mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}\left(-\frac{\mathrm{e}^{-2 x \sqrt{h^{2}-n^{2}}}}{2 \sqrt{h^{2}-n^{2}}}\right)\right)
\end{aligned}
\]

Initial conditions are used to solve for the constants of integration.
Looking at the above solution
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\frac{c_{2} \mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}}{2 \sqrt{h^{2}-n^{2}}} \tag{1}
\end{equation*}
\]

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting \(y=a\) and \(x=0\) in the above gives
\[
\begin{equation*}
a=\frac{2 c_{1} \sqrt{h^{2}-n^{2}}-c_{2}}{2 \sqrt{h^{2}-n^{2}}} \tag{1~A}
\end{equation*}
\]

Taking derivative of the solution gives
\[
y^{\prime}=c_{1}\left(-h+\sqrt{h^{2}-n^{2}}\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\frac{c_{2}\left(-h-\sqrt{h^{2}-n^{2}}\right) \mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}}{2 \sqrt{h^{2}-n^{2}}}
\]
substituting \(y^{\prime}=c\) and \(x=0\) in the above gives
\[
\begin{equation*}
c=\frac{-2 c_{1} h \sqrt{h^{2}-n^{2}}+2 c_{1} h^{2}-2 c_{1} n^{2}+c_{2} \sqrt{h^{2}-n^{2}}+c_{2} h}{2 \sqrt{h^{2}-n^{2}}} \tag{2~A}
\end{equation*}
\]

Equations \(\{1 \mathrm{~A}, 2 \mathrm{~A}\}\) are now solved for \(\left\{c_{1}, c_{2}\right\}\). Solving for the constants gives
\[
\begin{aligned}
& c_{1}=\frac{\sqrt{h^{2}-n^{2}} a+h a+c}{2 \sqrt{h^{2}-n^{2}}} \\
& c_{2}=-\sqrt{h^{2}-n^{2}} a+h a+c
\end{aligned}
\]

Substituting these values back in above solution results in
\(y=\frac{\sqrt{h^{2}-n^{2}} \mathrm{e}^{-\left(h-\sqrt{h^{2}-n^{2}}\right) x} a+\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x} \sqrt{h^{2}-n^{2}} a+\mathrm{e}^{-\left(h-\sqrt{h^{2}-n^{2}}\right) x} a h-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x} h a+\mathrm{e}^{-(h-}}{2 \sqrt{h^{2}-n^{2}}}\)

Which simplifies to
\[
y=\frac{\left(\sqrt{h^{2}-n^{2}} a+h a+c\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}\left(-\sqrt{h^{2}-n^{2}} a+h a+c\right)}{2 \sqrt{h^{2}-n^{2}}}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{\left(\sqrt{h^{2}-n^{2}} a+h a+c\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}\left(-\sqrt{h^{2}-n^{2}} a+h a+c\right)}{2 \sqrt{h^{2}-n^{2}}} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=\frac{\left(\sqrt{h^{2}-n^{2}} a+h a+c\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}\left(-\sqrt{h^{2}-n^{2}} a+h a+c\right)}{2 \sqrt{h^{2}-n^{2}}}
\]

Verified OK.

\subsection*{1.114.4 Maple step by step solution}

Let's solve
\(\left[y^{\prime \prime}+2 h y^{\prime}+y n^{2}=0, y(0)=a,\left.y^{\prime}\right|_{\{x=0\}}=c\right]\)
- Highest derivative means the order of the ODE is 2
\(y^{\prime \prime}\)
- Characteristic polynomial of ODE
\(2 h r+n^{2}+r^{2}=0\)
- Use quadratic formula to solve for \(r\)
\[
r=\frac{(-2 h) \pm\left(\sqrt{4 h^{2}-4 n^{2}}\right)}{2}
\]
- Roots of the characteristic polynomial
\[
r=\left(-h-\sqrt{h^{2}-n^{2}},-h+\sqrt{h^{2}-n^{2}}\right)
\]
- \(\quad 1\) st solution of the ODE
\[
y_{1}(x)=\mathrm{e}^{\left(-h-\sqrt{h^{2}-n^{2}}\right) x}
\]
- \(\quad 2 n d\) solution of the ODE
\[
y_{2}(x)=\mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}
\]
- General solution of the ODE
\[
y=c_{1} y_{1}(x)+c_{2} y_{2}(x)
\]
- \(\quad\) Substitute in solutions
\[
y=c_{1} \mathrm{e}^{\left(-h-\sqrt{h^{2}-n^{2}}\right) x}+c_{2} \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}
\]

Check validity of solution \(y=c_{1} \mathrm{e}^{\left(-h-\sqrt{h^{2}-n^{2}}\right) x}+c_{2} \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}\)
- Use initial condition \(y(0)=a\)
\[
a=c_{1}+c_{2}
\]
- Compute derivative of the solution
\[
y^{\prime}=c_{1}\left(-h-\sqrt{h^{2}-n^{2}}\right) \mathrm{e}^{\left(-h-\sqrt{h^{2}-n^{2}}\right) x}+c_{2}\left(-h+\sqrt{h^{2}-n^{2}}\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}
\]
- Use the initial condition \(\left.y^{\prime}\right|_{\{x=0\}}=c\)
\(c=c_{1}\left(-h-\sqrt{h^{2}-n^{2}}\right)+c_{2}\left(-h+\sqrt{h^{2}-n^{2}}\right)\)
- Solve for \(c_{1}\) and \(c_{2}\)
\(\left\{c_{1}=-\frac{-\sqrt{h^{2}-n^{2}} a+h a+c}{2 \sqrt{h^{2}-n^{2}}}, c_{2}=\frac{\sqrt{h^{2}-n^{2}} a+h a+c}{2 \sqrt{h^{2}-n^{2}}}\right\}\)
- Substitute constant values into general solution and simplify
\[
y=\frac{\left(\sqrt{h^{2}-n^{2}} a+h a+c\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}\left(-\sqrt{h^{2}-n^{2}} a+h a+c\right)}{2 \sqrt{h^{2}-n^{2}}}
\]
- \(\quad\) Solution to the IVP
\[
y=\frac{\left(\sqrt{h^{2}-n^{2}} a+h a+c\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}\left(-\sqrt{h^{2}-n^{2}} a+h a+c\right)}{2 \sqrt{h^{2}-n^{2}}}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.109 (sec). Leaf size: 93
dsolve([diff \(\left.(y(x), x \$ 2)+2 * h * \operatorname{diff}(y(x), x)+n^{\wedge} 2 * y(x)=0, y(0)=a, D(y)(0)=c\right], y(x)\), singsol=all)
\(y(x)=\frac{\left(\sqrt{h^{2}-n^{2}} a+h a+c\right) \mathrm{e}^{\left(-h+\sqrt{h^{2}-n^{2}}\right) x}-\mathrm{e}^{-\left(h+\sqrt{h^{2}-n^{2}}\right) x}\left(-\sqrt{h^{2}-n^{2}} a+h a+c\right)}{2 \sqrt{h^{2}-n^{2}}}\)
\(\checkmark\) Solution by Mathematica
Time used: 0.067 (sec). Leaf size: 123
```

DSolve[{y''[x]+2*h*y'[x]+n^2*y[x]==0,{y[0]==a, y'[0]==c}},y[x],x,IncludeSingularSolutions ->

```
\(y(x)\)
\(\rightarrow \frac{e^{-\left(x\left(\sqrt{h^{2}-n^{2}}+h\right)\right)}\left(a h\left(e^{2 x \sqrt{h^{2}-n^{2}}}-1\right)+a \sqrt{h^{2}-n^{2}}\left(e^{2 x \sqrt{h^{2}-n^{2}}}+1\right)+c\left(e^{2 x \sqrt{h^{2}-n^{2}}}-1\right)\right)}{2 \sqrt{h^{2}-n^{2}}}\)

\subsection*{1.115 problem 163}
1.115.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 1170
1.115.2 Solving as second order linear constant coeff ode . . . . . . . . 1171
1.115.3 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1175
1.115.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1180

Internal problem ID [12532]
Internal file name [OUTPUT/11184_Tuesday_October_17_2023_07_20_36_AM_96363/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 163.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second__order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _linear, _nonhomogeneous]]
\[
y^{\prime \prime}+y n^{2}=h \sin (r x)
\]

With initial conditions
\[
\left[y(0)=a, y^{\prime}(0)=c\right]
\]

\subsection*{1.115.1 Existence and uniqueness analysis}

This is a linear ODE. In canonical form it is written as
\[
y^{\prime \prime}+p(x) y^{\prime}+q(x) y=F
\]

Where here
\[
\begin{aligned}
p(x) & =0 \\
q(x) & =n^{2} \\
F & =h \sin (r x)
\end{aligned}
\]

Hence the ode is
\[
y^{\prime \prime}+y n^{2}=h \sin (r x)
\]

The domain of \(p(x)=0\) is
\[
\{-\infty<x<\infty\}
\]

And the point \(x_{0}=0\) is inside this domain. The domain of \(q(x)=n^{2}\) is
\[
\{-\infty<x<\infty\}
\]

And the point \(x_{0}=0\) is also inside this domain. The domain of \(F=h \sin (r x)\) is
\[
\{-\infty<x<\infty\}
\]

And the point \(x_{0}=0\) is also inside this domain. Hence solution exists and is unique.

\subsection*{1.115.2 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=0, C=n^{2}, f(x)=h \sin (r x)\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+y n^{2}=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=n^{2}\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+n^{2} \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+n^{2}=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=n^{2}\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)\left(n^{2}\right)} \\
& = \pm \sqrt{-n^{2}}
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=+\sqrt{-n^{2}} \\
& \lambda_{2}=-\sqrt{-n^{2}}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=\sqrt{-n^{2}} \\
& \lambda_{2}=-\sqrt{-n^{2}}
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{\left(\sqrt{-n^{2}}\right) x}+c_{2} e^{\left(-\sqrt{-n^{2}}\right) x}
\end{aligned}
\]

Or
\[
y=c_{1} \mathrm{e}^{\sqrt{-n^{2}} x}+c_{2} \mathrm{e}^{-\sqrt{-n^{2}} x}
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \mathrm{e}^{\sqrt{-n^{2}} x}+c_{2} \mathrm{e}^{-\sqrt{-n^{2}} x}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
h \sin (r x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{\cos (r x), \sin (r x)\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{\sqrt{-n^{2}} x}, \mathrm{e}^{-\sqrt{-n^{2}} x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \cos (r x)+A_{2} \sin (r x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-A_{1} r^{2} \cos (r x)-A_{2} r^{2} \sin (r x)+\left(A_{1} \cos (r x)+A_{2} \sin (r x)\right) n^{2}=h \sin (r x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=0, A_{2}=\frac{h}{n^{2}-r^{2}}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{h \sin (r x)}{n^{2}-r^{2}}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{\sqrt{-n^{2}} x}+c_{2} \mathrm{e}^{-\sqrt{-n^{2}} x}\right)+\left(\frac{h \sin (r x)}{n^{2}-r^{2}}\right)
\end{aligned}
\]

Initial conditions are used to solve for the constants of integration.
Looking at the above solution
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{\sqrt{-n^{2}} x}+c_{2} \mathrm{e}^{-\sqrt{-n^{2}} x}+\frac{h \sin (r x)}{n^{2}-r^{2}} \tag{1}
\end{equation*}
\]

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting \(y=a\) and \(x=0\) in the above gives
\[
\begin{equation*}
a=c_{1}+c_{2} \tag{1~A}
\end{equation*}
\]

Taking derivative of the solution gives
\[
y^{\prime}=c_{1} \sqrt{-n^{2}} \mathrm{e}^{\sqrt{-n^{2}} x}-c_{2} \sqrt{-n^{2}} \mathrm{e}^{-\sqrt{-n^{2}} x}+\frac{h r \cos (r x)}{n^{2}-r^{2}}
\]
substituting \(y^{\prime}=c\) and \(x=0\) in the above gives
\[
\begin{equation*}
c=\frac{(n+r)(n-r)\left(c_{1}-c_{2}\right) \sqrt{-n^{2}}+h r}{n^{2}-r^{2}} \tag{2~A}
\end{equation*}
\]

Equations \(\{1 \mathrm{~A}, 2 \mathrm{~A}\}\) are now solved for \(\left\{c_{1}, c_{2}\right\}\). Solving for the constants gives
\[
\begin{aligned}
& c_{1}=-\frac{-n^{4} a+n^{2} a r^{2}+\sqrt{-n^{2}} c n^{2}-\sqrt{-n^{2}} c r^{2}-\sqrt{-n^{2}} h r}{2 n^{2}\left(n^{2}-r^{2}\right)} \\
& c_{2}=-\frac{\sqrt{-n^{2}}\left(\sqrt{-n^{2}} n^{2} a-\sqrt{-n^{2}} r^{2} a-c n^{2}+c r^{2}+h r\right)}{2 n^{2}\left(n^{2}-r^{2}\right)}
\end{aligned}
\]

Substituting these values back in above solution results in
\(y=\frac{\mathrm{e}^{\sqrt{-n^{2}} x} a n^{4}-\mathrm{e}^{\sqrt{-n^{2}} x} a n^{2} r^{2}+\mathrm{e}^{-\sqrt{-n^{2}} x} a n^{4}-\mathrm{e}^{-\sqrt{-n^{2}} x} a n^{2} r^{2}-\mathrm{e}^{\sqrt{-n^{2}} x} \sqrt{-n^{2}} c n^{2}+\mathrm{e}^{\sqrt{-n^{2}} x} \sqrt{-n^{2}} c r^{2}}{2 n^{4}-2}\)
Which simplifies to
\(y\)
\(=\frac{\mathrm{e}^{-\sqrt{-n^{2}} x}\left(\left(\left(-c n^{2}+r(c r+h)\right) \sqrt{-n^{2}}+n^{4} a-n^{2} a r^{2}\right) \mathrm{e}^{2 \sqrt{-n^{2}} x}+2 \sin (r x) n^{2} h \mathrm{e}^{\sqrt{-n^{2}} x}+\left(c n^{2}-c r^{2}-l\right.\right.}{2 n^{4}-2 n^{2} r^{2}}\)
Summary
The solution(s) found are the following
\(y\)
\(=\frac{\mathrm{e}^{-\sqrt{-n^{2}} x}\left(\left(\left(-c n^{2}+r(c r+h)\right) \sqrt{-n^{2}}+n^{4} a-n^{2} a r^{2}\right) \mathrm{e}^{2 \sqrt{-n^{2}} x}+2 \sin (r x) n^{2} h \mathrm{e}^{\sqrt{-n^{2}} x}+\left(c n^{2}-c r^{2}-l\right.\right.}{2 n^{4}-2 n^{2} r^{2}}\)
Verification of solutions
\(y\)
\(=\frac{\mathrm{e}^{-\sqrt{-n^{2}} x}\left(\left(\left(-c n^{2}+r(c r+h)\right) \sqrt{-n^{2}}+n^{4} a-n^{2} a r^{2}\right) \mathrm{e}^{2 \sqrt{-n^{2}} x}+2 \sin (r x) n^{2} h \mathrm{e}^{\sqrt{-n^{2}} x}+\left(c n^{2}-c r^{2}-l\right.\right.}{2 n^{4}-2 n^{2} r^{2}}\)

\section*{Verified OK.}

\subsection*{1.115.3 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+y n^{2} & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=0  \tag{3}\\
& C=n^{2}
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{-n^{2}}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=-n^{2} \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=\left(-n^{2}\right) z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 179: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=-n^{2}\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{\sqrt{-n^{2}} x}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\mathrm{e}^{\sqrt{-n^{2} x}}
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{\sqrt{-n^{2}} x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\mathrm{e}^{\sqrt{-n^{2}} x} \int \frac{1}{\mathrm{e}^{2 \sqrt{-n^{2}} x}} d x \\
& =\mathrm{e}^{\sqrt{-n^{2}} x}\left(\frac{\sqrt{-n^{2}} \mathrm{e}^{-2 \sqrt{-n^{2}} x}}{2 n^{2}}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{\sqrt{-n^{2}} x}\right)+c_{2}\left(\mathrm{e}^{\sqrt{-n^{2}} x}\left(\frac{\sqrt{-n^{2}} \mathrm{e}^{-2 \sqrt{-n^{2}} x}}{2 n^{2}}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+y n^{2}=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \mathrm{e}^{\sqrt{-n^{2}} x}+\frac{c_{2} \sqrt{-n^{2}} \mathrm{e}^{-\sqrt{-n^{2}} x}}{2 n^{2}}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
h \sin (r x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{\cos (r x), \sin (r x)\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\frac{\sqrt{-n^{2}} \mathrm{e}^{-\sqrt{-n^{2}} x}}{2 n^{2}}, \mathrm{e}^{\sqrt{-n^{2}} x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \cos (r x)+A_{2} \sin (r x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-A_{1} r^{2} \cos (r x)-A_{2} r^{2} \sin (r x)+\left(A_{1} \cos (r x)+A_{2} \sin (r x)\right) n^{2}=h \sin (r x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=0, A_{2}=\frac{h}{n^{2}-r^{2}}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{h \sin (r x)}{n^{2}-r^{2}}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{\sqrt{-n^{2}} x}+\frac{c_{2} \sqrt{-n^{2}} \mathrm{e}^{-\sqrt{-n^{2}} x}}{2 n^{2}}\right)+\left(\frac{h \sin (r x)}{n^{2}-r^{2}}\right)
\end{aligned}
\]

Initial conditions are used to solve for the constants of integration.
Looking at the above solution
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{\sqrt{-n^{2}} x}+\frac{c_{2} \sqrt{-n^{2}} \mathrm{e}^{-\sqrt{-n^{2}} x}}{2 n^{2}}+\frac{h \sin (r x)}{n^{2}-r^{2}} \tag{1}
\end{equation*}
\]

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting \(y=a\) and \(x=0\) in the above gives
\[
\begin{equation*}
a=\frac{2 c_{1} n^{2}+\sqrt{-n^{2}} c_{2}}{2 n^{2}} \tag{1~A}
\end{equation*}
\]

Taking derivative of the solution gives
\[
y^{\prime}=c_{1} \sqrt{-n^{2}} \mathrm{e}^{\sqrt{-n^{2}} x}+\frac{c_{2} \mathrm{e}^{-\sqrt{-n^{2}} x}}{2}+\frac{h r \cos (r x)}{n^{2}-r^{2}}
\]
substituting \(y^{\prime}=c\) and \(x=0\) in the above gives
\[
\begin{equation*}
c=\frac{2 \sqrt{-n^{2}} c_{1} n^{2}-2 \sqrt{-n^{2}} c_{1} r^{2}+c_{2} n^{2}-c_{2} r^{2}+2 h r}{2 n^{2}-2 r^{2}} \tag{2~A}
\end{equation*}
\]

Equations \(\{1 \mathrm{~A}, 2 \mathrm{~A}\}\) are now solved for \(\left\{c_{1}, c_{2}\right\}\). Solving for the constants gives
\[
\begin{aligned}
& c_{1}=-\frac{-n^{4} a+n^{2} a r^{2}+\sqrt{-n^{2}} c n^{2}-\sqrt{-n^{2}} c r^{2}-\sqrt{-n^{2}} h r}{2 n^{2}\left(n^{2}-r^{2}\right)} \\
& c_{2}=\frac{-\sqrt{-n^{2}} n^{2} a+\sqrt{-n^{2}} r^{2} a+c n^{2}-c r^{2}-h r}{n^{2}-r^{2}}
\end{aligned}
\]

Substituting these values back in above solution results in
\(y=\frac{\mathrm{e}^{\sqrt{-n^{2}} x} a n^{4}-\mathrm{e}^{\sqrt{-n^{2}} x} a n^{2} r^{2}+\mathrm{e}^{-\sqrt{-n^{2}} x} a n^{4}-\mathrm{e}^{-\sqrt{-n^{2}} x} a n^{2} r^{2}-\mathrm{e}^{\sqrt{-n^{2}} x} \sqrt{-n^{2}} c n^{2}+\mathrm{e}^{\sqrt{-n^{2}} x} \sqrt{-n^{2}} c r^{2}}{2 n^{4}-2}\)
Which simplifies to
\(y\)
\(=\frac{\mathrm{e}^{-\sqrt{-n^{2}} x}\left(\left(\left(-c n^{2}+r(c r+h)\right) \sqrt{-n^{2}}+n^{4} a-n^{2} a r^{2}\right) \mathrm{e}^{2 \sqrt{-n^{2}} x}+2 \sin (r x) n^{2} h \mathrm{e}^{\sqrt{-n^{2}} x}+\left(c n^{2}-c r^{2}-l\right.\right.}{2 n^{4}-2 n^{2} r^{2}}\)

\section*{Summary}

The solution(s) found are the following
\(y\)
\(=\frac{\mathrm{e}^{-\sqrt{-n^{2}} x}\left(\left(\left(-c n^{2}+r(c r+h)\right) \sqrt{-n^{2}}+n^{4} a-n^{2} a r^{2}\right) \mathrm{e}^{2 \sqrt{-n^{2}} x}+2 \sin (r x) n^{2} h \mathrm{e}^{\sqrt{-n^{2}} x}+\left(c n^{2}-c r^{2}-l\right.\right.}{2 n^{4}-2 n^{2} r^{2}}\)
Verification of solutions
\(y\)
\(=\frac{\mathrm{e}^{-\sqrt{-n^{2}} x}\left(\left(\left(-c n^{2}+r(c r+h)\right) \sqrt{-n^{2}}+n^{4} a-n^{2} a r^{2}\right) \mathrm{e}^{2 \sqrt{-n^{2}} x}+2 \sin (r x) n^{2} h \mathrm{e}^{\sqrt{-n^{2}} x}+\left(c n^{2}-c r^{2}-l\right.\right.}{2 n^{4}-2 n^{2} r^{2}}\)
Verified OK.

\subsection*{1.115.4 Maple step by step solution}

Let's solve
\[
\left[y^{\prime \prime}+y n^{2}=h \sin (r x), y(0)=a,\left.y^{\prime}\right|_{\{x=0\}}=c\right]
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\[
n^{2}+r^{2}=0
\]
- Use quadratic formula to solve for \(r\)
\[
r=\frac{0 \pm\left(\sqrt{-4 n^{2}}\right)}{2}
\]
- Roots of the characteristic polynomial
\[
r=\left(\sqrt{-n^{2}},-\sqrt{-n^{2}}\right)
\]
- \(\quad 1\) st solution of the homogeneous ODE
\[
y_{1}(x)=\mathrm{e}^{\sqrt{-n^{2}} x}
\]
- \(\quad 2\) nd solution of the homogeneous ODE
\[
y_{2}(x)=\mathrm{e}^{-\sqrt{-n^{2}} x}
\]
- General solution of the ODE
\[
y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)
\]
- Substitute in solutions of the homogeneous ODE
\[
y=c_{1} \mathrm{e}^{\sqrt{-n^{2}} x}+c_{2} \mathrm{e}^{-\sqrt{-n^{2}} x}+y_{p}(x)
\]

Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\[
\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=h \sin (r x)\right]
\]
- Wronskian of solutions of the homogeneous equation
\[
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
\mathrm{e}^{\sqrt{-n^{2}} x} & \mathrm{e}^{-\sqrt{-n^{2}} x} \\
\sqrt{-n^{2}} \mathrm{e}^{\sqrt{-n^{2}} x} & -\sqrt{-n^{2}} \mathrm{e}^{-\sqrt{-n^{2}} x}
\end{array}\right]
\]
- Compute Wronskian
\[
W\left(y_{1}(x), y_{2}(x)\right)=-2 \sqrt{-n^{2}}
\]
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=\frac{h\left(\mathrm{e}^{\sqrt{-n^{2}} x}\left(\int \mathrm{e}^{-\sqrt{-n^{2}} x} \sin (r x) d x\right)-\mathrm{e}^{-\sqrt{-n^{2}} x}\left(\int \mathrm{e}^{\sqrt{-n^{2}} x} \sin (r x) d x\right)\right)}{2 \sqrt{-n^{2}}}
\]
- Compute integrals
\[
y_{p}(x)=\frac{h \sin (r x)}{n^{2}-r^{2}}
\]
- Substitute particular solution into general solution to ODE
\(y=c_{1} \mathrm{e}^{\sqrt{-n^{2}} x}+c_{2} \mathrm{e}^{-\sqrt{-n^{2}} x}+\frac{h \sin (r x)}{n^{2}-r^{2}}\)
Check validity of solution \(y=c_{1} \mathrm{e}^{\sqrt{-n^{2}} x}+c_{2} \mathrm{e}^{-\sqrt{-n^{2}} x}+\frac{h \sin (r x)}{n^{2}-r^{2}}\)
- Use initial condition \(y(0)=a\)
\[
a=c_{1}+c_{2}
\]
- Compute derivative of the solution
\[
y^{\prime}=c_{1} \sqrt{-n^{2}} \mathrm{e}^{\sqrt{-n^{2}} x}-c_{2} \sqrt{-n^{2}} \mathrm{e}^{-\sqrt{-n^{2}} x}+\frac{h r \cos (r x)}{n^{2}-r^{2}}
\]
- Use the initial condition \(\left.y^{\prime}\right|_{\{x=0\}}=c\)
\[
c=\sqrt{-n^{2}} c_{1}-\sqrt{-n^{2}} c_{2}+\frac{r h}{n^{2}-r^{2}}
\]
- Solve for \(c_{1}\) and \(c_{2}\)
\[
\left\{c_{1}=\frac{-n^{4} a+n^{2} a r^{2}+\sqrt{-n^{2}} c n^{2}-\sqrt{-n^{2}} c r^{2}-\sqrt{-n^{2}} h r}{2 n^{2}\left(-n^{2}+r^{2}\right)}, c_{2}=\frac{\sqrt{-n^{2}}\left(\sqrt{-n^{2}} n^{2} a-\sqrt{-n^{2}} r^{2} a-c n^{2}+c r^{2}+h r\right)}{2 n^{2}\left(-n^{2}+r^{2}\right)}\right\}
\]
- Substitute constant values into general solution and simplify
\(y=\frac{\mathrm{e}^{-\sqrt{-n^{2}} x}\left(\left(\left(-c n^{2}+r(c r+h)\right) \sqrt{-n^{2}}+n^{4} a-n^{2} a r^{2}\right) \mathrm{e}^{2 \sqrt{-n^{2}} x}+2 \sin (r x) n^{2} h \mathrm{e}^{\sqrt{-n^{2}} x}+\left(c n^{2}-c r^{2}-h r\right) \sqrt{-n^{2}}+n^{4} a-n^{2} a r^{2}\right)}{2 n^{4}-2 n^{2} r^{2}}\)
- \(\quad\) Solution to the IVP
\[
y=\frac{\mathrm{e}^{-\sqrt{-n^{2}} x}\left(\left(\left(-c n^{2}+r(c r+h)\right) \sqrt{-n^{2}}+n^{4} a-n^{2} a r^{2}\right) \mathrm{e}^{2 \sqrt{-n^{2}} x}+2 \sin (r x) n^{2} h \mathrm{e}^{\sqrt{-n^{2}} x}+\left(c n^{2}-c r^{2}-h r\right) \sqrt{-n^{2}}+n^{4} a-n^{2} a r^{2}\right)}{2 n^{4}-2 n^{2} r^{2}}
\]

Maple trace
- Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
checking if the LODE has constant coefficients
<- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`
\(\checkmark\) Solution by Maple
Time used: 0.047 (sec). Leaf size: 60
dsolve([diff \(\left.(y(x), x \$ 2)+n^{\wedge} 2 * y(x)=h * \sin (r * x), y(0)=a, D(y)(0)=c\right], y(x)\), singsol=all)
\[
y(x)=-\frac{\sin (x n)\left(-n^{2} c+c r^{2}+h r\right)}{n^{3}-n r^{2}}+\cos (x n) a+\frac{h \sin (r x)}{n^{2}-r^{2}}
\]

Solution by Mathematica
Time used: 0.066 (sec). Leaf size: 63
DSolve \(\left[\left\{y^{\prime \prime}[x]+n^{\wedge} 2 * y[x]==h * \operatorname{Sin}[r * x],\left\{y[0]==a, y^{\prime}[0]==c\right\}\right\}, y[x], x\right.\), IncludeSingularSolutions \(->\) I
\[
y(x) \rightarrow \frac{a n\left(n^{2}-r^{2}\right) \cos (n x)+\sin (n x)\left(c n^{2}-c r^{2}-h r\right)+h n \sin (r x)}{n^{3}-n r^{2}}
\]

\subsection*{1.116 problem 167}
1.116.1 Solving as second order linear constant coeff ode . . . . . . . . 1183
1.116.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1186
1.116.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1191

Internal problem ID [12533]
Internal file name [OUTPUT/11185_Tuesday_October_17_2023_07_20_37_AM_85751881/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 167.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _linear, _nonhomogeneous]]
\[
y^{\prime \prime}-7 y^{\prime}+6 y=\sin (x)
\]

\subsection*{1.116.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=-7, C=6, f(x)=\sin (x)\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-7 y^{\prime}+6 y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=-7, C=6\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-7 \lambda \mathrm{e}^{\lambda x}+6 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}-7 \lambda+6=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=-7, C=6\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{7}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-7^{2}-(4)(1)(6)} \\
& =\frac{7}{2} \pm \frac{5}{2}
\end{aligned}
\]

Hence
\[
\begin{aligned}
\lambda_{1} & =\frac{7}{2}+\frac{5}{2} \\
\lambda_{2} & =\frac{7}{2}-\frac{5}{2}
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=6 \\
& \lambda_{2}=1
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{(6) x}+c_{2} e^{(1) x}
\end{aligned}
\]

Or
\[
y=c_{1} \mathrm{e}^{6 x}+c_{2} \mathrm{e}^{x}
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \mathrm{e}^{6 x}+c_{2} \mathrm{e}^{x}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
\sin (x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{\cos (x), \sin (x)\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{x}, \mathrm{e}^{6 x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \cos (x)+A_{2} \sin (x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
5 A_{1} \cos (x)+5 A_{2} \sin (x)+7 A_{1} \sin (x)-7 A_{2} \cos (x)=\sin (x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{7}{74}, A_{2}=\frac{5}{74}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{6 x}+c_{2} \mathrm{e}^{x}\right)+\left(\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74}\right)
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{6 x}+c_{2} \mathrm{e}^{x}+\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74} \tag{1}
\end{equation*}
\]


Figure 185: Slope field plot

Verification of solutions
\[
y=c_{1} \mathrm{e}^{6 x}+c_{2} \mathrm{e}^{x}+\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74}
\]

Verified OK.

\subsection*{1.116.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}-7 y^{\prime}+6 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=-7  \tag{3}\\
& C=6
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{25}{4} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=25 \\
& t=4
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=\frac{25 z(x)}{4} \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi-
\end{tabular} & no condition \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\). & \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 181: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=\frac{25}{4}\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-\frac{5 x}{2}}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
\begin{aligned}
y_{1} & =z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x} \\
& =z_{1} e^{-\int \frac{1}{2} \frac{-7}{1} d x} \\
& =z_{1} e^{\frac{7 x}{2}} \\
& =z_{1}\left(\mathrm{e}^{\frac{7 x}{2}}\right)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Substituting gives
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{e^{\int-\frac{-7}{1} d x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1} \int \frac{e^{7 x}}{\left(y_{1}\right)^{2}} d x \\
& =y_{1}\left(\frac{\mathrm{e}^{5 x}}{5}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{x}\right)+c_{2}\left(\mathrm{e}^{x}\left(\frac{\mathrm{e}^{5 x}}{5}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-7 y^{\prime}+6 y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=\mathrm{e}^{x} c_{1}+\frac{c_{2} \mathrm{e}^{6 x}}{5}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
\sin (x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
[\{\cos (x), \sin (x)\}]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\frac{\mathrm{e}^{6 x}}{5}, \mathrm{e}^{x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \cos (x)+A_{2} \sin (x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
5 A_{1} \cos (x)+5 A_{2} \sin (x)+7 A_{1} \sin (x)-7 A_{2} \cos (x)=\sin (x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=\frac{7}{74}, A_{2}=\frac{5}{74}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(\mathrm{e}^{x} c_{1}+\frac{c_{2} \mathrm{e}^{6 x}}{5}\right)+\left(\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\mathrm{e}^{x} c_{1}+\frac{c_{2} \mathrm{e}^{6 x}}{5}+\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74} \tag{1}
\end{equation*}
\]


Figure 186: Slope field plot

\section*{Verification of solutions}
\[
y=\mathrm{e}^{x} c_{1}+\frac{c_{2} \mathrm{e}^{6 x}}{5}+\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74}
\]

Verified OK.

\subsection*{1.116.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}-7 y^{\prime}+6 y=\sin (x)
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\[
r^{2}-7 r+6=0
\]
- Factor the characteristic polynomial
\[
(r-1)(r-6)=0
\]
- Roots of the characteristic polynomial
\(r=(1,6)\)
- \(\quad 1\) st solution of the homogeneous ODE
\(y_{1}(x)=\mathrm{e}^{x}\)
- \(\quad 2 n d\) solution of the homogeneous ODE
\(y_{2}(x)=\mathrm{e}^{6 x}\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)\)
- \(\quad\) Substitute in solutions of the homogeneous ODE
\(y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{6 x}+y_{p}(x)\)
Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function \(\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=\sin (x)\right]\)
- Wronskian of solutions of the homogeneous equation
\(W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}\mathrm{e}^{x} & \mathrm{e}^{6 x} \\ \mathrm{e}^{x} & 6 \mathrm{e}^{6 x}\end{array}\right]\)
- Compute Wronskian
\(W\left(y_{1}(x), y_{2}(x)\right)=5 \mathrm{e}^{7 x}\)
- Substitute functions into equation for \(y_{p}(x)\)
\(y_{p}(x)=-\frac{\mathrm{e}^{x}\left(\int \mathrm{e}^{-x} \sin (x) d x\right)}{5}+\frac{\mathrm{e}^{6 x}\left(\int \sin (x) \mathrm{e}^{-6 x} d x\right)}{5}\)
- Compute integrals
\(y_{p}(x)=\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74}\)
- Substitute particular solution into general solution to ODE
\(y=\mathrm{e}^{x} c_{1}+c_{2} \mathrm{e}^{6 x}+\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74}\)

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 23
```

dsolve(diff(y(x),x\$2)-7*diff (y(x),x)+6*y(x)=sin(x),y(x), singsol=all)

```
\[
y(x)=\mathrm{e}^{6 x} c_{2}+c_{1} \mathrm{e}^{x}+\frac{7 \cos (x)}{74}+\frac{5 \sin (x)}{74}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.104 (sec). Leaf size: 32
DSolve[y''[x]-7*y'[x]+6*y[x]==Sin[x],y[x],x,IncludeSingularSolutions -> True]
\[
y(x) \rightarrow \frac{5 \sin (x)}{74}+\frac{7 \cos (x)}{74}+c_{1} e^{x}+c_{2} e^{6 x}
\]

\subsection*{1.117 problem 168}
1.117.1 Solving as second order linear constant coeff ode 1194
1.117.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1198
1.117.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1204

Internal problem ID [12534]
Internal file name [OUTPUT/11186_Tuesday_October_17_2023_07_20_38_AM_36289514/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 168.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _linear, _nonhomogeneous]]
\[
y^{\prime \prime}+y=\sec (x)
\]

\subsection*{1.117.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=0, C=1, f(x)=\sec (x)\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=1\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+\mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+1=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=1\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)(1)} \\
& = \pm i
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=+i \\
& \lambda_{2}=-i
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=i \\
& \lambda_{2}=-i
\end{aligned}
\]

Since roots are complex conjugate of each others, then let the roots be
\[
\lambda_{1,2}=\alpha \pm i \beta
\]

Where \(\alpha=0\) and \(\beta=1\). Therefore the final solution, when using Euler relation, can be written as
\[
y=e^{\alpha x}\left(c_{1} \cos (\beta x)+c_{2} \sin (\beta x)\right)
\]

Which becomes
\[
y=e^{0}\left(c_{1} \cos (x)+c_{2} \sin (x)\right)
\]

Or
\[
y=c_{1} \cos (x)+c_{2} \sin (x)
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \cos (x)+c_{2} \sin (x)
\]

The particular solution \(y_{p}\) can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on \(x\) as well. Let
\[
\begin{equation*}
y_{p}(x)=u_{1} y_{1}+u_{2} y_{2} \tag{1}
\end{equation*}
\]

Where \(u_{1}, u_{2}\) to be determined, and \(y_{1}, y_{2}\) are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as
\[
\begin{aligned}
& y_{1}=\cos (x) \\
& y_{2}=\sin (x)
\end{aligned}
\]

In the Variation of parameters \(u_{1}, u_{2}\) are found using
\[
\begin{align*}
& u_{1}=-\int \frac{y_{2} f(x)}{a W(x)}  \tag{2}\\
& u_{2}=\int \frac{y_{1} f(x)}{a W(x)} \tag{3}
\end{align*}
\]

Where \(W(x)\) is the Wronskian and \(a\) is the coefficient in front of \(y^{\prime \prime}\) in the given ODE. The Wronskian is given by \(W=\left|\begin{array}{ll}y_{1} & y_{2} \\ y_{1}^{\prime} & y_{2}^{\prime}\end{array}\right|\). Hence
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
\frac{d}{d x}(\cos (x)) & \frac{d}{d x}(\sin (x))
\end{array}\right|
\]

Which gives
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
-\sin (x) & \cos (x)
\end{array}\right|
\]

Therefore
\[
W=(\cos (x))(\cos (x))-(\sin (x))(-\sin (x))
\]

Which simplifies to
\[
W=\cos (x)^{2}+\sin (x)^{2}
\]

Which simplifies to
\[
W=1
\]

Therefore Eq. (2) becomes
\[
u_{1}=-\int \frac{\sin (x) \sec (x)}{1} d x
\]

Which simplifies to
\[
u_{1}=-\int \tan (x) d x
\]

Hence
\[
u_{1}=\ln (\cos (x))
\]

And Eq. (3) becomes
\[
u_{2}=\int \frac{\cos (x) \sec (x)}{1} d x
\]

Which simplifies to
\[
u_{2}=\int 1 d x
\]

Hence
\[
u_{2}=x
\]

Therefore the particular solution, from equation (1) is
\[
y_{p}(x)=\ln (\cos (x)) \cos (x)+x \sin (x)
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \cos (x)+c_{2} \sin (x)\right)+(\ln (\cos (x)) \cos (x)+x \sin (x))
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (x)+c_{2} \sin (x)+\ln (\cos (x)) \cos (x)+x \sin (x) \tag{1}
\end{equation*}
\]


Figure 187: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \cos (x)+c_{2} \sin (x)+\ln (\cos (x)) \cos (x)+x \sin (x)
\]

Verified OK.

\subsection*{1.117.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=0  \tag{3}\\
& C=1
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{-1}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=-1 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=-z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi-
\end{tabular} & no condition \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\). & \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 183: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=-1\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\cos (x)
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\cos (x)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\cos (x)
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\cos (x) \int \frac{1}{\cos (x)^{2}} d x \\
& =\cos (x)(\tan (x))
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}(\cos (x))+c_{2}(\cos (x)(\tan (x)))
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \cos (x)+c_{2} \sin (x)
\]

The particular solution \(y_{p}\) can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of
parameters will be used as it is more general and can be used when the coefficients of the ODE depend on \(x\) as well. Let
\[
\begin{equation*}
y_{p}(x)=u_{1} y_{1}+u_{2} y_{2} \tag{1}
\end{equation*}
\]

Where \(u_{1}, u_{2}\) to be determined, and \(y_{1}, y_{2}\) are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as
\[
\begin{aligned}
& y_{1}=\cos (x) \\
& y_{2}=\sin (x)
\end{aligned}
\]

In the Variation of parameters \(u_{1}, u_{2}\) are found using
\[
\begin{align*}
& u_{1}=-\int \frac{y_{2} f(x)}{a W(x)}  \tag{2}\\
& u_{2}=\int \frac{y_{1} f(x)}{a W(x)} \tag{3}
\end{align*}
\]

Where \(W(x)\) is the Wronskian and \(a\) is the coefficient in front of \(y^{\prime \prime}\) in the given ODE. The Wronskian is given by \(W=\left|\begin{array}{ll}y_{1} & y_{2} \\ y_{1}^{\prime} & y_{2}^{\prime}\end{array}\right|\). Hence
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
\frac{d}{d x}(\cos (x)) & \frac{d}{d x}(\sin (x))
\end{array}\right|
\]

Which gives
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
-\sin (x) & \cos (x)
\end{array}\right|
\]

Therefore
\[
W=(\cos (x))(\cos (x))-(\sin (x))(-\sin (x))
\]

Which simplifies to
\[
W=\cos (x)^{2}+\sin (x)^{2}
\]

Which simplifies to
\[
W=1
\]

Therefore Eq. (2) becomes
\[
u_{1}=-\int \frac{\sin (x) \sec (x)}{1} d x
\]

Which simplifies to
\[
u_{1}=-\int \tan (x) d x
\]

Hence
\[
u_{1}=\ln (\cos (x))
\]

And Eq. (3) becomes
\[
u_{2}=\int \frac{\cos (x) \sec (x)}{1} d x
\]

Which simplifies to
\[
u_{2}=\int 1 d x
\]

Hence
\[
u_{2}=x
\]

Therefore the particular solution, from equation (1) is
\[
y_{p}(x)=\ln (\cos (x)) \cos (x)+x \sin (x)
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \cos (x)+c_{2} \sin (x)\right)+(\ln (\cos (x)) \cos (x)+x \sin (x))
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (x)+c_{2} \sin (x)+\ln (\cos (x)) \cos (x)+x \sin (x) \tag{1}
\end{equation*}
\]


Figure 188: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \cos (x)+c_{2} \sin (x)+\ln (\cos (x)) \cos (x)+x \sin (x)
\]

Verified OK.

\subsection*{1.117.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}+y=\sec (x)
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\[
r^{2}+1=0
\]
- Use quadratic formula to solve for \(r\)
\[
r=\frac{0 \pm(\sqrt{-4})}{2}
\]
- Roots of the characteristic polynomial
\[
r=(-\mathrm{I}, \mathrm{I})
\]
- \(\quad 1\) st solution of the homogeneous ODE
\[
y_{1}(x)=\cos (x)
\]
- \(\quad 2 n d\) solution of the homogeneous ODE
\[
y_{2}(x)=\sin (x)
\]
- General solution of the ODE
\[
y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)
\]
- Substitute in solutions of the homogeneous ODE
\[
y=c_{1} \cos (x)+c_{2} \sin (x)+y_{p}(x)
\]

Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\[
\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=\sec (x)\right]
\]
- Wronskian of solutions of the homogeneous equation
\[
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
\cos (x) & \sin (x) \\
-\sin (x) & \cos (x)
\end{array}\right]
\]
- Compute Wronskian
\(W\left(y_{1}(x), y_{2}(x)\right)=1\)
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=-\cos (x)\left(\int \tan (x) d x\right)+\sin (x)\left(\int 1 d x\right)
\]
- Compute integrals
\[
y_{p}(x)=\ln (\cos (x)) \cos (x)+x \sin (x)
\]
- Substitute particular solution into general solution to ODE
\[
y=c_{1} \cos (x)+c_{2} \sin (x)+\ln (\cos (x)) \cos (x)+x \sin (x)
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 22
```

dsolve(diff(y(x),x\$2)+y(x)=sec(x),y(x), singsol=all)

```
\[
y(x)=-\ln (\sec (x)) \cos (x)+c_{1} \cos (x)+\sin (x)\left(c_{2}+x\right)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.037 (sec). Leaf size: 22
DSolve[y''[x]+y[x]==Sec[x],y[x],x,IncludeSingularSolutions -> True]
\[
y(x) \rightarrow\left(x+c_{2}\right) \sin (x)+\cos (x)\left(\log (\cos (x))+c_{1}\right)
\]

\subsection*{1.118 problem 169}
1.118.1 Solving as second order linear constant coeff ode . . . . . . . . 1207
1.118.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1212
1.118.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1218

Internal problem ID [12535]
Internal file name [OUTPUT/11187_Tuesday_October_17_2023_07_20_39_AM_97353951/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 169.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _linear, _nonhomogeneous]]
\[
y^{\prime \prime}+y=\frac{1}{\cos (2 x)^{\frac{3}{2}}}
\]

\subsection*{1.118.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=0, C=1, f(x)=\frac{1}{\cos (2 x)^{\frac{3}{2}}}\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=1\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+\mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\mathrm{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+1=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form.Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=1\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)(1)} \\
& = \pm i
\end{aligned}
\]

Hence
\[
\begin{gathered}
\lambda_{1}=+i \\
\lambda_{2}=-i
\end{gathered}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=i \\
& \lambda_{2}=-i
\end{aligned}
\]

Since roots are complex conjugate of each others, then let the roots be
\[
\lambda_{1,2}=\alpha \pm i \beta
\]

Where \(\alpha=0\) and \(\beta=1\). Therefore the final solution, when using Euler relation, can be written as
\[
y=e^{\alpha x}\left(c_{1} \cos (\beta x)+c_{2} \sin (\beta x)\right)
\]

Which becomes
\[
y=e^{0}\left(c_{1} \cos (x)+c_{2} \sin (x)\right)
\]

Or
\[
y=c_{1} \cos (x)+c_{2} \sin (x)
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \cos (x)+c_{2} \sin (x)
\]

The particular solution \(y_{p}\) can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on \(x\) as well. Let
\[
\begin{equation*}
y_{p}(x)=u_{1} y_{1}+u_{2} y_{2} \tag{1}
\end{equation*}
\]

Where \(u_{1}, u_{2}\) to be determined, and \(y_{1}, y_{2}\) are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as
\[
\begin{aligned}
& y_{1}=\cos (x) \\
& y_{2}=\sin (x)
\end{aligned}
\]

In the Variation of parameters \(u_{1}, u_{2}\) are found using
\[
\begin{align*}
& u_{1}=-\int \frac{y_{2} f(x)}{a W(x)}  \tag{2}\\
& u_{2}=\int \frac{y_{1} f(x)}{a W(x)} \tag{3}
\end{align*}
\]

Where \(W(x)\) is the Wronskian and \(a\) is the coefficient in front of \(y^{\prime \prime}\) in the given ODE. The Wronskian is given by \(W=\left|\begin{array}{ll}y_{1} & y_{2} \\ y_{1}^{\prime} & y_{2}^{\prime}\end{array}\right|\). Hence
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
\frac{d}{d x}(\cos (x)) & \frac{d}{d x}(\sin (x))
\end{array}\right|
\]

Which gives
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
-\sin (x) & \cos (x)
\end{array}\right|
\]

Therefore
\[
W=(\cos (x))(\cos (x))-(\sin (x))(-\sin (x))
\]

Which simplifies to
\[
W=\cos (x)^{2}+\sin (x)^{2}
\]

Which simplifies to
\[
W=1
\]

Therefore Eq. (2) becomes
\[
u_{1}=-\int \frac{\frac{\sin (x)}{\cos (2 x)^{\frac{3}{2}}}}{1} d x
\]

Which simplifies to
\[
u_{1}=-\int \frac{\sin (x)}{\cos (2 x)^{\frac{3}{2}}} d x
\]

Hence
\[
u_{1}=-\frac{\cos (x)}{\sqrt{2 \cos (x)^{2}-1}}
\]

And Eq. (3) becomes
\[
u_{2}=\int \frac{\frac{\cos (x)}{\cos (2 x)^{\frac{3}{2}}}}{1} d x
\]

Which simplifies to
\[
u_{2}=\int \frac{\cos (x)}{\cos (2 x)^{\frac{3}{2}}} d x
\]

Hence
\[
u_{2}=\frac{\sqrt{\left(2 \cos (x)^{2}-1\right) \sin (x)^{2}} \sin (x)}{\sqrt{-2 \sin (x)^{4}+\sin (x)^{2}} \sqrt{2 \cos (x)^{2}-1}}
\]

Which simplifies to
\[
\begin{aligned}
& u_{1}=-\frac{\cos (x)}{\sqrt{\cos (2 x)}} \\
& u_{2}=\frac{\sin (x)}{\sqrt{\cos (2 x)}}
\end{aligned}
\]

Therefore the particular solution, from equation (1) is
\[
y_{p}(x)=-\frac{\cos (x)^{2}}{\sqrt{\cos (2 x)}}+\frac{\sin (x)^{2}}{\sqrt{\cos (2 x)}}
\]

Which simplifies to
\[
y_{p}(x)=-\sqrt{\cos (2 x)}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \cos (x)+c_{2} \sin (x)\right)+(-\sqrt{\cos (2 x)})
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (x)+c_{2} \sin (x)-\sqrt{\cos (2 x)} \tag{1}
\end{equation*}
\]


Figure 189: Slope field plot

Verification of solutions
\[
y=c_{1} \cos (x)+c_{2} \sin (x)-\sqrt{\cos (2 x)}
\]

Verified OK.

\subsection*{1.118.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{array}{r}
y^{\prime \prime}+y=0 \\
A y^{\prime \prime}+B y^{\prime}+C y=0 \tag{2}
\end{array}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
A & =1 \\
B & =0  \tag{3}\\
C & =1
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{-1}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=-1 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=-z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 185: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=-1\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\cos (x)
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\cos (x)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\cos (x)
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\cos (x) \int \frac{1}{\cos (x)^{2}} d x \\
& =\cos (x)(\tan (x))
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}(\cos (x))+c_{2}(\cos (x)(\tan (x)))
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \cos (x)+c_{2} \sin (x)
\]

The particular solution \(y_{p}\) can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on \(x\) as well. Let
\[
\begin{equation*}
y_{p}(x)=u_{1} y_{1}+u_{2} y_{2} \tag{1}
\end{equation*}
\]

Where \(u_{1}, u_{2}\) to be determined, and \(y_{1}, y_{2}\) are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as
\[
\begin{aligned}
& y_{1}=\cos (x) \\
& y_{2}=\sin (x)
\end{aligned}
\]

In the Variation of parameters \(u_{1}, u_{2}\) are found using
\[
\begin{align*}
& u_{1}=-\int \frac{y_{2} f(x)}{a W(x)}  \tag{2}\\
& u_{2}=\int \frac{y_{1} f(x)}{a W(x)} \tag{3}
\end{align*}
\]

Where \(W(x)\) is the Wronskian and \(a\) is the coefficient in front of \(y^{\prime \prime}\) in the given ODE.
The Wronskian is given by \(W=\left|\begin{array}{ll}y_{1} & y_{2} \\ y_{1}^{\prime} & y_{2}^{\prime}\end{array}\right|\). Hence
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
\frac{d}{d x}(\cos (x)) & \frac{d}{d x}(\sin (x))
\end{array}\right|
\]

Which gives
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
-\sin (x) & \cos (x)
\end{array}\right|
\]

Therefore
\[
W=(\cos (x))(\cos (x))-(\sin (x))(-\sin (x))
\]

Which simplifies to
\[
W=\cos (x)^{2}+\sin (x)^{2}
\]

Which simplifies to
\[
W=1
\]

Therefore Eq. (2) becomes
\[
u_{1}=-\int \frac{\frac{\sin (x)}{\cos (2 x)^{\frac{3}{2}}}}{1} d x
\]

Which simplifies to
\[
u_{1}=-\int \frac{\sin (x)}{\cos (2 x)^{\frac{3}{2}}} d x
\]

Hence
\[
u_{1}=-\frac{\cos (x)}{\sqrt{2 \cos (x)^{2}-1}}
\]

And Eq. (3) becomes
\[
u_{2}=\int \frac{\frac{\cos (x)}{\cos (2 x)^{\frac{3}{2}}}}{1} d x
\]

Which simplifies to
\[
u_{2}=\int \frac{\cos (x)}{\cos (2 x)^{\frac{3}{2}}} d x
\]

Hence
\[
u_{2}=\frac{\sqrt{\left(2 \cos (x)^{2}-1\right) \sin (x)^{2}} \sin (x)}{\sqrt{-2 \sin (x)^{4}+\sin (x)^{2}} \sqrt{2 \cos (x)^{2}-1}}
\]

Which simplifies to
\[
\begin{aligned}
& u_{1}=-\frac{\cos (x)}{\sqrt{\cos (2 x)}} \\
& u_{2}=\frac{\sin (x)}{\sqrt{\cos (2 x)}}
\end{aligned}
\]

Therefore the particular solution, from equation (1) is
\[
y_{p}(x)=-\frac{\cos (x)^{2}}{\sqrt{\cos (2 x)}}+\frac{\sin (x)^{2}}{\sqrt{\cos (2 x)}}
\]

Which simplifies to
\[
y_{p}(x)=-\sqrt{\cos (2 x)}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \cos (x)+c_{2} \sin (x)\right)+(-\sqrt{\cos (2 x)})
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (x)+c_{2} \sin (x)-\sqrt{\cos (2 x)} \tag{1}
\end{equation*}
\]


Figure 190: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \cos (x)+c_{2} \sin (x)-\sqrt{\cos (2 x)}
\]

Verified OK.

\subsection*{1.118.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}+y=\frac{1}{\cos (2 x)^{\frac{3}{2}}}
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE \(r^{2}+1=0\)
- Use quadratic formula to solve for \(r\)
\[
r=\frac{0 \pm(\sqrt{-4})}{2}
\]
- Roots of the characteristic polynomial
\(r=(-\mathrm{I}, \mathrm{I})\)
- \(\quad 1\) st solution of the homogeneous ODE
\(y_{1}(x)=\cos (x)\)
- 2nd solution of the homogeneous ODE
\(y_{2}(x)=\sin (x)\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)\)
- \(\quad\) Substitute in solutions of the homogeneous ODE
\(y=c_{1} \cos (x)+c_{2} \sin (x)+y_{p}(x)\)
Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\(\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=\frac{1}{\cos (2 x)^{\frac{3}{2}}}\right]\)
- Wronskian of solutions of the homogeneous equation
\(W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}\cos (x) & \sin (x) \\ -\sin (x) & \cos (x)\end{array}\right]\)
- Compute Wronskian
\(W\left(y_{1}(x), y_{2}(x)\right)=1\)
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=-\cos (x)\left(\int \frac{\sin (x)}{\cos (2 x)^{\frac{3}{2}}} d x\right)+\sin (x)\left(\int \frac{\cos (x)}{\cos (2 x)^{\frac{3}{2}}} d x\right)
\]
- Compute integrals
\[
y_{p}(x)=-\sqrt{\cos (2 x)}
\]
- Substitute particular solution into general solution to ODE \(y=c_{1} \cos (x)+c_{2} \sin (x)-\sqrt{\cos (2 x)}\)

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.016 (sec). Leaf size: 21
```

dsolve(diff(y(x),x\$2)+y(x)=1/(cos(2*x)*sqrt(cos(2*x))),y(x), singsol=all)

```
\[
y(x)=\sin (x) c_{2}+c_{1} \cos (x)-\sqrt{\cos (2 x)}
\]

Solution by Mathematica
Time used: 0.123 (sec). Leaf size: 26
DSolve[y' \([x]+y[x]==1 /(\operatorname{Cos}[2 * x] * \operatorname{Sqrt}[\operatorname{Cos}[2 * x]]), y[x], x\), IncludeSingularSolutions \(\rightarrow\) True]
\[
y(x) \rightarrow-\sqrt{\cos (2 x)}+c_{1} \cos (x)+c_{2} \sin (x)
\]

\subsection*{1.119 problem 170}
1.119.1 Solution using Matrix exponential method . . . . . . . . . . . . 1221
1.119.2 Solution using explicit Eigenvalue and Eigenvector method . . . 1223

Internal problem ID [12536]
Internal file name [OUTPUT/11188_Tuesday_October_17_2023_07_20_39_AM_92642001/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 170.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "system of linear ODEs"
Solve
\[
\begin{aligned}
x^{\prime}(t) & =y(t)+1 \\
y^{\prime}(t) & =x(t)+1
\end{aligned}
\]

With initial conditions
\[
[x(0)=-2, y(0)=0]
\]

\subsection*{1.119.1 Solution using Matrix exponential method}

In this method, we will assume we have found the matrix exponential \(e^{A t}\) allready. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)+\vec{G}(t)
\]

Or
\[
\left[\begin{array}{l}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]+\left[\begin{array}{l}
1 \\
1
\end{array}\right]
\]

Since the system is nonhomogeneous, then the solution is given by
\[
\vec{x}(t)=\vec{x}_{h}(t)+\vec{x}_{p}(t)
\]

Where \(\vec{x}_{h}(t)\) is the homogeneous solution to \(\vec{x}^{\prime}(t)=A \vec{x}(t)\) and \(\vec{x}_{p}(t)\) is a particular solution to \(\vec{x}^{\prime}(t)=A \vec{x}(t)+\vec{G}(t)\). The particular solution will be found using variation
of parameters method applied to the fundamental matrix. For the above matrix \(A\), the matrix exponential can be found to be
\[
e^{A t}=\left[\begin{array}{cc}
\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2} & \frac{\mathrm{e}^{t}}{2}-\frac{\mathrm{e}^{-t}}{2} \\
\frac{\mathrm{e}^{t}}{2}-\frac{\mathrm{e}^{-t}}{2} & \frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2}
\end{array}\right]
\]

Therefore the homogeneous solution is
\[
\begin{aligned}
\vec{x}_{h}(t) & =e^{A t} \vec{x}_{0} \\
& =\left[\begin{array}{cc}
\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2} & \frac{\mathrm{e}^{t}}{2}-\frac{\mathrm{e}^{-t}}{2} \\
\frac{\mathrm{e}^{t}}{2}-\frac{\mathrm{e}^{-t}}{2} & \frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2}
\end{array}\right]\left[\begin{array}{c}
-2 \\
0
\end{array}\right] \\
& =\left[\begin{array}{c}
-\mathrm{e}^{-t}-\mathrm{e}^{t} \\
-\mathrm{e}^{t}+\mathrm{e}^{-t}
\end{array}\right]
\end{aligned}
\]

The particular solution given by
\[
\vec{x}_{p}(t)=e^{A t} \int e^{-A t} \vec{G}(t) d t
\]

But
\[
\begin{aligned}
e^{-A t} & =\left(e^{A t}\right)^{-1} \\
& =\left[\begin{array}{cc}
\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2} & -\frac{\mathrm{e}^{t}}{2}+\frac{\mathrm{e}^{-t}}{2} \\
-\frac{\mathrm{e}^{t}}{2}+\frac{\mathrm{e}^{-t}}{2} & \frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2}
\end{array}\right]
\end{aligned}
\]

Hence
\[
\begin{aligned}
\vec{x}_{p}(t) & =\left[\begin{array}{ll}
\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2} & \frac{\mathrm{e}^{t}}{2}-\frac{\mathrm{e}^{-t}}{2} \\
\frac{\mathrm{e}^{t}}{2}-\frac{\mathrm{e}^{-t}}{2} & \frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2}
\end{array}\right] \int\left[\begin{array}{cc}
\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2} & -\frac{\mathrm{e}^{t}}{2}+\frac{\mathrm{e}^{-t}}{2} \\
-\frac{\mathrm{e}^{t}}{2}+\frac{\mathrm{e}^{-t}}{2} & \frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2}
\end{array}\right]\left[\begin{array}{l}
1 \\
1
\end{array}\right] d t \\
& =\left[\begin{array}{ll}
\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2} & \frac{\mathrm{e}^{t}}{2}-\frac{\mathrm{e}^{-t}}{2} \\
\frac{\mathrm{e}^{t}}{2}-\frac{\mathrm{e}^{-t}}{2} & \frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{t}}{2}
\end{array}\right]\left[\begin{array}{c}
-\mathrm{e}^{-t} \\
-\mathrm{e}^{-t}
\end{array}\right] \\
& =\left[\begin{array}{l}
-1 \\
-1
\end{array}\right]
\end{aligned}
\]

Hence the complete solution is
\[
\begin{aligned}
\vec{x}(t) & =\vec{x}_{h}(t)+\vec{x}_{p}(t) \\
& =\left[\begin{array}{c}
-\mathrm{e}^{-t}-\mathrm{e}^{t}-1 \\
-\mathrm{e}^{t}+\mathrm{e}^{-t}-1
\end{array}\right]
\end{aligned}
\]

\subsection*{1.119.2 Solution using explicit Eigenvalue and Eigenvector method}

This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)+\vec{G}(t)
\]

Or
\[
\left[\begin{array}{c}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]+\left[\begin{array}{l}
1 \\
1
\end{array}\right]
\]

Since the system is nonhomogeneous, then the solution is given by
\[
\vec{x}(t)=\vec{x}_{h}(t)+\vec{x}_{p}(t)
\]

Where \(\vec{x}_{h}(t)\) is the homogeneous solution to \(\vec{x}^{\prime}(t)=A \vec{x}(t)\) and \(\vec{x}_{p}(t)\) is a particular solution to \(\vec{x}^{\prime}(t)=A \vec{x}(t)+\vec{G}(t)\). The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of \(A\). This is done by solving the following equation for the eigenvalues \(\lambda\)
\[
\operatorname{det}(A-\lambda I)=0
\]

Expanding gives
\[
\operatorname{det}\left(\left[\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right]-\lambda\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)=0
\]

Therefore
\[
\operatorname{det}\left(\left[\begin{array}{cc}
-\lambda & 1 \\
1 & -\lambda
\end{array}\right]\right)=0
\]

Which gives the characteristic equation
\[
\lambda^{2}-1=0
\]

The roots of the above are the eigenvalues.
\[
\begin{aligned}
\lambda_{1} & =1 \\
\lambda_{2} & =-1
\end{aligned}
\]

This table summarises the above result
\begin{tabular}{|l|l|l|}
\hline eigenvalue & algebraic multiplicity & type of eigenvalue \\
\hline-1 & 1 & real eigenvalue \\
\hline 1 & 1 & real eigenvalue \\
\hline
\end{tabular}

Now the eigenvector for each eigenvalue are found.
Considering the eigenvalue \(\lambda_{1}=-1\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
&\left(\left[\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right]-(-1)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
& {\left[\begin{array}{ll}
1 & 1 \\
1 & 1
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] }
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{ll|l}
1 & 1 & 0 \\
1 & 1 & 0
\end{array}\right]} \\
R_{2}=R_{2}-R_{1} \Longrightarrow\left[\begin{array}{ll|l}
1 & 1 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{ll}
1 & 1 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=-t\right\}\)
Hence the solution is
\[
\left[\begin{array}{c}
-t \\
t
\end{array}\right]=\left[\begin{array}{c}
-t \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
-t \\
t
\end{array}\right]=t\left[\begin{array}{c}
-1 \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{c}
-t \\
t
\end{array}\right]=\left[\begin{array}{c}
-1 \\
1
\end{array}\right]
\]

Considering the eigenvalue \(\lambda_{2}=1\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
\left(\left[\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right]-(1)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right] & =\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
{\left[\begin{array}{cc}
-1 & 1 \\
1 & -1
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right] } & =\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{cc|c}
-1 & 1 & 0 \\
1 & -1 & 0
\end{array}\right]} \\
R_{2}=R_{2}+R_{1} \Longrightarrow\left[\begin{array}{cc|c}
-1 & 1 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{cc}
-1 & 1 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=t\right\}\)

Hence the solution is
\[
\left[\begin{array}{l}
t \\
t
\end{array}\right]=\left[\begin{array}{l}
t \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
t \\
t
\end{array}\right]=t\left[\begin{array}{l}
1 \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{l}
t \\
t
\end{array}\right]=\left[\begin{array}{l}
1 \\
1
\end{array}\right]
\]

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity \(m\), and its geometric multiplicity \(k\) and the eigenvectors associated with the eigenvalue. If \(m>k\) then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity \(k\) ) does not equal the algebraic multiplicity \(m\), and we need to determine an additional \(m-k\) generalized eigenvectors for this eigenvalue.
\begin{tabular}{|c|c|c|c|c|}
\hline \multirow{2}{*}{ eigenvalue } & \multicolumn{2}{|c|}{ multiplicity } & & \\
\cline { 2 - 3 } & algebraic \(m\) & geometric \(k\) & defective? & eigenvectors \\
\hline 1 & 1 & 1 & No & {\(\left[\begin{array}{c}1 \\
1\end{array}\right]\)} \\
\hline-1 & 1 & 1 & No & {\(\left[\begin{array}{c}-1 \\
1\end{array}\right]\)} \\
\hline
\end{tabular}

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue 1 is real and distinct then the corresponding eigenvector solution is
\[
\begin{aligned}
\vec{x}_{1}(t) & =\vec{v}_{1} e^{t} \\
& =\left[\begin{array}{l}
1 \\
1
\end{array}\right] e^{t}
\end{aligned}
\]

Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is
\[
\begin{aligned}
\vec{x}_{2}(t) & =\vec{v}_{2} e^{-t} \\
& =\left[\begin{array}{c}
-1 \\
1
\end{array}\right] e^{-t}
\end{aligned}
\]

Therefore the homogeneous solution is
\[
\vec{x}_{h}(t)=c_{1} \vec{x}_{1}(t)+c_{2} \vec{x}_{2}(t)
\]

Which is written as
\[
\left[\begin{array}{c}
x(t) \\
y(t)
\end{array}\right]=c_{1}\left[\begin{array}{c}
\mathrm{e}^{t} \\
\mathrm{e}^{t}
\end{array}\right]+c_{2}\left[\begin{array}{c}
-\mathrm{e}^{-t} \\
\mathrm{e}^{-t}
\end{array}\right]
\]

Now that we found homogeneous solution above, we need to find a particular solution \(\vec{x}_{p}(t)\). We will use Variation of parameters. The fundamental matrix is
\[
\Phi=\left[\begin{array}{lll}
\vec{x}_{1} & \vec{x}_{2} & \cdots
\end{array}\right]
\]

Where \(\vec{x}_{i}\) are the solution basis found above. Therefore the fundamental matrix is
\[
\Phi(t)=\left[\begin{array}{cc}
\mathrm{e}^{t} & -\mathrm{e}^{-t} \\
\mathrm{e}^{t} & \mathrm{e}^{-t}
\end{array}\right]
\]

The particular solution is then given by
\[
\vec{x}_{p}(t)=\Phi \int \Phi^{-1} \vec{G}(t) d t
\]

But
\[
\Phi^{-1}=\left[\begin{array}{cc}
\frac{\mathrm{e}^{-t}}{2} & \frac{\mathrm{e}^{-t}}{2} \\
-\frac{\mathrm{e}^{t}}{2} & \frac{\mathrm{e}^{t}}{2}
\end{array}\right]
\]

Hence
\[
\begin{aligned}
\vec{x}_{p}(t) & =\left[\begin{array}{cc}
\mathrm{e}^{t} & -\mathrm{e}^{-t} \\
\mathrm{e}^{t} & \mathrm{e}^{-t}
\end{array}\right] \int\left[\begin{array}{cc}
\frac{\mathrm{e}^{-t}}{2} & \frac{\mathrm{e}^{-t}}{2} \\
-\frac{\mathrm{e}^{t}}{2} & \frac{\mathrm{e}^{t}}{2}
\end{array}\right]\left[\begin{array}{l}
1 \\
1
\end{array}\right] d t \\
& =\left[\begin{array}{cc}
\mathrm{e}^{t} & -\mathrm{e}^{-t} \\
\mathrm{e}^{t} & \mathrm{e}^{-t}
\end{array}\right] \int\left[\begin{array}{c}
\mathrm{e}^{-t} \\
0
\end{array}\right] d t \\
& =\left[\begin{array}{cc}
\mathrm{e}^{t} & -\mathrm{e}^{-t} \\
\mathrm{e}^{t} & \mathrm{e}^{-t}
\end{array}\right]\left[\begin{array}{c}
-\mathrm{e}^{-t} \\
0
\end{array}\right] \\
& =\left[\begin{array}{l}
-1 \\
-1
\end{array}\right]
\end{aligned}
\]

Now that we found particular solution, the final solution is
\[
\begin{aligned}
\vec{x}(t) & =\vec{x}_{h}(t)+\vec{x}_{p}(t) \\
{\left[\begin{array}{c}
x(t) \\
y(t)
\end{array}\right] } & =\left[\begin{array}{c}
c_{1} \mathrm{e}^{t} \\
c_{1} \mathrm{e}^{t}
\end{array}\right]+\left[\begin{array}{c}
-c_{2} \mathrm{e}^{-t} \\
c_{2} \mathrm{e}^{-t}
\end{array}\right]+\left[\begin{array}{c}
-1 \\
-1
\end{array}\right]
\end{aligned}
\]

Which becomes
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
c_{1} \mathrm{e}^{t}-c_{2} \mathrm{e}^{-t}-1 \\
c_{1} \mathrm{e}^{t}+c_{2} \mathrm{e}^{-t}-1
\end{array}\right]
\]

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions
\[
\left[\begin{array}{c}
x(0)=-2  \tag{1}\\
y(0)=0
\end{array}\right]
\]

Substituting initial conditions into the above solution at \(t=0\) gives
\[
\left[\begin{array}{c}
-2 \\
0
\end{array}\right]=\left[\begin{array}{c}
c_{1}-c_{2}-1 \\
c_{1}+c_{2}-1
\end{array}\right]
\]

Solving for the constants of integrations gives
\[
\left[\begin{array}{l}
c_{1}=0 \\
c_{2}=1
\end{array}\right]
\]

Substituting these constants back in original solution in Eq. (1) gives
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
-1-\mathrm{e}^{-t} \\
\mathrm{e}^{-t}-1
\end{array}\right]
\]

The following is the phase plot of the system.


Figure 191: Phase plot

The following are plots of each solution.


\(\checkmark\) Solution by Maple
Time used: 0.015 (sec). Leaf size: 22
dsolve([diff \((x(t), t)=y(t)+1, \operatorname{diff}(y(t), t)=x(t)+1, x(0)=-2, y(0)=0]\), singsol=all)
\[
\begin{aligned}
& x(t)=-1-\mathrm{e}^{-t} \\
& y(t)=-1+\mathrm{e}^{-t}
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.014 (sec). Leaf size: 24
DSolve \(\left[\left\{x^{\prime}[t]==y[t]+1, y^{\prime}[t]==x[t]+1\right\},\{x[0]==-2, y[0]==0\},\{x[t], y[t]\}, t\right.\), IncludeSingularSolutio
\[
\begin{aligned}
x(t) & \rightarrow-e^{-t}-1 \\
y(t) & \rightarrow e^{-t}-1
\end{aligned}
\]

\subsection*{1.120 problem 171}
1.120.1 Solution using Matrix exponential method . . . . . . . . . . . . 1231
1.120.2 Solution using explicit Eigenvalue and Eigenvector method . . . 1232

Internal problem ID [12537]
Internal file name [OUTPUT/11189_Tuesday_October_17_2023_07_20_40_AM_64705552/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 171.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "system of linear ODEs"
Solve
\[
\begin{aligned}
x^{\prime}(t) & =x(t)-2 y(t) \\
y^{\prime}(t) & =x(t)-y(t)
\end{aligned}
\]

With initial conditions
\[
[x(0)=1, y(0)=1]
\]

\subsection*{1.120.1 Solution using Matrix exponential method}

In this method, we will assume we have found the matrix exponential \(e^{A t}\) allready. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)
\]

Or
\[
\left[\begin{array}{c}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{ll}
1 & -2 \\
1 & -1
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]

For the above matrix \(A\), the matrix exponential can be found to be
\[
e^{A t}=\left[\begin{array}{cc}
\cos (t)+\sin (t) & -2 \sin (t) \\
\sin (t) & \cos (t)-\sin (t)
\end{array}\right]
\]

Therefore the homogeneous solution is
\[
\begin{aligned}
\vec{x}_{h}(t) & =e^{A t} \vec{x}_{0} \\
& =\left[\begin{array}{cc}
\cos (t)+\sin (t) & -2 \sin (t) \\
\sin (t) & \cos (t)-\sin (t)
\end{array}\right]\left[\begin{array}{l}
1 \\
1
\end{array}\right] \\
& =\left[\begin{array}{c}
\cos (t)-\sin (t) \\
\cos (t)
\end{array}\right]
\end{aligned}
\]

Since no forcing function is given, then the final solution is \(\vec{x}_{h}(t)\) above.

\subsection*{1.120.2 Solution using explicit Eigenvalue and Eigenvector method}

This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)
\]

Or
\[
\left[\begin{array}{c}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{ll}
1 & -2 \\
1 & -1
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]

The first step is find the homogeneous solution. We start by finding the eigenvalues of \(A\). This is done by solving the following equation for the eigenvalues \(\lambda\)
\[
\operatorname{det}(A-\lambda I)=0
\]

Expanding gives
\[
\operatorname{det}\left(\left[\begin{array}{ll}
1 & -2 \\
1 & -1
\end{array}\right]-\lambda\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)=0
\]

Therefore
\[
\operatorname{det}\left(\left[\begin{array}{cc}
1-\lambda & -2 \\
1 & -1-\lambda
\end{array}\right]\right)=0
\]

Which gives the characteristic equation
\[
\lambda^{2}+1=0
\]

The roots of the above are the eigenvalues.
\[
\begin{aligned}
\lambda_{1} & =i \\
\lambda_{2} & =-i
\end{aligned}
\]

This table summarises the above result
\begin{tabular}{|l|l|l|}
\hline eigenvalue & algebraic multiplicity & type of eigenvalue \\
\hline\(i\) & 1 & complex eigenvalue \\
\hline\(-i\) & 1 & complex eigenvalue \\
\hline
\end{tabular}

Now the eigenvector for each eigenvalue are found.
Considering the eigenvalue \(\lambda_{1}=-i\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
&\left(\left[\begin{array}{ll}
1 & -2 \\
1 & -1
\end{array}\right]-(-i)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
& {\left[\begin{array}{cc}
1+i & -2 \\
1 & -1+i
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] }
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{cc|c}
1+i & -2 & 0 \\
1 & -1+i & 0
\end{array}\right]} \\
R_{2}=R_{2}+\left(-\frac{1}{2}+\frac{i}{2}\right) R_{1} \Longrightarrow\left[\begin{array}{cc|c}
1+i & -2 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{cc}
1+i & -2 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=(1-i) t\right\}\)

Hence the solution is
\[
\left[\begin{array}{c}
(1-\mathrm{I}) t \\
t
\end{array}\right]=\left[\begin{array}{c}
(1-i) t \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
(1-\mathrm{I}) t \\
t
\end{array}\right]=t\left[\begin{array}{c}
1-i \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{c}
(1-\mathrm{I}) t \\
t
\end{array}\right]=\left[\begin{array}{c}
1-i \\
1
\end{array}\right]
\]

Considering the eigenvalue \(\lambda_{2}=i\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
&\left(\left[\begin{array}{ll}
1 & -2 \\
1 & -1
\end{array}\right]-(i)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
& {\left[\begin{array}{cc}
1-i & -2 \\
1 & -1-i
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] }
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{cc|c}
1-i & -2 & 0 \\
1 & -1-i & 0
\end{array}\right]} \\
R_{2}=R_{2}+\left(-\frac{1}{2}-\frac{i}{2}\right) R_{1} \Longrightarrow\left[\begin{array}{cc|c}
1-i & -2 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{cc}
1-i & -2 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=(1+i) t\right\}\)

Hence the solution is
\[
\left[\begin{array}{c}
(1+\mathrm{I}) t \\
t
\end{array}\right]=\left[\begin{array}{c}
(1+i) t \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
(1+\mathrm{I}) t \\
t
\end{array}\right]=t\left[\begin{array}{c}
1+i \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{c}
(1+\mathrm{I}) t \\
t
\end{array}\right]=\left[\begin{array}{c}
1+i \\
1
\end{array}\right]
\]

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity \(m\), and its geometric multiplicity \(k\) and the eigenvectors associated with the eigenvalue. If \(m>k\) then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity \(k\) ) does not equal the algebraic multiplicity \(m\), and we need to determine an additional \(m-k\) generalized eigenvectors for this eigenvalue.
\begin{tabular}{|c|c|c|c|c|}
\hline \multirow{2}{*}{ eigenvalue } & \multicolumn{2}{|c|}{ multiplicity } & & \\
\cline { 2 - 3 } & algebraic \(m\) & geometric \(k\) & defective? & eigenvectors \\
\hline\(i\) & 1 & 1 & No & {\(\left[\begin{array}{c}1+i \\
1\end{array}\right]\)} \\
\hline\(-i\) & 1 & 1 & No & {\(\left[\begin{array}{c}1-i \\
1\end{array}\right]\)} \\
\hline
\end{tabular}

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is
\[
\vec{x}_{h}(t)=c_{1} \vec{x}_{1}(t)+c_{2} \vec{x}_{2}(t)
\]

Which is written as
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=c_{1}\left[\begin{array}{c}
(1+i) \mathrm{e}^{i t} \\
\mathrm{e}^{i t}
\end{array}\right]+c_{2}\left[\begin{array}{c}
(1-i) \mathrm{e}^{-i t} \\
\mathrm{e}^{-i t}
\end{array}\right]
\]

Which becomes
\[
\left[\begin{array}{c}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
(1+i) c_{1} \mathrm{e}^{i t}+(1-i) c_{2} \mathrm{e}^{-i t} \\
c_{1} \mathrm{e}^{i t}+c_{2} \mathrm{e}^{-i t}
\end{array}\right]
\]

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions
\[
\left[\begin{array}{l}
x(0)=1  \tag{1}\\
y(0)=1
\end{array}\right]
\]

Substituting initial conditions into the above solution at \(t=0\) gives
\[
\left[\begin{array}{l}
1 \\
1
\end{array}\right]=\left[\begin{array}{c}
(1+i) c_{1}+(1-i) c_{2} \\
c_{1}+c_{2}
\end{array}\right]
\]

Solving for the constants of integrations gives
\[
\left[\begin{array}{l}
c_{1}=\frac{1}{2} \\
c_{2}=\frac{1}{2}
\end{array}\right]
\]

Substituting these constants back in original solution in Eq. (1) gives
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
\left(\frac{1}{2}-\frac{i}{2}\right) \mathrm{e}^{-i t}+\left(\frac{1}{2}+\frac{i}{2}\right) \mathrm{e}^{i t} \\
\frac{\mathrm{e}^{i t}}{2}+\frac{\mathrm{e}^{-i t}}{2}
\end{array}\right]
\]

The following is the phase plot of the system.


Figure 192: Phase plot

The following are plots of each solution.
\(\checkmark\) Solution by Maple
Time used: 0.015 (sec). Leaf size: 17
```

dsolve([diff(x(t),t) = x (t)-2*y(t), diff(y(t),t) = x (t)-y(t), x(0) = 1, y(0) = 1], singsol=a

```
\[
\begin{aligned}
& x(t)=-\sin (t)+\cos (t) \\
& y(t)=\cos (t)
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.009 (sec). Leaf size: 17
DSolve \(\left[\left\{x^{\prime}[t]==x[t]-2 * y[t], y^{\prime}[t]==x[t]-y[t]\right\},\{x[0]==1, y[0]==1\},\{x[t], y[t]\}, t\right.\), IncludeSingular
\[
\begin{aligned}
& x(t) \rightarrow \cos (t)-\sin (t) \\
& y(t) \rightarrow \cos (t)
\end{aligned}
\]

\subsection*{1.121 problem 172}
1.121.1 Solution using Matrix exponential method . . . . . . . . . . . . 1238
1.121.2 Solution using explicit Eigenvalue and Eigenvector method . . . 1240
1.121.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1245

Internal problem ID [12538]
Internal file name [OUTPUT/11190_Tuesday_October_17_2023_07_20_40_AM_61324329/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 172.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "system of linear ODEs"
Solve
\[
\begin{aligned}
x^{\prime}(t) & =-y(t)+\cos (t) \\
y^{\prime}(t) & =-4 y(t)+4 \cos (t)+3 x(t)-\sin (t)
\end{aligned}
\]

\subsection*{1.121.1 Solution using Matrix exponential method}

In this method, we will assume we have found the matrix exponential \(e^{A t}\) allready. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)+\vec{G}(t)
\]

Or
\[
\left[\begin{array}{c}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{ll}
0 & -1 \\
3 & -4
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]+\left[\begin{array}{c}
\cos (t) \\
4 \cos (t)-\sin (t)
\end{array}\right]
\]

Since the system is nonhomogeneous, then the solution is given by
\[
\vec{x}(t)=\vec{x}_{h}(t)+\vec{x}_{p}(t)
\]

Where \(\vec{x}_{h}(t)\) is the homogeneous solution to \(\vec{x}^{\prime}(t)=A \vec{x}(t)\) and \(\vec{x}_{p}(t)\) is a particular solution to \(\vec{x}^{\prime}(t)=A \vec{x}(t)+\vec{G}(t)\). The particular solution will be found using variation
of parameters method applied to the fundamental matrix. For the above matrix \(A\), the matrix exponential can be found to be
\[
e^{A t}=\left[\begin{array}{cc}
\frac{3 \mathrm{e}^{-t}}{2}-\frac{\mathrm{e}^{-3 t}}{2} & -\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{-3 t}}{2} \\
\frac{3 \mathrm{e}^{-t}}{2}-\frac{3 \mathrm{e}^{-3 t}}{2} & \frac{3 \mathrm{e}^{-3 t}}{2}-\frac{\mathrm{e}^{-t}}{2}
\end{array}\right]
\]

Therefore the homogeneous solution is
\[
\begin{aligned}
\vec{x}_{h}(t) & =e^{A t} \vec{c} \\
& =\left[\begin{array}{cc}
\frac{3 \mathrm{e}^{-t}}{2}-\frac{\mathrm{e}^{-3 t}}{2} & -\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{-3 t}}{2} \\
\frac{3 \mathrm{e}^{-t}}{2}-\frac{3 \mathrm{e}^{-3 t}}{2} & \frac{3 \mathrm{e}^{-3 t}}{2}-\frac{\mathrm{e}^{-t}}{2}
\end{array}\right]\left[\begin{array}{c}
c_{1} \\
c_{2}
\end{array}\right] \\
& =\left[\begin{array}{c}
\left(\frac{3 \mathrm{e}^{-t}}{2}-\frac{\mathrm{e}^{-3 t}}{2}\right) c_{1}+\left(-\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{-3 t}}{2}\right) c_{2} \\
\left(\frac{3 \mathrm{e}^{-t}}{2}-\frac{3 \mathrm{e}^{-3 t}}{2}\right) c_{1}+\left(\frac{3 \mathrm{e}^{-3 t}}{2}-\frac{\mathrm{e}^{-t}}{2}\right) c_{2}
\end{array}\right] \\
& =\left[\begin{array}{c}
\frac{\left(-c_{1}+c_{2}\right) \mathrm{e}^{-3 t}}{2}+\frac{3\left(c_{1}-\frac{c_{2}}{3}\right) \mathrm{e}^{-t}}{2} \\
\frac{\left(-3 c_{1}+3 c_{2}\right) \mathrm{e}^{-3 t}}{2}+\frac{3\left(c_{1}-\frac{c_{2}}{3}\right) \mathrm{e}^{-t}}{2}
\end{array}\right]
\end{aligned}
\]

The particular solution given by
\[
\vec{x}_{p}(t)=e^{A t} \int e^{-A t} \vec{G}(t) d t
\]

But
\[
\begin{aligned}
e^{-A t} & =\left(e^{A t}\right)^{-1} \\
& =\left[\begin{array}{cc}
-\frac{\left(\mathrm{e}^{2 t}-3\right) \mathrm{e}^{t}}{2} & \frac{\left(\mathrm{e}^{2 t}-1\right) \mathrm{e}^{t}}{2} \\
-\frac{3\left(\mathrm{e}^{2 t}-1\right) \mathrm{e}^{t}}{2} & \frac{\left(3 \mathrm{e}^{2 t}-1\right) \mathrm{e}^{t}}{2}
\end{array}\right]
\end{aligned}
\]

Hence
\[
\begin{aligned}
\vec{x}_{p}(t) & =\left[\begin{array}{cc}
\frac{3 \mathrm{e}^{-t}}{2}-\frac{\mathrm{e}^{-3 t}}{2} & -\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{-3 t}}{2} \\
\frac{3 \mathrm{e}^{-t}}{2}-\frac{3 \mathrm{e}^{-3 t}}{2} & \frac{3 \mathrm{e}^{-3 t}}{2}-\frac{\mathrm{e}^{-t}}{2}
\end{array}\right] \int\left[\begin{array}{cc}
-\frac{\left(\mathrm{e}^{2 t}-3\right) \mathrm{e}^{t}}{2} & \frac{\left(\mathrm{e}^{2 t}-1\right) \mathrm{e}^{t}}{2} \\
-\frac{3\left(\mathrm{e}^{2 t}-1\right) \mathrm{e}^{t}}{2} & \frac{\left(3 \mathrm{e}^{2 t}-1\right) \mathrm{e}^{t}}{2}
\end{array}\right]\left[\begin{array}{c}
\cos (t) \\
4 \cos (t)-\sin (t)
\end{array}\right] d t \\
& =\left[\begin{array}{cc}
\frac{3 \mathrm{e}^{-t}}{2}-\frac{\mathrm{e}^{-3 t}}{2} & -\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{-3 t}}{2} \\
\frac{3 \mathrm{e}^{-t}}{2}-\frac{3 \mathrm{e}^{-3 t}}{2} & \frac{3 \mathrm{e}^{-3 t}}{2}-\frac{\mathrm{e}^{-t}}{2}
\end{array}\right]\left[\begin{array}{c}
\frac{\left(\mathrm{e}^{2 t}-1\right) \mathrm{e}^{t} \cos (t)}{2} \\
\frac{\cos (t)\left(3 \mathrm{e}^{2 t}-1\right) \mathrm{e}^{t}}{2}
\end{array}\right] \\
& =\left[\begin{array}{c}
0 \\
\cos (t)
\end{array}\right]
\end{aligned}
\]

Hence the complete solution is
\[
\begin{aligned}
\vec{x}(t) & =\vec{x}_{h}(t)+\vec{x}_{p}(t) \\
& =\left[\begin{array}{c}
\frac{\left(-c_{1}+c_{2}\right) \mathrm{e}^{-3 t}}{2}+\frac{3\left(c_{1}-\frac{c_{2}}{3}\right) \mathrm{e}^{-t}}{2} \\
\frac{\left(-3 c_{1}+3 c_{2}\right) \mathrm{e}^{-3 t}}{2}+\frac{\left(3 c_{1}-c_{2}\right) \mathrm{e}^{-t}}{2}+\cos (t)
\end{array}\right]
\end{aligned}
\]

\subsection*{1.121.2 Solution using explicit Eigenvalue and Eigenvector method}

This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)+\vec{G}(t)
\]

Or
\[
\left[\begin{array}{c}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{ll}
0 & -1 \\
3 & -4
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]+\left[\begin{array}{c}
\cos (t) \\
4 \cos (t)-\sin (t)
\end{array}\right]
\]

Since the system is nonhomogeneous, then the solution is given by
\[
\vec{x}(t)=\vec{x}_{h}(t)+\vec{x}_{p}(t)
\]

Where \(\vec{x}_{h}(t)\) is the homogeneous solution to \(\vec{x}^{\prime}(t)=A \vec{x}(t)\) and \(\vec{x}_{p}(t)\) is a particular solution to \(\vec{x}^{\prime}(t)=A \vec{x}(t)+\vec{G}(t)\). The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of \(A\). This is done by solving the following equation for the eigenvalues \(\lambda\)
\[
\operatorname{det}(A-\lambda I)=0
\]

Expanding gives
\[
\operatorname{det}\left(\left[\begin{array}{ll}
0 & -1 \\
3 & -4
\end{array}\right]-\lambda\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)=0
\]

Therefore
\[
\operatorname{det}\left(\left[\begin{array}{cc}
-\lambda & -1 \\
3 & -4-\lambda
\end{array}\right]\right)=0
\]

Which gives the characteristic equation
\[
\lambda^{2}+4 \lambda+3=0
\]

The roots of the above are the eigenvalues.
\[
\begin{aligned}
& \lambda_{1}=-1 \\
& \lambda_{2}=-3
\end{aligned}
\]

This table summarises the above result
\begin{tabular}{|l|l|l|}
\hline eigenvalue & algebraic multiplicity & type of eigenvalue \\
\hline-1 & 1 & real eigenvalue \\
\hline-3 & 1 & real eigenvalue \\
\hline
\end{tabular}

Now the eigenvector for each eigenvalue are found.
Considering the eigenvalue \(\lambda_{1}=-3\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{array}{r}
\left(\left[\begin{array}{cc}
0 & -1 \\
3 & -4
\end{array}\right]-(-3)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
{\left[\begin{array}{ll}
3 & -1 \\
3 & -1
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]}
\end{array}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{ll|l}
3 & -1 & 0 \\
3 & -1 & 0
\end{array}\right]} \\
R_{2}=R_{2}-R_{1} \Longrightarrow\left[\begin{array}{cc|c}
3 & -1 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{cc}
3 & -1 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=\frac{t}{3}\right\}\)

Hence the solution is
\[
\left[\begin{array}{c}
\frac{t}{3} \\
t
\end{array}\right]=\left[\begin{array}{c}
\frac{t}{3} \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
\frac{t}{3} \\
t
\end{array}\right]=t\left[\begin{array}{c}
\frac{1}{3} \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{c}
\frac{t}{3} \\
t
\end{array}\right]=\left[\begin{array}{c}
\frac{1}{3} \\
1
\end{array}\right]
\]

Which is normalized to
\[
\left[\begin{array}{c}
\frac{t}{3} \\
t
\end{array}\right]=\left[\begin{array}{l}
1 \\
3
\end{array}\right]
\]

Considering the eigenvalue \(\lambda_{2}=-1\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
&\left(\left[\begin{array}{ll}
0 & -1 \\
3 & -4
\end{array}\right]-(-1)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
& {\left[\begin{array}{ll}
1 & -1 \\
3 & -3
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] }
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{ll|l}
1 & -1 & 0 \\
3 & -3 & 0
\end{array}\right]} \\
R_{2}=R_{2}-3 R_{1} \Longrightarrow\left[\begin{array}{cc|c}
1 & -1 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{cc}
1 & -1 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=t\right\}\)

Hence the solution is
\[
\left[\begin{array}{l}
t \\
t
\end{array}\right]=\left[\begin{array}{l}
t \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{l}
t \\
t
\end{array}\right]=t\left[\begin{array}{l}
1 \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{l}
t \\
t
\end{array}\right]=\left[\begin{array}{l}
1 \\
1
\end{array}\right]
\]

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity \(m\), and its geometric multiplicity \(k\) and the eigenvectors associated with the eigenvalue. If \(m>k\) then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity \(k\) ) does not equal the algebraic multiplicity \(m\), and we need to determine an additional \(m-k\) generalized eigenvectors for this eigenvalue.
\begin{tabular}{|c|c|c|c|c|}
\hline \multirow{2}{*}{ eigenvalue } & \multicolumn{2}{|c|}{ multiplicity } & & \\
\cline { 2 - 3 } & algebraic \(m\) & geometric \(k\) & defective? & eigenvectors \\
\hline-1 & 1 & 1 & No & {\(\left[\begin{array}{l}1 \\
1\end{array}\right]\)} \\
\hline-3 & 1 & 1 & No & {\(\left[\begin{array}{c}\frac{1}{3} \\
1\end{array}\right]\)} \\
\hline
\end{tabular}

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -1 is real and distinct then the
corresponding eigenvector solution is
\[
\begin{aligned}
\vec{x}_{1}(t) & =\vec{v}_{1} e^{-t} \\
& =\left[\begin{array}{l}
1 \\
1
\end{array}\right] e^{-t}
\end{aligned}
\]

Since eigenvalue -3 is real and distinct then the corresponding eigenvector solution is
\[
\begin{aligned}
\vec{x}_{2}(t) & =\vec{v}_{2} e^{-3 t} \\
& =\left[\begin{array}{c}
\frac{1}{3} \\
1
\end{array}\right] e^{-3 t}
\end{aligned}
\]

Therefore the homogeneous solution is
\[
\vec{x}_{h}(t)=c_{1} \vec{x}_{1}(t)+c_{2} \vec{x}_{2}(t)
\]

Which is written as
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=c_{1}\left[\begin{array}{l}
\mathrm{e}^{-t} \\
\mathrm{e}^{-t}
\end{array}\right]+c_{2}\left[\begin{array}{c}
\frac{\mathrm{e}^{-3 t}}{3} \\
\mathrm{e}^{-3 t}
\end{array}\right]
\]

Now that we found homogeneous solution above, we need to find a particular solution \(\vec{x}_{p}(t)\). We will use Variation of parameters. The fundamental matrix is
\[
\Phi=\left[\begin{array}{lll}
\vec{x}_{1} & \vec{x}_{2} & \cdots
\end{array}\right]
\]

Where \(\vec{x}_{i}\) are the solution basis found above. Therefore the fundamental matrix is
\[
\Phi(t)=\left[\begin{array}{cc}
\mathrm{e}^{-t} & \frac{\mathrm{e}^{-3 t}}{3} \\
\mathrm{e}^{-t} & \mathrm{e}^{-3 t}
\end{array}\right]
\]

The particular solution is then given by
\[
\vec{x}_{p}(t)=\Phi \int \Phi^{-1} \vec{G}(t) d t
\]

But
\[
\Phi^{-1}=\left[\begin{array}{cc}
\frac{3 \mathrm{e}^{t}}{2} & -\frac{\mathrm{e}^{t}}{2} \\
-\frac{3 \mathrm{e}^{3 t}}{2} & \frac{3 \mathrm{e}^{3 t}}{2}
\end{array}\right]
\]

Hence
\[
\begin{aligned}
\vec{x}_{p}(t) & =\left[\begin{array}{ll}
\mathrm{e}^{-t} & \frac{\mathrm{e}^{-3 t}}{3} \\
\mathrm{e}^{-t} & \mathrm{e}^{-3 t}
\end{array}\right] \int\left[\begin{array}{cc}
\frac{3 \mathrm{e}^{t}}{2} & -\frac{\mathrm{e}^{t}}{2} \\
-\frac{3 \mathrm{e}^{3 t}}{2} & \frac{3 \mathrm{e}^{3 t}}{2}
\end{array}\right]\left[\begin{array}{c}
\cos (t) \\
4 \cos (t)-\sin (t)
\end{array}\right] d t \\
& =\left[\begin{array}{ll}
\mathrm{e}^{-t} & \frac{\mathrm{e}^{-3 t}}{3} \\
\mathrm{e}^{-t} & \mathrm{e}^{-3 t}
\end{array}\right] \int\left[\begin{array}{c}
-\frac{(\cos (t)-\sin (t)) \mathrm{e}^{t}}{2} \\
\frac{3 \mathrm{e}^{3 t}(3 \cos (t)-\sin (t))}{2}
\end{array}\right] d t \\
& =\left[\begin{array}{cc}
\mathrm{e}^{-t} & \frac{\mathrm{e}^{-3 t}}{3} \\
\mathrm{e}^{-t} & \mathrm{e}^{-3 t}
\end{array}\right]\left[\begin{array}{c}
-\frac{\mathrm{e}^{t} \cos (t)}{2} \\
\frac{3 \mathrm{e}^{3 t} \cos (t)}{2}
\end{array}\right] \\
& =\left[\begin{array}{c}
0 \\
\cos (t)
\end{array}\right]
\end{aligned}
\]

Now that we found particular solution, the final solution is
\[
\begin{aligned}
\vec{x}(t) & =\vec{x}_{h}(t)+\vec{x}_{p}(t) \\
{\left[\begin{array}{c}
x(t) \\
y(t)
\end{array}\right] } & =\left[\begin{array}{c}
c_{1} \mathrm{e}^{-t} \\
c_{1} \mathrm{e}^{-t}
\end{array}\right]+\left[\begin{array}{c}
\frac{c_{2} \mathrm{e}^{-3 t}}{3} \\
c_{2} \mathrm{e}^{-3 t}
\end{array}\right]+\left[\begin{array}{c}
0 \\
\cos (t)
\end{array}\right]
\end{aligned}
\]

Which becomes
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
c_{1} \mathrm{e}^{-t}+\frac{c_{2} \mathrm{e}^{-3 t}}{3} \\
c_{1} \mathrm{e}^{-t}+c_{2} \mathrm{e}^{-3 t}+\cos (t)
\end{array}\right]
\]

\subsection*{1.121.3 Maple step by step solution}

Let's solve
\[
\left[x^{\prime}(t)=-y(t)+\cos (t), y^{\prime}(t)=-4 y(t)+4 \cos (t)+3 x(t)-\sin (t)\right]
\]
- Define vector
\[
\vec{x}(t)=\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]
- Convert system into a vector equation
\[
\vec{x}^{\prime}(t)=\left[\begin{array}{ll}
0 & -1 \\
3 & -4
\end{array}\right] \cdot \vec{x}(t)+\left[\begin{array}{c}
\cos (t) \\
4 \cos (t)-\sin (t)
\end{array}\right]
\]
- System to solve
\[
\vec{x}^{\prime}(t)=\left[\begin{array}{ll}
0 & -1 \\
3 & -4
\end{array}\right] \cdot \vec{x}(t)+\left[\begin{array}{c}
\cos (t) \\
4 \cos (t)-\sin (t)
\end{array}\right]
\]
- Define the forcing function
\(\vec{f}(t)=\left[\begin{array}{c}\cos (t) \\ 4 \cos (t)-\sin (t)\end{array}\right]\)
- Define the coefficient matrix
\[
A=\left[\begin{array}{ll}
0 & -1 \\
3 & -4
\end{array}\right]
\]
- Rewrite the system as
\(\vec{x}^{\prime}(t)=A \cdot \vec{x}(t)+\vec{f}\)
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- \(\quad\) Eigenpairs of \(A\)
\[
\left[\left[-3,\left[\begin{array}{c}
\frac{1}{3} \\
1
\end{array}\right]\right],\left[-1,\left[\begin{array}{l}
1 \\
1
\end{array}\right]\right]\right]
\]
- Consider eigenpair
\[
\left[-3,\left[\begin{array}{c}
\frac{1}{3} \\
1
\end{array}\right]\right]
\]
- Solution to homogeneous system from eigenpair
\[
\vec{x}_{1}=\mathrm{e}^{-3 t} \cdot\left[\begin{array}{c}
\frac{1}{3} \\
1
\end{array}\right]
\]
- Consider eigenpair
\[
\left[-1,\left[\begin{array}{l}
1 \\
1
\end{array}\right]\right]
\]
- Solution to homogeneous system from eigenpair
\(\vec{x}_{2}=\mathrm{e}^{-t} \cdot\left[\begin{array}{l}1 \\ 1\end{array}\right]\)
- General solution of the system of ODEs can be written in terms of the particular solution \(\vec{x}_{p}(\) \(\vec{x}(t)=c_{1} \vec{x}_{1}+c_{2} \vec{x}_{2}+\vec{x}_{p}(t)\)

\section*{Fundamental matrix}
- Let \(\phi(t)\) be the matrix whose columns are the independent solutions of the homogeneous syst \(\phi(t)=\left[\begin{array}{cc}\frac{\mathrm{e}^{-3 t}}{3} & \mathrm{e}^{-t} \\ \mathrm{e}^{-3 t} & \mathrm{e}^{-t}\end{array}\right]\)
- The fundamental matrix, \(\Phi(t)\) is a normalized version of \(\phi(t)\) satisfying \(\Phi(0)=I\) where \(I\) is th
\[
\Phi(t)=\phi(t) \cdot \frac{1}{\phi(0)}
\]
- Substitute the value of \(\phi(t)\) and \(\phi(0)\)
\[
\Phi(t)=\left[\begin{array}{cc}
\frac{\mathrm{e}^{-3 t}}{3} & \mathrm{e}^{-t} \\
\mathrm{e}^{-3 t} & \mathrm{e}^{-t}
\end{array}\right] \cdot \frac{1}{\left[\begin{array}{cc}
\frac{1}{3} & 1 \\
1 & 1
\end{array}\right]}
\]
- Evaluate and simplify to get the fundamental matrix
\[
\Phi(t)=\left[\begin{array}{cc}
\frac{3 \mathrm{e}^{-t}}{2}-\frac{\mathrm{e}^{-3 t}}{2} & -\frac{\mathrm{e}^{-t}}{2}+\frac{\mathrm{e}^{-3 t}}{2} \\
\frac{3 \mathrm{e}^{-t}}{2}-\frac{3 \mathrm{e}^{-3 t}}{2} & \frac{3 \mathrm{e}^{-3 t}}{2}-\frac{\mathrm{e}^{-t}}{2}
\end{array}\right]
\]

Find a particular solution of the system of ODEs using variation of parameters
- Let the particular solution be the fundamental matrix multiplied by \(\vec{v}(t)\) and solve for \(\vec{v}(t)\) \(\vec{x}_{p}(t)=\Phi(t) \cdot \vec{v}(t)\)
- Take the derivative of the particular solution
\(\vec{x}_{p}^{\prime}(t)=\Phi^{\prime}(t) \cdot \vec{v}(t)+\Phi(t) \cdot \vec{v}^{\prime}(t)\)
- Substitute particular solution and its derivative into the system of ODEs \(\Phi^{\prime}(t) \cdot \vec{v}(t)+\Phi(t) \cdot \vec{v}^{\prime}(t)=A \cdot \Phi(t) \cdot \vec{v}(t)+\vec{f}(t)\)
- The fundamental matrix has columns that are solutions to the homogeneous system so its der
\[
A \cdot \Phi(t) \cdot \vec{v}(t)+\Phi(t) \cdot \vec{v}^{\prime}(t)=A \cdot \Phi(t) \cdot \vec{v}(t)+\vec{f}(t)
\]
- Cancel like terms
\(\Phi(t) \cdot \vec{v}^{\prime}(t)=\vec{f}(t)\)
- Multiply by the inverse of the fundamental matrix
\[
\vec{v}^{\prime}(t)=\frac{1}{\Phi(t)} \cdot \vec{f}(t)
\]
- Integrate to solve for \(\vec{v}(t)\)
\[
\vec{v}(t)=\int_{0}^{t} \frac{1}{\Phi(s)} \cdot \vec{f}(s) d s
\]
- Plug \(\vec{v}(t)\) into the equation for the particular solution
\[
\vec{x}_{p}(t)=\Phi(t) \cdot\left(\int_{0}^{t} \frac{1}{\Phi(s)} \cdot \vec{f}(s) d s\right)
\]
- Plug in the fundamental matrix and the forcing function and compute
\[
\vec{x}_{p}(t)=\left[\begin{array}{c}
\frac{\mathrm{e}^{-t}}{2}-\frac{\mathrm{e}^{-3 t}}{2} \\
\cos (t)+\frac{\mathrm{e}^{-t}}{2}-\frac{3 \mathrm{e}^{-3 t}}{2}
\end{array}\right]
\]
- Plug particular solution back into general solution
\[
\vec{x}(t)=c_{1} \vec{x}_{1}+c_{2} \vec{x}_{2}+\left[\begin{array}{c}
\frac{\mathrm{e}^{-t}}{2}-\frac{\mathrm{e}^{-3 t}}{2} \\
\cos (t)+\frac{\mathrm{e}^{-t}}{2}-\frac{3 \mathrm{e}^{-3 t}}{2}
\end{array}\right]
\]
- \(\quad\) Substitute in vector of dependent variables
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
\frac{\left(2 c_{1}-3\right) \mathrm{e}^{-3 t}}{6}+\frac{\left(6 c_{2}+3\right) \mathrm{e}^{-t}}{6} \\
\frac{\left(2 c_{1}-3\right) \mathrm{e}^{-3 t}}{2}+\frac{\left(2 c_{2}+1\right) \mathrm{e}^{-t}}{2}+\cos (t)
\end{array}\right]
\]
- Solution to the system of ODEs
\[
\left\{x(t)=\frac{\left(2 c_{1}-3\right) \mathrm{e}^{-3 t}}{6}+\frac{\left(6 c_{2}+3\right) \mathrm{e}^{-t}}{6}, y(t)=\frac{\left(2 c_{1}-3\right) \mathrm{e}^{-3 t}}{2}+\frac{\left(2 c_{2}+1\right) \mathrm{e}^{-t}}{2}+\cos (t)\right\}
\]
\(\checkmark\) Solution by Maple
Time used: 0.078 (sec). Leaf size: 37
dsolve([4*diff \((x(t), t)-\operatorname{diff}(y(t), t)+3 * x(t)=\sin (t), \operatorname{diff}(x(t), t)+y(t)=\cos (t)]\), singsol=all)
\[
\begin{aligned}
& x(t)=\frac{c_{2} \mathrm{e}^{-3 t}}{3}+\mathrm{e}^{-t} c_{1} \\
& y(t)=c_{2} \mathrm{e}^{-3 t}+\mathrm{e}^{-t} c_{1}+\cos (t)
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.352 (sec). Leaf size: 76
DSolve \(\left[\left\{4 * x^{\prime}[t]-y^{\prime}[t]+3 * x[t]==\operatorname{Sin}[t], x^{\prime}[t]+y[t]==\operatorname{Cos}[t]\right\},\{x[t], y[t]\}, t\right.\), IncludeSingularSoluti
\[
\begin{aligned}
& x(t) \rightarrow \frac{1}{2} e^{-3 t}\left(c_{1}\left(3 e^{2 t}-1\right)-c_{2}\left(e^{2 t}-1\right)\right) \\
& y(t) \rightarrow \cos (t)+\frac{1}{2} e^{-3 t}\left(3 c_{1}\left(e^{2 t}-1\right)-c_{2}\left(e^{2 t}-3\right)\right)
\end{aligned}
\]

\subsection*{1.122 problem 181}
1.122.1 Solving as second order ode missing x ode
1.122.2 Maple step by step solution 1253

Internal problem ID [12539]
Internal file name [OUTPUT/11191_Wednesday_October_18_2023_03_46_59_AM_64086885/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 181.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "second_order_ode_missing_x"
Maple gives the following as the ode type
[[_2nd_order, _missing_x], [_2nd_order, _reducible, _mu_x_y1]]
\[
y y^{\prime \prime}-y^{\prime 2}=1
\]

\subsection*{1.122.1 Solving as second order ode missing \(x\) ode}

This is missing independent variable second order ode. Solved by reduction of order by using substitution which makes the dependent variable \(y\) an independent variable. Using
\[
y^{\prime}=p(y)
\]

Then
\[
\begin{aligned}
y^{\prime \prime} & =\frac{d p}{d x} \\
& =\frac{d y}{d x} \frac{d p}{d y} \\
& =p \frac{d p}{d y}
\end{aligned}
\]

Hence the ode becomes
\[
y p(y)\left(\frac{d}{d y} p(y)\right)-p(y)^{2}=1
\]

Which is now solved as first order ode for \(p(y)\). In canonical form the ODE is
\[
\begin{aligned}
p^{\prime} & =F(y, p) \\
& =f(y) g(p) \\
& =\frac{p^{2}+1}{y p}
\end{aligned}
\]

Where \(f(y)=\frac{1}{y}\) and \(g(p)=\frac{p^{2}+1}{p}\). Integrating both sides gives
\[
\begin{aligned}
\frac{1}{\frac{p^{2}+1}{p}} d p & =\frac{1}{y} d y \\
\int \frac{1}{\frac{p^{2}+1}{p}} d p & =\int \frac{1}{y} d y \\
\frac{\ln \left(p^{2}+1\right)}{2} & =\ln (y)+c_{1}
\end{aligned}
\]

Raising both side to exponential gives
\[
\sqrt{p^{2}+1}=\mathrm{e}^{\ln (y)+c_{1}}
\]

Which simplifies to
\[
\sqrt{p^{2}+1}=c_{2} y
\]

Which simplifies to
\[
\sqrt{p(y)^{2}+1}=c_{2} y \mathrm{e}^{c_{1}}
\]

The solution is
\[
\sqrt{p(y)^{2}+1}=c_{2} y \mathrm{e}^{c_{1}}
\]

For solution (1) found earlier, since \(p=y^{\prime}\) then we now have a new first order ode to solve which is
\[
\sqrt{y^{\prime 2}+1}=c_{2} y \mathrm{e}^{c_{1}}
\]

Solving the given ode for \(y^{\prime}\) results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are
\[
\begin{align*}
& y^{\prime}=\sqrt{-1+c_{2}^{2} \mathrm{e}^{2 c_{1}} y^{2}}  \tag{1}\\
& y^{\prime}=-\sqrt{-1+c_{2}^{2} \mathrm{e}^{2 c_{1}} y^{2}} \tag{2}
\end{align*}
\]

Now each one of the above ODE is solved.
Solving equation (1)
Integrating both sides gives
\[
\begin{aligned}
\int \frac{1}{\sqrt{-1+c_{2}^{2} \mathrm{e}^{2 c_{1}} y^{2}}} d y & =\int d x \\
\frac{\ln \left(\frac{c_{2}^{2} \mathrm{e}^{2 c_{1}} y}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}+\sqrt{-1+c_{2}^{2} \mathrm{e}^{2 c_{1}} y^{2}}\right)}{\sqrt{c_{2}^{2} \mathrm{e}^{\mathrm{e}_{1}}}} & =c_{3}+x
\end{aligned}
\]

Raising both side to exponential gives
\[
\mathrm{e}^{\frac{\ln \left(\frac{c_{2}^{2} \mathrm{e}^{2 c_{1}} y}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}+\sqrt{-1+c_{2}^{2} \mathrm{e}^{2 c_{1} y^{2}}}\right)}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}}=\mathrm{e}^{c_{3}+x}
\]

Which simplifies to
\[
\left(\frac{c_{2}^{2} \mathrm{e}^{2 c_{1}} y+\sqrt{-1+c_{2}^{2} \mathrm{e}^{2 c_{1}} y^{2}} \sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}\right)^{\frac{1}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}}=\mathrm{e}^{x} c_{4}
\]

Solving equation (2)
Integrating both sides gives
\[
\begin{aligned}
\int-\frac{1}{\sqrt{-1+c_{2}^{2} \mathrm{e}^{2 c_{1}} y^{2}}} d y & =\int d x \\
-\frac{\ln \left(\frac{c_{2}^{2} e^{2 c_{1}} y}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}+\sqrt{-1+c_{2}^{2} \mathrm{e}^{2 c_{1}} y^{2}}\right)}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}} & =x+c_{5}
\end{aligned}
\]

Raising both side to exponential gives
\[
\mathrm{e}^{-\frac{\ln \left(\frac{c_{2}^{2} \mathrm{e}^{2 c_{1} y}}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}+\sqrt{-1+c_{2}^{2} \mathrm{e}^{2 c_{1} y^{2}}}\right)}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}}=\mathrm{e}^{x+c_{5}}
\]

Which simplifies to
\[
\left(\frac{c_{2}^{2} \mathrm{e}^{2 c_{1}} y+\sqrt{-1+c_{2}^{2} \mathrm{e}^{2 c_{1}} y^{2}} \sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}\right)^{-\frac{1}{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}}=c_{6} \mathrm{e}^{x}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{align*}
& y=\frac{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}} \mathrm{e}^{-2 c_{1}}\left(\left(\mathrm{e}^{x} c_{4}\right)^{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}+\left(\mathrm{e}^{x} c_{4}\right)^{-\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}\right)}{2 c_{2}^{2}}  \tag{1}\\
& y=\frac{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}} \mathrm{e}^{-2 c_{1}}\left(\left(c_{6} \mathrm{e}^{x}\right)^{-\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}+\left(c_{6} \mathrm{e}^{x}\right)^{\sqrt{c_{2}^{2} \mathrm{e}^{c_{1}}}}\right)}{2 c_{2}^{2}} \tag{2}
\end{align*}
\]

Verification of solutions
\[
y=\frac{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}} \mathrm{e}^{-2 c_{1}}\left(\left(\mathrm{e}^{x} c_{4}\right)^{\sqrt{c_{2}^{2} \mathrm{e}^{c_{1}}}}+\left(\mathrm{e}^{x} c_{4}\right)^{-\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}\right)}{2 c_{2}^{2}}
\]

Verified OK.
\[
y=\frac{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}} \mathrm{e}^{-2 c_{1}}\left(\left(c_{6} \mathrm{e}^{x}\right)^{-\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}+\left(c_{6} \mathrm{e}^{x}\right)^{\sqrt{c_{2}^{2} \mathrm{e}^{2 c_{1}}}}\right)}{2 c_{2}^{2}}
\]

Verified OK.

\subsection*{1.122.2 Maple step by step solution}

Let's solve
\(y y^{\prime \prime}-y^{\prime 2}=1\)
- Highest derivative means the order of the ODE is 2
\(y^{\prime \prime}\)
- \(\quad\) Define new dependent variable \(u\)
\[
u(x)=y^{\prime}
\]
- Compute \(y^{\prime \prime}\)
\(u^{\prime}(x)=y^{\prime \prime}\)
- Use chain rule on the lhs
\(y^{\prime}\left(\frac{d}{d y} u(y)\right)=y^{\prime \prime}\)
- \(\quad\) Substitute in the definition of \(u\)
\(u(y)\left(\frac{d}{d y} u(y)\right)=y^{\prime \prime}\)
- Make substitutions \(y^{\prime}=u(y), y^{\prime \prime}=u(y)\left(\frac{d}{d y} u(y)\right)\) to reduce order of ODE \(y u(y)\left(\frac{d}{d y} u(y)\right)-u(y)^{2}=1\)
- Separate variables
\(\frac{\left(\frac{d}{d y} u(y)\right) u(y)}{u(y)^{2}+1}=\frac{1}{y}\)
- Integrate both sides with respect to \(y\)
\(\int \frac{\left(\frac{d}{d y} u(y)\right) u(y)}{u(y)^{2}+1} d y=\int \frac{1}{y} d y+c_{1}\)
- Evaluate integral
\(\frac{\ln \left(u(y)^{2}+1\right)}{2}=\ln (y)+c_{1}\)
- \(\quad\) Solve for \(u(y)\)
\(\left\{u(y)=\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}, u(y)=-\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}\right\}\)
- \(\quad\) Solve 1st ODE for \(u(y)\)
\(u(y)=\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}\)
- Revert to original variables with substitution \(u(y)=y^{\prime}, y=y\)
\(y^{\prime}=\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}\)
- Separate variables
\(\frac{y^{\prime}}{\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}}=1\)
- Integrate both sides with respect to \(x\)
\(\int \frac{y^{\prime}}{\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}} d x=\int 1 d x+c_{2}\)
- Evaluate integral
\(\frac{\ln \left(\frac{\left(\mathrm{e}^{c_{1}}\right)^{2} y}{\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}}+\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}\right)}{\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}}=x+c_{2}\)
- \(\quad\) Solve for \(y\)
\[
y=\frac{\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}\left(\left(\mathrm{e}^{c_{2} \sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}+x \sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}}\right)^{2}+1\right)}{2 \mathrm{e}^{c_{2} \sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}+x \sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}\left(\mathrm{e}^{c_{1}}\right)^{2}}}
\]
- \(\quad\) Solve 2nd ODE for \(u(y)\)
\[
u(y)=-\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}
\]
- Revert to original variables with substitution \(u(y)=y^{\prime}, y=y\)
\[
y^{\prime}=-\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}
\]
- \(\quad\) Separate variables
\[
\frac{y^{\prime}}{\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}}=-1
\]
- Integrate both sides with respect to \(x\)
\[
\int \frac{y^{\prime}}{\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}} d x=\int(-1) d x+c_{2}
\]
- Evaluate integral
\[
\frac{\ln \left(\frac{\left(e^{\left.c_{1}\right)^{2} y}\right.}{\sqrt{\left(\mathrm{e}_{1}\right)^{2}}}+\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2} y^{2}-1}\right)}{\sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}}=-x+c_{2}
\]
- \(\quad\) Solve for \(y\)
\[
y=\frac{\sqrt{\left(\mathrm{e}_{1}^{c_{1}}\right)^{2}}\left(\left(\mathrm{e}^{\left.\left.c_{2} \sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}-x \sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}\right)^{2}+1\right)}\right.\right.}{2 \mathrm{e}^{c_{2} \sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}-x \sqrt{\left(\mathrm{e}^{c_{1}}\right)^{2}}}\left(\mathrm{e}^{c_{1}}\right)^{2}}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying 2nd order Liouville trying 2nd order WeierstrassP trying 2nd order JacobiSN differential order: 2; trying a linearization to 3rd order trying 2nd order ODE linearizable_by_differentiation trying 2nd order, 2 integrating factors of the form mu(x,y) trying differential order: 2; missing variables `, `-> Computing symmetries using: way = 3 <- differential order: 2; canonical coordinates successful <- differential order 2; missing variables successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.11 (sec). Leaf size: 55
dsolve \((y(x) * \operatorname{diff}(y(x), x \$ 2)=1+\operatorname{diff}(y(x), x) \sim 2, y(x)\), singsol=all)
\[
\begin{aligned}
& y(x)=\frac{c_{1}\left(\mathrm{e}^{\frac{c_{2}+x}{c_{1}}}+\mathrm{e}^{\frac{-c_{2}-x}{c_{1}}}\right)}{2} \\
& y(x)=\frac{c_{1}\left(\mathrm{e}^{\frac{c_{2}+x}{c_{1}}}+\mathrm{e}^{\frac{-c_{2}-x}{c_{1}}}\right)}{2}
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 60.235 (sec). Leaf size: 80
DSolve[y[x]*y''[x]==1+(y'[x])~2,y[x],x,IncludeSingularSolutions -> True]
\[
\begin{aligned}
& y(x) \rightarrow-\frac{e^{-c_{1}} \tanh \left(e^{c_{1}}\left(x+c_{2}\right)\right)}{\sqrt{-\operatorname{sech}^{2}\left(e^{c_{1}}\left(x+c_{2}\right)\right)}} \\
& y(x) \rightarrow \frac{e^{-c_{1}} \tanh \left(e^{c_{1}}\left(x+c_{2}\right)\right)}{\sqrt{-\operatorname{sech}^{2}\left(e^{c_{1}}\left(x+c_{2}\right)\right)}}
\end{aligned}
\]

\subsection*{1.123 problem 182}
1.123.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 1257
1.123.2 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 1259
1.123.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 1260
1.123.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 1264
1.123.5 Solving as riccati ode . . . . . . . . . . . . . . . . . . . . . . . . 1268
1.123.6 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1270

Internal problem ID [12540]
Internal file name [OUTPUT/11192_Wednesday_October_18_2023_03_47_02_AM_50152952/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 182.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]
\[
\frac{x^{2} y^{\prime}}{(-y+x)^{2}}-\frac{y^{2}}{(-y+x)^{2}}=0
\]

\subsection*{1.123.1 Solving as separable ode}

In canonical form the ODE is
\[
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{y^{2}}{x^{2}}
\end{aligned}
\]

Where \(f(x)=\frac{1}{x^{2}}\) and \(g(y)=y^{2}\). Integrating both sides gives
\[
\begin{aligned}
\frac{1}{y^{2}} d y & =\frac{1}{x^{2}} d x \\
\int \frac{1}{y^{2}} d y & =\int \frac{1}{x^{2}} d x \\
-\frac{1}{y} & =-\frac{1}{x}+c_{1}
\end{aligned}
\]

Which results in
\[
y=-\frac{x}{c_{1} x-1}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=-\frac{x}{c_{1} x-1} \tag{1}
\end{equation*}
\]


Figure 193: Slope field plot

\section*{Verification of solutions}
\[
y=-\frac{x}{c_{1} x-1}
\]

Verified OK.

\subsection*{1.123.2 Solving as homogeneousTypeD2 ode}

Using the change of variables \(y=u(x) x\) on the above ode results in new ode in \(u(x)\)
\[
\frac{x^{2}\left(u^{\prime}(x) x+u(x)\right)}{(-u(x) x+x)^{2}}-\frac{u(x)^{2} x^{2}}{(-u(x) x+x)^{2}}=0
\]

In canonical form the ODE is
\[
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =\frac{u(u-1)}{x}
\end{aligned}
\]

Where \(f(x)=\frac{1}{x}\) and \(g(u)=u(u-1)\). Integrating both sides gives
\[
\begin{aligned}
\frac{1}{u(u-1)} d u & =\frac{1}{x} d x \\
\int \frac{1}{u(u-1)} d u & =\int \frac{1}{x} d x \\
\ln (u-1)-\ln (u) & =\ln (x)+c_{2}
\end{aligned}
\]

Raising both side to exponential gives
\[
\mathrm{e}^{\ln (u-1)-\ln (u)}=\mathrm{e}^{\ln (x)+c_{2}}
\]

Which simplifies to
\[
\frac{u-1}{u}=c_{3} x
\]

Therefore the solution \(y\) is
\[
\begin{aligned}
y & =x u \\
& =-\frac{x}{c_{3} x-1}
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=-\frac{x}{c_{3} x-1} \tag{1}
\end{equation*}
\]


Figure 194: Slope field plot

\section*{Verification of solutions}
\[
y=-\frac{x}{c_{3} x-1}
\]

Verified OK.

\subsection*{1.123.3 Solving as first order ode lie symmetry lookup ode}

Writing the ode as
\[
\begin{aligned}
y^{\prime} & =\frac{y^{2}}{x^{2}} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
\]

The condition of Lie symmetry is the linearized PDE given by
\[
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
\]

The type of this ode is known. It is of type separable. Therefore we do not need to solve the \(\operatorname{PDE}(\mathrm{A})\), and can just use the lookup table shown below to find \(\xi, \eta\)

Table 189: Lie symmetry infinitesimal lookup table for known first order ODE's
\begin{tabular}{|c|c|c|c|}
\hline ODE class & Form & \(\xi\) & \(\eta\) \\
\hline linear ode & \(y^{\prime}=f(x) y(x)+g(x)\) & 0 & \(e^{\int f d x}\) \\
\hline separable ode & \(y^{\prime}=f(x) g(y)\) & \(\frac{1}{f}\) & 0 \\
\hline quadrature ode & \(y^{\prime}=f(x)\) & 0 & 1 \\
\hline quadrature ode & \(y^{\prime}=g(y)\) & 1 & 0 \\
\hline homogeneous ODEs of Class A & \(y^{\prime}=f\left(\frac{y}{x}\right)\) & \(x\) & \(y\) \\
\hline homogeneous ODEs of Class C & \(y^{\prime}=(a+b x+c y)^{\frac{n}{m}}\) & 1 & \[
-\frac{b}{c}
\] \\
\hline homogeneous class D & \(y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)\) & \(x^{2}\) & \(x y\) \\
\hline First order special form ID 1 & \(y^{\prime}=g(x) e^{h(x)+b y}+f(x)\) & \[
\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}
\] & \[
\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}
\] \\
\hline polynomial type ode & \[
y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}
\] & \[
\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}
\] & \[
\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}
\] \\
\hline Bernoulli ode & \(y^{\prime}=f(x) y+g(x) y^{n}\) & 0 & \(e^{-\int(n-1) f(x) d x} y^{n}\) \\
\hline Reduced Riccati & \(y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}\) & 0 & \(e^{-\int f_{1} d x}\) \\
\hline
\end{tabular}

The above table shows that
\[
\begin{align*}
& \xi(x, y)=x^{2} \\
& \eta(x, y)=0 \tag{A1}
\end{align*}
\]

The next step is to determine the canonical coordinates \(R, S\). The canonical coordinates map \((x, y) \rightarrow(R, S)\) where \((R, S)\) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is
\[
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
\]

The above comes from the requirements that \(\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1\). Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable \(R\) in the canonical coordinates, where \(S(R)\). Since \(\eta=0\) then in this special case
\[
R=y
\]
\(S\) is found from
\[
\begin{aligned}
S & =\int \frac{1}{\xi} d x \\
& =\int \frac{1}{x^{2}} d x
\end{aligned}
\]

Which results in
\[
S=-\frac{1}{x}
\]

Now that \(R, S\) are found, we need to setup the ode in these coordinates. This is done by evaluating
\[
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
\]

Where in the above \(R_{x}, R_{y}, S_{x}, S_{y}\) are all partial derivatives and \(\omega(x, y)\) is the right hand side of the original ode given by
\[
\omega(x, y)=\frac{y^{2}}{x^{2}}
\]

Evaluating all the partial derivatives gives
\[
\begin{aligned}
R_{x} & =0 \\
R_{y} & =1 \\
S_{x} & =\frac{1}{x^{2}} \\
S_{y} & =0
\end{aligned}
\]

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.
\[
\begin{equation*}
\frac{d S}{d R}=\frac{1}{y^{2}} \tag{2~A}
\end{equation*}
\]

We now need to express the RHS as function of \(R\) only. This is done by solving for \(x, y\) in terms of \(R, S\) from the result obtained earlier and simplifying. This gives
\[
\frac{d S}{d R}=\frac{1}{R^{2}}
\]

The above is a quadrature ode．This is the whole point of Lie symmetry method． It converts an ode，no matter how complicated it is，to one that can be solved by integration when the ode is in the canonical coordiates \(R, S\) ．Integrating the above gives
\[
\begin{equation*}
S(R)=-\frac{1}{R}+c_{1} \tag{4}
\end{equation*}
\]

To complete the solution，we just need to transform（4）back to \(x, y\) coordinates．This results in
\[
-\frac{1}{x}=-\frac{1}{y}+c_{1}
\]

Which simplifies to
\[
-\frac{1}{x}=-\frac{1}{y}+c_{1}
\]

Which gives
\[
y=\frac{x}{c_{1} x+1}
\]

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown．
\begin{tabular}{|c|c|c|}
\hline Original ode in \(x, y\) coordinates & Canonical coordinates transformation & ODE in canonical coordinates
\[
(R, S)
\] \\
\hline \(\frac{d y}{d x}=\frac{y^{2}}{x^{2}}\) & & \(\frac{d S}{d R}=\frac{1}{R^{2}}\) \\
\hline  & &  \\
\hline  & &  \\
\hline  & & \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\infty]{ }+\) \\
\hline  & &  \\
\hline  & & \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow-\infty]{ }\) \\
\hline \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \infty]{ }\) & \(R=y\) & \(\rightarrow \pm 19\) \\
\hline  & &  \\
\hline  & & \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \infty]{ }\) \\
\hline  & & \(\rightarrow\) 为分新多多 \\
\hline  & & \(\rightarrow \pm 14\) \\
\hline  & & － \(4+\)＋ \\
\hline  & &  \\
\hline
\end{tabular}

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=\frac{x}{c_{1} x+1} \tag{1}
\end{equation*}
\]


Figure 195: Slope field plot
Verification of solutions
\[
y=\frac{x}{c_{1} x+1}
\]

Verified OK.

\subsection*{1.123.4 Solving as exact ode}

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form
\[
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
\]

We assume there exists a function \(\phi(x, y)=c\) where \(c\) is constant, that satisfies the ode. Taking derivative of \(\phi\) w.r.t. \(x\) gives
\[
\frac{d}{d x} \phi(x, y)=0
\]

Hence
\[
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
\]

Comparing ( \(\mathrm{A}, \mathrm{B}\) ) shows that
\[
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
\]

But since \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) then for the above to be valid, we require that
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

If the above condition is satisfied, then the original ode is called exact. We still need to determine \(\phi(x, y)\) but at least we know now that we can do that since the condition \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is
\[
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
\]

Therefore
\[
\begin{align*}
\left(\frac{1}{y^{2}}\right) \mathrm{d} y & =\left(\frac{1}{x^{2}}\right) \mathrm{d} x \\
\left(-\frac{1}{x^{2}}\right) \mathrm{d} x+\left(\frac{1}{y^{2}}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
\]

Comparing (1A) and (2A) shows that
\[
\begin{aligned}
& M(x, y)=-\frac{1}{x^{2}} \\
& N(x, y)=\frac{1}{y^{2}}
\end{aligned}
\]

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

Using result found above gives
\[
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{1}{x^{2}}\right) \\
& =0
\end{aligned}
\]

And
\[
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{1}{y^{2}}\right) \\
& =0
\end{aligned}
\]

Since \(\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}\), then the ODE is exact The following equations are now set up to solve for the function \(\phi(x, y)\)
\[
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
\]

Integrating (1) w.r.t. \(x\) gives
\[
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{1}{x^{2}} \mathrm{~d} x \\
\phi & =\frac{1}{x}+f(y) \tag{3}
\end{align*}
\]

Where \(f(y)\) is used for the constant of integration since \(\phi\) is a function of both \(x\) and \(y\). Taking derivative of equation (3) w.r.t \(y\) gives
\[
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
\]

But equation (2) says that \(\frac{\partial \phi}{\partial y}=\frac{1}{y^{2}}\). Therefore equation (4) becomes
\[
\begin{equation*}
\frac{1}{y^{2}}=0+f^{\prime}(y) \tag{5}
\end{equation*}
\]

Solving equation (5) for \(f^{\prime}(y)\) gives
\[
f^{\prime}(y)=\frac{1}{y^{2}}
\]

Integrating the above w.r.t \(y\) gives
\[
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{1}{y^{2}}\right) \mathrm{d} y \\
f(y) & =-\frac{1}{y}+c_{1}
\end{aligned}
\]

Where \(c_{1}\) is constant of integration. Substituting result found above for \(f(y)\) into equation (3) gives \(\phi\)
\[
\phi=\frac{1}{x}-\frac{1}{y}+c_{1}
\]

But since \(\phi\) itself is a constant function, then let \(\phi=c_{2}\) where \(c_{2}\) is new constant and combining \(c_{1}\) and \(c_{2}\) constants into new constant \(c_{1}\) gives the solution as
\[
c_{1}=\frac{1}{x}-\frac{1}{y}
\]

The solution becomes
\[
y=-\frac{x}{c_{1} x-1}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=-\frac{x}{c_{1} x-1} \tag{1}
\end{equation*}
\]


Figure 196: Slope field plot

Verification of solutions
\[
y=-\frac{x}{c_{1} x-1}
\]

Verified OK.

\subsection*{1.123.5 Solving as riccati ode}

In canonical form the ODE is
\[
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{y^{2}}{x^{2}}
\end{aligned}
\]

This is a Riccati ODE. Comparing the ODE to solve
\[
y^{\prime}=\frac{y^{2}}{x^{2}}
\]

With Riccati ODE standard form
\[
y^{\prime}=f_{0}(x)+f_{1}(x) y+f_{2}(x) y^{2}
\]

Shows that \(f_{0}(x)=0, f_{1}(x)=0\) and \(f_{2}(x)=\frac{1}{x^{2}}\). Let
\[
\begin{align*}
y & =\frac{-u^{\prime}}{f_{2} u} \\
& =\frac{-u^{\prime}}{\frac{u}{x^{2}}} \tag{1}
\end{align*}
\]

Using the above substitution in the given ODE results (after some simplification)in a second order ODE to solve for \(u(x)\) which is
\[
\begin{equation*}
f_{2} u^{\prime \prime}(x)-\left(f_{2}^{\prime}+f_{1} f_{2}\right) u^{\prime}(x)+f_{2}^{2} f_{0} u(x)=0 \tag{2}
\end{equation*}
\]

But
\[
\begin{aligned}
f_{2}^{\prime} & =-\frac{2}{x^{3}} \\
f_{1} f_{2} & =0 \\
f_{2}^{2} f_{0} & =0
\end{aligned}
\]

Substituting the above terms back in equation (2) gives
\[
\frac{u^{\prime \prime}(x)}{x^{2}}+\frac{2 u^{\prime}(x)}{x^{3}}=0
\]

Solving the above ODE (this ode solved using Maple, not this program), gives
\[
u(x)=c_{1}+\frac{c_{2}}{x}
\]

The above shows that
\[
u^{\prime}(x)=-\frac{c_{2}}{x^{2}}
\]

Using the above in (1) gives the solution
\[
y=\frac{c_{2}}{c_{1}+\frac{c_{2}}{x}}
\]

Dividing both numerator and denominator by \(c_{1}\) gives, after renaming the constant \(\frac{c_{2}}{c_{1}}=c_{3}\) the following solution
\[
y=\frac{1}{c_{3}+\frac{1}{x}}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=\frac{1}{c_{3}+\frac{1}{x}} \tag{1}
\end{equation*}
\]


Figure 197: Slope field plot

\section*{Verification of solutions}
\[
y=\frac{1}{c_{3}+\frac{1}{x}}
\]

Verified OK.

\subsection*{1.123.6 Maple step by step solution}

Let's solve
\[
\frac{x^{2} y^{\prime}}{(-y+x)^{2}}-\frac{y^{2}}{(-y+x)^{2}}=0
\]
- Highest derivative means the order of the ODE is 1 \(y^{\prime}\)
- Integrate both sides with respect to \(x\)
\[
\int\left(\frac{x^{2} y^{\prime}}{(-y+x)^{2}}-\frac{y^{2}}{(-y+x)^{2}}\right) d x=\int 0 d x+c_{1}
\]
- Evaluate integral
\[
\frac{x^{2}}{-y+x}-x=c_{1}
\]
- \(\quad\) Solve for \(y\)
\[
y=\frac{c_{1} x}{x+c_{1}}
\]

Maple trace
- Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 13
```

dsolve(x^2*diff(y(x),x)/(x-y(x))^2-y(x)^2/(x-y(x))^2=0,y(x), singsol=all)

```
\[
y(x)=\frac{x}{c_{1} x+1}
\]
\(\sqrt{ }\) Solution by Mathematica
Time used: 0.225 (sec). Leaf size: 21
DSolve \(\left[x^{\wedge} 2 * y^{\prime}[x] /(x-y[x])^{\wedge} 2-y[x] \sim 2 /(x-y[x]) \wedge 2==0, y[x], x\right.\), IncludeSingularSolutions \(\rightarrow\) True]
\[
\begin{aligned}
& y(x) \rightarrow \frac{x}{1-c_{1} x} \\
& y(x) \rightarrow 0
\end{aligned}
\]

\subsection*{1.124 problem 183}
1.124.1 Solving as dAlembert ode

Internal problem ID [12541]
Internal file name [OUTPUT/11193_Wednesday_October_18_2023_03_47_03_AM_12919399/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 183.
ODE order: 1.
ODE degree: 2.

The type(s) of ODE detected by this program : "dAlembert"
Maple gives the following as the ode type
```

[[_homogeneous, `class C`], _rational, _dAlembert]

```
\[
y-x y^{\prime 2}-y^{\prime 2}=0
\]

\subsection*{1.124.1 Solving as dAlembert ode}

Let \(p=y^{\prime}\) the ode becomes
\[
-x p^{2}-p^{2}+y=0
\]

Solving for \(y\) from the above results in
\[
\begin{equation*}
y=x p^{2}+p^{2} \tag{1A}
\end{equation*}
\]

This has the form
\[
\begin{equation*}
y=x f(p)+g(p) \tag{}
\end{equation*}
\]

Where \(f, g\) are functions of \(p=y^{\prime}(x)\). The above ode is dAlembert ode which is now solved. Taking derivative of \(\left({ }^{*}\right)\) w.r.t. \(x\) gives
\[
\begin{align*}
p & =f+\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \\
p-f & =\left(x f^{\prime}+g^{\prime}\right) \frac{d p}{d x} \tag{2}
\end{align*}
\]

Comparing the form \(y=x f+g\) to (1A) shows that
\[
\begin{aligned}
& f=p^{2} \\
& g=p^{2}
\end{aligned}
\]

Hence (2) becomes
\[
\begin{equation*}
-p^{2}+p=(2 x p+2 p) p^{\prime}(x) \tag{2~A}
\end{equation*}
\]

The singular solution is found by setting \(\frac{d p}{d x}=0\) in the above which gives
\[
-p^{2}+p=0
\]

Solving for \(p\) from the above gives
\[
\begin{aligned}
& p=0 \\
& p=1
\end{aligned}
\]

Substituting these in (1A) gives
\[
\begin{aligned}
& y=0 \\
& y=x+1
\end{aligned}
\]

The general solution is found when \(\frac{\mathrm{d} p}{\mathrm{~d} x} \neq 0\). From eq. (2A). This results in
\[
\begin{equation*}
p^{\prime}(x)=\frac{-p(x)^{2}+p(x)}{2 p(x) x+2 p(x)} \tag{3}
\end{equation*}
\]

This ODE is now solved for \(p(x)\).
Entering Linear first order ODE solver. In canonical form a linear first order is
\[
p^{\prime}(x)+p(x) p(x)=q(x)
\]

Where here
\[
\begin{aligned}
p(x) & =\frac{1}{2 x+2} \\
q(x) & =\frac{1}{2 x+2}
\end{aligned}
\]

Hence the ode is
\[
p^{\prime}(x)+\frac{p(x)}{2 x+2}=\frac{1}{2 x+2}
\]

The integrating factor \(\mu\) is
\[
\begin{aligned}
\mu & =\mathrm{e}^{\int \frac{1}{2 x+2} d x} \\
& =\sqrt{x+1}
\end{aligned}
\]

The ode becomes
\[
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu p) & =(\mu)\left(\frac{1}{2 x+2}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}(\sqrt{x+1} p) & =(\sqrt{x+1})\left(\frac{1}{2 x+2}\right) \\
\mathrm{d}(\sqrt{x+1} p) & =\left(\frac{1}{2 \sqrt{x+1}}\right) \mathrm{d} x
\end{aligned}
\]

Integrating gives
\[
\begin{aligned}
& \sqrt{x+1} p=\int \frac{1}{2 \sqrt{x+1}} \mathrm{~d} x \\
& \sqrt{x+1} p=\sqrt{x+1}+c_{1}
\end{aligned}
\]

Dividing both sides by the integrating factor \(\mu=\sqrt{x+1}\) results in
\[
p(x)=1+\frac{c_{1}}{\sqrt{x+1}}
\]

Substituing the above solution for \(p\) in (2A) gives
\[
y=x\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2}+\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2}
\]

Summary
The solution(s) found are the following
\[
\begin{align*}
& y=0  \tag{1}\\
& y=x+1  \tag{2}\\
& y=x\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2}+\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2} \tag{3}
\end{align*}
\]

\section*{Verification of solutions}
\[
y=0
\]

Verified OK.
\[
y=x+1
\]

Verified OK.
\[
y=x\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2}+\left(1+\frac{c_{1}}{\sqrt{x+1}}\right)^{2}
\]

Verified OK.
Maple trace
- Methods for first order ODEs:
*** Sublevel 2 ***
Methods for first order ODEs:
-> Solving 1st order ODE of high degree, 1st attempt
trying 1st order WeierstrassP solution for high degree ODE
trying 1st order WeierstrassPPrime solution for high degree ODE
trying 1st order JacobiSN solution for high degree ODE
trying 1st order ODE linearizable_by_differentiation
trying differential order: 1; missing variables
trying dAlembert
<- dAlembert successful`
\(\checkmark\) Solution by Maple
Time used: 0.016 (sec). Leaf size: 53
```

dsolve(y(x)=x*diff(y(x),x)^2+diff(y(x),x)^2,y(x), singsol=all)

```
\[
\begin{aligned}
& y(x)=0 \\
& y(x)=\frac{\left(x+1+\sqrt{(1+x)\left(c_{1}+1\right)}\right)^{2}}{1+x} \\
& y(x)=\frac{\left(-x-1+\sqrt{(1+x)\left(c_{1}+1\right)}\right)^{2}}{1+x}
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.108 (sec). Leaf size: 57
DSolve[y[x]==x*(y'[x]) \(2+(y \prime[x]) \sim 2, y[x], x\), IncludeSingularSolutions \(->\) True]
\[
\begin{aligned}
& y(x) \rightarrow x-c_{1} \sqrt{x+1}+1+\frac{c_{1}^{2}}{4} \\
& y(x) \rightarrow x+c_{1} \sqrt{x+1}+1+\frac{c_{1}^{2}}{4} \\
& y(x) \rightarrow 0
\end{aligned}
\]

\subsection*{1.125 problem 184}
1.125.1 Solving as second order linear constant coeff ode . . . . . . . . 1277
1.125.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1281
1.125.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1287

Internal problem ID [12542]
Internal file name [OUTPUT/11194_Wednesday_October_18_2023_03_47_04_AM_99036748/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 184.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _linear, _nonhomogeneous]]
\[
y^{\prime \prime}+y=\sec (x)
\]

\subsection*{1.125.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=0, C=1, f(x)=\sec (x)\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=1\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}+\mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}+1=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=1\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)(1)} \\
& = \pm i
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=+i \\
& \lambda_{2}=-i
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
& \lambda_{1}=i \\
& \lambda_{2}=-i
\end{aligned}
\]

Since roots are complex conjugate of each others, then let the roots be
\[
\lambda_{1,2}=\alpha \pm i \beta
\]

Where \(\alpha=0\) and \(\beta=1\). Therefore the final solution, when using Euler relation, can be written as
\[
y=e^{\alpha x}\left(c_{1} \cos (\beta x)+c_{2} \sin (\beta x)\right)
\]

Which becomes
\[
y=e^{0}\left(c_{1} \cos (x)+c_{2} \sin (x)\right)
\]

Or
\[
y=c_{1} \cos (x)+c_{2} \sin (x)
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \cos (x)+c_{2} \sin (x)
\]

The particular solution \(y_{p}\) can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on \(x\) as well. Let
\[
\begin{equation*}
y_{p}(x)=u_{1} y_{1}+u_{2} y_{2} \tag{1}
\end{equation*}
\]

Where \(u_{1}, u_{2}\) to be determined, and \(y_{1}, y_{2}\) are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as
\[
\begin{aligned}
& y_{1}=\cos (x) \\
& y_{2}=\sin (x)
\end{aligned}
\]

In the Variation of parameters \(u_{1}, u_{2}\) are found using
\[
\begin{align*}
& u_{1}=-\int \frac{y_{2} f(x)}{a W(x)}  \tag{2}\\
& u_{2}=\int \frac{y_{1} f(x)}{a W(x)} \tag{3}
\end{align*}
\]

Where \(W(x)\) is the Wronskian and \(a\) is the coefficient in front of \(y^{\prime \prime}\) in the given ODE. The Wronskian is given by \(W=\left|\begin{array}{ll}y_{1} & y_{2} \\ y_{1}^{\prime} & y_{2}^{\prime}\end{array}\right|\). Hence
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
\frac{d}{d x}(\cos (x)) & \frac{d}{d x}(\sin (x))
\end{array}\right|
\]

Which gives
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
-\sin (x) & \cos (x)
\end{array}\right|
\]

Therefore
\[
W=(\cos (x))(\cos (x))-(\sin (x))(-\sin (x))
\]

Which simplifies to
\[
W=\cos (x)^{2}+\sin (x)^{2}
\]

Which simplifies to
\[
W=1
\]

Therefore Eq. (2) becomes
\[
u_{1}=-\int \frac{\sin (x) \sec (x)}{1} d x
\]

Which simplifies to
\[
u_{1}=-\int \tan (x) d x
\]

Hence
\[
u_{1}=\ln (\cos (x))
\]

And Eq. (3) becomes
\[
u_{2}=\int \frac{\cos (x) \sec (x)}{1} d x
\]

Which simplifies to
\[
u_{2}=\int 1 d x
\]

Hence
\[
u_{2}=x
\]

Therefore the particular solution, from equation (1) is
\[
y_{p}(x)=\ln (\cos (x)) \cos (x)+x \sin (x)
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \cos (x)+c_{2} \sin (x)\right)+(\ln (\cos (x)) \cos (x)+x \sin (x))
\end{aligned}
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (x)+c_{2} \sin (x)+\ln (\cos (x)) \cos (x)+x \sin (x) \tag{1}
\end{equation*}
\]


Figure 198: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \cos (x)+c_{2} \sin (x)+\ln (\cos (x)) \cos (x)+x \sin (x)
\]

Verified OK.

\subsection*{1.125.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}+y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=0  \tag{3}\\
& C=1
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{-1}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{aligned}
& s=-1 \\
& t=1
\end{aligned}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=-z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi-
\end{tabular} & no condition \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\). & \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 192: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=-1\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\cos (x)
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\cos (x)
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\cos (x)
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\cos (x) \int \frac{1}{\cos (x)^{2}} d x \\
& =\cos (x)(\tan (x))
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}(\cos (x))+c_{2}(\cos (x)(\tan (x)))
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}+y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \cos (x)+c_{2} \sin (x)
\]

The particular solution \(y_{p}\) can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of
parameters will be used as it is more general and can be used when the coefficients of the ODE depend on \(x\) as well. Let
\[
\begin{equation*}
y_{p}(x)=u_{1} y_{1}+u_{2} y_{2} \tag{1}
\end{equation*}
\]

Where \(u_{1}, u_{2}\) to be determined, and \(y_{1}, y_{2}\) are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as
\[
\begin{aligned}
& y_{1}=\cos (x) \\
& y_{2}=\sin (x)
\end{aligned}
\]

In the Variation of parameters \(u_{1}, u_{2}\) are found using
\[
\begin{align*}
& u_{1}=-\int \frac{y_{2} f(x)}{a W(x)}  \tag{2}\\
& u_{2}=\int \frac{y_{1} f(x)}{a W(x)} \tag{3}
\end{align*}
\]

Where \(W(x)\) is the Wronskian and \(a\) is the coefficient in front of \(y^{\prime \prime}\) in the given ODE. The Wronskian is given by \(W=\left|\begin{array}{ll}y_{1} & y_{2} \\ y_{1}^{\prime} & y_{2}^{\prime}\end{array}\right|\). Hence
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
\frac{d}{d x}(\cos (x)) & \frac{d}{d x}(\sin (x))
\end{array}\right|
\]

Which gives
\[
W=\left|\begin{array}{cc}
\cos (x) & \sin (x) \\
-\sin (x) & \cos (x)
\end{array}\right|
\]

Therefore
\[
W=(\cos (x))(\cos (x))-(\sin (x))(-\sin (x))
\]

Which simplifies to
\[
W=\cos (x)^{2}+\sin (x)^{2}
\]

Which simplifies to
\[
W=1
\]

Therefore Eq. (2) becomes
\[
u_{1}=-\int \frac{\sin (x) \sec (x)}{1} d x
\]

Which simplifies to
\[
u_{1}=-\int \tan (x) d x
\]

Hence
\[
u_{1}=\ln (\cos (x))
\]

And Eq. (3) becomes
\[
u_{2}=\int \frac{\cos (x) \sec (x)}{1} d x
\]

Which simplifies to
\[
u_{2}=\int 1 d x
\]

Hence
\[
u_{2}=x
\]

Therefore the particular solution, from equation (1) is
\[
y_{p}(x)=\ln (\cos (x)) \cos (x)+x \sin (x)
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \cos (x)+c_{2} \sin (x)\right)+(\ln (\cos (x)) \cos (x)+x \sin (x))
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \cos (x)+c_{2} \sin (x)+\ln (\cos (x)) \cos (x)+x \sin (x) \tag{1}
\end{equation*}
\]


Figure 199: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \cos (x)+c_{2} \sin (x)+\ln (\cos (x)) \cos (x)+x \sin (x)
\]

Verified OK.

\subsection*{1.125.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}+y=\sec (x)
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\[
r^{2}+1=0
\]
- Use quadratic formula to solve for \(r\)
\[
r=\frac{0 \pm(\sqrt{-4})}{2}
\]
- Roots of the characteristic polynomial
\[
r=(-\mathrm{I}, \mathrm{I})
\]
- \(\quad 1\) st solution of the homogeneous ODE
\[
y_{1}(x)=\cos (x)
\]
- \(\quad 2 n d\) solution of the homogeneous ODE
\[
y_{2}(x)=\sin (x)
\]
- General solution of the ODE
\[
y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)
\]
- Substitute in solutions of the homogeneous ODE
\[
y=c_{1} \cos (x)+c_{2} \sin (x)+y_{p}(x)
\]

Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\[
\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=\sec (x)\right]
\]
- Wronskian of solutions of the homogeneous equation
\[
W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}
\cos (x) & \sin (x) \\
-\sin (x) & \cos (x)
\end{array}\right]
\]
- Compute Wronskian
\(W\left(y_{1}(x), y_{2}(x)\right)=1\)
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=-\cos (x)\left(\int \tan (x) d x\right)+\sin (x)\left(\int 1 d x\right)
\]
- Compute integrals
\[
y_{p}(x)=\ln (\cos (x)) \cos (x)+x \sin (x)
\]
- Substitute particular solution into general solution to ODE
\[
y=c_{1} \cos (x)+c_{2} \sin (x)+\ln (\cos (x)) \cos (x)+x \sin (x)
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 22
```

dsolve(diff(y(x),x\$2)+y(x)=sec(x),y(x), singsol=all)

```
\[
y(x)=-\ln (\sec (x)) \cos (x)+c_{1} \cos (x)+\sin (x)\left(c_{2}+x\right)
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.033 (sec). Leaf size: 22
DSolve[y''[x]+y[x]==Sec[x],y[x],x,IncludeSingularSolutions -> True]
\[
y(x) \rightarrow\left(x+c_{2}\right) \sin (x)+\cos (x)\left(\log (\cos (x))+c_{1}\right)
\]

\subsection*{1.126 problem 185}
1.126.1 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 1290
1.126.2 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 1292
1.126.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 1293
1.126.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 1296
1.126.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1300

Internal problem ID [12543]
Internal file name [OUTPUT/11195_Wednesday_October_18_2023_03_47_04_AM_50697412/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 185.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "linear", "homogeneousTypeD2", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry__lookup"

Maple gives the following as the ode type
[_linear]
\[
\left(x^{2}+1\right) y^{\prime}-y x=\alpha
\]

\subsection*{1.126.1 Solving as linear ode}

Entering Linear first order ODE solver. In canonical form a linear first order is
\[
y^{\prime}+p(x) y=q(x)
\]

Where here
\[
\begin{aligned}
& p(x)=-\frac{x}{x^{2}+1} \\
& q(x)=\frac{\alpha}{x^{2}+1}
\end{aligned}
\]

Hence the ode is
\[
y^{\prime}-\frac{x y}{x^{2}+1}=\frac{\alpha}{x^{2}+1}
\]

The integrating factor \(\mu\) is
\[
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{x}{x^{2}+1} d x} \\
& =\frac{1}{\sqrt{x^{2}+1}}
\end{aligned}
\]

The ode becomes
\[
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\frac{\alpha}{x^{2}+1}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{y}{\sqrt{x^{2}+1}}\right) & =\left(\frac{1}{\sqrt{x^{2}+1}}\right)\left(\frac{\alpha}{x^{2}+1}\right) \\
\mathrm{d}\left(\frac{y}{\sqrt{x^{2}+1}}\right) & =\left(\frac{\alpha}{\left(x^{2}+1\right)^{\frac{3}{2}}}\right) \mathrm{d} x
\end{aligned}
\]

Integrating gives
\[
\begin{aligned}
& \frac{y}{\sqrt{x^{2}+1}}=\int \frac{\alpha}{\left(x^{2}+1\right)^{\frac{3}{2}}} \mathrm{~d} x \\
& \frac{y}{\sqrt{x^{2}+1}}=\frac{x \alpha}{\sqrt{x^{2}+1}}+c_{1}
\end{aligned}
\]

Dividing both sides by the integrating factor \(\mu=\frac{1}{\sqrt{x^{2}+1}}\) results in
\[
y=x \alpha+c_{1} \sqrt{x^{2}+1}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=x \alpha+c_{1} \sqrt{x^{2}+1} \tag{1}
\end{equation*}
\]

\section*{Verification of solutions}
\[
y=x \alpha+c_{1} \sqrt{x^{2}+1}
\]

Verified OK.

\subsection*{1.126.2 Solving as homogeneousTypeD2 ode}

Using the change of variables \(y=u(x) x\) on the above ode results in new ode in \(u(x)\)
\[
\left(x^{2}+1\right)\left(u^{\prime}(x) x+u(x)\right)-u(x) x^{2}=\alpha
\]

In canonical form the ODE is
\[
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =\frac{-u+\alpha}{x\left(x^{2}+1\right)}
\end{aligned}
\]

Where \(f(x)=\frac{1}{x\left(x^{2}+1\right)}\) and \(g(u)=-u+\alpha\). Integrating both sides gives
\[
\begin{aligned}
\frac{1}{-u+\alpha} d u & =\frac{1}{x\left(x^{2}+1\right)} d x \\
\int \frac{1}{-u+\alpha} d u & =\int \frac{1}{x\left(x^{2}+1\right)} d x \\
-\ln (-u+\alpha) & =\ln (x)-\frac{\ln \left(x^{2}+1\right)}{2}+c_{2}
\end{aligned}
\]

Raising both side to exponential gives
\[
\frac{1}{-u+\alpha}=\mathrm{e}^{\ln (x)-\frac{\ln \left(x^{2}+1\right)}{2}+c_{2}}
\]

Which simplifies to
\[
\frac{1}{-u+\alpha}=c_{3} \mathrm{e}^{\ln (x)-\frac{\ln \left(x^{2}+1\right)}{2}}
\]

Which simplifies to
\[
u(x)=\frac{\left(\frac{c_{3} \mathrm{e}^{c_{2} x \alpha}}{\sqrt{x^{2}+1}}-1\right) \mathrm{e}^{-c_{2}} \sqrt{x^{2}+1}}{c_{3} x}
\]

Therefore the solution \(y\) is
\[
\begin{aligned}
y & =x u \\
& =\frac{\left(\frac{c_{3} e^{c_{2} x \alpha}}{\sqrt{x^{2}+1}}-1\right) \mathrm{e}^{-c_{2}} \sqrt{x^{2}+1}}{c_{3}}
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{\left(\frac{c_{3} e^{c_{2} x \alpha}}{\sqrt{x^{2}+1}}-1\right) \mathrm{e}^{-c_{2}} \sqrt{x^{2}+1}}{c_{3}} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=\frac{\left(\frac{c_{3} e^{c_{2}} x \alpha}{\sqrt{x^{2}+1}}-1\right) \mathrm{e}^{-c_{2}} \sqrt{x^{2}+1}}{c_{3}}
\]

Verified OK.

\subsection*{1.126.3 Solving as first order ode lie symmetry lookup ode}

Writing the ode as
\[
\begin{aligned}
y^{\prime} & =\frac{x y+\alpha}{x^{2}+1} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
\]

The condition of Lie symmetry is the linearized PDE given by
\[
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
\]

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find \(\xi, \eta\)

Table 194: Lie symmetry infinitesimal lookup table for known first order ODE's
\begin{tabular}{|l|l|l|l|}
\hline ODE class & Form & \(\xi\) & \(\eta\) \\
\hline \hline linear ode & \(y^{\prime}=f(x) y(x)+g(x)\) & 0 & \(e^{\int f d x}\) \\
\hline separable ode & \(y^{\prime}=f(x) g(y)\) & \(\frac{1}{f}\) & 0 \\
\hline quadrature ode & \(y^{\prime}=f(x)\) & 0 & 1 \\
\hline quadrature ode & \(y^{\prime}=g(y)\) & 1 & 0 \\
\hline \begin{tabular}{l} 
homogeneous ODEs of \\
Class A
\end{tabular} & \(y^{\prime}=f\left(\frac{y}{x}\right)\) & \(x\) & \(y\) \\
\hline \begin{tabular}{l} 
homogeneous ODEs of \\
Class C
\end{tabular} & \(y^{\prime}=(a+b x+c y)^{\frac{n}{m}}\) & 1 & \(-\frac{b}{c}\) \\
\hline homogeneous class D & \(y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)\) & \(x^{2}\) & \(x y\) \\
\hline \begin{tabular}{l} 
First order \\
form ID 1
\end{tabular} & \(y^{2}=g(x) e^{h(x)+b y}+f(x)\) & \(\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}\) & \(\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}\) \\
\hline polynomial type ode & \(y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}\) & \(\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}\) & \(\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}\) \\
\hline Bernoulli ode & \(y^{\prime}=f(x) y+g(x) y^{n}\) & 0 & \(e^{-\int(n-1) f(x) d x} y^{n}\) \\
\hline Reduced Riccati & \(y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}\) & 0 & \(e^{-\int f_{1} d x}\) \\
\hline
\end{tabular}

The above table shows that
\[
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\sqrt{x^{2}+1} \tag{A1}
\end{align*}
\]

The next step is to determine the canonical coordinates \(R, S\). The canonical coordinates \(\operatorname{map}(x, y) \rightarrow(R, S)\) where \((R, S)\) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is
\[
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
\]

The above comes from the requirements that \(\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1\). Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable \(R\) in the
canonical coordinates, where \(S(R)\). Since \(\xi=0\) then in this special case
\[
R=x
\]
\(S\) is found from
\[
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\sqrt{x^{2}+1}} d y
\end{aligned}
\]

Which results in
\[
S=\frac{y}{\sqrt{x^{2}+1}}
\]

Now that \(R, S\) are found, we need to setup the ode in these coordinates. This is done by evaluating
\[
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
\]

Where in the above \(R_{x}, R_{y}, S_{x}, S_{y}\) are all partial derivatives and \(\omega(x, y)\) is the right hand side of the original ode given by
\[
\omega(x, y)=\frac{x y+\alpha}{x^{2}+1}
\]

Evaluating all the partial derivatives gives
\[
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =-\frac{y x}{\left(x^{2}+1\right)^{\frac{3}{2}}} \\
S_{y} & =\frac{1}{\sqrt{x^{2}+1}}
\end{aligned}
\]

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.
\[
\begin{equation*}
\frac{d S}{d R}=\frac{\alpha}{\left(x^{2}+1\right)^{\frac{3}{2}}} \tag{2~A}
\end{equation*}
\]

We now need to express the RHS as function of \(R\) only. This is done by solving for \(x, y\) in terms of \(R, S\) from the result obtained earlier and simplifying. This gives
\[
\frac{d S}{d R}=\frac{\alpha}{\left(R^{2}+1\right)^{\frac{3}{2}}}
\]

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates \(R, S\). Integrating the above gives
\[
\begin{equation*}
S(R)=\frac{R \alpha}{\sqrt{R^{2}+1}}+c_{1} \tag{4}
\end{equation*}
\]

To complete the solution, we just need to transform (4) back to \(x, y\) coordinates. This results in
\[
\frac{y}{\sqrt{x^{2}+1}}=\frac{x \alpha}{\sqrt{x^{2}+1}}+c_{1}
\]

Which simplifies to
\[
\frac{y}{\sqrt{x^{2}+1}}=\frac{x \alpha}{\sqrt{x^{2}+1}}+c_{1}
\]

Which gives
\[
y=x \alpha+c_{1} \sqrt{x^{2}+1}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=x \alpha+c_{1} \sqrt{x^{2}+1} \tag{1}
\end{equation*}
\]

Verification of solutions
\[
y=x \alpha+c_{1} \sqrt{x^{2}+1}
\]

Verified OK.

\subsection*{1.126.4 Solving as exact ode}

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form
\[
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
\]

We assume there exists a function \(\phi(x, y)=c\) where \(c\) is constant, that satisfies the ode. Taking derivative of \(\phi\) w.r.t. \(x\) gives
\[
\frac{d}{d x} \phi(x, y)=0
\]

Hence
\[
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
\]

Comparing ( \(\mathrm{A}, \mathrm{B}\) ) shows that
\[
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
\]

But since \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) then for the above to be valid, we require that
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

If the above condition is satisfied, then the original ode is called exact. We still need to determine \(\phi(x, y)\) but at least we know now that we can do that since the condition \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is
\[
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
\]

Therefore
\[
\begin{align*}
\left(x^{2}+1\right) \mathrm{d} y & =(x y+\alpha) \mathrm{d} x \\
(-x y-\alpha) \mathrm{d} x+\left(x^{2}+1\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
\]

Comparing (1A) and (2A) shows that
\[
\begin{aligned}
M(x, y) & =-x y-\alpha \\
N(x, y) & =x^{2}+1
\end{aligned}
\]

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

Using result found above gives
\[
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}(-x y-\alpha) \\
& =-x
\end{aligned}
\]

And
\[
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(x^{2}+1\right) \\
& =2 x
\end{aligned}
\]

Since \(\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}\), then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let
\[
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =\frac{1}{x^{2}+1}((-x)-(2 x)) \\
& =-\frac{3 x}{x^{2}+1}
\end{aligned}
\]

Since \(A\) does not depend on \(y\), then it can be used to find an integrating factor. The integrating factor \(\mu\) is
\[
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-\frac{3 x}{x^{2}+1} \mathrm{~d} x}
\end{aligned}
\]

The result of integrating gives
\[
\begin{aligned}
\mu & =e^{-\frac{3 \ln \left(x^{2}+1\right)}{2}} \\
& =\frac{1}{\left(x^{2}+1\right)^{\frac{3}{2}}}
\end{aligned}
\]
\(M\) and \(N\) are multiplied by this integrating factor, giving new \(M\) and new \(N\) which are called \(\bar{M}\) and \(\bar{N}\) for now so not to confuse them with the original \(M\) and \(N\).
\[
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{\left(x^{2}+1\right)^{\frac{3}{2}}}(-x y-\alpha) \\
& =-\frac{x y+\alpha}{\left(x^{2}+1\right)^{\frac{3}{2}}}
\end{aligned}
\]

And
\[
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{\left(x^{2}+1\right)^{\frac{3}{2}}}\left(x^{2}+1\right) \\
& =\frac{1}{\sqrt{x^{2}+1}}
\end{aligned}
\]

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is
\[
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x} & =0 \\
\left(-\frac{x y+\alpha}{\left(x^{2}+1\right)^{\frac{3}{2}}}\right)+\left(\frac{1}{\sqrt{x^{2}+1}}\right) \frac{\mathrm{d} y}{\mathrm{~d} x} & =0
\end{aligned}
\]

The following equations are now set up to solve for the function \(\phi(x, y)\)
\[
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
\]

Integrating (1) w.r.t. \(x\) gives
\[
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{x y+\alpha}{\left(x^{2}+1\right)^{\frac{3}{2}}} \mathrm{~d} x \\
\phi & =\frac{-x \alpha+y}{\sqrt{x^{2}+1}}+f(y) \tag{3}
\end{align*}
\]

Where \(f(y)\) is used for the constant of integration since \(\phi\) is a function of both \(x\) and \(y\). Taking derivative of equation (3) w.r.t \(y\) gives
\[
\begin{equation*}
\frac{\partial \phi}{\partial y}=\frac{1}{\sqrt{x^{2}+1}}+f^{\prime}(y) \tag{4}
\end{equation*}
\]

But equation (2) says that \(\frac{\partial \phi}{\partial y}=\frac{1}{\sqrt{x^{2}+1}}\). Therefore equation (4) becomes
\[
\begin{equation*}
\frac{1}{\sqrt{x^{2}+1}}=\frac{1}{\sqrt{x^{2}+1}}+f^{\prime}(y) \tag{5}
\end{equation*}
\]

Solving equation (5) for \(f^{\prime}(y)\) gives
\[
f^{\prime}(y)=0
\]

Therefore
\[
f(y)=c_{1}
\]

Where \(c_{1}\) is constant of integration. Substituting this result for \(f(y)\) into equation (3) gives \(\phi\)
\[
\phi=\frac{-x \alpha+y}{\sqrt{x^{2}+1}}+c_{1}
\]

But since \(\phi\) itself is a constant function, then let \(\phi=c_{2}\) where \(c_{2}\) is new constant and combining \(c_{1}\) and \(c_{2}\) constants into new constant \(c_{1}\) gives the solution as
\[
c_{1}=\frac{-x \alpha+y}{\sqrt{x^{2}+1}}
\]

The solution becomes
\[
y=x \alpha+c_{1} \sqrt{x^{2}+1}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=x \alpha+c_{1} \sqrt{x^{2}+1} \tag{1}
\end{equation*}
\]

\section*{Verification of solutions}
\[
y=x \alpha+c_{1} \sqrt{x^{2}+1}
\]

Verified OK.

\subsection*{1.126.5 Maple step by step solution}

Let's solve
\[
\left(x^{2}+1\right) y^{\prime}-y x=\alpha
\]
- Highest derivative means the order of the ODE is 1
\(y^{\prime}\)
- Isolate the derivative
\(y^{\prime}=\frac{x y}{x^{2}+1}+\frac{\alpha}{x^{2}+1}\)
- Group terms with \(y\) on the lhs of the ODE and the rest on the rhs of the ODE \(y^{\prime}-\frac{x y}{x^{2}+1}=\frac{\alpha}{x^{2}+1}\)
- The ODE is linear; multiply by an integrating factor \(\mu(x)\)
\(\mu(x)\left(y^{\prime}-\frac{x y}{x^{2}+1}\right)=\frac{\mu(x) \alpha}{x^{2}+1}\)
- Assume the lhs of the ODE is the total derivative \(\frac{d}{d x}(\mu(x) y)\)
\(\mu(x)\left(y^{\prime}-\frac{x y}{x^{2}+1}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}\)
- Isolate \(\mu^{\prime}(x)\)
\(\mu^{\prime}(x)=-\frac{\mu(x) x}{x^{2}+1}\)
- Solve to find the integrating factor
\(\mu(x)=\frac{1}{\sqrt{x^{2}+1}}\)
- Integrate both sides with respect to \(x\)
\(\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int \frac{\mu(x) \alpha}{x^{2}+1} d x+c_{1}\)
- Evaluate the integral on the lhs
\(\mu(x) y=\int \frac{\mu(x) \alpha}{x^{2}+1} d x+c_{1}\)
- \(\quad\) Solve for \(y\)
\(y=\frac{\int \frac{\mu(x) \alpha}{x^{2}+1} d x+c_{1}}{\mu(x)}\)
- \(\quad\) Substitute \(\mu(x)=\frac{1}{\sqrt{x^{2}+1}}\)
\(y=\sqrt{x^{2}+1}\left(\int \frac{\alpha}{\left(x^{2}+1\right)^{\frac{3}{2}}} d x+c_{1}\right)\)
- Evaluate the integrals on the rhs
\(y=\sqrt{x^{2}+1}\left(\frac{x \alpha}{\sqrt{x^{2}+1}}+c_{1}\right)\)
- Simplify
\(y=x \alpha+c_{1} \sqrt{x^{2}+1}\)

Maple trace
```

`Methods for first order ODEs: --- Trying classification methods --- trying a quadrature trying 1st order linear <- 1st order linear successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.015 (sec). Leaf size: 17
dsolve( \(\left(1+x^{\wedge} 2\right) * \operatorname{diff}(y(x), x)-x * y(x)-\operatorname{alpha}=0, y(x)\), singsol=all)
\[
y(x)=c_{1} \sqrt{x^{2}+1}+\alpha x
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.114 (sec). Leaf size: 21
DSolve[(1+x^2)*y'[x]-x*y[x]-a==0,y[x],x,IncludeSingularSolutions -> True]
\[
y(x) \rightarrow a x+c_{1} \sqrt{x^{2}+1}
\]

\subsection*{1.127 problem 186}
1.127.1 Solving as homogeneousTypeD ode . . . . . . . . . . . . . . . . 1303
1.127.2 Solving as homogeneousTypeD2 ode . . . . . . . . . . . . . . . 1305
1.127.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 1307
1.127.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 1312

Internal problem ID [12544]
Internal file name [OUTPUT/11196_Wednesday_October_18_2023_03_47_05_AM_61324329/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 186.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "homogeneousTypeD", "homogeneousTypeD2", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[[_homogeneous, `class A`], _dAlembert]
\[
x \cos \left(\frac{y}{x}\right) y^{\prime}-y \cos \left(\frac{y}{x}\right)=-x
\]

\subsection*{1.127.1 Solving as homogeneousTypeD ode}

Writing the ode as
\[
\begin{equation*}
y^{\prime}=\frac{y}{x}-\frac{1}{\cos \left(\frac{y}{x}\right)} \tag{A}
\end{equation*}
\]

The given ode has the form
\[
\begin{equation*}
y^{\prime}=\frac{y}{x}+g(x) f\left(b \frac{y}{x}\right)^{\frac{n}{m}} \tag{1}
\end{equation*}
\]

Where \(b\) is scalar and \(g(x)\) is function of \(x\) and \(n, m\) are integers. The solution is given in Kamke page 20. Using the substitution \(y(x)=u(x) x\) then
\[
\frac{d y}{d x}=\frac{d u}{d x} x+u
\]

Hence the given ode becomes
\[
\begin{align*}
\frac{d u}{d x} x+u & =u+g(x) f(b u)^{\frac{n}{m}} \\
u^{\prime} & =\frac{1}{x} g(x) f(b u)^{\frac{n}{m}} \tag{2}
\end{align*}
\]

The above ode is always separable. This is easily solved for \(u\) assuming the integration can be resolved, and then the solution to the original ode becomes \(y=u x\). Comapring the given ode (A) with the form (1) shows that
\[
\begin{aligned}
g(x) & =-1 \\
b & =1 \\
f\left(\frac{b x}{y}\right) & =\cos \left(\frac{y}{x}\right)
\end{aligned}
\]

Substituting the above in (2) results in the \(u(x)\) ode as
\[
u^{\prime}(x)=-\frac{1}{x \cos (u(x))}
\]

Which is now solved as separable In canonical form the ODE is
\[
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =-\frac{\sec (u)}{x}
\end{aligned}
\]

Where \(f(x)=-\frac{1}{x}\) and \(g(u)=\sec (u)\). Integrating both sides gives
\[
\begin{aligned}
\frac{1}{\sec (u)} d u & =-\frac{1}{x} d x \\
\int \frac{1}{\sec (u)} d u & =\int-\frac{1}{x} d x \\
\sin (u) & =-\ln (x)+c_{1}
\end{aligned}
\]

The solution is
\[
\sin (u(x))+\ln (x)-c_{1}=0
\]

Therefore the solution is found using \(y=u x\). Hence
\[
\sin \left(\frac{y}{x}\right)+\ln (x)-c_{1}=0
\]

Summary
The solution(s) found are the following


Figure 200: Slope field plot

\section*{Verification of solutions}
\[
\sin \left(\frac{y}{x}\right)+\ln (x)-c_{1}=0
\]

Verified OK.

\subsection*{1.127.2 Solving as homogeneousTypeD2 ode}

Using the change of variables \(y=u(x) x\) on the above ode results in new ode in \(u(x)\)
\[
x \cos (u(x))\left(u^{\prime}(x) x+u(x)\right)-u(x) x \cos (u(x))=-x
\]

In canonical form the ODE is
\[
\begin{aligned}
u^{\prime} & =F(x, u) \\
& =f(x) g(u) \\
& =-\frac{\sec (u)}{x}
\end{aligned}
\]

Where \(f(x)=-\frac{1}{x}\) and \(g(u)=\sec (u)\). Integrating both sides gives
\[
\begin{aligned}
\frac{1}{\sec (u)} d u & =-\frac{1}{x} d x \\
\int \frac{1}{\sec (u)} d u & =\int-\frac{1}{x} d x \\
\sin (u) & =-\ln (x)+c_{2}
\end{aligned}
\]

The solution is
\[
\sin (u(x))+\ln (x)-c_{2}=0
\]

Replacing \(u(x)\) in the above solution by \(\frac{y}{x}\) results in the solution for \(y\) in implicit form
\[
\begin{aligned}
& \sin \left(\frac{y}{x}\right)+\ln (x)-c_{2}=0 \\
& \sin \left(\frac{y}{x}\right)+\ln (x)-c_{2}=0
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
\sin \left(\frac{y}{x}\right)+\ln (x)-c_{2}=0 \tag{1}
\end{equation*}
\]


Figure 201: Slope field plot

\section*{Verification of solutions}
\[
\sin \left(\frac{y}{x}\right)+\ln (x)-c_{2}=0
\]

Verified OK.

\subsection*{1.127.3 Solving as first order ode lie symmetry lookup ode}

Writing the ode as
\[
\begin{aligned}
& y^{\prime}=\frac{y \cos \left(\frac{y}{x}\right)-x}{x \cos \left(\frac{y}{x}\right)} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
\]

The condition of Lie symmetry is the linearized PDE given by
\[
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
\]

The type of this ode is known. It is of type homogeneous Type D. Therefore we do not need to solve the \(\operatorname{PDE}\) (A), and can just use the lookup table shown below to find \(\xi, \eta\)

Table 197: Lie symmetry infinitesimal lookup table for known first order ODE's
\begin{tabular}{|l|l|l|l|}
\hline ODE class & Form & \(\xi\) & \(\eta\) \\
\hline \hline linear ode & \(y^{\prime}=f(x) y(x)+g(x)\) & 0 & \(e^{\int f d x}\) \\
\hline separable ode & \(y^{\prime}=f(x) g(y)\) & \(\frac{1}{f}\) & 0 \\
\hline quadrature ode & \(y^{\prime}=f(x)\) & 0 & 1 \\
\hline quadrature ode & \(y^{\prime}=g(y)\) & 1 & 0 \\
\hline \begin{tabular}{l} 
homogeneous ODEs of \\
Class A
\end{tabular} & \(y^{\prime}=f\left(\frac{y}{x}\right)\) & \(x\) & \(y\) \\
\hline \begin{tabular}{l} 
homogeneous ODEs of \\
Class C
\end{tabular} & \(y^{\prime}=(a+b x+c y)^{\frac{n}{m}}\) & 1 & \(-\frac{b}{c}\) \\
\hline homogeneous class D & \(y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)\) & \(x^{2}\) & \(x y\) \\
\hline \begin{tabular}{l} 
First order \\
form ID 1
\end{tabular} & \(y^{2}=g(x) e^{h(x)+b y}+f(x)\) & \(\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}\) & \(\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}\) \\
\hline polynomial type ode & \(y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}\) & \(\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}\) & \(\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}\) \\
\hline Bernoulli ode & \(y^{\prime}=f(x) y+g(x) y^{n}\) & 0 & \(e^{-\int(n-1) f(x) d x} y^{n}\) \\
\hline Reduced Riccati & \(y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}\) & 0 & \(e^{-\int f_{1} d x}\) \\
\hline
\end{tabular}

The above table shows that
\[
\begin{align*}
& \xi(x, y)=x^{2} \\
& \eta(x, y)=x y \tag{A1}
\end{align*}
\]

The next step is to determine the canonical coordinates \(R, S\). The canonical coordinates map \((x, y) \rightarrow(R, S)\) where \((R, S)\) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is
\[
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
\]

The above comes from the requirements that \(\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1\). Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable \(R\) in the
canonical coordinates, where \(S(R)\). Therefore
\[
\begin{aligned}
\frac{d y}{d x} & =\frac{\eta}{\xi} \\
& =\frac{x y}{x^{2}} \\
& =\frac{y}{x}
\end{aligned}
\]

This is easily solved to give
\[
y=c_{1} x
\]

Where now the coordinate \(R\) is taken as the constant of integration. Hence
\[
R=\frac{y}{x}
\]

And \(S\) is found from
\[
\begin{aligned}
d S & =\frac{d x}{\xi} \\
& =\frac{d x}{x^{2}}
\end{aligned}
\]

Integrating gives
\[
\begin{aligned}
S & =\int \frac{d x}{T} \\
& =-\frac{1}{x}
\end{aligned}
\]

Where the constant of integration is set to zero as we just need one solution. Now that \(R, S\) are found, we need to setup the ode in these coordinates. This is done by evaluating
\[
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
\]

Where in the above \(R_{x}, R_{y}, S_{x}, S_{y}\) are all partial derivatives and \(\omega(x, y)\) is the right hand side of the original ode given by
\[
\omega(x, y)=\frac{y \cos \left(\frac{y}{x}\right)-x}{x \cos \left(\frac{y}{x}\right)}
\]

Evaluating all the partial derivatives gives
\[
\begin{aligned}
R_{x} & =-\frac{y}{x^{2}} \\
R_{y} & =\frac{1}{x} \\
S_{x} & =\frac{1}{x^{2}} \\
S_{y} & =0
\end{aligned}
\]

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.
\[
\begin{equation*}
\frac{d S}{d R}=-\frac{\cos \left(\frac{y}{x}\right)}{x} \tag{2~A}
\end{equation*}
\]

We now need to express the RHS as function of \(R\) only. This is done by solving for \(x, y\) in terms of \(R, S\) from the result obtained earlier and simplifying. This gives
\[
\frac{d S}{d R}=\cos (R) S(R)
\]

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates \(R, S\). Integrating the above gives
\[
\begin{equation*}
S(R)=c_{1} \mathrm{e}^{\sin (R)} \tag{4}
\end{equation*}
\]

To complete the solution, we just need to transform (4) back to \(x, y\) coordinates. This results in
\[
-\frac{1}{x}=c_{1} \mathrm{e}^{\sin \left(\frac{y}{x}\right)}
\]

Which simplifies to
\[
-\frac{1}{x}=c_{1} \mathrm{e}^{\sin \left(\frac{y}{x}\right)}
\]

Which gives
\[
y=\arcsin \left(\ln \left(-\frac{1}{c_{1} x}\right)\right) x
\]

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.
\begin{tabular}{|c|c|c|}
\hline Original ode in \(x, y\) coordinates & Canonical coordinates transformation & ODE in canonical coordinates
\[
(R, S)
\] \\
\hline \(\frac{d y}{d x}=\frac{y \cos \left(\frac{y}{x}\right)-x}{x \cos \left(\frac{y}{x}\right)}\) & & \(\frac{d S}{d R}=\cos (R) S(R)\) \\
\hline  & &  \\
\hline  & &  \\
\hline  & &  \\
\hline  & &  \\
\hline 为 & &  \\
\hline  & \(R=\frac{y}{x}\) &  \\
\hline  & &  \\
\hline  & \(S=-\frac{1}{x}\) &  \\
\hline  & \(x\) &  \\
\hline  & &  \\
\hline  & &  \\
\hline  & &  \\
\hline  & &  \\
\hline
\end{tabular}

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=\arcsin \left(\ln \left(-\frac{1}{c_{1} x}\right)\right) x \tag{1}
\end{equation*}
\]


Figure 202: Slope field plot

\section*{Verification of solutions}
\[
y=\arcsin \left(\ln \left(-\frac{1}{c_{1} x}\right)\right) x
\]

Verified OK.

\subsection*{1.127.4 Solving as exact ode}

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form
\[
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
\]

We assume there exists a function \(\phi(x, y)=c\) where \(c\) is constant, that satisfies the ode. Taking derivative of \(\phi\) w.r.t. \(x\) gives
\[
\frac{d}{d x} \phi(x, y)=0
\]

Hence
\[
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
\]

Comparing ( \(\mathrm{A}, \mathrm{B}\) ) shows that
\[
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
\]

But since \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) then for the above to be valid, we require that
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

If the above condition is satisfied, then the original ode is called exact. We still need to determine \(\phi(x, y)\) but at least we know now that we can do that since the condition \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is
\[
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
\]

Therefore
\[
\begin{align*}
\left(\cos \left(\frac{y}{x}\right) x\right) \mathrm{d} y & =\left(y \cos \left(\frac{y}{x}\right)-x\right) \mathrm{d} x \\
\left(-y \cos \left(\frac{y}{x}\right)+x\right) \mathrm{d} x+\left(\cos \left(\frac{y}{x}\right) x\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
\]

Comparing (1A) and (2A) shows that
\[
\begin{aligned}
M(x, y) & =-y \cos \left(\frac{y}{x}\right)+x \\
N(x, y) & =\cos \left(\frac{y}{x}\right) x
\end{aligned}
\]

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

Using result found above gives
\[
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-y \cos \left(\frac{y}{x}\right)+x\right) \\
& =-\cos \left(\frac{y}{x}\right)+\frac{y \sin \left(\frac{y}{x}\right)}{x}
\end{aligned}
\]

And
\[
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\cos \left(\frac{y}{x}\right) x\right) \\
& =\frac{y \sin \left(\frac{y}{x}\right)}{x}+\cos \left(\frac{y}{x}\right)
\end{aligned}
\]

Since \(\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}\), then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let
\[
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =\frac{\sec \left(\frac{y}{x}\right)}{x}\left(\left(-\cos \left(\frac{y}{x}\right)+\frac{y \sin \left(\frac{y}{x}\right)}{x}\right)-\left(\frac{y \sin \left(\frac{y}{x}\right)}{x}+\cos \left(\frac{y}{x}\right)\right)\right) \\
& =-\frac{2}{x}
\end{aligned}
\]

Since \(A\) does not depend on \(y\), then it can be used to find an integrating factor. The integrating factor \(\mu\) is
\[
\begin{aligned}
\mu & =e^{\int A \mathrm{~d} x} \\
& =e^{\int-\frac{2}{x} \mathrm{~d} x}
\end{aligned}
\]

The result of integrating gives
\[
\begin{aligned}
\mu & =e^{-2 \ln (x)} \\
& =\frac{1}{x^{2}}
\end{aligned}
\]
\(M\) and \(N\) are multiplied by this integrating factor, giving new \(M\) and new \(N\) which are called \(\bar{M}\) and \(\bar{N}\) for now so not to confuse them with the original \(M\) and \(N\).
\[
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{x^{2}}\left(-y \cos \left(\frac{y}{x}\right)+x\right) \\
& =\frac{-y \cos \left(\frac{y}{x}\right)+x}{x^{2}}
\end{aligned}
\]

And
\[
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{x^{2}}\left(\cos \left(\frac{y}{x}\right) x\right) \\
& =\frac{\cos \left(\frac{y}{x}\right)}{x}
\end{aligned}
\]

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is
\[
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x} & =0 \\
\left(\frac{-y \cos \left(\frac{y}{x}\right)+x}{x^{2}}\right)+\left(\frac{\cos \left(\frac{y}{x}\right)}{x}\right) \frac{\mathrm{d} y}{\mathrm{~d} x} & =0
\end{aligned}
\]

The following equations are now set up to solve for the function \(\phi(x, y)\)
\[
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
\]

Integrating (1) w.r.t. \(x\) gives
\[
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{-y \cos \left(\frac{y}{x}\right)+x}{x^{2}} \mathrm{~d} x \\
\phi & =\sin \left(\frac{y}{x}\right)-\ln \left(\frac{1}{x}\right)+f(y) \tag{3}
\end{align*}
\]

Where \(f(y)\) is used for the constant of integration since \(\phi\) is a function of both \(x\) and \(y\). Taking derivative of equation (3) w.r.t \(y\) gives
\[
\begin{equation*}
\frac{\partial \phi}{\partial y}=\frac{\cos \left(\frac{y}{x}\right)}{x}+f^{\prime}(y) \tag{4}
\end{equation*}
\]

But equation (2) says that \(\frac{\partial \phi}{\partial y}=\frac{\cos \left(\frac{y}{x}\right)}{x}\). Therefore equation (4) becomes
\[
\begin{equation*}
\frac{\cos \left(\frac{y}{x}\right)}{x}=\frac{\cos \left(\frac{y}{x}\right)}{x}+f^{\prime}(y) \tag{5}
\end{equation*}
\]

Solving equation (5) for \(f^{\prime}(y)\) gives
\[
f^{\prime}(y)=0
\]

Therefore
\[
f(y)=c_{1}
\]

Where \(c_{1}\) is constant of integration. Substituting this result for \(f(y)\) into equation (3) gives \(\phi\)
\[
\phi=\sin \left(\frac{y}{x}\right)-\ln \left(\frac{1}{x}\right)+c_{1}
\]

But since \(\phi\) itself is a constant function, then let \(\phi=c_{2}\) where \(c_{2}\) is new constant and combining \(c_{1}\) and \(c_{2}\) constants into new constant \(c_{1}\) gives the solution as
\[
c_{1}=\sin \left(\frac{y}{x}\right)-\ln \left(\frac{1}{x}\right)
\]

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
\sin \left(\frac{y}{x}\right)-\ln \left(\frac{1}{x}\right)=c_{1} \tag{1}
\end{equation*}
\]


Figure 203: Slope field plot
Verification of solutions
\[
\sin \left(\frac{y}{x}\right)-\ln \left(\frac{1}{x}\right)=c_{1}
\]

Verified OK.

Maple trace
```

`Methods for first order ODEs: --- Trying classification methods --- trying a quadrature trying 1st order linear trying Bernoulli trying separable trying inverse linear trying homogeneous types: trying homogeneous D <- homogeneous successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.0 (sec). Leaf size: 12
dsolve \((x * \cos (y(x) / x) * \operatorname{diff}(y(x), x)=y(x) * \cos (y(x) / x)-x, y(x)\), singsol=all)
\[
y(x)=-\arcsin \left(\ln (x)+c_{1}\right) x
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.628 (sec). Leaf size: 15
DSolve[x*Cos[y[x]/x]*y'[x]==y[x]*Cos[y[x]/x]-x,y[x],x,IncludeSingularSolutions -> True]
\[
y(x) \rightarrow x \arcsin \left(-\log (x)+c_{1}\right)
\]

\subsection*{1.128 problem 187}
1.128.1 Solving as second order linear constant coeff ode 1318
1.128.2 Solving using Kovacic algorithm . . . . . . . . . . . . . . . . . . 1321
1.128.3 Maple step by step solution 1326

Internal problem ID [12545]
Internal file name [OUTPUT/11197_Wednesday_October_18_2023_03_47_06_AM_21165008/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 187.
ODE order: 2.
ODE degree: 1 .

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type
[[_2nd_order, _linear, _nonhomogeneous]]
\[
y^{\prime \prime}-4 y=\mathrm{e}^{2 x} \sin (2 x)
\]

\subsection*{1.128.1 Solving as second order linear constant coeff ode}

This is second order non-homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)
\]

Where \(A=1, B=0, C=-4, f(x)=\mathrm{e}^{2 x} \sin (2 x)\). Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the non-homogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-4 y=0
\]

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is
\[
A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0
\]

Where in the above \(A=1, B=0, C=-4\). Let the solution be \(y=e^{\lambda x}\). Substituting this into the ODE gives
\[
\begin{equation*}
\lambda^{2} \mathrm{e}^{\lambda x}-4 \mathrm{e}^{\lambda x}=0 \tag{1}
\end{equation*}
\]

Since exponential function is never zero, then dividing \(\operatorname{Eq}(2)\) throughout by \(e^{\lambda x}\) gives
\[
\begin{equation*}
\lambda^{2}-4=0 \tag{2}
\end{equation*}
\]

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula
\[
\lambda_{1,2}=\frac{-B}{2 A} \pm \frac{1}{2 A} \sqrt{B^{2}-4 A C}
\]

Substituting \(A=1, B=0, C=-4\) into the above gives
\[
\begin{aligned}
\lambda_{1,2} & =\frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^{2}-(4)(1)(-4)} \\
& = \pm 2
\end{aligned}
\]

Hence
\[
\begin{aligned}
& \lambda_{1}=+2 \\
& \lambda_{2}=-2
\end{aligned}
\]

Which simplifies to
\[
\begin{aligned}
\lambda_{1} & =2 \\
\lambda_{2} & =-2
\end{aligned}
\]

Since roots are real and distinct, then the solution is
\[
\begin{aligned}
& y=c_{1} e^{\lambda_{1} x}+c_{2} e^{\lambda_{2} x} \\
& y=c_{1} e^{(2) x}+c_{2} e^{(-2) x}
\end{aligned}
\]

Or
\[
y=c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{-2 x}
\]

Therefore the homogeneous solution \(y_{h}\) is
\[
y_{h}=c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{-2 x}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
\mathrm{e}^{2 x} \sin (2 x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{2 x} \cos (2 x), \mathrm{e}^{2 x} \sin (2 x)\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\mathrm{e}^{-2 x}, \mathrm{e}^{2 x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{2 x} \cos (2 x)+A_{2} \mathrm{e}^{2 x} \sin (2 x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-8 A_{1} \mathrm{e}^{2 x} \sin (2 x)+8 A_{2} \mathrm{e}^{2 x} \cos (2 x)-4 A_{1} \mathrm{e}^{2 x} \cos (2 x)-4 A_{2} \mathrm{e}^{2 x} \sin (2 x)=\mathrm{e}^{2 x} \sin (2 x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-\frac{1}{10}, A_{2}=-\frac{1}{20}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=-\frac{\mathrm{e}^{2 x} \cos (2 x)}{10}-\frac{\mathrm{e}^{2 x} \sin (2 x)}{20}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{-2 x}\right)+\left(-\frac{\mathrm{e}^{2 x} \cos (2 x)}{10}-\frac{\mathrm{e}^{2 x} \sin (2 x)}{20}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{-2 x}-\frac{\mathrm{e}^{2 x} \cos (2 x)}{10}-\frac{\mathrm{e}^{2 x} \sin (2 x)}{20} \tag{1}
\end{equation*}
\]


Figure 204: Slope field plot

Verification of solutions
\[
y=c_{1} \mathrm{e}^{2 x}+c_{2} \mathrm{e}^{-2 x}-\frac{\mathrm{e}^{2 x} \cos (2 x)}{10}-\frac{\mathrm{e}^{2 x} \sin (2 x)}{20}
\]

Verified OK.

\subsection*{1.128.2 Solving using Kovacic algorithm}

Writing the ode as
\[
\begin{align*}
y^{\prime \prime}-4 y & =0  \tag{1}\\
A y^{\prime \prime}+B y^{\prime}+C y & =0 \tag{2}
\end{align*}
\]

Comparing (1) and (2) shows that
\[
\begin{align*}
& A=1 \\
& B=0  \tag{3}\\
& C=-4
\end{align*}
\]

Applying the Liouville transformation on the dependent variable gives
\[
z(x)=y e^{\int \frac{B}{2 A} d x}
\]

Then (2) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=r z(x) \tag{4}
\end{equation*}
\]

Where \(r\) is given by
\[
\begin{align*}
r & =\frac{s}{t}  \tag{5}\\
& =\frac{2 A B^{\prime}-2 B A^{\prime}+B^{2}-4 A C}{4 A^{2}}
\end{align*}
\]

Substituting the values of \(A, B, C\) from (3) in the above and simplifying gives
\[
\begin{equation*}
r=\frac{4}{1} \tag{6}
\end{equation*}
\]

Comparing the above to (5) shows that
\[
\begin{gathered}
s=4 \\
t=1
\end{gathered}
\]

Therefore eq. (4) becomes
\[
\begin{equation*}
z^{\prime \prime}(x)=4 z(x) \tag{7}
\end{equation*}
\]

Equation (7) is now solved. After finding \(z(x)\) then \(y\) is found using the inverse transformation
\[
y=z(x) e^{-\int \frac{B}{2 A} d x}
\]

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of \(r\) and the order of \(r\) at \(\infty\). The following table summarizes these cases.
\begin{tabular}{|l|l|l|}
\hline Case & Allowed pole order for \(r\) & Allowed value for \(\mathcal{O}(\infty)\) \\
\hline 1 & \(\{0,1,2,4,6,8, \cdots\}\) & \(\{\cdots,-6,-4,-2,0,2,3,4,5,6, \cdots\}\) \\
\hline 2 & \begin{tabular}{l} 
Need to have at least one pole that \\
is either order 2 or odd order greater \\
than 2. Any other pole order is \\
allowed as long as the above condi- \\
tion is satisfied. Hence the following \\
set of pole orders are all allowed. \\
\(\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}\).
\end{tabular} \\
\hline 3 & \(\{1,2\}\) & \(\{2,3,4,5,6,7, \cdots\}\) \\
\hline
\end{tabular}

Table 199: Necessary conditions for each Kovacic case

The order of \(r\) at \(\infty\) is the degree of \(t\) minus the degree of \(s\). Therefore
\[
\begin{aligned}
O(\infty) & =\operatorname{deg}(t)-\operatorname{deg}(s) \\
& =0-0 \\
& =0
\end{aligned}
\]

There are no poles in \(r\). Therefore the set of poles \(\Gamma\) is empty. Since there is no odd order pole larger than 2 and the order at \(\infty\) is 0 then the necessary conditions for case one are met. Therefore
\[
L=[1]
\]

Since \(r=4\) is not a function of \(x\), then there is no need run Kovacic algorithm to obtain a solution for transformed ode \(z^{\prime \prime}=r z\) as one solution is
\[
z_{1}(x)=\mathrm{e}^{-2 x}
\]

Using the above, the solution for the original ode can now be found. The first solution to the original ode in \(y\) is found from
\[
y_{1}=z_{1} e^{\int-\frac{1}{2} \frac{B}{A} d x}
\]

Since \(B=0\) then the above reduces to
\[
\begin{aligned}
y_{1} & =z_{1} \\
& =\mathrm{e}^{-2 x}
\end{aligned}
\]

Which simplifies to
\[
y_{1}=\mathrm{e}^{-2 x}
\]

The second solution \(y_{2}\) to the original ode is found using reduction of order
\[
y_{2}=y_{1} \int \frac{e^{\int-\frac{B}{A} d x}}{y_{1}^{2}} d x
\]

Since \(B=0\) then the above becomes
\[
\begin{aligned}
y_{2} & =y_{1} \int \frac{1}{y_{1}^{2}} d x \\
& =\mathrm{e}^{-2 x} \int \frac{1}{\mathrm{e}^{-4 x}} d x \\
& =\mathrm{e}^{-2 x}\left(\frac{\mathrm{e}^{4 x}}{4}\right)
\end{aligned}
\]

Therefore the solution is
\[
\begin{aligned}
y & =c_{1} y_{1}+c_{2} y_{2} \\
& =c_{1}\left(\mathrm{e}^{-2 x}\right)+c_{2}\left(\mathrm{e}^{-2 x}\left(\frac{\mathrm{e}^{4 x}}{4}\right)\right)
\end{aligned}
\]

This is second order nonhomogeneous ODE. Let the solution be
\[
y=y_{h}+y_{p}
\]

Where \(y_{h}\) is the solution to the homogeneous \(\operatorname{ODE} A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=0\), and \(y_{p}\) is a particular solution to the nonhomogeneous ODE \(A y^{\prime \prime}(x)+B y^{\prime}(x)+C y(x)=f(x)\). \(y_{h}\) is the solution to
\[
y^{\prime \prime}-4 y=0
\]

The homogeneous solution is found using the Kovacic algorithm which results in
\[
y_{h}=c_{1} \mathrm{e}^{-2 x}+\frac{c_{2} \mathrm{e}^{2 x}}{4}
\]

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is
\[
\mathrm{e}^{2 x} \sin (2 x)
\]

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is
\[
\left[\left\{\mathrm{e}^{2 x} \cos (2 x), \mathrm{e}^{2 x} \sin (2 x)\right\}\right]
\]

While the set of the basis functions for the homogeneous solution found earlier is
\[
\left\{\frac{\mathrm{e}^{2 x}}{4}, \mathrm{e}^{-2 x}\right\}
\]

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.
\[
y_{p}=A_{1} \mathrm{e}^{2 x} \cos (2 x)+A_{2} \mathrm{e}^{2 x} \sin (2 x)
\]

The unknowns \(\left\{A_{1}, A_{2}\right\}\) are found by substituting the above trial solution \(y_{p}\) into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives
\[
-8 A_{1} \mathrm{e}^{2 x} \sin (2 x)+8 A_{2} \mathrm{e}^{2 x} \cos (2 x)-4 A_{1} \mathrm{e}^{2 x} \cos (2 x)-4 A_{2} \mathrm{e}^{2 x} \sin (2 x)=\mathrm{e}^{2 x} \sin (2 x)
\]

Solving for the unknowns by comparing coefficients results in
\[
\left[A_{1}=-\frac{1}{10}, A_{2}=-\frac{1}{20}\right]
\]

Substituting the above back in the above trial solution \(y_{p}\), gives the particular solution
\[
y_{p}=-\frac{\mathrm{e}^{2 x} \cos (2 x)}{10}-\frac{\mathrm{e}^{2 x} \sin (2 x)}{20}
\]

Therefore the general solution is
\[
\begin{aligned}
y & =y_{h}+y_{p} \\
& =\left(c_{1} \mathrm{e}^{-2 x}+\frac{c_{2} \mathrm{e}^{2 x}}{4}\right)+\left(-\frac{\mathrm{e}^{2 x} \cos (2 x)}{10}-\frac{\mathrm{e}^{2 x} \sin (2 x)}{20}\right)
\end{aligned}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=c_{1} \mathrm{e}^{-2 x}+\frac{c_{2} \mathrm{e}^{2 x}}{4}-\frac{\mathrm{e}^{2 x} \cos (2 x)}{10}-\frac{\mathrm{e}^{2 x} \sin (2 x)}{20} \tag{1}
\end{equation*}
\]


Figure 205: Slope field plot

\section*{Verification of solutions}
\[
y=c_{1} \mathrm{e}^{-2 x}+\frac{c_{2} \mathrm{e}^{2 x}}{4}-\frac{\mathrm{e}^{2 x} \cos (2 x)}{10}-\frac{\mathrm{e}^{2 x} \sin (2 x)}{20}
\]

Verified OK.

\subsection*{1.128.3 Maple step by step solution}

Let's solve
\[
y^{\prime \prime}-4 y=\mathrm{e}^{2 x} \sin (2 x)
\]
- Highest derivative means the order of the ODE is 2
\[
y^{\prime \prime}
\]
- Characteristic polynomial of homogeneous ODE
\[
r^{2}-4=0
\]
- Factor the characteristic polynomial
\((r-2)(r+2)=0\)
- Roots of the characteristic polynomial
\(r=(-2,2)\)
- \(\quad 1\) st solution of the homogeneous ODE
\(y_{1}(x)=\mathrm{e}^{-2 x}\)
- \(\quad\) 2nd solution of the homogeneous ODE
\(y_{2}(x)=\mathrm{e}^{2 x}\)
- General solution of the ODE
\(y=c_{1} y_{1}(x)+c_{2} y_{2}(x)+y_{p}(x)\)
- Substitute in solutions of the homogeneous ODE
\(y=c_{1} \mathrm{e}^{-2 x}+c_{2} \mathrm{e}^{2 x}+y_{p}(x)\)
Find a particular solution \(y_{p}(x)\) of the ODE
- Use variation of parameters to find \(y_{p}\) here \(f(x)\) is the forcing function
\[
\left[y_{p}(x)=-y_{1}(x)\left(\int \frac{y_{2}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right)+y_{2}(x)\left(\int \frac{y_{1}(x) f(x)}{W\left(y_{1}(x), y_{2}(x)\right)} d x\right), f(x)=\mathrm{e}^{2 x} \sin (2 x)\right]
\]
- Wronskian of solutions of the homogeneous equation
\(W\left(y_{1}(x), y_{2}(x)\right)=\left[\begin{array}{cc}\mathrm{e}^{-2 x} & \mathrm{e}^{2 x} \\ -2 \mathrm{e}^{-2 x} & 2 \mathrm{e}^{2 x}\end{array}\right]\)
- Compute Wronskian
\(W\left(y_{1}(x), y_{2}(x)\right)=4\)
- Substitute functions into equation for \(y_{p}(x)\)
\[
y_{p}(x)=-\frac{\mathrm{e}^{-2 x}\left(\int \mathrm{e}^{4 x} \sin (2 x) d x\right)}{4}+\frac{\mathrm{e}^{2 x}\left(\int \sin (2 x) d x\right)}{4}
\]
- Compute integrals
\[
y_{p}(x)=-\frac{\mathrm{e}^{2 x}(\sin (2 x)+2 \cos (2 x))}{20}
\]
- Substitute particular solution into general solution to ODE
\[
y=c_{1} \mathrm{e}^{-2 x}+c_{2} \mathrm{e}^{2 x}-\frac{\mathrm{e}^{2 x}(\sin (2 x)+2 \cos (2 x))}{20}
\]

Maple trace
```

`Methods for second order ODEs: --- Trying classification methods --- trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE     checking if the LODE has constant coefficients     <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.016 (sec). Leaf size: 33
```

dsolve(diff(y(x),x\$2)-4*y(x)=exp(2*x)*sin(2*x),y(x), singsol=all)

```
\[
y(x)=\frac{\left(20 c_{2}-2 \cos (2 x)-\sin (2 x)\right) \mathrm{e}^{2 x}}{20}+\mathrm{e}^{-2 x} c_{1}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.152 (sec). Leaf size: 42
DSolve[y''[x]-4*y[x]==Exp[2*x]*Sin[2*x],y[x],x,IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow c_{1} e^{2 x}+c_{2} e^{-2 x}-\frac{1}{20} e^{2 x}(\sin (2 x)+2 \cos (2 x))
\]

\subsection*{1.129 problem 188}
1.129.1 Solving as first order ode lie symmetry lookup ode . . . . . . . 1329
1.129.2 Solving as bernoulli ode . . . . . . . . . . . . . . . . . . . . . . 1333
1.129.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 1337
1.129.4 Solving as riccati ode . . . . . . . . . . . . . . . . . . . . . . . . 1342

Internal problem ID [12546]
Internal file name [OUTPUT/11198_Wednesday_October_18_2023_03_47_06_AM_57167490/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 188.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "riccati", "bernoulli", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_Bernoulli]
\[
y^{\prime} x+y-y^{2} \ln (x)=0
\]

\subsection*{1.129.1 Solving as first order ode lie symmetry lookup ode}

Writing the ode as
\[
\begin{aligned}
& y^{\prime}=\frac{y(y \ln (x)-1)}{x} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
\]

The condition of Lie symmetry is the linearized PDE given by
\[
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
\]

The type of this ode is known. It is of type Bernoulli. Therefore we do not need to solve the \(\operatorname{PDE}(\mathrm{A})\), and can just use the lookup table shown below to find \(\xi, \eta\)

Table 201: Lie symmetry infinitesimal lookup table for known first order ODE's
\begin{tabular}{|l|l|l|l|}
\hline ODE class & Form & \(\xi\) & \(\eta\) \\
\hline \hline linear ode & \(y^{\prime}=f(x) y(x)+g(x)\) & 0 & \(e^{\int f d x}\) \\
\hline separable ode & \(y^{\prime}=f(x) g(y)\) & \(\frac{1}{f}\) & 0 \\
\hline quadrature ode & \(y^{\prime}=f(x)\) & 0 & 1 \\
\hline quadrature ode & \(y^{\prime}=g(y)\) & 1 & 0 \\
\hline \begin{tabular}{l} 
homogeneous ODEs of \\
Class A
\end{tabular} & \(y^{\prime}=f\left(\frac{y}{x}\right)\) & \(x\) & \(y\) \\
\hline \begin{tabular}{l} 
homogeneous ODEs of \\
Class C
\end{tabular} & \(y^{\prime}=(a+b x+c y)^{\frac{n}{m}}\) & 1 & \(-\frac{b}{c}\) \\
\hline homogeneous class D & \(y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)\) & \(x^{2}\) & \(x y\) \\
\hline \begin{tabular}{l} 
First order \\
form ID 1
\end{tabular} & \(y^{2}=g(x) e^{h(x)+b y}+f(x)\) & \(\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}\) & \(\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}\) \\
\hline polynomial type ode & \(y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}\) & \(\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}\) & \(\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}\) \\
\hline Bernoulli ode & \(y^{\prime}=f(x) y+g(x) y^{n}\) & 0 & \(e^{-\int(n-1) f(x) d x} y^{n}\) \\
\hline Reduced Riccati & \(y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}\) & 0 & \(e^{-\int f_{1} d x}\) \\
\hline
\end{tabular}

The above table shows that
\[
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=x y^{2} \tag{A1}
\end{align*}
\]

The next step is to determine the canonical coordinates \(R, S\). The canonical coordinates map \((x, y) \rightarrow(R, S)\) where \((R, S)\) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is
\[
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
\]

The above comes from the requirements that \(\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1\). Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable \(R\) in the
canonical coordinates, where \(S(R)\). Since \(\xi=0\) then in this special case
\[
R=x
\]
\(S\) is found from
\[
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{x y^{2}} d y
\end{aligned}
\]

Which results in
\[
S=-\frac{1}{x y}
\]

Now that \(R, S\) are found, we need to setup the ode in these coordinates. This is done by evaluating
\[
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
\]

Where in the above \(R_{x}, R_{y}, S_{x}, S_{y}\) are all partial derivatives and \(\omega(x, y)\) is the right hand side of the original ode given by
\[
\omega(x, y)=\frac{y(y \ln (x)-1)}{x}
\]

Evaluating all the partial derivatives gives
\[
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{1}{x^{2} y} \\
S_{y} & =\frac{1}{x y^{2}}
\end{aligned}
\]

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.
\[
\begin{equation*}
\frac{d S}{d R}=\frac{\ln (x)}{x^{2}} \tag{2~A}
\end{equation*}
\]

We now need to express the RHS as function of \(R\) only. This is done by solving for \(x, y\) in terms of \(R, S\) from the result obtained earlier and simplifying. This gives
\[
\frac{d S}{d R}=\frac{\ln (R)}{R^{2}}
\]

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates \(R, S\). Integrating the above gives
\[
\begin{equation*}
S(R)=-\frac{\ln (R)}{R}-\frac{1}{R}+c_{1} \tag{4}
\end{equation*}
\]

To complete the solution, we just need to transform (4) back to \(x, y\) coordinates. This results in
\[
-\frac{1}{y x}=-\frac{\ln (x)}{x}-\frac{1}{x}+c_{1}
\]

Which simplifies to
\[
\frac{-c_{1} x y+y \ln (x)+y-1}{x y}=0
\]

Which gives
\[
y=\frac{1}{-c_{1} x+\ln (x)+1}
\]

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.


\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=\frac{1}{-c_{1} x+\ln (x)+1} \tag{1}
\end{equation*}
\]


Figure 206: Slope field plot
Verification of solutions
\[
y=\frac{1}{-c_{1} x+\ln (x)+1}
\]

Verified OK.

\subsection*{1.129.2 Solving as bernoulli ode}

In canonical form, the ODE is
\[
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{y(y \ln (x)-1)}{x}
\end{aligned}
\]

This is a Bernoulli ODE.
\[
\begin{equation*}
y^{\prime}=-\frac{1}{x} y+\frac{\ln (x)}{x} y^{2} \tag{1}
\end{equation*}
\]

The standard Bernoulli ODE has the form
\[
\begin{equation*}
y^{\prime}=f_{0}(x) y+f_{1}(x) y^{n} \tag{2}
\end{equation*}
\]

The first step is to divide the above equation by \(y^{n}\) which gives
\[
\begin{equation*}
\frac{y^{\prime}}{y^{n}}=f_{0}(x) y^{1-n}+f_{1}(x) \tag{3}
\end{equation*}
\]

The next step is use the substitution \(w=y^{1-n}\) in equation (3) which generates a new ODE in \(w(x)\) which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution \(y(x)\) which is what we want.
This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that
\[
\begin{aligned}
f_{0}(x) & =-\frac{1}{x} \\
f_{1}(x) & =\frac{\ln (x)}{x} \\
n & =2
\end{aligned}
\]

Dividing both sides of ODE (1) by \(y^{n}=y^{2}\) gives
\[
\begin{equation*}
y^{\prime} \frac{1}{y^{2}}=-\frac{1}{x y}+\frac{\ln (x)}{x} \tag{4}
\end{equation*}
\]

Let
\[
\begin{align*}
w & =y^{1-n} \\
& =\frac{1}{y} \tag{5}
\end{align*}
\]

Taking derivative of equation (5) w.r.t \(x\) gives
\[
\begin{equation*}
w^{\prime}=-\frac{1}{y^{2}} y^{\prime} \tag{6}
\end{equation*}
\]

Substituting equations (5) and (6) into equation (4) gives
\[
\begin{align*}
-w^{\prime}(x) & =-\frac{w(x)}{x}+\frac{\ln (x)}{x} \\
w^{\prime} & =\frac{w}{x}-\frac{\ln (x)}{x} \tag{7}
\end{align*}
\]

The above now is a linear ODE in \(w(x)\) which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is
\[
w^{\prime}(x)+p(x) w(x)=q(x)
\]

Where here
\[
\begin{aligned}
p(x) & =-\frac{1}{x} \\
q(x) & =-\frac{\ln (x)}{x}
\end{aligned}
\]

Hence the ode is
\[
w^{\prime}(x)-\frac{w(x)}{x}=-\frac{\ln (x)}{x}
\]

The integrating factor \(\mu\) is
\[
\begin{aligned}
\mu & =\mathrm{e}^{\int-\frac{1}{x} d x} \\
& =\frac{1}{x}
\end{aligned}
\]

The ode becomes
\[
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu w) & =(\mu)\left(-\frac{\ln (x)}{x}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}\left(\frac{w}{x}\right) & =\left(\frac{1}{x}\right)\left(-\frac{\ln (x)}{x}\right) \\
\mathrm{d}\left(\frac{w}{x}\right) & =\left(-\frac{\ln (x)}{x^{2}}\right) \mathrm{d} x
\end{aligned}
\]

Integrating gives
\[
\begin{aligned}
& \frac{w}{x}=\int-\frac{\ln (x)}{x^{2}} \mathrm{~d} x \\
& \frac{w}{x}=\frac{\ln (x)}{x}+\frac{1}{x}+c_{1}
\end{aligned}
\]

Dividing both sides by the integrating factor \(\mu=\frac{1}{x}\) results in
\[
w(x)=x\left(\frac{\ln (x)}{x}+\frac{1}{x}\right)+c_{1} x
\]
which simplifies to
\[
w(x)=c_{1} x+\ln (x)+1
\]

Replacing \(w\) in the above by \(\frac{1}{y}\) using equation (5) gives the final solution.
\[
\frac{1}{y}=c_{1} x+\ln (x)+1
\]

Or
\[
y=\frac{1}{c_{1} x+\ln (x)+1}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{1}{c_{1} x+\ln (x)+1} \tag{1}
\end{equation*}
\]


Figure 207: Slope field plot

Verification of solutions
\[
y=\frac{1}{c_{1} x+\ln (x)+1}
\]

Verified OK.

\subsection*{1.129.3 Solving as exact ode}

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form
\[
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
\]

We assume there exists a function \(\phi(x, y)=c\) where \(c\) is constant, that satisfies the ode. Taking derivative of \(\phi\) w.r.t. \(x\) gives
\[
\frac{d}{d x} \phi(x, y)=0
\]

Hence
\[
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
\]

Comparing ( \(\mathrm{A}, \mathrm{B}\) ) shows that
\[
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
\]

But since \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) then for the above to be valid, we require that
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

If the above condition is satisfied, then the original ode is called exact. We still need to determine \(\phi(x, y)\) but at least we know now that we can do that since the condition \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is
\[
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
\]

Therefore
\[
\begin{align*}
(x) \mathrm{d} y & =\left(-y+y^{2} \ln (x)\right) \mathrm{d} x \\
\left(-y^{2} \ln (x)+y\right) \mathrm{d} x+(x) \mathrm{d} y & =0 \tag{2A}
\end{align*}
\]

Comparing (1A) and (2A) shows that
\[
\begin{aligned}
M(x, y) & =-y^{2} \ln (x)+y \\
N(x, y) & =x
\end{aligned}
\]

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

Using result found above gives
\[
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-y^{2} \ln (x)+y\right) \\
& =-2 y \ln (x)+1
\end{aligned}
\]

And
\[
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(x) \\
& =1
\end{aligned}
\]

Since \(\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}\), then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let
\[
\begin{aligned}
A & =\frac{1}{N}\left(\frac{\partial M}{\partial y}-\frac{\partial N}{\partial x}\right) \\
& =\frac{1}{x}((-2 y \ln (x)+1)-(1)) \\
& =-\frac{2 y \ln (x)}{x}
\end{aligned}
\]

Since \(A\) depends on \(y\), it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let
\[
\begin{aligned}
B & =\frac{1}{M}\left(\frac{\partial N}{\partial x}-\frac{\partial M}{\partial y}\right) \\
& =-\frac{1}{(y \ln (x)-1) y}((1)-(-2 y \ln (x)+1)) \\
& =-\frac{2 \ln (x)}{y \ln (x)-1}
\end{aligned}
\]

Since \(B\) depends on \(x\), it can not be used to obtain an integrating factor. We will now try a third method to find an integrating factor. Let
\[
R=\frac{\frac{\partial N}{\partial x}-\frac{\partial M}{\partial y}}{x M-y N}
\]
\(R\) is now checked to see if it is a function of only \(t=x y\). Therefore
\[
\begin{aligned}
R & =\frac{\frac{\partial N}{\partial x}-\frac{\partial M}{\partial y}}{x M-y N} \\
& =\frac{(1)-(-2 y \ln (x)+1)}{x\left(-y^{2} \ln (x)+y\right)-y(x)} \\
& =-\frac{2}{x y}
\end{aligned}
\]

Replacing all powers of terms \(x y\) by \(t\) gives
\[
R=-\frac{2}{t}
\]

Since \(R\) depends on \(t\) only, then it can be used to find an integrating factor. Let the integrating factor be \(\mu\) then
\[
\begin{aligned}
\mu & =e^{\int R \mathrm{~d} t} \\
& =e^{\int\left(-\frac{2}{t}\right) \mathrm{d} t}
\end{aligned}
\]

The result of integrating gives
\[
\begin{aligned}
\mu & =e^{-2 \ln (t)} \\
& =\frac{1}{t^{2}}
\end{aligned}
\]

Now \(t\) is replaced back with \(x y\) giving
\[
\mu=\frac{1}{x^{2} y^{2}}
\]

Multiplying \(M\) and \(N\) by this integrating factor gives new \(M\) and new \(N\) which are called \(\bar{M}\) and \(\bar{N}\) so not to confuse them with the original \(M\) and \(N\)
\[
\begin{aligned}
\bar{M} & =\mu M \\
& =\frac{1}{x^{2} y^{2}}\left(-y^{2} \ln (x)+y\right) \\
& =\frac{-y \ln (x)+1}{x^{2} y}
\end{aligned}
\]

And
\[
\begin{aligned}
\bar{N} & =\mu N \\
& =\frac{1}{x^{2} y^{2}}(x) \\
& =\frac{1}{x y^{2}}
\end{aligned}
\]

A modified ODE is now obtained from the original ODE, which is exact and can solved. The modified ODE is
\[
\begin{aligned}
\bar{M}+\bar{N} \frac{\mathrm{~d} y}{\mathrm{~d} x} & =0 \\
\left(\frac{-y \ln (x)+1}{x^{2} y}\right)+\left(\frac{1}{x y^{2}}\right) \frac{\mathrm{d} y}{\mathrm{~d} x} & =0
\end{aligned}
\]

The following equations are now set up to solve for the function \(\phi(x, y)\)
\[
\begin{align*}
& \frac{\partial \phi}{\partial x}=\bar{M}  \tag{1}\\
& \frac{\partial \phi}{\partial y}=\bar{N} \tag{2}
\end{align*}
\]

Integrating (1) w.r.t. \(x\) gives
\[
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \bar{M} \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int \frac{-y \ln (x)+1}{x^{2} y} \mathrm{~d} x \\
\phi & =\frac{y \ln (x)+y-1}{x y}+f(y) \tag{3}
\end{align*}
\]

Where \(f(y)\) is used for the constant of integration since \(\phi\) is a function of both \(x\) and \(y\). Taking derivative of equation (3) w.r.t \(y\) gives
\[
\begin{align*}
\frac{\partial \phi}{\partial y} & =\frac{\ln (x)+1}{x y}-\frac{y \ln (x)+y-1}{x y^{2}}+f^{\prime}(y)  \tag{4}\\
& =\frac{1}{x y^{2}}+f^{\prime}(y)
\end{align*}
\]

But equation (2) says that \(\frac{\partial \phi}{\partial y}=\frac{1}{x y^{2}}\). Therefore equation (4) becomes
\[
\begin{equation*}
\frac{1}{x y^{2}}=\frac{1}{x y^{2}}+f^{\prime}(y) \tag{5}
\end{equation*}
\]

Solving equation (5) for \(f^{\prime}(y)\) gives
\[
f^{\prime}(y)=0
\]

Therefore
\[
f(y)=c_{1}
\]

Where \(c_{1}\) is constant of integration. Substituting this result for \(f(y)\) into equation (3) gives \(\phi\)
\[
\phi=\frac{y \ln (x)+y-1}{x y}+c_{1}
\]

But since \(\phi\) itself is a constant function, then let \(\phi=c_{2}\) where \(c_{2}\) is new constant and combining \(c_{1}\) and \(c_{2}\) constants into new constant \(c_{1}\) gives the solution as
\[
c_{1}=\frac{y \ln (x)+y-1}{x y}
\]

The solution becomes
\[
y=\frac{1}{-c_{1} x+\ln (x)+1}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{1}{-c_{1} x+\ln (x)+1} \tag{1}
\end{equation*}
\]


Figure 208: Slope field plot

\section*{Verification of solutions}
\[
y=\frac{1}{-c_{1} x+\ln (x)+1}
\]

Verified OK.

\subsection*{1.129.4 Solving as riccati ode}

In canonical form the ODE is
\[
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =\frac{y(y \ln (x)-1)}{x}
\end{aligned}
\]

This is a Riccati ODE. Comparing the ODE to solve
\[
y^{\prime}=\frac{y^{2} \ln (x)}{x}-\frac{y}{x}
\]

With Riccati ODE standard form
\[
y^{\prime}=f_{0}(x)+f_{1}(x) y+f_{2}(x) y^{2}
\]

Shows that \(f_{0}(x)=0, f_{1}(x)=-\frac{1}{x}\) and \(f_{2}(x)=\frac{\ln (x)}{x}\). Let
\[
\begin{align*}
y & =\frac{-u^{\prime}}{f_{2} u} \\
& =\frac{-u^{\prime}}{\frac{\ln (x) u}{x}} \tag{1}
\end{align*}
\]

Using the above substitution in the given ODE results (after some simplification)in a second order ODE to solve for \(u(x)\) which is
\[
\begin{equation*}
f_{2} u^{\prime \prime}(x)-\left(f_{2}^{\prime}+f_{1} f_{2}\right) u^{\prime}(x)+f_{2}^{2} f_{0} u(x)=0 \tag{2}
\end{equation*}
\]

But
\[
\begin{aligned}
f_{2}^{\prime} & =-\frac{\ln (x)}{x^{2}}+\frac{1}{x^{2}} \\
f_{1} f_{2} & =-\frac{\ln (x)}{x^{2}} \\
f_{2}^{2} f_{0} & =0
\end{aligned}
\]

Substituting the above terms back in equation (2) gives
\[
\frac{\ln (x) u^{\prime \prime}(x)}{x}-\left(-\frac{2 \ln (x)}{x^{2}}+\frac{1}{x^{2}}\right) u^{\prime}(x)=0
\]

Solving the above ODE (this ode solved using Maple, not this program), gives
\[
u(x)=\frac{-\ln (x) c_{2}+c_{1} x-c_{2}}{x}
\]

The above shows that
\[
u^{\prime}(x)=\frac{\ln (x) c_{2}}{x^{2}}
\]

Using the above in (1) gives the solution
\[
y=-\frac{c_{2}}{-\ln (x) c_{2}+c_{1} x-c_{2}}
\]

Dividing both numerator and denominator by \(c_{1}\) gives, after renaming the constant \(\frac{c_{2}}{c_{1}}=c_{3}\) the following solution
\[
y=\frac{1}{-c_{3} x+\ln (x)+1}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{1}{-c_{3} x+\ln (x)+1} \tag{1}
\end{equation*}
\]


Figure 209: Slope field plot

Verification of solutions
\[
y=\frac{1}{-c_{3} x+\ln (x)+1}
\]

Verified OK.
Maple trace
```

`Methods for first order ODEs: --- Trying classification methods --- trying a quadrature trying 1st order linear trying Bernoulli <- Bernoulli successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.016 (sec). Leaf size: 13
dsolve( \(x * \operatorname{diff}(y(x), x)+y(x)-y(x)^{\wedge} 2 * \ln (x)=0, y(x)\), singsol=all)
\[
y(x)=\frac{1}{1+c_{1} x+\ln (x)}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.237 (sec). Leaf size: 20
DSolve \([x * y\) ' \([x]+y[x]-y[x] \sim 2 * \log [x]==0, y[x], x\), IncludeSingularSolutions \(\rightarrow\) True]
\[
\begin{aligned}
& y(x) \rightarrow \frac{1}{\log (x)+c_{1} x+1} \\
& y(x) \rightarrow 0
\end{aligned}
\]

\subsection*{1.130 problem 189}
1.130.1 Solving as first order ode lie symmetry calculated ode

Internal problem ID [12547]
Internal file name [OUTPUT/11199_Wednesday_October_18_2023_03_47_07_AM_2811439/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 189.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "first__order_ode__lie_symmetry_calculated"

Maple gives the following as the ode type
```

[[_homogeneous, `class C`], _rational, [_Abel, `2nd type`, `     class A`]]

```
\[
2 y+(x+y-2) y^{\prime}=-2 x+1
\]

\subsection*{1.130.1 Solving as first order ode lie symmetry calculated ode}

Writing the ode as
\[
\begin{aligned}
& y^{\prime}=-\frac{2 x+2 y-1}{x+y-2} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
\]

The condition of Lie symmetry is the linearized PDE given by
\[
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
\]

The type of this ode is not in the lookup table. To determine \(\xi, \eta\) then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives
\[
\begin{align*}
& \xi=x a_{2}+y a_{3}+a_{1}  \tag{1E}\\
& \eta=x b_{2}+y b_{3}+b_{1} \tag{2E}
\end{align*}
\]

Where the unknown coefficients are
\[
\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}
\]

Substituting equations (1E,2E) and \(\omega\) into (A) gives
\[
\begin{align*}
b_{2} & -\frac{(2 x+2 y-1)\left(b_{3}-a_{2}\right)}{x+y-2}-\frac{(2 x+2 y-1)^{2} a_{3}}{(x+y-2)^{2}} \\
& -\left(-\frac{2}{x+y-2}+\frac{2 x+2 y-1}{(x+y-2)^{2}}\right)\left(x a_{2}+y a_{3}+a_{1}\right)  \tag{5E}\\
& -\left(-\frac{2}{x+y-2}+\frac{2 x+2 y-1}{(x+y-2)^{2}}\right)\left(x b_{2}+y b_{3}+b_{1}\right)=0
\end{align*}
\]

Putting the above in normal form gives
\[
\begin{aligned}
& \frac{2 x^{2} a_{2}-4 x^{2} a_{3}+x^{2} b_{2}-2 x^{2} b_{3}+4 x y a_{2}-8 x y a_{3}+2 x y b_{2}-4 x y b_{3}+2 y^{2} a_{2}-4 y^{2} a_{3}+y^{2} b_{2}-2 y^{2} b_{3}-8 x a_{2}}{(x+y-2)^{2}} \\
& =0
\end{aligned}
\]

Setting the numerator to zero gives
\[
\begin{align*}
& 2 x^{2} a_{2}-4 x^{2} a_{3}+x^{2} b_{2}-2 x^{2} b_{3}+4 x y a_{2}-8 x y a_{3}+2 x y b_{2}-4 x y b_{3}  \tag{6E}\\
& \quad+2 y^{2} a_{2}-4 y^{2} a_{3}+y^{2} b_{2}-2 y^{2} b_{3}-8 x a_{2}+4 x a_{3}-7 x b_{2}+5 x b_{3} \\
& \quad-5 y a_{2}+y a_{3}-4 y b_{2}+2 y b_{3}-3 a_{1}+2 a_{2}-a_{3}-3 b_{1}+4 b_{2}-2 b_{3}=0
\end{align*}
\]

Looking at the above PDE shows the following are all the terms with \(\{x, y\}\) in them.
\[
\{x, y\}
\]

The following substitution is now made to be able to collect on all terms with \(\{x, y\}\) in them
\[
\left\{x=v_{1}, y=v_{2}\right\}
\]

The above PDE (6E) now becomes
\[
\begin{align*}
& 2 a_{2} v_{1}^{2}+4 a_{2} v_{1} v_{2}+2 a_{2} v_{2}^{2}-4 a_{3} v_{1}^{2}-8 a_{3} v_{1} v_{2}-4 a_{3} v_{2}^{2}+b_{2} v_{1}^{2}+2 b_{2} v_{1} v_{2}  \tag{7E}\\
& \quad+b_{2} v_{2}^{2}-2 b_{3} v_{1}^{2}-4 b_{3} v_{1} v_{2}-2 b_{3} v_{2}^{2}-8 a_{2} v_{1}-5 a_{2} v_{2}+4 a_{3} v_{1}+a_{3} v_{2} \\
& \quad-7 b_{2} v_{1}-4 b_{2} v_{2}+5 b_{3} v_{1}+2 b_{3} v_{2}-3 a_{1}+2 a_{2}-a_{3}-3 b_{1}+4 b_{2}-2 b_{3}=0
\end{align*}
\]

Collecting the above on the terms \(v_{i}\) introduced, and these are
\[
\left\{v_{1}, v_{2}\right\}
\]

Equation (7E) now becomes
\[
\begin{align*}
& \left(2 a_{2}-4 a_{3}+b_{2}-2 b_{3}\right) v_{1}^{2}+\left(4 a_{2}-8 a_{3}+2 b_{2}-4 b_{3}\right) v_{1} v_{2}  \tag{8E}\\
& \quad+\left(-8 a_{2}+4 a_{3}-7 b_{2}+5 b_{3}\right) v_{1}+\left(2 a_{2}-4 a_{3}+b_{2}-2 b_{3}\right) v_{2}^{2} \\
& \quad+\left(-5 a_{2}+a_{3}-4 b_{2}+2 b_{3}\right) v_{2}-3 a_{1}+2 a_{2}-a_{3}-3 b_{1}+4 b_{2}-2 b_{3}=0
\end{align*}
\]

Setting each coefficients in (8E) to zero gives the following equations to solve
\[
\begin{array}{r}
-8 a_{2}+4 a_{3}-7 b_{2}+5 b_{3}=0 \\
-5 a_{2}+a_{3}-4 b_{2}+2 b_{3}=0 \\
2 a_{2}-4 a_{3}+b_{2}-2 b_{3}=0 \\
4 a_{2}-8 a_{3}+2 b_{2}-4 b_{3}=0 \\
-3 a_{1}+2 a_{2}-a_{3}-3 b_{1}+4 b_{2}-2 b_{3}=0
\end{array}
\]

Solving the above equations for the unknowns gives
\[
\begin{aligned}
a_{1} & =-b_{1}-a_{2} \\
a_{2} & =a_{2} \\
a_{3} & =a_{2} \\
b_{1} & =b_{1} \\
b_{2} & =-2 a_{2} \\
b_{3} & =-2 a_{2}
\end{aligned}
\]

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives
\[
\begin{aligned}
& \xi=-1 \\
& \eta=1
\end{aligned}
\]

Shifting is now applied to make \(\xi=0\) in order to simplify the rest of the computation
\[
\begin{aligned}
\eta & =\eta-\omega(x, y) \xi \\
& =1-\left(-\frac{2 x+2 y-1}{x+y-2}\right)(-1) \\
& =\frac{-x-y-1}{x+y-2} \\
\xi & =0
\end{aligned}
\]

The next step is to determine the canonical coordinates \(R, S\). The canonical coordinates map \((x, y) \rightarrow(R, S)\) where \((R, S)\) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is
\[
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
\]

The above comes from the requirements that \(\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1\). Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable \(R\) in the canonical coordinates, where \(S(R)\). Since \(\xi=0\) then in this special case
\[
R=x
\]
\(S\) is found from
\[
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{-x-y-1}{x+y-2}} d y
\end{aligned}
\]

Which results in
\[
S=-y+3 \ln (x+y+1)
\]

Now that \(R, S\) are found, we need to setup the ode in these coordinates. This is done by evaluating
\[
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
\]

Where in the above \(R_{x}, R_{y}, S_{x}, S_{y}\) are all partial derivatives and \(\omega(x, y)\) is the right hand side of the original ode given by
\[
\omega(x, y)=-\frac{2 x+2 y-1}{x+y-2}
\]

Evaluating all the partial derivatives gives
\[
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =\frac{3}{x+y+1} \\
S_{y} & =-1+\frac{3}{x+y+1}
\end{aligned}
\]

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.
\[
\begin{equation*}
\frac{d S}{d R}=2 \tag{2~A}
\end{equation*}
\]

We now need to express the RHS as function of \(R\) only. This is done by solving for \(x, y\) in terms of \(R, S\) from the result obtained earlier and simplifying. This gives
\[
\frac{d S}{d R}=2
\]

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates \(R, S\). Integrating the above gives
\[
\begin{equation*}
S(R)=2 R+c_{1} \tag{4}
\end{equation*}
\]

To complete the solution, we just need to transform (4) back to \(x, y\) coordinates. This results in
\[
-y+3 \ln (x+y+1)=2 x+c_{1}
\]

Which simplifies to
\[
-y+3 \ln (x+y+1)=2 x+c_{1}
\]

Which gives
\[
y=-3 \operatorname{LambertW}\left(-\frac{\mathrm{e}^{\frac{x}{3}+\frac{c_{1}}{3}-\frac{1}{3}}}{3}\right)-x-1
\]

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.
\begin{tabular}{|c|c|c|}
\hline Original ode in \(x, y\) coordinates & Canonical coordinates transformation & ODE in canonical coordinates
\[
(R, S)
\] \\
\hline \(\frac{d y}{d x}=-\frac{2 x+2 y-1}{x+y-2}\) & & \(\frac{d S}{d R}=2\) \\
\hline  & &  \\
\hline  & &  \\
\hline 4, dy & &  \\
\hline  & &  \\
\hline  & &  \\
\hline  & &  \\
\hline  & \(R=x\) &  \\
\hline  & \(S=-y+3 \ln (x+\) & - \({ }^{\text {P4 }}\) \\
\hline  & &  \\
\hline  & &  \\
\hline \(t^{\text {a }}\) & &  \\
\hline  & &  \\
\hline  & &  \\
\hline
\end{tabular}

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
y=-3 \text { LambertW }\left(-\frac{\mathrm{e}^{\frac{x}{3}+\frac{c_{1}}{3}-\frac{1}{3}}}{3}\right)-x-1 \tag{1}
\end{equation*}
\]


Figure 210: Slope field plot

Verification of solutions
\[
y=-3 \operatorname{LambertW}\left(-\frac{\mathrm{e}^{\frac{x}{3}+\frac{c_{1}}{3}-\frac{1}{3}}}{3}\right)-x-1
\]

Verified OK.

Maple trace
```

`Methods for first order ODEs: --- Trying classification methods --- trying a quadrature trying 1st order linear trying Bernoulli trying separable trying inverse linear trying homogeneous types: trying homogeneous C 1st order, trying the canonical coordinates of the invariance group <- 1st order, canonical coordinates successful <- homogeneous successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.031 ( sec ). Leaf size: 21
```

dsolve((2*x+2*y(x)-1)+(x+y(x)-2)*diff (y(x),x)=0,y(x), singsol=all)

```
\[
y(x)=-x-3 \text { LambertW }\left(-\frac{c_{1} \mathrm{e}^{\frac{x}{3}-\frac{1}{3}}}{3}\right)-1
\]

Solution by Mathematica
Time used: 5.545 (sec). Leaf size: 35
```

DSolve[(2*x+2*y[x]-1)+(x+y[x]-2)*y'[x]==0,y[x],x,IncludeSingularSolutions -> True]

```
\[
\begin{aligned}
& y(x) \rightarrow-3 W\left(-e^{\frac{x}{3}-1+c_{1}}\right)-x-1 \\
& y(x) \rightarrow-x-1
\end{aligned}
\]

\subsection*{1.131 problem 190}
1.131.1 Solving as separable ode . . . . . . . . . . . . . . . . . . . . . . 1354
1.131.2 Solving as first order ode lie symmetry lookup ode . . . . . . . 1356
1.131.3 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 1360
1.131.4 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1364

Internal problem ID [12548]
Internal file name [OUTPUT/11200_Wednesday_October_18_2023_03_47_08_AM_1173136/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 190.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first__order__ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_separable]
\[
3 \mathrm{e}^{x} \tan (y)+\left(1-\mathrm{e}^{x}\right) \sec (y)^{2} y^{\prime}=0
\]

\subsection*{1.131.1 Solving as separable ode}

In canonical form the ODE is
\[
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =f(x) g(y) \\
& =\frac{3 \mathrm{e}^{x} \sin (2 y)}{2\left(-1+\mathrm{e}^{x}\right)}
\end{aligned}
\]

Where \(f(x)=\frac{3 \mathrm{e}^{x}}{-1+\mathrm{e}^{x}}\) and \(g(y)=\frac{\sin (2 y)}{2}\). Integrating both sides gives
\[
\begin{gathered}
\frac{1}{\frac{\sin (2 y)}{2}} d y=\frac{3 \mathrm{e}^{x}}{-1+\mathrm{e}^{x}} d x \\
\int \frac{1}{\frac{\sin (2 y)}{2}} d y=\int \frac{3 \mathrm{e}^{x}}{-1+\mathrm{e}^{x}} d x
\end{gathered}
\]
\[
\ln (\csc (2 y)-\cot (2 y))=3 \ln \left(-1+\mathrm{e}^{x}\right)+c_{1}
\]

Raising both side to exponential gives
\[
\csc (2 y)-\cot (2 y)=\mathrm{e}^{3 \ln \left(-1+\mathrm{e}^{x}\right)+c_{1}}
\]

Which simplifies to
\[
\csc (2 y)-\cot (2 y)=c_{2}\left(-1+\mathrm{e}^{x}\right)^{3}
\]

\section*{Summary}

The solution(s) found are the following

\section*{\(y\)}
\(=\frac{\arctan \left(\frac{2 \mathrm{e}^{c_{1}} c_{2}\left(\mathrm{e}^{3 x}-3 \mathrm{e}^{2 x}+3 \mathrm{e}^{x}-1\right)}{\mathrm{e}^{6 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-20 \mathrm{e}^{3 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{2 x} \mathrm{e}^{2 c_{1} c_{1}^{2}-6 \mathrm{e}^{x}} \mathrm{e}^{2 c_{1} c_{1}^{2}+c_{2}^{2} \mathrm{e}^{2 c_{1}}+1}},-\frac{\mathrm{e}^{6 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c} c_{1} c}{\mathrm{e}^{6 x} \mathrm{e}^{2 c} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c 1} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c_{1} c}}\right.}{2}\)


Figure 211: Slope field plot

\section*{Verification of solutions}
\(y\)
\(=\frac{\arctan \left(\frac{2 \mathrm{e}^{c_{1}} c_{2}\left(\mathrm{e}^{3 x}-3 \mathrm{e}^{2 x}+3 \mathrm{e}^{x}-1\right)}{\mathrm{e}^{6 x} \mathrm{x}^{2 c_{1}} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-20 \mathrm{e}^{3 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{2 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-6 \mathrm{e}^{x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+c_{2}^{2} \mathrm{e}^{2 c_{1}}+1},-\frac{\mathrm{e}^{6 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c_{1}} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c_{1}}}{\mathrm{e}^{6 x} \mathrm{x}^{2 c_{1}} c_{2}^{2}-6 \mathrm{e}^{5 x} \mathrm{e}^{2 c c_{1}} c_{2}^{2}+15 \mathrm{e}^{4 x} \mathrm{e}^{2 c_{1} c}}\right.}{2}\)
Verified OK.

\subsection*{1.131.2 Solving as first order ode lie symmetry lookup ode}

Writing the ode as
\[
\begin{aligned}
y^{\prime} & =\frac{3 \mathrm{e}^{x} \tan (y)}{\left(-1+\mathrm{e}^{x}\right) \sec (y)^{2}} \\
y^{\prime} & =\omega(x, y)
\end{aligned}
\]

The condition of Lie symmetry is the linearized PDE given by
\[
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
\]

The type of this ode is known. It is of type separable. Therefore we do not need to solve the \(\operatorname{PDE}\) (A), and can just use the lookup table shown below to find \(\xi, \eta\)

Table 203: Lie symmetry infinitesimal lookup table for known first order ODE's
\begin{tabular}{|c|c|c|c|}
\hline ODE class & Form & \(\xi\) & \(\eta\) \\
\hline linear ode & \(y^{\prime}=f(x) y(x)+g(x)\) & 0 & \(e^{\int f d x}\) \\
\hline separable ode & \(y^{\prime}=f(x) g(y)\) & \(\frac{1}{f}\) & 0 \\
\hline quadrature ode & \(y^{\prime}=f(x)\) & 0 & 1 \\
\hline quadrature ode & \(y^{\prime}=g(y)\) & 1 & 0 \\
\hline homogeneous ODEs of Class A & \(y^{\prime}=f\left(\frac{y}{x}\right)\) & \(x\) & \(y\) \\
\hline homogeneous ODEs of Class C & \(y^{\prime}=(a+b x+c y)^{\frac{n}{m}}\) & 1 & \[
-\frac{b}{c}
\] \\
\hline homogeneous class D & \(y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)\) & \(x^{2}\) & \(x y\) \\
\hline First order special form ID 1 & \(y^{\prime}=g(x) e^{h(x)+b y}+f(x)\) & \[
\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}
\] & \[
\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}
\] \\
\hline polynomial type ode & \(y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}\) & \[
\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}
\] & \[
\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}
\] \\
\hline Bernoulli ode & \(y^{\prime}=f(x) y+g(x) y^{n}\) & 0 & \(e^{-\int(n-1) f(x) d x} y^{n}\) \\
\hline Reduced Riccati & \(y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}\) & 0 & \(e^{-\int f_{1} d x}\) \\
\hline
\end{tabular}

The above table shows that
\[
\begin{align*}
& \xi(x, y)=\frac{\mathrm{e}^{-x}\left(-1+\mathrm{e}^{x}\right)}{3} \\
& \eta(x, y)=0 \tag{A1}
\end{align*}
\]

The next step is to determine the canonical coordinates \(R, S\). The canonical coordinates \(\operatorname{map}(x, y) \rightarrow(R, S)\) where \((R, S)\) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is
\[
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
\]

The above comes from the requirements that \(\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1\). Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable \(R\) in the canonical coordinates, where \(S(R)\). Since \(\eta=0\) then in this special case
\[
R=y
\]
\(S\) is found from
\[
\begin{aligned}
S & =\int \frac{1}{\xi} d x \\
& =\int \frac{1}{\frac{\mathrm{e}^{-x}\left(-1+\mathrm{e}^{x}\right)}{3}} d x
\end{aligned}
\]

Which results in
\[
S=3 \ln \left(-1+\mathrm{e}^{x}\right)
\]

Now that \(R, S\) are found, we need to setup the ode in these coordinates. This is done by evaluating
\[
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
\]

Where in the above \(R_{x}, R_{y}, S_{x}, S_{y}\) are all partial derivatives and \(\omega(x, y)\) is the right hand side of the original ode given by
\[
\omega(x, y)=\frac{3 \mathrm{e}^{x} \tan (y)}{\left(-1+\mathrm{e}^{x}\right) \sec (y)^{2}}
\]

Evaluating all the partial derivatives gives
\[
\begin{aligned}
R_{x} & =0 \\
R_{y} & =1 \\
S_{x} & =\frac{3 \mathrm{e}^{x}}{-1+\mathrm{e}^{x}} \\
S_{y} & =0
\end{aligned}
\]

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.
\[
\begin{equation*}
\frac{d S}{d R}=\sec (y) \csc (y) \tag{2~A}
\end{equation*}
\]

We now need to express the RHS as function of \(R\) only. This is done by solving for \(x, y\) in terms of \(R, S\) from the result obtained earlier and simplifying. This gives
\[
\frac{d S}{d R}=\sec (R) \csc (R)
\]

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates \(R, S\). Integrating the above gives
\[
\begin{equation*}
S(R)=\ln (\tan (R))+c_{1} \tag{4}
\end{equation*}
\]

To complete the solution, we just need to transform (4) back to \(x, y\) coordinates. This results in
\[
3 \ln \left(-1+\mathrm{e}^{x}\right)=\ln (\tan (y))+c_{1}
\]

Which simplifies to
\[
3 \ln \left(-1+\mathrm{e}^{x}\right)=\ln (\tan (y))+c_{1}
\]

Which gives
\[
y=\arctan \left(\mathrm{e}^{-c_{1}}\left(-1+\mathrm{e}^{x}\right)^{3}\right)
\]

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.
\begin{tabular}{|c|c|c|}
\hline Original ode in \(x, y\) coordinates & Canonical coordinates transformation & ODE in canonical coordinates
\[
(R, S)
\] \\
\hline \(\frac{d y}{d x}=\frac{3 \mathrm{e}^{x} \tan (y)}{\left(-1+\mathrm{e}^{x}\right) \sec (y)^{2}}\) & & \(\frac{d S}{d R}=\sec (R) \csc (R)\) \\
\hline  & &  \\
\hline \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-\infty]{\rightarrow \rightarrow \rightarrow+\infty}\) & & 星: \\
\hline  & &  \\
\hline \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \infty]{ }\) & &  \\
\hline \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\infty]{ }\) & &  \\
\hline \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+\infty]{ }\) & \(R=y\) & +iditativatidiapti \\
\hline  & \[
S=3 \ln \left(-1+\mathrm{e}^{x}\right)
\] &  \\
\hline  & &  \\
\hline  & &  \\
\hline  & &  \\
\hline  & &  \\
\hline \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-\infty]{ }\) & &  \\
\hline
\end{tabular}

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\arctan \left(\mathrm{e}^{-c_{1}}\left(-1+\mathrm{e}^{x}\right)^{3}\right) \tag{1}
\end{equation*}
\]


Figure 212: Slope field plot

Verification of solutions
\[
y=\arctan \left(\mathrm{e}^{-c_{1}}\left(-1+\mathrm{e}^{x}\right)^{3}\right)
\]

Verified OK.

\subsection*{1.131.3 Solving as exact ode}

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form
\[
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
\]

We assume there exists a function \(\phi(x, y)=c\) where \(c\) is constant, that satisfies the ode. Taking derivative of \(\phi\) w.r.t. \(x\) gives
\[
\frac{d}{d x} \phi(x, y)=0
\]

Hence
\[
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
\]

Comparing ( \(\mathrm{A}, \mathrm{B}\) ) shows that
\[
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
\]

But since \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) then for the above to be valid, we require that
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

If the above condition is satisfied, then the original ode is called exact. We still need to determine \(\phi(x, y)\) but at least we know now that we can do that since the condition \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is
\[
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
\]

Therefore
\[
\begin{align*}
\left(\frac{\sec (y)^{2}}{3 \tan (y)}\right) \mathrm{d} y & =\left(\frac{\mathrm{e}^{x}}{-1+\mathrm{e}^{x}}\right) \mathrm{d} x \\
\left(-\frac{\mathrm{e}^{x}}{-1+\mathrm{e}^{x}}\right) \mathrm{d} x+\left(\frac{\sec (y)^{2}}{3 \tan (y)}\right) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
\]

Comparing (1A) and (2A) shows that
\[
\begin{aligned}
& M(x, y)=-\frac{\mathrm{e}^{x}}{-1+\mathrm{e}^{x}} \\
& N(x, y)=\frac{\sec (y)^{2}}{3 \tan (y)}
\end{aligned}
\]

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

Using result found above gives
\[
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-\frac{\mathrm{e}^{x}}{-1+\mathrm{e}^{x}}\right) \\
& =0
\end{aligned}
\]

And
\[
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}\left(\frac{\sec (y)^{2}}{3 \tan (y)}\right) \\
& =0
\end{aligned}
\]

Since \(\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}\), then the ODE is exact The following equations are now set up to solve for the function \(\phi(x, y)\)
\[
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
\]

Integrating (1) w.r.t. \(x\) gives
\[
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-\frac{\mathrm{e}^{x}}{-1+\mathrm{e}^{x}} \mathrm{~d} x \\
\phi & =-\ln \left(-1+\mathrm{e}^{x}\right)+f(y) \tag{3}
\end{align*}
\]

Where \(f(y)\) is used for the constant of integration since \(\phi\) is a function of both \(x\) and \(y\). Taking derivative of equation (3) w.r.t \(y\) gives
\[
\begin{equation*}
\frac{\partial \phi}{\partial y}=0+f^{\prime}(y) \tag{4}
\end{equation*}
\]

But equation (2) says that \(\frac{\partial \phi}{\partial y}=\frac{\sec (y)^{2}}{3 \tan (y)}\). Therefore equation (4) becomes
\[
\begin{equation*}
\frac{\sec (y)^{2}}{3 \tan (y)}=0+f^{\prime}(y) \tag{5}
\end{equation*}
\]

Solving equation (5) for \(f^{\prime}(y)\) gives
\[
\begin{aligned}
f^{\prime}(y) & =\frac{\sec (y)^{2}}{3 \tan (y)} \\
& =\frac{\sec (y) \csc (y)}{3}
\end{aligned}
\]

Integrating the above w.r.t \(y\) results in
\[
\begin{aligned}
\int f^{\prime}(y) \mathrm{d} y & =\int\left(\frac{\sec (y) \csc (y)}{3}\right) \mathrm{d} y \\
f(y) & =\frac{\ln (\tan (y))}{3}+c_{1}
\end{aligned}
\]

Where \(c_{1}\) is constant of integration. Substituting result found above for \(f(y)\) into equation (3) gives \(\phi\)
\[
\phi=-\ln \left(-1+\mathrm{e}^{x}\right)+\frac{\ln (\tan (y))}{3}+c_{1}
\]

But since \(\phi\) itself is a constant function, then let \(\phi=c_{2}\) where \(c_{2}\) is new constant and combining \(c_{1}\) and \(c_{2}\) constants into new constant \(c_{1}\) gives the solution as
\[
c_{1}=-\ln \left(-1+\mathrm{e}^{x}\right)+\frac{\ln (\tan (y))}{3}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
-\ln \left(-1+\mathrm{e}^{x}\right)+\frac{\ln (\tan (y))}{3}=c_{1} \tag{1}
\end{equation*}
\]


Figure 213: Slope field plot

Verification of solutions
\[
-\ln \left(-1+\mathrm{e}^{x}\right)+\frac{\ln (\tan (y))}{3}=c_{1}
\]

Verified OK.

\subsection*{1.131.4 Maple step by step solution}

Let's solve
\(3 \mathrm{e}^{x} \tan (y)+\left(1-\mathrm{e}^{x}\right) \sec (y)^{2} y^{\prime}=0\)
- Highest derivative means the order of the ODE is 1 \(y^{\prime}\)
- \(\quad\) Separate variables
\(\frac{y^{\prime} \sec (y)^{2}}{\tan (y)}=-\frac{3 \mathrm{e}^{x}}{1-\mathrm{e}^{x}}\)
- Integrate both sides with respect to \(x\)
\[
\int \frac{y^{\prime} \sec (y)^{2}}{\tan (y)} d x=\int-\frac{3 \mathrm{e}^{x}}{1-\mathrm{e}^{x}} d x+c_{1}
\]
- Evaluate integral
\[
\ln (\tan (y))=3 \ln \left(1-\mathrm{e}^{x}\right)+c_{1}
\]
- \(\quad\) Solve for \(y\)
\[
y=-\arctan \left(\mathrm{e}^{c_{1}}\left(-1+\mathrm{e}^{x}\right)^{3}\right)
\]

Maple trace
```

`Methods for first order ODEs: --- Trying classification methods --- trying a quadrature trying 1st order linear trying Bernoulli trying separable <- separable successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.032 (sec). Leaf size: 205
```

dsolve(3*exp(x)*\operatorname{tan}(y(x))+(1-exp(x))*sec(y(x))^2*diff (y(x),x)=0,y(x), singsol=all)

```
\(y(x)\)
\(=\frac{\arctan \left(-\frac{2 c_{1}\left(\mathrm{e}^{3 x}-3 \mathrm{e}^{2 x}+3 \mathrm{e}^{x}-1\right)}{-c_{1}^{2} \mathrm{e}^{6 x}+6 c_{1}^{2} 5^{5 x}-15 c_{1}^{2} \mathrm{e}^{4 x}+20 c_{1}^{2} \mathrm{e}^{3 x}-15 c_{1}^{2} \mathrm{e}^{2 x}+6 c_{1}^{2} \mathrm{e}^{x}-c_{1}^{2}-1}, \frac{c_{1}^{2} e^{6 x}-6 c_{1}^{2} 5^{5 x}+15 c_{1}^{2} \mathrm{e}^{4 x}-20 c_{1}^{2} \mathrm{e}^{3 x}+15 c_{1}^{2} \mathrm{e}^{2 x}-6 c_{1}^{2} \mathrm{e}^{x}+c_{1}^{2}-1}{-c_{1}^{2} \mathrm{e}^{6 x}+6 c_{1}^{2} \mathrm{e}^{5 x}-15 c_{1}^{2} \mathrm{e}^{4 x}+20 c_{1}^{2} \mathrm{e}^{3 x}-15 c_{1}^{2} \mathrm{e}^{2 x}+6 c_{1}^{2} \mathrm{e}^{x}-c_{1}^{2}-1}\right)}{2}\)
\(\checkmark\) Solution by Mathematica
Time used: 1.847 (sec). Leaf size: 74
DSolve \(\left[3 * \operatorname{Exp}[x] * \operatorname{Tan}[y[x]]+(1-\operatorname{Exp}[x]) * \operatorname{Sec}[y[x]]^{\wedge} 2 * y^{\prime}[x]==0, y[x], x\right.\), IncludeSingularSolutions \(\rightarrow\)
\[
\begin{aligned}
& y(x) \rightarrow-\frac{1}{2} \arccos \left(-\tanh \left(3 \log \left(e^{x}-1\right)+2 c_{1}\right)\right) \\
& y(x) \rightarrow \frac{1}{2} \arccos \left(-\tanh \left(3 \log \left(e^{x}-1\right)+2 c_{1}\right)\right) \\
& y(x) \rightarrow 0 \\
& y(x) \rightarrow-\frac{\pi}{2} \\
& y(x) \rightarrow \frac{\pi}{2}
\end{aligned}
\]

\subsection*{1.132 problem 191}
1.132.1 Solution using Matrix exponential method . . . . . . . . . . . . 1366
1.132.2 Solution using explicit Eigenvalue and Eigenvector method . . . 1367
1.132.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1372

Internal problem ID [12549]
Internal file name [OUTPUT/11201_Wednesday_October_18_2023_03_47_10_AM_83868436/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 191.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "system of linear ODEs"
Solve
\[
\begin{aligned}
x^{\prime}(t) & =2 x(t)-3 y(t) \\
y^{\prime}(t) & =5 x(t)+6 y(t)
\end{aligned}
\]

\subsection*{1.132.1 Solution using Matrix exponential method}

In this method, we will assume we have found the matrix exponential \(e^{A t}\) allready. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)
\]

Or
\[
\left[\begin{array}{c}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{cc}
2 & -3 \\
5 & 6
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]

For the above matrix \(A\), the matrix exponential can be found to be
\[
\begin{aligned}
e^{A t} & =\left[\begin{array}{cc}
\mathrm{e}^{4 t} \cos (\sqrt{11} t)-\frac{2 \mathrm{e}^{4 t} \sin (\sqrt{11} t) \sqrt{11}}{11} & -\frac{3 \mathrm{e}^{4 t} \sin (\sqrt{11} t) \sqrt{11}}{11} \\
\frac{5 \mathrm{e}^{4 t} \sin (\sqrt{11} t) \sqrt{11}}{11} & \mathrm{e}^{4 t} \cos (\sqrt{11} t)+\frac{2 \mathrm{e}^{4 t} \sin (\sqrt{11} t) \sqrt{11}}{11}
\end{array}\right] \\
& =\left[\begin{array}{cc}
\mathrm{e}^{4 t}\left(\cos (\sqrt{11} t)-\frac{2 \sin (\sqrt{11} t) \sqrt{11}}{11}\right) & -\frac{3 \mathrm{e}^{4 t} \sin (\sqrt{11} t) \sqrt{11}}{11} \\
\frac{5 \mathrm{e}^{4 t} \sin (\sqrt{11} t) \sqrt{11}}{11} & \frac{\mathrm{e}^{4 t}(2 \sin (\sqrt{11} t) \sqrt{11}+11 \cos (\sqrt{11} t))}{11}
\end{array}\right]
\end{aligned}
\]

Therefore the homogeneous solution is
\[
\begin{aligned}
\vec{x}_{h}(t) & =e^{A t} \vec{c} \\
& =\left[\begin{array}{cc}
\mathrm{e}^{4 t}\left(\cos (\sqrt{11} t)-\frac{2 \sin (\sqrt{11} t) \sqrt{11}}{11}\right) & -\frac{3 \mathrm{e}^{4 t} \sin (\sqrt{11} t) \sqrt{11}}{11} \\
\frac{5 \mathrm{e}^{4 t} \sin (\sqrt{11} t) \sqrt{11}}{11} & \frac{\mathrm{e}^{4 t}(2 \sin (\sqrt{11} t) \sqrt{11}+11 \cos (\sqrt{11} t))}{11}
\end{array}\right]\left[\begin{array}{l}
c_{1} \\
c_{2}
\end{array}\right] \\
& =\left[\begin{array}{c}
\mathrm{e}^{4 t\left(\cos (\sqrt{11} t)-\frac{2 \sin (\sqrt{11} t) \sqrt{11}}{11}\right) c_{1}-\frac{3 \mathrm{e}^{4 t} \sin (\sqrt{11} t) \sqrt{11} c_{2}}{11}} \\
\frac{5 \mathrm{e}^{4 t} \sin (\sqrt{11} t) \sqrt{11} c_{1}}{11}+\frac{\mathrm{e}^{4 t}(2 \sin (\sqrt{11} t) \sqrt{11}+11 \cos (\sqrt{11} t)) c_{2}}{11}
\end{array}\right] \\
& =\left[\begin{array}{c}
-\frac{2\left(\sqrt{11}\left(c_{1}+\frac{3 c_{2}}{2}\right) \sin (\sqrt{11} t)-\frac{11 c_{1} \cos (\sqrt{11} t)}{2}\right) \mathrm{e}^{4 t}}{11} \\
\frac{5 \mathrm{e}^{4 t}\left(\sqrt{11}\left(c_{1}+\frac{2 c_{2}}{5}\right) \sin (\sqrt{11} t)+\frac{11 \cos (\sqrt{11} t) c_{2}}{5}\right)}{11}
\end{array}\right]
\end{aligned}
\]

Since no forcing function is given, then the final solution is \(\vec{x}_{h}(t)\) above.

\subsection*{1.132.2 Solution using explicit Eigenvalue and Eigenvector method}

This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)
\]

Or
\[
\left[\begin{array}{c}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{cc}
2 & -3 \\
5 & 6
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]

The first step is find the homogeneous solution. We start by finding the eigenvalues of \(A\). This is done by solving the following equation for the eigenvalues \(\lambda\)
\[
\operatorname{det}(A-\lambda I)=0
\]

Expanding gives
\[
\operatorname{det}\left(\left[\begin{array}{cc}
2 & -3 \\
5 & 6
\end{array}\right]-\lambda\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)=0
\]

Therefore
\[
\operatorname{det}\left(\left[\begin{array}{cc}
2-\lambda & -3 \\
5 & 6-\lambda
\end{array}\right]\right)=0
\]

Which gives the characteristic equation
\[
\lambda^{2}-8 \lambda+27=0
\]

The roots of the above are the eigenvalues.
\[
\begin{aligned}
& \lambda_{1}=4+i \sqrt{11} \\
& \lambda_{2}=4-i \sqrt{11}
\end{aligned}
\]

This table summarises the above result
\begin{tabular}{|l|l|l|}
\hline eigenvalue & algebraic multiplicity & type of eigenvalue \\
\hline \(4+i \sqrt{11}\) & 1 & complex eigenvalue \\
\hline \(4-i \sqrt{11}\) & 1 & complex eigenvalue \\
\hline
\end{tabular}

Now the eigenvector for each eigenvalue are found.
Considering the eigenvalue \(\lambda_{1}=4-i \sqrt{11}\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
\left(\left[\begin{array}{cc}
2 & -3 \\
5 & 6
\end{array}\right]-(4-i \sqrt{11})\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right] & =\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
{\left[\begin{array}{cc}
-2+i \sqrt{11} & -3 \\
5 & 2+i \sqrt{11}
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right] } & =\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{cc|c}
-2+i \sqrt{11} & -3 & 0 \\
5 & 2+i \sqrt{11} & 0
\end{array}\right]} \\
R_{2}=R_{2}-\frac{5 R_{1}}{-2+i \sqrt{11}} \Longrightarrow\left[\begin{array}{cc|c}
-2+i \sqrt{11} & -3 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{cc}
-2+i \sqrt{11} & -3 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=\frac{3 t}{-2+i \sqrt{11}}\right\}\)
Hence the solution is
\[
\left[\begin{array}{c}
\frac{3 t}{-2+\mathrm{I} \sqrt{11}} \\
t
\end{array}\right]=\left[\begin{array}{c}
\frac{3 t}{-2+i \sqrt{11}} \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
\frac{3 t}{-2+\mathrm{I} \sqrt{11}} \\
t
\end{array}\right]=t\left[\begin{array}{c}
\frac{3}{-2+i \sqrt{11}} \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{c}
\frac{3 t}{-2+\mathrm{I} \sqrt{11}} \\
t
\end{array}\right]=\left[\begin{array}{c}
\frac{3}{-2+i \sqrt{11}} \\
1
\end{array}\right]
\]

Which is normalized to
\[
\left[\begin{array}{c}
\frac{3 t}{-2+\mathrm{I} \sqrt{11}} \\
t
\end{array}\right]=\left[\begin{array}{c}
\frac{3}{-2+i \sqrt{11}} \\
1
\end{array}\right]
\]

Considering the eigenvalue \(\lambda_{2}=4+i \sqrt{11}\)

We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
&\left(\left[\begin{array}{cc}
2 & -3 \\
5 & 6
\end{array}\right]-(4+i \sqrt{11})\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
& {\left[\begin{array}{cc}
-2-i \sqrt{11} & -3 \\
5 & 2-i \sqrt{11}
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] }
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{cc|c}
-2-i \sqrt{11} & -3 & 0 \\
5 & 2-i \sqrt{11} & 0
\end{array}\right]} \\
R_{2}=R_{2}-\frac{5 R_{1}}{-2-i \sqrt{11}} \Longrightarrow\left[\begin{array}{cc|c}
-2-i \sqrt{11} & -3 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{cc}
-2-i \sqrt{11} & -3 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=-\frac{3 t}{2+i \sqrt{11}}\right\}\)
Hence the solution is
\[
\left[\begin{array}{c}
-\frac{3 t}{2+\mathrm{I} \sqrt{11}} \\
t
\end{array}\right]=\left[\begin{array}{c}
-\frac{3 t}{2+i \sqrt{11}} \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
-\frac{3 t}{2+\mathrm{I} \sqrt{11}} \\
t
\end{array}\right]=t\left[\begin{array}{c}
-\frac{3}{2+i \sqrt{11}} \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{c}
-\frac{3 t}{2+\mathrm{I} \sqrt{11}} \\
t
\end{array}\right]=\left[\begin{array}{c}
-\frac{3}{2+i \sqrt{11}} \\
1
\end{array}\right]
\]

Which is normalized to
\[
\left[\begin{array}{c}
-\frac{3 t}{2+\mathrm{I} \sqrt{11}} \\
t
\end{array}\right]=\left[\begin{array}{c}
-\frac{3}{2+i \sqrt{11}} \\
1
\end{array}\right]
\]

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity \(m\), and its geometric multiplicity \(k\) and the eigenvectors associated with the eigenvalue. If \(m>k\) then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity \(k\) ) does not equal the algebraic multiplicity \(m\), and we need to determine an additional \(m-k\) generalized eigenvectors for this eigenvalue.
\begin{tabular}{|c|c|c|c|c|}
\hline \multirow{2}{*}{ eigenvalue } & \multicolumn{2}{|c|}{ multiplicity } & \multirow{2}{*}{} & \\
\cline { 2 - 3 } & algebraic \(m\) & geometric \(k\) & defective? & eigenvectors \\
\hline \(4+i \sqrt{11}\) & 1 & 1 & No & {\(\left[\begin{array}{c}-\frac{3}{2+i \sqrt{11}} \\
1\end{array}\right]\)} \\
\hline \(4-i \sqrt{11}\) & 1 & 1 & No & {\(\left[\begin{array}{c}-\frac{3}{2-i \sqrt{11}} \\
1\end{array}\right]\)} \\
\hline
\end{tabular}

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is
\[
\vec{x}_{h}(t)=c_{1} \vec{x}_{1}(t)+c_{2} \vec{x}_{2}(t)
\]

Which is written as
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=c_{1}\left[\begin{array}{c}
-\frac{3 \mathrm{e}^{(4+i \sqrt{11}) t}}{2+i \sqrt{11}} \\
\mathrm{e}^{(4+i \sqrt{11}) t}
\end{array}\right]+c_{2}\left[\begin{array}{c}
-\frac{3 \mathrm{e}^{(4-i \sqrt{11}) t}}{2-i \sqrt{11}} \\
\mathrm{e}^{(4-i \sqrt{11}) t}
\end{array}\right]
\]

Which becomes
\[
\left[\begin{array}{c}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
\frac{2 i\left(i-\frac{\sqrt{11}}{2}\right) c_{2} \mathrm{e}^{-(i \sqrt{11}-4) t}}{5}+\frac{2 i e^{(4+i \sqrt{11}) t} c_{1}\left(i+\frac{\sqrt{11}}{2}\right)}{5} \\
c_{1} \mathrm{e}^{(4+i \sqrt{11}) t}+c_{2} \mathrm{e}^{-(i \sqrt{11}-4) t}
\end{array}\right]
\]

The following is the phase plot of the system.


Figure 214: Phase plot

\subsection*{1.132.3 Maple step by step solution}

Let's solve
\(\left[x^{\prime}(t)=2 x(t)-3 y(t), y^{\prime}(t)=5 x(t)+6 y(t)\right]\)
- Define vector
\(\vec{x}(t)=\left[\begin{array}{l}x(t) \\ y(t)\end{array}\right]\)
- Convert system into a vector equation
\(\vec{x}^{\prime}(t)=\left[\begin{array}{cc}2 & -3 \\ 5 & 6\end{array}\right] \cdot \vec{x}(t)+\left[\begin{array}{l}0 \\ 0\end{array}\right]\)
- System to solve
\(\vec{x}^{\prime}(t)=\left[\begin{array}{cc}2 & -3 \\ 5 & 6\end{array}\right] \cdot \vec{x}(t)\)
- Define the coefficient matrix
\[
A=\left[\begin{array}{cc}
2 & -3 \\
5 & 6
\end{array}\right]
\]
- Rewrite the system as
\(\vec{x}^{\prime}(t)=A \cdot \vec{x}(t)\)
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- \(\quad\) Eigenpairs of \(A\)
\[
\left[\left[4-\mathrm{I} \sqrt{11},\left[\begin{array}{c}
-\frac{3}{2-\mathrm{I} \sqrt{11}} \\
1
\end{array}\right]\right],\left[4+\mathrm{I} \sqrt{11},\left[\begin{array}{c}
-\frac{3}{2+\mathrm{I} \sqrt{11}} \\
1
\end{array}\right]\right]\right]
\]
- Consider complex eigenpair, complex conjugate eigenvalue can be ignored
\[
\left[4-\mathrm{I} \sqrt{11},\left[\begin{array}{c}
-\frac{3}{2-\mathrm{I} \sqrt{11}} \\
1
\end{array}\right]\right]
\]
- \(\quad\) Solution from eigenpair
\(\mathrm{e}^{(4-\mathrm{I} \sqrt{11}) t} \cdot\left[\begin{array}{c}-\frac{3}{2-\mathrm{I} \sqrt{11}} \\ 1\end{array}\right]\)
- Use Euler identity to write solution in terms of sin and cos
\[
\mathrm{e}^{4 t} \cdot(\cos (\sqrt{11} t)-\mathrm{I} \sin (\sqrt{11} t)) \cdot\left[\begin{array}{c}
-\frac{3}{2-\mathrm{I} \sqrt{11}} \\
1
\end{array}\right]
\]
- Simplify expression
\[
\mathrm{e}^{4 t} \cdot\left[\begin{array}{c}
-\frac{3(\cos (\sqrt{11} t)-\mathrm{I} \sin (\sqrt{11} t))}{2-\mathrm{I} \sqrt{11}} \\
\cos (\sqrt{11} t)-\mathrm{I} \sin (\sqrt{11} t)
\end{array}\right]
\]
- Both real and imaginary parts are solutions to the homogeneous system
\[
\left[\vec{x}_{1}(t)=\mathrm{e}^{4 t} \cdot\left[-\frac{2 \cos (\sqrt{11} t)}{5}-\frac{\sin (\sqrt{11} t) \sqrt{11}}{5}\right], \vec{x}_{2}(t)=\mathrm{e}^{4 t} \cdot\left[\begin{array}{c}
-\frac{\cos (\sqrt{11} t) \sqrt{11}}{5}+\frac{2 \sin (\sqrt{11} t)}{5} \\
\cos (\sqrt{11} t)
\end{array}\right]\right]
\]
- General solution to the system of ODEs
\[
\vec{x}=c_{1} \vec{x}_{1}(t)+c_{2} \vec{x}_{2}(t)
\]
- Substitute solutions into the general solution
\[
\vec{x}=c_{1} \mathrm{e}^{4 t} \cdot\left[\begin{array}{c}
-\frac{2 \cos (\sqrt{11} t)}{5}-\frac{\sin (\sqrt{11} t) \sqrt{11}}{5} \\
\cos (\sqrt{11} t)
\end{array}\right]+c_{2} \mathrm{e}^{4 t} \cdot\left[\begin{array}{c}
-\frac{\cos (\sqrt{11} t) \sqrt{11}}{5}+\frac{2 \sin (\sqrt{11} t)}{5} \\
-\sin (\sqrt{11} t)
\end{array}\right]
\]
- Substitute in vector of dependent variables
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
-\frac{\mathrm{e}^{4 t}\left(\left(c_{2} \sqrt{11}+2 c_{1}\right) \cos (\sqrt{11} t)+\sin (\sqrt{11} t)\left(\sqrt{11} c_{1}-2 c_{2}\right)\right)}{5} \\
\mathrm{e}^{4 t}\left(\cos (\sqrt{11} t) c_{1}-\sin (\sqrt{11} t) c_{2}\right)
\end{array}\right]
\]
- \(\quad\) Solution to the system of ODEs
\[
\left\{x(t)=-\frac{\mathrm{e}^{4 t}\left(\left(c_{2} \sqrt{11}+2 c_{1}\right) \cos (\sqrt{11} t)+\sin (\sqrt{11} t)\left(\sqrt{11} c_{1}-2 c_{2}\right)\right)}{5}, y(t)=\mathrm{e}^{4 t}\left(\cos (\sqrt{11} t) c_{1}-\sin (\sqrt{11} t) c_{2}\right)\right.
\]
\(\checkmark\) Solution by Maple
Time used: 0.031 (sec). Leaf size: 78
```

dsolve([diff(x(t),t)=2*x(t)-3*y(t), diff(y(t),t)=5*x(t)+6*y(t)],singsol=all)

```
\(x(t)=\mathrm{e}^{4 t}\left(\sin (\sqrt{11} t) c_{1}+\cos (\sqrt{11} t) c_{2}\right)\)
\(y(t)=\frac{\mathrm{e}^{4 t}\left(\sin (\sqrt{11} t) \sqrt{11} c_{2}-\cos (\sqrt{11} t) \sqrt{11} c_{1}-2 \sin (\sqrt{11} t) c_{1}-2 \cos (\sqrt{11} t) c_{2}\right)}{3}\)
Solution by Mathematica
Time used: 0.033 (sec). Leaf size: 99
DSolve \(\left[\left\{x^{\prime}[t]==2 * x[t]-3 * y[t], y^{\prime}[t]==5 * x[t]+6 * y[t]\right\},\{x[t], y[t]\}, t\right.\), IncludeSingularSolutions \(\rightarrow\)
\[
\begin{aligned}
& x(t) \rightarrow c_{1} e^{4 t} \cos (\sqrt{11} t)-\frac{\left(2 c_{1}+3 c_{2}\right) e^{4 t} \sin (\sqrt{11} t)}{\sqrt{11}} \\
& y(t) \rightarrow c_{2} e^{4 t} \cos (\sqrt{11} t)+\frac{\left(5 c_{1}+2 c_{2}\right) e^{4 t} \sin (\sqrt{11} t)}{\sqrt{11}}
\end{aligned}
\]

\subsection*{1.133 problem 192}
1.133.1 Solution using Matrix exponential method . . . . . . . . . . . . 1375
1.133.2 Solution using explicit Eigenvalue and Eigenvector method . . . 1376
1.133.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1380

Internal problem ID [12550]
Internal file name [OUTPUT/11202_Wednesday_October_18_2023_03_47_10_AM_76593534/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 192.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "system of linear ODEs"
Solve
\[
\begin{aligned}
x^{\prime}(t) & =-4 x(t)-10 y(t) \\
y^{\prime}(t) & =x(t)-2 y(t)
\end{aligned}
\]

\subsection*{1.133.1 Solution using Matrix exponential method}

In this method, we will assume we have found the matrix exponential \(e^{A t}\) allready. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)
\]

Or
\[
\left[\begin{array}{l}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{cc}
-4 & -10 \\
1 & -2
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]

For the above matrix \(A\), the matrix exponential can be found to be
\[
\begin{aligned}
e^{A t} & =\left[\begin{array}{cc}
\mathrm{e}^{-3 t} \cos (3 t)-\frac{\mathrm{e}^{-3 t} \sin (3 t)}{3} & -\frac{10 \mathrm{e}^{-3 t} \sin (3 t)}{3} \\
\frac{\mathrm{e}^{-3 t} \sin (3 t)}{3} & \mathrm{e}^{-3 t} \cos (3 t)+\frac{\mathrm{e}^{-3 t} \sin (3 t)}{3}
\end{array}\right] \\
& =\left[\begin{array}{cc}
\frac{\mathrm{e}^{-3 t}(3 \cos (3 t)-\sin (3 t))}{3} & -\frac{10 \mathrm{e}^{-3 t} \sin (3 t)}{3} \\
\frac{\mathrm{e}^{-3 t} \sin (3 t)}{3} & \frac{\mathrm{e}^{-3 t}(3 \cos (3 t)+\sin (3 t))}{3}
\end{array}\right]
\end{aligned}
\]

Therefore the homogeneous solution is
\[
\begin{aligned}
\vec{x}_{h}(t) & =e^{A t} \vec{c} \\
& =\left[\begin{array}{cc}
\frac{\mathrm{e}^{-3 t}(3 \cos (3 t)-\sin (3 t))}{3} & -\frac{10 \mathrm{e}^{-3 t} \sin (3 t)}{3} \\
\frac{\mathrm{e}^{-3 t} \sin (3 t)}{3} & \frac{\mathrm{e}^{-3 t}(3 \cos (3 t)+\sin (3 t))}{3}
\end{array}\right]\left[\begin{array}{l}
c_{1} \\
c_{2}
\end{array}\right] \\
& =\left[\begin{array}{c}
\frac{\mathrm{e}^{-3 t}(3 \cos (3 t)-\sin (3 t)) c_{1}}{3}-\frac{10 \mathrm{e}^{-3 t} \sin (3 t) c_{2}}{3} \\
\frac{\mathrm{e}^{-3 t} \sin (3 t) c_{1}}{3}+\frac{\mathrm{e}^{-3 t}(3 \cos (3 t)+\sin (3 t)) c_{2}}{3}
\end{array}\right] \\
& =\left[\begin{array}{c}
\frac{\mathrm{e}^{-3 t}\left(3 c_{1} \cos (3 t)-c_{1} \sin (3 t)-10 \sin (3 t) c_{2}\right)}{3} \\
\frac{\mathrm{e}^{-3 t}\left(\left(c_{1}+c_{2}\right) \sin (3 t)+3 c_{2} \cos (3 t)\right)}{3}
\end{array}\right]
\end{aligned}
\]

Since no forcing function is given, then the final solution is \(\vec{x}_{h}(t)\) above.

\subsection*{1.133.2 Solution using explicit Eigenvalue and Eigenvector method}

This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)
\]

Or
\[
\left[\begin{array}{c}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{cc}
-4 & -10 \\
1 & -2
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]

The first step is find the homogeneous solution. We start by finding the eigenvalues of \(A\). This is done by solving the following equation for the eigenvalues \(\lambda\)
\[
\operatorname{det}(A-\lambda I)=0
\]

Expanding gives
\[
\operatorname{det}\left(\left[\begin{array}{cc}
-4 & -10 \\
1 & -2
\end{array}\right]-\lambda\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)=0
\]

Therefore
\[
\operatorname{det}\left(\left[\begin{array}{cc}
-4-\lambda & -10 \\
1 & -2-\lambda
\end{array}\right]\right)=0
\]

Which gives the characteristic equation
\[
\lambda^{2}+6 \lambda+18=0
\]

The roots of the above are the eigenvalues.
\[
\begin{gathered}
\lambda_{1}=-3+3 i \\
\lambda_{2}=-3-3 i
\end{gathered}
\]

This table summarises the above result
\begin{tabular}{|l|l|l|}
\hline eigenvalue & algebraic multiplicity & type of eigenvalue \\
\hline\(-3-3 i\) & 1 & complex eigenvalue \\
\hline\(-3+3 i\) & 1 & complex eigenvalue \\
\hline
\end{tabular}

Now the eigenvector for each eigenvalue are found.
Considering the eigenvalue \(\lambda_{1}=-3-3 i\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
&\left(\left[\begin{array}{cc}
-4 & -10 \\
1 & -2
\end{array}\right]-(-3-3 i)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
& {\left[\begin{array}{cc}
-1+3 i & -10 \\
1 & 1+3 i
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] }
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{cc|c}
-1+3 i & -10 & 0 \\
1 & 1+3 i & 0
\end{array}\right]} \\
R_{2}=R_{2}+\left(\frac{1}{10}+\frac{3 i}{10}\right) R_{1} \Longrightarrow\left[\begin{array}{cc|c}
-1+3 i & -10 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{cc}
-1+3 i & -10 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=(-1-3 i) t\right\}\)

Hence the solution is
\[
\left[\begin{array}{c}
(-1-3 \mathrm{I}) t \\
t
\end{array}\right]=\left[\begin{array}{c}
(-1-3 i) t \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
(-1-3 \mathrm{I}) t \\
t
\end{array}\right]=t\left[\begin{array}{c}
-1-3 i \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{c}
(-1-3 \mathrm{I}) t \\
t
\end{array}\right]=\left[\begin{array}{c}
-1-3 i \\
1
\end{array}\right]
\]

Considering the eigenvalue \(\lambda_{2}=-3+3 i\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
&\left(\left[\begin{array}{cc}
-4 & -10 \\
1 & -2
\end{array}\right]-(-3+3 i)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
& {\left[\begin{array}{cc}
-1-3 i & -10 \\
1 & 1-3 i
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] }
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{cc|c}
-1-3 i & -10 & 0 \\
1 & 1-3 i & 0
\end{array}\right]} \\
R_{2}=R_{2}+\left(\frac{1}{10}-\frac{3 i}{10}\right) R_{1} \Longrightarrow\left[\begin{array}{cc|c}
-1-3 i & -10 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{cc}
-1-3 i & -10 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=(-1+3 i) t\right\}\)

Hence the solution is
\[
\left[\begin{array}{c}
(-1+3 \mathrm{I}) t \\
t
\end{array}\right]=\left[\begin{array}{c}
(-1+3 i) t \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
(-1+3 \mathrm{I}) t \\
t
\end{array}\right]=t\left[\begin{array}{c}
-1+3 i \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{c}
(-1+3 \mathrm{I}) t \\
t
\end{array}\right]=\left[\begin{array}{c}
-1+3 i \\
1
\end{array}\right]
\]

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity \(m\), and its geometric multiplicity \(k\) and the eigenvectors associated with the eigenvalue. If \(m>k\) then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity \(k\) ) does not equal the algebraic multiplicity \(m\), and we need to determine an additional \(m-k\) generalized eigenvectors for this eigenvalue.
\begin{tabular}{|c|c|c|c|c|}
\hline \multirow{2}{*}{ eigenvalue } & \multicolumn{2}{|c|}{ multiplicity } & & \\
\cline { 2 - 3 } & algebraic \(m\) & geometric \(k\) & defective? & eigenvectors \\
\hline\(-3+3 i\) & 1 & 1 & No & {\(\left[\begin{array}{c}-1+3 i \\
1\end{array}\right]\)} \\
\hline\(-3-3 i\) & 1 & 1 & No & {\(\left[\begin{array}{c}-1-3 i \\
1\end{array}\right]\)} \\
\hline
\end{tabular}

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is
\[
\vec{x}_{h}(t)=c_{1} \vec{x}_{1}(t)+c_{2} \vec{x}_{2}(t)
\]

Which is written as
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=c_{1}\left[\begin{array}{c}
(-1+3 i) \mathrm{e}^{(-3+3 i) t} \\
\mathrm{e}^{(-3+3 i) t}
\end{array}\right]+c_{2}\left[\begin{array}{c}
(-1-3 i) \mathrm{e}^{(-3-3 i) t} \\
\mathrm{e}^{(-3-3 i) t}
\end{array}\right]
\]

Which becomes
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
(-1+3 i) c_{1} \mathrm{e}^{(-3+3 i) t}+(-1-3 i) c_{2} \mathrm{e}^{(-3-3 i) t} \\
c_{1} \mathrm{e}^{(-3+3 i) t}+c_{2} \mathrm{e}^{(-3-3 i) t}
\end{array}\right]
\]

The following is the phase plot of the system.


Figure 215: Phase plot

\subsection*{1.133.3 Maple step by step solution}

Let's solve
\[
\left[x^{\prime}(t)=-4 x(t)-10 y(t), y^{\prime}(t)=x(t)-2 y(t)\right]
\]
- Define vector
\(\vec{x}(t)=\left[\begin{array}{l}x(t) \\ y(t)\end{array}\right]\)
- Convert system into a vector equation
\(\vec{x}^{\prime}(t)=\left[\begin{array}{cc}-4 & -10 \\ 1 & -2\end{array}\right] \cdot \vec{x}(t)+\left[\begin{array}{l}0 \\ 0\end{array}\right]\)
- System to solve
\[
\vec{x}^{\prime}(t)=\left[\begin{array}{cc}
-4 & -10 \\
1 & -2
\end{array}\right] \cdot \vec{x}(t)
\]
- Define the coefficient matrix
\[
A=\left[\begin{array}{cc}
-4 & -10 \\
1 & -2
\end{array}\right]
\]
- Rewrite the system as
\[
\vec{x}^{\prime}(t)=A \cdot \vec{x}(t)
\]
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- \(\quad\) Eigenpairs of \(A\)
\[
\left[\left[\left[-3-3 \mathrm{I},\left[\begin{array}{c}
-1-3 \mathrm{I} \\
1
\end{array}\right]\right],\left[-3+3 \mathrm{I},\left[\begin{array}{c}
-1+3 \mathrm{I} \\
1
\end{array}\right]\right]\right]\right.
\]
- Consider complex eigenpair, complex conjugate eigenvalue can be ignored
\[
\left[-3-3 \mathrm{I},\left[\begin{array}{c}
-1-3 \mathrm{I} \\
1
\end{array}\right]\right]
\]
- Solution from eigenpair
\[
\mathrm{e}^{(-3-3 \mathrm{I}) t} \cdot\left[\begin{array}{c}
-1-3 \mathrm{I} \\
1
\end{array}\right]
\]
- Use Euler identity to write solution in terms of sin and cos
\[
\mathrm{e}^{-3 t} \cdot(\cos (3 t)-\mathrm{I} \sin (3 t)) \cdot\left[\begin{array}{c}
-1-3 \mathrm{I} \\
1
\end{array}\right]
\]
- Simplify expression
\[
\mathrm{e}^{-3 t} \cdot\left[\begin{array}{c}
(-1-3 \mathrm{I})(\cos (3 t)-\mathrm{I} \sin (3 t)) \\
\cos (3 t)-\mathrm{I} \sin (3 t)
\end{array}\right]
\]
- Both real and imaginary parts are solutions to the homogeneous system
\[
\left[\vec{x}_{1}(t)=\mathrm{e}^{-3 t} \cdot\left[\begin{array}{c}
-\cos (3 t)-3 \sin (3 t) \\
\cos (3 t)
\end{array}\right], \vec{x}_{2}(t)=\mathrm{e}^{-3 t} \cdot\left[\begin{array}{c}
-3 \cos (3 t)+\sin (3 t) \\
-\sin (3 t)
\end{array}\right]\right]
\]
- General solution to the system of ODEs
\[
\vec{x}=c_{1} \vec{x}_{1}(t)+c_{2} \vec{x}_{2}(t)
\]
- Substitute solutions into the general solution
\[
\vec{x}=\mathrm{e}^{-3 t} c_{1} \cdot\left[\begin{array}{c}
-\cos (3 t)-3 \sin (3 t) \\
\cos (3 t)
\end{array}\right]+c_{2} \mathrm{e}^{-3 t} \cdot\left[\begin{array}{c}
-3 \cos (3 t)+\sin (3 t) \\
-\sin (3 t)
\end{array}\right]
\]
- \(\quad\) Substitute in vector of dependent variables
\[
\left[\begin{array}{c}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
-\mathrm{e}^{-3 t}\left(\left(c_{1}+3 c_{2}\right) \cos (3 t)+3\left(c_{1}-\frac{c_{2}}{3}\right) \sin (3 t)\right) \\
\mathrm{e}^{-3 t}\left(c_{1} \cos (3 t)-c_{2} \sin (3 t)\right)
\end{array}\right]
\]
- \(\quad\) Solution to the system of ODEs
\[
\left\{x(t)=-\mathrm{e}^{-3 t}\left(\left(c_{1}+3 c_{2}\right) \cos (3 t)+3\left(c_{1}-\frac{c_{2}}{3}\right) \sin (3 t)\right), y(t)=\mathrm{e}^{-3 t}\left(c_{1} \cos (3 t)-c_{2} \sin (3 t)\right)\right\}
\]

\section*{Solution by Maple}

Time used: 0.031 (sec). Leaf size: 59
```

dsolve([diff(x(t),t)=-4*x(t)-10*y(t), diff (y(t),t)=x(t)-2*y(t)],singsol=all)

```
\[
\begin{aligned}
& x(t)=\mathrm{e}^{-3 t}\left(c_{1} \sin (3 t)+c_{2} \cos (3 t)\right) \\
& y(t)=-\frac{\mathrm{e}^{-3 t}\left(c_{1} \sin (3 t)-3 c_{2} \sin (3 t)+3 c_{1} \cos (3 t)+c_{2} \cos (3 t)\right)}{10}
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.009 (sec). Leaf size: 67
DSolve[\{x' \(\left.[t]==-4 * x[t]-10 * y[t], y^{\prime}[t]==x[t]-2 * y[t]\right\},\{x[t], y[t]\}, t\), IncludeSingularSolutions
\[
\begin{aligned}
& x(t) \rightarrow \frac{1}{3} e^{-3 t}\left(3 c_{1} \cos (3 t)-\left(c_{1}+10 c_{2}\right) \sin (3 t)\right) \\
& y(t) \rightarrow \frac{1}{3} e^{-3 t}\left(3 c_{2} \cos (3 t)+\left(c_{1}+c_{2}\right) \sin (3 t)\right)
\end{aligned}
\]

\subsection*{1.134 problem 193}
1.134.1 Solution using Matrix exponential method . . . . . . . . . . . . 1384
1.134.2 Solution using explicit Eigenvalue and Eigenvector method . . . 1385
1.134.3 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1390

Internal problem ID [12551]
Internal file name [OUTPUT/11203_Wednesday_October_18_2023_03_47_11_AM_532605/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 193.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "system of linear ODEs"
Solve
\[
\begin{aligned}
x^{\prime}(t) & =12 x(t)+18 y(t) \\
y^{\prime}(t) & =-8 x(t)-12 y(t)
\end{aligned}
\]

\subsection*{1.134.1 Solution using Matrix exponential method}

In this method, we will assume we have found the matrix exponential \(e^{A t}\) allready. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)
\]

Or
\[
\left[\begin{array}{l}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{cc}
12 & 18 \\
-8 & -12
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]

For the above matrix \(A\), the matrix exponential can be found to be
\[
e^{A t}=\left[\begin{array}{cc}
1+12 t & 18 t \\
-8 t & 1-12 t
\end{array}\right]
\]

Therefore the homogeneous solution is
\[
\begin{aligned}
\vec{x}_{h}(t) & =e^{A t} \vec{c} \\
& =\left[\begin{array}{cc}
1+12 t & 18 t \\
-8 t & 1-12 t
\end{array}\right]\left[\begin{array}{l}
c_{1} \\
c_{2}
\end{array}\right] \\
& =\left[\begin{array}{c}
(1+12 t) c_{1}+18 t c_{2} \\
-8 t c_{1}+(1-12 t) c_{2}
\end{array}\right] \\
& =\left[\begin{array}{c}
\left(12 c_{1}+18 c_{2}\right) t+c_{1} \\
\left(-8 c_{1}-12 c_{2}\right) t+c_{2}
\end{array}\right]
\end{aligned}
\]

Since no forcing function is given, then the final solution is \(\vec{x}_{h}(t)\) above.

\subsection*{1.134.2 Solution using explicit Eigenvalue and Eigenvector method}

This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)
\]

Or
\[
\left[\begin{array}{c}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{cc}
12 & 18 \\
-8 & -12
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]

The first step is find the homogeneous solution. We start by finding the eigenvalues of \(A\). This is done by solving the following equation for the eigenvalues \(\lambda\)
\[
\operatorname{det}(A-\lambda I)=0
\]

Expanding gives
\[
\operatorname{det}\left(\left[\begin{array}{cc}
12 & 18 \\
-8 & -12
\end{array}\right]-\lambda\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)=0
\]

Therefore
\[
\operatorname{det}\left(\left[\begin{array}{cc}
12-\lambda & 18 \\
-8 & -12-\lambda
\end{array}\right]\right)=0
\]

Which gives the characteristic equation
\[
\lambda^{2}=0
\]

The roots of the above are the eigenvalues.
\[
\lambda_{1}=0
\]

This table summarises the above result
\begin{tabular}{|l|l|l|}
\hline eigenvalue & algebraic multiplicity & type of eigenvalue \\
\hline 0 & 1 & real eigenvalue \\
\hline
\end{tabular}

Now the eigenvector for each eigenvalue are found.
Considering the eigenvalue \(\lambda_{1}=0\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
&\left(\left[\begin{array}{cc}
12 & 18 \\
-8 & -12
\end{array}\right]-(0)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
& {\left[\begin{array}{cc}
12 & 18 \\
-8 & -12
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] }
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{cc|c}
12 & 18 & 0 \\
-8 & -12 & 0
\end{array}\right]} \\
R_{2}=R_{2}+\frac{2 R_{1}}{3} \Longrightarrow\left[\begin{array}{cc|c}
12 & 18 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{cc}
12 & 18 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=-\frac{3 t}{2}\right\}\)

Hence the solution is
\[
\left[\begin{array}{c}
-\frac{3 t}{2} \\
t
\end{array}\right]=\left[\begin{array}{c}
-\frac{3 t}{2} \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
-\frac{3 t}{2} \\
t
\end{array}\right]=t\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{c}
-\frac{3 t}{2} \\
t
\end{array}\right]=\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right]
\]

Which is normalized to
\[
\left[\begin{array}{c}
-\frac{3 t}{2} \\
t
\end{array}\right]=\left[\begin{array}{c}
-3 \\
2
\end{array}\right]
\]

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity \(m\), and its geometric multiplicity \(k\) and the eigenvectors associated with the eigenvalue. If \(m>k\) then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity \(k\) ) does not equal the algebraic multiplicity \(m\), and we need to determine an additional \(m-k\) generalized eigenvectors for this eigenvalue.
\begin{tabular}{|c|c|c|c|c|}
\hline \multirow{2}{*}{ eigenvalue } & \multicolumn{2}{|c|}{ multiplicity } & & \\
\cline { 2 - 3 } & algebraic \(m\) & geometric \(k\) & defective? & eigenvectors \\
\hline 0 & 2 & 1 & Yes & {\(\left[\begin{array}{c}-\frac{3}{2} \\
1\end{array}\right]\)} \\
\hline
\end{tabular}

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue 0 is real and repated eigenvalue of multiplicity 2.There are two possible cases that can happen. This is illustrated in this diagram


Figure 216: Possible case for repeated \(\lambda\) of multiplicity 2

This eigenvalue has algebraic multiplicity of 2 , and geometric multiplicity 1 , therefore this is defective eigenvalue. The defect is 1 . This falls into case 2 shown above. We need to generate the missing additonal generalized eigevector \(\vec{v}_{2}\) by solving
\[
(A-\lambda I) \vec{v}_{2}=\vec{v}_{1}
\]

Where \(\vec{v}_{1}\) is the normal (rank 1) eigenvector found above. Hence we need to solve
\[
\begin{aligned}
&\left(\left[\begin{array}{cc}
12 & 18 \\
-8 & -12
\end{array}\right]-(0)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right] \\
& {\left[\begin{array}{cc}
12 & 18 \\
-8 & -12
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right] }
\end{aligned}
\]

Solving for \(\vec{v}_{2}\) gives
\[
\vec{v}_{2}=\left[\begin{array}{c}
-\frac{13}{8} \\
1
\end{array}\right]
\]

We have found two generalized eigenvectors for eigenvalue 0 . Therefore the two basis solution associated with this eigenvalue are
\[
\begin{aligned}
\vec{x}_{1}(t) & =\vec{v}_{1} e^{\lambda t} \\
& =\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right] 1 \\
& =\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right]
\end{aligned}
\]

And
\[
\begin{aligned}
\vec{x}_{2}(t) & =\left(\vec{v}_{1} t+\vec{v}_{2}\right) e^{\lambda t} \\
& =\left(\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right] t+\left[\begin{array}{c}
-\frac{13}{8} \\
1
\end{array}\right]\right) 1 \\
& =\left[\begin{array}{c}
-\frac{3 t}{2}-\frac{13}{8} \\
1+t
\end{array}\right]
\end{aligned}
\]

Therefore the final solution is
\[
\vec{x}_{h}(t)=c_{1} \vec{x}_{1}(t)+c_{2} \vec{x}_{2}(t)
\]

Which is written as
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=c_{1}\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right]+c_{2}\left[\begin{array}{c}
-\frac{3 t}{2}-\frac{13}{8} \\
1+t
\end{array}\right]
\]

Which becomes
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
\frac{(-12 t-13) c_{2}}{8}-\frac{3 c_{1}}{2} \\
c_{2} t+c_{1}+c_{2}
\end{array}\right]
\]

The following is the phase plot of the system.


Figure 217: Phase plot

\subsection*{1.134.3 Maple step by step solution}

Let's solve
\(\left[x^{\prime}(t)=12 x(t)+18 y(t), y^{\prime}(t)=-8 x(t)-12 y(t)\right]\)
- Define vector
\(\vec{x}(t)=\left[\begin{array}{l}x(t) \\ y(t)\end{array}\right]\)
- Convert system into a vector equation
\(\vec{x}^{\prime}(t)=\left[\begin{array}{cc}12 & 18 \\ -8 & -12\end{array}\right] \cdot \vec{x}(t)+\left[\begin{array}{l}0 \\ 0\end{array}\right]\)
- \(\quad\) System to solve
\(\vec{x}^{\prime}(t)=\left[\begin{array}{cc}12 & 18 \\ -8 & -12\end{array}\right] \cdot \vec{x}(t)\)
- Define the coefficient matrix
\[
A=\left[\begin{array}{cc}
12 & 18 \\
-8 & -12
\end{array}\right]
\]
- Rewrite the system as
\[
\vec{x}^{\prime}(t)=A \cdot \vec{x}(t)
\]
- To solve the system, find the eigenvalues and eigenvectors of \(A\)
- \(\quad\) Eigenpairs of \(A\)
\[
\left[\left[0,\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right]\right],\left[0,\left[\begin{array}{l}
0 \\
0
\end{array}\right]\right]\right]
\]
- Consider eigenpair
\[
\left[0,\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right]\right]
\]
- Solution to homogeneous system from eigenpair
\[
\vec{x}_{1}=\left[\begin{array}{c}
-\frac{3}{2} \\
1
\end{array}\right]
\]
- Consider eigenpair
\(\left[0,\left[\begin{array}{l}0 \\ 0\end{array}\right]\right]\)
- Solution to homogeneous system from eigenpair
\(\vec{x}_{2}=\left[\begin{array}{l}0 \\ 0\end{array}\right]\)
- General solution to the system of ODEs
\[
\vec{x}=c_{1} \vec{x}_{1}+c_{2} \vec{x}_{2}
\]
- Substitute solutions into the general solution
\[
\vec{x}=\left[\begin{array}{c}
-\frac{3 c_{1}}{2} \\
c_{1}
\end{array}\right]
\]
- Substitute in vector of dependent variables
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
-\frac{3 c_{1}}{2} \\
c_{1}
\end{array}\right]
\]
- \(\quad\) Solution to the system of ODEs
\[
\left\{x(t)=-\frac{3 c_{1}}{2}, y(t)=c_{1}\right\}
\]
\(\checkmark\) Solution by Maple
Time used: 0.015 (sec). Leaf size: 24
```

dsolve([diff(x(t),t)=12*x(t)+18*y(t), diff (y(t),t)=-8*x(t)-12*y(t)],singsol=all)

```
\[
\begin{aligned}
& x(t)=c_{1} t+c_{2} \\
& y(t)=\frac{1}{18} c_{1}-\frac{2}{3} c_{1} t-\frac{2}{3} c_{2}
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.003 (sec). Leaf size: 36
DSolve \(\left[\left\{x^{\prime}[t]==12 * x[t]+18 * y[t], y^{\prime}[t]==-8 * x[t]-12 * y[t]\right\},\{x[t], y[t]\}, t\right.\), IncludeSingularSolution
\[
\begin{aligned}
& x(t) \rightarrow 12 c_{1} t+18 c_{2} t+c_{1} \\
& y(t) \rightarrow c_{2}-4\left(2 c_{1}+3 c_{2}\right) t
\end{aligned}
\]

\subsection*{1.135 problem 194}
1.135.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 1393
1.135.2 Solving as riccati ode . . . . . . . . . . . . . . . . . . . . . . . . 1394

Internal problem ID [12552]
Internal file name [OUTPUT/11204_Wednesday_October_18_2023_03_47_11_AM_17442311/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 194.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "riccati"
Maple gives the following as the ode type
[[_Riccati, _special]]
\[
y^{\prime}-y^{2}=x
\]

With initial conditions
\[
[y(0)=1]
\]

\subsection*{1.135.1 Existence and uniqueness analysis}

This is non linear first order ODE. In canonical form it is written as
\[
\begin{aligned}
y^{\prime} & =f(x, y) \\
& =y^{2}+x
\end{aligned}
\]

The \(x\) domain of \(f(x, y)\) when \(y=1\) is
\[
\{-\infty<x<\infty\}
\]

And the point \(x_{0}=0\) is inside this domain. The \(y\) domain of \(f(x, y)\) when \(x=0\) is
\[
\{-\infty<y<\infty\}
\]

And the point \(y_{0}=1\) is inside this domain. Now we will look at the continuity of
\[
\begin{aligned}
\frac{\partial f}{\partial y} & =\frac{\partial}{\partial y}\left(y^{2}+x\right) \\
& =2 y
\end{aligned}
\]

The \(y\) domain of \(\frac{\partial f}{\partial y}\) when \(x=0\) is
\[
\{-\infty<y<\infty\}
\]

And the point \(y_{0}=1\) is inside this domain. Therefore solution exists and is unique.

\subsection*{1.135.2 Solving as riccati ode}

In canonical form the ODE is
\[
\begin{aligned}
y^{\prime} & =F(x, y) \\
& =y^{2}+x
\end{aligned}
\]

This is a Riccati ODE. Comparing the ODE to solve
\[
y^{\prime}=y^{2}+x
\]

With Riccati ODE standard form
\[
y^{\prime}=f_{0}(x)+f_{1}(x) y+f_{2}(x) y^{2}
\]

Shows that \(f_{0}(x)=x, f_{1}(x)=0\) and \(f_{2}(x)=1\). Let
\[
\begin{align*}
y & =\frac{-u^{\prime}}{f_{2} u} \\
& =\frac{-u^{\prime}}{u} \tag{1}
\end{align*}
\]

Using the above substitution in the given ODE results (after some simplification)in a second order ODE to solve for \(u(x)\) which is
\[
\begin{equation*}
f_{2} u^{\prime \prime}(x)-\left(f_{2}^{\prime}+f_{1} f_{2}\right) u^{\prime}(x)+f_{2}^{2} f_{0} u(x)=0 \tag{2}
\end{equation*}
\]

But
\[
\begin{aligned}
f_{2}^{\prime} & =0 \\
f_{1} f_{2} & =0 \\
f_{2}^{2} f_{0} & =x
\end{aligned}
\]

Substituting the above terms back in equation (2) gives
\[
u^{\prime \prime}(x)+x u(x)=0
\]

Solving the above ODE (this ode solved using Maple, not this program), gives
\[
u(x)=c_{1} \operatorname{AiryAi}(-x)+c_{2} \operatorname{AiryBi}(-x)
\]

The above shows that
\[
u^{\prime}(x)=-c_{1} \operatorname{AiryAi}(1,-x)-c_{2} \operatorname{AiryBi}(1,-x)
\]

Using the above in (1) gives the solution
\[
y=-\frac{-c_{1} \operatorname{Airy} \operatorname{Ai}(1,-x)-c_{2} \operatorname{AiryBi}(1,-x)}{c_{1} \operatorname{Airy} \operatorname{Ai}(-x)+c_{2} \operatorname{AiryBi}(-x)}
\]

Dividing both numerator and denominator by \(c_{1}\) gives, after renaming the constant \(\frac{c_{2}}{c_{1}}=c_{3}\) the following solution
\[
y=\frac{c_{3} \operatorname{AiryAi}(1,-x)+\operatorname{AiryBi}(1,-x)}{c_{3} \operatorname{AiryAi}(-x)+\operatorname{AiryBi}(-x)}
\]

Initial conditions are used to solve for \(c_{3}\). Substituting \(x=0\) and \(y=1\) in the above solution gives an equation to solve for the constant of integration.
\[
\begin{gathered}
1=\frac{3 \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{2}{3}}-3 \Gamma\left(\frac{2}{3}\right)^{2} c_{3} 3^{\frac{1}{6}}}{23^{\frac{5}{6}} \pi+2 \pi c_{3} 3^{\frac{1}{3}}} \\
c_{3}=\frac{-23^{\frac{5}{6}} \pi+3 \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{2}{3}}}{3 \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{1}{6}}+2 \pi 3^{\frac{1}{3}}}
\end{gathered}
\]

Substituting \(c_{3}\) found above in the general solution gives
\(y=\frac{-2 \operatorname{AiryAi}(1,-x) \pi 3^{\frac{5}{6}}+3 \operatorname{AiryAi}(1,-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{2}{3}}+3 \operatorname{AiryBi}(1,-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{1}{6}}+2 \operatorname{AiryBi}(1,-x) \pi 3}{-2 \operatorname{AiryAi}(-x) \pi 3^{\frac{5}{6}}+3 \operatorname{AiryAi}(-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{2}{3}}+3 \operatorname{AiryBi}(-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{1}{6}}+2 \operatorname{AiryBi}(-x) \pi 3^{\frac{1}{3}}}\)

\section*{Summary}

The solution(s) found are the following
\[
\begin{equation*}
=\frac{-2 \operatorname{AiryAi}(1,-x) \pi 3^{\frac{5}{6}}+3 \operatorname{AiryAi}(1,-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{2}{3}}+3 \operatorname{AiryBi}(1,-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{1}{6}}+2 \operatorname{AiryBi}(1,-x) \pi 3^{\frac{1}{3}}}{-2 \operatorname{Airy} \operatorname{Ai}(-x) \pi 3^{\frac{5}{6}}+3 \operatorname{Airy} \operatorname{Ai}(-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{2}{3}}+3 \operatorname{AiryBi}(-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{1}{6}}+2 \operatorname{AiryBi}(-x) \pi 3^{\frac{1}{3}}} \tag{1}
\end{equation*}
\]

(a) Solution plot (b) Slope field plot


\section*{Verification of solutions}
\(y\)
\(=\frac{-2 \operatorname{AiryAi}(1,-x) \pi 3^{\frac{5}{6}}+3 \operatorname{AiryAi}(1,-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{2}{3}}+3 \operatorname{AiryBi}(1,-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{1}{6}}+2 \operatorname{AiryBi}(1,-x) \pi 3^{\frac{1}{3}}}{-2 \operatorname{Airy} \operatorname{Ai}(-x) \pi 3^{\frac{5}{6}}+3 \operatorname{AiryAi}(-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{2}{3}}+3 \operatorname{AiryBi}(-x) \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{1}{6}}+2 \operatorname{AiryBi}(-x) \pi 3^{\frac{1}{3}}}\)
Verified OK.

Maple trace
```

`Methods for first order ODEs: --- Trying classification methods --- trying a quadrature trying 1st order linear trying Bernoulli trying separable trying inverse linear trying homogeneous types: trying Chini differential order: 1; looking for linear symmetries trying exact Looking for potential symmetries trying Riccati trying Riccati Special <- Riccati Special successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.109 (sec). Leaf size: 97
```

dsolve([diff(y(x),x)=y(x)^2+x,y(0) = 1],y(x), singsol=all)

```
\[
y(x)=\frac{\left(-23^{\frac{5}{6}} \pi+3 \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{2}{3}}\right) \operatorname{AiryAi}(1,-x)+\operatorname{AiryBi}(1,-x)\left(33^{\frac{1}{6}} \Gamma\left(\frac{2}{3}\right)^{2}+2 \pi 3^{\frac{1}{3}}\right)}{\left(-23^{\frac{5}{6}} \pi+3 \Gamma\left(\frac{2}{3}\right)^{2} 3^{\frac{2}{3}}\right) \operatorname{AiryAi}(-x)+\operatorname{AiryBi}(-x)\left(33^{\frac{1}{6}} \Gamma\left(\frac{2}{3}\right)^{2}+2 \pi 3^{\frac{1}{3}}\right)}
\]
\(\checkmark\) Solution by Mathematica
Time used: 1.986 (sec). Leaf size: 145
```

DSolve[{y'[x]==y[x]^2+x,{y[0]==1}},y[x],x,IncludeSingularSolutions -> True]

```
\(y(x)\)
\(\rightarrow \frac{\sqrt[3]{3} \operatorname{Gamma}\left(\frac{2}{3}\right)\left(x^{3 / 2} \operatorname{BesselJ}\left(-\frac{4}{3}, \frac{2 x^{3 / 2}}{3}\right)-x^{3 / 2} \operatorname{BesselJ}\left(\frac{2}{3}, \frac{2 x^{3 / 2}}{3}\right)+\operatorname{BesselJ}\left(-\frac{1}{3}, \frac{2 x^{3 / 2}}{3}\right)\right)-2 x^{3 / 2} \mathrm{Gat}}{2 x\left(\operatorname{Gamma}\left(\frac{1}{3}\right) \operatorname{BesselJ}\left(\frac{1}{3}, \frac{2 x^{3 / 2}}{3}\right)-\sqrt[3]{3} \operatorname{Gamma}\left(\frac{2}{3}\right) \operatorname{BesselJ}\left(-\frac{1}{3}, \frac{2 x^{3 / 2}}{3}\right)\right)}\)

\subsection*{1.136 problem 195}
1.136.1 Existence and uniqueness analysis . . . . . . . . . . . . . . . . . 1398
1.136.2 Solving as linear ode . . . . . . . . . . . . . . . . . . . . . . . . 1399
1.136.3 Solving as first order ode lie symmetry lookup ode . . . . . . . 1401
1.136.4 Solving as exact ode . . . . . . . . . . . . . . . . . . . . . . . . 1405
1.136.5 Maple step by step solution . . . . . . . . . . . . . . . . . . . . 1409

Internal problem ID [12553]
Internal file name [OUTPUT/11205_Wednesday_October_18_2023_03_47_12_AM_61480679/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR
PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 195.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "exact", "linear", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type
[_linear]
\[
y^{\prime}+\frac{y}{x}=\mathrm{e}^{x}
\]

With initial conditions
\[
[y(1)=1]
\]

\subsection*{1.136.1 Existence and uniqueness analysis}

This is a linear ODE. In canonical form it is written as
\[
y^{\prime}+p(x) y=q(x)
\]

Where here
\[
\begin{aligned}
& p(x)=\frac{1}{x} \\
& q(x)=\mathrm{e}^{x}
\end{aligned}
\]

Hence the ode is
\[
y^{\prime}+\frac{y}{x}=\mathrm{e}^{x}
\]

The domain of \(p(x)=\frac{1}{x}\) is
\[
\{x<0 \vee 0<x\}
\]

And the point \(x_{0}=1\) is inside this domain. The domain of \(q(x)=\mathrm{e}^{x}\) is
\[
\{-\infty<x<\infty\}
\]

And the point \(x_{0}=1\) is also inside this domain. Hence solution exists and is unique.

\subsection*{1.136.2 Solving as linear ode}

Entering Linear first order ODE solver. The integrating factor \(\mu\) is
\[
\begin{aligned}
& \mu=\mathrm{e}^{\int \frac{1}{x} d x} \\
& =x
\end{aligned}
\]

The ode becomes
\[
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} x}(\mu y) & =(\mu)\left(\mathrm{e}^{x}\right) \\
\frac{\mathrm{d}}{\mathrm{~d} x}(x y) & =(x)\left(\mathrm{e}^{x}\right) \\
\mathrm{d}(x y) & =\left(x \mathrm{e}^{x}\right) \mathrm{d} x
\end{aligned}
\]

Integrating gives
\[
\begin{aligned}
& x y=\int x \mathrm{e}^{x} \mathrm{~d} x \\
& x y=\mathrm{e}^{x}(x-1)+c_{1}
\end{aligned}
\]

Dividing both sides by the integrating factor \(\mu=x\) results in
\[
y=\frac{\mathrm{e}^{x}(x-1)}{x}+\frac{c_{1}}{x}
\]
which simplifies to
\[
y=\frac{\mathrm{e}^{x}(x-1)+c_{1}}{x}
\]

Initial conditions are used to solve for \(c_{1}\). Substituting \(x=1\) and \(y=1\) in the above solution gives an equation to solve for the constant of integration.
\[
\begin{aligned}
& 1=c_{1} \\
& c_{1}=1
\end{aligned}
\]

Substituting \(c_{1}\) found above in the general solution gives
\[
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+1}{x}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+1}{x} \tag{1}
\end{equation*}
\]

(a) Solution plot
(b) Slope field plot

\section*{Verification of solutions}
\[
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+1}{x}
\]

Verified OK.

\subsection*{1.136.3 Solving as first order ode lie symmetry lookup ode}

Writing the ode as
\[
\begin{aligned}
& y^{\prime}=-\frac{-x \mathrm{e}^{x}+y}{x} \\
& y^{\prime}=\omega(x, y)
\end{aligned}
\]

The condition of Lie symmetry is the linearized PDE given by
\[
\begin{equation*}
\eta_{x}+\omega\left(\eta_{y}-\xi_{x}\right)-\omega^{2} \xi_{y}-\omega_{x} \xi-\omega_{y} \eta=0 \tag{A}
\end{equation*}
\]

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find \(\xi, \eta\)

Table 209: Lie symmetry infinitesimal lookup table for known first order ODE's
\begin{tabular}{|c|c|c|c|}
\hline ODE class & Form & \(\xi\) & \(\eta\) \\
\hline linear ode & \(y^{\prime}=f(x) y(x)+g(x)\) & 0 & \(e^{\int f d x}\) \\
\hline separable ode & \(y^{\prime}=f(x) g(y)\) & \(\frac{1}{f}\) & 0 \\
\hline quadrature ode & \(y^{\prime}=f(x)\) & 0 & 1 \\
\hline quadrature ode & \(y^{\prime}=g(y)\) & 1 & 0 \\
\hline homogeneous ODEs of Class A & \(y^{\prime}=f\left(\frac{y}{x}\right)\) & \(x\) & \(y\) \\
\hline homogeneous ODEs of Class C & \(y^{\prime}=(a+b x+c y)^{\frac{n}{m}}\) & 1 & \[
-\frac{b}{c}
\] \\
\hline homogeneous class D & \(y^{\prime}=\frac{y}{x}+g(x) F\left(\frac{y}{x}\right)\) & \(x^{2}\) & \(x y\) \\
\hline First order special form ID 1 & \(y^{\prime}=g(x) e^{h(x)+b y}+f(x)\) & \[
\frac{e^{-\int b f(x) d x-h(x)}}{g(x)}
\] & \[
\frac{f(x) e^{-\int b f(x) d x-h(x)}}{g(x)}
\] \\
\hline polynomial type ode & \(y^{\prime}=\frac{a_{1} x+b_{1} y+c_{1}}{a_{2} x+b_{2} y+c_{2}}\) & \[
\frac{a_{1} b_{2} x-a_{2} b_{1} x-b_{1} c_{2}+b_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}
\] & \[
\frac{a_{1} b_{2} y-a_{2} b_{1} y-a_{1} c_{2}-a_{2} c_{1}}{a_{1} b_{2}-a_{2} b_{1}}
\] \\
\hline Bernoulli ode & \(y^{\prime}=f(x) y+g(x) y^{n}\) & 0 & \(e^{-\int(n-1) f(x) d x} y^{n}\) \\
\hline Reduced Riccati & \(y^{\prime}=f_{1}(x) y+f_{2}(x) y^{2}\) & 0 & \(e^{-\int f_{1} d x}\) \\
\hline
\end{tabular}

The above table shows that
\[
\begin{align*}
& \xi(x, y)=0 \\
& \eta(x, y)=\frac{1}{x} \tag{A1}
\end{align*}
\]

The next step is to determine the canonical coordinates \(R, S\). The canonical coordinates \(\operatorname{map}(x, y) \rightarrow(R, S)\) where \((R, S)\) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is
\[
\begin{equation*}
\frac{d x}{\xi}=\frac{d y}{\eta}=d S \tag{1}
\end{equation*}
\]

The above comes from the requirements that \(\left(\xi \frac{\partial}{\partial x}+\eta \frac{\partial}{\partial y}\right) S(x, y)=1\). Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable \(R\) in the canonical coordinates, where \(S(R)\). Since \(\xi=0\) then in this special case
\[
R=x
\]
\(S\) is found from
\[
\begin{aligned}
S & =\int \frac{1}{\eta} d y \\
& =\int \frac{1}{\frac{1}{x}} d y
\end{aligned}
\]

Which results in
\[
S=x y
\]

Now that \(R, S\) are found, we need to setup the ode in these coordinates. This is done by evaluating
\[
\begin{equation*}
\frac{d S}{d R}=\frac{S_{x}+\omega(x, y) S_{y}}{R_{x}+\omega(x, y) R_{y}} \tag{2}
\end{equation*}
\]

Where in the above \(R_{x}, R_{y}, S_{x}, S_{y}\) are all partial derivatives and \(\omega(x, y)\) is the right hand side of the original ode given by
\[
\omega(x, y)=-\frac{-x \mathrm{e}^{x}+y}{x}
\]

Evaluating all the partial derivatives gives
\[
\begin{aligned}
R_{x} & =1 \\
R_{y} & =0 \\
S_{x} & =y \\
S_{y} & =x
\end{aligned}
\]

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.
\[
\begin{equation*}
\frac{d S}{d R}=x \mathrm{e}^{x} \tag{2~A}
\end{equation*}
\]

We now need to express the RHS as function of \(R\) only. This is done by solving for \(x, y\) in terms of \(R, S\) from the result obtained earlier and simplifying. This gives
\[
\frac{d S}{d R}=R \mathrm{e}^{R}
\]

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates \(R, S\). Integrating the above gives
\[
\begin{equation*}
S(R)=(R-1) \mathrm{e}^{R}+c_{1} \tag{4}
\end{equation*}
\]

To complete the solution, we just need to transform (4) back to \(x, y\) coordinates. This results in
\[
y x=\mathrm{e}^{x}(x-1)+c_{1}
\]

Which simplifies to
\[
y x=\mathrm{e}^{x}(x-1)+c_{1}
\]

Which gives
\[
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+c_{1}}{x}
\]

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.
\begin{tabular}{|c|c|c|}
\hline Original ode in \(x, y\) coordinates & \[
\begin{gathered}
\text { Canonical } \\
\text { coordinates } \\
\text { transformation }
\end{gathered}
\] & ODE in canonical coordinates
\[
(R, S)
\] \\
\hline \(\frac{d y}{d x}=-\frac{-x \mathrm{e}^{x}+y}{x}\) & & \(\frac{d S}{d R}=R \mathrm{e}^{R}\) \\
\hline  & &  \\
\hline  & &  \\
\hline  & & \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow-S T R]{ }\) \\
\hline  & & \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow+]{\rightarrow \rightarrow+}\) \\
\hline  & \(R=x\) & \(\xrightarrow[\rightarrow \rightarrow \rightarrow \rightarrow \rightarrow \infty]{\rightarrow \rightarrow-\infty}\) \\
\hline  & &  \\
\hline  & \(S=x y\) &  \\
\hline  & &  \\
\hline  & &  \\
\hline  & &  \\
\hline  & &  \\
\hline
\end{tabular}

Initial conditions are used to solve for \(c_{1}\). Substituting \(x=1\) and \(y=1\) in the above solution gives an equation to solve for the constant of integration.
\[
\begin{aligned}
& 1=c_{1} \\
& c_{1}=1
\end{aligned}
\]

Substituting \(c_{1}\) found above in the general solution gives
\[
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+1}{x}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+1}{x} \tag{1}
\end{equation*}
\]

(a) Solution plot

\section*{Verification of solutions}
\[
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+1}{x}
\]

Verified OK.

\subsection*{1.136.4 Solving as exact ode}

Entering Exact first order ODE solver. (Form one type)
To solve an ode of the form
\[
\begin{equation*}
M(x, y)+N(x, y) \frac{d y}{d x}=0 \tag{A}
\end{equation*}
\]

We assume there exists a function \(\phi(x, y)=c\) where \(c\) is constant, that satisfies the ode. Taking derivative of \(\phi\) w.r.t. \(x\) gives
\[
\frac{d}{d x} \phi(x, y)=0
\]

Hence
\[
\begin{equation*}
\frac{\partial \phi}{\partial x}+\frac{\partial \phi}{\partial y} \frac{d y}{d x}=0 \tag{B}
\end{equation*}
\]

Comparing ( \(\mathrm{A}, \mathrm{B}\) ) shows that
\[
\begin{aligned}
& \frac{\partial \phi}{\partial x}=M \\
& \frac{\partial \phi}{\partial y}=N
\end{aligned}
\]

But since \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) then for the above to be valid, we require that
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

If the above condition is satisfied, then the original ode is called exact. We still need to determine \(\phi(x, y)\) but at least we know now that we can do that since the condition \(\frac{\partial^{2} \phi}{\partial x \partial y}=\frac{\partial^{2} \phi}{\partial y \partial x}\) is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is
\[
\begin{equation*}
M(x, y) \mathrm{d} x+N(x, y) \mathrm{d} y=0 \tag{1A}
\end{equation*}
\]

Therefore
\[
\begin{align*}
(x) \mathrm{d} y & =\left(x \mathrm{e}^{x}-y\right) \mathrm{d} x \\
\left(-x \mathrm{e}^{x}+y\right) \mathrm{d} x+(x) \mathrm{d} y & =0 \tag{2~A}
\end{align*}
\]

Comparing (1A) and (2A) shows that
\[
\begin{aligned}
M(x, y) & =-x \mathrm{e}^{x}+y \\
N(x, y) & =x
\end{aligned}
\]

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied
\[
\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}
\]

Using result found above gives
\[
\begin{aligned}
\frac{\partial M}{\partial y} & =\frac{\partial}{\partial y}\left(-x \mathrm{e}^{x}+y\right) \\
& =1
\end{aligned}
\]

And
\[
\begin{aligned}
\frac{\partial N}{\partial x} & =\frac{\partial}{\partial x}(x) \\
& =1
\end{aligned}
\]

Since \(\frac{\partial M}{\partial y}=\frac{\partial N}{\partial x}\), then the ODE is exact The following equations are now set up to solve for the function \(\phi(x, y)\)
\[
\begin{align*}
& \frac{\partial \phi}{\partial x}=M  \tag{1}\\
& \frac{\partial \phi}{\partial y}=N \tag{2}
\end{align*}
\]

Integrating (1) w.r.t. \(x\) gives
\[
\begin{align*}
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int M \mathrm{~d} x \\
\int \frac{\partial \phi}{\partial x} \mathrm{~d} x & =\int-x \mathrm{e}^{x}+y \mathrm{~d} x \\
\phi & =x y-x \mathrm{e}^{x}+\mathrm{e}^{x}+f(y) \tag{3}
\end{align*}
\]

Where \(f(y)\) is used for the constant of integration since \(\phi\) is a function of both \(x\) and \(y\). Taking derivative of equation (3) w.r.t \(y\) gives
\[
\begin{equation*}
\frac{\partial \phi}{\partial y}=x+f^{\prime}(y) \tag{4}
\end{equation*}
\]

But equation (2) says that \(\frac{\partial \phi}{\partial y}=x\). Therefore equation (4) becomes
\[
\begin{equation*}
x=x+f^{\prime}(y) \tag{5}
\end{equation*}
\]

Solving equation (5) for \(f^{\prime}(y)\) gives
\[
f^{\prime}(y)=0
\]

Therefore
\[
f(y)=c_{1}
\]

Where \(c_{1}\) is constant of integration. Substituting this result for \(f(y)\) into equation (3) gives \(\phi\)
\[
\phi=x y-x \mathrm{e}^{x}+\mathrm{e}^{x}+c_{1}
\]

But since \(\phi\) itself is a constant function, then let \(\phi=c_{2}\) where \(c_{2}\) is new constant and combining \(c_{1}\) and \(c_{2}\) constants into new constant \(c_{1}\) gives the solution as
\[
c_{1}=x y-x \mathrm{e}^{x}+\mathrm{e}^{x}
\]

The solution becomes
\[
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+c_{1}}{x}
\]

Initial conditions are used to solve for \(c_{1}\). Substituting \(x=1\) and \(y=1\) in the above solution gives an equation to solve for the constant of integration.
\[
\begin{aligned}
& 1=c_{1} \\
& c_{1}=1
\end{aligned}
\]

Substituting \(c_{1}\) found above in the general solution gives
\[
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+1}{x}
\]

Summary
The solution(s) found are the following
\[
\begin{equation*}
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+1}{x} \tag{1}
\end{equation*}
\]

(a) Solution plot
(b) Slope field plot

\section*{Verification of solutions}
\[
y=\frac{x \mathrm{e}^{x}-\mathrm{e}^{x}+1}{x}
\]

Verified OK.

\subsection*{1.136.5 Maple step by step solution}

Let's solve
\[
\left[y^{\prime}+\frac{y}{x}=\mathrm{e}^{x}, y(1)=1\right]
\]
- Highest derivative means the order of the ODE is 1
\(y^{\prime}\)
- Isolate the derivative
\(y^{\prime}=-\frac{y}{x}+\mathrm{e}^{x}\)
- Group terms with \(y\) on the lhs of the ODE and the rest on the rhs of the ODE \(y^{\prime}+\frac{y}{x}=\mathrm{e}^{x}\)
- The ODE is linear; multiply by an integrating factor \(\mu(x)\)
\(\mu(x)\left(y^{\prime}+\frac{y}{x}\right)=\mu(x) \mathrm{e}^{x}\)
- Assume the lhs of the ODE is the total derivative \(\frac{d}{d x}(\mu(x) y)\)
\(\mu(x)\left(y^{\prime}+\frac{y}{x}\right)=\mu^{\prime}(x) y+\mu(x) y^{\prime}\)
- Isolate \(\mu^{\prime}(x)\)
\(\mu^{\prime}(x)=\frac{\mu(x)}{x}\)
- Solve to find the integrating factor
\(\mu(x)=x\)
- Integrate both sides with respect to \(x\)
\(\int\left(\frac{d}{d x}(\mu(x) y)\right) d x=\int \mu(x) \mathrm{e}^{x} d x+c_{1}\)
- Evaluate the integral on the lhs
\(\mu(x) y=\int \mu(x) \mathrm{e}^{x} d x+c_{1}\)
- \(\quad\) Solve for \(y\)
\(y=\frac{\int \mu(x) \mathrm{e}^{x} d x+c_{1}}{\mu(x)}\)
- \(\quad\) Substitute \(\mu(x)=x\)
\(y=\frac{\int x \mathrm{e}^{x} d x+c_{1}}{x}\)
- Evaluate the integrals on the rhs
\(y=\frac{\mathrm{e}^{x}(x-1)+c_{1}}{x}\)
- Use initial condition \(y(1)=1\)
\(1=c_{1}\)
- \(\quad\) Solve for \(c_{1}\)
\(c_{1}=1\)
- Substitute \(c_{1}=1\) into general solution and simplify
\(y=\frac{\mathrm{e}^{x}(x-1)+1}{x}\)
- Solution to the IVP
\(y=\frac{\mathrm{e}^{x}(x-1)+1}{x}\)

Maple trace
```

`Methods for first order ODEs: --- Trying classification methods --- trying a quadrature trying 1st order linear <- 1st order linear successful`

```
\(\checkmark\) Solution by Maple
Time used: 0.016 (sec). Leaf size: 16
```

dsolve([diff(y(x),x)+1/x*y(x)=exp(x),y(1) = 1],y(x), singsol=all)

```
\[
y(x)=\frac{(-1+x) \mathrm{e}^{x}+1}{x}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.082 (sec). Leaf size: 18
DSolve[\{y' \([x]+1 / x * y[x]==\operatorname{Exp}[x],\{y[1]==1\}\}, y[x], x\), IncludeSingularSolutions \(->\) True]
\[
y(x) \rightarrow \frac{e^{x}(x-1)+1}{x}
\]

\subsection*{1.137 problem 196}
1.137.1 Solution using Matrix exponential method . . . . . . . . . . . . 1411
1.137.2 Solution using explicit Eigenvalue and Eigenvector method . . . 1412

Internal problem ID [12554]
Internal file name [OUTPUT/11206_Wednesday_October_18_2023_03_47_13_AM_92578670/index.tex]
Book: DIFFERENTIAL and INTEGRAL CALCULUS. VOL I. by N. PISKUNOV. MIR PUBLISHERS, Moscow 1969.
Section: Chapter 8. Differential equations. Exercises page 595
Problem number: 196.
ODE order: 1.
ODE degree: 1 .

The type(s) of ODE detected by this program : "system of linear ODEs"
Solve
\[
\begin{aligned}
x^{\prime}(t) & =-x(t)+y(t) \\
y^{\prime}(t) & =-x(t)-3 y(t)
\end{aligned}
\]

With initial conditions
\[
[x(1)=0, y(1)=1]
\]

\subsection*{1.137.1 Solution using Matrix exponential method}

In this method, we will assume we have found the matrix exponential \(e^{A t}\) allready. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)
\]

Or
\[
\left[\begin{array}{l}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{cc}
-1 & 1 \\
-1 & -3
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]

For the above matrix \(A\), the matrix exponential can be found to be
\[
e^{A t}=\left[\begin{array}{cc}
\mathrm{e}^{-2 t}(1+t) & t \mathrm{e}^{-2 t} \\
-t \mathrm{e}^{-2 t} & \mathrm{e}^{-2 t}(1-t)
\end{array}\right]
\]

Therefore the homogeneous solution is
\[
\begin{aligned}
\vec{x}_{h}(t) & =e^{A t} \vec{x}_{0} \\
& =\left[\begin{array}{cc}
\mathrm{e}^{-2 t}(1+t) & t \mathrm{e}^{-2 t} \\
-t \mathrm{e}^{-2 t} & \mathrm{e}^{-2 t}(1-t)
\end{array}\right]\left[\begin{array}{l}
0 \\
1
\end{array}\right] \\
& =\left[\begin{array}{c}
t \mathrm{e}^{-2 t} \\
\mathrm{e}^{-2 t}(1-t)
\end{array}\right]
\end{aligned}
\]

Since no forcing function is given, then the final solution is \(\vec{x}_{h}(t)\) above.

\subsection*{1.137.2 Solution using explicit Eigenvalue and Eigenvector method}

This is a system of linear ODE's given as
\[
\vec{x}^{\prime}(t)=A \vec{x}(t)
\]

Or
\[
\left[\begin{array}{c}
x^{\prime}(t) \\
y^{\prime}(t)
\end{array}\right]=\left[\begin{array}{cc}
-1 & 1 \\
-1 & -3
\end{array}\right]\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]
\]

The first step is find the homogeneous solution. We start by finding the eigenvalues of \(A\). This is done by solving the following equation for the eigenvalues \(\lambda\)
\[
\operatorname{det}(A-\lambda I)=0
\]

Expanding gives
\[
\operatorname{det}\left(\left[\begin{array}{cc}
-1 & 1 \\
-1 & -3
\end{array}\right]-\lambda\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)=0
\]

Therefore
\[
\operatorname{det}\left(\left[\begin{array}{cc}
-1-\lambda & 1 \\
-1 & -3-\lambda
\end{array}\right]\right)=0
\]

Which gives the characteristic equation
\[
\lambda^{2}+4 \lambda+4=0
\]

The roots of the above are the eigenvalues.
\[
\lambda_{1}=-2
\]

This table summarises the above result
\begin{tabular}{|l|l|l|}
\hline eigenvalue & algebraic multiplicity & type of eigenvalue \\
\hline-2 & 1 & real eigenvalue \\
\hline
\end{tabular}

Now the eigenvector for each eigenvalue are found.
Considering the eigenvalue \(\lambda_{1}=-2\)
We need to solve \(A \vec{v}=\lambda \vec{v}\) or \((A-\lambda I) \vec{v}=\overrightarrow{0}\) which becomes
\[
\begin{aligned}
&\left(\left[\begin{array}{cc}
-1 & 1 \\
-1 & -3
\end{array}\right]-(-2)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right)\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] \\
& {\left[\begin{array}{cc}
1 & 1 \\
-1 & -1
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right] }
\end{aligned}
\]

Now forward elimination is applied to solve for the eigenvector \(\vec{v}\). The augmented matrix is
\[
\begin{gathered}
{\left[\begin{array}{cc|c}
1 & 1 & 0 \\
-1 & -1 & 0
\end{array}\right]} \\
R_{2}=R_{2}+R_{1} \Longrightarrow\left[\begin{array}{ll|l}
1 & 1 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{gathered}
\]

Therefore the system in Echelon form is
\[
\left[\begin{array}{ll}
1 & 1 \\
0 & 0
\end{array}\right]\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
\]

The free variables are \(\left\{v_{2}\right\}\) and the leading variables are \(\left\{v_{1}\right\}\). Let \(v_{2}=t\). Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation \(\left\{v_{1}=-t\right\}\)

Hence the solution is
\[
\left[\begin{array}{c}
-t \\
t
\end{array}\right]=\left[\begin{array}{c}
-t \\
t
\end{array}\right]
\]

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as
\[
\left[\begin{array}{c}
-t \\
t
\end{array}\right]=t\left[\begin{array}{c}
-1 \\
1
\end{array}\right]
\]

Let \(t=1\) the eigenvector becomes
\[
\left[\begin{array}{c}
-t \\
t
\end{array}\right]=\left[\begin{array}{c}
-1 \\
1
\end{array}\right]
\]

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity \(m\), and its geometric multiplicity \(k\) and the eigenvectors associated with the eigenvalue. If \(m>k\) then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity \(k\) ) does not equal the algebraic multiplicity \(m\), and we need to determine an additional \(m-k\) generalized eigenvectors for this eigenvalue.
\begin{tabular}{|c|c|c|c|c|}
\hline \multirow{2}{*}{ eigenvalue } & \multicolumn{2}{|c|}{ multiplicity } & & \\
\cline { 2 - 3 } & algebraic \(m\) & geometric \(k\) & defective? & eigenvectors \\
\hline-2 & 2 & 1 & Yes & {\(\left[\begin{array}{c}-1 \\
1\end{array}\right]\)} \\
\hline
\end{tabular}

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue -2 is real and repated eigenvalue of multiplicity 2.There are two possible cases that can happen. This is illustrated in this diagram


Figure 222: Possible case for repeated \(\lambda\) of multiplicity 2

This eigenvalue has algebraic multiplicity of 2 , and geometric multiplicity 1 , therefore this is defective eigenvalue. The defect is 1 . This falls into case 2 shown above. We need to generate the missing additonal generalized eigevector \(\vec{v}_{2}\) by solving
\[
(A-\lambda I) \vec{v}_{2}=\vec{v}_{1}
\]

Where \(\vec{v}_{1}\) is the normal (rank 1) eigenvector found above. Hence we need to solve
\[
\begin{aligned}
\left(\left[\begin{array}{cc}
-1 & 1 \\
-1 & -3
\end{array}\right]-(-2)\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]\right) & {\left[\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right] }
\end{aligned}=\left[\begin{array}{c}
-1 \\
1
\end{array}\right]
\]

Solving for \(\vec{v}_{2}\) gives
\[
\vec{v}_{2}=\left[\begin{array}{c}
-2 \\
1
\end{array}\right]
\]

We have found two generalized eigenvectors for eigenvalue -2 . Therefore the two basis solution associated with this eigenvalue are
\[
\begin{aligned}
\vec{x}_{1}(t) & =\vec{v}_{1} e^{\lambda t} \\
& =\left[\begin{array}{c}
-1 \\
1
\end{array}\right] \mathrm{e}^{-2 t} \\
& =\left[\begin{array}{c}
-\mathrm{e}^{-2 t} \\
\mathrm{e}^{-2 t}
\end{array}\right]
\end{aligned}
\]

And
\[
\begin{aligned}
\vec{x}_{2}(t) & =\left(\vec{v}_{1} t+\vec{v}_{2}\right) e^{\lambda t} \\
& =\left(\left[\begin{array}{c}
-1 \\
1
\end{array}\right] t+\left[\begin{array}{c}
-2 \\
1
\end{array}\right]\right) \mathrm{e}^{-2 t} \\
& =\left[\begin{array}{c}
-\mathrm{e}^{-2 t}(t+2) \\
\mathrm{e}^{-2 t}(1+t)
\end{array}\right]
\end{aligned}
\]

Therefore the final solution is
\[
\vec{x}_{h}(t)=c_{1} \vec{x}_{1}(t)+c_{2} \vec{x}_{2}(t)
\]

Which is written as
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=c_{1}\left[\begin{array}{c}
-\mathrm{e}^{-2 t} \\
\mathrm{e}^{-2 t}
\end{array}\right]+c_{2}\left[\begin{array}{c}
\mathrm{e}^{-2 t}(-t-2) \\
\mathrm{e}^{-2 t}(1+t)
\end{array}\right]
\]

Which becomes
\[
\left[\begin{array}{c}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
-\left((t+2) c_{2}+c_{1}\right) \mathrm{e}^{-2 t} \\
\mathrm{e}^{-2 t}\left(c_{2} t+c_{1}+c_{2}\right)
\end{array}\right]
\]

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions
\[
\left[\begin{array}{l}
x(1)=0  \tag{1}\\
y(1)=1
\end{array}\right]
\]

Substituting initial conditions into the above solution at \(t=1\) gives
\[
\left[\begin{array}{l}
0 \\
1
\end{array}\right]=\left[\begin{array}{c}
-\left(3 c_{2}+c_{1}\right) \mathrm{e}^{-2} \\
\mathrm{e}^{-2}\left(c_{1}+2 c_{2}\right)
\end{array}\right]
\]

Solving for the constants of integrations gives
\[
\left[\begin{array}{c}
c_{1}=3 \mathrm{e}^{2} \\
c_{2}=-\mathrm{e}^{2}
\end{array}\right]
\]

Substituting these constants back in original solution in Eq. (1) gives
\[
\left[\begin{array}{l}
x(t) \\
y(t)
\end{array}\right]=\left[\begin{array}{c}
-\left(-(t+2) \mathrm{e}^{2}+3 \mathrm{e}^{2}\right) \mathrm{e}^{-2 t} \\
\mathrm{e}^{-2 t}\left(-\mathrm{e}^{2} t+2 \mathrm{e}^{2}\right)
\end{array}\right]
\]

The following is the phase plot of the system.


Figure 223: Phase plot

The following are plots of each solution.

\(\checkmark\) Solution by Maple
Time used: 0.031 (sec). Leaf size: 45
```

dsolve([diff(x(t),t) = y(t)-x(t), diff(y(t),t) = -x(t)-3*y(t), x(1) = 0, y(1) = 1], singsol=

```
\[
\begin{aligned}
x(t) & =\mathrm{e}^{-2 t}\left(t \mathrm{e}^{2}-\mathrm{e}^{2}\right) \\
y(t) & =-\mathrm{e}^{-2 t}\left(t \mathrm{e}^{2}-2 \mathrm{e}^{2}\right)
\end{aligned}
\]
\(\checkmark\) Solution by Mathematica
Time used: 0.006 (sec). Leaf size: 31
```

DSolve[{x'[t]==y[t]-x[t], y'[t]==-x[t]-3*y[t]},{x[1]==0,y[1]==1},{x[t],y[t]},t, IncludeSingula

```
\[
\begin{aligned}
& x(t) \rightarrow e^{2-2 t}(t-1) \\
& y(t) \rightarrow-e^{2-2 t}(t-2)
\end{aligned}
\]```

