## HW 3

# Math 5587 <br> Elementary Partial Differential Equations I 

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## 1 Problem 3.1.2

Find all separable eigensolutions to the heat equation $u_{t}=u_{x x}$ on $0 \leq x \leq \pi$ subject to (a) homogeneous boundary conditions $u(t, 0)=0, u(t, \pi)=0$. (b) mixed boundary conditions $u(t, 0)=0, u_{x}(t, \pi)=0$
solution
Using separation of variables, let $u(t, x)=T(t) X(x)$. Substituting this into $u_{t}=u_{x x}$ gives $T^{\prime} X=T X^{\prime \prime}$. Dividing by $X T \neq 0$ results in

$$
\frac{T^{\prime}}{T}=\frac{X^{\prime \prime}}{X}=-\lambda
$$

Where $\lambda$ is the seperation constant. The above gives the following ODE's to solve

$$
\begin{aligned}
X^{\prime \prime}(x)+\lambda X(x) & =0 \\
T^{\prime}(t)+\lambda T(t) & =0
\end{aligned}
$$

The boundary and initial conditions are transfered from the PDE to the ODE as shown below.

### 1.1 Part (a)

Using $u(t, 0)=0, u(t, \pi)=0$. Starting with the spatial ODE, and transferring the boundary conditions to the ODE results in

$$
\begin{aligned}
X^{\prime \prime}(x)+\lambda X(x) & =0 \\
X(0) & =0 \\
X(\pi) & =0
\end{aligned}
$$

This is an eigenvalue boundary value ODE. The solution to the spatial ODE is

$$
\begin{equation*}
X(x)=c_{1} e^{\sqrt{-\lambda} x}+c_{2} e^{-\sqrt{-\lambda} x} \tag{1}
\end{equation*}
$$

case $\lambda<0$
Since $\lambda<0$, then $-\lambda$ is positive. Let $\mu=-\lambda$, where $\mu$ is positive. The above solution becomes

$$
X(x)=c_{1} e^{\sqrt{\mu} x}+c_{2} e^{-\sqrt{\mu} x}
$$

Which can be written as

$$
X(x)=c_{1} \cosh (\sqrt{\mu} x)+c_{2} \sinh (\sqrt{\mu} x)
$$

At $x=0$ this gives

$$
0=c_{1}
$$

The solution now reduces to $X(x)=c_{2} \sinh (\sqrt{\mu} x)$. At $x=\pi$ this gives

$$
0=c_{2} \sinh (\sqrt{\mu} \pi)
$$

But sinh is only zero when its argument is zero. Since $\mu \neq 0$, then the only choice is that $c_{2}=0$ also. But this gives trivial solution therefore $\lambda<0$ is not an eigenvalue.
case $\lambda=0$
In this case the solution is $X(x)=c_{1}+c_{2} x$. At $x=0$ this gives $0=c_{1}$. The solution becomes $X(x)=c_{2} x$. At $x=\pi$, this gives $0=c_{2} \pi$. Therefore $c_{2}=0$ also. This also gives the trivial solution. Hence $\lambda=0$ is not an eigenvalue.
case $\lambda>0$
The solution in this case is

$$
\begin{aligned}
X(x) & =c_{1} e^{\sqrt{-\lambda} x}+c_{2} e^{-\sqrt{-\lambda} x} \\
& =c_{1} e^{i \sqrt{\lambda} x}+c_{2} e^{-i \sqrt{\lambda} x}
\end{aligned}
$$

Which can be rewritten as (the constants $c_{1}, c_{2}$ below will be different than the above $c_{1}, c_{2}$, but kept the same name for simplicity).

$$
X(x)=c_{1} \cos (\sqrt{\lambda} x)+c_{2} \sin (\sqrt{\lambda} x)
$$

At $x=0$ this gives

$$
0=c_{1}
$$

The solution now reduces to

$$
X(x)=c_{2} \sin (\sqrt{\lambda} x)
$$

At $x=\pi$ this gives

$$
0=c_{2} \sin (\sqrt{\lambda} \pi)
$$

non-trivial solution requires that $\sin (\sqrt{\lambda} \pi)=0$ which implies that $\sqrt{\lambda} \pi=n \pi, n=1,2,3, \cdots$. Hence eigenvalues are

$$
\lambda_{n}=n^{2} \quad n=1,2,3, \cdots
$$

And corresponding eigenfunctions are

$$
X_{n}(x)=\sin (n x) \quad n=1,2,3, \cdots
$$

Now that the eigenvalues and eigenfunction are found, the time ODE can be solved. The time ODE now becomes

$$
T^{\prime}(t)+n^{2} T(t)=0
$$

This is linear first order ode. The solution is $T_{n}(t)=C_{n} e^{-n^{2} t}$. Therefore the fundamental solution is

$$
\begin{aligned}
u_{n}(t, x) & =C_{n} T_{n}(t) X_{n}(x) \\
& =C_{n} e^{-n^{2} t} \sin (n x)
\end{aligned}
$$

Since this is a linear PDE, a linear combination of all fundamental solutions is a solution. Hence the general solution is

$$
u(t, x)=\sum_{n=1}^{\infty} C_{n} e^{-n^{2} t} \sin (n x)
$$

The constant $C_{n}$ can be found if initial conditions are given.

### 1.2 Part (b)

Using $u(t, 0)=0, u_{x}(t, \pi)=0$. Starting with the spatial ODE, and transferring the boundary condition to $X$, it becomes

$$
\begin{aligned}
X^{\prime \prime}(x)+\lambda X(x) & =0 \\
X(0) & =0 \\
X^{\prime}(\pi) & =0
\end{aligned}
$$

This is an eigenvalue boundary value problem. The solution to the spatial ODE is

$$
\begin{equation*}
X(x)=c_{1} e^{\sqrt{-\lambda} x}+c_{2} e^{-\sqrt{-\lambda} x} \tag{1}
\end{equation*}
$$

case $\lambda<0$
Since $\lambda<0$, then $-\lambda$ is positive. Let $\mu=-\lambda$, where $\mu$ is positive. The solution becomes

$$
X(x)=c_{1} e^{\sqrt{\mu} x}+c_{2} e^{-\sqrt{\mu} x}
$$

The above can be written as

$$
X(x)=c_{1} \cosh (\sqrt{\mu} x)+c_{2} \sinh (\sqrt{\mu} x)
$$

At $x=0$ this gives

$$
0=c_{1}
$$

Hence the solution now becomes

$$
X(x)=c_{2} \sinh (\sqrt{\mu} x)
$$

Taking derivative gives

$$
X^{\prime}(x)=c_{2} \sqrt{\mu} \cosh (\sqrt{\mu} x)
$$

And at $x=\pi$ the above gives

$$
0=c_{2} \sqrt{\mu} \cosh (\sqrt{\mu} \pi)
$$

But $\mu \neq 0$ and cosh is never zero for any argument. Hence the only choice is that $c_{2}=0$. This gives the trivial solution. Hence $\lambda<0$ is not an eigenvalue.
case $\lambda=0$
In this case the solution is $X(x)=c_{1}+c_{2} x$. At $x=0$ this results in $0=c_{1}$. The solution becomes $X(x)=c_{2} x$. Hence $X^{\prime}(x)=c_{2}$. At $x=\pi$, this implies $0=c_{2} \pi$. Therefore $c_{2}=0$ also. This gives the trivial solution. Hence $\lambda=0$ is not an eigenvalue.
case $\lambda>0$
The solution in this case is

$$
\begin{aligned}
X(x) & =c_{1} e^{\sqrt{-\lambda} x}+c_{2} e^{-\sqrt{-\lambda} x} \\
& =c_{1} e^{i \sqrt{\lambda} x}+c_{2} e^{-i \sqrt{\lambda} x}
\end{aligned}
$$

Which can be rewritten as (the constants $c_{1}, c_{2}$ below will be different than the above $c_{1}, c_{2}$, but kept the same name for simplicity).

$$
X(x)=c_{1} \cos (\sqrt{\lambda} x)+c_{2} \sin (\sqrt{\lambda} x)
$$

At $x=0$ this gives

$$
0=c_{1}
$$

The solution now reduces to

$$
X(x)=c_{2} \sin (\sqrt{\lambda} x)
$$

Therefore

$$
X^{\prime}(x)=\sqrt{\lambda} c_{2} \cos (\sqrt{\lambda} x)
$$

At $x=\pi$

$$
0=\sqrt{\lambda} c_{2} \cos (\sqrt{\lambda} \pi)
$$

Non-trivial solution requires that $\cos (\sqrt{\lambda} \pi)=0$, which implies $\sqrt{\lambda} \pi=\frac{n \pi}{2}, n=1,3,5, \cdots$. or $\sqrt{\lambda}=\frac{n}{2}, n=1,3,5, \cdots$. Therefore the eigenvalues are

$$
\lambda_{n}=\left(\frac{n}{2}\right)^{2} \quad n=1,3,5, \cdots
$$

Or

$$
\lambda_{n}=\left(\frac{2 n-1}{2}\right)^{2} \quad n=1,2,3, \cdots
$$

Few eigenvalues are $\lambda=\left\{\frac{1}{4}, \frac{9}{4}, \frac{25}{4}, \cdots\right\}$. The corresponding eigenfunctions are

$$
X_{n}(x)=\sin \left(\frac{2 n-1}{2} x\right) \quad n=1,2,3, \cdots
$$

Now that the eigenvalues and eigenfunction are found, the time ODE is solved. The time ODE now becomes

$$
T^{\prime}(t)+\left(\frac{2 n-1}{2}\right)^{2} T(t)=0
$$

This is linear first order ode. The solution is $T_{n}(t)=C_{n} e^{-\left(\frac{2 n-1}{2}\right)^{2} t}$. Therefore the fundamental solution is

$$
\begin{aligned}
u_{n}(t, x) & =C_{n} T_{n}(t) X_{n}(x) \\
& =C_{n} e^{-\left(\frac{2 n-1}{2}\right)^{2} t} \sin \left(\frac{2 n-1}{2} x\right)
\end{aligned}
$$

A linear combination of all fundamental solution is a solution (due to linearity). Hence the general solution is

$$
u(t, x)=\sum_{n=1}^{\infty} C_{n} e^{-\left(\frac{2 n-1}{2}\right)^{2} t} \sin \left(\frac{2 n-1}{2} x\right)
$$

## 2 Problem 3.1.5

(a) Find the real eigensolutions to the damped heat equation $u_{t}=u_{x x}-u$. (b) Which solutions satisfy the periodic boundary conditions $u(t,-\pi)=u(t, \pi), u_{x}(t,-\pi)=u_{x}(t, \pi)$ ?
solution

### 2.1 Part (a)

Using separation of variables, Let $u(t, x)=T(t) X(x)$. Substituting this into $u_{t}+u=u_{x x}$ gives $T^{\prime} X+T X=T X^{\prime \prime}$. Dividing by $X T \neq 0$ gives

$$
\frac{T^{\prime}}{T}+1=\frac{X^{\prime \prime}}{X}=-\lambda
$$

Where $\lambda$ is the separation constant. This gives the following ODE's to solve

$$
\begin{aligned}
X^{\prime \prime}(x)+\lambda X(x) & =0 \\
T^{\prime}(t)+(\lambda+1) T(t) & =0
\end{aligned}
$$

Eigenfunctions are solutions to the spatial ODE.

$$
\begin{equation*}
X(x)=c_{1} e^{\sqrt{-\lambda} x}+c_{2} e^{-\sqrt{-\lambda} x} \tag{1}
\end{equation*}
$$

To determine the actual eigenfunctions and eigenvalues, boundary conditions are used. This is part b below.

### 2.2 Part (b)

Using $u(t,-\pi)=u(t, \pi), u_{x}(t,-\pi)=u_{x}(t, \pi)$. Starting with the spatial ODE above, and transferring the boundary condition to $X$ gives

$$
\begin{aligned}
X^{\prime \prime}(x)+\lambda X(x) & =0 \\
X(-\pi) & =X(\pi) \\
X^{\prime}(-\pi) & =X^{\prime}(\pi)
\end{aligned}
$$

This is an eigenvalue boundary value problem. The solution is

$$
\begin{equation*}
X(x)=c_{1} e^{\sqrt{-\lambda} x}+c_{2} e^{-\sqrt{-\lambda} x} \tag{1}
\end{equation*}
$$

case $\lambda<0$
Since $\lambda<0$, then $-\lambda$ is positive. Let $\mu=-\lambda$, where $\mu$ is now positive. The solution (1) becomes

$$
X(x)=c_{1} e^{\sqrt{\mu} x}+c_{2} e^{-\sqrt{\mu} x}
$$

The above can be written as

$$
\begin{equation*}
X(x)=c_{1} \cosh (\sqrt{\mu} x)+c_{2} \sinh (\sqrt{\mu} x) \tag{2}
\end{equation*}
$$

Applying first B.C. $X(-\pi)=X(\pi)$ using (2) gives

$$
\begin{aligned}
c_{1} \cosh (\sqrt{\mu} \pi)+c_{2} \sinh (-\sqrt{\mu} \pi) & =c_{1} \cosh (\sqrt{\mu} \pi)+c_{2} \sinh (\sqrt{\mu} \pi) \\
c_{2} \sinh (-\sqrt{\mu} \pi) & =c_{2} \sinh (\sqrt{\mu} \pi)
\end{aligned}
$$

But sinh is only zero when its argument is zero which is not the case here. Therefore the above implies that $c_{2}=0$ as only possibility to satisfy the above equation. The solution (2) now reduces to

$$
\begin{equation*}
X(x)=c_{1} \cosh (\sqrt{\mu} x) \tag{3}
\end{equation*}
$$

Taking derivative

$$
\begin{equation*}
X^{\prime}(x)=c_{1} \sqrt{\mu} \sinh (\sqrt{\mu} x) \tag{4}
\end{equation*}
$$

Applying the second $\mathrm{BC} X^{\prime}(-\pi)=X^{\prime}(\pi)$ gives

$$
c_{1} \sqrt{\mu} \sinh (-\sqrt{\mu} \pi)=c_{1} \sqrt{\mu} \sinh (\sqrt{\mu} x)
$$

But sinh is only zero when its argument is zero which is not the case here. Therefore the above implies that $c_{1}=0$. This means a trivial solution. Therefore $\lambda<0$ is not an eigenvalue.
case $\lambda=0$

In this case the solution is $X(x)=c_{1}+c_{2} x$. Applying first $\mathrm{BC} X(-\pi)=X(\pi)$ gives

$$
\begin{aligned}
c_{1}-c_{2} \pi & =c_{1}+c_{2} \pi \\
-c_{2} \pi & =c_{2} \pi
\end{aligned}
$$

This gives $c_{2}=0$. The solution now becomes

$$
X(x)=c_{1}
$$

Therefore $X^{\prime}(x)=0$. Applying the second boundary conditions $X^{\prime}(-\pi)=X^{\prime}(\pi)$ is now satisfied for any $c_{1}$, since it gives $(0=0)$. Therefore $\lambda=0$ is an eigenvalue with eigenfunction $X_{0}(0)=1$ (selecting $c_{1}=1$ since any arbitrary constant will work).
case $\lambda>0$
The solution in this case is

$$
\begin{aligned}
X(x) & =c_{1} e^{\sqrt{-\lambda} x}+c_{2} e^{-\sqrt{-\lambda} x} \\
& =c_{1} e^{i \sqrt{\lambda} x}+c_{2} e^{-i \sqrt{\lambda} x}
\end{aligned}
$$

Which can be rewritten as (the constants $c_{1}, c_{2}$ below will be different than the above $c_{1}, c_{2}$, but kept the same name for simplicity).

$$
\begin{equation*}
X(x)=c_{1} \cos (\sqrt{\lambda} x)+c_{2} \sin (\sqrt{\lambda} x) \tag{5}
\end{equation*}
$$

Applying first B.C. $X(-\pi)=X(\pi)$ using the above gives

$$
\begin{aligned}
c_{1} \cos (\sqrt{\lambda} \pi)+c_{2} \sin (-\sqrt{\lambda} \pi) & =c_{1} \cos (\sqrt{\lambda} \pi)+c_{2} \sin (\sqrt{\lambda} \pi) \\
c_{2} \sin (-\sqrt{\lambda} \pi) & =c_{2} \sin (\sqrt{\lambda} \pi)
\end{aligned}
$$

There are two choices here. Either $c_{2}=0$ or $\sqrt{\lambda} \pi=n \pi, n=1,2,3, \cdots$. Using the second choice for now, which implies that

$$
\lambda_{n}=n^{2} \quad n=1,2,3, \cdots
$$

And now we will now look to see what happens using the second BC with the above choice. The solution (5) now becomes

$$
X(x)=c_{1} \cos (n x)+c_{2} \sin (n x) \quad n=1,2,3, \cdots
$$

Therefore

$$
X^{\prime}(x)=-c_{1} n \sin (n x)+c_{2} n \cos (n x)
$$

Applying the second $\mathrm{BC} \mathrm{X}^{\prime}(-\pi)=X^{\prime}(\pi)$ using the above gives

$$
\begin{aligned}
c_{1} n \sin (n \pi)+c_{2} n \cos (n \pi) & =-c_{1} n \sin (n \pi)+c_{2} n \cos (n \pi) \\
c_{1} n \sin (n \pi) & =-c_{1} n \sin (n \pi) \\
0 & =0
\end{aligned}
$$

Since $n$ is integer.
Therefore this means that using the choice $\lambda_{n}=n^{2}$ satisfied both boundary conditions with $c_{2} \neq 0, c_{1} \neq 0$. This means the solution (5) is

$$
X_{n}(x)=A_{n} \cos (n x)+B_{n} \sin (n x) \quad n=1,2,3, \cdots
$$

The above says that there are two eigenfunctions in this case. They are

$$
X_{n}(x)=\left\{\begin{array}{l}
\cos (n x) \\
\sin (n x)
\end{array}\right.
$$

Recalling that there is also a zero eigenvalue with constant as its eigenfunction, then the complete set of eigenfunctions is

$$
X_{n}(x)=\left\{\begin{array}{c}
1 \\
\cos (n x) \\
\sin (n x)
\end{array}\right.
$$

Now that the eigenvalues are found, the solution to the time ODE can be found. The time ODE from above was found to be

$$
T^{\prime}(t)+(\lambda+1) T(t)=0
$$

For the zero eigenvalue case, the above reduces to $T^{\prime}(t)+T(t)=0$ which has the solution $T_{0}(t)=C_{0} e^{-t}$. For non zero eigenvalues $\lambda_{n}=n^{2}$, the ODE becomes $T^{\prime}(t)+\left(n^{2}+1\right) T(t)=0$, whose solution is $T_{0}(t)=C_{n} e^{-\left(n^{2}+1\right) t}$.

Putting all the above together, gives the fundamental solution as

$$
u_{n}(t, x)= \begin{cases}\multicolumn{2}{c}{C_{0} e^{-t}} \\ C_{n} \cos (n x) e^{-\left(n^{2}+1\right) t} & n=1,2,3, \cdots \\ B_{n} \sin (n x) e^{-\left(n^{2}+1\right) t} & n=1,2,3, \cdots\end{cases}
$$

The complete solution is the sum of the above solutions

$$
u(t, x)=C_{0} e^{-t}+\sum_{n=1}^{\infty} e^{-\left(n^{2}+1\right) t}\left(C_{n} \cos (n x)+B_{n} \sin (n x)\right)
$$

The constants $C_{0}, C_{n}, B_{n}$ can be found from initial conditions.

## 3 Problem 3.2.1

(d) Find the Fourier series of the following functions $f(x)=x^{2}$ (using $-\pi \leq x \leq \pi$ ) solution

The Fourier series is given by

$$
x^{2} \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos \left(\frac{2 \pi}{T} n x\right)+a_{n} \sin \left(\frac{2 \pi}{T} n x\right)
$$

Where $T$ is the period of $f(x)$. Taking this period to be $2 \pi$, the above simplifies to

$$
x^{2} \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos (n x)+b_{n} \sin (n x)
$$

The function $x^{2}$ is even, hence all $b_{n}$ are zero. The above becomes

$$
\begin{equation*}
x^{2} \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos (n x) \tag{1}
\end{equation*}
$$

But

$$
\begin{aligned}
a_{0} & =\frac{1}{\pi} \int_{-\pi}^{\pi} x^{2} d x \\
& =\frac{2}{\pi} \int_{0}^{\pi} x^{2} d x \\
& =\frac{2}{\pi}\left[\frac{x^{3}}{3}\right]_{0}^{\pi} \\
& =\frac{2}{3 \pi} \pi^{3} \\
& =\frac{2}{3} \pi^{2}
\end{aligned}
$$

And

$$
\begin{align*}
a_{n} & =\frac{1}{\pi} \int_{-\pi}^{\pi} x^{2} \cos (n x) d x \\
& =\frac{2}{\pi} \int_{0}^{\pi} x^{2} \cos (n x) d x \tag{1A}
\end{align*}
$$

Let $I=\int_{0}^{\pi} x^{2} \cos (n x) d x$. Using integration by parts $\int u d v=u v-\int v d u$. Let $u=x^{2}, d v=$ $\cos (n x)$. Then $d u=2 x, v=\frac{\sin (n x)}{n}$. Hence

$$
\begin{aligned}
I & =\left[x^{2} \frac{\sin (n x)}{n}\right]_{0}^{\pi}-2 \int_{0}^{\pi} x \frac{\sin (n x)}{n} d x \\
& =\overbrace{\frac{1}{n}\left[x^{2} \sin (n x)\right]_{0}^{\pi}}^{0}-\frac{2}{n} \int_{0}^{\pi} x \sin (n x) d x \\
& =-\frac{2}{n} \int_{0}^{\pi} x \sin (n x) d x
\end{aligned}
$$

Integration by parts again. $u=x, d v=\sin (n x)$, then $d u=1, v=-\frac{\cos (n x)}{n}$. The above becomes

$$
\begin{aligned}
I & =-\frac{2}{n}\left(\left[-x \frac{\cos (n x)}{n}\right]_{0}^{\pi}-\int_{0}^{\pi}-\frac{\cos (n x)}{n} d x\right) \\
& =-\frac{2}{n}\left(-\frac{1}{n}[x \cos (n x)]_{0}^{\pi}+\frac{1}{n} \int_{0}^{\pi} \cos (n x) d x\right) \\
& =\frac{2}{n^{2}}\left([x \cos (n x)]_{0}^{\pi}-\int_{0}^{\pi} \cos (n x) d x\right) \\
& =\frac{2}{n^{2}}\left([\pi \cos (n \pi)]-\left[\frac{\sin (n x)}{n}\right]_{0}^{\pi}\right) \\
& =\frac{2 \pi}{n^{2}} \cos (n \pi) \\
& =\frac{2 \pi}{n^{2}}(-1)^{n}
\end{aligned}
$$

The above is $I$. Substituting this result back in (1A) gives

$$
\begin{aligned}
a_{n} & =\frac{2}{\pi} I \\
& =\frac{2}{\pi} \frac{2 \pi}{n^{2}}(-1)^{n} \\
& =\frac{4}{n^{2}}(-1)^{n}
\end{aligned}
$$

Therefore (1) becomes

$$
x^{2} \sim \frac{1}{3} \pi^{2}+4 \sum_{n=1}^{\infty} \frac{(-1)^{n}}{n^{2}} \cos (n x)
$$

To verify this result, the Fourier series was compared to $x^{2}$ for an increasing number of terms to see if it converged to $x^{2}$. Here is the result. This shows the convergence is fast, after 6 terms only the approximation (in red color) is almost the same as the original function $x^{2}$.


Figure 1: Fourier series of $x^{2}$

```
\(\mathrm{fs}\left[x_{-}, \max _{-}\right]:=\frac{1}{3} \pi^{2}+4 \operatorname{Sum}\left[\frac{(-1)^{n}}{\mathrm{n}^{2}} \operatorname{Cos}[\mathrm{n} x],\{n, 1, \max \}\right]\)
makePlot \(\left[n_{-}\right]:=P l o t\left[\left\{x^{\wedge} 2, f s[x, n]\right\},\{x,-\operatorname{Pi}, \operatorname{Pi}\}\right.\),
            PlotStyle \(\rightarrow\) \{ Gray, Red, AxesLabel \(\rightarrow\{" x "\), None \(\}\),
            PlotLabel \(\rightarrow\) Row[\{"Fourier series approx using ", n, " terms"\}],
            ImageSize \(\rightarrow 300\)
        ];
Grid[Partition[Table[makePlot[n], \{n, \{1, 2, 3, 4, 5, 6\}\}], 2],
    Frame \(\rightarrow\) All]
```

Figure 2: Code used for the above plot
the following plot shows how the Fourier series approximation to $x^{2}$ when it is periodically extended to outside $[-\pi, \pi]$. This uses the range $[-3 \pi, 3 \pi]$ by adding one period to left and one period to the right.

```
\(\left.\ln [\cdot]=\mathrm{fs}\left[x_{-}, \max \right]\right]:=\frac{1}{3} \pi^{2}+4 \operatorname{Sum}\left[\frac{(-1)^{n}}{n^{2}} \operatorname{Cos}[n x],\{n, 1, \max \}\right]\)
fx[x_]:=Piecewise[\{
            \(\left\{(x+2 \mathrm{Pi})^{\wedge} 2, x<-\mathrm{Pi}\right\}\),
            \(\left\{x^{\wedge} 2,-\mathrm{Pi}<x<\mathrm{Pi}\right\}\),
            \(\left.\left\{(x-2 \mathrm{Pi})^{\wedge} 2, x>\operatorname{Pi}\right\}\right\} ;\)
makePlot \(\left[n_{-}\right]:=\operatorname{Plot}[\{f x[x], f s[x, n]\},\{x,-3 P i, 3 P i\}\),
            PlotStyle \(\rightarrow\) \{ Gray, Red \}, AxesLabel \(\rightarrow\) \{" x ", None \(\}\),
            PlotLabel \(\rightarrow\) Row[\{"Fourier series approx using ", \(n\), " terms"\}],
            ImageSize \(\rightarrow 300\)
            ];
Grid[Partition[Table[makePlot[n], \(\{n,\{1,2,3,4,5,6\}\}], 2]\),
    Frame \(\rightarrow\) All]
```

Figure 3: Code used for the above plot

## 4 Problem 3.2.2

(d) Find the Fourier series of the following function $f(x)=\left\{\begin{array}{cc}x & |x|<\frac{\pi}{2} \\ 0 & \text { otherwise }\end{array}\right.$ solution

This is plot showing $f(x)$


Figure 4: Plot of $f(x)$

The Fourier series is given by

$$
f(x) \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos \left(\frac{2 \pi}{T} n x\right)+a_{n} \sin \left(\frac{2 \pi}{T} n x\right)
$$

Where $T$ is the period of the function to be approximated. Taking this period to be $2 \pi$, the above simplifies to

$$
f(x) \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos (n x)+b_{n} \sin (n x)
$$

The function $f(x)$ is odd then all $a_{n}$ will zero. The above simplifies to

$$
f(x) \sim \sum_{n=1}^{\infty} b_{n} \sin (n x)
$$

Where

$$
\begin{aligned}
b_{n} & =\frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin (n x) d x \\
& =\frac{1}{\pi} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} x \sin (n x) d x
\end{aligned}
$$

But $x$ is odd and $\sin (x)$ is odd, hence the product is even. The above simplifies to

$$
b_{n}=\frac{2}{\pi} \int_{0}^{\frac{\pi}{2}} x \sin (n x) d x
$$

Using integration by parts $\int u d v=u v-\int v d u$. Let $x=u, d u=1, d v=\sin (n x), v=\frac{-\cos (n x)}{n}$, the above gives

$$
\begin{aligned}
b_{n} & =\frac{2}{\pi}\left(\frac{-1}{n}[x \cos (n x)]_{0}^{\frac{\pi}{2}}+\frac{1}{n} \int_{0}^{\frac{\pi}{2}} \cos (n x) d x\right) \\
& =\frac{2}{\pi n}\left(-[x \cos (n x)]_{0}^{\frac{\pi}{2}}+\int_{0}^{\frac{\pi}{2}} \cos (n x) d x\right) \\
& =\frac{2}{\pi n}\left(-\left[\frac{\pi}{2} \cos \left(n \frac{\pi}{2}\right)\right]+\frac{1}{n}[\sin (n x)]_{0}^{\frac{\pi}{2}}\right) \\
& =\frac{2}{\pi n}\left(-\left[\frac{\pi}{2} \cos \left(n \frac{\pi}{2}\right)\right]+\frac{1}{n}\left[\sin \left(n \frac{\pi}{2}\right)\right]\right) \\
& =\frac{2}{\pi n^{2}}\left(\sin \left(n \frac{\pi}{2}\right)-\frac{n \pi}{2} \cos \left(n \frac{\pi}{2}\right)\right)
\end{aligned}
$$

Therefore the Fourier series becomes

$$
f(x) \sim \sum_{n=1}^{\infty} \frac{2}{\pi n^{2}}\left(\sin \left(\frac{n \pi}{2}\right)-\frac{1}{2} n \pi \cos \left(\frac{n \pi}{2}\right)\right) \sin (n x)
$$

To verify this result, the Fourier series was compared to $f(x)$ for increasing number of terms to see if it converges to $x^{2}$. Here is the result. This shows the convergence is fast, but not as fast as last problem due to jump discontinuity in $f(x) .10$ terms are used below.


Figure 5: Fourier series approximation of $f(x)$

```
In[-]:= fs[\mp@subsup{x}{-}{\prime},\operatorname{max}_]:= Sum[\frac{2}{\mp@subsup{n}{}{2}\pi}(\operatorname{Sin}[\frac{n\pi}{2}]-\frac{1}{2}n\pi\operatorname{cos}[\frac{n\pi}{2}])\operatorname{Sin}[nx],{n,1,\operatorname{max}}];
    f[x_]:= Piecewise[{{x, Abs[x]<Pi/2},{0, True}}];
    makePlot[n_] := Plot[{f[x], fs[x, n]}, {x, -Pi, Pi},
            PlotStyle }->\mathrm{ {Blue, Red}, AxesLabel }->{"x", None}
            PlotLabel }->\mathrm{ Row[{"Fourier series approx using ", n, " terms"}],
            ImageSize }->\mathrm{ 300,
            Ticks }->\mathrm{ {Range[-Pi, Pi, Pi/ 2], Automatic}
            ];
Grid[Partition[Table[makePlot[n], {n, {1, 2, 3, 4, 5, 6, 7, 8, 9, 10}}], 2],
    Frame }->\mathrm{ All]
```

Figure 6: Code used for the above plot
the following plot shows how the Fourier series approximate $f(x)$ when it is periodically
extended to outside $[-\pi, \pi]$. This uses the range $[-3 \pi, 3 \pi]$ by adding one more period to left and to the right.


Figure 7: Fourier series of periodic extension $f(x)$

```
\(\ln [\cdot]=\mathrm{fs}\left[x_{-}, \max x_{-}\right]:=\operatorname{Sum}\left[\frac{2}{n^{2} \pi}\left(\operatorname{Sin}\left[\frac{n \pi}{2}\right]-\frac{1}{2} n \pi \operatorname{Cos}\left[\frac{n \pi}{2}\right]\right) \operatorname{Sin}[n x],\{n, 1, \max \}\right] ;\)
    \(\mathrm{f}\left[x_{-}\right]:=\)Piecewise[ \(\{\)
            \(\{0, x<-5 / 2 \mathrm{Pi}\}\),
            \(\{x+2 \mathrm{Pi},-5 / 2 \mathrm{Pi}<x<-3 / 2 \mathrm{Pi}\}\),
            \(\{0,-3 / 2 \mathrm{Pi}<x<-\mathrm{Pi} / 2\}\),
            \(\{x,-\mathrm{Pi} / 2<x<\mathrm{Pi} / 2\}\),
            \(\{0, \mathrm{Pi} / 2<x<3 / 2 \mathrm{Pi}\}\),
            \(\{x-2 \mathrm{Pi}, 2 / 3 \mathrm{Pi}<x<5 / 2 \mathrm{Pi}\}\),
            \(\{0,5 / 2 \mathrm{Pi}<x<3 \mathrm{Pi}\}\}] ;\)
makePlot[n_]:=Plot[\{f[x],fs[x,n]\},\{x,-3Pi,3Pi\},
            PlotStyle \(\rightarrow\) \{ Blue, Red\}, AxesLabel \(\rightarrow\) \{" \(x\) ", None \},
            PlotLabel \(\rightarrow\) Row [\{"Fourier series approx using ", \(n\), " terms"\}],
            ImageSize \(\rightarrow\) 300,
            Ticks \(\rightarrow\) \{Range[-Pi, Pi, Pi/2], Automatic \(\}\)
        ];
Grid [Partition[Table[makePlot[n], \(\{n,\{1,2,3,4,5,6,7,8,9,10\}\}], 2]\),
    Frame \(\rightarrow\) All]
```

Figure 8: Code used for the above plot

## 5 Problem 3.2.3

Find the Fourier series of $\sin ^{2} x$ and $\cos ^{2} x$ without directly calculating the Fourier coefficients.
solution
Using the known trig identity

$$
\begin{equation*}
\sin ^{2} x=\frac{1}{2}-\frac{1}{2} \cos (2 x) \tag{1}
\end{equation*}
$$

And comparing the the above to the Fourier series expansion
$\sin ^{2} x=\frac{a_{0}}{2}+\left(a_{1} \cos (x)+a_{2} \cos (2 x)+a_{3} \cos (3 x)+\cdots\right)+\left(b_{1} \sin (x)+b_{2} \sin (2 x)+b_{3} \sin (3 x)+\cdots\right)$

Shows that $\frac{a_{0}}{2}=\frac{1}{2}$ and $a_{2}=\frac{-1}{2}$ and all other terms are zero. Because the Fourier series is unique for a function, then (1) is the Fourier series for $\sin ^{2} x$.

Similarly, Using the known trig identity

$$
\begin{equation*}
\cos ^{2} x=\frac{1}{2}+\frac{1}{2} \cos (2 x) \tag{2}
\end{equation*}
$$

And comparing the the above to the Fourier series expansion (A), shows that $\frac{a_{0}}{2}=\frac{1}{2}$ and $a_{2}=\frac{1}{2}$ and all other terms are zero. Therefore (2) is the Fourier series expansion for $\cos ^{2} x$.

## 6 Problem 3.2.6

Graph the $2 \pi$ periodic extension of each of the following functions (h) $f(x)=\frac{1}{x}$. Which extension are continuous? Differentiable?
solution

### 6.1 Part (h)

The original function $f(x)=\frac{1}{x}$ is always taken from $-\pi \leq x \leq \pi$ (before extending it periodically). At $x=0$ the function is not defined.


Figure 9: Plot of $f(x)=\frac{1}{x}$

Periodically extending it, it becomes (showing one extra period to the left and right) then following


Figure 10: Plot of periodic extension of $f(x)=\frac{1}{x}$

$$
\ln [\cdot]=\mathbf{f}\left[x_{-}\right]:=\text {Piecewise }[\{
$$

$\{1 /(x+2 \mathrm{Pi}), x<-\mathrm{Pi}\}$,
$\{1 / x,-\mathrm{Pi}<x<\mathrm{Pi}\}$,
$\{1 /(x-2 \mathrm{Pi}), \mathrm{Pi}<x\}$
\}];
Plot $[f[x],\{x,-3$ Pi, 3 Pi $\}$, Ticks $\rightarrow\{$ Range [-3Pi, 3 Pi, Pi], Automatic $\}$,
AxesLabel $\rightarrow\{$ " $x$ ", " $1 / x$ extended" $\}$,
GridLines $\rightarrow$ \{Range[-3 Pi, $3 \mathrm{Pi}, \mathrm{Pi}]$, Automatic $\}$,
GridLinesStyle $\rightarrow$ LightGray, PlotStyle $\rightarrow$ Red, AspectRatio $\rightarrow$ Automatic]
Figure 11: Code for the above plot

Looking at the above plot shows the extension is not continuous and also not Differentiable due to jump discontinuities.

## 7 Problem 3.2.9

Suppose that $f(x)$ is periodic with period $T$ (using $T$ instead of $l$ as in book as it is more clear). Prove that for any $a$ (a) $\int_{a}^{a+T} f(x) d x=\int_{0}^{T} f(x) d x$. (b) $\int_{0}^{T} f(x+a) d x=\int_{0}^{T} f(x) d x$ solution

### 7.1 Part (a)

$$
\int_{a}^{a+T} f(x) d x-\int_{0}^{T} f(x) d x=\overbrace{\left(\int_{a}^{T} f(x) d x+\int_{T}^{a+T} f(x) d x\right)}^{\int_{a}^{a+T} f(x) d x}-\overbrace{\left(\int_{0}^{a} f(x) d x+\int_{a}^{T} f(x) d x\right)}^{\int_{0}^{T} f(x) d x}
$$

Simplifying the RHS above gives

$$
\begin{equation*}
\int_{a}^{a+T} f(x) d x-\int_{0}^{T} f(x) d x=\int_{T}^{a+T} f(x) d x-\int_{0}^{a} f(x) d x \tag{1}
\end{equation*}
$$

But

$$
\begin{equation*}
\int_{T}^{a+T} f(x) d x=\int_{0}^{a} f(x+T) d x \tag{2}
\end{equation*}
$$

To show how Eq(2) was derived: Let $u=x-T$. Then $d u=d x$. When $x=T$ then $u=0$. When $x=a+T$ then $u=a$. Hence $\int_{T}^{a+T} f(x) d x=\int_{0}^{a} f(u+T) d u$. But $u$ is arbitrary integral variable. Renaming it back to $x$ gives that $\int_{T}^{a+T} f(x) d x=\int_{0}^{a} f(x+T) d x$.

Now, substituting (2) back into RHS of (1) gives

$$
\begin{aligned}
\int_{a}^{a+T} f(x) d x-\int_{0}^{T} f(x) d x & =\int_{0}^{a} f(x+T) d x-\int_{0}^{a} f(x) d x \\
& =\int_{0}^{a} f(x+T)-f(x) d x
\end{aligned}
$$

But since $f(x)$ is periodic, then $f(x+T)=f(x)$. Therefore the RHS above is zero.

$$
\begin{aligned}
\int_{a}^{a+T} f(x) d x-\int_{0}^{T} f(x) d x & =0 \\
\int_{a}^{a+T} f(x) d x & =\int_{0}^{T} f(x) d x
\end{aligned}
$$

Which is what the problem is asking to show.

### 7.2 Part (b)

Starting by rewriting $\int_{0}^{T} f(x+a) d x$ as the following. Let $u=x+a$. Hence $d u=d x$. When $x=0, u=a$ and when $x=T, u=a+T$. The integral becomes $\int_{a}^{a+T} f(u) d u$. But now $u$ is arbitrary integration variable. Renaming is back to $x$ then we obtain that

$$
\begin{equation*}
\int_{0}^{T} f(x+a) d x=\int_{a}^{a+T} f(x) d x \tag{1}
\end{equation*}
$$

Now, to show that main result, considering

$$
\int_{0}^{T} f(x+a) d x-\int_{0}^{T} f(x) d x=\int_{a}^{a+T} f(x) d x-\int_{0}^{T} f(x) d x
$$

Where in the above, (1) was used to obtain RHS. The above can now be written as

$$
\int_{0}^{T} f(x+a) d x-\int_{0}^{T} f(x) d x=\overbrace{\left(\int_{a}^{T} f(x) d x+\int_{T}^{T+a} f(x) d x\right)}^{\int_{a}^{a+T} f(x) d x}-\int_{0}^{T} f(x) d x
$$

But $\int_{T}^{T+a} f(x) d x=\int_{0}^{a} f(x) d x$ since $f(x)$ is periodic with period $T$. The above now becomes

$$
\begin{aligned}
\int_{0}^{T} f(x+a) d x-\int_{0}^{T} f(x) d x & =\left(\int_{a}^{T} f(x) d x+\int_{0}^{a} f(x) d x\right)-\int_{0}^{T} f(x) d x \\
& =\int_{0}^{T} f(x) d x-\int_{0}^{T} f(x) d x \\
& =0
\end{aligned}
$$

Therefore $\int_{0}^{T} f(x+a) d x=\int_{0}^{T} f(x) d x$ which is what the problem is asking to show.

## 8 Problem 3.2.25

(a) Sketch the $2 \pi$ periodic half-wave $f(x)=\left\{\begin{array}{cc}\sin x & 0<x \leq \pi \\ 0 & -\pi \leq x<0\end{array}\right.$. (b) Find its Fourier series. (c) Graph the first five Fourier sums and compare the function. (d) Discuss convergence of the Fourier series.

## solution

### 8.1 Part (a)



Figure 12: Plot of $f(x)$

```
ln[v]:= f[x_] := Piecewise[{{Sin[x], 0<x\leq Pi}, {0,-Pi\leqx<0}}];
Plot[f[x], {x, -Pi, Pi}, Ticks }->\mathrm{ {Range[-Pi, Pi, Pi / 2], Automatic},
    AxesLabel }->\mathrm{ {"x", "f(x)"},
    GridLines }->\mathrm{ {Range [-Pi, Pi, Pi / 2], Automatic},
    GridLinesStyle }->\mathrm{ LightGray, PlotStyle }->\mathrm{ Red]
```

Figure 13: Code for the above plot

### 8.2 Part (b)

The Fourier series is given by

$$
f(x) \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos \left(\frac{2 \pi}{T} n x\right)+a_{n} \sin \left(\frac{2 \pi}{T} n x\right)
$$

Where $T$ is the period of the function to be approximated. Taking this period to be $2 \pi$, the above simplifies to

$$
f(x) \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos (n x)+b_{n} \sin (n x)
$$

## Hence

$$
\begin{aligned}
a_{0} & =\frac{1}{\pi} \int_{-\pi}^{\pi} f(x) d x \\
& =\frac{1}{\pi} \int_{0}^{\pi} \sin (x) d x \\
& =\frac{1}{\pi}[-\cos (x)]_{0}^{\pi} \\
& =\frac{-1}{\pi}[\cos (x)]_{0}^{\pi} \\
& =\frac{-1}{\pi}[\cos (\pi)-1] \\
& =\frac{-1}{\pi}[-1-1] \\
& =\frac{2}{\pi}
\end{aligned}
$$

And

$$
\begin{aligned}
a_{n} & =\frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos (n x) d x \\
& =\frac{1}{\pi} \int_{0}^{\pi} \sin (x) \cos (n x) d x
\end{aligned}
$$

For $n=1$

$$
\begin{aligned}
a_{n} & =\frac{1}{\pi} \int_{0}^{\pi} \sin (x) \cos (x) d x \\
& =0
\end{aligned}
$$

And for $n>1$

$$
a_{n}=\frac{1}{\pi} \int_{0}^{\pi} \sin (x) \cos (n x) d x
$$

Using $\sin A \cos B=\frac{1}{2}(\sin (A-B)+\sin (A+b))$, then $\sin (x) \cos (n x)=\frac{1}{2}(\sin (x-n x)+\sin (x+n x))$. The above becomes

$$
\begin{aligned}
a_{n} & =\frac{1}{2 \pi} \int_{0}^{\pi} \sin (x-n x)+\sin (x+n x) d x \\
& =\frac{1}{2 \pi}\left(\int_{0}^{\pi} \sin (x-n x) d x+\int_{0}^{\pi} \sin (x+n x) d x\right) \\
& =\frac{1}{2 \pi}\left(-\frac{1}{1-n}[\cos (x-n x)]_{0}^{\pi}-\frac{1}{1+n}[\cos (x+n x)]_{0}^{\pi}\right) \\
& =\frac{-1}{2 \pi}\left(\frac{1}{1-n}[\cos (\pi-n \pi)-1]+\frac{1}{1+n}[\cos (\pi-n \pi)-1]\right)
\end{aligned}
$$

But $\cos (\pi-n \pi)=-\cos (n \pi)$. The above becomes

$$
\begin{aligned}
a_{n} & =\frac{-1}{2 \pi}\left(\frac{1}{1-n}[-\cos (n \pi)-1]+\frac{1}{1+n}[-\cos (n \pi)-1]\right) \\
& =\frac{1}{2 \pi}\left(\frac{\cos (n \pi)+1}{1-n}+\frac{\cos (n \pi)+1}{1+n}\right) \\
& =\frac{1}{2 \pi}\left(\frac{(1+n)(\cos (n \pi)+1)+(1-n)(\cos (n \pi)+1)}{(1-n)(1+n)}\right) \\
& =\frac{1}{2 \pi}\left(\frac{(1+n)(\cos (n \pi)+1)+(1-n)(\cos (n \pi)+1)}{\left(1-n^{2}\right)}\right) \\
& =\frac{1}{2 \pi\left(1-n^{2}\right)}((1+n)(\cos (n \pi)+1)+(1-n)(\cos (n \pi)+1)) \\
& =\frac{1}{2 \pi\left(1-n^{2}\right)}(2 \cos (\pi n)+2) \\
& =\frac{1}{\pi\left(1-n^{2}\right)}(\cos (\pi n)+1) \\
& =\frac{1+(-1)^{n}}{\pi\left(1-n^{2}\right)}
\end{aligned}
$$

For odd $n=3,5, \cdots$ then $a_{n}=0$. For even $n$ the above can be written as

$$
a_{n}=\frac{2}{\pi\left(1-n^{2}\right)} \quad n=2,4,6, \cdots
$$

Now $b_{n}$ is found

$$
\begin{aligned}
b_{n} & =\frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin (n x) d x \\
& =\frac{1}{\pi} \int_{0}^{\pi} \sin (x) \sin (n x) d x
\end{aligned}
$$

Consider case $n=1$ first. The above gives

$$
\begin{aligned}
b_{1} & =\frac{1}{\pi} \int_{0}^{\pi} \sin ^{2}(x) d x \\
& =\frac{1}{\pi} \int_{0}^{\pi} \frac{1}{2}-\frac{1}{2} \cos (2 x) d x \\
& =\frac{1}{\pi}\left(\int_{0}^{\pi} \frac{1}{2} d x-\frac{1}{2} \int_{0}^{\pi} \cos (2 x) d x\right) \\
& =\frac{1}{\pi}\left(\frac{1}{2} \pi-\frac{1}{2}\left[\frac{\sin (2 x)}{2}\right]_{0}^{\pi}\right) \\
& =\frac{1}{2}
\end{aligned}
$$

For $n>1$

$$
\begin{aligned}
b_{n} & =\frac{1}{\pi} \int_{0}^{\pi} \sin x \sin (n x) d x \\
& =\frac{1}{\pi} \frac{\sin (n \pi)}{n^{2}-1} \\
& =0
\end{aligned}
$$

Therefore the Fourier series is

$$
\begin{aligned}
f(x) & \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos (n x)+b_{n} \sin (n x) \\
& =\frac{1}{\pi}+\frac{1}{2} \sin (x)+\frac{2}{\pi} \sum_{n=2,4,6, \cdots}^{\infty} \frac{1}{1-n^{2}} \cos (n x) \\
& =\frac{1}{\pi}+\frac{1}{2} \sin (x)+\frac{2}{\pi} \sum_{n=1}^{\infty} \frac{1}{1-(2 n)^{2}} \cos (2 n x)
\end{aligned}
$$

### 8.3 Part (c)



Figure 14: Plot of Fourier series approximation and $f(x)$

```
In[0]:= fs[\mp@subsup{x}{-}{\prime},\operatorname{max}]]:=\frac{1}{\pi}+\frac{1}{2}\operatorname{Sin}[x]+\frac{2}{\pi}\operatorname{Sum}[\frac{1}{1-(2n\mp@subsup{)}{}{2}}\operatorname{Cos[2nx],{n,1, max}];}
f[x_] := Piecewise[{{Sin[x],0<x\leq Pi}, {0, - Pi \leq x < 0}}];
makePlot[n_] := Plot[{f[x], fs[x, n]}, {x, -Pi, Pi},
            PlotStyle }->\mathrm{ { Blue, Red}, AxesLabel }->\mathrm{ {"x", None},
            PlotLabel }->\mathrm{ Row[{"Fourier series approx using ", n, " terms"}],
            ImageSize }->\mathrm{ 300,
            Ticks }->\mathrm{ {Range[-Pi, Pi, Pi / 2], Automatic}
        ];
Grid[Partition[Table[makePlot[n], {n, 0, 5}], 2],
    Frame }->\mathrm{ All]
```

Figure 15: Code for the above plot

### 8.4 Part (d)

The function $f(x)$ is piecewise $C^{1}$ continuous over $-\pi \leq x \leq \pi$. Therefore the $2 \pi$ periodic extension is also piecewise $C^{1}$ continuous over all $x$. There are no jump discontinues (only corner points). The Fourier series converges to $f(x)$ at each $x \in \mathfrak{R}$. (If there was a jump discontinuity at some $x$, then the Fourier series will converge to the average of $f(x)$ at that $x$, but this is not the case here).

## $9 \quad$ Problem 3.2.27

(a) Find the Fourier series of $f(x)=e^{x}$. (b) For which values of $x$ does the Fourier series converges? Is the convergence uniform? (c) Graph the function it converges to.

## solution

### 9.1 Part (a)

For generality, the Fourier series for $e^{a x}$ is found, then at the end $a$ is set to be one. It is assumed the period is $2 \pi$.

$$
e^{a x} \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos \left(\frac{2 \pi}{T} n x\right)+b_{n} \sin \left(\frac{2 \pi}{T} n x\right)
$$

But $T=2 \pi$ and the above becomes

$$
e^{a x} \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos (n x)+b_{n} \sin (n x)
$$

Where

$$
\begin{aligned}
a_{0} & =\frac{1}{\frac{T}{2}} \int_{-\frac{T}{2}}^{\frac{T}{2}} f(x) d x \\
& =\frac{1}{\pi} \int_{-\pi}^{\pi} e^{a x} d x \\
& =\frac{1}{\pi}\left[\frac{e^{a x}}{a}\right]_{-\pi}^{\pi} \\
& =\frac{1}{\pi a}\left(e^{a \pi}-e^{-a \pi}\right)
\end{aligned}
$$

But $\frac{e^{a \pi}-e^{-a \pi}}{2}=\sinh (a \pi)$ hence the above simplifies to

$$
a_{0}=\frac{2}{\pi a} \sinh (a \pi)
$$

And for $n>0$

$$
\begin{align*}
a_{n} & =\frac{1}{\frac{T}{2}} \int_{-\frac{T}{2}}^{\frac{T}{2}} f(x) \cos \left(\frac{2 \pi}{T} n x\right) d x \\
& =\frac{1}{\pi} \int_{-\pi}^{\pi} e^{a x} \cos (n x) d x \tag{1}
\end{align*}
$$

Let $I=\int_{-\pi}^{\pi} e^{a x} \cos (n x) d x$. Using integration by parts, $\int u d v=u v-\int v d u$. Let $u=\cos n x, d v=$ $e^{a x}$ then $v=\frac{e^{a x}}{a}, d u=-n \sin (n x)$. Hence

$$
\begin{aligned}
I & =u v-\int v d u \\
& =\left[\cos (n x) \frac{e^{a x}}{a}\right]_{-\pi}^{\pi}+\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x \\
& =\left[\cos (n \pi) \frac{e^{a \pi}}{a}-\cos (n \pi) \frac{e^{-a \pi}}{a}\right]+\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x \\
& =(-1)^{n}\left[\frac{e^{a \pi}-e^{-a \pi}}{a}\right]+\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x \\
& =\frac{2(-1)^{n}}{a}\left[\frac{e^{a \pi}-e^{-a \pi}}{2}\right]+\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x \\
& =\frac{2(-1)^{n}}{a} \sinh (a \pi)+\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x
\end{aligned}
$$

Applying integration by parts again on the integral above. Let $u=\sin n x, d v=e^{a x}$ then $v=\frac{e^{a x}}{a}, d u=n \cos (n x)$ and the above becomes

$$
\begin{aligned}
I & =\frac{2(-1)^{n}}{a} \sinh (a \pi)+\frac{n}{a}\left(\left(\sin n x \frac{e^{a x}}{a}\right)_{-\pi}^{\pi}-\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \cos (n x) d x\right) \\
& =\frac{2(-1)^{n}}{a} \sinh (a \pi)+\frac{n}{a}(\frac{1}{a} \overbrace{\left(\sin (n \pi) e^{a \pi}+\sin (n \pi) e^{-a \pi}\right)}^{a}-\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \cos (n x) d x) \\
& =\frac{2(-1)^{n}}{a} \sinh (a \pi)-\frac{n^{2}}{a^{2}} \int_{-\pi}^{\pi} e^{a x} \cos (n x) d x
\end{aligned}
$$

But $\int_{-\pi}^{\pi} e^{a x} \cos (n x) d x=I$, the original integral we are solving for. Hence solving for $I$ from the above gives gives

$$
\begin{align*}
I & =\frac{2(-1)^{n}}{a} \sinh (a \pi)-\frac{n^{2}}{a^{2}} I \\
I+\frac{n^{2}}{a^{2}} I & =\frac{2(-1)^{n}}{a} \sinh (a \pi) \\
I\left(1+\frac{n^{2}}{a^{2}}\right) & =\frac{2(-1)^{n}}{a} \sinh (a \pi) \\
I & =\frac{\frac{2(-1)^{n}}{a} \sinh (a \pi)}{1+\frac{n^{2}}{a^{2}}} \\
& =\frac{2 a(-1)^{n} \sinh (a \pi)}{a^{2}+n^{2}} \tag{2}
\end{align*}
$$

Using (2) in (1) gives

$$
\begin{align*}
a_{n} & =\frac{1}{\pi} \int_{-\pi}^{\pi} e^{a x} \cos (n x) d x \\
& =\frac{a}{\pi} \frac{2(-1)^{n} \sinh (a \pi)}{a^{2}+n^{2}} \tag{3}
\end{align*}
$$

Now we will do the same to find $b_{n}$

$$
\begin{align*}
b_{n} & =\frac{1}{\frac{T}{2}} \int_{-\frac{T}{2}}^{\frac{T}{2}} f(x) \sin \left(\frac{2 \pi}{T} n x\right) d x \\
& =\frac{1}{\pi} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x \tag{4}
\end{align*}
$$

Let $I=\int_{-\pi}^{\pi} e^{a x} \sin (n x) d x$. Using integration by parts, $\int u d v=u v-\int v d u$. Let $u=\sin (n x)$, $d v=$ $e^{a x}$ then $v=\frac{e^{a x}}{a}, d u=n \cos (n x)$. Hence

$$
\begin{aligned}
I & =u v-\int v d u \\
& =\left[\sin (n x) \frac{e^{a x}}{a}\right]_{-\pi}^{\pi}-\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \cos (n x) d x \\
& =\overbrace{\left[\sin (n \pi) \frac{e^{a \pi}}{a}-\sin (n \pi) \frac{e^{-a \pi}}{a}\right]}^{0}-\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \cos (n x) d x \\
& =-\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \cos (n x) d x
\end{aligned}
$$

Now we apply integration by parts again on the integral above. Let $u=\cos n x, d v=e^{a x}$ then $v=\frac{e^{a x}}{a}, d u=-n \sin (n x)$ and the above becomes

$$
\begin{aligned}
I & =-\frac{n}{a}\left(\left(\cos (n x) \frac{e^{a x}}{a}\right)_{-\pi}^{\pi}+\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x\right) \\
& =-\frac{n}{a}\left(\frac{1}{a}\left(\cos (n \pi) e^{a \pi}-\cos (n \pi) e^{-a \pi}\right)+\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x\right) \\
& =-\frac{n}{a}\left(\frac{1}{a} \cos (n \pi)\left(e^{a \pi}-e^{-a \pi}\right)+\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x\right) \\
& =-\frac{n}{a}\left(\frac{2}{a} \cos (n \pi)\left(\frac{e^{a \pi}-e^{-a \pi}}{2}\right)+\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x\right) \\
& =-\frac{n}{a}\left(\frac{2}{a} \cos (n \pi) \sinh (a \pi)+\frac{n}{a} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x\right) \\
& =-\frac{2 n}{a^{2}}(-1)^{n} \sinh (a \pi)-\frac{n^{2}}{a^{2}} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x
\end{aligned}
$$

But $\int_{-\pi}^{\pi} e^{a x} \sin (n x) d x=I$. Hence solving for $I$ gives

$$
\begin{align*}
I & =-\frac{2 n}{a^{2}}(-1)^{n} \sinh (a \pi)-\frac{n^{2}}{a^{2}} I \\
I+\frac{n^{2}}{a^{2}} I & =-\frac{2 n}{a^{2}}(-1)^{n} \sinh (a \pi) \\
I\left(1+\frac{n^{2}}{a^{2}}\right) & =-\frac{2 n}{a^{2}}(-1)^{n} \sinh (a \pi) \\
I & =-\frac{\frac{2 n}{a^{2}}(-1)^{n} \sinh (a \pi)}{1+\frac{n^{2}}{a^{2}}} \\
I & =-\frac{2 n(-1)^{n}}{a^{2}+n^{2}} \sinh (a \pi) \tag{5}
\end{align*}
$$

Using (5) in (4) gives

$$
\begin{aligned}
b_{n} & =\frac{1}{\pi} \int_{-\pi}^{\pi} e^{a x} \sin (n x) d x \\
& =-\frac{1}{\pi} \frac{2 n(-1)^{n}}{a^{2}+n^{2}} \sinh (a \pi)
\end{aligned}
$$

Now that we found $a_{0}, a_{n}, b_{n}$ then the Fourier series is

$$
\begin{aligned}
e^{a x} & \sim \frac{a_{0}}{2}+\sum_{n=1}^{\infty} a_{n} \cos (n x)+b_{n} \sin (n x) \\
& \sim \frac{\frac{2}{\pi a} \sinh (a \pi)}{2}+\sum_{n=1}^{\infty} \frac{a}{\pi} \frac{2(-1)^{n} \sinh (a \pi)}{a^{2}+n^{2}} \cos (n x)-\frac{1}{\pi} \frac{2 n(-1)^{n}}{a^{2}+n^{2}} \sinh (a \pi) \sin (n x) \\
& \sim \frac{\sinh (a \pi)}{\pi a}+\frac{1}{\pi} \sinh (a \pi) \sum_{n=1}^{\infty} \frac{2(-1)^{n}}{a^{2}+n^{2}}(a \cos (n x)-n \sin (n x)) \\
& \sim \sinh (a \pi)\left(\frac{1}{\pi a}+\frac{1}{\pi} \sum_{n=1}^{\infty} \frac{2(-1)^{n}}{a^{2}+n^{2}}(a \cos (n x)-n \sin (n x))\right) \\
& \sim \frac{2 \sinh (a \pi)}{\pi}\left(\frac{1}{2 a}+\sum_{n=1}^{\infty} \frac{(-1)^{n}}{a^{2}+n^{2}}(a \cos (n x)-n \sin (n x))\right)
\end{aligned}
$$

When $a=1$ the above becomes

$$
e^{x} \sim \frac{2 \sinh (\pi)}{\pi}\left(\frac{1}{2}+\sum_{n=1}^{\infty} \frac{(-1)^{n}}{1+n^{2}}(\cos (n x)-n \sin (n x))\right)
$$

### 9.2 Part (b)

The $2 \pi$ periodic extended function shows it piecewise $C^{1}$ over all points except at the points $x=\cdots,-5 \pi,-3 \pi, \pi, 3 \pi, 5 \pi, \cdots$. These are points at the ends of the original domain. At these points, there is a jump discontinuity. Therefore the Fourier series at these points will converge to the average of the $2 \pi$ periodic extended $e^{x}$. Due to the jump discontinuity Gibbs phenomena shows up at these points. This also means that the convergence is
not uniform.

### 9.3 Part (c)

The following is a plot showing the convergence using different number of terms in the above sum. This shows the Fourier series converges to $e^{x}$ at all points inside the interval, except at the end points $x=-\pi, \pi$ where it converges to the average of $f(x)$.


Figure 16: Plot of Fourier series approximation and $f(x)$

```
In[v]:= padIt2[v_, f_List] := AccountingForm[v,f, NumberSigns }->\mathrm{ {"", ""},
    NumberPadding }->\mathrm{ {" ", " "}, SignPadding }->\mathrm{ True];
fs[\mp@subsup{x}{-}{\prime},\operatorname{max}_]:=\frac{2\operatorname{Sinh[Pi]}}{\textrm{Pi}}(\frac{1}{2}+\operatorname{Sum}[\frac{(-1\mp@subsup{)}{}{n}}{1+\mp@subsup{n}{}{2}}(\operatorname{Cos}[nx]-n\operatorname{Sin}[nx]),{n,1,\operatorname{max}}]);
f[x_] := Exp[x];
fp[x_] := Piecewise[{{f[x+2 Pi], x < - Pi}, {f[x], - Pi<x < Pi}, {f[x-2Pi], x > Pi}}];
makePlot[n_] := Plot[{f[x], fs[x, n]}, {x, -Pi, Pi},
    PlotStyle }->\mathrm{ { Blue, Red}, AxesLabel }->\mathrm{ {"x", None},
    PlotLabel }->\mathrm{ Row[{"Fourier series approx using ", n, " terms"}],
    ImageSize }->\mathrm{ 300,
    Ticks }->\mathrm{ {Range[-Pi, Pi, Pi], Automatic},
    PlotRange }->\mathrm{ {{-1.1 Pi, 1.1 Pi}, {-4, 25}},
    GridLines }->\mathrm{ {Range [-Pi, Pi, Pi], Automatic}, GridLinesStyle }->\mathrm{ LightGray
        ];
    Grid[Partition[Table[makePlot[n] , {n, {0, 3, 6, 9, 12, 15}}], 2],
    Frame }->\mathrm{ All]
```

Figure 17: Code for the above plot

## 10 Problem 3.2.30

Suppose $a_{k}, b_{k}$ are the Fourier coefficients of the function $f(x)$. (a) To which function does the Fourier series

$$
\frac{a_{0}}{2}+\sum_{k=1}^{\infty} a_{k} \cos (2 k x)+b_{k} \sin (2 k x)
$$

Converge? (b) Test your answer with the Fourier series (3.37) for $f(x)=x$.

$$
\begin{equation*}
x \sim 2\left(\sin x-\frac{\sin 2 x}{2}+\frac{\sin 3 x}{3}-\frac{\sin 4 x}{4}+\cdots\right) \tag{3.37}
\end{equation*}
$$

solution

### 10.1 Part (a)

Let

$$
\begin{aligned}
& g(x) \sim \frac{a_{0}}{2}+\sum_{k=1}^{\infty} a_{k} \cos (2 k x)+b_{k} \sin (2 k x) \\
& f(x) \sim \frac{a_{0}}{2}+\sum_{k=1}^{\infty} a_{k} \cos (k x)+b_{k} \sin (k x)
\end{aligned}
$$

Then $g(x)$ has as its period half the period of $f(x)$. This is because when $2 k x=\frac{2 \pi}{T} k x$ then $T=\pi$ and when $k x=\frac{2 \pi}{T} k x$ then $T=2 \pi$.
Therefore, if $f(x)$ has fundamental period as $-\pi<x<\pi$, then $g(x)$ has a fundamental period as $-\frac{\pi}{2}<x<\frac{\pi}{2}$. And since $f(x), g(x)$ have the same Fourier series coefficients, then $g(x)$ converges to $2 f(x)$ but only over $-\frac{\pi}{2}<x<\frac{\pi}{2}$.

### 10.2 Part (b)

Let $f(x)=x$ whose we are given that its Fourier series is

$$
\begin{aligned}
f(x) & \sim 2\left(\sin x-\frac{\sin 2 x}{2}+\frac{\sin 3 x}{3}-\frac{\sin 4 x}{4}+\cdots\right) \\
& =2 \sin x-\sin 2 x+\frac{2}{3} \sin 3 x-\frac{1}{2} \sin 4 x+\cdots
\end{aligned}
$$

The above says that $a_{k}=0$ and $b_{k}=\frac{2(-1)^{k+1}}{k}$. Hence

$$
f(x) \sim \sum_{k=1}^{\infty} \frac{2(-1)^{k+1}}{k} \sin (k x)
$$

Therefore $g(x)$ will converge to

$$
\begin{aligned}
g(x) & \sim \frac{a_{0}}{2}+\sum_{k=1}^{\infty} a_{k} \cos (2 k x)+b_{k} \sin (2 k x) \\
& =\sum_{k=1}^{\infty} \frac{2(-1)^{k+1}}{k} \sin (2 k x) \\
& =2(+1) \sin (2 x)+\frac{-2}{2} \sin (4 x)+\frac{2(+1)}{3} \sin (6 x)+\frac{-2}{4} \sin (8 x)+\cdots \\
& =2 \sin (2 x)-\sin (4 x)+\frac{2}{3} \sin (6 x)-\frac{1}{2} \sin (8 x)+\cdots
\end{aligned}
$$

Over $-\frac{\pi}{2}<x<\frac{\pi}{2}$. To verify the above, we will now find $a_{k}, b_{k}$ directly for $x$ but using $T=\pi$ and not $T=2 \pi$ to see if the above Fourier series is obtained.

$$
\begin{aligned}
a_{0} & =\frac{2}{\pi} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} x d x \\
& =0
\end{aligned}
$$

And

$$
a_{k}=\frac{2}{\pi} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} x \cos (2 k x) d x
$$

Since $x$ is odd function and cos is even, the product is odd. Hence $a_{k}=0$.

$$
\begin{aligned}
b_{k} & =\frac{2}{\pi} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} x \sin (2 k x) d x \\
& =\frac{4}{\pi} \int_{0}^{\frac{\pi}{2}} x \sin (2 k x) d x \\
& =\frac{4}{\pi}\left(\frac{-k \pi \cos (k \pi)+\sin (k \pi)}{4 k^{2}}\right) \\
& =\frac{1}{\pi k^{2}}(-k \pi \cos (k \pi)) \\
& =\frac{-1}{k} \cos (k \pi) \\
& =\frac{-1}{k}(-1)^{k} \\
& =\frac{(-1)^{k+1}}{k}
\end{aligned}
$$

Therefore

$$
\begin{aligned}
g(x) & \sim \frac{a_{0}}{2}+\sum_{k=1}^{\infty} a_{k} \cos (2 k x)+b_{k} \sin (2 k x) \\
& =\sum_{k=1}^{\infty} \frac{(-1)^{k}}{k} \sin (2 k x)
\end{aligned}
$$

