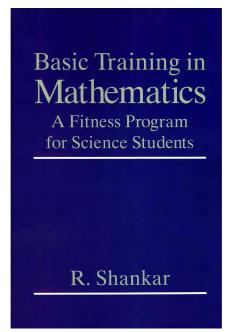
A Solution Manual For

# Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995



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#### 1.1 problem 10.2.4

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Internal problem ID [5045] Internal file name [OUTPUT/4538\_Sunday\_June\_05\_2022\_03\_00\_32\_PM\_16368218/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.2, ODEs with constant Coefficients. page 307
Problem number: 10.2.4.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second\_order\_linear\_constant\_coeff", "second\_order\_ode\_can\_be\_made\_integrable"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$x'' - \omega^2 x = 0$$

#### 1.1.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ax''(t) + Bx'(t) + Cx(t) = 0$$

Where in the above  $A = 1, B = 0, C = -\omega^2$ . Let the solution be  $x = e^{\lambda t}$ . Substituting this into the ODE gives

$$\lambda^2 \mathrm{e}^{\lambda t} - \omega^2 \mathrm{e}^{\lambda t} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by  $e^{\lambda t}$  gives

$$\lambda^2 - \omega^2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting  $A = 1, B = 0, C = -\omega^2$  into the above gives

$$\lambda_{1,2} = \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(-\omega^2)} \\ = \pm \sqrt{\omega^2}$$

Hence

$$\lambda_1 = +\sqrt{\omega^2}$$
$$\lambda_2 = -\sqrt{\omega^2}$$

Which simplifies to

$$\lambda_1 = \sqrt{\omega^2}$$
$$\lambda_2 = -\sqrt{\omega^2}$$

Since roots are real and distinct, then the solution is

$$x = c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t}$$
$$x = c_1 e^{\left(\sqrt{\omega^2}\right)t} + c_2 e^{\left(-\sqrt{\omega^2}\right)t}$$

Or

$$x = c_1 \mathrm{e}^{\sqrt{\omega^2}t} + c_2 \mathrm{e}^{-\sqrt{\omega^2}t}$$

Summary

The solution(s) found are the following

$$x = c_1 e^{\sqrt{\omega^2 t}} + c_2 e^{-\sqrt{\omega^2 t}}$$
 (1)

Verification of solutions

$$x = c_1 \mathrm{e}^{\sqrt{\omega^2} t} + c_2 \mathrm{e}^{-\sqrt{\omega^2} t}$$

Verified OK.

# 1.1.2 Solving as second order ode can be made integrable ode

Multiplying the ode by x' gives

$$x'x'' - \omega^2 x'x = 0$$

Integrating the above w.r.t t gives

$$\int \left(x'x'' - \omega^2 x'x\right) dt = 0$$
$$\frac{x'^2}{2} - \frac{\omega^2 x^2}{2} = c_2$$

Which is now solved for x. Solving the given ode for x' results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$x' = \sqrt{\omega^2 x^2 + 2c_1} \tag{1}$$

$$x' = -\sqrt{\omega^2 x^2 + 2c_1} \tag{2}$$

Now each one of the above ODE is solved.

Solving equation (1)

Integrating both sides gives

$$\int \frac{1}{\sqrt{\omega^2 x^2 + 2c_1}} dx = \int dt$$
$$\frac{\ln\left(\frac{\omega^2 x}{\sqrt{\omega^2}} + \sqrt{\omega^2 x^2 + 2c_1}\right)}{\sqrt{\omega^2}} = t + c_2$$

Raising both side to exponential gives

$$\mathrm{e}^{\frac{\ln\left(\frac{\omega^2 x}{\sqrt{\omega^2}} + \sqrt{\omega^2 x^2 + 2c_1}\right)}{\sqrt{\omega^2}}} = \mathrm{e}^{t+c_2}$$

Which simplifies to

$$\left(\omega x \operatorname{csgn}(\omega) + \sqrt{\omega^2 x^2 + 2c_1}\right)^{\frac{1}{\sqrt{\omega^2}}} = c_3 \mathrm{e}^t$$
  
$$\operatorname{n} x = \frac{\operatorname{csgn}(\omega)\left((c_3 \mathrm{e}^t)^{\operatorname{csgn}(\omega)\omega} - 2(c_3 \mathrm{e}^t)^{-\operatorname{csgn}(\omega)\omega} c_1\right)}{2\omega} \text{ to } x = \frac{(c_3 \mathrm{e}^t)^{\omega} - 2(c_3 \mathrm{e}^t)^{-\omega} c_1}{2\omega}$$

Simplifying the solution aSolving equation (2)

Integrating both sides gives

$$\int -\frac{1}{\sqrt{\omega^2 x^2 + 2c_1}} dx = \int dt$$
$$\frac{\ln\left(\frac{\omega^2 x}{\sqrt{\omega^2}} + \sqrt{\omega^2 x^2 + 2c_1}\right)}{\sqrt{\omega^2}} = t + c_4$$

Raising both side to exponential gives

$$\mathrm{e}^{-\frac{\ln\left(\frac{\omega^2 x}{\sqrt{\omega^2}} + \sqrt{\omega^2 x^2 + 2c_1}\right)}{\sqrt{\omega^2}}} = \mathrm{e}^{t+c_4}$$

Which simplifies to

$$\left(\omega x \operatorname{csgn}(\omega) + \sqrt{\omega^2 x^2 + 2c_1}\right)^{-\frac{\operatorname{csgn}(\omega)}{\omega}} = c_5 e^t$$

Simplifying the solution  $x = -\frac{\operatorname{csgn}(\omega)\left(2(c_5e^t)^{\operatorname{csgn}(\omega)\omega}c_1 - (c_5e^t)^{-\operatorname{csgn}(\omega)\omega}\right)}{2\omega}$  to  $x = -\frac{2(c_5e^t)^{\omega}c_1 - (c_5e^t)^{-\omega}}{2\omega}$ Summary

The solution(s) found are the following

$$x = \frac{(c_3 e^t)^{\omega} - 2(c_3 e^t)^{-\omega} c_1}{2\omega}$$
(1)

$$x = -\frac{2(c_5 e^t)^{\omega} c_1 - (c_5 e^t)^{-\omega}}{2\omega}$$
(2)

Verification of solutions

$$x = \frac{(c_3 e^t)^{\omega} - 2(c_3 e^t)^{-\omega} c_1}{2\omega}$$

Verified OK.

$$x = -rac{2(c_5 \mathrm{e}^t)^{\omega} c_1 - (c_5 \mathrm{e}^t)^{-\omega}}{2\omega}$$

Verified OK.

#### 1.1.3 Solving using Kovacic algorithm

Writing the ode as

$$x'' - \omega^2 x = 0 \tag{1}$$

$$Ax'' + Bx' + Cx = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$
  

$$B = 0$$

$$C = -\omega^2$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(t) = x e^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{\omega^2}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = \omega^2$$
$$t = 1$$

Therefore eq. (4) becomes

$$z''(t) = \left(\omega^2\right) z(t) \tag{7}$$

Equation (7) is now solved. After finding z(t) then x is found using the inverse transformation

$$x = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 1: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 0 - 0$$
$$= 0$$

There are no poles in r. Therefore the set of poles  $\Gamma$  is empty. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since  $r = \omega^2$  is not a function of t, then there is no need run Kovacic algorithm to obtain a solution for transformed ode z'' = rz as one solution is

$$z_1(t) = \mathrm{e}^{\sqrt{\omega^2} t}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in x is found from

$$x_1 = z_1 e^{\int -\frac{1}{2}\frac{B}{A}\,dt}$$

Since B = 0 then the above reduces to

$$\begin{aligned} x_1 &= z_1 \\ &= \mathrm{e}^{\sqrt{\omega^2} t} \end{aligned}$$

Which simplifies to

$$x_1 = \mathrm{e}^{\sqrt{\omega^2} t}$$

The second solution  $x_2$  to the original ode is found using reduction of order

$$x_2 = x_1 \int \frac{e^{\int -\frac{B}{A} dt}}{x_1^2} dt$$

Since B = 0 then the above becomes

$$\begin{aligned} x_2 &= x_1 \int \frac{1}{x_1^2} dt \\ &= e^{\sqrt{\omega^2} t} \int \frac{1}{e^{2\sqrt{\omega^2} t}} dt \\ &= e^{\sqrt{\omega^2} t} \left( -\frac{\operatorname{csgn}(\omega) e^{-2\operatorname{csgn}(\omega)\omega t}}{2\omega} \right) \end{aligned}$$

Therefore the solution is

$$x = c_1 x_1 + c_2 x_2$$
  
=  $c_1 \left( e^{\sqrt{\omega^2} t} \right) + c_2 \left( e^{\sqrt{\omega^2} t} \left( -\frac{\operatorname{csgn}(\omega) e^{-2 \operatorname{csgn}(\omega) \omega t}}{2\omega} \right) \right)$ 

Simplifying the solution  $x = c_1 e^{\sqrt{\omega^2 t}} - \frac{c_2 \operatorname{csgn}(\omega)e^{-\operatorname{csgn}(\omega)\omega t}}{2\omega}$  to  $x = c_1 e^{\sqrt{\omega^2 t}} - \frac{c_2 e^{-\omega t}}{2\omega}$ Summary

The solution(s) found are the following

$$x = c_1 \mathrm{e}^{\sqrt{\omega^2}t} - \frac{c_2 \mathrm{e}^{-\omega t}}{2\omega} \tag{1}$$

Verification of solutions

$$x = c_1 \mathrm{e}^{\sqrt{\omega^2} t} - \frac{c_2 \mathrm{e}^{-\omega t}}{2\omega}$$

Verified OK.

#### 1.1.4 Maple step by step solution

Let's solve  
$$x'' - \omega^2 x = 0$$

- Highest derivative means the order of the ODE is 2 x''
- Characteristic polynomial of ODE  $-\omega^2 + r^2 = 0$
- Factor the characteristic polynomial  $-(\omega - r) (\omega + r) = 0$
- Roots of the characteristic polynomial  $r = (\omega, -\omega)$
- 1st solution of the ODE  $x_1(t) = e^{\omega t}$

• 2nd solution of the ODE 
$$x_2(t) = e^{-\omega t}$$

• General solution of the ODE

 $x = c_1 x_1(t) + c_2 x_2(t)$ 

• Substitute in solutions

$$x = c_1 \mathrm{e}^{\omega t} + c_2 \mathrm{e}^{-\omega t}$$

#### Maple trace

`Methods for second order ODEs: --- Trying classification methods --trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`</pre>

Solution by Maple Time used: 0.0 (sec). Leaf size: 18

dsolve(diff(x(t),t\$2)-omega^2\*x(t)=0,x(t), singsol=all)

$$x(t) = c_1 \mathrm{e}^{\omega t} + c_2 \mathrm{e}^{-\omega t}$$

Solution by Mathematica Time used: 0.013 (sec). Leaf size: 23

DSolve[x''[t]-\[Omega]^2\*x[t]==0,x[t],t,IncludeSingularSolutions -> True]

 $x(t) \rightarrow c_1 e^{t\omega} + c_2 e^{-t\omega}$ 

#### 1.2 problem 10.2.5

Internal problem ID [5046]

 $Internal\,file\,name\,[\texttt{OUTPUT/4539}\_\texttt{Sunday}\_\texttt{June}\_\texttt{05}\_\texttt{2022}\_\texttt{03}\_\texttt{00}\_\texttt{33}\_\texttt{PM}\_\texttt{16585502}/\texttt{index}.\texttt{tex}]$ 

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.2, ODEs with constant Coefficients. page 307
Problem number: 10.2.5.
ODE order: 3.
ODE degree: 1.

The type(s) of ODE detected by this program : "higher\_order\_linear\_constant\_coefficients\_ODE"

Maple gives the following as the ode type

```
[[_3rd_order, _missing_x]]
```

$$x''' - x'' + x' - x = 0$$

The characteristic equation is

$$\lambda^3 - \lambda^2 + \lambda - 1 = 0$$

The roots of the above equation are

$$\lambda_1 = 1$$
  
 $\lambda_2 = i$   
 $\lambda_3 = -i$ 

Therefore the homogeneous solution is

$$x_h(t) = c_1 e^t + e^{-it} c_2 + e^{it} c_3$$

The fundamental set of solutions for the homogeneous solution are the following

$$x_1 = e^t$$
$$x_2 = e^{-it}$$
$$x_3 = e^{it}$$

 $\frac{Summary}{The solution(s)}$  found are the following

$$x = c_1 e^t + e^{-it} c_2 + e^{it} c_3 \tag{1}$$

Verification of solutions

$$x = c_1 \mathrm{e}^t + \mathrm{e}^{-it} c_2 + \mathrm{e}^{it} c_3$$

Verified OK.

#### 1.2.1 Maple step by step solution

Let's solve

x''' - x'' + x' - x = 0

- Highest derivative means the order of the ODE is 3 x'''
- $\Box$  Convert linear ODE into a system of first order ODEs
  - Define new variable  $x_1(t)$  $x_1(t) = x$
  - Define new variable  $x_2(t)$

$$x_2(t) = x^t$$

• Define new variable  $x_3(t)$ 

$$x_3(t) = x^{\prime\prime}$$

• Isolate for  $x'_3(t)$  using original ODE

 $x_3'(t) = x_3(t) - x_2(t) + x_1(t)$ 

Convert linear ODE into a system of first order ODEs

$$[x_2(t) = x_1'(t), x_3(t) = x_2'(t), x_3'(t) = x_3(t) - x_2(t) + x_1(t)]$$

• Define vector

$$ec{x}(t) = \left[egin{array}{c} x_1(t) \ x_2(t) \ x_3(t) \end{array}
ight]$$

• System to solve

$$\vec{x}'(t) = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & -1 & 1 \end{bmatrix} \cdot \vec{x}(t)$$

• Define the coefficient matrix

$$A = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & -1 & 1 \end{bmatrix}$$

• Rewrite the system as

$$\vec{x}'(t) = A \cdot \vec{x}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[ \left[ 1, \left[ \begin{array}{c} 1\\1\\1\\1 \end{array} \right] \right], \left[ -I, \left[ \begin{array}{c} -1\\I\\1\\1 \end{array} \right] \right], \left[ I, \left[ \begin{array}{c} -1\\-I\\1\\1 \end{array} \right] \right] \right]$$

• Consider eigenpair

$$\left[ \begin{matrix} 1 \\ 1 \\ 1 \\ 1 \end{matrix} \right]$$

• Solution to homogeneous system from eigenpair

$$\vec{x}_1 = \mathbf{e}^t \cdot \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

• Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\begin{bmatrix} -I \\ -I \\ I \\ 1 \end{bmatrix}$$

• Solution from eigenpair

$$e^{-It} \cdot \begin{bmatrix} -1 \\ I \\ 1 \end{bmatrix}$$

• Use Euler identity to write solution in terms of sin and cos

$$(\cos(t) - \operatorname{I}\sin(t)) \cdot \begin{bmatrix} -1 \\ \mathrm{I} \\ 1 \end{bmatrix}$$

• Simplify expression

$$-\cos(t) + I\sin(t)$$
$$I(\cos(t) - I\sin(t))$$
$$\cos(t) - I\sin(t)$$

• Both real and imaginary parts are solutions to the homogeneous system

$$\begin{bmatrix} \vec{x}_2(t) = \begin{bmatrix} -\cos(t) \\ \sin(t) \\ \cos(t) \end{bmatrix}, \vec{x}_3(t) = \begin{bmatrix} \sin(t) \\ \cos(t) \\ -\sin(t) \end{bmatrix} \end{bmatrix}$$

• General solution to the system of ODEs

$$\vec{x} = c_1 \vec{x}_1 + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

• Substitute solutions into the general solution

$$\vec{x} = c_1 e^t \cdot \begin{bmatrix} 1\\ 1\\ 1 \end{bmatrix} + \begin{bmatrix} -c_2 \cos(t) + c_3 \sin(t) \\ c_2 \sin(t) + c_3 \cos(t) \\ c_2 \cos(t) - c_3 \sin(t) \end{bmatrix}$$

• First component of the vector is the solution to the ODE  $x = c_1 e^t + c_3 \sin(t) - c_2 \cos(t)$ 

#### Maple trace

`Methods for third order ODEs: --- Trying classification methods --trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`</pre>

Solution by Maple Time used: 0.015 (sec). Leaf size: 17

dsolve(diff(x(t),t3)-diff(x(t),t2)+diff(x(t),t)-x(t)=0,x(t), singsol=all)

 $x(t) = c_1 e^t + c_2 \sin(t) + c_3 \cos(t)$ 

Solution by Mathematica Time used: 0.003 (sec). Leaf size: 22

DSolve[x'''[t]-x''[t]+x'[t]-x[t]==0,x[t],t,IncludeSingularSolutions -> True]

 $x(t) \rightarrow c_3 e^t + c_1 \cos(t) + c_2 \sin(t)$ 

# 1.3 problem 10.2.8 part(1)

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	Solving as second order linear constant coeff ode

Internal problem ID [5047] Internal file name [OUTPUT/4540\_Sunday\_June\_05\_2022\_03\_00\_33\_PM\_66076883/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.2, ODEs with constant Coefficients. page 307
Problem number: 10.2.8 part(1).
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second\_order\_linear\_constant\_coeff"

Maple gives the following as the ode type

[[\_2nd\_order, \_missing\_x]]

$$x'' + 42x' + x = 0$$

With initial conditions

$$[x(0) = 1, x'(0) = 0]$$

#### 1.3.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$x'' + p(t)x' + q(t)x = F$$

Where here

$$p(t) = 42$$
$$q(t) = 1$$
$$F = 0$$

Hence the ode is

$$x'' + 42x' + x = 0$$

The domain of p(t) = 42 is

$$\{-\infty < t < \infty\}$$

And the point  $t_0 = 0$  is inside this domain. The domain of q(t) = 1 is

$$\{-\infty < t < \infty\}$$

And the point  $t_0 = 0$  is also inside this domain. Hence solution exists and is unique.

#### 1.3.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ax''(t) + Bx'(t) + Cx(t) = 0$$

Where in the above A = 1, B = 42, C = 1. Let the solution be  $x = e^{\lambda t}$ . Substituting this into the ODE gives

$$\lambda^2 \mathrm{e}^{\lambda t} + 42\lambda \,\mathrm{e}^{\lambda t} + \mathrm{e}^{\lambda t} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by  $e^{\lambda t}$  gives

$$\lambda^2 + 42\lambda + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting A = 1, B = 42, C = 1 into the above gives

$$\lambda_{1,2} = \frac{-42}{(2)(1)} \pm \frac{1}{(2)(1)}\sqrt{42^2 - (4)(1)(1)}$$
$$= -21 \pm 2\sqrt{110}$$

Hence

$$\lambda_1 = -21 + 2\sqrt{110}$$
  
 $\lambda_2 = -21 - 2\sqrt{110}$ 

Which simplifies to

$$\lambda_1 = -21 + 2\sqrt{110}$$
$$\lambda_2 = -21 - 2\sqrt{110}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned} x &= c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t} \\ x &= c_1 e^{\left(-21 + 2\sqrt{110}\right)t} + c_2 e^{\left(-21 - 2\sqrt{110}\right)t} \end{aligned}$$

Or

$$x = c_1 e^{\left(-21 + 2\sqrt{110}\right)t} + c_2 e^{\left(-21 - 2\sqrt{110}\right)t}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$x = c_1 e^{\left(-21 + 2\sqrt{110}\right)t} + c_2 e^{\left(-21 - 2\sqrt{110}\right)t}$$
(1)

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting x = 1 and t = 0 in the above gives

$$1 = c_1 + c_2 \tag{1A}$$

Taking derivative of the solution gives

$$x' = c_1 \left(-21 + 2\sqrt{110}\right) e^{\left(-21 + 2\sqrt{110}\right)t} + c_2 \left(-21 - 2\sqrt{110}\right) e^{\left(-21 - 2\sqrt{110}\right)t}$$

substituting x' = 0 and t = 0 in the above gives

$$0 = (2c_1 - 2c_2)\sqrt{110} - 21c_1 - 21c_2$$
(2A)

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

$$c_1 = \frac{1}{2} + \frac{21\sqrt{110}}{440}$$
$$c_2 = \frac{1}{2} - \frac{21\sqrt{110}}{440}$$

Substituting these values back in above solution results in

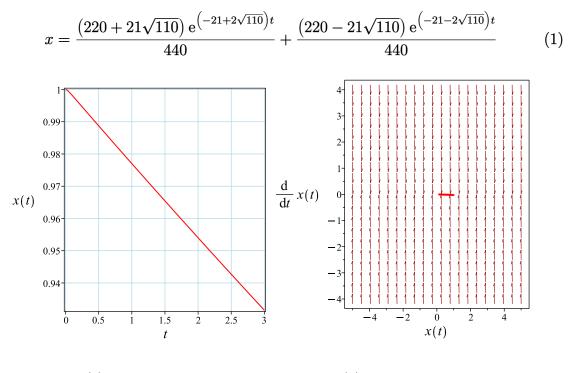
$$x = \frac{e^{\left(-21+2\sqrt{110}\right)t}}{2} + \frac{21e^{\left(-21+2\sqrt{110}\right)t}\sqrt{110}}{440} + \frac{e^{-\left(21+2\sqrt{110}\right)t}}{2} - \frac{21e^{-\left(21+2\sqrt{110}\right)t}\sqrt{110}}{440}$$

Which simplifies to

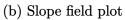
$$x = \frac{\left(220 + 21\sqrt{110}\right)e^{\left(-21 + 2\sqrt{110}\right)t}}{440} + \frac{\left(220 - 21\sqrt{110}\right)e^{\left(-21 - 2\sqrt{110}\right)t}}{440}$$

## Summary

The solution(s) found are the following



(a) Solution plot



Verification of solutions

$$x = \frac{\left(220 + 21\sqrt{110}\right)e^{\left(-21 + 2\sqrt{110}\right)t}}{440} + \frac{\left(220 - 21\sqrt{110}\right)e^{\left(-21 - 2\sqrt{110}\right)t}}{440}$$

Verified OK.

### 1.3.3 Solving using Kovacic algorithm

Writing the ode as

$$x'' + 42x' + x = 0 \tag{1}$$

$$Ax'' + Bx' + Cx = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$
  

$$B = 42$$
 (3)  

$$C = 1$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = x e^{\int rac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{440}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 440$$
  
 $t = 1$ 

Therefore eq. (4) becomes

$$z''(t) = 440z(t)$$
(7)

Equation (7) is now solved. After finding z(t) then x is found using the inverse transformation

$$x = z(t) e^{-\int rac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 4: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 0 - 0$$
$$= 0$$

There are no poles in r. Therefore the set of poles  $\Gamma$  is empty. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since r = 440 is not a function of t, then there is no need run Kovacic algorithm to obtain a solution for transformed ode z'' = rz as one solution is

$$z_1(t) = e^{-2t\sqrt{110}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in x is found from

$$x_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt}$$
$$= z_1 e^{-\int \frac{1}{2} \frac{42}{1} dt}$$

$$= z_1 e^{-21t}$$
$$= z_1 \left( e^{-21t} \right)$$

Which simplifies to

$$x_1 = \mathrm{e}^{\left(-21 - 2\sqrt{110}\right)t}$$

The second solution  $x_2$  to the original ode is found using reduction of order

$$x_2 = x_1 \int \frac{e^{\int -\frac{B}{A} dt}}{x_1^2} dt$$

Substituting gives

$$x_{2} = x_{1} \int \frac{e^{\int -\frac{42}{1} dt}}{(x_{1})^{2}} dt$$
$$= x_{1} \int \frac{e^{-42t}}{(x_{1})^{2}} dt$$
$$= x_{1} \left(\frac{\sqrt{110} e^{4t\sqrt{110}}}{440}\right)$$

Therefore the solution is

$$x = c_1 x_1 + c_2 x_2$$
  
=  $c_1 \left( e^{\left(-21 - 2\sqrt{110}\right)t} \right) + c_2 \left( e^{\left(-21 - 2\sqrt{110}\right)t} \left( \frac{\sqrt{110} e^{4t\sqrt{110}}}{440} \right) \right)$ 

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$x = c_1 e^{\left(-21 - 2\sqrt{110}\right)t} + \frac{c_2 e^{\left(-21 + 2\sqrt{110}\right)t}\sqrt{110}}{440}$$
(1)

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting x = 1 and t = 0 in the above gives

$$1 = c_1 + \frac{c_2\sqrt{110}}{440} \tag{1A}$$

Taking derivative of the solution gives

$$x' = c_1 \left(-21 - 2\sqrt{110}\right) e^{\left(-21 - 2\sqrt{110}\right)t} + \frac{c_2 \left(-21 + 2\sqrt{110}\right) e^{\left(-21 + 2\sqrt{110}\right)t} \sqrt{110}}{440}$$

substituting x' = 0 and t = 0 in the above gives

$$0 = \frac{(-880c_1 - 21c_2)\sqrt{110}}{440} - 21c_1 + \frac{c_2}{2}$$
(2A)

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

$$c_1 = \frac{1}{2} - \frac{21\sqrt{110}}{440}$$
$$c_2 = 21 + 2\sqrt{110}$$

Substituting these values back in above solution results in

$$x = \frac{e^{\left(-21+2\sqrt{110}\right)t}}{2} + \frac{21e^{\left(-21+2\sqrt{110}\right)t}\sqrt{110}}{440} + \frac{e^{-\left(21+2\sqrt{110}\right)t}}{2} - \frac{21e^{-\left(21+2\sqrt{110}\right)t}\sqrt{110}}{440}$$

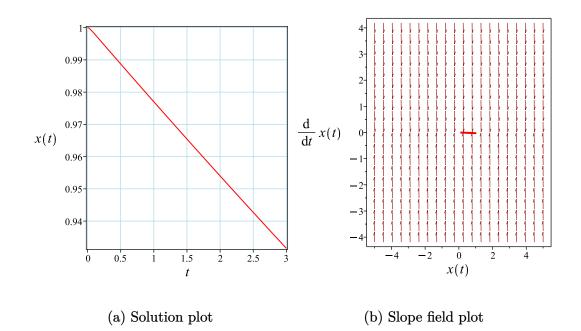
Which simplifies to

$$x = \frac{\left(220 + 21\sqrt{110}\right)e^{\left(-21 + 2\sqrt{110}\right)t}}{440} + \frac{\left(220 - 21\sqrt{110}\right)e^{\left(-21 - 2\sqrt{110}\right)t}}{440}$$

Summary

The solution(s) found are the following

$$x = \frac{\left(220 + 21\sqrt{110}\right)e^{\left(-21 + 2\sqrt{110}\right)t}}{440} + \frac{\left(220 - 21\sqrt{110}\right)e^{\left(-21 - 2\sqrt{110}\right)t}}{440} \tag{1}$$



Verification of solutions

$$x = \frac{\left(220 + 21\sqrt{110}\right)e^{\left(-21 + 2\sqrt{110}\right)t}}{440} + \frac{\left(220 - 21\sqrt{110}\right)e^{\left(-21 - 2\sqrt{110}\right)t}}{440}$$

Verified OK.

#### 1.3.4 Maple step by step solution

Let's solve  

$$\begin{bmatrix} x'' + 42x' + x = 0, x(0) = 1, x' \Big|_{\{t=0\}} = 0 \end{bmatrix}$$
Highest derivative means the order of the O

- Highest derivative means the order of the ODE is 2 x''
- Characteristic polynomial of ODE  $r^2 + 42r + 1 = 0$
- Use quadratic formula to solve for r $r = \frac{(-42) \pm \left(\sqrt{1760}\right)}{2}$
- Roots of the characteristic polynomial  $r = \left(-21 2\sqrt{110}, -21 + 2\sqrt{110}\right)$

- 1st solution of the ODE  $x_1(t) = e^{\left(-21 2\sqrt{110}\right)t}$
- 2nd solution of the ODE $x_2(t) = e^{\left(-21+2\sqrt{110}\right)t}$
- General solution of the ODE

$$x = c_1 x_1(t) + c_2 x_2(t)$$

• Substitute in solutions  $x = c_1 e^{\left(-21 - 2\sqrt{110}\right)t} + c_2 e^{\left(-21 + 2\sqrt{110}\right)t}$ 

 $\Box$  Check validity of solution  $x = c_1 e^{\left(-21 - 2\sqrt{110}\right)t} + c_2 e^{\left(-21 + 2\sqrt{110}\right)t}$ 

 $\circ \quad \text{Use initial condition } x(0) = 1$ 

 $1 = c_1 + c_2$ 

• Compute derivative of the solution

$$x' = c_1 \left(-21 - 2\sqrt{110}\right) e^{\left(-21 - 2\sqrt{110}\right)t} + c_2 \left(-21 + 2\sqrt{110}\right) e^{\left(-21 + 2\sqrt{110}\right)t}$$

• Use the initial condition 
$$x'\Big|_{\{t=0\}} = 0$$
  
 $0 = c_1(-21 - 2\sqrt{110}) + (-21 + 2\sqrt{110}) c_2$ 

• Solve for  $c_1$  and  $c_2$ 

$$\left\{c_1 = \frac{1}{2} - \frac{21\sqrt{110}}{440}, c_2 = \frac{1}{2} + \frac{21\sqrt{110}}{440}\right\}$$

• Substitute constant values into general solution and simplify

$$x = \frac{\left(220 + 21\sqrt{110}\right)e^{\left(-21 + 2\sqrt{110}\right)t}}{440} + \frac{\left(220 - 21\sqrt{110}\right)e^{\left(-21 - 2\sqrt{110}\right)t}}{440}$$

• Solution to the IVP

$$x = \frac{\left(220 + 21\sqrt{110}\right)e^{\left(-21 + 2\sqrt{110}\right)t}}{440} + \frac{\left(220 - 21\sqrt{110}\right)e^{\left(-21 - 2\sqrt{110}\right)t}}{440}$$

#### Maple trace

`Methods for second order ODEs: --- Trying classification methods --trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`</pre>

Solution by Maple Time used: 0.078 (sec). Leaf size: 43

dsolve([diff(x(t),t\$2)+42\*diff(x(t),t)+x(t)=0,x(0) = 1, D(x)(0) = 0],x(t), singsol=all)

$$x(t) = \frac{\left(220 + 21\sqrt{110}\right)e^{\left(-21 + 2\sqrt{110}\right)t}}{440} + \frac{\left(220 - 21\sqrt{110}\right)e^{\left(-21 - 2\sqrt{110}\right)t}}{440}$$

Solution by Mathematica Time used: 0.03 (sec). Leaf size: 53

DSolve[{x''[t]+42\*x'[t]+x[t]==0,{x[0]==1,x'[0]==0}},x[t],t,IncludeSingularSolutions -> True]

$$x(t) \to \frac{e^{-\left(\left(21+2\sqrt{110}\right)t\right)}\left(\left(881+84\sqrt{110}\right)e^{4\sqrt{110}t}-1\right)}{880+84\sqrt{110}}$$

# 1.4 problem 10.2.8 part(2)

Internal problem ID [5048]

 $Internal file name \left[ \texttt{OUTPUT/4541}\_\texttt{Sunday}\_\texttt{June}\_\texttt{05}\_\texttt{2022}\_\texttt{03}\_\texttt{00}\_\texttt{34}\_\texttt{PM}\_\texttt{44292396}/\texttt{index}.\texttt{tex} \right]$ 

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.2, ODEs with constant Coefficients. page 307
Problem number: 10.2.8 part(2).
ODE order: 4.
ODE degree: 1.

The type(s) of ODE detected by this program : "higher\_order\_linear\_constant\_coefficients\_ODE"

Maple gives the following as the ode type

```
[[_high_order, _missing_x]]
```

$$x'''' + x = 0$$

The characteristic equation is

$$\lambda^4 + 1 = 0$$

The roots of the above equation are

$$\lambda_{1} = \frac{\sqrt{2}}{2} + \frac{i\sqrt{2}}{2}$$
$$\lambda_{2} = -\frac{\sqrt{2}}{2} + \frac{i\sqrt{2}}{2}$$
$$\lambda_{3} = -\frac{\sqrt{2}}{2} - \frac{i\sqrt{2}}{2}$$
$$\lambda_{4} = \frac{\sqrt{2}}{2} - \frac{i\sqrt{2}}{2}$$

Therefore the homogeneous solution is

$$x_h(t) = e^{\left(-\frac{\sqrt{2}}{2} + \frac{i\sqrt{2}}{2}\right)t} c_1 + e^{\left(\frac{\sqrt{2}}{2} + \frac{i\sqrt{2}}{2}\right)t} c_2 + e^{\left(-\frac{\sqrt{2}}{2} - \frac{i\sqrt{2}}{2}\right)t} c_3 + e^{\left(\frac{\sqrt{2}}{2} - \frac{i\sqrt{2}}{2}\right)t} c_4$$

The fundamental set of solutions for the homogeneous solution are the following

$$\begin{aligned} x_1 &= \mathrm{e}^{\left(-\frac{\sqrt{2}}{2} + \frac{i\sqrt{2}}{2}\right)t} \\ x_2 &= \mathrm{e}^{\left(\frac{\sqrt{2}}{2} + \frac{i\sqrt{2}}{2}\right)t} \\ x_3 &= \mathrm{e}^{\left(-\frac{\sqrt{2}}{2} - \frac{i\sqrt{2}}{2}\right)t} \\ x_4 &= \mathrm{e}^{\left(\frac{\sqrt{2}}{2} - \frac{i\sqrt{2}}{2}\right)t} \end{aligned}$$

#### Summary

The solution(s) found are the following

$$x = e^{\left(-\frac{\sqrt{2}}{2} + \frac{i\sqrt{2}}{2}\right)t}c_1 + e^{\left(\frac{\sqrt{2}}{2} + \frac{i\sqrt{2}}{2}\right)t}c_2 + e^{\left(-\frac{\sqrt{2}}{2} - \frac{i\sqrt{2}}{2}\right)t}c_3 + e^{\left(\frac{\sqrt{2}}{2} - \frac{i\sqrt{2}}{2}\right)t}c_4$$
(1)

Verification of solutions

$$x = e^{\left(-\frac{\sqrt{2}}{2} + \frac{i\sqrt{2}}{2}\right)t}c_1 + e^{\left(\frac{\sqrt{2}}{2} + \frac{i\sqrt{2}}{2}\right)t}c_2 + e^{\left(-\frac{\sqrt{2}}{2} - \frac{i\sqrt{2}}{2}\right)t}c_3 + e^{\left(\frac{\sqrt{2}}{2} - \frac{i\sqrt{2}}{2}\right)t}c_4$$

Verified OK.

#### 1.4.1 Maple step by step solution

Let's solve

x'''' + x = 0

- Highest derivative means the order of the ODE is 4 x''''
- $\Box$  Convert linear ODE into a system of first order ODEs
  - Define new variable  $x_1(t)$  $x_1(t) = x$
  - Define new variable  $x_2(t)$

$$x_2(t) = x'$$

• Define new variable  $x_3(t)$  $x_3(t) = x''$ 

$$x_3(t) = x'$$

• Define new variable  $x_4(t)$  $x_4(t) = x'''$ 

$$x_4(t) = x^{\prime\prime}$$

 $\circ~$  Isolate for  $x_4'(t)~$  using original ODE  $x_4'(t)=-x_1(t)$ 

Convert linear ODE into a system of first order ODEs

$$[x_2(t) = x'_1(t), x_3(t) = x'_2(t), x_4(t) = x'_3(t), x'_4(t) = -x_1(t)]$$

• Define vector

$$ec{x}(t) = egin{bmatrix} x_1(t) \ x_2(t) \ x_3(t) \ x_4(t) \end{bmatrix}$$

• System to solve

$$\vec{x}'(t) = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -1 & 0 & 0 & 0 \end{bmatrix} \cdot \vec{x}(t)$$

• Define the coefficient matrix

$$A = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -1 & 0 & 0 & 0 \end{bmatrix}$$

• Rewrite the system as

$$\vec{x}'(t) = A \cdot \vec{x}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\begin{bmatrix} \begin{bmatrix} \frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{1\sqrt{2}}{2}\right)^{3}} \\ \frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{1\sqrt{2}}{2}\right)^{2}} \\ \frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{1\sqrt{2}}{2}\right)^{2}} \\ \frac{1}{-\frac{\sqrt{2}}{2}-\frac{1\sqrt{2}}{2}} \\ 1 \end{bmatrix} \end{bmatrix}, \begin{bmatrix} \frac{1}{\left(-\frac{\sqrt{2}}{2}+\frac{1\sqrt{2}}{2}\right)^{3}} \\ \frac{1}{\left(-\frac{\sqrt{2}}{2}+\frac{1\sqrt{2}}{2}\right)^{2}} \\ \frac{1}{-\frac{\sqrt{2}}{2}+\frac{1\sqrt{2}}{2}} \\ 1 \end{bmatrix} \end{bmatrix}, \begin{bmatrix} \frac{1}{\left(\frac{\sqrt{2}}{2}-\frac{1\sqrt{2}}{2}\right)^{3}} \\ \frac{1}{\left(\frac{\sqrt{2}}{2}-\frac{1\sqrt{2}}{2}\right)^{2}} \\ \frac{1}{\frac{\sqrt{2}}{2}-\frac{1\sqrt{2}}{2}} \\ 1 \end{bmatrix} \end{bmatrix}$$

• Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\begin{bmatrix} -\frac{\sqrt{2}}{2} - \frac{I\sqrt{2}}{2}, \begin{bmatrix} \frac{1}{\left(-\frac{\sqrt{2}}{2} - \frac{I\sqrt{2}}{2}\right)^3} \\ \frac{1}{\left(-\frac{\sqrt{2}}{2} - \frac{I\sqrt{2}}{2}\right)^2} \\ \frac{-\frac{1}{\sqrt{2}} - \frac{I\sqrt{2}}{2}}{1} \end{bmatrix}$$

• Solution from eigenpair

$$e^{\left(-\frac{\sqrt{2}}{2}-\frac{I\sqrt{2}}{2}\right)t} \cdot \begin{bmatrix} \frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{I\sqrt{2}}{2}\right)^3} \\ \frac{1}{\left(-\frac{\sqrt{2}}{2}-\frac{I\sqrt{2}}{2}\right)^2} \\ \frac{1}{-\frac{\sqrt{2}}{2}-\frac{I\sqrt{2}}{2}} \\ 1 \end{bmatrix}$$

• Use Euler identity to write solution in terms of sin and cos

$$\mathrm{e}^{-\frac{t\sqrt{2}}{2}} \cdot \left( \cos\left(\frac{t\sqrt{2}}{2}\right) - \mathrm{I}\sin\left(\frac{t\sqrt{2}}{2}\right) \right) \cdot \begin{bmatrix} \frac{1}{\left(-\frac{\sqrt{2}}{2} - \frac{1\sqrt{2}}{2}\right)^3} \\ \frac{1}{\left(-\frac{\sqrt{2}}{2} - \frac{1\sqrt{2}}{2}\right)^2} \\ \frac{1}{-\frac{\sqrt{2}}{2} - \frac{1\sqrt{2}}{2}} \\ 1 \end{bmatrix}$$

• Simplify expression

$$\mathbf{e}^{-\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right) - \mathrm{I}\sin\left(\frac{t\sqrt{2}}{2}\right)}{\left(-\frac{\sqrt{2}}{2} - \frac{\mathrm{I}\sqrt{2}}{2}\right)^3} \\ \frac{\cos\left(\frac{t\sqrt{2}}{2}\right) - \mathrm{I}\sin\left(\frac{t\sqrt{2}}{2}\right)}{\left(-\frac{\sqrt{2}}{2} - \frac{\mathrm{I}\sqrt{2}}{2}\right)^2} \\ \frac{\cos\left(\frac{t\sqrt{2}}{2}\right) - \mathrm{I}\sin\left(\frac{t\sqrt{2}}{2}\right)}{-\frac{\sqrt{2}}{2} - \frac{\mathrm{I}\sqrt{2}}{2}} \\ \cos\left(\frac{t\sqrt{2}}{2}\right) - \mathrm{I}\sin\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix}$$

• Both real and imaginary parts are solutions to the homogeneous system

$$\begin{bmatrix} \vec{x}_{1}(t) = e^{-\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\sin\left(\frac{t\sqrt{2}}{2}\right) \\ -\frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ \cos\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix}, \vec{x}_{2}(t) = e^{-\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} - \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\cos\left(\frac{t\sqrt{2}}{2}\right) \\ \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ \cos\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix}$$

• Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\begin{bmatrix} \frac{1}{\left(\frac{\sqrt{2}}{2} - \frac{\mathrm{I}\sqrt{2}}{2}\right)^3} \\ \frac{1}{\left(\frac{\sqrt{2}}{2} - \frac{\mathrm{I}\sqrt{2}}{2}\right)^2} \\ \frac{1}{\frac{\sqrt{2}}{2} - \frac{\mathrm{I}\sqrt{2}}{2}} \\ \frac{1}{\frac{\sqrt{2}}{2} - \frac{\mathrm{I}\sqrt{2}}{2}} \\ 1 \end{bmatrix}$$

• Solution from eigenpair

$$e^{\left(\frac{\sqrt{2}}{2} - \frac{I\sqrt{2}}{2}\right)t} \cdot \begin{bmatrix} \frac{1}{\left(\frac{\sqrt{2}}{2} - \frac{I\sqrt{2}}{2}\right)^3} \\ \frac{1}{\left(\frac{\sqrt{2}}{2} - \frac{I\sqrt{2}}{2}\right)^2} \\ \frac{1}{\frac{\sqrt{2}}{2} - \frac{I\sqrt{2}}{2}} \\ 1 \end{bmatrix}$$

• Use Euler identity to write solution in terms of sin and cos

$$e^{\frac{t\sqrt{2}}{2}} \cdot \left(\cos\left(\frac{t\sqrt{2}}{2}\right) - I\sin\left(\frac{t\sqrt{2}}{2}\right)\right) \cdot \begin{bmatrix} \frac{1}{\left(\frac{\sqrt{2}}{2} - \frac{1\sqrt{2}}{2}\right)^3} \\ \frac{1}{\left(\frac{\sqrt{2}}{2} - \frac{1\sqrt{2}}{2}\right)^2} \\ \frac{1}{\frac{\sqrt{2}}{2} - \frac{1\sqrt{2}}{2}} \\ 1 \end{bmatrix}$$

• Simplify expression

$$\mathbf{e}^{\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right) - \mathrm{I}\sin\left(\frac{t\sqrt{2}}{2}\right)}{\left(\frac{\sqrt{2}}{2} - \frac{\mathrm{I}\sqrt{2}}{2}\right)^3} \\ \frac{\cos\left(\frac{t\sqrt{2}}{2}\right) - \mathrm{I}\sin\left(\frac{t\sqrt{2}}{2}\right)}{\left(\frac{\sqrt{2}}{2} - \frac{\mathrm{I}\sqrt{2}}{2}\right)^2} \\ \frac{\cos\left(\frac{t\sqrt{2}}{2}\right) - \mathrm{I}\sin\left(\frac{t\sqrt{2}}{2}\right)}{\frac{\sqrt{2}}{2} - \frac{\mathrm{I}\sqrt{2}}{2}} \\ \cos\left(\frac{t\sqrt{2}}{2}\right) - \mathrm{I}\sin\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix}$$

• Both real and imaginary parts are solutions to the homogeneous system

$$\begin{bmatrix} \vec{x}_{3}(t) = e^{\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} -\frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ \sin\left(\frac{t\sqrt{2}}{2}\right) \\ \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ \cos\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix}, \vec{x}_{4}(t) = e^{\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ \cos\left(\frac{t\sqrt{2}}{2}\right) \\ \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} - \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\sin\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix} \end{bmatrix}$$

• General solution to the system of ODEs

$$\vec{x} = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t) + c_4 \vec{x}_4(t)$$

• Substitute solutions into the general solution

$$\vec{x} = c_1 e^{-\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\sin\left(\frac{t\sqrt{2}}{2}\right) \\ -\frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ \cos\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix} + c_2 e^{-\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} - \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\cos\left(\frac{t\sqrt{2}}{2}\right) \\ \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\sin\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix} + c_3 e^{\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} - \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} + \frac{\sin\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\sin\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix} + c_3 e^{\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\sin\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix} + c_3 e^{\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\sin\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix} \end{bmatrix} + c_3 e^{\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\sin\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix} + c_3 e^{\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\sin\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix} \end{bmatrix} + c_3 e^{\frac{t\sqrt{2}}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{t\sqrt{2}}{2}\right)\sqrt{2}}{2} \\ -\sin\left(\frac{t\sqrt{2}}{2}\right) \end{bmatrix} \end{bmatrix}$$

• First component of the vector is the solution to the ODE

$$x = \frac{\left(\left((c_1+c_2)e^{-\frac{t\sqrt{2}}{2}} - e^{\frac{t\sqrt{2}}{2}}(c_3-c_4)\right)\cos\left(\frac{t\sqrt{2}}{2}\right) + \sin\left(\frac{t\sqrt{2}}{2}\right)\left((c_1-c_2)e^{-\frac{t\sqrt{2}}{2}} + e^{\frac{t\sqrt{2}}{2}}(c_3+c_4)\right)\right)\sqrt{2}}{2}$$

#### Maple trace

`Methods for high order ODEs: --- Trying classification methods --trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`</pre>

Solution by Maple Time used: 0.015 (sec). Leaf size: 61

dsolve(diff(x(t),t\$4)+x(t)=0,x(t), singsol=all)

$$x(t) = \left(-c_1 \mathrm{e}^{-\frac{\sqrt{2}t}{2}} - c_2 \mathrm{e}^{\frac{\sqrt{2}t}{2}}\right) \sin\left(\frac{\sqrt{2}t}{2}\right) + \left(c_3 \mathrm{e}^{-\frac{\sqrt{2}t}{2}} + c_4 \mathrm{e}^{\frac{\sqrt{2}t}{2}}\right) \cos\left(\frac{\sqrt{2}t}{2}\right)$$

Solution by Mathematica Time used: 0.003 (sec). Leaf size: 65

DSolve[x'''[t]+x[t]==0,x[t],t,IncludeSingularSolutions -> True]

$$x(t) \to e^{-\frac{t}{\sqrt{2}}} \left( \left( c_1 e^{\sqrt{2}t} + c_2 \right) \cos\left(\frac{t}{\sqrt{2}}\right) + \left( c_4 e^{\sqrt{2}t} + c_3 \right) \sin\left(\frac{t}{\sqrt{2}}\right) \right)$$

# 1.5 problem 10.2.8 part(3)

1.5.1 Maple step by step solution  $\ldots \ldots \ldots \ldots \ldots \ldots \ldots 35$ 

Internal problem ID [5049]

 $Internal file name \left[ \texttt{OUTPUT/4542\_Sunday\_June\_05\_2022\_03\_00\_35\_PM\_42275176/index.tex} \right]$ 

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.2, ODEs with constant Coefficients. page 307
Problem number: 10.2.8 part(3).
ODE order: 3.
ODE degree: 1.

The type(s) of ODE detected by this program : "higher\_order\_linear\_constant\_coefficients\_ODE"

Maple gives the following as the ode type

```
[[_3rd_order, _missing_x]]
```

$$x''' - 3x'' - 9x' - 5x = 0$$

The characteristic equation is

$$\lambda^3 - 3\lambda^2 - 9\lambda - 5 = 0$$

The roots of the above equation are

$$\lambda_1 = 5$$
$$\lambda_2 = -1$$
$$\lambda_3 = -1$$

Therefore the homogeneous solution is

$$x_h(t) = e^{-t}c_1 + t e^{-t}c_2 + e^{5t}c_3$$

The fundamental set of solutions for the homogeneous solution are the following

$$x_1 = e^{-t}$$
$$x_2 = t e^{-t}$$
$$x_3 = e^{5t}$$

 $\frac{\text{Summary}}{\text{The solution(s) found are the following}}$ 

$$x = e^{-t}c_1 + t e^{-t}c_2 + e^{5t}c_3$$
(1)

Verification of solutions

$$x = e^{-t}c_1 + t e^{-t}c_2 + e^{5t}c_3$$

Verified OK.

#### 1.5.1 Maple step by step solution

Let's solve

x''' - 3x'' - 9x' - 5x = 0

- Highest derivative means the order of the ODE is 3 x'''
- $\Box$  Convert linear ODE into a system of first order ODEs
  - Define new variable  $x_1(t)$  $x_1(t) = x$

$$x_1(\iota) = x$$

• Define new variable  $x_2(t)$ 

$$x_2(t) = x'$$

 $\circ$  Define new variable  $x_3(t)$  $x_3(t) = x''$ 

• Isolate for 
$$x'_3(t)$$
 using original ODE

 $x_3'(t) = 3x_3(t) + 9x_2(t) + 5x_1(t)$ 

Convert linear ODE into a system of first order ODEs

$$[x_2(t) = x'_1(t), x_3(t) = x'_2(t), x'_3(t) = 3x_3(t) + 9x_2(t) + 5x_1(t)]$$

• Define vector

$$ec{x}(t) = \left[egin{array}{c} x_1(t) \ x_2(t) \ x_3(t) \end{array}
ight]$$

• System to solve

$$\vec{x}'(t) = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 5 & 9 & 3 \end{bmatrix} \cdot \vec{x}(t)$$

• Define the coefficient matrix

$$A = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 5 & 9 & 3 \end{bmatrix}$$

• Rewrite the system as

$$\vec{x}'(t) = A \cdot \vec{x}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[ \left[ -1, \left[ \begin{array}{c} 1\\ -1\\ 1 \end{array} \right] \right], \left[ -1, \left[ \begin{array}{c} 0\\ 0\\ 0 \end{array} \right] \right], \left[ 5, \left[ \begin{array}{c} \frac{1}{25}\\ \frac{1}{5}\\ 1 \end{array} \right] \right] \right]$$

• Consider eigenpair, with eigenvalue of algebraic multiplicity 2

$$\begin{bmatrix} -1, \begin{bmatrix} 1\\ -1\\ 1 \end{bmatrix}$$

• First solution from eigenvalue -1

$$\vec{x}_1(t) = e^{-t} \cdot \begin{bmatrix} 1\\ -1\\ 1 \end{bmatrix}$$

- Form of the 2nd homogeneous solution where  $\vec{p}$  is to be solved for,  $\lambda = -1$  is the eigenvalue, a  $\vec{x}_2(t) = e^{\lambda t} \left( t \vec{v} + \vec{p} \right)$
- Note that the t multiplying  $\vec{v}$  makes this solution linearly independent to the 1st solution obta

• Substitute 
$$\vec{x}_2(t)$$
 into the homogeneous system  
 $\lambda e^{\lambda t} \left( t \vec{v} + \vec{p} \right) + e^{\lambda t} \vec{v} = (e^{\lambda t} A) \cdot \left( t \vec{v} + \vec{p} \right)$ 

- Use the fact that  $\vec{v}$  is an eigenvector of A $\lambda e^{\lambda t} \left( t \vec{v} + \vec{p} \right) + e^{\lambda t} \vec{v} = e^{\lambda t} \left( \lambda t \vec{v} + A \cdot \vec{p} \right)$
- Simplify equation  $\lambda \vec{p} + \vec{v} = A \cdot \vec{p}$

• Make use of the identity matrix I  

$$(\lambda \cdot I) \cdot \overrightarrow{p} + \overrightarrow{v} = A \cdot \overrightarrow{p}$$

- Condition  $\vec{p}$  must meet for  $\vec{x}_2(t)$  to be a solution to the homogeneous system  $(A - \lambda \cdot I) \cdot \vec{p} = \vec{v}$
- Choose  $\overrightarrow{p}$  to use in the second solution to the homogeneous system from eigenvalue -1

$$\left( \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 5 & 9 & 3 \end{bmatrix} - (-1) \cdot \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \cdot \vec{p} = \begin{bmatrix} 1 \\ -1 \\ 1 \end{bmatrix}$$

• Choice of  $\overrightarrow{p}$ 

$$\vec{p} = \begin{bmatrix} 1\\0\\0 \end{bmatrix}$$

• Second solution from eigenvalue -1

$$\vec{x}_{2}(t) = \mathrm{e}^{-t} \cdot \left( t \cdot \begin{bmatrix} 1 \\ -1 \\ 1 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} \right)$$

• Consider eigenpair

$$\begin{bmatrix} 5, \begin{bmatrix} \frac{1}{25} \\ \frac{1}{5} \\ 1 \end{bmatrix}$$

• Solution to homogeneous system from eigenpair

$$\vec{x}_3 = \mathrm{e}^{5t} \cdot \begin{bmatrix} \frac{1}{25} \\ \frac{1}{5} \\ 1 \end{bmatrix}$$

• General solution to the system of ODEs

$$\vec{x} = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + c_3 \vec{x}_3$$

• Substitute solutions into the general solution

$$\vec{x} = \mathbf{e}^{-t}c_1 \cdot \begin{bmatrix} 1\\ -1\\ 1 \end{bmatrix} + c_2 \mathbf{e}^{-t} \cdot \left( t \cdot \begin{bmatrix} 1\\ -1\\ 1 \end{bmatrix} + \begin{bmatrix} 1\\ 0\\ 0 \end{bmatrix} \right) + \mathbf{e}^{5t}c_3 \cdot \begin{vmatrix} \frac{1}{25}\\ \frac{1}{5}\\ 1 \end{vmatrix}$$

• First component of the vector is the solution to the ODE  $x = ((t+1)c_2 + c_1)e^{-t} + \frac{e^{5t}c_3}{25}$ 

#### Maple trace

`Methods for third order ODEs: --- Trying classification methods --trying a quadrature checking if the LODE has constant coefficients <- constant coefficients successful`</pre>

Solution by Maple Time used: 0.0 (sec). Leaf size: 21

dsolve(diff(x(t),t\$3)-3\*diff(x(t),t\$2)-9\*diff(x(t),t)-5\*x(t)=0,x(t), singsol=all)

$$x(t) = (c_3 t + c_2) e^{-t} + c_1 e^{5t}$$

Solution by Mathematica Time used: 0.003 (sec). Leaf size: 26

DSolve[x'''[t]-3\*x''[t]-9\*x'[t]-5\*x[t]==0,x[t],t,IncludeSingularSolutions -> True]

$$x(t) \rightarrow e^{-t} (c_2 t + c_3 e^{6t} + c_1)$$

#### 1.6 problem 10.2.10

1.6.1	Solving as second order linear constant coeff ode	39
1.6.2	Solving using Kovacic algorithm	42
1.6.3	Maple step by step solution	47

Internal problem ID [5050]

Internal file name [OUTPUT/4543\_Sunday\_June\_05\_2022\_03\_00\_36\_PM\_76937184/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.2, ODEs with constant Coefficients. page 307
Problem number: 10.2.10.
ODE order: 2.

**ODE degree**: 1.

The type(s) of ODE detected by this program : "kovacic", "second\_order\_linear\_constant\_coeff"

Maple gives the following as the ode type

```
[[_2nd_order, _linear, _nonhomogeneous]]
```

 $x'' + 2\gamma x' + \omega_0 x = F \cos\left(\omega t\right)$ 

#### 1.6.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ax''(t) + Bx'(t) + Cx(t) = f(t)$$

Where  $A = 1, B = 2\gamma, C = \omega_0, f(t) = F \cos(\omega t)$ . Let the solution be

$$x = x_h + x_p$$

Where  $x_h$  is the solution to the homogeneous ODE Ax''(t) + Bx'(t) + Cx(t) = 0, and  $x_p$  is a particular solution to the non-homogeneous ODE Ax''(t) + Bx'(t) + Cx(t) = f(t).  $x_h$  is the solution to

$$x'' + 2\gamma x' + \omega_0 x = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ax''(t) + Bx'(t) + Cx(t) = 0$$

Where in the above  $A = 1, B = 2\gamma, C = \omega_0$ . Let the solution be  $x = e^{\lambda t}$ . Substituting this into the ODE gives

$$\lambda^2 \mathrm{e}^{\lambda t} + 2\gamma \lambda \,\mathrm{e}^{\lambda t} + \omega_0 \mathrm{e}^{\lambda t} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by  $e^{\lambda t}$  gives

$$2\gamma\lambda + \lambda^2 + \omega_0 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting  $A = 1, B = 2\gamma, C = \omega_0$  into the above gives

$$\lambda_{1,2} = \frac{-2\gamma}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{2\gamma^2 - (4)(1)(\omega_0)} \\ = -\gamma \pm \sqrt{\gamma^2 - \omega_0}$$

Hence

$$\lambda_1 = -\gamma + \sqrt{\gamma^2 - \omega_0}$$
$$\lambda_2 = -\gamma - \sqrt{\gamma^2 - \omega_0}$$

Which simplifies to

$$\lambda_1 = -\gamma + \sqrt{\gamma^2 - \omega_0}$$
$$\lambda_2 = -\gamma - \sqrt{\gamma^2 - \omega_0}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned} x &= c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t} \\ x &= c_1 e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} + c_2 e^{\left(-\gamma - \sqrt{\gamma^2 - \omega_0}\right)t} \end{aligned}$$

Or

$$x = c_1 e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} + c_2 e^{\left(-\gamma - \sqrt{\gamma^2 - \omega_0}\right)t}$$

Therefore the homogeneous solution  $x_h$  is

$$x_h = c_1 \mathrm{e}^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} + c_2 \mathrm{e}^{\left(-\gamma - \sqrt{\gamma^2 - \omega_0}\right)t}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

 $F\cos\left(\omega t\right)$ 

Shows that the corresponding undetermined set of the basis functions (UC\_set) for the trial solution is

$$\left[\left\{\cos\left(\omega t\right),\sin\left(\omega t\right)\right\}\right]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ e^{\left(-\gamma - \sqrt{\gamma^2 - \omega_0}\right)t}, e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} \right\}$$

Since there is no duplication between the basis function in the UC\_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC\_set.

$$x_p = A_1 \cos\left(\omega t\right) + A_2 \sin\left(\omega t\right)$$

The unknowns  $\{A_1, A_2\}$  are found by substituting the above trial solution  $x_p$  into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-A_1\omega^2\cos(\omega t) - A_2\omega^2\sin(\omega t) + 2\gamma(-A_1\omega\sin(\omega t) + A_2\omega\cos(\omega t)) + \omega_0(A_1\cos(\omega t) + A_2\sin(\omega t)) = F\cos(\omega t)$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_{1} = -\frac{(\omega^{2} - \omega_{0}) F}{\omega^{4} + (4\gamma^{2} - 2\omega_{0}) \omega^{2} + \omega_{0}^{2}}, A_{2} = \frac{2\gamma F\omega}{\omega^{4} + (4\gamma^{2} - 2\omega_{0}) \omega^{2} + \omega_{0}^{2}}\right]$$

Substituting the above back in the above trial solution  $x_p$ , gives the particular solution

$$x_p = -\frac{\left(\omega^2 - \omega_0\right)F\cos\left(\omega t\right)}{\omega^4 + \left(4\gamma^2 - 2\omega_0\right)\omega^2 + \omega_0^2} + \frac{2\gamma F\omega\sin\left(\omega t\right)}{\omega^4 + \left(4\gamma^2 - 2\omega_0\right)\omega^2 + \omega_0^2}$$

Therefore the general solution is

$$\begin{aligned} x &= x_h + x_p \\ &= \left( c_1 e^{\left( -\gamma + \sqrt{\gamma^2 - \omega_0} \right)t} + c_2 e^{\left( -\gamma - \sqrt{\gamma^2 - \omega_0} \right)t} \right) \\ &+ \left( -\frac{\left( \omega^2 - \omega_0 \right)F\cos\left(\omega t \right)}{\omega^4 + \left( 4\gamma^2 - 2\omega_0 \right)\omega^2 + \omega_0^2} + \frac{2\gamma F\omega\sin\left(\omega t \right)}{\omega^4 + \left( 4\gamma^2 - 2\omega_0 \right)\omega^2 + \omega_0^2} \right) \end{aligned}$$

Which simplifies to

$$x = c_1 e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} + c_2 e^{-\left(\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} - \frac{(\omega^2 - \omega_0) F \cos(\omega t)}{\omega^4 + (4\gamma^2 - 2\omega_0) \omega^2 + \omega_0^2} + \frac{2\gamma F \omega \sin(\omega t)}{\omega^4 + (4\gamma^2 - 2\omega_0) \omega^2 + \omega_0^2}$$

Summary

The solution(s) found are the following

$$x = c_1 e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} + c_2 e^{-\left(\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} - \frac{(\omega^2 - \omega_0) F \cos(\omega t)}{\omega^4 + (4\gamma^2 - 2\omega_0)\omega^2 + \omega_0^2} + \frac{2\gamma F \omega \sin(\omega t)}{\omega^4 + (4\gamma^2 - 2\omega_0)\omega^2 + \omega_0^2}$$
(1)

Verification of solutions

$$x = c_1 e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} + c_2 e^{-\left(\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} - \frac{(\omega^2 - \omega_0) F \cos(\omega t)}{\omega^4 + (4\gamma^2 - 2\omega_0) \omega^2 + \omega_0^2} + \frac{2\gamma F \omega \sin(\omega t)}{\omega^4 + (4\gamma^2 - 2\omega_0) \omega^2 + \omega_0^2}$$

Verified OK.

## 1.6.2 Solving using Kovacic algorithm

Writing the ode as

$$x'' + 2\gamma x' + \omega_0 x = 0 \tag{1}$$

$$Ax'' + Bx' + Cx = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$
  

$$B = 2\gamma$$

$$C = \omega_0$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(t) = x e^{\int rac{B}{2A} \, dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{\gamma^2 - \omega_0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = \gamma^2 - \omega_0$$
$$t = 1$$

Therefore eq. (4) becomes

$$z''(t) = \left(\gamma^2 - \omega_0\right) z(t) \tag{7}$$

Equation (7) is now solved. After finding z(t) then x is found using the inverse transformation

$$x = z(t) e^{-\int rac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

#### Table 8: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 0 - 0$$
$$= 0$$

There are no poles in r. Therefore the set of poles  $\Gamma$  is empty. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since  $r = \gamma^2 - \omega_0$  is not a function of t, then there is no need run Kovacic algorithm to obtain a solution for transformed ode z'' = rz as one solution is

$$z_1(t) = \mathrm{e}^{t\sqrt{\gamma^2 - \omega_0}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in x is found from

$$\begin{aligned} x_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{2\gamma}{1} dt} \\ &= z_1 e^{-t\gamma} \\ &= z_1 \left( e^{-t\gamma} \right) \end{aligned}$$

Which simplifies to

$$x_1 = \mathrm{e}^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}
ight)t}$$

The second solution  $x_2$  to the original ode is found using reduction of order

$$x_2 = x_1 \int \frac{e^{\int -\frac{B}{A}\,dt}}{x_1^2}\,dt$$

Substituting gives

$$x_2 = x_1 \int \frac{e^{\int -\frac{2\gamma}{1} dt}}{(x_1)^2} dt$$
$$= x_1 \int \frac{e^{-2t\gamma}}{(x_1)^2} dt$$
$$= x_1 \left( -\frac{e^{-2t\sqrt{\gamma^2 - \omega_0}}}{2\sqrt{\gamma^2 - \omega_0}} \right)$$

Therefore the solution is

$$\begin{aligned} x &= c_1 x_1 + c_2 x_2 \\ &= c_1 \left( e^{\left( -\gamma + \sqrt{\gamma^2 - \omega_0} \right) t} \right) + c_2 \left( e^{\left( -\gamma + \sqrt{\gamma^2 - \omega_0} \right) t} \left( -\frac{e^{-2t\sqrt{\gamma^2 - \omega_0}}}{2\sqrt{\gamma^2 - \omega_0}} \right) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$x = x_h + x_p$$

Where  $x_h$  is the solution to the homogeneous ODE Ax''(t) + Bx'(t) + Cx(t) = 0, and  $x_p$  is a particular solution to the nonhomogeneous ODE Ax''(t) + Bx'(t) + Cx(t) = f(t).  $x_h$  is the solution to

$$x'' + 2\gamma x' + \omega_0 x = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$x_h = c_1 \mathrm{e}^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} - \frac{c_2 \mathrm{e}^{-\left(\gamma + \sqrt{\gamma^2 - \omega_0}\right)t}}{2\sqrt{\gamma^2 - \omega_0}}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$F\cos\left(\omega t\right)$$

Shows that the corresponding undetermined set of the basis functions (UC\_set) for the trial solution is

$$\left[\left\{\cos\left(\omega t\right),\sin\left(\omega t\right)\right\}\right]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{-\frac{\mathrm{e}^{-\left(\gamma+\sqrt{\gamma^{2}-\omega_{0}}\right)t}}{2\sqrt{\gamma^{2}-\omega_{0}}},\mathrm{e}^{\left(-\gamma+\sqrt{\gamma^{2}-\omega_{0}}\right)t}\right\}$$

Since there is no duplication between the basis function in the UC\_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC\_set.

$$x_p = A_1 \cos\left(\omega t\right) + A_2 \sin\left(\omega t\right)$$

The unknowns  $\{A_1, A_2\}$  are found by substituting the above trial solution  $x_p$  into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-A_1\omega^2\cos(\omega t) - A_2\omega^2\sin(\omega t) + 2\gamma(-A_1\omega\sin(\omega t) + A_2\omega\cos(\omega t)) + \omega_0(A_1\cos(\omega t) + A_2\sin(\omega t)) = F\cos(\omega t)$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_{1} = -\frac{(\omega^{2} - \omega_{0}) F}{\omega^{4} + (4\gamma^{2} - 2\omega_{0}) \omega^{2} + \omega_{0}^{2}}, A_{2} = \frac{2\gamma F\omega}{\omega^{4} + (4\gamma^{2} - 2\omega_{0}) \omega^{2} + \omega_{0}^{2}}\right]$$

Substituting the above back in the above trial solution  $x_p$ , gives the particular solution

$$x_p = -\frac{\left(\omega^2 - \omega_0\right)F\cos\left(\omega t\right)}{\omega^4 + \left(4\gamma^2 - 2\omega_0\right)\omega^2 + \omega_0^2} + \frac{2\gamma F\omega\sin\left(\omega t\right)}{\omega^4 + \left(4\gamma^2 - 2\omega_0\right)\omega^2 + \omega_0^2}$$

Therefore the general solution is

$$\begin{aligned} x &= x_h + x_p \\ &= \left( c_1 e^{\left( -\gamma + \sqrt{\gamma^2 - \omega_0} \right)t} - \frac{c_2 e^{-\left( \gamma + \sqrt{\gamma^2 - \omega_0} \right)t}}{2\sqrt{\gamma^2 - \omega_0}} \right) \\ &+ \left( -\frac{\left( \omega^2 - \omega_0 \right) F \cos\left(\omega t \right)}{\omega^4 + \left( 4\gamma^2 - 2\omega_0 \right) \omega^2 + \omega_0^2} + \frac{2\gamma F \omega \sin\left(\omega t \right)}{\omega^4 + \left( 4\gamma^2 - 2\omega_0 \right) \omega^2 + \omega_0^2} \right) \end{aligned}$$

#### Summary

The solution(s) found are the following

$$x = c_1 e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} - \frac{c_2 e^{-\left(\gamma + \sqrt{\gamma^2 - \omega_0}\right)t}}{2\sqrt{\gamma^2 - \omega_0}} - \frac{(\omega^2 - \omega_0) F \cos(\omega t)}{\omega^4 + (4\gamma^2 - 2\omega_0)\omega^2 + \omega_0^2} + \frac{2\gamma F \omega \sin(\omega t)}{\omega^4 + (4\gamma^2 - 2\omega_0)\omega^2 + \omega_0^2}$$
(1)

Verification of solutions

$$x = c_1 e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} - \frac{c_2 e^{-\left(\gamma + \sqrt{\gamma^2 - \omega_0}\right)t}}{2\sqrt{\gamma^2 - \omega_0}} - \frac{(\omega^2 - \omega_0) F \cos(\omega t)}{\omega^4 + (4\gamma^2 - 2\omega_0) \omega^2 + \omega_0^2} + \frac{2\gamma F \omega \sin(\omega t)}{\omega^4 + (4\gamma^2 - 2\omega_0) \omega^2 + \omega_0^2}$$

Verified OK.

#### 1.6.3 Maple step by step solution

Let's solve

 $x'' + 2\gamma x' + \omega_0 x = F \cos\left(\omega t\right)$ 

- Highest derivative means the order of the ODE is 2 x''
- Characteristic polynomial of homogeneous ODE

$$2\gamma r + r^2 + \omega_0 = 0$$

• Use quadratic formula to solve for r

$$r=rac{(-2\gamma)\pm\left(\sqrt{4\gamma^2-4\omega_0}
ight)}{2}$$

- Roots of the characteristic polynomial  $r = (-\gamma - \sqrt{\gamma^2 - \omega_0}, -\gamma + \sqrt{\gamma^2 - \omega_0})$
- 1st solution of the homogeneous ODE

$$x_1(t) = \mathrm{e}^{\left(-\gamma - \sqrt{\gamma^2 - \omega_0}
ight)t}$$

- 2nd solution of the homogeneous ODE $x_2(t) = \mathrm{e}^{\left(-\gamma + \sqrt{\gamma^2 \omega_0}
  ight)t}$
- General solution of the ODE  $x = c_1 x_1(t) + c_2 x_2(t) + x_p(t)$
- Substitute in solutions of the homogeneous ODE

$$x = c_1 \mathrm{e}^{\left(-\gamma - \sqrt{\gamma^2 - \omega_0}\right)t} + c_2 \mathrm{e}^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} + x_p(t)$$

- $\Box$  Find a particular solution  $x_p(t)$  of the ODE
  - Use variation of parameters to find  $x_p$  here f(t) is the forcing function

$$x_p(t) = -x_1(t) \left( \int \frac{x_2(t)f(t)}{W(x_1(t), x_2(t))} dt \right) + x_2(t) \left( \int \frac{x_1(t)f(t)}{W(x_1(t), x_2(t))} dt \right), f(t) = F \cos\left(\omega t\right) \right]$$

 $\circ$   $\,$  Wronskian of solutions of the homogeneous equation

$$W(x_1(t), x_2(t)) = \begin{bmatrix} e^{\left(-\gamma - \sqrt{\gamma^2 - \omega_0}\right)t} & e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} \\ \left(-\gamma - \sqrt{\gamma^2 - \omega_0}\right)e^{\left(-\gamma - \sqrt{\gamma^2 - \omega_0}\right)t} & \left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} \end{bmatrix}$$

• Compute Wronskian

 $W(x_1(t), x_2(t)) = 2\sqrt{\gamma^2 - \omega_0} e^{-2t\gamma}$ 

• Substitute functions into equation for  $x_p(t)$ 

$$x_p(t) = \frac{F\left(e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t}\left(\int \cos(\omega t)e^{\left(\gamma - \sqrt{\gamma^2 - \omega_0}\right)t}dt\right) - e^{-\left(\gamma + \sqrt{\gamma^2 - \omega_0}\right)t}\left(\int \cos(\omega t)e^{\left(\gamma + \sqrt{\gamma^2 - \omega_0}\right)t}dt\right)\right)}{2\sqrt{\gamma^2 - \omega_0}}$$

• Compute integrals

 $x_p(t) = rac{F((-\omega^2+\omega_0)\cos(\omega t)+2\gamma\omega\sin(\omega t))}{\omega^4+2(2\gamma^2-\omega_0)\omega^2+\omega_0^2}$ 

• Substitute particular solution into general solution to ODE

$$x = c_1 e^{\left(-\gamma - \sqrt{\gamma^2 - \omega_0}\right)t} + c_2 e^{\left(-\gamma + \sqrt{\gamma^2 - \omega_0}\right)t} + \frac{F\left(\left(-\omega^2 + \omega_0\right)\cos(\omega t) + 2\gamma\omega\sin(\omega t)\right)}{\omega^4 + 2(2\gamma^2 - \omega_0)\omega^2 + \omega_0^2}$$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
checking if the LODE has constant coefficients
<- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`</pre>
```

Solution by Maple Time used: 0.015 (sec). Leaf size: 111

dsolve(diff(x(t),t\$2)+2\*gamma\*diff(x(t),t)+omega\_\_0\*x(t)=F\*cos(omega\*t),x(t), singsol=all)

x(t)

$$=\frac{-F(\omega^{2}-\omega_{0})\cos(\omega t)+2F\sin(\omega t)\gamma\omega+4\left(\frac{\omega^{4}}{4}+\left(\gamma^{2}-\frac{\omega_{0}}{2}\right)\omega^{2}+\frac{\omega_{0}^{2}}{4}\right)\left(e^{-\left(\gamma+\sqrt{\gamma^{2}-\omega_{0}}\right)t}c_{1}+e^{\left(-\gamma+\sqrt{\gamma^{2}-\omega_{0}}\right)t}c_{2}+\frac{\omega^{2}}{4}+\left(4\gamma^{2}-2\omega_{0}\right)\omega^{2}+\omega_{0}^{2}\right)}{\omega^{4}+\left(4\gamma^{2}-2\omega_{0}\right)\omega^{2}+\omega_{0}^{2}}$$

# ✓ Solution by Mathematica Time used: 0.509 (sec). Leaf size: 108

DSolve[x''[t]+2\*\[Gamma]\*x'[t]+Subscript[\[Omega],0]\*x[t]==F\*Cos[\[Omega]\*t],x[t],t,IncludeS

$$x(t) \rightarrow \frac{F(\omega(2\gamma\sin(t\omega) - \omega\cos(t\omega)) + \omega_0\cos(t\omega))}{4\gamma^2\omega^2 + \omega^4 - 2\omega_0\omega^2 + \omega_0^2} + c_1e^{-t\left(\sqrt{\gamma^2 - \omega_0} + \gamma\right)} + c_2e^{t\left(\sqrt{\gamma^2 - \omega_0} - \gamma\right)}$$

## 1.7 problem 10.2.11 (i)

1.7.1	Existence and uniqueness analysis	50
1.7.2	Solving as second order linear constant coeff ode	51
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Internal problem ID [5051] Internal file name [OUTPUT/4544\_Sunday\_June\_05\_2022\_03\_00\_37\_PM\_98654921/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.2, ODEs with constant Coefficients. page 307
Problem number: 10.2.11 (i).
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second\_order\_linear\_constant\_coeff"

Maple gives the following as the ode type

[[\_2nd\_order, \_with\_linear\_symmetries]]

$$y'' - y' - 2y = \mathrm{e}^{2x}$$

With initial conditions

$$[y(0) = 1, y'(0) = 0]$$

#### 1.7.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = -1$$
$$q(x) = -2$$
$$F = e^{2x}$$

Hence the ode is

$$y'' - y' - 2y = e^{2x}$$

The domain of p(x) = -1 is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is inside this domain. The domain of q(x) = -2 is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is also inside this domain. The domain of  $F = e^{2x}$  is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is also inside this domain. Hence solution exists and is unique.

#### 1.7.2 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = f(x)$$

Where  $A = 1, B = -1, C = -2, f(x) = e^{2x}$ . Let the solution be

$$y = y_h + y_p$$

Where  $y_h$  is the solution to the homogeneous ODE Ay''(x) + By'(x) + Cy(x) = 0, and  $y_p$  is a particular solution to the non-homogeneous ODE Ay''(x) + By'(x) + Cy(x) = f(x).  $y_h$  is the solution to

$$y'' - y' - 2y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above A = 1, B = -1, C = -2. Let the solution be  $y = e^{\lambda x}$ . Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - \lambda e^{\lambda x} - 2 e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by  $e^{\lambda x}$  gives

$$\lambda^2 - \lambda - 2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting A = 1, B = -1, C = -2 into the above gives

$$\lambda_{1,2} = \frac{1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-1^2 - (4)(1)(-2)}$$
$$= \frac{1}{2} \pm \frac{3}{2}$$

Hence

$$\lambda_1 = \frac{1}{2} + \frac{3}{2}$$
$$\lambda_2 = \frac{1}{2} - \frac{3}{2}$$

Which simplifies to

$$\lambda_1 = 2$$
$$\lambda_2 = -1$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$
  
 $y = c_1 e^{(2)x} + c_2 e^{(-1)x}$ 

Or

$$y = c_1 e^{2x} + c_2 e^{-x}$$

Therefore the homogeneous solution  $y_h$  is

$$y_h = c_1 e^{2x} + c_2 e^{-x}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

 $e^{2x}$ 

Shows that the corresponding undetermined set of the basis functions (UC\_set) for the trial solution is

 $[\{\mathrm{e}^{2x}\}]$ 

While the set of the basis functions for the homogeneous solution found earlier is

$$\{e^{-x}, e^{2x}\}$$

Since  $e^{2x}$  is duplicated in the UC\_set, then this basis is multiplied by extra x. The UC\_set becomes

$$[\{e^{2x}x\}]$$

Since there was duplication between the basis functions in the UC\_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC\_set.

$$y_p = A_1 e^{2x} x$$

The unknowns  $\{A_1\}$  are found by substituting the above trial solution  $y_p$  into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

 $3A_1e^{2x} = e^{2x}$ 

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = \frac{1}{3}\right]$$

Substituting the above back in the above trial solution  $y_p$ , gives the particular solution

$$y_p = \frac{\mathrm{e}^{2x}x}{3}$$

Therefore the general solution is

$$y = y_h + y_p$$
$$= \left(c_1 e^{2x} + c_2 e^{-x}\right) + \left(\frac{e^{2x} x}{3}\right)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{2x} + c_2 e^{-x} + \frac{e^{2x} x}{3}$$
(1)

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting y = 1 and x = 0 in the above gives

$$1 = c_1 + c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = 2c_1 e^{2x} - c_2 e^{-x} + \frac{2 e^{2x} x}{3} + \frac{e^{2x}}{3}$$

substituting y' = 0 and x = 0 in the above gives

$$0 = 2c_1 - c_2 + \frac{1}{3} \tag{2A}$$

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

$$c_1 = \frac{2}{9}$$
$$c_2 = \frac{7}{9}$$

Substituting these values back in above solution results in

$$y = \frac{2 e^{2x}}{9} + \frac{7 e^{-x}}{9} + \frac{e^{2x} x}{3}$$

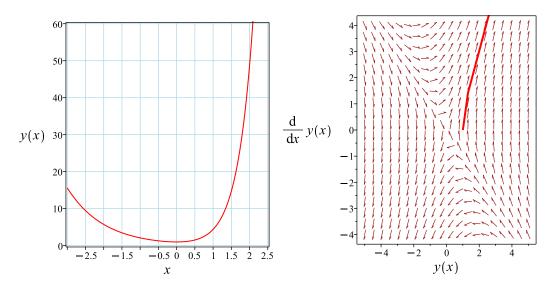
Which simplifies to

$$y = \frac{(2+3x)e^{2x}}{9} + \frac{7e^{-x}}{9}$$

Summary

The solution(s) found are the following

$$y = \frac{(2+3x)e^{2x}}{9} + \frac{7e^{-x}}{9}$$
(1)



(a) Solution plot

(b) Slope field plot

Verification of solutions

$$y = \frac{(2+3x)e^{2x}}{9} + \frac{7e^{-x}}{9}$$

Verified OK.

### 1.7.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' - y' - 2y = 0 (1)$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$
  

$$B = -1$$

$$C = -2$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{9}{4} \tag{6}$$

Comparing the above to (5) shows that

s = 9t = 4

Therefore eq. (4) becomes

$$z''(x) = \frac{9z(x)}{4}$$
(7)

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 10: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 0 - 0$$
$$= 0$$

There are no poles in r. Therefore the set of poles  $\Gamma$  is empty. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since  $r = \frac{9}{4}$  is not a function of x, then there is no need run Kovacic algorithm to obtain a solution for transformed ode z'' = rz as one solution is

$$z_1(x) = \mathrm{e}^{-\frac{3x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_{1} = z_{1}e^{\int -\frac{1}{2}\frac{B}{A}dx}$$
  
=  $z_{1}e^{-\int \frac{1}{2}\frac{-1}{1}dx}$   
=  $z_{1}e^{\frac{x}{2}}$   
=  $z_{1}(e^{\frac{x}{2}})$ 

Which simplifies to

$$y_1 = e^{-x}$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} \, dx}}{y_1^2} \, dx$$

Substituting gives

$$y_2 = y_1 \int \frac{e^{\int -\frac{-1}{1} dx}}{(y_1)^2} dx$$
$$= y_1 \int \frac{e^x}{(y_1)^2} dx$$
$$= y_1 \left(\frac{e^{3x}}{3}\right)$$

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
  
=  $c_1 (e^{-x}) + c_2 \left( e^{-x} \left( \frac{e^{3x}}{3} \right) \right)$ 

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where  $y_h$  is the solution to the homogeneous ODE Ay''(x) + By'(x) + Cy(x) = 0, and  $y_p$  is a particular solution to the nonhomogeneous ODE Ay''(x) + By'(x) + Cy(x) = f(x).  $y_h$  is the solution to

$$y'' - y' - 2y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 \mathrm{e}^{-x} + \frac{c_2 \mathrm{e}^{2x}}{3}$$

The particular solution  $y_p$  can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where  $u_1, u_2$  to be determined, and  $y_1, y_2$  are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = e^{-x}$$
$$y_2 = \frac{e^{2x}}{3}$$

In the Variation of parameters  $u_1, u_2$  are found using

$$u_1 = -\int \frac{y_2 f(x)}{a W(x)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(x)}{a W(x)} \tag{3}$$

Where W(x) is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by  $W = \begin{vmatrix} y_1 & y_2 \\ y'_1 & y'_2 \end{vmatrix}$ . Hence

 $W = \begin{vmatrix} e^{-x} & \frac{e^{2x}}{3} \\ \frac{d}{dx}(e^{-x}) & \frac{d}{dx}\left(\frac{e^{2x}}{3}\right) \end{vmatrix}$ 

Which gives

 $W = egin{bmatrix} {
m e}^{-x} & {
m e}^{2x} \ {
m -e}^{-x} & {
m 2}{
m e}^{2x} \ {
m -e}^{2x} \ {
m 3} \end{bmatrix}$ 

 $W = \left(\mathrm{e}^{-x}
ight)\left(rac{2\,\mathrm{e}^{2x}}{3}
ight) - \left(rac{\mathrm{e}^{2x}}{3}
ight)\left(-\mathrm{e}^{-x}
ight)$ 

Therefore

Which simplifies to

Which simplifies to

Therefore Eq. (2) becomes

Which simplifies to

Hence

And Eq. (3) becomes

$$\int e^{-x}e^{2x}$$

$$u_2 = \int \frac{\mathrm{e}^{-x} \mathrm{e}^{2x}}{\mathrm{e}^x} \, dx$$

$$W = e^x$$

 $W = e^{-x}e^{2x}$ 

$$u_1 = -\int rac{\mathrm{e}^{4x}}{\mathrm{e}^x} \, dx$$

$$u_1 = -\int \frac{\mathrm{e}^{3x}}{3} dx$$

$$u_1 = -\frac{\mathrm{e}^{3x}}{9}$$

Which simplifies to

$$u_2 = \int 1 dx$$

Hence

$$u_2 = x$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = -\frac{\mathrm{e}^{-x}\mathrm{e}^{3x}}{9} + \frac{\mathrm{e}^{2x}x}{3}$$

Which simplifies to

$$y_p(x) = \frac{(3x-1)e^{2x}}{9}$$

Therefore the general solution is

$$y = y_h + y_p$$
  
=  $\left(c_1 e^{-x} + \frac{c_2 e^{2x}}{3}\right) + \left(\frac{(3x-1)e^{2x}}{9}\right)$ 

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-x} + \frac{c_2 e^{2x}}{3} + \frac{(3x-1)e^{2x}}{9}$$
(1)

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting y = 1 and x = 0 in the above gives

$$1 = c_1 + \frac{c_2}{3} - \frac{1}{9} \tag{1A}$$

Taking derivative of the solution gives

$$y' = -c_1 e^{-x} + \frac{2c_2 e^{2x}}{3} + \frac{e^{2x}}{3} + \frac{2(3x-1)e^{2x}}{9}$$

substituting y' = 0 and x = 0 in the above gives

$$0 = -c_1 + \frac{2c_2}{3} + \frac{1}{9} \tag{2A}$$

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

$$c_1 = \frac{7}{9}$$
$$c_2 = 1$$

Substituting these values back in above solution results in

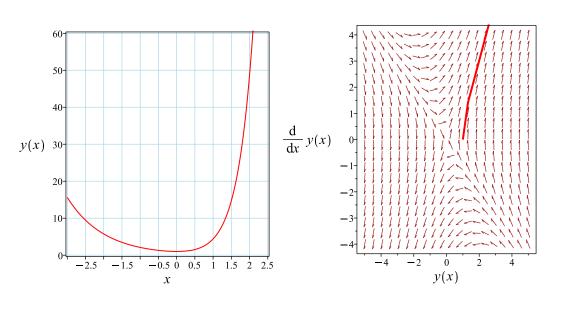
$$y = \frac{2e^{2x}}{9} + \frac{7e^{-x}}{9} + \frac{e^{2x}x}{3}$$

Which simplifies to

$$y = \frac{(2+3x)e^{2x}}{9} + \frac{7e^{-x}}{9}$$

 $\frac{Summary}{The solution(s) found are the following}$ 

$$y = \frac{(2+3x)e^{2x}}{9} + \frac{7e^{-x}}{9}$$
(1)



(a) Solution plot

(b) Slope field plot

Verification of solutions

$$y = \frac{(2+3x)e^{2x}}{9} + \frac{7e^{-x}}{9}$$

Verified OK.

#### 1.7.4 Maple step by step solution

Let's solve

$$\left[y'' - y' - 2y = e^{2x}, y(0) = 1, y'\Big|_{\{x=0\}} = 0\right]$$

- Highest derivative means the order of the ODE is 2 y''
- Characteristic polynomial of homogeneous ODE

$$r^2 - r - 2 = 0$$

- Factor the characteristic polynomial (r+1)(r-2) = 0
- Roots of the characteristic polynomial

$$r = (-1, 2)$$

- 1st solution of the homogeneous ODE  $y_1(x) = e^{-x}$
- 2nd solution of the homogeneous ODE

$$y_2(x) = e^{2x}$$

• General solution of the ODE

 $y = c_1 y_1(x) + c_2 y_2(x) + y_p(x)$ 

• Substitute in solutions of the homogeneous ODE

 $y = c_1 e^{-x} + c_2 e^{2x} + y_p(x)$ 

- $\Box$  Find a particular solution  $y_p(x)$  of the ODE
  - Use variation of parameters to find  $y_p$  here f(x) is the forcing function

$$\int y_p(x) = -y_1(x) \left(\int rac{y_2(x)f(x)}{W(y_1(x),y_2(x))} dx
ight) + y_2(x) \left(\int rac{y_1(x)f(x)}{W(y_1(x),y_2(x))} dx
ight), f(x) = \mathrm{e}^{2x}$$

 $\circ$  Wronskian of solutions of the homogeneous equation

$$W(y_1(x), y_2(x)) = \begin{bmatrix} e^{-x} & e^{2x} \\ -e^{-x} & 2e^{2x} \end{bmatrix}$$

• Compute Wronskian

$$W(y_1(x), y_2(x)) = 3 \operatorname{e}^x$$

• Substitute functions into equation for  $y_p(x)$ 

$$y_p(x) = -rac{\mathrm{e}^{-x}(\int \mathrm{e}^{3x} dx)}{3} + rac{\mathrm{e}^{2x}(\int 1 dx)}{3}$$

• Compute integrals

$$y_p(x) = \frac{(3x-1)\mathrm{e}^{2x}}{9}$$

- Substitute particular solution into general solution to ODE  $y = c_1 e^{-x} + c_2 e^{2x} + \frac{(3x-1)e^{2x}}{9}$
- $\Box \qquad \text{Check validity of solution } y = c_1 e^{-x} + c_2 e^{2x} + \frac{(3x-1)e^{2x}}{9}$ 
  - Use initial condition y(0) = 1

$$1 = c_1 + c_2 - \frac{1}{9}$$

• Compute derivative of the solution

$$y' = -c_1 e^{-x} + 2c_2 e^{2x} + \frac{e^{2x}}{3} + \frac{2(3x-1)e^{2x}}{9}$$

 $\circ \quad \text{Use the initial condition } y'\Big|_{\{x=0\}} = 0$ 

$$0 = -c_1 + 2c_2 + \frac{1}{9}$$

• Solve for  $c_1$  and  $c_2$ 

$$\left\{c_1 = \frac{7}{9}, c_2 = \frac{1}{3}\right\}$$

• Substitute constant values into general solution and simplify

$$y = \frac{(2+3x)e^{2x}}{9} + \frac{7e^{-x}}{9}$$

• Solution to the IVP $y = \frac{(2+3x)e^{2x}}{9} + \frac{7e^{-x}}{9}$ 

#### Maple trace

`Methods for second order ODEs: --- Trying classification methods --trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE checking if the LODE has constant coefficients <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`</pre> Solution by Maple Time used: 0.016 (sec). Leaf size: 22

dsolve([diff(y(x),x\$2)-diff(y(x),x)-2\*y(x)=exp(2\*x),y(0) = 1, D(y)(0) = 0],y(x), singsol=all

$$y(x) = \frac{(3x+2)e^{2x}}{9} + \frac{7e^{-x}}{9}$$

✓ Solution by Mathematica

Time used: 0.024 (sec). Leaf size: 27

DSolve[{y''[x]-y'[x]-2\*y[x]==Exp[2\*x],{y[0]==1,y'[0]==0}},y[x],x,IncludeSingularSolutions ->

$$y(x) \to \frac{1}{9}e^{-x} (e^{3x}(3x+2)+7)$$

## 1.8 problem 10.2.11 (ii)

1		Existence and uniqueness analysis	66
1		Solving as second order linear constant coeff ode	66
1	8.3	Solving as linear second order ode solved by an integrating factor	
		ode	70
1		Solving using Kovacic algorithm	72
1	8.5	Maple step by step solution	77
Internal pr	roblem	ID [5052]	
Internal file	e name	[OUTPUT/4545_Sunday_June_05_2022_03_00_38_PM_77665157/index.	tex]
Book: Ba	isic Tra	aining in Mathematics. By R. Shankar. Plenum Press. NY. 1995	
Section:	Chapte	er 10, Differential equations. Section 10.2, ODEs with constant Coefficient	nts.
page 307			
Problem	num	<b>ber</b> : 10.2.11 (ii).	
ODE ord	<b>ler</b> : 2.		
ODE deg	gree: 1	1.	

The type(s) of ODE detected by this program : "kovacic", "second\_order\_linear\_constant\_coeff", "linear\_second\_order\_ode\_solved\_by\_an\_integrating\_factor"

Maple gives the following as the ode type

[[\_2nd\_order, \_linear, \_nonhomogeneous]]

$$y'' - 2y' + y = 2\cos(x)$$

With initial conditions

$$[y(0) = 1, y'(0) = 0]$$

#### **1.8.1** Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = -2$$
$$q(x) = 1$$
$$F = 2\cos(x)$$

Hence the ode is

$$y'' - 2y' + y = 2\cos\left(x\right)$$

The domain of p(x) = -2 is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is inside this domain. The domain of q(x) = 1 is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is also inside this domain. The domain of  $F = 2\cos(x)$  is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is also inside this domain. Hence solution exists and is unique.

#### 1.8.2 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = f(x)$$

Where  $A = 1, B = -2, C = 1, f(x) = 2\cos(x)$ . Let the solution be

$$y = y_h + y_p$$

Where  $y_h$  is the solution to the homogeneous ODE Ay''(x) + By'(x) + Cy(x) = 0, and  $y_p$  is a particular solution to the non-homogeneous ODE Ay''(x) + By'(x) + Cy(x) = f(x).  $y_h$  is the solution to

$$y'' - 2y' + y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above A = 1, B = -2, C = 1. Let the solution be  $y = e^{\lambda x}$ . Substituting this into the ODE gives

$$\lambda^2 \mathrm{e}^{\lambda x} - 2\lambda \,\mathrm{e}^{\lambda x} + \mathrm{e}^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by  $e^{\lambda x}$  gives

$$\lambda^2 - 2\lambda + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting A = 1, B = -2, C = 1 into the above gives

$$\lambda_{1,2} = \frac{2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(-2)^2 - (4)(1)(1)}$$
  
= 1

Hence this is the case of a double root  $\lambda_{1,2} = -1$ . Therefore the solution is

$$y = c_1 \mathrm{e}^x + c_2 \mathrm{e}^x x \tag{1}$$

Therefore the homogeneous solution  $y_h$  is

$$y_h = c_1 \mathrm{e}^x + c_2 x \, \mathrm{e}^x$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

```
2\cos(x)
```

Shows that the corresponding undetermined set of the basis functions (UC\_set) for the trial solution is

$$\left[\left\{\cos\left(x\right),\sin\left(x\right)\right\}\right]$$

While the set of the basis functions for the homogeneous solution found earlier is

 $\{e^x x, e^x\}$ 

Since there is no duplication between the basis function in the UC\_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC\_set.

$$y_p = A_1 \cos\left(x\right) + A_2 \sin\left(x\right)$$

The unknowns  $\{A_1, A_2\}$  are found by substituting the above trial solution  $y_p$  into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$2A_1 \sin(x) - 2A_2 \cos(x) = 2\cos(x)$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = 0, A_2 = -1]$$

Substituting the above back in the above trial solution  $y_p$ , gives the particular solution

$$y_p = -\sin\left(x\right)$$

Therefore the general solution is

$$y = y_h + y_p$$
  
=  $(c_1 e^x + c_2 x e^x) + (-\sin(x))$ 

Which simplifies to

$$y = e^{x}(c_{2}x + c_{1}) - \sin(x)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = e^{x}(c_{2}x + c_{1}) - \sin(x)$$
(1)

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting y = 1 and x = 0 in the above gives

$$1 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = e^x(c_2x + c_1) + c_2e^x - \cos(x)$$

substituting y' = 0 and x = 0 in the above gives

$$0 = -1 + c_1 + c_2 \tag{2A}$$

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

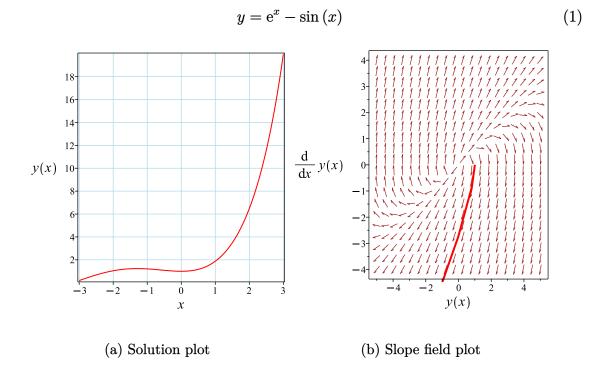
$$c_1 = 1$$
  
 $c_2 = 0$ 

Substituting these values back in above solution results in

$$y = e^x - \sin\left(x\right)$$

Summary

The solution(s) found are the following



Verification of solutions

$$y = e^x - \sin\left(x\right)$$

Verified OK.

# 1.8.3 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(x)y' + \frac{(p(x)^2 + p'(x))y}{2} = f(x)$$

Where p(x) = -2. Therefore, there is an integrating factor given by

$$M(x) = e^{\frac{1}{2} \int p \, dx}$$
$$= e^{\int -2 \, dx}$$
$$= e^{-x}$$

Multiplying both sides of the ODE by the integrating factor M(x) makes the left side of the ODE a complete differential

$$(M(x)y)'' = 2e^{-x}\cos(x)$$
  
 $(e^{-x}y)'' = 2e^{-x}\cos(x)$ 

Integrating once gives

$$(e^{-x}y)' = -e^{-x}(-\sin(x) + \cos(x)) + c_1$$

Integrating again gives

$$(e^{-x}y) = c_1 x - e^{-x} \sin(x) + c_2$$

Hence the solution is

$$y = \frac{c_1 x - e^{-x} \sin(x) + c_2}{e^{-x}}$$

Or

$$y = c_1 x \operatorname{e}^x + c_2 \operatorname{e}^x - \sin\left(x\right)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 x e^x + c_2 e^x - \sin(x)$$
(1)

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting y = 1 and x = 0 in the above gives

$$1 = c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = c_1 e^x + c_1 x e^x + c_2 e^x - \cos(x)$$

substituting y' = 0 and x = 0 in the above gives

$$0 = -1 + c_1 + c_2 \tag{2A}$$

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

$$c_1 = 0$$
  
 $c_2 = 1$ 

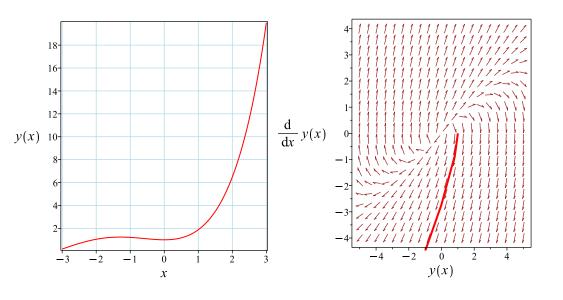
Substituting these values back in above solution results in

$$y = e^x - \sin\left(x\right)$$

Summary

The solution(s) found are the following

$$y = e^x - \sin\left(x\right) \tag{1}$$



(a) Solution plot

(b) Slope field plot

Verification of solutions

$$y = e^x - \sin\left(x\right)$$

Verified OK.

# 1.8.4 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 2y' + y = 0 (1)$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$
  

$$B = -2$$

$$C = 1$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int rac{B}{2A} \, dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 0$$
  
 $t = 1$ 

Therefore eq. (4) becomes

$$z''(x) = 0 \tag{7}$$

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 12: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 0 - -\infty$$
$$= \infty$$

There are no poles in r. Therefore the set of poles  $\Gamma$  is empty. Since there is no odd order pole larger than 2 and the order at  $\infty$  is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since r = 0 is not a function of x, then there is no need run Kovacic algorithm to obtain a solution for transformed ode z'' = rz as one solution is

$$z_1(x) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$
$$= z_1 e^{-\int \frac{1}{2} -\frac{2}{1} dx}$$

$$= z_1 e^x$$
$$= z_1 (e^x)$$

Which simplifies to

 $y_1 = e^x$ 

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2 = y_1 \int rac{e^{\int -rac{B}{A}\,dx}}{y_1^2}\,dx$$

Substituting gives

$$y_{2} = y_{1} \int \frac{e^{\int -\frac{-2}{1} dx}}{(y_{1})^{2}} dx$$
$$= y_{1} \int \frac{e^{2x}}{(y_{1})^{2}} dx$$
$$= y_{1}(x)$$

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
  
=  $c_1(e^x) + c_2(e^x(x))$ 

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where  $y_h$  is the solution to the homogeneous ODE Ay''(x) + By'(x) + Cy(x) = 0, and  $y_p$  is a particular solution to the nonhomogeneous ODE Ay''(x) + By'(x) + Cy(x) = f(x).  $y_h$  is the solution to

$$y'' - 2y' + y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 \mathrm{e}^x + c_2 x \, \mathrm{e}^x$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$2\cos(x)$$

Shows that the corresponding undetermined set of the basis functions (UC\_set) for the trial solution is

$$\left[\left\{\cos\left(x\right),\sin\left(x\right)\right\}\right]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\{\mathrm{e}^x x,\mathrm{e}^x\}$$

Since there is no duplication between the basis function in the UC\_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC\_set.

$$y_p = A_1 \cos\left(x\right) + A_2 \sin\left(x\right)$$

The unknowns  $\{A_1, A_2\}$  are found by substituting the above trial solution  $y_p$  into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$2A_{1}\sin(x) - 2A_{2}\cos(x) = 2\cos(x)$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = 0, A_2 = -1]$$

Substituting the above back in the above trial solution  $y_p$ , gives the particular solution

$$y_p = -\sin\left(x\right)$$

Therefore the general solution is

$$y = y_h + y_p$$
  
=  $(c_1 e^x + c_2 x e^x) + (-\sin(x))$ 

Which simplifies to

$$y = e^x(c_2x + c_1) - \sin\left(x\right)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = e^{x}(c_{2}x + c_{1}) - \sin(x)$$
(1)

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting y = 1 and x = 0 in the above gives

$$1 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = e^{x}(c_{2}x + c_{1}) + c_{2}e^{x} - \cos(x)$$

substituting y' = 0 and x = 0 in the above gives

$$0 = -1 + c_1 + c_2 \tag{2A}$$

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

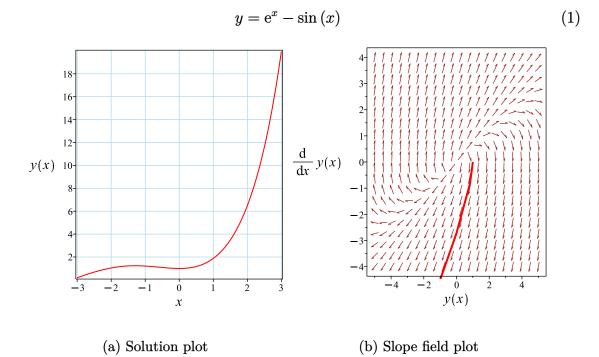
$$c_1 = 1$$
  
 $c_2 = 0$ 

Substituting these values back in above solution results in

$$y = e^x - \sin\left(x\right)$$

Summary

The solution(s) found are the following



Verification of solutions

$$y = e^x - \sin\left(x\right)$$

Verified OK.

## 1.8.5 Maple step by step solution

Let's solve

$$\left[y'' - 2y' + y = 2\cos(x), y(0) = 1, y'\Big|_{\{x=0\}} = 0\right]$$

- Highest derivative means the order of the ODE is 2 y''
- Characteristic polynomial of homogeneous ODE  $r^2 2r + 1 = 0$
- Factor the characteristic polynomial

$$(r-1)^2 = 0$$

- Root of the characteristic polynomial r = 1
- 1st solution of the homogeneous ODE

$$y_1(x) = e^x$$

- Repeated root, multiply  $y_1(x)$  by x to ensure linear independence  $y_2(x) = e^x x$
- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x) + y_p(x)$$

• Substitute in solutions of the homogeneous ODE

$$y = c_1 \mathrm{e}^x + c_2 x \, \mathrm{e}^x + y_p(x)$$

 $\Box$  Find a particular solution  $y_p(x)$  of the ODE

• Use variation of parameters to find  $y_p$  here f(x) is the forcing function

$$\left[y_p(x) = -y_1(x) \left(\int \frac{y_2(x)f(x)}{W(y_1(x),y_2(x))} dx\right) + y_2(x) \left(\int \frac{y_1(x)f(x)}{W(y_1(x),y_2(x))} dx\right), f(x) = 2\cos\left(x\right)\right]$$

• Wronskian of solutions of the homogeneous equation

$$W(y_1(x)\,,y_2(x))=\left[egin{array}{cc} \mathrm{e}^x & \mathrm{e}^x x\ \mathrm{e}^x & \mathrm{e}^x x+\mathrm{e}^x\end{array}
ight.$$

- $\circ$  Compute Wronskian  $W(y_1(x)\,,y_2(x))={\rm e}^{2x}$
- Substitute functions into equation for  $y_p(x)$  $y_p(x) = 2 e^x \left( -\left( \int \cos(x) x e^{-x} dx \right) + x \left( \int e^{-x} \cos(x) dx \right) \right)$
- Compute integrals

 $y_p(x) = -\sin\left(x\right)$ 

- Substitute particular solution into general solution to ODE  $y = c_2 x e^x + c_1 e^x - \sin(x)$
- $\Box \qquad \text{Check validity of solution } y = c_2 x e^x + c_1 e^x \sin(x)$ 
  - Use initial condition y(0) = 1

$$1 = c_1$$

 $\circ \quad {\rm Compute \ derivative \ of \ the \ solution}$ 

$$y' = c_2 e^x + c_2 x e^x + c_1 e^x - \cos(x)$$

 $\circ \quad \text{Use the initial condition } y'\Big|_{\{x=0\}} = 0$ 

$$0 = -1 + c_1 + c_2$$

• Solve for  $c_1$  and  $c_2$ 

$$\{c_1 = 1, c_2 = 0\}$$

- $\circ$   $\;$  Substitute constant values into general solution and simplify
  - $y = e^x \sin\left(x\right)$
- Solution to the IVP

$$y = e^x - \sin\left(x\right)$$

#### Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`</pre>
```

Solution by Maple Time used: 0.015 (sec). Leaf size: 11

dsolve([diff(y(x),x\$2)-2\*diff(y(x),x)+y(x)=2\*cos(x),y(0) = 1, D(y)(0) = 0],y(x), singsol=all

 $y(x) = e^x - \sin(x)$ 

Solution by Mathematica Time used: 0.019 (sec). Leaf size: 13

DSolve[{y''[x]-2\*y'[x]+y[x]==2\*Cos[x],{y[0]==1,y'[0]==0}},y[x],x,IncludeSingularSolutions ->

 $y(x) \to e^x - \sin(x)$ 

# 1.9 problem 10.2.11 (iii)

1.9.1	Existence and uniqueness analysis	80
1.9.2	Solving as second order linear constant coeff ode	81
1.9.3	Solving using Kovacic algorithm	85
1.9.4	Maple step by step solution	90

Internal problem ID [5053] Internal file name [OUTPUT/4546\_Sunday\_June\_05\_2022\_03\_00\_39\_PM\_97533883/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.2, ODEs with constant Coefficients. page 307
Problem number: 10.2.11 (iii).
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second\_order\_linear\_constant\_coeff"

Maple gives the following as the ode type

[[\_2nd\_order, \_linear, \_nonhomogeneous]]

$$y'' + 16y = 16\cos\left(4x\right)$$

With initial conditions

$$[y(0) = 1, y'(0) = 0]$$

### 1.9.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 0$$
  

$$q(x) = 16$$
  

$$F = 16 \cos (4x)$$

Hence the ode is

$$y'' + 16y = 16\cos\left(4x\right)$$

The domain of p(x) = 0 is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is inside this domain. The domain of q(x) = 16 is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is also inside this domain. The domain of  $F = 16 \cos(4x)$  is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is also inside this domain. Hence solution exists and is unique.

#### 1.9.2 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = f(x)$$

Where  $A = 1, B = 0, C = 16, f(x) = 16 \cos(4x)$ . Let the solution be

 $y = y_h + y_p$ 

Where  $y_h$  is the solution to the homogeneous ODE Ay''(x) + By'(x) + Cy(x) = 0, and  $y_p$  is a particular solution to the non-homogeneous ODE Ay''(x) + By'(x) + Cy(x) = f(x).  $y_h$  is the solution to

$$y'' + 16y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above A = 1, B = 0, C = 16. Let the solution be  $y = e^{\lambda x}$ . Substituting this into the ODE gives

$$\lambda^2 \mathrm{e}^{\lambda x} + 16 \,\mathrm{e}^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by  $e^{\lambda x}$  gives

$$\lambda^2 + 16 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting A = 1, B = 0, C = 16 into the above gives

$$\lambda_{1,2} = \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(16)}$$
$$= \pm 4i$$

Hence

$$\lambda_1 = +4i$$
$$\lambda_2 = -4i$$

Which simplifies to

$$\lambda_1 = 4i$$
$$\lambda_2 = -4i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where  $\alpha = 0$  and  $\beta = 4$ . Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x} (c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^0(c_1 \cos{(4x)} + c_2 \sin{(4x)})$$

Or

$$y = c_1 \cos\left(4x\right) + c_2 \sin\left(4x\right)$$

Therefore the homogeneous solution  $y_h$  is

$$y_h = c_1 \cos(4x) + c_2 \sin(4x)$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$16\cos(4x)$$

Shows that the corresponding undetermined set of the basis functions (UC\_set) for the trial solution is

$$\left[\left\{\cos\left(4x\right),\sin\left(4x\right)\right\}\right]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ \cos\left(4x\right),\sin\left(4x\right) \right\}$$

Since  $\cos(4x)$  is duplicated in the UC\_set, then this basis is multiplied by extra x. The UC\_set becomes

$$\left[\left\{x\cos\left(4x\right),x\sin\left(4x\right)\right\}\right]$$

Since there was duplication between the basis functions in the UC\_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC\_set.

$$y_p = A_1 x \cos\left(4x\right) + A_2 x \sin\left(4x\right)$$

The unknowns  $\{A_1, A_2\}$  are found by substituting the above trial solution  $y_p$  into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-8A_1\sin(4x) + 8A_2\cos(4x) = 16\cos(4x)$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = 0, A_2 = 2]$$

Substituting the above back in the above trial solution  $y_p$ , gives the particular solution

$$y_p = 2x\sin\left(4x\right)$$

Therefore the general solution is

$$y = y_h + y_p$$
  
=  $(c_1 \cos (4x) + c_2 \sin (4x)) + (2x \sin (4x))$ 

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 \cos(4x) + c_2 \sin(4x) + 2x \sin(4x) \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting y = 1 and x = 0 in the above gives

$$1 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -4c_1 \sin(4x) + 4c_2 \cos(4x) + 2\sin(4x) + 8x \cos(4x)$$

substituting y' = 0 and x = 0 in the above gives

$$0 = 4c_2 \tag{2A}$$

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

$$c_1 = 1$$
  
 $c_2 = 0$ 

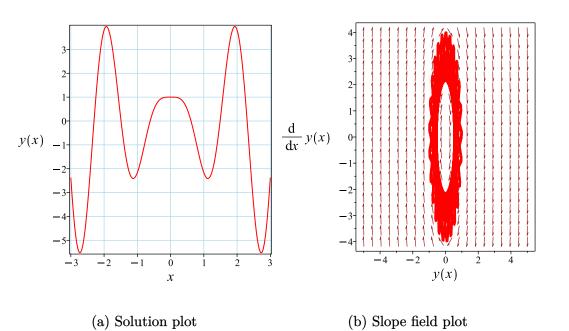
Substituting these values back in above solution results in

$$y = 2x\sin\left(4x\right) + \cos\left(4x\right)$$

Summary

The solution(s) found are the following

$$y = 2x\sin\left(4x\right) + \cos\left(4x\right) \tag{1}$$



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Verification of solutions

$$y = 2x\sin\left(4x\right) + \cos\left(4x\right)$$

Verified OK.

### 1.9.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 16y = 0 (1)$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$
  

$$B = 0$$

$$C = 16$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int \frac{B}{2A} \, dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-16}{1} \tag{6}$$

Comparing the above to (5) shows that

s = -16t = 1

Therefore eq. (4) becomes

$$z''(x) = -16z(x)$$
(7)

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 14: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 0 - 0$$
$$= 0$$

There are no poles in r. Therefore the set of poles  $\Gamma$  is empty. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since r = -16 is not a function of x, then there is no need run Kovacic algorithm to obtain a solution for transformed ode z'' = rz as one solution is

$$z_1(x) = \cos\left(4x\right)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -rac{1}{2}rac{B}{A}\,dx}$$

Since B = 0 then the above reduces to

$$y_1 = z_1$$
$$= \cos(4x)$$

Which simplifies to

$$y_1 = \cos\left(4x\right)$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2=y_1\int rac{e^{\int -rac{B}{A}\,dx}}{y_1^2}\,dx$$

Since B = 0 then the above becomes

$$y_2 = y_1 \int \frac{1}{y_1^2} dx$$
$$= \cos(4x) \int \frac{1}{\cos(4x)^2} dx$$
$$= \cos(4x) \left(\frac{\tan(4x)}{4}\right)$$

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
  
=  $c_1(\cos(4x)) + c_2\left(\cos(4x)\left(\frac{\tan(4x)}{4}\right)\right)$ 

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where  $y_h$  is the solution to the homogeneous ODE Ay''(x) + By'(x) + Cy(x) = 0, and  $y_p$  is a particular solution to the nonhomogeneous ODE Ay''(x) + By'(x) + Cy(x) = f(x).  $y_h$  is the solution to

$$y'' + 16y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 \cos(4x) + \frac{c_2 \sin(4x)}{4}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

```
16\cos(4x)
```

Shows that the corresponding undetermined set of the basis functions (UC\_set) for the trial solution is

$$\left[\left\{\cos\left(4x\right),\sin\left(4x\right)\right\}\right]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{\frac{\sin\left(4x\right)}{4},\cos\left(4x\right)\right\}$$

Since  $\cos(4x)$  is duplicated in the UC\_set, then this basis is multiplied by extra x. The UC\_set becomes

$$\left[\left\{x\cos\left(4x\right),x\sin\left(4x\right)\right\}\right]$$

Since there was duplication between the basis functions in the UC\_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC\_set.

$$y_p = A_1 x \cos\left(4x\right) + A_2 x \sin\left(4x\right)$$

The unknowns  $\{A_1, A_2\}$  are found by substituting the above trial solution  $y_p$  into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-8A_1\sin(4x) + 8A_2\cos(4x) = 16\cos(4x)$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = 0, A_2 = 2]$$

Substituting the above back in the above trial solution  $y_p$ , gives the particular solution

$$y_p = 2x\sin\left(4x\right)$$

Therefore the general solution is

$$y = y_h + y_p$$
  
=  $\left(c_1 \cos(4x) + \frac{c_2 \sin(4x)}{4}\right) + (2x \sin(4x))$ 

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 \cos(4x) + \frac{c_2 \sin(4x)}{4} + 2x \sin(4x) \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting y = 1 and x = 0 in the above gives

$$1 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -4c_1 \sin(4x) + c_2 \cos(4x) + 2\sin(4x) + 8x \cos(4x)$$

substituting y' = 0 and x = 0 in the above gives

$$0 = c_2 \tag{2A}$$

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

$$c_1 = 1$$
$$c_2 = 0$$

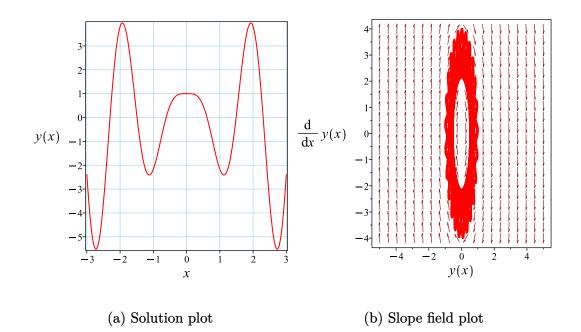
Substituting these values back in above solution results in

$$y = 2x\sin\left(4x\right) + \cos\left(4x\right)$$

Summary

The solution(s) found are the following

$$y = 2x\sin\left(4x\right) + \cos\left(4x\right) \tag{1}$$



Verification of solutions

 $y = 2x\sin\left(4x\right) + \cos\left(4x\right)$ 

Verified OK.

# 1.9.4 Maple step by step solution

Let's solve

$$y'' + 16y = 16\cos(4x), y(0) = 1, y'\Big|_{\{x=0\}} = 0$$

• Highest derivative means the order of the ODE is 2

- Characteristic polynomial of homogeneous ODE  $r^2 + 16 = 0$
- Use quadratic formula to solve for *r*

$$r = \frac{0 \pm (\sqrt{-64})}{2}$$

- Roots of the characteristic polynomial r = (-4 I, 4 I)
- 1st solution of the homogeneous ODE

 $y_1(x) = \cos\left(4x\right)$ 

- 2nd solution of the homogeneous ODE  $y_2(x) = \sin(4x)$
- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x) + y_p(x)$$

- Substitute in solutions of the homogeneous ODE  $y = c_1 \cos(4x) + c_2 \sin(4x) + y_p(x)$
- $\Box$  Find a particular solution  $y_p(x)$  of the ODE
  - Use variation of parameters to find  $y_p$  here f(x) is the forcing function

$$\left[y_p(x) = -y_1(x) \left(\int \frac{y_2(x)f(x)}{W(y_1(x),y_2(x))} dx\right) + y_2(x) \left(\int \frac{y_1(x)f(x)}{W(y_1(x),y_2(x))} dx\right), f(x) = 16\cos\left(4x\right)\right]$$

• Wronskian of solutions of the homogeneous equation

$$W(y_1(x), y_2(x)) = \begin{bmatrix} \cos(4x) & \sin(4x) \\ -4\sin(4x) & 4\cos(4x) \end{bmatrix}$$

 $\circ \quad \text{Compute Wronskian} \\$ 

$$W(y_1(x), y_2(x)) = 4$$

- Substitute functions into equation for  $y_p(x)$  $y_p(x) = -2\cos(4x) \left(\int \sin(8x) dx\right) + 2\sin(4x) \left(\int (1 + \cos(8x)) dx\right)$
- Compute integrals

$$y_p(x) = \frac{\cos(4x)}{4} + 2x\sin(4x)$$

• Substitute particular solution into general solution to ODE

 $y = c_1 \cos(4x) + c_2 \sin(4x) + \frac{\cos(4x)}{4} + 2x \sin(4x)$ 

- $\Box \qquad \text{Check validity of solution } y = c_1 \cos(4x) + c_2 \sin(4x) + \frac{\cos(4x)}{4} + 2x \sin(4x)$ 
  - Use initial condition y(0) = 1

$$1 = \frac{1}{4} + c_1$$

• Compute derivative of the solution

$$y' = -4c_1 \sin(4x) + 4c_2 \cos(4x) + \sin(4x) + 8x \cos(4x)$$

• Use the initial condition 
$$y'\Big|_{\{x=0\}} = 0$$
  
 $0 = 4c_2$ 

• Solve for  $c_1$  and  $c_2$ 

 $\left\{c_1 = \frac{3}{4}, c_2 = 0\right\}$ 

• Substitute constant values into general solution and simplify

 $y = 2x\sin\left(4x\right) + \cos\left(4x\right)$ 

• Solution to the IVP

 $y = 2x\sin\left(4x\right) + \cos\left(4x\right)$ 

### Maple trace

`Methods for second order ODEs: --- Trying classification methods --trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE checking if the LODE has constant coefficients <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`</pre>

Solution by Maple Time used: 0.015 (sec). Leaf size: 16

dsolve([diff(y(x),x\$2)+16\*y(x)=16\*cos(4\*x),y(0) = 1, D(y)(0) = 0],y(x), singsol=all)

 $y(x) = \cos\left(4x\right) + 2\sin\left(4x\right)x$ 

Solution by Mathematica Time used: 0.036 (sec). Leaf size: 17

DSolve[{y''[x]+16\*y[x]==16\*Cos[4\*x],{y[0]==1,y'[0]==0}},y[x],x,IncludeSingularSolutions -> T

 $y(x) \rightarrow 2x\sin(4x) + \cos(4x)$ 

# 1.10 problem 10.2.11 (iv)

1.10.1	Existence and uniqueness analysis	93
1.10.2	Solving as second order linear constant coeff ode	94
1.10.3	Solving using Kovacic algorithm	99
1.10.4	Maple step by step solution	106

Internal problem ID [5054] Internal file name [OUTPUT/4547\_Sunday\_June\_05\_2022\_03\_00\_40\_PM\_40071015/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.2, ODEs with constant Coefficients. page 307
Problem number: 10.2.11 (iv).
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second\_order\_linear\_constant\_coeff"

Maple gives the following as the ode type

[[\_2nd\_order, \_linear, \_nonhomogeneous]]

$$y'' - y = \cosh\left(x\right)$$

With initial conditions

[y(0) = 1, y'(0) = 0]

### 1.10.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 0$$
$$q(x) = -1$$
$$F = \cosh(x)$$

Hence the ode is

$$y'' - y = \cosh\left(x\right)$$

The domain of p(x) = 0 is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is inside this domain. The domain of q(x) = -1 is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is also inside this domain. The domain of  $F = \cosh(x)$  is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is also inside this domain. Hence solution exists and is unique.

#### 1.10.2 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = f(x)$$

Where  $A = 1, B = 0, C = -1, f(x) = \cosh(x)$ . Let the solution be

 $y = y_h + y_p$ 

Where  $y_h$  is the solution to the homogeneous ODE Ay''(x) + By'(x) + Cy(x) = 0, and  $y_p$  is a particular solution to the non-homogeneous ODE Ay''(x) + By'(x) + Cy(x) = f(x).  $y_h$  is the solution to

$$y'' - y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above A = 1, B = 0, C = -1. Let the solution be  $y = e^{\lambda x}$ . Substituting this into the ODE gives

$$\lambda^2 \mathrm{e}^{\lambda x} - \mathrm{e}^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by  $e^{\lambda x}$  gives

$$\lambda^2 - 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting A = 1, B = 0, C = -1 into the above gives

$$\lambda_{1,2} = \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(-1)}$$
$$= \pm 1$$

Hence

$$\lambda_1 = +1$$
$$\lambda_2 = -1$$

Which simplifies to

$$\lambda_1 = 1$$
$$\lambda_2 = -1$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$
  
 $y = c_1 e^{(1)x} + c_2 e^{(-1)x}$ 

Or

 $y = c_1 \mathrm{e}^x + c_2 \mathrm{e}^{-x}$ 

Therefore the homogeneous solution  $y_h$  is

$$y_h = c_1 \mathrm{e}^x + c_2 \mathrm{e}^{-x}$$

The particular solution  $y_p$  can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where  $u_1, u_2$  to be determined, and  $y_1, y_2$  are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = e^x$$
$$y_2 = e^{-x}$$

In the Variation of parameters  $u_1, u_2$  are found using

$$u_1 = -\int \frac{y_2 f(x)}{a W(x)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(x)}{a W(x)} \tag{3}$$

Where W(x) is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by  $W = \begin{vmatrix} y_1 & y_2 \\ y'_1 & y'_2 \end{vmatrix}$ . Hence $W = \begin{vmatrix} e^x & e^{-x} \\ \frac{d}{dx}(e^x) & \frac{d}{dx}(e^{-x}) \end{vmatrix}$ Which gives

$$W = \begin{vmatrix} \mathbf{e}^x & \mathbf{e}^{-x} \\ \mathbf{e}^x & -\mathbf{e}^{-x} \end{vmatrix}$$

Therefore

$$W = (e^{x})(-e^{-x}) - (e^{-x})(e^{x})$$

Which simplifies to

$$W = -2 \,\mathrm{e}^x \mathrm{e}^{-x}$$

Which simplifies to

$$W = -2$$

Therefore Eq. (2) becomes

$$u_1 = -\int \frac{\mathrm{e}^{-x} \cosh\left(x\right)}{-2} \, dx$$

Which simplifies to

$$u_1 = -\int -\frac{\mathrm{e}^{-x}\cosh\left(x\right)}{2}dx$$

Hence

$$u_1 = \frac{x}{4} + \frac{\sinh(2x)}{8} - \frac{\cosh(2x)}{8}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\mathrm{e}^x \cosh\left(x\right)}{-2} \, dx$$

Which simplifies to

$$u_2 = \int -\frac{\mathrm{e}^x \cosh\left(x\right)}{2} dx$$

Hence

$$u_{2} = -rac{\cosh(x)^{2}}{4} - rac{\cosh(x)\sinh(x)}{4} - rac{x}{4}$$

Which simplifies to

$$u_1 = \frac{x}{4} + \frac{\sinh(2x)}{8} - \frac{\cosh(2x)}{8}$$
$$u_2 = -\frac{x}{4} - \frac{\sinh(2x)}{8} - \frac{\cosh(2x)}{8} - \frac{1}{8}$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = \left(\frac{x}{4} + \frac{\sinh(2x)}{8} - \frac{\cosh(2x)}{8}\right)e^x + \left(-\frac{x}{4} - \frac{\sinh(2x)}{8} - \frac{\cosh(2x)}{8} - \frac{1}{8}\right)e^{-x}$$

Which simplifies to

$$y_p(x) = \frac{\left(-2x - 1 - \cosh\left(2x\right) - \sinh\left(2x\right)\right) e^{-x}}{8} + \frac{\left(x + \frac{\sinh(2x)}{2} - \frac{\cosh(2x)}{2}\right) e^x}{4}$$

Therefore the general solution is

$$y = y_h + y_p$$
  
=  $(c_1 e^x + c_2 e^{-x}) + \left(\frac{(-2x - 1 - \cosh(2x) - \sinh(2x))e^{-x}}{8} + \frac{\left(x + \frac{\sinh(2x)}{2} - \frac{\cosh(2x)}{2}\right)e^x}{4}\right)$ 

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^x + c_2 e^{-x} + \frac{\left(-2x - 1 - \cosh\left(2x\right) - \sinh\left(2x\right)\right) e^{-x}}{8} + \frac{\left(x + \frac{\sinh(2x)}{2} - \frac{\cosh(2x)}{2}\right) e^x}{4}$$
(1)

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting y = 1 and x = 0 in the above gives

$$1 = c_1 + c_2 - \frac{3}{8} \tag{1A}$$

Taking derivative of the solution gives

$$y' = c_1 e^x - c_2 e^{-x} + \frac{(-2 - 2\sinh(2x) - 2\cosh(2x))e^{-x}}{8} - \frac{(-2x - 1 - \cosh(2x) - \sinh(2x))e^{-x}}{8} + \frac{(1 + 2)\exp(2x) - 2\cosh(2x) - 2\cosh(2x)}{8} + \frac{(-2x - 2)\exp(2x) - 2\cosh(2x)}{8} + \frac{(-2x - 2$$

substituting y' = 0 and x = 0 in the above gives

$$0 = c_1 - c_2 + \frac{1}{8} \tag{2A}$$

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

$$c_1 = \frac{5}{8}$$
$$c_2 = \frac{3}{4}$$

Substituting these values back in above solution results in

$$y = \frac{5e^x}{8} + \frac{5e^{-x}}{8} - \frac{e^{-x}\sinh(2x)}{8} - \frac{e^{-x}\cosh(2x)}{8} - \frac{xe^{-x}}{4} + \frac{e^xx}{4} + \frac{e^x\sinh(2x)}{8} - \frac{e^x\cosh(2x)}{8} - \frac{e^$$

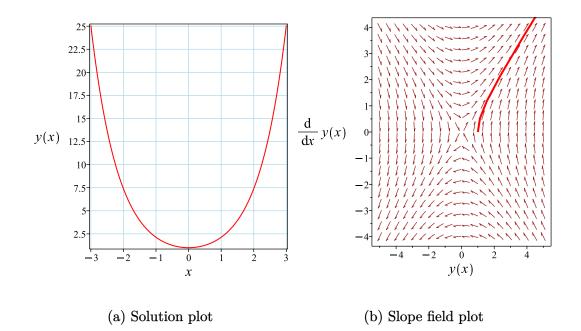
Which simplifies to

$$y = \frac{\left(-2x - \cosh\left(2x\right) - \sinh\left(2x\right) + 5\right)e^{-x}}{8} + \frac{e^{x}\left(x - \frac{\cosh(2x)}{2} + \frac{\sinh(2x)}{2} + \frac{5}{2}\right)}{4}$$

#### Summary

The solution(s) found are the following

$$y = \frac{\left(-2x - \cosh\left(2x\right) - \sinh\left(2x\right) + 5\right)e^{-x}}{8} + \frac{e^{x}\left(x - \frac{\cosh(2x)}{2} + \frac{\sinh(2x)}{2} + \frac{5}{2}\right)}{4}$$
(1)



Verification of solutions

$$y = \frac{\left(-2x - \cosh\left(2x\right) - \sinh\left(2x\right) + 5\right)e^{-x}}{8} + \frac{e^{x}\left(x - \frac{\cosh(2x)}{2} + \frac{\sinh(2x)}{2} + \frac{5}{2}\right)}{4}$$

Verified OK.

# 1.10.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' - y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$
  

$$B = 0$$

$$C = -1$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{1} \tag{6}$$

Comparing the above to (5) shows that

s = 1t = 1

Therefore eq. (4) becomes

$$z''(x) = z(x) \tag{7}$$

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 16: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 0 - 0$$
$$= 0$$

There are no poles in r. Therefore the set of poles  $\Gamma$  is empty. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since r = 1 is not a function of x, then there is no need run Kovacic algorithm to obtain a solution for transformed ode z'' = rz as one solution is

$$z_1(x) = \mathrm{e}^{-x}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2}\frac{B}{A}\,dx}$$

Since B = 0 then the above reduces to

$$y_1 = z_1$$
$$= e^{-x}$$

Which simplifies to

$$y_1 = e^{-x}$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2 = y_1 \int rac{e^{\int -rac{B}{A}\,dx}}{y_1^2}\,dx$$

Since B = 0 then the above becomes

$$y_2 = y_1 \int \frac{1}{y_1^2} dx$$
$$= e^{-x} \int \frac{1}{e^{-2x}} dx$$
$$= e^{-x} \left(\frac{e^{2x}}{2}\right)$$

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
  
=  $c_1 (e^{-x}) + c_2 \left( e^{-x} \left( \frac{e^{2x}}{2} \right) \right)$ 

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where  $y_h$  is the solution to the homogeneous ODE Ay''(x) + By'(x) + Cy(x) = 0, and  $y_p$  is a particular solution to the nonhomogeneous ODE Ay''(x) + By'(x) + Cy(x) = f(x).  $y_h$  is the solution to

$$y'' - y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 \mathrm{e}^{-x} + \frac{c_2 \mathrm{e}^x}{2}$$

The particular solution  $y_p$  can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where  $u_1, u_2$  to be determined, and  $y_1, y_2$  are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = e^{-x}$$
$$y_2 = \frac{e^x}{2}$$

In the Variation of parameters  $u_1, u_2$  are found using

$$u_1 = -\int \frac{y_2 f(x)}{a W(x)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(x)}{a W(x)} \tag{3}$$

Where W(x) is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by  $W = \begin{vmatrix} y_1 & y_2 \\ y'_1 & y'_2 \end{vmatrix}$ . Hence

$$W = \begin{vmatrix} e^{-x} & \frac{e^x}{2} \\ \frac{d}{dx}(e^{-x}) & \frac{d}{dx}\left(\frac{e^x}{2}\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \mathrm{e}^{-x} & \frac{\mathrm{e}^{x}}{2} \\ -\mathrm{e}^{-x} & \frac{\mathrm{e}^{x}}{2} \end{vmatrix}$$

Therefore

$$W = \left(e^{-x}\right) \left(\frac{e^x}{2}\right) - \left(\frac{e^x}{2}\right) \left(-e^{-x}\right)$$

Which simplifies to

$$W = e^x e^{-x}$$

Which simplifies to

W = 1

Therefore Eq. (2) becomes

$$u_1 = -\int \frac{\frac{\mathrm{e}^x \cosh(x)}{2}}{1} \, dx$$

Which simplifies to

$$u_1 = -\int \frac{\mathrm{e}^x \cosh\left(x\right)}{2} dx$$

Hence

$$u_{1} = -\frac{\cosh(x)^{2}}{4} - \frac{\cosh(x)\sinh(x)}{4} - \frac{x}{4}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\mathrm{e}^{-x} \cosh\left(x\right)}{1} \, dx$$

Which simplifies to

$$u_2 = \int e^{-x} \cosh\left(x\right) dx$$

Hence

$$u_2 = \frac{x}{2} + \frac{\sinh(2x)}{4} - \frac{\cosh(2x)}{4}$$

Which simplifies to

$$u_{1} = -\frac{x}{4} - \frac{\sinh(2x)}{8} - \frac{\cosh(2x)}{8} - \frac{1}{8}$$
$$u_{2} = \frac{x}{2} + \frac{\sinh(2x)}{4} - \frac{\cosh(2x)}{4}$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = \left(-\frac{x}{4} - \frac{\sinh\left(2x\right)}{8} - \frac{\cosh\left(2x\right)}{8} - \frac{1}{8}\right)e^{-x} + \frac{\left(\frac{x}{2} + \frac{\sinh\left(2x\right)}{4} - \frac{\cosh\left(2x\right)}{4}\right)e^x}{2}$$

Which simplifies to

$$y_p(x) = \frac{\left(-2x - 1 - \cosh\left(2x\right) - \sinh\left(2x\right)\right)e^{-x}}{8} + \frac{\left(x + \frac{\sinh(2x)}{2} - \frac{\cosh(2x)}{2}\right)e^x}{4}$$

Therefore the general solution is

$$y = y_h + y_p$$

$$= \left(c_1 e^{-x} + \frac{c_2 e^x}{2}\right)$$

$$+ \left(\frac{\left(-2x - 1 - \cosh\left(2x\right) - \sinh\left(2x\right)\right) e^{-x}}{8} + \frac{\left(x + \frac{\sinh(2x)}{2} - \frac{\cosh(2x)}{2}\right) e^x}{4}\right)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-x} + \frac{c_2 e^x}{2} + \frac{\left(-2x - 1 - \cosh\left(2x\right) - \sinh\left(2x\right)\right) e^{-x}}{8} + \frac{\left(x + \frac{\sinh(2x)}{2} - \frac{\cosh(2x)}{2}\right) e^x}{4}$$
(1)

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting y = 1 and x = 0 in the above gives

$$1 = c_1 + \frac{c_2}{2} - \frac{3}{8} \tag{1A}$$

Taking derivative of the solution gives

$$y' = -c_1 e^{-x} + \frac{c_2 e^x}{2} + \frac{(-2 - 2\sinh(2x) - 2\cosh(2x))e^{-x}}{8} - \frac{(-2x - 1 - \cosh(2x) - \sinh(2x))e^{-x}}{8} + \frac{(1 - 2x)e^{-x}}{8} + \frac{(-2x)e^{-x}}{8} + \frac{(-2x)e^$$

substituting y' = 0 and x = 0 in the above gives

$$0 = -c_1 + \frac{c_2}{2} + \frac{1}{8} \tag{2A}$$

Equations  $\{1A, 2A\}$  are now solved for  $\{c_1, c_2\}$ . Solving for the constants gives

$$c_1 = \frac{3}{4}$$
$$c_2 = \frac{5}{4}$$

Substituting these values back in above solution results in

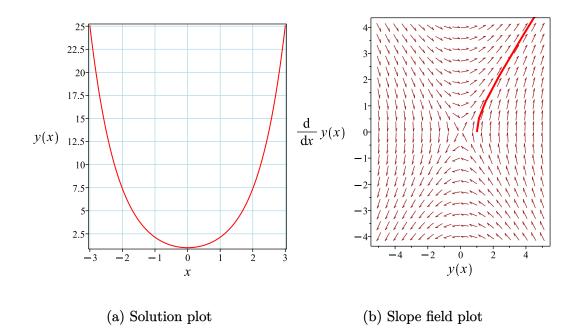
$$y = \frac{5e^x}{8} + \frac{5e^{-x}}{8} - \frac{e^{-x}\sinh(2x)}{8} - \frac{e^{-x}\cosh(2x)}{8} - \frac{xe^{-x}}{4} + \frac{e^xx}{4} + \frac{e^x\sinh(2x)}{8} - \frac{e^x\cosh(2x)}{8} - \frac{e^$$

Which simplifies to

$$y = \frac{\left(-2x - \cosh\left(2x\right) - \sinh\left(2x\right) + 5\right)e^{-x}}{8} + \frac{e^{x}\left(x - \frac{\cosh(2x)}{2} + \frac{\sinh(2x)}{2} + \frac{5}{2}\right)}{4}$$

 $\frac{\text{Summary}}{\text{The solution(s) found are the following}}$ 

$$y = \frac{\left(-2x - \cosh\left(2x\right) - \sinh\left(2x\right) + 5\right)e^{-x}}{8} + \frac{e^{x}\left(x - \frac{\cosh\left(2x\right)}{2} + \frac{\sinh\left(2x\right)}{2} + \frac{5}{2}\right)}{4} \quad (1)$$

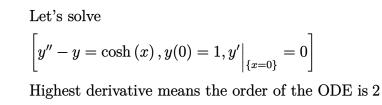


Verification of solutions

$$y = \frac{\left(-2x - \cosh\left(2x\right) - \sinh\left(2x\right) + 5\right)e^{-x}}{8} + \frac{e^{x}\left(x - \frac{\cosh(2x)}{2} + \frac{\sinh(2x)}{2} + \frac{5}{2}\right)}{4}$$

Verified OK.

## 1.10.4 Maple step by step solution



- Highest derivative means the order of the ODE is 2 y''
- Characteristic polynomial of homogeneous ODE  $r^2 1 = 0$
- Factor the characteristic polynomial

(r-1)(r+1) = 0

- Roots of the characteristic polynomial r = (-1, 1)
- 1st solution of the homogeneous ODE

 $y_1(x) = \mathrm{e}^{-x}$ 

- 2nd solution of the homogeneous ODE  $y_2(x) = e^x$
- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x) + y_p(x)$$

• Substitute in solutions of the homogeneous ODE  $y = c_1 e^{-x} + c_2 e^x + y_p(x)$ 

 $\Box$  Find a particular solution  $y_p(x)$  of the ODE

• Use variation of parameters to find  $y_p$  here f(x) is the forcing function

$$\left[y_p(x) = -y_1(x) \left(\int \frac{y_2(x)f(x)}{W(y_1(x), y_2(x))} dx\right) + y_2(x) \left(\int \frac{y_1(x)f(x)}{W(y_1(x), y_2(x))} dx\right), f(x) = \cosh\left(x\right)\right]$$

• Wronskian of solutions of the homogeneous equation

$$W(y_1(x)\,,y_2(x))=\left[egin{array}{cc} \mathrm{e}^{-x}&\mathrm{e}^x\ -\mathrm{e}^{-x}&\mathrm{e}^x\end{array}
ight]$$

• Compute Wronskian

$$W(y_1(x), y_2(x)) = 2$$

- Substitute functions into equation for  $y_p(x)$  $y_p(x) = -\frac{e^{-x}(\int e^x \cosh(x)dx)}{2} + \frac{e^x(\int e^{-x} \cosh(x)dx)}{2}$
- Compute integrals

$$y_p(x) = rac{(-2x-1-\cosh(2x)-\sinh(2x))\mathrm{e}^{-x}}{8} + rac{\left(x+rac{\sinh(2x)}{2}-rac{\cosh(2x)}{2}
ight)\mathrm{e}^x}{4}$$

• Substitute particular solution into general solution to ODE

$$y = c_1 e^{-x} + c_2 e^x + \frac{(-2x - 1 - \cosh(2x) - \sinh(2x))e^{-x}}{8} + \frac{\left(x + \frac{\sinh(2x)}{2} - \frac{\cosh(2x)}{2}\right)e^x}{4}$$

Check validity of solution  $y = c_1 e^{-x} + c_2 e^x + \frac{(-2x - 1 - \cosh(2x) - \sinh(2x))e^{-x}}{8} + \frac{\left(x + \frac{\sinh(2x)}{2} - \frac{\cosh(2x)}{2}\right)e^x}{4}$ 

• Use initial condition y(0) = 1

$$1 = c_1 + c_2 - \frac{3}{8}$$

 $\circ \quad {\rm Compute \ derivative \ of \ the \ solution}$ 

$$y' = -c_1 e^{-x} + c_2 e^x + \frac{(-2-2\sinh(2x)-2\cosh(2x))e^{-x}}{8} - \frac{(-2x-1-\cosh(2x)-\sinh(2x))e^{-x}}{8} + \frac{(1+\cosh(2x)-\sinh(2x))e^x}{4}$$

 $\circ \quad \text{Use the initial condition } y'\Big|_{\{x=0\}} = 0$ 

 $0 = -c_1 + c_2 + \frac{1}{8}$ 

• Solve for  $c_1$  and  $c_2$ 

 $\left\{c_1 = \frac{3}{4}, c_2 = \frac{5}{8}\right\}$ 

• Substitute constant values into general solution and simplify

 $y = \frac{(-2x - \cosh(2x) - \sinh(2x) + 5)e^{-x}}{8} + \frac{e^x \left(x - \frac{\cosh(2x)}{2} + \frac{\sinh(2x)}{2} + \frac{5}{2}\right)}{4}$ 

• Solution to the IVP

$$y = \frac{(-2x - \cosh(2x) - \sinh(2x) + 5)e^{-x}}{8} + \frac{e^x \left(x - \frac{\cosh(2x)}{2} + \frac{\sinh(2x)}{2} + \frac{5}{2}\right)}{4}$$

Maple trace

`Methods for second order ODEs: --- Trying classification methods --trying a quadrature trying high order exact linear fully integrable trying differential order: 2; linear nonhomogeneous with symmetry [0,1] trying a double symmetry of the form [xi=0, eta=F(x)] -> Try solving first the homogeneous part of the ODE checking if the LODE has constant coefficients <- constant coefficients successful <- solving first the homogeneous part of the ODE successful`</pre>

Solution by Maple Time used: 0.032 (sec). Leaf size: 23

dsolve([diff(y(x),x\$2)-y(x)=cosh(x),y(0) = 1, D(y)(0) = 0],y(x), singsol=all)

$$y(x) = \frac{(-x+2)e^{-x}}{4} + \frac{e^x(x+2)}{4}$$

✓ Solution by Mathematica

Time used: 0.041 (sec). Leaf size: 28

DSolve[{y''[x]-y[x]==Cosh[x],{y[0]==1,y'[0]==0}},y[x],x,IncludeSingularSolutions -> True]

$$y(x) \to \frac{1}{4}e^{-x}(-x+e^{2x}(x+2)+2)$$

# 2 Chapter 10, Differential equations. Section 10.3, ODEs with variable Coefficients. First order. page 315

2.1	problem	10.3.2		•	•			•	•	•	•	•	•	•	•	•	•	•			•	•		•			•	•	•	•	•	•	•	110
2.2	problem	10.3.3		•	•			•	•	•	•	•		•	•	•	•	•			•			•	•		•	•		•	•	•	•	123
2.3	problem	10.3.4		•	•			•	•	•	•	•		•	•	•	•	•			•			•	•		•	•		•	•	•	•	138
2.4	$\operatorname{problem}$	10.3.5	•	•	•				•	•	•	•		•	•	•	•	•	•			•		•	•		•	•		•	•	•		152
2.5	$\operatorname{problem}$	10.3.6	•	•	•				•	•	•	•		•	•	•	•	•	•			•		•	•		•	•		•	•	•		165
2.6	problem	10.3.7		•	•	•	•	•	•	•	•	•	•	•	•	•	•	•	• •	 •	•	•		•	•		•	•	•	•	•	•	•	178
2.7	problem	10.3.8		•	•			•	•	•	•	•		•	•	•	•	•	• •		•	•		•	•		•	•		•	•	•	•	191
2.8	problem	10.3.9	(a)		•	•		•	•	•	•			•	•	•	•	•	• •		•			•	•		•	•		•	•	•	•	206
2.9	$\operatorname{problem}$	10.3.9	(b)		•	•	•	•	•	•	•	•	•	•	•	•	•	•	•	 •	•	•	•	•	•	•	•	•	•	•	•	•	•	223

# 2.1 problem 10.3.2

2.1.1	Solving as linear ode	110
2.1.2	Solving as first order ode lie symmetry lookup ode	112
2.1.3	Solving as exact ode	116
2.1.4	Maple step by step solution	120

Internal problem ID [5055] Internal file name [OUTPUT/4548\_Sunday\_June\_05\_2022\_03\_00\_42\_PM\_28220764/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.3, ODEs with variable Coefficients.
First order. page 315
Problem number: 10.3.2.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[[\_linear, `class A`]]

$$y' - y = e^{2x}$$

#### 2.1.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -1$$
$$q(x) = e^{2x}$$

Hence the ode is

 $y' - y = e^{2x}$ 

The integrating factor  $\mu$  is

$$\mu = e^{\int (-1)dx}$$
$$= e^{-x}$$

The ode becomes

$$\begin{aligned} \frac{\mathrm{d}}{\mathrm{d}x}(\mu y) &= (\mu) \left( \mathrm{e}^{2x} \right) \\ \frac{\mathrm{d}}{\mathrm{d}x} \left( \mathrm{e}^{-x} y \right) &= \left( \mathrm{e}^{-x} \right) \left( \mathrm{e}^{2x} \right) \\ \mathrm{d} \left( \mathrm{e}^{-x} y \right) &= \mathrm{e}^{x} \, \mathrm{d}x \end{aligned}$$

Integrating gives

$$e^{-x}y = \int e^x dx$$
  
 $e^{-x}y = e^x + c_1$ 

Dividing both sides by the integrating factor  $\mu = e^{-x}$  results in

$$y = e^{2x} + c_1 e^x$$

 $\frac{Summary}{The solution(s) found are the following}$ 

$$y = e^{2x} + c_1 e^x \tag{1}$$

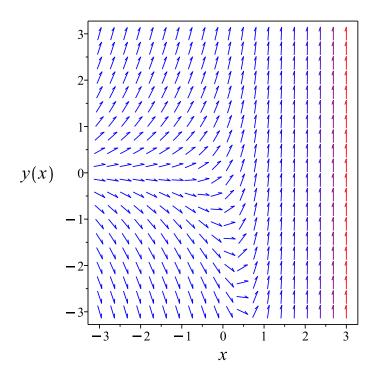


Figure 12: Slope field plot

Verification of solutions

$$y = e^{2x} + c_1 e^x$$

Verified OK.

# 2.1.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = y + e^{2x}$$
  
 $y' = \omega(x, y)$ 

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
 (A)

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi$ ,  $\eta$ 

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f(rac{y}{x})$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x + b_1y + c_1}{a_2x + b_2y + c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x) y + f_2(x) y^2$	0	$e^{-\int f_1 dx}$

Table 18: Lie symmetry infinitesimal lookup table for known first order ODE's

The above table shows that

$$\xi(x,y) = 0$$
  

$$\eta(x,y) = e^x$$
(A1)

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where S(R). Since  $\xi = 0$  then in this special case

$$R = x$$

 ${\cal S}$  is found from

$$S = \int rac{1}{\eta} dy \ = \int rac{1}{\mathrm{e}^x} dy$$

Which results in

$$S = e^{-x}y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}$$
(2)

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = y + e^{2x}$$

Evaluating all the partial derivatives gives

$$R_x = 1$$
  

$$R_y = 0$$
  

$$S_x = -e^{-x}y$$
  

$$S_y = e^{-x}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^x \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = e^R + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$y e^{-x} = e^x + c_1$$

Which simplifies to

$$y e^{-x} = e^x + c_1$$

Which gives

$$y = e^x(e^x + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	$ODE  ext{ in canonical coordinates} \ (R,S)$
$\frac{dy}{dx} = y + e^{2x}$	$R = x$ $S = e^{-x}y$	$\frac{dS}{dR} = e^{R}$

#### Summary

The solution(s) found are the following

$$y = e^x(e^x + c_1) \tag{1}$$

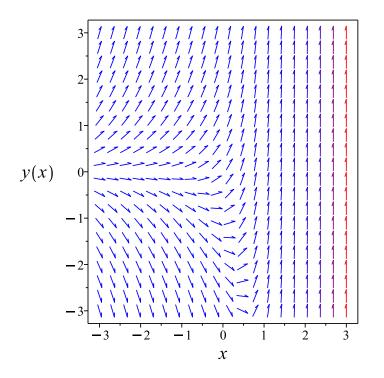


Figure 13: Slope field plot

Verification of solutions

$$y = e^x(e^x + c_1)$$

Verified OK.

#### 2.1.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x,y) + N(x,y)\frac{dy}{dx} = 0$$
 (A)

We assume there exists a function  $\phi(x, y) = c$  where c is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t. x gives

$$\frac{d}{dx}\phi(x,y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x,y) dx + N(x,y) dy = 0$$
(1A)

Therefore

$$dy = (y + e^{2x}) dx$$
$$(-y - e^{2x}) dx + dy = 0$$
(2A)

Comparing (1A) and (2A) shows that

$$M(x, y) = -y - e^{2x}$$
  
 $N(x, y) = 1$ 

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y} \left( -y - e^{2x} \right)$$
$$= -1$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x}(1)$$
$$= 0$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is <u>not exact</u>. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$A = \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right)$$
$$= 1((-1) - (0))$$
$$= -1$$

Since A does not depend on y, then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\mu = e^{\int A \, \mathrm{d}x}$$
$$= e^{\int -1 \, \mathrm{d}x}$$

The result of integrating gives

$$\mu = e^{-x}$$
$$= e^{-x}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called  $\overline{M}$  and  $\overline{N}$  for now so not to confuse them with the original M and N.

$$\overline{M} = \mu M$$
  
= e<sup>-x</sup> (-y - e<sup>2x</sup>)  
= -e<sup>-x</sup>y - e<sup>x</sup>

And

$$\overline{N} = \mu N$$
  
=  $e^{-x}(1)$   
=  $e^{-x}$ 

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$\overline{M} + \overline{N} rac{\mathrm{d}y}{\mathrm{d}x} = 0$$
  
 $\left(-\mathrm{e}^{-x}y - \mathrm{e}^{x}\right) + \left(\mathrm{e}^{-x}\right) rac{\mathrm{d}y}{\mathrm{d}x} = 0$ 

The following equations are now set up to solve for the function  $\phi(x, y)$ 

$$\frac{\partial \phi}{\partial x} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \overline{M} dx$$
$$\int \frac{\partial \phi}{\partial x} dx = \int -e^{-x}y - e^{x} dx$$
$$\phi = e^{-x}y - e^{x} + f(y)$$
(3)

Where f(y) is used for the constant of integration since  $\phi$  is a function of both x and y. Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-x} + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = e^{-x}$ . Therefore equation (4) becomes

$$e^{-x} = e^{-x} + f'(y)$$
 (5)

Solving equation (5) for f'(y) gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for f(y) into equation (3) gives  $\phi$ 

$$\phi = \mathrm{e}^{-x}y - \mathrm{e}^x + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \mathrm{e}^{-x}y - \mathrm{e}^x$$

The solution becomes

$$y = e^x(e^x + c_1)$$

Summary

The solution(s) found are the following

$$y = e^x(e^x + c_1) \tag{1}$$

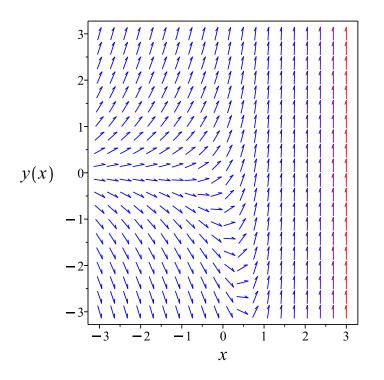


Figure 14: Slope field plot

# Verification of solutions

$$y = e^x(e^x + c_1)$$

Verified OK.

#### 2.1.4 Maple step by step solution

Let's solve

$$y' - y = e^{2x}$$

- Highest derivative means the order of the ODE is 1 y'
- Isolate the derivative

$$y' = y + e^{2x}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE  $y' y = e^{2x}$
- The ODE is linear; multiply by an integrating factor  $\mu(x)$  $\mu(x) (y' - y) = \mu(x) e^{2x}$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$  $\mu(x)(y'-y) = \mu'(x)y + \mu(x)y'$
- Isolate  $\mu'(x)$

 $\mu'(x) = -\mu(x)$ 

- Solve to find the integrating factor  $\mu(x) = \mathrm{e}^{-x}$
- Integrate both sides with respect to x

$$\int \left( rac{d}{dx} (\mu(x) \, y) 
ight) dx = \int \mu(x) \, \mathrm{e}^{2x} dx + c_1$$

• Evaluate the integral on the lhs

$$\mu(x) y = \int \mu(x) e^{2x} dx + c_1$$

• Solve for y

$$y = \frac{\int \mu(x) \mathrm{e}^{2x} dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = e^{-x}$  $y = \frac{\int e^{-x} e^{2x} dx + c_1}{e^{-x}}$
- Evaluate the integrals on the rhs  $y = \frac{e^x + c_1}{e^{-x}}$
- Simplify

$$y = e^x(e^x + c_1)$$

# Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`</pre>
```

Solution by Maple Time used: 0.0 (sec). Leaf size: 11

dsolve(diff(y(x),x)-y(x)=exp(2\*x),y(x), singsol=all)

$$y(x) = (e^x + c_1) e^x$$

✓ Solution by Mathematica Time used: 0.04 (sec). Leaf size: 15

DSolve[y'[x]-y[x]==Exp[2\*x],y[x],x,IncludeSingularSolutions -> True]

$$y(x) \to e^x(e^x + c_1)$$

# 2.2 problem 10.3.3

2	2.2.1	Existence and uniqueness analysis	123						
2	2.2.2	Solving as linear ode	124						
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Internal pr	roblem	ID [5056]							
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Book: Ba	sic Tra	ining in Mathematics. By R. Shankar. Plenum Press. NY. 1995							
Section:	Chapte	er 10, Differential equations. Section 10.3, ODEs with variable Coefficie	ents.						
First order	First order. page 315								
Problem	numl	ber: 10.3.3.							
ODE order: 1.									
ODE deg	gree: 1	1.							

The type(s) of ODE detected by this program : "exact", "linear", "differentialType", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_linear]

$$x^2y' + 2xy = x - 1$$

With initial conditions

$$[y(1) = 0]$$

#### 2.2.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{2}{x}$$
$$q(x) = \frac{x-1}{x^2}$$

Hence the ode is

$$y' + \frac{2y}{x} = \frac{x-1}{x^2}$$

The domain of  $p(x) = \frac{2}{x}$  is

$$\{x < 0 \lor 0 < x\}$$

And the point  $x_0 = 1$  is inside this domain. The domain of  $q(x) = \frac{x-1}{x^2}$  is

$${x < 0 \lor 0 < x}$$

And the point  $x_0 = 1$  is also inside this domain. Hence solution exists and is unique.

#### 2.2.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor  $\mu$  is

$$\mu = \mathrm{e}^{\int rac{2}{x} dx}$$
  
=  $x^2$ 

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu y) = (\mu) \left(\frac{x-1}{x^2}\right)$$
$$\frac{\mathrm{d}}{\mathrm{d}x}(y x^2) = (x^2) \left(\frac{x-1}{x^2}\right)$$
$$\mathrm{d}(y x^2) = (x-1) \mathrm{d}x$$

Integrating gives

$$y x^{2} = \int x - 1 dx$$
$$y x^{2} = \frac{1}{2}x^{2} - x + c_{1}$$

Dividing both sides by the integrating factor  $\mu=x^2$  results in

$$y = \frac{\frac{1}{2}x^2 - x}{x^2} + \frac{c_1}{x^2}$$

which simplifies to

$$y = \frac{x^2 + 2c_1 - 2x}{2x^2}$$

Initial conditions are used to solve for  $c_1$ . Substituting x = 1 and y = 0 in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{1}{2} + c_1$$

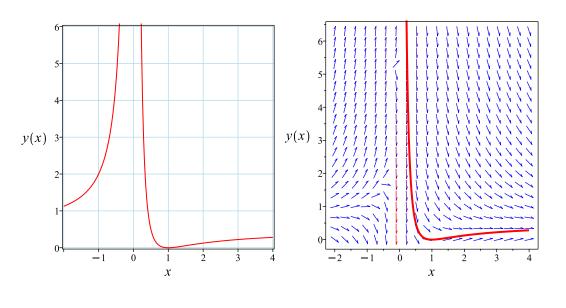
 $c_1 = \frac{1}{2}$ 

Substituting  $c_1$  found above in the general solution gives

$$y = \frac{x^2 - 2x + 1}{2x^2}$$

 $\frac{Summary}{The solution(s) found are the following}$ 

$$y = \frac{x^2 - 2x + 1}{2x^2} \tag{1}$$



(a) Solution plot

(b) Slope field plot

Verification of solutions

$$y = \frac{x^2 - 2x + 1}{2x^2}$$

Verified OK.

#### 2.2.3 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-2xy + x - 1}{x^2}$$
(1)

Which becomes

$$0 = (-x^2) dy + (-2xy + x - 1) dx$$
(2)

But the RHS is complete differential because

$$(-x^2) dy + (-2xy + x - 1) dx = d\left(-y x^2 + \frac{1}{2}x^2 - x\right)$$

Hence (2) becomes

$$0 = d\left(-y\,x^2 + \frac{1}{2}x^2 - x\right)$$

Integrating both sides gives gives these solutions

$$y = \frac{x^2 + 2c_1 - 2x}{2x^2} + c_1$$

Initial conditions are used to solve for  $c_1$ . Substituting x = 1 and y = 0 in the above solution gives an equation to solve for the constant of integration.

$$0 = 2c_1 - \frac{1}{2}$$
$$c_1 = \frac{1}{4}$$

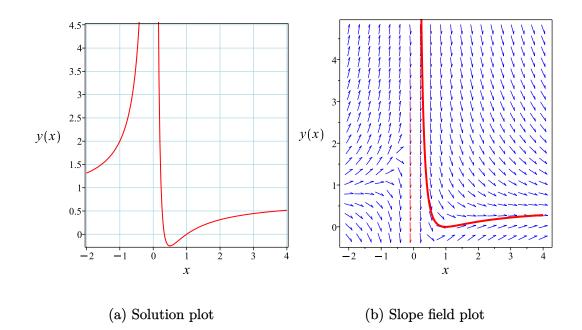
Substituting  $c_1$  found above in the general solution gives

$$y = \frac{3x^2 - 4x + 1}{4x^2}$$

Summary

The solution(s) found are the following

$$y = \frac{3x^2 - 4x + 1}{4x^2} \tag{1}$$



Verification of solutions

$$y = \frac{3x^2 - 4x + 1}{4x^2}$$

Verified OK.

#### 2.2.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{2xy - x + 1}{x^2}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
 (A)

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi$ ,  $\eta$ 

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f(rac{y}{x})$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y'=rac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x) y + f_2(x) y^2$	0	$e^{-\int f_1 dx}$

Table 21: Lie symmetry infinitesimal lookup table for known first order ODE's

The above table shows that

$$\xi(x,y) = 0$$
  

$$\eta(x,y) = \frac{1}{x^2}$$
(A1)

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where S(R). Since  $\xi = 0$  then in this special case

$$R = x$$

 ${\cal S}$  is found from

$$S = \int \frac{1}{\eta} dy$$
$$= \int \frac{1}{\frac{1}{x^2}} dy$$

Which results in

$$S = y x^2$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \tag{2}$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = -\frac{2xy - x + 1}{x^2}$$

Evaluating all the partial derivatives gives

$$egin{aligned} R_x &= 1 \ R_y &= 0 \ S_x &= 2xy \ S_y &= x^2 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = x - 1 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R - 1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = \frac{1}{2}R^2 - R + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$yx^2 = \frac{1}{2}x^2 - x + c_1$$

Which simplifies to

$$yx^2 = \frac{1}{2}x^2 - x + c_1$$

Which gives

$$y = \frac{x^2 + 2c_1 - 2x}{2x^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	$ODE  ext{ in canonical coordinates} \ (R,S)$
$\frac{dy}{dx} = -\frac{2xy - x + 1}{x^2}$	$R = x$ $S = y x^2$	$\frac{dS}{dR} = R - 1$

Initial conditions are used to solve for  $c_1$ . Substituting x = 1 and y = 0 in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{1}{2} + c_1$$

$$c_1 = \frac{1}{2}$$

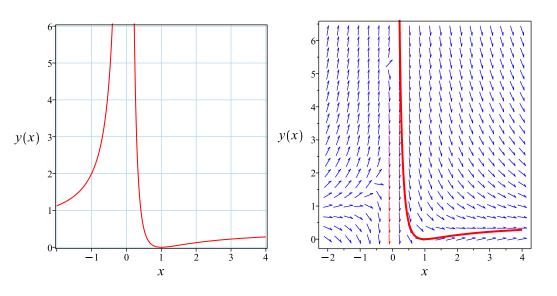
Substituting  $c_1$  found above in the general solution gives

$$y=\frac{x^2-2x+1}{2x^2}$$

Summary

The solution(s) found are the following

$$y = \frac{x^2 - 2x + 1}{2x^2} \tag{1}$$



(a) Solution plot

(b) Slope field plot

Verification of solutions

$$y = \frac{x^2 - 2x + 1}{2x^2}$$

Verified OK.

#### 2.2.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x,y) + N(x,y)\frac{dy}{dx} = 0$$
 (A)

We assume there exists a function  $\phi(x, y) = c$  where c is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t. x gives

$$\frac{d}{dx}\phi(x,y) = 0$$

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y}\frac{dy}{dx} = 0$$
(B)

Comparing (A,B) shows that

Hence

$$rac{\partial \phi}{\partial x} = M$$
  
 $rac{\partial \phi}{\partial y} = N$ 

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x,y) dx + N(x,y) dy = 0$$
(1A)

Therefore

$$(x^{2}) dy = (-2xy + x - 1) dx$$
$$(2xy - x + 1) dx + (x^{2}) dy = 0$$
(2A)

Comparing (1A) and (2A) shows that

$$M(x, y) = 2xy - x + 1$$
$$N(x, y) = x^{2}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y}(2xy - x + 1)$$
$$= 2x$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x} (x^2)$$
$$= 2x$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is <u>exact</u> The following equations are now set up to solve for the function  $\phi(x, y)$ 

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int 2xy - x + 1 dx$$

$$\phi = \frac{(2y - 1)x^2}{2} + x + f(y)$$
(3)

Where f(y) is used for the constant of integration since  $\phi$  is a function of both x and y. Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x^2 + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = x^2$ . Therefore equation (4) becomes

$$x^2 = x^2 + f'(y) (5)$$

Solving equation (5) for f'(y) gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for f(y) into equation (3) gives  $\phi$ 

$$\phi = \frac{(2y-1)x^2}{2} + x + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \frac{(2y-1)x^2}{2} + x$$

The solution becomes

$$y = \frac{x^2 + 2c_1 - 2x}{2x^2}$$

Initial conditions are used to solve for  $c_1$ . Substituting x = 1 and y = 0 in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{1}{2} + c_1$$
$$c_1 = \frac{1}{2}$$

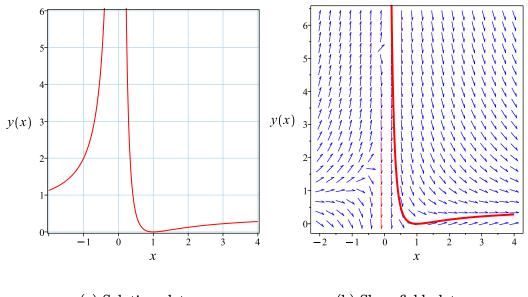
Substituting  $c_1$  found above in the general solution gives

$$y = \frac{x^2 - 2x + 1}{2x^2}$$

Summary

The solution(s) found are the following

$$y = \frac{x^2 - 2x + 1}{2x^2} \tag{1}$$



(a) Solution plot

(b) Slope field plot

Verification of solutions

$$y = \frac{x^2 - 2x + 1}{2x^2}$$

Verified OK.

#### 2.2.6 Maple step by step solution

Let's solve

 $[x^2y' + 2xy = x - 1, y(1) = 0]$ 

- Highest derivative means the order of the ODE is 1 y'
- Isolate the derivative

$$y' = -\frac{2y}{x} + \frac{x-1}{x^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE  $y' + \frac{2y}{x} = \frac{x-1}{x^2}$
- The ODE is linear; multiply by an integrating factor  $\mu(x)$  $\mu(x) \left(y' + \frac{2y}{x}\right) = \frac{\mu(x)(x-1)}{x^2}$
- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$  $\mu(x)\left(y'+\frac{2y}{x}\right) = \mu'(x)y + \mu(x)y'$

- Isolate  $\mu'(x)$  $\mu'(x) = \frac{2\mu(x)}{x}$
- Solve to find the integrating factor  $\mu(x) = x^2$
- Integrate both sides with respect to x $\int \left(\frac{d}{dx}(\mu(x) y)\right) dx = \int \frac{\mu(x)(x-1)}{x^2} dx + c_1$
- Evaluate the integral on the lhs

$$\mu(x) y = \int \frac{\mu(x)(x-1)}{x^2} dx + c_1$$

• Solve for y

$$y = rac{\int rac{\mu(x)(x-1)}{x^2} dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = x^2$  $y = rac{\int (x-1)dx + c_1}{x^2}$
- Evaluate the integrals on the rhs

$$y = rac{rac{1}{2}x^2 - x + c_1}{x^2}$$

• Simplify

$$y = \frac{x^2 + 2c_1 - 2x}{2x^2}$$

• Use initial condition y(1) = 0

$$0 = -\frac{1}{2} + c_1$$

• Solve for  $c_1$ 

$$c_1 = \frac{1}{2}$$

• Substitute  $c_1 = \frac{1}{2}$  into general solution and simplify

$$y = \frac{(x-1)^2}{2x^2}$$

• Solution to the IVP

$$y = \frac{(x-1)^2}{2x^2}$$

#### Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

✓ Solution by Maple Time used: 0.015 (sec). Leaf size: 14

dsolve([x^2\*diff(y(x),x)+2\*x\*y(x)-x+1=0,y(1) = 0],y(x), singsol=all)

$$y(x) = \frac{(x-1)^2}{2x^2}$$

Solution by Mathematica  $\checkmark$ 

Time used: 0.028 (sec). Leaf size: 17

DSolve[{x^2\*y'[x]+2\*x\*y[x]-x+1==0,{y[1]==0}},y[x],x,IncludeSingularSolutions -> True]

$$y(x) \rightarrow \frac{(x-1)^2}{2x^2}$$

# 2.3 problem 10.3.4

2.3.1	Existence and uniqueness analysis	138
2.3.2	Solving as linear ode	139
2.3.3	Solving as first order ode lie symmetry lookup ode	141
2.3.4	Solving as exact ode	145
2.3.5	Maple step by step solution	149

Internal problem ID [5057] Internal file name [OUTPUT/4550\_Sunday\_June\_05\_2022\_03\_00\_43\_PM\_3682306/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.3, ODEs with variable Coefficients.
First order. page 315
Problem number: 10.3.4.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[[\_linear, `class A`]]

$$y + y' = (x+1)^2$$

With initial conditions

$$[y(0) = 0]$$

#### 2.3.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = 1$$
$$q(x) = (x+1)^2$$

Hence the ode is

$$y + y' = (x+1)^2$$

The domain of p(x) = 1 is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is inside this domain. The domain of  $q(x) = (x+1)^2$  is

$$\{-\infty < x < \infty\}$$

And the point  $x_0 = 0$  is also inside this domain. Hence solution exists and is unique.

### 2.3.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor  $\mu$  is

$$\mu = e^{\int 1dx}$$
$$= e^x$$

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu y) = (\mu) \left( (x+1)^2 \right)$$
$$\frac{\mathrm{d}}{\mathrm{d}x}(y \,\mathrm{e}^x) = (\mathrm{e}^x) \left( (x+1)^2 \right)$$
$$\mathrm{d}(y \,\mathrm{e}^x) = \left( (x+1)^2 \,\mathrm{e}^x \right) \,\mathrm{d}x$$

Integrating gives

$$y e^{x} = \int (x+1)^{2} e^{x} dx$$
$$y e^{x} = (x^{2}+1) e^{x} + c_{1}$$

Dividing both sides by the integrating factor  $\mu = e^x$  results in

$$y = e^{-x} (x^2 + 1) e^x + c_1 e^{-x}$$

which simplifies to

$$y = x^2 + 1 + c_1 e^{-x}$$

Initial conditions are used to solve for  $c_1$ . Substituting x = 0 and y = 0 in the above solution gives an equation to solve for the constant of integration.

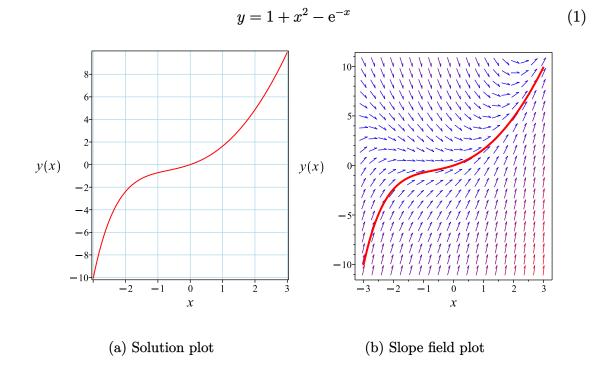
$$0 = c_1 + 1$$

$$c_1 = -1$$

Substituting  $c_1$  found above in the general solution gives

$$y = 1 + x^2 - e^{-x}$$

 $\frac{\text{Summary}}{\text{The solution(s) found are the following}}$ 



Verification of solutions

$$y = 1 + x^2 - \mathrm{e}^{-x}$$

Verified OK.

#### 2.3.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = x^{2} + 2x - y + 1$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
 (A)

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi$ ,  $\eta$ 

Table 24: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f(rac{y}{x})$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y'=rac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x) y + f_2(x) y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned} \xi(x,y) &= 0\\ \eta(x,y) &= \mathrm{e}^{-x} \end{aligned} \tag{A1}$$

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where S(R). Since  $\xi = 0$  then in this special case

$$R = x$$

 ${\cal S}$  is found from

$$S = \int rac{1}{\eta} dy \ = \int rac{1}{\mathrm{e}^{-x}} dy$$

Which results in

 $S = y e^x$ 

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \tag{2}$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = x^2 + 2x - y + 1$$

Evaluating all the partial derivatives gives

$$R_x = 1$$
$$R_y = 0$$
$$S_x = y e^x$$
$$S_y = e^x$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = (x+1)^2 e^x \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \left(R+1\right)^2 \mathbf{e}^R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = (R^2 + 1) e^R + c_1$$
(4)

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\mathrm{e}^{x}y = \left(x^{2}+1\right)\mathrm{e}^{x}+c_{1}$$

Which simplifies to

$$\mathrm{e}^{x}y = \left(x^{2}+1\right)\mathrm{e}^{x}+c_{1}$$

Which gives

$$y = \left(x^2 \mathrm{e}^x + \mathrm{e}^x + c_1\right) \mathrm{e}^{-x}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = x^2 + 2x - y + 1$	$R = x$ $S = y e^x$	$\frac{dS}{dR} = (R+1)^2 e^R$

Initial conditions are used to solve for  $c_1$ . Substituting x = 0 and y = 0 in the above solution gives an equation to solve for the constant of integration.

 $0 = c_1 + 1$ 

 $c_1 = -1$ 

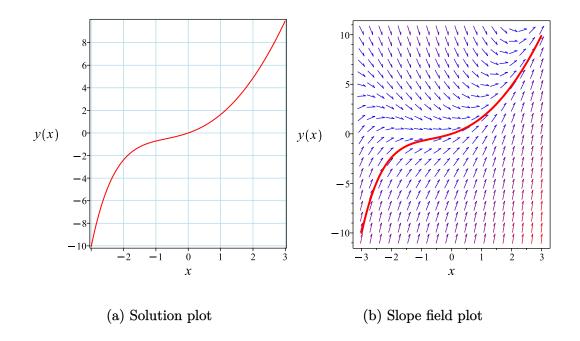
Substituting  $c_1$  found above in the general solution gives

$$y = 1 + x^2 - \mathrm{e}^{-x}$$

Summary

The solution(s) found are the following

$$y = 1 + x^2 - e^{-x} \tag{1}$$



Verification of solutions

$$y = 1 + x^2 - e^{-x}$$

Verified OK.

## 2.3.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x,y) + N(x,y)\frac{dy}{dx} = 0$$
 (A)

We assume there exists a function  $\phi(x, y) = c$  where c is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t. x gives

$$\frac{d}{dx}\phi(x,y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y}\frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$
$$\frac{\partial \phi}{\partial y} = N$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x,y) dx + N(x,y) dy = 0$$
(1A)

Therefore

$$dy = (-y + (x+1)^2) dx$$
$$(y - (x+1)^2) dx + dy = 0$$
(2A)

Comparing (1A) and (2A) shows that

$$M(x, y) = y - (x + 1)^{2}$$
$$N(x, y) = 1$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$rac{\partial M}{\partial y} = rac{\partial}{\partial y} (y - (x+1)^2)$$
  
= 1

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x}(1)$$
$$= 0$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is <u>not exact</u>. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$A = \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right)$$
$$= 1((1) - (0))$$
$$= 1$$

Since A does not depend on y, then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\mu = e^{\int A \, \mathrm{d}x}$$
$$= e^{\int 1 \, \mathrm{d}x}$$

The result of integrating gives

 $\begin{aligned} \mu &= e^x \\ &= e^x \end{aligned}$ 

M and N are multiplied by this integrating factor, giving new M and new N which are called  $\overline{M}$  and  $\overline{N}$  for now so not to confuse them with the original M and N.

$$\overline{M} = \mu M$$
  
=  $e^x (y - (x+1)^2)$   
=  $-e^x (x^2 + 2x - y + 1)$ 

And

$$egin{aligned} N &= \mu N \ &= \mathrm{e}^x(1) \ &= \mathrm{e}^x \end{aligned}$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$\overline{M} + \overline{N}\frac{\mathrm{d}y}{\mathrm{d}x} = 0$$
$$\left(-\mathrm{e}^{x}\left(x^{2} + 2x - y + 1\right)\right) + \left(\mathrm{e}^{x}\right)\frac{\mathrm{d}y}{\mathrm{d}x} = 0$$

The following equations are now set up to solve for the function  $\phi(x, y)$ 

$$\frac{\partial \phi}{\partial x} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \overline{M} dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -e^x (x^2 + 2x - y + 1) dx$$

$$\phi = -(x^2 - y + 1) e^x + f(y)$$
(3)

Where f(y) is used for the constant of integration since  $\phi$  is a function of both x and y. Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^x + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = e^x$ . Therefore equation (4) becomes

$$\mathbf{e}^x = \mathbf{e}^x + f'(y) \tag{5}$$

Solving equation (5) for f'(y) gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for f(y) into equation (3) gives  $\phi$ 

$$\phi = -\left(x^2 - y + 1\right)e^x + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\left(x^2 - y + 1\right) \mathrm{e}^x$$

The solution becomes

$$y = \left(x^2 \mathbf{e}^x + \mathbf{e}^x + c_1\right) \mathbf{e}^{-x}$$

Initial conditions are used to solve for  $c_1$ . Substituting x = 0 and y = 0 in the above solution gives an equation to solve for the constant of integration.

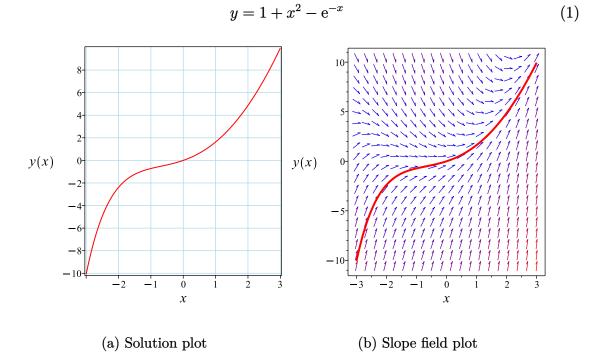
$$0 = c_1 + 1$$
  
 $c_1 = -1$ 

Substituting  $c_1$  found above in the general solution gives

$$y = 1 + x^2 - e^{-x}$$

### Summary

The solution(s) found are the following



### Verification of solutions

$$y = 1 + x^2 - e^{-x}$$

Verified OK.

# 2.3.5 Maple step by step solution

Let's solve

$$[y + y' = (x + 1)^2, y(0) = 0]$$

• Highest derivative means the order of the ODE is 1

y'

• Isolate the derivative

 $y' = -y + (x+1)^2$ 

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE  $y + y' = (x + 1)^2$
- The ODE is linear; multiply by an integrating factor  $\mu(x)$

 $\mu(x) (y + y') = \mu(x) (x + 1)^2$ 

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x) y)$  $\mu(x) (y + y') = \mu'(x) y + \mu(x) y'$
- Isolate  $\mu'(x)$  $\mu'(x) = \mu(x)$
- Solve to find the integrating factor  $\mu(x) = e^x$
- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x) y)\right) dx = \int \mu(x) (x+1)^2 dx + c_1$$

- Evaluate the integral on the lhs $\mu(x) \, y = \int \mu(x) \, (x+1)^2 \, dx + c_1$
- Solve for y

$$y = \frac{\int \mu(x)(x+1)^2 dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = e^x$  $y = \frac{\int (x+1)^2 e^x dx + c_1}{e^x}$
- Evaluate the integrals on the rhs  $y = \frac{(x^2+1)\mathrm{e}^x + c_1}{\mathrm{e}^x}$
- Simplify

$$y = x^2 + 1 + c_1 e^{-x}$$

• Use initial condition y(0) = 0

$$0 = c_1 + 1$$

• Solve for  $c_1$ 

$$c_1 = -1$$

- Substitute  $c_1 = -1$  into general solution and simplify  $y = 1 + x^2 - e^{-x}$
- Solution to the IVP

$$y = 1 + x^2 - e^{-x}$$

### Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`</pre>
```

Solution by Maple Time used: 0.0 (sec). Leaf size: 15

 $dsolve([diff(y(x),x)+y(x)=(x+1)^2,y(0) = 0],y(x), singsol=all)$ 

$$y(x) = x^2 + 1 - e^{-x}$$

✓ Solution by Mathematica

Time used: 0.102 (sec). Leaf size: 17

DSolve[{y'[x]+y[x]==(x+1)^2,{y[0]==0}},y[x],x,IncludeSingularSolutions -> True]

$$y(x) \to x^2 - e^{-x} + 1$$

# 2.4 problem 10.3.5

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Internal problem	ID [5058]	
Internal file name	[OUTPUT/4551_Sunday_June_05_2022_03_00_45_PM_20469628/index.	tex]
Book: Basic Tra	aining in Mathematics. By R. Shankar. Plenum Press. NY. 1995	

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.3, ODEs with variable Coefficients.
First order. page 315
Problem number: 10.3.5.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "linear", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_linear]

$$x^2y' + 2xy = \sinh\left(x\right)$$

With initial conditions

[y(1) = 2]

### 2.4.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{2}{x}$$
$$q(x) = \frac{\sinh(x)}{x^2}$$

Hence the ode is

$$y' + \frac{2y}{x} = \frac{\sinh\left(x\right)}{x^2}$$

The domain of  $p(x) = \frac{2}{x}$  is

$$\{x < 0 \lor 0 < x\}$$

And the point  $x_0 = 1$  is inside this domain. The domain of  $q(x) = \frac{\sinh(x)}{x^2}$  is

$$\{x < 0 \lor 0 < x\}$$

And the point  $x_0 = 1$  is also inside this domain. Hence solution exists and is unique.

## 2.4.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor  $\mu$  is

$$\mu = e^{\int \frac{2}{x} dx}$$
$$= x^2$$

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu y) = (\mu) \left(\frac{\sinh(x)}{x^2}\right)$$
$$\frac{\mathrm{d}}{\mathrm{d}x}(y x^2) = (x^2) \left(\frac{\sinh(x)}{x^2}\right)$$
$$\mathrm{d}(y x^2) = \sinh(x) \mathrm{d}x$$

Integrating gives

$$y x^{2} = \int \sinh(x) dx$$
$$y x^{2} = \cosh(x) + c_{1}$$

Dividing both sides by the integrating factor  $\mu = x^2$  results in

$$y = \frac{\cosh\left(x\right)}{x^2} + \frac{c_1}{x^2}$$

which simplifies to

$$y = \frac{\cosh\left(x\right) + c_1}{x^2}$$

Initial conditions are used to solve for  $c_1$ . Substituting x = 1 and y = 2 in the above solution gives an equation to solve for the constant of integration.

$$2 = \cosh\left(1\right) + c_1$$

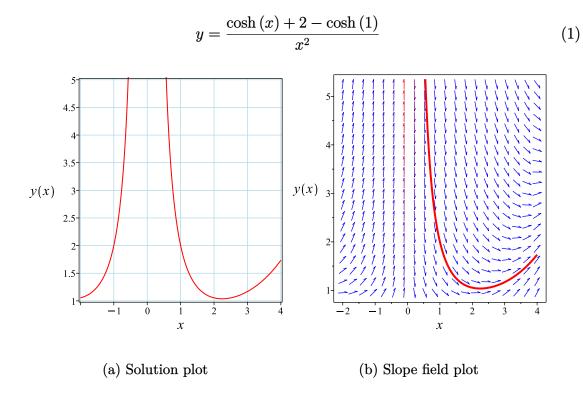
$$c_1 = -\cosh\left(1\right) + 2$$

Substituting  $c_1$  found above in the general solution gives

$$y = \frac{\cosh\left(x\right) + 2 - \cosh\left(1\right)}{x^2}$$

Summary

The solution(s) found are the following



Verification of solutions

$$y = \frac{\cosh\left(x\right) + 2 - \cosh\left(1\right)}{x^2}$$

Verified OK.

## 2.4.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-2xy + \sinh(x)}{x^2}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
 (A)

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi$ ,  $\eta$ 

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f\left(rac{y}{x} ight)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx - h(x)}}{g(x)}$
polynomial type ode	$y'=rac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x) y + f_2(x) y^2$	0	$e^{-\int f_1 dx}$

Table 27: Lie symmetry infinitesimal lookup table for known first order ODE's

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The above table shows that

$$\xi(x,y) = 0$$
  

$$\eta(x,y) = \frac{1}{x^2}$$
(A1)

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where S(R). Since  $\xi = 0$  then in this special case

R = x

 ${\cal S}$  is found from

$$S = \int rac{1}{\eta} dy \ = \int rac{1}{rac{1}{x^2}} dy$$

Which results in

 $S = y x^2$ 

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \tag{2}$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = \frac{-2xy + \sinh\left(x\right)}{x^2}$$

Evaluating all the partial derivatives gives

$$egin{aligned} R_x &= 1 \ R_y &= 0 \ S_x &= 2xy \ S_y &= x^2 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \sinh\left(x\right) \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \sinh\left(R\right)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = \cosh\left(R\right) + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$yx^2 = \cosh\left(x\right) + c_1$$

Which simplifies to

$$yx^2 = \cosh\left(x\right) + c_1$$

Which gives

$$y = \frac{\cosh\left(x\right) + c_1}{x^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	$ODE  ext{ in canonical coordinates} \ (R,S)$
$\frac{dy}{dx} = \frac{-2xy + \sinh(x)}{x^2}$	$R = x$ $S = y x^2$	$\frac{dS}{dR} = \sinh(R)$

Initial conditions are used to solve for  $c_1$ . Substituting x = 1 and y = 2 in the above solution gives an equation to solve for the constant of integration.

 $2 = \cosh\left(1\right) + c_1$ 

$$c_1 = -\cosh\left(1\right) + 2$$

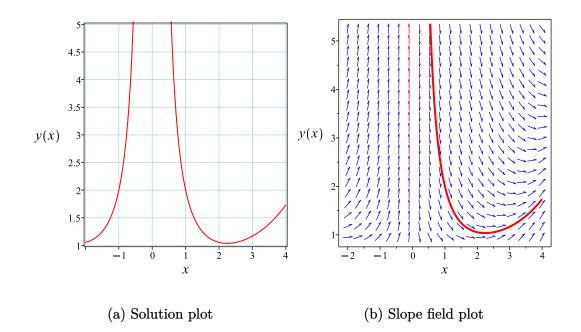
Substituting  $c_1$  found above in the general solution gives

$$y = \frac{\cosh\left(x\right) + 2 - \cosh\left(1\right)}{x^2}$$

Summary

The solution(s) found are the following

$$y = \frac{\cosh{(x)} + 2 - \cosh{(1)}}{x^2}$$
(1)



Verification of solutions

$$y = \frac{\cosh\left(x\right) + 2 - \cosh\left(1\right)}{x^2}$$

Verified OK.

### 2.4.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x,y) + N(x,y)\frac{dy}{dx} = 0$$
 (A)

We assume there exists a function  $\phi(x, y) = c$  where c is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t. x gives

$$\frac{d}{dx}\phi(x,y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y}\frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$
$$\frac{\partial \phi}{\partial y} = N$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x,y) dx + N(x,y) dy = 0$$
(1A)

Therefore

$$(x^{2}) dy = (-2xy + \sinh(x)) dx$$
$$(2xy - \sinh(x)) dx + (x^{2}) dy = 0$$
(2A)

Comparing (1A) and (2A) shows that

$$M(x, y) = 2xy - \sinh(x)$$
  
 $N(x, y) = x^2$ 

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y} (2xy - \sinh(x))$$
$$= 2x$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x} (x^2)$$
$$= 2x$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is <u>exact</u> The following equations are now set up to solve for the function  $\phi(x, y)$ 

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int 2xy - \sinh(x) dx$$

$$\phi = y x^2 - \cosh(x) + f(y)$$
(3)

Where f(y) is used for the constant of integration since  $\phi$  is a function of both x and y. Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x^2 + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = x^2$ . Therefore equation (4) becomes

$$x^2 = x^2 + f'(y) (5)$$

Solving equation (5) for f'(y) gives

f'(y) = 0

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for f(y) into equation (3) gives  $\phi$ 

$$\phi = y x^2 - \cosh\left(x\right) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = y \, x^2 - \cosh\left(x\right)$$

The solution becomes

$$y = \frac{\cosh\left(x\right) + c_1}{x^2}$$

Initial conditions are used to solve for  $c_1$ . Substituting x = 1 and y = 2 in the above solution gives an equation to solve for the constant of integration.

$$2 = \cosh\left(1\right) + c_1$$

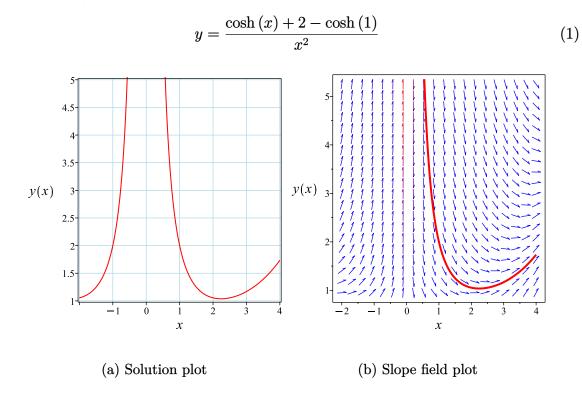
$$c_1 = -\cosh\left(1\right) + 2$$

Substituting  $c_1$  found above in the general solution gives

$$y = \frac{\cosh\left(x\right) + 2 - \cosh\left(1\right)}{x^2}$$

Summary

The solution(s) found are the following



Verification of solutions

$$y = \frac{\cosh\left(x\right) + 2 - \cosh\left(1\right)}{x^2}$$

Verified OK.

#### 2.4.5Maple step by step solution

Let's solve

 $[x^{2}y' + 2xy = \sinh(x), y(1) = 2]$ 

Highest derivative means the order of the ODE is 1 *y*′

Isolate the derivative

$$y' = -\frac{2y}{x} + \frac{\sinh(x)}{x^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE .  $y' + \frac{2y}{x} = \frac{\sinh(x)}{x^2}$
- The ODE is linear; multiply by an integrating factor  $\mu(x)$ .  $\mu(x)\left(y'+rac{2y}{x}
  ight)=rac{\mu(x)\sinh(x)}{x^2}$
- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$

$$\mu(x)\left(y' + \frac{2y}{x}\right) = \mu'(x)y + \mu(x)y'$$

- Isolate  $\mu'(x)$  $\mu'(x) = \frac{2\mu(x)}{x}$
- Solve to find the integrating factor

$$\mu(x) = x^2$$

Integrate both sides with respect to x

$$\int \left( rac{d}{dx}(\mu(x)\,y) 
ight) dx = \int rac{\mu(x)\sinh(x)}{x^2} dx + c_1$$

Evaluate the integral on the lhs

$$\mu(x) \, y = \int rac{\mu(x) \sinh(x)}{x^2} dx + c_1$$

Solve for y

$$y=rac{\intrac{\mu(x)\sinh(x)}{x^2}dx+c_1}{\mu(x)}$$

- Substitute  $\mu(x) = x^2$ •  $y = rac{\int \sinh(x)dx + c_1}{x^2}$
- Evaluate the integrals on the rhs  $y = \frac{\cosh(x) + c_1}{x^2}$
- Use initial condition y(1) = 2•

$$2 = \cosh\left(1\right) + c_1$$

• Solve for  $c_1$ 

 $c_1 = -\cosh\left(1\right) + 2$ 

- Substitute  $c_1 = -\cosh(1) + 2$  into general solution and simplify  $y = \frac{\cosh(x) + 2 - \cosh(1)}{x^2}$
- Solution to the IVP  $y = \frac{\cosh(x) + 2 - \cosh(1)}{x^2}$

Maple trace

`Methods for first order ODEs: --- Trying classification methods --trying a quadrature trying 1st order linear <- 1st order linear successful`</pre>

Solution by Maple Time used: 0.031 (sec). Leaf size: 16

 $dsolve([x^2*diff(y(x),x)+2*x*y(x)=sinh(x),y(1) = 2],y(x), singsol=all)$ 

$$y(x) = \frac{\cosh\left(x\right) + 2 - \cosh\left(1\right)}{x^2}$$

Solution by Mathematica

Time used: 0.038 (sec). Leaf size: 17

DSolve[{x^2\*y'[x]+2\*x\*y[x]==Sinh[x],{y[1]==2}},y[x],x,IncludeSingularSolutions -> True]

$$y(x) 
ightarrow rac{\cosh(x) + 2 - \cosh(1)}{x^2}$$

# 2.5 problem 10.3.6

2.5.1	Solving as linear ode	165
2.5.2	Solving as first order ode lie symmetry lookup ode	167
2.5.3	Solving as exact ode	171
2.5.4	Maple step by step solution	176

Internal problem ID [5059] Internal file name [OUTPUT/4552\_Sunday\_June\_05\_2022\_03\_00\_46\_PM\_11495574/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.3, ODEs with variable Coefficients.
First order. page 315
Problem number: 10.3.6.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_linear]

$$y' + \frac{y}{1-x} = x^2 - 2x$$

### 2.5.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{1}{x-1}$$
$$q(x) = x(-2+x)$$

Hence the ode is

$$y' - \frac{y}{x-1} = x(-2+x)$$

The integrating factor  $\mu$  is

$$\mu = e^{\int -\frac{1}{x-1}dx}$$
$$= \frac{1}{x-1}$$

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu y) = (\mu) \left(x(-2+x)\right)$$
$$\frac{\mathrm{d}}{\mathrm{d}x}\left(\frac{y}{x-1}\right) = \left(\frac{1}{x-1}\right) \left(x(-2+x)\right)$$
$$\mathrm{d}\left(\frac{y}{x-1}\right) = \left(\frac{x(-2+x)}{x-1}\right) \,\mathrm{d}x$$

Integrating gives

$$\frac{y}{x-1} = \int \frac{x(-2+x)}{x-1} \, \mathrm{d}x$$
$$\frac{y}{x-1} = \frac{x^2}{2} - x - \ln(x-1) + c_1$$

Dividing both sides by the integrating factor  $\mu = \frac{1}{x-1}$  results in

$$y = (x-1)\left(\frac{x^2}{2} - x - \ln(x-1)\right) + c_1(x-1)$$

which simplifies to

$$y = \frac{(x-1)(x^2 - 2x - 2\ln(x-1) + 2c_1)}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{(x-1)\left(x^2 - 2x - 2\ln\left(x - 1\right) + 2c_1\right)}{2} \tag{1}$$

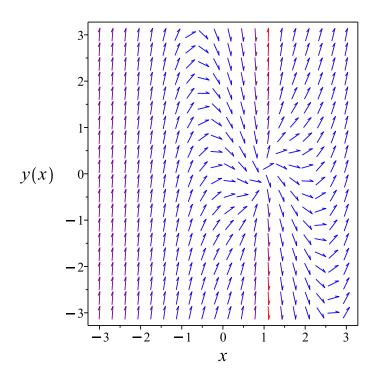


Figure 25: Slope field plot

Verification of solutions

$$y = \frac{(x-1)(x^2 - 2x - 2\ln(x-1) + 2c_1)}{2}$$

Verified OK.

## 2.5.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x^3 - 3x^2 + 2x + y}{x - 1}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
 (A)

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi$ ,  $\eta$ 

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f(rac{y}{x})$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x + b_1y + c_1}{a_2x + b_2y + c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x) y + f_2(x) y^2$	0	$e^{-\int f_1 dx}$

Table 30: Lie symmetry infinitesimal lookup table for known first order ODE's

The above table shows that

$$\xi(x, y) = 0$$
  

$$\eta(x, y) = x - 1$$
(A1)

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where S(R). Since  $\xi = 0$  then in this special case

R = x

 ${\cal S}$  is found from

$$S = \int \frac{1}{\eta} dy$$
$$= \int \frac{1}{x - 1} dy$$

Which results in

$$S = \frac{y}{x-1}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \tag{2}$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = \frac{x^3 - 3x^2 + 2x + y}{x - 1}$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 0$$

$$S_x = -\frac{y}{(x-1)^2}$$

$$S_y = \frac{1}{x-1}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{x(-2+x)}{x-1} \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{R(-2+R)}{R-1}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = \frac{R^2}{2} - R - \ln(R - 1) + c_1 \tag{4}$$

Which gives

$$y = -\frac{(x-1)(-x^2 + 2\ln(x-1) - 2c_1 + 2x)}{2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	${ m ODE} { m in \ canonical \ coordinates} \ (R,S)$
$\frac{dy}{dx} = \frac{x^3 - 3x^2 + 2x + y}{x - 1}$	$R = x$ $S = \frac{y}{x - 1}$	$\frac{dS}{dR} = \frac{R(-2+R)}{R-1}$

Summary

The solution(s) found are the following

$$y = -\frac{(x-1)\left(-x^2 + 2\ln(x-1) - 2c_1 + 2x\right)}{2} \tag{1}$$

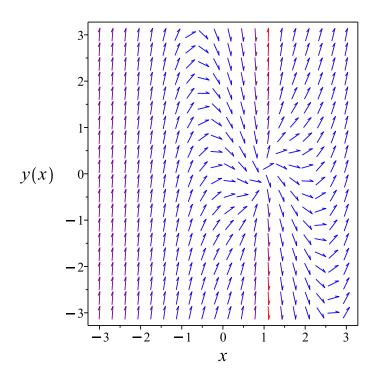


Figure 26: Slope field plot

Verification of solutions

$$y = -\frac{(x-1)(-x^2 + 2\ln(x-1) - 2c_1 + 2x)}{2}$$

Verified OK.

## 2.5.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x,y) + N(x,y)\frac{dy}{dx} = 0$$
(A)

We assume there exists a function  $\phi(x, y) = c$  where c is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t. x gives

.

$$\frac{d}{dx}\phi(x,y) = 0$$

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y}\frac{dy}{dx} = 0$$
(B)

Hence

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$
$$\frac{\partial \phi}{\partial y} = N$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0$$
(1A)

Therefore

$$dy = \left(-\frac{y}{1-x} - 2x + x^2\right) dx$$
$$\left(\frac{y}{1-x} + 2x - x^2\right) dx + dy = 0$$
(2A)

Comparing (1A) and (2A) shows that

$$M(x,y) = \frac{y}{1-x} + 2x - x^2$$
$$N(x,y) = 1$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y} \left( \frac{y}{1-x} + 2x - x^2 \right)$$
$$= \frac{1}{1-x}$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x}(1)$$
$$= 0$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is <u>not exact</u>. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$A = \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right)$$
$$= 1 \left( \left( \frac{1}{1-x} \right) - (0) \right)$$
$$= \frac{1}{1-x}$$

Since A does not depend on y, then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\mu = e^{\int A \, \mathrm{d}x}$$
$$= e^{\int \frac{1}{1-x} \, \mathrm{d}x}$$

The result of integrating gives

$$\mu = e^{-\ln(1-x)}$$
$$= \frac{1}{1-x}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called  $\overline{M}$  and  $\overline{N}$  for now so not to confuse them with the original M and N.

$$\overline{M} = \mu M$$
$$= \frac{1}{1-x} \left( \frac{y}{1-x} + 2x - x^2 \right)$$
$$= \frac{x^3 - 3x^2 + 2x + y}{\left(x-1\right)^2}$$

And

$$N = \mu N$$
$$= \frac{1}{1-x}(1)$$
$$= \frac{1}{1-x}$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$\overline{M} + \overline{N}\frac{\mathrm{d}y}{\mathrm{d}x} = 0$$
$$\left(\frac{x^3 - 3x^2 + 2x + y}{\left(x - 1\right)^2}\right) + \left(\frac{1}{1 - x}\right)\frac{\mathrm{d}y}{\mathrm{d}x} = 0$$

The following equations are now set up to solve for the function  $\phi(x, y)$ 

$$\frac{\partial \phi}{\partial x} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \overline{M} dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int \frac{x^3 - 3x^2 + 2x + y}{(x - 1)^2} dx$$

$$\phi = \frac{x^2}{2} - x - \ln(x - 1) - \frac{y}{x - 1} + f(y)$$
(3)

Where f(y) is used for the constant of integration since  $\phi$  is a function of both x and y. Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = -\frac{1}{x-1} + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{1}{1-x}$ . Therefore equation (4) becomes

$$\frac{1}{1-x} = -\frac{1}{x-1} + f'(y) \tag{5}$$

Solving equation (5) for f'(y) gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for f(y) into equation (3) gives  $\phi$ 

$$\phi = \frac{x^2}{2} - x - \ln(x - 1) - \frac{y}{x - 1} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \frac{x^2}{2} - x - \ln(x - 1) - \frac{y}{x - 1}$$

The solution becomes

$$y = -\frac{\left(-x^2 + 2\ln(x-1) + 2c_1 + 2x\right)(x-1)}{2}$$

Summary

The solution(s) found are the following

$$y = -\frac{(-x^{2} + 2\ln(x - 1) + 2c_{1} + 2x)(x - 1)}{2}$$
(1)

Figure 27: Slope field plot

Verification of solutions

$$y = -\frac{(-x^2 + 2\ln(x-1) + 2c_1 + 2x)(x-1)}{2}$$

Verified OK.

### 2.5.4 Maple step by step solution

Let's solve

$$y' + \frac{y}{1-x} = x^2 - 2x$$

• Highest derivative means the order of the ODE is 1

y'

• Isolate the derivative

$$y' = \frac{y}{x-1} + x^2 - 2x$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE  $y' \frac{y}{x-1} = x^2 2x$
- The ODE is linear; multiply by an integrating factor  $\mu(x)$  $\mu(x) \left(y' - \frac{y}{x-1}\right) = \mu(x) \left(x^2 - 2x\right)$
- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x) y)$  $\mu(x) \left(y' - \frac{y}{x-1}\right) = \mu'(x) y + \mu(x) y'$
- Isolate  $\mu'(x)$  $\mu'(x) = -\frac{\mu(x)}{x-1}$
- Solve to find the integrating factor  $\mu(x) = \frac{1}{x-1}$
- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x)y)\right) dx = \int \mu(x) \left(x^2 - 2x\right) dx + c_1$$

• Evaluate the integral on the lhs

$$\mu(x) y = \int \mu(x) (x^2 - 2x) dx + c_1$$

- Solve for y $y = \frac{\int \mu(x)(x^2 - 2x)dx + c_1}{\mu(x)}$
- Substitute  $\mu(x) = \frac{1}{x-1}$

$$y = (x-1)\left(\int rac{x^2-2x}{x-1}dx + c_1
ight)$$

• Evaluate the integrals on the rhs

$$y = (x - 1)\left(\frac{x^2}{2} - x - \ln(x - 1) + c_1\right)$$

• Simplify  $y = \frac{(x-1)(x^2-2x-2\ln(x-1)+2c_1)}{2}$ 

## Maple trace

`Methods for first order ODEs: --- Trying classification methods --trying a quadrature trying 1st order linear <- 1st order linear successful`</pre>

Solution by Maple Time used: 0.0 (sec). Leaf size: 25

 $dsolve(diff(y(x),x)+y(x)/(1-x)+2*x-x^2=0,y(x), singsol=all)$ 

$$y(x) = \frac{(x^2 - 2x - 2\ln(x - 1) + 2c_1)(x - 1)}{2}$$

✓ Solution by Mathematica

Time used: 0.034 (sec). Leaf size: 27

DSolve[y'[x]+y[x]/(1-x)+2\*x-x^2==0,y[x],x,IncludeSingularSolutions -> True]

$$y(x) \to (x-1)\left(\frac{1}{2}(x-1)^2 - \log(x-1) + c_1\right)$$

# 2.6 problem 10.3.7

2.6.1	Solving as linear ode	178
2.6.2	Solving as first order ode lie symmetry lookup ode	180
2.6.3	Solving as exact ode	184
2.6.4	Maple step by step solution	189

Internal problem ID [5060] Internal file name [OUTPUT/4553\_Sunday\_June\_05\_2022\_03\_00\_47\_PM\_35483196/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.3, ODEs with variable Coefficients.
First order. page 315
Problem number: 10.3.7.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_linear]

$$y' + \frac{y}{1-x} = x^2 - x$$

### 2.6.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{1}{x-1}$$
$$q(x) = x(x-1)$$

Hence the ode is

$$y' - \frac{y}{x-1} = x(x-1)$$

The integrating factor  $\mu$  is

$$\mu = e^{\int -\frac{1}{x-1}dx}$$
$$= \frac{1}{x-1}$$

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu y) = (\mu) \left(x(x-1)\right)$$
$$\frac{\mathrm{d}}{\mathrm{d}x}\left(\frac{y}{x-1}\right) = \left(\frac{1}{x-1}\right) \left(x(x-1)\right)$$
$$\mathrm{d}\left(\frac{y}{x-1}\right) = x \,\mathrm{d}x$$

Integrating gives

$$\frac{y}{x-1} = \int x \, \mathrm{d}x$$
$$\frac{y}{x-1} = \frac{x^2}{2} + c_1$$

Dividing both sides by the integrating factor  $\mu = \frac{1}{x-1}$  results in

$$y = \frac{x^2(x-1)}{2} + c_1(x-1)$$

which simplifies to

$$y = \frac{(x-1)(x^2 + 2c_1)}{2}$$

 $\frac{Summary}{The solution(s) found are the following}$ 

$$y = \frac{(x-1)(x^2+2c_1)}{2} \tag{1}$$

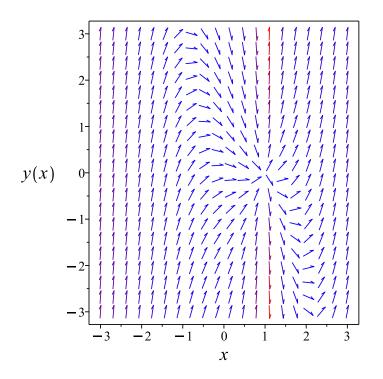


Figure 28: Slope field plot

Verification of solutions

$$y = \frac{(x-1)(x^2 + 2c_1)}{2}$$

Verified OK.

#### 2.6.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x^3 - 2x^2 + x + y}{x - 1}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
 (A)

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi$ ,  $\eta$ 

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f(rac{y}{x})$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x + b_1y + c_1}{a_2x + b_2y + c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x) y + f_2(x) y^2$	0	$e^{-\int f_1 dx}$

Table 33: Lie symmetry infinitesimal lookup table for known first order ODE's

The above table shows that

$$\xi(x, y) = 0$$
  

$$\eta(x, y) = x - 1$$
(A1)

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where S(R). Since  $\xi = 0$  then in this special case

R = x

 ${\cal S}$  is found from

$$S = \int \frac{1}{\eta} dy$$
$$= \int \frac{1}{x - 1} dy$$

Which results in

$$S = \frac{y}{x-1}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}$$
(2)

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = \frac{x^3 - 2x^2 + x + y}{x - 1}$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 0$$

$$S_x = -\frac{y}{(x-1)^2}$$

$$S_y = \frac{1}{x-1}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = x \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = \frac{R^2}{2} + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{y}{x-1} = \frac{x^2}{2} + c_1$$

Which simplifies to

$$\frac{y}{x-1} = \frac{x^2}{2} + c_1$$

Which gives

$$y = \frac{(x-1)(x^2 + 2c_1)}{2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	$ODE  ext{ in canonical coordinates} \ (R,S)$
$\frac{dy}{dx} = \frac{x^3 - 2x^2 + x + y}{x - 1}$	$R = x$ $S = \frac{y}{x - 1}$	$\frac{dS}{dR} = R$

# $\frac{Summary}{The solution(s) found are the following}$

$$y = \frac{(x-1)(x^2+2c_1)}{2} \tag{1}$$

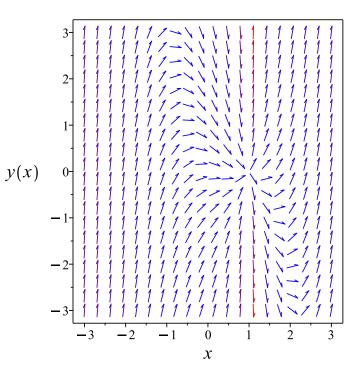


Figure 29: Slope field plot

Verification of solutions

$$y = \frac{(x-1)(x^2 + 2c_1)}{2}$$

Verified OK.

#### 2.6.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x,y) + N(x,y)\frac{dy}{dx} = 0$$
 (A)

We assume there exists a function  $\phi(x, y) = c$  where c is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t. x gives

$$\frac{d}{dx}\phi(x,y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y}\frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$
$$\frac{\partial \phi}{\partial y} = N$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0$$
(1A)

Therefore

$$dy = \left(-\frac{y}{1-x} - x + x^2\right) dx$$
$$\left(\frac{y}{1-x} + x - x^2\right) dx + dy = 0$$
(2A)

Comparing (1A) and (2A) shows that

$$M(x,y) = \frac{y}{1-x} + x - x^2$$
$$N(x,y) = 1$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$egin{aligned} &rac{\partial M}{\partial y} = rac{\partial}{\partial y}igg(rac{y}{1-x}+x-x^2igg)\ &= rac{1}{1-x} \end{aligned}$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x}(1)$$
$$= 0$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is <u>not exact</u>. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$A = \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right)$$
$$= 1 \left( \left( \frac{1}{1-x} \right) - (0) \right)$$
$$= \frac{1}{1-x}$$

Since A does not depend on y, then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\mu = e^{\int A \, \mathrm{d}x}$$
$$= e^{\int \frac{1}{1-x} \, \mathrm{d}x}$$

The result of integrating gives

$$\mu = e^{-\ln(1-x)}$$
$$= \frac{1}{1-x}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called  $\overline{M}$  and  $\overline{N}$  for now so not to confuse them with the original M and N.

$$\overline{M} = \mu M$$
$$= \frac{1}{1-x} \left( \frac{y}{1-x} + x - x^2 \right)$$
$$= \frac{x^3 - 2x^2 + x + y}{\left(x-1\right)^2}$$

And

$$N = \mu N$$
$$= \frac{1}{1-x}(1)$$
$$= \frac{1}{1-x}$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$\overline{M} + \overline{N}\frac{\mathrm{d}y}{\mathrm{d}x} = 0$$
$$\left(\frac{x^3 - 2x^2 + x + y}{\left(x - 1\right)^2}\right) + \left(\frac{1}{1 - x}\right)\frac{\mathrm{d}y}{\mathrm{d}x} = 0$$

The following equations are now set up to solve for the function  $\phi(x, y)$ 

$$\frac{\partial \phi}{\partial x} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \overline{M} dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int \frac{x^3 - 2x^2 + x + y}{(x - 1)^2} dx$$

$$\phi = \frac{x^2}{2} - \frac{y}{x - 1} + f(y)$$
(3)

Where f(y) is used for the constant of integration since  $\phi$  is a function of both x and y. Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = -\frac{1}{x-1} + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{1}{1-x}$ . Therefore equation (4) becomes

$$\frac{1}{1-x} = -\frac{1}{x-1} + f'(y) \tag{5}$$

Solving equation (5) for f'(y) gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for f(y) into equation (3) gives  $\phi$ 

$$\phi = \frac{x^2}{2} - \frac{y}{x-1} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \frac{x^2}{2} - \frac{y}{x - 1}$$

The solution becomes

$$y = -\frac{(-x^2 + 2c_1)(x-1)}{2}$$

Summary

The solution(s) found are the following

$$y = -\frac{(-x^2 + 2c_1)(x-1)}{2} \tag{1}$$

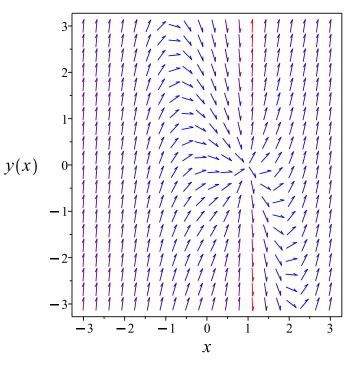


Figure 30: Slope field plot

Verification of solutions

$$y = -\frac{(-x^2 + 2c_1)(x-1)}{2}$$

Verified OK.

#### 2.6.4 Maple step by step solution

Let's solve

 $y' + \tfrac{y}{1-x} = x^2 - x$ 

• Highest derivative means the order of the ODE is 1

y'

• Isolate the derivative

$$y' = \frac{y}{x-1} + x^2 - x$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE  $y' \frac{y}{x-1} = x^2 x$
- The ODE is linear; multiply by an integrating factor  $\mu(x)$  $\mu(x) \left(y' - \frac{y}{x-1}\right) = \mu(x) \left(x^2 - x\right)$
- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x) y)$  $\mu(x) \left(y' - \frac{y}{x-1}\right) = \mu'(x) y + \mu(x) y'$
- Isolate  $\mu'(x)$  $\mu'(x) = -\frac{\mu(x)}{x-1}$
- Solve to find the integrating factor  $\mu(x) = \frac{1}{x-1}$
- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x) y)\right) dx = \int \mu(x) \left(x^2 - x\right) dx + c_1$$

• Evaluate the integral on the lhs

$$\mu(x) y = \int \mu(x) (x^2 - x) dx + c_1$$

- Solve for y $y = \frac{\int \mu(x)(x^2 - x)dx + c_1}{\mu(x)}$
- Substitute  $\mu(x) = \frac{1}{x-1}$

$$y = (x-1)\left(\int \frac{x^2-x}{x-1}dx + c_1\right)$$

• Evaluate the integrals on the rhs

$$y = \left(\frac{x^2}{2} + c_1\right)(x - 1)$$

• Simplify

$$y = \frac{(x-1)(x^2+2c_1)}{2}$$

## Maple trace

`Methods for first order ODEs: --- Trying classification methods --trying a quadrature trying 1st order linear <- 1st order linear successful`</pre>

Solution by Maple Time used: 0.0 (sec). Leaf size: 16

 $dsolve(diff(y(x),x)+y(x)/(1-x)+x-x^2=0,y(x), singsol=all)$ 

$$y(x) = \frac{(x^2 + 2c_1)(x - 1)}{2}$$

Solution by Mathematica

Time used: 0.027 (sec). Leaf size: 20

DSolve[y'[x]+y[x]/(1-x)+x-x^2==0,y[x],x,IncludeSingularSolutions -> True]

$$y(x)
ightarrow rac{1}{2}(x-1)\left(x^2+2c_1
ight)$$

## 2.7 problem 10.3.8

2.7.1	Solving as linear ode	191
2.7.2	Solving as homogeneousTypeD2 ode	193
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2.7.5	Maple step by step solution $\ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots$	204
Internal problem	ID [5061]	
Internal file name	[OUTPUT/4554_Sunday_June_05_2022_03_00_48_PM_6281316/index.	tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.3, ODEs with variable Coefficients.
First order. page 315
Problem number: 10.3.8.
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "homogeneousTypeD2", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_linear]

$$\left(x^2+1\right)y'-xy=1$$

#### 2.7.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{x}{x^2 + 1}$$
$$q(x) = \frac{1}{x^2 + 1}$$

Hence the ode is

$$y' - \frac{xy}{x^2 + 1} = \frac{1}{x^2 + 1}$$

The integrating factor  $\mu$  is

$$\mu = e^{\int -\frac{x}{x^2+1}dx}$$
$$= \frac{1}{\sqrt{x^2+1}}$$

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu y) = (\mu) \left(\frac{1}{x^2 + 1}\right)$$
$$\frac{\mathrm{d}}{\mathrm{d}x}\left(\frac{y}{\sqrt{x^2 + 1}}\right) = \left(\frac{1}{\sqrt{x^2 + 1}}\right) \left(\frac{1}{x^2 + 1}\right)$$
$$\mathrm{d}\left(\frac{y}{\sqrt{x^2 + 1}}\right) = \frac{1}{\left(x^2 + 1\right)^{\frac{3}{2}}} \,\mathrm{d}x$$

Integrating gives

$$\frac{y}{\sqrt{x^2 + 1}} = \int \frac{1}{(x^2 + 1)^{\frac{3}{2}}} dx$$
$$\frac{y}{\sqrt{x^2 + 1}} = \frac{x}{\sqrt{x^2 + 1}} + c_1$$

Dividing both sides by the integrating factor  $\mu = \frac{1}{\sqrt{x^2+1}}$  results in

$$y = c_1 \sqrt{x^2 + 1} + x$$

Summary

 $\overline{\text{The solution}(s)}$  found are the following

$$y = c_1 \sqrt{x^2 + 1} + x \tag{1}$$

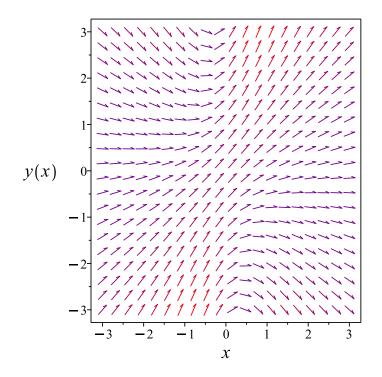


Figure 31: Slope field plot

Verification of solutions

$$y = c_1 \sqrt{x^2 + 1} + x$$

Verified OK.

### 2.7.2 Solving as homogeneousTypeD2 ode

Using the change of variables y = u(x) x on the above ode results in new ode in u(x)

$$(x^{2}+1)(u'(x)x+u(x)) - x^{2}u(x) = 1$$

In canonical form the ODE is

$$u' = F(x, u)$$
$$= f(x)g(u)$$
$$= \frac{-u+1}{x(x^2+1)}$$

Where  $f(x) = \frac{1}{x(x^2+1)}$  and g(u) = -u + 1. Integrating both sides gives

$$\frac{1}{-u+1}\,du=\frac{1}{x\left(x^2+1\right)}\,dx$$

$$\int \frac{1}{-u+1} \, du = \int \frac{1}{x \, (x^2+1)} \, dx$$
$$-\ln\left(u-1\right) = -\frac{\ln\left(x^2+1\right)}{2} + \ln\left(x\right) + c_2$$

Raising both side to exponential gives

$$\frac{1}{u-1} = e^{-\frac{\ln(x^2+1)}{2} + \ln(x) + c_2}$$

Which simplifies to

$$\frac{1}{u-1} = c_3 \mathrm{e}^{-\frac{\ln\left(x^2+1\right)}{2} + \ln(x)}$$

Which simplifies to

$$u(x) = rac{\left(rac{c_3\mathrm{e}^{c_2}x}{\sqrt{x^2+1}}+1
ight)\mathrm{e}^{-c_2}\sqrt{x^2+1}}{c_3x}$$

Therefore the solution y is

$$y = xu$$
  
=  $\frac{\left(\frac{c_3 e^{c_2} x}{\sqrt{x^2 + 1}} + 1\right) e^{-c_2} \sqrt{x^2 + 1}}{c_3}$ 

 $\frac{Summary}{The solution(s) found are the following}$ 

$$y = \frac{\left(\frac{c_3 e^{c_2} x}{\sqrt{x^2 + 1}} + 1\right) e^{-c_2} \sqrt{x^2 + 1}}{c_3} \tag{1}$$

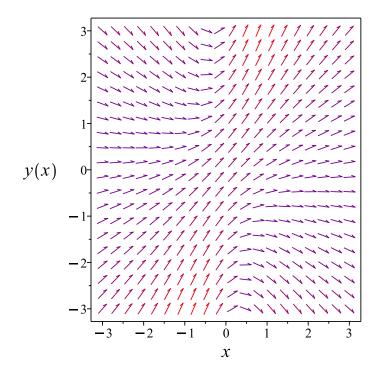


Figure 32: Slope field plot

Verification of solutions

$$y = \frac{\left(\frac{c_{3}e^{c_2}x}{\sqrt{x^2+1}} + 1\right)e^{-c_2}\sqrt{x^2+1}}{c_3}$$

Verified OK.

### 2.7.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{xy+1}{x^2+1}$$
$$y' = \omega(x,y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
(A)

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$ 

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f(rac{y}{x})$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x + b_1y + c_1}{a_2x + b_2y + c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x) y + f_2(x) y^2$	0	$e^{-\int f_1 dx}$

Table 36: Lie symmetry infinitesimal lookup table for known first order ODE's

The above table shows that

$$\xi(x,y) = 0$$
  

$$\eta(x,y) = \sqrt{x^2 + 1}$$
(A1)

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where S(R). Since  $\xi = 0$  then in this special case

R = x

 ${\cal S}$  is found from

$$S = \int \frac{1}{\eta} dy$$
$$= \int \frac{1}{\sqrt{x^2 + 1}} dy$$

Which results in

$$S = \frac{y}{\sqrt{x^2 + 1}}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}$$
(2)

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = \frac{xy+1}{x^2+1}$$

Evaluating all the partial derivatives gives

$$R_{x} = 1$$

$$R_{y} = 0$$

$$S_{x} = -\frac{yx}{(x^{2}+1)^{\frac{3}{2}}}$$

$$S_{y} = \frac{1}{\sqrt{x^{2}+1}}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{(x^2 + 1)^{\frac{3}{2}}}$$
(2A)

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{\left(R^2 + 1\right)^{\frac{3}{2}}}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = \frac{R}{\sqrt{R^2 + 1}} + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{y}{\sqrt{x^2+1}} = \frac{x}{\sqrt{x^2+1}} + c_1$$

Which simplifies to

$$\frac{y}{\sqrt{x^2+1}} = \frac{x}{\sqrt{x^2+1}} + c_1$$

Which gives

$$y = c_1 \sqrt{x^2 + 1} + x$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	$ODE  ext{ in canonical coordinates} \ (R,S)$
$\frac{dy}{dx} = \frac{xy+1}{x^2+1}$	$R = x$ $S = \frac{y}{\sqrt{x^2 + 1}}$	$\frac{dS}{dR} = \frac{1}{(R^2+1)^{\frac{3}{2}}}$

# $\frac{\text{Summary}}{\text{The solution(s) found are the following}}$

$$y = c_1 \sqrt{x^2 + 1} + x \tag{1}$$

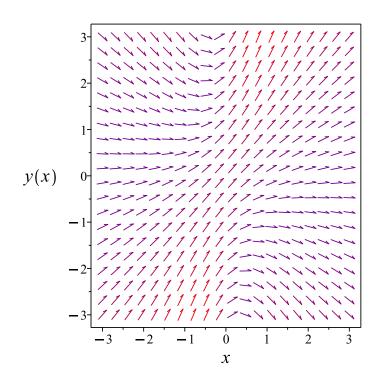


Figure 33: Slope field plot

Verification of solutions

$$y = c_1 \sqrt{x^2 + 1} + x$$

Verified OK.

#### 2.7.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x,y) + N(x,y)\frac{dy}{dx} = 0$$
 (A)

We assume there exists a function  $\phi(x, y) = c$  where c is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t. x gives

$$\frac{d}{dx}\phi(x,y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y}\frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$
$$\frac{\partial \phi}{\partial y} = N$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x,y) dx + N(x,y) dy = 0$$
(1A)

Therefore

$$(x^{2} + 1) dy = (xy + 1) dx$$
$$(-xy - 1) dx + (x^{2} + 1) dy = 0$$
(2A)

Comparing (1A) and (2A) shows that

$$M(x, y) = -xy - 1$$
$$N(x, y) = x^{2} + 1$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y}(-xy-1)$$
$$= -x$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x} (x^2 + 1)$$
$$= 2x$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is <u>not exact</u>. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$A = \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right)$$
$$= \frac{1}{x^2 + 1} ((-x) - (2x))$$
$$= -\frac{3x}{x^2 + 1}$$

Since A does not depend on y, then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\mu = e^{\int A \, \mathrm{d}x}$$
$$= e^{\int -\frac{3x}{x^2 + 1} \, \mathrm{d}x}$$

The result of integrating gives

$$\mu = e^{-\frac{3\ln(x^2+1)}{2}}$$
$$= \frac{1}{(x^2+1)^{\frac{3}{2}}}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called  $\overline{M}$  and  $\overline{N}$  for now so not to confuse them with the original M and N.

$$\overline{M} = \mu M$$
  
=  $\frac{1}{(x^2 + 1)^{\frac{3}{2}}}(-xy - 1)$   
=  $-\frac{xy + 1}{(x^2 + 1)^{\frac{3}{2}}}$ 

And

$$\begin{split} \overline{N} &= \mu N \\ &= \frac{1}{(x^2 + 1)^{\frac{3}{2}}} \big( x^2 + 1 \big) \\ &= \frac{1}{\sqrt{x^2 + 1}} \end{split}$$

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$\overline{M} + \overline{N} \frac{\mathrm{d}y}{\mathrm{d}x} = 0$$
$$\left(-\frac{xy+1}{\left(x^2+1\right)^{\frac{3}{2}}}\right) + \left(\frac{1}{\sqrt{x^2+1}}\right) \frac{\mathrm{d}y}{\mathrm{d}x} = 0$$

The following equations are now set up to solve for the function  $\phi(x, y)$ 

$$\frac{\partial \phi}{\partial x} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \overline{M} dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -\frac{xy+1}{(x^2+1)^{\frac{3}{2}}} dx$$

$$\phi = \frac{-x+y}{\sqrt{x^2+1}} + f(y)$$
(3)

Where f(y) is used for the constant of integration since  $\phi$  is a function of both x and y. Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = \frac{1}{\sqrt{x^2 + 1}} + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{1}{\sqrt{x^2+1}}$ . Therefore equation (4) becomes

$$\frac{1}{\sqrt{x^2+1}} = \frac{1}{\sqrt{x^2+1}} + f'(y) \tag{5}$$

Solving equation (5) for f'(y) gives

f'(y) = 0

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for f(y) into equation (3) gives  $\phi$ 

$$\phi = \frac{-x+y}{\sqrt{x^2+1}} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \frac{-x+y}{\sqrt{x^2+1}}$$

The solution becomes

$$y = c_1 \sqrt{x^2 + 1} + x$$

Summary

The solution(s) found are the following

$$y = c_1 \sqrt{x^2 + 1} + x \tag{1}$$

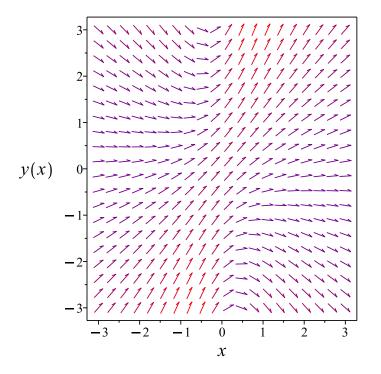


Figure 34: Slope field plot

Verification of solutions

$$y = c_1 \sqrt{x^2 + 1} + x$$

Verified OK.

#### 2.7.5 Maple step by step solution

Let's solve

$$(x^2 + 1) y' - xy = 1$$

- Highest derivative means the order of the ODE is 1 y'
- Isolate the derivative

$$y' = \frac{xy}{x^2+1} + \frac{1}{x^2+1}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE  $y' \frac{xy}{x^2+1} = \frac{1}{x^2+1}$
- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x)\left(y'-rac{xy}{x^2+1}
ight)=rac{\mu(x)}{x^2+1}$$

• Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$ 

$$\mu(x)\left(y' - \frac{xy}{x^2+1}\right) = \mu'(x)\,y + \mu(x)\,y$$

• Isolate  $\mu'(x)$ 

$$\mu'(x) = -\frac{\mu(x)x}{x^2+1}$$

• Solve to find the integrating factor

$$\mu(x) = \frac{1}{\sqrt{x^2 + 1}}$$

• Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x) y)\right) dx = \int \frac{\mu(x)}{x^2+1} dx + c_1$$

• Evaluate the integral on the lhs

$$\mu(x) y = \int rac{\mu(x)}{x^2+1} dx + c_1$$

• Solve for y

$$y = rac{\int rac{\mu(x)}{x^2+1} dx + c_1}{\mu(x)}$$

• Substitute  $\mu(x) = \frac{1}{\sqrt{x^2+1}}$ 

$$y = \sqrt{x^2 + 1} \left( \int \frac{1}{(x^2 + 1)^{\frac{3}{2}}} dx + c_1 \right)$$

- Evaluate the integrals on the rhs  $y = \sqrt{x^2 + 1} \left( \frac{x}{\sqrt{x^2 + 1}} + c_1 \right)$
- Simplify  $y = c_1 \sqrt{x^2 + 1} + x$

#### Maple trace

`Methods for first order ODEs: --- Trying classification methods --trying a quadrature trying 1st order linear <- 1st order linear successful`</pre>

Solution by Maple Time used: 0.0 (sec). Leaf size: 15

 $dsolve((1+x^2)*diff(y(x),x)=1+x*y(x),y(x), singsol=all)$ 

$$y(x) = \sqrt{x^2 + 1} c_1 + x$$

✓ Solution by Mathematica

Time used: 0.031 (sec). Leaf size: 19

DSolve[(1+x^2)\*y'[x]==1+x\*y[x],y[x],x,IncludeSingularSolutions -> True]

$$y(x) \to x + c_1 \sqrt{x^2 + 1}$$

# 2.8 problem 10.3.9 (a)

Solving as separable ode
Solving as first order ode lie symmetry lookup ode 208
Solving as bernoulli ode 212
Solving as exact ode
Solving as riccati ode
Maple step by step solution
n ID [5062]
e [OUTPUT/4555_Sunday_June_05_2022_03_00_49_PM_941248/index.tex]
aining in Mathematics. By R. Shankar. Plenum Press. NY. 1995
ter 10, Differential equations. Section 10.3, ODEs with variable Coefficients.
e 315
<b>iber</b> : 10.3.9 (a).
1.

The type(s) of ODE detected by this program : "exact", "riccati", "bernoulli", "separable", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_separable]

$$y' + xy - xy^2 = 0$$

#### 2.8.1 Solving as separable ode

In canonical form the ODE is

$$y' = F(x, y)$$
$$= f(x)g(y)$$
$$= xy(y-1)$$

Where f(x) = x and g(y) = y(y - 1). Integrating both sides gives

$$\frac{1}{y\left(y-1\right)}\,dy = x\,dx$$

$$\int \frac{1}{y(y-1)} \, dy = \int x \, dx$$
$$\ln(y-1) - \ln(y) = \frac{x^2}{2} + c_1$$

Raising both side to exponential gives

$$e^{\ln(y-1)-\ln(y)} = e^{\frac{x^2}{2}+c_1}$$

Which simplifies to

$$\frac{y-1}{y} = c_2 \mathrm{e}^{\frac{x^2}{2}}$$

Summary

The solution(s) found are the following

$$y = -\frac{1}{-1 + c_2 e^{\frac{x^2}{2}}} \tag{1}$$

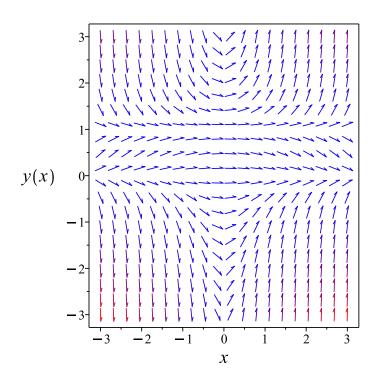


Figure 35: Slope field plot

Verification of solutions

$$y = -\frac{1}{-1 + c_2 \mathrm{e}^{\frac{x^2}{2}}}$$

Verified OK.

#### 2.8.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = y^2 x - xy$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
 (A)

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$ 

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f(rac{y}{x})$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x + b_1y + c_1}{a_2x + b_2y + c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x) y + f_2(x) y^2$	0	$e^{-\int f_1 dx}$

Table 39: Lie symmetry infinitesimal lookup table for known first order ODE's

The above table shows that

$$\xi(x,y) = \frac{1}{x}$$
  

$$\eta(x,y) = 0$$
(A1)

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where S(R). Since  $\eta = 0$  then in this special case

$$R = y$$

 ${\cal S}$  is found from

$$S = \int \frac{1}{\xi} dx$$
$$= \int \frac{1}{\frac{1}{x}} dx$$

Which results in

$$S=\frac{x^2}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}$$
(2)

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = y^2 x - xy$$

Evaluating all the partial derivatives gives

$$R_x = 0$$
$$R_y = 1$$
$$S_x = x$$
$$S_y = 0$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{y\left(y-1\right)} \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R\left(R-1\right)}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = \ln (R - 1) - \ln (R) + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{x^2}{2} = \ln{(y-1)} - \ln{(y)} + c_1$$

Which simplifies to

$$\frac{x^2}{2} = \ln{(y-1)} - \ln{(y)} + c_1$$

Which gives

$$y = \frac{\mathrm{e}^{-\frac{x^2}{2} + c_1}}{-1 + \mathrm{e}^{-\frac{x^2}{2} + c_1}}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	$ODE  ext{ in canonical coordinates} (R, S)$
$\frac{dy}{dx} = y^2 x - xy$	$R = y$ $S = \frac{x^2}{2}$	$\frac{dS}{dR} = \frac{1}{R(R-1)}$

# $\frac{Summary}{The solution(s) found are the following}$

$$y = \frac{1}{-1 + e^{-\frac{x^2}{2} + c_1}}$$

 $e^{-\frac{x^2}{2}+c_1}$ 

(1)

Figure 36: Slope field plot

Verification of solutions

$$y = \frac{\mathrm{e}^{-\frac{x^2}{2} + c_1}}{-1 + \mathrm{e}^{-\frac{x^2}{2} + c_1}}$$

Verified OK.

## 2.8.3 Solving as bernoulli ode

In canonical form, the ODE is

$$y' = F(x, y)$$
$$= y^2 x - xy$$

This is a Bernoulli ODE.

$$y' = -xy + xy^2 \tag{1}$$

The standard Bernoulli ODE has the form

$$y' = f_0(x)y + f_1(x)y^n$$
(2)

The first step is to divide the above equation by  $y^n$  which gives

$$\frac{y'}{y^n} = f_0(x)y^{1-n} + f_1(x) \tag{3}$$

The next step is use the substitution  $w = y^{1-n}$  in equation (3) which generates a new ODE in w(x) which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution y(x) which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$f_0(x) = -x$$
$$f_1(x) = x$$
$$n = 2$$

Dividing both sides of ODE (1) by  $y^n = y^2$  gives

 $y'\frac{1}{y^2} = -\frac{x}{y} + x$  (4)

Let

$$w = y^{1-n}$$
$$= \frac{1}{y}$$
(5)

Taking derivative of equation (5) w.r.t x gives

$$w' = -\frac{1}{y^2}y' \tag{6}$$

Substituting equations (5) and (6) into equation (4) gives

$$-w'(x) = -w(x) x + x$$
  

$$w' = xw - x$$
(7)

The above now is a linear ODE in w(x) which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(x) + p(x)w(x) = q(x)$$

Where here

$$p(x) = -x$$
$$q(x) = -x$$

Hence the ode is

$$w'(x) - w(x) x = -x$$

The integrating factor  $\mu$  is

$$\mu = e^{\int -xdx}$$
$$= e^{-\frac{x^2}{2}}$$

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu w) = (\mu) (-x)$$
$$\frac{\mathrm{d}}{\mathrm{d}x} \left( \mathrm{e}^{-\frac{x^2}{2}} w \right) = \left( \mathrm{e}^{-\frac{x^2}{2}} \right) (-x)$$
$$\mathrm{d} \left( \mathrm{e}^{-\frac{x^2}{2}} w \right) = \left( -x \, \mathrm{e}^{-\frac{x^2}{2}} \right) \, \mathrm{d}x$$

Integrating gives

$$e^{-\frac{x^2}{2}}w = \int -x e^{-\frac{x^2}{2}} dx$$
$$e^{-\frac{x^2}{2}}w = e^{-\frac{x^2}{2}} + c_1$$

Dividing both sides by the integrating factor  $\mu = e^{-\frac{x^2}{2}}$  results in

$$w(x) = e^{\frac{x^2}{2}} e^{-\frac{x^2}{2}} + c_1 e^{\frac{x^2}{2}}$$

which simplifies to

$$w(x) = 1 + c_1 \mathrm{e}^{\frac{x^2}{2}}$$

Replacing w in the above by  $\frac{1}{y}$  using equation (5) gives the final solution.

$$\frac{1}{y} = 1 + c_1 \mathrm{e}^{\frac{x^2}{2}}$$

Or

$$y = \frac{1}{1 + c_1 \mathrm{e}^{\frac{x^2}{2}}}$$

# $\frac{Summary}{The solution(s) found are the following}$

$$y = \frac{1}{1 + c_1 e^{\frac{x^2}{2}}}$$

(1)

Figure 37: Slope field plot

Verification of solutions

$$y = \frac{1}{1 + c_1 \mathrm{e}^{\frac{x^2}{2}}}$$

Verified OK.

### 2.8.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x,y) + N(x,y)\frac{dy}{dx} = 0$$
 (A)

We assume there exists a function  $\phi(x,y) = c$  where c is constant, that satisfies the

ode. Taking derivative of  $\phi$  w.r.t. x gives

$$\frac{d}{dx}\phi(x,y) = 0$$

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y}\frac{dy}{dx} = 0$$
(B)

Hence

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$
$$\frac{\partial \phi}{\partial y} = N$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x,y) dx + N(x,y) dy = 0$$
(1A)

Therefore

$$\left(\frac{1}{y(y-1)}\right) dy = (x) dx$$
$$(-x) dx + \left(\frac{1}{y(y-1)}\right) dy = 0$$
(2A)

Comparing (1A) and (2A) shows that

$$\begin{split} M(x,y) &= -x \\ N(x,y) &= \frac{1}{y \left(y-1\right)} \end{split}$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y}(-x)$$
$$= 0$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x} \left( \frac{1}{y(y-1)} \right)$$
$$= 0$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is <u>exact</u> The following equations are now set up to solve for the function  $\phi(x, y)$ 

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -x dx$$

$$\phi = -\frac{x^2}{2} + f(y)$$
(3)

Where f(y) is used for the constant of integration since  $\phi$  is a function of both x and y. Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{1}{y(y-1)}$ . Therefore equation (4) becomes

$$\frac{1}{y(y-1)} = 0 + f'(y) \tag{5}$$

Solving equation (5) for f'(y) gives

$$f'(y) = \frac{1}{y(y-1)}$$

Integrating the above w.r.t y gives

$$\int f'(y) \, \mathrm{d}y = \int \left(\frac{1}{y(y-1)}\right) \mathrm{d}y$$
$$f(y) = \ln(y-1) - \ln(y) + c_1$$

Where  $c_1$  is constant of integration. Substituting result found above for f(y) into equation (3) gives  $\phi$ 

$$\phi = -\frac{x^2}{2} + \ln{(y-1)} - \ln{(y)} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1=-rac{x^2}{2}+\ln{\left(y-1
ight)}-\ln{\left(y
ight)}$$

The solution becomes

$$y = -\frac{1}{\mathrm{e}^{\frac{x^2}{2} + c_1} - 1}$$

 $\frac{Summary}{The solution(s) found are the following}$ 

$$y = -\frac{1}{\mathrm{e}^{\frac{x^2}{2} + c_1} - 1} \tag{1}$$

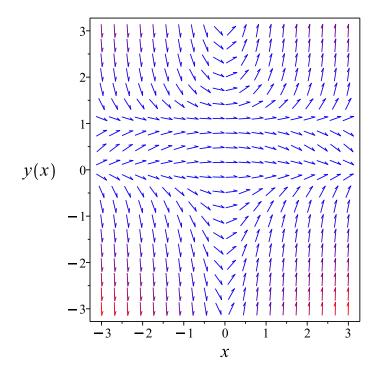


Figure 38: Slope field plot

Verification of solutions

$$y = -\frac{1}{\mathrm{e}^{\frac{x^2}{2} + c_1} - 1}$$

Verified OK.

#### 2.8.5 Solving as riccati ode

In canonical form the ODE is

$$y' = F(x, y)$$
$$= y^2 x - xy$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = y^2 x - xy$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that  $f_0(x) = 0$ ,  $f_1(x) = -x$  and  $f_2(x) = x$ . Let

$$y = \frac{-u'}{f_2 u}$$
$$= \frac{-u'}{x u}$$
(1)

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for u(x) which is

$$f_2 u''(x) - (f'_2 + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0$$
<sup>(2)</sup>

But

$$f_2' = 1 \ f_1 f_2 = -x^2 \ f_2^2 f_0 = 0$$

Substituting the above terms back in equation (2) gives

$$xu''(x) - (-x^2 + 1) u'(x) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 + e^{-\frac{x^2}{2}}c_2$$

The above shows that

$$u'(x) = -x e^{-\frac{x^2}{2}} c_2$$

Using the above in (1) gives the solution

$$y = \frac{\mathrm{e}^{-\frac{x^2}{2}}c_2}{c_1 + \mathrm{e}^{-\frac{x^2}{2}}c_2}$$

Dividing both numerator and denominator by  $c_1$  gives, after renaming the constant  $\frac{c_2}{c_1} = c_3$  the following solution

$$y = \frac{e^{-\frac{x^2}{2}}}{c_3 + e^{-\frac{x^2}{2}}}$$

# $\frac{Summary}{The solution(s) found are the following}$

$$y = \frac{e^{-\frac{x^2}{2}}}{c_3 + e^{-\frac{x^2}{2}}} \tag{1}$$

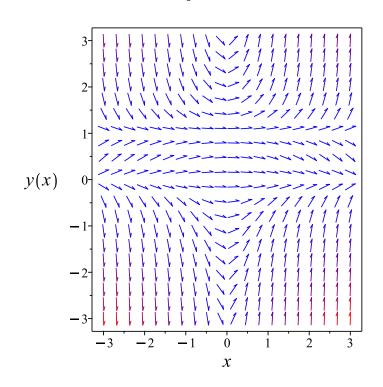


Figure 39: Slope field plot

Verification of solutions

$$y = \frac{e^{-\frac{x^2}{2}}}{c_3 + e^{-\frac{x^2}{2}}}$$

Verified OK.

### 2.8.6 Maple step by step solution

Let's solve  $y' + xy - xy^2 = 0$ Highest derivative means the order of the ODE is 1 y' • Separate variables

$$\frac{y'}{y(y-1)} = x$$

• Integrate both sides with respect to x

$$\int rac{y'}{y(y-1)} dx = \int x dx + c_1$$

• Evaluate integral

$$\ln(y-1) - \ln(y) = \frac{x^2}{2} + c_1$$

• Solve for y $y = -\frac{1}{1}$ 

$$y = -\frac{1}{e^{rac{x^2}{2} + c_1} - 1}$$

Maple trace

`Methods for first order ODEs: --- Trying classification methods --trying a quadrature trying 1st order linear trying Bernoulli <- Bernoulli successful`</pre>

Solution by Maple Time used: 0.0 (sec). Leaf size: 16

 $dsolve(diff(y(x),x)+x*y(x)=x*y(x)^2,y(x), singsol=all)$ 

$$y(x) = rac{1}{1 + \mathrm{e}^{rac{x^2}{2}}c_1}$$

✓ Solution by Mathematica

Time used: 0.25 (sec). Leaf size: 31

DSolve[y'[x]+x\*y[x]==x\*y[x]^2,y[x],x,IncludeSingularSolutions -> True]

$$y(x) \rightarrow \frac{1}{1 + e^{\frac{x^2}{2} + c_1}}$$
$$y(x) \rightarrow 0$$
$$y(x) \rightarrow 1$$

#### 2.9 problem 10.3.9 (b)

2.9.1	Solving as first order ode lie symmetry lookup ode 223							
2.9.2	Solving as bernoulli ode							
Internal problem	n ID [5063]							
Internal file name [OUTPUT/4556_Sunday_June_05_2022_03_00_50_PM_48381868/index.tex]								
Book: Basic Tr	Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995							
Section: Chapter 10, Differential equations. Section 10.3, ODEs with variable Coefficients.								
First order. page 315								
Problem num	<b>aber</b> : 10.3.9 (b).							
<b>ODE order</b> : 1								

**ODE degree**: 1.

The type(s) of ODE detected by this program : "bernoulli", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

```
[[_homogeneous, `class G`], _rational, _Bernoulli]
```

$$\boxed{3xy' + y + y^4x^2 = 0}$$

#### 2.9.1 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{y(y^3x^2 + 1)}{3x}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
 (A)

The type of this ode is known. It is of type Bernoulli. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$ 

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f(rac{y}{x})$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x + b_1y + c_1}{a_2x + b_2y + c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x) y + f_2(x) y^2$	0	$e^{-\int f_1 dx}$

Table 42: Lie symmetry infinitesimal lookup table for known first order ODE's

The above table shows that

$$\begin{aligned} \xi(x,y) &= 0\\ \eta(x,y) &= y^4 x \end{aligned} \tag{A1}$$

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where S(R). Since  $\xi = 0$  then in this special case

$$R = x$$

 ${\cal S}$  is found from

$$S = \int \frac{1}{\eta} dy$$
$$= \int \frac{1}{y^4 x} dy$$

Which results in

$$S = -\frac{1}{3x \, y^3}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}$$
(2)

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = -\frac{y(y^3x^2+1)}{3x}$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 0$$

$$S_x = \frac{1}{3x^2y^3}$$

$$S_y = \frac{1}{y^4x}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{3} \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{3}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = -\frac{R}{3} + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-\frac{1}{3y^3x} = -\frac{x}{3} + c_1$$

Which simplifies to

$$-\frac{1}{3y^3x} = -\frac{x}{3} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	$ODE  ext{ in canonical coordinates} (R, S)$
$\frac{dy}{dx} = -\frac{y(y^3x^2+1)}{3x}$	$R = x$ $S = -\frac{1}{3x y^3}$	$\frac{dS}{dR} = -\frac{1}{3}$

#### Summary

The solution(s) found are the following

$$-\frac{1}{3y^3x} = -\frac{x}{3} + c_1 \tag{1}$$

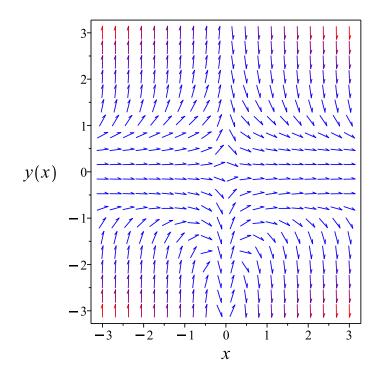


Figure 40: Slope field plot

Verification of solutions

$$-\frac{1}{3y^3x} = -\frac{x}{3} + c_1$$

Verified OK.

#### 2.9.2 Solving as bernoulli ode

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x,y) \\ &= -\frac{y(y^3x^2+1)}{3x} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = -\frac{1}{3x}y - \frac{x}{3}y^4$$
 (1)

The standard Bernoulli ODE has the form

$$y' = f_0(x)y + f_1(x)y^n$$
(2)

The first step is to divide the above equation by  $y^n$  which gives

$$\frac{y'}{y^n} = f_0(x)y^{1-n} + f_1(x) \tag{3}$$

The next step is use the substitution  $w = y^{1-n}$  in equation (3) which generates a new ODE in w(x) which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution y(x) which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$f_0(x) = -\frac{1}{3x}$$
$$f_1(x) = -\frac{x}{3}$$
$$n = 4$$

Dividing both sides of ODE (1) by  $y^n = y^4$  gives

$$y'\frac{1}{y^4} = -\frac{1}{3x\,y^3} - \frac{x}{3} \tag{4}$$

Let

$$w = y^{1-n}$$
$$= \frac{1}{y^3}$$
(5)

Taking derivative of equation (5) w.r.t x gives

$$w' = -\frac{3}{y^4}y' \tag{6}$$

Substituting equations (5) and (6) into equation (4) gives

$$-\frac{w'(x)}{3} = -\frac{w(x)}{3x} - \frac{x}{3}$$
$$w' = \frac{w}{x} + x \tag{7}$$

The above now is a linear ODE in w(x) which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(x) + p(x)w(x) = q(x)$$

Where here

$$p(x) = -\frac{1}{x}$$
$$q(x) = x$$

Hence the ode is

$$w'(x) - \frac{w(x)}{x} = x$$

The integrating factor  $\mu$  is

$$\mu = e^{\int -\frac{1}{x}dx}$$
$$= \frac{1}{x}$$

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu w) = (\mu) (x)$$
$$\frac{\mathrm{d}}{\mathrm{d}x}\left(\frac{w}{x}\right) = \left(\frac{1}{x}\right) (x)$$
$$\mathrm{d}\left(\frac{w}{x}\right) = \mathrm{d}x$$

Integrating gives

$$\frac{w}{x} = \int \mathrm{d}x$$
$$\frac{w}{x} = x + c_1$$

Dividing both sides by the integrating factor  $\mu = \frac{1}{x}$  results in

$$w(x) = c_1 x + x^2$$

which simplifies to

$$w(x) = x(x + c_1)$$

Replacing w in the above by  $\frac{1}{y^3}$  using equation (5) gives the final solution.

$$\frac{1}{y^3} = x(x+c_1)$$

Solving for y gives

$$y(x) = \frac{\left(x^2(x+c_1)^2\right)^{\frac{1}{3}}}{x(x+c_1)}$$
$$y(x) = \frac{\left(x^2(x+c_1)^2\right)^{\frac{1}{3}}\left(i\sqrt{3}-1\right)}{2x(x+c_1)}$$
$$y(x) = -\frac{\left(x^2(x+c_1)^2\right)^{\frac{1}{3}}\left(1+i\sqrt{3}\right)}{2x(x+c_1)}$$

### Summary

 $\overline{\text{The solution}(s)}$  found are the following

$$y = \frac{\left(x^2(x+c_1)^2\right)^{\frac{1}{3}}}{x\left(x+c_1\right)} \tag{1}$$

$$y = \frac{\left(x^2(x+c_1)^2\right)^{\frac{1}{3}}\left(i\sqrt{3}-1\right)}{2x\left(x+c_1\right)} \tag{2}$$

$$y = -\frac{\left(x^2(x+c_1)^2\right)^{\frac{1}{3}}\left(1+i\sqrt{3}\right)}{2x\left(x+c_1\right)}$$
(3)

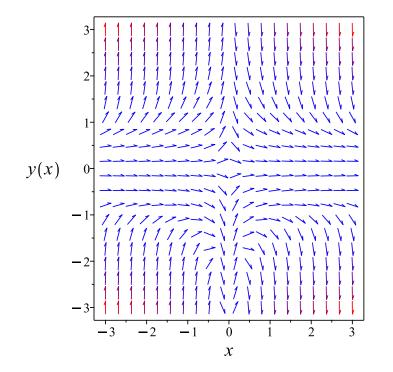


Figure 41: Slope field plot

Verification of solutions

$$y = \frac{\left(x^{2}(x+c_{1})^{2}\right)^{\frac{1}{3}}}{x(x+c_{1})}$$

Verified OK.

$$y = rac{\left(x^2(x+c_1)^2
ight)^{rac{1}{3}}\left(i\sqrt{3}-1
ight)}{2x\left(x+c_1
ight)}$$

Verified OK.

$$y = -\frac{\left(x^2(x+c_1)^2\right)^{\frac{1}{3}}\left(1+i\sqrt{3}\right)}{2x\left(x+c_1\right)}$$

Verified OK.

Maple trace

`Methods for first order ODEs: --- Trying classification methods --trying a quadrature trying 1st order linear trying Bernoulli <- Bernoulli successful`</pre>

Solution by Maple Time used: 0.0 (sec). Leaf size: 88

 $dsolve(3*x*diff(y(x),x)+y(x)+x^2*y(x)^4=0,y(x), singsol=all)$ 

$$\begin{split} y(x) &= \frac{\left((x+c_1)^2 \, x^2\right)^{\frac{1}{3}}}{(x+c_1) \, x} \\ y(x) &= -\frac{\left((x+c_1)^2 \, x^2\right)^{\frac{1}{3}} \left(1+i\sqrt{3}\right)}{2 \, (x+c_1) \, x} \\ y(x) &= \frac{\left((x+c_1)^2 \, x^2\right)^{\frac{1}{3}} \left(i\sqrt{3}-1\right)}{2 \, (x+c_1) \, x} \end{split}$$

## Solution by Mathematica

Time used: 0.3 (sec). Leaf size: 61

DSolve[3\*x\*y'[x]+y[x]+x^2\*y[x]^4==0,y[x],x,IncludeSingularSolutions -> True]

$$\begin{split} y(x) &\to \frac{1}{\sqrt[3]{x(x+c_1)}} \\ y(x) &\to -\frac{\sqrt[3]{-1}}{\sqrt[3]{x(x+c_1)}} \\ y(x) &\to \frac{(-1)^{2/3}}{\sqrt[3]{x(x+c_1)}} \\ y(x) &\to 0 \end{split}$$

## 3 Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients. Second order and Homogeneous. page 318

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<b>3.1</b> problem 10.4.8 (a)         3.1.1       Solving using Kovacic algorithm         3.1.2       Maple step by step solution						
Internal file name [OUTPUT/4557_Sunday_June_05_2022_03_00_51_PM_97582610/index.tex]						
<ul> <li>Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995</li> <li>Section: Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients.</li> <li>Second order and Homogeneous. page 318</li> <li>Problem number: 10.4.8 (a).</li> <li>ODE order: 2.</li> <li>ODE degree: 1.</li> </ul>						

The type(s) of ODE detected by this program : "kovacic"

Maple gives the following as the ode type

[[\_2nd\_order, \_with\_linear\_symmetries]]

 $x(x+1)^{2} y'' + (-x^{2}+1) y' + (x-1) y = 0$ 

#### 3.1.1 Solving using Kovacic algorithm

Writing the ode as

$$x(x+1)^{2} y'' + (-x^{2}+1) y' + (x-1) y = 0$$
<sup>(1)</sup>

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = x(x+1)^{2}$$
  

$$B = -x^{2} + 1$$
  

$$C = x - 1$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4x^2} \tag{6}$$

Comparing the above to (5) shows that

$$s = -1$$
$$t = 4x^2$$

Therefore eq. (4) becomes

$$z''(x) = \left(-\frac{1}{4x^2}\right)z(x) \tag{7}$$

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 44: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 2 - 0$$
$$= 2$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of  $t = 4x^2$ . There is a pole at x = 0 of order 2. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case n = 1.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{4x^2}$$

For the pole at x = 0 let b be the coefficient of  $\frac{1}{x^2}$  in the partial fractions decomposition of r given above. Therefore  $b = -\frac{1}{4}$ . Hence

$$\begin{split} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1+4b} = \frac{1}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1+4b} = \frac{1}{2} \end{split}$$

Since the order of r at  $\infty$  is 2 then  $[\sqrt{r}]_{\infty} = 0$ . Let b be the coefficient of  $\frac{1}{x^2}$  in the Laurent series expansion of r at  $\infty$ , which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{1}{4x^2}$$

Since the gcd(s,t) = 1. This gives  $b = -\frac{1}{4}$ . Hence

$$\begin{split} [\sqrt{r}]_{\infty} &= 0\\ \alpha_{\infty}^{+} &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2}\\ \alpha_{\infty}^{-} &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{split}$$

The following table summarizes the findings so far for poles and for the order of r at  $\infty$  where r is 1

		r =	$\overline{4x^2}$				
pole	e $c$ location	pole	order	[	$\sqrt{r}]_c$	$\alpha_c^+$	$lpha_c^-$
	0	4	2		0	$\frac{1}{2}$	$\frac{1}{2}$
Order of $r$ at $\propto$			$[\sqrt{r}]_{c}$	Ø	$lpha_\infty^+$	$lpha_{\infty}^{-}$	
	0		$\frac{1}{2}$	$\frac{1}{2}$			

Now that the all  $[\sqrt{r}]_c$  and its associated  $\alpha_c^{\pm}$  have been determined for all the poles in the set  $\Gamma$  and  $[\sqrt{r}]_{\infty}$  and its associated  $\alpha_{\infty}^{\pm}$  have also been found, the next step is to determine possible non negative integer d from these using

$$d=\alpha_\infty^{s(\infty)}-\sum_{c\in\Gamma}\alpha_c^{s(c)}$$

Where s(c) is either + or - and  $s(\infty)$  is the sign of  $\alpha_{\infty}^{\pm}$ . This is done by trial over all set of families  $s = (s(c))_{c \in \Gamma \cup \infty}$  until such d is found to work in finding candidate  $\omega$ . Trying  $\alpha_{\infty}^{-} = \frac{1}{2}$  then

$$d = \alpha_{\infty}^{-} - (\alpha_{c_1}^{-})$$
$$= \frac{1}{2} - \left(\frac{1}{2}\right)$$
$$= 0$$

Since d an integer and  $d \geq 0$  then it can be used to find  $\omega$  using

$$\omega = \sum_{c \in \Gamma} \left( s(c) [\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{x - c} \right) + s(\infty) [\sqrt{r}]_{\infty}$$

The above gives

$$\begin{split} \omega &= \left( (-) [\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}}{x - c_1} \right) + (-) [\sqrt{r}]_{\infty} \\ &= \frac{1}{2x} + (-) (0) \\ &= \frac{1}{2x} \\ &= \frac{1}{2x} \end{split}$$

Now that  $\omega$  is determined, the next step is find a corresponding minimal polynomial p(x) of degree d = 0 to solve the ode. The polynomial p(x) needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0$$
(1A)

Let

$$p(x) = 1 \tag{2A}$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(\frac{1}{2x}\right)(0) + \left(\left(-\frac{1}{2x^2}\right) + \left(\frac{1}{2x}\right)^2 - \left(-\frac{1}{4x^2}\right)\right) = 0$$
$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode z'' = rz is

$$z_1(x) = p e^{\int \omega \, dx}$$
$$= e^{\int \frac{1}{2x} \, dx}$$
$$= \sqrt{x}$$

The first solution to the original ode in y is found from

$$y_{1} = z_{1}e^{\int -\frac{1}{2}\frac{B}{A}dx}$$
  
=  $z_{1}e^{-\int \frac{1}{2}\frac{-x^{2}+1}{x(x+1)^{2}}dx}$   
=  $z_{1}e^{\ln(x+1)-\frac{\ln(x)}{2}}$   
=  $z_{1}\left(\frac{x+1}{\sqrt{x}}\right)$ 

Which simplifies to

$$y_1 = x + 1$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} \, dx}}{y_1^2} \, dx$$

Substituting gives

$$y_{2} = y_{1} \int \frac{e^{\int -\frac{-x^{2}+1}{x(x+1)^{2}} dx}}{(y_{1})^{2}} dx$$
$$= y_{1} \int \frac{e^{2\ln(x+1)-\ln(x)}}{(y_{1})^{2}} dx$$
$$= y_{1}(\ln(x))$$

Therefore the solution is

$$egin{aligned} y &= c_1 y_1 + c_2 y_2 \ &= c_1 (x+1) + c_2 (x+1 (\ln{(x)})) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = (x+1)c_1 + c_2(x+1)\ln(x)$$
(1)

Verification of solutions

$$y = (x+1)c_1 + c_2(x+1)\ln(x)$$

Verified OK.

#### 3.1.2 Maple step by step solution

Let's solve

 $x(x+1)^{2} y'' + (-x^{2}+1) y' + (x-1) y = 0$ 

- Highest derivative means the order of the ODE is 2 y''
- Isolate 2nd derivative

$$y'' = -\frac{(x-1)y}{x(x+1)^2} + \frac{(x-1)y'}{x(x+1)}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear  $y'' \frac{(x-1)y'}{x(x+1)} + \frac{(x-1)y}{x(x+1)^2} = 0$
- $\Box$  Check to see if  $x_0$  is a regular singular point
  - Define functions

$$\begin{split} \left| P_2(x) &= -\frac{x-1}{(x+1)x}, P_3(x) = \frac{x-1}{x(x+1)^2} \right] \\ \circ \quad (x+1) \cdot P_2(x) \text{ is analytic at } x = -1 \\ ((x+1) \cdot P_2(x)) \right|_{x=-1} &= -2 \\ \circ \quad (x+1)^2 \cdot P_3(x) \text{ is analytic at } x = -1 \\ ((x+1)^2 \cdot P_3(x)) \left|_{x=-1} = 2 \end{split}$$

 $\circ \quad x = -1 \text{is a regular singular point}$ 

Check to see if  $x_0$  is a regular singular point

$$x_0 = -1$$

 $x(x+1)^{2} y'' - (x-1) (x+1) y' + (x-1) y = 0$ 

• Change variables using x = u - 1 so that the regular singular point is at u = 0

$$(u^{3} - u^{2})\left(\frac{d^{2}}{du^{2}}y(u)\right) + (-u^{2} + 2u)\left(\frac{d}{du}y(u)\right) + (u - 2)y(u) = 0$$

• Assume series solution for y(u)

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

 $\Box$  Rewrite ODE with series expansions

 $\circ \quad \text{Convert} \; u^m \cdot y(u) \text{ to series expansion for } m = 0..1$ 

$$u^m \cdot y(u) = \sum_{k=0}^{\infty} a_k u^{k+r+m}$$

• Shift index using k - > k - m

$$u^m \cdot y(u) = \sum_{k=m}^{\infty} a_{k-m} u^{k+r}$$

- Convert  $u^m \cdot \left(\frac{d}{du}y(u)\right)$  to series expansion for m = 1..2 $u^m \cdot \left(\frac{d}{du}y(u)\right) = \sum_{k=0}^{\infty} a_k(k+r) u^{k+r-1+m}$
- $\circ \quad \text{Shift index using } k->k+1-m$

$$u^{m} \cdot \left(\frac{d}{du}y(u)\right) = \sum_{k=-1+m}^{\infty} a_{k+1-m}(k+1-m+r) u^{k+r}$$

• Convert  $u^m \cdot \left(\frac{d^2}{du^2}y(u)\right)$  to series expansion for m = 2..3

$$u^m \cdot \left(\frac{d^2}{du^2}y(u)\right) = \sum_{k=0}^{\infty} a_k(k+r)\left(k+r-1\right)u^{k+r-2+m}$$

 $\circ \quad \text{Shift index using } k->\!\!k+2-m$ 

$$u^{m} \cdot \left(\frac{d^{2}}{du^{2}}y(u)\right) = \sum_{k=-2+m}^{\infty} a_{k+2-m}(k+2-m+r)\left(k+1-m+r\right)u^{k+r}$$

Rewrite ODE with series expansions

$$-a_0(-1+r)\left(-2+r\right)u^r + \left(\sum_{k=1}^{\infty} \left(-a_k(k+r-1)\left(k+r-2\right) + a_{k-1}(k+r-2)^2\right)u^{k+r}\right) = 0$$

•  $a_0$  cannot be 0 by assumption, giving the indicial equation

$$-(-1+r)(-2+r) = 0$$

- Values of **r** that satisfy the indicial equation  $r \in \{1, 2\}$
- Each term in the series must be 0, giving the recursion relation  $(l + l + l)(l + l + l) = (l + l + l)^2 = 0$

$$-a_k(k+r-1)(k+r-2) + a_{k-1}(k+r-2)^2 = 0$$

• Shift index using k - >k + 1

$$-a_{k+1}(k+r)(k+r-1) + a_k(k+r-1)^2 = 0$$

- Recursion relation that defines series solution to ODE  $a_{k+1} = \frac{a_k(k+r-1)}{k+r}$
- Recursion relation for r = 1

$$a_{k+1} = \frac{a_k k}{k+1}$$

• Solution for r = 1

$$\left[y(u)=\sum\limits_{k=0}^{\infty}a_{k}u^{k+1},a_{k+1}=rac{a_{k}k}{k+1}
ight]$$

• Revert the change of variables u = x + 1

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^{k+1}, a_{k+1} = \frac{a_k k}{k+1}\right]$$

• Recursion relation for r = 2

$$a_{k+1} = \frac{a_k(k+1)}{k+2}$$

• Solution for r = 2

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+2}, a_{k+1} = rac{a_k(k+1)}{k+2}
ight]$$

• Revert the change of variables u = x + 1

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^{k+2}, a_{k+1} = rac{a_k (k+1)}{k+2}
ight]$$

Combine solutions and rename parameters  $\left[y = \left(\sum_{k=0}^{\infty} a_k (x+1)^{k+1}\right) + \left(\sum_{k=0}^{\infty} b_k (x+1)^{k+2}\right), a_{k+1} = \frac{a_k k}{k+1}, b_{k+1} = \frac{b_k (k+1)}{k+2}\right]$ 

#### Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
A Liouvillian solution exists
Reducible group (found an exponential solution)
Group is reducible, not completely reducible
<- Kovacics algorithm successful`</pre>
```

Solution by Maple Time used: 0.016 (sec). Leaf size: 14

 $dsolve(x*(x+1)^2*diff(y(x),x$2)+(1-x^2)*diff(y(x),x)+(x-1)*y(x)=0,y(x), singsol=all)$ 

 $y(x) = (x + 1) (c_2 \ln (x) + c_1)$ 

Solution by Mathematica

Time used: 0.029 (sec). Leaf size: 17

DSolve[x\*(x+1)^2\*y''[x]+(1-x^2)\*y'[x]+(x-1)\*y[x]==0,y[x],x,IncludeSingularSolutions -> True]

 $y(x) \to (x+1)(c_2 \log(x) + c_1)$ 

#### 3.2 problem 10.4.8 (b)

3.	.2.1	Solving as second order change of variable on y method 1 ode $\ .$	243					
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3.	.2.3	Solving as second order integrable as is ode	247					
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3.	.2.7	Maple step by step solution	255					
Internal pro	oblem	ID [5065]						
Internal file	name	[OUTPUT/4558_Sunday_June_05_2022_03_00_52_PM_84210290/index	.tex]					
Book: Bas	sic Tra	ining in Mathematics. By R. Shankar. Plenum Press. NY. 1995						
Section: Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients.								
Second order and Homogeneous. page 318								
<b>Problem number</b> : 10.4.8 (b).								

**ODE order**: 2.

**ODE degree**: 1.

The type(s) of ODE detected by this program : "kovacic", "exact linear second order ode", "second\_order\_integrable\_as\_is", "second\_order\_change\_of\_variable\_on\_y\_method\_1", "second\_order\_change\_of\_variable\_on\_y\_method\_2"

Maple gives the following as the ode type

[[\_2nd\_order, \_exact, \_linear, \_homogeneous]]

$$x(1-x)y'' + 2(1-2x)y' - 2y = 0$$

#### 3.2.1 Solving as second order change of variable on y method 1 ode

In normal form the given ode is written as

$$y'' + p(x) y' + q(x) y = 0$$
(2)

Where

$$p(x) = \frac{-4x+2}{-x^2+x} q(x) = -\frac{2}{-x^2+x}$$

Calculating the Liouville ode invariant Q given by

$$\begin{aligned} Q &= q - \frac{p'}{2} - \frac{p^2}{4} \\ &= -\frac{2}{-x^2 + x} - \frac{\left(\frac{-4x + 2}{-x^2 + x}\right)'}{2} - \frac{\left(\frac{-4x + 2}{-x^2 + x}\right)^2}{4} \\ &= -\frac{2}{-x^2 + x} - \frac{\left(-\frac{4}{-x^2 + x} - \frac{\left(-4x + 2\right)\left(1 - 2x\right)}{\left(-x^2 + x\right)^2}\right)}{2} - \frac{\left(\frac{\left(-4x + 2\right)^2}{\left(-x^2 + x\right)^2}\right)}{4} \\ &= -\frac{2}{-x^2 + x} - \left(-\frac{2}{-x^2 + x} - \frac{\left(-4x + 2\right)\left(1 - 2x\right)}{2\left(-x^2 + x\right)^2}\right) - \frac{\left(-4x + 2\right)^2}{4\left(-x^2 + x\right)^2} \\ &= 0 \end{aligned}$$

Since the Liouville ode invariant does not depend on the independent variable x then the transformation

$$y = v(x) z(x) \tag{3}$$

is used to change the original ode to a constant coefficients ode in v. In (3) the term z(x) is given by

$$z(x) = e^{-\left(\int \frac{p(x)}{2} dx\right)} = e^{-\int \frac{-4x+2}{-x^2+x}} = \frac{1}{x(x-1)}$$
(5)

Hence (3) becomes

$$y = \frac{v(x)}{x(x-1)} \tag{4}$$

Applying this change of variable to the original ode results in

$$-v''(x) = 0$$

Which is now solved for v(x) Integrating twice gives the solution

$$v(x) = c_1 x + c_2$$

Now that v(x) is known, then

$$y = v(x) z(x) = (c_1 x + c_2) (z(x))$$
(7)

But from (5)

$$z(x) = \frac{1}{x(x-1)}$$

Hence (7) becomes

$$y = \frac{c_1 x + c_2}{x \left(x - 1\right)}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 x + c_2}{x \left(x - 1\right)} \tag{1}$$

Verification of solutions

$$y = \frac{c_1 x + c_2}{x \left(x - 1\right)}$$

Verified OK.

#### 3.2.2 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$(-x^{2} + x) y'' + (-4x + 2) y' - 2y = 0$$
(1)

Becomes

$$y'' + p(x) y' + q(x) y = 0$$
(2)

Where

$$p(x) = \frac{4x - 2}{x(x - 1)}$$
$$q(x) = \frac{2}{x(x - 1)}$$

Applying change of variables on the dependent variable  $y = v(x) x^n$  to (2) gives the following ode where the dependent variables is v(x) and not y.

$$v''(x) + \left(\frac{2n}{x} + p\right)v'(x) + \left(\frac{n(n-1)}{x^2} + \frac{np}{x} + q\right)v(x) = 0$$
(3)

Let the coefficient of v(x) above be zero. Hence

$$\frac{n(n-1)}{x^2} + \frac{np}{x} + q = 0 \tag{4}$$

Substituting the earlier values found for p(x) and q(x) into (4) gives

$$\frac{n(n-1)}{x^2} + \frac{n(4x-2)}{x^2(x-1)} + \frac{2}{x(x-1)} = 0$$
(5)

Solving (5) for n gives

$$n = -1 \tag{6}$$

Substituting this value in (3) gives

$$v''(x) + \left(-\frac{2}{x} + \frac{4x - 2}{x(x - 1)}\right)v'(x) = 0$$
$$v''(x) + \frac{2v'(x)}{x - 1} = 0$$
(7)

Using the substitution

$$u(x) = v'(x)$$

Then (7) becomes

$$u'(x) + \frac{2u(x)}{x-1} = 0 \tag{8}$$

The above is now solved for u(x). In canonical form the ODE is

$$u' = F(x, u)$$
$$= f(x)g(u)$$
$$= -\frac{2u}{x-1}$$

Where  $f(x) = -\frac{2}{x-1}$  and g(u) = u. Integrating both sides gives

$$\frac{1}{u} du = -\frac{2}{x-1} dx$$
$$\int \frac{1}{u} du = \int -\frac{2}{x-1} dx$$
$$\ln(u) = -2\ln(x-1) + c_1$$
$$u = e^{-2\ln(x-1)+c_1}$$
$$= \frac{c_1}{(x-1)^2}$$

Now that u(x) is known, then

$$v'(x) = u(x)$$
$$v(x) = \int u(x) dx + c_2$$
$$= -\frac{c_1}{x - 1} + c_2$$

Hence

$$y = v(x) x^{n}$$
  
=  $\frac{-\frac{c_{1}}{x-1} + c_{2}}{x}$   
=  $\frac{-\frac{c_{1}}{x-1} + c_{2}}{x}$ 

Summary

The solution(s) found are the following

$$y = \frac{-\frac{c_1}{x-1} + c_2}{x}$$
(1)

Verification of solutions

$$y = \frac{-\frac{c_1}{x-1} + c_2}{x}$$

Verified OK.

#### 3.2.3 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t x gives

$$\int \left( \left( -x^2 + x \right) y'' + \left( -4x + 2 \right) y' - 2y \right) dx = 0$$
$$-(2x - 1) y - \left( x^2 - x \right) y' = c_1$$

Which is now solved for y.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{1-2x}{x(x-1)}$$
$$q(x) = -\frac{c_1}{x(x-1)}$$

Hence the ode is

$$y' - \frac{(1-2x)y}{x(x-1)} = -\frac{c_1}{x(x-1)}$$

The integrating factor  $\mu$  is

$$\mu = e^{\int -\frac{1-2x}{x(x-1)}dx}$$
$$= x(x-1)$$

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu y) = (\mu) \left(-\frac{c_1}{x(x-1)}\right)$$
$$\frac{\mathrm{d}}{\mathrm{d}x}(yx(x-1)) = (x(x-1)) \left(-\frac{c_1}{x(x-1)}\right)$$
$$\mathrm{d}(yx(x-1)) = (-c_1) \mathrm{d}x$$

Integrating gives

$$yx(x-1) = \int -c_1 \, \mathrm{d}x$$
$$yx(x-1) = -c_1x + c_2$$

Dividing both sides by the integrating factor  $\mu = x(x-1)$  results in

$$y = -\frac{c_1}{x-1} + \frac{c_2}{x(x-1)}$$

which simplifies to

$$y = \frac{-c_1 x + c_2}{x \left(x - 1\right)}$$

Summary

The solution(s) found are the following

$$y = \frac{-c_1 x + c_2}{x \left(x - 1\right)} \tag{1}$$

Verification of solutions

$$y = \frac{-c_1 x + c_2}{x (x - 1)}$$

Verified OK.

# 3.2.4 Solving as type second\_order\_integrable\_as\_is (not using ABC version)

Writing the ode as

$$(-x^{2} + x) y'' + (-4x + 2) y' - 2y = 0$$

Integrating both sides of the ODE w.r.t x gives

$$\int \left( \left( -x^2 + x \right) y'' + \left( -4x + 2 \right) y' - 2y \right) dx = 0$$
$$-(2x - 1) y - \left( x^2 - x \right) y' = c_1$$

Which is now solved for y.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{1-2x}{x(x-1)}$$
$$q(x) = -\frac{c_1}{x(x-1)}$$

Hence the ode is

$$y' - \frac{(1-2x)y}{x(x-1)} = -\frac{c_1}{x(x-1)}$$

The integrating factor  $\mu$  is

$$\mu = e^{\int -\frac{1-2x}{x(x-1)}dx}$$
$$= x(x-1)$$

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu y) = (\mu) \left(-\frac{c_1}{x(x-1)}\right)$$
$$\frac{\mathrm{d}}{\mathrm{d}x}(yx(x-1)) = (x(x-1)) \left(-\frac{c_1}{x(x-1)}\right)$$
$$\mathrm{d}(yx(x-1)) = (-c_1) \mathrm{d}x$$

Integrating gives

$$yx(x-1) = \int -c_1 \, \mathrm{d}x$$
$$yx(x-1) = -c_1x + c_2$$

Dividing both sides by the integrating factor  $\mu = x(x-1)$  results in

$$y = -\frac{c_1}{x - 1} + \frac{c_2}{x (x - 1)}$$

which simplifies to

$$y = \frac{-c_1 x + c_2}{x \left(x - 1\right)}$$

Summary

The solution(s) found are the following

$$y = \frac{-c_1 x + c_2}{x \left(x - 1\right)} \tag{1}$$

Verification of solutions

$$y = \frac{-c_1 x + c_2}{x \left(x - 1\right)}$$

Verified OK.

#### 3.2.5 Solving using Kovacic algorithm

Writing the ode as

$$(-x^{2} + x) y'' + (-4x + 2) y' - 2y = 0$$
(1)

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = -x^{2} + x$$
  

$$B = -4x + 2$$
  

$$C = -2$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 0$$
  
 $t = 1$ 

Therefore eq. (4) becomes

$$z''(x) = 0 \tag{7}$$

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 46: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 0 - -\infty$$
$$= \infty$$

There are no poles in r. Therefore the set of poles  $\Gamma$  is empty. Since there is no odd order pole larger than 2 and the order at  $\infty$  is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since r = 0 is not a function of x, then there is no need run Kovacic algorithm to obtain a solution for transformed ode z'' = rz as one solution is

$$z_1(x) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_{1} = z_{1}e^{\int -\frac{1}{2}\frac{B}{A}dx}$$
  
=  $z_{1}e^{-\int \frac{1}{2}\frac{-4x+2}{-x^{2}+x}dx}$   
=  $z_{1}e^{-\ln(x(x-1))}$   
=  $z_{1}\left(\frac{1}{x(x-1)}\right)$ 

Which simplifies to

$$y_1 = \frac{1}{x\left(x-1\right)}$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} \, dx}}{y_1^2} \, dx$$

Substituting gives

$$y_2 = y_1 \int rac{e^{\int -rac{-4x+2}{-x^2+x} \, dx}}{\left(y_1
ight)^2} \, dx$$
 $= y_1 \int rac{e^{-2\ln(x(x-1))}}{\left(y_1
ight)^2} \, dx$ 
 $= y_1(x)$ 

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
  
=  $c_1 \left( \frac{1}{x (x - 1)} \right) + c_2 \left( \frac{1}{x (x - 1)} (x) \right)$ 

Summary

The solution(s) found are the following

$$y = \frac{c_1}{x(x-1)} + \frac{c_2}{x-1} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{x(x-1)} + \frac{c_2}{x-1}$$

Verified OK.

### 3.2.6 Solving as exact linear second order ode ode

An ode of the form

$$p(x) y'' + q(x) y' + r(x) y = s(x)$$

is exact if

$$p''(x) - q'(x) + r(x) = 0$$
(1)

For the given ode we have

$$p(x) = -x^{2} + x$$
$$q(x) = -4x + 2$$
$$r(x) = -2$$
$$s(x) = 0$$

Hence

$$p''(x) = -2$$
$$q'(x) = -4$$

Therefore (1) becomes

$$-2 - (-4) + (-2) = 0$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$(p(x) y' + (q(x) - p'(x)) y)' = s(x)$$

Integrating gives

$$p(x) y' + (q(x) - p'(x)) y = \int s(x) dx$$

Substituting the above values for p, q, r, s gives

$$(-x^{2} + x) y' + (1 - 2x) y = c_{1}$$

We now have a first order ode to solve which is

$$(-x^{2}+x) y' + (1-2x) y = c_{1}$$

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{1-2x}{x(x-1)}$$
$$q(x) = -\frac{c_1}{x(x-1)}$$

Hence the ode is

$$y' - rac{(1-2x)y}{x(x-1)} = -rac{c_1}{x(x-1)}$$

The integrating factor  $\mu$  is

$$\mu = e^{\int -\frac{1-2x}{x(x-1)}dx}$$
$$= x(x-1)$$

The ode becomes

$$\frac{\mathrm{d}}{\mathrm{d}x}(\mu y) = (\mu) \left(-\frac{c_1}{x(x-1)}\right)$$
$$\frac{\mathrm{d}}{\mathrm{d}x}(yx(x-1)) = (x(x-1)) \left(-\frac{c_1}{x(x-1)}\right)$$
$$\mathrm{d}(yx(x-1)) = (-c_1) \mathrm{d}x$$

Integrating gives

$$yx(x-1) = \int -c_1 \, \mathrm{d}x$$
$$yx(x-1) = -c_1x + c_2$$

Dividing both sides by the integrating factor  $\mu = x(x-1)$  results in

$$y = -\frac{c_1}{x - 1} + \frac{c_2}{x (x - 1)}$$

which simplifies to

$$y = \frac{-c_1 x + c_2}{x \left(x - 1\right)}$$

Summary

The solution(s) found are the following

$$y = \frac{-c_1 x + c_2}{x \left(x - 1\right)} \tag{1}$$

Verification of solutions

$$y = \frac{-c_1 x + c_2}{x \left(x - 1\right)}$$

Verified OK.

#### 3.2.7 Maple step by step solution

Let's solve

$$(-x^{2} + x) y'' + (-4x + 2) y' - 2y = 0$$

- Highest derivative means the order of the ODE is 2 y''
- Isolate 2nd derivative

$$y'' = -\frac{2y}{x(x-1)} - \frac{2(2x-1)y'}{x(x-1)}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear  $y'' + \frac{2(2x-1)y'}{x(x-1)} + \frac{2y}{x(x-1)} = 0$
- $\Box$  Check to see if  $x_0$  is a regular singular point
  - $\circ$  Define functions

$$\left[P_2(x) = \frac{2(2x-1)}{x(x-1)}, P_3(x) = \frac{2}{x(x-1)}\right]$$

•  $x \cdot P_2(x)$  is analytic at x = 0

$$(x \cdot P_2(x)) \bigg|_{x=0} = 2$$

 $\circ \quad x^2 \cdot P_3(x) \text{ is analytic at } x = 0$ 

$$\left(x^2 \cdot P_3(x)\right) \bigg|_{x=0} = 0$$

 $\circ \quad x = 0 \text{is a regular singular point}$ 

Check to see if  $x_0$  is a regular singular point

 $x_0 = 0$ 

• Multiply by denominators

$$y''x(x-1) + (4x-2)y' + 2y = 0$$

• Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^{k+r}$$

 $\Box$  Rewrite ODE with series expansions

• Convert  $x^m \cdot y'$  to series expansion for m = 0..1

$$x^m \cdot y' = \sum_{k=0}^{\infty} a_k(k+r) x^{k+r-1+m}$$

• Shift index using k - >k + 1 - m

$$x^m \cdot y' = \sum_{k=-1+m}^{\infty} a_{k+1-m}(k+1-m+r) \, x^{k+r}$$

• Convert  $x^m \cdot y''$  to series expansion for m = 1..2

$$x^m\cdot y'' = \sum\limits_{k=0}^\infty a_k(k+r) \left(k+r-1\right) x^{k+r-2+m}$$

 $\circ \quad \text{Shift index using } k->k+2-m$ 

$$x^{m} \cdot y'' = \sum_{k=-2+m}^{\infty} a_{k+2-m}(k+2-m+r) \left(k+1-m+r\right) x^{k+r}$$

Rewrite ODE with series expansions

$$-a_0 r(1+r) x^{-1+r} + \left(\sum_{k=0}^{\infty} \left(-a_{k+1}(k+r+1) \left(k+r+2\right) + a_k(k+r+2) \left(k+r+1\right)\right) x^{k+r}\right) = 0$$

•  $a_0$  cannot be 0 by assumption, giving the indicial equation

-r(1+r) = 0

• Values of r that satisfy the indicial equation

 $r\in\{-1,0\}$ 

- Each term in the series must be 0, giving the recursion relation  $(k+r+2)(k+r+1)(-a_{k+1}+a_k) = 0$
- Recursion relation that defines series solution to ODE

 $a_{k+1} = a_k$ 

• Recursion relation for r = -1

 $a_{k+1} = a_k$ 

• Solution for r = -1

$$\left[y=\sum\limits_{k=0}^{\infty}a_kx^{k-1},a_{k+1}=a_k
ight]$$

• Recursion relation for r = 0

 $a_{k+1} = a_k$ 

• Solution for r = 0

$$igg[y=\sum\limits_{k=0}^\infty a_k x^k,a_{k+1}=a_kigg]$$

• Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k x^{k-1}\right) + \left(\sum_{k=0}^{\infty} b_k x^k\right), a_{k+1} = a_k, b_{k+1} = b_k\right]$$

#### Maple trace

`Methods for second order ODEs: --- Trying classification methods --trying a quadrature checking if the LODE has constant coefficients checking if the LODE is of Euler type trying a symmetry of the form [xi=0, eta=F(x)] <- linear\_1 successful`</pre> Solution by Maple Time used: 0.0 (sec). Leaf size: 18

dsolve(x\*(1-x)\*diff(y(x),x\$2)+2\*(1-2\*x)\*diff(y(x),x)-2\*y(x)=0,y(x), singsol=all)

$$y(x) = rac{c_1 x + c_2}{x (x - 1)}$$

Solution by Mathematica Time used: 0.035 (sec). Leaf size: 22

DSolve[x\*(1-x)\*y''[x]+2\*(1-2\*x)\*y'[x]-2\*y[x]==0,y[x],x,IncludeSingularSolutions -> True]

$$y(x) \to \frac{c_2 x + c_1}{x - x^2}$$

# 3.3 problem 10.4.8 (c)

-						
3.3.1	Solving as second order euler ode ode					
3.3.2 Solving as second order change of variable on x method 2 ode .						
3.3.3	Solving as second order change of variable on x method 1 ode $\therefore$ 263					
3.3.4	Solving as second order change of variable on y method 2 ode $\ . \ 265$					
3.3.5	Solving using Kovacic algorithm					
3.3.6	Maple step by step solution					
Internal problem ID [5066]						
Internal file name [OUTPUT/4559_Sunday_June_05_2022_03_00_53_PM_83484169/index.tex]						
Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995						
Section: Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients.						
Second order and Homogeneous. page 318						
<b>Problem number</b> : 10.4.8 (c).						
ODE order: 2.						
ODE dogmos: 1						

**ODE degree**: 1.

The type(s) of ODE detected by this program : "kovacic", "second\_order\_euler\_ode", "second\_order\_change\_of\_variable\_on\_x\_method\_1", "second\_order\_change\_of\_variable\_on\_x\_method\_2", "second\_order\_change\_of\_variable\_on\_y\_method\_2"

Maple gives the following as the ode type

```
[[_Emden, _Fowler], [_2nd_order, _linear, `_with_symmetry_[0,F(
    x)]`]]
```

$$x^2y'' + xy' - 9y = 0$$

### 3.3.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be  $y = x^r$ , then  $y' = rx^{r-1}$  and  $y'' = r(r-1)x^{r-2}$ . Substituting these back into the given ODE gives

$$x^{2}(r(r-1))x^{r-2} + xrx^{r-1} - 9x^{r} = 0$$

Simplifying gives

$$r(r-1) x^r + r x^r - 9x^r = 0$$

Since  $x^r \neq 0$  then dividing throughout by  $x^r$  gives

$$r(r-1) + r - 9 = 0$$

$$r^2 - 9 = 0$$
(1)

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

 $r_1 = -3$  $r_2 = 3$ 

Since the roots are real and distinct, then the general solution is

$$y = c_1 y_1 + c_2 y_2$$

Where  $y_1 = x^{r_1}$  and  $y_2 = x^{r_2}$ . Hence

$$y = \frac{c_1}{x^3} + c_2 x^3$$

Summary

Or

The solution(s) found are the following

$$y = \frac{c_1}{x^3} + c_2 x^3 \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{x^3} + c_2 x^3$$

Verified OK.

### 3.3.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$x^2y'' + xy' - 9y = 0 \tag{1}$$

Becomes

$$y'' + p(x) y' + q(x) y = 0$$
(2)

Where

$$p(x) = \frac{1}{x}$$
$$q(x) = -\frac{9}{x^2}$$

Applying change of variables  $\tau = g(x)$  to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0$$
(3)

Where  $\tau$  is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(x) + p(x)\,\tau'(x)}{\tau'(x)^2} \tag{4}$$

$$q_1(\tau) = \frac{q(x)}{\tau'(x)^2} \tag{5}$$

Let  $p_1 = 0$ . Eq (4) simplifies to

$$\tau''(x) + p(x)\,\tau'(x) = 0$$

This ode is solved resulting in

$$\tau = \int e^{-(\int p(x)dx)} dx$$
  
=  $\int e^{-(\int \frac{1}{x}dx)} dx$   
=  $\int e^{-\ln(x)} dx$   
=  $\int \frac{1}{x}dx$   
=  $\ln(x)$  (6)

Using (6) to evaluate  $q_1$  from (5) gives

$$q_{1}(\tau) = \frac{q(x)}{\tau'(x)^{2}} = \frac{-\frac{9}{x^{2}}}{\frac{1}{x^{2}}} = -9$$
(7)

Substituting the above in (3) and noting that now  $p_1 = 0$  results in

$$egin{aligned} &rac{d^2}{d au^2}y( au)+q_1y( au)=0\ &rac{d^2}{d au^2}y( au)-9y( au)=0 \end{aligned}$$

The above ode is now solved for  $y(\tau)$ . This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(\tau) + By'(\tau) + Cy(\tau) = 0$$

Where in the above A = 1, B = 0, C = -9. Let the solution be  $y(\tau) = e^{\lambda \tau}$ . Substituting this into the ODE gives

$$\lambda^2 \mathrm{e}^{\lambda\tau} - 9\,\mathrm{e}^{\lambda\tau} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by  $e^{\lambda \tau}$  gives

$$\lambda^2 - 9 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting A = 1, B = 0, C = -9 into the above gives

$$\lambda_{1,2} = \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(-9)}$$
$$= \pm 3$$

Hence

$$\lambda_1 = +3$$
$$\lambda_2 = -3$$

Which simplifies to

$$\lambda_1 = 3$$
$$\lambda_2 = -3$$

Since roots are real and distinct, then the solution is

$$y(\tau) = c_1 e^{\lambda_1 \tau} + c_2 e^{\lambda_2 \tau}$$
$$y(\tau) = c_1 e^{(3)\tau} + c_2 e^{(-3)\tau}$$

Or

$$y(\tau) = c_1 e^{3\tau} + c_2 e^{-3\tau}$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{c_1 x^6 + c_2}{x^3}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 x^6 + c_2}{x^3} \tag{1}$$

Verification of solutions

$$y = \frac{c_1 x^6 + c_2}{x^3}$$

Verified OK.

### 3.3.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$x^2y'' + xy' - 9y = 0 \tag{1}$$

Becomes

$$y'' + p(x) y' + q(x) y = 0$$
(2)

Where

$$p(x) = \frac{1}{x}$$
$$q(x) = -\frac{9}{x^2}$$

Applying change of variables  $\tau = g(x)$  to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0$$
(3)

Where  $\tau$  is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(x) + p(x)\,\tau'(x)}{\tau'(x)^2} \tag{4}$$

$$q_1(\tau) = \frac{q(x)}{\tau'(x)^2} \tag{5}$$

Let  $q_1 = c^2$  where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c}\sqrt{q}$$

$$= \frac{3\sqrt{-\frac{1}{x^2}}}{c}$$

$$\tau'' = \frac{3}{c\sqrt{-\frac{1}{x^2}}x^3}$$
(6)

Substituting the above into (4) results in

$$p_{1}(\tau) = \frac{\tau''(x) + p(x)\tau'(x)}{\tau'(x)^{2}}$$
$$= \frac{\frac{3}{c\sqrt{-\frac{1}{x^{2}}}x^{3}} + \frac{1}{x}\frac{3\sqrt{-\frac{1}{x^{2}}}}{c}}{\left(\frac{3\sqrt{-\frac{1}{x^{2}}}}{c}\right)^{2}}$$
$$= 0$$

Therefore ode (3) now becomes

$$y(\tau)'' + p_1 y(\tau)' + q_1 y(\tau) = 0$$
  
$$\frac{d^2}{d\tau^2} y(\tau) + c^2 y(\tau) = 0$$
(7)

The above ode is now solved for  $y(\tau)$ . Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = c_1 \cos\left(c\tau\right) + c_2 \sin\left(c\tau\right)$$

Now from (6)

$$\tau = \int \frac{1}{c} \sqrt{q} \, dx$$
$$= \frac{\int 3\sqrt{-\frac{1}{x^2}} \, dx}{\frac{c}{3\sqrt{-\frac{1}{x^2}} x \ln(x)}}$$

Substituting the above into the solution obtained gives

$$y = c_1 \cosh \left(3 \ln \left(x\right)\right) + i c_2 \sinh \left(3 \ln \left(x\right)\right)$$

 $\frac{\text{Summary}}{\text{The solution(s) found are the following}}$ 

$$y = c_1 \cosh(3\ln(x)) + ic_2 \sinh(3\ln(x))$$
(1)

Verification of solutions

$$y = c_1 \cosh\left(3\ln\left(x\right)\right) + ic_2 \sinh\left(3\ln\left(x\right)\right)$$

Verified OK.

#### 3.3.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$x^2y'' + xy' - 9y = 0 \tag{1}$$

Becomes

$$y'' + p(x) y' + q(x) y = 0$$
(2)

Where

$$p(x) = \frac{1}{x}$$
$$q(x) = -\frac{9}{x^2}$$

Applying change of variables on the dependent variable  $y = v(x) x^n$  to (2) gives the following ode where the dependent variables is v(x) and not y.

$$v''(x) + \left(\frac{2n}{x} + p\right)v'(x) + \left(\frac{n(n-1)}{x^2} + \frac{np}{x} + q\right)v(x) = 0$$
(3)

Let the coefficient of v(x) above be zero. Hence

$$\frac{n(n-1)}{x^2} + \frac{np}{x} + q = 0 \tag{4}$$

Substituting the earlier values found for p(x) and q(x) into (4) gives

$$\frac{n(n-1)}{x^2} + \frac{n}{x^2} - \frac{9}{x^2} = 0 \tag{5}$$

Solving (5) for n gives

$$n = 3 \tag{6}$$

Substituting this value in (3) gives

$$v''(x) + \frac{7v'(x)}{x} = 0$$
  
$$v''(x) + \frac{7v'(x)}{x} = 0$$
 (7)

Using the substitution

$$u(x) = v'(x)$$

Then (7) becomes

$$u'(x) + \frac{7u(x)}{x} = 0$$
 (8)

The above is now solved for u(x). In canonical form the ODE is

$$u' = F(x, u)$$
  
=  $f(x)g(u)$   
=  $-\frac{7u}{x}$ 

Where  $f(x) = -\frac{7}{x}$  and g(u) = u. Integrating both sides gives

$$\frac{1}{u} du = -\frac{7}{x} dx$$
$$\int \frac{1}{u} du = \int -\frac{7}{x} dx$$
$$\ln(u) = -7\ln(x) + c_1$$
$$u = e^{-7\ln(x) + c_1}$$
$$= \frac{c_1}{x^7}$$

Now that u(x) is known, then

$$egin{aligned} v'(x) &= u(x) \ v(x) &= \int u(x) \, dx + c_2 \ &= -rac{c_1}{6x^6} + c_2 \end{aligned}$$

Hence

$$y = v(x) x^{n}$$
  
=  $\left(-\frac{c_{1}}{6x^{6}} + c_{2}\right) x^{3}$   
=  $\frac{6c_{2}x^{6} - c_{1}}{6x^{3}}$ 

 $\frac{Summary}{The solution(s) found are the following}$ 

$$y = \left(-\frac{c_1}{6x^6} + c_2\right)x^3$$
 (1)

Verification of solutions

$$y = \left(-\frac{c_1}{6x^6} + c_2\right)x^3$$

Verified OK.

# 3.3.5 Solving using Kovacic algorithm

Writing the ode as

$$x^2y'' + xy' - 9y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = x^{2}$$
  

$$B = x$$
  

$$C = -9$$
  
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int rac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{35}{4x^2} \tag{6}$$

Comparing the above to (5) shows that

$$s = 35$$
$$t = 4x^2$$

Therefore eq. (4) becomes

$$z''(x) = \left(\frac{35}{4x^2}\right)z(x) \tag{7}$$

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) \, e^{-\int rac{B}{2A} \, dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 48: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 2 - 0$$
$$= 2$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of  $t = 4x^2$ . There is a pole at x = 0 of order 2. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case n = 1.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{35}{4x^2}$$

For the pole at x = 0 let b be the coefficient of  $\frac{1}{x^2}$  in the partial fractions decomposition of r given above. Therefore  $b = \frac{35}{4}$ . Hence

$$\begin{split} [\sqrt{r}]_c &= 0\\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1+4b} = \frac{7}{2}\\ \alpha_c^- &= \frac{1}{2} - \sqrt{1+4b} = -\frac{5}{2} \end{split}$$

Since the order of r at  $\infty$  is 2 then  $[\sqrt{r}]_{\infty} = 0$ . Let b be the coefficient of  $\frac{1}{x^2}$  in the Laurent series expansion of r at  $\infty$ , which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = \frac{35}{4x^2}$$

Since the gcd(s,t) = 1. This gives  $b = \frac{35}{4}$ . Hence

$$\begin{split} [\sqrt{r}]_{\infty} &= 0 \\ \alpha_{\infty}^{+} &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{7}{2} \\ \alpha_{\infty}^{-} &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{5}{2} \end{split}$$

The following table summarizes the findings so far for poles and for the order of r at  $\infty$  where r is

$$r = \frac{35}{4x^2}$$

pole $c$ location		pole order		۱]	$\sqrt{r}]_c$	$\alpha_c^+$	$\alpha_c^-$
0		2		0		$\frac{7}{2}$	$-\frac{5}{2}$
	Order of $r$ at $\infty$		$[\sqrt{r}]_{c}$	$\infty$	$lpha^+_\infty$	$\alpha_{\infty}^{-}$	
	2		0		$\frac{7}{2}$	$-\frac{5}{2}$	

Now that the all  $[\sqrt{r}]_c$  and its associated  $\alpha_c^{\pm}$  have been determined for all the poles in the set  $\Gamma$  and  $[\sqrt{r}]_{\infty}$  and its associated  $\alpha_{\infty}^{\pm}$  have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_{\infty}^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where s(c) is either + or - and  $s(\infty)$  is the sign of  $\alpha_{\infty}^{\pm}$ . This is done by trial over all set of families  $s = (s(c))_{c \in \Gamma \cup \infty}$  until such d is found to work in finding candidate  $\omega$ . Trying  $\alpha_{\infty}^{-} = -\frac{5}{2}$  then

$$d = \alpha_{\infty}^{-} - (\alpha_{c_1}^{-})$$
$$= -\frac{5}{2} - \left(-\frac{5}{2}\right)$$
$$= 0$$

Since d an integer and  $d \ge 0$  then it can be used to find  $\omega$  using

$$\omega = \sum_{c \in \Gamma} \left( s(c) [\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{x - c} \right) + s(\infty) [\sqrt{r}]_{\infty}$$

The above gives

$$\begin{split} \omega &= \left( (-) [\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}}{x - c_1} \right) + (-) [\sqrt{r}]_{\infty} \\ &= -\frac{5}{2x} + (-) (0) \\ &= -\frac{5}{2x} \\ &= -\frac{5}{2x} \end{split}$$

Now that  $\omega$  is determined, the next step is find a corresponding minimal polynomial p(x) of degree d = 0 to solve the ode. The polynomial p(x) needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0$$
(1A)

Let

$$p(x) = 1 \tag{2A}$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(-\frac{5}{2x}\right)(0) + \left(\left(\frac{5}{2x^2}\right) + \left(-\frac{5}{2x}\right)^2 - \left(\frac{35}{4x^2}\right)\right) = 0$$
$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode  $z^{\prime\prime}=rz$  is

$$egin{aligned} z_1(x) &= p e^{\int \omega \, dx} \ &= e^{\int -rac{5}{2x} dx} \ &= rac{1}{x^{rac{5}{2}}} \end{aligned}$$

The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$
$$= z_1 e^{-\int \frac{1}{2} \frac{x}{x^2} dx}$$
$$= z_1 e^{-\frac{\ln(x)}{2}}$$
$$= z_1 \left(\frac{1}{\sqrt{x}}\right)$$

Which simplifies to

$$y_1 = \frac{1}{x^3}$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} \, dx}}{y_1^2} \, dx$$

Substituting gives

$$y_{2} = y_{1} \int \frac{e^{\int -\frac{x}{x^{2}} dx}}{(y_{1})^{2}} dx$$
$$= y_{1} \int \frac{e^{-\ln(x)}}{(y_{1})^{2}} dx$$
$$= y_{1} \left(\frac{x^{6}}{6}\right)$$

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
$$= c_1 \left(\frac{1}{x^3}\right) + c_2 \left(\frac{1}{x^3} \left(\frac{x^6}{6}\right)\right)$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{x^3} + \frac{c_2 x^3}{6} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{x^3} + \frac{c_2 x^3}{6}$$

Verified OK.

### 3.3.6 Maple step by step solution

Let's solve

 $x^2y'' + xy' - 9y = 0$ 

• Highest derivative means the order of the ODE is 2

• Isolate 2nd derivative

$$y'' = -\frac{y'}{x} + \frac{9y}{x^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear  $y'' + \frac{y'}{x} \frac{9y}{x^2} = 0$
- Multiply by denominators of the ODE

$$x^2y^{\prime\prime} + xy^\prime - 9y = 0$$

• Make a change of variables

$$t = \ln\left(x\right)$$

- $\Box$  Substitute the change of variables back into the ODE
  - $\circ~$  Calculate the 1st derivative of y with respect to x , using the chain rule  $y' = \left( \frac{d}{dt} y(t) \right) t'(x)$

• Compute derivative

$$y' = rac{rac{d}{dt}y(t)}{x}$$

- Calculate the 2nd derivative of y with respect to x , using the chain rule  $y'' = \left(\frac{d^2}{dt^2}y(t)\right)t'(x)^2 + t''(x)\left(\frac{d}{dt}y(t)\right)$
- Compute derivative

$$y'' = rac{rac{d^2}{dt^2}y(t)}{x^2} - rac{rac{d}{dt}y(t)}{x^2}$$

Substitute the change of variables back into the ODE

$$x^{2}\left(\frac{\frac{d^{2}}{dt^{2}}y(t)}{x^{2}} - \frac{\frac{d}{dt}y(t)}{x^{2}}\right) + \frac{d}{dt}y(t) - 9y(t) = 0$$

• Simplify

$$\frac{d^2}{dt^2}y(t) - 9y(t) = 0$$

• Characteristic polynomial of ODE

$$r^2 - 9 = 0$$

• Factor the characteristic polynomial

$$(r-3)(r+3) = 0$$

- Roots of the characteristic polynomial r = (-3, 3)
- 1st solution of the ODE

$$y_1(t) = e^{-3t}$$

• 2nd solution of the ODE

$$y_2(t) = e^{3t}$$

• General solution of the ODE

$$y(t) = c_1 y_1(t) + c_2 y_2(t)$$

• Substitute in solutions

$$y(t) = c_1 e^{-3t} + c_2 e^{3t}$$

• Change variables back using  $t = \ln(x)$ 

$$y = \frac{c_1}{x^3} + c_2 x$$

• Simplify

$$y = \frac{c_1}{x^3} + c_2 x^3$$

#### Maple trace

`Methods for second order ODEs: --- Trying classification methods --trying a quadrature checking if the LODE has constant coefficients checking if the LODE is of Euler type <- LODE of Euler type successful`</pre>

Solution by Maple Time used: 0.0 (sec). Leaf size: 15

 $dsolve(x^2*diff(y(x),x^2)+x*diff(y(x),x)-9*y(x)=0,y(x), singsol=all)$ 

$$y(x) = \frac{c_2 x^6 + c_1}{x^3}$$

Solution by Mathematica Time used: 0.011 (sec). Leaf size: 18

DSolve[x^2\*y''[x]+x\*y'[x]-9\*y[x]==0,y[x],x,IncludeSingularSolutions -> True]

$$y(x) \to \frac{c_2 x^6 + c_1}{x^3}$$

# 3.4 problem 10.4.8 (d)

-						
3.4.1	Solving as second order change of variable on x method 2 ode $\therefore$ 275					
3.4.2 Solving as second order change of variable on x method 1 ode . 278						
3.4.3	3.4.3 Solving as second order bessel ode ode					
3.4.4	Solving using Kovacic algorithm					
3.4.5	Maple step by step solution					
Internal problem ID [5067]						
$Internal file name \left[ \texttt{OUTPUT/4560\_Sunday\_June\_05\_2022\_03\_00\_54\_PM\_41197975/index.tex} \right]$						
Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995						
Section: Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients.						
Second order and Homogeneous. page 318						
Problem number: 10.4.8 (d).						
ODE order: 2.						
ODE degree: 1.						

The type(s) of ODE detected by this program : "kovacic", "second\_order\_bessel\_ode", "second\_order\_change\_of\_variable\_on\_x\_method\_1", "second\_order\_change\_of\_variable\_on\_x\_method\_2"

Maple gives the following as the ode type

$$xy'' + \frac{y'}{2} + 2y = 0$$

### 3.4.1 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$xy'' + \frac{y'}{2} + 2y = 0 \tag{1}$$

Becomes

$$y'' + p(x) y' + q(x) y = 0$$
(2)

Where

$$p(x) = \frac{1}{2x}$$
$$q(x) = \frac{2}{x}$$

Applying change of variables  $\tau = g(x)$  to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0$$
(3)

Where  $\tau$  is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(x) + p(x)\tau'(x)}{\tau'(x)^2}$$
(4)

$$q_1(\tau) = \frac{q(x)}{\tau'(x)^2} \tag{5}$$

Let  $p_1 = 0$ . Eq (4) simplifies to

$$\tau''(x) + p(x)\tau'(x) = 0$$

This ode is solved resulting in

$$\tau = \int e^{-(\int p(x)dx)} dx$$
  
=  $\int e^{-(\int \frac{1}{2x}dx)} dx$   
=  $\int e^{-\frac{\ln(x)}{2}} dx$   
=  $\int \frac{1}{\sqrt{x}} dx$   
=  $2\sqrt{x}$  (6)

Using (6) to evaluate  $q_1$  from (5) gives

$$q_{1}(\tau) = \frac{q(x)}{\tau'(x)^{2}}$$
$$= \frac{\frac{2}{x}}{\frac{1}{x}}$$
$$= 2$$
(7)

Substituting the above in (3) and noting that now  $p_1 = 0$  results in

$$egin{aligned} &rac{d^2}{d au^2}y( au)+q_1y( au)=0\ &rac{d^2}{d au^2}y( au)+2y( au)=0 \end{aligned}$$

The above ode is now solved for  $y(\tau)$ . This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(\tau) + By'(\tau) + Cy(\tau) = 0$$

Where in the above A = 1, B = 0, C = 2. Let the solution be  $y(\tau) = e^{\lambda \tau}$ . Substituting this into the ODE gives

$$\lambda^2 \mathrm{e}^{\lambda \tau} + 2 \, \mathrm{e}^{\lambda \tau} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by  $e^{\lambda \tau}$  gives

$$\lambda^2 + 2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting A = 1, B = 0, C = 2 into the above gives

$$\lambda_{1,2} = \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(2)}$$
$$= \pm i\sqrt{2}$$

Hence

$$\lambda_1 = +i\sqrt{2}$$
$$\lambda_2 = -i\sqrt{2}$$

Which simplifies to

$$\lambda_1 = i\sqrt{2}$$
$$\lambda_2 = -i\sqrt{2}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where  $\alpha = 0$  and  $\beta = \sqrt{2}$ . Therefore the final solution, when using Euler relation, can be written as

$$y(\tau) = e^{\alpha \tau} (c_1 \cos(\beta \tau) + c_2 \sin(\beta \tau))$$

Which becomes

$$y(\tau) = e^0 \left( c_1 \cos \left( \sqrt{2} \tau \right) + c_2 \sin \left( \sqrt{2} \tau \right) \right)$$

Or

$$y(\tau) = c_1 \cos\left(\sqrt{2}\,\tau\right) + c_2 \sin\left(\sqrt{2}\,\tau\right)$$

The above solution is now transformed back to y using (6) which results in

$$y = c_1 \cos\left(2\sqrt{2}\sqrt{x}\right) + c_2 \sin\left(2\sqrt{2}\sqrt{x}\right)$$

Summary

The solution(s) found are the following

$$y = c_1 \cos\left(2\sqrt{2}\sqrt{x}\right) + c_2 \sin\left(2\sqrt{2}\sqrt{x}\right) \tag{1}$$

Verification of solutions

$$y = c_1 \cos\left(2\sqrt{2}\sqrt{x}\right) + c_2 \sin\left(2\sqrt{2}\sqrt{x}\right)$$

Verified OK.

# 3.4.2 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$xy'' + \frac{y'}{2} + 2y = 0 \tag{1}$$

Becomes

$$y'' + p(x) y' + q(x) y = 0$$
(2)

Where

$$p(x) = \frac{1}{2x}$$
$$q(x) = \frac{2}{x}$$

Applying change of variables  $\tau = g(x)$  to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0$$
(3)

Where  $\tau$  is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(x) + p(x)\,\tau'(x)}{\tau'(x)^2} \tag{4}$$

$$q_1(\tau) = \frac{q(x)}{\tau'(x)^2} \tag{5}$$

Let  $q_1 = c^2$  where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c}\sqrt{q}$$

$$= \frac{\sqrt{2}\sqrt{\frac{1}{x}}}{c}$$

$$\tau'' = -\frac{\sqrt{2}}{2c\sqrt{\frac{1}{x}}x^2}$$
(6)

Substituting the above into (4) results in

$$p_{1}(\tau) = \frac{\tau''(x) + p(x) \tau'(x)}{\tau'(x)^{2}}$$
$$= \frac{-\frac{\sqrt{2}}{2c\sqrt{\frac{1}{x}x^{2}}} + \frac{1}{2x}\frac{\sqrt{2}\sqrt{\frac{1}{x}}}{c}}{\left(\frac{\sqrt{2}\sqrt{\frac{1}{x}}}{c}\right)^{2}}$$
$$= 0$$

Therefore ode (3) now becomes

$$y(\tau)'' + p_1 y(\tau)' + q_1 y(\tau) = 0$$
  
$$\frac{d^2}{d\tau^2} y(\tau) + c^2 y(\tau) = 0$$
 (7)

The above ode is now solved for  $y(\tau)$ . Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = c_1 \cos(c\tau) + c_2 \sin(c\tau)$$

Now from (6)

$$\tau = \int \frac{1}{c} \sqrt{q} \, dx$$
$$= \frac{\int \sqrt{2} \sqrt{\frac{1}{x}} dx}{c}$$
$$= \frac{2x\sqrt{2}\sqrt{\frac{1}{x}}}{c}$$

Substituting the above into the solution obtained gives

$$y = c_1 \cos\left(2\sqrt{2}\sqrt{x}\right) + c_2 \sin\left(2\sqrt{2}\sqrt{x}\right)$$

#### Summary

The solution(s) found are the following

$$y = c_1 \cos\left(2\sqrt{2}\sqrt{x}\right) + c_2 \sin\left(2\sqrt{2}\sqrt{x}\right) \tag{1}$$

Verification of solutions

$$y = c_1 \cos\left(2\sqrt{2}\sqrt{x}\right) + c_2 \sin\left(2\sqrt{2}\sqrt{x}\right)$$

Verified OK.

### 3.4.3 Solving as second order bessel ode ode

Writing the ode as

$$x^2y'' + \frac{xy'}{2} + 2xy = 0 \tag{1}$$

Bessel ode has the form

$$x^{2}y'' + xy' + (-n^{2} + x^{2})y = 0$$
<sup>(2)</sup>

The generalized form of Bessel ode is given by Bowman (1958) as the following

$$x^{2}y'' + (1 - 2\alpha)xy' + \left(\beta^{2}\gamma^{2}x^{2\gamma} - n^{2}\gamma^{2} + \alpha^{2}\right)y = 0$$
(3)

With the standard solution

$$y = x^{\alpha}(c_1 \operatorname{BesselJ}(n, \beta x^{\gamma}) + c_2 \operatorname{BesselY}(n, \beta x^{\gamma}))$$
(4)

Comparing (3) to (1) and solving for  $\alpha,\beta,n,\gamma$  gives

$$\alpha = \frac{1}{4}$$
$$\beta = 2\sqrt{2}$$
$$n = \frac{1}{2}$$
$$\gamma = \frac{1}{2}$$

Substituting all the above into (4) gives the solution as

$$y = \frac{c_1 x^{\frac{1}{4}} \sin\left(2\sqrt{2}\sqrt{x}\right)}{\sqrt{\pi}\sqrt{\sqrt{2}\sqrt{x}}} - \frac{c_2 x^{\frac{1}{4}} \cos\left(2\sqrt{2}\sqrt{x}\right)}{\sqrt{\pi}\sqrt{\sqrt{2}\sqrt{x}}}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 x^{\frac{1}{4}} \sin\left(2\sqrt{2}\sqrt{x}\right)}{\sqrt{\pi}\sqrt{\sqrt{2}\sqrt{x}}} - \frac{c_2 x^{\frac{1}{4}} \cos\left(2\sqrt{2}\sqrt{x}\right)}{\sqrt{\pi}\sqrt{\sqrt{2}\sqrt{x}}}$$
(1)

Verification of solutions

$$y = \frac{c_1 x^{\frac{1}{4}} \sin\left(2\sqrt{2}\sqrt{x}\right)}{\sqrt{\pi}\sqrt{\sqrt{2}\sqrt{x}}} - \frac{c_2 x^{\frac{1}{4}} \cos\left(2\sqrt{2}\sqrt{x}\right)}{\sqrt{\pi}\sqrt{\sqrt{2}\sqrt{x}}}$$

Verified OK.

## 3.4.4 Solving using Kovacic algorithm

Writing the ode as

$$xy'' + \frac{y'}{2} + 2y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = x$$
  

$$B = \frac{1}{2}$$

$$C = 2$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-32x - 3}{16x^2} \tag{6}$$

Comparing the above to (5) shows that

$$s = -32x - 3$$
$$t = 16x^2$$

Therefore eq. (4) becomes

$$z''(x) = \left(\frac{-32x - 3}{16x^2}\right) z(x)$$
(7)

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 50: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 2 - 1$$
$$= 1$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of  $t = 16x^2$ . There is a pole at x = 0 of order 2. Since there is a pole of order 2 then necessary conditions for case two are met. Therefore

L = [2]

Attempting to find a solution using case n = 2.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{2}{x} - \frac{3}{16x^2}$$

For the pole at x = 0 let b be the coefficient of  $\frac{1}{x^2}$  in the partial fractions decomposition of r given above. Therefore  $b = -\frac{3}{16}$ . Hence

$$E_c = \{2, 2 + 2\sqrt{1+4b}, 2 - 2\sqrt{1+4b}\}\$$
  
=  $\{1, 2, 3\}$ 

Since the order of r at  $\infty$  is 1 < 2 then

 $E_{\infty} = \{1\}$ 

The following table summarizes the findings so far for poles and for the order of r at  $\infty$  for case 2 of Kovacic algorithm.

pole $c$ location		pole order		$E_c$	
0		2		$\{1, 2, 3\}$	
Order of $r$ at $\infty$			$E_{c}$	$\infty$	
	1		{1	.}	

Using the family  $\{e_1, e_2, \ldots, e_\infty\}$  given by

$$e_1 = 1, \ e_\infty = 1$$

Gives a non negative integer d (the degree of the polynomial p(x)), which is generated using

$$d = \frac{1}{2} \left( e_{\infty} - \sum_{c \in \Gamma} e_c \right)$$
$$= \frac{1}{2} (1 - (1))$$
$$= 0$$

We now form the following rational function

$$\theta = \frac{1}{2} \sum_{c \in \Gamma} \frac{e_c}{x - c}$$
$$= \frac{1}{2} \left( \frac{1}{(x - (0))} \right)$$
$$= \frac{1}{2x}$$

Now we search for a monic polynomial p(x) of degree d = 0 such that

$$p''' + 3\theta p'' + (3\theta^2 + 3\theta' - 4r) p' + (\theta'' + 3\theta\theta' + \theta^3 - 4r\theta - 2r') p = 0$$
(1A)

Since d = 0, then letting

$$p = 1 \tag{2A}$$

Substituting p and  $\theta$  into Eq. (1A) gives

0 = 0

And solving for p gives

p = 1

Now that p(x) is found let

$$\phi = \theta + \frac{p'}{p}$$
$$= \frac{1}{2x}$$

Let  $\omega$  be the solution of

$$\omega^2 - \phi\omega + \left(\frac{1}{2}\phi' + \frac{1}{2}\phi^2 - r\right) = 0$$

Substituting the values for  $\phi$  and r into the above equation gives

$$w^2 - \frac{w}{2x} + \frac{1+32x}{16x^2} = 0$$

Solving for  $\omega$  gives

$$\omega = \frac{1 + 4\sqrt{2}\sqrt{-x}}{4x}$$

Therefore the first solution to the ode z'' = rz is

$$egin{aligned} z_1(x) &= e^{\int \omega \, dx} \ &= \mathrm{e}^{\int rac{1+4\sqrt{2}\sqrt{-x}}{4x}dx} \ &= x^{rac{1}{4}}\mathrm{e}^{2\sqrt{2}\sqrt{-x}} \end{aligned}$$

The first solution to the original ode in y is found from

$$y_{1} = z_{1}e^{\int -\frac{1}{2}\frac{B}{A}dx}$$
  
=  $z_{1}e^{-\int \frac{1}{2}\frac{1}{x}dx}$   
=  $z_{1}e^{-\frac{\ln(x)}{4}}$   
=  $z_{1}\left(\frac{1}{x^{\frac{1}{4}}}\right)$ 

Which simplifies to

$$y_1 = e^{2\sqrt{2}\sqrt{-x}}$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} \, dx}}{y_1^2} \, dx$$

Substituting gives

$$y_{2} = y_{1} \int \frac{e^{\int -\frac{1}{x} dx}}{(y_{1})^{2}} dx$$
  
=  $y_{1} \int \frac{e^{-\frac{\ln(x)}{2}}}{(y_{1})^{2}} dx$   
=  $y_{1} \left( \frac{\sqrt{2}\sqrt{-x} \left(-1 + e^{-4\sqrt{2}\sqrt{-x}}\right)}{4\sqrt{x}} \right)$ 

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
  
=  $c_1 \left( e^{2\sqrt{2}\sqrt{-x}} \right) + c_2 \left( e^{2\sqrt{2}\sqrt{-x}} \left( \frac{\sqrt{2}\sqrt{-x} \left( -1 + e^{-4\sqrt{2}\sqrt{-x}} \right)}{4\sqrt{x}} \right) \right)$ 

# Summary

The solution(s) found are the following

$$y = c_1 e^{2\sqrt{2}\sqrt{-x}} - \frac{c_2\sqrt{2}\sqrt{-x}\left(e^{2\sqrt{2}\sqrt{-x}} - e^{-2\sqrt{2}\sqrt{-x}}\right)}{4\sqrt{x}}$$
(1)

Verification of solutions

$$y = c_1 e^{2\sqrt{2}\sqrt{-x}} - \frac{c_2\sqrt{2}\sqrt{-x} \left(e^{2\sqrt{2}\sqrt{-x}} - e^{-2\sqrt{2}\sqrt{-x}}\right)}{4\sqrt{x}}$$

Verified OK.

# 3.4.5 Maple step by step solution

Let's solve

$$y''x + \frac{y'}{2} + 2y = 0$$

- Highest derivative means the order of the ODE is 2 y''
- Isolate 2nd derivative

$$y'' = -\frac{y'}{2x} - \frac{2y}{x}$$

• Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear  $y'' + \frac{y'}{2x} + \frac{2y}{x} = 0$ 

0

Define functions

$$\left[P_2(x)=rac{1}{2x},P_3(x)=rac{2}{x}
ight]$$

 $\circ \quad x \cdot P_2(x) \text{ is analytic at } x = 0$ 

$$(x \cdot P_2(x)) \bigg|_{x=0} = \frac{1}{2}$$

- $\circ \quad x^2 \cdot P_3(x) ext{ is analytic at } x = 0$  $(x^2 \cdot P_3(x)) \Big|_{x=0} = 0$
- $\circ$  x = 0 is a regular singular point

Check to see if  $x_0 = 0$  is a regular singular point

Check to see if  $x_0 = 0$  is a regular singular point

 $x_0 = 0$ 

• Multiply by denominators

2y''x + 4y + y' = 0

• Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^{k+r}$$

- $\Box$  Rewrite ODE with series expansions
  - Convert y' to series expansion

$$y' = \sum_{k=0}^{\infty} a_k(k+r) x^{k+r-1}$$

• Shift index using k - >k + 1

$$y' = \sum_{k=-1}^{\infty} a_{k+1}(k+1+r) x^{k+r}$$

 $\circ \quad \text{Convert} \; x \cdot y'' \text{ to series expansion}$ 

$$x \cdot y'' = \sum_{k=0}^{\infty} a_k(k+r) (k+r-1) x^{k+r-1}$$

• Shift index using k - > k + 1

$$x \cdot y'' = \sum_{k=-1}^{\infty} a_{k+1}(k+1+r)(k+r)x^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r(-1+2r) x^{-1+r} + \left(\sum_{k=0}^{\infty} \left(a_{k+1}(k+1+r) \left(2k+1+2r\right) + 4a_k\right) x^{k+r}\right) = 0$$

- $a_0$  cannot be 0 by assumption, giving the indicial equation r(-1+2r) = 0
- Values of r that satisfy the indicial equation  $r \in \left\{0, \frac{1}{2}\right\}$
- Each term in the series must be 0, giving the recursion relation  $2(k + \frac{1}{2} + r)(k + 1 + r)a_{k+1} + 4a_k = 0$
- Recursion relation that defines series solution to ODE

$$a_{k+1} = -\frac{4a_k}{(2k+1+2r)(k+1+r)}$$

• Recursion relation for r = 0

$$a_{k+1} = -\frac{4a_k}{(2k+1)(k+1)}$$

• Solution for r = 0

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+1} = -rac{4a_k}{(2k+1)(k+1)}
ight]$$

• Recursion relation for  $r = \frac{1}{2}$ 

$$a_{k+1} = -rac{4a_k}{(2k+2)(k+rac{3}{2})}$$

• Solution for  $r = \frac{1}{2}$ 

$$\left[y = \sum_{k=0}^{\infty} a_k x^{k+rac{1}{2}}, a_{k+1} = -rac{4a_k}{(2k+2)(k+rac{3}{2})}
ight]$$

• Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k x^k\right) + \left(\sum_{k=0}^{\infty} b_k x^{k+\frac{1}{2}}\right), a_{k+1} = -\frac{4a_k}{(2k+1)(k+1)}, b_{k+1} = -\frac{4b_k}{(2k+2)(k+\frac{3}{2})}\right]$$

#### Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
<- linear_1 successful`</pre>
```

Solution by Maple Time used: 0.0 (sec). Leaf size: 27

dsolve(x\*diff(y(x),x\$2)+1/2\*diff(y(x),x)+2\*y(x)=0,y(x), singsol=all)

$$y(x) = c_1 \sin\left(2\sqrt{x}\sqrt{2}\right) + c_2 \cos\left(2\sqrt{x}\sqrt{2}\right)$$

Solution by Mathematica Time used: 0.024 (sec). Leaf size: 38

DSolve[x\*y''[x]+1/2\*y'[x]+2\*y[x]==0,y[x],x,IncludeSingularSolutions -> True]

$$y(x) \rightarrow c_1 \cos\left(2\sqrt{2}\sqrt{x}\right) + c_2 \sin\left(2\sqrt{2}\sqrt{x}\right)$$

# 3.5 problem 10.4.8 (e)

3.5.	1 Solving as second order euler ode ode
3.5.	2 Solving as second order change of variable on x method 2 ode $\therefore$ 292
3.5.	3 Solving as second order change of variable on x method 1 ode $.294$
3.5.	4 Solving as second order change of variable on y method 2 ode . 296
3.5.	5 Solving as second order ode non constant coeff transformation
	on B ode
3.5.	6 Solving using Kovacic algorithm
3.5.	7 Maple step by step solution
Internal probl	em ID [5068]
Internal file na	$\mathrm{Ime}\left[ \texttt{OUTPUT/4561} \texttt{Sunday} \texttt{June} \texttt{05} \texttt{2022} \texttt{03} \texttt{00} \texttt{55} \texttt{PM} \texttt{87042529} \texttt{index.tex} \right]$
Book: Basic	Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Ch	apter 10, Differential equations. Section 10.4, ODEs with variable Coefficients.
Second order	and Homogeneous. page 318
Problem nu	<b>imber</b> : 10.4.8 (e).
~~~ .	

ODE order: 2. ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second\_order\_euler\_ode", "second\_order\_change\_of\_variable\_on\_x\_method\_1", "second\_order\_change\_of\_variable\_on\_x\_method\_2", "second\_order\_change\_of\_variable\_on\_y\_method\_2", "second\_order\_ode\_non\_constant\_coeff\_transformation\_on\_B"

Maple gives the following as the ode type

[[\_Emden, \_Fowler]]

$$x^2y'' - xy' + y = 0$$

#### 3.5.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be  $y = x^r$ , then  $y' = rx^{r-1}$  and  $y'' = r(r-1)x^{r-2}$ . Substituting these back into the given ODE gives

$$x^{2}(r(r-1))x^{r-2} - xrx^{r-1} + x^{r} = 0$$

Simplifying gives

$$r(r-1)x^{r} - rx^{r} + x^{r} = 0$$

Since  $x^r \neq 0$  then dividing throughout by  $x^r$  gives

$$r(r-1) - r + 1 = 0$$

Or

$$r^2 - 2r + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = 1$$
  
 $r_2 = 1$ 

Since the roots are equal, then the general solution is

$$y = c_1 y_1 + c_2 y_2$$

Where  $y_1 = x^r$  and  $y_2 = x^r \ln(x)$ . Hence

$$y = c_1 x + c_2 x \ln\left(x\right)$$

Summary

The solution(s) found are the following

$$y = c_1 x + c_2 x \ln\left(x\right) \tag{1}$$

Verification of solutions

 $y = c_1 x + c_2 x \ln\left(x\right)$ 

Verified OK.

# 3.5.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$x^2y'' - xy' + y = 0 (1)$$

Becomes

$$y'' + p(x) y' + q(x) y = 0$$
(2)

Where

$$p(x) = -\frac{1}{x}$$
$$q(x) = \frac{1}{x^2}$$

Applying change of variables  $\tau = g(x)$  to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0$$
(3)

Where  $\tau$  is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(x) + p(x)\,\tau'(x)}{\tau'(x)^2} \tag{4}$$

$$q_1(\tau) = \frac{q(x)}{\tau'(x)^2} \tag{5}$$

Let  $p_1 = 0$ . Eq (4) simplifies to

$$\tau''(x) + p(x)\,\tau'(x) = 0$$

This ode is solved resulting in

$$\tau = \int e^{-(\int p(x)dx)} dx$$
  
=  $\int e^{-(\int -\frac{1}{x}dx)} dx$   
=  $\int e^{\ln(x)} dx$   
=  $\int xdx$   
=  $\frac{x^2}{2}$  (6)

Using (6) to evaluate  $q_1$  from (5) gives

$$q_{1}(\tau) = \frac{q(x)}{\tau'(x)^{2}}$$
$$= \frac{\frac{1}{x^{2}}}{x^{2}}$$
$$= \frac{1}{x^{4}}$$
(7)

Substituting the above in (3) and noting that now  $p_1 = 0$  results in

$$egin{aligned} &rac{d^2}{d au^2}y( au) + q_1y( au) = 0 \ &rac{d^2}{d au^2}y( au) + rac{y( au)}{x^4} = 0 \end{aligned}$$

But in terms of  $\tau$ 

$$\frac{1}{x^4} = \frac{1}{4\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{y(\tau)}{4\tau^2} = 0$$

The above ode is now solved for  $y(\tau)$ . The ode can be written as

$$4\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be  $y(\tau) = \tau^r$ , then  $y' = r\tau^{r-1}$  and  $y'' = r(r-1)\tau^{r-2}$ . Substituting these back into the given ODE gives

$$4\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + \tau^r = 0$$

Simplifying gives

$$4r(r-1)\,\tau^r + 0\,\tau^r + \tau^r = 0$$

Since  $\tau^r \neq 0$  then dividing throughout by  $\tau^r$  gives

$$4r(r-1) + 0 + 1 = 0$$

Or

$$4r^2 - 4r + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2}$$
$$r_2 = \frac{1}{2}$$

Since the roots are equal, then the general solution is

$$y(\tau) = c_1 y_1 + c_2 y_2$$

Where  $y_1 = \tau^r$  and  $y_2 = \tau^r \ln(\tau)$ . Hence

$$y(\tau) = c_1 \sqrt{\tau} + c_2 \sqrt{\tau} \ln(\tau)$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{x\sqrt{2}(c_1 + 2c_2\ln(x) - c_2\ln(2))}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{x\sqrt{2}\left(c_1 + 2c_2\ln(x) - c_2\ln(2)\right)}{2} \tag{1}$$

Verification of solutions

$$y = \frac{x\sqrt{2}(c_1 + 2c_2\ln(x) - c_2\ln(2))}{2}$$

Verified OK.

#### 3.5.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$x^2y'' - xy' + y = 0 (1)$$

Becomes

$$y'' + p(x) y' + q(x) y = 0$$
(2)

Where

$$p(x) = -\frac{1}{x}$$
$$q(x) = \frac{1}{x^2}$$

Applying change of variables  $\tau = g(x)$  to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0$$
(3)

Where  $\tau$  is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(x) + p(x)\,\tau'(x)}{\tau'(x)^2} \tag{4}$$

$$q_1(\tau) = \frac{q(x)}{\tau'(x)^2} \tag{5}$$

Let  $q_1 = c^2$  where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c}\sqrt{q}$$

$$= \frac{\sqrt{\frac{1}{x^2}}}{c}$$

$$\tau'' = -\frac{1}{c\sqrt{\frac{1}{x^2}}x^3}$$
(6)

Substituting the above into (4) results in

$$p_{1}(\tau) = \frac{\tau''(x) + p(x)\tau'(x)}{\tau'(x)^{2}}$$
$$= \frac{-\frac{1}{c\sqrt{\frac{1}{x^{2}}x^{3}}} - \frac{1}{x}\frac{\sqrt{\frac{1}{x^{2}}}}{c}}{\left(\frac{\sqrt{\frac{1}{x^{2}}}}{c}\right)^{2}}$$
$$= -2c$$

Therefore ode (3) now becomes

$$y(\tau)'' + p_1 y(\tau)' + q_1 y(\tau) = 0$$
  
$$\frac{d^2}{d\tau^2} y(\tau) - 2c \left(\frac{d}{d\tau} y(\tau)\right) + c^2 y(\tau) = 0$$
(7)

The above ode is now solved for  $y(\tau)$ . Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = \mathrm{e}^{c\tau} c_1$$

Now from (6)

$$\tau = \int \frac{1}{c} \sqrt{q} \, dx$$
$$= \frac{\int \sqrt{\frac{1}{x^2}} dx}{c}$$
$$= \frac{\sqrt{\frac{1}{x^2}} x \ln(x)}{c}$$

Substituting the above into the solution obtained gives

$$y = c_1 x$$

 $\frac{\text{Summary}}{\text{The solution(s) found are the following}}$ 

$$y = c_1 x \tag{1}$$

Verification of solutions

 $y = c_1 x$ 

Verified OK.

#### 3.5.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$x^2y'' - xy' + y = 0 (1)$$

Becomes

$$y'' + p(x) y' + q(x) y = 0$$
(2)

Where

$$p(x) = -\frac{1}{x}$$
$$q(x) = \frac{1}{x^2}$$

Applying change of variables on the dependent variable  $y = v(x) x^n$  to (2) gives the following ode where the dependent variables is v(x) and not y.

$$v''(x) + \left(\frac{2n}{x} + p\right)v'(x) + \left(\frac{n(n-1)}{x^2} + \frac{np}{x} + q\right)v(x) = 0$$
(3)

Let the coefficient of v(x) above be zero. Hence

$$\frac{n(n-1)}{x^2} + \frac{np}{x} + q = 0 \tag{4}$$

Substituting the earlier values found for p(x) and q(x) into (4) gives

$$\frac{n(n-1)}{x^2} - \frac{n}{x^2} + \frac{1}{x^2} = 0$$
(5)

Solving (5) for n gives

$$n = 1 \tag{6}$$

Substituting this value in (3) gives

$$v''(x) + \frac{v'(x)}{x} = 0$$
  
$$v''(x) + \frac{v'(x)}{x} = 0$$
 (7)

Using the substitution

$$u(x) = v'(x)$$

Then (7) becomes

$$u'(x) + \frac{u(x)}{x} = 0$$
 (8)

The above is now solved for u(x). In canonical form the ODE is

$$u' = F(x, u)$$
  
=  $f(x)g(u)$   
=  $-\frac{u}{x}$ 

Where  $f(x) = -\frac{1}{x}$  and g(u) = u. Integrating both sides gives

$$\frac{1}{u} du = -\frac{1}{x} dx$$
$$\int \frac{1}{u} du = \int -\frac{1}{x} dx$$
$$\ln(u) = -\ln(x) + c_1$$
$$u = e^{-\ln(x) + c_1}$$
$$= \frac{c_1}{x}$$

Now that u(x) is known, then

$$egin{aligned} v'(x) &= u(x) \ v(x) &= \int u(x) \, dx + c_2 \ &= c_1 \ln{(x)} + c_2 \end{aligned}$$

Hence

$$y = v(x) x^{n}$$
  
= (c<sub>1</sub> ln (x) + c<sub>2</sub>) x  
= (c<sub>1</sub> ln (x) + c<sub>2</sub>) x

 $\frac{\text{Summary}}{\text{The solution(s) found are the following}}$ 

$$y = (c_1 \ln (x) + c_2) x$$
(1)

Verification of solutions

$$y = \left(c_1 \ln \left(x\right) + c_2\right) x$$

Verified OK.

# 3.5.5 Solving as second order ode non constant coeff transformation on B ode

Given an ode of the form

$$Ay'' + By' + Cy = F(x)$$

This method reduces the order ode the ODE by one by applying the transformation

$$y = Bv$$

This results in

$$y' = B'v + v'B$$
  

$$y'' = B''v + B'v' + v''B + v'B'$$
  

$$= v''B + 2v' + B' + B''v$$

And now the original ode becomes

$$A(v''B + 2v'B' + B''v) + B(B'v + v'B) + CBv = 0$$
  

$$ABv'' + (2AB' + B^2)v' + (AB'' + BB' + CB)v = 0$$
(1)

If the term AB'' + BB' + CB is zero, then this method works and can be used to solve

$$ABv'' + (2AB' + B^2) v' = 0$$

By Using u = v' which reduces the order of the above ode to one. The new ode is

$$ABu' + \left(2AB' + B^2\right)u = 0$$

The above ode is first order ode which is solved for u. Now a new ode v' = u is solved for v as first order ode. Then the final solution is obtain from y = Bv.

This method works only if the term AB'' + BB' + CB is zero. The given ODE shows that

$$A = x^{2}$$
$$B = -x$$
$$C = 1$$
$$F = 0$$

The above shows that for this ode

$$AB'' + BB' + CB = (x^2) (0) + (-x) (-1) + (1) (-x)$$
  
= 0

Hence the ode in v given in (1) now simplifies to

$$-x^{3}v'' + (-x^{2})v' = 0$$

Now by applying v' = u the above becomes

$$-x^{2}(u'(x) x + u(x)) = 0$$

Which is now solved for u. In canonical form the ODE is

$$u' = F(x, u)$$
  
=  $f(x)g(u)$   
=  $-\frac{u}{x}$ 

Where  $f(x) = -\frac{1}{x}$  and g(u) = u. Integrating both sides gives

$$\frac{1}{u} du = -\frac{1}{x} dx$$
$$\int \frac{1}{u} du = \int -\frac{1}{x} dx$$
$$\ln(u) = -\ln(x) + c_1$$
$$u = e^{-\ln(x) + c_1}$$
$$= \frac{c_1}{x}$$

The ode for v now becomes

$$v' = u$$
$$= \frac{c_1}{x}$$

Which is now solved for v. Integrating both sides gives

$$egin{aligned} v(x) &= \int rac{c_1}{x} \, \mathrm{d}x \ &= c_1 \ln{(x)} + c_2 \end{aligned}$$

Therefore the solution is

$$y(x) = Bv$$
  
= (-x) (c<sub>1</sub> ln (x) + c<sub>2</sub>)  
= -(c<sub>1</sub> ln (x) + c<sub>2</sub>) x

Summary

The solution(s) found are the following

$$y = -(c_1 \ln(x) + c_2) x \tag{1}$$

Verification of solutions

$$y = -(c_1 \ln \left(x\right) + c_2) x$$

Verified OK.

#### 3.5.6 Solving using Kovacic algorithm

Writing the ode as

$$x^2y'' - xy' + y = 0 (1)$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = x^{2}$$
  

$$B = -x$$
  

$$C = 1$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4x^2} \tag{6}$$

Comparing the above to (5) shows that

$$s = -1$$
$$t = 4x^2$$

Therefore eq. (4) becomes

$$z''(x) = \left(-\frac{1}{4x^2}\right)z(x) \tag{7}$$

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 52: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 2 - 0$$
$$= 2$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of  $t = 4x^2$ . There is a pole at x = 0 of order 2. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case n = 1.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{4x^2}$$

For the pole at x = 0 let b be the coefficient of  $\frac{1}{x^2}$  in the partial fractions decomposition of r given above. Therefore  $b = -\frac{1}{4}$ . Hence

$$\begin{split} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1+4b} = \frac{1}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1+4b} = \frac{1}{2} \end{split}$$

Since the order of r at  $\infty$  is 2 then  $[\sqrt{r}]_{\infty} = 0$ . Let b be the coefficient of  $\frac{1}{x^2}$  in the Laurent series expansion of r at  $\infty$ , which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{1}{4x^2}$$

Since the gcd(s,t) = 1. This gives  $b = -\frac{1}{4}$ . Hence

$$\begin{split} [\sqrt{r}]_{\infty} &= 0\\ \alpha_{\infty}^{+} &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2}\\ \alpha_{\infty}^{-} &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{split}$$

The following table summarizes the findings so far for poles and for the order of r at  $\infty$  where r is 1

		r =	$\overline{4x^2}$				
pole $c$ location		pole order		·]	$\sqrt{r}]_c$	$\alpha_c^+$	$\alpha_c^-$
0		4	2 0		0	$\frac{1}{2}$	$\frac{1}{2}$
	Order of $r$ at $\infty$		$[\sqrt{r}]_{\circ}$	×	$lpha_\infty^+$	$lpha_{\infty}^{-}$	
	2		0		$\frac{1}{2}$	$\frac{1}{2}$	]

Now that the all  $[\sqrt{r}]_c$  and its associated  $\alpha_c^{\pm}$  have been determined for all the poles in the set  $\Gamma$  and  $[\sqrt{r}]_{\infty}$  and its associated  $\alpha_{\infty}^{\pm}$  have also been found, the next step is to determine possible non negative integer d from these using

$$d=\alpha_\infty^{s(\infty)}-\sum_{c\in\Gamma}\alpha_c^{s(c)}$$

Where s(c) is either + or - and  $s(\infty)$  is the sign of  $\alpha_{\infty}^{\pm}$ . This is done by trial over all set of families  $s = (s(c))_{c \in \Gamma \cup \infty}$  until such d is found to work in finding candidate  $\omega$ . Trying  $\alpha_{\infty}^{-} = \frac{1}{2}$  then

$$d = \alpha_{\infty}^{-} - (\alpha_{c_1}^{-})$$
$$= \frac{1}{2} - \left(\frac{1}{2}\right)$$
$$= 0$$

Since d an integer and  $d \ge 0$  then it can be used to find  $\omega$  using

$$\omega = \sum_{c \in \Gamma} \left( s(c) [\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{x - c} \right) + s(\infty) [\sqrt{r}]_{\infty}$$

The above gives

$$\begin{split} \omega &= \left( (-) [\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}}{x - c_1} \right) + (-) [\sqrt{r}]_{\infty} \\ &= \frac{1}{2x} + (-) (0) \\ &= \frac{1}{2x} \\ &= \frac{1}{2x} \end{split}$$

Now that  $\omega$  is determined, the next step is find a corresponding minimal polynomial p(x) of degree d = 0 to solve the ode. The polynomial p(x) needs to satisfy the equation

$$p'' + 2\omega p' + \left(\omega' + \omega^2 - r\right)p = 0 \tag{1A}$$

Let

$$p(x) = 1 \tag{2A}$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(\frac{1}{2x}\right)(0) + \left(\left(-\frac{1}{2x^2}\right) + \left(\frac{1}{2x}\right)^2 - \left(-\frac{1}{4x^2}\right)\right) = 0$$
$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode z'' = rz is

$$z_1(x) = p e^{\int \omega \, dx}$$
$$= e^{\int \frac{1}{2x} \, dx}$$
$$= \sqrt{x}$$

The first solution to the original ode in y is found from

$$y_{1} = z_{1}e^{\int -\frac{1}{2}\frac{B}{A} dx}$$
  
=  $z_{1}e^{-\int \frac{1}{2}\frac{-x}{x^{2}} dx}$   
=  $z_{1}e^{\frac{\ln(x)}{2}}$   
=  $z_{1}(\sqrt{x})$ 

Which simplifies to

$$y_1 = x$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} \, dx}}{y_1^2} \, dx$$

Substituting gives

$$egin{aligned} y_2 &= y_1 \int rac{e^{\int -rac{-x}{x^2}\,dx}}{\left(y_1
ight)^2}\,dx \ &= y_1 \int rac{e^{\ln(x)}}{\left(y_1
ight)^2}\,dx \ &= y_1(\ln{(x)}) \end{aligned}$$

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
  
=  $c_1(x) + c_2(x(\ln(x)))$ 

 $\frac{Summary}{The solution(s) found are the following}$ 

$$y = c_1 x + c_2 x \ln\left(x\right) \tag{1}$$

Verification of solutions

$$y = c_1 x + c_2 x \ln\left(x\right)$$

Verified OK.

### 3.5.7 Maple step by step solution

Let's solve

 $x^2y^{\prime\prime}-xy^\prime+y=0$ 

- Highest derivative means the order of the ODE is 2 y''
- Isolate 2nd derivative

$$y'' = \frac{y'}{x} - \frac{y}{x^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear  $y'' \frac{y'}{x} + \frac{y}{x^2} = 0$
- Multiply by denominators of the ODE

$$x^2y'' - xy' + y = 0$$

• Make a change of variables

 $t = \ln\left(x\right)$ 

 $\Box$  Substitute the change of variables back into the ODE

- $\circ~$  Calculate the 1st derivative of y with respect to x , using the chain rule  $y' = \left( \tfrac{d}{dt} y(t) \right) t'(x)$
- Compute derivative

$$y' = \frac{\frac{d}{dt}y(t)}{x}$$

- Calculate the 2nd derivative of **y** with respect to **x** , using the chain rule  $y'' = \left(\frac{d^2}{dt^2}y(t)\right)t'(x)^2 + t''(x)\left(\frac{d}{dt}y(t)\right)$
- Compute derivative

$$y'' = \frac{\frac{d^2}{dt^2}y(t)}{x^2} - \frac{\frac{d}{dt}y(t)}{x^2}$$

Substitute the change of variables back into the ODE

$$x^2 \left( \frac{\frac{d^2}{dt^2} y(t)}{x^2} - \frac{\frac{d}{dt} y(t)}{x^2} \right) - \frac{d}{dt} y(t) + y(t) = 0$$

• Simplify

$$\frac{d^2}{dt^2}y(t) - 2\frac{d}{dt}y(t) + y(t) = 0$$

• Characteristic polynomial of ODE

$$r^2 - 2r + 1 = 0$$

• Factor the characteristic polynomial

$$\left(r-1\right)^2 = 0$$

• Root of the characteristic polynomial

r = 1

• 1st solution of the ODE

$$y_1(t) = e^t$$

- Repeated root, multiply  $y_1(t)$  by t to ensure linear independence  $y_2(t) = t e^t$
- General solution of the ODE

$$y(t) = c_1 y_1(t) + c_2 y_2(t)$$

• Substitute in solutions

$$y(t) = c_1 \mathrm{e}^t + c_2 t \, \mathrm{e}^t$$

• Change variables back using  $t = \ln(x)$ 

 $y = c_1 x + c_2 x \ln\left(x\right)$ 

• Simplify

$$y = x(c_2 \ln (x) + c_1)$$

#### Maple trace

`Methods for second order ODEs: --- Trying classification methods --trying a quadrature checking if the LODE has constant coefficients checking if the LODE is of Euler type <- LODE of Euler type successful`</pre>

Solution by Maple Time used: 0.0 (sec). Leaf size: 12

 $dsolve(x^2*diff(y(x),x$2)-x*diff(y(x),x)+y(x)=0,y(x), singsol=all)$ 

$$y(x) = x(c_2 \ln (x) + c_1)$$

Solution by Mathematica Time used: 0.018 (sec). Leaf size: 15

DSolve[x^2\*y''[x]-x\*y'[x]+y[x]==0,y[x],x,IncludeSingularSolutions -> True]

 $y(x) \rightarrow x(c_2 \log(x) + c_1)$ 

# 3.6 problem 10.4.8 (f)

3.6.1	Solving as second order bessel ode ode	308		
3.6.2	Solving using Kovacic algorithm	309		
3.6.3	Maple step by step solution $\ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots$	314		
Internal problem ID [5069]				

Internal file name [OUTPUT/4562\_Sunday\_June\_05\_2022\_03\_00\_56\_PM\_86182655/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients. Second order and Homogeneous. page 318
Problem number: 10.4.8 (f).
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second\_order\_bessel\_ode" Maple gives the following as the ode type

[[\_Emden, \_Fowler]]

$$2xy'' - y' + 2y = 0$$

#### 3.6.1 Solving as second order bessel ode ode

Writing the ode as

$$x^2y'' - \frac{xy'}{2} + xy = 0 \tag{1}$$

Bessel ode has the form

$$x^{2}y'' + xy' + \left(-n^{2} + x^{2}\right)y = 0$$
<sup>(2)</sup>

The generalized form of Bessel ode is given by Bowman (1958) as the following

$$x^{2}y'' + (1 - 2\alpha)xy' + (\beta^{2}\gamma^{2}x^{2\gamma} - n^{2}\gamma^{2} + \alpha^{2})y = 0$$
(3)

With the standard solution

$$y = x^{\alpha}(c_1 \operatorname{BesselJ}(n, \beta x^{\gamma}) + c_2 \operatorname{BesselY}(n, \beta x^{\gamma}))$$
(4)

Comparing (3) to (1) and solving for  $\alpha, \beta, n, \gamma$  gives

$$\begin{aligned} \alpha &= \frac{3}{4} \\ \beta &= 2 \\ n &= \frac{3}{2} \\ \gamma &= \frac{1}{2} \end{aligned}$$

Substituting all the above into (4) gives the solution as

$$y = -\frac{c_1(2\cos(2\sqrt{x})\sqrt{x} - \sin(2\sqrt{x}))}{2\sqrt{\pi}} - \frac{c_2(2\sin(2\sqrt{x})\sqrt{x} + \cos(2\sqrt{x}))}{2\sqrt{\pi}}$$

Summary

 $\overline{\text{The solution}}(s)$  found are the following

$$y = -\frac{c_1(2\cos(2\sqrt{x})\sqrt{x} - \sin(2\sqrt{x}))}{2\sqrt{\pi}} - \frac{c_2(2\sin(2\sqrt{x})\sqrt{x} + \cos(2\sqrt{x}))}{2\sqrt{\pi}}$$
(1)

Verification of solutions

$$y = -\frac{c_1 \left(2\cos\left(2\sqrt{x}\right)\sqrt{x} - \sin\left(2\sqrt{x}\right)\right)}{2\sqrt{\pi}} - \frac{c_2 \left(2\sin\left(2\sqrt{x}\right)\sqrt{x} + \cos\left(2\sqrt{x}\right)\right)}{2\sqrt{\pi}}$$

Verified OK.

### 3.6.2 Solving using Kovacic algorithm

Writing the ode as

$$2xy'' - y' + 2y = 0 (1)$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 2x$$
  

$$B = -1$$

$$C = 2$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{5 - 16x}{16x^2} \tag{6}$$

Comparing the above to (5) shows that

$$s = 5 - 16x$$
$$t = 16x^2$$

Therefore eq. (4) becomes

$$z''(x) = \left(\frac{5-16x}{16x^2}\right)z(x) \tag{7}$$

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 54: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 2 - 1$$
$$= 1$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of  $t = 16x^2$ . There is a pole at x = 0 of order 2. Since there is a pole of order 2 then necessary conditions for case two are met. Therefore

$$L = [2]$$

Attempting to find a solution using case n = 2.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{x} + \frac{5}{16x^2}$$

For the pole at x = 0 let b be the coefficient of  $\frac{1}{x^2}$  in the partial fractions decomposition of r given above. Therefore  $b = \frac{5}{16}$ . Hence

$$E_c = \{2, 2 + 2\sqrt{1+4b}, 2 - 2\sqrt{1+4b}\}\$$
  
= \{-1, 2, 5\}

Since the order of r at  $\infty$  is 1 < 2 then

$$E_{\infty} = \{1\}$$

The following table summarizes the findings so far for poles and for the order of r at  $\infty$  for case 2 of Kovacic algorithm.

pole $c$ location		pole order			$E_c$
0		2		{-	$-1, 2, 5\}$
	Order o	of $r$ at $\infty$	E	$\overline{f}_{\infty}$	
		1	{	1}	

Using the family  $\{e_1, e_2, \ldots, e_\infty\}$  given by

$$e_1 = -1, \ e_\infty = 1$$

Gives a non negative integer d (the degree of the polynomial p(x)), which is generated using

$$d = \frac{1}{2} \left( e_{\infty} - \sum_{c \in \Gamma} e_c \right)$$
$$= \frac{1}{2} (1 - (-1))$$
$$= 1$$

We now form the following rational function

$$\begin{split} \theta &= \frac{1}{2} \sum_{c \in \Gamma} \frac{e_c}{x - c} \\ &= \frac{1}{2} \left( \frac{-1}{(x - (0))} \right) \\ &= -\frac{1}{2x} \end{split}$$

Now we search for a monic polynomial p(x) of degree d = 1 such that

$$p''' + 3\theta p'' + (3\theta^2 + 3\theta' - 4r) p' + (\theta'' + 3\theta\theta' + \theta^3 - 4r\theta - 2r') p = 0$$
(1A)

Since d = 1, then letting

$$p = x + a_0 \tag{2A}$$

Substituting p and  $\theta$  into Eq. (1A) gives

$$\frac{1-4a_0}{x^2} = 0$$

And solving for p gives

$$p = x + \frac{1}{4}$$

Now that p(x) is found let

$$\phi = \theta + \frac{p'}{p}$$
$$= \frac{1}{x + \frac{1}{4}} - \frac{1}{2x}$$

Let  $\omega$  be the solution of

$$\omega^2 - \phi\omega + \left(\frac{1}{2}\phi' + \frac{1}{2}\phi^2 - r\right) = 0$$

Substituting the values for  $\phi$  and r into the above equation gives

$$w^{2} - \left(\frac{1}{x + \frac{1}{4}} - \frac{1}{2x}\right)w + \frac{64x^{2} - 12x + 1}{64x^{3} + 16x^{2}} = 0$$

Solving for  $\omega$  gives

$$\omega = \frac{16x\sqrt{-x} + 4x - 1}{4\left(4x + 1\right)x}$$

Therefore the first solution to the ode z'' = rz is

$$z_1(x) = e^{\int \omega \, dx}$$
  
=  $e^{\int \frac{16x\sqrt{-x}+4x-1}{4(4x+1)x} dx}$   
=  $\frac{\sqrt{2\sqrt{-x}-1}\sqrt{4x+1}e^{2\sqrt{-x}}}{\sqrt{2\sqrt{-x}+1}x^{\frac{1}{4}}}$ 

The first solution to the original ode in y is found from

$$egin{aligned} y_1 &= z_1 e^{\int -rac{1}{2} rac{B}{A} \, dx} \ &= z_1 e^{-\int rac{1}{2} rac{-1}{2x} \, dx} \ &= z_1 e^{rac{\ln(x)}{4}} \ &= z_1 \Big( x^rac{1}{4} \Big) \end{aligned}$$

Which simplifies to

$$y_1 = \frac{e^{2\sqrt{-x}}\sqrt{(2\sqrt{-x}-1)(4x+1)}}{\sqrt{2\sqrt{-x}+1}}$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} \, dx}}{y_1^2} \, dx$$

Substituting gives

$$y_{2} = y_{1} \int \frac{e^{\int -\frac{-1}{2x} dx}}{(y_{1})^{2}} dx$$
  
=  $y_{1} \int \frac{e^{\frac{\ln(x)}{2}}}{(y_{1})^{2}} dx$   
=  $y_{1} \left( \int \frac{\sqrt{x} e^{-4\sqrt{-x}} (2\sqrt{-x} + 1)}{(2\sqrt{-x} - 1) (4x + 1)} dx \right)$ 

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
  
=  $c_1 \left( \frac{e^{2\sqrt{-x}} \sqrt{(2\sqrt{-x} - 1)(4x + 1)}}{\sqrt{2\sqrt{-x} + 1}} \right)$   
+  $c_2 \left( \frac{e^{2\sqrt{-x}} \sqrt{(2\sqrt{-x} - 1)(4x + 1)}}{\sqrt{2\sqrt{-x} + 1}} \left( \int \frac{\sqrt{x} e^{-4\sqrt{-x}}(2\sqrt{-x} + 1)}{(2\sqrt{-x} - 1)(4x + 1)} dx \right) \right)$ 

Summary

The solution(s) found are the following

$$y = \frac{c_1 e^{2\sqrt{-x}} \sqrt{\left(2\sqrt{-x} - 1\right)\left(4x + 1\right)}}{\sqrt{2\sqrt{-x} + 1}} + \frac{c_2 e^{2\sqrt{-x}} \sqrt{\left(2\sqrt{-x} - 1\right)\left(4x + 1\right)} \left(\int \frac{\sqrt{x} e^{-4\sqrt{-x}}(2\sqrt{-x} + 1)}{(2\sqrt{-x} - 1)(4x + 1)} dx\right)}{\sqrt{2\sqrt{-x} + 1}}$$
(1)

Verification of solutions

$$y = \frac{c_1 e^{2\sqrt{-x}} \sqrt{\left(2\sqrt{-x} - 1\right)\left(4x + 1\right)}}{\sqrt{2\sqrt{-x} + 1}} + \frac{c_2 e^{2\sqrt{-x}} \sqrt{\left(2\sqrt{-x} - 1\right)\left(4x + 1\right)} \left(\int \frac{\sqrt{x} e^{-4\sqrt{-x}}(2\sqrt{-x} + 1)}{(2\sqrt{-x} - 1)(4x + 1)} dx\right)}{\sqrt{2\sqrt{-x} + 1}}$$

Verified OK.

# 3.6.3 Maple step by step solution

Let's solve

2y''x - y' + 2y = 0

• Highest derivative means the order of the ODE is 2

• Isolate 2nd derivative

$$y'' = \frac{y'}{2x} - \frac{y}{x}$$

• Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{y'}{2x} + \frac{y}{x} = 0$$

 $\Box$  Check to see if  $x_0 = 0$  is a regular singular point

• Define functions

$$\left[P_2(x) = -\frac{1}{2x}, P_3(x) = \frac{1}{x}\right]$$

 $\circ \quad x \cdot P_2(x)$  is analytic at x = 0

$$(x \cdot P_2(x))\Big|_{x=0} = -\frac{1}{2}$$

 $\circ \quad x^2 \cdot P_3(x)$  is analytic at x = 0

$$\left(x^2 \cdot P_3(x)\right)\Big|_{x=0} = 0$$

 $\circ \quad x = 0 \text{is a regular singular point}$ 

Check to see if  $x_0 = 0$  is a regular singular point

 $x_0 = 0$ 

• Multiply by denominators

$$2y''x - y' + 2y = 0$$

• Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^{k+r}$$

- $\Box$  Rewrite ODE with series expansions
  - Convert y' to series expansion

$$y' = \sum\limits_{k=0}^{\infty} a_k(k+r) \, x^{k+r-1}$$

• Shift index using k - >k + 1

$$y' = \sum_{k=-1}^{\infty} a_{k+1}(k+1+r) x^{k+r}$$

 $\circ \quad \text{Convert} \; x \cdot y'' \text{ to series expansion}$ 

$$x \cdot y'' = \sum_{k=0}^{\infty} a_k(k+r) (k+r-1) x^{k+r-1}$$

• Shift index using k - >k + 1

$$x \cdot y'' = \sum_{k=-1}^{\infty} a_{k+1}(k+1+r) (k+r) x^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r(-3+2r) x^{-1+r} + \left(\sum_{k=0}^{\infty} \left(a_{k+1}(k+1+r) \left(2k-1+2r\right)+2a_k\right) x^{k+r}\right) = 0$$

- $a_0$  cannot be 0 by assumption, giving the indicial equation r(-3+2r) = 0
- Values of r that satisfy the indicial equation  $r \in \left\{0, \frac{3}{2}\right\}$
- Each term in the series must be 0, giving the recursion relation  $2(k - \frac{1}{2} + r)(k + 1 + r)a_{k+1} + 2a_k = 0$
- Recursion relation that defines series solution to ODE

$$a_{k+1} = -\frac{2a_k}{(2k-1+2r)(k+1+r)}$$

• Recursion relation for r = 0

$$a_{k+1} = -\frac{2a_k}{(2k-1)(k+1)}$$

• Solution for r = 0

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+1} = -rac{2a_k}{(2k-1)(k+1)}
ight]$$

• Recursion relation for  $r = \frac{3}{2}$ 

$$a_{k+1} = -\frac{2a_k}{(2k+2)(k+\frac{5}{2})}$$

• Solution for  $r = \frac{3}{2}$ 

$$\left[y = \sum_{k=0}^{\infty} a_k x^{k+rac{3}{2}}, a_{k+1} = -rac{2a_k}{(2k+2)(k+rac{5}{2})}
ight]$$

• Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k x^k\right) + \left(\sum_{k=0}^{\infty} b_k x^{k+\frac{3}{2}}\right), a_{k+1} = -\frac{2a_k}{(2k-1)(k+1)}, b_{k+1} = -\frac{2b_k}{(2k+2)(k+\frac{5}{2})}\right]$$

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
        A Liouvillian solution exists
        Group is reducible or imprimitive
<- Kovacics algorithm successful`</pre>
```

Solution by Maple Time used: 0.032 (sec). Leaf size: 36

dsolve(2\*x\*diff(y(x),x\$2)-diff(y(x),x)+2\*y(x)=0,y(x), singsol=all)

$$y(x) = (2\sqrt{x}c_1 + c_2)\cos(2\sqrt{x}) - \sin(2\sqrt{x})(-2\sqrt{x}c_2 + c_1)$$

Solution by Mathematica Time used: 0.199 (sec). Leaf size: 59

DSolve[2\*x\*y''[x]-y'[x]+2\*y[x]==0,y[x],x,IncludeSingularSolutions -> True]

$$y(x) \rightarrow c_1 e^{2i\sqrt{x}} \left(2\sqrt{x}+i\right) + \frac{1}{8}c_2 e^{-2i\sqrt{x}} \left(1+2i\sqrt{x}\right)$$

# 3.7 problem 10.4.8 (g)

Internal problem ID [5070]

Internal file name [OUTPUT/4563\_Sunday\_June\_05\_2022\_03\_00\_57\_PM\_84765408/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients.
Second order and Homogeneous. page 318
Problem number: 10.4.8 (g).
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic"

Maple gives the following as the ode type

```
[[_2nd_order, _with_linear_symmetries]]
```

$$xy'' + xy' - 2y = 0$$

#### 3.7.1 Solving using Kovacic algorithm

Writing the ode as

$$xy'' + xy' - 2y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = x$$
  

$$B = x$$

$$C = -2$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int rac{B}{2A} \, dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{8+x}{4x} \tag{6}$$

Comparing the above to (5) shows that

$$s = 8 + x$$
$$t = 4x$$

Therefore eq. (4) becomes

$$z''(x) = \left(\frac{8+x}{4x}\right)z(x) \tag{7}$$

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 56: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 1 - 1$$
$$= 0$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of t = 4x. There is a pole at x = 0 of order 1. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Attempting to find a solution using case n = 1.

Looking at poles of order 1. For the pole at x = 0 of order 1 then

$$[\sqrt{r}]_c = 0$$
  
 $\alpha_c^+ = 1$   
 $\alpha_c^- = 1$ 

Since the order of r at  $\infty$  is  $O_r(\infty) = 0$  then

$$v = \frac{-O_r(\infty)}{2} = \frac{0}{2} = 0$$

 $[\sqrt{r}]_{\infty}$  is the sum of terms involving  $x^i$  for  $0 \le i \le v$  in the Laurent series for  $\sqrt{r}$  at  $\infty$ . Therefore

$$[\sqrt{r}]_{\infty} = \sum_{i=0}^{v} a_i x^i$$
$$= \sum_{i=0}^{0} a_i x^i$$
(8)

Let a be the coefficient of  $x^v = x^0$  in the above sum. The Laurent series of  $\sqrt{r}$  at  $\infty$  is

$$\sqrt{r} \approx \frac{1}{2} + \frac{2}{x} - \frac{4}{x^2} + \frac{16}{x^3} - \frac{80}{x^4} + \frac{448}{x^5} - \frac{2688}{x^6} + \frac{16896}{x^7} + \dots$$
(9)

Comparing Eq. (9) with Eq. (8) shows that

$$a = \frac{1}{2}$$

From Eq. (9) the sum up to v = 0 gives

$$[\sqrt{r}]_{\infty} = \sum_{i=0}^{0} a_i x^i$$
$$= \frac{1}{2}$$
(10)

Now we need to find b, where b be the coefficient of  $x^{v-1} = x^{-1} = \frac{1}{x}$  in r minus the coefficient of same term but in  $([\sqrt{r}]_{\infty})^2$  where  $[\sqrt{r}]_{\infty}$  was found above in Eq (10). Hence

$$\left( [\sqrt{r}]_{\infty} \right)^2 = \frac{1}{4}$$

This shows that the coefficient of  $\frac{1}{x}$  in the above is 0. Now we need to find the coefficient of  $\frac{1}{x}$  in r. How this is done depends on if v = 0 or not. Since v = 0 then starting from  $r = \frac{s}{t}$  and doing long division in the form

$$r = Q + \frac{R}{t}$$

Where Q is the quotient and R is the remainder. Then the coefficient of  $\frac{1}{x}$  in r will be the coefficient in R of the term in x of degree of t minus one, divided by the leading coefficient in t. Doing long division gives

$$r = \frac{s}{t}$$
$$= \frac{8+x}{4x}$$
$$= Q + \frac{R}{4x}$$
$$= \left(\frac{1}{4}\right) + \left(\frac{2}{x}\right)$$
$$= \frac{1}{4} + \frac{2}{x}$$

Since the degree of t is 1, then we see that the coefficient of the term 1 in the remainder R is 8. Dividing this by leading coefficient in t which is 4 gives 2. Now b can be found.

$$b = (2) - (0)$$
  
= 2

Hence

$$\begin{split} [\sqrt{r}]_{\infty} &= \frac{1}{2} \\ \alpha_{\infty}^{+} &= \frac{1}{2} \left( \frac{b}{a} - v \right) = \frac{1}{2} \left( \frac{2}{\frac{1}{2}} - 0 \right) = 2 \\ \alpha_{\infty}^{-} &= \frac{1}{2} \left( -\frac{b}{a} - v \right) = \frac{1}{2} \left( -\frac{2}{\frac{1}{2}} - 0 \right) = -2 \end{split}$$

The following table summarizes the findings so far for poles and for the order of r at  $\infty$  where r is

$$r = \frac{8+x}{4x}$$

$r = rac{8+x}{4x}$							
pol	e $c$ location	pole	order	$[\sqrt{r}]_c$		$lpha_c^+$	$lpha_c^-$
0		1		0		0	1
	Order of $r$ at $\infty$		$[\sqrt{r}]_{\circ}$	ø	$lpha^+_\infty$	$\alpha_{\infty}^{-}$	]
	0		$\frac{1}{2}$		2	-2	

Now that the all  $[\sqrt{r}]_c$  and its associated  $\alpha_c^{\pm}$  have been determined for all the poles in the set  $\Gamma$  and  $[\sqrt{r}]_{\infty}$  and its associated  $\alpha_{\infty}^{\pm}$  have also been found, the next step is to determine possible non negative integer d from these using

$$d = lpha_{\infty}^{s(\infty)} - \sum_{c \in \Gamma} lpha_{c}^{s(c)}$$

Where s(c) is either + or - and  $s(\infty)$  is the sign of  $\alpha_{\infty}^{\pm}$ . This is done by trial over all set of families  $s = (s(c))_{c \in \Gamma \cup \infty}$  until such d is found to work in finding candidate  $\omega$ . Trying  $\alpha_{\infty}^+ = 2$  then

$$d = \alpha_{\infty}^{+} - (\alpha_{c_{1}}^{-})$$
$$= 2 - (1)$$
$$= 1$$

Since d an integer and  $d \geq 0$  then it can be used to find  $\omega$  using

$$\omega = \sum_{c \in \Gamma} \left( s(c) [\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{x - c} \right) + s(\infty) [\sqrt{r}]_{\infty}$$

Substituting the above values in the above results in

$$\begin{split} \omega &= \left( (-) [\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}}{x - c_1} \right) + (+) [\sqrt{r}]_{\infty} \\ &= \frac{1}{x} + \left(\frac{1}{2}\right) \\ &= \frac{1}{x} + \frac{1}{2} \\ &= \frac{1}{x} + \frac{1}{2} \end{split}$$

Now that  $\omega$  is determined, the next step is find a corresponding minimal polynomial p(x) of degree d = 1 to solve the ode. The polynomial p(x) needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0$$
(1A)

Let

$$p(x) = x + a_0 \tag{2A}$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(\frac{1}{x} + \frac{1}{2}\right)(1) + \left(\left(-\frac{1}{x^2}\right) + \left(\frac{1}{x} + \frac{1}{2}\right)^2 - \left(\frac{8+x}{4x}\right)\right) = 0$$
$$\frac{2-a_0}{x} = 0$$

Solving for the coefficients  $a_i$  in the above using method of undetermined coefficients gives

$$\{a_0 = 2\}$$

Substituting these coefficients in p(x) in eq. (2A) results in

$$p(x) = x + 2$$

Therefore the first solution to the ode z'' = rz is

$$z_1(x) = p e^{\int \omega \, dx}$$
  
=  $(x+2) e^{\int (\frac{1}{x} + \frac{1}{2}) dx}$   
=  $(x+2) e^{\frac{x}{2} + \ln(x)}$   
=  $(x+2) x e^{\frac{x}{2}}$ 

The first solution to the original ode in y is found from

$$egin{aligned} y_1 &= z_1 e^{\int -rac{1}{2}rac{B}{A}\,dx} \ &= z_1 e^{-\int rac{1}{2}rac{x}{x}\,dx} \ &= z_1 e^{-rac{x}{2}} \ &= z_1 (\mathrm{e}^{-rac{x}{2}}) \end{aligned}$$

Which simplifies to

$$y_1 = x(x+2)$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2=y_1\int rac{e^{\int -rac{B}{A}\,dx}}{y_1^2}\,dx$$

Substituting gives

$$y_{2} = y_{1} \int \frac{e^{\int -\frac{x}{x} dx}}{(y_{1})^{2}} dx$$
  
=  $y_{1} \int \frac{e^{-x}}{(y_{1})^{2}} dx$   
=  $y_{1} \left( \frac{(-1-x)e^{-x} + (x+2)x \text{ expIntegral}_{1}(x)}{2(x+2)x} \right)$ 

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
  
=  $c_1(x(x+2)) + c_2\left(x(x+2)\left(\frac{(-1-x)e^{-x} + (x+2)x \text{ expIntegral}_1(x)}{2(x+2)x}\right)\right)$ 

Summary

The solution(s) found are the following

$$y = c_1 x(x+2) + c_2 \left(\frac{(-1-x)e^{-x}}{2} + \frac{(x+2)x \operatorname{expIntegral}_1(x)}{2}\right)$$
(1)

Verification of solutions

$$y = c_1 x(x+2) + c_2 \left(\frac{(-1-x)e^{-x}}{2} + \frac{(x+2)x \operatorname{expIntegral}_1(x)}{2}\right)$$

Verified OK.

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
A Liouvillian solution exists
Reducible group (found an exponential solution)
Group is reducible, not completely reducible
<- Kovacics algorithm successful`</pre>
```

Solution by Maple Time used: 0.015 (sec). Leaf size: 28

dsolve(x\*diff(y(x),x\$2)+x\*diff(y(x),x)-2\*y(x)=0,y(x), singsol=all)

$$y(x) = -\frac{(x+1)c_2e^{-x}}{2} + (x+2)x\left(c_1 + \frac{\text{expIntegral}_1(x)c_2}{2}\right)$$

Solution by Mathematica Time used: 0.125 (sec). Leaf size: 39

DSolve[x\*y''[x]+x\*y'[x]-2\*y[x]==0,y[x],x,IncludeSingularSolutions -> True]

$$y(x) \rightarrow c_1 x(x+2) - \frac{1}{2} c_2 e^{-x} (e^x (x+2) x \operatorname{ExpIntegralEi}(-x) + x + 1)$$

# 3.8 problem 10.4.8 (h)

3.8.1	Solving as second order bessel ode ode	326
3.8.2	Solving using Kovacic algorithm	327
3.8.3	Maple step by step solution	332
1 11		

Internal problem ID [5071] Internal file name [OUTPUT/4564\_Sunday\_June\_05\_2022\_03\_00\_58\_PM\_93789482/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients.
Second order and Homogeneous. page 318
Problem number: 10.4.8 (h).
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second\_order\_bessel\_ode" Maple gives the following as the ode type

```
[[_2nd_order, _with_linear_symmetries]]
```

$$x(x-1)^2 \, y'' - 2y = 0$$

#### 3.8.1 Solving as second order bessel ode ode

Writing the ode as

$$x^2 y'' - \frac{2y}{x} = 0 (1)$$

Bessel ode has the form

$$x^{2}y'' + xy' + (-n^{2} + x^{2})y = 0$$
(2)

The generalized form of Bessel ode is given by Bowman (1958) as the following

$$x^{2}y'' + (1 - 2\alpha)xy' + (\beta^{2}\gamma^{2}x^{2\gamma} - n^{2}\gamma^{2} + \alpha^{2})y = 0$$
(3)

With the standard solution

$$y = x^{\alpha}(c_1 \operatorname{BesselJ}(n, \beta x^{\gamma}) + c_2 \operatorname{BesselY}(n, \beta x^{\gamma}))$$
(4)

Comparing (3) to (1) and solving for  $\alpha,\beta,n,\gamma$  gives

$$\alpha = \frac{1}{2}$$
$$\beta = 2i\sqrt{2}$$
$$n = 1$$
$$\gamma = -\frac{1}{2}$$

Substituting all the above into (4) gives the solution as

$$y = ic_1\sqrt{x} \text{ BesselI}\left(1, \frac{2\sqrt{2}}{\sqrt{x}}\right) + c_2\sqrt{x} \text{ BesselY}\left(1, \frac{2i\sqrt{2}}{\sqrt{x}}\right)$$

Summary

The solution(s) found are the following

$$y = ic_1\sqrt{x} \text{ BesselI}\left(1, \frac{2\sqrt{2}}{\sqrt{x}}\right) + c_2\sqrt{x} \text{ BesselY}\left(1, \frac{2i\sqrt{2}}{\sqrt{x}}\right)$$
 (1)

\_

Verification of solutions

$$y = ic_1\sqrt{x} \text{ BesselI}\left(1, \frac{2\sqrt{2}}{\sqrt{x}}\right) + c_2\sqrt{x} \text{ BesselY}\left(1, \frac{2i\sqrt{2}}{\sqrt{x}}\right)$$

Verified OK.

# 3.8.2 Solving using Kovacic algorithm

Writing the ode as

$$x(x-1)^2 y'' - 2y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = x(x-1)^{2}$$
  

$$B = 0$$

$$C = -2$$
(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = y e^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t}$$
(5)  
=  $\frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$ 

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{2}{x \left(x - 1\right)^2}$$
(6)

Comparing the above to (5) shows that

$$s = 2$$
$$t = x(x-1)^2$$

Therefore eq. (4) becomes

$$z''(x) = \left(\frac{2}{x(x-1)^2}\right)z(x) \tag{7}$$

Equation (7) is now solved. After finding z(x) then y is found using the inverse transformation

$$y = z(x) \, e^{-\int rac{B}{2A} \, dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at  $\infty$ . The following table summarizes these cases.

Case	Allowed pole order for $r$	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \cdots\}$	$\{\cdots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \cdots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condi- tion is satisfied. Hence the following set of pole orders are all allowed. $\{1,2\},\{1,3\},\{2\},\{3\},\{3,4\},\{1,2,5\}.$	no condition
3	{1,2}	$\{2, 3, 4, 5, 6, 7, \cdots\}$

Table 57: Necessary conditions for each Kovacic case

The order of r at  $\infty$  is the degree of t minus the degree of s. Therefore

$$O(\infty) = \deg(t) - \deg(s)$$
$$= 3 - 0$$
$$= 3$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of  $t = x(x-1)^2$ . There is a pole at x = 0 of order 1. There is a pole at x = 1 of order 2. Since there is no odd order pole larger than 2 and the order at  $\infty$  is 3 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 3 then the necessary conditions for case two are met. Since pole order is not larger than 2 and the order at  $\infty$  is 3 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case n = 1.

Looking at poles of order 1. For the pole at x = 0 of order 1 then

$$[\sqrt{r}]_c = 0$$
$$\alpha_c^+ = 1$$
$$\alpha_c^- = 1$$

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{2}{(x-1)^2} + \frac{2}{x} - \frac{2}{x-1}$$

For the pole at x = 1 let b be the coefficient of  $\frac{1}{(x-1)^2}$  in the partial fractions decomposition of r given above. Therefore b = 2. Hence

$$\begin{split} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1+4b} = 2 \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1+4b} = -1 \end{split}$$

Since the order of r at  $\infty$  is 3 > 2 then

$$[\sqrt{r}]_{\infty} = 0$$
  
 $lpha_{\infty}^{+} = 0$   
 $lpha_{\infty}^{-} = 1$ 

The following table summarizes the findings so far for poles and for the order of r at  $\infty$  where r is 2

$r = \frac{1}{x \left( x - 1 \right)^2}$							
pole $c$	location	pole	order	[1	$\sqrt{r}]_c$	$lpha_c^+$	$lpha_c^-$
	0	]	L		0	0	1
	1	6 2	2		0	2	-1
0	rder of r	at $\infty$	$\sqrt{r}$	ø	$\alpha^+_\infty$	$\alpha_{\infty}^{-}$	
	0		0		0	<b>–</b>	

Now that the all  $[\sqrt{r}]_c$  and its associated  $\alpha_c^{\pm}$  have been determined for all the poles in the set  $\Gamma$  and  $[\sqrt{r}]_{\infty}$  and its associated  $\alpha_{\infty}^{\pm}$  have also been found, the next step is to determine possible non negative integer d from these using

$$d = lpha_{\infty}^{s(\infty)} - \sum_{c \in \Gamma} lpha_{c}^{s(c)}$$

Where s(c) is either + or - and  $s(\infty)$  is the sign of  $\alpha_{\infty}^{\pm}$ . This is done by trial over all set of families  $s = (s(c))_{c \in \Gamma \cup \infty}$  until such d is found to work in finding candidate  $\omega$ . Trying  $\alpha_{\infty}^{+} = 0$  then

$$d = \alpha_{\infty}^{+} - (\alpha_{c_{1}}^{-} + \alpha_{c_{2}}^{-})$$
  
= 0 - (0)  
= 0

Since d an integer and  $d \ge 0$  then it can be used to find  $\omega$  using

$$\omega = \sum_{c \in \Gamma} \left( s(c) [\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{x - c} \right) + s(\infty) [\sqrt{r}]_{\infty}$$

Substituting the above values in the above results in

$$\begin{split} \omega &= \left( (-) [\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}}{x - c_1} \right) + \left( (-) [\sqrt{r}]_{c_2} + \frac{\alpha_{c_2}}{x - c_2} \right) + (+) [\sqrt{r}]_{\infty} \\ &= \frac{1}{x} - \frac{1}{x - 1} + (0) \\ &= \frac{1}{x} - \frac{1}{x - 1} \\ &= -\frac{1}{x (x - 1)} \end{split}$$

Now that  $\omega$  is determined, the next step is find a corresponding minimal polynomial p(x) of degree d = 0 to solve the ode. The polynomial p(x) needs to satisfy the equation

$$p'' + 2\omega p' + \left(\omega' + \omega^2 - r\right)p = 0 \tag{1A}$$

Let

$$p(x) = 1 \tag{2A}$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(\frac{1}{x} - \frac{1}{x-1}\right)(0) + \left(\left(-\frac{1}{x^2} + \frac{1}{(x-1)^2}\right) + \left(\frac{1}{x} - \frac{1}{x-1}\right)^2 - \left(\frac{2}{x(x-1)^2}\right)\right) = 0$$
$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode z'' = rz is

$$z_1(x) = p e^{\int \omega \, dx}$$
$$= e^{\int \left(\frac{1}{x} - \frac{1}{x-1}\right) dx}$$
$$= \frac{x}{x-1}$$

The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2}\frac{B}{A}\,dx}$$

Since B = 0 then the above reduces to

$$y_1 = z_1$$
$$= \frac{x}{x-1}$$

Which simplifies to

$$y_1 = \frac{x}{x-1}$$

The second solution  $y_2$  to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} \, dx}}{y_1^2} \, dx$$

Since B = 0 then the above becomes

$$y_{2} = y_{1} \int \frac{1}{y_{1}^{2}} dx$$
$$= \frac{x}{x-1} \int \frac{1}{\frac{x^{2}}{(x-1)^{2}}} dx$$
$$= \frac{x}{x-1} \left( x - 2\ln(x) - \frac{1}{x} \right)$$

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$
  
=  $c_1 \left( \frac{x}{x-1} \right) + c_2 \left( \frac{x}{x-1} \left( x - 2\ln(x) - \frac{1}{x} \right) \right)$ 

Summary

The solution(s) found are the following

$$y = \frac{c_1 x}{x - 1} + \frac{c_2 (-2\ln(x) x + x^2 - 1)}{x - 1} \tag{1}$$

Verification of solutions

$$y = \frac{c_1 x}{x - 1} + \frac{c_2(-2\ln(x)x + x^2 - 1)}{x - 1}$$

Verified OK.

### 3.8.3 Maple step by step solution

Let's solve

$$x(x-1)^2 y'' - 2y = 0$$

- Highest derivative means the order of the ODE is 2 y''
- Isolate 2nd derivative

$$y'' = \frac{2y}{x(x-1)^2}$$

• Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \tfrac{2y}{x(x-1)^2} = 0$$

Check to see if  $x_0$  is a regular singular point

• Define functions

$$\left[P_2(x) = 0, P_3(x) = -\frac{2}{x(x-1)^2}\right]$$

- $\circ \quad x \cdot P_2(x) ext{ is analytic at } x = 0$  $(x \cdot P_2(x)) \Big|_{x=0} = 0$
- $x^2 \cdot P_3(x)$  is analytic at x = 0

$$\left(x^2 \cdot P_3(x)\right) \bigg|_{x=0} = 0$$

 $\circ \quad x = 0$  is a regular singular point

Check to see if  $x_0$  is a regular singular point

 $x_0 = 0$ 

• Multiply by denominators

$$x(x-1)^2 \, y'' - 2y = 0$$

• Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^{k+r}$$

 $\Box$  Rewrite ODE with series expansions

• Convert  $x^m \cdot y''$  to series expansion for m = 1..3

$$x^m \cdot y'' = \sum_{k=0}^{\infty} a_k (k+r) (k+r-1) x^{k+r-2+m}$$

 $\circ \quad \text{Shift index using } k->k+2-m$ 

$$x^{m} \cdot y'' = \sum_{k=-2+m}^{\infty} a_{k+2-m}(k+2-m+r) \left(k+1-m+r\right) x^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r(-1+r) x^{-1+r} + (a_1(1+r) r - 2a_0(r^2 - r + 1)) x^r + \left(\sum_{k=1}^{\infty} (a_{k+1}(k+1+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k^2 - r + 1)) x^r + (a_1(1+r) (k+r) (k+r) - 2a_k(k+r) (k+r) - 2a_k(k+r) (k+r) - 2a_k(k+r) (k+r) - 2a_k(k+r) (k+r) (k+r) - 2a_k(k+r) (k+r) (k+r) (k+r) - 2a_k(k+r) (k+r) (k+r) (k+r) (k+r) - 2a_k(k+r) (k+r) (k+r) (k+r) (k+r) (k+r) - 2a_k(k+r) (k+r) (k+r)$$

- $a_0$  cannot be 0 by assumption, giving the indicial equation r(-1+r) = 0
- Values of r that satisfy the indicial equation

 $r \in \{0,1\}$ 

• Each term must be 0

$$a_1(1+r)r - 2a_0(r^2 - r + 1) = 0$$

• Each term in the series must be 0, giving the recursion relation  $(-2a_k + a_{k-1} + a_{k+1})k^2 + ((-4a_k + 2a_{k-1} + 2a_{k+1})r + 2a_k - 3a_{k-1} + a_{k+1})k + (-2a_k + a_{k-1} + a_{k+1})r + (-2a_k + a_{k+1})r +$ 

• Shift index using 
$$k - > k + 1$$
  
 $(-2a_{k+1} + a_k + a_{k+2})(k+1)^2 + ((-4a_{k+1} + 2a_k + 2a_{k+2})r + 2a_{k+1} - 3a_k + a_{k+2})(k+1) + (-a_{k+1} + a_k + a_{k+2})(k+1) + (-a_{k+1} + a_{k+2})($ 

• Recursion relation that defines series solution to ODE

$$a_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} + 2kra_k - 4kra_{k+1} + r^2 a_k - 2r^2 a_{k+1} - ka_k - 2ka_{k+1} - ra_k - 2ra_{k+1} - 2a_{k+1}}{k^2 + 2kr + r^2 + 3k + 3r + 2}$$

• Recursion relation for r = 0 $a_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - ka_k - 2ka_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}$ 

• Solution for 
$$r = 0$$

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - ka_k - 2ka_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0\right]$$

• Recursion relation for r = 1

$$a_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} + ka_k - 6ka_{k+1} - 6a_{k+1}}{k^2 + 5k + 6}$$

• Solution for 
$$r = 1$$

$$\left[y = \sum_{k=0}^{\infty} a_k x^{k+1}, a_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} + ka_k - 6ka_{k+1} - 6a_{k+1}}{k^2 + 5k + 6}, 2a_1 - 2a_0 = 0\right]$$

• Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k x^k\right) + \left(\sum_{k=0}^{\infty} b_k x^{k+1}\right), a_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - ka_k - 2ka_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - ka_k - 2ka_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - ka_k - 2ka_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - ka_k - 2ka_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - ka_k - 2ka_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - 2a_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+1} - 2a_{k+1}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = 0, b_{k+2} = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = -\frac{k^2 a_k - 2k^2 a_{k+2}}{k^2 + 3k + 2}, -2a_0 = -\frac{k^2 a$$

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
A Liouvillian solution exists
Reducible group (found an exponential solution)
Group is reducible, not completely reducible
<- Kovacics algorithm successful`</pre>
```

Solution by Maple Time used: 0.0 (sec). Leaf size: 27

 $dsolve(x*(x-1)^2*diff(y(x),x$2)-2*y(x)=0,y(x), singsol=all)$ 

$$y(x) = \frac{2c_2x\ln(x) - c_2x^2 + c_1x + c_2}{x - 1}$$

Solution by Mathematica Time used: 0.034 (sec). Leaf size: 33

DSolve[x\*(x-1)^2\*y''[x]-2\*y[x]==0,y[x],x,IncludeSingularSolutions -> True]

$$y(x) \rightarrow \frac{-c_2 x^2 - c_1 x + 2c_2 x \log(x) + c_2}{x - 1}$$

# 3.9 problem 10.4.9 (i)

3.9.1	Solving as linear ode	336
3.9.2	Solving as first order ode lie symmetry lookup ode	338
3.9.3	Solving as exact ode	342
3.9.4	Maple step by step solution	347

Internal problem ID [5072] Internal file name [OUTPUT/4565\_Sunday\_June\_05\_2022\_03\_00\_59\_PM\_16506391/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients.
Second order and Homogeneous. page 318
Problem number: 10.4.9 (i).
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_linear]

$$y' - \frac{2y}{x} = x^2$$

### 3.9.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{2}{x}$$
$$q(x) = x^2$$

Hence the ode is

$$y' - \frac{2y}{x} = x^2$$

The integrating factor  $\mu$  is

$$\mu = \mathrm{e}^{\int -\frac{2}{x}dx}$$
 $= rac{1}{x^2}$ 

The ode becomes

$$\begin{split} \frac{\mathrm{d}}{\mathrm{d}x}(\mu y) &= (\mu) \left(x^2\right) \\ \frac{\mathrm{d}}{\mathrm{d}x} \left(\frac{y}{x^2}\right) &= \left(\frac{1}{x^2}\right) \left(x^2\right) \\ \mathrm{d} \left(\frac{y}{x^2}\right) &= \mathrm{d}x \end{split}$$

Integrating gives

$$\frac{y}{x^2} = \int \mathrm{d}x$$
$$\frac{y}{x^2} = x + c_1$$

Dividing both sides by the integrating factor  $\mu = \frac{1}{x^2}$  results in

$$y = c_1 x^2 + x^3$$

which simplifies to

$$y = x^2(x + c_1)$$

Summary

The solution(s) found are the following

$$y = x^2(x + c_1) \tag{1}$$

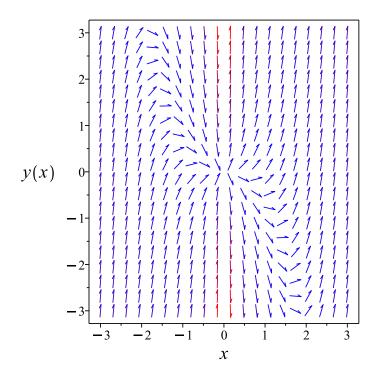


Figure 42: Slope field plot

Verification of solutions

$$y = x^2(x + c_1)$$

Verified OK.

# 3.9.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x^3 + 2y}{x}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
(A)

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi$ ,  $\eta$ 

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f(rac{y}{x})$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y'=rac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)  y + f_2(x)  y^2$	0	$e^{-\int f_1 dx}$

Table 59: Lie symmetry infinitesimal lookup table for known first order ODE's

The above table shows that

$$\begin{aligned} \xi(x,y) &= 0\\ \eta(x,y) &= x^2 \end{aligned} \tag{A1}$$

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where S(R). Since  $\xi = 0$  then in this special case

$$R = x$$

 ${\cal S}$  is found from

$$S = \int rac{1}{\eta} dy \ = \int rac{1}{x^2} dy$$

Which results in

$$S = \frac{y}{x^2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}$$
(2)

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = \frac{x^3 + 2y}{x}$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 0$$

$$S_x = -\frac{2y}{x^3}$$

$$S_y = \frac{1}{x^2}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 1 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = R + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{y}{x^2} = x + c_1$$

Which simplifies to

$$\frac{y}{x^2} = x + c_1$$

Which gives

$$y = x^2(x + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	$ODE  ext{ in canonical coordinates} \ (R,S)$
$\frac{dy}{dx} = \frac{x^3 + 2y}{x}$	$R = x$ $S = \frac{y}{x^2}$	$\frac{dS}{dR} = 1$

Summary

The solution(s) found are the following

$$y = x^2(x + c_1) \tag{1}$$

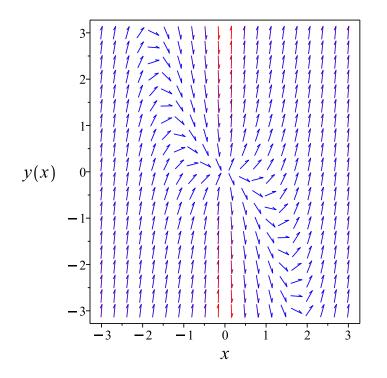


Figure 43: Slope field plot

Verification of solutions

$$y = x^2(x + c_1)$$

Verified OK.

### 3.9.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x,y) + N(x,y)\frac{dy}{dx} = 0$$
 (A)

We assume there exists a function  $\phi(x, y) = c$  where c is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t. x gives

$$\frac{d}{dx}\phi(x,y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$
$$\frac{\partial \phi}{\partial y} = N$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0$$
(1A)

Therefore

$$dy = \left(\frac{2y}{x} + x^2\right) dx$$
$$\left(-\frac{2y}{x} - x^2\right) dx + dy = 0$$
(2A)

Comparing (1A) and (2A) shows that

$$\begin{split} M(x,y) &= -\frac{2y}{x} - x^2 \\ N(x,y) &= 1 \end{split}$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$rac{\partial M}{\partial y} = rac{\partial}{\partial y} \left( -rac{2y}{x} - x^2 
ight) \ = -rac{2}{x}$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x}(1)$$
$$= 0$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is <u>not exact</u>. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$A = \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right)$$
$$= 1 \left( \left( -\frac{2}{x} \right) - (0) \right)$$
$$= -\frac{2}{x}$$

Since A does not depend on y, then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\mu = e^{\int A \, \mathrm{d}x}$$
$$= e^{\int -\frac{2}{x} \, \mathrm{d}x}$$

The result of integrating gives

$$\mu = e^{-2\ln(x)}$$
$$= \frac{1}{x^2}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called  $\overline{M}$  and  $\overline{N}$  for now so not to confuse them with the original M and N.

$$egin{aligned} \overline{M} &= \mu M \ &= rac{1}{x^2}igg(-rac{2y}{x}-x^2igg) \ &= rac{-x^3-2y}{x^3} \end{aligned}$$

And

$$\overline{N} = \mu N$$
  
 $= rac{1}{x^2}(1)$   
 $= rac{1}{x^2}$ 

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$\overline{M} + \overline{N} \frac{\mathrm{d}y}{\mathrm{d}x} = 0$$
$$\left(\frac{-x^3 - 2y}{x^3}\right) + \left(\frac{1}{x^2}\right) \frac{\mathrm{d}y}{\mathrm{d}x} = 0$$

The following equations are now set up to solve for the function  $\phi(x, y)$ 

$$\frac{\partial \phi}{\partial x} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \overline{M} dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int \frac{-x^3 - 2y}{x^3} dx$$

$$\phi = -x + \frac{y}{x^2} + f(y)$$
(3)

Where f(y) is used for the constant of integration since  $\phi$  is a function of both x and y. Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = \frac{1}{x^2} + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{1}{x^2}$ . Therefore equation (4) becomes

$$\frac{1}{x^2} = \frac{1}{x^2} + f'(y) \tag{5}$$

Solving equation (5) for f'(y) gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for f(y) into equation (3) gives  $\phi$ 

$$\phi = -x + \frac{y}{x^2} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -x + \frac{y}{x^2}$$

The solution becomes

$$y = x^2(x + c_1)$$

Summary

The solution(s) found are the following

$$y = x^2(x + c_1) \tag{1}$$

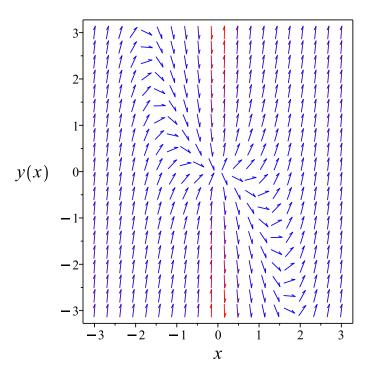


Figure 44: Slope field plot

Verification of solutions

$$y = x^2(x + c_1)$$

Verified OK.

### 3.9.4 Maple step by step solution

Let's solve

 $y' - \frac{2y}{x} = x^2$ 

- Highest derivative means the order of the ODE is 1
  - y'
- Isolate the derivative

$$y' = \frac{2y}{x} + x^2$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE  $y' \frac{2y}{x} = x^2$
- The ODE is linear; multiply by an integrating factor  $\mu(x)$  $\mu(x) \left(y' - \frac{2y}{x}\right) = \mu(x) x^2$
- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$

$$\mu(x)\left(y'-rac{2y}{x}
ight)=\mu'(x)\,y+\mu(x)\,y'$$

- Isolate  $\mu'(x)$  $\mu'(x) = -\frac{2\mu(x)}{x}$
- Solve to find the integrating factor

$$\mu(x) = \frac{1}{x^2}$$

• Integrate both sides with respect to x

$$\int \left( rac{d}{dx}(\mu(x)\,y) 
ight) dx = \int \mu(x)\,x^2 dx + c_1$$

• Evaluate the integral on the lhs

$$\mu(x) y = \int \mu(x) x^2 dx + c_1$$

• Solve for y

$$y = rac{\int \mu(x)x^2 dx + c_1}{\mu(x)}$$

• Substitute  $\mu(x) = \frac{1}{x^2}$ 

$$y = x^2 \left( \int 1 dx + c_1 \right)$$

• Evaluate the integrals on the rhs

$$y = x^2(x + c_1)$$

#### Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`</pre>
```

Solution by Maple Time used: 0.0 (sec). Leaf size: 11

 $dsolve(diff(y(x),x)-2*y(x)/x-x^2=0,y(x), singsol=all)$ 

$$y(x) = (x + c_1) x^2$$

✓ Solution by Mathematica

Time used:  $0.\overline{026}$  (sec). Leaf size: 13

DSolve[y'[x]-2\*y[x]/x-x^2==0,y[x],x,IncludeSingularSolutions -> True]

$$y(x) \rightarrow x^2(x+c_1)$$

# 3.10 problem 10.4.9 (ii)

Solving as linear ode	349
2 Solving as first order ode lie symmetry lookup ode	351
3 Solving as exact ode	355
4 Maple step by step solution	360
	<ol> <li>Solving as linear ode</li></ol>

Internal problem ID [5073] Internal file name [OUTPUT/4566\_Sunday\_June\_05\_2022\_03\_01\_00\_PM\_73400741/index.tex]

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients.
Second order and Homogeneous. page 318
Problem number: 10.4.9 (ii).
ODE order: 1.
ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_linear]

$$y' + \frac{2y}{x} = x^3$$

#### 3.10.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{2}{x}$$
$$q(x) = x^3$$

Hence the ode is

$$y' + \frac{2y}{x} = x^3$$

The integrating factor  $\mu$  is

$$\mu = e^{\int \frac{2}{x} dx}$$
$$= x^2$$

The ode becomes

$$egin{aligned} &rac{\mathrm{d}}{\mathrm{d}x}(\mu y)=(\mu)\left(x^3
ight)\ &rac{\mathrm{d}}{\mathrm{d}x}(y\,x^2)=\left(x^2
ight)\left(x^3
ight)\ &\mathrm{d}\left(y\,x^2
ight)=x^5\,\mathrm{d}x \end{aligned}$$

Integrating gives

Dividing both sides by the integrating factor 
$$\mu = x^2$$
 results in

$$y = \frac{x^4}{6} + \frac{c_1}{x^2}$$

 $y x^2 = \int x^5 dx$  $y x^2 = \frac{x^6}{6} + c_1$ 

Summary

The solution(s) found are the following

$$y = \frac{x^4}{6} + \frac{c_1}{x^2} \tag{1}$$

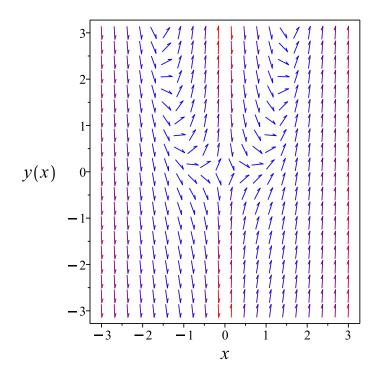


Figure 45: Slope field plot

Verification of solutions

$$y = \frac{x^4}{6} + \frac{c_1}{x^2}$$

Verified OK.

## 3.10.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{-x^4 + 2y}{x}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0$$
 (A)

The type of this ode is known. It is of type linear. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$ 

ODE class	Form	ξ	η
linear ode	y' = f(x)y(x) + g(x)	0	$e^{\int f dx}$
separable ode	y' = f(x) g(y)	$\frac{1}{f}$	0
quadrature ode	y' = f(x)	0	1
quadrature ode	y' = g(y)	1	0
homogeneous ODEs of Class A	$y' = f(rac{y}{x})$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x) F\left(\frac{y}{x}\right)$	$x^2$	xy
First order special form ID 1	$y' = g(x) e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y'=rac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x - a_2b_1x - b_1c_2 + b_2c_1}{a_1b_2 - a_2b_1}$	$\frac{a_1b_2y - a_2b_1y - a_1c_2 - a_2c_1}{a_1b_2 - a_2b_1}$
Bernoulli ode	$y' = f(x) y + g(x) y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)  y + f_2(x)  y^2$	0	$e^{-\int f_1 dx}$

Table 62: Lie symmetry infinitesimal lookup table for known first order ODE's

The above table shows that

$$\xi(x,y) = 0$$
  

$$\eta(x,y) = \frac{1}{x^2}$$
(A1)

The next step is to determine the canonical coordinates R, S. The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where S(R). Since  $\xi = 0$  then in this special case

$$R = x$$

 ${\cal S}$  is found from

$$S = \int rac{1}{\eta} dy \ = \int rac{1}{rac{1}{x^2}} dy$$

Which results in

$$S = y x^2$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \tag{2}$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x,y) = -\frac{-x^4 + 2y}{x}$$

Evaluating all the partial derivatives gives

$$egin{aligned} R_x &= 1 \ R_y &= 0 \ S_x &= 2xy \ S_y &= x^2 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = x^5 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R^5$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordiates R, S. Integrating the above gives

$$S(R) = \frac{R^6}{6} + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$yx^2 = rac{x^6}{6} + c_1$$

Which simplifies to

$$yx^2 = \frac{x^6}{6} + c_1$$

Which gives

$$y = \frac{x^6 + 6c_1}{6x^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	$ODE  ext{ in canonical coordinates} \ (R,S)$
$\frac{dy}{dx} = -\frac{-x^4 + 2y}{x}$	$R = x$ $S = y x^2$	$\frac{dS}{dR} = R^5$

#### Summary

The solution(s) found are the following

$$y = \frac{x^6 + 6c_1}{6x^2} \tag{1}$$

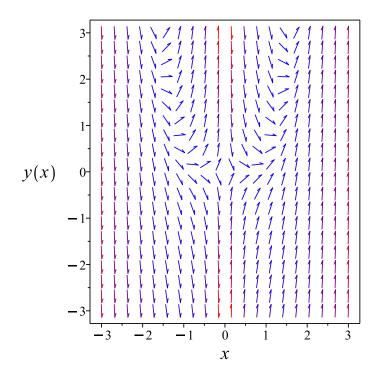


Figure 46: Slope field plot

Verification of solutions

$$y = \frac{x^6 + 6c_1}{6x^2}$$

Verified OK.

### 3.10.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x,y) + N(x,y)\frac{dy}{dx} = 0$$
 (A)

We assume there exists a function  $\phi(x, y) = c$  where c is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t. x gives

$$\frac{d}{dx}\phi(x,y) = 0$$

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y}\frac{dy}{dx} = 0$$
(B)

Hence

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x,y) dx + N(x,y) dy = 0$$
(1A)

Therefore

$$dy = \left(-\frac{2y}{x} + x^3\right) dx$$
$$\left(-x^3 + \frac{2y}{x}\right) dx + dy = 0$$
(2A)

Comparing (1A) and (2A) shows that

$$\begin{split} M(x,y) &= -x^3 + \frac{2y}{x} \\ N(x,y) &= 1 \end{split}$$

The next step is to determine if the ODE is is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$rac{\partial M}{\partial y} = rac{\partial}{\partial y} \left( -x^3 + rac{2y}{x} 
ight) = rac{2}{x}$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x}(1)$$
$$= 0$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is <u>not exact</u>. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$A = \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right)$$
$$= 1 \left( \left( \frac{2}{x} \right) - (0) \right)$$
$$= \frac{2}{x}$$

Since A does not depend on y, then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\mu = e^{\int A \, \mathrm{d}x}$$
$$= e^{\int \frac{2}{x} \, \mathrm{d}x}$$

The result of integrating gives

$$\mu = e^{2\ln(x)}$$
$$= x^2$$

M and N are multiplied by this integrating factor, giving new M and new N which are called  $\overline{M}$  and  $\overline{N}$  for now so not to confuse them with the original M and N.

$$\overline{M} = \mu M$$
  
=  $x^2 \left( -x^3 + \frac{2y}{x} \right)$   
=  $-x \left( x^4 - 2y \right)$ 

And

$$\overline{N} = \mu N$$
  
=  $x^2(1)$   
=  $x^2$ 

Now a modified ODE is ontained from the original ODE, which is exact and can be solved. The modified ODE is

$$\overline{M} + \overline{N} rac{\mathrm{d}y}{\mathrm{d}x} = 0$$
  
 $\left(-x \left(x^4 - 2y\right)\right) + \left(x^2\right) rac{\mathrm{d}y}{\mathrm{d}x} = 0$ 

The following equations are now set up to solve for the function  $\phi(x, y)$ 

$$\frac{\partial \phi}{\partial x} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \overline{M} dx$$
$$\int \frac{\partial \phi}{\partial x} dx = \int -x(x^4 - 2y) dx$$
$$\phi = -\frac{1}{6}x^6 + yx^2 + f(y)$$
(3)

Where f(y) is used for the constant of integration since  $\phi$  is a function of both x and y. Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x^2 + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = x^2$ . Therefore equation (4) becomes

$$x^2 = x^2 + f'(y) (5)$$

Solving equation (5) for f'(y) gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for f(y) into equation (3) gives  $\phi$ 

$$\phi = -\frac{1}{6}x^6 + y\,x^2 + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\frac{1}{6}x^6 + y\,x^2$$

The solution becomes

$$y = \frac{x^6 + 6c_1}{6x^2}$$

Summary

The solution(s) found are the following

$$y = \frac{x^6 + 6c_1}{6x^2} \tag{1}$$

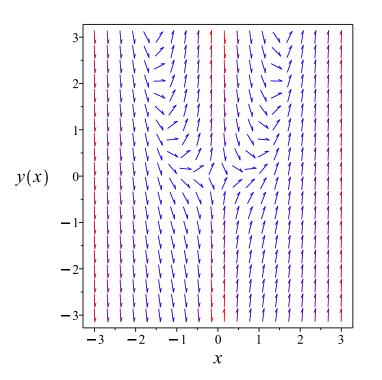


Figure 47: Slope field plot

Verification of solutions

$$y = \frac{x^6 + 6c_1}{6x^2}$$

Verified OK.

### 3.10.4 Maple step by step solution

Let's solve

 $y' + \frac{2y}{x} = x^3$ 

• Highest derivative means the order of the ODE is 1

• Isolate the derivative

$$y' = -\frac{2y}{x} + x^3$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE  $y' + \frac{2y}{x} = x^3$
- The ODE is linear; multiply by an integrating factor  $\mu(x)$  $\mu(x) \left(y' + \frac{2y}{x}\right) = \mu(x) x^3$
- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$

$$\mu(x)\left(y'+\frac{2y}{x}\right) = \mu'(x)\,y + \mu(x)\,y'$$

- Isolate  $\mu'(x)$  $\mu'(x) = \frac{2\mu(x)}{x}$
- Solve to find the integrating factor  $\mu(x) = x^2$
- Integrate both sides with respect to x $\int \left(\frac{d}{dx}(\mu(x) y)\right) dx = \int \mu(x) x^3 dx + c_1$
- Evaluate the integral on the lhs

 $\mu(x) y = \int \mu(x) x^3 dx + c_1$ 

- Solve for y $y = rac{\int \mu(x) x^3 dx + c_1}{\mu(x)}$
- Substitute  $\mu(x) = x^2$

$$y = rac{\int x^5 dx + c_1}{x^2}$$

• Evaluate the integrals on the rhs

$$y = \frac{\frac{x^6}{6} + c_1}{x^2}$$

• Simplify  $y = \frac{x^6 + 6c_1}{6x^2}$ 

#### Maple trace

`Methods for first order ODEs: --- Trying classification methods --trying a quadrature trying 1st order linear <- 1st order linear successful`</pre>

Solution by Maple Time used: 0.015 (sec). Leaf size: 16

 $dsolve(diff(y(x),x)+2*y(x)/x-x^3=0,y(x), singsol=all)$ 

$$y(x) = \frac{x^6 + 6c_1}{6x^2}$$

✓ Solution by Mathematica

Time used: 0.026 (sec). Leaf size: 13

DSolve[y'[x]-2\*y[x]/x-x^2==0,y[x],x,IncludeSingularSolutions -> True]

 $y(x) \rightarrow x^2(x+c_1)$ 

## 3.11 problem 10.4.10

3.11.1 Maple step by step solution  $\ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots 362$ 

Internal problem ID [5074]

 $Internal\,file\,name\,[\texttt{OUTPUT/4567}\_\texttt{Sunday}\_\texttt{June}\_\texttt{05}\_\texttt{2022}\_\texttt{03}\_\texttt{01}\_\texttt{01}\_\texttt{PM}\_\texttt{94712491}/\texttt{index}.\texttt{tex}]$ 

Book: Basic Training in Mathematics. By R. Shankar. Plenum Press. NY. 1995
Section: Chapter 10, Differential equations. Section 10.4, ODEs with variable Coefficients.
Second order and Homogeneous. page 318
Problem number: 10.4.10.
ODE order: 2.
ODE degree: 1.

The type(s) of ODE detected by this program : "unknown"

Maple gives the following as the ode type

[\_Laguerre]

Unable to solve or complete the solution.

$$xy'' + (1 - x)y' + my = 0$$

### 3.11.1 Maple step by step solution

Let's solve

y''x + (1-x)y' + my = 0

- Highest derivative means the order of the ODE is 2 y''
- Isolate 2nd derivative

$$y'' = \frac{(x-1)y'}{x} - \frac{my}{x}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear  $y'' \frac{(x-1)y'}{x} + \frac{my}{x} = 0$
- $\Box$  Check to see if  $x_0 = 0$  is a regular singular point
  - Define functions

$$\left[P_2(x) = -\frac{x-1}{x}, P_3(x) = \frac{m}{x}\right]$$

 $\circ \quad x \cdot P_2(x) \text{ is analytic at } x = 0$ 

$$(x \cdot P_2(x)) \bigg|_{x=0} = 1$$

- $x^2 \cdot P_3(x)$  is analytic at x = 0 $(x^2 \cdot P_3(x)) \Big|_{x=0} = 0$
- $\circ \quad x = 0 \text{is a regular singular point}$

Check to see if  $x_0 = 0$  is a regular singular point

$$x_0 = 0$$

• Multiply by denominators

$$y''x + (1-x)y' + my = 0$$

• Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^{k+r}$$

 $\Box$  Rewrite ODE with series expansions

• Convert  $x^m \cdot y'$  to series expansion for m = 0..1

$$x^m \cdot y' = \sum_{k=0}^{\infty} a_k(k+r) \, x^{k+r-1+m}$$

 $\circ \quad \text{Shift index using } k->k+1-m$ 

$$x^m \cdot y' = \sum_{k=-1+m}^{\infty} a_{k+1-m}(k+1-m+r) x^{k+r}$$

 $\circ \quad \text{Convert} \; x \cdot y'' \text{ to series expansion}$ 

$$x \cdot y'' = \sum_{k=0}^{\infty} a_k (k+r) (k+r-1) x^{k+r-1}$$

• Shift index using k - > k + 1

$$x \cdot y'' = \sum_{k=-1}^{\infty} a_{k+1}(k+1+r) (k+r) x^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r^2 x^{-1+r} + \left(\sum_{k=0}^{\infty} \left(a_{k+1}(k+1+r)^2 - a_k(-m+k+r)\right) x^{k+r}\right) = 0$$

•  $a_0$  cannot be 0 by assumption, giving the indicial equation  $r^2 = 0$  • Values of r that satisfy the indicial equation

r = 0

- Each term in the series must be 0, giving the recursion relation  $a_{k+1}(k+1)^2 a_k(-m+k) = 0$
- Recursion relation that defines series solution to ODE  $a_{k+1} = \frac{a_k(-m+k)}{(k+1)^2}$
- Recursion relation for r = 0

$$a_{k+1} = \frac{a_k(-m+k)}{(k+1)^2}$$

• Solution for r = 0

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+1} = \frac{a_k(-m+k)}{(k+1)^2}\right]$$

#### Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
<- No Liouvillian solutions exists
-> Trying a solution in terms of special functions:
  -> Bessel
  -> elliptic
   -> Legendre
  -> Kummer
      -> hyper3: Equivalence to 1F1 under a power @ Moebius
      <- hyper3 successful: received ODE is equivalent to the 1F1 ODE
   <- Kummer successful
<- special function solution successful`
```

Solution by Maple Time used: 0.047 (sec). Leaf size: 21

dsolve(x\*diff(y(x),x\$2)+(1-x)\*diff(y(x),x)+m\*y(x)=0,y(x), singsol=all)

 $y(x) = c_1 \operatorname{KummerM}(-m, 1, x) + c_2 \operatorname{KummerU}(-m, 1, x)$ 

Solution by Mathematica Time used: 0.021 (sec). Leaf size: 21

DSolve[x\*y''[x]+(1-x)\*y'[x]+m\*y[x]==0,y[x],x,IncludeSingularSolutions -> True]

 $y(x) \rightarrow c_1$  Hypergeometric  $U(-m, 1, x) + c_2$  Laguerre L(m, x)