

HW # 7

MAE 170

Winter 2005

49.5/50

by NASSER ABBASI

HW#7

Problem B-8-2

Consider system whose closed loop TF is  $\frac{C(s)}{R(s)} = \frac{K(T_2s+1)}{T_1s+1}$

Obtain steady state output of the system

when it is subjected to input  $r(t) = R \sin \omega t$ .

Solution

let  $G(s) = \frac{C(s)}{R(s)}$

since input is sinusoidal, and system is LTI, then output is  $R |G(s)|_{s=j\omega} \sin(\omega t + \phi)$

where  $\phi = \angle G(s)|_{s=j\omega}$

$$G(s)|_{s=j\omega} = \frac{K(T_2j\omega + 1)}{T_1j\omega + 1} = \frac{j(KT_2\omega) + K}{j(T_1\omega) + 1}$$

$$\therefore |G(s)|_{s=j\omega} = \frac{\sqrt{(KT_2\omega)^2 + K^2}}{\sqrt{T_1^2\omega^2 + 1}} = \frac{K \sqrt{T_2^2\omega^2 + 1}}{\sqrt{T_1^2\omega^2 + 1}}$$

$$\angle G(j\omega) = \tan^{-1} \frac{KT_2\omega}{K} - \tan^{-1} T_1\omega$$

$$= \tan^{-1} T_2\omega - \tan^{-1} T_1\omega$$

so  $C(s) = R \frac{K \sqrt{T_2^2\omega^2 + 1}}{\sqrt{T_1^2\omega^2 + 1}} \sin(\omega t + (\tan^{-1} T_2\omega - \tan^{-1} T_1\omega))$

10

The above is the answer.

side question why can't I seem to be able to solve above using Final value theorem? what Am I doing wrong below?

$C(s) = X(s)G(s)$  . but  $X(s) = R \frac{\omega}{s^2 + \omega^2}$  .

hence  $C(s) = \frac{R\omega}{(s^2 + \omega^2)} \cdot \frac{K(T_2s+1)}{(T_1s+1)} \Rightarrow C(s) = \lim_{s \rightarrow 0} s C(s)$

i.e  $C(s) = \lim_{s \rightarrow 0} s \cdot \frac{R\omega K(T_2s+1)}{(s^2 + \omega^2)(T_1s+1)} = \lim_{s \rightarrow 0} \frac{s^2 R\omega K T_2 + s R\omega K}{T_1 s^3 + s^2 + T_1 \omega^2 s + \omega^2}$

$= \lim_{s \rightarrow 0} \frac{R\omega K T_2 + \frac{R\omega K}{s}}{T_1 s + 1 + T_1 \omega^2/s + \omega^2/s} \rightarrow \frac{R\omega K T_2}{\infty} = 0$

?? how is this equivalent to  $\infty$  oscillation?

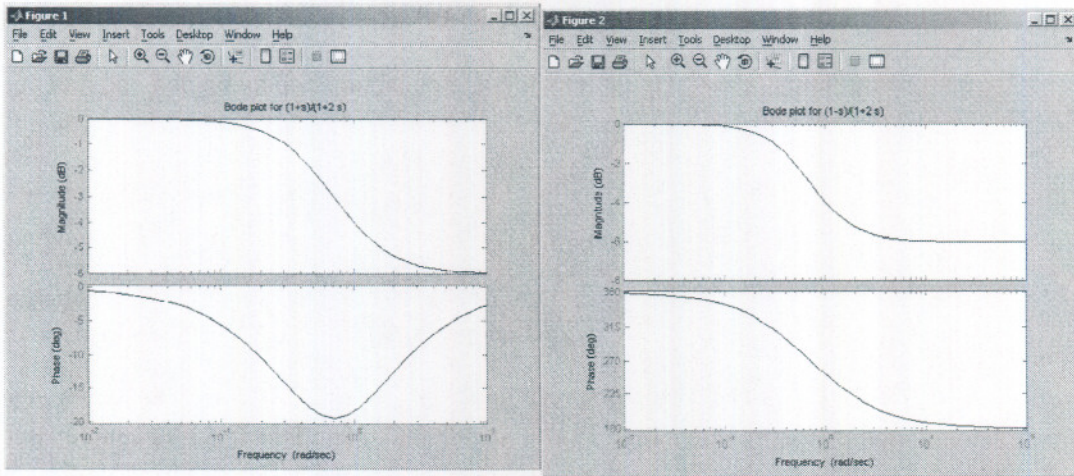
see pg. 25  
FVT only applies when f(t) settles down to a definite value for  $t \rightarrow \infty$   
 $\infty$  oscillation?

```
close all;
clear all;
%problem B-8-3, by Nasser Abbasi
```

```
s=tf('s');
sys=(1+s)/(1+2*s);
bode(sys);
title('Bode plot for (1+s)/(1+2 s)');
```

Lo

```
sys=(1-s)/(1+2*s);
figure;
bode(sys);
title('Bode plot for (1-s)/(1+2 s)');
```



$$\frac{1+s}{1+2s}$$

$$\frac{1-s}{1+2s}$$

HW#7

Problem B-8-4

Ketch Bode diagram of

a)  $G(s) = \frac{T_1 s + 1}{T_2 s + 1}$   $T_1 > T_2 > 0$

b)  $G(s) = \frac{T_1 s - 1}{T_2 s + 1}$   $T_1 > T_2 > 0$

c)  $G(s) = \frac{-T_1 s + 1}{T_2 s + 1}$   $T_1 > T_2 > 0$

note:  
after doing this I see that  $|G(s)|$  for parts a, b, c is the same only phase diagrams changed. I should have taken advantage of this to reduce work needed.  
yup

9.5

Answer

a) let  $G(s) = G_1(s) G_2(s)$  where  $G_1(s) = T_1 s + 1$ ,  $G_2(s) = \frac{1}{T_2 s + 1}$

now plot bode plot for  $G_1$  and  $G_2$  separately and combine plots.

Case  $T_1 s + 1$ , then  $G_1(j\omega) = T_1 j\omega + 1$ ,  $|G_1(j\omega)| = \sqrt{(T_1 \omega)^2 + 1}$

so  $\log |G_1(j\omega)| = \frac{1}{2} \log (1 + (T_1 \omega)^2)$

For small  $\omega$ ,  $\log |G_1| \rightarrow \frac{1}{2} \log (1) = 0$

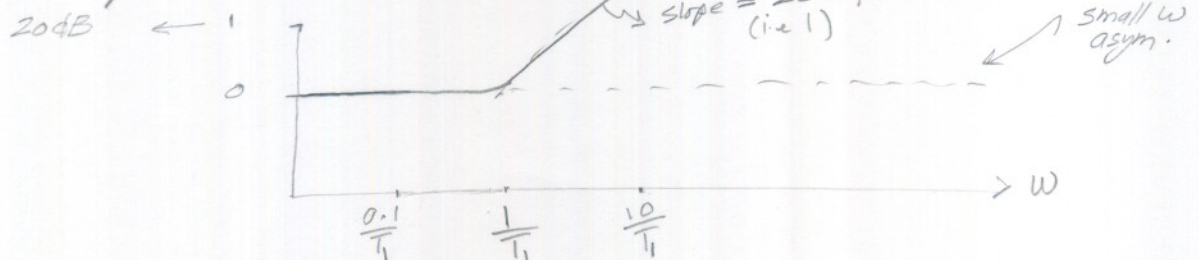
for large  $\omega$ ,  $\log |G_1| \rightarrow \frac{1}{2} \log (T_1 \omega)^2 = \log T_1 \omega$

$\angle G_1 = \tan^{-1} T_1 \omega$

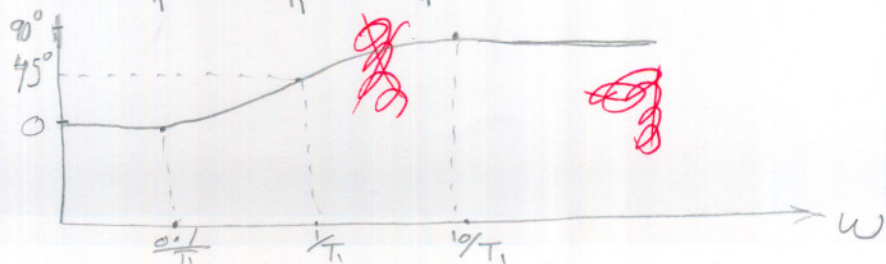
For small  $\omega$ ,  $\angle G_1 = 0^\circ$

for large  $\omega$ ,  $\angle G_1 = 90^\circ$

so Bode plot for  $G_1$  only is



next I do bode for  $G_2$



$$G_2 = \frac{1}{T_2 s + 1} \quad \text{so} \quad G_2(j\omega) = \frac{1}{T_2 j\omega + 1}$$

$$|G_2| = \frac{1}{\sqrt{1 + (T_2\omega)^2}} \rightarrow \log|G_2| = \log(1) - \log\sqrt{1 + (T_2\omega)^2} \\ = 0 - \frac{1}{2} \log(1 + (T_2\omega)^2)$$

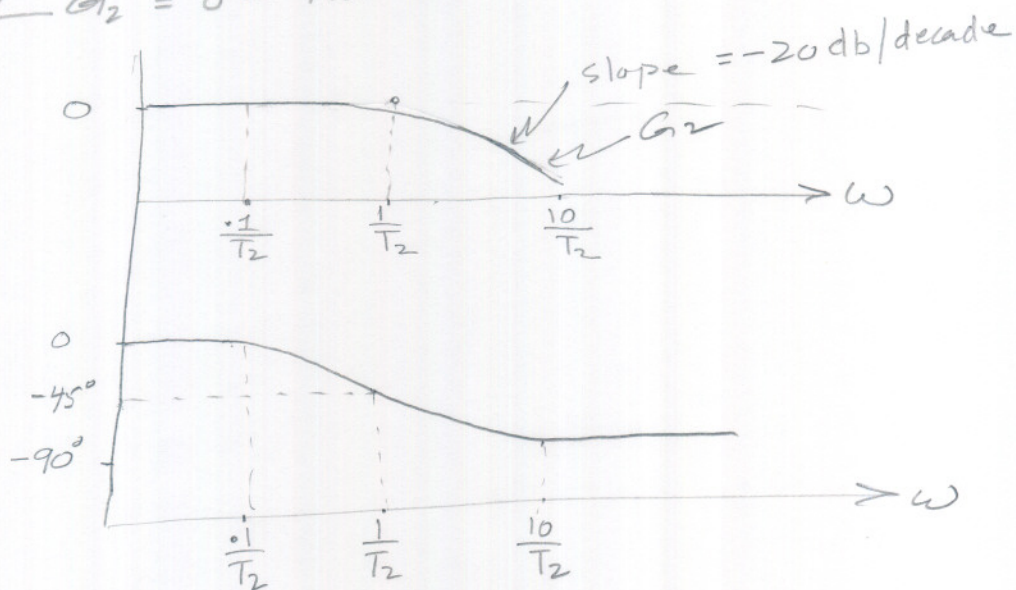
for small  $\omega$ ,  $\log|G_2| \rightarrow -\frac{1}{2} \log(1) = 0$

for large  $\omega$ ,  $\log|G_2| \rightarrow -\log T_2\omega$

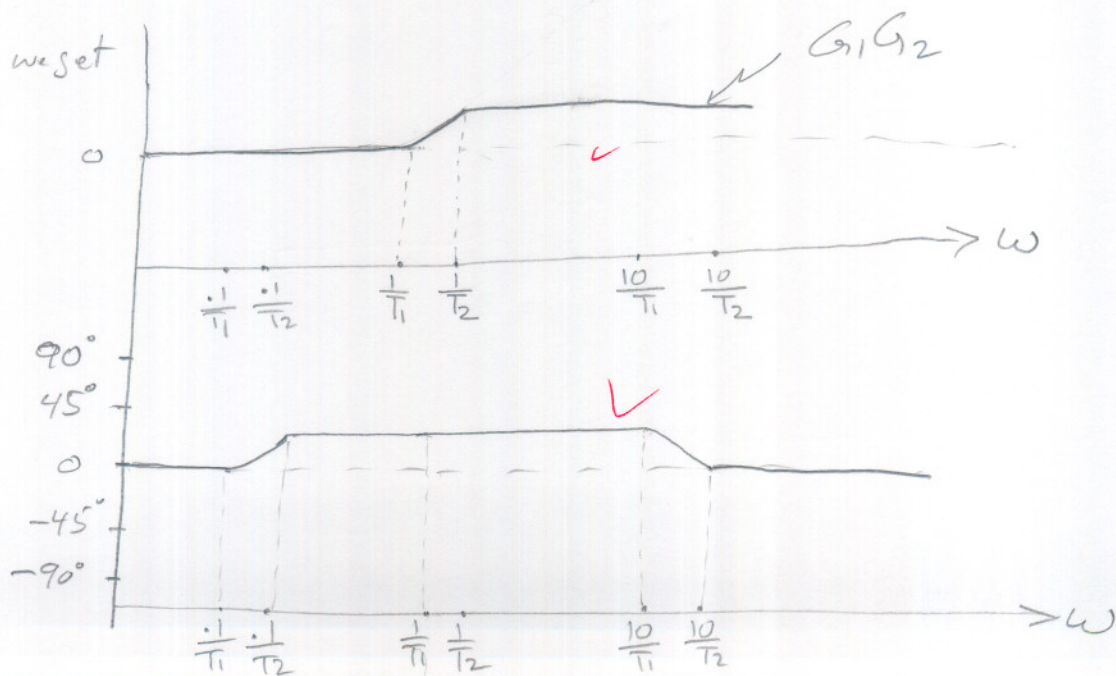
$$\angle G_2(j\omega) = 0 - \tan^{-1} T_2\omega$$

for small  $\omega$ ,  $\angle G_2 = 0 - \tan^{-1} 0 = 0$

for large  $\omega$ ,  $\angle G_2 = 0 - \tan^{-1} \infty = -90^\circ$



so now add  $G_1, G_2$  we get



Part b

$$G(s) = \frac{T_1 s - 1}{T_2 s + 1}$$

$$T_1 > T_2 > 0$$

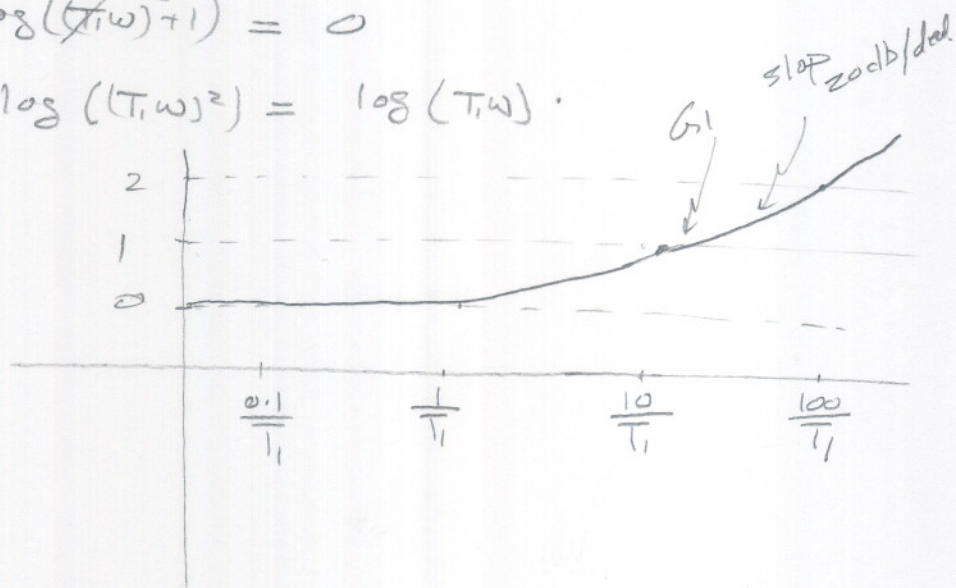
let  $G(s) = G_1 G_2$  where  $G_1 = T_1 s - 1$ ,  $G_2 = \frac{1}{T_2 s + 1}$

$G_1$  let  $G_1(j\omega) = T_1 j\omega - 1$

$$|G_1(j\omega)| = \sqrt{(T_1 \omega)^2 + 1}, \quad \log |G_1(j\omega)| = \frac{1}{2} \log((T_1 \omega)^2 + 1)$$

small  $\omega$ ,  $\log |G_1(j\omega)| \rightarrow \frac{1}{2} \log((T_1 \omega)^2 + 1) = 0$

large  $\omega$ ,  $\log |G_1(j\omega)| \rightarrow \frac{1}{2} \log((T_1 \omega)^2) = \log(T_1 \omega)$

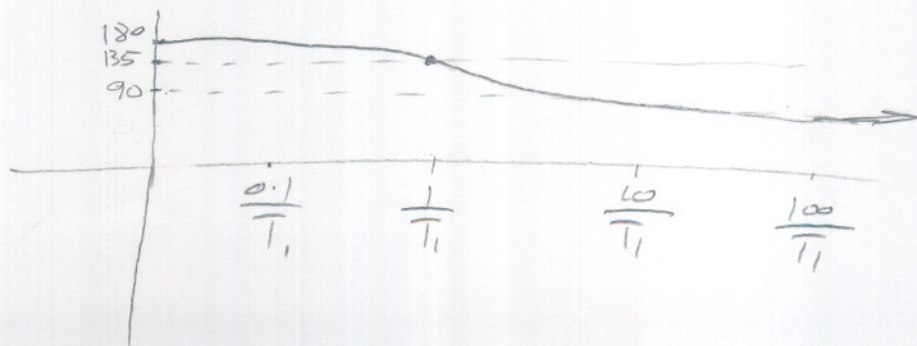


To Find Phase:

$$\angle G_1 = \tan^{-1} \frac{T_1 \omega}{-1}$$



so for small  $\omega$ , we see that  $\angle G_1$  goes to  $180^\circ$  and for large  $\omega$ , we see that  $\angle G_1$  goes to  $90^\circ$   
 when  $\omega = \frac{1}{T_1}$ ,  $\angle G_1 = \tan^{-1} -1$  i.e.  $135^\circ$   
 when  $\omega = \frac{10}{T_1}$ ,  $\angle G_1 = \tan^{-1} -10$ , etc.



next I do  $G_2 \rightarrow$

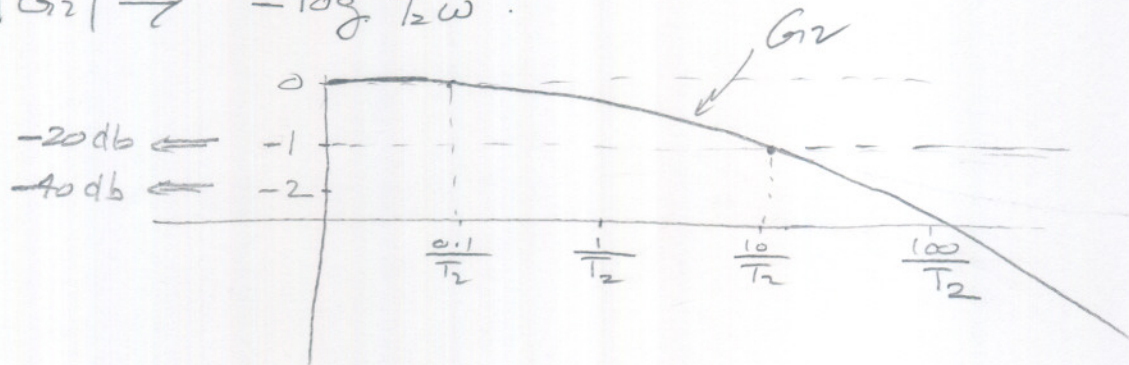
$$G_2 = \frac{1}{T_2 s + 1}$$

$$G_2(j\omega) = \frac{1}{T_2 j\omega + 1} \quad |G_2(j\omega)| = \frac{1}{\sqrt{(T_2\omega)^2 + 1}}$$

$$\log |G_2| = 0 - \log \sqrt{(T_2\omega)^2 + 1} = -\frac{1}{2} \log ((T_2\omega)^2 + 1)$$

For small  $\omega$ ,  $\log |G_2| \rightarrow 0$

For large  $\omega$ ,  $\log |G_2| \rightarrow -\log T_2\omega$

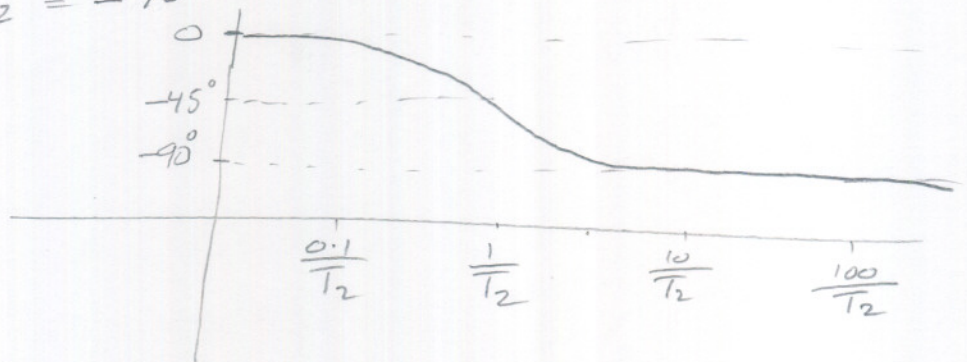


For phase

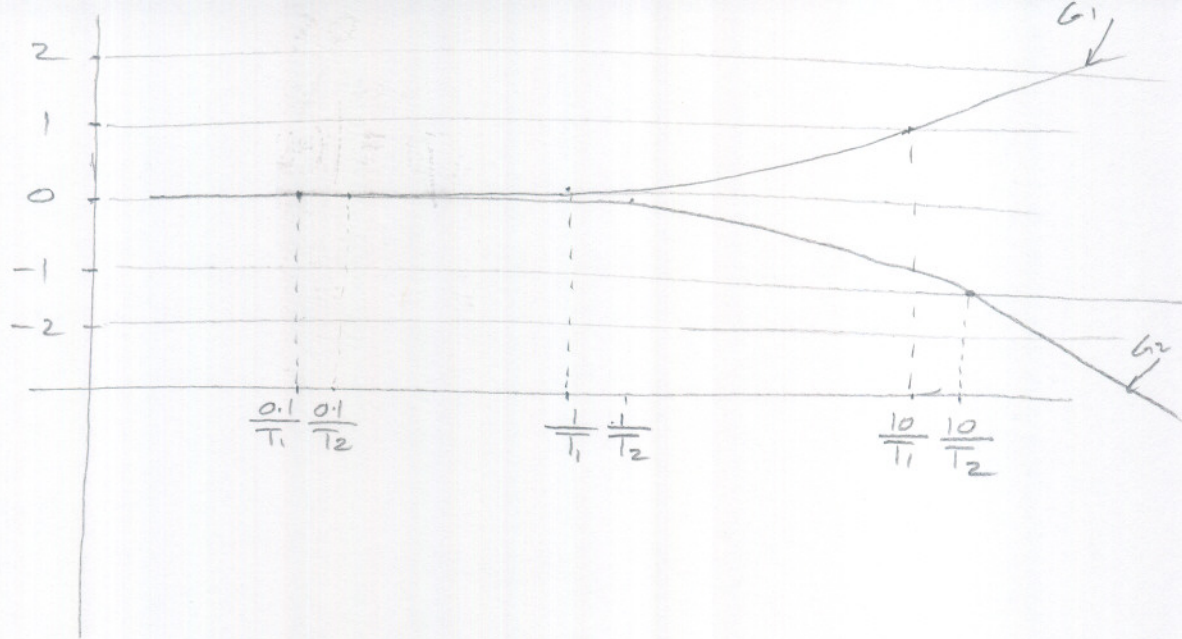
$$\angle G_2 = \tan^{-1} 0 - \tan^{-1}(T_2\omega) = -\tan^{-1} T_2\omega$$

For small  $\omega$ ,  $\angle G_2 = 0$

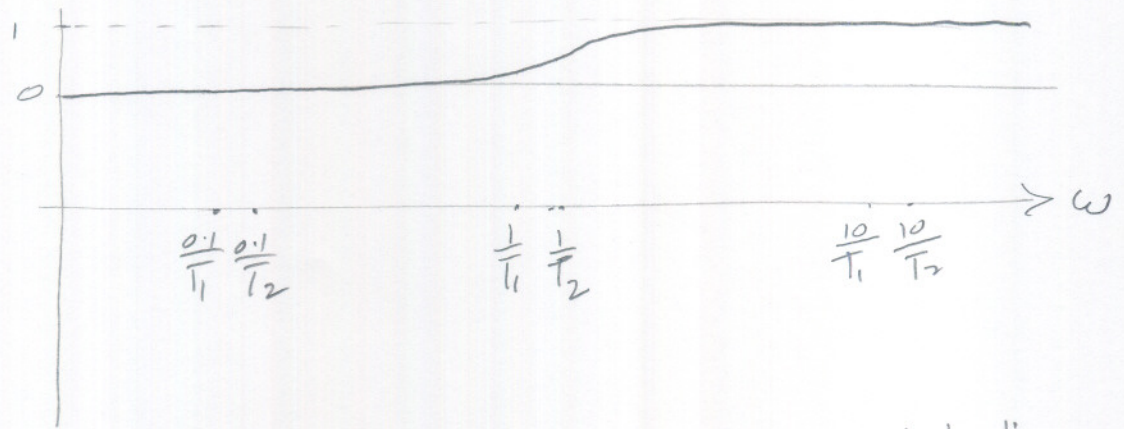
For large  $\omega$ ,  $\angle G_2 = -90^\circ$



so now I put  $G_1, G_2$  together  $\Rightarrow$

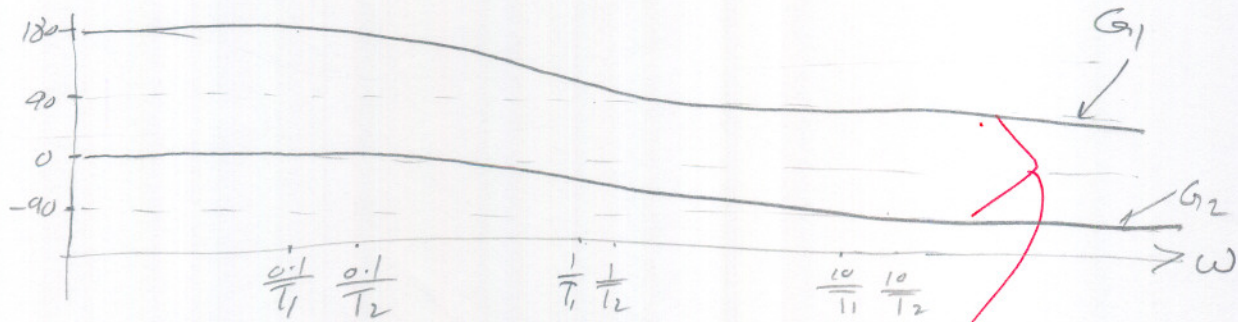


combine  $\Rightarrow$



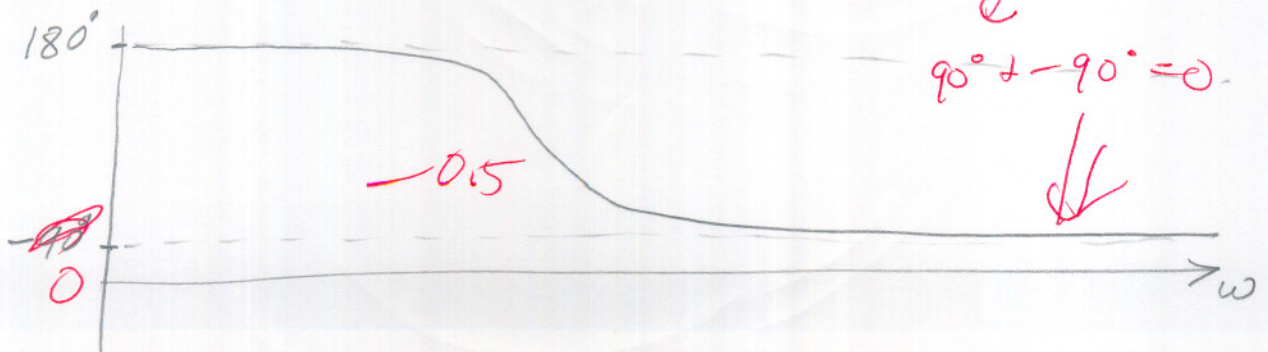
at large  $\omega$ ,  $G_1$  slope cancel  $G_2$ , we get straight line.  
 at low  $\omega$ ,  $G_1$  rises before  $G_2$  since  $T_1 > T_2$ .

for phase:



$90^\circ + -90^\circ = 0$

combine





part c

$$G(s) = \frac{-T_1 s + 1}{T_2 s + 1} \quad T_1 > T_2 > 0$$

let  $G_1(s) = -T_1 s + 1$ ,  $G_2(s) = \frac{1}{T_2 s + 1}$

now do Bode plot for each  $G_i$  separately and combine results.

$G_1$

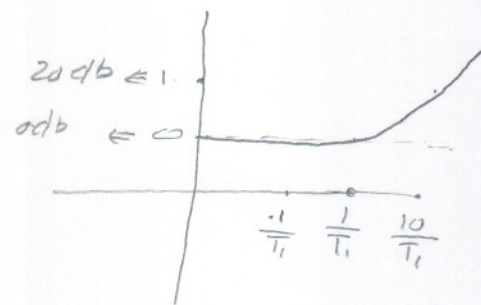
$$G_1(j\omega) = -T_1 j\omega + 1$$

$$|G_1| = \sqrt{(T_1 \omega)^2 + 1}$$

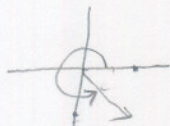
$$\log |G_1| = \frac{1}{2} \log [(T_1 \omega)^2 + 1]$$

for large  $\omega$ ,  $\log |G_1| \rightarrow \log (T_1 \omega)$

for small  $\omega$ ,  $\log |G_1| \rightarrow 0$

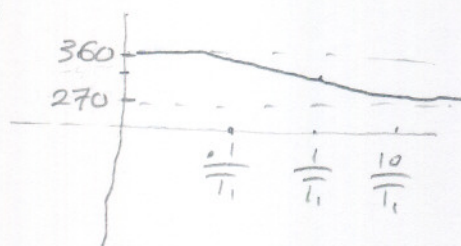


$$\angle G_1 = \tan^{-1} -T_1 \omega$$



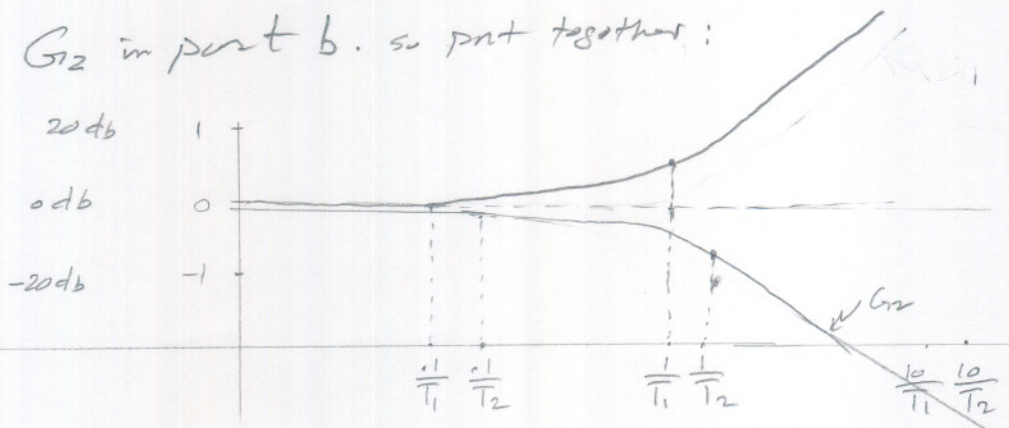
so for small  $\omega$ ,  $\angle G_1 = 360^\circ$  (could also say  $0^\circ$ )

for large  $\omega$ ,  $\angle G_1 = 270^\circ$  (could also say  $-90^\circ$ )

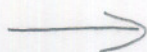


Now do  $G_2$

This is the same as  $G_2$  in part b. so put together:



Combine, we get

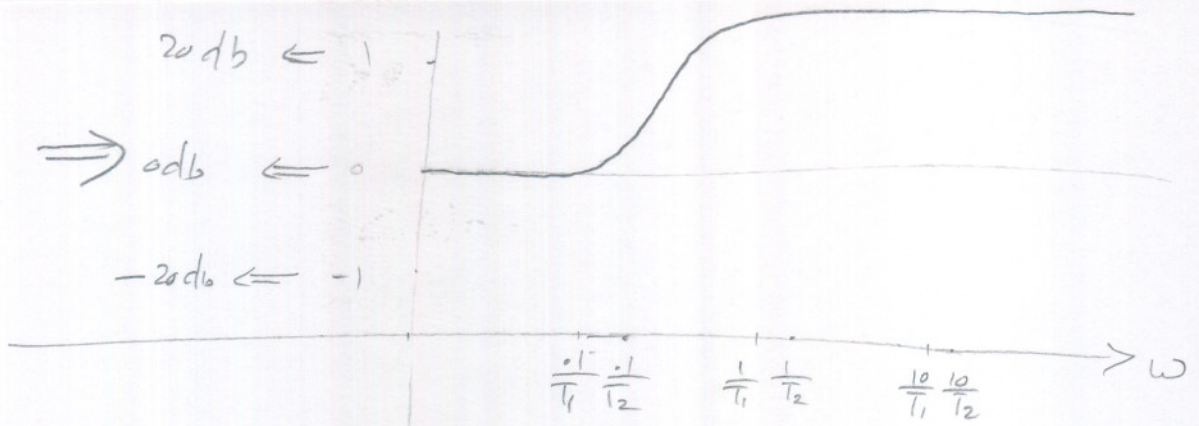


$$G(s) = \frac{-T_1 s + 1}{T_2 s + 1}$$

$$\Rightarrow 0 \text{ dB} \leftarrow 0$$

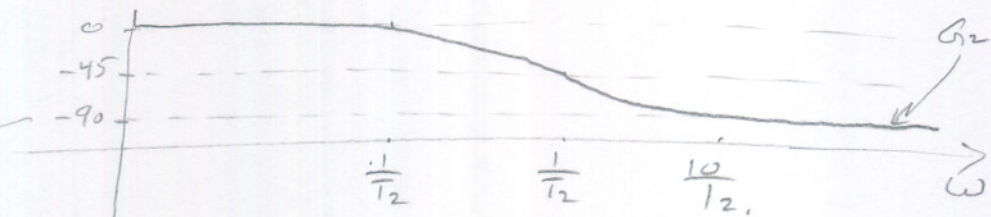
$$-20 \text{ dB} \leftarrow -1$$

$$20 \text{ dB} \leftarrow 1$$

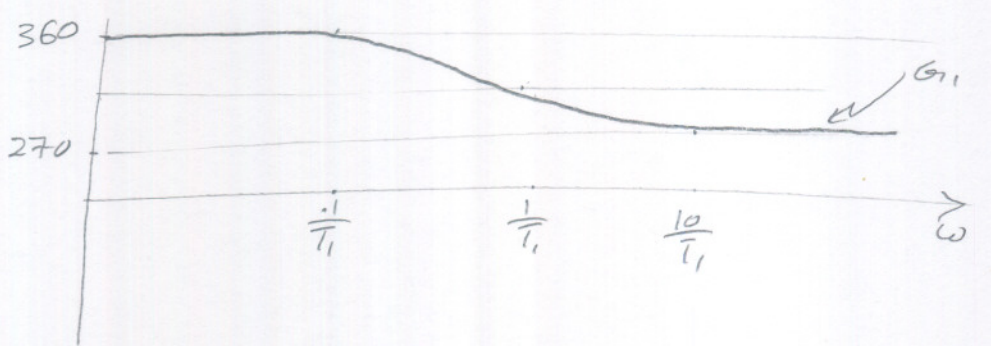


How do phase

$\angle$  for  $G_2$  is same as part (b)  $G_2$  which is

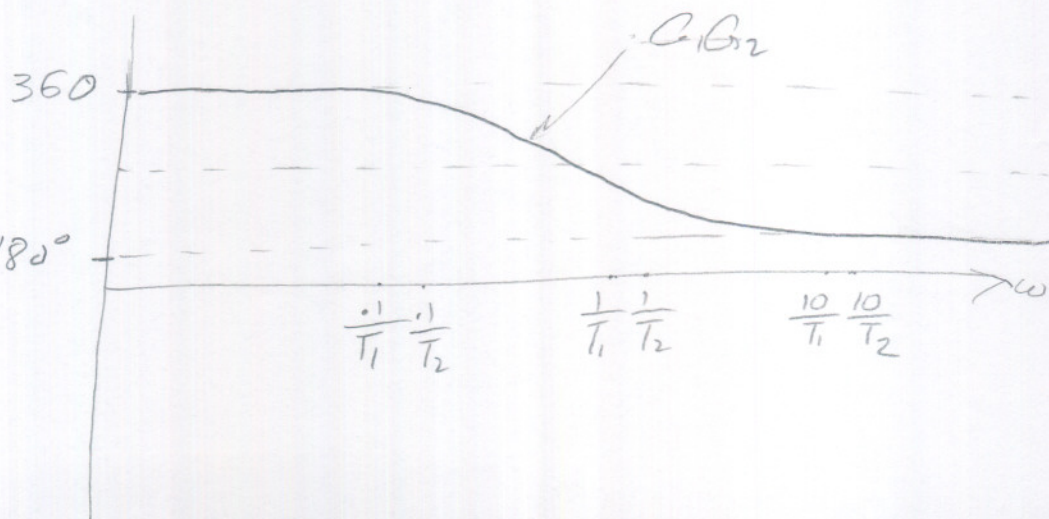


$$270 - 90 = 180^\circ$$



combine :

$$\text{This } (270 - 90) \leftarrow 180^\circ$$



HW#7

Problem B-8-6

show that  $|G(j\omega_n)| = \frac{1}{2\zeta}$

when  $G(s) = \frac{\omega_n^2}{s^2 + 2\zeta\omega_n s + \omega_n^2}$

Solution

$$G(j\omega) = \frac{\omega_n^2}{-\omega^2 + 2\zeta\omega_n j\omega + \omega_n^2}$$

$$|G(j\omega)| = \frac{\omega_n^2}{\sqrt{(-\omega^2 + \omega_n^2)^2 + (2\zeta\omega_n\omega)^2}}$$

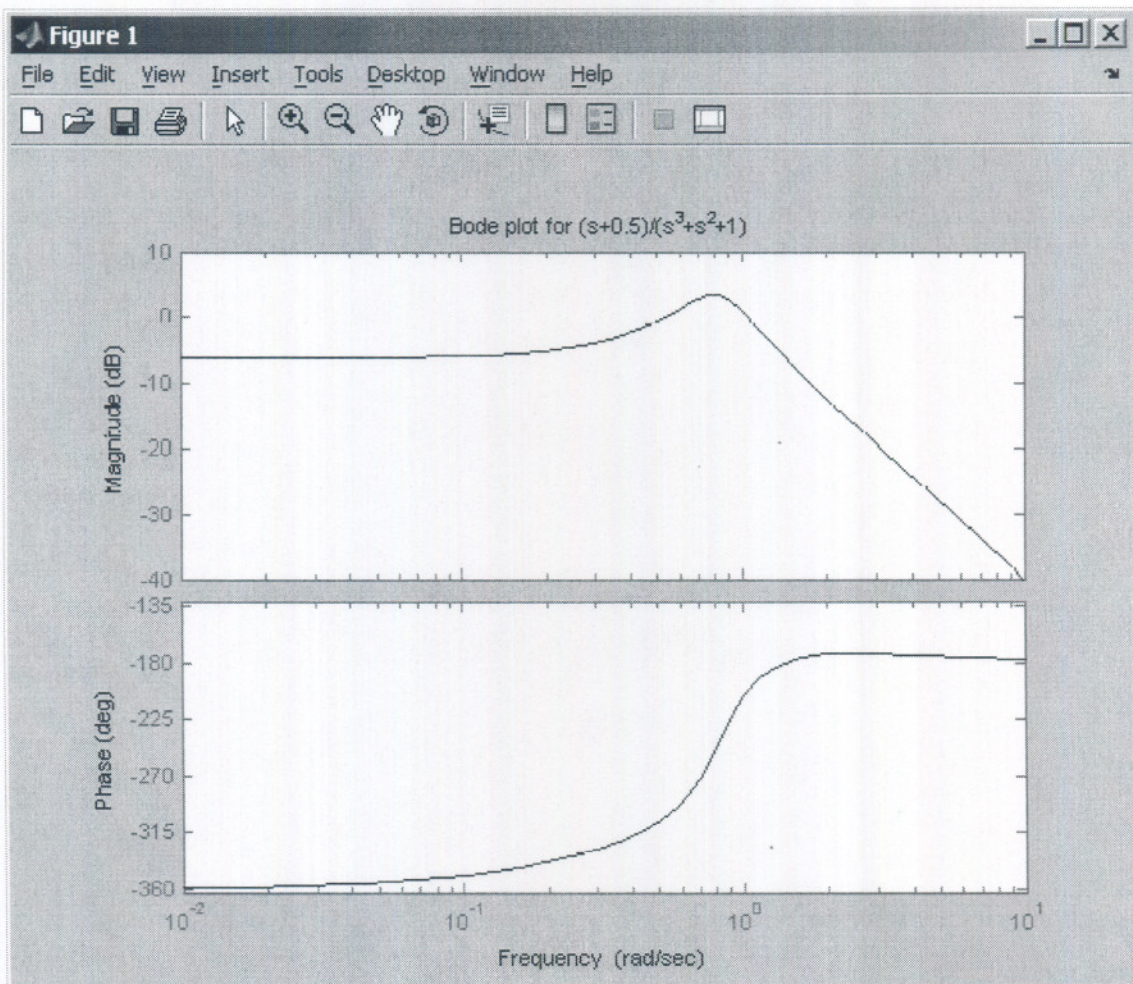
let  $\omega = \omega_n$  as required

$$|G(j\omega_n)| = \frac{\omega_n^2}{2\zeta\omega_n^2} = \boxed{\frac{1}{2\zeta}}$$

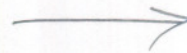
10

```
close all;
clear all;
%problem B-8-7, by Nasser Abbasi

s=tf('s');
sys=(s+0.5)/(s^3+s^2+1);
bode(sys);
title('Bode plot for (s+0.5)/(s^3+s^2+1)');
```



(10)



now need to explain phase diagram.

Phase as shown by matlab starts from  $-360^\circ$  and ends at  $-180^\circ$ . This is the same as starting from  $0^\circ$  and ending at  $180^\circ$  as per problem statement.

$$G(j\omega) = \frac{j\omega + 1}{-j\omega^3 - \omega^2 + 1}$$

$$\begin{aligned}\angle G(j\omega) &= \angle j\omega + 1 - \angle -j\omega^3 - \omega^2 + 1 \\ &= \tan^{-1} \omega - \tan^{-1} \left( \frac{-\omega^3}{-\omega^2 + 1} \right).\end{aligned}$$

For small  $\omega$ ,  $\frac{-\omega^3}{-\omega^2 + 1} \rightarrow \frac{-\omega}{1} \rightarrow 0$

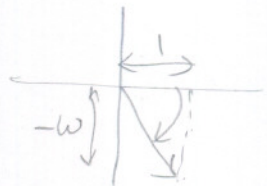
So we get in the limit

$$\angle G(j\omega) \underset{\omega \rightarrow 0}{=} \tan^{-1} \omega - \tan^{-1} (-\omega) = 0 - 360^\circ = -360^\circ$$

which is the result given by Matlab. But  $-360^\circ = 0^\circ$  also.

Now I consider what happens when  $\omega \rightarrow \infty$ .

$$\begin{aligned}\lim_{\omega \rightarrow \infty} \angle G(j\omega) &= \tan^{-1} \omega \underset{\omega \rightarrow \infty}{=} \tan^{-1} \left( \frac{-\omega^3}{-\omega^2 + 1} \right) \underset{\omega \rightarrow \infty}{=} \\ &= 90^\circ - \tan^{-1} (-\omega) \underset{\omega \rightarrow \infty}{=} \\ &= 90^\circ - (-90^\circ) \\ &= 180^\circ\end{aligned}$$



This explains the asymptotic behavior given.