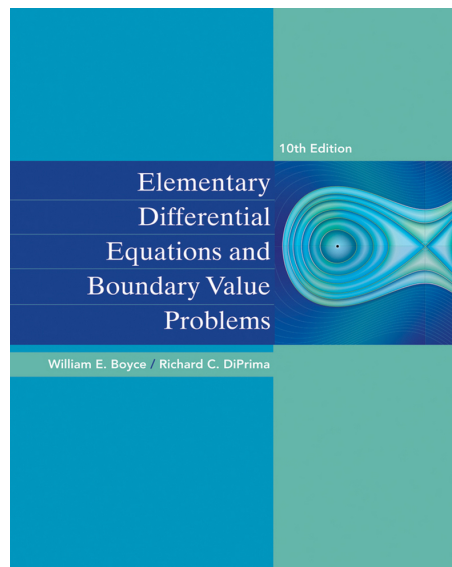


A Solution Manual For

**Elementary differential equations and
boundary value problems, 10th ed.,
Boyce and DiPrima**



Nasser M. Abbasi

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1.1 problem 1

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Internal problem ID [448]

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Section: Section 2.1. Page 40

Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$3y + y' = e^{-2t} + t$$

1.1.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = 3$$
$$q(t) = e^{-2t} + t$$

Hence the ode is

$$3y + y' = e^{-2t} + t$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int 3dt} \\ &= e^{3t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (e^{-2t} + t) \\ \frac{d}{dt}(e^{3t}y) &= (e^{3t}) (e^{-2t} + t) \\ d(e^{3t}y) &= ((e^{2t}t + 1) e^t) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{3t}y &= \int (e^{2t}t + 1) e^t dt \\ e^{3t}y &= \frac{e^{3t}t}{3} - \frac{e^{3t}}{9} + e^t + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{3t}$ results in

$$y = e^{-3t} \left(\frac{e^{3t}t}{3} - \frac{e^{3t}}{9} + e^t \right) + c_1 e^{-3t}$$

which simplifies to

$$y = \frac{t}{3} - \frac{1}{9} + e^{-2t} + c_1 e^{-3t}$$

Summary

The solution(s) found are the following

$$y = \frac{t}{3} - \frac{1}{9} + e^{-2t} + c_1 e^{-3t} \tag{1}$$

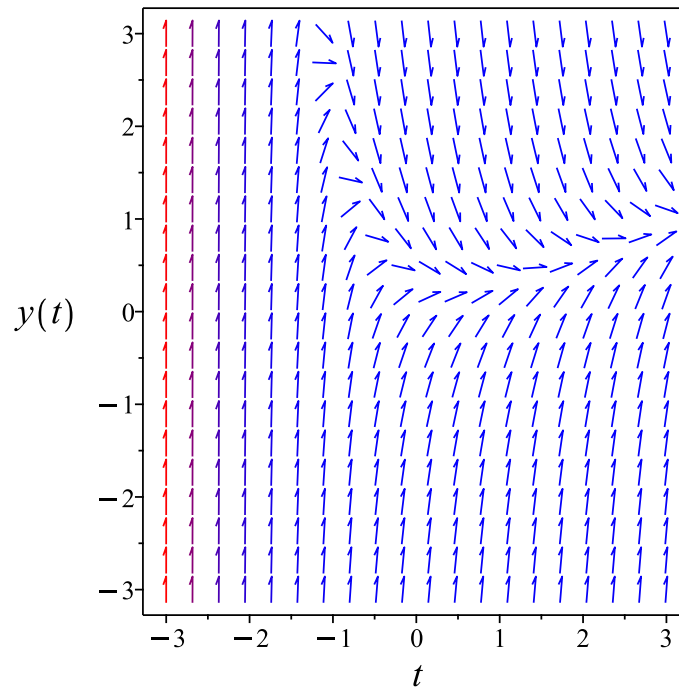


Figure 1: Slope field plot

Verification of solutions

$$y = \frac{t}{3} - \frac{1}{9} + e^{-2t} + c_1 e^{-3t}$$

Verified OK.

1.1.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -3y + e^{-2t} + t$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 1: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-3t}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-3t}} dy \end{aligned}$$

Which results in

$$S = e^{3t}y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -3y + e^{-2t} + t$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= 3e^{3t}y \\ S_y &= e^{3t} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^t + e^{3t}t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^R + e^{3R}R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{e^{3R}R}{3} - \frac{e^{3R}}{9} + e^R + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{3t}y = \frac{e^{3t}t}{3} - \frac{e^{3t}}{9} + e^t + c_1$$

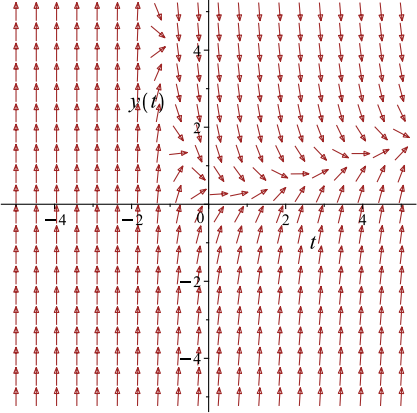
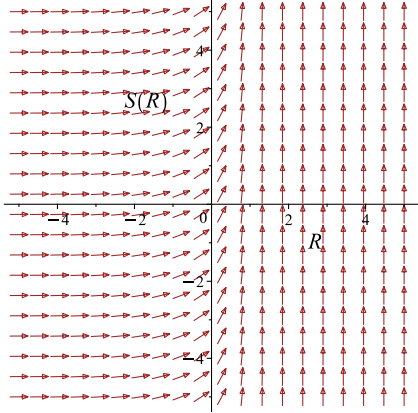
Which simplifies to

$$e^{3t}y = \frac{e^{3t}t}{3} - \frac{e^{3t}}{9} + e^t + c_1$$

Which gives

$$y = \frac{(3e^{3t}t - e^{3t} + 9e^t + 9c_1)e^{-3t}}{9}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -3y + e^{-2t} + t$ 	$R = t$ $S = e^{3t}y$	$\frac{dS}{dR} = e^R + e^{3R}R$ 

Summary

The solution(s) found are the following

$$y = \frac{(3e^{3t}t - e^{3t} + 9e^t + 9c_1)e^{-3t}}{9} \quad (1)$$

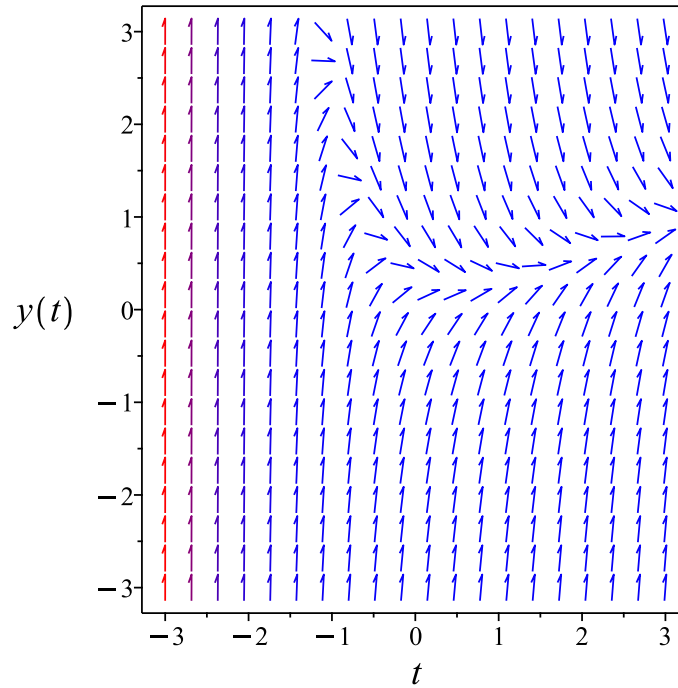


Figure 2: Slope field plot

Verification of solutions

$$y = \frac{(3e^{3t}t - e^{3t} + 9e^t + 9c_1)e^{-3t}}{9}$$

Verified OK.

1.1.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}dy &= (-3y + e^{-2t} + t) dt \\ (3y - e^{-2t} - t) dt + dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= 3y - e^{-2t} - t \\ N(t, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(3y - e^{-2t} - t) \\ &= 3\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1((3) - (0)) \\ &= 3 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int 3 dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{3t} \\ &= e^{3t} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{3t}(3y - e^{-2t} - t) \\ &= -(1 + (-3y + t)e^{2t})e^t \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{3t}(1) \\ &= e^{3t} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ (-1 + (-3y + t)e^{2t})e^t + (e^{3t}) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \overline{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -(1 + (-3y + t) e^{2t}) e^t dt \\ \phi &= \frac{(-3t + 9y + 1) e^{3t}}{9} - e^t + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{3t} + f'(y)\quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{3t}$. Therefore equation (4) becomes

$$e^{3t} = e^{3t} + f'(y)\quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{(-3t + 9y + 1) e^{3t}}{9} - e^t + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{(-3t + 9y + 1) e^{3t}}{9} - e^t$$

The solution becomes

$$y = \frac{(3 e^{3t} t - e^{3t} + 9 e^t + 9 c_1) e^{-3t}}{9}$$

Summary

The solution(s) found are the following

$$y = \frac{(3e^{3t}t - e^{3t} + 9e^t + 9c_1)e^{-3t}}{9} \quad (1)$$

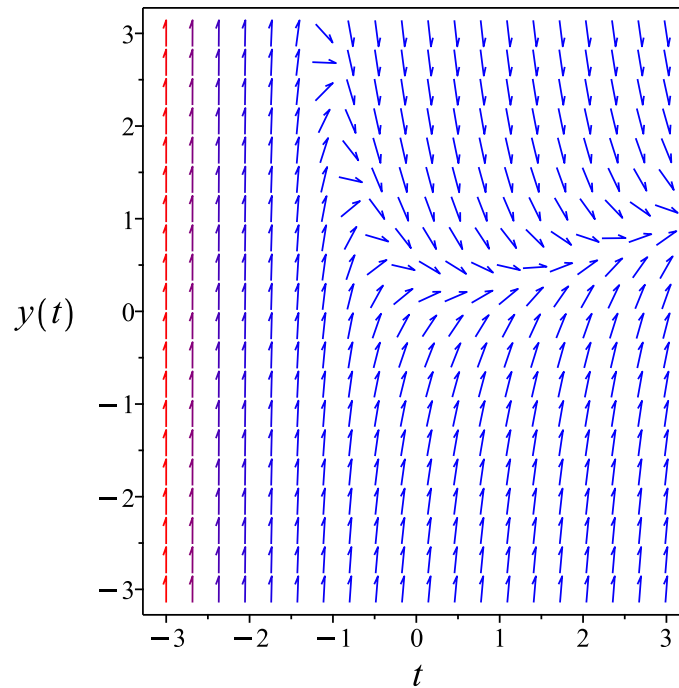


Figure 3: Slope field plot

Verification of solutions

$$y = \frac{(3e^{3t}t - e^{3t} + 9e^t + 9c_1)e^{-3t}}{9}$$

Verified OK.

1.1.4 Maple step by step solution

Let's solve

$$3y + y' = e^{-2t} + t$$

- Highest derivative means the order of the ODE is 1

y'

- Isolate the derivative

$$y' = -3y + e^{-2t} + t$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$3y + y' = e^{-2t} + t$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) (3y + y') = \mu(t) (e^{-2t} + t)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t) y)$

$$\mu(t) (3y + y') = \mu'(t) y + \mu(t) y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = 3\mu(t)$$

- Solve to find the integrating factor

$$\mu(t) = e^{3t}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y) \right) dt = \int \mu(t) (e^{-2t} + t) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t) y = \int \mu(t) (e^{-2t} + t) dt + c_1$$

- Solve for y

$$y = \frac{\int \mu(t) (e^{-2t} + t) dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{3t}$

$$y = \frac{\int (e^{-2t} + t) e^{3t} dt + c_1}{e^{3t}}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{(e^t)^3}{3} - \frac{(e^t)^3}{9} + e^t + c_1}{e^{3t}}$$

- Simplify

$$y = \frac{(3e^{3t}t - e^{3t} + 9e^t + 9c_1)e^{-3t}}{9}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(3*y(t)+diff(y(t),t) = exp(-2*t)+t,y(t), singsol=all)
```

$$y(t) = \frac{t}{3} - \frac{1}{9} + e^{-2t} + c_1 e^{-3t}$$

✓ Solution by Mathematica

Time used: 0.086 (sec). Leaf size: 27

```
DSolve[3*y[t]+y'[t] == Exp[-2*t]+t,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{t}{3} + e^{-2t} + c_1 e^{-3t} - \frac{1}{9}$$

1.2 problem 2

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Internal problem ID [449]

Internal file name [OUTPUT/449_Sunday_June_05_2022_01_41_40_AM_16129924/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$-2y + y' = e^{2t}t^2$$

1.2.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -2$$
$$q(t) = e^{2t}t^2$$

Hence the ode is

$$-2y + y' = e^{2t}t^2$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int(-2)dt} \\ &= e^{-2t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (e^{2t}t^2) \\ \frac{d}{dt}(e^{-2t}y) &= (e^{-2t}) (e^{2t}t^2) \\ d(e^{-2t}y) &= t^2 dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-2t}y &= \int t^2 dt \\ e^{-2t}y &= \frac{t^3}{3} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-2t}$ results in

$$y = \frac{e^{2t}t^3}{3} + c_1e^{2t}$$

which simplifies to

$$y = e^{2t} \left(\frac{t^3}{3} + c_1 \right)$$

Summary

The solution(s) found are the following

$$y = e^{2t} \left(\frac{t^3}{3} + c_1 \right) \tag{1}$$

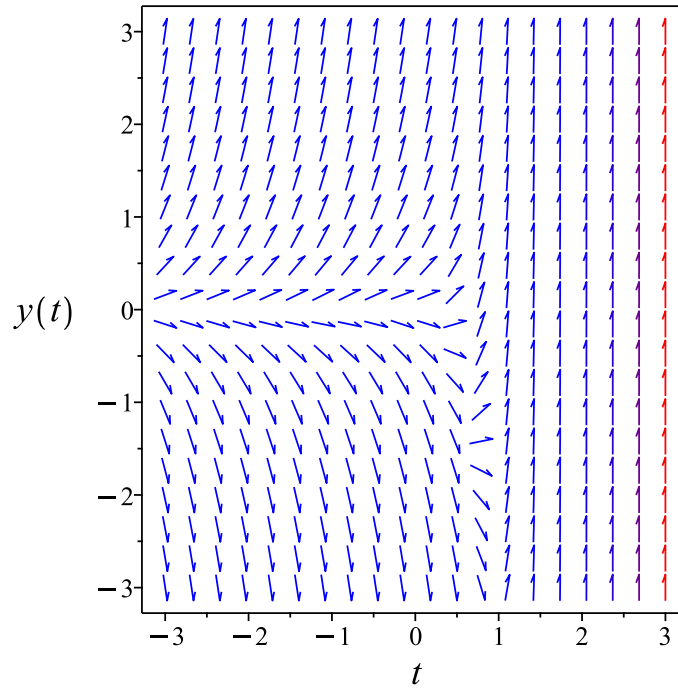


Figure 4: Slope field plot

Verification of solutions

$$y = e^{2t} \left(\frac{t^3}{3} + c_1 \right)$$

Verified OK.

1.2.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = 2y + e^{2t}t^2$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 4: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{2t}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{2t}} dy \end{aligned}$$

Which results in

$$S = e^{-2t}y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = 2y + e^{2t}t^2$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -2e^{-2t}y \\ S_y &= e^{-2t} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = t^2 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R^2$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^3}{3} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{-2t}y = \frac{t^3}{3} + c_1$$

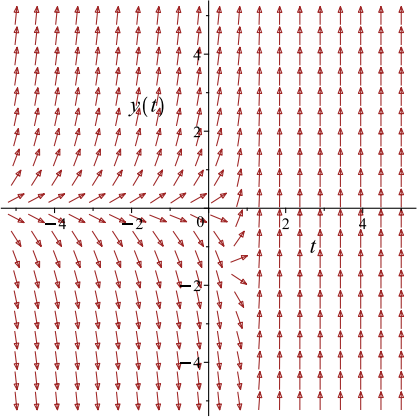
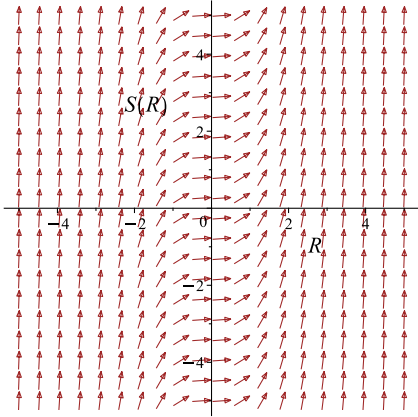
Which simplifies to

$$e^{-2t}y = \frac{t^3}{3} + c_1$$

Which gives

$$y = \frac{e^{2t}(t^3 + 3c_1)}{3}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = 2y + e^{2t}t^2$ 	$R = t$ $S = e^{-2t}y$	$\frac{dS}{dR} = R^2$ 

Summary

The solution(s) found are the following

$$y = \frac{e^{2t}(t^3 + 3c_1)}{3} \quad (1)$$

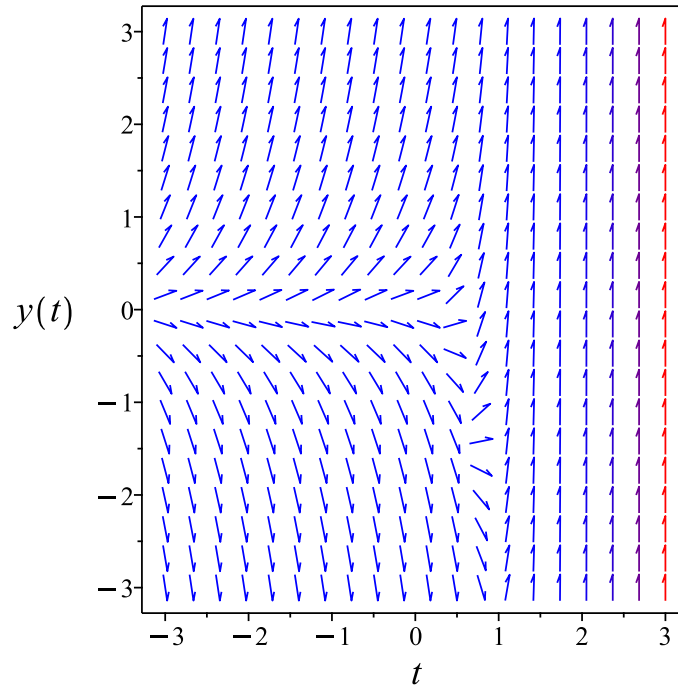


Figure 5: Slope field plot

Verification of solutions

$$y = \frac{e^{2t}(t^3 + 3c_1)}{3}$$

Verified OK.

1.2.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}dy &= (2y + e^{2t}t^2) dt \\ (-2y - e^{2t}t^2) dt + dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -2y - e^{2t}t^2 \\ N(t, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-2y - e^{2t}t^2) \\ &= -2\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1((-2) - (0)) \\ &= -2 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int -2 dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-2t} \\ &= e^{-2t} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{-2t}(-2y - e^{2t}t^2) \\ &= -2e^{-2t}y - t^2 \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{-2t}(1) \\ &= e^{-2t} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ (-2e^{-2t}y - t^2) + (e^{-2t}) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -2e^{-2t}y - t^2 dt \\ \phi &= -\frac{t^3}{3} + e^{-2t}y + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-2t} + f'(y)\quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-2t}$. Therefore equation (4) becomes

$$e^{-2t} = e^{-2t} + f'(y)\quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{t^3}{3} + e^{-2t}y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{t^3}{3} + e^{-2t}y$$

The solution becomes

$$y = \frac{e^{2t}(t^3 + 3c_1)}{3}$$

Summary

The solution(s) found are the following

$$y = \frac{e^{2t}(t^3 + 3c_1)}{3} \quad (1)$$

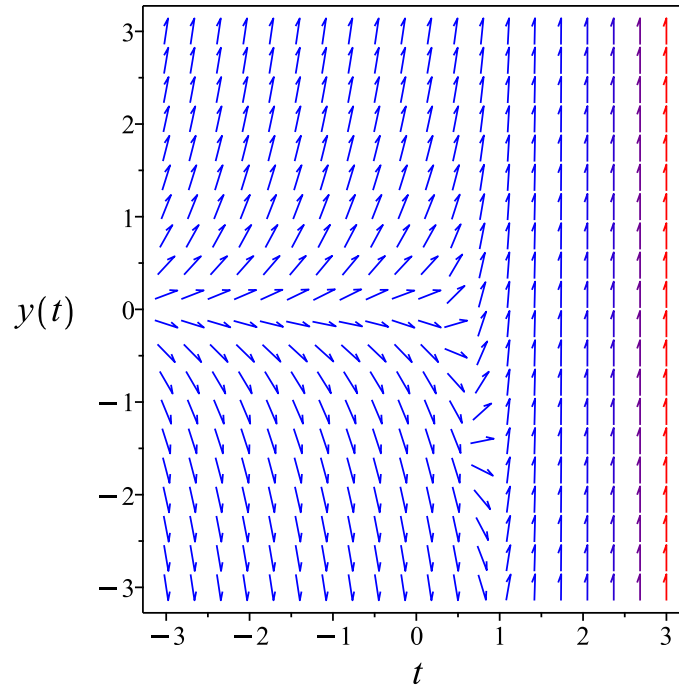


Figure 6: Slope field plot

Verification of solutions

$$y = \frac{e^{2t}(t^3 + 3c_1)}{3}$$

Verified OK.

1.2.4 Maple step by step solution

Let's solve

$$-2y + y' = e^{2t}t^2$$

- Highest derivative means the order of the ODE is 1

y'

- Isolate the derivative

$$y' = 2y + e^{2t}t^2$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$-2y + y' = e^{2t}t^2$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) (-2y + y') = \mu(t) e^{2t}t^2$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t) y)$

$$\mu(t) (-2y + y') = \mu'(t) y + \mu(t) y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -2\mu(t)$$

- Solve to find the integrating factor

$$\mu(t) = e^{-2t}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y) \right) dt = \int \mu(t) e^{2t}t^2 dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t) y = \int \mu(t) e^{2t}t^2 dt + c_1$$

- Solve for y

$$y = \frac{\int \mu(t) e^{2t}t^2 dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{-2t}$

$$y = \frac{\int e^{2t}t^2 e^{-2t} dt + c_1}{e^{-2t}}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{t^3}{3} + c_1}{e^{-2t}}$$

- Simplify

$$y = \frac{e^{2t}(t^3 + 3c_1)}{3}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(-2*y(t)+diff(y(t),t) = exp(2*t)*t^2,y(t), singsol=all)
```

$$y(t) = \frac{(t^3 + 3c_1)e^{2t}}{3}$$

✓ Solution by Mathematica

Time used: 0.043 (sec). Leaf size: 22

```
DSolve[-2*y[t]+y'[t]== Exp[2*t]*t^2,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{3}e^{2t}(t^3 + 3c_1)$$

1.3 problem 3

1.3.1	Solving as linear ode	30
1.3.2	Solving as first order ode lie symmetry lookup ode	32
1.3.3	Solving as exact ode	36
1.3.4	Maple step by step solution	40

Internal problem ID [450]

Internal file name [OUTPUT/450_Sunday_June_05_2022_01_41_41_AM_58171829/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 3.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "linear", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$y + y' = 1 + te^{-t}$$

1.3.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = 1$$
$$q(t) = 1 + te^{-t}$$

Hence the ode is

$$y + y' = 1 + te^{-t}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int 1 dt} \\ &= e^t\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (1 + t e^{-t}) \\ \frac{d}{dt}(y e^t) &= (e^t) (1 + t e^{-t}) \\ d(y e^t) &= (e^t + t) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}y e^t &= \int e^t + t dt \\ y e^t &= \frac{t^2}{2} + e^t + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^t$ results in

$$y = e^{-t} \left(\frac{t^2}{2} + e^t \right) + c_1 e^{-t}$$

which simplifies to

$$y = 1 + \frac{(t^2 + 2c_1) e^{-t}}{2}$$

Summary

The solution(s) found are the following

$$y = 1 + \frac{(t^2 + 2c_1) e^{-t}}{2} \tag{1}$$

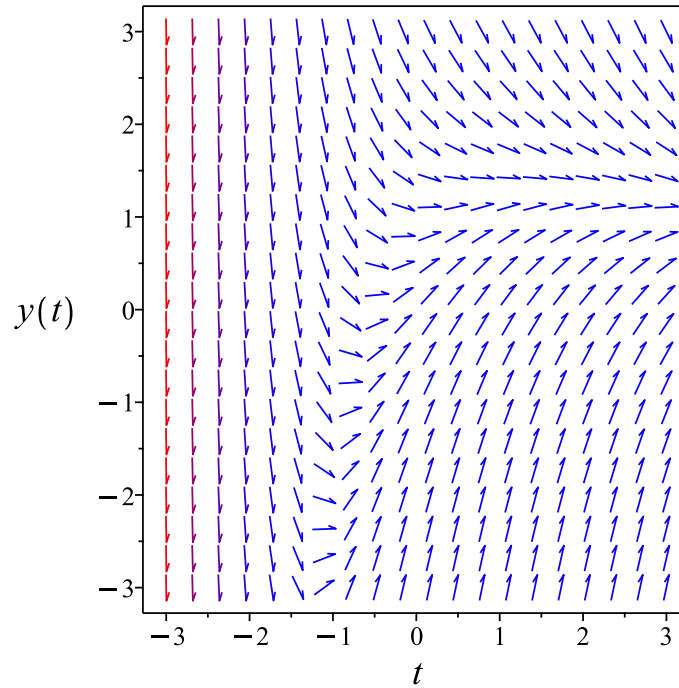


Figure 7: Slope field plot

Verification of solutions

$$y = 1 + \frac{(t^2 + 2c_1) e^{-t}}{2}$$

Verified OK.

1.3.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -(y e^t - e^t - t) e^{-t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 7: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-t}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-t}} dy \end{aligned}$$

Which results in

$$S = y e^t$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -(y e^t - e^t - t) e^{-t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= y e^t \\ S_y &= e^t \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^t + t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^R + R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^2}{2} + e^R + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$y e^t = \frac{t^2}{2} + e^t + c_1$$

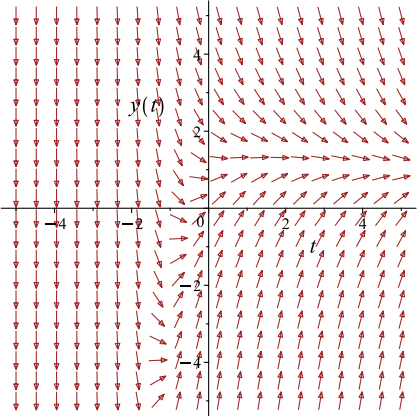
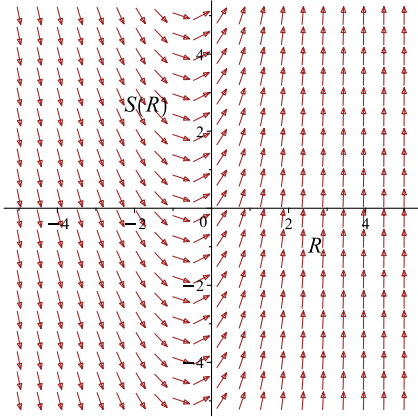
Which simplifies to

$$y e^t = \frac{t^2}{2} + e^t + c_1$$

Which gives

$$y = \frac{(t^2 + 2e^t + 2c_1)e^{-t}}{2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -(y e^t - e^t - t) e^{-t}$ 	$R = t$ $S = y e^t$	$\frac{dS}{dR} = e^R + R$ 

Summary

The solution(s) found are the following

$$y = \frac{(t^2 + 2e^t + 2c_1)e^{-t}}{2} \quad (1)$$

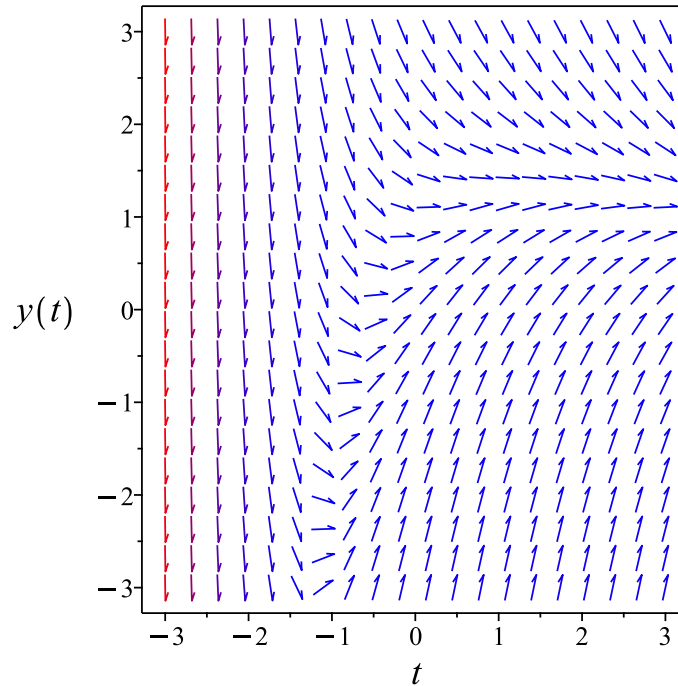


Figure 8: Slope field plot

Verification of solutions

$$y = \frac{(t^2 + 2e^t + 2c_1)e^{-t}}{2}$$

Verified OK.

1.3.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(e^t) dy &= (-y e^t + e^t + t) dt \\ (y e^t - e^t - t) dt + (e^t) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= y e^t - e^t - t \\ N(t, y) &= e^t\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (y e^t - e^t - t) \\ &= e^t\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (e^t) \\ &= e^t\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int M dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int y e^t - e^t - t dt$$

$$\phi = e^t(y - 1) - \frac{t^2}{2} + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^t + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^t$. Therefore equation (4) becomes

$$e^t = e^t + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = e^t(y - 1) - \frac{t^2}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = e^t(y - 1) - \frac{t^2}{2}$$

The solution becomes

$$y = \frac{(t^2 + 2e^t + 2c_1)e^{-t}}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{(t^2 + 2e^t + 2c_1)e^{-t}}{2} \tag{1}$$

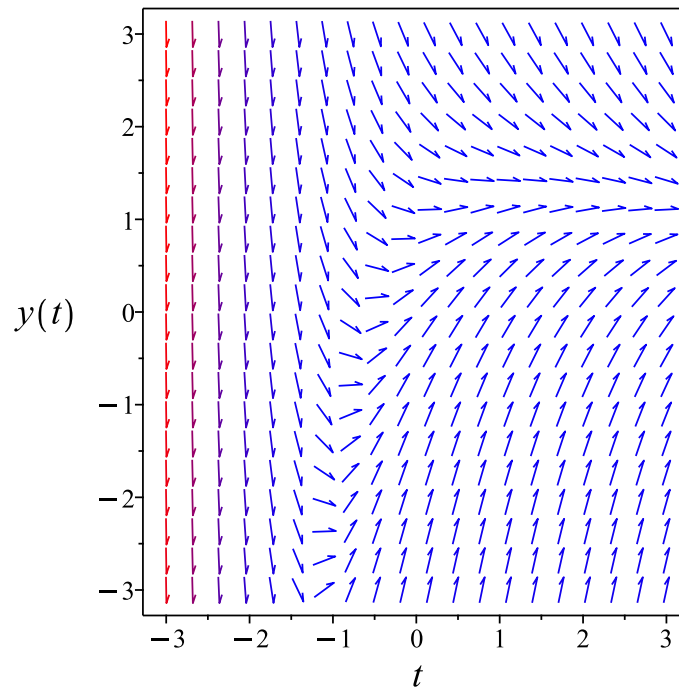


Figure 9: Slope field plot

Verification of solutions

$$y = \frac{(t^2 + 2e^t + 2c_1)e^{-t}}{2}$$

Verified OK.

1.3.4 Maple step by step solution

Let's solve

$$y + y' = 1 + \frac{t}{e^t}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -y + \frac{e^t+t}{e^t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y + y' = \frac{e^t+t}{e^t}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t)(y + y') = \frac{\mu(t)(e^t+t)}{e^t}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t)(y + y') = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \mu(t)$$

- Solve to find the integrating factor

$$\mu(t) = (e^t)^2 e^{-t}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)(e^t+t)}{e^t} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)(e^t+t)}{e^t} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)(e^t+t)}{e^t} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = (e^t)^2 e^{-t}$

$$y = \frac{\int e^t e^{-t} (e^t+t) dt + c_1}{(e^t)^2 e^{-t}}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{t^2}{2} + e^t + c_1}{(e^t)^2 e^{-t}}$$

- Simplify

$$y = 1 + \frac{(t^2 + 2c_1)e^{-t}}{2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(y(t)+diff(y(t),t) = 1+t/exp(t),y(t), singsol=all)
```

$$y(t) = 1 + \frac{(t^2 + 2c_1)e^{-t}}{2}$$

✓ Solution by Mathematica

Time used: 0.052 (sec). Leaf size: 27

```
DSolve[y[t]+y'[t] == 1+t/Exp[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{2}e^{-t}(t^2 + 2e^t + 2c_1)$$

1.4 problem 4

1.4.1	Solving as linear ode	42
1.4.2	Solving as first order ode lie symmetry lookup ode	44
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1.4.4	Maple step by step solution	52

Internal problem ID [451]

Internal file name [OUTPUT/451_Sunday_June_05_2022_01_41_42_AM_48741648/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 4.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**exact**", "**linear**", "**first_order_ode_lie_symmetry_lookup**"

Maple gives the following as the ode type

`[_linear]`

$$\frac{y}{t} + y' = 3 \cos(2t)$$

1.4.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{1}{t}$$
$$q(t) = 3 \cos(2t)$$

Hence the ode is

$$\frac{y}{t} + y' = 3 \cos(2t)$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{1}{t} dt} \\ &= t\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (3 \cos (2t)) \\ \frac{d}{dt}(ty) &= (t) (3 \cos (2t)) \\ d(ty) &= (3 \cos (2t) t) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}ty &= \int 3 \cos (2t) t dt \\ ty &= \frac{3 \cos (2t)}{4} + \frac{3 \sin (2t) t}{2} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t$ results in

$$y = \frac{\frac{3 \cos (2t)}{4} + \frac{3 \sin (2t) t}{2}}{t} + \frac{c_1}{t}$$

which simplifies to

$$y = \frac{6 \sin (2t) t + 3 \cos (2t) + 4c_1}{4t}$$

Summary

The solution(s) found are the following

$$y = \frac{6 \sin (2t) t + 3 \cos (2t) + 4c_1}{4t} \tag{1}$$

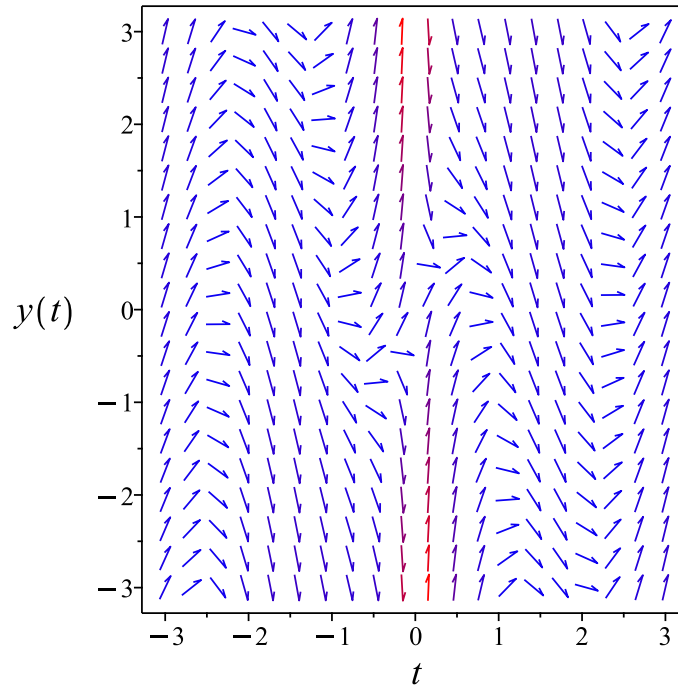


Figure 10: Slope field plot

Verification of solutions

$$y = \frac{6 \sin(2t) t + 3 \cos(2t) + 4c_1}{4t}$$

Verified OK.

1.4.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-y + 3 \cos(2t) t}{t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 10: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \frac{1}{t}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{t}} dy \end{aligned}$$

Which results in

$$S = ty$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{-y + 3 \cos(2t) t}{t}$$

Evaluating all the partial derivatives gives

$$R_t = 1$$

$$R_y = 0$$

$$S_t = y$$

$$S_y = t$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 3 \cos(2t) t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 3 \cos(2R) R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{3 \cos(2R)}{4} + \frac{3 \sin(2R) R}{2} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$yt = \frac{3 \cos(2t)}{4} + \frac{3 \sin(2t) t}{2} + c_1$$

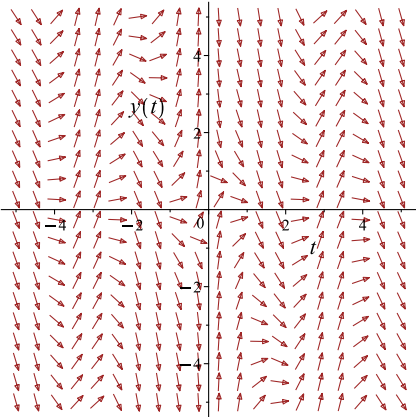
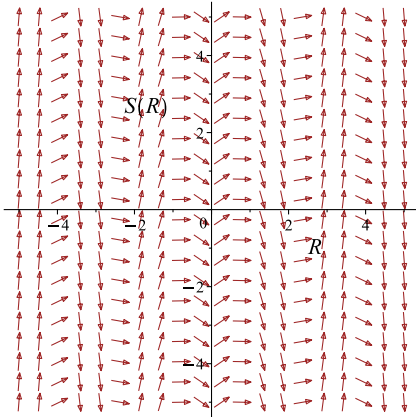
Which simplifies to

$$yt = \frac{3 \cos(2t)}{4} + \frac{3 \sin(2t) t}{2} + c_1$$

Which gives

$$y = \frac{6 \sin(2t) t + 3 \cos(2t) + 4c_1}{4t}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{-y + 3 \cos(2t)t}{t}$ 	$R = t$ $S = ty$	$\frac{dS}{dR} = 3 \cos(2R) R$ 

Summary

The solution(s) found are the following

$$y = \frac{6 \sin(2t) t + 3 \cos(2t) + 4c_1}{4t} \quad (1)$$

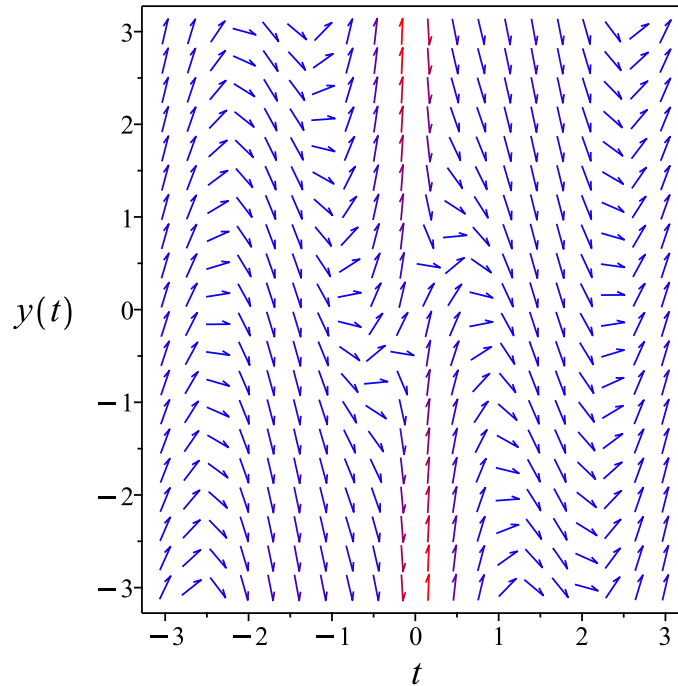


Figure 11: Slope field plot

Verification of solutions

$$y = \frac{6 \sin(2t) t + 3 \cos(2t) + 4c_1}{4t}$$

Verified OK.

1.4.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(t) dy &= (-y + 3 \cos(2t) t) dt \\ (y - 3 \cos(2t) t) dt + (t) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= y - 3 \cos(2t) t \\ N(t, y) &= t\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y - 3 \cos(2t) t) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(t) \\ &= 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int M dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int y - 3 \cos(2t) t dt$$

$$\phi = ty - \frac{3 \cos(2t)}{4} - \frac{3 \sin(2t) t}{2} + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t$. Therefore equation (4) becomes

$$t = t + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = ty - \frac{3 \cos(2t)}{4} - \frac{3 \sin(2t) t}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = ty - \frac{3 \cos(2t)}{4} - \frac{3 \sin(2t) t}{2}$$

The solution becomes

$$y = \frac{6 \sin(2t) t + 3 \cos(2t) + 4c_1}{4t}$$

Summary

The solution(s) found are the following

$$y = \frac{6 \sin(2t) t + 3 \cos(2t) + 4c_1}{4t} \tag{1}$$

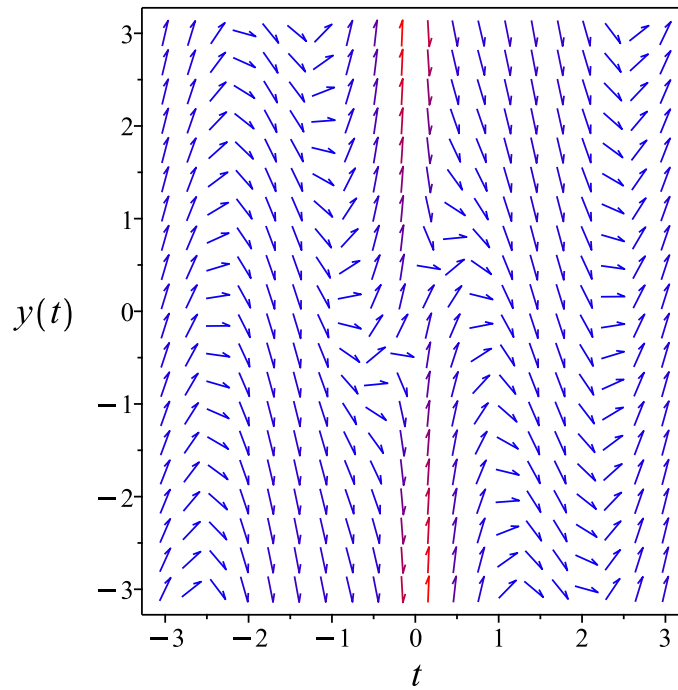


Figure 12: Slope field plot

Verification of solutions

$$y = \frac{6 \sin(2t) t + 3 \cos(2t) + 4c_1}{4t}$$

Verified OK.

1.4.4 Maple step by step solution

Let's solve

$$\frac{y}{t} + y' = 3 \cos(2t)$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{y}{t} + 3 \cos(2t)$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$\frac{y}{t} + y' = 3 \cos(2t)$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(\frac{y}{t} + y' \right) = 3\mu(t) \cos(2t)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(\frac{y}{t} + y' \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = t$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int 3\mu(t) \cos(2t) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int 3\mu(t) \cos(2t) dt + c_1$$

- Solve for y

$$y = \frac{\int 3\mu(t) \cos(2t) dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = t$

$$y = \frac{\int 3 \cos(2t) t dt + c_1}{t}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{3 \cos(2t)}{4} + \frac{3 \sin(2t)t}{2} + c_1}{t}$$

- Simplify

$$y = \frac{6 \sin(2t)t + 3 \cos(2t) + 4c_1}{4t}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 26

```
dsolve(y(t)/t+diff(y(t),t) = 3*cos(2*t),y(t), singsol=all)
```

$$y(t) = \frac{4c_1 + 6 \sin(2t)t + 3 \cos(2t)}{4t}$$

✓ Solution by Mathematica

Time used: 0.073 (sec). Leaf size: 30

```
DSolve[y[t]/t+y'[t] == 3*Cos[2*t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{6t \sin(2t) + 3 \cos(2t) + 4c_1}{4t}$$

1.5 problem 5

1.5.1	Solving as linear ode	54
1.5.2	Solving as first order ode lie symmetry lookup ode	56
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1.5.4	Maple step by step solution	64

Internal problem ID [452]

Internal file name [OUTPUT/452_Sunday_June_05_2022_01_41_43_AM_90446891/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 5.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$-2y + y' = 3e^t$$

1.5.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -2$$

$$q(t) = 3e^t$$

Hence the ode is

$$-2y + y' = 3e^t$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int (-2) dt} \\ &= e^{-2t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (3 e^t) \\ \frac{d}{dt}(e^{-2t}y) &= (e^{-2t}) (3 e^t) \\ d(e^{-2t}y) &= (3 e^{-t}) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-2t}y &= \int 3 e^{-t} dt \\ e^{-2t}y &= -3 e^{-t} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-2t}$ results in

$$y = -3 e^{2t} e^{-t} + c_1 e^{2t}$$

which simplifies to

$$y = -3 e^t + c_1 e^{2t}$$

Summary

The solution(s) found are the following

$$y = -3 e^t + c_1 e^{2t} \tag{1}$$

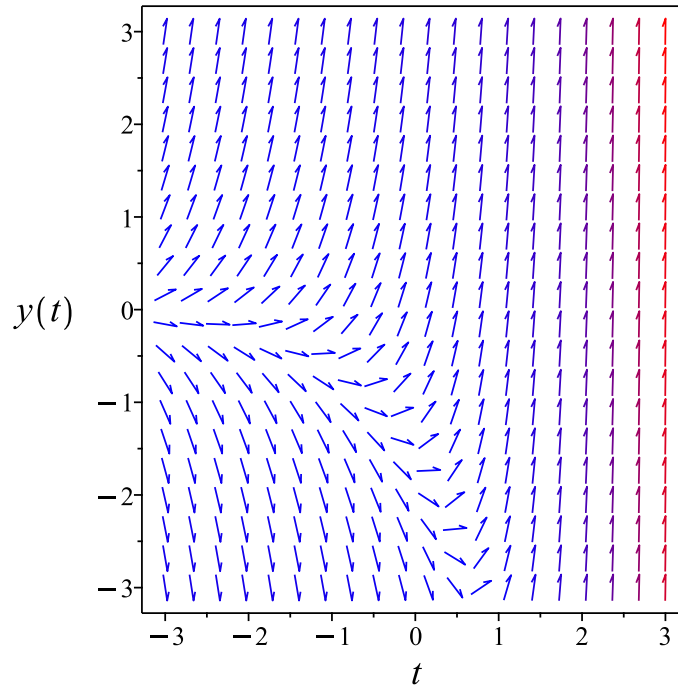


Figure 13: Slope field plot

Verification of solutions

$$y = -3e^t + c_1e^{2t}$$

Verified OK.

1.5.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = 2y + 3e^t$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2\xi_y - \omega_t\xi - \omega_y\eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 13: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{2t}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{2t}} dy \end{aligned}$$

Which results in

$$S = e^{-2t} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = 2y + 3e^t$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -2e^{-2t}y \\ S_y &= e^{-2t} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 3e^{-t} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 3e^{-R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -3e^{-R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{-2t}y = -3e^{-t} + c_1$$

Which simplifies to

$$e^{-2t}y = -3e^{-t} + c_1$$

Which gives

$$y = -(3e^{-t} - c_1)e^{2t}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = 2y + 3e^t$	$R = t$ $S = e^{-2t}y$	$\frac{dS}{dR} = 3e^{-R}$

Summary

The solution(s) found are the following

$$y = -(3e^{-t} - c_1)e^{2t} \quad (1)$$

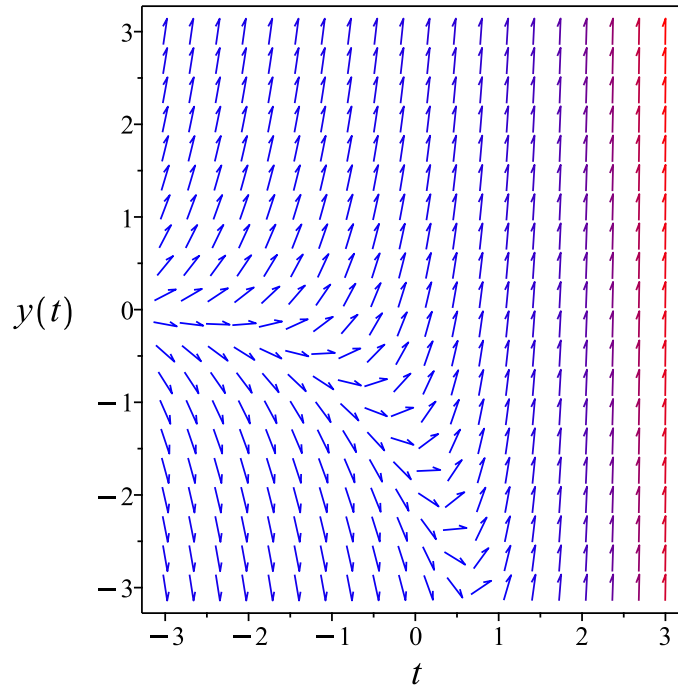


Figure 14: Slope field plot

Verification of solutions

$$y = -(3e^{-t} - c_1)e^{2t}$$

Verified OK.

1.5.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dy &= (2y + 3e^t) dt \\ (-2y - 3e^t) dt + dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -2y - 3e^t \\ N(t, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-2y - 3e^t) \\ &= -2\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1((-2) - (0)) \\ &= -2 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int -2 dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-2t} \\ &= e^{-2t} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{-2t}(-2y - 3e^t) \\ &= (-2y - 3e^t)e^{-2t} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{-2t}(1) \\ &= e^{-2t} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ ((-2y - 3e^t)e^{-2t}) + (e^{-2t}) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int (-2y - 3e^t) e^{-2t} dt \\ \phi &= 3e^{-t} + e^{-2t}y + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-2t} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-2t}$. Therefore equation (4) becomes

$$e^{-2t} = e^{-2t} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = 3e^{-t} + e^{-2t}y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = 3e^{-t} + e^{-2t}y$$

The solution becomes

$$y = -(3e^{-t} - c_1) e^{2t}$$

Summary

The solution(s) found are the following

$$y = -(3e^{-t} - c_1) e^{2t} \quad (1)$$

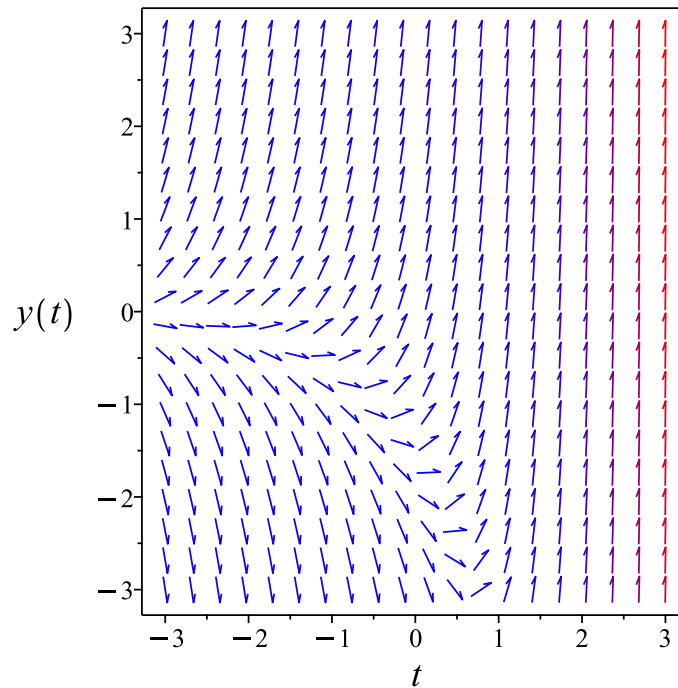


Figure 15: Slope field plot

Verification of solutions

$$y = -(3e^{-t} - c_1)e^{2t}$$

Verified OK.

1.5.4 Maple step by step solution

Let's solve

$$-2y + y' = 3e^t$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = 2y + 3e^t$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$-2y + y' = 3e^t$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t)(-2y + y') = 3\mu(t)e^t$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t) y)$

$$\mu(t) (-2y + y') = \mu'(t) y + \mu(t) y'$$
- Isolate $\mu'(t)$

$$\mu'(t) = -2\mu(t)$$
- Solve to find the integrating factor

$$\mu(t) = e^{-2t}$$
- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y) \right) dt = \int 3\mu(t) e^t dt + c_1$$
- Evaluate the integral on the lhs

$$\mu(t) y = \int 3\mu(t) e^t dt + c_1$$
- Solve for y

$$y = \frac{\int 3\mu(t) e^t dt + c_1}{\mu(t)}$$
- Substitute $\mu(t) = e^{-2t}$

$$y = \frac{\int 3e^t e^{-2t} dt + c_1}{e^{-2t}}$$
- Evaluate the integrals on the rhs

$$y = \frac{-3e^{-t} + c_1}{e^{-2t}}$$
- Simplify

$$y = -3e^t + c_1 e^{2t}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve(-2*y(t)+diff(y(t),t) = 3*exp(t),y(t), singsol=all)
```

$$y(t) = -3e^t + c_1e^{2t}$$

✓ Solution by Mathematica

Time used: 0.046 (sec). Leaf size: 17

```
DSolve[-2*y[t]+y'[t] == 3*Exp[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^t(-3 + c_1e^t)$$

1.6 problem 6

1.6.1	Solving as linear ode	67
1.6.2	Solving as first order ode lie symmetry lookup ode	69
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1.6.4	Maple step by step solution	78

Internal problem ID [453]

Internal file name [OUTPUT/453_Sunday_June_05_2022_01_41_43_AM_13381902/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 6.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$2y + ty' = \sin(t)$$

1.6.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{2}{t}$$
$$q(t) = \frac{\sin(t)}{t}$$

Hence the ode is

$$y' + \frac{2y}{t} = \frac{\sin(t)}{t}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{2}{t} dt} \\ &= t^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{\sin(t)}{t} \right) \\ \frac{d}{dt}(t^2 y) &= (t^2) \left(\frac{\sin(t)}{t} \right) \\ d(t^2 y) &= (t \sin(t)) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t^2 y &= \int t \sin(t) dt \\ t^2 y &= -t \cos(t) + \sin(t) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t^2$ results in

$$y = \frac{-t \cos(t) + \sin(t)}{t^2} + \frac{c_1}{t^2}$$

which simplifies to

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2}$$

Summary

The solution(s) found are the following

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2} \tag{1}$$

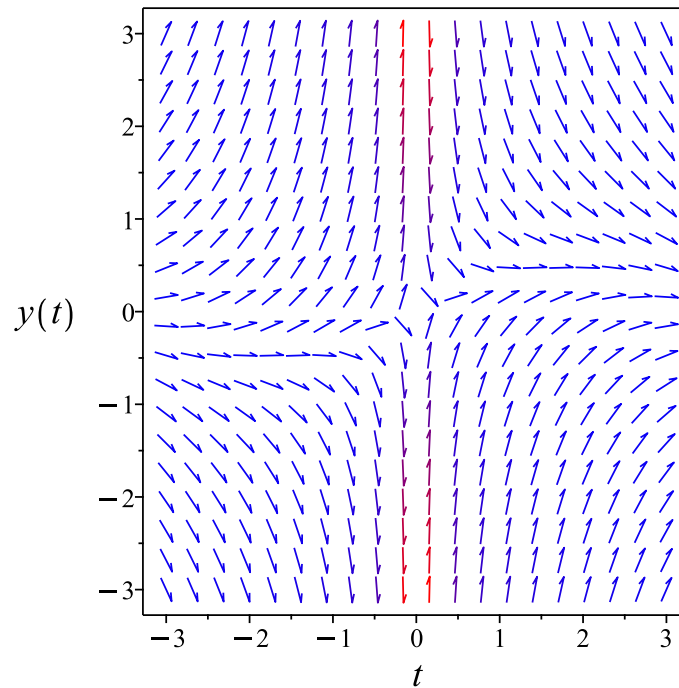


Figure 16: Slope field plot

Verification of solutions

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2}$$

Verified OK.

1.6.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-2y + \sin(t)}{t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 16: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \frac{1}{t^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{t^2}} dy \end{aligned}$$

Which results in

$$S = t^2 y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{-2y + \sin(t)}{t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= 2ty \\ S_y &= t^2 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = t \sin(t) \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R \sin(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -R \cos(R) + \sin(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$yt^2 = -t \cos(t) + \sin(t) + c_1$$

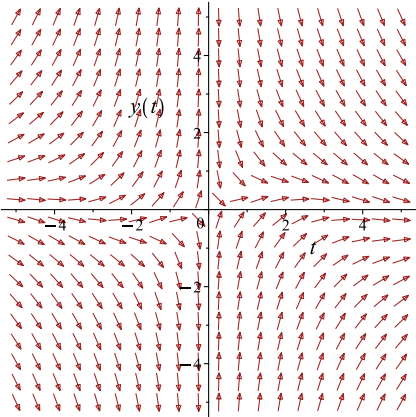
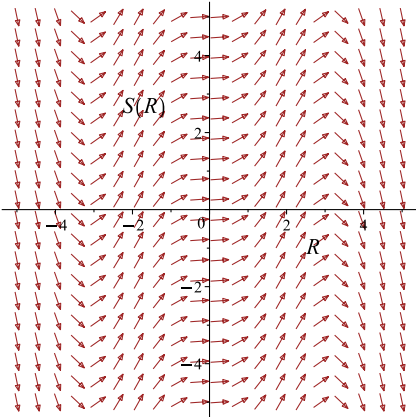
Which simplifies to

$$yt^2 = -t \cos(t) + \sin(t) + c_1$$

Which gives

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{-2y + \sin(t)}{t}$ 	$R = t$ $S = t^2 y$	$\frac{dS}{dR} = R \sin(R)$ 

Summary

The solution(s) found are the following

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2} \quad (1)$$

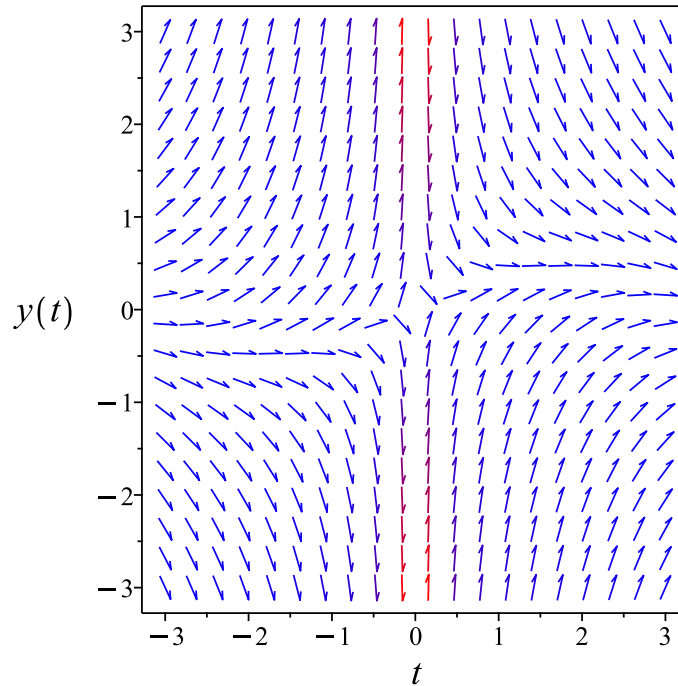


Figure 17: Slope field plot

Verification of solutions

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2}$$

Verified OK.

1.6.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(t) dy &= (-2y + \sin(t)) dt \\ (2y - \sin(t)) dt + (t) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= 2y - \sin(t) \\ N(t, y) &= t\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2y - \sin(t)) \\ &= 2\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(t) \\ &= 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{t} ((2) - (1)) \\ &= \frac{1}{t} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int \frac{1}{t} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\ln(t)} \\ &= t \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \overline{M} &= \mu M \\ &= t(2y - \sin(t)) \\ &= (2y - \sin(t)) t \end{aligned}$$

And

$$\begin{aligned} \overline{N} &= \mu N \\ &= t(t) \\ &= t^2 \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \overline{M} + \overline{N} \frac{dy}{dt} &= 0 \\ ((2y - \sin(t)) t) + (t^2) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int \bar{M} dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int (2y - \sin(t)) t dt$$

$$\phi = t^2 y + t \cos(t) - \sin(t) + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t^2 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t^2$. Therefore equation (4) becomes

$$t^2 = t^2 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = t^2 y + t \cos(t) - \sin(t) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = t^2 y + t \cos(t) - \sin(t)$$

The solution becomes

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2}$$

Summary

The solution(s) found are the following

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2} \tag{1}$$

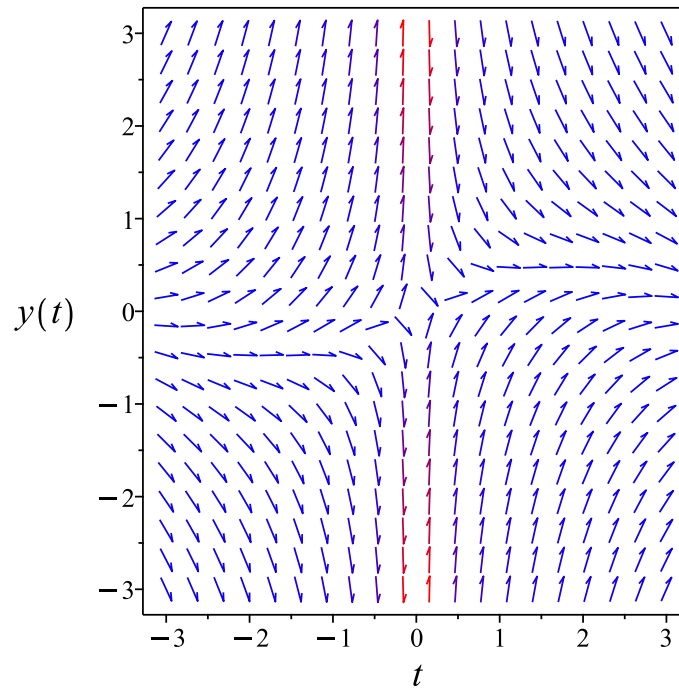


Figure 18: Slope field plot

Verification of solutions

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2}$$

Verified OK.

1.6.4 Maple step by step solution

Let's solve

$$2y + ty' = \sin(t)$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{2y}{t} + \frac{\sin(t)}{t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{2y}{t} = \frac{\sin(t)}{t}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \frac{\mu(t) \sin(t)}{t}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t) y)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \mu'(t) y + \mu(t) y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{2\mu(t)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = t^2$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y) \right) dt = \int \frac{\mu(t) \sin(t)}{t} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t) y = \int \frac{\mu(t) \sin(t)}{t} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t) \sin(t)}{t} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = t^2$

$$y = \frac{\int t \sin(t) dt + c_1}{t^2}$$

- Evaluate the integrals on the rhs

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(2*y(t)+t*diff(y(t),t) = sin(t),y(t), singsol=all)
```

$$y(t) = \frac{\sin(t) - \cos(t)t + c_1}{t^2}$$

✓ Solution by Mathematica

Time used: 0.039 (sec). Leaf size: 19

```
DSolve[2*y[t]+t*y'[t]== Sin[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{\sin(t) - t \cos(t) + c_1}{t^2}$$

1.7 problem 7

1.7.1	Solving as linear ode	80
1.7.2	Solving as first order ode lie symmetry lookup ode	82
1.7.3	Solving as exact ode	86
1.7.4	Maple step by step solution	90

Internal problem ID [454]

Internal file name [OUTPUT/454_Sunday_June_05_2022_01_41_44_AM_75941329/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 7.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "linear", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

`[_linear]`

$$2yt + y' = 2t e^{-t^2}$$

1.7.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = 2t$$

$$q(t) = 2t e^{-t^2}$$

Hence the ode is

$$2yt + y' = 2t e^{-t^2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int 2t dt} \\ &= e^{t^2}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (2t e^{-t^2}) \\ \frac{d}{dt}(y e^{t^2}) &= (e^{t^2}) (2t e^{-t^2}) \\ d(y e^{t^2}) &= (2t) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}y e^{t^2} &= \int 2t dt \\ y e^{t^2} &= t^2 + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{t^2}$ results in

$$y = t^2 e^{-t^2} + c_1 e^{-t^2}$$

which simplifies to

$$y = e^{-t^2} (t^2 + c_1)$$

Summary

The solution(s) found are the following

$$y = e^{-t^2} (t^2 + c_1) \tag{1}$$

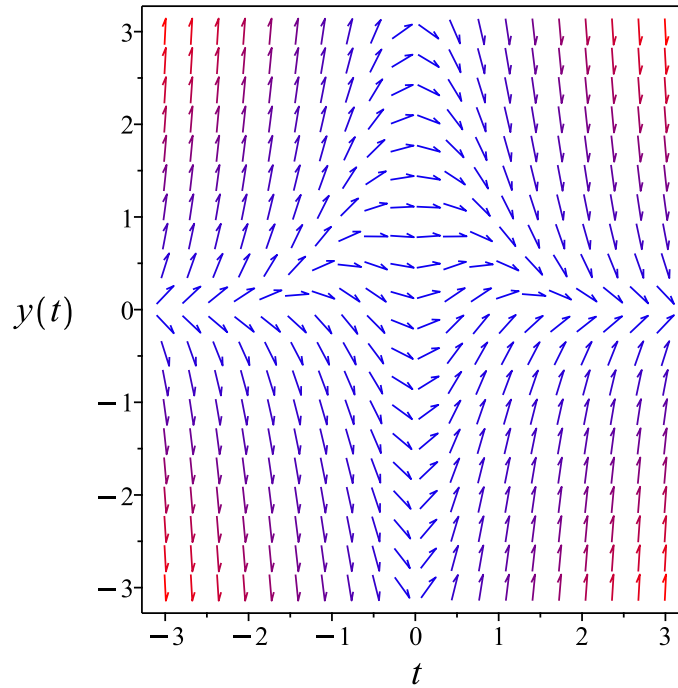


Figure 19: Slope field plot

Verification of solutions

$$y = e^{-t^2} (t^2 + c_1)$$

Verified OK.

1.7.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -2t(y e^{t^2} - 1) e^{-t^2}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 19: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-t^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-t^2}} dy \end{aligned}$$

Which results in

$$S = y e^{t^2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -2t(y e^{t^2} - 1) e^{-t^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= 2ty e^{t^2} \\ S_y &= e^{t^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 2t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 2R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R^2 + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$y e^{t^2} = t^2 + c_1$$

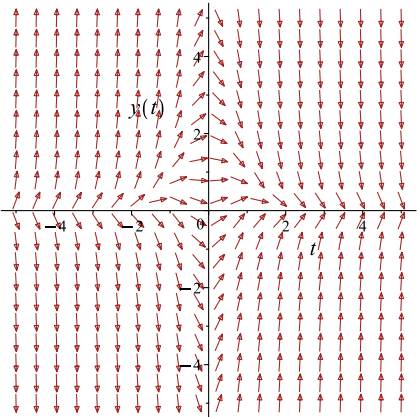
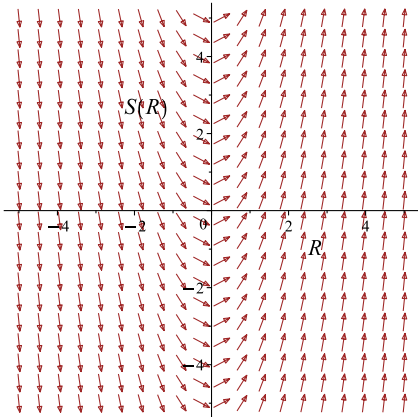
Which simplifies to

$$y e^{t^2} = t^2 + c_1$$

Which gives

$$y = e^{-t^2} (t^2 + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -2t(y e^{t^2} - 1) e^{-t^2}$ 	$R = t$ $S = y e^{t^2}$	$\frac{dS}{dR} = 2R$ 

Summary

The solution(s) found are the following

$$y = e^{-t^2} (t^2 + c_1) \quad (1)$$

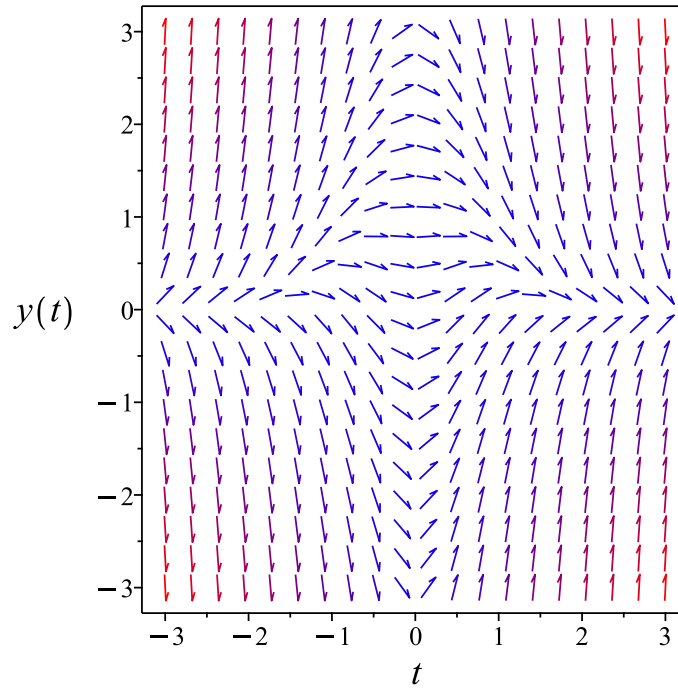


Figure 20: Slope field plot

Verification of solutions

$$y = e^{-t^2} (t^2 + c_1)$$

Verified OK.

1.7.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} & \left(e^{t^2} \right) dy = \left(-2t \left(y e^{t^2} - 1 \right) \right) dt \\ \left(2t \left(y e^{t^2} - 1 \right) \right) dt + \left(e^{t^2} \right) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= 2t \left(y e^{t^2} - 1 \right) \\ N(t, y) &= e^{t^2} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(2t \left(y e^{t^2} - 1 \right) \right) \\ &= 2t e^{t^2}\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (e^{t^2}) \\ &= 2t e^{t^2}\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int 2t(y e^{t^2} - 1) dt \\ \phi &= -t^2 + y e^{t^2} + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{t^2} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{t^2}$. Therefore equation (4) becomes

$$e^{t^2} = e^{t^2} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -t^2 + y e^{t^2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -t^2 + y e^{t^2}$$

The solution becomes

$$y = e^{-t^2} (t^2 + c_1)$$

Summary

The solution(s) found are the following

$$y = e^{-t^2} (t^2 + c_1) \tag{1}$$

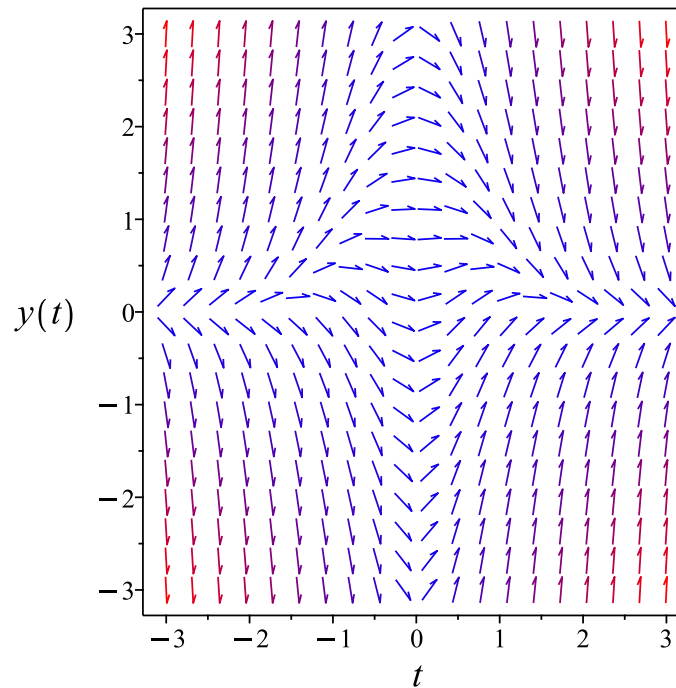


Figure 21: Slope field plot

Verification of solutions

$$y = e^{-t^2} (t^2 + c_1)$$

Verified OK.

1.7.4 Maple step by step solution

Let's solve

$$2yt + y' = \frac{2t}{e^{t^2}}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -2yt + \frac{2t}{e^{t^2}}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$2yt + y' = \frac{2t}{e^{t^2}}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) (2yt + y') = \frac{2\mu(t)t}{e^{t^2}}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t) y)$

$$\mu(t) (2yt + y') = \mu'(t) y + \mu(t) y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = 2\mu(t) t$$

- Solve to find the integrating factor

$$\mu(t) = \left(e^{t^2}\right)^2 e^{-t^2}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y)\right) dt = \int \frac{2\mu(t)t}{e^{t^2}} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t) y = \int \frac{2\mu(t)t}{e^{t^2}} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{2\mu(t)t}{e^{t^2}} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = \left(e^{t^2}\right)^2 e^{-t^2}$

$$y = \frac{\int 2t e^{-t^2} e^{t^2} dt + c_1}{\left(e^{t^2}\right)^2 e^{-t^2}}$$

- Evaluate the integrals on the rhs

$$y = \frac{t^2 + c_1}{(e^{t^2})^2 e^{-t^2}}$$

- Simplify

$$y = e^{-t^2} (t^2 + c_1)$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 16

```
dsolve(2*t*y(t)+diff(y(t),t) = 2*t/exp(t^2),y(t), singsol=all)
```

$$y(t) = (t^2 + c_1) e^{-t^2}$$

✓ Solution by Mathematica

Time used: 0.055 (sec). Leaf size: 19

```
DSolve[2*t*y[t]+y'[t] == 2*t/Exp[t^2],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{-t^2} (t^2 + c_1)$$

1.8 problem 8

1.8.1	Solving as linear ode	92
1.8.2	Solving as first order ode lie symmetry lookup ode	94
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Internal problem ID [455]

Internal file name [OUTPUT/455_Sunday_June_05_2022_01_41_45_AM_39090059/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 8.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$4yt + (t^2 + 1)y' = \frac{1}{(t^2 + 1)^2}$$

1.8.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{4t}{t^2 + 1}$$
$$q(t) = \frac{1}{(t^2 + 1)^3}$$

Hence the ode is

$$y' + \frac{4yt}{t^2 + 1} = \frac{1}{(t^2 + 1)^3}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{-4t}{t^2+1} dt} \\ &= (t^2 + 1)^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{1}{(t^2 + 1)^3} \right) \\ \frac{d}{dt}((t^2 + 1)^2 y) &= ((t^2 + 1)^2) \left(\frac{1}{(t^2 + 1)^3} \right) \\ d((t^2 + 1)^2 y) &= \frac{1}{t^2 + 1} dt\end{aligned}$$

Integrating gives

$$\begin{aligned}(t^2 + 1)^2 y &= \int \frac{1}{t^2 + 1} dt \\ (t^2 + 1)^2 y &= \arctan(t) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = (t^2 + 1)^2$ results in

$$y = \frac{\arctan(t)}{(t^2 + 1)^2} + \frac{c_1}{(t^2 + 1)^2}$$

which simplifies to

$$y = \frac{\arctan(t) + c_1}{(t^2 + 1)^2}$$

Summary

The solution(s) found are the following

$$y = \frac{\arctan(t) + c_1}{(t^2 + 1)^2} \tag{1}$$

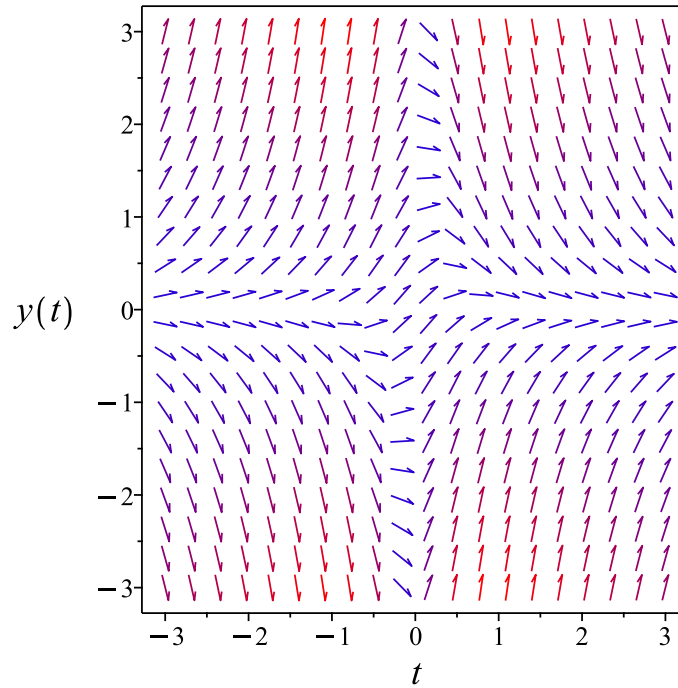


Figure 22: Slope field plot

Verification of solutions

$$y = \frac{\arctan(t) + c_1}{(t^2 + 1)^2}$$

Verified OK.

1.8.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{4t^5y + 8t^3y + 4ty - 1}{(t^2 + 1)^3}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2\xi_y - \omega_t\xi - \omega_y\eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 22: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \frac{1}{(t^2 + 1)^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{(t^2+1)^2}} dy \end{aligned}$$

Which results in

$$S = (t^2 + 1)^2 y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{4t^5 y + 8t^3 y + 4ty - 1}{(t^2 + 1)^3}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= 4(t^2 + 1) yt \\ S_y &= (t^2 + 1)^2 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{t^2 + 1} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R^2 + 1}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \arctan(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$(t^2 + 1)^2 y = \arctan(t) + c_1$$

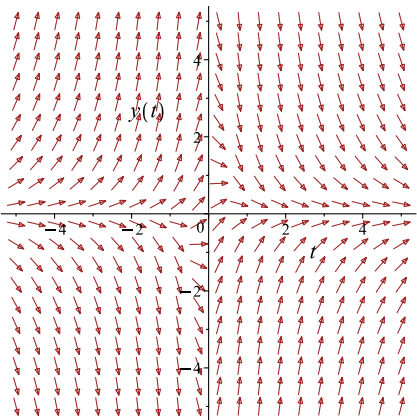
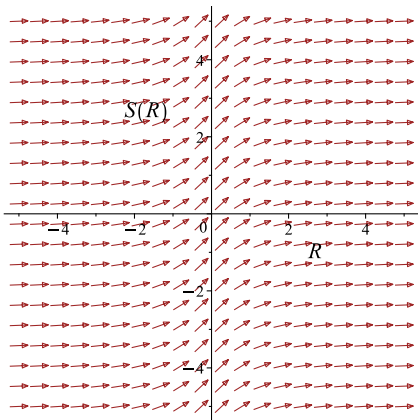
Which simplifies to

$$(t^2 + 1)^2 y = \arctan(t) + c_1$$

Which gives

$$y = \frac{\arctan(t) + c_1}{(t^2 + 1)^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{4t^5 y + 8t^3 y + 4ty - 1}{(t^2 + 1)^3}$ 	$R = t$ $S = (t^2 + 1)^2 y$	$\frac{dS}{dR} = \frac{1}{R^2 + 1}$ 

Summary

The solution(s) found are the following

$$y = \frac{\arctan(t) + c_1}{(t^2 + 1)^2} \quad (1)$$

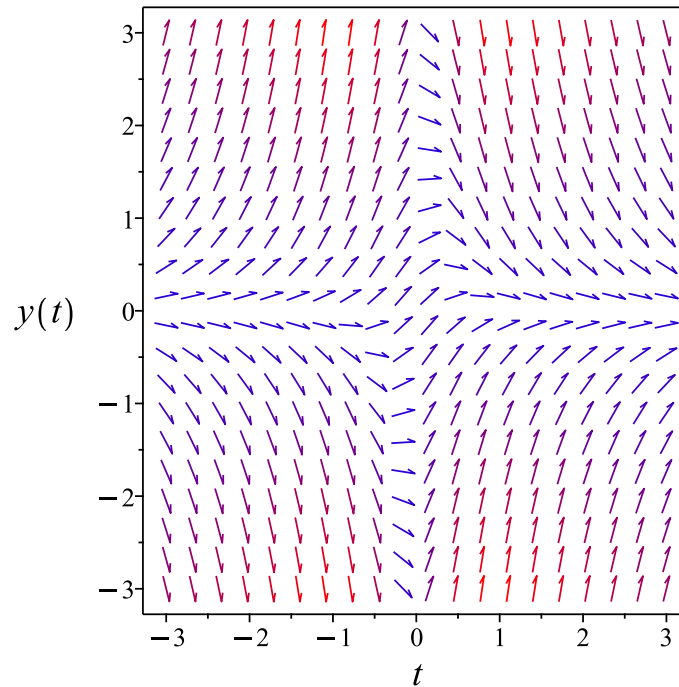


Figure 23: Slope field plot

Verification of solutions

$$y = \frac{\arctan(t) + c_1}{(t^2 + 1)^2}$$

Verified OK.

1.8.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the

ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (t^2 + 1) dy &= \left(-4ty + \frac{1}{(t^2 + 1)^2} \right) dt \\ \left(4ty - \frac{1}{(t^2 + 1)^2} \right) dt + (t^2 + 1) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= 4ty - \frac{1}{(t^2 + 1)^2} \\ N(t, y) &= t^2 + 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(4ty - \frac{1}{(t^2 + 1)^2} \right) \\ &= 4t\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (t^2 + 1) \\ &= 2t\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{t^2 + 1} ((4t) - (2t)) \\ &= \frac{2t}{t^2 + 1}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int \frac{2t}{t^2+1} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{\ln(t^2+1)} \\ &= t^2 + 1\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= t^2 + 1 \left(4ty - \frac{1}{(t^2 + 1)^2} \right) \\ &= \left(4ty - \frac{1}{(t^2 + 1)^2} \right) (t^2 + 1)\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= t^2 + 1(t^2 + 1) \\ &= (t^2 + 1)^2\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left(\left(4ty - \frac{1}{(t^2 + 1)^2} \right) (t^2 + 1) \right) + \left((t^2 + 1)^2 \right) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int \left(4ty - \frac{1}{(t^2 + 1)^2} \right) (t^2 + 1) dt \\ \phi &= t^4 y + 2t^2 y - \arctan(t) + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t^4 + 2t^2 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = (t^2 + 1)^2$. Therefore equation (4) becomes

$$(t^2 + 1)^2 = t^4 + 2t^2 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 1$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (1) dy$$

$$f(y) = y + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = t^4 y + 2t^2 y - \arctan(t) + y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = t^4 y + 2t^2 y - \arctan(t) + y$$

Summary

The solution(s) found are the following

$$yt^4 + 2yt^2 - \arctan(t) + y = c_1 \tag{1}$$

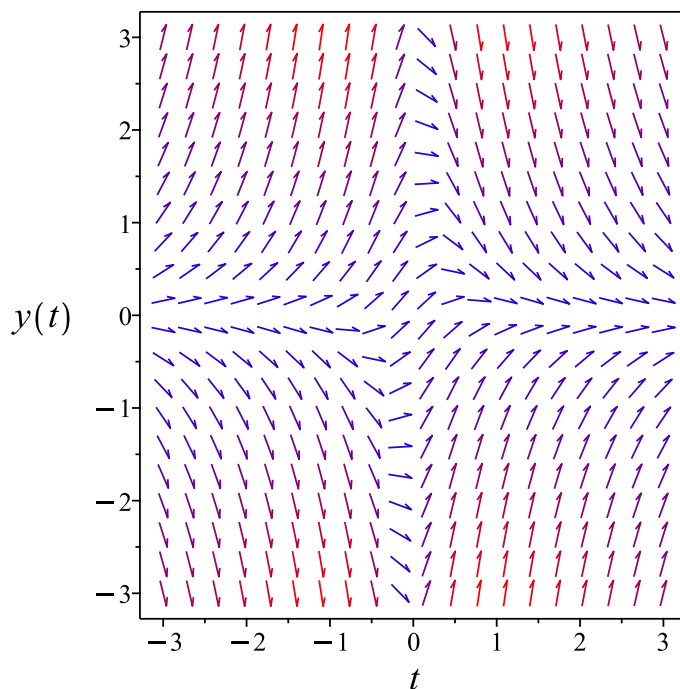


Figure 24: Slope field plot

Verification of solutions

$$yt^4 + 2yt^2 - \arctan(t) + y = c_1$$

Verified OK.

1.8.4 Maple step by step solution

Let's solve

$$4yt + (t^2 + 1)y' = \frac{1}{(t^2+1)^2}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{4yt}{t^2+1} + \frac{1}{(t^2+1)^3}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{4yt}{t^2+1} = \frac{1}{(t^2+1)^3}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{4yt}{t^2+1} \right) = \frac{\mu(t)}{(t^2+1)^3}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{4yt}{t^2+1} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{4\mu(t)t}{t^2+1}$$

- Solve to find the integrating factor

$$\mu(t) = (t^2 + 1)^2$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)}{(t^2+1)^3} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)}{(t^2+1)^3} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)}{(t^2+1)^3} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = (t^2 + 1)^2$

$$y = \frac{\int \frac{1}{t^2+1} dt + c_1}{(t^2+1)^2}$$
- Evaluate the integrals on the rhs

$$y = \frac{\arctan(t) + c_1}{(t^2+1)^2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 16

```
dsolve(4*t*y(t)+(t^2+1)*diff(y(t),t) = 1/(t^2+1)^2,y(t), singsol=all)
```

$$y(t) = \frac{\arctan(t) + c_1}{(t^2 + 1)^2}$$

✓ Solution by Mathematica

Time used: 0.041 (sec). Leaf size: 18

```
DSolve[4*t*y[t]+(t^2+1)*y'[t] == 1/(t^2+1)^2,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{\arctan(t) + c_1}{(t^2 + 1)^2}$$

1.9 problem 9

1.9.1	Solving as linear ode	105
1.9.2	Solving as first order ode lie symmetry lookup ode	107
1.9.3	Solving as exact ode	111
1.9.4	Maple step by step solution	116

Internal problem ID [456]

Internal file name [OUTPUT/456_Sunday_June_05_2022_01_41_46_AM_4296471/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 9.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$y + 2y' = 3t$$

1.9.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{1}{2}$$
$$q(t) = \frac{3t}{2}$$

Hence the ode is

$$y' + \frac{y}{2} = \frac{3t}{2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{1}{2} dt} \\ &= e^{\frac{t}{2}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{3t}{2} \right) \\ \frac{d}{dt} \left(e^{\frac{t}{2}} y \right) &= \left(e^{\frac{t}{2}} \right) \left(\frac{3t}{2} \right) \\ d \left(e^{\frac{t}{2}} y \right) &= \left(\frac{3t e^{\frac{t}{2}}}{2} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{\frac{t}{2}} y &= \int \frac{3t e^{\frac{t}{2}}}{2} dt \\ e^{\frac{t}{2}} y &= 3(t-2) e^{\frac{t}{2}} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{\frac{t}{2}}$ results in

$$y = 3 e^{-\frac{t}{2}} (t-2) e^{\frac{t}{2}} + c_1 e^{-\frac{t}{2}}$$

which simplifies to

$$y = 3t - 6 + c_1 e^{-\frac{t}{2}}$$

Summary

The solution(s) found are the following

$$y = 3t - 6 + c_1 e^{-\frac{t}{2}} \tag{1}$$

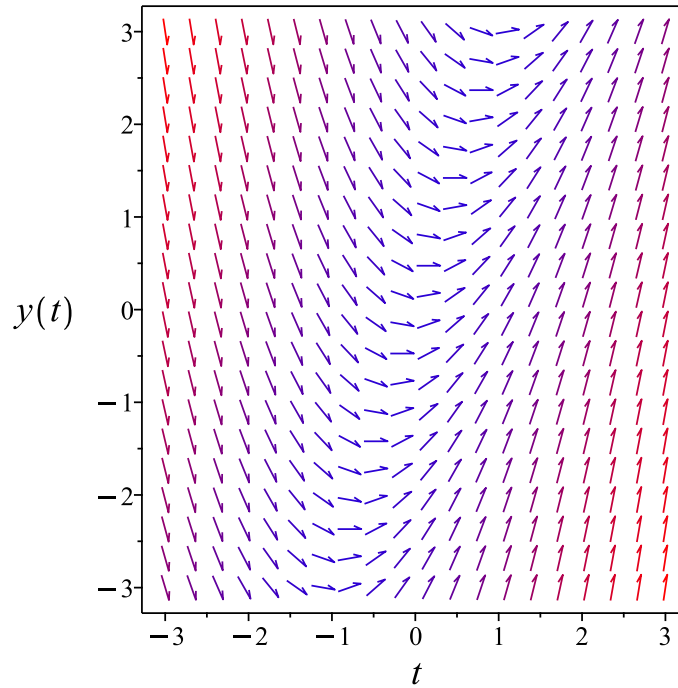


Figure 25: Slope field plot

Verification of solutions

$$y = 3t - 6 + c_1 e^{-\frac{t}{2}}$$

Verified OK.

1.9.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{y}{2} + \frac{3t}{2}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 25: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-\frac{t}{2}}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-\frac{t}{2}}} dy \end{aligned}$$

Which results in

$$S = e^{\frac{t}{2}} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{y}{2} + \frac{3t}{2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= \frac{e^{\frac{t}{2}} y}{2} \\ S_y &= e^{\frac{t}{2}} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{3t e^{\frac{t}{2}}}{2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{3R e^{\frac{R}{2}}}{2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = 3(R - 2)e^{\frac{R}{2}} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{\frac{t}{2}}y = 3(t - 2)e^{\frac{t}{2}} + c_1$$

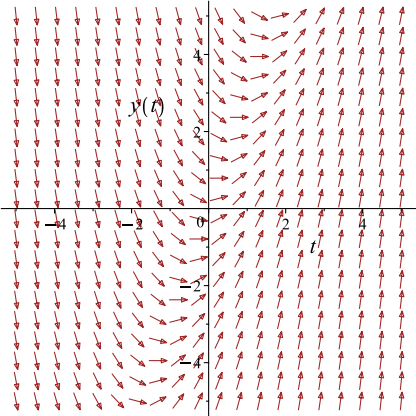
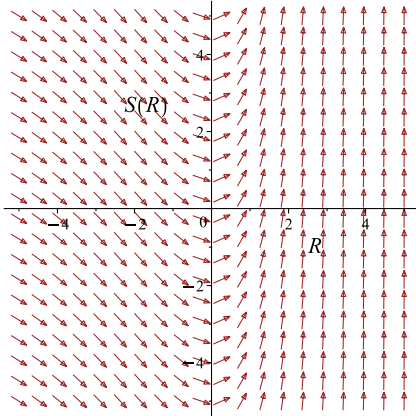
Which simplifies to

$$(-3t + y + 6)e^{\frac{t}{2}} - c_1 = 0$$

Which gives

$$y = \left(3te^{\frac{t}{2}} - 6e^{\frac{t}{2}} + c_1\right)e^{-\frac{t}{2}}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{y}{2} + \frac{3t}{2}$ 	$R = t$ $S = e^{\frac{t}{2}}y$	$\frac{dS}{dR} = \frac{3Re^{\frac{R}{2}}}{2}$ 

Summary

The solution(s) found are the following

$$y = \left(3te^{\frac{t}{2}} - 6e^{\frac{t}{2}} + c_1\right)e^{-\frac{t}{2}} \quad (1)$$

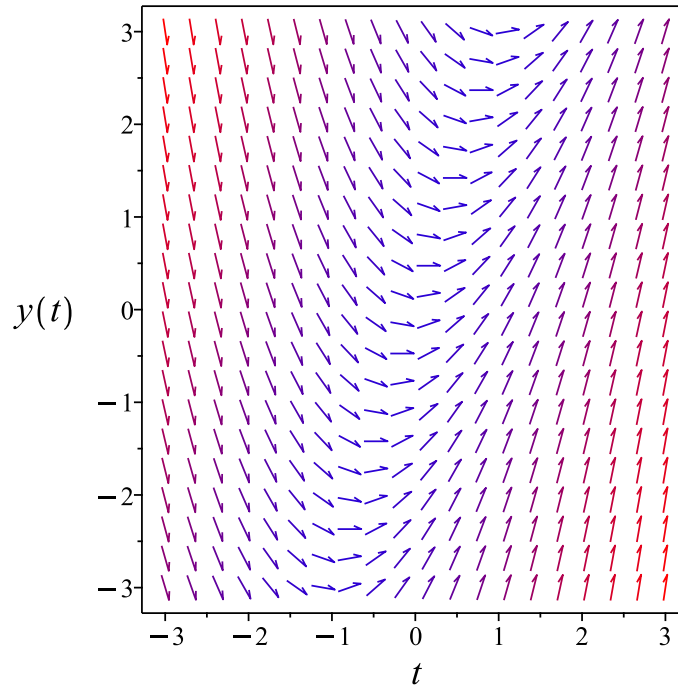


Figure 26: Slope field plot

Verification of solutions

$$y = \left(3t e^{\frac{t}{2}} - 6 e^{\frac{t}{2}} + c_1 \right) e^{-\frac{t}{2}}$$

Verified OK.

1.9.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(2) dy &= (-y + 3t) dt \\ (y - 3t) dt + (2) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= y - 3t \\ N(t, y) &= 2\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y - 3t) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(2) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{2} ((1) - (0)) \\ &= \frac{1}{2} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int \frac{1}{2} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\frac{t}{2}} \\ &= e^{\frac{t}{2}} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{\frac{t}{2}}(y - 3t) \\ &= (y - 3t) e^{\frac{t}{2}} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{\frac{t}{2}}(2) \\ &= 2 e^{\frac{t}{2}} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left((y - 3t) e^{\frac{t}{2}} \right) + \left(2 e^{\frac{t}{2}} \right) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int \overline{M} dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int (y - 3t) e^{\frac{t}{2}} dt$$

$$\phi = (-6t + 2y + 12) e^{\frac{t}{2}} + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 2 e^{\frac{t}{2}} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 2 e^{\frac{t}{2}}$. Therefore equation (4) becomes

$$2 e^{\frac{t}{2}} = 2 e^{\frac{t}{2}} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = (-6t + 2y + 12) e^{\frac{t}{2}} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = (-6t + 2y + 12) e^{\frac{t}{2}}$$

The solution becomes

$$y = \frac{\left(6t e^{\frac{t}{2}} - 12 e^{\frac{t}{2}} + c_1\right) e^{-\frac{t}{2}}}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(6t e^{\frac{t}{2}} - 12 e^{\frac{t}{2}} + c_1\right) e^{-\frac{t}{2}}}{2} \tag{1}$$

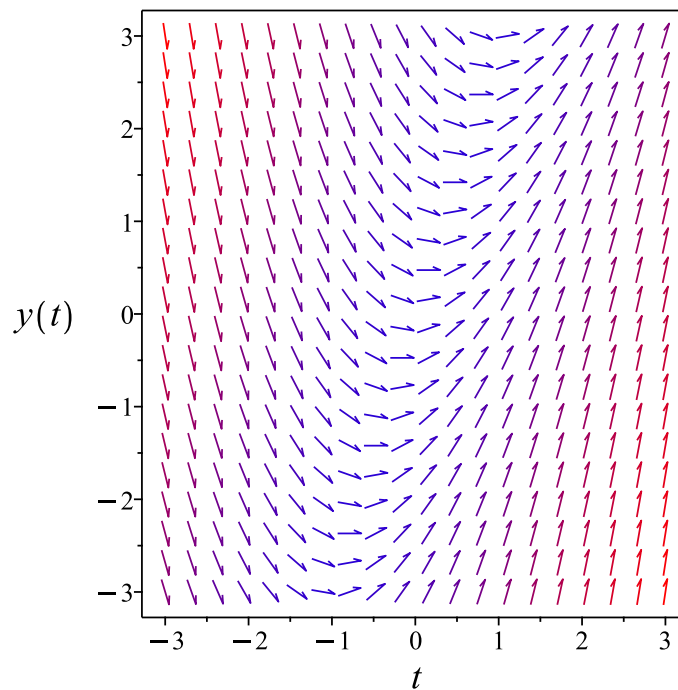


Figure 27: Slope field plot

Verification of solutions

$$y = \frac{\left(6t e^{\frac{t}{2}} - 12 e^{\frac{t}{2}} + c_1\right) e^{-\frac{t}{2}}}{2}$$

Verified OK.

1.9.4 Maple step by step solution

Let's solve

$$y + 2y' = 3t$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{y}{2} + \frac{3t}{2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{y}{2} = \frac{3t}{2}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{y}{2} \right) = \frac{3\mu(t)t}{2}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{y}{2} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)}{2}$$

- Solve to find the integrating factor

$$\mu(t) = e^{\frac{t}{2}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{3\mu(t)t}{2} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{3\mu(t)t}{2} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{3\mu(t)t}{2} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{\frac{t}{2}}$

$$y = \frac{\int \frac{3te^{\frac{t}{2}}}{2} dt + c_1}{e^{\frac{t}{2}}}$$

- Evaluate the integrals on the rhs

$$y = \frac{3(t-2)e^{\frac{t}{2}} + c_1}{e^{\frac{t}{2}}}$$

- Simplify

$$y = 3t - 6 + c_1 e^{-\frac{t}{2}}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve(y(t)+2*diff(y(t),t) = 3*t,y(t), singsol=all)
```

$$y(t) = 3t - 6 + e^{-\frac{t}{2}}c_1$$

✓ Solution by Mathematica

Time used: 0.045 (sec). Leaf size: 20

```
DSolve[y[t]+2*y'[t] == 3*t,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow 3t + c_1 e^{-t/2} - 6$$

1.10 problem 10

1.10.1 Solving as linear ode	118
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Internal problem ID [457]

Internal file name [OUTPUT/457_Sunday_June_05_2022_01_41_47_AM_99710545/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 10.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear"**, **"homogeneousTypeD2"**, **"exactWithIntegrationFactor"**, **"first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$ty' - y = t^2e^{-t}$$

1.10.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{1}{t}$$
$$q(t) = te^{-t}$$

Hence the ode is

$$y' - \frac{y}{t} = te^{-t}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{1}{t} dt} \\ &= \frac{1}{t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu)(t e^{-t}) \\ \frac{d}{dt}\left(\frac{y}{t}\right) &= \left(\frac{1}{t}\right)(t e^{-t}) \\ d\left(\frac{y}{t}\right) &= e^{-t} dt\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y}{t} &= \int e^{-t} dt \\ \frac{y}{t} &= -e^{-t} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = \frac{1}{t}$ results in

$$y = -t e^{-t} + c_1 t$$

which simplifies to

$$y = t(-e^{-t} + c_1)$$

Summary

The solution(s) found are the following

$$y = t(-e^{-t} + c_1) \tag{1}$$

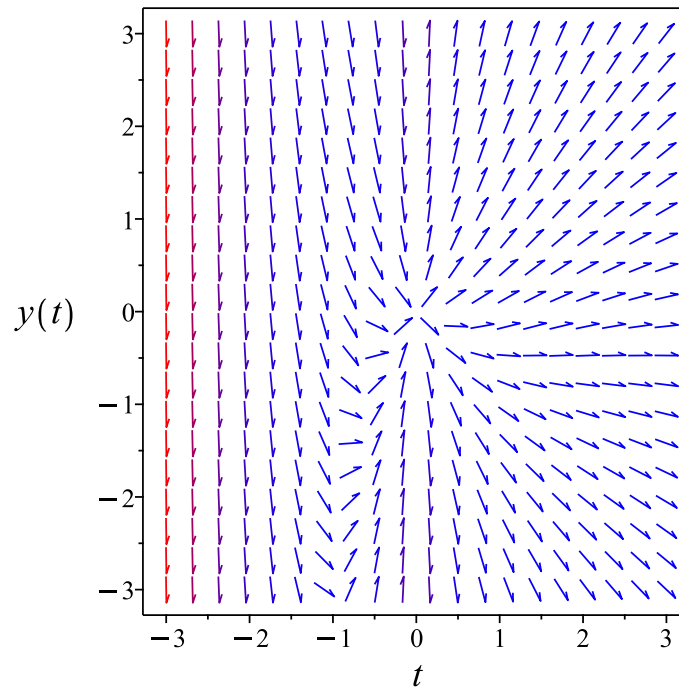


Figure 28: Slope field plot

Verification of solutions

$$y = t(-e^{-t} + c_1)$$

Verified OK.

1.10.2 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(t)t$ on the above ode results in new ode in $u(t)$

$$t(u'(t)t + u(t)) - u(t)t = t^2e^{-t}$$

Integrating both sides gives

$$\begin{aligned} u(t) &= \int e^{-t} dt \\ &= -e^{-t} + c_2 \end{aligned}$$

Therefore the solution y is

$$\begin{aligned} y &= ut \\ &= t(-e^{-t} + c_2) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = t(-e^{-t} + c_2) \quad (1)$$

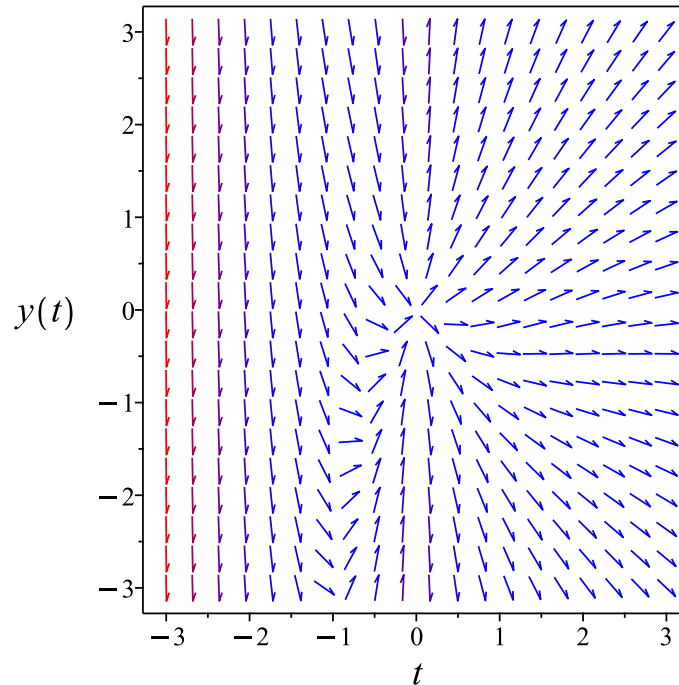


Figure 29: Slope field plot

Verification of solutions

$$y = t(-e^{-t} + c_2)$$

Verified OK.

1.10.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{(y e^t + t^2) e^{-t}}{t}$$
$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (A)$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 28: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= t\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{t} dy \end{aligned}$$

Which results in

$$S = \frac{y}{t}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{(y e^t + t^2) e^{-t}}{t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -\frac{y}{t^2} \\ S_y &= \frac{1}{t} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^{-t} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^{-R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -e^{-R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$\frac{y}{t} = -e^{-t} + c_1$$

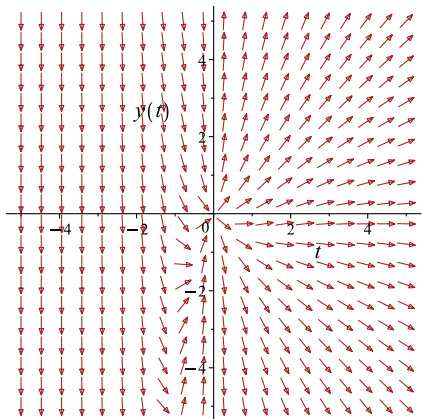
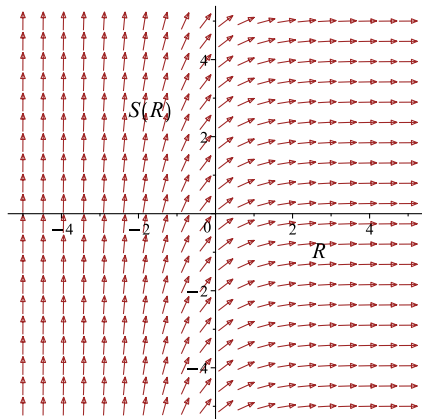
Which simplifies to

$$\frac{y}{t} = -e^{-t} + c_1$$

Which gives

$$y = -t(e^{-t} - c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{(ye^t + t^2)e^{-t}}{t}$ 	$R = t$ $S = \frac{y}{t}$	$\frac{dS}{dR} = e^{-R}$ 

Summary

The solution(s) found are the following

$$y = -t(e^{-t} - c_1) \quad (1)$$

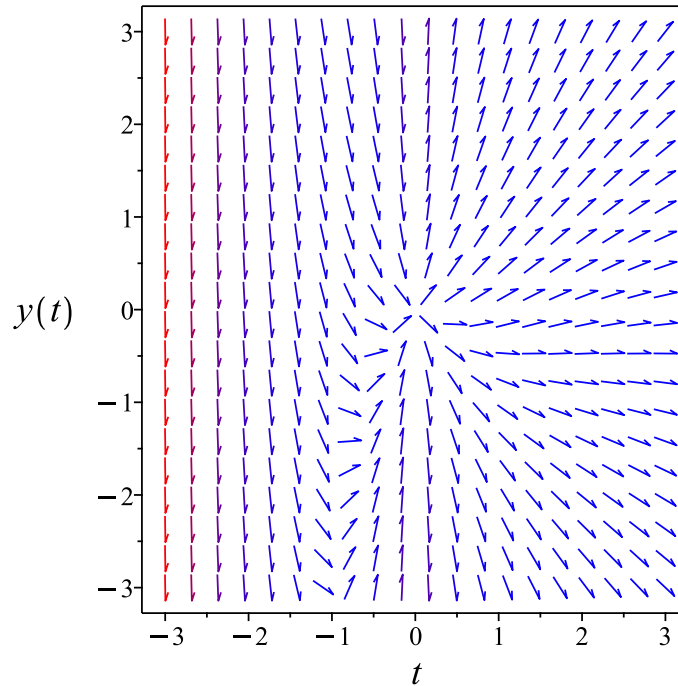


Figure 30: Slope field plot

Verification of solutions

$$y = -t(e^{-t} - c_1)$$

Verified OK.

1.10.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(t) dy &= (t^2 e^{-t} + y) dt \\ (-t^2 e^{-t} - y) dt + (t) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -t^2 e^{-t} - y \\ N(t, y) &= t\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (-t^2 e^{-t} - y) \\ &= -1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (t) \\ &= 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{t} ((-1) - (1)) \\ &= -\frac{2}{t} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int -\frac{2}{t} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-2 \ln(t)} \\ &= \frac{1}{t^2} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{1}{t^2} (-t^2 e^{-t} - y) \\ &= \frac{-t^2 e^{-t} - y}{t^2} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \frac{1}{t^2} (t) \\ &= \frac{1}{t} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left(\frac{-t^2 e^{-t} - y}{t^2} \right) + \left(\frac{1}{t} \right) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial t} dt &= \int \overline{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int \frac{-t^2 e^{-t} - y}{t^2} dt \\ \phi &= \frac{t e^{-t} + y}{t} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = \frac{1}{t} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{t}$. Therefore equation (4) becomes

$$\frac{1}{t} = \frac{1}{t} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{t e^{-t} + y}{t} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{t e^{-t} + y}{t}$$

The solution becomes

$$y = -t(e^{-t} - c_1)$$

Summary

The solution(s) found are the following

$$y = -t(e^{-t} - c_1) \tag{1}$$

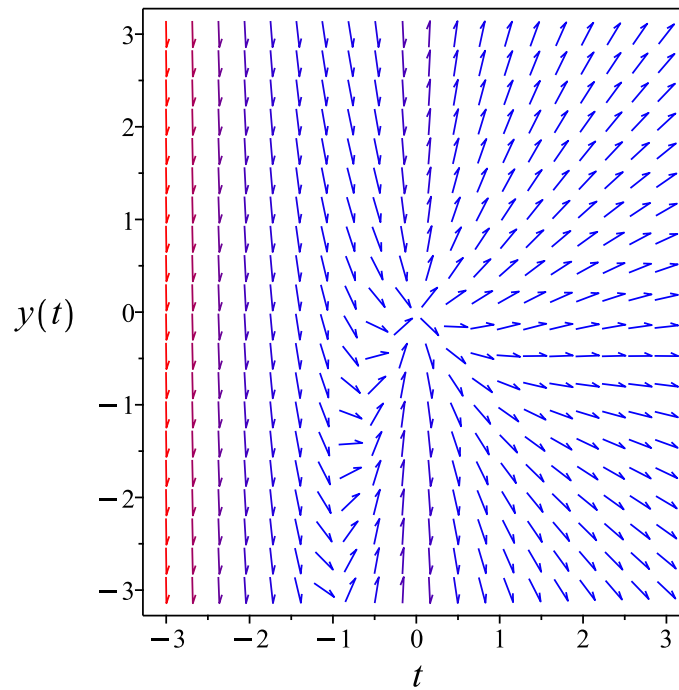


Figure 31: Slope field plot

Verification of solutions

$$y = -t(e^{-t} - c_1)$$

Verified OK.

1.10.5 Maple step by step solution

Let's solve

$$ty' - y = \frac{t^2}{e^t}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = \frac{y}{t} + \frac{t}{e^t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' - \frac{y}{t} = \frac{t}{e^t}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' - \frac{y}{t} \right) = \frac{\mu(t)t}{e^t}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' - \frac{y}{t} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -\frac{\mu(t)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = \frac{e^{-t}e^t}{t}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)t}{e^t} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)t}{e^t} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)t}{e^t} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = \frac{e^{-t}e^t}{t}$

$$y = \frac{t(\int e^{-t} dt + c_1)}{e^{-t}e^t}$$

- Evaluate the integrals on the rhs

$$y = \frac{t(-e^{-t} + c_1)}{e^{-t}e^t}$$

- Simplify

$$y = t(-e^{-t} + c_1)$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(-y(t)+t*diff(y(t),t) = t^2/exp(t),y(t), singsol=all)
```

$$y(t) = (-e^{-t} + c_1) t$$

✓ Solution by Mathematica

Time used: 0.053 (sec). Leaf size: 17

```
DSolve[-y[t]+t*y'[t] == t^2/Exp[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t(-e^{-t} + c_1)$$

1.11 problem 11

1.11.1 Solving as linear ode	132
1.11.2 Solving as first order ode lie symmetry lookup ode	134
1.11.3 Solving as exact ode	138
1.11.4 Maple step by step solution	142

Internal problem ID [458]

Internal file name [OUTPUT/458_Sunday_June_05_2022_01_41_48_AM_82720145/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 11.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$y + y' = 5 \sin(2t)$$

1.11.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = 1$$

$$q(t) = 5 \sin(2t)$$

Hence the ode is

$$y + y' = 5 \sin(2t)$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int 1 dt} \\ &= e^t\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (5 \sin(2t)) \\ \frac{d}{dt}(y e^t) &= (e^t) (5 \sin(2t)) \\ d(y e^t) &= (5 e^t \sin(2t)) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}y e^t &= \int 5 e^t \sin(2t) dt \\ y e^t &= -2 e^t \cos(2t) + e^t \sin(2t) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^t$ results in

$$y = e^{-t}(-2 e^t \cos(2t) + e^t \sin(2t)) + c_1 e^{-t}$$

which simplifies to

$$y = \sin(2t) - 2 \cos(2t) + c_1 e^{-t}$$

Summary

The solution(s) found are the following

$$y = \sin(2t) - 2 \cos(2t) + c_1 e^{-t} \tag{1}$$

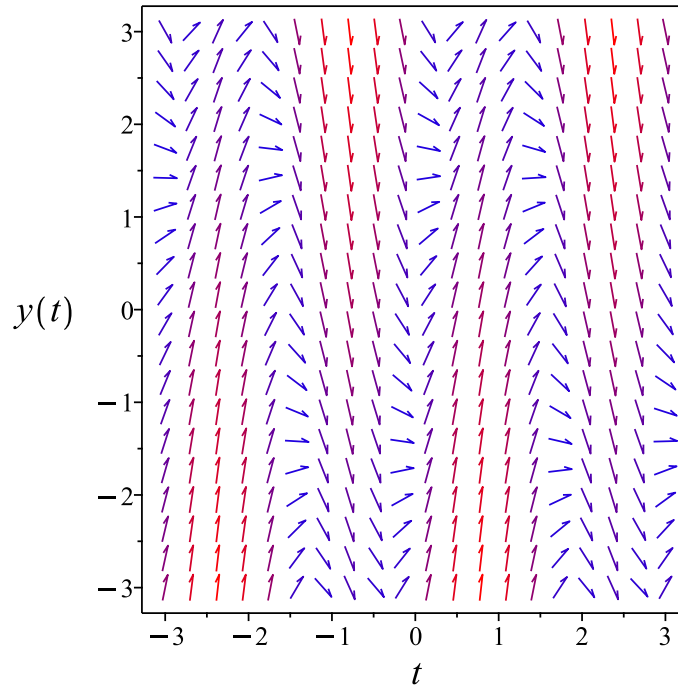


Figure 32: Slope field plot

Verification of solutions

$$y = \sin(2t) - 2 \cos(2t) + c_1 e^{-t}$$

Verified OK.

1.11.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -y + 5 \sin(2t)$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 31: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-t}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-t}} dy \end{aligned}$$

Which results in

$$S = y e^t$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -y + 5 \sin(2t)$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= y e^t \\ S_y &= e^t \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 5 e^t \sin(2t) \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 5 e^R \sin(2R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 - e^R(2 \cos(2R) - \sin(2R)) \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$y e^t = c_1 - e^t(2 \cos(2t) - \sin(2t))$$

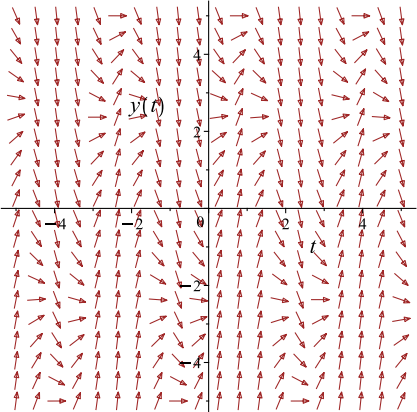
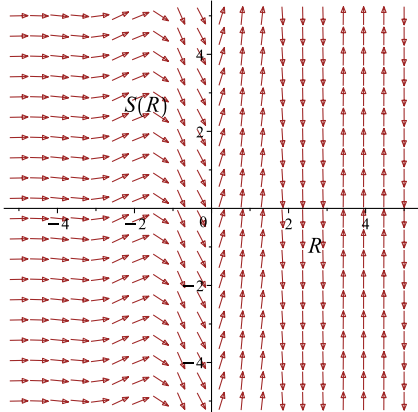
Which simplifies to

$$y e^t = c_1 - e^t(2 \cos(2t) - \sin(2t))$$

Which gives

$$y = e^{-t}(e^t \sin(2t) - 2 e^t \cos(2t) + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -y + 5 \sin(2t)$ 	$R = t$ $S = y e^t$	$\frac{dS}{dR} = 5 e^R \sin(2R)$ 

Summary

The solution(s) found are the following

$$y = e^{-t}(e^t \sin(2t) - 2 e^t \cos(2t) + c_1) \quad (1)$$

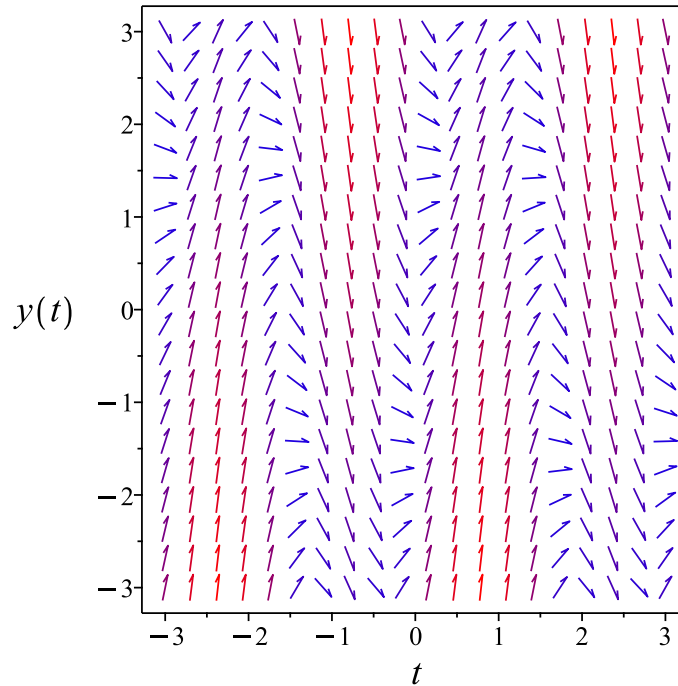


Figure 33: Slope field plot

Verification of solutions

$$y = e^{-t}(e^t \sin(2t) - 2e^t \cos(2t) + c_1)$$

Verified OK.

1.11.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dy &= (-y + 5 \sin(2t)) dt \\ (y - 5 \sin(2t)) dt + dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= y - 5 \sin(2t) \\ N(t, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y - 5 \sin(2t)) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1((1) - (0)) \\ &= 1 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int 1 dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^t \\ &= e^t \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^t(y - 5 \sin(2t)) \\ &= (y - 5 \sin(2t)) e^t \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^t(1) \\ &= e^t \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ ((y - 5 \sin(2t)) e^t) + (e^t) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int (y - 5 \sin(2t)) e^t dt \\ \phi &= e^t(y + 2 \cos(2t) - \sin(2t)) + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^t + f'(y)\quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^t$. Therefore equation (4) becomes

$$e^t = e^t + f'(y)\quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = e^t(y + 2 \cos(2t) - \sin(2t)) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = e^t(y + 2 \cos(2t) - \sin(2t))$$

The solution becomes

$$y = e^{-t}(e^t \sin(2t) - 2 e^t \cos(2t) + c_1)$$

Summary

The solution(s) found are the following

$$y = e^{-t}(e^t \sin(2t) - 2e^t \cos(2t) + c_1) \quad (1)$$

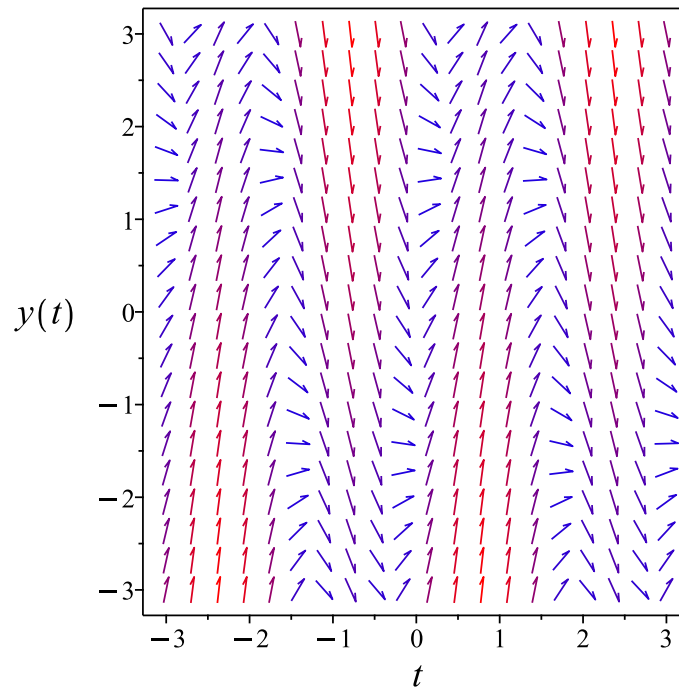


Figure 34: Slope field plot

Verification of solutions

$$y = e^{-t}(e^t \sin(2t) - 2e^t \cos(2t) + c_1)$$

Verified OK.

1.11.4 Maple step by step solution

Let's solve

$$y + y' = 5 \sin(2t)$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -y + 5 \sin(2t)$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y + y' = 5 \sin(2t)$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t)(y + y') = 5\mu(t) \sin(2t)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t)(y + y') = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \mu(t)$$

- Solve to find the integrating factor

$$\mu(t) = e^t$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int 5\mu(t) \sin(2t) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int 5\mu(t) \sin(2t) dt + c_1$$

- Solve for y

$$y = \frac{\int 5\mu(t) \sin(2t) dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^t$

$$y = \frac{\int 5e^t \sin(2t) dt + c_1}{e^t}$$

- Evaluate the integrals on the rhs

$$y = \frac{e^t \sin(2t) - 2e^t \cos(2t) + c_1}{e^t}$$

- Simplify

$$y = \sin(2t) - 2 \cos(2t) + c_1 e^{-t}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```


✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 21

```
dsolve(y(t)+diff(y(t),t) = 5*sin(2*t),y(t), singsol=all)
```

$$y(t) = \sin(2t) - 2 \cos(2t) + e^{-t}c_1$$

✓ Solution by Mathematica

Time used: 0.087 (sec). Leaf size: 24

```
DSolve[y[t]+y'[t] == 5*Sin[2*t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \sin(2t) - 2 \cos(2t) + c_1 e^{-t}$$

1.12 problem 12

1.12.1 Solving as linear ode	145
1.12.2 Solving as first order ode lie symmetry lookup ode	147
1.12.3 Solving as exact ode	151
1.12.4 Maple step by step solution	156

Internal problem ID [459]

Internal file name [OUTPUT/459_Sunday_June_05_2022_01_41_49_AM_20042804/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 12.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$y + 2y' = 3t^2$$

1.12.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{1}{2}$$
$$q(t) = \frac{3t^2}{2}$$

Hence the ode is

$$y' + \frac{y}{2} = \frac{3t^2}{2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{1}{2} dt} \\ &= e^{\frac{t}{2}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{3t^2}{2} \right) \\ \frac{d}{dt} \left(e^{\frac{t}{2}} y \right) &= \left(e^{\frac{t}{2}} \right) \left(\frac{3t^2}{2} \right) \\ d \left(e^{\frac{t}{2}} y \right) &= \left(\frac{3t^2 e^{\frac{t}{2}}}{2} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{\frac{t}{2}} y &= \int \frac{3t^2 e^{\frac{t}{2}}}{2} dt \\ e^{\frac{t}{2}} y &= 3(t^2 - 4t + 8) e^{\frac{t}{2}} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{\frac{t}{2}}$ results in

$$y = 3e^{-\frac{t}{2}}(t^2 - 4t + 8) e^{\frac{t}{2}} + c_1 e^{-\frac{t}{2}}$$

which simplifies to

$$y = 3t^2 - 12t + 24 + c_1 e^{-\frac{t}{2}}$$

Summary

The solution(s) found are the following

$$y = 3t^2 - 12t + 24 + c_1 e^{-\frac{t}{2}} \tag{1}$$

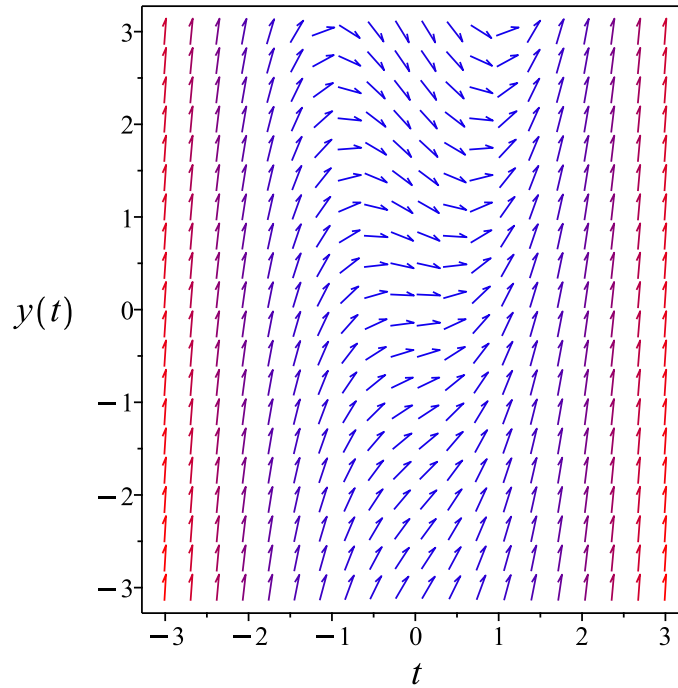


Figure 35: Slope field plot

Verification of solutions

$$y = 3t^2 - 12t + 24 + c_1 e^{-\frac{t}{2}}$$

Verified OK.

1.12.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{y}{2} + \frac{3t^2}{2}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 34: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-\frac{t}{2}}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-\frac{t}{2}}} dy \end{aligned}$$

Which results in

$$S = e^{\frac{t}{2}} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{y}{2} + \frac{3t^2}{2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= \frac{e^{\frac{t}{2}} y}{2} \\ S_y &= e^{\frac{t}{2}} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{3t^2 e^{\frac{t}{2}}}{2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{3R^2 e^{\frac{R}{2}}}{2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = 3(R^2 - 4R + 8) e^{\frac{R}{2}} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{\frac{t}{2}} y = 3(t^2 - 4t + 8) e^{\frac{t}{2}} + c_1$$

Which simplifies to

$$(-3t^2 + 12t + y - 24) e^{\frac{t}{2}} - c_1 = 0$$

Which gives

$$y = \left(3t^2 e^{\frac{t}{2}} - 12t e^{\frac{t}{2}} + 24 e^{\frac{t}{2}} + c_1 \right) e^{-\frac{t}{2}}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{y}{2} + \frac{3t^2}{2}$	$R = t$ $S = e^{\frac{t}{2}} y$	$\frac{dS}{dR} = \frac{3R^2 e^{\frac{R}{2}}}{2}$

Summary

The solution(s) found are the following

$$y = \left(3t^2 e^{\frac{t}{2}} - 12t e^{\frac{t}{2}} + 24 e^{\frac{t}{2}} + c_1 \right) e^{-\frac{t}{2}} \quad (1)$$

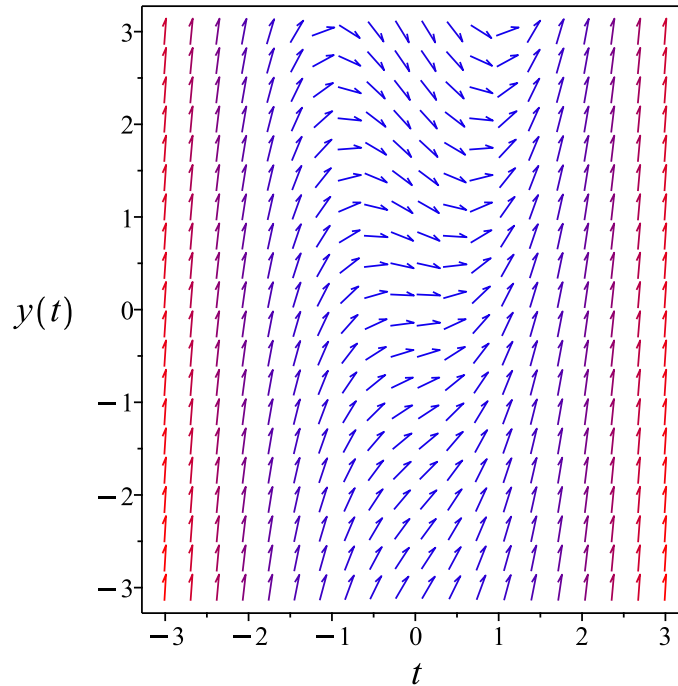


Figure 36: Slope field plot

Verification of solutions

$$y = \left(3t^2 e^{\frac{t}{2}} - 12t e^{\frac{t}{2}} + 24 e^{\frac{t}{2}} + c_1 \right) e^{-\frac{t}{2}}$$

Verified OK.

1.12.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(2) dy &= (3t^2 - y) dt \\ (-3t^2 + y) dt + (2) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -3t^2 + y \\ N(t, y) &= 2\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-3t^2 + y) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(2) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{2} ((1) - (0)) \\ &= \frac{1}{2} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int \frac{1}{2} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\frac{t}{2}} \\ &= e^{\frac{t}{2}} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \overline{M} &= \mu M \\ &= e^{\frac{t}{2}} (-3t^2 + y) \\ &= (-3t^2 + y) e^{\frac{t}{2}} \end{aligned}$$

And

$$\begin{aligned} \overline{N} &= \mu N \\ &= e^{\frac{t}{2}} (2) \\ &= 2 e^{\frac{t}{2}} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \overline{M} + \overline{N} \frac{dy}{dt} &= 0 \\ \left((-3t^2 + y) e^{\frac{t}{2}} \right) + \left(2 e^{\frac{t}{2}} \right) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int \overline{M} dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int (-3t^2 + y) e^{\frac{t}{2}} dt$$

$$\phi = -6 \left(t^2 - 4t - \frac{1}{3}y + 8 \right) e^{\frac{t}{2}} + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 2e^{\frac{t}{2}} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 2e^{\frac{t}{2}}$. Therefore equation (4) becomes

$$2e^{\frac{t}{2}} = 2e^{\frac{t}{2}} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -6 \left(t^2 - 4t - \frac{1}{3}y + 8 \right) e^{\frac{t}{2}} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -6 \left(t^2 - 4t - \frac{1}{3}y + 8 \right) e^{\frac{t}{2}}$$

The solution becomes

$$y = \frac{\left(6t^2 e^{\frac{t}{2}} - 24t e^{\frac{t}{2}} + 48 e^{\frac{t}{2}} + c_1\right) e^{-\frac{t}{2}}}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(6t^2 e^{\frac{t}{2}} - 24t e^{\frac{t}{2}} + 48 e^{\frac{t}{2}} + c_1\right) e^{-\frac{t}{2}}}{2} \quad (1)$$

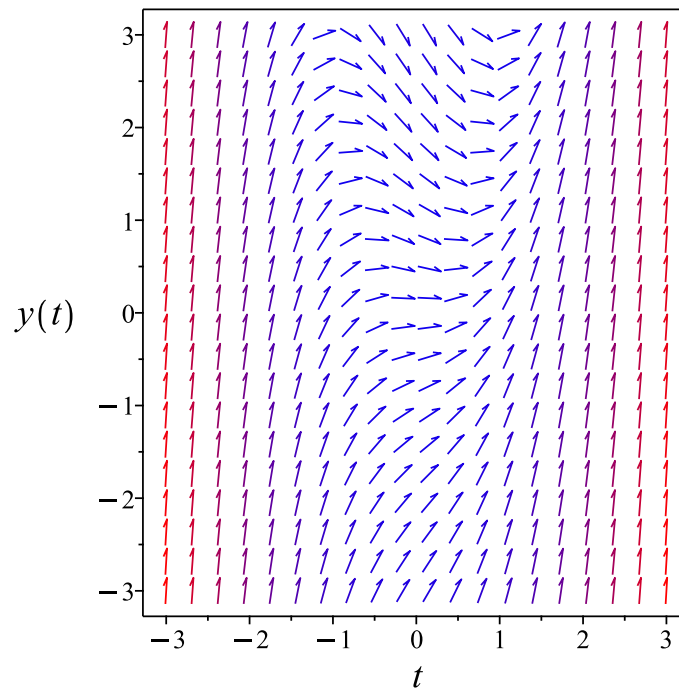


Figure 37: Slope field plot

Verification of solutions

$$y = \frac{\left(6t^2 e^{\frac{t}{2}} - 24t e^{\frac{t}{2}} + 48 e^{\frac{t}{2}} + c_1\right) e^{-\frac{t}{2}}}{2}$$

Verified OK.

1.12.4 Maple step by step solution

Let's solve

$$y + 2y' = 3t^2$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{y}{2} + \frac{3t^2}{2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{y}{2} = \frac{3t^2}{2}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{y}{2} \right) = \frac{3\mu(t)t^2}{2}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{y}{2} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)}{2}$$

- Solve to find the integrating factor

$$\mu(t) = e^{\frac{t}{2}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{3\mu(t)t^2}{2} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{3\mu(t)t^2}{2} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{3\mu(t)t^2}{2} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{\frac{t}{2}}$

$$y = \frac{\int \frac{3t^2 e^{\frac{t}{2}}}{2} dt + c_1}{e^{\frac{t}{2}}}$$

- Evaluate the integrals on the rhs

$$y = \frac{3(t^2 - 4t + 8)e^{\frac{t}{2}} + c_1}{e^{\frac{t}{2}}}$$

- Simplify

$$y = 3t^2 - 12t + 24 + c_1 e^{-\frac{t}{2}}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 20

```
dsolve(y(t)+2*diff(y(t),t) = 3*t^2,y(t), singsol=all)
```

$$y(t) = 3t^2 - 12t + 24 + e^{-\frac{t}{2}}c_1$$

✓ Solution by Mathematica

Time used: 0.047 (sec). Leaf size: 25

```
DSolve[y[t]+2*y'[t] == 3*t^2,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow 3t^2 - 12t + c_1 e^{-t/2} + 24$$

1.13 problem 13

1.13.1 Existence and uniqueness analysis	158
1.13.2 Solving as linear ode	159
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Internal problem ID [460]

Internal file name [OUTPUT/460_Sunday_June_05_2022_01_41_49_AM_58508263/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 13.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$-y + y' = 2e^{2t}t$$

With initial conditions

$$[y(0) = 1]$$

1.13.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -1$$
$$q(t) = 2e^{2t}t$$

Hence the ode is

$$-y + y' = 2e^{2t}$$

The domain of $p(t) = -1$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = 2e^{2t}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

1.13.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int(-1)dt} \\ &= e^{-t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (2e^{2t}) \\ \frac{d}{dt}(e^{-t}y) &= (e^{-t}) (2e^{2t}) \\ d(e^{-t}y) &= (2te^t) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-t}y &= \int 2te^t dt \\ e^{-t}y &= 2e^t(-1+t) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-t}$ results in

$$y = 2e^{2t}(-1+t) + c_1e^t$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = c_1 - 2$$

$$c_1 = 3$$

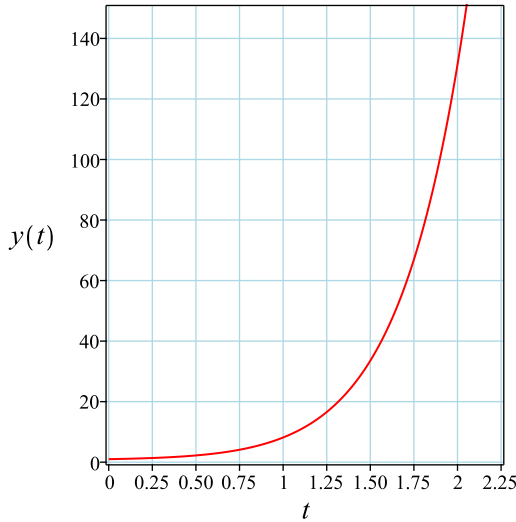
Substituting c_1 found above in the general solution gives

$$y = 2e^{2t}t - 2e^{2t} + 3e^t$$

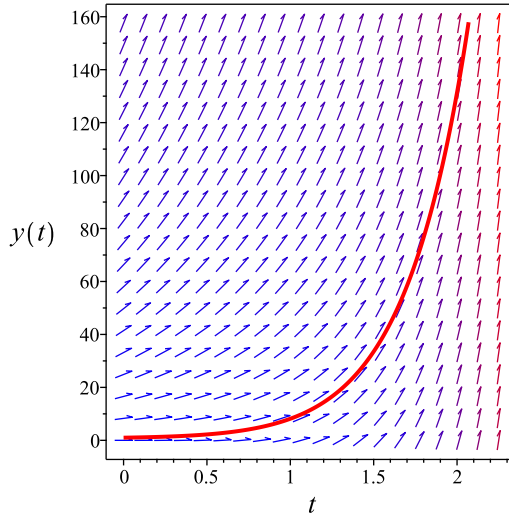
Summary

The solution(s) found are the following

$$y = 2e^{2t}t - 2e^{2t} + 3e^t \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = 2e^{2t}t - 2e^{2t} + 3e^t$$

Verified OK.

1.13.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = y + 2e^{2t}t$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (A)$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 37: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^t\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^t} dy \end{aligned}$$

Which results in

$$S = e^{-t}y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = y + 2e^{2t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -e^{-t}y \\ S_y &= e^{-t} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 2t e^t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 2R e^R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = 2(R - 1)e^R + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{-t}y = 2e^t(-1 + t) + c_1$$

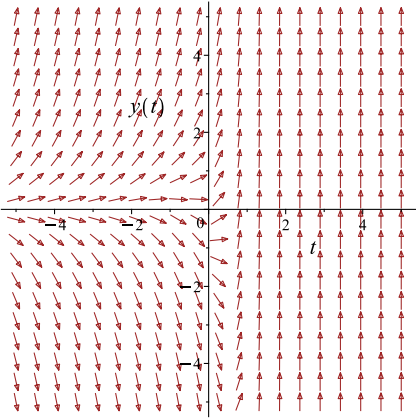
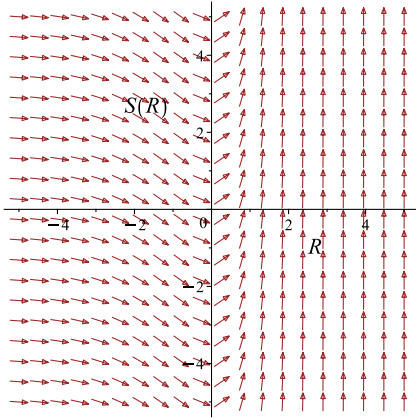
Which simplifies to

$$e^{-t}y = 2e^t(-1 + t) + c_1$$

Which gives

$$y = e^t(2te^t - 2e^t + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = y + 2e^{2t}t$ 	$R = t$ $S = e^{-t}y$	$\frac{dS}{dR} = 2Re^R$ 

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = c_1 - 2$$

$$c_1 = 3$$

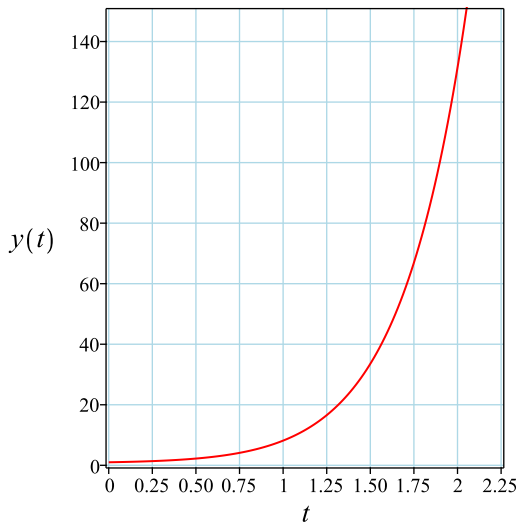
Substituting c_1 found above in the general solution gives

$$y = 2e^{2t}t - 2e^{2t} + 3e^t$$

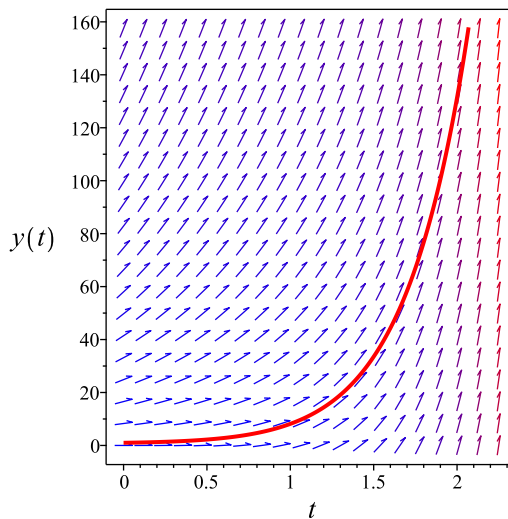
Summary

The solution(s) found are the following

$$y = 2e^{2t}t - 2e^{2t} + 3e^t \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = 2e^{2t}t - 2e^{2t} + 3e^t$$

Verified OK.

1.13.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} dy &= (y + 2e^{2t}) dt \\ (-y - 2e^{2t}) dt + dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= -y - 2e^{2t} \\ N(t, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-y - 2e^{2t}) \\ &= -1 \end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1((-1) - (0)) \\ &= -1\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int -1 dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-t} \\ &= e^{-t}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\overline{M} &= \mu M \\ &= e^{-t}(-y - 2e^{2t}t) \\ &= -e^{-t}y - 2te^t\end{aligned}$$

And

$$\begin{aligned}\overline{N} &= \mu N \\ &= e^{-t}(1) \\ &= e^{-t}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\overline{M} + \overline{N} \frac{dy}{dt} &= 0 \\ (-e^{-t}y - 2te^t) + (e^{-t}) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int \bar{M} dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int -e^{-t}y - 2te^t dt$$

$$\phi = e^{-t}y - 2e^t(-1 + t) + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-t} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-t}$. Therefore equation (4) becomes

$$e^{-t} = e^{-t} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = e^{-t}y - 2e^t(-1 + t) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = e^{-t}y - 2e^t(-1 + t)$$

The solution becomes

$$y = e^t(2te^t - 2e^t + c_1)$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = c_1 - 2$$

$$c_1 = 3$$

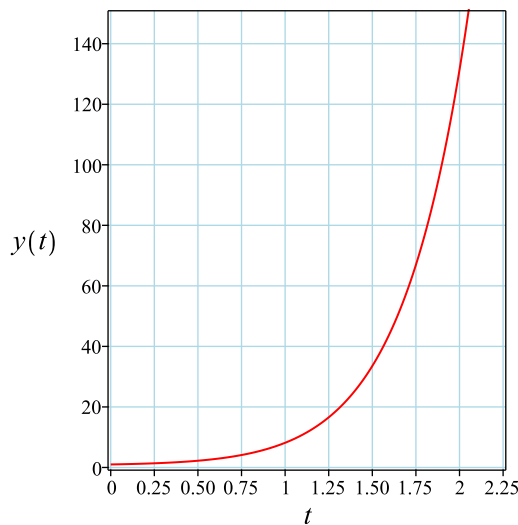
Substituting c_1 found above in the general solution gives

$$y = 2e^{2t}t - 2e^{2t} + 3e^t$$

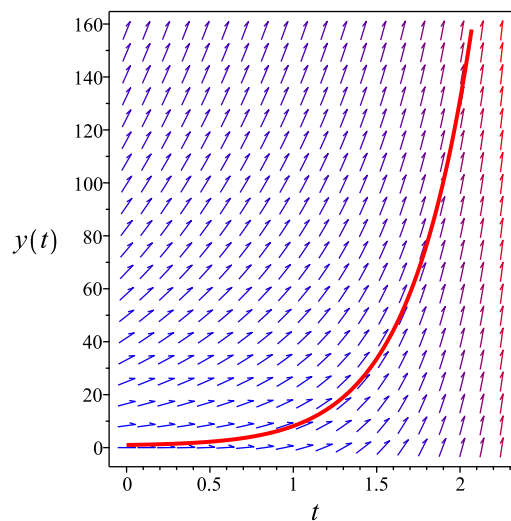
Summary

The solution(s) found are the following

$$y = 2e^{2t}t - 2e^{2t} + 3e^t \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = 2e^{2t}t - 2e^{2t} + 3e^t$$

Verified OK.

1.13.5 Maple step by step solution

Let's solve

$$[-y + y' = 2e^{2t}t, y(0) = 1]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = y + 2e^{2t}t$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$-y + y' = 2e^{2t}t$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t)(-y + y') = 2\mu(t)e^{2t}t$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t)(-y + y') = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -\mu(t)$$

- Solve to find the integrating factor

$$\mu(t) = e^{-t}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int 2\mu(t)e^{2t}t dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int 2\mu(t)e^{2t}t dt + c_1$$

- Solve for y

$$y = \frac{\int 2\mu(t)e^{2t}t dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{-t}$

$$y = \frac{\int 2e^{2t}te^{-t} dt + c_1}{e^{-t}}$$

- Evaluate the integrals on the rhs

$$y = \frac{2e^t(-1+t) + c_1}{e^{-t}}$$

- Simplify

$$y = (2t - 2)(e^t)^2 + c_1e^t$$

- Use initial condition $y(0) = 1$
 $1 = c_1 - 2$
- Solve for c_1
 $c_1 = 3$
- Substitute $c_1 = 3$ into general solution and simplify
 $y = (2t - 2)e^{2t} + 3e^t$
- Solution to the IVP
 $y = (2t - 2)e^{2t} + 3e^t$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 16

```
dsolve([-y(t)+diff(y(t),t) = 2*exp(2*t)*t,y(0) = 1],y(t), singsol=all)
```

$$y(t) = (2t - 2)e^{2t} + 3e^t$$

✓ Solution by Mathematica

Time used: 0.053 (sec). Leaf size: 19

```
DSolve[{-y[t]+y'[t] == 2*Exp[2*t]*t,y[0]==1},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^t(2e^t(t - 1) + 3)$$

1.14 problem 14

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Internal problem ID [461]

Internal file name [OUTPUT/461_Sunday_June_05_2022_01_41_50_AM_10145323/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 14.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "linear", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$2y + y' = te^{-2t}$$

With initial conditions

$$[y(1) = 0]$$

1.14.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$\begin{aligned} p(t) &= 2 \\ q(t) &= te^{-2t} \end{aligned}$$

Hence the ode is

$$2y + y' = t e^{-2t}$$

The domain of $p(t) = 2$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 1$ is inside this domain. The domain of $q(t) = t e^{-2t}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 1$ is also inside this domain. Hence solution exists and is unique.

1.14.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int 2dt} \\ &= e^{2t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (t e^{-2t}) \\ \frac{d}{dt}(y e^{2t}) &= (e^{2t}) (t e^{-2t}) \\ d(y e^{2t}) &= t dt\end{aligned}$$

Integrating gives

$$\begin{aligned}y e^{2t} &= \int t dt \\ y e^{2t} &= \frac{t^2}{2} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{2t}$ results in

$$y = \frac{e^{-2t}t^2}{2} + c_1 e^{-2t}$$

which simplifies to

$$y = e^{-2t} \left(\frac{t^2}{2} + c_1 \right)$$

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{e^{-2}(2c_1 + 1)}{2}$$

$$c_1 = -\frac{1}{2}$$

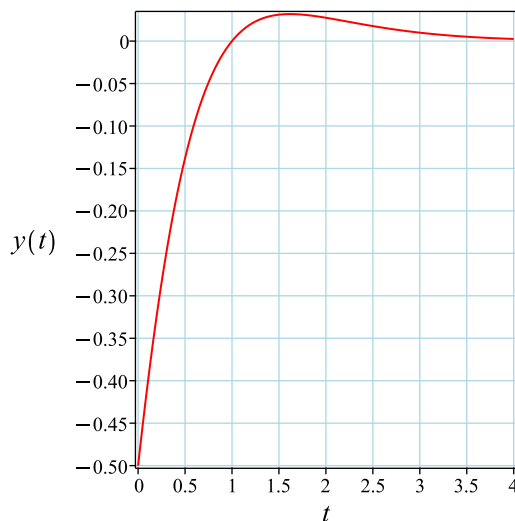
Substituting c_1 found above in the general solution gives

$$y = \frac{e^{-2t}(t^2 - 1)}{2}$$

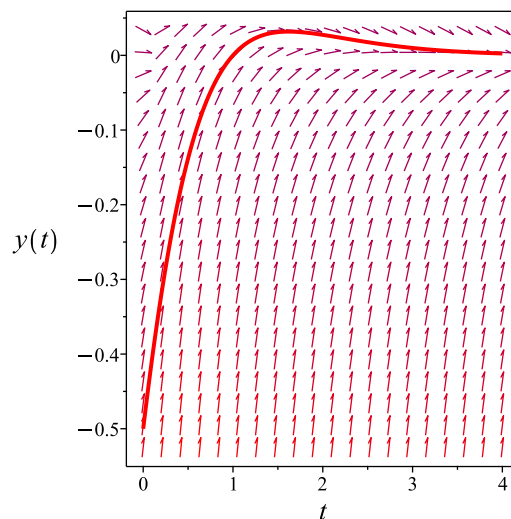
Summary

The solution(s) found are the following

$$y = \frac{e^{-2t}(t^2 - 1)}{2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^{-2t}(t^2 - 1)}{2}$$

Verified OK.

1.14.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -(2y e^{2t} - t) e^{-2t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 40: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-2t}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-2t}} dy\end{aligned}$$

Which results in

$$S = y e^{2t}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -(2y e^{2t} - t) e^{-2t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 1 \\ R_y &= 0 \\ S_t &= 2y e^{2t} \\ S_y &= e^{2t}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = t \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^2}{2} + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$y e^{2t} = \frac{t^2}{2} + c_1$$

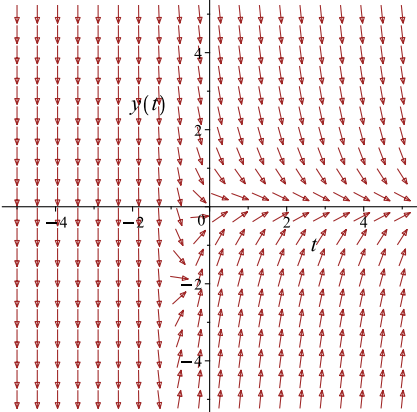
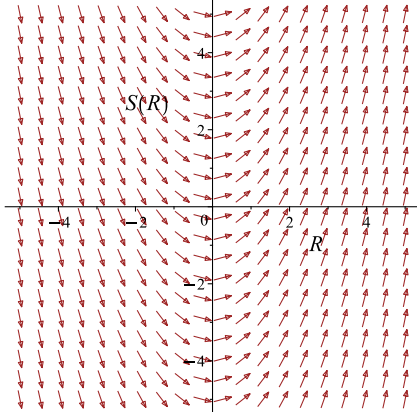
Which simplifies to

$$y e^{2t} = \frac{t^2}{2} + c_1$$

Which gives

$$y = \frac{e^{-2t}(t^2 + 2c_1)}{2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -(2ye^{2t} - t)e^{-2t}$ 	$R = t$ $S = ye^{2t}$	$\frac{dS}{dR} = R$ 

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = e^{-2}c_1 + \frac{e^{-2}}{2}$$

$$c_1 = -\frac{1}{2}$$

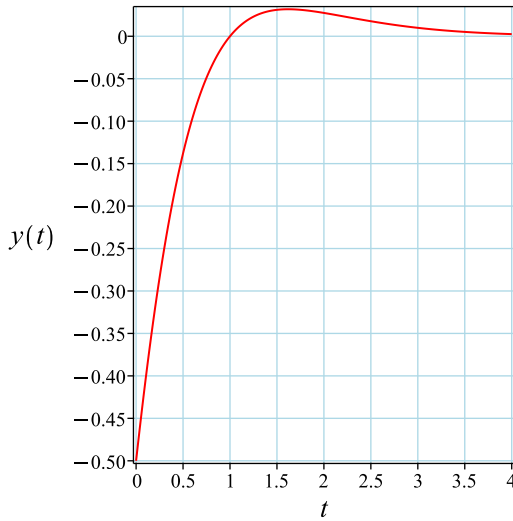
Substituting c_1 found above in the general solution gives

$$y = \frac{e^{-2t}t^2}{2} - \frac{e^{-2t}}{2}$$

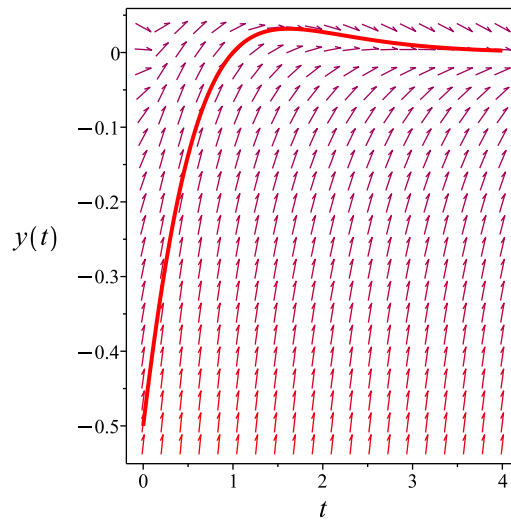
Summary

The solution(s) found are the following

$$y = \frac{e^{-2t}t^2}{2} - \frac{e^{-2t}}{2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^{-2t}t^2}{2} - \frac{e^{-2t}}{2}$$

Verified OK.

1.14.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (e^{2t}) dy &= (-2y e^{2t} + t) dt \\ (2y e^{2t} - t) dt + (e^{2t}) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= 2y e^{2t} - t \\ N(t, y) &= e^{2t} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (2y e^{2t} - t) \\ &= 2 e^{2t} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (e^{2t}) \\ &= 2 e^{2t} \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int 2y e^{2t} - t dt \\ \phi &= -\frac{t^2}{2} + y e^{2t} + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{2t} + f'(y)\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{2t}$. Therefore equation (4) becomes

$$e^{2t} = e^{2t} + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{t^2}{2} + y e^{2t} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{t^2}{2} + y e^{2t}$$

The solution becomes

$$y = \frac{e^{-2t}(t^2 + 2c_1)}{2}$$

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = e^{-2}c_1 + \frac{e^{-2}}{2}$$

$$c_1 = -\frac{1}{2}$$

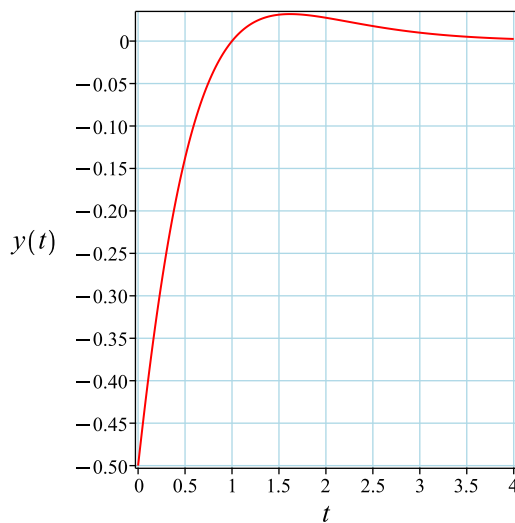
Substituting c_1 found above in the general solution gives

$$y = \frac{e^{-2t}t^2}{2} - \frac{e^{-2t}}{2}$$

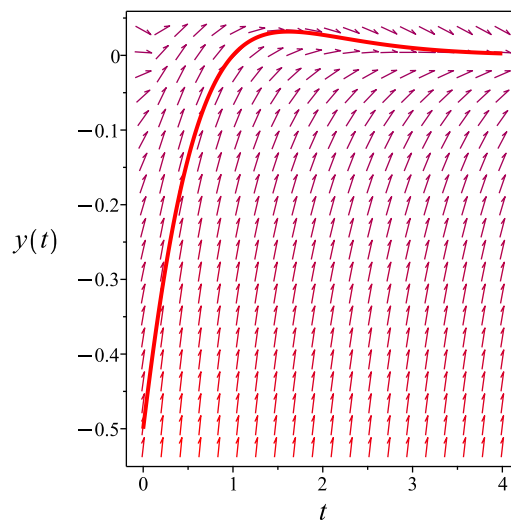
Summary

The solution(s) found are the following

$$y = \frac{e^{-2t}t^2}{2} - \frac{e^{-2t}}{2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^{-2t}t^2}{2} - \frac{e^{-2t}}{2}$$

Verified OK.

1.14.5 Maple step by step solution

Let's solve

$$[2y + y' = \frac{t}{e^{2t}}, y(1) = 0]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -2y + \frac{t}{e^{2t}}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$2y + y' = \frac{t}{e^{2t}}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) (2y + y') = \frac{\mu(t)t}{e^{2t}}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t) y)$

$$\mu(t) (2y + y') = \mu'(t) y + \mu(t) y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = 2\mu(t)$$

- Solve to find the integrating factor

$$\mu(t) = (e^{2t})^2 e^{-2t}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y) \right) dt = \int \frac{\mu(t)t}{e^{2t}} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t) y = \int \frac{\mu(t)t}{e^{2t}} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)t}{e^{2t}} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = (e^{2t})^2 e^{-2t}$

$$y = \frac{\int t e^{-2t} e^{2t} dt + c_1}{(e^{2t})^2 e^{-2t}}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{t^2}{2} + c_1}{(e^{2t})^2 e^{-2t}}$$

- Simplify

$$y = \frac{e^{-2t}(t^2+2c_1)}{2}$$

- Use initial condition $y(1) = 0$

$$0 = \frac{e^{-2}(2c_1+1)}{2}$$

- Solve for c_1

$$c_1 = -\frac{1}{2}$$

- Substitute $c_1 = -\frac{1}{2}$ into general solution and simplify

$$y = \frac{e^{-2t}(t^2-1)}{2}$$

- Solution to the IVP

$$y = \frac{e^{-2t}(t^2-1)}{2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve([2*y(t)+diff(y(t),t) = t/exp(2*t),y(1) = 0],y(t), singsol=all)
```

$$y(t) = \frac{(t^2 - 1)e^{-2t}}{2}$$

✓ Solution by Mathematica

Time used: 0.054 (sec). Leaf size: 19

```
DSolve[{2*y[t]+y'[t] == t/Exp[2*t],y[1]==0},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{2}e^{-2t}(t^2 - 1)$$

1.15 problem 15

1.15.1 Existence and uniqueness analysis	184
1.15.2 Solving as linear ode	185
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Internal problem ID [462]

Internal file name [OUTPUT/462_Sunday_June_05_2022_01_41_51_AM_24139145/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 15.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**linear**", "**exactWithIntegrationFactor**", "**first_order_ode_lie_symmetry_lookup**"

Maple gives the following as the ode type

[_linear]

$$2y + ty' = t^2 - t + 1$$

With initial conditions

$$\left[y(1) = \frac{1}{2} \right]$$

1.15.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{2}{t}$$
$$q(t) = \frac{t^2 - t + 1}{t}$$

Hence the ode is

$$y' + \frac{2y}{t} = \frac{t^2 - t + 1}{t}$$

The domain of $p(t) = \frac{2}{t}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = 1$ is inside this domain. The domain of $q(t) = \frac{t^2 - t + 1}{t}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = 1$ is also inside this domain. Hence solution exists and is unique.

1.15.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{2}{t} dt} \\ &= t^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{t^2 - t + 1}{t} \right) \\ \frac{d}{dt}(t^2 y) &= (t^2) \left(\frac{t^2 - t + 1}{t} \right) \\ d(t^2 y) &= ((t^2 - t + 1) t) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t^2 y &= \int (t^2 - t + 1) t dt \\ t^2 y &= \frac{1}{4}t^4 - \frac{1}{3}t^3 + \frac{1}{2}t^2 + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t^2$ results in

$$y = \frac{\frac{1}{4}t^4 - \frac{1}{3}t^3 + \frac{1}{2}t^2}{t^2} + \frac{c_1}{t^2}$$

which simplifies to

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 12c_1}{12t^2}$$

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = \frac{1}{2}$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{1}{2} = c_1 + \frac{5}{12}$$

$$c_1 = \frac{1}{12}$$

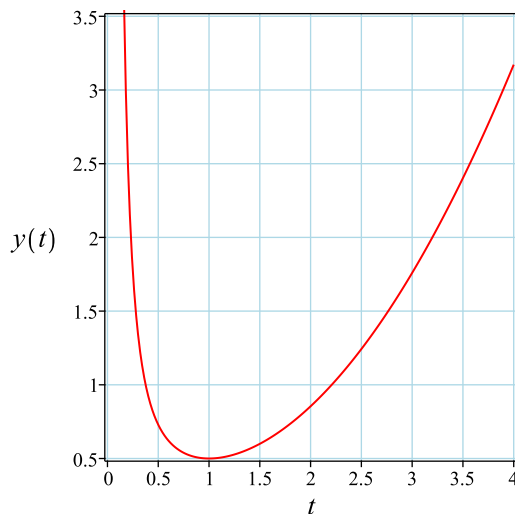
Substituting c_1 found above in the general solution gives

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2}$$

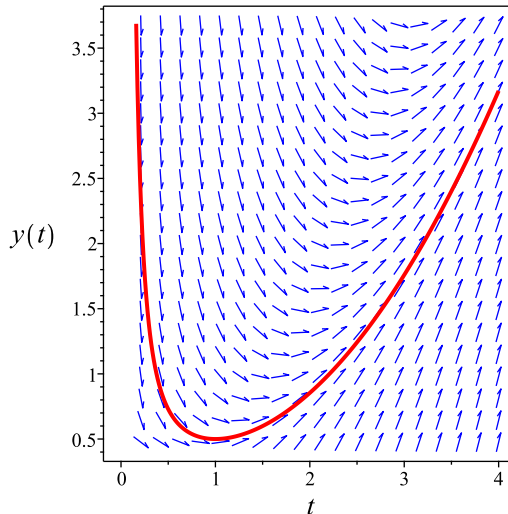
Summary

The solution(s) found are the following

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2}$$

Verified OK.

1.15.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{-t^2 + t + 2y - 1}{t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 43: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \frac{1}{t^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{t^2}} dy\end{aligned}$$

Which results in

$$S = t^2 y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{-t^2 + t + 2y - 1}{t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 1 \\R_y &= 0 \\S_t &= 2ty \\S_y &= t^2\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = t^3 - t^2 + t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R^3 - R^2 + R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{1}{4}R^4 - \frac{1}{3}R^3 + \frac{1}{2}R^2 + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$yt^2 = \frac{1}{4}t^4 - \frac{1}{3}t^3 + \frac{1}{2}t^2 + c_1$$

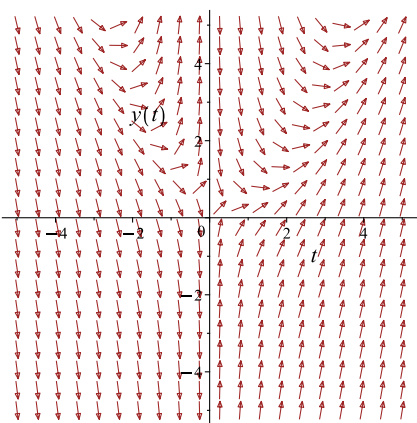
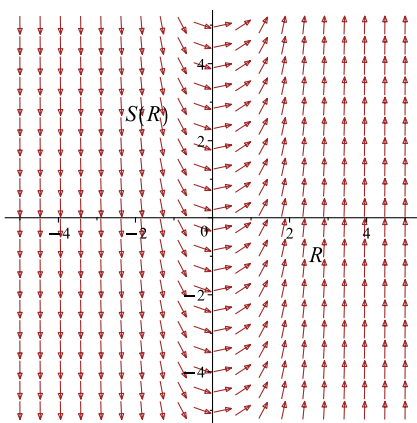
Which simplifies to

$$yt^2 = \frac{1}{4}t^4 - \frac{1}{3}t^3 + \frac{1}{2}t^2 + c_1$$

Which gives

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 12c_1}{12t^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{-t^2+t+2y-1}{t}$ 	$R = t$ $S = t^2y$	$\frac{dS}{dR} = R^3 - R^2 + R$ 

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = \frac{1}{2}$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{1}{2} = c_1 + \frac{5}{12}$$

$$c_1 = \frac{1}{12}$$

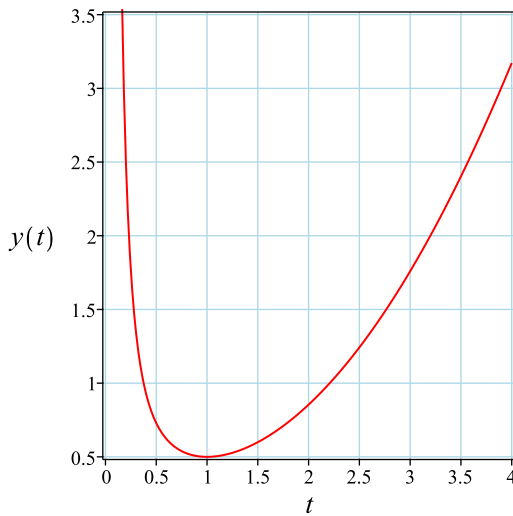
Substituting c_1 found above in the general solution gives

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2}$$

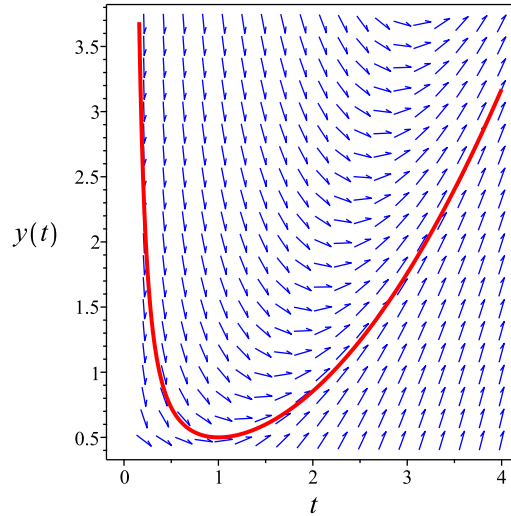
Summary

The solution(s) found are the following

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2}$$

Verified OK.

1.15.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (t) dy &= (t^2 - t - 2y + 1) dt \\ (-t^2 + t + 2y - 1) dt + (t) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= -t^2 + t + 2y - 1 \\ N(t, y) &= t \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (-t^2 + t + 2y - 1) \\ &= 2 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (t) \\ &= 1 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{t} ((2) - (1)) \\ &= \frac{1}{t} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int \frac{1}{t} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{\ln(t)} \\ &= t\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= t(-t^2 + t + 2y - 1) \\ &= -t(t^2 - t - 2y + 1)\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= t(t) \\ &= t^2\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ (-t(t^2 - t - 2y + 1)) + (t^2) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -t(t^2 - t - 2y + 1) dt \\ \phi &= -\frac{t^2(t^2 - \frac{4}{3}t - 4y + 2)}{4} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t^2 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t^2$. Therefore equation (4) becomes

$$t^2 = t^2 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{t^2(t^2 - \frac{4}{3}t - 4y + 2)}{4} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{t^2(t^2 - \frac{4}{3}t - 4y + 2)}{4}$$

The solution becomes

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 12c_1}{12t^2}$$

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = \frac{1}{2}$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{1}{2} = c_1 + \frac{5}{12}$$

$$c_1 = \frac{1}{12}$$

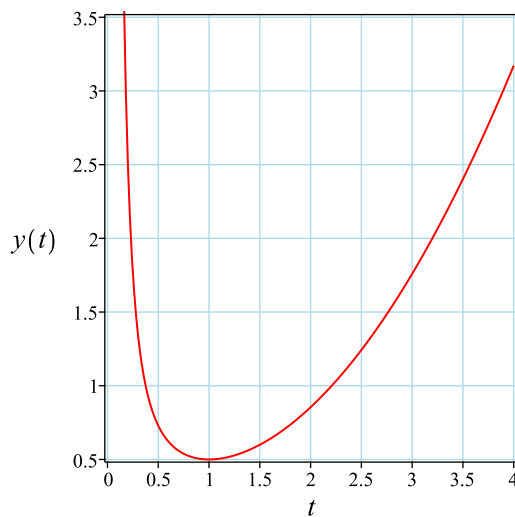
Substituting c_1 found above in the general solution gives

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2}$$

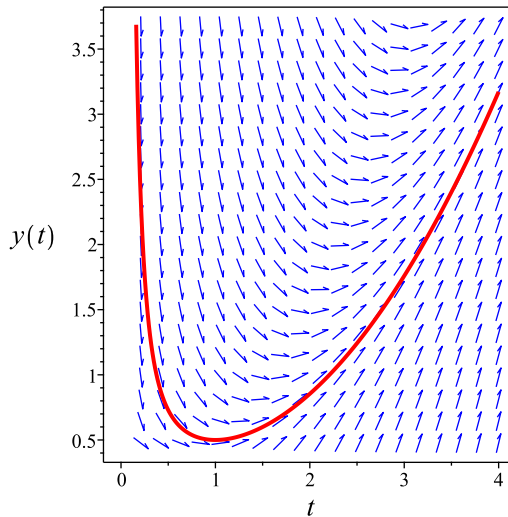
Summary

The solution(s) found are the following

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2}$$

Verified OK.

1.15.5 Maple step by step solution

Let's solve

$$[2y + ty' = t^2 - t + 1, y(1) = \frac{1}{2}]$$

- Highest derivative means the order of the ODE is 1

y'

- Isolate the derivative

$$y' = -\frac{2y}{t} + \frac{t^2-t+1}{t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{2y}{t} = \frac{t^2-t+1}{t}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \frac{\mu(t)(t^2-t+1)}{t}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{2\mu(t)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = t^2$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)(t^2-t+1)}{t} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)(t^2-t+1)}{t} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)(t^2-t+1)}{t} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = t^2$

$$y = \frac{\int (t^2-t+1)tdt + c_1}{t^2}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{1}{4}t^4 - \frac{1}{3}t^3 + \frac{1}{2}t^2 + c_1}{t^2}$$

- Simplify

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 12c_1}{12t^2}$$

- Use initial condition $y(1) = \frac{1}{2}$

$$\frac{1}{2} = c_1 + \frac{5}{12}$$

- Solve for c_1

$$c_1 = \frac{1}{12}$$

- Substitute $c_1 = \frac{1}{12}$ into general solution and simplify

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2}$$

- Solution to the IVP

$$y = \frac{3t^4 - 4t^3 + 6t^2 + 1}{12t^2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve([2*y(t)+t*diff(y(t),t) = t^2-t+1,y(1) = 1/2],y(t), singsol=all)
```

$$y(t) = \frac{t^2}{4} - \frac{t}{3} + \frac{1}{2} + \frac{1}{12t^2}$$

✓ Solution by Mathematica

Time used: 0.028 (sec). Leaf size: 22

```
DSolve[{2*y[t]+t*y'[t] == t^2-t+1,y[1]==1/2},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{12} \left(3t^2 + \frac{1}{t^2} - 4t + 6 \right)$$

1.16 problem 16

1.16.1 Existence and uniqueness analysis	198
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1.16.4 Solving as exact ode	205
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Internal problem ID [463]

Internal file name [OUTPUT/463_Sunday_June_05_2022_01_41_52_AM_63516619/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 16.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "linear", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$y' + \frac{2y}{t} = \frac{\cos(t)}{t^2}$$

With initial conditions

$$[y(\pi) = 0]$$

1.16.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{2}{t}$$
$$q(t) = \frac{\cos(t)}{t^2}$$

Hence the ode is

$$y' + \frac{2y}{t} = \frac{\cos(t)}{t^2}$$

The domain of $p(t) = \frac{2}{t}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = \pi$ is inside this domain. The domain of $q(t) = \frac{\cos(t)}{t^2}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = \pi$ is also inside this domain. Hence solution exists and is unique.

1.16.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{2}{t} dt} \\ &= t^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{\cos(t)}{t^2} \right) \\ \frac{d}{dt}(t^2 y) &= (t^2) \left(\frac{\cos(t)}{t^2} \right) \\ d(t^2 y) &= \cos(t) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t^2 y &= \int \cos(t) dt \\ t^2 y &= \sin(t) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t^2$ results in

$$y = \frac{\sin(t)}{t^2} + \frac{c_1}{t^2}$$

which simplifies to

$$y = \frac{\sin(t) + c_1}{t^2}$$

Initial conditions are used to solve for c_1 . Substituting $t = \pi$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{c_1}{\pi^2}$$

$$c_1 = 0$$

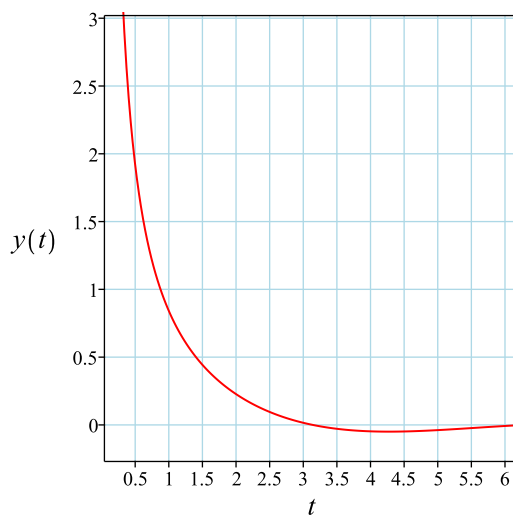
Substituting c_1 found above in the general solution gives

$$y = \frac{\sin(t)}{t^2}$$

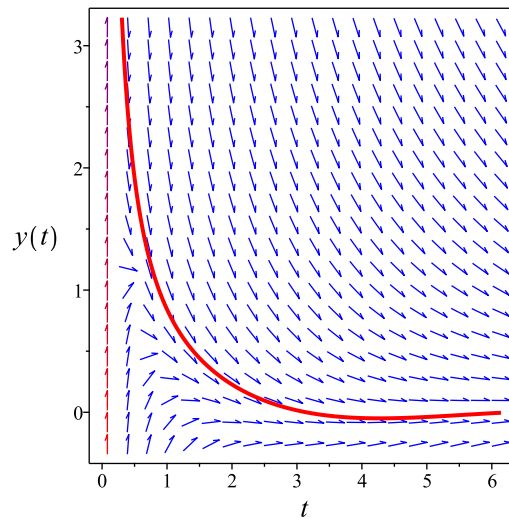
Summary

The solution(s) found are the following

$$y = \frac{\sin(t)}{t^2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{\sin(t)}{t^2}$$

Verified OK.

1.16.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-2ty + \cos(t)}{t^2}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 46: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \frac{1}{t^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{t^2}} dy\end{aligned}$$

Which results in

$$S = t^2 y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{-2ty + \cos(t)}{t^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 1 \\R_y &= 0 \\S_t &= 2ty \\S_y &= t^2\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \cos(t) \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \cos(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \sin(R) + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$yt^2 = \sin(t) + c_1$$

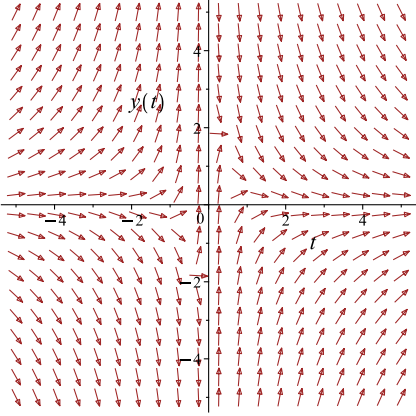
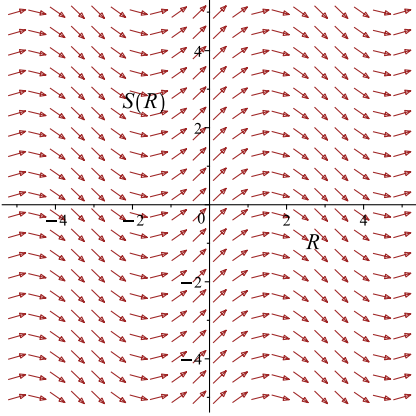
Which simplifies to

$$yt^2 = \sin(t) + c_1$$

Which gives

$$y = \frac{\sin(t) + c_1}{t^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{-2ty + \cos(t)}{t^2}$ 	$R = t$ $S = t^2 y$	$\frac{dS}{dR} = \cos(R)$ 

Initial conditions are used to solve for c_1 . Substituting $t = \pi$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{c_1}{\pi^2}$$

$$c_1 = 0$$

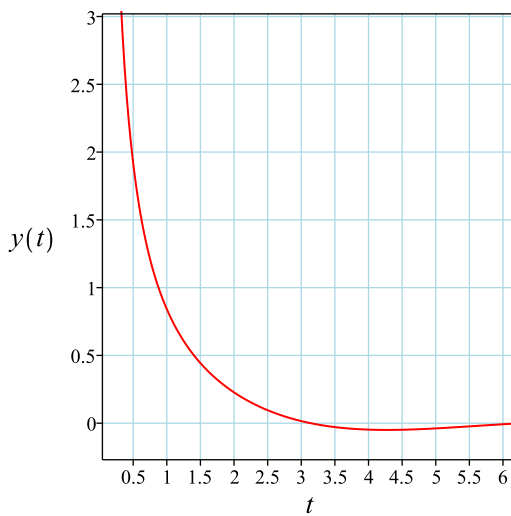
Substituting c_1 found above in the general solution gives

$$y = \frac{\sin(t)}{t^2}$$

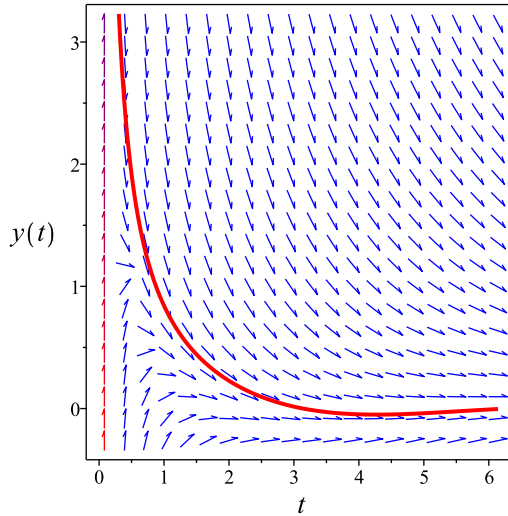
Summary

The solution(s) found are the following

$$y = \frac{\sin(t)}{t^2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{\sin(t)}{t^2}$$

Verified OK.

1.16.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (t^2) dy &= (-2ty + \cos(t)) dt \\ (2ty - \cos(t)) dt + (t^2) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= 2ty - \cos(t) \\ N(t, y) &= t^2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2ty - \cos(t)) \\ &= 2t \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(t^2) \\ &= 2t \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int 2ty - \cos(t) dt \\ \phi &= t^2y - \sin(t) + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t^2 + f'(y)\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t^2$. Therefore equation (4) becomes

$$t^2 = t^2 + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = t^2y - \sin(t) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = t^2y - \sin(t)$$

The solution becomes

$$y = \frac{\sin(t) + c_1}{t^2}$$

Initial conditions are used to solve for c_1 . Substituting $t = \pi$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{c_1}{\pi^2}$$

$$c_1 = 0$$

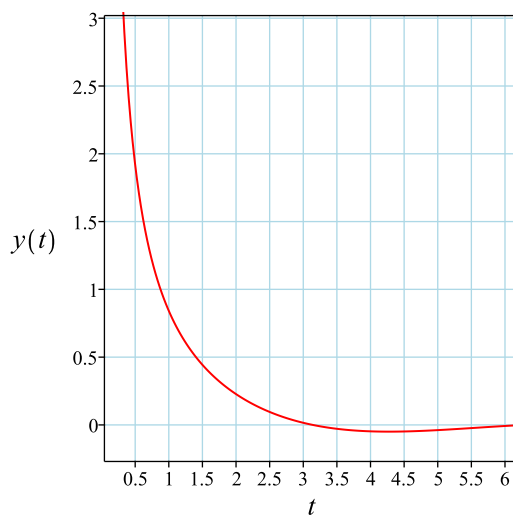
Substituting c_1 found above in the general solution gives

$$y = \frac{\sin(t)}{t^2}$$

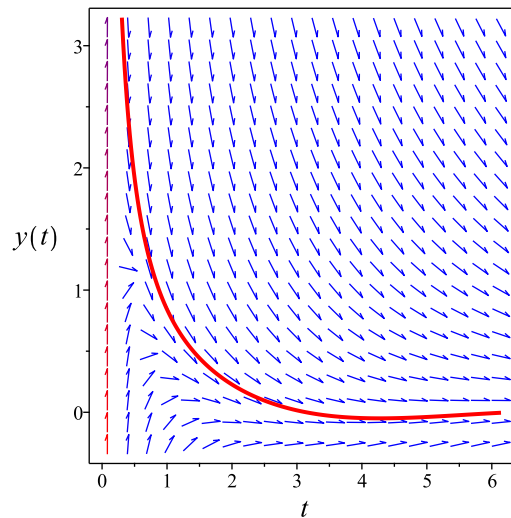
Summary

The solution(s) found are the following

$$y = \frac{\sin(t)}{t^2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{\sin(t)}{t^2}$$

Verified OK.

1.16.5 Maple step by step solution

Let's solve

$$\left[y' + \frac{2y}{t} = \frac{\cos(t)}{t^2}, y(\pi) = 0 \right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{2y}{t} + \frac{\cos(t)}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{2y}{t} = \frac{\cos(t)}{t^2}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \frac{\mu(t)\cos(t)}{t^2}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{2\mu(t)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = t^2$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)\cos(t)}{t^2} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)\cos(t)}{t^2} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)\cos(t)}{t^2} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = t^2$

$$y = \frac{\int \cos(t) dt + c_1}{t^2}$$

- Evaluate the integrals on the rhs

$$y = \frac{\sin(t) + c_1}{t^2}$$

- Use initial condition $y(\pi) = 0$

$$0 = \frac{c_1}{\pi^2}$$

- Solve for c_1

$$c_1 = 0$$

- Substitute $c_1 = 0$ into general solution and simplify

$$y = \frac{\sin(t)}{t^2}$$

- Solution to the IVP

$$y = \frac{\sin(t)}{t^2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 10

```
dsolve([2*y(t)/t+diff(y(t),t) = cos(t)/t^2,y(Pi) = 0],y(t), singsol=all)
```

$$y(t) = \frac{\sin(t)}{t^2}$$

✓ Solution by Mathematica

Time used: 0.041 (sec). Leaf size: 11

```
DSolve[{2*y[t]/t+y'[t] == Cos[t]/t^2,y[Pi]==0},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{\sin(t)}{t^2}$$

1.17 problem 17

1.17.1 Existence and uniqueness analysis	211
1.17.2 Solving as linear ode	212
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Internal problem ID [464]

Internal file name [OUTPUT/464_Sunday_June_05_2022_01_41_53_AM_28340272/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 17.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$-2y + y' = e^{2t}$$

With initial conditions

$$[y(0) = 2]$$

1.17.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -2$$

$$q(t) = e^{2t}$$

Hence the ode is

$$-2y + y' = e^{2t}$$

The domain of $p(t) = -2$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = e^{2t}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

1.17.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int (-2)dt} \\ &= e^{-2t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu)(e^{2t}) \\ \frac{d}{dt}(e^{-2t}y) &= (e^{-2t})(e^{2t}) \\ d(e^{-2t}y) &= dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-2t}y &= \int dt \\ e^{-2t}y &= t + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-2t}$ results in

$$y = e^{2t}t + c_1e^{2t}$$

which simplifies to

$$y = e^{2t}(t + c_1)$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = 2$ in the above solution gives an equation to solve for the constant of integration.

$$2 = c_1$$

$$c_1 = 2$$

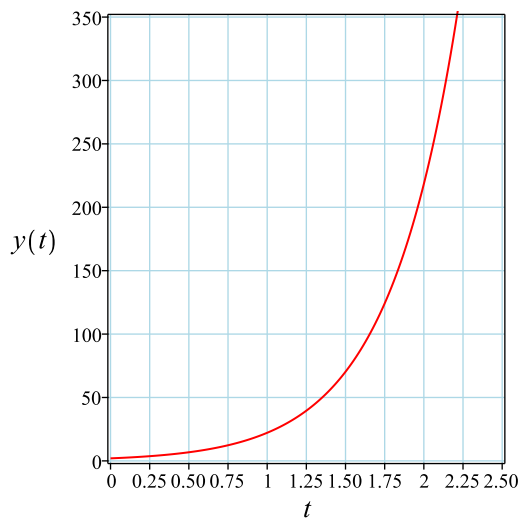
Substituting c_1 found above in the general solution gives

$$y = e^{2t}(2 + t)$$

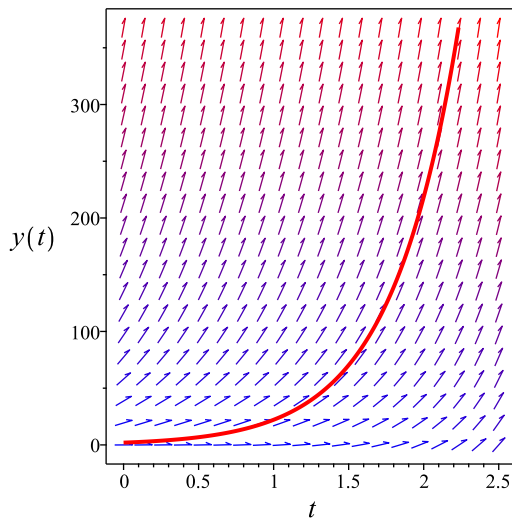
Summary

The solution(s) found are the following

$$y = e^{2t}(2 + t) \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{2t}(2 + t)$$

Verified OK.

1.17.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = 2y + e^{2t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 49: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int(n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{2t}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{2t}} dy\end{aligned}$$

Which results in

$$S = e^{-2t}y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = 2y + e^{2t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 1 \\ R_y &= 0 \\ S_t &= -2e^{-2t}y \\ S_y &= e^{-2t}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 1 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{-2t}y = t + c_1$$

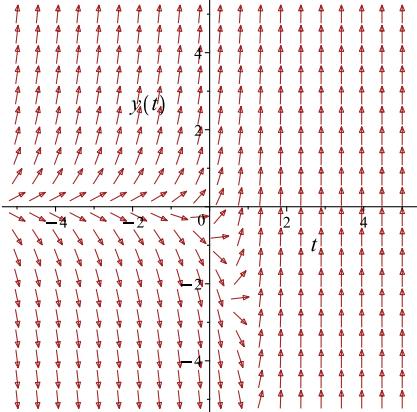
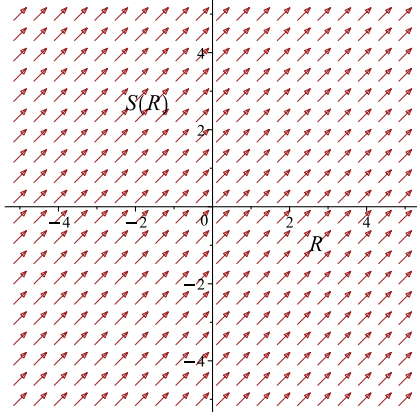
Which simplifies to

$$e^{-2t}y = t + c_1$$

Which gives

$$y = e^{2t}(t + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = 2y + e^{2t}$ 	$R = t$ $S = e^{-2t}y$	$\frac{dS}{dR} = 1$ 

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = 2$ in the above solution gives an equation to solve for the constant of integration.

$$2 = c_1$$

$$c_1 = 2$$

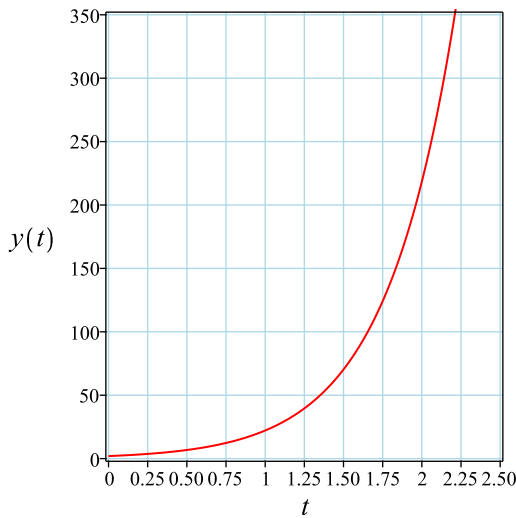
Substituting c_1 found above in the general solution gives

$$y = e^{2t}t + 2e^{2t}$$

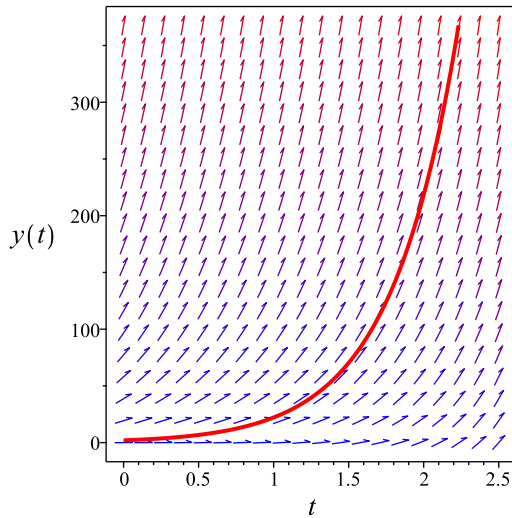
Summary

The solution(s) found are the following

$$y = e^{2t}t + 2e^{2t} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{2t}t + 2e^{2t}$$

Verified OK.

1.17.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} dy &= (2y + e^{2t}) dt \\ (-2y - e^{2t}) dt + dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= -2y - e^{2t} \\ N(t, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (-2y - e^{2t}) \\ &= -2 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (1) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1((-2) - (0)) \\ &= -2 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int -2 dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-2t} \\ &= e^{-2t}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= e^{-2t}(-2y - e^{2t}) \\ &= -2e^{-2t}y - 1\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^{-2t}(1) \\ &= e^{-2t}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ (-2e^{-2t}y - 1) + (e^{-2t}) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -2e^{-2t}y - 1 dt \\ \phi &= -t + e^{-2t}y + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-2t} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-2t}$. Therefore equation (4) becomes

$$e^{-2t} = e^{-2t} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -t + e^{-2t}y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -t + e^{-2t}y$$

The solution becomes

$$y = e^{2t}(t + c_1)$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = 2$ in the above solution gives an equation to solve for the constant of integration.

$$2 = c_1$$

$$c_1 = 2$$

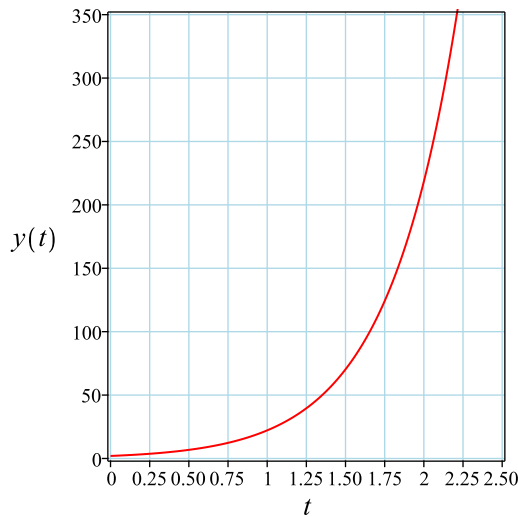
Substituting c_1 found above in the general solution gives

$$y = e^{2t}t + 2e^{2t}$$

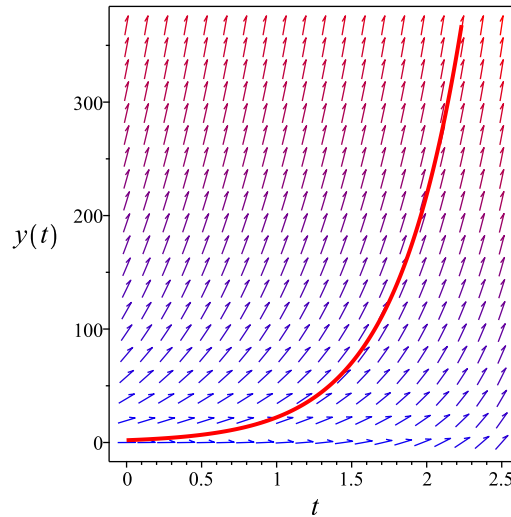
Summary

The solution(s) found are the following

$$y = e^{2t}t + 2e^{2t} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{2t}t + 2e^{2t}$$

Verified OK.

1.17.5 Maple step by step solution

Let's solve

$$[-2y + y' = e^{2t}, y(0) = 2]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = 2y + e^{2t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$-2y + y' = e^{2t}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t)(-2y + y') = \mu(t)e^{2t}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t)(-2y + y') = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -2\mu(t)$$

- Solve to find the integrating factor

$$\mu(t) = e^{-2t}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \mu(t)e^{2t} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \mu(t)e^{2t} dt + c_1$$

- Solve for y

$$y = \frac{\int \mu(t)e^{2t} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{-2t}$

$$y = \frac{\int e^{-2t}e^{2t} dt + c_1}{e^{-2t}}$$

- Evaluate the integrals on the rhs

$$y = \frac{t + c_1}{e^{-2t}}$$

- Simplify

$$y = e^{2t}(t + c_1)$$

- Use initial condition $y(0) = 2$

$$2 = c_1$$

- Solve for c_1

$$c_1 = 2$$

- Substitute $c_1 = 2$ into general solution and simplify

$$y = e^{2t}(2 + t)$$

- Solution to the IVP

$$y = e^{2t}(2 + t)$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 12

```
dsolve([-2*y(t)+diff(y(t),t) = exp(2*t),y(0) = 2],y(t), singsol=all)
```

$$y(t) = (2 + t)e^{2t}$$

✓ Solution by Mathematica

Time used: 0.039 (sec). Leaf size: 14

```
DSolve[{-2*y[t]+y'[t] == Exp[2*t],y[0]==2},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{2t}(t + 2)$$

1.18 problem 18

1.18.1 Existence and uniqueness analysis	225
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Internal problem ID [465]

Internal file name [OUTPUT/465_Sunday_June_05_2022_01_41_54_AM_98318111/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 18.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**linear**", "**exactWithIntegrationFactor**", "**first_order_ode_lie_symmetry_lookup**"

Maple gives the following as the ode type

[_linear]

$$2y + ty' = \sin(t)$$

With initial conditions

$$\left[y\left(\frac{\pi}{2}\right) = 1 \right]$$

1.18.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{2}{t}$$
$$q(t) = \frac{\sin(t)}{t}$$

Hence the ode is

$$y' + \frac{2y}{t} = \frac{\sin(t)}{t}$$

The domain of $p(t) = \frac{2}{t}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = \frac{\pi}{2}$ is inside this domain. The domain of $q(t) = \frac{\sin(t)}{t}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = \frac{\pi}{2}$ is also inside this domain. Hence solution exists and is unique.

1.18.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{2}{t} dt} \\ &= t^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{\sin(t)}{t} \right) \\ \frac{d}{dt}(t^2 y) &= (t^2) \left(\frac{\sin(t)}{t} \right) \\ d(t^2 y) &= (t \sin(t)) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t^2 y &= \int t \sin(t) dt \\ t^2 y &= -t \cos(t) + \sin(t) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t^2$ results in

$$y = \frac{-t \cos(t) + \sin(t)}{t^2} + \frac{c_1}{t^2}$$

which simplifies to

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2}$$

Initial conditions are used to solve for c_1 . Substituting $t = \frac{\pi}{2}$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{4 + 4c_1}{\pi^2}$$

$$c_1 = \frac{\pi^2}{4} - 1$$

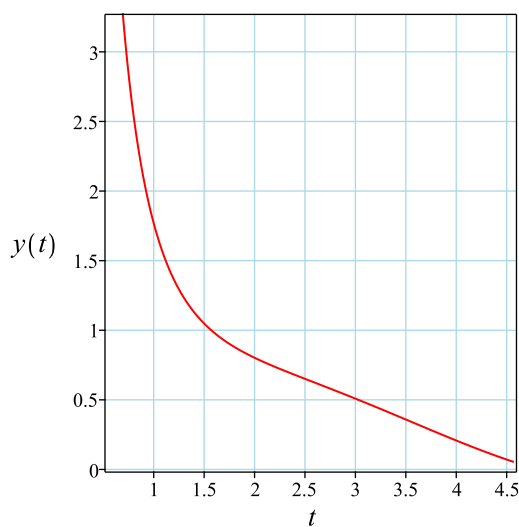
Substituting c_1 found above in the general solution gives

$$y = \frac{-4t \cos(t) + \pi^2 + 4 \sin(t) - 4}{4t^2}$$

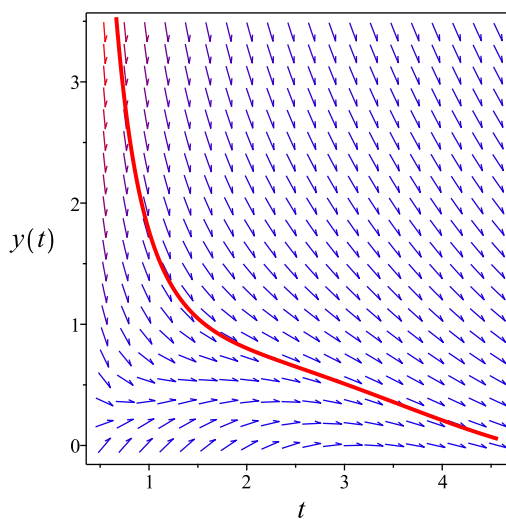
Summary

The solution(s) found are the following

$$y = \frac{-4t \cos(t) + \pi^2 + 4 \sin(t) - 4}{4t^2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{-4t \cos(t) + \pi^2 + 4 \sin(t) - 4}{4t^2}$$

Verified OK.

1.18.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{2y - \sin(t)}{t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 52: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \frac{1}{t^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{t^2}} dy\end{aligned}$$

Which results in

$$S = t^2 y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{2y - \sin(t)}{t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 1 \\R_y &= 0 \\S_t &= 2ty \\S_y &= t^2\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = t \sin(t) \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R \sin(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -R \cos(R) + \sin(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$yt^2 = -t \cos(t) + \sin(t) + c_1$$

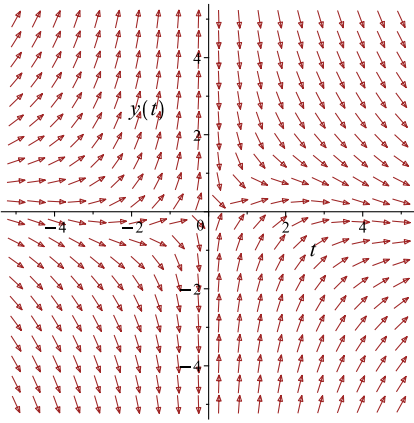
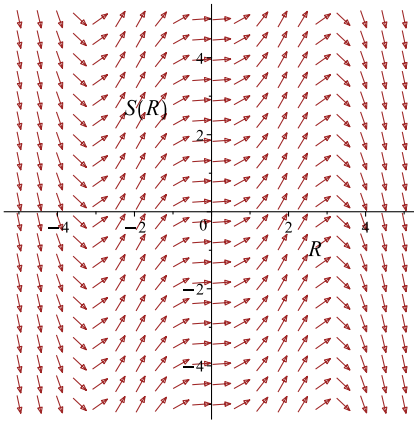
Which simplifies to

$$yt^2 = -t \cos(t) + \sin(t) + c_1$$

Which gives

$$y = -\frac{t \cos(t) - \sin(t) - c_1}{t^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{2y - \sin(t)}{t}$ 	$R = t$ $S = t^2 y$	$\frac{dS}{dR} = R \sin(R)$ 

Initial conditions are used to solve for c_1 . Substituting $t = \frac{\pi}{2}$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{4 + 4c_1}{\pi^2}$$

$$c_1 = \frac{\pi^2}{4} - 1$$

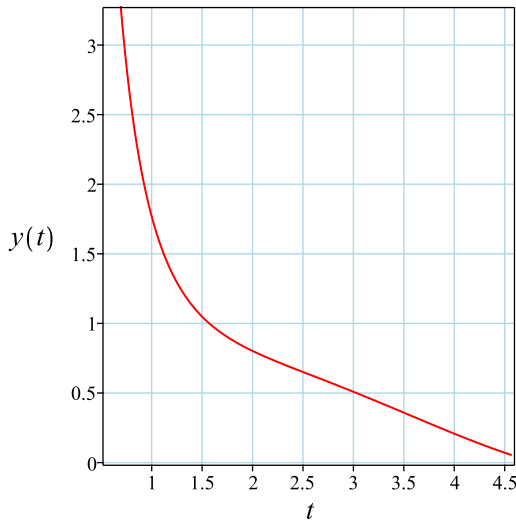
Substituting c_1 found above in the general solution gives

$$y = \frac{-4t \cos(t) + \pi^2 + 4 \sin(t) - 4}{4t^2}$$

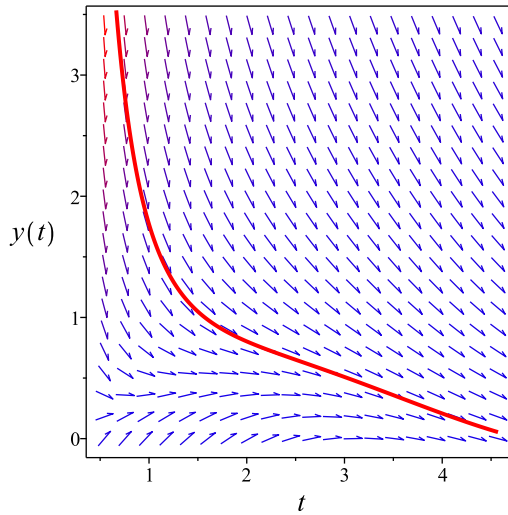
Summary

The solution(s) found are the following

$$y = \frac{-4t \cos(t) + \pi^2 + 4 \sin(t) - 4}{4t^2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{-4t \cos(t) + \pi^2 + 4 \sin(t) - 4}{4t^2}$$

Verified OK.

1.18.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (t) dy &= (-2y + \sin(t)) dt \\ (2y - \sin(t)) dt + (t) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= 2y - \sin(t) \\ N(t, y) &= t \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2y - \sin(t)) \\ &= 2 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(t) \\ &= 1 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{t} ((2) - (1)) \\ &= \frac{1}{t} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int \frac{1}{t} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{\ln(t)} \\ &= t\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= t(2y - \sin(t)) \\ &= (2y - \sin(t)) t\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= t(t) \\ &= t^2\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ ((2y - \sin(t)) t) + (t^2) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int (2y - \sin(t)) t dt \\ \phi &= t^2 y + t \cos(t) - \sin(t) + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t^2 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t^2$. Therefore equation (4) becomes

$$t^2 = t^2 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = t^2 y + t \cos(t) - \sin(t) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = t^2 y + t \cos(t) - \sin(t)$$

The solution becomes

$$y = -\frac{t \cos(t) - \sin(t) - c_1}{t^2}$$

Initial conditions are used to solve for c_1 . Substituting $t = \frac{\pi}{2}$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{4 + 4c_1}{\pi^2}$$

$$c_1 = \frac{\pi^2}{4} - 1$$

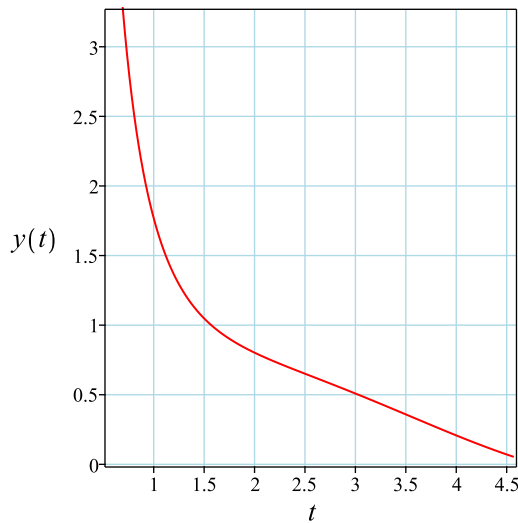
Substituting c_1 found above in the general solution gives

$$y = \frac{-4t \cos(t) + \pi^2 + 4 \sin(t) - 4}{4t^2}$$

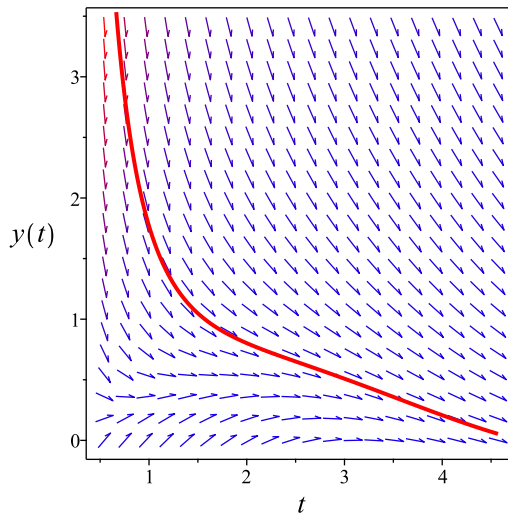
Summary

The solution(s) found are the following

$$y = \frac{-4t \cos(t) + \pi^2 + 4 \sin(t) - 4}{4t^2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{-4t \cos(t) + \pi^2 + 4 \sin(t) - 4}{4t^2}$$

Verified OK.

1.18.5 Maple step by step solution

Let's solve

$$[2y + ty' = \sin(t), y(\frac{\pi}{2}) = 1]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{2y}{t} + \frac{\sin(t)}{t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{2y}{t} = \frac{\sin(t)}{t}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \frac{\mu(t)\sin(t)}{t}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{2\mu(t)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = t^2$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)\sin(t)}{t} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)\sin(t)}{t} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)\sin(t)}{t} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = t^2$

$$y = \frac{\int t \sin(t) dt + c_1}{t^2}$$

- Evaluate the integrals on the rhs

$$y = \frac{-t \cos(t) + \sin(t) + c_1}{t^2}$$

- Use initial condition $y\left(\frac{\pi}{2}\right) = 1$

$$1 = \frac{4(1+c_1)}{\pi^2}$$

- Solve for c_1

$$c_1 = \frac{\pi^2}{4} - 1$$

- Substitute $c_1 = \frac{\pi^2}{4} - 1$ into general solution and simplify

$$y = \frac{-t \cos(t) + \sin(t) + \frac{\pi^2}{4} - 1}{t^2}$$

- Solution to the IVP

$$y = \frac{-t \cos(t) + \sin(t) + \frac{\pi^2}{4} - 1}{t^2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 22

```
dsolve([2*y(t)+t*diff(y(t),t) = sin(t),y(1/2*Pi) = 1],y(t), singsol=all)
```

$$y(t) = \frac{\sin(t) - \cos(t)t + \frac{\pi^2}{4} - 1}{t^2}$$

✓ Solution by Mathematica

Time used: 0.034 (sec). Leaf size: 26

```
DSolve[{2*y[t]+t*y'[t] == Sin[t],y[Pi/2]==1},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{4 \sin(t) - 4t \cos(t) + \pi^2 - 4}{4t^2}$$

1.19 problem 19

1.19.1 Existence and uniqueness analysis	239
1.19.2 Solving as linear ode	240
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Internal problem ID [466]

Internal file name [OUTPUT/466_Sunday_June_05_2022_01_41_55_AM_9788455/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 19.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$4yt^2 + y't^3 = e^{-t}$$

With initial conditions

$$[y(-1) = 0]$$

1.19.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{4}{t}$$
$$q(t) = \frac{e^{-t}}{t^3}$$

Hence the ode is

$$y' + \frac{4y}{t} = \frac{e^{-t}}{t^3}$$

The domain of $p(t) = \frac{4}{t}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = -1$ is inside this domain. The domain of $q(t) = \frac{e^{-t}}{t^3}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = -1$ is also inside this domain. Hence solution exists and is unique.

1.19.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{4}{t} dt} \\ &= t^4\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{e^{-t}}{t^3} \right) \\ \frac{d}{dt}(t^4 y) &= (t^4) \left(\frac{e^{-t}}{t^3} \right) \\ d(t^4 y) &= (t e^{-t}) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t^4 y &= \int t e^{-t} dt \\ t^4 y &= -e^{-t}(t+1) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t^4$ results in

$$y = -\frac{e^{-t}(t+1)}{t^4} + \frac{c_1}{t^4}$$

which simplifies to

$$y = \frac{(-t-1)e^{-t} + c_1}{t^4}$$

Initial conditions are used to solve for c_1 . Substituting $t = -1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = c_1$$

$$c_1 = 0$$

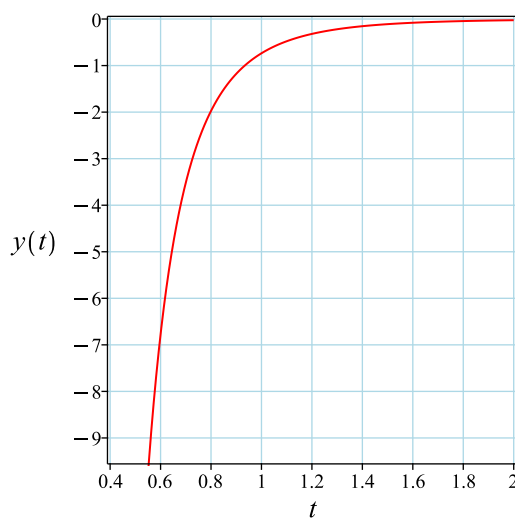
Substituting c_1 found above in the general solution gives

$$y = -\frac{e^{-t}(t+1)}{t^4}$$

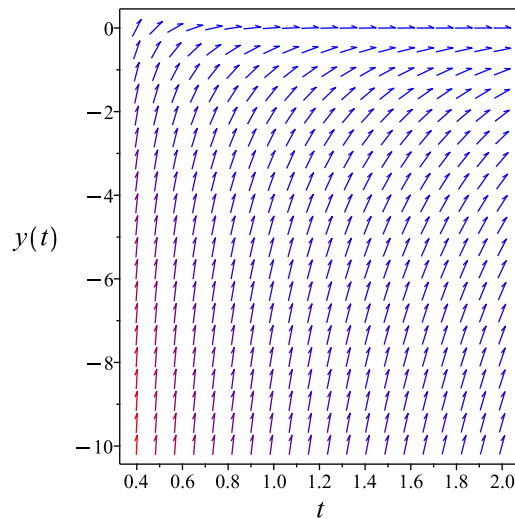
Summary

The solution(s) found are the following

$$y = -\frac{e^{-t}(t+1)}{t^4} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{e^{-t}(t+1)}{t^4}$$

Verified OK.

1.19.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-4t^2y + e^{-t}}{t^3}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 55: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \frac{1}{t^4}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{t^4}} dy\end{aligned}$$

Which results in

$$S = t^4 y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{-4t^2 y + e^{-t}}{t^3}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 1 \\R_y &= 0 \\S_t &= 4t^3y \\S_y &= t^4\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = t e^{-t} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R e^{-R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -(R + 1) e^{-R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$yt^4 = -e^{-t}(t + 1) + c_1$$

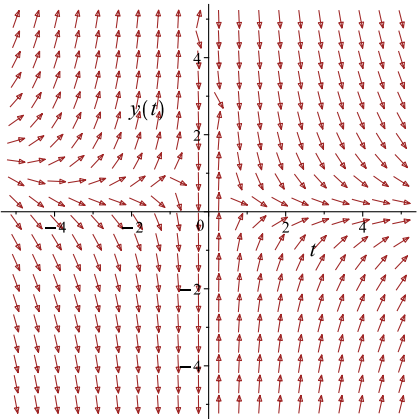
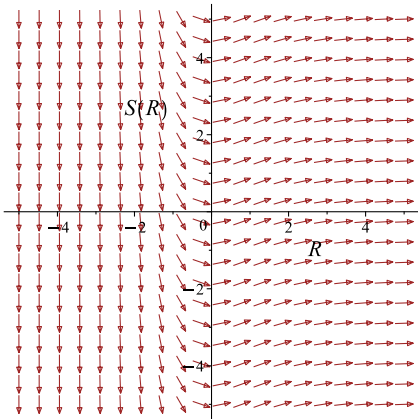
Which simplifies to

$$yt^4 = -e^{-t}(t + 1) + c_1$$

Which gives

$$y = -\frac{te^{-t} + e^{-t} - c_1}{t^4}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{-4t^2y + e^{-t}}{t^3}$ 	$R = t$ $S = t^4 y$	$\frac{dS}{dR} = R e^{-R}$ 

Initial conditions are used to solve for c_1 . Substituting $t = -1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = c_1$$

$$c_1 = 0$$

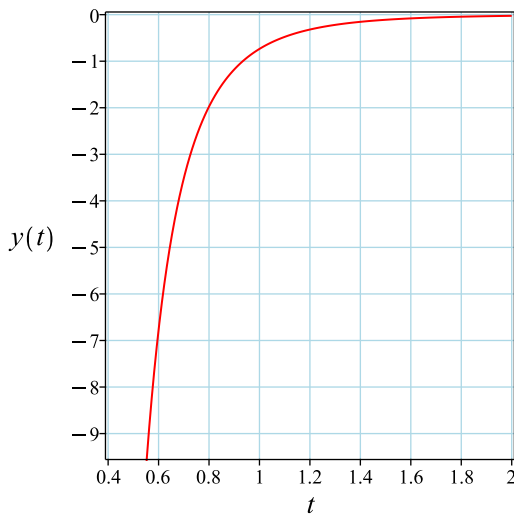
Substituting c_1 found above in the general solution gives

$$y = -\frac{e^{-t}(t + 1)}{t^4}$$

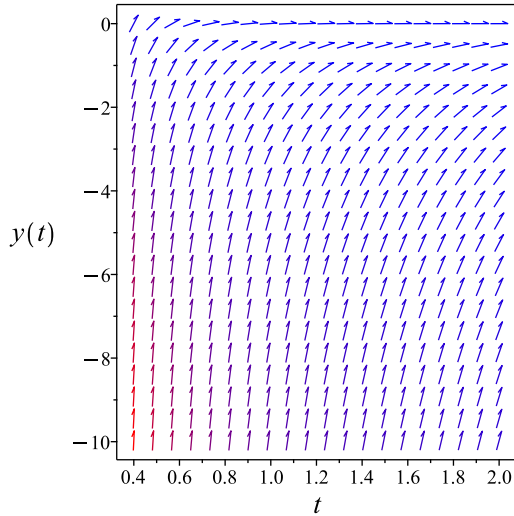
Summary

The solution(s) found are the following

$$y = -\frac{e^{-t}(t + 1)}{t^4} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{e^{-t}(t+1)}{t^4}$$

Verified OK.

1.19.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (t^3) dy &= (-4t^2y + e^{-t}) dt \\ (4t^2y - e^{-t}) dt + (t^3) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= 4t^2y - e^{-t} \\ N(t, y) &= t^3 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (4t^2y - e^{-t}) \\ &= 4t^2 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (t^3) \\ &= 3t^2 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{t^3} ((4t^2) - (3t^2)) \\ &= \frac{1}{t} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int \frac{1}{t} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{\ln(t)} \\ &= t\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= t(4t^2y - e^{-t}) \\ &= 4t^3y - te^{-t}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= t(t^3) \\ &= t^4\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ (4t^3y - te^{-t}) + (t^4) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int 4t^3y - te^{-t} dt \\ \phi &= te^{-t} + e^{-t} + t^4y + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t^4 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t^4$. Therefore equation (4) becomes

$$t^4 = t^4 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = t e^{-t} + e^{-t} + t^4 y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = t e^{-t} + e^{-t} + t^4 y$$

The solution becomes

$$y = -\frac{t e^{-t} + e^{-t} - c_1}{t^4}$$

Initial conditions are used to solve for c_1 . Substituting $t = -1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = c_1$$

$$c_1 = 0$$

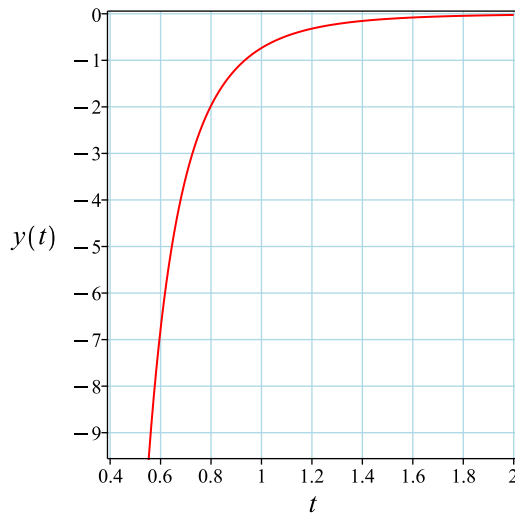
Substituting c_1 found above in the general solution gives

$$y = -\frac{e^{-t}(t+1)}{t^4}$$

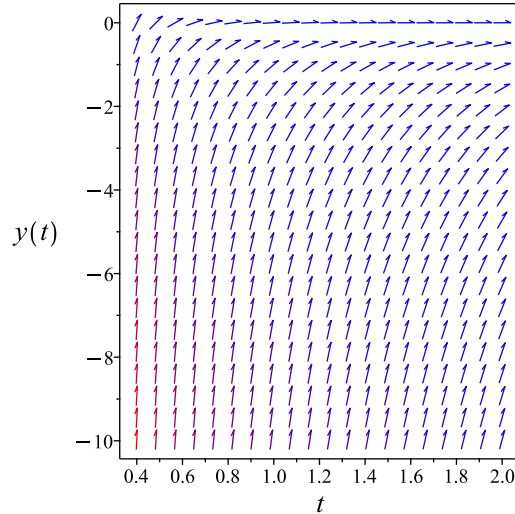
Summary

The solution(s) found are the following

$$y = -\frac{e^{-t}(t+1)}{t^4} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{e^{-t}(t+1)}{t^4}$$

Verified OK.

1.19.5 Maple step by step solution

Let's solve

$$[4yt^2 + y't^3 = e^{-t}, y(-1) = 0]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{4y}{t} + \frac{e^{-t}}{t^3}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{4y}{t} = \frac{e^{-t}}{t^3}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{4y}{t} \right) = \frac{\mu(t)e^{-t}}{t^3}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{4y}{t} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{4\mu(t)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = t^4$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)e^{-t}}{t^3} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)e^{-t}}{t^3} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)e^{-t}}{t^3} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = t^4$

$$y = \frac{\int t e^{-t} dt + c_1}{t^4}$$

- Evaluate the integrals on the rhs

$$y = \frac{-e^{-t}(t+1) + c_1}{t^4}$$

- Simplify

$$y = \frac{(-t-1)e^{-t} + c_1}{t^4}$$

- Use initial condition $y(-1) = 0$

$$0 = c_1$$

- Solve for c_1

$$c_1 = 0$$

- Substitute $c_1 = 0$ into general solution and simplify

$$y = -\frac{e^{-t}(t+1)}{t^4}$$

- Solution to the IVP

$$y = -\frac{e^{-t}(t+1)}{t^4}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 16

```
dsolve([4*t^2*y(t)+t^3*diff(y(t),t) = exp(-t),y(-1) = 0],y(t), singsol=all)
```

$$y(t) = -\frac{(t+1)e^{-t}}{t^4}$$

✓ Solution by Mathematica

Time used: 0.067 (sec). Leaf size: 18

```
DSolve[{4*t^2*y[t]+t^3*y'[t] == Exp[-t],y[-1]==0},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -\frac{e^{-t}(t+1)}{t^4}$$

1.20 problem 20

1.20.1 Existence and uniqueness analysis	253
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1.20.5 Maple step by step solution	264

Internal problem ID [467]

Internal file name [OUTPUT/467_Sunday_June_05_2022_01_41_56_AM_59241620/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 20.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**linear**", "**exactWithIntegrationFactor**", "**first_order_ode_lie_symmetry_lookup**"

Maple gives the following as the ode type

`[_linear]`

$$(t + 1)y + ty' = t$$

With initial conditions

$$[y(\ln(2)) = 1]$$

1.20.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{-t-1}{t}$$

$$q(t) = 1$$

Hence the ode is

$$y' - \frac{(-t-1)y}{t} = 1$$

The domain of $p(t) = -\frac{-t-1}{t}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = \ln(2)$ is inside this domain. The domain of $q(t) = 1$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = \ln(2)$ is also inside this domain. Hence solution exists and is unique.

1.20.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{-t-1}{t} dt} \\ &= e^{t+\ln(t)}\end{aligned}$$

Which simplifies to

$$\mu = t e^t$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= \mu \\ \frac{d}{dt}(t e^t y) &= t e^t \\ d(t e^t y) &= t e^t dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t e^t y &= \int t e^t dt \\ t e^t y &= e^t(-1+t) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t e^t$ results in

$$y = \frac{e^{-t}e^t(-1+t)}{t} + \frac{c_1 e^{-t}}{t}$$

which simplifies to

$$y = \frac{c_1 e^{-t} + t - 1}{t}$$

Initial conditions are used to solve for c_1 . Substituting $t = \ln(2)$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{c_1 + 2 \ln(2) - 2}{2 \ln(2)}$$

$$c_1 = 2$$

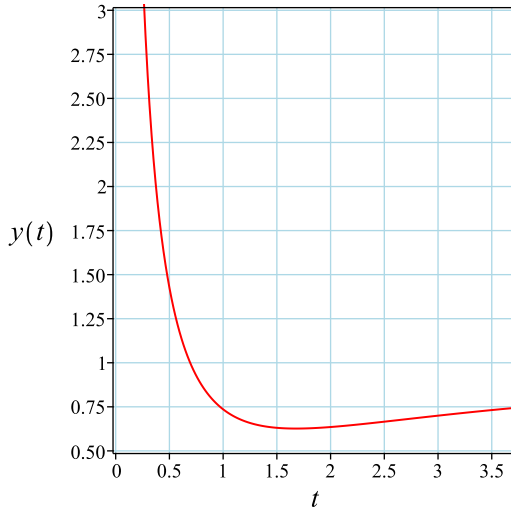
Substituting c_1 found above in the general solution gives

$$y = \frac{-1 + 2e^{-t} + t}{t}$$

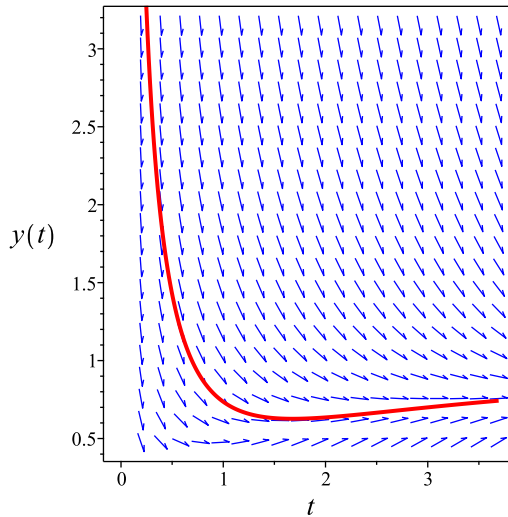
Summary

The solution(s) found are the following

$$y = \frac{-1 + 2e^{-t} + t}{t} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{-1 + 2e^{-t} + t}{t}$$

Verified OK.

1.20.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{ty - t + y}{t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 58: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-t-\ln(t)}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-t-\ln(t)}} dy\end{aligned}$$

Which results in

$$S = t e^t y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y) S_y}{R_t + \omega(t, y) R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{ty - t + y}{t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 1 \\ R_y &= 0 \\ S_t &= y e^t (t + 1) \\ S_y &= t e^t\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = t e^t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R e^R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = (R - 1) e^R + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$t e^t y = e^t (-1 + t) + c_1$$

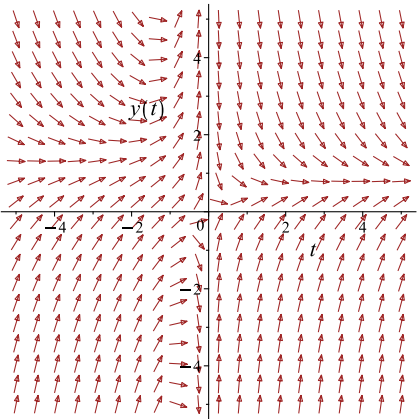
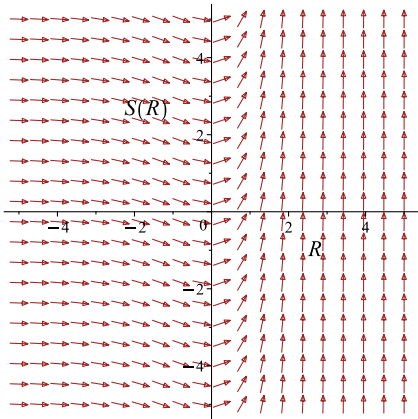
Which simplifies to

$$t e^t y = e^t (-1 + t) + c_1$$

Which gives

$$y = \frac{(t e^t - e^t + c_1) e^{-t}}{t}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{ty-t+y}{t}$ 	$R = t$ $S = te^t y$	$\frac{dS}{dR} = R e^R$ 

Initial conditions are used to solve for c_1 . Substituting $t = \ln(2)$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{c_1 + 2 \ln(2) - 2}{2 \ln(2)}$$

$$c_1 = 2$$

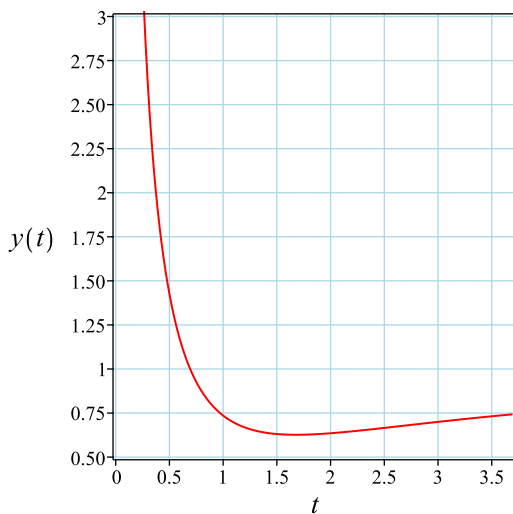
Substituting c_1 found above in the general solution gives

$$y = \frac{e^{-t}e^{tt} - e^{-t}e^t + 2e^{-t}}{t}$$

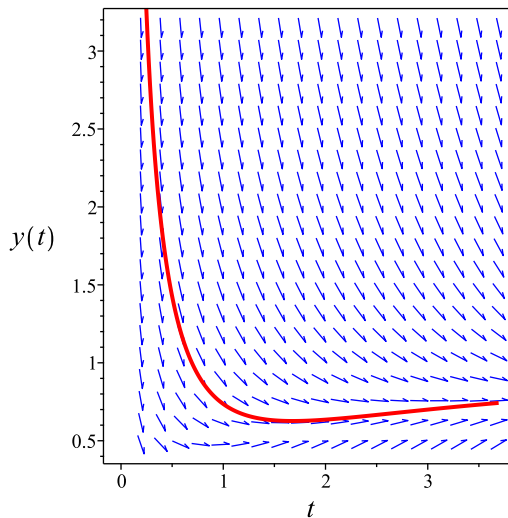
Summary

The solution(s) found are the following

$$y = \frac{e^{-t}e^{tt} - e^{-t}e^t + 2e^{-t}}{t} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^{-t}e^{tt} - e^{-t}e^t + 2e^{-t}}{t}$$

Verified OK.

1.20.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (t) dy &= -(t+1)y + t) dt \\ ((t+1)y - t) dt + (t) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= (t+1)y - t \\ N(t, y) &= t \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}((t+1)y - t) \\ &= t + 1 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(t) \\ &= 1 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{t} ((t+1) - (1)) \\ &= 1 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int 1 dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^t \\ &= e^t\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= e^t((t+1)y - t) \\ &= ((y-1)t + y)e^t\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^t(t) \\ &= te^t\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ (((y-1)t + y)e^t) + (te^t) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int ((y-1)t + y)e^t dt \\ \phi &= (1 + (y-1)t)e^t + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t e^t + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t e^t$. Therefore equation (4) becomes

$$t e^t = t e^t + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = (1 + (y - 1)t) e^t + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = (1 + (y - 1)t) e^t$$

The solution becomes

$$y = \frac{(t e^t - e^t + c_1) e^{-t}}{t}$$

Initial conditions are used to solve for c_1 . Substituting $t = \ln(2)$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{c_1 + 2 \ln(2) - 2}{2 \ln(2)}$$

$$c_1 = 2$$

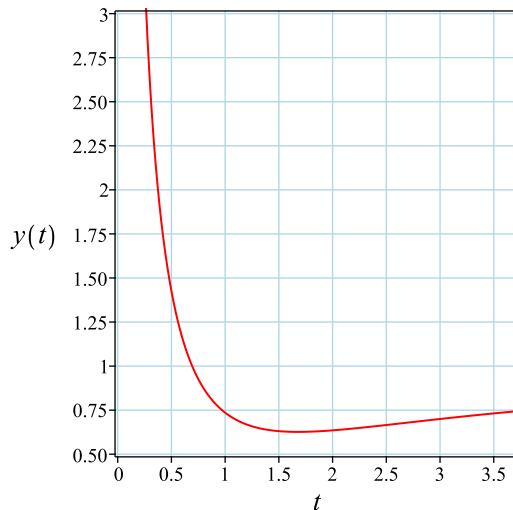
Substituting c_1 found above in the general solution gives

$$y = \frac{e^{-t} e^t t - e^{-t} e^t + 2 e^{-t}}{t}$$

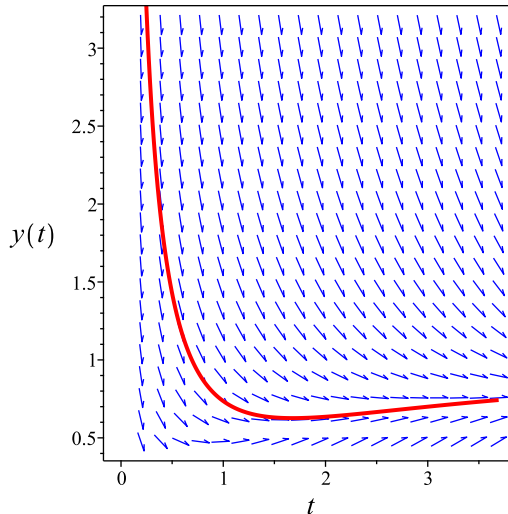
Summary

The solution(s) found are the following

$$y = \frac{e^{-t}e^t t - e^{-t}e^t + 2e^{-t}}{t} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^{-t}e^t t - e^{-t}e^t + 2e^{-t}}{t}$$

Verified OK.

1.20.5 Maple step by step solution

Let's solve

$$[(t+1)y + ty' = t, y(\ln(2)) = 1]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = 1 - \frac{(t+1)y}{t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{(t+1)y}{t} = 1$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{(t+1)y}{t} \right) = \mu(t)$$
- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{(t+1)y}{t} \right) = \mu'(t)y + \mu(t)y'$$
- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)(t+1)}{t}$$
- Solve to find the integrating factor

$$\mu(t) = t e^t$$
- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \mu(t) dt + c_1$$
- Evaluate the integral on the lhs

$$\mu(t)y = \int \mu(t) dt + c_1$$
- Solve for y

$$y = \frac{\int \mu(t) dt + c_1}{\mu(t)}$$
- Substitute $\mu(t) = t e^t$

$$y = \frac{\int t e^t dt + c_1}{t e^t}$$
- Evaluate the integrals on the rhs

$$y = \frac{e^t(-1+t) + c_1}{t e^t}$$
- Simplify

$$y = \frac{c_1 e^{-t} + t - 1}{t}$$
- Use initial condition $y(\ln(2)) = 1$

$$1 = \frac{c_1 + \ln(2) - 1}{\ln(2)}$$
- Solve for c_1

$$c_1 = 2$$
- Substitute $c_1 = 2$ into general solution and simplify

$$y = \frac{-1 + 2e^{-t} + t}{t}$$
- Solution to the IVP

$$y = \frac{-1 + 2e^{-t} + t}{t}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 17

```
dsolve([(1+t)*y(t)+t*diff(y(t),t) = t,y(ln(2)) = 1],y(t), singsol=all)
```

$$y(t) = \frac{t - 1 + 2e^{-t}}{t}$$

✓ Solution by Mathematica

Time used: 0.034 (sec). Leaf size: 23

```
DSolve[{(1+t)*y[t]+t*y'[t]== t,y[Log[2]]==1},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{e^{-t}(e^t(t-1)+2)}{t}$$

1.21 problem 21

1.21.1 Existence and uniqueness analysis	267
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Internal problem ID [468]

Internal file name [OUTPUT/468_Sunday_June_05_2022_01_41_57_AM_10339982/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 21.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$-\frac{y}{2} + y' = 2 \cos(t)$$

With initial conditions

$$[y(0) = a]$$

1.21.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{1}{2}$$
$$q(t) = 2 \cos(t)$$

Hence the ode is

$$-\frac{y}{2} + y' = 2 \cos(t)$$

The domain of $p(t) = -\frac{1}{2}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = 2 \cos(t)$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

1.21.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{1}{2} dt} \\ &= e^{-\frac{t}{2}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu)(2 \cos(t)) \\ \frac{d}{dt}\left(e^{-\frac{t}{2}} y\right) &= \left(e^{-\frac{t}{2}}\right)(2 \cos(t)) \\ d\left(e^{-\frac{t}{2}} y\right) &= \left(2 \cos(t) e^{-\frac{t}{2}}\right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-\frac{t}{2}} y &= \int 2 \cos(t) e^{-\frac{t}{2}} dt \\ e^{-\frac{t}{2}} y &= -\frac{4 \cos(t) e^{-\frac{t}{2}}}{5} + \frac{8 \sin(t) e^{-\frac{t}{2}}}{5} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-\frac{t}{2}}$ results in

$$y = e^{\frac{t}{2}} \left(-\frac{4 \cos(t) e^{-\frac{t}{2}}}{5} + \frac{8 \sin(t) e^{-\frac{t}{2}}}{5} \right) + c_1 e^{\frac{t}{2}}$$

which simplifies to

$$y = c_1 e^{\frac{t}{2}} + \frac{8 \sin(t)}{5} - \frac{4 \cos(t)}{5}$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = c_1 - \frac{4}{5}$$

$$c_1 = \frac{4}{5} + a$$

Substituting c_1 found above in the general solution gives

$$y = \left(\frac{4}{5} + a \right) e^{\frac{t}{2}} + \frac{8 \sin(t)}{5} - \frac{4 \cos(t)}{5}$$

Summary

The solution(s) found are the following

$$y = \left(\frac{4}{5} + a \right) e^{\frac{t}{2}} + \frac{8 \sin(t)}{5} - \frac{4 \cos(t)}{5} \quad (1)$$

Verification of solutions

$$y = \left(\frac{4}{5} + a \right) e^{\frac{t}{2}} + \frac{8 \sin(t)}{5} - \frac{4 \cos(t)}{5}$$

Verified OK.

1.21.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{y}{2} + 2 \cos(t)$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (A)$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 61: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{\frac{t}{2}}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{\frac{t}{2}}} dy \end{aligned}$$

Which results in

$$S = e^{-\frac{t}{2}} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{y}{2} + 2 \cos(t)$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -\frac{e^{-\frac{t}{2}} y}{2} \\ S_y &= e^{-\frac{t}{2}} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 2 \cos(t) e^{-\frac{t}{2}} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 2 \cos(R) e^{-\frac{R}{2}}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 - \frac{4 e^{-\frac{R}{2}} (\cos (R) - 2 \sin (R))}{5} \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{-\frac{t}{2}} y = c_1 - \frac{4 e^{-\frac{t}{2}} (\cos (t) - 2 \sin (t))}{5}$$

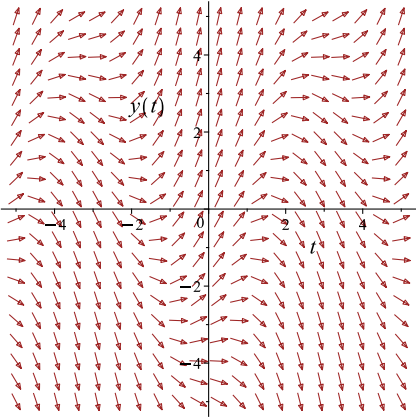
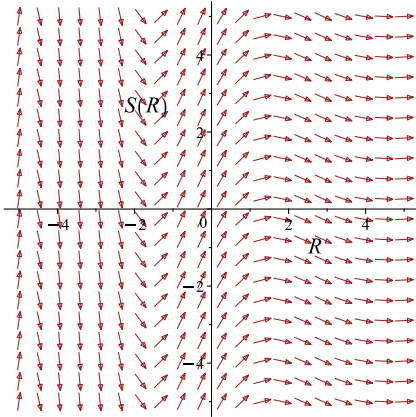
Which simplifies to

$$e^{-\frac{t}{2}} y = c_1 - \frac{4 e^{-\frac{t}{2}} (\cos (t) - 2 \sin (t))}{5}$$

Which gives

$$y = -\frac{e^{\frac{t}{2}} (4 \cos (t) e^{-\frac{t}{2}} - 8 \sin (t) e^{-\frac{t}{2}} - 5c_1)}{5}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{y}{2} + 2 \cos (t)$ 	$R = t$ $S = e^{-\frac{t}{2}} y$	$\frac{dS}{dR} = 2 \cos (R) e^{-\frac{R}{2}}$ 

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = c_1 - \frac{4}{5}$$

$$c_1 = \frac{4}{5} + a$$

Substituting c_1 found above in the general solution gives

$$y = -\frac{\left(4 \cos(t) e^{-\frac{t}{2}} - 8 \sin(t) e^{-\frac{t}{2}} - 4 - 5a\right) e^{\frac{t}{2}}}{5}$$

Summary

The solution(s) found are the following

$$y = -\frac{\left(4 \cos(t) e^{-\frac{t}{2}} - 8 \sin(t) e^{-\frac{t}{2}} - 4 - 5a\right) e^{\frac{t}{2}}}{5} \quad (1)$$

Verification of solutions

$$y = -\frac{\left(4 \cos(t) e^{-\frac{t}{2}} - 8 \sin(t) e^{-\frac{t}{2}} - 4 - 5a\right) e^{\frac{t}{2}}}{5}$$

Verified OK.

1.21.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}dy &= \left(\frac{y}{2} + 2 \cos(t) \right) dt \\ \left(-\frac{y}{2} - 2 \cos(t) \right) dt + dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -\frac{y}{2} - 2 \cos(t) \\ N(t, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{y}{2} - 2 \cos(t) \right) \\ &= -\frac{1}{2}\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1 \left(\left(-\frac{1}{2} \right) - (0) \right) \\ &= -\frac{1}{2}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int -\frac{1}{2} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\frac{t}{2}} \\ &= e^{-\frac{t}{2}}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= e^{-\frac{t}{2}} \left(-\frac{y}{2} - 2 \cos(t) \right) \\ &= -\frac{(y + 4 \cos(t)) e^{-\frac{t}{2}}}{2}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^{-\frac{t}{2}}(1) \\ &= e^{-\frac{t}{2}}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \overline{M} + \overline{N} \frac{dy}{dt} &= 0 \\ \left(-\frac{(y + 4 \cos(t)) e^{-\frac{t}{2}}}{2} \right) + \left(e^{-\frac{t}{2}} \right) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial t} dt &= \int \overline{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -\frac{(y + 4 \cos(t)) e^{-\frac{t}{2}}}{2} dt \\ \phi &= \frac{(5y + 4 \cos(t) - 8 \sin(t)) e^{-\frac{t}{2}}}{5} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-\frac{t}{2}} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-\frac{t}{2}}$. Therefore equation (4) becomes

$$e^{-\frac{t}{2}} = e^{-\frac{t}{2}} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{(5y + 4 \cos(t) - 8 \sin(t)) e^{-\frac{t}{2}}}{5} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{(5y + 4 \cos(t) - 8 \sin(t)) e^{-\frac{t}{2}}}{5}$$

The solution becomes

$$y = -\frac{e^{\frac{t}{2}} \left(4 \cos(t) e^{-\frac{t}{2}} - 8 \sin(t) e^{-\frac{t}{2}} - 5c_1 \right)}{5}$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = c_1 - \frac{4}{5}$$

$$c_1 = \frac{4}{5} + a$$

Substituting c_1 found above in the general solution gives

$$y = -\frac{\left(4 \cos(t) e^{-\frac{t}{2}} - 8 \sin(t) e^{-\frac{t}{2}} - 4 - 5a \right) e^{\frac{t}{2}}}{5}$$

Summary

The solution(s) found are the following

$$y = -\frac{\left(4 \cos(t) e^{-\frac{t}{2}} - 8 \sin(t) e^{-\frac{t}{2}} - 4 - 5a \right) e^{\frac{t}{2}}}{5} \quad (1)$$

Verification of solutions

$$y = -\frac{\left(4 \cos(t) e^{-\frac{t}{2}} - 8 \sin(t) e^{-\frac{t}{2}} - 4 - 5a \right) e^{\frac{t}{2}}}{5}$$

Verified OK.

1.21.5 Maple step by step solution

Let's solve

$$\left[-\frac{y}{2} + y' = 2 \cos(t), y(0) = a\right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = \frac{y}{2} + 2 \cos(t)$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$-\frac{y}{2} + y' = 2 \cos(t)$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(-\frac{y}{2} + y'\right) = 2\mu(t) \cos(t)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(-\frac{y}{2} + y'\right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -\frac{\mu(t)}{2}$$

- Solve to find the integrating factor

$$\mu(t) = e^{-\frac{t}{2}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y)\right) dt = \int 2\mu(t) \cos(t) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int 2\mu(t) \cos(t) dt + c_1$$

- Solve for y

$$y = \frac{\int 2\mu(t) \cos(t) dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{-\frac{t}{2}}$

$$y = \frac{\int 2 \cos(t) e^{-\frac{t}{2}} dt + c_1}{e^{-\frac{t}{2}}}$$

- Evaluate the integrals on the rhs

$$y = \frac{-\frac{4 \cos(t) e^{-\frac{t}{2}}}{5} + \frac{8 \sin(t) e^{-\frac{t}{2}}}{5} + c_1}{e^{-\frac{t}{2}}}$$

- Simplify

$$y = c_1 e^{\frac{t}{2}} + \frac{8 \sin(t)}{5} - \frac{4 \cos(t)}{5}$$

- Use initial condition $y(0) = a$

$$a = c_1 - \frac{4}{5}$$

- Solve for c_1

$$c_1 = \frac{4}{5} + a$$

- Substitute $c_1 = \frac{4}{5} + a$ into general solution and simplify

$$y = e^{\frac{t}{2}} a - \frac{4 \cos(t)}{5} + \frac{8 \sin(t)}{5} + \frac{4 e^{\frac{t}{2}}}{5}$$

- Solution to the IVP

$$y = e^{\frac{t}{2}} a - \frac{4 \cos(t)}{5} + \frac{8 \sin(t)}{5} + \frac{4 e^{\frac{t}{2}}}{5}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 21

```
dsolve([-1/2*y(t)+diff(y(t),t) = 2*cos(t),y(0) = a],y(t), singsol=all)
```

$$y(t) = -\frac{4 \cos(t)}{5} + \frac{8 \sin(t)}{5} + e^{\frac{t}{2}} a + \frac{4 e^{\frac{t}{2}}}{5}$$

✓ Solution by Mathematica

Time used: 0.054 (sec). Leaf size: 31

```
DSolve[{-1/2*y[t]+y'[t] == 2*Cos[t],y[0]==a},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{5}((5a + 4)e^{t/2} + 8 \sin(t) - 4 \cos(t))$$

1.22 problem 22

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Internal problem ID [469]

Internal file name [OUTPUT/469_Sunday_June_05_2022_01_41_59_AM_35727842/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 22.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$-y + 2y' = e^{\frac{t}{3}}$$

With initial conditions

$$[y(0) = a]$$

1.22.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{1}{2}$$
$$q(t) = \frac{e^{\frac{t}{3}}}{2}$$

Hence the ode is

$$-\frac{y}{2} + y' = \frac{e^{\frac{t}{3}}}{2}$$

The domain of $p(t) = -\frac{1}{2}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = \frac{e^{\frac{t}{3}}}{2}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

1.22.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{1}{2} dt} \\ &= e^{-\frac{t}{2}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{e^{\frac{t}{3}}}{2} \right) \\ \frac{d}{dt} \left(e^{-\frac{t}{2}} y \right) &= \left(e^{-\frac{t}{2}} \right) \left(\frac{e^{\frac{t}{3}}}{2} \right) \\ d \left(e^{-\frac{t}{2}} y \right) &= \left(\frac{e^{-\frac{t}{6}}}{2} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-\frac{t}{2}} y &= \int \frac{e^{-\frac{t}{6}}}{2} dt \\ e^{-\frac{t}{2}} y &= -3 e^{-\frac{t}{6}} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-\frac{t}{2}}$ results in

$$y = -3 e^{\frac{t}{2}} e^{-\frac{t}{6}} + c_1 e^{\frac{t}{2}}$$

which simplifies to

$$y = -3e^{\frac{t}{3}} + c_1 e^{\frac{t}{2}}$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = -3 + c_1$$

$$c_1 = 3 + a$$

Substituting c_1 found above in the general solution gives

$$y = -3e^{\frac{t}{3}} + (3 + a)e^{\frac{t}{2}}$$

Summary

The solution(s) found are the following

$$y = -3e^{\frac{t}{3}} + (3 + a)e^{\frac{t}{2}} \quad (1)$$

Verification of solutions

$$y = -3e^{\frac{t}{3}} + (3 + a)e^{\frac{t}{2}}$$

Verified OK.

1.22.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{y}{2} + \frac{e^{\frac{t}{3}}}{2}$$
$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (A)$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 64: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{\frac{t}{2}}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{\frac{t}{2}}} dy \end{aligned}$$

Which results in

$$S = e^{-\frac{t}{2}} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{y}{2} + \frac{e^{\frac{t}{3}}}{2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -\frac{e^{-\frac{t}{2}} y}{2} \\ S_y &= e^{-\frac{t}{2}} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{e^{-\frac{t}{6}}}{2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{e^{-\frac{R}{6}}}{2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -3e^{-\frac{R}{6}} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{-\frac{t}{2}}y = -3e^{-\frac{t}{6}} + c_1$$

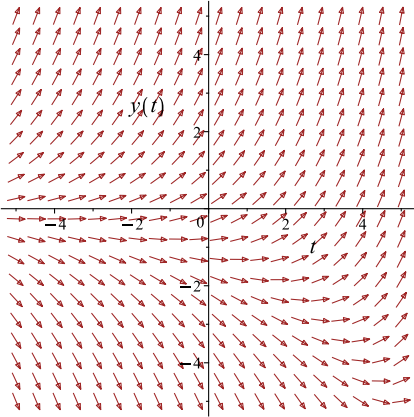
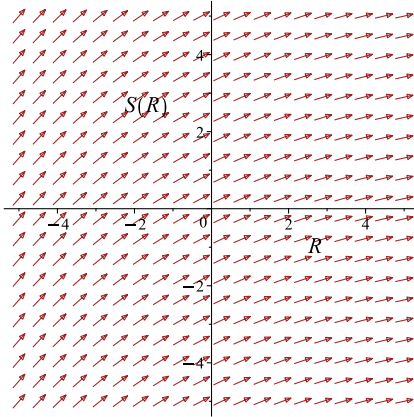
Which simplifies to

$$e^{-\frac{t}{2}}y = -3e^{-\frac{t}{6}} + c_1$$

Which gives

$$y = -\left(3e^{-\frac{t}{6}} - c_1\right)e^{\frac{t}{2}}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{y}{2} + \frac{e^{\frac{t}{3}}}{2}$ 	$R = t$ $S = e^{-\frac{t}{2}}y$	$\frac{dS}{dR} = \frac{e^{-\frac{R}{6}}}{2}$ 

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = -3 + c_1$$

$$c_1 = 3 + a$$

Substituting c_1 found above in the general solution gives

$$y = -\left(3e^{-\frac{t}{6}} - 3 - a\right) e^{\frac{t}{2}}$$

Summary

The solution(s) found are the following

$$y = -\left(3e^{-\frac{t}{6}} - 3 - a\right) e^{\frac{t}{2}} \quad (1)$$

Verification of solutions

$$y = -\left(3e^{-\frac{t}{6}} - 3 - a\right) e^{\frac{t}{2}}$$

Verified OK.

1.22.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (2) dy &= \left(y + e^{\frac{t}{3}} \right) dt \\ \left(-y - e^{\frac{t}{3}} \right) dt + (2) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= -y - e^{\frac{t}{3}} \\ N(t, y) &= 2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-y - e^{\frac{t}{3}} \right) \\ &= -1 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (2) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{2} ((-1) - (0)) \\ &= -\frac{1}{2} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int -\frac{1}{2} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\frac{t}{2}} \\ &= e^{-\frac{t}{2}}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= e^{-\frac{t}{2}} \left(-y - e^{\frac{t}{3}} \right) \\ &= - \left(y + e^{\frac{t}{3}} \right) e^{-\frac{t}{2}}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^{-\frac{t}{2}} (2) \\ &= 2 e^{-\frac{t}{2}}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left(- \left(y + e^{\frac{t}{3}} \right) e^{-\frac{t}{2}} \right) + \left(2 e^{-\frac{t}{2}} \right) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int - \left(y + e^{\frac{t}{3}} \right) e^{-\frac{t}{2}} dt \\ \phi &= 2 e^{-\frac{t}{2}} y + 6 e^{-\frac{t}{6}} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 2e^{-\frac{t}{2}} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 2e^{-\frac{t}{2}}$. Therefore equation (4) becomes

$$2e^{-\frac{t}{2}} = 2e^{-\frac{t}{2}} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = 2e^{-\frac{t}{2}}y + 6e^{-\frac{t}{6}} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = 2e^{-\frac{t}{2}}y + 6e^{-\frac{t}{6}}$$

The solution becomes

$$y = -\frac{(6e^{-\frac{t}{6}} - c_1)e^{\frac{t}{2}}}{2}$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = -3 + \frac{c_1}{2}$$

$$c_1 = 6 + 2a$$

Substituting c_1 found above in the general solution gives

$$y = -\frac{\left(6e^{-\frac{t}{6}} - 6 - 2a\right)e^{\frac{t}{2}}}{2}$$

Summary

The solution(s) found are the following

$$y = -\frac{\left(6e^{-\frac{t}{6}} - 6 - 2a\right)e^{\frac{t}{2}}}{2} \quad (1)$$

Verification of solutions

$$y = -\frac{\left(6e^{-\frac{t}{6}} - 6 - 2a\right)e^{\frac{t}{2}}}{2}$$

Verified OK.

1.22.5 Maple step by step solution

Let's solve

$$\left[-y + 2y' = e^{\frac{t}{3}}, y(0) = a\right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = \frac{y}{2} + \frac{e^{\frac{t}{3}}}{2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$-\frac{y}{2} + y' = \frac{e^{\frac{t}{3}}}{2}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(-\frac{y}{2} + y'\right) = \frac{\mu(t)e^{\frac{t}{3}}}{2}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(-\frac{y}{2} + y'\right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -\frac{\mu(t)}{2}$$

- Solve to find the integrating factor

$$\mu(t) = e^{-\frac{t}{2}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)e^{\frac{t}{3}}}{2} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)e^{\frac{t}{3}}}{2} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)e^{\frac{t}{3}}}{2} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{-\frac{t}{2}}$

$$y = \frac{\int \frac{e^{\frac{t}{3}}e^{-\frac{t}{2}}}{2} dt + c_1}{e^{-\frac{t}{2}}}$$

- Evaluate the integrals on the rhs

$$y = \frac{-3e^{-\frac{t}{6}} + c_1}{e^{-\frac{t}{2}}}$$

- Simplify

$$y = e^{\frac{t}{3}} \left(e^{\frac{t}{6}} c_1 - 3 \right)$$

- Use initial condition $y(0) = a$

$$a = -3 + c_1$$

- Solve for c_1

$$c_1 = 3 + a$$

- Substitute $c_1 = 3 + a$ into general solution and simplify

$$y = e^{\frac{t}{3}} \left(e^{\frac{t}{6}} (3 + a) - 3 \right)$$

- Solution to the IVP

$$y = e^{\frac{t}{3}} \left(e^{\frac{t}{6}} (3 + a) - 3 \right)$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 18

```
dsolve([-y(t)+2*diff(y(t),t) = exp(1/3*t),y(0) = a],y(t), singsol=all)
```

$$y(t) = e^{\frac{t}{3}} \left(-3 + (a + 3) e^{\frac{t}{6}} \right)$$

✓ Solution by Mathematica

Time used: 0.052 (sec). Leaf size: 26

```
DSolve[{-y[t]+2*y'[t] == Exp[1/3*t],y[0]==a},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{t/3} \left((a + 3) e^{t/6} - 3 \right)$$

1.23 problem 23

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Internal problem ID [470]

Internal file name [OUTPUT/470_Sunday_June_05_2022_01_41_59_AM_31829908/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 23.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$-2y + 3y' = e^{-\frac{\pi t}{2}}$$

With initial conditions

$$[y(0) = a]$$

1.23.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{2}{3}$$
$$q(t) = \frac{e^{-\frac{\pi t}{2}}}{3}$$

Hence the ode is

$$y' - \frac{2y}{3} = \frac{e^{-\frac{\pi t}{2}}}{3}$$

The domain of $p(t) = -\frac{2}{3}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = \frac{e^{-\frac{\pi t}{2}}}{3}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

1.23.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{2}{3} dt} \\ &= e^{-\frac{2t}{3}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{e^{-\frac{\pi t}{2}}}{3} \right) \\ \frac{d}{dt} \left(e^{-\frac{2t}{3}} y \right) &= \left(e^{-\frac{2t}{3}} \right) \left(\frac{e^{-\frac{\pi t}{2}}}{3} \right) \\ d \left(e^{-\frac{2t}{3}} y \right) &= \left(\frac{e^{t(-\frac{\pi}{2} - \frac{2}{3})}}{3} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-\frac{2t}{3}} y &= \int \frac{e^{t(-\frac{\pi}{2} - \frac{2}{3})}}{3} dt \\ e^{-\frac{2t}{3}} y &= -\frac{2 e^{-\frac{t(3\pi+4)}{6}}}{3\pi+4} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-\frac{2t}{3}}$ results in

$$y = -\frac{2 e^{\frac{2t}{3}} e^{-\frac{t(3\pi+4)}{6}}}{3\pi+4} + c_1 e^{\frac{2t}{3}}$$

which simplifies to

$$y = \frac{-2e^{-\frac{\pi t}{2}} + (3\pi + 4)c_1 e^{\frac{2t}{3}}}{3\pi + 4}$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = \frac{3\pi c_1 + 4c_1 - 2}{3\pi + 4}$$

$$c_1 = \frac{3\pi a + 4a + 2}{3\pi + 4}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{-2e^{-\frac{\pi t}{2}} + (3\pi a + 4a + 2)e^{\frac{2t}{3}}}{3\pi + 4}$$

Summary

The solution(s) found are the following

$$y = \frac{-2e^{-\frac{\pi t}{2}} + (3\pi a + 4a + 2)e^{\frac{2t}{3}}}{3\pi + 4} \quad (1)$$

Verification of solutions

$$y = \frac{-2e^{-\frac{\pi t}{2}} + (3\pi a + 4a + 2)e^{\frac{2t}{3}}}{3\pi + 4}$$

Verified OK.

1.23.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{2y}{3} + \frac{e^{-\frac{\pi t}{2}}}{3}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (A)$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 67: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{\frac{2t}{3}}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{\frac{2t}{3}}} dy \end{aligned}$$

Which results in

$$S = e^{-\frac{2t}{3}} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{2y}{3} + \frac{e^{-\frac{\pi t}{2}}}{3}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -\frac{2e^{-\frac{2t}{3}}y}{3} \\ S_y &= e^{-\frac{2t}{3}} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{e^{t(-\frac{\pi}{2} - \frac{2}{3})}}{3} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{e^{R(-\frac{\pi}{2} - \frac{2}{3})}}{3}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{2e^{-\frac{R(3\pi+4)}{6}}}{3\pi+4} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{-\frac{2t}{3}}y = -\frac{2e^{-\frac{t(3\pi+4)}{6}}}{3\pi+4} + c_1$$

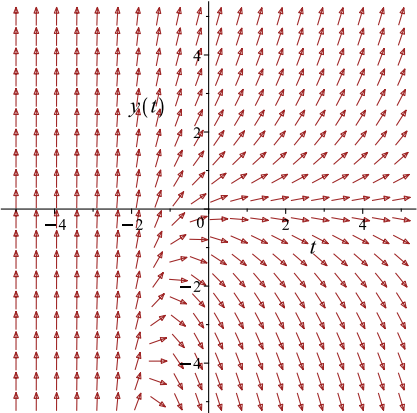
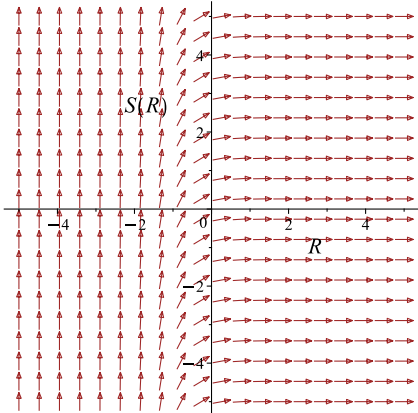
Which simplifies to

$$e^{-\frac{2t}{3}}y = -\frac{2e^{-\frac{t(3\pi+4)}{6}}}{3\pi+4} + c_1$$

Which gives

$$y = \frac{(3\pi c_1 - 2e^{-\frac{t(3\pi+4)}{6}} + 4c_1)e^{\frac{2t}{3}}}{3\pi+4}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{2y}{3} + \frac{e^{-\frac{\pi t}{2}}}{3}$ 	$R = t$ $S = e^{-\frac{2t}{3}}y$	$\frac{dS}{dR} = \frac{e^{R(-\frac{\pi}{2}-\frac{2}{3})}}{3}$ 

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = \frac{3\pi c_1 + 4c_1 - 2}{3\pi + 4}$$

$$c_1 = \frac{3\pi a + 4a + 2}{3\pi + 4}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{\left(\frac{3\pi(3\pi a + 4a + 2)}{3\pi + 4} - 2e^{-\frac{t(3\pi + 4)}{6}} + \frac{12\pi a + 16a + 8}{3\pi + 4} \right) e^{\frac{2t}{3}}}{3\pi + 4}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(\frac{3\pi(3\pi a + 4a + 2)}{3\pi + 4} - 2e^{-\frac{t(3\pi + 4)}{6}} + \frac{12\pi a + 16a + 8}{3\pi + 4} \right) e^{\frac{2t}{3}}}{3\pi + 4} \quad (1)$$

Verification of solutions

$$y = \frac{\left(\frac{3\pi(3\pi a + 4a + 2)}{3\pi + 4} - 2e^{-\frac{t(3\pi + 4)}{6}} + \frac{12\pi a + 16a + 8}{3\pi + 4} \right) e^{\frac{2t}{3}}}{3\pi + 4}$$

Verified OK.

1.23.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(3) dy &= \left(2y + e^{-\frac{\pi t}{2}}\right) dt \\ \left(-2y - e^{-\frac{\pi t}{2}}\right) dt + (3) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -2y - e^{-\frac{\pi t}{2}} \\ N(t, y) &= 3\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-2y - e^{-\frac{\pi t}{2}}\right) \\ &= -2\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (3) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{3}((-2) - (0)) \\ &= -\frac{2}{3} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int -\frac{2}{3} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-\frac{2t}{3}} \\ &= e^{-\frac{2t}{3}} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \overline{M} &= \mu M \\ &= e^{-\frac{2t}{3}} \left(-2y - e^{-\frac{\pi t}{2}} \right) \\ &= - \left(2y + e^{-\frac{\pi t}{2}} \right) e^{-\frac{2t}{3}} \end{aligned}$$

And

$$\begin{aligned} \overline{N} &= \mu N \\ &= e^{-\frac{2t}{3}} (3) \\ &= 3e^{-\frac{2t}{3}} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \overline{M} + \overline{N} \frac{dy}{dt} &= 0 \\ \left(- \left(2y + e^{-\frac{\pi t}{2}} \right) e^{-\frac{2t}{3}} \right) + \left(3e^{-\frac{2t}{3}} \right) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial t} dt &= \int \overline{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -\left(2y + e^{-\frac{\pi t}{2}}\right) e^{-\frac{2t}{3}} dt \\ \phi &= \frac{6e^{t(-\frac{\pi}{2}-\frac{2}{3})} + 3y(3\pi + 4)e^{-\frac{2t}{3}}}{3\pi + 4} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 3e^{-\frac{2t}{3}} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 3e^{-\frac{2t}{3}}$. Therefore equation (4) becomes

$$3e^{-\frac{2t}{3}} = 3e^{-\frac{2t}{3}} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{6e^{t(-\frac{\pi}{2}-\frac{2}{3})} + 3y(3\pi + 4)e^{-\frac{2t}{3}}}{3\pi + 4} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{6e^{t(-\frac{\pi}{2}-\frac{2}{3})} + 3y(3\pi + 4)e^{-\frac{2t}{3}}}{3\pi + 4}$$

The solution becomes

$$y = \frac{\left(3\pi c_1 - 6 e^{-\frac{t(3\pi+4)}{6}} + 4c_1\right) e^{\frac{2t}{3}}}{9\pi + 12}$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = \frac{3\pi c_1 + 4c_1 - 6}{9\pi + 12}$$

$$c_1 = \frac{9\pi a + 12a + 6}{3\pi + 4}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{\left(\frac{9\pi(3\pi a+4a+2)}{3\pi+4} - 6 e^{-\frac{t(3\pi+4)}{6}} + \frac{36\pi a+48a+24}{3\pi+4}\right) e^{\frac{2t}{3}}}{9\pi + 12}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(\frac{9\pi(3\pi a+4a+2)}{3\pi+4} - 6 e^{-\frac{t(3\pi+4)}{6}} + \frac{36\pi a+48a+24}{3\pi+4}\right) e^{\frac{2t}{3}}}{9\pi + 12} \quad (1)$$

Verification of solutions

$$y = \frac{\left(\frac{9\pi(3\pi a+4a+2)}{3\pi+4} - 6 e^{-\frac{t(3\pi+4)}{6}} + \frac{36\pi a+48a+24}{3\pi+4}\right) e^{\frac{2t}{3}}}{9\pi + 12}$$

Verified OK.

1.23.5 Maple step by step solution

Let's solve

$$\left[-2y + 3y' = e^{-\frac{\pi t}{2}}, y(0) = a\right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = \frac{2y}{3} + \frac{e^{-\frac{\pi t}{2}}}{3}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' - \frac{2y}{3} = e^{-\frac{\pi t}{2}}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' - \frac{2y}{3} \right) = \frac{\mu(t)e^{-\frac{\pi t}{2}}}{3}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' - \frac{2y}{3} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -\frac{2\mu(t)}{3}$$

- Solve to find the integrating factor

$$\mu(t) = e^{-\frac{2t}{3}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)e^{-\frac{\pi t}{2}}}{3} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)e^{-\frac{\pi t}{2}}}{3} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)e^{-\frac{\pi t}{2}}}{3} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{-\frac{2t}{3}}$

$$y = \frac{\int \frac{e^{-\frac{\pi t}{2}} e^{-\frac{2t}{3}}}{3} dt + c_1}{e^{-\frac{2t}{3}}}$$

- Evaluate the integrals on the rhs

$$y = \frac{-\frac{2}{3} e^{-\frac{1}{2}\pi t - \frac{2}{3}t}}{3\pi + 4} + c_1}{e^{-\frac{2t}{3}}}$$

- Simplify

$$y = \frac{\left(3\pi c_1 - 2 e^{t\left(-\frac{\pi}{2} - \frac{2}{3}\right)} + 4c_1 \right) e^{\frac{2t}{3}}}{3\pi + 4}$$

- Use initial condition $y(0) = a$

$$a = \frac{3\pi c_1 + 4c_1 - 2}{3\pi + 4}$$

- Solve for c_1

$$c_1 = \frac{3\pi a + 4a + 2}{3\pi + 4}$$

- Substitute $c_1 = \frac{3\pi a + 4a + 2}{3\pi + 4}$ into general solution and simplify

$$y = \frac{\left(3\pi a - 2e^{t\left(-\frac{\pi}{2} - \frac{2}{3}\right)} + 4a + 2\right)e^{\frac{2t}{3}}}{3\pi + 4}$$

- Solution to the IVP

$$y = \frac{\left(3\pi a - 2e^{t\left(-\frac{\pi}{2} - \frac{2}{3}\right)} + 4a + 2\right)e^{\frac{2t}{3}}}{3\pi + 4}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 35

```
dsolve([-2*y(t)+3*diff(y(t),t) = exp(-1/2*Pi*t),y(0) = a],y(t), singsol=all)
```

$$y(t) = \frac{\left(3\pi a - 2e^{t\left(-\frac{\pi}{2} - \frac{2}{3}\right)} + 4a + 2\right)e^{\frac{2t}{3}}}{3\pi + 4}$$

✓ Solution by Mathematica

Time used: 0.083 (sec). Leaf size: 43

```
DSolve[{-2*y[t]+3*y'[t] == Exp[-1/2*Pi*t],y[0]==a},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{e^{2t/3} \left((4 + 3\pi)a - 2e^{-\frac{1}{6}(4+3\pi)t} + 2 \right)}{4 + 3\pi}$$

1.24 problem 24

1.24.1 Existence and uniqueness analysis	306
1.24.2 Solving as linear ode	307
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Internal problem ID [471]

Internal file name [OUTPUT/471_Sunday_June_05_2022_01_42_00_AM_12055452/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 24.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "linear", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$(t + 1)y + ty' = 2te^{-t}$$

With initial conditions

$$[y(1) = a]$$

1.24.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{-t-1}{t}$$
$$q(t) = 2e^{-t}$$

Hence the ode is

$$y' - \frac{(-t-1)y}{t} = 2e^{-t}$$

The domain of $p(t) = -\frac{-t-1}{t}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = 1$ is inside this domain. The domain of $q(t) = 2e^{-t}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 1$ is also inside this domain. Hence solution exists and is unique.

1.24.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{-t-1}{t} dt} \\ &= e^{t+\ln(t)}\end{aligned}$$

Which simplifies to

$$\mu = t e^t$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (2e^{-t}) \\ \frac{d}{dt}(t e^t y) &= (t e^t) (2e^{-t}) \\ d(t e^t y) &= (2t) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t e^t y &= \int 2t dt \\ t e^t y &= t^2 + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t e^t$ results in

$$y = t e^{-t} + \frac{c_1 e^{-t}}{t}$$

which simplifies to

$$y = \frac{e^{-t}(t^2 + c_1)}{t}$$

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = e^{-1}(1 + c_1)$$

$$c_1 = -(e^{-1} - a) e$$

Substituting c_1 found above in the general solution gives

$$y = \frac{e^{-t}(t^2 - (e^{-1} - a) e)}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{e^{-t}(t^2 - (e^{-1} - a) e)}{t} \tag{1}$$

Verification of solutions

$$y = \frac{e^{-t}(t^2 - (e^{-1} - a) e)}{t}$$

Verified OK.

1.24.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{(t e^t y + y e^t - 2t) e^{-t}}{t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 70: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-t-\ln(t)}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-t-\ln(t)}} dy \end{aligned}$$

Which results in

$$S = t e^t y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{(t e^t y + y e^t - 2t) e^{-t}}{t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= y e^t (t + 1) \\ S_y &= t e^t \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 2t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 2R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R^2 + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$t e^t y = t^2 + c_1$$

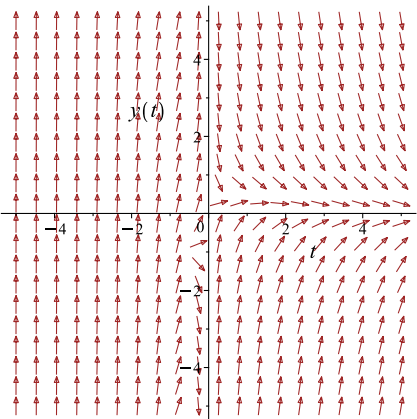
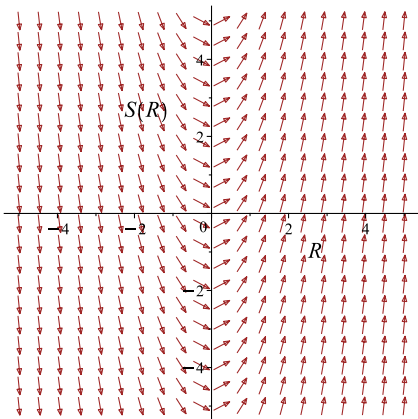
Which simplifies to

$$t e^t y = t^2 + c_1$$

Which gives

$$y = \frac{e^{-t}(t^2 + c_1)}{t}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{(t e^t y + y e^t - 2t) e^{-t}}{t}$ 	$R = t$ $S = t e^t y$	$\frac{dS}{dR} = 2R$ 

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = e^{-1} c_1 + e^{-1}$$

$$c_1 = -(e^{-1} - a) e$$

Substituting c_1 found above in the general solution gives

$$y = \frac{e^{-t}(t^2 - (e^{-1} - a) e)}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{e^{-t}(t^2 - (e^{-1} - a) e)}{t} \quad (1)$$

Verification of solutions

$$y = \frac{e^{-t}(t^2 - (e^{-1} - a) e)}{t}$$

Verified OK.

1.24.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (t e^t) dy &= (-t e^t y - y e^t + 2t) dt \\ (t e^t y + y e^t - 2t) dt + (t e^t) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= t e^t y + y e^t - 2t \\ N(t, y) &= t e^t \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (t e^t y + y e^t - 2t) \\ &= e^t (t + 1) \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (t e^t) \\ &= e^t (t + 1) \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int t e^t y + y e^t - 2t dt \\ \phi &= -t(-y e^t + t) + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t e^t + f'(y)\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t e^t$. Therefore equation (4) becomes

$$t e^t = t e^t + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -t(-y e^t + t) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -t(-y e^t + t)$$

The solution becomes

$$y = \frac{e^{-t}(t^2 + c_1)}{t}$$

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = e^{-1}c_1 + e^{-1}$$

$$c_1 = -(e^{-1} - a) e$$

Substituting c_1 found above in the general solution gives

$$y = \frac{e^{-t}(t^2 - (e^{-1} - a) e)}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{e^{-t}(t^2 - (e^{-1} - a) e)}{t} \tag{1}$$

Verification of solutions

$$y = \frac{e^{-t}(t^2 - (e^{-1} - a) e)}{t}$$

Verified OK.

1.24.5 Maple step by step solution

Let's solve

$$[(t + 1)y + ty' = \frac{2t}{e^t}, y(1) = a]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{(t+1)y}{t} + \frac{2}{e^t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{(t+1)y}{t} = \frac{2}{e^t}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{(t+1)y}{t} \right) = \frac{2\mu(t)}{e^t}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{(t+1)y}{t} \right) = \mu'(t) y + \mu(t) y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)(t+1)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = t(e^t)^2 e^{-t}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y) \right) dt = \int \frac{2\mu(t)}{e^t} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t) y = \int \frac{2\mu(t)}{e^t} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{2\mu(t)}{e^t} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = t(e^t)^2 e^{-t}$

$$y = \frac{\int 2e^{-t} e^t t dt + c_1}{t(e^t)^2 e^{-t}}$$

- Evaluate the integrals on the rhs

$$y = \frac{t^2 + c_1}{t(e^t)^2 e^{-t}}$$

- Simplify

$$y = \frac{e^{-t}(t^2 + c_1)}{t}$$

- Use initial condition $y(1) = a$

$$a = e^{-1}(1 + c_1)$$

- Solve for c_1

$$c_1 = -\frac{e^{-1} - a}{e^{-1}}$$

- Substitute $c_1 = -\frac{e^{-1} - a}{e^{-1}}$ into general solution and simplify

$$y = \frac{e^{-t}(-1 + ea + t^2)}{t}$$

- Solution to the IVP

$$y = \frac{e^{-t}(-1 + ea + t^2)}{t}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 21

```
dsolve([(1+t)*y(t)+t*diff(y(t),t) = 2*t/exp(t),y(1) = a],y(t), singsol=all)
```

$$y(t) = \frac{(t^2 + a e - 1) e^{-t}}{t}$$

✓ Solution by Mathematica

Time used: 0.065 (sec). Leaf size: 22

```
DSolve[{(1+t)*y[t]+t*y'[t] == 2*t/Exp[t],y[1]==a},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{e^{-t}(ea + t^2 - 1)}{t}$$

1.25 problem 25

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Internal problem ID [472]

Internal file name [OUTPUT/472_Sunday_June_05_2022_01_42_01_AM_79347504/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 25.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "linear", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$2y + ty' = \frac{\sin(t)}{t}$$

With initial conditions

$$\left[y\left(-\frac{\pi}{2}\right) = a \right]$$

1.25.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{2}{t}$$
$$q(t) = \frac{\sin(t)}{t^2}$$

Hence the ode is

$$y' + \frac{2y}{t} = \frac{\sin(t)}{t^2}$$

The domain of $p(t) = \frac{2}{t}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = -\frac{\pi}{2}$ is inside this domain. The domain of $q(t) = \frac{\sin(t)}{t^2}$ is

$$\{t < 0 \vee 0 < t\}$$

And the point $t_0 = -\frac{\pi}{2}$ is also inside this domain. Hence solution exists and is unique.

1.25.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{2}{t} dt} \\ &= t^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{\sin(t)}{t^2} \right) \\ \frac{d}{dt}(t^2 y) &= (t^2) \left(\frac{\sin(t)}{t^2} \right) \\ d(t^2 y) &= \sin(t) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t^2 y &= \int \sin(t) dt \\ t^2 y &= -\cos(t) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t^2$ results in

$$y = -\frac{\cos(t)}{t^2} + \frac{c_1}{t^2}$$

which simplifies to

$$y = \frac{-\cos(t) + c_1}{t^2}$$

Initial conditions are used to solve for c_1 . Substituting $t = -\frac{\pi}{2}$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = \frac{4c_1}{\pi^2}$$

$$c_1 = \frac{\pi^2 a}{4}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{-\cos(t) + \frac{\pi^2 a}{4}}{t^2}$$

Summary

The solution(s) found are the following

$$y = \frac{-\cos(t) + \frac{\pi^2 a}{4}}{t^2} \tag{1}$$

Verification of solutions

$$y = \frac{-\cos(t) + \frac{\pi^2 a}{4}}{t^2}$$

Verified OK.

1.25.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-2ty + \sin(t)}{t^2}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 73: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int(n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \frac{1}{t^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{t^2}} dy \end{aligned}$$

Which results in

$$S = t^2 y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{-2ty + \sin(t)}{t^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= 2ty \\ S_y &= t^2 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \sin(t) \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \sin(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\cos(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$yt^2 = -\cos(t) + c_1$$

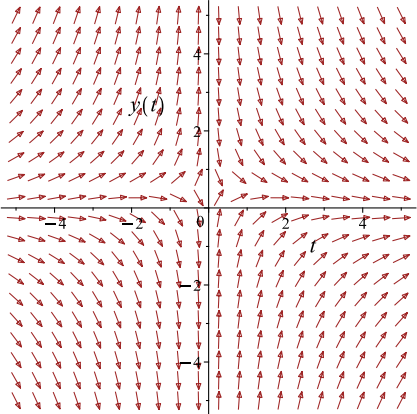
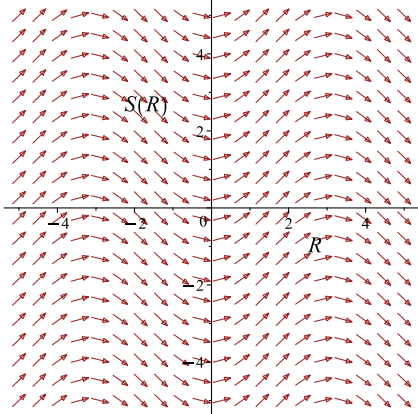
Which simplifies to

$$yt^2 = -\cos(t) + c_1$$

Which gives

$$y = -\frac{\cos(t) - c_1}{t^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{-2ty + \sin(t)}{t^2}$ 	$R = t$ $S = t^2 y$	$\frac{dS}{dR} = \sin(R)$ 

Initial conditions are used to solve for c_1 . Substituting $t = -\frac{\pi}{2}$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = \frac{4c_1}{\pi^2}$$

$$c_1 = \frac{\pi^2 a}{4}$$

Substituting c_1 found above in the general solution gives

$$y = -\frac{\cos(t) - \frac{\pi^2 a}{4}}{t^2}$$

Summary

The solution(s) found are the following

$$y = -\frac{\cos(t) - \frac{\pi^2 a}{4}}{t^2} \quad (1)$$

Verification of solutions

$$y = -\frac{\cos(t) - \frac{\pi^2 a}{4}}{t^2}$$

Verified OK.

1.25.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (t^2) dy &= (-2ty + \sin(t)) dt \\ (2ty - \sin(t)) dt + (t^2) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= 2ty - \sin(t) \\ N(t, y) &= t^2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (2ty - \sin(t)) \\ &= 2t \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (t^2) \\ &= 2t \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int 2ty - \sin(t) dt \\ \phi &= t^2y + \cos(t) + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t^2 + f'(y)\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t^2$. Therefore equation (4) becomes

$$t^2 = t^2 + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = t^2y + \cos(t) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = t^2y + \cos(t)$$

The solution becomes

$$y = -\frac{\cos(t) - c_1}{t^2}$$

Initial conditions are used to solve for c_1 . Substituting $t = -\frac{\pi}{2}$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = \frac{4c_1}{\pi^2}$$

$$c_1 = \frac{\pi^2 a}{4}$$

Substituting c_1 found above in the general solution gives

$$y = -\frac{\cos(t) - \frac{\pi^2 a}{4}}{t^2}$$

Summary

The solution(s) found are the following

$$y = -\frac{\cos(t) - \frac{\pi^2 a}{4}}{t^2} \quad (1)$$

Verification of solutions

$$y = -\frac{\cos(t) - \frac{\pi^2 a}{4}}{t^2}$$

Verified OK.

1.25.5 Maple step by step solution

Let's solve

$$\left[2y + ty' = \frac{\sin(t)}{t}, y\left(-\frac{\pi}{2}\right) = a \right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{2y}{t} + \frac{\sin(t)}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{2y}{t} = \frac{\sin(t)}{t^2}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \frac{\mu(t)\sin(t)}{t^2}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{2\mu(t)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = t^2$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)\sin(t)}{t^2} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)\sin(t)}{t^2} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)\sin(t)}{t^2} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = t^2$

$$y = \frac{\int \sin(t) dt + c_1}{t^2}$$

- Evaluate the integrals on the rhs

$$y = \frac{-\cos(t) + c_1}{t^2}$$

- Use initial condition $y\left(-\frac{\pi}{2}\right) = a$

$$a = \frac{4c_1}{\pi^2}$$

- Solve for c_1

$$c_1 = \frac{\pi^2 a}{4}$$

- Substitute $c_1 = \frac{\pi^2 a}{4}$ into general solution and simplify

$$y = \frac{-\cos(t) + \frac{\pi^2 a}{4}}{t^2}$$

- Solution to the IVP

$$y = \frac{-\cos(t) + \frac{\pi^2 a}{4}}{t^2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 19

```
dsolve([2*y(t)+t*diff(y(t),t) = sin(t)/t,y(-1/2*Pi) = a],y(t), singsol=all)
```

$$y(t) = \frac{-\cos(t) + \frac{a\pi^2}{4}}{t^2}$$

✓ Solution by Mathematica

Time used: 0.042 (sec). Leaf size: 22

```
DSolve[{2*y[t]+t*y'[t] == Sin[t]/t,y[-Pi/2]==a},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{\pi^2 a - 4 \cos(t)}{4t^2}$$

1.26 problem 26

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1.26.4 Maple step by step solution	339

Internal problem ID [473]

Internal file name [OUTPUT/473_Sunday_June_05_2022_01_42_02_AM_28951275/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 26.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "linear", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_linear]

$$\cos(t) y + \sin(t) y' = e^t$$

With initial conditions

$$[y(1) = a]$$

1.26.1 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \cot(t) dt} \\ &= \sin(t)\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) (\csc(t) e^t) \\ \frac{d}{dt}(\sin(t) y) &= (\sin(t)) (\csc(t) e^t) \\ d(\sin(t) y) &= e^t dt\end{aligned}$$

Integrating gives

$$\sin(t) y = \int e^t dt$$

$$\sin(t) y = e^t + c_1$$

Dividing both sides by the integrating factor $\mu = \sin(t)$ results in

$$y = \csc(t) e^t + c_1 \csc(t)$$

which simplifies to

$$y = \csc(t) (e^t + c_1)$$

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = \frac{e + c_1}{\sin(1)}$$

$$c_1 = -\frac{\csc(1) e - a}{\csc(1)}$$

Substituting c_1 found above in the general solution gives

$$y = \csc(t) \left(e^t - \frac{\csc(1) e - a}{\csc(1)} \right)$$

Summary

The solution(s) found are the following

$$y = \csc(t) \left(e^t - \frac{\csc(1) e - a}{\csc(1)} \right) \quad (1)$$

Verification of solutions

$$y = \csc(t) \left(e^t - \frac{\csc(1) e - a}{\csc(1)} \right)$$

Verified OK.

1.26.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{\cos(t)y - e^t}{\sin(t)}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 76: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \frac{1}{\sin(t)}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{\sin(t)}} dy\end{aligned}$$

Which results in

$$S = \sin(t) y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{\cos(t) y - e^t}{\sin(t)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 1 \\R_y &= 0 \\S_t &= \cos(t) y \\S_y &= \sin(t)\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^t \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = e^R + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$\sin(t) y = e^t + c_1$$

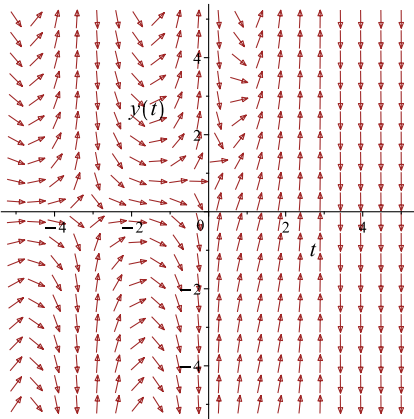
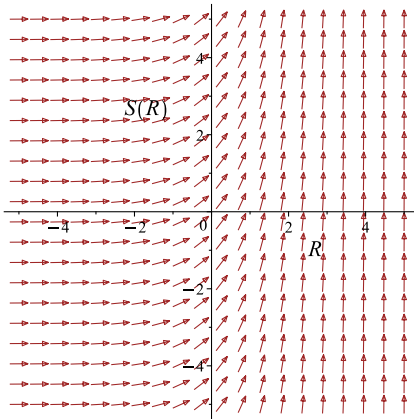
Which simplifies to

$$\sin(t) y = e^t + c_1$$

Which gives

$$y = \frac{e^t + c_1}{\sin(t)}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{\cos(t)y - e^t}{\sin(t)}$ 	$R = t$ $S = \sin(t) y$	$\frac{dS}{dR} = e^R$ 

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = \frac{e + c_1}{\sin(1)}$$

$$c_1 = -\frac{\csc(1)e - a}{\csc(1)}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{e^t - \frac{\csc(1)e - a}{\csc(1)}}{\sin(t)}$$

Summary

The solution(s) found are the following

$$y = \frac{e^t - \frac{\csc(1)e - a}{\csc(1)}}{\sin(t)} \quad (1)$$

Verification of solutions

$$y = \frac{e^t - \frac{\csc(1)e^{-a}}{\csc(1)}}{\sin(t)}$$

Verified OK.

1.26.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned}(\sin(t)) dy &= (-\cos(t)y + e^t) dt \\(\cos(t)y - e^t) dt + (\sin(t)) dy &= 0\end{aligned}\tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= \cos(t)y - e^t \\N(t, y) &= \sin(t)\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(\cos(t)y - e^t) \\&= \cos(t)\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(\sin(t)) \\&= \cos(t)\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M\tag{1}$$

$$\frac{\partial \phi}{\partial y} = N\tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int \cos(t)y - e^t dt \\ \phi &= \sin(t)y - e^t + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = \sin(t) + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \sin(t)$. Therefore equation (4) becomes

$$\sin(t) = \sin(t) + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \sin(t)y - e^t + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \sin(t)y - e^t$$

The solution becomes

$$y = \frac{e^t + c_1}{\sin(t)}$$

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = a$ in the above solution gives an equation to solve for the constant of integration.

$$a = \frac{e + c_1}{\sin(1)}$$

$$c_1 = -\frac{\csc(1)e - a}{\csc(1)}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{e^t - \frac{\csc(1)e^{-a}}{\csc(1)}}{\sin(t)}$$

Summary

The solution(s) found are the following

$$y = \frac{e^t - \frac{\csc(1)e^{-a}}{\csc(1)}}{\sin(t)} \quad (1)$$

Verification of solutions

$$y = \frac{e^t - \frac{\csc(1)e^{-a}}{\csc(1)}}{\sin(t)}$$

Verified OK.

1.26.4 Maple step by step solution

Let's solve

$$[\cos(t)y + \sin(t)y' = e^t, y(1) = a]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{\cos(t)y}{\sin(t)} + \frac{e^t}{\sin(t)}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{\cos(t)y}{\sin(t)} = \frac{e^t}{\sin(t)}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{\cos(t)y}{\sin(t)} \right) = \frac{\mu(t)e^t}{\sin(t)}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{\cos(t)y}{\sin(t)} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)\cos(t)}{\sin(t)}$$

- Solve to find the integrating factor

$$\mu(t) = \sin(t)$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)e^t}{\sin(t)} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)e^t}{\sin(t)} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)e^t}{\sin(t)} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = \sin(t)$

$$y = \frac{\int e^t dt + c_1}{\sin(t)}$$

- Evaluate the integrals on the rhs

$$y = \frac{e^t + c_1}{\sin(t)}$$

- Simplify

$$y = \csc(t) (e^t + c_1)$$

- Use initial condition $y(1) = a$

$$a = \csc(1) (e + c_1)$$

- Solve for c_1

$$c_1 = -\frac{\csc(1)e - a}{\csc(1)}$$

- Substitute $c_1 = -\frac{\csc(1)e - a}{\csc(1)}$ into general solution and simplify

$$y = \csc(t) (a \sin(1) + e^t - e)$$

- Solution to the IVP

$$y = \csc(t) (a \sin(1) + e^t - e)$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 18

```
dsolve([cos(t)*y(t)+sin(t)*diff(y(t),t) = exp(t),y(1) = a],y(t), singsol=all)
```

$$y(t) = \csc(t) (e^t + a \sin(1) - e)$$

✓ Solution by Mathematica

Time used: 0.062 (sec). Leaf size: 19

```
DSolve[{Cos[t]*y[t]+Sin[t]*y'[t] == Exp[t],y[1]==a},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \csc(t) (a \sin(1) + e^t - e)$$

1.27 problem 27

1.27.1 Existence and uniqueness analysis	342
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Internal problem ID [474]

Internal file name [OUTPUT/474_Sunday_June_05_2022_01_42_03_AM_40812399/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 27.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$y' + \frac{y}{2} = 2 \cos(t)$$

With initial conditions

$$[y(0) = -1]$$

1.27.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{1}{2}$$

$$q(t) = 2 \cos(t)$$

Hence the ode is

$$y' + \frac{y}{2} = 2 \cos(t)$$

The domain of $p(t) = \frac{1}{2}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = 2 \cos(t)$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

1.27.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{1}{2} dt} \\ &= e^{\frac{t}{2}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu)(2 \cos(t)) \\ \frac{d}{dt}\left(e^{\frac{t}{2}} y\right) &= \left(e^{\frac{t}{2}}\right)(2 \cos(t)) \\ d\left(e^{\frac{t}{2}} y\right) &= \left(2 \cos(t) e^{\frac{t}{2}}\right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{\frac{t}{2}} y &= \int 2 \cos(t) e^{\frac{t}{2}} dt \\ e^{\frac{t}{2}} y &= \frac{4 \cos(t) e^{\frac{t}{2}}}{5} + \frac{8 \sin(t) e^{\frac{t}{2}}}{5} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{\frac{t}{2}}$ results in

$$y = e^{-\frac{t}{2}} \left(\frac{4 \cos(t) e^{\frac{t}{2}}}{5} + \frac{8 \sin(t) e^{\frac{t}{2}}}{5} \right) + c_1 e^{-\frac{t}{2}}$$

which simplifies to

$$y = \frac{8 \sin(t)}{5} + \frac{4 \cos(t)}{5} + c_1 e^{-\frac{t}{2}}$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = -1$ in the above solution gives an equation to solve for the constant of integration.

$$-1 = \frac{4}{5} + c_1$$

$$c_1 = -\frac{9}{5}$$

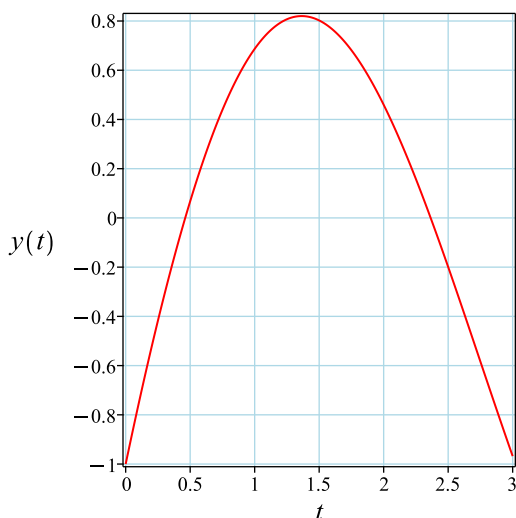
Substituting c_1 found above in the general solution gives

$$y = \frac{8 \sin(t)}{5} + \frac{4 \cos(t)}{5} - \frac{9 e^{-\frac{t}{2}}}{5}$$

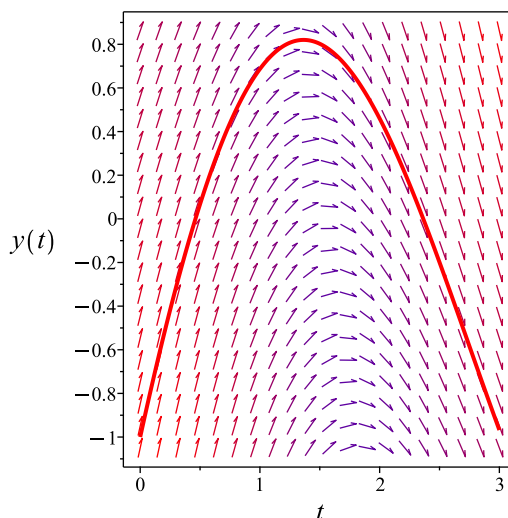
Summary

The solution(s) found are the following

$$y = \frac{8 \sin(t)}{5} + \frac{4 \cos(t)}{5} - \frac{9 e^{-\frac{t}{2}}}{5} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{8 \sin(t)}{5} + \frac{4 \cos(t)}{5} - \frac{9 e^{-\frac{t}{2}}}{5}$$

Verified OK.

1.27.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{y}{2} + 2 \cos(t)$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 79: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-\frac{t}{2}}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-\frac{t}{2}}} dy\end{aligned}$$

Which results in

$$S = e^{\frac{t}{2}} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{y}{2} + 2 \cos(t)$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 1 \\ R_y &= 0 \\ S_t &= \frac{e^{\frac{t}{2}} y}{2} \\ S_y &= e^{\frac{t}{2}}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 2 \cos (t) e^{\frac{t}{2}} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 2 \cos (R) e^{\frac{R}{2}}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 + \frac{4 e^{\frac{R}{2}} (\cos (R) + 2 \sin (R))}{5} \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{\frac{t}{2}} y = \frac{4 e^{\frac{t}{2}} (\cos (t) + 2 \sin (t))}{5} + c_1$$

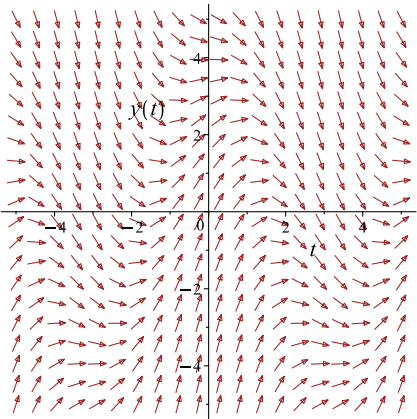
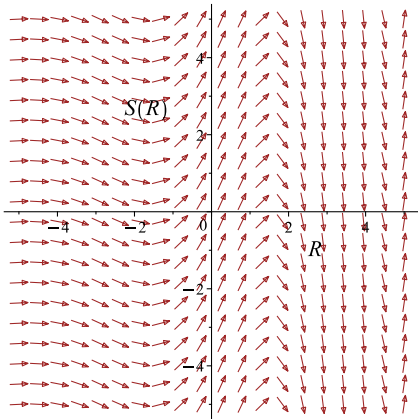
Which simplifies to

$$e^{\frac{t}{2}} y = \frac{4 e^{\frac{t}{2}} (\cos (t) + 2 \sin (t))}{5} + c_1$$

Which gives

$$y = \frac{e^{-\frac{t}{2}} \left(4 \cos (t) e^{\frac{t}{2}} + 8 \sin (t) e^{\frac{t}{2}} + 5c_1 \right)}{5}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{y}{2} + 2 \cos(t)$ 	$R = t$ $S = e^{\frac{t}{2}} y$	$\frac{dS}{dR} = 2 \cos(R) e^{\frac{R}{2}}$ 

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = -1$ in the above solution gives an equation to solve for the constant of integration.

$$-1 = \frac{4}{5} + c_1$$

$$c_1 = -\frac{9}{5}$$

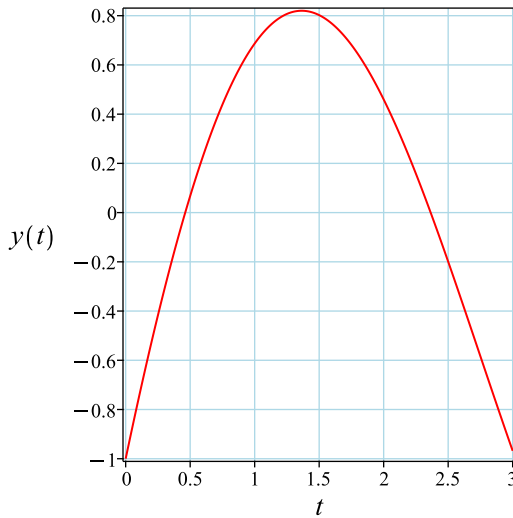
Substituting c_1 found above in the general solution gives

$$y = \frac{8 \sin(t)}{5} + \frac{4 \cos(t)}{5} - \frac{9 e^{-\frac{t}{2}}}{5}$$

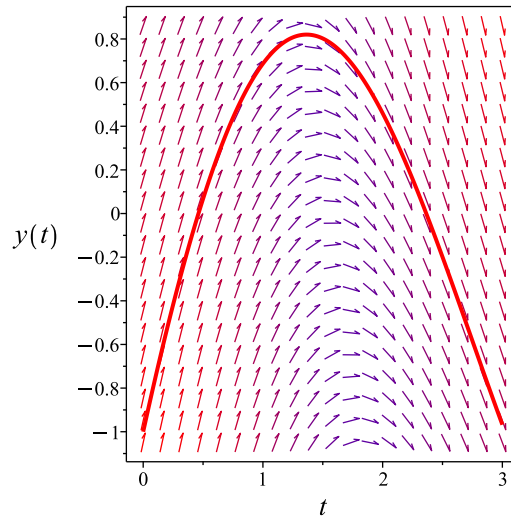
Summary

The solution(s) found are the following

$$y = \frac{8 \sin(t)}{5} + \frac{4 \cos(t)}{5} - \frac{9 e^{-\frac{t}{2}}}{5} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{8 \sin(t)}{5} + \frac{4 \cos(t)}{5} - \frac{9 e^{-\frac{t}{2}}}{5}$$

Verified OK.

1.27.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} dy &= \left(-\frac{y}{2} + 2 \cos(t) \right) dt \\ \left(\frac{y}{2} - 2 \cos(t) \right) dt + dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= \frac{y}{2} - 2 \cos(t) \\ N(t, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{y}{2} - 2 \cos(t) \right) \\ &= \frac{1}{2} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (1) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1 \left(\left(\frac{1}{2} \right) - (0) \right) \\ &= \frac{1}{2} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int \frac{1}{2} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\frac{t}{2}} \\ &= e^{\frac{t}{2}} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{\frac{t}{2}} \left(\frac{y}{2} - 2 \cos(t) \right) \\ &= \frac{(y - 4 \cos(t)) e^{\frac{t}{2}}}{2} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{\frac{t}{2}} (1) \\ &= e^{\frac{t}{2}} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left(\frac{(y - 4 \cos(t)) e^{\frac{t}{2}}}{2} \right) + \left(e^{\frac{t}{2}} \right) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int \overline{M} dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int \frac{(y - 4 \cos(t)) e^{\frac{t}{2}}}{2} dt$$

$$\phi = -\frac{e^{\frac{t}{2}}(-5y + 4 \cos(t) + 8 \sin(t))}{5} + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{\frac{t}{2}} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{\frac{t}{2}}$. Therefore equation (4) becomes

$$e^{\frac{t}{2}} = e^{\frac{t}{2}} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{e^{\frac{t}{2}}(-5y + 4 \cos(t) + 8 \sin(t))}{5} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{e^{\frac{t}{2}}(-5y + 4 \cos(t) + 8 \sin(t))}{5}$$

The solution becomes

$$y = \frac{e^{-\frac{t}{2}} \left(4 \cos(t) e^{\frac{t}{2}} + 8 \sin(t) e^{\frac{t}{2}} + 5c_1 \right)}{5}$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = -1$ in the above solution gives an equation to solve for the constant of integration.

$$-1 = \frac{4}{5} + c_1$$

$$c_1 = -\frac{9}{5}$$

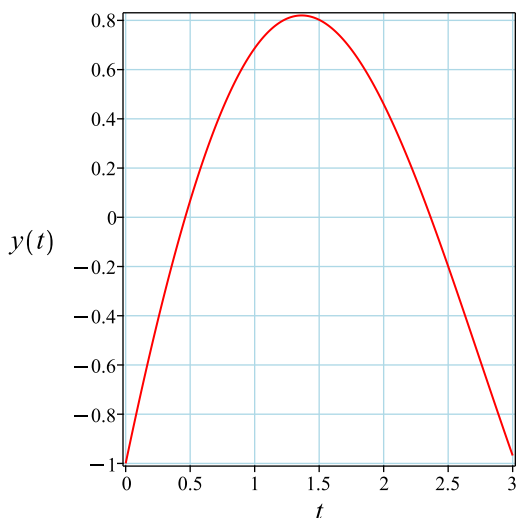
Substituting c_1 found above in the general solution gives

$$y = \frac{8 \sin(t)}{5} + \frac{4 \cos(t)}{5} - \frac{9 e^{-\frac{t}{2}}}{5}$$

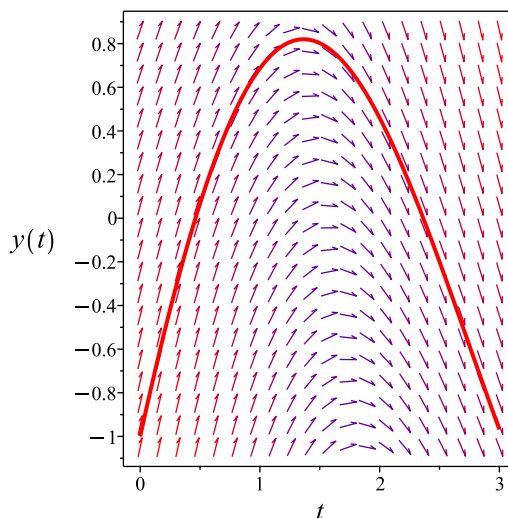
Summary

The solution(s) found are the following

$$y = \frac{8 \sin(t)}{5} + \frac{4 \cos(t)}{5} - \frac{9 e^{-\frac{t}{2}}}{5} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{8 \sin(t)}{5} + \frac{4 \cos(t)}{5} - \frac{9 e^{-\frac{t}{2}}}{5}$$

Verified OK.

1.27.5 Maple step by step solution

Let's solve

$$\left[y' + \frac{y}{2} = 2 \cos(t), y(0) = -1 \right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{y}{2} + 2 \cos(t)$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{y}{2} = 2 \cos(t)$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{y}{2} \right) = 2\mu(t) \cos(t)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{y}{2} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)}{2}$$

- Solve to find the integrating factor

$$\mu(t) = e^{\frac{t}{2}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int 2\mu(t) \cos(t) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int 2\mu(t) \cos(t) dt + c_1$$

- Solve for y

$$y = \frac{\int 2\mu(t) \cos(t) dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{\frac{t}{2}}$

$$y = \frac{\int 2 \cos(t) e^{\frac{t}{2}} dt + c_1}{e^{\frac{t}{2}}}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{4 \cos(t) e^{\frac{t}{2}}}{5} + \frac{8 \sin(t) e^{\frac{t}{2}}}{5} + c_1}{e^{\frac{t}{2}}}$$

- Simplify

$$y = \frac{8\sin(t)}{5} + \frac{4\cos(t)}{5} + c_1 e^{-\frac{t}{2}}$$

- Use initial condition $y(0) = -1$

$$-1 = \frac{4}{5} + c_1$$

- Solve for c_1

$$c_1 = -\frac{9}{5}$$

- Substitute $c_1 = -\frac{9}{5}$ into general solution and simplify

$$y = \frac{8\sin(t)}{5} + \frac{4\cos(t)}{5} - \frac{9e^{-\frac{t}{2}}}{5}$$

- Solution to the IVP

$$y = \frac{8\sin(t)}{5} + \frac{4\cos(t)}{5} - \frac{9e^{-\frac{t}{2}}}{5}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve([1/2*y(t)+diff(y(t),t) = 2*cos(t),y(0) = -1],y(t), singsol=all)
```

$$y(t) = \frac{4\cos(t)}{5} + \frac{8\sin(t)}{5} - \frac{9e^{-\frac{t}{2}}}{5}$$

✓ Solution by Mathematica

Time used: 0.049 (sec). Leaf size: 27

```
DSolve[{1/2*y[t]+y'[t] == 2*Cos[t],y[0]==-1},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{5}(-9e^{-t/2} + 8\sin(t) + 4\cos(t))$$

1.28 problem 28

1.28.1 Solving as linear ode	356
1.28.2 Solving as first order ode lie symmetry lookup ode	358
1.28.3 Solving as exact ode	362
1.28.4 Maple step by step solution	367

Internal problem ID [475]

Internal file name [OUTPUT/475_Sunday_June_05_2022_01_42_04_AM_90970000/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 28.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$\frac{2y}{3} + y' = -\frac{t}{2} + 1$$

1.28.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{2}{3}$$

$$q(t) = -\frac{t}{2} + 1$$

Hence the ode is

$$\frac{2y}{3} + y' = -\frac{t}{2} + 1$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{2}{3} dt} \\ &= e^{\frac{2t}{3}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(-\frac{t}{2} + 1 \right) \\ \frac{d}{dt} \left(e^{\frac{2t}{3}} y \right) &= \left(e^{\frac{2t}{3}} \right) \left(-\frac{t}{2} + 1 \right) \\ d \left(e^{\frac{2t}{3}} y \right) &= \left(-\frac{(t-2)e^{\frac{2t}{3}}}{2} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{\frac{2t}{3}} y &= \int -\frac{(t-2)e^{\frac{2t}{3}}}{2} dt \\ e^{\frac{2t}{3}} y &= -\frac{3(2t-7)e^{\frac{2t}{3}}}{8} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{\frac{2t}{3}}$ results in

$$y = -\frac{3e^{-\frac{2t}{3}}(2t-7)e^{\frac{2t}{3}}}{8} + c_1 e^{-\frac{2t}{3}}$$

which simplifies to

$$y = -\frac{3t}{4} + \frac{21}{8} + c_1 e^{-\frac{2t}{3}}$$

Summary

The solution(s) found are the following

$$y = -\frac{3t}{4} + \frac{21}{8} + c_1 e^{-\frac{2t}{3}} \quad (1)$$

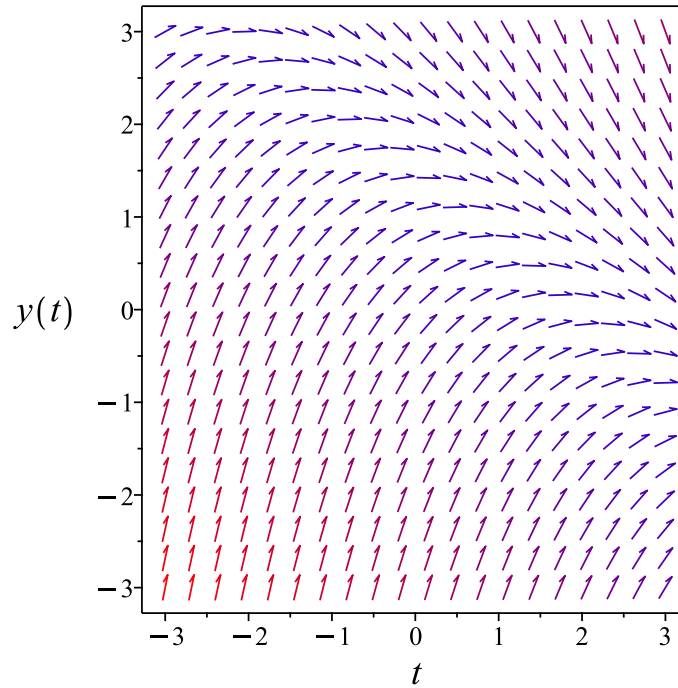


Figure 65: Slope field plot

Verification of solutions

$$y = -\frac{3t}{4} + \frac{21}{8} + c_1 e^{-\frac{2t}{3}}$$

Verified OK.

1.28.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{2y}{3} - \frac{t}{2} + 1$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 82: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-\frac{2t}{3}}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-\frac{2t}{3}}} dy \end{aligned}$$

Which results in

$$S = e^{\frac{2t}{3}} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{2y}{3} - \frac{t}{2} + 1$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= \frac{2e^{\frac{2t}{3}}y}{3} \\ S_y &= e^{\frac{2t}{3}} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \left(-\frac{t}{2} + 1 \right) e^{\frac{2t}{3}} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \left(-\frac{R}{2} + 1 \right) e^{\frac{2R}{3}}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{3(2R - 7)e^{\frac{2R}{3}}}{8} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{\frac{2t}{3}}y = -\frac{3(2t - 7)e^{\frac{2t}{3}}}{8} + c_1$$

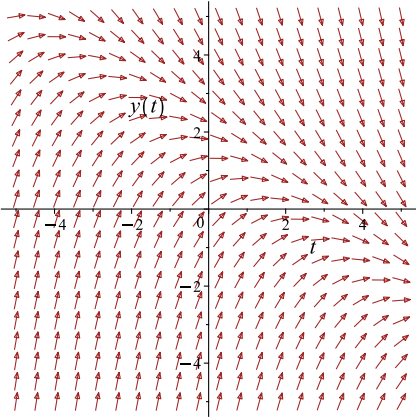
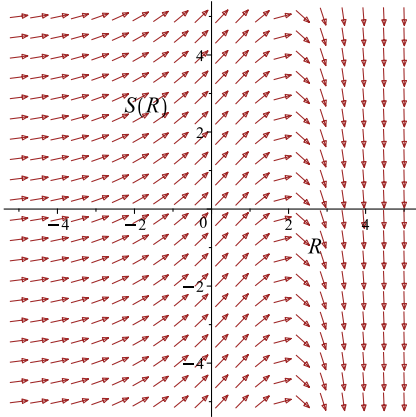
Which simplifies to

$$e^{\frac{2t}{3}}y = -\frac{3(2t - 7)e^{\frac{2t}{3}}}{8} + c_1$$

Which gives

$$y = -\frac{(6e^{\frac{2t}{3}}t - 21e^{\frac{2t}{3}} - 8c_1)e^{-\frac{2t}{3}}}{8}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{2y}{3} - \frac{t}{2} + 1$ 	$R = t$ $S = e^{\frac{2t}{3}}y$	$\frac{dS}{dR} = \left(-\frac{R}{2} + 1\right)e^{\frac{2R}{3}}$ 

Summary

The solution(s) found are the following

$$y = -\frac{\left(6 e^{\frac{2t}{3}} t - 21 e^{\frac{2t}{3}} - 8c_1\right) e^{-\frac{2t}{3}}}{8} \quad (1)$$

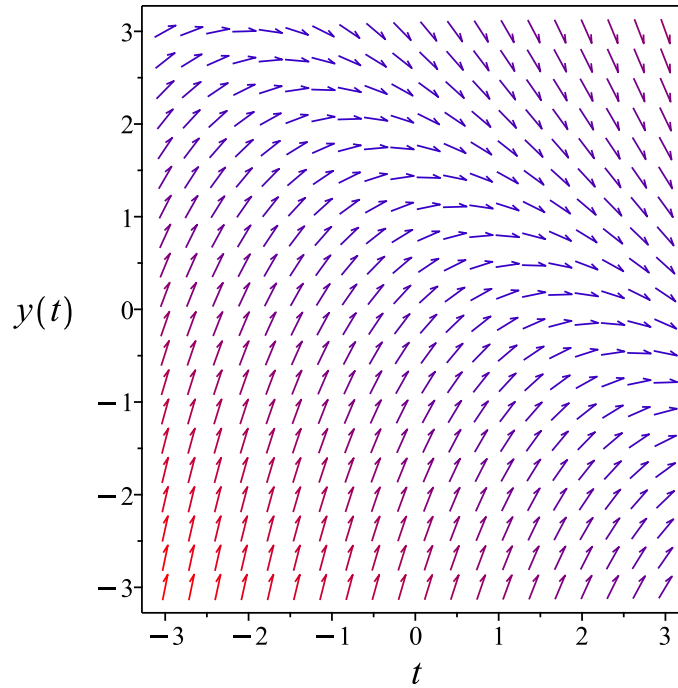


Figure 66: Slope field plot

Verification of solutions

$$y = -\frac{\left(6 e^{\frac{2t}{3}} t - 21 e^{\frac{2t}{3}} - 8c_1\right) e^{-\frac{2t}{3}}}{8}$$

Verified OK.

1.28.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the

ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} dy &= \left(-\frac{2y}{3} - \frac{t}{2} + 1 \right) dt \\ \left(\frac{2y}{3} + \frac{t}{2} - 1 \right) dt + dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= \frac{2y}{3} + \frac{t}{2} - 1 \\ N(t, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{2y}{3} + \frac{t}{2} - 1 \right) \\ &= \frac{2}{3}\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1 \left(\left(\frac{2}{3} \right) - (0) \right) \\ &= \frac{2}{3}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int \frac{2}{3} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{\frac{2t}{3}} \\ &= e^{\frac{2t}{3}}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= e^{\frac{2t}{3}} \left(\frac{2y}{3} + \frac{t}{2} - 1 \right) \\ &= \frac{(4y + 3t - 6) e^{\frac{2t}{3}}}{6}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^{\frac{2t}{3}}(1) \\ &= e^{\frac{2t}{3}}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left(\frac{(4y + 3t - 6) e^{\frac{2t}{3}}}{6} \right) + \left(e^{\frac{2t}{3}} \right) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int \frac{(4y + 3t - 6) e^{\frac{2t}{3}}}{6} dt \\ \phi &= \frac{(6t + 8y - 21) e^{\frac{2t}{3}}}{8} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{\frac{2t}{3}} + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{\frac{2t}{3}}$. Therefore equation (4) becomes

$$e^{\frac{2t}{3}} = e^{\frac{2t}{3}} + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{(6t + 8y - 21)e^{\frac{2t}{3}}}{8} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{(6t + 8y - 21)e^{\frac{2t}{3}}}{8}$$

The solution becomes

$$y = -\frac{\left(6e^{\frac{2t}{3}}t - 21e^{\frac{2t}{3}} - 8c_1\right)e^{-\frac{2t}{3}}}{8}$$

Summary

The solution(s) found are the following

$$y = -\frac{\left(6e^{\frac{2t}{3}}t - 21e^{\frac{2t}{3}} - 8c_1\right)e^{-\frac{2t}{3}}}{8} \quad (1)$$

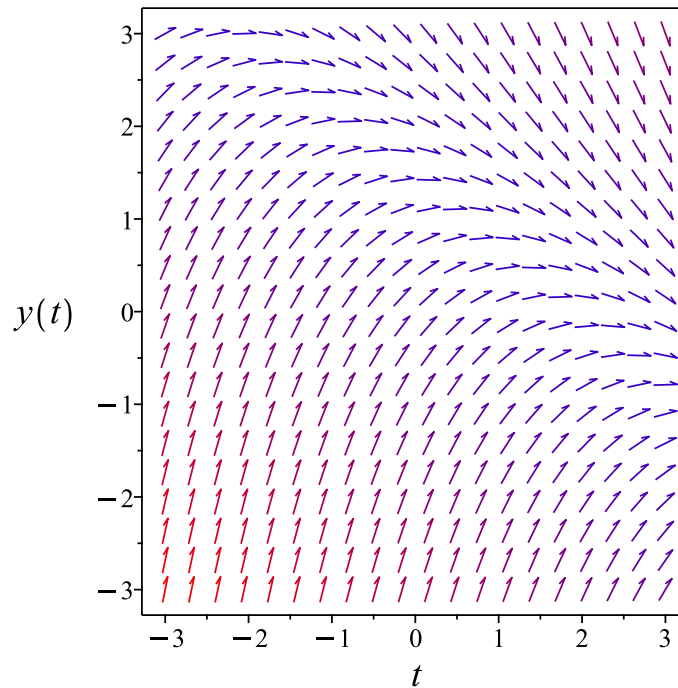


Figure 67: Slope field plot

Verification of solutions

$$y = -\frac{\left(6 e^{\frac{2t}{3}} t - 21 e^{\frac{2t}{3}} - 8c_1\right) e^{-\frac{2t}{3}}}{8}$$

Verified OK.

1.28.4 Maple step by step solution

Let's solve

$$\frac{2y}{3} + y' = -\frac{t}{2} + 1$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{2y}{3} - \frac{t}{2} + 1$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$\frac{2y}{3} + y' = -\frac{t}{2} + 1$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(\frac{2y}{3} + y' \right) = \mu(t) \left(-\frac{t}{2} + 1 \right)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t) y)$

$$\mu(t) \left(\frac{2y}{3} + y' \right) = \mu'(t) y + \mu(t) y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{2\mu(t)}{3}$$

- Solve to find the integrating factor

$$\mu(t) = e^{\frac{2t}{3}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y) \right) dt = \int \mu(t) \left(-\frac{t}{2} + 1 \right) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t) y = \int \mu(t) \left(-\frac{t}{2} + 1 \right) dt + c_1$$

- Solve for y

$$y = \frac{\int \mu(t) \left(-\frac{t}{2} + 1 \right) dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{\frac{2t}{3}}$

$$y = \frac{\int \left(-\frac{t}{2} + 1 \right) e^{\frac{2t}{3}} dt + c_1}{e^{\frac{2t}{3}}}$$

- Evaluate the integrals on the rhs

$$y = \frac{-\frac{3(2t-7)e^{\frac{2t}{3}}}{8} + c_1}{e^{\frac{2t}{3}}}$$

- Simplify

$$y = -\frac{3t}{4} + \frac{21}{8} + c_1 e^{-\frac{2t}{3}}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 15

```
dsolve(2/3*y(t)+diff(y(t),t) = 1-1/2*t,y(t), singsol=all)
```

$$y(t) = -\frac{3t}{4} + \frac{21}{8} + e^{-\frac{2t}{3}} c_1$$

✓ Solution by Mathematica

Time used: 0.061 (sec). Leaf size: 24

```
DSolve[2/3*y[t]+y'[t] == 1-1/2*t,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -\frac{3t}{4} + c_1 e^{-2t/3} + \frac{21}{8}$$

1.29 problem 29

1.29.1 Existence and uniqueness analysis	370
1.29.2 Solving as linear ode	371
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1.29.4 Solving as exact ode	377
1.29.5 Maple step by step solution	382

Internal problem ID [476]

Internal file name [OUTPUT/476_Sunday_June_05_2022_01_42_05_AM_51193873/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 29.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$\frac{y}{4} + y' = 3 + 2 \cos(2t)$$

With initial conditions

$$[y(0) = 0]$$

1.29.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{1}{4}$$
$$q(t) = 3 + 2 \cos(2t)$$

Hence the ode is

$$\frac{y}{4} + y' = 3 + 2 \cos(2t)$$

The domain of $p(t) = \frac{1}{4}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = 3 + 2 \cos(2t)$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

1.29.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{1}{4} dt} \\ &= e^{\frac{t}{4}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu)(3 + 2 \cos(2t)) \\ \frac{d}{dt}\left(e^{\frac{t}{4}} y\right) &= \left(e^{\frac{t}{4}}\right)(3 + 2 \cos(2t)) \\ d\left(e^{\frac{t}{4}} y\right) &= \left((3 + 2 \cos(2t)) e^{\frac{t}{4}}\right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{\frac{t}{4}} y &= \int (3 + 2 \cos(2t)) e^{\frac{t}{4}} dt \\ e^{\frac{t}{4}} y &= 12 e^{\frac{t}{4}} + \frac{8 \cos(2t) e^{\frac{t}{4}}}{65} + \frac{64 \sin(2t) e^{\frac{t}{4}}}{65} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{\frac{t}{4}}$ results in

$$y = e^{-\frac{t}{4}} \left(12 e^{\frac{t}{4}} + \frac{8 \cos(2t) e^{\frac{t}{4}}}{65} + \frac{64 \sin(2t) e^{\frac{t}{4}}}{65} \right) + c_1 e^{-\frac{t}{4}}$$

which simplifies to

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 + c_1 e^{-\frac{t}{4}}$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{788}{65} + c_1$$

$$c_1 = -\frac{788}{65}$$

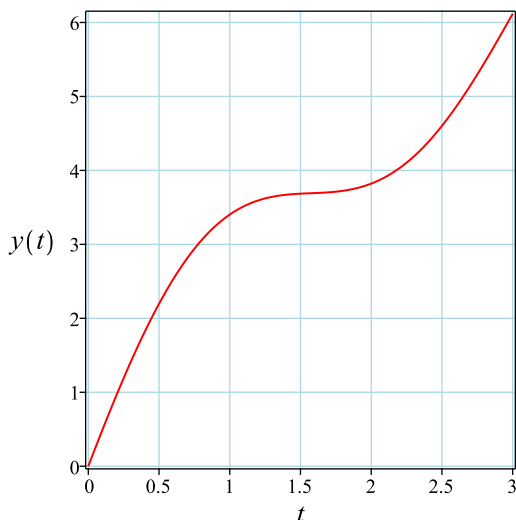
Substituting c_1 found above in the general solution gives

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65}$$

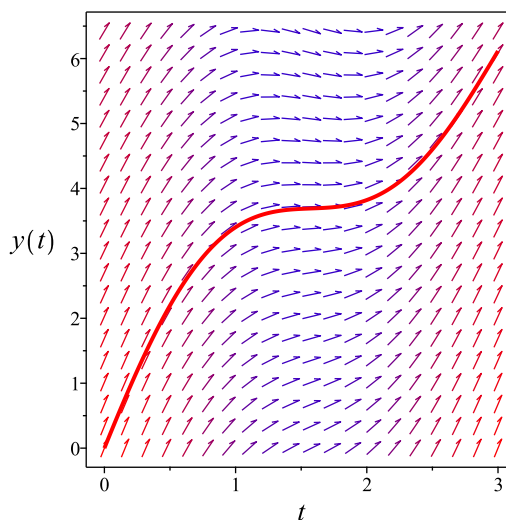
Summary

The solution(s) found are the following

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65}$$

Verified OK.

1.29.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{y}{4} + 3 + 2 \cos(2t)$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 85: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-\frac{t}{4}}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-\frac{t}{4}}} dy\end{aligned}$$

Which results in

$$S = e^{\frac{t}{4}} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{y}{4} + 3 + 2 \cos(2t)$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 1 \\ R_y &= 0 \\ S_t &= \frac{e^{\frac{t}{4}} y}{4} \\ S_y &= e^{\frac{t}{4}}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = (3 + 2 \cos(2t)) e^{\frac{t}{4}} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = (3 + 2 \cos(2R)) e^{\frac{R}{4}}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = 12 e^{\frac{R}{4}} + c_1 + \frac{8 e^{\frac{R}{4}} (\cos(2R) + 8 \sin(2R))}{65} \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{\frac{t}{4}} y = 12 e^{\frac{t}{4}} + c_1 + \frac{8 e^{\frac{t}{4}} (\cos(2t) + 8 \sin(2t))}{65}$$

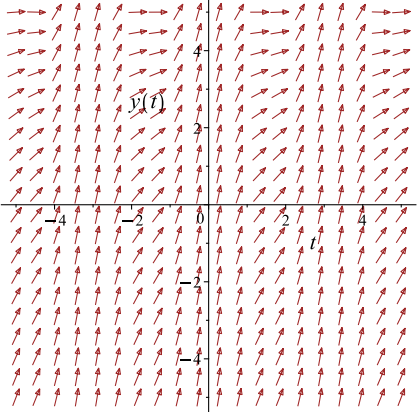
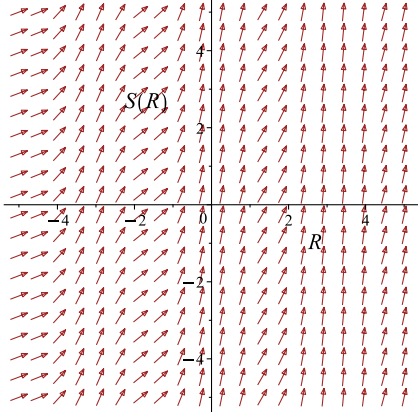
Which simplifies to

$$\frac{(65y - 8 \cos(2t) - 64 \sin(2t) - 780) e^{\frac{t}{4}}}{65} - c_1 = 0$$

Which gives

$$y = \frac{e^{-\frac{t}{4}} (8 \cos(2t) e^{\frac{t}{4}} + 64 \sin(2t) e^{\frac{t}{4}} + 780 e^{\frac{t}{4}} + 65c_1)}{65}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{y}{4} + 3 + 2 \cos(2t)$ 	$R = t$ $S = e^{\frac{t}{4}} y$	$\frac{dS}{dR} = (3 + 2 \cos(2R)) e^{\frac{R}{4}}$ 

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{788}{65} + c_1$$

$$c_1 = -\frac{788}{65}$$

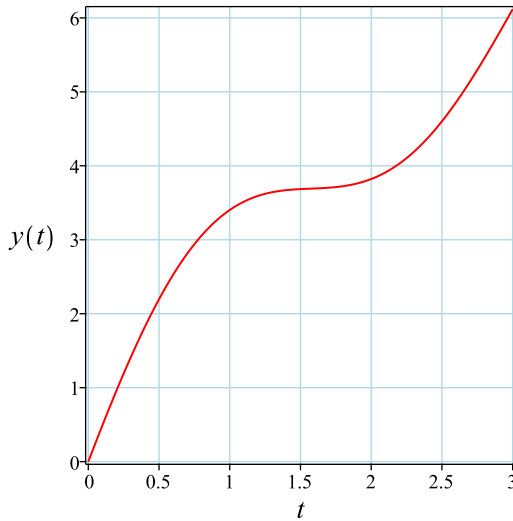
Substituting c_1 found above in the general solution gives

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65}$$

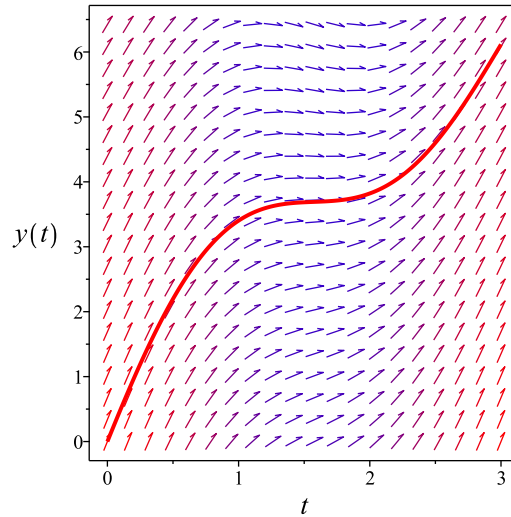
Summary

The solution(s) found are the following

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65}$$

Verified OK.

1.29.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} dy &= \left(-\frac{y}{4} + 3 + 2 \cos(2t) \right) dt \\ \left(\frac{y}{4} - 3 - 2 \cos(2t) \right) dt + dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= \frac{y}{4} - 3 - 2 \cos(2t) \\ N(t, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{y}{4} - 3 - 2 \cos(2t) \right) \\ &= \frac{1}{4} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (1) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1 \left(\left(\frac{1}{4} \right) - (0) \right) \\ &= \frac{1}{4} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int \frac{1}{4} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\frac{t}{4}} \\ &= e^{\frac{t}{4}} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{\frac{t}{4}} \left(\frac{y}{4} - 3 - 2 \cos(2t) \right) \\ &= \frac{(y - 12 - 8 \cos(2t)) e^{\frac{t}{4}}}{4} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{\frac{t}{4}} (1) \\ &= e^{\frac{t}{4}} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left(\frac{(y - 12 - 8 \cos(2t)) e^{\frac{t}{4}}}{4} \right) + \left(e^{\frac{t}{4}} \right) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int \overline{M} dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int \frac{(y - 12 - 8 \cos(2t)) e^{\frac{t}{4}}}{4} dt$$

$$\phi = -\frac{e^{\frac{t}{4}}(780 - 65y + 8 \cos(2t) + 64 \sin(2t))}{65} + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{\frac{t}{4}} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{\frac{t}{4}}$. Therefore equation (4) becomes

$$e^{\frac{t}{4}} = e^{\frac{t}{4}} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{e^{\frac{t}{4}}(780 - 65y + 8 \cos(2t) + 64 \sin(2t))}{65} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{e^{\frac{t}{4}}(780 - 65y + 8 \cos(2t) + 64 \sin(2t))}{65}$$

The solution becomes

$$y = \frac{e^{-\frac{t}{4}} \left(8 \cos(2t) e^{\frac{t}{4}} + 64 \sin(2t) e^{\frac{t}{4}} + 780 e^{\frac{t}{4}} + 65c_1 \right)}{65}$$

Initial conditions are used to solve for c_1 . Substituting $t = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{788}{65} + c_1$$

$$c_1 = -\frac{788}{65}$$

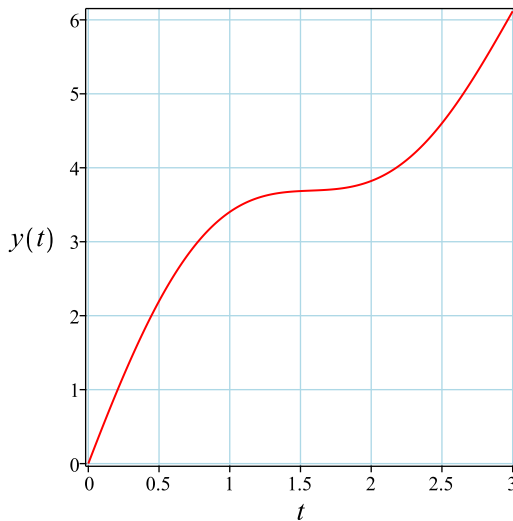
Substituting c_1 found above in the general solution gives

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65}$$

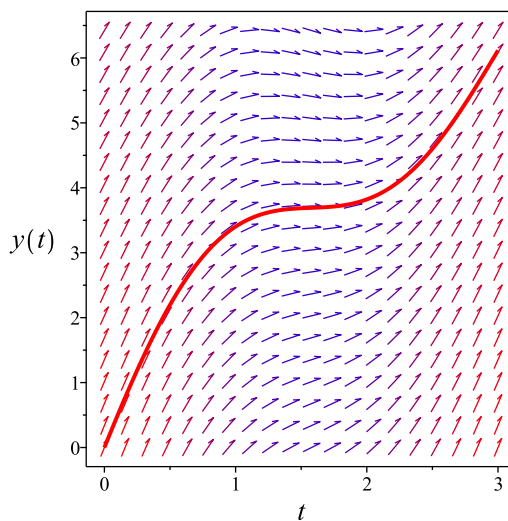
Summary

The solution(s) found are the following

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65}$$

Verified OK.

1.29.5 Maple step by step solution

Let's solve

$$\left[\frac{y}{4} + y' = 3 + 2 \cos(2t), y(0) = 0 \right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{y}{4} + 3 + 2 \cos(2t)$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$\frac{y}{4} + y' = 3 + 2 \cos(2t)$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(\frac{y}{4} + y' \right) = \mu(t) (3 + 2 \cos(2t))$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t) y)$

$$\mu(t) \left(\frac{y}{4} + y' \right) = \mu'(t) y + \mu(t) y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)}{4}$$

- Solve to find the integrating factor

$$\mu(t) = e^{\frac{t}{4}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y) \right) dt = \int \mu(t) (3 + 2 \cos(2t)) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t) y = \int \mu(t) (3 + 2 \cos(2t)) dt + c_1$$

- Solve for y

$$y = \frac{\int \mu(t)(3+2 \cos(2t)) dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{\frac{t}{4}}$

$$y = \frac{\int (3+2 \cos(2t)) e^{\frac{t}{4}} dt + c_1}{e^{\frac{t}{4}}}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{8 \cos(2t) e^{\frac{t}{4}}}{65} + \frac{64 \sin(2t) e^{\frac{t}{4}}}{65} + 12 e^{\frac{t}{4}} + c_1}{e^{\frac{t}{4}}}$$

- Simplify

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 + c_1 e^{-\frac{t}{4}}$$

- Use initial condition $y(0) = 0$

$$0 = \frac{788}{65} + c_1$$

- Solve for c_1

$$c_1 = -\frac{788}{65}$$

- Substitute $c_1 = -\frac{788}{65}$ into general solution and simplify

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65}$$

- Solution to the IVP

$$y = \frac{64 \sin(2t)}{65} + \frac{8 \cos(2t)}{65} + 12 - \frac{788 e^{-\frac{t}{4}}}{65}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 24

```
dsolve([1/4*y(t)+diff(y(t),t) = 3+2*cos(2*t),y(0) = 0],y(t), singsol=all)
```

$$y(t) = 12 + \frac{8 \cos(2t)}{65} + \frac{64 \sin(2t)}{65} - \frac{788 e^{-\frac{t}{4}}}{65}$$

✓ Solution by Mathematica

Time used: 0.158 (sec). Leaf size: 32

```
DSolve[{1/4*y[t]+y'[t] == 3+2*Cos[2*t],y[0]==0},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{4}{65}(-197e^{-t/4} + 16 \sin(2t) + 2 \cos(2t) + 195)$$

1.30 problem 30

1.30.1 Solving as linear ode	384
1.30.2 Solving as first order ode lie symmetry lookup ode	386
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1.30.4 Maple step by step solution	394

Internal problem ID [477]

Internal file name [OUTPUT/477_Sunday_June_05_2022_01_42_06_AM_27968450/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 30.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$-y + y' = 1 + 3 \sin(t)$$

1.30.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -1$$

$$q(t) = 1 + 3 \sin(t)$$

Hence the ode is

$$-y + y' = 1 + 3 \sin(t)$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int(-1)dt} \\ &= e^{-t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu)(1 + 3 \sin(t)) \\ \frac{d}{dt}(e^{-t}y) &= (e^{-t})(1 + 3 \sin(t)) \\ d(e^{-t}y) &= ((1 + 3 \sin(t))e^{-t}) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-t}y &= \int (1 + 3 \sin(t)) e^{-t} dt \\ e^{-t}y &= -e^{-t} - \frac{3e^{-t} \cos(t)}{2} - \frac{3e^{-t} \sin(t)}{2} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-t}$ results in

$$y = e^t \left(-e^{-t} - \frac{3e^{-t} \cos(t)}{2} - \frac{3e^{-t} \sin(t)}{2} \right) + c_1 e^t$$

which simplifies to

$$y = -1 + c_1 e^t - \frac{3 \cos(t)}{2} - \frac{3 \sin(t)}{2}$$

Summary

The solution(s) found are the following

$$y = -1 + c_1 e^t - \frac{3 \cos(t)}{2} - \frac{3 \sin(t)}{2} \tag{1}$$

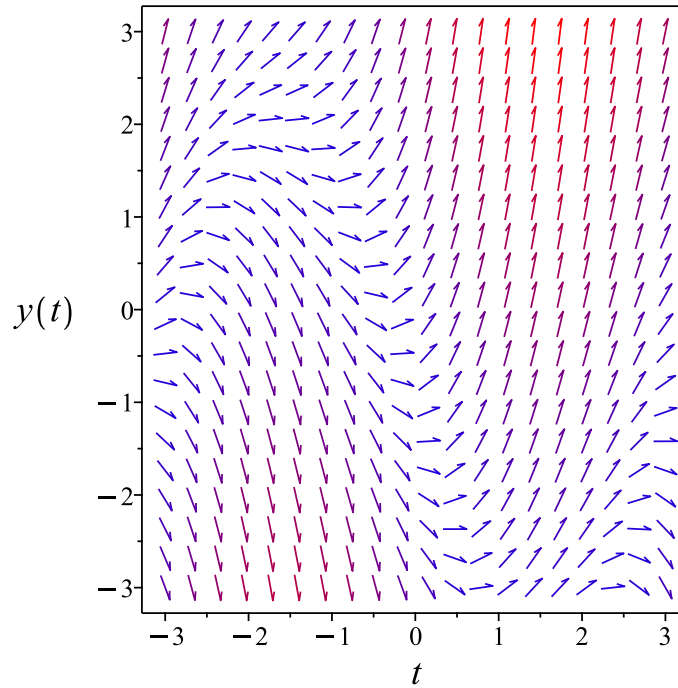


Figure 71: Slope field plot

Verification of solutions

$$y = -1 + c_1 e^t - \frac{3 \cos(t)}{2} - \frac{3 \sin(t)}{2}$$

Verified OK.

1.30.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$\begin{aligned} y' &= y + 1 + 3 \sin(t) \\ y' &= \omega(t, y) \end{aligned}$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 88: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^t\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^t} dy \end{aligned}$$

Which results in

$$S = e^{-t}y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = y + 1 + 3 \sin(t)$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -e^{-t}y \\ S_y &= e^{-t} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = (1 + 3 \sin(t)) e^{-t} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = (1 + 3 \sin(R)) e^{-R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -e^{-R} + c_1 - \frac{3e^{-R}(\cos(R) + \sin(R))}{2} \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{-t}y = -e^{-t} + c_1 - \frac{3e^{-t}(\cos(t) + \sin(t))}{2}$$

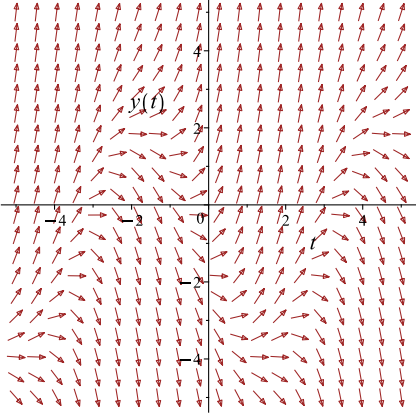
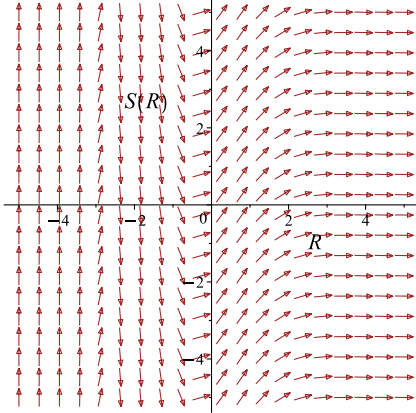
Which simplifies to

$$\frac{(2y + 3\cos(t) + 3\sin(t) + 2)e^{-t}}{2} - c_1 = 0$$

Which gives

$$y = -\frac{(3e^{-t}\sin(t) + 3e^{-t}\cos(t) + 2e^{-t} - 2c_1)e^t}{2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = y + 1 + 3\sin(t)$ 	$R = t$ $S = e^{-t}y$	$\frac{dS}{dR} = (1 + 3\sin(R))e^{-R}$ 

Summary

The solution(s) found are the following

$$y = -\frac{(3e^{-t}\sin(t) + 3e^{-t}\cos(t) + 2e^{-t} - 2c_1)e^t}{2} \quad (1)$$

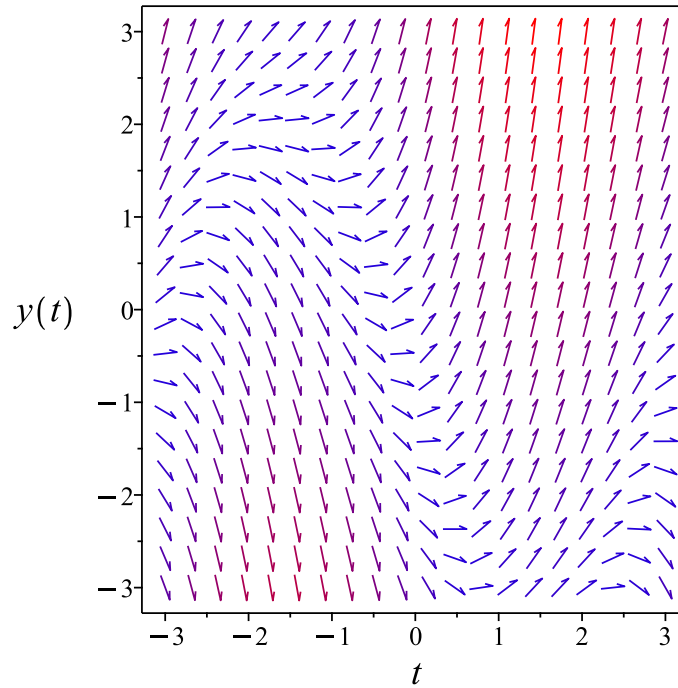


Figure 72: Slope field plot

Verification of solutions

$$y = -\frac{(3 e^{-t} \sin (t) + 3 e^{-t} \cos (t) + 2 e^{-t} - 2c_1) e^t}{2}$$

Verified OK.

1.30.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dy &= (y + 1 + 3 \sin(t)) dt \\ (-y - 1 - 3 \sin(t)) dt + dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -y - 1 - 3 \sin(t) \\ N(t, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-y - 1 - 3 \sin(t)) \\ &= -1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1((-1) - (0)) \\ &= -1 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int -1 dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-t} \\ &= e^{-t} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{-t}(-y - 1 - 3 \sin(t)) \\ &= -e^{-t}(y + 1 + 3 \sin(t)) \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{-t}(1) \\ &= e^{-t} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ (-e^{-t}(y + 1 + 3 \sin(t))) + (e^{-t}) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -e^{-t}(y + 1 + 3 \sin(t)) dt \\ \phi &= \frac{(2y + 3 \cos(t) + 3 \sin(t) + 2) e^{-t}}{2} + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-t} + f'(y)\quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-t}$. Therefore equation (4) becomes

$$e^{-t} = e^{-t} + f'(y)\quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{(2y + 3 \cos(t) + 3 \sin(t) + 2) e^{-t}}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{(2y + 3 \cos(t) + 3 \sin(t) + 2) e^{-t}}{2}$$

The solution becomes

$$y = -\frac{(3 e^{-t} \sin(t) + 3 e^{-t} \cos(t) + 2 e^{-t} - 2c_1) e^t}{2}$$

Summary

The solution(s) found are the following

$$y = -\frac{(3 e^{-t} \sin (t) + 3 e^{-t} \cos (t) + 2 e^{-t} - 2 c_1) e^t}{2} \quad (1)$$

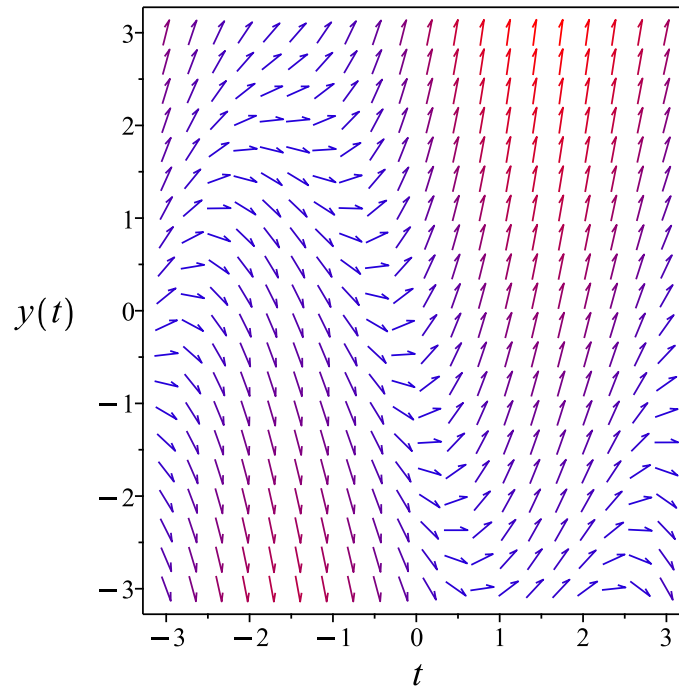


Figure 73: Slope field plot

Verification of solutions

$$y = -\frac{(3 e^{-t} \sin (t) + 3 e^{-t} \cos (t) + 2 e^{-t} - 2 c_1) e^t}{2}$$

Verified OK.

1.30.4 Maple step by step solution

Let's solve

$$-y + y' = 1 + 3 \sin (t)$$

- Highest derivative means the order of the ODE is 1

y'

- Isolate the derivative

$$y' = y + 1 + 3 \sin(t)$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$-y + y' = 1 + 3 \sin(t)$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t)(-y + y') = \mu(t)(1 + 3 \sin(t))$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t)(-y + y') = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -\mu(t)$$

- Solve to find the integrating factor

$$\mu(t) = e^{-t}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \mu(t)(1 + 3 \sin(t)) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \mu(t)(1 + 3 \sin(t)) dt + c_1$$

- Solve for y

$$y = \frac{\int \mu(t)(1+3 \sin(t))dt+c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{-t}$

$$y = \frac{\int (1+3 \sin(t))e^{-t}dt+c_1}{e^{-t}}$$

- Evaluate the integrals on the rhs

$$y = \frac{-\frac{1}{e^t} - \frac{3e^{-t}\cos(t)}{2} - \frac{3e^{-t}\sin(t)}{2} + c_1}{e^{-t}}$$

- Simplify

$$y = -1 + c_1e^t - \frac{3 \cos(t)}{2} - \frac{3 \sin(t)}{2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 18

```
dsolve(-y(t)+diff(y(t),t) = 1+3*sin(t),y(t), singsol=all)
```

$$y(t) = -1 - \frac{3 \cos(t)}{2} - \frac{3 \sin(t)}{2} + e^t c_1$$

✓ Solution by Mathematica

Time used: 0.072 (sec). Leaf size: 25

```
DSolve[-y[t]+y'[t] == 1+3*Sin[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -\frac{3 \sin(t)}{2} - \frac{3 \cos(t)}{2} + c_1 e^t - 1$$

1.31 problem 31

1.31.1 Solving as linear ode	397
1.31.2 Solving as first order ode lie symmetry lookup ode	399
1.31.3 Solving as exact ode	403
1.31.4 Maple step by step solution	408

Internal problem ID [478]

Internal file name [OUTPUT/478_Sunday_June_05_2022_01_42_07_AM_9283464/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.1. Page 40

Problem number: 31.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$-\frac{3y}{2} + y' = 2e^t + 3t$$

1.31.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{3}{2}$$
$$q(t) = 2e^t + 3t$$

Hence the ode is

$$-\frac{3y}{2} + y' = 2e^t + 3t$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{3}{2} dt} \\ &= e^{-\frac{3t}{2}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu)(2e^t + 3t) \\ \frac{d}{dt}(e^{-\frac{3t}{2}} y) &= (e^{-\frac{3t}{2}})(2e^t + 3t) \\ d(e^{-\frac{3t}{2}} y) &= ((2e^t + 3t)e^{-\frac{3t}{2}}) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-\frac{3t}{2}} y &= \int (2e^t + 3t)e^{-\frac{3t}{2}} dt \\ e^{-\frac{3t}{2}} y &= -2e^{-\frac{3t}{2}} t - \frac{4e^{-\frac{3t}{2}}}{3} - 4e^{-\frac{t}{2}} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-\frac{3t}{2}}$ results in

$$y = e^{\frac{3t}{2}} \left(-2e^{-\frac{3t}{2}} t - \frac{4e^{-\frac{3t}{2}}}{3} - 4e^{-\frac{t}{2}} \right) + c_1 e^{\frac{3t}{2}}$$

which simplifies to

$$y = -2t - \frac{4}{3} - 4e^t + c_1 e^{\frac{3t}{2}}$$

Summary

The solution(s) found are the following

$$y = -2t - \frac{4}{3} - 4e^t + c_1 e^{\frac{3t}{2}} \quad (1)$$

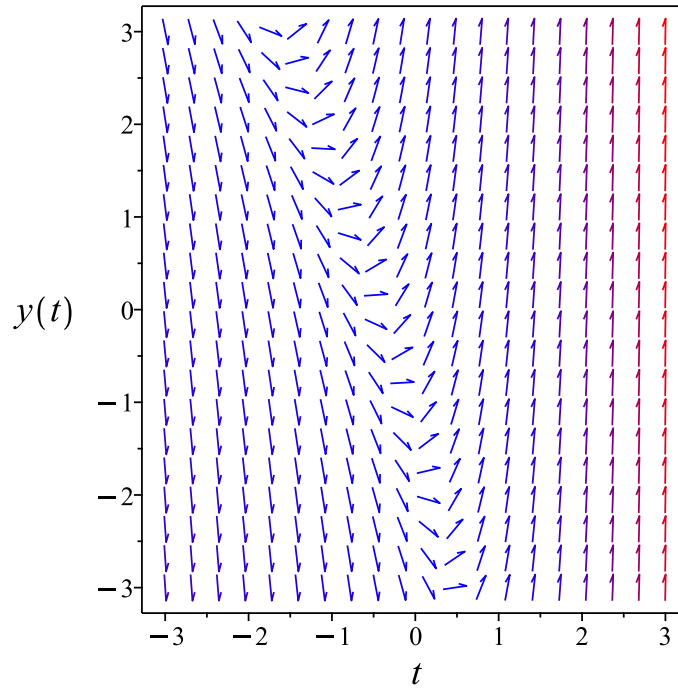


Figure 74: Slope field plot

Verification of solutions

$$y = -2t - \frac{4}{3} - 4e^t + c_1 e^{\frac{3t}{2}}$$

Verified OK.

1.31.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{3y}{2} + 2e^t + 3t$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 91: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{\frac{3t}{2}}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{\frac{3t}{2}}} dy \end{aligned}$$

Which results in

$$S = e^{-\frac{3t}{2}} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{3y}{2} + 2e^t + 3t$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -\frac{3e^{-\frac{3t}{2}}y}{2} \\ S_y &= e^{-\frac{3t}{2}} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 2e^{-\frac{t}{2}} + 3e^{-\frac{3t}{2}}t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 2e^{-\frac{R}{2}} + 3e^{-\frac{3R}{2}}R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -2e^{-\frac{3R}{2}}R - \frac{4e^{-\frac{3R}{2}}}{3} - 4e^{-\frac{R}{2}} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{-\frac{3t}{2}}y = -2e^{-\frac{3t}{2}}t - \frac{4e^{-\frac{3t}{2}}}{3} - 4e^{-\frac{t}{2}} + c_1$$

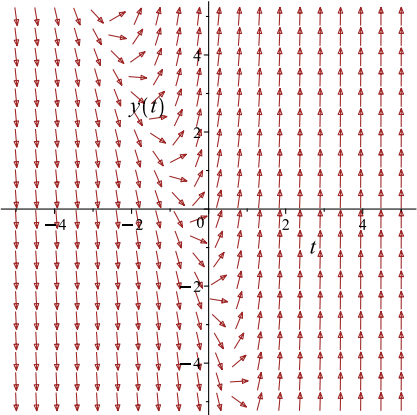
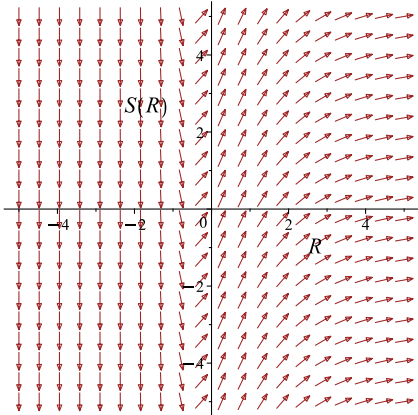
Which simplifies to

$$\frac{(6t + 3y + 4)e^{-\frac{3t}{2}}}{3} - c_1 + 4e^{-\frac{t}{2}} = 0$$

Which gives

$$y = -\frac{(6e^{-\frac{3t}{2}}t + 4e^{-\frac{3t}{2}} + 12e^{-\frac{t}{2}} - 3c_1)e^{\frac{3t}{2}}}{3}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{3y}{2} + 2e^t + 3t$ 	$R = t$ $S = e^{-\frac{3t}{2}}y$	$\frac{dS}{dR} = 2e^{-\frac{R}{2}} + 3e^{-\frac{3R}{2}}R$ 

Summary

The solution(s) found are the following

$$y = -\frac{\left(6e^{-\frac{3t}{2}}t + 4e^{-\frac{3t}{2}} + 12e^{-\frac{t}{2}} - 3c_1\right)e^{\frac{3t}{2}}}{3} \quad (1)$$

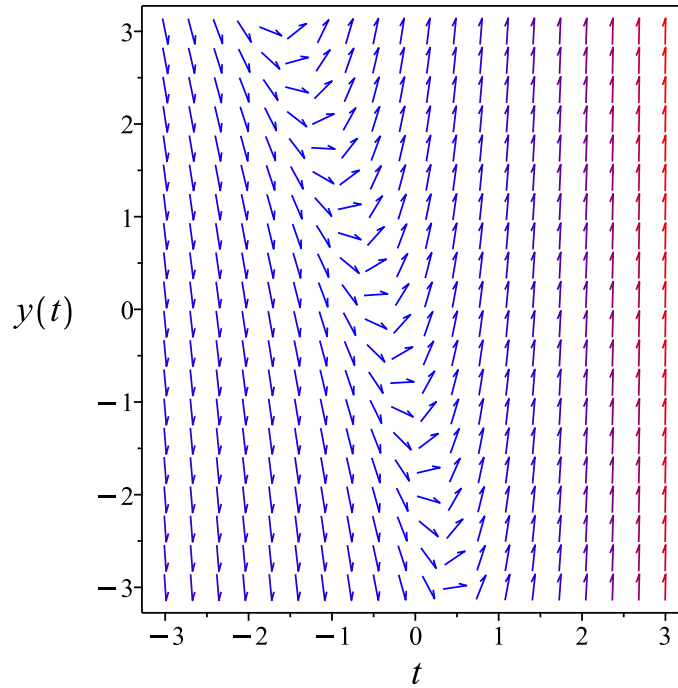


Figure 75: Slope field plot

Verification of solutions

$$y = -\frac{\left(6e^{-\frac{3t}{2}}t + 4e^{-\frac{3t}{2}} + 12e^{-\frac{t}{2}} - 3c_1\right)e^{\frac{3t}{2}}}{3}$$

Verified OK.

1.31.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the

ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} dy &= \left(\frac{3y}{2} + 2e^t + 3t \right) dt \\ \left(-\frac{3y}{2} - 2e^t - 3t \right) dt + dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= -\frac{3y}{2} - 2e^t - 3t \\ N(t, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{3y}{2} - 2e^t - 3t \right) \\ &= -\frac{3}{2}\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1 \left(\left(-\frac{3}{2} \right) - (0) \right) \\ &= -\frac{3}{2}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int -\frac{3}{2} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\frac{3t}{2}} \\ &= e^{-\frac{3t}{2}}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\overline{M} &= \mu M \\ &= e^{-\frac{3t}{2}} \left(-\frac{3y}{2} - 2e^t - 3t \right) \\ &= -\frac{(3y + 4e^t + 6t) e^{-\frac{3t}{2}}}{2}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^{-\frac{3t}{2}}(1) \\ &= e^{-\frac{3t}{2}}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left(-\frac{(3y + 4e^t + 6t)e^{-\frac{3t}{2}}}{2} \right) + \left(e^{-\frac{3t}{2}} \right) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -\frac{(3y + 4e^t + 6t)e^{-\frac{3t}{2}}}{2} dt \\ \phi &= \frac{(6t + 3y + 4)e^{-\frac{3t}{2}}}{3} + 4e^{-\frac{t}{2}} + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-\frac{3t}{2}} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-\frac{3t}{2}}$. Therefore equation (4) becomes

$$e^{-\frac{3t}{2}} = e^{-\frac{3t}{2}} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{(6t + 3y + 4)e^{-\frac{3t}{2}}}{3} + 4e^{-\frac{t}{2}} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{(6t + 3y + 4)e^{-\frac{3t}{2}}}{3} + 4e^{-\frac{t}{2}}$$

The solution becomes

$$y = -\frac{\left(6e^{-\frac{3t}{2}}t + 4e^{-\frac{3t}{2}} + 12e^{-\frac{t}{2}} - 3c_1\right)e^{\frac{3t}{2}}}{3}$$

Summary

The solution(s) found are the following

$$y = -\frac{\left(6e^{-\frac{3t}{2}}t + 4e^{-\frac{3t}{2}} + 12e^{-\frac{t}{2}} - 3c_1\right)e^{\frac{3t}{2}}}{3} \quad (1)$$

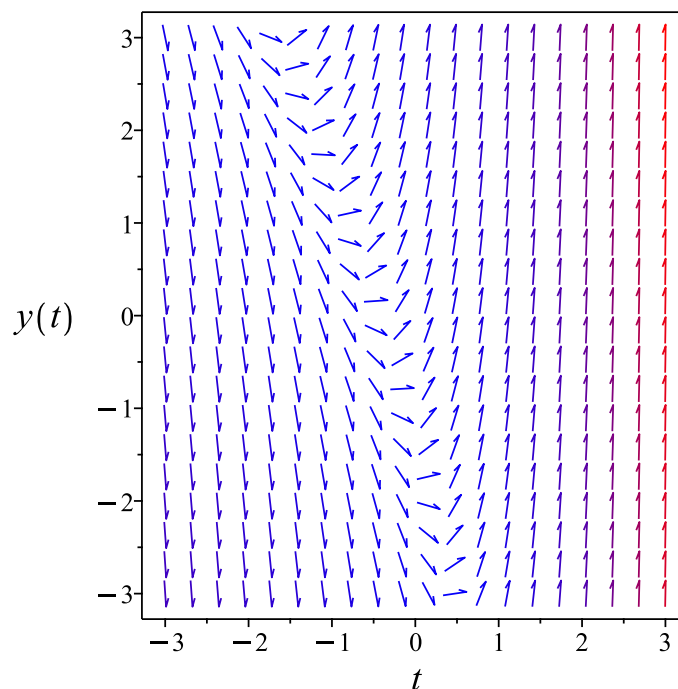


Figure 76: Slope field plot

Verification of solutions

$$y = -\frac{\left(6e^{-\frac{3t}{2}}t + 4e^{-\frac{3t}{2}} + 12e^{-\frac{t}{2}} - 3c_1\right)e^{\frac{3t}{2}}}{3}$$

Verified OK.

1.31.4 Maple step by step solution

Let's solve

$$-\frac{3y}{2} + y' = 2e^t + 3t$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = \frac{3y}{2} + 2e^t + 3t$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$-\frac{3y}{2} + y' = 2e^t + 3t$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(-\frac{3y}{2} + y'\right) = \mu(t) (2e^t + 3t)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(-\frac{3y}{2} + y'\right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -\frac{3\mu(t)}{2}$$

- Solve to find the integrating factor

$$\mu(t) = e^{-\frac{3t}{2}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y)\right) dt = \int \mu(t) (2e^t + 3t) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \mu(t) (2e^t + 3t) dt + c_1$$

- Solve for y

$$y = \frac{\int \mu(t)(2e^t+3t)dt+c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{-\frac{3t}{2}}$

$$y = \frac{\int (2e^t+3t)e^{-\frac{3t}{2}} dt+c_1}{e^{-\frac{3t}{2}}}$$

- Evaluate the integrals on the rhs

$$y = \frac{-2e^{-\frac{3t}{2}}t - \frac{4e^{-\frac{3t}{2}}}{3} - 4e^{-\frac{t}{2}} + c_1}{e^{-\frac{3t}{2}}}$$

- Simplify

$$y = -2t - \frac{4}{3} - 4e^t + c_1e^{\frac{3t}{2}}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(-3/2*y(t)+diff(y(t),t) = 2*exp(t)+3*t,y(t), singsol=all)
```

$$y(t) = -2t - \frac{4}{3} - 4e^t + e^{\frac{3t}{2}}c_1$$

✓ Solution by Mathematica

Time used: 0.1 (sec). Leaf size: 27

```
DSolve[-3/2*y[t]+y'[t] == 2*Exp[t]+3*t,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -2t - 4e^t + c_1e^{3t/2} - \frac{4}{3}$$

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2.1 problem 1

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Internal problem ID [479]

Internal file name [OUTPUT/479_Sunday_June_05_2022_01_42_08_AM_65104678/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{x^2}{y} = 0$$

2.1.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{x^2}{y}\end{aligned}$$

Where $f(x) = x^2$ and $g(y) = \frac{1}{y}$. Integrating both sides gives

$$\frac{1}{y} dy = x^2 dx$$

$$\int \frac{1}{y} dy = \int x^2 dx$$

$$\frac{y^2}{2} = \frac{x^3}{3} + c_1$$

Which results in

$$y = \frac{\sqrt{6x^3 + 18c_1}}{3}$$

$$y = -\frac{\sqrt{6x^3 + 18c_1}}{3}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{6x^3 + 18c_1}}{3} \tag{1}$$

$$y = -\frac{\sqrt{6x^3 + 18c_1}}{3} \tag{2}$$

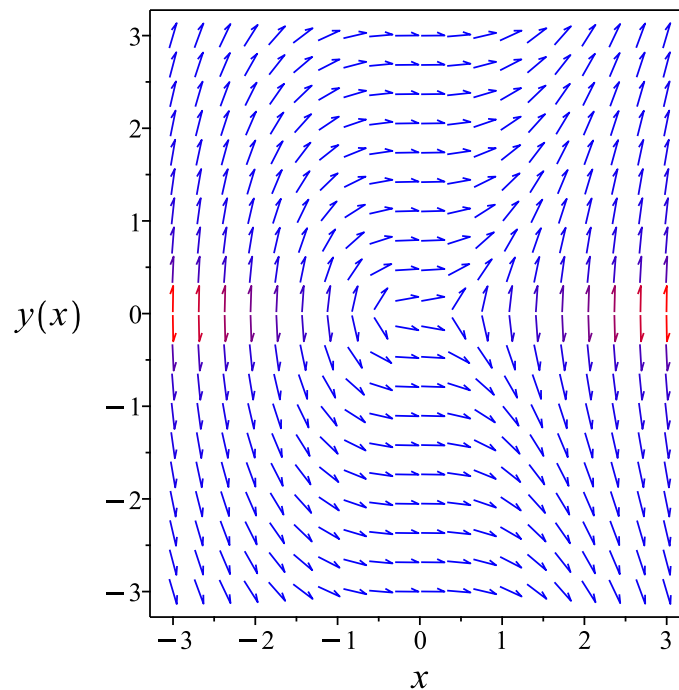


Figure 77: Slope field plot

Verification of solutions

$$y = \frac{\sqrt{6x^3 + 18c_1}}{3}$$

Verified OK.

$$y = -\frac{\sqrt{6x^3 + 18c_1}}{3}$$

Verified OK.

2.1.2 Solving as differentialType ode

Writing the ode as

$$y' = \frac{x^2}{y} \tag{1}$$

Which becomes

$$(y) dy = (x^2) dx \tag{2}$$

But the RHS is complete differential because

$$(x^2) dx = d\left(\frac{x^3}{3}\right)$$

Hence (2) becomes

$$(y) dy = d\left(\frac{x^3}{3}\right)$$

Integrating both sides gives gives these solutions

$$y = \frac{\sqrt{6x^3 + 18c_1}}{3} + c_1$$

$$y = -\frac{\sqrt{6x^3 + 18c_1}}{3} + c_1$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{6x^3 + 18c_1}}{3} + c_1 \tag{1}$$

$$y = -\frac{\sqrt{6x^3 + 18c_1}}{3} + c_1 \tag{2}$$

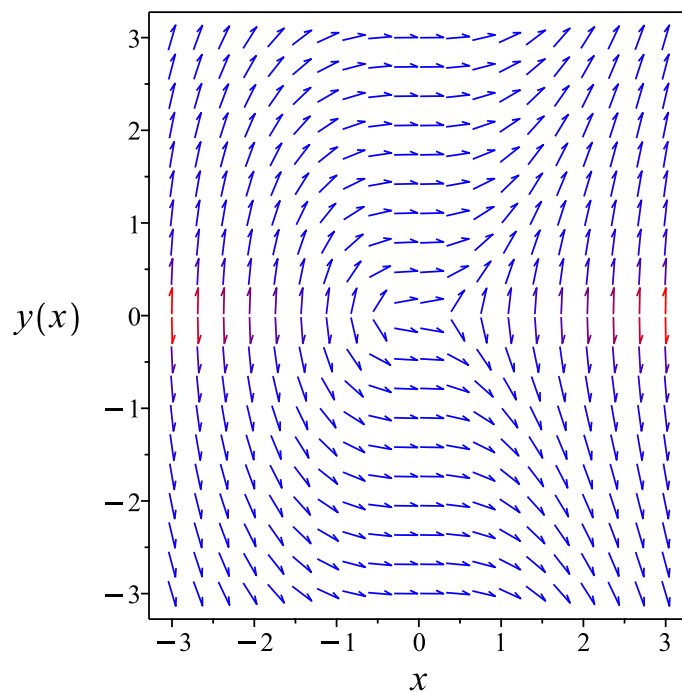


Figure 78: Slope field plot

Verification of solutions

$$y = \frac{\sqrt{6x^3 + 18c_1}}{3} + c_1$$

Verified OK.

$$y = -\frac{\sqrt{6x^3 + 18c_1}}{3} + c_1$$

Verified OK.

2.1.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x^2}{y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 94: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{x^2} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{x^2}} dx \end{aligned}$$

Which results in

$$S = \frac{x^3}{3}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x^2}{y}$$

Evaluating all the partial derivatives gives

$$R_x = 0$$

$$R_y = 1$$

$$S_x = x^2$$

$$S_y = 0$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^2}{2} + c_1 \quad (4)$$

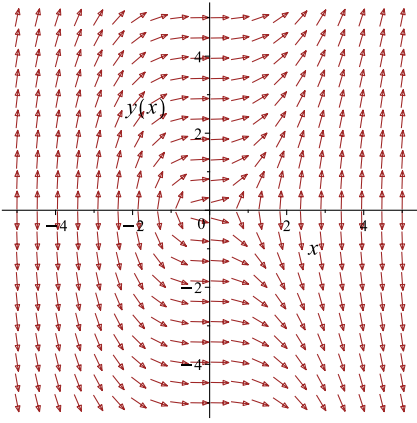
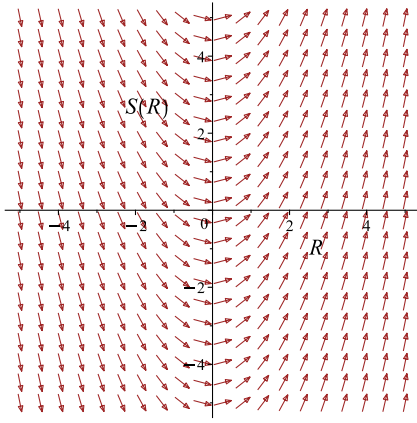
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{x^3}{3} = \frac{y^2}{2} + c_1$$

Which simplifies to

$$\frac{x^3}{3} = \frac{y^2}{2} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{x^2}{y}$ 	$R = y$ $S = \frac{x^3}{3}$	$\frac{dS}{dR} = R$ 

Summary

The solution(s) found are the following

$$\frac{x^3}{3} = \frac{y^2}{2} + c_1 \quad (1)$$

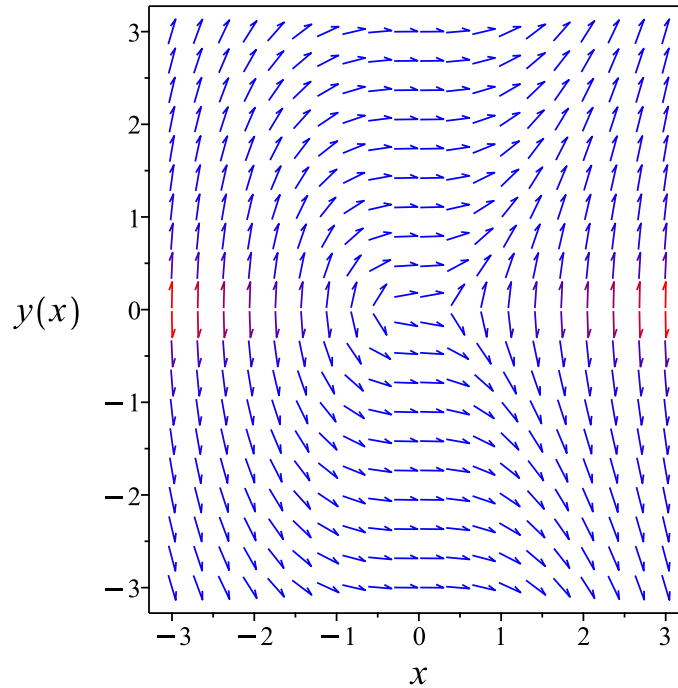


Figure 79: Slope field plot

Verification of solutions

$$\frac{x^3}{3} = \frac{y^2}{2} + c_1$$

Verified OK.

2.1.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(y) dy &= (x^2) dx \\ (-x^2) dx + (y) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -x^2 \\ N(x, y) &= y\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x^2) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(y) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x^2 dx \\ \phi &= -\frac{x^3}{3} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = y$. Therefore equation (4) becomes

$$y = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = y$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int (y) dy \\ f(y) &= \frac{y^2}{2} + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{x^3}{3} + \frac{y^2}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{x^3}{3} + \frac{y^2}{2}$$

Summary

The solution(s) found are the following

$$-\frac{x^3}{3} + \frac{y^2}{2} = c_1 \tag{1}$$

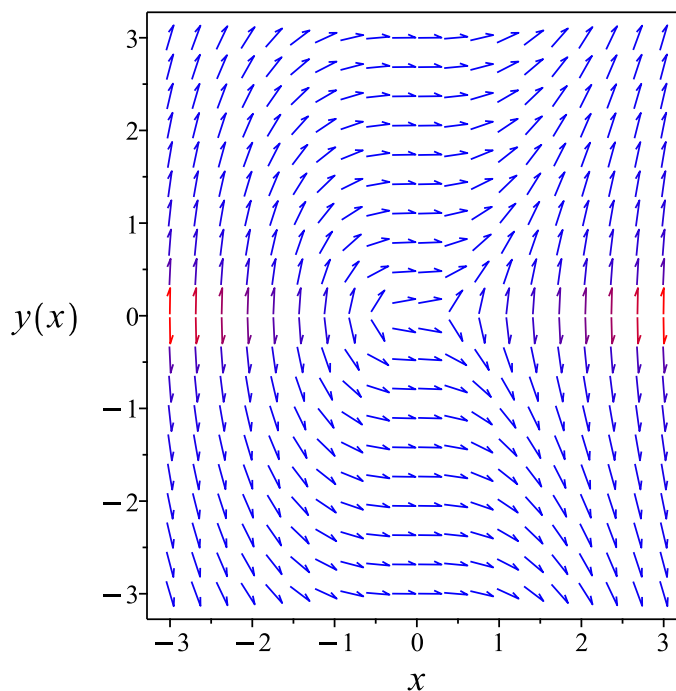


Figure 80: Slope field plot

Verification of solutions

$$-\frac{x^3}{3} + \frac{y^2}{2} = c_1$$

Verified OK.

2.1.5 Maple step by step solution

Let's solve

$$y' - \frac{x^2}{y} = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$yy' = x^2$$

- Integrate both sides with respect to x

$$\int yy'dx = \int x^2 dx + c_1$$

- Evaluate integral

$$\frac{y^2}{2} = \frac{x^3}{3} + c_1$$

- Solve for y

$$\left\{ y = -\frac{\sqrt{6x^3+18c_1}}{3}, y = \frac{\sqrt{6x^3+18c_1}}{3} \right\}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 33

```
dsolve(diff(y(x),x) = x^2/y(x),y(x), singsol=all)
```

$$y(x) = -\frac{\sqrt{6x^3 + 9c_1}}{3}$$
$$y(x) = \frac{\sqrt{6x^3 + 9c_1}}{3}$$

✓ Solution by Mathematica

Time used: 0.093 (sec). Leaf size: 50

```
DSolve[y'[x] == x^2/y[x], y[x], x, IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\sqrt{\frac{2}{3}}\sqrt{x^3 + 3c_1}$$

$$y(x) \rightarrow \sqrt{\frac{2}{3}}\sqrt{x^3 + 3c_1}$$

2.2 problem 2

2.2.1	Solving as separable ode	425
2.2.2	Solving as first order ode lie symmetry lookup ode	427
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2.2.4	Maple step by step solution	435

Internal problem ID [480]

Internal file name [OUTPUT/480_Sunday_June_05_2022_01_42_09_AM_72882121/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{x^2}{(x^3 + 1)y} = 0$$

2.2.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{x^2}{(x^3 + 1)y}\end{aligned}$$

Where $f(x) = \frac{x^2}{x^3+1}$ and $g(y) = \frac{1}{y}$. Integrating both sides gives

$$\frac{1}{y} dy = \frac{x^2}{x^3 + 1} dx$$

$$\int \frac{1}{y} dy = \int \frac{x^2}{x^3 + 1} dx$$

$$\frac{y^2}{2} = \frac{\ln(x^3 + 1)}{3} + c_1$$

Which results in

$$y = \frac{\sqrt{6 \ln(x^3 + 1) + 18c_1}}{3}$$

$$y = -\frac{\sqrt{6 \ln(x^3 + 1) + 18c_1}}{3}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{6 \ln(x^3 + 1) + 18c_1}}{3} \tag{1}$$

$$y = -\frac{\sqrt{6 \ln(x^3 + 1) + 18c_1}}{3} \tag{2}$$

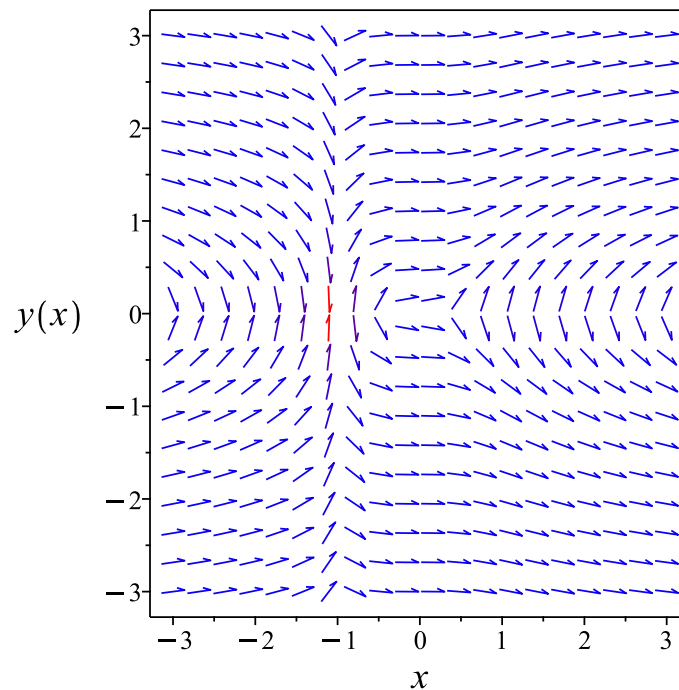


Figure 81: Slope field plot

Verification of solutions

$$y = \frac{\sqrt{6 \ln(x^3 + 1) + 18c_1}}{3}$$

Verified OK.

$$y = -\frac{\sqrt{6 \ln(x^3 + 1) + 18c_1}}{3}$$

Verified OK.

2.2.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x^2}{(x^3 + 1)y}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 97: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{x^3 + 1}{x^2} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{x^3+1}{x^2}} dx \end{aligned}$$

Which results in

$$S = \frac{\ln(x^3 + 1)}{3}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x^2}{(x^3 + 1)y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= \frac{x^2}{(x^2 - x + 1)(x + 1)} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^2}{2} + c_1 \quad (4)$$

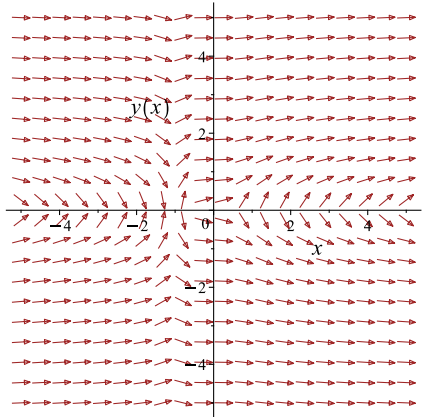
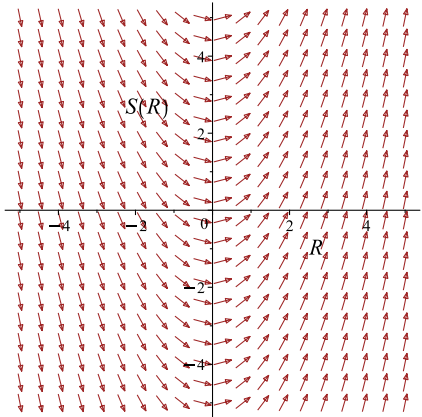
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\ln(x+1)}{3} + \frac{\ln(x^2-x+1)}{3} = \frac{y^2}{2} + c_1$$

Which simplifies to

$$\frac{\ln(x+1)}{3} + \frac{\ln(x^2-x+1)}{3} = \frac{y^2}{2} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{x^2}{(x^3+1)y}$ 	$R = y$ $S = \frac{\ln(x+1)}{3} + \frac{\ln(x^2-x+1)}{3}$	$\frac{dS}{dR} = R$ 

Summary

The solution(s) found are the following

$$\frac{\ln(x+1)}{3} + \frac{\ln(x^2-x+1)}{3} = \frac{y^2}{2} + c_1 \quad (1)$$

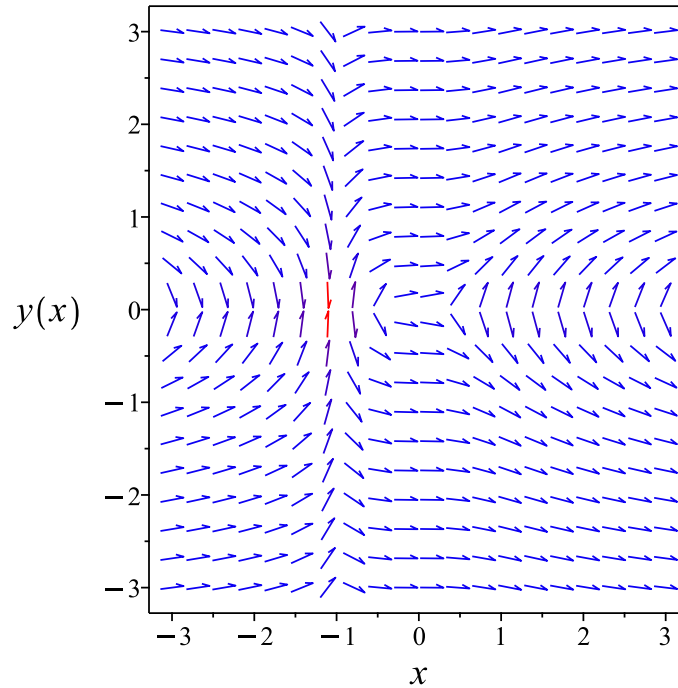


Figure 82: Slope field plot

Verification of solutions

$$\frac{\ln(x+1)}{3} + \frac{\ln(x^2 - x + 1)}{3} = \frac{y^2}{2} + c_1$$

Verified OK.

2.2.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(y) dy &= \left(\frac{x^2}{x^3 + 1} \right) dx \\ \left(-\frac{x^2}{x^3 + 1} \right) dx + (y) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{x^2}{x^3 + 1} \\ N(x, y) &= y\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{x^2}{x^3 + 1} \right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(y) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{x^2}{x^3 + 1} dx \\ \phi &= -\frac{\ln(x^3 + 1)}{3} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = y$. Therefore equation (4) becomes

$$y = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = y$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (y) dy \\ f(y) &= \frac{y^2}{2} + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{\ln(x^3 + 1)}{3} + \frac{y^2}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{\ln(x^3 + 1)}{3} + \frac{y^2}{2}$$

Summary

The solution(s) found are the following

$$-\frac{\ln(x^3 + 1)}{3} + \frac{y^2}{2} = c_1 \quad (1)$$

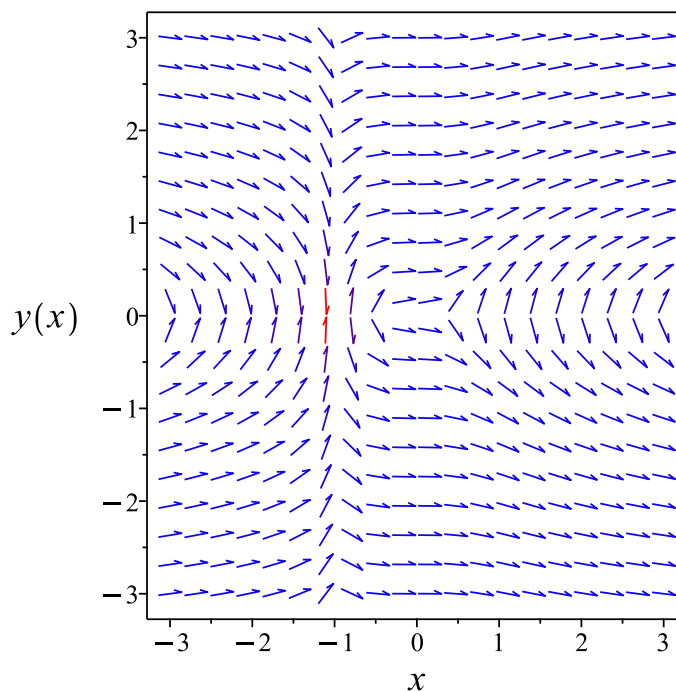


Figure 83: Slope field plot

Verification of solutions

$$-\frac{\ln(x^3 + 1)}{3} + \frac{y^2}{2} = c_1$$

Verified OK.

2.2.4 Maple step by step solution

Let's solve

$$y' - \frac{x^2}{(x^3+1)y} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$yy' = \frac{x^2}{x^3+1}$$

- Integrate both sides with respect to x

$$\int yy' dx = \int \frac{x^2}{x^3+1} dx + c_1$$

- Evaluate integral

$$\frac{y^2}{2} = \frac{\ln(x^3+1)}{3} + c_1$$

- Solve for y

$$\left\{ y = -\frac{\sqrt{6\ln(x^3+1)+18c_1}}{3}, y = \frac{\sqrt{6\ln(x^3+1)+18c_1}}{3} \right\}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 39

```
dsolve(diff(y(x),x) = x^2/(x^3+1)/y(x),y(x), singsol=all)
```

$$y(x) = -\frac{\sqrt{6\ln(x^3+1)+9c_1}}{3}$$
$$y(x) = \frac{\sqrt{6\ln(x^3+1)+9c_1}}{3}$$

✓ Solution by Mathematica

Time used: 0.12 (sec). Leaf size: 56

```
DSolve[y'[x] == x^2/(x^3+1)/y[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\sqrt{\frac{2}{3}}\sqrt{\log(x^3 + 1) + 3c_1}$$

$$y(x) \rightarrow \sqrt{\frac{2}{3}}\sqrt{\log(x^3 + 1) + 3c_1}$$

2.3 problem 3

2.3.1	Solving as separable ode	437
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Internal problem ID [481]

Internal file name [OUTPUT/481_Sunday_June_05_2022_01_42_10_AM_4186121/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 3.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$\sin(x)y^2 + y' = 0$$

2.3.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\sin(x)y^2\end{aligned}$$

Where $f(x) = -\sin(x)$ and $g(y) = y^2$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y^2} dy &= -\sin(x) dx \\ \int \frac{1}{y^2} dy &= \int -\sin(x) dx\end{aligned}$$

$$-\frac{1}{y} = \cos(x) + c_1$$

Which results in

$$y = -\frac{1}{\cos(x) + c_1}$$

Summary

The solution(s) found are the following

$$y = -\frac{1}{\cos(x) + c_1} \tag{1}$$

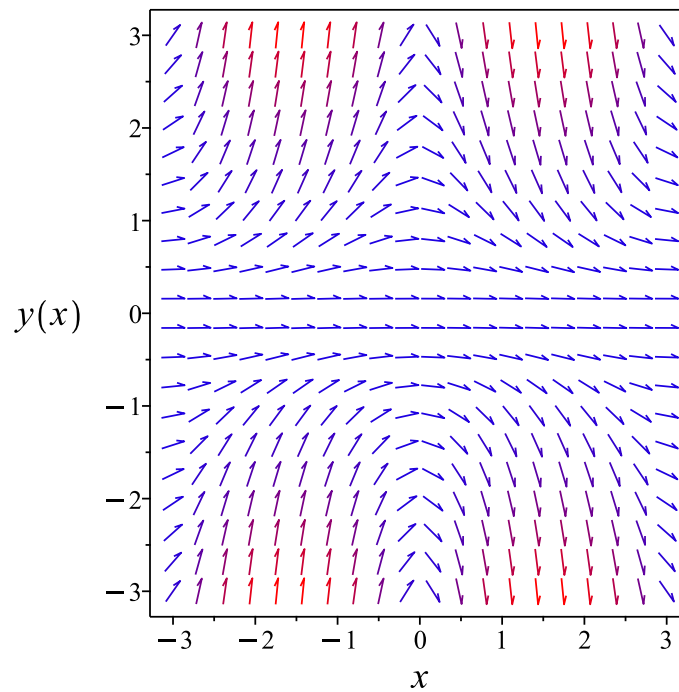


Figure 84: Slope field plot

Verification of solutions

$$y = -\frac{1}{\cos(x) + c_1}$$

Verified OK.

2.3.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\sin(x) y^2$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 100: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int(n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= -\frac{1}{\sin(x)} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{-\frac{1}{\sin(x)}} dx\end{aligned}$$

Which results in

$$S = \cos(x)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\sin(x) y^2$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= -\sin(x) \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{y^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{1}{R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\cos(x) = -\frac{1}{y} + c_1$$

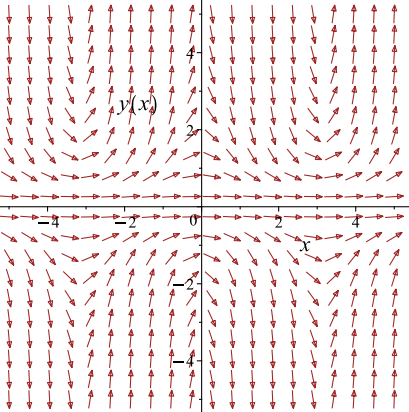
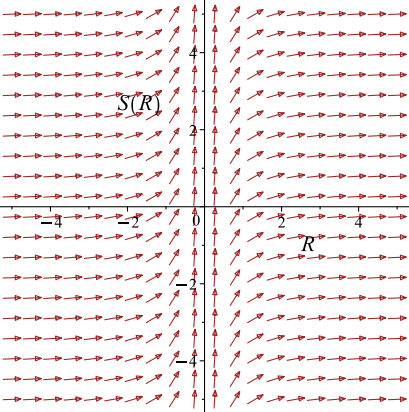
Which simplifies to

$$\cos(x) = -\frac{1}{y} + c_1$$

Which gives

$$y = -\frac{1}{\cos(x) - c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\sin(x) y^2$ 	$R = y$ $S = \cos(x)$	$\frac{dS}{dR} = \frac{1}{R^2}$ 

Summary

The solution(s) found are the following

$$y = -\frac{1}{\cos(x) - c_1} \tag{1}$$

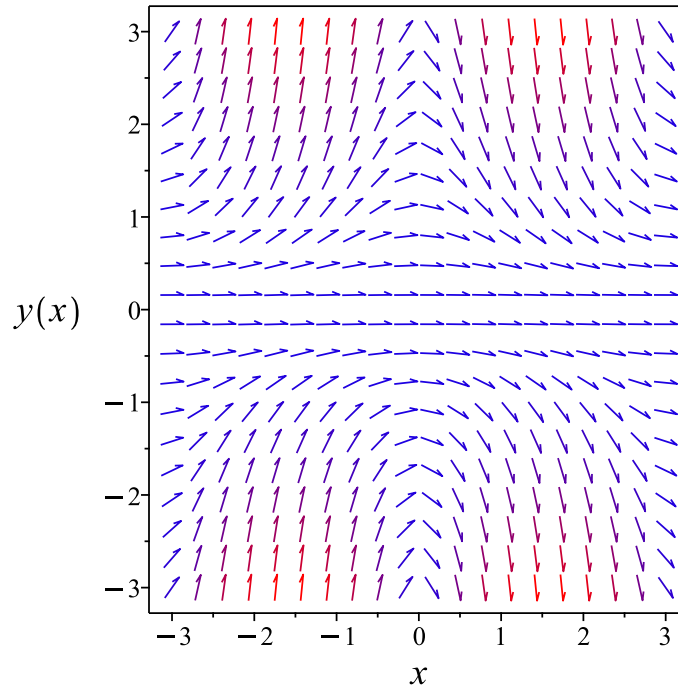


Figure 85: Slope field plot

Verification of solutions

$$y = -\frac{1}{\cos(x) - c_1}$$

Verified OK.

2.3.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}\left(-\frac{1}{y^2}\right) dy &= (\sin(x)) dx \\ (-\sin(x)) dx + \left(-\frac{1}{y^2}\right) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\sin(x) \\ N(x, y) &= -\frac{1}{y^2}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-\sin(x)) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(-\frac{1}{y^2} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\sin(x) dx \\ \phi &= \cos(x) + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{1}{y^2}$. Therefore equation (4) becomes

$$-\frac{1}{y^2} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{1}{y^2}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(-\frac{1}{y^2} \right) dy \\ f(y) &= \frac{1}{y} + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = \cos(x) + \frac{1}{y} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \cos(x) + \frac{1}{y}$$

The solution becomes

$$y = -\frac{1}{\cos(x) - c_1}$$

Summary

The solution(s) found are the following

$$y = -\frac{1}{\cos(x) - c_1} \tag{1}$$

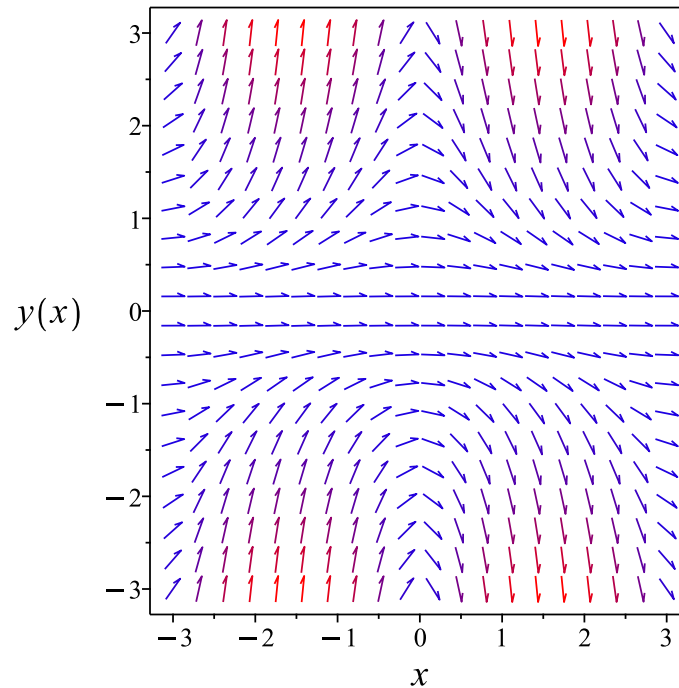


Figure 86: Slope field plot

Verification of solutions

$$y = -\frac{1}{\cos(x) - c_1}$$

Verified OK.

2.3.4 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= -\sin(x) y^2\end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = -\sin(x) y^2$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that $f_0(x) = 0$, $f_1(x) = 0$ and $f_2(x) = -\sin(x)$. Let

$$\begin{aligned}y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{-\sin(x) u}\end{aligned}\tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0\tag{2}$$

But

$$\begin{aligned}f_2' &= -\cos(x) \\ f_1 f_2 &= 0 \\ f_2^2 f_0 &= 0\end{aligned}$$

Substituting the above terms back in equation (2) gives

$$-\sin(x) u''(x) + \cos(x) u'(x) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 + \cos(x) c_2$$

The above shows that

$$u'(x) = -c_2 \sin(x)$$

Using the above in (1) gives the solution

$$y = -\frac{c_2}{c_1 + \cos(x) c_2}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = -\frac{1}{c_3 + \cos(x)}$$

Summary

The solution(s) found are the following

$$y = -\frac{1}{c_3 + \cos(x)} \tag{1}$$

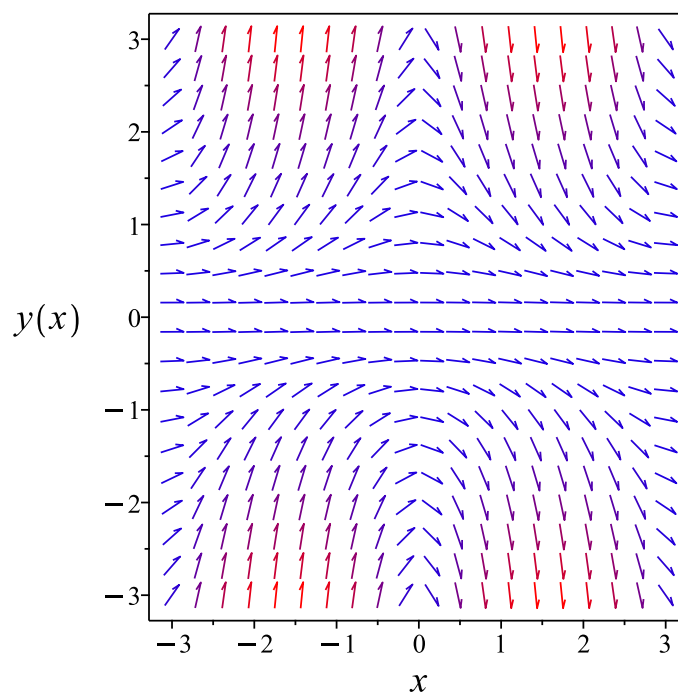


Figure 87: Slope field plot

Verification of solutions

$$y = -\frac{1}{c_3 + \cos(x)}$$

Verified OK.

2.3.5 Maple step by step solution

Let's solve

$$\sin(x)y^2 + y' = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{y^2} = -\sin(x)$$

- Integrate both sides with respect to x

$$\int \frac{y'}{y^2} dx = \int -\sin(x) dx + c_1$$

- Evaluate integral

$$-\frac{1}{y} = \cos(x) + c_1$$

- Solve for y

$$y = -\frac{1}{\cos(x) + c_1}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 12

```
dsolve(sin(x)*y(x)^2+diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = \frac{1}{-\cos(x) + c_1}$$

✓ Solution by Mathematica

Time used: 0.133 (sec). Leaf size: 19

```
DSolve[Sin[x]*y[x]^2+y'[x]== 0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{1}{\cos(x) + c_1}$$
$$y(x) \rightarrow 0$$

2.4 problem 4

2.4.1	Solving as separable ode	451
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Internal problem ID [482]

Internal file name [OUTPUT/482_Sunday_June_05_2022_01_42_11_AM_75746612/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 4.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "differential-Type", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{3x^2 - 1}{3 + 2y} = 0$$

2.4.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{3x^2 - 1}{3 + 2y}\end{aligned}$$

Where $f(x) = 3x^2 - 1$ and $g(y) = \frac{1}{3+2y}$. Integrating both sides gives

$$\frac{1}{\frac{1}{3+2y}} dy = 3x^2 - 1 dx$$

$$\int \frac{1}{\frac{1}{3+2y}} dy = \int 3x^2 - 1 dx$$

$$y^2 + 3y = x^3 + c_1 - x$$

Which results in

$$y = -\frac{3}{2} + \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2}$$

$$y = -\frac{3}{2} - \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2}$$

Summary

The solution(s) found are the following

$$y = -\frac{3}{2} + \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2} \quad (1)$$

$$y = -\frac{3}{2} - \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2} \quad (2)$$

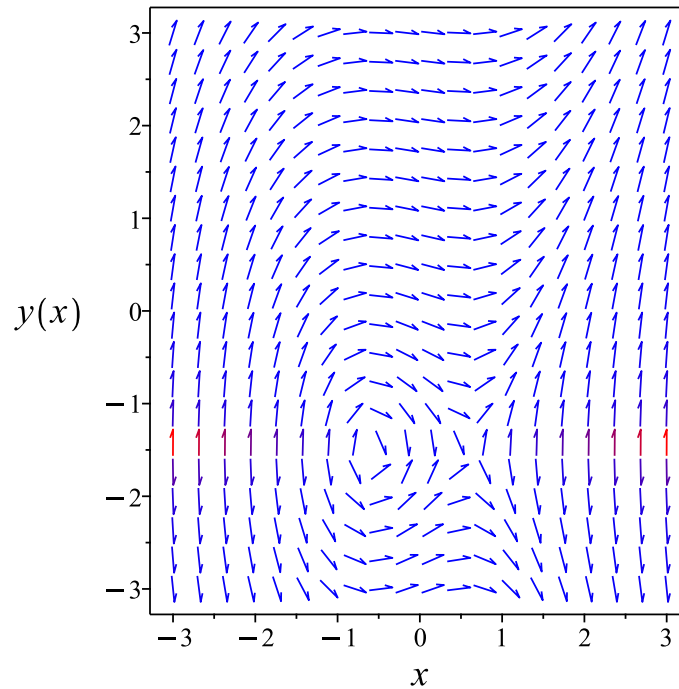


Figure 88: Slope field plot

Verification of solutions

$$y = -\frac{3}{2} + \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2}$$

Verified OK.

$$y = -\frac{3}{2} - \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2}$$

Verified OK.

2.4.2 Solving as differential Type ode

Writing the ode as

$$y' = \frac{3x^2 - 1}{3 + 2y} \quad (1)$$

Which becomes

$$(3 + 2y) dy = (3x^2 - 1) dx \quad (2)$$

But the RHS is complete differential because

$$(3x^2 - 1) dx = d(x^3 - x)$$

Hence (2) becomes

$$(3 + 2y) dy = d(x^3 - x)$$

Integrating both sides gives these solutions

$$y = -\frac{3}{2} + \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2} + c_1$$

$$y = -\frac{3}{2} - \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2} + c_1$$

Summary

The solution(s) found are the following

$$y = -\frac{3}{2} + \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2} + c_1 \quad (1)$$

$$y = -\frac{3}{2} - \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2} + c_1 \quad (2)$$

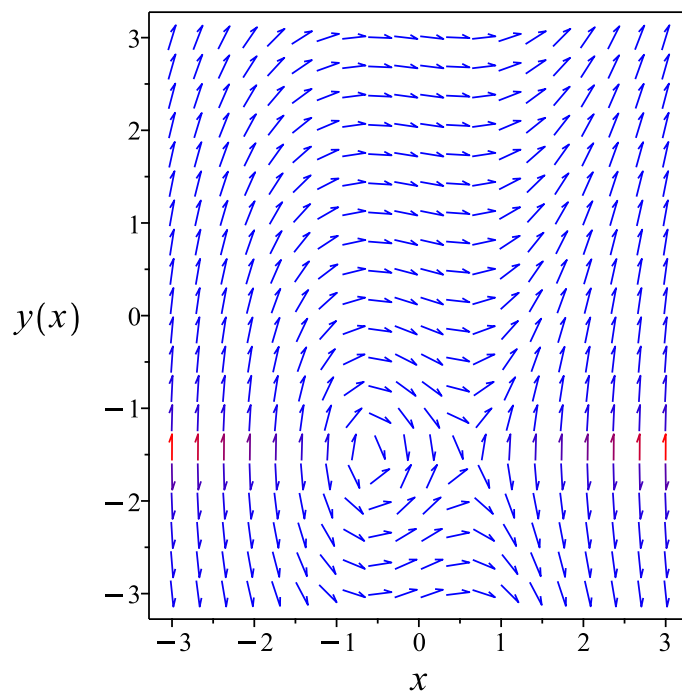


Figure 89: Slope field plot

Verification of solutions

$$y = -\frac{3}{2} + \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2} + c_1$$

Verified OK.

$$y = -\frac{3}{2} - \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2} + c_1$$

Verified OK.

2.4.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{3x^2 - 1}{3 + 2y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 103: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{3x^2 - 1} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{3x^2-1}} dx \end{aligned}$$

Which results in

$$S = x^3 - x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{3x^2 - 1}{3 + 2y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= 3x^2 - 1 \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 3 + 2y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 3 + 2R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R^2 + 3R + c_1 \quad (4)$$

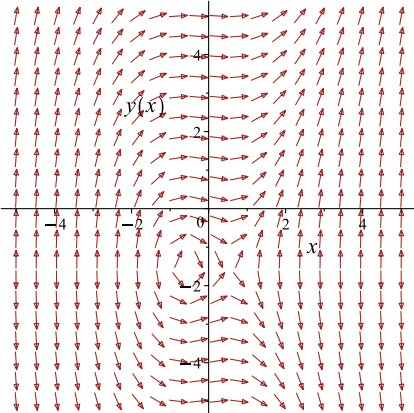
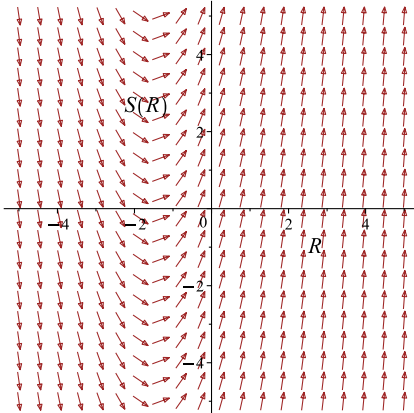
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$x^3 - x = y^2 + c_1 + 3y$$

Which simplifies to

$$x^3 - x = y^2 + c_1 + 3y$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{3x^2-1}{3+2y}$ 	$R = y$ $S = x^3 - x$	$\frac{dS}{dR} = 3 + 2R$ 

Summary

The solution(s) found are the following

$$x^3 - x = y^2 + c_1 + 3y \quad (1)$$

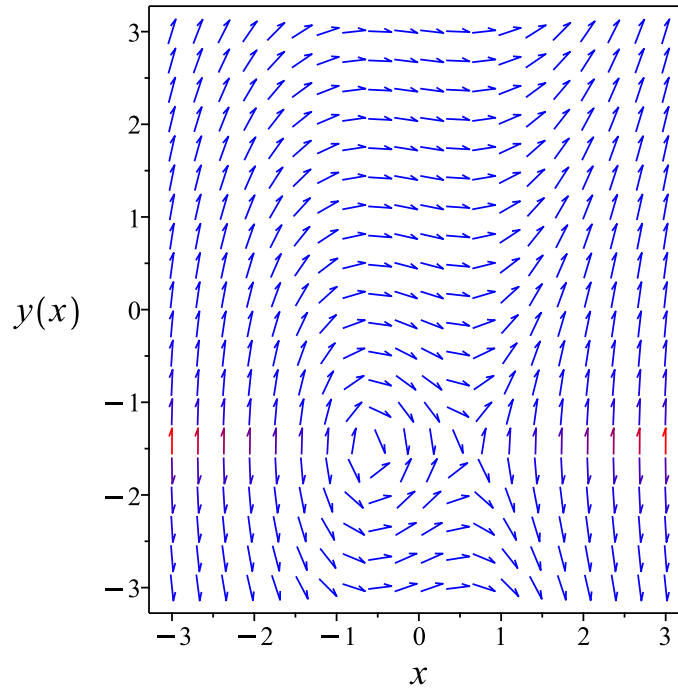


Figure 90: Slope field plot

Verification of solutions

$$x^3 - x = y^2 + c_1 + 3y$$

Verified OK.

2.4.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(3 + 2y) dy &= (3x^2 - 1) dx \\ (-3x^2 + 1) dx + (3 + 2y) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -3x^2 + 1 \\ N(x, y) &= 3 + 2y\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-3x^2 + 1) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(3 + 2y) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -3x^2 + 1 dx$$

$$\phi = -x^3 + x + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 3 + 2y$. Therefore equation (4) becomes

$$3 + 2y = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 3 + 2y$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (3 + 2y) dy$$

$$f(y) = y^2 + 3y + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x^3 + y^2 + x + 3y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x^3 + y^2 + x + 3y$$

Summary

The solution(s) found are the following

$$-x^3 + y^2 + 3y + x = c_1 \tag{1}$$

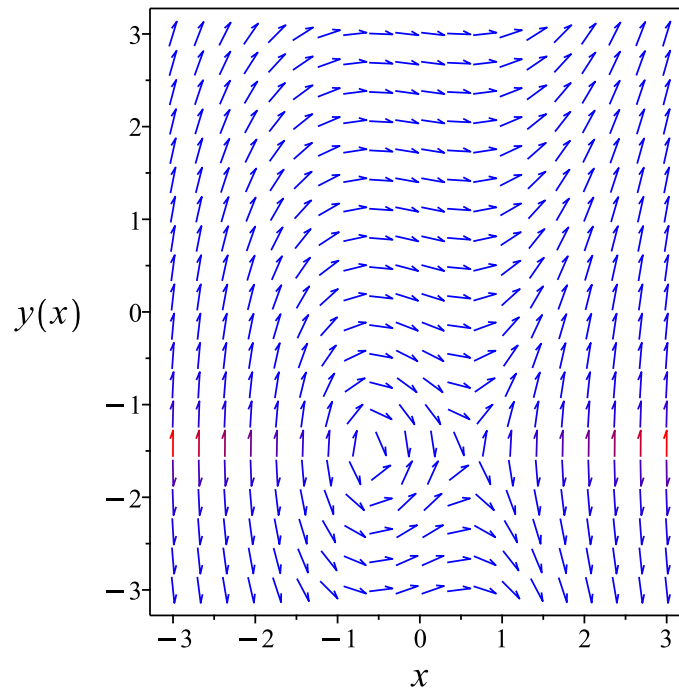


Figure 91: Slope field plot

Verification of solutions

$$-x^3 + y^2 + 3y + x = c_1$$

Verified OK.

2.4.5 Maple step by step solution

Let's solve

$$y' - \frac{3x^2-1}{3+2y} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$(3 + 2y) y' = 3x^2 - 1$$

- Integrate both sides with respect to x

$$\int (3 + 2y) y' dx = \int (3x^2 - 1) dx + c_1$$

- Evaluate integral

$$y^2 + 3y = x^3 + c_1 - x$$

- Solve for y

$$\left\{ y = -\frac{3}{2} - \frac{\sqrt{4x^3+4c_1-4x+9}}{2}, y = -\frac{3}{2} + \frac{\sqrt{4x^3+4c_1-4x+9}}{2} \right\}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 45

```
dsolve(diff(y(x),x) = (3*x^2-1)/(3+2*y(x)),y(x), singsol=all)
```

$$y(x) = -\frac{3}{2} - \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2}$$
$$y(x) = -\frac{3}{2} + \frac{\sqrt{4x^3 + 4c_1 - 4x + 9}}{2}$$

✓ Solution by Mathematica

Time used: 0.121 (sec). Leaf size: 59

```
DSolve[y'[x] == (3*x^2-1)/(3+2*y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2} \left(-3 - \sqrt{4x^3 - 4x + 9 + 4c_1} \right)$$

$$y(x) \rightarrow \frac{1}{2} \left(-3 + \sqrt{4x^3 - 4x + 9 + 4c_1} \right)$$

2.5 problem 5

2.5.1	Solving as separable ode	464
2.5.2	Solving as first order ode lie symmetry lookup ode	466
2.5.3	Solving as exact ode	470
2.5.4	Maple step by step solution	474

Internal problem ID [483]

Internal file name [OUTPUT/483_Sunday_June_05_2022_01_42_12_AM_72792709/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 5.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y' - \cos(x)^2 \cos(2y)^2 = 0$$

2.5.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \cos(x)^2 \cos(2y)^2\end{aligned}$$

Where $f(x) = \cos(x)^2$ and $g(y) = \cos(2y)^2$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\cos(2y)^2} dy &= \cos(x)^2 dx \\ \int \frac{1}{\cos(2y)^2} dy &= \int \cos(x)^2 dx \\ \frac{\tan(2y)}{2} &= \frac{\cos(x) \sin(x)}{2} + \frac{x}{2} + c_1\end{aligned}$$

Which results in

$$y = \frac{\arctan(\cos(x) \sin(x) + 2c_1 + x)}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{\arctan(\cos(x) \sin(x) + 2c_1 + x)}{2} \quad (1)$$

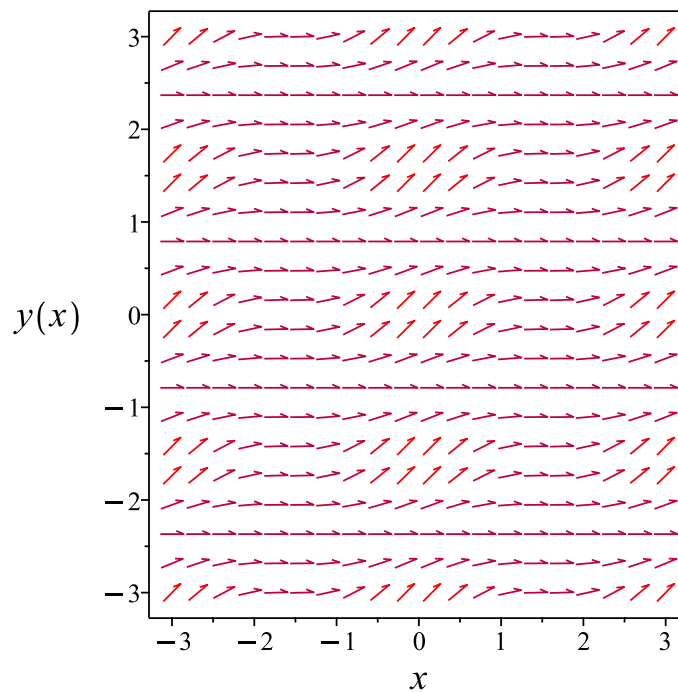


Figure 92: Slope field plot

Verification of solutions

$$y = \frac{\arctan(\cos(x) \sin(x) + 2c_1 + x)}{2}$$

Verified OK.

2.5.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \cos(x)^2 \cos(2y)^2$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 106: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int(n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{\cos(x)^2} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{\cos(x)^2}} dx\end{aligned}$$

Which results in

$$S = \frac{\cos(x) \sin(x)}{2} + \frac{x}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \cos(x)^2 \cos(2y)^2$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= \frac{\cos(2x)}{2} + \frac{1}{2} \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \sec(2y)^2 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \sec(2R)^2$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{\tan(2R)}{2} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\sin(2x)}{4} + \frac{x}{2} = \frac{\tan(2y)}{2} + c_1$$

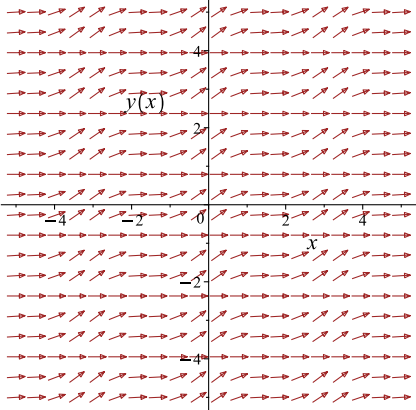
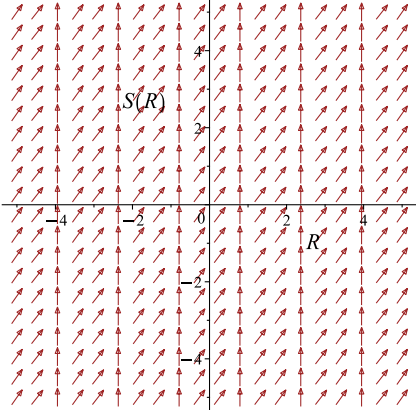
Which simplifies to

$$\frac{\sin(2x)}{4} + \frac{x}{2} = \frac{\tan(2y)}{2} + c_1$$

Which gives

$$y = -\frac{\arctan\left(-\frac{\sin(2x)}{2} - x + 2c_1\right)}{2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \cos(x)^2 \cos(2y)^2$ 	$R = y$ $S = \frac{\sin(2x)}{4} + \frac{x}{2}$	$\frac{dS}{dR} = \sec(2R)^2$ 

Summary

The solution(s) found are the following

$$y = -\frac{\arctan\left(-\frac{\sin(2x)}{2} - x + 2c_1\right)}{2} \quad (1)$$

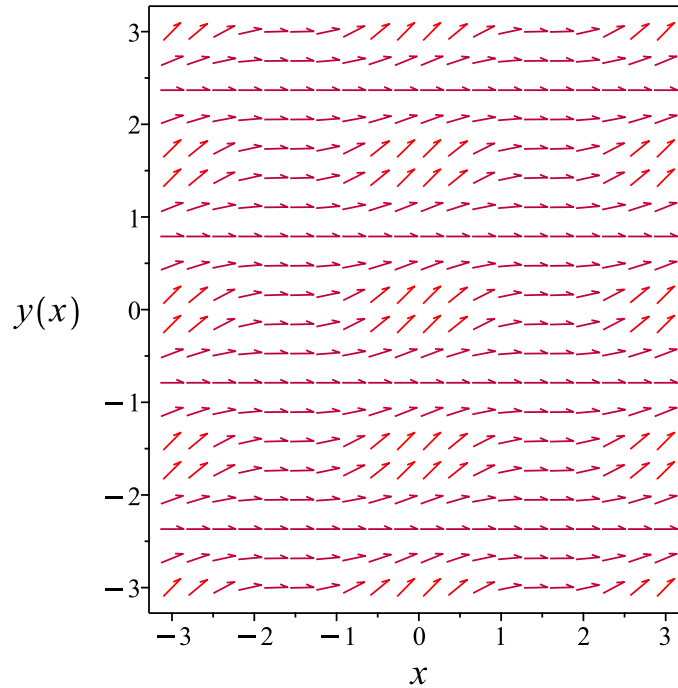


Figure 93: Slope field plot

Verification of solutions

$$y = -\frac{\arctan\left(-\frac{\sin(2x)}{2} - x + 2c_1\right)}{2}$$

Verified OK.

2.5.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}\left(\frac{1}{\cos(2y)^2}\right) dy &= (\cos(x)^2) dx \\ (-\cos(x)^2) dx + \left(\frac{1}{\cos(2y)^2}\right) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\cos(x)^2 \\ N(x, y) &= \frac{1}{\cos(2y)^2}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-\cos(x)^2) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{1}{\cos(2y)^2} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\cos(x)^2 dx \\ \phi &= -\frac{\sin(2x)}{4} - \frac{x}{2} + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{\cos(2y)^2}$. Therefore equation (4) becomes

$$\frac{1}{\cos(2y)^2} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{1}{\cos(2y)^2}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (\sec(2y))^2 dy$$
$$f(y) = \frac{\tan(2y)}{2} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{\sin(2x)}{4} - \frac{x}{2} + \frac{\tan(2y)}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{\sin(2x)}{4} - \frac{x}{2} + \frac{\tan(2y)}{2}$$

Summary

The solution(s) found are the following

$$-\frac{x}{2} - \frac{\sin(2x)}{4} + \frac{\tan(2y)}{2} = c_1 \quad (1)$$

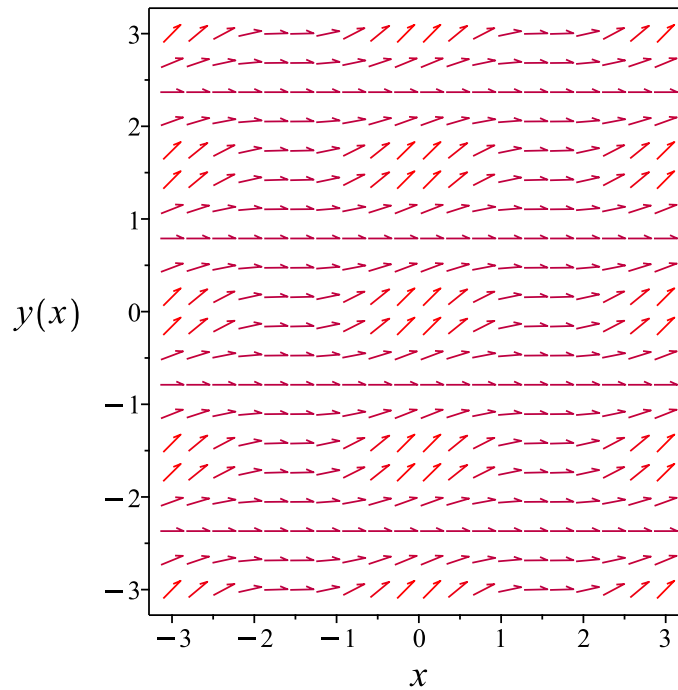


Figure 94: Slope field plot

Verification of solutions

$$-\frac{x}{2} - \frac{\sin(2x)}{4} + \frac{\tan(2y)}{2} = c_1$$

Verified OK.

2.5.4 Maple step by step solution

Let's solve

$$y' - \cos(x)^2 \cos(2y)^2 = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'}{\cos(2y)^2} = \cos(x)^2$$

- Integrate both sides with respect to x

$$\int \frac{y'}{\cos(2y)^2} dx = \int \cos(x)^2 dx + c_1$$

- Evaluate integral

$$\frac{\tan(2y)}{2} = \frac{\cos(x)\sin(x)}{2} + \frac{x}{2} + c_1$$

- Solve for y

$$y = \frac{\arctan(\cos(x)\sin(x)+2c_1+x)}{2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 18

```
dsolve(diff(y(x),x) = cos(x)^2*cos(2*y(x))^2,y(x), singsol=all)
```

$$y(x) = \frac{\arctan\left(x + 2c_1 + \frac{\sin(2x)}{2}\right)}{2}$$

✓ Solution by Mathematica

Time used: 1.312 (sec). Leaf size: 63

```
DSolve[y'[x] == Cos[x]^2*Cos[2*y[x]]^2,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2} \arctan\left(x + \sin(x)\cos(x) + \frac{c_1}{4}\right)$$

$$y(x) \rightarrow \frac{1}{2} \arctan\left(x + \sin(x)\cos(x) + \frac{c_1}{4}\right)$$

$$y(x) \rightarrow -\frac{\pi}{4}$$

$$y(x) \rightarrow \frac{\pi}{4}$$

2.6 problem 6

2.6.1	Solving as separable ode	476
2.6.2	Solving as first order ode lie symmetry lookup ode	478
2.6.3	Solving as exact ode	482
2.6.4	Maple step by step solution	486

Internal problem ID [484]

Internal file name [OUTPUT/484_Sunday_June_05_2022_01_42_13_AM_90091053/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 6.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y'x - \sqrt{1 - y^2} = 0$$

2.6.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{\sqrt{-y^2 + 1}}{x}\end{aligned}$$

Where $f(x) = \frac{1}{x}$ and $g(y) = \sqrt{-y^2 + 1}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\sqrt{-y^2 + 1}} dy &= \frac{1}{x} dx \\ \int \frac{1}{\sqrt{-y^2 + 1}} dy &= \int \frac{1}{x} dx \\ \arcsin(y) &= \ln(x) + c_1\end{aligned}$$

Which results in

$$y = \sin(\ln(x) + c_1)$$

Summary

The solution(s) found are the following

$$y = \sin(\ln(x) + c_1) \tag{1}$$

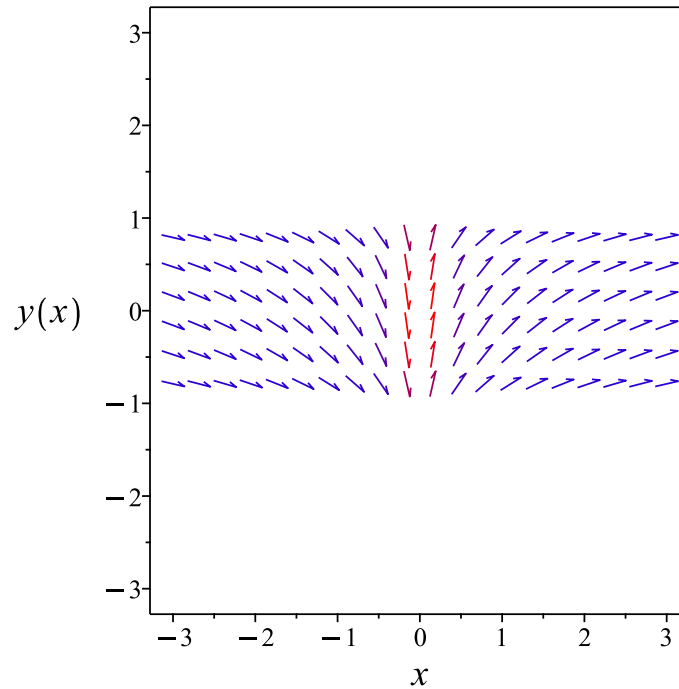


Figure 95: Slope field plot

Verification of solutions

$$y = \sin(\ln(x) + c_1)$$

Verified OK.

2.6.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{\sqrt{-y^2 + 1}}{x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 109: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= x \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{x} dx\end{aligned}$$

Which results in

$$S = \ln(x)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{\sqrt{-y^2 + 1}}{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= \frac{1}{x} \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{\sqrt{-y^2 + 1}} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{\sqrt{-R^2 + 1}}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \arcsin(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\ln(x) = \arcsin(y) + c_1$$

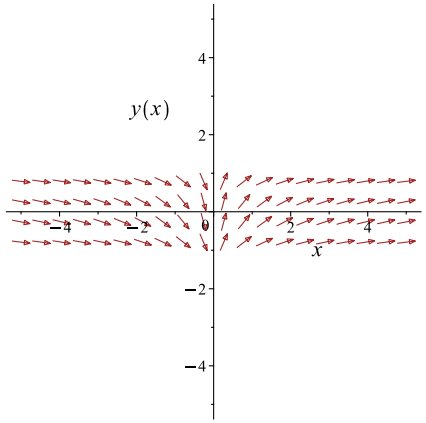
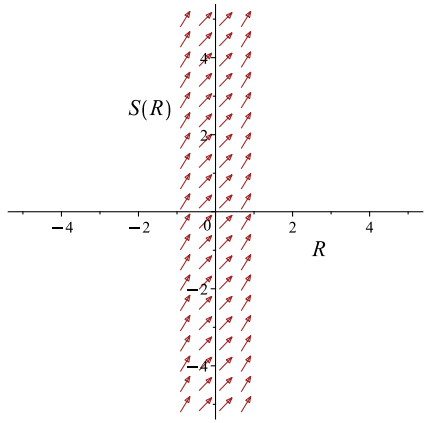
Which simplifies to

$$\ln(x) = \arcsin(y) + c_1$$

Which gives

$$y = -\sin(-\ln(x) + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{\sqrt{-y^2+1}}{x}$ 	$R = y$ $S = \ln(x)$	$\frac{dS}{dR} = \frac{1}{\sqrt{-R^2+1}}$ 

Summary

The solution(s) found are the following

$$y = -\sin(-\ln(x) + c_1) \tag{1}$$

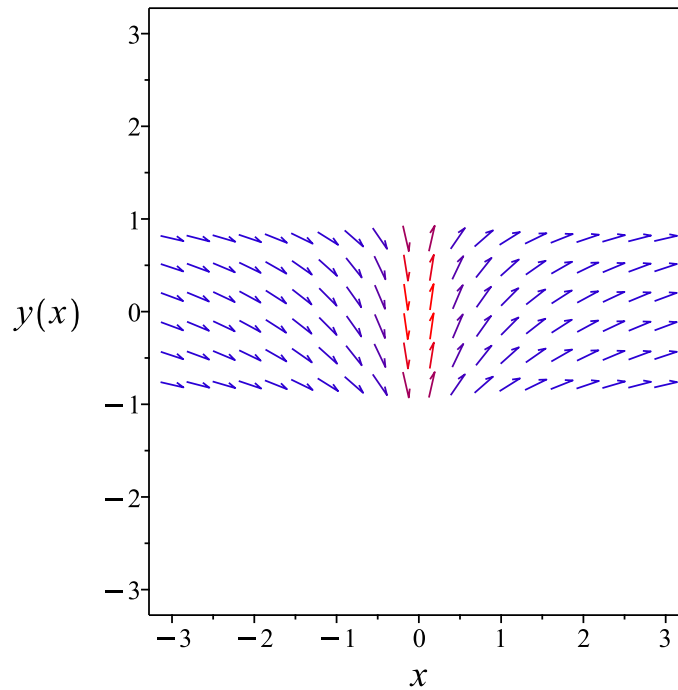


Figure 96: Slope field plot

Verification of solutions

$$y = -\sin(-\ln(x) + c_1)$$

Verified OK.

2.6.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}\left(\frac{1}{\sqrt{-y^2+1}}\right) dy &= \left(\frac{1}{x}\right) dx \\ \left(-\frac{1}{x}\right) dx + \left(\frac{1}{\sqrt{-y^2+1}}\right) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{1}{x} \\ N(x, y) &= \frac{1}{\sqrt{-y^2+1}}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{1}{x}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{1}{\sqrt{-y^2+1}} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{1}{x} dx \\ \phi &= -\ln(x) + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{\sqrt{-y^2+1}}$. Therefore equation (4) becomes

$$\frac{1}{\sqrt{-y^2+1}} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{1}{\sqrt{-y^2+1}}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(\frac{1}{\sqrt{-y^2+1}} \right) dy \\ f(y) &= \arcsin(y) + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\ln(x) + \arcsin(y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\ln(x) + \arcsin(y)$$

The solution becomes

$$y = \sin(\ln(x) + c_1)$$

Summary

The solution(s) found are the following

$$y = \sin(\ln(x) + c_1) \tag{1}$$

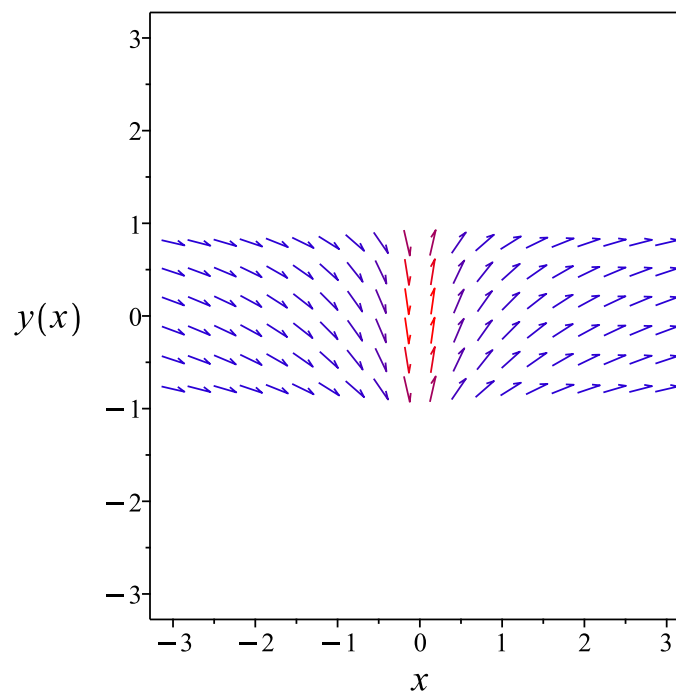


Figure 97: Slope field plot

Verification of solutions

$$y = \sin(\ln(x) + c_1)$$

Verified OK.

2.6.4 Maple step by step solution

Let's solve

$$y'x - \sqrt{1-y^2} = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{\sqrt{1-y^2}} = \frac{1}{x}$$

- Integrate both sides with respect to x

$$\int \frac{y'}{\sqrt{1-y^2}} dx = \int \frac{1}{x} dx + c_1$$

- Evaluate integral

$$\arcsin(y) = \ln(x) + c_1$$

- Solve for y

$$y = \sin(\ln(x) + c_1)$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 9

```
dsolve(x*diff(y(x),x) = (1-y(x)^2)^(1/2),y(x), singsol=all)
```

$$y(x) = \sin(\ln(x) + c_1)$$

✓ Solution by Mathematica

Time used: 0.206 (sec). Leaf size: 29

```
DSolve[x*y'[x] == (1-y[x]^2)^(1/2),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \cos(\log(x) + c_1)$$

$$y(x) \rightarrow -1$$

$$y(x) \rightarrow 1$$

$$y(x) \rightarrow \text{Interval}[\{-1, 1\}]$$

2.7 problem 7

Internal problem ID [485]

Internal file name [OUTPUT/485_Sunday_June_05_2022_01_42_14_AM_30092807/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 7.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**unknown**"

Maple gives the following as the ode type

[`y=_G(x,y')`]

Unable to solve or complete the solution.

$$y' - \frac{-e^{-x} + x}{e^y + x} = 0$$

Unable to determine ODE type.

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
Looking for potential symmetries
trying inverse_Riccati
trying an equivalence to an Abel ODE
differential order: 1; trying a linearization to 2nd order
--- trying a change of variables {x -> y(x), y(x) -> x}
differential order: 1; trying a linearization to 2nd order
trying 1st order ODE linearizable_by_differentiation
--- Trying Lie symmetry methods, 1st order ---
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = 4
`, `-> Computing symmetries using: way = 5
trying symmetry patterns for 1st order ODEs
-> trying a symmetry pattern of the form [F(x)*G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)*G(y)]
-> trying symmetry patterns of the forms [F(x),G(y)] and [G(y),F(x)]
`, `-> Computing symmetries using: way = HINT
-> Calling odsolve with the ODE`, diff(y(x), x) = -y(x)/x, y(x)` *** Sublevel 2 ***
  Methods for first order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  trying 1st order linear
  <- 1st order linear successful
-> Calling odsolve with the ODE`, diff(y(x), x), y(x)` *** Sublevel 2 ***
  Methods for first order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  trying 1st order linear
  <- 1st order linear successful
-> Calling odsolve with the ODE`, diff(y(x), x)+y(x), y(x)` *** Sublevel 2 ***
  Methods for first order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  trying 1st order linear
  <- 1st order linear successful
```

X Solution by Maple

```
dsolve(diff(y(x),x) = (-exp(-x)+x)/(exp(y(x))+x),y(x), singsol=all)
```

No solution found

X Solution by Mathematica

Time used: 0.0 (sec). Leaf size: 0

```
DSolve[y'[x] == (-Exp[-x]+x)/(Exp[y[x]]+x),y[x],x,IncludeSingularSolutions -> True]
```

Not solved

2.8 problem 8

2.8.1	Solving as separable ode	491
2.8.2	Solving as differentialType ode	495
2.8.3	Solving as first order ode lie symmetry lookup ode	499
2.8.4	Solving as exact ode	503
2.8.5	Maple step by step solution	507

Internal problem ID [486]

Internal file name [OUTPUT/486_Sunday_June_05_2022_01_42_16_AM_5493618/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 8.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "differential-Type", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{x^2}{1+y^2} = 0$$

2.8.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{x^2}{y^2 + 1}\end{aligned}$$

Where $f(x) = x^2$ and $g(y) = \frac{1}{y^2+1}$. Integrating both sides gives

$$\frac{1}{y^2+1} dy = x^2 dx$$

$$\int \frac{1}{y^2+1} dy = \int x^2 dx$$

$$\frac{1}{3}y^3 + y = \frac{x^3}{3} + c_1$$

Which results in

$$y = \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} - \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}$$

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{1}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} + \frac{i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} \right)}{2}$$

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{1}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} - \frac{i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} \right)}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} \quad (1)$$

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{1} \quad (2)$$

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} \right)}{2} \quad (3)$$

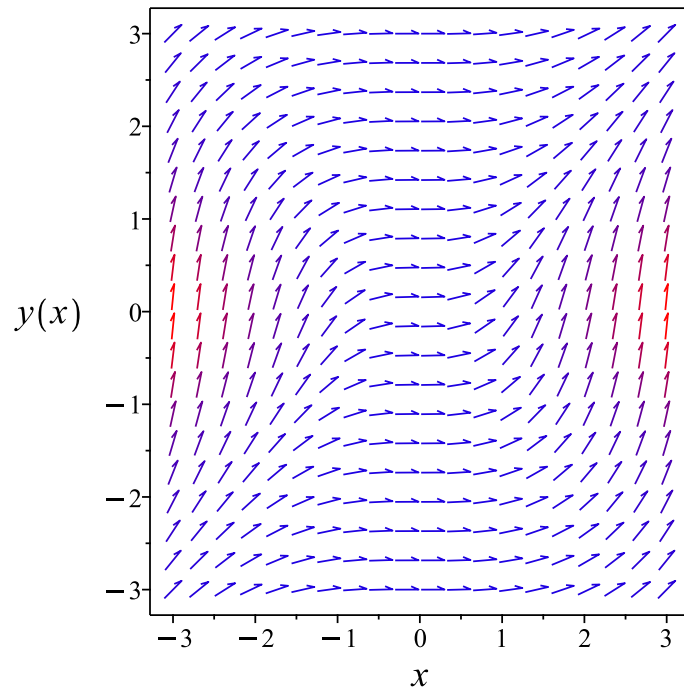


Figure 98: Slope field plot

Verification of solutions

$$y = \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} - \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}$$

Verified OK.

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{1}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} + i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} \right)$$

Verified OK.

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{1}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} - i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} \right)$$

Verified OK.

2.8.2 Solving as differential Type ode

Writing the ode as

$$y' = \frac{x^2}{1 + y^2} \tag{1}$$

Which becomes

$$(y^2 + 1) dy = (x^2) dx \tag{2}$$

But the RHS is complete differential because

$$(x^2) dx = d\left(\frac{x^3}{3}\right)$$

Hence (2) becomes

$$(y^2 + 1) dy = d\left(\frac{x^3}{3}\right)$$

Integrating both sides gives gives these solutions

$$y = \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} - \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} + c_1$$

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{1}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} + \frac{i\sqrt{3}}{4} \left(\frac{4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}\right)$$

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{1}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} - \frac{i\sqrt{3}}{4} \left(\frac{4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}\right)$$

Summary

The solution(s) found are the following

$$y = \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} - \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + c_1 \quad (1)$$

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{1} \quad (2)$$

$$+ \frac{i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} \right)}{2} + c_1$$

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{1} \quad (3)$$

$$- \frac{i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} \right)}{2} + c_1$$

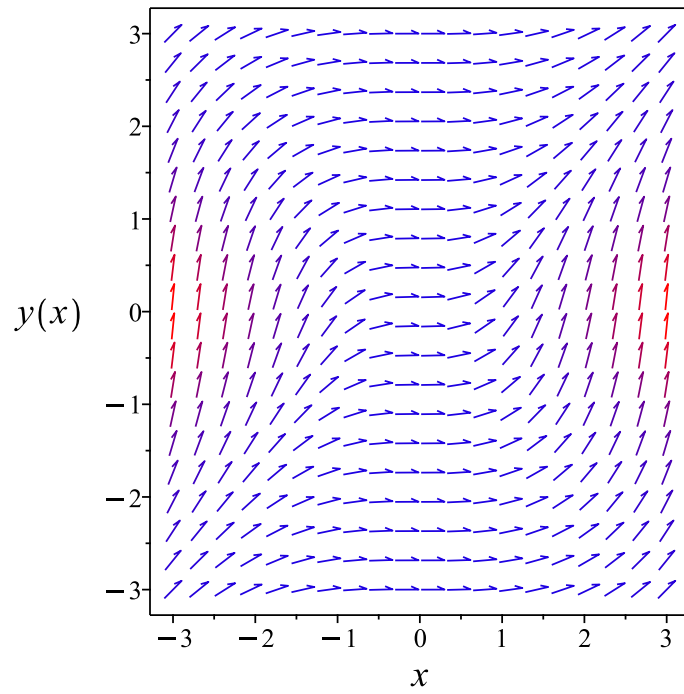


Figure 99: Slope field plot

Verification of solutions

$$y = \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} - \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + c_1$$

Verified OK.

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{1} + i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} \right) + c_1$$

Verified OK.

$$y = -\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{1} + i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}} \right) - \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} + c_1$$

Verified OK.

2.8.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x^2}{y^2 + 1}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 112: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned} \xi(x, y) &= \frac{1}{x^2} \\ \eta(x, y) &= 0 \end{aligned} \quad (\text{A1})$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{x^2} dx \end{aligned}$$

Which results in

$$S = \frac{x^3}{3}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x^2}{y^2 + 1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= x^2 \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y^2 + 1 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R^2 + 1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{1}{3}R^3 + R + c_1 \quad (4)$$

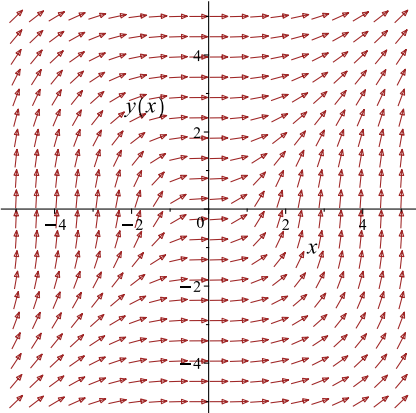
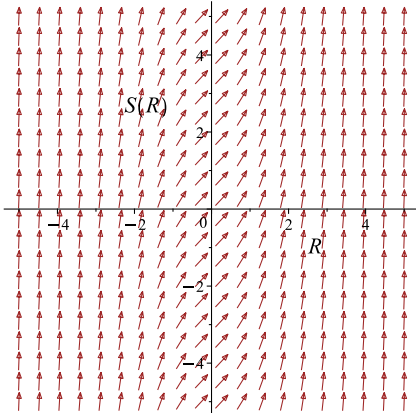
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{x^3}{3} = \frac{y^3}{3} + y + c_1$$

Which simplifies to

$$\frac{x^3}{3} = \frac{y^3}{3} + y + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{x^2}{y^2+1}$ 	$R = y$ $S = \frac{x^3}{3}$	$\frac{dS}{dR} = R^2 + 1$ 

Summary

The solution(s) found are the following

$$\frac{x^3}{3} = \frac{y^3}{3} + y + c_1 \quad (1)$$

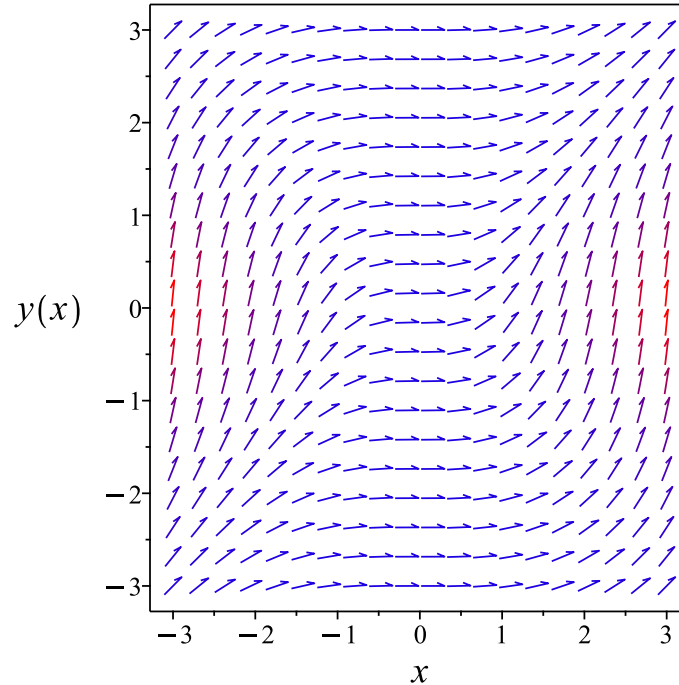


Figure 100: Slope field plot

Verification of solutions

$$\frac{x^3}{3} = \frac{y^3}{3} + y + c_1$$

Verified OK.

2.8.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (y^2 + 1) dy &= (x^2) dx \\ (-x^2) dx + (y^2 + 1) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x^2 \\ N(x, y) &= y^2 + 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x^2) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(y^2 + 1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x^2 dx \\ \phi &= -\frac{x^3}{3} + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = y^2 + 1$. Therefore equation (4) becomes

$$y^2 + 1 = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = y^2 + 1$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (y^2 + 1) dy \\ f(y) &= \frac{1}{3}y^3 + y + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{1}{3}x^3 + \frac{1}{3}y^3 + y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{1}{3}x^3 + \frac{1}{3}y^3 + y$$

Summary

The solution(s) found are the following

$$-\frac{x^3}{3} + \frac{y^3}{3} + y = c_1 \tag{1}$$

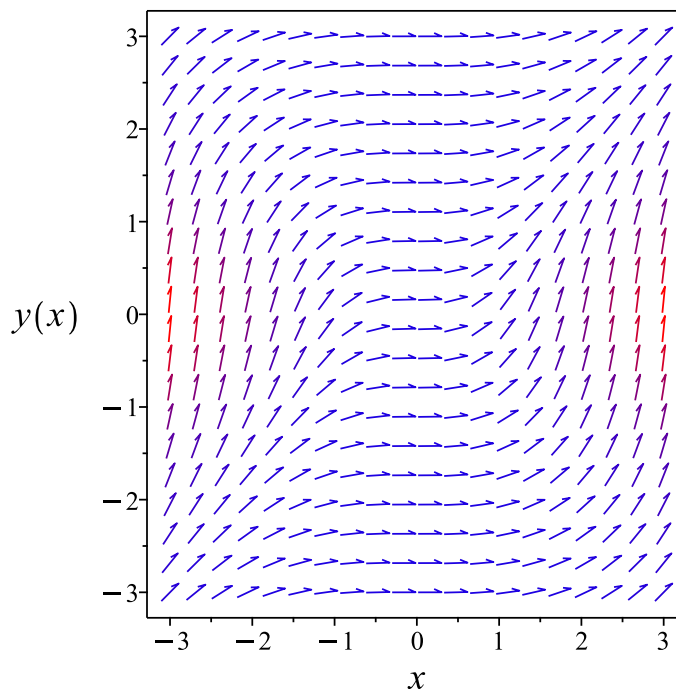


Figure 101: Slope field plot

Verification of solutions

$$-\frac{x^3}{3} + \frac{y^3}{3} + y = c_1$$

Verified OK.

2.8.5 Maple step by step solution

Let's solve

$$y' - \frac{x^2}{1+y^2} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$(1 + y^2) y' = x^2$$

- Integrate both sides with respect to x

$$\int (1 + y^2) y' dx = \int x^2 dx + c_1$$

- Evaluate integral

$$\frac{y^3}{3} + y = \frac{x^3}{3} + c_1$$

- Solve for y

$$y = \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}{2} - \frac{2}{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 268

```
dsolve(diff(y(x),x) = x^2/(1+y(x)^2),y(x), singsol=all)
```

$$y(x) = \frac{\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{2}{3}} - 4}{2\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}$$

$$y(x) = -\frac{(1 + i\sqrt{3})\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{2}{3}} + 4i\sqrt{3} - 4}{4\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}$$

$$y(x) = \frac{i\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{2}{3}}\sqrt{3} + 4i\sqrt{3} - \left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{2}{3}} + 4}{4\left(4x^3 + 12c_1 + 4\sqrt{x^6 + 6c_1x^3 + 9c_1^2 + 4}\right)^{\frac{1}{3}}}$$

✓ Solution by Mathematica

Time used: 2.246 (sec). Leaf size: 307

```
DSolve[y'[x]== x^2/(1+y[x]^2),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{-2 + \sqrt[3]{2}(x^3 + \sqrt{x^6 + 6c_1x^3 + 4 + 9c_1^2 + 3c_1})^{2/3}}{2^{2/3}\sqrt[3]{x^3 + \sqrt{x^6 + 6c_1x^3 + 4 + 9c_1^2 + 3c_1}}}$$

$$y(x) \rightarrow \frac{i(\sqrt{3} + i)\sqrt[3]{x^3 + \sqrt{x^6 + 6c_1x^3 + 4 + 9c_1^2 + 3c_1}}}{2\sqrt[3]{2}} + \frac{1 + i\sqrt{3}}{2^{2/3}\sqrt[3]{x^3 + \sqrt{x^6 + 6c_1x^3 + 4 + 9c_1^2 + 3c_1}}}$$

$$y(x) \rightarrow \frac{1 - i\sqrt{3}}{2^{2/3}\sqrt[3]{x^3 + \sqrt{x^6 + 6c_1x^3 + 4 + 9c_1^2 + 3c_1}}} - \frac{(1 + i\sqrt{3})\sqrt[3]{x^3 + \sqrt{x^6 + 6c_1x^3 + 4 + 9c_1^2 + 3c_1}}}{2\sqrt[3]{2}}$$

2.9 problem 9

2.9.1	Existence and uniqueness analysis	509
2.9.2	Solving as separable ode	510
2.9.3	Solving as first order ode lie symmetry lookup ode	512
2.9.4	Solving as exact ode	516
2.9.5	Solving as riccati ode	520
2.9.6	Maple step by step solution	523

Internal problem ID [487]

Internal file name [OUTPUT/487_Sunday_June_05_2022_01_42_17_AM_24604466/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 9.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - (1 - 2x)y^2 = 0$$

With initial conditions

$$\left[y(0) = -\frac{1}{6} \right]$$

2.9.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned} y' &= f(x, y) \\ &= -y^2(2x - 1) \end{aligned}$$

The x domain of $f(x, y)$ when $y = -\frac{1}{6}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = -\frac{1}{6}$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y}(-y^2(2x - 1)) \\ &= -2(2x - 1)y\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = -\frac{1}{6}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = -\frac{1}{6}$ is inside this domain. Therefore solution exists and is unique.

2.9.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= (1 - 2x)y^2\end{aligned}$$

Where $f(x) = 1 - 2x$ and $g(y) = y^2$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y^2} dy &= 1 - 2x dx \\ \int \frac{1}{y^2} dy &= \int 1 - 2x dx \\ -\frac{1}{y} &= -x^2 + c_1 + x\end{aligned}$$

Which results in

$$y = -\frac{1}{-x^2 + c_1 + x}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -\frac{1}{6}$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{1}{6} = -\frac{1}{c_1}$$

$$c_1 = 6$$

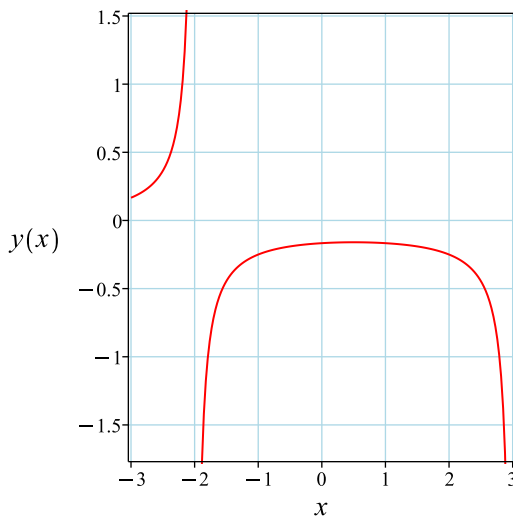
Substituting c_1 found above in the general solution gives

$$y = \frac{1}{x^2 - x - 6}$$

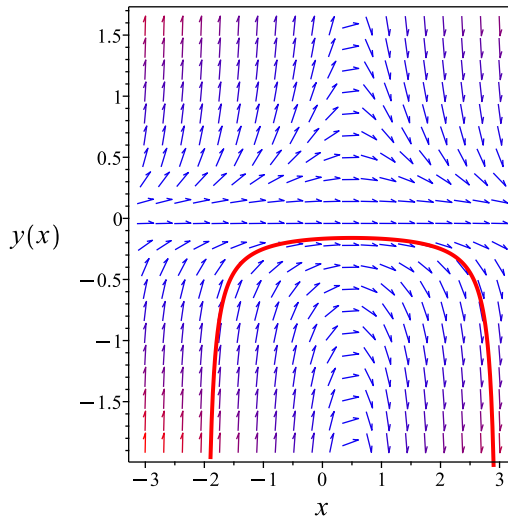
Summary

The solution(s) found are the following

$$y = \frac{1}{x^2 - x - 6} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{1}{x^2 - x - 6}$$

Verified OK.

2.9.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -y^2(2x - 1)$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 115: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{1 - 2x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{1-2x}} dx\end{aligned}$$

Which results in

$$S = -x^2 + x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -y^2(2x - 1)$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\ R_y &= 1 \\ S_x &= 1 - 2x \\ S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{y^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{1}{R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-x^2 + x = -\frac{1}{y} + c_1$$

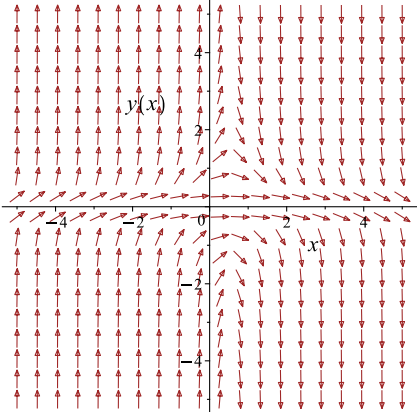
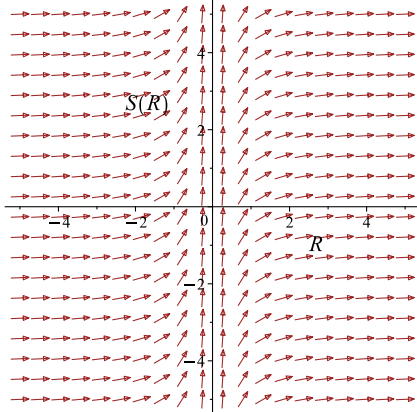
Which simplifies to

$$-x^2 + x = -\frac{1}{y} + c_1$$

Which gives

$$y = \frac{1}{x^2 + c_1 - x}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -y^2(2x - 1)$ 	$R = y$ $S = -x^2 + x$	$\frac{dS}{dR} = \frac{1}{R^2}$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -\frac{1}{6}$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{1}{6} = \frac{1}{c_1}$$

$$c_1 = -6$$

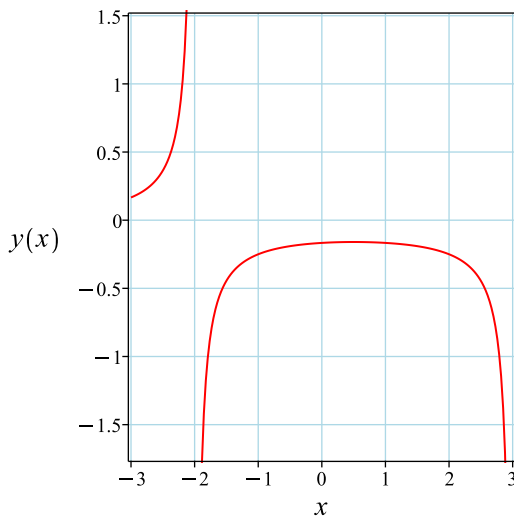
Substituting c_1 found above in the general solution gives

$$y = \frac{1}{x^2 - x - 6}$$

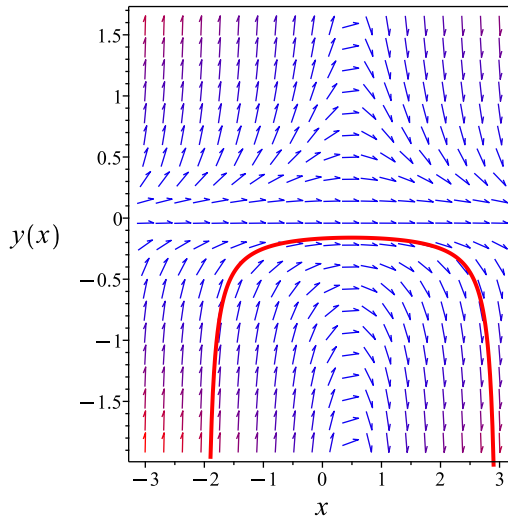
Summary

The solution(s) found are the following

$$y = \frac{1}{x^2 - x - 6} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{1}{x^2 - x - 6}$$

Verified OK.

2.9.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left(-\frac{1}{y^2}\right) dy &= (2x - 1) dx \\ (1 - 2x) dx + \left(-\frac{1}{y^2}\right) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 1 - 2x \\ N(x, y) &= -\frac{1}{y^2} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(1 - 2x) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}\left(-\frac{1}{y^2}\right) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 1 - 2x dx \\ \phi &= -x^2 + x + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{1}{y^2}$. Therefore equation (4) becomes

$$-\frac{1}{y^2} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{1}{y^2}$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int \left(-\frac{1}{y^2}\right) dy \\ f(y) &= \frac{1}{y} + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x^2 + x + \frac{1}{y} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x^2 + x + \frac{1}{y}$$

The solution becomes

$$y = \frac{1}{x^2 + c_1 - x}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -\frac{1}{6}$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{1}{6} = \frac{1}{c_1}$$

$$c_1 = -6$$

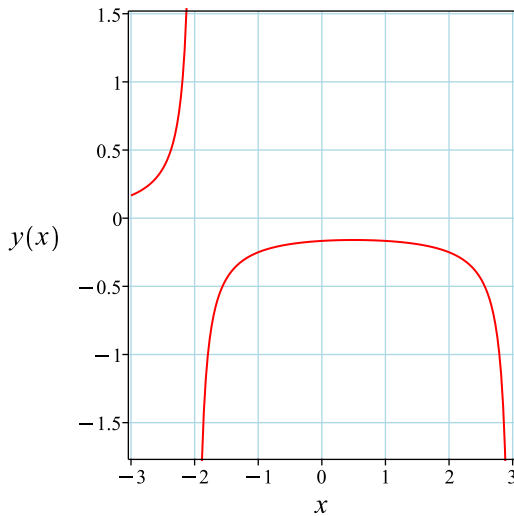
Substituting c_1 found above in the general solution gives

$$y = \frac{1}{x^2 - x - 6}$$

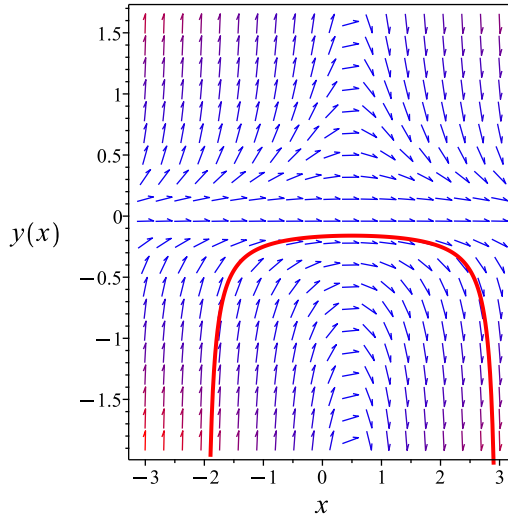
Summary

The solution(s) found are the following

$$y = \frac{1}{x^2 - x - 6} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{1}{x^2 - x - 6}$$

Verified OK.

2.9.5 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= -y^2(2x - 1) \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = -2xy^2 + y^2$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that $f_0(x) = 0$, $f_1(x) = 0$ and $f_2(x) = 1 - 2x$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{(1 - 2x)u} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \quad (2)$$

But

$$\begin{aligned} f_2' &= -2 \\ f_1 f_2 &= 0 \\ f_2^2 f_0 &= 0 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$(1 - 2x) u''(x) + 2u'(x) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 + c_2 \left(x - \frac{1}{2} \right)^2$$

The above shows that

$$u'(x) = c_2(2x - 1)$$

Using the above in (1) gives the solution

$$y = -\frac{c_2(2x - 1)}{(1 - 2x) \left(c_1 + c_2 \left(x - \frac{1}{2} \right)^2 \right)}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = \frac{4}{4x^2 + 4c_3 - 4x + 1}$$

Initial conditions are used to solve for c_3 . Substituting $x = 0$ and $y = -\frac{1}{6}$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{1}{6} = \frac{4}{4c_3 + 1}$$

$$c_3 = -\frac{25}{4}$$

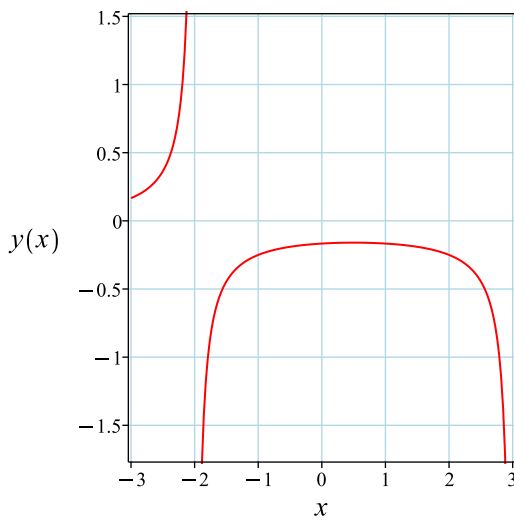
Substituting c_3 found above in the general solution gives

$$y = \frac{1}{x^2 - x - 6}$$

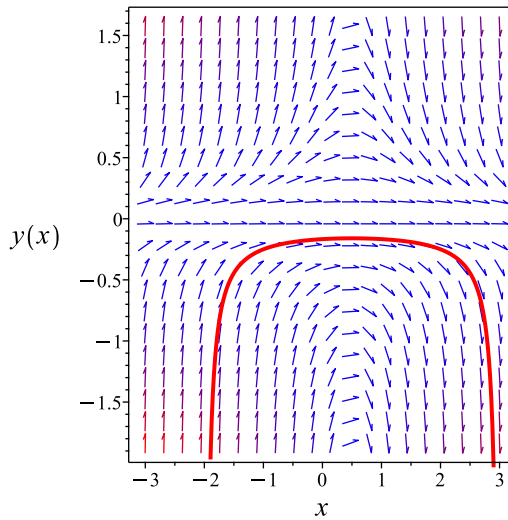
Summary

The solution(s) found are the following

$$y = \frac{1}{x^2 - x - 6} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{1}{x^2 - x - 6}$$

Verified OK.

2.9.6 Maple step by step solution

Let's solve

$$[y' - (1 - 2x)y^2 = 0, y(0) = -\frac{1}{6}]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{y^2} = 1 - 2x$$

- Integrate both sides with respect to x

$$\int \frac{y'}{y^2} dx = \int (1 - 2x) dx + c_1$$

- Evaluate integral

$$-\frac{1}{y} = -x^2 + c_1 + x$$

- Solve for y

$$y = -\frac{1}{-x^2 + c_1 + x}$$

- Use initial condition $y(0) = -\frac{1}{6}$

$$-\frac{1}{6} = -\frac{1}{c_1}$$

- Solve for c_1

$$c_1 = 6$$

- Substitute $c_1 = 6$ into general solution and simplify

$$y = \frac{1}{x^2 - x - 6}$$

- Solution to the IVP

$$y = \frac{1}{x^2 - x - 6}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.062 (sec). Leaf size: 14

```
dsolve([diff(y(x),x) = (1-2*x)*y(x)^2,y(0) = -1/6],y(x), singsol=all)
```

$$y(x) = \frac{1}{x^2 - x - 6}$$

✓ Solution by Mathematica

Time used: 0.131 (sec). Leaf size: 15

```
DSolve[{y'[x] == (1-2*x)*y[x]^2,y[0]==-1/6},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{x^2 - x - 6}$$

2.10 problem 10

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Internal problem ID [488]

Internal file name [OUTPUT/488_Sunday_June_05_2022_01_42_18_AM_32830898/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 10.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "homogeneousTypeMapleC", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{1 - 2x}{y} = 0$$

With initial conditions

$$[y(1) = -2]$$

2.10.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= -\frac{2x-1}{y}\end{aligned}$$

The x domain of $f(x, y)$ when $y = -2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 1$ is inside this domain. The y domain of $f(x, y)$ when $x = 1$ is

$$\{y < 0 \vee 0 < y\}$$

And the point $y_0 = -2$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{2x-1}{y} \right) \\ &= \frac{2x-1}{y^2}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = -2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 1$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 1$ is

$$\{y < 0 \vee 0 < y\}$$

And the point $y_0 = -2$ is inside this domain. Therefore solution exists and is unique.

2.10.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{1-2x}{y}\end{aligned}$$

Where $f(x) = 1 - 2x$ and $g(y) = \frac{1}{y}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y} dy &= 1 - 2x dx \\ \int \frac{1}{y} dy &= \int 1 - 2x dx \\ \frac{y^2}{2} &= -x^2 + c_1 + x\end{aligned}$$

Which results in

$$\begin{aligned}y &= \sqrt{-2x^2 + 2c_1 + 2x} \\ y &= -\sqrt{-2x^2 + 2c_1 + 2x}\end{aligned}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = -\sqrt{c_1} \sqrt{2}$$

$$c_1 = 2$$

Substituting c_1 found above in the general solution gives

$$y = -\sqrt{-2x^2 + 2x + 4}$$

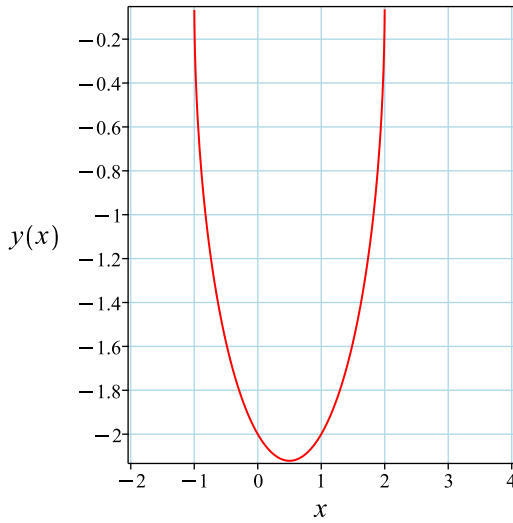
Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = \sqrt{c_1} \sqrt{2}$$

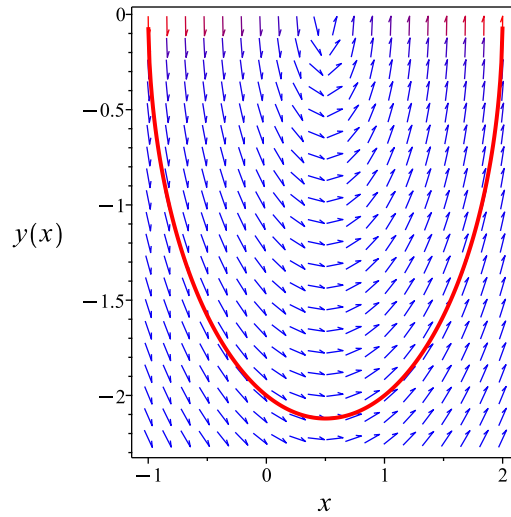
Summary

Warning: Unable to solve for constant of integration. The solution(s) found are the following

$$y = -\sqrt{-2x^2 + 2x + 4}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\sqrt{-2x^2 + 2x + 4}$$

Verified OK.

2.10.3 Solving as differentialType ode

Writing the ode as

$$y' = \frac{1 - 2x}{y} \tag{1}$$

Which becomes

$$(y) dy = (1 - 2x) dx \tag{2}$$

But the RHS is complete differential because

$$(1 - 2x) dx = d(-x^2 + x)$$

Hence (2) becomes

$$(y) dy = d(-x^2 + x)$$

Integrating both sides gives these solutions

$$y = \sqrt{-2x^2 + 2c_1 + 2x + c_1}$$

$$y = -\sqrt{-2x^2 + 2c_1 + 2x + c_1}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = -\sqrt{c_1} \sqrt{2} + c_1$$

$$c_1 = \sqrt{2} \left(\frac{\sqrt{2}}{2} + \frac{i\sqrt{6}}{2} \right) - 2$$

Substituting c_1 found above in the general solution gives

$$y = -\sqrt{-2x^2 - 2 + 2i\sqrt{3} + 2x - 1 + i\sqrt{3}}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = \sqrt{c_1} \sqrt{2} + c_1$$

Summary

The solution(s) found are the following

Warning: Unable to solve for constant of integration.

$$y = -\sqrt{-2x^2 - 2 + 2i\sqrt{3} + 2x - 1 + i\sqrt{3}}$$

Verification of solutions

$$y = -\sqrt{-2x^2 - 2 + 2i\sqrt{3} + 2x - 1 + i\sqrt{3}}$$

Verified OK.

2.10.4 Solving as homogeneousTypeMapleC ode

Let $Y = y + y_0$ and $X = x + x_0$ then the above is transformed to new ode in $Y(X)$

$$\frac{d}{dX}Y(X) = -\frac{2X + 2x_0 - 1}{Y(X) + y_0}$$

Solving for possible values of x_0 and y_0 which makes the above ode a homogeneous ode results in

$$x_0 = \frac{1}{2}$$

$$y_0 = 0$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$\frac{d}{dX}Y(X) = -\frac{2X}{Y(X)}$$

In canonical form, the ODE is

$$\begin{aligned} Y' &= F(X, Y) \\ &= -\frac{2X}{Y} \end{aligned} \tag{1}$$

An ode of the form $Y' = \frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order n if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both $M = -2X$ and $N = Y$ are both homogeneous and of the same order $n = 1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{Y}{X}$, or $Y = uX$. Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation $Y = uX$ to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dX}X + u &= -\frac{2}{u} \\ \frac{du}{dX} &= \frac{-\frac{2}{u(X)} - u(X)}{X} \end{aligned}$$

Or

$$\frac{d}{dX}u(X) - \frac{-\frac{2}{u(X)} - u(X)}{X} = 0$$

Or

$$\left(\frac{d}{dX}u(X) \right) u(X) X + u(X)^2 + 2 = 0$$

Which is now solved as separable in $u(X)$. Which is now solved in $u(X)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(X, u) \\ &= f(X)g(u) \\ &= -\frac{u^2 + 2}{uX} \end{aligned}$$

Where $f(X) = -\frac{1}{X}$ and $g(u) = \frac{u^2+2}{u}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{u^2+2}{u}} du &= -\frac{1}{X} dX \\ \int \frac{1}{\frac{u^2+2}{u}} du &= \int -\frac{1}{X} dX \\ \frac{\ln(u^2 + 2)}{2} &= -\ln(X) + c_2\end{aligned}$$

Raising both side to exponential gives

$$\sqrt{u^2 + 2} = e^{-\ln(X)+c_2}$$

Which simplifies to

$$\sqrt{u^2 + 2} = \frac{c_3}{X}$$

Which simplifies to

$$\sqrt{u(X)^2 + 2} = \frac{c_3 e^{c_2}}{X}$$

The solution is

$$\sqrt{u(X)^2 + 2} = \frac{c_3 e^{c_2}}{X}$$

Now u in the above solution is replaced back by Y using $u = \frac{Y}{X}$ which results in the solution

$$\sqrt{\frac{Y(X)^2}{X^2} + 2} = \frac{c_3 e^{c_2}}{X}$$

Using the solution for $Y(X)$

$$\sqrt{\frac{Y(X)^2 + 2X^2}{X^2}} = \frac{c_3 e^{c_2}}{X}$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y$$
$$X = x + \frac{1}{2}$$

Then the solution in y becomes

$$\sqrt{\frac{y^2 + 2\left(x - \frac{1}{2}\right)^2}{\left(x - \frac{1}{2}\right)^2}} = \frac{c_3 e^{c_2}}{x - \frac{1}{2}}$$

Initial conditions are used to solve for c_2 . Substituting $x = 1$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$3\sqrt{2} = 2c_3 e^{c_2}$$

$$c_2 = \frac{\ln\left(\frac{9}{2c_3^2}\right)}{2}$$

Substituting c_2 found above in the general solution gives

$$\sqrt{\frac{y^2 + 2\left(x - \frac{1}{2}\right)^2}{\left(x - \frac{1}{2}\right)^2}} = \frac{3c_3\sqrt{2}\sqrt{\frac{1}{c_3^2}}}{2x - 1}$$

The above simplifies to

$$\sqrt{2}\left(2\sqrt{\frac{4x^2 + 2y^2 - 4x + 1}{(2x - 1)^2}}x - 3c_3\sqrt{\frac{1}{c_3^2}} - \sqrt{\frac{4x^2 + 2y^2 - 4x + 1}{(2x - 1)^2}}\right) = 0$$

Summary

The solution(s) found are the following

$$\sqrt{2}\left(2\sqrt{\frac{4x^2 + 2y^2 - 4x + 1}{(2x - 1)^2}}x - 3 \operatorname{csgn}\left(\frac{1}{c_3}\right) - \sqrt{\frac{4x^2 + 2y^2 - 4x + 1}{(2x - 1)^2}}\right) = 0 \quad (1)$$

Verification of solutions

$$\sqrt{2}\left(2\sqrt{\frac{4x^2 + 2y^2 - 4x + 1}{(2x - 1)^2}}x - 3 \operatorname{csgn}\left(\frac{1}{c_3}\right) - \sqrt{\frac{4x^2 + 2y^2 - 4x + 1}{(2x - 1)^2}}\right) = 0$$

Verified OK.

2.10.5 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{2x-1}{y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 118: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{1 - 2x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{1-2x}} dx\end{aligned}$$

Which results in

$$S = -x^2 + x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{2x - 1}{y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= 1 - 2x \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^2}{2} + c_1 \tag{4}$$

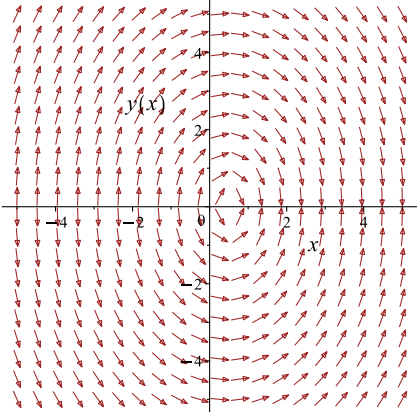
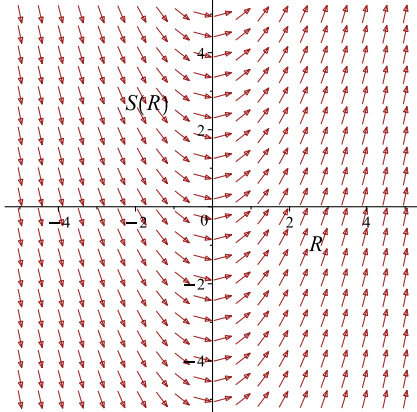
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-x^2 + x = \frac{y^2}{2} + c_1$$

Which simplifies to

$$-x^2 + x = \frac{y^2}{2} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{2x-1}{y}$ 	$R = y$ $S = -x^2 + x$	$\frac{dS}{dR} = R$ 

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$0 = c_1 + 2$$

$$c_1 = -2$$

Substituting c_1 found above in the general solution gives

$$-x^2 + x = \frac{y^2}{2} - 2$$

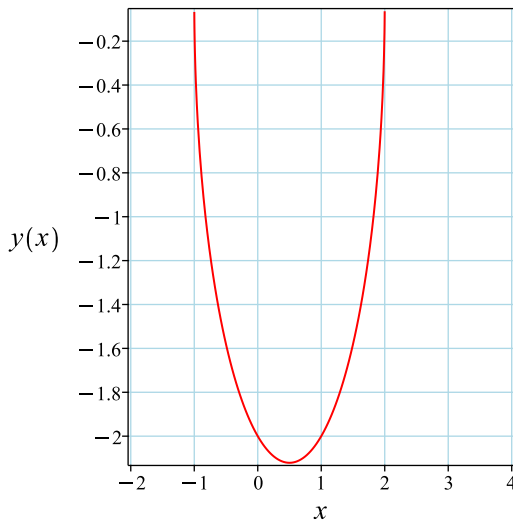
Solving for y from the above gives

$$y = -\sqrt{-2x^2 + 2x + 4}$$

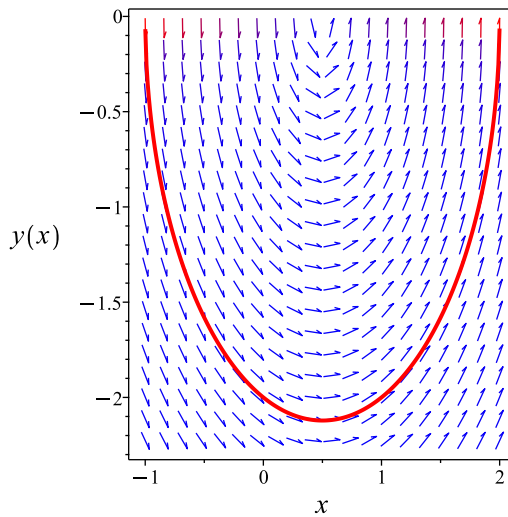
Summary

The solution(s) found are the following

$$y = -\sqrt{-2x^2 + 2x + 4} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\sqrt{-2x^2 + 2x + 4}$$

Verified OK. {positive}

2.10.6 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (-y) dy &= (2x - 1) dx \\ (1 - 2x) dx + (-y) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 1 - 2x \\ N(x, y) &= -y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(1 - 2x) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-y) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 1 - 2x dx \\ \phi &= -x^2 + x + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y)\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -y$. Therefore equation (4) becomes

$$-y = 0 + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -y$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (-y) dy \\ f(y) &= -\frac{y^2}{2} + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x^2 - \frac{1}{2}y^2 + x + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x^2 - \frac{1}{2}y^2 + x$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = c_1$$

$$c_1 = -2$$

Substituting c_1 found above in the general solution gives

$$-x^2 - \frac{1}{2}y^2 + x = -2$$

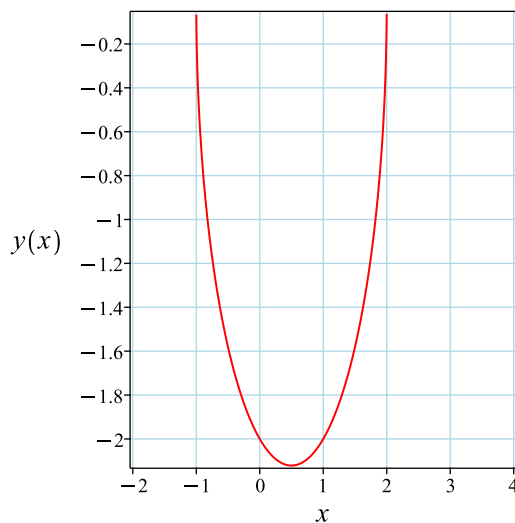
Solving for y from the above gives

$$y = -\sqrt{-2x^2 + 2x + 4}$$

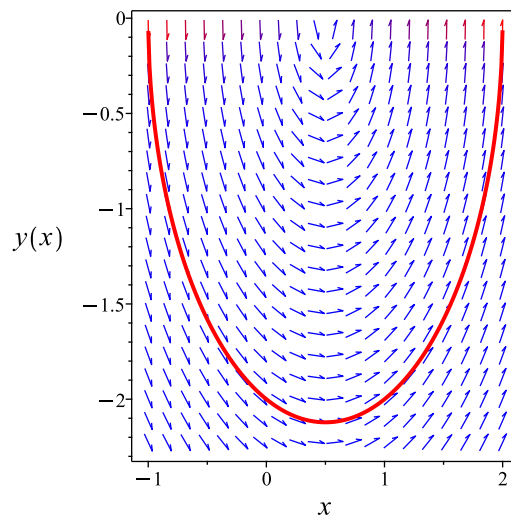
Summary

The solution(s) found are the following

$$y = -\sqrt{-2x^2 + 2x + 4} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\sqrt{-2x^2 + 2x + 4}$$

Verified OK. {positive}

2.10.7 Maple step by step solution

Let's solve

$$\left[y' - \frac{1-2x}{y} = 0, y(1) = -2 \right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$yy' = 1 - 2x$$

- Integrate both sides with respect to x

$$\int yy' dx = \int (1 - 2x) dx + c_1$$

- Evaluate integral

$$\frac{y^2}{2} = -x^2 + c_1 + x$$

- Solve for y

$$\{y = \sqrt{-2x^2 + 2c_1 + 2x}, y = -\sqrt{-2x^2 + 2c_1 + 2x}\}$$

- Use initial condition $y(1) = -2$

$$-2 = \sqrt{c_1} \sqrt{2}$$

- Solution does not satisfy initial condition

- Use initial condition $y(1) = -2$

$$-2 = -\sqrt{c_1} \sqrt{2}$$

- Solve for c_1

$$c_1 = 2$$

- Substitute $c_1 = 2$ into general solution and simplify

$$y = -\sqrt{-2x^2 + 2x + 4}$$

- Solution to the IVP

$$y = -\sqrt{-2x^2 + 2x + 4}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 18

```
dsolve([diff(y(x),x) = (1-2*x)/y(x),y(1) = -2],y(x), singsol=all)
```

$$y(x) = -\sqrt{-2x^2 + 2x + 4}$$

✓ Solution by Mathematica

Time used: 0.094 (sec). Leaf size: 24

```
DSolve[{y'[x] == (1-2*x)/y[x],y[1]==-2},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\sqrt{2}\sqrt{-x^2 + x + 2}$$

2.11 problem 11

2.11.1 Existence and uniqueness analysis	543
2.11.2 Solving as separable ode	544
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2.11.4 Solving as exact ode	551
2.11.5 Maple step by step solution	554

Internal problem ID [489]

Internal file name [OUTPUT/489_Sunday_June_05_2022_01_42_19_AM_46023698/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 11.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y'e^{-x}y = -x$$

With initial conditions

$$[y(0) = 1]$$

2.11.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$y' = f(x, y) \\ = -\frac{x e^x}{y}$$

The x domain of $f(x, y)$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{x e^x}{y} \right) \\ &= \frac{x e^x}{y^2}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Therefore solution exists and is unique.

2.11.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\frac{x e^x}{y}\end{aligned}$$

Where $f(x) = -x e^x$ and $g(y) = \frac{1}{y}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y} dy &= -x e^x dx \\ \int \frac{1}{y} dy &= \int -x e^x dx \\ \frac{y^2}{2} &= -(x-1)e^x + c_1\end{aligned}$$

Which results in

$$y = \sqrt{-2x e^x + 2 e^x + 2c_1}$$

$$y = -\sqrt{-2x e^x + 2 e^x + 2c_1}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\sqrt{2 + 2c_1}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \sqrt{2 + 2c_1}$$

$$c_1 = -\frac{1}{2}$$

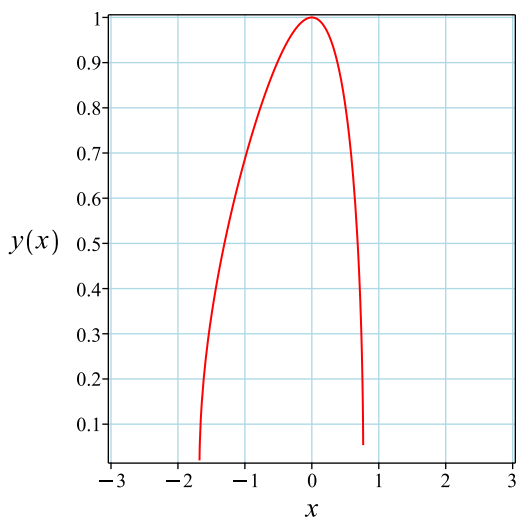
Substituting c_1 found above in the general solution gives

$$y = \sqrt{-2x e^x + 2 e^x - 1}$$

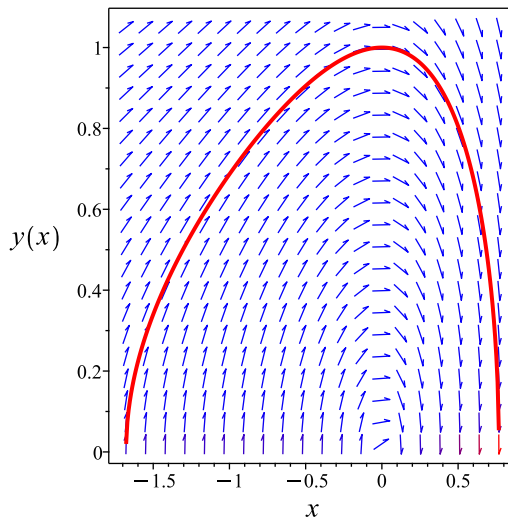
Summary

The solution(s) found are the following

$$y = \sqrt{-2x e^x + 2 e^x - 1} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \sqrt{-2x e^x + 2e^x - 1}$$

Verified OK.

2.11.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{x e^x}{y}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 121: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= -\frac{e^{-x}}{x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{-\frac{e^{-x}}{x}} dx \end{aligned}$$

Which results in

$$S = -(x - 1) e^x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{x e^x}{y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= -x e^x \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^2}{2} + c_1 \quad (4)$$

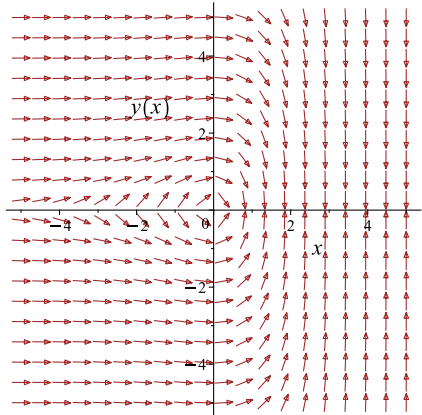
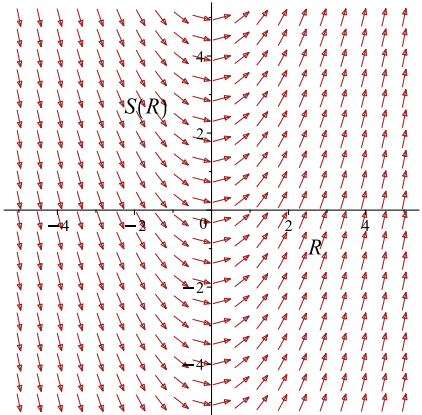
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-(x - 1) e^x = \frac{y^2}{2} + c_1$$

Which simplifies to

$$-(x - 1) e^x = \frac{y^2}{2} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{x e^x}{y}$ 	$R = y$ $S = -(x - 1) e^x$	$\frac{dS}{dR} = R$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{1}{2} + c_1$$

$$c_1 = \frac{1}{2}$$

Substituting c_1 found above in the general solution gives

$$-(x - 1) e^x = \frac{y^2}{2} + \frac{1}{2}$$

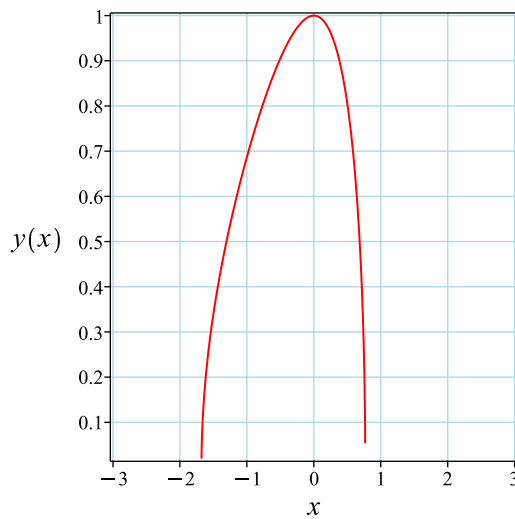
Solving for y from the above gives

$$y = \sqrt{-2x e^x + 2 e^x - 1}$$

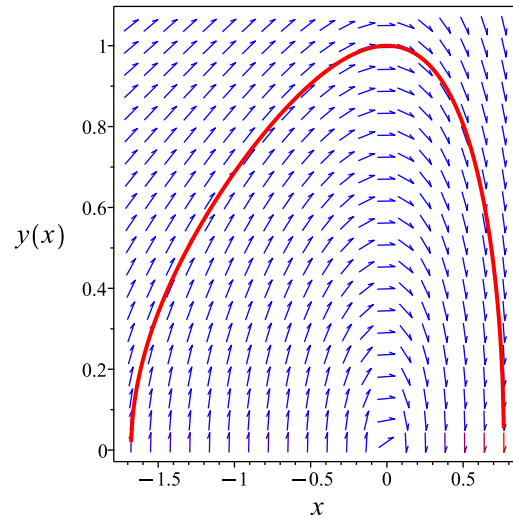
Summary

The solution(s) found are the following

$$y = \sqrt{-2x e^x + 2 e^x - 1} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \sqrt{-2x e^x + 2 e^x - 1}$$

Verified OK.

2.11.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (-y) dy &= (x e^x) dx \\ (-x e^x) dx + (-y) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x e^x \\ N(x, y) &= -y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x e^x) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-y) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x e^x dx \\ \phi &= -(x - 1) e^x + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -y$. Therefore equation (4) becomes

$$-y = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -y$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (-y) dy$$

$$f(y) = -\frac{y^2}{2} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -(x-1)e^x - \frac{y^2}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -(x-1)e^x - \frac{y^2}{2}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{1}{2} = c_1$$

$$c_1 = \frac{1}{2}$$

Substituting c_1 found above in the general solution gives

$$-(x-1)e^x - \frac{y^2}{2} = \frac{1}{2}$$

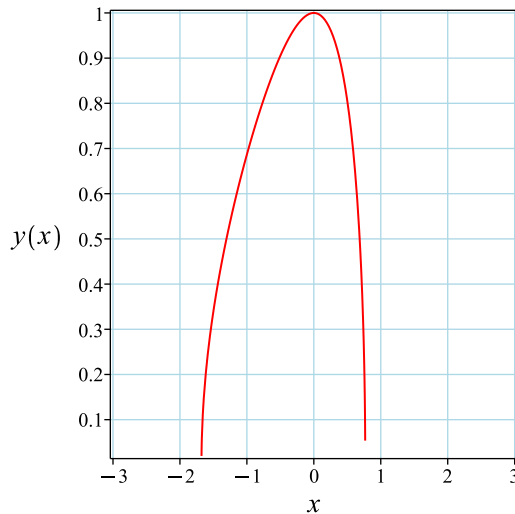
Solving for y from the above gives

$$y = \sqrt{-2xe^x + 2e^x - 1}$$

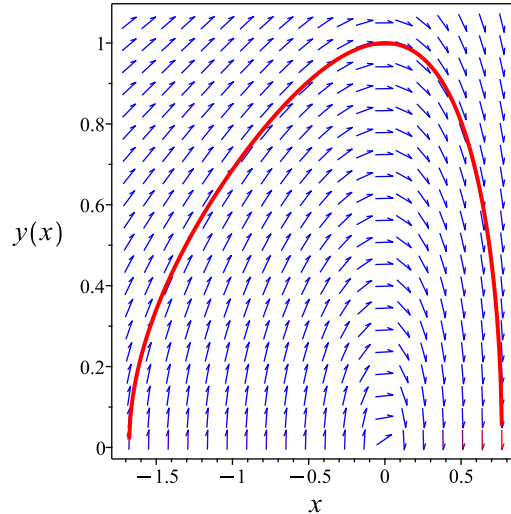
Summary

The solution(s) found are the following

$$y = \sqrt{-2x e^x + 2 e^x - 1} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \sqrt{-2x e^x + 2 e^x - 1}$$

Verified OK.

2.11.5 Maple step by step solution

Let's solve

$$\left[\frac{yy'}{e^x} = -x, y(0) = 1 \right]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$yy' = -x e^x$$

- Integrate both sides with respect to x

$$\int yy' dx = \int -x e^x dx + c_1$$

- Evaluate integral

$$\frac{y^2}{2} = -(x-1)e^x + c_1$$

- Solve for y
 $\{y = \sqrt{-2xe^x + 2e^x + 2c_1}, y = -\sqrt{-2xe^x + 2e^x + 2c_1}\}$
- Use initial condition $y(0) = 1$
 $1 = \sqrt{2 + 2c_1}$
- Solve for c_1
 $c_1 = -\frac{1}{2}$
- Substitute $c_1 = -\frac{1}{2}$ into general solution and simplify
 $y = \sqrt{-2xe^x + 2e^x - 1}$
- Use initial condition $y(0) = 1$
 $1 = -\sqrt{2 + 2c_1}$
- Solution does not satisfy initial condition
- Solution to the IVP
 $y = \sqrt{-2xe^x + 2e^x - 1}$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`

```

✓ Solution by Maple

Time used: 0.078 (sec). Leaf size: 17

```
dsolve([x+y(x)*diff(y(x),x)/exp(x) = 0,y(0) = 1],y(x), singsol=all)
```

$$y(x) = \sqrt{-1 - 2xe^x + 2e^x}$$

✓ Solution by Mathematica

Time used: 1.763 (sec). Leaf size: 19

```
DSolve[{x+y[x]*y'[x]/Exp[x] == 0,y[0]==1},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \sqrt{-2e^x(x-1)} - 1$$

2.12 problem 12

2.12.1 Existence and uniqueness analysis	557
2.12.2 Solving as separable ode	558
2.12.3 Solving as first order ode lie symmetry lookup ode	560
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2.12.6 Maple step by step solution	570

Internal problem ID [490]

Internal file name [OUTPUT/490_Sunday_June_05_2022_01_42_20_AM_13382328/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 12.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$r' - \frac{r^2}{x} = 0$$

With initial conditions

$$[r(1) = 2]$$

2.12.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned} r' &= f(x, r) \\ &= \frac{r^2}{x} \end{aligned}$$

The x domain of $f(x, r)$ when $r = 2$ is

$$\{x < 0 \vee 0 < x\}$$

And the point $x_0 = 1$ is inside this domain. The r domain of $f(x, r)$ when $x = 1$ is

$$\{-\infty < r < \infty\}$$

And the point $r_0 = 2$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial r} &= \frac{\partial}{\partial r} \left(\frac{r^2}{x} \right) \\ &= \frac{2r}{x}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial r}$ when $r = 2$ is

$$\{x < 0 \vee 0 < x\}$$

And the point $x_0 = 1$ is inside this domain. The r domain of $\frac{\partial f}{\partial r}$ when $x = 1$ is

$$\{-\infty < r < \infty\}$$

And the point $r_0 = 2$ is inside this domain. Therefore solution exists and is unique.

2.12.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}r' &= F(x, r) \\ &= f(x)g(r) \\ &= \frac{r^2}{x}\end{aligned}$$

Where $f(x) = \frac{1}{x}$ and $g(r) = r^2$. Integrating both sides gives

$$\begin{aligned}\frac{1}{r^2} dr &= \frac{1}{x} dx \\ \int \frac{1}{r^2} dr &= \int \frac{1}{x} dx \\ -\frac{1}{r} &= \ln(x) + c_1\end{aligned}$$

Which results in

$$r = -\frac{1}{\ln(x) + c_1}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $r = 2$ in the above solution gives an equation to solve for the constant of integration.

$$2 = -\frac{1}{c_1}$$

$$c_1 = -\frac{1}{2}$$

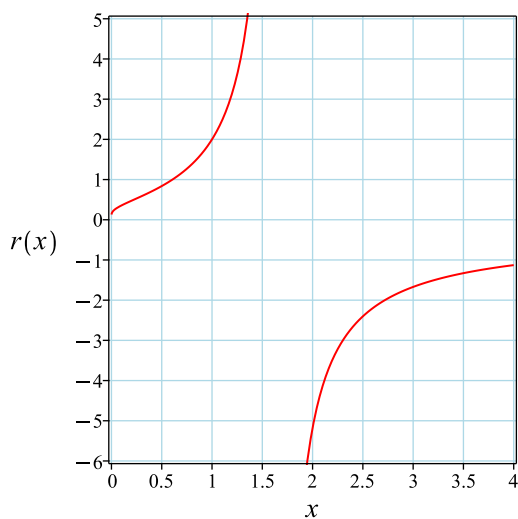
Substituting c_1 found above in the general solution gives

$$r = -\frac{2}{2\ln(x) - 1}$$

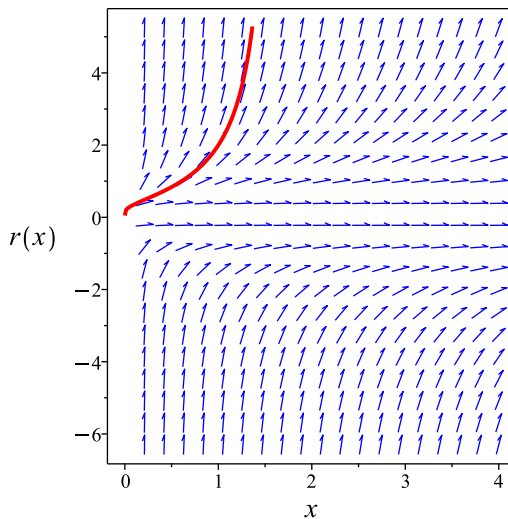
Summary

The solution(s) found are the following

$$r = -\frac{2}{2\ln(x) - 1} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$r = -\frac{2}{2\ln(x) - 1}$$

Verified OK.

2.12.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$r' = \frac{r^2}{x}$$

$$r' = \omega(x, r)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_r - \xi_x) - \omega^2 \xi_r - \omega_x \xi - \omega_r \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 124: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, r) &= x \\ \eta(x, r) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, r) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dr}{\eta} = dS\tag{1}$$

The above comes from the requirements that $(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial r}) S(x, r) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = r$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{x} dx\end{aligned}$$

Which results in

$$S = \ln(x)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, r)S_r}{R_x + \omega(x, r)R_r}\tag{2}$$

Where in the above R_x, R_r, S_x, S_r are all partial derivatives and $\omega(x, r)$ is the right hand side of the original ode given by

$$\omega(x, r) = \frac{r^2}{x}$$

Evaluating all the partial derivatives gives

$$R_x = 0$$

$$R_r = 1$$

$$S_x = \frac{1}{x}$$

$$S_r = 0$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{r^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, r in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{1}{R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, r coordinates. This results in

$$\ln(x) = -\frac{1}{r} + c_1$$

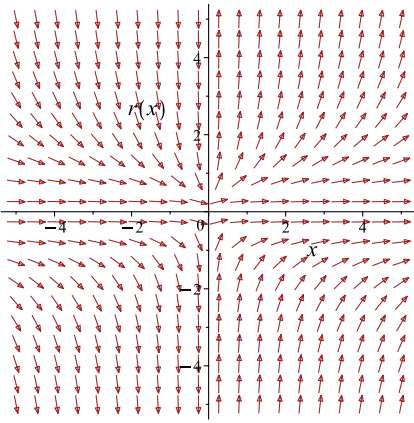
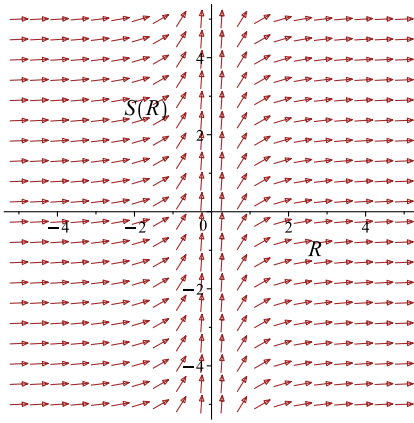
Which simplifies to

$$\ln(x) = -\frac{1}{r} + c_1$$

Which gives

$$r = -\frac{1}{\ln(x) - c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, r coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dr}{dx} = \frac{r^2}{x}$ 	$R = r$ $S = \ln(x)$	$\frac{dS}{dR} = \frac{1}{R^2}$ 

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $r = 2$ in the above solution gives an equation to solve for the constant of integration.

$$2 = \frac{1}{c_1}$$

$$c_1 = \frac{1}{2}$$

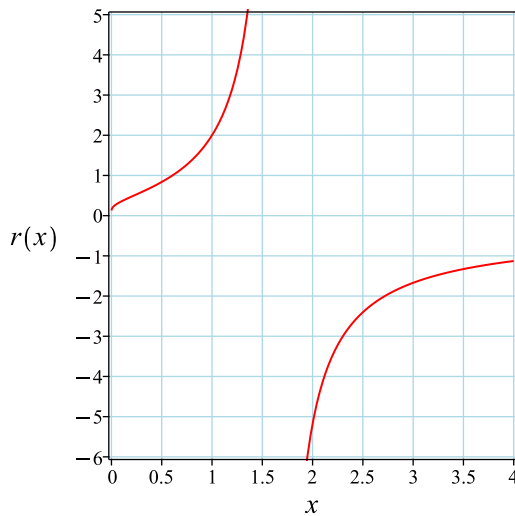
Substituting c_1 found above in the general solution gives

$$r = -\frac{2}{2 \ln(x) - 1}$$

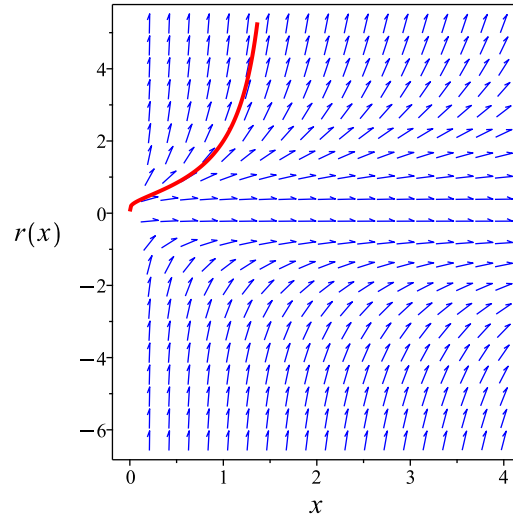
Summary

The solution(s) found are the following

$$r = -\frac{2}{2 \ln(x) - 1} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$r = -\frac{2}{2 \ln(x) - 1}$$

Verified OK.

2.12.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, r) dx + N(x, r) dr = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left(\frac{1}{r^2}\right) dr &= \left(\frac{1}{x}\right) dx \\ \left(-\frac{1}{x}\right) dx + \left(\frac{1}{r^2}\right) dr &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, r) &= -\frac{1}{x} \\ N(x, r) &= \frac{1}{r^2} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial r} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial r} &= \frac{\partial}{\partial r} \left(-\frac{1}{x}\right) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{1}{r^2}\right) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial r} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, r)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial r} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -\frac{1}{x} dx$$

$$\phi = -\ln(x) + f(r) \quad (3)$$

Where $f(r)$ is used for the constant of integration since ϕ is a function of both x and r . Taking derivative of equation (3) w.r.t r gives

$$\frac{\partial \phi}{\partial r} = 0 + f'(r) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial r} = \frac{1}{r^2}$. Therefore equation (4) becomes

$$\frac{1}{r^2} = 0 + f'(r) \quad (5)$$

Solving equation (5) for $f'(r)$ gives

$$f'(r) = \frac{1}{r^2}$$

Integrating the above w.r.t r gives

$$\int f'(r) dr = \int \left(\frac{1}{r^2}\right) dr$$

$$f(r) = -\frac{1}{r} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(r)$ into equation (3) gives ϕ

$$\phi = -\ln(x) - \frac{1}{r} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\ln(x) - \frac{1}{r}$$

The solution becomes

$$r = -\frac{1}{\ln(x) + c_1}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $r = 2$ in the above solution gives an equation to solve for the constant of integration.

$$2 = -\frac{1}{c_1}$$

$$c_1 = -\frac{1}{2}$$

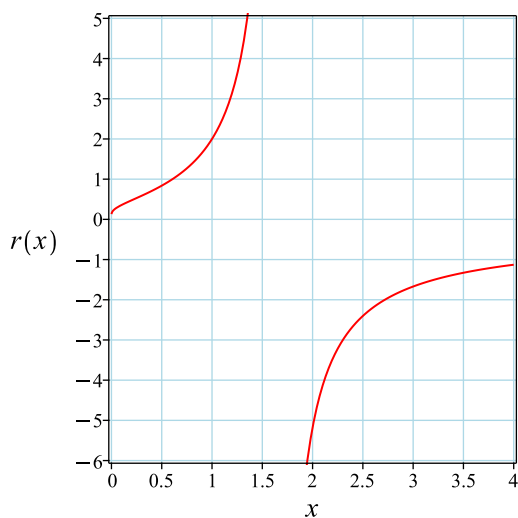
Substituting c_1 found above in the general solution gives

$$r = -\frac{2}{2\ln(x) - 1}$$

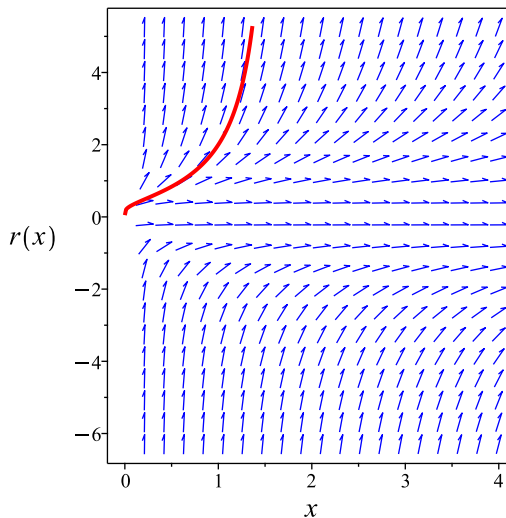
Summary

The solution(s) found are the following

$$r = -\frac{2}{2\ln(x) - 1} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$r = -\frac{2}{2 \ln(x) - 1}$$

Verified OK.

2.12.5 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} r' &= F(x, r) \\ &= \frac{r^2}{x} \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$r' = \frac{r^2}{x}$$

With Riccati ODE standard form

$$r' = f_0(x) + f_1(x)r + f_2(x)r^2$$

Shows that $f_0(x) = 0$, $f_1(x) = 0$ and $f_2(x) = \frac{1}{x}$. Let

$$\begin{aligned} r &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{\frac{u}{x}} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \tag{2}$$

But

$$\begin{aligned} f_2' &= -\frac{1}{x^2} \\ f_1 f_2 &= 0 \\ f_2^2 f_0 &= 0 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$\frac{u''(x)}{x} + \frac{u'(x)}{x^2} = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_2 \ln(x) + c_1$$

The above shows that

$$u'(x) = \frac{c_2}{x}$$

Using the above in (1) gives the solution

$$r = -\frac{c_2}{c_2 \ln(x) + c_1}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$r = -\frac{1}{\ln(x) + c_3}$$

Initial conditions are used to solve for c_3 . Substituting $x = 1$ and $r = 2$ in the above solution gives an equation to solve for the constant of integration.

$$2 = -\frac{1}{c_3}$$

$$c_3 = -\frac{1}{2}$$

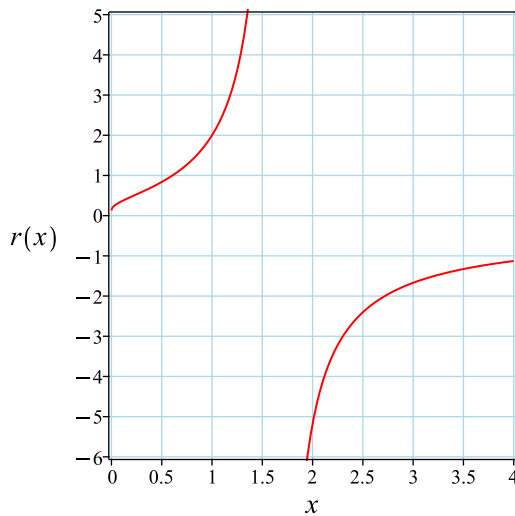
Substituting c_3 found above in the general solution gives

$$r = -\frac{2}{2 \ln(x) - 1}$$

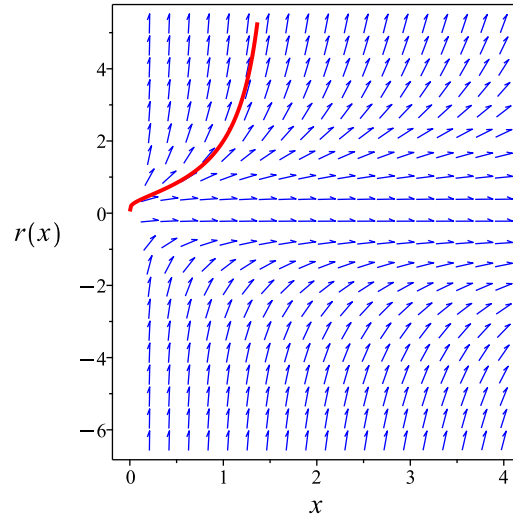
Summary

The solution(s) found are the following

$$r = -\frac{2}{2 \ln(x) - 1} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$r = -\frac{2}{2 \ln(x) - 1}$$

Verified OK.

2.12.6 Maple step by step solution

Let's solve

$$\left[r' - \frac{r^2}{x} = 0, r(1) = 2 \right]$$

- Highest derivative means the order of the ODE is 1

r'

- Separate variables

$$\frac{r'}{r^2} = \frac{1}{x}$$

- Integrate both sides with respect to x

$$\int \frac{r'}{r^2} dx = \int \frac{1}{x} dx + c_1$$

- Evaluate integral

$$-\frac{1}{r} = \ln(x) + c_1$$

- Solve for r

$$r = -\frac{1}{\ln(x)+c_1}$$

- Use initial condition $r(1) = 2$

$$2 = -\frac{1}{c_1}$$

- Solve for c_1

$$c_1 = -\frac{1}{2}$$

- Substitute $c_1 = -\frac{1}{2}$ into general solution and simplify

$$r = -\frac{2}{2\ln(x)-1}$$

- Solution to the IVP

$$r = -\frac{2}{2\ln(x)-1}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 14

```
dsolve([diff(r(x),x) = r(x)^2/x,r(1) = 2],r(x), singsol=all)
```

$$r(x) = -\frac{2}{2\ln(x) - 1}$$

✓ Solution by Mathematica

Time used: 0.128 (sec). Leaf size: 15

```
DSolve[{r'[x] == r[x]^2/x,r[1]==2},r[x],x,IncludeSingularSolutions -> True]
```

$$r(x) \rightarrow \frac{2}{1 - 2\log(x)}$$

2.13 problem 13

2.13.1 Existence and uniqueness analysis	572
2.13.2 Solving as separable ode	573
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Internal problem ID [491]

Internal file name [OUTPUT/491_Sunday_June_05_2022_01_42_21_AM_11702963/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 13.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{2x}{y + x^2y} = 0$$

With initial conditions

$$[y(0) = -2]$$

2.13.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned} y' &= f(x, y) \\ &= \frac{2x}{y(x^2 + 1)} \end{aligned}$$

The x domain of $f(x, y)$ when $y = -2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = -2$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{2x}{y(x^2 + 1)} \right) \\ &= -\frac{2x}{y^2(x^2 + 1)}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = -2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = -2$ is inside this domain. Therefore solution exists and is unique.

2.13.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{2x}{y(x^2 + 1)}\end{aligned}$$

Where $f(x) = \frac{2x}{x^2+1}$ and $g(y) = \frac{1}{y}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{1}{y}} dy &= \frac{2x}{x^2 + 1} dx \\ \int \frac{1}{\frac{1}{y}} dy &= \int \frac{2x}{x^2 + 1} dx \\ \frac{y^2}{2} &= \ln(x^2 + 1) + c_1\end{aligned}$$

Which results in

$$y = \sqrt{2 \ln(x^2 + 1) + 2c_1}$$

$$y = -\sqrt{2 \ln(x^2 + 1) + 2c_1}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = -\sqrt{c_1} \sqrt{2}$$

$$c_1 = 2$$

Substituting c_1 found above in the general solution gives

$$y = -\sqrt{2 \ln(x^2 + 1) + 4}$$

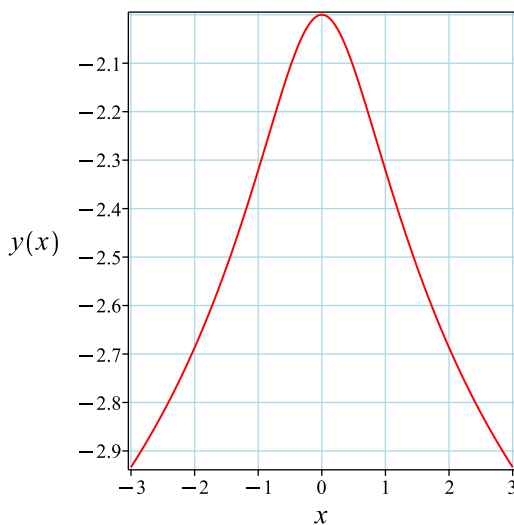
Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = \sqrt{c_1} \sqrt{2}$$

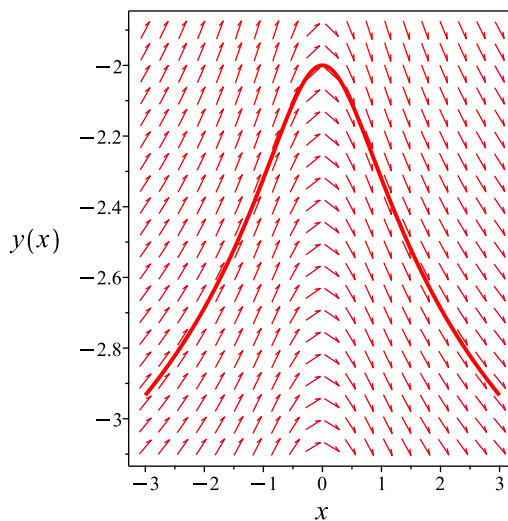
Summary

Warning: Unable to solve for constant of integration. The solution(s) found are the following

$$y = -\sqrt{2 \ln(x^2 + 1)}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\sqrt{2 \ln(x^2 + 1) + 4}$$

Verified OK.

2.13.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{2x}{y(x^2 + 1)}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 127: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{x^2 + 1}{2x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{x^2+1}{2x}} dx \end{aligned}$$

Which results in

$$S = \ln(x^2 + 1)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{2x}{y(x^2 + 1)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= \frac{2x}{x^2 + 1} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^2}{2} + c_1 \quad (4)$$

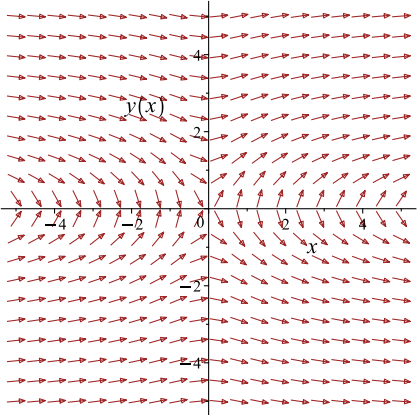
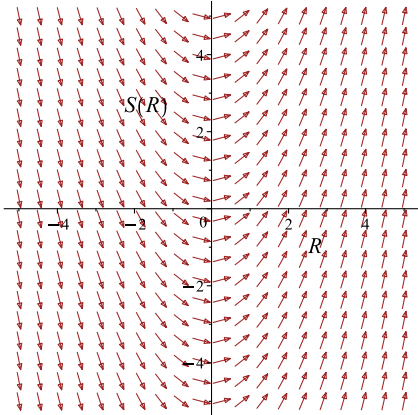
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\ln(x^2 + 1) = \frac{y^2}{2} + c_1$$

Which simplifies to

$$\ln(x^2 + 1) = \frac{y^2}{2} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{2x}{y(x^2+1)}$ 	$R = y$ $S = \ln(x^2 + 1)$	$\frac{dS}{dR} = R$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$0 = c_1 + 2$$

$$c_1 = -2$$

Substituting c_1 found above in the general solution gives

$$\ln(x^2 + 1) = \frac{y^2}{2} - 2$$

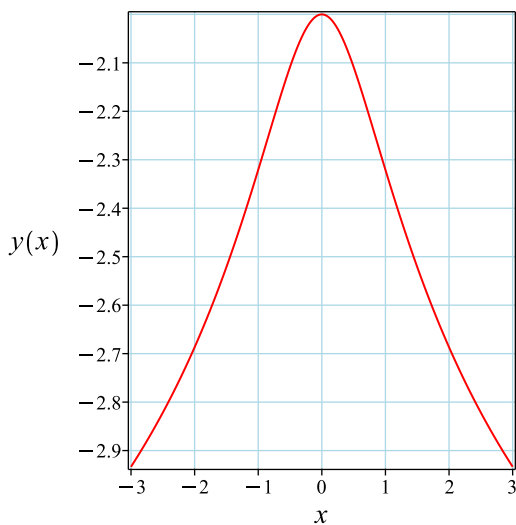
Solving for y from the above gives

$$y = -\sqrt{2 \ln(x^2 + 1) + 4}$$

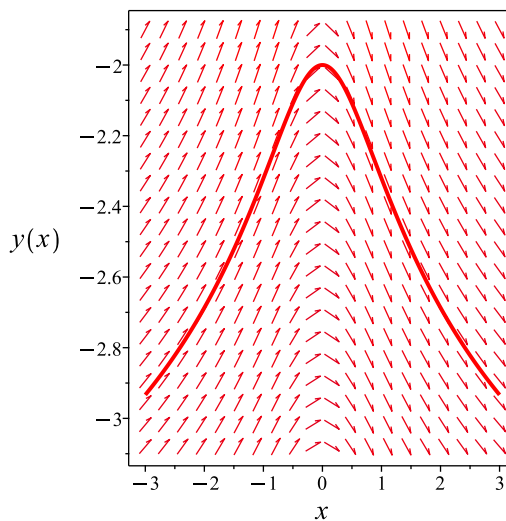
Summary

The solution(s) found are the following

$$y = -\sqrt{2 \ln(x^2 + 1) + 4} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\sqrt{2 \ln(x^2 + 1) + 4}$$

Verified OK.

2.13.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} \left(\frac{y}{2}\right) dy &= \left(\frac{x}{x^2+1}\right) dx \\ \left(-\frac{x}{x^2+1}\right) dx + \left(\frac{y}{2}\right) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$M(x, y) = -\frac{x}{x^2 + 1}$$
$$N(x, y) = \frac{y}{2}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y} \left(-\frac{x}{x^2 + 1} \right)$$
$$= 0$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x} \left(\frac{y}{2} \right)$$
$$= 0$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -\frac{x}{x^2 + 1} dx$$

$$\phi = -\frac{\ln(x^2 + 1)}{2} + f(y) \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{y}{2}$. Therefore equation (4) becomes

$$\frac{y}{2} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{y}{2}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(\frac{y}{2}\right) dy$$
$$f(y) = \frac{y^2}{4} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{\ln(x^2 + 1)}{2} + \frac{y^2}{4} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{\ln(x^2 + 1)}{2} + \frac{y^2}{4}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$1 = c_1$$

$$c_1 = 1$$

Substituting c_1 found above in the general solution gives

$$-\frac{\ln(x^2 + 1)}{2} + \frac{y^2}{4} = 1$$

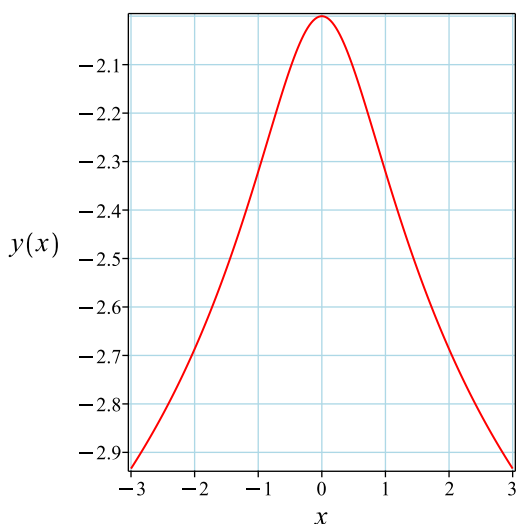
Solving for y from the above gives

$$y = -\sqrt{2 \ln(x^2 + 1)} + 4$$

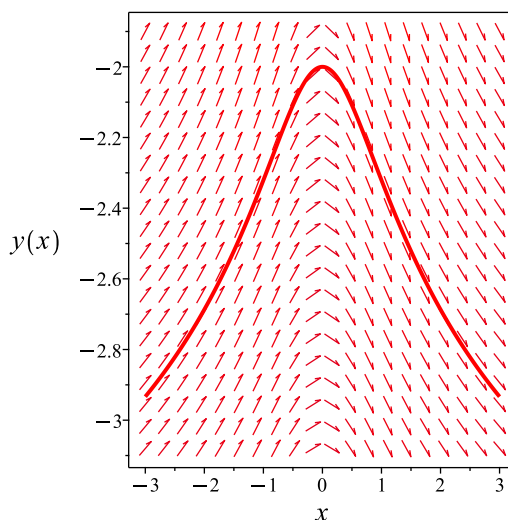
Summary

The solution(s) found are the following

$$y = -\sqrt{2 \ln(x^2 + 1)} + 4 \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\sqrt{2 \ln(x^2 + 1)} + 4$$

Verified OK.

2.13.5 Maple step by step solution

Let's solve

$$\left[y' - \frac{2x}{y+x^2y} = 0, y(0) = -2 \right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$yy' = \frac{2x}{x^2+1}$$

- Integrate both sides with respect to x

$$\int yy' dx = \int \frac{2x}{x^2+1} dx + c_1$$
- Evaluate integral

$$\frac{y^2}{2} = \ln(x^2 + 1) + c_1$$
- Solve for y

$$\left\{ y = \sqrt{2 \ln(x^2 + 1) + 2c_1}, y = -\sqrt{2 \ln(x^2 + 1) + 2c_1} \right\}$$
- Use initial condition $y(0) = -2$

$$-2 = \sqrt{c_1} \sqrt{2}$$
- Solution does not satisfy initial condition
- Use initial condition $y(0) = -2$

$$-2 = -\sqrt{c_1} \sqrt{2}$$
- Solve for c_1

$$c_1 = 2$$
- Substitute $c_1 = 2$ into general solution and simplify

$$y = -\sqrt{2 \ln(x^2 + 1) + 4}$$
- Solution to the IVP

$$y = -\sqrt{2 \ln(x^2 + 1) + 4}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`

```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 18

```
dsolve([diff(y(x),x) = 2*x/(y(x)+x^2*y(x)),y(0) = -2],y(x), singsol=all)
```

$$y(x) = -\sqrt{2 \ln(x^2 + 1) + 4}$$

✓ Solution by Mathematica

Time used: 0.091 (sec). Leaf size: 24

```
DSolve[{y'[x] == 2*x/(y[x]+x^2*y[x]),y[0]==-2},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\sqrt{2} \sqrt{\log(x^2 + 1) + 2}$$

2.14 problem 14

2.14.1 Existence and uniqueness analysis	587
2.14.2 Solving as separable ode	587
2.14.3 Solving as first order ode lie symmetry lookup ode	589
2.14.4 Solving as exact ode	594
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2.14.6 Maple step by step solution	600

Internal problem ID [492]

Internal file name [OUTPUT/492_Sunday_June_05_2022_01_42_22_AM_57989922/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 14.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{xy^2}{\sqrt{x^2 + 1}} = 0$$

With initial conditions

$$[y(0) = 1]$$

2.14.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= \frac{xy^2}{\sqrt{x^2 + 1}}\end{aligned}$$

The x domain of $f(x, y)$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{xy^2}{\sqrt{x^2 + 1}} \right) \\ &= \frac{2xy}{\sqrt{x^2 + 1}}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Therefore solution exists and is unique.

2.14.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{xy^2}{\sqrt{x^2 + 1}}\end{aligned}$$

Where $f(x) = \frac{x}{\sqrt{x^2+1}}$ and $g(y) = y^2$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y^2} dy &= \frac{x}{\sqrt{x^2+1}} dx \\ \int \frac{1}{y^2} dy &= \int \frac{x}{\sqrt{x^2+1}} dx \\ -\frac{1}{y} &= \sqrt{x^2+1} + c_1\end{aligned}$$

Which results in

$$y = -\frac{1}{\sqrt{x^2+1} + c_1}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{1}{1 + c_1}$$

$$c_1 = -2$$

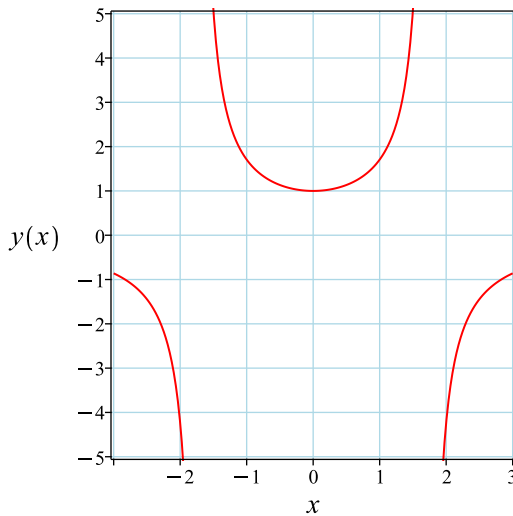
Substituting c_1 found above in the general solution gives

$$y = -\frac{1}{\sqrt{x^2+1} - 2}$$

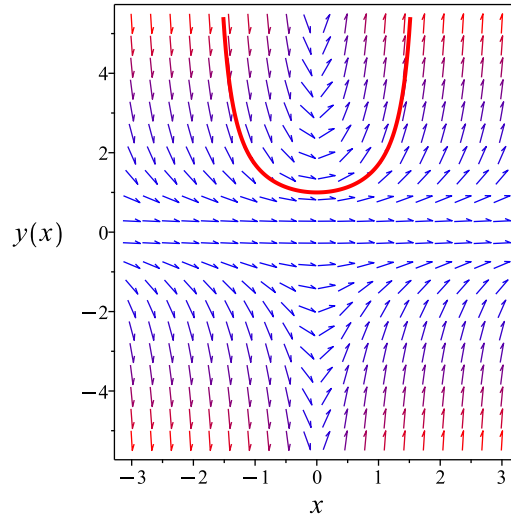
Summary

The solution(s) found are the following

$$y = -\frac{1}{\sqrt{x^2+1} - 2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{1}{\sqrt{x^2 + 1} - 2}$$

Verified OK.

2.14.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x y^2}{\sqrt{x^2 + 1}}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 130: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{\sqrt{x^2 + 1}}{x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{\sqrt{x^2+1}}{x}} dx \end{aligned}$$

Which results in

$$S = \sqrt{x^2 + 1}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{xy^2}{\sqrt{x^2 + 1}}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= \frac{x}{\sqrt{x^2 + 1}} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{y^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{1}{R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\sqrt{x^2 + 1} = -\frac{1}{y} + c_1$$

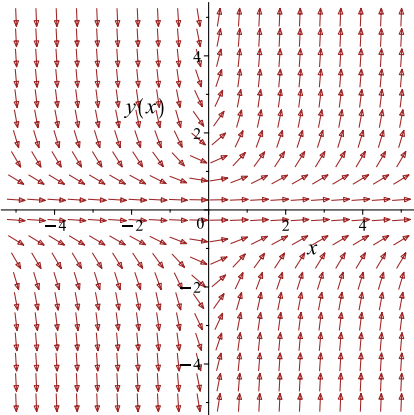
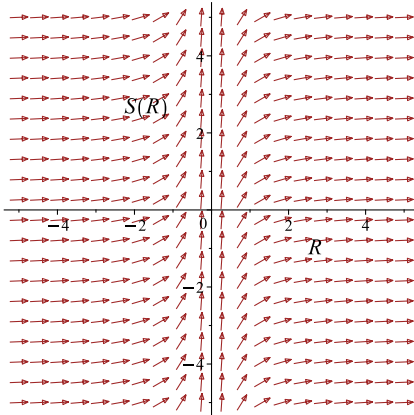
Which simplifies to

$$\sqrt{x^2 + 1} = -\frac{1}{y} + c_1$$

Which gives

$$y = -\frac{1}{\sqrt{x^2 + 1} - c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{xy^2}{\sqrt{x^2+1}}$ 	$R = y$ $S = \sqrt{x^2 + 1}$	$\frac{dS}{dR} = \frac{1}{R^2}$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{1}{c_1 - 1}$$

$$c_1 = 2$$

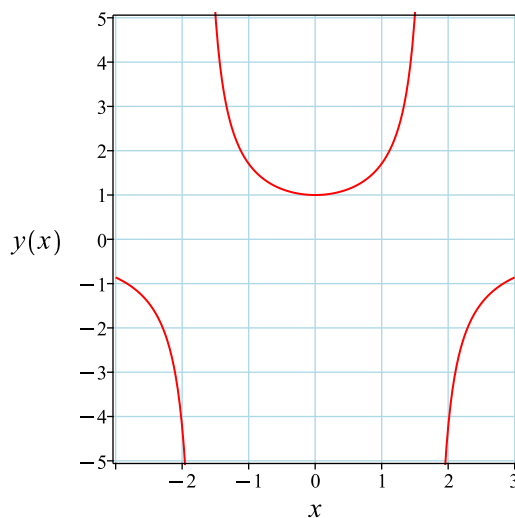
Substituting c_1 found above in the general solution gives

$$y = -\frac{1}{\sqrt{x^2 + 1} - 2}$$

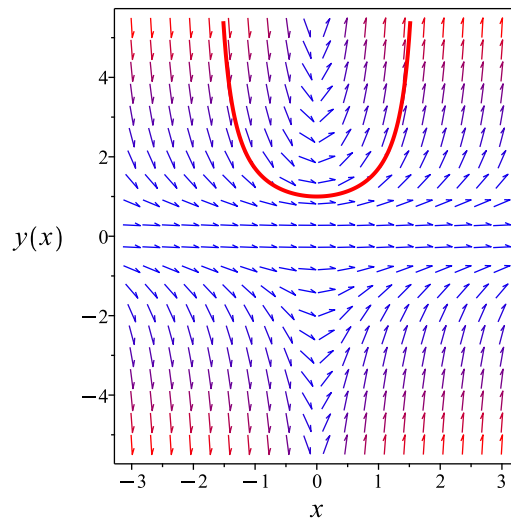
Summary

The solution(s) found are the following

$$y = -\frac{1}{\sqrt{x^2 + 1} - 2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{1}{\sqrt{x^2 + 1} - 2}$$

Verified OK.

2.14.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} \left(\frac{1}{y^2}\right) dy &= \left(\frac{x}{\sqrt{x^2+1}}\right) dx \\ \left(-\frac{x}{\sqrt{x^2+1}}\right) dx &+ \left(\frac{1}{y^2}\right) dy = 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$M(x, y) = -\frac{x}{\sqrt{x^2 + 1}}$$
$$N(x, y) = \frac{1}{y^2}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y} \left(-\frac{x}{\sqrt{x^2 + 1}} \right)$$
$$= 0$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x} \left(\frac{1}{y^2} \right)$$
$$= 0$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$
$$\int \frac{\partial \phi}{\partial x} dx = \int -\frac{x}{\sqrt{x^2 + 1}} dx$$
$$\phi = -\sqrt{x^2 + 1} + f(y) \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{y^2}$. Therefore equation (4) becomes

$$\frac{1}{y^2} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{1}{y^2}$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) \, dy &= \int \left(\frac{1}{y^2} \right) \, dy \\ f(y) &= -\frac{1}{y} + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\sqrt{x^2 + 1} - \frac{1}{y} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\sqrt{x^2 + 1} - \frac{1}{y}$$

The solution becomes

$$y = -\frac{1}{\sqrt{x^2 + 1} + c_1}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{1}{1 + c_1}$$

$$c_1 = -2$$

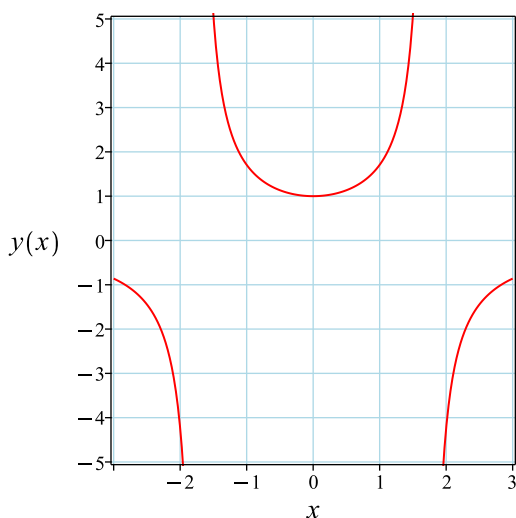
Substituting c_1 found above in the general solution gives

$$y = -\frac{1}{\sqrt{x^2 + 1} - 2}$$

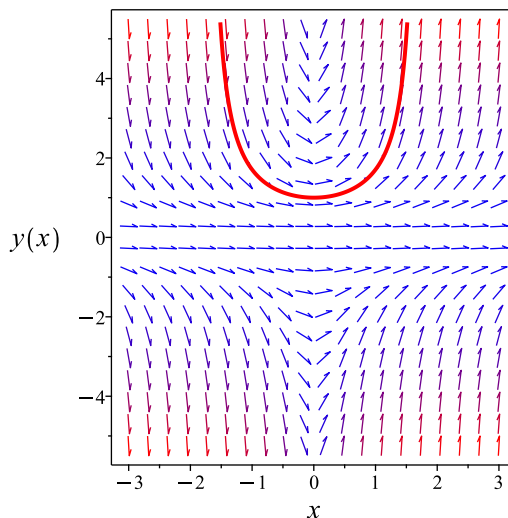
Summary

The solution(s) found are the following

$$y = -\frac{1}{\sqrt{x^2 + 1} - 2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{1}{\sqrt{x^2 + 1} - 2}$$

Verified OK.

2.14.5 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{xy^2}{\sqrt{x^2 + 1}} \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = \frac{xy^2}{\sqrt{x^2 + 1}}$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that $f_0(x) = 0$, $f_1(x) = 0$ and $f_2(x) = \frac{x}{\sqrt{x^2+1}}$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{\frac{xu}{\sqrt{x^2+1}}} \end{aligned} \quad (1)$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \quad (2)$$

But

$$\begin{aligned} f_2' &= \frac{1}{\sqrt{x^2+1}} - \frac{x^2}{(x^2+1)^{\frac{3}{2}}} \\ f_1 f_2 &= 0 \\ f_2^2 f_0 &= 0 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$\frac{xu''(x)}{\sqrt{x^2+1}} - \left(\frac{1}{\sqrt{x^2+1}} - \frac{x^2}{(x^2+1)^{\frac{3}{2}}} \right) u'(x) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 + \sqrt{x^2+1} c_2$$

The above shows that

$$u'(x) = \frac{c_2 x}{\sqrt{x^2+1}}$$

Using the above in (1) gives the solution

$$y = -\frac{c_2}{c_1 + \sqrt{x^2+1} c_2}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = -\frac{1}{c_3 + \sqrt{x^2 + 1}}$$

Initial conditions are used to solve for c_3 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{1}{c_3 + 1}$$

$$c_3 = -2$$

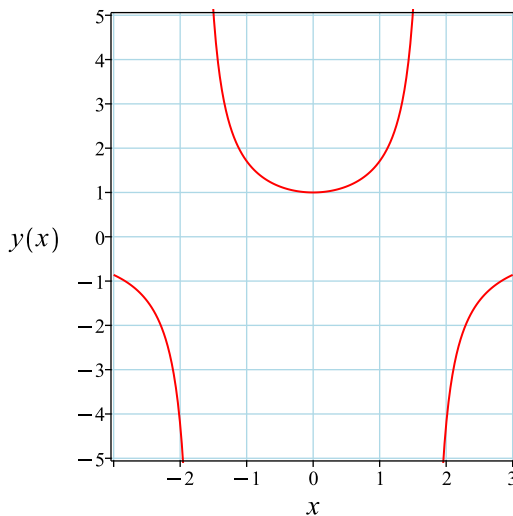
Substituting c_3 found above in the general solution gives

$$y = -\frac{1}{\sqrt{x^2 + 1} - 2}$$

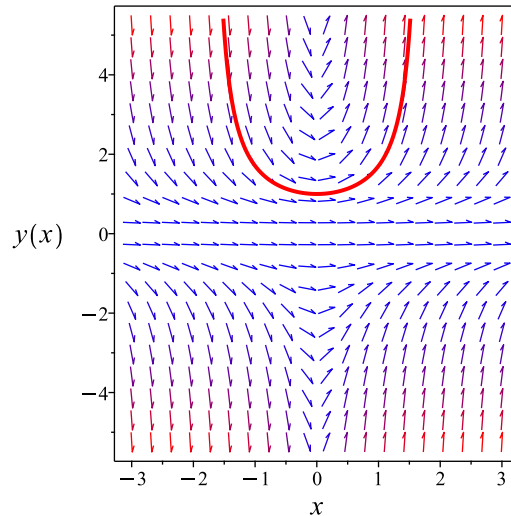
Summary

The solution(s) found are the following

$$y = -\frac{1}{\sqrt{x^2 + 1} - 2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{1}{\sqrt{x^2 + 1} - 2}$$

Verified OK.

2.14.6 Maple step by step solution

Let's solve

$$\left[y' - \frac{xy^2}{\sqrt{x^2+1}} = 0, y(0) = 1 \right]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'}{y^2} = \frac{x}{\sqrt{x^2+1}}$$

- Integrate both sides with respect to x

$$\int \frac{y'}{y^2} dx = \int \frac{x}{\sqrt{x^2+1}} dx + c_1$$

- Evaluate integral

$$-\frac{1}{y} = \sqrt{x^2+1} + c_1$$

- Solve for y

$$y = -\frac{1}{\sqrt{x^2+1} + c_1}$$

- Use initial condition $y(0) = 1$

$$1 = -\frac{1}{1+c_1}$$

- Solve for c_1

$$c_1 = -2$$

- Substitute $c_1 = -2$ into general solution and simplify

$$y = -\frac{1}{\sqrt{x^2+1}-2}$$

- Solution to the IVP

$$y = -\frac{1}{\sqrt{x^2+1}-2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.032 (sec). Leaf size: 17

```
dsolve([diff(y(x),x) = x*y(x)^2/(x^2+1)^(1/2),y(0) = 1],y(x), singsol=all)
```

$$y(x) = -\frac{1}{\sqrt{x^2 + 1} - 2}$$

✓ Solution by Mathematica

Time used: 0.183 (sec). Leaf size: 20

```
DSolve[{y'[x] == x*y[x]^2/(x^2+1)^(1/2),y[0]==1},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2 - \sqrt{x^2 + 1}}$$

2.15 problem 15

2.15.1 Existence and uniqueness analysis	603
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Internal problem ID [493]

Internal file name [OUTPUT/493_Sunday_June_05_2022_01_42_24_AM_81327021/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 15.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "homogeneousTypeMapleC", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{2x}{1+2y} = 0$$

With initial conditions

$$[y(2) = 0]$$

2.15.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= \frac{2x}{1 + 2y}\end{aligned}$$

The x domain of $f(x, y)$ when $y = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 2$ is inside this domain. The y domain of $f(x, y)$ when $x = 2$ is

$$\left\{y < -\frac{1}{2} \vee -\frac{1}{2} < y\right\}$$

And the point $y_0 = 0$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{2x}{1 + 2y} \right) \\ &= -\frac{4x}{(1 + 2y)^2}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 2$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 2$ is

$$\left\{y < -\frac{1}{2} \vee -\frac{1}{2} < y\right\}$$

And the point $y_0 = 0$ is inside this domain. Therefore solution exists and is unique.

2.15.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{2x}{1+2y}\end{aligned}$$

Where $f(x) = 2x$ and $g(y) = \frac{1}{1+2y}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{1}{1+2y}} dy &= 2x dx \\ \int \frac{1}{\frac{1}{1+2y}} dy &= \int 2x dx \\ y^2 + y &= x^2 + c_1\end{aligned}$$

Which results in

$$\begin{aligned}y &= -\frac{1}{2} + \frac{\sqrt{4x^2 + 4c_1 + 1}}{2} \\ y &= -\frac{1}{2} - \frac{\sqrt{4x^2 + 4c_1 + 1}}{2}\end{aligned}$$

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{1}{2} - \frac{\sqrt{17 + 4c_1}}{2}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{1}{2} + \frac{\sqrt{17 + 4c_1}}{2}$$

$$c_1 = -4$$

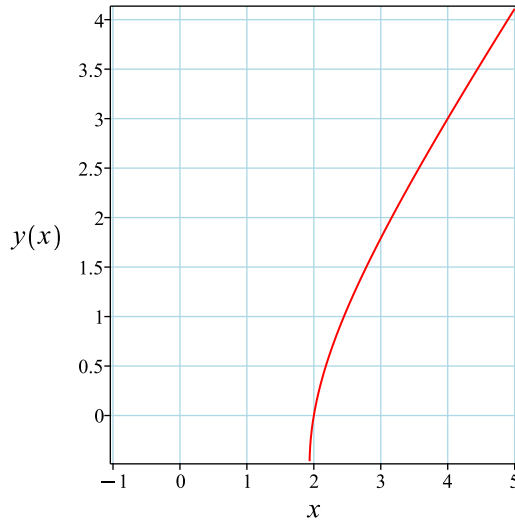
Substituting c_1 found above in the general solution gives

$$y = -\frac{1}{2} + \frac{\sqrt{4x^2 - 15}}{2}$$

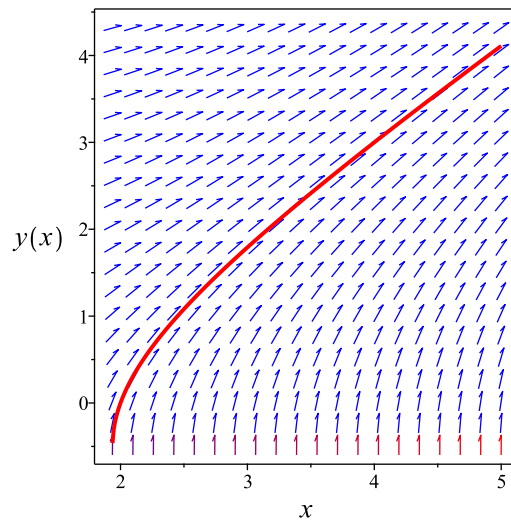
Summary

The solution(s) found are the following

$$y = -\frac{1}{2} + \frac{\sqrt{4x^2 - 15}}{2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{1}{2} + \frac{\sqrt{4x^2 - 15}}{2}$$

Verified OK.

2.15.3 Solving as differentialType ode

Writing the ode as

$$y' = \frac{2x}{1 + 2y} \quad (1)$$

Which becomes

$$(1 + 2y) dy = (2x) dx \quad (2)$$

But the RHS is complete differential because

$$(2x) dx = d(x^2)$$

Hence (2) becomes

$$(1 + 2y) dy = d(x^2)$$

Integrating both sides gives gives these solutions

$$y = -\frac{1}{2} + \frac{\sqrt{4x^2 + 4c_1 + 1}}{2} + c_1$$

$$y = -\frac{1}{2} - \frac{\sqrt{4x^2 + 4c_1 + 1}}{2} + c_1$$

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{1}{2} - \frac{\sqrt{17 + 4c_1}}{2} + c_1$$

$$c_1 = \sqrt{5} + 1$$

Substituting c_1 found above in the general solution gives

$$y = \frac{1}{2} - \frac{\sqrt{4x^2 + 4\sqrt{5} + 5}}{2} + \sqrt{5}$$

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{1}{2} + \frac{\sqrt{17 + 4c_1}}{2} + c_1$$

$$c_1 = -\sqrt{5} + 1$$

Substituting c_1 found above in the general solution gives

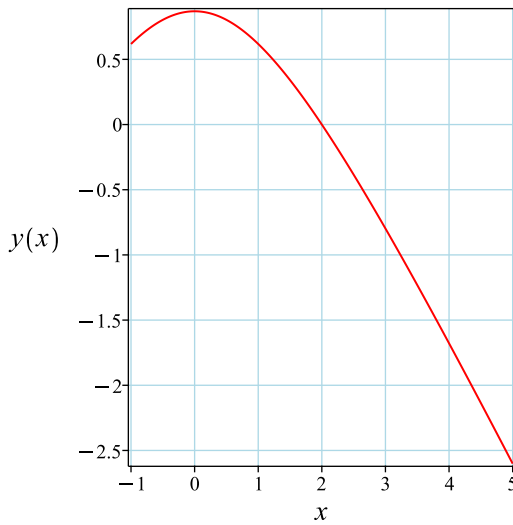
$$y = \frac{1}{2} + \frac{\sqrt{4x^2 - 4\sqrt{5} + 5}}{2} - \sqrt{5}$$

Summary

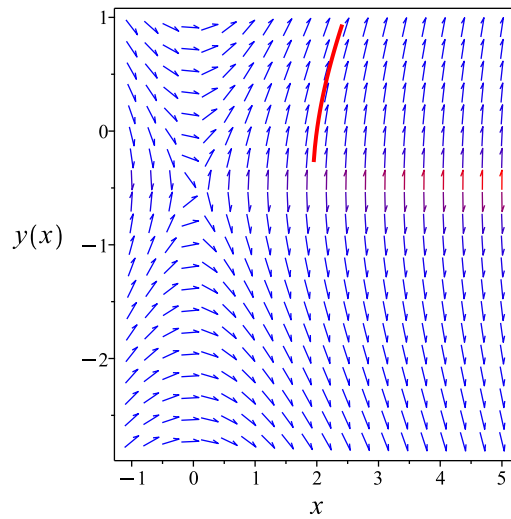
The solution(s) found are the following

$$y = \frac{1}{2} + \frac{\sqrt{4x^2 - 4\sqrt{5} + 5}}{2} - \sqrt{5} \quad (1)$$

$$y = \frac{1}{2} - \frac{\sqrt{4x^2 + 4\sqrt{5} + 5}}{2} + \sqrt{5} \quad (2)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{1}{2} + \frac{\sqrt{4x^2 - 4\sqrt{5} + 5}}{2} - \sqrt{5}$$

Verified OK.

$$y = \frac{1}{2} - \frac{\sqrt{4x^2 + 4\sqrt{5} + 5}}{2} + \sqrt{5}$$

Verified OK.

2.15.4 Solving as homogeneousTypeMapleC ode

Let $Y = y + y_0$ and $X = x + x_0$ then the above is transformed to new ode in $Y(X)$

$$\frac{d}{dX}Y(X) = \frac{2X + 2x_0}{1 + 2Y(X) + 2y_0}$$

Solving for possible values of x_0 and y_0 which makes the above ode a homogeneous ode results in

$$\begin{aligned} x_0 &= 0 \\ y_0 &= -\frac{1}{2} \end{aligned}$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$\frac{d}{dX}Y(X) = \frac{X}{Y(X)}$$

In canonical form, the ODE is

$$\begin{aligned} Y' &= F(X, Y) \\ &= \frac{X}{Y} \end{aligned} \tag{1}$$

An ode of the form $Y' = \frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order n if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both $M = X$ and $N = Y$ are both homogeneous and of the same order $n = 1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{Y}{X}$, or $Y = uX$. Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation $Y = uX$ to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dX}X + u &= \frac{1}{u} \\ \frac{du}{dX} &= \frac{\frac{1}{u(X)} - u(X)}{X} \end{aligned}$$

Or

$$\frac{d}{dX}u(X) - \frac{\frac{1}{u(X)} - u(X)}{X} = 0$$

Or

$$\left(\frac{d}{dX}u(X) \right) u(X) X + u(X)^2 - 1 = 0$$

Which is now solved as separable in $u(X)$. Which is now solved in $u(X)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(X, u) \\ &= f(X)g(u) \\ &= -\frac{u^2 - 1}{uX} \end{aligned}$$

Where $f(X) = -\frac{1}{X}$ and $g(u) = \frac{u^2-1}{u}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{u^2-1}{u}} du &= -\frac{1}{X} dX \\ \int \frac{1}{\frac{u^2-1}{u}} du &= \int -\frac{1}{X} dX \\ \frac{\ln(u-1)}{2} + \frac{\ln(u+1)}{2} &= -\ln(X) + c_2\end{aligned}$$

The above can be written as

$$\begin{aligned}\left(\frac{1}{2}\right) (\ln(u-1) + \ln(u+1)) &= -\ln(X) + 2c_2 \\ \ln(u-1) + \ln(u+1) &= (2)(-\ln(X) + 2c_2) \\ &= -2\ln(X) + 4c_2\end{aligned}$$

Raising both side to exponential gives

$$e^{\ln(u-1)+\ln(u+1)} = e^{-2\ln(X)+2c_2}$$

Which simplifies to

$$\begin{aligned}u^2 - 1 &= \frac{2c_2}{X^2} \\ &= \frac{c_3}{X^2}\end{aligned}$$

The solution is

$$u(X)^2 - 1 = \frac{c_3}{X^2}$$

Now u in the above solution is replaced back by Y using $u = \frac{Y}{X}$ which results in the solution

$$\frac{Y(X)^2}{X^2} - 1 = \frac{c_3}{X^2}$$

Which simplifies to

$$-(X - Y(X))(X + Y(X)) = c_3$$

Using the solution for $Y(X)$

$$-(X - Y(X))(X + Y(X)) = c_3$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y - \frac{1}{2}$$

$$X = x$$

Then the solution in y becomes

$$-\left(x - y - \frac{1}{2}\right)\left(x + y + \frac{1}{2}\right) = c_3$$

Initial conditions are used to solve for c_3 . Substituting $x = 2$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{15}{4} = c_3$$

$$c_3 = -\frac{15}{4}$$

Substituting c_3 found above in the general solution gives

$$-\left(x - y - \frac{1}{2}\right)\left(x + y + \frac{1}{2}\right) = -\frac{15}{4}$$

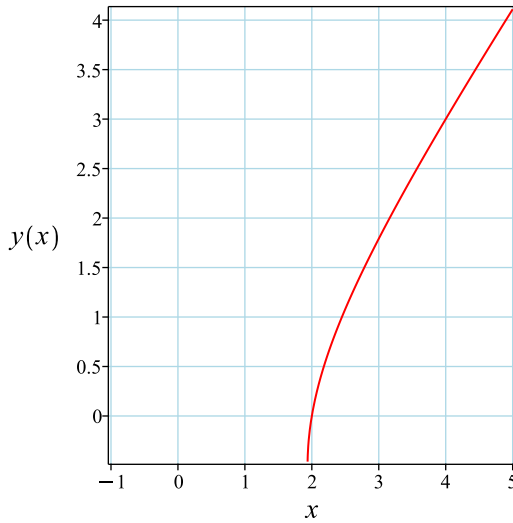
Solving for y from the above gives

$$y = -\frac{1}{2} + \frac{\sqrt{4x^2 - 15}}{2}$$

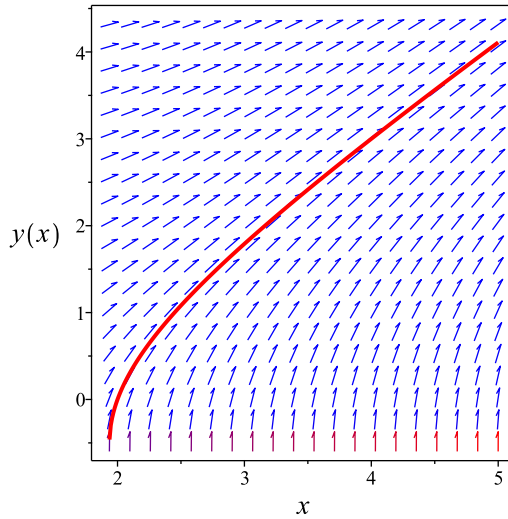
Summary

The solution(s) found are the following

$$y = -\frac{1}{2} + \frac{\sqrt{4x^2 - 15}}{2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{1}{2} + \frac{\sqrt{4x^2 - 15}}{2}$$

Verified OK.

2.15.5 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{2x}{1 + 2y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 133: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{2x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{2x}} dx \end{aligned}$$

Which results in

$$S = x^2$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{2x}{1 + 2y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= 2x \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 1 + 2y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 1 + 2R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R^2 + R + c_1 \quad (4)$$

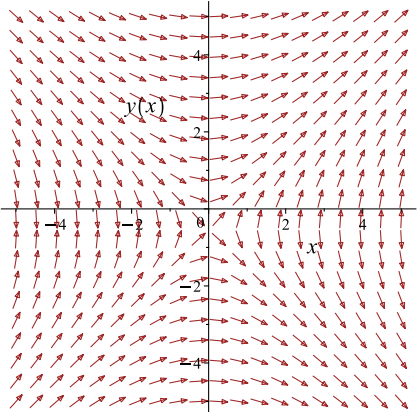
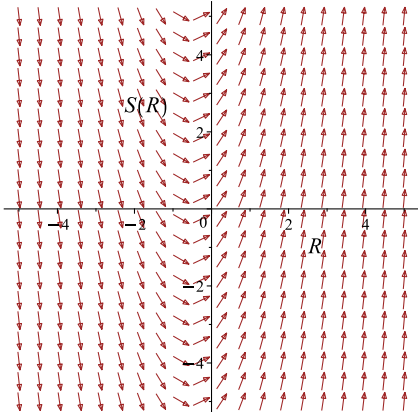
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$x^2 = y^2 + c_1 + y$$

Which simplifies to

$$x^2 = y^2 + c_1 + y$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{2x}{1+2y}$ 	$R = y$ $S = x^2$	$\frac{dS}{dR} = 1 + 2R$ 

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$4 = c_1$$

$$c_1 = 4$$

Substituting c_1 found above in the general solution gives

$$x^2 = y^2 + y + 4$$

Summary

The solution(s) found are the following

$$x^2 = y^2 + y + 4 \quad (1)$$

Verification of solutions

$$x^2 = y^2 + y + 4$$

Verified OK.

2.15.6 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work

and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left(y + \frac{1}{2}\right) dy &= (x) dx \\ (-x) dx + \left(y + \frac{1}{2}\right) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x \\ N(x, y) &= y + \frac{1}{2} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}\left(y + \frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x dx \\ \phi &= -\frac{x^2}{2} + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y)\quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = y + \frac{1}{2}$. Therefore equation (4) becomes

$$y + \frac{1}{2} = 0 + f'(y)\quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = y + \frac{1}{2}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(y + \frac{1}{2}\right) dy \\ f(y) &= \frac{1}{2}y^2 + \frac{1}{2}y + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{1}{2}x^2 + \frac{1}{2}y^2 + \frac{1}{2}y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{1}{2}x^2 + \frac{1}{2}y^2 + \frac{1}{2}y$$

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = c_1$$

$$c_1 = -2$$

Substituting c_1 found above in the general solution gives

$$-\frac{1}{2}x^2 + \frac{1}{2}y^2 + \frac{1}{2}y = -2$$

Summary

The solution(s) found are the following

$$-\frac{x^2}{2} + \frac{y^2}{2} + \frac{y}{2} = -2 \quad (1)$$

Verification of solutions

$$-\frac{x^2}{2} + \frac{y^2}{2} + \frac{y}{2} = -2$$

Verified OK.

2.15.7 Maple step by step solution

Let's solve

$$\left[y' - \frac{2x}{1+2y} = 0, y(2) = 0 \right]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$y'(1 + 2y) = 2x$$

- Integrate both sides with respect to x

$$\int y'(1 + 2y) dx = \int 2x dx + c_1$$

- Evaluate integral

$$y^2 + y = x^2 + c_1$$

- Solve for y

$$\left\{ y = -\frac{1}{2} - \frac{\sqrt{4x^2+4c_1+1}}{2}, y = -\frac{1}{2} + \frac{\sqrt{4x^2+4c_1+1}}{2} \right\}$$

- Use initial condition $y(2) = 0$

$$0 = -\frac{1}{2} - \frac{\sqrt{17+4c_1}}{2}$$

- Solution does not satisfy initial condition

- Use initial condition $y(2) = 0$

$$0 = -\frac{1}{2} + \frac{\sqrt{17+4c_1}}{2}$$

- Solve for c_1

$$c_1 = -4$$

- Substitute $c_1 = -4$ into general solution and simplify

$$y = -\frac{1}{2} + \frac{\sqrt{4x^2-15}}{2}$$

- Solution to the IVP

$$y = -\frac{1}{2} + \frac{\sqrt{4x^2-15}}{2}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 17

```
dsolve([diff(y(x),x) = 2*x/(1+2*y(x)),y(2) = 0],y(x), singsol=all)
```

$$y(x) = -\frac{1}{2} + \frac{\sqrt{4x^2 - 15}}{2}$$

✓ Solution by Mathematica

Time used: 0.115 (sec). Leaf size: 22

```
DSolve[{y'[x] == 2*x/(1+2*y[x]), y[2]==0}, y[x], x, IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2} \left(\sqrt{4x^2 - 15} - 1 \right)$$

2.16 problem 16

2.16.1 Existence and uniqueness analysis	622
2.16.2 Solving as separable ode	622
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Internal problem ID [494]

Internal file name [OUTPUT/494_Sunday_June_05_2022_01_42_24_AM_46498650/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 16.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{x(x^2 + 1)}{4y^3} = 0$$

With initial conditions

$$\left[y(0) = -\frac{\sqrt{2}}{2} \right]$$

2.16.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= \frac{x(x^2 + 1)}{4y^3}\end{aligned}$$

The x domain of $f(x, y)$ when $y = -\frac{\sqrt{2}}{2}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = -\frac{\sqrt{2}}{2}$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{x(x^2 + 1)}{4y^3} \right) \\ &= -\frac{3x(x^2 + 1)}{4y^4}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = -\frac{\sqrt{2}}{2}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = -\frac{\sqrt{2}}{2}$ is inside this domain. Therefore solution exists and is unique.

2.16.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{x(x^2 + 1)}{4y^3}\end{aligned}$$

Where $f(x) = \frac{x(x^2+1)}{4}$ and $g(y) = \frac{1}{y^3}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{1}{y^3}} dy &= \frac{x(x^2+1)}{4} dx \\ \int \frac{1}{\frac{1}{y^3}} dy &= \int \frac{x(x^2+1)}{4} dx \\ \frac{y^4}{4} &= \frac{(x^2+1)^2}{16} + c_1\end{aligned}$$

The solution is

$$\frac{y^4}{4} - \frac{(x^2+1)^2}{16} - c_1 = 0$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -\frac{\sqrt{2}}{2}$ in the above solution gives an equation to solve for the constant of integration.

$$-c_1 = 0$$

$$c_1 = 0$$

Substituting c_1 found above in the general solution gives

$$\frac{1}{4}y^4 - \frac{1}{16}x^4 - \frac{1}{8}x^2 - \frac{1}{16} = 0$$

Summary

The solution(s) found are the following

$$\frac{y^4}{4} - \frac{x^4}{16} - \frac{x^2}{8} - \frac{1}{16} = 0 \tag{1}$$

Verification of solutions

$$\frac{y^4}{4} - \frac{x^4}{16} - \frac{x^2}{8} - \frac{1}{16} = 0$$

Verified OK.

2.16.3 Solving as differential Type ode

Writing the ode as

$$y' = \frac{x(x^2 + 1)}{4y^3} \quad (1)$$

Which becomes

$$(4y^3) dy = (x(x^2 + 1)) dx \quad (2)$$

But the RHS is complete differential because

$$(x(x^2 + 1)) dx = d\left(\frac{1}{2}x^2 + \frac{1}{4}x^4\right)$$

Hence (2) becomes

$$(4y^3) dy = d\left(\frac{1}{2}x^2 + \frac{1}{4}x^4\right)$$

Integrating both sides gives gives these solutions

$$y = \left(\frac{1}{2}x^2 + \frac{1}{4}x^4 + c_1\right)^{\frac{1}{4}} + c_1$$

$$y = i\left(\frac{1}{2}x^2 + \frac{1}{4}x^4 + c_1\right)^{\frac{1}{4}} + c_1$$

$$y = -\left(\frac{1}{2}x^2 + \frac{1}{4}x^4 + c_1\right)^{\frac{1}{4}} + c_1$$

$$y = -i\left(\frac{1}{2}x^2 + \frac{1}{4}x^4 + c_1\right)^{\frac{1}{4}} + c_1$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -\frac{\sqrt{2}}{2}$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{\sqrt{2}}{2} = -ic_1^{\frac{1}{4}} + c_1$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -\frac{\sqrt{2}}{2}$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{\sqrt{2}}{2} = -c_1^{\frac{1}{4}} + c_1$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -\frac{\sqrt{2}}{2}$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{\sqrt{2}}{2} = ic_1^{\frac{1}{4}} + c_1$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -\frac{\sqrt{2}}{2}$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{\sqrt{2}}{2} = c_1^{\frac{1}{4}} + c_1$$

Warning: Unable to solve for constant of integration.

Verification of solutions N/A

2.16.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x(x^2 + 1)}{4y^3}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 136: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{4}{(x^2 + 1)x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{4}{(x^2+1)x}} dx \end{aligned}$$

Which results in

$$S = \frac{(x^2 + 1)^2}{16}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x(x^2 + 1)}{4y^3}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= \frac{x(x^2 + 1)}{4} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y^3 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R^3$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^4}{4} + c_1 \quad (4)$$

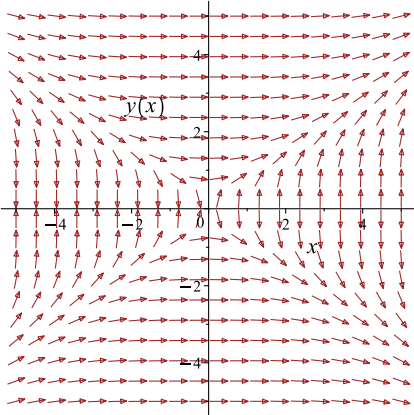
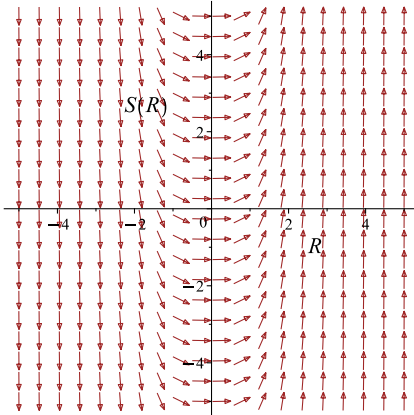
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{(x^2 + 1)^2}{16} = \frac{y^4}{4} + c_1$$

Which simplifies to

$$\frac{(x^2 + 1)^2}{16} = \frac{y^4}{4} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{x(x^2+1)}{4y^3}$ 	$R = y$ $S = \frac{(x^2 + 1)^2}{16}$	$\frac{dS}{dR} = R^3$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -\frac{\sqrt{2}}{2}$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{1}{16} = \frac{1}{16} + c_1$$

$$c_1 = 0$$

Substituting c_1 found above in the general solution gives

$$\frac{(x^2 + 1)^2}{16} = \frac{y^4}{4}$$

Summary

The solution(s) found are the following

$$\frac{(x^2 + 1)^2}{16} = \frac{y^4}{4} \quad (1)$$

Verification of solutions

$$\frac{(x^2 + 1)^2}{16} = \frac{y^4}{4}$$

Verified OK.

2.16.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (4y^3) dy &= (x(x^2 + 1)) dx \\ (-x(x^2 + 1)) dx + (4y^3) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x(x^2 + 1) \\ N(x, y) &= 4y^3 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x(x^2 + 1)) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(4y^3) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x(x^2 + 1) dx \\ \phi &= -\frac{(x^2 + 1)^2}{4} + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y)\quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 4y^3$. Therefore equation (4) becomes

$$4y^3 = 0 + f'(y)\quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 4y^3$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (4y^3) dy \\ f(y) &= y^4 + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{(x^2 + 1)^2}{4} + y^4 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{(x^2 + 1)^2}{4} + y^4$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -\frac{\sqrt{2}}{2}$ in the above solution gives an equation to solve for the constant of integration.

$$0 = c_1$$

$$c_1 = 0$$

Substituting c_1 found above in the general solution gives

$$-\frac{(x^2 + 1)^2}{4} + y^4 = 0$$

Summary

The solution(s) found are the following

$$-\frac{(x^2 + 1)^2}{4} + y^4 = 0 \tag{1}$$

Verification of solutions

$$-\frac{(x^2 + 1)^2}{4} + y^4 = 0$$

Verified OK.

2.16.6 Maple step by step solution

Let's solve

$$\left[y' - \frac{x(x^2+1)}{4y^3} = 0, y(0) = -\frac{\sqrt{2}}{2} \right]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$y'y^3 = \frac{x(x^2+1)}{4}$$

- Integrate both sides with respect to x

$$\int y'y^3 dx = \int \frac{x(x^2+1)}{4} dx + c_1$$

- Evaluate integral

$$\frac{y^4}{4} = \frac{(x^2+1)^2}{16} + c_1$$

- Solve for y

$$\left\{ y = \left(\frac{1}{4}x^4 + \frac{1}{2}x^2 + 4c_1 + \frac{1}{4}\right)^{\frac{1}{4}}, y = -\left(\frac{1}{4}x^4 + \frac{1}{2}x^2 + 4c_1 + \frac{1}{4}\right)^{\frac{1}{4}} \right\}$$

- Use initial condition $y(0) = -\frac{\sqrt{2}}{2}$
- $-\frac{\sqrt{2}}{2} = \left(\frac{1}{4} + 4c_1\right)^{\frac{1}{4}}$
- Solution does not satisfy initial condition
- Use initial condition $y(0) = -\frac{\sqrt{2}}{2}$
- $-\frac{\sqrt{2}}{2} = -\left(\frac{1}{4} + 4c_1\right)^{\frac{1}{4}}$
- Solve for c_1
- $c_1 = 0$
- Substitute $c_1 = 0$ into general solution and simplify

$$y = -\frac{\sqrt{2}((x^2+1)^2)^{\frac{1}{4}}}{2}$$

- Solution to the IVP

$$y = -\frac{\sqrt{2}((x^2+1)^2)^{\frac{1}{4}}}{2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`

```

✓ Solution by Maple

Time used: 0.078 (sec). Leaf size: 15

```
dsolve([diff(y(x),x) = 1/4*x*(x^2+1)/y(x)^3,y(0) = -1/sqrt(2)],y(x), singsol=all)
```

$$y(x) = -\frac{\sqrt{2x^2 + 2}}{2}$$

✓ Solution by Mathematica

Time used: 0.23 (sec). Leaf size: 23

```
DSolve[{y'[x] == 1/4*x*(x^2+1)/y[x]^3, y[0]==-(1/Sqrt[2])}, y[x], x, IncludeSingularSolutions ->
```

$$y(x) \rightarrow -\frac{\sqrt[4]{(x^2 + 1)^2}}{\sqrt{2}}$$

2.17 problem 17

2.17.1 Existence and uniqueness analysis	635
2.17.2 Solving as separable ode	636
2.17.3 Solving as first order ode lie symmetry lookup ode	638
2.17.4 Solving as exact ode	642
2.17.5 Maple step by step solution	645

Internal problem ID [495]

Internal file name [OUTPUT/495_Sunday_June_05_2022_01_42_26_AM_66277094/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 17.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{-e^x + 3x^2}{-5 + 2y} = 0$$

With initial conditions

$$[y(0) = 1]$$

2.17.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned} y' &= f(x, y) \\ &= -\frac{-3x^2 + e^x}{-5 + 2y} \end{aligned}$$

The x domain of $f(x, y)$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\left\{ y < \frac{5}{2} \vee \frac{5}{2} < y \right\}$$

And the point $y_0 = 1$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned} \frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{3x^2 + e^x}{-5 + 2y} \right) \\ &= \frac{-6x^2 + 2e^x}{(-5 + 2y)^2} \end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\left\{ y < \frac{5}{2} \vee \frac{5}{2} < y \right\}$$

And the point $y_0 = 1$ is inside this domain. Therefore solution exists and is unique.

2.17.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{-e^x + 3x^2}{-5 + 2y} \end{aligned}$$

Where $f(x) = -e^x + 3x^2$ and $g(y) = \frac{1}{-5+2y}$. Integrating both sides gives

$$\begin{aligned} \frac{1}{-5+2y} dy &= -e^x + 3x^2 dx \\ \int \frac{1}{-5+2y} dy &= \int -e^x + 3x^2 dx \\ y^2 - 5y &= x^3 - e^x + c_1 \end{aligned}$$

Which results in

$$y = \frac{5}{2} + \frac{\sqrt{25 + 4x^3 - 4e^x + 4c_1}}{2}$$

$$y = \frac{5}{2} - \frac{\sqrt{25 + 4x^3 - 4e^x + 4c_1}}{2}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{5}{2} - \frac{\sqrt{21 + 4c_1}}{2}$$

$$c_1 = -3$$

Substituting c_1 found above in the general solution gives

$$y = \frac{5}{2} - \frac{\sqrt{13 + 4x^3 - 4e^x}}{2}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

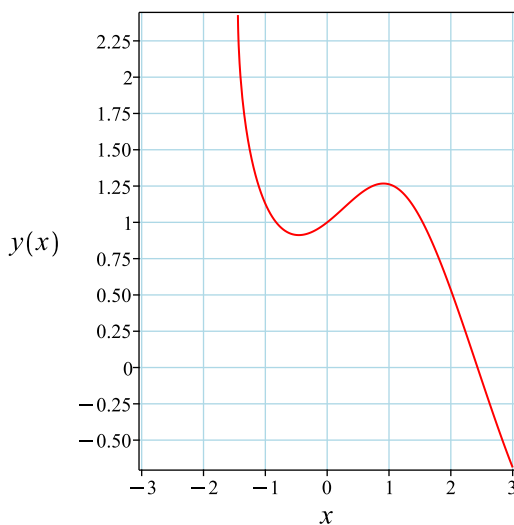
$$1 = \frac{5}{2} + \frac{\sqrt{21 + 4c_1}}{2}$$

Summary

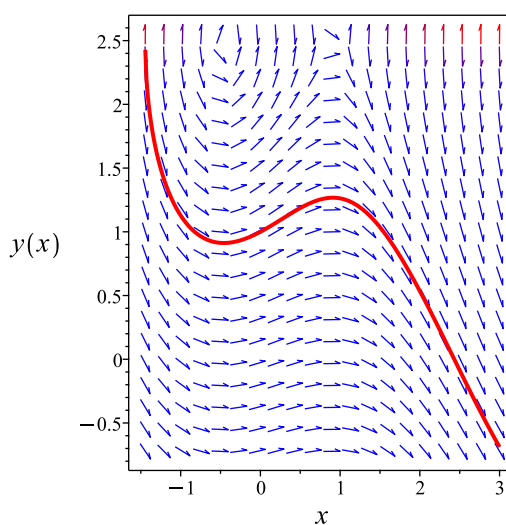
The solution(s) found are the following

Warning: Unable to solve for constant of integration.

$$y = \frac{5}{2} - \frac{\sqrt{13 + 4x^3 - 4e^x}}{2}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{5}{2} - \frac{\sqrt{13 + 4x^3 - 4e^x}}{2}$$

Verified OK.

2.17.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{-3x^2 + e^x}{-5 + 2y}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 139: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{-e^x + 3x^2} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{-e^x + 3x^2}} dx \end{aligned}$$

Which results in

$$S = x^3 - e^x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{-3x^2 + e^x}{-5 + 2y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= -e^x + 3x^2 \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -5 + 2y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -5 + 2R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R^2 - 5R + c_1 \quad (4)$$

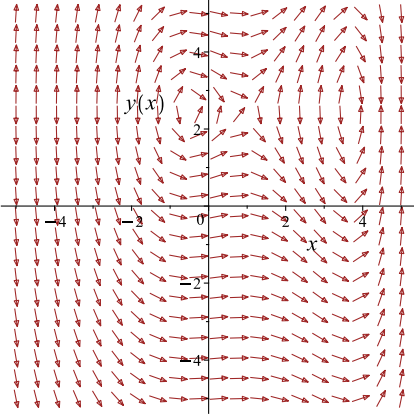
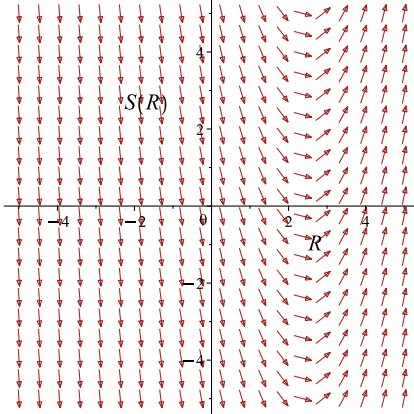
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$x^3 - e^x = y^2 + c_1 - 5y$$

Which simplifies to

$$x^3 - e^x = y^2 + c_1 - 5y$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{-3x^2 + e^x}{-5 + 2y}$ 	$R = y$ $S = x^3 - e^x$	$\frac{dS}{dR} = -5 + 2R$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$-1 = -4 + c_1$$

$$c_1 = 3$$

Substituting c_1 found above in the general solution gives

$$x^3 - e^x = y^2 - 5y + 3$$

Summary

The solution(s) found are the following

$$x^3 - e^x = y^2 - 5y + 3 \quad (1)$$

Verification of solutions

$$x^3 - e^x = y^2 - 5y + 3$$

Verified OK.

2.17.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work

and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (5 - 2y) dy &= (-3x^2 + e^x) dx \\ (-e^x + 3x^2) dx + (5 - 2y) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -e^x + 3x^2 \\ N(x, y) &= 5 - 2y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-e^x + 3x^2) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(5 - 2y) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -e^x + 3x^2 dx \\ \phi &= x^3 - e^x + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial\phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial\phi}{\partial y} = 5 - 2y$. Therefore equation (4) becomes

$$5 - 2y = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 5 - 2y$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) \, dy &= \int (5 - 2y) \, dy \\ f(y) &= -y^2 + 5y + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = x^3 - y^2 - e^x + 5y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = x^3 - y^2 - e^x + 5y$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$3 = c_1$$

$$c_1 = 3$$

Substituting c_1 found above in the general solution gives

$$x^3 - y^2 - e^x + 5y = 3$$

Summary

The solution(s) found are the following

$$x^3 - y^2 - e^x + 5y = 3 \quad (1)$$

Verification of solutions

$$x^3 - y^2 - e^x + 5y = 3$$

Verified OK.

2.17.5 Maple step by step solution

Let's solve

$$\left[y' - \frac{-e^x + 3x^2}{-5 + 2y} = 0, y(0) = 1 \right]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$y'(-5 + 2y) = -e^x + 3x^2$$

- Integrate both sides with respect to x

$$\int y'(-5 + 2y) dx = \int (-e^x + 3x^2) dx + c_1$$

- Evaluate integral

$$y^2 - 5y = x^3 - e^x + c_1$$

- Solve for y

$$\left\{ y = \frac{5}{2} - \frac{\sqrt{25 + 4x^3 - 4e^x + 4c_1}}{2}, y = \frac{5}{2} + \frac{\sqrt{25 + 4x^3 - 4e^x + 4c_1}}{2} \right\}$$

- Use initial condition $y(0) = 1$

$$1 = \frac{5}{2} - \frac{\sqrt{21 + 4c_1}}{2}$$

- Solve for c_1

$$c_1 = -3$$

- Substitute $c_1 = -3$ into general solution and simplify

$$y = \frac{5}{2} - \frac{\sqrt{13 + 4x^3 - 4e^x}}{2}$$

- Use initial condition $y(0) = 1$

$$1 = \frac{5}{2} + \frac{\sqrt{21 + 4c_1}}{2}$$

- Solution does not satisfy initial condition

- Solution to the IVP

$$y = \frac{5}{2} - \frac{\sqrt{13+4x^3-4e^x}}{2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.172 (sec). Leaf size: 21

```
dsolve([diff(y(x),x) = (-exp(x)+3*x^2)/(-5+2*y(x)),y(0) = 1],y(x), singsol=all)
```

$$y(x) = \frac{5}{2} - \frac{\sqrt{13 + 4x^3 - 4e^x}}{2}$$

✓ Solution by Mathematica

Time used: 0.891 (sec). Leaf size: 29

```
DSolve[{y'[x] == (-Exp[x]+3*x^2)/(-5+2*y[x]),y[0]==1},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2} \left(5 - \sqrt{4x^3 - 4e^x + 13} \right)$$

2.18 problem 18

2.18.1 Existence and uniqueness analysis	647
2.18.2 Solving as separable ode	648
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Internal problem ID [496]

Internal file name [OUTPUT/496_Sunday_June_05_2022_01_42_27_AM_16414608/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 18.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{-e^x + e^{-x}}{3 + 4y} = 0$$

With initial conditions

$$[y(0) = 1]$$

2.18.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned} y' &= f(x, y) \\ &= -\frac{e^x - e^{-x}}{3 + 4y} \end{aligned}$$

The x domain of $f(x, y)$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{e^x - e^{-x}}{3 + 4y} \right) \\ &= \frac{4e^x - 4e^{-x}}{(3 + 4y)^2}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Therefore solution exists and is unique.

2.18.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{-e^x + e^{-x}}{3 + 4y}\end{aligned}$$

Where $f(x) = -e^x + e^{-x}$ and $g(y) = \frac{1}{3+4y}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{3+4y} dy &= -e^x + e^{-x} dx \\ \int \frac{1}{3+4y} dy &= \int -e^x + e^{-x} dx \\ 2y^2 + 3y &= -e^x - e^{-x} + c_1\end{aligned}$$

Which results in

$$y = -\frac{(3e^x - \sqrt{-8e^{3x} + 8c_1e^{2x} + 9e^{2x} - 8e^x})e^{-x}}{4}$$

$$y = -\frac{(3e^x + \sqrt{-8e^{3x} + 8c_1e^{2x} + 9e^{2x} - 8e^x})e^{-x}}{4}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{3}{4} - \frac{\sqrt{-7 + 8c_1}}{4}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{3}{4} + \frac{\sqrt{-7 + 8c_1}}{4}$$

$$c_1 = 7$$

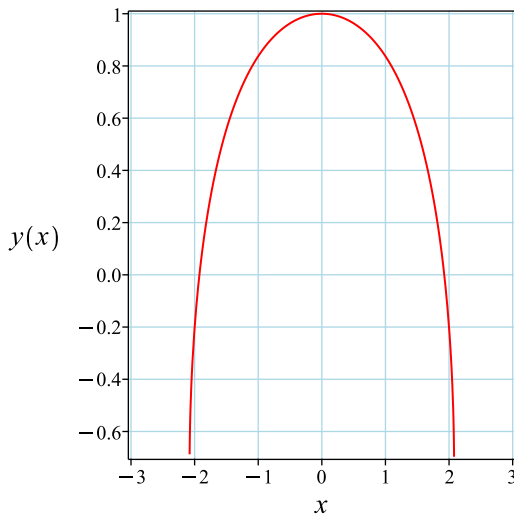
Substituting c_1 found above in the general solution gives

$$y = -\frac{3}{4} + \frac{e^{-x}\sqrt{65e^{2x} - 8e^xe^{2x} - 8e^x}}{4}$$

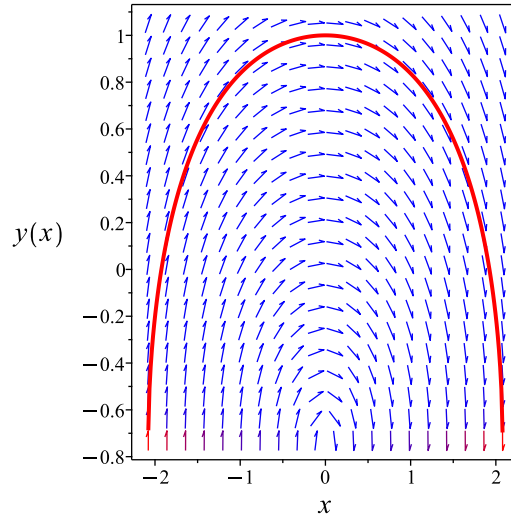
Summary

The solution(s) found are the following

$$y = -\frac{3}{4} + \frac{e^{-x}\sqrt{65e^{2x} - 8e^xe^{2x} - 8e^x}}{4} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{3}{4} + \frac{e^{-x}\sqrt{65e^{2x} - 8e^xe^{2x} - 8e^x}}{4}$$

Verified OK.

2.18.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{e^x - e^{-x}}{3 + 4y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2\xi_y - \omega_x\xi - \omega_y\eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 142: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{-e^x + e^{-x}} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{-e^x + e^{-x}}} dx \end{aligned}$$

Which results in

$$S = -e^x - e^{-x}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{e^x - e^{-x}}{3 + 4y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= -e^x + e^{-x} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 3 + 4y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 3 + 4R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = 2R^2 + 3R + c_1 \quad (4)$$

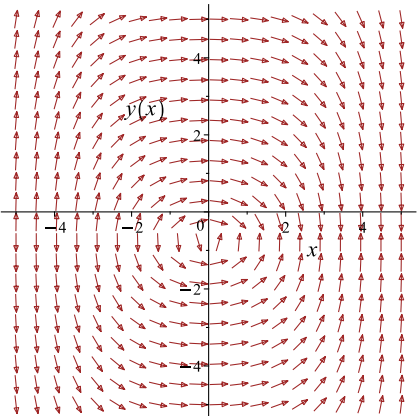
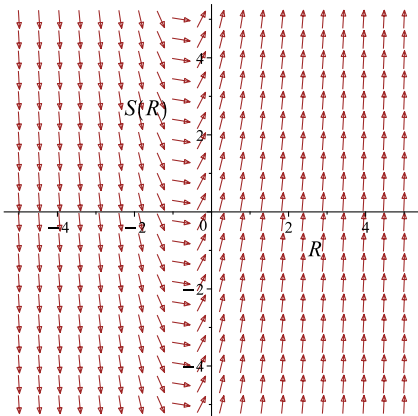
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-e^x - e^{-x} = 2y^2 + c_1 + 3y$$

Which simplifies to

$$-e^x - e^{-x} = 2y^2 + c_1 + 3y$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{e^x - e^{-x}}{3 + 4y}$ 	$R = y$ $S = -e^x - e^{-x}$	$\frac{dS}{dR} = 3 + 4R$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = 5 + c_1$$

$$c_1 = -7$$

Substituting c_1 found above in the general solution gives

$$-e^x - e^{-x} = 2y^2 + 3y - 7$$

Summary

The solution(s) found are the following

$$-e^x - e^{-x} = 2y^2 + 3y - 7 \quad (1)$$

Verification of solutions

$$-e^x - e^{-x} = 2y^2 + 3y - 7$$

Verified OK.

2.18.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work

and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (-3 - 4y) dy &= (e^x - e^{-x}) dx \\ (-e^x + e^{-x}) dx + (-3 - 4y) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -e^x + e^{-x} \\ N(x, y) &= -3 - 4y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (-e^x + e^{-x}) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (-3 - 4y) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -e^x + e^{-x} dx \\ \phi &= -e^x - e^{-x} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -3 - 4y$. Therefore equation (4) becomes

$$-3 - 4y = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -3 - 4y$$

Integrating the above w.r.t y gives

$$\int f'(y) \, dy = \int (-3 - 4y) \, dy$$

$$f(y) = -2y^2 - 3y + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -2y^2 - e^x - 3y - e^{-x} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -2y^2 - e^x - 3y - e^{-x}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$-7 = c_1$$

$$c_1 = -7$$

Substituting c_1 found above in the general solution gives

$$-2y^2 - e^x - 3y - e^{-x} = -7$$

Summary

The solution(s) found are the following

$$-2y^2 - e^x - 3y - e^{-x} = -7 \quad (1)$$

Verification of solutions

$$-2y^2 - e^x - 3y - e^{-x} = -7$$

Verified OK.

2.18.5 Maple step by step solution

Let's solve

$$\left[y' - \frac{-e^x + e^{-x}}{3 + 4y} = 0, y(0) = 1 \right]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$y'(3 + 4y) = -e^x + e^{-x}$$

- Integrate both sides with respect to x

$$\int y'(3 + 4y) dx = \int (-e^x + e^{-x}) dx + c_1$$

- Evaluate integral

$$2y^2 + 3y = -e^x - e^{-x} + c_1$$

- Solve for y

$$\left\{ y = -\frac{3e^x - \sqrt{-8(e^x)^3 + 8c_1(e^x)^2 + 9(e^x)^2 - 8e^x}}{4e^x}, y = -\frac{3e^x + \sqrt{-8(e^x)^3 + 8c_1(e^x)^2 + 9(e^x)^2 - 8e^x}}{4e^x} \right\}$$

- Use initial condition $y(0) = 1$

$$1 = -\frac{3}{4} + \frac{\sqrt{-7 + 8c_1}}{4}$$

- Solve for c_1

$$c_1 = 7$$

- Substitute $c_1 = 7$ into general solution and simplify

$$y = -\frac{3}{4} + \frac{e^{-x} \sqrt{(-8e^{2x} + 65e^x - 8)e^x}}{4}$$

- Use initial condition $y(0) = 1$

$$1 = -\frac{3}{4} - \frac{\sqrt{-7 + 8c_1}}{4}$$

- Solution does not satisfy initial condition
- Solution to the IVP

$$y = -\frac{3}{4} + \frac{e^{-x}\sqrt{(-8e^{2x}+65e^x-8)e^x}}{4}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.235 (sec). Leaf size: 29

```
dsolve([diff(y(x),x) = (exp(-x)-exp(x))/(3+4*y(x)),y(0) = 1],y(x), singsol=all)
```

$$y(x) = -\frac{3}{4} + \frac{\sqrt{e^x(-8e^{2x} + 65e^x - 8)}e^{-x}}{4}$$

✓ Solution by Mathematica

Time used: 1.347 (sec). Leaf size: 29

```
DSolve[{y'[x] == (Exp[-x]-Exp[x])/(3+4*y[x]),y[0]==1},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{4} \left(\sqrt{-8e^{-x} - 8e^x + 65} - 3 \right)$$

2.19 problem 19

2.19.1 Existence and uniqueness analysis	659
2.19.2 Solving as separable ode	660
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Internal problem ID [497]

Internal file name [OUTPUT/497_Sunday_June_05_2022_01_42_28_AM_60921212/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 19.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$\cos(3y)y' = -\sin(2x)$$

With initial conditions

$$\left[y\left(\frac{\pi}{2}\right) = 0 \right]$$

2.19.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned} y' &= f(x, y) \\ &= -\frac{\sin(2x)}{\cos(3y)} \end{aligned}$$

The x domain of $f(x, y)$ when $y = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = \frac{\pi}{2}$ is inside this domain. The y domain of $f(x, y)$ when $x = \frac{\pi}{2}$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 0$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{\sin(2x)}{\cos(3y)} \right) \\ &= -\frac{3 \sin(2x) \sin(3y)}{\cos(3y)^2}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = \frac{\pi}{2}$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = \frac{\pi}{2}$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 0$ is inside this domain. Therefore solution exists and is unique.

2.19.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\sin(2x) \sec(3y)\end{aligned}$$

Where $f(x) = -\sin(2x)$ and $g(y) = \sec(3y)$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\sec(3y)} dy &= -\sin(2x) dx \\ \int \frac{1}{\sec(3y)} dy &= \int -\sin(2x) dx \\ \frac{\sin(3y)}{3} &= \frac{\cos(2x)}{2} + c_1\end{aligned}$$

Which results in

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + 3c_1\right)}{3}$$

Initial conditions are used to solve for c_1 . Substituting $x = \frac{\pi}{2}$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{\arcsin\left(-\frac{3}{2} + 3c_1\right)}{3}$$

$$c_1 = \frac{1}{2}$$

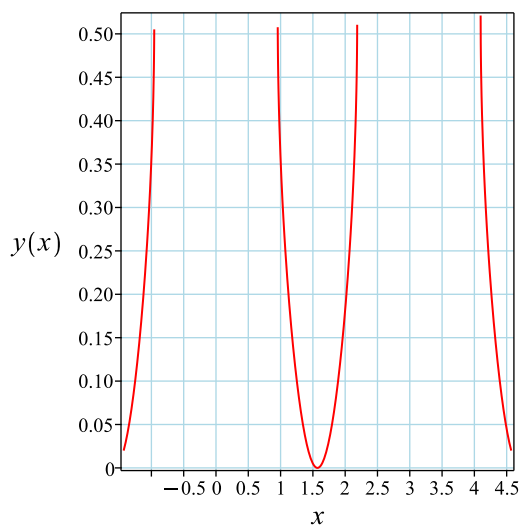
Substituting c_1 found above in the general solution gives

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3}$$

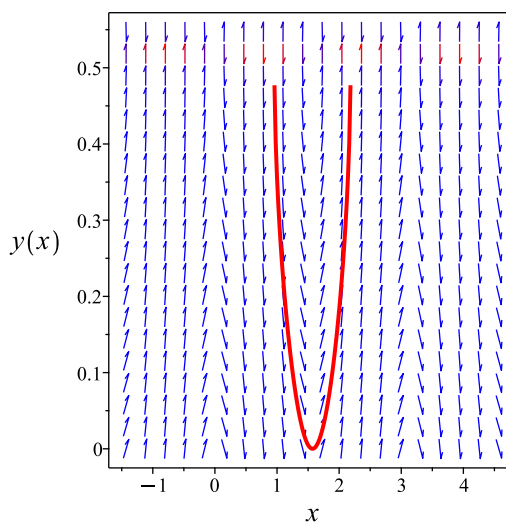
Summary

The solution(s) found are the following

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3}$$

Verified OK.

2.19.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{\sin(2x)}{\cos(3y)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 145: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= -\frac{1}{\sin(2x)} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{-\frac{1}{\sin(2x)}} dx\end{aligned}$$

Which results in

$$S = \frac{\cos(2x)}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{\sin(2x)}{\cos(3y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= -\sin(2x) \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \cos(3y) \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \cos(3R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{\sin(3R)}{3} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\cos(2x)}{2} = \frac{\sin(3y)}{3} + c_1$$

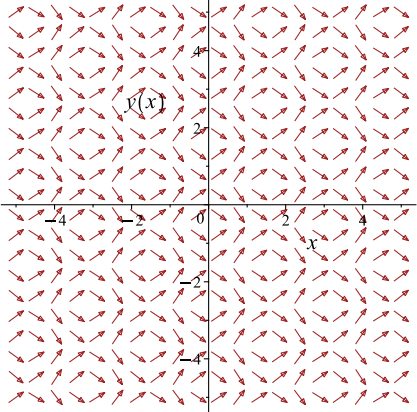
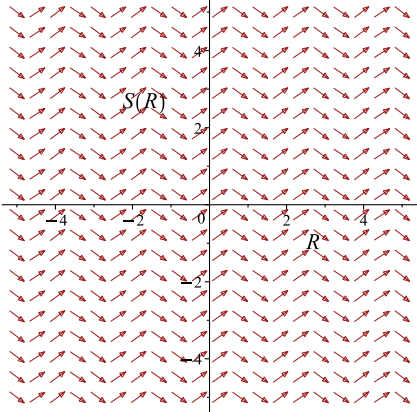
Which simplifies to

$$\frac{\cos(2x)}{2} = \frac{\sin(3y)}{3} + c_1$$

Which gives

$$y = -\frac{\arcsin\left(-\frac{3\cos(2x)}{2} + 3c_1\right)}{3}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{\sin(2x)}{\cos(3y)}$ 	$R = y$ $S = \frac{\cos(2x)}{2}$	$\frac{dS}{dR} = \cos(3R)$ 

Initial conditions are used to solve for c_1 . Substituting $x = \frac{\pi}{2}$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{\arcsin\left(\frac{3}{2} + 3c_1\right)}{3}$$

$$c_1 = -\frac{1}{2}$$

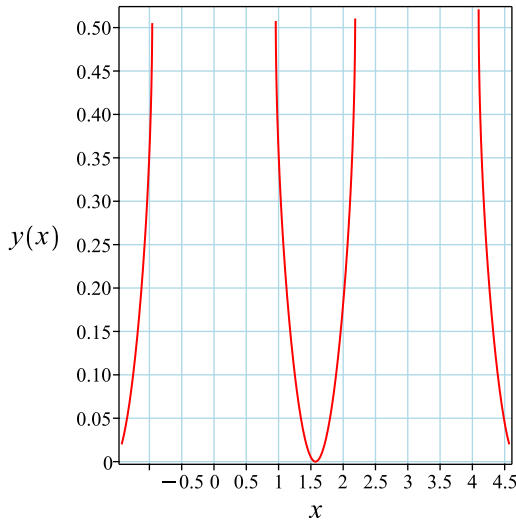
Substituting c_1 found above in the general solution gives

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3}$$

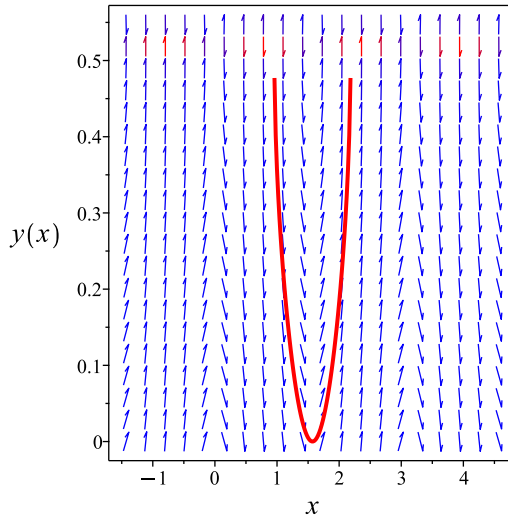
Summary

The solution(s) found are the following

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3}$$

Verified OK.

2.19.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (-\cos(3y)) dy &= (\sin(2x)) dx \\ (-\sin(2x)) dx + (-\cos(3y)) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -\sin(2x) \\ N(x, y) &= -\cos(3y) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-\sin(2x)) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-\cos(3y)) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\sin(2x) dx \\ \phi &= \frac{\cos(2x)}{2} + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y)\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\cos(3y)$. Therefore equation (4) becomes

$$-\cos(3y) = 0 + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\cos(3y)$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (-\cos(3y)) dy \\ f(y) &= -\frac{\sin(3y)}{3} + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{\cos(2x)}{2} - \frac{\sin(3y)}{3} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{\cos(2x)}{2} - \frac{\sin(3y)}{3}$$

Initial conditions are used to solve for c_1 . Substituting $x = \frac{\pi}{2}$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{1}{2} = c_1$$

$$c_1 = -\frac{1}{2}$$

Substituting c_1 found above in the general solution gives

$$\frac{\cos(2x)}{2} - \frac{\sin(3y)}{3} = -\frac{1}{2}$$

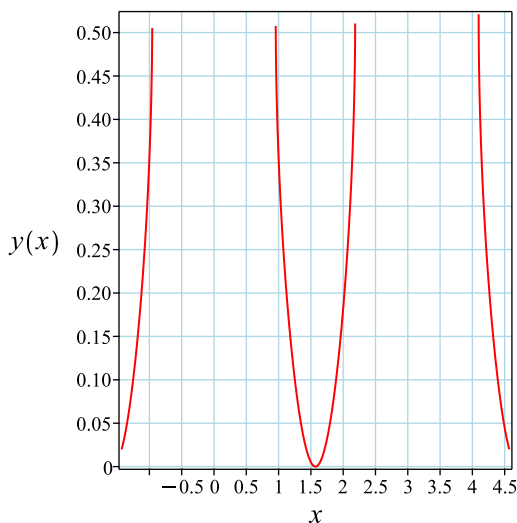
Solving for y from the above gives

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3}$$

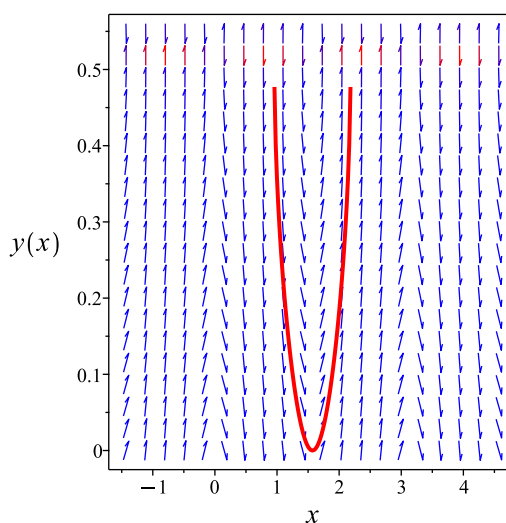
Summary

The solution(s) found are the following

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3}$$

Verified OK.

2.19.5 Maple step by step solution

Let's solve

$$[\cos(3y)y' = -\sin(2x), y(\frac{\pi}{2}) = 0]$$

- Highest derivative means the order of the ODE is 1

y'

- Integrate both sides with respect to x

$$\int \cos(3y)y'dx = \int -\sin(2x)dx + c_1$$

- Evaluate integral

$$\frac{\sin(3y)}{3} = \frac{\cos(2x)}{2} + c_1$$

- Solve for y

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + 3c_1\right)}{3}$$

- Use initial condition $y(\frac{\pi}{2}) = 0$

$$0 = \frac{\arcsin\left(-\frac{3}{2} + 3c_1\right)}{3}$$

- Solve for c_1

$$c_1 = \frac{1}{2}$$

- Substitute $c_1 = \frac{1}{2}$ into general solution and simplify

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3}$$

- Solution to the IVP

$$y = \frac{\arcsin\left(\frac{3\cos(2x)}{2} + \frac{3}{2}\right)}{3}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.297 (sec). Leaf size: 15

```
dsolve([sin(2*x)+cos(3*y(x))*diff(y(x),x) = 0,y(1/2*Pi) = 0],y(x), singsol=all)
```

$$y(x) = \frac{\arcsin\left(\frac{3}{2} + \frac{3\cos(2x)}{2}\right)}{3}$$

✓ Solution by Mathematica

Time used: 0.614 (sec). Leaf size: 16

```
DSolve[{Sin[2*x]+Cos[3*y[x]]*y'[x] == 0,y[Pi/2]==0},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{3} \arcsin(3 \cos^2(x))$$

2.20 problem 20

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Internal problem ID [498]

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Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 20.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$\sqrt{-x^2 + 1} y^2 y' = \arcsin(x)$$

With initial conditions

$$[y(0) = 1]$$

2.20.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned} y' &= f(x, y) \\ &= \frac{\arcsin(x)}{\sqrt{-x^2 + 1} y^2} \end{aligned}$$

The x domain of $f(x, y)$ when $y = 1$ is

$$\{-1 < x < 1\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{\arcsin(x)}{\sqrt{-x^2+1}y^2} \right) \\ &= -\frac{2\arcsin(x)}{\sqrt{-x^2+1}y^3}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 1$ is

$$\{-1 < x < 1\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Therefore solution exists and is unique.

2.20.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{\arcsin(x)}{\sqrt{-x^2+1}y^2}\end{aligned}$$

Where $f(x) = \frac{\arcsin(x)}{\sqrt{-x^2+1}}$ and $g(y) = \frac{1}{y^2}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{1}{y^2}} dy &= \frac{\arcsin(x)}{\sqrt{-x^2+1}} dx \\ \int \frac{1}{\frac{1}{y^2}} dy &= \int \frac{\arcsin(x)}{\sqrt{-x^2+1}} dx \\ \frac{y^3}{3} &= \frac{\arcsin(x)^2}{2} + c_1\end{aligned}$$

Which results in

$$y = \frac{(12 \arcsin(x)^2 + 24c_1)^{\frac{1}{3}}}{2}$$

$$y = -\frac{(12 \arcsin(x)^2 + 24c_1)^{\frac{1}{3}}}{4} + \frac{i\sqrt{3}(12 \arcsin(x)^2 + 24c_1)^{\frac{1}{3}}}{4}$$

$$y = -\frac{(12 \arcsin(x)^2 + 24c_1)^{\frac{1}{3}}}{4} - \frac{i\sqrt{3}(12 \arcsin(x)^2 + 24c_1)^{\frac{1}{3}}}{4}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{ic_1^{\frac{1}{3}}24^{\frac{1}{3}}\sqrt{3}}{4} - \frac{c_1^{\frac{1}{3}}24^{\frac{1}{3}}}{4}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{ic_1^{\frac{1}{3}}24^{\frac{1}{3}}\sqrt{3}}{4} - \frac{c_1^{\frac{1}{3}}24^{\frac{1}{3}}}{4}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{c_1^{\frac{1}{3}}24^{\frac{1}{3}}}{2}$$

$$c_1 = \frac{1}{3}$$

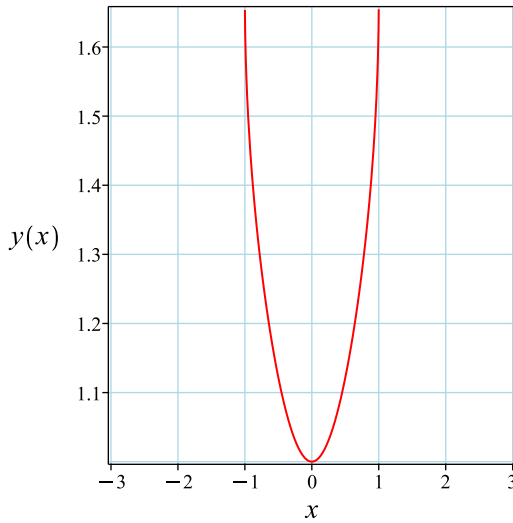
Substituting c_1 found above in the general solution gives

$$y = \frac{(12 \arcsin(x)^2 + 8)^{\frac{1}{3}}}{2}$$

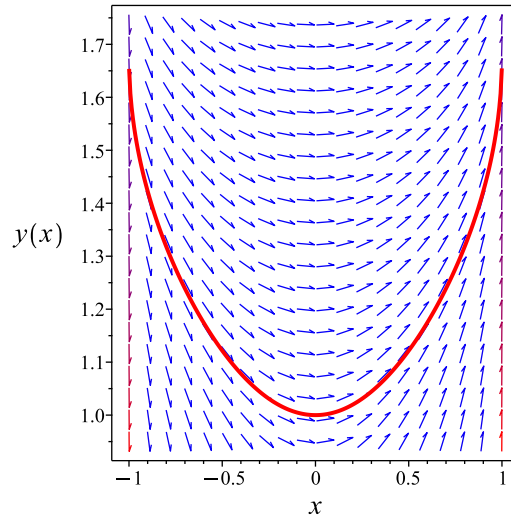
Summary

The solution(s) found are the following

$$y = \frac{(12 \arcsin(x)^2 + 8)^{\frac{1}{3}}}{2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{(12 \arcsin(x)^2 + 8)^{\frac{1}{3}}}{2}$$

Verified OK.

2.20.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{\arcsin(x)}{\sqrt{-x^2 + 1} y^2}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 148: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{\sqrt{-x^2 + 1}}{\arcsin(x)} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{\sqrt{-x^2+1}}{\arcsin(x)}} dx \end{aligned}$$

Which results in

$$S = \frac{\arcsin(x)^2}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{\arcsin(x)}{\sqrt{-x^2+1}y^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= \frac{\arcsin(x)}{\sqrt{-x^2+1}} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y^2 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R^2$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^3}{3} + c_1 \quad (4)$$

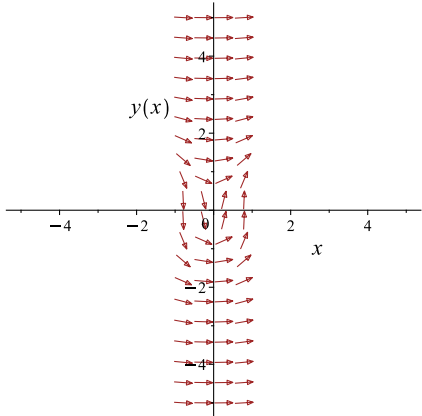
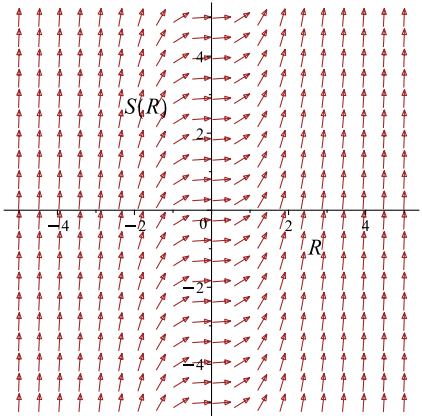
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\arcsin(x)^2}{2} = \frac{y^3}{3} + c_1$$

Which simplifies to

$$\frac{\arcsin(x)^2}{2} = \frac{y^3}{3} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{\arcsin(x)}{\sqrt{-x^2+1}y^2}$ 	$R = y$ $S = \frac{\arcsin(x)^2}{2}$	$\frac{dS}{dR} = R^2$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{1}{3} + c_1$$

$$c_1 = -\frac{1}{3}$$

Substituting c_1 found above in the general solution gives

$$\frac{\arcsin(x)^2}{2} = \frac{y^3}{3} - \frac{1}{3}$$

Summary

The solution(s) found are the following

$$\frac{\arcsin(x)^2}{2} = \frac{y^3}{3} - \frac{1}{3} \quad (1)$$

Verification of solutions

$$\frac{\arcsin(x)^2}{2} = \frac{y^3}{3} - \frac{1}{3}$$

Verified OK.

2.20.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (y^2) dy &= \left(\frac{\arcsin(x)}{\sqrt{-x^2 + 1}} \right) dx \\ \left(-\frac{\arcsin(x)}{\sqrt{-x^2 + 1}} \right) dx + (y^2) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -\frac{\arcsin(x)}{\sqrt{-x^2 + 1}} \\ N(x, y) &= y^2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{\arcsin(x)}{\sqrt{-x^2 + 1}} \right) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (y^2) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{\arcsin(x)}{\sqrt{-x^2+1}} dx \\ \phi &= -\frac{\arcsin(x)^2}{2} + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = y^2$. Therefore equation (4) becomes

$$y^2 = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = y^2$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (y^2) dy \\ f(y) &= \frac{y^3}{3} + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{\arcsin(x)^2}{2} + \frac{y^3}{3} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{\arcsin(x)^2}{2} + \frac{y^3}{3}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{1}{3} = c_1$$

$$c_1 = \frac{1}{3}$$

Substituting c_1 found above in the general solution gives

$$-\frac{\arcsin(x)^2}{2} + \frac{y^3}{3} = \frac{1}{3}$$

Summary

The solution(s) found are the following

$$-\frac{\arcsin(x)^2}{2} + \frac{y^3}{3} = \frac{1}{3} \quad (1)$$

Verification of solutions

$$-\frac{\arcsin(x)^2}{2} + \frac{y^3}{3} = \frac{1}{3}$$

Verified OK.

2.20.5 Maple step by step solution

Let's solve

$$[\sqrt{-x^2 + 1} y^2 y' = \arcsin(x), y(0) = 1]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$y' y^2 = \frac{\arcsin(x)}{\sqrt{-x^2+1}}$$

- Integrate both sides with respect to x

$$\int y' y^2 dx = \int \frac{\arcsin(x)}{\sqrt{-x^2+1}} dx + c_1$$

- Evaluate integral

$$\frac{y^3}{3} = \frac{\arcsin(x)^2}{2} + c_1$$

- Solve for y

$$y = \frac{(12 \arcsin(x)^2 + 24c_1)^{\frac{1}{3}}}{2}$$

- Use initial condition $y(0) = 1$

$$1 = \frac{c_1^{\frac{1}{3}} 24^{\frac{1}{3}}}{2}$$

- Solve for c_1

$$c_1 = \frac{1}{3}$$

- Substitute $c_1 = \frac{1}{3}$ into general solution and simplify

$$y = \frac{(12 \arcsin(x)^2 + 8)^{\frac{1}{3}}}{2}$$

- Solution to the IVP

$$y = \frac{(12 \arcsin(x)^2 + 8)^{\frac{1}{3}}}{2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`

```

✓ Solution by Maple

Time used: 0.109 (sec). Leaf size: 16

```
dsolve([(-x^2+1)^(1/2)*y(x)^2*diff(y(x),x) = arcsin(x),y(0) = 1],y(x), singsol=all)
```

$$y(x) = \frac{(8 + 12 \arcsin(x)^2)^{\frac{1}{3}}}{2}$$

✓ Solution by Mathematica

Time used: 0.527 (sec). Leaf size: 19

```
DSolve[{-x^2+1)^(1/2)*y[x]^2*y'[x] == ArcSin[x],y[0]==1},y[x],x,IncludeSingularSolutions ->
```

$$y(x) \rightarrow \sqrt[3]{\frac{3 \arcsin(x)^2}{2} + 1}$$

2.21 problem 21

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Internal problem ID [499]

Internal file name [OUTPUT/499_Sunday_June_05_2022_01_42_31_AM_97879552/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 21.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{3x^2 + 1}{-6y + 3y^2} = 0$$

With initial conditions

$$[y(0) = 1]$$

2.21.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= \frac{3x^2 + 1}{3y(y - 2)}\end{aligned}$$

The x domain of $f(x, y)$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty \leq y < 0, 0 < y < 2, 2 < y \leq \infty\}$$

And the point $y_0 = 1$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{3x^2 + 1}{3y(y - 2)} \right) \\ &= -\frac{3x^2 + 1}{3y^2(y - 2)} - \frac{3x^2 + 1}{3y(y - 2)^2}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty \leq y < 0, 0 < y < 2, 2 < y \leq \infty\}$$

And the point $y_0 = 1$ is inside this domain. Therefore solution exists and is unique.

2.21.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{x^2 + \frac{1}{3}}{y(y - 2)}\end{aligned}$$

Where $f(x) = x^2 + \frac{1}{3}$ and $g(y) = \frac{1}{y(y-2)}$. Integrating both sides gives

$$\frac{1}{\frac{1}{y(y-2)}} dy = x^2 + \frac{1}{3} dx$$

$$\int \frac{1}{\frac{1}{y(y-2)}} dy = \int x^2 + \frac{1}{3} dx$$

$$\frac{1}{3}y^3 - y^2 = \frac{1}{3}x^3 + \frac{1}{3}x + c_1$$

Which results in

$$y = \frac{\left(8 + 4x^3 + 12c_1 + 4x + 4\sqrt{x^6 + 6c_1x^3 + 2x^4 + 4x^3 + 9c_1^2 + 6c_1x + x^2 + 12c_1 + 4x}\right)^{\frac{1}{3}}}{2} + \frac{\left(8 + 4x^3 + 12c_1 + 4x + 4\sqrt{x^6 + 6c_1x^3 + 2x^4 + 4x^3 + 9c_1^2 + 6c_1x + x^2 + 12c_1 + 4x}\right)^{\frac{1}{3}}}{2} + 1$$

$$y = \frac{\left(8 + 4x^3 + 12c_1 + 4x + 4\sqrt{x^6 + 6c_1x^3 + 2x^4 + 4x^3 + 9c_1^2 + 6c_1x + x^2 + 12c_1 + 4x}\right)^{\frac{1}{3}}}{4} - \frac{\left(8 + 4x^3 + 12c_1 + 4x + 4\sqrt{x^6 + 6c_1x^3 + 2x^4 + 4x^3 + 9c_1^2 + 6c_1x + x^2 + 12c_1 + 4x}\right)^{\frac{1}{3}}}{1} + 1$$

$$+ i\sqrt{3} \left(\frac{\left(8 + 4x^3 + 12c_1 + 4x + 4\sqrt{x^6 + 6c_1x^3 + 2x^4 + 4x^3 + 9c_1^2 + 6c_1x + x^2 + 12c_1 + 4x}\right)^{\frac{1}{3}}}{2} - \frac{2}{\left(8 + 4x^3 + 12c_1 + 4x + 4\sqrt{x^6 + 6c_1x^3 + 2x^4 + 4x^3 + 9c_1^2 + 6c_1x + x^2 + 12c_1 + 4x}\right)^{\frac{1}{3}}} \right)$$

$$y = \frac{\left(8 + 4x^3 + 12c_1 + 4x + 4\sqrt{x^6 + 6c_1x^3 + 2x^4 + 4x^3 + 9c_1^2 + 6c_1x + x^2 + 12c_1 + 4x}\right)^{\frac{1}{3}}}{\frac{4}{1}} - \frac{\left(8 + 4x^3 + 12c_1 + 4x + 4\sqrt{x^6 + 6c_1x^3 + 2x^4 + 4x^3 + 9c_1^2 + 6c_1x + x^2 + 12c_1 + 4x}\right)^{\frac{1}{3}}}{+1} + i\sqrt{3} \left(\frac{\left(8 + 4x^3 + 12c_1 + 4x + 4\sqrt{x^6 + 6c_1x^3 + 2x^4 + 4x^3 + 9c_1^2 + 6c_1x + x^2 + 12c_1 + 4x}\right)^{\frac{1}{3}}}{2} - \frac{2}{\left(8 + 4x^3 + 12c_1 + 4x + 4\sqrt{x^6 + 6c_1x^3 + 2x^4 + 4x^3 + 9c_1^2 + 6c_1x + x^2 + 12c_1 + 4x}\right)^{\frac{1}{3}}} \right)$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{-i\sqrt{3} \left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{2}{3}} + 4i\sqrt{3} - \left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{2}{3}} + 4\left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{1}{3}}}{4 \left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{i\sqrt{3} \left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{2}{3}} - 4i\sqrt{3} - \left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{2}{3}} + 4\left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{1}{3}}}{4 \left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{\left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{2}{3}} + 2\left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{1}{3}} + 4}{2 \left(8 + 12c_1 + 4\sqrt{3} \sqrt{3c_1^2 + 4c_1}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration.

Verification of solutions N/A

2.21.3 Solving as differential Type ode

Writing the ode as

$$y' = \frac{3x^2 + 1}{-6y + 3y^2} \quad (1)$$

Which becomes

$$(3y^2 - 6y) dy = (3x^2 + 1) dx \quad (2)$$

But the RHS is complete differential because

$$(3x^2 + 1) dx = d(x^3 + x)$$

Hence (2) becomes

$$(3y^2 - 6y) dy = d(x^3 + x)$$

Integrating both sides gives these solutions

$$y = \frac{\left(8 + 4x^3 + 4c_1 + 4x + 4\sqrt{x^6 + 2c_1x^3 + 2x^4 + 4x^3 + c_1^2 + 2c_1x + x^2 + 4c_1 + 4x}\right)^{\frac{1}{3}}}{2} + \frac{\dots}{(8 + 4x^3 + 4c_1)}$$

$$y = -\frac{\left(8 + 4x^3 + 4c_1 + 4x + 4\sqrt{x^6 + 2c_1x^3 + 2x^4 + 4x^3 + c_1^2 + 2c_1x + x^2 + 4c_1 + 4x}\right)^{\frac{1}{3}}}{4} - \frac{\dots}{(8 + 4x^3 + 4c_1)}$$

$$y = -\frac{\left(8 + 4x^3 + 4c_1 + 4x + 4\sqrt{x^6 + 2c_1x^3 + 2x^4 + 4x^3 + c_1^2 + 2c_1x + x^2 + 4c_1 + 4x}\right)^{\frac{1}{3}}}{4} - \frac{\dots}{(8 + 4x^3 + 4c_1)}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{-i\sqrt{3} \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{2}{3}} + 4i\sqrt{3} - \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{2}{3}} + 4c_1 \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{1}{3}}}{4 \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{i\sqrt{3} \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{2}{3}} - 4i\sqrt{3} - \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{2}{3}} + 4c_1 \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{1}{3}}}{4 \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{\left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{2}{3}} + 2c_1 \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{1}{3}} + 2 \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{1}{3}} + 4}{2 \left(8 + 4c_1 + 4\sqrt{c_1^2 + 4c_1}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration.

Verification of solutions N/A

2.21.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{3x^2 + 1}{3y(y - 2)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 151: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{x^2 + \frac{1}{3}} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{x^2 + \frac{1}{3}}} dx \end{aligned}$$

Which results in

$$S = \frac{1}{3}x^3 + \frac{1}{3}x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{3x^2 + 1}{3y(y - 2)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= x^2 + \frac{1}{3} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y(y - 2) \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R(R - 2)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{1}{3}R^3 - R^2 + c_1 \quad (4)$$

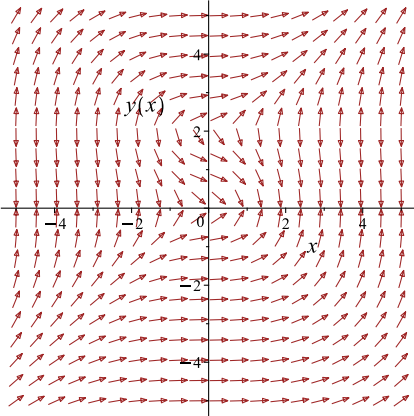
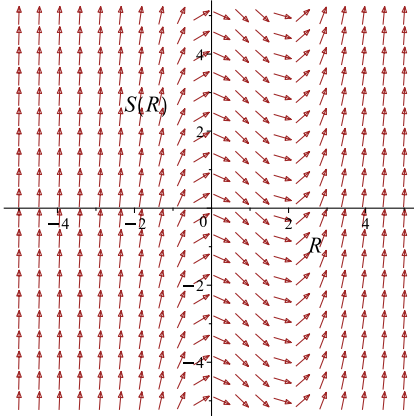
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{1}{3}x^3 + \frac{1}{3}x = \frac{y^3}{3} - y^2 + c_1$$

Which simplifies to

$$\frac{1}{3}x^3 + \frac{1}{3}x = \frac{y^3}{3} - y^2 + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{3x^2+1}{3y(y-2)}$ 	$R = y$ $S = \frac{1}{3}x^3 + \frac{1}{3}x$	$\frac{dS}{dR} = R(R - 2)$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{2}{3} + c_1$$

$$c_1 = \frac{2}{3}$$

Substituting c_1 found above in the general solution gives

$$\frac{1}{3}x^3 + \frac{1}{3}x = \frac{1}{3}y^3 - y^2 + \frac{2}{3}$$

Summary

The solution(s) found are the following

$$\frac{1}{3}x^3 + \frac{1}{3}x = \frac{y^3}{3} - y^2 + \frac{2}{3} \quad (1)$$

Verification of solutions

$$\frac{1}{3}x^3 + \frac{1}{3}x = \frac{y^3}{3} - y^2 + \frac{2}{3}$$

Verified OK.

2.21.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (3y(y - 2)) dy &= (3x^2 + 1) dx \\ (-3x^2 - 1) dx + (3y(y - 2)) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -3x^2 - 1 \\ N(x, y) &= 3y(y - 2) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (-3x^2 - 1) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (3y(y - 2)) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -3x^2 - 1 dx \\ \phi &= -x^3 - x + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y)\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 3y(y - 2)$. Therefore equation (4) becomes

$$3y(y - 2) = 0 + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 3y(y - 2)$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (3y(y - 2)) dy \\ f(y) &= y^3 - 3y^2 + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x^3 + y^3 - 3y^2 - x + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x^3 + y^3 - 3y^2 - x$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = c_1$$

$$c_1 = -2$$

Substituting c_1 found above in the general solution gives

$$-x^3 + y^3 - 3y^2 - x = -2$$

Summary

The solution(s) found are the following

$$-x^3 + y^3 - 3y^2 - x = -2 \quad (1)$$

Verification of solutions

$$-x^3 + y^3 - 3y^2 - x = -2$$

Verified OK.

2.21.6 Maple step by step solution

Let's solve

$$\left[y' - \frac{3x^2+1}{-6y+3y^2} = 0, y(0) = 1 \right]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$y'(-6y + 3y^2) = 3x^2 + 1$$

- Integrate both sides with respect to x

$$\int y'(-6y + 3y^2) dx = \int (3x^2 + 1) dx + c_1$$

- Evaluate integral

$$y^3 - 3y^2 = x^3 + c_1 + x$$

- Solve for y

$$y = \frac{\left(8+4x^3+4c_1+4x+4\sqrt{x^6+2c_1x^3+2x^4+4x^3+c_1^2+2c_1x+x^2+4c_1+4x}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(8+4x^3+4c_1+4x+4\sqrt{x^6+2c_1x^3+2x^4+4x^3+c_1^2+2c_1x+x^2+4c_1+4x}\right)^{\frac{1}{3}}}$$

- Use initial condition $y(0) = 1$

$$1 = \frac{\left(8+4c_1+4\sqrt{c_1^2+4c_1}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(8+4c_1+4\sqrt{c_1^2+4c_1}\right)^{\frac{1}{3}}} + 1$$

- Solve for c_1

$$c_1 = \text{RootOf}\left(\left(8+4_Z+4\sqrt{-Z^2+4_Z}\right)^{\frac{2}{3}}+4\right)$$

- Substitute $c_1 = \text{RootOf}\left(\left(8+4_Z+4\sqrt{-Z^2+4_Z}\right)^{\frac{2}{3}}+4\right)$ into general solution and simplify

$$y = \frac{\left(8+4x^3+4\text{RootOf}\left(\left(8+4_Z+4\sqrt{-Z(-Z+4)}\right)^{\frac{2}{3}}+4\right)+4x+4\sqrt{\left(x^3+x+\text{RootOf}\left(\left(8+4_Z+4\sqrt{-Z(-Z+4)}\right)^{\frac{2}{3}}+4\right)+4\right)}}{2\left(8+4x^3+4\text{RootOf}\left(\left(8+4_Z+4\sqrt{-Z(-Z+4)}\right)^{\frac{2}{3}}+4\right)+4x+4\sqrt{\left(x^3+x+\text{RootOf}\left(\left(8+4_Z+4\sqrt{-Z(-Z+4)}\right)^{\frac{2}{3}}+4\right)+4\right)}}\right)}$$

- Solution to the IVP

$$y = \frac{\left(8+4x^3+4\text{RootOf}\left(\left(8+4_Z+4\sqrt{-Z(-Z+4)}\right)^{\frac{2}{3}}+4\right)+4x+4\sqrt{\left(x^3+x+\text{RootOf}\left(\left(8+4_Z+4\sqrt{-Z(-Z+4)}\right)^{\frac{2}{3}}+4\right)+4\right)}}{2\left(8+4x^3+4\text{RootOf}\left(\left(8+4_Z+4\sqrt{-Z(-Z+4)}\right)^{\frac{2}{3}}+4\right)+4x+4\sqrt{\left(x^3+x+\text{RootOf}\left(\left(8+4_Z+4\sqrt{-Z(-Z+4)}\right)^{\frac{2}{3}}+4\right)+4\right)}}\right)}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.11 (sec). Leaf size: 109

```
dsolve([diff(y(x),x) = (3*x^2+1)/(-6*y(x)+3*y(x)^2),y(0) = 1],y(x), singsol=all)
```

$$y(x) = \frac{(1 + i\sqrt{3}) (4x^3 + 4x + 4\sqrt{x^6 + 2x^4 + x^2 - 4})^{\frac{2}{3}} - 4i\sqrt{3} - 4(4x^3 + 4x + 4\sqrt{x^6 + 2x^4 + x^2 - 4})^{\frac{1}{3}} + 4}{4 (4x^3 + 4x + 4\sqrt{x^6 + 2x^4 + x^2 - 4})^{\frac{1}{3}}}$$

✓ Solution by Mathematica

Time used: 4.019 (sec). Leaf size: 158

```
DSolve[{y'[x] == (3*x^2+1)/(-6*y[x]+3*y[x]^2),y[0]==1},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{-i2^{2/3}\sqrt{3}(x^3 + \sqrt{x^6 + 2x^4 + x^2 - 4} + x)^{2/3} - 2^{2/3}(x^3 + \sqrt{x^6 + 2x^4 + x^2 - 4} + x)^{2/3} + 4\sqrt[3]{x^3 + \sqrt{x^6 + 2x^4 + x^2 - 4}}}{4\sqrt[3]{x^3 + \sqrt{x^6 + 2x^4 + x^2 - 4} + x}}$$

2.22 problem 22

2.22.1 Solving as separable ode	700
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Internal problem ID [500]

Internal file name [OUTPUT/500_Sunday_June_05_2022_01_42_32_AM_94693732/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 22.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{3x^2}{-4 + 3y^2} = 0$$

With initial conditions

$$[y(1) = 0]$$

2.22.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{3x^2}{3y^2 - 4}\end{aligned}$$

Where $f(x) = 3x^2$ and $g(y) = \frac{1}{3y^2-4}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{1}{3y^2-4}} dy &= 3x^2 dx \\ \int \frac{1}{\frac{1}{3y^2-4}} dy &= \int 3x^2 dx \\ y^3 - 4y &= x^3 + c_1\end{aligned}$$

Which results in

$$\begin{aligned}y &= \frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{\frac{6}{8}} \\ &+ \frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}} \\ y &= -\frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{\frac{12}{4}} \\ &- \frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}} \\ &+ \frac{i\sqrt{3} \left(\frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{6} - \frac{8}{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}} \right)}{2} \\ y &= -\frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{\frac{12}{4}} \\ &- \frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}} \\ &+ \frac{i\sqrt{3} \left(\frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{6} - \frac{8}{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}} \right)}{2}\end{aligned}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{-i\sqrt{3} \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{2}{3}} + 48i\sqrt{3} - \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{2}{3}}}{12 \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}}}$$

$$c_1 = -1$$

Substituting c_1 found above in the general solution gives

$$y = \frac{-i\sqrt{3} (108x^3 - 108 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{2}{3}} - (108x^3 - 108 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{2}{3}} + 48i\sqrt{3}}{12 (108x^3 - 108 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{1}{3}}}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{i\sqrt{3} \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{2}{3}} - 48i\sqrt{3} - \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{2}{3}}}{12 \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}}}$$

Unable to solve for constant of integration. Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{48 + \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{2}{3}}}{6 \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}}}$$

Summary

The solution(s) found are the following

Warning: Unable to solve for constant of integration. y

$$= \frac{-i\sqrt{3} (108x^3 - 108 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{2}{3}} - (108x^3 - 108 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{2}{3}} + 48i\sqrt{3}}{12 (108x^3 - 108 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{1}{3}}}$$

Verification of solutions

$$y = \frac{-i\sqrt{3} (108x^3 - 108 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{2}{3}} - (108x^3 - 108 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{2}{3}} + 48i\sqrt{3}}{12 (108x^3 - 108 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{1}{3}}}$$

Verified OK.

2.22.2 Solving as differential Type ode

Writing the ode as

$$y' = \frac{3x^2}{-4 + 3y^2} \quad (1)$$

Which becomes

$$(3y^2 - 4) dy = (3x^2) dx \quad (2)$$

But the RHS is complete differential because

$$(3x^2) dx = d(x^3)$$

Hence (2) becomes

$$(3y^2 - 4) dy = d(x^3)$$

Integrating both sides gives these solutions

$$y = \frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{6} + \frac{8}{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}$$

$$y = -\frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{12} - \frac{4}{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}$$

$$y = -\frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{12} - \frac{4}{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{-i\sqrt{3} \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{2}{3}} + 48i\sqrt{3} - \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}}}{12 \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{i\sqrt{3} \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{2}{3}} - 48i\sqrt{3} - \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}}}{12 \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{\left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{2}{3}} + 6c_1 \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}} + 48}{6 \left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration.

Verification of solutions N/A

2.22.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{3x^2}{3y^2 - 4}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 154: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{3x^2} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{3x^2} dx \end{aligned}$$

Which results in

$$S = x^3$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{3x^2}{3y^2 - 4}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= 3x^2 \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 3y^2 - 4 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 3R^2 - 4$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R^3 - 4R + c_1 \quad (4)$$

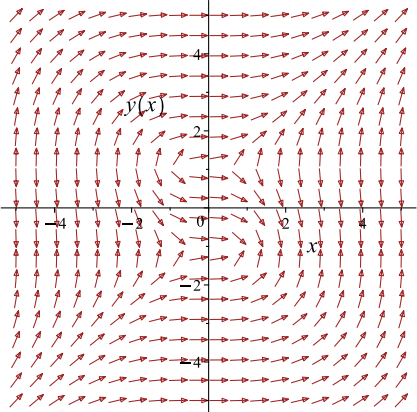
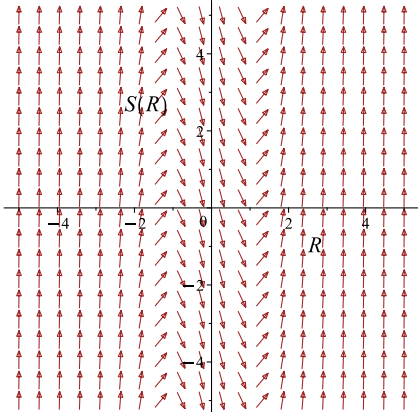
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$x^3 = y^3 + c_1 - 4y$$

Which simplifies to

$$x^3 = y^3 + c_1 - 4y$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{3x^2}{3y^2 - 4}$ 	$R = y$ $S = x^3$	$\frac{dS}{dR} = 3R^2 - 4$ 

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$1 = c_1$$

$$c_1 = 1$$

Substituting c_1 found above in the general solution gives

$$x^3 = y^3 - 4y + 1$$

Summary

The solution(s) found are the following

$$x^3 = y^3 - 4y + 1 \quad (1)$$

Verification of solutions

$$x^3 = y^3 - 4y + 1$$

Verified OK.

2.22.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work

and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left(y^2 - \frac{4}{3}\right) dy &= (x^2) dx \\ (-x^2) dx + \left(y^2 - \frac{4}{3}\right) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x^2 \\ N(x, y) &= y^2 - \frac{4}{3} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x^2) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}\left(y^2 - \frac{4}{3}\right) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x^2 dx \\ \phi &= -\frac{x^3}{3} + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = y^2 - \frac{4}{3}$. Therefore equation (4) becomes

$$y^2 - \frac{4}{3} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = y^2 - \frac{4}{3}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(y^2 - \frac{4}{3} \right) dy \\ f(y) &= \frac{1}{3}y^3 - \frac{4}{3}y + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{1}{3}x^3 + \frac{1}{3}y^3 - \frac{4}{3}y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{1}{3}x^3 + \frac{1}{3}y^3 - \frac{4}{3}y$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$-\frac{1}{3} = c_1$$

$$c_1 = -\frac{1}{3}$$

Substituting c_1 found above in the general solution gives

$$-\frac{1}{3}x^3 + \frac{1}{3}y^3 - \frac{4}{3}y = -\frac{1}{3}$$

Summary

The solution(s) found are the following

$$-\frac{x^3}{3} + \frac{y^3}{3} - \frac{4y}{3} = -\frac{1}{3} \quad (1)$$

Verification of solutions

$$-\frac{x^3}{3} + \frac{y^3}{3} - \frac{4y}{3} = -\frac{1}{3}$$

Verified OK.

2.22.5 Maple step by step solution

Let's solve

$$\left[y' - \frac{3x^2}{-4+3y^2} = 0, y(1) = 0 \right]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$y'(-4 + 3y^2) = 3x^2$$

- Integrate both sides with respect to x

$$\int y'(-4 + 3y^2) dx = \int 3x^2 dx + c_1$$

- Evaluate integral

$$y^3 - 4y = x^3 + c_1$$

- Solve for y

$$y = \frac{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}{6} + \frac{8}{\left(108x^3 + 108c_1 + 12\sqrt{81x^6 + 162c_1x^3 + 81c_1^2 - 768}\right)^{\frac{1}{3}}}$$

- Use initial condition $y(1) = 0$

$$0 = \frac{\left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}}}{6} + \frac{8}{\left(108 + 108c_1 + 12\sqrt{81c_1^2 + 162c_1 - 687}\right)^{\frac{1}{3}}}$$

- Solution does not satisfy initial condition

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.094 (sec). Leaf size: 73

```
dsolve([diff(y(x),x) = 3*x^2/(-4+3*y(x)^2),y(1) = 0],y(x), singsol=all)
```

$$y(x) = -\frac{(1 + i\sqrt{3}) (-108 + 108x^3 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{2}{3}} - 48i\sqrt{3} + 48}{12 (-108 + 108x^3 + 12\sqrt{81x^6 - 162x^3 - 687})^{\frac{1}{3}}}$$

✓ Solution by Mathematica

Time used: 9.526 (sec). Leaf size: 137

```
DSolve[{y'[x]== 3*x^2/(-4+3*y[x]^2),y[1]==0},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{-i\sqrt[3]{2}2^{2/3}(9x^3 + \sqrt{81x^6 - 162x^3 - 687} - 9)^{2/3} - \sqrt[3]{2}\sqrt[6]{3}(9x^3 + \sqrt{81x^6 - 162x^3 - 687} - 9)^{2/3} - 8\sqrt{3} + 48}{2 \cdot 2^{2/3}3^{5/6}\sqrt[3]{9x^3 + \sqrt{81x^6 - 162x^3 - 687} - 9}}$$

2.23 problem 23

2.23.1 Existence and uniqueness analysis	713
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2.23.4 Solving as exact ode	720
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Internal problem ID [501]

Internal file name [OUTPUT/501_Sunday_June_05_2022_01_42_33_AM_10464111/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 23.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - 2y^2 - xy^2 = 0$$

With initial conditions

$$[y(0) = 1]$$

2.23.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= xy^2 + 2y^2\end{aligned}$$

The x domain of $f(x, y)$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y}(xy^2 + 2y^2) \\ &= 2yx + 4y\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Therefore solution exists and is unique.

2.23.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= y^2(2 + x)\end{aligned}$$

Where $f(x) = 2 + x$ and $g(y) = y^2$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y^2} dy &= 2 + x dx \\ \int \frac{1}{y^2} dy &= \int 2 + x dx \\ -\frac{1}{y} &= 2x + \frac{1}{2}x^2 + c_1\end{aligned}$$

Which results in

$$y = -\frac{2}{x^2 + 2c_1 + 4x}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{1}{c_1}$$

$$c_1 = -1$$

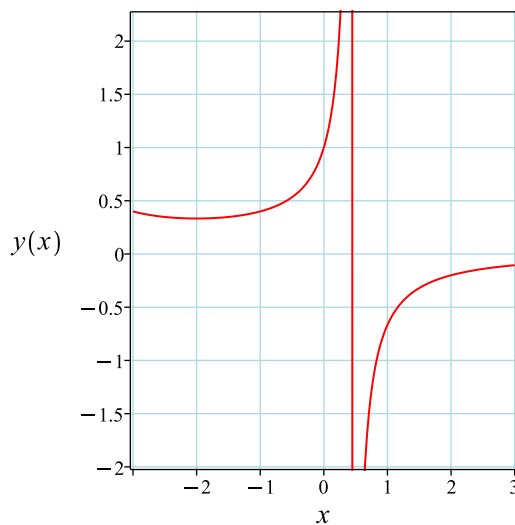
Substituting c_1 found above in the general solution gives

$$y = -\frac{2}{x^2 + 4x - 2}$$

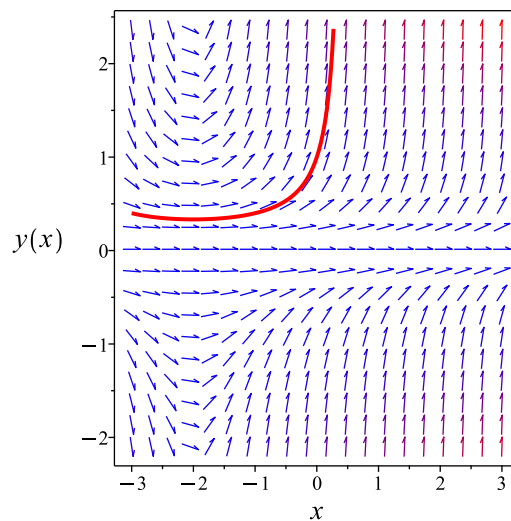
Summary

The solution(s) found are the following

$$y = -\frac{2}{x^2 + 4x - 2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{2}{x^2 + 4x - 2}$$

Verified OK.

2.23.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = x y^2 + 2y^2$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 157: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{2+x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{2+x}} dx\end{aligned}$$

Which results in

$$S = 2x + \frac{1}{2}x^2$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = x y^2 + 2y^2$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\ R_y &= 1 \\ S_x &= 2 + x \\ S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{y^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{1}{R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$2x + \frac{1}{2}x^2 = -\frac{1}{y} + c_1$$

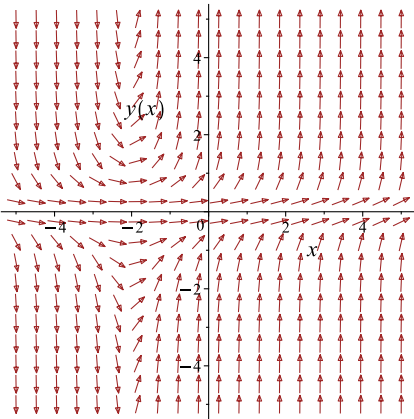
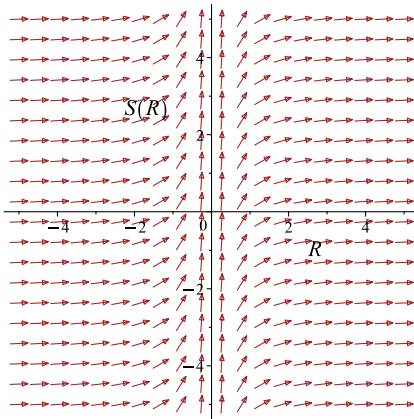
Which simplifies to

$$2x + \frac{1}{2}x^2 = -\frac{1}{y} + c_1$$

Which gives

$$y = \frac{2}{-x^2 + 2c_1 - 4x}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = x y^2 + 2y^2$ 	$R = y$ $S = 2x + \frac{1}{2}x^2$	$\frac{dS}{dR} = \frac{1}{R^2}$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{1}{c_1}$$

$$c_1 = 1$$

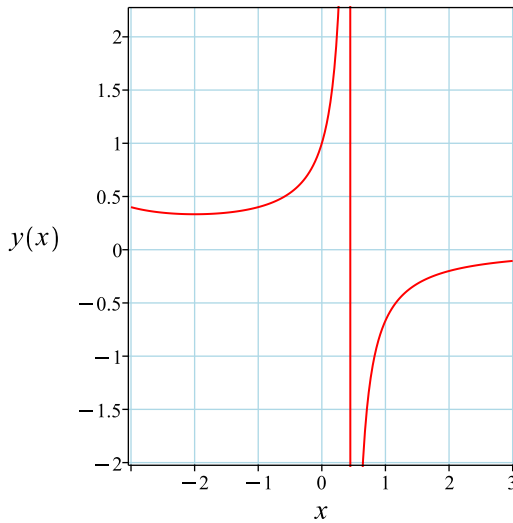
Substituting c_1 found above in the general solution gives

$$y = -\frac{2}{x^2 + 4x - 2}$$

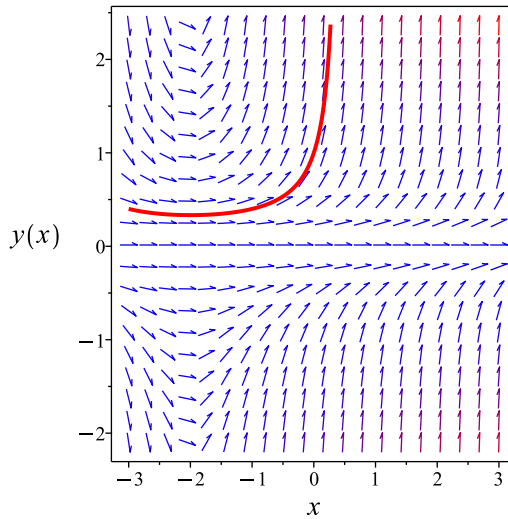
Summary

The solution(s) found are the following

$$y = -\frac{2}{x^2 + 4x - 2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{2}{x^2 + 4x - 2}$$

Verified OK.

2.23.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} \left(\frac{1}{y^2}\right) dy &= (2 + x) dx \\ (-x - 2) dx + \left(\frac{1}{y^2}\right) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x - 2 \\ N(x, y) &= \frac{1}{y^2} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x - 2) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{1}{y^2}\right) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -x - 2 dx$$

$$\phi = -\frac{1}{2}x^2 - 2x + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{y^2}$. Therefore equation (4) becomes

$$\frac{1}{y^2} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{1}{y^2}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(\frac{1}{y^2} \right) dy$$

$$f(y) = -\frac{1}{y} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{x^2}{2} - 2x - \frac{1}{y} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{x^2}{2} - 2x - \frac{1}{y}$$

The solution becomes

$$y = -\frac{2}{x^2 + 2c_1 + 4x}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{1}{c_1}$$

$$c_1 = -1$$

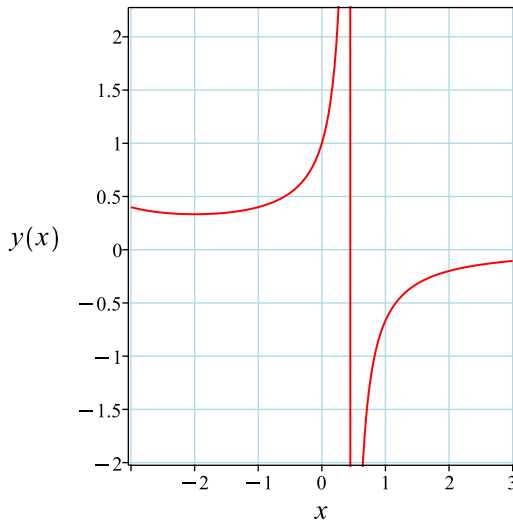
Substituting c_1 found above in the general solution gives

$$y = -\frac{2}{x^2 + 4x - 2}$$

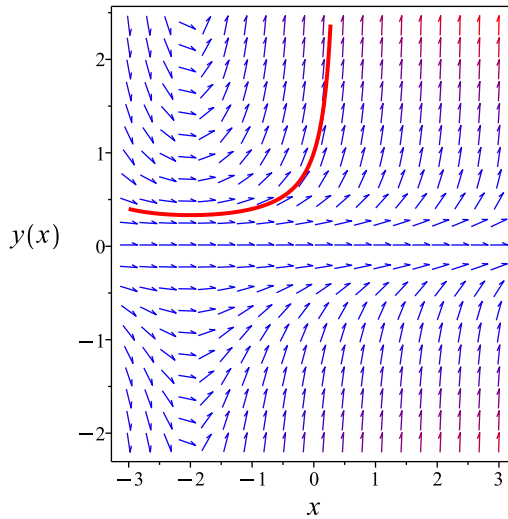
Summary

The solution(s) found are the following

$$y = -\frac{2}{x^2 + 4x - 2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{2}{x^2 + 4x - 2}$$

Verified OK.

2.23.5 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= xy^2 + 2y^2 \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = xy^2 + 2y^2$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that $f_0(x) = 0$, $f_1(x) = 0$ and $f_2(x) = 2 + x$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{(2+x)u} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \quad (2)$$

But

$$f_2' = 1$$

$$f_1 f_2 = 0$$

$$f_2^2 f_0 = 0$$

Substituting the above terms back in equation (2) gives

$$(2 + x) u''(x) - u'(x) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 + c_2(2 + x)^2$$

The above shows that

$$u'(x) = 2(2 + x) c_2$$

Using the above in (1) gives the solution

$$y = -\frac{2c_2}{c_1 + c_2(2 + x)^2}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = -\frac{2}{x^2 + c_3 + 4x + 4}$$

Initial conditions are used to solve for c_3 . Substituting $x = 0$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{2}{c_3 + 4}$$

$$c_3 = -6$$

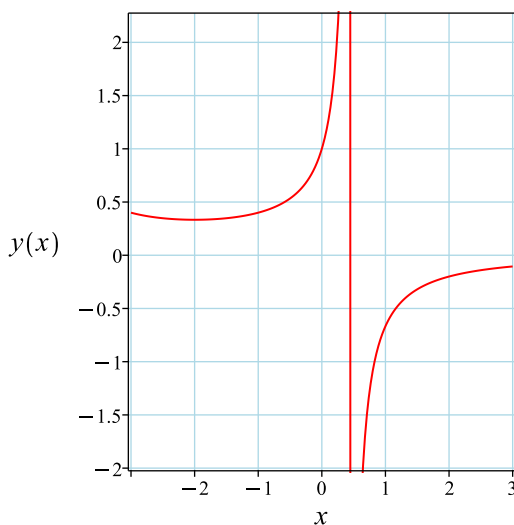
Substituting c_3 found above in the general solution gives

$$y = -\frac{2}{x^2 + 4x - 2}$$

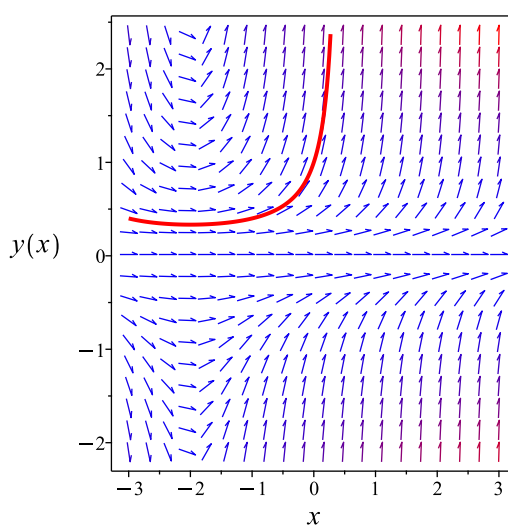
Summary

The solution(s) found are the following

$$y = -\frac{2}{x^2 + 4x - 2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{2}{x^2 + 4x - 2}$$

Verified OK.

2.23.6 Maple step by step solution

Let's solve

$$[y' - 2y^2 - xy^2 = 0, y(0) = 1]$$

- Highest derivative means the order of the ODE is 1
- y'
- Separate variables

$$\frac{y'}{y^2} = 2 + x$$

- Integrate both sides with respect to x

$$\int \frac{y'}{y^2} dx = \int (2 + x) dx + c_1$$

- Evaluate integral

$$-\frac{1}{y} = 2x + \frac{1}{2}x^2 + c_1$$

- Solve for y

$$y = -\frac{2}{x^2 + 2c_1 + 4x}$$

- Use initial condition $y(0) = 1$

$$1 = -\frac{1}{c_1}$$

- Solve for c_1

$$c_1 = -1$$

- Substitute $c_1 = -1$ into general solution and simplify

$$y = -\frac{2}{x^2 + 4x - 2}$$

- Solution to the IVP

$$y = -\frac{2}{x^2 + 4x - 2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.063 (sec). Leaf size: 16

```
dsolve([diff(y(x),x) = 2*y(x)^2+x*y(x)^2,y(0) = 1],y(x), singsol=all)
```

$$y(x) = -\frac{2}{x^2 + 4x - 2}$$

✓ Solution by Mathematica

Time used: 0.127 (sec). Leaf size: 17

```
DSolve[{y'[x] == 2*y[x]^2+x*y[x]^2,y[0]==1},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{2}{x^2 + 4x - 2}$$

2.24 problem 24

2.24.1 Existence and uniqueness analysis	729
2.24.2 Solving as separable ode	730
2.24.3 Solving as first order ode lie symmetry lookup ode	732
2.24.4 Solving as exact ode	736
2.24.5 Maple step by step solution	739

Internal problem ID [502]

Internal file name [OUTPUT/502_Sunday_June_05_2022_01_42_34_AM_42187153/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 24.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{2 - e^x}{3 + 2y} = 0$$

With initial conditions

$$[y(0) = 0]$$

2.24.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned} y' &= f(x, y) \\ &= -\frac{-2 + e^x}{3 + 2y} \end{aligned}$$

The x domain of $f(x, y)$ when $y = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\left\{ y < -\frac{3}{2} \vee -\frac{3}{2} < y \right\}$$

And the point $y_0 = 0$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned} \frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{-2 + e^x}{3 + 2y} \right) \\ &= \frac{-4 + 2e^x}{(3 + 2y)^2} \end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\left\{ y < -\frac{3}{2} \vee -\frac{3}{2} < y \right\}$$

And the point $y_0 = 0$ is inside this domain. Therefore solution exists and is unique.

2.24.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{2 - e^x}{3 + 2y} \end{aligned}$$

Where $f(x) = 2 - e^x$ and $g(y) = \frac{1}{3+2y}$. Integrating both sides gives

$$\begin{aligned} \frac{1}{3+2y} dy &= 2 - e^x dx \\ \int \frac{1}{3+2y} dy &= \int 2 - e^x dx \\ y^2 + 3y &= 2x - e^x + c_1 \end{aligned}$$

Which results in

$$y = -\frac{3}{2} + \frac{\sqrt{9 - 4e^x + 4c_1 + 8x}}{2}$$

$$y = -\frac{3}{2} - \frac{\sqrt{9 - 4e^x + 4c_1 + 8x}}{2}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{3}{2} - \frac{\sqrt{5 + 4c_1}}{2}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -\frac{3}{2} + \frac{\sqrt{5 + 4c_1}}{2}$$

$$c_1 = 1$$

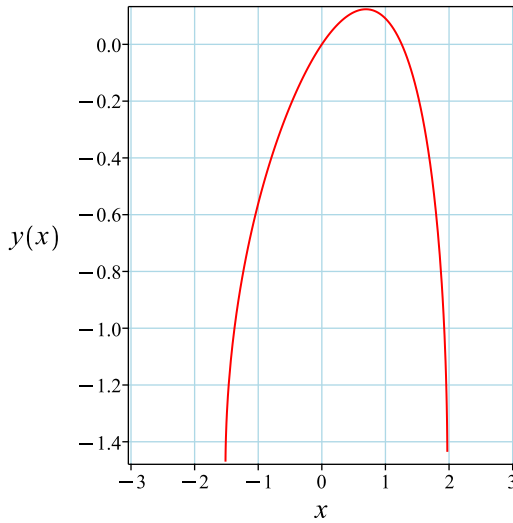
Substituting c_1 found above in the general solution gives

$$y = -\frac{3}{2} + \frac{\sqrt{13 - 4e^x + 8x}}{2}$$

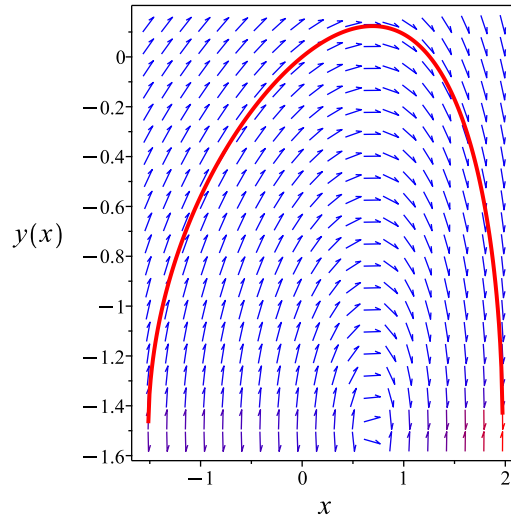
Summary

The solution(s) found are the following

$$y = -\frac{3}{2} + \frac{\sqrt{13 - 4e^x + 8x}}{2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{3}{2} + \frac{\sqrt{13 - 4e^x + 8x}}{2}$$

Verified OK.

2.24.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{-2 + e^x}{3 + 2y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 160: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{2 - e^x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{2-e^x}} dx \end{aligned}$$

Which results in

$$S = 2x - e^x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{-2 + e^x}{3 + 2y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= 2 - e^x \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 3 + 2y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 3 + 2R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R^2 + 3R + c_1 \quad (4)$$

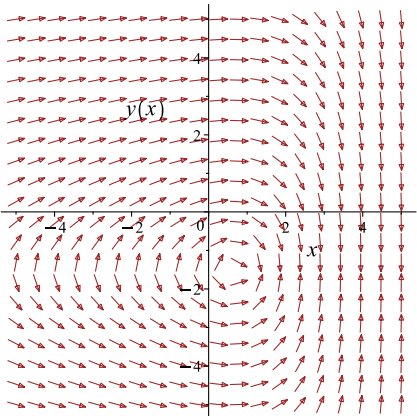
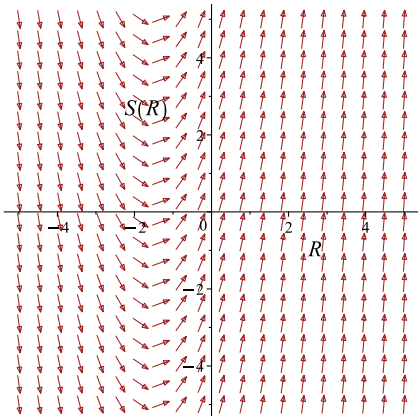
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$2x - e^x = y^2 + c_1 + 3y$$

Which simplifies to

$$2x - e^x = y^2 + c_1 + 3y$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{-2+e^x}{3+2y}$ 	$R = y$ $S = 2x - e^x$	$\frac{dS}{dR} = 3 + 2R$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$-1 = c_1$$

$$c_1 = -1$$

Substituting c_1 found above in the general solution gives

$$2x - e^x = y^2 + 3y - 1$$

Summary

The solution(s) found are the following

$$2x - e^x = y^2 + 3y - 1 \quad (1)$$

Verification of solutions

$$2x - e^x = y^2 + 3y - 1$$

Verified OK.

2.24.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work

and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (-3 - 2y) dy &= (-2 + e^x) dx \\ (2 - e^x) dx + (-3 - 2y) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 2 - e^x \\ N(x, y) &= -3 - 2y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2 - e^x) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-3 - 2y) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 2 - e^x dx \\ \phi &= 2x - e^x + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial\phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial\phi}{\partial y} = -3 - 2y$. Therefore equation (4) becomes

$$-3 - 2y = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -3 - 2y$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) \, dy &= \int (-3 - 2y) \, dy \\ f(y) &= -y^2 - 3y + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -y^2 - e^x + 2x - 3y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -y^2 - e^x + 2x - 3y$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$-1 = c_1$$

$$c_1 = -1$$

Substituting c_1 found above in the general solution gives

$$-y^2 - e^x + 2x - 3y = -1$$

Summary

The solution(s) found are the following

$$-y^2 - e^x + 2x - 3y = -1 \quad (1)$$

Verification of solutions

$$-y^2 - e^x + 2x - 3y = -1$$

Verified OK.

2.24.5 Maple step by step solution

Let's solve

$$\left[y' - \frac{2-e^x}{3+2y} = 0, y(0) = 0 \right]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$(3 + 2y) y' = 2 - e^x$$

- Integrate both sides with respect to x

$$\int (3 + 2y) y' dx = \int (2 - e^x) dx + c_1$$

- Evaluate integral

$$y^2 + 3y = 2x - e^x + c_1$$

- Solve for y

$$\left\{ y = -\frac{3}{2} - \frac{\sqrt{9-4e^x+4c_1+8x}}{2}, y = -\frac{3}{2} + \frac{\sqrt{9-4e^x+4c_1+8x}}{2} \right\}$$

- Use initial condition $y(0) = 0$

$$0 = -\frac{3}{2} - \frac{\sqrt{5+4c_1}}{2}$$

- Solution does not satisfy initial condition

- Use initial condition $y(0) = 0$

$$0 = -\frac{3}{2} + \frac{\sqrt{5+4c_1}}{2}$$

- Solve for c_1

$$c_1 = 1$$

- Substitute $c_1 = 1$ into general solution and simplify

$$y = -\frac{3}{2} + \frac{\sqrt{13-4e^x+8x}}{2}$$

- Solution to the IVP

$$y = -\frac{3}{2} + \frac{\sqrt{13-4e^x+8x}}{2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.141 (sec). Leaf size: 19

```
dsolve([diff(y(x),x) = (2-exp(x))/(3+2*y(x)),y(0) = 0],y(x), singsol=all)
```

$$y(x) = -\frac{3}{2} + \frac{\sqrt{13-4e^x+8x}}{2}$$

✓ Solution by Mathematica

Time used: 0.737 (sec). Leaf size: 25

```
DSolve[{y'[x] == (2-Exp[x])/(3+2*y[x]),y[0]==0},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2}(\sqrt{8x-4e^x+13}-3)$$

2.25 problem 25

2.25.1 Existence and uniqueness analysis	741
2.25.2 Solving as separable ode	742
2.25.3 Solving as first order ode lie symmetry lookup ode	744
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2.25.5 Maple step by step solution	751

Internal problem ID [503]

Internal file name [OUTPUT/503_Sunday_June_05_2022_01_42_35_AM_21918699/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 25.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{2 \cos(2x)}{3 + 2y} = 0$$

With initial conditions

$$[y(0) = -1]$$

2.25.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned} y' &= f(x, y) \\ &= \frac{2 \cos(2x)}{3 + 2y} \end{aligned}$$

The x domain of $f(x, y)$ when $y = -1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\left\{ y < -\frac{3}{2} \vee -\frac{3}{2} < y \right\}$$

And the point $y_0 = -1$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned} \frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{2 \cos(2x)}{3 + 2y} \right) \\ &= -\frac{4 \cos(2x)}{(3 + 2y)^2} \end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = -1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\left\{ y < -\frac{3}{2} \vee -\frac{3}{2} < y \right\}$$

And the point $y_0 = -1$ is inside this domain. Therefore solution exists and is unique.

2.25.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{2 \cos(2x)}{3 + 2y} \end{aligned}$$

Where $f(x) = 2 \cos(2x)$ and $g(y) = \frac{1}{3+2y}$. Integrating both sides gives

$$\begin{aligned} \frac{1}{3+2y} dy &= 2 \cos(2x) dx \\ \int \frac{1}{3+2y} dy &= \int 2 \cos(2x) dx \\ y^2 + 3y &= \sin(2x) + c_1 \end{aligned}$$

Which results in

$$y = -\frac{3}{2} + \frac{\sqrt{9 + 4 \sin(2x) + 4c_1}}{2}$$

$$y = -\frac{3}{2} - \frac{\sqrt{9 + 4 \sin(2x) + 4c_1}}{2}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -1$ in the above solution gives an equation to solve for the constant of integration.

$$-1 = -\frac{3}{2} - \frac{\sqrt{9 + 4c_1}}{2}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -1$ in the above solution gives an equation to solve for the constant of integration.

$$-1 = -\frac{3}{2} + \frac{\sqrt{9 + 4c_1}}{2}$$

$$c_1 = -2$$

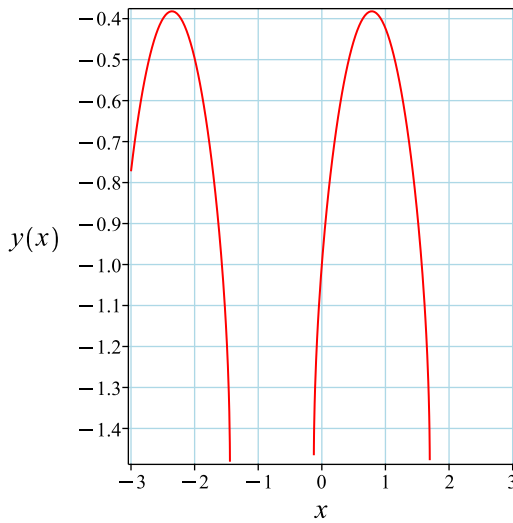
Substituting c_1 found above in the general solution gives

$$y = -\frac{3}{2} + \frac{\sqrt{1 + 4 \sin(2x)}}{2}$$

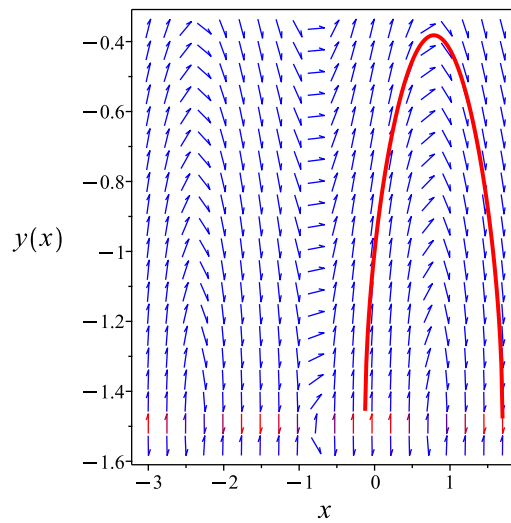
Summary

The solution(s) found are the following

$$y = -\frac{3}{2} + \frac{\sqrt{1 + 4 \sin(2x)}}{2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{3}{2} + \frac{\sqrt{1 + 4 \sin(2x)}}{2}$$

Verified OK.

2.25.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{2 \cos(2x)}{3 + 2y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 163: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\xi(x, y) = \frac{1}{2 \cos(2x)}$$

$$\eta(x, y) = 0 \tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{2 \cos(2x)}} dx \end{aligned}$$

Which results in

$$S = \sin(2x)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{2 \cos(2x)}{3 + 2y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= 2 \cos(2x) \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 3 + 2y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 3 + 2R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R^2 + 3R + c_1 \quad (4)$$

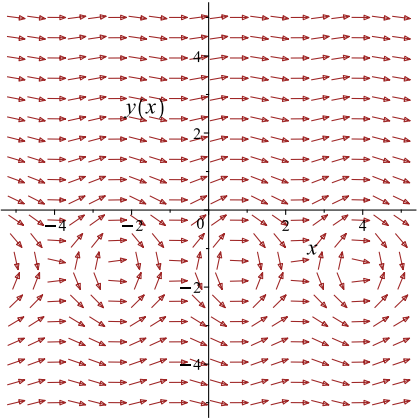
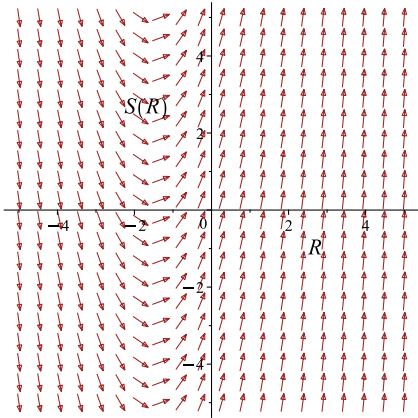
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\sin(2x) = y^2 + c_1 + 3y$$

Which simplifies to

$$\sin(2x) = y^2 + c_1 + 3y$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{2 \cos(2x)}{3+2y}$ 	$R = y$ $S = \sin(2x)$	$\frac{dS}{dR} = 3 + 2R$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -1$ in the above solution gives an equation to solve for the constant of integration.

$$0 = c_1 - 2$$

$$c_1 = 2$$

Substituting c_1 found above in the general solution gives

$$\sin(2x) = y^2 + 3y + 2$$

Summary

The solution(s) found are the following

$$\sin(2x) = y^2 + 3y + 2 \quad (1)$$

Verification of solutions

$$\sin(2x) = y^2 + 3y + 2$$

Verified OK.

2.25.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work

and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left(\frac{3}{2} + y\right) dy &= (\cos(2x)) dx \\ (-\cos(2x)) dx + \left(\frac{3}{2} + y\right) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -\cos(2x) \\ N(x, y) &= \frac{3}{2} + y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-\cos(2x)) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}\left(\frac{3}{2} + y\right) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\cos(2x) dx \\ \phi &= -\frac{\sin(2x)}{2} + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{3}{2} + y$. Therefore equation (4) becomes

$$\frac{3}{2} + y = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{3}{2} + y$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(\frac{3}{2} + y\right) dy \\ f(y) &= \frac{3}{2}y + \frac{1}{2}y^2 + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{\sin(2x)}{2} + \frac{3y}{2} + \frac{y^2}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{\sin(2x)}{2} + \frac{3y}{2} + \frac{y^2}{2}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = -1$ in the above solution gives an equation to solve for the constant of integration.

$$-1 = c_1$$

$$c_1 = -1$$

Substituting c_1 found above in the general solution gives

$$-\frac{\sin(2x)}{2} + \frac{3y}{2} + \frac{y^2}{2} = -1$$

Summary

The solution(s) found are the following

$$-\frac{\sin(2x)}{2} + \frac{3y}{2} + \frac{y^2}{2} = -1 \quad (1)$$

Verification of solutions

$$-\frac{\sin(2x)}{2} + \frac{3y}{2} + \frac{y^2}{2} = -1$$

Verified OK.

2.25.5 Maple step by step solution

Let's solve

$$\left[y' - \frac{2 \cos(2x)}{3+2y} = 0, y(0) = -1 \right]$$

- Highest derivative means the order of the ODE is 1
- y'
- Separate variables
- $(3 + 2y) y' = 2 \cos(2x)$
- Integrate both sides with respect to x
- $\int (3 + 2y) y' dx = \int 2 \cos(2x) dx + c_1$
- Evaluate integral
- $y^2 + 3y = \sin(2x) + c_1$
- Solve for y

$$\left\{ y = -\frac{3}{2} - \frac{\sqrt{9+4\sin(2x)+4c_1}}{2}, y = -\frac{3}{2} + \frac{\sqrt{9+4\sin(2x)+4c_1}}{2} \right\}$$

- Use initial condition $y(0) = -1$

$$-1 = -\frac{3}{2} - \frac{\sqrt{9+4c_1}}{2}$$
- Solution does not satisfy initial condition
- Use initial condition $y(0) = -1$

$$-1 = -\frac{3}{2} + \frac{\sqrt{9+4c_1}}{2}$$
- Solve for c_1

$$c_1 = -2$$
- Substitute $c_1 = -2$ into general solution and simplify

$$y = -\frac{3}{2} + \frac{\sqrt{1+4\sin(2x)}}{2}$$
- Solution to the IVP

$$y = -\frac{3}{2} + \frac{\sqrt{1+4\sin(2x)}}{2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.25 (sec). Leaf size: 18

```
dsolve([diff(y(x),x) = 2*cos(2*x)/(3+2*y(x)),y(0) = -1],y(x), singsol=all)
```

$$y(x) = -\frac{3}{2} + \frac{\sqrt{1+4\sin(2x)}}{2}$$

✓ Solution by Mathematica

Time used: 0.175 (sec). Leaf size: 23

```
DSolve[{y'[x] == 2*Cos[2*x]/(3+2*y[x]), y[0]==-1}, y[x], x, IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2} \left(\sqrt{4 \sin(2x) + 1} - 3 \right)$$

2.26 problem 26

2.26.1 Existence and uniqueness analysis	754
2.26.2 Solving as separable ode	755
2.26.3 Solving as first order ode lie symmetry lookup ode	757
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2.26.5 Solving as riccati ode	765
2.26.6 Maple step by step solution	767

Internal problem ID [504]

Internal file name [OUTPUT/504_Sunday_June_05_2022_01_42_37_AM_55464902/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 26.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - 2(x + 1)(1 + y^2) = 0$$

With initial conditions

$$[y(0) = 0]$$

2.26.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= 2(x + 1)(y^2 + 1)\end{aligned}$$

The x domain of $f(x, y)$ when $y = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 0$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y}(2(x+1)(y^2+1)) \\ &= 4(x+1)y\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 0$ is inside this domain. Therefore solution exists and is unique.

2.26.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= (x+1)(2y^2+2)\end{aligned}$$

Where $f(x) = x+1$ and $g(y) = 2y^2+2$. Integrating both sides gives

$$\begin{aligned}\frac{1}{2y^2+2} dy &= x+1 dx \\ \int \frac{1}{2y^2+2} dy &= \int x+1 dx \\ \frac{\arctan(y)}{2} &= \frac{1}{2}x^2 + x + c_1\end{aligned}$$

Which results in

$$y = \tan(x^2 + 2c_1 + 2x)$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \tan(2c_1)$$

$$c_1 = 0$$

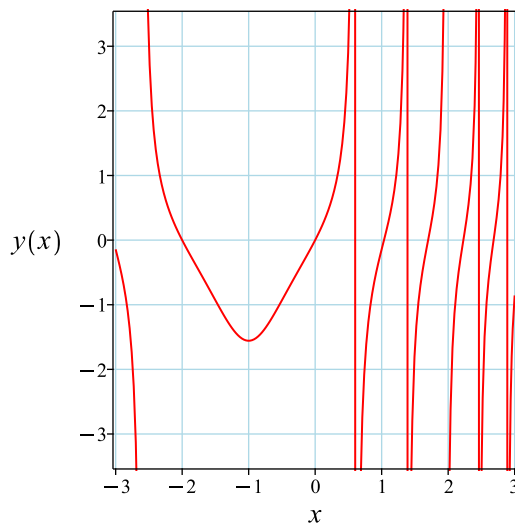
Substituting c_1 found above in the general solution gives

$$y = \tan(x^2 + 2x)$$

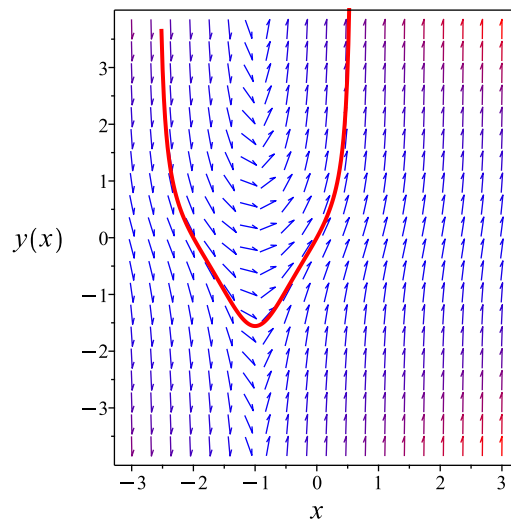
Summary

The solution(s) found are the following

$$y = \tan(x^2 + 2x) \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \tan(x^2 + 2x)$$

Verified OK.

2.26.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = 2(x+1)(y^2+1)$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 166: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int(n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{x+1} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{x+1}} dx\end{aligned}$$

Which results in

$$S = \frac{1}{2}x^2 + x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = 2(x+1)(y^2+1)$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\ R_y &= 1 \\ S_x &= x+1 \\ S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{2y^2 + 2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{2R^2 + 2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{\arctan(R)}{2} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{1}{2}x^2 + x = \frac{\arctan(y)}{2} + c_1$$

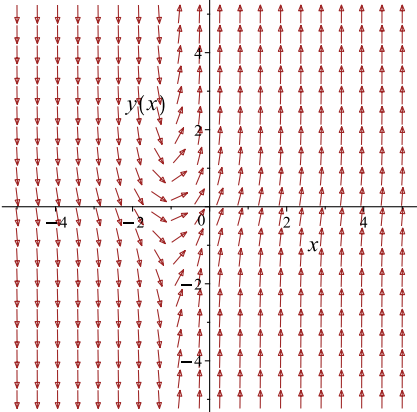
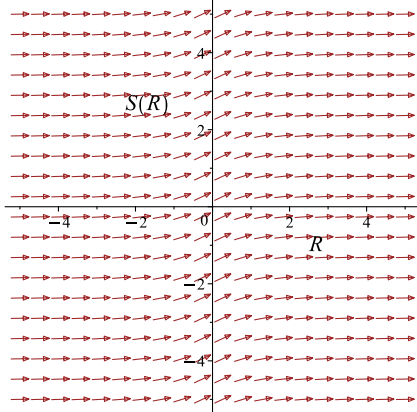
Which simplifies to

$$\frac{1}{2}x^2 + x = \frac{\arctan(y)}{2} + c_1$$

Which gives

$$y = -\tan(-x^2 + 2c_1 - 2x)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = 2(x+1)(y^2+1)$ 	$R = y$ $S = \frac{1}{2}x^2 + x$	$\frac{dS}{dR} = \frac{1}{2R^2+2}$ 

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -\tan(2c_1)$$

$$c_1 = 0$$

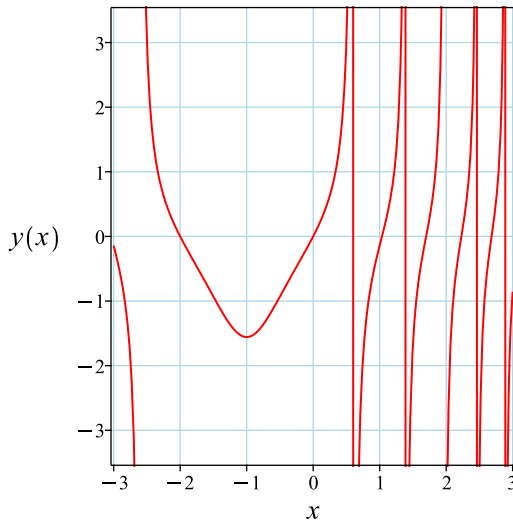
Substituting c_1 found above in the general solution gives

$$y = \tan(x^2 + 2x)$$

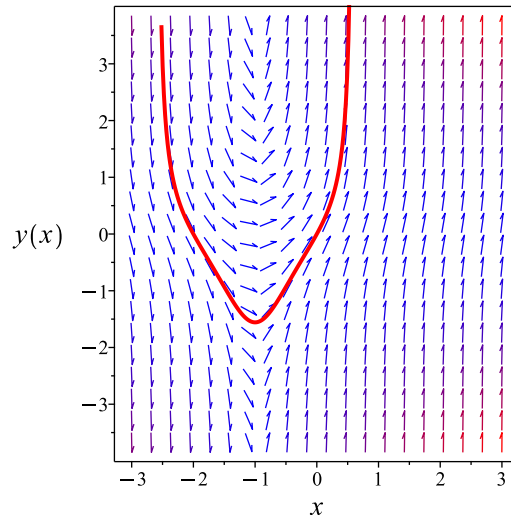
Summary

The solution(s) found are the following

$$y = \tan(x^2 + 2x) \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \tan(x^2 + 2x)$$

Verified OK.

2.26.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left(\frac{1}{2y^2 + 2}\right) dy &= (x + 1) dx \\ (-x - 1) dx + \left(\frac{1}{2y^2 + 2}\right) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x - 1 \\ N(x, y) &= \frac{1}{2y^2 + 2} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x - 1) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{1}{2y^2 + 2}\right) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x - 1 dx \\ \phi &= -\frac{1}{2}x^2 - x + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{2y^2+2}$. Therefore equation (4) becomes

$$\frac{1}{2y^2+2} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{1}{2y^2+2}$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int \left(\frac{1}{2y^2+2} \right) dy \\ f(y) &= \frac{\arctan(y)}{2} + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{x^2}{2} - x + \frac{\arctan(y)}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{x^2}{2} - x + \frac{\arctan(y)}{2}$$

The solution becomes

$$y = \tan(x^2 + 2c_1 + 2x)$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \tan(2c_1)$$

$$c_1 = 0$$

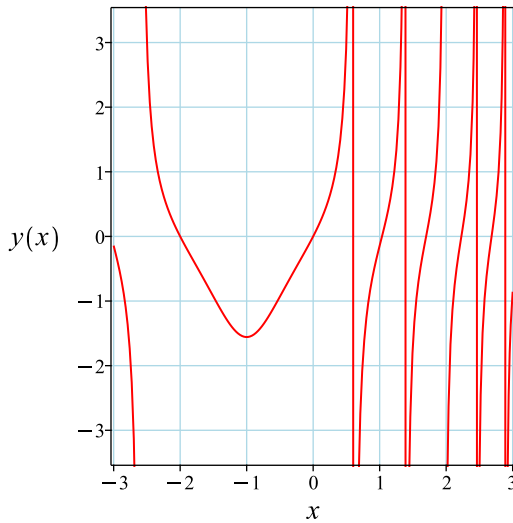
Substituting c_1 found above in the general solution gives

$$y = \tan(x^2 + 2x)$$

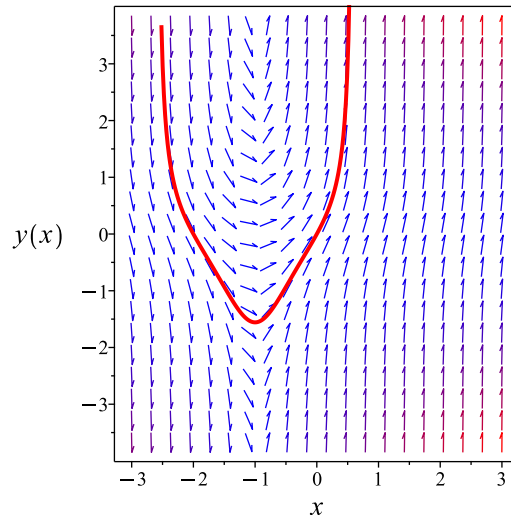
Summary

The solution(s) found are the following

$$y = \tan(x^2 + 2x) \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \tan(x^2 + 2x)$$

Verified OK.

2.26.5 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= 2(x+1)(y^2+1) \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = 2x y^2 + 2y^2 + 2x + 2$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that $f_0(x) = 2 + 2x$, $f_1(x) = 0$ and $f_2(x) = 2 + 2x$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{(2+2x)u} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \quad (2)$$

But

$$\begin{aligned} f_2' &= 2 \\ f_1 f_2 &= 0 \\ f_2^2 f_0 &= (2 + 2x)^3 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$(2 + 2x) u''(x) - 2u'(x) + (2 + 2x)^3 u(x) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 \sin(x^2 + 2x) + c_2 \cos(x^2 + 2x)$$

The above shows that

$$u'(x) = -2(x + 1) (c_2 \sin(x^2 + 2x) - c_1 \cos(x^2 + 2x))$$

Using the above in (1) gives the solution

$$y = \frac{2(x + 1) (c_2 \sin(x^2 + 2x) - c_1 \cos(x^2 + 2x))}{(2 + 2x) (c_1 \sin(x^2 + 2x) + c_2 \cos(x^2 + 2x))}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = \frac{\sin(x^2 + 2x) - c_3 \cos(x^2 + 2x)}{c_3 \sin(x^2 + 2x) + \cos(x^2 + 2x)}$$

Initial conditions are used to solve for c_3 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -c_3$$

$$c_3 = 0$$

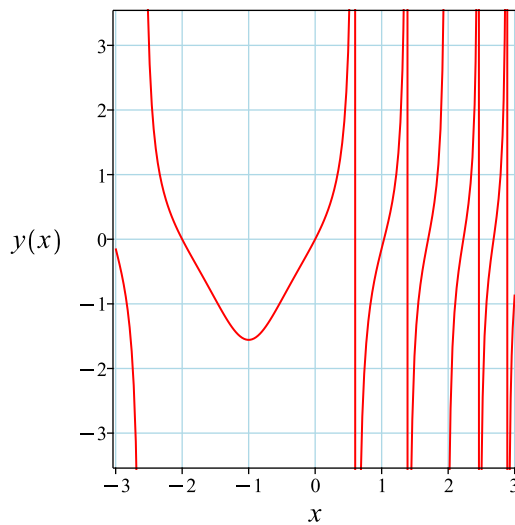
Substituting c_3 found above in the general solution gives

$$y = \frac{\sin(x(2+x))}{\cos(x(2+x))}$$

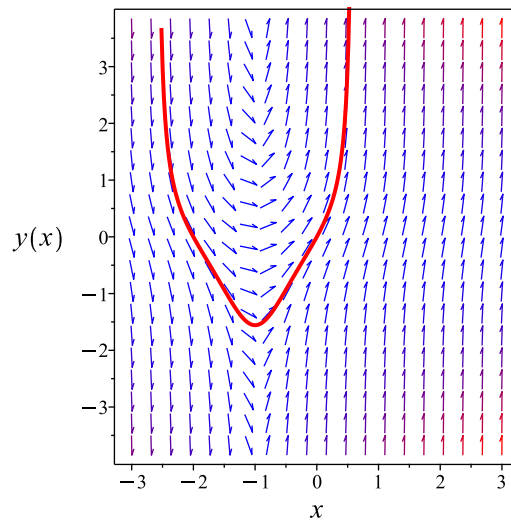
Summary

The solution(s) found are the following

$$y = \frac{\sin(x(2+x))}{\cos(x(2+x))} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{\sin(x(2+x))}{\cos(x(2+x))}$$

Verified OK.

2.26.6 Maple step by step solution

Let's solve

$$[y' - 2(x+1)(1+y^2) = 0, y(0) = 0]$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'}{1+y^2} = 2 + 2x$$
- Integrate both sides with respect to x

$$\int \frac{y'}{1+y^2} dx = \int (2 + 2x) dx + c_1$$
- Evaluate integral

$$\arctan(y) = x^2 + c_1 + 2x$$
- Solve for y

$$y = \tan(x^2 + c_1 + 2x)$$
- Use initial condition $y(0) = 0$

$$0 = \tan(c_1)$$
- Solve for c_1

$$c_1 = 0$$
- Substitute $c_1 = 0$ into general solution and simplify

$$y = \tan(x^2 + 2x)$$
- Solution to the IVP

$$y = \tan(x^2 + 2x)$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 12

```
dsolve([diff(y(x),x) = 2*(1+x)*(1+y(x)^2),y(0) = 0],y(x), singsol=all)
```

$$y(x) = \tan(x^2 + 2x)$$

✓ Solution by Mathematica

Time used: 0.199 (sec). Leaf size: 11

```
DSolve[{y'[x] == 2*(1+x)*(1+y[x]^2), y[0]==0}, y[x], x, IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \tan(x(x + 2))$$

2.27 problem 27

2.27.1 Solving as separable ode	770
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Internal problem ID [505]

Internal file name [OUTPUT/505_Sunday_June_05_2022_01_42_38_AM_39711965/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 27.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "bernoulli", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{t(4-y)y}{3} = 0$$

2.27.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(t, y) \\ &= f(t)g(y) \\ &= -\frac{ty(y-4)}{3}\end{aligned}$$

Where $f(t) = -\frac{t}{3}$ and $g(y) = y(y - 4)$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y(y-4)} dy &= -\frac{t}{3} dt \\ \int \frac{1}{y(y-4)} dy &= \int -\frac{t}{3} dt \\ \frac{\ln(y-4)}{4} - \frac{\ln(y)}{4} &= -\frac{t^2}{6} + c_1\end{aligned}$$

The above can be written as

$$\begin{aligned}\left(\frac{1}{4}\right) (\ln(y-4) - \ln(y)) &= -\frac{t^2}{6} + 2c_1 \\ \ln(y-4) - \ln(y) &= (4) \left(-\frac{t^2}{6} + 2c_1\right) \\ &= -\frac{2t^2}{3} + 8c_1\end{aligned}$$

Raising both side to exponential gives

$$e^{\ln(y-4) - \ln(y)} = e^{-\frac{2t^2}{3} + 8c_1}$$

Which simplifies to

$$\begin{aligned}\frac{y-4}{y} &= 4c_1 e^{-\frac{2t^2}{3}} \\ &= c_2 e^{-\frac{2t^2}{3}}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = -\frac{4}{-1 + c_2 e^{-\frac{2t^2}{3}}} \quad (1)$$

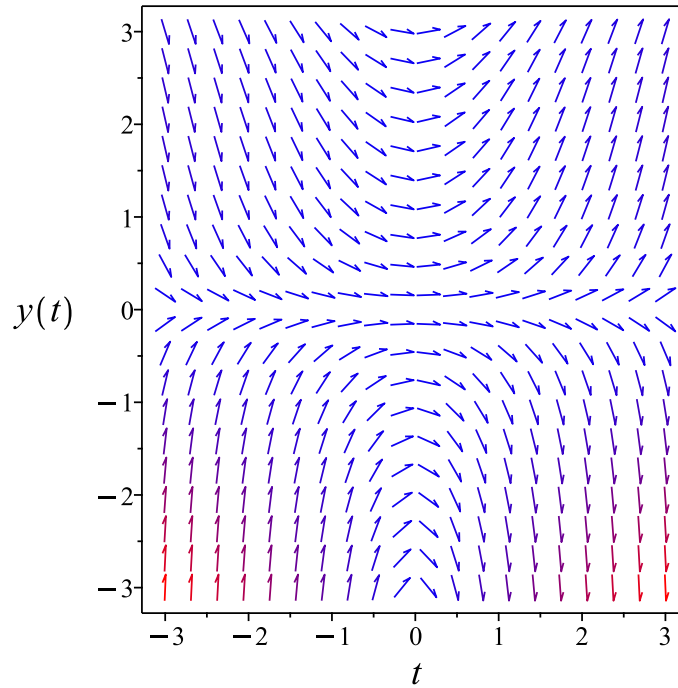


Figure 142: Slope field plot

Verification of solutions

$$y = -\frac{4}{-1 + c_2 e^{-\frac{2t^2}{3}}}$$

Verified OK.

2.27.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{ty(y-4)}{3}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 169: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int(n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= -\frac{3}{t} \\ \eta(t, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dt \\ &= \int \frac{1}{-\frac{3}{t}} dt \end{aligned}$$

Which results in

$$S = -\frac{t^2}{6}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{ty(y-4)}{3}$$

Evaluating all the partial derivatives gives

$$R_t = 0$$

$$R_y = 1$$

$$S_t = -\frac{t}{3}$$

$$S_y = 0$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{y(y-4)} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R(R-4)}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{\ln(R-4)}{4} - \frac{\ln(R)}{4} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$-\frac{t^2}{6} = \frac{\ln(y-4)}{4} - \frac{\ln(y)}{4} + c_1$$

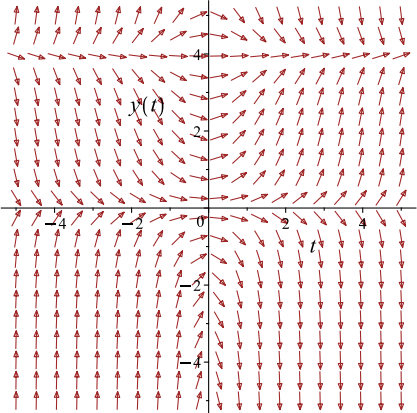
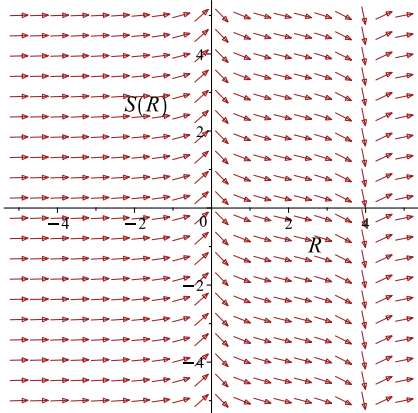
Which simplifies to

$$-\frac{t^2}{6} = \frac{\ln(y-4)}{4} - \frac{\ln(y)}{4} + c_1$$

Which gives

$$y = \frac{4e^{\frac{2t^2}{3} + 4c_1}}{-1 + e^{\frac{2t^2}{3} + 4c_1}}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{ty(y-4)}{3}$ 	$R = y$ $S = -\frac{t^2}{6}$	$\frac{dS}{dR} = \frac{1}{R(R-4)}$ 

Summary

The solution(s) found are the following

$$y = \frac{4 e^{\frac{2t^2}{3} + 4c_1}}{-1 + e^{\frac{2t^2}{3} + 4c_1}} \quad (1)$$

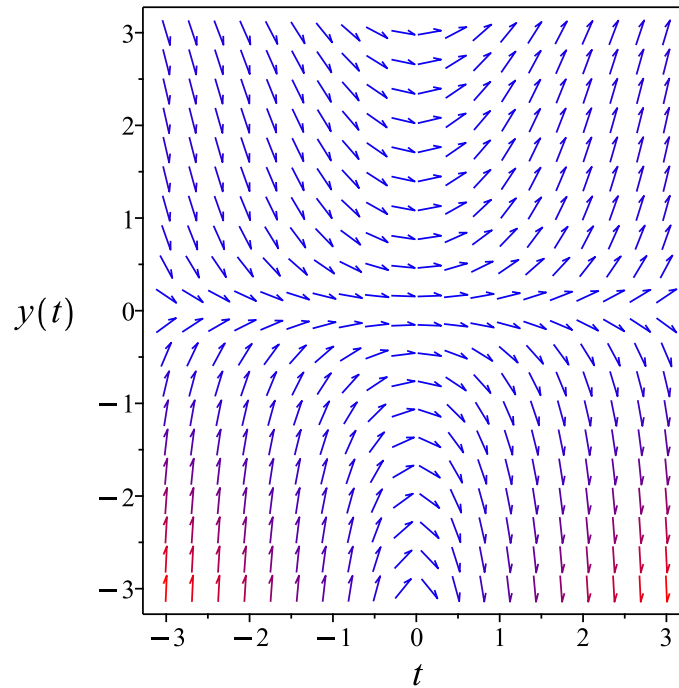


Figure 143: Slope field plot

Verification of solutions

$$y = \frac{4 e^{\frac{2t^2}{3} + 4c_1}}{-1 + e^{\frac{2t^2}{3} + 4c_1}}$$

Verified OK.

2.27.3 Solving as bernoulli ode

In canonical form, the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= -\frac{ty(y-4)}{3} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = \frac{4t}{3}y - \frac{t}{3}y^2 \quad (1)$$

The standard Bernoulli ODE has the form

$$y' = f_0(t)y + f_1(t)y^n \quad (2)$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(t)y^{1-n} + f_1(t) \quad (3)$$

The next step is use the substitution $w = y^{1-n}$ in equation (3) which generates a new ODE in $w(t)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(t)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(t) &= \frac{4t}{3} \\ f_1(t) &= -\frac{t}{3} \\ n &= 2 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = y^2$ gives

$$y' \frac{1}{y^2} = \frac{4t}{3y} - \frac{t}{3} \quad (4)$$

Let

$$\begin{aligned} w &= y^{1-n} \\ &= \frac{1}{y} \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t t gives

$$w' = -\frac{1}{y^2}y' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} -w'(t) &= \frac{4w(t)t}{3} - \frac{t}{3} \\ w' &= -\frac{4}{3}tw + \frac{1}{3}t \end{aligned} \quad (7)$$

The above now is a linear ODE in $w(t)$ which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(t) + p(t)w(t) = q(t)$$

Where here

$$p(t) = \frac{4t}{3}$$
$$q(t) = \frac{t}{3}$$

Hence the ode is

$$w'(t) + \frac{4w(t)t}{3} = \frac{t}{3}$$

The integrating factor μ is

$$\mu = e^{\int \frac{4t}{3} dt}$$
$$= e^{\frac{2t^2}{3}}$$

The ode becomes

$$\frac{d}{dt}(\mu w) = (\mu) \left(\frac{t}{3} \right)$$
$$\frac{d}{dt} \left(e^{\frac{2t^2}{3}} w \right) = \left(e^{\frac{2t^2}{3}} \right) \left(\frac{t}{3} \right)$$
$$d \left(e^{\frac{2t^2}{3}} w \right) = \left(\frac{t e^{\frac{2t^2}{3}}}{3} \right) dt$$

Integrating gives

$$e^{\frac{2t^2}{3}} w = \int \frac{t e^{\frac{2t^2}{3}}}{3} dt$$
$$e^{\frac{2t^2}{3}} w = \frac{e^{\frac{2t^2}{3}}}{4} + c_1$$

Dividing both sides by the integrating factor $\mu = e^{\frac{2t^2}{3}}$ results in

$$w(t) = \frac{e^{-\frac{2t^2}{3}} e^{\frac{2t^2}{3}}}{4} + c_1 e^{-\frac{2t^2}{3}}$$

which simplifies to

$$w(t) = \frac{1}{4} + c_1 e^{-\frac{2t^2}{3}}$$

Replacing w in the above by $\frac{1}{y}$ using equation (5) gives the final solution.

$$\frac{1}{y} = \frac{1}{4} + c_1 e^{-\frac{2t^2}{3}}$$

Or

$$y = \frac{1}{\frac{1}{4} + c_1 e^{-\frac{2t^2}{3}}}$$

Summary

The solution(s) found are the following

$$y = \frac{1}{\frac{1}{4} + c_1 e^{-\frac{2t^2}{3}}} \quad (1)$$

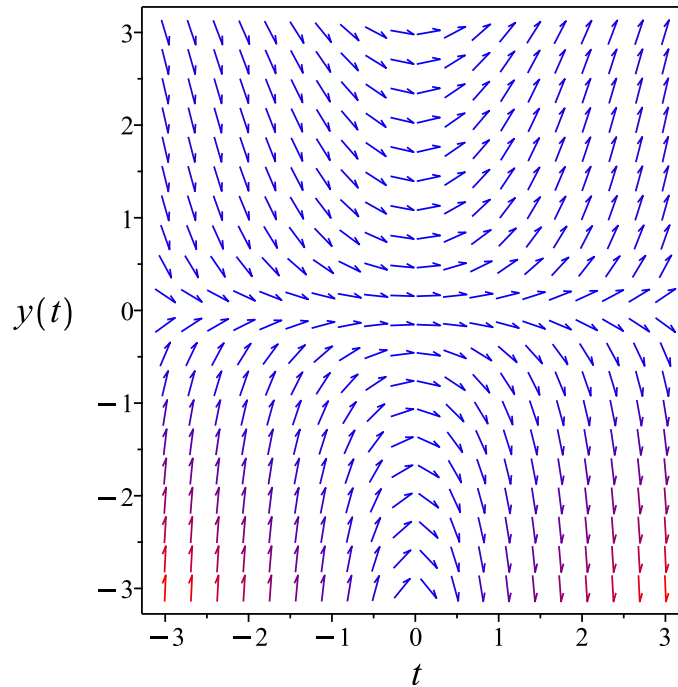


Figure 144: Slope field plot

Verification of solutions

$$y = \frac{1}{\frac{1}{4} + c_1 e^{-\frac{2t^2}{3}}}$$

Verified OK.

2.27.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} &\left(-\frac{3}{y(y-4)}\right) dy = (t) dt \\ (-t) dt + \left(-\frac{3}{y(y-4)}\right) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$M(t, y) = -t$$
$$N(t, y) = -\frac{3}{y(y-4)}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y}(-t)$$
$$= 0$$

And

$$\frac{\partial N}{\partial t} = \frac{\partial}{\partial t} \left(-\frac{3}{y(y-4)} \right)$$
$$= 0$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int M dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int -t dt$$

$$\phi = -\frac{t^2}{2} + f(y) \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{3}{y(y-4)}$. Therefore equation (4) becomes

$$-\frac{3}{y(y-4)} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{3}{y(y-4)}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(-\frac{3}{y(y-4)} \right) dy$$

$$f(y) = -\frac{3 \ln(y-4)}{4} + \frac{3 \ln(y)}{4} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{t^2}{2} - \frac{3 \ln(y-4)}{4} + \frac{3 \ln(y)}{4} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{t^2}{2} - \frac{3 \ln(y-4)}{4} + \frac{3 \ln(y)}{4}$$

The solution becomes

$$y = \frac{4 e^{\frac{2t^2}{3} + \frac{4c_1}{3}}}{-1 + e^{\frac{2t^2}{3} + \frac{4c_1}{3}}}$$

Summary

The solution(s) found are the following

$$y = \frac{4 e^{\frac{2t^2}{3} + \frac{4c_1}{3}}}{-1 + e^{\frac{2t^2}{3} + \frac{4c_1}{3}}} \quad (1)$$

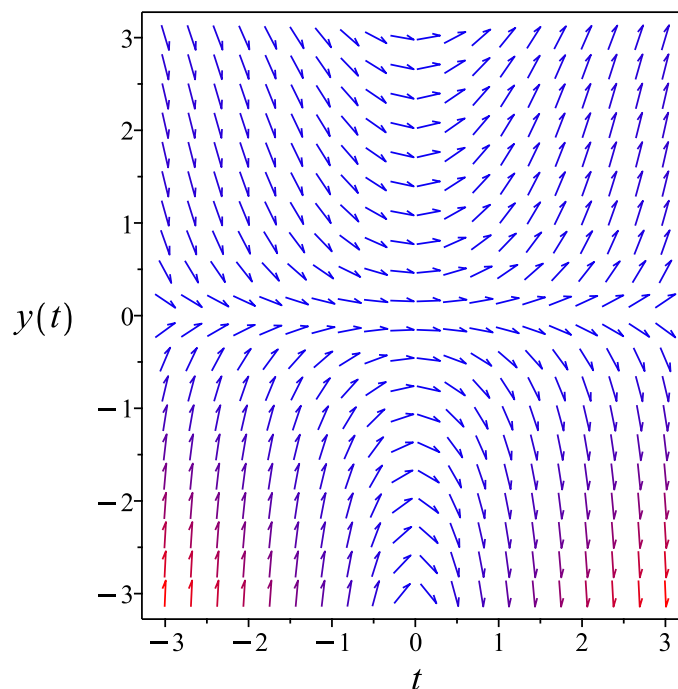


Figure 145: Slope field plot

Verification of solutions

$$y = \frac{4e^{\frac{2t^2}{3} + \frac{4c_1}{3}}}{-1 + e^{\frac{2t^2}{3} + \frac{4c_1}{3}}}$$

Verified OK.

2.27.5 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= -\frac{ty(y-4)}{3} \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = -\frac{1}{3}ty^2 + \frac{4}{3}ty$$

With Riccati ODE standard form

$$y' = f_0(t) + f_1(t)y + f_2(t)y^2$$

Shows that $f_0(t) = 0$, $f_1(t) = \frac{4t}{3}$ and $f_2(t) = -\frac{t}{3}$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{-\frac{tu}{3}} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(t) - (f_2' + f_1 f_2) u'(t) + f_2^2 f_0 u(t) = 0 \tag{2}$$

But

$$\begin{aligned} f_2' &= -\frac{1}{3} \\ f_1 f_2 &= -\frac{4t^2}{9} \\ f_2^2 f_0 &= 0 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$-\frac{tu''(t)}{3} - \left(-\frac{1}{3} - \frac{4t^2}{9}\right) u'(t) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(t) = c_1 + e^{\frac{2t^2}{3}} c_2$$

The above shows that

$$u'(t) = \frac{4t e^{\frac{2t^2}{3}} c_2}{3}$$

Using the above in (1) gives the solution

$$y = \frac{4 e^{\frac{2t^2}{3}} c_2}{c_1 + e^{\frac{2t^2}{3}} c_2}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = \frac{4e^{\frac{2t^2}{3}}}{c_3 + e^{\frac{2t^2}{3}}}$$

Summary

The solution(s) found are the following

$$y = \frac{4e^{\frac{2t^2}{3}}}{c_3 + e^{\frac{2t^2}{3}}} \quad (1)$$

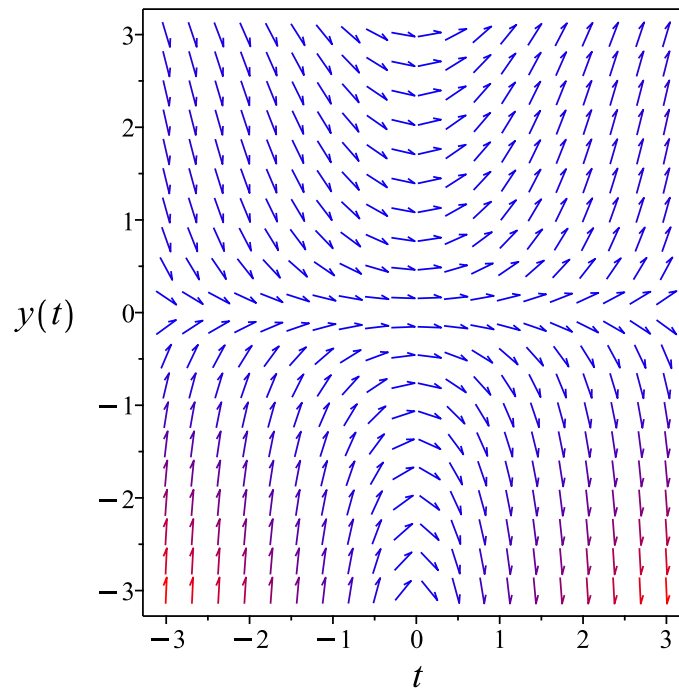


Figure 146: Slope field plot

Verification of solutions

$$y = \frac{4e^{\frac{2t^2}{3}}}{c_3 + e^{\frac{2t^2}{3}}}$$

Verified OK.

2.27.6 Maple step by step solution

Let's solve

$$y' - \frac{t(4-y)y}{3} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'}{(4-y)y} = \frac{t}{3}$$

- Integrate both sides with respect to t

$$\int \frac{y'}{(4-y)y} dt = \int \frac{t}{3} dt + c_1$$

- Evaluate integral

$$-\frac{\ln(y-4)}{4} + \frac{\ln(y)}{4} = \frac{t^2}{6} + c_1$$

- Solve for y

$$y = \frac{4e^{\frac{2t^2}{3} + 4c_1}}{-1 + e^{\frac{2t^2}{3} + 4c_1}}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 19

```
dsolve(diff(y(t),t) = 1/3*t*(4-y(t))*y(t),y(t), singsol=all)
```

$$y(t) = \frac{4}{1 + 4e^{-\frac{2t^2}{3}} c_1}$$

✓ Solution by Mathematica

Time used: 0.248 (sec). Leaf size: 44

```
DSolve[y'[t]== 1/3*t*(4-y[t])*y[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{4e^{\frac{2t^2}{3}}}{e^{\frac{2t^2}{3}} + e^{4c_1}}$$

$$y(t) \rightarrow 0$$

$$y(t) \rightarrow 4$$

2.28 problem 28

2.28.1 Solving as separable ode	788
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Internal problem ID [506]

Internal file name [OUTPUT/506_Sunday_June_05_2022_01_42_39_AM_99102865/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 28.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "bernoulli", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{ty(4-y)}{t+1} = 0$$

2.28.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(t, y) \\ &= f(t)g(y) \\ &= -\frac{ty(y-4)}{t+1}\end{aligned}$$

Where $f(t) = -\frac{t}{t+1}$ and $g(y) = y(y-4)$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y(y-4)} dy &= -\frac{t}{t+1} dt \\ \int \frac{1}{y(y-4)} dy &= \int -\frac{t}{t+1} dt \\ \frac{\ln(y-4)}{4} - \frac{\ln(y)}{4} &= -t + \ln(t+1) + c_1\end{aligned}$$

The above can be written as

$$\begin{aligned}\left(\frac{1}{4}\right) (\ln(y-4) - \ln(y)) &= -t + \ln(t+1) + 2c_1 \\ \ln(y-4) - \ln(y) &= (4)(-t + \ln(t+1) + 2c_1) \\ &= -4t + 4\ln(t+1) + 8c_1\end{aligned}$$

Raising both side to exponential gives

$$e^{\ln(y-4) - \ln(y)} = e^{-4t + 4\ln(t+1) + 4c_1}$$

Which simplifies to

$$\begin{aligned}\frac{y-4}{y} &= 4c_1 e^{-4t + 4\ln(t+1)} \\ &= c_2 e^{-4t + 4\ln(t+1)}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = -\frac{4}{-1 + c_2 e^{-4t + 4\ln(t+1)}} \quad (1)$$

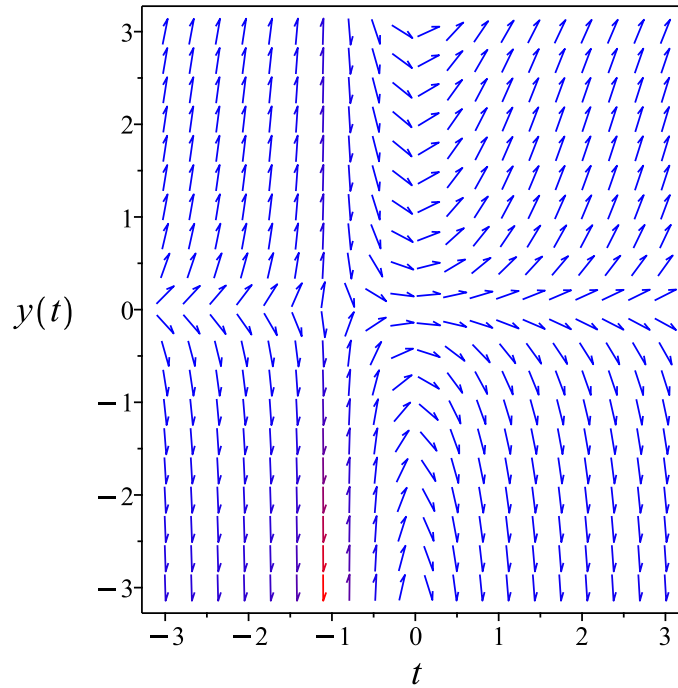


Figure 147: Slope field plot

Verification of solutions

$$y = -\frac{4}{-1 + c_2 e^{-4t+4\ln(t+1)}}$$

Verified OK.

2.28.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{ty(y-4)}{t+1}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 172: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= -\frac{t+1}{t} \\ \eta(t, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dt \\ &= \int \frac{1}{-\frac{t+1}{t}} dt \end{aligned}$$

Which results in

$$S = -t + \ln(t + 1)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{ty(y - 4)}{t + 1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 0 \\ R_y &= 1 \\ S_t &= -\frac{t}{t + 1} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{y(y - 4)} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R(R - 4)}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{\ln(R-4)}{4} - \frac{\ln(R)}{4} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$-t + \ln(t+1) = \frac{\ln(y-4)}{4} - \frac{\ln(y)}{4} + c_1$$

Which simplifies to

$$-t + \ln(t+1) = \frac{\ln(y-4)}{4} - \frac{\ln(y)}{4} + c_1$$

Which gives

$$y = \frac{4e^{4t+4c_1}}{-1-t^4-4t^3-6t^2+e^{4t+4c_1}-4t}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{ty(y-4)}{t+1}$	$R = y$ $S = -t + \ln(t+1)$	$\frac{dS}{dR} = \frac{1}{R(R-4)}$

Summary

The solution(s) found are the following

$$y = \frac{4e^{4t+4c_1}}{-1 - t^4 - 4t^3 - 6t^2 + e^{4t+4c_1} - 4t} \quad (1)$$

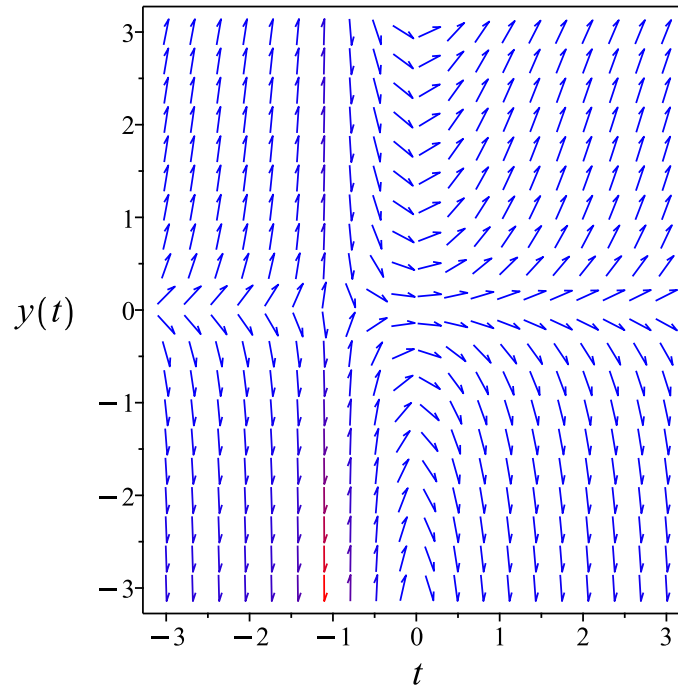


Figure 148: Slope field plot

Verification of solutions

$$y = \frac{4e^{4t+4c_1}}{-1 - t^4 - 4t^3 - 6t^2 + e^{4t+4c_1} - 4t}$$

Verified OK.

2.28.3 Solving as bernoulli ode

In canonical form, the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= -\frac{ty(y-4)}{t+1} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = \frac{4t}{t+1}y - \frac{t}{t+1}y^2 \quad (1)$$

The standard Bernoulli ODE has the form

$$y' = f_0(t)y + f_1(t)y^n \quad (2)$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(t)y^{1-n} + f_1(t) \quad (3)$$

The next step is use the substitution $w = y^{1-n}$ in equation (3) which generates a new ODE in $w(t)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(t)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(t) &= \frac{4t}{t+1} \\ f_1(t) &= -\frac{t}{t+1} \\ n &= 2 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = y^2$ gives

$$y' \frac{1}{y^2} = \frac{4t}{(t+1)y} - \frac{t}{t+1} \quad (4)$$

Let

$$\begin{aligned} w &= y^{1-n} \\ &= \frac{1}{y} \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t t gives

$$w' = -\frac{1}{y^2}y' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} -w'(t) &= \frac{4tw(t)}{t+1} - \frac{t}{t+1} \\ w' &= -\frac{4tw}{t+1} + \frac{t}{t+1} \end{aligned} \quad (7)$$

The above now is a linear ODE in $w(t)$ which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(t) + p(t)w(t) = q(t)$$

Where here

$$p(t) = \frac{4t}{t+1}$$
$$q(t) = \frac{t}{t+1}$$

Hence the ode is

$$w'(t) + \frac{4tw(t)}{t+1} = \frac{t}{t+1}$$

The integrating factor μ is

$$\mu = e^{\int \frac{4t}{t+1} dt}$$
$$= e^{4t-4\ln(t+1)}$$

Which simplifies to

$$\mu = \frac{e^{4t}}{(t+1)^4}$$

The ode becomes

$$\frac{d}{dt}(\mu w) = (\mu) \left(\frac{t}{t+1} \right)$$
$$\frac{d}{dt} \left(\frac{e^{4t}w}{(t+1)^4} \right) = \left(\frac{e^{4t}}{(t+1)^4} \right) \left(\frac{t}{t+1} \right)$$
$$d \left(\frac{e^{4t}w}{(t+1)^4} \right) = \left(\frac{t e^{4t}}{(t+1)^5} \right) dt$$

Integrating gives

$$\frac{e^{4t}w}{(t+1)^4} = \int \frac{t e^{4t}}{(t+1)^5} dt$$
$$\frac{e^{4t}w}{(t+1)^4} = \frac{e^{4t}}{4(t+1)^4} + c_1$$

Dividing both sides by the integrating factor $\mu = \frac{e^{4t}}{(t+1)^4}$ results in

$$w(t) = \frac{e^{-4t}e^{4t}}{4} + c_1 e^{-4t}(t+1)^4$$

which simplifies to

$$w(t) = \frac{1}{4} + c_1 e^{-4t} (t+1)^4$$

Replacing w in the above by $\frac{1}{y}$ using equation (5) gives the final solution.

$$\frac{1}{y} = \frac{1}{4} + c_1 e^{-4t} (t+1)^4$$

Or

$$y = \frac{1}{\frac{1}{4} + c_1 e^{-4t} (t+1)^4}$$

Summary

The solution(s) found are the following

$$y = \frac{1}{\frac{1}{4} + c_1 e^{-4t} (t+1)^4} \tag{1}$$

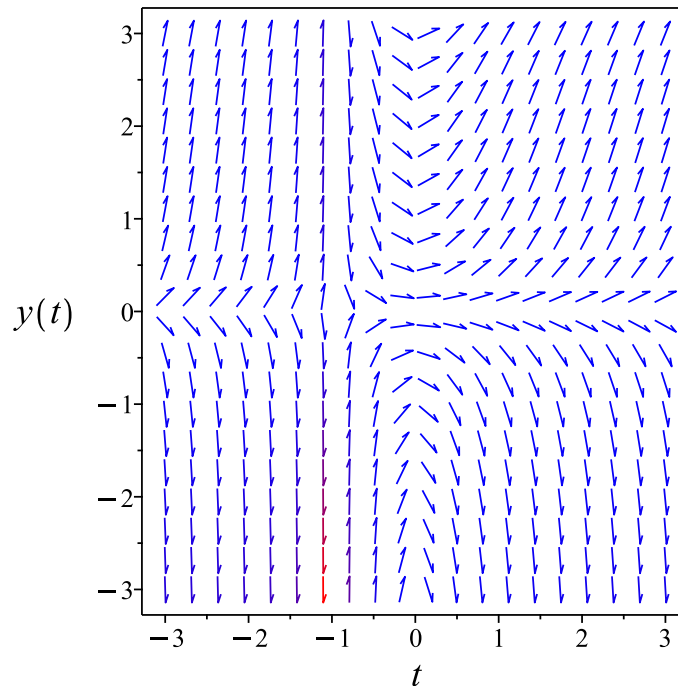


Figure 149: Slope field plot

Verification of solutions

$$y = \frac{1}{\frac{1}{4} + c_1 e^{-4t} (t+1)^4}$$

Verified OK.

2.28.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} \left(-\frac{1}{y(y-4)} \right) dy &= \left(\frac{t}{t+1} \right) dt \\ \left(-\frac{t}{t+1} \right) dt + \left(-\frac{1}{y(y-4)} \right) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$M(t, y) = -\frac{t}{t+1}$$
$$N(t, y) = -\frac{1}{y(y-4)}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y} \left(-\frac{t}{t+1} \right)$$
$$= 0$$

And

$$\frac{\partial N}{\partial t} = \frac{\partial}{\partial t} \left(-\frac{1}{y(y-4)} \right)$$
$$= 0$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int M dt$$
$$\int \frac{\partial \phi}{\partial t} dt = \int -\frac{t}{t+1} dt$$
$$\phi = -t + \ln(t+1) + f(y) \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{1}{y(y-4)}$. Therefore equation (4) becomes

$$-\frac{1}{y(y-4)} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{1}{y(y-4)}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(-\frac{1}{y(y-4)} \right) dy$$

$$f(y) = -\frac{\ln(y-4)}{4} + \frac{\ln(y)}{4} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -t + \ln(t+1) - \frac{\ln(y-4)}{4} + \frac{\ln(y)}{4} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -t + \ln(t+1) - \frac{\ln(y-4)}{4} + \frac{\ln(y)}{4}$$

The solution becomes

$$y = \frac{4e^{4t+4c_1}}{-1 - t^4 - 4t^3 - 6t^2 + e^{4t+4c_1} - 4t}$$

Summary

The solution(s) found are the following

$$y = \frac{4e^{4t+4c_1}}{-1 - t^4 - 4t^3 - 6t^2 + e^{4t+4c_1} - 4t} \quad (1)$$

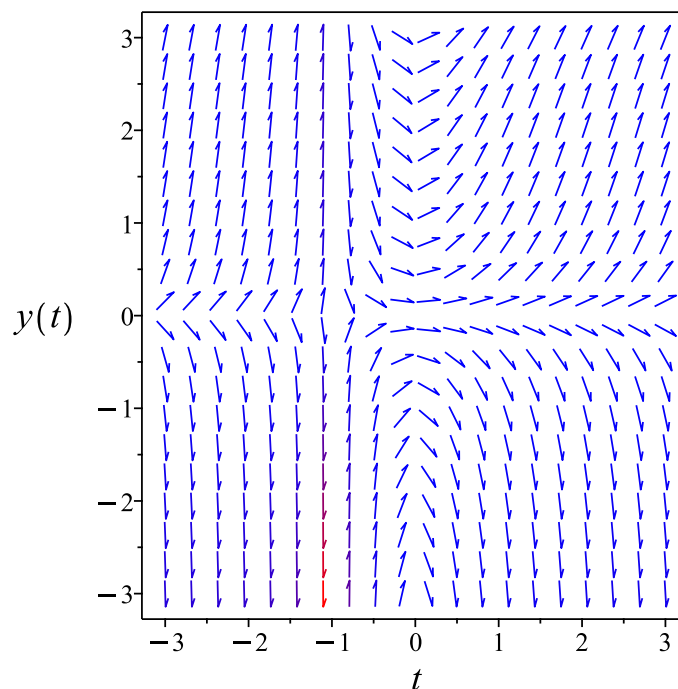


Figure 150: Slope field plot

Verification of solutions

$$y = \frac{4e^{4t+4c_1}}{-1 - t^4 - 4t^3 - 6t^2 + e^{4t+4c_1} - 4t}$$

Verified OK.

2.28.5 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= -\frac{ty(y-4)}{t+1} \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = \frac{4ty}{t+1} - \frac{ty^2}{t+1}$$

With Riccati ODE standard form

$$y' = f_0(t) + f_1(t)y + f_2(t)y^2$$

Shows that $f_0(t) = 0$, $f_1(t) = \frac{4t}{t+1}$ and $f_2(t) = -\frac{t}{t+1}$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{-\frac{tu}{t+1}} \end{aligned} \quad (1)$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(t) - (f_2' + f_1 f_2) u'(t) + f_2^2 f_0 u(t) = 0 \quad (2)$$

But

$$\begin{aligned} f_2' &= -\frac{1}{t+1} + \frac{t}{(t+1)^2} \\ f_1 f_2 &= -\frac{4t^2}{(t+1)^2} \\ f_2^2 f_0 &= 0 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$-\frac{tu''(t)}{t+1} - \left(-\frac{1}{t+1} + \frac{t}{(t+1)^2} - \frac{4t^2}{(t+1)^2} \right) u'(t) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(t) = \frac{c_2 e^{4t} + (t+1)^4 c_1}{(t+1)^4}$$

The above shows that

$$u'(t) = \frac{4e^{4t} c_2 t}{(t+1)^5}$$

Using the above in (1) gives the solution

$$y = \frac{4e^{4t} c_2}{c_2 e^{4t} + (t+1)^4 c_1}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = \frac{4 e^{4t}}{e^{4t} + c_3 (t^4 + 4t^3 + 6t^2 + 4t + 1)}$$

Summary

The solution(s) found are the following

$$y = \frac{4 e^{4t}}{e^{4t} + c_3 (t^4 + 4t^3 + 6t^2 + 4t + 1)} \quad (1)$$

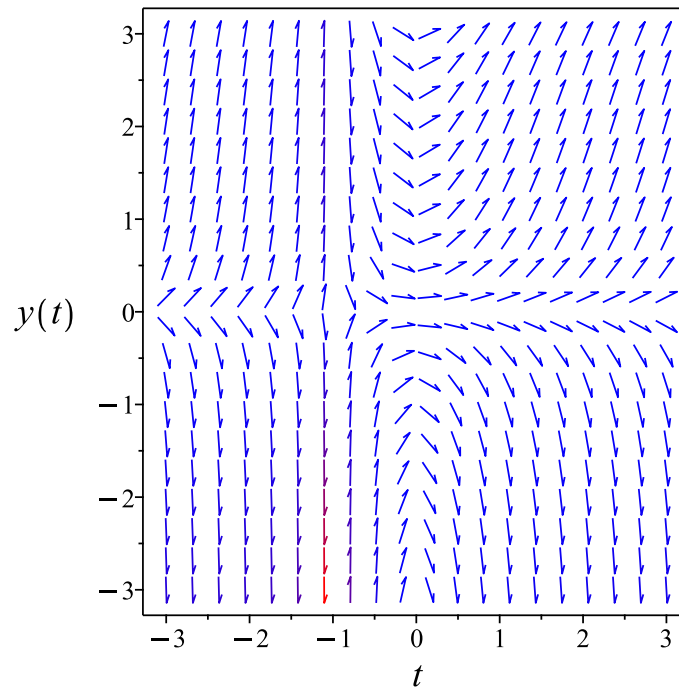


Figure 151: Slope field plot

Verification of solutions

$$y = \frac{4 e^{4t}}{e^{4t} + c_3 (t^4 + 4t^3 + 6t^2 + 4t + 1)}$$

Verified OK.

2.28.6 Maple step by step solution

Let's solve

$$y' - \frac{ty(4-y)}{t+1} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'}{(4-y)y} = \frac{t}{t+1}$$

- Integrate both sides with respect to t

$$\int \frac{y'}{(4-y)y} dt = \int \frac{t}{t+1} dt + c_1$$

- Evaluate integral

$$-\frac{\ln(y-4)}{4} + \frac{\ln(y)}{4} = t - \ln(t+1) + c_1$$

- Solve for y

$$y = \frac{4e^{4t+4c_1}}{-1-t^4-4t^3-6t^2+e^{4t+4c_1}-4t}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 22

```
dsolve(diff(y(t),t) = t*y(t)*(4-y(t))/(1+t),y(t), singsol=all)
```

$$y(t) = \frac{4}{1 + 4e^{-4t}(t+1)^4 c_1}$$

✓ Solution by Mathematica

Time used: 1.337 (sec). Leaf size: 42

```
DSolve[y'[t] == t*y[t]*(4-y[t])/(1+t),y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{4e^{4t}}{e^{4t} + e^{4c_1}(t+1)^4}$$

$$y(t) \rightarrow 0$$

$$y(t) \rightarrow 4$$

2.29 problem 29

2.29.1 Solving as quadrature ode	806
2.29.2 Maple step by step solution	807

Internal problem ID [507]

Internal file name [OUTPUT/507_Sunday_June_05_2022_01_42_40_AM_35330923/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 29.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"quadrature"**

Maple gives the following as the ode type

`[_quadrature]`

$$y' - \frac{b + ay}{d + cy} = 0$$

2.29.1 Solving as quadrature ode

Integrating both sides gives

$$\int \frac{cy + d}{ay + b} dy = x + c_1$$

$$\frac{cy}{a} + \frac{(ad - bc) \ln (ay + b)}{a^2} = x + c_1$$

Solving for y gives these solutions

Summary

The solution(s) found are the following

$$y = \frac{c_1 a^2 + x a^2 - \left(-\text{LambertW} \left(-\frac{c e^{\frac{c_1 a^2 + x a^2 + bc}{ad - bc}}}{-ad + bc} \right) + \frac{c_1 a^2 + x a^2 + bc}{ad - bc} \right) ad + \left(-\text{LambertW} \left(-\frac{c e^{\frac{c_1 a^2 + x a^2 + bc}{ad - bc}}}{-ad + bc} \right) \right)}{ac} \quad (1)$$

Verification of solutions

y

$$= \frac{c_1 a^2 + x a^2 - \left(-\text{LambertW} \left(-\frac{c e^{\frac{c_1 a^2 + x a^2 + bc}{ad-bc}}}{-ad+bc} \right) + \frac{c_1 a^2 + x a^2 + bc}{ad-bc} \right) ad + \left(-\text{LambertW} \left(-\frac{c e^{\frac{c_1 a^2 + x a^2 + bc}{ad-bc}}}{-ad+bc} \right) + \frac{c_1 a^2 + x a^2 + bc}{ad-bc} \right) bc}{ac}$$

Verified OK.

2.29.2 Maple step by step solution

Let's solve

$$y' - \frac{b+ay}{d+cy} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'(d+cy)}{b+ay} = 1$$

- Integrate both sides with respect to x

$$\int \frac{y'(d+cy)}{b+ay} dx = \int 1 dx + c_1$$

- Evaluate integral

$$\frac{cy}{a} + \frac{(ad-bc)\ln(b+ay)}{a^2} = x + c_1$$

- Solve for y

$$y = \frac{c_1 a^2 + x a^2 - \left(-\text{LambertW} \left(-\frac{c e^{\frac{c_1 a^2 + x a^2 + bc}{ad-bc}}}{-ad+bc} \right) + \frac{c_1 a^2 + x a^2 + bc}{ad-bc} \right) ad + \left(-\text{LambertW} \left(-\frac{c e^{\frac{c_1 a^2 + x a^2 + bc}{ad-bc}}}{-ad+bc} \right) + \frac{c_1 a^2 + x a^2 + bc}{ad-bc} \right) bc}{ac}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```


✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 61

```
dsolve(diff(y(x),x) = (b+a*y(x))/(d+c*y(x)),y(x), singsol=all)
```

$$y(x) = \frac{(ad - bc) \operatorname{LambertW}\left(\frac{ce^{\frac{(c_1+x)a^2+bc}{ad-bc}}}{ad-bc}\right) - bc}{ac}$$

✓ Solution by Mathematica

Time used: 16.166 (sec). Leaf size: 83

```
DSolve[y'[x] == (b+a*y[x])/(d+c*y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{-bc + (ad - bc)W\left(-\frac{c\left(e^{-1-\frac{a^2(x+c_1)}{bc}}\right)^{\frac{bc}{bc-ad}}}{bc-ad}\right)}{ac}$$
$$y(x) \rightarrow -\frac{b}{a}$$

2.30 problem 31

2.30.1 Solving as homogeneousTypeD2 ode	809
2.30.2 Solving as first order ode lie symmetry calculated ode	811
2.30.3 Solving as riccati ode	816

Internal problem ID [508]

Internal file name [OUTPUT/508_Sunday_June_05_2022_01_42_41_AM_3383851/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 31.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "riccati", "homogeneousTypeD2", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _Riccati]
```

$$y' - \frac{x^2 + yx + y^2}{x^2} = 0$$

2.30.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u'(x)x + u(x) - \frac{x^2 + u(x)x^2 + u(x)^2x^2}{x^2} = 0$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{u^2 + 1}{x}\end{aligned}$$

Where $f(x) = \frac{1}{x}$ and $g(u) = u^2 + 1$. Integrating both sides gives

$$\frac{1}{u^2 + 1} du = \frac{1}{x} dx$$

$$\int \frac{1}{u^2 + 1} du = \int \frac{1}{x} dx$$

$$\arctan(u) = \ln(x) + c_2$$

The solution is

$$\arctan(u(x)) - \ln(x) - c_2 = 0$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\arctan\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0$$

$$\arctan\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0$$

Summary

The solution(s) found are the following

$$\arctan\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0 \tag{1}$$

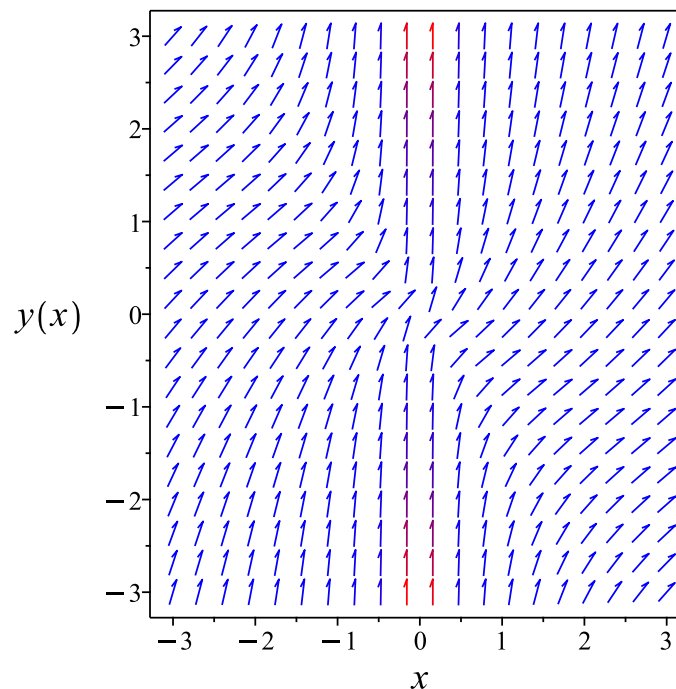


Figure 152: Slope field plot

Verification of solutions

$$\arctan\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0$$

Verified OK.

2.30.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{x^2 + yx + y^2}{x^2}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + \frac{(x^2 + yx + y^2)(b_3 - a_2)}{x^2} - \frac{(x^2 + yx + y^2)^2 a_3}{x^4} - \left(\frac{2x + y}{x^2} - \frac{2(x^2 + yx + y^2)}{x^3} \right) (xa_2 + ya_3 + a_1) - \frac{(x + 2y)(xb_2 + yb_3 + b_1)}{x^2} = 0 \quad (\text{5E})$$

Putting the above in normal form gives

$$\frac{x^4 a_2 + x^4 a_3 - x^4 b_3 + 2x^3 y a_3 + 2x^3 y b_2 - x^2 y^2 a_2 + 2x^2 y^2 a_3 + x^2 y^2 b_3 + y^4 a_3 + x^3 b_1 - x^2 y a_1 + 2x^2 y b_1 -}{x^4} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} & -x^4a_2 - x^4a_3 + x^4b_3 - 2x^3ya_3 - 2x^3yb_2 + x^2y^2a_2 - 2x^2y^2a_3 \\ & - x^2y^2b_3 - y^4a_3 - x^3b_1 + x^2ya_1 - 2x^2yb_1 + 2xy^2a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & -a_2v_1^4 + a_2v_1^2v_2^2 - a_3v_1^4 - 2a_3v_1^3v_2 - 2a_3v_1^2v_2^2 - a_3v_2^4 - 2b_2v_1^3v_2 \\ & + b_3v_1^4 - b_3v_1^2v_2^2 + a_1v_1^2v_2 + 2a_1v_1v_2^2 - b_1v_1^3 - 2b_1v_1^2v_2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & (-a_2 - a_3 + b_3)v_1^4 + (-2a_3 - 2b_2)v_1^3v_2 - b_1v_1^3 \\ & + (a_2 - 2a_3 - b_3)v_1^2v_2^2 + (a_1 - 2b_1)v_1^2v_2 + 2a_1v_1v_2^2 - a_3v_2^4 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 2a_1 &= 0 \\ -a_3 &= 0 \\ -b_1 &= 0 \\ a_1 - 2b_1 &= 0 \\ -2a_3 - 2b_2 &= 0 \\ -a_2 - a_3 + b_3 &= 0 \\ a_2 - 2a_3 - b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$a_1 = 0$$

$$a_2 = b_3$$

$$a_3 = 0$$

$$b_1 = 0$$

$$b_2 = 0$$

$$b_3 = b_3$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\xi = x$$

$$\eta = y$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= y - \left(\frac{x^2 + yx + y^2}{x^2} \right) (x) \\ &= \frac{-x^2 - y^2}{x} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-x^2 - y^2}{x}} dy \end{aligned}$$

Which results in

$$S = -\arctan\left(\frac{y}{x}\right)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x^2 + yx + y^2}{x^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{y}{x^2 + y^2} \\ S_y &= -\frac{x}{x^2 + y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{x} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\ln(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-\arctan\left(\frac{y}{x}\right) = -\ln(x) + c_1$$

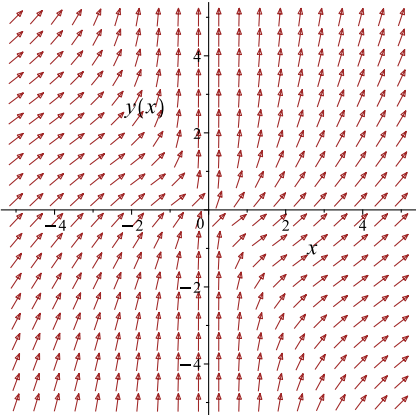
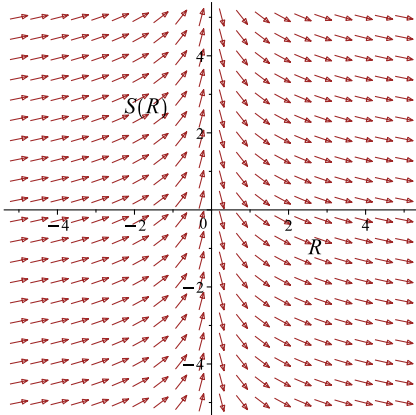
Which simplifies to

$$-\arctan\left(\frac{y}{x}\right) = -\ln(x) + c_1$$

Which gives

$$y = -\tan(-\ln(x) + c_1)x$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{x^2 + yx + y^2}{x^2}$ 	$R = x$ $S = -\arctan\left(\frac{y}{x}\right)$	$\frac{dS}{dR} = -\frac{1}{R}$ 

Summary

The solution(s) found are the following

$$y = -\tan(-\ln(x) + c_1)x \tag{1}$$

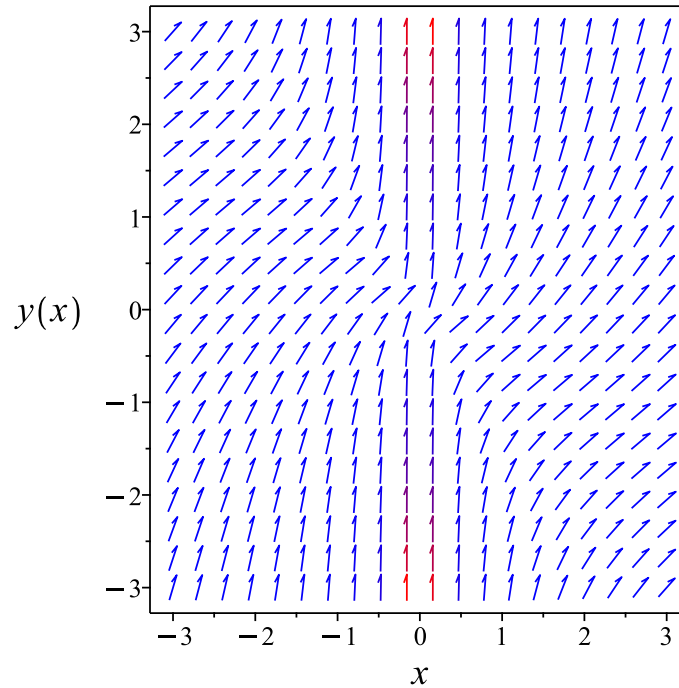


Figure 153: Slope field plot

Verification of solutions

$$y = -\tan(-\ln(x) + c_1)x$$

Verified OK.

2.30.3 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{x^2 + yx + y^2}{x^2} \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = 1 + \frac{y}{x} + \frac{y^2}{x^2}$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that $f_0(x) = 1$, $f_1(x) = \frac{1}{x}$ and $f_2(x) = \frac{1}{x^2}$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{\frac{u}{x^2}} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \tag{2}$$

But

$$\begin{aligned} f_2' &= -\frac{2}{x^3} \\ f_1 f_2 &= \frac{1}{x^3} \\ f_2^2 f_0 &= \frac{1}{x^4} \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$\frac{u''(x)}{x^2} + \frac{u'(x)}{x^3} + \frac{u(x)}{x^4} = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 \sin(\ln(x)) + c_2 \cos(\ln(x))$$

The above shows that

$$u'(x) = \frac{c_1 \cos(\ln(x)) - c_2 \sin(\ln(x))}{x}$$

Using the above in (1) gives the solution

$$y = -\frac{x(c_1 \cos(\ln(x)) - c_2 \sin(\ln(x)))}{c_1 \sin(\ln(x)) + c_2 \cos(\ln(x))}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = \frac{(-c_3 \cos(\ln(x)) + \sin(\ln(x)))x}{c_3 \sin(\ln(x)) + \cos(\ln(x))}$$

Summary

The solution(s) found are the following

$$y = \frac{(-c_3 \cos(\ln(x)) + \sin(\ln(x)))x}{c_3 \sin(\ln(x)) + \cos(\ln(x))} \quad (1)$$

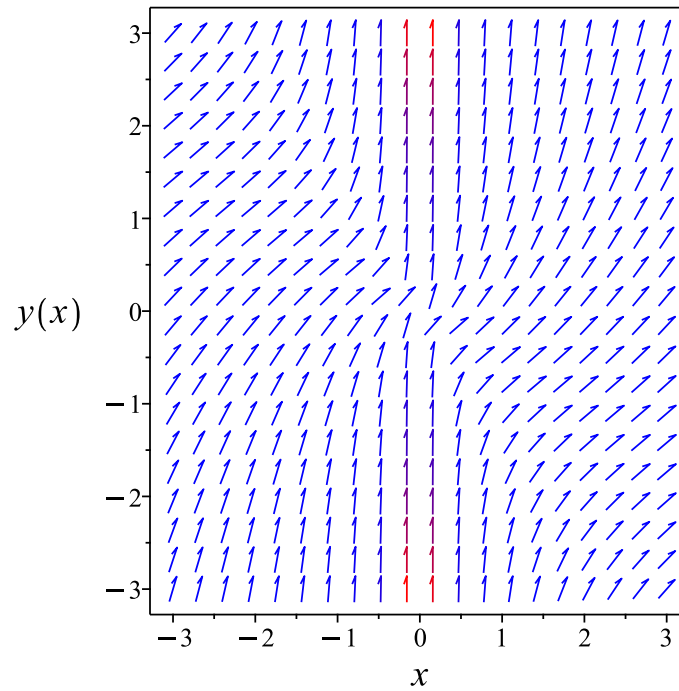


Figure 154: Slope field plot

Verification of solutions

$$y = \frac{(-c_3 \cos(\ln(x)) + \sin(\ln(x)))x}{c_3 \sin(\ln(x)) + \cos(\ln(x))}$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 11

```
dsolve(diff(y(x),x) = (x^2+x*y(x)+y(x)^2)/x^2,y(x), singsol=all)
```

$$y(x) = \tan(\ln(x) + c_1) x$$

✓ Solution by Mathematica

Time used: 0.199 (sec). Leaf size: 13

```
DSolve[y'[x] == (x^2+x*y[x]+y[x]^2)/x^2,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow x \tan(\log(x) + c_1)$$

2.31 problem 32

2.31.1 Solving as homogeneousTypeD2 ode	820
2.31.2 Solving as first order ode lie symmetry lookup ode	822
2.31.3 Solving as bernoulli ode	826
2.31.4 Solving as exact ode	829

Internal problem ID [509]

Internal file name [OUTPUT/509_Sunday_June_05_2022_01_42_42_AM_58248688/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 32.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**bernoulli**", "**homogeneousTypeD2**", "**exactWithIntegrationFactor**", "**first_order_ode_lie_symmetry_lookup**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _Bernoulli]
```

$$y' - \frac{x^2 + 3y^2}{2xy} = 0$$

2.31.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u'(x)x + u(x) - \frac{x^2 + 3u(x)^2x^2}{2x^2u(x)} = 0$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{u^2 + 1}{2xu}\end{aligned}$$

Where $f(x) = \frac{1}{2x}$ and $g(u) = \frac{u^2+1}{u}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{u^2+1}{u}} du &= \frac{1}{2x} dx \\ \int \frac{1}{\frac{u^2+1}{u}} du &= \int \frac{1}{2x} dx \\ \frac{\ln(u^2 + 1)}{2} &= \frac{\ln(x)}{2} + c_2\end{aligned}$$

Raising both side to exponential gives

$$\sqrt{u^2 + 1} = e^{\frac{\ln(x)}{2} + c_2}$$

Which simplifies to

$$\sqrt{u^2 + 1} = c_3\sqrt{x}$$

Which simplifies to

$$\sqrt{u(x)^2 + 1} = c_3\sqrt{x} e^{c_2}$$

The solution is

$$\sqrt{u(x)^2 + 1} = c_3\sqrt{x} e^{c_2}$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\begin{aligned}\sqrt{\frac{y^2}{x^2} + 1} &= c_3\sqrt{x} e^{c_2} \\ \sqrt{\frac{x^2 + y^2}{x^2}} &= c_3\sqrt{x} e^{c_2}\end{aligned}$$

Summary

The solution(s) found are the following

$$\sqrt{\frac{x^2 + y^2}{x^2}} = c_3\sqrt{x} e^{c_2} \quad (1)$$

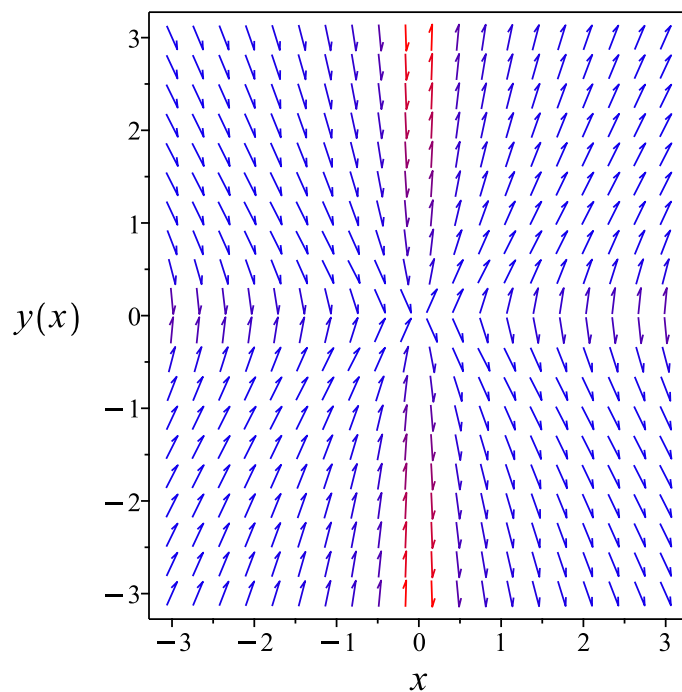


Figure 155: Slope field plot

Verification of solutions

$$\sqrt{\frac{x^2 + y^2}{x^2}} = c_3 \sqrt{x} e^{c_2}$$

Verified OK.

2.31.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x^2 + 3y^2}{2xy}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **Bernoulli**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 176: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{x^3}{y}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{x^3}{y}} dy \end{aligned}$$

Which results in

$$S = \frac{y^2}{2x^3}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x^2 + 3y^2}{2xy}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -\frac{3y^2}{2x^4} \\ S_y &= \frac{y}{x^3} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{2x^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{2R^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{1}{2R} + c_1 \quad (4)$$

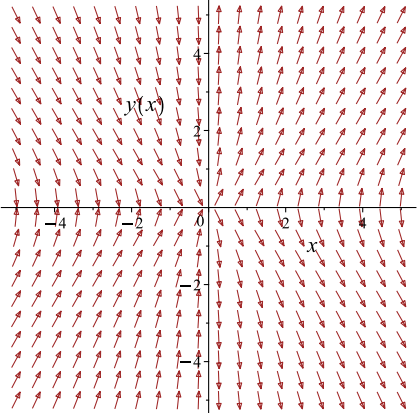
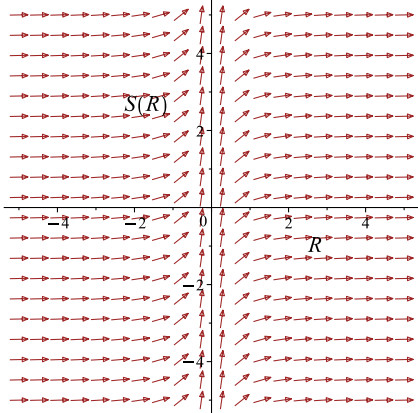
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{y^2}{2x^3} = -\frac{1}{2x} + c_1$$

Which simplifies to

$$\frac{y^2}{2x^3} = -\frac{1}{2x} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{x^2 + 3y^2}{2xy}$ 	$R = x$ $S = \frac{y^2}{2x^3}$	$\frac{dS}{dR} = \frac{1}{2R^2}$ 

Summary

The solution(s) found are the following

$$\frac{y^2}{2x^3} = -\frac{1}{2x} + c_1 \quad (1)$$

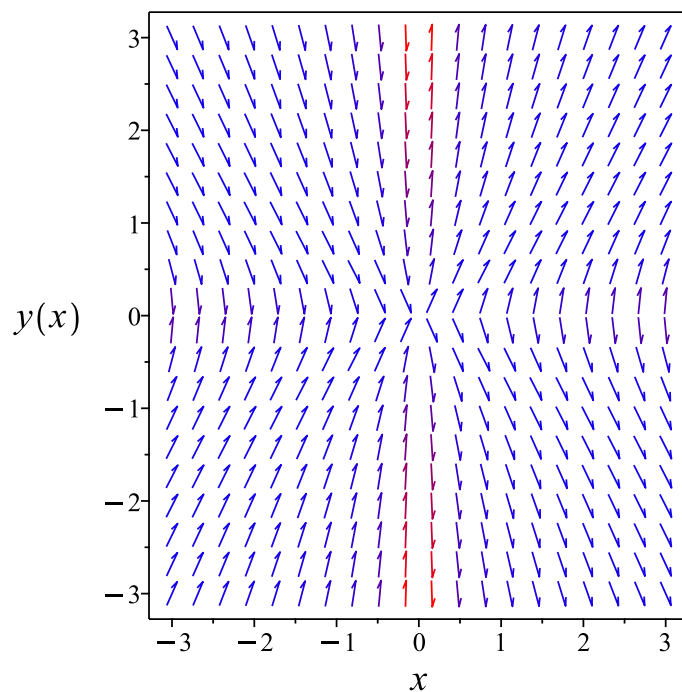


Figure 156: Slope field plot

Verification of solutions

$$\frac{y^2}{2x^3} = -\frac{1}{2x} + c_1$$

Verified OK.

2.31.3 Solving as bernoulli ode

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{x^2 + 3y^2}{2xy} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = \frac{3}{2x}y + \frac{x}{2} \frac{1}{y} \tag{1}$$

The standard Bernoulli ODE has the form

$$y' = f_0(x)y + f_1(x)y^n \tag{2}$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(x)y^{1-n} + f_1(x) \tag{3}$$

The next step is use the substitution $w = y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(x) &= \frac{3}{2x} \\ f_1(x) &= \frac{x}{2} \\ n &= -1 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = \frac{1}{y}$ gives

$$y'y = \frac{3y^2}{2x} + \frac{x}{2} \quad (4)$$

Let

$$\begin{aligned} w &= y^{1-n} \\ &= y^2 \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t x gives

$$w' = 2yy' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} \frac{w'(x)}{2} &= \frac{3w(x)}{2x} + \frac{x}{2} \\ w' &= \frac{3w}{x} + x \end{aligned} \quad (7)$$

The above now is a linear ODE in $w(x)$ which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(x) + p(x)w(x) = q(x)$$

Where here

$$\begin{aligned} p(x) &= -\frac{3}{x} \\ q(x) &= x \end{aligned}$$

Hence the ode is

$$w'(x) - \frac{3w(x)}{x} = x$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{3}{x} dx} \\ &= \frac{1}{x^3}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu w) &= (\mu)(x) \\ \frac{d}{dx}\left(\frac{w}{x^3}\right) &= \left(\frac{1}{x^3}\right)(x) \\ d\left(\frac{w}{x^3}\right) &= \frac{1}{x^2} dx\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{w}{x^3} &= \int \frac{1}{x^2} dx \\ \frac{w}{x^3} &= -\frac{1}{x} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = \frac{1}{x^3}$ results in

$$w(x) = c_1 x^3 - x^2$$

Replacing w in the above by y^2 using equation (5) gives the final solution.

$$y^2 = c_1 x^3 - x^2$$

Solving for y gives

$$\begin{aligned}y(x) &= \sqrt{c_1 x - 1} x \\ y(x) &= -\sqrt{c_1 x - 1} x\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \sqrt{c_1 x - 1} x \tag{1}$$

$$y = -\sqrt{c_1 x - 1} x \tag{2}$$

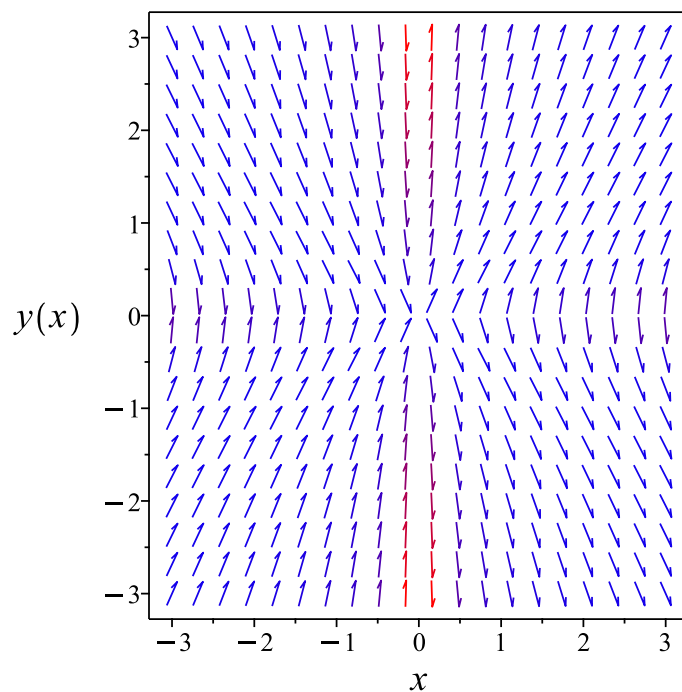


Figure 157: Slope field plot

Verification of solutions

$$y = \sqrt{c_1 x - 1} x$$

Verified OK.

$$y = -\sqrt{c_1 x - 1} x$$

Verified OK.

2.31.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (2yx) dy &= (x^2 + 3y^2) dx \\ (-x^2 - 3y^2) dx + (2yx) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x^2 - 3y^2 \\ N(x, y) &= 2yx \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (-x^2 - 3y^2) \\ &= -6y \end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(2yx) \\ &= 2y\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{2yx} ((-6y) - (2y)) \\ &= -\frac{4}{x}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int -\frac{4}{x} dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-4 \ln(x)} \\ &= \frac{1}{x^4}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{x^4}(-x^2 - 3y^2) \\ &= \frac{-x^2 - 3y^2}{x^4}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{x^4}(2yx) \\ &= \frac{2y}{x^3}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \overline{M} + \overline{N} \frac{dy}{dx} &= 0 \\ \left(\frac{-x^2 - 3y^2}{x^4} \right) + \left(\frac{2y}{x^3} \right) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{-x^2 - 3y^2}{x^4} dx \\ \phi &= \frac{x^2 + y^2}{x^3} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = \frac{2y}{x^3} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{2y}{x^3}$. Therefore equation (4) becomes

$$\frac{2y}{x^3} = \frac{2y}{x^3} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{x^2 + y^2}{x^3} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{x^2 + y^2}{x^3}$$

Summary

The solution(s) found are the following

$$\frac{x^2 + y^2}{x^3} = c_1 \tag{1}$$

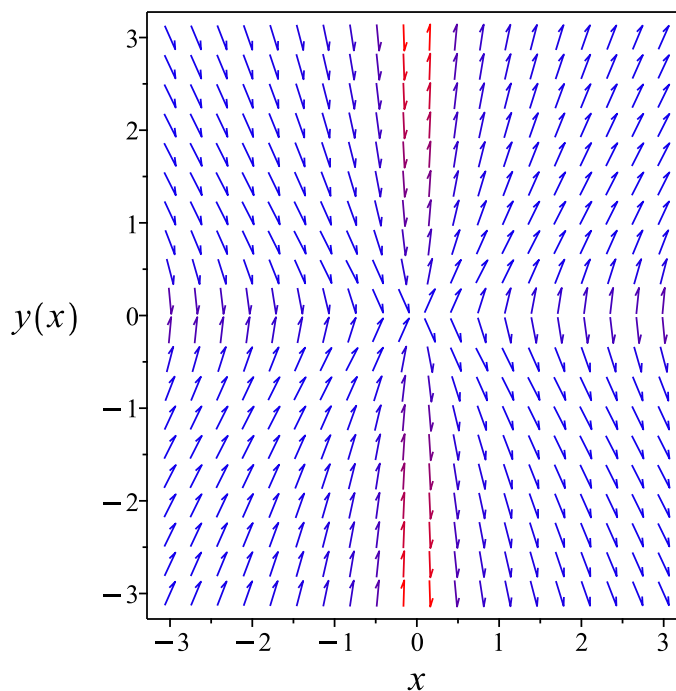


Figure 158: Slope field plot

Verification of solutions

$$\frac{x^2 + y^2}{x^3} = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 26

```
dsolve(diff(y(x),x) = (x^2+3*y(x)^2)/(2*x*y(x)),y(x), singsol=all)
```

$$y(x) = \sqrt{c_1 x - 1} x$$
$$y(x) = -\sqrt{c_1 x - 1} x$$

✓ Solution by Mathematica

Time used: 0.251 (sec). Leaf size: 34

```
DSolve[y'[x] == (x^2+3*y[x]^2)/(2*x*y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -x\sqrt{-1 + c_1 x}$$
$$y(x) \rightarrow x\sqrt{-1 + c_1 x}$$

2.32 problem 33

- 2.32.1 Solving as homogeneousTypeD2 ode 835
- 2.32.2 Solving as first order ode lie symmetry calculated ode 837

Internal problem ID [510]

Internal file name [OUTPUT/510_Sunday_June_05_2022_01_42_44_AM_10559292/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 33.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD2**", "**first_order_ode_lie_symmetry_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$y' - \frac{4y - 3x}{2x - y} = 0$$

2.32.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u'(x)x + u(x) - \frac{4u(x)x - 3x}{2x - u(x)x} = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{u^2 + 2u - 3}{x(u - 2)} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(u) = \frac{u^2+2u-3}{u-2}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{u^2+2u-3}{u-2}} du &= -\frac{1}{x} dx \\ \int \frac{1}{\frac{u^2+2u-3}{u-2}} du &= \int -\frac{1}{x} dx \\ -\frac{\ln(u-1)}{4} + \frac{5\ln(u+3)}{4} &= -\ln(x) + c_2\end{aligned}$$

The above can be written as

$$\begin{aligned}\frac{-\ln(u-1) + 5\ln(u+3)}{4} &= -\ln(x) + c_2 \\ -\ln(u-1) + 5\ln(u+3) &= (4)(-\ln(x) + c_2) \\ &= -4\ln(x) + 4c_2\end{aligned}$$

Raising both side to exponential gives

$$e^{-\ln(u-1)+5\ln(u+3)} = e^{-4\ln(x)+4c_2}$$

Which simplifies to

$$\begin{aligned}\frac{(u+3)^5}{u-1} &= \frac{4c_2}{x^4} \\ &= \frac{c_3}{x^4}\end{aligned}$$

Which simplifies to

$$u(x) = \text{RootOf}\left(_Z^5 + 15_Z^4 + 90_Z^3 + 270_Z^2 + \left(-\frac{c_3 e^{4c_2}}{x^4} + 405\right)_Z + \frac{c_3 e^{4c_2}}{x^4} + 243\right)$$

Therefore the solution y is

$$\begin{aligned}y &= ux \\ &= x \text{RootOf}\left(_Z^5 x^4 + 15x^4_Z^4 + 90_Z^3 x^4 + 270x^4_Z^2 + (-c_3 e^{4c_2} + 405x^4)_Z + c_3 e^{4c_2} + 243x^4\right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = x \text{RootOf}\left(_Z^5 x^4 + 15x^4_Z^4 + 90_Z^3 x^4 + 270x^4_Z^2 + (-c_3 e^{4c_2} + 405x^4)_Z + c_3 e^{4c_2} + 243x^4\right)$$

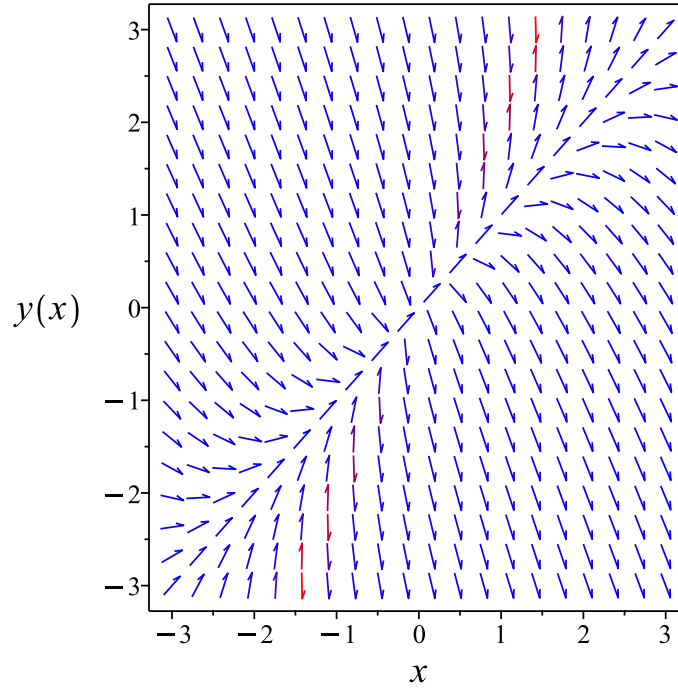


Figure 159: Slope field plot

Verification of solutions

$$y = x \text{ RootOf } (_Z^5 x^4 + 15x^4 _Z^4 + 90 _Z^3 x^4 + 270x^4 _Z^2 + (-c_3 e^{4c_2} + 405x^4) _Z + c_3 e^{4c_2} + 243x^4)$$

Verified OK.

2.32.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{4y - 3x}{-2x + y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 - \frac{(4y - 3x)(b_3 - a_2)}{-2x + y} - \frac{(4y - 3x)^2 a_3}{(-2x + y)^2} \\ - \left(\frac{3}{-2x + y} - \frac{2(4y - 3x)}{(-2x + y)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left(-\frac{4}{-2x + y} + \frac{4y - 3x}{(-2x + y)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{6x^2a_2 - 9x^2a_3 - x^2b_2 - 6x^2b_3 - 6xya_2 + 24xya_3 - 4xyb_2 + 6xyb_3 + 4y^2a_2 - 11y^2a_3 + y^2b_2 - 4y^2b_3 - 5xb_1 - 5ya_1}{(2x - y)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} 6x^2a_2 - 9x^2a_3 - x^2b_2 - 6x^2b_3 - 6xya_2 + 24xya_3 - 4xyb_2 \\ + 6xyb_3 + 4y^2a_2 - 11y^2a_3 + y^2b_2 - 4y^2b_3 - 5xb_1 + 5ya_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} 6a_2v_1^2 - 6a_2v_1v_2 + 4a_2v_2^2 - 9a_3v_1^2 + 24a_3v_1v_2 - 11a_3v_2^2 - b_2v_1^2 \\ - 4b_2v_1v_2 + b_2v_2^2 - 6b_3v_1^2 + 6b_3v_1v_2 - 4b_3v_2^2 + 5a_1v_2 - 5b_1v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$(6a_2 - 9a_3 - b_2 - 6b_3)v_1^2 + (-6a_2 + 24a_3 - 4b_2 + 6b_3)v_1v_2 - 5b_1v_1 + (4a_2 - 11a_3 + b_2 - 4b_3)v_2^2 + 5a_1v_2 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 5a_1 &= 0 \\ -5b_1 &= 0 \\ -6a_2 + 24a_3 - 4b_2 + 6b_3 &= 0 \\ 4a_2 - 11a_3 + b_2 - 4b_3 &= 0 \\ 6a_2 - 9a_3 - b_2 - 6b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= 2a_3 + b_3 \\ a_3 &= a_3 \\ b_1 &= 0 \\ b_2 &= 3a_3 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= y - \left(-\frac{4y - 3x}{-2x + y} \right) (x) \\ &= \frac{3x^2 - 2yx - y^2}{2x - y} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{3x^2 - 2yx - y^2}{2x - y}} dy \end{aligned}$$

Which results in

$$S = \frac{5 \ln(3x + y)}{4} - \frac{\ln(-x + y)}{4}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{4y - 3x}{-2x + y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{-4y + 3x}{(3x + y)(x - y)} \\ S_y &= \frac{2x - y}{(3x + y)(x - y)} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

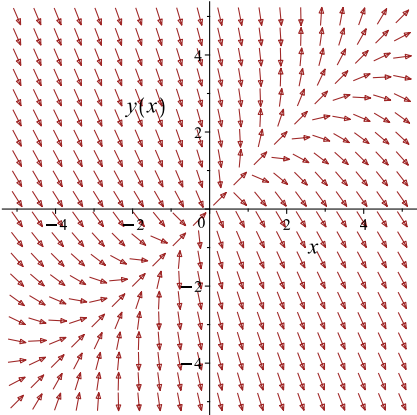
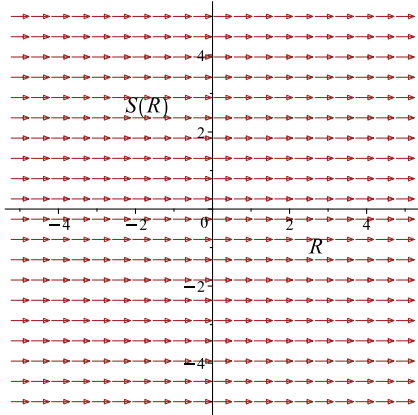
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{5 \ln(3x + y)}{4} - \frac{\ln(-x + y)}{4} = c_1$$

Which simplifies to

$$\frac{5 \ln(3x + y)}{4} - \frac{\ln(-x + y)}{4} = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{4y-3x}{-2x+y}$ 	$R = x$ $S = \frac{5 \ln(3x + y)}{4} - \frac{\ln(-x + y)}{4}$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$\frac{5 \ln(3x + y)}{4} - \frac{\ln(-x + y)}{4} = c_1 \quad (1)$$

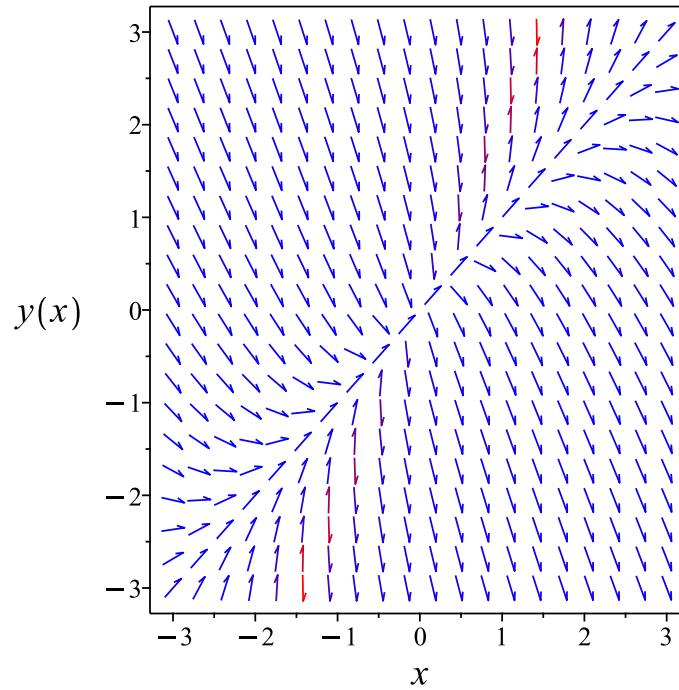


Figure 160: Slope field plot

Verification of solutions

$$\frac{5 \ln(3x + y)}{4} - \frac{\ln(-x + y)}{4} = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 0.375 (sec). Leaf size: 26

```
dsolve(diff(y(x),x) = (4*y(x)-3*x)/(2*x-y(x)),y(x), singsol=all)
```

$$y(x) = x \left(-3 + \text{RootOf} \left(_Z^{20} c_1 x^4 - _Z^4 + 4 \right)^4 \right)$$

✓ Solution by Mathematica

Time used: 3.328 (sec). Leaf size: 336

```
DSolve[y'[x] == (4*y[x]-3*x)/(2*x-y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$\begin{aligned} y(x) &\rightarrow \text{Root} \left[\#1^5 + 15\#1^4 x + 90\#1^3 x^2 + 270\#1^2 x^3 + \#1(405x^4 - e^{4c_1}) + 243x^5 \right. \\ &\quad \left. + e^{4c_1} x \&, 1 \right] \\ y(x) &\rightarrow \text{Root} \left[\#1^5 + 15\#1^4 x + 90\#1^3 x^2 + 270\#1^2 x^3 + \#1(405x^4 - e^{4c_1}) + 243x^5 \right. \\ &\quad \left. + e^{4c_1} x \&, 2 \right] \\ y(x) &\rightarrow \text{Root} \left[\#1^5 + 15\#1^4 x + 90\#1^3 x^2 + 270\#1^2 x^3 + \#1(405x^4 - e^{4c_1}) + 243x^5 \right. \\ &\quad \left. + e^{4c_1} x \&, 3 \right] \\ y(x) &\rightarrow \text{Root} \left[\#1^5 + 15\#1^4 x + 90\#1^3 x^2 + 270\#1^2 x^3 + \#1(405x^4 - e^{4c_1}) + 243x^5 \right. \\ &\quad \left. + e^{4c_1} x \&, 4 \right] \\ y(x) &\rightarrow \text{Root} \left[\#1^5 + 15\#1^4 x + 90\#1^3 x^2 + 270\#1^2 x^3 + \#1(405x^4 - e^{4c_1}) + 243x^5 \right. \\ &\quad \left. + e^{4c_1} x \&, 5 \right] \end{aligned}$$

2.33 problem 34

- 2.33.1 Solving as homogeneousTypeD2 ode 844
- 2.33.2 Solving as first order ode lie symmetry calculated ode 847

Internal problem ID [511]

Internal file name [OUTPUT/511_Sunday_June_05_2022_01_42_45_AM_86261302/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 34.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD2**", "**first_order_ode_lie_symmetry_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$y' + \frac{4x + 3y}{2x + y} = 0$$

2.33.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u'(x)x + u(x) + \frac{4x + 3u(x)x}{2x + u(x)x} = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{u^2 + 5u + 4}{x(u + 2)} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(u) = \frac{u^2+5u+4}{u+2}$. Integrating both sides gives

$$\frac{1}{\frac{u^2+5u+4}{u+2}} du = -\frac{1}{x} dx$$

$$\int \frac{1}{\frac{u^2+5u+4}{u+2}} du = \int -\frac{1}{x} dx$$

$$\frac{2 \ln(u+4)}{3} + \frac{\ln(u+1)}{3} = -\ln(x) + c_2$$

The above can be written as

$$\frac{2 \ln(u+4) + \ln(u+1)}{3} = -\ln(x) + c_2$$

$$2 \ln(u+4) + \ln(u+1) = (3)(-\ln(x) + c_2)$$

$$= -3 \ln(x) + 3c_2$$

Raising both side to exponential gives

$$e^{2 \ln(u+4) + \ln(u+1)} = e^{-3 \ln(x) + 3c_2}$$

Which simplifies to

$$(u+4)^2 (u+1) = \frac{3c_2}{x^3}$$

$$= \frac{c_3}{x^3}$$

Which simplifies to

$$(u(x)+4)^2 (u(x)+1) = \frac{c_3 e^{3c_2}}{x^3}$$

The solution is

$$(u(x)+4)^2 (u(x)+1) = \frac{c_3 e^{3c_2}}{x^3}$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\left(\frac{y}{x} + 4\right)^2 \left(1 + \frac{y}{x}\right) = \frac{c_3 e^{3c_2}}{x^3}$$

$$\frac{(4x+y)^2 (x+y)}{x^3} = \frac{c_3 e^{3c_2}}{x^3}$$

Which simplifies to

$$(4x + y)^2 (x + y) = c_3 e^{3c_2}$$

Summary

The solution(s) found are the following

$$(4x + y)^2 (x + y) = c_3 e^{3c_2} \tag{1}$$

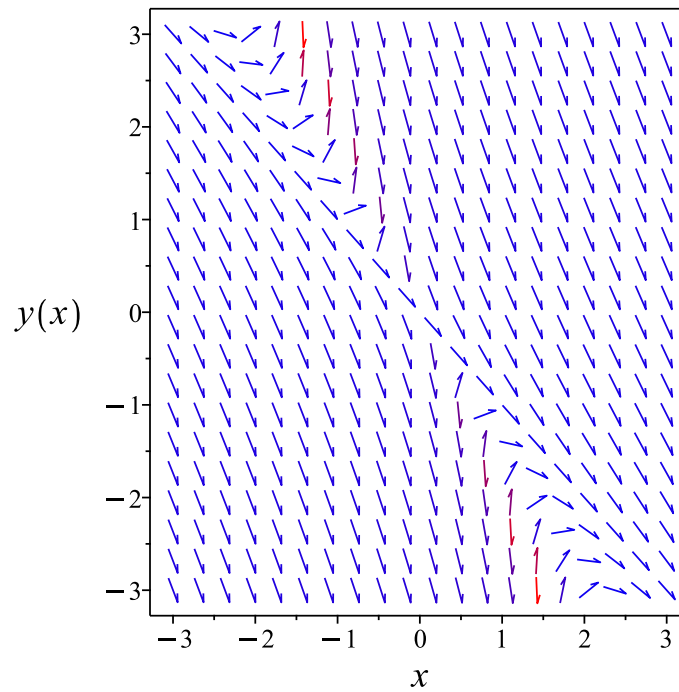


Figure 161: Slope field plot

Verification of solutions

$$(4x + y)^2 (x + y) = c_3 e^{3c_2}$$

Verified OK.

2.33.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{4x + 3y}{2x + y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 - \frac{(4x + 3y)(b_3 - a_2)}{2x + y} - \frac{(4x + 3y)^2 a_3}{(2x + y)^2}$$

$$- \left(-\frac{4}{2x + y} + \frac{8x + 6y}{(2x + y)^2} \right) (xa_2 + ya_3 + a_1) \quad (\text{5E})$$

$$- \left(-\frac{3}{2x + y} + \frac{4x + 3y}{(2x + y)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{8x^2 a_2 - 16x^2 a_3 + 6x^2 b_2 - 8x^2 b_3 + 8xy a_2 - 24xy a_3 + 4xy b_2 - 8xy b_3 + 3y^2 a_2 - 11y^2 a_3 + y^2 b_2 - 3y^2 b_3 + 2xb_1 - 2ya_1}{(2x + y)^2}$$

$$= 0$$

Setting the numerator to zero gives

$$8x^2 a_2 - 16x^2 a_3 + 6x^2 b_2 - 8x^2 b_3 + 8xy a_2 - 24xy a_3 + 4xy b_2$$

$$- 8xy b_3 + 3y^2 a_2 - 11y^2 a_3 + y^2 b_2 - 3y^2 b_3 + 2xb_1 - 2ya_1 = 0 \quad (\text{6E})$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} 8a_2v_1^2 + 8a_2v_1v_2 + 3a_2v_2^2 - 16a_3v_1^2 - 24a_3v_1v_2 - 11a_3v_2^2 + 6b_2v_1^2 \\ + 4b_2v_1v_2 + b_2v_2^2 - 8b_3v_1^2 - 8b_3v_1v_2 - 3b_3v_2^2 - 2a_1v_2 + 2b_1v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} (8a_2 - 16a_3 + 6b_2 - 8b_3)v_1^2 + (8a_2 - 24a_3 + 4b_2 - 8b_3)v_1v_2 \\ + 2b_1v_1 + (3a_2 - 11a_3 + b_2 - 3b_3)v_2^2 - 2a_1v_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$-2a_1 = 0$$

$$2b_1 = 0$$

$$3a_2 - 11a_3 + b_2 - 3b_3 = 0$$

$$8a_2 - 24a_3 + 4b_2 - 8b_3 = 0$$

$$8a_2 - 16a_3 + 6b_2 - 8b_3 = 0$$

Solving the above equations for the unknowns gives

$$a_1 = 0$$

$$a_2 = 5a_3 + b_3$$

$$a_3 = a_3$$

$$b_1 = 0$$

$$b_2 = -4a_3$$

$$b_3 = b_3$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= x \\ \eta &= y\end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left(-\frac{4x + 3y}{2x + y} \right) (x) \\ &= \frac{4x^2 + 5yx + y^2}{2x + y} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{4x^2 + 5yx + y^2}{2x + y}} dy\end{aligned}$$

Which results in

$$S = \frac{2 \ln(4x + y)}{3} + \frac{\ln(x + y)}{3}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y) S_y}{R_x + \omega(x, y) R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{4x + 3y}{2x + y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{4x + 3y}{(x + y)(4x + y)} \\ S_y &= \frac{2x + y}{(x + y)(4x + y)} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

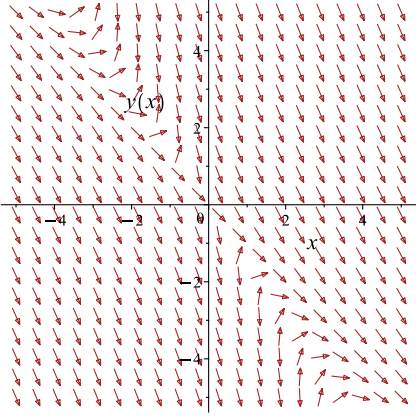
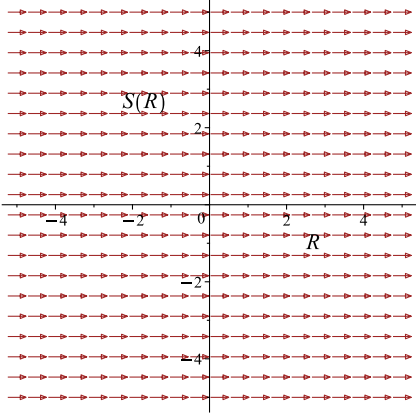
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{2 \ln(4x + y)}{3} + \frac{\ln(x + y)}{3} = c_1$$

Which simplifies to

$$\frac{2 \ln(4x + y)}{3} + \frac{\ln(x + y)}{3} = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{4x+3y}{2x+y}$ 	$R = x$ $S = \frac{2 \ln(4x + y)}{3} + \frac{\ln(x)}{\xi}$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$\frac{2 \ln(4x + y)}{3} + \frac{\ln(x + y)}{3} = c_1 \tag{1}$$

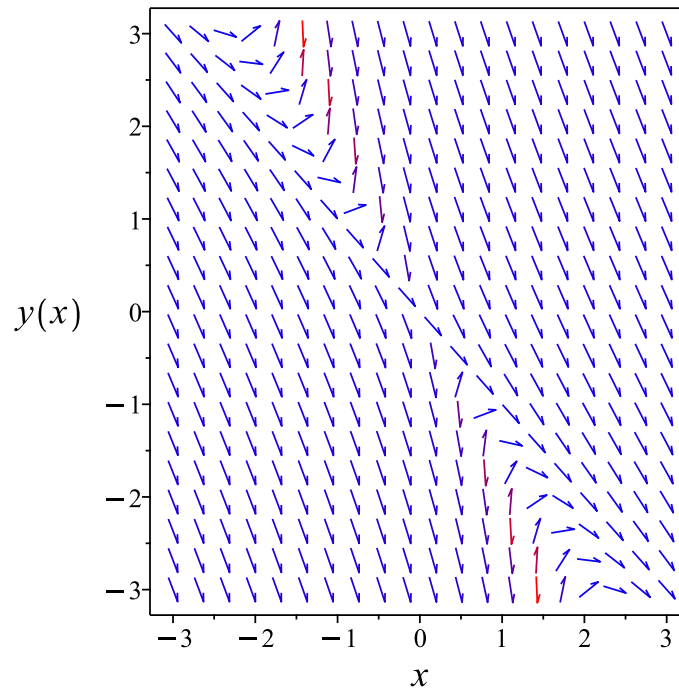


Figure 162: Slope field plot

Verification of solutions

$$\frac{2 \ln(4x + y)}{3} + \frac{\ln(x + y)}{3} = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.125 (sec). Leaf size: 1228

`dsolve(diff(y(x),x) = - (4*x+3*y(x))/(2*x+y(x)),y(x), singsol=all)`

$$y(x) = \frac{\left(-4c_1x^3 + \left(4c_1x^3 + 4\sqrt{x^6c_1^2(4c_1x^3+1)}\right)^{\frac{2}{3}}\right)^2}{4\left(4c_1x^3 + 4\sqrt{x^6c_1^2(4c_1x^3+1)}\right)^{\frac{2}{3}}c_1} - x^3$$

$$y(x) = \frac{-3x^3\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{2}{3}}c_1 + \left(c_1x^3 + \sqrt{4c_1^3x^9 + x^6c_1^2}\right)\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{1}{3}} + 4x^6c_1^2}{\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{2}{3}}c_1x^2}$$

$$y(x) = \frac{-3x^3\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{2}{3}}c_1 + \left(c_1x^3 + \sqrt{4c_1^3x^9 + x^6c_1^2}\right)\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{1}{3}} + 4x^6c_1^2}{\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{2}{3}}c_1x^2}$$

$$y(x) = -\frac{\left(4\sqrt{3}c_1x^3 + \sqrt{3}\left(4c_1x^3 + 4\sqrt{x^6c_1^2(4c_1x^3+1)}\right)^{\frac{2}{3}} + 4ic_1x^3 - i\left(4c_1x^3 + 4\sqrt{x^6c_1^2(4c_1x^3+1)}\right)^{\frac{2}{3}}\right)^2}{16\left(4c_1x^3 + 4\sqrt{x^6c_1^2(4c_1x^3+1)}\right)^{\frac{2}{3}}c_1} + x^3$$

$$y(x) = -\frac{\left(4\sqrt{3}c_1x^3 + \sqrt{3}\left(4c_1x^3 + 4\sqrt{x^6c_1^2(4c_1x^3+1)}\right)^{\frac{2}{3}} - 4ic_1x^3 + i\left(4c_1x^3 + 4\sqrt{x^6c_1^2(4c_1x^3+1)}\right)^{\frac{2}{3}}\right)^2}{16\left(4c_1x^3 + 4\sqrt{x^6c_1^2(4c_1x^3+1)}\right)^{\frac{2}{3}}c_1} + x^3$$

$$y(x) = \frac{2\left(\frac{3x^3\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{2}{3}}c_1}{2} - \frac{\left(c_1x^3 + \sqrt{4c_1^3x^9 + x^6c_1^2}\right)(i\sqrt{3}-1)\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{1}{3}}}{4} + x^6(1+i\sqrt{3})c_1^2\right)}{\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{2}{3}}c_1x^2}$$

$$y(x) = \frac{-3x^3\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{2}{3}}c_1 - \frac{\left(c_1x^3 + \sqrt{4c_1^3x^9 + x^6c_1^2}\right)(1+i\sqrt{3})\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{1}{3}}}{2} + 2x^6(i\sqrt{3}-1)c_1^2}{\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{2}{3}}c_1x^2}$$

$$y(x) = \frac{2\left(\frac{3x^3\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{2}{3}}c_1}{2} - \frac{\left(c_1x^3 + \sqrt{4c_1^3x^9 + x^6c_1^2}\right)(i\sqrt{3}-1)\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{1}{3}}}{4} + x^6(1+i\sqrt{3})c_1^2\right)}{853\left(4c_1x^3 + 4\sqrt{4c_1^3x^9 + x^6c_1^2}\right)^{\frac{2}{3}}c_1x^2}$$

✓ Solution by Mathematica

Time used: 20.375 (sec). Leaf size: 484

`DSolve[y'[x] == - (4*x+3*y[x])/(2*x+y[x]), y[x], x, IncludeSingularSolutions -> True]`

$$y(x) \rightarrow \frac{\sqrt[3]{2x^3 + \sqrt{4e^{3c_1}x^3 + e^{6c_1}} + e^{3c_1}}}{\sqrt[3]{2}} + \frac{\sqrt[3]{2}x^2}{\sqrt[3]{2x^3 + \sqrt{4e^{3c_1}x^3 + e^{6c_1}} + e^{3c_1}}} - 3x$$

$$y(x) \rightarrow \frac{i(\sqrt{3} + i) \sqrt[3]{2x^3 + \sqrt{4e^{3c_1}x^3 + e^{6c_1}} + e^{3c_1}}}{2\sqrt[3]{2}} - \frac{(1 + i\sqrt{3})x^2}{2^{2/3} \sqrt[3]{2x^3 + \sqrt{4e^{3c_1}x^3 + e^{6c_1}} + e^{3c_1}}} - 3x$$

$$y(x) \rightarrow -\frac{(1 + i\sqrt{3}) \sqrt[3]{2x^3 + \sqrt{4e^{3c_1}x^3 + e^{6c_1}} + e^{3c_1}}}{2\sqrt[3]{2}} + \frac{i(\sqrt{3} + i)x^2}{2^{2/3} \sqrt[3]{2x^3 + \sqrt{4e^{3c_1}x^3 + e^{6c_1}} + e^{3c_1}}} - 3x$$

$$y(x) \rightarrow \sqrt[3]{x^3} + \frac{(x^3)^{2/3}}{x} - 3x$$

$$y(x) \rightarrow \frac{1}{2} \left(i(\sqrt{3} + i) \sqrt[3]{x^3} + \frac{(-1 - i\sqrt{3})(x^3)^{2/3}}{x} - 6x \right)$$

$$y(x) \rightarrow \frac{1}{2} \left((-1 - i\sqrt{3}) \sqrt[3]{x^3} + \frac{i(\sqrt{3} + i)(x^3)^{2/3}}{x} - 6x \right)$$

2.34 problem 35

- 2.34.1 Solving as homogeneousTypeD2 ode 855
- 2.34.2 Solving as first order ode lie symmetry calculated ode 857

Internal problem ID [512]

Internal file name [OUTPUT/512_Sunday_June_05_2022_01_42_48_AM_92797294/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 35.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD2**", "**first_order_ode_lie_symmetry_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$y' - \frac{x + 3y}{x - y} = 0$$

2.34.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u'(x)x + u(x) - \frac{x + 3u(x)x}{x - u(x)x} = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{(u + 1)^2}{x(u - 1)} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(u) = \frac{(u+1)^2}{u-1}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{(u+1)^2}{u-1}} du &= -\frac{1}{x} dx \\ \int \frac{1}{\frac{(u+1)^2}{u-1}} du &= \int -\frac{1}{x} dx \\ \ln(u+1) + \frac{2}{u+1} &= -\ln(x) + c_2\end{aligned}$$

The solution is

$$\ln(u(x)+1) + \frac{2}{u(x)+1} + \ln(x) - c_2 = 0$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\begin{aligned}\ln\left(1 + \frac{y}{x}\right) + \frac{2}{1 + \frac{y}{x}} + \ln(x) - c_2 &= 0 \\ \ln\left(\frac{x+y}{x}\right) + \frac{2x}{x+y} + \ln(x) - c_2 &= 0\end{aligned}$$

Summary

The solution(s) found are the following

$$\ln\left(\frac{x+y}{x}\right) + \frac{2x}{x+y} + \ln(x) - c_2 = 0 \quad (1)$$

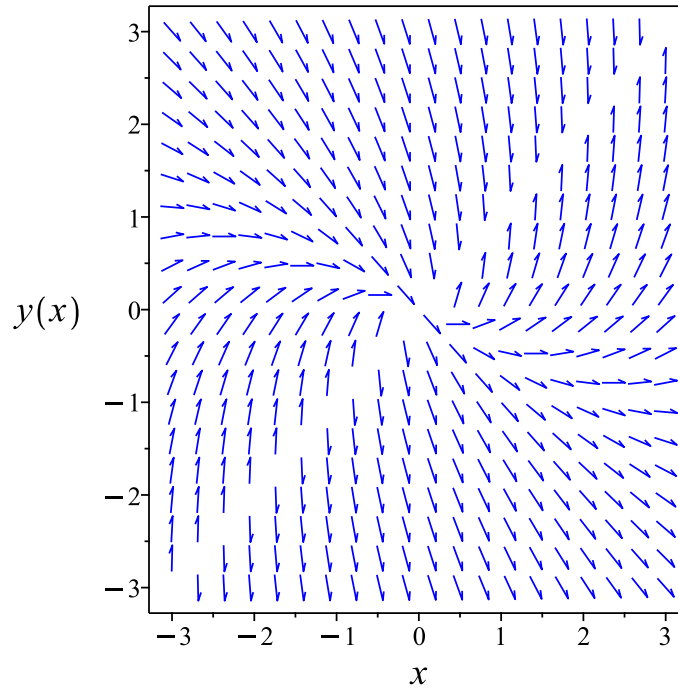


Figure 163: Slope field plot

Verification of solutions

$$\ln\left(\frac{x+y}{x}\right) + \frac{2x}{x+y} + \ln(x) - c_2 = 0$$

Verified OK.

2.34.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{3y+x}{-x+y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 - \frac{(3y+x)(b_3-a_2)}{-x+y} - \frac{(3y+x)^2 a_3}{(-x+y)^2} \\ - \left(-\frac{1}{-x+y} - \frac{3y+x}{(-x+y)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left(-\frac{3}{-x+y} + \frac{3y+x}{(-x+y)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{x^2 a_2 + x^2 a_3 + 3x^2 b_2 - x^2 b_3 - 2xy a_2 + 6xy a_3 + 2xy b_2 + 2xy b_3 - 3y^2 a_2 + 5y^2 a_3 - y^2 b_2 + 3y^2 b_3 + 4xb_1 - 4yb_1}{(x-y)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} -x^2 a_2 - x^2 a_3 - 3x^2 b_2 + x^2 b_3 + 2xy a_2 - 6xy a_3 - 2xy b_2 \\ - 2xy b_3 + 3y^2 a_2 - 5y^2 a_3 + y^2 b_2 - 3y^2 b_3 - 4xb_1 + 4yb_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} -a_2 v_1^2 + 2a_2 v_1 v_2 + 3a_2 v_2^2 - a_3 v_1^2 - 6a_3 v_1 v_2 - 5a_3 v_2^2 - 3b_2 v_1^2 \\ - 2b_2 v_1 v_2 + b_2 v_2^2 + b_3 v_1^2 - 2b_3 v_1 v_2 - 3b_3 v_2^2 + 4a_1 v_2 - 4b_1 v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} &(-a_2 - a_3 - 3b_2 + b_3)v_1^2 + (2a_2 - 6a_3 - 2b_2 - 2b_3)v_1v_2 \\ &- 4b_1v_1 + (3a_2 - 5a_3 + b_2 - 3b_3)v_2^2 + 4a_1v_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 4a_1 &= 0 \\ -4b_1 &= 0 \\ -a_2 - a_3 - 3b_2 + b_3 &= 0 \\ 2a_2 - 6a_3 - 2b_2 - 2b_3 &= 0 \\ 3a_2 - 5a_3 + b_2 - 3b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= -2b_2 + b_3 \\ a_3 &= -b_2 \\ b_1 &= 0 \\ b_2 &= b_2 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y)\xi \\ &= y - \left(-\frac{3y+x}{-x+y}\right)(x) \\ &= \frac{-x^2 - 2yx - y^2}{x-y} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-x^2 - 2yx - y^2}{x-y}} dy \end{aligned}$$

Which results in

$$S = \frac{2x}{x+y} + \ln(x+y)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{3y+x}{-x+y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{3y+x}{(x+y)^2} \\ S_y &= \frac{-x+y}{(x+y)^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{(x+y) \ln(x+y) + 2x}{x+y} = c_1$$

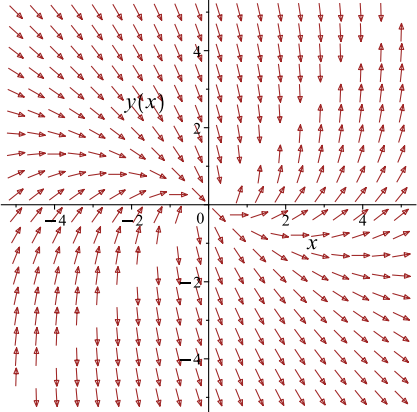
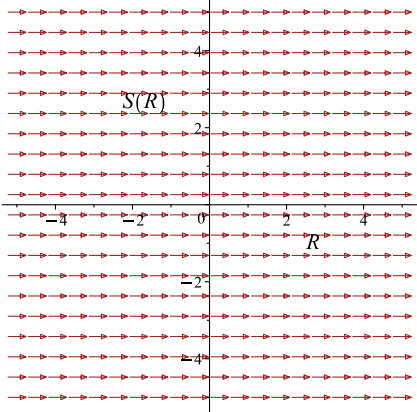
Which simplifies to

$$\frac{(x+y) \ln(x+y) + 2x}{x+y} = c_1$$

Which gives

$$y = e^{\text{LambertW}(-2x e^{-c_1}) + c_1} - x$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{3y+x}{-x+y}$ 	$R = x$ $S = \frac{(x+y) \ln(x+y) + x^2 - y^2}{x+y}$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$y = e^{\text{LambertW}(-2x e^{-c_1}) + c_1} - x \tag{1}$$

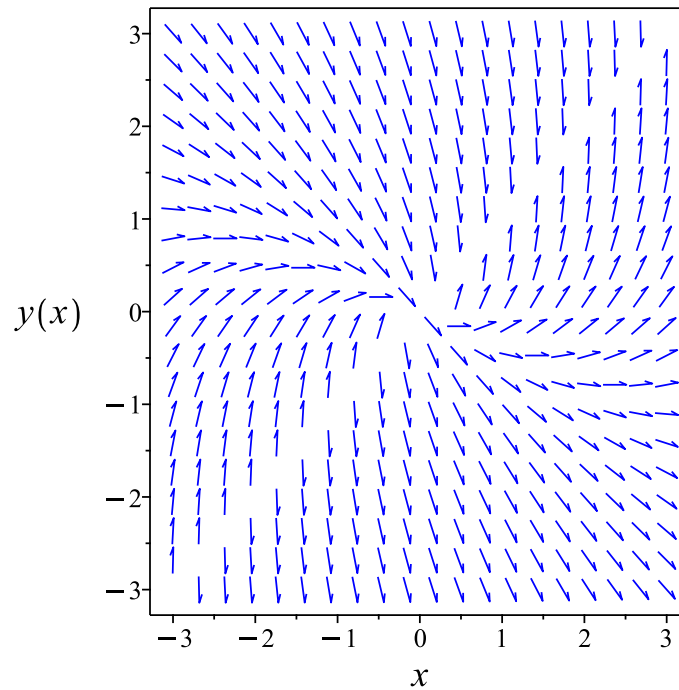


Figure 164: Slope field plot

Verification of solutions

$$y = e^{\text{LambertW}(-2x e^{-c_1}) + c_1} - x$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```


✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 21

```
dsolve(diff(y(x),x) = (x+3*y(x))/(x-y(x)),y(x), singsol=all)
```

$$y(x) = -\frac{x(\text{LambertW}(-2c_1x) + 2)}{\text{LambertW}(-2c_1x)}$$

✓ Solution by Mathematica

Time used: 0.113 (sec). Leaf size: 33

```
DSolve[y'[x] == (x+3*y[x])/(x-y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve} \left[\frac{2}{\frac{y(x)}{x} + 1} + \log \left(\frac{y(x)}{x} + 1 \right) = -\log(x) + c_1, y(x) \right]$$

2.35 problem 36

2.35.1 Solving as homogeneousTypeD2 ode	865
2.35.2 Solving as first order ode lie symmetry calculated ode	867
2.35.3 Solving as riccati ode	873

Internal problem ID [513]

Internal file name [OUTPUT/513_Sunday_June_05_2022_01_42_49_AM_6394563/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 36.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**riccati**", "**homogeneousTypeD2**", "**first_order_ode_lie_symmetry_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _Riccati]
```

$$3yx + y^2 - y'x^2 = -x^2$$

2.35.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$3u(x)x^2 + u(x)^2x^2 - (u'(x)x + u(x))x^2 = -x^2$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{u^2 + 2u + 1}{x}\end{aligned}$$

Where $f(x) = \frac{1}{x}$ and $g(u) = u^2 + 2u + 1$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u^2 + 2u + 1} du &= \frac{1}{x} dx \\ \int \frac{1}{u^2 + 2u + 1} du &= \int \frac{1}{x} dx \\ -\frac{1}{u + 1} &= \ln(x) + c_2\end{aligned}$$

The solution is

$$-\frac{1}{u(x) + 1} - \ln(x) - c_2 = 0$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\begin{aligned}-\frac{1}{1 + \frac{y}{x}} - \ln(x) - c_2 &= 0 \\ \frac{(-c_2 - \ln(x))y - x(c_2 + \ln(x) + 1)}{x + y} &= 0\end{aligned}$$

Which simplifies to

$$-\frac{\ln(x)y + c_2y + x \ln(x) + c_2x + x}{x + y} = 0$$

Summary

The solution(s) found are the following

$$-\frac{\ln(x)y + c_2y + x \ln(x) + c_2x + x}{x + y} = 0 \tag{1}$$

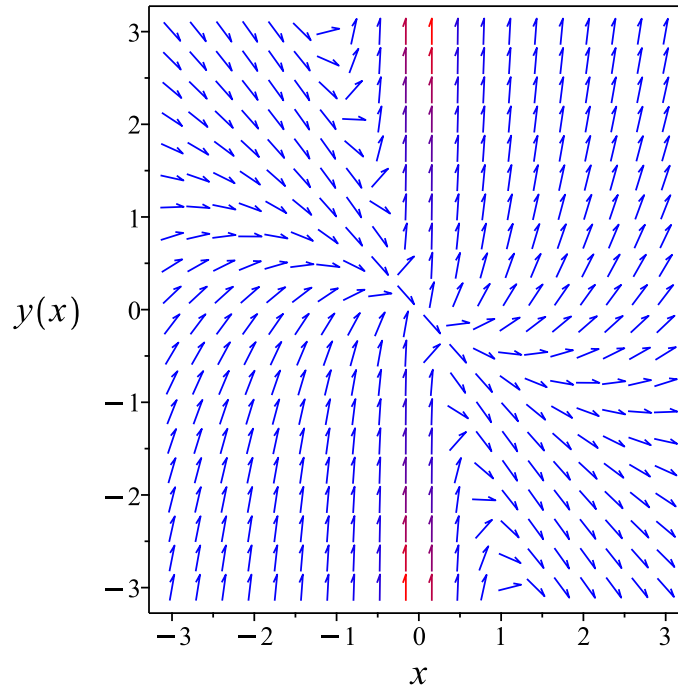


Figure 165: Slope field plot

Verification of solutions

$$\frac{\ln(x)y + c_2y + x \ln(x) + c_2x + x}{x + y} = 0$$

Verified OK.

2.35.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{x^2 + 3yx + y^2}{x^2}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 + \frac{(x^2 + 3yx + y^2)(b_3 - a_2)}{x^2} - \frac{(x^2 + 3yx + y^2)^2 a_3}{x^4} \\ - \left(\frac{2x + 3y}{x^2} - \frac{2(x^2 + 3yx + y^2)}{x^3} \right) (xa_2 + ya_3 + a_1) \\ - \frac{(3x + 2y)(xb_2 + yb_3 + b_1)}{x^2} = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{x^4 a_2 + x^4 a_3 + 2b_2 x^4 - x^4 b_3 + 6x^3 y a_3 + 2x^3 y b_2 - x^2 y^2 a_2 + 8x^2 y^2 a_3 + x^2 y^2 b_3 + 4x y^3 a_3 + y^4 a_3 + 3x^3 b_1 - x^2 y^2 b_3 - 4x y^3 a_3 - y^4 a_3 - 3x^3 b_1 + 3x^2 y a_1 - 2x^2 y b_1 + 2x y^2 a_1}{x^4} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} -x^4 a_2 - x^4 a_3 - 2b_2 x^4 + x^4 b_3 - 6x^3 y a_3 - 2x^3 y b_2 + x^2 y^2 a_2 - 8x^2 y^2 a_3 \\ - x^2 y^2 b_3 - 4x y^3 a_3 - y^4 a_3 - 3x^3 b_1 + 3x^2 y a_1 - 2x^2 y b_1 + 2x y^2 a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} -a_2 v_1^4 + a_2 v_1^2 v_2^2 - a_3 v_1^4 - 6a_3 v_1^3 v_2 - 8a_3 v_1^2 v_2^2 - 4a_3 v_1 v_2^3 - a_3 v_2^4 - 2b_2 v_1^4 \\ - 2b_2 v_1^3 v_2 + b_3 v_1^4 - b_3 v_1^2 v_2^2 + 3a_1 v_1^2 v_2 + 2a_1 v_1 v_2^2 - 3b_1 v_1^3 - 2b_1 v_1^2 v_2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} &(-a_2 - a_3 - 2b_2 + b_3) v_1^4 + (-6a_3 - 2b_2) v_1^3 v_2 - 3b_1 v_1^3 + (a_2 - 8a_3 - b_3) v_1^2 v_2^2 \quad (8E) \\ &+ (3a_1 - 2b_1) v_1^2 v_2 - 4a_3 v_1 v_2^3 + 2a_1 v_1 v_2^2 - a_3 v_2^4 = 0 \end{aligned}$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 2a_1 &= 0 \\ -4a_3 &= 0 \\ -a_3 &= 0 \\ -3b_1 &= 0 \\ 3a_1 - 2b_1 &= 0 \\ -6a_3 - 2b_2 &= 0 \\ a_2 - 8a_3 - b_3 &= 0 \\ -a_2 - a_3 - 2b_2 + b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left(\frac{x^2 + 3yx + y^2}{x^2} \right) (x) \\ &= \frac{-x^2 - 2yx - y^2}{x} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-x^2 - 2yx - y^2}{x}} dy\end{aligned}$$

Which results in

$$S = \frac{x}{x + y}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y) S_y}{R_x + \omega(x, y) R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x^2 + 3yx + y^2}{x^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= \frac{y}{(x+y)^2} \\S_y &= -\frac{x}{(x+y)^2}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{x} \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\ln(R) + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{x}{x+y} = -\ln(x) + c_1$$

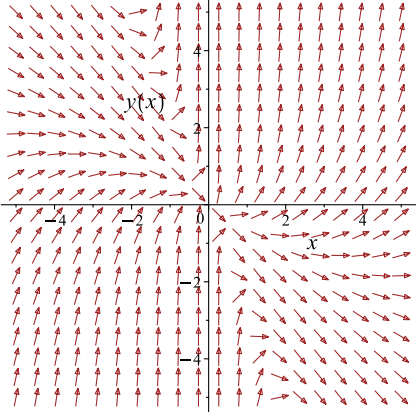
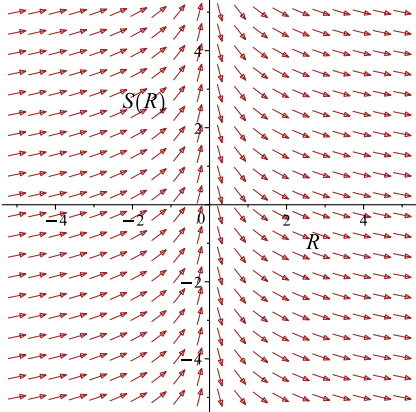
Which simplifies to

$$\frac{x}{x+y} = -\ln(x) + c_1$$

Which gives

$$y = -\frac{x(\ln(x) - c_1 + 1)}{\ln(x) - c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{x^2 + 3yx + y^2}{x^2}$ 	$R = x$ $S = \frac{x}{x + y}$	$\frac{dS}{dR} = -\frac{1}{R}$ 

Summary

The solution(s) found are the following

$$y = -\frac{x(\ln(x) - c_1 + 1)}{\ln(x) - c_1} \tag{1}$$

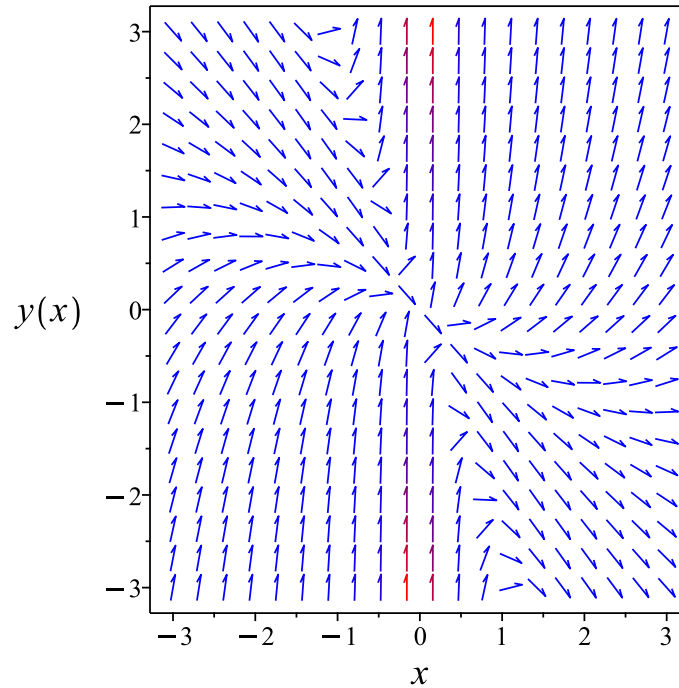


Figure 166: Slope field plot

Verification of solutions

$$y = -\frac{x(\ln(x) - c_1 + 1)}{\ln(x) - c_1}$$

Verified OK.

2.35.3 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{x^2 + 3yx + y^2}{x^2} \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = 1 + \frac{3y}{x} + \frac{y^2}{x^2}$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that $f_0(x) = 1$, $f_1(x) = \frac{3}{x}$ and $f_2(x) = \frac{1}{x^2}$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{\frac{u}{x^2}} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \tag{2}$$

But

$$\begin{aligned} f_2' &= -\frac{2}{x^3} \\ f_1 f_2 &= \frac{3}{x^3} \\ f_2^2 f_0 &= \frac{1}{x^4} \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$\frac{u''(x)}{x^2} - \frac{u'(x)}{x^3} + \frac{u(x)}{x^4} = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = x(c_2 \ln(x) + c_1)$$

The above shows that

$$u'(x) = c_2 \ln(x) + c_1 + c_2$$

Using the above in (1) gives the solution

$$y = -\frac{(c_2 \ln(x) + c_1 + c_2) x}{c_2 \ln(x) + c_1}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = -\frac{(\ln(x) + c_3 + 1) x}{\ln(x) + c_3}$$

Summary

The solution(s) found are the following

$$y = -\frac{(\ln(x) + c_3 + 1)x}{\ln(x) + c_3} \quad (1)$$

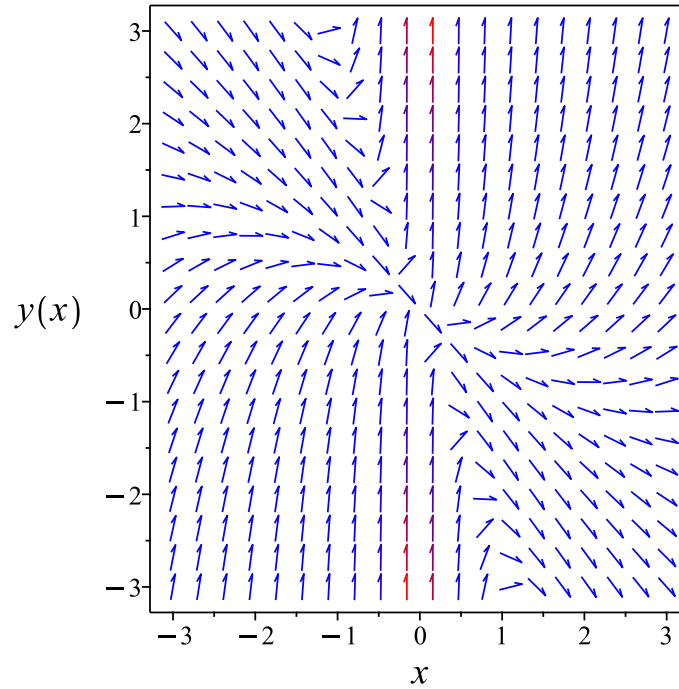


Figure 167: Slope field plot

Verification of solutions

$$y = -\frac{(\ln(x) + c_3 + 1)x}{\ln(x) + c_3}$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 18

```
dsolve((x^2+3*x*y(x)+y(x)^2)-x^2* diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = -\frac{x(\ln(x) + c_1 + 1)}{\ln(x) + c_1}$$

✓ Solution by Mathematica

Time used: 0.141 (sec). Leaf size: 28

```
DSolve[(x^2+3*x*y[x]+y[x]^2)-x^2* y'[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{x(\log(x) + 1 + c_1)}{\log(x) + c_1}$$
$$y(x) \rightarrow -x$$

2.36 problem 37

2.36.1 Solving as homogeneousTypeD2 ode	877
2.36.2 Solving as first order ode lie symmetry lookup ode	879
2.36.3 Solving as bernoulli ode	883
2.36.4 Solving as exact ode	886

Internal problem ID [514]

Internal file name [OUTPUT/514_Sunday_June_05_2022_01_42_50_AM_58581911/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 37.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**bernoulli**", "**homogeneousTypeD2**", "**exactWithIntegrationFactor**", "**first_order_ode_lie_symmetry_lookup**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _Bernoulli]
```

$$y' - \frac{x^2 - 3y^2}{2yx} = 0$$

2.36.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u'(x)x + u(x) - \frac{x^2 - 3u(x)^2 x^2}{2u(x)x^2} = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{5u^2 - 1}{2xu} \end{aligned}$$

Where $f(x) = -\frac{1}{2x}$ and $g(u) = \frac{5u^2-1}{u}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{5u^2-1}{u}} du &= -\frac{1}{2x} dx \\ \int \frac{1}{\frac{5u^2-1}{u}} du &= \int -\frac{1}{2x} dx \\ \frac{\ln(5u^2-1)}{10} &= -\frac{\ln(x)}{2} + c_2\end{aligned}$$

Raising both side to exponential gives

$$(5u^2 - 1)^{\frac{1}{10}} = e^{-\frac{\ln(x)}{2} + c_2}$$

Which simplifies to

$$(5u^2 - 1)^{\frac{1}{10}} = \frac{c_3}{\sqrt{x}}$$

Which simplifies to

$$(5u(x)^2 - 1)^{\frac{1}{10}} = \frac{c_3 e^{c_2}}{\sqrt{x}}$$

The solution is

$$(5u(x)^2 - 1)^{\frac{1}{10}} = \frac{c_3 e^{c_2}}{\sqrt{x}}$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\begin{aligned}\left(\frac{5y^2}{x^2} - 1\right)^{\frac{1}{10}} &= \frac{c_3 e^{c_2}}{\sqrt{x}} \\ \left(\frac{5y^2 - x^2}{x^2}\right)^{\frac{1}{10}} &= \frac{c_3 e^{c_2}}{\sqrt{x}}\end{aligned}$$

Which simplifies to

$$\left(-\frac{5y^2 + x^2}{x^2}\right)^{\frac{1}{10}} = \frac{c_3 e^{c_2}}{\sqrt{x}}$$

Summary

The solution(s) found are the following

$$\left(-\frac{5y^2 + x^2}{x^2}\right)^{\frac{1}{10}} = \frac{c_3 e^{c_2}}{\sqrt{x}} \quad (1)$$

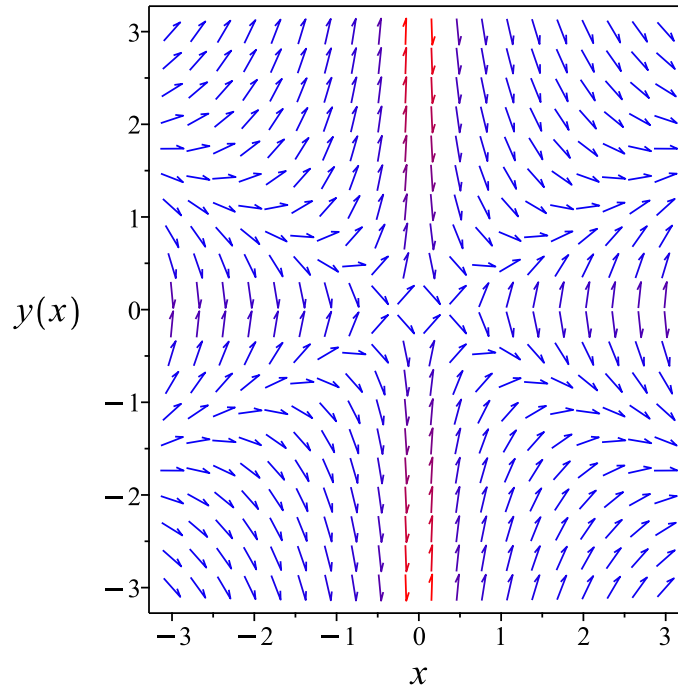


Figure 168: Slope field plot

Verification of solutions

$$\left(-\frac{5y^2 + x^2}{x^2} \right)^{\frac{1}{10}} = \frac{c_3 e^{c_2}}{\sqrt{x}}$$

Verified OK.

2.36.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{-x^2 + 3y^2}{2yx}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **Bernoulli**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 178: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{x^3y}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{x^3 y}} dy \end{aligned}$$

Which results in

$$S = \frac{x^3 y^2}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{-x^2 + 3y^2}{2yx}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{3x^2 y^2}{2} \\ S_y &= y x^3 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{x^4}{2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{R^4}{2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^5}{10} + c_1 \quad (4)$$

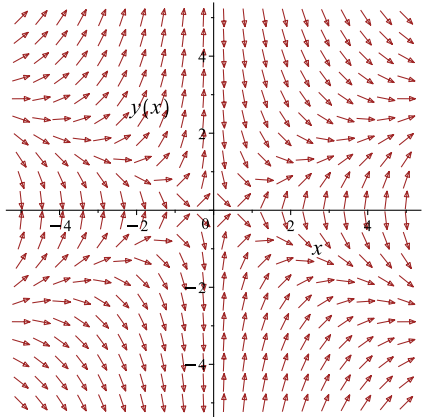
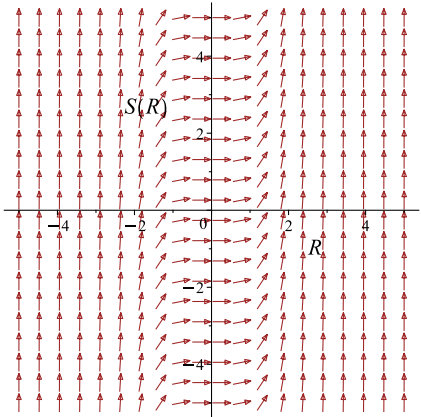
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{x^3 y^2}{2} = \frac{x^5}{10} + c_1$$

Which simplifies to

$$\frac{x^3 y^2}{2} = \frac{x^5}{10} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{-x^2 + 3y^2}{2yx}$ 	$R = x$ $S = \frac{x^3 y^2}{2}$	$\frac{dS}{dR} = \frac{R^4}{2}$ 

Summary

The solution(s) found are the following

$$\frac{x^3 y^2}{2} = \frac{x^5}{10} + c_1 \quad (1)$$

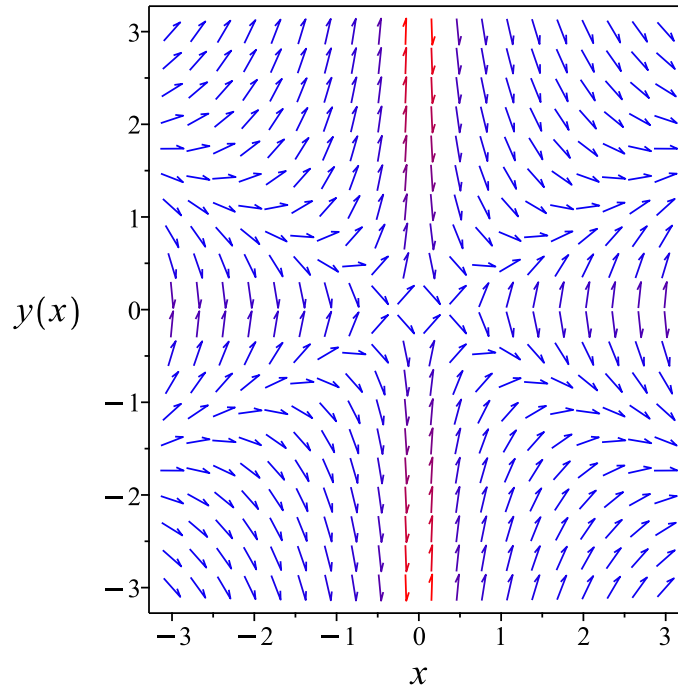


Figure 169: Slope field plot

Verification of solutions

$$\frac{x^3 y^2}{2} = \frac{x^5}{10} + c_1$$

Verified OK.

2.36.3 Solving as bernoulli ode

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= -\frac{-x^2 + 3y^2}{2yx} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = -\frac{3}{2x}y + \frac{x}{2} \frac{1}{y} \quad (1)$$

The standard Bernoulli ODE has the form

$$y' = f_0(x)y + f_1(x)y^n \quad (2)$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(x)y^{1-n} + f_1(x) \quad (3)$$

The next step is use the substitution $w = y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(x) &= -\frac{3}{2x} \\ f_1(x) &= \frac{x}{2} \\ n &= -1 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = \frac{1}{y}$ gives

$$y'y = -\frac{3y^2}{2x} + \frac{x}{2} \quad (4)$$

Let

$$\begin{aligned} w &= y^{1-n} \\ &= y^2 \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t x gives

$$w' = 2yy' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} \frac{w'(x)}{2} &= -\frac{3w(x)}{2x} + \frac{x}{2} \\ w' &= -\frac{3w}{x} + x \end{aligned} \quad (7)$$

The above now is a linear ODE in $w(x)$ which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(x) + p(x)w(x) = q(x)$$

Where here

$$\begin{aligned} p(x) &= \frac{3}{x} \\ q(x) &= x \end{aligned}$$

Hence the ode is

$$w'(x) + \frac{3w(x)}{x} = x$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{3}{x} dx} \\ &= x^3\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu w) &= (\mu)(x) \\ \frac{d}{dx}(x^3 w) &= (x^3)(x) \\ d(x^3 w) &= x^4 dx\end{aligned}$$

Integrating gives

$$\begin{aligned}x^3 w &= \int x^4 dx \\ x^3 w &= \frac{x^5}{5} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = x^3$ results in

$$w(x) = \frac{x^2}{5} + \frac{c_1}{x^3}$$

Replacing w in the above by y^2 using equation (5) gives the final solution.

$$y^2 = \frac{x^2}{5} + \frac{c_1}{x^3}$$

Solving for y gives

$$\begin{aligned}y(x) &= \frac{\sqrt{5} \sqrt{x(x^5 + 5c_1)}}{5x^2} \\ y(x) &= -\frac{\sqrt{5} \sqrt{x(x^5 + 5c_1)}}{5x^2}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{5} \sqrt{x(x^5 + 5c_1)}}{5x^2} \quad (1)$$

$$y = -\frac{\sqrt{5} \sqrt{x(x^5 + 5c_1)}}{5x^2} \quad (2)$$

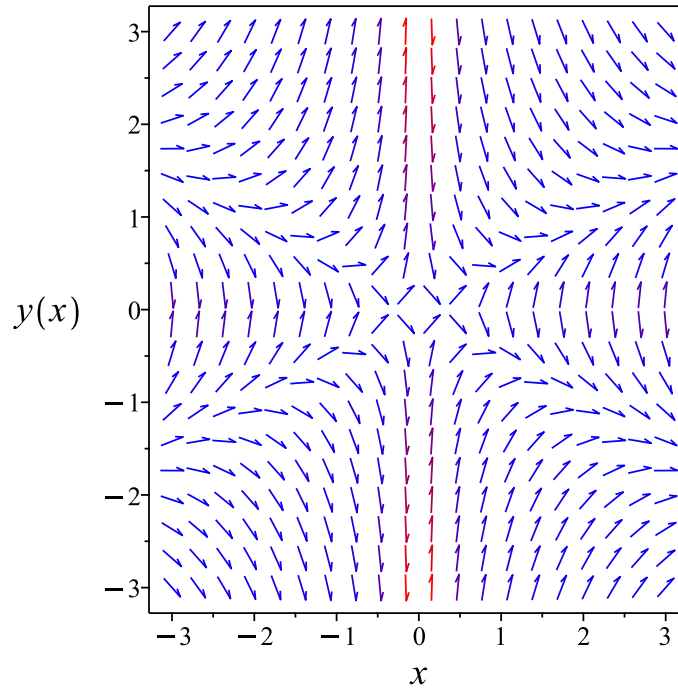


Figure 170: Slope field plot

Verification of solutions

$$y = \frac{\sqrt{5} \sqrt{x(x^5 + 5c_1)}}{5x^2}$$

Verified OK.

$$y = -\frac{\sqrt{5} \sqrt{x(x^5 + 5c_1)}}{5x^2}$$

Verified OK.

2.36.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (2yx) dy &= (x^2 - 3y^2) dx \\ (-x^2 + 3y^2) dx + (2yx) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x^2 + 3y^2 \\ N(x, y) &= 2yx \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (-x^2 + 3y^2) \\ &= 6y \end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(2yx) \\ &= 2y\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{2yx} ((6y) - (2y)) \\ &= \frac{2}{x}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int \frac{2}{x} dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{2\ln(x)} \\ &= x^2\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= x^2(-x^2 + 3y^2) \\ &= -x^2(x^2 - 3y^2)\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= x^2(2yx) \\ &= 2y x^3\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (-x^2(x^2 - 3y^2)) + (2y x^3) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x^2(x^2 - 3y^2) dx \\ \phi &= -\frac{1}{5}x^5 + x^3y^2 + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 2y x^3 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 2y x^3$. Therefore equation (4) becomes

$$2y x^3 = 2y x^3 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{1}{5}x^5 + x^3y^2 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{1}{5}x^5 + x^3y^2$$

Summary

The solution(s) found are the following

$$-\frac{x^5}{5} + x^3y^2 = c_1 \quad (1)$$

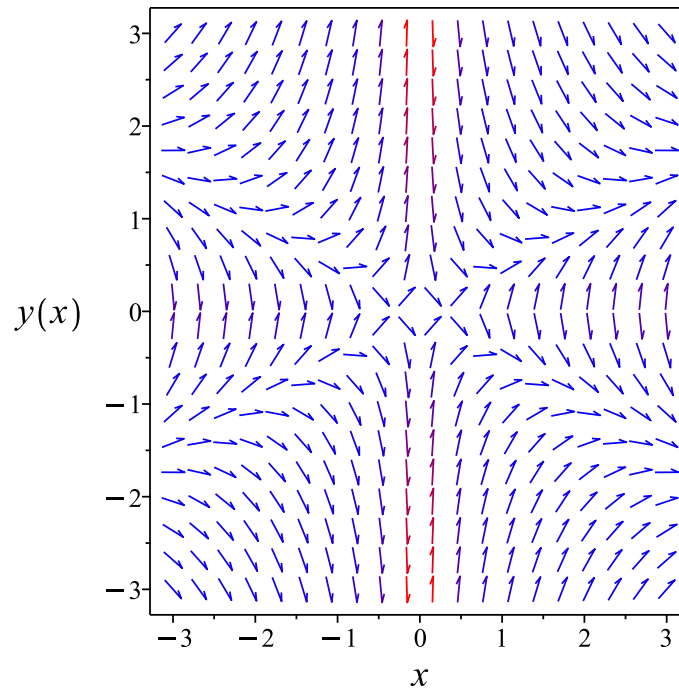


Figure 171: Slope field plot

Verification of solutions

$$-\frac{x^5}{5} + x^3y^2 = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 45

```
dsolve(diff(y(x),x) = (x^2-3*y(x)^2)/(2*x*y(x)),y(x), singsol=all)
```

$$y(x) = -\frac{\sqrt{5} \sqrt{x(x^5 + 5c_1)}}{5x^2}$$
$$y(x) = \frac{\sqrt{5} \sqrt{x(x^5 + 5c_1)}}{5x^2}$$

✓ Solution by Mathematica

Time used: 0.211 (sec). Leaf size: 50

```
DSolve[y'[x] == (x^2-3*y[x]^2)/(2*x*y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{\sqrt{\frac{x^5}{5} + c_1}}{x^{3/2}}$$
$$y(x) \rightarrow \frac{\sqrt{\frac{x^5}{5} + c_1}}{x^{3/2}}$$

2.37 problem 38

2.37.1 Solving as homogeneousTypeD2 ode	892
2.37.2 Solving as first order ode lie symmetry lookup ode	894
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Internal problem ID [515]

Internal file name [OUTPUT/515_Sunday_June_05_2022_01_42_52_AM_69066053/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.2. Page 48

Problem number: 38.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**bernoulli**", "**homogeneousTypeD2**", "**exactWithIntegrationFactor**", "**first_order_ode_lie_symmetry_lookup**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _Bernoulli]
```

$$y' - \frac{3y^2 - x^2}{2yx} = 0$$

2.37.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u'(x)x + u(x) - \frac{3u(x)^2x^2 - x^2}{2u(x)x^2} = 0$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{u^2 - 1}{2xu}\end{aligned}$$

Where $f(x) = \frac{1}{2x}$ and $g(u) = \frac{u^2-1}{u}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{u^2-1}{u}} du &= \frac{1}{2x} dx \\ \int \frac{1}{\frac{u^2-1}{u}} du &= \int \frac{1}{2x} dx \\ \frac{\ln(u-1)}{2} + \frac{\ln(u+1)}{2} &= \frac{\ln(x)}{2} + c_2\end{aligned}$$

The above can be written as

$$\begin{aligned}\left(\frac{1}{2}\right) (\ln(u-1) + \ln(u+1)) &= \frac{\ln(x)}{2} + 2c_2 \\ \ln(u-1) + \ln(u+1) &= (2) \left(\frac{\ln(x)}{2} + 2c_2\right) \\ &= \ln(x) + 4c_2\end{aligned}$$

Raising both side to exponential gives

$$e^{\ln(u-1)+\ln(u+1)} = e^{\ln(x)+2c_2}$$

Which simplifies to

$$\begin{aligned}u^2 - 1 &= 2c_2x \\ &= c_3x\end{aligned}$$

The solution is

$$u(x)^2 - 1 = c_3x$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\begin{aligned}\frac{y^2}{x^2} - 1 &= c_3x \\ \frac{y^2}{x^2} - 1 &= c_3x\end{aligned}$$

Summary

The solution(s) found are the following

$$\frac{y^2}{x^2} - 1 = c_3x \tag{1}$$

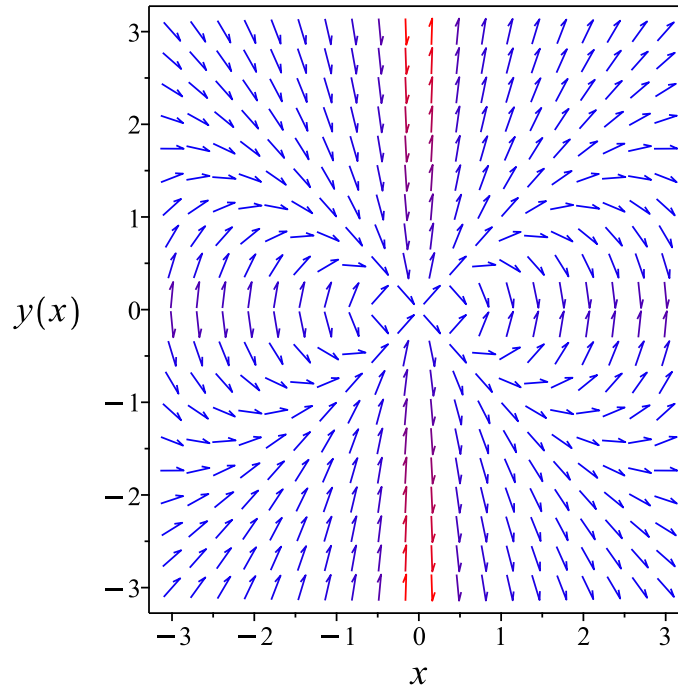


Figure 172: Slope field plot

Verification of solutions

$$\frac{y^2}{x^2} - 1 = c_3 x$$

Verified OK.

2.37.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-x^2 + 3y^2}{2yx}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **Bernoulli**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 180: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{x^3}{y}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{x^3}{y}} dy \end{aligned}$$

Which results in

$$S = \frac{y^2}{2x^3}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{-x^2 + 3y^2}{2yx}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -\frac{3y^2}{2x^4} \\ S_y &= \frac{y}{x^3} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{2x^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{2R^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{1}{2R} + c_1 \quad (4)$$

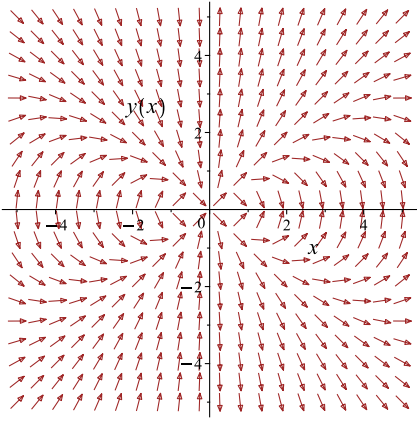
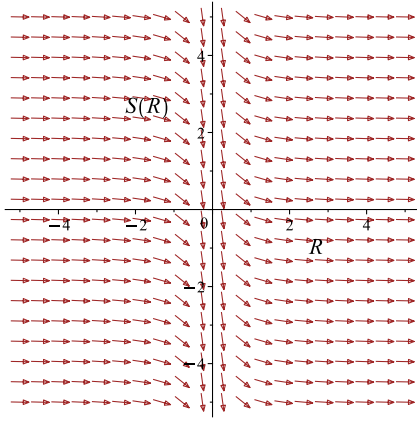
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{y^2}{2x^3} = \frac{1}{2x} + c_1$$

Which simplifies to

$$\frac{y^2}{2x^3} = \frac{1}{2x} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{-x^2 + 3y^2}{2yx}$ 	$R = x$ $S = \frac{y^2}{2x^3}$	$\frac{dS}{dR} = -\frac{1}{2R^2}$ 

Summary

The solution(s) found are the following

$$\frac{y^2}{2x^3} = \frac{1}{2x} + c_1 \quad (1)$$

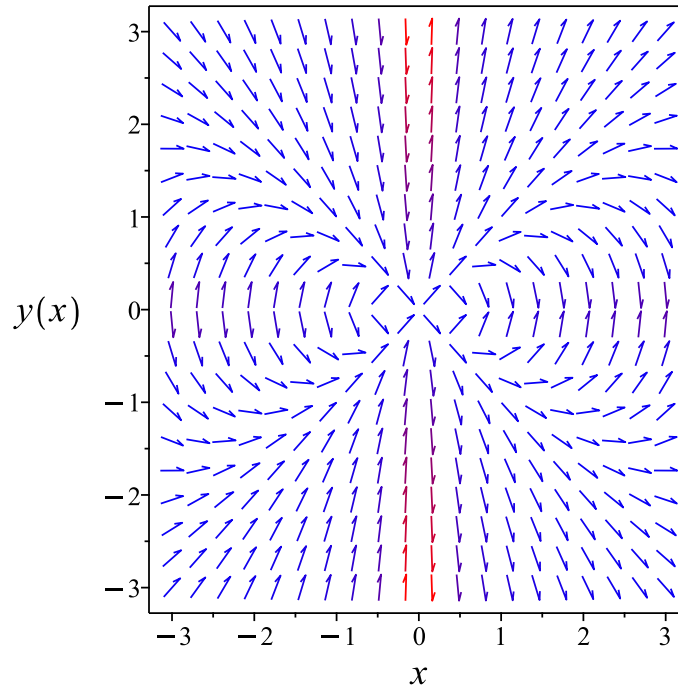


Figure 173: Slope field plot

Verification of solutions

$$\frac{y^2}{2x^3} = \frac{1}{2x} + c_1$$

Verified OK.

2.37.3 Solving as bernoulli ode

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{-x^2 + 3y^2}{2yx} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = \frac{3}{2x}y - \frac{x}{2} \frac{1}{y} \tag{1}$$

The standard Bernoulli ODE has the form

$$y' = f_0(x)y + f_1(x)y^n \tag{2}$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(x)y^{1-n} + f_1(x) \tag{3}$$

The next step is use the substitution $w = y^{1-n}$ in equation (3) which generates a new ODE in $w(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(x) &= \frac{3}{2x} \\ f_1(x) &= -\frac{x}{2} \\ n &= -1 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = \frac{1}{y}$ gives

$$y'y = \frac{3y^2}{2x} - \frac{x}{2} \quad (4)$$

Let

$$\begin{aligned} w &= y^{1-n} \\ &= y^2 \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t x gives

$$w' = 2yy' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} \frac{w'(x)}{2} &= \frac{3w(x)}{2x} - \frac{x}{2} \\ w' &= \frac{3w}{x} - x \end{aligned} \quad (7)$$

The above now is a linear ODE in $w(x)$ which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(x) + p(x)w(x) = q(x)$$

Where here

$$\begin{aligned} p(x) &= -\frac{3}{x} \\ q(x) &= -x \end{aligned}$$

Hence the ode is

$$w'(x) - \frac{3w(x)}{x} = -x$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{3}{x} dx} \\ &= \frac{1}{x^3}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu w) &= (\mu)(-x) \\ \frac{d}{dx}\left(\frac{w}{x^3}\right) &= \left(\frac{1}{x^3}\right)(-x) \\ d\left(\frac{w}{x^3}\right) &= \left(-\frac{1}{x^2}\right) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{w}{x^3} &= \int -\frac{1}{x^2} dx \\ \frac{w}{x^3} &= \frac{1}{x} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = \frac{1}{x^3}$ results in

$$w(x) = c_1 x^3 + x^2$$

Replacing w in the above by y^2 using equation (5) gives the final solution.

$$y^2 = c_1 x^3 + x^2$$

Solving for y gives

$$\begin{aligned}y(x) &= \sqrt{c_1 x + 1} x \\ y(x) &= -\sqrt{c_1 x + 1} x\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \sqrt{c_1 x + 1} x \tag{1}$$

$$y = -\sqrt{c_1 x + 1} x \tag{2}$$

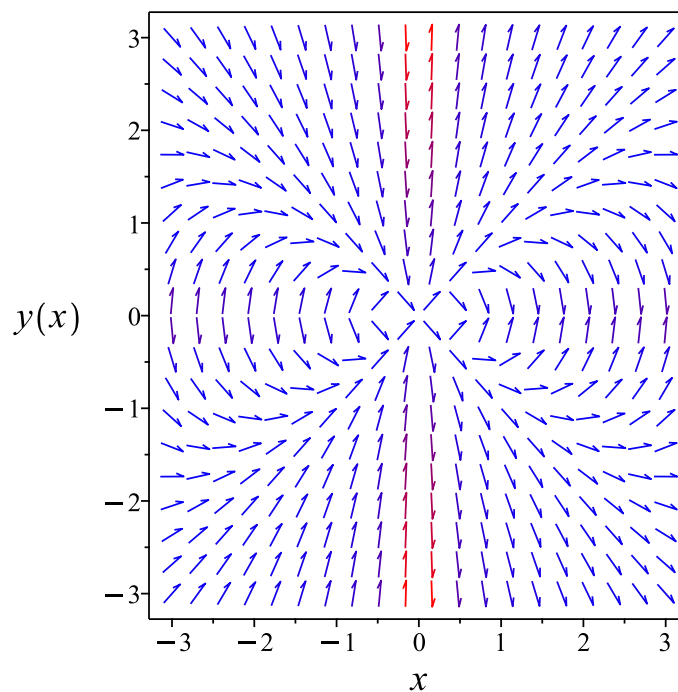


Figure 174: Slope field plot

Verification of solutions

$$y = \sqrt{c_1 x + 1} x$$

Verified OK.

$$y = -\sqrt{c_1 x + 1} x$$

Verified OK.

2.37.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (2yx) dy &= (-x^2 + 3y^2) dx \\ (x^2 - 3y^2) dx + (2yx) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= x^2 - 3y^2 \\ N(x, y) &= 2yx \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(x^2 - 3y^2) \\ &= -6y \end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(2yx) \\ &= 2y\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{2yx} ((-6y) - (2y)) \\ &= -\frac{4}{x}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A \, dx} \\ &= e^{\int -\frac{4}{x} \, dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-4 \ln(x)} \\ &= \frac{1}{x^4}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{x^4}(x^2 - 3y^2) \\ &= \frac{x^2 - 3y^2}{x^4}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{x^4}(2yx) \\ &= \frac{2y}{x^3}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left(\frac{x^2 - 3y^2}{x^4} \right) + \left(\frac{2y}{x^3} \right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{x^2 - 3y^2}{x^4} dx \\ \phi &= -\frac{1}{x} + \frac{y^2}{x^3} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = \frac{2y}{x^3} + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{2y}{x^3}$. Therefore equation (4) becomes

$$\frac{2y}{x^3} = \frac{2y}{x^3} + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{1}{x} + \frac{y^2}{x^3} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{1}{x} + \frac{y^2}{x^3}$$

Summary

The solution(s) found are the following

$$-\frac{1}{x} + \frac{y^2}{x^3} = c_1 \quad (1)$$

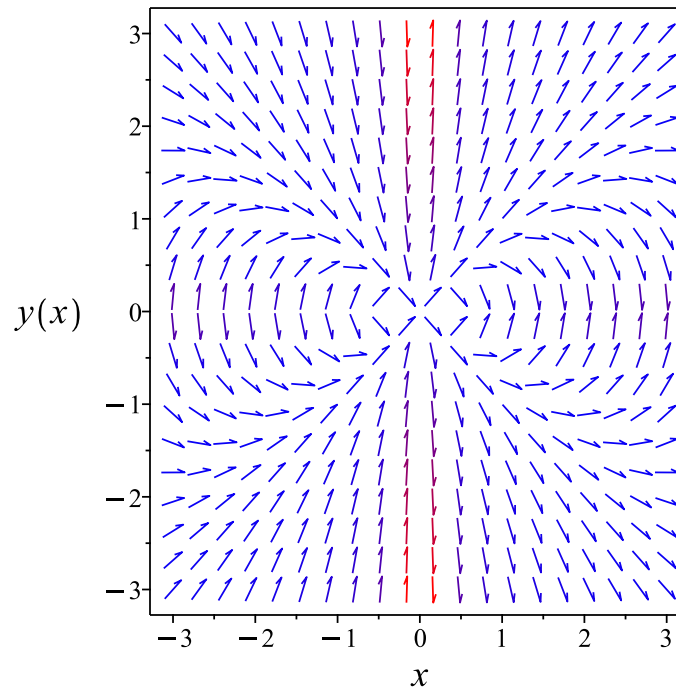


Figure 175: Slope field plot

Verification of solutions

$$-\frac{1}{x} + \frac{y^2}{x^3} = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 26

```
dsolve(diff(y(x),x) = (3*y(x)^2-x^2)/(2*x*y(x)),y(x), singsol=all)
```

$$y(x) = \sqrt{c_1x + 1}x$$
$$y(x) = -\sqrt{c_1x + 1}x$$

✓ Solution by Mathematica

Time used: 0.209 (sec). Leaf size: 34

```
DSolve[y'[x] == (3*y[x]^2-x^2)/(2*x*y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -x\sqrt{1 + c_1x}$$
$$y(x) \rightarrow x\sqrt{1 + c_1x}$$

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3.1 problem 1

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Internal problem ID [516]

Internal file name [OUTPUT/516_Sunday_June_05_2022_01_42_53_AM_57257413/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_linear]

$$\ln(t)y + (t - 3)y' = 2t$$

3.1.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{\ln(t)}{t - 3}$$

$$q(t) = \frac{2t}{t - 3}$$

Hence the ode is

$$y' + \frac{\ln(t)y}{t - 3} = \frac{2t}{t - 3}$$

The integrating factor μ is

$$\mu = e^{\int \frac{\ln(t)}{t-3} dt}$$

The ode becomes

$$\begin{aligned} \frac{d}{dt}(\mu y) &= (\mu) \left(\frac{2t}{t-3} \right) \\ \frac{d}{dt} \left(e^{\int \frac{\ln(t)}{t-3} dt} y \right) &= \left(e^{\int \frac{\ln(t)}{t-3} dt} \right) \left(\frac{2t}{t-3} \right) \\ d \left(e^{\int \frac{\ln(t)}{t-3} dt} y \right) &= \left(\frac{2t e^{\int \frac{\ln(t)}{t-3} dt}}{t-3} \right) dt \end{aligned}$$

Integrating gives

$$\begin{aligned} e^{\int \frac{\ln(t)}{t-3} dt} y &= \int \frac{2t e^{\int \frac{\ln(t)}{t-3} dt}}{t-3} dt \\ e^{\int \frac{\ln(t)}{t-3} dt} y &= \int \frac{2t e^{\int \frac{\ln(t)}{t-3} dt}}{t-3} dt + c_1 \end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{\int \frac{\ln(t)}{t-3} dt}$ results in

$$y = e^{-\left(\int \frac{\ln(t)}{t-3} dt\right)} \left(\int \frac{2t e^{\int \frac{\ln(t)}{t-3} dt}}{t-3} dt \right) + c_1 e^{-\left(\int \frac{\ln(t)}{t-3} dt\right)}$$

which simplifies to

$$y = e^{-\left(\int \frac{\ln(t)}{t-3} dt\right)} \left(2 \left(\int \frac{t e^{\int \frac{\ln(t)}{t-3} dt}}{t-3} dt \right) + c_1 \right)$$

Which can be simplified to become

$$y = e^{\int -\frac{\ln(t)}{t-3} dt} \left(2 \left(\int \frac{t e^{\int \frac{\ln(t)}{t-3} dt}}{t-3} dt \right) + c_1 \right)$$

Summary

The solution(s) found are the following

$$y = e^{\int -\frac{\ln(t)}{t-3} dt} \left(2 \left(\int \frac{t e^{\int \frac{\ln(t)}{t-3} dt}}{t-3} dt \right) + c_1 \right) \quad (1)$$

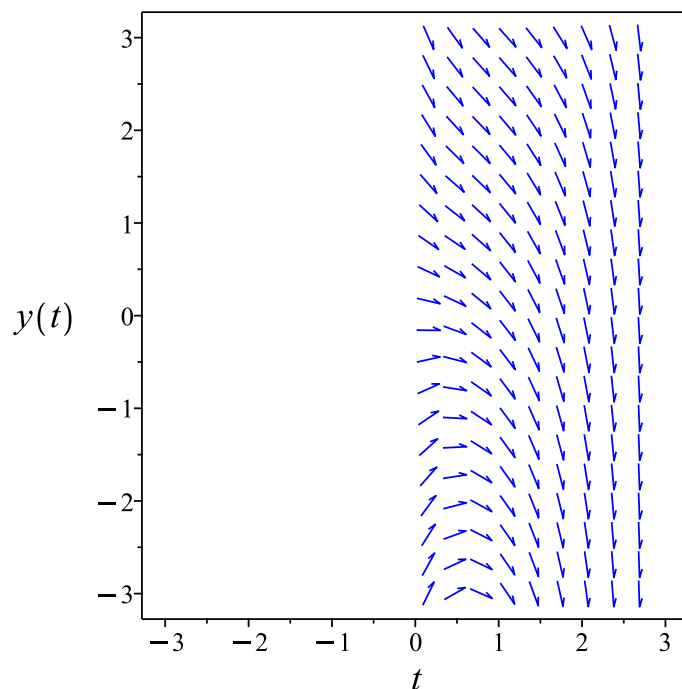


Figure 176: Slope field plot

Verification of solutions

$$y = e^{\int -\frac{\ln(t)}{t-3} dt} \left(2 \left(\int \frac{t e^{\int \frac{\ln(t)}{t-3} dt}}{t-3} dt \right) + c_1 \right)$$

Verified OK.

3.1.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{\ln(t)y - 2t}{t-3}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 182: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int(n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-(\ln(t)-\ln(\frac{t}{3}))\ln(-\frac{t}{3}+1)+\text{dilog}(\frac{t}{3})}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-(\ln(t) - \ln(\frac{t}{3})) \ln(-\frac{t}{3} + 1) + \text{dilog}(\frac{t}{3})}} dy \end{aligned}$$

Which results in

$$S = e^{-\ln(\frac{1}{3}) \ln(-\frac{t}{3} + 1) - \text{dilog}(\frac{t}{3})} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y) S_y}{R_t + \omega(t, y) R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{\ln(t) y - 2t}{t - 3}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -(-t + 3)^{-1 + \ln(3)} y \ln(t) e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} \\ S_y &= (-t + 3)^{\ln(3)} e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -2 e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} (-t + 3)^{-1 + \ln(3)} t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -2 e^{-\ln(3)^2 - \text{dilog}(\frac{R}{3})} (-R + 3)^{-1 + \ln(3)} R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \int -2 e^{-\ln(3)^2 - \operatorname{dilog}\left(\frac{R}{3}\right)} (-R + 3)^{-1 + \ln(3)} R dR + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$y(-t + 3)^{\ln(3)} e^{-\ln(3)^2 - \operatorname{dilog}\left(\frac{t}{3}\right)} = \int -2 e^{-\ln(3)^2 - \operatorname{dilog}\left(\frac{t}{3}\right)} (-t + 3)^{-1 + \ln(3)} t dt + c_1$$

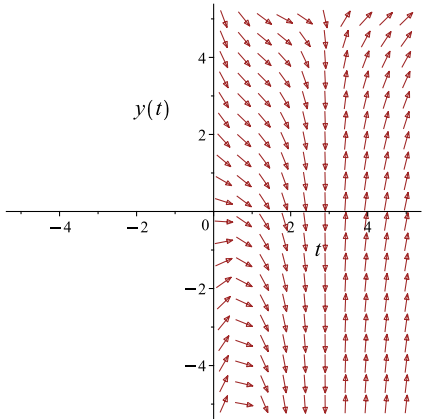
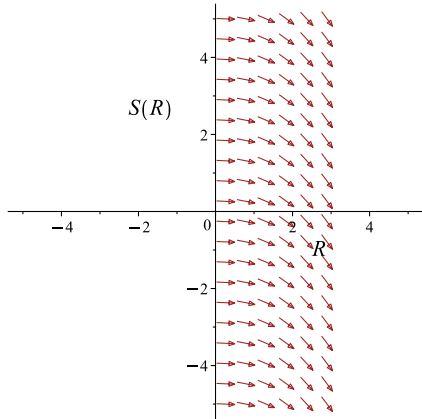
Which simplifies to

$$y(-t + 3)^{\ln(3)} e^{-\ln(3)^2 - \operatorname{dilog}\left(\frac{t}{3}\right)} = \int -2 e^{-\ln(3)^2 - \operatorname{dilog}\left(\frac{t}{3}\right)} (-t + 3)^{-1 + \ln(3)} t dt + c_1$$

Which gives

$$y = \left(\int -2 e^{-\ln(3)^2 - \operatorname{dilog}\left(\frac{t}{3}\right)} (-t + 3)^{-1 + \ln(3)} t dt + c_1 \right) (-t + 3)^{-\ln(3)} e^{\ln(3)^2 + \operatorname{dilog}\left(\frac{t}{3}\right)}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{\ln(t)y - 2t}{t - 3}$ 	$R = t$ $S = y(-t + 3)^{\ln(3)} e^{-\ln(3)^2}$	$\frac{dS}{dR} = -2 e^{-\ln(3)^2 - \operatorname{dilog}\left(\frac{R}{3}\right)} (-R + 3)^{-1 + \ln(3)} R$ 

Summary

The solution(s) found are the following

$$y = \left(\int -2 e^{-\ln(3)^2 - \operatorname{dilog}\left(\frac{t}{3}\right)} (-t + 3)^{-1 + \ln(3)} t dt + c_1 \right) (-t + 3)^{-\ln(3)} e^{\ln(3)^2 + \operatorname{dilog}\left(\frac{t}{3}\right)} \quad (1)$$

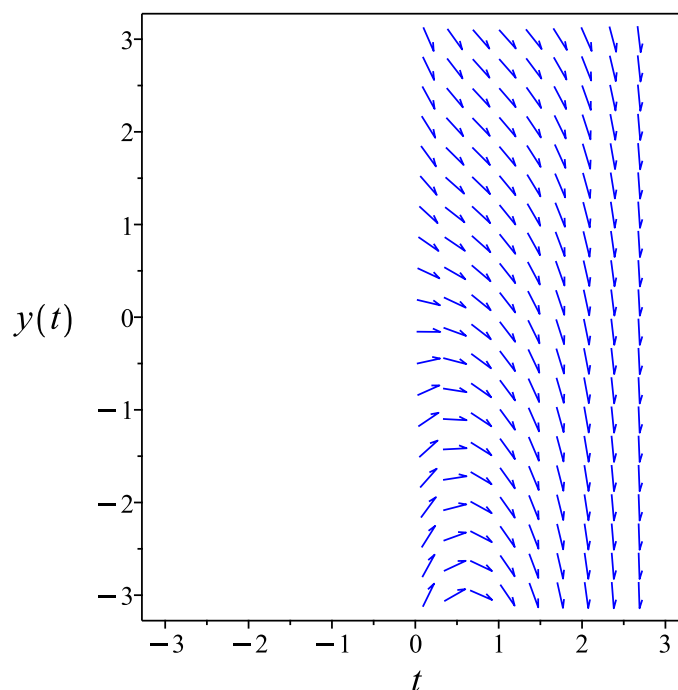


Figure 177: Slope field plot

Verification of solutions

$$y = \left(\int -2 e^{-\ln(3)^2 - \operatorname{dilog}\left(\frac{t}{3}\right)} (-t + 3)^{-1 + \ln(3)} t dt + c_1 \right) (-t + 3)^{-\ln(3)} e^{\ln(3)^2 + \operatorname{dilog}\left(\frac{t}{3}\right)}$$

Verified OK.

3.1.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the

ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (t - 3) dy &= (-\ln(t) y + 2t) dt \\ (\ln(t) y - 2t) dt + (t - 3) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= \ln(t) y - 2t \\ N(t, y) &= t - 3 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(\ln(t)y - 2t) \\ &= \ln(t)\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(t - 3) \\ &= 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{t-3} ((\ln(t)) - (1)) \\ &= \frac{\ln(t) - 1}{t-3}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int \frac{\ln(t)-1}{t-3} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{(\ln(t)-\ln(\frac{t}{3})) \ln(-\frac{t}{3}+1) - \text{dilog}(\frac{t}{3}) - \ln(t-3)} \\ &= -(-t+3)^{-1+\ln(3)} e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= -(-t+3)^{-1+\ln(3)} e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} (\ln(t)y - 2t) \\ &= (-t+3)^{-1+\ln(3)} (-\ln(t)y + 2t) e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= -(-t+3)^{-1+\ln(3)} e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} (t-3) \\ &= (-t+3)^{\ln(3)} e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\overline{M} + \overline{N} \frac{dy}{dt} = 0$$

$$\left((-t + 3)^{-1+\ln(3)} (-\ln(t)y + 2t) e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} \right) + \left((-t + 3)^{\ln(3)} e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} \right) \frac{dy}{dt} = 0$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int \overline{M} dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int (-t + 3)^{-1+\ln(3)} (-\ln(t)y + 2t) e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} dt$$

$$\phi = \int^t (-a + 3)^{-1+\ln(3)} (-\ln(a)y + 2a) e^{-\ln(3)^2 - \text{dilog}(\frac{a}{3})} da + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = - \left(\int^t (-a + 3)^{-1+\ln(3)} \ln(a) e^{-\ln(3)^2 - \text{dilog}(\frac{a}{3})} da \right) + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = (-t + 3)^{\ln(3)} e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})}$. Therefore equation (4) becomes

$$(-t + 3)^{\ln(3)} e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} = - \left(\int^t (-a + 3)^{-1+\ln(3)} \ln(a) e^{-\ln(3)^2 - \text{dilog}(\frac{a}{3})} da \right) + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = (-t + 3)^{\ln(3)} e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} + \int^t (-a + 3)^{-1+\ln(3)} \ln(a) e^{-\ln(3)^2 - \text{dilog}(\frac{a}{3})} da$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left((-t+3)^{\ln(3)} e^{-\ln(3)^2 - \operatorname{dilog}(\frac{t}{3})} \right. \\ \left. + \int^t (-_a+3)^{-1+\ln(3)} \ln(_a) e^{-\ln(3)^2 - \operatorname{dilog}(\frac{-a}{3})} d_a \right) dy \\ f(y) = \left((-t+3)^{\ln(3)} e^{-\ln(3)^2 - \operatorname{dilog}(\frac{t}{3})} \right. \\ \left. + \int^t (-_a+3)^{-1+\ln(3)} \ln(_a) e^{-\ln(3)^2 - \operatorname{dilog}(\frac{-a}{3})} d_a \right) y + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = \int^t (-_a+3)^{-1+\ln(3)} (-\ln(_a)y + 2_a) e^{-\ln(3)^2 - \operatorname{dilog}(\frac{-a}{3})} d_a \\ + \left((-t+3)^{\ln(3)} e^{-\ln(3)^2 - \operatorname{dilog}(\frac{t}{3})} \right. \\ \left. + \int^t (-_a+3)^{-1+\ln(3)} \ln(_a) e^{-\ln(3)^2 - \operatorname{dilog}(\frac{-a}{3})} d_a \right) y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \int^t (-_a+3)^{-1+\ln(3)} (-\ln(_a)y + 2_a) e^{-\ln(3)^2 - \operatorname{dilog}(\frac{-a}{3})} d_a \\ + \left((-t+3)^{\ln(3)} e^{-\ln(3)^2 - \operatorname{dilog}(\frac{t}{3})} \right. \\ \left. + \int^t (-_a+3)^{-1+\ln(3)} \ln(_a) e^{-\ln(3)^2 - \operatorname{dilog}(\frac{-a}{3})} d_a \right) y$$

Summary

The solution(s) found are the following

$$\int^t (-_a+3)^{-1+\ln(3)} (-\ln(_a)y + 2_a) e^{-\ln(3)^2 - \operatorname{dilog}(\frac{-a}{3})} d_a \\ + \left((-t+3)^{\ln(3)} e^{-\ln(3)^2 - \operatorname{dilog}(\frac{t}{3})} \right. \\ \left. + \int^t (-_a+3)^{-1+\ln(3)} \ln(_a) e^{-\ln(3)^2 - \operatorname{dilog}(\frac{-a}{3})} d_a \right) y = c_1 \quad (1)$$

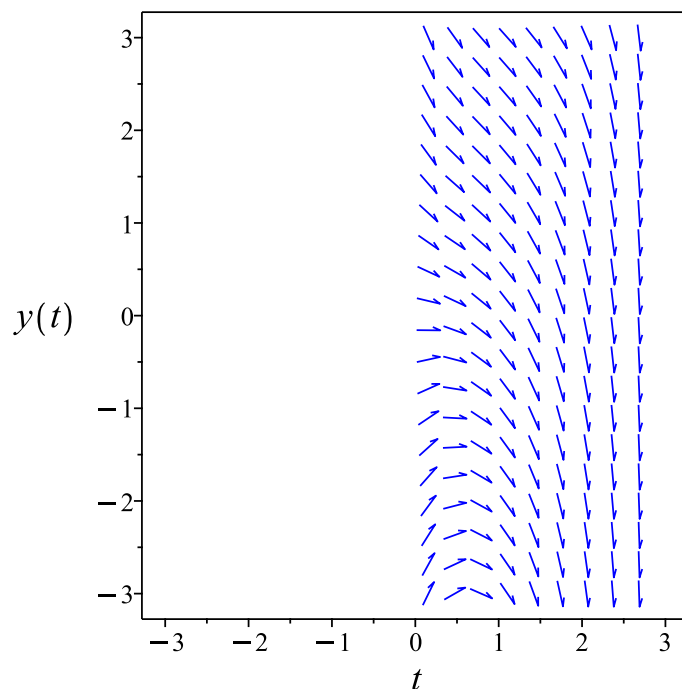


Figure 178: Slope field plot

Verification of solutions

$$\int^t (-_a + 3)^{-1+\ln(3)} (-\ln(_a) y + 2_a) e^{-\ln(3)^2 - \text{dilog}(\frac{a}{3})} d_a$$

$$+ \left((-t + 3)^{\ln(3)} e^{-\ln(3)^2 - \text{dilog}(\frac{t}{3})} \right.$$

$$\left. + \int^t (-_a + 3)^{-1+\ln(3)} \ln(_a) e^{-\ln(3)^2 - \text{dilog}(\frac{a}{3})} d_a \right) y = c_1$$

Verified OK.

3.1.4 Maple step by step solution

Let's solve

$$\ln(t)y + (t-3)y' = 2t$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{\ln(t)y}{t-3} + \frac{2t}{t-3}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{\ln(t)y}{t-3} = \frac{2t}{t-3}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{\ln(t)y}{t-3} \right) = \frac{2\mu(t)t}{t-3}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{\ln(t)y}{t-3} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)\ln(t)}{t-3}$$

- Solve to find the integrating factor

$$\mu(t) = 3^{\ln(-t+3)} e^{-\ln(3)^2 - \text{dilog}\left(\frac{t}{3}\right)}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{2\mu(t)t}{t-3} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{2\mu(t)t}{t-3} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{2\mu(t)t}{t-3} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = 3^{\ln(-t+3)} e^{-\ln(3)^2 - \text{dilog}\left(\frac{t}{3}\right)}$

$$y = \frac{\int \frac{2 \cdot 3^{\ln(-t+3)} e^{-\ln(3)^2 - \text{dilog}\left(\frac{t}{3}\right)} t}{t-3} dt + c_1}{3^{\ln(-t+3)} e^{-\ln(3)^2 - \text{dilog}\left(\frac{t}{3}\right)}}$$

- Simplify

$$y = \left(-2 \left(\int e^{-\ln(3)^2 - \text{dilog}\left(\frac{t}{3}\right)} (-t+3)^{-1+\ln(3)} t dt \right) + c_1 \right) (-t+3)^{-\ln(3)} e^{\ln(3)^2 + \text{dilog}\left(\frac{t}{3}\right)}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 57

```
dsolve(ln(t)*y(t)+(-3+t)*diff(y(t),t) = 2*t,y(t), singsol=all)
```

$$y(t) = e^{\ln(3)^2 + \operatorname{dilog}\left(\frac{t}{3}\right)} (-t + 3)^{-\ln(3)} \left(-2 \left(\int t (-t + 3)^{-1 + \ln(3)} e^{-\ln(3)^2 - \operatorname{dilog}\left(\frac{t}{3}\right)} dt \right) + c_1 \right)$$

✓ Solution by Mathematica

Time used: 0.188 (sec). Leaf size: 69

```
DSolve[Log[t]*y[t]+(-3+t)*y'[t] == 2*t,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{\operatorname{PolyLog}\left(2, 1 - \frac{t}{3}\right) - \log(3) \log(t-3)} \left(\int_1^t \frac{2e^{\log(3) \log(K[1]-3) - \operatorname{PolyLog}\left(2, 1 - \frac{K[1]}{3}\right)}}{K[1] - 3} K[1] dK[1] + c_1 \right)$$

3.2 problem 2

3.2.1	Existence and uniqueness analysis	923
3.2.2	Solving as separable ode	923
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3.2.4	Solving as homogeneousTypeD2 ode	926
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Internal problem ID [517]

Internal file name [OUTPUT/517_Sunday_June_05_2022_01_42_55_AM_26177264/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

`[_separable]`

$$y + (t - 4) ty' = 0$$

With initial conditions

$$[y(2) = 1]$$

3.2.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{1}{(t-4)t}$$

$$q(t) = 0$$

Hence the ode is

$$y' + \frac{y}{(t-4)t} = 0$$

The domain of $p(t) = \frac{1}{(t-4)t}$ is

$$\{-\infty \leq t < 0, 0 < t < 4, 4 < t \leq \infty\}$$

And the point $t_0 = 2$ is inside this domain. Hence solution exists and is unique.

3.2.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= f(t)g(y) \\ &= -\frac{y}{(t-4)t} \end{aligned}$$

Where $f(t) = -\frac{1}{(t-4)t}$ and $g(y) = y$. Integrating both sides gives

$$\begin{aligned} \frac{1}{y} dy &= -\frac{1}{(t-4)t} dt \\ \int \frac{1}{y} dy &= \int -\frac{1}{(t-4)t} dt \\ \ln(y) &= -\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4} + c_1 \\ y &= e^{-\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4} + c_1} \\ &= c_1 e^{-\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4}} \end{aligned}$$

Which can be simplified to become

$$y = \frac{c_1 t^{\frac{1}{4}}}{(t-4)^{\frac{1}{4}}}$$

Initial conditions are used to solve for c_1 . Substituting $t = 2$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{\sqrt{2} c_1}{2} - \frac{i\sqrt{2} c_1}{2}$$

$$c_1 = \left(\frac{1}{2} + \frac{i}{2}\right) \sqrt{2}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{\sqrt{2} t^{\frac{1}{4}} + i\sqrt{2} t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{2} t^{\frac{1}{4}} + i\sqrt{2} t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}} \quad (1)$$

Verification of solutions

$$y = \frac{\sqrt{2} t^{\frac{1}{4}} + i\sqrt{2} t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}}$$

Verified OK.

3.2.3 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int \frac{1}{(t-4)t} dt} \\ &= e^{\frac{\ln(t-4)}{4} - \frac{\ln(t)}{4}} \end{aligned}$$

Which simplifies to

$$\mu = \frac{(t-4)^{\frac{1}{4}}}{t^{\frac{1}{4}}}$$

The ode becomes

$$\frac{d}{dt} \mu y = 0$$
$$\frac{d}{dt} \left(\frac{(t-4)^{\frac{1}{4}} y}{t^{\frac{1}{4}}} \right) = 0$$

Integrating gives

$$\frac{(t-4)^{\frac{1}{4}} y}{t^{\frac{1}{4}}} = c_1$$

Dividing both sides by the integrating factor $\mu = \frac{(t-4)^{\frac{1}{4}}}{t^{\frac{1}{4}}}$ results in

$$y = \frac{c_1 t^{\frac{1}{4}}}{(t-4)^{\frac{1}{4}}}$$

Initial conditions are used to solve for c_1 . Substituting $t = 2$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{\sqrt{2} c_1}{2} - \frac{i\sqrt{2} c_1}{2}$$

$$c_1 = \left(\frac{1}{2} + \frac{i}{2} \right) \sqrt{2}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{\sqrt{2} t^{\frac{1}{4}} + i\sqrt{2} t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{2} t^{\frac{1}{4}} + i\sqrt{2} t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}} \tag{1}$$

Verification of solutions

$$y = \frac{\sqrt{2} t^{\frac{1}{4}} + i\sqrt{2} t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}}$$

Verified OK.

3.2.4 Solving as homogeneous Type D2 ode

Using the change of variables $y = u(t)t$ on the above ode results in new ode in $u(t)$

$$u(t)t + (t-4)t(u'(t)t + u(t)) = 0$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{u(t-3)}{t(t-4)}\end{aligned}$$

Where $f(t) = -\frac{t-3}{(t-4)t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= -\frac{t-3}{(t-4)t} dt \\ \int \frac{1}{u} du &= \int -\frac{t-3}{(t-4)t} dt \\ \ln(u) &= -\frac{\ln(t-4)}{4} - \frac{3\ln(t)}{4} + c_2 \\ u &= e^{-\frac{\ln(t-4)}{4} - \frac{3\ln(t)}{4} + c_2} \\ &= c_2 e^{-\frac{\ln(t-4)}{4} - \frac{3\ln(t)}{4}}\end{aligned}$$

Which simplifies to

$$u(t) = \frac{c_2}{(t-4)^{\frac{1}{4}} t^{\frac{3}{4}}}$$

Therefore the solution y is

$$\begin{aligned}y &= ut \\ &= \frac{t^{\frac{1}{4}} c_2}{(t-4)^{\frac{1}{4}}}\end{aligned}$$

Initial conditions are used to solve for c_2 . Substituting $t = 2$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$\begin{aligned}1 &= \frac{\sqrt{2} c_2}{2} - \frac{i\sqrt{2} c_2}{2} \\ c_2 &= \left(\frac{1}{2} + \frac{i}{2}\right) \sqrt{2}\end{aligned}$$

Substituting c_2 found above in the general solution gives

$$y = \frac{\sqrt{2}t^{\frac{1}{4}} + i\sqrt{2}t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{2}t^{\frac{1}{4}} + i\sqrt{2}t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}} \quad (1)$$

Verification of solutions

$$y = \frac{\sqrt{2}t^{\frac{1}{4}} + i\sqrt{2}t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}}$$

Verified OK.

3.2.5 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{y}{(t-4)t}$$
$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2\xi_y - \omega_t\xi - \omega_y\eta = 0 \quad (A)$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 185: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4}}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4}}} dy \end{aligned}$$

Which results in

$$S = e^{\ln((t-4)^{\frac{1}{4}}) + \ln\left(\frac{1}{t^{\frac{1}{4}}}\right)} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{y}{(t-4)t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= \frac{y}{(t-4)^{\frac{3}{4}} t^{\frac{5}{4}}} \\ S_y &= \frac{(t-4)^{\frac{1}{4}}}{t^{\frac{1}{4}}} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$\frac{(t-4)^{\frac{1}{4}} y}{t^{\frac{1}{4}}} = c_1$$

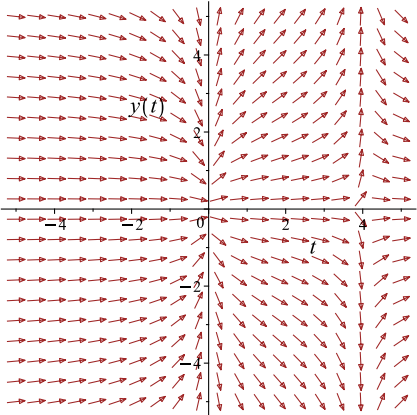
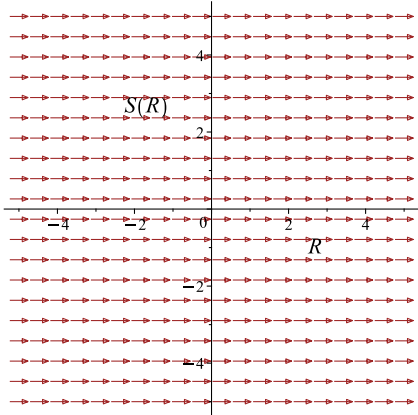
Which simplifies to

$$\frac{(t-4)^{\frac{1}{4}} y}{t^{\frac{1}{4}}} = c_1$$

Which gives

$$y = \frac{c_1 t^{\frac{1}{4}}}{(t-4)^{\frac{1}{4}}}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{y}{(t-4)t}$ 	$R = t$ $S = \frac{(t-4)^{\frac{1}{4}} y}{t^{\frac{1}{4}}}$	$\frac{dS}{dR} = 0$ 

Initial conditions are used to solve for c_1 . Substituting $t = 2$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{\sqrt{2} c_1}{2} - \frac{i\sqrt{2} c_1}{2}$$

$$c_1 = \left(\frac{1}{2} + \frac{i}{2} \right) \sqrt{2}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{\sqrt{2} t^{\frac{1}{4}} + i\sqrt{2} t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{2} t^{\frac{1}{4}} + i\sqrt{2} t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}} \quad (1)$$

Verification of solutions

$$y = \frac{\sqrt{2} t^{\frac{1}{4}} + i\sqrt{2} t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}}$$

Verified OK.

3.2.6 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}\left(-\frac{1}{y}\right) dy &= \left(\frac{1}{(t-4)t}\right) dt \\ \left(-\frac{1}{(t-4)t}\right) dt + \left(-\frac{1}{y}\right) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -\frac{1}{(t-4)t} \\ N(t, y) &= -\frac{1}{y}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{1}{(t-4)t}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} \left(-\frac{1}{y} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -\frac{1}{(t-4)t} dt \\ \phi &= -\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{1}{y}$. Therefore equation (4) becomes

$$-\frac{1}{y} = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{1}{y}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(-\frac{1}{y} \right) dy \\ f(y) &= -\ln(y) + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4} - \ln(y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4} - \ln(y)$$

The solution becomes

$$y = e^{-\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4} - c_1}$$

Initial conditions are used to solve for c_1 . Substituting $t = 2$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{\sqrt{2}e^{-c_1}}{2} - \frac{i\sqrt{2}e^{-c_1}}{2}$$

$$c_1 = -\frac{i\pi}{4}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{\sqrt{2}t^{\frac{1}{4}} + i\sqrt{2}t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{2}t^{\frac{1}{4}} + i\sqrt{2}t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}} \quad (1)$$

Verification of solutions

$$y = \frac{\sqrt{2}t^{\frac{1}{4}} + i\sqrt{2}t^{\frac{1}{4}}}{2(t-4)^{\frac{1}{4}}}$$

Verified OK.

3.2.7 Maple step by step solution

Let's solve

$$[y + (t - 4)ty' = 0, y(2) = 1]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{y} = -\frac{1}{(t-4)t}$$

- Integrate both sides with respect to t

$$\int \frac{y'}{y} dt = \int -\frac{1}{(t-4)t} dt + c_1$$

- Evaluate integral

$$\ln(y) = -\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4} + c_1$$

- Use initial condition $y(2) = 1$

$$0 = -\frac{I\pi}{4} + c_1$$

- Solve for c_1

$$c_1 = \frac{I}{4}\pi$$

- Substitute $c_1 = \frac{I}{4}\pi$ into general solution and simplify

$$\ln(y) = -\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4} + \frac{I\pi}{4}$$

- Solution to the IVP

$$\ln(y) = -\frac{\ln(t-4)}{4} + \frac{\ln(t)}{4} + \frac{I\pi}{4}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```


✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 19

```
dsolve([y(t)+(-4+t)*t*diff(y(t),t) = 0,y(2) = 1],y(t), singsol=all)
```

$$y(t) = \frac{\left(\frac{1}{2} + \frac{i}{2}\right) \sqrt{2} t^{\frac{1}{4}}}{(-4 + t)^{\frac{1}{4}}}$$

✓ Solution by Mathematica

Time used: 0.033 (sec). Leaf size: 20

```
DSolve[{y[t]+(-4+t)*t*y'[t] == 0,y[2]==1},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{\sqrt[4]{t}}{\sqrt[4]{4-t}}$$

3.3 problem 3

3.3.1	Existence and uniqueness analysis	937
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3.3.5	Maple step by step solution	947

Internal problem ID [518]

Internal file name [OUTPUT/518_Sunday_June_05_2022_01_42_55_AM_79033667/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 3.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_linear]

$$\tan(t)y + y' = \sin(t)$$

With initial conditions

$$[y(\pi) = 0]$$

3.3.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \tan(t)$$

$$q(t) = \sin(t)$$

Hence the ode is

$$\tan(t) y + y' = \sin(t)$$

The domain of $p(t) = \tan(t)$ is

$$\left\{ t < \frac{1}{2}\pi + \pi_{-Z8} \vee \frac{1}{2}\pi + \pi_{-Z8} < t \right\}$$

And the point $t_0 = \pi$ is inside this domain. The domain of $q(t) = \sin(t)$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = \pi$ is also inside this domain. Hence solution exists and is unique.

3.3.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int \tan(t) dt} \\ &= \frac{1}{\cos(t)} \end{aligned}$$

Which simplifies to

$$\mu = \sec(t)$$

The ode becomes

$$\begin{aligned} \frac{d}{dt}(\mu y) &= (\mu) (\sin(t)) \\ \frac{d}{dt}(\sec(t) y) &= (\sec(t)) (\sin(t)) \\ d(\sec(t) y) &= \tan(t) dt \end{aligned}$$

Integrating gives

$$\begin{aligned} \sec(t) y &= \int \tan(t) dt \\ \sec(t) y &= -\ln(\cos(t)) + c_1 \end{aligned}$$

Dividing both sides by the integrating factor $\mu = \sec(t)$ results in

$$y = -\cos(t) \ln(\cos(t)) + c_1 \cos(t)$$

which simplifies to

$$y = \cos(t) (-\ln(\cos(t)) + c_1)$$

Initial conditions are used to solve for c_1 . Substituting $t = \pi$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = i\pi - c_1$$

$$c_1 = i\pi$$

Substituting c_1 found above in the general solution gives

$$y = i \cos(t) \pi - \cos(t) \ln(\cos(t))$$

Summary

The solution(s) found are the following

$$y = i \cos(t) \pi - \cos(t) \ln(\cos(t)) \quad (1)$$

Verification of solutions

$$y = i \cos(t) \pi - \cos(t) \ln(\cos(t))$$

Verified OK.

3.3.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$\begin{aligned} y' &= -\tan(t) y + \sin(t) \\ y' &= \omega(t, y) \end{aligned}$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (A)$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 188: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \cos(t)\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\cos(t)} dy \end{aligned}$$

Which results in

$$S = \frac{y}{\cos(t)}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\tan(t)y + \sin(t)$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= \sec(t) \tan(t) y \\ S_y &= \sec(t) \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \tan(t) \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \tan(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\ln(\cos(R)) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$\sec(t) y = -\ln(\cos(t)) + c_1$$

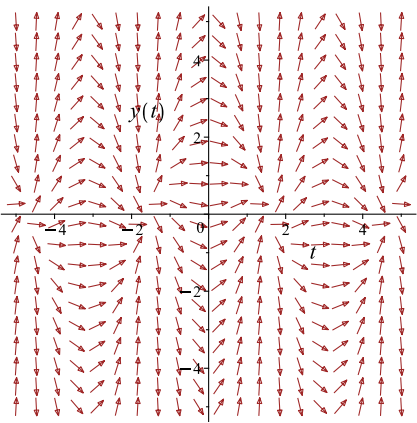
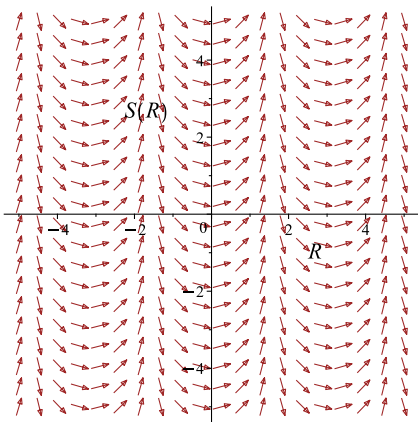
Which simplifies to

$$\sec(t) y = -\ln(\cos(t)) + c_1$$

Which gives

$$y = -\frac{\ln(\cos(t)) - c_1}{\sec(t)}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\tan(t) y + \sin(t)$ 	$R = t$ $S = \sec(t) y$	$\frac{dS}{dR} = \tan(R)$ 

Initial conditions are used to solve for c_1 . Substituting $t = \pi$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = i\pi - c_1$$

$$c_1 = i\pi$$

Substituting c_1 found above in the general solution gives

$$y = i \cos(t) \pi - \cos(t) \ln(\cos(t))$$

Summary

The solution(s) found are the following

$$y = i \cos(t) \pi - \cos(t) \ln(\cos(t)) \quad (1)$$

Verification of solutions

$$y = i \cos(t) \pi - \cos(t) \ln(\cos(t))$$

Verified OK.

3.3.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition

$\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} dy &= (-\tan(t)y + \sin(t)) dt \\ (\tan(t)y - \sin(t)) dt + dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= \tan(t)y - \sin(t) \\ N(t, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(\tan(t)y - \sin(t)) \\ &= \tan(t) \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(1) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= 1((\tan(t)) - (0)) \\ &= \tan(t) \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int \tan(t) dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\ln(\cos(t))} \\ &= \sec(t)\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\overline{M} &= \mu M \\ &= \sec(t) (\tan(t) y - \sin(t)) \\ &= \tan(t) (\sec(t) y - 1)\end{aligned}$$

And

$$\begin{aligned}\overline{N} &= \mu N \\ &= \sec(t) (1) \\ &= \sec(t)\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\overline{M} + \overline{N} \frac{dy}{dt} &= 0 \\ (\tan(t) (\sec(t) y - 1)) + (\sec(t)) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \overline{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int \tan(t) (\sec(t) y - 1) dt \\ \phi &= \sec(t) y - \ln(\sec(t)) + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = \sec(t) + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \sec(t)$. Therefore equation (4) becomes

$$\sec(t) = \sec(t) + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \sec(t) y - \ln(\sec(t)) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \sec(t) y - \ln(\sec(t))$$

The solution becomes

$$y = \frac{\ln(\sec(t)) + c_1}{\sec(t)}$$

Initial conditions are used to solve for c_1 . Substituting $t = \pi$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = -i\pi - c_1$$

$$c_1 = -i\pi$$

Substituting c_1 found above in the general solution gives

$$y = -i \cos(t) \pi + \ln(\sec(t)) \cos(t)$$

Summary

The solution(s) found are the following

$$y = -i \cos(t) \pi + \ln(\sec(t)) \cos(t) \quad (1)$$

Verification of solutions

$$y = -i \cos(t) \pi + \ln(\sec(t)) \cos(t)$$

Verified OK.

3.3.5 Maple step by step solution

Let's solve

$$[\tan(t) y + y' = \sin(t), y(\pi) = 0]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\tan(t) y + \sin(t)$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$\tan(t) y + y' = \sin(t)$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) (\tan(t) y + y') = \mu(t) \sin(t)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t) y)$

$$\mu(t) (\tan(t) y + y') = \mu'(t) y + \mu(t) y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \mu(t) \tan(t)$$

- Solve to find the integrating factor

$$\mu(t) = \frac{1}{\cos(t)}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y) \right) dt = \int \mu(t) \sin(t) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t) y = \int \mu(t) \sin(t) dt + c_1$$

- Solve for y

$$y = \frac{\int \mu(t) \sin(t) dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = \frac{1}{\cos(t)}$

$$y = \cos(t) \left(\int \frac{\sin(t)}{\cos(t)} dt + c_1 \right)$$

- Evaluate the integrals on the rhs

$$y = \cos(t) (-\ln(\cos(t)) + c_1)$$

- Use initial condition $y(\pi) = 0$

$$0 = I\pi - c_1$$

- Solve for c_1

$$c_1 = I\pi$$

- Substitute $c_1 = I\pi$ into general solution and simplify

$$y = (-\ln(\cos(t)) + I\pi) \cos(t)$$

- Solution to the IVP

$$y = (-\ln(\cos(t)) + I\pi) \cos(t)$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 17

```
dsolve([tan(t)*y(t)+diff(y(t),t) = sin(t),y(Pi) = 0],y(t), singsol=all)
```

$$y(t) = (-\ln(\cos(t)) + i\pi) \cos(t)$$

✓ Solution by Mathematica

Time used: 0.073 (sec). Leaf size: 20

```
DSolve[{Tan[t]*y[t]+y'[t] == Sin[t],y[Pi]==0},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow i \cos(t)(\pi + i \log(\cos(t)))$$

3.4 problem 4

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Internal problem ID [519]

Internal file name [OUTPUT/519_Sunday_June_05_2022_01_42_56_AM_17528492/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 4.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$2yt + (-t^2 + 4)y' = 3t^2$$

With initial conditions

$$[y(-3) = 1]$$

3.4.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{2t}{t^2 - 4}$$
$$q(t) = -\frac{3t^2}{t^2 - 4}$$

Hence the ode is

$$y' - \frac{2ty}{t^2 - 4} = -\frac{3t^2}{t^2 - 4}$$

The domain of $p(t) = -\frac{2t}{t^2-4}$ is

$$\{-\infty \leq t < -2, -2 < t < 2, 2 < t \leq \infty\}$$

And the point $t_0 = -3$ is inside this domain. The domain of $q(t) = -\frac{3t^2}{t^2-4}$ is

$$\{-\infty \leq t < -2, -2 < t < 2, 2 < t \leq \infty\}$$

And the point $t_0 = -3$ is also inside this domain. Hence solution exists and is unique.

3.4.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{2t}{t^2-4} dt} \\ &= \frac{1}{t^2 - 4}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(-\frac{3t^2}{t^2 - 4} \right) \\ \frac{d}{dt} \left(\frac{y}{t^2 - 4} \right) &= \left(\frac{1}{t^2 - 4} \right) \left(-\frac{3t^2}{t^2 - 4} \right) \\ d \left(\frac{y}{t^2 - 4} \right) &= \left(-\frac{3t^2}{(t^2 - 4)^2} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y}{t^2 - 4} &= \int -\frac{3t^2}{(t^2 - 4)^2} dt \\ \frac{y}{t^2 - 4} &= \frac{3}{4(t-2)} - \frac{3 \ln(t-2)}{8} + \frac{3}{4(2+t)} + \frac{3 \ln(2+t)}{8} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = \frac{1}{t^2-4}$ results in

$$y = (t^2 - 4) \left(\frac{3}{4(t-2)} - \frac{3 \ln(t-2)}{8} + \frac{3}{4(2+t)} + \frac{3 \ln(2+t)}{8} \right) + c_1(t^2 - 4)$$

which simplifies to

$$y = \frac{3(t^2 - 4) \ln(2 + t)}{8} + c_1 t^2 - \frac{3 \ln(t - 2) t^2}{8} + \frac{3t}{2} - 4c_1 + \frac{3 \ln(t - 2)}{2}$$

Initial conditions are used to solve for c_1 . Substituting $t = -3$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{9}{2} + 5c_1 - \frac{15 \ln(5)}{8}$$

$$c_1 = \frac{11}{10} + \frac{3 \ln(5)}{8}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{3 \ln(2 + t) t^2}{8} - \frac{3 \ln(2 + t)}{2} + \frac{11t^2}{10} + \frac{3t^2 \ln(5)}{8} - \frac{3 \ln(t - 2) t^2}{8} + \frac{3t}{2} - \frac{22}{5} - \frac{3 \ln(5)}{2} + \frac{3 \ln(t - 2)}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{3 \ln(2 + t) t^2}{8} - \frac{3 \ln(2 + t)}{2} + \frac{11t^2}{10} + \frac{3t^2 \ln(5)}{8} - \frac{3 \ln(t - 2) t^2}{8} + \frac{3t}{2} - \frac{22}{5} - \frac{3 \ln(5)}{2} + \frac{3 \ln(t - 2)}{2} \quad (1)$$

Verification of solutions

$$y = \frac{3 \ln(2 + t) t^2}{8} - \frac{3 \ln(2 + t)}{2} + \frac{11t^2}{10} + \frac{3t^2 \ln(5)}{8} - \frac{3 \ln(t - 2) t^2}{8} + \frac{3t}{2} - \frac{22}{5} - \frac{3 \ln(5)}{2} + \frac{3 \ln(t - 2)}{2}$$

Verified OK.

3.4.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{t(-3t + 2y)}{t^2 - 4}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 191: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= t^2 - 4\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{t^2 - 4} dy \end{aligned}$$

Which results in

$$S = \frac{y}{t^2 - 4}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{t(-3t + 2y)}{t^2 - 4}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -\frac{2yt}{(t^2 - 4)^2} \\ S_y &= \frac{1}{t^2 - 4} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{3t^2}{(t^2 - 4)^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{3R^2}{(R^2 - 4)^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{3}{4(R-2)} - \frac{3 \ln(R-2)}{8} + \frac{3}{4(R+2)} + \frac{3 \ln(R+2)}{8} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$\frac{y}{t^2 - 4} = \frac{3}{4(t-2)} - \frac{3 \ln(t-2)}{8} + \frac{3}{4(2+t)} + \frac{3 \ln(2+t)}{8} + c_1$$

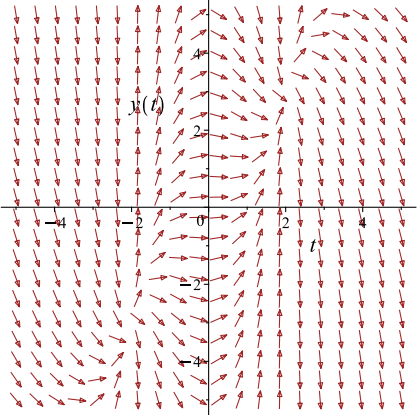
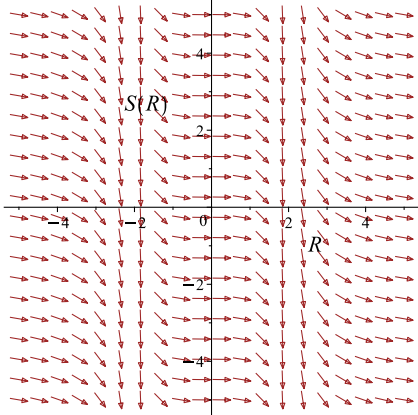
Which simplifies to

$$\frac{y}{t^2 - 4} = \frac{3}{4(t-2)} - \frac{3 \ln(t-2)}{8} + \frac{3}{4(2+t)} + \frac{3 \ln(2+t)}{8} + c_1$$

Which gives

$$y = -\frac{3 \ln(t-2)t^2}{8} + \frac{3 \ln(2+t)t^2}{8} + c_1 t^2 + \frac{3 \ln(t-2)}{2} - \frac{3 \ln(2+t)}{2} - 4c_1 + \frac{3t}{2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{t(-3t+2y)}{t^2-4}$ 	$R = t$ $S = \frac{y}{t^2 - 4}$	$\frac{dS}{dR} = -\frac{3R^2}{(R^2-4)^2}$ 

Initial conditions are used to solve for c_1 . Substituting $t = -3$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{9}{2} + 5c_1 - \frac{15 \ln(5)}{8}$$

$$c_1 = \frac{11}{10} + \frac{3 \ln(5)}{8}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{3 \ln(2+t)t^2}{8} - \frac{3 \ln(2+t)}{2} + \frac{11t^2}{10} + \frac{3t^2 \ln(5)}{8} - \frac{3 \ln(t-2)t^2}{8} + \frac{3t}{2} - \frac{22}{5} - \frac{3 \ln(5)}{2} + \frac{3 \ln(t-2)}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{3 \ln(2+t)t^2}{8} - \frac{3 \ln(2+t)}{2} + \frac{11t^2}{10} + \frac{3t^2 \ln(5)}{8} - \frac{3 \ln(t-2)t^2}{8} + \frac{3t}{2} - \frac{22}{5} - \frac{3 \ln(5)}{2} + \frac{3 \ln(t-2)}{2} \quad (1)$$

Verification of solutions

$$y = \frac{3 \ln(2+t)t^2}{8} - \frac{3 \ln(2+t)}{2} + \frac{11t^2}{10} + \frac{3t^2 \ln(5)}{8} - \frac{3 \ln(t-2)t^2}{8} + \frac{3t}{2} - \frac{22}{5} - \frac{3 \ln(5)}{2} + \frac{3 \ln(t-2)}{2}$$

Verified OK.

3.4.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(-t^2 + 4) dy &= (3t^2 - 2ty) dt \\ (-3t^2 + 2ty) dt + (-t^2 + 4) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -3t^2 + 2ty \\ N(t, y) &= -t^2 + 4\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-3t^2 + 2ty) \\ &= 2t\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(-t^2 + 4) \\ &= -2t\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= -\frac{1}{t^2 - 4} ((2t) - (-2t)) \\ &= -\frac{4t}{t^2 - 4} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int -\frac{4t}{t^2-4} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-2 \ln(t^2-4)} \\ &= \frac{1}{(t^2 - 4)^2} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{1}{(t^2 - 4)^2} (-3t^2 + 2ty) \\ &= \frac{-3t^2 + 2ty}{(t^2 - 4)^2} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \frac{1}{(t^2 - 4)^2} (-t^2 + 4) \\ &= -\frac{1}{t^2 - 4} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left(\frac{-3t^2 + 2ty}{(t^2 - 4)^2} \right) + \left(-\frac{1}{t^2 - 4} \right) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int \frac{-3t^2 + 2ty}{(t^2 - 4)^2} dt \\ \phi &= -\frac{3 \ln(t-2)}{8} + \frac{3-y}{4t-8} + \frac{3 \ln(2+t)}{8} + \frac{3+y}{8+4t} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned} \frac{\partial \phi}{\partial y} &= -\frac{1}{4t-8} + \frac{1}{8+4t} + f'(y) \\ &= -\frac{1}{t^2-4} + f'(y) \end{aligned} \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{1}{t^2-4}$. Therefore equation (4) becomes

$$-\frac{1}{t^2-4} = -\frac{1}{t^2-4} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{3 \ln(t-2)}{8} + \frac{3-y}{4t-8} + \frac{3 \ln(2+t)}{8} + \frac{3+y}{8+4t} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{3 \ln(t-2)}{8} + \frac{3-y}{4t-8} + \frac{3 \ln(2+t)}{8} + \frac{3+y}{8+4t}$$

The solution becomes

$$y = -\frac{3 \ln(t-2) t^2}{8} + \frac{3 \ln(2+t) t^2}{8} - c_1 t^2 + \frac{3 \ln(t-2)}{2} - \frac{3 \ln(2+t)}{2} + 4c_1 + \frac{3t}{2}$$

Initial conditions are used to solve for c_1 . Substituting $t = -3$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{15 \ln(5)}{8} - \frac{9}{2} - 5c_1$$

$$c_1 = -\frac{11}{10} - \frac{3 \ln(5)}{8}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{3 \ln(2+t) t^2}{8} - \frac{3 \ln(2+t)}{2} + \frac{11t^2}{10} + \frac{3t^2 \ln(5)}{8} - \frac{3 \ln(t-2) t^2}{8} + \frac{3t}{2} - \frac{22}{5} - \frac{3 \ln(5)}{2} + \frac{3 \ln(t-2)}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{3 \ln(2+t) t^2}{8} - \frac{3 \ln(2+t)}{2} + \frac{11t^2}{10} + \frac{3t^2 \ln(5)}{8} - \frac{3 \ln(t-2) t^2}{8} + \frac{3t}{2} - \frac{22}{5} - \frac{3 \ln(5)}{2} + \frac{3 \ln(t-2)}{2} \quad (1)$$

Verification of solutions

$$y = \frac{3 \ln(2+t) t^2}{8} - \frac{3 \ln(2+t)}{2} + \frac{11t^2}{10} + \frac{3t^2 \ln(5)}{8} - \frac{3 \ln(t-2) t^2}{8} + \frac{3t}{2} - \frac{22}{5} - \frac{3 \ln(5)}{2} + \frac{3 \ln(t-2)}{2}$$

Verified OK.

3.4.5 Maple step by step solution

Let's solve

$$[2ty + (-t^2 + 4)y' = 3t^2, y(-3) = 1]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = \frac{2ty}{t^2-4} - \frac{3t^2}{t^2-4}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' - \frac{2ty}{t^2-4} = -\frac{3t^2}{t^2-4}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' - \frac{2ty}{t^2-4} \right) = -\frac{3\mu(t)t^2}{t^2-4}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' - \frac{2ty}{t^2-4} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -\frac{2\mu(t)t}{t^2-4}$$

- Solve to find the integrating factor

$$\mu(t) = \frac{1}{(t-2)(2+t)}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int -\frac{3\mu(t)t^2}{t^2-4} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int -\frac{3\mu(t)t^2}{t^2-4} dt + c_1$$

- Solve for y

$$y = \frac{\int -\frac{3\mu(t)t^2}{t^2-4} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = \frac{1}{(t-2)(2+t)}$

$$y = (t-2)(2+t) \left(\int -\frac{3t^2}{(t-2)(2+t)(t^2-4)} dt + c_1 \right)$$

- Evaluate the integrals on the rhs

$$y = (t-2)(2+t) \left(\frac{3}{4(t-2)} - \frac{3\ln(t-2)}{8} + \frac{3}{4(2+t)} + \frac{3\ln(2+t)}{8} + c_1 \right)$$

- Simplify

$$y = \frac{3(t^2-4)\ln(2+t)}{8} + c_1 t^2 - \frac{3\ln(t-2)t^2}{8} + \frac{3t}{2} - 4c_1 + \frac{3\ln(t-2)}{2}$$

- Use initial condition $y(-3) = 1$

$$1 = -\frac{9}{2} + 5c_1 - \frac{15\ln(5)}{8}$$

- Solve for c_1

$$c_1 = \frac{11}{10} + \frac{3\ln(5)}{8}$$

- Substitute $c_1 = \frac{11}{10} + \frac{3\ln(5)}{8}$ into general solution and simplify

$$y = -\frac{22}{5} + \frac{3(t^2-4)\ln(2+t)}{8} + \frac{(44+15\ln(5))t^2}{40} - \frac{3\ln(t-2)t^2}{8} + \frac{3t}{2} - \frac{3\ln(5)}{2} + \frac{3\ln(t-2)}{2}$$

- Solution to the IVP

$$y = -\frac{22}{5} + \frac{3(t^2-4)\ln(2+t)}{8} + \frac{(44+15\ln(5))t^2}{40} - \frac{3\ln(t-2)t^2}{8} + \frac{3t}{2} - \frac{3\ln(5)}{2} + \frac{3\ln(t-2)}{2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.032 (sec). Leaf size: 42

```
dsolve([2*t*y(t)+(-t^2+4)*diff(y(t),t) = 3*t^2,y(-3) = 1],y(t), singsol=all)
```

$$y(t) = \frac{3t}{2} + \frac{3\ln(2+t)t^2}{8} - \frac{3\ln(2+t)}{2} - \frac{3\ln(t-2)t^2}{8} + \frac{3\ln(t-2)}{2} + \frac{11t^2}{10} - \frac{22}{5} + \frac{3\ln(5)t^2}{8} - \frac{3\ln(5)}{2}$$

✓ Solution by Mathematica

Time used: 0.049 (sec). Leaf size: 67

```
DSolve[{2*t*y[t]+(-t^2+4)*y'[t] == 3*t^2,y[-3]==1},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{40} \left(-15i\pi t^2 + 44t^2 + 15t^2 \log(5) - 15(t^2 - 4) \log(2 - t) + 15(t^2 - 4) \log(t + 2) \right. \\ \left. + 60t + 60i\pi - 176 - 60 \log(5) \right)$$

3.5 problem 5

3.5.1	Existence and uniqueness analysis	964
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Internal problem ID [520]

Internal file name [OUTPUT/520_Sunday_June_05_2022_01_42_58_AM_84257707/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 5.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**linear**", "**exactWithIntegrationFactor**", "**first_order_ode_lie_symmetry_lookup**"

Maple gives the following as the ode type

`[_linear]`

$$2yt + (-t^2 + 4)y' = 3t^2$$

With initial conditions

$$[y(1) = -3]$$

3.5.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{2t}{t^2 - 4}$$
$$q(t) = -\frac{3t^2}{t^2 - 4}$$

Hence the ode is

$$y' - \frac{2ty}{t^2 - 4} = -\frac{3t^2}{t^2 - 4}$$

The domain of $p(t) = -\frac{2t}{t^2-4}$ is

$$\{-\infty \leq t < -2, -2 < t < 2, 2 < t \leq \infty\}$$

And the point $t_0 = 1$ is inside this domain. The domain of $q(t) = -\frac{3t^2}{t^2-4}$ is

$$\{-\infty \leq t < -2, -2 < t < 2, 2 < t \leq \infty\}$$

And the point $t_0 = 1$ is also inside this domain. Hence solution exists and is unique.

3.5.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{2t}{t^2-4} dt} \\ &= \frac{1}{t^2 - 4}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(-\frac{3t^2}{t^2 - 4} \right) \\ \frac{d}{dt} \left(\frac{y}{t^2 - 4} \right) &= \left(\frac{1}{t^2 - 4} \right) \left(-\frac{3t^2}{t^2 - 4} \right) \\ d \left(\frac{y}{t^2 - 4} \right) &= \left(-\frac{3t^2}{(t^2 - 4)^2} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y}{t^2 - 4} &= \int -\frac{3t^2}{(t^2 - 4)^2} dt \\ \frac{y}{t^2 - 4} &= \frac{3}{4(t-2)} - \frac{3 \ln(t-2)}{8} + \frac{3}{4(2+t)} + \frac{3 \ln(2+t)}{8} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = \frac{1}{t^2-4}$ results in

$$y = (t^2 - 4) \left(\frac{3}{4(t-2)} - \frac{3 \ln(t-2)}{8} + \frac{3}{4(2+t)} + \frac{3 \ln(2+t)}{8} \right) + c_1(t^2 - 4)$$

which simplifies to

$$y = \frac{3(t^2 - 4) \ln(2 + t)}{8} + c_1 t^2 - \frac{3 \ln(t - 2) t^2}{8} + \frac{3t}{2} - 4c_1 + \frac{3 \ln(t - 2)}{2}$$

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = -3$ in the above solution gives an equation to solve for the constant of integration.

$$-3 = -\frac{9 \ln(3)}{8} - 3c_1 + \frac{9i\pi}{8} + \frac{3}{2}$$

$$c_1 = -\frac{3 \ln(3)}{8} + \frac{3i\pi}{8} + \frac{3}{2}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{3 \ln(2 + t) t^2}{8} - \frac{3 \ln(2 + t)}{2} - \frac{3t^2 \ln(3)}{8} + \frac{3it^2\pi}{8} + \frac{3t^2}{2} - \frac{3 \ln(t - 2) t^2}{8} + \frac{3t}{2} + \frac{3 \ln(3)}{2} - \frac{3i\pi}{2} - 6 + \frac{3 \ln(t - 2)}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{3 \ln(2 + t) t^2}{8} - \frac{3 \ln(2 + t)}{2} - \frac{3t^2 \ln(3)}{8} + \frac{3it^2\pi}{8} + \frac{3t^2}{2} - \frac{3 \ln(t - 2) t^2}{8} + \frac{3t}{2} + \frac{3 \ln(3)}{2} - \frac{3i\pi}{2} - 6 + \frac{3 \ln(t - 2)}{2} \quad (1)$$

Verification of solutions

$$y = \frac{3 \ln(2 + t) t^2}{8} - \frac{3 \ln(2 + t)}{2} - \frac{3t^2 \ln(3)}{8} + \frac{3it^2\pi}{8} + \frac{3t^2}{2} - \frac{3 \ln(t - 2) t^2}{8} + \frac{3t}{2} + \frac{3 \ln(3)}{2} - \frac{3i\pi}{2} - 6 + \frac{3 \ln(t - 2)}{2}$$

Verified OK.

3.5.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{t(-3t + 2y)}{t^2 - 4}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (A)$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 194: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= t^2 - 4\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{t^2 - 4} dy \end{aligned}$$

Which results in

$$S = \frac{y}{t^2 - 4}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{t(-3t + 2y)}{t^2 - 4}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -\frac{2yt}{(t^2 - 4)^2} \\ S_y &= \frac{1}{t^2 - 4} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{3t^2}{(t^2 - 4)^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{3R^2}{(R^2 - 4)^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{3}{4(R-2)} - \frac{3 \ln(R-2)}{8} + \frac{3}{4(R+2)} + \frac{3 \ln(R+2)}{8} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$\frac{y}{t^2 - 4} = \frac{3}{4(t-2)} - \frac{3 \ln(t-2)}{8} + \frac{3}{4(2+t)} + \frac{3 \ln(2+t)}{8} + c_1$$

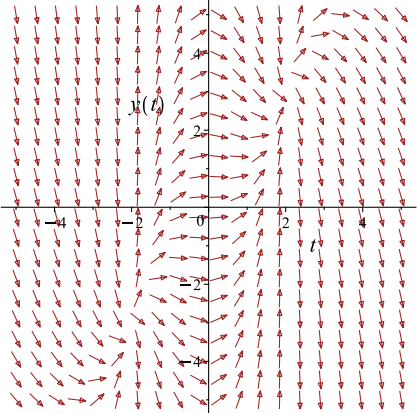
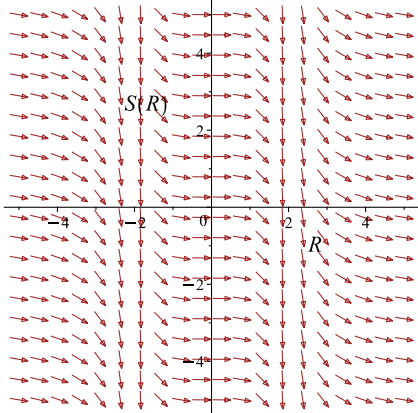
Which simplifies to

$$\frac{y}{t^2 - 4} = \frac{3}{4(t-2)} - \frac{3 \ln(t-2)}{8} + \frac{3}{4(2+t)} + \frac{3 \ln(2+t)}{8} + c_1$$

Which gives

$$y = -\frac{3 \ln(t-2)t^2}{8} + \frac{3 \ln(2+t)t^2}{8} + c_1 t^2 + \frac{3 \ln(t-2)}{2} - \frac{3 \ln(2+t)}{2} - 4c_1 + \frac{3t}{2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{t(-3t+2y)}{t^2-4}$ 	$R = t$ $S = \frac{y}{t^2 - 4}$	$\frac{dS}{dR} = -\frac{3R^2}{(R^2-4)^2}$ 

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = -3$ in the above solution gives an equation to solve for the constant of integration.

$$-3 = -\frac{9 \ln(3)}{8} - 3c_1 + \frac{9i\pi}{8} + \frac{3}{2}$$

$$c_1 = -\frac{3 \ln(3)}{8} + \frac{3i\pi}{8} + \frac{3}{2}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{3 \ln(2+t)t^2}{8} - \frac{3 \ln(2+t)}{2} - \frac{3t^2 \ln(3)}{8} + \frac{3it^2\pi}{8} + \frac{3t^2}{2} - \frac{3 \ln(t-2)t^2}{8} + \frac{3t}{2} + \frac{3 \ln(3)}{2} - \frac{3i\pi}{2} - 6 + \frac{3}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{3 \ln(2+t)t^2}{8} - \frac{3 \ln(2+t)}{2} - \frac{3t^2 \ln(3)}{8} + \frac{3it^2\pi}{8} + \frac{3t^2}{2} - \frac{3 \ln(t-2)t^2}{8} + \frac{3t}{2} + \frac{3 \ln(3)}{2} - \frac{3i\pi}{2} - 6 + \frac{3 \ln(t-2)}{2} \quad (1)$$

Verification of solutions

$$y = \frac{3 \ln(2+t)t^2}{8} - \frac{3 \ln(2+t)}{2} - \frac{3t^2 \ln(3)}{8} + \frac{3it^2\pi}{8} + \frac{3t^2}{2} - \frac{3 \ln(t-2)t^2}{8} + \frac{3t}{2} + \frac{3 \ln(3)}{2} - \frac{3i\pi}{2} - 6 + \frac{3 \ln(t-2)}{2}$$

Verified OK.

3.5.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (B)$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(-t^2 + 4) dy &= (3t^2 - 2ty) dt \\ (-3t^2 + 2ty) dt + (-t^2 + 4) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -3t^2 + 2ty \\ N(t, y) &= -t^2 + 4\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-3t^2 + 2ty) \\ &= 2t\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(-t^2 + 4) \\ &= -2t\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= -\frac{1}{t^2 - 4} ((2t) - (-2t)) \\ &= -\frac{4t}{t^2 - 4} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int -\frac{4t}{t^2-4} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-2 \ln(t^2-4)} \\ &= \frac{1}{(t^2 - 4)^2} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{1}{(t^2 - 4)^2} (-3t^2 + 2ty) \\ &= \frac{-3t^2 + 2ty}{(t^2 - 4)^2} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \frac{1}{(t^2 - 4)^2} (-t^2 + 4) \\ &= -\frac{1}{t^2 - 4} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left(\frac{-3t^2 + 2ty}{(t^2 - 4)^2} \right) + \left(-\frac{1}{t^2 - 4} \right) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int \frac{-3t^2 + 2ty}{(t^2 - 4)^2} dt \\ \phi &= -\frac{3 \ln(t-2)}{8} + \frac{3-y}{4t-8} + \frac{3 \ln(2+t)}{8} + \frac{3+y}{8+4t} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned} \frac{\partial \phi}{\partial y} &= -\frac{1}{4t-8} + \frac{1}{8+4t} + f'(y) \\ &= -\frac{1}{t^2-4} + f'(y) \end{aligned} \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{1}{t^2-4}$. Therefore equation (4) becomes

$$-\frac{1}{t^2-4} = -\frac{1}{t^2-4} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{3 \ln(t-2)}{8} + \frac{3-y}{4t-8} + \frac{3 \ln(2+t)}{8} + \frac{3+y}{8+4t} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{3 \ln(t-2)}{8} + \frac{3-y}{4t-8} + \frac{3 \ln(2+t)}{8} + \frac{3+y}{8+4t}$$

The solution becomes

$$y = -\frac{3 \ln(t-2) t^2}{8} + \frac{3 \ln(2+t) t^2}{8} - c_1 t^2 + \frac{3 \ln(t-2)}{2} - \frac{3 \ln(2+t)}{2} + 4c_1 + \frac{3t}{2}$$

Initial conditions are used to solve for c_1 . Substituting $t = 1$ and $y = -3$ in the above solution gives an equation to solve for the constant of integration.

$$-3 = \frac{9i\pi}{8} - \frac{9 \ln(3)}{8} + 3c_1 + \frac{3}{2}$$

$$c_1 = \frac{3 \ln(3)}{8} - \frac{3i\pi}{8} - \frac{3}{2}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{3 \ln(2+t) t^2}{8} - \frac{3 \ln(2+t)}{2} - \frac{3t^2 \ln(3)}{8} + \frac{3it^2\pi}{8} + \frac{3t^2}{2} - \frac{3 \ln(t-2) t^2}{8} + \frac{3t}{2} + \frac{3 \ln(3)}{2} - \frac{3i\pi}{2} - 6 + \frac{3 \ln(t-2)}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{3 \ln(2+t) t^2}{8} - \frac{3 \ln(2+t)}{2} - \frac{3t^2 \ln(3)}{8} + \frac{3it^2\pi}{8} + \frac{3t^2}{2} - \frac{3 \ln(t-2) t^2}{8} + \frac{3t}{2} + \frac{3 \ln(3)}{2} - \frac{3i\pi}{2} - 6 + \frac{3 \ln(t-2)}{2} \quad (1)$$

Verification of solutions

$$y = \frac{3 \ln(2+t) t^2}{8} - \frac{3 \ln(2+t)}{2} - \frac{3t^2 \ln(3)}{8} + \frac{3it^2\pi}{8} + \frac{3t^2}{2} - \frac{3 \ln(t-2) t^2}{8} + \frac{3t}{2} + \frac{3 \ln(3)}{2} - \frac{3i\pi}{2} - 6 + \frac{3 \ln(t-2)}{2}$$

Verified OK.

3.5.5 Maple step by step solution

Let's solve

$$[2yt + (-t^2 + 4)y' = 3t^2, y(1) = -3]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = \frac{2ty}{t^2-4} - \frac{3t^2}{t^2-4}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' - \frac{2ty}{t^2-4} = -\frac{3t^2}{t^2-4}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' - \frac{2ty}{t^2-4} \right) = -\frac{3\mu(t)t^2}{t^2-4}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' - \frac{2ty}{t^2-4} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = -\frac{2\mu(t)t}{t^2-4}$$

- Solve to find the integrating factor

$$\mu(t) = \frac{1}{(t-2)(2+t)}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int -\frac{3\mu(t)t^2}{t^2-4} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int -\frac{3\mu(t)t^2}{t^2-4} dt + c_1$$

- Solve for y

$$y = \frac{\int -\frac{3\mu(t)t^2}{t^2-4} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = \frac{1}{(t-2)(2+t)}$

$$y = (t-2)(2+t) \left(\int -\frac{3t^2}{(t-2)(2+t)(t^2-4)} dt + c_1 \right)$$

- Evaluate the integrals on the rhs

$$y = (t-2)(2+t) \left(\frac{3}{4(t-2)} - \frac{3\ln(t-2)}{8} + \frac{3}{4(2+t)} + \frac{3\ln(2+t)}{8} + c_1 \right)$$

- Simplify

$$y = \frac{3(t^2-4)\ln(2+t)}{8} + c_1 t^2 - \frac{3\ln(t-2)t^2}{8} + \frac{3t}{2} - 4c_1 + \frac{3\ln(t-2)}{2}$$

- Use initial condition $y(1) = -3$

$$-3 = -\frac{9\ln(3)}{8} - 3c_1 + \frac{9\pi}{8} + \frac{3}{2}$$

- Solve for c_1

$$c_1 = -\frac{3\ln(3)}{8} + \frac{3\pi}{8} + \frac{3}{2}$$

- Substitute $c_1 = -\frac{3\ln(3)}{8} + \frac{3\pi}{8} + \frac{3}{2}$ into general solution and simplify

$$y = -6 + \frac{3(t^2-4)\ln(2+t)}{8} + \frac{3\pi t^2}{8} - \frac{3t^2\ln(3)}{8} - \frac{3\ln(t-2)t^2}{8} - \frac{3\pi}{2} + \frac{3t^2}{2} + \frac{3t}{2} + \frac{3\ln(3)}{2} + \frac{3\ln(t-2)}{2}$$

- Solution to the IVP

$$y = -6 + \frac{3(t^2-4)\ln(2+t)}{8} + \frac{3\pi t^2}{8} - \frac{3t^2\ln(3)}{8} - \frac{3\ln(t-2)t^2}{8} - \frac{3\pi}{2} + \frac{3t^2}{2} + \frac{3t}{2} + \frac{3\ln(3)}{2} + \frac{3\ln(t-2)}{2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 46

```
dsolve([2*t*y(t)+(-t^2+4)*diff(y(t),t) = 3*t^2,y(1) = -3],y(t), singsol=all)
```

$$y(t) = -6 + \frac{3(t^2-4)\ln(2+t)}{8} + \frac{3\pi t^2}{8} - \frac{3\ln(3)t^2}{8} - \frac{3\ln(t-2)t^2}{8} - \frac{3\pi}{2} + \frac{3t^2}{2} + \frac{3t}{2} + \frac{3\ln(3)}{2} + \frac{3\ln(t-2)}{2}$$

✓ Solution by Mathematica

Time used: 0.043 (sec). Leaf size: 52

```
DSolve[{2*t*y[t]+(-t^2+4)*y'[t] == 3*t^2,y[1]==-3},y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -\frac{3}{8}(-4t^2 + t^2 \log(3) + (t^2 - 4) \log(2 - t) - (t^2 - 4) \log(t + 2) - 4t + 16 - 4 \log(3))$$

3.6 problem 6

3.6.1	Solving as linear ode	978
3.6.2	Solving as first order ode lie symmetry lookup ode	980
3.6.3	Solving as exact ode	984
3.6.4	Maple step by step solution	989

Internal problem ID [521]

Internal file name [OUTPUT/521_Sunday_June_05_2022_01_42_59_AM_5852861/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 6.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$y + \ln(t) y' = \cot(t)$$

3.6.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{1}{\ln(t)}$$
$$q(t) = \frac{\cot(t)}{\ln(t)}$$

Hence the ode is

$$y' + \frac{y}{\ln(t)} = \frac{\cot(t)}{\ln(t)}$$

The integrating factor μ is

$$\mu = e^{\int \frac{1}{\ln(t)} dt}$$

The ode becomes

$$\begin{aligned} \frac{d}{dt}(\mu y) &= (\mu) \left(\frac{\cot(t)}{\ln(t)} \right) \\ \frac{d}{dt} \left(e^{\int \frac{1}{\ln(t)} dt} y \right) &= \left(e^{\int \frac{1}{\ln(t)} dt} \right) \left(\frac{\cot(t)}{\ln(t)} \right) \\ d \left(e^{\int \frac{1}{\ln(t)} dt} y \right) &= \left(\frac{\cot(t) e^{\int \frac{1}{\ln(t)} dt}}{\ln(t)} \right) dt \end{aligned}$$

Integrating gives

$$\begin{aligned} e^{\int \frac{1}{\ln(t)} dt} y &= \int \frac{\cot(t) e^{\int \frac{1}{\ln(t)} dt}}{\ln(t)} dt \\ e^{\int \frac{1}{\ln(t)} dt} y &= \int \frac{\cot(t) e^{\int \frac{1}{\ln(t)} dt}}{\ln(t)} dt + c_1 \end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{\int \frac{1}{\ln(t)} dt}$ results in

$$y = e^{-\left(\int \frac{1}{\ln(t)} dt\right)} \left(\int \frac{\cot(t) e^{\int \frac{1}{\ln(t)} dt}}{\ln(t)} dt \right) + c_1 e^{-\left(\int \frac{1}{\ln(t)} dt\right)}$$

which simplifies to

$$y = e^{-\left(\int \frac{1}{\ln(t)} dt\right)} \left(\int \frac{\cot(t) e^{\int \frac{1}{\ln(t)} dt}}{\ln(t)} dt + c_1 \right)$$

Which can be simplified to become

$$y = e^{\int -\frac{1}{\ln(t)} dt} \left(\int \frac{\cot(t) e^{\int \frac{1}{\ln(t)} dt}}{\ln(t)} dt + c_1 \right)$$

Summary

The solution(s) found are the following

$$y = e^{\int -\frac{1}{\ln(t)} dt} \left(\int \frac{\cot(t) e^{\int \frac{1}{\ln(t)} dt}}{\ln(t)} dt + c_1 \right) \quad (1)$$

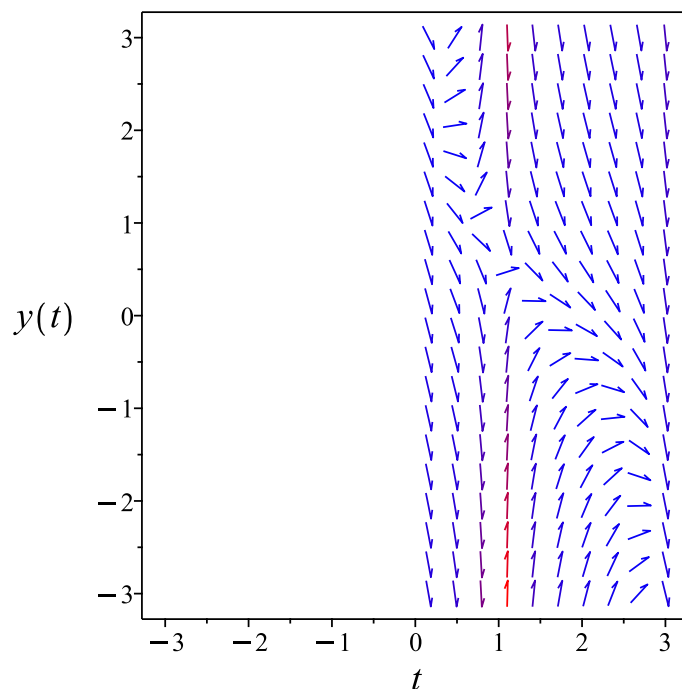


Figure 179: Slope field plot

Verification of solutions

$$y = e^{\int -\frac{1}{\ln(t)} dt} \left(\int \frac{\cot(t) e^{\int \frac{1}{\ln(t)} dt}}{\ln(t)} dt + c_1 \right)$$

Verified OK.

3.6.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-y + \cot(t)}{\ln(t)}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 197: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{\text{expIntegral}_1(-\ln(t))}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{\text{expIntegral}_1(-\ln(t))}} dy \end{aligned}$$

Which results in

$$S = e^{-\text{expIntegral}_1(-\ln(t))} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{-y + \cot(t)}{\ln(t)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= \frac{e^{-\text{expIntegral}_1(-\ln(t))} y}{\ln(t)} \\ S_y &= e^{-\text{expIntegral}_1(-\ln(t))} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{e^{-\text{expIntegral}_1(-\ln(t))} \cot(t)}{\ln(t)} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{e^{-\text{expIntegral}_1(-\ln(R))} \cot(R)}{\ln(R)}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \int \frac{e^{-\exp\text{Integral}_1(-\ln(R))} \cot(R)}{\ln(R)} dR + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$e^{-\exp\text{Integral}_1(-\ln(t))} y = \int \frac{e^{-\exp\text{Integral}_1(-\ln(t))} \cot(t)}{\ln(t)} dt + c_1$$

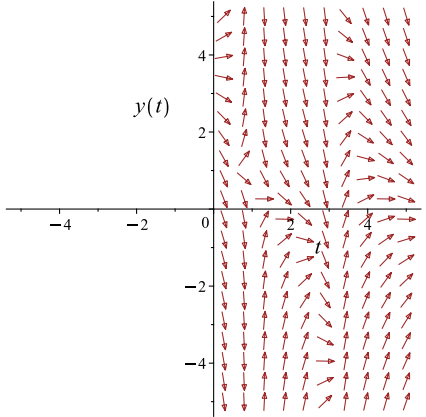
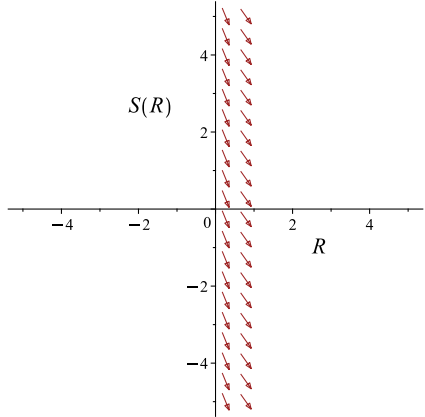
Which simplifies to

$$e^{-\exp\text{Integral}_1(-\ln(t))} y = \int \frac{e^{-\exp\text{Integral}_1(-\ln(t))} \cot(t)}{\ln(t)} dt + c_1$$

Which gives

$$y = \left(\int \frac{e^{-\exp\text{Integral}_1(-\ln(t))} \cot(t)}{\ln(t)} dt + c_1 \right) e^{\exp\text{Integral}_1(-\ln(t))}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{-y + \cot(t)}{\ln(t)}$ 	$R = t$ $S = e^{-\exp\text{Integral}_1(-\ln(t))} y$	$\frac{dS}{dR} = \frac{e^{-\exp\text{Integral}_1(-\ln(R))} \cot(R)}{\ln(R)}$ 

Summary

The solution(s) found are the following

$$y = \left(\int \frac{e^{-\exp\text{Integral}_1(-\ln(t))} \cot(t)}{\ln(t)} dt + c_1 \right) e^{\exp\text{Integral}_1(-\ln(t))} \quad (1)$$

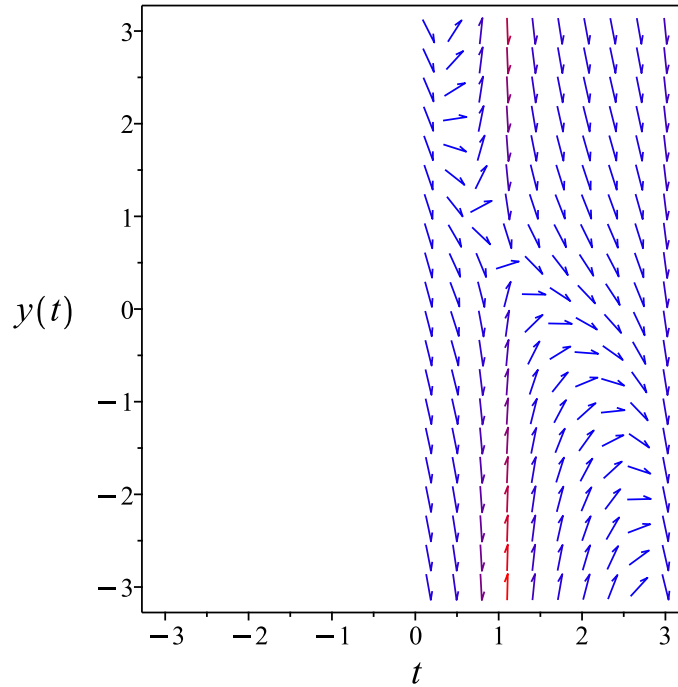


Figure 180: Slope field plot

Verification of solutions

$$y = \left(\int \frac{e^{-\exp\text{Integral}_1(-\ln(t))} \cot(t)}{\ln(t)} dt + c_1 \right) e^{\exp\text{Integral}_1(-\ln(t))}$$

Verified OK.

3.6.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the

ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (\ln(t)) dy &= (-y + \cot(t)) dt \\ (y - \cot(t)) dt + (\ln(t)) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, y) &= y - \cot(t) \\ N(t, y) &= \ln(t) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y - \cot(t)) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(\ln(t)) \\ &= \frac{1}{t}\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{\ln(t)} \left((1) - \left(\frac{1}{t} \right) \right) \\ &= \frac{-1 + t}{t \ln(t)}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int \frac{-1+t}{t \ln(t)} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\text{expIntegral}_1(-\ln(t)) - \ln(\ln(t))} \\ &= \frac{e^{-\text{expIntegral}_1(-\ln(t))}}{\ln(t)}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\overline{M} &= \mu M \\ &= \frac{e^{-\text{expIntegral}_1(-\ln(t))}}{\ln(t)}(y - \cot(t)) \\ &= \frac{(y - \cot(t)) e^{-\text{expIntegral}_1(-\ln(t))}}{\ln(t)}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{e^{-\exp\text{Integral}_1(-\ln(t))}}{\ln(t)} (\ln(t)) \\ &= e^{-\exp\text{Integral}_1(-\ln(t))}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ \left(\frac{(y - \cot(t)) e^{-\exp\text{Integral}_1(-\ln(t))}}{\ln(t)} \right) + (e^{-\exp\text{Integral}_1(-\ln(t))}) \frac{dy}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int \frac{(y - \cot(t)) e^{-\exp\text{Integral}_1(-\ln(t))}}{\ln(t)} dt \\ \phi &= \int^t \frac{(y - \cot(_a)) e^{-\exp\text{Integral}_1(-\ln(_a))}}{\ln(_a)} d_a + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-\exp\text{Integral}_1(-\ln(t))} + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-\exp\text{Integral}_1(-\ln(t))}$. Therefore equation (4) becomes

$$e^{-\exp\text{Integral}_1(-\ln(t))} = e^{-\exp\text{Integral}_1(-\ln(t))} + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \int^t \frac{(y - \cot(_a)) e^{-\expIntegral_1(-\ln(_a))}}{\ln(_a)} d_a + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \int^t \frac{(y - \cot(_a)) e^{-\expIntegral_1(-\ln(_a))}}{\ln(_a)} d_a$$

Summary

The solution(s) found are the following

$$\int^t \frac{(y - \cot(_a)) e^{-\expIntegral_1(-\ln(_a))}}{\ln(_a)} d_a = c_1 \quad (1)$$

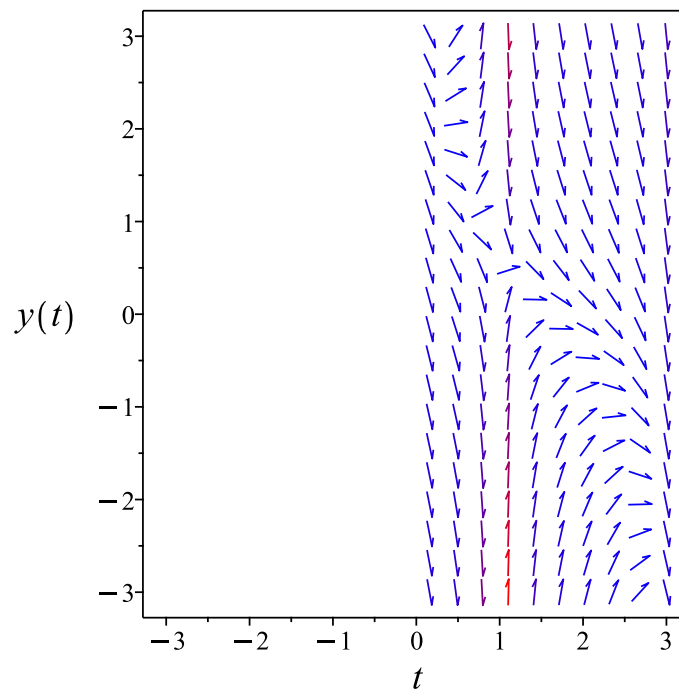


Figure 181: Slope field plot

Verification of solutions

$$\int^t \frac{(y - \cot(\ln a)) e^{-\exp \int_1^t (-\ln(\ln a))}}{\ln(\ln a)} d \ln a = c_1$$

Verified OK.

3.6.4 Maple step by step solution

Let's solve

$$y + \ln(t) y' = \cot(t)$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{y}{\ln(t)} + \frac{\cot(t)}{\ln(t)}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{y}{\ln(t)} = \frac{\cot(t)}{\ln(t)}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{y}{\ln(t)} \right) = \frac{\mu(t) \cot(t)}{\ln(t)}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t) y)$

$$\mu(t) \left(y' + \frac{y}{\ln(t)} \right) = \mu'(t) y + \mu(t) y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)}{\ln(t)}$$

- Solve to find the integrating factor

$$\mu(t) = e^{-\text{Ei}_1(-\ln(t))}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t) y) \right) dt = \int \frac{\mu(t) \cot(t)}{\ln(t)} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t) y = \int \frac{\mu(t) \cot(t)}{\ln(t)} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t) \cot(t)}{\ln(t)} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = e^{-\text{Ei}_1(-\ln(t))}$

$$y = \frac{\int \frac{e^{-\text{Ei}_1(-\ln(t))} \cot(t)}{\ln(t)} dt + c_1}{e^{-\text{Ei}_1(-\ln(t))}}$$

- Simplify

$$y = \left(\int \frac{e^{-\text{Ei}_1(-\ln(t))} \cot(t)}{\ln(t)} dt + c_1 \right) e^{\text{Ei}_1(-\ln(t))}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 32

```
dsolve(y(t)+ln(t)*diff(y(t),t) = cot(t),y(t), singsol=all)
```

$$y(t) = \left(\int \frac{\cot(t) e^{-\text{expIntegral}_1(-\ln(t))}}{\ln(t)} dt + c_1 \right) e^{\text{expIntegral}_1(-\ln(t))}$$

✓ Solution by Mathematica

Time used: 0.142 (sec). Leaf size: 36

```
DSolve[y[t]+Log[t]*y'[t] == Cot[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{-\text{LogIntegral}(t)} \left(\int_1^t \frac{e^{\text{LogIntegral}(K[1])} \cot(K[1])}{\log(K[1])} dK[1] + c_1 \right)$$

3.7 problem 11

3.7.1	Solving as separable ode	991
3.7.2	Solving as differentialType ode	996
3.7.3	Solving as first order ode lie symmetry lookup ode	1000
3.7.4	Solving as exact ode	1004
3.7.5	Maple step by step solution	1008

Internal problem ID [522]

Internal file name [OUTPUT/522_Sunday_June_05_2022_01_43_00_AM_26088981/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 11.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{t^2 + 1}{3y - y^2} = 0$$

3.7.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(t, y) \\ &= f(t)g(y) \\ &= \frac{-t^2 - 1}{y(-3 + y)}\end{aligned}$$

Where $f(t) = -t^2 - 1$ and $g(y) = \frac{1}{y(-3+y)}$. Integrating both sides gives

$$\frac{1}{\frac{1}{y(-3+y)}} dy = -t^2 - 1 dt$$

$$\int \frac{1}{\frac{1}{y(-3+y)}} dy = \int -t^2 - 1 dt$$

$$\frac{1}{3}y^3 - \frac{3}{2}y^2 = -\frac{1}{3}t^3 - t + c_1$$

Which results in

$$y = \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} + \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2}$$

$$y = \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{\frac{4}{9}} - \frac{4\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2} + i\sqrt{3} \left(\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} - \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3}\right)^{\frac{1}{3}}}{2} \right)$$

$$y = \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{\frac{4}{9}} - \frac{4\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2} + i\sqrt{3} \left(\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} - \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3}\right)^{\frac{1}{3}}}{2} \right)$$

Summary

The solution(s) found are the following

$$y = \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} \quad (1)$$

$$+ \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2}$$

$$y = \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{4} \quad (2)$$

$$- \frac{4\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2}$$

$$+ i\sqrt{3} \left(\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} - \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3}\right)^{\frac{1}{3}}}{2} \right)$$

$$y = \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{4} \quad (3)$$

$$- \frac{4\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2}$$

$$+ i\sqrt{3} \left(\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} - \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3}\right)^{\frac{1}{3}}}{2} \right)$$

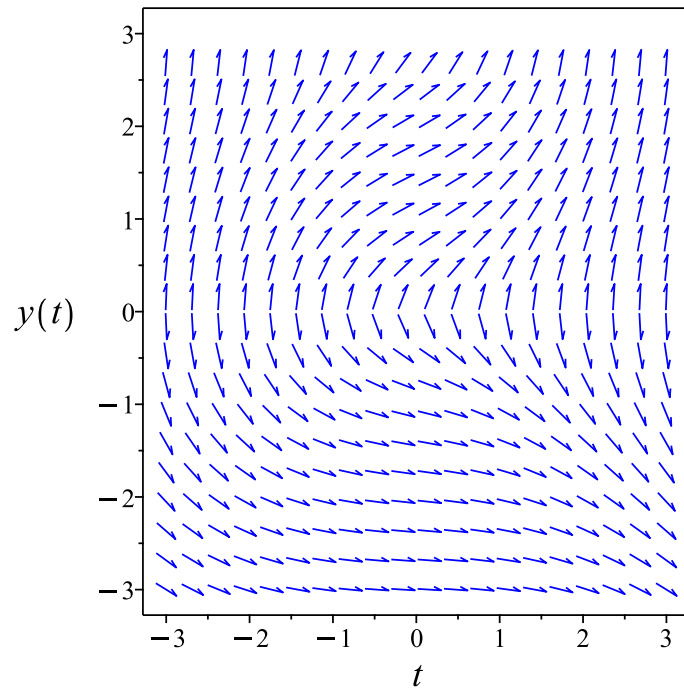


Figure 182: Slope field plot

Verification of solutions

y

$$= \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} + \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2}$$

Verified OK.

$y =$

$$- \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{4} - \frac{4\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2} + i\sqrt{3} \left(\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} - \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} \right)$$

Verified OK.

$y =$

$$- \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{4} - \frac{4\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2} + i\sqrt{3} \left(\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} - \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} \right)$$

Verified OK.

3.7.2 Solving as differentialType ode

Writing the ode as

$$y' = \frac{t^2 + 1}{3y - y^2} \quad (1)$$

Which becomes

$$(y^2 - 3y) dy = (-t^2 - 1) dt \quad (2)$$

But the RHS is complete differential because

$$(-t^2 - 1) dt = d\left(-\frac{1}{3}t^3 - t\right)$$

Hence (2) becomes

$$(y^2 - 3y) dy = d\left(-\frac{1}{3}t^3 - t\right)$$

Integrating both sides gives gives these solutions

$$y = \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} + \frac{t}{2}$$

$$y = -\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{4}$$

$$y = -\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{4}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} \quad (1)$$

$$+ \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2} + c_1$$

$$y = \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{4} \quad (2)$$

$$+ \frac{4\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2}$$

$$+ i\sqrt{3} \left(\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} - \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3}\right)^{\frac{1}{3}}}{2} \right) + c_1$$

$$y = \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{4} \quad (3)$$

$$+ \frac{4\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2}$$

$$+ i\sqrt{3} \left(\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} - \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3}\right)^{\frac{1}{3}}}{2} \right) + c_1$$

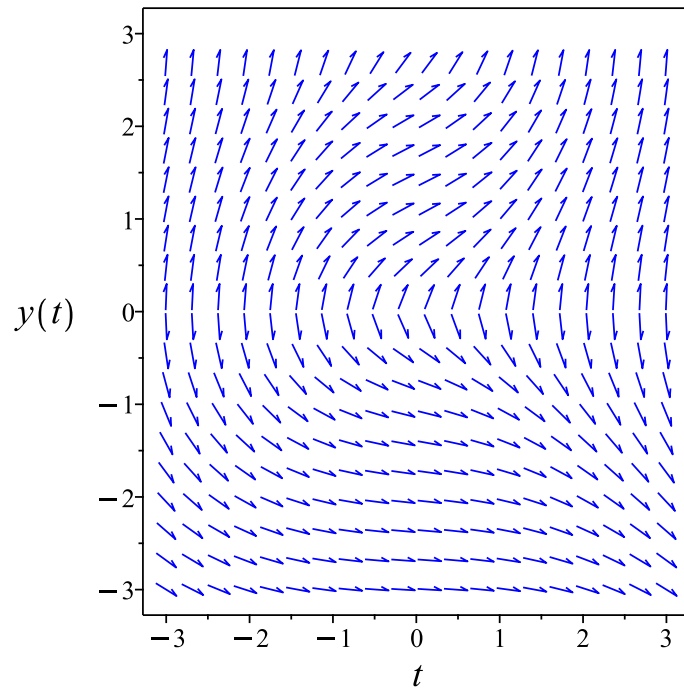


Figure 183: Slope field plot

Verification of solutions

y

$$= \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} + \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2} + c_1$$

Verified OK.

$y =$

$$- \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{4} - \frac{4\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2} + i\sqrt{3} \left(\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} - \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3}\right)^{\frac{1}{3}}}{2} \right) + c_1$$

Verified OK.

$y =$

$$- \frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{4} - \frac{4\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{9} + \frac{3}{2} + i\sqrt{3} \left(\frac{\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 - 72c_1t + 36t^2 + 162c_1 - 162t}\right)^{\frac{1}{3}}}{2} - \frac{2\left(27 - 4t^3 + 12c_1 - 12t + 2\sqrt{4t^6 - 24c_1t^3}\right)^{\frac{1}{3}}}{2} \right) + c_1$$

Verified OK.

3.7.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{t^2 + 1}{y(-3 + y)}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 200: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= \frac{1}{-t^2 - 1} \\ \eta(t, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dt \\ &= \int \frac{1}{\frac{1}{-t^2-1}} dt\end{aligned}$$

Which results in

$$S = -\frac{1}{3}t^3 - t$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{t^2 + 1}{y(-3 + y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 0 \\R_y &= 1 \\S_t &= -t^2 - 1 \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y(-3 + y) \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R(-3 + R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{1}{3}R^3 - \frac{3}{2}R^2 + c_1 \quad (4)$$

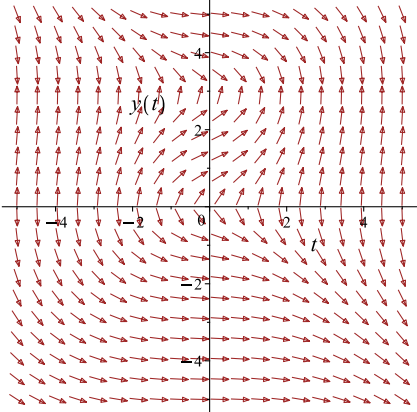
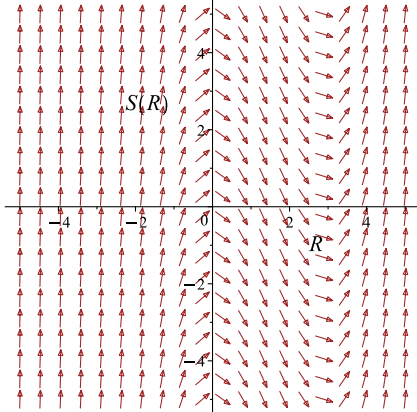
To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$-\frac{1}{3}t^3 - t = \frac{y^3}{3} - \frac{3y^2}{2} + c_1$$

Which simplifies to

$$-\frac{1}{3}t^3 - t = \frac{y^3}{3} - \frac{3y^2}{2} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{t^2+1}{y(-3+y)}$ 	$R = y$ $S = -\frac{1}{3}t^3 - t$	$\frac{dS}{dR} = R(-3 + R)$ 

Summary

The solution(s) found are the following

$$-\frac{1}{3}t^3 - t = \frac{y^3}{3} - \frac{3y^2}{2} + c_1 \quad (1)$$

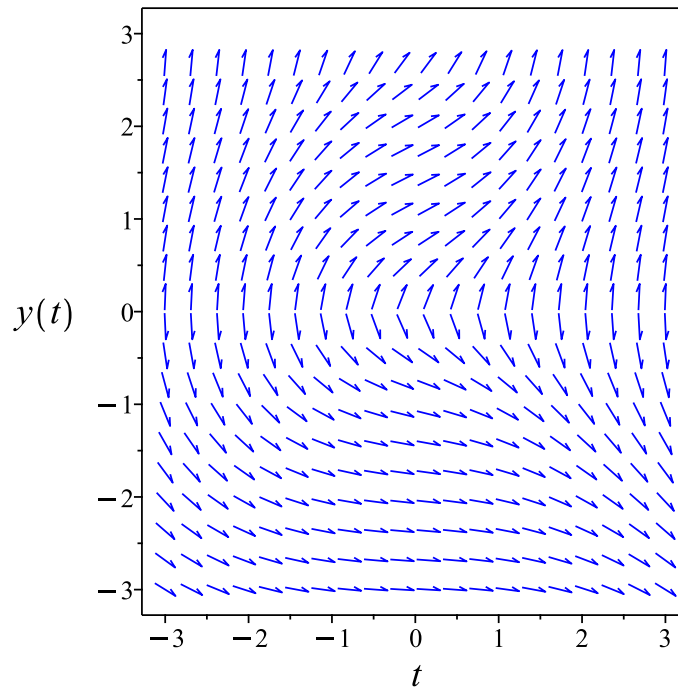


Figure 184: Slope field plot

Verification of solutions

$$-\frac{1}{3}t^3 - t = \frac{y^3}{3} - \frac{3y^2}{2} + c_1$$

Verified OK.

3.7.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(-y(-3 + y)) dy &= (t^2 + 1) dt \\ (-t^2 - 1) dt + (-y(-3 + y)) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -t^2 - 1 \\ N(t, y) &= -y(-3 + y)\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-t^2 - 1) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(-y(-3 + y)) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -t^2 - 1 dt \\ \phi &= -\frac{1}{3}t^3 - t + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -y(-3 + y)$. Therefore equation (4) becomes

$$-y(-3 + y) = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -y(-3 + y)$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int (-y(-3 + y)) dy \\ f(y) &= -\frac{1}{3}y^3 + \frac{3}{2}y^2 + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{1}{3}t^3 - t - \frac{1}{3}y^3 + \frac{3}{2}y^2 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{1}{3}t^3 - t - \frac{1}{3}y^3 + \frac{3}{2}y^2$$

Summary

The solution(s) found are the following

$$-\frac{t^3}{3} - \frac{y^3}{3} + \frac{3y^2}{2} - t = c_1 \quad (1)$$

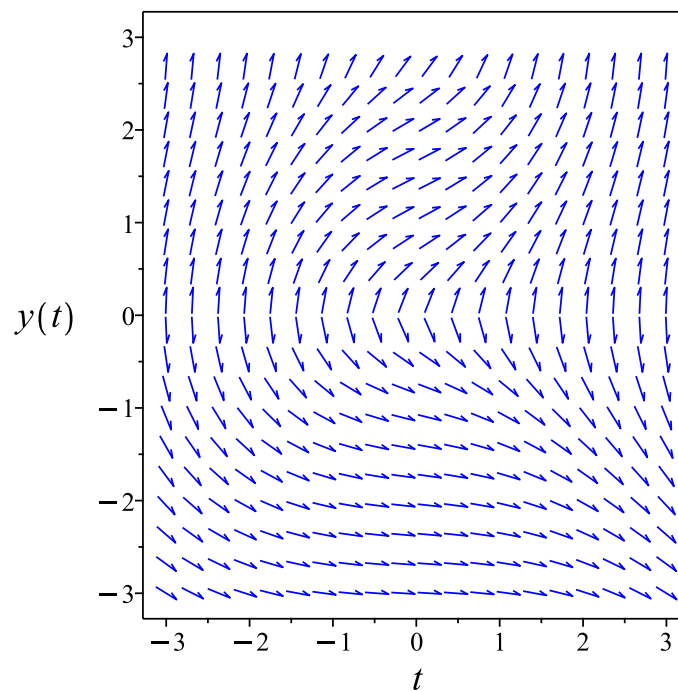


Figure 185: Slope field plot

Verification of solutions

$$-\frac{t^3}{3} - \frac{y^3}{3} + \frac{3y^2}{2} - t = c_1$$

Verified OK.

3.7.5 Maple step by step solution

Let's solve

$$y' - \frac{t^2+1}{3y-y^2} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$(3y - y^2) y' = t^2 + 1$$

- Integrate both sides with respect to t

$$\int (3y - y^2) y' dt = \int (t^2 + 1) dt + c_1$$

- Evaluate integral

$$-\frac{y^3}{3} + \frac{3y^2}{2} = \frac{1}{3}t^3 + c_1 + t$$

- Solve for y

$$y = \frac{\left(27-4t^3-12c_1-12t+2\sqrt{4t^6+24c_1t^3+24t^4-54t^3+36c_1^2+72c_1t+36t^2-162c_1-162t}\right)^{\frac{1}{3}}}{2} + \frac{1}{2\left(27-4t^3-12c_1-12t+2\sqrt{4t^6+24c_1t^3}\right)^{\frac{1}{3}}}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 444

```
dsolve(diff(y(t),t) = (t^2+1)/(3*y(t)-y(t)^2),y(t), singsol=all)
```

$y(t)$

$$= \frac{\left(27 - 4t^3 - 12c_1 - 12t + 2\sqrt{4t^6 + 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 + 72c_1t + 36t^2 - 162c_1 - 162t}\right)^{\frac{1}{3}}}{2 \cdot 9}$$

$$+ \frac{2 \left(27 - 4t^3 - 12c_1 - 12t + 2\sqrt{4t^6 + 24c_1t^3 + 24t^4 - 54t^3 + 36c_1^2 + 72c_1t + 36t^2 - 162c_1 - 162t}\right)^{\frac{1}{3}}}{2}$$

$y(t) =$

$$= \frac{(1 + i\sqrt{3}) \left(27 - 4t^3 - 12c_1 - 12t + 2\sqrt{4} \sqrt{\left(t^3 + 3t + 3c_1 - \frac{27}{2}\right) (t^3 + 3c_1 + 3t)}\right)^{\frac{2}{3}} - 9i\sqrt{3} - 6 \left(27 - 4t^3 - 12c_1 - 12t + 2\sqrt{4} \sqrt{\left(t^3 + 3t + 3c_1 - \frac{27}{2}\right) (t^3 + 3c_1 + 3t)}\right)^{\frac{1}{3}}}{4 \left(27 - 4t^3 - 12c_1 - 12t + 2\sqrt{4} \sqrt{\left(t^3 + 3t + 3c_1 - \frac{27}{2}\right) (t^3 + 3c_1 + 3t)}\right)^{\frac{2}{3}} - 9i\sqrt{3} - 6 \left(27 - 4t^3 - 12c_1 - 12t + 2\sqrt{4} \sqrt{\left(t^3 + 3t + 3c_1 - \frac{27}{2}\right) (t^3 + 3c_1 + 3t)}\right)^{\frac{1}{3}}}$$

$y(t)$

$$= \frac{(i\sqrt{3} - 1) \left(27 - 4t^3 - 12c_1 - 12t + 2\sqrt{4} \sqrt{\left(t^3 + 3t + 3c_1 - \frac{27}{2}\right) (t^3 + 3c_1 + 3t)}\right)^{\frac{2}{3}} - 9i\sqrt{3} + 6 \left(27 - 4t^3 - 12c_1 - 12t + 2\sqrt{4} \sqrt{\left(t^3 + 3t + 3c_1 - \frac{27}{2}\right) (t^3 + 3c_1 + 3t)}\right)^{\frac{1}{3}}}{4 \left(27 - 4t^3 - 12c_1 - 12t + 2\sqrt{4} \sqrt{\left(t^3 + 3t + 3c_1 - \frac{27}{2}\right) (t^3 + 3c_1 + 3t)}\right)^{\frac{2}{3}} - 9i\sqrt{3} + 6 \left(27 - 4t^3 - 12c_1 - 12t + 2\sqrt{4} \sqrt{\left(t^3 + 3t + 3c_1 - \frac{27}{2}\right) (t^3 + 3c_1 + 3t)}\right)^{\frac{1}{3}}}$$

✓ Solution by Mathematica

Time used: 3.185 (sec). Leaf size: 343

`DSolve[y'[t] == (t^2+1)/(3*y[t]-y[t]^2),y[t],t,IncludeSingularSolutions -> True]`

$$y(t) \rightarrow \frac{1}{2} \left(\sqrt[3]{-4t^3 + \sqrt{-729 + (4t^3 + 12t - 3(9 + 4c_1))^2} - 12t + 27 + 12c_1} + \frac{9}{\sqrt[3]{-4t^3 + \sqrt{-729 + (4t^3 + 12t - 3(9 + 4c_1))^2} - 12t + 27 + 12c_1}} + 3 \right)$$

$$y(t) \rightarrow \frac{1}{4} \left(i(\sqrt{3} + i) \sqrt[3]{-4t^3 + \sqrt{-729 + (4t^3 + 12t - 3(9 + 4c_1))^2} - 12t + 27 + 12c_1} - \frac{9(1 + i\sqrt{3})}{\sqrt[3]{-4t^3 + \sqrt{-729 + (4t^3 + 12t - 3(9 + 4c_1))^2} - 12t + 27 + 12c_1}} + 6 \right)$$

$$y(t) \rightarrow \frac{1}{4} \left(- \left((1 + i\sqrt{3}) \sqrt[3]{-4t^3 + \sqrt{-729 + (4t^3 + 12t - 3(9 + 4c_1))^2} - 12t + 27 + 12c_1} + \frac{9i(\sqrt{3} + i)}{\sqrt[3]{-4t^3 + \sqrt{-729 + (4t^3 + 12t - 3(9 + 4c_1))^2} - 12t + 27 + 12c_1}} + 6 \right) \right)$$

3.8 problem 12

3.8.1	Solving as separable ode	1011
3.8.2	Solving as first order ode lie symmetry lookup ode	1013
3.8.3	Solving as exact ode	1017
3.8.4	Maple step by step solution	1021

Internal problem ID [523]

Internal file name [OUTPUT/523_Sunday_June_05_2022_01_43_01_AM_79042422/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 12.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{\cot(t)y}{1+y} = 0$$

3.8.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(t, y) \\ &= f(t)g(y) \\ &= \frac{\cot(t)y}{y+1}\end{aligned}$$

Where $f(t) = \cot(t)$ and $g(y) = \frac{y}{y+1}$. Integrating both sides gives

$$\frac{1}{\frac{y}{y+1}} dy = \cot(t) dt$$

$$\int \frac{1}{\frac{y}{y+1}} dy = \int \cot(t) dt$$

$$y + \ln(y) = \ln(\sin(t)) + c_1$$

Which results in

$$y = \text{LambertW}(e^{c_1} \sin(t))$$

Since c_1 is constant, then exponential powers of this constant are constants also, and these can be simplified to just c_1 in the above solution. Which simplifies to

$$y = \text{LambertW}(e^{c_1} \sin(t))$$

gives

$$y = \text{LambertW}(c_1 \sin(t))$$

Summary

The solution(s) found are the following

$$y = \text{LambertW}(c_1 \sin(t)) \tag{1}$$

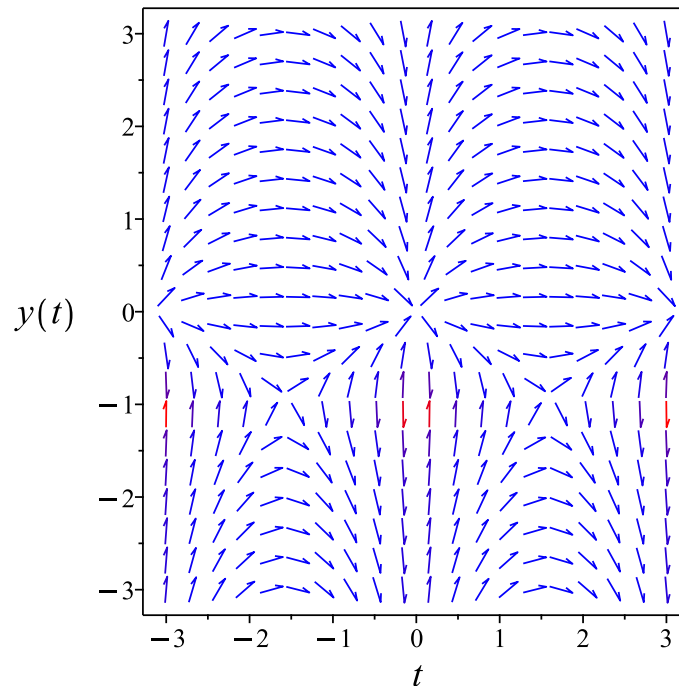


Figure 186: Slope field plot

Verification of solutions

$$y = \text{LambertW}(c_1 \sin(t))$$

Verified OK.

3.8.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{\cot(t)y}{y+1}$$
$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 203: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= \frac{1}{\cot(t)} \\ \eta(t, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dt \\ &= \int \frac{1}{\frac{1}{\cot(t)}} dt \end{aligned}$$

Which results in

$$S = \ln(\sin(t))$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{\cot(t)y}{y+1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 0 \\ R_y &= 1 \\ S_t &= \cot(t) \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{y+1}{y} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{R+1}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R + \ln(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$\ln(\sin(t)) = y + \ln(y) + c_1$$

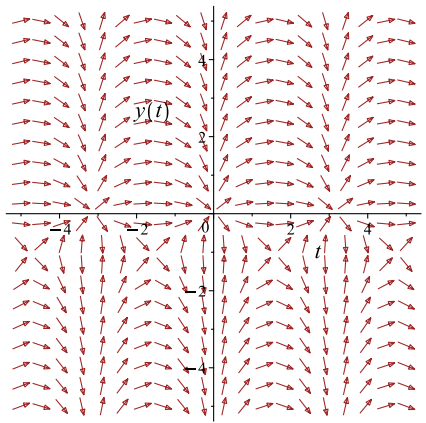
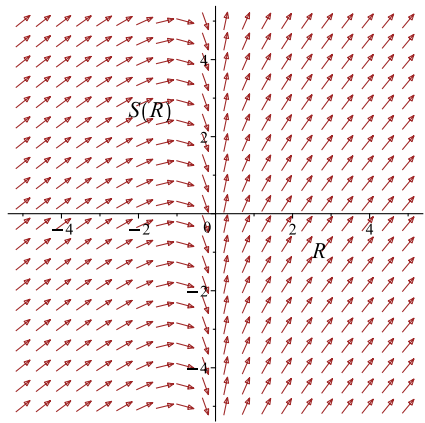
Which simplifies to

$$\ln(\sin(t)) = y + \ln(y) + c_1$$

Which gives

$$y = \text{LambertW}(e^{-c_1} \sin(t))$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{\cot(t)y}{y+1}$ 	$R = y$ $S = \ln(\sin(t))$	$\frac{dS}{dR} = \frac{R+1}{R}$ 

Summary

The solution(s) found are the following

$$y = \text{LambertW}(e^{-c_1} \sin(t)) \quad (1)$$

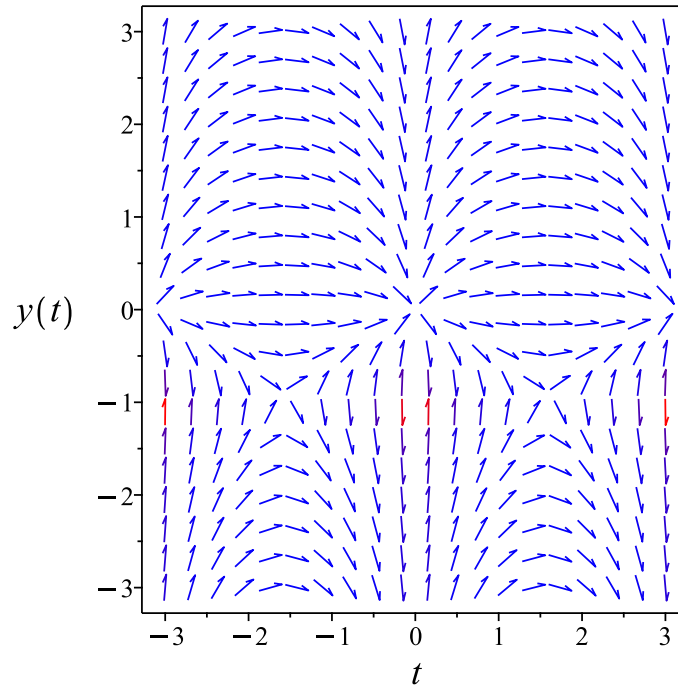


Figure 187: Slope field plot

Verification of solutions

$$y = \text{LambertW}(e^{-c_1} \sin(t))$$

Verified OK.

3.8.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(\frac{y+1}{y}\right) dy &= (\cot(t)) dt \\ (-\cot(t)) dt + \left(\frac{y+1}{y}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -\cot(t) \\ N(t, y) &= \frac{y+1}{y}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-\cot(t)) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} \left(\frac{y+1}{y} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -\cot(t) dt \\ \phi &= -\ln(\sin(t)) + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{y+1}{y}$. Therefore equation (4) becomes

$$\frac{y+1}{y} = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{y+1}{y}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(\frac{y+1}{y} \right) dy \\ f(y) &= y + \ln(y) + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\ln(\sin(t)) + y + \ln(y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\ln(\sin(t)) + y + \ln(y)$$

The solution becomes

$$y = \text{LambertW}(e^{c_1} \sin(t))$$

Summary

The solution(s) found are the following

$$y = \text{LambertW}(e^{c_1} \sin(t)) \tag{1}$$

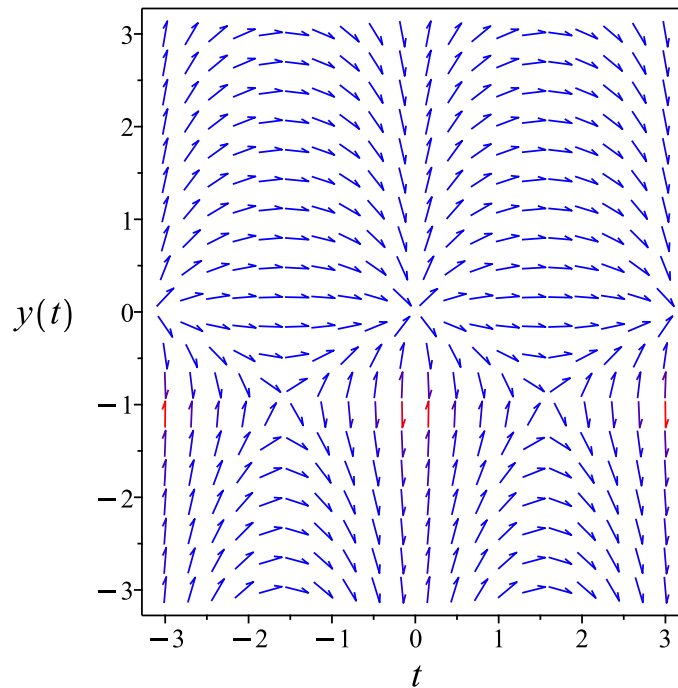


Figure 188: Slope field plot

Verification of solutions

$$y = \text{LambertW}(e^{c_1} \sin(t))$$

Verified OK.

3.8.4 Maple step by step solution

Let's solve

$$y' - \frac{\cot(t)y}{1+y} = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'(1+y)}{y} = \cot(t)$$

- Integrate both sides with respect to t

$$\int \frac{y'(1+y)}{y} dt = \int \cot(t) dt + c_1$$

- Evaluate integral

$$y + \ln(y) = \ln(\sin(t)) + c_1$$

- Solve for y

$$y = \text{LambertW}(e^{c_1} \sin(t))$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.032 (sec). Leaf size: 9

```
dsolve(diff(y(t),t) = cot(t)*y(t)/(1+y(t)),y(t), singsol=all)
```

$$y(t) = \text{LambertW}(c_1 \sin(t))$$

✓ Solution by Mathematica

Time used: 1.602 (sec). Leaf size: 18

```
DSolve[y'[t] == Cot[t]*y[t]/(1+y[t]),y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow W(e^{c_1} \sin(t))$$

$$y(t) \rightarrow 0$$

3.9 problem 13

3.9.1	Solving as separable ode	1023
3.9.2	Solving as homogeneousTypeD2 ode	1025
3.9.3	Solving as differentialType ode	1027
3.9.4	Solving as first order ode lie symmetry lookup ode	1028
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3.9.6	Maple step by step solution	1036

Internal problem ID [524]

Internal file name [OUTPUT/524_Sunday_June_05_2022_01_43_02_AM_61510835/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 13.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' + \frac{4t}{y} = 0$$

3.9.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(t, y) \\ &= f(t)g(y) \\ &= -\frac{4t}{y}\end{aligned}$$

Where $f(t) = -4t$ and $g(y) = \frac{1}{y}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y} dy &= -4t dt \\ \int \frac{1}{y} dy &= \int -4t dt \\ \frac{y^2}{2} &= -2t^2 + c_1\end{aligned}$$

Which results in

$$\begin{aligned}y &= \sqrt{-4t^2 + 2c_1} \\ y &= -\sqrt{-4t^2 + 2c_1}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \sqrt{-4t^2 + 2c_1} \tag{1}$$

$$y = -\sqrt{-4t^2 + 2c_1} \tag{2}$$

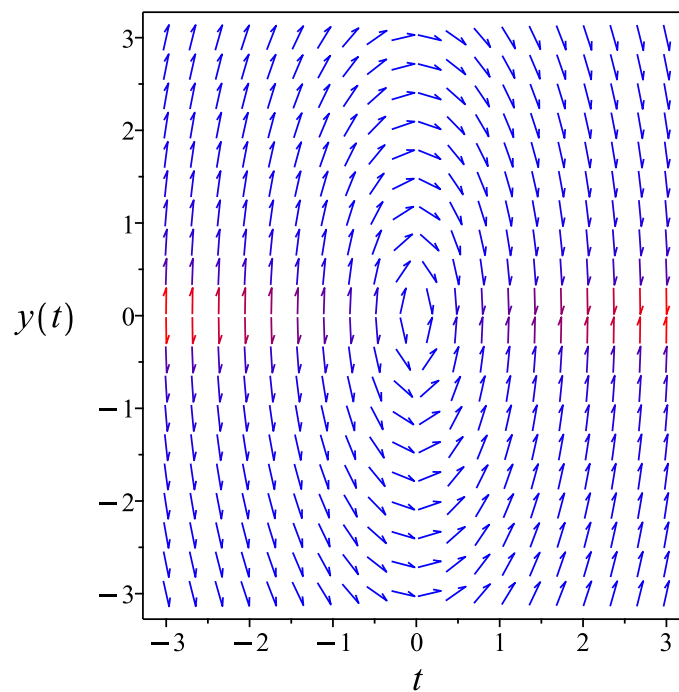


Figure 189: Slope field plot

Verification of solutions

$$y = \sqrt{-4t^2 + 2c_1}$$

Verified OK.

$$y = -\sqrt{-4t^2 + 2c_1}$$

Verified OK.

3.9.2 Solving as homogeneous Type D2 ode

Using the change of variables $y = u(t)t$ on the above ode results in new ode in $u(t)$

$$u'(t)t + u(t) + \frac{4}{u(t)} = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{u^2 + 4}{tu} \end{aligned}$$

Where $f(t) = -\frac{1}{t}$ and $g(u) = \frac{u^2+4}{u}$. Integrating both sides gives

$$\begin{aligned} \frac{1}{\frac{u^2+4}{u}} du &= -\frac{1}{t} dt \\ \int \frac{1}{\frac{u^2+4}{u}} du &= \int -\frac{1}{t} dt \\ \frac{\ln(u^2 + 4)}{2} &= -\ln(t) + c_2 \end{aligned}$$

Raising both side to exponential gives

$$\sqrt{u^2 + 4} = e^{-\ln(t)+c_2}$$

Which simplifies to

$$\sqrt{u^2 + 4} = \frac{c_3}{t}$$

Which simplifies to

$$\sqrt{u(t)^2 + 4} = \frac{c_3 e^{c_2}}{t}$$

The solution is

$$\sqrt{u(t)^2 + 4} = \frac{c_3 e^{c_2}}{t}$$

Replacing $u(t)$ in the above solution by $\frac{y}{t}$ results in the solution for y in implicit form

$$\begin{aligned}\sqrt{\frac{y^2}{t^2} + 4} &= \frac{c_3 e^{c_2}}{t} \\ \sqrt{\frac{y^2 + 4t^2}{t^2}} &= \frac{c_3 e^{c_2}}{t}\end{aligned}$$

Summary

The solution(s) found are the following

$$\sqrt{\frac{y^2 + 4t^2}{t^2}} = \frac{c_3 e^{c_2}}{t} \quad (1)$$

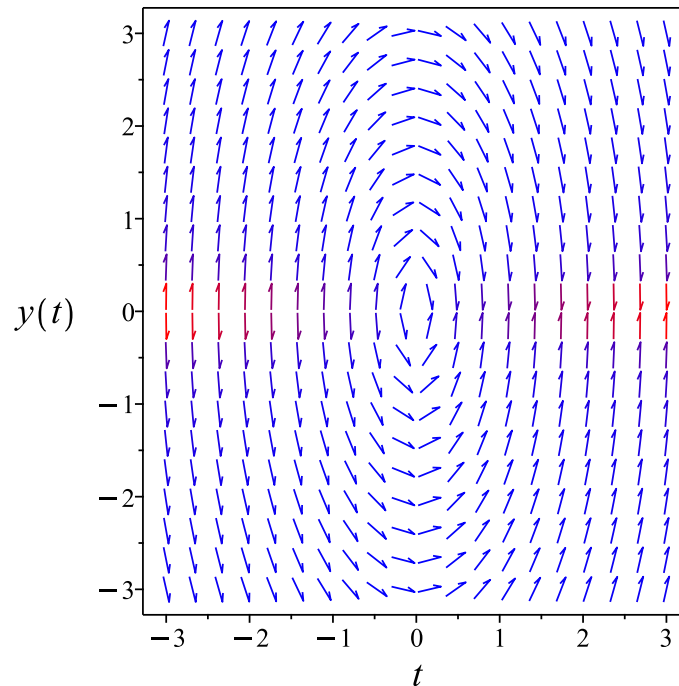


Figure 190: Slope field plot

Verification of solutions

$$\sqrt{\frac{y^2 + 4t^2}{t^2}} = \frac{c_3 e^{c_2}}{t}$$

Verified OK.

3.9.3 Solving as differentialType ode

Writing the ode as

$$y' = -\frac{4t}{y} \quad (1)$$

Which becomes

$$(y) dy = (-4t) dt \quad (2)$$

But the RHS is complete differential because

$$(-4t) dt = d(-2t^2)$$

Hence (2) becomes

$$(y) dy = d(-2t^2)$$

Integrating both sides gives gives these solutions

$$y = \sqrt{-4t^2 + 2c_1} + c_1$$

$$y = -\sqrt{-4t^2 + 2c_1} + c_1$$

Summary

The solution(s) found are the following

$$y = \sqrt{-4t^2 + 2c_1} + c_1 \quad (1)$$

$$y = -\sqrt{-4t^2 + 2c_1} + c_1 \quad (2)$$

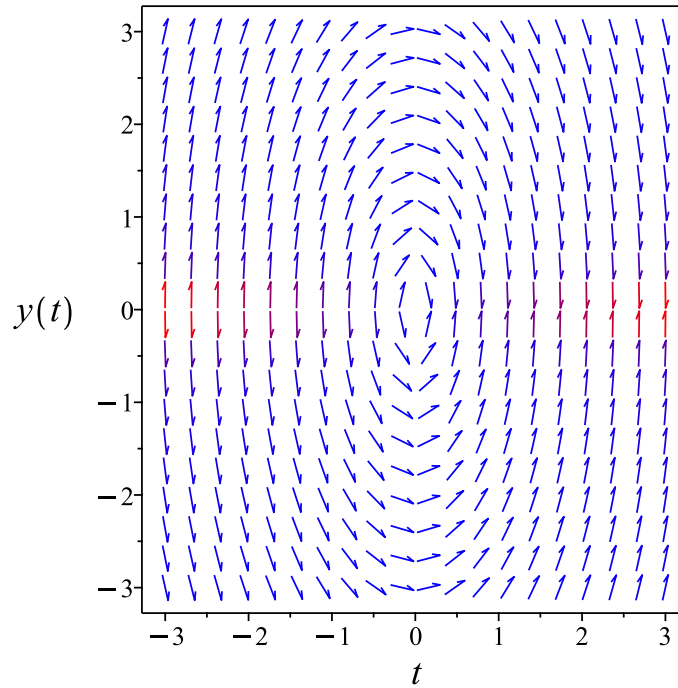


Figure 191: Slope field plot

Verification of solutions

$$y = \sqrt{-4t^2 + 2c_1} + c_1$$

Verified OK.

$$y = -\sqrt{-4t^2 + 2c_1} + c_1$$

Verified OK.

3.9.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{4t}{y}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 206: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= -\frac{1}{4t} \\ \eta(t, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dt \\ &= \int \frac{1}{-\frac{1}{4t}} dt \end{aligned}$$

Which results in

$$S = -2t^2$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{4t}{y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 0 \\ R_y &= 1 \\ S_t &= -4t \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^2}{2} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$-2t^2 = \frac{y^2}{2} + c_1$$

Which simplifies to

$$-2t^2 = \frac{y^2}{2} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{4t}{y}$	$R = y$ $S = -2t^2$	$\frac{dS}{dR} = R$

Summary

The solution(s) found are the following

$$-2t^2 = \frac{y^2}{2} + c_1 \quad (1)$$

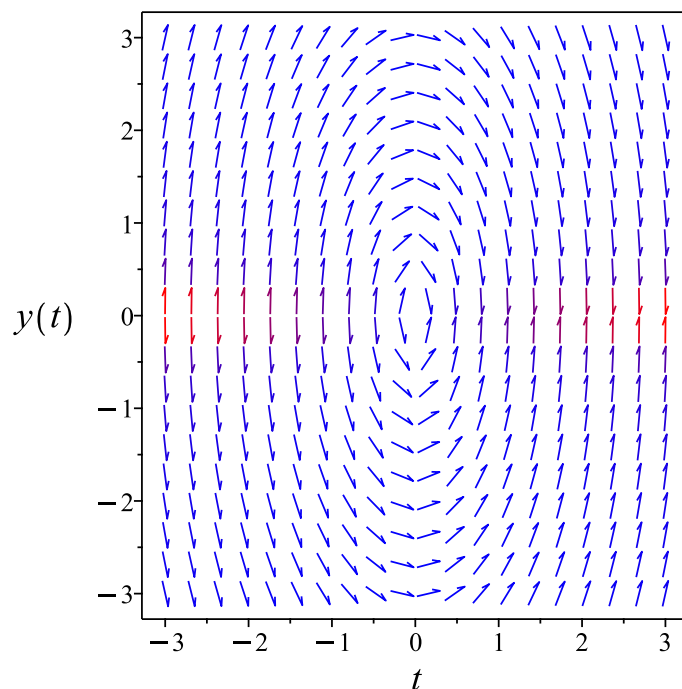


Figure 192: Slope field plot

Verification of solutions

$$-2t^2 = \frac{y^2}{2} + c_1$$

Verified OK.

3.9.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(-\frac{y}{4}\right) dy &= (t) dt \\ (-t) dt + \left(-\frac{y}{4}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -t \\ N(t, y) &= -\frac{y}{4}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-t) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} \left(-\frac{y}{4} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -t dt \\ \phi &= -\frac{t^2}{2} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{y}{4}$. Therefore equation (4) becomes

$$-\frac{y}{4} = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{y}{4}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(-\frac{y}{4} \right) dy \\ f(y) &= -\frac{y^2}{8} + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{t^2}{2} - \frac{y^2}{8} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{t^2}{2} - \frac{y^2}{8}$$

Summary

The solution(s) found are the following

$$-\frac{t^2}{2} - \frac{y^2}{8} = c_1 \tag{1}$$

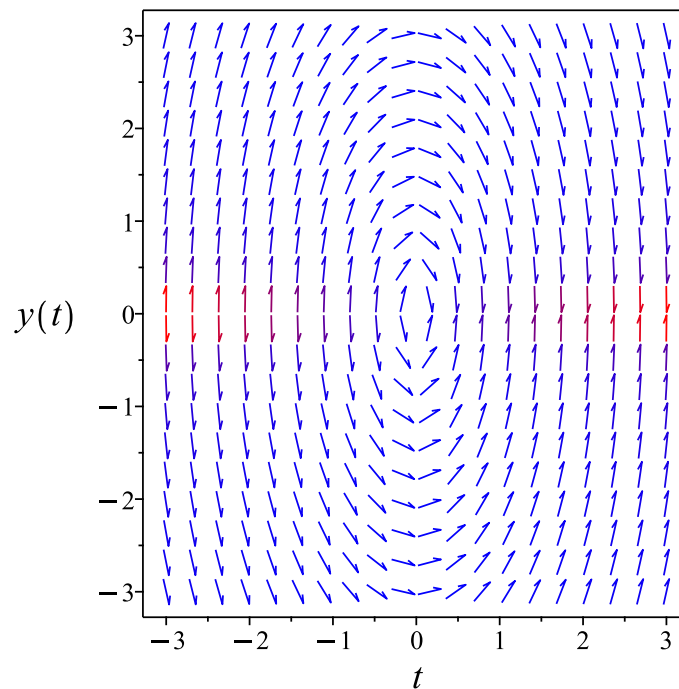


Figure 193: Slope field plot

Verification of solutions

$$-\frac{t^2}{2} - \frac{y^2}{8} = c_1$$

Verified OK.

3.9.6 Maple step by step solution

Let's solve

$$y' + \frac{4t}{y} = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$y'y = -4t$$

- Integrate both sides with respect to t

$$\int y'y dt = \int -4t dt + c_1$$

- Evaluate integral

$$\frac{y^2}{2} = -2t^2 + c_1$$

- Solve for y

$$\{y = \sqrt{-4t^2 + 2c_1}, y = -\sqrt{-4t^2 + 2c_1}\}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 27

```
dsolve(diff(y(t),t) = -4*t/y(t),y(t), singsol=all)
```

$$y(t) = \sqrt{-4t^2 + c_1}$$
$$y(t) = -\sqrt{-4t^2 + c_1}$$

✓ Solution by Mathematica

Time used: 0.087 (sec). Leaf size: 46

```
DSolve[y'[t]== -4*t/y[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -\sqrt{2}\sqrt{-2t^2 + c_1}$$

$$y(t) \rightarrow \sqrt{2}\sqrt{-2t^2 + c_1}$$

3.10 problem 14

3.10.1 Solving as separable ode	1038
3.10.2 Solving as first order ode lie symmetry lookup ode	1040
3.10.3 Solving as exact ode	1044
3.10.4 Solving as riccati ode	1048
3.10.5 Maple step by step solution	1050

Internal problem ID [525]

Internal file name [OUTPUT/525_Sunday_June_05_2022_01_43_03_AM_244180/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 14.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - 2ty^2 = 0$$

3.10.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(t, y) \\ &= f(t)g(y) \\ &= 2t y^2\end{aligned}$$

Where $f(t) = 2t$ and $g(y) = y^2$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y^2} dy &= 2t dt \\ \int \frac{1}{y^2} dy &= \int 2t dt\end{aligned}$$

$$-\frac{1}{y} = t^2 + c_1$$

Which results in

$$y = -\frac{1}{t^2 + c_1}$$

Summary

The solution(s) found are the following

$$y = -\frac{1}{t^2 + c_1} \tag{1}$$

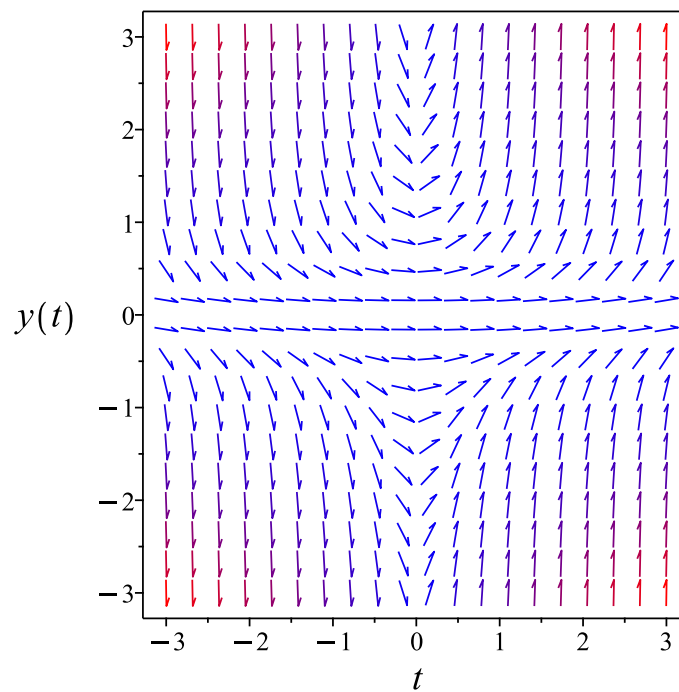


Figure 194: Slope field plot

Verification of solutions

$$y = -\frac{1}{t^2 + c_1}$$

Verified OK.

3.10.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = 2ty^2$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 209: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= \frac{1}{2t} \\ \eta(t, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dt \\ &= \int \frac{1}{\frac{1}{2t}} dt\end{aligned}$$

Which results in

$$S = t^2$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y}\tag{2}$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = 2t y^2$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_t &= 0 \\ R_y &= 1 \\ S_t &= 2t \\ S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{y^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{1}{R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$t^2 = -\frac{1}{y} + c_1$$

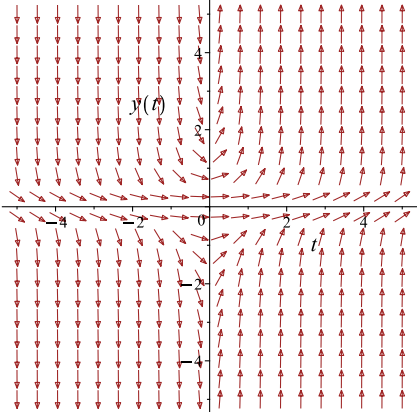
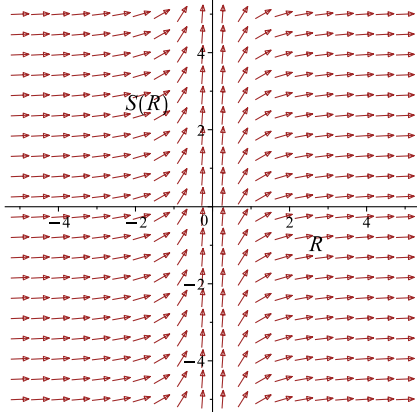
Which simplifies to

$$t^2 = -\frac{1}{y} + c_1$$

Which gives

$$y = \frac{1}{-t^2 + c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = 2t y^2$ 	$R = y$ $S = t^2$	$\frac{dS}{dR} = \frac{1}{R^2}$ 

Summary

The solution(s) found are the following

$$y = \frac{1}{-t^2 + c_1} \tag{1}$$

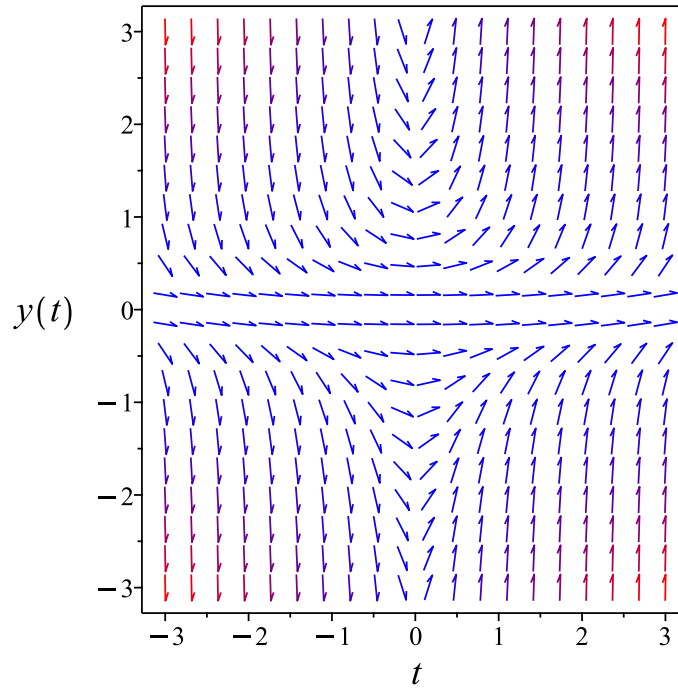


Figure 195: Slope field plot

Verification of solutions

$$y = \frac{1}{-t^2 + c_1}$$

Verified OK.

3.10.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(\frac{1}{2y^2}\right) dy &= (t) dt \\ (-t) dt + \left(\frac{1}{2y^2}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -t \\ N(t, y) &= \frac{1}{2y^2}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-t) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} \left(\frac{1}{2y^2} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -t dt \\ \phi &= -\frac{t^2}{2} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{2y^2}$. Therefore equation (4) becomes

$$\frac{1}{2y^2} = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{1}{2y^2}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(\frac{1}{2y^2} \right) dy$$
$$f(y) = -\frac{1}{2y} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{t^2}{2} - \frac{1}{2y} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{t^2}{2} - \frac{1}{2y}$$

The solution becomes

$$y = -\frac{1}{t^2 + 2c_1}$$

Summary

The solution(s) found are the following

$$y = -\frac{1}{t^2 + 2c_1} \tag{1}$$

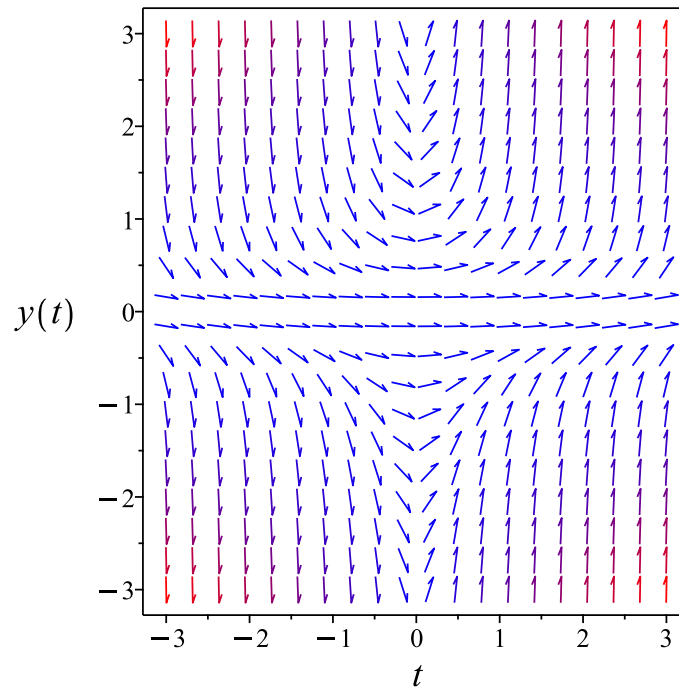


Figure 196: Slope field plot

Verification of solutions

$$y = -\frac{1}{t^2 + 2c_1}$$

Verified OK.

3.10.4 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= 2ty^2 \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = 2ty^2$$

With Riccati ODE standard form

$$y' = f_0(t) + f_1(t)y + f_2(t)y^2$$

Shows that $f_0(t) = 0$, $f_1(t) = 0$ and $f_2(t) = 2t$. Let

$$\begin{aligned}y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{2tu}\end{aligned}\tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(t) - (f_2' + f_1 f_2) u'(t) + f_2^2 f_0 u(t) = 0\tag{2}$$

But

$$\begin{aligned}f_2' &= 2 \\ f_1 f_2 &= 0 \\ f_2^2 f_0 &= 0\end{aligned}$$

Substituting the above terms back in equation (2) gives

$$2t u''(t) - 2u'(t) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(t) = c_2 t^2 + c_1$$

The above shows that

$$u'(t) = 2c_2 t$$

Using the above in (1) gives the solution

$$y = -\frac{c_2}{c_2 t^2 + c_1}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = -\frac{1}{t^2 + c_3}$$

Summary

The solution(s) found are the following

$$y = -\frac{1}{t^2 + c_3} \quad (1)$$

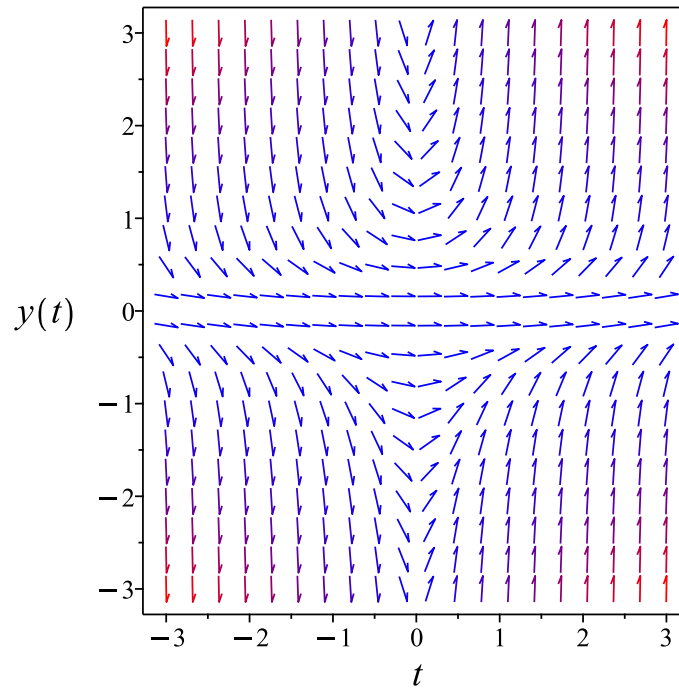


Figure 197: Slope field plot

Verification of solutions

$$y = -\frac{1}{t^2 + c_3}$$

Verified OK.

3.10.5 Maple step by step solution

Let's solve

$$y' - 2ty^2 = 0$$

- Highest derivative means the order of the ODE is 1
- y'
- Separate variables

$$\frac{y'}{y^2} = 2t$$

- Integrate both sides with respect to t

$$\int \frac{y'}{y^2} dt = \int 2t dt + c_1$$

- Evaluate integral

$$-\frac{1}{y} = t^2 + c_1$$

- Solve for y

$$y = -\frac{1}{t^2 + c_1}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 13

```
dsolve(diff(y(t),t) = 2*t*y(t)^2,y(t), singsol=all)
```

$$y(t) = \frac{1}{-t^2 + c_1}$$

✓ Solution by Mathematica

Time used: 0.111 (sec). Leaf size: 20

```
DSolve[y'[t] == 2*t*y[t]^2,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -\frac{1}{t^2 + c_1}$$
$$y(t) \rightarrow 0$$

3.11 problem 15

3.11.1 Solving as quadrature ode 1052

3.11.2 Maple step by step solution 1053

Internal problem ID [526]

Internal file name [OUTPUT/526_Sunday_June_05_2022_01_43_04_AM_90328302/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 15.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**quadrature**"

Maple gives the following as the ode type

[_quadrature]

$$y^3 + y' = 0$$

3.11.1 Solving as quadrature ode

Integrating both sides gives

$$\int -\frac{1}{y^3} dy = t + c_1$$
$$\frac{1}{2y^2} = t + c_1$$

Solving for y gives these solutions

$$y_1 = \frac{1}{\sqrt{2c_1 + 2t}}$$
$$y_2 = -\frac{1}{\sqrt{2c_1 + 2t}}$$

Summary

The solution(s) found are the following

$$y = \frac{1}{\sqrt{2c_1 + 2t}} \tag{1}$$

$$y = -\frac{1}{\sqrt{2c_1 + 2t}} \tag{2}$$

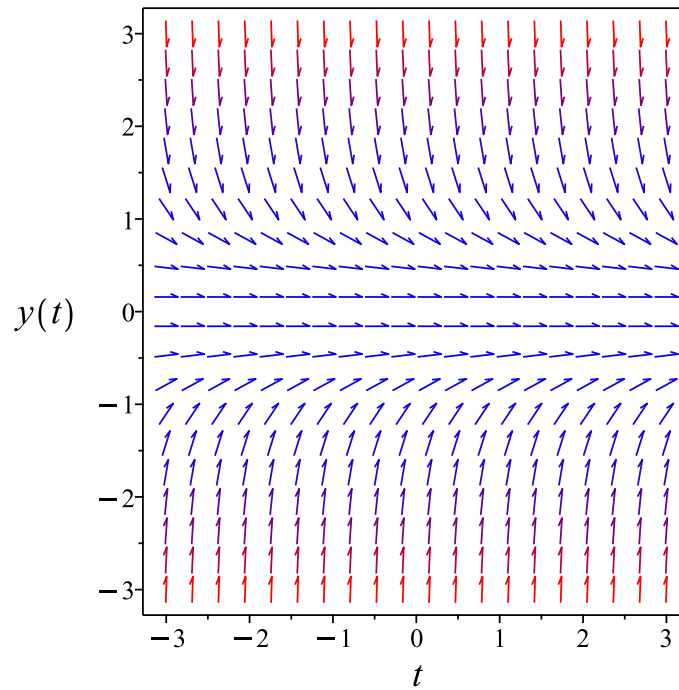


Figure 198: Slope field plot

Verification of solutions

$$y = \frac{1}{\sqrt{2c_1 + 2t}}$$

Verified OK.

$$y = -\frac{1}{\sqrt{2c_1 + 2t}}$$

Verified OK.

3.11.2 Maple step by step solution

Let's solve

$$y^3 + y' = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{y^3} = -1$$

- Integrate both sides with respect to t

$$\int \frac{y'}{y^3} dt = \int (-1) dt + c_1$$

- Evaluate integral

$$-\frac{1}{2y^2} = -t + c_1$$

- Solve for y

$$\left\{ y = \frac{1}{\sqrt{-2c_1+2t}}, y = -\frac{1}{\sqrt{-2c_1+2t}} \right\}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 23

```
dsolve(y(t)^3+diff(y(t),t) = 0,y(t), singsol=all)
```

$$y(t) = \frac{1}{\sqrt{2t + c_1}}$$

$$y(t) = -\frac{1}{\sqrt{2t + c_1}}$$

✓ Solution by Mathematica

Time used: 0.124 (sec). Leaf size: 40

```
DSolve[y[t]^3+y'[t] == 0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -\frac{1}{\sqrt{2t - 2c_1}}$$

$$y(t) \rightarrow \frac{1}{\sqrt{2t - 2c_1}}$$

$$y(t) \rightarrow 0$$

3.12 problem 16

3.12.1 Solving as separable ode	1055
3.12.2 Solving as first order ode lie symmetry lookup ode	1057
3.12.3 Solving as exact ode	1061
3.12.4 Maple step by step solution	1065

Internal problem ID [527]

Internal file name [OUTPUT/527_Sunday_June_05_2022_01_43_05_AM_3895157/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 16.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{t^2}{(t^3 + 1)y} = 0$$

3.12.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(t, y) \\ &= f(t)g(y) \\ &= \frac{t^2}{(t^3 + 1)y}\end{aligned}$$

Where $f(t) = \frac{t^2}{t^3+1}$ and $g(y) = \frac{1}{y}$. Integrating both sides gives

$$\frac{1}{y} dy = \frac{t^2}{t^3 + 1} dt$$

$$\int \frac{1}{\frac{1}{y}} dy = \int \frac{t^2}{t^3 + 1} dt$$

$$\frac{y^2}{2} = \frac{\ln(t^3 + 1)}{3} + c_1$$

Which results in

$$y = \frac{\sqrt{6 \ln(t^3 + 1) + 18c_1}}{3}$$

$$y = -\frac{\sqrt{6 \ln(t^3 + 1) + 18c_1}}{3}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{6 \ln(t^3 + 1) + 18c_1}}{3} \tag{1}$$

$$y = -\frac{\sqrt{6 \ln(t^3 + 1) + 18c_1}}{3} \tag{2}$$

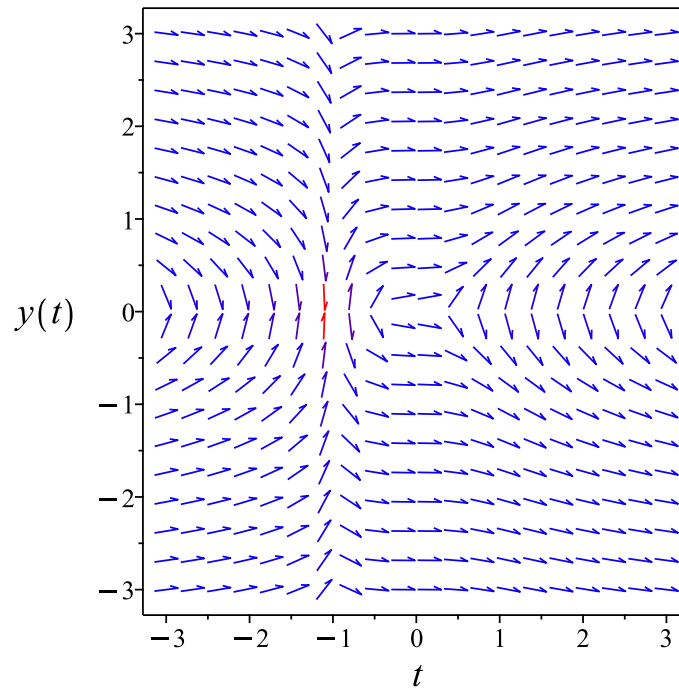


Figure 199: Slope field plot

Verification of solutions

$$y = \frac{\sqrt{6 \ln(t^3 + 1) + 18c_1}}{3}$$

Verified OK.

$$y = -\frac{\sqrt{6 \ln(t^3 + 1) + 18c_1}}{3}$$

Verified OK.

3.12.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{t^2}{(t^3 + 1)y}$$
$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 213: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= \frac{t^3 + 1}{t^2} \\ \eta(t, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dt \\ &= \int \frac{1}{\frac{t^3+1}{t^2}} dt \end{aligned}$$

Which results in

$$S = \frac{\ln(t^3 + 1)}{3}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{t^2}{(t^3 + 1)y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 0 \\ R_y &= 1 \\ S_t &= \frac{t^2}{(t^2 - t + 1)(t + 1)} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^2}{2} + c_1 \quad (4)$$

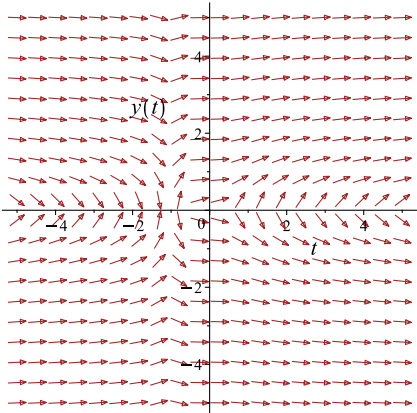
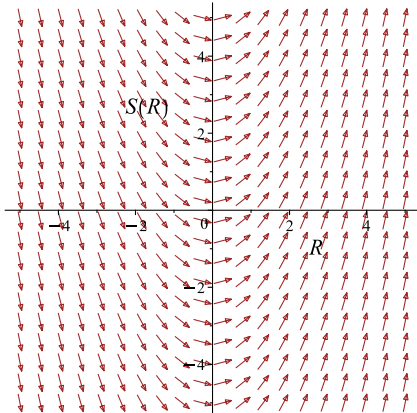
To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$\frac{\ln(t+1)}{3} + \frac{\ln(t^2-t+1)}{3} = \frac{y^2}{2} + c_1$$

Which simplifies to

$$\frac{\ln(t+1)}{3} + \frac{\ln(t^2-t+1)}{3} = \frac{y^2}{2} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{t^2}{(t^3+1)y}$ 	$R = y$ $S = \frac{\ln(t+1)}{3} + \frac{\ln(t^2-t+1)}{3}$	$\frac{dS}{dR} = R$ 

Summary

The solution(s) found are the following

$$\frac{\ln(t+1)}{3} + \frac{\ln(t^2-t+1)}{3} = \frac{y^2}{2} + c_1 \quad (1)$$

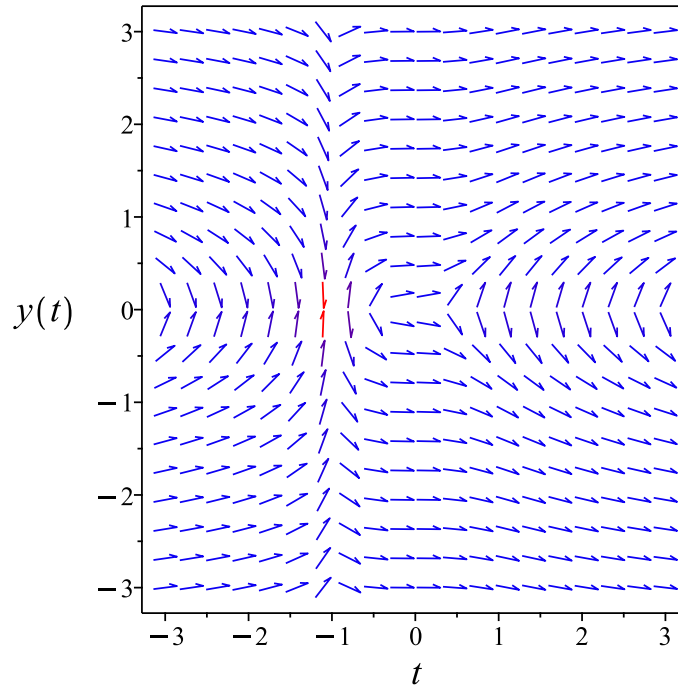


Figure 200: Slope field plot

Verification of solutions

$$\frac{\ln(t+1)}{3} + \frac{\ln(t^2-t+1)}{3} = \frac{y^2}{2} + c_1$$

Verified OK.

3.12.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(y) dy &= \left(\frac{t^2}{t^3 + 1}\right) dt \\ \left(-\frac{t^2}{t^3 + 1}\right) dt + (y) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -\frac{t^2}{t^3 + 1} \\ N(t, y) &= y\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{t^2}{t^3 + 1}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(y) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -\frac{t^2}{t^3 + 1} dt \\ \phi &= -\frac{\ln(t^3 + 1)}{3} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = y$. Therefore equation (4) becomes

$$y = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = y$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (y) dy \\ f(y) &= \frac{y^2}{2} + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{\ln(t^3 + 1)}{3} + \frac{y^2}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{\ln(t^3 + 1)}{3} + \frac{y^2}{2}$$

Summary

The solution(s) found are the following

$$-\frac{\ln(t^3 + 1)}{3} + \frac{y^2}{2} = c_1 \quad (1)$$

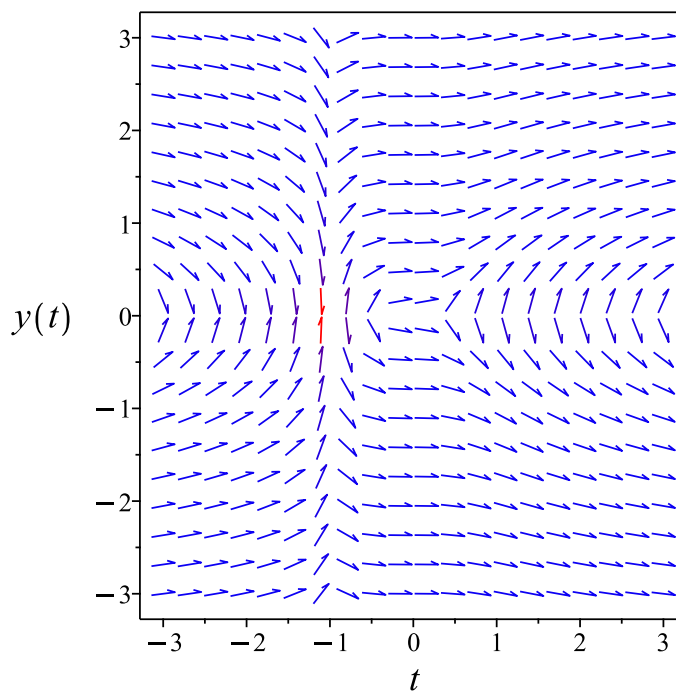


Figure 201: Slope field plot

Verification of solutions

$$-\frac{\ln(t^3 + 1)}{3} + \frac{y^2}{2} = c_1$$

Verified OK.

3.12.4 Maple step by step solution

Let's solve

$$y' - \frac{t^2}{(t^3+1)y} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$y'y = \frac{t^2}{t^3+1}$$

- Integrate both sides with respect to t

$$\int y'y dt = \int \frac{t^2}{t^3+1} dt + c_1$$

- Evaluate integral

$$\frac{y^2}{2} = \frac{\ln(t^3+1)}{3} + c_1$$

- Solve for y

$$\left\{ y = -\frac{\sqrt{6\ln(t^3+1)+18c_1}}{3}, y = \frac{\sqrt{6\ln(t^3+1)+18c_1}}{3} \right\}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 39

```
dsolve(diff(y(t),t) = t^2/(t^3+1)/y(t),y(t), singsol=all)
```

$$y(t) = -\frac{\sqrt{6\ln(t^3+1)+9c_1}}{3}$$
$$y(t) = \frac{\sqrt{6\ln(t^3+1)+9c_1}}{3}$$

✓ Solution by Mathematica

Time used: 0.123 (sec). Leaf size: 56

```
DSolve[y'[t] == t^2/(t^3+1)/y[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -\sqrt{\frac{2}{3}}\sqrt{\log(t^3 + 1) + 3c_1}$$

$$y(t) \rightarrow \sqrt{\frac{2}{3}}\sqrt{\log(t^3 + 1) + 3c_1}$$

3.13 problem 17

3.13.1 Solving as separable ode	1067
3.13.2 Solving as first order ode lie symmetry lookup ode	1069
3.13.3 Solving as bernoulli ode	1073
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3.13.6 Maple step by step solution	1082

Internal problem ID [528]

Internal file name [OUTPUT/528_Sunday_June_05_2022_01_43_06_AM_89383031/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 17.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "bernoulli", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - t(3 - y)y = 0$$

3.13.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(t, y) \\ &= f(t)g(y) \\ &= -ty(-3 + y)\end{aligned}$$

Where $f(t) = -t$ and $g(y) = y(-3 + y)$. Integrating both sides gives

$$\frac{1}{y(-3 + y)} dy = -t dt$$

$$\int \frac{1}{y(-3+y)} dy = \int -t dt$$

$$-\frac{\ln(y)}{3} + \frac{\ln(-3+y)}{3} = -\frac{t^2}{2} + c_1$$

The above can be written as

$$\left(-\frac{1}{3}\right)(\ln(y) - \ln(-3+y)) = -\frac{t^2}{2} + 2c_1$$

$$\ln(y) - \ln(-3+y) = (-3) \left(-\frac{t^2}{2} + 2c_1\right)$$

$$= \frac{3t^2}{2} - 6c_1$$

Raising both side to exponential gives

$$e^{\ln(y) - \ln(-3+y)} = e^{\frac{3t^2}{2} - 6c_1}$$

Which simplifies to

$$\frac{y}{-3+y} = -3c_1 e^{\frac{3t^2}{2}}$$

$$= c_2 e^{\frac{3t^2}{2}}$$

Summary

The solution(s) found are the following

$$y = \frac{3c_2 e^{\frac{3t^2}{2}}}{-1 + c_2 e^{\frac{3t^2}{2}}} \quad (1)$$

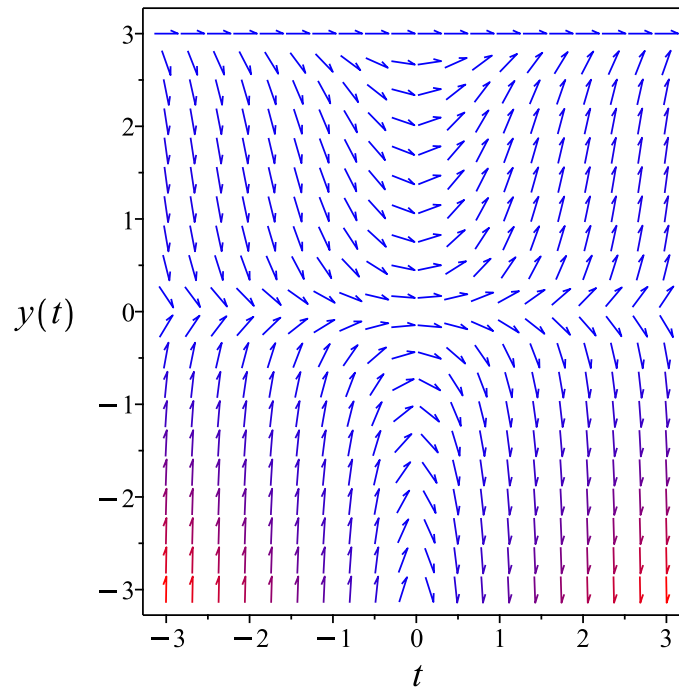


Figure 202: Slope field plot

Verification of solutions

$$y = \frac{3c_2 e^{\frac{3t^2}{2}}}{-1 + c_2 e^{\frac{3t^2}{2}}}$$

Verified OK.

3.13.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -ty(-3 + y)$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 216: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= -\frac{1}{t} \\ \eta(t, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dt \\ &= \int \frac{1}{-\frac{1}{t}} dt \end{aligned}$$

Which results in

$$S = -\frac{t^2}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -ty(-3 + y)$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 0 \\ R_y &= 1 \\ S_t &= -t \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{y(-3 + y)} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R(-3 + R)}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{\ln(R)}{3} + \frac{\ln(-3+R)}{3} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$-\frac{t^2}{2} = -\frac{\ln(y)}{3} + \frac{\ln(-3+y)}{3} + c_1$$

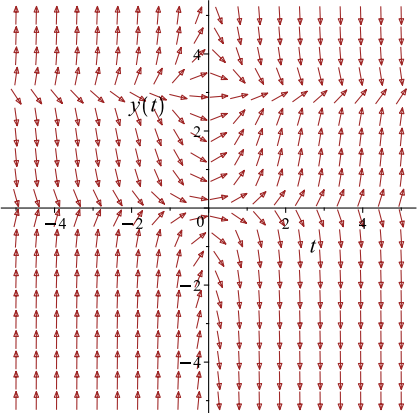
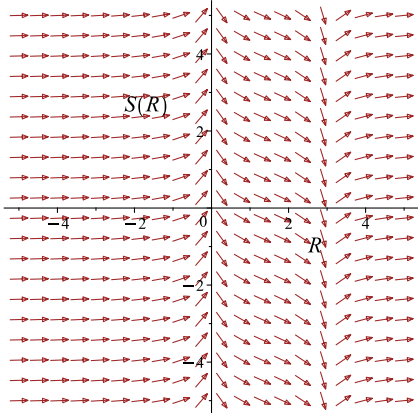
Which simplifies to

$$-\frac{t^2}{2} = -\frac{\ln(y)}{3} + \frac{\ln(-3+y)}{3} + c_1$$

Which gives

$$y = \frac{3e^{\frac{3t^2}{2} + 3c_1}}{-1 + e^{\frac{3t^2}{2} + 3c_1}}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -ty(-3+y)$ 	$R = y$ $S = -\frac{t^2}{2}$	$\frac{dS}{dR} = \frac{1}{R(-3+R)}$ 

Summary

The solution(s) found are the following

$$y = \frac{3e^{\frac{3t^2}{2} + 3c_1}}{-1 + e^{\frac{3t^2}{2} + 3c_1}} \quad (1)$$

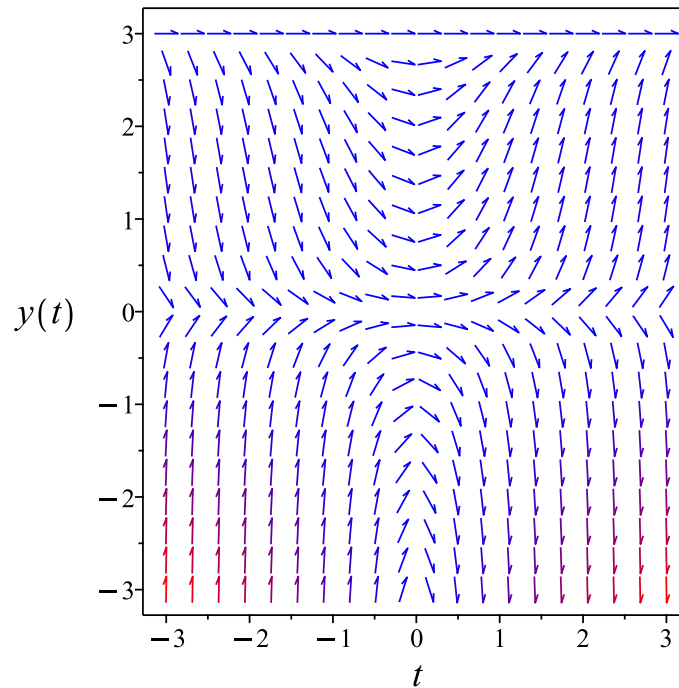


Figure 203: Slope field plot

Verification of solutions

$$y = \frac{3e^{\frac{3t^2}{2} + 3c_1}}{-1 + e^{\frac{3t^2}{2} + 3c_1}}$$

Verified OK.

3.13.3 Solving as bernoulli ode

In canonical form, the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= -ty(-3 + y) \end{aligned}$$

This is a Bernoulli ODE.

$$y' = 3ty - ty^2 \quad (1)$$

The standard Bernoulli ODE has the form

$$y' = f_0(t)y + f_1(t)y^n \quad (2)$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(t)y^{1-n} + f_1(t) \quad (3)$$

The next step is use the substitution $w = y^{1-n}$ in equation (3) which generates a new ODE in $w(t)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(t)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(t) &= 3t \\ f_1(t) &= -t \\ n &= 2 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = y^2$ gives

$$y' \frac{1}{y^2} = \frac{3t}{y} - t \quad (4)$$

Let

$$\begin{aligned} w &= y^{1-n} \\ &= \frac{1}{y} \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t t gives

$$w' = -\frac{1}{y^2}y' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} -w'(t) &= 3w(t)t - t \\ w' &= -3tw + t \end{aligned} \quad (7)$$

The above now is a linear ODE in $w(t)$ which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(t) + p(t)w(t) = q(t)$$

Where here

$$p(t) = 3t$$

$$q(t) = t$$

Hence the ode is

$$w'(t) + 3w(t)t = t$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int 3tdt} \\ &= e^{\frac{3t^2}{2}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu w) &= (\mu)(t) \\ \frac{d}{dt}\left(e^{\frac{3t^2}{2}} w\right) &= \left(e^{\frac{3t^2}{2}}\right)(t) \\ d\left(e^{\frac{3t^2}{2}} w\right) &= \left(t e^{\frac{3t^2}{2}}\right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{\frac{3t^2}{2}} w &= \int t e^{\frac{3t^2}{2}} dt \\ e^{\frac{3t^2}{2}} w &= \frac{e^{\frac{3t^2}{2}}}{3} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{\frac{3t^2}{2}}$ results in

$$w(t) = \frac{e^{-\frac{3t^2}{2}} e^{\frac{3t^2}{2}}}{3} + c_1 e^{-\frac{3t^2}{2}}$$

which simplifies to

$$w(t) = \frac{1}{3} + c_1 e^{-\frac{3t^2}{2}}$$

Replacing w in the above by $\frac{1}{y}$ using equation (5) gives the final solution.

$$\frac{1}{y} = \frac{1}{3} + c_1 e^{-\frac{3t^2}{2}}$$

Or

$$y = \frac{1}{\frac{1}{3} + c_1 e^{-\frac{3t^2}{2}}}$$

Summary

The solution(s) found are the following

$$y = \frac{1}{\frac{1}{3} + c_1 e^{-\frac{3t^2}{2}}} \tag{1}$$

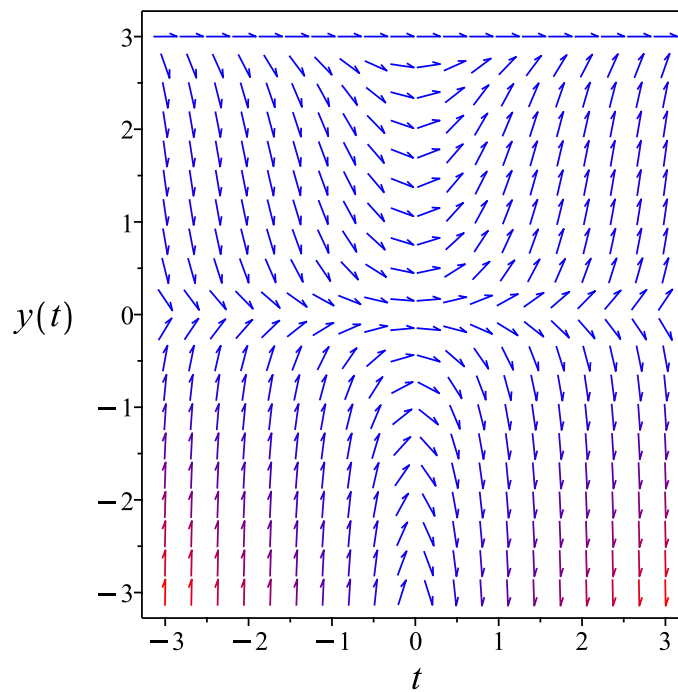


Figure 204: Slope field plot

Verification of solutions

$$y = \frac{1}{\frac{1}{3} + c_1 e^{-\frac{3t^2}{2}}}$$

Verified OK.

3.13.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} \left(-\frac{1}{y(-3+y)} \right) dy &= (t) dt \\ (-t) dt + \left(-\frac{1}{y(-3+y)} \right) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$M(t, y) = -t$$
$$N(t, y) = -\frac{1}{y(-3 + y)}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y}(-t)$$
$$= 0$$

And

$$\frac{\partial N}{\partial t} = \frac{\partial}{\partial t} \left(-\frac{1}{y(-3 + y)} \right)$$
$$= 0$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int M dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int -t dt$$

$$\phi = -\frac{t^2}{2} + f(y) \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{1}{y(-3+y)}$. Therefore equation (4) becomes

$$-\frac{1}{y(-3+y)} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{1}{y(-3+y)}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(-\frac{1}{y(-3+y)} \right) dy$$

$$f(y) = \frac{\ln(y)}{3} - \frac{\ln(-3+y)}{3} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{t^2}{2} + \frac{\ln(y)}{3} - \frac{\ln(-3+y)}{3} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{t^2}{2} + \frac{\ln(y)}{3} - \frac{\ln(-3+y)}{3}$$

The solution becomes

$$y = \frac{3e^{\frac{3t^2}{2} + 3c_1}}{-1 + e^{\frac{3t^2}{2} + 3c_1}}$$

Summary

The solution(s) found are the following

$$y = \frac{3e^{\frac{3t^2}{2} + 3c_1}}{-1 + e^{\frac{3t^2}{2} + 3c_1}} \quad (1)$$

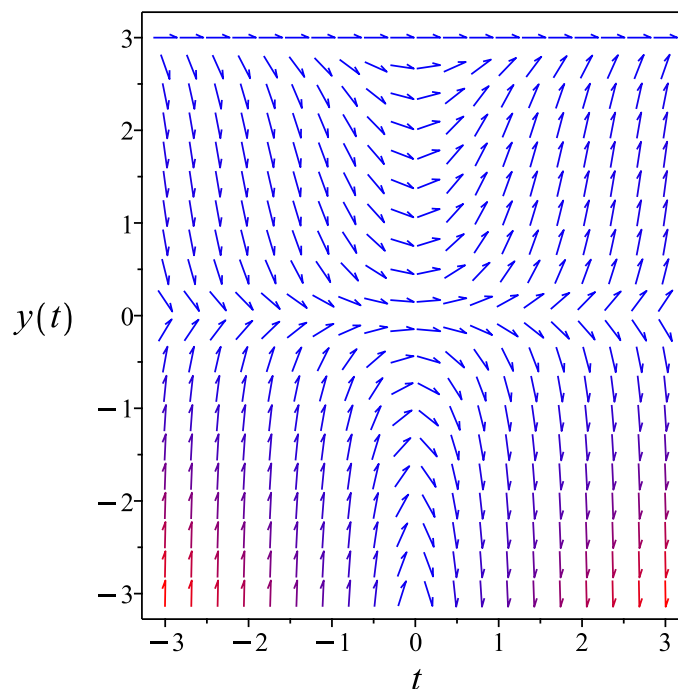


Figure 205: Slope field plot

Verification of solutions

$$y = \frac{3e^{\frac{3t^2}{2} + 3c_1}}{-1 + e^{\frac{3t^2}{2} + 3c_1}}$$

Verified OK.

3.13.5 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= -ty(-3 + y) \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = -ty^2 + 3ty$$

With Riccati ODE standard form

$$y' = f_0(t) + f_1(t)y + f_2(t)y^2$$

Shows that $f_0(t) = 0$, $f_1(t) = 3t$ and $f_2(t) = -t$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{-tu} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(t) - (f_2' + f_1 f_2) u'(t) + f_2^2 f_0 u(t) = 0 \tag{2}$$

But

$$\begin{aligned} f_2' &= -1 \\ f_1 f_2 &= -3t^2 \\ f_2^2 f_0 &= 0 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$-t u''(t) - (-3t^2 - 1) u'(t) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(t) = c_1 + c_2 e^{\frac{3t^2}{2}}$$

The above shows that

$$u'(t) = 3c_2 t e^{\frac{3t^2}{2}}$$

Using the above in (1) gives the solution

$$y = \frac{3c_2 e^{\frac{3t^2}{2}}}{c_1 + c_2 e^{\frac{3t^2}{2}}}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = \frac{3 e^{\frac{3t^2}{2}}}{c_3 + e^{\frac{3t^2}{2}}}$$

Summary

The solution(s) found are the following

$$y = \frac{3e^{\frac{3t^2}{2}}}{c_3 + e^{\frac{3t^2}{2}}} \quad (1)$$

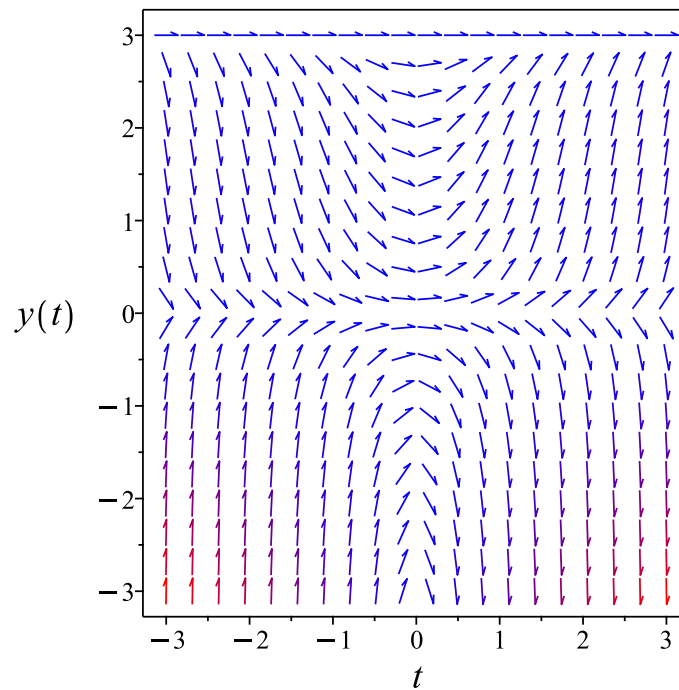


Figure 206: Slope field plot

Verification of solutions

$$y = \frac{3e^{\frac{3t^2}{2}}}{c_3 + e^{\frac{3t^2}{2}}}$$

Verified OK.

3.13.6 Maple step by step solution

Let's solve

$$y' - t(3 - y)y = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'}{(3-y)y} = t$$

- Integrate both sides with respect to t

$$\int \frac{y'}{(3-y)y} dt = \int t dt + c_1$$

- Evaluate integral

$$\frac{\ln(y)}{3} - \frac{\ln(-3+y)}{3} = \frac{t^2}{2} + c_1$$

- Solve for y

$$y = \frac{3e^{\frac{3t^2}{2} + 3c_1}}{-1 + e^{\frac{3t^2}{2} + 3c_1}}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(diff(y(t),t) = t*(3-y(t))*y(t),y(t), singsol=all)
```

$$y(t) = \frac{3}{1 + 3e^{-\frac{3t^2}{2}} c_1}$$

✓ Solution by Mathematica

Time used: 0.266 (sec). Leaf size: 44

```
DSolve[y'[t] == t*(3-y[t])*y[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{3e^{\frac{3t^2}{2}}}{e^{\frac{3t^2}{2}} + e^{3c_1}}$$

$$y(t) \rightarrow 0$$

$$y(t) \rightarrow 3$$

3.14 problem 18

3.14.1 Solving as first order ode lie symmetry lookup ode	1085
3.14.2 Solving as bernoulli ode	1089
3.14.3 Solving as riccati ode	1093

Internal problem ID [529]

Internal file name [OUTPUT/529_Sunday_June_05_2022_01_43_07_AM_62544468/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 18.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**riccati**", "**bernoulli**", "**first_order_ode_lie_symmetry_lookup**"

Maple gives the following as the ode type

[_Bernoulli]

$$y' - y(3 - yt) = 0$$

3.14.1 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -y(ty - 3)$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type **Bernoulli**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 219: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= y^2 e^{-3t}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{y^2 e^{-3t}} dy \end{aligned}$$

Which results in

$$S = -\frac{e^{3t}}{y}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -y(ty - 3)$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= -\frac{3e^{3t}}{y} \\ S_y &= \frac{e^{3t}}{y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -e^{3t}t \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -e^{3R}R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{(3R - 1)e^{3R}}{9} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$-\frac{e^{3t}}{y} = -\frac{(-1 + 3t)e^{3t}}{9} + c_1$$

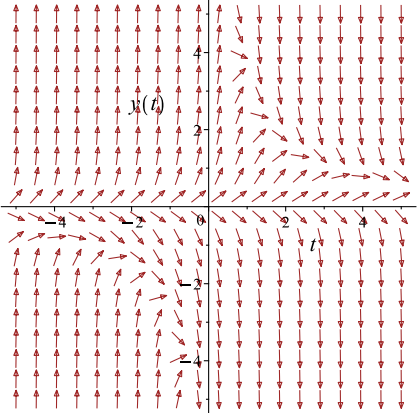
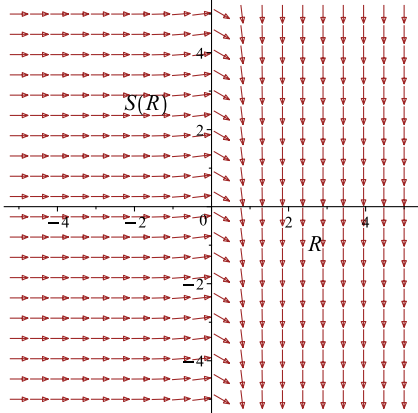
Which simplifies to

$$-\frac{e^{3t}}{y} = -\frac{(-1 + 3t)e^{3t}}{9} + c_1$$

Which gives

$$y = \frac{9e^{3t}}{3e^{3t}t - e^{3t} - 9c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -y(ty - 3)$ 	$R = t$ $S = -\frac{e^{3t}}{y}$	$\frac{dS}{dR} = -e^{3R}R$ 

Summary

The solution(s) found are the following

$$y = \frac{9e^{3t}}{3e^{3t}t - e^{3t} - 9c_1} \quad (1)$$

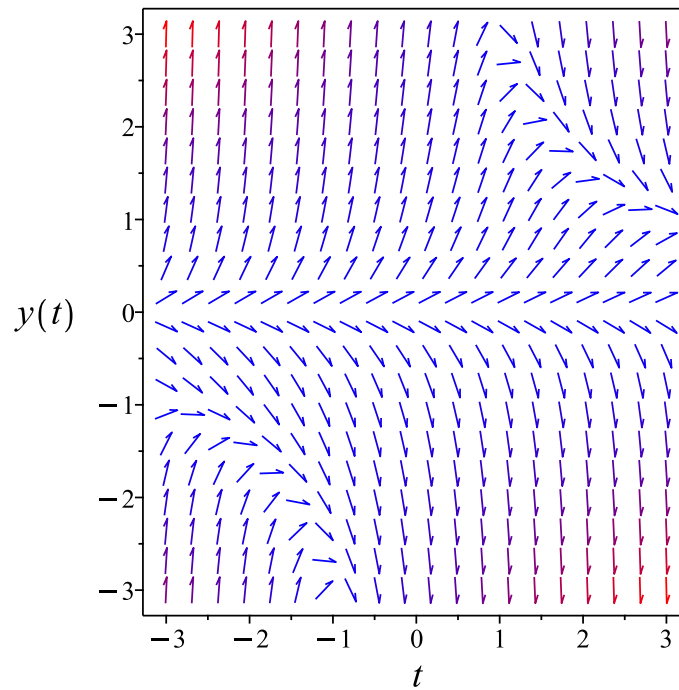


Figure 207: Slope field plot

Verification of solutions

$$y = \frac{9e^{3t}}{3e^{3t}t - e^{3t} - 9c_1}$$

Verified OK.

3.14.2 Solving as bernoulli ode

In canonical form, the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= -y(ty - 3) \end{aligned}$$

This is a Bernoulli ODE.

$$y' = 3y - ty^2 \quad (1)$$

The standard Bernoulli ODE has the form

$$y' = f_0(t)y + f_1(t)y^n \quad (2)$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(t)y^{1-n} + f_1(t) \quad (3)$$

The next step is use the substitution $w = y^{1-n}$ in equation (3) which generates a new ODE in $w(t)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(t)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(t) &= 3 \\ f_1(t) &= -t \\ n &= 2 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = y^2$ gives

$$y' \frac{1}{y^2} = \frac{3}{y} - t \quad (4)$$

Let

$$\begin{aligned} w &= y^{1-n} \\ &= \frac{1}{y} \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t t gives

$$w' = -\frac{1}{y^2}y' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} -w'(t) &= 3w(t) - t \\ w' &= -3w + t \end{aligned} \quad (7)$$

The above now is a linear ODE in $w(t)$ which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(t) + p(t)w(t) = q(t)$$

Where here

$$\begin{aligned}p(t) &= 3 \\q(t) &= t\end{aligned}$$

Hence the ode is

$$w'(t) + 3w(t) = t$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int 3dt} \\ &= e^{3t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu w) &= (\mu)(t) \\ \frac{d}{dt}(e^{3t}w) &= (e^{3t})(t) \\ d(e^{3t}w) &= (e^{3t}t) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{3t}w &= \int e^{3t}t dt \\ e^{3t}w &= \frac{(-1 + 3t)e^{3t}}{9} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{3t}$ results in

$$w(t) = \frac{e^{-3t}(-1 + 3t)e^{3t}}{9} + c_1e^{-3t}$$

which simplifies to

$$w(t) = \frac{t}{3} - \frac{1}{9} + c_1e^{-3t}$$

Replacing w in the above by $\frac{1}{y}$ using equation (5) gives the final solution.

$$\frac{1}{y} = \frac{t}{3} - \frac{1}{9} + c_1e^{-3t}$$

Or

$$y = \frac{1}{\frac{t}{3} - \frac{1}{9} + c_1e^{-3t}}$$

Which is simplified to

$$y = \frac{9}{9c_1e^{-3t} + 3t - 1}$$

Summary

The solution(s) found are the following

$$y = \frac{9}{9c_1e^{-3t} + 3t - 1} \tag{1}$$

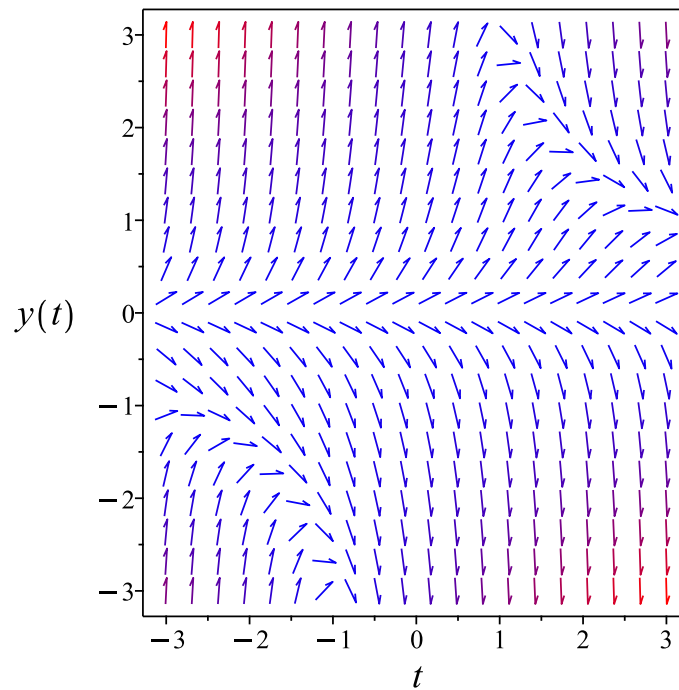


Figure 208: Slope field plot

Verification of solutions

$$y = \frac{9}{9c_1e^{-3t} + 3t - 1}$$

Verified OK.

3.14.3 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(t, y) \\ &= -y(ty - 3)\end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = -ty^2 + 3y$$

With Riccati ODE standard form

$$y' = f_0(t) + f_1(t)y + f_2(t)y^2$$

Shows that $f_0(t) = 0$, $f_1(t) = 3$ and $f_2(t) = -t$. Let

$$\begin{aligned}y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{-tu}\end{aligned}\tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(t) - (f_2' + f_1 f_2) u'(t) + f_2^2 f_0 u(t) = 0\tag{2}$$

But

$$\begin{aligned}f_2' &= -1 \\ f_1 f_2 &= -3t \\ f_2^2 f_0 &= 0\end{aligned}$$

Substituting the above terms back in equation (2) gives

$$-tu''(t) - (-3t - 1)u'(t) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(t) = c_1 + (-1 + 3t)e^{3t}c_2$$

The above shows that

$$u'(t) = 9c_2 e^{3t}t$$

Using the above in (1) gives the solution

$$y = \frac{9c_2 e^{3t}}{c_1 + (-1 + 3t) e^{3t} c_2}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = \frac{9 e^{3t}}{c_3 + (-1 + 3t) e^{3t}}$$

Summary

The solution(s) found are the following

$$y = \frac{9 e^{3t}}{c_3 + (-1 + 3t) e^{3t}} \tag{1}$$

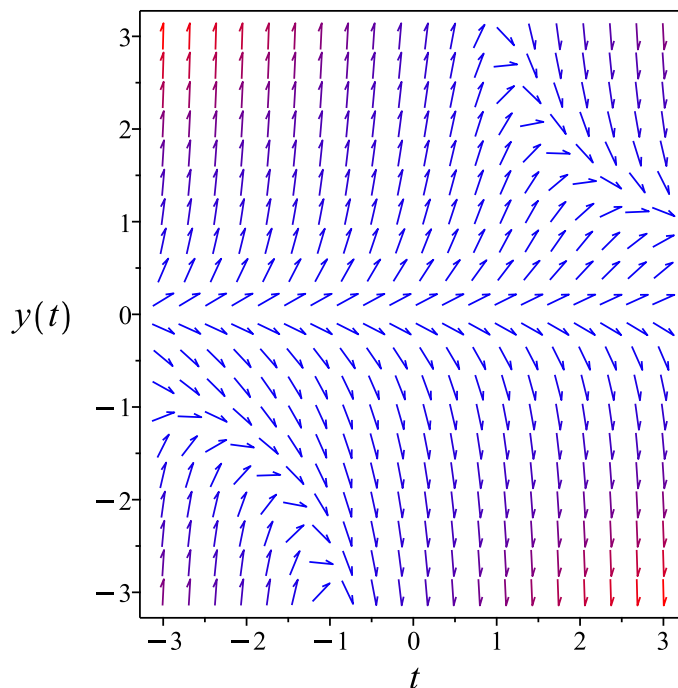


Figure 209: Slope field plot

Verification of solutions

$$y = \frac{9 e^{3t}}{c_3 + (-1 + 3t) e^{3t}}$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

```
dsolve(diff(y(t),t) = y(t)*(3-t*y(t)),y(t), singsol=all)
```

$$y(t) = \frac{9}{-1 + 9c_1e^{-3t} + 3t}$$

✓ Solution by Mathematica

Time used: 0.118 (sec). Leaf size: 35

```
DSolve[y'[t] == y[t]*(3-t*y[t]),y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{9e^{3t}}{e^{3t}(3t - 1) + 9c_1}$$
$$y(t) \rightarrow 0$$

3.15 problem 19

3.15.1 Solving as first order ode lie symmetry lookup ode	1096
3.15.2 Solving as bernoulli ode	1100
3.15.3 Solving as riccati ode	1104

Internal problem ID [530]

Internal file name [OUTPUT/530_Sunday_June_05_2022_01_43_08_AM_17571426/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 19.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "riccati", "bernoulli", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_Bernoulli]

$$y' + y(3 - yt) = 0$$

3.15.1 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = y(ty - 3)$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type `Bernoulli`. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 221: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= y^2 e^{3t}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{y^2 e^{3t}} dy \end{aligned}$$

Which results in

$$S = -\frac{e^{-3t}}{y}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = y(ty - 3)$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= \frac{3e^{-3t}}{y} \\ S_y &= \frac{e^{-3t}}{y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = t e^{-3t} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R e^{-3R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{(3R + 1)e^{-3R}}{9} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$-\frac{e^{-3t}}{y} = -\frac{(1 + 3t)e^{-3t}}{9} + c_1$$

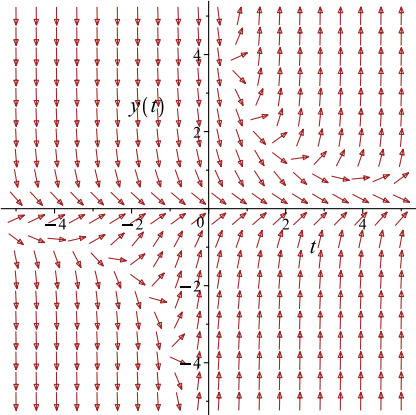
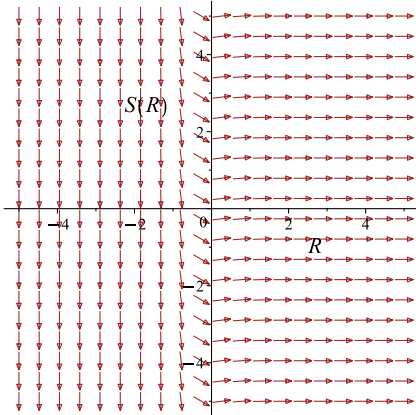
Which simplifies to

$$-\frac{e^{-3t}}{y} = -\frac{(1 + 3t)e^{-3t}}{9} + c_1$$

Which gives

$$y = \frac{9e^{-3t}}{3te^{-3t} + e^{-3t} - 9c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = y(ty - 3)$ 	$R = t$ $S = -\frac{e^{-3t}}{y}$	$\frac{dS}{dR} = R e^{-3R}$ 

Summary

The solution(s) found are the following

$$y = \frac{9e^{-3t}}{3te^{-3t} + e^{-3t} - 9c_1} \quad (1)$$

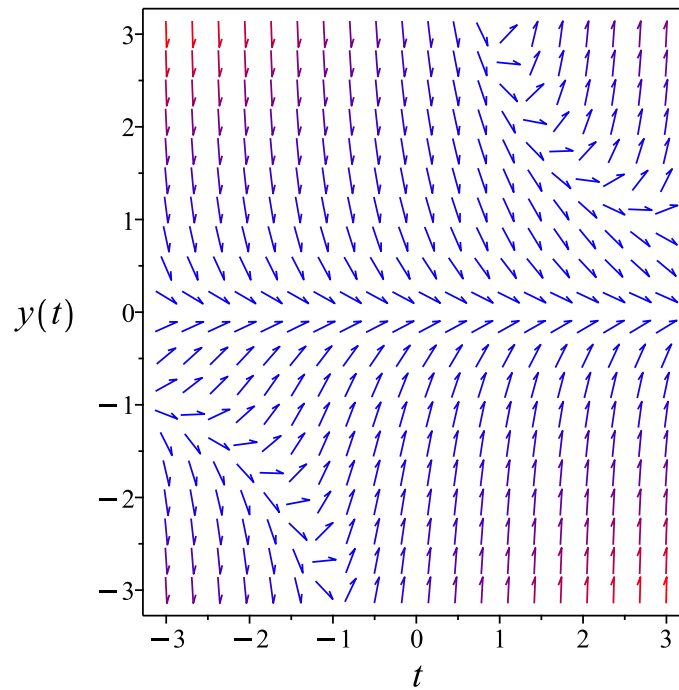


Figure 210: Slope field plot

Verification of solutions

$$y = \frac{9e^{-3t}}{3te^{-3t} + e^{-3t} - 9c_1}$$

Verified OK.

3.15.2 Solving as bernoulli ode

In canonical form, the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= y(ty - 3) \end{aligned}$$

This is a Bernoulli ODE.

$$y' = -3y + ty^2 \quad (1)$$

The standard Bernoulli ODE has the form

$$y' = f_0(t)y + f_1(t)y^n \quad (2)$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(t)y^{1-n} + f_1(t) \quad (3)$$

The next step is use the substitution $w = y^{1-n}$ in equation (3) which generates a new ODE in $w(t)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(t)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(t) &= -3 \\ f_1(t) &= t \\ n &= 2 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = y^2$ gives

$$y' \frac{1}{y^2} = -\frac{3}{y} + t \quad (4)$$

Let

$$\begin{aligned} w &= y^{1-n} \\ &= \frac{1}{y} \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t t gives

$$w' = -\frac{1}{y^2}y' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} -w'(t) &= -3w(t) + t \\ w' &= 3w - t \end{aligned} \quad (7)$$

The above now is a linear ODE in $w(t)$ which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(t) + p(t)w(t) = q(t)$$

Where here

$$\begin{aligned}p(t) &= -3 \\q(t) &= -t\end{aligned}$$

Hence the ode is

$$w'(t) - 3w(t) = -t$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int(-3)dt} \\ &= e^{-3t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu w) &= (\mu)(-t) \\ \frac{d}{dt}(e^{-3t}w) &= (e^{-3t})(-t) \\ d(e^{-3t}w) &= (-te^{-3t}) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-3t}w &= \int -te^{-3t} dt \\ e^{-3t}w &= \frac{(1+3t)e^{-3t}}{9} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-3t}$ results in

$$w(t) = \frac{e^{3t}(1+3t)e^{-3t}}{9} + c_1e^{3t}$$

which simplifies to

$$w(t) = \frac{1}{9} + \frac{t}{3} + c_1e^{3t}$$

Replacing w in the above by $\frac{1}{y}$ using equation (5) gives the final solution.

$$\frac{1}{y} = \frac{1}{9} + \frac{t}{3} + c_1e^{3t}$$

Or

$$y = \frac{1}{\frac{1}{9} + \frac{t}{3} + c_1e^{3t}}$$

Which is simplified to

$$y = \frac{9}{9c_1e^{3t} + 3t + 1}$$

Summary

The solution(s) found are the following

$$y = \frac{9}{9c_1e^{3t} + 3t + 1} \tag{1}$$

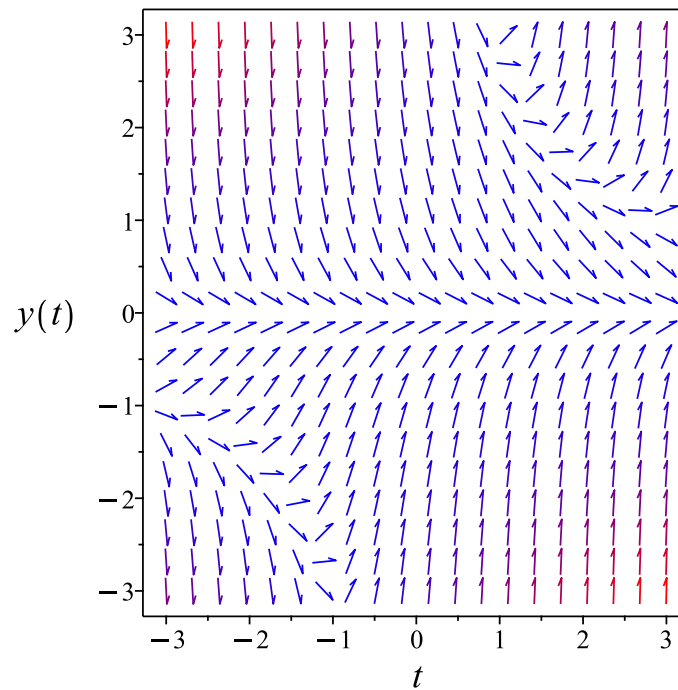


Figure 211: Slope field plot

Verification of solutions

$$y = \frac{9}{9c_1e^{3t} + 3t + 1}$$

Verified OK.

3.15.3 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(t, y) \\ &= y(ty - 3)\end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = ty^2 - 3y$$

With Riccati ODE standard form

$$y' = f_0(t) + f_1(t)y + f_2(t)y^2$$

Shows that $f_0(t) = 0$, $f_1(t) = -3$ and $f_2(t) = t$. Let

$$\begin{aligned}y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{tu}\end{aligned}\tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(t) - (f_2' + f_1 f_2) u'(t) + f_2^2 f_0 u(t) = 0\tag{2}$$

But

$$\begin{aligned}f_2' &= 1 \\ f_1 f_2 &= -3t \\ f_2^2 f_0 &= 0\end{aligned}$$

Substituting the above terms back in equation (2) gives

$$tu''(t) - (1 - 3t) u'(t) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(t) = c_1 + (1 + 3t) e^{-3t} c_2$$

The above shows that

$$u'(t) = -9c_2 e^{-3t} t$$

Using the above in (1) gives the solution

$$y = \frac{9c_2 e^{-3t}}{c_1 + (1 + 3t) e^{-3t} c_2}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = \frac{9 e^{-3t}}{3t e^{-3t} + e^{-3t} + c_3}$$

Summary

The solution(s) found are the following

$$y = \frac{9 e^{-3t}}{3t e^{-3t} + e^{-3t} + c_3} \quad (1)$$

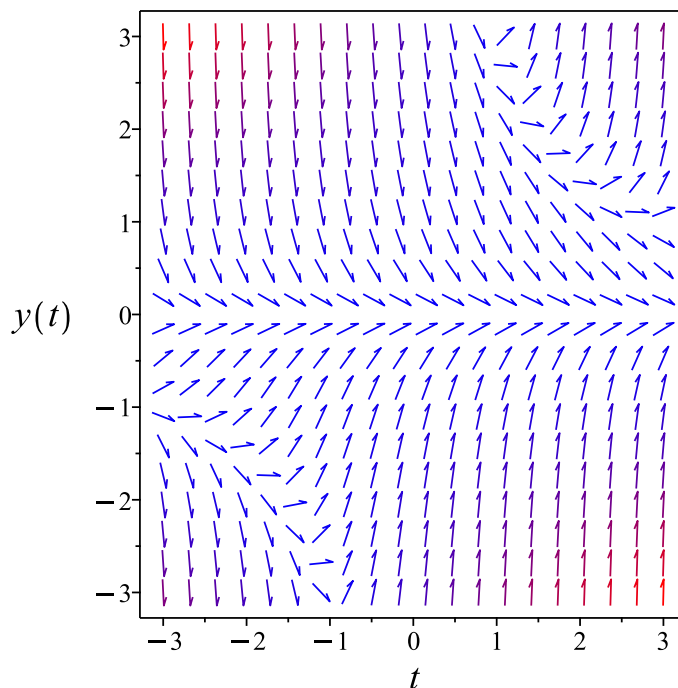


Figure 212: Slope field plot

Verification of solutions

$$y = \frac{9 e^{-3t}}{3t e^{-3t} + e^{-3t} + c_3}$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

```
dsolve(diff(y(t),t) = -y(t)*(3-t*y(t)),y(t), singsol=all)
```

$$y(t) = \frac{9}{1 + 9c_1e^{3t} + 3t}$$

✓ Solution by Mathematica

Time used: 0.126 (sec). Leaf size: 28

```
DSolve[y'[t] == -y[t]*(3-t*y[t]),y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{9}{3t + 9c_1e^{3t} + 1}$$
$$y(t) \rightarrow 0$$

3.16 problem 20

3.16.1 Solving as riccati ode 1107

Internal problem ID [531]

Internal file name [OUTPUT/531_Sunday_June_05_2022_01_43_09_AM_59273030/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.4. Page 76

Problem number: 20.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**riccati**"

Maple gives the following as the ode type

[_Riccati]

$$y' + y^2 = -1 + t$$

3.16.1 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(t, y) \\ &= -y^2 + t - 1 \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = -y^2 + t - 1$$

With Riccati ODE standard form

$$y' = f_0(t) + f_1(t)y + f_2(t)y^2$$

Shows that $f_0(t) = -1 + t$, $f_1(t) = 0$ and $f_2(t) = -1$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{-u} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(t) - (f_2' + f_1 f_2) u'(t) + f_2^2 f_0 u(t) = 0 \quad (2)$$

But

$$\begin{aligned} f_2' &= 0 \\ f_1 f_2 &= 0 \\ f_2^2 f_0 &= -1 + t \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$-u''(t) + (-1 + t) u(t) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(t) = c_1 \text{AiryAi}(-1 + t) + c_2 \text{AiryBi}(-1 + t)$$

The above shows that

$$u'(t) = c_1 \text{AiryAi}(1, -1 + t) + c_2 \text{AiryBi}(1, -1 + t)$$

Using the above in (1) gives the solution

$$y = \frac{c_1 \text{AiryAi}(1, -1 + t) + c_2 \text{AiryBi}(1, -1 + t)}{c_1 \text{AiryAi}(-1 + t) + c_2 \text{AiryBi}(-1 + t)}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = \frac{c_3 \text{AiryAi}(1, -1 + t) + \text{AiryBi}(1, -1 + t)}{c_3 \text{AiryAi}(-1 + t) + \text{AiryBi}(-1 + t)}$$

Summary

The solution(s) found are the following

$$y = \frac{c_3 \text{AiryAi}(1, -1 + t) + \text{AiryBi}(1, -1 + t)}{c_3 \text{AiryAi}(-1 + t) + \text{AiryBi}(-1 + t)} \quad (1)$$

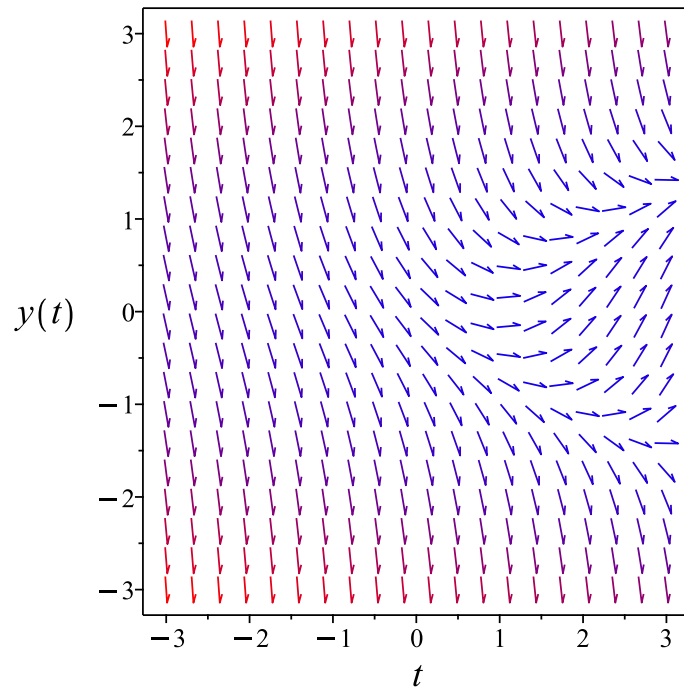


Figure 213: Slope field plot

Verification of solutions

$$y = \frac{c_3 \text{AiryAi}(1, -1 + t) + \text{AiryBi}(1, -1 + t)}{c_3 \text{AiryAi}(-1 + t) + \text{AiryBi}(-1 + t)}$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
Looking for potential symmetries  
trying Riccati  
trying Riccati Special  
trying Riccati sub-methods:  
  <- Abel AIR successful: ODE belongs to the OF1 0-parameter (Airy type) class`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 31

```
dsolve(diff(y(t),t) = t-1-y(t)^2,y(t), singsol=all)
```

$$y(t) = \frac{\text{AiryAi}(1, t - 1) c_1 + \text{AiryBi}(1, t - 1)}{\text{AiryAi}(t - 1) c_1 + \text{AiryBi}(t - 1)}$$

✓ Solution by Mathematica

Time used: 0.123 (sec). Leaf size: 47

```
DSolve[y'[t] == t-1-y[t]^2,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{\text{AiryBiPrime}(t - 1) + c_1 \text{AiryAiPrime}(t - 1)}{\text{AiryBi}(t - 1) + c_1 \text{AiryAi}(t - 1)}$$
$$y(t) \rightarrow \frac{\text{AiryAiPrime}(t - 1)}{\text{AiryAi}(t - 1)}$$

4 Section 2.5. Page 88

4.1	problem 1	1112
4.2	problem 3	1115
4.3	problem 4	1119
4.4	problem 5	1122
4.5	problem 6	1125
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4.1 problem 1

- 4.1.1 Solving as quadrature ode 1112
- 4.1.2 Maple step by step solution 1113

Internal problem ID [532]

Internal file name [OUTPUT/532_Sunday_June_05_2022_01_43_10_AM_54801064/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"quadrature"**

Maple gives the following as the ode type

`[_quadrature]`

$$y' - ay - by^2 = 0$$

4.1.1 Solving as quadrature ode

Integrating both sides gives

$$\int \frac{1}{by^2 + ay} dy = \int dx$$
$$\frac{\ln(y)}{a} - \frac{\ln(by + a)}{a} = x + c_1$$

The above can be written as

$$\left(\frac{1}{a}\right) (\ln(y) - \ln(by + a)) = x + c_1$$
$$\ln(y) - \ln(by + a) = (a)(x + c_1)$$
$$= a(x + c_1)$$

Raising both side to exponential gives

$$e^{\ln(y) - \ln(by + a)} = ac_1 e^{ax}$$

Which simplifies to

$$\frac{y}{by + a} = c_2 e^{ax}$$

Summary

The solution(s) found are the following

$$y = -\frac{e^{ax} c_2 a}{-1 + e^{ax} c_2 b} \quad (1)$$

Verification of solutions

$$y = -\frac{e^{ax} c_2 a}{-1 + e^{ax} c_2 b}$$

Verified OK.

4.1.2 Maple step by step solution

Let's solve

$$y' - ay - by^2 = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{ay+by^2} = 1$$

- Integrate both sides with respect to x

$$\int \frac{y'}{ay+by^2} dx = \int 1 dx + c_1$$

- Evaluate integral

$$\frac{\ln(y)}{a} - \frac{\ln(by+a)}{a} = x + c_1$$

- Solve for y

$$y = -\frac{e^{c_1 a + ax} a}{-1 + b e^{c_1 a + ax}}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

```
dsolve(diff(y(x),x) = a*y(x)+b*y(x)^2,y(x), singsol=all)
```

$$y(x) = \frac{a}{e^{-ax}c_1a - b}$$

✓ Solution by Mathematica

Time used: 0.888 (sec). Leaf size: 45

```
DSolve[y'[x]== a*y[x]+b*y[x]^2,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{ae^{a(x+c_1)}}{-1 + be^{a(x+c_1)}}$$
$$y(x) \rightarrow 0$$
$$y(x) \rightarrow -\frac{a}{b}$$

4.2 problem 3

4.2.1 Solving as quadrature ode	1115
4.2.2 Maple step by step solution	1116

Internal problem ID [533]

Internal file name [OUTPUT/533_Sunday_June_05_2022_01_43_11_AM_86686196/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 3.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"quadrature"**

Maple gives the following as the ode type

[_quadrature]

$$y' - y(-2 + y)(-1 + y) = 0$$

4.2.1 Solving as quadrature ode

Integrating both sides gives

$$\int \frac{1}{y(y-2)(y-1)} dy = \int dt$$
$$\frac{\ln(y)}{2} - \ln(y-1) + \frac{\ln(y-2)}{2} = t + c_1$$

Raising both side to exponential gives

$$e^{\frac{\ln(y)}{2} - \ln(y-1) + \frac{\ln(y-2)}{2}} = e^{t+c_1}$$

Which simplifies to

$$\frac{\sqrt{y} \sqrt{y-2}}{y-1} = c_2 e^t$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{-c_2^2 e^{2t} + 1} - 1}{\sqrt{-c_2^2 e^{2t} + 1}} \quad (1)$$

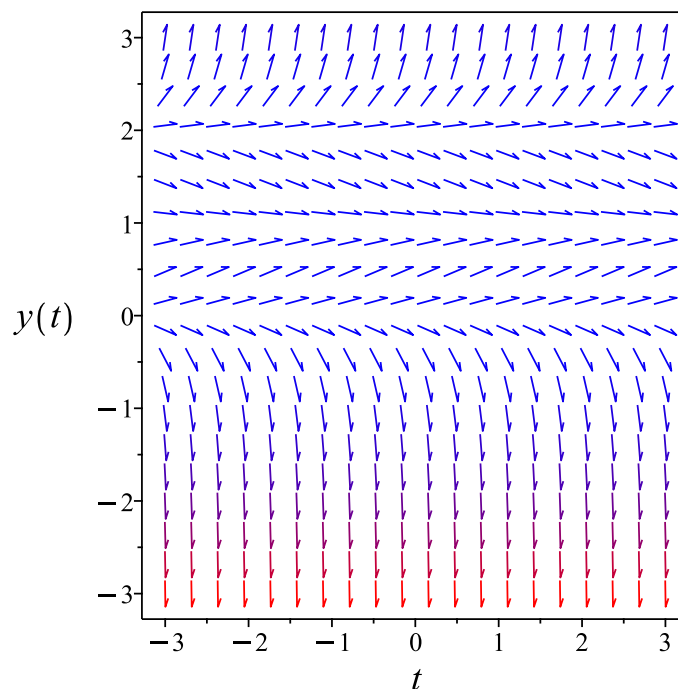


Figure 214: Slope field plot

Verification of solutions

$$y = \frac{\sqrt{-c_2^2 e^{2t} + 1} - 1}{\sqrt{-c_2^2 e^{2t} + 1}}$$

Verified OK.

4.2.2 Maple step by step solution

Let's solve

$$y' - y(-2 + y)(-1 + y) = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'}{y(-2+y)(-1+y)} = 1$$

- Integrate both sides with respect to t

$$\int \frac{y'}{y(-2+y)(-1+y)} dt = \int 1 dt + c_1$$

- Evaluate integral

$$\frac{\ln(y)}{2} - \ln(-1 + y) + \frac{\ln(-2+y)}{2} = t + c_1$$

- Solve for y

$$\left\{ y = \frac{(e^{c_1})^2 (e^t)^2}{(e^{c_1})^2 (e^t)^2 - \sqrt{-(e^{c_1})^2 (e^t)^2 + 1} - 1}, y = \frac{(e^{c_1})^2 (e^t)^2}{(e^{c_1})^2 (e^t)^2 + \sqrt{-(e^{c_1})^2 (e^t)^2 + 1} - 1} \right\}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.125 (sec). Leaf size: 75

```
dsolve(diff(y(t),t) = y(t)*(-2+y(t))*(-1+y(t)),y(t), singsol=all)
```

$$y(t) = \frac{e^{2t}c_1}{(-1 - \sqrt{-c_1e^{2t} + 1})\sqrt{-c_1e^{2t} + 1}}$$

$$y(t) = \frac{e^{2t}c_1}{(1 - \sqrt{-c_1e^{2t} + 1})\sqrt{-c_1e^{2t} + 1}}$$

✓ Solution by Mathematica

Time used: 11.055 (sec). Leaf size: 100

```
DSolve[y'[t] == y[t]*(-2+y[t])*(-1+y[t]),y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{-\sqrt{1 + e^{2(t+c_1)}} + e^{2(t+c_1)} + 1}{1 + e^{2(t+c_1)}}$$

$$y(t) \rightarrow \frac{\sqrt{1 + e^{2(t+c_1)}} + e^{2(t+c_1)} + 1}{1 + e^{2(t+c_1)}}$$

$$y(t) \rightarrow 0$$

$$y(t) \rightarrow 1$$

$$y(t) \rightarrow 2$$

4.3 problem 4

4.3.1 Solving as quadrature ode	1119
4.3.2 Maple step by step solution	1120

Internal problem ID [534]

Internal file name [OUTPUT/534_Sunday_June_05_2022_01_43_47_AM_54468129/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 4.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**quadrature**"

Maple gives the following as the ode type

[_quadrature]

$$y' - e^y = -1$$

4.3.1 Solving as quadrature ode

Integrating both sides gives

$$\int \frac{1}{-1 + e^y} dy = \int dt$$
$$-\ln(e^y) + \ln(-1 + e^y) = t + c_1$$

Raising both side to exponential gives

$$e^{-\ln(e^y) + \ln(-1 + e^y)} = e^{t + c_1}$$

Which simplifies to

$$-e^{-y} + 1 = c_2 e^t$$

Summary

The solution(s) found are the following

$$y = -\ln(-c_2 e^t + 1) \tag{1}$$

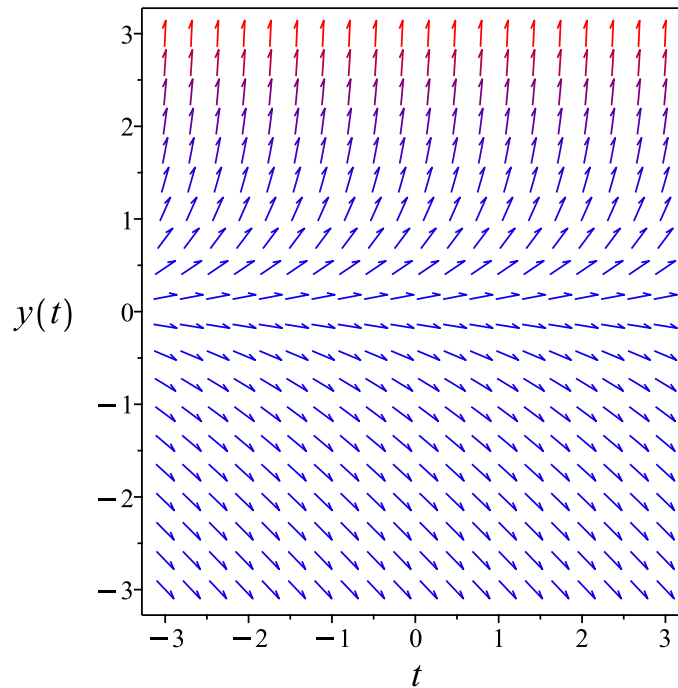


Figure 215: Slope field plot

Verification of solutions

$$y = -\ln(-c_2 e^t + 1)$$

Verified OK.

4.3.2 Maple step by step solution

Let's solve

$$y' - e^y = -1$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{-1+e^y} = 1$$

- Integrate both sides with respect to t

$$\int \frac{y'}{-1+e^y} dt = \int 1 dt + c_1$$

- Evaluate integral

$$-\ln(e^y) + \ln(-1 + e^y) = t + c_1$$

- Solve for y

$$y = \ln\left(-\frac{1}{e^{t+c_1}-1}\right)$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve(diff(y(t),t) = -1+exp(y(t)),y(t), singsol=all)
```

$$y(t) = \ln\left(-\frac{1}{e^{t+c_1}-1}\right)$$

✓ Solution by Mathematica

Time used: 0.817 (sec). Leaf size: 28

```
DSolve[y'[t]== -1+Exp[y[t]],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \log\left(\frac{1}{2}\left(1 - \tanh\left(\frac{t+c_1}{2}\right)\right)\right)$$

$$y(t) \rightarrow 0$$

4.4 problem 5

4.4.1 Solving as quadrature ode	1122
4.4.2 Maple step by step solution	1123

Internal problem ID [535]

Internal file name [OUTPUT/535_Sunday_June_05_2022_01_43_49_AM_20194512/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 5.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"quadrature"**

Maple gives the following as the ode type

`[_quadrature]`

$$y' - e^{-y} = -1$$

4.4.1 Solving as quadrature ode

Integrating both sides gives

$$\int \frac{1}{-1 + e^{-y}} dy = \int dt$$
$$\ln(e^{-y}) - \ln(-1 + e^{-y}) = t + c_1$$

Raising both side to exponential gives

$$e^{\ln(e^{-y}) - \ln(-1 + e^{-y})} = e^{t+c_1}$$

Which simplifies to

$$\frac{e^{-y}}{-1 + e^{-y}} = c_2 e^t$$

Summary

The solution(s) found are the following

$$y = -\ln\left(\frac{c_2}{c_2 e^t - 1}\right) - t \tag{1}$$

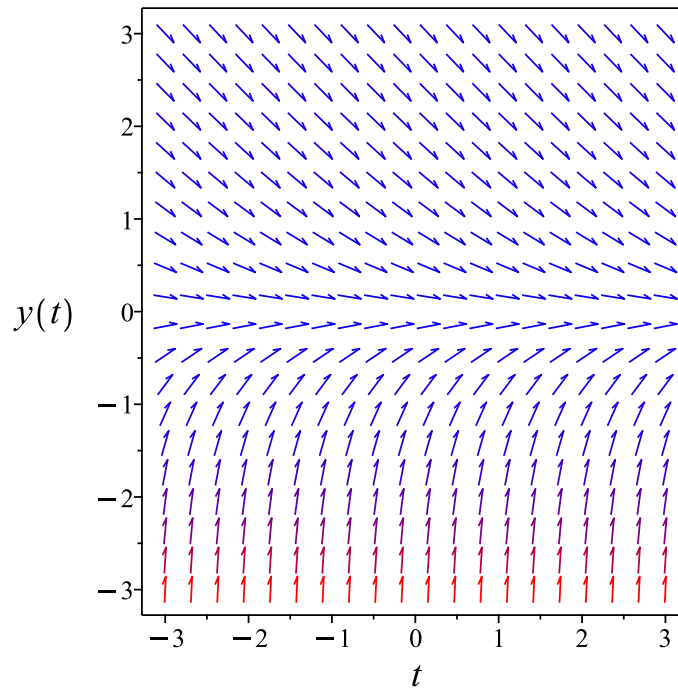


Figure 216: Slope field plot

Verification of solutions

$$y = -\ln\left(\frac{c_2}{c_2 e^t - 1}\right) - t$$

Verified OK.

4.4.2 Maple step by step solution

Let's solve

$$y' - e^{-y} = -1$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{-1+e^{-y}} = 1$$

- Integrate both sides with respect to t

$$\int \frac{y'}{-1+e^{-y}} dt = \int 1 dt + c_1$$

- Evaluate integral

$$\ln(e^{-y}) - \ln(-1 + e^{-y}) = t + c_1$$

- Solve for y

$$y = \ln(e^{t+c_1} - 1) - t - c_1$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 18

```
dsolve(diff(y(t),t) = -1+exp(-y(t)),y(t), singsol=all)
```

$$y(t) = -t + \ln(e^{t+c_1} - 1) - c_1$$

✓ Solution by Mathematica

Time used: 0.853 (sec). Leaf size: 21

```
DSolve[y'[t] == -1+Exp[-y[t]],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \log(1 + e^{-t+c_1})$$

$$y(t) \rightarrow 0$$

4.5 problem 6

4.5.1 Solving as quadrature ode	1125
4.5.2 Maple step by step solution	1126

Internal problem ID [536]

Internal file name [OUTPUT/536_Sunday_June_05_2022_01_43_50_AM_22554720/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 6.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"quadrature"**

Maple gives the following as the ode type

`[_quadrature]`

$$y' + \frac{2 \arctan(y)}{1 + y^2} = 0$$

4.5.1 Solving as quadrature ode

Integrating both sides gives

$$\int -\frac{y^2 + 1}{2 \arctan(y)} dy = \int dt$$
$$\int^y -\frac{-a^2 + 1}{2 \arctan(-a)} d-a = t + c_1$$

Summary

The solution(s) found are the following

$$\int^y -\frac{-a^2 + 1}{2 \arctan(-a)} d-a = t + c_1 \tag{1}$$

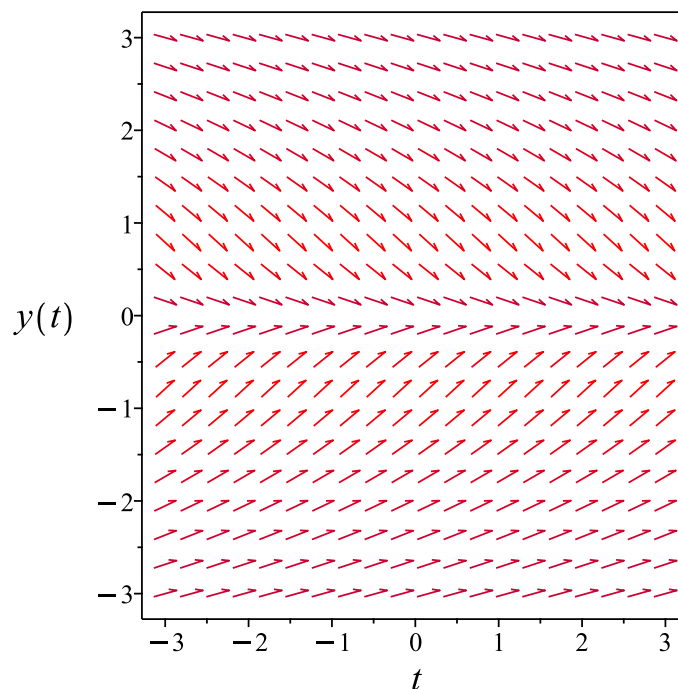


Figure 217: Slope field plot

Verification of solutions

$$\int^y -\frac{-a^2 + 1}{2 \arctan(-a)} d_a = t + c_1$$

Verified OK.

4.5.2 Maple step by step solution

Let's solve

$$y' + \frac{2 \arctan(y)}{1+y^2} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{(1+y^2)y'}{\arctan(y)} = -2$$

- Integrate both sides with respect to t

$$\int \frac{(1+y^2)y'}{\arctan(y)} dt = \int (-2) dt + c_1$$

- Cannot compute integral

$$\int \frac{(1+y^2)y'}{\arctan(y)} dt = -2t + c_1$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 23

```
dsolve(diff(y(t),t) = -2*arctan(y(t))/(1+y(t)^2),y(t), singsol=all)
```

$$t + \frac{\left(\int^{y(t)} \frac{a^2+1}{\arctan(a)} da \right)}{2} + c_1 = 0$$

✓ Solution by Mathematica

Time used: 1.013 (sec). Leaf size: 38

```
DSolve[y'[t] == -2*ArcTan[y[t]]/(1+y[t]^2),y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \text{InverseFunction} \left[\int_1^{\#1} \frac{K[1]^2 + 1}{\arctan(K[1])} dK[1] \& \right] [-2t + c_1]$$

$$y(t) \rightarrow 0$$

4.6 problem 7

4.6.1 Solving as quadrature ode	1128
4.6.2 Maple step by step solution	1129

Internal problem ID [537]

Internal file name [OUTPUT/537_Sunday_June_05_2022_01_43_52_AM_26451284/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 7.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"quadrature"**

Maple gives the following as the ode type

[_quadrature]

$$y' + k(-1 + y)^2 = 0$$

4.6.1 Solving as quadrature ode

Integrating both sides gives

$$\int -\frac{1}{k(y-1)^2} dy = t + c_1$$
$$\frac{1}{k(y-1)} = t + c_1$$

Solving for y gives these solutions

$$y_1 = \frac{c_1 k + t k + 1}{k(t + c_1)}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 k + t k + 1}{k(t + c_1)} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 k + tk + 1}{k(t + c_1)}$$

Verified OK.

4.6.2 Maple step by step solution

Let's solve

$$y' + k(-1 + y)^2 = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{(-1+y)^2} = -k$$

- Integrate both sides with respect to t

$$\int \frac{y'}{(-1+y)^2} dt = \int -k dt + c_1$$

- Evaluate integral

$$-\frac{1}{-1+y} = -tk + c_1$$

- Solve for y

$$y = \frac{-tk + c_1 - 1}{-tk + c_1}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

```
dsolve(diff(y(t),t) = -k*(-1+y(t))^2,y(t), singsol=all)
```

$$y(t) = \frac{1 + k(t + c_1)}{k(t + c_1)}$$

✓ Solution by Mathematica

Time used: 0.139 (sec). Leaf size: 30

```
DSolve[y'[t]== -k*(-1+y[t])^2,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{kt + 1 - c_1}{kt - c_1}$$
$$y(t) \rightarrow 1$$

4.7 problem 9

4.7.1 Solving as quadrature ode	1131
4.7.2 Maple step by step solution	1132

Internal problem ID [538]

Internal file name [OUTPUT/538_Sunday_June_05_2022_01_43_53_AM_86819030/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 9.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"quadrature"**

Maple gives the following as the ode type

`[_quadrature]`

$$y' - y^2(y^2 - 1) = 0$$

4.7.1 Solving as quadrature ode

Integrating both sides gives

$$\int \frac{1}{y^2(y^2 - 1)} dy = \int dt$$
$$\int \frac{1}{-a^2(-a^2 - 1)} d-a = t + c_1$$

Summary

The solution(s) found are the following

$$\int \frac{1}{-a^2(-a^2 - 1)} d-a = t + c_1 \tag{1}$$

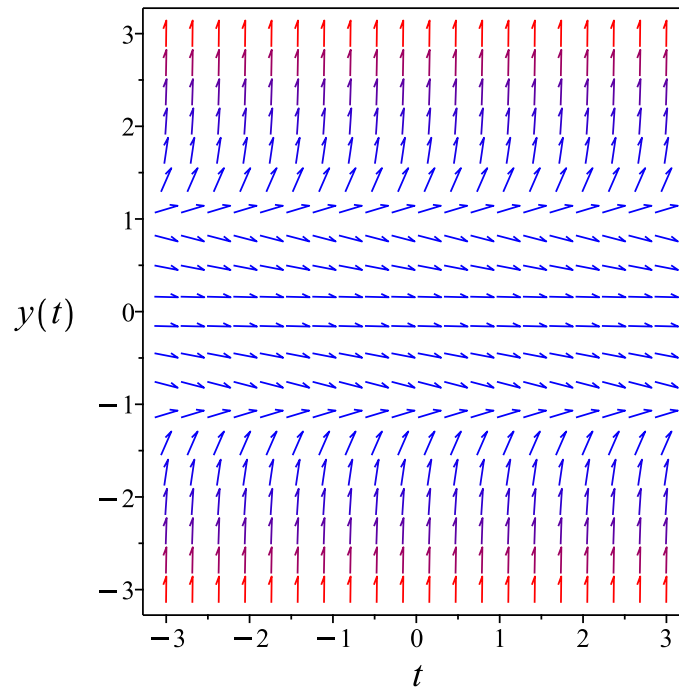


Figure 218: Slope field plot

Verification of solutions

$$\int^y \frac{1}{-a^2(-a^2-1)} da = t + c_1$$

Verified OK.

4.7.2 Maple step by step solution

Let's solve

$$y' - y^2(y^2 - 1) = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'}{y^2(y^2-1)} = 1$$

- Integrate both sides with respect to t

$$\int \frac{y'}{y^2(y^2-1)} dt = \int 1 dt + c_1$$

- Evaluate integral

$$-\frac{\ln(1+y)}{2} + \frac{1}{y} + \frac{\ln(-1+y)}{2} = t + c_1$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.282 (sec). Leaf size: 47

```
dsolve(diff(y(t),t) = y(t)^2*(y(t)^2-1),y(t), singsol=all)
```

$$y(t) = e^{\text{RootOf}(-\ln(e^{-Z}-2)e^{-Z}+2c_1e^{-Z}+Ze^{-Z}+2te^{-Z}+\ln(e^{-Z}-2)-2c_1-_-Z-2t-2)} - 1$$

✓ Solution by Mathematica

Time used: 0.246 (sec). Leaf size: 51

```
DSolve[y'[t] == y[t]^2*(y[t]^2-1),y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \text{InverseFunction} \left[\frac{1}{\#1} + \frac{1}{2} \log(1 - \#1) - \frac{1}{2} \log(\#1 + 1) \& \right] [t + c_1]$$

$$y(t) \rightarrow -1$$

$$y(t) \rightarrow 0$$

$$y(t) \rightarrow 1$$

4.8 problem 10

4.8.1 Solving as quadrature ode	1134
4.8.2 Maple step by step solution	1135

Internal problem ID [539]

Internal file name [OUTPUT/539_Sunday_June_05_2022_01_43_54_AM_75461840/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 10.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"quadrature"**

Maple gives the following as the ode type

`[_quadrature]`

$$y' - y(1 - y^2) = 0$$

4.8.1 Solving as quadrature ode

Integrating both sides gives

$$\int -\frac{1}{y(y^2 - 1)} dy = \int dt$$
$$\ln(y) - \frac{\ln(y+1)}{2} - \frac{\ln(y-1)}{2} = t + c_1$$

Raising both side to exponential gives

$$e^{\ln(y) - \frac{\ln(y+1)}{2} - \frac{\ln(y-1)}{2}} = e^{t+c_1}$$

Which simplifies to

$$\frac{y}{\sqrt{y+1}\sqrt{y-1}} = c_2 e^t$$

Summary

The solution(s) found are the following

$$y = \frac{-c_2^2 e^{2t} + \sqrt{e^{4t} c_2^4 - c_2^2 e^{2t}} + 1}{c_2^2 e^{2t} - 1} + 1 \quad (1)$$

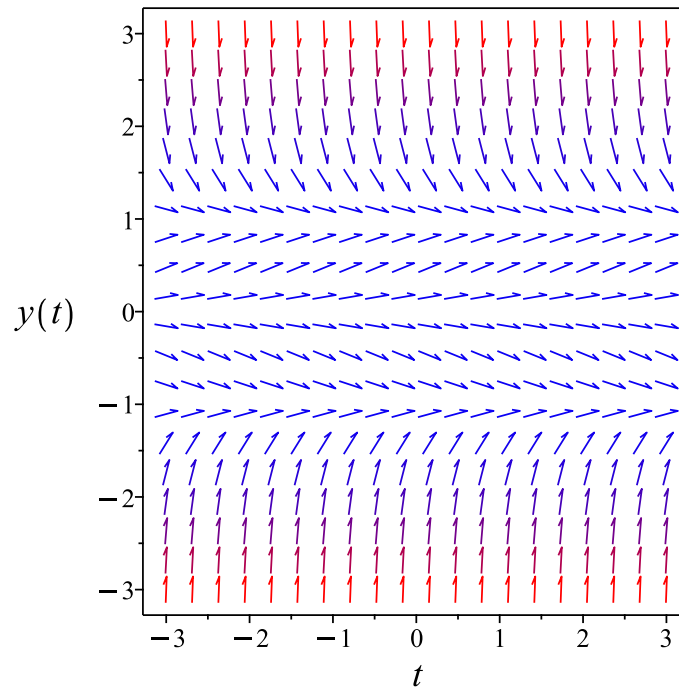


Figure 219: Slope field plot

Verification of solutions

$$y = \frac{-c_2^2 e^{2t} + \sqrt{e^{4t} c_2^4 - c_2^2 e^{2t} + 1}}{c_2^2 e^{2t} - 1} + 1$$

Verified OK.

4.8.2 Maple step by step solution

Let's solve

$$y' - y(1 - y^2) = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{y(1-y^2)} = 1$$

- Integrate both sides with respect to t

$$\int \frac{y'}{y(1-y^2)} dt = \int 1 dt + c_1$$

- Evaluate integral

$$\ln(y) - \frac{\ln(1+y)}{2} - \frac{\ln(-1+y)}{2} = t + c_1$$

- Solve for y

$$\left\{ y = \frac{\sqrt{(e^{2c_1+2t}-1)e^{2c_1+2t}}}{e^{2c_1+2t}-1}, y = -\frac{\sqrt{(e^{2c_1+2t}-1)e^{2c_1+2t}}}{e^{2c_1+2t}-1} \right\}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 29

```
dsolve(diff(y(t),t) = y(t)*(1-y(t)^2),y(t), singsol=all)
```

$$y(t) = \frac{1}{\sqrt{c_1 e^{-2t} + 1}}$$
$$y(t) = -\frac{1}{\sqrt{c_1 e^{-2t} + 1}}$$

✓ Solution by Mathematica

Time used: 0.676 (sec). Leaf size: 100

```
DSolve[y'[t]== y[t]*(1-y[t]^2),y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -\frac{e^t}{\sqrt{e^{2t} + e^{2c_1}}}$$
$$y(t) \rightarrow \frac{e^t}{\sqrt{e^{2t} + e^{2c_1}}}$$
$$y(t) \rightarrow -1$$
$$y(t) \rightarrow 0$$
$$y(t) \rightarrow 1$$
$$y(t) \rightarrow -\frac{e^t}{\sqrt{e^{2t}}}$$
$$y(t) \rightarrow \frac{e^t}{\sqrt{e^{2t}}}$$

4.9 problem 11

4.9.1 Solving as quadrature ode	1138
4.9.2 Maple step by step solution	1139

Internal problem ID [540]

Internal file name [OUTPUT/540_Sunday_June_05_2022_01_43_57_AM_7870853/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 11.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"quadrature"**

Maple gives the following as the ode type

`[_quadrature]`

$$y' + b\sqrt{y} - ay = 0$$

4.9.1 Solving as quadrature ode

Integrating both sides gives

$$\int \frac{1}{-b\sqrt{y} + ay} dy = \int dt$$
$$\frac{\ln(a^2y - b^2)}{a} - \frac{\ln(a\sqrt{y} + b)}{a} + \frac{\ln(a\sqrt{y} - b)}{a} = t + c_1$$

The above can be written as

$$\left(\frac{1}{a}\right) (\ln(a^2y - b^2) - \ln(a\sqrt{y} + b) + \ln(a\sqrt{y} - b)) = t + c_1$$
$$\ln(a^2y - b^2) - \ln(a\sqrt{y} + b) + \ln(a\sqrt{y} - b) = (a)(t + c_1)$$
$$= a(t + c_1)$$

Raising both side to exponential gives

$$e^{\ln(a^2y - b^2) - \ln(a\sqrt{y} + b) + \ln(a\sqrt{y} - b)} = ac_1 e^{at}$$

Which simplifies to

$$\frac{(a^2y - b^2)(a\sqrt{y} - b)}{a\sqrt{y} + b} = c_2e^{at}$$

Summary

The solution(s) found are the following

$$y = \frac{2b(b - \sqrt{c_2e^{at}}) + c_2e^{at} - b^2}{a^2} \quad (1)$$

Verification of solutions

$$y = \frac{2b(b - \sqrt{c_2e^{at}}) + c_2e^{at} - b^2}{a^2}$$

Verified OK.

4.9.2 Maple step by step solution

Let's solve

$$y' + b\sqrt{y} - ay = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{-b\sqrt{y}+ay} = 1$$

- Integrate both sides with respect to t

$$\int \frac{y'}{-b\sqrt{y}+ay} dt = \int 1 dt + c_1$$

- Evaluate integral

$$\frac{\ln(ya^2-b^2)}{a} - \frac{\ln(b+a\sqrt{y})}{a} + \frac{\ln(a\sqrt{y}-b)}{a} = t + c_1$$

- Solve for y

$$\left\{ y = \frac{2b(b - e^{\frac{1}{2}c_1 a + \frac{1}{2}at}) - b^2 + e^{c_1 a + at}}{a^2}, y = \frac{2b(b + e^{\frac{1}{2}c_1 a + \frac{1}{2}at}) - b^2 + e^{c_1 a + at}}{a^2} \right\}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 26

```
dsolve(diff(y(t),t) = -b*y(t)^(1/2)+a*y(t),y(t), singsol=all)
```

$$\frac{-e^{\frac{at}{2}}c_1a + \sqrt{y(t)}a - b}{a} = 0$$

✓ Solution by Mathematica

Time used: 0.844 (sec). Leaf size: 55

```
DSolve[y'[t] == -b*y[t]^(1/2)+a*y[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{e^{-ac_1} \left(e^{\frac{at}{2}} - be^{\frac{ac_1}{2}} \right)^2}{a^2}$$
$$y(t) \rightarrow 0$$
$$y(t) \rightarrow \frac{b^2}{a^2}$$

4.10 problem 12

4.10.1 Solving as quadrature ode	1141
4.10.2 Maple step by step solution	1142

Internal problem ID [541]

Internal file name [OUTPUT/541_Sunday_June_05_2022_01_43_58_AM_15077121/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 12.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"quadrature"**

Maple gives the following as the ode type

`[_quadrature]`

$$y' - y^2(4 - y^2) = 0$$

4.10.1 Solving as quadrature ode

Integrating both sides gives

$$\int -\frac{1}{y^2(y^2 - 4)} dy = \int dt$$
$$\int -\frac{1}{a^2(a^2 - 4)} da = t + c_1$$

Summary

The solution(s) found are the following

$$\int -\frac{1}{a^2(a^2 - 4)} da = t + c_1 \tag{1}$$

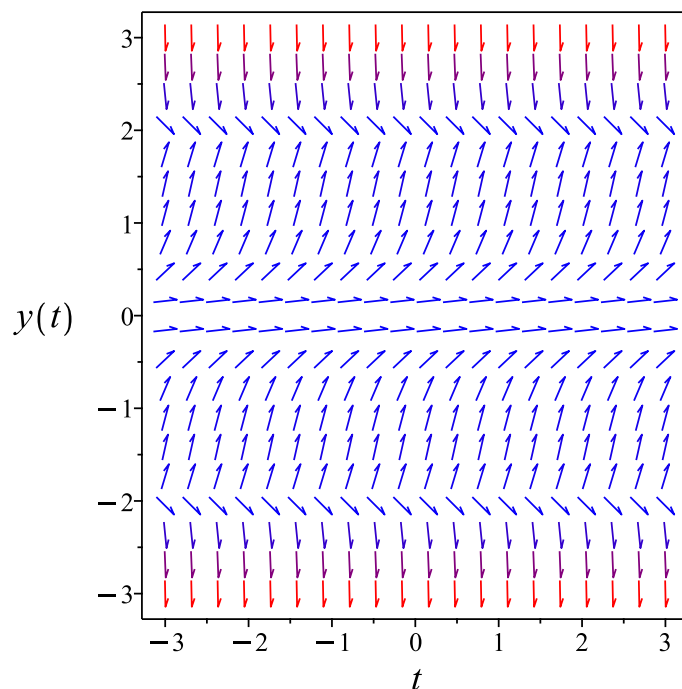


Figure 220: Slope field plot

Verification of solutions

$$\int^y -\frac{1}{a^2(a^2 - 4)} da = t + c_1$$

Verified OK.

4.10.2 Maple step by step solution

Let's solve

$$y' - y^2(4 - y^2) = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'}{y^2(4-y^2)} = 1$$

- Integrate both sides with respect to t

$$\int \frac{y'}{y^2(4-y^2)} dt = \int 1 dt + c_1$$

- Evaluate integral

$$-\frac{1}{4y} - \frac{\ln(-2+y)}{16} + \frac{\ln(2+y)}{16} = t + c_1$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.032 (sec). Leaf size: 49

```
dsolve(diff(y(t),t) = y(t)^2*(4-y(t)^2),y(t), singsol=all)
```

$$y(t) = e^{\text{RootOf}(\ln(e^{-Z}-4)e^{-Z}+16c_1e^{-Z}-_Ze^{-Z}+16t e^{-Z}-2\ln(e^{-Z}-4)-32c_1+2_Z-32t+4) - 2}$$

✓ Solution by Mathematica

Time used: 0.247 (sec). Leaf size: 57

```
DSolve[y'[t] == y[t]^2*(4-y[t]^2),y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \text{InverseFunction} \left[\frac{1}{4\#1} + \frac{1}{16} \log(2 - \#1) - \frac{1}{16} \log(\#1 + 2) \& \right] [-t + c_1]$$

$$y(t) \rightarrow -2$$

$$y(t) \rightarrow 0$$

$$y(t) \rightarrow 2$$

4.11 problem 13

- 4.11.1 Solving as quadrature ode 1144
- 4.11.2 Maple step by step solution 1145

Internal problem ID [542]

Internal file name [OUTPUT/542_Sunday_June_05_2022_01_43_59_AM_906908/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.5. Page 88

Problem number: 13.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**quadrature**"

Maple gives the following as the ode type

[_quadrature]

$$y' - (1 - y)^2 y^2 = 0$$

4.11.1 Solving as quadrature ode

Integrating both sides gives

$$\int \frac{1}{y^2 (y - 1)^2} dy = \int dt$$
$$\int^y \frac{1}{_a^2 (_a - 1)^2} d_a = t + c_1$$

Summary

The solution(s) found are the following

$$\int^y \frac{1}{_a^2 (_a - 1)^2} d_a = t + c_1 \tag{1}$$

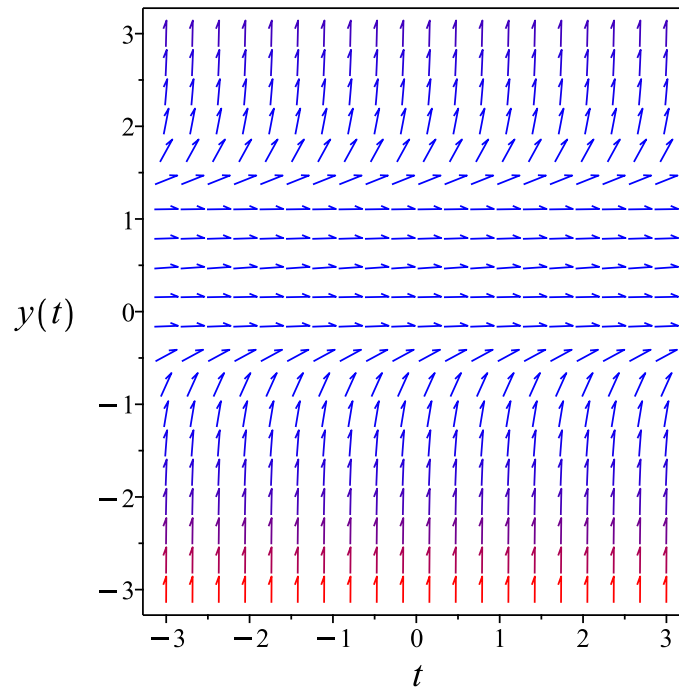


Figure 221: Slope field plot

Verification of solutions

$$\int \frac{1}{-a^2 (-a - 1)^2} da = t + c_1$$

Verified OK.

4.11.2 Maple step by step solution

Let's solve

$$y' - (1 - y)^2 y^2 = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'}{(1-y)^2 y^2} = 1$$

- Integrate both sides with respect to t

$$\int \frac{y'}{(1-y)^2 y^2} dt = \int 1 dt + c_1$$

- Evaluate integral

$$-\frac{1}{y} + 2 \ln(y) - \frac{1}{-1+y} - 2 \ln(-1+y) = t + c_1$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 66

```
dsolve(diff(y(t),t) = (1-y(t))^2*y(t)^2,y(t), singsol=all)
```

$$y(t) = e^{\text{RootOf}(-2 \ln(e^{-Z}+1)e^{2-Z}+c_1e^{2-Z}+2_Ze^{2-Z}+te^{2-Z}-2 \ln(e^{-Z}+1)e^{-Z}+c_1e^{-Z}+2_Ze^{-Z}+te^{-Z}+2e^{-Z}+1)} + 1$$

✓ Solution by Mathematica

Time used: 0.365 (sec). Leaf size: 50

```
DSolve[y'[t] == (1-y[t])^2*y[t]^2,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \text{InverseFunction} \left[-\frac{1}{\#1-1} - \frac{1}{\#1} - 2 \log(1-\#1) + 2 \log(\#1) \& \right] [t + c_1]$$

$$y(t) \rightarrow 0$$

$$y(t) \rightarrow 1$$

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5.1 problem 1

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Internal problem ID [543]

Internal file name [OUTPUT/543_Sunday_June_05_2022_01_44_01_AM_47710901/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "homogeneousTypeMapleC", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$(-2 + 2y)y' = -3 - 2x$$

5.1.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{-\frac{3}{2} - x}{y - 1}\end{aligned}$$

Where $f(x) = -\frac{3}{2} - x$ and $g(y) = \frac{1}{y-1}$. Integrating both sides gives

$$\frac{1}{y-1} dy = -\frac{3}{2} - x dx$$

$$\int \frac{1}{\frac{1}{y-1}} dy = \int -\frac{3}{2} - x dx$$

$$\frac{1}{2}y^2 - y = -\frac{3}{2}x - \frac{1}{2}x^2 + c_1$$

Which results in

$$y = 1 + \sqrt{-x^2 + 2c_1 - 3x + 1}$$

$$y = 1 - \sqrt{-x^2 + 2c_1 - 3x + 1}$$

Summary

The solution(s) found are the following

$$y = 1 + \sqrt{-x^2 + 2c_1 - 3x + 1} \tag{1}$$

$$y = 1 - \sqrt{-x^2 + 2c_1 - 3x + 1} \tag{2}$$

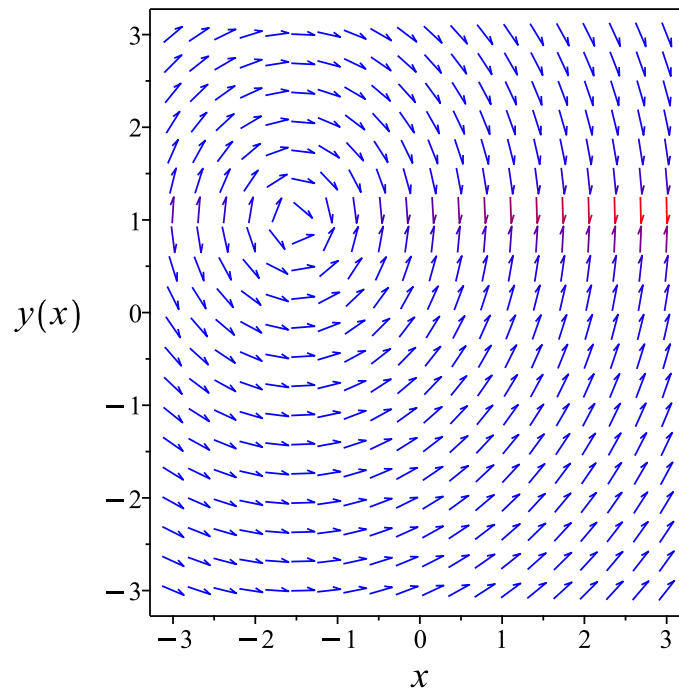


Figure 222: Slope field plot

Verification of solutions

$$y = 1 + \sqrt{-x^2 + 2c_1 - 3x + 1}$$

Verified OK.

$$y = 1 - \sqrt{-x^2 + 2c_1 - 3x + 1}$$

Verified OK.

5.1.2 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-3 - 2x}{-2 + 2y} \quad (1)$$

Which becomes

$$(2y - 2) dy = (-3 - 2x) dx \quad (2)$$

But the RHS is complete differential because

$$(-3 - 2x) dx = d(-x^2 - 3x)$$

Hence (2) becomes

$$(2y - 2) dy = d(-x^2 - 3x)$$

Integrating both sides gives gives these solutions

$$y = 1 + \sqrt{-x^2 + c_1 - 3x + 1} + c_1$$

$$y = 1 - \sqrt{-x^2 + c_1 - 3x + 1} + c_1$$

Summary

The solution(s) found are the following

$$y = 1 + \sqrt{-x^2 + c_1 - 3x + 1} + c_1 \quad (1)$$

$$y = 1 - \sqrt{-x^2 + c_1 - 3x + 1} + c_1 \quad (2)$$

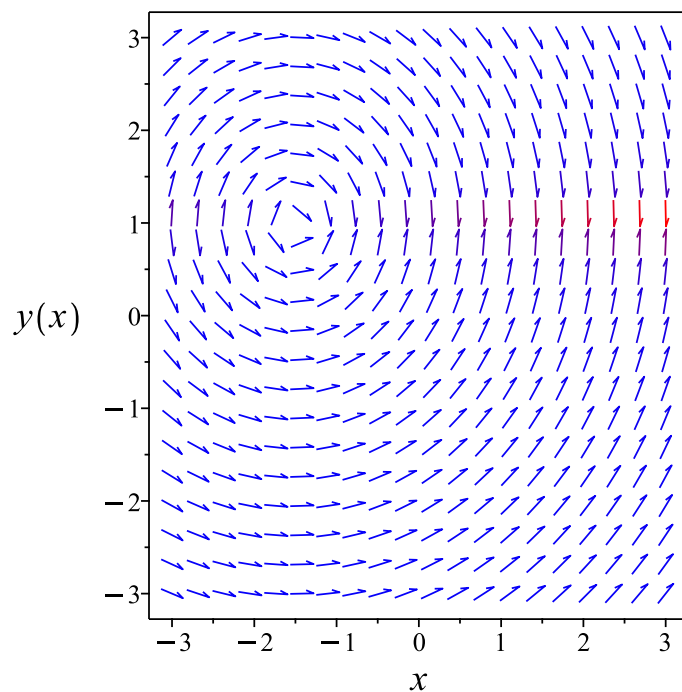


Figure 223: Slope field plot

Verification of solutions

$$y = 1 + \sqrt{-x^2 + c_1 - 3x + 1} + c_1$$

Verified OK.

$$y = 1 - \sqrt{-x^2 + c_1 - 3x + 1} + c_1$$

Verified OK.

5.1.3 Solving as homogeneousTypeMapleC ode

Let $Y = y + y_0$ and $X = x + x_0$ then the above is transformed to new ode in $Y(X)$

$$\frac{d}{dX}Y(X) = -\frac{2X + 2x_0 + 3}{2(Y(X) + y_0 - 1)}$$

Solving for possible values of x_0 and y_0 which makes the above ode a homogeneous ode results in

$$\begin{aligned} x_0 &= -\frac{3}{2} \\ y_0 &= 1 \end{aligned}$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$\frac{d}{dX}Y(X) = -\frac{X}{Y(X)}$$

In canonical form, the ODE is

$$\begin{aligned} Y' &= F(X, Y) \\ &= -\frac{X}{Y} \end{aligned} \tag{1}$$

An ode of the form $Y' = \frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order n if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both $M = -X$ and $N = Y$ are both homogeneous and of the same order $n = 1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{Y}{X}$, or $Y = uX$. Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation $Y = uX$ to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dX}X + u &= -\frac{1}{u} \\ \frac{du}{dX} &= \frac{-\frac{1}{u(X)} - u(X)}{X} \end{aligned}$$

Or

$$\frac{d}{dX}u(X) - \frac{-\frac{1}{u(X)} - u(X)}{X} = 0$$

Or

$$\left(\frac{d}{dX}u(X)\right)u(X)X + u(X)^2 + 1 = 0$$

Which is now solved as separable in $u(X)$. Which is now solved in $u(X)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(X, u) \\ &= f(X)g(u) \\ &= -\frac{u^2 + 1}{uX} \end{aligned}$$

Where $f(X) = -\frac{1}{X}$ and $g(u) = \frac{u^2+1}{u}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{u^2+1}{u}} du &= -\frac{1}{X} dX \\ \int \frac{1}{\frac{u^2+1}{u}} du &= \int -\frac{1}{X} dX \\ \frac{\ln(u^2 + 1)}{2} &= -\ln(X) + c_2\end{aligned}$$

Raising both side to exponential gives

$$\sqrt{u^2 + 1} = e^{-\ln(X)+c_2}$$

Which simplifies to

$$\sqrt{u^2 + 1} = \frac{c_3}{X}$$

Which simplifies to

$$\sqrt{u(X)^2 + 1} = \frac{c_3 e^{c_2}}{X}$$

The solution is

$$\sqrt{u(X)^2 + 1} = \frac{c_3 e^{c_2}}{X}$$

Now u in the above solution is replaced back by Y using $u = \frac{Y}{X}$ which results in the solution

$$\sqrt{\frac{Y(X)^2}{X^2} + 1} = \frac{c_3 e^{c_2}}{X}$$

Using the solution for $Y(X)$

$$\sqrt{\frac{Y(X)^2 + X^2}{X^2}} = \frac{c_3 e^{c_2}}{X}$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y + 1$$
$$X = x - \frac{3}{2}$$

Then the solution in y becomes

$$\sqrt{\frac{(y-1)^2 + (x + \frac{3}{2})^2}{(x + \frac{3}{2})^2}} = \frac{c_3 e^{c_2}}{x + \frac{3}{2}}$$

Summary

The solution(s) found are the following

$$\sqrt{\frac{(y-1)^2 + (x + \frac{3}{2})^2}{(x + \frac{3}{2})^2}} = \frac{c_3 e^{c_2}}{x + \frac{3}{2}} \quad (1)$$

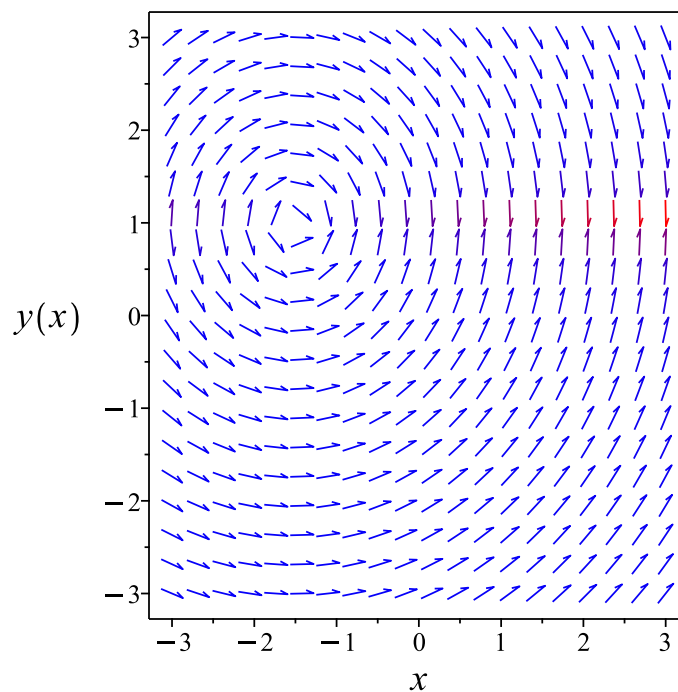


Figure 224: Slope field plot

Verification of solutions

$$\sqrt{\frac{(y-1)^2 + (x + \frac{3}{2})^2}{(x + \frac{3}{2})^2}} = \frac{c_3 e^{c_2}}{x + \frac{3}{2}}$$

Verified OK.

5.1.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{2x+3}{2(y-1)}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 234: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{-\frac{3}{2} - x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{-\frac{3}{2}-x} dx \end{aligned}$$

Which results in

$$S = -\frac{3}{2}x - \frac{1}{2}x^2$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{2x + 3}{2(y - 1)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= -\frac{3}{2} - x \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y - 1 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R - 1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{1}{2}R^2 - R + c_1 \quad (4)$$

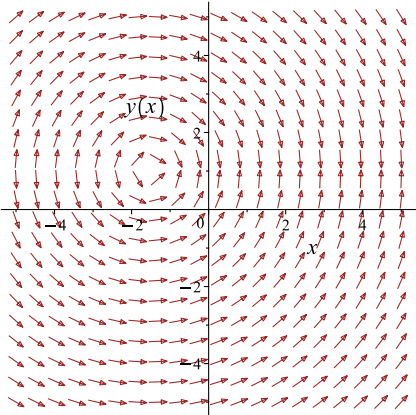
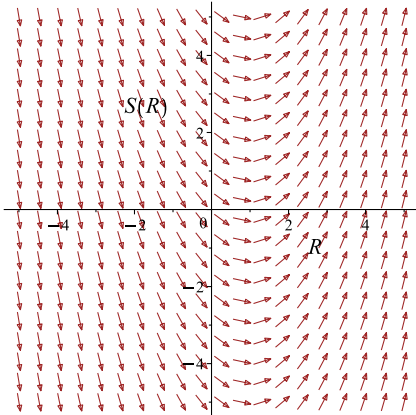
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-\frac{3}{2}x - \frac{1}{2}x^2 = \frac{y^2}{2} - y + c_1$$

Which simplifies to

$$-\frac{3}{2}x - \frac{1}{2}x^2 = \frac{y^2}{2} - y + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{2x+3}{2(y-1)}$ 	$R = y$ $S = -\frac{3}{2}x - \frac{1}{2}x^2$	$\frac{dS}{dR} = R - 1$ 

Summary

The solution(s) found are the following

$$-\frac{3}{2}x - \frac{1}{2}x^2 = \frac{y^2}{2} - y + c_1 \quad (1)$$

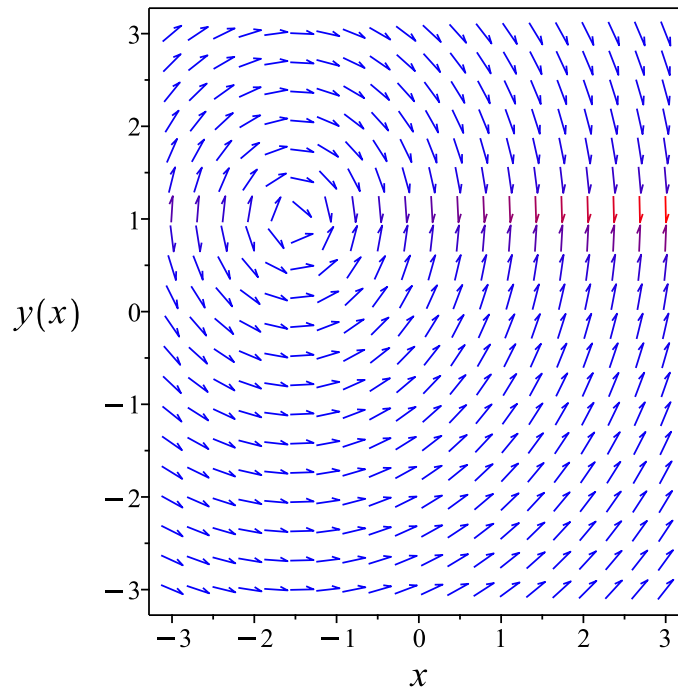


Figure 225: Slope field plot

Verification of solutions

$$-\frac{3}{2}x - \frac{1}{2}x^2 = \frac{y^2}{2} - y + c_1$$

Verified OK.

5.1.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(-2y + 2) dy &= (2x + 3) dx \\ (-3 - 2x) dx + (-2y + 2) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -3 - 2x \\ N(x, y) &= -2y + 2\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-3 - 2x) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-2y + 2) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -3 - 2x dx \\ \phi &= -x^2 - 3x + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -2y + 2$. Therefore equation (4) becomes

$$-2y + 2 = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -2y + 2$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int (-2y + 2) dy \\ f(y) &= -y^2 + 2y + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x^2 - y^2 - 3x + 2y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x^2 - y^2 - 3x + 2y$$

Summary

The solution(s) found are the following

$$-y^2 - x^2 + 2y - 3x = c_1 \tag{1}$$

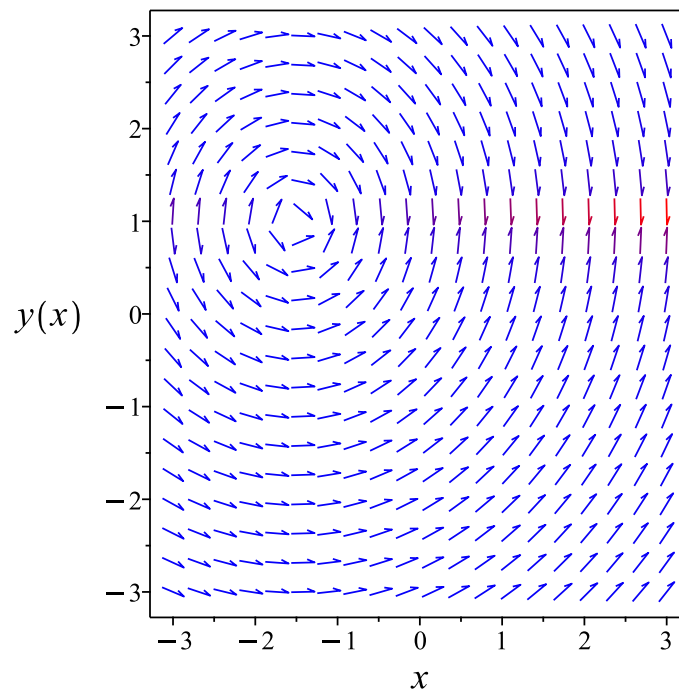


Figure 226: Slope field plot

Verification of solutions

$$-y^2 - x^2 + 2y - 3x = c_1$$

Verified OK.

5.1.6 Maple step by step solution

Let's solve

$$(-2 + 2y)y' = -3 - 2x$$

- Highest derivative means the order of the ODE is 1

y'

- Integrate both sides with respect to x

$$\int (-2 + 2y)y'dx = \int (-3 - 2x)dx + c_1$$

- Evaluate integral

$$y^2 - 2y = -x^2 + c_1 - 3x$$

- Solve for y

$$\{y = 1 - \sqrt{-x^2 + c_1 - 3x + 1}, y = 1 + \sqrt{-x^2 + c_1 - 3x + 1}\}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 43

```
dsolve(3+2*x+(-2+2*y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = 1 - \sqrt{-x^2 - c_1 - 3x + 1}$$

$$y(x) = 1 + \sqrt{-x^2 - c_1 - 3x + 1}$$

✓ Solution by Mathematica

Time used: 0.115 (sec). Leaf size: 51

```
DSolve[3+2*x+(-2+2*y[x])*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow 1 - \sqrt{-x^2 - 3x + 1 + 2c_1}$$

$$y(x) \rightarrow 1 + \sqrt{-x^2 - 3x + 1 + 2c_1}$$

5.2 problem 2

- 5.2.1 Solving as homogeneousTypeD2 ode 1165
- 5.2.2 Solving as first order ode lie symmetry calculated ode 1167

Internal problem ID [544]

Internal file name [OUTPUT/544_Sunday_June_05_2022_01_44_02_AM_13505422/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD2**", "**first_order_ode_lie_symmetry_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$4y + (2x - 2y)y' = -2x$$

5.2.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$4u(x)x + (2x - 2u(x)x)(u'(x)x + u(x)) = -2x$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{u^2 - 3u - 1}{(u - 1)x} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(u) = \frac{u^2-3u-1}{u-1}$. Integrating both sides gives

$$\frac{1}{\frac{u^2-3u-1}{u-1}} du = -\frac{1}{x} dx$$

$$\int \frac{1}{\frac{u^2-3u-1}{u-1}} du = \int -\frac{1}{x} dx$$

$$\frac{\ln(u^2 - 3u - 1)}{2} - \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(2u-3)\sqrt{13}}{13}\right)}{13} = -\ln(x) + c_2$$

The solution is

$$\frac{\ln(u(x)^2 - 3u(x) - 1)}{2} - \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(2u(x)-3)\sqrt{13}}{13}\right)}{13} + \ln(x) - c_2 = 0$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\frac{\ln\left(\frac{y^2}{x^2} - \frac{3y}{x} - 1\right)}{2} - \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{\left(\frac{2y}{x}-3\right)\sqrt{13}}{13}\right)}{13} + \ln(x) - c_2 = 0$$

$$\frac{\ln\left(\frac{y^2}{x^2} - \frac{3y}{x} - 1\right)}{2} - \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(2y-3x)\sqrt{13}}{13x}\right)}{13} + \ln(x) - c_2 = 0$$

Summary

The solution(s) found are the following

$$\frac{\ln\left(\frac{y^2}{x^2} - \frac{3y}{x} - 1\right)}{2} - \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(2y-3x)\sqrt{13}}{13x}\right)}{13} + \ln(x) - c_2 = 0 \quad (1)$$

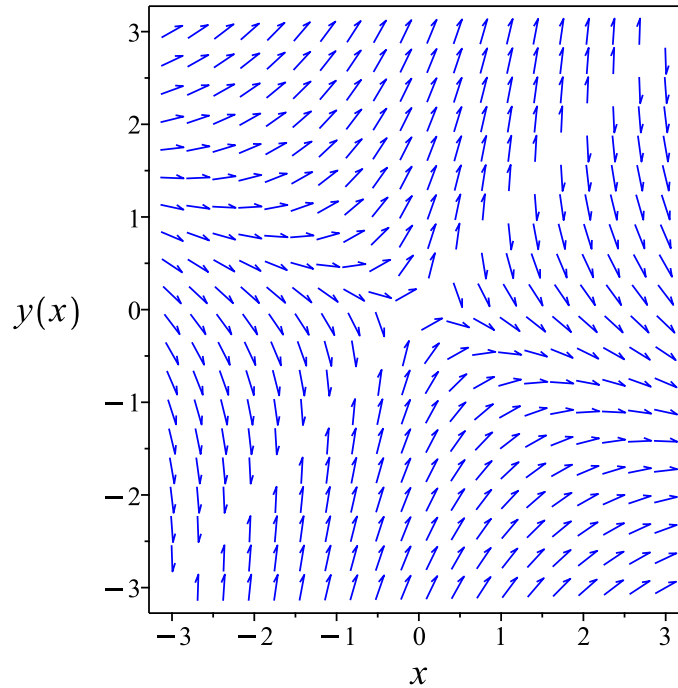


Figure 227: Slope field plot

Verification of solutions

$$\frac{\ln\left(\frac{y^2}{x^2} - \frac{3y}{x} - 1\right)}{2} - \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(2y-3x)\sqrt{13}}{13x}\right)}{13} + \ln(x) - c_2 = 0$$

Verified OK.

5.2.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{x + 2y}{-x + y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 + \frac{(x+2y)(b_3 - a_2)}{-x+y} - \frac{(x+2y)^2 a_3}{(-x+y)^2} \\ - \left(\frac{1}{-x+y} + \frac{x+2y}{(-x+y)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left(\frac{2}{-x+y} - \frac{x+2y}{(-x+y)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{x^2 a_2 - x^2 a_3 + 4x^2 b_2 - x^2 b_3 - 2xy a_2 - 4xy a_3 - 2xy b_2 + 2xy b_3 - 2y^2 a_2 - 7y^2 a_3 + y^2 b_2 + 2y^2 b_3 + 3xb_1 - 3ya_1}{(x-y)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} x^2 a_2 - x^2 a_3 + 4x^2 b_2 - x^2 b_3 - 2xy a_2 - 4xy a_3 - 2xy b_2 \\ + 2xy b_3 - 2y^2 a_2 - 7y^2 a_3 + y^2 b_2 + 2y^2 b_3 + 3xb_1 - 3ya_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} a_2 v_1^2 - 2a_2 v_1 v_2 - 2a_2 v_2^2 - a_3 v_1^2 - 4a_3 v_1 v_2 - 7a_3 v_2^2 + 4b_2 v_1^2 \\ - 2b_2 v_1 v_2 + b_2 v_2^2 - b_3 v_1^2 + 2b_3 v_1 v_2 + 2b_3 v_2^2 - 3a_1 v_2 + 3b_1 v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$(a_2 - a_3 + 4b_2 - b_3) v_1^2 + (-2a_2 - 4a_3 - 2b_2 + 2b_3) v_1 v_2 + 3b_1 v_1 + (-2a_2 - 7a_3 + b_2 + 2b_3) v_2^2 - 3a_1 v_2 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -3a_1 &= 0 \\ 3b_1 &= 0 \\ -2a_2 - 7a_3 + b_2 + 2b_3 &= 0 \\ -2a_2 - 4a_3 - 2b_2 + 2b_3 &= 0 \\ a_2 - a_3 + 4b_2 - b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= -3a_3 + b_3 \\ a_3 &= a_3 \\ b_1 &= 0 \\ b_2 &= a_3 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= y - \left(\frac{x + 2y}{-x + y} \right) (x) \\ &= \frac{x^2 + 3yx - y^2}{x - y} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{x^2 + 3yx - y^2}{x - y}} dy \end{aligned}$$

Which results in

$$S = \frac{\ln(-x^2 - 3yx + y^2)}{2} - \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(-3x+2y)\sqrt{13}}{13x}\right)}{13}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x + 2y}{-x + y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{x + 2y}{x^2 + 3yx - y^2} \\ S_y &= \frac{x - y}{x^2 + 3yx - y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \quad (4)$$

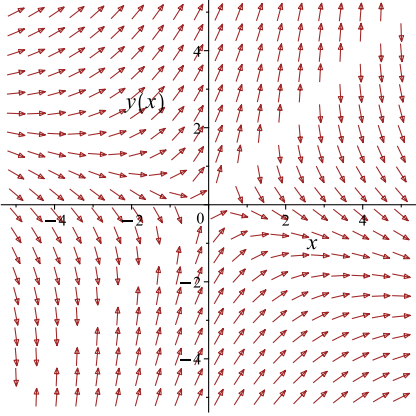
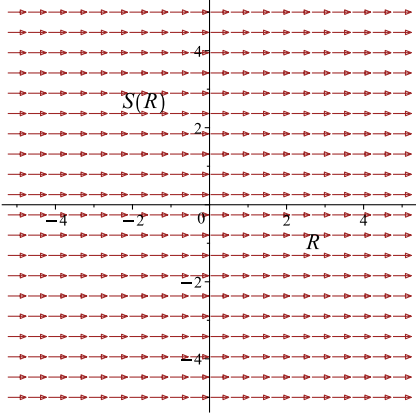
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\ln(y^2 - 3yx - x^2)}{2} + \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(-2y+3x)\sqrt{13}}{13x}\right)}{13} = c_1$$

Which simplifies to

$$\frac{\ln(y^2 - 3yx - x^2)}{2} + \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(-2y+3x)\sqrt{13}}{13x}\right)}{13} = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{x+2y}{-x+y}$ 	$R = x$ $S = \frac{\ln(-x^2 - 3yx + y^2)}{2}$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$\frac{\ln(y^2 - 3yx - x^2)}{2} + \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(-2y+3x)\sqrt{13}}{13x}\right)}{13} = c_1 \quad (1)$$

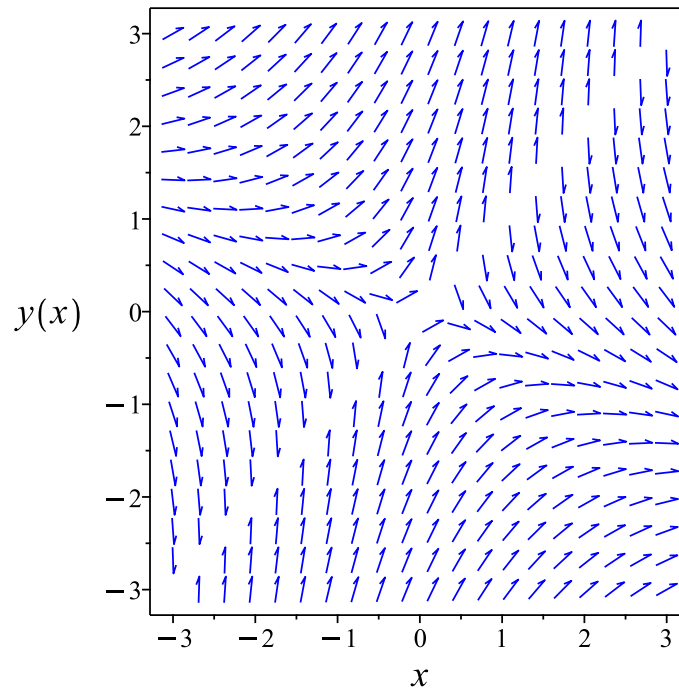


Figure 228: Slope field plot

Verification of solutions

$$\frac{\ln(y^2 - 3yx - x^2)}{2} + \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(-2y+3x)\sqrt{13}}{13x}\right)}{13} = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.14 (sec). Leaf size: 55

```
dsolve(2*x+4*y(x)+(2*x-2*y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$-\frac{\ln\left(\frac{-x^2-3xy(x)+y(x)^2}{x^2}\right)}{2} + \frac{\sqrt{13} \operatorname{arctanh}\left(\frac{(2y(x)-3x)\sqrt{13}}{13x}\right)}{13} - \ln(x) - c_1 = 0$$

✓ Solution by Mathematica

Time used: 0.065 (sec). Leaf size: 63

```
DSolve[2*x+4*y[x]+(2*x-2*y[x])*y'[x]== 0,y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve}\left[\frac{1}{26}\left(\left(13 + \sqrt{13}\right)\log\left(-\frac{2y(x)}{x} + \sqrt{13} + 3\right) - \left(\sqrt{13} - 13\right)\log\left(\frac{2y(x)}{x} + \sqrt{13} - 3\right)\right) = -\log(x) + c_1, y(x)\right]$$

5.3 problem 3

5.3.1 Solving as differentialType ode	1175
5.3.2 Solving as exact ode	1180
5.3.3 Maple step by step solution	1183

Internal problem ID [545]

Internal file name [OUTPUT/545_Sunday_June_05_2022_01_44_04_AM_28395296/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 3.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "differentialType"

Maple gives the following as the ode type

`[_exact, _rational]`

$$-2yx + (3 - x^2 + 6y^2) y' = -3x^2 - 2$$

5.3.1 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-2 - 3x^2 + 2yx}{3 - x^2 + 6y^2} \tag{1}$$

Which becomes

$$(-6y^2 - 3) dy = (-x^2) dy + (3x^2 - 2yx + 2) dx \tag{2}$$

But the RHS is complete differential because

$$(-x^2) dy + (3x^2 - 2yx + 2) dx = d(x^3 - yx^2 + 2x)$$

Hence (2) becomes

$$(-6y^2 - 3) dy = d(x^3 - yx^2 + 2x)$$

Integrating both sides gives gives these solutions

$$y = \frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{6} - \frac{1}{(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162})^{\frac{1}{3}}}$$

$$y = -\frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{12} + \frac{1}{(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162})^{\frac{1}{3}}}$$

$$y = -\frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{12} + \frac{1}{(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162})^{\frac{1}{3}}}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{6\left(-\frac{x^2}{6} + \frac{1}{2}\right)} \quad (1)$$

$$y = \frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{6\left(-\frac{x^2}{6} + \frac{1}{2}\right)} + c_1 \quad (2)$$

$$y = \frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{6\left(-\frac{x^2}{6} + \frac{1}{2}\right)} + \frac{i\sqrt{3}\left(\frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{6} + \frac{-x^2 + 3}{2}\right)}{2} + c_1$$

$$y = \frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{6\left(-\frac{x^2}{6} + \frac{1}{2}\right)} + \frac{i\sqrt{3}\left(\frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{6} + \frac{-x^2 + 3}{2}\right)}{2} + c_1 \quad (3)$$

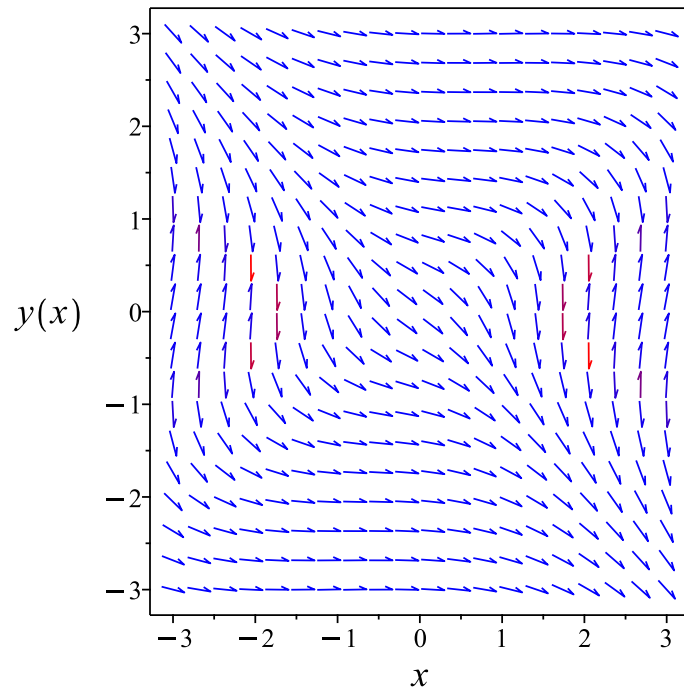


Figure 229: Slope field plot

Verification of solutions

y

$$= \frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{6\left(-\frac{x^2}{6} + \frac{1}{2}\right)} + c_1$$

Verified OK.

$y =$

$$\frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{12} + \frac{-\frac{x^2}{2} + \frac{3}{2}}{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}} + \frac{i\sqrt{3}\left(\frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{6}\right)}{2} + c_1$$

Verified OK.

$y =$

$$\frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{12} + \frac{-\frac{x^2}{2} + \frac{3}{2}}{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}} + \frac{i\sqrt{3}\left(\frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{6}\right)}{2} + c_1$$

Verified OK.

5.3.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (-x^2 + 6y^2 + 3) dy &= (-3x^2 + 2yx - 2) dx \\ (3x^2 - 2yx + 2) dx + (-x^2 + 6y^2 + 3) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 3x^2 - 2yx + 2 \\ N(x, y) &= -x^2 + 6y^2 + 3 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(3x^2 - 2yx + 2) \\ &= -2x\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x^2 + 6y^2 + 3) \\ &= -2x\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 3x^2 - 2yx + 2 dx \\ \phi &= x^3 - yx^2 + 2x + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = -x^2 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -x^2 + 6y^2 + 3$. Therefore equation (4) becomes

$$-x^2 + 6y^2 + 3 = -x^2 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 6y^2 + 3$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (6y^2 + 3) dy$$

$$f(y) = 2y^3 + 3y + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = x^3 - yx^2 + 2y^3 + 2x + 3y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = x^3 - yx^2 + 2y^3 + 2x + 3y$$

Summary

The solution(s) found are the following

$$x^3 - x^2y + 2y^3 + 2x + 3y = c_1 \tag{1}$$

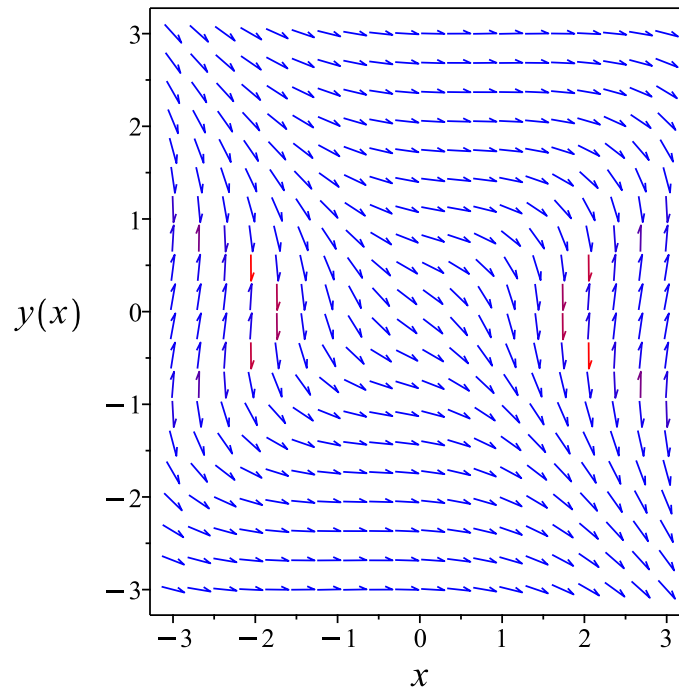


Figure 230: Slope field plot

Verification of solutions

$$x^3 - x^2y + 2y^3 + 2x + 3y = c_1$$

Verified OK.

5.3.3 Maple step by step solution

Let's solve

$$-2yx + (3 - x^2 + 6y^2) y' = -3x^2 - 2$$

- Highest derivative means the order of the ODE is 1
 y'
- Check if ODE is exact
 - ODE is exact if the lhs is the total derivative of a C^2 function
 $F'(x, y) = 0$
 - Compute derivative of lhs
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$

- Evaluate derivatives

$$-2x = -2x$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$\left[F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$

- Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to x

$$F(x, y) = \int (3x^2 - 2yx + 2) dx + f_1(y)$$

- Evaluate integral

$$F(x, y) = x^3 - yx^2 + 2x + f_1(y)$$

- Take derivative of $F(x, y)$ with respect to y

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$

- Compute derivative

$$-x^2 + 6y^2 + 3 = -x^2 + \frac{d}{dy} f_1(y)$$

- Isolate for $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = 6y^2 + 3$$

- Solve for $f_1(y)$

$$f_1(y) = 2y^3 + 3y$$

- Substitute $f_1(y)$ into equation for $F(x, y)$

$$F(x, y) = x^3 - yx^2 + 2y^3 + 2x + 3y$$

- Substitute $F(x, y)$ into the solution of the ODE

$$x^3 - yx^2 + 2y^3 + 2x + 3y = c_1$$

- Solve for y

$$\left\{ y = \frac{\left(-54x^3 + 54c_1 - 108x + 6\sqrt{75x^6 - 162c_1x^3 + 378x^4 + 81c_1^2 - 324c_1x + 162x^2 + 162} \right)^{\frac{1}{3}}}{6} - \frac{6\left(-54x^3 + 54c_1 - 108x + 6\sqrt{75x^6 - 162c_1x^3 + 378x^4 + 81c_1^2 - 324c_1x + 162x^2 + 162} \right)^{\frac{1}{3}}}{\left(-54x^3 + 54c_1 - 108x + 6\sqrt{75x^6 - 162c_1x^3 + 378x^4 + 81c_1^2 - 324c_1x + 162x^2 + 162} \right)^{\frac{1}{3}}} \right.$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 372

```
dsolve(2+3*x^2-2*x*y(x)+(3-x^2+6*y(x)^2)*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = \frac{\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{2}{3}} + 6x^2 - 18}{6\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}$$
$$y(x) = \frac{(-1 - i\sqrt{3})\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{\frac{12}{(x^2 - 3)(i\sqrt{3} - 1)}} + \frac{2\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}{2\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}$$
$$y(x) = \frac{(i\sqrt{3} - 1)\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{2}{3}}}{6} + \frac{(-1 - i\sqrt{3})(x^2 - 3)}{2\left(-54x^3 - 54c_1 - 108x + 6\sqrt{75x^6 + 162c_1x^3 + 378x^4 + 81c_1^2 + 324c_1x + 162x^2 + 162}\right)^{\frac{1}{3}}}$$

✓ Solution by Mathematica

Time used: 8.724 (sec). Leaf size: 421

DSolve[2+3*x^2-2*x*y[x]+(3-x^2+6*y[x]^2)*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]

$$y(x) \rightarrow -\frac{x^2 - 3}{\frac{\sqrt[3]{6}\sqrt[3]{9x^3 + \sqrt{3}\sqrt{-2(x^2 - 3)^3 + 27(x^3 + 2x + c_1)^2 + 18x + 9c_1}}}{\sqrt[3]{9x^3 + \sqrt{3}\sqrt{-2(x^2 - 3)^3 + 27(x^3 + 2x + c_1)^2 + 18x + 9c_1}}}} - \frac{6^{2/3}}{6^{2/3}}$$

$$y(x) \rightarrow \frac{\sqrt[3]{6}(1 + i\sqrt{3})(x^2 - 3) + (1 - i\sqrt{3})\left(9x^3 + \sqrt{3}\sqrt{-2(x^2 - 3)^3 + 27(x^3 + 2x + c_1)^2 + 18x + 9c_1}\right)^{2/3}}{2 \cdot 6^{2/3}\sqrt[3]{9x^3 + \sqrt{3}\sqrt{-2(x^2 - 3)^3 + 27(x^3 + 2x + c_1)^2 + 18x + 9c_1}}}$$

$$y(x) \rightarrow \frac{\sqrt[3]{6}(1 - i\sqrt{3})(x^2 - 3) + (1 + i\sqrt{3})\left(9x^3 + \sqrt{3}\sqrt{-2(x^2 - 3)^3 + 27(x^3 + 2x + c_1)^2 + 18x + 9c_1}\right)^{2/3}}{2 \cdot 6^{2/3}\sqrt[3]{9x^3 + \sqrt{3}\sqrt{-2(x^2 - 3)^3 + 27(x^3 + 2x + c_1)^2 + 18x + 9c_1}}}$$

5.4 problem 4

5.4.1	Solving as separable ode	1187
5.4.2	Solving as linear ode	1189
5.4.3	Solving as homogeneousTypeD2 ode	1190
5.4.4	Solving as differentialType ode	1192
5.4.5	Solving as first order ode lie symmetry lookup ode	1193
5.4.6	Solving as exact ode	1197
5.4.7	Maple step by step solution	1201

Internal problem ID [546]

Internal file name [OUTPUT/546_Sunday_June_05_2022_01_44_05_AM_3993535/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 4.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "linear", "separable", "differentialType", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

`[_separable]`

$$2y + 2xy^2 + (2x + 2x^2y) y' = 0$$

5.4.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\frac{y}{x} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(y) = y$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y} dy &= -\frac{1}{x} dx \\ \int \frac{1}{y} dy &= \int -\frac{1}{x} dx \\ \ln(y) &= -\ln(x) + c_1 \\ y &= e^{-\ln(x)+c_1} \\ &= \frac{c_1}{x}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{x} \tag{1}$$

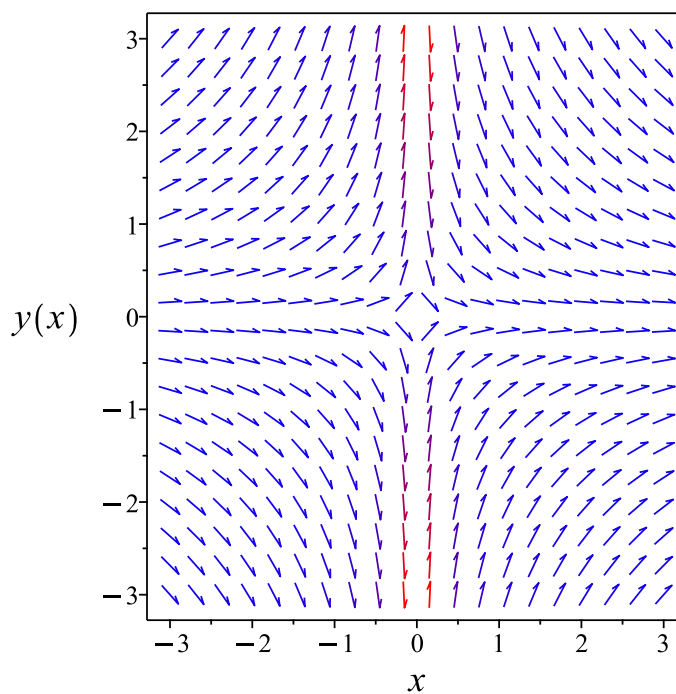


Figure 231: Slope field plot

Verification of solutions

$$y = \frac{c_1}{x}$$

Verified OK.

5.4.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{1}{x}$$
$$q(x) = 0$$

Hence the ode is

$$y' + \frac{y}{x} = 0$$

The integrating factor μ is

$$\mu = e^{\int \frac{1}{x} dx}$$
$$= x$$

The ode becomes

$$\frac{d}{dx} \mu y = 0$$
$$\frac{d}{dx} (yx) = 0$$

Integrating gives

$$yx = c_1$$

Dividing both sides by the integrating factor $\mu = x$ results in

$$y = \frac{c_1}{x}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{x} \tag{1}$$

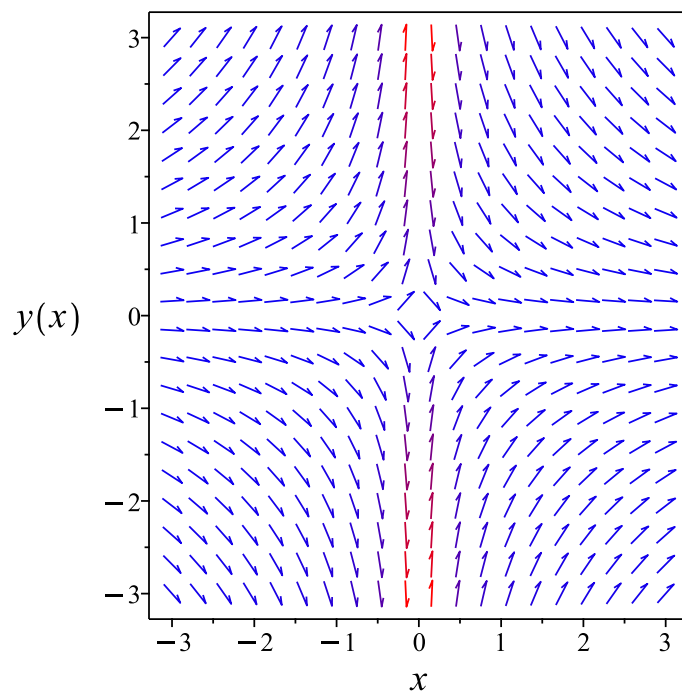


Figure 232: Slope field plot

Verification of solutions

$$y = \frac{c_1}{x}$$

Verified OK.

5.4.3 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$2u(x)x + 2x^3u(x)^2 + (2x + 2x^3u(x))(u'(x)x + u(x)) = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{2u}{x} \end{aligned}$$

Where $f(x) = -\frac{2}{x}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= -\frac{2}{x} dx \\ \int \frac{1}{u} du &= \int -\frac{2}{x} dx \\ \ln(u) &= -2 \ln(x) + c_2 \\ u &= e^{-2 \ln(x) + c_2} \\ &= \frac{c_2}{x^2}\end{aligned}$$

Therefore the solution y is

$$\begin{aligned}y &= ux \\ &= \frac{c_2}{x}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_2}{x} \tag{1}$$

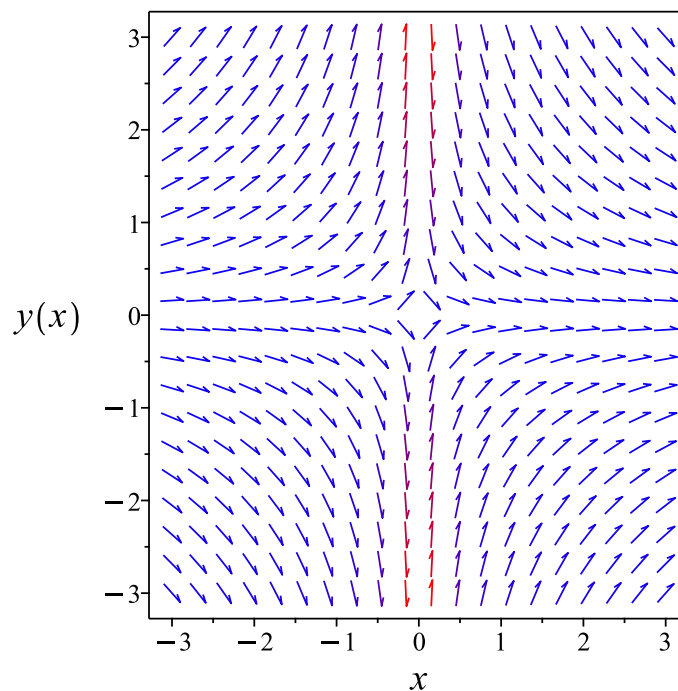


Figure 233: Slope field plot

Verification of solutions

$$y = \frac{c_2}{x}$$

Verified OK.

5.4.4 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-2y - 2xy^2}{2x + 2x^2y} \quad (1)$$

Which becomes

$$0 = (-x) dy + (-y) dx \quad (2)$$

But the RHS is complete differential because

$$(-x) dy + (-y) dx = d(-yx)$$

Hence (2) becomes

$$0 = d(-yx)$$

Integrating both sides gives gives these solutions

$$y = \frac{c_1}{x} + c_1$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{x} + c_1 \quad (1)$$

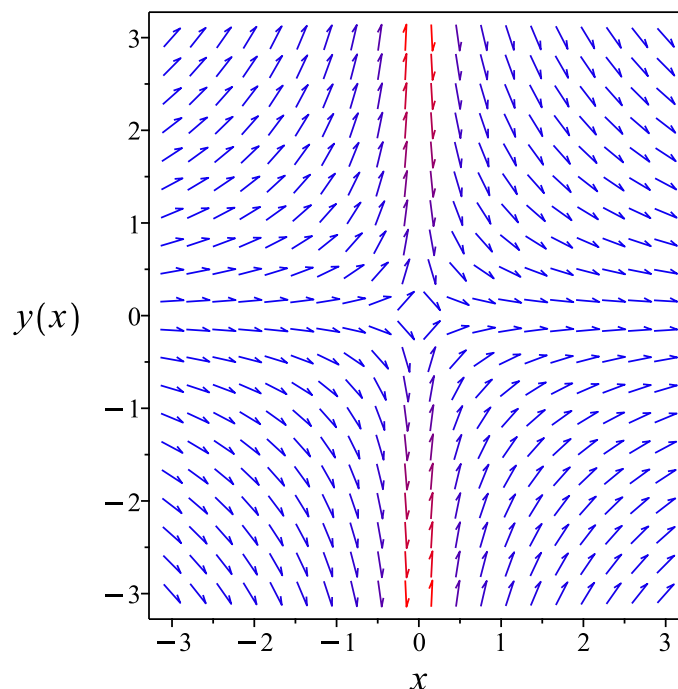


Figure 234: Slope field plot

Verification of solutions

$$y = \frac{c_1}{x} + c_2$$

Verified OK.

5.4.5 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{y}{x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 238: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int(n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{x}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{x}} dy \end{aligned}$$

Which results in

$$S = yx$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y}{x}$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 0$$

$$S_x = y$$

$$S_y = x$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$yx = c_1$$

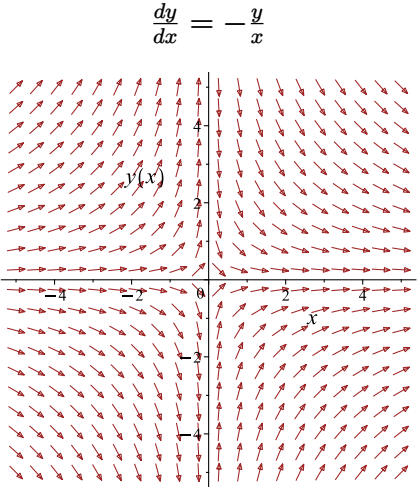
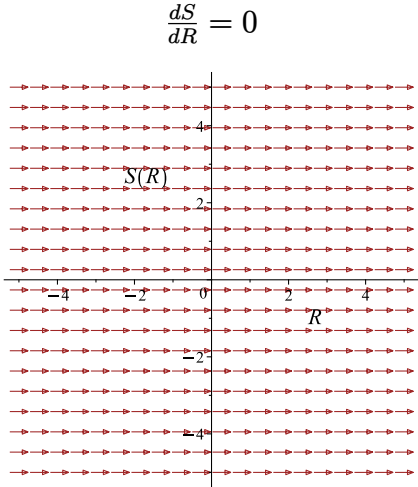
Which simplifies to

$$y = \frac{c_1}{x}$$

Which gives

$$y = \frac{c_1}{x}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{y}{x}$ 	$R = x$ $S = yx$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$y = \frac{c_1}{x} \quad (1)$$

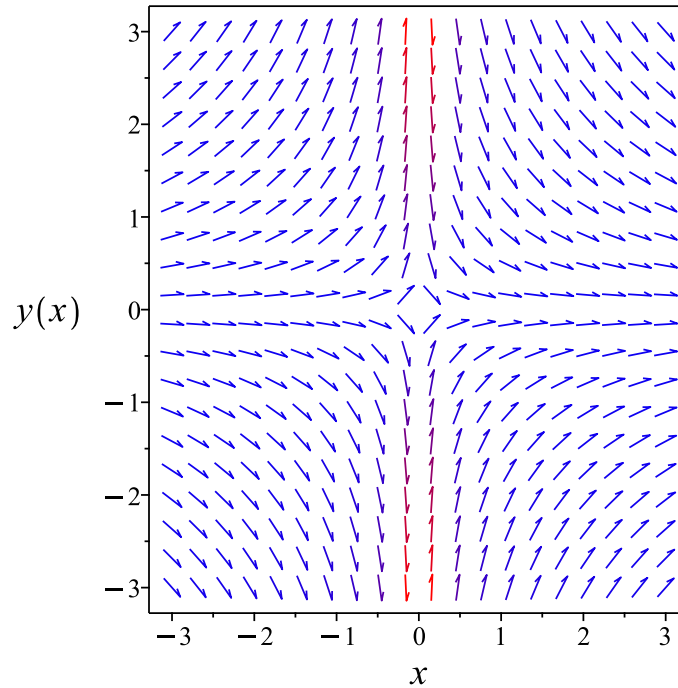


Figure 235: Slope field plot

Verification of solutions

$$y = \frac{c_1}{x}$$

Verified OK.

5.4.6 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(-\frac{1}{y}\right) dy &= \left(\frac{1}{x}\right) dx \\ \left(-\frac{1}{x}\right) dx + \left(-\frac{1}{y}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{1}{x} \\ N(x, y) &= -\frac{1}{y}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{1}{x}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(-\frac{1}{y} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{1}{x} dx \\ \phi &= -\ln(x) + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{1}{y}$. Therefore equation (4) becomes

$$-\frac{1}{y} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{1}{y}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(-\frac{1}{y} \right) dy \\ f(y) &= -\ln(y) + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\ln(x) - \ln(y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\ln(x) - \ln(y)$$

The solution becomes

$$y = \frac{e^{-c_1}}{x}$$

Summary

The solution(s) found are the following

$$y = \frac{e^{-c_1}}{x} \tag{1}$$

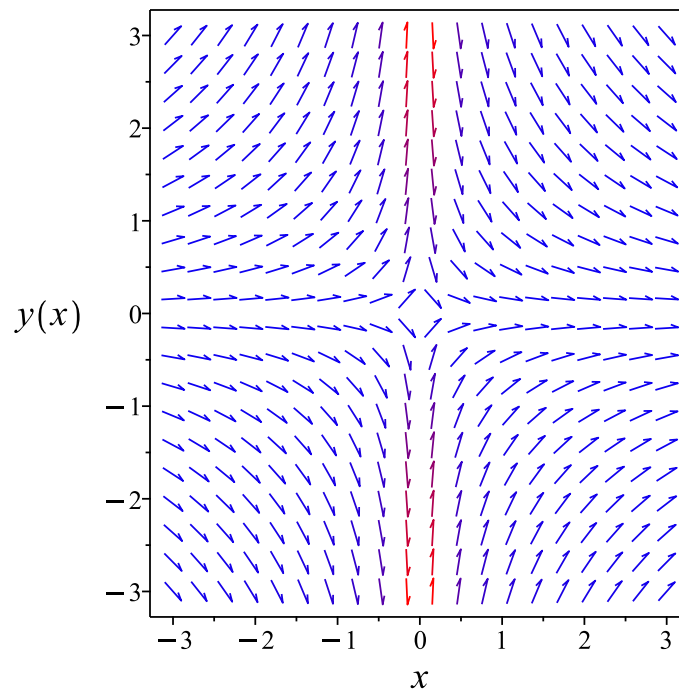


Figure 236: Slope field plot

Verification of solutions

$$y = \frac{e^{-c_1}}{x}$$

Verified OK.

5.4.7 Maple step by step solution

Let's solve

$$2y + 2xy^2 + (2x + 2x^2y) y' = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Integrate both sides with respect to x

$$\int (2y + 2xy^2 + (2x + 2x^2y) y') dx = \int 0 dx + c_1$$

- Evaluate integral

$$2x \left(\frac{xy^2}{2} + y \right) = c_1$$

- Solve for y

$$\left\{ y = \frac{-1 - \sqrt{1 + c_1}}{x}, y = \frac{-1 + \sqrt{1 + c_1}}{x} \right\}$$

Maple trace

```
`Classification methods on request
Methods to be used are: [exact]
-----
* Tackling ODE using method: exact
--- Trying classification methods ---
trying exact
<- exact successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 31

```
dsolve(2*y(x)+2*x*y(x)^2+(2*x+2*x^2*y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = -\frac{1}{x}$$
$$y(x) = \frac{-1 - c_1}{x}$$
$$y(x) = \frac{c_1 - 1}{x}$$

✓ Solution by Mathematica

Time used: 0.023 (sec). Leaf size: 29

```
DSolve[2*y[x]+2*x*y[x]^2+(2*x+2*x^2*y[x])*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{1}{x}$$
$$y(x) \rightarrow \frac{c_1}{x}$$
$$y(x) \rightarrow -\frac{1}{x}$$

5.5 problem 5

5.5.1	Solving as homogeneousTypeD2 ode	1203
5.5.2	Solving as first order ode lie symmetry calculated ode	1205
5.5.3	Solving as exact ode	1209

Internal problem ID [547]

Internal file name [OUTPUT/547_Sunday_June_05_2022_01_44_06_AM_24958851/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 5.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "homogeneousTypeD2", "first_order_ode_lie_symmetry_calculated"**

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$y' - \frac{-ax - by}{bx + cy} = 0$$

5.5.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u'(x)x + u(x) - \frac{-ax - bu(x)x}{bx + cu(x)x} = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{cu^2 + 2bu + a}{x(cu + b)} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(u) = \frac{cu^2+2bu+a}{cu+b}$. Integrating both sides gives

$$\frac{1}{\frac{cu^2+2bu+a}{cu+b}} du = -\frac{1}{x} dx$$

$$\int \frac{1}{\frac{cu^2+2bu+a}{cu+b}} du = \int -\frac{1}{x} dx$$

$$\frac{\ln(cu^2 + 2bu + a)}{2} = -\ln(x) + c_2$$

Raising both side to exponential gives

$$\sqrt{cu^2 + 2bu + a} = e^{-\ln(x)+c_2}$$

Which simplifies to

$$\sqrt{cu^2 + 2bu + a} = \frac{c_3}{x}$$

Which simplifies to

$$\sqrt{u(x)^2 c + 2u(x)b + a} = \frac{c_3 e^{c_2}}{x}$$

The solution is

$$\sqrt{u(x)^2 c + 2u(x)b + a} = \frac{c_3 e^{c_2}}{x}$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\sqrt{\frac{y^2 c}{x^2} + \frac{2yb}{x} + a} = \frac{c_3 e^{c_2}}{x}$$

$$\sqrt{\frac{cy^2 + 2byx + ax^2}{x^2}} = \frac{c_3 e^{c_2}}{x}$$

Summary

The solution(s) found are the following

$$\sqrt{\frac{cy^2 + 2byx + ax^2}{x^2}} = \frac{c_3 e^{c_2}}{x} \quad (1)$$

Verification of solutions

$$\sqrt{\frac{cy^2 + 2byx + ax^2}{x^2}} = \frac{c_3 e^{c_2}}{x}$$

Verified OK.

5.5.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{ax + by}{bx + cy}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 - \frac{(ax + by)(b_3 - a_2)}{bx + cy} - \frac{(ax + by)^2 a_3}{(bx + cy)^2}$$

$$- \left(-\frac{a}{bx + cy} + \frac{(ax + by)b}{(bx + cy)^2} \right) (xa_2 + ya_3 + a_1) \quad (\text{5E})$$

$$- \left(-\frac{b}{bx + cy} + \frac{(ax + by)c}{(bx + cy)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{a^2 x^2 a_3 - ab x^2 a_2 + ab x^2 b_3 + 2abxya_3 + ac x^2 b_2 - 2acxya_2 + 2acxyb_3 - ac y^2 a_3 - 2b^2 x^2 b_2 + 2b^2 y^2 a_3 - 2b^2 xy b_3 + 2bcxyb_2 + bc y^2 a_2 - bc y^2 b_3 + c^2 y^2 b_2 - acxb_1 + acya_1 + b^2 xb_1 - b^2 ya_1}{(bx + cy)^2} = 0$$

Setting the numerator to zero gives

$$-a^2 x^2 a_3 + ab x^2 a_2 - ab x^2 b_3 - 2abxya_3 - ac x^2 b_2 + 2acxya_2$$

$$- 2acxyb_3 + ac y^2 a_3 + 2b^2 x^2 b_2 - 2b^2 y^2 a_3 + 2bcxyb_2 + bc y^2 a_2$$

$$- bc y^2 b_3 + c^2 y^2 b_2 - acxb_1 + acya_1 + b^2 xb_1 - b^2 ya_1 = 0 \quad (\text{6E})$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & -a^2 a_3 v_1^2 + a b a_2 v_1^2 - 2 a b a_3 v_1 v_2 - a b b_3 v_1^2 + 2 a c a_2 v_1 v_2 + a c a_3 v_2^2 \\ & - a c b_2 v_1^2 - 2 a c b_3 v_1 v_2 - 2 b^2 a_3 v_2^2 + 2 b^2 b_2 v_1^2 + b c a_2 v_2^2 + 2 b c b_2 v_1 v_2 \\ & - b c b_3 v_2^2 + c^2 b_2 v_2^2 + a c a_1 v_2 - a c b_1 v_1 - b^2 a_1 v_2 + b^2 b_1 v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & (-a^2 a_3 + a b a_2 - a b b_3 - a c b_2 + 2 b^2 b_2) v_1^2 + (-2 a b a_3 + 2 a c a_2 - 2 a c b_3 + 2 b c b_2) v_1 v_2 \\ & + (-a c b_1 + b^2 b_1) v_1 + (a c a_3 - 2 b^2 a_3 + b c a_2 - b c b_3 + c^2 b_2) v_2^2 + (a c a_1 - b^2 a_1) v_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} & a c a_1 - b^2 a_1 = 0 \\ & -a c b_1 + b^2 b_1 = 0 \\ & -2 a b a_3 + 2 a c a_2 - 2 a c b_3 + 2 b c b_2 = 0 \\ & a c a_3 - 2 b^2 a_3 + b c a_2 - b c b_3 + c^2 b_2 = 0 \\ & -a^2 a_3 + a b a_2 - a b b_3 - a c b_2 + 2 b^2 b_2 = 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= \frac{a b_3 - 2 b b_2}{a} \\ a_3 &= -\frac{c b_2}{a} \\ b_1 &= 0 \\ b_2 &= b_2 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= x \\ \eta &= y\end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left(-\frac{ax + by}{bx + cy} \right) (x) \\ &= \frac{ax^2 + 2bxy + cy^2}{bx + cy} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{ax^2 + 2bxy + cy^2}{bx + cy}} dy\end{aligned}$$

Which results in

$$S = \frac{\ln(ax^2 + 2bxy + cy^2)}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{ax + by}{bx + cy}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{ax + by}{ax^2 + 2bxy + cy^2} \\ S_y &= \frac{bx + cy}{ax^2 + 2bxy + cy^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\ln(cy^2 + 2bxy + ax^2)}{2} = c_1$$

Which simplifies to

$$\frac{\ln(cy^2 + 2bxy + ax^2)}{2} = c_1$$

Summary

The solution(s) found are the following

$$\frac{\ln(cy^2 + 2bxy + ax^2)}{2} = c_1 \tag{1}$$

Verification of solutions

$$\frac{\ln (cy^2 + 2byx + ax^2)}{2} = c_1$$

Verified OK.

5.5.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (bx + cy) dy &= (-ax - by) dx \\ (ax + by) dx + (bx + cy) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$M(x, y) = ax + by$$

$$N(x, y) = bx + cy$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(ax + by) \\ &= b\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(bx + cy) \\ &= b\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int ax + by dx$$

$$\phi = \frac{1}{2}ax^2 + bxy + f(y) \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = bx + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = bx + cy$. Therefore equation (4) becomes

$$bx + cy = bx + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = cy$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (cy) dy$$
$$f(y) = \frac{cy^2}{2} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{1}{2}ax^2 + bxy + \frac{1}{2}cy^2 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{1}{2}ax^2 + bxy + \frac{1}{2}cy^2$$

Summary

The solution(s) found are the following

$$\frac{ax^2}{2} + bxy + \frac{cy^2}{2} = c_1 \quad (1)$$

Verification of solutions

$$\frac{ax^2}{2} + bxy + \frac{cy^2}{2} = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 77

```
dsolve(diff(y(x),x) = (-a*x-b*y(x))/(b*x+c*y(x)),y(x), singsol=all)
```

$$y(x) = \frac{-bxc_1 + \sqrt{-x^2(ac - b^2)c_1^2 + c}}{cc_1}$$
$$y(x) = \frac{-bxc_1 - \sqrt{-x^2(ac - b^2)c_1^2 + c}}{cc_1}$$

✓ Solution by Mathematica

Time used: 17.783 (sec). Leaf size: 139

```
DSolve[y'[x]== (-a*x-b*y[x])/(b*x+c*y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{bx + \sqrt{-acx^2 + b^2x^2 + ce^{2c_1}}}{c}$$
$$y(x) \rightarrow \frac{-bx + \sqrt{b^2x^2 + c(-ax^2 + e^{2c_1})}}{c}$$
$$y(x) \rightarrow -\frac{\sqrt{x^2(b^2 - ac)} + bx}{c}$$
$$y(x) \rightarrow \frac{\sqrt{x^2(b^2 - ac)} - bx}{c}$$

5.6 problem 6

- 5.6.1 Solving as homogeneousTypeD2 ode 1213
- 5.6.2 Solving as first order ode lie symmetry calculated ode 1215

Internal problem ID [548]

Internal file name [OUTPUT/548_Sunday_June_05_2022_01_44_07_AM_46494958/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 6.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD2**", "**first_order_ode_lie_symmetry_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$y' - \frac{-ax + by}{bx - cy} = 0$$

5.6.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u'(x)x + u(x) - \frac{-ax + bu(x)x}{bx - cu(x)x} = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{-cu^2 + a}{x(-cu + b)} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(u) = \frac{-cu^2+a}{-cu+b}$. Integrating both sides gives

$$\frac{1}{\frac{-cu^2+a}{-cu+b}} du = -\frac{1}{x} dx$$

$$\int \frac{1}{\frac{-cu^2+a}{-cu+b}} du = \int -\frac{1}{x} dx$$

$$\frac{\ln(cu^2 - a)}{2} + \frac{b \operatorname{arctanh}\left(\frac{uc}{\sqrt{ac}}\right)}{\sqrt{ac}} = -\ln(x) + c_2$$

The solution is

$$\frac{\ln(u(x)^2 c - a)}{2} + \frac{b \operatorname{arctanh}\left(\frac{u(x)c}{\sqrt{ac}}\right)}{\sqrt{ac}} + \ln(x) - c_2 = 0$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\frac{\ln\left(\frac{y^2 c}{x^2} - a\right)}{2} + \frac{b \operatorname{arctanh}\left(\frac{yc}{x\sqrt{ac}}\right)}{\sqrt{ac}} + \ln(x) - c_2 = 0$$

$$\frac{\ln\left(\frac{y^2 c}{x^2} - a\right)}{2} + \frac{b \operatorname{arctanh}\left(\frac{yc}{x\sqrt{ac}}\right)}{\sqrt{ac}} + \ln(x) - c_2 = 0$$

Summary

The solution(s) found are the following

$$\frac{\ln\left(\frac{y^2 c}{x^2} - a\right)}{2} + \frac{b \operatorname{arctanh}\left(\frac{yc}{x\sqrt{ac}}\right)}{\sqrt{ac}} + \ln(x) - c_2 = 0 \quad (1)$$

Verification of solutions

$$\frac{\ln\left(\frac{y^2 c}{x^2} - a\right)}{2} + \frac{b \operatorname{arctanh}\left(\frac{yc}{x\sqrt{ac}}\right)}{\sqrt{ac}} + \ln(x) - c_2 = 0$$

Verified OK.

5.6.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{-ax + by}{-bx + cy}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (A)$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1E)$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2E)$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 - \frac{(-ax + by)(b_3 - a_2)}{-bx + cy} - \frac{(-ax + by)^2 a_3}{(-bx + cy)^2}$$

$$- \left(\frac{a}{-bx + cy} - \frac{(-ax + by)b}{(-bx + cy)^2} \right) (xa_2 + ya_3 + a_1) \quad (5E)$$

$$- \left(-\frac{b}{-bx + cy} + \frac{(-ax + by)c}{(-bx + cy)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{-a^2 x^2 a_3 - ab x^2 a_2 + ab x^2 b_3 - 2abxy a_3 - ac x^2 b_2 + 2acxy a_2 - 2acxy b_3 + ac y^2 a_3 + 2bcxy b_2 - bc y^2 a_2 + b_1}{(bx - cy)^2}$$

$$= 0$$

Setting the numerator to zero gives

$$-a^2 x^2 a_3 + ab x^2 a_2 - ab x^2 b_3 + 2abxy a_3 + ac x^2 b_2 - 2acxy a_2 + 2acxy b_3 - ac y^2 a_3 \quad (6E)$$

$$- 2bcxy b_2 + bc y^2 a_2 - bc y^2 b_3 + c^2 y^2 b_2 + acx b_1 - acy a_1 - b^2 x b_1 + b^2 y a_1 = 0$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & -a^2 a_3 v_1^2 + a b a_2 v_1^2 + 2 a b a_3 v_1 v_2 - a b b_3 v_1^2 - 2 a c a_2 v_1 v_2 \\ & - a c a_3 v_2^2 + a c b_2 v_1^2 + 2 a c b_3 v_1 v_2 + b c a_2 v_2^2 - 2 b c b_2 v_1 v_2 \\ & - b c b_3 v_2^2 + c^2 b_2 v_2^2 - a c a_1 v_2 + a c b_1 v_1 + b^2 a_1 v_2 - b^2 b_1 v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & (-a^2 a_3 + a b a_2 - a b b_3 + a c b_2) v_1^2 + (2 a b a_3 - 2 a c a_2 + 2 a c b_3 - 2 b c b_2) v_1 v_2 \\ & + (a c b_1 - b^2 b_1) v_1 + (-a c a_3 + b c a_2 - b c b_3 + c^2 b_2) v_2^2 + (-a c a_1 + b^2 a_1) v_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} & -a c a_1 + b^2 a_1 = 0 \\ & a c b_1 - b^2 b_1 = 0 \\ & -a c a_3 + b c a_2 - b c b_3 + c^2 b_2 = 0 \\ & 2 a b a_3 - 2 a c a_2 + 2 a c b_3 - 2 b c b_2 = 0 \\ & -a^2 a_3 + a b a_2 - a b b_3 + a c b_2 = 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= a_3 \\ b_1 &= 0 \\ b_2 &= \frac{a_3 a}{c} \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= x \\ \eta &= y\end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left(\frac{-ax + by}{-bx + cy} \right) (x) \\ &= \frac{ax^2 - cy^2}{bx - cy} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{ax^2 - cy^2}{bx - cy}} dy\end{aligned}$$

Which results in

$$S = \frac{\ln(-ax^2 + cy^2)}{2} + \frac{b \operatorname{arctanh}\left(\frac{yc}{x\sqrt{ac}}\right)}{\sqrt{ac}}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{-ax + by}{-bx + cy}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{ax - by}{ax^2 - cy^2} \\ S_y &= \frac{bx - cy}{ax^2 - cy^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\ln(cy^2 - ax^2) \sqrt{a} \sqrt{c} + 2b \operatorname{arctanh}\left(\frac{y\sqrt{c}}{x\sqrt{a}}\right)}{2\sqrt{a} \sqrt{c}} = c_1$$

Which simplifies to

$$\frac{\ln (cy^2 - a x^2) \sqrt{a} \sqrt{c} + 2b \operatorname{arctanh} \left(\frac{y\sqrt{c}}{x\sqrt{a}} \right)}{2\sqrt{a} \sqrt{c}} = c_1$$

Summary

The solution(s) found are the following

$$\frac{\ln (cy^2 - a x^2) \sqrt{a} \sqrt{c} + 2b \operatorname{arctanh} \left(\frac{y\sqrt{c}}{x\sqrt{a}} \right)}{2\sqrt{a} \sqrt{c}} = c_1 \quad (1)$$

Verification of solutions

$$\frac{\ln (cy^2 - a x^2) \sqrt{a} \sqrt{c} + 2b \operatorname{arctanh} \left(\frac{y\sqrt{c}}{x\sqrt{a}} \right)}{2\sqrt{a} \sqrt{c}} = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.422 (sec). Leaf size: 47

```
dsolve(diff(y(x),x) = (-a*x+b*y(x))/(b*x-c*y(x)),y(x), singsol=all)
```

$$y(x) = \operatorname{RootOf} \left(c_Z^2 - a - e^{\operatorname{RootOf} \left(e^{-Z} \cosh \left(\frac{\sqrt{ac} (2c_1 + _Z + 2 \ln(x))}{2b} \right)^2 + a \right)} \right) x$$

✓ Solution by Mathematica

Time used: 0.072 (sec). Leaf size: 58

```
DSolve[y'[x] == (-a*x+b*y[x])/(b*x-c*y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve} \left[-\frac{\text{arctanh}\left(\frac{\sqrt{cy(x)}}{\sqrt{ax}}\right)}{\sqrt{a}\sqrt{c}} - \frac{1}{2} \log\left(\frac{cy(x)^2}{x^2} - a\right) = \log(x) + c_1, y(x) \right]$$

5.7 problem 7

5.7.1 Solving as exact ode	1221
5.7.2 Maple step by step solution	1224

Internal problem ID [549]

Internal file name [OUTPUT/549_Sunday_June_05_2022_01_44_09_AM_4332274/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 7.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact"**

Maple gives the following as the ode type

[_exact]

$$e^x \sin(y) - 2 \sin(x) y + (2 \cos(x) + e^x \cos(y)) y' = 0$$

5.7.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (2 \cos(x) + e^x \cos(y)) dy &= (-e^x \sin(y) + 2 \sin(x) y) dx \\ (e^x \sin(y) - 2 \sin(x) y) dx &+ (2 \cos(x) + e^x \cos(y)) dy = 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= e^x \sin(y) - 2 \sin(x) y \\ N(x, y) &= 2 \cos(x) + e^x \cos(y) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (e^x \sin(y) - 2 \sin(x) y) \\ &= -2 \sin(x) + e^x \cos(y) \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (2 \cos(x) + e^x \cos(y)) \\ &= -2 \sin(x) + e^x \cos(y) \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int e^x \sin(y) - 2 \sin(x) y dx \\ \phi &= e^x \sin(y) + 2 \cos(x) y + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 2 \cos(x) + e^x \cos(y) + f'(y)\quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 2 \cos(x) + e^x \cos(y)$. Therefore equation (4) becomes

$$2 \cos(x) + e^x \cos(y) = 2 \cos(x) + e^x \cos(y) + f'(y)\quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = e^x \sin(y) + 2 \cos(x) y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = e^x \sin(y) + 2 \cos(x) y$$

Summary

The solution(s) found are the following

$$e^x \sin(y) + 2 \cos(x) y = c_1\quad (1)$$

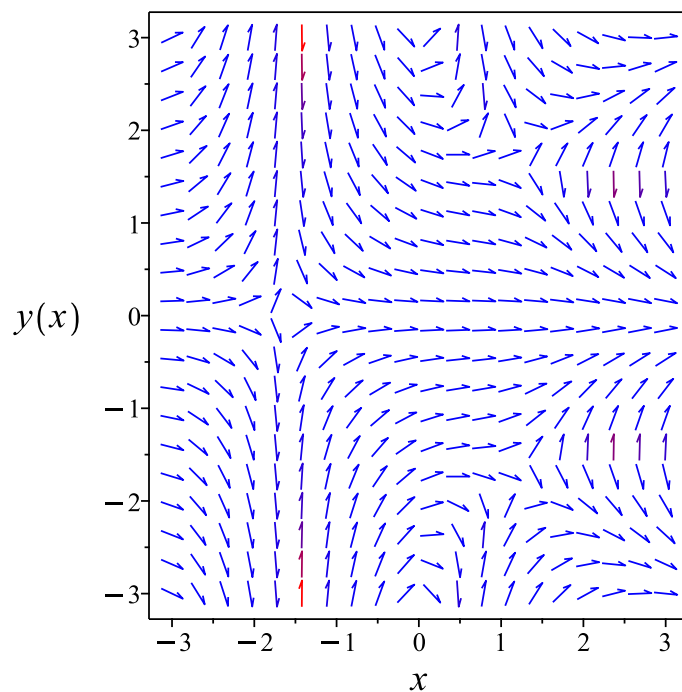


Figure 237: Slope field plot

Verification of solutions

$$e^x \sin(y) + 2 \cos(x) y = c_1$$

Verified OK.

5.7.2 Maple step by step solution

Let's solve

$$e^x \sin(y) - 2 \sin(x) y + (2 \cos(x) + e^x \cos(y)) y' = 0$$

- Highest derivative means the order of the ODE is 1
 y'
- Check if ODE is exact
 - ODE is exact if the lhs is the total derivative of a C^2 function
 $F'(x, y) = 0$
 - Compute derivative of lhs
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$

- Evaluate derivatives

$$-2 \sin(x) + e^x \cos(y) = -2 \sin(x) + e^x \cos(y)$$
- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$\left[F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$
- Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to x

$$F(x, y) = \int (e^x \sin(y) - 2 \sin(x) y) dx + f_1(y)$$
- Evaluate integral

$$F(x, y) = e^x \sin(y) + 2 \cos(x) y + f_1(y)$$
- Take derivative of $F(x, y)$ with respect to y

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$
- Compute derivative

$$2 \cos(x) + e^x \cos(y) = e^x \cos(y) + 2 \cos(x) + \frac{d}{dy} f_1(y)$$
- Isolate for $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = 0$$
- Solve for $f_1(y)$

$$f_1(y) = 0$$
- Substitute $f_1(y)$ into equation for $F(x, y)$

$$F(x, y) = e^x \sin(y) + 2 \cos(x) y$$
- Substitute $F(x, y)$ into the solution of the ODE

$$e^x \sin(y) + 2 \cos(x) y = c_1$$
- Solve for y

$$y = \text{RootOf}(-e^x \sin(_Z) - 2_Z \cos(x) + c_1)$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 17

```
dsolve(exp(x)*sin(y(x))-2*sin(x)*y(x)+(2*cos(x)+exp(x)*cos(y(x)))*diff(y(x),x) = 0,y(x), sin
```

$$e^x \sin(y(x)) + 2 \cos(x) y(x) + c_1 = 0$$

✓ Solution by Mathematica

Time used: 0.282 (sec). Leaf size: 20

```
DSolve[Exp[x]*Sin[y[x]]-2*Ssin[x]*y[x]+(2*Cos[x]+Exp[x]*Cos[y[x]])*y'[x] == 0,y[x],x,IncludeS
```

$$\text{Solve}[e^x \sin(y(x)) + 2y(x) \cos(x) = c_1, y(x)]$$

5.8 problem 8

Internal problem ID [550]

Internal file name [OUTPUT/550_Sunday_June_05_2022_01_44_16_AM_39437276/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 8.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**unknown**"

Maple gives the following as the ode type

[`x=_G(y,y')`]

Unable to solve or complete the solution.

$$e^x \sin(y) + 3y - (3x - e^x \sin(y)) y' = 0$$

Unable to determine ODE type.

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
Looking for potential symmetries
trying inverse_Riccati
trying an equivalence to an Abel ODE
differential order: 1; trying a linearization to 2nd order
--- trying a change of variables {x -> y(x), y(x) -> x}
differential order: 1; trying a linearization to 2nd order
trying 1st order ODE linearizable_by_differentiation
--- Trying Lie symmetry methods, 1st order ---
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = 4
`, `-> Computing symmetries using: way = 5
trying symmetry patterns for 1st order ODEs
-> trying a symmetry pattern of the form [F(x)*G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)*G(y)]
-> trying symmetry patterns of the forms [F(x),G(y)] and [G(y),F(x)]
`, `-> Computing symmetries using: way = HINT
-> Calling odsolve with the ODE`, diff(y(x), x) = 0, y(x)` *** Sublevel 2 ***
  Methods for first order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  trying 1st order linear
  <- 1st order linear successful
-> Calling odsolve with the ODE`, diff(y(x), x) = y(x)/x, y(x)` *** Sublevel 2 ***
  Methods for first order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  trying 1st order linear
  <- 1st order linear successful
-> Calling odsolve with the ODE`, diff(y(x), x)-y(x)/x, y(x)` *** Sublevel 2 ***
  Methods for first order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  trying 1st order linear
  <- 1st order linear successful
```

X Solution by Maple

```
dsolve(exp(x)*sin(y(x))+3*y(x)-(3*x-exp(x)*sin(y(x)))*diff(y(x),x) = 0,y(x), singsol=all)
```

No solution found

X Solution by Mathematica

Time used: 0.0 (sec). Leaf size: 0

```
DSolve[Exp[x]*Sin[y[x]]+3*y[x]-(3*x-Exp[x]*Sin[y[x]])*y'[x] == 0,y[x],x,IncludeSingularSolut
```

Not solved

5.9 problem 9

5.9.1 Solving as exact ode	1230
5.9.2 Maple step by step solution	1234

Internal problem ID [551]

Internal file name [OUTPUT/551_Sunday_June_05_2022_01_44_22_AM_93677906/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 9.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact"**

Maple gives the following as the ode type

[_exact]

$$-2 e^{yx} \sin(2x) + e^{yx} \cos(2x) y + (-3 + e^{yx} x \cos(2x)) y' = -2x$$

5.9.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (-3 + e^{yx} x \cos(2x)) dy &= (-2x + 2 e^{yx} \sin(2x) - e^{yx} \cos(2x) y) \\ (2x - 2 e^{yx} \sin(2x) + e^{yx} \cos(2x) y) dx &+ (-3 + e^{yx} x \cos(2x)) dy = 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 2x - 2 e^{yx} \sin(2x) + e^{yx} \cos(2x) y \\ N(x, y) &= -3 + e^{yx} x \cos(2x) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (2x - 2 e^{yx} \sin(2x) + e^{yx} \cos(2x) y) \\ &= e^{yx} (yx \cos(2x) + \cos(2x) - 2x \sin(2x)) \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (-3 + e^{yx} x \cos(2x)) \\ &= e^{yx} (yx \cos(2x) + \cos(2x) - 2x \sin(2x)) \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 2x - 2e^{yx} \sin(2x) + e^{yx} \cos(2x) y dx \\ \phi &= e^{yx} \cos(2x) + x^2 + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{yx} x \cos(2x) + f'(y)\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -3 + e^{yx} x \cos(2x)$. Therefore equation (4) becomes

$$-3 + e^{yx} x \cos(2x) = e^{yx} x \cos(2x) + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -3$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (-3) dy \\ f(y) &= -3y + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = e^{yx} \cos(2x) + x^2 - 3y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = e^{yx} \cos(2x) + x^2 - 3y$$

The solution becomes

$$y = -\frac{-x^3 + c_1x + 3 \operatorname{LambertW}\left(-\frac{x \cos(2x)e^{\frac{1}{3}x^3 - \frac{1}{3}c_1x}}{3}\right)}{3x}$$

Summary

The solution(s) found are the following

$$y = -\frac{-x^3 + c_1x + 3 \operatorname{LambertW}\left(-\frac{x \cos(2x)e^{\frac{1}{3}x^3 - \frac{1}{3}c_1x}}{3}\right)}{3x} \tag{1}$$

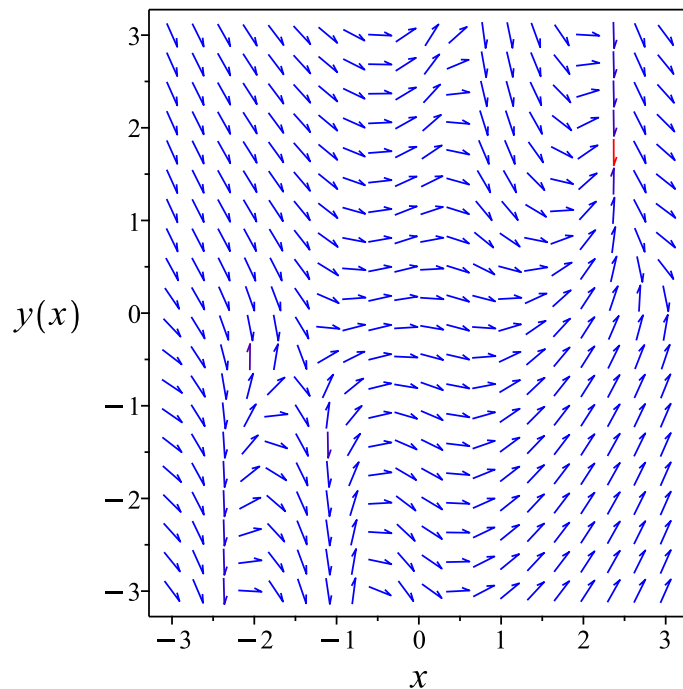


Figure 238: Slope field plot

Verification of solutions

$$y = -\frac{-x^3 + c_1x + 3 \operatorname{LambertW}\left(-\frac{x \cos(2x)e^{\frac{1}{3}x^3 - \frac{1}{3}c_1x}}{3}\right)}{3x}$$

Verified OK.

5.9.2 Maple step by step solution

Let's solve

$$-2 e^{yx} \sin(2x) + e^{yx} \cos(2x) y + (-3 + e^{yx} x \cos(2x)) y' = -2x$$

- Highest derivative means the order of the ODE is 1

y'

- Check if ODE is exact

- ODE is exact if the lhs is the total derivative of a C^2 function

$$F'(x, y) = 0$$

- Compute derivative of lhs

$$F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y) \right) y' = 0$$

- Evaluate derivatives

$$e^{yx} y x \cos(2x) + e^{yx} \cos(2x) - 2 e^{yx} x \sin(2x) = e^{yx} y x \cos(2x) + e^{yx} \cos(2x) - 2 e^{yx} x \sin(2x)$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$\left[F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$

- Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to x

$$F(x, y) = \int (2x - 2 e^{yx} \sin(2x) + e^{yx} \cos(2x) y) dx + f_1(y)$$

- Evaluate integral

$$F(x, y) = y \left(\frac{y e^{yx} \cos(2x)}{y^2+4} + \frac{2 e^{yx} \sin(2x)}{y^2+4} \right) + x^2 + \frac{4 e^{yx} \cos(2x)}{y^2+4} - \frac{2y e^{yx} \sin(2x)}{y^2+4} + f_1(y)$$

- Take derivative of $F(x, y)$ with respect to y

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$

- Compute derivative

$$-3 + e^{yx} x \cos(2x) = \frac{y e^{yx} \cos(2x)}{y^2+4} + y \left(\frac{e^{yx} \cos(2x)}{y^2+4} - \frac{2y^2 e^{yx} \cos(2x)}{(y^2+4)^2} + \frac{y e^{yx} x \cos(2x)}{y^2+4} - \frac{4 e^{yx} \sin(2x) y}{(y^2+4)^2} + \frac{2 e^{yx} x \sin(2x)}{y^2+4} \right)$$

- Isolate for $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = -3 + e^{yx} x \cos(2x) - \frac{y e^{yx} \cos(2x)}{y^2+4} - y \left(\frac{e^{yx} \cos(2x)}{y^2+4} - \frac{2y^2 e^{yx} \cos(2x)}{(y^2+4)^2} + \frac{y e^{yx} x \cos(2x)}{y^2+4} - \frac{4 e^{yx} \sin(2x)}{(y^2+4)^2} \right)$$

- Solve for $f_1(y)$

$$f_1(y) = -3y$$

- Substitute $f_1(y)$ into equation for $F(x, y)$

$$F(x, y) = y \left(\frac{y e^{yx} \cos(2x)}{y^2+4} + \frac{2 e^{yx} \sin(2x)}{y^2+4} \right) + x^2 + \frac{4 e^{yx} \cos(2x)}{y^2+4} - \frac{2y e^{yx} \sin(2x)}{y^2+4} - 3y$$

- Substitute $F(x, y)$ into the solution of the ODE

$$y \left(\frac{y e^{yx} \cos(2x)}{y^2+4} + \frac{2 e^{yx} \sin(2x)}{y^2+4} \right) + x^2 + \frac{4 e^{yx} \cos(2x)}{y^2+4} - \frac{2y e^{yx} \sin(2x)}{y^2+4} - 3y = c_1$$

- Solve for y

$$y = -\frac{-x^3 + c_1 x + 3 \operatorname{LambertW}\left(-\frac{x \cos(2x) e^{\frac{1}{3} x^3 - \frac{1}{3} c_1 x}}{3}\right)}{3x}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 35

```
dsolve(2*x-2*exp(x*y(x))*sin(2*x)+exp(x*y(x))*cos(2*x)*y(x)+(-3+exp(x*y(x))*x*cos(2*x))*diff
```

$$y(x) = \frac{x^3 + c_1 x - 3 \operatorname{LambertW}\left(-\frac{x \cos(2x) e^{\frac{x(x^2+c_1)}{3}}}{3}\right)}{3x}$$

✓ Solution by Mathematica

Time used: 5.197 (sec). Leaf size: 48

```
DSolve [2*x-2*Exp [x*y [x]] *Sin [2*x]+Exp [x*y [x]] *Cos [2*x] *y [x]+(-3+Exp [x*y [x]] *x *Cos [2*x]) *y' [x]
```

$$y(x) \rightarrow \frac{-3W\left(-\frac{1}{3}xe^{\frac{1}{3}x(x^2-c_1)} \cos(2x)\right) + x^3 - c_1x}{3x}$$

5.10 problem 10

5.10.1 Solving as linear ode	1237
5.10.2 Solving as first order ode lie symmetry lookup ode	1239
5.10.3 Solving as exact ode	1243
5.10.4 Maple step by step solution	1247

Internal problem ID [552]

Internal file name [OUTPUT/552_Sunday_June_05_2022_01_44_25_AM_13151512/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 10.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**exact**", "**linear**", "**first_order_ode_lie_symmetry_lookup**"

Maple gives the following as the ode type

`[_linear]`

$$\frac{y}{x} + (\ln(x) - 2)y' = -6x$$

5.10.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{1}{(\ln(x) - 2)x}$$
$$q(x) = -\frac{6x}{\ln(x) - 2}$$

Hence the ode is

$$y' + \frac{y}{(\ln(x) - 2)x} = -\frac{6x}{\ln(x) - 2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{1}{(\ln(x)-2)x} dx} \\ &= \ln(x) - 2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) \left(-\frac{6x}{\ln(x) - 2} \right) \\ \frac{d}{dx}((\ln(x) - 2)y) &= (\ln(x) - 2) \left(-\frac{6x}{\ln(x) - 2} \right) \\ d((\ln(x) - 2)y) &= (-6x) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}(\ln(x) - 2)y &= \int -6x dx \\ (\ln(x) - 2)y &= -3x^2 + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = \ln(x) - 2$ results in

$$y = -\frac{3x^2}{\ln(x) - 2} + \frac{c_1}{\ln(x) - 2}$$

which simplifies to

$$y = \frac{-3x^2 + c_1}{\ln(x) - 2}$$

Summary

The solution(s) found are the following

$$y = \frac{-3x^2 + c_1}{\ln(x) - 2} \tag{1}$$

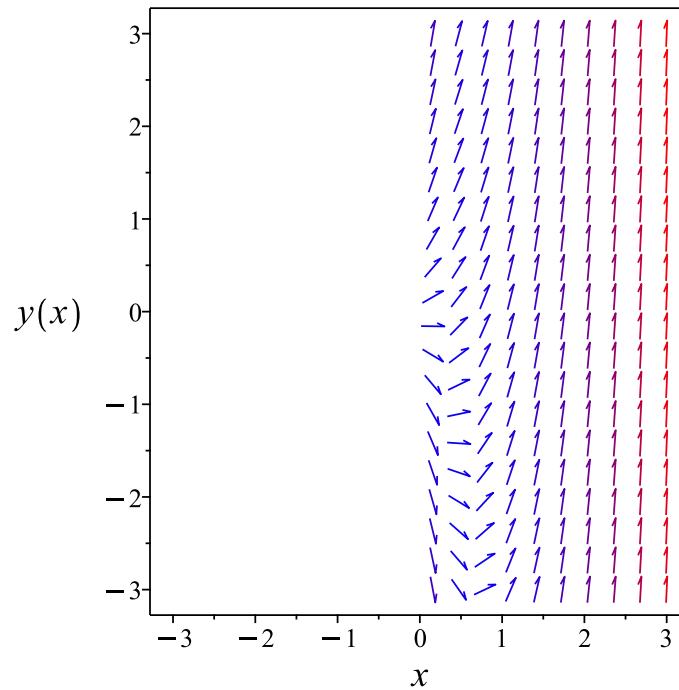


Figure 239: Slope field plot

Verification of solutions

$$y = \frac{-3x^2 + c_1}{\ln(x) - 2}$$

Verified OK.

5.10.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{6x^2 + y}{(\ln(x) - 2)x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 243: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned} \xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{\ln(x) - 2} \end{aligned} \tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{\ln(x)-2}} dy \end{aligned}$$

Which results in

$$S = (\ln(x) - 2) y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{6x^2 + y}{(\ln(x) - 2)x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{y}{x} \\ S_y &= \ln(x) - 2 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -6x \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -6R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -3R^2 + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$(\ln(x) - 2)y = -3x^2 + c_1$$

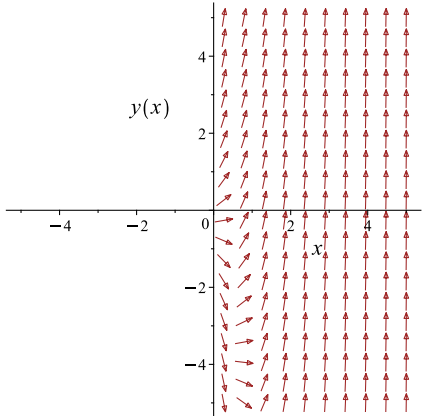
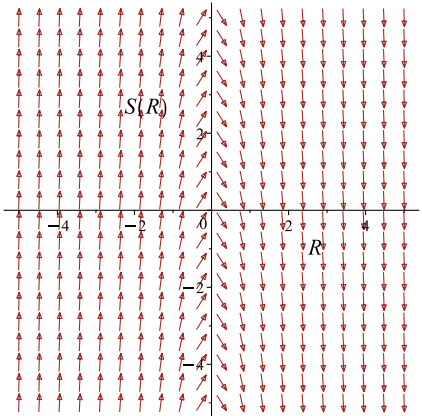
Which simplifies to

$$(\ln(x) - 2)y = -3x^2 + c_1$$

Which gives

$$y = \frac{-3x^2 + c_1}{\ln(x) - 2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{6x^2+y}{(\ln(x)-2)x}$ 	$R = x$ $S = (\ln(x) - 2)y$	$\frac{dS}{dR} = -6R$ 

Summary

The solution(s) found are the following

$$y = \frac{-3x^2 + c_1}{\ln(x) - 2} \quad (1)$$

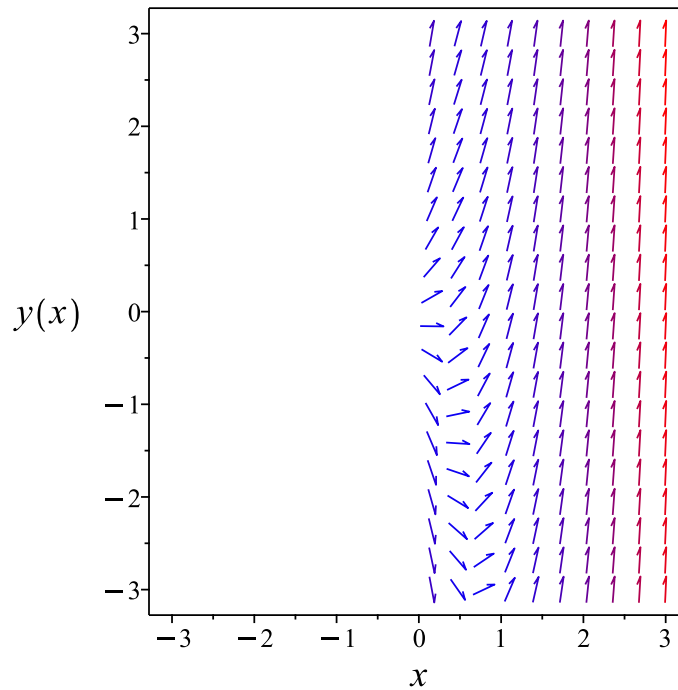


Figure 240: Slope field plot

Verification of solutions

$$y = \frac{-3x^2 + c_1}{\ln(x) - 2}$$

Verified OK.

5.10.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(\ln(x) - 2) dy &= \left(-\frac{y}{x} - 6x\right) dx \\ \left(\frac{y}{x} + 6x\right) dx + (\ln(x) - 2) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= \frac{y}{x} + 6x \\ N(x, y) &= \ln(x) - 2\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{y}{x} + 6x\right) \\ &= \frac{1}{x}\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(\ln(x) - 2) \\ &= \frac{1}{x}\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{y}{x} + 6x dx \\ \phi &= \ln(x)y + 3x^2 + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = \ln(x) + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \ln(x) - 2$. Therefore equation (4) becomes

$$\ln(x) - 2 = \ln(x) + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -2$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (-2) dy \\ f(y) &= -2y + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = \ln(x)y + 3x^2 - 2y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \ln(x)y + 3x^2 - 2y$$

The solution becomes

$$y = \frac{-3x^2 + c_1}{\ln(x) - 2}$$

Summary

The solution(s) found are the following

$$y = \frac{-3x^2 + c_1}{\ln(x) - 2} \tag{1}$$

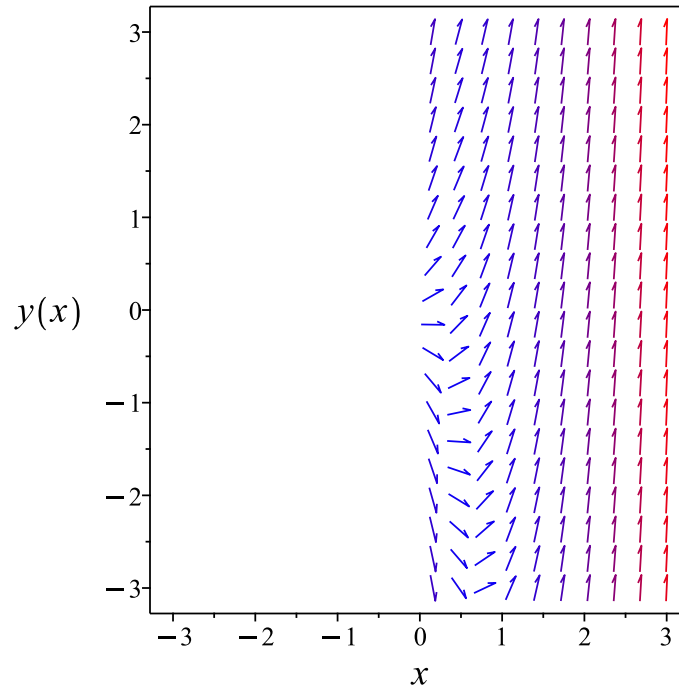


Figure 241: Slope field plot

Verification of solutions

$$y = \frac{-3x^2 + c_1}{\ln(x) - 2}$$

Verified OK.

5.10.4 Maple step by step solution

Let's solve

$$\frac{y}{x} + (\ln(x) - 2) y' = -6x$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{y}{(\ln(x)-2)x} - \frac{6x}{\ln(x)-2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{y}{(\ln(x)-2)x} = -\frac{6x}{\ln(x)-2}$$

- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$\mu(x) \left(y' + \frac{y}{(\ln(x)-2)x} \right) = -\frac{6\mu(x)x}{\ln(x)-2}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dx}(\mu(x)y)$

$$\mu(x) \left(y' + \frac{y}{(\ln(x)-2)x} \right) = \mu'(x)y + \mu(x)y'$$

- Isolate $\mu'(x)$

$$\mu'(x) = \frac{\mu(x)}{(\ln(x)-2)x}$$

- Solve to find the integrating factor

$$\mu(x) = \ln(x) - 2$$

- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x)y) \right) dx = \int -\frac{6\mu(x)x}{\ln(x)-2} dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int -\frac{6\mu(x)x}{\ln(x)-2} dx + c_1$$

- Solve for y

$$y = \frac{\int -\frac{6\mu(x)x}{\ln(x)-2} dx + c_1}{\mu(x)}$$

- Substitute $\mu(x) = \ln(x) - 2$

$$y = \frac{\int -6x dx + c_1}{\ln(x) - 2}$$

- Evaluate the integrals on the rhs

$$y = \frac{-3x^2 + c_1}{\ln(x) - 2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.219 (sec). Leaf size: 18

```
dsolve((y(x)/x+6*x)+(ln(x)-2)*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = \frac{-3x^2 + c_1}{\ln(x) - 2}$$

✓ Solution by Mathematica

Time used: 0.079 (sec). Leaf size: 20

```
DSolve[(y[x]/x+6*x)+(Log[x]-2)*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{-3x^2 + c_1}{\log(x) - 2}$$

5.11 problem 11

Internal problem ID [553]

Internal file name [OUTPUT/553_Sunday_June_05_2022_01_44_26_AM_51220061/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 11.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**unknown**"

Maple gives the following as the ode type

```
[[_Abel, `2nd type`, `class B`]]
```

Unable to solve or complete the solution.

$$yx + (\ln(x)y + yx)y' = -x \ln(x)$$

Unable to determine ODE type.

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
trying Abel
Looking for potential symmetries
Looking for potential symmetries
Looking for potential symmetries
trying inverse_Riccati
trying an equivalence to an Abel ODE
differential order: 1; trying a linearization to 2nd order
--- trying a change of variables {x -> y(x), y(x) -> x}
differential order: 1; trying a linearization to 2nd order
trying 1st order ODE linearizable_by_differentiation
--- Trying Lie symmetry methods, 1st order ---
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = 4
`, `-> Computing symmetries using: way = 2
trying symmetry patterns for 1st order ODEs
-> trying a symmetry pattern of the form [F(x)*G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)*G(y)]
-> trying symmetry patterns of the forms [F(x),G(y)] and [G(y),F(x)]
`, `-> Computing symmetries using: way = HINT
-> Calling odsolve with the ODE`, diff(y(x), x)+y(x)*(ln(x)-1)/(x*(ln(x)+x)), y(x)`
  Methods for first order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  trying 1st order linear
  <- 1st order linear successful
-> Calling odsolve with the ODE`, diff(y(x), x)+y(x)*(ln(x)^2+x)/(x*(ln(x)+x)*ln(x)), y(x)`
  Methods for first order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  trying 1st order linear
  <- 1st order linear successful 1250
`, `-> Computing symmetries using: way = HINT
-> trying a symmetry pattern of the form [F(x),G(x)]
-> trying a symmetry pattern of the form [F(x),G(x)]
```

X Solution by Maple

```
dsolve((x*ln(x)+x*y(x))+(y(x)*ln(x)+x*y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

No solution found

X Solution by Mathematica

Time used: 0.0 (sec). Leaf size: 0

```
DSolve[(x*Log[x]+x*y[x])+(y[x]*Log[x]+x*y[x])*y'[x] == 0,y[x],x,IncludeSingularSolutions ->
```

Not solved

5.12 problem 12

5.12.1 Solving as separable ode	1252
5.12.2 Solving as homogeneousTypeD2 ode	1254
5.12.3 Solving as differentialType ode	1256
5.12.4 Solving as first order ode lie symmetry lookup ode	1257
5.12.5 Solving as exact ode	1261
5.12.6 Maple step by step solution	1265

Internal problem ID [554]

Internal file name [OUTPUT/554_Sunday_June_05_2022_01_44_30_AM_73717968/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 12.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$\frac{x}{(x^2 + y^2)^{\frac{3}{2}}} + \frac{yy'}{(x^2 + y^2)^{\frac{3}{2}}} = 0$$

5.12.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\frac{x}{y}\end{aligned}$$

Where $f(x) = -x$ and $g(y) = \frac{1}{y}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y} dy &= -x dx \\ \int \frac{1}{y} dy &= \int -x dx \\ \frac{y^2}{2} &= -\frac{x^2}{2} + c_1\end{aligned}$$

Which results in

$$\begin{aligned}y &= \sqrt{-x^2 + 2c_1} \\ y &= -\sqrt{-x^2 + 2c_1}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \sqrt{-x^2 + 2c_1} \tag{1}$$

$$y = -\sqrt{-x^2 + 2c_1} \tag{2}$$

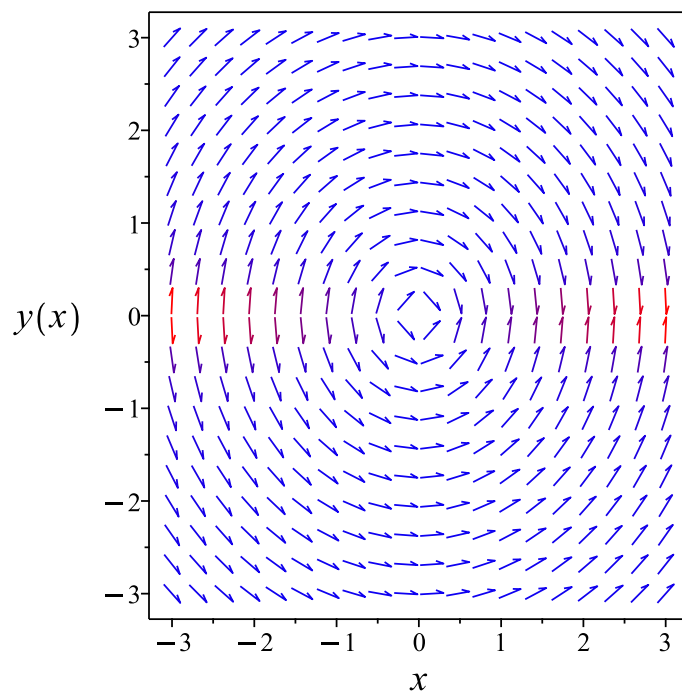


Figure 242: Slope field plot

Verification of solutions

$$y = \sqrt{-x^2 + 2c_1}$$

Verified OK.

$$y = -\sqrt{-x^2 + 2c_1}$$

Verified OK.

5.12.2 Solving as homogeneous Type D2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$\frac{x}{(x^2 + u(x)^2 x^2)^{\frac{3}{2}}} + \frac{u(x)x(u'(x)x + u(x))}{(x^2 + u(x)^2 x^2)^{\frac{3}{2}}} = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{u^2 + 1}{ux} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(u) = \frac{u^2+1}{u}$. Integrating both sides gives

$$\begin{aligned} \frac{1}{\frac{u^2+1}{u}} du &= -\frac{1}{x} dx \\ \int \frac{1}{\frac{u^2+1}{u}} du &= \int -\frac{1}{x} dx \\ \frac{\ln(u^2 + 1)}{2} &= -\ln(x) + c_2 \end{aligned}$$

Raising both side to exponential gives

$$\sqrt{u^2 + 1} = e^{-\ln(x)+c_2}$$

Which simplifies to

$$\sqrt{u^2 + 1} = \frac{c_3}{x}$$

Which simplifies to

$$\sqrt{u(x)^2 + 1} = \frac{c_3 e^{c_2}}{x}$$

The solution is

$$\sqrt{u(x)^2 + 1} = \frac{c_3 e^{c_2}}{x}$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\begin{aligned}\sqrt{\frac{y^2}{x^2} + 1} &= \frac{c_3 e^{c_2}}{x} \\ \sqrt{\frac{x^2 + y^2}{x^2}} &= \frac{c_3 e^{c_2}}{x}\end{aligned}$$

Summary

The solution(s) found are the following

$$\sqrt{\frac{x^2 + y^2}{x^2}} = \frac{c_3 e^{c_2}}{x} \quad (1)$$

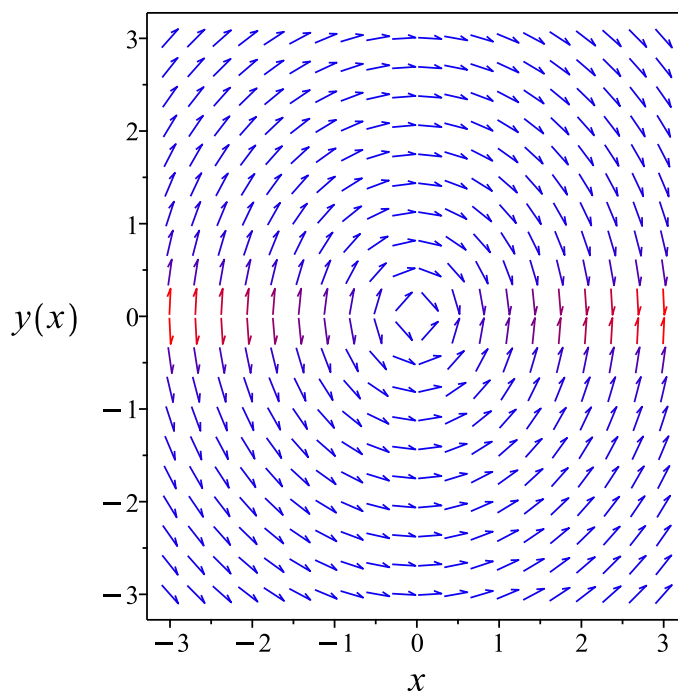


Figure 243: Slope field plot

Verification of solutions

$$\sqrt{\frac{x^2 + y^2}{x^2}} = \frac{c_3 e^{c_2}}{x}$$

Verified OK.

5.12.3 Solving as differential Type ode

Writing the ode as

$$y' = -\frac{x}{y} \quad (1)$$

Which becomes

$$(y) dy = (-x) dx \quad (2)$$

But the RHS is complete differential because

$$(-x) dx = d\left(-\frac{x^2}{2}\right)$$

Hence (2) becomes

$$(y) dy = d\left(-\frac{x^2}{2}\right)$$

Integrating both sides gives gives these solutions

$$y = \sqrt{-x^2 + 2c_1} + c_1$$

$$y = -\sqrt{-x^2 + 2c_1} + c_1$$

Summary

The solution(s) found are the following

$$y = \sqrt{-x^2 + 2c_1} + c_1 \quad (1)$$

$$y = -\sqrt{-x^2 + 2c_1} + c_1 \quad (2)$$

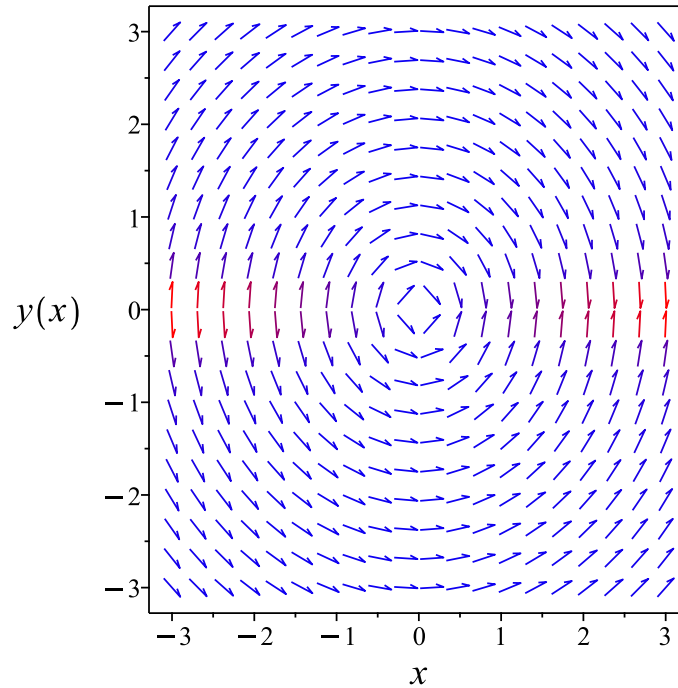


Figure 244: Slope field plot

Verification of solutions

$$y = \sqrt{-x^2 + 2c_1} + c_1$$

Verified OK.

$$y = -\sqrt{-x^2 + 2c_1} + c_1$$

Verified OK.

5.12.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{x}{y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 246: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= -\frac{1}{x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{-\frac{1}{x}} dx \end{aligned}$$

Which results in

$$S = -\frac{x^2}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{x}{y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= -x \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^2}{2} + c_1 \quad (4)$$

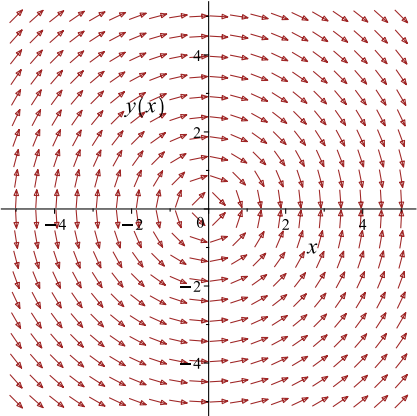
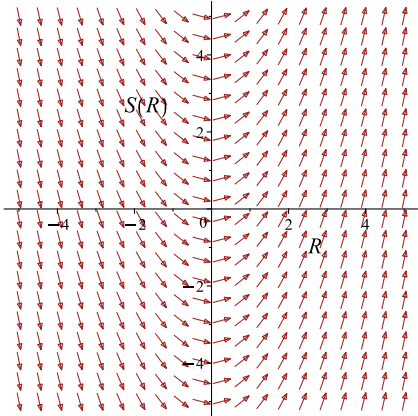
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-\frac{x^2}{2} = \frac{y^2}{2} + c_1$$

Which simplifies to

$$-\frac{x^2}{2} = \frac{y^2}{2} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{x}{y}$ 	$R = y$ $S = -\frac{x^2}{2}$	$\frac{dS}{dR} = R$ 

Summary

The solution(s) found are the following

$$-\frac{x^2}{2} = \frac{y^2}{2} + c_1 \quad (1)$$

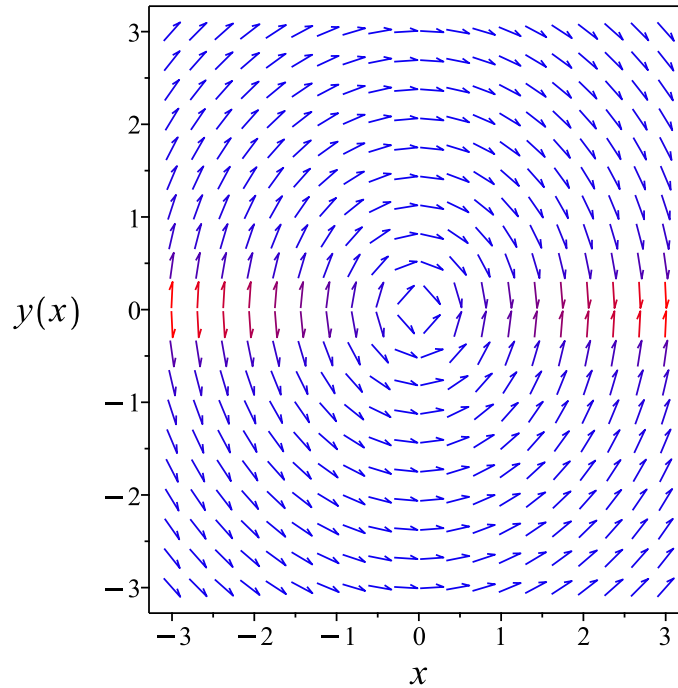


Figure 245: Slope field plot

Verification of solutions

$$-\frac{x^2}{2} = \frac{y^2}{2} + c_1$$

Verified OK.

5.12.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(-y) dy &= (x) dx \\ (-x) dx + (-y) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -x \\ N(x, y) &= -y\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-y) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -x dx$$

$$\phi = -\frac{x^2}{2} + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -y$. Therefore equation (4) becomes

$$-y = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -y$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (-y) dy$$

$$f(y) = -\frac{y^2}{2} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{x^2}{2} - \frac{y^2}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{x^2}{2} - \frac{y^2}{2}$$

Summary

The solution(s) found are the following

$$-\frac{x^2}{2} - \frac{y^2}{2} = c_1 \tag{1}$$

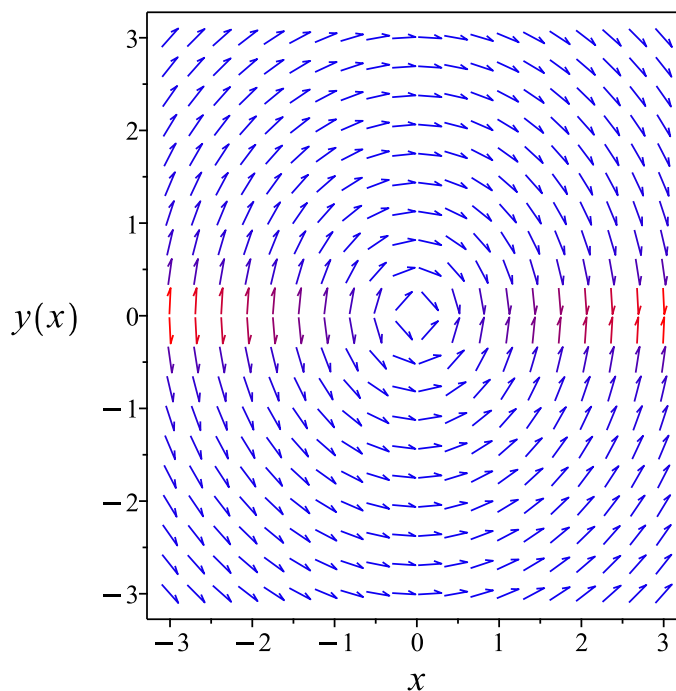


Figure 246: Slope field plot

Verification of solutions

$$-\frac{x^2}{2} - \frac{y^2}{2} = c_1$$

Verified OK.

5.12.6 Maple step by step solution

Let's solve

$$\frac{x}{(x^2+y^2)^{\frac{3}{2}}} + \frac{yy'}{(x^2+y^2)^{\frac{3}{2}}} = 0$$

- Highest derivative means the order of the ODE is 1
 y'

- Integrate both sides with respect to x

$$\int \left(\frac{x}{(x^2+y^2)^{\frac{3}{2}}} + \frac{yy'}{(x^2+y^2)^{\frac{3}{2}}} \right) dx = \int 0 dx + c_1$$

- Evaluate integral

$$-\frac{1}{\sqrt{x^2+y^2}} = c_1$$

- Solve for y

$$\left\{ y = \frac{\sqrt{-c_1^2 x^2 + 1}}{c_1}, y = -\frac{\sqrt{-c_1^2 x^2 + 1}}{c_1} \right\}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 27

```
dsolve(x/(x^2+y(x)^2)^(3/2)+y(x)*diff(y(x),x)/(x^2+y(x)^2)^(3/2) = 0,y(x), singsol=all)
```

$$y(x) = \sqrt{-x^2 + c_1}$$
$$y(x) = -\sqrt{-x^2 + c_1}$$

✓ Solution by Mathematica

Time used: 0.09 (sec). Leaf size: 39

```
DSolve[x/(x^2+y[x]^2)^(3/2)+y[x]*y'[x]/(x^2+y[x]^2)^(3/2) == 0,y[x],x,IncludeSingularSolutio
```

$$y(x) \rightarrow -\sqrt{-x^2 + 2c_1}$$

$$y(x) \rightarrow \sqrt{-x^2 + 2c_1}$$

5.13 problem 13

5.13.1 Existence and uniqueness analysis	1267
5.13.2 Solving as homogeneousTypeD2 ode	1268
5.13.3 Solving as differentialType ode	1270
5.13.4 Solving as first order ode lie symmetry calculated ode	1271
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5.13.6 Maple step by step solution	1279

Internal problem ID [555]

Internal file name [OUTPUT/555_Sunday_June_05_2022_01_44_31_AM_94833341/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 13.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "differentialType", "homogeneousTypeD2", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _exact, _rational, [_Abel, `2nd
  type`, `class A`]]
```

$$-y + (-x + 2y)y' = -2x$$

With initial conditions

$$[y(1) = 3]$$

5.13.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= \frac{-2x + y}{-x + 2y}\end{aligned}$$

The x domain of $f(x, y)$ when $y = 3$ is

$$\{x < 6 \vee 6 < x\}$$

And the point $x_0 = 1$ is inside this domain. The y domain of $f(x, y)$ when $x = 1$ is

$$\left\{y < \frac{1}{2} \vee \frac{1}{2} < y\right\}$$

And the point $y_0 = 3$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned} \frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{-2x + y}{-x + 2y} \right) \\ &= \frac{1}{-x + 2y} - \frac{2(-2x + y)}{(-x + 2y)^2} \end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 3$ is

$$\{x < 6 \vee 6 < x\}$$

And the point $x_0 = 1$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 1$ is

$$\left\{y < \frac{1}{2} \vee \frac{1}{2} < y\right\}$$

And the point $y_0 = 3$ is inside this domain. Therefore solution exists and is unique.

5.13.2 Solving as homogeneous TypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$-u(x)x + (-x + 2u(x)x)(u'(x)x + u(x)) = -2x$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{2(u^2 - u + 1)}{x(2u - 1)} \end{aligned}$$

Where $f(x) = -\frac{2}{x}$ and $g(u) = \frac{u^2-u+1}{2u-1}$. Integrating both sides gives

$$\frac{1}{\frac{u^2-u+1}{2u-1}} du = -\frac{2}{x} dx$$
$$\int \frac{1}{\frac{u^2-u+1}{2u-1}} du = \int -\frac{2}{x} dx$$
$$\ln(u^2 - u + 1) = -2 \ln(x) + c_2$$

Raising both side to exponential gives

$$u^2 - u + 1 = e^{-2 \ln(x) + c_2}$$

Which simplifies to

$$u^2 - u + 1 = \frac{c_3}{x^2}$$

Which simplifies to

$$u(x)^2 - u(x) + 1 = \frac{c_3 e^{c_2}}{x^2}$$

The solution is

$$u(x)^2 - u(x) + 1 = \frac{c_3 e^{c_2}}{x^2}$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\frac{y^2}{x^2} - \frac{y}{x} + 1 = \frac{c_3 e^{c_2}}{x^2}$$
$$\frac{y^2}{x^2} - \frac{y}{x} + 1 = \frac{c_3 e^{c_2}}{x^2}$$

Which simplifies to

$$y^2 - yx + x^2 = c_3 e^{c_2}$$

Substituting initial conditions and solving for c_2 gives $c_2 = \ln\left(\frac{7}{c_3}\right)$. Hence the solution

Summary
becomes The solution(s) found are the following

$$y^2 - yx + x^2 = 7 \tag{1}$$

Verification of solutions

$$y^2 - yx + x^2 = 7$$

Verified OK.

5.13.3 Solving as differential Type ode

Writing the ode as

$$y' = \frac{-2x + y}{-x + 2y} \quad (1)$$

Which becomes

$$(-2y) dy = (-x) dy + (2x - y) dx \quad (2)$$

But the RHS is complete differential because

$$(-x) dy + (2x - y) dx = d(x^2 - yx)$$

Hence (2) becomes

$$(-2y) dy = d(x^2 - yx)$$

Integrating both sides gives gives these solutions

$$y = \frac{x}{2} + \frac{\sqrt{-3x^2 - 4c_1}}{2} + c_1$$
$$y = \frac{x}{2} - \frac{\sqrt{-3x^2 - 4c_1}}{2} + c_1$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 3$ in the above solution gives an equation to solve for the constant of integration.

$$3 = \frac{1}{2} - \frac{\sqrt{-3 - 4c_1}}{2} + c_1$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 3$ in the above solution gives an equation to solve for the constant of integration.

$$3 = \frac{1}{2} + \frac{\sqrt{-3 - 4c_1}}{2} + c_1$$

$$c_1 = 2 - i\sqrt{3}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{x}{2} + \frac{\sqrt{-3x^2 - 8 + 4i\sqrt{3}}}{2} + 2 - i\sqrt{3}$$

Summary

The solution(s) found are the following

$$y = \frac{x}{2} + \frac{\sqrt{-3x^2 - 8 + 4i\sqrt{3}}}{2} + 2 - i\sqrt{3} \quad (1)$$

Verification of solutions

$$y = \frac{x}{2} + \frac{\sqrt{-3x^2 - 8 + 4i\sqrt{3}}}{2} + 2 - i\sqrt{3}$$

Verified OK.

5.13.4 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{-2x + y}{-x + 2y}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (A)$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1E)$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2E)$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + \frac{(-2x + y)(b_3 - a_2)}{-x + 2y} - \frac{(-2x + y)^2 a_3}{(-x + 2y)^2}$$
$$- \left(-\frac{2}{-x + 2y} + \frac{-2x + y}{(-x + 2y)^2} \right) (xa_2 + ya_3 + a_1) \quad (5E)$$
$$- \left(\frac{1}{-x + 2y} - \frac{2(-2x + y)}{(-x + 2y)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{2x^2a_2 + 4x^2a_3 + 2x^2b_2 - 2x^2b_3 - 8xya_2 - 4xya_3 + 4xyb_2 + 8xyb_3 + 2y^2a_2 - 2y^2a_3 - 4y^2b_2 - 2y^2b_3 + 3a_1}{(x - 2y)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} -2x^2a_2 - 4x^2a_3 - 2x^2b_2 + 2x^2b_3 + 8xya_2 + 4xya_3 - 4xyb_2 \\ - 8xyb_3 - 2y^2a_2 + 2y^2a_3 + 4y^2b_2 + 2y^2b_3 - 3a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} -2a_2v_1^2 + 8a_2v_1v_2 - 2a_2v_2^2 - 4a_3v_1^2 + 4a_3v_1v_2 + 2a_3v_2^2 - 2b_2v_1^2 \\ - 4b_2v_1v_2 + 4b_2v_2^2 + 2b_3v_1^2 - 8b_3v_1v_2 + 2b_3v_2^2 + 3a_1v_2 - 3b_1v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} (-2a_2 - 4a_3 - 2b_2 + 2b_3)v_1^2 + (8a_2 + 4a_3 - 4b_2 - 8b_3)v_1v_2 \\ - 3b_1v_1 + (-2a_2 + 2a_3 + 4b_2 + 2b_3)v_2^2 + 3a_1v_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 3a_1 &= 0 \\ -3b_1 &= 0 \\ -2a_2 - 4a_3 - 2b_2 + 2b_3 &= 0 \\ -2a_2 + 2a_3 + 4b_2 + 2b_3 &= 0 \\ 8a_2 + 4a_3 - 4b_2 - 8b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_2 + b_3 \\ a_3 &= -b_2 \\ b_1 &= 0 \\ b_2 &= b_2 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= y - \left(\frac{-2x + y}{-x + 2y} \right) (x) \\ &= \frac{-2x^2 + 2yx - 2y^2}{x - 2y} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-2x^2 + 2yx - 2y^2}{x - 2y}} dy \end{aligned}$$

Which results in

$$S = \frac{\ln(x^2 - yx + y^2)}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{-2x + y}{-x + 2y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{2x - y}{2x^2 - 2yx + 2y^2} \\ S_y &= \frac{-x + 2y}{2x^2 - 2yx + 2y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \quad (4)$$

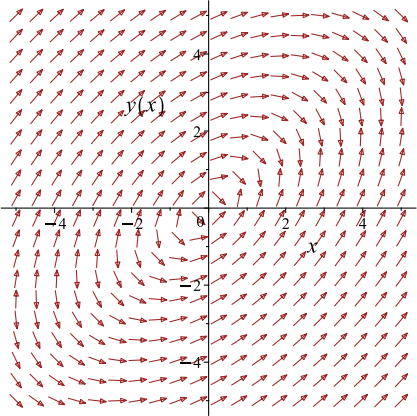
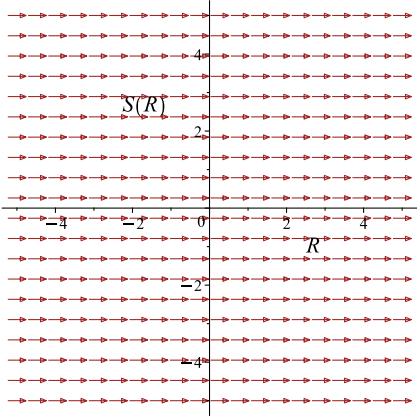
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\ln(y^2 - yx + x^2)}{2} = c_1$$

Which simplifies to

$$\frac{\ln(y^2 - yx + x^2)}{2} = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{-2x+y}{-x+2y}$ 	$R = x$ $S = \frac{\ln(x^2 - yx + y^2)}{2}$	$\frac{dS}{dR} = 0$ 

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 3$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{\ln(7)}{2} = c_1$$

$$c_1 = \frac{\ln(7)}{2}$$

Substituting c_1 found above in the general solution gives

$$\frac{\ln(x^2 - yx + y^2)}{2} = \frac{\ln(7)}{2}$$

Summary

The solution(s) found are the following

$$\frac{\ln(y^2 - yx + x^2)}{2} = \frac{\ln(7)}{2} \tag{1}$$

Verification of solutions

$$\frac{\ln(y^2 - yx + x^2)}{2} = \frac{\ln(7)}{2}$$

Verified OK.

5.13.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (-x + 2y) dy &= (-2x + y) dx \\ (2x - y) dx + (-x + 2y) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 2x - y \\N(x, y) &= -x + 2y\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2x - y) \\ &= -1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x + 2y) \\ &= -1\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 2x - y dx \\ \phi &= x(x - y) + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = -x + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -x + 2y$. Therefore equation (4) becomes

$$-x + 2y = -x + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 2y$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (2y) dy$$
$$f(y) = y^2 + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = x(x - y) + y^2 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = x(x - y) + y^2$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 3$ in the above solution gives an equation to solve for the constant of integration.

$$7 = c_1$$

$$c_1 = 7$$

Substituting c_1 found above in the general solution gives

$$x(x - y) + y^2 = 7$$

Summary

The solution(s) found are the following

$$y^2 - yx + x^2 = 7 \quad (1)$$

Verification of solutions

$$y^2 - yx + x^2 = 7$$

Verified OK.

The solution

$$\frac{\ln(y^2 - yx + x^2)}{2} = \frac{\ln(7)}{2}$$

can be simplified to

$$\ln(y^2 - yx + x^2) = \ln(7)$$

5.13.6 Maple step by step solution

Let's solve

$$[-y + (-x + 2y)y' = -2x, y(1) = 3]$$

- Highest derivative means the order of the ODE is 1
 y'
- Check if ODE is exact
 - ODE is exact if the lhs is the total derivative of a C^2 function
 $F'(x, y) = 0$
 - Compute derivative of lhs
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$
 - Evaluate derivatives
 $-1 = -1$
 - Condition met, ODE is exact
- Exact ODE implies solution will be of this form
 $\left[F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y)\right]$
- Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to x
 $F(x, y) = \int (2x - y) dx + f_1(y)$
- Evaluate integral
 $F(x, y) = x^2 - yx + f_1(y)$
- Take derivative of $F(x, y)$ with respect to y
 $N(x, y) = \frac{\partial}{\partial y} F(x, y)$

- Compute derivative

$$-x + 2y = -x + \frac{d}{dy}f_1(y)$$
- Isolate for $\frac{d}{dy}f_1(y)$

$$\frac{d}{dy}f_1(y) = 2y$$
- Solve for $f_1(y)$

$$f_1(y) = y^2$$
- Substitute $f_1(y)$ into equation for $F(x, y)$

$$F(x, y) = x^2 - yx + y^2$$
- Substitute $F(x, y)$ into the solution of the ODE

$$x^2 - yx + y^2 = c_1$$
- Solve for y

$$\left\{ y = \frac{x}{2} - \frac{\sqrt{-3x^2+4c_1}}{2}, y = \frac{x}{2} + \frac{\sqrt{-3x^2+4c_1}}{2} \right\}$$
- Use initial condition $y(1) = 3$

$$3 = \frac{1}{2} - \frac{\sqrt{4c_1-3}}{2}$$
- Solution does not satisfy initial condition
- Use initial condition $y(1) = 3$

$$3 = \frac{1}{2} + \frac{\sqrt{4c_1-3}}{2}$$
- Solve for c_1

$$c_1 = 7$$
- Substitute $c_1 = 7$ into general solution and simplify

$$y = \frac{x}{2} + \frac{\sqrt{-3x^2+28}}{2}$$
- Solution to the IVP

$$y = \frac{x}{2} + \frac{\sqrt{-3x^2+28}}{2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 0.094 (sec). Leaf size: 19

```
dsolve([2*x-y(x)+(-x+2*y(x))*diff(y(x),x) = 0,y(1) = 3],y(x), singsol=all)
```

$$y(x) = \frac{x}{2} + \frac{\sqrt{-3x^2 + 28}}{2}$$

✓ Solution by Mathematica

Time used: 0.456 (sec). Leaf size: 22

```
DSolve[{2*x-y[x]+(-x+2*y[x])*y'[x] == 0,y[1]==3},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2} \left(\sqrt{28 - 3x^2} + x \right)$$

5.14 problem 14

5.14.1 Existence and uniqueness analysis	1282
5.14.2 Solving as differentialType ode	1283
5.14.3 Solving as exact ode	1285
5.14.4 Maple step by step solution	1288

Internal problem ID [556]

Internal file name [OUTPUT/556_Sunday_June_05_2022_01_44_32_AM_62551946/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 14.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "differentialType"

Maple gives the following as the ode type

```
[_exact, _rational, [_1st_order, `_with_symmetry_[F(x),G(x)]`],  
[_Abel, `2nd type`, `class A`]]
```

$$y + (x - 4y)y' = -9x^2 + 1$$

With initial conditions

$$[y(1) = 0]$$

5.14.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= \frac{9x^2 + y - 1}{-x + 4y}\end{aligned}$$

The x domain of $f(x, y)$ when $y = 0$ is

$$\{x < 0 \vee 0 < x\}$$

And the point $x_0 = 1$ is inside this domain. The y domain of $f(x, y)$ when $x = 1$ is

$$\left\{ y < \frac{1}{4} \vee \frac{1}{4} < y \right\}$$

And the point $y_0 = 0$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned} \frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{9x^2 + y - 1}{-x + 4y} \right) \\ &= \frac{1}{-x + 4y} - \frac{4(9x^2 + y - 1)}{(-x + 4y)^2} \end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 0$ is

$$\{x < 0 \vee 0 < x\}$$

And the point $x_0 = 1$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 1$ is

$$\left\{ y < \frac{1}{4} \vee \frac{1}{4} < y \right\}$$

And the point $y_0 = 0$ is inside this domain. Therefore solution exists and is unique.

5.14.2 Solving as differentialType ode

Writing the ode as

$$y' = \frac{1 - 9x^2 - y}{x - 4y} \quad (1)$$

Which becomes

$$(-4y) dy = (-x) dy + (-9x^2 - y + 1) dx \quad (2)$$

But the RHS is complete differential because

$$(-x) dy + (-9x^2 - y + 1) dx = d(-3x^3 - yx + x)$$

Hence (2) becomes

$$(-4y) dy = d(-3x^3 - yx + x)$$

Integrating both sides gives gives these solutions

$$y = \frac{x}{4} + \frac{\sqrt{24x^3 + x^2 - 8c_1 - 8x}}{4} + c_1$$

$$y = \frac{x}{4} - \frac{\sqrt{24x^3 + x^2 - 8c_1 - 8x}}{4} + c_1$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{1}{4} - \frac{\sqrt{17 - 8c_1}}{4} + c_1$$

$$c_1 = \frac{\sqrt{5}}{2} - \frac{1}{2}$$

Substituting c_1 found above in the general solution gives

$$y = \frac{x}{4} - \frac{\sqrt{24x^3 + x^2 - 4\sqrt{5} + 4 - 8x}}{4} + \frac{\sqrt{5}}{2} - \frac{1}{2}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{1}{4} + \frac{\sqrt{17 - 8c_1}}{4} + c_1$$

$$c_1 = -\frac{1}{2} - \frac{\sqrt{5}}{2}$$

Substituting c_1 found above in the general solution gives

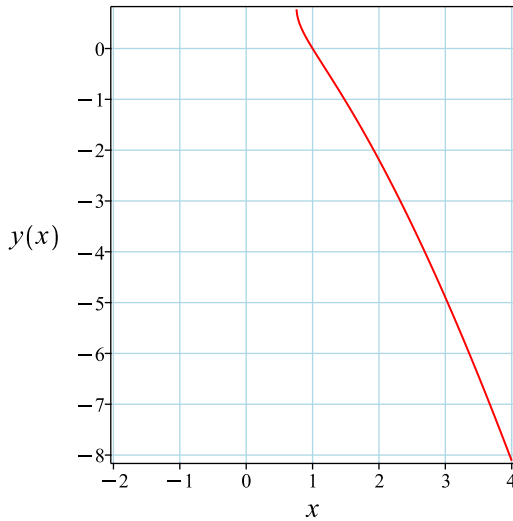
$$y = \frac{x}{4} + \frac{\sqrt{24x^3 + x^2 + 4 + 4\sqrt{5} - 8x}}{4} - \frac{1}{2} - \frac{\sqrt{5}}{2}$$

Summary

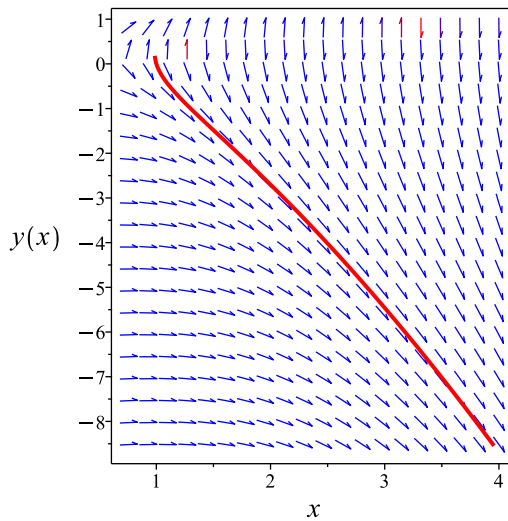
The solution(s) found are the following

$$y = \frac{x}{4} + \frac{\sqrt{24x^3 + x^2 + 4 + 4\sqrt{5} - 8x}}{4} - \frac{1}{2} - \frac{\sqrt{5}}{2} \quad (1)$$

$$y = \frac{x}{4} - \frac{\sqrt{24x^3 + x^2 - 4\sqrt{5} + 4 - 8x}}{4} + \frac{\sqrt{5}}{2} - \frac{1}{2} \quad (2)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{x}{4} + \frac{\sqrt{24x^3 + x^2 + 4 + 4\sqrt{5} - 8x}}{4} - \frac{1}{2} - \frac{\sqrt{5}}{2}$$

Verified OK.

$$y = \frac{x}{4} - \frac{\sqrt{24x^3 + x^2 - 4\sqrt{5} + 4 - 8x}}{4} + \frac{\sqrt{5}}{2} - \frac{1}{2}$$

Verified OK.

5.14.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(x - 4y) dy &= (-9x^2 - y + 1) dx \\ (9x^2 + y - 1) dx + (x - 4y) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 9x^2 + y - 1 \\ N(x, y) &= x - 4y\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(9x^2 + y - 1) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x - 4y) \\ &= 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 9x^2 + y - 1 dx \\ \phi &= x(3x^2 + y - 1) + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = x - 4y$. Therefore equation (4) becomes

$$x - 4y = x + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -4y$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int (-4y) dy \\ f(y) &= -2y^2 + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = x(3x^2 + y - 1) - 2y^2 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = x(3x^2 + y - 1) - 2y^2$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$2 = c_1$$

$$c_1 = 2$$

Substituting c_1 found above in the general solution gives

$$x(3x^2 + y - 1) - 2y^2 = 2$$

Summary

The solution(s) found are the following

$$3x^3 + (y - 1)x - 2y^2 = 2 \tag{1}$$

Verification of solutions

$$3x^3 + (y - 1)x - 2y^2 = 2$$

Verified OK.

5.14.4 Maple step by step solution

Let's solve

$$[y + (x - 4y)y' = -9x^2 + 1, y(1) = 0]$$

- Highest derivative means the order of the ODE is 1
 y'
- Check if ODE is exact
 - ODE is exact if the lhs is the total derivative of a C^2 function
 $F'(x, y) = 0$
 - Compute derivative of lhs
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$

- Evaluate derivatives
 $1 = 1$
- Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$\left[F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$
- Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to x

$$F(x, y) = \int (9x^2 + y - 1) dx + f_1(y)$$
- Evaluate integral

$$F(x, y) = 3x^3 + yx - x + f_1(y)$$
- Take derivative of $F(x, y)$ with respect to y

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$
- Compute derivative

$$x - 4y = x + \frac{d}{dy} f_1(y)$$
- Isolate for $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = -4y$$
- Solve for $f_1(y)$

$$f_1(y) = -2y^2$$
- Substitute $f_1(y)$ into equation for $F(x, y)$

$$F(x, y) = 3x^3 + yx - 2y^2 - x$$
- Substitute $F(x, y)$ into the solution of the ODE

$$3x^3 + yx - 2y^2 - x = c_1$$
- Solve for y

$$\left\{ y = \frac{x}{4} - \frac{\sqrt{24x^3 + x^2 - 8c_1 - 8x}}{4}, y = \frac{x}{4} + \frac{\sqrt{24x^3 + x^2 - 8c_1 - 8x}}{4} \right\}$$
- Use initial condition $y(1) = 0$

$$0 = \frac{1}{4} - \frac{\sqrt{17 - 8c_1}}{4}$$
- Solve for c_1

$$c_1 = 2$$
- Substitute $c_1 = 2$ into general solution and simplify

$$y = \frac{x}{4} - \frac{\sqrt{24x^3 + x^2 - 8x - 16}}{4}$$

- Use initial condition $y(1) = 0$

$$0 = \frac{1}{4} + \frac{\sqrt{17 - 8c_1}}{4}$$

- Solution does not satisfy initial condition
- Solution to the IVP

$$y = \frac{x}{4} - \frac{\sqrt{24x^3 + x^2 - 8x - 16}}{4}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 25

```
dsolve([-1+9*x^2+y(x)+(x-4*y(x))*diff(y(x),x) = 0,y(1) = 0],y(x), singsol=all)
```

$$y(x) = \frac{x}{4} - \frac{\sqrt{24x^3 + x^2 - 8x - 16}}{4}$$

✓ Solution by Mathematica

Time used: 0.148 (sec). Leaf size: 34

```
DSolve[{-1+9*x^2+y[x]+(x-4*y[x])*y'[x] == 0,y[1]==0},y[x],x,IncludeSingularSolutions -> True
```

$$y(x) \rightarrow \frac{1}{4} \left(x + i\sqrt{-24x^3 - x^2 + 8x + 16} \right)$$

5.15 problem 19

5.15.1 Solving as separable ode	1291
5.15.2 Solving as first order ode lie symmetry lookup ode	1293
5.15.3 Solving as exact ode	1297
5.15.4 Maple step by step solution	1301

Internal problem ID [557]

Internal file name [OUTPUT/557_Sunday_June_05_2022_01_44_33_AM_13501171/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 19.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y^3 x^2 + x(1 + y^2) y' = 0$$

5.15.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\frac{y^3 x}{y^2 + 1} \end{aligned}$$

Where $f(x) = -x$ and $g(y) = \frac{y^3}{y^2+1}$. Integrating both sides gives

$$\begin{aligned} \frac{1}{\frac{y^3}{y^2+1}} dy &= -x dx \\ \int \frac{1}{\frac{y^3}{y^2+1}} dy &= \int -x dx \end{aligned}$$

$$\ln(y) - \frac{1}{2y^2} = -\frac{x^2}{2} + c_1$$

Which results in

$$y = e^{\frac{\text{LambertW}(e^{x^2-2c_1})}{2} - \frac{x^2}{2} + c_1}$$

Summary

The solution(s) found are the following

$$y = e^{\frac{\text{LambertW}(e^{x^2-2c_1})}{2} - \frac{x^2}{2} + c_1} \quad (1)$$

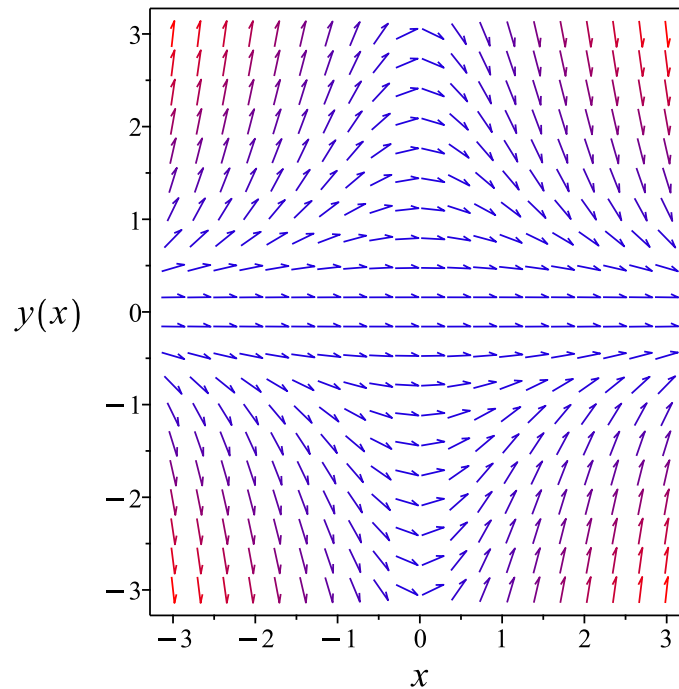


Figure 248: Slope field plot

Verification of solutions

$$y = e^{\frac{\text{LambertW}(e^{x^2-2c_1})}{2} - \frac{x^2}{2} + c_1}$$

Verified OK.

5.15.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{y^3 x}{y^2 + 1}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 251: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= -\frac{1}{x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{-\frac{1}{x}} dx\end{aligned}$$

Which results in

$$S = -\frac{x^2}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y^3 x}{y^2 + 1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= -x \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{y^2 + 1}{y^3} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{R^2 + 1}{R^3}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \ln(R) - \frac{1}{2R^2} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-\frac{x^2}{2} = \ln(y) - \frac{1}{2y^2} + c_1$$

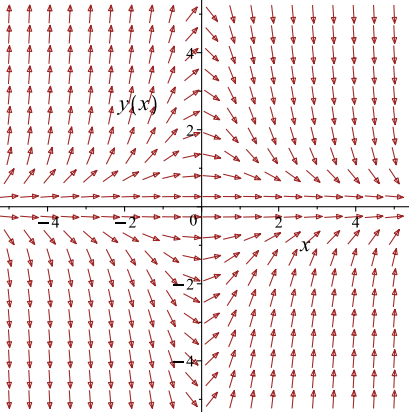
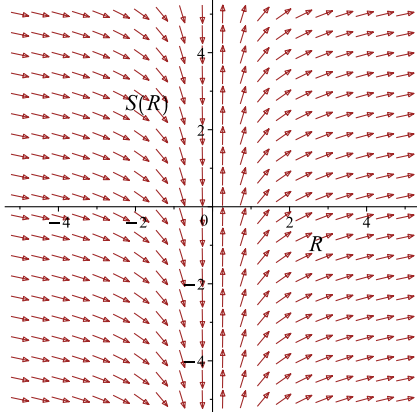
Which simplifies to

$$-\frac{x^2}{2} = \ln(y) - \frac{1}{2y^2} + c_1$$

Which gives

$$y = e^{\frac{\text{LambertW}(e^{x^2+2c_1})}{2} - \frac{x^2}{2} - c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{y^3 x}{y^2 + 1}$ 	$R = y$ $S = -\frac{x^2}{2}$	$\frac{dS}{dR} = \frac{R^2 + 1}{R^3}$ 

Summary

The solution(s) found are the following

$$y = e^{\frac{\text{LambertW}(e^{x^2+2c_1})}{2} - \frac{x^2}{2} - c_1} \tag{1}$$

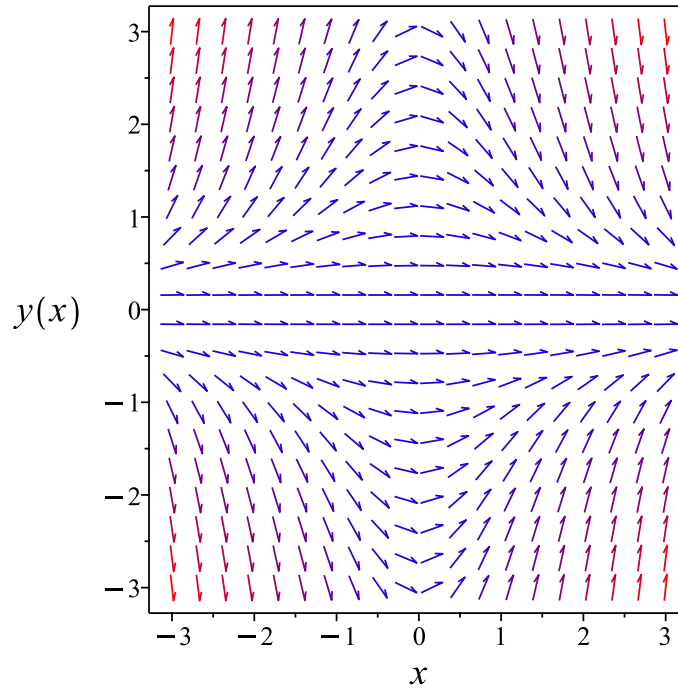


Figure 249: Slope field plot

Verification of solutions

$$y = e^{\frac{\text{LambertW}(e^{x^2+2c_1})}{2} - \frac{x^2}{2} - c_1}$$

Verified OK.

5.15.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(-\frac{y^2 + 1}{y^3}\right) dy &= (x) dx \\ (-x) dx + \left(-\frac{y^2 + 1}{y^3}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -x \\ N(x, y) &= -\frac{y^2 + 1}{y^3}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(-\frac{y^2 + 1}{y^3} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x dx \\ \phi &= -\frac{x^2}{2} + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{y^2+1}{y^3}$. Therefore equation (4) becomes

$$-\frac{y^2 + 1}{y^3} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{y^2 + 1}{y^3}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(\frac{-y^2 - 1}{y^3} \right) dy$$
$$f(y) = -\ln(y) + \frac{1}{2y^2} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{x^2}{2} - \ln(y) + \frac{1}{2y^2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{x^2}{2} - \ln(y) + \frac{1}{2y^2}$$

The solution becomes

$$y = e^{\frac{\text{LambertW}(e^{x^2+2c_1})}{2} - \frac{x^2}{2} - c_1}$$

Summary

The solution(s) found are the following

$$y = e^{\frac{\text{LambertW}(e^{x^2+2c_1})}{2} - \frac{x^2}{2} - c_1} \quad (1)$$

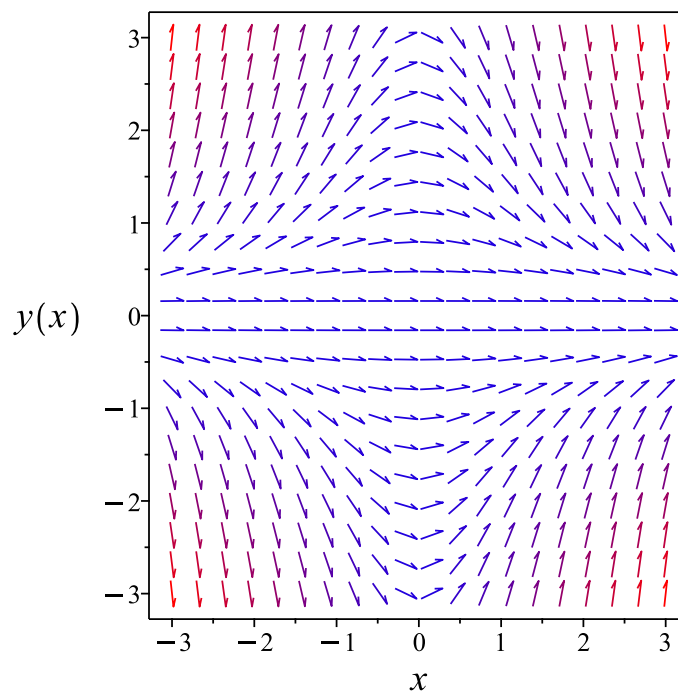


Figure 250: Slope field plot

Verification of solutions

$$y = e^{\frac{\text{LambertW}(e^{x^2+2c_1})}{2} - \frac{x^2}{2} - c_1}$$

Verified OK.

5.15.4 Maple step by step solution

Let's solve

$$y^3 x^2 + x(1 + y^2) y' = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y'(1+y^2)}{y^3} = -x$$

- Integrate both sides with respect to x

$$\int \frac{y'(1+y^2)}{y^3} dx = \int -x dx + c_1$$

- Evaluate integral

$$\ln(y) - \frac{1}{2y^2} = -\frac{x^2}{2} + c_1$$

- Solve for y

$$y = e^{\frac{\text{LambertW}(e^{x^2-2c_1})}{2} - \frac{x^2}{2} + c_1}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 37

```
dsolve(x^2*y(x)^3+x*(1+y(x)^2)*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-\frac{x^2}{2} - c_1} \sqrt{\frac{e^{x^2+2c_1}}{\text{LambertW}(e^{x^2+2c_1})}}$$

✓ Solution by Mathematica

Time used: 4.095 (sec). Leaf size: 46

```
DSolve[x^2*y[x]^3+x*(1+y[x]^2)*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{1}{\sqrt{W(e^{x^2-2c_1})}}$$

$$y(x) \rightarrow \frac{1}{\sqrt{W(e^{x^2-2c_1})}}$$

$$y(x) \rightarrow 0$$

5.16 problem 21

5.16.1 Solving as exact ode 1303

Internal problem ID [558]

Internal file name [OUTPUT/558_Sunday_June_05_2022_01_44_35_AM_40301990/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 21.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**exactWithIntegrationFactor**"

Maple gives the following as the ode type

```
[[_1st_order , `_with_symmetry_[F(x)*G(y),0]`]]
```

$$y + (2x - e^y y) y' = 0$$

5.16.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (2x - e^y y) dy &= (-y) dx \\ (y) dx + (2x - e^y y) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= y \\ N(x, y) &= 2x - e^y y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y) \\ &= 1 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(2x - e^y y) \\ &= 2 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{2x - e^y y} ((1) - (2)) \\ &= \frac{1}{e^y y - 2x} \end{aligned}$$

Since A depends on y , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= \frac{1}{y} ((2) - (1)) \\ &= \frac{1}{y} \end{aligned}$$

Since B does not depend on x , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned} \mu &= e^{\int B \, dy} \\ &= e^{\int \frac{1}{y} \, dy} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\ln(y)} \\ &= y \end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= y(y) \\ &= y^2 \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= y(2x - e^y y) \\ &= -y^2 e^y + 2yx \end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (y^2) + (-y^2 e^y + 2yx) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int y^2 dx \\ \phi &= x y^2 + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 2yx + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -y^2 e^y + 2yx$. Therefore equation (4) becomes

$$-y^2 e^y + 2yx = 2yx + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -y^2 e^y$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int (-y^2 e^y) dy \\ f(y) &= -(y^2 - 2y + 2) e^y + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = x y^2 - (y^2 - 2y + 2) e^y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = x y^2 - (y^2 - 2y + 2) e^y$$

Summary

The solution(s) found are the following

$$x y^2 - (y^2 - 2y + 2) e^y = c_1 \tag{1}$$

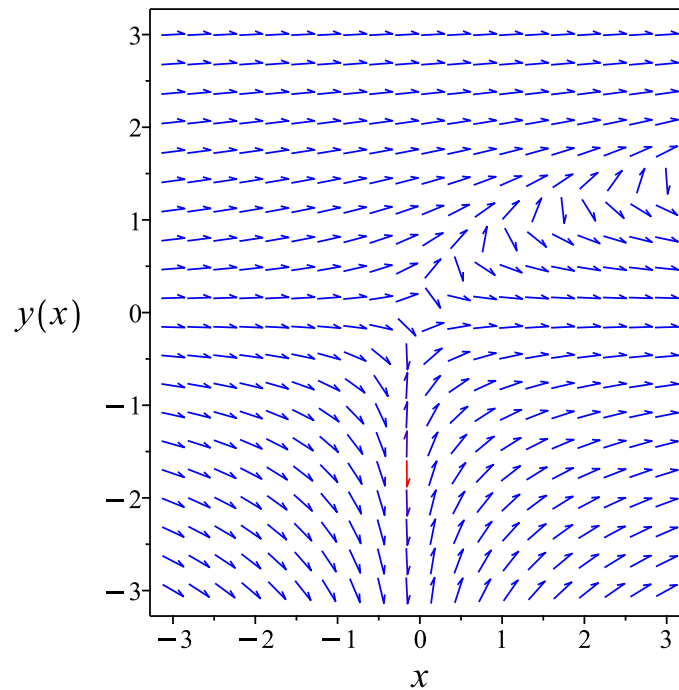


Figure 251: Slope field plot

Verification of solutions

$$x y^2 - (y^2 - 2y + 2) e^y = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
<- 1st order linear successful  
<- inverse linear successful`
```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 34

```
dsolve(y(x)+(2*x-exp(y(x))*y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$\frac{(-y(x)^2 + 2y(x) - 2) e^{y(x)} + xy(x)^2 - c_1}{y(x)^2} = 0$$

✓ Solution by Mathematica

Time used: 0.225 (sec). Leaf size: 32

```
DSolve[y[x]+(2*x-Exp[y[x]]*y[x])*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve} \left[x = \frac{e^{y(x)}(y(x)^2 - 2y(x) + 2)}{y(x)^2} + \frac{c_1}{y(x)^2}, y(x) \right]$$

5.17 problem 22

5.17.1 Solving as separable ode	1309
5.17.2 Solving as first order ode lie symmetry lookup ode	1311
5.17.3 Solving as exact ode	1315
5.17.4 Maple step by step solution	1319

Internal problem ID [559]

Internal file name [OUTPUT/559_Sunday_June_05_2022_01_44_36_AM_63277223/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 22.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$(2 + x) \sin(y) + x \cos(y) y' = 0$$

5.17.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\frac{(2+x)\tan(y)}{x} \end{aligned}$$

Where $f(x) = -\frac{2+x}{x}$ and $g(y) = \tan(y)$. Integrating both sides gives

$$\begin{aligned} \frac{1}{\tan(y)} dy &= -\frac{2+x}{x} dx \\ \int \frac{1}{\tan(y)} dy &= \int -\frac{2+x}{x} dx \\ \ln(\sin(y)) &= -x - 2 \ln(x) + c_1 \end{aligned}$$

Raising both side to exponential gives

$$\sin(y) = e^{-x-2\ln(x)+c_1}$$

Which simplifies to

$$\sin(y) = c_2 e^{-x-2\ln(x)}$$

Summary

The solution(s) found are the following

$$y = \arcsin\left(\frac{c_2 e^{-x+c_1}}{x^2}\right) \quad (1)$$

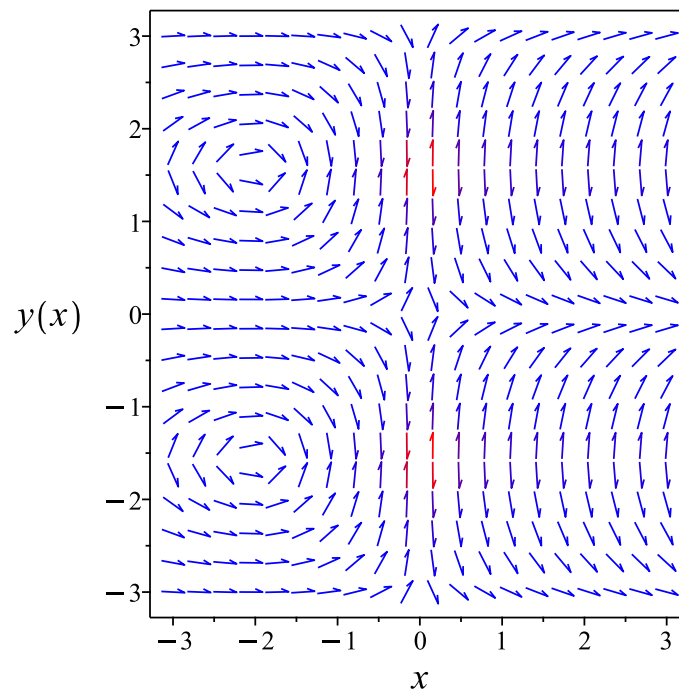


Figure 252: Slope field plot

Verification of solutions

$$y = \arcsin\left(\frac{c_2 e^{-x+c_1}}{x^2}\right)$$

Verified OK.

5.17.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{(2+x)\sin(y)}{x\cos(y)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 254: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= -\frac{x}{2+x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{-\frac{x}{2+x}} dx\end{aligned}$$

Which results in

$$S = -x - 2 \ln(x)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{(2+x) \sin(y)}{x \cos(y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= \frac{-x-2}{x} \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \cot(y) \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \cot(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \ln(\sin(R)) + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-x - 2 \ln(x) = \ln(\sin(y)) + c_1$$

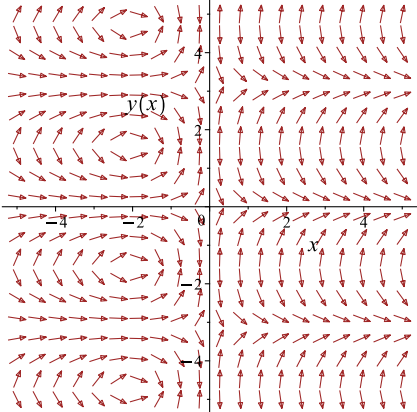
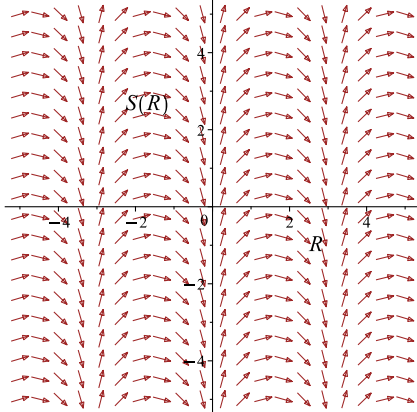
Which simplifies to

$$-x - 2 \ln(x) = \ln(\sin(y)) + c_1$$

Which gives

$$y = \arcsin\left(\frac{e^{-x-c_1}}{x^2}\right)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{(2+x)\sin(y)}{x\cos(y)}$ 	$R = y$ $S = -x - 2 \ln(x)$	$\frac{dS}{dR} = \cot(R)$ 

Summary

The solution(s) found are the following

$$y = \arcsin\left(\frac{e^{-x-c_1}}{x^2}\right) \tag{1}$$

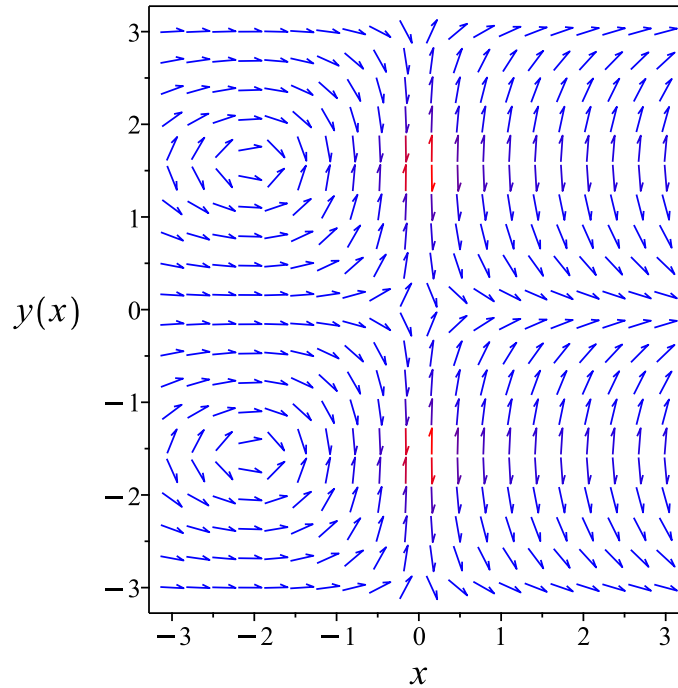


Figure 253: Slope field plot

Verification of solutions

$$y = \arcsin\left(\frac{e^{-x-c_1}}{x^2}\right)$$

Verified OK.

5.17.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}\left(-\frac{\cos(y)}{\sin(y)}\right) dy &= \left(\frac{2+x}{x}\right) dx \\ \left(-\frac{2+x}{x}\right) dx + \left(-\frac{\cos(y)}{\sin(y)}\right) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{2+x}{x} \\ N(x, y) &= -\frac{\cos(y)}{\sin(y)}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{2+x}{x}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(-\frac{\cos(y)}{\sin(y)} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{2+x}{x} dx \\ \phi &= -x - 2 \ln(x) + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{\cos(y)}{\sin(y)}$. Therefore equation (4) becomes

$$-\frac{\cos(y)}{\sin(y)} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$\begin{aligned}f'(y) &= -\frac{\cos(y)}{\sin(y)} \\ &= -\cot(y)\end{aligned}$$

Integrating the above w.r.t y results in

$$\int f'(y) dy = \int (-\cot(y)) dy$$

$$f(y) = -\ln(\sin(y)) + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x - 2\ln(x) - \ln(\sin(y)) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x - 2\ln(x) - \ln(\sin(y))$$

Summary

The solution(s) found are the following

$$-x - 2\ln(x) - \ln(\sin(y)) = c_1 \tag{1}$$

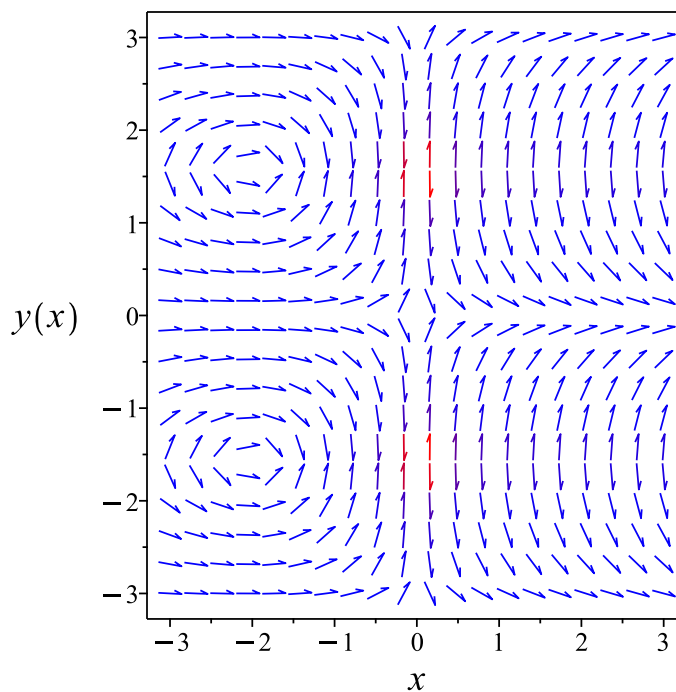


Figure 254: Slope field plot

Verification of solutions

$$-x - 2 \ln(x) - \ln(\sin(y)) = c_1$$

Verified OK.

5.17.4 Maple step by step solution

Let's solve

$$(2 + x) \sin(y) + x \cos(y) y' = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$\frac{y' \cos(y)}{\sin(y)} = -\frac{2+x}{x}$$

- Integrate both sides with respect to x

$$\int \frac{y' \cos(y)}{\sin(y)} dx = \int -\frac{2+x}{x} dx + c_1$$

- Evaluate integral

$$\ln(\sin(y)) = -x - 2 \ln(x) + c_1$$

- Solve for y

$$y = \arcsin\left(\frac{e^{-x+c_1}}{x^2}\right)$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```


✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 16

```
dsolve((2+x)*sin(y(x))+x*cos(y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = \arcsin\left(\frac{e^{-x}}{c_1 x^2}\right)$$

✓ Solution by Mathematica

Time used: 51.022 (sec). Leaf size: 23

```
DSolve[(2+x)*Sin[y[x]]+x*Cos[y[x]]*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \arcsin\left(\frac{e^{-x+c_1}}{x^2}\right)$$

$$y(x) \rightarrow 0$$

5.18 problem 25

- 5.18.1 Solving as homogeneousTypeD2 ode 1321
5.18.2 Solving as exact ode 1323

Internal problem ID [560]

Internal file name [OUTPUT/560_Sunday_June_05_2022_01_44_37_AM_81876546/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 25.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD2**", "**exactWithIntegrationFactor**"

Maple gives the following as the ode type

```
[[_homogeneous, `class D`], _rational]
```

$$2yx + 3x^2y + y^3 + (x^2 + y^2) y' = 0$$

5.18.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$2u(x)x^2 + 3x^3u(x) + u(x)^3x^3 + (x^2 + u(x)^2x^2)(u'(x)x + u(x)) = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{u(u^2 + 3)(x + 1)}{x(u^2 + 1)} \end{aligned}$$

Where $f(x) = -\frac{x+1}{x}$ and $g(u) = \frac{(u^2+3)u}{u^2+1}$. Integrating both sides gives

$$\frac{1}{\frac{(u^2+3)u}{u^2+1}} du = -\frac{x+1}{x} dx$$

$$\int \frac{1}{\frac{(u^2+3)u}{u^2+1}} du = \int -\frac{x+1}{x} dx$$

$$\frac{\ln((u^2+3)u)}{3} = -x - \ln(x) + c_2$$

Raising both side to exponential gives

$$((u^2+3)u)^{\frac{1}{3}} = e^{-x-\ln(x)+c_2}$$

Which simplifies to

$$((u^2+3)u)^{\frac{1}{3}} = c_3 e^{-x-\ln(x)}$$

Which simplifies to

$$((u(x)^2+3)u(x))^{\frac{1}{3}} = \frac{c_3 e^{-x} e^{c_2}}{x}$$

The solution is

$$((u(x)^2+3)u(x))^{\frac{1}{3}} = \frac{c_3 e^{-x} e^{c_2}}{x}$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\left(\frac{\left(\frac{y^2}{x^2} + 3 \right) y}{x} \right)^{\frac{1}{3}} = \frac{c_3 e^{-x} e^{c_2}}{x}$$

$$\left(\frac{(y^2 + 3x^2) y}{x^3} \right)^{\frac{1}{3}} = \frac{c_3 e^{-x+c_2}}{x}$$

Summary

The solution(s) found are the following

$$\left(\frac{(y^2 + 3x^2) y}{x^3} \right)^{\frac{1}{3}} = \frac{c_3 e^{-x+c_2}}{x} \quad (1)$$

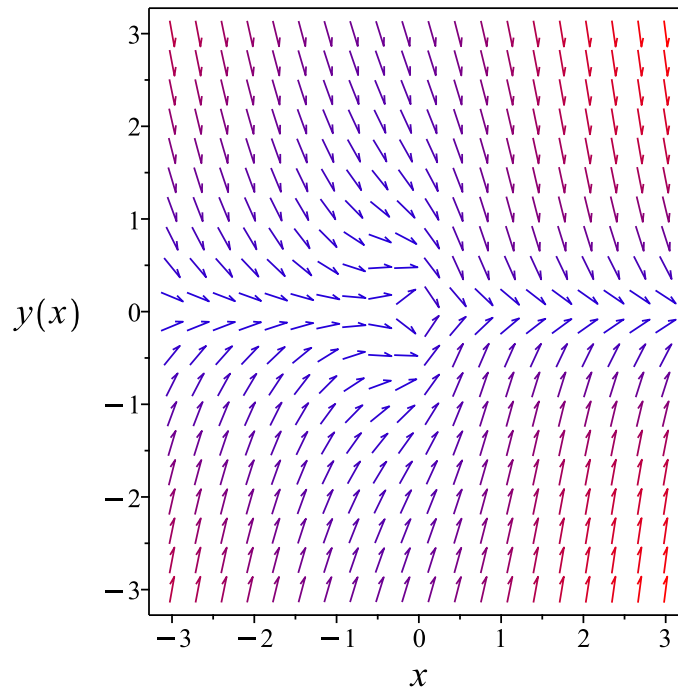


Figure 255: Slope field plot

Verification of solutions

$$\left(\frac{(y^2 + 3x^2)y}{x^3} \right)^{\frac{1}{3}} = \frac{c_3 e^{-x+c_2}}{x}$$

Verified OK.

5.18.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(x^2 + y^2) dy &= (-3y x^2 - y^3 - 2yx) dx \\ (3y x^2 + y^3 + 2yx) dx + (x^2 + y^2) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 3y x^2 + y^3 + 2yx \\ N(x, y) &= x^2 + y^2\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (3y x^2 + y^3 + 2yx) \\ &= 3x^2 + 3y^2 + 2x\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (x^2 + y^2) \\ &= 2x\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{x^2 + y^2} ((3x^2 + 3y^2 + 2x) - (2x)) \\ &= 3 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A \, dx} \\ &= e^{\int 3 \, dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{3x} \\ &= e^{3x} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{3x} (3y x^2 + y^3 + 2yx) \\ &= y(3x^2 + y^2 + 2x) e^{3x} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{3x} (x^2 + y^2) \\ &= (x^2 + y^2) e^{3x} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (y(3x^2 + y^2 + 2x) e^{3x}) + ((x^2 + y^2) e^{3x}) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int y(3x^2 + y^2 + 2x) e^{3x} dx \\ \phi &= \frac{(3x^2 + y^2) y e^{3x}}{3} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned} \frac{\partial \phi}{\partial y} &= \frac{2y^2 e^{3x}}{3} + \frac{(3x^2 + y^2) e^{3x}}{3} + f'(y) \\ &= (x^2 + y^2) e^{3x} + f'(y) \end{aligned} \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = (x^2 + y^2) e^{3x}$. Therefore equation (4) becomes

$$(x^2 + y^2) e^{3x} = (x^2 + y^2) e^{3x} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{(3x^2 + y^2) y e^{3x}}{3} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{(3x^2 + y^2) y e^{3x}}{3}$$

Summary

The solution(s) found are the following

$$\frac{(y^2 + 3x^2) y e^{3x}}{3} = c_1 \quad (1)$$

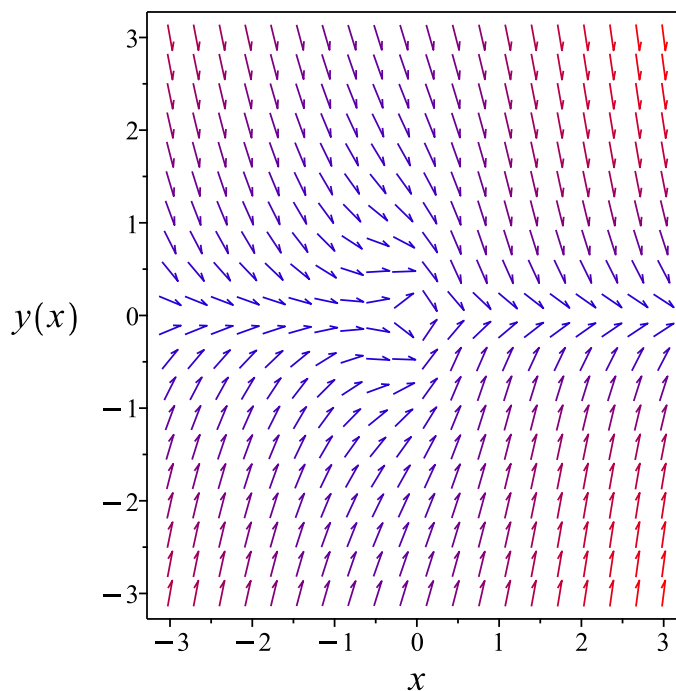


Figure 256: Slope field plot

Verification of solutions

$$\frac{(y^2 + 3x^2) y e^{3x}}{3} = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 297

```
dsolve(2*x*y(x)+3*x^2*y(x)+y(x)^3+(x^2+y(x)^2)*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = -\frac{\left(x^2 e^{6x} c_1^2 - \frac{2^{\frac{1}{3}} \left(\left(1 + \sqrt{4x^6 e^{6x} c_1^2 + 1}\right) e^{6x} c_1^2\right)^{\frac{2}{3}}}{2}\right) 2^{\frac{1}{3}} e^{-3x}}{\left(\left(1 + \sqrt{4x^6 e^{6x} c_1^2 + 1}\right) e^{6x} c_1^2\right)^{\frac{1}{3}} c_1}$$
$$y(x) = -\frac{e^{-3x} 2^{\frac{1}{3}} \left(2x^2 (i\sqrt{3} - 1) e^{6x} c_1^2 + 2^{\frac{1}{3}} (1 + i\sqrt{3}) \left(\left(1 + \sqrt{4x^6 e^{6x} c_1^2 + 1}\right) e^{6x} c_1^2\right)^{\frac{2}{3}}\right)}{4 \left(\left(1 + \sqrt{4x^6 e^{6x} c_1^2 + 1}\right) e^{6x} c_1^2\right)^{\frac{1}{3}} c_1}$$
$$y(x) = \frac{\left(2x^2 (1 + i\sqrt{3}) e^{6x} c_1^2 + 2^{\frac{1}{3}} (i\sqrt{3} - 1) \left(\left(1 + \sqrt{4x^6 e^{6x} c_1^2 + 1}\right) e^{6x} c_1^2\right)^{\frac{2}{3}}\right) e^{-3x} 2^{\frac{1}{3}}}{4 \left(\left(1 + \sqrt{4x^6 e^{6x} c_1^2 + 1}\right) e^{6x} c_1^2\right)^{\frac{1}{3}} c_1}$$

✓ Solution by Mathematica

Time used: 60.305 (sec). Leaf size: 383

`DSolve [2*x*y [x]+3*x^2*y [x]+y [x]^3+(x^2+y [x]^2)*y' [x] == 0,y [x],x,IncludeSingularSolutions ->`

$$y(x) \rightarrow \frac{e^{-3x} \left(-2e^{6x}x^2 + \sqrt[3]{2} \left(\sqrt{4e^{18x}x^6 + e^{6(2x+c_1)}} + e^{6x+3c_1} \right)^{2/3} \right)}{2^{2/3} \sqrt[3]{\sqrt{4e^{18x}x^6 + e^{6(2x+c_1)}} + e^{6x+3c_1}}}$$

$$y(x) \rightarrow \frac{i(\sqrt{3} + i) e^{-3x} \sqrt[3]{\sqrt{4e^{18x}x^6 + e^{6(2x+c_1)}} + e^{6x+3c_1}}}{2\sqrt[3]{2}} + \frac{(1 + i\sqrt{3}) e^{3x}x^2}{2^{2/3} \sqrt[3]{\sqrt{4e^{18x}x^6 + e^{6(2x+c_1)}} + e^{6x+3c_1}}}$$

$$y(x) \rightarrow \frac{(1 - i\sqrt{3}) e^{3x}x^2}{2^{2/3} \sqrt[3]{\sqrt{4e^{18x}x^6 + e^{6(2x+c_1)}} + e^{6x+3c_1}}} - \frac{(1 + i\sqrt{3}) e^{-3x} \sqrt[3]{\sqrt{4e^{18x}x^6 + e^{6(2x+c_1)}} + e^{6x+3c_1}}}{2\sqrt[3]{2}}$$

5.19 problem 26

5.19.1 Solving as linear ode	1330
5.19.2 Solving as first order ode lie symmetry lookup ode	1332
5.19.3 Solving as exact ode	1336
5.19.4 Maple step by step solution	1340

Internal problem ID [561]

Internal file name [OUTPUT/561_Sunday_June_05_2022_01_44_39_AM_71606817/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 26.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$y' - y = -1 + e^{2x}$$

5.19.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$\begin{aligned} p(x) &= -1 \\ q(x) &= -1 + e^{2x} \end{aligned}$$

Hence the ode is

$$y' - y = -1 + e^{2x}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int (-1) dx} \\ &= e^{-x}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) (-1 + e^{2x}) \\ \frac{d}{dx}(e^{-x}y) &= (e^{-x}) (-1 + e^{2x}) \\ d(e^{-x}y) &= (e^x - e^{-x}) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-x}y &= \int e^x - e^{-x} dx \\ e^{-x}y &= e^x + e^{-x} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-x}$ results in

$$y = e^x(e^x + e^{-x}) + c_1e^x$$

which simplifies to

$$y = e^{2x} + 1 + c_1e^x$$

Summary

The solution(s) found are the following

$$y = e^{2x} + 1 + c_1e^x \tag{1}$$

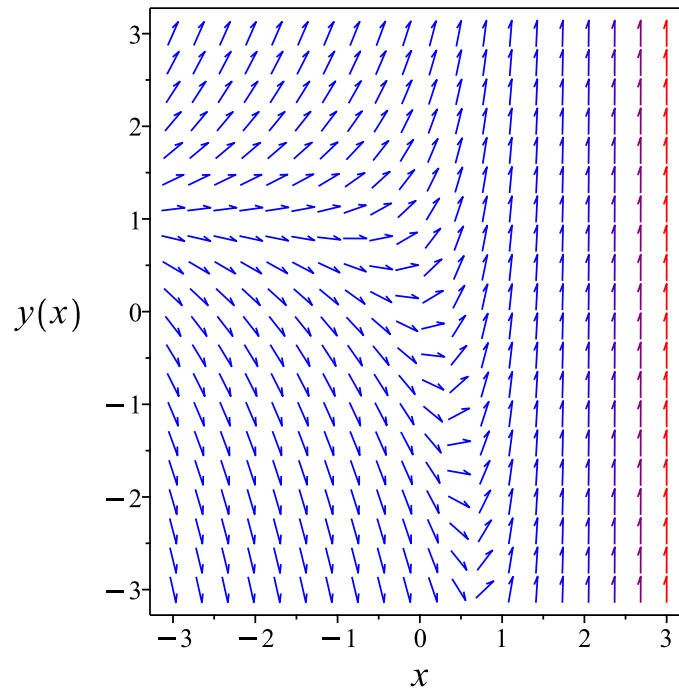


Figure 257: Slope field plot

Verification of solutions

$$y = e^{2x} + 1 + c_1 e^x$$

Verified OK.

5.19.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -1 + e^{2x} + y$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 257: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= e^x\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^x} dy \end{aligned}$$

Which results in

$$S = e^{-x}y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -1 + e^{2x} + y$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -e^{-x}y \\ S_y &= e^{-x} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^x - e^{-x} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^R - e^{-R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = e^R + e^{-R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$y e^{-x} = e^x + e^{-x} + c_1$$

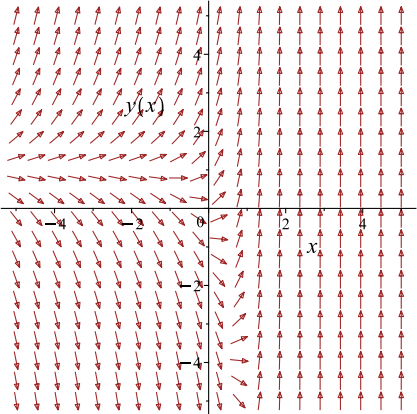
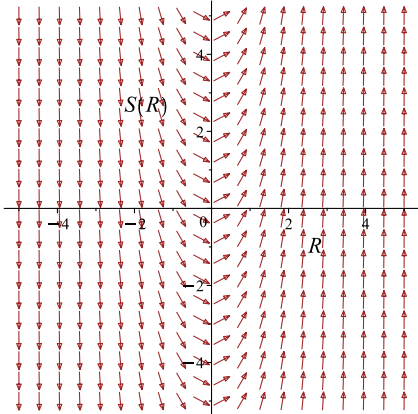
Which simplifies to

$$y e^{-x} = e^x + e^{-x} + c_1$$

Which gives

$$y = (e^x + e^{-x} + c_1) e^x$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -1 + e^{2x} + y$ 	$R = x$ $S = e^{-x}y$	$\frac{dS}{dR} = e^R - e^{-R}$ 

Summary

The solution(s) found are the following

$$y = (e^x + e^{-x} + c_1) e^x \quad (1)$$

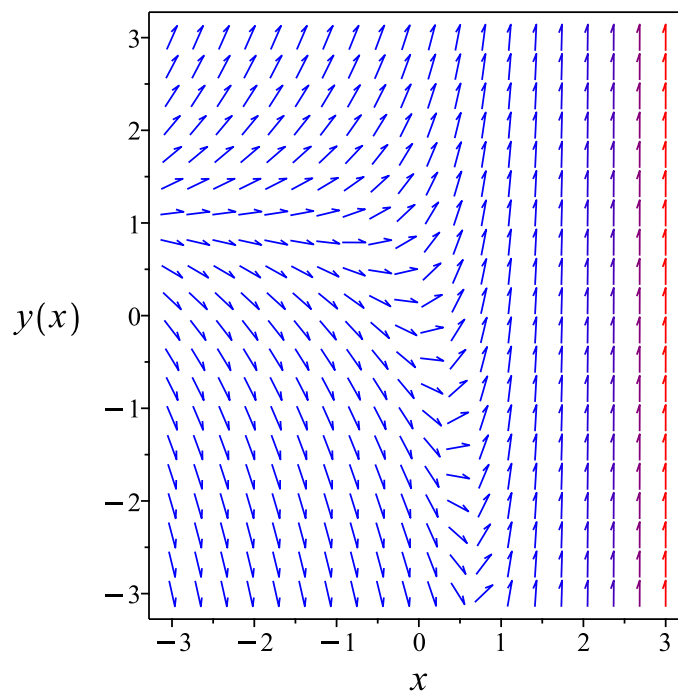


Figure 258: Slope field plot

Verification of solutions

$$y = (e^x + e^{-x} + c_1) e^x$$

Verified OK.

5.19.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}dy &= (-1 + e^{2x} + y) dx \\ (1 - e^{2x} - y) dx + dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 1 - e^{2x} - y \\ N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(1 - e^{2x} - y) \\ &= -1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((-1) - (0)) \\ &= -1 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dx} \\ &= e^{\int -1 dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-x} \\ &= e^{-x} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{-x}(1 - e^{2x} - y) \\ &= e^{-x} - e^x - e^{-x}y \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{-x}(1) \\ &= e^{-x} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (e^{-x} - e^x - e^{-x}y) + (e^{-x}) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int e^{-x} - e^x - e^{-x}y dx \\ \phi &= (y - 1)e^{-x} - e^x + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-x} + f'(y)\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-x}$. Therefore equation (4) becomes

$$e^{-x} = e^{-x} + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = (y - 1)e^{-x} - e^x + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = (y - 1)e^{-x} - e^x$$

The solution becomes

$$y = (e^x + e^{-x} + c_1) e^x$$

Summary

The solution(s) found are the following

$$y = (e^x + e^{-x} + c_1) e^x\tag{1}$$

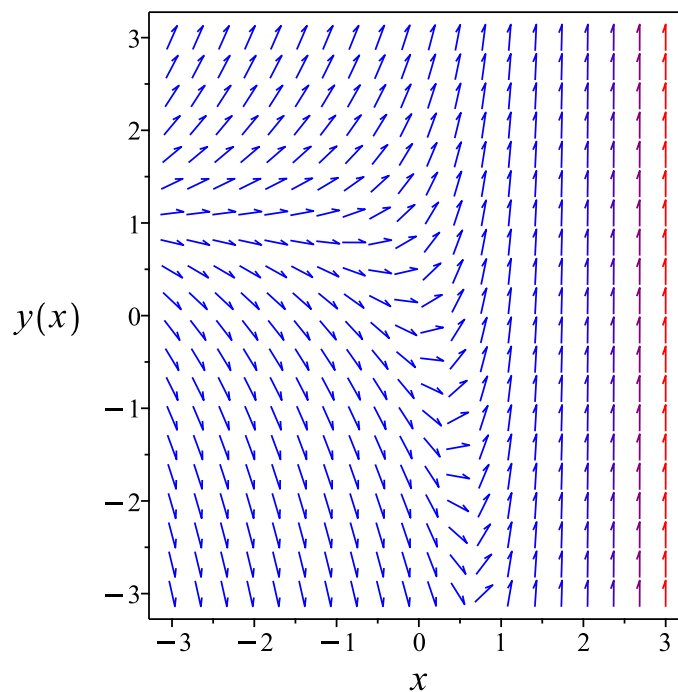


Figure 259: Slope field plot

Verification of solutions

$$y = (e^x + e^{-x} + c_1) e^x$$

Verified OK.

5.19.4 Maple step by step solution

Let's solve

$$y' - y = -1 + e^{2x}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -1 + e^{2x} + y$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' - y = -1 + e^{2x}$$

- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$\mu(x)(y' - y) = \mu(x)(-1 + e^{2x})$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dx}(\mu(x) y)$

$$\mu(x) (y' - y) = \mu'(x) y + \mu(x) y'$$
- Isolate $\mu'(x)$

$$\mu'(x) = -\mu(x)$$
- Solve to find the integrating factor

$$\mu(x) = e^{-x}$$
- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x) y) \right) dx = \int \mu(x) (-1 + e^{2x}) dx + c_1$$
- Evaluate the integral on the lhs

$$\mu(x) y = \int \mu(x) (-1 + e^{2x}) dx + c_1$$
- Solve for y

$$y = \frac{\int \mu(x) (-1 + e^{2x}) dx + c_1}{\mu(x)}$$
- Substitute $\mu(x) = e^{-x}$

$$y = \frac{\int e^{-x} (-1 + e^{2x}) dx + c_1}{e^{-x}}$$
- Evaluate the integrals on the rhs

$$y = \frac{e^x + \frac{1}{e^x} + c_1}{e^{-x}}$$
- Simplify

$$y = e^{2x} + 1 + c_1 e^x$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(diff(y(x),x) = -1+exp(2*x)+y(x),y(x), singsol=all)
```

$$y(x) = e^{2x} + 1 + e^x c_1$$

✓ Solution by Mathematica

Time used: 0.047 (sec). Leaf size: 18

```
DSolve[y'[x] == -1+Exp[2*x]+y[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{2x} + c_1 e^x + 1$$

5.20 problem 27

5.20.1 Solving as differentialType ode 1343

5.20.2 Solving as exact ode 1345

Internal problem ID [562]

Internal file name [OUTPUT/562_Sunday_June_05_2022_01_44_40_AM_44632801/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 27.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "differentialType"

Maple gives the following as the ode type

```
[[_1st_order , ` _with_symmetry_[F(x)*G(y),0] `]]
```

$$\left(-\sin(y) + \frac{x}{y}\right) y' = -1$$

5.20.1 Solving as differentialType ode

Writing the ode as

$$y' = -\frac{1}{-\sin(y) + \frac{x}{y}} \quad (1)$$

Which becomes

$$(y \sin(y)) dy = (x) dy + (y) dx \quad (2)$$

But the RHS is complete differential because

$$(x) dy + (y) dx = d(yx)$$

Hence (2) becomes

$$(y \sin(y)) dy = d(yx)$$

Integrating both sides gives gives the solution as

$$\sin(y) - y \cos(y) = yx + c_1$$

Summary

The solution(s) found are the following

$$\sin(y) - y \cos(y) = yx + c_1 \tag{1}$$

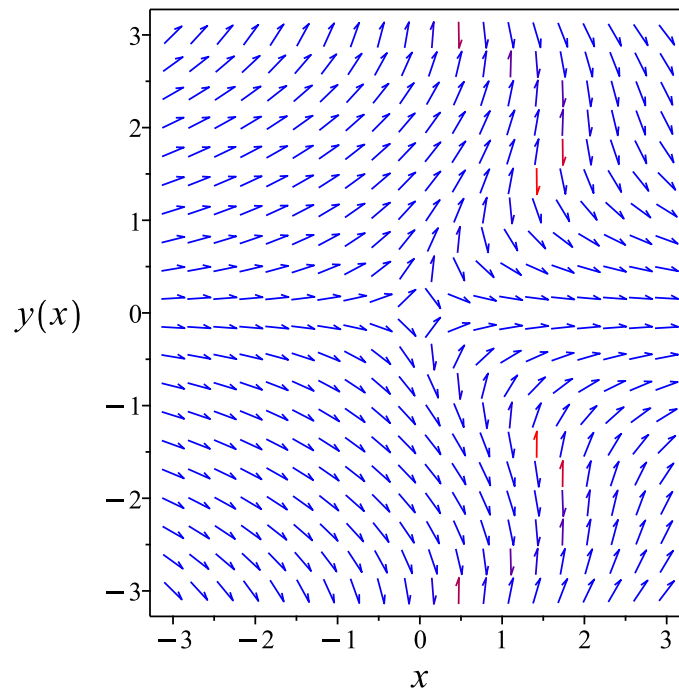


Figure 260: Slope field plot

Verification of solutions

$$\sin(y) - y \cos(y) = yx + c_1$$

Verified OK.

5.20.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (y \sin(y) - x) dy &= (y) dx \\ (-y) dx + (y \sin(y) - x) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -y \\ N(x, y) &= y \sin(y) - x \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-y) \\ &= -1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(y \sin(y) - x) \\ &= -1\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -y dx \\ \phi &= -yx + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = -x + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = y \sin(y) - x$. Therefore equation (4) becomes

$$y \sin(y) - x = -x + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = y \sin (y)$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int (y \sin (y)) dy \\ f(y) &= \sin (y) - y \cos (y) + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -yx + \sin (y) - y \cos (y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -yx + \sin (y) - y \cos (y)$$

Summary

The solution(s) found are the following

$$\sin (y) - y \cos (y) - yx = c_1 \tag{1}$$

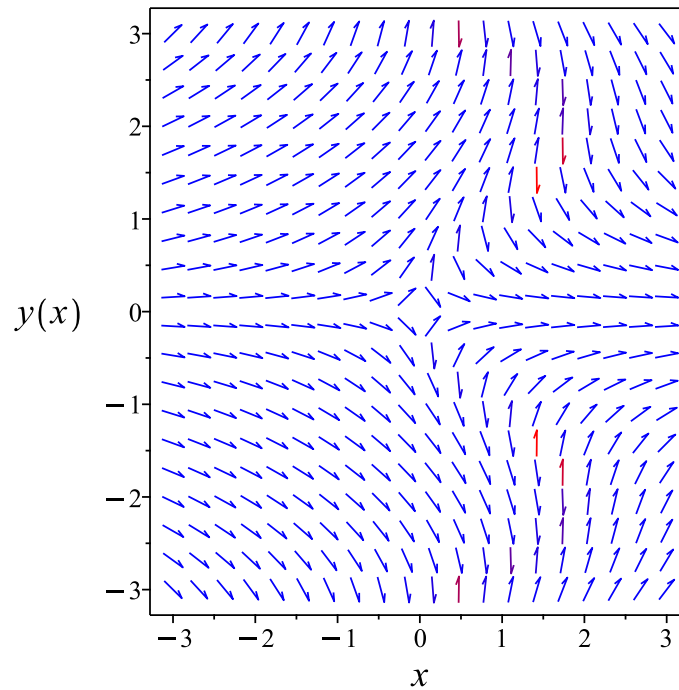


Figure 261: Slope field plot

Verification of solutions

$$\sin(y) - y \cos(y) - yx = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
<- 1st order linear successful
<- inverse linear successful`

```

✓ Solution by Maple

Time used: 0.078 (sec). Leaf size: 25

```
dsolve(1+(-sin(y(x))+x/y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$x + \frac{y(x) \cos(y(x)) - \sin(y(x)) - c_1}{y(x)} = 0$$

✓ Solution by Mathematica

Time used: 0.143 (sec). Leaf size: 29

```
DSolve[1+(-Sin[y[x]]+x/y[x])*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve} \left[x = \frac{\sin(y(x)) - y(x) \cos(y(x))}{y(x)} + \frac{c_1}{y(x)}, y(x) \right]$$

5.21 problem 28

5.21.1 Solving as first order ode lie symmetry calculated ode 1350

5.21.2 Solving as exact ode 1357

Internal problem ID [563]

Internal file name [OUTPUT/563_Sunday_June_05_2022_01_44_42_AM_83508478/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 28.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exactWithIntegrationFactor", "first_order_ode_lie_symmetry_calculated"**

Maple gives the following as the ode type

```
[[_1st_order , _with_exponential_symmetries]]
```

$$y + (-e^{-2y} + 2yx) y' = 0$$

5.21.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{y}{-e^{-2y} + 2yx}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 2 to use as anstaz gives

$$\xi = x^2 a_4 + y x a_5 + y^2 a_6 + x a_2 + y a_3 + a_1 \quad (1\text{E})$$

$$\eta = x^2 b_4 + y x b_5 + y^2 b_6 + x b_2 + y b_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, a_4, a_5, a_6, b_1, b_2, b_3, b_4, b_5, b_6\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} & 2xb_4 + yb_5 + b_2 - \frac{y(-2xa_4 + xb_5 - ya_5 + 2yb_6 - a_2 + b_3)}{-e^{-2y} + 2yx} \\ & - \frac{y^2(xa_5 + 2ya_6 + a_3)}{(-e^{-2y} + 2yx)^2} - \frac{2y^2(x^2a_4 + yxa_5 + y^2a_6 + xa_2 + ya_3 + a_1)}{(-e^{-2y} + 2yx)^2} \quad (5E) \\ & - \left(-\frac{1}{-e^{-2y} + 2yx} + \frac{y(2e^{-2y} + 2x)}{(-e^{-2y} + 2yx)^2} \right) (x^2b_4 \\ & + yxb_5 + y^2b_6 + xb_2 + yb_3 + b_1) = 0 \end{aligned}$$

Putting the above in normal form gives

$$\begin{aligned} & \frac{-10e^{-2y}x^2yb_4 - 6e^{-2y}xy^2b_5 - 2e^{-2y}xya_4 - 6e^{-2y}xyb_2 + e^{-4y}yb_5 + 8x^3y^2b_4 + 4x^2y^3b_5 - 2y^4a_6 - 2y^3a_6 - 4xy^3b_6 + 2x^2y^2a_4 - 2x^2y^2b_5 - y^2a_3 + 4x^2y^2b_2 - 2xy^2b_3 - 2e^{-2y}y^2b_3 - e^{-2y}xb_2 - e^{-2y}ya_2 - 2e^{-2y}yb_1 - 2y^3a_3 - 2y^2a_1 - e^{-2y}b_1 + 2e^{-4y}xb_4 - xy^2a_5 - 2e^{-2y}y^3b_6 - e^{-2y}x^2b_4 - e^{-2y}y^2a_5 + e^{-2y}y^2b_6 + e^{-4y}b_2}{=} = 0 \end{aligned}$$

Setting the numerator to zero gives

$$\begin{aligned} & -10e^{-2y}x^2yb_4 - 6e^{-2y}xy^2b_5 - 2e^{-2y}xya_4 - 6e^{-2y}xyb_2 + e^{-4y}yb_5 \\ & + 8x^3y^2b_4 + 4x^2y^3b_5 - 2y^4a_6 - 2y^3a_6 - 4xy^3b_6 + 2x^2y^2a_4 - 2x^2y^2b_5 \quad (6E) \\ & - y^2a_3 + 4x^2y^2b_2 - 2xy^2b_3 - 2e^{-2y}y^2b_3 - e^{-2y}xb_2 - e^{-2y}ya_2 \\ & - 2e^{-2y}yb_1 - 2y^3a_3 - 2y^2a_1 - e^{-2y}b_1 + 2e^{-4y}xb_4 - xy^2a_5 \\ & - 2e^{-2y}y^3b_6 - e^{-2y}x^2b_4 - e^{-2y}y^2a_5 + e^{-2y}y^2b_6 + e^{-4y}b_2 = 0 \end{aligned}$$

Simplifying the above gives

$$\begin{aligned} & -10e^{-2y}x^2yb_4 - 6e^{-2y}xy^2b_5 - 2e^{-2y}xya_4 - 6e^{-2y}xyb_2 + e^{-4y}yb_5 \\ & + 8x^3y^2b_4 + 4x^2y^3b_5 - 2y^4a_6 - 2y^3a_6 - 4xy^3b_6 + 2x^2y^2a_4 - 2x^2y^2b_5 \quad (6E) \\ & - y^2a_3 + 4x^2y^2b_2 - 2xy^2b_3 - 2e^{-2y}y^2b_3 - e^{-2y}xb_2 - e^{-2y}ya_2 \\ & - 2e^{-2y}yb_1 - 2y^3a_3 - 2y^2a_1 - e^{-2y}b_1 + 2e^{-4y}xb_4 - xy^2a_5 \\ & - 2e^{-2y}y^3b_6 - e^{-2y}x^2b_4 - e^{-2y}y^2a_5 + e^{-2y}y^2b_6 + e^{-4y}b_2 = 0 \end{aligned}$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y, e^{-4y}, e^{-2y}\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2, e^{-4y} = v_3, e^{-2y} = v_4\}$$

The above PDE (6E) now becomes

$$\begin{aligned} &8v_1^3v_2^2b_4 + 4v_1^2v_2^3b_5 + 2v_1^2v_2^2a_4 - 2v_2^4a_6 + 4v_1^2v_2^2b_2 - 10v_4v_1^2v_2b_4 - 2v_1^2v_2^2b_5 \\ &- 6v_4v_1v_2^2b_5 - 4v_1v_2^3b_6 - 2v_4v_2^3b_6 - 2v_2^3a_3 - 2v_4v_1v_2a_4 - v_1v_2^2a_5 - v_4v_2^2a_5 \quad (7E) \\ &- 2v_2^3a_6 - 6v_4v_1v_2b_2 - 2v_1v_2^2b_3 - 2v_4v_2^2b_3 - v_4v_1^2b_4 + v_4v_2^2b_6 - 2v_2^2a_1 \\ &- v_4v_2a_2 - v_2^2a_3 - 2v_4v_2b_1 - v_4v_1b_2 + 2v_3v_1b_4 + v_3v_2b_5 - v_4b_1 + v_3b_2 = 0 \end{aligned}$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2, v_3, v_4\}$$

Equation (7E) now becomes

$$\begin{aligned} &8v_1^3v_2^2b_4 + 4v_1^2v_2^3b_5 + (2a_4 + 4b_2 - 2b_5)v_1^2v_2^2 - 10v_4v_1^2v_2b_4 - v_4v_1^2b_4 \\ &- 4v_1v_2^3b_6 - 6v_4v_1v_2^2b_5 + (-a_5 - 2b_3)v_1v_2^2 + (-2a_4 - 6b_2)v_1v_2v_4 + 2v_3v_1b_4 \quad (8E) \\ &- v_4v_1b_2 - 2v_2^4a_6 - 2v_4v_2^3b_6 + (-2a_3 - 2a_6)v_2^3 + (-a_5 - 2b_3 + b_6)v_2^2v_4 \\ &+ (-2a_1 - a_3)v_2^2 + v_3v_2b_5 + (-a_2 - 2b_1)v_2v_4 + v_3b_2 - v_4b_1 = 0 \end{aligned}$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$b_2 = 0$$

$$b_5 = 0$$

$$-2a_6 = 0$$

$$-b_1 = 0$$

$$-b_2 = 0$$

$$-10b_4 = 0$$

$$-b_4 = 0$$

$$2b_4 = 0$$

$$8b_4 = 0$$

$$-6b_5 = 0$$

$$4b_5 = 0$$

$$-4b_6 = 0$$

$$-2b_6 = 0$$

$$-2a_1 - a_3 = 0$$

$$-a_2 - 2b_1 = 0$$

$$-2a_3 - 2a_6 = 0$$

$$-2a_4 - 6b_2 = 0$$

$$-a_5 - 2b_3 = 0$$

$$2a_4 + 4b_2 - 2b_5 = 0$$

$$-a_5 - 2b_3 + b_6 = 0$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= 0 \\
 a_2 &= 0 \\
 a_3 &= 0 \\
 a_4 &= 0 \\
 a_5 &= -2b_3 \\
 a_6 &= 0 \\
 b_1 &= 0 \\
 b_2 &= 0 \\
 b_3 &= b_3 \\
 b_4 &= 0 \\
 b_5 &= 0 \\
 b_6 &= 0
 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}
 \xi &= -2yx \\
 \eta &= y
 \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}
 \eta &= \eta - \omega(x, y) \xi \\
 &= y - \left(-\frac{y}{-e^{-2y} + 2yx} \right) (-2yx) \\
 &= \frac{y e^{-2y}}{-2yx + e^{-2y}} \\
 \xi &= 0
 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{y e^{-2y}}{-2yx + e^{-2y}}} dy \end{aligned}$$

Which results in

$$S = \ln(y) - e^{2y}x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y}{-e^{-2y} + 2yx}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -e^{2y} \\ S_y &= \frac{-2e^{2y}xy + 1}{y} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

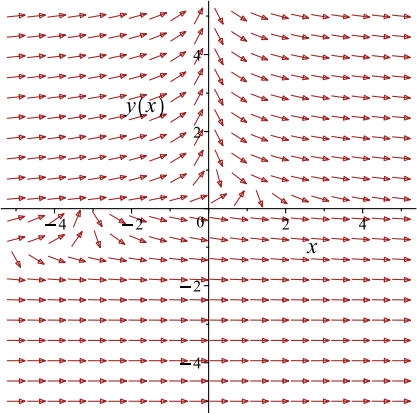
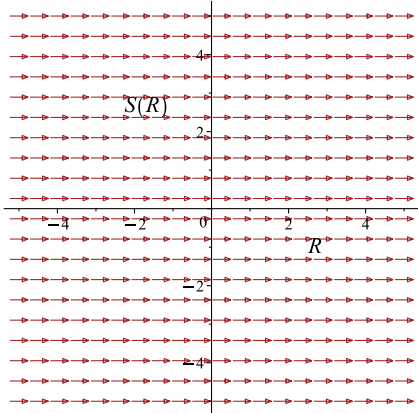
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\ln(y) - e^{2y}x = c_1$$

Which simplifies to

$$\ln(y) - e^{2y}x = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{y}{-e^{-2y} + 2yx}$ 	$R = x$ $S = \ln(y) - e^{2y}x$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$\ln(y) - e^{2y}x = c_1 \tag{1}$$

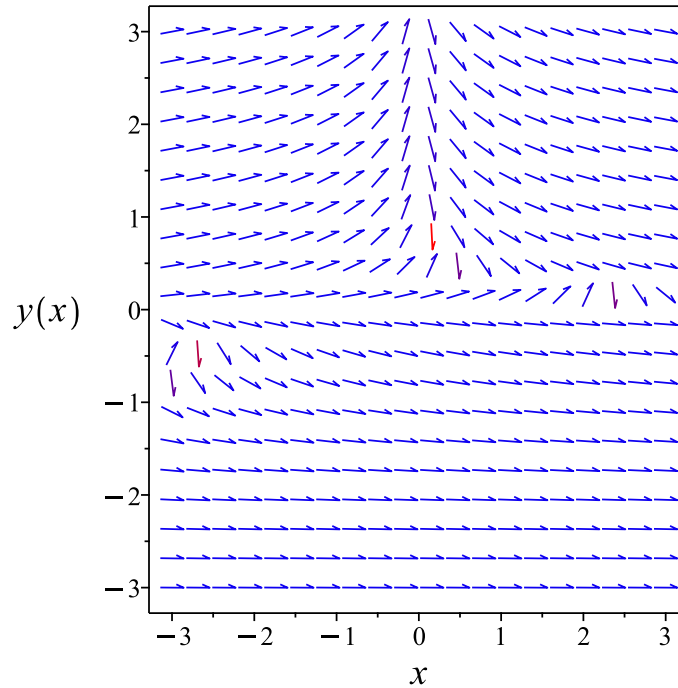


Figure 262: Slope field plot

Verification of solutions

$$\ln(y) - e^{2y}x = c_1$$

Verified OK.

5.21.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(-e^{-2y} + 2yx) dy &= (-y) dx \\ (y) dx + (-e^{-2y} + 2yx) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= y \\ N(x, y) &= -e^{-2y} + 2yx\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-e^{-2y} + 2yx) \\ &= 2y\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{-e^{-2y} + 2yx} ((1) - (2y)) \\ &= \frac{-1 + 2y}{-2yx + e^{-2y}} \end{aligned}$$

Since A depends on y , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= \frac{1}{y} ((2y) - (1)) \\ &= \frac{-1 + 2y}{y} \end{aligned}$$

Since B does not depend on x , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned} \mu &= e^{\int B \, dy} \\ &= e^{\int \frac{-1+2y}{y} \, dy} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{2y - \ln(y)} \\ &= \frac{e^{2y}}{y} \end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{e^{2y}}{y} (y) \\ &= e^{2y} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \frac{e^{2y}}{y} (-e^{-2y} + 2yx) \\ &= \frac{2e^{2y}xy - 1}{y} \end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned}\overline{M} + \overline{N} \frac{dy}{dx} &= 0 \\ (e^{2y}) + \left(\frac{2e^{2y}xy - 1}{y} \right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int e^{2y} dx \\ \phi &= e^{2y}x + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 2e^{2y}x + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{2e^{2y}xy - 1}{y}$. Therefore equation (4) becomes

$$\frac{2e^{2y}xy - 1}{y} = 2e^{2y}x + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{1}{y}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(-\frac{1}{y} \right) dy \\ f(y) &= -\ln(y) + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = e^{2y}x - \ln(y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = e^{2y}x - \ln(y)$$

Summary

The solution(s) found are the following

$$-\ln(y) + e^{2y}x = c_1 \tag{1}$$

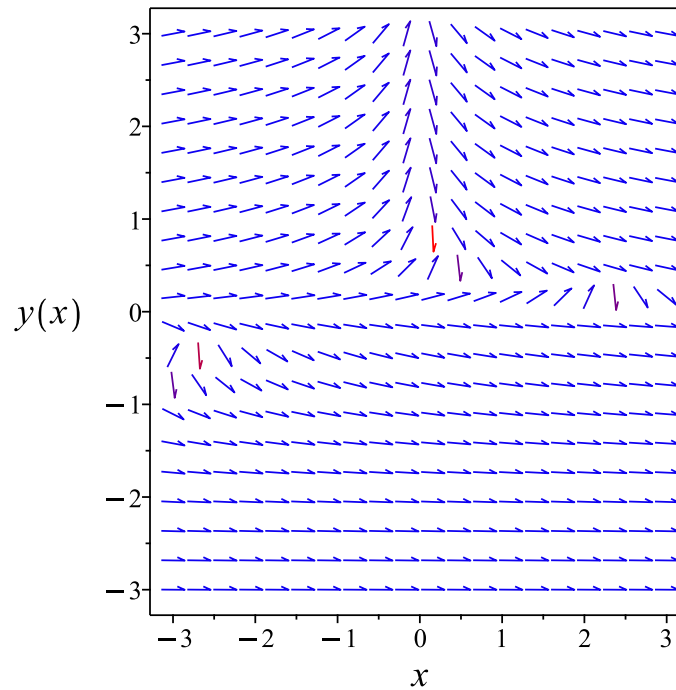


Figure 263: Slope field plot

Verification of solutions

$$-\ln(y) + e^{2y}x = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
<- 1st order linear successful  
<- inverse linear successful`
```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 24

```
dsolve(y(x)+(-exp(-2*y(x))+2*x*y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = e^{\text{RootOf}(c_1 e^{-2e^{-Z}} + Ze^{-2e^{-Z}} - x)}$$

✓ Solution by Mathematica

Time used: 0.243 (sec). Leaf size: 25

```
DSolve[y[x]+(-Exp[-2*y[x]]+2*x*y[x])*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve}[x = e^{-2y(x)} \log(y(x)) + c_1 e^{-2y(x)}, y(x)]$$

5.22 problem 29

5.22.1 Solving as exact ode 1363

Internal problem ID [564]

Internal file name [OUTPUT/564_Sunday_June_05_2022_01_44_43_AM_26805416/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 29.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exactWithIntegrationFactor"**

Maple gives the following as the ode type

```
[[_1st_order , `_with_symmetry_[F(x)*G(y),0]`]]
```

$$(e^x \cot(y) + 2 \csc(y) y) y' = -e^x$$

5.22.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (e^x \cot(y) + 2 \csc(y) y) dy &= (-e^x) dx \\ (e^x) dx + (e^x \cot(y) + 2 \csc(y) y) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= e^x \\ N(x, y) &= e^x \cot(y) + 2 \csc(y) y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(e^x) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(e^x \cot(y) + 2 \csc(y) y) \\ &= e^x \cot(y) \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{\sin(y)}{e^x \cos(y) + 2y} ((0) - (e^x \cot(y))) \\ &= -\frac{\cos(y) e^x}{e^x \cos(y) + 2y} \end{aligned}$$

Since A depends on y , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= e^{-x} ((e^x \cot(y)) - (0)) \\ &= \cot(y) \end{aligned}$$

Since B does not depend on x , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned} \mu &= e^{\int B \, dy} \\ &= e^{\int \cot(y) \, dy} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\ln(\sin(y))} \\ &= \sin(y) \end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \sin(y) (e^x) \\ &= e^x \sin(y) \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \sin(y) (e^x \cot(y) + 2 \csc(y) y) \\ &= e^x \cos(y) + 2y \end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned} \overline{M} + \overline{N} \frac{dy}{dx} &= 0 \\ (e^x \sin(y)) + (e^x \cos(y) + 2y) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int e^x \sin(y) dx \\ \phi &= e^x \sin(y) + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^x \cos(y) + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^x \cos(y) + 2y$. Therefore equation (4) becomes

$$e^x \cos(y) + 2y = e^x \cos(y) + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 2y$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int (2y) dy \\ f(y) &= y^2 + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = e^x \sin(y) + y^2 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = e^x \sin(y) + y^2$$

Summary

The solution(s) found are the following

$$e^x \sin(y) + y^2 = c_1 \tag{1}$$

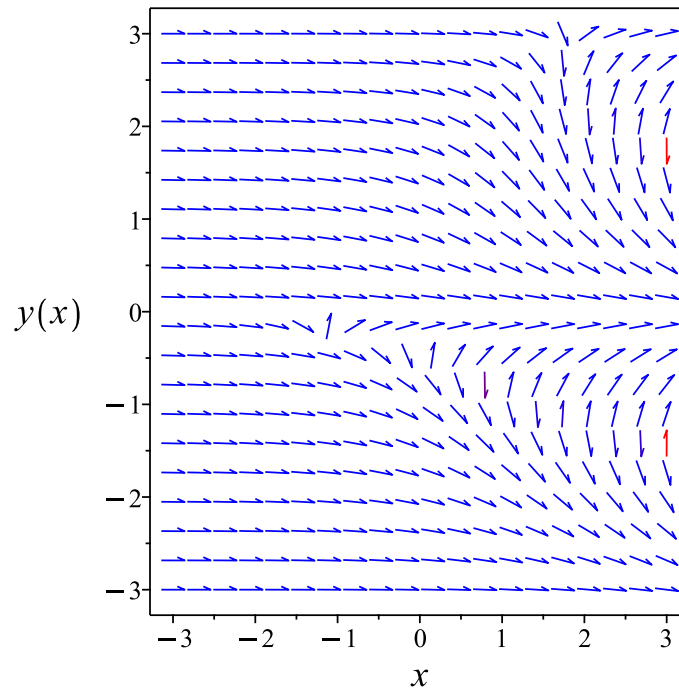


Figure 264: Slope field plot

Verification of solutions

$$e^x \sin(y) + y^2 = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 15

```
dsolve(exp(x)+(exp(x)*cot(y(x))+2*csc(y(x))*y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$e^x \sin(y(x)) + y(x)^2 + c_1 = 0$$

✓ Solution by Mathematica

Time used: 0.312 (sec). Leaf size: 18

```
DSolve[Exp[x]+(Exp[x]*Cot[y[x]]+2*Csc[y[x]]*y[x])*y'[x] == 0,y[x],x,IncludeSingularSolutions
```

$$\text{Solve}[y(x)^2 + e^x \sin(y(x)) = c_1, y(x)]$$

5.23 problem 30

- 5.23.1 Solving as differentialType ode 1369
5.23.2 Solving as exact ode 1371

Internal problem ID [565]

Internal file name [OUTPUT/565_Sunday_June_05_2022_01_44_50_AM_82120052/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 30.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact"**, **"differentialType"**

Maple gives the following as the ode type

`[_rational]`

$$\frac{4x^3}{y^2} + \frac{3}{y} + \left(\frac{3x}{y^2} + 4y\right) y' = 0$$

5.23.1 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-\frac{4x^3}{y^2} - \frac{3}{y}}{\frac{3x}{y^2} + 4y} \quad (1)$$

Which becomes

$$(4y^3) dy = (-3x) dy + (-4x^3 - 3y) dx \quad (2)$$

But the RHS is complete differential because

$$(-3x) dy + (-4x^3 - 3y) dx = d(-x^4 - 3yx)$$

Hence (2) becomes

$$(4y^3) dy = d(-x^4 - 3yx)$$

Integrating both sides gives gives the solution as

$$y^4 = -x^4 - 3yx + c_1$$

Summary

The solution(s) found are the following

$$y^4 = -x^4 - 3yx + c_1 \tag{1}$$

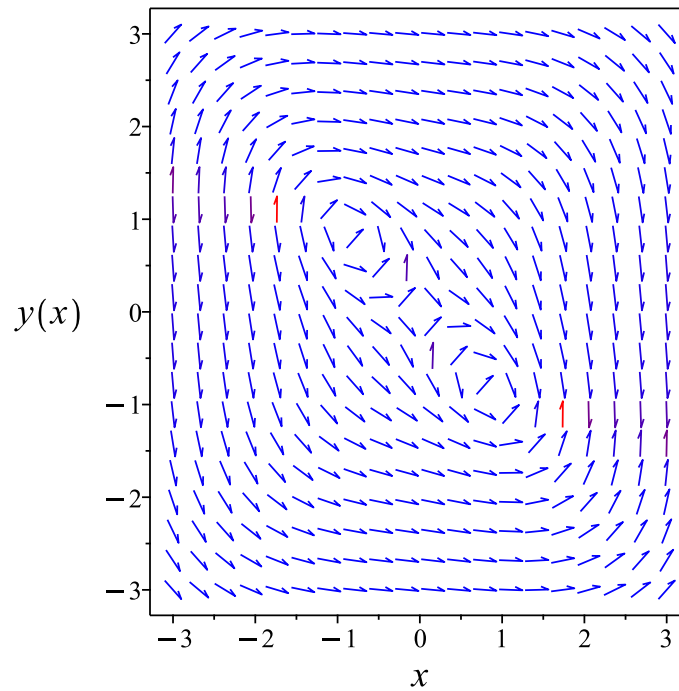


Figure 265: Slope field plot

Verification of solutions

$$y^4 = -x^4 - 3yx + c_1$$

Verified OK.

5.23.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (4y^3 + 3x) dy &= (-4x^3 - 3y) dx \\ (4x^3 + 3y) dx + (4y^3 + 3x) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 4x^3 + 3y \\ N(x, y) &= 4y^3 + 3x \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(4x^3 + 3y) \\ &= 3\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(4y^3 + 3x) \\ &= 3\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 4x^3 + 3y dx \\ \phi &= x(x^3 + 3y) + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 3x + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 4y^3 + 3x$. Therefore equation (4) becomes

$$4y^3 + 3x = 3x + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 4y^3$$

Integrating the above w.r.t y gives

$$\int f'(y) \, dy = \int (4y^3) \, dy$$

$$f(y) = y^4 + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = x(x^3 + 3y) + y^4 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = x(x^3 + 3y) + y^4$$

Summary

The solution(s) found are the following

$$x(3y + x^3) + y^4 = c_1 \tag{1}$$

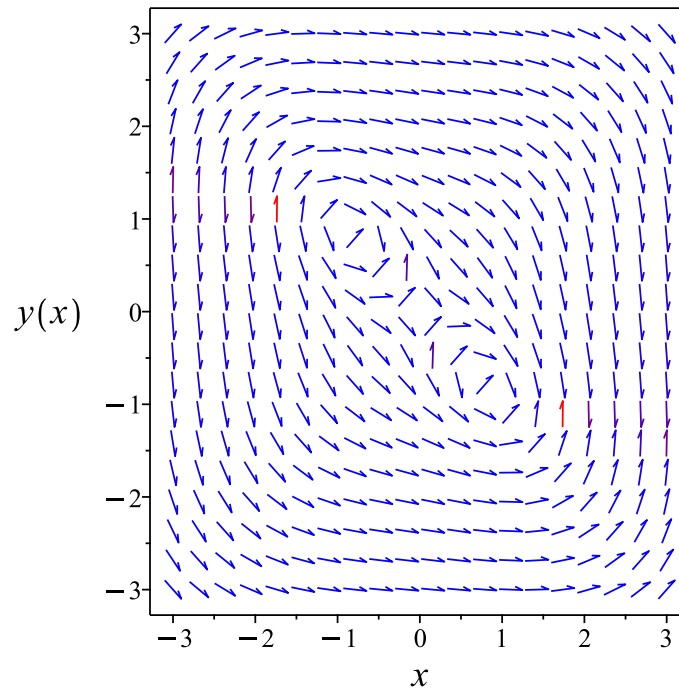


Figure 266: Slope field plot

Verification of solutions

$$x(3y + x^3) + y^4 = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(4*x^3/y(x)^2+3/y(x)+(3*x/y(x)^2+4*y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$x^4 + y(x)^4 + 3xy(x) + c_1 = 0$$

✓ Solution by Mathematica

Time used: 60.148 (sec). Leaf size: 1181

`DSolve[4*x^3/y[x]^2+3/y[x]+(3*x/y[x]^2+4*y[x])*y'[x]== 0,y[x],x,IncludeSingularSolutions ->`

$y(x) \rightarrow$

$$-\frac{1}{2} \sqrt{\frac{6x}{\sqrt{\frac{4\sqrt[3]{2}(x^4-c_1)}{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}} + \frac{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}{3\sqrt[3]{2}}}} - \sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}}$$

$$-\frac{1}{2} \sqrt{\frac{4\sqrt[3]{2}(x^4 - c_1)}{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}} + \frac{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}{3\sqrt[3]{2}}}}$$

$y(x)$

$$\rightarrow \frac{1}{2} \sqrt{\frac{6x}{\sqrt{\frac{4\sqrt[3]{2}(x^4-c_1)}{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}} + \frac{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}{3\sqrt[3]{2}}}} - \sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}}$$

$$-\frac{1}{2} \sqrt{\frac{4\sqrt[3]{2}(x^4 - c_1)}{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}} + \frac{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}{3\sqrt[3]{2}}}}$$

$y(x)$

$$\rightarrow \frac{1}{2} \sqrt{\frac{4\sqrt[3]{2}(x^4 - c_1)}{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}} + \frac{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}{3\sqrt[3]{2}}}}$$

$$-\frac{1}{2} \sqrt{\frac{6x}{\sqrt{\frac{4\sqrt[3]{2}(x^4-c_1)}{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}} + \frac{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}{3\sqrt[3]{2}}}} - \sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}}$$

$y(x)$

$$\rightarrow \frac{1}{2} \sqrt{\frac{6x}{\sqrt{\frac{4\sqrt[3]{2}(x^4-c_1)}{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}} + \frac{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}{3\sqrt[3]{2}}}} - \sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}}$$

$$+\frac{1}{2} \sqrt{\frac{4\sqrt[3]{2}(x^4 - c_1) \quad 1376}{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}} + \frac{\sqrt[3]{243x^2 + \sqrt{59049x^4 - 6912(x^4 - c_1)^3}}}{3\sqrt[3]{2}}}}$$

5.24 problem 30

- 5.24.1 Solving as differentialType ode 1377
5.24.2 Solving as exact ode 1382

Internal problem ID [566]

Internal file name [OUTPUT/566_Sunday_June_05_2022_01_44_51_AM_23911975/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 30.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "differentialType"

Maple gives the following as the ode type

[_rational]

$$\frac{6}{y} + \left(\frac{x^2}{y} + \frac{3y}{x} \right) y' = -3x$$

5.24.1 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-3x - \frac{6}{y}}{\frac{x^2}{y} + \frac{3y}{x}} \quad (1)$$

Which becomes

$$(3y^2) dy = (-x^3) dy + (-3x(yx + 2)) dx \quad (2)$$

But the RHS is complete differential because

$$(-x^3) dy + (-3x(yx + 2)) dx = d(-y x^3 - 3x^2)$$

Hence (2) becomes

$$(3y^2) dy = d(-y x^3 - 3x^2)$$

Integrating both sides gives gives these solutions

$$y = \frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{6} - \frac{2x^3}{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}$$

$$y = -\frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{12} + \frac{x^3}{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}$$

$$y = -\frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{12} + \frac{x^3}{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{6 \cdot 2x^3} + c_1 \quad (1)$$

$$y = -\frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{12x^3} + \frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{2}$$

$$+ \frac{i\sqrt{3} \left(\frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{6} + \frac{2x^3}{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}} \right)}{2} + c_1$$

$$y = -\frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{12x^3} + \frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{2} + c_1 \quad (3)$$

$$+ \frac{i\sqrt{3} \left(\frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{6} + \frac{2x^3}{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}} \right)}{2}$$

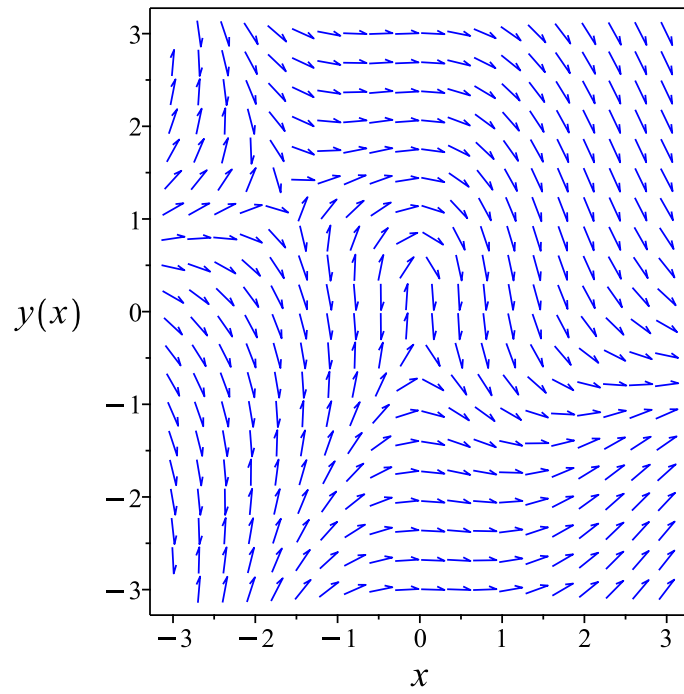


Figure 267: Slope field plot

Verification of solutions

$$y = \frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{6} - \frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{2x^3} + c_1$$

Verified OK.

$$y = -\frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{12} + \frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{x^3} + i\sqrt{3} \left(\frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{6} + \frac{2x^3}{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}} \right) + c_1$$

Verified OK.

$$y = -\frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{12} + \frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{x^3} + i\sqrt{3} \left(\frac{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{6} + \frac{2x^3}{\left(-324x^2 + 108c_1 + 12\sqrt{12x^9 + 729x^4 - 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}} \right) - c_1$$

Verified OK.

5.24.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (x^3 + 3y^2) dy &= (-3x(yx + 2)) dx \\ (3x(yx + 2)) dx + (x^3 + 3y^2) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 3x(yx + 2) \\ N(x, y) &= x^3 + 3y^2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(3x(yx + 2)) \\ &= 3x^2\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x^3 + 3y^2) \\ &= 3x^2\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 3x(yx + 2) dx \\ \phi &= yx^3 + 3x^2 + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x^3 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = x^3 + 3y^2$. Therefore equation (4) becomes

$$x^3 + 3y^2 = x^3 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 3y^2$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (3y^2) dy$$

$$f(y) = y^3 + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = yx^3 + y^3 + 3x^2 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = yx^3 + y^3 + 3x^2$$

Summary

The solution(s) found are the following

$$yx^3 + y^3 + 3x^2 = c_1 \tag{1}$$

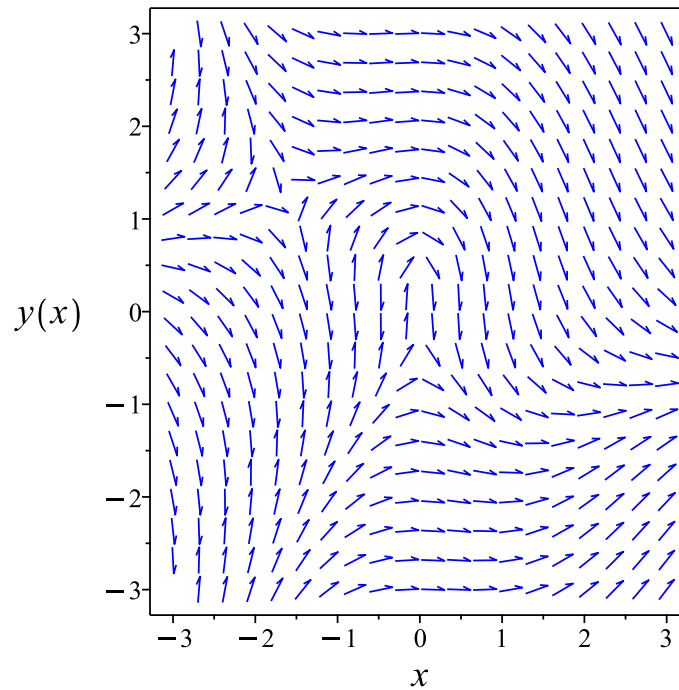


Figure 268: Slope field plot

Verification of solutions

$$yx^3 + y^3 + 3x^2 = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 326

```
dsolve(3*x+6/y(x)+(x^2/y(x)+3*y(x)/x)*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = \frac{-12x^3 + \left(-324x^2 - 108c_1 + 12\sqrt{12x^9 + 729x^4 + 486c_1x^2 + 81c_1^2}\right)^{\frac{2}{3}}}{6 \left(-324x^2 - 108c_1 + 12\sqrt{12x^9 + 729x^4 + 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}$$

$$y(x) = -\frac{(1 + i\sqrt{3}) \left(-324x^2 - 108c_1 + 12\sqrt{12x^9 + 729x^4 + 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{x^3(i\sqrt{3} - 1)} - \frac{12}{\left(-324x^2 - 108c_1 + 12\sqrt{12x^9 + 729x^4 + 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}$$

$$y(x) = \frac{12i\sqrt{3}x^3 + i\sqrt{3} \left(-324x^2 - 108c_1 + 12\sqrt{12x^9 + 729x^4 + 486c_1x^2 + 81c_1^2}\right)^{\frac{2}{3}} + 12x^3 - \left(-324x^2 - 108c_1 + 12\sqrt{12x^9 + 729x^4 + 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}{12 \left(-324x^2 - 108c_1 + 12\sqrt{12x^9 + 729x^4 + 486c_1x^2 + 81c_1^2}\right)^{\frac{1}{3}}}$$

✓ Solution by Mathematica

Time used: 4.558 (sec). Leaf size: 331

`DSolve[3*x+6/y[x]+(x^2/y[x]+3*y[x]/x)*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]`

$$\begin{aligned}
 y(x) &\rightarrow \frac{\sqrt[3]{-81x^2 + \sqrt{108x^9 + 729(-3x^2 + c_1)^2 + 27c_1}}}{3\sqrt[3]{2}\sqrt[3]{2x^3}} \\
 &\quad - \frac{\sqrt[3]{-81x^2 + \sqrt{108x^9 + 729(-3x^2 + c_1)^2 + 27c_1}}}{\sqrt[3]{2x^3}} \\
 y(x) &\rightarrow \frac{(-1 + i\sqrt{3})\sqrt[3]{-81x^2 + \sqrt{108x^9 + 729(-3x^2 + c_1)^2 + 27c_1}}}{6\sqrt[3]{2}} \\
 &\quad + \frac{(1 + i\sqrt{3})x^3}{2^{2/3}\sqrt[3]{-81x^2 + \sqrt{108x^9 + 729(-3x^2 + c_1)^2 + 27c_1}}} \\
 y(x) &\rightarrow \frac{(1 - i\sqrt{3})x^3}{2^{2/3}\sqrt[3]{-81x^2 + \sqrt{108x^9 + 729(-3x^2 + c_1)^2 + 27c_1}}} \\
 &\quad - \frac{(1 + i\sqrt{3})\sqrt[3]{-81x^2 + \sqrt{108x^9 + 729(-3x^2 + c_1)^2 + 27c_1}}}{6\sqrt[3]{2}}
 \end{aligned}$$

5.25 problem 32

5.25.1 Solving as homogeneousTypeD2 ode	1388
5.25.2 Solving as first order ode lie symmetry calculated ode	1390
5.25.3 Solving as exact ode	1396

Internal problem ID [567]

Internal file name [OUTPUT/567_Sunday_June_05_2022_01_44_53_AM_42477074/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Section 2.6. Page 100

Problem number: 32.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD2**", "**exactWithIntegrationFactor**", "**first_order_ode_lie_symmetry_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class B`]]
```

$$3yx + y^2 + (x^2 + yx)y' = 0$$

5.25.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$3u(x)x^2 + u(x)^2x^2 + (x^2 + u(x)x^2)(u'(x)x + u(x)) = 0$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{2u(u+2)}{x(u+1)}\end{aligned}$$

Where $f(x) = -\frac{2}{x}$ and $g(u) = \frac{u(u+2)}{u+1}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{u(u+2)}{u+1}} du &= -\frac{2}{x} dx \\ \int \frac{1}{\frac{u(u+2)}{u+1}} du &= \int -\frac{2}{x} dx \\ \frac{\ln(u(u+2))}{2} &= -2 \ln(x) + c_2\end{aligned}$$

Raising both side to exponential gives

$$\sqrt{u(u+2)} = e^{-2 \ln(x) + c_2}$$

Which simplifies to

$$\sqrt{u(u+2)} = \frac{c_3}{x^2}$$

Which simplifies to

$$\sqrt{u(x)(u(x)+2)} = \frac{c_3 e^{c_2}}{x^2}$$

The solution is

$$\sqrt{u(x)(u(x)+2)} = \frac{c_3 e^{c_2}}{x^2}$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\begin{aligned}\sqrt{\frac{y(\frac{y}{x}+2)}{x}} &= \frac{c_3 e^{c_2}}{x^2} \\ \sqrt{\frac{y(2x+y)}{x^2}} &= \frac{c_3 e^{c_2}}{x^2}\end{aligned}$$

Summary

The solution(s) found are the following

$$\sqrt{\frac{y(2x+y)}{x^2}} = \frac{c_3 e^{c_2}}{x^2} \quad (1)$$

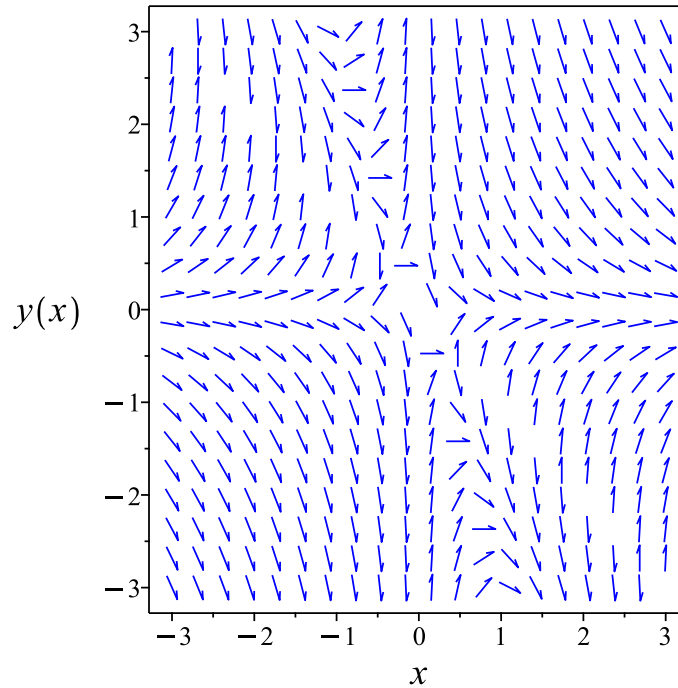


Figure 269: Slope field plot

Verification of solutions

$$\sqrt{\frac{y(2x+y)}{x^2}} = \frac{c_3 e^{c_2}}{x^2}$$

Verified OK.

5.25.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{y(3x+y)}{x(x+y)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \tag{1E}$$

$$\eta = xb_2 + yb_3 + b_1 \tag{2E}$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 - \frac{y(3x+y)(b_3-a_2)}{x(x+y)} - \frac{y^2(3x+y)^2 a_3}{x^2(x+y)^2} \\ - \left(-\frac{3y}{x(x+y)} + \frac{y(3x+y)}{x^2(x+y)} + \frac{y(3x+y)}{x(x+y)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left(-\frac{3x+y}{x(x+y)} - \frac{y}{x(x+y)} + \frac{y(3x+y)}{x(x+y)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{4x^4b_2 + 4x^3yb_2 + 2x^2y^2a_2 - 12x^2y^2a_3 + 2x^2y^2b_2 - 2x^2y^2b_3 - 8xy^3a_3 - 2y^4a_3 + 3x^3b_1 - 3x^2ya_1 + 2x^2yb_1}{x^2(x+y)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} 4x^4b_2 + 4x^3yb_2 + 2x^2y^2a_2 - 12x^2y^2a_3 + 2x^2y^2b_2 - 2x^2y^2b_3 - 8xy^3a_3 \\ - 2y^4a_3 + 3x^3b_1 - 3x^2ya_1 + 2x^2yb_1 - 2xy^2a_1 + xy^2b_1 - y^3a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} 2a_2v_1^2v_2^2 - 12a_3v_1^2v_2^2 - 8a_3v_1v_2^3 - 2a_3v_2^4 + 4b_2v_1^4 + 4b_2v_1^3v_2 + 2b_2v_1^2v_2^2 \\ - 2b_3v_1^2v_2^2 - 3a_1v_1^2v_2 - 2a_1v_1v_2^2 - a_1v_2^3 + 3b_1v_1^3 + 2b_1v_1^2v_2 + b_1v_1v_2^2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$4b_2v_1^4 + 4b_2v_1^3v_2 + 3b_1v_1^3 + (2a_2 - 12a_3 + 2b_2 - 2b_3)v_1^2v_2^2 + (-3a_1 + 2b_1)v_1^2v_2 - 8a_3v_1v_2^3 + (-2a_1 + b_1)v_1v_2^2 - 2a_3v_2^4 - a_1v_2^3 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -a_1 &= 0 \\ -8a_3 &= 0 \\ -2a_3 &= 0 \\ 3b_1 &= 0 \\ 4b_2 &= 0 \\ -3a_1 + 2b_1 &= 0 \\ -2a_1 + b_1 &= 0 \\ 2a_2 - 12a_3 + 2b_2 - 2b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left(-\frac{y(3x + y)}{x(x + y)} \right) (x) \\ &= \frac{4yx + 2y^2}{x + y} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{4yx + 2y^2}{x + y}} dy\end{aligned}$$

Which results in

$$S = \frac{\ln(y(2x + y))}{4}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y(3x + y)}{x(x + y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= \frac{1}{4x + 2y} \\S_y &= \frac{x + y}{2y(2x + y)}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{2R} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{2R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{\ln(R)}{2} + c_1 \quad (4)$$

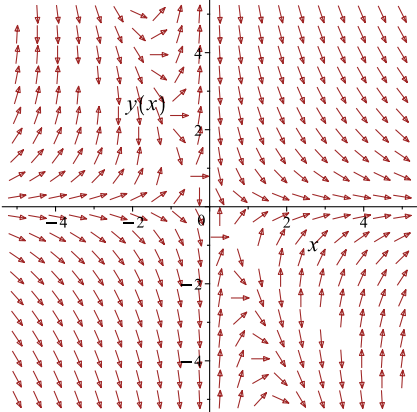
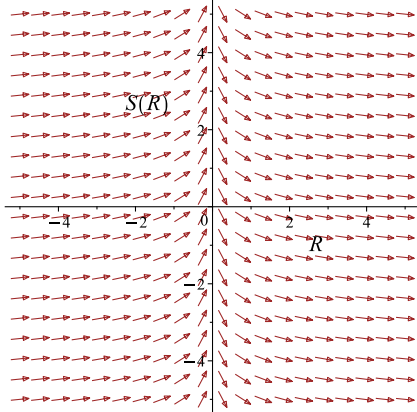
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\ln(y)}{4} + \frac{\ln(2x + y)}{4} = -\frac{\ln(x)}{2} + c_1$$

Which simplifies to

$$\frac{\ln(y)}{4} + \frac{\ln(2x + y)}{4} = -\frac{\ln(x)}{2} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{y(3x+y)}{x(x+y)}$ 	$R = x$ $S = \frac{\ln(y)}{4} + \frac{\ln(2x+y)}{4}$	$\frac{dS}{dR} = -\frac{1}{2R}$ 

Summary

The solution(s) found are the following

$$\frac{\ln(y)}{4} + \frac{\ln(2x+y)}{4} = -\frac{\ln(x)}{2} + c_1 \tag{1}$$

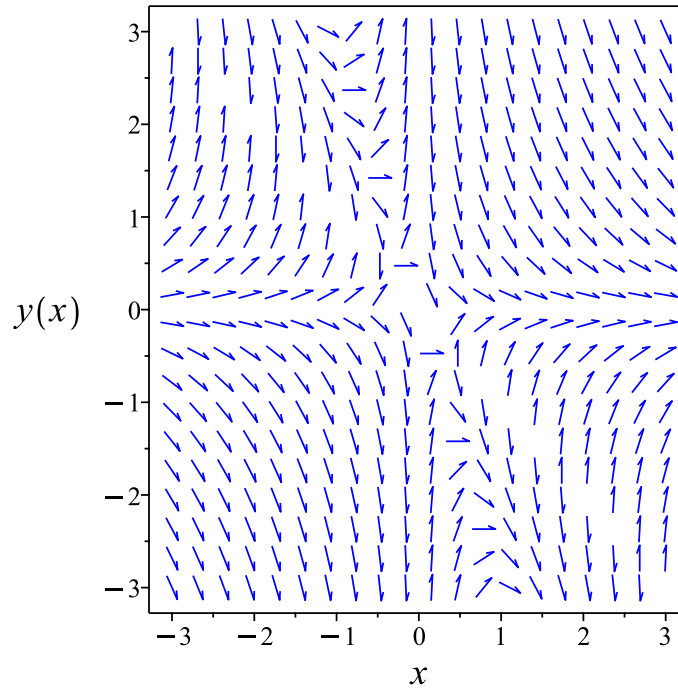


Figure 270: Slope field plot

Verification of solutions

$$\frac{\ln(y)}{4} + \frac{\ln(2x+y)}{4} = -\frac{\ln(x)}{2} + c_1$$

Verified OK.

5.25.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(x^2 + yx) dy &= (-3yx - y^2) dx \\ (3yx + y^2) dx + (x^2 + yx) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 3yx + y^2 \\ N(x, y) &= x^2 + yx\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(3yx + y^2) \\ &= 3x + 2y\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x^2 + yx) \\ &= 2x + y\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{x(x+y)} ((3x+2y) - (2x+y)) \\ &= \frac{1}{x} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dx} \\ &= e^{\int \frac{1}{x} dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\ln(x)} \\ &= x \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= x(3yx + y^2) \\ &= y(3x + y)x \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= x(x^2 + yx) \\ &= x^2(x + y) \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (y(3x + y)x) + (x^2(x + y)) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int y(3x + y) x dx \\ \phi &= \frac{y x^2(2x + y)}{2} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned} \frac{\partial \phi}{\partial y} &= \frac{x^2(2x + y)}{2} + \frac{y x^2}{2} + f'(y) \\ &= x^2(x + y) + f'(y) \end{aligned} \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = x^2(x + y)$. Therefore equation (4) becomes

$$x^2(x + y) = x^2(x + y) + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{y x^2(2x + y)}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{y x^2(2x + y)}{2}$$

Summary

The solution(s) found are the following

$$\frac{y x^2(2x + y)}{2} = c_1 \quad (1)$$

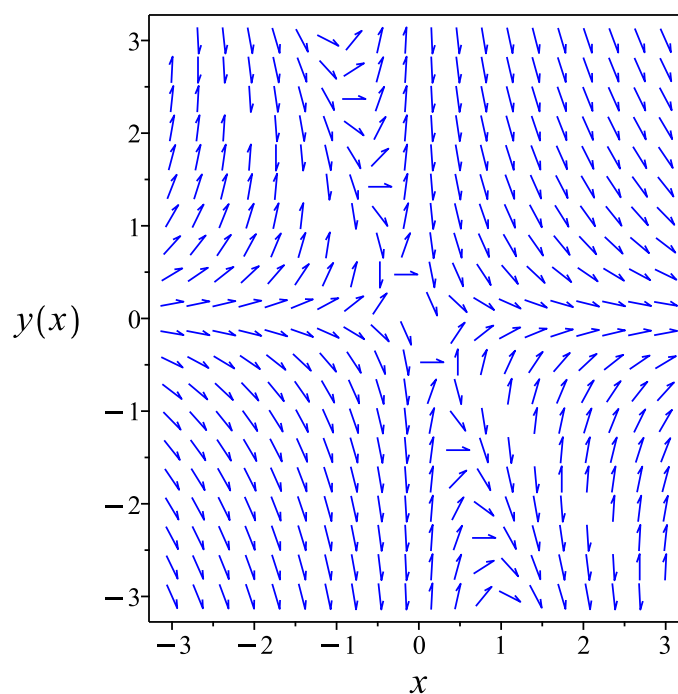


Figure 271: Slope field plot

Verification of solutions

$$\frac{y x^2(2x + y)}{2} = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 59

```
dsolve(3*x*y(x)+y(x)^2+(x^2+x*y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = \frac{-c_1 x^2 - \sqrt{c_1^2 x^4 + 1}}{c_1 x}$$
$$y(x) = \frac{-c_1 x^2 + \sqrt{c_1^2 x^4 + 1}}{c_1 x}$$

✓ Solution by Mathematica

Time used: 0.661 (sec). Leaf size: 93

```
DSolve[3*x*y[x]+y[x]^2+(x^2+x*y[x])*y'[x]== 0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{x^2 + \sqrt{x^4 + e^{2c_1}}}{x}$$
$$y(x) \rightarrow -x + \frac{\sqrt{x^4 + e^{2c_1}}}{x}$$
$$y(x) \rightarrow -\frac{\sqrt{x^4 + x^2}}{x}$$
$$y(x) \rightarrow \frac{\sqrt{x^4}}{x} - x$$

6 Miscellaneous problems, end of chapter 2. Page 133

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6.1 problem 1

6.1.1	Solving as linear ode	1403
6.1.2	Solving as first order ode lie symmetry lookup ode	1405
6.1.3	Solving as exact ode	1409
6.1.4	Maple step by step solution	1414

Internal problem ID [568]

Internal file name [OUTPUT/568_Sunday_June_05_2022_01_44_54_AM_67792014/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$y' - \frac{x^3 - 2y}{x} = 0$$

6.1.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{2}{x}$$
$$q(x) = x^2$$

Hence the ode is

$$y' + \frac{2y}{x} = x^2$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{2}{x} dx} \\ &= x^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) (x^2) \\ \frac{d}{dx}(y x^2) &= (x^2) (x^2) \\ d(y x^2) &= x^4 dx\end{aligned}$$

Integrating gives

$$\begin{aligned}y x^2 &= \int x^4 dx \\ y x^2 &= \frac{x^5}{5} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = x^2$ results in

$$y = \frac{x^3}{5} + \frac{c_1}{x^2}$$

Summary

The solution(s) found are the following

$$y = \frac{x^3}{5} + \frac{c_1}{x^2} \tag{1}$$

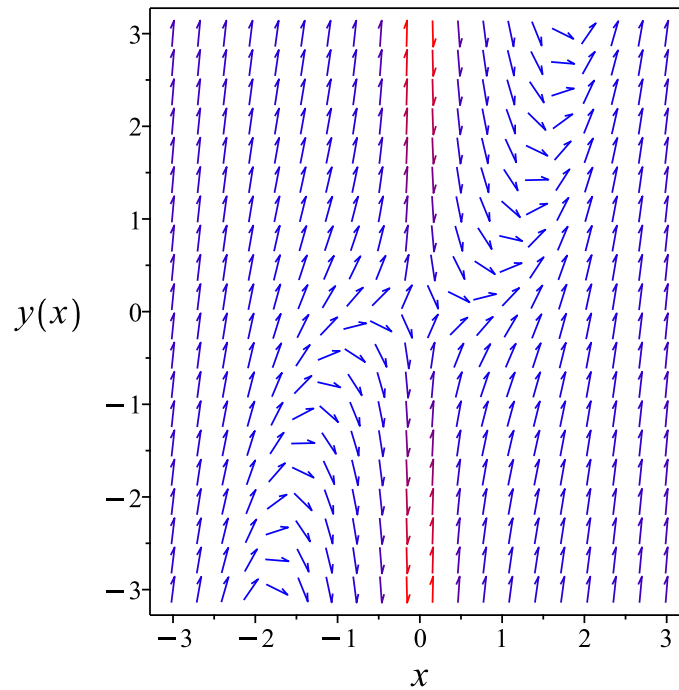


Figure 272: Slope field plot

Verification of solutions

$$y = \frac{x^3}{5} + \frac{c_1}{x^2}$$

Verified OK.

6.1.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{-x^3 + 2y}{x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 260: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int(n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{x^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{x^2}} dy \end{aligned}$$

Which results in

$$S = y x^2$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{-x^3 + 2y}{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= 2yx \\ S_y &= x^2 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = x^4 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R^4$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{R^5}{5} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$x^2 y = \frac{x^5}{5} + c_1$$

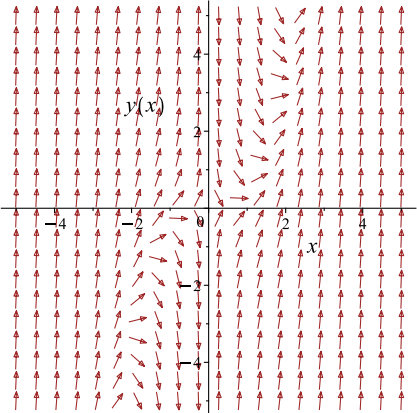
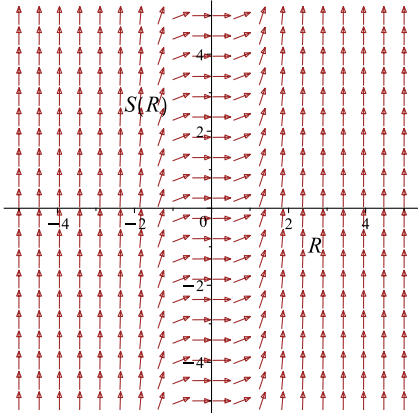
Which simplifies to

$$x^2 y = \frac{x^5}{5} + c_1$$

Which gives

$$y = \frac{x^5 + 5c_1}{5x^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{-x^3+2y}{x}$ 	$R = x$ $S = y x^2$	$\frac{dS}{dR} = R^4$ 

Summary

The solution(s) found are the following

$$y = \frac{x^5 + 5c_1}{5x^2} \quad (1)$$

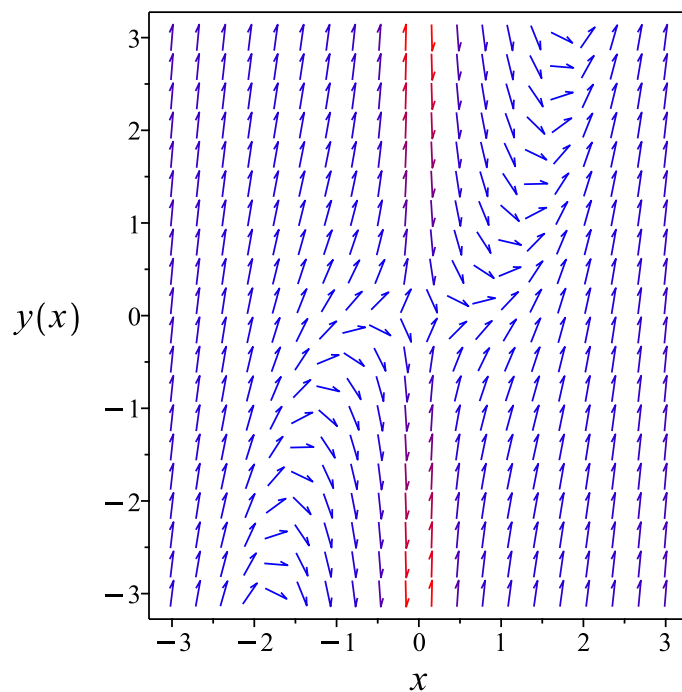


Figure 273: Slope field plot

Verification of solutions

$$y = \frac{x^5 + 5c_1}{5x^2}$$

Verified OK.

6.1.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}dy &= \left(\frac{x^3 - 2y}{x} \right) dx \\ \left(-\frac{x^3 - 2y}{x} \right) dx + dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{x^3 - 2y}{x} \\ N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{x^3 - 2y}{x} \right) \\ &= \frac{2}{x}\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1 \left(\left(\frac{2}{x} \right) - (0) \right) \\ &= \frac{2}{x}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int \frac{2}{x} dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{2\ln(x)} \\ &= x^2\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\overline{M} &= \mu M \\ &= x^2 \left(-\frac{x^3 - 2y}{x} \right) \\ &= -x(x^3 - 2y)\end{aligned}$$

And

$$\begin{aligned}\overline{N} &= \mu N \\ &= x^2(1) \\ &= x^2\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\overline{M} + \overline{N} \frac{dy}{dx} &= 0 \\ (-x(x^3 - 2y)) + (x^2) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x(x^3 - 2y) dx \\ \phi &= -\frac{1}{5}x^5 + yx^2 + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x^2 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = x^2$. Therefore equation (4) becomes

$$x^2 = x^2 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{1}{5}x^5 + yx^2 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{1}{5}x^5 + yx^2$$

The solution becomes

$$y = \frac{x^5 + 5c_1}{5x^2}$$

Summary

The solution(s) found are the following

$$y = \frac{x^5 + 5c_1}{5x^2} \tag{1}$$

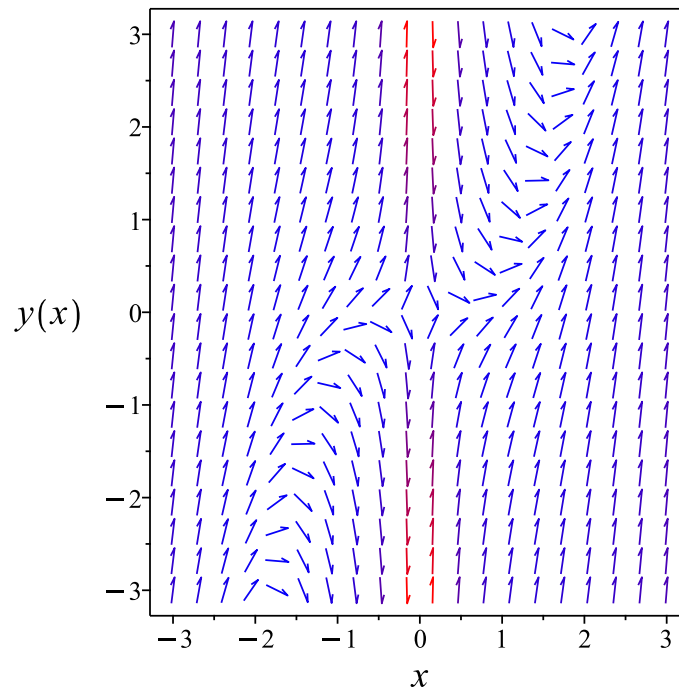


Figure 274: Slope field plot

Verification of solutions

$$y = \frac{x^5 + 5c_1}{5x^2}$$

Verified OK.

6.1.4 Maple step by step solution

Let's solve

$$y' - \frac{x^3 - 2y}{x} = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{2y}{x} + x^2$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{2y}{x} = x^2$$

- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$\mu(x) \left(y' + \frac{2y}{x} \right) = \mu(x) x^2$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dx}(\mu(x)y)$

$$\mu(x) \left(y' + \frac{2y}{x} \right) = \mu'(x)y + \mu(x)y'$$

- Isolate $\mu'(x)$

$$\mu'(x) = \frac{2\mu(x)}{x}$$

- Solve to find the integrating factor

$$\mu(x) = x^2$$

- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x)y) \right) dx = \int \mu(x) x^2 dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int \mu(x) x^2 dx + c_1$$

- Solve for y

$$y = \frac{\int \mu(x) x^2 dx + c_1}{\mu(x)}$$

- Substitute $\mu(x) = x^2$

$$y = \frac{\int x^4 dx + c_1}{x^2}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{x^5}{5} + c_1}{x^2}$$

- Simplify

$$y = \frac{x^5 + 5c_1}{5x^2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 16

```
dsolve(diff(y(x),x) = (x^3-2*y(x))/x,y(x), singsol=all)
```

$$y(x) = \frac{x^5 + 5c_1}{5x^2}$$

✓ Solution by Mathematica

Time used: 0.024 (sec). Leaf size: 19

```
DSolve[y'[x]== (x^3-2*y[x])/x,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{x^3}{5} + \frac{c_1}{x^2}$$

6.2 problem 2

6.2.1	Solving as separable ode	1416
6.2.2	Solving as first order ode lie symmetry lookup ode	1418
6.2.3	Solving as exact ode	1422
6.2.4	Maple step by step solution	1426

Internal problem ID [569]

Internal file name [OUTPUT/569_Sunday_June_05_2022_01_44_55_AM_90905761/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{\cos(x) + 1}{2 - \sin(y)} = 0$$

6.2.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{-\cos(x) - 1}{-2 + \sin(y)}\end{aligned}$$

Where $f(x) = -\cos(x) - 1$ and $g(y) = \frac{1}{-2 + \sin(y)}$. Integrating both sides gives

$$\frac{1}{-2 + \sin(y)} dy = -\cos(x) - 1 dx$$

$$\int \frac{1}{-2+\sin(y)} dy = \int -\cos(x) - 1 dx$$

$$-2y - \cos(y) = -x - \sin(x) + c_1$$

Which results in

$$y = \text{RootOf}(2_Z + \cos(_Z) - x - \sin(x) + c_1)$$

Summary

The solution(s) found are the following

$$y = \text{RootOf}(2_Z + \cos(_Z) - x - \sin(x) + c_1) \quad (1)$$

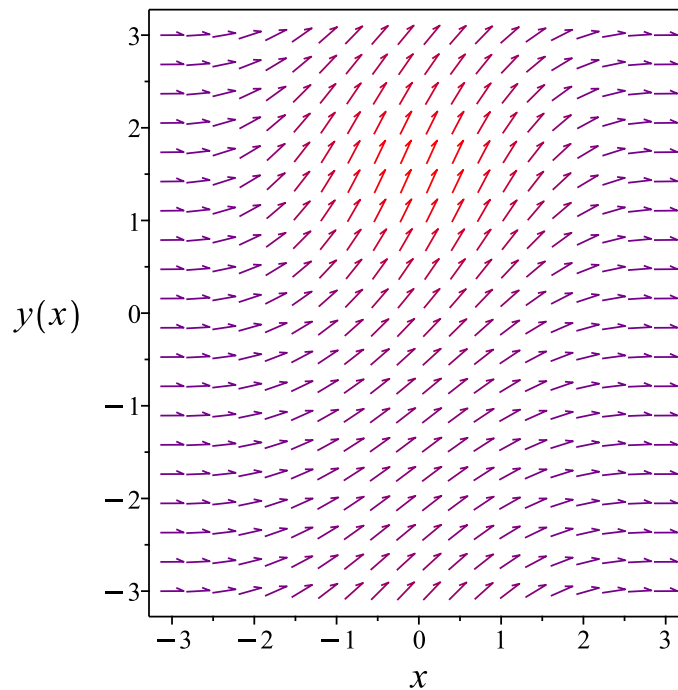


Figure 275: Slope field plot

Verification of solutions

$$y = \text{RootOf}(2_Z + \cos(_Z) - x - \sin(x) + c_1)$$

Verified OK.

6.2.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{\cos(x) + 1}{-2 + \sin(y)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 263: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{-\cos(x) - 1} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{-\cos(x)-1} dx\end{aligned}$$

Which results in

$$S = -x - \sin(x)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{\cos(x) + 1}{-2 + \sin(y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= -\cos(x) - 1 \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -2 + \sin(y) \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -2 + \sin(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\cos(R) - 2R + c_1 \quad (4)$$

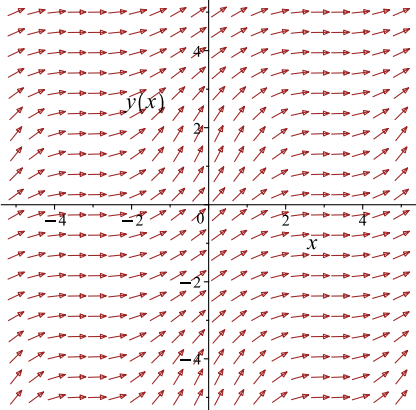
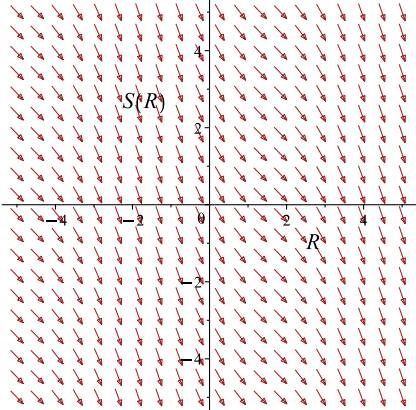
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-x - \sin(x) = -\cos(y) - 2y + c_1$$

Which simplifies to

$$-x - \sin(x) = -\cos(y) - 2y + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{\cos(x)+1}{-2+\sin(y)}$ 	$R = y$ $S = -x - \sin(x)$	$\frac{dS}{dR} = -2 + \sin(R)$ 

Summary

The solution(s) found are the following

$$-x - \sin(x) = -\cos(y) - 2y + c_1 \tag{1}$$

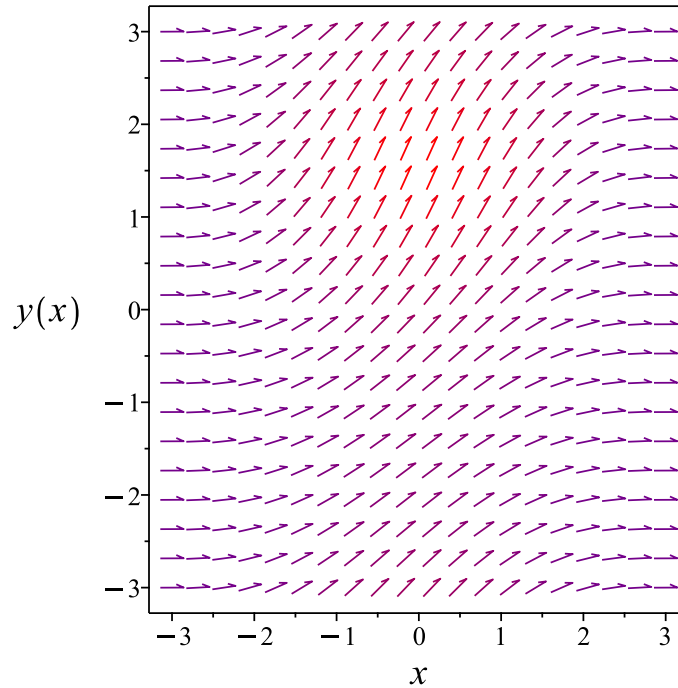


Figure 276: Slope field plot

Verification of solutions

$$-x - \sin(x) = -\cos(y) - 2y + c_1$$

Verified OK.

6.2.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(2 - \sin(y)) dy &= (\cos(x) + 1) dx \\ (-\cos(x) - 1) dx + (2 - \sin(y)) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\cos(x) - 1 \\ N(x, y) &= 2 - \sin(y)\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-\cos(x) - 1) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(2 - \sin(y)) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -\cos(x) - 1 dx$$

$$\phi = -x - \sin(x) + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 2 - \sin(y)$. Therefore equation (4) becomes

$$2 - \sin(y) = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 2 - \sin(y)$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (2 - \sin(y)) dy$$

$$f(y) = 2y + \cos(y) + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x - \sin(x) + 2y + \cos(y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x - \sin(x) + 2y + \cos(y)$$

Summary

The solution(s) found are the following

$$2y + \cos(y) - x - \sin(x) = c_1 \quad (1)$$

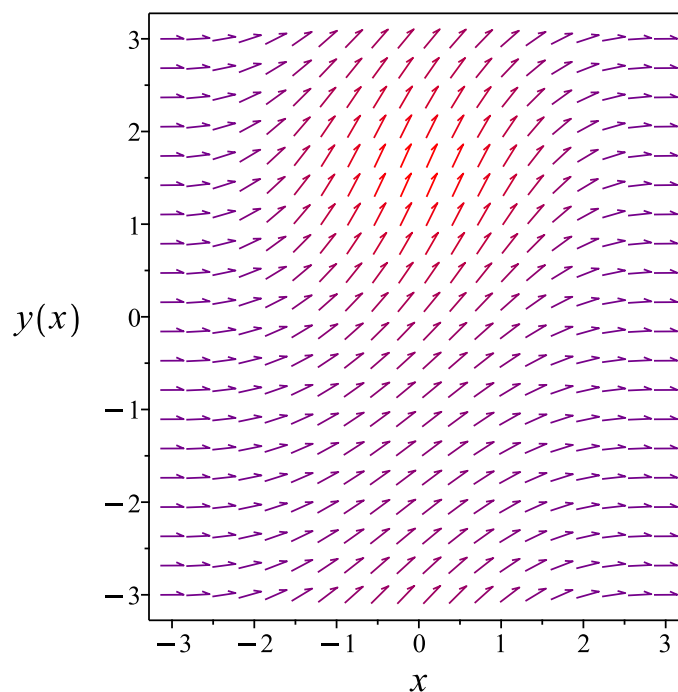


Figure 277: Slope field plot

Verification of solutions

$$2y + \cos(y) - x - \sin(x) = c_1$$

Verified OK.

6.2.4 Maple step by step solution

Let's solve

$$y' - \frac{\cos(x)+1}{2-\sin(y)} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$(2 - \sin(y)) y' = \cos(x) + 1$$

- Integrate both sides with respect to x

$$\int (2 - \sin(y)) y' dx = \int (\cos(x) + 1) dx + c_1$$

- Evaluate integral

$$2y + \cos(y) = x + \sin(x) + c_1$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(diff(y(x),x) = (1+cos(x))/(2-sin(y(x))),y(x), singsol=all)
```

$$x + \sin(x) - 2y(x) - \cos(y(x)) + c_1 = 0$$

✓ Solution by Mathematica

Time used: 0.372 (sec). Leaf size: 27

```
DSolve[y'[x] == (1+Cos[x])/(2-Sin[y[x]]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \text{InverseFunction}[-2\#1 - \cos(\#1)\&][-x - \sin(x) + c_1]$$

6.3 problem 3

6.3.1	Existence and uniqueness analysis	1427
6.3.2	Solving as differentialType ode	1428
6.3.3	Solving as exact ode	1433

Internal problem ID [570]

Internal file name [OUTPUT/570_Sunday_June_05_2022_01_44_56_AM_17223756/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 3.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "differentialType"**

Maple gives the following as the ode type

`[_rational]`

$$y' - \frac{2x + y}{3 - x + 3y^2} = 0$$

With initial conditions

$$[y(0) = 0]$$

6.3.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned} y' &= f(x, y) \\ &= \frac{2x + y}{3y^2 - x + 3} \end{aligned}$$

The x domain of $f(x, y)$ when $y = 0$ is

$$\{x < 3 \vee 3 < x\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $f(x, y)$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 0$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{2x + y}{3y^2 - x + 3} \right) \\ &= \frac{1}{3y^2 - x + 3} - \frac{6(2x + y)y}{(3y^2 - x + 3)^2}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 0$ is

$$\{x < 3 \vee 3 < x\}$$

And the point $x_0 = 0$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 0$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 0$ is inside this domain. Therefore solution exists and is unique.

6.3.2 Solving as differentialType ode

Writing the ode as

$$y' = \frac{2x + y}{3 - x + 3y^2} \quad (1)$$

Which becomes

$$(-3y^2 - 3) dy = (-x) dy + (-2x - y) dx \quad (2)$$

But the RHS is complete differential because

$$(-x) dy + (-2x - y) dx = d(-x^2 - yx)$$

Hence (2) becomes

$$(-3y^2 - 3) dy = d(-x^2 - yx)$$

Integrating both sides gives gives these solutions

$$y = \frac{\left(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324}\right)^{\frac{1}{3}}}{6} - \frac{\left(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324}\right)^{\frac{1}{3}}}{12}$$

$$y = -\frac{\left(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324}\right)^{\frac{1}{3}}}{12} + \frac{\left(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324}\right)^{\frac{1}{3}}}{12}$$

$$y = -\frac{\left(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324}\right)^{\frac{1}{3}}}{12} + \frac{\left(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324}\right)^{\frac{1}{3}}}{12}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{-i\sqrt{3}\left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{2}{3}} - 4i\sqrt{3} - \left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{2}{3}} + 4c_1\left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{1}{3}} + 4}{4\left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{1}{3}}}$$

$$c_1 = \frac{i(2\sqrt{3} - 2i)\sqrt{3} + 4i\sqrt{3} + 2\sqrt{3} - 4 - 2i}{4\sqrt{2\sqrt{3} - 2i}}$$

Substituting c_1 found above in the general solution gives

$$y = \lim_{c_1 \rightarrow \frac{i(2\sqrt{3}-2i)\sqrt{3}+4i\sqrt{3}+2\sqrt{3}-4-2i}{4\sqrt{2\sqrt{3}-2i}}} \left(-\frac{\left(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324}\right)^{\frac{1}{3}}}{12} + \frac{\left(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324}\right)^{\frac{1}{3}}}{12} \right)$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{i\sqrt{3}\left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{2}{3}} + 4i\sqrt{3} - \left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{2}{3}} + 4c_1\left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{1}{3}} + 4}{4\left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{1}{3}}}$$

$$c_1 = -\frac{i(2i + 2\sqrt{3})\sqrt{3} + 4i\sqrt{3} + 4 - 2i - 2\sqrt{3}}{4\sqrt{2i + 2\sqrt{3}}}$$

Substituting c_1 found above in the general solution gives

$$y = \text{Expression too large to display}$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = \frac{\left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{2}{3}} + 2c_1\left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{1}{3}} - 4}{2\left(-4c_1 + 4\sqrt{c_1^2 + 4}\right)^{\frac{1}{3}}}$$

$$c_1 = 0$$

Substituting c_1 found above in the general solution gives

$$y = \frac{\left(108x^2 + 12\sqrt{81x^4 - 12x^3 + 108x^2 - 324x + 324}\right)^{\frac{2}{3}} + 12x - 36}{6\left(108x^2 + 12\sqrt{81x^4 - 12x^3 + 108x^2 - 324x + 324}\right)^{\frac{1}{3}}}$$

Summary

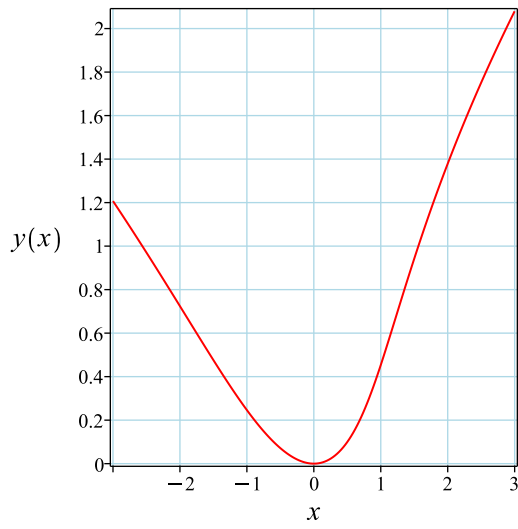
The solution(s) found are the following

$$y = \frac{(108x^2 + 12\sqrt{81x^4 - 12x^3 + 108x^2 - 324x + 324})^{\frac{2}{3}} + 12x - 36}{6(108x^2 + 12\sqrt{81x^4 - 12x^3 + 108x^2 - 324x + 324})^{\frac{1}{3}}} \quad (1)$$

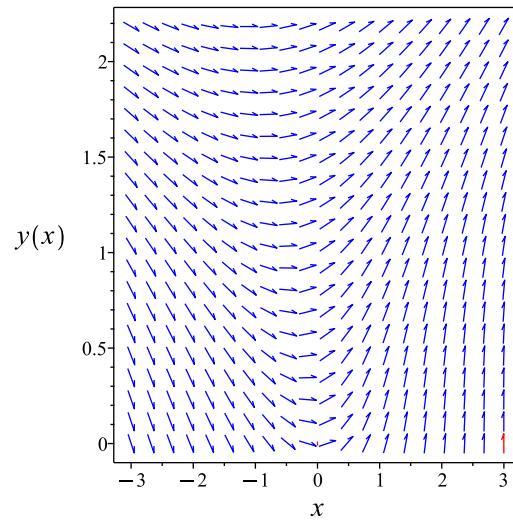
Expression too large to display (2)

y (3)

$$= \lim_{c_1 \rightarrow \frac{i(2\sqrt{3}-2i)\sqrt{3+4i\sqrt{3}+2\sqrt{3}-4-2i}}{4\sqrt{2\sqrt{3}-2i}}} \left(\frac{(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324})^{\frac{2}{3}} + 12x - 36}{12} \right. \\ + \frac{3 - x}{(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324})^{\frac{1}{3}}} \\ \left. + i\sqrt{3} \left(\frac{(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324})^{\frac{1}{3}}}{6} + \frac{-2x + 6}{(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324})^{\frac{1}{3}}} \right) \right) + c_1$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{(108x^2 + 12\sqrt{81x^4 - 12x^3 + 108x^2 - 324x + 324})^{\frac{2}{3}} + 12x - 36}{6(108x^2 + 12\sqrt{81x^4 - 12x^3 + 108x^2 - 324x + 324})^{\frac{1}{3}}}$$

Verified OK.

Expression too large to display

Warning, solution could not be verified

y

$$= \lim_{c_1 \rightarrow \frac{i(2\sqrt{3}-2i)\sqrt{3+4i}\sqrt{3+2\sqrt{3}-4-2i}}{4\sqrt{2\sqrt{3}-2i}}} \left(- \frac{(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324})^{\frac{2}{3}}}{12} \right. \\ \left. + \frac{3-x}{(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324})^{\frac{1}{3}}} \right. \\ \left. + i\sqrt{3} \left(\frac{(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324})^{\frac{1}{3}}}{6} + \frac{-2x+6}{(108x^2 - 108c_1 + 12\sqrt{81x^4 - 162c_1x^2 - 12x^3 + 81c_1^2 + 108x^2 - 324x + 324})^{\frac{1}{3}}} \right) \right. \\ \left. - \frac{\dots}{2} \right) + c_1$$

Warning, solution could not be verified

6.3.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the

ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (3y^2 - x + 3) dy &= (2x + y) dx \\ (-2x - y) dx + (3y^2 - x + 3) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -2x - y \\ N(x, y) &= 3y^2 - x + 3 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-2x - y) \\ &= -1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(3y^2 - x + 3) \\ &= -1\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -2x - y dx \\ \phi &= -x(x + y) + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = -x + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 3y^2 - x + 3$. Therefore equation (4) becomes

$$3y^2 - x + 3 = -x + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 3y^2 + 3$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (3y^2 + 3) dy$$
$$f(y) = y^3 + 3y + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x(x + y) + y^3 + 3y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x(x + y) + y^3 + 3y$$

Initial conditions are used to solve for c_1 . Substituting $x = 0$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = c_1$$

$$c_1 = 0$$

Substituting c_1 found above in the general solution gives

$$-x(x + y) + y^3 + 3y = 0$$

Summary

The solution(s) found are the following

$$y^3 + (3 - x)y - x^2 = 0 \tag{1}$$

Verification of solutions

$$y^3 + (3 - x)y - x^2 = 0$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

✓ Solution by Maple

Time used: 0.078 (sec). Leaf size: 75

```
dsolve([diff(y(x),x) = (2*x+y(x))/(3-x+3*y(x)^2),y(0) = 0],y(x), singsol=all)
```

$$y(x) = \frac{(108x^2 + 12\sqrt{81x^4 - 12x^3 + 108x^2 - 324x + 324})^{\frac{2}{3}} + 12x - 36}{6(108x^2 + 12\sqrt{81x^4 - 12x^3 + 108x^2 - 324x + 324})^{\frac{1}{3}}}$$

✓ Solution by Mathematica

Time used: 5.408 (sec). Leaf size: 114

```
DSolve[{y'[x] == (2*x+y[x])/(3-x+3*y[x]^2),y[0]==0},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{\sqrt[3]{2}(\sqrt{3}\sqrt{27x^4 - 4x^3 + 36x^2 - 108x + 108} - 9x^2)^{2/3} + 2\sqrt[3]{3}x - 6\sqrt[3]{3}}{6^{2/3}\sqrt[3]{\sqrt{3}\sqrt{27x^4 - 4x^3 + 36x^2 - 108x + 108} - 9x^2}}$$

6.4 problem 4

6.4.1	Solving as separable ode	1438
6.4.2	Solving as linear ode	1440
6.4.3	Solving as first order ode lie symmetry lookup ode	1441
6.4.4	Solving as exact ode	1445
6.4.5	Maple step by step solution	1449

Internal problem ID [571]

Internal file name [OUTPUT/571_Sunday_June_05_2022_01_44_58_AM_16995004/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 4.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "linear", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - y + 2yx = -6x + 3$$

6.4.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= (2x - 1)(-y - 3)\end{aligned}$$

Where $f(x) = 2x - 1$ and $g(y) = -y - 3$. Integrating both sides gives

$$\begin{aligned}\frac{1}{-y - 3} dy &= 2x - 1 dx \\ \int \frac{1}{-y - 3} dy &= \int 2x - 1 dx \\ -\ln(3 + y) &= x^2 + c_1 - x\end{aligned}$$

Raising both side to exponential gives

$$\frac{1}{3+y} = e^{x^2+c_1-x}$$

Which simplifies to

$$\frac{1}{3+y} = c_2 e^{x^2-x}$$

Summary

The solution(s) found are the following

$$y = -\frac{(3c_2 e^{x^2+c_1-x} - 1) e^{-x^2-c_1+x}}{c_2} \quad (1)$$

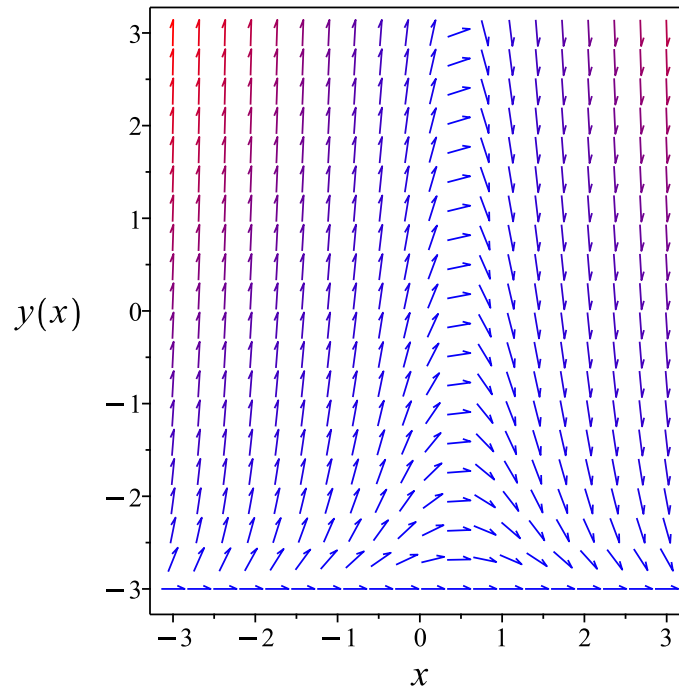


Figure 279: Slope field plot

Verification of solutions

$$y = -\frac{(3c_2 e^{x^2+c_1-x} - 1) e^{-x^2-c_1+x}}{c_2}$$

Verified OK.

6.4.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$\begin{aligned}p(x) &= 2x - 1 \\q(x) &= -6x + 3\end{aligned}$$

Hence the ode is

$$y' + y(2x - 1) = -6x + 3$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int(2x-1)dx} \\&= e^{x^2-x}\end{aligned}$$

Which simplifies to

$$\mu = e^{x(x-1)}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu)(-6x + 3) \\ \frac{d}{dx}(e^{x(x-1)}y) &= (e^{x(x-1)})(-6x + 3) \\ d(e^{x(x-1)}y) &= ((-6x + 3)e^{x(x-1)}) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{x(x-1)}y &= \int (-6x + 3)e^{x(x-1)} dx \\ e^{x(x-1)}y &= -3e^{x(x-1)} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{x(x-1)}$ results in

$$y = -3e^{-x(x-1)}e^{x(x-1)} + c_1e^{-x(x-1)}$$

which simplifies to

$$y = -3 + c_1e^{-x(x-1)}$$

Summary

The solution(s) found are the following

$$y = -3 + c_1 e^{-x(x-1)} \quad (1)$$

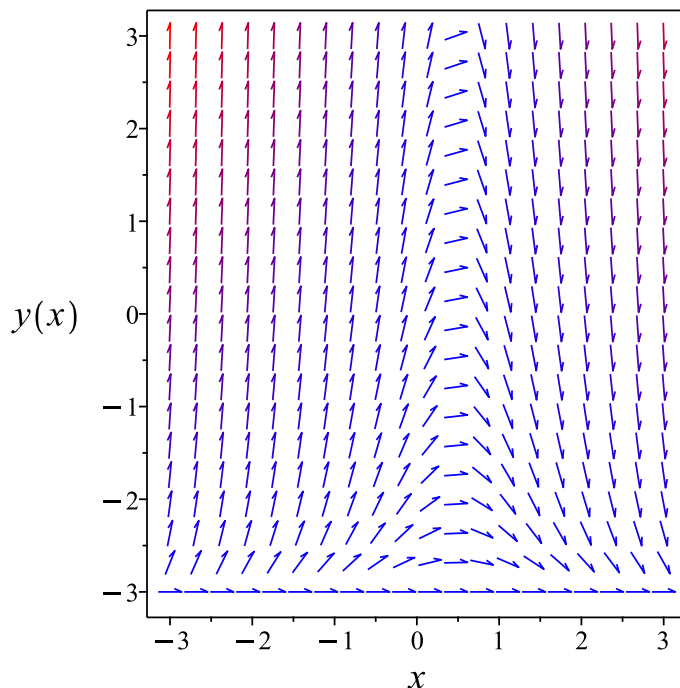


Figure 280: Slope field plot

Verification of solutions

$$y = -3 + c_1 e^{-x(x-1)}$$

Verified OK.

6.4.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$\begin{aligned} y' &= -2yx - 6x + y + 3 \\ y' &= \omega(x, y) \end{aligned}$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 266: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int(n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= e^{-x^2+xy}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-x^2+x}} dy \end{aligned}$$

Which results in

$$S = e^{x^2-x} y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -2yx - 6x + y + 3$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= (2x - 1) e^{x(x-1)} y \\ S_y &= e^{x(x-1)} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = (-6x + 3) e^{x(x-1)} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = (-6R + 3) e^{R(R-1)}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -3e^{R(R-1)} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$e^{x(x-1)}y = -3e^{x(x-1)} + c_1$$

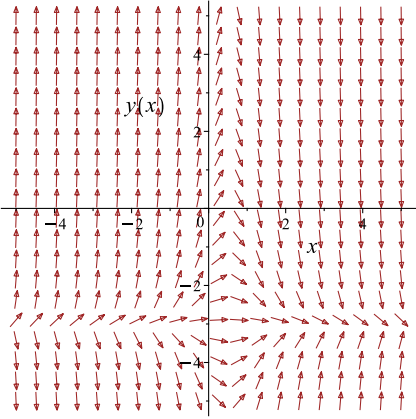
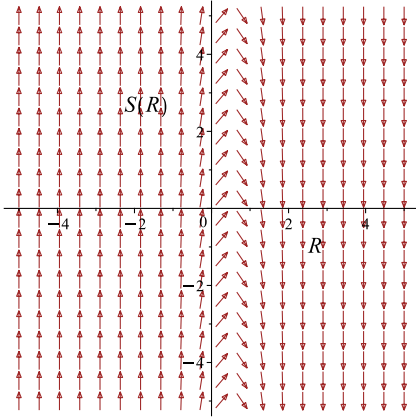
Which simplifies to

$$(3 + y) e^{x(x-1)} - c_1 = 0$$

Which gives

$$y = -(3e^{x(x-1)} - c_1) e^{-x(x-1)}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -2yx - 6x + y + 3$ 	$R = x$ $S = e^{x(x-1)}y$	$\frac{dS}{dR} = (-6R + 3) e^{R(R-1)}$ 

Summary

The solution(s) found are the following

$$y = -(3e^{x(x-1)} - c_1) e^{-x(x-1)} \quad (1)$$

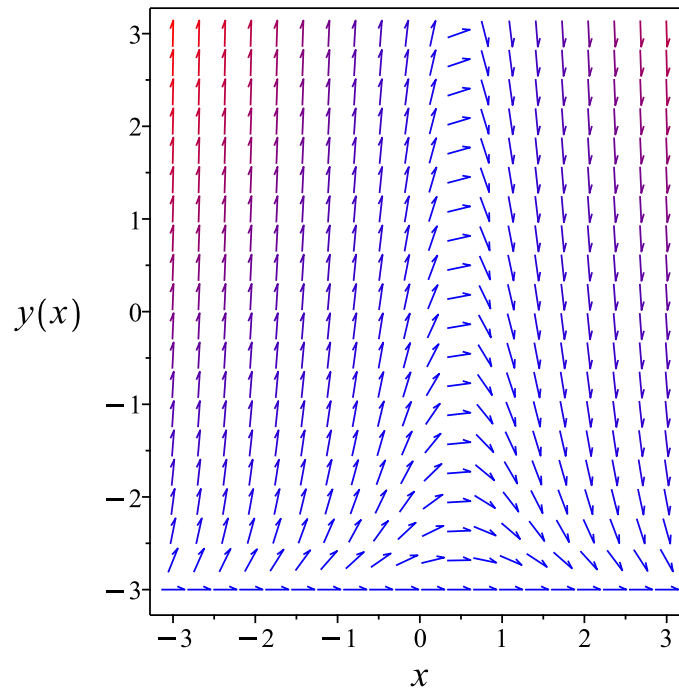


Figure 281: Slope field plot

Verification of solutions

$$y = -(3e^{x(x-1)} - c_1) e^{-x(x-1)}$$

Verified OK.

6.4.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(\frac{1}{-y-3}\right) dy &= (2x-1) dx \\ (1-2x) dx + \left(\frac{1}{-y-3}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 1 - 2x \\ N(x, y) &= \frac{1}{-y-3}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(1-2x) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{1}{-y-3} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 1 - 2x dx \\ \phi &= -x^2 + x + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{-y-3}$. Therefore equation (4) becomes

$$\frac{1}{-y-3} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{1}{3+y}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(-\frac{1}{3+y} \right) dy \\ f(y) &= -\ln(3+y) + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x^2 + x - \ln(3 + y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x^2 + x - \ln(3 + y)$$

The solution becomes

$$y = e^{-x^2 - c_1 + x} - 3$$

Summary

The solution(s) found are the following

$$y = e^{-x^2 - c_1 + x} - 3 \quad (1)$$

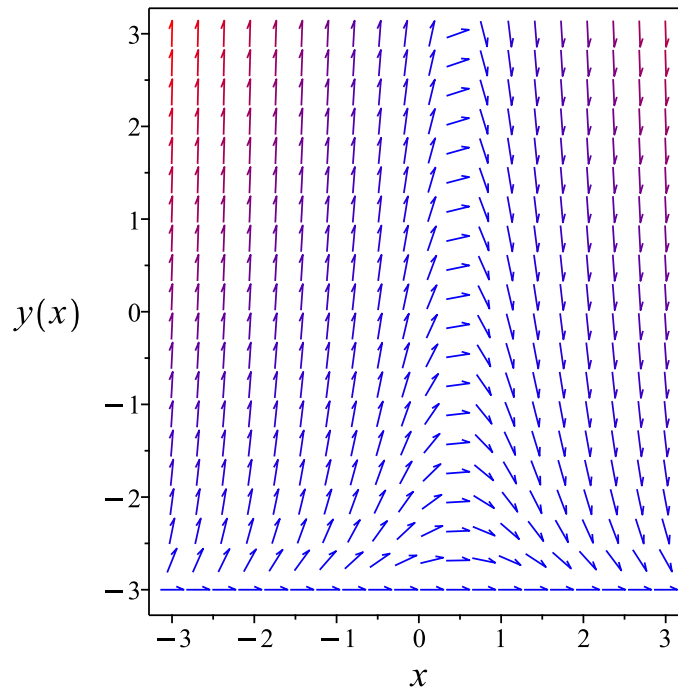


Figure 282: Slope field plot

Verification of solutions

$$y = e^{-x^2 - c_1 + x} - 3$$

Verified OK.

6.4.5 Maple step by step solution

Let's solve

$$y' - y + 2yx = -6x + 3$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{3+y} = 1 - 2x$$

- Integrate both sides with respect to x

$$\int \frac{y'}{3+y} dx = \int (1 - 2x) dx + c_1$$

- Evaluate integral

$$\ln(3 + y) = -x^2 + c_1 + x$$

- Solve for y

$$y = e^{-x^2+c_1+x} - 3$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 15

```
dsolve(diff(y(x),x) = 3-6*x+y(x)-2*x*y(x),y(x), singsol=all)
```

$$y(x) = -3 + e^{-x(x-1)}c_1$$

✓ Solution by Mathematica

Time used: 0.048 (sec). Leaf size: 24

```
DSolve[y'[x] == 3-6*x+y[x]-2*x*y[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -3 + c_1 e^{x-x^2}$$

$$y(x) \rightarrow -3$$

6.5 problem 5

6.5.1 Solving as exact ode 1451

Internal problem ID [572]

Internal file name [OUTPUT/572_Sunday_June_05_2022_01_44_59_AM_84609071/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 5.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact"

Maple gives the following as the ode type

```
[_rational, [_Abel, `2nd type`, `class B`]]
```

$$y' - \frac{-1 - 2yx - y^2}{x^2 + 2yx} = 0$$

6.5.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (x(x + 2y)) dy &= (-2yx - y^2 - 1) dx \\ (2yx + y^2 + 1) dx + (x(x + 2y)) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 2yx + y^2 + 1 \\ N(x, y) &= x(x + 2y) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (2yx + y^2 + 1) \\ &= 2x + 2y \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (x(x + 2y)) \\ &= 2x + 2y \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 2yx + y^2 + 1 dx \\ \phi &= yx^2 + xy^2 + x + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned}\frac{\partial \phi}{\partial y} &= x^2 + 2yx + f'(y) \\ &= x(x + 2y) + f'(y)\end{aligned}\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = x(x + 2y)$. Therefore equation (4) becomes

$$x(x + 2y) = x(x + 2y) + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = yx^2 + xy^2 + x + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = yx^2 + xy^2 + x$$

Summary

The solution(s) found are the following

$$x^2y + xy^2 + x = c_1\tag{1}$$

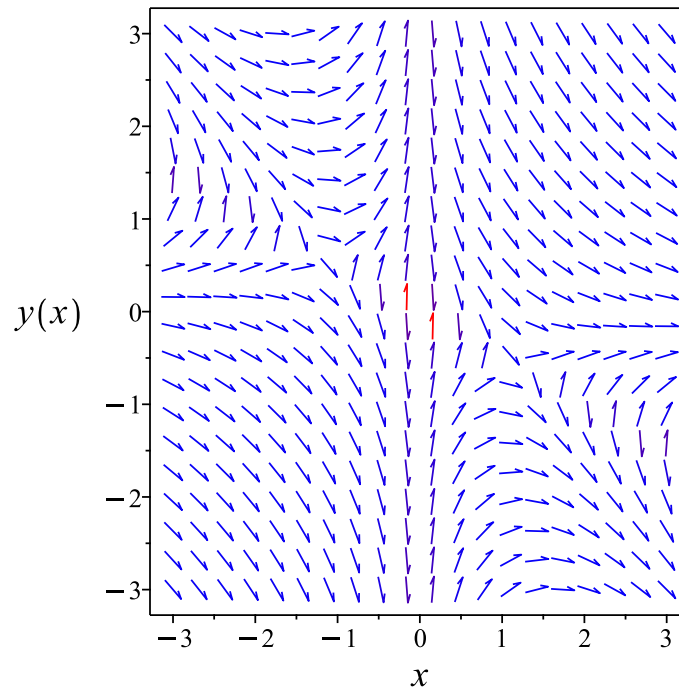


Figure 283: Slope field plot

Verification of solutions

$$x^2y + xy^2 + x = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 59

```
dsolve(diff(y(x),x) = (-1-2*x*y(x)-y(x)^2)/(x^2+2*x*y(x)),y(x), singsol=all)
```

$$y(x) = \frac{-x^2 + \sqrt{x(x^3 - 4c_1 - 4x)}}{2x}$$
$$y(x) = \frac{-x^2 - \sqrt{x(x^3 - 4c_1 - 4x)}}{2x}$$

✓ Solution by Mathematica

Time used: 0.502 (sec). Leaf size: 67

```
DSolve[y'[x] == (-1-2*x*y[x]-y[x]^2)/(x^2+2*x*y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{x^2 + \sqrt{x(x^3 - 4x + 4c_1)}}{2x}$$
$$y(x) \rightarrow \frac{-x^2 + \sqrt{x(x^3 - 4x + 4c_1)}}{2x}$$

6.6 problem 6

6.6.1	Existence and uniqueness analysis	1456
6.6.2	Solving as linear ode	1457
6.6.3	Solving as first order ode lie symmetry lookup ode	1459
6.6.4	Solving as exact ode	1463
6.6.5	Maple step by step solution	1467

Internal problem ID [573]

Internal file name [OUTPUT/573_Sunday_June_05_2022_01_45_00_AM_94379850/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 6.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$yx + y'x + y = 1$$

With initial conditions

$$[y(1) = 0]$$

6.6.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{-x-1}{x}$$
$$q(x) = \frac{1}{x}$$

Hence the ode is

$$y' - \frac{(-x-1)y}{x} = \frac{1}{x}$$

The domain of $p(x) = -\frac{-x-1}{x}$ is

$$\{x < 0 \vee 0 < x\}$$

And the point $x_0 = 1$ is inside this domain. The domain of $q(x) = \frac{1}{x}$ is

$$\{x < 0 \vee 0 < x\}$$

And the point $x_0 = 1$ is also inside this domain. Hence solution exists and is unique.

6.6.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{-x-1}{x} dx} \\ &= e^{x+\ln(x)}\end{aligned}$$

Which simplifies to

$$\mu = x e^x$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) \left(\frac{1}{x}\right) \\ \frac{d}{dx}(x e^x y) &= (x e^x) \left(\frac{1}{x}\right) \\ d(x e^x y) &= e^x dx\end{aligned}$$

Integrating gives

$$\begin{aligned}x e^x y &= \int e^x dx \\ x e^x y &= e^x + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = x e^x$ results in

$$y = \frac{e^{-x} e^x}{x} + \frac{c_1 e^{-x}}{x}$$

which simplifies to

$$y = \frac{c_1 e^{-x} + 1}{x}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = e^{-1}c_1 + 1$$

$$c_1 = -e$$

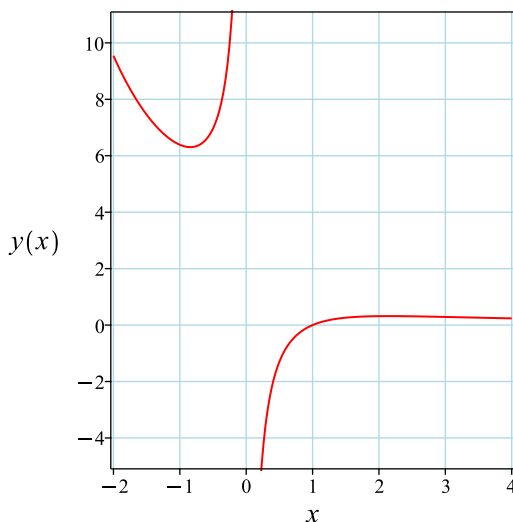
Substituting c_1 found above in the general solution gives

$$y = \frac{-e^{1-x} + 1}{x}$$

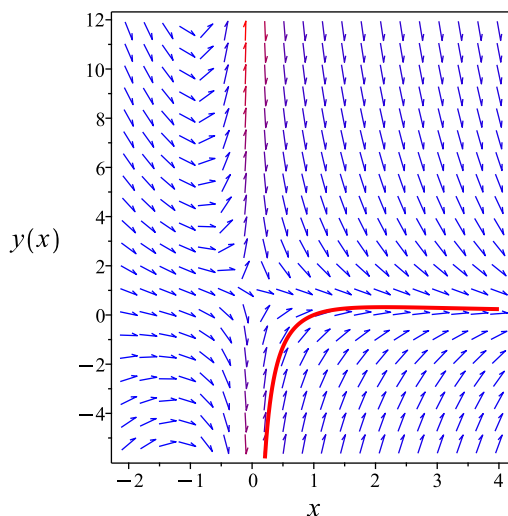
Summary

The solution(s) found are the following

$$y = \frac{-e^{1-x} + 1}{x} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{-e^{1-x} + 1}{x}$$

Verified OK.

6.6.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{yx + y - 1}{x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 269: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= e^{-x-\ln(x)}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-x-\ln(x)}} dy\end{aligned}$$

Which results in

$$S = x e^x y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{yx + y - 1}{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\ R_y &= 0 \\ S_x &= e^x y (x + 1) \\ S_y &= x e^x\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^x \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = e^R + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$y e^x x = e^x + c_1$$

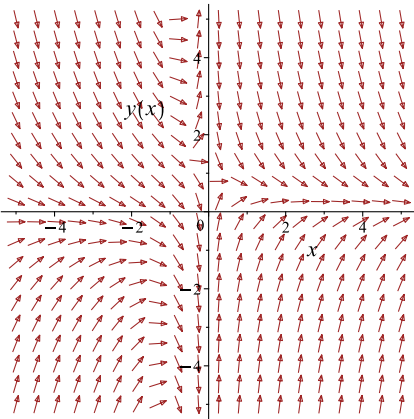
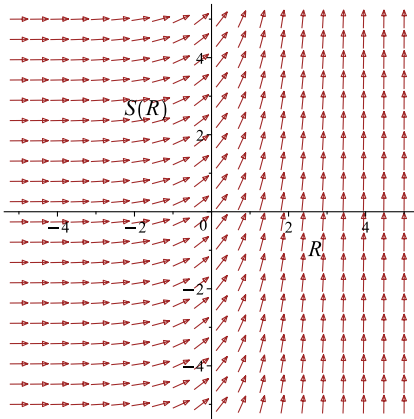
Which simplifies to

$$y e^x x = e^x + c_1$$

Which gives

$$y = \frac{(e^x + c_1) e^{-x}}{x}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{yx+y-1}{x}$ 	$R = x$ $S = x e^x y$	$\frac{dS}{dR} = e^R$ 

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = e^{-1}c_1 + 1$$

$$c_1 = -e$$

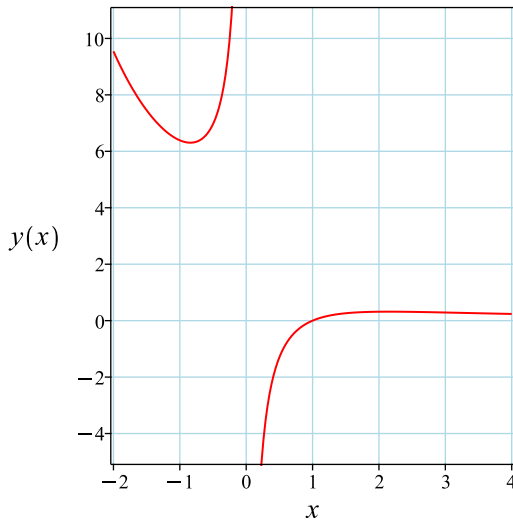
Substituting c_1 found above in the general solution gives

$$y = \frac{e^{-x}e^x - e e^{-x}}{x}$$

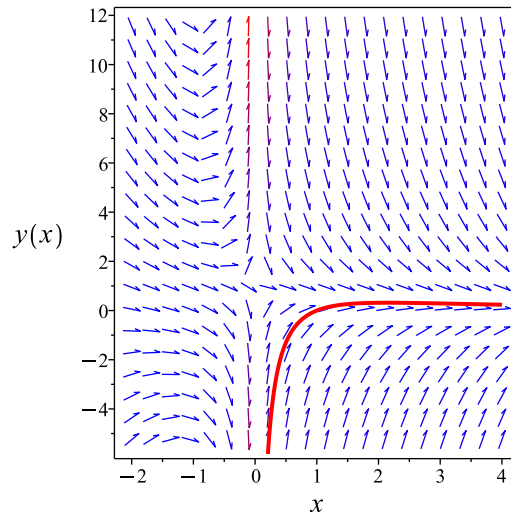
Summary

The solution(s) found are the following

$$y = \frac{e^{-x}e^x - e e^{-x}}{x} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^{-x}e^x - ee^{-x}}{x}$$

Verified OK.

6.6.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (x) dy &= (-yx - y + 1) dx \\ (yx + y - 1) dx + (x) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= yx + y - 1 \\ N(x, y) &= x \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(yx + y - 1) \\ &= x + 1 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x) \\ &= 1 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{x} ((x + 1) - (1)) \\ &= 1 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int 1 dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^x \\ &= e^x\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= e^x(yx + y - 1) \\ &= e^x(yx + y - 1)\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^x(x) \\ &= x e^x\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (e^x(yx + y - 1)) + (x e^x) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int e^x(yx + y - 1) dx \\ \phi &= (yx - 1) e^x + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x e^x + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = x e^x$. Therefore equation (4) becomes

$$x e^x = x e^x + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = (yx - 1) e^x + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = (yx - 1) e^x$$

The solution becomes

$$y = \frac{(e^x + c_1) e^{-x}}{x}$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = 0$ in the above solution gives an equation to solve for the constant of integration.

$$0 = e^{-1} c_1 + 1$$

$$c_1 = -e$$

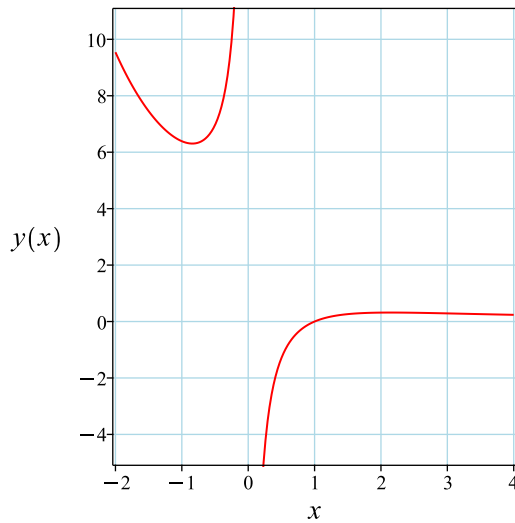
Substituting c_1 found above in the general solution gives

$$y = \frac{e^{-x} e^x - e e^{-x}}{x}$$

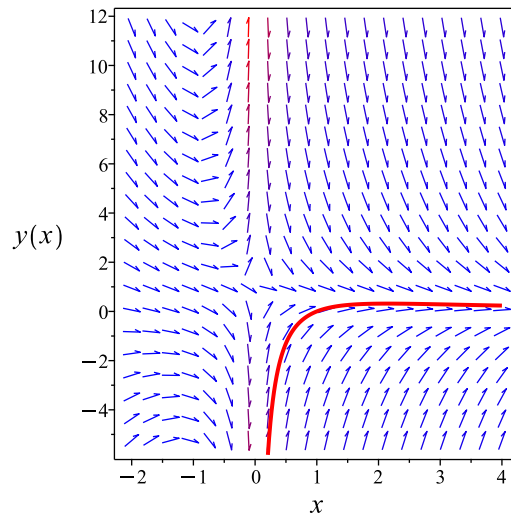
Summary

The solution(s) found are the following

$$y = \frac{e^{-x}e^x - e e^{-x}}{x} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^{-x}e^x - e e^{-x}}{x}$$

Verified OK.

6.6.5 Maple step by step solution

Let's solve

$$[yx + y'x + y = 1, y(1) = 0]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{(x+1)y}{x} + \frac{1}{x}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{(x+1)y}{x} = \frac{1}{x}$$

- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$\mu(x) \left(y' + \frac{(x+1)y}{x} \right) = \frac{\mu(x)}{x}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dx}(\mu(x)y)$

$$\mu(x) \left(y' + \frac{(x+1)y}{x} \right) = \mu'(x)y + \mu(x)y'$$

- Isolate $\mu'(x)$

$$\mu'(x) = \frac{\mu(x)(x+1)}{x}$$

- Solve to find the integrating factor

$$\mu(x) = x e^x$$

- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x)y) \right) dx = \int \frac{\mu(x)}{x} dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int \frac{\mu(x)}{x} dx + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(x)}{x} dx + c_1}{\mu(x)}$$

- Substitute $\mu(x) = x e^x$

$$y = \frac{\int e^x dx + c_1}{x e^x}$$

- Evaluate the integrals on the rhs

$$y = \frac{e^x + c_1}{e^x x}$$

- Simplify

$$y = \frac{c_1 e^{-x} + 1}{x}$$

- Use initial condition $y(1) = 0$

$$0 = e^{-1}c_1 + 1$$

- Solve for c_1

$$c_1 = -\frac{1}{e^{-1}}$$

- Substitute $c_1 = -\frac{1}{e^{-1}}$ into general solution and simplify

$$y = \frac{-e^{1-x} + 1}{x}$$

- Solution to the IVP

$$y = \frac{-e^{1-x} + 1}{x}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 18

```
dsolve([x*y(x)+x*diff(y(x),x) = 1-y(x),y(1) = 0],y(x), singsol=all)
```

$$y(x) = \frac{1 - e^{1-x}}{x}$$

✓ Solution by Mathematica

Time used: 0.03 (sec). Leaf size: 20

```
DSolve[{x*y[x]+x*y'[x] == 1-y[x],y[1]==0},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1 - e^{1-x}}{x}$$

6.7 problem 7

6.7.1	Solving as separable ode	1470
6.7.2	Solving as differentialType ode	1475
6.7.3	Solving as first order ode lie symmetry lookup ode	1479
6.7.4	Solving as exact ode	1483
6.7.5	Maple step by step solution	1487

Internal problem ID [574]

Internal file name [OUTPUT/574_Sunday_June_05_2022_01_45_01_AM_69403732/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 7.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "differential-Type", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$y' - \frac{4x^3 + 1}{y(2 + 3y)} = 0$$

6.7.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{4x^3 + 1}{y(2 + 3y)}\end{aligned}$$

Where $f(x) = 4x^3 + 1$ and $g(y) = \frac{1}{y(2+3y)}$. Integrating both sides gives

$$\frac{1}{\frac{1}{y(2+3y)}} dy = 4x^3 + 1 dx$$

$$\int \frac{1}{\frac{1}{y(2+3y)}} dy = \int 4x^3 + 1 dx$$

$$y^3 + y^2 = x^4 + c_1 + x$$

Which results in

y

$$= \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{6} - \frac{1}{3}$$

$$+ \frac{3\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{1}{3}$$

$y =$

$$\frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{1}{3}$$

$$+ \frac{3\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{1}{3}$$

$$+ i\sqrt{3} \left(\frac{\left(\frac{-8+108x^4+108c_1+108x+12\sqrt{81x^8+162c_1x^4+162x^5-12x^4+81c_1^2+162c_1x+81x^2-12c_1-12x}}{6}\right)^{\frac{1}{3}}}{3\left(-8+108x^4+108c_1+108x+12\sqrt{81x^8+162c_1x^4+162x^5-12x^4+81c_1^2+162c_1x+81x^2-12c_1-12x}\right)} \right) - \frac{1}{3}$$

$y =$

$$\frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{1}{3}$$

$$+ \frac{3\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{1}{3}$$

$$+ i\sqrt{3} \left(\frac{\left(\frac{-8+108x^4+108c_1+108x+12\sqrt{81x^8+162c_1x^4+162x^5-12x^4+81c_1^2+162c_1x+81x^2-12c_1-12x}}{6}\right)^{\frac{1}{3}}}{3\left(-8+108x^4+108c_1+108x+12\sqrt{81x^8+162c_1x^4+162x^5-12x^4+81c_1^2+162c_1x+81x^2-12c_1-12x}\right)} \right) - \frac{1}{3}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{6} \quad (1)$$

$$y = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{6} - \frac{1}{3} \quad (2)$$

$$y = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{6} - \frac{12}{1} \quad (3)$$

$$y = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{6} - \frac{1}{3} + i\sqrt{3} \left(\frac{\left(\frac{-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}}{6}\right)^{\frac{1}{3}}}{3(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x})} \right) - \frac{12}{1} \quad (3)$$

$$y = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{6} - \frac{1}{3} + i\sqrt{3} \left(\frac{\left(\frac{-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}}{6}\right)^{\frac{1}{3}}}{3(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x})} \right) - \frac{12}{1} \quad (3)$$

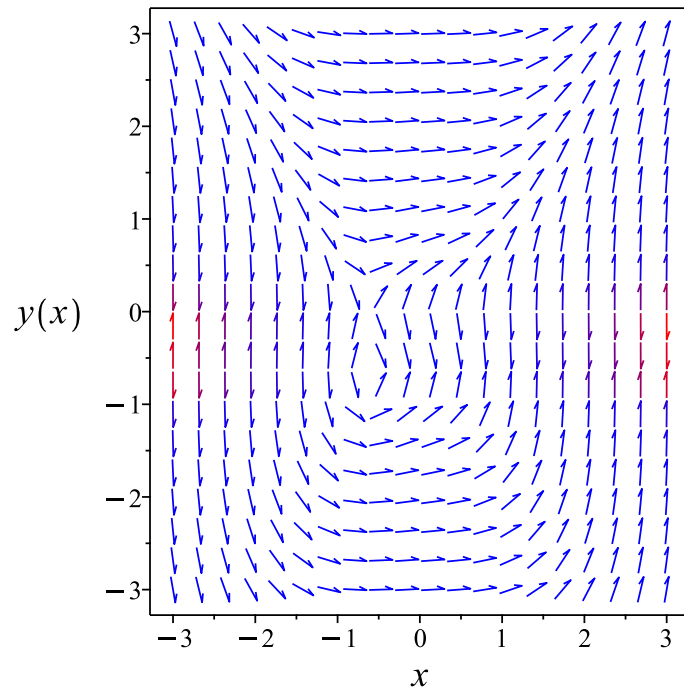


Figure 287: Slope field plot

Verification of solutions

y

$$= \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{6} - \frac{1}{3} + \frac{3\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{2}$$

Verified OK.

$y =$

$$- \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{1}{3} + \frac{3\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{2} + i\sqrt{3} \left(\frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)^{\frac{1}{3}}}{6} - \frac{3\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{2} \right)$$

Verified OK.

$y =$

$$- \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{1}{3} + \frac{3\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{2} + i\sqrt{3} \left(\frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)^{\frac{1}{3}}}{6} - \frac{3\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{2} \right)$$

Verified OK.

6.7.2 Solving as differentialType ode

Writing the ode as

$$y' = \frac{4x^3 + 1}{y(2 + 3y)} \quad (1)$$

Which becomes

$$(3y^2 + 2y) dy = (4x^3 + 1) dx \quad (2)$$

But the RHS is complete differential because

$$(4x^3 + 1) dx = d(x^4 + x)$$

Hence (2) becomes

$$(3y^2 + 2y) dy = d(x^4 + x)$$

Integrating both sides gives gives these solutions

$$y = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1}\right)}{6}$$

$$y = -\frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1}\right)}{12}$$

$$y = -\frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1}\right)}{12}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{6} + \frac{1}{3} + c_1 \tag{1}$$

$$y = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{1}{3} + c_1 \tag{2}$$

$$y = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{1}{3} + i\sqrt{3} \left(\frac{\left(\frac{-8+108x^4+108c_1+108x+12\sqrt{81x^8+162c_1x^4+162x^5-12x^4+81c_1^2+162c_1x+81x^2-12c_1-12x}}{6}\right)^{\frac{1}{3}}}{3(-8+108x^4+108c_1+108x+12\sqrt{81x^8+162c_1x^4+162x^5-12x^4+81c_1^2+162c_1x+81x^2-12c_1-12x})} \right) + c_1 \tag{3}$$

$$y = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{1}{3} + i\sqrt{3} \left(\frac{\left(\frac{-8+108x^4+108c_1+108x+12\sqrt{81x^8+162c_1x^4+162x^5-12x^4+81c_1^2+162c_1x+81x^2-12c_1-12x}}{6}\right)^{\frac{1}{3}}}{3(-8+108x^4+108c_1+108x+12\sqrt{81x^8+162c_1x^4+162x^5-12x^4+81c_1^2+162c_1x+81x^2-12c_1-12x})} \right) + c_1$$

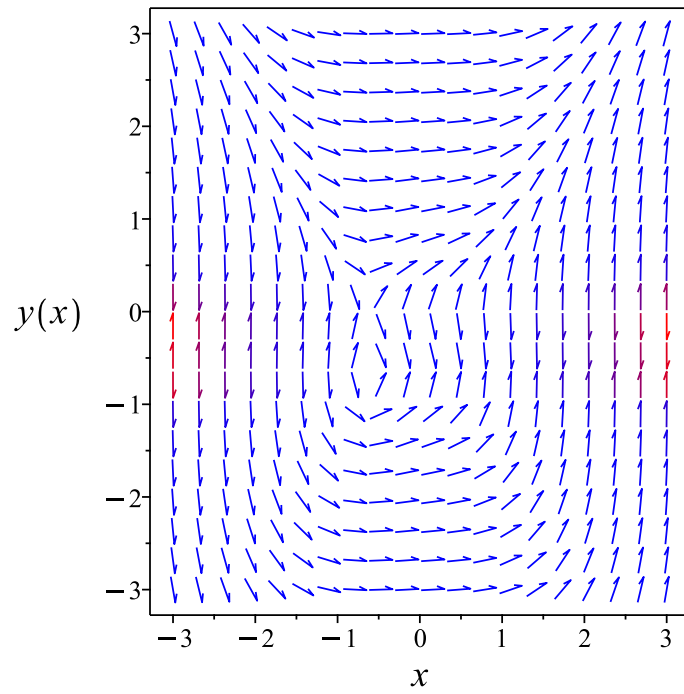


Figure 288: Slope field plot

Verification of solutions

y

$$= \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{6} + \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{2} - \frac{1}{3} + c_1$$

Verified OK.

$y =$

$$- \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{1} - \frac{3\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{3} - \frac{1}{3} + i\sqrt{3} \left(\frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)^{\frac{1}{3}}}{6} - \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)^{\frac{1}{3}}}{3} \right) + c_1$$

Verified OK.

$y =$

$$- \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{12} - \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{1} - \frac{3\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)}{3} - \frac{1}{3} + i\sqrt{3} \left(\frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)^{\frac{1}{3}}}{6} - \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81x^8 + 162c_1x^4 + 162x^5 - 12x^4 + 81c_1^2 + 162c_1x + 81x^2 - 12c_1 - 12x}\right)^{\frac{1}{3}}}{3} \right) + c_1$$

Verified OK.

6.7.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{4x^3 + 1}{y(2 + 3y)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 272: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{4x^3 + 1} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{4x^3+1}} dx\end{aligned}$$

Which results in

$$S = x^4 + x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{4x^3 + 1}{y(2 + 3y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= 4x^3 + 1 \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 3y^2 + 2y \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 3R^2 + 2R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R^3 + R^2 + c_1 \quad (4)$$

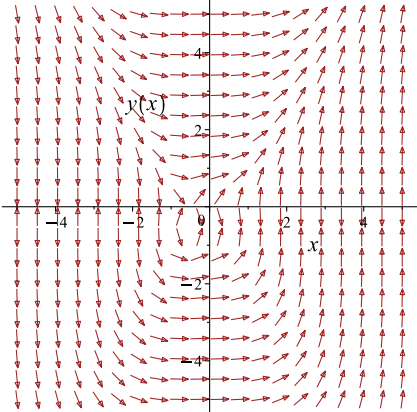
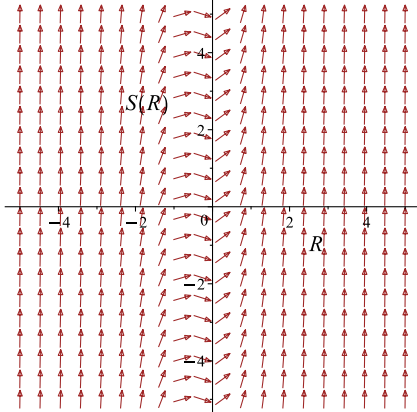
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$x^4 + x = y^3 + y^2 + c_1$$

Which simplifies to

$$x^4 + x = y^3 + y^2 + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{4x^3+1}{y(2+3y)}$ 	$R = y$ $S = x^4 + x$	$\frac{dS}{dR} = 3R^2 + 2R$ 

Summary

The solution(s) found are the following

$$x^4 + x = y^3 + y^2 + c_1 \tag{1}$$

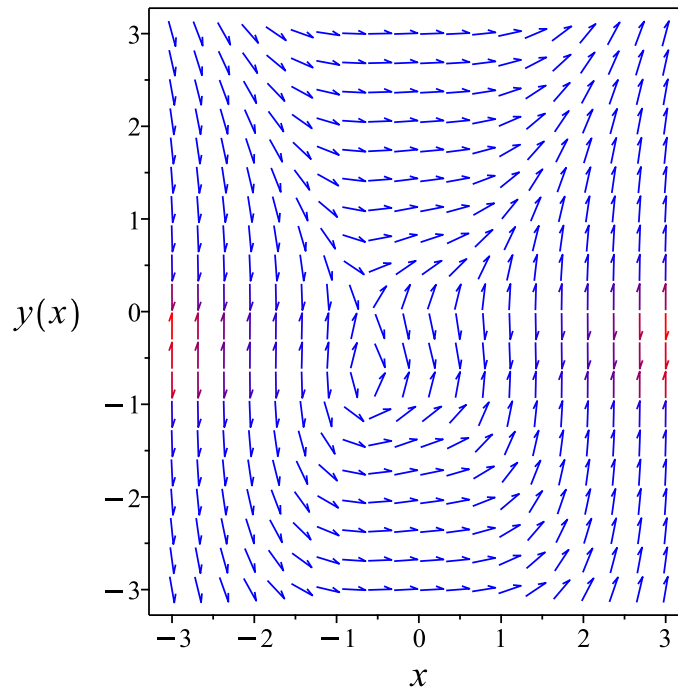


Figure 289: Slope field plot

Verification of solutions

$$x^4 + x = y^3 + y^2 + c_1$$

Verified OK.

6.7.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(y(2 + 3y)) dy &= (4x^3 + 1) dx \\ (-4x^3 - 1) dx + (y(2 + 3y)) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -4x^3 - 1 \\ N(x, y) &= y(2 + 3y)\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-4x^3 - 1) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(y(2 + 3y)) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -4x^3 - 1 dx \\ \phi &= -x^4 - x + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = y(2 + 3y)$. Therefore equation (4) becomes

$$y(2 + 3y) = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = y(2 + 3y)$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int (3y^2 + 2y) dy \\ f(y) &= y^3 + y^2 + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x^4 + y^3 + y^2 - x + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x^4 + y^3 + y^2 - x$$

Summary

The solution(s) found are the following

$$y^3 + y^2 - x^4 - x = c_1 \tag{1}$$

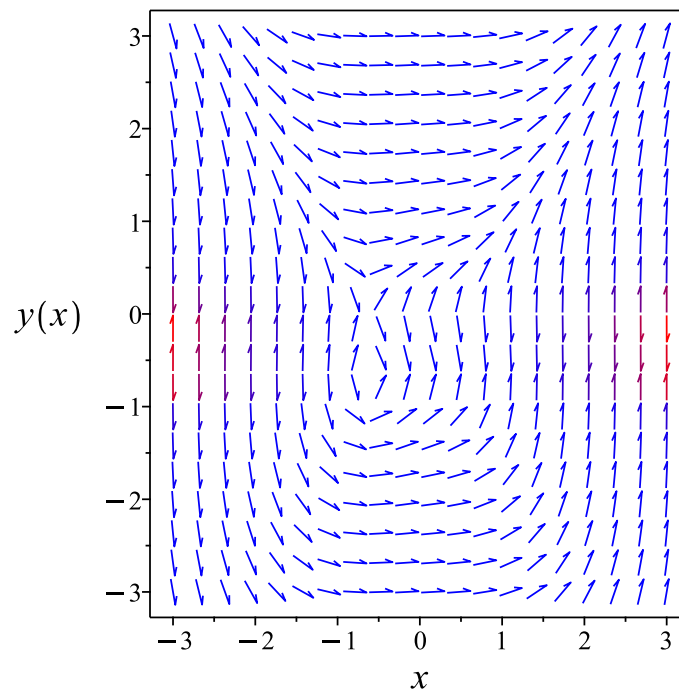


Figure 290: Slope field plot

Verification of solutions

$$y^3 + y^2 - x^4 - x = c_1$$

Verified OK.

6.7.5 Maple step by step solution

Let's solve

$$y' - \frac{4x^3+1}{y(2+3y)} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Separate variables

$$y'y(2+3y) = 4x^3 + 1$$

- Integrate both sides with respect to x

$$\int y'y(2+3y) dx = \int (4x^3 + 1) dx + c_1$$

- Evaluate integral

$$y^3 + y^2 = x^4 + c_1 + x$$

- Solve for y

$$y = \frac{\left(-8+108x^4+108c_1+108x+12\sqrt{81x^8+162c_1x^4+162x^5-12x^4+81c_1^2+162c_1x+81x^2-12c_1-12x}\right)^{\frac{1}{3}}}{6} + \frac{\sqrt[3]{-8+108x^4+108c_1+108x+12\sqrt{81x^8+162c_1x^4+162x^5-12x^4+81c_1^2+162c_1x+81x^2-12c_1-12x}}}{3}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```


✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 382

```
dsolve(diff(y(x),x) = (4*x^3+1)/(y(x)*(2+3*y(x))),y(x), singsol=all)
```

$$y(x) = \frac{\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81} \sqrt{(x^4 + c_1 + x) \left(x^4 + c_1 + x - \frac{4}{27}\right)}\right)^{\frac{2}{3}} - 2\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81} \sqrt{(x^4 + c_1 + x) \left(x^4 + c_1 + x - \frac{4}{27}\right)}\right)^{\frac{2}{3}}}{6\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81} \sqrt{(x^4 + c_1 + x) \left(x^4 + c_1 + x - \frac{4}{27}\right)}\right)^{\frac{2}{3}}}$$

$$y(x) = \frac{(1 + i\sqrt{3})\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81} \sqrt{(x^4 + c_1 + x) \left(x^4 + c_1 + x - \frac{4}{27}\right)}\right)^{\frac{2}{3}} - 4i\sqrt{3} + 4\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81} \sqrt{(x^4 + c_1 + x) \left(x^4 + c_1 + x - \frac{4}{27}\right)}\right)^{\frac{2}{3}}}{12\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81} \sqrt{(x^4 + c_1 + x) \left(x^4 + c_1 + x - \frac{4}{27}\right)}\right)^{\frac{2}{3}}}$$

$$y(x) = \frac{(i\sqrt{3} - 1)\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81} \sqrt{(x^4 + c_1 + x) \left(x^4 + c_1 + x - \frac{4}{27}\right)}\right)^{\frac{2}{3}} - 4i\sqrt{3} - 4\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81} \sqrt{(x^4 + c_1 + x) \left(x^4 + c_1 + x - \frac{4}{27}\right)}\right)^{\frac{2}{3}}}{12\left(-8 + 108x^4 + 108c_1 + 108x + 12\sqrt{81} \sqrt{(x^4 + c_1 + x) \left(x^4 + c_1 + x - \frac{4}{27}\right)}\right)^{\frac{2}{3}}}$$

✓ Solution by Mathematica

Time used: 4.502 (sec). Leaf size: 356

`DSolve[y'[x]== (4*x^3+1)/(y[x]*(2+3*y[x])),y[x],x,IncludeSingularSolutions -> True]`

$$y(x) \rightarrow \frac{1}{6} \left(2^{2/3} \sqrt[3]{27x^4 + \sqrt{-4 + (27x^4 + 27x - 2 + 27c_1)^2 + 27x - 2 + 27c_1}} \right. \\ \left. + \frac{2\sqrt[3]{2}}{\sqrt[3]{27x^4 + \sqrt{-4 + (27x^4 + 27x - 2 + 27c_1)^2 + 27x - 2 + 27c_1}}} - 2 \right)$$

$$y(x) \rightarrow \frac{1}{12} \left(i 2^{2/3} (\sqrt{3} + i) \sqrt[3]{27x^4 + \sqrt{-4 + (27x^4 + 27x - 2 + 27c_1)^2 + 27x - 2 + 27c_1}} \right. \\ \left. - \frac{2\sqrt[3]{2}(1 + i\sqrt{3})}{\sqrt[3]{27x^4 + \sqrt{-4 + (27x^4 + 27x - 2 + 27c_1)^2 + 27x - 2 + 27c_1}}} - 4 \right)$$

$$y(x) \rightarrow \frac{1}{12} \left(-2^{2/3} (1 + i\sqrt{3}) \sqrt[3]{27x^4 + \sqrt{-4 + (27x^4 + 27x - 2 + 27c_1)^2 + 27x - 2 + 27c_1}} \right. \\ \left. + \frac{2i\sqrt[3]{2}(\sqrt{3} + i)}{\sqrt[3]{27x^4 + \sqrt{-4 + (27x^4 + 27x - 2 + 27c_1)^2 + 27x - 2 + 27c_1}}} - 4 \right)$$

6.8 problem 8

6.8.1	Existence and uniqueness analysis	1490
6.8.2	Solving as linear ode	1491
6.8.3	Solving as first order ode lie symmetry lookup ode	1493
6.8.4	Solving as exact ode	1497
6.8.5	Maple step by step solution	1501

Internal problem ID [575]

Internal file name [OUTPUT/575_Sunday_June_05_2022_01_45_02_AM_24783477/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 8.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "linear", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$2y + y'x = \frac{\sin(x)}{x}$$

With initial conditions

$$[y(2) = 1]$$

6.8.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{2}{x}$$
$$q(x) = \frac{\sin(x)}{x^2}$$

Hence the ode is

$$y' + \frac{2y}{x} = \frac{\sin(x)}{x^2}$$

The domain of $p(x) = \frac{2}{x}$ is

$$\{x < 0 \vee 0 < x\}$$

And the point $x_0 = 2$ is inside this domain. The domain of $q(x) = \frac{\sin(x)}{x^2}$ is

$$\{x < 0 \vee 0 < x\}$$

And the point $x_0 = 2$ is also inside this domain. Hence solution exists and is unique.

6.8.2 Solving as linear ode

Entering Linear first order ODE solver. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{2}{x} dx} \\ &= x^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) \left(\frac{\sin(x)}{x^2} \right) \\ \frac{d}{dx}(y x^2) &= (x^2) \left(\frac{\sin(x)}{x^2} \right) \\ d(y x^2) &= \sin(x) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}y x^2 &= \int \sin(x) dx \\ y x^2 &= -\cos(x) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = x^2$ results in

$$y = -\frac{\cos(x)}{x^2} + \frac{c_1}{x^2}$$

which simplifies to

$$y = \frac{-\cos(x) + c_1}{x^2}$$

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{\cos(2)}{4} + \frac{c_1}{4}$$

$$c_1 = \cos(2) + 4$$

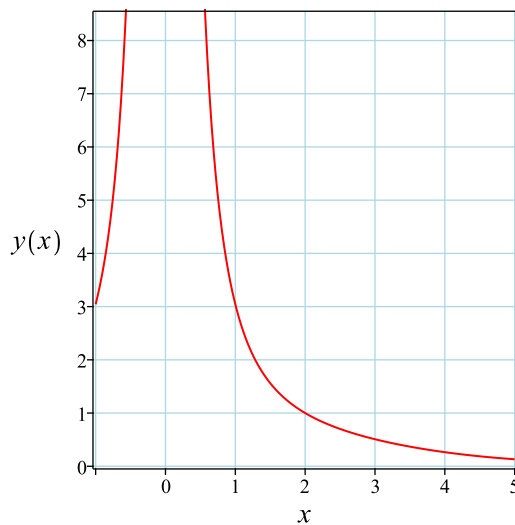
Substituting c_1 found above in the general solution gives

$$y = -\frac{\cos(x) - 4 - \cos(2)}{x^2}$$

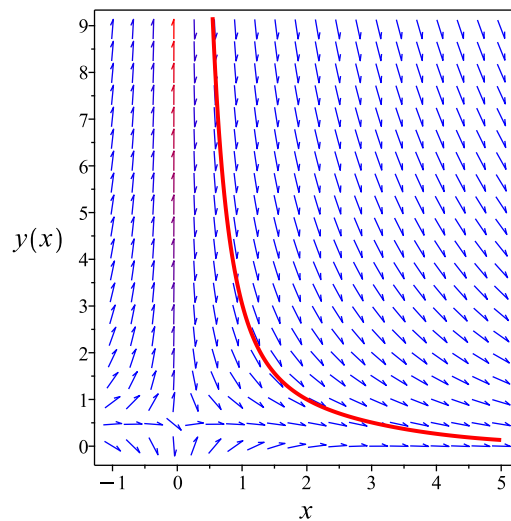
Summary

The solution(s) found are the following

$$y = -\frac{\cos(x) - 4 - \cos(2)}{x^2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{\cos(x) - 4 - \cos(2)}{x^2}$$

Verified OK.

6.8.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-2yx + \sin(x)}{x^2}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 275: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{x^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{x^2}} dy\end{aligned}$$

Which results in

$$S = y x^2$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{-2yx + \sin(x)}{x^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= 2yx \\S_y &= x^2\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \sin(x) \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \sin(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\cos(R) + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$x^2y = -\cos(x) + c_1$$

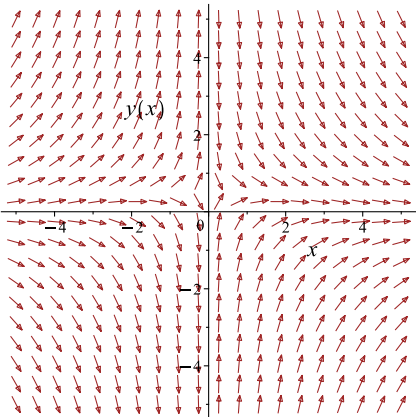
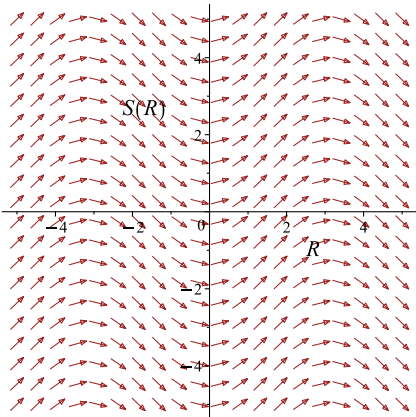
Which simplifies to

$$x^2y = -\cos(x) + c_1$$

Which gives

$$y = -\frac{\cos(x) - c_1}{x^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{-2yx + \sin(x)}{x^2}$ 	$R = x$ $S = yx^2$	$\frac{dS}{dR} = \sin(R)$ 

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{\cos(2)}{4} + \frac{c_1}{4}$$

$$c_1 = \cos(2) + 4$$

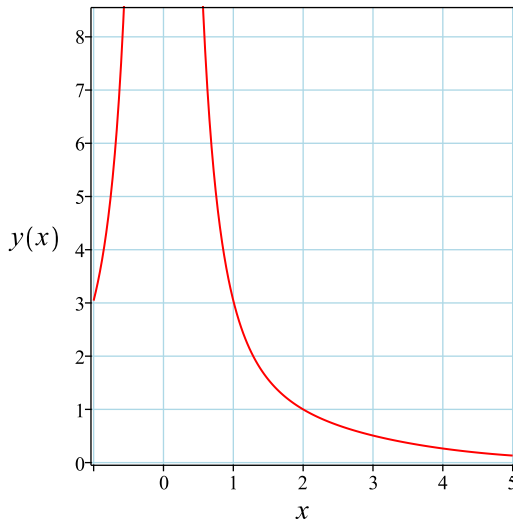
Substituting c_1 found above in the general solution gives

$$y = -\frac{\cos(x) - 4 - \cos(2)}{x^2}$$

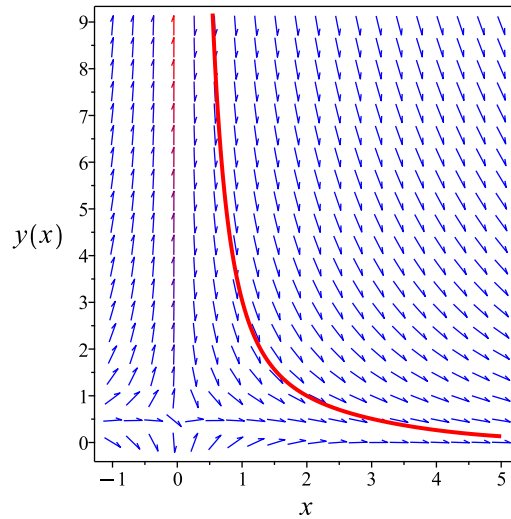
Summary

The solution(s) found are the following

$$y = -\frac{\cos(x) - 4 - \cos(2)}{x^2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{\cos(x) - 4 - \cos(2)}{x^2}$$

Verified OK.

6.8.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (x^2) dy &= (-2yx + \sin(x)) dx \\ (2yx - \sin(x)) dx + (x^2) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 2yx - \sin(x) \\ N(x, y) &= x^2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2yx - \sin(x)) \\ &= 2x \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x^2) \\ &= 2x \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 2yx - \sin(x) dx \\ \phi &= yx^2 + \cos(x) + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x^2 + f'(y)\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = x^2$. Therefore equation (4) becomes

$$x^2 = x^2 + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = yx^2 + \cos(x) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = yx^2 + \cos(x)$$

The solution becomes

$$y = -\frac{\cos(x) - c_1}{x^2}$$

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = -\frac{\cos(2)}{4} + \frac{c_1}{4}$$

$$c_1 = \cos(2) + 4$$

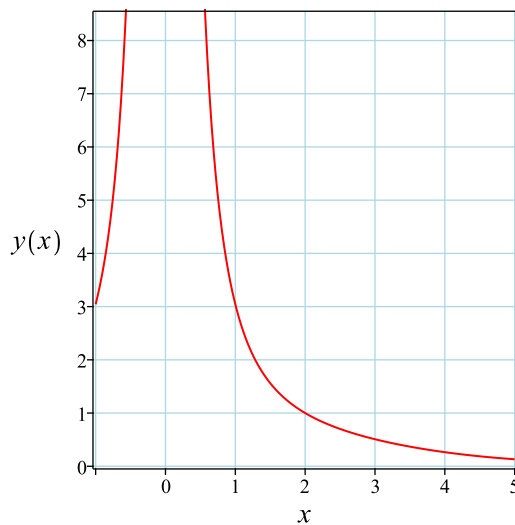
Substituting c_1 found above in the general solution gives

$$y = -\frac{\cos(x) - 4 - \cos(2)}{x^2}$$

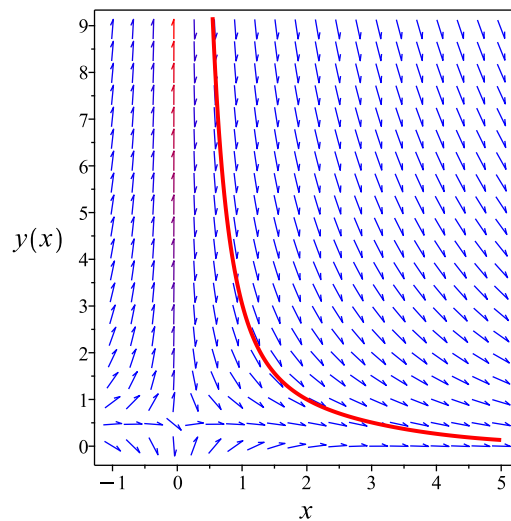
Summary

The solution(s) found are the following

$$y = -\frac{\cos(x) - 4 - \cos(2)}{x^2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{\cos(x) - 4 - \cos(2)}{x^2}$$

Verified OK.

6.8.5 Maple step by step solution

Let's solve

$$\left[2y + y'x = \frac{\sin(x)}{x}, y(2) = 1 \right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{2y}{x} + \frac{\sin(x)}{x^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{2y}{x} = \frac{\sin(x)}{x^2}$$

- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$\mu(x) \left(y' + \frac{2y}{x} \right) = \frac{\mu(x)\sin(x)}{x^2}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dx}(\mu(x)y)$

$$\mu(x) \left(y' + \frac{2y}{x} \right) = \mu'(x)y + \mu(x)y'$$

- Isolate $\mu'(x)$

$$\mu'(x) = \frac{2\mu(x)}{x}$$

- Solve to find the integrating factor

$$\mu(x) = x^2$$

- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x)y) \right) dx = \int \frac{\mu(x)\sin(x)}{x^2} dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int \frac{\mu(x)\sin(x)}{x^2} dx + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(x)\sin(x)}{x^2} dx + c_1}{\mu(x)}$$

- Substitute $\mu(x) = x^2$

$$y = \frac{\int \sin(x) dx + c_1}{x^2}$$

- Evaluate the integrals on the rhs

$$y = \frac{-\cos(x) + c_1}{x^2}$$

- Use initial condition $y(2) = 1$

$$1 = -\frac{\cos(2)}{4} + \frac{c_1}{4}$$

- Solve for c_1

$$c_1 = \cos(2) + 4$$

- Substitute $c_1 = \cos(2) + 4$ into general solution and simplify

$$y = \frac{-\cos(x) + 4 + \cos(2)}{x^2}$$

- Solution to the IVP

$$y = \frac{-\cos(x) + 4 + \cos(2)}{x^2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 16

```
dsolve([2*y(x)+x*diff(y(x),x) = sin(x)/x,y(2) = 1],y(x), singsol=all)
```

$$y(x) = \frac{-\cos(x) + 4 + \cos(2)}{x^2}$$

✓ Solution by Mathematica

Time used: 0.039 (sec). Leaf size: 17

```
DSolve[{2*y[x]+x*y'[x] == Sin[x]/x,y[2]==1},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{-\cos(x) + 4 + \cos(2)}{x^2}$$

6.9 problem 9

- 6.9.1 Solving as differentialType ode 1503
6.9.2 Solving as exact ode 1505

Internal problem ID [576]

Internal file name [OUTPUT/576_Sunday_June_05_2022_01_45_03_AM_51375781/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 9.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "differentialType"

Maple gives the following as the ode type

```
[_rational, [_1st_order, `_with_symmetry_[F(x),G(x)]`], [_Abel,  
`2nd type`, `class A`]]
```

$$y' - \frac{-1 - 2yx}{x^2 + 2y} = 0$$

6.9.1 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-1 - 2yx}{x^2 + 2y} \tag{1}$$

Which becomes

$$(2y) dy = (-x^2) dy + (-2yx - 1) dx \tag{2}$$

But the RHS is complete differential because

$$(-x^2) dy + (-2yx - 1) dx = d(-y x^2 - x)$$

Hence (2) becomes

$$(2y) dy = d(-y x^2 - x)$$

Integrating both sides gives gives these solutions

$$y = -\frac{x^2}{2} + \frac{\sqrt{x^4 + 4c_1 - 4x}}{2} + c_1$$

$$y = -\frac{x^2}{2} - \frac{\sqrt{x^4 + 4c_1 - 4x}}{2} + c_1$$

Summary

The solution(s) found are the following

$$y = -\frac{x^2}{2} + \frac{\sqrt{x^4 + 4c_1 - 4x}}{2} + c_1 \quad (1)$$

$$y = -\frac{x^2}{2} - \frac{\sqrt{x^4 + 4c_1 - 4x}}{2} + c_1 \quad (2)$$

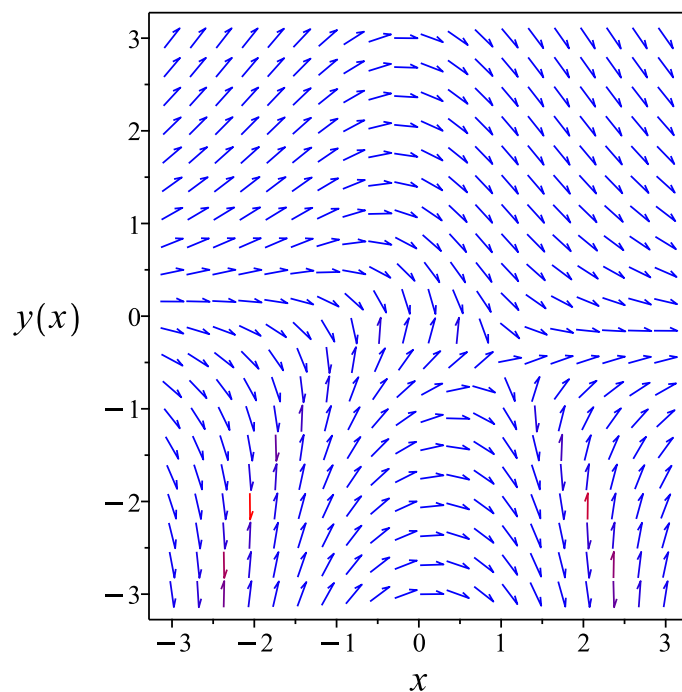


Figure 294: Slope field plot

Verification of solutions

$$y = -\frac{x^2}{2} + \frac{\sqrt{x^4 + 4c_1 - 4x}}{2} + c_1$$

Verified OK.

$$y = -\frac{x^2}{2} - \frac{\sqrt{x^4 + 4c_1 - 4x}}{2} + c_1$$

Verified OK.

6.9.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned}(x^2 + 2y) dy &= (-2yx - 1) dx \\ (2yx + 1) dx + (x^2 + 2y) dy &= 0\end{aligned}\tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 2yx + 1 \\ N(x, y) &= x^2 + 2y\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2yx + 1) \\ &= 2x\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x^2 + 2y) \\ &= 2x\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M\tag{1}$$

$$\frac{\partial \phi}{\partial y} = N\tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 2yx + 1 dx \\ \phi &= yx^2 + x + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial\phi}{\partial y} = x^2 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial\phi}{\partial y} = x^2 + 2y$. Therefore equation (4) becomes

$$x^2 + 2y = x^2 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 2y$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (2y) dy$$
$$f(y) = y^2 + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = yx^2 + y^2 + x + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = yx^2 + y^2 + x$$

Summary

The solution(s) found are the following

$$x^2y + y^2 + x = c_1 \quad (1)$$

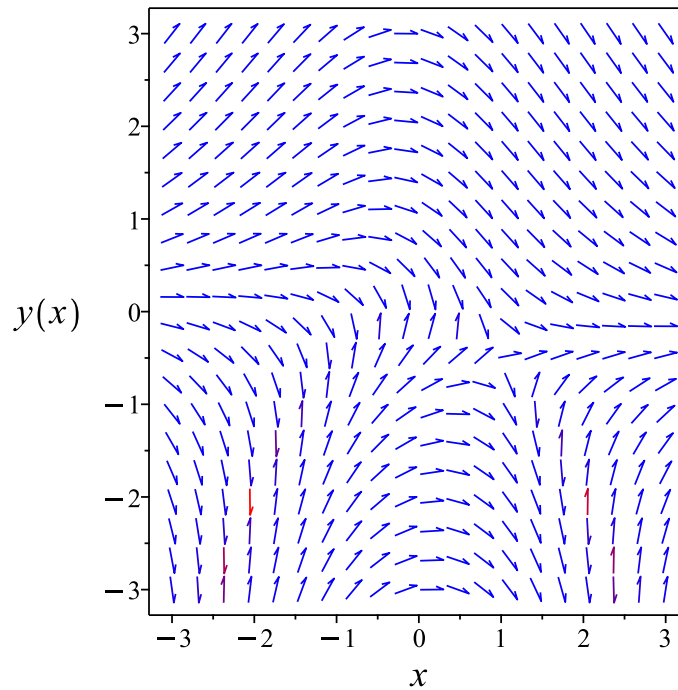


Figure 295: Slope field plot

Verification of solutions

$$x^2y + y^2 + x = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 47

```
dsolve(diff(y(x),x) = (-1-2*x*y(x))/(x^2+2*y(x)),y(x), singsol=all)
```

$$y(x) = -\frac{x^2}{2} - \frac{\sqrt{x^4 - 4c_1 - 4x}}{2}$$
$$y(x) = -\frac{x^2}{2} + \frac{\sqrt{x^4 - 4c_1 - 4x}}{2}$$

✓ Solution by Mathematica

Time used: 0.15 (sec). Leaf size: 61

```
DSolve[y'[x]== (-1-2*x*y[x])/(x^2+2*y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2} \left(-x^2 - \sqrt{x^4 - 4x + 4c_1} \right)$$
$$y(x) \rightarrow \frac{1}{2} \left(-x^2 + \sqrt{x^4 - 4x + 4c_1} \right)$$

6.10 problem 10

6.10.1 Solving as separable ode	1510
6.10.2 Solving as first order ode lie symmetry lookup ode	1512
6.10.3 Solving as exact ode	1517
6.10.4 Maple step by step solution	1520

Internal problem ID [577]

Internal file name [OUTPUT/577_Sunday_June_05_2022_01_45_05_AM_92187482/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 10.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$\frac{yy'}{y-2} = -\frac{-x^2+x+1}{x^2}$$

6.10.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{(x^2 - x - 1)(y - 2)}{y x^2}\end{aligned}$$

Where $f(x) = \frac{x^2-x-1}{x^2}$ and $g(y) = \frac{y-2}{y}$. Integrating both sides gives

$$\frac{1}{\frac{y-2}{y}} dy = \frac{x^2 - x - 1}{x^2} dx$$

$$\int \frac{1}{\frac{y-2}{y}} dy = \int \frac{x^2 - x - 1}{x^2} dx$$

$$y + 2 \ln(y - 2) = x - \ln(x) + \frac{1}{x} + c_1$$

Which results in

$$y = e^{-\frac{x \ln(x) + 2 \operatorname{LambertW}\left(\frac{e^{-\frac{-1 - c_1 x - x^2 + x \ln(x) + 2x}{2x}}}{2}\right) x - c_1 x - x^2 + 2x - 1}{2x}} + 2$$

Summary

The solution(s) found are the following

$$y = e^{-\frac{x \ln(x) + 2 \operatorname{LambertW}\left(\frac{e^{-\frac{-1 - c_1 x - x^2 + x \ln(x) + 2x}{2x}}}{2}\right) x - c_1 x - x^2 + 2x - 1}{2x}} + 2 \quad (1)$$

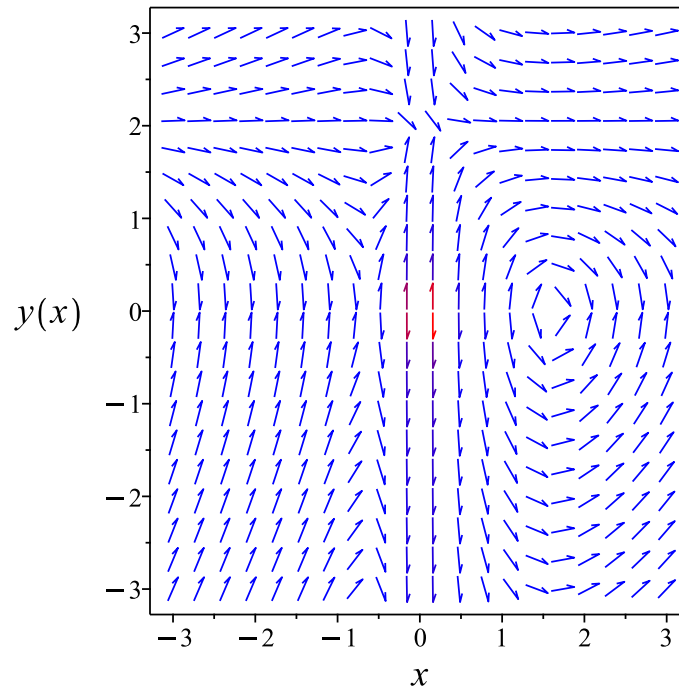


Figure 296: Slope field plot

Verification of solutions

$$y = e^{-\frac{x \ln(x)+2 \operatorname{LambertW}\left(\frac{e^{-1-c_1 x-x^2+x \ln(x)+2x}}{2x}\right)}{2x} x-c_1 x-x^2+2x-1} + 2$$

Verified OK.

6.10.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{y x^2 - 2x^2 - yx + 2x - y + 2}{y x^2}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 278: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{x^2}{x^2 - x - 1} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{x^2}{x^2-x-1}} dx \end{aligned}$$

Which results in

$$S = x - \ln(x) + \frac{1}{x}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y x^2 - 2x^2 - yx + 2x - y + 2}{y x^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= \frac{x^2 - x - 1}{x^2} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{y}{y - 2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{R}{R - 2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = R + 2 \ln(R - 2) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{-x \ln(x) + x^2 + 1}{x} = y + 2 \ln(y - 2) + c_1$$

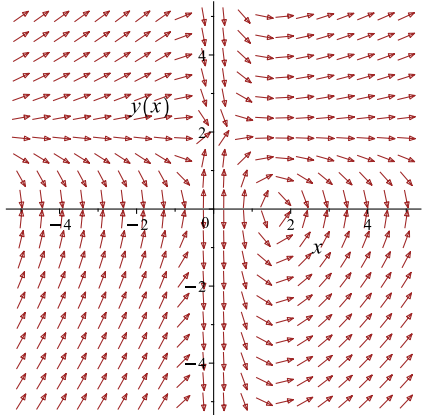
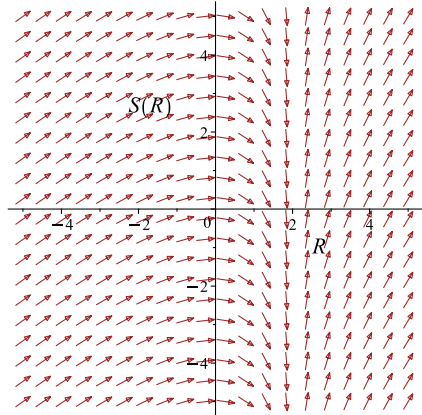
Which simplifies to

$$\frac{-x \ln(x) + x^2 + 1}{x} = y + 2 \ln(y - 2) + c_1$$

Which gives

$$y = 2 \text{LambertW} \left(\frac{e^{-\frac{x \ln(x) + c_1 x - x^2 + 2x - 1}{2x}}}{2} \right) + 2$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{y x^2 - 2x^2 - yx + 2x - y + 2}{y x^2}$ 	$R = y$ $S = \frac{-x \ln(x) + x^2 + 1}{x}$	$\frac{dS}{dR} = \frac{R}{R-2}$ 

Summary

The solution(s) found are the following

$$y = 2 \operatorname{LambertW} \left(\frac{e^{-\frac{x \ln(x) + c_1 x - x^2 + 2x - 1}{2x}}}{2} \right) + 2 \quad (1)$$

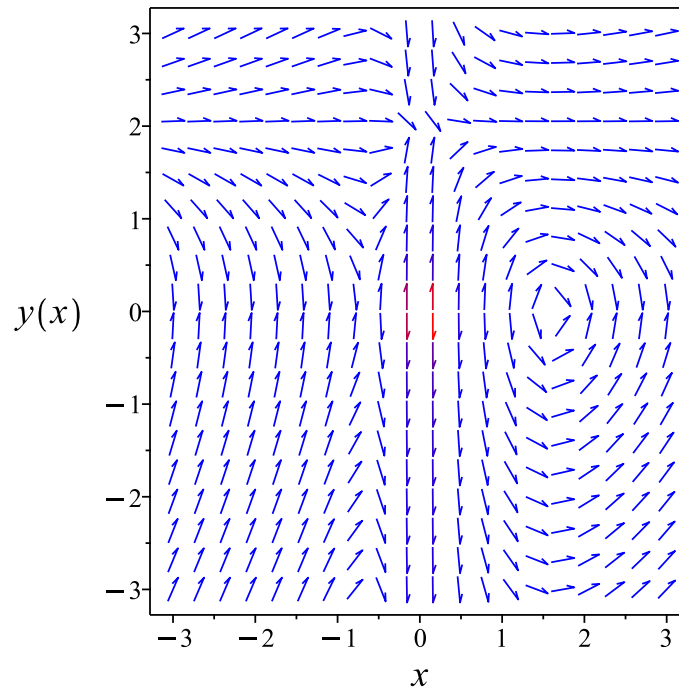


Figure 297: Slope field plot

Verification of solutions

$$y = 2 \operatorname{LambertW} \left(\frac{e^{-\frac{x \ln(x) + c_1 x - x^2 + 2x - 1}{2x}}}{2} \right) + 2$$

Verified OK.

6.10.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} \left(\frac{y}{y-2}\right) dy &= \left(\frac{x^2 - x - 1}{x^2}\right) dx \\ \left(-\frac{x^2 - x - 1}{x^2}\right) dx &+ \left(\frac{y}{y-2}\right) dy = 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$M(x, y) = -\frac{x^2 - x - 1}{x^2}$$
$$N(x, y) = \frac{y}{y - 2}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y} \left(-\frac{x^2 - x - 1}{x^2} \right)$$
$$= 0$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x} \left(\frac{y}{y - 2} \right)$$
$$= 0$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$
$$\int \frac{\partial \phi}{\partial x} dx = \int -\frac{x^2 - x - 1}{x^2} dx$$
$$\phi = -x + \ln(x) - \frac{1}{x} + f(y) \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{y}{y-2}$. Therefore equation (4) becomes

$$\frac{y}{y-2} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{y}{y-2}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(\frac{y}{y-2} \right) dy$$

$$f(y) = y + 2 \ln(y-2) + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x + \ln(x) - \frac{1}{x} + y + 2 \ln(y-2) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x + \ln(x) - \frac{1}{x} + y + 2 \ln(y-2)$$

The solution becomes

$$y = e^{-\frac{x \ln(x) + 2 \operatorname{LambertW}\left(\frac{e^{-1-c_1 x - x^2 + x \ln(x) + 2x}}{2x}\right) x - c_1 x - x^2 + 2x - 1}{2x}} + 2$$

Summary

The solution(s) found are the following

$$y = e^{-\frac{x \ln(x) + 2 \operatorname{LambertW}\left(\frac{e^{-1-c_1 x - x^2 + x \ln(x) + 2x}}{2x}\right) x - c_1 x - x^2 + 2x - 1}{2x}} + 2 \quad (1)$$

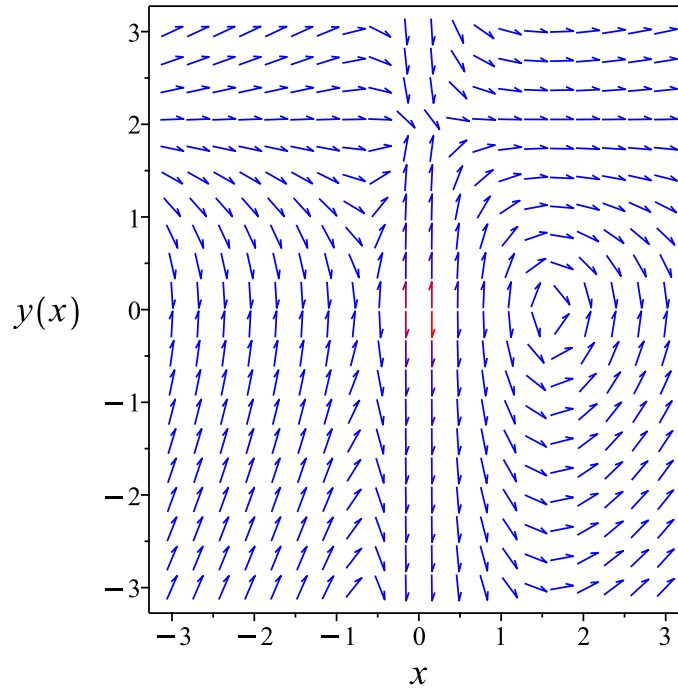


Figure 298: Slope field plot

Verification of solutions

$$y = e^{-\frac{x \ln(x)+2 \operatorname{LambertW}\left(\frac{-1-c_1 x-x^2+x \ln(x)+2 x}{2 x}\right)}{2 x}} x^{-c_1 x-x^2+2 x-1} + 2$$

Verified OK.

6.10.4 Maple step by step solution

Let's solve

$$\frac{yy'}{y-2} = -\frac{-x^2+x+1}{x^2}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Integrate both sides with respect to x

$$\int \frac{yy'}{y-2} dx = \int -\frac{-x^2+x+1}{x^2} dx + c_1$$

- Evaluate integral

$$y + 2 \ln(y - 2) = x - \ln(x) + \frac{1}{x} + c_1$$

- Solve for y

$$y = e^{-\frac{x \ln(x) + 2 \operatorname{LambertW}\left(\frac{e^{-1 - c_1 x - x^2 + x \ln(x) + 2x}}{2x}\right)}{2x} - c_1 x - x^2 + 2x - 1} + 2$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 26

```
dsolve((-x^2+x+1)/x^2+y(x)*diff(y(x),x)/(-2+y(x)) = 0,y(x), singsol=all)
```

$$y(x) = 2 \operatorname{LambertW}\left(\frac{c_1 e^{\frac{(x-1)^2}{2x}}}{2\sqrt{x}}\right) + 2$$

✓ Solution by Mathematica

Time used: 60.036 (sec). Leaf size: 63

```
DSolve[(-x^2+x+1)/x^2+y[x]*y'[x]/(-2+y[x]) == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow 2 \left(1 + W\left(-\frac{1}{2} \sqrt{\frac{e^{x+\frac{1}{x}-2+c_1}}{x}}\right) \right)$$

$$y(x) \rightarrow 2 \left(1 + W\left(\frac{1}{2} \sqrt{\frac{e^{x+\frac{1}{x}-2+c_1}}{x}}\right) \right)$$

6.11 problem 11

- 6.11.1 Solving as differentialType ode 1522
- 6.11.2 Solving as exact ode 1524
- 6.11.3 Maple step by step solution 1527

Internal problem ID [578]

Internal file name [OUTPUT/578_Sunday_June_05_2022_01_45_06_AM_15814802/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 11.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**exact**", "**differentialType**"

Maple gives the following as the ode type

[_exact]

$$y + (e^y + x) y' = -x^2$$

6.11.1 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-x^2 - y}{e^y + x} \quad (1)$$

Which becomes

$$(e^y) dy = (-x) dy + (-x^2 - y) dx \quad (2)$$

But the RHS is complete differential because

$$(-x) dy + (-x^2 - y) dx = d\left(-\frac{1}{3}x^3 - yx\right)$$

Hence (2) becomes

$$(e^y) dy = d\left(-\frac{1}{3}x^3 - yx\right)$$

Integrating both sides gives gives these solutions

$$y = -\text{LambertW}\left(\frac{e^{-x^3+3c_1}}{x}\right) + \frac{-x^3 + 3c_1}{3x} + c_1$$

Summary

The solution(s) found are the following

$$y = -\text{LambertW}\left(\frac{e^{-x^3+3c_1}}{x}\right) + \frac{-x^3 + 3c_1}{3x} + c_1 \quad (1)$$

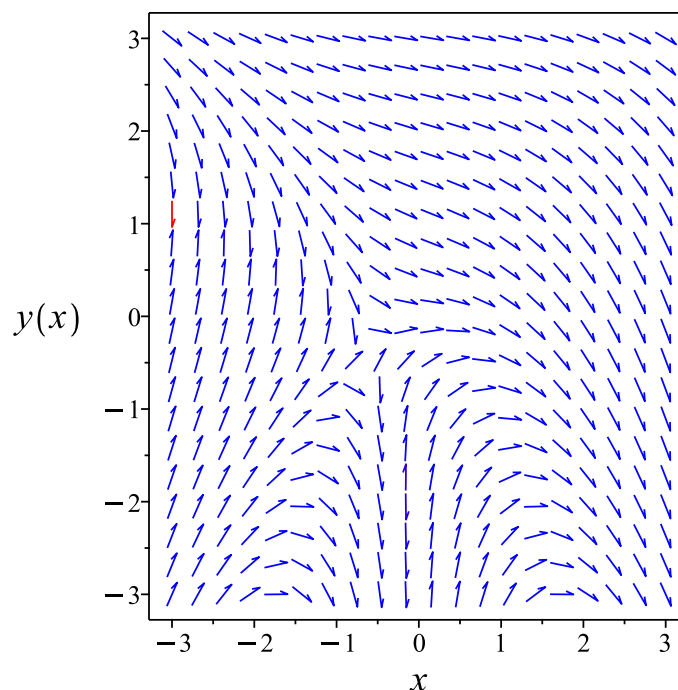


Figure 299: Slope field plot

Verification of solutions

$$y = -\text{LambertW}\left(\frac{e^{-x^3+3c_1}}{x}\right) + \frac{-x^3 + 3c_1}{3x} + c_1$$

Verified OK.

6.11.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (e^y + x) dy &= (-x^2 - y) dx \\ (x^2 + y) dx + (e^y + x) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= x^2 + y \\ N(x, y) &= e^y + x \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(x^2 + y) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(e^y + x) \\ &= 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int x^2 + y dx$$

$$\phi = \frac{1}{3}x^3 + yx + f(y) \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^y + x$. Therefore equation (4) becomes

$$e^y + x = x + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = e^y$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (e^y) dy$$

$$f(y) = e^y + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{x^3}{3} + yx + e^y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{x^3}{3} + yx + e^y$$

The solution becomes

$$y = -\text{LambertW}\left(\frac{e^{\frac{-x^3+3c_1}{3x}}}{x}\right) + \frac{-x^3 + 3c_1}{3x}$$

Summary

The solution(s) found are the following

$$y = -\text{LambertW}\left(\frac{e^{\frac{-x^3+3c_1}{3x}}}{x}\right) + \frac{-x^3 + 3c_1}{3x} \quad (1)$$

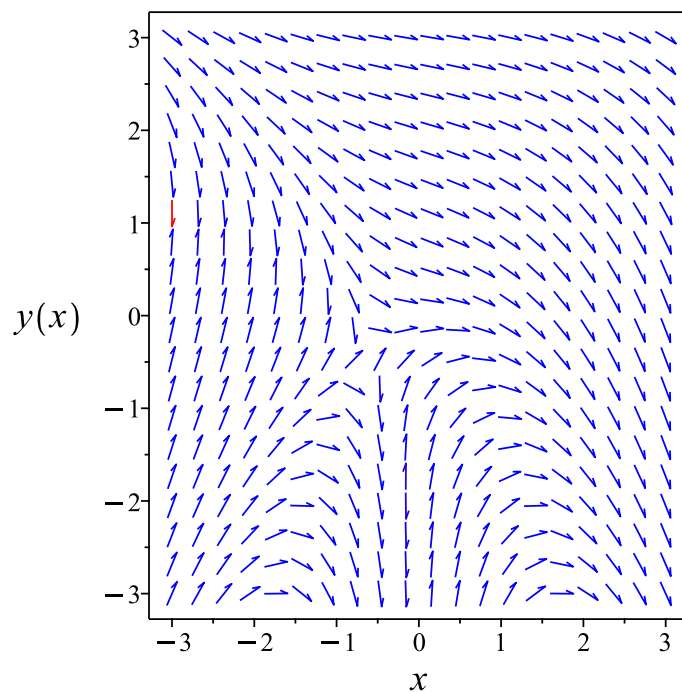


Figure 300: Slope field plot

Verification of solutions

$$y = -\text{LambertW}\left(\frac{e^{\frac{-x^3+3c_1}{3x}}}{x}\right) + \frac{-x^3 + 3c_1}{3x}$$

Verified OK.

6.11.3 Maple step by step solution

Let's solve

$$y + (e^y + x)y' = -x^2$$

- Highest derivative means the order of the ODE is 1
- Check if ODE is exact
 - ODE is exact if the lhs is the total derivative of a C^2 function
 - $F'(x, y) = 0$
 - Compute derivative of lhs

$$F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y) \right) y' = 0$$

- Evaluate derivatives

$$1 = 1$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$\left[F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$

- Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to x

$$F(x, y) = \int (x^2 + y) dx + f_1(y)$$

- Evaluate integral

$$F(x, y) = \frac{x^3}{3} + yx + f_1(y)$$

- Take derivative of $F(x, y)$ with respect to y

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$

- Compute derivative

$$e^y + x = x + \frac{d}{dy} f_1(y)$$

- Isolate for $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = e^y$$

- Solve for $f_1(y)$

$$f_1(y) = e^y$$

- Substitute $f_1(y)$ into equation for $F(x, y)$

$$F(x, y) = \frac{x^3}{3} + yx + e^y$$

- Substitute $F(x, y)$ into the solution of the ODE

$$\frac{x^3}{3} + yx + e^y = c_1$$

- Solve for y

$$y = -\text{LambertW}\left(\frac{e^{\frac{-x^3+3c_1}{3x}}}{x}\right) + \frac{-x^3+3c_1}{3x}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 39

```
dsolve(x^2+y(x)+(exp(y(x))+x)*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = \frac{-x^3 - 3x \operatorname{LambertW}\left(\frac{e^{-\frac{x^3+3c_1}{3x}}}{x}\right) - 3c_1}{3x}$$

✓ Solution by Mathematica

Time used: 3.877 (sec). Leaf size: 42

```
DSolve[x^2+y[x]+(Exp[y[x]]+x)*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -W\left(\frac{e^{-\frac{x^2+c_1}{3}}}{x}\right) - \frac{x^2}{3} + \frac{c_1}{x}$$

6.12 problem 12

6.12.1 Solving as linear ode	1530
6.12.2 Solving as first order ode lie symmetry lookup ode	1532
6.12.3 Solving as exact ode	1536
6.12.4 Maple step by step solution	1541

Internal problem ID [579]

Internal file name [OUTPUT/579_Sunday_June_05_2022_01_45_07_AM_95196417/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 12.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$y' + y = \frac{1}{1 + e^x}$$

6.12.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = 1$$
$$q(x) = \frac{1}{1 + e^x}$$

Hence the ode is

$$y' + y = \frac{1}{1 + e^x}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int 1 dx} \\ &= e^x\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) \left(\frac{1}{1 + e^x} \right) \\ \frac{d}{dx}(e^x y) &= (e^x) \left(\frac{1}{1 + e^x} \right) \\ d(e^x y) &= \left(\frac{e^x}{1 + e^x} \right) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}e^x y &= \int \frac{e^x}{1 + e^x} dx \\ e^x y &= \ln(1 + e^x) + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^x$ results in

$$y = e^{-x} \ln(1 + e^x) + c_1 e^{-x}$$

which simplifies to

$$y = e^{-x}(\ln(1 + e^x) + c_1)$$

Summary

The solution(s) found are the following

$$y = e^{-x}(\ln(1 + e^x) + c_1) \tag{1}$$

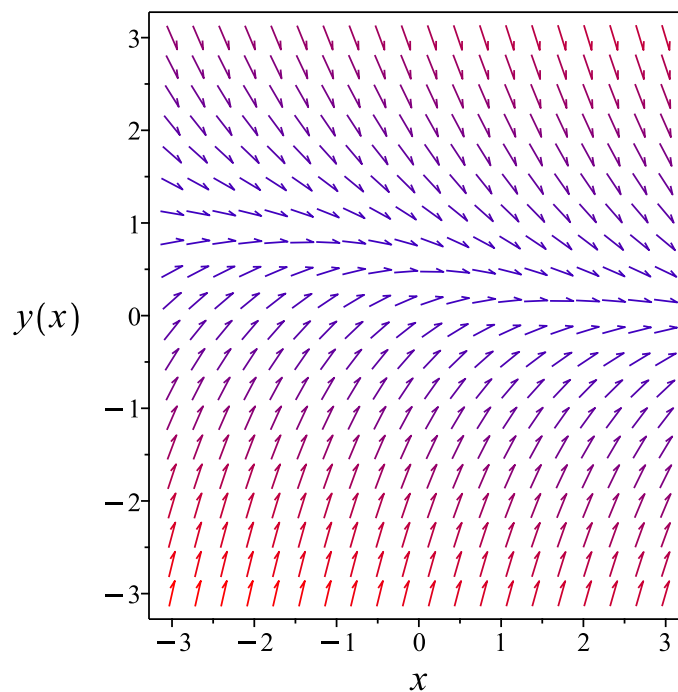


Figure 301: Slope field plot

Verification of solutions

$$y = e^{-x}(\ln(1 + e^x) + c_1)$$

Verified OK.

6.12.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{e^x y + y - 1}{1 + e^x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 282: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= e^{-x}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-x}} dy \end{aligned}$$

Which results in

$$S = e^x y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{e^x y + y - 1}{1 + e^x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= e^x y \\ S_y &= e^x \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{e^x}{1 + e^x} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{e^R}{1 + e^R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \ln(1 + e^R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$e^x y = \ln(1 + e^x) + c_1$$

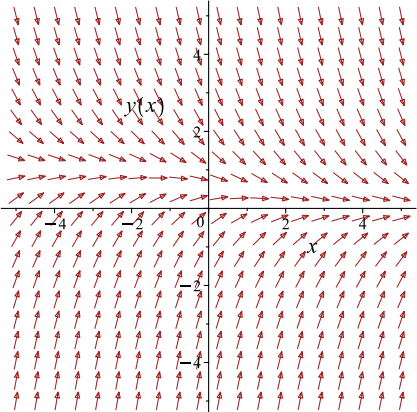
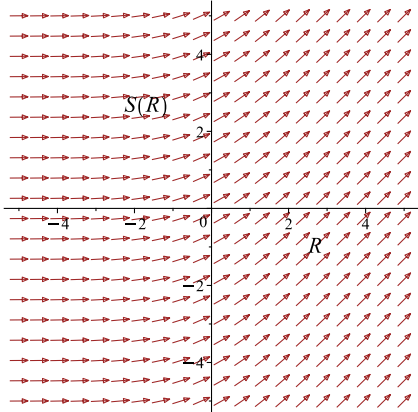
Which simplifies to

$$e^x y = \ln(1 + e^x) + c_1$$

Which gives

$$y = e^{-x}(\ln(1 + e^x) + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{e^x y + y - 1}{1 + e^x}$ 	$R = x$ $S = e^x y$	$\frac{dS}{dR} = \frac{e^R}{1 + e^R}$ 

Summary

The solution(s) found are the following

$$y = e^{-x}(\ln(1 + e^x) + c_1) \quad (1)$$

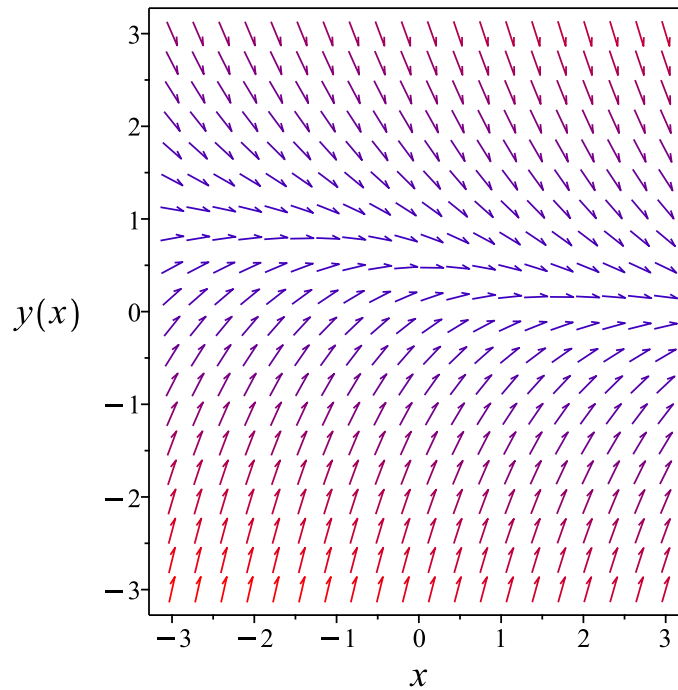


Figure 302: Slope field plot

Verification of solutions

$$y = e^{-x}(\ln(1 + e^x) + c_1)$$

Verified OK.

6.12.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dy &= \left(-y + \frac{1}{1 + e^x}\right) dx \\ \left(y - \frac{1}{1 + e^x}\right) dx + dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= y - \frac{1}{1 + e^x} \\ N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(y - \frac{1}{1 + e^x}\right) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((1) - (0)) \\ &= 1\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int 1 dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^x \\ &= e^x\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= e^x \left(y - \frac{1}{1 + e^x} \right) \\ &= \frac{e^x(e^x y + y - 1)}{1 + e^x}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^x(1) \\ &= e^x\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left(\frac{e^x(e^x y + y - 1)}{1 + e^x} \right) + (e^x) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{e^x(e^x y + y - 1)}{1 + e^x} dx \\ \phi &= e^x y - \ln(1 + e^x) + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^x + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^x$. Therefore equation (4) becomes

$$e^x = e^x + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = e^x y - \ln(1 + e^x) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = e^x y - \ln(1 + e^x)$$

The solution becomes

$$y = e^{-x}(\ln(1 + e^x) + c_1)$$

Summary

The solution(s) found are the following

$$y = e^{-x}(\ln(1 + e^x) + c_1) \tag{1}$$

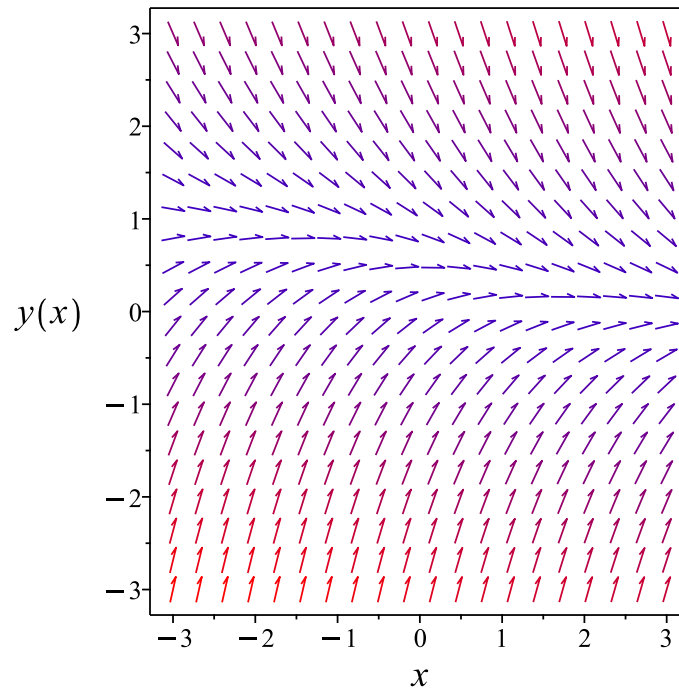


Figure 303: Slope field plot

Verification of solutions

$$y = e^{-x}(\ln(1 + e^x) + c_1)$$

Verified OK.

6.12.4 Maple step by step solution

Let's solve

$$y' + y = \frac{1}{1+e^x}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -y + \frac{1}{1+e^x}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + y = \frac{1}{1+e^x}$$

- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$\mu(x)(y' + y) = \frac{\mu(x)}{1+e^x}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dx}(\mu(x)y)$

$$\mu(x)(y' + y) = \mu'(x)y + \mu(x)y'$$

- Isolate $\mu'(x)$

$$\mu'(x) = \mu(x)$$

- Solve to find the integrating factor

$$\mu(x) = e^x$$

- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x)y) \right) dx = \int \frac{\mu(x)}{1+e^x} dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int \frac{\mu(x)}{1+e^x} dx + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(x)}{1+e^x} dx + c_1}{\mu(x)}$$

- Substitute $\mu(x) = e^x$

$$y = \frac{\int \frac{e^x}{1+e^x} dx + c_1}{e^x}$$

- Evaluate the integrals on the rhs

$$y = \frac{\ln(1+e^x) + c_1}{e^x}$$

- Simplify

$$y = e^{-x}(\ln(1 + e^x) + c_1)$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 16

```
dsolve(y(x)+diff(y(x),x) = 1/(1+exp(x)),y(x), singsol=all)
```

$$y(x) = (\ln(1 + e^x) + c_1)e^{-x}$$

✓ Solution by Mathematica

Time used: 0.067 (sec). Leaf size: 20

```
DSolve[y[x]+y'[x] == 1/(1+Exp[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x}(\log(e^x + 1) + c_1)$$

6.13 problem 13

6.13.1 Solving as separable ode	1543
6.13.2 Solving as first order ode lie symmetry lookup ode	1545
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6.13.5 Maple step by step solution	1555

Internal problem ID [580]

Internal file name [OUTPUT/580_Sunday_June_05_2022_01_45_08_AM_16506274/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 13.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "riccati", "separable", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - y^2 - 2xy^2 = 1 + 2x$$

6.13.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= (y^2 + 1)(1 + 2x)\end{aligned}$$

Where $f(x) = 1 + 2x$ and $g(y) = y^2 + 1$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y^2 + 1} dy &= 1 + 2x dx \\ \int \frac{1}{y^2 + 1} dy &= \int 1 + 2x dx \\ \arctan(y) &= x^2 + c_1 + x\end{aligned}$$

Which results in

$$y = \tan(x^2 + c_1 + x)$$

Summary

The solution(s) found are the following

$$y = \tan(x^2 + c_1 + x) \tag{1}$$

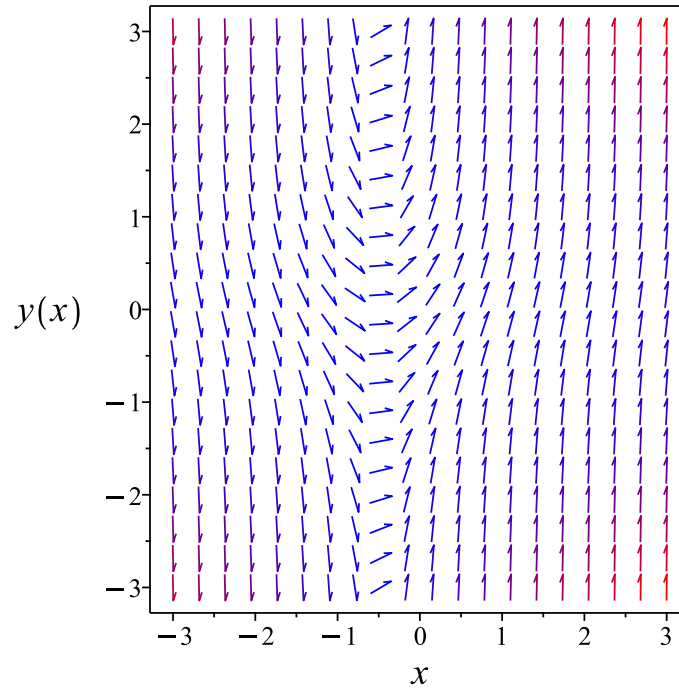


Figure 304: Slope field plot

Verification of solutions

$$y = \tan(x^2 + c_1 + x)$$

Verified OK.

6.13.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = 2xy^2 + y^2 + 2x + 1$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 285: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{1 + 2x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{1+2x}} dx\end{aligned}$$

Which results in

$$S = x^2 + x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = 2x y^2 + y^2 + 2x + 1$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\ R_y &= 1 \\ S_x &= 1 + 2x \\ S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{y^2 + 1} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R^2 + 1}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \arctan(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$x^2 + x = \arctan(y) + c_1$$

Which simplifies to

$$x^2 + x = \arctan(y) + c_1$$

Which gives

$$y = -\tan(-x^2 + c_1 - x)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = 2x y^2 + y^2 + 2x + 1$	$R = y$ $S = x^2 + x$	$\frac{dS}{dR} = \frac{1}{R^2+1}$

Summary

The solution(s) found are the following

$$y = -\tan(-x^2 + c_1 - x) \tag{1}$$

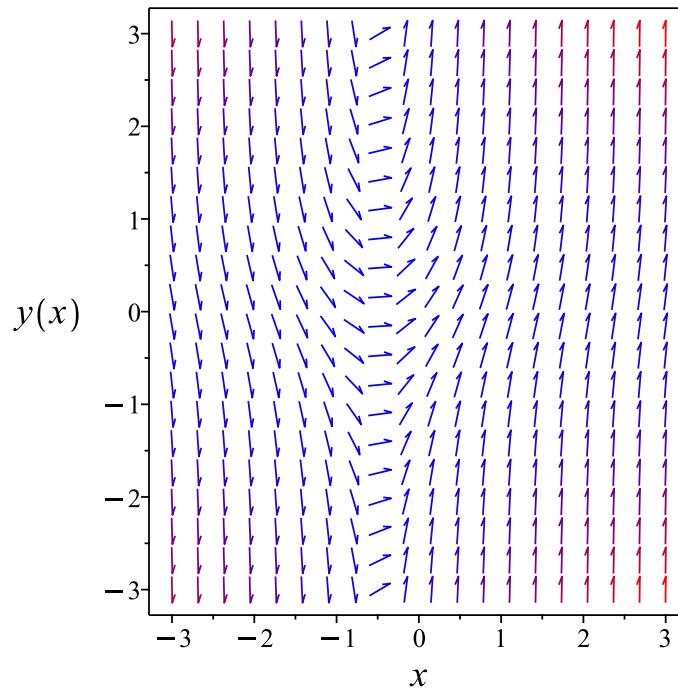


Figure 305: Slope field plot

Verification of solutions

$$y = -\tan(-x^2 + c_1 - x)$$

Verified OK.

6.13.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(\frac{1}{y^2 + 1}\right) dy &= (1 + 2x) dx \\ (-1 - 2x) dx + \left(\frac{1}{y^2 + 1}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -1 - 2x \\ N(x, y) &= \frac{1}{y^2 + 1}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-1 - 2x) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{1}{y^2 + 1} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -1 - 2x dx \\ \phi &= -x^2 - x + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{y^2+1}$. Therefore equation (4) becomes

$$\frac{1}{y^2 + 1} = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{1}{y^2 + 1}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(\frac{1}{y^2 + 1} \right) dy \\ f(y) &= \arctan(y) + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x^2 - x + \arctan(y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x^2 - x + \arctan(y)$$

The solution becomes

$$y = \tan(x^2 + c_1 + x)$$

Summary

The solution(s) found are the following

$$y = \tan(x^2 + c_1 + x) \tag{1}$$

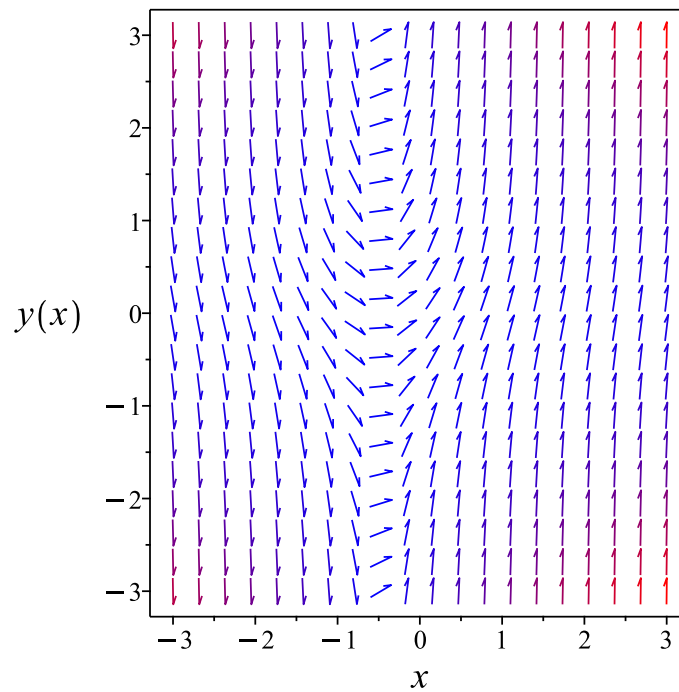


Figure 306: Slope field plot

Verification of solutions

$$y = \tan(x^2 + c_1 + x)$$

Verified OK.

6.13.4 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= 2x y^2 + y^2 + 2x + 1\end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = 2x y^2 + y^2 + 2x + 1$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that $f_0(x) = 1 + 2x$, $f_1(x) = 0$ and $f_2(x) = 1 + 2x$. Let

$$\begin{aligned}y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{(1 + 2x) u}\end{aligned}\tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0\tag{2}$$

But

$$\begin{aligned}f_2' &= 2 \\ f_1 f_2 &= 0 \\ f_2^2 f_0 &= (1 + 2x)^3\end{aligned}$$

Substituting the above terms back in equation (2) gives

$$(1 + 2x) u''(x) - 2u'(x) + (1 + 2x)^3 u(x) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 \sin(x^2 + x) + c_2 \cos(x^2 + x)$$

The above shows that

$$u'(x) = (1 + 2x) (c_1 \cos(x^2 + x) - c_2 \sin(x^2 + x))$$

Using the above in (1) gives the solution

$$y = -\frac{c_1 \cos(x^2 + x) - c_2 \sin(x^2 + x)}{c_1 \sin(x^2 + x) + c_2 \cos(x^2 + x)}$$

Dividing both numerator and denominator by c_1 gives, after renaming the constant $\frac{c_2}{c_1} = c_3$ the following solution

$$y = \frac{-c_3 \cos(x^2 + x) + \sin(x^2 + x)}{c_3 \sin(x^2 + x) + \cos(x^2 + x)}$$

Summary

The solution(s) found are the following

$$y = \frac{-c_3 \cos(x^2 + x) + \sin(x^2 + x)}{c_3 \sin(x^2 + x) + \cos(x^2 + x)} \quad (1)$$

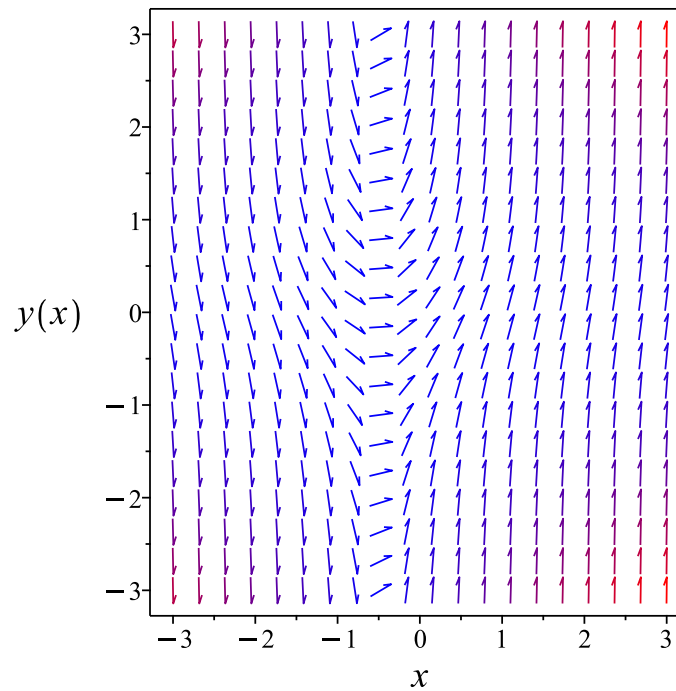


Figure 307: Slope field plot

Verification of solutions

$$y = \frac{-c_3 \cos(x^2 + x) + \sin(x^2 + x)}{c_3 \sin(x^2 + x) + \cos(x^2 + x)}$$

Verified OK.

6.13.5 Maple step by step solution

Let's solve

$$y' - y^2 - 2xy^2 = 1 + 2x$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{1+y^2} = 1 + 2x$$

- Integrate both sides with respect to x

$$\int \frac{y'}{1+y^2} dx = \int (1 + 2x) dx + c_1$$

- Evaluate integral

$$\arctan(y) = x^2 + c_1 + x$$

- Solve for y

$$y = \tan(x^2 + c_1 + x)$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 11

```
dsolve(diff(y(x),x) = 1+2*x+y(x)^2+2*x*y(x)^2,y(x), singsol=all)
```

$$y(x) = \tan(x^2 + c_1 + x)$$

✓ Solution by Mathematica

Time used: 0.178 (sec). Leaf size: 13

```
DSolve[y'[x] == 1+2*x+y[x]^2+2*x*y[x]^2,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \tan(x^2 + x + c_1)$$

6.14 problem 14

6.14.1 Existence and uniqueness analysis	1557
6.14.2 Solving as homogeneousTypeD2 ode	1558
6.14.3 Solving as differentialType ode	1560
6.14.4 Solving as first order ode lie symmetry calculated ode	1562
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6.14.6 Maple step by step solution	1570

Internal problem ID [581]

Internal file name [OUTPUT/581_Sunday_June_05_2022_01_45_09_AM_47757679/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 14.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "differentialType", "homogeneousTypeD2", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _exact, _rational, [_Abel, `2nd
type`, `class A`]]
```

$$y + (x + 2y)y' = -x$$

With initial conditions

$$[y(2) = 3]$$

6.14.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= -\frac{x + y}{x + 2y}\end{aligned}$$

The x domain of $f(x, y)$ when $y = 3$ is

$$\{x < -6 \vee -6 < x\}$$

And the point $x_0 = 2$ is inside this domain. The y domain of $f(x, y)$ when $x = 2$ is

$$\{y < -1 \vee -1 < y\}$$

And the point $y_0 = 3$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned} \frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{x+y}{x+2y} \right) \\ &= -\frac{1}{x+2y} + \frac{2x+2y}{(x+2y)^2} \end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 3$ is

$$\{x < -6 \vee -6 < x\}$$

And the point $x_0 = 2$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = 2$ is

$$\{y < -1 \vee -1 < y\}$$

And the point $y_0 = 3$ is inside this domain. Therefore solution exists and is unique.

6.14.2 Solving as homogeneous TypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u(x)x + (x + 2u(x)x)(u'(x)x + u(x)) = -x$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{2u^2 + 2u + 1}{x(2u + 1)} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(u) = \frac{2u^2+2u+1}{2u+1}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{2u^2+2u+1}{2u+1}} du &= -\frac{1}{x} dx \\ \int \frac{1}{\frac{2u^2+2u+1}{2u+1}} du &= \int -\frac{1}{x} dx \\ \frac{\ln(2u^2 + 2u + 1)}{2} &= -\ln(x) + c_2\end{aligned}$$

Raising both side to exponential gives

$$\sqrt{2u^2 + 2u + 1} = e^{-\ln(x)+c_2}$$

Which simplifies to

$$\sqrt{2u^2 + 2u + 1} = \frac{c_3}{x}$$

Which simplifies to

$$\sqrt{2u(x)^2 + 2u(x) + 1} = \frac{c_3 e^{c_2}}{x}$$

The solution is

$$\sqrt{2u(x)^2 + 2u(x) + 1} = \frac{c_3 e^{c_2}}{x}$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\begin{aligned}\sqrt{\frac{2y^2}{x^2} + \frac{2y}{x} + 1} &= \frac{c_3 e^{c_2}}{x} \\ \sqrt{\frac{2y^2 + 2yx + x^2}{x^2}} &= \frac{c_3 e^{c_2}}{x}\end{aligned}$$

Substituting initial conditions and solving for c_2 gives $c_2 = \frac{\ln\left(\frac{34}{c_3^2}\right)}{2}$. Hence the solution becomes Initial conditions are used to solve for c_3 . Substituting $x = 2$ and $y = 3$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{\sqrt{34}}{2} = \frac{c_3 \sqrt{34}}{2} \sqrt{\frac{1}{c_3^2}}$$

This solution is valid for any c_3 . Hence there are infinite number of solutions.

Summary

The solution(s) found are the following

$$\sqrt{\frac{2y^2 + 2yx + x^2}{x^2}} = \frac{c_3 \sqrt{34} \sqrt{\frac{1}{c_3^2}}}{x} \quad (1)$$

Verification of solutions

$$\sqrt{\frac{2y^2 + 2yx + x^2}{x^2}} = \frac{c_3 \sqrt{34} \sqrt{\frac{1}{c_3^2}}}{x}$$

Verified OK.

6.14.3 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-x - y}{x + 2y} \quad (1)$$

Which becomes

$$(2y) dy = (-x) dy + (-x - y) dx \quad (2)$$

But the RHS is complete differential because

$$(-x) dy + (-x - y) dx = d\left(-\frac{1}{2}x^2 - yx\right)$$

Hence (2) becomes

$$(2y) dy = d\left(-\frac{1}{2}x^2 - yx\right)$$

Integrating both sides gives gives these solutions

$$y = -\frac{x}{2} + \frac{\sqrt{-x^2 + 4c_1}}{2} + c_1$$
$$y = -\frac{x}{2} - \frac{\sqrt{-x^2 + 4c_1}}{2} + c_1$$

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 3$ in the above solution gives an equation to solve for the constant of integration.

$$3 = -1 - \sqrt{c_1 - 1} + c_1$$

$$c_1 = \frac{9}{2} + \frac{\sqrt{13}}{2}$$

Substituting c_1 found above in the general solution gives

$$y = -\frac{x}{2} - \frac{\sqrt{-x^2 + 18 + 2\sqrt{13}}}{2} + \frac{9}{2} + \frac{\sqrt{13}}{2}$$

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 3$ in the above solution gives an equation to solve for the constant of integration.

$$3 = -1 + \sqrt{c_1 - 1} + c_1$$

$$c_1 = \frac{9}{2} - \frac{\sqrt{13}}{2}$$

Substituting c_1 found above in the general solution gives

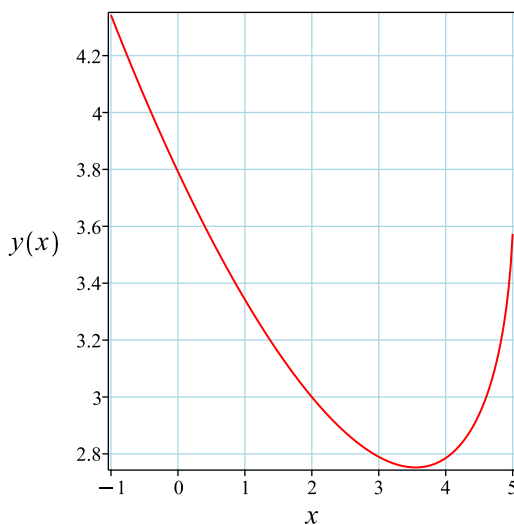
$$y = -\frac{x}{2} + \frac{\sqrt{-x^2 + 18 - 2\sqrt{13}}}{2} + \frac{9}{2} - \frac{\sqrt{13}}{2}$$

Summary

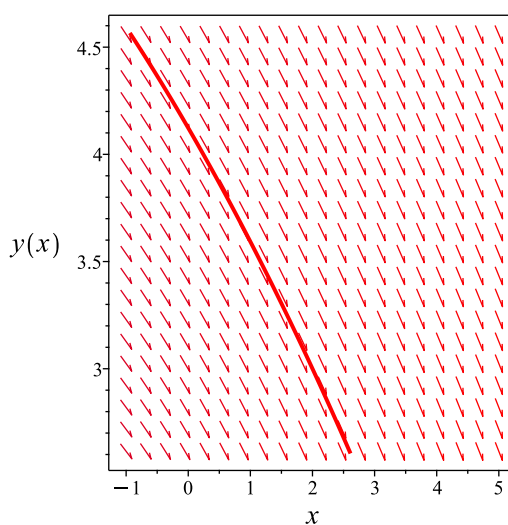
The solution(s) found are the following

$$y = -\frac{x}{2} + \frac{\sqrt{-x^2 + 18 - 2\sqrt{13}}}{2} + \frac{9}{2} - \frac{\sqrt{13}}{2} \quad (1)$$

$$y = -\frac{x}{2} - \frac{\sqrt{-x^2 + 18 + 2\sqrt{13}}}{2} + \frac{9}{2} + \frac{\sqrt{13}}{2} \quad (2)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{x}{2} + \frac{\sqrt{-x^2 + 18 - 2\sqrt{13}}}{2} + \frac{9}{2} - \frac{\sqrt{13}}{2}$$

Verified OK. {positive}

$$y = -\frac{x}{2} - \frac{\sqrt{-x^2 + 18 + 2\sqrt{13}}}{2} + \frac{9}{2} + \frac{\sqrt{13}}{2}$$

Verified OK. {positive}

6.14.4 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{x+y}{x+2y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 - \frac{(x+y)(b_3 - a_2)}{x+2y} - \frac{(x+y)^2 a_3}{(x+2y)^2} - \left(-\frac{1}{x+2y} + \frac{x+y}{(x+2y)^2} \right) (xa_2 + ya_3 + a_1) \quad (\text{5E})$$

$$- \left(-\frac{1}{x+2y} + \frac{2x+2y}{(x+2y)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{x^2 a_2 - x^2 a_3 - x^2 b_3 + 4xy a_2 - 2xy a_3 + 4xy b_2 - 4xy b_3 + 2y^2 a_2 + 4y^2 b_2 - 2y^2 b_3 - x b_1 + y a_1}{(x + 2y)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} x^2 a_2 - x^2 a_3 - x^2 b_3 + 4xy a_2 - 2xy a_3 + 4xy b_2 \\ - 4xy b_3 + 2y^2 a_2 + 4y^2 b_2 - 2y^2 b_3 - x b_1 + y a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} a_2 v_1^2 + 4a_2 v_1 v_2 + 2a_2 v_2^2 - a_3 v_1^2 - 2a_3 v_1 v_2 + 4b_2 v_1 v_2 \\ + 4b_2 v_2^2 - b_3 v_1^2 - 4b_3 v_1 v_2 - 2b_3 v_2^2 + a_1 v_2 - b_1 v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$(a_2 - a_3 - b_3) v_1^2 + (4a_2 - 2a_3 + 4b_2 - 4b_3) v_1 v_2 - b_1 v_1 + (2a_2 + 4b_2 - 2b_3) v_2^2 + a_1 v_2 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} a_1 &= 0 \\ -b_1 &= 0 \\ a_2 - a_3 - b_3 &= 0 \\ 2a_2 + 4b_2 - 2b_3 &= 0 \\ 4a_2 - 2a_3 + 4b_2 - 4b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= 0 \\
 a_2 &= -2b_2 + b_3 \\
 a_3 &= -2b_2 \\
 b_1 &= 0 \\
 b_2 &= b_2 \\
 b_3 &= b_3
 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}
 \xi &= x \\
 \eta &= y
 \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}
 \eta &= \eta - \omega(x, y) \xi \\
 &= y - \left(-\frac{x+y}{x+2y} \right) (x) \\
 &= \frac{x^2 + 2yx + 2y^2}{x+2y} \\
 \xi &= 0
 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}
 S &= \int \frac{1}{\eta} dy \\
 &= \int \frac{1}{\frac{x^2 + 2yx + 2y^2}{x+2y}} dy
 \end{aligned}$$

Which results in

$$S = \frac{\ln(x^2 + 2yx + 2y^2)}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{x + y}{x + 2y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{x + y}{x^2 + 2yx + 2y^2} \\ S_y &= \frac{x + 2y}{x^2 + 2yx + 2y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \quad (4)$$

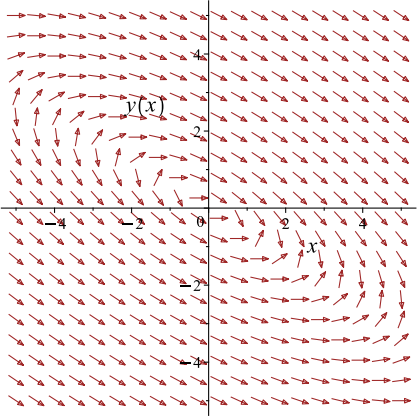
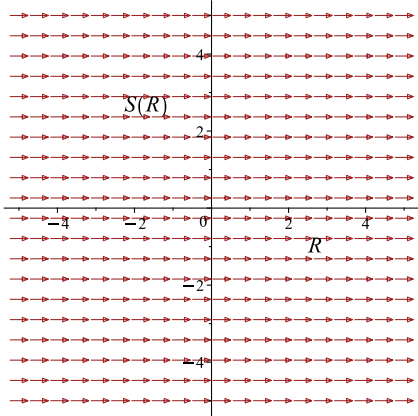
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\ln(2y^2 + 2yx + x^2)}{2} = c_1$$

Which simplifies to

$$\frac{\ln(2y^2 + 2yx + x^2)}{2} = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{x+y}{x+2y}$ 	$R = x$ $S = \frac{\ln(x^2 + 2yx + 2y^2)}{2}$	$\frac{dS}{dR} = 0$ 

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 3$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{\ln(2)}{2} + \frac{\ln(17)}{2} = c_1$$

$$c_1 = \frac{\ln(2)}{2} + \frac{\ln(17)}{2}$$

Substituting c_1 found above in the general solution gives

$$\frac{\ln(x^2 + 2yx + 2y^2)}{2} = \frac{\ln(2)}{2} + \frac{\ln(17)}{2}$$

Summary

The solution(s) found are the following

$$\frac{\ln(2y^2 + 2yx + x^2)}{2} = \frac{\ln(2)}{2} + \frac{\ln(17)}{2} \quad (1)$$

Verification of solutions

$$\frac{\ln(2y^2 + 2yx + x^2)}{2} = \frac{\ln(2)}{2} + \frac{\ln(17)}{2}$$

Verified OK.

6.14.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (x + 2y) dy &= (-x - y) dx \\ (x + y) dx + (x + 2y) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= x + y \\N(x, y) &= x + 2y\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(x + y) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x + 2y) \\ &= 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int x + y dx \\ \phi &= \frac{x(x + 2y)}{2} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = x + 2y$. Therefore equation (4) becomes

$$x + 2y = x + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 2y$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (2y) dy$$
$$f(y) = y^2 + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{x(x + 2y)}{2} + y^2 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{x(x + 2y)}{2} + y^2$$

Initial conditions are used to solve for c_1 . Substituting $x = 2$ and $y = 3$ in the above solution gives an equation to solve for the constant of integration.

$$17 = c_1$$

$$c_1 = 17$$

Substituting c_1 found above in the general solution gives

$$\frac{x(x + 2y)}{2} + y^2 = 17$$

Summary

The solution(s) found are the following

$$\frac{x^2}{2} + yx + y^2 = 17 \quad (1)$$

Verification of solutions

$$\frac{x^2}{2} + yx + y^2 = 17$$

Verified OK.

6.14.6 Maple step by step solution

Let's solve

$$[y + (x + 2y)y' = -x, y(2) = 3]$$

- Highest derivative means the order of the ODE is 1
 y'
- Check if ODE is exact
 - ODE is exact if the lhs is the total derivative of a C^2 function
 $F'(x, y) = 0$
 - Compute derivative of lhs
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$
 - Evaluate derivatives
 $1 = 1$
 - Condition met, ODE is exact
- Exact ODE implies solution will be of this form
 $\left[F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y)\right]$
- Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to x
 $F(x, y) = \int (x + y) dx + f_1(y)$
- Evaluate integral
 $F(x, y) = \frac{x^2}{2} + yx + f_1(y)$
- Take derivative of $F(x, y)$ with respect to y
 $N(x, y) = \frac{\partial}{\partial y} F(x, y)$
- Compute derivative
 $x + 2y = x + \frac{d}{dy} f_1(y)$
- Isolate for $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = 2y$$

- Solve for $f_1(y)$

$$f_1(y) = y^2$$

- Substitute $f_1(y)$ into equation for $F(x, y)$

$$F(x, y) = \frac{1}{2}x^2 + yx + y^2$$

- Substitute $F(x, y)$ into the solution of the ODE

$$\frac{1}{2}x^2 + yx + y^2 = c_1$$

- Solve for y

$$\left\{ y = -\frac{x}{2} - \frac{\sqrt{-x^2+4c_1}}{2}, y = -\frac{x}{2} + \frac{\sqrt{-x^2+4c_1}}{2} \right\}$$

- Use initial condition $y(2) = 3$

$$3 = -1 - \frac{\sqrt{-4+4c_1}}{2}$$

- Solution does not satisfy initial condition

- Use initial condition $y(2) = 3$

$$3 = -1 + \frac{\sqrt{-4+4c_1}}{2}$$

- Solve for c_1

$$c_1 = 17$$

- Substitute $c_1 = 17$ into general solution and simplify

$$y = -\frac{x}{2} + \frac{\sqrt{-x^2+68}}{2}$$

- Solution to the IVP

$$y = -\frac{x}{2} + \frac{\sqrt{-x^2+68}}{2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 0.093 (sec). Leaf size: 19

```
dsolve([x+y(x)+(x+2*y(x))*diff(y(x),x) = 0,y(2) = 3],y(x), singsol=all)
```

$$y(x) = -\frac{x}{2} + \frac{\sqrt{-x^2 + 68}}{2}$$

✓ Solution by Mathematica

Time used: 0.458 (sec). Leaf size: 24

```
DSolve[{x+y[x]+(x+2*y[x])*y'[x] == 0,y[2]==3},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2}(\sqrt{68 - x^2} - x)$$

6.15 problem 15

6.15.1 Solving as separable ode	1573
6.15.2 Solving as linear ode	1575
6.15.3 Solving as homogeneousTypeD2 ode	1576
6.15.4 Solving as first order ode lie symmetry lookup ode	1578
6.15.5 Solving as exact ode	1582
6.15.6 Maple step by step solution	1586

Internal problem ID [582]

Internal file name [OUTPUT/582_Sunday_June_05_2022_01_45_11_AM_3677461/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 15.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$(1 + e^x)y' - y + e^xy = 0$$

6.15.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\frac{y(e^x - 1)}{1 + e^x}\end{aligned}$$

Where $f(x) = -\frac{e^x-1}{1+e^x}$ and $g(y) = y$. Integrating both sides gives

$$\begin{aligned}\frac{1}{y} dy &= -\frac{e^x - 1}{1 + e^x} dx \\ \int \frac{1}{y} dy &= \int -\frac{e^x - 1}{1 + e^x} dx \\ \ln(y) &= \ln(e^x) - 2 \ln(1 + e^x) + c_1 \\ y &= e^{\ln(e^x) - 2 \ln(1 + e^x) + c_1} \\ &= c_1 e^{\ln(e^x) - 2 \ln(1 + e^x)}\end{aligned}$$

Which simplifies to

$$y = \frac{c_1 e^x}{(1 + e^x)^2}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 e^x}{(1 + e^x)^2} \tag{1}$$

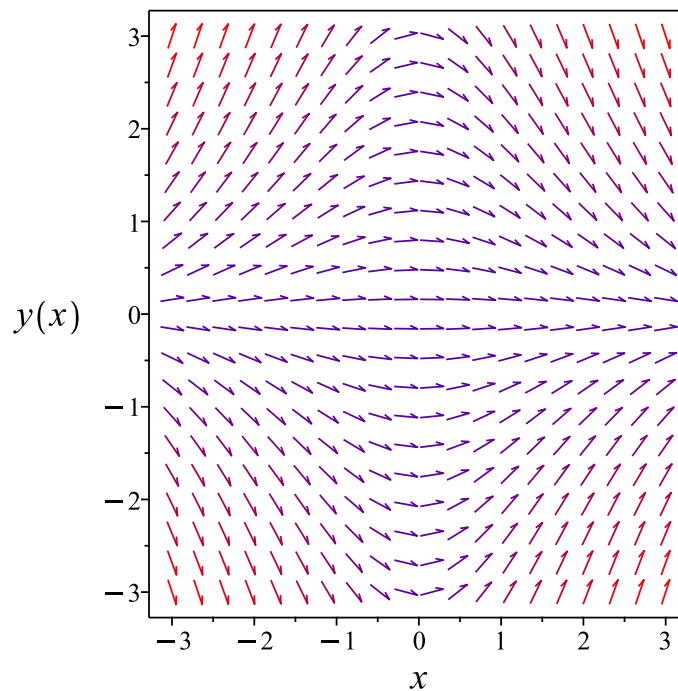


Figure 309: Slope field plot

Verification of solutions

$$y = \frac{c_1 e^x}{(1 + e^x)^2}$$

Verified OK.

6.15.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{-e^x + 1}{1 + e^x}$$

$$q(x) = 0$$

Hence the ode is

$$y' - \frac{(-e^x + 1)y}{1 + e^x} = 0$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{-e^x + 1}{1 + e^x} dx} \\ &= e^{-\ln(e^x) + 2\ln(1 + e^x)}\end{aligned}$$

Which simplifies to

$$\mu = e^{-x}(1 + e^x)^2$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}\mu y &= 0 \\ \frac{d}{dx}(e^{-x}(1 + e^x)^2 y) &= 0\end{aligned}$$

Integrating gives

$$e^{-x}(1 + e^x)^2 y = c_1$$

Dividing both sides by the integrating factor $\mu = e^{-x}(1 + e^x)^2$ results in

$$y = \frac{c_1 e^x}{(1 + e^x)^2}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 e^x}{(1 + e^x)^2} \quad (1)$$

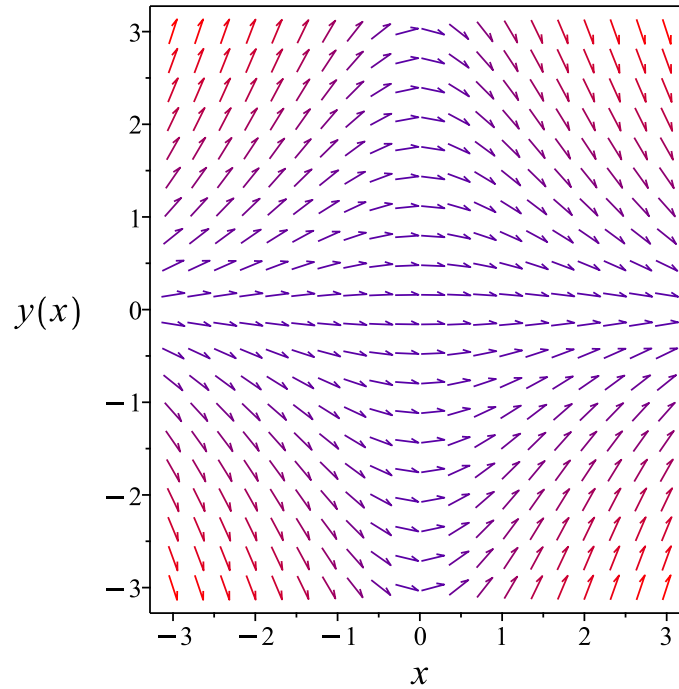


Figure 310: Slope field plot

Verification of solutions

$$y = \frac{c_1 e^x}{(1 + e^x)^2}$$

Verified OK.

6.15.3 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)$ on the above ode results in new ode in $u(x)$

$$(1 + e^x) (u'(x) x + u(x)) - u(x) x + e^x u(x) x = 0$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{u(x e^x + e^x - x + 1)}{(1 + e^x)x}\end{aligned}$$

Where $f(x) = -\frac{x e^x + e^x - x + 1}{x(1 + e^x)}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= -\frac{x e^x + e^x - x + 1}{x(1 + e^x)} dx \\ \int \frac{1}{u} du &= \int -\frac{x e^x + e^x - x + 1}{x(1 + e^x)} dx \\ \ln(u) &= x - \ln(x) - 2 \ln(1 + e^x) + c_2 \\ u &= e^{x - \ln(x) - 2 \ln(1 + e^x) + c_2} \\ &= c_2 e^{x - \ln(x) - 2 \ln(1 + e^x)}\end{aligned}$$

Which simplifies to

$$u(x) = \frac{c_2 e^x}{x(1 + e^x)^2}$$

Therefore the solution y is

$$\begin{aligned}y &= ux \\ &= \frac{c_2 e^x}{(1 + e^x)^2}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_2 e^x}{(1 + e^x)^2} \tag{1}$$

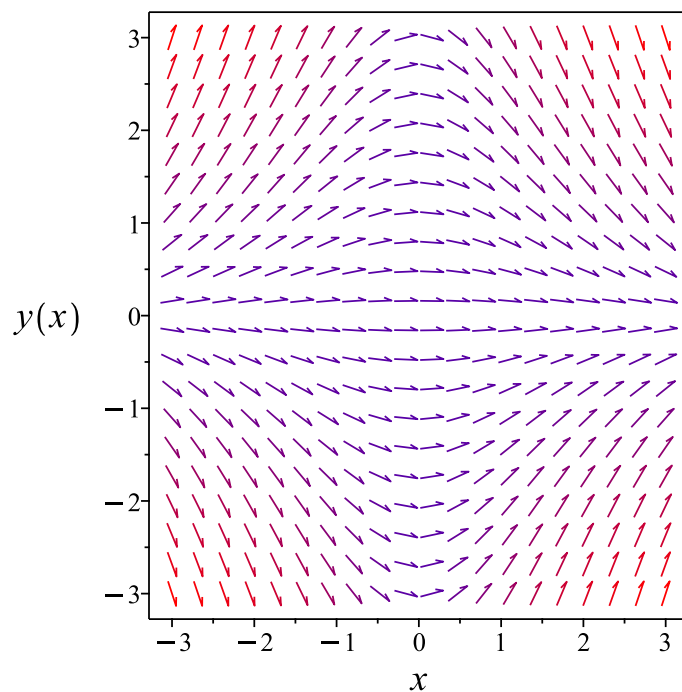


Figure 311: Slope field plot

Verification of solutions

$$y = \frac{c_2 e^x}{(1 + e^x)^2}$$

Verified OK.

6.15.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{y(e^x - 1)}{1 + e^x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 289: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= e^{\ln(e^x) - 2\ln(1+e^x)}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{\ln(e^x) - 2\ln(1+e^x)}} dy \end{aligned}$$

Which results in

$$S = e^{-x}(1 + e^x)^2 y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y(e^x - 1)}{1 + e^x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= y(e^x - e^{-x}) \\ S_y &= e^{-x}(1 + e^x)^2 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$e^{-x}(1 + e^x)^2 y = c_1$$

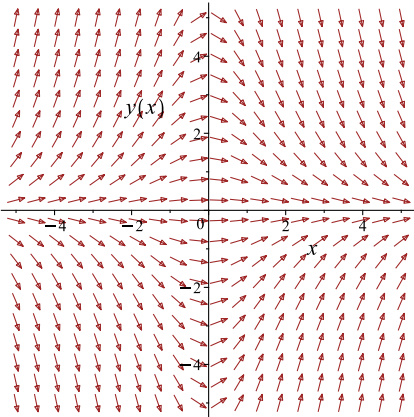
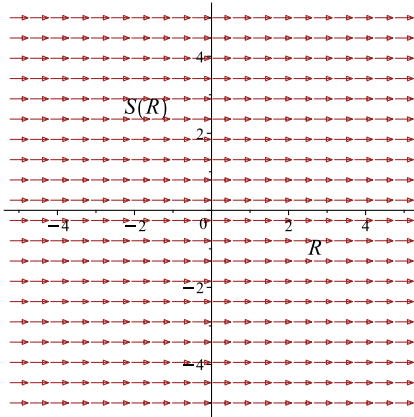
Which simplifies to

$$e^{-x}(1 + e^x)^2 y = c_1$$

Which gives

$$y = \frac{c_1 e^x}{e^{2x} + 2e^x + 1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{y(e^x - 1)}{1 + e^x}$ 	$R = x$ $S = e^{-x}(1 + e^x)^2 y$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$y = \frac{c_1 e^x}{e^{2x} + 2e^x + 1} \quad (1)$$

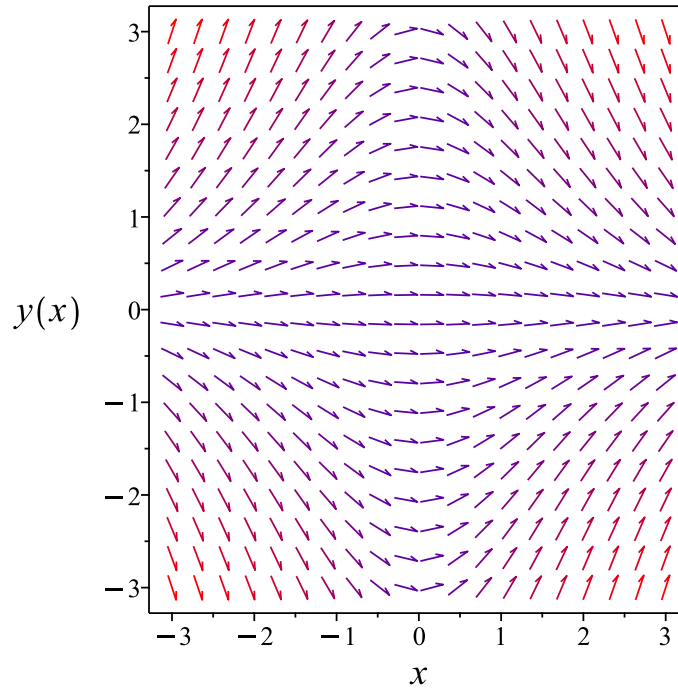


Figure 312: Slope field plot

Verification of solutions

$$y = \frac{c_1 e^x}{e^{2x} + 2e^x + 1}$$

Verified OK.

6.15.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(-\frac{1}{y}\right) dy &= \left(\frac{e^x - 1}{1 + e^x}\right) dx \\ \left(-\frac{e^x - 1}{1 + e^x}\right) dx + \left(-\frac{1}{y}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{e^x - 1}{1 + e^x} \\ N(x, y) &= -\frac{1}{y}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{e^x - 1}{1 + e^x}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(-\frac{1}{y} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{e^x - 1}{1 + e^x} dx \\ \phi &= \ln(e^x) - 2 \ln(1 + e^x) + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{1}{y}$. Therefore equation (4) becomes

$$-\frac{1}{y} = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = -\frac{1}{y}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(-\frac{1}{y} \right) dy \\ f(y) &= -\ln(y) + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = \ln(e^x) - 2\ln(1 + e^x) - \ln(y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \ln(e^x) - 2\ln(1 + e^x) - \ln(y)$$

The solution becomes

$$y = \frac{e^{x-c_1}}{(1 + e^x)^2}$$

Summary

The solution(s) found are the following

$$y = \frac{e^{x-c_1}}{(1 + e^x)^2} \tag{1}$$

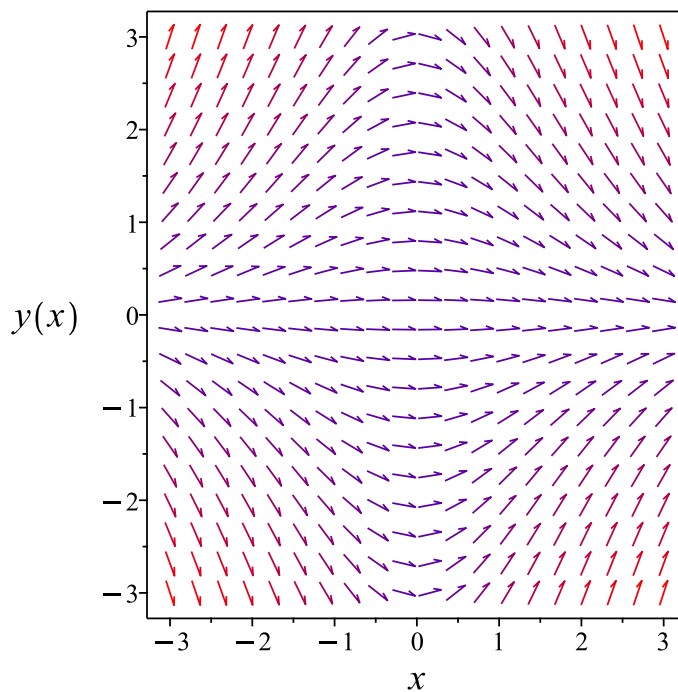


Figure 313: Slope field plot

Verification of solutions

$$y = \frac{e^{x-c_1}}{(1+e^x)^2}$$

Verified OK.

6.15.6 Maple step by step solution

Let's solve

$$(1+e^x)y' - y + e^xy = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{y} = -\frac{e^x-1}{1+e^x}$$

- Integrate both sides with respect to x

$$\int \frac{y'}{y} dx = \int -\frac{e^x-1}{1+e^x} dx + c_1$$

- Evaluate integral

$$\ln(y) = \ln(e^x) - 2\ln(1+e^x) + c_1$$

- Solve for y

$$y = \frac{e^{x+c_1}}{(1+e^x)^2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve((1+exp(x))*diff(y(x),x) = y(x)-exp(x)*y(x),y(x), singsol=all)
```

$$y(x) = \frac{c_1 e^x}{(1 + e^x)^2}$$

✓ Solution by Mathematica

Time used: 0.077 (sec). Leaf size: 23

```
DSolve[(1+Exp[x])*y'[x]== y[x]-Exp[x]*y[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{c_1 e^x}{(e^x + 1)^2}$$
$$y(x) \rightarrow 0$$

6.16 problem 16

6.16.1 Solving as exact ode 1588

Internal problem ID [583]

Internal file name [OUTPUT/583_Sunday_June_05_2022_01_45_12_AM_88473674/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 16.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exactWithIntegrationFactor"**

Maple gives the following as the ode type

[NONE]

$$y' - \frac{-e^{2y} \cos(x) + \cos(y) e^{-x}}{2 e^{2y} \sin(x) - \sin(y) e^{-x}} = 0$$

6.16.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (2 e^{2y} \sin(x) e^x - \sin(y)) dy &= (-e^{2y} \cos(x) e^x + \cos(y)) dx \\ (e^{2y} \cos(x) e^x - \cos(y)) dx &+ (2 e^{2y} \sin(x) e^x - \sin(y)) dy = 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= e^{2y} \cos(x) e^x - \cos(y) \\ N(x, y) &= 2 e^{2y} \sin(x) e^x - \sin(y) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (e^{2y} \cos(x) e^x - \cos(y)) \\ &= 2 \cos(x) e^{x+2y} + \sin(y) \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (2 e^{2y} \sin(x) e^x - \sin(y)) \\ &= 2(\cos(x) + \sin(x)) e^{x+2y} \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{2 \sin(x) e^{x+2y} - \sin(y)} \left((2 e^{2y} \cos(x) e^x + \sin(y)) - (2 e^{2y} \cos(x) e^x + 2 e^{2y} \sin(x) e^x) \right) \\ &= -1 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int -1 dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-x} \\ &= e^{-x}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= e^{-x} (e^{2y} \cos(x) e^x - \cos(y)) \\ &= (\cos(x) e^{x+2y} - \cos(y)) e^{-x}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^{-x} (2 e^{2y} \sin(x) e^x - \sin(y)) \\ &= (2 \sin(x) e^{x+2y} - \sin(y)) e^{-x}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ ((\cos(x) e^{x+2y} - \cos(y)) e^{-x}) + ((2 \sin(x) e^{x+2y} - \sin(y)) e^{-x}) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int (\cos(x) e^{x+2y} - \cos(y)) e^{-x} dx \\ \phi &= e^{2y} \sin(x) + \cos(y) e^{-x} + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 2 e^{2y} \sin(x) - \sin(y) e^{-x} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = (2 \sin(x) e^{x+2y} - \sin(y)) e^{-x}$. Therefore equation (4) becomes

$$(2 \sin(x) e^{x+2y} - \sin(y)) e^{-x} = 2 e^{2y} \sin(x) - \sin(y) e^{-x} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = e^{2y} \sin(x) + \cos(y) e^{-x} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = e^{2y} \sin(x) + \cos(y) e^{-x}$$

Summary

The solution(s) found are the following

$$e^{2y} \sin(x) + \cos(y) e^{-x} = c_1 \quad (1)$$

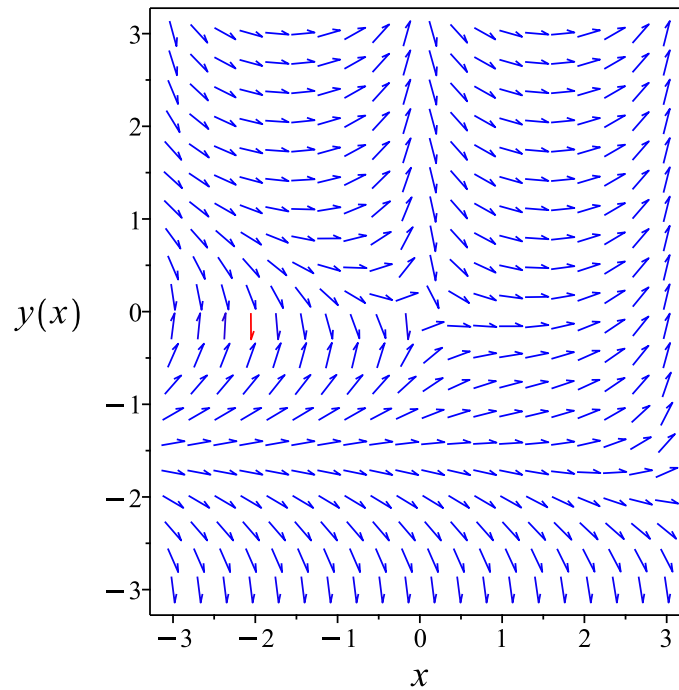


Figure 314: Slope field plot

Verification of solutions

$$e^{2y} \sin(x) + \cos(y) e^{-x} = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 21

```
dsolve(diff(y(x),x) = (-exp(2*y(x))*cos(x)+cos(y(x))/exp(x))/(2*exp(2*y(x))*sin(x)-sin(y(x)))
```

$$c_1 + \cos(y(x)) e^{-x} + e^{2y(x)} \sin(x) = 0$$

✓ Solution by Mathematica

Time used: 0.473 (sec). Leaf size: 25

```
DSolve[y' [x] == (-Exp [2*y [x]]*Cos [x]+Cos [y [x]]/Exp [x])/(2*Exp [2*y [x]]*Sin [x]-Sin [y [x]]/Exp [x]
```

$$\text{Solve}[e^{2y(x)} \sin(x) + e^{-x} \cos(y(x)) = c_1, y(x)]$$

6.17 problem 17

6.17.1 Solving as linear ode	1594
6.17.2 Solving as first order ode lie symmetry lookup ode	1596
6.17.3 Solving as exact ode	1600
6.17.4 Maple step by step solution	1604

Internal problem ID [584]

Internal file name [OUTPUT/584_Sunday_June_05_2022_01_45_16_AM_73508669/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 17.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$y' - 3y = e^{2x}$$

6.17.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -3$$
$$q(x) = e^{2x}$$

Hence the ode is

$$y' - 3y = e^{2x}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int (-3) dx} \\ &= e^{-3x}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) (e^{2x}) \\ \frac{d}{dx}(e^{-3x}y) &= (e^{-3x}) (e^{2x}) \\ d(e^{-3x}y) &= e^{-x} dx\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{-3x}y &= \int e^{-x} dx \\ e^{-3x}y &= -e^{-x} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{-3x}$ results in

$$y = -e^{3x}e^{-x} + c_1e^{3x}$$

which simplifies to

$$y = -e^{2x} + c_1e^{3x}$$

Summary

The solution(s) found are the following

$$y = -e^{2x} + c_1e^{3x} \tag{1}$$

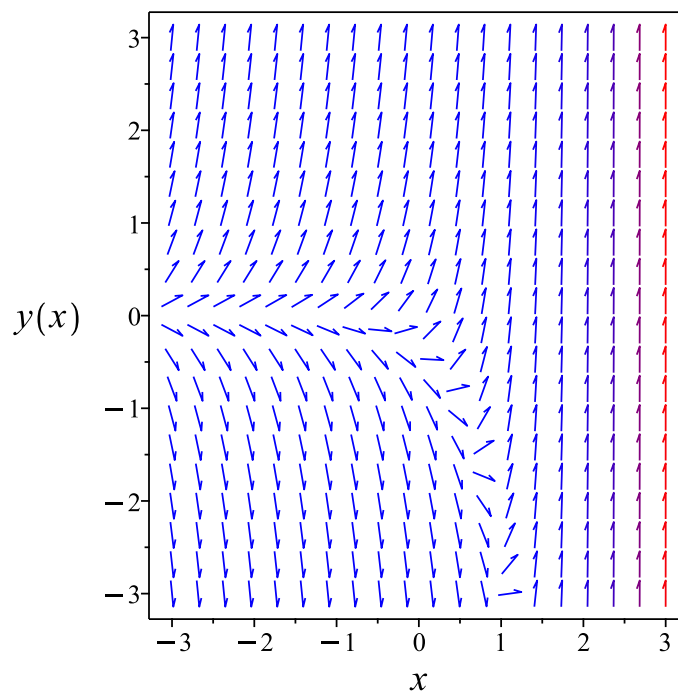


Figure 315: Slope field plot

Verification of solutions

$$y = -e^{2x} + c_1 e^{3x}$$

Verified OK.

6.17.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = e^{2x} + 3y$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 292: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= e^{3x}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{3x}} dy \end{aligned}$$

Which results in

$$S = e^{-3x}y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = e^{2x} + 3y$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -3e^{-3x}y \\ S_y &= e^{-3x} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^{-x} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^{-R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -e^{-R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$y e^{-3x} = -e^{-x} + c_1$$

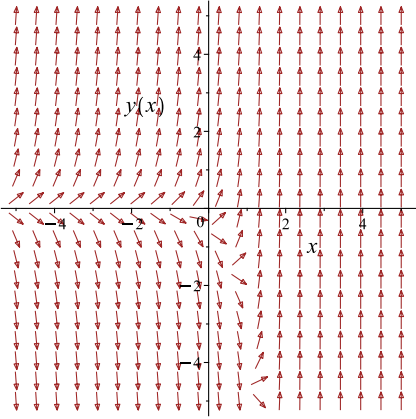
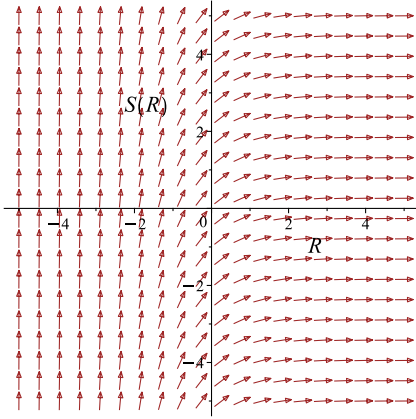
Which simplifies to

$$y e^{-3x} = -e^{-x} + c_1$$

Which gives

$$y = -(e^{-x} - c_1) e^{3x}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = e^{2x} + 3y$ 	$R = x$ $S = e^{-3x}y$	$\frac{dS}{dR} = e^{-R}$ 

Summary

The solution(s) found are the following

$$y = -(e^{-x} - c_1) e^{3x} \quad (1)$$

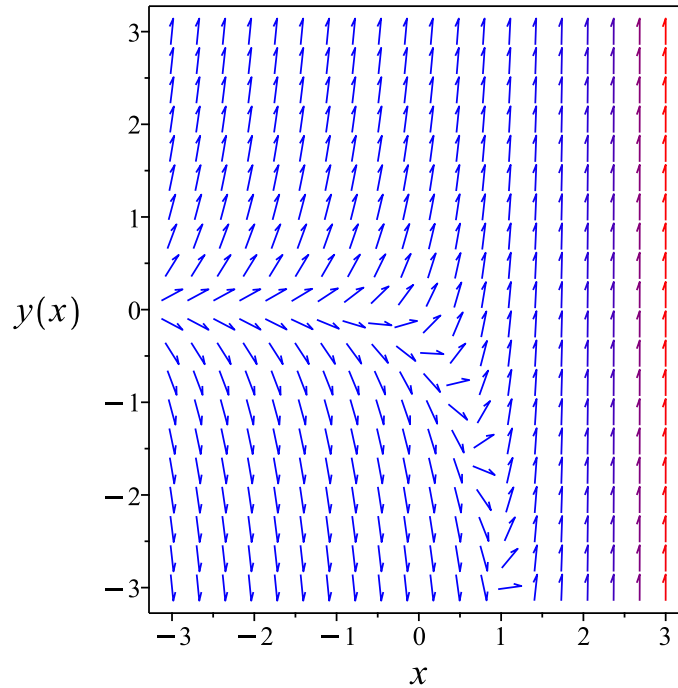


Figure 316: Slope field plot

Verification of solutions

$$y = -(e^{-x} - c_1) e^{3x}$$

Verified OK.

6.17.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}dy &= (e^{2x} + 3y) dx \\ (-e^{2x} - 3y) dx + dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -e^{2x} - 3y \\ N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-e^{2x} - 3y) \\ &= -3\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((-3) - (0)) \\ &= -3 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dx} \\ &= e^{\int -3 dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-3x} \\ &= e^{-3x} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{-3x}(-e^{2x} - 3y) \\ &= (-e^{2x} - 3y)e^{-3x} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{-3x}(1) \\ &= e^{-3x} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ ((-e^{2x} - 3y)e^{-3x}) + (e^{-3x}) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int (-e^{2x} - 3y) e^{-3x} dx \\ \phi &= e^{-x} + e^{-3x}y + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{-3x} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-3x}$. Therefore equation (4) becomes

$$e^{-3x} = e^{-3x} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = e^{-x} + e^{-3x}y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = e^{-x} + e^{-3x}y$$

The solution becomes

$$y = -(e^{-x} - c_1) e^{3x}$$

Summary

The solution(s) found are the following

$$y = -(e^{-x} - c_1) e^{3x} \quad (1)$$

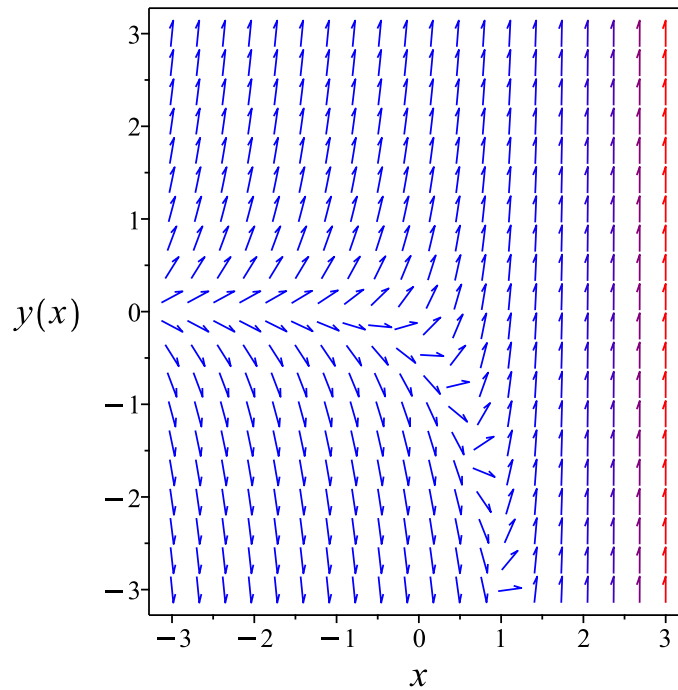


Figure 317: Slope field plot

Verification of solutions

$$y = -(e^{-x} - c_1) e^{3x}$$

Verified OK.

6.17.4 Maple step by step solution

Let's solve

$$y' - 3y = e^{2x}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = e^{2x} + 3y$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' - 3y = e^{2x}$$

- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$\mu(x) (y' - 3y) = \mu(x) e^{2x}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dx}(\mu(x)y)$

$$\mu(x)(y' - 3y) = \mu'(x)y + \mu(x)y'$$
- Isolate $\mu'(x)$

$$\mu'(x) = -3\mu(x)$$
- Solve to find the integrating factor

$$\mu(x) = e^{-3x}$$
- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x)y) \right) dx = \int \mu(x) e^{2x} dx + c_1$$
- Evaluate the integral on the lhs

$$\mu(x)y = \int \mu(x) e^{2x} dx + c_1$$
- Solve for y

$$y = \frac{\int \mu(x) e^{2x} dx + c_1}{\mu(x)}$$
- Substitute $\mu(x) = e^{-3x}$

$$y = \frac{\int e^{2x} e^{-3x} dx + c_1}{e^{-3x}}$$
- Evaluate the integrals on the rhs

$$y = \frac{-e^{-x} + c_1}{e^{-3x}}$$
- Simplify

$$y = e^{2x}(c_1 e^x - 1)$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve(diff(y(x),x) = exp(2*x)+3*y(x),y(x), singsol=all)
```

$$y(x) = (e^x c_1 - 1) e^{2x}$$

✓ Solution by Mathematica

Time used: 0.04 (sec). Leaf size: 19

```
DSolve[y'[x]== Exp[2*x]+3*y[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{2x}(-1 + c_1 e^x)$$

6.18 problem 18

6.18.1 Solving as linear ode	1607
6.18.2 Solving as first order ode lie symmetry lookup ode	1609
6.18.3 Solving as exact ode	1613
6.18.4 Maple step by step solution	1617

Internal problem ID [585]

Internal file name [OUTPUT/585_Sunday_June_05_2022_01_45_17_AM_81100757/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 18.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$2y + y' = e^{-x^2-2x}$$

6.18.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = 2$$
$$q(x) = e^{-x(2+x)}$$

Hence the ode is

$$2y + y' = e^{-x(2+x)}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int 2dx} \\ &= e^{2x}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) (e^{-x(2+x)}) \\ \frac{d}{dx}(e^{2x}y) &= (e^{2x}) (e^{-x(2+x)}) \\ d(e^{2x}y) &= e^{-x^2} dx\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{2x}y &= \int e^{-x^2} dx \\ e^{2x}y &= \frac{\sqrt{\pi} \operatorname{erf}(x)}{2} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = e^{2x}$ results in

$$y = \frac{e^{-2x} \sqrt{\pi} \operatorname{erf}(x)}{2} + c_1 e^{-2x}$$

which simplifies to

$$y = e^{-2x} \left(\frac{\sqrt{\pi} \operatorname{erf}(x)}{2} + c_1 \right)$$

Summary

The solution(s) found are the following

$$y = e^{-2x} \left(\frac{\sqrt{\pi} \operatorname{erf}(x)}{2} + c_1 \right) \tag{1}$$

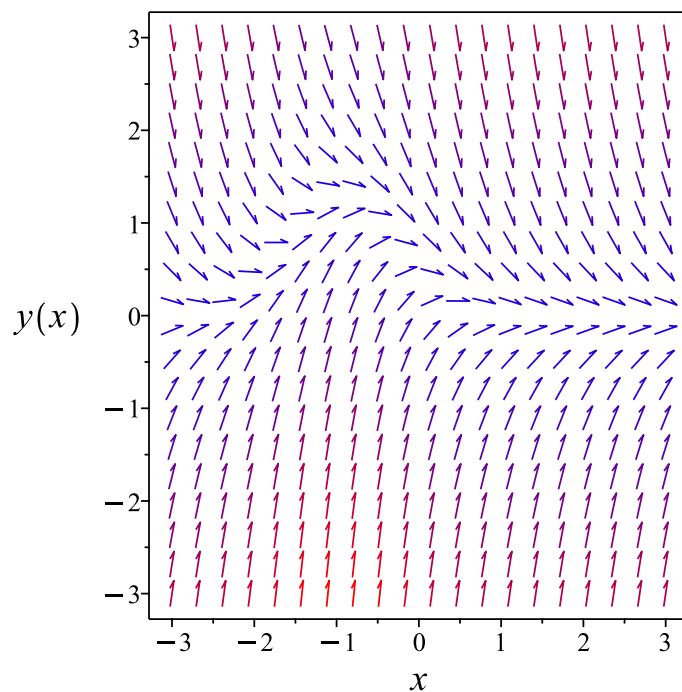


Figure 318: Slope field plot

Verification of solutions

$$y = e^{-2x} \left(\frac{\sqrt{\pi} \operatorname{erf}(x)}{2} + c_1 \right)$$

Verified OK.

6.18.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -2y + e^{-x^2-2x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 295: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= e^{-2x}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-2x}} dy \end{aligned}$$

Which results in

$$S = e^{2x}y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -2y + e^{-x^2-2x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= 2e^{2x}y \\ S_y &= e^{2x} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^{-x^2} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^{-R^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{\sqrt{\pi} \operatorname{erf}(R)}{2} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$y e^{2x} = \frac{\sqrt{\pi} \operatorname{erf}(x)}{2} + c_1$$

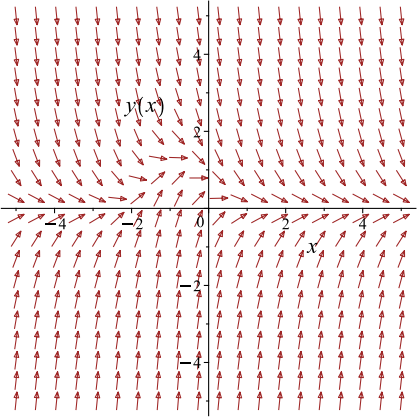
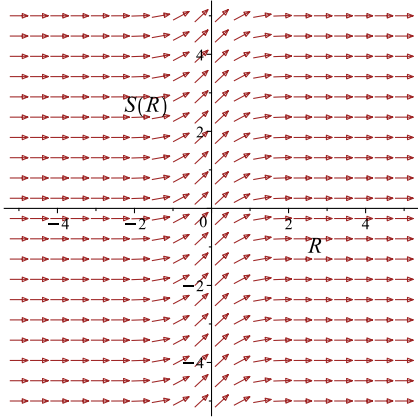
Which simplifies to

$$y e^{2x} = \frac{\sqrt{\pi} \operatorname{erf}(x)}{2} + c_1$$

Which gives

$$y = \frac{e^{-2x} (\sqrt{\pi} \operatorname{erf}(x) + 2c_1)}{2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -2y + e^{-x^2-2x}$ 	$R = x$ $S = e^{2x} y$	$\frac{dS}{dR} = e^{-R^2}$ 

Summary

The solution(s) found are the following

$$y = \frac{e^{-2x} (\sqrt{\pi} \operatorname{erf}(x) + 2c_1)}{2} \quad (1)$$

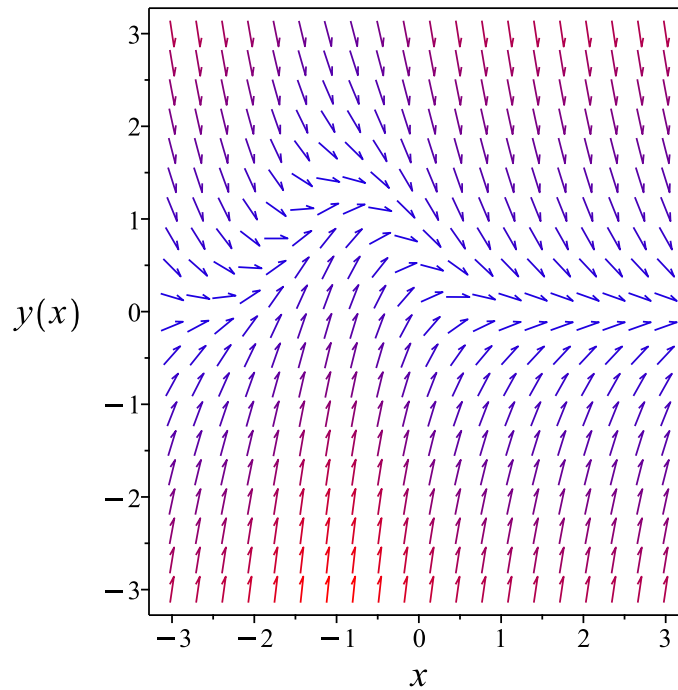


Figure 319: Slope field plot

Verification of solutions

$$y = \frac{e^{-2x}(\sqrt{\pi} \operatorname{erf}(x) + 2c_1)}{2}$$

Verified OK.

6.18.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dy &= \left(-2y + e^{-x^2-2x}\right) dx \\ \left(2y - e^{-x^2-2x}\right) dx + dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 2y - e^{-x^2-2x} \\ N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(2y - e^{-x^2-2x}\right) \\ &= 2\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((2) - (0)) \\ &= 2 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dx} \\ &= e^{\int 2 dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{2x} \\ &= e^{2x} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{2x} (2y - e^{-x^2-2x}) \\ &= (2y - e^{-x(2+x)}) e^{2x} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{2x} (1) \\ &= e^{2x} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ ((2y - e^{-x(2+x)}) e^{2x}) + (e^{2x}) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \bar{M} dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int (2y - e^{-x(2+x)}) e^{2x} dx$$

$$\phi = -\frac{\sqrt{\pi} \operatorname{erf}(x)}{2} + e^{2x}y + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = e^{2x} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{2x}$. Therefore equation (4) becomes

$$e^{2x} = e^{2x} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{\sqrt{\pi} \operatorname{erf}(x)}{2} + e^{2x}y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{\sqrt{\pi} \operatorname{erf}(x)}{2} + e^{2x}y$$

The solution becomes

$$y = \frac{e^{-2x}(\sqrt{\pi} \operatorname{erf}(x) + 2c_1)}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{e^{-2x}(\sqrt{\pi} \operatorname{erf}(x) + 2c_1)}{2} \quad (1)$$

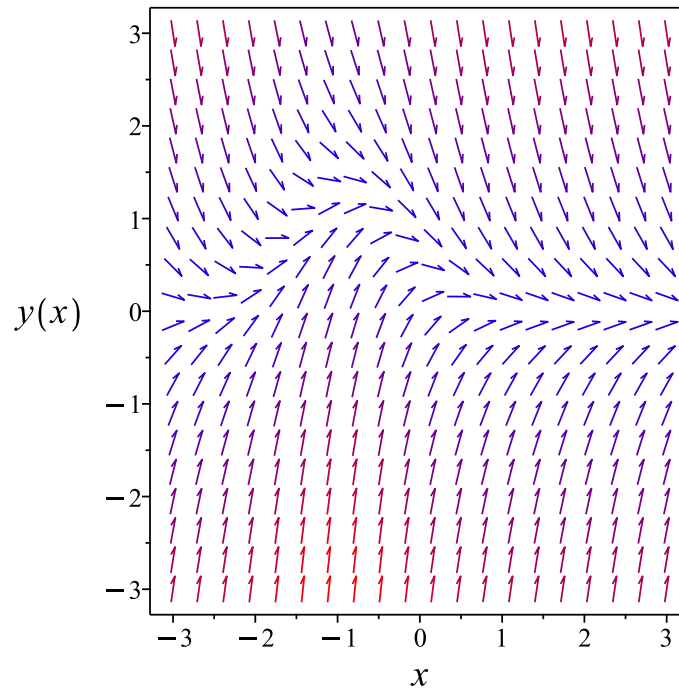


Figure 320: Slope field plot

Verification of solutions

$$y = \frac{e^{-2x}(\sqrt{\pi} \operatorname{erf}(x) + 2c_1)}{2}$$

Verified OK.

6.18.4 Maple step by step solution

Let's solve

$$2y + y' = e^{-x^2-2x}$$

- Highest derivative means the order of the ODE is 1
- y'
- Isolate the derivative

$$y' = -2y + e^{-x^2-2x}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$2y + y' = e^{-x^2-2x}$$

- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$\mu(x) (2y + y') = \mu(x) e^{-x^2-2x}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dx}(\mu(x) y)$

$$\mu(x) (2y + y') = \mu'(x) y + \mu(x) y'$$

- Isolate $\mu'(x)$

$$\mu'(x) = 2\mu(x)$$

- Solve to find the integrating factor

$$\mu(x) = e^{2x}$$

- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(\mu(x) y) \right) dx = \int \mu(x) e^{-x^2-2x} dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x) y = \int \mu(x) e^{-x^2-2x} dx + c_1$$

- Solve for y

$$y = \frac{\int \mu(x) e^{-x^2-2x} dx + c_1}{\mu(x)}$$

- Substitute $\mu(x) = e^{2x}$

$$y = \frac{\int e^{2x} e^{-x^2-2x} dx + c_1}{e^{2x}}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{\sqrt{\pi} \operatorname{erf}(x)}{2} + c_1}{e^{2x}}$$

- Simplify

$$y = \frac{e^{-2x} (\sqrt{\pi} \operatorname{erf}(x) + 2c_1)}{2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

```
dsolve(2*y(x)+diff(y(x),x) = exp(-x^2-2*x),y(x), singsol=all)
```

$$y(x) = \frac{(\sqrt{\pi} \operatorname{erf}(x) + 2c_1) e^{-2x}}{2}$$

✓ Solution by Mathematica

Time used: 0.09 (sec). Leaf size: 27

```
DSolve[2*y[x]+y'[x] == Exp[-x^2-2*x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2} e^{-2x} (\sqrt{\pi} \operatorname{erf}(x) + 2c_1)$$

6.19 problem 19

6.19.1 Solving as exact ode 1620

Internal problem ID [586]

Internal file name [OUTPUT/586_Sunday_June_05_2022_01_45_18_AM_94339622/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 19.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact"**

Maple gives the following as the ode type

`[_rational]`

$$y' - \frac{3x^2 - 2y - y^3}{2x + 3xy^2} = 0$$

6.19.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (x(3y^2 + 2)) dy &= (-y^3 + 3x^2 - 2y) dx \\ (y^3 - 3x^2 + 2y) dx + (x(3y^2 + 2)) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= y^3 - 3x^2 + 2y \\ N(x, y) &= x(3y^2 + 2) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (y^3 - 3x^2 + 2y) \\ &= 3y^2 + 2 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (x(3y^2 + 2)) \\ &= 3y^2 + 2 \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int y^3 - 3x^2 + 2y dx \\ \phi &= -x(-y^3 + x^2 - 2y) + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned}\frac{\partial \phi}{\partial y} &= -x(-3y^2 - 2) + f'(y) \\ &= x(3y^2 + 2) + f'(y)\end{aligned}\quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = x(3y^2 + 2)$. Therefore equation (4) becomes

$$x(3y^2 + 2) = x(3y^2 + 2) + f'(y)\quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -x(-y^3 + x^2 - 2y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -x(-y^3 + x^2 - 2y)$$

Summary

The solution(s) found are the following

$$-x(-y^3 + x^2 - 2y) = c_1\quad (1)$$

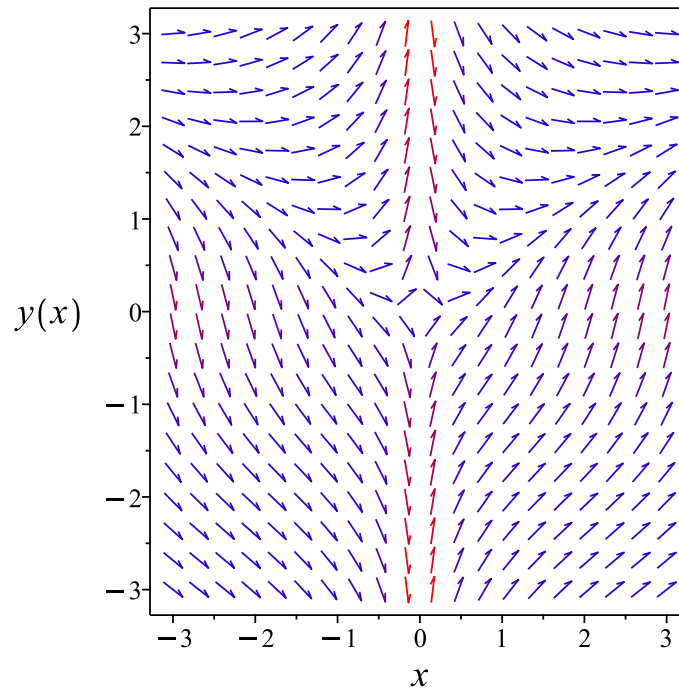


Figure 321: Slope field plot

Verification of solutions

$$-x(-y^3 + x^2 - 2y) = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```


✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 409

```
dsolve(diff(y(x),x) = (3*x^2-2*y(x)-y(x)^3)/(2*x+3*x*y(x)^2),y(x), singsol=all)
```

$$y(x) = -\frac{12^{\frac{1}{3}} \left(x^2 12^{\frac{1}{3}} - \frac{\left((9x^3 + \sqrt{3} \sqrt{27x^6 - 54c_1x^3 + 27c_1^2 + 32x^2 - 9c_1}) x^2 \right)^{\frac{2}{3}}}{2} \right)}{3 \left((9x^3 + \sqrt{3} \sqrt{27x^6 - 54c_1x^3 + 27c_1^2 + 32x^2 - 9c_1}) x^2 \right)^{\frac{1}{3}} x}$$

$$y(x) = \frac{2^{\frac{2}{3}} 3^{\frac{1}{3}} \left(2i 2^{\frac{2}{3}} 3^{\frac{5}{6}} x^2 - 2x^2 2^{\frac{2}{3}} 3^{\frac{1}{3}} + i \left((9x^3 + \sqrt{3} \sqrt{27x^6 - 54c_1x^3 + 27c_1^2 + 32x^2 - 9c_1}) x^2 \right)^{\frac{2}{3}} \sqrt{3} + \left((9x^3 + \sqrt{3} \sqrt{27x^6 - 54c_1x^3 + 27c_1^2 + 32x^2 - 9c_1}) x^2 \right)^{\frac{2}{3}} \right)}{12x \left((9x^3 + \sqrt{3} \sqrt{27x^6 - 54c_1x^3 + 27c_1^2 + 32x^2 - 9c_1}) x^2 \right)^{\frac{1}{3}}}$$

$$y(x) = \frac{2^{\frac{2}{3}} 3^{\frac{1}{3}} \left(2 \left(i 3^{\frac{5}{6}} + 3^{\frac{1}{3}} \right) x^2 2^{\frac{2}{3}} + (i\sqrt{3} - 1) \left((9x^3 + \sqrt{3} \sqrt{27x^6 - 54c_1x^3 + 27c_1^2 + 32x^2 - 9c_1}) x^2 \right)^{\frac{2}{3}} \right)}{12 \left((9x^3 + \sqrt{3} \sqrt{27x^6 - 54c_1x^3 + 27c_1^2 + 32x^2 - 9c_1}) x^2 \right)^{\frac{1}{3}} x}$$

✓ Solution by Mathematica

Time used: 32.075 (sec). Leaf size: 358

`DSolve[y'[x] == (3*x^2-2*y[x]-y[x]^3)/(2*x+3*x*y[x]^2),y[x],x,IncludeSingularSolutions -> Tr`

$$y(x) \rightarrow \frac{\sqrt[3]{27x^5 + 27c_1x^2 + \sqrt{864x^6 + 729x^4(x^3 + c_1)^2}}}{3\sqrt[3]{2x}} - \frac{2\sqrt[3]{2x}}{\sqrt[3]{27x^5 + 27c_1x^2 + \sqrt{864x^6 + 729x^4(x^3 + c_1)^2}}}$$

$$y(x) \rightarrow \frac{\sqrt[3]{27x^5 + 27c_1x^2 + \sqrt{864x^6 + 729x^4(x^3 + c_1)^2}}}{\sqrt[3]{2}(1 + i\sqrt{3})x} - \frac{(1 - i\sqrt{3})\sqrt[3]{27x^5 + 27c_1x^2 + \sqrt{864x^6 + 729x^4(x^3 + c_1)^2}}}{6\sqrt[3]{2x}}$$

$$y(x) \rightarrow \frac{\sqrt[3]{27x^5 + 27c_1x^2 + \sqrt{864x^6 + 729x^4(x^3 + c_1)^2}}}{\sqrt[3]{2}(1 - i\sqrt{3})x} - \frac{(1 + i\sqrt{3})\sqrt[3]{27x^5 + 27c_1x^2 + \sqrt{864x^6 + 729x^4(x^3 + c_1)^2}}}{6\sqrt[3]{2x}}$$

6.20 problem 20

6.20.1 Solving as separable ode	1626
6.20.2 Solving as first order special form ID 1 ode	1628
6.20.3 Solving as first order ode lie symmetry lookup ode	1629
6.20.4 Solving as exact ode	1633
6.20.5 Maple step by step solution	1637

Internal problem ID [587]

Internal file name [OUTPUT/587_Sunday_June_05_2022_01_45_20_AM_44183904/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 20.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "first order special form ID 1", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - e^{x+y} = 0$$

6.20.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= e^x e^y\end{aligned}$$

Where $f(x) = e^x$ and $g(y) = e^y$. Integrating both sides gives

$$\begin{aligned}\frac{1}{e^y} dy &= e^x dx \\ \int \frac{1}{e^y} dy &= \int e^x dx \\ -e^{-y} &= e^x + c_1\end{aligned}$$

Which results in

$$y = \ln \left(-\frac{1}{e^x + c_1} \right)$$

Summary

The solution(s) found are the following

$$y = \ln \left(-\frac{1}{e^x + c_1} \right) \tag{1}$$

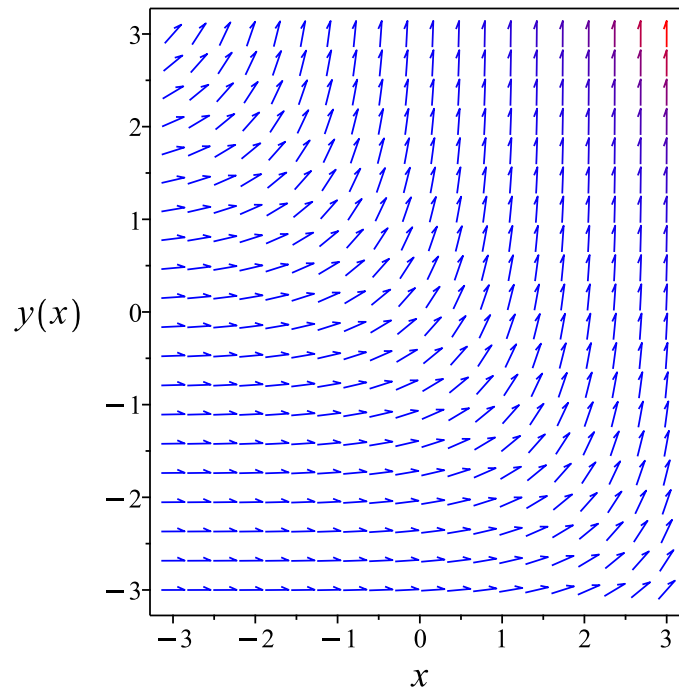


Figure 322: Slope field plot

Verification of solutions

$$y = \ln \left(-\frac{1}{e^x + c_1} \right)$$

Verified OK.

6.20.2 Solving as first order special form ID 1 ode

Writing the ode as

$$y' = e^{x+y} \quad (1)$$

And using the substitution $u = e^{-y}$ then

$$u' = -y'e^{-y}$$

The above shows that

$$\begin{aligned} y' &= -u'(x) e^y \\ &= -\frac{u'(x)}{u} \end{aligned}$$

Substituting this in (1) gives

$$-\frac{u'(x)}{u} = \frac{e^x}{u}$$

The above simplifies to

$$u'(x) = -e^x \quad (2)$$

Now ode (2) is solved for $u(x)$ Integrating both sides gives

$$\begin{aligned} u(x) &= \int -e^x dx \\ &= -e^x + c_1 \end{aligned}$$

Substituting the solution found for $u(x)$ in $u = e^{-y}$ gives

$$\begin{aligned} y &= -\ln(u(x)) \\ &= -\ln(-e^x + c_1) \\ &= -\ln(-e^x + c_1) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = -\ln(-e^x + c_1) \quad (1)$$

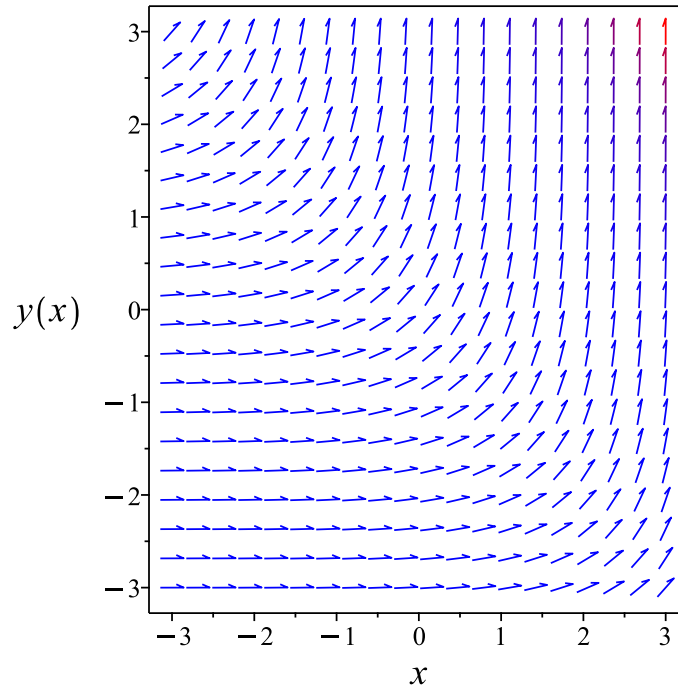


Figure 323: Slope field plot

Verification of solutions

$$y = -\ln(-e^x + c_1)$$

Verified OK.

6.20.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = e^{x+y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 298: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= e^{-x} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{e^{-x}} dx \end{aligned}$$

Which results in

$$S = e^x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = e^{x+y}$$

Evaluating all the partial derivatives gives

$$R_x = 0$$

$$R_y = 1$$

$$S_x = e^x$$

$$S_y = 0$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^{-y} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^{-R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -e^{-R} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$e^x = -e^{-y} + c_1$$

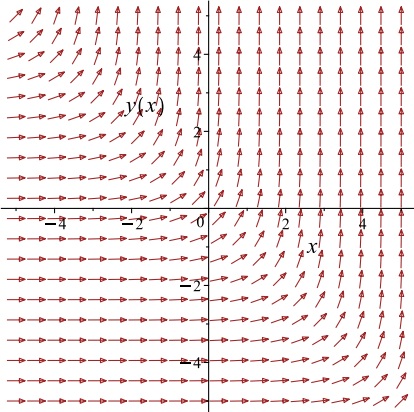
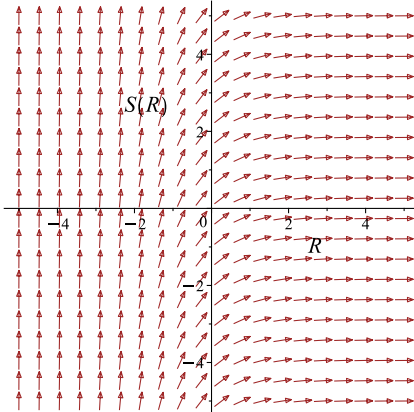
Which simplifies to

$$e^x = -e^{-y} + c_1$$

Which gives

$$y = -\ln(-e^x + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = e^{x+y}$ 	$R = y$ $S = e^x$	$\frac{dS}{dR} = e^{-R}$ 

Summary

The solution(s) found are the following

$$y = -\ln(-e^x + c_1) \quad (1)$$

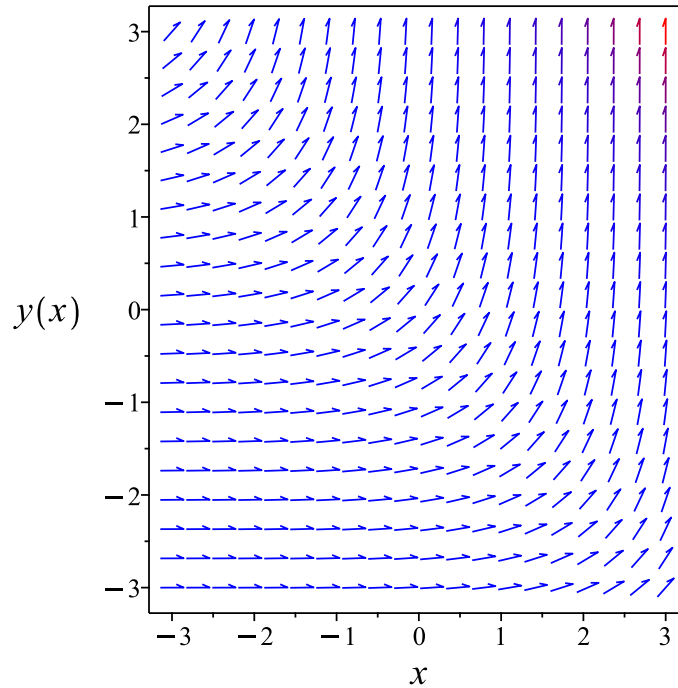


Figure 324: Slope field plot

Verification of solutions

$$y = -\ln(-e^x + c_1)$$

Verified OK.

6.20.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(e^{-y}) dy &= (e^x) dx \\ (-e^x) dx + (e^{-y}) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -e^x \\ N(x, y) &= e^{-y}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-e^x) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(e^{-y}) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -e^x dx \\ \phi &= -e^x + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = e^{-y}$. Therefore equation (4) becomes

$$e^{-y} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = e^{-y}$$

Integrating the above w.r.t y gives

$$\begin{aligned} \int f'(y) dy &= \int (e^{-y}) dy \\ f(y) &= -e^{-y} + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -e^x - e^{-y} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -e^x - e^{-y}$$

The solution becomes

$$y = -\ln(-e^x - c_1)$$

Summary

The solution(s) found are the following

$$y = -\ln(-e^x - c_1) \tag{1}$$

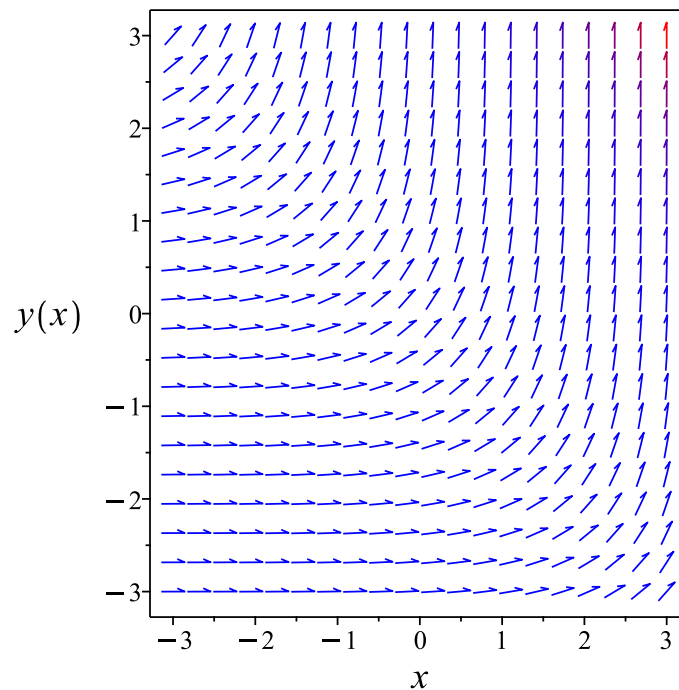


Figure 325: Slope field plot

Verification of solutions

$$y = -\ln(-e^x - c_1)$$

Verified OK.

6.20.5 Maple step by step solution

Let's solve

$$y' - e^{x+y} = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{e^y} = e^x$$

- Integrate both sides with respect to x

$$\int \frac{y'}{e^y} dx = \int e^x dx + c_1$$

- Evaluate integral

$$-\frac{1}{e^y} = e^x + c_1$$

- Solve for y

$$y = \ln\left(-\frac{1}{e^x + c_1}\right)$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 13

```
dsolve(diff(y(x),x) = exp(x+y(x)),y(x), singsol=all)
```

$$y(x) = \ln\left(-\frac{1}{e^x + c_1}\right)$$

✓ Solution by Mathematica

Time used: 0.752 (sec). Leaf size: 18

```
DSolve[y'[x] == Exp[x+y[x]],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\log(-e^x - c_1)$$

6.21 problem 21

6.21.1 Solving as exact ode 1639

Internal problem ID [588]

Internal file name [OUTPUT/588_Sunday_June_05_2022_01_45_20_AM_85449594/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 21.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact"**

Maple gives the following as the ode type

[_rational]

$$\frac{-4 + 6yx + 2y^2}{3x^2 + 4yx + 3y^2} + y' = 0$$

6.21.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (3x^2 + 4yx + 3y^2) dy &= (-6yx - 2y^2 + 4) dx \\ (6yx + 2y^2 - 4) dx + (3x^2 + 4yx + 3y^2) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 6yx + 2y^2 - 4 \\ N(x, y) &= 3x^2 + 4yx + 3y^2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (6yx + 2y^2 - 4) \\ &= 6x + 4y \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (3x^2 + 4yx + 3y^2) \\ &= 6x + 4y \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 6yx + 2y^2 - 4 dx \\ \phi &= 3y x^2 + 2x y^2 - 4x + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned}\frac{\partial \phi}{\partial y} &= 3x^2 + 4yx + f'(y) \\ &= x(3x + 4y) + f'(y)\end{aligned}\tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = 3x^2 + 4yx + 3y^2$. Therefore equation (4) becomes

$$3x^2 + 4yx + 3y^2 = x(3x + 4y) + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 3y^2$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int (3y^2) dy \\ f(y) &= y^3 + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = 3y x^2 + 2x y^2 + y^3 - 4x + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = 3y x^2 + 2x y^2 + y^3 - 4x$$

Summary

The solution(s) found are the following

$$3x^2y + 2xy^2 + y^3 - 4x = c_1 \quad (1)$$

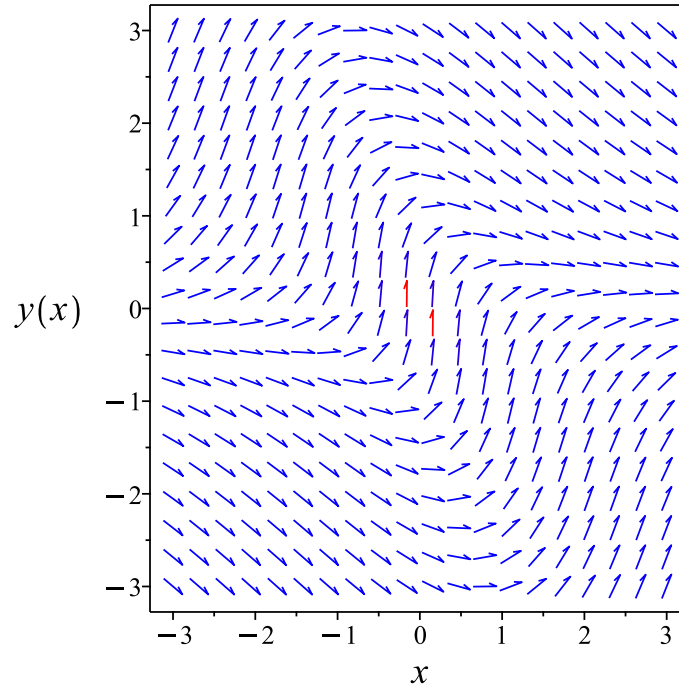


Figure 326: Slope field plot

Verification of solutions

$$3x^2y + 2xy^2 + y^3 - 4x = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 517

`dsolve((-4+6*x*y(x)+2*y(x)^2)/(3*x^2+4*x*y(x)+3*y(x)^2)+diff(y(x),x)=0,y(x), singsol=all)`

$y(x)$

$$= \frac{\left(152x^3 - 108c_1 + 432x + 12\sqrt{216x^6 - 228c_1x^3 + 912x^4 + 81c_1^2 - 648c_1x + 1296x^2}\right)^{\frac{1}{3}}}{6 \cdot 10x^2}$$

$$- \frac{3 \left(152x^3 - 108c_1 + 432x + 12\sqrt{216x^6 - 228c_1x^3 + 912x^4 + 81c_1^2 - 648c_1x + 1296x^2}\right)^{\frac{1}{3}}}{\frac{2x}{3}}$$

$y(x) =$

$$\frac{(1 + i\sqrt{3}) \left(152x^3 - 108c_1 + 432x + 12\sqrt{216x^6 - 228c_1x^3 + 912x^4 + 81c_1^2 - 648c_1x + 1296x^2}\right)^{\frac{1}{3}}}{12}$$

$$- \frac{5x \left(i\sqrt{3}x - x + \frac{2 \left(152x^3 - 108c_1 + 432x + 12\sqrt{216x^6 - 228c_1x^3 + 912x^4 + 81c_1^2 - 648c_1x + 1296x^2}\right)^{\frac{1}{3}}}{5} \right)}{3 \left(152x^3 - 108c_1 + 432x + 12\sqrt{216x^6 - 228c_1x^3 + 912x^4 + 81c_1^2 - 648c_1x + 1296x^2}\right)^{\frac{1}{3}}}$$

$y(x)$

$$= \frac{i \left(152x^3 - 108c_1 + 432x + 12\sqrt{216x^6 - 228c_1x^3 + 912x^4 + 81c_1^2 - 648c_1x + 1296x^2}\right)^{\frac{2}{3}} \sqrt{3} + 20i\sqrt{3}x^2}{12}$$

12

✓ Solution by Mathematica

Time used: 4.748 (sec). Leaf size: 383

`DSolve[(-4+6*x*y[x]+2*y[x]^2)/(3*x^2+4*x*y[x]+3*y[x]^2)+y'[x]==0,y[x],x,IncludeSingularSolut`

$$y(x) \rightarrow \frac{1}{6} \left(2^{2/3} \sqrt[3]{38x^3 + \sqrt{500x^6 + (38x^3 + 108x + 27c_1)^2} + 108x + 27c_1} \right. \\ \left. - \frac{10\sqrt[3]{2}x^2}{\sqrt[3]{38x^3 + \sqrt{500x^6 + (38x^3 + 108x + 27c_1)^2} + 108x + 27c_1}} - 4x \right)$$

$$y(x) \rightarrow \frac{1}{12} \left(i2^{2/3}(\sqrt{3} + i) \sqrt[3]{38x^3 + \sqrt{500x^6 + (38x^3 + 108x + 27c_1)^2} + 108x + 27c_1} \right. \\ \left. + \frac{10\sqrt[3]{2}(1 + i\sqrt{3})x^2}{\sqrt[3]{38x^3 + \sqrt{500x^6 + (38x^3 + 108x + 27c_1)^2} + 108x + 27c_1}} - 8x \right)$$

$$y(x) \rightarrow \frac{1}{12} \left(-2^{2/3}(1 + i\sqrt{3}) \sqrt[3]{38x^3 + \sqrt{500x^6 + (38x^3 + 108x + 27c_1)^2} + 108x + 27c_1} \right. \\ \left. + \frac{10\sqrt[3]{2}(1 - i\sqrt{3})x^2}{\sqrt[3]{38x^3 + \sqrt{500x^6 + (38x^3 + 108x + 27c_1)^2} + 108x + 27c_1}} - 8x \right)$$

6.22 problem 22

6.22.1 Existence and uniqueness analysis	1647
6.22.2 Solving as separable ode	1647
6.22.3 Solving as differentialType ode	1650
6.22.4 Solving as first order ode lie symmetry lookup ode	1652
6.22.5 Solving as exact ode	1656
6.22.6 Maple step by step solution	1659

Internal problem ID [589]

Internal file name [OUTPUT/589_Sunday_June_05_2022_01_45_22_AM_65389013/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 22.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact", "separable", "differential-Type", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

[_separable]

$$y' - \frac{x^2 - 1}{1 + y^2} = 0$$

With initial conditions

$$[y(-1) = 1]$$

6.22.1 Existence and uniqueness analysis

This is non linear first order ODE. In canonical form it is written as

$$\begin{aligned}y' &= f(x, y) \\ &= \frac{x^2 - 1}{y^2 + 1}\end{aligned}$$

The x domain of $f(x, y)$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = -1$ is inside this domain. The y domain of $f(x, y)$ when $x = -1$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Now we will look at the continuity of

$$\begin{aligned}\frac{\partial f}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{x^2 - 1}{y^2 + 1} \right) \\ &= -\frac{2(x^2 - 1)y}{(y^2 + 1)^2}\end{aligned}$$

The x domain of $\frac{\partial f}{\partial y}$ when $y = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = -1$ is inside this domain. The y domain of $\frac{\partial f}{\partial y}$ when $x = -1$ is

$$\{-\infty < y < \infty\}$$

And the point $y_0 = 1$ is inside this domain. Therefore solution exists and is unique.

6.22.2 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{x^2 - 1}{y^2 + 1}\end{aligned}$$

Where $f(x) = x^2 - 1$ and $g(y) = \frac{1}{y^2+1}$. Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{1}{y^2+1}} dy &= x^2 - 1 dx \\ \int \frac{1}{\frac{1}{y^2+1}} dy &= \int x^2 - 1 dx \\ \frac{1}{3}y^3 + y &= \frac{1}{3}x^3 - x + c_1\end{aligned}$$

Which results in

$$\begin{aligned}y &= \frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{2} \\ &\quad - \frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{2} \\ y &= -\frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{4} \\ &\quad + \frac{1}{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}} \\ &\quad + \frac{i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{2}\right)}{2} + \frac{2}{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}} \\ y &= -\frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{4} \\ &\quad + \frac{1}{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}} \\ &\quad + \frac{i\sqrt{3} \left(\frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{2}\right)}{2} + \frac{2}{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}\end{aligned}$$

Initial conditions are used to solve for c_1 . Substituting $x = -1$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{-i\sqrt{3} \left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{2}{3}} - 4i\sqrt{3} - \left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{2}{3}} + 4}{4 \left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}}}$$

Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = -1$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{i\sqrt{3} \left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{2}{3}} - \left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{2}{3}} + 4i\sqrt{3} + 4}{4 \left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}}}$$

Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = -1$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{2}{3}} - 4}{2 \left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}}}$$

$$c_1 = \frac{2}{3}$$

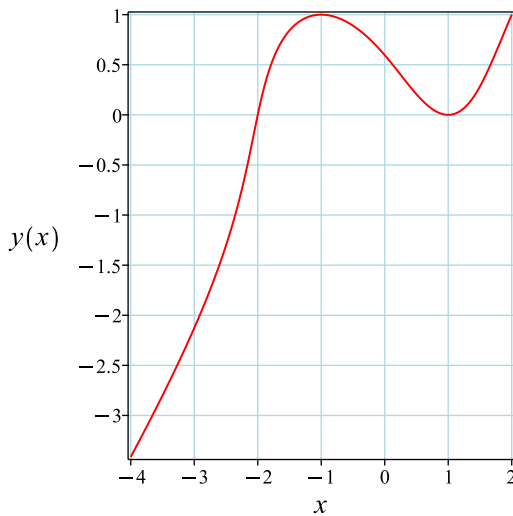
Substituting c_1 found above in the general solution gives

$$y = \frac{(4x^3 + 8 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8})^{\frac{2}{3}} - 4}{2(4x^3 + 8 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8})^{\frac{1}{3}}}$$

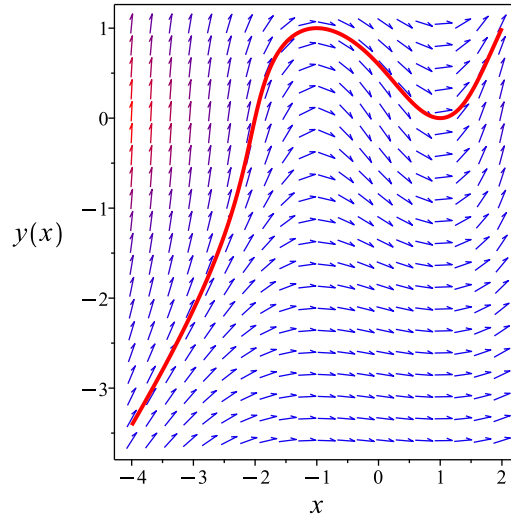
Summary

The solution(s) found are the following

$$y = \frac{(4x^3 + 8 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8})^{\frac{2}{3}} - 4}{2(4x^3 + 8 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8})^{\frac{1}{3}}} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{(4x^3 + 8 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8})^{\frac{2}{3}} - 4}{2(4x^3 + 8 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8})^{\frac{1}{3}}}$$

Verified OK.

6.22.3 Solving as differential Type ode

Writing the ode as

$$y' = \frac{x^2 - 1}{1 + y^2} \quad (1)$$

Which becomes

$$(y^2 + 1) dy = (x^2 - 1) dx \quad (2)$$

But the RHS is complete differential because

$$(x^2 - 1) dx = d\left(\frac{1}{3}x^3 - x\right)$$

Hence (2) becomes

$$(y^2 + 1) dy = d\left(\frac{1}{3}x^3 - x\right)$$

Integrating both sides gives gives these solutions

$$y = \frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{2} - \frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}$$

$$y = -\frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}$$

$$y = -\frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{4} + \frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}$$

Initial conditions are used to solve for c_1 . Substituting $x = -1$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{-i\sqrt{3}\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{2}{3}} - \left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{2}{3}} + 4c_1\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}}}{4\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = -1$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{i\sqrt{3}\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{2}{3}} - \left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{2}{3}} + 4c_1\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}}}{4\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration. Initial conditions are used to solve for c_1 . Substituting $x = -1$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$1 = \frac{\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{2}{3}} + 2c_1\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}} - 4}{2\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}}}$$

Warning: Unable to solve for constant of integration.

Verification of solutions N/A

6.22.4 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x^2 - 1}{y^2 + 1}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 301: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= \frac{1}{x^2 - 1} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{\frac{1}{x^2-1}} dx\end{aligned}$$

Which results in

$$S = \frac{1}{3}x^3 - x$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x^2 - 1}{y^2 + 1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= x^2 - 1 \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = y^2 + 1 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R^2 + 1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{1}{3}R^3 + R + c_1 \tag{4}$$

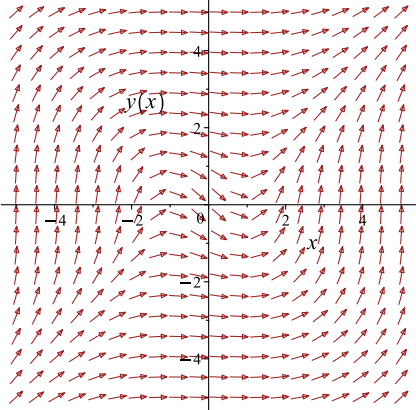
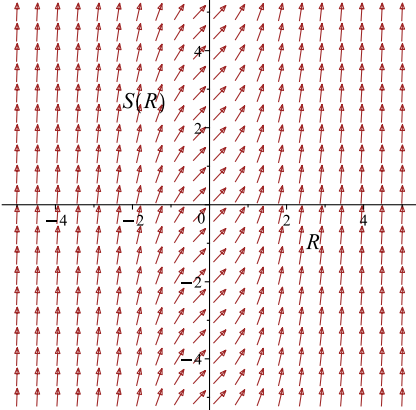
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{1}{3}x^3 - x = \frac{y^3}{3} + y + c_1$$

Which simplifies to

$$\frac{1}{3}x^3 - x = \frac{y^3}{3} + y + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{x^2-1}{y^2+1}$ 	$R = y$ $S = \frac{1}{3}x^3 - x$	$\frac{dS}{dR} = R^2 + 1$ 

Initial conditions are used to solve for c_1 . Substituting $x = -1$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{2}{3} = c_1 + \frac{4}{3}$$

$$c_1 = -\frac{2}{3}$$

Substituting c_1 found above in the general solution gives

$$\frac{1}{3}x^3 - x = \frac{1}{3}y^3 + y - \frac{2}{3}$$

Summary

The solution(s) found are the following

$$\frac{1}{3}x^3 - x = \frac{y^3}{3} + y - \frac{2}{3} \tag{1}$$

Verification of solutions

$$\frac{1}{3}x^3 - x = \frac{y^3}{3} + y - \frac{2}{3}$$

Verified OK.

6.22.5 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (y^2 + 1) dy &= (x^2 - 1) dx \\ (-x^2 + 1) dx + (y^2 + 1) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x^2 + 1 \\ N(x, y) &= y^2 + 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x^2 + 1) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(y^2 + 1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -x^2 + 1 dx$$

$$\phi = -\frac{1}{3}x^3 + x + f(y) \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = y^2 + 1$. Therefore equation (4) becomes

$$y^2 + 1 = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = y^2 + 1$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (y^2 + 1) dy$$

$$f(y) = \frac{1}{3}y^3 + y + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{1}{3}x^3 + x + \frac{1}{3}y^3 + y + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{1}{3}x^3 + x + \frac{1}{3}y^3 + y$$

Initial conditions are used to solve for c_1 . Substituting $x = -1$ and $y = 1$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{2}{3} = c_1$$

$$c_1 = \frac{2}{3}$$

Substituting c_1 found above in the general solution gives

$$-\frac{1}{3}x^3 + x + \frac{1}{3}y^3 + y = \frac{2}{3}$$

Summary

The solution(s) found are the following

$$-\frac{x^3}{3} + x + \frac{y^3}{3} + y = \frac{2}{3} \quad (1)$$

Verification of solutions

$$-\frac{x^3}{3} + x + \frac{y^3}{3} + y = \frac{2}{3}$$

Verified OK.

6.22.6 Maple step by step solution

Let's solve

$$\left[y' - \frac{x^2-1}{1+y^2} = 0, y(-1) = 1 \right]$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$y'(1 + y^2) = x^2 - 1$$

- Integrate both sides with respect to x

$$\int y'(1 + y^2) dx = \int (x^2 - 1) dx + c_1$$

- Evaluate integral

$$\frac{y^3}{3} + y = \frac{1}{3}x^3 - x + c_1$$

- Solve for y

$$y = \frac{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}{2} - \frac{2}{\left(4x^3 + 12c_1 - 12x + 4\sqrt{x^6 + 6c_1x^3 - 6x^4 + 9c_1^2 - 18c_1x + 9x^2 + 4}\right)^{\frac{1}{3}}}$$

- Use initial condition $y(-1) = 1$

$$1 = \frac{\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}}}{2} - \frac{2}{\left(8 + 12c_1 + 4\sqrt{9c_1^2 + 12c_1 + 8}\right)^{\frac{1}{3}}}$$

- Solve for c_1

$$c_1 = \frac{2}{3}$$

- Substitute $c_1 = \frac{2}{3}$ into general solution and simplify

$$y = \frac{\left(4x^3 + 8 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8}\right)^{\frac{2}{3}} - 4}{2\left(4x^3 + 8 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8}\right)^{\frac{1}{3}}}$$

- Solution to the IVP

$$y = \frac{\left(4x^3 + 8 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8}\right)^{\frac{2}{3}} - 4}{2\left(4x^3 + 8 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8}\right)^{\frac{1}{3}}}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```

✓ Solution by Maple

Time used: 0.11 (sec). Leaf size: 87

```
dsolve([diff(y(x),x) = (x^2-1)/(1+y(x)^2),y(-1) = 1],y(x), singsol=all)
```

$$y(x) = \frac{(8 + 4x^3 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8})^{\frac{2}{3}} - 4}{2(8 + 4x^3 - 12x + 4\sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8})^{\frac{1}{3}}}$$

✓ Solution by Mathematica

Time used: 2.95 (sec). Leaf size: 97

```
DSolve[{y'[x]== (x^2-1)/(1+y[x]^2),y[-1]==1},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{\sqrt[3]{2}(x^3 + \sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8} - 3x + 2)^{2/3} - 2}{2^{2/3}\sqrt[3]{x^3 + \sqrt{x^6 - 6x^4 + 4x^3 + 9x^2 - 12x + 8} - 3x + 2}}$$

6.23 problem 23

6.23.1 Solving as linear ode	1661
6.23.2 Solving as first order ode lie symmetry lookup ode	1663
6.23.3 Solving as exact ode	1667
6.23.4 Maple step by step solution	1671

Internal problem ID [590]

Internal file name [OUTPUT/590_Sunday_June_05_2022_01_45_23_AM_42038206/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 23.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$(t + 1)y + ty' = e^{2t}$$

6.23.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{-t - 1}{t}$$

$$q(t) = \frac{e^{2t}}{t}$$

Hence the ode is

$$y' - \frac{(-t - 1)y}{t} = \frac{e^{2t}}{t}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{t-1}{t} dt} \\ &= e^{t+\ln(t)}\end{aligned}$$

Which simplifies to

$$\mu = t e^t$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{e^{2t}}{t} \right) \\ \frac{d}{dt}(t e^t y) &= (t e^t) \left(\frac{e^{2t}}{t} \right) \\ d(t e^t y) &= e^{3t} dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t e^t y &= \int e^{3t} dt \\ t e^t y &= \frac{e^{3t}}{3} + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t e^t$ results in

$$y = \frac{e^{-t} e^{3t}}{3t} + \frac{c_1 e^{-t}}{t}$$

which simplifies to

$$y = \frac{3c_1 e^{-t} + e^{2t}}{3t}$$

Summary

The solution(s) found are the following

$$y = \frac{3c_1 e^{-t} + e^{2t}}{3t} \tag{1}$$

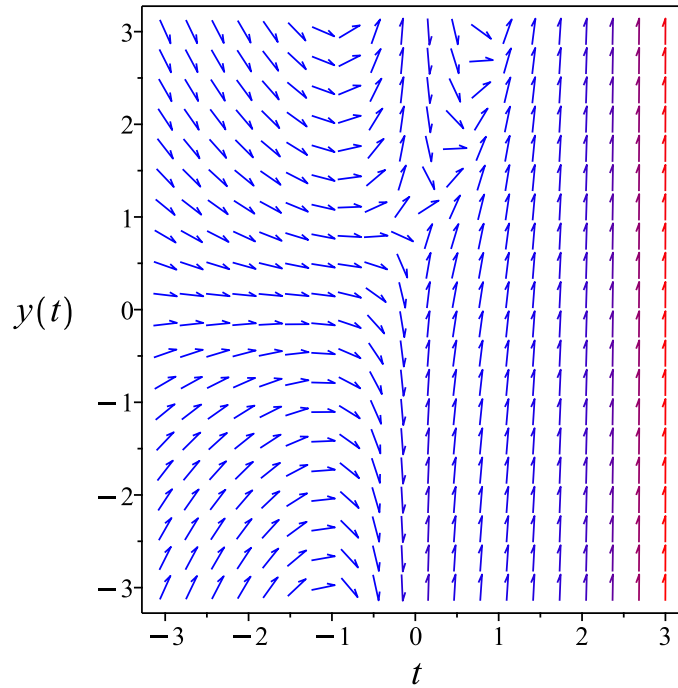


Figure 328: Slope field plot

Verification of solutions

$$y = \frac{3c_1 e^{-t} + e^{2t}}{3t}$$

Verified OK.

6.23.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-ty + e^{2t} - y}{t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 304: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= e^{-t-\ln(t)}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-t-\ln(t)}} dy \end{aligned}$$

Which results in

$$S = t e^t y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = \frac{-ty + e^{2t} - y}{t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= y e^t (t + 1) \\ S_y &= t e^t \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = e^{3t} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = e^{3R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \frac{e^{3R}}{3} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$t e^t y = \frac{e^{3t}}{3} + c_1$$

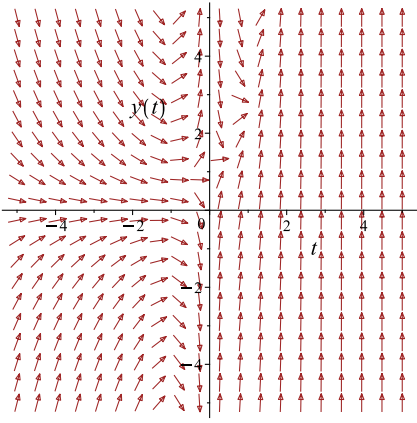
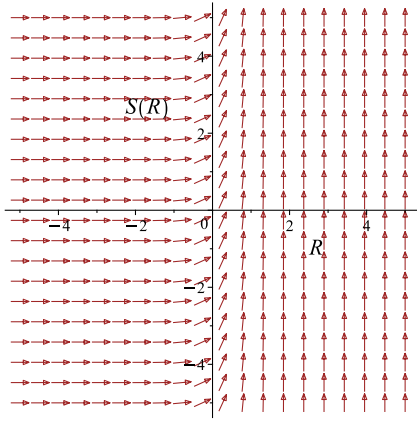
Which simplifies to

$$t e^t y = \frac{e^{3t}}{3} + c_1$$

Which gives

$$y = \frac{(e^{3t} + 3c_1) e^{-t}}{3t}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = \frac{-ty + e^{2t} - y}{t}$ 	$R = t$ $S = t e^t y$	$\frac{dS}{dR} = e^{3R}$ 

Summary

The solution(s) found are the following

$$y = \frac{(e^{3t} + 3c_1) e^{-t}}{3t} \quad (1)$$

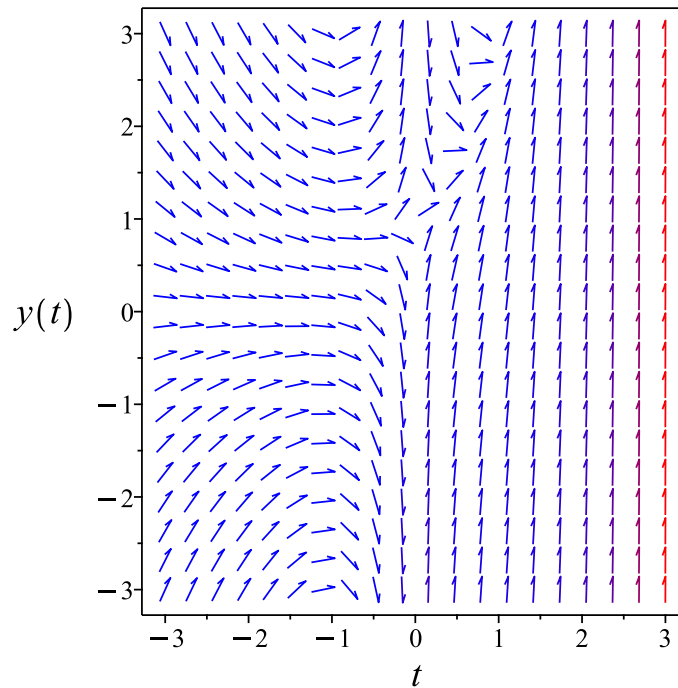


Figure 329: Slope field plot

Verification of solutions

$$y = \frac{(e^{3t} + 3c_1) e^{-t}}{3t}$$

Verified OK.

6.23.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(t) dy &= (-t + 1)y + e^{2t} dt \\ (-e^{2t} + (t + 1)y) dt + (t) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= -e^{2t} + (t + 1)y \\ N(t, y) &= t\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-e^{2t} + (t + 1)y) \\ &= t + 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(t) \\ &= 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{t} ((t+1) - (1)) \\ &= 1 \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int 1 dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^t \\ &= e^t \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^t (-e^{2t} + (t+1)y) \\ &= (-e^{2t} + (t+1)y) e^t \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^t(t) \\ &= t e^t \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ ((-e^{2t} + (t+1)y) e^t) + (t e^t) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\int \frac{\partial \phi}{\partial t} dt = \int \bar{M} dt$$

$$\int \frac{\partial \phi}{\partial t} dt = \int (-e^{2t} + (t+1)y) e^t dt$$

$$\phi = t e^t y - \frac{e^{3t}}{3} + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t e^t + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t e^t$. Therefore equation (4) becomes

$$t e^t = t e^t + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = t e^t y - \frac{e^{3t}}{3} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = t e^t y - \frac{e^{3t}}{3}$$

The solution becomes

$$y = \frac{(e^{3t} + 3c_1) e^{-t}}{3t}$$

Summary

The solution(s) found are the following

$$y = \frac{(e^{3t} + 3c_1) e^{-t}}{3t} \quad (1)$$

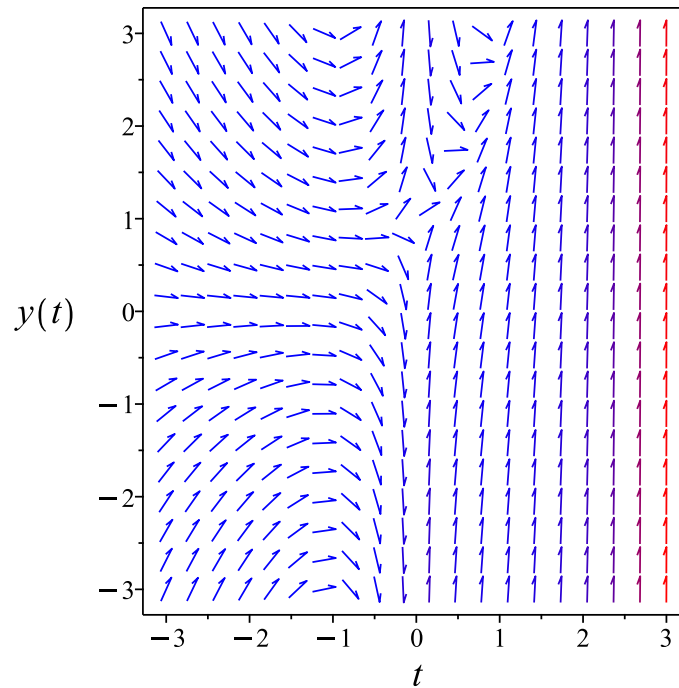


Figure 330: Slope field plot

Verification of solutions

$$y = \frac{(e^{3t} + 3c_1) e^{-t}}{3t}$$

Verified OK.

6.23.4 Maple step by step solution

Let's solve

$$(t + 1)y + ty' = e^{2t}$$

- Highest derivative means the order of the ODE is 1

y'

- Isolate the derivative

$$y' = -\frac{(t+1)y}{t} + \frac{e^{2t}}{t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{(t+1)y}{t} = \frac{e^{2t}}{t}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{(t+1)y}{t} \right) = \frac{\mu(t)e^{2t}}{t}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{(t+1)y}{t} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)(t+1)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = t e^t$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int \frac{\mu(t)e^{2t}}{t} dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int \frac{\mu(t)e^{2t}}{t} dt + c_1$$

- Solve for y

$$y = \frac{\int \frac{\mu(t)e^{2t}}{t} dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = t e^t$

$$y = \frac{\int e^{2t} e^t dt + c_1}{t e^t}$$

- Evaluate the integrals on the rhs

$$y = \frac{\frac{e^{3t}}{3} + c_1}{t e^t}$$

- Simplify

$$y = \frac{3c_1 e^{-t} + e^{2t}}{3t}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 21

```
dsolve((1+t)*y(t)+t*diff(y(t),t) = exp(2*t),y(t), singsol=all)
```

$$y(t) = \frac{e^{2t} + 3e^{-t}c_1}{3t}$$

✓ Solution by Mathematica

Time used: 0.053 (sec). Leaf size: 27

```
DSolve[(1+t)*y[t]+t*y'[t] == Exp[2*t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{e^{2t} + 3c_1e^{-t}}{3t}$$

6.24 problem 24

6.24.1 Solving as separable ode	1674
6.24.2 Solving as first order ode lie symmetry lookup ode	1676
6.24.3 Solving as exact ode	1680
6.24.4 Maple step by step solution	1684

Internal problem ID [591]

Internal file name [OUTPUT/591_Sunday_June_05_2022_01_45_24_AM_63570909/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 24.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact", "separable", "first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_separable]`

$$2 \cos(x) \sin(x) \sin(y) + \cos(y) \sin(x)^2 y' = 0$$

6.24.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\frac{2 \cos(x) \tan(y)}{\sin(x)} \end{aligned}$$

Where $f(x) = -\frac{2 \cos(x)}{\sin(x)}$ and $g(y) = \tan(y)$. Integrating both sides gives

$$\begin{aligned} \frac{1}{\tan(y)} dy &= -\frac{2 \cos(x)}{\sin(x)} dx \\ \int \frac{1}{\tan(y)} dy &= \int -\frac{2 \cos(x)}{\sin(x)} dx \\ \ln(\sin(y)) &= -2 \ln(\sin(x)) + c_1 \end{aligned}$$

Raising both side to exponential gives

$$\sin(y) = e^{-2\ln(\sin(x))+c_1}$$

Which simplifies to

$$\sin(y) = \frac{c_2}{\sin(x)^2}$$

Summary

The solution(s) found are the following

$$y = \arcsin\left(\frac{c_2 e^{c_1}}{\sin(x)^2}\right) \quad (1)$$

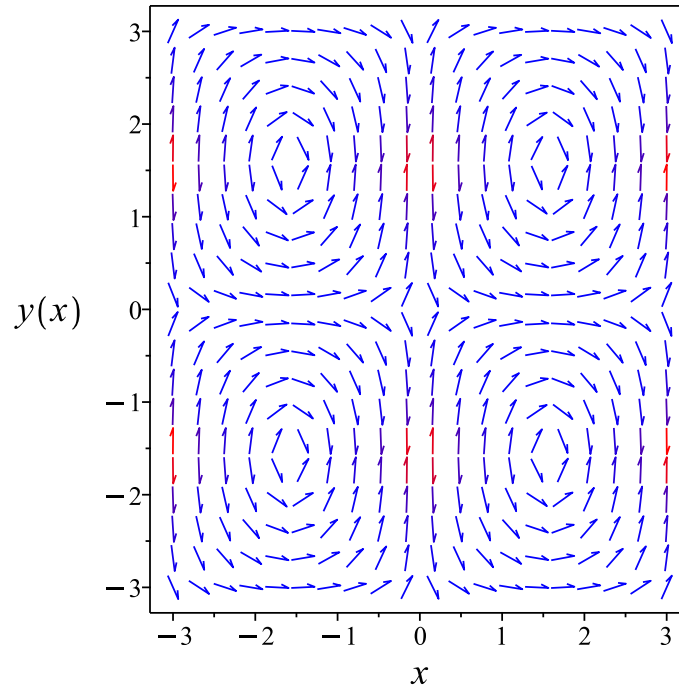


Figure 331: Slope field plot

Verification of solutions

$$y = \arcsin\left(\frac{c_2 e^{c_1}}{\sin(x)^2}\right)$$

Verified OK.

6.24.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{2 \cos(x) \sin(y)}{\sin(x) \cos(y)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 307: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= -\frac{\sin(x)}{2 \cos(x)} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\eta = 0$ then in this special case

$$R = y$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{-\frac{\sin(x)}{2 \cos(x)}} dx\end{aligned}$$

Which results in

$$S = -2 \ln(\sin(x))$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{2 \cos(x) \sin(y)}{\sin(x) \cos(y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 0 \\R_y &= 1 \\S_x &= -2 \cot(x) \\S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \cot(y) \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \cot(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = \ln(\sin(R)) + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-2 \ln(\sin(x)) = \ln(\sin(y)) + c_1$$

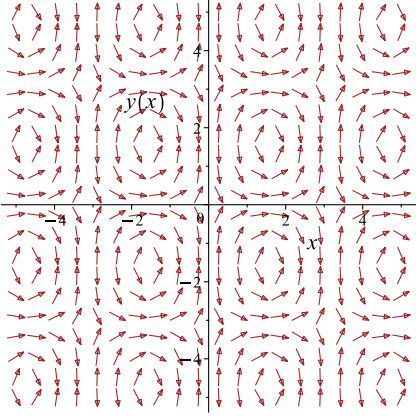
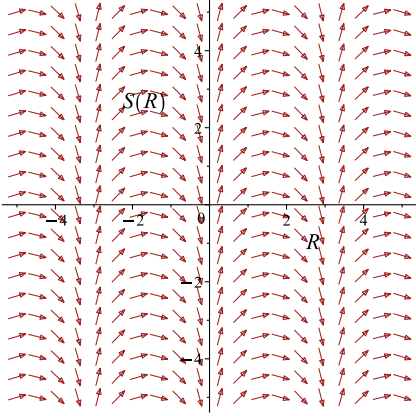
Which simplifies to

$$-2 \ln(\sin(x)) = \ln(\sin(y)) + c_1$$

Which gives

$$y = \arcsin\left(\frac{e^{-c_1}}{\sin(x)^2}\right)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{2 \cos(x) \sin(y)}{\sin(x) \cos(y)}$ 	$R = y$ $S = -2 \ln(\sin(x))$	$\frac{dS}{dR} = \cot(R)$ 

Summary

The solution(s) found are the following

$$y = \arcsin\left(\frac{e^{-c_1}}{\sin(x)^2}\right) \tag{1}$$

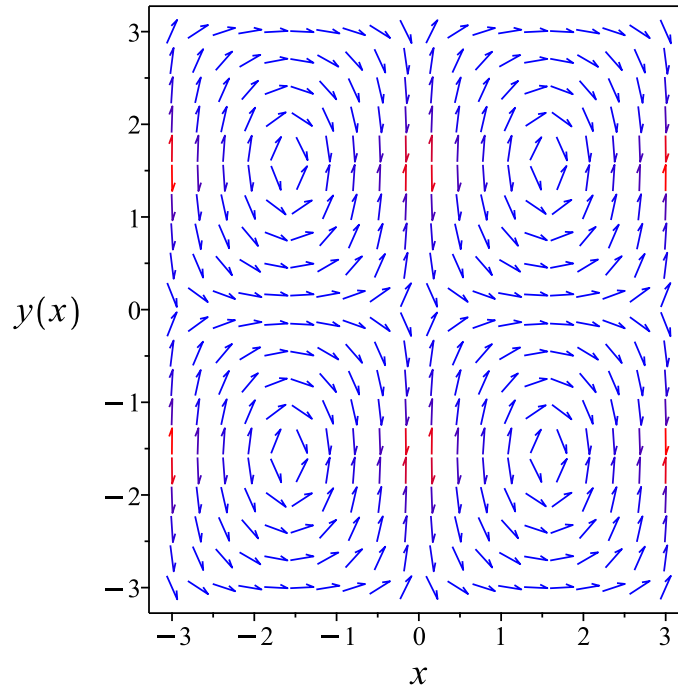


Figure 332: Slope field plot

Verification of solutions

$$y = \arcsin\left(\frac{e^{-c_1}}{\sin(x)^2}\right)$$

Verified OK.

6.24.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(-\frac{\cos(y)}{2 \sin(y)}\right) dy &= \left(\frac{\cos(x)}{\sin(x)}\right) dx \\ \left(-\frac{\cos(x)}{\sin(x)}\right) dx + \left(-\frac{\cos(y)}{2 \sin(y)}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{\cos(x)}{\sin(x)} \\ N(x, y) &= -\frac{\cos(y)}{2 \sin(y)}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{\cos(x)}{\sin(x)}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(-\frac{\cos(y)}{2\sin(y)} \right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{\cos(x)}{\sin(x)} dx \\ \phi &= -\ln(\sin(x)) + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{\cos(y)}{2\sin(y)}$. Therefore equation (4) becomes

$$-\frac{\cos(y)}{2\sin(y)} = 0 + f'(y) \tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$\begin{aligned}f'(y) &= -\frac{\cos(y)}{2\sin(y)} \\ &= -\frac{\cot(y)}{2}\end{aligned}$$

Integrating the above w.r.t y results in

$$\int f'(y) dy = \int \left(-\frac{\cot(y)}{2} \right) dy$$
$$f(y) = -\frac{\ln(\sin(y))}{2} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\ln(\sin(x)) - \frac{\ln(\sin(y))}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\ln(\sin(x)) - \frac{\ln(\sin(y))}{2}$$

Summary

The solution(s) found are the following

$$-\ln(\sin(x)) - \frac{\ln(\sin(y))}{2} = c_1 \tag{1}$$

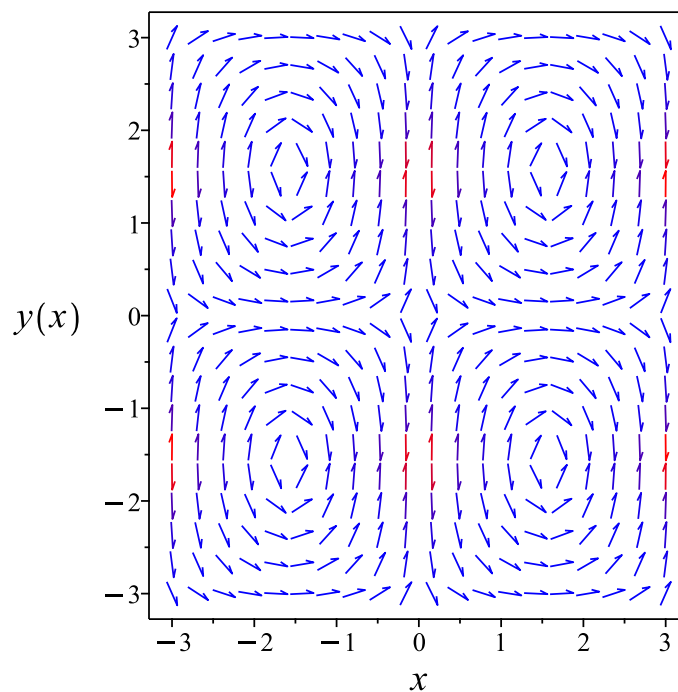


Figure 333: Slope field plot

Verification of solutions

$$-\ln(\sin(x)) - \frac{\ln(\sin(y))}{2} = c_1$$

Verified OK.

6.24.4 Maple step by step solution

Let's solve

$$2 \cos(x) \sin(x) \sin(y) + \cos(y) \sin(x)^2 y' = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Integrate both sides with respect to x

$$\int (2 \cos(x) \sin(x) \sin(y) + \cos(y) \sin(x)^2 y') dx = \int 0 dx + c_1$$

- Evaluate integral

$$\sin(x)^2 \sin(y) = c_1$$

- Solve for y

$$y = \arcsin\left(\frac{c_1}{\sin(x)^2}\right)$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

✓ Solution by Maple

Time used: 0.328 (sec). Leaf size: 18

```
dsolve(2*cos(x)*sin(x)*sin(y(x))+cos(y(x))*sin(x)^2*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = -\arcsin\left(\frac{2c_1}{-1 + \cos(2x)}\right)$$

✓ Solution by Mathematica

Time used: 5.176 (sec). Leaf size: 21

```
DSolve[2*Cos[x]*Sin[x]*Sin[y[x]]+Cos[y[x]]*Sin[x]^2*y'[x] == 0,y[x],x,IncludeSingularSolutio
```

$$y(x) \rightarrow \arcsin\left(\frac{1}{2}c_1 \csc^2(x)\right)$$

$$y(x) \rightarrow 0$$

6.25 problem 25

- 6.25.1 Solving as exact ode 1686
- 6.25.2 Maple step by step solution 1690

Internal problem ID [592]

Internal file name [OUTPUT/592_Sunday_June_05_2022_01_45_26_AM_93723432/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 25.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exact"**

Maple gives the following as the ode type

`[_exact, _rational]`

$$\frac{2x}{y} - \frac{y}{x^2 + y^2} + \left(-\frac{x^2}{y^2} + \frac{x}{x^2 + y^2} \right) y' = 0$$

6.25.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(-\frac{x^2}{y^2} + \frac{x}{x^2 + y^2}\right) dy &= \left(-\frac{2x}{y} + \frac{y}{x^2 + y^2}\right) dx \\ \left(\frac{2x}{y} - \frac{y}{x^2 + y^2}\right) dx + \left(-\frac{x^2}{y^2} + \frac{x}{x^2 + y^2}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= \frac{2x}{y} - \frac{y}{x^2 + y^2} \\ N(x, y) &= -\frac{x^2}{y^2} + \frac{x}{x^2 + y^2}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{2x}{y} - \frac{y}{x^2 + y^2} \right) \\ &= -\frac{2x}{y^2} - \frac{1}{x^2 + y^2} + \frac{2y^2}{(x^2 + y^2)^2}\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(-\frac{x^2}{y^2} + \frac{x}{x^2 + y^2} \right) \\ &= -\frac{2x}{y^2} - \frac{2x^2}{(x^2 + y^2)^2} + \frac{1}{x^2 + y^2}\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{2x}{y} - \frac{y}{x^2 + y^2} dx \\ \phi &= \frac{x^2}{y} - \arctan\left(\frac{x}{y}\right) + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = -\frac{x^2}{y^2} + \frac{x}{y^2 \left(\frac{x^2}{y^2} + 1 \right)} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{x^2}{y^2} + \frac{x}{x^2 + y^2}$. Therefore equation (4) becomes

$$-\frac{x^2}{y^2} + \frac{x}{x^2 + y^2} = -\frac{((x-1)y^2 + x^3)x}{y^2(x^2 + y^2)} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{x^2}{y} - \arctan\left(\frac{x}{y}\right) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{x^2}{y} - \arctan\left(\frac{x}{y}\right)$$

Summary

The solution(s) found are the following

$$\frac{x^2}{y} - \arctan\left(\frac{x}{y}\right) = c_1 \tag{1}$$

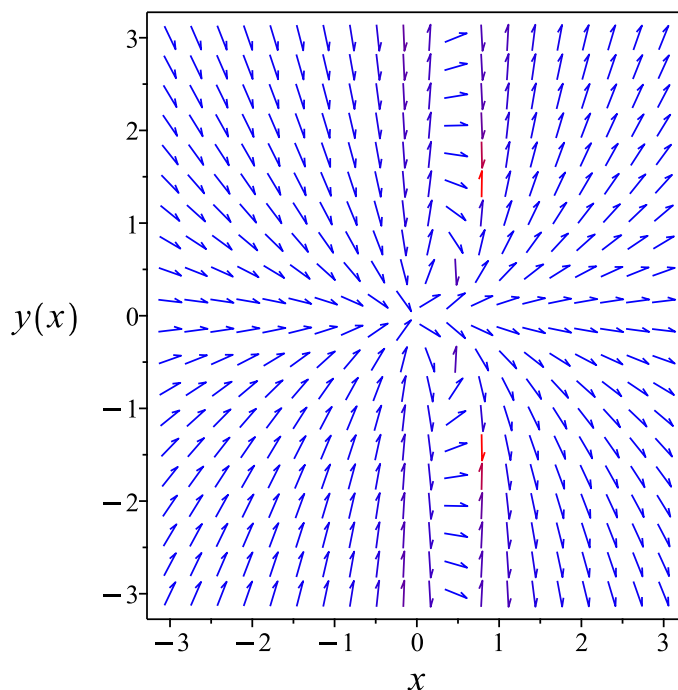


Figure 334: Slope field plot

Verification of solutions

$$\frac{x^2}{y} - \arctan\left(\frac{x}{y}\right) = c_1$$

Verified OK.

6.25.2 Maple step by step solution

Let's solve

$$\frac{2x}{y} - \frac{y}{x^2+y^2} + \left(-\frac{x^2}{y^2} + \frac{x}{x^2+y^2}\right) y' = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Check if ODE is exact

- ODE is exact if the lhs is the total derivative of a C^2 function

$$F'(x, y) = 0$$

- Compute derivative of lhs

$$F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$$

- Evaluate derivatives

$$-\frac{2x}{y^2} - \frac{1}{x^2+y^2} + \frac{2y^2}{(x^2+y^2)^2} = -\frac{2x}{y^2} - \frac{2x^2}{(x^2+y^2)^2} + \frac{1}{x^2+y^2}$$

- Simplify

$$-\frac{2x}{y^2} - \frac{1}{x^2+y^2} + \frac{2y^2}{(x^2+y^2)^2} = -\frac{2x}{y^2} - \frac{2x^2}{(x^2+y^2)^2} + \frac{1}{x^2+y^2}$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$\left[F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$

- Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to x

$$F(x, y) = \int \left(\frac{2x}{y} - \frac{y}{x^2+y^2} \right) dx + f_1(y)$$

- Evaluate integral

$$F(x, y) = \frac{x^2}{y} - \arctan\left(\frac{x}{y}\right) + f_1(y)$$

- Take derivative of $F(x, y)$ with respect to y

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$

- Compute derivative

$$-\frac{x^2}{y^2} + \frac{x}{x^2+y^2} = -\frac{x^2}{y^2} + \frac{x}{y^2\left(\frac{x^2}{y^2}+1\right)} + \frac{d}{dy} f_1(y)$$

- Isolate for $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = \frac{x}{x^2+y^2} - \frac{x}{y^2\left(\frac{x^2}{y^2}+1\right)}$$

- Solve for $f_1(y)$

$$f_1(y) = 0$$

- Substitute $f_1(y)$ into equation for $F(x, y)$

$$F(x, y) = \frac{x^2}{y} - \arctan\left(\frac{x}{y}\right)$$

- Substitute $F(x, y)$ into the solution of the ODE

$$\frac{x^2}{y} - \arctan\left(\frac{x}{y}\right) = c_1$$

- Solve for y

$$y = \frac{x}{\tan(\text{RootOf}(-Z - x \tan(Z) + c_1))}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 17

```
dsolve(2*x/y(x)-y(x)/(x^2+y(x)^2)+(-x^2/y(x)^2+x/(x^2+y(x)^2))*diff(y(x),x) = 0,y(x), singso
```

$$y(x) = \cot(\text{RootOf}(-Z + x \tan(Z) + c_1)) x$$

✓ Solution by Mathematica

Time used: 0.255 (sec). Leaf size: 23

```
DSolve[2*x/y[x]-y[x]/(x^2+y[x]^2)+(-x^2/y[x]^2+x/(x^2+y[x]^2))*y'[x] == 0,y[x],x,IncludeSing
```

$$\text{Solve} \left[\arctan \left(\frac{x}{y(x)} \right) - \frac{x^2}{y(x)} = c_1, y(x) \right]$$

6.26 problem 26

- 6.26.1 Solving as homogeneousTypeD ode 1693
- 6.26.2 Solving as homogeneousTypeD2 ode 1695
- 6.26.3 Solving as first order ode lie symmetry lookup ode 1697

Internal problem ID [593]

Internal file name [OUTPUT/593_Sunday_June_05_2022_01_45_28_AM_41374823/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 26.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "homogeneousTypeD", "homogeneousTypeD2", "first_order_ode_lie_symmetry_lookup"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _dAlembert]
```

$$y'x - e^{\frac{y}{x}}x - y = 0$$

6.26.1 Solving as homogeneousTypeD ode

Writing the ode as

$$y' = e^{\frac{y}{x}} + \frac{y}{x} \tag{A}$$

The given ode has the form

$$y' = \frac{y}{x} + g(x) f\left(b\frac{y}{x}\right)^{\frac{n}{m}} \tag{1}$$

Where b is scalar and $g(x)$ is function of x and n, m are integers. The solution is given in Kamke page 20. Using the substitution $y(x) = u(x)x$ then

$$\frac{dy}{dx} = \frac{du}{dx}x + u$$

Hence the given ode becomes

$$\begin{aligned}\frac{du}{dx}x + u &= u + g(x) f(bu)^{\frac{n}{m}} \\ u' &= \frac{1}{x}g(x) f(bu)^{\frac{n}{m}}\end{aligned}\quad (2)$$

The above ode is always separable. This is easily solved for u assuming the integration can be resolved, and then the solution to the original ode becomes $y = ux$. Comparing the given ode (A) with the form (1) shows that

$$\begin{aligned}g(x) &= 1 \\ b &= 1 \\ f\left(\frac{bx}{y}\right) &= e^{\frac{y}{x}}\end{aligned}$$

Substituting the above in (2) results in the $u(x)$ ode as

$$u'(x) = \frac{e^{u(x)}}{x}$$

Which is now solved as separable In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{e^u}{x}\end{aligned}$$

Where $f(x) = \frac{1}{x}$ and $g(u) = e^u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{e^u} du &= \frac{1}{x} dx \\ \int \frac{1}{e^u} du &= \int \frac{1}{x} dx \\ -e^{-u} &= \ln(x) + c_1\end{aligned}$$

The solution is

$$-e^{-u(x)} - \ln(x) - c_1 = 0$$

Therefore the solution is found using $y = ux$. Hence

$$-e^{-\frac{y}{x}} - \ln(x) - c_1 = 0$$

Summary

The solution(s) found are the following

$$-e^{-\frac{y}{x}} - \ln(x) - c_1 = 0 \quad (1)$$

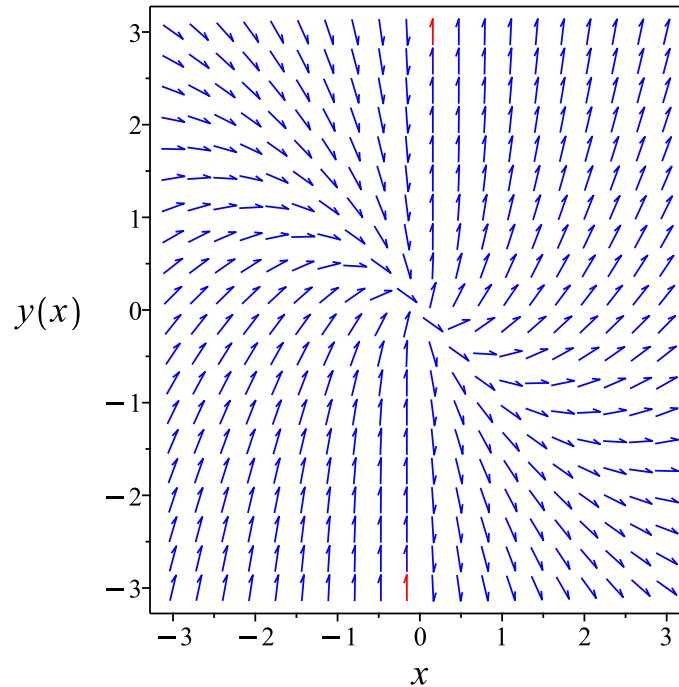


Figure 335: Slope field plot

Verification of solutions

$$-e^{-\frac{y}{x}} - \ln(x) - c_1 = 0$$

Verified OK.

6.26.2 Solving as homogeneous TypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$(u'(x)x + u(x))x - e^{u(x)}x - u(x)x = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{e^u}{x} \end{aligned}$$

Where $f(x) = \frac{1}{x}$ and $g(u) = e^u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{e^u} du &= \frac{1}{x} dx \\ \int \frac{1}{e^u} du &= \int \frac{1}{x} dx \\ -e^{-u} &= \ln(x) + c_2\end{aligned}$$

The solution is

$$-e^{-u(x)} - \ln(x) - c_2 = 0$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\begin{aligned}-e^{-\frac{y}{x}} - \ln(x) - c_2 &= 0 \\ -e^{-\frac{y}{x}} - \ln(x) - c_2 &= 0\end{aligned}$$

Summary

The solution(s) found are the following

$$-e^{-\frac{y}{x}} - \ln(x) - c_2 = 0 \tag{1}$$

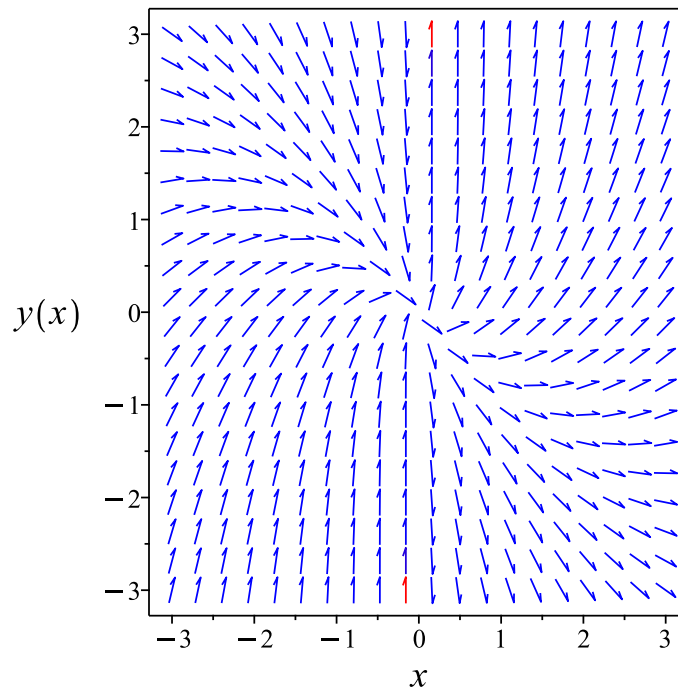


Figure 336: Slope field plot

Verification of solutions

$$-e^{-\frac{y}{x}} - \ln(x) - c_2 = 0$$

Verified OK.

6.26.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{e^{\frac{y}{x}}x + y}{x}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2\xi_y - \omega_x\xi - \omega_y\eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **homogeneous Type D**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 311: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= x^2 \\ \eta(x, y) &= yx\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Therefore

$$\begin{aligned}\frac{dy}{dx} &= \frac{\eta}{\xi} \\ &= \frac{yx}{x^2} \\ &= \frac{y}{x}\end{aligned}$$

This is easily solved to give

$$y = c_1 x$$

Where now the coordinate R is taken as the constant of integration. Hence

$$R = \frac{y}{x}$$

And S is found from

$$\begin{aligned}dS &= \frac{dx}{\xi} \\ &= \frac{dx}{x^2}\end{aligned}$$

Integrating gives

$$\begin{aligned}S &= \int \frac{dx}{T} \\ &= -\frac{1}{x}\end{aligned}$$

Where the constant of integration is set to zero as we just need one solution. Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{e^{\frac{y}{x}}x + y}{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= -\frac{y}{x^2} \\ R_y &= \frac{1}{x} \\ S_x &= \frac{1}{x^2} \\ S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{e^{-\frac{y}{x}}}{x} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -S(R) e^{-R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 e^{e^{-R}} \quad (4)$$

To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-\frac{1}{x} = c_1 e^{e^{-\frac{y}{x}}}$$

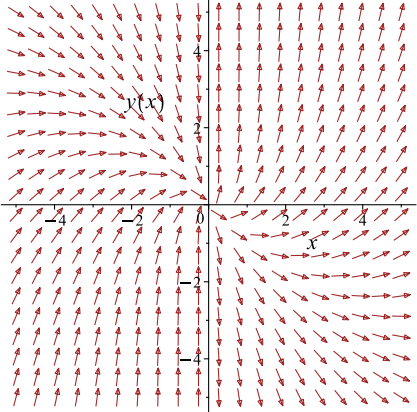
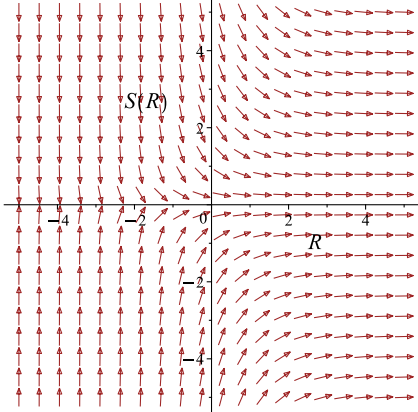
Which simplifies to

$$-\frac{1}{x} = c_1 e^{e^{-\frac{y}{x}}}$$

Which gives

$$y = -\ln \left(\ln \left(-\frac{1}{c_1 x} \right) \right) x$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{e^{\frac{y}{x}} x + y}{x}$ 	$R = \frac{y}{x}$ $S = -\frac{1}{x}$	$\frac{dS}{dR} = -S(R) e^{-R}$ 

Summary

The solution(s) found are the following

$$y = -\ln \left(\ln \left(-\frac{1}{c_1 x} \right) \right) x \tag{1}$$

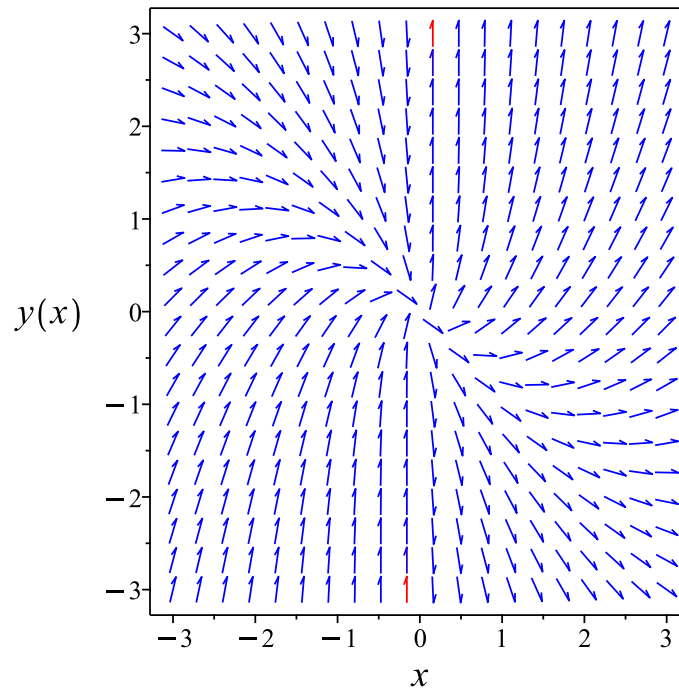


Figure 337: Slope field plot

Verification of solutions

$$y = -\ln\left(\ln\left(-\frac{1}{c_1 x}\right)\right) x$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve(x*diff(y(x),x) = exp(y(x)/x)*x+y(x),y(x), singsol=all)
```

$$y(x) = \ln\left(-\frac{1}{\ln(x) + c_1}\right)x$$

✓ Solution by Mathematica

Time used: 0.316 (sec). Leaf size: 18

```
DSolve[x*y'[x] == Exp[y[x]/x]*x+y[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -x \log(-\log(x) - c_1)$$

6.27 problem 27

6.27.1 Solving as exact ode 1704

Internal problem ID [594]

Internal file name [OUTPUT/594_Sunday_June_05_2022_01_45_29_AM_32116103/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 27.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"exactWithIntegrationFactor"**

Maple gives the following as the ode type

```
[_rational, [_1st_order, ` _with_symmetry_[F(x)*G(y),0] `]]
```

$$y' - \frac{x}{x^2 + y + y^3} = 0$$

6.27.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (y^3 + x^2 + y) dy &= (x) dx \\ (-x) dx + (y^3 + x^2 + y) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x \\ N(x, y) &= y^3 + x^2 + y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(y^3 + x^2 + y) \\ &= 2x \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{y^3 + x^2 + y} ((0) - (2x)) \\ &= -\frac{2x}{y^3 + x^2 + y} \end{aligned}$$

Since A depends on y , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= -\frac{1}{x} ((2x) - (0)) \\ &= -2 \end{aligned}$$

Since B does not depend on x , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned} \mu &= e^{\int B \, dy} \\ &= e^{\int -2 \, dy} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-2y} \\ &= e^{-2y} \end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{-2y}(-x) \\ &= -x e^{-2y} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{-2y}(y^3 + x^2 + y) \\ &= (y^3 + x^2 + y) e^{-2y} \end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (-x e^{-2y}) + ((y^3 + x^2 + y) e^{-2y}) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x e^{-2y} dx \\ \phi &= -\frac{x^2 e^{-2y}}{2} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x^2 e^{-2y} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = (y^3 + x^2 + y) e^{-2y}$. Therefore equation (4) becomes

$$(y^3 + x^2 + y) e^{-2y} = x^2 e^{-2y} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$\begin{aligned} f'(y) &= e^{-2y} y^3 + e^{-2y} y \\ &= e^{-2y} (y^3 + y) \end{aligned}$$

Integrating the above w.r.t y results in

$$\begin{aligned} \int f'(y) dy &= \int (e^{-2y} (y^3 + y)) dy \\ f(y) &= -\frac{e^{-2y} (4y^3 + 6y^2 + 10y + 5)}{8} + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{x^2 e^{-2y}}{2} - \frac{e^{-2y}(4y^3 + 6y^2 + 10y + 5)}{8} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{x^2 e^{-2y}}{2} - \frac{e^{-2y}(4y^3 + 6y^2 + 10y + 5)}{8}$$

Summary

The solution(s) found are the following

$$-\frac{x^2 e^{-2y}}{2} - \frac{e^{-2y}(4y^3 + 6y^2 + 10y + 5)}{8} = c_1 \quad (1)$$

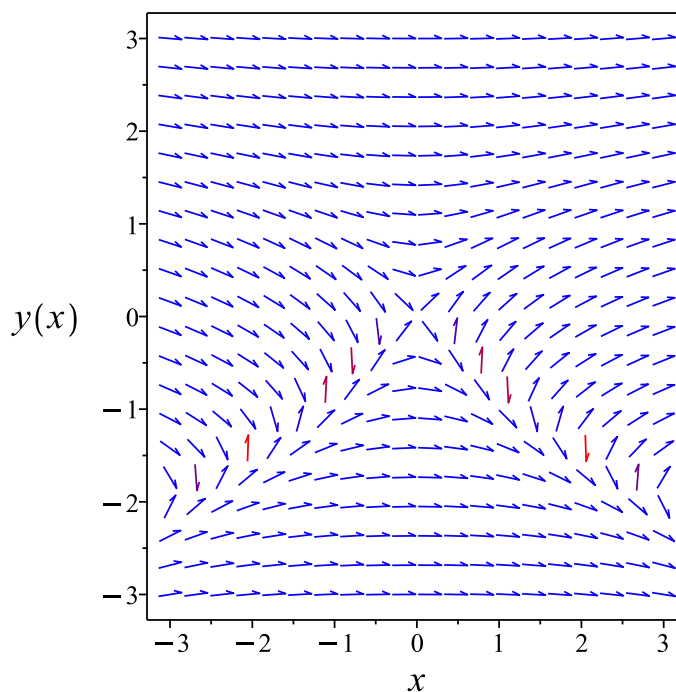


Figure 338: Slope field plot

Verification of solutions

$$-\frac{x^2 e^{-2y}}{2} - \frac{e^{-2y}(4y^3 + 6y^2 + 10y + 5)}{8} = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 35

```
dsolve(diff(y(x),x) = x/(x^2+y(x)+y(x)^3),y(x), singsol=all)
```

$$\frac{(-4y(x)^3 - 4x^2 - 6y(x)^2 - 10y(x) - 5) e^{-2y(x)}}{4} + c_1 = 0$$

✓ Solution by Mathematica

Time used: 0.173 (sec). Leaf size: 48

```
DSolve[y'[x] == x/(x^2+y[x]+y[x]^3),y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve} \left[-\frac{1}{2}x^2 e^{-2y(x)} - \frac{1}{8}e^{-2y(x)} (4y(x)^3 + 6y(x)^2 + 10y(x) + 5) = c_1, y(x) \right]$$

6.28 problem 28

6.28.1 Solving as linear ode	1710
6.28.2 Solving as homogeneousTypeD2 ode	1712
6.28.3 Solving as first order ode lie symmetry lookup ode	1714
6.28.4 Solving as exact ode	1718
6.28.5 Maple step by step solution	1723

Internal problem ID [595]

Internal file name [OUTPUT/595_Sunday_June_05_2022_01_45_31_AM_56902748/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 28.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"linear"**, **"homogeneousTypeD2"**, **"exactWithIntegrationFactor"**, **"first_order_ode_lie_symmetry_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$2y + ty' = -3t$$

6.28.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{2}{t}$$
$$q(t) = -3$$

Hence the ode is

$$y' + \frac{2y}{t} = -3$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{2}{t} dt} \\ &= t^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu)(-3) \\ \frac{d}{dt}(t^2 y) &= (t^2)(-3) \\ d(t^2 y) &= (-3t^2) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t^2 y &= \int -3t^2 dt \\ t^2 y &= -t^3 + c_1\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t^2$ results in

$$y = -t + \frac{c_1}{t^2}$$

Summary

The solution(s) found are the following

$$y = -t + \frac{c_1}{t^2} \tag{1}$$

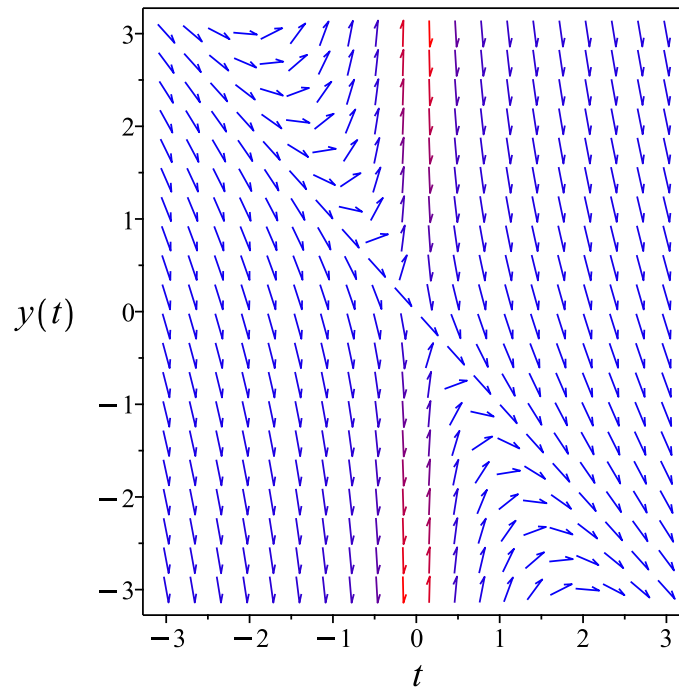


Figure 339: Slope field plot

Verification of solutions

$$y = -t + \frac{c_1}{t^2}$$

Verified OK.

6.28.2 Solving as homogeneous TypeD2 ode

Using the change of variables $y = u(t)t$ on the above ode results in new ode in $u(t)$

$$2u(t)t + t(u'(t)t + u(t)) = -3t$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{-3u - 3}{t} \end{aligned}$$

Where $f(t) = \frac{1}{t}$ and $g(u) = -3u - 3$. Integrating both sides gives

$$\frac{1}{-3u - 3} du = \frac{1}{t} dt$$

$$\int \frac{1}{-3u-3} du = \int \frac{1}{t} dt$$
$$-\frac{\ln(u+1)}{3} = \ln(t) + c_2$$

Raising both side to exponential gives

$$\frac{1}{(u+1)^{\frac{1}{3}}} = e^{\ln(t)+c_2}$$

Which simplifies to

$$\frac{1}{(u+1)^{\frac{1}{3}}} = c_3 t$$

Therefore the solution y is

$$y = ut$$
$$= -\frac{(c_3^3 e^{3c_2} t^3 - 1) e^{-3c_2}}{t^2 c_3^3}$$

Summary

The solution(s) found are the following

$$y = -\frac{(c_3^3 e^{3c_2} t^3 - 1) e^{-3c_2}}{t^2 c_3^3} \quad (1)$$

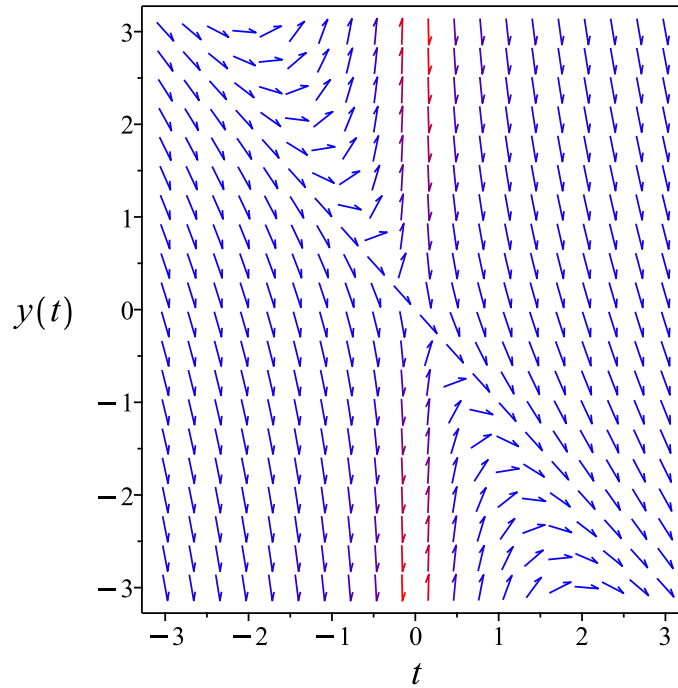


Figure 340: Slope field plot

Verification of solutions

$$y = -\frac{(c_3^3 e^{3c_2 t^3} - 1) e^{-3c_2}}{t^2 c_3^3}$$

Verified OK.

6.28.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{3t + 2y}{t}$$

$$y' = \omega(t, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_y - \xi_t) - \omega^2 \xi_y - \omega_t \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find ξ, η

Table 313: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	ξ	η
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	x	y
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	x^2	xy
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(t, y) &= 0 \\ \eta(t, y) &= \frac{1}{t^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial y}\right) S(t, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{t^2}} dy \end{aligned}$$

Which results in

$$S = t^2 y$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, y)S_y}{R_t + \omega(t, y)R_y} \quad (2)$$

Where in the above R_t, R_y, S_t, S_y are all partial derivatives and $\omega(t, y)$ is the right hand side of the original ode given by

$$\omega(t, y) = -\frac{3t + 2y}{t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_y &= 0 \\ S_t &= 2ty \\ S_y &= t^2 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -3t^2 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -3R^2$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -R^3 + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to t, y coordinates. This results in

$$yt^2 = -t^3 + c_1$$

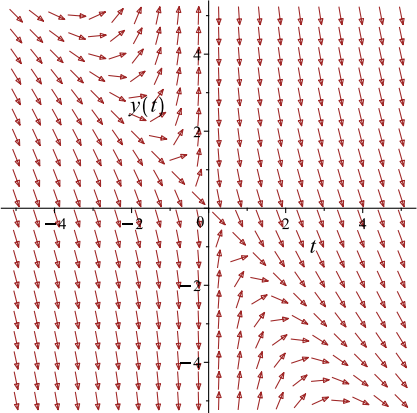
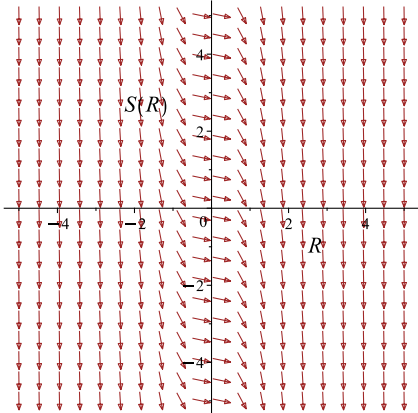
Which simplifies to

$$yt^2 = -t^3 + c_1$$

Which gives

$$y = \frac{-t^3 + c_1}{t^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dt} = -\frac{3t+2y}{t}$ 	$R = t$ $S = t^2y$	$\frac{dS}{dR} = -3R^2$ 

Summary

The solution(s) found are the following

$$y = \frac{-t^3 + c_1}{t^2} \quad (1)$$

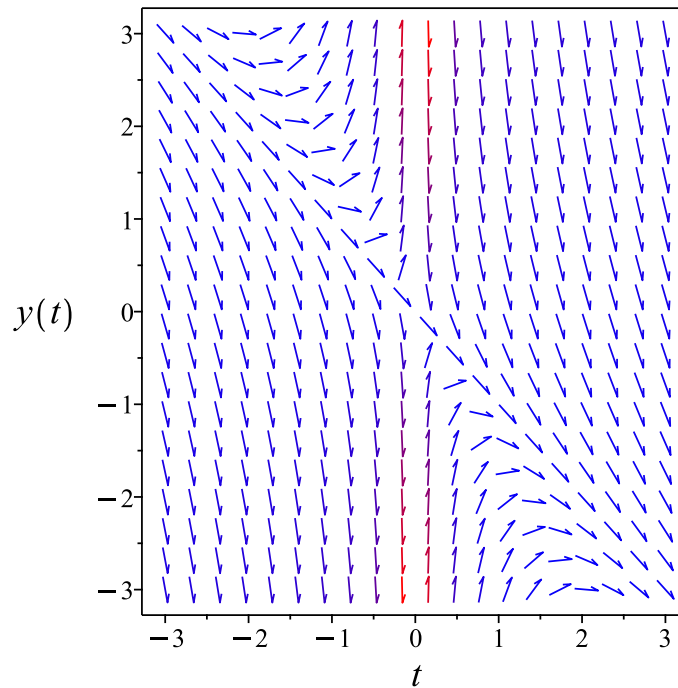


Figure 341: Slope field plot

Verification of solutions

$$y = \frac{-t^3 + c_1}{t^2}$$

Verified OK.

6.28.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, y) dt + N(t, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(t) dy &= (-3t - 2y) dt \\ (3t + 2y) dt + (t) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, y) &= 3t + 2y \\ N(t, y) &= t\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(3t + 2y) \\ &= 2\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(t) \\ &= 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{t} ((2) - (1)) \\ &= \frac{1}{t} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int \frac{1}{t} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\ln(t)} \\ &= t \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= t(3t + 2y) \\ &= t(3t + 2y) \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= t(t) \\ &= t^2 \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dt} &= 0 \\ (t(3t + 2y)) + (t^2) \frac{dy}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, y)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int t(3t + 2y) dt \\ \phi &= t^2(t + y) + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both t and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = t^2 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = t^2$. Therefore equation (4) becomes

$$t^2 = t^2 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = t^2(t + y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = t^2(t + y)$$

The solution becomes

$$y = \frac{-t^3 + c_1}{t^2}$$

Summary

The solution(s) found are the following

$$y = \frac{-t^3 + c_1}{t^2} \tag{1}$$

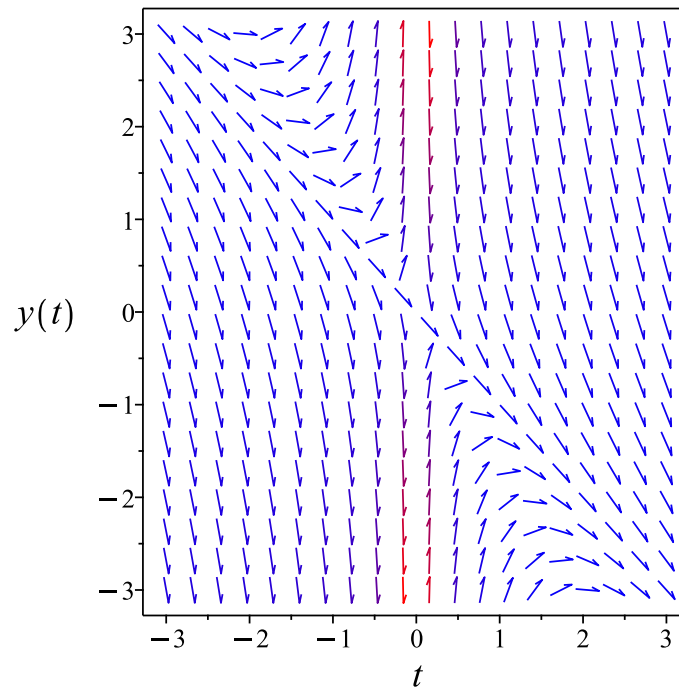


Figure 342: Slope field plot

Verification of solutions

$$y = \frac{-t^3 + c_1}{t^2}$$

Verified OK.

6.28.5 Maple step by step solution

Let's solve

$$2y + ty' = -3t$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -3 - \frac{2y}{t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{2y}{t} = -3$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = -3\mu(t)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(\mu(t)y)$

$$\mu(t) \left(y' + \frac{2y}{t} \right) = \mu'(t)y + \mu(t)y'$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{2\mu(t)}{t}$$

- Solve to find the integrating factor

$$\mu(t) = t^2$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(\mu(t)y) \right) dt = \int -3\mu(t) dt + c_1$$

- Evaluate the integral on the lhs

$$\mu(t)y = \int -3\mu(t) dt + c_1$$

- Solve for y

$$y = \frac{\int -3\mu(t)dt + c_1}{\mu(t)}$$

- Substitute $\mu(t) = t^2$

$$y = \frac{\int -3t^2 dt + c_1}{t^2}$$

- Evaluate the integrals on the rhs

$$y = \frac{-t^3 + c_1}{t^2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 13

```
dsolve(3*t+2*y(t) = -t*diff(y(t),t),y(t), singsol=all)
```

$$y(t) = -t + \frac{c_1}{t^2}$$

✓ Solution by Mathematica

Time used: 0.023 (sec). Leaf size: 15

```
DSolve[3*t+2*y[t] == -t*y'[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -t + \frac{c_1}{t^2}$$

6.29 problem 29

- 6.29.1 Solving as homogeneousTypeD2 ode 1725
- 6.29.2 Solving as first order ode lie symmetry calculated ode 1727
- 6.29.3 Solving as exact ode 1732

Internal problem ID [596]

Internal file name [OUTPUT/596_Sunday_June_05_2022_01_45_32_AM_62306633/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 29.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exactByInspection", "homogeneousTypeD2", "first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$y' - \frac{x+y}{x-y} = 0$$

6.29.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$u'(x)x + u(x) - \frac{x + u(x)x}{x - u(x)x} = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{u^2 + 1}{x(u - 1)} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(u) = \frac{u^2+1}{u-1}$. Integrating both sides gives

$$\frac{1}{\frac{u^2+1}{u-1}} du = -\frac{1}{x} dx$$

$$\int \frac{1}{\frac{u^2+1}{u-1}} du = \int -\frac{1}{x} dx$$

$$\frac{\ln(u^2 + 1)}{2} - \arctan(u) = -\ln(x) + c_2$$

The solution is

$$\frac{\ln(u(x)^2 + 1)}{2} - \arctan(u(x)) + \ln(x) - c_2 = 0$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} - \arctan\left(\frac{y}{x}\right) + \ln(x) - c_2 = 0$$

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} - \arctan\left(\frac{y}{x}\right) + \ln(x) - c_2 = 0$$

Summary

The solution(s) found are the following

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} - \arctan\left(\frac{y}{x}\right) + \ln(x) - c_2 = 0 \quad (1)$$

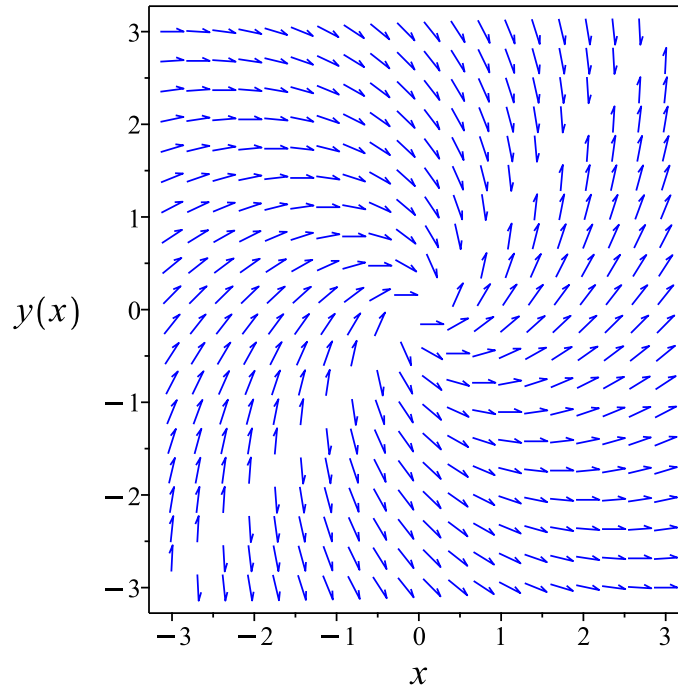


Figure 343: Slope field plot

Verification of solutions

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} - \arctan\left(\frac{y}{x}\right) + \ln(x) - c_2 = 0$$

Verified OK.

6.29.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{x+y}{-x+y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 - \frac{(x+y)(b_3 - a_2)}{-x+y} - \frac{(x+y)^2 a_3}{(-x+y)^2} \\ - \left(-\frac{1}{-x+y} - \frac{x+y}{(-x+y)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left(-\frac{1}{-x+y} + \frac{x+y}{(-x+y)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{x^2 a_2 + x^2 a_3 + x^2 b_2 - x^2 b_3 - 2xy a_2 + 2xy a_3 + 2xy b_2 + 2xy b_3 - y^2 a_2 - y^2 a_3 - y^2 b_2 + y^2 b_3 + 2xb_1 - 2ya_1}{(x-y)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} -x^2 a_2 - x^2 a_3 - x^2 b_2 + x^2 b_3 + 2xy a_2 - 2xy a_3 - 2xy b_2 \\ - 2xy b_3 + y^2 a_2 + y^2 a_3 + y^2 b_2 - y^2 b_3 - 2xb_1 + 2ya_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} -a_2 v_1^2 + 2a_2 v_1 v_2 + a_2 v_2^2 - a_3 v_1^2 - 2a_3 v_1 v_2 + a_3 v_2^2 - b_2 v_1^2 \\ - 2b_2 v_1 v_2 + b_2 v_2^2 + b_3 v_1^2 - 2b_3 v_1 v_2 - b_3 v_2^2 + 2a_1 v_2 - 2b_1 v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} (-a_2 - a_3 - b_2 + b_3) v_1^2 + (2a_2 - 2a_3 - 2b_2 - 2b_3) v_1 v_2 \\ - 2b_1 v_1 + (a_2 + a_3 + b_2 - b_3) v_2^2 + 2a_1 v_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 2a_1 &= 0 \\ -2b_1 &= 0 \\ -a_2 - a_3 - b_2 + b_3 &= 0 \\ a_2 + a_3 + b_2 - b_3 &= 0 \\ 2a_2 - 2a_3 - 2b_2 - 2b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= -b_2 \\ b_1 &= 0 \\ b_2 &= b_2 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= y - \left(-\frac{x+y}{-x+y} \right) (x) \\ &= \frac{-x^2 - y^2}{x-y} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-x^2-y^2}{x-y}} dy \end{aligned}$$

Which results in

$$S = \frac{\ln(x^2 + y^2)}{2} - \arctan\left(\frac{y}{x}\right)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{x+y}{-x+y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{x+y}{x^2+y^2} \\ S_y &= \frac{-x+y}{x^2+y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

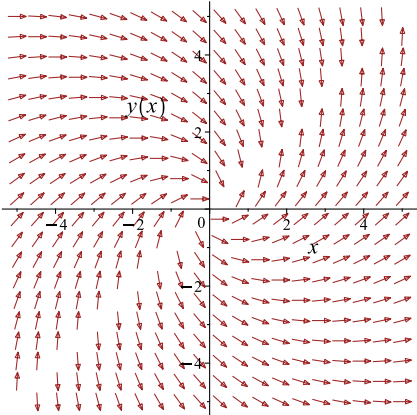
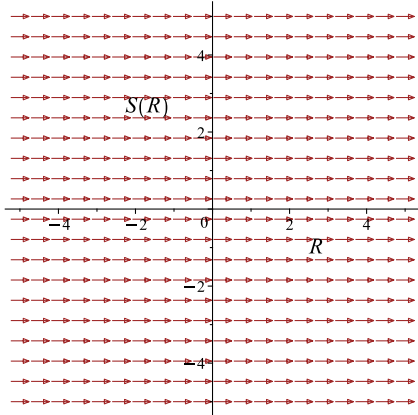
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\ln(x^2 + y^2)}{2} - \arctan\left(\frac{y}{x}\right) = c_1$$

Which simplifies to

$$\frac{\ln(x^2 + y^2)}{2} - \arctan\left(\frac{y}{x}\right) = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{x+y}{-x+y}$ 	$R = x$ $S = \frac{\ln(x^2 + y^2)}{2} - \arctan\left(\frac{y}{x}\right)$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$\frac{\ln(x^2 + y^2)}{2} - \arctan\left(\frac{y}{x}\right) = c_1 \quad (1)$$

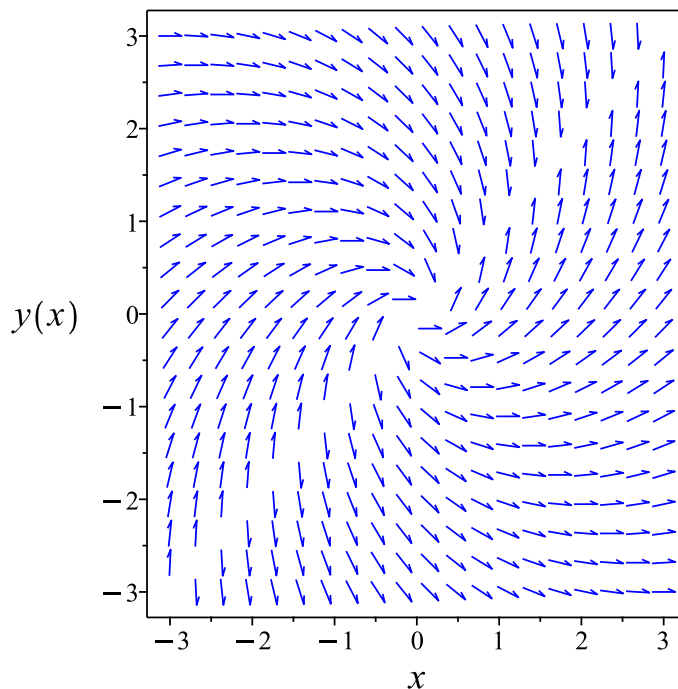


Figure 344: Slope field plot

Verification of solutions

$$\frac{\ln(x^2 + y^2)}{2} - \arctan\left(\frac{y}{x}\right) = c_1$$

Verified OK.

6.29.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (-x + y) dy &= (-x - y) dx \\ (x + y) dx + (-x + y) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= x + y \\ N(x, y) &= -x + y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(x + y) \\ &= 1 \end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x + y) \\ &= -1\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. By inspection $\frac{1}{x^2+y^2}$ is an integrating factor. Therefore by multiplying $M = x + y$ and $N = -x + y$ by this integrating factor the ode becomes exact. The new M, N are

$$\begin{aligned}M &= \frac{x + y}{x^2 + y^2} \\ N &= \frac{-x + y}{x^2 + y^2}\end{aligned}$$

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might

or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left(\frac{-x+y}{x^2+y^2}\right) dy &= \left(-\frac{x+y}{x^2+y^2}\right) dx \\ \left(\frac{x+y}{x^2+y^2}\right) dx + \left(\frac{-x+y}{x^2+y^2}\right) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= \frac{x+y}{x^2+y^2} \\ N(x, y) &= \frac{-x+y}{x^2+y^2} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{x+y}{x^2+y^2} \right) \\ &= \frac{x^2 - 2yx - y^2}{(x^2+y^2)^2} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{-x+y}{x^2+y^2} \right) \\ &= \frac{x^2 - 2yx - y^2}{(x^2+y^2)^2} \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{x+y}{x^2+y^2} dx \\ \phi &= \frac{\ln(x^2+y^2)}{2} + \arctan\left(\frac{x}{y}\right) + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned}\frac{\partial \phi}{\partial y} &= \frac{y}{x^2+y^2} - \frac{x}{y^2\left(\frac{x^2}{y^2}+1\right)} + f'(y) \\ &= \frac{-x+y}{x^2+y^2} + f'(y)\end{aligned}\quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{-x+y}{x^2+y^2}$. Therefore equation (4) becomes

$$\frac{-x+y}{x^2+y^2} = \frac{-x+y}{x^2+y^2} + f'(y)\quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{\ln(x^2+y^2)}{2} + \arctan\left(\frac{x}{y}\right) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = \frac{\ln(x^2+y^2)}{2} + \arctan\left(\frac{x}{y}\right)$$

Summary

The solution(s) found are the following

$$\frac{\ln(x^2 + y^2)}{2} + \arctan\left(\frac{x}{y}\right) = c_1 \quad (1)$$

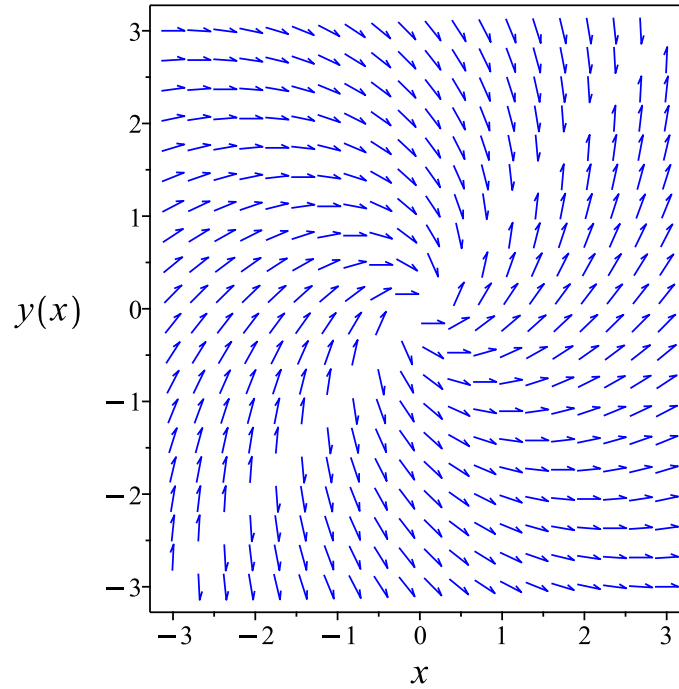


Figure 345: Slope field plot

Verification of solutions

$$\frac{\ln(x^2 + y^2)}{2} + \arctan\left(\frac{x}{y}\right) = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 24

```
dsolve(diff(y(x),x) = (x+y(x))/(x-y(x)),y(x), singsol=all)
```

$$y(x) = \tan \left(\text{RootOf} \left(-2_Z + \ln \left(\sec \left(_Z \right)^2 \right) + 2 \ln (x) + 2c_1 \right) \right) x$$

✓ Solution by Mathematica

Time used: 0.033 (sec). Leaf size: 36

```
DSolve[y'[x] == (x+y[x])/(x-y[x]),y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve} \left[\frac{1}{2} \log \left(\frac{y(x)^2}{x^2} + 1 \right) - \arctan \left(\frac{y(x)}{x} \right) = -\log(x) + c_1, y(x) \right]$$

6.30 problem 30

6.30.1 Solving as homogeneousTypeD2 ode	1739
6.30.2 Solving as first order ode lie symmetry calculated ode	1741
6.30.3 Solving as exact ode	1747

Internal problem ID [597]

Internal file name [OUTPUT/597_Sunday_June_05_2022_01_45_33_AM_62087244/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 30.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD2**", "**exactWithIntegrationFactor**", "**first_order_ode_lie_symmetry_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class B`]]
```

$$2yx + 3y^2 - (x^2 + 2yx) y' = 0$$

6.30.1 Solving as homogeneousTypeD2 ode

Using the change of variables $y = u(x)x$ on the above ode results in new ode in $u(x)$

$$2u(x)x^2 + 3u(x)^2x^2 - (x^2 + 2u(x)x^2)(u'(x)x + u(x)) = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{u(u+1)}{(2u+1)x} \end{aligned}$$

Where $f(x) = \frac{1}{x}$ and $g(u) = \frac{u(u+1)}{2u+1}$. Integrating both sides gives

$$\frac{1}{\frac{u(u+1)}{2u+1}} du = \frac{1}{x} dx$$
$$\int \frac{1}{\frac{u(u+1)}{2u+1}} du = \int \frac{1}{x} dx$$
$$\ln(u(u+1)) = \ln(x) + c_2$$

Raising both side to exponential gives

$$u(u+1) = e^{\ln(x)+c_2}$$

Which simplifies to

$$u(u+1) = c_3 x$$

Which simplifies to

$$u(x)(u(x)+1) = c_3 e^{c_2} x$$

The solution is

$$u(x)(u(x)+1) = c_3 e^{c_2} x$$

Replacing $u(x)$ in the above solution by $\frac{y}{x}$ results in the solution for y in implicit form

$$\frac{y(1 + \frac{y}{x})}{x} = c_3 e^{c_2} x$$
$$\frac{y(x+y)}{x^2} = c_3 e^{c_2} x$$

Summary

The solution(s) found are the following

$$\frac{y(x+y)}{x^2} = c_3 e^{c_2} x \quad (1)$$

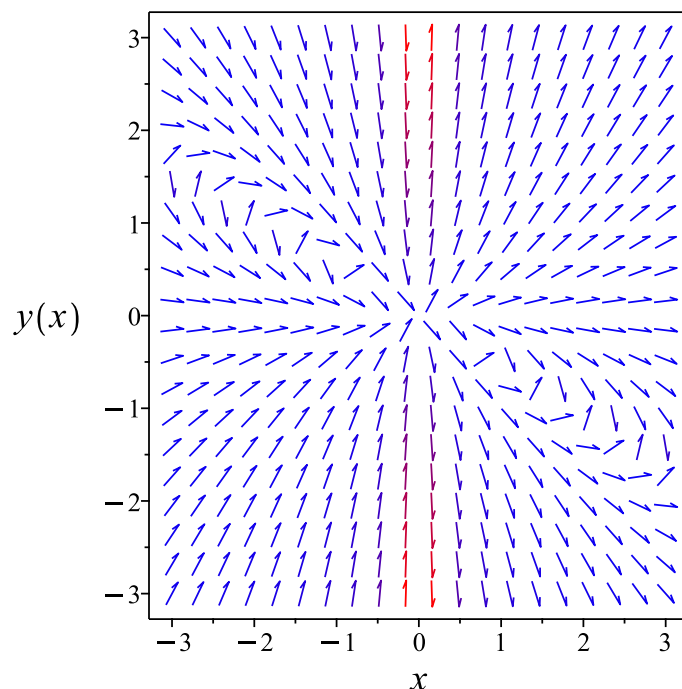


Figure 346: Slope field plot

Verification of solutions

$$\frac{y(x+y)}{x^2} = c_3 e^{c_2 x}$$

Verified OK.

6.30.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{y(2x+3y)}{x(x+2y)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 + \frac{y(2x+3y)(b_3-a_2)}{x(x+2y)} - \frac{y^2(2x+3y)^2 a_3}{x^2(x+2y)^2} \\ - \left(\frac{2y}{x(x+2y)} - \frac{y(2x+3y)}{x^2(x+2y)} - \frac{y(2x+3y)}{x(x+2y)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left(\frac{2x+3y}{x(x+2y)} + \frac{3y}{x(x+2y)} - \frac{2y(2x+3y)}{x(x+2y)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{x^4 b_2 + 2x^3 y b_2 + x^2 y^2 a_2 + 2x^2 y^2 a_3 + 2x^2 y^2 b_2 - x^2 y^2 b_3 + 6x y^3 a_3 + 3y^4 a_3 + 2x^3 b_1 - 2x^2 y a_1 + 6x^2 y b_1 - \dots}{x^2 (x+2y)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} -x^4 b_2 - 2x^3 y b_2 - x^2 y^2 a_2 - 2x^2 y^2 a_3 - 2x^2 y^2 b_2 + x^2 y^2 b_3 - 6x y^3 a_3 \\ - 3y^4 a_3 - 2x^3 b_1 + 2x^2 y a_1 - 6x^2 y b_1 + 6x y^2 a_1 - 6x y^2 b_1 + 6y^3 a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} -a_2 v_1^2 v_2^2 - 2a_3 v_1^2 v_2^2 - 6a_3 v_1 v_2^3 - 3a_3 v_2^4 - b_2 v_1^4 - 2b_2 v_1^3 v_2 - 2b_2 v_1^2 v_2^2 \\ + b_3 v_1^2 v_2^2 + 2a_1 v_1^2 v_2 + 6a_1 v_1 v_2^2 + 6a_1 v_2^3 - 2b_1 v_1^3 - 6b_1 v_1^2 v_2 - 6b_1 v_1 v_2^2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & -b_2v_1^4 - 2b_2v_1^3v_2 - 2b_1v_1^3 + (-a_2 - 2a_3 - 2b_2 + b_3)v_1^2v_2^2 \\ & + (2a_1 - 6b_1)v_1^2v_2 - 6a_3v_1v_2^3 + (6a_1 - 6b_1)v_1v_2^2 - 3a_3v_2^4 + 6a_1v_2^3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 6a_1 &= 0 \\ -6a_3 &= 0 \\ -3a_3 &= 0 \\ -2b_1 &= 0 \\ -2b_2 &= 0 \\ -b_2 &= 0 \\ 2a_1 - 6b_1 &= 0 \\ 6a_1 - 6b_1 &= 0 \\ -a_2 - 2a_3 - 2b_2 + b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left(\frac{y(2x + 3y)}{x(x + 2y)} \right) (x) \\ &= \frac{-yx - y^2}{x + 2y} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-yx - y^2}{x + 2y}} dy\end{aligned}$$

Which results in

$$S = -\ln(y(x + y))$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y(2x + 3y)}{x(x + 2y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= -\frac{1}{x+y} \\S_y &= \frac{-2y-x}{y(x+y)}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{3}{x} \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{3}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -3 \ln(R) + c_1 \tag{4}$$

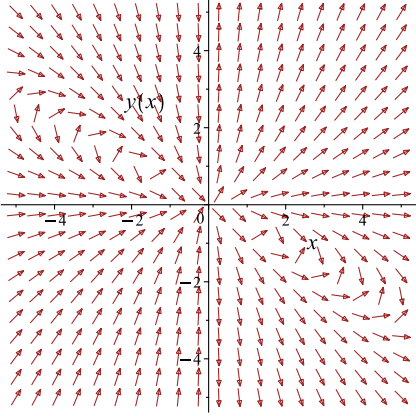
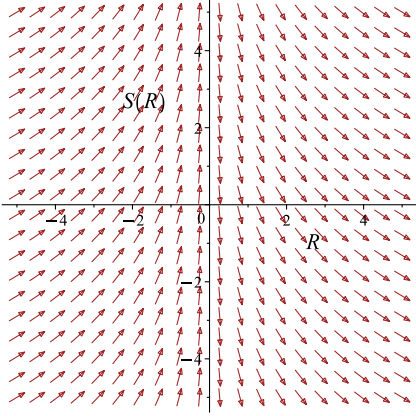
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$-\ln(y) - \ln(x+y) = -3 \ln(x) + c_1$$

Which simplifies to

$$-\ln(y) - \ln(x+y) = -3 \ln(x) + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{y(2x+3y)}{x(x+2y)}$ 	$R = x$ $S = -\ln(y) - \ln(x+y)$	$\frac{dS}{dR} = -\frac{3}{R}$ 

Summary

The solution(s) found are the following

$$-\ln(y) - \ln(x+y) = -3\ln(x) + c_1 \tag{1}$$

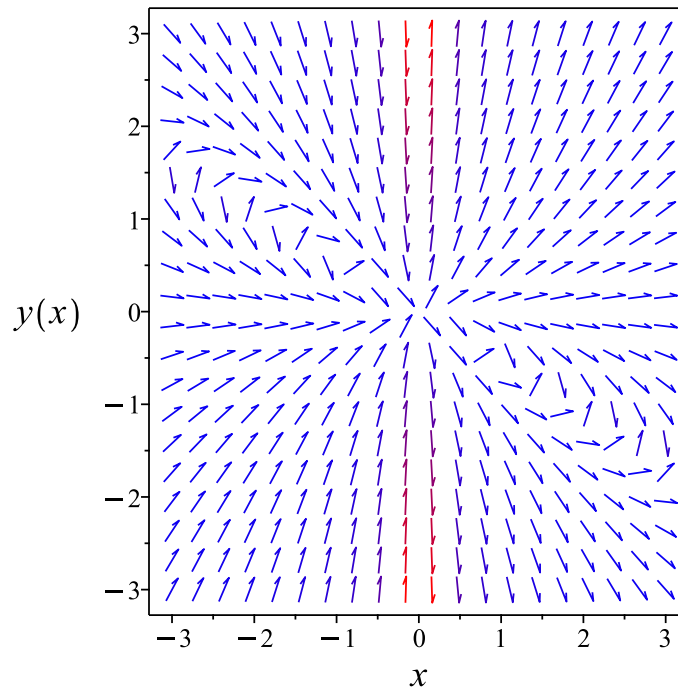


Figure 347: Slope field plot

Verification of solutions

$$-\ln(y) - \ln(x + y) = -3 \ln(x) + c_1$$

Verified OK.

6.30.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(-x^2 - 2yx) dy &= (-2yx - 3y^2) dx \\ (2yx + 3y^2) dx + (-x^2 - 2yx) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 2yx + 3y^2 \\ N(x, y) &= -x^2 - 2yx\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2yx + 3y^2) \\ &= 2x + 6y\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x^2 - 2yx) \\ &= -2x - 2y\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= -\frac{1}{x(x+2y)} ((2x+6y) - (-2x-2y)) \\ &= -\frac{4}{x} \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A \, dx} \\ &= e^{\int -\frac{4}{x} \, dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-4 \ln(x)} \\ &= \frac{1}{x^4} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{1}{x^4} (2yx + 3y^2) \\ &= \frac{y(2x + 3y)}{x^4} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \frac{1}{x^4} (-x^2 - 2yx) \\ &= \frac{-2y - x}{x^3} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left(\frac{y(2x + 3y)}{x^4} \right) + \left(\frac{-2y - x}{x^3} \right) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{y(2x + 3y)}{x^4} dx \\ \phi &= -\frac{y(x + y)}{x^3} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned} \frac{\partial \phi}{\partial y} &= -\frac{x + y}{x^3} - \frac{y}{x^3} + f'(y) \\ &= \frac{-2y - x}{x^3} + f'(y) \end{aligned} \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{-2y-x}{x^3}$. Therefore equation (4) becomes

$$\frac{-2y - x}{x^3} = \frac{-2y - x}{x^3} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{y(x + y)}{x^3} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = -\frac{y(x+y)}{x^3}$$

Summary

The solution(s) found are the following

$$-\frac{y(x+y)}{x^3} = c_1 \tag{1}$$

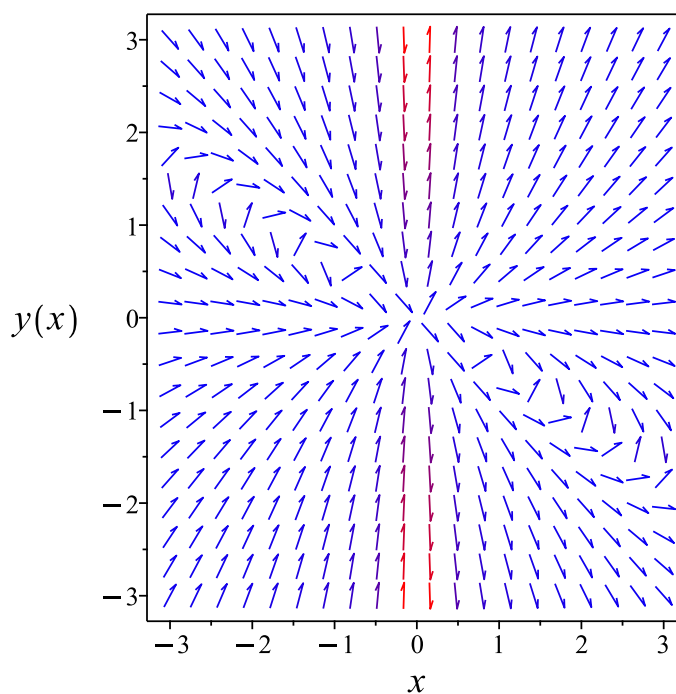


Figure 348: Slope field plot

Verification of solutions

$$-\frac{y(x+y)}{x^3} = c_1$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 33

```
dsolve(2*x*y(x)+3*y(x)^2-(x^2+2*x*y(x))*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = -\frac{(1 + \sqrt{4c_1x + 1})x}{2}$$
$$y(x) = \frac{(-1 + \sqrt{4c_1x + 1})x}{2}$$

✓ Solution by Mathematica

Time used: 0.408 (sec). Leaf size: 61

```
DSolve[2*x*y[x]+3*y[x]^2-(x^2+2*x*y[x])*y'[x] == 0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{1}{2}x(1 + \sqrt{1 + 4e^{c_1x}})$$
$$y(x) \rightarrow \frac{1}{2}x(-1 + \sqrt{1 + 4e^{c_1x}})$$
$$y(x) \rightarrow 0$$
$$y(x) \rightarrow -x$$

6.31 problem 31

6.31.1 Solving as first order ode lie symmetry calculated ode 1753

6.31.2 Solving as exact ode 1759

Internal problem ID [598]

Internal file name [OUTPUT/598_Sunday_June_05_2022_01_45_35_AM_42029649/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Miscellaneous problems, end of chapter 2. Page 133

Problem number: 31.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "exactWithIntegrationFactor",
"first_order_ode_lie_symmetry_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class G`], _rational, [_Abel, `2nd type`, `class B`]]
```

$$y' - \frac{-3x^2y - y^2}{2x^3 + 3yx} = 0$$

With initial conditions

$$[y(1) = -2]$$

6.31.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{y(3x^2 + y)}{x(2x^2 + 3y)}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1E)$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2E)$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 - \frac{y(3x^2 + y)(b_3 - a_2)}{x(2x^2 + 3y)} - \frac{y^2(3x^2 + y)^2 a_3}{x^2(2x^2 + 3y)^2} \\ - \left(-\frac{6y}{2x^2 + 3y} + \frac{y(3x^2 + y)}{x^2(2x^2 + 3y)} + \frac{4y(3x^2 + y)}{(2x^2 + 3y)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left(-\frac{3x^2 + y}{x(2x^2 + 3y)} - \frac{y}{x(2x^2 + 3y)} + \frac{3y(3x^2 + y)}{x(2x^2 + 3y)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{10x^6b_2 - 15x^4y^2a_3 + 6x^5b_1 - 6x^4ya_1 + 16x^4yb_2 + 14x^3y^2a_2 - 7x^3y^2b_3 - 3x^2y^3a_3 + 4x^3yb_1 + 3x^2y^2a_1 + 12x^2y^2b_2 - 4y^4a_3 + 3xy^2b_1 - 3y^3a_1}{(2x^2 + 3y)^2 x^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} 10x^6b_2 - 15x^4y^2a_3 + 6x^5b_1 - 6x^4ya_1 + 16x^4yb_2 + 14x^3y^2a_2 - 7x^3y^2b_3 \\ - 3x^2y^3a_3 + 4x^3yb_1 + 3x^2y^2a_1 + 12x^2y^2b_2 - 4y^4a_3 + 3xy^2b_1 - 3y^3a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & -15a_3v_1^4v_2^2 + 10b_2v_1^6 - 6a_1v_1^4v_2 + 14a_2v_1^3v_2^2 - 3a_3v_1^2v_2^3 + 6b_1v_1^5 + 16b_2v_1^4v_2 \\ & - 7b_3v_1^3v_2^2 + 3a_1v_1^2v_2^3 - 4a_3v_2^4 + 4b_1v_1^3v_2 + 12b_2v_1^2v_2^2 - 3a_1v_2^3 + 3b_1v_1v_2^2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & 10b_2v_1^6 + 6b_1v_1^5 - 15a_3v_1^4v_2^2 + (-6a_1 + 16b_2)v_1^4v_2 + (14a_2 - 7b_3)v_1^3v_2^2 \\ & + 4b_1v_1^3v_2 - 3a_3v_1^2v_2^3 + (3a_1 + 12b_2)v_1^2v_2^2 + 3b_1v_1v_2^2 - 4a_3v_2^4 - 3a_1v_2^3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} & -3a_1 = 0 \\ & -15a_3 = 0 \\ & -4a_3 = 0 \\ & -3a_3 = 0 \\ & 3b_1 = 0 \\ & 4b_1 = 0 \\ & 6b_1 = 0 \\ & 10b_2 = 0 \\ & -6a_1 + 16b_2 = 0 \\ & 3a_1 + 12b_2 = 0 \\ & 14a_2 - 7b_3 = 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} & a_1 = 0 \\ & a_2 = a_2 \\ & a_3 = 0 \\ & b_1 = 0 \\ & b_2 = 0 \\ & b_3 = 2a_2 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= x \\ \eta &= 2y\end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= 2y - \left(-\frac{y(3x^2 + y)}{x(2x^2 + 3y)} \right) (x) \\ &= \frac{7yx^2 + 7y^2}{2x^2 + 3y} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{7yx^2 + 7y^2}{2x^2 + 3y}} dy\end{aligned}$$

Which results in

$$S = \frac{\ln(x^2 + y)}{7} + \frac{2 \ln(y)}{7}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y) S_y}{R_x + \omega(x, y) R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y(3x^2 + y)}{x(2x^2 + 3y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{2x}{7x^2 + 7y} \\ S_y &= \frac{1}{7x^2 + 7y} + \frac{2}{7y} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{7x} \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{7R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S . Integrating the above gives

$$S(R) = -\frac{\ln(R)}{7} + c_1 \tag{4}$$

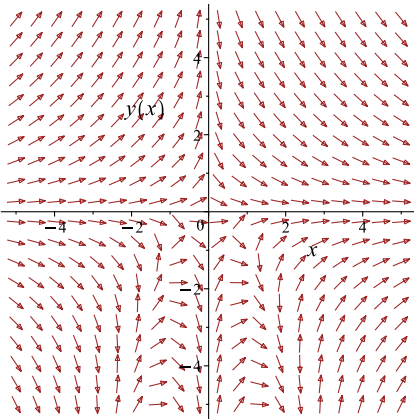
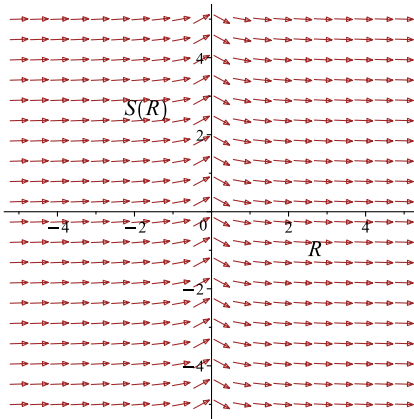
To complete the solution, we just need to transform (4) back to x, y coordinates. This results in

$$\frac{\ln(x^2 + y)}{7} + \frac{2\ln(y)}{7} = -\frac{\ln(x)}{7} + c_1$$

Which simplifies to

$$\frac{\ln(x^2 + y)}{7} + \frac{2\ln(y)}{7} = -\frac{\ln(x)}{7} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{y(3x^2+y)}{x(2x^2+3y)}$ 	$R = x$ $S = \frac{\ln(x^2 + y)}{7} + \frac{2 \ln(y)}{7}$	$\frac{dS}{dR} = -\frac{1}{7R}$ 

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$\frac{3i\pi}{7} + \frac{2 \ln(2)}{7} = c_1$$

$$c_1 = \frac{3i\pi}{7} + \frac{2 \ln(2)}{7}$$

Substituting c_1 found above in the general solution gives

$$\frac{\ln(x^2 + y)}{7} + \frac{2 \ln(y)}{7} = -\frac{\ln(x)}{7} + \frac{3i\pi}{7} + \frac{2 \ln(2)}{7}$$

Summary

The solution(s) found are the following

$$\frac{\ln(x^2 + y)}{7} + \frac{2 \ln(y)}{7} = -\frac{\ln(x)}{7} + \frac{3i\pi}{7} + \frac{2 \ln(2)}{7} \tag{1}$$

Verification of solutions

$$\frac{\ln(x^2 + y)}{7} + \frac{2 \ln(y)}{7} = -\frac{\ln(x)}{7} + \frac{3i\pi}{7} + \frac{2 \ln(2)}{7}$$

Verified OK.

6.31.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (x(2x^2 + 3y)) dy &= (-y(3x^2 + y)) dx \\ (y(3x^2 + y)) dx &+ (x(2x^2 + 3y)) dy = 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= y(3x^2 + y) \\ N(x, y) &= x(2x^2 + 3y) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y(3x^2 + y)) \\ &= 3x^2 + 2y\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x(2x^2 + 3y)) \\ &= 6x^2 + 3y\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{2x^3 + 3yx} ((3x^2 + 2y) - (6x^2 + 3y)) \\ &= \frac{-3x^2 - y}{2x^3 + 3yx}\end{aligned}$$

Since A depends on y , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned}B &= \frac{1}{M} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= \frac{1}{3yx^2 + y^2} ((6x^2 + 3y) - (3x^2 + 2y)) \\ &= \frac{1}{y}\end{aligned}$$

Since B does not depend on x , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned}\mu &= e^{\int B \, dy} \\ &= e^{\int \frac{1}{y} \, dy}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{\ln(y)} \\ &= y\end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} so not to confuse them with the original M and N .

$$\begin{aligned}\overline{M} &= \mu M \\ &= y(y(3x^2 + y)) \\ &= y^2(3x^2 + y)\end{aligned}$$

And

$$\begin{aligned}\overline{N} &= \mu N \\ &= y(x(2x^2 + 3y)) \\ &= yx(2x^2 + 3y)\end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned}\overline{M} + \overline{N} \frac{dy}{dx} &= 0 \\ (y^2(3x^2 + y)) + (yx(2x^2 + 3y)) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int y^2(3x^2 + y) dx \\ \phi &= y^2x(x^2 + y) + f(y)\end{aligned} \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned}\frac{\partial\phi}{\partial y} &= 2yx(x^2 + y) + xy^2 + f'(y) \\ &= yx(2x^2 + 3y) + f'(y)\end{aligned}\tag{4}$$

But equation (2) says that $\frac{\partial\phi}{\partial y} = yx(2x^2 + 3y)$. Therefore equation (4) becomes

$$yx(2x^2 + 3y) = yx(2x^2 + 3y) + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = y^2x(x^2 + y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into new constant c_1 gives the solution as

$$c_1 = y^2x(x^2 + y)$$

Initial conditions are used to solve for c_1 . Substituting $x = 1$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$-4 = c_1$$

$$c_1 = -4$$

Substituting c_1 found above in the general solution gives

$$y^2x(x^2 + y) = -4$$

Summary

The solution(s) found are the following

$$y^2x(x^2 + y) = -4\tag{1}$$

Verification of solutions

$$y^2x(x^2 + y) = -4$$

Verified OK.

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous G  
<- homogeneous successful`
```

✓ Solution by Maple

Time used: 1.266 (sec). Leaf size: 111

```
dsolve([diff(y(x),x) = (-3*x^2*y(x)-y(x)^2)/(2*x^3+3*x*y(x)),y(1) = -2],y(x), singsol=all)
```

$$y(x) = \frac{(i\sqrt{3} - 1) \left(-(x^7 - 6\sqrt{3}\sqrt{x^7 + 27} + 54) x^2 \right)^{\frac{2}{3}} - x^3 \left(i\sqrt{3} x^3 + x^3 + 2 \left(-(x^7 - 6\sqrt{3}\sqrt{x^7 + 27} + 54) x^2 \right)^{\frac{1}{3}} \right)}{6 \left(-(x^7 - 6\sqrt{3}\sqrt{x^7 + 27} + 54) x^2 \right)^{\frac{1}{3}} x}$$

✓ Solution by Mathematica

Time used: 40.923 (sec). Leaf size: 136

```
DSolve[{y'[x]== (-3*x^2*y[x]-y[x]^2)/(2*x^3+3*x*y[x]),y[1]==-2},y[x],x,IncludeSingularSoluti
```

$$y(x) \rightarrow \frac{i \left((\sqrt{3} + i) x^3 - (\sqrt{3} - i) x^3 + (\sqrt{3} + i) \sqrt[3]{-x^9 - 54x^2 + 6\sqrt{3}\sqrt{x^4(x^7 + 27)}} - \frac{(\sqrt{3} - i)x^6}{\sqrt[3]{-x^9 - 54x^2 + 6\sqrt{3}}} \right)}{6x}$$

7 Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

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7.3	problem 3	1781
7.4	problem 4	1789
7.5	problem 5	1797
7.6	problem 6	1812
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7.1 problem 1

- 7.1.1 Solving as second order linear constant coeff ode 1765
- 7.1.2 Solving using Kovacic algorithm 1767
- 7.1.3 Maple step by step solution 1771

Internal problem ID [599]

Internal file name [OUTPUT/599_Sunday_June_05_2022_01_45_39_AM_9710673/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 1.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 2y' - 3y = 0$$

7.1.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 2, C = -3$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 2\lambda e^{\lambda x} - 3e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 2\lambda - 3 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 2, C = -3$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{-2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{2^2 - (4)(1)(-3)} \\ &= -1 \pm 2\end{aligned}$$

Hence

$$\lambda_1 = -1 + 2$$

$$\lambda_2 = -1 - 2$$

Which simplifies to

$$\lambda_1 = 1$$

$$\lambda_2 = -3$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(1)x} + c_2 e^{(-3)x}$$

Or

$$y = c_1 e^x + e^{-3x} c_2$$

Summary

The solution(s) found are the following

$$y = c_1 e^x + e^{-3x} c_2 \tag{1}$$

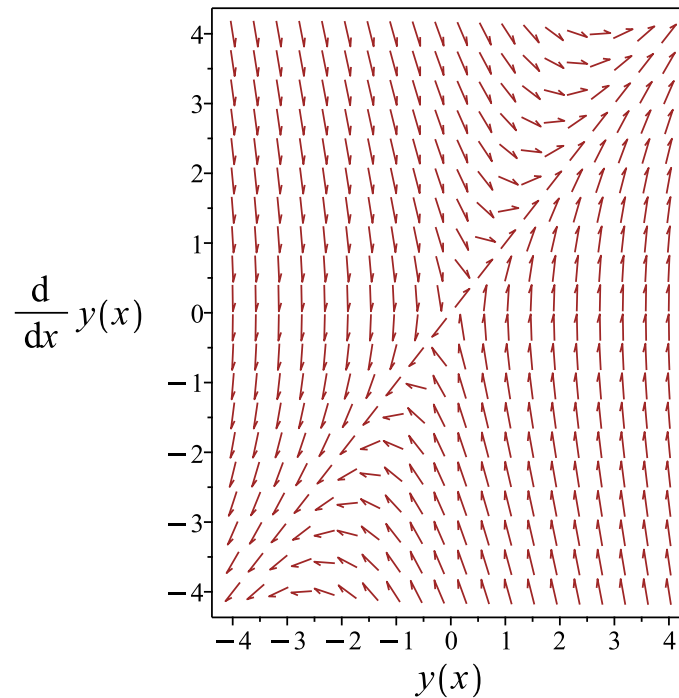


Figure 349: Slope field plot

Verification of solutions

$$y = c_1 e^x + e^{-3x} c_2$$

Verified OK.

7.1.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 2y' - 3y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 2 \\ C &= -3 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{4}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 4$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = 4z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 316: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 4$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-2x}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{2}{1} dx} \\ &= z_1 e^{-x} \\ &= z_1 (e^{-x})\end{aligned}$$

Which simplifies to

$$y_1 = e^{-3x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-2x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{e^{4x}}{4} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-3x}) + c_2 \left(e^{-3x} \left(\frac{e^{4x}}{4} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-3x} + \frac{c_2 e^x}{4} \quad (1)$$

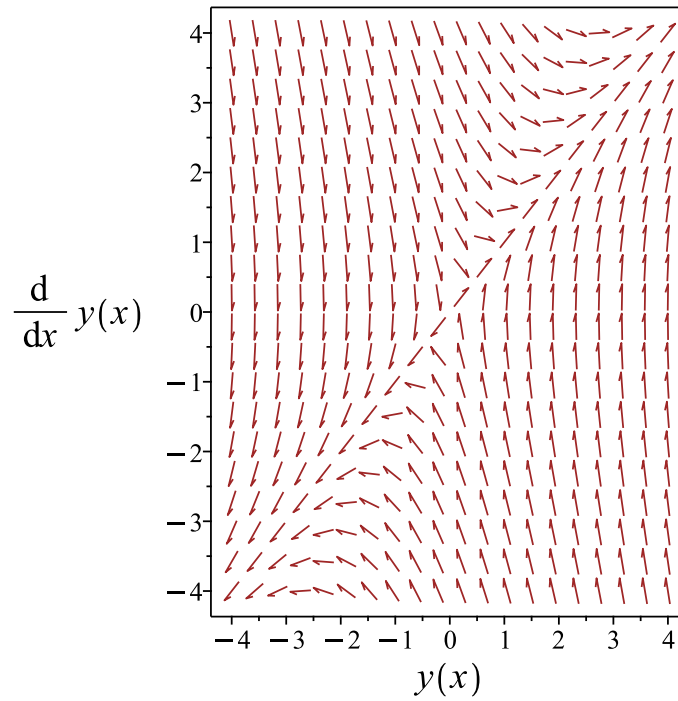


Figure 350: Slope field plot

Verification of solutions

$$y = c_1 e^{-3x} + \frac{c_2 e^x}{4}$$

Verified OK.

7.1.3 Maple step by step solution

Let's solve

$$y'' + 2y' - 3y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 2r - 3 = 0$$

- Factor the characteristic polynomial

$$(r + 3)(r - 1) = 0$$

- Roots of the characteristic polynomial

$$r = (-3, 1)$$

- 1st solution of the ODE

$$y_1(x) = e^{-3x}$$

- 2nd solution of the ODE

$$y_2(x) = e^x$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^{-3x} + c_2e^x$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 17

```
dsolve(diff(y(x),x$2) +2*diff(y(x),x)-3*y(x) = 0,y(x), singsol=all)
```

$$y(x) = (c_1 e^{4x} + c_2) e^{-3x}$$

✓ Solution by Mathematica

Time used: 0.011 (sec). Leaf size: 20

```
DSolve[y''[x]+2*y'[x]-3*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow c_1 e^{-3x} + c_2 e^x$$

7.2 problem 2

7.2.1 Solving as second order linear constant coeff ode	1773
7.2.2 Solving using Kovacic algorithm	1775
7.2.3 Maple step by step solution	1779

Internal problem ID [600]

Internal file name [OUTPUT/600_Sunday_June_05_2022_01_45_39_AM_98928217/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 2.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 3y' + 2y = 0$$

7.2.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 3, C = 2$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 3\lambda e^{\lambda x} + 2e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 3\lambda + 2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 3, C = 2$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{-3}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{3^2 - (4)(1)(2)} \\ &= -\frac{3}{2} \pm \frac{1}{2}\end{aligned}$$

Hence

$$\begin{aligned}\lambda_1 &= -\frac{3}{2} + \frac{1}{2} \\ \lambda_2 &= -\frac{3}{2} - \frac{1}{2}\end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= -1 \\ \lambda_2 &= -2\end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned}y &= c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x} \\ y &= c_1 e^{(-1)x} + c_2 e^{(-2)x}\end{aligned}$$

Or

$$y = c_1 e^{-x} + c_2 e^{-2x}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-x} + c_2 e^{-2x} \tag{1}$$

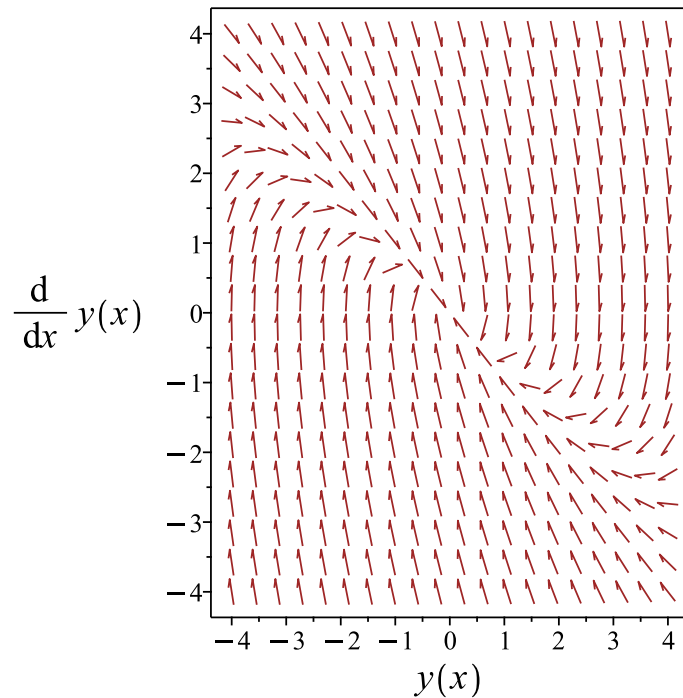


Figure 351: Slope field plot

Verification of solutions

$$y = c_1 e^{-x} + c_2 e^{-2x}$$

Verified OK.

7.2.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 3y' + 2y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 3 \\ C &= 2 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{4} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{z(x)}{4} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 318: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{3}{1} dx} \\ &= z_1 e^{-\frac{3x}{2}} \\ &= z_1 \left(e^{-\frac{3x}{2}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = e^{-2x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{3}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-3x}}{(y_1)^2} dx \\ &= y_1(e^x)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1(e^{-2x}) + c_2(e^{-2x}(e^x))\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-2x} + c_2 e^{-x} \tag{1}$$

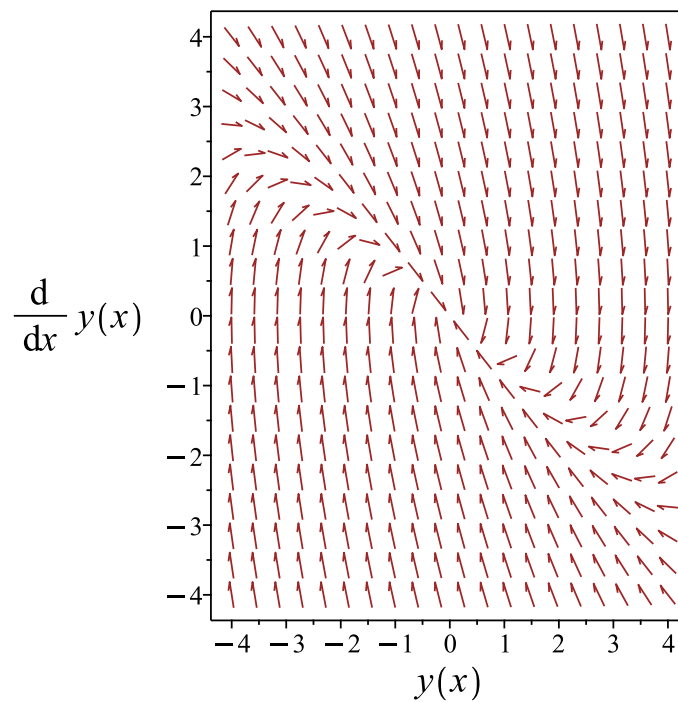


Figure 352: Slope field plot

Verification of solutions

$$y = c_1 e^{-2x} + c_2 e^{-x}$$

Verified OK.

7.2.3 Maple step by step solution

Let's solve

$$y'' + 3y' + 2y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 3r + 2 = 0$$

- Factor the characteristic polynomial

$$(r + 2)(r + 1) = 0$$

- Roots of the characteristic polynomial

$$r = (-2, -1)$$

- 1st solution of the ODE

$$y_1(x) = e^{-2x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{-x}$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^{-2x} + c_2e^{-x}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(diff(y(x),x$2) +3*diff(y(x),x)+2*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-2x}c_1 + c_2e^{-x}$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 20

```
DSolve[y''[x]+3*y'[x]+2*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-2x}(c_2e^x + c_1)$$

7.3 problem 3

7.3.1 Solving as second order linear constant coeff ode	1781
7.3.2 Solving using Kovacic algorithm	1783
7.3.3 Maple step by step solution	1787

Internal problem ID [601]

Internal file name [OUTPUT/601_Sunday_June_05_2022_01_45_40_AM_67843365/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 3.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$6y'' - y' - y = 0$$

7.3.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 6, B = -1, C = -1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$6\lambda^2 e^{\lambda x} - \lambda e^{\lambda x} - e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$6\lambda^2 - \lambda - 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 6, B = -1, C = -1$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{1}{(2)(6)} \pm \frac{1}{(2)(6)} \sqrt{-1^2 - (4)(6)(-1)} \\ &= \frac{1}{12} \pm \frac{5}{12}\end{aligned}$$

Hence

$$\begin{aligned}\lambda_1 &= \frac{1}{12} + \frac{5}{12} \\ \lambda_2 &= \frac{1}{12} - \frac{5}{12}\end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= \frac{1}{2} \\ \lambda_2 &= -\frac{1}{3}\end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned}y &= c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x} \\ y &= c_1 e^{(\frac{1}{2})x} + c_2 e^{(-\frac{1}{3})x}\end{aligned}$$

Or

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{-\frac{x}{3}}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{-\frac{x}{3}} \quad (1)$$

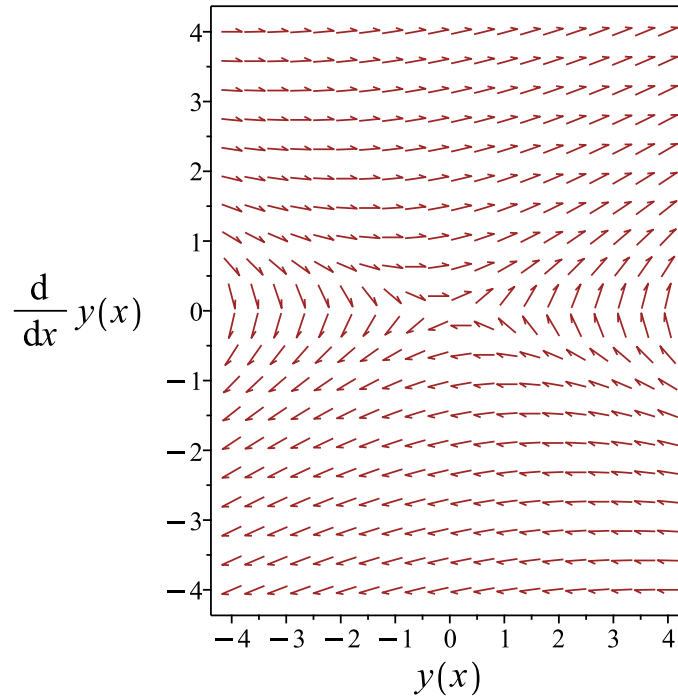


Figure 353: Slope field plot

Verification of solutions

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{-\frac{x}{3}}$$

Verified OK.

7.3.2 Solving using Kovacic algorithm

Writing the ode as

$$6y'' - y' - y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 6 \\ B &= -1 \\ C &= -1 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{25}{144} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 25 \\ t &= 144 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{25z(x)}{144} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 320: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{25}{144}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{5x}{12}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-1}{6} dx} \\ &= z_1 e^{\frac{x}{12}} \\ &= z_1 \left(e^{\frac{x}{12}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{x}{3}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-1}{6} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{\frac{x}{6}}}{(y_1)^2} dx \\ &= y_1 \left(\frac{6 e^{\frac{5x}{6}}}{5} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-\frac{x}{3}} \right) + c_2 \left(e^{-\frac{x}{3}} \left(\frac{6 e^{\frac{5x}{6}}}{5} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{x}{3}} + \frac{6c_2 e^{\frac{x}{2}}}{5} \quad (1)$$

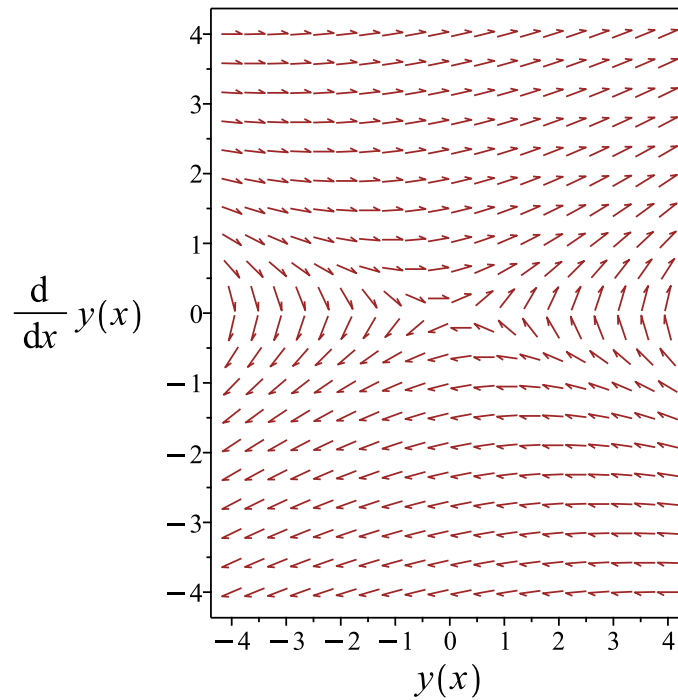


Figure 354: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{x}{3}} + \frac{6c_2 e^{\frac{x}{2}}}{5}$$

Verified OK.

7.3.3 Maple step by step solution

Let's solve

$$6y'' - y' - y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{y'}{6} + \frac{y}{6}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{y'}{6} - \frac{y}{6} = 0$$

- Characteristic polynomial of ODE

$$r^2 - \frac{1}{6}r - \frac{1}{6} = 0$$

- Factor the characteristic polynomial

$$\frac{(3r+1)(2r-1)}{6} = 0$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{3}, \frac{1}{2}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{x}{3}}$$

- 2nd solution of the ODE

$$y_2(x) = e^{\frac{x}{2}}$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^{-\frac{x}{3}} + c_2e^{\frac{x}{2}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(6*diff(y(x),x$2) -diff(y(x),x)-y(x) = 0,y(x), singsol=all)
```

$$y(x) = \left(c_1 e^{\frac{5x}{6}} + c_2\right) e^{-\frac{x}{3}}$$

✓ Solution by Mathematica

Time used: 0.014 (sec). Leaf size: 26

```
DSolve[6*y''[x]-y'[x]-y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x/3} (c_2 e^{5x/6} + c_1)$$

7.4 problem 4

7.4.1	Solving as second order linear constant coeff ode	1789
7.4.2	Solving using Kovacic algorithm	1791
7.4.3	Maple step by step solution	1795

Internal problem ID [602]

Internal file name [OUTPUT/602_Sunday_June_05_2022_01_45_41_AM_23423378/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 4.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$2y'' - 3y' + y = 0$$

7.4.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 2, B = -3, C = 1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$2\lambda^2 e^{\lambda x} - 3\lambda e^{\lambda x} + e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$2\lambda^2 - 3\lambda + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 2, B = -3, C = 1$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{3}{(2)(2)} \pm \frac{1}{(2)(2)} \sqrt{-3^2 - (4)(2)(1)} \\ &= \frac{3}{4} \pm \frac{1}{4}\end{aligned}$$

Hence

$$\lambda_1 = \frac{3}{4} + \frac{1}{4}$$

$$\lambda_2 = \frac{3}{4} - \frac{1}{4}$$

Which simplifies to

$$\lambda_1 = 1$$

$$\lambda_2 = \frac{1}{2}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(1)x} + c_2 e^{(\frac{1}{2})x}$$

Or

$$y = c_1 e^x + c_2 e^{\frac{x}{2}}$$

Summary

The solution(s) found are the following

$$y = c_1 e^x + c_2 e^{\frac{x}{2}} \quad (1)$$

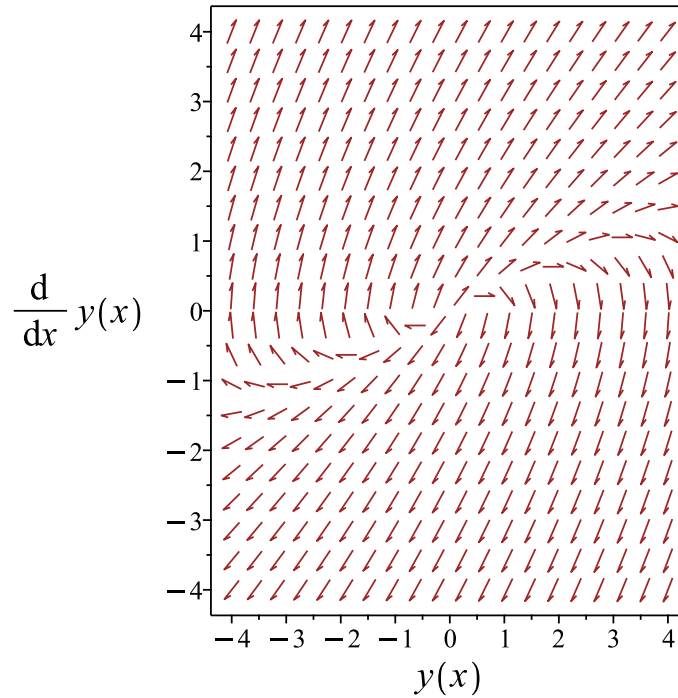


Figure 355: Slope field plot

Verification of solutions

$$y = c_1 e^x + c_2 e^{\frac{x}{2}}$$

Verified OK.

7.4.2 Solving using Kovacic algorithm

Writing the ode as

$$2y'' - 3y' + y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 2 \\ B &= -3 \\ C &= 1 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{16} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 16 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{z(x)}{16} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 322: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{16}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x}{4}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-3}{2} dx} \\ &= z_1 e^{\frac{3x}{4}} \\ &= z_1 \left(e^{\frac{3x}{4}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{\frac{x}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-3}{2} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{\frac{3x}{2}}}{(y_1)^2} dx \\ &= y_1 (2 e^{\frac{x}{2}}) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{\frac{x}{2}} \right) + c_2 \left(e^{\frac{x}{2}} (2e^{\frac{x}{2}}) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{\frac{x}{2}} + 2c_2 e^x \tag{1}$$

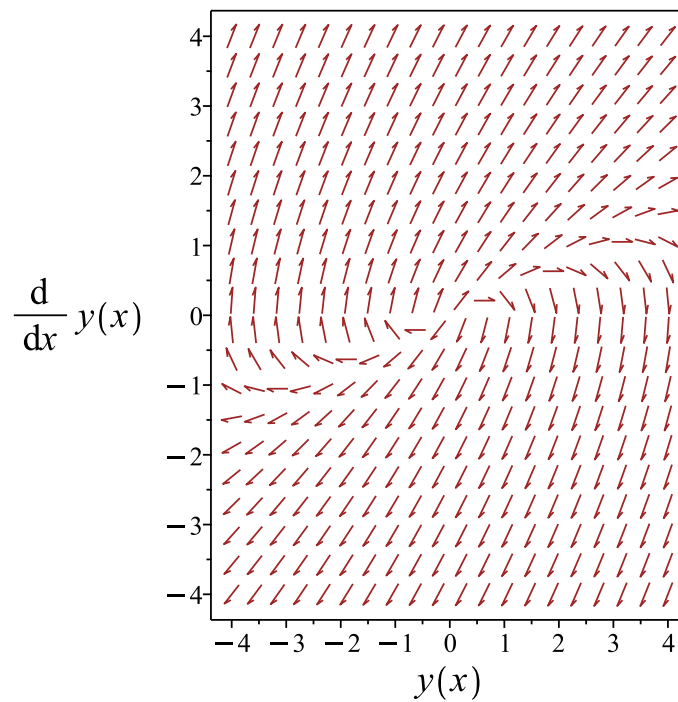


Figure 356: Slope field plot

Verification of solutions

$$y = c_1 e^{\frac{x}{2}} + 2c_2 e^x$$

Verified OK.

7.4.3 Maple step by step solution

Let's solve

$$2y'' - 3y' + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{3y'}{2} - \frac{y}{2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{3y'}{2} + \frac{y}{2} = 0$$

- Characteristic polynomial of ODE

$$r^2 - \frac{3}{2}r + \frac{1}{2} = 0$$

- Factor the characteristic polynomial

$$\frac{(2r-1)(r-1)}{2} = 0$$

- Roots of the characteristic polynomial

$$r = \left(1, \frac{1}{2}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^x$$

- 2nd solution of the ODE

$$y_2(x) = e^{\frac{x}{2}}$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^x + c_2e^{\frac{x}{2}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve(2*diff(y(x),x$2) -3*diff(y(x),x)+y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^x c_1 + c_2 e^{\frac{x}{2}}$$

✓ Solution by Mathematica

Time used: 0.021 (sec). Leaf size: 35

```
DSolve[y''[x]-3*y'[x]+y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-\frac{1}{2}(\sqrt{5}-3)x} \left(c_2 e^{\sqrt{5}x} + c_1 \right)$$

7.5 problem 5

7.5.1	Solving as second order linear constant coeff ode	1797
7.5.2	Solving as second order integrable as is ode	1799
7.5.3	Solving as second order ode missing y ode	1801
7.5.4	Solving as type second_order_integrable_as_is (not using ABC version)	1802
7.5.5	Solving using Kovacic algorithm	1804
7.5.6	Solving as exact linear second order ode ode	1807
7.5.7	Maple step by step solution	1810

Internal problem ID [603]

Internal file name [OUTPUT/603_Sunday_June_05_2022_01_45_42_AM_87737464/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 5.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "exact linear second order ode", "second_order_integrable_as_is", "second_order_ode_missing_y", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$y'' + 5y' = 0$$

7.5.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 5, C = 0$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 5\lambda e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 5\lambda = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 5, C = 0$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-5}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{5^2 - (4)(1)(0)} \\ &= -\frac{5}{2} \pm \frac{5}{2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= -\frac{5}{2} + \frac{5}{2} \\ \lambda_2 &= -\frac{5}{2} - \frac{5}{2} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= 0 \\ \lambda_2 &= -5 \end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned} y &= c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x} \\ y &= c_1 e^{(0)x} + c_2 e^{(-5)x} \end{aligned}$$

Or

$$y = c_1 + c_2 e^{-5x}$$

Summary

The solution(s) found are the following

$$y = c_1 + c_2 e^{-5x} \quad (1)$$

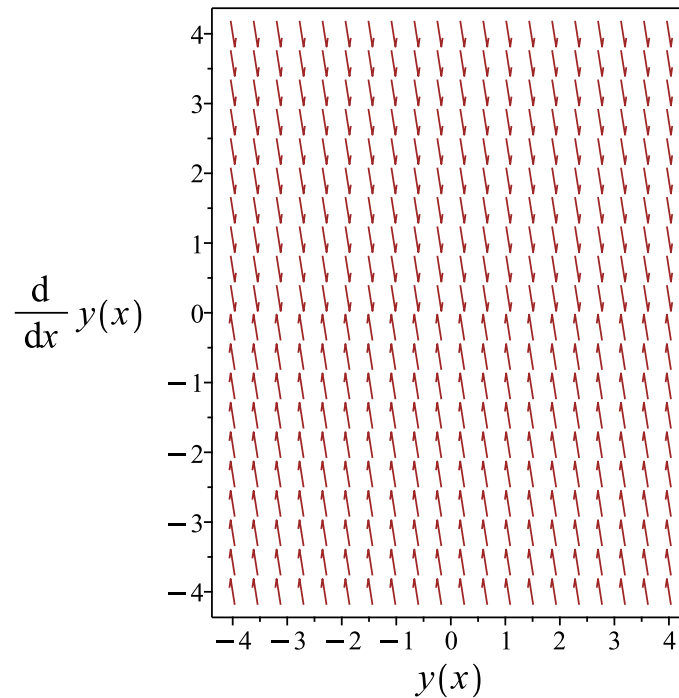


Figure 357: Slope field plot

Verification of solutions

$$y = c_1 + c_2 e^{-5x}$$

Verified OK.

7.5.2 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t x gives

$$\int (y'' + 5y') dx = 0$$

$$5y + y' = c_1$$

Which is now solved for y . Integrating both sides gives

$$\int \frac{1}{-5y + c_1} dy = \int dx$$

$$-\frac{\ln(-5y + c_1)}{5} = x + c_2$$

Raising both side to exponential gives

$$\frac{1}{(-5y + c_1)^{\frac{1}{5}}} = e^{x+c_2}$$

Which simplifies to

$$\frac{1}{(-5y + c_1)^{\frac{1}{5}}} = c_3 e^x$$

Summary

The solution(s) found are the following

$$y = -\frac{e^{-5x}}{5c_3^5} + \frac{c_1}{5} \quad (1)$$

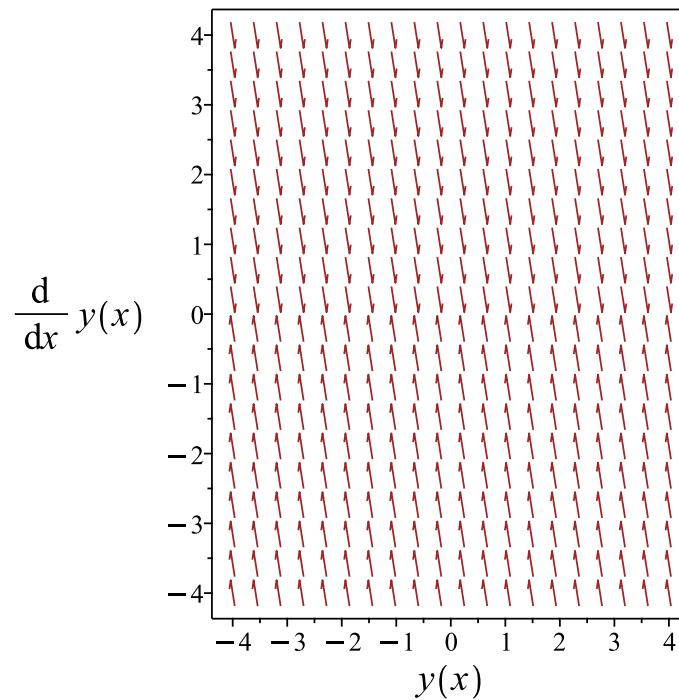


Figure 358: Slope field plot

Verification of solutions

$$y = -\frac{e^{-5x}}{5c_3^5} + \frac{c_1}{5}$$

Verified OK.

7.5.3 Solving as second order ode missing y ode

This is second order ode with missing dependent variable y . Let

$$p(x) = y'$$

Then

$$p'(x) = y''$$

Hence the ode becomes

$$p'(x) + 5p(x) = 0$$

Which is now solve for $p(x)$ as first order ode. Integrating both sides gives

$$\int -\frac{1}{5p} dp = \int dx$$
$$-\frac{\ln(p)}{5} = x + c_1$$

Raising both side to exponential gives

$$\frac{1}{p^{\frac{1}{5}}} = e^{x+c_1}$$

Which simplifies to

$$\frac{1}{p^{\frac{1}{5}}} = c_2 e^x$$

Since $p = y'$ then the new first order ode to solve is

$$y' = \frac{e^{-5x}}{c_2^5}$$

Integrating both sides gives

$$y = \int \frac{e^{-5x}}{c_2^5} dx$$
$$= -\frac{e^{-5x}}{5c_2^5} + c_3$$

Summary

The solution(s) found are the following

$$y = -\frac{e^{-5x}}{5c_2^5} + c_3 \quad (1)$$

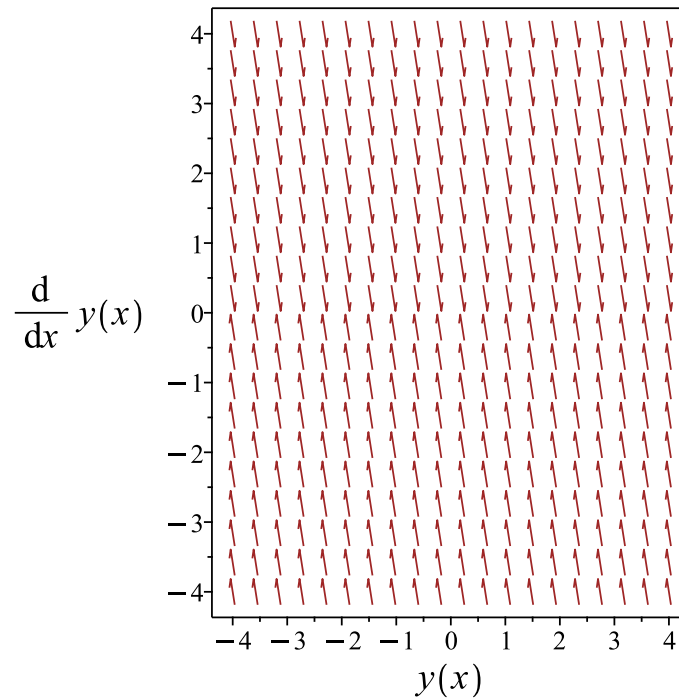


Figure 359: Slope field plot

Verification of solutions

$$y = -\frac{e^{-5x}}{5c_2^5} + c_3$$

Verified OK.

7.5.4 Solving as type second_order_integrable_as_is (not using ABC version)

Writing the ode as

$$y'' + 5y' = 0$$

Integrating both sides of the ODE w.r.t x gives

$$\int (y'' + 5y') dx = 0$$

$$5y + y' = c_1$$

Which is now solved for y . Integrating both sides gives

$$\int \frac{1}{-5y + c_1} dy = \int dx$$

$$-\frac{\ln(-5y + c_1)}{5} = x + c_2$$

Raising both side to exponential gives

$$\frac{1}{(-5y + c_1)^{\frac{1}{5}}} = e^{x+c_2}$$

Which simplifies to

$$\frac{1}{(-5y + c_1)^{\frac{1}{5}}} = c_3 e^x$$

Summary

The solution(s) found are the following

$$y = -\frac{e^{-5x}}{5c_3^5} + \frac{c_1}{5} \tag{1}$$

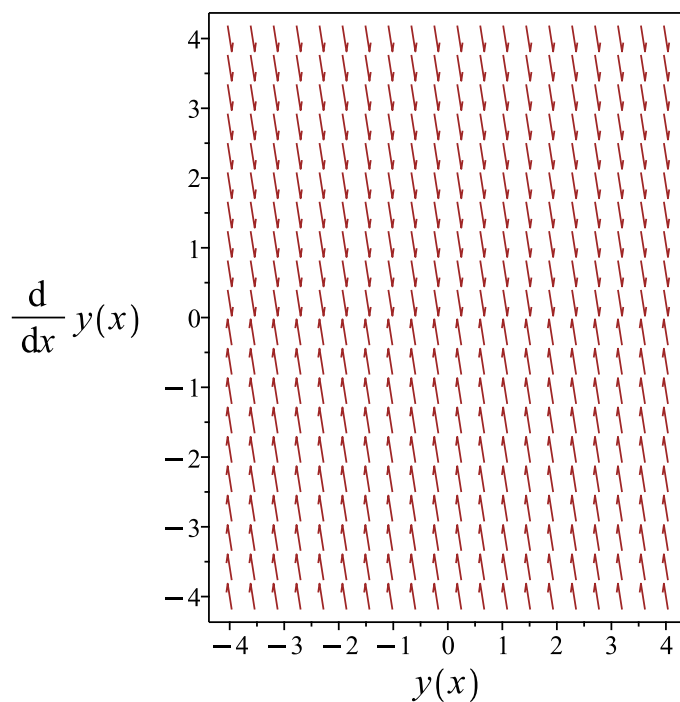


Figure 360: Slope field plot

Verification of solutions

$$y = -\frac{e^{-5x}}{5c_3^5} + \frac{c_1}{5}$$

Verified OK.

7.5.5 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 5y' = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 5 \tag{3}$$

$$C = 0$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} \tag{5}$$

$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{25}{4} \tag{6}$$

Comparing the above to (5) shows that

$$s = 25$$

$$t = 4$$

Therefore eq. (4) becomes

$$z''(x) = \frac{25z(x)}{4} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 324: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{25}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{5x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\&= z_1 e^{-\int \frac{1}{2} \frac{5}{1} dx} \\&= z_1 e^{-\frac{5x}{2}} \\&= z_1 \left(e^{-\frac{5x}{2}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = e^{-5x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{5}{1} dx}}{(y_1)^2} dx \\&= y_1 \int \frac{e^{-5x}}{(y_1)^2} dx \\&= y_1 \left(\frac{e^{5x}}{5} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 (e^{-5x}) + c_2 \left(e^{-5x} \left(\frac{e^{5x}}{5} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-5x} + \frac{c_2}{5} \quad (1)$$

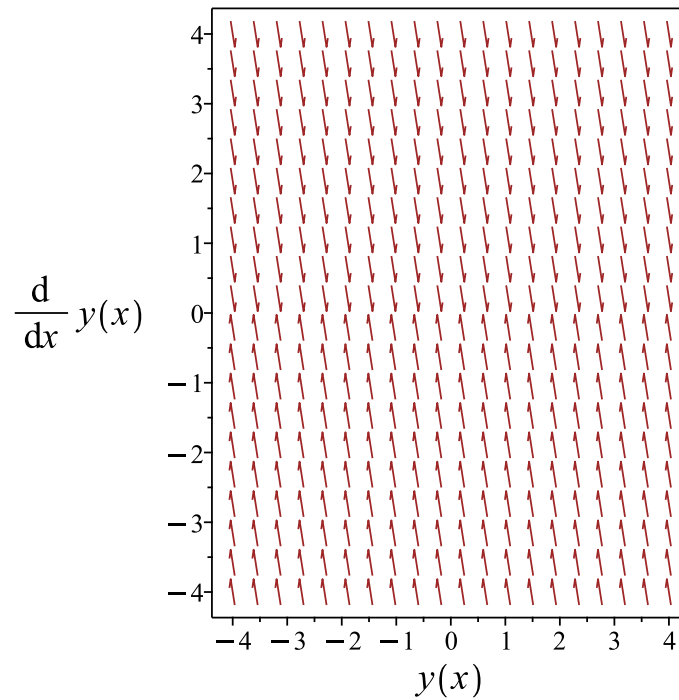


Figure 361: Slope field plot

Verification of solutions

$$y = c_1 e^{-5x} + \frac{c_2}{5}$$

Verified OK.

7.5.6 Solving as exact linear second order ode ode

An ode of the form

$$p(x) y'' + q(x) y' + r(x) y = s(x)$$

is exact if

$$p''(x) - q'(x) + r(x) = 0 \quad (1)$$

For the given ode we have

$$p(x) = 1$$

$$q(x) = 5$$

$$r(x) = 0$$

$$s(x) = 0$$

Hence

$$p''(x) = 0$$

$$q'(x) = 0$$

Therefore (1) becomes

$$0 - (0) + (0) = 0$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$(p(x)y' + (q(x) - p'(x))y)' = s(x)$$

Integrating gives

$$p(x)y' + (q(x) - p'(x))y = \int s(x) dx$$

Substituting the above values for p, q, r, s gives

$$5y + y' = c_1$$

We now have a first order ode to solve which is

$$5y + y' = c_1$$

Integrating both sides gives

$$\int \frac{1}{-5y + c_1} dy = \int dx$$
$$-\frac{\ln(-5y + c_1)}{5} = x + c_2$$

Raising both side to exponential gives

$$\frac{1}{(-5y + c_1)^{\frac{1}{5}}} = e^{x+c_2}$$

Which simplifies to

$$\frac{1}{(-5y + c_1)^{\frac{1}{5}}} = c_3 e^x$$

Summary

The solution(s) found are the following

$$y = -\frac{e^{-5x}}{5c_3^5} + \frac{c_1}{5} \quad (1)$$

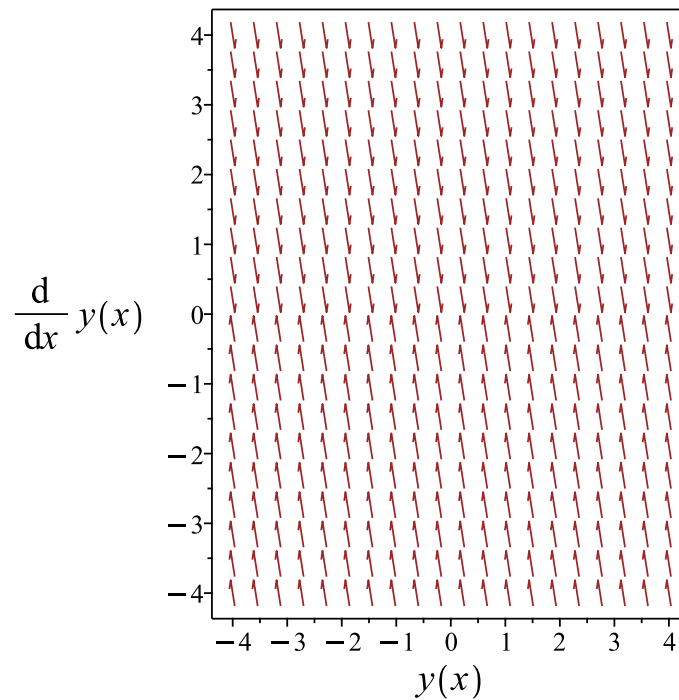


Figure 362: Slope field plot

Verification of solutions

$$y = -\frac{e^{-5x}}{5c_3^5} + \frac{c_1}{5}$$

Verified OK.

7.5.7 Maple step by step solution

Let's solve

$$y'' + 5y' = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 5r = 0$$

- Factor the characteristic polynomial

$$r(r + 5) = 0$$

- Roots of the characteristic polynomial

$$r = (-5, 0)$$

- 1st solution of the ODE

$$y_1(x) = e^{-5x}$$

- 2nd solution of the ODE

$$y_2(x) = 1$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^{-5x} + c_2$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 12

```
dsolve(diff(y(x),x$2) +5*diff(y(x),x) = 0,y(x), singsol=all)
```

$$y(x) = c_1 + c_2 e^{-5x}$$

✓ Solution by Mathematica

Time used: 0.01 (sec). Leaf size: 19

```
DSolve[y''[x]+5*y'[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow c_2 - \frac{1}{5}c_1 e^{-5x}$$

7.6 problem 6

7.6.1	Solving as second order linear constant coeff ode	1812
7.6.2	Solving as second order ode can be made integrable ode	1814
7.6.3	Solving using Kovacic algorithm	1817
7.6.4	Maple step by step solution	1820

Internal problem ID [604]

Internal file name [OUTPUT/604_Sunday_June_05_2022_01_45_42_AM_9328719/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 6.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "second_order_ode_can_be_made_integrable"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$4y'' - 9y = 0$$

7.6.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 4, B = 0, C = -9$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$4\lambda^2 e^{\lambda x} - 9e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$4\lambda^2 - 9 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 4, B = 0, C = -9$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(4)} \pm \frac{1}{(2)(4)} \sqrt{0^2 - (4)(4)(-9)} \\ &= \pm \frac{3}{2} \end{aligned}$$

Hence

$$\lambda_1 = +\frac{3}{2}$$

$$\lambda_2 = -\frac{3}{2}$$

Which simplifies to

$$\lambda_1 = \frac{3}{2}$$

$$\lambda_2 = -\frac{3}{2}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(\frac{3}{2})x} + c_2 e^{(-\frac{3}{2})x}$$

Or

$$y = c_1 e^{\frac{3x}{2}} + c_2 e^{-\frac{3x}{2}}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{\frac{3x}{2}} + c_2 e^{-\frac{3x}{2}} \quad (1)$$

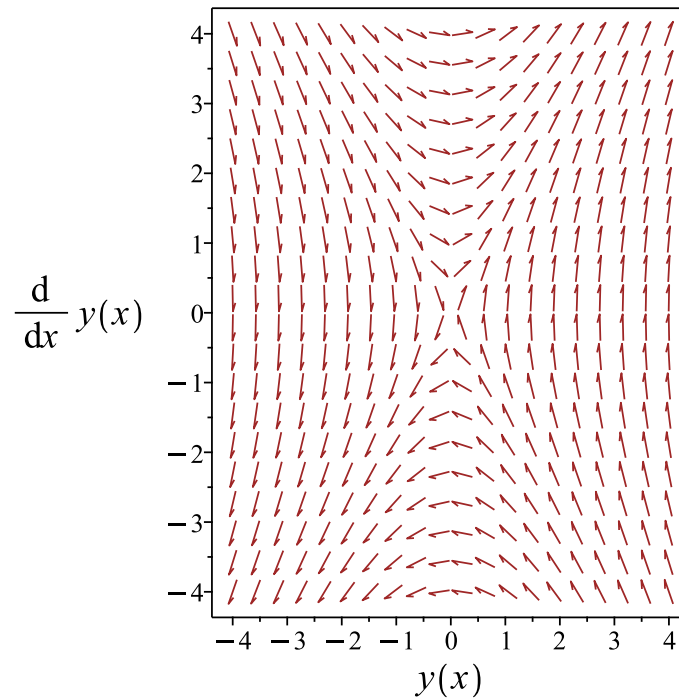


Figure 363: Slope field plot

Verification of solutions

$$y = c_1 e^{\frac{3x}{2}} + c_2 e^{-\frac{3x}{2}}$$

Verified OK.

7.6.2 Solving as second order ode can be made integrable ode

Multiplying the ode by y' gives

$$4y'y'' - 9yy' = 0$$

Integrating the above w.r.t x gives

$$\int (4y'y'' - 9yy') dx = 0$$

$$2y'^2 - \frac{9y^2}{2} = c_2$$

Which is now solved for y . Solving the given ode for y' results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$y' = \frac{\sqrt{9y^2 + 2c_1}}{2} \quad (1)$$

$$y' = -\frac{\sqrt{9y^2 + 2c_1}}{2} \quad (2)$$

Now each one of the above ODE is solved.

Solving equation (1)

Integrating both sides gives

$$\int \frac{2}{\sqrt{9y^2 + 2c_1}} dy = \int dx$$

$$\frac{2 \ln (y\sqrt{9} + \sqrt{9y^2 + 2c_1}) \sqrt{9}}{9} = x + c_2$$

Raising both side to exponential gives

$$e^{\frac{2 \ln (y\sqrt{9} + \sqrt{9y^2 + 2c_1}) \sqrt{9}}{9}} = e^{x+c_2}$$

Which simplifies to

$$\left(3y + \sqrt{9y^2 + 2c_1}\right)^{\frac{2}{3}} = c_3 e^x$$

Solving equation (2)

Integrating both sides gives

$$\int -\frac{2}{\sqrt{9y^2 + 2c_1}} dy = \int dx$$

$$-\frac{2 \ln (y\sqrt{9} + \sqrt{9y^2 + 2c_1}) \sqrt{9}}{9} = x + c_4$$

Raising both side to exponential gives

$$e^{-\frac{2 \ln (y\sqrt{9} + \sqrt{9y^2 + 2c_1}) \sqrt{9}}{9}} = e^{x+c_4}$$

Which simplifies to

$$\frac{1}{\left(3y + \sqrt{9y^2 + 2c_1}\right)^{\frac{2}{3}}} = c_5 e^x$$

Summary

The solution(s) found are the following

$$y = \frac{(c_3^3 e^{3x} - 2c_1) e^{-x}}{6c_3 \sqrt{c_3 e^x}} \quad (1)$$

$$y = -\frac{(2c_1 c_5^3 e^{3x} - 1) e^{-x}}{6c_5 \sqrt{c_5 e^x}} \quad (2)$$

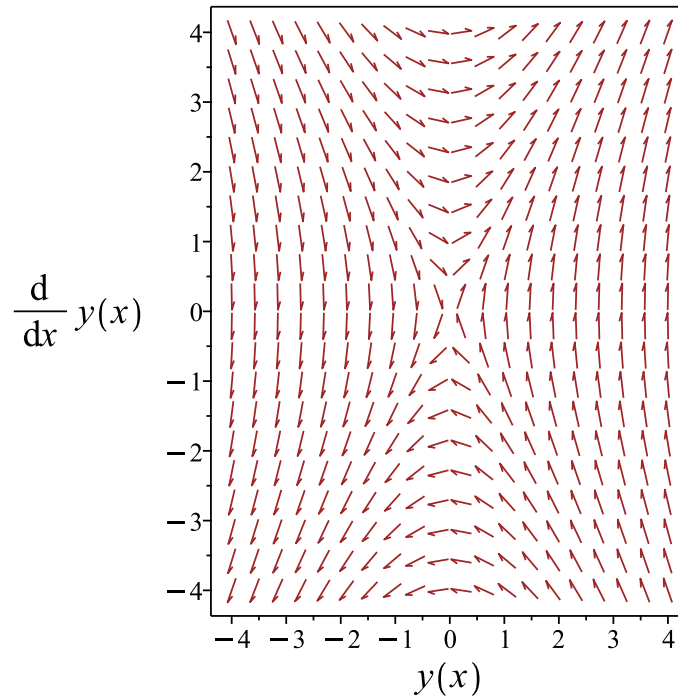


Figure 364: Slope field plot

Verification of solutions

$$y = \frac{(c_3^3 e^{3x} - 2c_1) e^{-x}}{6c_3 \sqrt{c_3 e^x}}$$

Verified OK.

$$y = -\frac{(2c_1 c_5^3 e^{3x} - 1) e^{-x}}{6c_5 \sqrt{c_5 e^x}}$$

Verified OK.

7.6.3 Solving using Kovacic algorithm

Writing the ode as

$$4y'' - 9y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 4$$

$$B = 0 \quad (3)$$

$$C = -9$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \quad (5)$$
$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{9}{4} \quad (6)$$

Comparing the above to (5) shows that

$$s = 9$$

$$t = 4$$

Therefore eq. (4) becomes

$$z''(x) = \frac{9z(x)}{4} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 326: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{9}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{3x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}y_1 &= z_1 \\ &= e^{-\frac{3x}{2}}\end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{3x}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Since $B = 0$ then the above becomes

$$\begin{aligned}y_2 &= y_1 \int \frac{1}{y_1^2} dx \\ &= e^{-\frac{3x}{2}} \int \frac{1}{e^{-3x}} dx \\ &= e^{-\frac{3x}{2}} \left(\frac{e^{3x}}{3} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-\frac{3x}{2}} \right) + c_2 \left(e^{-\frac{3x}{2}} \left(\frac{e^{3x}}{3} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{3x}{2}} + \frac{c_2 e^{\frac{3x}{2}}}{3} \quad (1)$$

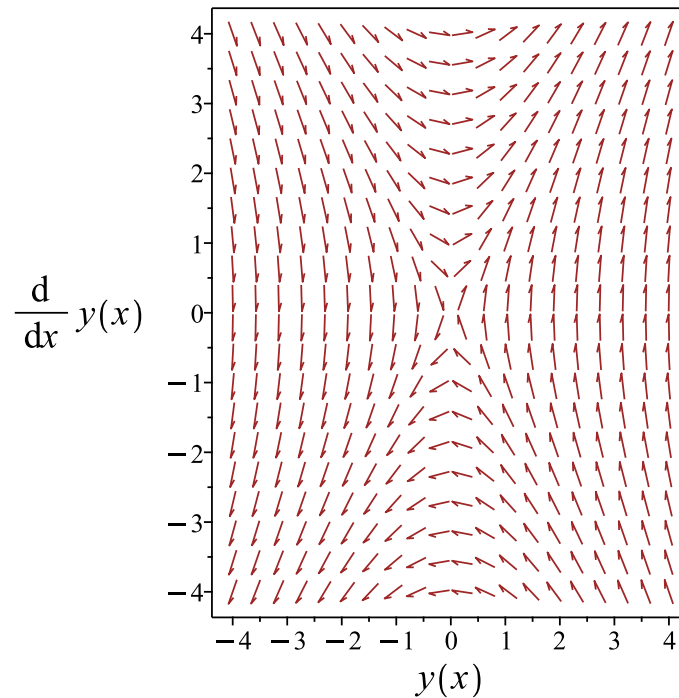


Figure 365: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{3x}{2}} + \frac{c_2 e^{\frac{3x}{2}}}{3}$$

Verified OK.

7.6.4 Maple step by step solution

Let's solve

$$4y'' - 9y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{9y}{4}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{9y}{4} = 0$$

- Characteristic polynomial of ODE

$$r^2 - \frac{9}{4} = 0$$

- Factor the characteristic polynomial

$$\frac{(2r-3)(2r+3)}{4} = 0$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{3}{2}, \frac{3}{2}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{3x}{2}}$$

- 2nd solution of the ODE

$$y_2(x) = e^{\frac{3x}{2}}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-\frac{3x}{2}} + c_2 e^{\frac{3x}{2}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(4*diff(y(x),x$2) -9*y(x) = 0,y(x), singsol=all)
```

$$y(x) = c_1 e^{-\frac{3x}{2}} + c_2 e^{\frac{3x}{2}}$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 24

```
DSolve[4*y'[x]-9*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-3x/2}(c_1 e^{3x} + c_2)$$

7.7 problem 7

7.7.1	Solving as second order linear constant coeff ode	1823
7.7.2	Solving using Kovacic algorithm	1825
7.7.3	Maple step by step solution	1829

Internal problem ID [605]

Internal file name [OUTPUT/605_Sunday_June_05_2022_01_45_43_AM_24728919/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 7.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - 9y' + 9y = 0$$

7.7.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -9, C = 9$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - 9\lambda e^{\lambda x} + 9e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - 9\lambda + 9 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -9, C = 9$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{9}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-9^2 - (4)(1)(9)} \\ &= \frac{9}{2} \pm \frac{3\sqrt{5}}{2}\end{aligned}$$

Hence

$$\lambda_1 = \frac{9}{2} + \frac{3\sqrt{5}}{2}$$

$$\lambda_2 = \frac{9}{2} - \frac{3\sqrt{5}}{2}$$

Which simplifies to

$$\lambda_1 = \frac{9}{2} + \frac{3\sqrt{5}}{2}$$

$$\lambda_2 = \frac{9}{2} - \frac{3\sqrt{5}}{2}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{\left(\frac{9}{2} + \frac{3\sqrt{5}}{2}\right)x} + c_2 e^{\left(\frac{9}{2} - \frac{3\sqrt{5}}{2}\right)x}$$

Or

$$y = c_1 e^{\left(\frac{9}{2} + \frac{3\sqrt{5}}{2}\right)x} + c_2 e^{\left(\frac{9}{2} - \frac{3\sqrt{5}}{2}\right)x}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{\left(\frac{9}{2} + \frac{3\sqrt{5}}{2}\right)x} + c_2 e^{\left(\frac{9}{2} - \frac{3\sqrt{5}}{2}\right)x} \quad (1)$$

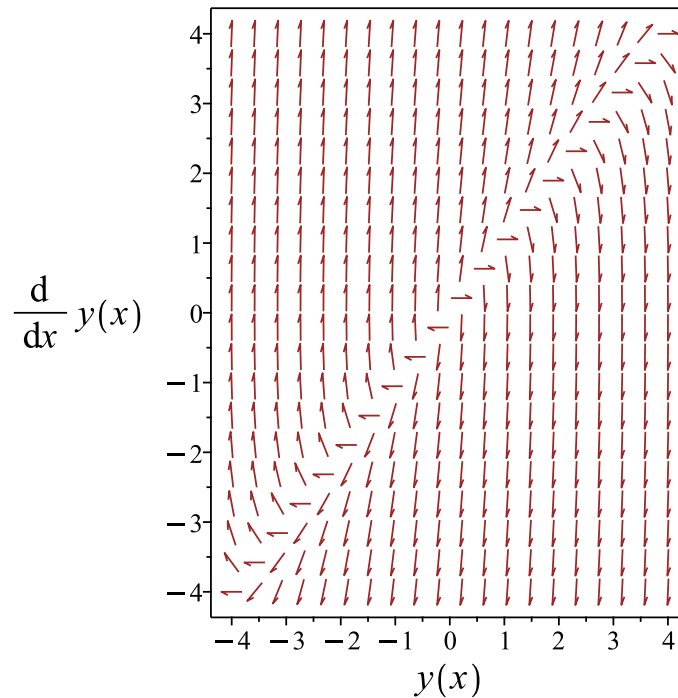


Figure 366: Slope field plot

Verification of solutions

$$y = c_1 e^{\left(\frac{9}{2} + \frac{3\sqrt{5}}{2}\right)x} + c_2 e^{\left(\frac{9}{2} - \frac{3\sqrt{5}}{2}\right)x}$$

Verified OK.

7.7.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 9y' + 9y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = -9 \tag{3}$$

$$C = 9$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{45}{4} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 45 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{45z(x)}{4} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 328: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{45}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{3x\sqrt{5}}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-9}{1} dx} \\ &= z_1 e^{\frac{9x}{2}} \\ &= z_1 \left(e^{\frac{9x}{2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{3(\sqrt{5}-3)x}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-9}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{9x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{\sqrt{5} e^{3x\sqrt{5}}}{15} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-\frac{3(\sqrt{5}-3)x}{2}} \right) + c_2 \left(e^{-\frac{3(\sqrt{5}-3)x}{2}} \left(\frac{\sqrt{5} e^{3x\sqrt{5}}}{15} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{3(\sqrt{5}-3)x}{2}} + \frac{c_2 \sqrt{5} e^{\frac{3(3+\sqrt{5})x}{2}}}{15} \quad (1)$$

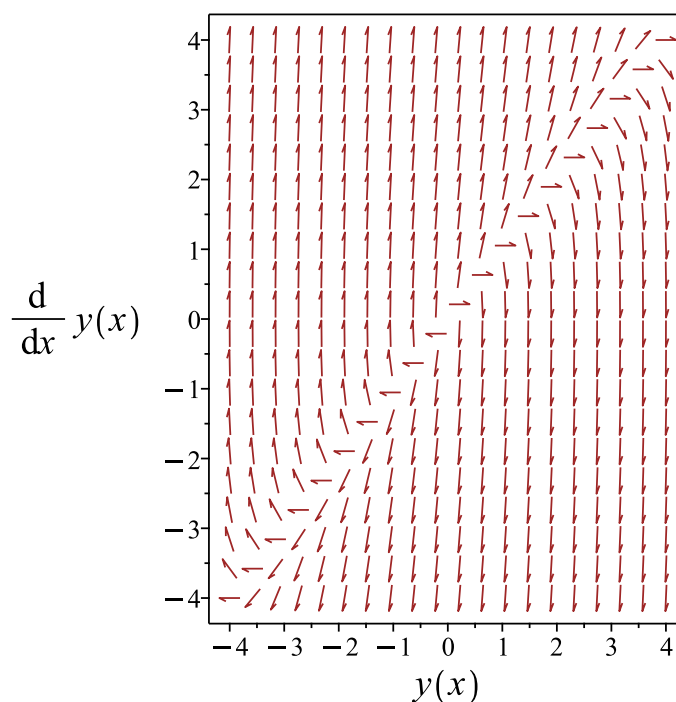


Figure 367: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{3(\sqrt{5}-3)x}{2}} + \frac{c_2 \sqrt{5} e^{\frac{3(3+\sqrt{5})x}{2}}}{15}$$

Verified OK.

7.7.3 Maple step by step solution

Let's solve

$$y'' - 9y' + 9y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 - 9r + 9 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{9 \pm (\sqrt{45})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(\frac{9}{2} - \frac{3\sqrt{5}}{2}, \frac{9}{2} + \frac{3\sqrt{5}}{2} \right)$$

- 1st solution of the ODE

$$y_1(x) = e^{\left(\frac{9}{2} - \frac{3\sqrt{5}}{2}\right)x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{\left(\frac{9}{2} + \frac{3\sqrt{5}}{2}\right)x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{\left(\frac{9}{2} - \frac{3\sqrt{5}}{2}\right)x} + c_2 e^{\left(\frac{9}{2} + \frac{3\sqrt{5}}{2}\right)x}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 27

```
dsolve(diff(y(x),x$2) -9*diff(y(x),x)+9*y(x) = 0,y(x), singsol=all)
```

$$y(x) = c_1 e^{\frac{3(3+\sqrt{5})x}{2}} + c_2 e^{-\frac{3(\sqrt{5}-3)x}{2}}$$

✓ Solution by Mathematica

Time used: 0.022 (sec). Leaf size: 36

```
DSolve[y''[x]-9*y'[x]+9*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-\frac{3}{2}(\sqrt{5}-3)x} \left(c_2 e^{3\sqrt{5}x} + c_1 \right)$$

7.8 problem 8

7.8.1 Solving as second order linear constant coeff ode	1831
7.8.2 Solving using Kovacic algorithm	1833
7.8.3 Maple step by step solution	1837

Internal problem ID [606]

Internal file name [OUTPUT/606_Sunday_June_05_2022_01_45_44_AM_43508000/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 8.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - 2y' - 2y = 0$$

7.8.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -2, C = -2$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - 2\lambda e^{\lambda x} - 2e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - 2\lambda - 2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -2, C = -2$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-2^2 - (4)(1)(-2)} \\ &= 1 \pm \sqrt{3}\end{aligned}$$

Hence

$$\lambda_1 = 1 + \sqrt{3}$$

$$\lambda_2 = 1 - \sqrt{3}$$

Which simplifies to

$$\lambda_1 = 1 + \sqrt{3}$$

$$\lambda_2 = -\sqrt{3} + 1$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(1+\sqrt{3})x} + c_2 e^{(-\sqrt{3}+1)x}$$

Or

$$y = c_1 e^{(1+\sqrt{3})x} + c_2 e^{(-\sqrt{3}+1)x}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{(1+\sqrt{3})x} + c_2 e^{(-\sqrt{3}+1)x} \quad (1)$$

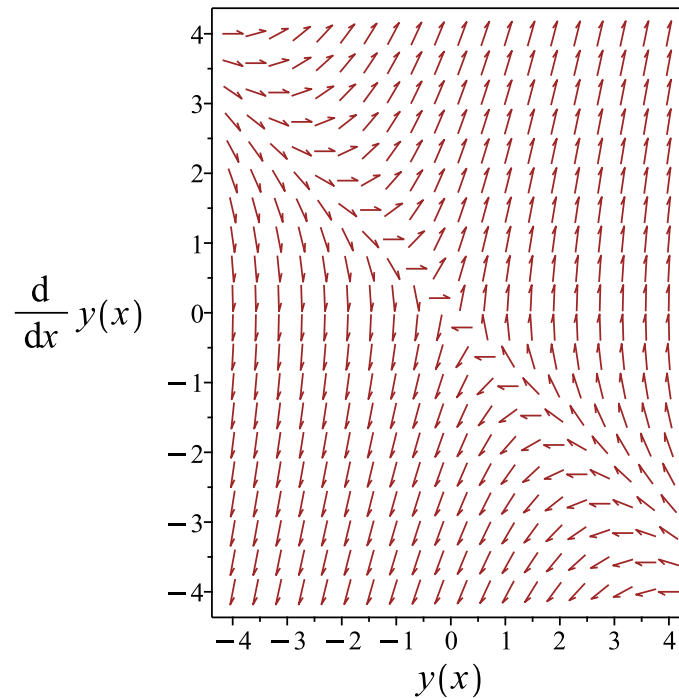


Figure 368: Slope field plot

Verification of solutions

$$y = c_1 e^{(1+\sqrt{3})x} + c_2 e^{(-\sqrt{3}+1)x}$$

Verified OK.

7.8.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 2y' - 2y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= -2 \\ C &= -2 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{3}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 3$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = 3z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 330: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 3$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\sqrt{3}x}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-2}{1} dx} \\ &= z_1 e^x \\ &= z_1(e^x) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-(\sqrt{3}-1)x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{2x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{\sqrt{3} e^{2\sqrt{3}x}}{6} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-(\sqrt{3}-1)x} \right) + c_2 \left(e^{-(\sqrt{3}-1)x} \left(\frac{\sqrt{3} e^{2\sqrt{3}x}}{6} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-(\sqrt{3}-1)x} + \frac{c_2 \sqrt{3} e^{(1+\sqrt{3})x}}{6} \quad (1)$$

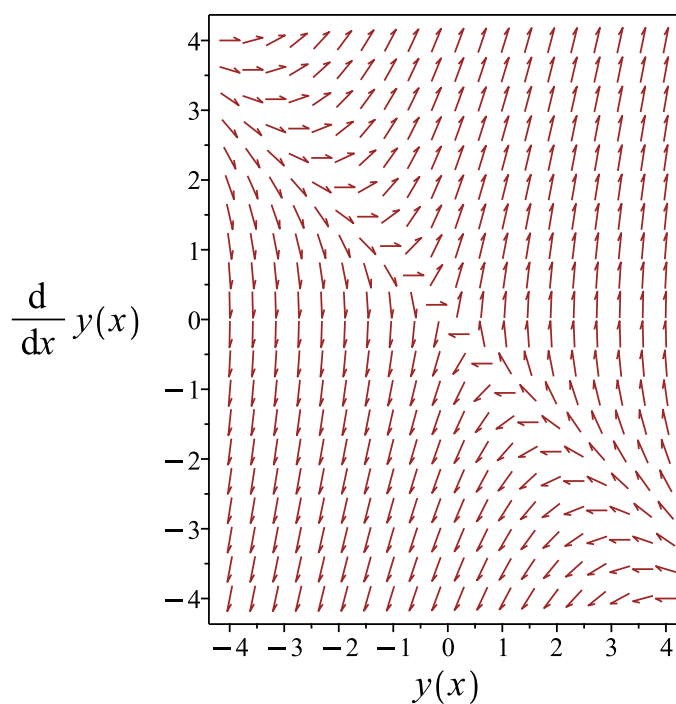


Figure 369: Slope field plot

Verification of solutions

$$y = c_1 e^{-(\sqrt{3}-1)x} + \frac{c_2 \sqrt{3} e^{(1+\sqrt{3})x}}{6}$$

Verified OK.

7.8.3 Maple step by step solution

Let's solve

$$y'' - 2y' - 2y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 - 2r - 2 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{2 \pm (\sqrt{12})}{2}$$

- Roots of the characteristic polynomial

$$r = (1 + \sqrt{3}, -\sqrt{3} + 1)$$

- 1st solution of the ODE

$$y_1(x) = e^{(1+\sqrt{3})x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{(-\sqrt{3}+1)x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{(1+\sqrt{3})x} + c_2 e^{(-\sqrt{3}+1)x}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 26

```
dsolve(diff(y(x),x$2) -2*diff(y(x),x)-2*y(x) = 0,y(x), singsol=all)
```

$$y(x) = c_1 e^{(1+\sqrt{3})x} + c_2 e^{-(\sqrt{3}-1)x}$$

✓ Solution by Mathematica

Time used: 0.019 (sec). Leaf size: 34

```
DSolve[y''[x]-2*y'[x]-2*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{x-\sqrt{3}x} \left(c_2 e^{2\sqrt{3}x} + c_1 \right)$$

7.9 problem 9

7.9.1	Existence and uniqueness analysis	1839
7.9.2	Solving as second order linear constant coeff ode	1840
7.9.3	Solving using Kovacic algorithm	1842
7.9.4	Maple step by step solution	1847

Internal problem ID [607]

Internal file name [OUTPUT/607_Sunday_June_05_2022_01_45_44_AM_60102257/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 9.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + y' - 2y = 0$$

With initial conditions

$$[y(0) = 0, y'(0) = 1]$$

7.9.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 1$$

$$q(x) = -2$$

$$F = 0$$

Hence the ode is

$$y'' + y' - 2y = 0$$

The domain of $p(x) = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = -2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.9.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 1, C = -2$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + \lambda e^{\lambda x} - 2 e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + \lambda - 2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 1, C = -2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{1^2 - (4)(1)(-2)} \\ &= -\frac{1}{2} \pm \frac{3}{2} \end{aligned}$$

Hence

$$\lambda_1 = -\frac{1}{2} + \frac{3}{2}$$

$$\lambda_2 = -\frac{1}{2} - \frac{3}{2}$$

Which simplifies to

$$\lambda_1 = 1$$

$$\lambda_2 = -2$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(1)x} + c_2 e^{(-2)x}$$

Or

$$y = c_1 e^x + c_2 e^{-2x}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^x + c_2 e^{-2x} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $x = 0$ in the above gives

$$0 = c_1 + c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = c_1 e^x - 2c_2 e^{-2x}$$

substituting $y' = 1$ and $x = 0$ in the above gives

$$1 = c_1 - 2c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{1}{3}$$

$$c_2 = -\frac{1}{3}$$

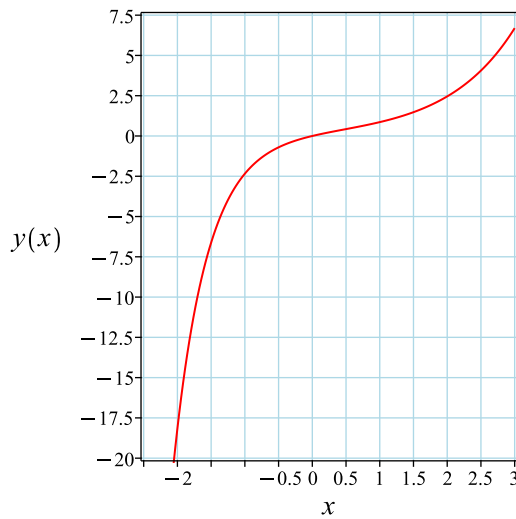
Substituting these values back in above solution results in

$$y = \frac{e^x}{3} - \frac{e^{-2x}}{3}$$

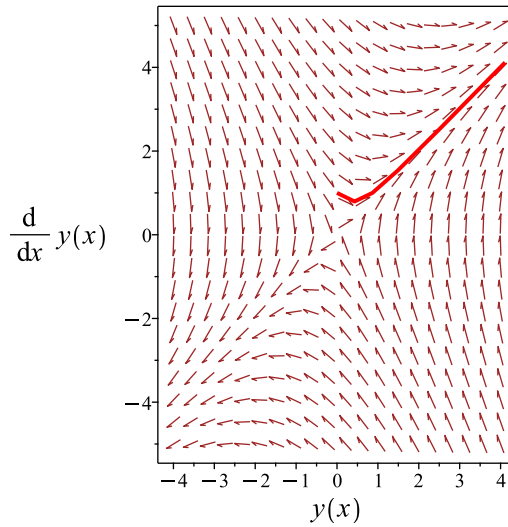
Summary

The solution(s) found are the following

$$y = \frac{e^x}{3} - \frac{e^{-2x}}{3} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^x}{3} - \frac{e^{-2x}}{3}$$

Verified OK.

7.9.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + y' - 2y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 1$$

$$C = -2$$

(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{9}{4} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 9 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{9z(x)}{4} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 332: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{9}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{3x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}
 y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\
 &= z_1 e^{-\int \frac{1}{2} \frac{1}{1} dx} \\
 &= z_1 e^{-\frac{x}{2}} \\
 &= z_1 (e^{-\frac{x}{2}})
 \end{aligned}$$

Which simplifies to

$$y_1 = e^{-2x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{1}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{e^{3x}}{3} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-2x}) + c_2 \left(e^{-2x} \left(\frac{e^{3x}}{3} \right) \right) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-2x} + \frac{c_2 e^x}{3} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $x = 0$ in the above gives

$$0 = c_1 + \frac{c_2}{3} \tag{1A}$$

Taking derivative of the solution gives

$$y' = -2c_1 e^{-2x} + \frac{c_2 e^x}{3}$$

substituting $y' = 1$ and $x = 0$ in the above gives

$$1 = -2c_1 + \frac{c_2}{3} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -\frac{1}{3}$$
$$c_2 = 1$$

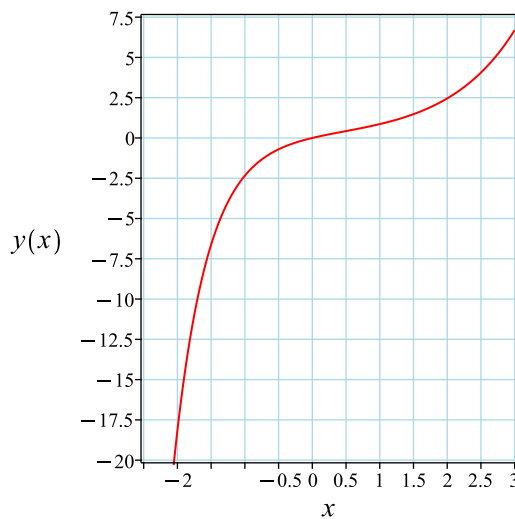
Substituting these values back in above solution results in

$$y = \frac{e^x}{3} - \frac{e^{-2x}}{3}$$

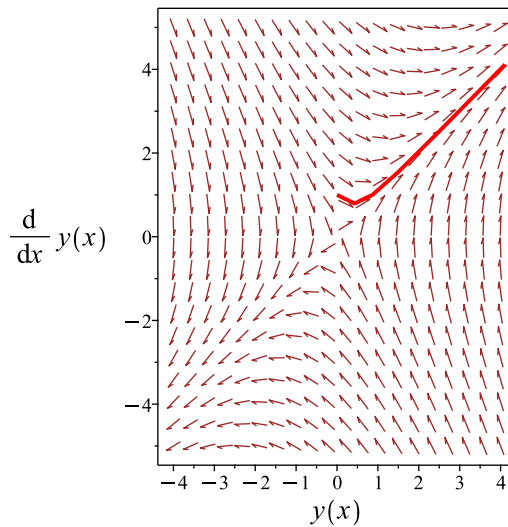
Summary

The solution(s) found are the following

$$y = \frac{e^x}{3} - \frac{e^{-2x}}{3} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^x}{3} - \frac{e^{-2x}}{3}$$

Verified OK.

7.9.4 Maple step by step solution

Let's solve

$$\left[y'' + y' - 2y = 0, y(0) = 0, y' \Big|_{\{x=0\}} = 1 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + r - 2 = 0$$

- Factor the characteristic polynomial

$$(r + 2)(r - 1) = 0$$

- Roots of the characteristic polynomial

$$r = (-2, 1)$$

- 1st solution of the ODE

$$y_1(x) = e^{-2x}$$

- 2nd solution of the ODE

$$y_2(x) = e^x$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-2x} + c_2 e^x$$

- Check validity of solution $y = c_1 e^{-2x} + c_2 e^x$

- Use initial condition $y(0) = 0$

$$0 = c_1 + c_2$$

- Compute derivative of the solution

$$y' = -2c_1 e^{-2x} + c_2 e^x$$

- Use the initial condition $y' \Big|_{\{x=0\}} = 1$

$$1 = -2c_1 + c_2$$

- Solve for c_1 and c_2

$$\left\{ c_1 = -\frac{1}{3}, c_2 = \frac{1}{3} \right\}$$

- Substitute constant values into general solution and simplify

$$y = \frac{(e^{3x}-1)e^{-2x}}{3}$$

- Solution to the IVP

$$y = \frac{(e^{3x}-1)e^{-2x}}{3}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 16

```
dsolve([diff(y(x),x$2) +diff(y(x),x)-2*y(x) = 0,y(0) = 0, D(y)(0) = 1],y(x), singsol=all)
```

$$y(x) = \frac{(e^{3x} - 1)e^{-2x}}{3}$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 21

```
DSolve[{y''[x]+y'[x]-2*y[x]==0,{y[0]==0,y'[0]==1}},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{3}e^{-2x}(e^{3x} - 1)$$

7.10 problem 10

7.10.1 Existence and uniqueness analysis	1849
7.10.2 Solving as second order linear constant coeff ode	1850
7.10.3 Solving using Kovacic algorithm	1852
7.10.4 Maple step by step solution	1857

Internal problem ID [608]

Internal file name [OUTPUT/608_Sunday_June_05_2022_01_45_46_AM_30342769/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 10.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$y'' + 4y' + 3y = 0$$

With initial conditions

$$[y(0) = 2, y'(0) = -1]$$

7.10.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 4$$

$$q(x) = 3$$

$$F = 0$$

Hence the ode is

$$y'' + 4y' + 3y = 0$$

The domain of $p(x) = 4$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = 3$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.10.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 4, C = 3$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 4\lambda e^{\lambda x} + 3e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 4\lambda + 3 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 4, C = 3$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-4}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{4^2 - (4)(1)(3)} \\ &= -2 \pm 1 \end{aligned}$$

Hence

$$\lambda_1 = -2 + 1$$

$$\lambda_2 = -2 - 1$$

Which simplifies to

$$\lambda_1 = -1$$

$$\lambda_2 = -3$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(-1)x} + c_2 e^{(-3)x}$$

Or

$$y = c_1 e^{-x} + e^{-3x} c_2$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-x} + e^{-3x} c_2 \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = c_1 + c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -c_1 e^{-x} - 3 e^{-3x} c_2$$

substituting $y' = -1$ and $x = 0$ in the above gives

$$-1 = -c_1 - 3c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{5}{2}$$

$$c_2 = -\frac{1}{2}$$

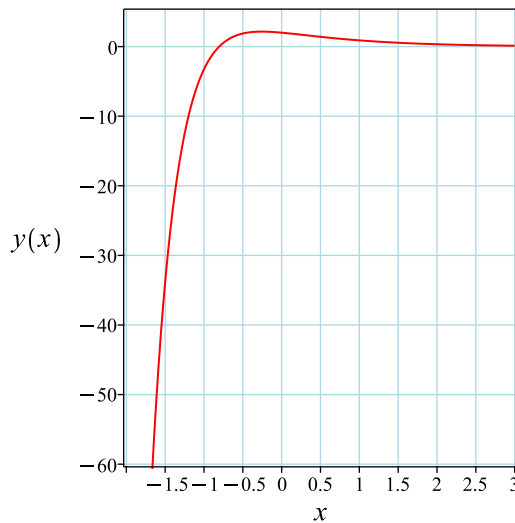
Substituting these values back in above solution results in

$$y = \frac{5 e^{-x}}{2} - \frac{e^{-3x}}{2}$$

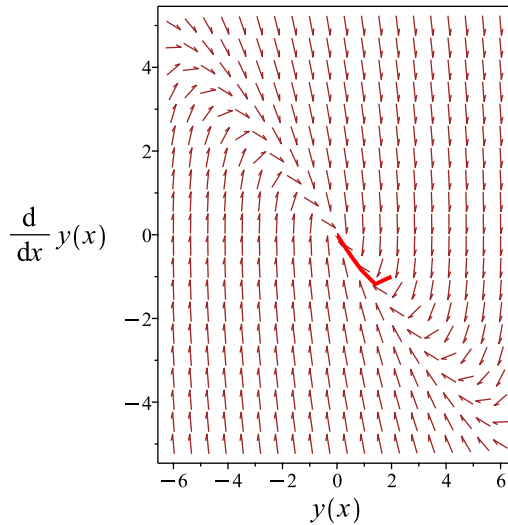
Summary

The solution(s) found are the following

$$y = \frac{5e^{-x}}{2} - \frac{e^{-3x}}{2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{5e^{-x}}{2} - \frac{e^{-3x}}{2}$$

Verified OK.

7.10.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 4y' + 3y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 4$$

$$C = 3$$

(3)

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 334: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-x}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}
 y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\
 &= z_1 e^{-\int \frac{1}{2} \frac{4}{1} dx} \\
 &= z_1 e^{-2x} \\
 &= z_1 (e^{-2x})
 \end{aligned}$$

Which simplifies to

$$y_1 = e^{-3x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{4}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-4x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{e^{2x}}{2} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-3x}) + c_2 \left(e^{-3x} \left(\frac{e^{2x}}{2} \right) \right) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-3x} + \frac{c_2 e^{-x}}{2} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = c_1 + \frac{c_2}{2} \tag{1A}$$

Taking derivative of the solution gives

$$y' = -3c_1 e^{-3x} - \frac{c_2 e^{-x}}{2}$$

substituting $y' = -1$ and $x = 0$ in the above gives

$$-1 = -3c_1 - \frac{c_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -\frac{1}{2}$$
$$c_2 = 5$$

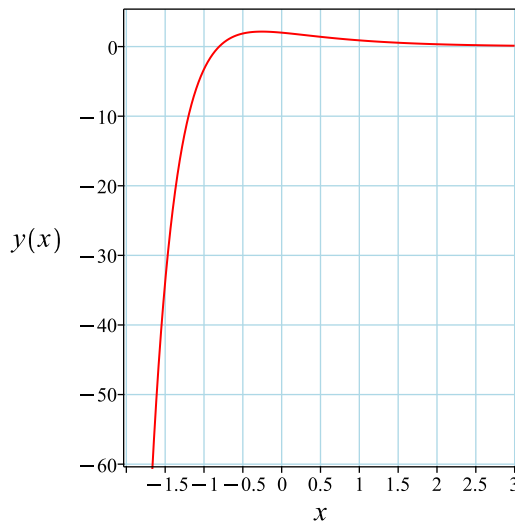
Substituting these values back in above solution results in

$$y = \frac{5e^{-x}}{2} - \frac{e^{-3x}}{2}$$

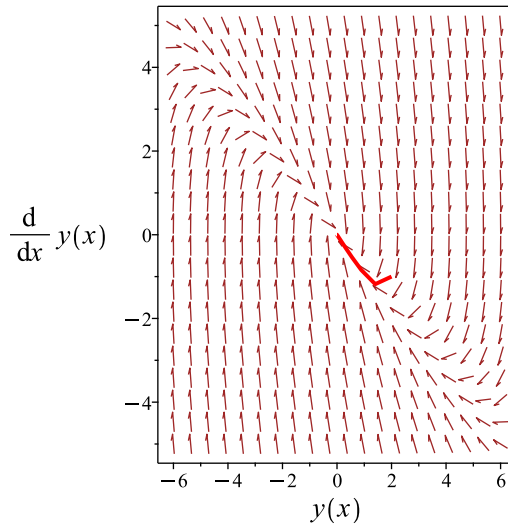
Summary

The solution(s) found are the following

$$y = \frac{5e^{-x}}{2} - \frac{e^{-3x}}{2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{5e^{-x}}{2} - \frac{e^{-3x}}{2}$$

Verified OK.

7.10.4 Maple step by step solution

Let's solve

$$\left[y'' + 4y' + 3y = 0, y(0) = 2, y' \Big|_{\{x=0\}} = -1 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 4r + 3 = 0$$

- Factor the characteristic polynomial

$$(r + 3)(r + 1) = 0$$

- Roots of the characteristic polynomial

$$r = (-3, -1)$$

- 1st solution of the ODE

$$y_1(x) = e^{-3x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{-x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-3x} + c_2 e^{-x}$$

- Check validity of solution $y = c_1 e^{-3x} + c_2 e^{-x}$

- Use initial condition $y(0) = 2$

$$2 = c_1 + c_2$$

- Compute derivative of the solution

$$y' = -3c_1 e^{-3x} - c_2 e^{-x}$$

- Use the initial condition $y' \Big|_{\{x=0\}} = -1$

$$-1 = -3c_1 - c_2$$

- Solve for c_1 and c_2

$$\left\{ c_1 = -\frac{1}{2}, c_2 = \frac{5}{2} \right\}$$

- Substitute constant values into general solution and simplify

$$y = \frac{5e^{-x}}{2} - \frac{e^{-3x}}{2}$$

- Solution to the IVP

$$y = \frac{5e^{-x}}{2} - \frac{e^{-3x}}{2}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 17

```
dsolve([diff(y(x),x$2) +4*diff(y(x),x)+3*y(x) = 0,y(0) = 2, D(y)(0) = -1],y(x), singsol=all)
```

$$y(x) = -\frac{e^{-3x}}{2} + \frac{5e^{-x}}{2}$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 23

```
DSolve[{y''[x]+4*y'[x]+3*y[x]==0,{y[0]==2,y'[0]==-1}},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2}e^{-3x}(5e^{2x} - 1)$$

7.11 problem 11

7.11.1 Existence and uniqueness analysis	1859
7.11.2 Solving as second order linear constant coeff ode	1860
7.11.3 Solving using Kovacic algorithm	1862
7.11.4 Maple step by step solution	1866

Internal problem ID [609]

Internal file name [OUTPUT/609_Sunday_June_05_2022_01_45_47_AM_47206086/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 11.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$6y'' - 5y' + y = 0$$

With initial conditions

$$[y(0) = 4, y'(0) = 0]$$

7.11.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = -\frac{5}{6}$$
$$q(x) = \frac{1}{6}$$
$$F = 0$$

Hence the ode is

$$y'' - \frac{5y'}{6} + \frac{y}{6} = 0$$

The domain of $p(x) = -\frac{5}{6}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = \frac{1}{6}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.11.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 6, B = -5, C = 1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$6\lambda^2 e^{\lambda x} - 5\lambda e^{\lambda x} + e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$6\lambda^2 - 5\lambda + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 6, B = -5, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{5}{(2)(6)} \pm \frac{1}{(2)(6)} \sqrt{-5^2 - (4)(6)(1)} \\ &= \frac{5}{12} \pm \frac{1}{12} \end{aligned}$$

Hence

$$\lambda_1 = \frac{5}{12} + \frac{1}{12}$$

$$\lambda_2 = \frac{5}{12} - \frac{1}{12}$$

Which simplifies to

$$\lambda_1 = \frac{1}{2}$$

$$\lambda_2 = \frac{1}{3}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(\frac{1}{2})x} + c_2 e^{(\frac{1}{3})x}$$

Or

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{\frac{x}{3}}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{\frac{x}{3}} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 4$ and $x = 0$ in the above gives

$$4 = c_1 + c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = \frac{c_1 e^{\frac{x}{2}}}{2} + \frac{c_2 e^{\frac{x}{3}}}{3}$$

substituting $y' = 0$ and $x = 0$ in the above gives

$$0 = \frac{c_1}{2} + \frac{c_2}{3} \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -8$$

$$c_2 = 12$$

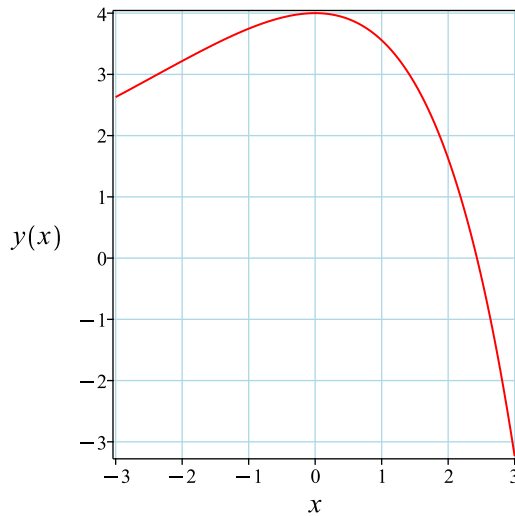
Substituting these values back in above solution results in

$$y = -8 e^{\frac{x}{2}} + 12 e^{\frac{x}{3}}$$

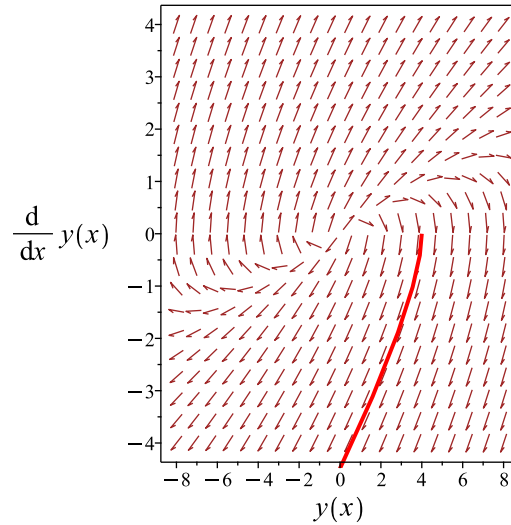
Summary

The solution(s) found are the following

$$y = -8e^{\frac{x}{2}} + 12e^{\frac{x}{3}} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -8e^{\frac{x}{2}} + 12e^{\frac{x}{3}}$$

Verified OK.

7.11.3 Solving using Kovacic algorithm

Writing the ode as

$$6y'' - 5y' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 6$$

$$B = -5 \quad (3)$$

$$C = 1$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{144} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 144 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{z(x)}{144} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 336: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{144}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x}{12}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-5}{6} dx} \\ &= z_1 e^{\frac{5x}{12}} \\ &= z_1 \left(e^{\frac{5x}{12}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{\frac{x}{3}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-5}{6} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{\frac{5x}{6}}}{(y_1)^2} dx \\ &= y_1 (6 e^{\frac{x}{6}}) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{\frac{x}{3}}) + c_2 (e^{\frac{x}{3}} (6 e^{\frac{x}{6}}))\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\frac{x}{3}} + 6c_2 e^{\frac{x}{2}} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 4$ and $x = 0$ in the above gives

$$4 = c_1 + 6c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{c_1 e^{\frac{x}{3}}}{3} + 3c_2 e^{\frac{x}{2}}$$

substituting $y' = 0$ and $x = 0$ in the above gives

$$0 = \frac{c_1}{3} + 3c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}c_1 &= 12 \\ c_2 &= -\frac{4}{3}\end{aligned}$$

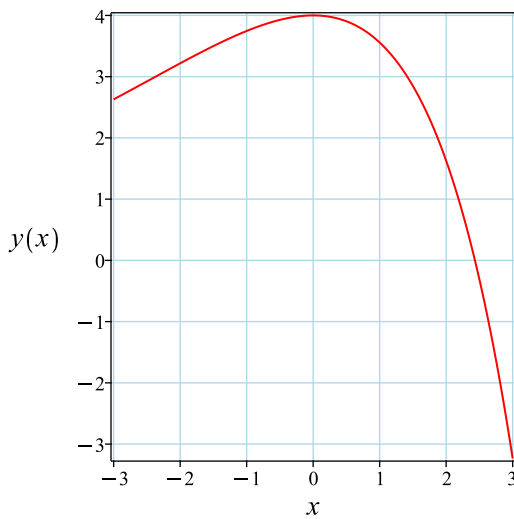
Substituting these values back in above solution results in

$$y = -8 e^{\frac{x}{2}} + 12 e^{\frac{x}{3}}$$

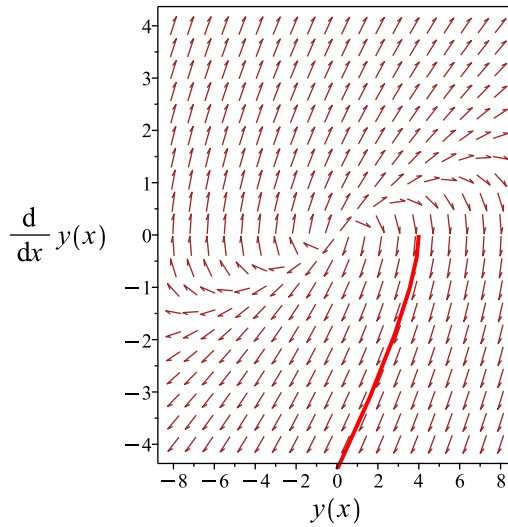
Summary

The solution(s) found are the following

$$y = -8 e^{\frac{x}{2}} + 12 e^{\frac{x}{3}} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -8e^{\frac{x}{2}} + 12e^{\frac{x}{3}}$$

Verified OK.

7.11.4 Maple step by step solution

Let's solve

$$\left[6y'' - 5y' + y = 0, y(0) = 4, y'|_{\{x=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{5y'}{6} - \frac{y}{6}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{5y'}{6} + \frac{y}{6} = 0$$

- Characteristic polynomial of ODE

$$r^2 - \frac{5}{6}r + \frac{1}{6} = 0$$

- Factor the characteristic polynomial

$$\frac{(3r-1)(2r-1)}{6} = 0$$

- Roots of the characteristic polynomial

$$r = \left(\frac{1}{2}, \frac{1}{3}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{\frac{x}{2}}$$

- 2nd solution of the ODE

$$y_2(x) = e^{\frac{x}{3}}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{\frac{x}{3}}$$

- Check validity of solution $y = c_1 e^{\frac{x}{2}} + c_2 e^{\frac{x}{3}}$

- Use initial condition $y(0) = 4$

$$4 = c_1 + c_2$$

- Compute derivative of the solution

$$y' = \frac{c_1 e^{\frac{x}{2}}}{2} + \frac{c_2 e^{\frac{x}{3}}}{3}$$

- Use the initial condition $y' \Big|_{\{x=0\}} = 0$

$$0 = \frac{c_1}{2} + \frac{c_2}{3}$$

- Solve for c_1 and c_2

$$\{c_1 = -8, c_2 = 12\}$$

- Substitute constant values into general solution and simplify

$$y = -8 e^{\frac{x}{2}} + 12 e^{\frac{x}{3}}$$

- Solution to the IVP

$$y = -8 e^{\frac{x}{2}} + 12 e^{\frac{x}{3}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 17

```
dsolve([6*diff(y(x),x$2) -5*diff(y(x),x)+y(x) = 0,y(0) = 4, D(y)(0) = 0],y(x), singsol=all)
```

$$y(x) = -8e^{\frac{x}{2}} + 12e^{\frac{x}{3}}$$

✓ Solution by Mathematica

Time used: 0.025 (sec). Leaf size: 48

```
DSolve[{6*y'[x]-5*y'[x]+2*y[x]==0,{y[0]==4,y'[0]==0}},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{4}{23} e^{5x/12} \left(23 \cos\left(\frac{\sqrt{23}x}{12}\right) - 5\sqrt{23} \sin\left(\frac{\sqrt{23}x}{12}\right) \right)$$

7.12 problem 12

7.12.1 Existence and uniqueness analysis	1870
7.12.2 Solving as second order linear constant coeff ode	1870
7.12.3 Solving as second order integrable as is ode	1872
7.12.4 Solving as second order ode missing y ode	1873
7.12.5 Solving as type second_order_integrable_as_is (not using ABC version)	1875
7.12.6 Solving using Kovacic algorithm	1877
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7.12.8 Maple step by step solution	1883

Internal problem ID [610]

Internal file name [OUTPUT/610_Sunday_June_05_2022_01_45_47_AM_93318101/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 12.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "exact linear second order ode", "second_order_integrable_as_is", "second_order_ode_missing_y", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$y'' + 3y' = 0$$

With initial conditions

$$[y(0) = -2, y'(0) = 3]$$

7.12.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 3$$

$$q(x) = 0$$

$$F = 0$$

Hence the ode is

$$y'' + 3y' = 0$$

The domain of $p(x) = 3$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. Hence solution exists and is unique.

7.12.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 3, C = 0$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 3\lambda e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 3\lambda = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 3, C = 0$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-3}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{3^2 - (4)(1)(0)} \\ &= -\frac{3}{2} \pm \frac{3}{2} \end{aligned}$$

Hence

$$\lambda_1 = -\frac{3}{2} + \frac{3}{2}$$

$$\lambda_2 = -\frac{3}{2} - \frac{3}{2}$$

Which simplifies to

$$\lambda_1 = 0$$

$$\lambda_2 = -3$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(0)x} + c_2 e^{(-3)x}$$

Or

$$y = c_1 + e^{-3x} c_2$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 + e^{-3x} c_2 \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = -2$ and $x = 0$ in the above gives

$$-2 = c_1 + c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -3e^{-3x} c_2$$

substituting $y' = 3$ and $x = 0$ in the above gives

$$3 = -3c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -1$$

$$c_2 = -1$$

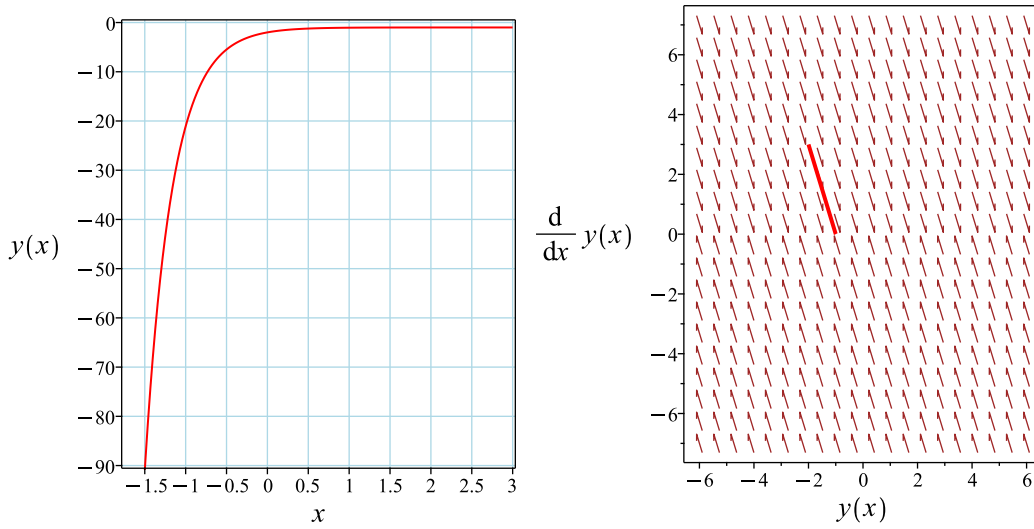
Substituting these values back in above solution results in

$$y = -1 - e^{-3x}$$

Summary

The solution(s) found are the following

$$y = -1 - e^{-3x} \tag{1}$$



(a) Solution plot

(b) Slope field plot

Verification of solutions

$$y = -1 - e^{-3x}$$

Verified OK.

7.12.3 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t x gives

$$\int (y'' + 3y') dx = 0$$

$$3y + y' = c_1$$

Which is now solved for y . Integrating both sides gives

$$\int \frac{1}{-3y + c_1} dy = \int dx$$

$$-\frac{\ln(-3y + c_1)}{3} = x + c_2$$

Raising both side to exponential gives

$$\frac{1}{(-3y + c_1)^{\frac{1}{3}}} = e^{x+c_2}$$

Which simplifies to

$$\frac{1}{(-3y + c_1)^{\frac{1}{3}}} = c_3 e^x$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = -\frac{e^{-3x}}{3c_3^3} + \frac{c_1}{3} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = -2$ and $x = 0$ in the above gives

$$-2 = \frac{c_1 c_3^3 - 1}{3c_3^3} \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{e^{-3x}}{c_3^3}$$

substituting $y' = 3$ and $x = 0$ in the above gives

$$3 = \frac{1}{c_3^3} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_3\}$. Warning, unable to solve for constants of integrations.

Verification of solutions N/A

7.12.4 Solving as second order ode missing y ode

This is second order ode with missing dependent variable y . Let

$$p(x) = y'$$

Then

$$p'(x) = y''$$

Hence the ode becomes

$$p'(x) + 3p(x) = 0$$

Which is now solve for $p(x)$ as first order ode. Integrating both sides gives

$$\int -\frac{1}{3p} dp = \int dx$$
$$-\frac{\ln(p)}{3} = x + c_1$$

Raising both side to exponential gives

$$\frac{1}{p^{\frac{1}{3}}} = e^{x+c_1}$$

Which simplifies to

$$\frac{1}{p^{\frac{1}{3}}} = c_2 e^x$$

Initial conditions are used to solve for c_2 . Substituting $x = 0$ and $p = 3$ in the above solution gives an equation to solve for the constant of integration.

$$3 = \frac{1}{c_2^3}$$

$$c_2 = \frac{3^{\frac{2}{3}}}{3}$$

Substituting c_2 found above in the general solution gives

$$p(x) = 3e^{-3x}$$

Since $p = y'$ then the new first order ode to solve is

$$y' = 3e^{-3x}$$

Integrating both sides gives

$$y = \int 3e^{-3x} dx$$
$$= -e^{-3x} + c_3$$

Initial conditions are used to solve for c_3 . Substituting $x = 0$ and $y = -2$ in the above solution gives an equation to solve for the constant of integration.

$$-2 = -1 + c_3$$

$$c_3 = -1$$

Substituting c_3 found above in the general solution gives

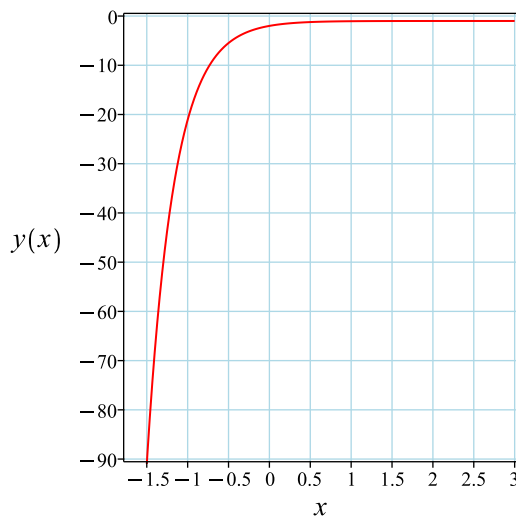
$$y = -1 - e^{-3x}$$

Initial conditions are used to solve for the constants of integration.

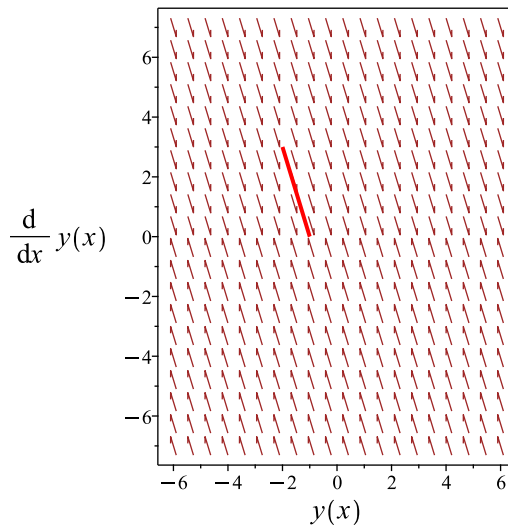
Summary

The solution(s) found are the following

$$y = -1 - e^{-3x} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -1 - e^{-3x}$$

Verified OK.

7.12.5 Solving as type second_order_integrable_as_is (not using ABC version)

Writing the ode as

$$y'' + 3y' = 0$$

Integrating both sides of the ODE w.r.t x gives

$$\int (y'' + 3y') dx = 0$$
$$3y + y' = c_1$$

Which is now solved for y . Integrating both sides gives

$$\int \frac{1}{-3y + c_1} dy = \int dx$$
$$-\frac{\ln(-3y + c_1)}{3} = x + c_2$$

Raising both side to exponential gives

$$\frac{1}{(-3y + c_1)^{\frac{1}{3}}} = e^{x+c_2}$$

Which simplifies to

$$\frac{1}{(-3y + c_1)^{\frac{1}{3}}} = c_3 e^x$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = -\frac{e^{-3x}}{3c_3^3} + \frac{c_1}{3} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = -2$ and $x = 0$ in the above gives

$$-2 = \frac{c_1 c_3^3 - 1}{3c_3^3} \tag{1A}$$

Taking derivative of the solution gives

$$y' = \frac{e^{-3x}}{c_3^3}$$

substituting $y' = 3$ and $x = 0$ in the above gives

$$3 = \frac{1}{c_3^3} \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_3\}$. Warning, unable to solve for constants of integrations.

7.12.6 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 3y' = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 3 \quad (3)$$

$$C = 0$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \quad (5)$$
$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{9}{4} \quad (6)$$

Comparing the above to (5) shows that

$$s = 9$$

$$t = 4$$

Therefore eq. (4) becomes

$$z''(x) = \frac{9z(x)}{4} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 338: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{9}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{3x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{3}{1} dx} \end{aligned}$$

$$\begin{aligned}
&= z_1 e^{-\frac{3x}{2}} \\
&= z_1 \left(e^{-\frac{3x}{2}} \right)
\end{aligned}$$

Which simplifies to

$$y_1 = e^{-3x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}
y_2 &= y_1 \int \frac{e^{\int -\frac{3}{1} dx}}{(y_1)^2} dx \\
&= y_1 \int \frac{e^{-3x}}{(y_1)^2} dx \\
&= y_1 \left(\frac{e^{3x}}{3} \right)
\end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 (e^{-3x}) + c_2 \left(e^{-3x} \left(\frac{e^{3x}}{3} \right) \right)
\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-3x} + \frac{c_2}{3} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = -2$ and $x = 0$ in the above gives

$$-2 = c_1 + \frac{c_2}{3} \tag{1A}$$

Taking derivative of the solution gives

$$y' = -3c_1 e^{-3x}$$

substituting $y' = 3$ and $x = 0$ in the above gives

$$3 = -3c_1 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -1$$

$$c_2 = -3$$

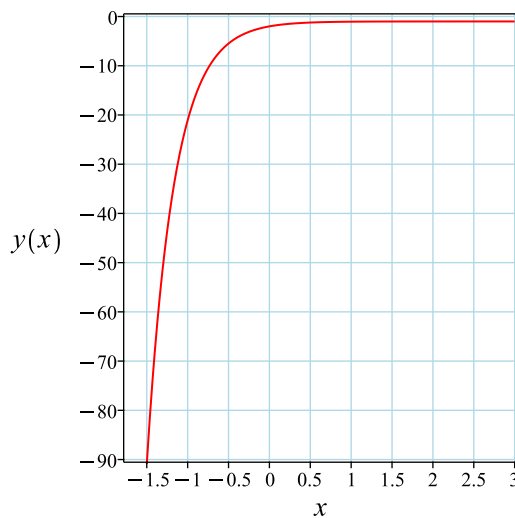
Substituting these values back in above solution results in

$$y = -1 - e^{-3x}$$

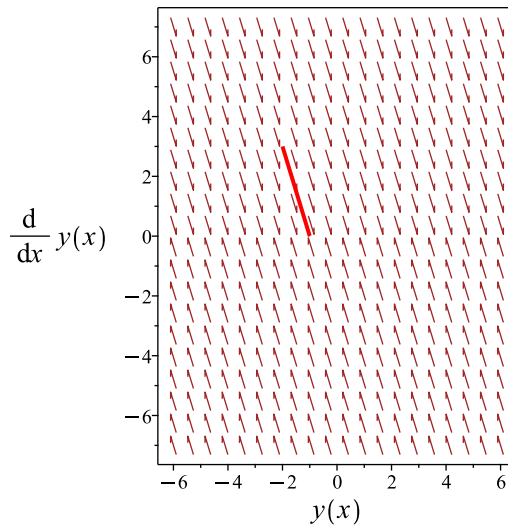
Summary

The solution(s) found are the following

$$y = -1 - e^{-3x} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -1 - e^{-3x}$$

Verified OK.

7.12.7 Solving as exact linear second order ode

An ode of the form

$$p(x)y'' + q(x)y' + r(x)y = s(x)$$

is exact if

$$p''(x) - q'(x) + r(x) = 0 \tag{1}$$

For the given ode we have

$$p(x) = 1$$

$$q(x) = 3$$

$$r(x) = 0$$

$$s(x) = 0$$

Hence

$$p''(x) = 0$$

$$q'(x) = 0$$

Therefore (1) becomes

$$0 - (0) + (0) = 0$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$(p(x)y' + (q(x) - p'(x))y)' = s(x)$$

Integrating gives

$$p(x)y' + (q(x) - p'(x))y = \int s(x) dx$$

Substituting the above values for p, q, r, s gives

$$3y + y' = c_1$$

We now have a first order ode to solve which is

$$3y + y' = c_1$$

Integrating both sides gives

$$\int \frac{1}{-3y + c_1} dy = \int dx$$
$$-\frac{\ln(-3y + c_1)}{3} = x + c_2$$

Raising both side to exponential gives

$$\frac{1}{(-3y + c_1)^{\frac{1}{3}}} = e^{x+c_2}$$

Which simplifies to

$$\frac{1}{(-3y + c_1)^{\frac{1}{3}}} = c_3 e^x$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = -\frac{e^{-3x}}{3c_3^3} + \frac{c_1}{3} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = -2$ and $x = 0$ in the above gives

$$-2 = \frac{c_1 c_3^3 - 1}{3c_3^3} \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{e^{-3x}}{c_3^3}$$

substituting $y' = 3$ and $x = 0$ in the above gives

$$3 = \frac{1}{c_3^3} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_3\}$. Warning, unable to solve for constants of integrations.

Verification of solutions N/A

7.12.8 Maple step by step solution

Let's solve

$$\left[y'' + 3y' = 0, y(0) = -2, y' \Big|_{\{x=0\}} = 3 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 3r = 0$$

- Factor the characteristic polynomial

$$r(r + 3) = 0$$

- Roots of the characteristic polynomial

$$r = (-3, 0)$$

- 1st solution of the ODE

$$y_1(x) = e^{-3x}$$

- 2nd solution of the ODE

$$y_2(x) = 1$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-3x} + c_2$$

- Check validity of solution $y = c_1 e^{-3x} + c_2$

- Use initial condition $y(0) = -2$

$$-2 = c_1 + c_2$$

- Compute derivative of the solution

$$y' = -3c_1 e^{-3x}$$

- Use the initial condition $y' \Big|_{\{x=0\}} = 3$

$$3 = -3c_1$$

- Solve for c_1 and c_2

$$\{c_1 = -1, c_2 = -1\}$$

- Substitute constant values into general solution and simplify

$$y = -1 - e^{-3x}$$

- Solution to the IVP

$$y = -1 - e^{-3x}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 12

```
dsolve([diff(y(x),x$2) +3*diff(y(x),x) = 0,y(0) = -2, D(y)(0) = 3],y(x), singsol=all)
```

$$y(x) = -1 - e^{-3x}$$

✓ Solution by Mathematica

Time used: 0.014 (sec). Leaf size: 14

```
DSolve[{y'[x]+3*y'[x]==0,{y[0]==-2,y'[0]==3}},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -e^{-3x} - 1$$

7.13 problem 13

7.13.1 Existence and uniqueness analysis	1885
7.13.2 Solving as second order linear constant coeff ode	1886
7.13.3 Solving using Kovacic algorithm	1889
7.13.4 Maple step by step solution	1893

Internal problem ID [611]

Internal file name [OUTPUT/611_Sunday_June_05_2022_01_45_48_AM_49802737/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 13.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$y'' + 5y' + 3y = 0$$

With initial conditions

$$[y(0) = 1, y'(0) = 0]$$

7.13.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 5$$

$$q(x) = 3$$

$$F = 0$$

Hence the ode is

$$y'' + 5y' + 3y = 0$$

The domain of $p(x) = 5$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = 3$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.13.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 5, C = 3$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 5\lambda e^{\lambda x} + 3e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 5\lambda + 3 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 5, C = 3$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-5}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{5^2 - (4)(1)(3)} \\ &= -\frac{5}{2} \pm \frac{\sqrt{13}}{2} \end{aligned}$$

Hence

$$\lambda_1 = -\frac{5}{2} + \frac{\sqrt{13}}{2}$$

$$\lambda_2 = -\frac{5}{2} - \frac{\sqrt{13}}{2}$$

Which simplifies to

$$\lambda_1 = -\frac{5}{2} + \frac{\sqrt{13}}{2}$$

$$\lambda_2 = -\frac{5}{2} - \frac{\sqrt{13}}{2}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{\left(-\frac{5}{2} + \frac{\sqrt{13}}{2}\right)x} + c_2 e^{\left(-\frac{5}{2} - \frac{\sqrt{13}}{2}\right)x}$$

Or

$$y = c_1 e^{\left(-\frac{5}{2} + \frac{\sqrt{13}}{2}\right)x} + c_2 e^{\left(-\frac{5}{2} - \frac{\sqrt{13}}{2}\right)x}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\left(-\frac{5}{2} + \frac{\sqrt{13}}{2}\right)x} + c_2 e^{\left(-\frac{5}{2} - \frac{\sqrt{13}}{2}\right)x} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = 0$ in the above gives

$$1 = c_1 + c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = c_1 \left(-\frac{5}{2} + \frac{\sqrt{13}}{2} \right) e^{\left(-\frac{5}{2} + \frac{\sqrt{13}}{2}\right)x} + c_2 \left(-\frac{5}{2} - \frac{\sqrt{13}}{2} \right) e^{\left(-\frac{5}{2} - \frac{\sqrt{13}}{2}\right)x}$$

substituting $y' = 0$ and $x = 0$ in the above gives

$$0 = \frac{(c_1 - c_2)\sqrt{13}}{2} - \frac{5c_1}{2} - \frac{5c_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{1}{2} + \frac{5\sqrt{13}}{26}$$

$$c_2 = \frac{1}{2} - \frac{5\sqrt{13}}{26}$$

Substituting these values back in above solution results in

$$y = \frac{e^{\frac{(-5+\sqrt{13})x}{2}}}{2} + \frac{5e^{\frac{(-5+\sqrt{13})x}{2}}\sqrt{13}}{26} + \frac{e^{-\frac{(5+\sqrt{13})x}{2}}}{2} - \frac{5e^{-\frac{(5+\sqrt{13})x}{2}}\sqrt{13}}{26}$$

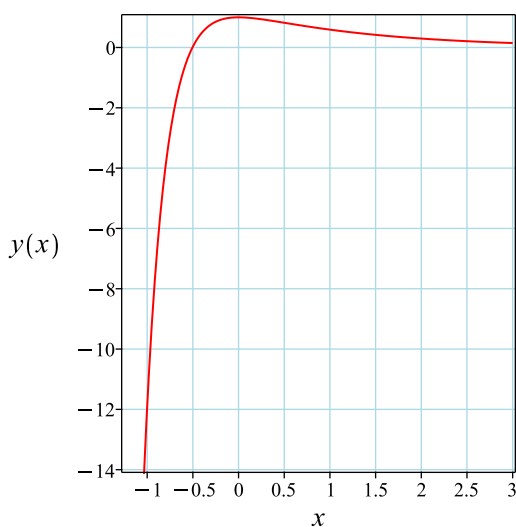
Which simplifies to

$$y = \frac{(13 + 5\sqrt{13})e^{\frac{(-5+\sqrt{13})x}{2}}}{26} + \frac{(13 - 5\sqrt{13})e^{-\frac{(5+\sqrt{13})x}{2}}}{26}$$

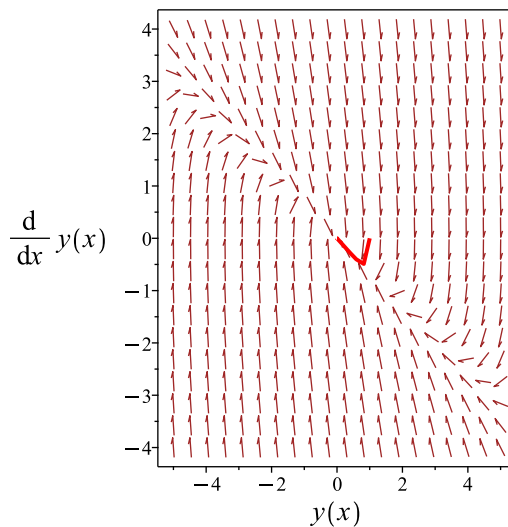
Summary

The solution(s) found are the following

$$y = \frac{(13 + 5\sqrt{13})e^{\frac{(-5+\sqrt{13})x}{2}}}{26} + \frac{(13 - 5\sqrt{13})e^{-\frac{(5+\sqrt{13})x}{2}}}{26} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{(13 + 5\sqrt{13})e^{\frac{(-5+\sqrt{13})x}{2}}}{26} + \frac{(13 - 5\sqrt{13})e^{-\frac{(5+\sqrt{13})x}{2}}}{26}$$

Verified OK.

7.13.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 5y' + 3y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 5 \quad (3)$$

$$C = 3$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \quad (5)$$
$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{13}{4} \quad (6)$$

Comparing the above to (5) shows that

$$s = 13$$

$$t = 4$$

Therefore eq. (4) becomes

$$z''(x) = \frac{13z(x)}{4} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 340: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{13}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x\sqrt{13}}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{5}{1} dx} \end{aligned}$$

$$\begin{aligned}
&= z_1 e^{-\frac{5x}{2}} \\
&= z_1 \left(e^{-\frac{5x}{2}} \right)
\end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{(5+\sqrt{13})x}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}
y_2 &= y_1 \int \frac{e^{\int -\frac{5}{1} dx}}{(y_1)^2} dx \\
&= y_1 \int \frac{e^{-5x}}{(y_1)^2} dx \\
&= y_1 \left(\frac{\sqrt{13} e^{x\sqrt{13}}}{13} \right)
\end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 \left(e^{-\frac{(5+\sqrt{13})x}{2}} \right) + c_2 \left(e^{-\frac{(5+\sqrt{13})x}{2}} \left(\frac{\sqrt{13} e^{x\sqrt{13}}}{13} \right) \right)
\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-\frac{(5+\sqrt{13})x}{2}} + \frac{c_2 e^{\frac{(-5+\sqrt{13})x}{2}} \sqrt{13}}{13} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = 0$ in the above gives

$$1 = c_1 + \frac{c_2 \sqrt{13}}{13} \quad (1A)$$

Taking derivative of the solution gives

$$y' = c_1 \left(-\frac{5}{2} - \frac{\sqrt{13}}{2} \right) e^{-\frac{(5+\sqrt{13})x}{2}} + \frac{c_2 \left(-\frac{5}{2} + \frac{\sqrt{13}}{2} \right) e^{-\frac{(-5+\sqrt{13})x}{2}} \sqrt{13}}{13}$$

substituting $y' = 0$ and $x = 0$ in the above gives

$$0 = \frac{(-13c_1 - 5c_2) \sqrt{13}}{26} - \frac{5c_1}{2} + \frac{c_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{1}{2} - \frac{5\sqrt{13}}{26}$$

$$c_2 = \frac{5}{2} + \frac{\sqrt{13}}{2}$$

Substituting these values back in above solution results in

$$y = \frac{e^{\frac{(-5+\sqrt{13})x}{2}}}{2} + \frac{5e^{\frac{(-5+\sqrt{13})x}{2}} \sqrt{13}}{26} + \frac{e^{-\frac{(5+\sqrt{13})x}{2}}}{2} - \frac{5e^{-\frac{(5+\sqrt{13})x}{2}} \sqrt{13}}{26}$$

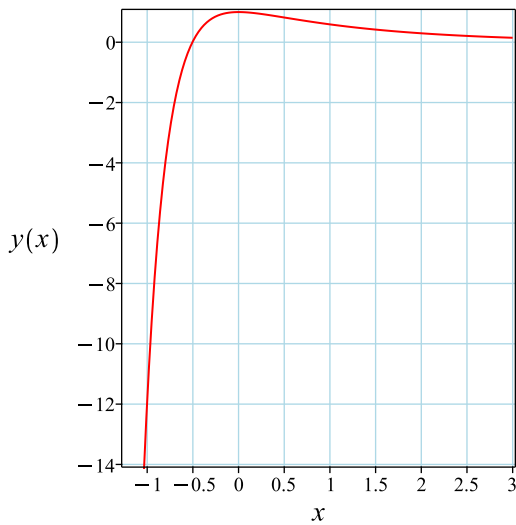
Which simplifies to

$$y = \frac{(13 + 5\sqrt{13}) e^{\frac{(-5+\sqrt{13})x}{2}}}{26} + \frac{(13 - 5\sqrt{13}) e^{-\frac{(5+\sqrt{13})x}{2}}}{26}$$

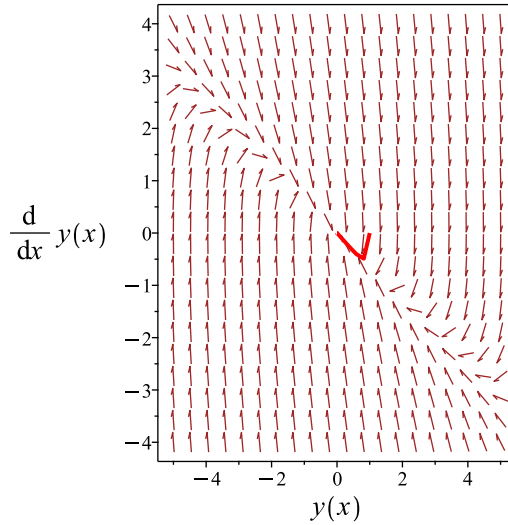
Summary

The solution(s) found are the following

$$y = \frac{(13 + 5\sqrt{13}) e^{\frac{(-5+\sqrt{13})x}{2}}}{26} + \frac{(13 - 5\sqrt{13}) e^{-\frac{(5+\sqrt{13})x}{2}}}{26} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{(13 + 5\sqrt{13}) e^{\frac{(-5+\sqrt{13})x}{2}}}{26} + \frac{(13 - 5\sqrt{13}) e^{-\frac{(5+\sqrt{13})x}{2}}}{26}$$

Verified OK.

7.13.4 Maple step by step solution

Let's solve

$$\left[y'' + 5y' + 3y = 0, y(0) = 1, y'|_{\{x=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

y''

- Characteristic polynomial of ODE

$$r^2 + 5r + 3 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-5) \pm (\sqrt{13})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{5}{2} - \frac{\sqrt{13}}{2}, -\frac{5}{2} + \frac{\sqrt{13}}{2} \right)$$

- 1st solution of the ODE

$$y_1(x) = e^{\left(-\frac{5}{2} - \frac{\sqrt{13}}{2}\right)x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{\left(-\frac{5}{2} + \frac{\sqrt{13}}{2}\right)x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{\left(-\frac{5}{2} - \frac{\sqrt{13}}{2}\right)x} + c_2 e^{\left(-\frac{5}{2} + \frac{\sqrt{13}}{2}\right)x}$$

- Check validity of solution $y = c_1 e^{\left(-\frac{5}{2} - \frac{\sqrt{13}}{2}\right)x} + c_2 e^{\left(-\frac{5}{2} + \frac{\sqrt{13}}{2}\right)x}$

- Use initial condition $y(0) = 1$

$$1 = c_1 + c_2$$

- Compute derivative of the solution

$$y' = c_1 \left(-\frac{5}{2} - \frac{\sqrt{13}}{2}\right) e^{\left(-\frac{5}{2} - \frac{\sqrt{13}}{2}\right)x} + c_2 \left(-\frac{5}{2} + \frac{\sqrt{13}}{2}\right) e^{\left(-\frac{5}{2} + \frac{\sqrt{13}}{2}\right)x}$$

- Use the initial condition $y'|_{\{x=0\}} = 0$

$$0 = c_1 \left(-\frac{5}{2} - \frac{\sqrt{13}}{2}\right) + \left(-\frac{5}{2} + \frac{\sqrt{13}}{2}\right) c_2$$

- Solve for c_1 and c_2

$$\left\{ c_1 = \frac{1}{2} - \frac{5\sqrt{13}}{26}, c_2 = \frac{1}{2} + \frac{5\sqrt{13}}{26} \right\}$$

- Substitute constant values into general solution and simplify

$$y = \frac{(13+5\sqrt{13})e^{\frac{(-5+\sqrt{13})x}{2}}}{26} + \frac{(13-5\sqrt{13})e^{-\frac{(5+\sqrt{13})x}{2}}}{26}$$

- Solution to the IVP

$$y = \frac{(13+5\sqrt{13})e^{\frac{(-5+\sqrt{13})x}{2}}}{26} + \frac{(13-5\sqrt{13})e^{-\frac{(5+\sqrt{13})x}{2}}}{26}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 39

```
dsolve([diff(y(x),x$2) +5*diff(y(x),x)+3*y(x) = 0,y(0) = 1, D(y)(0) = 0],y(x), singsol=all)
```

$$y(x) = \frac{(13 + 5\sqrt{13}) e^{\frac{(-5+\sqrt{13})x}{2}}}{26} + \frac{(13 - 5\sqrt{13}) e^{-\frac{(5+\sqrt{13})x}{2}}}{26}$$

✓ Solution by Mathematica

Time used: 0.023 (sec). Leaf size: 51

```
DSolve[{y''[x]+5*y'[x]+3*y[x]==0,{y[0]==1,y'[0]==0}},y[x],x,IncludeSingularSolutions -> True
```

$$y(x) \rightarrow \frac{1}{26} e^{-\frac{1}{2}(5+\sqrt{13})x} \left((13 + 5\sqrt{13}) e^{\sqrt{13}x} + 13 - 5\sqrt{13} \right)$$

7.14 problem 14

7.14.1 Existence and uniqueness analysis	1896
7.14.2 Solving as second order linear constant coeff ode	1897
7.14.3 Solving using Kovacic algorithm	1900
7.14.4 Maple step by step solution	1904

Internal problem ID [612]

Internal file name [OUTPUT/612_Sunday_June_05_2022_01_45_50_AM_78937181/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 14.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$2y'' + y' - 4y = 0$$

With initial conditions

$$[y(0) = 0, y'(0) = 1]$$

7.14.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$\begin{aligned} p(x) &= \frac{1}{2} \\ q(x) &= -2 \\ F &= 0 \end{aligned}$$

Hence the ode is

$$y'' + \frac{y'}{2} - 2y = 0$$

The domain of $p(x) = \frac{1}{2}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = -2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.14.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 2, B = 1, C = -4$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$2\lambda^2 e^{\lambda x} + \lambda e^{\lambda x} - 4e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$2\lambda^2 + \lambda - 4 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 2, B = 1, C = -4$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-1}{(2)(2)} \pm \frac{1}{(2)(2)} \sqrt{1^2 - (4)(2)(-4)} \\ &= -\frac{1}{4} \pm \frac{\sqrt{33}}{4} \end{aligned}$$

Hence

$$\lambda_1 = -\frac{1}{4} + \frac{\sqrt{33}}{4}$$

$$\lambda_2 = -\frac{1}{4} - \frac{\sqrt{33}}{4}$$

Which simplifies to

$$\lambda_1 = -\frac{1}{4} + \frac{\sqrt{33}}{4}$$

$$\lambda_2 = -\frac{1}{4} - \frac{\sqrt{33}}{4}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{\left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right)x} + c_2 e^{\left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right)x}$$

Or

$$y = c_1 e^{\left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right)x} + c_2 e^{\left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right)x}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right)x} + c_2 e^{\left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right)x} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $x = 0$ in the above gives

$$0 = c_1 + c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = c_1 \left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right) e^{\left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right)x} + c_2 \left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right) e^{\left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right)x}$$

substituting $y' = 1$ and $x = 0$ in the above gives

$$1 = \frac{(c_1 - c_2)\sqrt{33}}{4} - \frac{c_1}{4} - \frac{c_2}{4} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{2\sqrt{33}}{33}$$
$$c_2 = -\frac{2\sqrt{33}}{33}$$

Substituting these values back in above solution results in

$$y = \frac{2\sqrt{33} e^{\frac{(-1+\sqrt{33})x}{4}}}{33} - \frac{2\sqrt{33} e^{-\frac{(1+\sqrt{33})x}{4}}}{33}$$

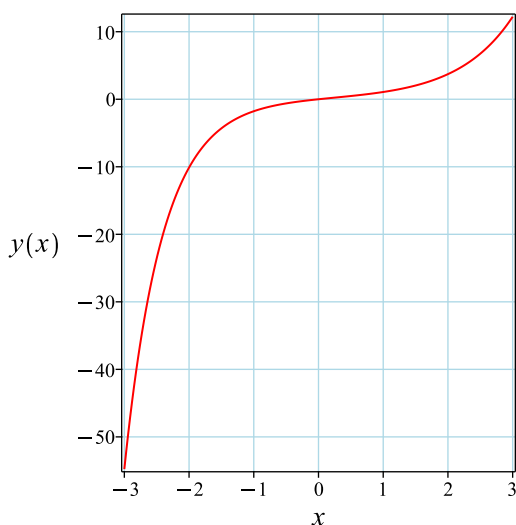
Which simplifies to

$$y = \frac{2\sqrt{33} \left(e^{\frac{(-1+\sqrt{33})x}{4}} - e^{-\frac{(1+\sqrt{33})x}{4}} \right)}{33}$$

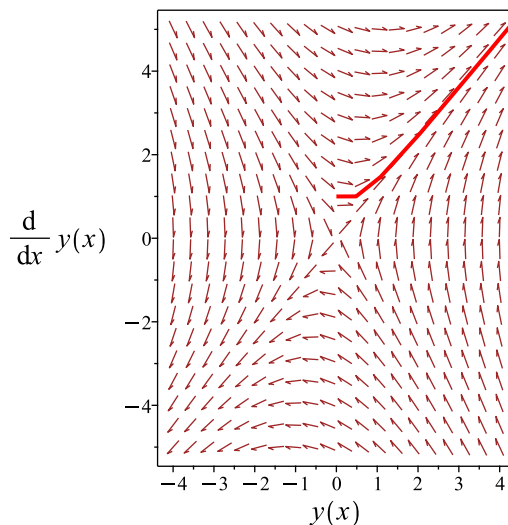
Summary

The solution(s) found are the following

$$y = \frac{2\sqrt{33} \left(e^{\frac{(-1+\sqrt{33})x}{4}} - e^{-\frac{(1+\sqrt{33})x}{4}} \right)}{33} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{2\sqrt{33} \left(e^{\frac{(-1+\sqrt{33})x}{4}} - e^{-\frac{(1+\sqrt{33})x}{4}} \right)}{33}$$

Verified OK.

7.14.3 Solving using Kovacic algorithm

Writing the ode as

$$2y'' + y' - 4y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 2$$

$$B = 1 \quad (3)$$

$$C = -4$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \quad (5)$$
$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{33}{16} \quad (6)$$

Comparing the above to (5) shows that

$$s = 33$$

$$t = 16$$

Therefore eq. (4) becomes

$$z''(x) = \frac{33z(x)}{16} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 342: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{33}{16}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x\sqrt{33}}{4}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{1}{2} dx} \end{aligned}$$

$$\begin{aligned}
&= z_1 e^{-\frac{x}{4}} \\
&= z_1 \left(e^{-\frac{x}{4}} \right)
\end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{(1+\sqrt{33})x}{4}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}
y_2 &= y_1 \int \frac{e^{\int -\frac{1}{2} dx}}{(y_1)^2} dx \\
&= y_1 \int \frac{e^{-\frac{x}{2}}}{(y_1)^2} dx \\
&= y_1 \left(\frac{2\sqrt{33} e^{\frac{x\sqrt{33}}{2}}}{33} \right)
\end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 \left(e^{-\frac{(1+\sqrt{33})x}{4}} \right) + c_2 \left(e^{-\frac{(1+\sqrt{33})x}{4}} \left(\frac{2\sqrt{33} e^{\frac{x\sqrt{33}}{2}}}{33} \right) \right)
\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-\frac{(1+\sqrt{33})x}{4}} + \frac{2c_2\sqrt{33} e^{\frac{(-1+\sqrt{33})x}{4}}}{33} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $x = 0$ in the above gives

$$0 = c_1 + \frac{2c_2\sqrt{33}}{33} \quad (1A)$$

Taking derivative of the solution gives

$$y' = c_1 \left(-\frac{1}{4} - \frac{\sqrt{33}}{4} \right) e^{-\frac{(1+\sqrt{33})x}{4}} + \frac{2c_2\sqrt{33} \left(-\frac{1}{4} + \frac{\sqrt{33}}{4} \right) e^{\frac{(-1+\sqrt{33})x}{4}}}{33}$$

substituting $y' = 1$ and $x = 0$ in the above gives

$$1 = \frac{(-33c_1 - 2c_2)\sqrt{33}}{132} - \frac{c_1}{4} + \frac{c_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -\frac{2\sqrt{33}}{33}$$

$$c_2 = 1$$

Substituting these values back in above solution results in

$$y = \frac{2\sqrt{33} e^{\frac{(-1+\sqrt{33})x}{4}}}{33} - \frac{2\sqrt{33} e^{-\frac{(1+\sqrt{33})x}{4}}}{33}$$

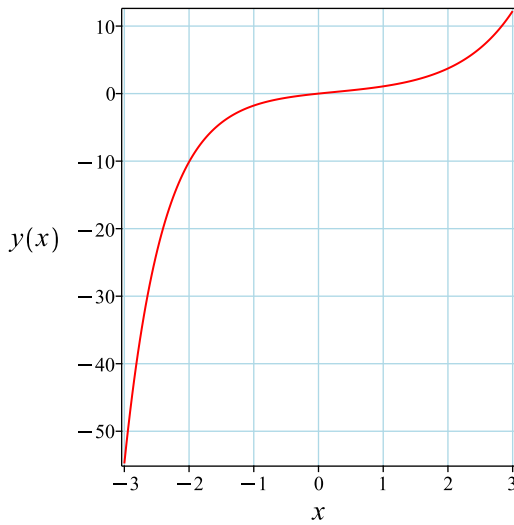
Which simplifies to

$$y = \frac{2\sqrt{33} \left(e^{\frac{(-1+\sqrt{33})x}{4}} - e^{-\frac{(1+\sqrt{33})x}{4}} \right)}{33}$$

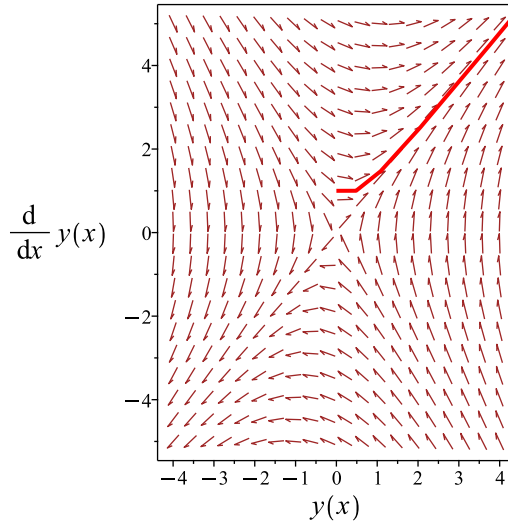
Summary

The solution(s) found are the following

$$y = \frac{2\sqrt{33} \left(e^{\frac{(-1+\sqrt{33})x}{4}} - e^{-\frac{(1+\sqrt{33})x}{4}} \right)}{33} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{2\sqrt{33} \left(e^{\frac{(-1+\sqrt{33})x}{4}} - e^{-\frac{(1+\sqrt{33})x}{4}} \right)}{33}$$

Verified OK.

7.14.4 Maple step by step solution

Let's solve

$$\left[2y'' + y' - 4y = 0, y(0) = 0, y'|_{\{x=0\}} = 1 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{y'}{2} + 2y$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{y'}{2} - 2y = 0$$

- Characteristic polynomial of ODE

$$r^2 + \frac{1}{2}r - 2 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-\frac{1}{2}) \pm (\sqrt{\frac{33}{4}})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{4} - \frac{\sqrt{33}}{4}, -\frac{1}{4} + \frac{\sqrt{33}}{4}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{\left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right)x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{\left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right)x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{\left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right)x} + c_2 e^{\left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right)x}$$

- Check validity of solution $y = c_1 e^{\left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right)x} + c_2 e^{\left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right)x}$

- Use initial condition $y(0) = 0$

$$0 = c_1 + c_2$$

- Compute derivative of the solution

$$y' = c_1 \left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right) e^{\left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right)x} + c_2 \left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right) e^{\left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right)x}$$

- Use the initial condition $y' \Big|_{\{x=0\}} = 1$

$$1 = \left(-\frac{1}{4} - \frac{\sqrt{33}}{4}\right) c_1 + \left(-\frac{1}{4} + \frac{\sqrt{33}}{4}\right) c_2$$

- Solve for c_1 and c_2

$$\left\{ c_1 = -\frac{2\sqrt{33}}{33}, c_2 = \frac{2\sqrt{33}}{33} \right\}$$

- Substitute constant values into general solution and simplify

$$y = \frac{2\sqrt{33} \left(e^{\frac{(-1+\sqrt{33})x}{4}} - e^{-\frac{(1+\sqrt{33})x}{4}} \right)}{33}$$

- Solution to the IVP

$$y = \frac{2\sqrt{33} \left(e^{\frac{(-1+\sqrt{33})x}{4}} - e^{-\frac{(1+\sqrt{33})x}{4}} \right)}{33}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 30

```
dsolve([2*diff(y(x),x$2) +diff(y(x),x)-4*y(x) = 0,y(0) = 0, D(y)(0) = 1],y(x), singsol=all)
```

$$y(x) = \frac{2 \left(e^{\frac{(-1+\sqrt{33})x}{4}} - e^{-\frac{(1+\sqrt{33})x}{4}} \right) \sqrt{33}}{33}$$

✓ Solution by Mathematica

Time used: 0.027 (sec). Leaf size: 40

```
DSolve[{2*y''[x]+y'[x]-4*y[x]==0,{y[0]==0,y'[0]==1}},y[x],x,IncludeSingularSolutions -> True
```

$$y(x) \rightarrow \frac{2e^{-\frac{1}{4}(1+\sqrt{33})x} \left(e^{\frac{\sqrt{33}x}{2}} - 1 \right)}{\sqrt{33}}$$

7.15 problem 15

7.15.1 Existence and uniqueness analysis	1907
7.15.2 Solving as second order linear constant coeff ode	1908
7.15.3 Solving using Kovacic algorithm	1910
7.15.4 Maple step by step solution	1915

Internal problem ID [613]

Internal file name [OUTPUT/613_Sunday_June_05_2022_01_45_51_AM_15862367/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 15.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 8y' - 9y = 0$$

With initial conditions

$$[y(1) = 1, y'(1) = 0]$$

7.15.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 8$$

$$q(x) = -9$$

$$F = 0$$

Hence the ode is

$$y'' + 8y' - 9y = 0$$

The domain of $p(x) = 8$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 1$ is inside this domain. The domain of $q(x) = -9$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 1$ is also inside this domain. Hence solution exists and is unique.

7.15.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 8, C = -9$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 8\lambda e^{\lambda x} - 9e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 8\lambda - 9 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 8, C = -9$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-8}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{8^2 - (4)(1)(-9)} \\ &= -4 \pm 5 \end{aligned}$$

Hence

$$\lambda_1 = -4 + 5$$

$$\lambda_2 = -4 - 5$$

Which simplifies to

$$\lambda_1 = 1$$

$$\lambda_2 = -9$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(1)x} + c_2 e^{(-9)x}$$

Or

$$y = c_1 e^x + c_2 e^{-9x}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^x + c_2 e^{-9x} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = 1$ in the above gives

$$1 = (e^{10} c_1 + c_2) e^{-9} \quad (1A)$$

Taking derivative of the solution gives

$$y' = c_1 e^x - 9c_2 e^{-9x}$$

substituting $y' = 0$ and $x = 1$ in the above gives

$$0 = (e^{10} c_1 - 9c_2) e^{-9} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{9 e^{-1}}{10}$$

$$c_2 = \frac{e^9}{10}$$

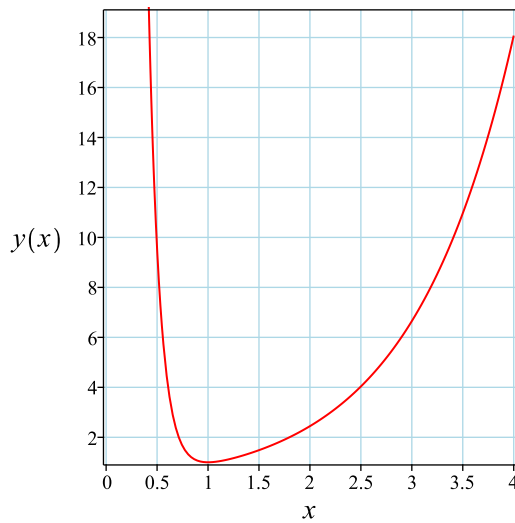
Substituting these values back in above solution results in

$$y = \frac{9 e^x e^{-1}}{10} + \frac{e^{-9x} e^9}{10}$$

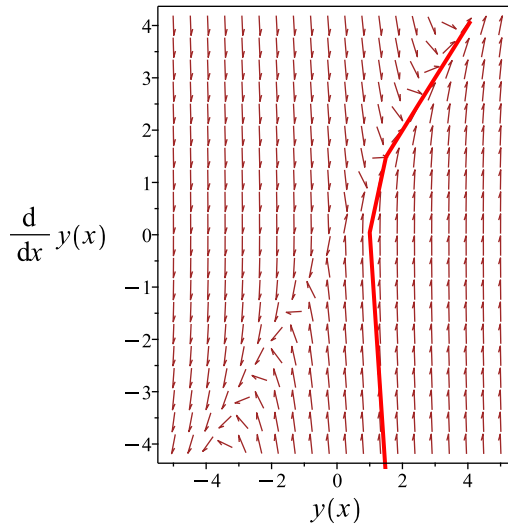
Summary

The solution(s) found are the following

$$y = \frac{9e^xe^{-1}}{10} + \frac{e^{-9x}e^9}{10} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{9e^xe^{-1}}{10} + \frac{e^{-9x}e^9}{10}$$

Verified OK.

7.15.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 8y' - 9y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 8 \\ C &= -9 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{25}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 25 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = 25z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 344: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 25$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-5x}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}
 y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\
 &= z_1 e^{-\int \frac{1}{2} \frac{8}{1} dx} \\
 &= z_1 e^{-4x} \\
 &= z_1 (e^{-4x})
 \end{aligned}$$

Which simplifies to

$$y_1 = e^{-9x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{8}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-8x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{e^{10x}}{10} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-9x}) + c_2 \left(e^{-9x} \left(\frac{e^{10x}}{10} \right) \right) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-9x} + \frac{c_2 e^x}{10} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = 1$ in the above gives

$$1 = \frac{(e^{10} c_2 + 10 c_1) e^{-9}}{10} \tag{1A}$$

Taking derivative of the solution gives

$$y' = -9c_1 e^{-9x} + \frac{c_2 e^x}{10}$$

substituting $y' = 0$ and $x = 1$ in the above gives

$$0 = \frac{(e^{10}c_2 - 90c_1)e^{-9}}{10} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{e^9}{10}$$
$$c_2 = 9e^{-1}$$

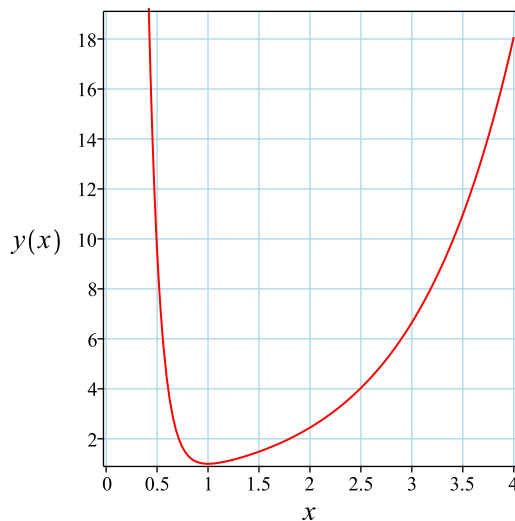
Substituting these values back in above solution results in

$$y = \frac{9e^xe^{-1}}{10} + \frac{e^{-9x}e^9}{10}$$

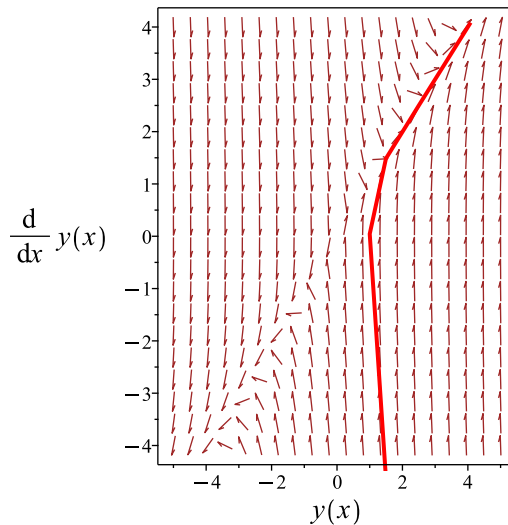
Summary

The solution(s) found are the following

$$y = \frac{9e^xe^{-1}}{10} + \frac{e^{-9x}e^9}{10} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{9e^xe^{-1}}{10} + \frac{e^{-9x}e^9}{10}$$

Verified OK.

7.15.4 Maple step by step solution

Let's solve

$$\left[y'' + 8y' - 9y = 0, y(1) = 1, y' \Big|_{\{x=1\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 8r - 9 = 0$$

- Factor the characteristic polynomial

$$(r + 9)(r - 1) = 0$$

- Roots of the characteristic polynomial

$$r = (-9, 1)$$

- 1st solution of the ODE

$$y_1(x) = e^{-9x}$$

- 2nd solution of the ODE

$$y_2(x) = e^x$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-9x} + c_2 e^x$$

- Check validity of solution $y = c_1 e^{-9x} + c_2 e^x$

- Use initial condition $y(1) = 1$

$$1 = c_1 e^{-9} + c_2$$

- Compute derivative of the solution

$$y' = -9c_1 e^{-9x} + c_2 e^x$$

- Use the initial condition $y' \Big|_{\{x=1\}} = 0$

$$0 = -9c_1 e^{-9} + c_2$$

- Solve for c_1 and c_2

$$\left\{ c_1 = \frac{1}{10e^{-9}}, c_2 = \frac{9}{10e} \right\}$$

- Substitute constant values into general solution and simplify

$$y = \frac{9e^{x-1}}{10} + \frac{e^{-9x+9}}{10}$$

- Solution to the IVP

$$y = \frac{9e^{x-1}}{10} + \frac{e^{-9x+9}}{10}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 19

```
dsolve([diff(y(x),x$2) +8*diff(y(x),x)-9*y(x) = 0,y(1) = 1, D(y)(1) = 0],y(x), singsol=all)
```

$$y(x) = \frac{e^{9-9x}}{10} + \frac{9e^{x-1}}{10}$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 26

```
DSolve[{y''[x]+8*y'[x]-9*y[x]==0,{y[1]==1,y'[1]==0}},y[x],x,IncludeSingularSolutions -> True
```

$$y(x) \rightarrow \frac{1}{10}e^{9-9x} + \frac{9e^{x-1}}{10}$$

7.16 problem 16

7.16.1 Existence and uniqueness analysis	1917
7.16.2 Solving as second order linear constant coeff ode	1918
7.16.3 Solving as second order ode can be made integrable ode	1920
7.16.4 Solving using Kovacic algorithm	1924
7.16.5 Maple step by step solution	1928

Internal problem ID [614]

Internal file name [OUTPUT/614_Sunday_June_05_2022_01_45_52_AM_59244929/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 16.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "second_order_ode_can_be_made_integrable"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$4y'' - y = 0$$

With initial conditions

$$[y(-2) = 1, y'(-2) = -1]$$

7.16.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 0$$

$$q(x) = -\frac{1}{4}$$

$$F = 0$$

Hence the ode is

$$y'' - \frac{y}{4} = 0$$

The domain of $p(x) = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = -2$ is inside this domain. The domain of $q(x) = -\frac{1}{4}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = -2$ is also inside this domain. Hence solution exists and is unique.

7.16.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 4, B = 0, C = -1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$4\lambda^2 e^{\lambda x} - e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$4\lambda^2 - 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 4, B = 0, C = -1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(4)} \pm \frac{1}{(2)(4)} \sqrt{0^2 - (4)(4)(-1)} \\ &= \pm \frac{1}{2} \end{aligned}$$

Hence

$$\lambda_1 = +\frac{1}{2}$$

$$\lambda_2 = -\frac{1}{2}$$

Which simplifies to

$$\lambda_1 = \frac{1}{2}$$

$$\lambda_2 = -\frac{1}{2}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(\frac{1}{2})x} + c_2 e^{(-\frac{1}{2})x}$$

Or

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{-\frac{x}{2}}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{-\frac{x}{2}} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = -2$ in the above gives

$$1 = e^{-1}c_1 + ec_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{c_1 e^{\frac{x}{2}}}{2} - \frac{c_2 e^{-\frac{x}{2}}}{2}$$

substituting $y' = -1$ and $x = -2$ in the above gives

$$-1 = \frac{e^{-1}c_1}{2} - \frac{ec_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -\frac{e}{2}$$
$$c_2 = \frac{3e^{-1}}{2}$$

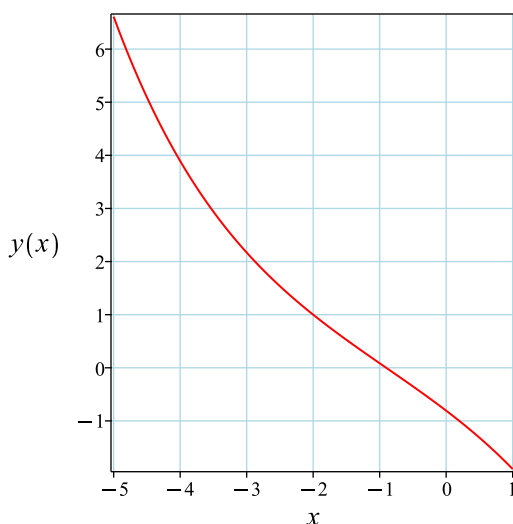
Substituting these values back in above solution results in

$$y = -\frac{e^{\frac{x}{2}}e}{2} + \frac{3e^{-\frac{x}{2}}e^{-1}}{2}$$

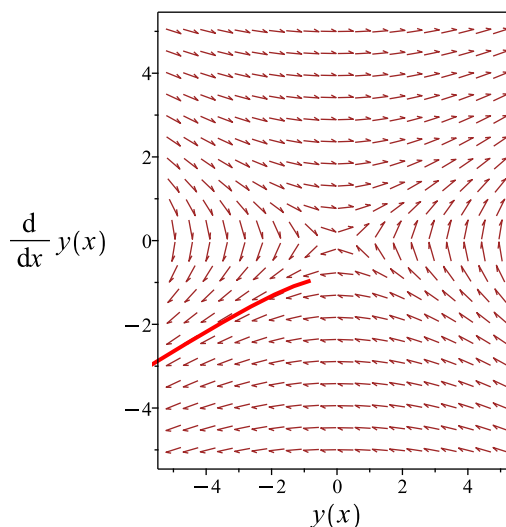
Summary

The solution(s) found are the following

$$y = -\frac{e^{\frac{x}{2}}e}{2} + \frac{3e^{-\frac{x}{2}}e^{-1}}{2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{e^{\frac{x}{2}}e}{2} + \frac{3e^{-\frac{x}{2}}e^{-1}}{2}$$

Verified OK.

7.16.3 Solving as second order ode can be made integrable ode

Multiplying the ode by y' gives

$$4y'y'' - yy' = 0$$

Integrating the above w.r.t x gives

$$\int (4y'y'' - yy') dx = 0$$

$$2y'^2 - \frac{y^2}{2} = c_2$$

Which is now solved for y . Solving the given ode for y' results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$y' = \frac{\sqrt{y^2 + 2c_1}}{2} \quad (1)$$

$$y' = -\frac{\sqrt{y^2 + 2c_1}}{2} \quad (2)$$

Now each one of the above ODE is solved.

Solving equation (1)

Integrating both sides gives

$$\int \frac{2}{\sqrt{y^2 + 2c_1}} dy = \int dx$$
$$2 \ln \left(y + \sqrt{y^2 + 2c_1} \right) = x + c_2$$

Raising both side to exponential gives

$$\left(y + \sqrt{y^2 + 2c_1} \right)^2 = e^{x+c_2}$$

Which simplifies to

$$\left(y + \sqrt{y^2 + 2c_1} \right)^2 = c_3 e^x$$

Solving equation (2)

Integrating both sides gives

$$\int -\frac{2}{\sqrt{y^2 + 2c_1}} dy = \int dx$$
$$-2 \ln \left(y + \sqrt{y^2 + 2c_1} \right) = x + c_4$$

Raising both side to exponential gives

$$\frac{1}{\left(y + \sqrt{y^2 + 2c_1} \right)^2} = e^{x+c_4}$$

Which simplifies to

$$\frac{1}{\left(y + \sqrt{y^2 + 2c_1} \right)^2} = c_5 e^x$$

Initial conditions are used to solve for the constants of integration.

Looking at the First solution

$$y = \frac{c_3 e^x - 2c_1}{2\sqrt{c_3 e^x}} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = -2$ in the above gives

$$1 = -\frac{(e^2 c_1 - \frac{c_3}{2}) e^{-1}}{\sqrt{c_3}} \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{c_3 e^x}{2\sqrt{c_3 e^x}} - \frac{(c_3 e^x - 2c_1) c_3 e^x}{4(c_3 e^x)^{\frac{3}{2}}}$$

substituting $y' = -1$ and $x = -2$ in the above gives

$$-1 = \frac{2e c_1 + c_3 e^{-1}}{4\sqrt{c_3}} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_3\}$. There is no solution for the constants of integrations. This solution is removed.

Looking at the Second solution

$$y = -\frac{2c_1 c_5 e^x - 1}{2\sqrt{c_5 e^x}} \quad (2)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = -2$ in the above gives

$$1 = \frac{-2c_1 c_5 e^{-1} + e}{2\sqrt{c_5}} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -\frac{c_5 e^x c_1}{\sqrt{c_5 e^x}} + \frac{(2c_1 c_5 e^x - 1) c_5 e^x}{4(c_5 e^x)^{\frac{3}{2}}}$$

substituting $y' = -1$ and $x = -2$ in the above gives

$$-1 = -\frac{c_1 c_5 e^{-1} + \frac{e}{2}}{2\sqrt{c_5}} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_5\}$. Solving for the constants gives

$$c_1 = \frac{3}{2}$$

$$c_5 = \frac{e^2}{9}$$

Substituting these values back in above solution results in

$$y = \frac{-e^{-1}e^{2+x} + 3e^{-1}}{2\sqrt{e^x}}$$

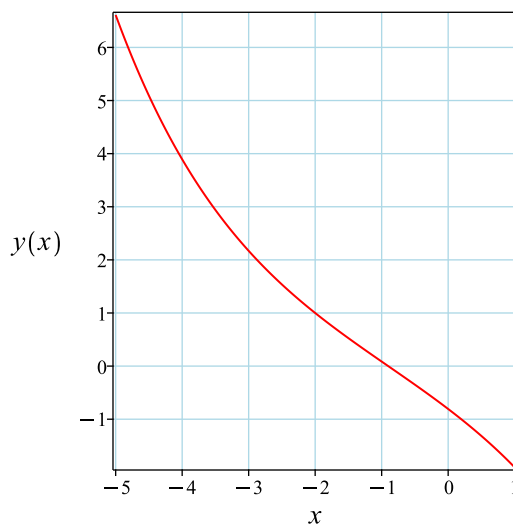
Which simplifies to

$$y = -\frac{(e^{2+x} - 3)e^{-1}}{2\sqrt{e^x}}$$

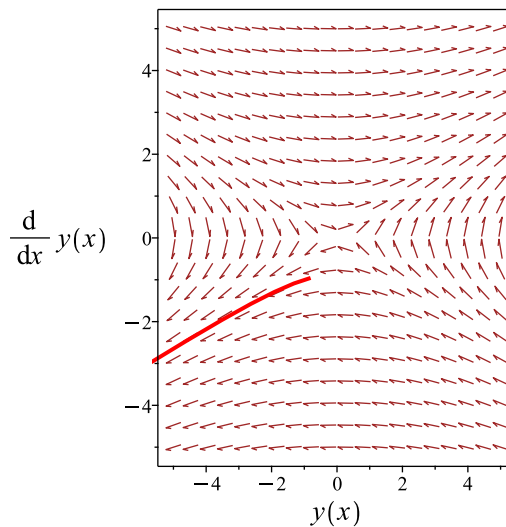
Summary

The solution(s) found are the following

$$y = -\frac{(e^{2+x} - 3)e^{-1}}{2\sqrt{e^x}} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{(e^{2+x} - 3)e^{-1}}{2\sqrt{e^x}}$$

Verified OK.

7.16.4 Solving using Kovacic algorithm

Writing the ode as

$$4y'' - y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 4$$

$$B = 0 \quad (3)$$

$$C = -1$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \quad (5)$$
$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{4} \quad (6)$$

Comparing the above to (5) shows that

$$s = 1$$

$$t = 4$$

Therefore eq. (4) becomes

$$z''(x) = \frac{z(x)}{4} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 346: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}y_1 &= z_1 \\ &= e^{-\frac{x}{2}}\end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{x}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Since $B = 0$ then the above becomes

$$\begin{aligned}y_2 &= y_1 \int \frac{1}{y_1^2} dx \\ &= e^{-\frac{x}{2}} \int \frac{1}{e^{-x}} dx \\ &= e^{-\frac{x}{2}} (e^x)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-\frac{x}{2}}) + c_2 (e^{-\frac{x}{2}} (e^x))\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-\frac{x}{2}} + c_2 e^{\frac{x}{2}} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = -2$ in the above gives

$$1 = e c_1 + c_2 e^{-1} \tag{1A}$$

Taking derivative of the solution gives

$$y' = -\frac{c_1 e^{-\frac{x}{2}}}{2} + \frac{c_2 e^{\frac{x}{2}}}{2}$$

substituting $y' = -1$ and $x = -2$ in the above gives

$$-1 = -\frac{ec_1}{2} + \frac{c_2 e^{-1}}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{3e^{-1}}{2}$$

$$c_2 = -\frac{e}{2}$$

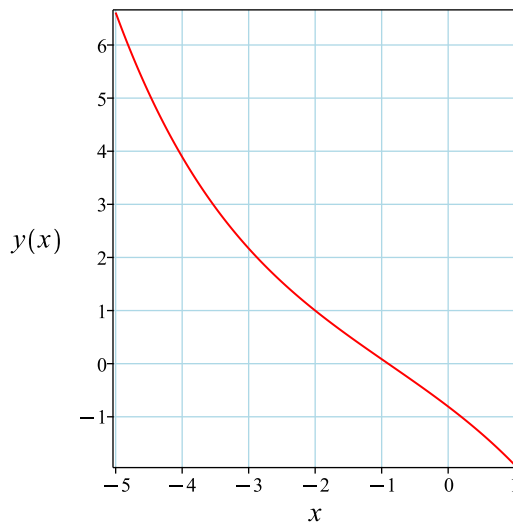
Substituting these values back in above solution results in

$$y = -\frac{e^{\frac{x}{2}} e}{2} + \frac{3e^{-\frac{x}{2}} e^{-1}}{2}$$

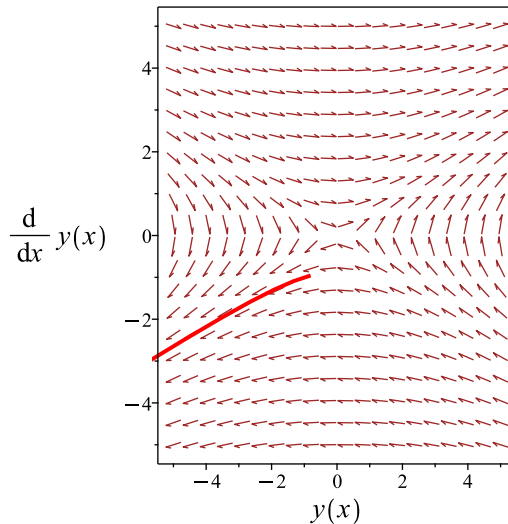
Summary

The solution(s) found are the following

$$y = -\frac{e^{\frac{x}{2}} e}{2} + \frac{3e^{-\frac{x}{2}} e^{-1}}{2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{e^{\frac{x}{2}}e}{2} + \frac{3e^{-\frac{x}{2}}e^{-1}}{2}$$

Verified OK.

7.16.5 Maple step by step solution

Let's solve

$$\left[4y'' - y = 0, y(-2) = 1, y' \Big|_{\{x=-2\}} = -1 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{y}{4}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{y}{4} = 0$$

- Characteristic polynomial of ODE

$$r^2 - \frac{1}{4} = 0$$

- Factor the characteristic polynomial

$$\frac{(2r-1)(2r+1)}{4} = 0$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{2}, \frac{1}{2}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{x}{2}}$$

- 2nd solution of the ODE

$$y_2(x) = e^{\frac{x}{2}}$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^{-\frac{x}{2}} + c_2e^{\frac{x}{2}}$$

- Check validity of solution $y = c_1e^{-\frac{x}{2}} + c_2e^{\frac{x}{2}}$

- Use initial condition $y(-2) = 1$

$$1 = ec_1 + c_2e^{-1}$$
- Compute derivative of the solution
$$y' = -\frac{c_1e^{-\frac{x}{2}}}{2} + \frac{c_2e^{\frac{x}{2}}}{2}$$
- Use the initial condition $y' \Big|_{\{x=-2\}} = -1$

$$-1 = -\frac{ec_1}{2} + \frac{c_2e^{-1}}{2}$$
- Solve for c_1 and c_2

$$\left\{ c_1 = \frac{3}{2e}, c_2 = -\frac{1}{2e^{-1}} \right\}$$
- Substitute constant values into general solution and simplify
$$y = -\frac{e^{\frac{x}{2}+1}}{2} + \frac{3e^{-\frac{x}{2}-1}}{2}$$
- Solution to the IVP
$$y = -\frac{e^{\frac{x}{2}+1}}{2} + \frac{3e^{-\frac{x}{2}-1}}{2}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 21

```
dsolve([4*diff(y(x),x$2) -y(x) = 0,y(-2) = 1, D(y)(-2) = -1],y(x), singsol=all)
```

$$y(x) = -\frac{e^{1+\frac{x}{2}}}{2} + \frac{3e^{-1-\frac{x}{2}}}{2}$$

✓ Solution by Mathematica

Time used: 0.012 (sec). Leaf size: 25

```
DSolve[{4*y'[x]-y[x]==0,{y[-2]==1,y'[-2]==-1}},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{1}{2}e^{-\frac{x}{2}-1}(e^{x+2} - 3)$$

7.17 problem 19

7.17.1 Existence and uniqueness analysis	1931
7.17.2 Solving as second order linear constant coeff ode	1932
7.17.3 Solving as second order ode can be made integrable ode	1934
7.17.4 Solving using Kovacic algorithm	1938
7.17.5 Maple step by step solution	1942

Internal problem ID [615]

Internal file name [OUTPUT/615_Sunday_June_05_2022_01_45_53_AM_95572567/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 19.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "second_order_ode_can_be_made_integrable"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - y = 0$$

With initial conditions

$$\left[y(0) = \frac{5}{4}, y'(0) = -\frac{3}{4} \right]$$

7.17.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 0$$

$$q(x) = -1$$

$$F = 0$$

Hence the ode is

$$y'' - y = 0$$

The domain of $p(x) = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = -1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.17.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 0, C = -1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = -1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(-1)} \\ &= \pm 1 \end{aligned}$$

Hence

$$\lambda_1 = +1$$

$$\lambda_2 = -1$$

Which simplifies to

$$\lambda_1 = 1$$

$$\lambda_2 = -1$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(1)x} + c_2 e^{(-1)x}$$

Or

$$y = c_1 e^x + c_2 e^{-x}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^x + c_2 e^{-x} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = \frac{5}{4}$ and $x = 0$ in the above gives

$$\frac{5}{4} = c_1 + c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = c_1 e^x - c_2 e^{-x}$$

substituting $y' = -\frac{3}{4}$ and $x = 0$ in the above gives

$$-\frac{3}{4} = c_1 - c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{1}{4}$$

$$c_2 = 1$$

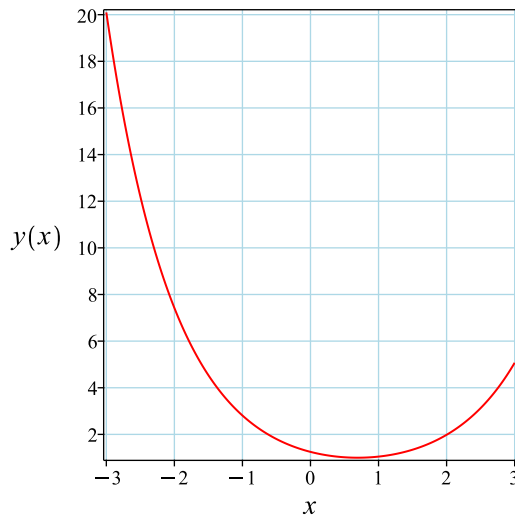
Substituting these values back in above solution results in

$$y = \frac{e^x}{4} + e^{-x}$$

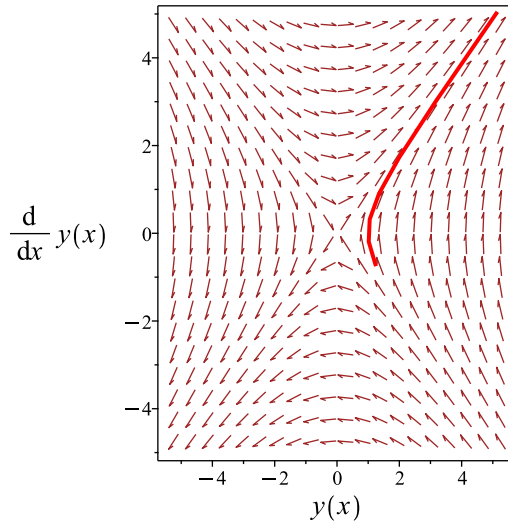
Summary

The solution(s) found are the following

$$y = \frac{e^x}{4} + e^{-x} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^x}{4} + e^{-x}$$

Verified OK.

7.17.3 Solving as second order ode can be made integrable ode

Multiplying the ode by y' gives

$$y'y'' - yy' = 0$$

Integrating the above w.r.t x gives

$$\int (y'y'' - yy') dx = 0$$
$$\frac{y'^2}{2} - \frac{y^2}{2} = c_2$$

Which is now solved for y . Solving the given ode for y' results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$y' = \sqrt{y^2 + 2c_1} \quad (1)$$

$$y' = -\sqrt{y^2 + 2c_1} \quad (2)$$

Now each one of the above ODE is solved.

Solving equation (1)

Integrating both sides gives

$$\int \frac{1}{\sqrt{y^2 + 2c_1}} dy = \int dx$$

$$\ln(y + \sqrt{y^2 + 2c_1}) = x + c_2$$

Raising both side to exponential gives

$$y + \sqrt{y^2 + 2c_1} = e^{x+c_2}$$

Which simplifies to

$$y + \sqrt{y^2 + 2c_1} = c_3 e^x$$

Solving equation (2)

Integrating both sides gives

$$\int -\frac{1}{\sqrt{y^2 + 2c_1}} dy = \int dx$$

$$-\ln(y + \sqrt{y^2 + 2c_1}) = x + c_4$$

Raising both side to exponential gives

$$\frac{1}{y + \sqrt{y^2 + 2c_1}} = e^{x+c_4}$$

Which simplifies to

$$\frac{1}{y + \sqrt{y^2 + 2c_1}} = c_5 e^x$$

Initial conditions are used to solve for the constants of integration.

Looking at the First solution

$$y = \frac{(e^{2x}c_3^2 - 2c_1)e^{-x}}{2c_3} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = \frac{5}{4}$ and $x = 0$ in the above gives

$$\frac{5}{4} = \frac{c_3^2 - 2c_1}{2c_3} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -\frac{(e^{2x}c_3^2 - 2c_1)e^{-x}}{2c_3} + c_3e^x$$

substituting $y' = -\frac{3}{4}$ and $x = 0$ in the above gives

$$-\frac{3}{4} = \frac{c_3^2 + 2c_1}{2c_3} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_3\}$. Solving for the constants gives

$$c_1 = -\frac{1}{2}$$

$$c_3 = \frac{1}{2}$$

Substituting these values back in above solution results in

$$y = \frac{e^{-x}(4 + e^{2x})}{4}$$

Which simplifies to

$$y = \frac{e^x}{4} + e^{-x}$$

Looking at the Second solution

$$y = -\frac{(2c_1c_5^2e^{2x} - 1)e^{-x}}{2c_5} \quad (2)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = \frac{5}{4}$ and $x = 0$ in the above gives

$$\frac{5}{4} = \frac{-2c_1c_5^2 + 1}{2c_5} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -2c_1c_5e^x + \frac{(2c_1c_5^2e^{2x} - 1)e^{-x}}{2c_5}$$

substituting $y' = -\frac{3}{4}$ and $x = 0$ in the above gives

$$-\frac{3}{4} = \frac{-2c_1c_5^2 - 1}{2c_5} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_5\}$. Solving for the constants gives

$$c_1 = -\frac{1}{2}$$

$$c_5 = \frac{1}{2}$$

Substituting these values back in above solution results in

$$y = \frac{e^{-x}(4 + e^{2x})}{4}$$

Which simplifies to

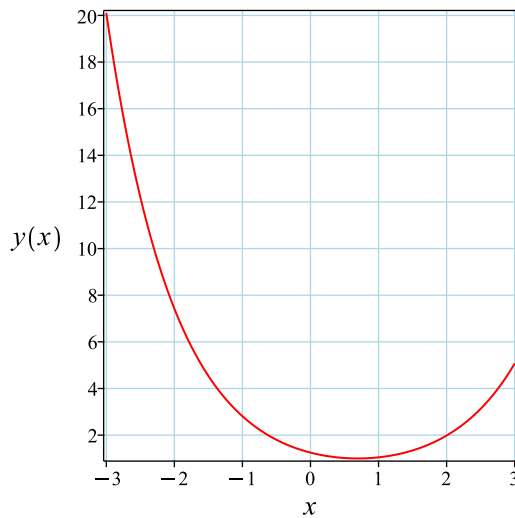
$$y = \frac{e^x}{4} + e^{-x}$$

Summary

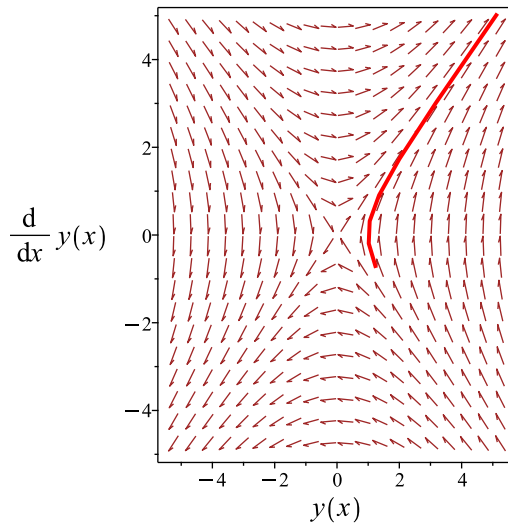
The solution(s) found are the following

$$y = \frac{e^x}{4} + e^{-x} \quad (1)$$

$$y = \frac{e^x}{4} + e^{-x} \quad (2)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^x}{4} + e^{-x}$$

Verified OK.

$$y = \frac{e^x}{4} + e^{-x}$$

Verified OK.

7.17.4 Solving using Kovacic algorithm

Writing the ode as

$$y'' - y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 0 \tag{3}$$

$$C = -1$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} \tag{5}$$
$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 1$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 348: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-x}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned} y_1 &= z_1 \\ &= e^{-x} \end{aligned}$$

Which simplifies to

$$y_1 = e^{-x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Since $B = 0$ then the above becomes

$$\begin{aligned} y_2 &= y_1 \int \frac{1}{y_1^2} dx \\ &= e^{-x} \int \frac{1}{e^{-2x}} dx \\ &= e^{-x} \left(\frac{e^{2x}}{2} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-x}) + c_2 \left(e^{-x} \left(\frac{e^{2x}}{2} \right) \right) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-x} + \frac{c_2 e^x}{2} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = \frac{5}{4}$ and $x = 0$ in the above gives

$$\frac{5}{4} = c_1 + \frac{c_2}{2} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -c_1 e^{-x} + \frac{c_2 e^x}{2}$$

substituting $y' = -\frac{3}{4}$ and $x = 0$ in the above gives

$$-\frac{3}{4} = -c_1 + \frac{c_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 1$$

$$c_2 = \frac{1}{2}$$

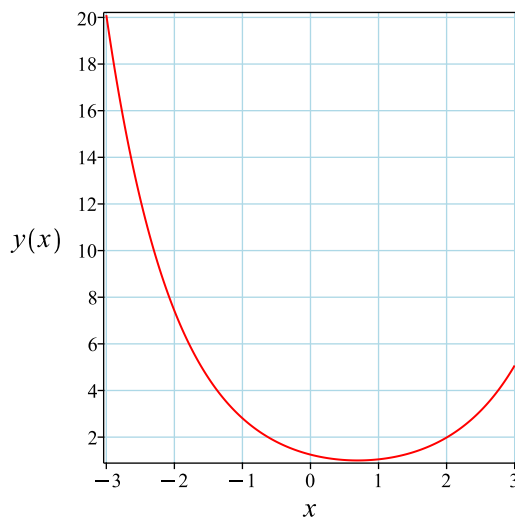
Substituting these values back in above solution results in

$$y = \frac{e^x}{4} + e^{-x}$$

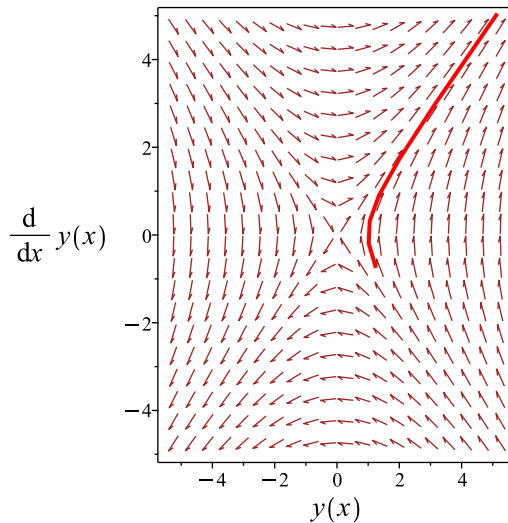
Summary

The solution(s) found are the following

$$y = \frac{e^x}{4} + e^{-x} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^x}{4} + e^{-x}$$

Verified OK.

7.17.5 Maple step by step solution

Let's solve

$$\left[y'' - y = 0, y(0) = \frac{5}{4}, y' \Big|_{\{x=0\}} = -\frac{3}{4} \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 - 1 = 0$$

- Factor the characteristic polynomial

$$(r - 1)(r + 1) = 0$$

- Roots of the characteristic polynomial

$$r = (-1, 1)$$

- 1st solution of the ODE

$$y_1(x) = e^{-x}$$

- 2nd solution of the ODE

$$y_2(x) = e^x$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-x} + c_2 e^x$$

- Check validity of solution $y = c_1 e^{-x} + c_2 e^x$

- Use initial condition $y(0) = \frac{5}{4}$

$$\frac{5}{4} = c_1 + c_2$$

- Compute derivative of the solution

$$y' = -c_1 e^{-x} + c_2 e^x$$

- Use the initial condition $y' \Big|_{\{x=0\}} = -\frac{3}{4}$

$$-\frac{3}{4} = -c_1 + c_2$$
- Solve for c_1 and c_2

$$\{c_1 = 1, c_2 = \frac{1}{4}\}$$
- Substitute constant values into general solution and simplify
$$y = \frac{e^x}{4} + e^{-x}$$
- Solution to the IVP
$$y = \frac{e^x}{4} + e^{-x}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 13

```
dsolve([diff(y(x),x$2) -y(x) = 0,y(0) = 5/4, D(y)(0) = -3/4],y(x), singsol=all)
```

$$y(x) = \frac{e^x}{4} + e^{-x}$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 18

```
DSolve[{y'[x]-y[x]==0,{y[0]==5/4,y'[0]==-3/4}},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x} + \frac{e^x}{4}$$

7.18 problem 20

7.18.1 Existence and uniqueness analysis	1944
7.18.2 Solving as second order linear constant coeff ode	1945
7.18.3 Solving using Kovacic algorithm	1947
7.18.4 Maple step by step solution	1951

Internal problem ID [616]

Internal file name [OUTPUT/616_Sunday_June_05_2022_01_45_54_AM_32300670/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 20.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$2y'' - 3y' + y = 0$$

With initial conditions

$$\left[y(0) = 2, y'(0) = \frac{1}{2} \right]$$

7.18.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = -\frac{3}{2}$$
$$q(x) = \frac{1}{2}$$
$$F = 0$$

Hence the ode is

$$y'' - \frac{3y'}{2} + \frac{y}{2} = 0$$

The domain of $p(x) = -\frac{3}{2}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = \frac{1}{2}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.18.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 2, B = -3, C = 1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$2\lambda^2 e^{\lambda x} - 3\lambda e^{\lambda x} + e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$2\lambda^2 - 3\lambda + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 2, B = -3, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{3}{(2)(2)} \pm \frac{1}{(2)(2)} \sqrt{-3^2 - (4)(2)(1)} \\ &= \frac{3}{4} \pm \frac{1}{4} \end{aligned}$$

Hence

$$\lambda_1 = \frac{3}{4} + \frac{1}{4}$$

$$\lambda_2 = \frac{3}{4} - \frac{1}{4}$$

Which simplifies to

$$\lambda_1 = 1$$

$$\lambda_2 = \frac{1}{2}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(1)x} + c_2 e^{(\frac{1}{2})x}$$

Or

$$y = c_1 e^x + c_2 e^{\frac{x}{2}}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^x + c_2 e^{\frac{x}{2}} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = c_1 + c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = c_1 e^x + \frac{c_2 e^{\frac{x}{2}}}{2}$$

substituting $y' = \frac{1}{2}$ and $x = 0$ in the above gives

$$\frac{1}{2} = c_1 + \frac{c_2}{2} \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -1$$

$$c_2 = 3$$

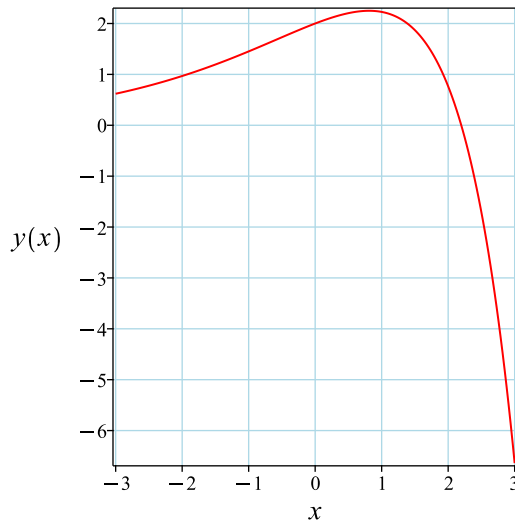
Substituting these values back in above solution results in

$$y = -e^x + 3e^{\frac{x}{2}}$$

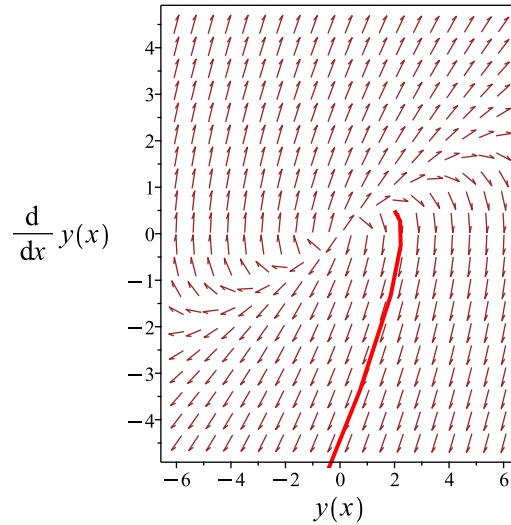
Summary

The solution(s) found are the following

$$y = -e^x + 3e^{\frac{x}{2}} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -e^x + 3e^{\frac{x}{2}}$$

Verified OK.

7.18.3 Solving using Kovacic algorithm

Writing the ode as

$$2y'' - 3y' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 2$$

$$B = -3 \quad (3)$$

$$C = 1$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{16} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 16 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{z(x)}{16} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 350: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{16}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x}{4}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-3}{2} dx} \\ &= z_1 e^{\frac{3x}{4}} \\ &= z_1 \left(e^{\frac{3x}{4}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{\frac{x}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-3}{2} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{\frac{3x}{2}}}{(y_1)^2} dx \\ &= y_1 (2 e^{\frac{x}{2}}) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{\frac{x}{2}} \right) + c_2 \left(e^{\frac{x}{2}} \left(2 e^{\frac{x}{2}} \right) \right)\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\frac{x}{2}} + 2c_2 e^x \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = c_1 + 2c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{c_1 e^{\frac{x}{2}}}{2} + 2c_2 e^x$$

substituting $y' = \frac{1}{2}$ and $x = 0$ in the above gives

$$\frac{1}{2} = \frac{c_1}{2} + 2c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}c_1 &= 3 \\ c_2 &= -\frac{1}{2}\end{aligned}$$

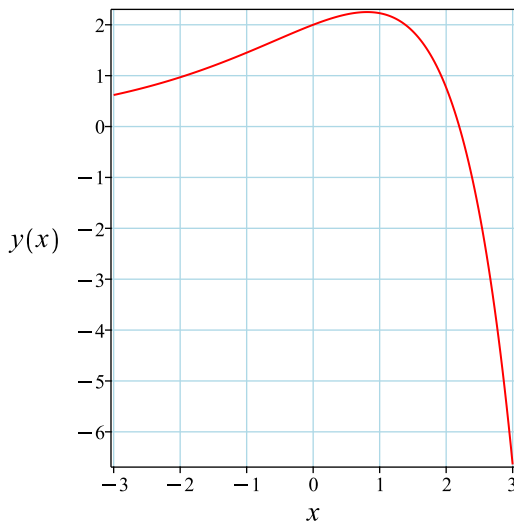
Substituting these values back in above solution results in

$$y = -e^x + 3e^{\frac{x}{2}}$$

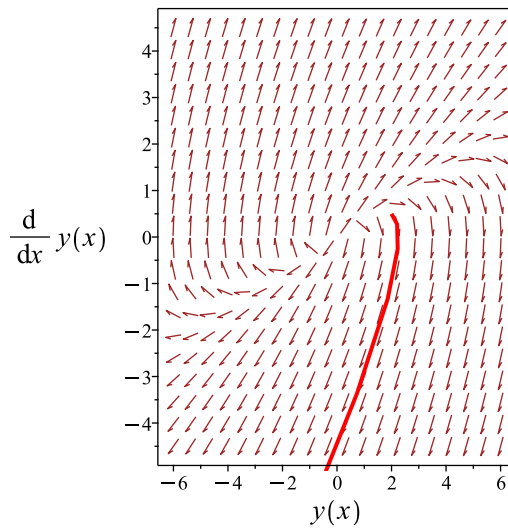
Summary

The solution(s) found are the following

$$y = -e^x + 3e^{\frac{x}{2}} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -e^x + 3e^{\frac{x}{2}}$$

Verified OK.

7.18.4 Maple step by step solution

Let's solve

$$\left[2y'' - 3y' + y = 0, y(0) = 2, y' \Big|_{\{x=0\}} = \frac{1}{2} \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{3y'}{2} - \frac{y}{2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{3y'}{2} + \frac{y}{2} = 0$$

- Characteristic polynomial of ODE

$$r^2 - \frac{3}{2}r + \frac{1}{2} = 0$$

- Factor the characteristic polynomial

$$\frac{(2r-1)(r-1)}{2} = 0$$

- Roots of the characteristic polynomial

$$r = \left(1, \frac{1}{2}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^x$$

- 2nd solution of the ODE

$$y_2(x) = e^{\frac{x}{2}}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^x + c_2 e^{\frac{x}{2}}$$

- Check validity of solution $y = c_1 e^x + c_2 e^{\frac{x}{2}}$

- Use initial condition $y(0) = 2$

$$2 = c_1 + c_2$$

- Compute derivative of the solution

$$y' = c_1 e^x + \frac{c_2 e^{\frac{x}{2}}}{2}$$

- Use the initial condition $y' \Big|_{\{x=0\}} = \frac{1}{2}$

$$\frac{1}{2} = c_1 + \frac{c_2}{2}$$

- Solve for c_1 and c_2

$$\{c_1 = -1, c_2 = 3\}$$

- Substitute constant values into general solution and simplify

$$y = -e^x + 3e^{\frac{x}{2}}$$

- Solution to the IVP

$$y = -e^x + 3e^{\frac{x}{2}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 15

```
dsolve([2*diff(y(x),x$2) -3*diff(y(x),x)+y(x) = 0,y(0) = 2, D(y)(0) = 1/2],y(x), singsol=all
```

$$y(x) = -e^x + 3e^{\frac{x}{2}}$$

✓ Solution by Mathematica

Time used: 0.012 (sec). Leaf size: 20

```
DSolve[{2*y'[x]-3*y'[x]+y[x]==0,{y[0]==2,y'[0]==1/2}},y[x],x,IncludeSingularSolutions -> Tr
```

$$y(x) \rightarrow 3e^{x/2} - e^x$$

7.19 problem 21

7.19.1 Existence and uniqueness analysis	1954
7.19.2 Solving as second order linear constant coeff ode	1955
7.19.3 Solving using Kovacic algorithm	1957
7.19.4 Maple step by step solution	1961

Internal problem ID [617]

Internal file name [OUTPUT/617_Sunday_June_05_2022_01_45_55_AM_95430660/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 21.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - y' - 2y = 0$$

With initial conditions

$$[y(0) = \alpha, y'(0) = 2]$$

7.19.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = -1$$

$$q(x) = -2$$

$$F = 0$$

Hence the ode is

$$y'' - y' - 2y = 0$$

The domain of $p(x) = -1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = -2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.19.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -1, C = -2$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - \lambda e^{\lambda x} - 2 e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - \lambda - 2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -1, C = -2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-1^2 - (4)(1)(-2)} \\ &= \frac{1}{2} \pm \frac{3}{2} \end{aligned}$$

Hence

$$\lambda_1 = \frac{1}{2} + \frac{3}{2}$$

$$\lambda_2 = \frac{1}{2} - \frac{3}{2}$$

Which simplifies to

$$\lambda_1 = 2$$

$$\lambda_2 = -1$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(2)x} + c_2 e^{(-1)x}$$

Or

$$y = c_1 e^{2x} + c_2 e^{-x}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{2x} + c_2 e^{-x} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = \alpha$ and $x = 0$ in the above gives

$$\alpha = c_1 + c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = 2c_1 e^{2x} - c_2 e^{-x}$$

substituting $y' = 2$ and $x = 0$ in the above gives

$$2 = 2c_1 - c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{\alpha}{3} + \frac{2}{3}$$
$$c_2 = -\frac{2}{3} + \frac{2\alpha}{3}$$

Substituting these values back in above solution results in

$$y = \frac{e^{2x}\alpha}{3} + \frac{2e^{2x}}{3} - \frac{2e^{-x}}{3} + \frac{2e^{-x}\alpha}{3}$$

Which simplifies to

$$y = \frac{(2\alpha - 2) e^{-x}}{3} + \frac{e^{2x}(\alpha + 2)}{3}$$

Summary

The solution(s) found are the following

$$y = \frac{(2\alpha - 2) e^{-x}}{3} + \frac{e^{2x}(\alpha + 2)}{3} \quad (1)$$

Verification of solutions

$$y = \frac{(2\alpha - 2) e^{-x}}{3} + \frac{e^{2x}(\alpha + 2)}{3}$$

Verified OK.

7.19.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' - y' - 2y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= -1 \\ C &= -2 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{9}{4} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 9 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{9z(x)}{4} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 352: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{9}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{3x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-1}{1} dx} \\ &= z_1 e^{\frac{x}{2}} \\ &= z_1 (e^{\frac{x}{2}}) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-1}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^x}{(y_1)^2} dx \\ &= y_1 \left(\frac{e^{3x}}{3} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 (e^{-x}) + c_2 \left(e^{-x} \left(\frac{e^{3x}}{3} \right) \right)
\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-x} + \frac{c_2 e^{2x}}{3} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = \alpha$ and $x = 0$ in the above gives

$$\alpha = c_1 + \frac{c_2}{3} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -c_1 e^{-x} + \frac{2c_2 e^{2x}}{3}$$

substituting $y' = 2$ and $x = 0$ in the above gives

$$2 = -c_1 + \frac{2c_2}{3} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}
c_1 &= -\frac{2}{3} + \frac{2\alpha}{3} \\
c_2 &= \alpha + 2
\end{aligned}$$

Substituting these values back in above solution results in

$$y = \frac{e^{2x}\alpha}{3} + \frac{2e^{2x}}{3} - \frac{2e^{-x}}{3} + \frac{2e^{-x}\alpha}{3}$$

Which simplifies to

$$y = \frac{(2\alpha - 2)e^{-x}}{3} + \frac{e^{2x}(\alpha + 2)}{3}$$

Summary

The solution(s) found are the following

$$y = \frac{(2\alpha - 2)e^{-x}}{3} + \frac{e^{2x}(\alpha + 2)}{3} \quad (1)$$

Verification of solutions

$$y = \frac{(2\alpha - 2)e^{-x}}{3} + \frac{e^{2x}(\alpha + 2)}{3}$$

Verified OK.

7.19.4 Maple step by step solution

Let's solve

$$\left[y'' - y' - 2y = 0, y(0) = \alpha, y' \Big|_{\{x=0\}} = 2 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 - r - 2 = 0$$

- Factor the characteristic polynomial

$$(r + 1)(r - 2) = 0$$

- Roots of the characteristic polynomial

$$r = (-1, 2)$$

- 1st solution of the ODE

$$y_1(x) = e^{-x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{2x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-x} + c_2 e^{2x}$$

- Check validity of solution $y = c_1 e^{-x} + c_2 e^{2x}$

- Use initial condition $y(0) = \alpha$

$$\alpha = c_1 + c_2$$

- Compute derivative of the solution

$$y' = -c_1 e^{-x} + 2c_2 e^{2x}$$

- Use the initial condition $y' \Big|_{\{x=0\}} = 2$

$$2 = -c_1 + 2c_2$$

- Solve for c_1 and c_2

$$\left\{ c_1 = -\frac{2}{3} + \frac{2\alpha}{3}, c_2 = \frac{\alpha}{3} + \frac{2}{3} \right\}$$

- Substitute constant values into general solution and simplify

$$y = \frac{(2\alpha-2)e^{-x}}{3} + \frac{e^{2x}(\alpha+2)}{3}$$

- Solution to the IVP

$$y = \frac{(2\alpha-2)e^{-x}}{3} + \frac{e^{2x}(\alpha+2)}{3}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 25

```
dsolve([diff(y(x),x$2) -diff(y(x),x)-2*y(x) = 0,y(0) = alpha, D(y)(0) = 2],y(x), singsol=all
```

$$y(x) = \frac{(2\alpha - 2)e^{-x}}{3} + \frac{e^{2x}(\alpha + 2)}{3}$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 29

```
DSolve[{y''[x]-y'[x]-2*y[x]==0,{y[0]==\[Alpha],y'[0]==2}},y[x],x,IncludeSingularSolutions ->
```

$$y(x) \rightarrow \frac{1}{3}e^{-x}(2(\alpha - 1) + (\alpha + 2)e^{3x})$$

7.20 problem 22

7.20.1 Existence and uniqueness analysis	1963
7.20.2 Solving as second order linear constant coeff ode	1964
7.20.3 Solving as second order ode can be made integrable ode	1966
7.20.4 Solving using Kovacic algorithm	1969
7.20.5 Maple step by step solution	1973

Internal problem ID [618]

Internal file name [OUTPUT/618_Sunday_June_05_2022_01_45_56_AM_92637636/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 22.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "second_order_ode_can_be_made_integrable"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$4y'' - y = 0$$

With initial conditions

$$[y(0) = 2, y'(0) = \beta]$$

7.20.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 0$$

$$q(x) = -\frac{1}{4}$$

$$F = 0$$

Hence the ode is

$$y'' - \frac{y}{4} = 0$$

The domain of $p(x) = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = -\frac{1}{4}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.20.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 4, B = 0, C = -1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$4\lambda^2 e^{\lambda x} - e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$4\lambda^2 - 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 4, B = 0, C = -1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(4)} \pm \frac{1}{(2)(4)} \sqrt{0^2 - (4)(4)(-1)} \\ &= \pm \frac{1}{2} \end{aligned}$$

Hence

$$\lambda_1 = +\frac{1}{2}$$

$$\lambda_2 = -\frac{1}{2}$$

Which simplifies to

$$\lambda_1 = \frac{1}{2}$$

$$\lambda_2 = -\frac{1}{2}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(\frac{1}{2})x} + c_2 e^{(-\frac{1}{2})x}$$

Or

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{-\frac{x}{2}}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{-\frac{x}{2}} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = c_1 + c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{c_1 e^{\frac{x}{2}}}{2} - \frac{c_2 e^{-\frac{x}{2}}}{2}$$

substituting $y' = \beta$ and $x = 0$ in the above gives

$$\beta = \frac{c_1}{2} - \frac{c_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 1 + \beta$$

$$c_2 = -\beta + 1$$

Substituting these values back in above solution results in

$$y = e^{\frac{x}{2}}\beta - e^{-\frac{x}{2}}\beta + e^{\frac{x}{2}} + e^{-\frac{x}{2}}$$

Which simplifies to

$$y = (1 + \beta) e^{\frac{x}{2}} - (\beta - 1) e^{-\frac{x}{2}}$$

Summary

The solution(s) found are the following

$$y = (1 + \beta) e^{\frac{x}{2}} - (\beta - 1) e^{-\frac{x}{2}} \quad (1)$$

Verification of solutions

$$y = (1 + \beta) e^{\frac{x}{2}} - (\beta - 1) e^{-\frac{x}{2}}$$

Verified OK.

7.20.3 Solving as second order ode can be made integrable ode

Multiplying the ode by y' gives

$$4y'y'' - yy' = 0$$

Integrating the above w.r.t x gives

$$\int (4y'y'' - yy') dx = 0$$
$$2y'^2 - \frac{y^2}{2} = c_2$$

Which is now solved for y . Solving the given ode for y' results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$y' = \frac{\sqrt{y^2 + 2c_1}}{2} \quad (1)$$

$$y' = -\frac{\sqrt{y^2 + 2c_1}}{2} \quad (2)$$

Now each one of the above ODE is solved.

Solving equation (1)

Integrating both sides gives

$$\int \frac{2}{\sqrt{y^2 + 2c_1}} dy = \int dx$$
$$2 \ln \left(y + \sqrt{y^2 + 2c_1} \right) = x + c_2$$

Raising both side to exponential gives

$$\left(y + \sqrt{y^2 + 2c_1}\right)^2 = e^{x+c_2}$$

Which simplifies to

$$\left(y + \sqrt{y^2 + 2c_1}\right)^2 = c_3 e^x$$

Solving equation (2)

Integrating both sides gives

$$\int -\frac{2}{\sqrt{y^2 + 2c_1}} dy = \int dx$$
$$-2 \ln \left(y + \sqrt{y^2 + 2c_1}\right) = x + c_4$$

Raising both side to exponential gives

$$\frac{1}{\left(y + \sqrt{y^2 + 2c_1}\right)^2} = e^{x+c_4}$$

Which simplifies to

$$\frac{1}{\left(y + \sqrt{y^2 + 2c_1}\right)^2} = c_5 e^x$$

Initial conditions are used to solve for the constants of integration.

Looking at the First solution

$$y = \frac{c_3 e^x - 2c_1}{2\sqrt{c_3 e^x}} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = \frac{c_3 - 2c_1}{2\sqrt{c_3}} \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{c_3 e^x}{2\sqrt{c_3 e^x}} - \frac{(c_3 e^x - 2c_1) c_3 e^x}{4(c_3 e^x)^{\frac{3}{2}}}$$

substituting $y' = \beta$ and $x = 0$ in the above gives

$$\beta = \frac{c_3 + 2c_1}{4\sqrt{c_3}} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_3\}$. Solving for the constants gives

$$\begin{aligned} c_1 &= -4 \operatorname{csgn}(1 + \beta) \beta + 2\beta^2 - 4 \operatorname{csgn}(1 + \beta) + 4\beta + 2 \\ c_3 &= 4(1 + \beta)^2 \end{aligned}$$

Substituting these values back in above solution results in

$$y = \frac{e^x \beta^2 + 2 \operatorname{csgn}(1 + \beta) \beta + 2 e^x \beta - \beta^2 + 2 \operatorname{csgn}(1 + \beta) + e^x - 2\beta - 1}{\sqrt{e^x \beta^2 + 2 e^x \beta + e^x}}$$

Which simplifies to

$$y = \frac{(1 + \beta) (2 \operatorname{csgn}(1 + \beta) + (e^x - 1) (1 + \beta))}{\sqrt{(1 + \beta)^2 e^x}}$$

Looking at the Second solution

$$y = -\frac{2c_1 c_5 e^x - 1}{2\sqrt{c_5 e^x}} \quad (2)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = \frac{-2c_1 c_5 + 1}{2\sqrt{c_5}} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -\frac{c_5 e^x c_1}{\sqrt{c_5 e^x}} + \frac{(2c_1 c_5 e^x - 1) c_5 e^x}{4 (c_5 e^x)^{\frac{3}{2}}}$$

substituting $y' = \beta$ and $x = 0$ in the above gives

$$\beta = \frac{-2c_1 c_5 - 1}{4\sqrt{c_5}} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_5\}$. Solving for the constants gives

$$\begin{aligned} c_1 &= 2\beta^2 - 2 \\ c_5 &= \frac{1}{4(\beta - 1)^2} \end{aligned}$$

Substituting these values back in above solution results in

$$y = \frac{-e^x \beta - e^x + \beta - 1}{\sqrt{\frac{e^x}{\beta^2 - 2\beta + 1}} \beta - \sqrt{\frac{e^x}{\beta^2 - 2\beta + 1}}}$$

Which simplifies to

$$y = \frac{(-1 - \beta) e^x + \beta - 1}{\sqrt{\frac{e^x}{(\beta - 1)^2}} (\beta - 1)}$$

Summary

The solution(s) found are the following

$$y = \frac{(1 + \beta) (2 \operatorname{csgn}(1 + \beta) + (e^x - 1) (1 + \beta))}{\sqrt{(1 + \beta)^2 e^x}} \quad (1)$$

$$y = \frac{(-1 - \beta) e^x + \beta - 1}{\sqrt{\frac{e^x}{(\beta - 1)^2}} (\beta - 1)} \quad (2)$$

Verification of solutions

$$y = \frac{(1 + \beta) (2 \operatorname{csgn}(1 + \beta) + (e^x - 1) (1 + \beta))}{\sqrt{(1 + \beta)^2 e^x}}$$

Verified OK.

$$y = \frac{(-1 - \beta) e^x + \beta - 1}{\sqrt{\frac{e^x}{(\beta - 1)^2}} (\beta - 1)}$$

Verified OK.

7.20.4 Solving using Kovacic algorithm

Writing the ode as

$$4y'' - y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 4 \\ B &= 0 \\ C &= -1 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{4} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{z(x)}{4} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 354: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}
 y_1 &= z_1 \\
 &= e^{-\frac{x}{2}}
 \end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{x}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Since $B = 0$ then the above becomes

$$\begin{aligned} y_2 &= y_1 \int \frac{1}{y_1^2} dx \\ &= e^{-\frac{x}{2}} \int \frac{1}{e^{-x}} dx \\ &= e^{-\frac{x}{2}} (e^x) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-\frac{x}{2}}) + c_2 (e^{-\frac{x}{2}} (e^x)) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-\frac{x}{2}} + c_2 e^{\frac{x}{2}} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = c_1 + c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = -\frac{c_1 e^{-\frac{x}{2}}}{2} + \frac{c_2 e^{\frac{x}{2}}}{2}$$

substituting $y' = \beta$ and $x = 0$ in the above gives

$$\beta = -\frac{c_1}{2} + \frac{c_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -\beta + 1$$

$$c_2 = 1 + \beta$$

Substituting these values back in above solution results in

$$y = e^{\frac{x}{2}}\beta - e^{-\frac{x}{2}}\beta + e^{\frac{x}{2}} + e^{-\frac{x}{2}}$$

Which simplifies to

$$y = (1 + \beta) e^{\frac{x}{2}} - (\beta - 1) e^{-\frac{x}{2}}$$

Summary

The solution(s) found are the following

$$y = (1 + \beta) e^{\frac{x}{2}} - (\beta - 1) e^{-\frac{x}{2}} \quad (1)$$

Verification of solutions

$$y = (1 + \beta) e^{\frac{x}{2}} - (\beta - 1) e^{-\frac{x}{2}}$$

Verified OK.

7.20.5 Maple step by step solution

Let's solve

$$\left[4y'' - y = 0, y(0) = 2, y' \Big|_{\{x=0\}} = \beta \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{y}{4}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{y}{4} = 0$$

- Characteristic polynomial of ODE

$$r^2 - \frac{1}{4} = 0$$

- Factor the characteristic polynomial

$$\frac{(2r-1)(2r+1)}{4} = 0$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{2}, \frac{1}{2}\right)$$
- 1st solution of the ODE

$$y_1(x) = e^{-\frac{x}{2}}$$
- 2nd solution of the ODE

$$y_2(x) = e^{\frac{x}{2}}$$
- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$
- Substitute in solutions

$$y = c_1 e^{-\frac{x}{2}} + c_2 e^{\frac{x}{2}}$$
- Check validity of solution $y = c_1 e^{-\frac{x}{2}} + c_2 e^{\frac{x}{2}}$
 - Use initial condition $y(0) = 2$

$$2 = c_1 + c_2$$
 - Compute derivative of the solution

$$y' = -\frac{c_1 e^{-\frac{x}{2}}}{2} + \frac{c_2 e^{\frac{x}{2}}}{2}$$
 - Use the initial condition $y' \Big|_{\{x=0\}} = \beta$

$$\beta = -\frac{c_1}{2} + \frac{c_2}{2}$$
 - Solve for c_1 and c_2

$$\{c_1 = -\beta + 1, c_2 = 1 + \beta\}$$
 - Substitute constant values into general solution and simplify

$$y = (1 + \beta) e^{\frac{x}{2}} - (\beta - 1) e^{-\frac{x}{2}}$$
- Solution to the IVP

$$y = (1 + \beta) e^{\frac{x}{2}} - (\beta - 1) e^{-\frac{x}{2}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 22

```
dsolve([4*diff(y(x),x$2) -y(x) = 0,y(0) = 2, D(y)(0) = beta],y(x), singsol=all)
```

$$y(x) = (1 + \beta) e^{\frac{x}{2}} - (\beta - 1) e^{-\frac{x}{2}}$$

✓ Solution by Mathematica

Time used: 0.012 (sec). Leaf size: 25

```
DSolve[{4*y'[x]-y[x]==0,{y[0]==2,y'[0]==\[Beta]}],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x/2}(-\beta + (\beta + 1)e^x + 1)$$

7.21 problem 23

7.21.1 Solving as second order linear constant coeff ode	1976
7.21.2 Solving using Kovacic algorithm	1978
7.21.3 Maple step by step solution	1981

Internal problem ID [619]

Internal file name [OUTPUT/619_Sunday_June_05_2022_01_45_56_AM_20917638/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 23.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - (2\alpha - 1)y' + \alpha(\alpha - 1)y = 0$$

7.21.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -2\alpha + 1, C = \alpha^2 - \alpha$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + (-2\alpha + 1)\lambda e^{\lambda x} + (\alpha^2 - \alpha)e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + (-2\alpha + 1)\lambda + \alpha^2 - \alpha = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -2\alpha + 1, C = \alpha^2 - \alpha$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{2\alpha - 1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-2\alpha + 1^2 - (4)(1)(\alpha^2 - \alpha)} \\ &= \alpha - \frac{1}{2} \pm \frac{1}{2} \end{aligned}$$

Hence

$$\lambda_1 = \alpha - \frac{1}{2} + \frac{1}{2}$$

$$\lambda_2 = \alpha - \frac{1}{2} - \frac{1}{2}$$

Which simplifies to

$$\lambda_1 = \alpha$$

$$\lambda_2 = \alpha - 1$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(\alpha)x} + c_2 e^{(\alpha-1)x}$$

Or

$$y = c_1 e^{\alpha x} + c_2 e^{(\alpha-1)x}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{\alpha x} + c_2 e^{(\alpha-1)x} \tag{1}$$

Verification of solutions

$$y = c_1 e^{\alpha x} + c_2 e^{(\alpha-1)x}$$

Verified OK.

7.21.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + (-2\alpha + 1)y' + (\alpha^2 - \alpha)y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= -2\alpha + 1 \\ C &= \alpha^2 - \alpha \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{4} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{z(x)}{4} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 356: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-2\alpha+1}{1} dx} \end{aligned}$$

$$\begin{aligned}
&= z_1 e^{(\alpha - \frac{1}{2})x} \\
&= z_1 \left(e^{(\alpha - \frac{1}{2})x} \right)
\end{aligned}$$

Which simplifies to

$$y_1 = e^{(\alpha - 1)x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}
y_2 &= y_1 \int \frac{e^{\int -\frac{-2\alpha + 1}{1} dx}}{(y_1)^2} dx \\
&= y_1 \int \frac{e^{(2\alpha - 1)x}}{(y_1)^2} dx \\
&= y_1 (e^x)
\end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 (e^{(\alpha - 1)x}) + c_2 (e^{(\alpha - 1)x} (e^x))
\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{(\alpha - 1)x} + c_2 e^{\alpha x} \tag{1}$$

Verification of solutions

$$y = c_1 e^{(\alpha - 1)x} + c_2 e^{\alpha x}$$

Verified OK.

7.21.3 Maple step by step solution

Let's solve

$$y'' + (-2\alpha + 1)y' + (\alpha^2 - \alpha)y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + (-2\alpha + 1)r + \alpha^2 - \alpha = 0$$

- Factor the characteristic polynomial

$$(\alpha - r)(\alpha - r - 1) = 0$$

- Roots of the characteristic polynomial

$$r = (\alpha, \alpha - 1)$$

- 1st solution of the ODE

$$y_1(x) = e^{\alpha x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{(\alpha-1)x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{\alpha x} + c_2 e^{(\alpha-1)x}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(diff(y(x),x$2) -(2*alpha-1)*diff(y(x),x)+alpha*(alpha-1)*y(x) = 0,y(x), singsol=all)
```

$$y(x) = c_1 e^{\alpha x} + c_2 e^{(\alpha-1)x}$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 24

```
DSolve[y''[x]-(2*[Alpha]-1)*y'[x]+[Alpha]*(\ [Alpha]-1)*y[x]==0,y[x],x,IncludeSingularSolut
```

$$y(x) \rightarrow c_1 e^{(\alpha-1)x} + c_2 e^{\alpha x}$$

7.22 problem 24

7.22.1 Solving as second order linear constant coeff ode	1983
7.22.2 Solving using Kovacic algorithm	1985
7.22.3 Maple step by step solution	1988

Internal problem ID [620]

Internal file name [OUTPUT/620_Sunday_June_05_2022_01_45_57_AM_13631438/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 24.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + (3 - \alpha)y' - 2(\alpha - 1)y = 0$$

7.22.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 3 - \alpha, C = -2\alpha + 2$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + (3 - \alpha)\lambda e^{\lambda x} + (-2\alpha + 2)e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + (3 - \alpha)\lambda - 2\alpha + 2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 3 - \alpha, C = -2\alpha + 2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{\alpha - 3}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{3 - \alpha^2 - (4)(1)(-2\alpha + 2)} \\ &= -\frac{3}{2} + \frac{\alpha}{2} \pm \frac{\sqrt{(\alpha + 1)^2}}{2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= -\frac{3}{2} + \frac{\alpha}{2} + \frac{\sqrt{(\alpha + 1)^2}}{2} \\ \lambda_2 &= -\frac{3}{2} + \frac{\alpha}{2} - \frac{\sqrt{(\alpha + 1)^2}}{2} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= -\frac{3}{2} + \frac{\alpha}{2} + \frac{\sqrt{(\alpha + 1)^2}}{2} \\ \lambda_2 &= -\frac{3}{2} + \frac{\alpha}{2} - \frac{\sqrt{(\alpha + 1)^2}}{2} \end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned} y &= c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x} \\ y &= c_1 e^{\left(-\frac{3}{2} + \frac{\alpha}{2} + \frac{\sqrt{(\alpha + 1)^2}}{2}\right)x} + c_2 e^{\left(-\frac{3}{2} + \frac{\alpha}{2} - \frac{\sqrt{(\alpha + 1)^2}}{2}\right)x} \end{aligned}$$

Or

$$y = c_1 e^{\left(-\frac{3}{2} + \frac{\alpha}{2} + \frac{\sqrt{(\alpha + 1)^2}}{2}\right)x} + c_2 e^{\left(-\frac{3}{2} + \frac{\alpha}{2} - \frac{\sqrt{(\alpha + 1)^2}}{2}\right)x}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{\left(-\frac{3}{2} + \frac{\alpha}{2} + \frac{\sqrt{(\alpha + 1)^2}}{2}\right)x} + c_2 e^{\left(-\frac{3}{2} + \frac{\alpha}{2} - \frac{\sqrt{(\alpha + 1)^2}}{2}\right)x} \quad (1)$$

Verification of solutions

$$y = c_1 e^{\left(-\frac{3}{2} + \frac{\alpha}{2} + \frac{\sqrt{(\alpha+1)^2}}{2}\right)x} + c_2 e^{\left(-\frac{3}{2} + \frac{\alpha}{2} - \frac{\sqrt{(\alpha+1)^2}}{2}\right)x}$$

Verified OK.

7.22.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + (3 - \alpha)y' + (-2\alpha + 2)y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 3 - \alpha \\ C &= -2\alpha + 2 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{\alpha^2 + 2\alpha + 1}{4} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= \alpha^2 + 2\alpha + 1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \left(\frac{1}{4}\alpha^2 + \frac{1}{2}\alpha + \frac{1}{4} \right) z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 358: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{4}\alpha^2 + \frac{1}{2}\alpha + \frac{1}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{\frac{\sqrt{(\alpha+1)^2}x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{3-\alpha}{1} dx} \\ &= z_1 e^{(-\frac{3}{2} + \frac{\alpha}{2})x} \\ &= z_1 \left(e^{\frac{(\alpha-3)x}{2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{\frac{(\text{csgn}(\alpha+1)\alpha + \text{csgn}(\alpha+1) + \alpha - 3)x}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{3-\alpha}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{(\alpha-3)x}}{(y_1)^2} dx \\ &= y_1 \left(-\frac{\text{csgn}(\alpha+1) e^{-\text{csgn}(\alpha+1)(\alpha+1)x}}{\alpha+1} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{\frac{(\text{csgn}(\alpha+1)\alpha + \text{csgn}(\alpha+1) + \alpha - 3)x}{2}} \right) \\ &\quad + c_2 \left(e^{\frac{(\text{csgn}(\alpha+1)\alpha + \text{csgn}(\alpha+1) + \alpha - 3)x}{2}} \left(-\frac{\text{csgn}(\alpha+1) e^{-\text{csgn}(\alpha+1)(\alpha+1)x}}{\alpha+1} \right) \right) \end{aligned}$$

Simplifying the solution $y = c_1 e^{\frac{(\operatorname{csgn}(\alpha+1)\alpha + \operatorname{csgn}(\alpha+1) + \alpha - 3)x}{2}} - \frac{c_2 \operatorname{csgn}(\alpha+1) e^{-\frac{(\operatorname{csgn}(\alpha+1)\alpha + \operatorname{csgn}(\alpha+1) - \alpha + 3)x}{2}}}{\alpha + 1}$

Summary

The solution(s) found are the following
to $y = c_1 e^{\frac{(2\alpha-2)x}{2}} - \frac{c_2 e^{-2x}}{\alpha+1}$

$$y = c_1 e^{\frac{(2\alpha-2)x}{2}} - \frac{c_2 e^{-2x}}{\alpha + 1}$$

Verification of solutions

$$y = c_1 e^{\frac{(2\alpha-2)x}{2}} - \frac{c_2 e^{-2x}}{\alpha + 1}$$

Verified OK.

7.22.3 Maple step by step solution

Let's solve

$$y'' + (3 - \alpha)y' + (-2\alpha + 2)y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + (3 - \alpha)r - 2\alpha + 2 = 0$$

- Factor the characteristic polynomial

$$-(r + 2)(-r + \alpha - 1) = 0$$

- Roots of the characteristic polynomial

$$r = (-2, \alpha - 1)$$

- 1st solution of the ODE

$$y_1(x) = e^{-2x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{(\alpha-1)x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-2x} + c_2 e^{(\alpha-1)x}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 19

```
dsolve(diff(y(x),x$2) +(3-alpha)*diff(y(x),x)-2*(alpha-1)*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-2x}c_1 + c_2e^{(\alpha-1)x}$$

✓ Solution by Mathematica

Time used: 0.014 (sec). Leaf size: 24

```
DSolve[y''[x]+(3-\[Alpha])*y'[x]-2*(\[Alpha]-1)*y[x]==0,y[x],x,IncludeSingularSolutions -> T
```

$$y(x) \rightarrow e^{-2x}(c_1e^{\alpha x+x} + c_2)$$

7.23 problem 25

7.23.1 Existence and uniqueness analysis	1990
7.23.2 Solving as second order linear constant coeff ode	1991
7.23.3 Solving using Kovacic algorithm	1993
7.23.4 Maple step by step solution	1997

Internal problem ID [621]

Internal file name [OUTPUT/621_Sunday_June_05_2022_01_45_58_AM_931008/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 25.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$2y'' + 3y' - 2y = 0$$

With initial conditions

$$[y(0) = 1, y'(0) = -\beta]$$

7.23.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$\begin{aligned} p(x) &= \frac{3}{2} \\ q(x) &= -1 \\ F &= 0 \end{aligned}$$

Hence the ode is

$$y'' + \frac{3y'}{2} - y = 0$$

The domain of $p(x) = \frac{3}{2}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = -1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.23.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 2, B = 3, C = -2$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$2\lambda^2 e^{\lambda x} + 3\lambda e^{\lambda x} - 2e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$2\lambda^2 + 3\lambda - 2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 2, B = 3, C = -2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-3}{(2)(2)} \pm \frac{1}{(2)(2)} \sqrt{3^2 - (4)(2)(-2)} \\ &= -\frac{3}{4} \pm \frac{5}{4} \end{aligned}$$

Hence

$$\lambda_1 = -\frac{3}{4} + \frac{5}{4}$$

$$\lambda_2 = -\frac{3}{4} - \frac{5}{4}$$

Which simplifies to

$$\lambda_1 = \frac{1}{2}$$

$$\lambda_2 = -2$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(\frac{1}{2})x} + c_2 e^{(-2)x}$$

Or

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{-2x}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\frac{x}{2}} + c_2 e^{-2x} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = 0$ in the above gives

$$1 = c_1 + c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{c_1 e^{\frac{x}{2}}}{2} - 2c_2 e^{-2x}$$

substituting $y' = -\beta$ and $x = 0$ in the above gives

$$-\beta = \frac{c_1}{2} - 2c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{4}{5} - \frac{2\beta}{5}$$

$$c_2 = \frac{2\beta}{5} + \frac{1}{5}$$

Substituting these values back in above solution results in

$$y = \frac{4e^{\frac{x}{2}}}{5} - \frac{2e^{\frac{x}{2}}\beta}{5} + \frac{2e^{-2x}\beta}{5} + \frac{e^{-2x}}{5}$$

Which simplifies to

$$y = -\frac{\left(2e^{\frac{5x}{2}}\beta - 4e^{\frac{5x}{2}} - 2\beta - 1\right)e^{-2x}}{5}$$

Summary

The solution(s) found are the following

$$y = -\frac{\left(2e^{\frac{5x}{2}}\beta - 4e^{\frac{5x}{2}} - 2\beta - 1\right)e^{-2x}}{5} \quad (1)$$

Verification of solutions

$$y = -\frac{\left(2e^{\frac{5x}{2}}\beta - 4e^{\frac{5x}{2}} - 2\beta - 1\right)e^{-2x}}{5}$$

Verified OK.

7.23.3 Solving using Kovacic algorithm

Writing the ode as

$$2y'' + 3y' - 2y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 2 \\ B &= 3 \\ C &= -2 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{25}{16} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 25 \\ t &= 16 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{25z(x)}{16} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 360: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{25}{16}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{5x}{4}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{3}{2} dx} \\ &= z_1 e^{-\frac{3x}{4}} \\ &= z_1 \left(e^{-\frac{3x}{4}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-2x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{3}{2} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-\frac{3x}{2}}}{(y_1)^2} dx \\ &= y_1 \left(\frac{2 e^{\frac{5x}{2}}}{5} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 (e^{-2x}) + c_2 \left(e^{-2x} \left(\frac{2e^{\frac{5x}{2}}}{5} \right) \right)
\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-2x} + \frac{2c_2 e^{\frac{x}{2}}}{5} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = 0$ in the above gives

$$1 = c_1 + \frac{2c_2}{5} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -2c_1 e^{-2x} + \frac{c_2 e^{\frac{x}{2}}}{5}$$

substituting $y' = -\beta$ and $x = 0$ in the above gives

$$-\beta = -2c_1 + \frac{c_2}{5} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}
c_1 &= \frac{2\beta}{5} + \frac{1}{5} \\
c_2 &= 2 - \beta
\end{aligned}$$

Substituting these values back in above solution results in

$$y = \frac{4e^{\frac{x}{2}}}{5} - \frac{2e^{\frac{x}{2}}\beta}{5} + \frac{2e^{-2x}\beta}{5} + \frac{e^{-2x}}{5}$$

Which simplifies to

$$y = -\frac{\left(2e^{\frac{5x}{2}}\beta - 4e^{\frac{5x}{2}} - 2\beta - 1\right)e^{-2x}}{5}$$

Summary

The solution(s) found are the following

$$y = -\frac{\left(2e^{\frac{5x}{2}}\beta - 4e^{\frac{5x}{2}} - 2\beta - 1\right)e^{-2x}}{5} \quad (1)$$

Verification of solutions

$$y = -\frac{\left(2e^{\frac{5x}{2}}\beta - 4e^{\frac{5x}{2}} - 2\beta - 1\right)e^{-2x}}{5}$$

Verified OK.

7.23.4 Maple step by step solution

Let's solve

$$\left[2y'' + 3y' - 2y = 0, y(0) = 1, y'|_{\{x=0\}} = -\beta\right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{3y'}{2} + y$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{3y'}{2} - y = 0$$

- Characteristic polynomial of ODE

$$r^2 + \frac{3}{2}r - 1 = 0$$

- Factor the characteristic polynomial

$$\frac{(r+2)(2r-1)}{2} = 0$$

- Roots of the characteristic polynomial

$$r = \left(-2, \frac{1}{2}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-2x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{\frac{x}{2}}$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^{-2x} + c_2e^{\frac{x}{2}}$$

- Check validity of solution $y = c_1e^{-2x} + c_2e^{\frac{x}{2}}$

- Use initial condition $y(0) = 1$

$$1 = c_1 + c_2$$

- Compute derivative of the solution

$$y' = -2c_1e^{-2x} + \frac{c_2e^{\frac{x}{2}}}{2}$$

- Use the initial condition $y'|_{\{x=0\}} = -\beta$

$$-\beta = -2c_1 + \frac{c_2}{2}$$

- Solve for c_1 and c_2

$$\left\{ c_1 = \frac{2\beta}{5} + \frac{1}{5}, c_2 = \frac{4}{5} - \frac{2\beta}{5} \right\}$$

- Substitute constant values into general solution and simplify

$$y = -\frac{(2e^{\frac{5x}{2}}\beta - 4e^{\frac{5x}{2}} - 2\beta - 1)e^{-2x}}{5}$$

- Solution to the IVP

$$y = -\frac{(2e^{\frac{5x}{2}}\beta - 4e^{\frac{5x}{2}} - 2\beta - 1)e^{-2x}}{5}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 28

```
dsolve([2*dif(y(x),x$2) +3*dif(y(x),x)-2*y(x) = 0,y(0) = 1, D(y)(0) = -beta],y(x), singsol
```

$$y(x) = -\frac{(2e^{\frac{5x}{2}}\beta - 4e^{\frac{5x}{2}} - 2\beta - 1)e^{-2x}}{5}$$

✓ Solution by Mathematica

Time used: 0.024 (sec). Leaf size: 67

```
DSolve[{y''[x]+3*y'[x]-2*y[x]==0,{y[0]==1,y'[0]==-\[Beta]}},y[x],x,IncludeSingularSolutions
```

$$y(x) \rightarrow \frac{1}{34} e^{-\frac{1}{2}(3+\sqrt{17})x} \left(2\sqrt{17}\beta + (-2\sqrt{17}\beta + 3\sqrt{17} + 17) e^{\sqrt{17}x} - 3\sqrt{17} + 17 \right)$$

7.24 problem 26

7.24.1 Existence and uniqueness analysis	2000
7.24.2 Solving as second order linear constant coeff ode	2001
7.24.3 Solving using Kovacic algorithm	2003
7.24.4 Maple step by step solution	2007

Internal problem ID [622]

Internal file name [OUTPUT/622_Sunday_June_05_2022_01_45_59_AM_64820801/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.1 Homogeneous Equations with Constant Coefficients, page 144

Problem number: 26.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 5y' + 6y = 0$$

With initial conditions

$$[y(0) = 2, y'(0) = \beta]$$

7.24.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 5$$

$$q(x) = 6$$

$$F = 0$$

Hence the ode is

$$y'' + 5y' + 6y = 0$$

The domain of $p(x) = 5$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = 6$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

7.24.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 5, C = 6$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 5\lambda e^{\lambda x} + 6e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 5\lambda + 6 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 5, C = 6$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-5}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{5^2 - (4)(1)(6)} \\ &= -\frac{5}{2} \pm \frac{1}{2} \end{aligned}$$

Hence

$$\lambda_1 = -\frac{5}{2} + \frac{1}{2}$$

$$\lambda_2 = -\frac{5}{2} - \frac{1}{2}$$

Which simplifies to

$$\lambda_1 = -2$$

$$\lambda_2 = -3$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(-2)x} + c_2 e^{(-3)x}$$

Or

$$y = c_1 e^{-2x} + e^{-3x} c_2$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-2x} + e^{-3x} c_2 \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = c_1 + c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = -2c_1 e^{-2x} - 3e^{-3x} c_2$$

substituting $y' = \beta$ and $x = 0$ in the above gives

$$\beta = -2c_1 - 3c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 6 + \beta$$

$$c_2 = -\beta - 4$$

Substituting these values back in above solution results in

$$y = e^{-2x} \beta - e^{-3x} \beta + 6e^{-2x} - 4e^{-3x}$$

Which simplifies to

$$y = (-\beta - 4) e^{-3x} + (6 + \beta) e^{-2x}$$

Summary

The solution(s) found are the following

$$y = (-\beta - 4) e^{-3x} + (6 + \beta) e^{-2x} \quad (1)$$

Verification of solutions

$$y = (-\beta - 4) e^{-3x} + (6 + \beta) e^{-2x}$$

Verified OK.

7.24.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 5y' + 6y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 5 \quad (3)$$

$$C = 6$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{4} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{z(x)}{4} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 362: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{5}{1} dx} \\ &= z_1 e^{-\frac{5x}{2}} \\ &= z_1 \left(e^{-\frac{5x}{2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-3x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{5}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-5x}}{(y_1)^2} dx \\ &= y_1 (e^x) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-3x}) + c_2 (e^{-3x} (e^x)) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-3x} + c_2 e^{-2x} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = c_1 + c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = -3c_1 e^{-3x} - 2c_2 e^{-2x}$$

substituting $y' = \beta$ and $x = 0$ in the above gives

$$\beta = -3c_1 - 2c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -\beta - 4$$

$$c_2 = 6 + \beta$$

Substituting these values back in above solution results in

$$y = e^{-2x} \beta - e^{-3x} \beta + 6 e^{-2x} - 4 e^{-3x}$$

Which simplifies to

$$y = (-\beta - 4) e^{-3x} + (6 + \beta) e^{-2x}$$

Summary

The solution(s) found are the following

$$y = (-\beta - 4) e^{-3x} + (6 + \beta) e^{-2x} \quad (1)$$

Verification of solutions

$$y = (-\beta - 4) e^{-3x} + (6 + \beta) e^{-2x}$$

Verified OK.

7.24.4 Maple step by step solution

Let's solve

$$\left[y'' + 5y' + 6y = 0, y(0) = 2, y' \Big|_{\{x=0\}} = \beta \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 5r + 6 = 0$$

- Factor the characteristic polynomial

$$(r + 3)(r + 2) = 0$$

- Roots of the characteristic polynomial

$$r = (-3, -2)$$

- 1st solution of the ODE

$$y_1(x) = e^{-3x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{-2x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-3x} + c_2 e^{-2x}$$

- Check validity of solution $y = c_1 e^{-3x} + c_2 e^{-2x}$

- Use initial condition $y(0) = 2$

$$2 = c_1 + c_2$$

- Compute derivative of the solution

$$y' = -3c_1 e^{-3x} - 2c_2 e^{-2x}$$

- Use the initial condition $y' \Big|_{\{x=0\}} = \beta$

$$\beta = -3c_1 - 2c_2$$

- Solve for c_1 and c_2

$$\{c_1 = -\beta - 4, c_2 = 6 + \beta\}$$

- Substitute constant values into general solution and simplify

$$y = (-\beta - 4)e^{-3x} + (6 + \beta)e^{-2x}$$

- Solution to the IVP

$$y = (-\beta - 4)e^{-3x} + (6 + \beta)e^{-2x}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 23

```
dsolve([diff(y(x),x$2) +5*diff(y(x),x)+6*y(x) = 0,y(0) = 2, D(y)(0) = beta],y(x), singsol=all)
```

$$y(x) = e^{-2x}(6 + \beta) + (-\beta - 4)e^{-3x}$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 23

```
DSolve[{y'[x]+5*y'[x]+6*y[x]==0,{y[0]==2,y'[0]==\[Beta]}},y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow e^{-3x}(-\beta + (\beta + 6)e^x - 4)$$

8 Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

8.1	problem 7	2010
8.2	problem 8	2018
8.3	problem 9	2026
8.4	problem 10	2034
8.5	problem 11	2042
8.6	problem 12	2050
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8.1 problem 7

- 8.1.1 Solving as second order linear constant coeff ode 2010
- 8.1.2 Solving using Kovacic algorithm 2012
- 8.1.3 Maple step by step solution 2016

Internal problem ID [623]

Internal file name [OUTPUT/623_Sunday_June_05_2022_01_46_00_AM_10820559/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 7.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - 2y' + 2y = 0$$

8.1.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -2, C = 2$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - 2\lambda e^{\lambda x} + 2e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - 2\lambda + 2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -2, C = 2$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-2^2 - (4)(1)(2)} \\ &= 1 \pm i\end{aligned}$$

Hence

$$\lambda_1 = 1 + i$$

$$\lambda_2 = 1 - i$$

Which simplifies to

$$\lambda_1 = 1 + i$$

$$\lambda_2 = 1 - i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 1$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^x(c_1 \cos(x) + c_2 \sin(x))$$

Summary

The solution(s) found are the following

$$y = e^x(c_1 \cos(x) + c_2 \sin(x)) \quad (1)$$

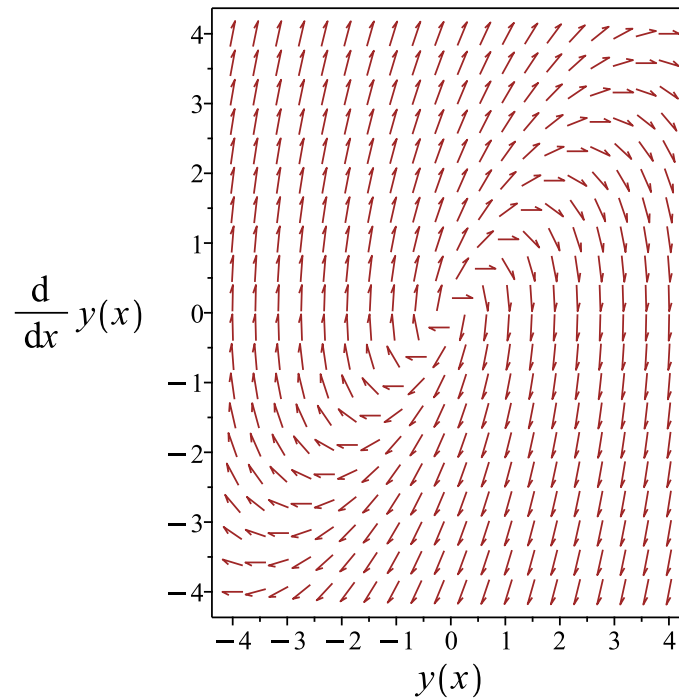


Figure 393: Slope field plot

Verification of solutions

$$y = e^x (c_1 \cos(x) + c_2 \sin(x))$$

Verified OK.

8.1.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 2y' + 2y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= -2 \\ C &= 2 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 364: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-2}{1} dx} \\ &= z_1 e^x \\ &= z_1(e^x)\end{aligned}$$

Which simplifies to

$$y_1 = \cos(x) e^x$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{2x}}{(y_1)^2} dx \\ &= y_1(\tan(x))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1(\cos(x) e^x) + c_2(\cos(x) e^x (\tan(x)))\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(x) e^x + c_2 \sin(x) e^x \quad (1)$$

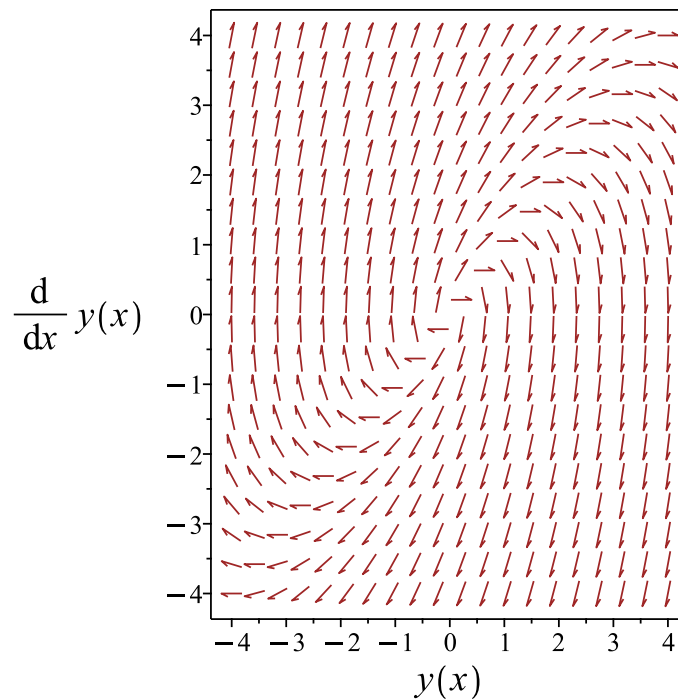


Figure 394: Slope field plot

Verification of solutions

$$y = c_1 \cos(x) e^x + c_2 \sin(x) e^x$$

Verified OK.

8.1.3 Maple step by step solution

Let's solve

$$y'' - 2y' + 2y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 - 2r + 2 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{2 \pm (\sqrt{-4})}{2}$$

- Roots of the characteristic polynomial

$$r = (1 - I, 1 + I)$$

- 1st solution of the ODE

$$y_1(x) = \cos(x) e^x$$

- 2nd solution of the ODE

$$y_2(x) = \sin(x) e^x$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 \cos(x) e^x + c_2 \sin(x) e^x$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 16

```
dsolve(diff(y(x),x$2) -2*diff(y(x),x)+2*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^x(c_1 \sin(x) + c_2 \cos(x))$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 20

```
DSolve[y''[x]-2*y'[x]+2*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^x(c_2 \cos(x) + c_1 \sin(x))$$

8.2 problem 8

- 8.2.1 Solving as second order linear constant coeff ode 2018
- 8.2.2 Solving using Kovacic algorithm 2020
- 8.2.3 Maple step by step solution 2024

Internal problem ID [624]

Internal file name [OUTPUT/624_Sunday_June_05_2022_01_46_01_AM_78108325/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 8.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - 2y' + 6y = 0$$

8.2.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -2, C = 6$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - 2\lambda e^{\lambda x} + 6 e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - 2\lambda + 6 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -2, C = 6$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-2^2 - (4)(1)(6)} \\ &= 1 \pm i\sqrt{5}\end{aligned}$$

Hence

$$\begin{aligned}\lambda_1 &= 1 + i\sqrt{5} \\ \lambda_2 &= 1 - i\sqrt{5}\end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= 1 + i\sqrt{5} \\ \lambda_2 &= 1 - i\sqrt{5}\end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 1$ and $\beta = \sqrt{5}$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x} (c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^x (c_1 \cos(x\sqrt{5}) + c_2 \sin(x\sqrt{5}))$$

Summary

The solution(s) found are the following

$$y = e^x (c_1 \cos(x\sqrt{5}) + c_2 \sin(x\sqrt{5})) \quad (1)$$

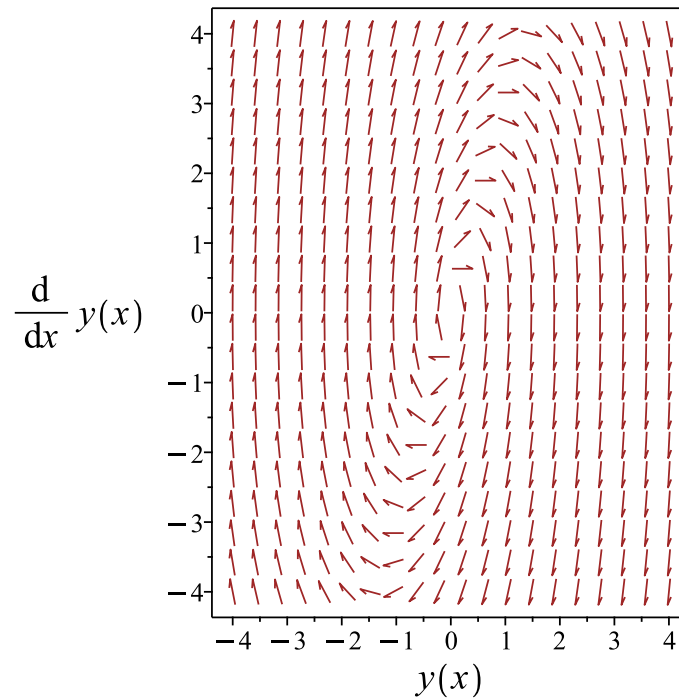


Figure 395: Slope field plot

Verification of solutions

$$y = e^x \left(c_1 \cos \left(x\sqrt{5} \right) + c_2 \sin \left(x\sqrt{5} \right) \right)$$

Verified OK.

8.2.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 2y' + 6y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= -2 \\ C &= 6 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-5}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -5 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -5z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 366: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -5$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x\sqrt{5})$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-2}{1} dx} \\ &= z_1 e^x \\ &= z_1 (e^x) \end{aligned}$$

Which simplifies to

$$y_1 = e^x \cos(x\sqrt{5})$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{2x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{\sqrt{5} \tan(x\sqrt{5})}{5} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^x \cos(x\sqrt{5}) \right) + c_2 \left(e^x \cos(x\sqrt{5}) \left(\frac{\sqrt{5} \tan(x\sqrt{5})}{5} \right) \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^x \cos(x\sqrt{5}) + \frac{c_2 \sin(x\sqrt{5}) e^x \sqrt{5}}{5} \quad (1)$$

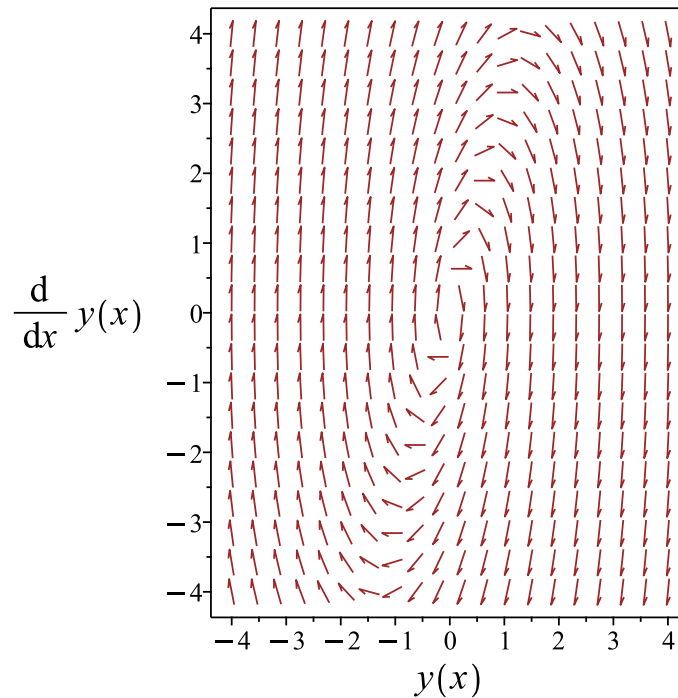


Figure 396: Slope field plot

Verification of solutions

$$y = c_1 e^x \cos(x\sqrt{5}) + \frac{c_2 \sin(x\sqrt{5}) e^x \sqrt{5}}{5}$$

Verified OK.

8.2.3 Maple step by step solution

Let's solve

$$y'' - 2y' + 6y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 - 2r + 6 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{2 \pm (\sqrt{-20})}{2}$$

- Roots of the characteristic polynomial

$$r = (1 - i\sqrt{5}, 1 + i\sqrt{5})$$

- 1st solution of the ODE

$$y_1(x) = e^x \cos(x\sqrt{5})$$

- 2nd solution of the ODE

$$y_2(x) = e^x \sin(x\sqrt{5})$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^x \cos(x\sqrt{5}) + c_2 e^x \sin(x\sqrt{5})$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 24

```
dsolve(diff(y(x),x$2) -2*diff(y(x),x)+6*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^x \left(c_1 \sin(\sqrt{5}x) + c_2 \cos(\sqrt{5}x) \right)$$

✓ Solution by Mathematica

Time used: 0.02 (sec). Leaf size: 32

```
DSolve[y''[x]-2*y'[x]+6*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^x \left(c_2 \cos(\sqrt{5}x) + c_1 \sin(\sqrt{5}x) \right)$$

8.3 problem 9

8.3.1 Solving as second order linear constant coeff ode	2026
8.3.2 Solving using Kovacic algorithm	2028
8.3.3 Maple step by step solution	2032

Internal problem ID [625]

Internal file name [OUTPUT/625_Sunday_June_05_2022_01_46_02_AM_10842904/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 9.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 2y' - 8y = 0$$

8.3.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 2, C = -8$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 2\lambda e^{\lambda x} - 8 e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 2\lambda - 8 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 2, C = -8$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{-2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{2^2 - (4)(1)(-8)} \\ &= -1 \pm 3\end{aligned}$$

Hence

$$\lambda_1 = -1 + 3$$

$$\lambda_2 = -1 - 3$$

Which simplifies to

$$\lambda_1 = 2$$

$$\lambda_2 = -4$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(2)x} + c_2 e^{(-4)x}$$

Or

$$y = c_1 e^{2x} + c_2 e^{-4x}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{2x} + c_2 e^{-4x} \tag{1}$$

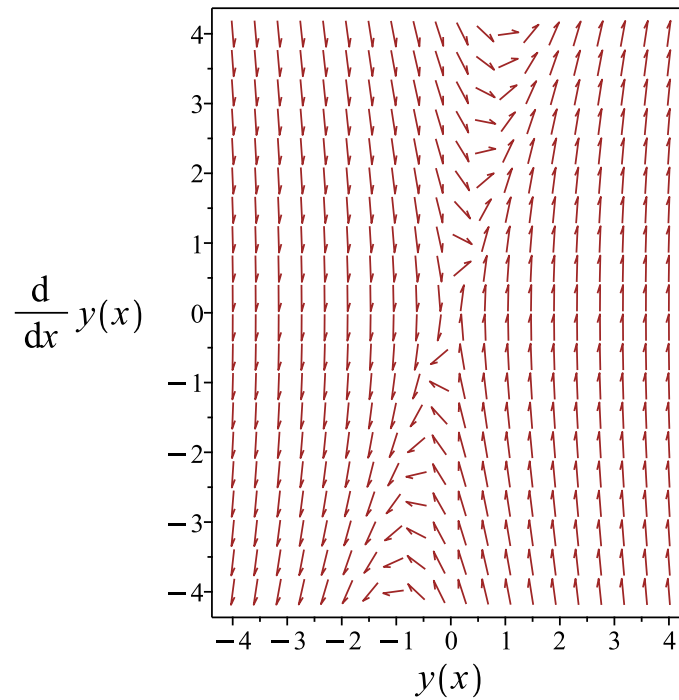


Figure 397: Slope field plot

Verification of solutions

$$y = c_1 e^{2x} + c_2 e^{-4x}$$

Verified OK.

8.3.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 2y' - 8y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 2 \tag{3}$$

$$C = -8$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{9}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 9$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = 9z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 368: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 9$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-3x}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{2}{1} dx} \\ &= z_1 e^{-x} \\ &= z_1 (e^{-x})\end{aligned}$$

Which simplifies to

$$y_1 = e^{-4x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-2x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{e^{6x}}{6} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-4x}) + c_2 \left(e^{-4x} \left(\frac{e^{6x}}{6} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-4x} + \frac{c_2 e^{2x}}{6} \quad (1)$$

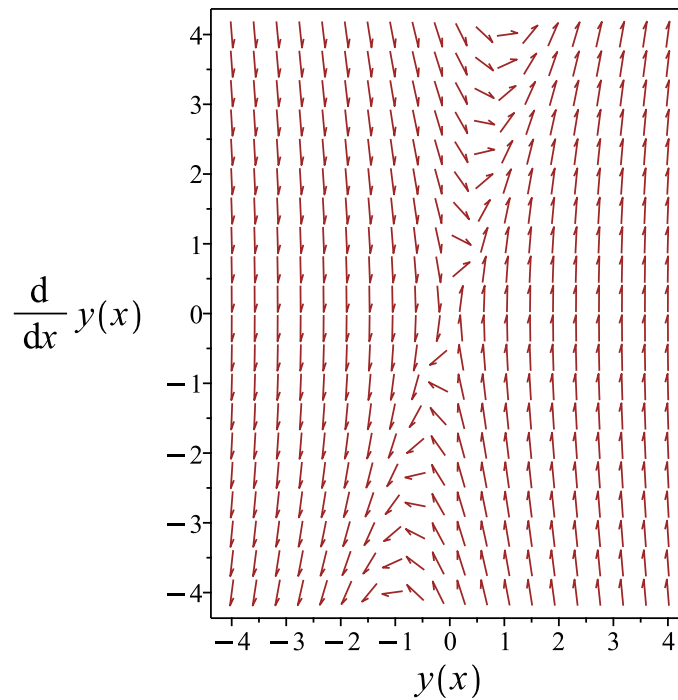


Figure 398: Slope field plot

Verification of solutions

$$y = c_1 e^{-4x} + \frac{c_2 e^{2x}}{6}$$

Verified OK.

8.3.3 Maple step by step solution

Let's solve

$$y'' + 2y' - 8y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 2r - 8 = 0$$

- Factor the characteristic polynomial

$$(r + 4)(r - 2) = 0$$

- Roots of the characteristic polynomial

$$r = (-4, 2)$$

- 1st solution of the ODE

$$y_1(x) = e^{-4x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{2x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-4x} + c_2 e^{2x}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 17

```
dsolve(diff(y(x),x$2) +2*diff(y(x),x)-8*y(x) = 0,y(x), singsol=all)
```

$$y(x) = (e^{6x}c_1 + c_2) e^{-4x}$$

✓ Solution by Mathematica

Time used: 0.012 (sec). Leaf size: 22

```
DSolve[y''[x]+2*y'[x]-8*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-4x}(c_2 e^{6x} + c_1)$$

8.4 problem 10

- 8.4.1 Solving as second order linear constant coeff ode 2034
- 8.4.2 Solving using Kovacic algorithm 2036
- 8.4.3 Maple step by step solution 2040

Internal problem ID [626]

Internal file name [OUTPUT/626_Sunday_June_05_2022_01_46_02_AM_81884660/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 10.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 2y' + 2y = 0$$

8.4.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 2, C = 2$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 2\lambda e^{\lambda x} + 2e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 2\lambda + 2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 2, C = 2$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{-2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{2^2 - (4)(1)(2)} \\ &= -1 \pm i\end{aligned}$$

Hence

$$\begin{aligned}\lambda_1 &= -1 + i \\ \lambda_2 &= -1 - i\end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= -1 + i \\ \lambda_2 &= -1 - i\end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -1$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-x}(c_1 \cos(x) + c_2 \sin(x))$$

Summary

The solution(s) found are the following

$$y = e^{-x}(c_1 \cos(x) + c_2 \sin(x)) \tag{1}$$

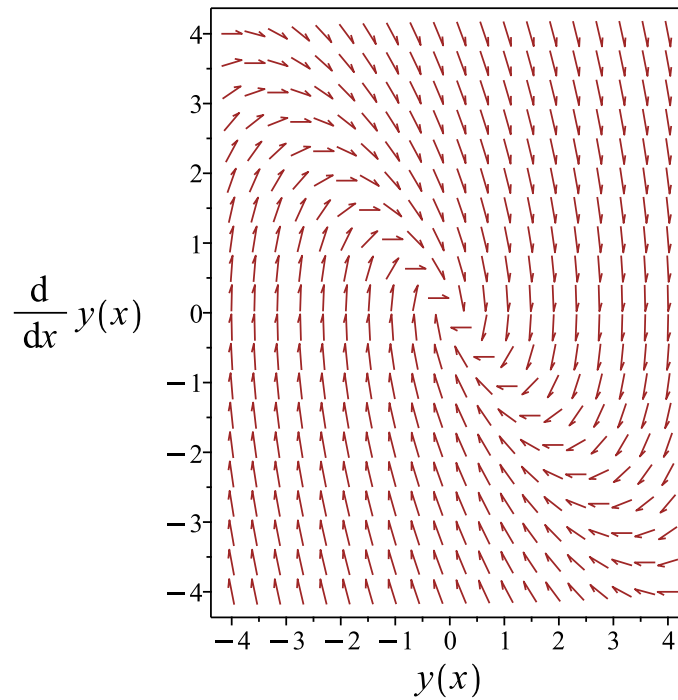


Figure 399: Slope field plot

Verification of solutions

$$y = e^{-x}(c_1 \cos(x) + c_2 \sin(x))$$

Verified OK.

8.4.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 2y' + 2y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 2 \tag{3}$$

$$C = 2$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 370: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{2}{1} dx} \\ &= z_1 e^{-x} \\ &= z_1 (e^{-x})\end{aligned}$$

Which simplifies to

$$y_1 = \cos(x) e^{-x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-2x}}{(y_1)^2} dx \\ &= y_1 (\tan(x))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (\cos(x) e^{-x}) + c_2 (\cos(x) e^{-x} (\tan(x)))\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(x) e^{-x} + c_2 \sin(x) e^{-x} \quad (1)$$

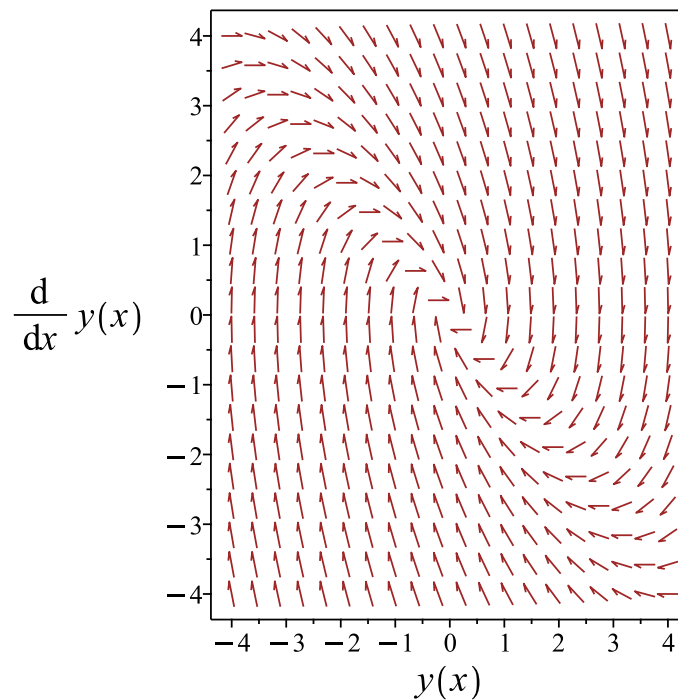


Figure 400: Slope field plot

Verification of solutions

$$y = c_1 \cos(x) e^{-x} + c_2 \sin(x) e^{-x}$$

Verified OK.

8.4.3 Maple step by step solution

Let's solve

$$y'' + 2y' + 2y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 2r + 2 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-2) \pm (\sqrt{-4})}{2}$$

- Roots of the characteristic polynomial

$$r = (-1 - I, -1 + I)$$

- 1st solution of the ODE

$$y_1(x) = \cos(x) e^{-x}$$

- 2nd solution of the ODE

$$y_2(x) = \sin(x) e^{-x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 \cos(x) e^{-x} + c_2 \sin(x) e^{-x}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 18

```
dsolve(diff(y(x),x$2) +2*diff(y(x),x)+2*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-x}(c_1 \sin(x) + c_2 \cos(x))$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 22

```
DSolve[y''[x]+2*y'[x]+2*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x}(c_2 \cos(x) + c_1 \sin(x))$$

8.5 problem 11

- 8.5.1 Solving as second order linear constant coeff ode 2042
- 8.5.2 Solving using Kovacic algorithm 2044
- 8.5.3 Maple step by step solution 2048

Internal problem ID [627]

Internal file name [OUTPUT/627_Sunday_June_05_2022_01_46_03_AM_30401613/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 11.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 6y' + 13y = 0$$

8.5.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 6, C = 13$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 6\lambda e^{\lambda x} + 13e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 6\lambda + 13 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 6, C = 13$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{-6}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{6^2 - (4)(1)(13)} \\ &= -3 \pm 2i\end{aligned}$$

Hence

$$\begin{aligned}\lambda_1 &= -3 + 2i \\ \lambda_2 &= -3 - 2i\end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= -3 + 2i \\ \lambda_2 &= -3 - 2i\end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -3$ and $\beta = 2$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-3x}(c_1 \cos(2x) + c_2 \sin(2x))$$

Summary

The solution(s) found are the following

$$y = e^{-3x}(c_1 \cos(2x) + c_2 \sin(2x)) \quad (1)$$

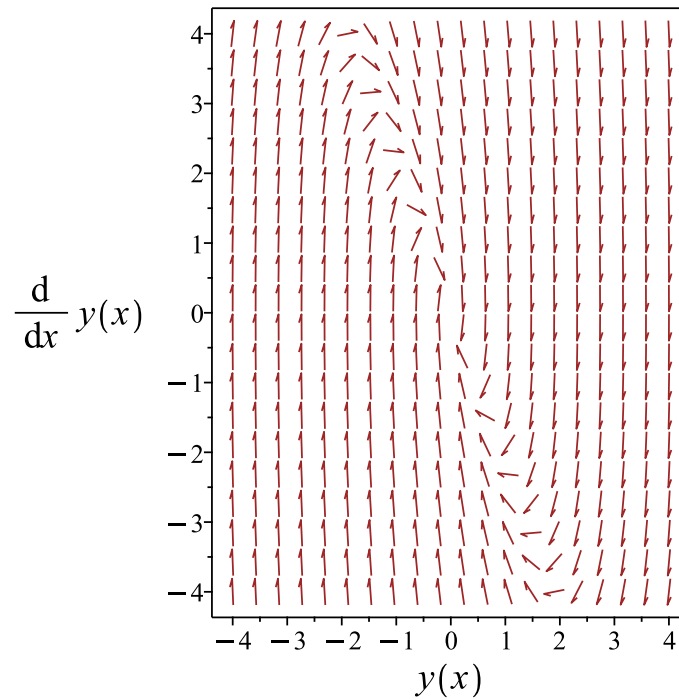


Figure 401: Slope field plot

Verification of solutions

$$y = e^{-3x}(c_1 \cos(2x) + c_2 \sin(2x))$$

Verified OK.

8.5.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 6y' + 13y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 6 \\ C &= 13 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-4}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -4 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -4z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 372: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -4$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(2x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{6}{1} dx} \\ &= z_1 e^{-3x} \\ &= z_1 (e^{-3x})\end{aligned}$$

Which simplifies to

$$y_1 = e^{-3x} \cos(2x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{6}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-6x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{\tan(2x)}{2} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-3x} \cos(2x)) + c_2 \left(e^{-3x} \cos(2x) \left(\frac{\tan(2x)}{2} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-3x} \cos(2x) + \frac{c_2 e^{-3x} \sin(2x)}{2} \quad (1)$$

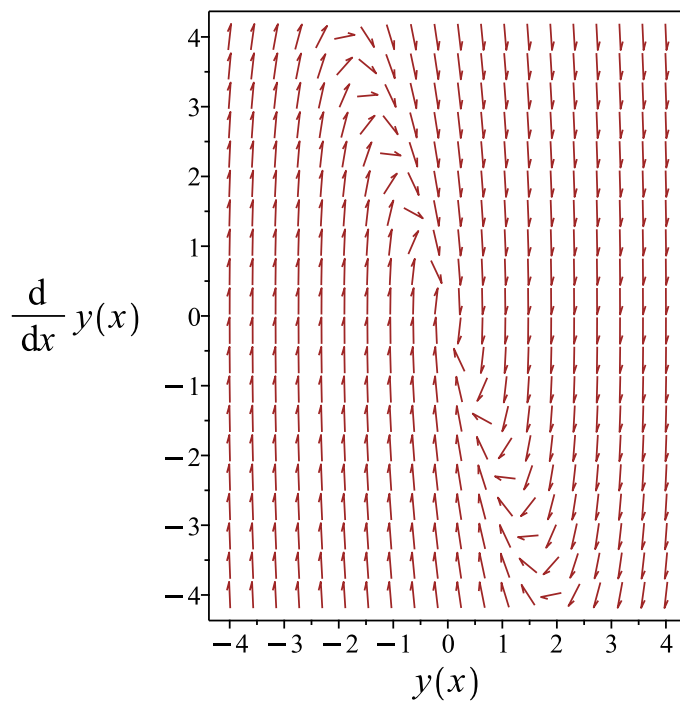


Figure 402: Slope field plot

Verification of solutions

$$y = c_1 e^{-3x} \cos(2x) + \frac{c_2 e^{-3x} \sin(2x)}{2}$$

Verified OK.

8.5.3 Maple step by step solution

Let's solve

$$y'' + 6y' + 13y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 6r + 13 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-6) \pm (\sqrt{-16})}{2}$$

- Roots of the characteristic polynomial

$$r = (-3 - 2I, -3 + 2I)$$

- 1st solution of the ODE

$$y_1(x) = e^{-3x} \cos(2x)$$

- 2nd solution of the ODE

$$y_2(x) = e^{-3x} \sin(2x)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-3x} \cos(2x) + c_2 e^{-3x} \sin(2x)$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 22

```
dsolve(diff(y(x),x$2) +6*diff(y(x),x)+13*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-3x}(\sin(2x)c_1 + c_2 \cos(2x))$$

✓ Solution by Mathematica

Time used: 0.018 (sec). Leaf size: 26

```
DSolve[y''[x]+6*y'[x]+13*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-3x}(c_2 \cos(2x) + c_1 \sin(2x))$$

8.6 problem 12

8.6.1	Solving as second order linear constant coeff ode	2050
8.6.2	Solving as second order ode can be made integrable ode	2052
8.6.3	Solving using Kovacic algorithm	2054
8.6.4	Maple step by step solution	2058

Internal problem ID [628]

Internal file name [OUTPUT/628_Sunday_June_05_2022_01_46_04_AM_31108206/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 12.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "second_order_ode_can_be_made_integrable"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$4y'' + 9y = 0$$

8.6.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 4, B = 0, C = 9$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$4\lambda^2 e^{\lambda x} + 9e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$4\lambda^2 + 9 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 4, B = 0, C = 9$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(4)} \pm \frac{1}{(2)(4)} \sqrt{0^2 - (4)(4)(9)} \\ &= \pm \frac{3i}{2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= +\frac{3i}{2} \\ \lambda_2 &= -\frac{3i}{2} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= \frac{3i}{2} \\ \lambda_2 &= -\frac{3i}{2} \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = \frac{3}{2}$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x} (c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^0 \left(c_1 \cos \left(\frac{3x}{2} \right) + c_2 \sin \left(\frac{3x}{2} \right) \right)$$

Or

$$y = c_1 \cos \left(\frac{3x}{2} \right) + c_2 \sin \left(\frac{3x}{2} \right)$$

Summary

The solution(s) found are the following

$$y = c_1 \cos\left(\frac{3x}{2}\right) + c_2 \sin\left(\frac{3x}{2}\right) \quad (1)$$

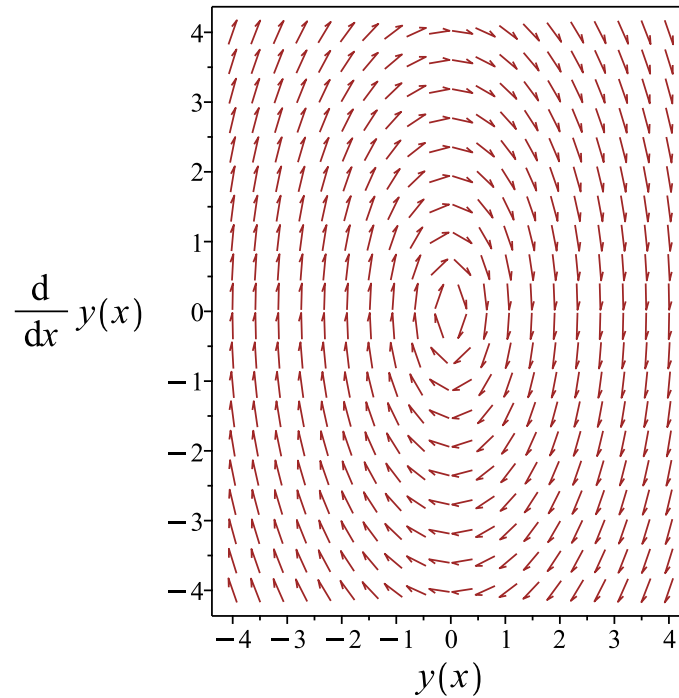


Figure 403: Slope field plot

Verification of solutions

$$y = c_1 \cos\left(\frac{3x}{2}\right) + c_2 \sin\left(\frac{3x}{2}\right)$$

Verified OK.

8.6.2 Solving as second order ode can be made integrable ode

Multiplying the ode by y' gives

$$4y'y'' + 9yy' = 0$$

Integrating the above w.r.t x gives

$$\int (4y'y'' + 9yy') dx = 0$$
$$2y'^2 + \frac{9y^2}{2} = c_2$$

Which is now solved for y . Solving the given ode for y' results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$y' = \frac{\sqrt{-9y^2 + 2c_1}}{2} \quad (1)$$

$$y' = -\frac{\sqrt{-9y^2 + 2c_1}}{2} \quad (2)$$

Now each one of the above ODE is solved.

Solving equation (1)

Integrating both sides gives

$$\int \frac{2}{\sqrt{-9y^2 + 2c_1}} dy = \int dx$$
$$\frac{2 \arctan\left(\frac{3y}{\sqrt{-9y^2 + 2c_1}}\right)}{3} = x + c_2$$

Solving equation (2)

Integrating both sides gives

$$\int -\frac{2}{\sqrt{-9y^2 + 2c_1}} dy = \int dx$$
$$-\frac{2 \arctan\left(\frac{3y}{\sqrt{-9y^2 + 2c_1}}\right)}{3} = c_3 + x$$

Summary

The solution(s) found are the following

$$\frac{2 \arctan\left(\frac{3y}{\sqrt{-9y^2 + 2c_1}}\right)}{3} = x + c_2 \quad (1)$$

$$-\frac{2 \arctan\left(\frac{3y}{\sqrt{-9y^2 + 2c_1}}\right)}{3} = c_3 + x \quad (2)$$

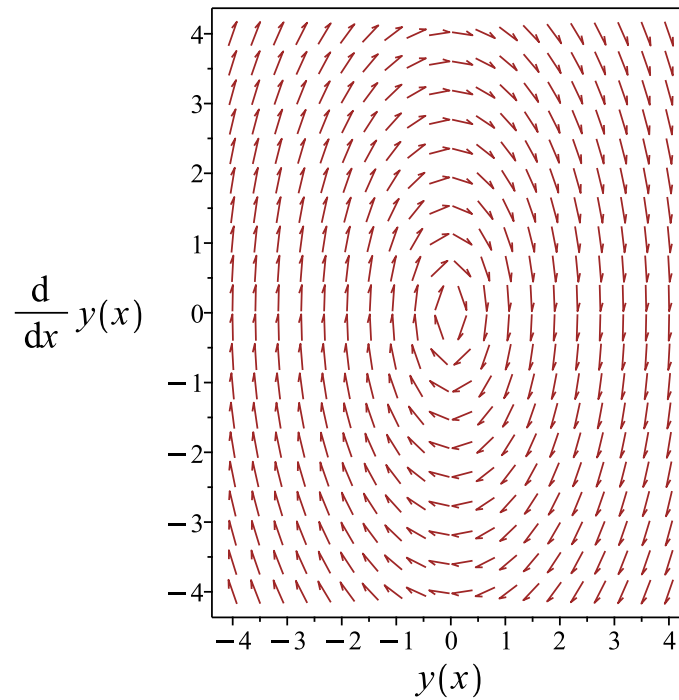


Figure 404: Slope field plot

Verification of solutions

$$\frac{2 \arctan\left(\frac{3y}{\sqrt{-9y^2+2c_1}}\right)}{3} = x + c_2$$

Verified OK.

$$-\frac{2 \arctan\left(\frac{3y}{\sqrt{-9y^2+2c_1}}\right)}{3} = c_3 + x$$

Verified OK.

8.6.3 Solving using Kovacic algorithm

Writing the ode as

$$4y'' + 9y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= 4 \\B &= 0 \\C &= 9\end{aligned}\tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x)\tag{4}$$

Where r is given by

$$\begin{aligned}r &= \frac{s}{t} \\&= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}\end{aligned}\tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-9}{4}\tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned}s &= -9 \\t &= 4\end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -\frac{9z(x)}{4}\tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 374: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -\frac{9}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos\left(\frac{3x}{2}\right)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}
 y_1 &= z_1 \\
 &= \cos\left(\frac{3x}{2}\right)
 \end{aligned}$$

Which simplifies to

$$y_1 = \cos\left(\frac{3x}{2}\right)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Since $B = 0$ then the above becomes

$$\begin{aligned} y_2 &= y_1 \int \frac{1}{y_1^2} dx \\ &= \cos\left(\frac{3x}{2}\right) \int \frac{1}{\cos\left(\frac{3x}{2}\right)^2} dx \\ &= \cos\left(\frac{3x}{2}\right) \left(\frac{2 \tan\left(\frac{3x}{2}\right)}{3}\right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(\cos\left(\frac{3x}{2}\right)\right) + c_2 \left(\cos\left(\frac{3x}{2}\right) \left(\frac{2 \tan\left(\frac{3x}{2}\right)}{3}\right)\right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 \cos\left(\frac{3x}{2}\right) + \frac{2c_2 \sin\left(\frac{3x}{2}\right)}{3} \quad (1)$$

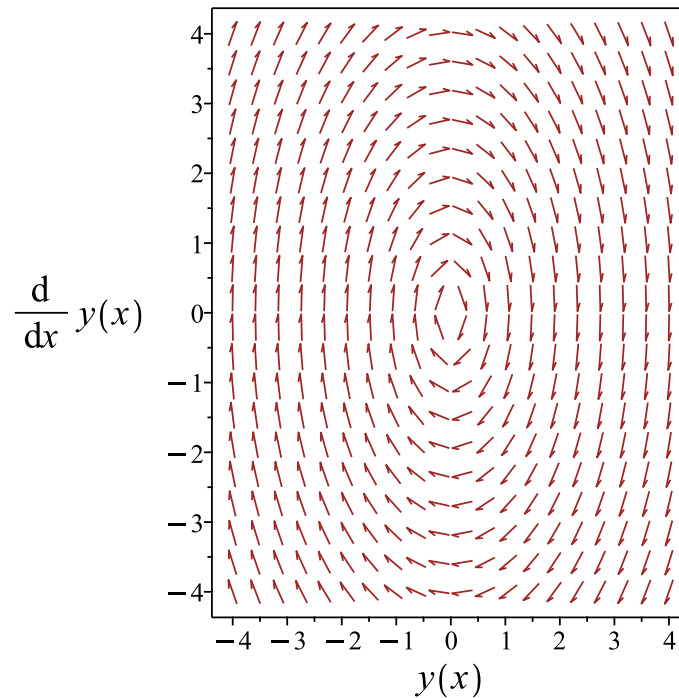


Figure 405: Slope field plot

Verification of solutions

$$y = c_1 \cos\left(\frac{3x}{2}\right) + \frac{2c_2 \sin\left(\frac{3x}{2}\right)}{3}$$

Verified OK.

8.6.4 Maple step by step solution

Let's solve

$$4y'' + 9y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{9y}{4}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{9y}{4} = 0$$

- Characteristic polynomial of ODE

$$r^2 + \frac{9}{4} = 0$$

- Use quadratic formula to solve for r

$$r = \frac{0 \pm (\sqrt{-9})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{3i}{2}, \frac{3i}{2}\right)$$

- 1st solution of the ODE

$$y_1(x) = \cos\left(\frac{3x}{2}\right)$$

- 2nd solution of the ODE

$$y_2(x) = \sin\left(\frac{3x}{2}\right)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 \cos\left(\frac{3x}{2}\right) + c_2 \sin\left(\frac{3x}{2}\right)$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(4*diff(y(x),x$2) +9*y(x) = 0,y(x), singsol=all)
```

$$y(x) = c_1 \sin\left(\frac{3x}{2}\right) + c_2 \cos\left(\frac{3x}{2}\right)$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 20

```
DSolve[y''[x]+9*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow c_1 \cos(3x) + c_2 \sin(3x)$$

8.7 problem 13

- 8.7.1 Solving as second order linear constant coeff ode 2061
- 8.7.2 Solving using Kovacic algorithm 2063
- 8.7.3 Maple step by step solution 2067

Internal problem ID [629]

Internal file name [OUTPUT/629_Sunday_June_05_2022_01_46_05_AM_47762062/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 13.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 2y' + \frac{5y}{4} = 0$$

8.7.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 2, C = \frac{5}{4}$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 2\lambda e^{\lambda x} + \frac{5 e^{\lambda x}}{4} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 2\lambda + \frac{5}{4} = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 2, C = \frac{5}{4}$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{2^2 - (4)(1)\left(\frac{5}{4}\right)} \\ &= -1 \pm \frac{i}{2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= -1 + \frac{i}{2} \\ \lambda_2 &= -1 - \frac{i}{2} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= -1 + \frac{i}{2} \\ \lambda_2 &= -1 - \frac{i}{2} \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -1$ and $\beta = \frac{1}{2}$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x} (c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-x} \left(c_1 \cos\left(\frac{x}{2}\right) + c_2 \sin\left(\frac{x}{2}\right) \right)$$

Summary

The solution(s) found are the following

$$y = e^{-x} \left(c_1 \cos\left(\frac{x}{2}\right) + c_2 \sin\left(\frac{x}{2}\right) \right) \quad (1)$$

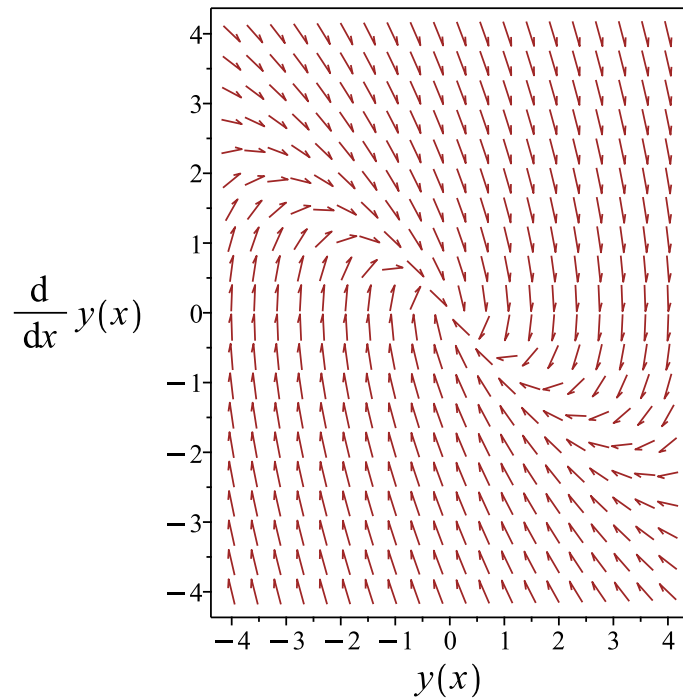


Figure 406: Slope field plot

Verification of solutions

$$y = e^{-x} \left(c_1 \cos \left(\frac{x}{2} \right) + c_2 \sin \left(\frac{x}{2} \right) \right)$$

Verified OK.

8.7.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 2y' + \frac{5y}{4} = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 2 \\ C &= \frac{5}{4} \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -\frac{z(x)}{4} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 376: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -\frac{1}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos\left(\frac{x}{2}\right)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} dx} \\ &= z_1 e^{-\frac{1}{2}x} \\ &= z_1 e^{-x} \\ &= z_1 (e^{-x}) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-x} \cos\left(\frac{x}{2}\right)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-2x}}{(y_1)^2} dx \\ &= y_1 \left(2 \tan\left(\frac{x}{2}\right) \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-x} \cos \left(\frac{x}{2} \right) \right) + c_2 \left(e^{-x} \cos \left(\frac{x}{2} \right) \left(2 \tan \left(\frac{x}{2} \right) \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-x} \cos \left(\frac{x}{2} \right) + 2c_2 e^{-x} \sin \left(\frac{x}{2} \right) \quad (1)$$

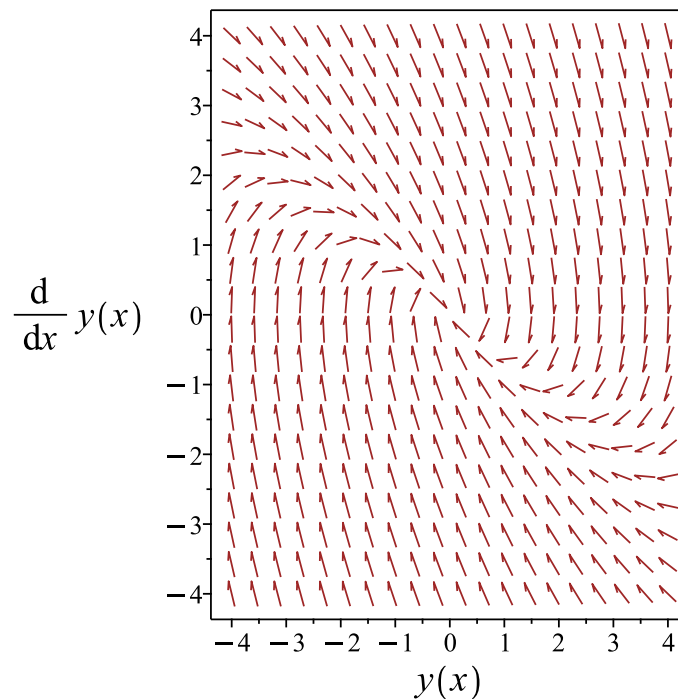


Figure 407: Slope field plot

Verification of solutions

$$y = c_1 e^{-x} \cos \left(\frac{x}{2} \right) + 2c_2 e^{-x} \sin \left(\frac{x}{2} \right)$$

Verified OK.

8.7.3 Maple step by step solution

Let's solve

$$y'' + 2y' + \frac{5y}{4} = 0$$

- Highest derivative means the order of the ODE is 2

y''

- Characteristic polynomial of ODE

$$r^2 + 2r + \frac{5}{4} = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-2) \pm (\sqrt{-1})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-1 - \frac{1}{2}, -1 + \frac{1}{2}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-x} \cos\left(\frac{x}{2}\right)$$

- 2nd solution of the ODE

$$y_2(x) = e^{-x} \sin\left(\frac{x}{2}\right)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-x} \cos\left(\frac{x}{2}\right) + c_2 e^{-x} \sin\left(\frac{x}{2}\right)$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 22

```
dsolve(diff(y(x),x$2) +2*diff(y(x),x)+125/100*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-x} \left(c_1 \sin \left(\frac{x}{2} \right) + c_2 \cos \left(\frac{x}{2} \right) \right)$$

✓ Solution by Mathematica

Time used: 0.018 (sec). Leaf size: 30

```
DSolve[y''[x]+2*y'[x]+125/100*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x} \left(c_2 \cos \left(\frac{x}{2} \right) + c_1 \sin \left(\frac{x}{2} \right) \right)$$

8.8 problem 14

8.8.1 Solving as second order linear constant coeff ode	2069
8.8.2 Solving using Kovacic algorithm	2071
8.8.3 Maple step by step solution	2075

Internal problem ID [630]

Internal file name [OUTPUT/630_Sunday_June_05_2022_01_46_06_AM_17484274/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 14.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$9y'' + 9y' - 4y = 0$$

8.8.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 9, B = 9, C = -4$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$9\lambda^2 e^{\lambda x} + 9\lambda e^{\lambda x} - 4e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$9\lambda^2 + 9\lambda - 4 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 9, B = 9, C = -4$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{-9}{(2)(9)} \pm \frac{1}{(2)(9)} \sqrt{9^2 - (4)(9)(-4)} \\ &= -\frac{1}{2} \pm \frac{5}{6}\end{aligned}$$

Hence

$$\lambda_1 = -\frac{1}{2} + \frac{5}{6}$$

$$\lambda_2 = -\frac{1}{2} - \frac{5}{6}$$

Which simplifies to

$$\lambda_1 = \frac{1}{3}$$

$$\lambda_2 = -\frac{4}{3}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(\frac{1}{3})x} + c_2 e^{(-\frac{4}{3})x}$$

Or

$$y = c_1 e^{\frac{x}{3}} + c_2 e^{-\frac{4x}{3}}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{\frac{x}{3}} + c_2 e^{-\frac{4x}{3}} \tag{1}$$

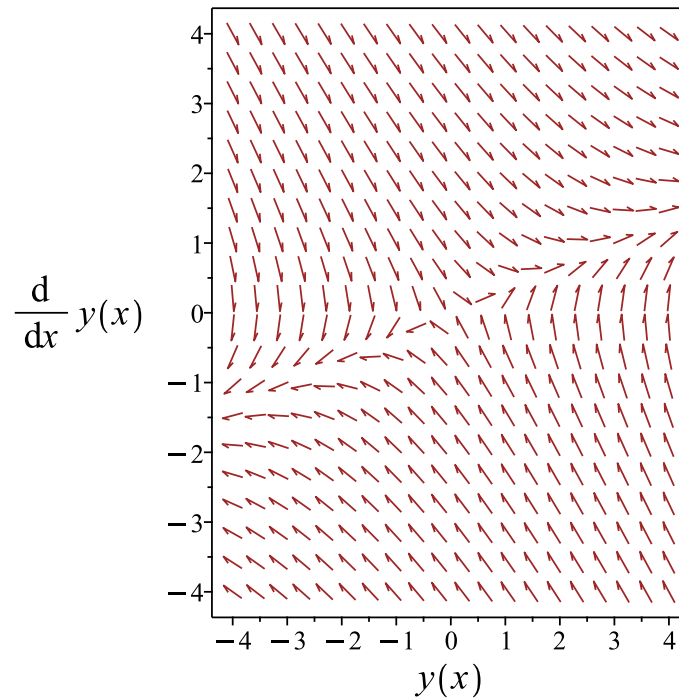


Figure 408: Slope field plot

Verification of solutions

$$y = c_1 e^{\frac{x}{3}} + c_2 e^{-\frac{4x}{3}}$$

Verified OK.

8.8.2 Solving using Kovacic algorithm

Writing the ode as

$$9y'' + 9y' - 4y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 9 \\ B &= 9 \\ C &= -4 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{25}{36} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 25 \\ t &= 36 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{25z(x)}{36} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 378: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{25}{36}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{5x}{6}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{9}{9} dx} \\ &= z_1 e^{-\frac{x}{2}} \\ &= z_1 \left(e^{-\frac{x}{2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{4x}{3}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{9}{9} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{3e^{\frac{5x}{3}}}{5} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-\frac{4x}{3}} \right) + c_2 \left(e^{-\frac{4x}{3}} \left(\frac{3 e^{\frac{5x}{3}}}{5} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{4x}{3}} + \frac{3c_2 e^{\frac{x}{3}}}{5} \quad (1)$$

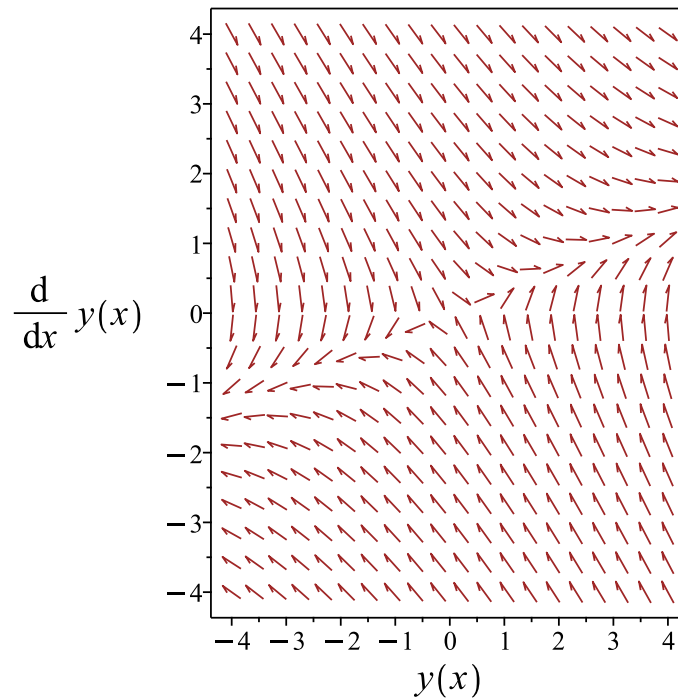


Figure 409: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{4x}{3}} + \frac{3c_2 e^{\frac{x}{3}}}{5}$$

Verified OK.

8.8.3 Maple step by step solution

Let's solve

$$9y'' + 9y' - 4y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -y' + \frac{4y}{9}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + y' - \frac{4y}{9} = 0$$

- Characteristic polynomial of ODE

$$r^2 + r - \frac{4}{9} = 0$$

- Factor the characteristic polynomial

$$\frac{(3r+4)(3r-1)}{9} = 0$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{4}{3}, \frac{1}{3}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{4x}{3}}$$

- 2nd solution of the ODE

$$y_2(x) = e^{\frac{x}{3}}$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^{-\frac{4x}{3}} + c_2e^{\frac{x}{3}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(9*diff(y(x),x$2) +9*diff(y(x),x)-4*y(x) = 0,y(x), singsol=all)
```

$$y(x) = \left(c_2 e^{\frac{5x}{3}} + c_1 \right) e^{-\frac{4x}{3}}$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 26

```
DSolve[9*y''[x]+9*y'[x]-4*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-4x/3} (c_2 e^{5x/3} + c_1)$$

8.9 problem 15

8.9.1	Solving as second order linear constant coeff ode	2077
8.9.2	Solving using Kovacic algorithm	2079
8.9.3	Maple step by step solution	2083

Internal problem ID [631]

Internal file name [OUTPUT/631_Sunday_June_05_2022_01_46_06_AM_45655324/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 15.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + y' + \frac{5y}{4} = 0$$

8.9.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 1, C = \frac{5}{4}$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + \lambda e^{\lambda x} + \frac{5 e^{\lambda x}}{4} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + \lambda + \frac{5}{4} = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 1, C = \frac{5}{4}$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{1^2 - (4)(1)\left(\frac{5}{4}\right)} \\ &= -\frac{1}{2} \pm i \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= -\frac{1}{2} + i \\ \lambda_2 &= -\frac{1}{2} - i \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= -\frac{1}{2} + i \\ \lambda_2 &= -\frac{1}{2} - i \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -\frac{1}{2}$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x} (c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-\frac{x}{2}} (c_1 \cos(x) + c_2 \sin(x))$$

Summary

The solution(s) found are the following

$$y = e^{-\frac{x}{2}} (c_1 \cos(x) + c_2 \sin(x)) \quad (1)$$

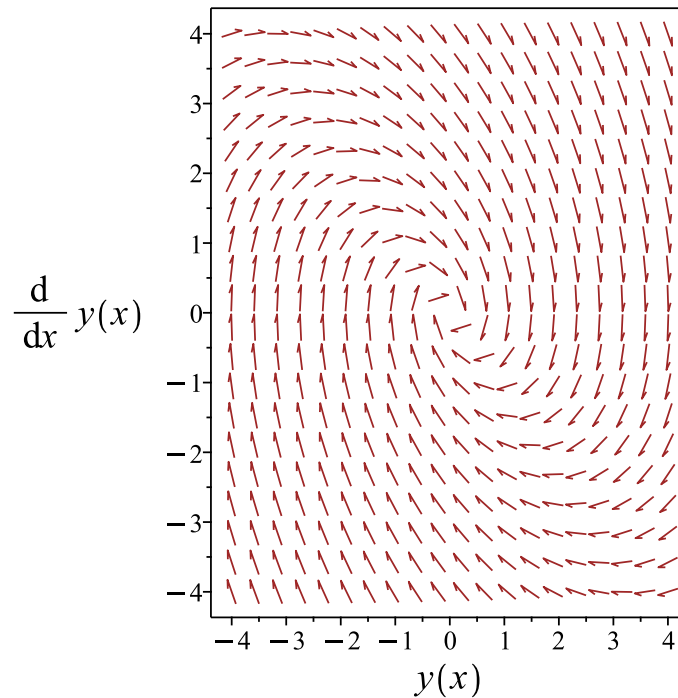


Figure 410: Slope field plot

Verification of solutions

$$y = e^{-\frac{x}{2}}(c_1 \cos(x) + c_2 \sin(x))$$

Verified OK.

8.9.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + y' + \frac{5y}{4} = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 1 \tag{3}$$

$$C = \frac{5}{4}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = -1$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = -z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 380: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} dx} \\ &= z_1 e^{-\frac{1}{2}x} \\ &= z_1 e^{-\frac{x}{2}} \\ &= z_1 (e^{-\frac{x}{2}})\end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{x}{2}} \cos(x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{1}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-x}}{(y_1)^2} dx \\ &= y_1 (\tan(x))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-\frac{x}{2}} \cos(x)) + c_2 (e^{-\frac{x}{2}} \cos(x) (\tan(x)))\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{x}{2}} \cos(x) + c_2 e^{-\frac{x}{2}} \sin(x) \quad (1)$$

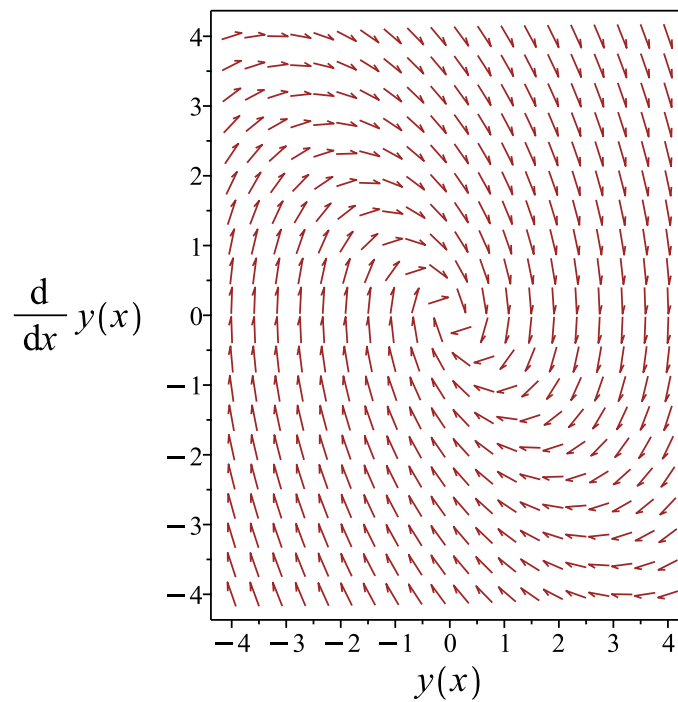


Figure 411: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{x}{2}} \cos(x) + c_2 e^{-\frac{x}{2}} \sin(x)$$

Verified OK.

8.9.3 Maple step by step solution

Let's solve

$$y'' + y' + \frac{5y}{4} = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + r + \frac{5}{4} = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-1) \pm (\sqrt{-4})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{2} - I, -\frac{1}{2} + I\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{x}{2}} \cos(x)$$

- 2nd solution of the ODE

$$y_2(x) = e^{-\frac{x}{2}} \sin(x)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-\frac{x}{2}} \cos(x) + c_2 e^{-\frac{x}{2}} \sin(x)$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 18

```
dsolve(diff(y(x),x$2) +diff(y(x),x)+125/100*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-\frac{x}{2}}(c_1 \sin(x) + c_2 \cos(x))$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 24

```
DSolve[y''[x]+y'[x]+125/100*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x/2}(c_2 \cos(x) + c_1 \sin(x))$$

8.10 problem 16

8.10.1 Solving as second order linear constant coeff ode	2085
8.10.2 Solving using Kovacic algorithm	2087
8.10.3 Maple step by step solution	2091

Internal problem ID [632]

Internal file name [OUTPUT/632_Sunday_June_05_2022_01_46_07_AM_6524054/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 16.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 4y' + \frac{25y}{4} = 0$$

8.10.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 4, C = \frac{25}{4}$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 4\lambda e^{\lambda x} + \frac{25 e^{\lambda x}}{4} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 4\lambda + \frac{25}{4} = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 4, C = \frac{25}{4}$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-4}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{4^2 - (4)(1) \left(\frac{25}{4}\right)} \\ &= -2 \pm \frac{3i}{2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= -2 + \frac{3i}{2} \\ \lambda_2 &= -2 - \frac{3i}{2} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= -2 + \frac{3i}{2} \\ \lambda_2 &= -2 - \frac{3i}{2} \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -2$ and $\beta = \frac{3}{2}$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x} (c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-2x} \left(c_1 \cos \left(\frac{3x}{2} \right) + c_2 \sin \left(\frac{3x}{2} \right) \right)$$

Summary

The solution(s) found are the following

$$y = e^{-2x} \left(c_1 \cos \left(\frac{3x}{2} \right) + c_2 \sin \left(\frac{3x}{2} \right) \right) \quad (1)$$

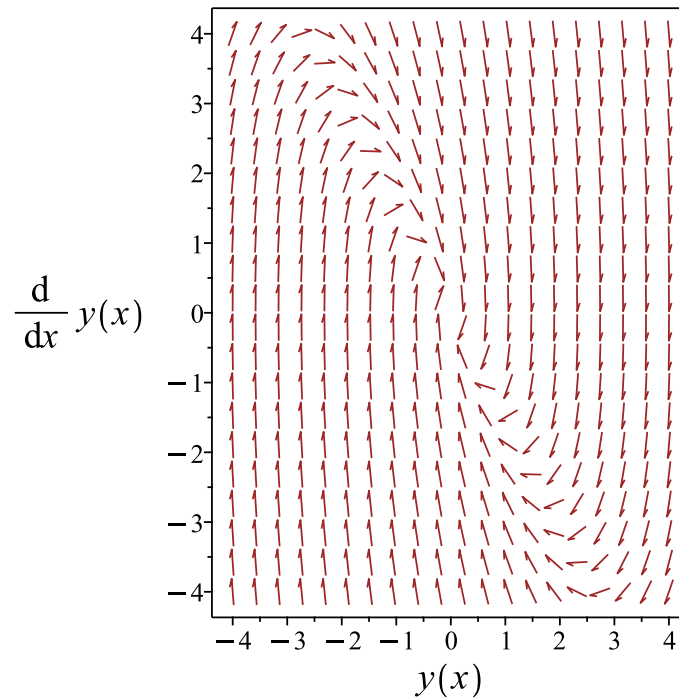


Figure 412: Slope field plot

Verification of solutions

$$y = e^{-2x} \left(c_1 \cos \left(\frac{3x}{2} \right) + c_2 \sin \left(\frac{3x}{2} \right) \right)$$

Verified OK.

8.10.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 4y' + \frac{25y}{4} = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 4 \\ C &= \frac{25}{4} \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-9}{4} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -9 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -\frac{9z(x)}{4} \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 382: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -\frac{9}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos\left(\frac{3x}{2}\right)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}
 y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\
 &= z_1 e^{-\int \frac{1}{2} \frac{4}{1} dx} \\
 &= z_1 e^{-2x} \\
 &= z_1 (e^{-2x})
 \end{aligned}$$

Which simplifies to

$$y_1 = e^{-2x} \cos\left(\frac{3x}{2}\right)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{4}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-4x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{2 \tan\left(\frac{3x}{2}\right)}{3} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-2x} \cos\left(\frac{3x}{2}\right) \right) + c_2 \left(e^{-2x} \cos\left(\frac{3x}{2}\right) \left(\frac{2 \tan\left(\frac{3x}{2}\right)}{3} \right) \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-2x} \cos\left(\frac{3x}{2}\right) + \frac{2c_2 e^{-2x} \sin\left(\frac{3x}{2}\right)}{3} \quad (1)$$

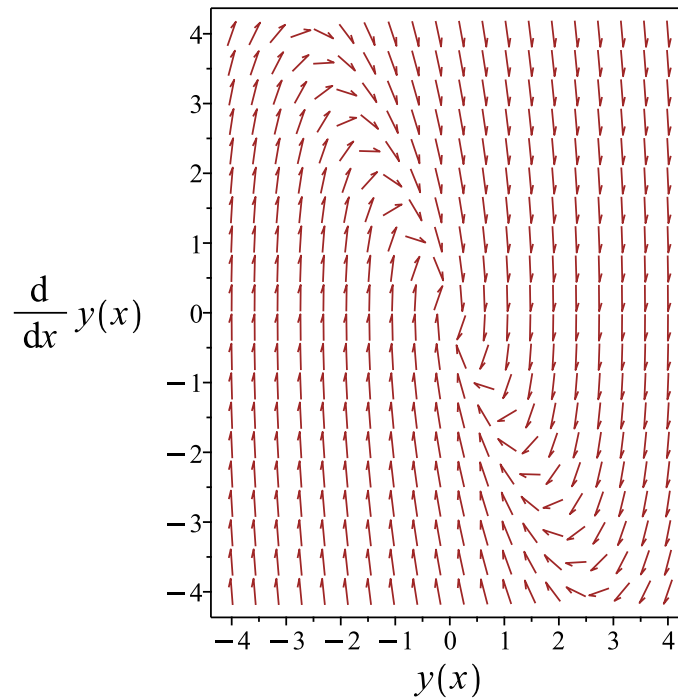


Figure 413: Slope field plot

Verification of solutions

$$y = c_1 e^{-2x} \cos\left(\frac{3x}{2}\right) + \frac{2c_2 e^{-2x} \sin\left(\frac{3x}{2}\right)}{3}$$

Verified OK.

8.10.3 Maple step by step solution

Let's solve

$$y'' + 4y' + \frac{25y}{4} = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 4r + \frac{25}{4} = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-4) \pm (\sqrt{-9})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-2 - \frac{3I}{2}, -2 + \frac{3I}{2}\right)$$
- 1st solution of the ODE

$$y_1(x) = e^{-2x} \cos\left(\frac{3x}{2}\right)$$
- 2nd solution of the ODE

$$y_2(x) = e^{-2x} \sin\left(\frac{3x}{2}\right)$$
- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$
- Substitute in solutions

$$y = c_1 e^{-2x} \cos\left(\frac{3x}{2}\right) + c_2 e^{-2x} \sin\left(\frac{3x}{2}\right)$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 22

```
dsolve(diff(y(x),x$2)+ 4*diff(y(x),x)+625/100*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-2x} \left(c_1 \sin\left(\frac{3x}{2}\right) + c_2 \cos\left(\frac{3x}{2}\right) \right)$$

✓ Solution by Mathematica

Time used: 0.018 (sec). Leaf size: 30

```
DSolve[y''[x]+4*y'[x]+625/100*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-2x} \left(c_2 \cos\left(\frac{3x}{2}\right) + c_1 \sin\left(\frac{3x}{2}\right) \right)$$

8.11 problem 17

8.11.1 Existence and uniqueness analysis	2093
8.11.2 Solving as second order linear constant coeff ode	2094
8.11.3 Solving as second order ode can be made integrable ode	2096
8.11.4 Solving using Kovacic algorithm	2099
8.11.5 Maple step by step solution	2103

Internal problem ID [633]

Internal file name [OUTPUT/633_Sunday_June_05_2022_01_46_08_AM_45312836/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 17.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "second_order_ode_can_be_made_integrable"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 4y = 0$$

With initial conditions

$$[y(0) = 0, y'(0) = 1]$$

8.11.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 0$$

$$q(x) = 4$$

$$F = 0$$

Hence the ode is

$$y'' + 4y = 0$$

The domain of $p(x) = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = 4$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

8.11.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 0, C = 4$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 4e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 4 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 4$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(4)} \\ &= \pm 2i \end{aligned}$$

Hence

$$\lambda_1 = +2i$$

$$\lambda_2 = -2i$$

Which simplifies to

$$\lambda_1 = 2i$$

$$\lambda_2 = -2i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 2$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^0(c_1 \cos(2x) + c_2 \sin(2x))$$

Or

$$y = c_1 \cos(2x) + c_2 \sin(2x)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 \cos(2x) + c_2 \sin(2x) \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $x = 0$ in the above gives

$$0 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -2c_1 \sin(2x) + 2c_2 \cos(2x)$$

substituting $y' = 1$ and $x = 0$ in the above gives

$$1 = 2c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 0$$

$$c_2 = \frac{1}{2}$$

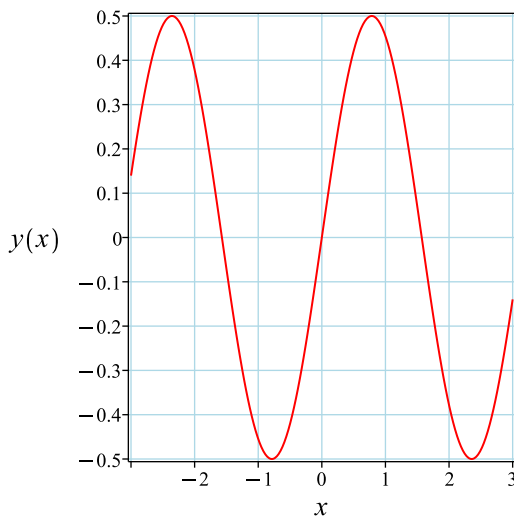
Substituting these values back in above solution results in

$$y = \frac{\sin(2x)}{2}$$

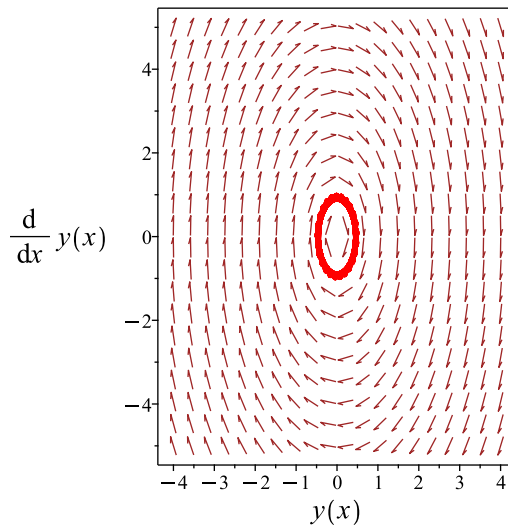
Summary

The solution(s) found are the following

$$y = \frac{\sin(2x)}{2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{\sin(2x)}{2}$$

Verified OK.

8.11.3 Solving as second order ode can be made integrable ode

Multiplying the ode by y' gives

$$y'y'' + 4yy' = 0$$

Integrating the above w.r.t x gives

$$\int (y'y'' + 4yy') dx = 0$$

$$\frac{y'^2}{2} + 2y^2 = c_2$$

Which is now solved for y . Solving the given ode for y' results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$y' = \sqrt{-4y^2 + 2c_1} \quad (1)$$

$$y' = -\sqrt{-4y^2 + 2c_1} \quad (2)$$

Now each one of the above ODE is solved.

Solving equation (1)

Integrating both sides gives

$$\int \frac{1}{\sqrt{-4y^2 + 2c_1}} dy = \int dx$$

$$\frac{\arctan\left(\frac{2y}{\sqrt{-4y^2 + 2c_1}}\right)}{2} = x + c_2$$

Solving equation (2)

Integrating both sides gives

$$\int -\frac{1}{\sqrt{-4y^2 + 2c_1}} dy = \int dx$$

$$-\frac{\arctan\left(\frac{2y}{\sqrt{-4y^2 + 2c_1}}\right)}{2} = c_3 + x$$

Initial conditions are used to solve for the constants of integration.

Looking at the First solution

$$\frac{\arctan\left(\frac{2y}{\sqrt{-4y^2 + 2c_1}}\right)}{2} = x + c_2 \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $x = 0$ in the above gives

$$0 = c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{(2 \tan(2x + 2c_2)^2 + 2) \sqrt{2} \sqrt{\frac{c_1}{\tan(2x + 2c_2)^2 + 1}}}{2} - \frac{\tan(2x + 2c_2)^2 \sqrt{2} c_1 (2 \tan(2x + 2c_2)^2 + 2)}{2 \sqrt{\frac{c_1}{\tan(2x + 2c_2)^2 + 1}} (\tan(2x + 2c_2)^2 + 1)^2}$$

substituting $y' = 1$ and $x = 0$ in the above gives

$$1 = \frac{\cos(2c_2)^2 \sqrt{2} c_1}{\sqrt{\cos(2c_2)^2 c_1}} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{1}{2}$$

$$c_2 = 0$$

Substituting these values back in above solution results in

$$\frac{\arctan\left(\frac{2y}{\sqrt{-4y^2+1}}\right)}{2} = x$$

Looking at the Second solution

$$-\frac{\arctan\left(\frac{2y}{\sqrt{-4y^2+2c_1}}\right)}{2} = c_3 + x \quad (2)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $x = 0$ in the above gives

$$0 = c_3 \quad (1A)$$

Taking derivative of the solution gives

$$y' = -\frac{(2 \tan(2c_3 + 2x)^2 + 2) \sqrt{2} \sqrt{\frac{c_1}{\tan(2c_3+2x)^2+1}}}{2} + \frac{\tan(2c_3 + 2x)^2 \sqrt{2} c_1 (2 \tan(2c_3 + 2x)^2 + 2)}{2 \sqrt{\frac{c_1}{\tan(2c_3+2x)^2+1}} (\tan(2c_3 + 2x)^2 + 1)^2}$$

substituting $y' = 1$ and $x = 0$ in the above gives

$$1 = -\frac{\cos(2c_3)^2 \sqrt{2} c_1}{\sqrt{\cos(2c_3)^2 c_1}} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_3\}$. Warning, unable to solve for constants of

Summary

The solution(s) found are the following integrations.

$$\frac{\arctan\left(\frac{2y}{\sqrt{-4y^2+1}}\right)}{2} = x \quad (1)$$

Verification of solutions

$$\frac{\arctan\left(\frac{2y}{\sqrt{-4y^2+1}}\right)}{2} = x$$

Verified OK.

8.11.4 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 4y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 0 \\ C &= 4 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-4}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -4 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -4z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 384: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -4$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(2x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned} y_1 &= z_1 \\ &= \cos(2x) \end{aligned}$$

Which simplifies to

$$y_1 = \cos(2x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Since $B = 0$ then the above becomes

$$\begin{aligned} y_2 &= y_1 \int \frac{1}{y_1^2} dx \\ &= \cos(2x) \int \frac{1}{\cos(2x)^2} dx \\ &= \cos(2x) \left(\frac{\tan(2x)}{2} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (\cos(2x)) + c_2 \left(\cos(2x) \left(\frac{\tan(2x)}{2} \right) \right) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 \cos(2x) + \frac{c_2 \sin(2x)}{2} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $x = 0$ in the above gives

$$0 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -2c_1 \sin(2x) + c_2 \cos(2x)$$

substituting $y' = 1$ and $x = 0$ in the above gives

$$1 = c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 0$$

$$c_2 = 1$$

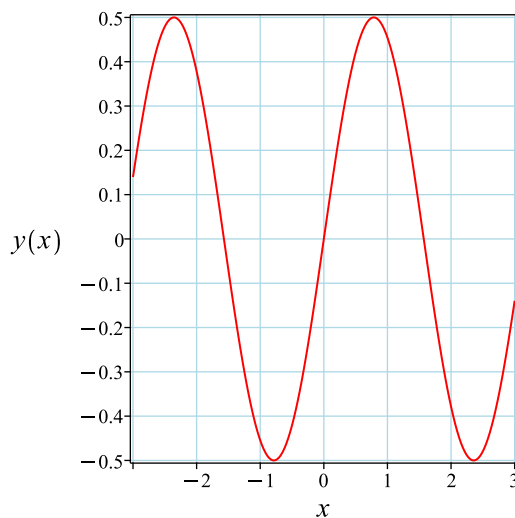
Substituting these values back in above solution results in

$$y = \frac{\sin(2x)}{2}$$

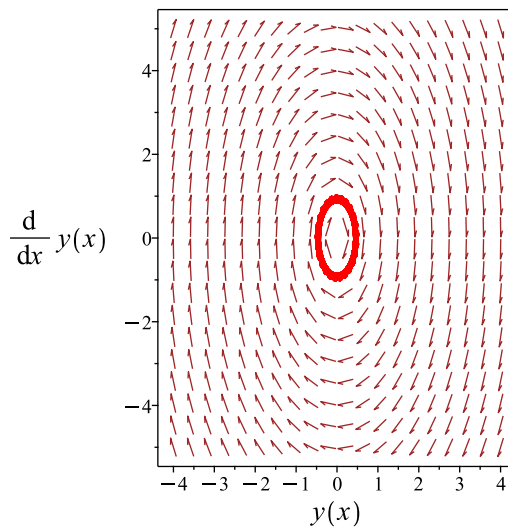
Summary

The solution(s) found are the following

$$y = \frac{\sin(2x)}{2} \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{\sin(2x)}{2}$$

Verified OK.

8.11.5 Maple step by step solution

Let's solve

$$\left[y'' + 4y = 0, y(0) = 0, y' \Big|_{\{x=0\}} = 1 \right]$$

- Highest derivative means the order of the ODE is 2
 y''
- Characteristic polynomial of ODE
 $r^2 + 4 = 0$
- Use quadratic formula to solve for r
 $r = \frac{0 \pm (\sqrt{-16})}{2}$
- Roots of the characteristic polynomial
 $r = (-2i, 2i)$
- 1st solution of the ODE
 $y_1(x) = \cos(2x)$
- 2nd solution of the ODE
 $y_2(x) = \sin(2x)$
- General solution of the ODE
 $y = c_1 y_1(x) + c_2 y_2(x)$
- Substitute in solutions
 $y = c_1 \cos(2x) + c_2 \sin(2x)$
- Check validity of solution $y = c_1 \cos(2x) + c_2 \sin(2x)$
 - Use initial condition $y(0) = 0$
 $0 = c_1$
 - Compute derivative of the solution
 $y' = -2c_1 \sin(2x) + 2c_2 \cos(2x)$

- Use the initial condition $y' \Big|_{\{x=0\}} = 1$

$$1 = 2c_2$$

- Solve for c_1 and c_2

$$\{c_1 = 0, c_2 = \frac{1}{2}\}$$

- Substitute constant values into general solution and simplify

$$y = \frac{\sin(2x)}{2}$$

- Solution to the IVP

$$y = \frac{\sin(2x)}{2}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 10

```
dsolve([diff(y(x),x$2)+ 4*y(x) = 0,y(0) = 0, D(y)(0) = 1],y(x), singsol=all)
```

$$y(x) = \frac{\sin(2x)}{2}$$

✓ Solution by Mathematica

Time used: 0.012 (sec). Leaf size: 10

```
DSolve[{y''[x]+4*y[x]==0,{y[0]==0,y'[0]==1}},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \sin(x) \cos(x)$$

8.12 problem 18

8.12.1 Existence and uniqueness analysis	2105
8.12.2 Solving as second order linear constant coeff ode	2106
8.12.3 Solving using Kovacic algorithm	2108
8.12.4 Maple step by step solution	2112

Internal problem ID [634]

Internal file name [OUTPUT/634_Sunday_June_05_2022_01_46_09_AM_49749350/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 18.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 4y' + 5y = 0$$

With initial conditions

$$[y(0) = 1, y'(0) = 0]$$

8.12.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 4$$

$$q(x) = 5$$

$$F = 0$$

Hence the ode is

$$y'' + 4y' + 5y = 0$$

The domain of $p(x) = 4$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = 5$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

8.12.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 4, C = 5$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 4\lambda e^{\lambda x} + 5e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 4\lambda + 5 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 4, C = 5$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-4}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{4^2 - (4)(1)(5)} \\ &= -2 \pm i \end{aligned}$$

Hence

$$\lambda_1 = -2 + i$$

$$\lambda_2 = -2 - i$$

Which simplifies to

$$\lambda_1 = -2 + i$$

$$\lambda_2 = -2 - i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -2$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-2x}(c_1 \cos(x) + c_2 \sin(x))$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = e^{-2x}(c_1 \cos(x) + c_2 \sin(x)) \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = 0$ in the above gives

$$1 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -2e^{-2x}(c_1 \cos(x) + c_2 \sin(x)) + e^{-2x}(-c_1 \sin(x) + c_2 \cos(x))$$

substituting $y' = 0$ and $x = 0$ in the above gives

$$0 = -2c_1 + c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 1$$

$$c_2 = 2$$

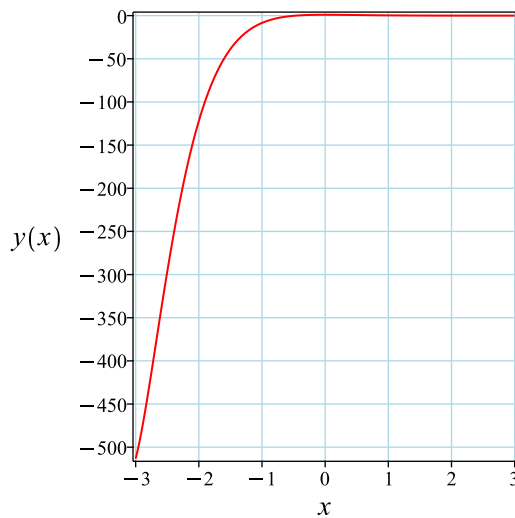
Substituting these values back in above solution results in

$$y = e^{-2x}(\cos(x) + 2 \sin(x))$$

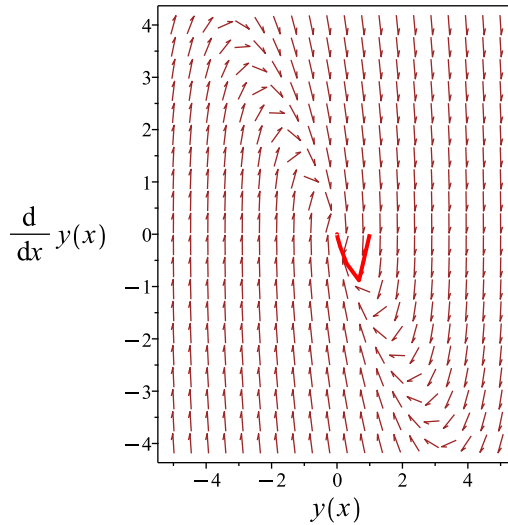
Summary

The solution(s) found are the following

$$y = e^{-2x}(\cos(x) + 2 \sin(x)) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{-2x}(\cos(x) + 2 \sin(x))$$

Verified OK.

8.12.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 4y' + 5y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 4 \quad (3)$$

$$C = 5$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = -1$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = -z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 386: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{4}{1} dx} \\ &= z_1 e^{-2x} \\ &= z_1 (e^{-2x})\end{aligned}$$

Which simplifies to

$$y_1 = e^{-2x} \cos(x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{4}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-4x}}{(y_1)^2} dx \\ &= y_1 (\tan(x))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-2x} \cos(x)) + c_2 (e^{-2x} \cos(x) (\tan(x)))\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-2x} \cos(x) + c_2 e^{-2x} \sin(x) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = 0$ in the above gives

$$1 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$y' = -2c_1 e^{-2x} \cos(x) - c_1 e^{-2x} \sin(x) - 2c_2 e^{-2x} \sin(x) + c_2 e^{-2x} \cos(x)$$

substituting $y' = 0$ and $x = 0$ in the above gives

$$0 = -2c_1 + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 1$$

$$c_2 = 2$$

Substituting these values back in above solution results in

$$y = e^{-2x} \cos(x) + 2 e^{-2x} \sin(x)$$

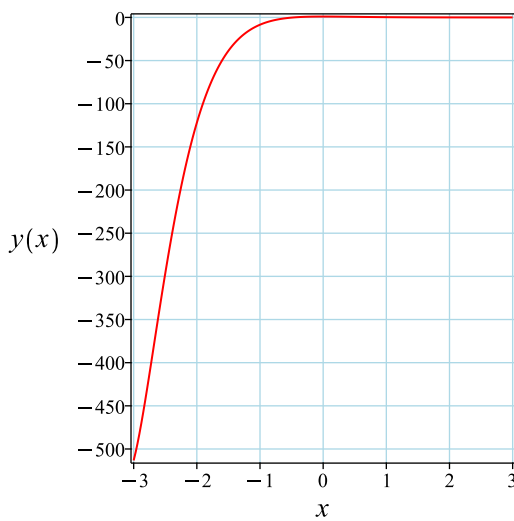
Which simplifies to

$$y = e^{-2x} (\cos(x) + 2 \sin(x))$$

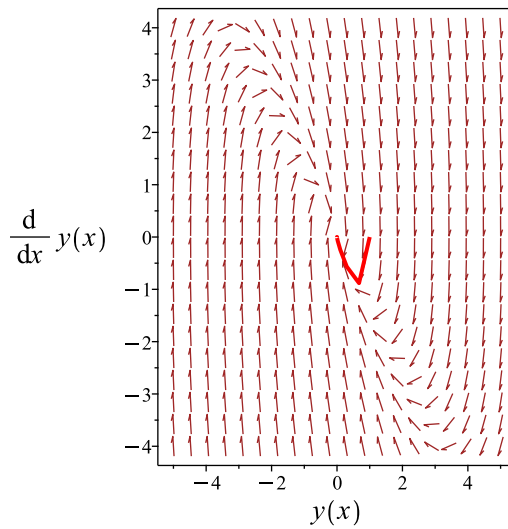
Summary

The solution(s) found are the following

$$y = e^{-2x} (\cos(x) + 2 \sin(x)) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{-2x}(\cos(x) + 2 \sin(x))$$

Verified OK.

8.12.4 Maple step by step solution

Let's solve

$$\left[y'' + 4y' + 5y = 0, y(0) = 1, y'|_{\{x=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 4r + 5 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-4) \pm (\sqrt{-4})}{2}$$

- Roots of the characteristic polynomial

$$r = (-2 - I, -2 + I)$$

- 1st solution of the ODE

$$y_1(x) = e^{-2x} \cos(x)$$

- 2nd solution of the ODE

$$y_2(x) = e^{-2x} \sin(x)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-2x} \cos(x) + c_2 e^{-2x} \sin(x)$$

- Check validity of solution $y = c_1 e^{-2x} \cos(x) + c_2 e^{-2x} \sin(x)$

- Use initial condition $y(0) = 1$

$$1 = c_1$$

- Compute derivative of the solution

$$y' = -2c_1 e^{-2x} \cos(x) - c_1 e^{-2x} \sin(x) - 2c_2 e^{-2x} \sin(x) + c_2 e^{-2x} \cos(x)$$

- Use the initial condition $y' \Big|_{\{x=0\}} = 0$

$$0 = -2c_1 + c_2$$

- Solve for c_1 and c_2

$$\{c_1 = 1, c_2 = 2\}$$

- Substitute constant values into general solution and simplify

$$y = e^{-2x}(\cos(x) + 2 \sin(x))$$

- Solution to the IVP

$$y = e^{-2x}(\cos(x) + 2 \sin(x))$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 16

```
dsolve([diff(y(x),x$2)+ 4*diff(y(x),x)+5*y(x) = 0,y(0) = 1, D(y)(0) = 0],y(x), singsol=all)
```

$$y(x) = e^{-2x}(2 \sin(x) + \cos(x))$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 18

```
DSolve[{y''[x]+4*y'[x]+5*y[x]==0,{y[0]==1,y'[0]==0}},y[x],x,IncludeSingularSolutions -> True
```

$$y(x) \rightarrow e^{-2x}(2 \sin(x) + \cos(x))$$

8.13 problem 19

8.13.1 Existence and uniqueness analysis	2115
8.13.2 Solving as second order linear constant coeff ode	2116
8.13.3 Solving using Kovacic algorithm	2118
8.13.4 Maple step by step solution	2122

Internal problem ID [635]

Internal file name [OUTPUT/635_Sunday_June_05_2022_01_46_10_AM_46194284/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 19.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - 2y' + 5y = 0$$

With initial conditions

$$\left[y\left(\frac{\pi}{2}\right) = 0, y'\left(\frac{\pi}{2}\right) = 2 \right]$$

8.13.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = -2$$

$$q(x) = 5$$

$$F = 0$$

Hence the ode is

$$y'' - 2y' + 5y = 0$$

The domain of $p(x) = -2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = \frac{\pi}{2}$ is inside this domain. The domain of $q(x) = 5$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = \frac{\pi}{2}$ is also inside this domain. Hence solution exists and is unique.

8.13.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -2, C = 5$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - 2\lambda e^{\lambda x} + 5 e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - 2\lambda + 5 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -2, C = 5$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-2^2 - (4)(1)(5)} \\ &= 1 \pm 2i \end{aligned}$$

Hence

$$\lambda_1 = 1 + 2i$$

$$\lambda_2 = 1 - 2i$$

Which simplifies to

$$\lambda_1 = 1 + 2i$$

$$\lambda_2 = 1 - 2i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 1$ and $\beta = 2$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^x(c_1 \cos(2x) + c_2 \sin(2x))$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = e^x(c_1 \cos(2x) + c_2 \sin(2x)) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $x = \frac{\pi}{2}$ in the above gives

$$0 = -c_1 e^{\frac{\pi}{2}} \quad (1A)$$

Taking derivative of the solution gives

$$y' = e^x(c_1 \cos(2x) + c_2 \sin(2x)) + e^x(-2c_1 \sin(2x) + 2c_2 \cos(2x))$$

substituting $y' = 2$ and $x = \frac{\pi}{2}$ in the above gives

$$2 = (-c_1 - 2c_2) e^{\frac{\pi}{2}} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 0$$

$$c_2 = -e^{-\frac{\pi}{2}}$$

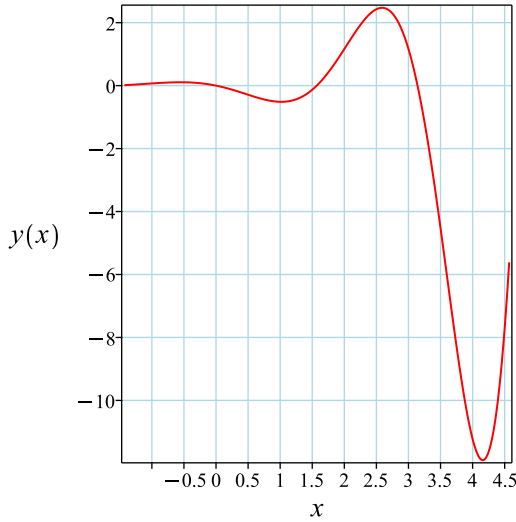
Substituting these values back in above solution results in

$$y = -\sin(2x) e^{-\frac{\pi}{2}+x}$$

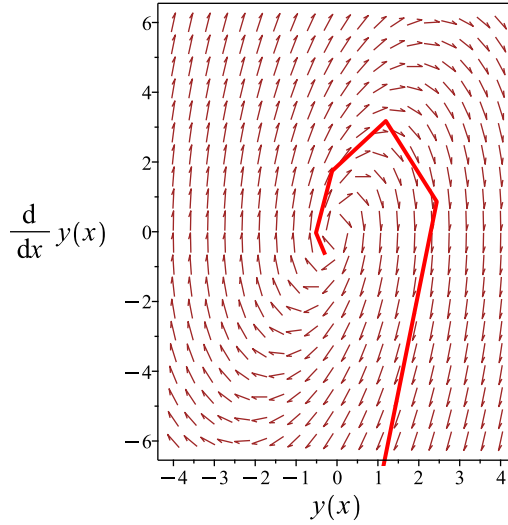
Summary

The solution(s) found are the following

$$y = -\sin(2x) e^{-\frac{\pi}{2}+x} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\sin(2x) e^{-\frac{\pi}{2}+x}$$

Verified OK.

8.13.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 2y' + 5y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = -2 \quad (3)$$

$$C = 5$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-4}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -4 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -4z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 388: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -4$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(2x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-2}{1} dx} \\ &= z_1 e^x \\ &= z_1 (e^x)\end{aligned}$$

Which simplifies to

$$y_1 = e^x \cos(2x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{2x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{\tan(2x)}{2} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^x \cos(2x)) + c_2 \left(e^x \cos(2x) \left(\frac{\tan(2x)}{2} \right) \right)\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^x \cos(2x) + \frac{e^x c_2 \sin(2x)}{2} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $x = \frac{\pi}{2}$ in the above gives

$$0 = -c_1 e^{\frac{\pi}{2}} \quad (1A)$$

Taking derivative of the solution gives

$$y' = c_1 e^x \cos(2x) - 2c_1 e^x \sin(2x) + \frac{e^x c_2 \sin(2x)}{2} + e^x c_2 \cos(2x)$$

substituting $y' = 2$ and $x = \frac{\pi}{2}$ in the above gives

$$2 = (-c_1 - c_2) e^{\frac{\pi}{2}} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}c_1 &= 0 \\ c_2 &= -2 e^{-\frac{\pi}{2}}\end{aligned}$$

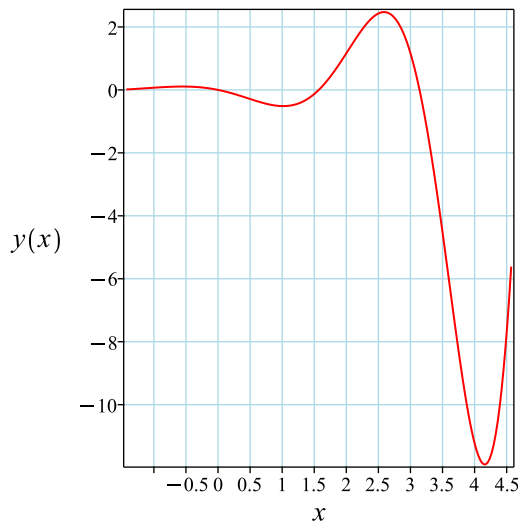
Substituting these values back in above solution results in

$$y = -\sin(2x) e^{-\frac{\pi}{2}+x}$$

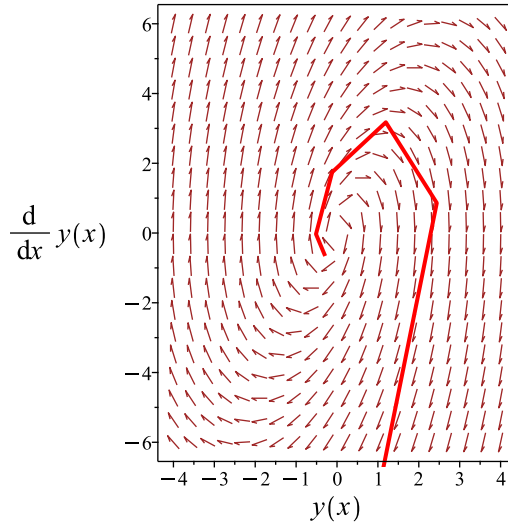
Summary

The solution(s) found are the following

$$y = -\sin(2x) e^{-\frac{\pi}{2}+x} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\sin(2x) e^{-\frac{\pi}{2}+x}$$

Verified OK.

8.13.4 Maple step by step solution

Let's solve

$$\left[y'' - 2y' + 5y = 0, y\left(\frac{\pi}{2}\right) = 0, y'\Big|_{\{x=\frac{\pi}{2}\}} = 2 \right]$$

- Highest derivative means the order of the ODE is 2
 y''
- Characteristic polynomial of ODE
 $r^2 - 2r + 5 = 0$
- Use quadratic formula to solve for r
 $r = \frac{2 \pm (\sqrt{-16})}{2}$
- Roots of the characteristic polynomial
 $r = (1 - 2I, 1 + 2I)$
- 1st solution of the ODE

$$y_1(x) = e^x \cos(2x)$$

- 2nd solution of the ODE

$$y_2(x) = e^x \sin(2x)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^x \cos(2x) + e^x c_2 \sin(2x)$$

- Check validity of solution $y = c_1 e^x \cos(2x) + e^x c_2 \sin(2x)$

- Use initial condition $y\left(\frac{\pi}{2}\right) = 0$

$$0 = -c_1 e^{\frac{\pi}{2}}$$

- Compute derivative of the solution

$$y' = c_1 e^x \cos(2x) - 2c_1 e^x \sin(2x) + e^x c_2 \sin(2x) + 2e^x c_2 \cos(2x)$$

- Use the initial condition $y' \Big|_{\{x=\frac{\pi}{2}\}} = 2$

$$2 = -c_1 e^{\frac{\pi}{2}} - 2e^{\frac{\pi}{2}} c_2$$

- Solve for c_1 and c_2

$$\left\{ c_1 = 0, c_2 = -\frac{1}{e^{\frac{\pi}{2}}} \right\}$$

- Substitute constant values into general solution and simplify

$$y = -\sin(2x) e^{-\frac{\pi}{2}+x}$$

- Solution to the IVP

$$y = -\sin(2x) e^{-\frac{\pi}{2}+x}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 16

```
dsolve([diff(y(x),x$2)- 2*diff(y(x),x)+5*y(x) = 0,y(1/2*Pi) = 0, D(y)(1/2*Pi) = 2],y(x), sin
```

$$y(x) = -\sin(2x)e^{-\frac{\pi}{2}+x}$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 20

```
DSolve[{y''[x]-2*y'[x]+5*y[x]==0,{y[Pi/2]==0,y'[Pi/2]==2}},y[x],x,IncludeSingularSolutions -
```

$$y(x) \rightarrow -e^{x-\frac{\pi}{2}} \sin(2x)$$

8.14 problem 20

8.14.1 Existence and uniqueness analysis	2125
8.14.2 Solving as second order linear constant coeff ode	2126
8.14.3 Solving as second order ode can be made integrable ode	2129
8.14.4 Solving using Kovacic algorithm	2131
8.14.5 Maple step by step solution	2135

Internal problem ID [636]

Internal file name [OUTPUT/636_Sunday_June_05_2022_01_46_11_AM_56093143/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 20.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "second_order_ode_can_be_made_integrable"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + y = 0$$

With initial conditions

$$\left[y\left(\frac{\pi}{3}\right) = 2, y'\left(\frac{\pi}{3}\right) = -4 \right]$$

8.14.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 0$$

$$q(x) = 1$$

$$F = 0$$

Hence the ode is

$$y'' + y = 0$$

The domain of $p(x) = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = \frac{\pi}{3}$ is inside this domain. The domain of $q(x) = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = \frac{\pi}{3}$ is also inside this domain. Hence solution exists and is unique.

8.14.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 0, C = 1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(1)} \\ &= \pm i \end{aligned}$$

Hence

$$\lambda_1 = +i$$

$$\lambda_2 = -i$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= i \\ \lambda_2 &= -i\end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^0(c_1 \cos(x) + c_2 \sin(x))$$

Or

$$y = c_1 \cos(x) + c_2 \sin(x)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 \cos(x) + c_2 \sin(x) \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = \frac{\pi}{3}$ in the above gives

$$2 = \frac{c_1}{2} + \frac{\sqrt{3}c_2}{2} \tag{1A}$$

Taking derivative of the solution gives

$$y' = -c_1 \sin(x) + c_2 \cos(x)$$

substituting $y' = -4$ and $x = \frac{\pi}{3}$ in the above gives

$$-4 = -\frac{\sqrt{3}c_1}{2} + \frac{c_2}{2} \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 2\sqrt{3} + 1$$

$$c_2 = \sqrt{3} - 2$$

Substituting these values back in above solution results in

$$y = 2 \cos(x) \sqrt{3} + \sin(x) \sqrt{3} + \cos(x) - 2 \sin(x)$$

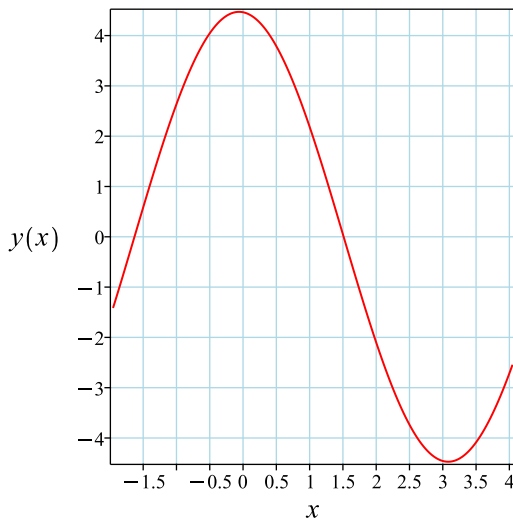
Which simplifies to

$$y = (2 \cos(x) + \sin(x)) \sqrt{3} + \cos(x) - 2 \sin(x)$$

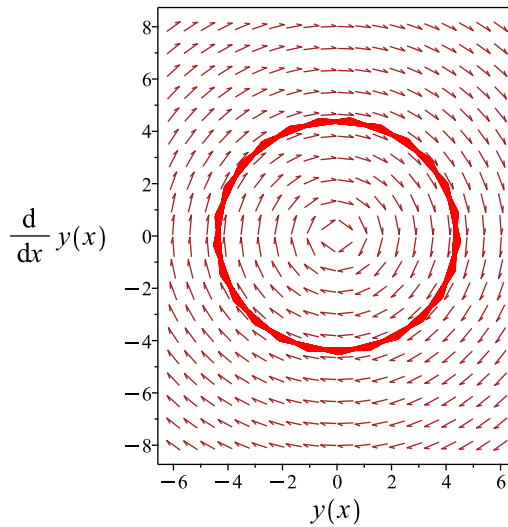
Summary

The solution(s) found are the following

$$y = (2 \cos(x) + \sin(x)) \sqrt{3} + \cos(x) - 2 \sin(x) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = (2 \cos(x) + \sin(x)) \sqrt{3} + \cos(x) - 2 \sin(x)$$

Verified OK.

8.14.3 Solving as second order ode can be made integrable ode

Multiplying the ode by y' gives

$$y'y'' + yy' = 0$$

Integrating the above w.r.t x gives

$$\int (y'y'' + yy') dx = 0$$
$$\frac{y'^2}{2} + \frac{y^2}{2} = c_2$$

Which is now solved for y . Solving the given ode for y' results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$y' = \sqrt{-y^2 + 2c_1} \quad (1)$$

$$y' = -\sqrt{-y^2 + 2c_1} \quad (2)$$

Now each one of the above ODE is solved.

Solving equation (1)

Integrating both sides gives

$$\int \frac{1}{\sqrt{-y^2 + 2c_1}} dy = \int dx$$
$$\arctan\left(\frac{y}{\sqrt{-y^2 + 2c_1}}\right) = x + c_2$$

Solving equation (2)

Integrating both sides gives

$$\int -\frac{1}{\sqrt{-y^2 + 2c_1}} dy = \int dx$$
$$-\arctan\left(\frac{y}{\sqrt{-y^2 + 2c_1}}\right) = c_3 + x$$

Initial conditions are used to solve for the constants of integration.

Looking at the First solution

$$\arctan\left(\frac{y}{\sqrt{-y^2 + 2c_1}}\right) = x + c_2 \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = \frac{\pi}{3}$ in the above gives

$$\arctan\left(\frac{2}{\sqrt{-4+2c_1}}\right) = \frac{\pi}{3} + c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = (\tan(x+c_2)^2 + 1) \sqrt{2} \sqrt{\frac{c_1}{\tan(x+c_2)^2 + 1}} - \frac{\tan(x+c_2)^2 \sqrt{2} c_1}{\sqrt{\frac{c_1}{\tan(x+c_2)^2 + 1}} (\tan(x+c_2)^2 + 1)}$$

substituting $y' = -4$ and $x = \frac{\pi}{3}$ in the above gives

$$-4 = \frac{\cos\left(\frac{\pi}{3} + c_2\right)^2 \sqrt{2} c_1}{\sqrt{\cos\left(\frac{\pi}{3} + c_2\right)^2 c_1}} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. There is no solution for the constants of integrations. This solution is removed.

Looking at the Second solution

$$-\arctan\left(\frac{y}{\sqrt{-y^2+2c_1}}\right) = c_3 + x \quad (2)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = \frac{\pi}{3}$ in the above gives

$$-\arctan\left(\frac{2}{\sqrt{-4+2c_1}}\right) = c_3 + \frac{\pi}{3} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -(\tan(c_3+x)^2 + 1) \sqrt{2} \sqrt{\frac{c_1}{\tan(c_3+x)^2 + 1}} + \frac{\tan(c_3+x)^2 \sqrt{2} c_1}{\sqrt{\frac{c_1}{\tan(c_3+x)^2 + 1}} (\tan(c_3+x)^2 + 1)}$$

substituting $y' = -4$ and $x = \frac{\pi}{3}$ in the above gives

$$-4 = -\frac{\cos\left(c_3 + \frac{\pi}{3}\right)^2 \sqrt{2} c_1}{\sqrt{\cos\left(c_3 + \frac{\pi}{3}\right)^2 c_1}} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_3\}$. Solving for the constants gives

$$c_1 = 10$$

$$c_3 = -\arctan\left(\frac{1}{2}\right) - \frac{\pi}{3}$$

Substituting these values back in above solution results in

$$-\arctan\left(\frac{y}{\sqrt{-y^2+20}}\right) = -\arctan\left(\frac{1}{2}\right) - \frac{\pi}{3} + x$$

Summary

The solution(s) found are the following

$$-\arctan\left(\frac{y}{\sqrt{-y^2+20}}\right) = -\arctan\left(\frac{1}{2}\right) - \frac{\pi}{3} + x \quad (1)$$

Verification of solutions

$$-\arctan\left(\frac{y}{\sqrt{-y^2+20}}\right) = -\arctan\left(\frac{1}{2}\right) - \frac{\pi}{3} + x$$

Verified OK.

8.14.4 Solving using Kovacic algorithm

Writing the ode as

$$y'' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 0 \quad (3)$$

$$C = 1$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \tag{5}$$

$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = -1$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = -z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 390: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}y_1 &= z_1 \\ &= \cos(x)\end{aligned}$$

Which simplifies to

$$y_1 = \cos(x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Since $B = 0$ then the above becomes

$$\begin{aligned}y_2 &= y_1 \int \frac{1}{y_1^2} dx \\ &= \cos(x) \int \frac{1}{\cos(x)^2} dx \\ &= \cos(x) (\tan(x))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1(\cos(x)) + c_2(\cos(x)(\tan(x)))\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 \cos(x) + c_2 \sin(x) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = \frac{\pi}{3}$ in the above gives

$$2 = \frac{c_1}{2} + \frac{\sqrt{3}c_2}{2} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -c_1 \sin(x) + c_2 \cos(x)$$

substituting $y' = -4$ and $x = \frac{\pi}{3}$ in the above gives

$$-4 = -\frac{\sqrt{3}c_1}{2} + \frac{c_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}c_1 &= 2\sqrt{3} + 1 \\ c_2 &= \sqrt{3} - 2\end{aligned}$$

Substituting these values back in above solution results in

$$y = 2 \cos(x) \sqrt{3} + \sin(x) \sqrt{3} + \cos(x) - 2 \sin(x)$$

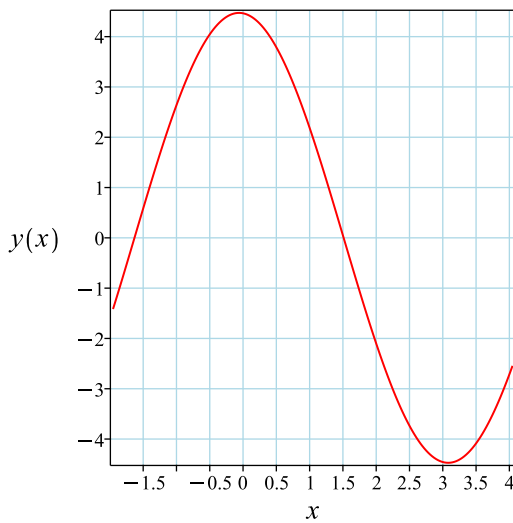
Which simplifies to

$$y = (2 \cos(x) + \sin(x)) \sqrt{3} + \cos(x) - 2 \sin(x)$$

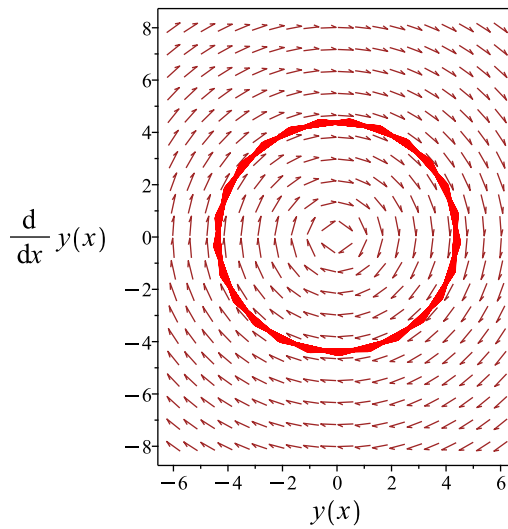
Summary

The solution(s) found are the following

$$y = (2 \cos(x) + \sin(x)) \sqrt{3} + \cos(x) - 2 \sin(x) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = (2 \cos(x) + \sin(x)) \sqrt{3} + \cos(x) - 2 \sin(x)$$

Verified OK.

8.14.5 Maple step by step solution

Let's solve

$$\left[y'' + y = 0, y\left(\frac{\pi}{3}\right) = 2, y' \Big|_{\{x=\frac{\pi}{3}\}} = -4 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 1 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{0 \pm (\sqrt{-4})}{2}$$

- Roots of the characteristic polynomial

$$r = (-I, I)$$

- 1st solution of the ODE

$$y_1(x) = \cos(x)$$

- 2nd solution of the ODE

$$y_2(x) = \sin(x)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 \cos(x) + c_2 \sin(x)$$

- Check validity of solution $y = c_1 \cos(x) + c_2 \sin(x)$

- Use initial condition $y\left(\frac{\pi}{3}\right) = 2$

$$2 = \frac{c_1}{2} + \frac{\sqrt{3}c_2}{2}$$

- Compute derivative of the solution

$$y' = -c_1 \sin(x) + c_2 \cos(x)$$

- Use the initial condition $y' \Big|_{\{x=\frac{\pi}{3}\}} = -4$

$$-4 = -\frac{\sqrt{3}c_1}{2} + \frac{c_2}{2}$$

- Solve for c_1 and c_2

$$\{c_1 = 2\sqrt{3} + 1, c_2 = \sqrt{3} - 2\}$$

- Substitute constant values into general solution and simplify

$$y = (2 \cos(x) + \sin(x)) \sqrt{3} + \cos(x) - 2 \sin(x)$$

- Solution to the IVP

$$y = (2 \cos(x) + \sin(x)) \sqrt{3} + \cos(x) - 2 \sin(x)$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 23

```
dsolve([diff(y(x),x$2)+y(x) = 0,y(1/3*Pi) = 2, D(y)(1/3*Pi) = -4],y(x), singsol=all)
```

$$y(x) = (\sin(x) + 2 \cos(x)) \sqrt{3} + \cos(x) - 2 \sin(x)$$

✓ Solution by Mathematica

Time used: 0.012 (sec). Leaf size: 28

```
DSolve[{y'[x]+y[x]==0,{y[Pi/3]==2,y'[Pi/3]==-4}},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow (\sqrt{3} - 2) \sin(x) + (1 + 2\sqrt{3}) \cos(x)$$

8.15 problem 21

8.15.1 Existence and uniqueness analysis	2138
8.15.2 Solving as second order linear constant coeff ode	2139
8.15.3 Solving using Kovacic algorithm	2142
8.15.4 Maple step by step solution	2146

Internal problem ID [637]

Internal file name [OUTPUT/637_Sunday_June_05_2022_01_46_12_AM_12825249/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 21.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + y' + \frac{5y}{4} = 0$$

With initial conditions

$$[y(0) = 3, y'(0) = 1]$$

8.15.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 1$$

$$q(x) = \frac{5}{4}$$

$$F = 0$$

Hence the ode is

$$y'' + y' + \frac{5y}{4} = 0$$

The domain of $p(x) = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = \frac{5}{4}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

8.15.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 1, C = \frac{5}{4}$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + \lambda e^{\lambda x} + \frac{5 e^{\lambda x}}{4} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + \lambda + \frac{5}{4} = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 1, C = \frac{5}{4}$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{1^2 - (4)(1)\left(\frac{5}{4}\right)} \\ &= -\frac{1}{2} \pm i \end{aligned}$$

Hence

$$\lambda_1 = -\frac{1}{2} + i$$
$$\lambda_2 = -\frac{1}{2} - i$$

Which simplifies to

$$\lambda_1 = -\frac{1}{2} + i$$
$$\lambda_2 = -\frac{1}{2} - i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -\frac{1}{2}$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-\frac{x}{2}}(c_1 \cos(x) + c_2 \sin(x))$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = e^{-\frac{x}{2}}(c_1 \cos(x) + c_2 \sin(x)) \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 3$ and $x = 0$ in the above gives

$$3 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -\frac{e^{-\frac{x}{2}}(c_1 \cos(x) + c_2 \sin(x))}{2} + e^{-\frac{x}{2}}(-c_1 \sin(x) + c_2 \cos(x))$$

substituting $y' = 1$ and $x = 0$ in the above gives

$$1 = -\frac{c_1}{2} + c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 3$$
$$c_2 = \frac{5}{2}$$

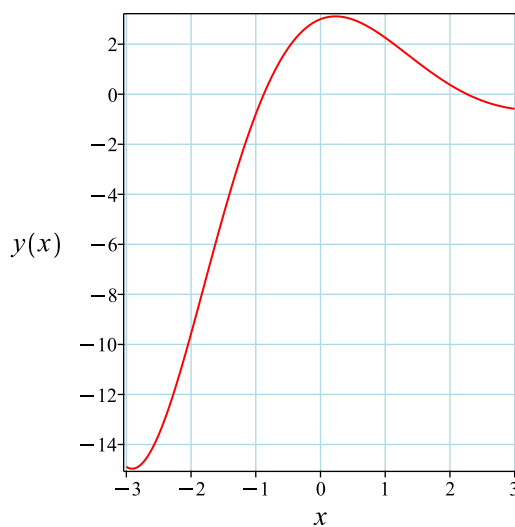
Substituting these values back in above solution results in

$$y = \frac{e^{-\frac{x}{2}}(6 \cos(x) + 5 \sin(x))}{2}$$

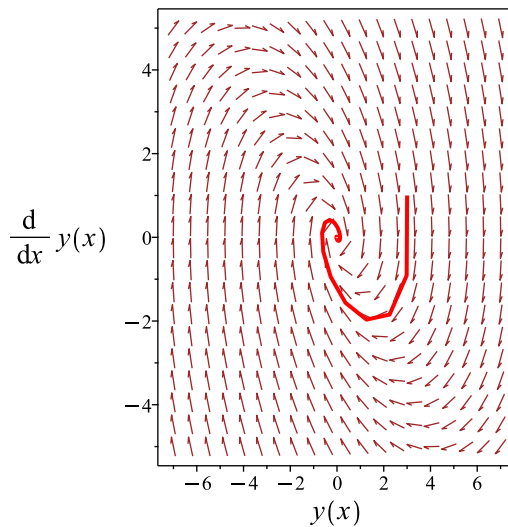
Summary

The solution(s) found are the following

$$y = \frac{e^{-\frac{x}{2}}(6 \cos(x) + 5 \sin(x))}{2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^{-\frac{x}{2}}(6 \cos(x) + 5 \sin(x))}{2}$$

Verified OK.

8.15.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + y' + \frac{5y}{4} = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 1 \\ C &= \frac{5}{4} \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 392: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{1}{1} dx} \end{aligned}$$

$$\begin{aligned}
&= z_1 e^{-\frac{x}{2}} \\
&= z_1 (e^{-\frac{x}{2}})
\end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{x}{2}} \cos(x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}
y_2 &= y_1 \int \frac{e^{\int -\frac{1}{1} dx}}{(y_1)^2} dx \\
&= y_1 \int \frac{e^{-x}}{(y_1)^2} dx \\
&= y_1 (\tan(x))
\end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 (e^{-\frac{x}{2}} \cos(x)) + c_2 (e^{-\frac{x}{2}} \cos(x) (\tan(x)))
\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-\frac{x}{2}} \cos(x) + c_2 e^{-\frac{x}{2}} \sin(x) \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 3$ and $x = 0$ in the above gives

$$3 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -\frac{c_1 e^{-\frac{x}{2}} \cos(x)}{2} - c_1 e^{-\frac{x}{2}} \sin(x) - \frac{c_2 e^{-\frac{x}{2}} \sin(x)}{2} + c_2 e^{-\frac{x}{2}} \cos(x)$$

substituting $y' = 1$ and $x = 0$ in the above gives

$$1 = -\frac{c_1}{2} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 3$$
$$c_2 = \frac{5}{2}$$

Substituting these values back in above solution results in

$$y = \frac{5 e^{-\frac{x}{2}} \sin(x)}{2} + 3 e^{-\frac{x}{2}} \cos(x)$$

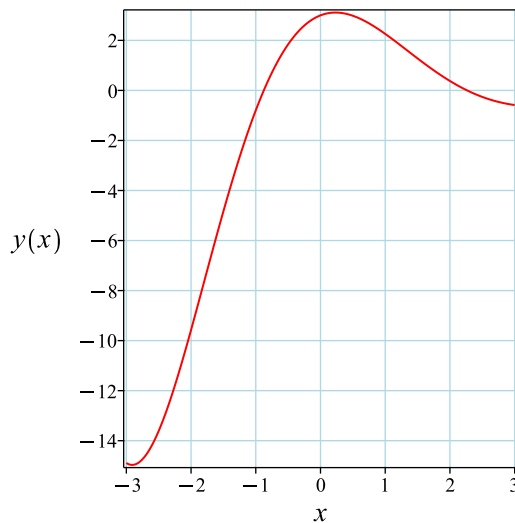
Which simplifies to

$$y = \frac{e^{-\frac{x}{2}} (6 \cos(x) + 5 \sin(x))}{2}$$

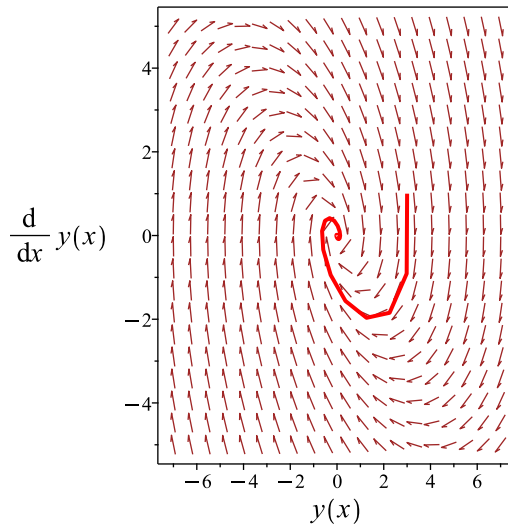
Summary

The solution(s) found are the following

$$y = \frac{e^{-\frac{x}{2}} (6 \cos(x) + 5 \sin(x))}{2} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \frac{e^{-\frac{x}{2}} (6 \cos(x) + 5 \sin(x))}{2}$$

Verified OK.

8.15.4 Maple step by step solution

Let's solve

$$\left[y'' + y' + \frac{5y}{4} = 0, y(0) = 3, y' \Big|_{\{x=0\}} = 1 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + r + \frac{5}{4} = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-1) \pm (\sqrt{-4})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{2} - I, -\frac{1}{2} + I\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{x}{2}} \cos(x)$$

- 2nd solution of the ODE

$$y_2(x) = e^{-\frac{x}{2}} \sin(x)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-\frac{x}{2}} \cos(x) + c_2 e^{-\frac{x}{2}} \sin(x)$$

- Check validity of solution $y = c_1 e^{-\frac{x}{2}} \cos(x) + c_2 e^{-\frac{x}{2}} \sin(x)$

- Use initial condition $y(0) = 3$

$$3 = c_1$$

- Compute derivative of the solution

$$y' = -\frac{c_1 e^{-\frac{x}{2}} \cos(x)}{2} - c_1 e^{-\frac{x}{2}} \sin(x) - \frac{c_2 e^{-\frac{x}{2}} \sin(x)}{2} + c_2 e^{-\frac{x}{2}} \cos(x)$$

- Use the initial condition $y' \Big|_{\{x=0\}} = 1$

$$1 = -\frac{c_1}{2} + c_2$$

- Solve for c_1 and c_2

$$\{c_1 = 3, c_2 = \frac{5}{2}\}$$

- Substitute constant values into general solution and simplify

$$y = \frac{e^{-\frac{x}{2}}(6 \cos(x) + 5 \sin(x))}{2}$$

- Solution to the IVP

$$y = \frac{e^{-\frac{x}{2}}(6 \cos(x) + 5 \sin(x))}{2}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 19

```
dsolve([diff(y(x),x$2)+ diff(y(x),x)+125/100*y(x) = 0,y(0) = 3, D(y)(0) = 1],y(x), singsol=a
```

$$y(x) = \frac{e^{-\frac{x}{2}}(5 \sin(x) + 6 \cos(x))}{2}$$

✓ Solution by Mathematica

Time used: 0.016 (sec). Leaf size: 25

```
DSolve[{y''[x]+y'[x]+125/100*y[x]==0,{y[0]==3,y'[0]==1}},y[x],x,IncludeSingularSolutions ->
```

$$y(x) \rightarrow \frac{1}{2}e^{-x/2}(5 \sin(x) + 6 \cos(x))$$

8.16 problem 22

8.16.1 Existence and uniqueness analysis	2148
8.16.2 Solving as second order linear constant coeff ode	2149
8.16.3 Solving using Kovacic algorithm	2152
8.16.4 Maple step by step solution	2156

Internal problem ID [638]

Internal file name [OUTPUT/638_Sunday_June_05_2022_01_46_13_AM_47679787/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 22.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$y'' + 2y' + 2y = 0$$

With initial conditions

$$\left[y\left(\frac{\pi}{4}\right) = 2, y'\left(\frac{\pi}{4}\right) = -2 \right]$$

8.16.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 2$$

$$q(x) = 2$$

$$F = 0$$

Hence the ode is

$$y'' + 2y' + 2y = 0$$

The domain of $p(x) = 2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = \frac{\pi}{4}$ is inside this domain. The domain of $q(x) = 2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = \frac{\pi}{4}$ is also inside this domain. Hence solution exists and is unique.

8.16.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 2, C = 2$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 2\lambda e^{\lambda x} + 2e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 2\lambda + 2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 2, C = 2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{2^2 - (4)(1)(2)} \\ &= -1 \pm i \end{aligned}$$

Hence

$$\lambda_1 = -1 + i$$

$$\lambda_2 = -1 - i$$

Which simplifies to

$$\lambda_1 = -1 + i$$

$$\lambda_2 = -1 - i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -1$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-x}(c_1 \cos(x) + c_2 \sin(x))$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = e^{-x}(c_1 \cos(x) + c_2 \sin(x)) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = \frac{\pi}{4}$ in the above gives

$$2 = \frac{(c_1 + c_2)\sqrt{2}e^{-\frac{\pi}{4}}}{2} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -e^{-x}(c_1 \cos(x) + c_2 \sin(x)) + e^{-x}(-c_1 \sin(x) + c_2 \cos(x))$$

substituting $y' = -2$ and $x = \frac{\pi}{4}$ in the above gives

$$-2 = -\sqrt{2}e^{-\frac{\pi}{4}}c_1 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \sqrt{2}e^{\frac{\pi}{4}}$$

$$c_2 = \sqrt{2}e^{\frac{\pi}{4}}$$

Substituting these values back in above solution results in

$$y = \sqrt{2} \cos(x) e^{\frac{\pi}{4}-x} + \sqrt{2} \sin(x) e^{\frac{\pi}{4}-x}$$

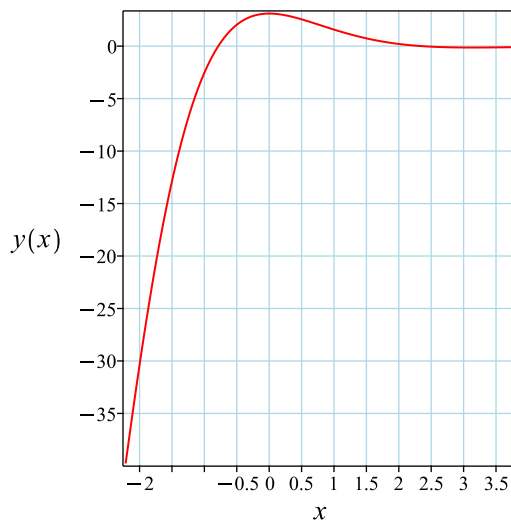
Which simplifies to

$$y = \sqrt{2} e^{\frac{\pi}{4}-x} (\cos(x) + \sin(x))$$

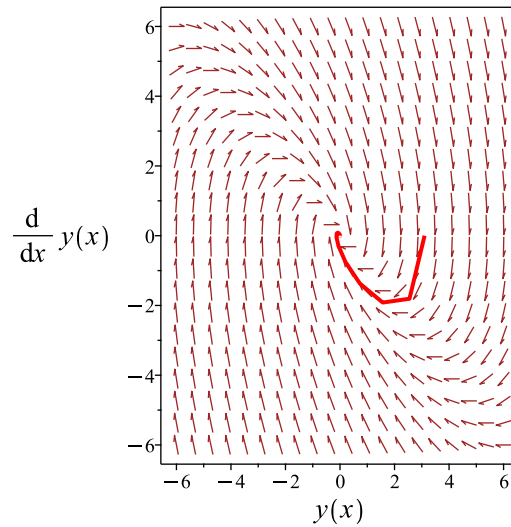
Summary

The solution(s) found are the following

$$y = \sqrt{2} e^{\frac{\pi}{4}-x} (\cos(x) + \sin(x)) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \sqrt{2} e^{\frac{\pi}{4}-x} (\cos(x) + \sin(x))$$

Verified OK.

8.16.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 2y' + 2y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 2 \\ C &= 2 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 394: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{2}{1} dx} \end{aligned}$$

$$\begin{aligned}
&= z_1 e^{-x} \\
&= z_1 (e^{-x})
\end{aligned}$$

Which simplifies to

$$y_1 = \cos(x) e^{-x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}
y_2 &= y_1 \int \frac{e^{\int -\frac{2}{1} dx}}{(y_1)^2} dx \\
&= y_1 \int \frac{e^{-2x}}{(y_1)^2} dx \\
&= y_1 (\tan(x))
\end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 (\cos(x) e^{-x}) + c_2 (\cos(x) e^{-x} (\tan(x)))
\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 \cos(x) e^{-x} + c_2 \sin(x) e^{-x} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = \frac{\pi}{4}$ in the above gives

$$2 = \frac{(c_1 + c_2) \sqrt{2} e^{-\frac{\pi}{4}}}{2} \quad (1A)$$

Taking derivative of the solution gives

$$y' = -c_1 \sin(x) e^{-x} - c_1 \cos(x) e^{-x} + c_2 \cos(x) e^{-x} - c_2 \sin(x) e^{-x}$$

substituting $y' = -2$ and $x = \frac{\pi}{4}$ in the above gives

$$-2 = -\sqrt{2} e^{-\frac{\pi}{4}} c_1 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \sqrt{2} e^{\frac{\pi}{4}}$$

$$c_2 = \sqrt{2} e^{\frac{\pi}{4}}$$

Substituting these values back in above solution results in

$$y = \sqrt{2} \cos(x) e^{\frac{\pi}{4}-x} + \sqrt{2} \sin(x) e^{\frac{\pi}{4}-x}$$

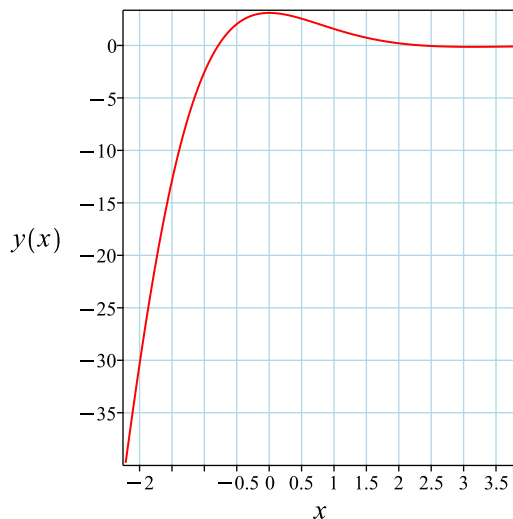
Which simplifies to

$$y = \sqrt{2} e^{\frac{\pi}{4}-x} (\cos(x) + \sin(x))$$

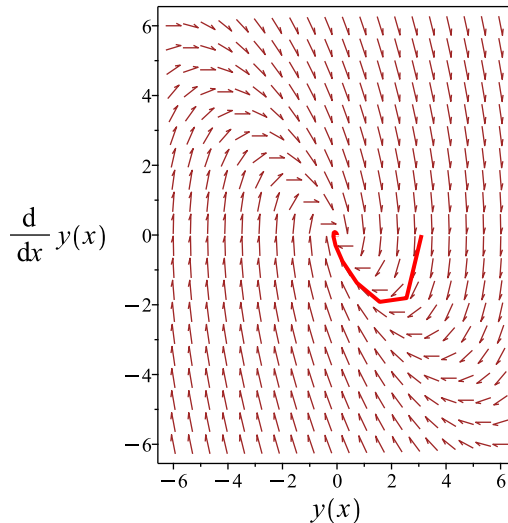
Summary

The solution(s) found are the following

$$y = \sqrt{2} e^{\frac{\pi}{4}-x} (\cos(x) + \sin(x)) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = \sqrt{2} e^{\frac{\pi}{4}-x} (\cos(x) + \sin(x))$$

Verified OK.

8.16.4 Maple step by step solution

Let's solve

$$\left[y'' + 2y' + 2y = 0, y\left(\frac{\pi}{4}\right) = 2, y'\Big|_{\{x=\frac{\pi}{4}\}} = -2 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 2r + 2 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-2) \pm (\sqrt{-4})}{2}$$

- Roots of the characteristic polynomial

$$r = (-1 - I, -1 + I)$$

- 1st solution of the ODE

$$y_1(x) = \cos(x) e^{-x}$$

- 2nd solution of the ODE

$$y_2(x) = \sin(x) e^{-x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 \cos(x) e^{-x} + c_2 \sin(x) e^{-x}$$

- Check validity of solution $y = c_1 \cos(x) e^{-x} + c_2 \sin(x) e^{-x}$

- Use initial condition $y\left(\frac{\pi}{4}\right) = 2$

$$2 = \frac{\sqrt{2} e^{-\frac{\pi}{4}} c_1}{2} + \frac{\sqrt{2} e^{-\frac{\pi}{4}} c_2}{2}$$

- Compute derivative of the solution

$$y' = -c_1 \sin(x) e^{-x} - c_1 \cos(x) e^{-x} + c_2 \cos(x) e^{-x} - c_2 \sin(x) e^{-x}$$

- Use the initial condition $y'\Big|_{\{x=\frac{\pi}{4}\}} = -2$

$$-2 = -\sqrt{2} e^{-\frac{\pi}{4}} c_1$$

- Solve for c_1 and c_2

$$\left\{ c_1 = \frac{\sqrt{2}}{e^{-\frac{\pi}{4}}}, c_2 = \frac{\sqrt{2}}{e^{-\frac{\pi}{4}}} \right\}$$

- Substitute constant values into general solution and simplify

$$y = \sqrt{2} e^{\frac{\pi}{4}-x} (\cos(x) + \sin(x))$$

- Solution to the IVP

$$y = \sqrt{2} e^{\frac{\pi}{4}-x} (\cos(x) + \sin(x))$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 21

```
dsolve([diff(y(x),x$2)+ 2*diff(y(x),x)+2*y(x) = 0,y(1/4*Pi) = 2, D(y)(1/4*Pi) = -2],y(x), si
```

$$y(x) = \sqrt{2} e^{-x+\frac{\pi}{4}} (\sin(x) + \cos(x))$$

✓ Solution by Mathematica

Time used: 0.016 (sec). Leaf size: 27

```
DSolve[{y''[x]+2*y'[x]+2*y[x]==0,{y[Pi/4]==2,y'[Pi/4]==-2}},y[x],x,IncludeSingularSolutions
```

$$y(x) \rightarrow \sqrt{2} e^{\frac{\pi}{4}-x} (\sin(x) + \cos(x))$$

8.17 problem 23

8.17.1 Existence and uniqueness analysis	2158
8.17.2 Solving as second order linear constant coeff ode	2159
8.17.3 Solving using Kovacic algorithm	2162
8.17.4 Maple step by step solution	2166

Internal problem ID [639]

Internal file name [OUTPUT/639_Sunday_June_05_2022_01_46_15_AM_93629760/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 23.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$u'' - u' + 2u = 0$$

With initial conditions

$$[u(0) = 2, u'(0) = 0]$$

8.17.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$u'' + p(x)u' + q(x)u = F$$

Where here

$$p(x) = -1$$

$$q(x) = 2$$

$$F = 0$$

Hence the ode is

$$u'' - u' + 2u = 0$$

The domain of $p(x) = -1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = 2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

8.17.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Au''(x) + Bu'(x) + Cu(x) = 0$$

Where in the above $A = 1, B = -1, C = 2$. Let the solution be $u = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - \lambda e^{\lambda x} + 2e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - \lambda + 2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -1, C = 2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-1^2 - (4)(1)(2)} \\ &= \frac{1}{2} \pm \frac{i\sqrt{7}}{2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= \frac{1}{2} + \frac{i\sqrt{7}}{2} \\ \lambda_2 &= \frac{1}{2} - \frac{i\sqrt{7}}{2} \end{aligned}$$

Which simplifies to

$$\lambda_1 = \frac{1}{2} + \frac{i\sqrt{7}}{2}$$
$$\lambda_2 = \frac{1}{2} - \frac{i\sqrt{7}}{2}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = \frac{1}{2}$ and $\beta = \frac{\sqrt{7}}{2}$. Therefore the final solution, when using Euler relation, can be written as

$$u = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$u = e^{\frac{x}{2}} \left(c_1 \cos \left(\frac{\sqrt{7}x}{2} \right) + c_2 \sin \left(\frac{\sqrt{7}x}{2} \right) \right)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = e^{\frac{x}{2}} \left(c_1 \cos \left(\frac{\sqrt{7}x}{2} \right) + c_2 \sin \left(\frac{\sqrt{7}x}{2} \right) \right) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 2$ and $x = 0$ in the above gives

$$2 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$u' = \frac{e^{\frac{x}{2}} \left(c_1 \cos \left(\frac{\sqrt{7}x}{2} \right) + c_2 \sin \left(\frac{\sqrt{7}x}{2} \right) \right)}{2} + e^{\frac{x}{2}} \left(-\frac{c_1 \sqrt{7} \sin \left(\frac{\sqrt{7}x}{2} \right)}{2} + \frac{c_2 \sqrt{7} \cos \left(\frac{\sqrt{7}x}{2} \right)}{2} \right)$$

substituting $u' = 0$ and $x = 0$ in the above gives

$$0 = \frac{c_1}{2} + \frac{\sqrt{7}c_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 2$$

$$c_2 = -\frac{2\sqrt{7}}{7}$$

Substituting these values back in above solution results in

$$u = -\frac{2 \sin\left(\frac{\sqrt{7}x}{2}\right) e^{\frac{x}{2}} \sqrt{7}}{7} + 2 e^{\frac{x}{2}} \cos\left(\frac{\sqrt{7}x}{2}\right)$$

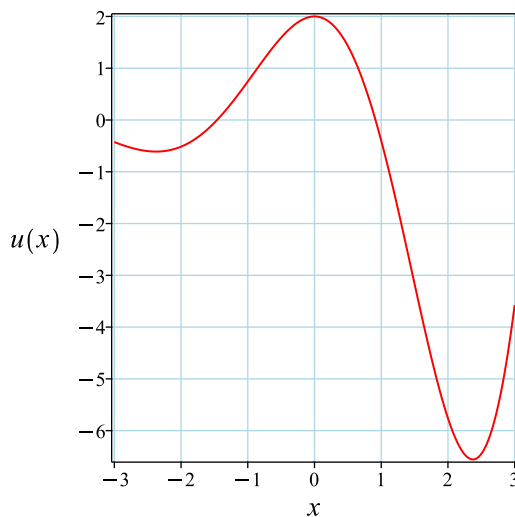
Which simplifies to

$$u = -\frac{2\left(\sqrt{7} \sin\left(\frac{\sqrt{7}x}{2}\right) - 7 \cos\left(\frac{\sqrt{7}x}{2}\right)\right) e^{\frac{x}{2}}}{7}$$

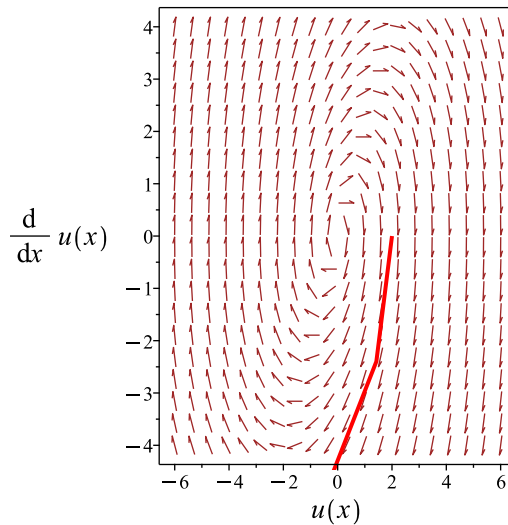
Summary

The solution(s) found are the following

$$u = -\frac{2\left(\sqrt{7} \sin\left(\frac{\sqrt{7}x}{2}\right) - 7 \cos\left(\frac{\sqrt{7}x}{2}\right)\right) e^{\frac{x}{2}}}{7} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = -\frac{2\left(\sqrt{7} \sin\left(\frac{\sqrt{7}x}{2}\right) - 7 \cos\left(\frac{\sqrt{7}x}{2}\right)\right) e^{\frac{x}{2}}}{7}$$

Verified OK.

8.17.3 Solving using Kovacic algorithm

Writing the ode as

$$u'' - u' + 2u = 0 \quad (1)$$

$$Au'' + Bu' + Cu = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = -1 \quad (3)$$

$$C = 2$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ue^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \quad (5)$$
$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-7}{4} \quad (6)$$

Comparing the above to (5) shows that

$$s = -7$$

$$t = 4$$

Therefore eq. (4) becomes

$$z''(x) = -\frac{7z(x)}{4} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then u is found using the inverse transformation

$$u = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 396: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -\frac{7}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos\left(\frac{\sqrt{7}x}{2}\right)$$

Using the above, the solution for the original ode can now be found. The first solution

to the original ode in u is found from

$$\begin{aligned}u_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\&= z_1 e^{-\int \frac{1}{2} \frac{-1}{1} dx} \\&= z_1 e^{\frac{x}{2}} \\&= z_1 \left(e^{\frac{x}{2}} \right)\end{aligned}$$

Which simplifies to

$$u_1 = e^{\frac{x}{2}} \cos \left(\frac{\sqrt{7} x}{2} \right)$$

The second solution u_2 to the original ode is found using reduction of order

$$u_2 = u_1 \int \frac{e^{\int -\frac{B}{A} dx}}{u_1^2} dx$$

Substituting gives

$$\begin{aligned}u_2 &= u_1 \int \frac{e^{\int -\frac{-1}{1} dx}}{(u_1)^2} dx \\&= u_1 \int \frac{e^x}{(u_1)^2} dx \\&= u_1 \left(\frac{2\sqrt{7} \tan \left(\frac{\sqrt{7} x}{2} \right)}{7} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}u &= c_1 u_1 + c_2 u_2 \\&= c_1 \left(e^{\frac{x}{2}} \cos \left(\frac{\sqrt{7} x}{2} \right) \right) + c_2 \left(e^{\frac{x}{2}} \cos \left(\frac{\sqrt{7} x}{2} \right) \left(\frac{2\sqrt{7} \tan \left(\frac{\sqrt{7} x}{2} \right)}{7} \right) \right)\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = c_1 e^{\frac{x}{2}} \cos\left(\frac{\sqrt{7}x}{2}\right) + \frac{2c_2 \sin\left(\frac{\sqrt{7}x}{2}\right) e^{\frac{x}{2}} \sqrt{7}}{7} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 2$ and $x = 0$ in the above gives

$$2 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$u' = \frac{c_1 e^{\frac{x}{2}} \cos\left(\frac{\sqrt{7}x}{2}\right)}{2} - \frac{c_1 e^{\frac{x}{2}} \sqrt{7} \sin\left(\frac{\sqrt{7}x}{2}\right)}{2} + c_2 \cos\left(\frac{\sqrt{7}x}{2}\right) e^{\frac{x}{2}} + \frac{c_2 \sin\left(\frac{\sqrt{7}x}{2}\right) e^{\frac{x}{2}} \sqrt{7}}{7}$$

substituting $u' = 0$ and $x = 0$ in the above gives

$$0 = \frac{c_1}{2} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned} c_1 &= 2 \\ c_2 &= -1 \end{aligned}$$

Substituting these values back in above solution results in

$$u = -\frac{2 \sin\left(\frac{\sqrt{7}x}{2}\right) e^{\frac{x}{2}} \sqrt{7}}{7} + 2 e^{\frac{x}{2}} \cos\left(\frac{\sqrt{7}x}{2}\right)$$

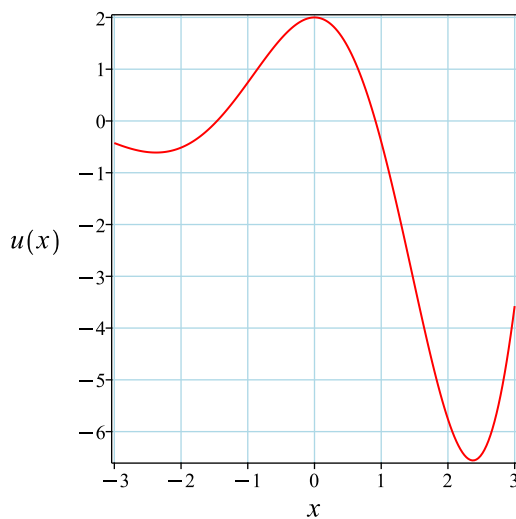
Which simplifies to

$$u = -\frac{2\left(\sqrt{7} \sin\left(\frac{\sqrt{7}x}{2}\right) - 7 \cos\left(\frac{\sqrt{7}x}{2}\right)\right) e^{\frac{x}{2}}}{7}$$

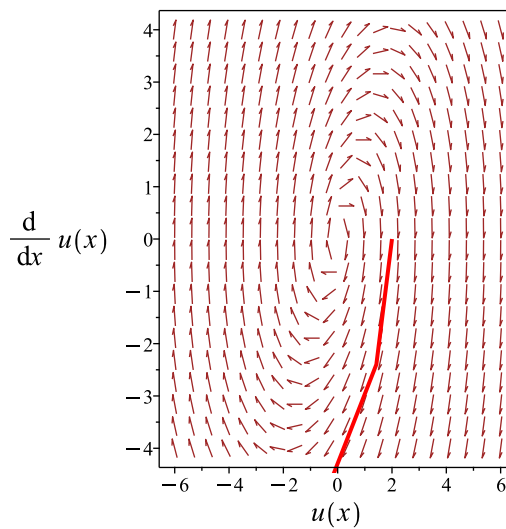
Summary

The solution(s) found are the following

$$u = -\frac{2\left(\sqrt{7} \sin\left(\frac{\sqrt{7}x}{2}\right) - 7 \cos\left(\frac{\sqrt{7}x}{2}\right)\right) e^{\frac{x}{2}}}{7} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = -\frac{2\left(\sqrt{7} \sin\left(\frac{\sqrt{7}x}{2}\right) - 7 \cos\left(\frac{\sqrt{7}x}{2}\right)\right) e^{\frac{x}{2}}}{7}$$

Verified OK.

8.17.4 Maple step by step solution

Let's solve

$$\left[u'' - u' + 2u = 0, u(0) = 2, u'|_{\{x=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

$$u''$$

- Characteristic polynomial of ODE

$$r^2 - r + 2 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{1 \pm (\sqrt{-7})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(\frac{1}{2} - \frac{i\sqrt{7}}{2}, \frac{1}{2} + \frac{i\sqrt{7}}{2} \right)$$

- 1st solution of the ODE

$$u_1(x) = e^{\frac{x}{2}} \cos\left(\frac{\sqrt{7}x}{2}\right)$$

- 2nd solution of the ODE

$$u_2(x) = e^{\frac{x}{2}} \sin\left(\frac{\sqrt{7}x}{2}\right)$$

- General solution of the ODE

$$u = c_1 u_1(x) + c_2 u_2(x)$$

- Substitute in solutions

$$u = c_1 e^{\frac{x}{2}} \cos\left(\frac{\sqrt{7}x}{2}\right) + \sin\left(\frac{\sqrt{7}x}{2}\right) e^{\frac{x}{2}} c_2$$

- Check validity of solution $u = c_1 e^{\frac{x}{2}} \cos\left(\frac{\sqrt{7}x}{2}\right) + \sin\left(\frac{\sqrt{7}x}{2}\right) e^{\frac{x}{2}} c_2$

- Use initial condition $u(0) = 2$

$$2 = c_1$$

- Compute derivative of the solution

$$u' = \frac{c_1 e^{\frac{x}{2}} \cos\left(\frac{\sqrt{7}x}{2}\right)}{2} - \frac{c_1 e^{\frac{x}{2}} \sqrt{7} \sin\left(\frac{\sqrt{7}x}{2}\right)}{2} + \frac{\sqrt{7} \cos\left(\frac{\sqrt{7}x}{2}\right) e^{\frac{x}{2}} c_2}{2} + \frac{\sin\left(\frac{\sqrt{7}x}{2}\right) e^{\frac{x}{2}} c_2}{2}$$

- Use the initial condition $u' \Big|_{\{x=0\}} = 0$

$$0 = \frac{c_1}{2} + \frac{\sqrt{7} c_2}{2}$$

- Solve for c_1 and c_2

$$\left\{ c_1 = 2, c_2 = -\frac{2\sqrt{7}}{7} \right\}$$

- Substitute constant values into general solution and simplify

$$u = -\frac{2\left(\sqrt{7} \sin\left(\frac{\sqrt{7}x}{2}\right) - 7 \cos\left(\frac{\sqrt{7}x}{2}\right)\right) e^{\frac{x}{2}}}{7}$$

- Solution to the IVP

$$u = -\frac{2\left(\sqrt{7} \sin\left(\frac{\sqrt{7}x}{2}\right) - 7 \cos\left(\frac{\sqrt{7}x}{2}\right)\right) e^{\frac{x}{2}}}{7}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 31

```
dsolve([diff(u(x),x$2)- diff(u(x),x)+2*u(x) = 0,u(0) = 2, D(u)(0) = 0],u(x), singsol=all)
```

$$u(x) = -\frac{2e^{\frac{x}{2}}\left(\sqrt{7}\sin\left(\frac{\sqrt{7}x}{2}\right) - 7\cos\left(\frac{\sqrt{7}x}{2}\right)\right)}{7}$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 19

```
DSolve[{u''[x]+4*u'[x]+5*u[x]==0,{u[0]==2,u'[0]==0}},u[x],x,IncludeSingularSolutions -> True
```

$$u(x) \rightarrow 2e^{-2x}(2\sin(x) + \cos(x))$$

8.18 problem 24

8.18.1 Existence and uniqueness analysis	2169
8.18.2 Solving as second order linear constant coeff ode	2170
8.18.3 Solving using Kovacic algorithm	2173
8.18.4 Maple step by step solution	2177

Internal problem ID [640]

Internal file name [OUTPUT/640_Sunday_June_05_2022_01_46_16_AM_5872921/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 24.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$5u'' + 2u' + 7u = 0$$

With initial conditions

$$[u(0) = 2, u'(0) = 1]$$

8.18.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$u'' + p(x)u' + q(x)u = F$$

Where here

$$p(x) = \frac{2}{5}$$
$$q(x) = \frac{7}{5}$$
$$F = 0$$

Hence the ode is

$$u'' + \frac{2u'}{5} + \frac{7u}{5} = 0$$

The domain of $p(x) = \frac{2}{5}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = \frac{7}{5}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

8.18.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Au''(x) + Bu'(x) + Cu(x) = 0$$

Where in the above $A = 5, B = 2, C = 7$. Let the solution be $u = e^{\lambda x}$. Substituting this into the ODE gives

$$5\lambda^2 e^{\lambda x} + 2\lambda e^{\lambda x} + 7e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$5\lambda^2 + 2\lambda + 7 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 5, B = 2, C = 7$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-2}{(2)(5)} \pm \frac{1}{(2)(5)} \sqrt{2^2 - (4)(5)(7)} \\ &= -\frac{1}{5} \pm \frac{i\sqrt{34}}{5} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= -\frac{1}{5} + \frac{i\sqrt{34}}{5} \\ \lambda_2 &= -\frac{1}{5} - \frac{i\sqrt{34}}{5} \end{aligned}$$

Which simplifies to

$$\lambda_1 = -\frac{1}{5} + \frac{i\sqrt{34}}{5}$$

$$\lambda_2 = -\frac{1}{5} - \frac{i\sqrt{34}}{5}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -\frac{1}{5}$ and $\beta = \frac{\sqrt{34}}{5}$. Therefore the final solution, when using Euler relation, can be written as

$$u = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$u = e^{-\frac{x}{5}} \left(c_1 \cos \left(\frac{\sqrt{34}x}{5} \right) + c_2 \sin \left(\frac{\sqrt{34}x}{5} \right) \right)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = e^{-\frac{x}{5}} \left(c_1 \cos \left(\frac{\sqrt{34}x}{5} \right) + c_2 \sin \left(\frac{\sqrt{34}x}{5} \right) \right) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 2$ and $x = 0$ in the above gives

$$2 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$u' = -\frac{e^{-\frac{x}{5}} \left(c_1 \cos \left(\frac{\sqrt{34}x}{5} \right) + c_2 \sin \left(\frac{\sqrt{34}x}{5} \right) \right)}{5} + e^{-\frac{x}{5}} \left(-\frac{c_1 \sqrt{34} \sin \left(\frac{\sqrt{34}x}{5} \right)}{5} + \frac{c_2 \sqrt{34} \cos \left(\frac{\sqrt{34}x}{5} \right)}{5} \right)$$

substituting $u' = 1$ and $x = 0$ in the above gives

$$1 = -\frac{c_1}{5} + \frac{\sqrt{34}c_2}{5} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 2$$

$$c_2 = \frac{7\sqrt{34}}{34}$$

Substituting these values back in above solution results in

$$u = \frac{7 \sin\left(\frac{\sqrt{34}x}{5}\right) e^{-\frac{x}{5}} \sqrt{34}}{34} + 2 e^{-\frac{x}{5}} \cos\left(\frac{\sqrt{34}x}{5}\right)$$

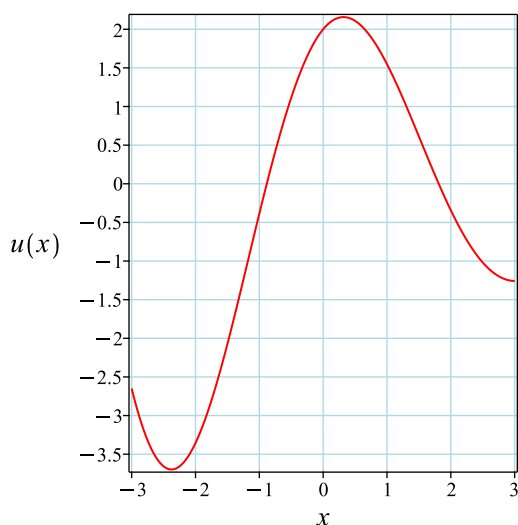
Which simplifies to

$$u = \frac{\left(7\sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right) + 68 \cos\left(\frac{\sqrt{34}x}{5}\right)\right) e^{-\frac{x}{5}}}{34}$$

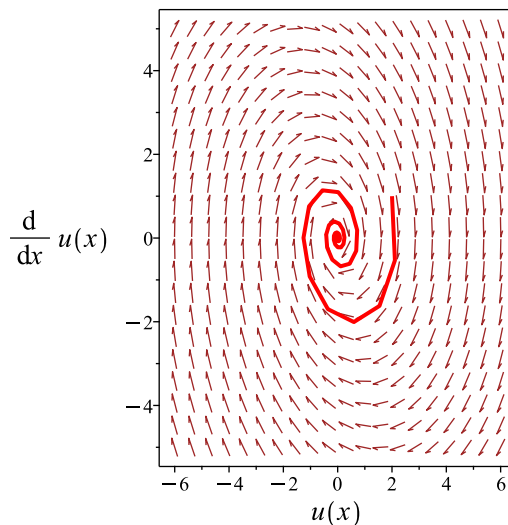
Summary

The solution(s) found are the following

$$u = \frac{\left(7\sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right) + 68 \cos\left(\frac{\sqrt{34}x}{5}\right)\right) e^{-\frac{x}{5}}}{34} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = \frac{\left(7\sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right) + 68 \cos\left(\frac{\sqrt{34}x}{5}\right)\right) e^{-\frac{x}{5}}}{34}$$

Verified OK.

8.18.3 Solving using Kovacic algorithm

Writing the ode as

$$5u'' + 2u' + 7u = 0 \quad (1)$$

$$Au'' + Bu' + Cu = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 5$$

$$B = 2 \quad (3)$$

$$C = 7$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ue^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \quad (5)$$
$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-34}{25} \quad (6)$$

Comparing the above to (5) shows that

$$s = -34$$

$$t = 25$$

Therefore eq. (4) becomes

$$z''(x) = -\frac{34z(x)}{25} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then u is found using the inverse transformation

$$u = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 398: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -\frac{34}{25}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos\left(\frac{\sqrt{34}x}{5}\right)$$

Using the above, the solution for the original ode can now be found. The first solution

to the original ode in u is found from

$$\begin{aligned}u_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\&= z_1 e^{-\int \frac{1}{2} \frac{2}{5} dx} \\&= z_1 e^{-\frac{x}{5}} \\&= z_1 \left(e^{-\frac{x}{5}} \right)\end{aligned}$$

Which simplifies to

$$u_1 = e^{-\frac{x}{5}} \cos \left(\frac{\sqrt{34} x}{5} \right)$$

The second solution u_2 to the original ode is found using reduction of order

$$u_2 = u_1 \int \frac{e^{\int -\frac{B}{A} dx}}{u_1^2} dx$$

Substituting gives

$$\begin{aligned}u_2 &= u_1 \int \frac{e^{\int -\frac{2}{5} dx}}{(u_1)^2} dx \\&= u_1 \int \frac{e^{-\frac{2x}{5}}}{(u_1)^2} dx \\&= u_1 \left(\frac{5\sqrt{34} \tan \left(\frac{\sqrt{34} x}{5} \right)}{34} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}u &= c_1 u_1 + c_2 u_2 \\&= c_1 \left(e^{-\frac{x}{5}} \cos \left(\frac{\sqrt{34} x}{5} \right) \right) + c_2 \left(e^{-\frac{x}{5}} \cos \left(\frac{\sqrt{34} x}{5} \right) \left(\frac{5\sqrt{34} \tan \left(\frac{\sqrt{34} x}{5} \right)}{34} \right) \right)\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = c_1 e^{-\frac{x}{5}} \cos\left(\frac{\sqrt{34}x}{5}\right) + \frac{5c_2 \sin\left(\frac{\sqrt{34}x}{5}\right) e^{-\frac{x}{5}} \sqrt{34}}{34} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 2$ and $x = 0$ in the above gives

$$2 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$u' = -\frac{c_1 e^{-\frac{x}{5}} \cos\left(\frac{\sqrt{34}x}{5}\right)}{5} - \frac{c_1 e^{-\frac{x}{5}} \sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right)}{5} + c_2 \cos\left(\frac{\sqrt{34}x}{5}\right) e^{-\frac{x}{5}} - \frac{c_2 \sin\left(\frac{\sqrt{34}x}{5}\right) e^{-\frac{x}{5}} \sqrt{34}}{34}$$

substituting $u' = 1$ and $x = 0$ in the above gives

$$1 = -\frac{c_1}{5} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned} c_1 &= 2 \\ c_2 &= \frac{7}{5} \end{aligned}$$

Substituting these values back in above solution results in

$$u = \frac{7 \sin\left(\frac{\sqrt{34}x}{5}\right) e^{-\frac{x}{5}} \sqrt{34}}{34} + 2 e^{-\frac{x}{5}} \cos\left(\frac{\sqrt{34}x}{5}\right)$$

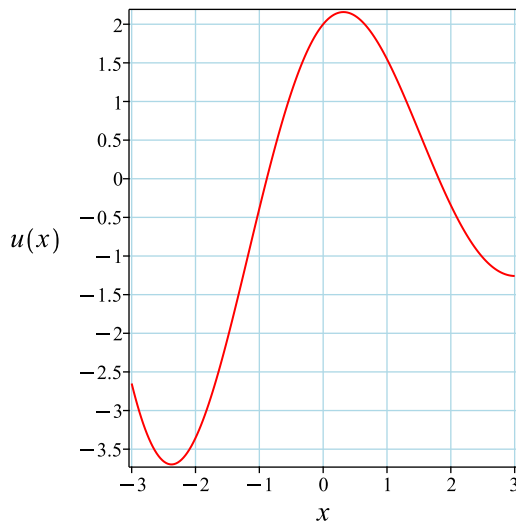
Which simplifies to

$$u = \frac{\left(7\sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right) + 68 \cos\left(\frac{\sqrt{34}x}{5}\right)\right) e^{-\frac{x}{5}}}{34}$$

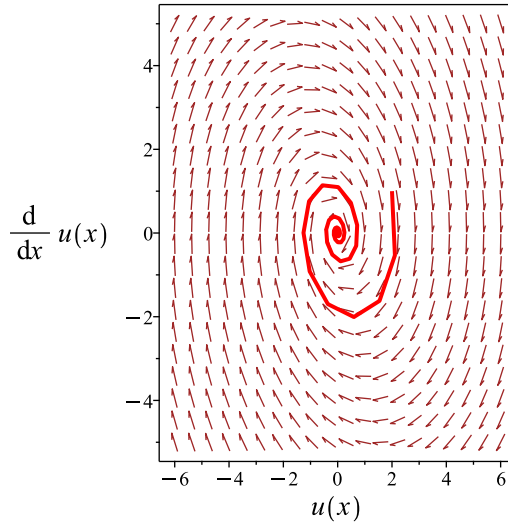
Summary

The solution(s) found are the following

$$u = \frac{\left(7\sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right) + 68 \cos\left(\frac{\sqrt{34}x}{5}\right)\right) e^{-\frac{x}{5}}}{34} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = \frac{\left(7\sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right) + 68 \cos\left(\frac{\sqrt{34}x}{5}\right)\right) e^{-\frac{x}{5}}}{34}$$

Verified OK.

8.18.4 Maple step by step solution

Let's solve

$$\left[5u'' + 2u' + 7u = 0, u(0) = 2, u' \Big|_{\{x=0\}} = 1\right]$$

- Highest derivative means the order of the ODE is 2

$$u''$$

- Isolate 2nd derivative

$$u'' = -\frac{2u'}{5} - \frac{7u}{5}$$

- Group terms with u on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$u'' + \frac{2u'}{5} + \frac{7u}{5} = 0$$

- Characteristic polynomial of ODE

$$r^2 + \frac{2}{5}r + \frac{7}{5} = 0$$

- Use quadratic formula to solve for r

$$r = \frac{\left(-\frac{2}{5}\right) \pm \left(\sqrt{-\frac{136}{25}}\right)}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{5} - \frac{i\sqrt{34}}{5}, -\frac{1}{5} + \frac{i\sqrt{34}}{5}\right)$$

- 1st solution of the ODE

$$u_1(x) = e^{-\frac{x}{5}} \cos\left(\frac{\sqrt{34}x}{5}\right)$$

- 2nd solution of the ODE

$$u_2(x) = e^{-\frac{x}{5}} \sin\left(\frac{\sqrt{34}x}{5}\right)$$

- General solution of the ODE

$$u = c_1 u_1(x) + c_2 u_2(x)$$

- Substitute in solutions

$$u = c_1 e^{-\frac{x}{5}} \cos\left(\frac{\sqrt{34}x}{5}\right) + \sin\left(\frac{\sqrt{34}x}{5}\right) e^{-\frac{x}{5}} c_2$$

- Check validity of solution $u = c_1 e^{-\frac{x}{5}} \cos\left(\frac{\sqrt{34}x}{5}\right) + \sin\left(\frac{\sqrt{34}x}{5}\right) e^{-\frac{x}{5}} c_2$

- Use initial condition $u(0) = 2$

$$2 = c_1$$

- Compute derivative of the solution

$$u' = -\frac{c_1 e^{-\frac{x}{5}} \cos\left(\frac{\sqrt{34}x}{5}\right)}{5} - \frac{c_1 e^{-\frac{x}{5}} \sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right)}{5} + \frac{\sqrt{34} \cos\left(\frac{\sqrt{34}x}{5}\right) e^{-\frac{x}{5}} c_2}{5} - \frac{\sin\left(\frac{\sqrt{34}x}{5}\right) e^{-\frac{x}{5}} c_2}{5}$$

- Use the initial condition $u' \Big|_{\{x=0\}} = 1$

$$1 = -\frac{c_1}{5} + \frac{\sqrt{34}c_2}{5}$$

- Solve for c_1 and c_2

$$\left\{ c_1 = 2, c_2 = \frac{7\sqrt{34}}{34} \right\}$$

- Substitute constant values into general solution and simplify

$$u = \frac{\left(7\sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right) + 68 \cos\left(\frac{\sqrt{34}x}{5}\right)\right) e^{-\frac{x}{5}}}{34}$$

- Solution to the IVP

$$u = \frac{\left(7\sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right) + 68 \cos\left(\frac{\sqrt{34}x}{5}\right)\right) e^{-\frac{x}{5}}}{34}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 32

```
dsolve([5*diff(u(x),x$2)+ 2*diff(u(x),x)+7*u(x) = 0,u(0) = 2, D(u)(0) = 1],u(x), singsol=all
```

$$u(x) = \frac{e^{-\frac{x}{5}} \left(7\sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right) + 68 \cos\left(\frac{\sqrt{34}x}{5}\right) \right)}{34}$$

✓ Solution by Mathematica

Time used: 0.025 (sec). Leaf size: 48

```
DSolve[{5*u''[x]+2*u'[x]+7*u[x]==0,{u[0]==2,u'[0]==1}},u[x],x,IncludeSingularSolutions -> Tr
```

$$u(x) \rightarrow \frac{1}{34} e^{-x/5} \left(7\sqrt{34} \sin\left(\frac{\sqrt{34}x}{5}\right) + 68 \cos\left(\frac{\sqrt{34}x}{5}\right) \right)$$

8.19 problem 25

8.19.1 Existence and uniqueness analysis	2180
8.19.2 Solving as second order linear constant coeff ode	2181
8.19.3 Solving using Kovacic algorithm	2183
8.19.4 Maple step by step solution	2187

Internal problem ID [641]

Internal file name [OUTPUT/641_Sunday_June_05_2022_01_46_17_AM_9066687/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 25.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 2y' + 6y = 0$$

With initial conditions

$$[y(0) = 2, y'(0) = \alpha]$$

8.19.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 2$$

$$q(x) = 6$$

$$F = 0$$

Hence the ode is

$$y'' + 2y' + 6y = 0$$

The domain of $p(x) = 2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = 6$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

8.19.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 2, C = 6$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 2\lambda e^{\lambda x} + 6e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 2\lambda + 6 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 2, C = 6$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{2^2 - (4)(1)(6)} \\ &= -1 \pm i\sqrt{5} \end{aligned}$$

Hence

$$\lambda_1 = -1 + i\sqrt{5}$$

$$\lambda_2 = -1 - i\sqrt{5}$$

Which simplifies to

$$\lambda_1 = -1 + i\sqrt{5}$$

$$\lambda_2 = -1 - i\sqrt{5}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -1$ and $\beta = \sqrt{5}$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-x}(c_1 \cos(x\sqrt{5}) + c_2 \sin(x\sqrt{5}))$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = e^{-x}(c_1 \cos(x\sqrt{5}) + c_2 \sin(x\sqrt{5})) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$y' = -e^{-x}(c_1 \cos(x\sqrt{5}) + c_2 \sin(x\sqrt{5})) + e^{-x}(-c_1\sqrt{5} \sin(x\sqrt{5}) + c_2\sqrt{5} \cos(x\sqrt{5}))$$

substituting $y' = \alpha$ and $x = 0$ in the above gives

$$\alpha = -c_1 + \sqrt{5}c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 2$$
$$c_2 = \frac{\sqrt{5}(\alpha + 2)}{5}$$

Substituting these values back in above solution results in

$$y = \frac{e^{-x} \sin(x\sqrt{5}) \sqrt{5}(\alpha + 2)}{5} + 2e^{-x} \cos(x\sqrt{5})$$

Which simplifies to

$$y = \frac{(\sin(x\sqrt{5}) \sqrt{5}(\alpha + 2) + 10 \cos(x\sqrt{5})) e^{-x}}{5}$$

Summary

The solution(s) found are the following

$$y = \frac{(\sin(x\sqrt{5}) \sqrt{5}(\alpha + 2) + 10 \cos(x\sqrt{5})) e^{-x}}{5} \quad (1)$$

Verification of solutions

$$y = \frac{(\sin(x\sqrt{5}) \sqrt{5}(\alpha + 2) + 10 \cos(x\sqrt{5})) e^{-x}}{5}$$

Verified OK.

8.19.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 2y' + 6y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 2 \\ C &= 6 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \tag{5}$$

$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-5}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = -5$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = -5z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 400: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -5$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x\sqrt{5})$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{2}{1} dx} \\ &= z_1 e^{-x} \\ &= z_1 (e^{-x}) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-x} \cos(x\sqrt{5})$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-2x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{\sqrt{5} \tan(x\sqrt{5})}{5} \right) \end{aligned}$$

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$

$$= c_1 \left(e^{-x} \cos(x\sqrt{5}) \right) + c_2 \left(e^{-x} \cos(x\sqrt{5}) \left(\frac{\sqrt{5} \tan(x\sqrt{5})}{5} \right) \right)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-x} \cos(x\sqrt{5}) + \frac{c_2 e^{-x} \sin(x\sqrt{5}) \sqrt{5}}{5} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = 0$ in the above gives

$$2 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$y' = -c_1 e^{-x} \cos(x\sqrt{5}) - c_1 e^{-x} \sqrt{5} \sin(x\sqrt{5}) - \frac{c_2 e^{-x} \sin(x\sqrt{5}) \sqrt{5}}{5} + c_2 e^{-x} \cos(x\sqrt{5})$$

substituting $y' = \alpha$ and $x = 0$ in the above gives

$$\alpha = -c_1 + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 2$$

$$c_2 = \alpha + 2$$

Substituting these values back in above solution results in

$$y = \frac{e^{-x} \sin(x\sqrt{5}) \sqrt{5} (\alpha + 2)}{5} + 2 e^{-x} \cos(x\sqrt{5})$$

Which simplifies to

$$y = \frac{(\sin(x\sqrt{5}) \sqrt{5} (\alpha + 2) + 10 \cos(x\sqrt{5})) e^{-x}}{5}$$

Summary

The solution(s) found are the following

$$y = \frac{(\sin(x\sqrt{5})\sqrt{5}(\alpha + 2) + 10\cos(x\sqrt{5}))e^{-x}}{5} \quad (1)$$

Verification of solutions

$$y = \frac{(\sin(x\sqrt{5})\sqrt{5}(\alpha + 2) + 10\cos(x\sqrt{5}))e^{-x}}{5}$$

Verified OK.

8.19.4 Maple step by step solution

Let's solve

$$\left[y'' + 2y' + 6y = 0, y(0) = 2, y'|_{\{x=0\}} = \alpha \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 + 2r + 6 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-2) \pm (\sqrt{-20})}{2}$$

- Roots of the characteristic polynomial

$$r = (-1 - I\sqrt{5}, -1 + I\sqrt{5})$$

- 1st solution of the ODE

$$y_1(x) = e^{-x} \cos(x\sqrt{5})$$

- 2nd solution of the ODE

$$y_2(x) = e^{-x} \sin(x\sqrt{5})$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-x} \cos(x\sqrt{5}) + e^{-x} \sin(x\sqrt{5}) c_2$$

- Check validity of solution $y = c_1 e^{-x} \cos(x\sqrt{5}) + e^{-x} \sin(x\sqrt{5}) c_2$

- Use initial condition $y(0) = 2$

$$2 = c_1$$

- Compute derivative of the solution

$$y' = -c_1 e^{-x} \cos(x\sqrt{5}) - c_1 e^{-x} \sqrt{5} \sin(x\sqrt{5}) - e^{-x} \sin(x\sqrt{5}) c_2 + e^{-x} \sqrt{5} \cos(x\sqrt{5}) c_2$$

- Use the initial condition $y' \Big|_{\{x=0\}} = \alpha$

$$\alpha = -c_1 + \sqrt{5} c_2$$

- Solve for c_1 and c_2

$$\left\{ c_1 = 2, c_2 = \frac{\sqrt{5}(\alpha+2)}{5} \right\}$$

- Substitute constant values into general solution and simplify

$$y = \frac{(\sin(x\sqrt{5})\sqrt{5}(\alpha+2) + 10 \cos(x\sqrt{5})) e^{-x}}{5}$$

- Solution to the IVP

$$y = \frac{(\sin(x\sqrt{5})\sqrt{5}(\alpha+2) + 10 \cos(x\sqrt{5})) e^{-x}}{5}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 32

```
dsolve([diff(y(x),x$2)+ 2*diff(y(x),x)+6*y(x) = 0,y(0) = 2, D(y)(0) = alpha],y(x), singsol=a
```

$$y(x) = \frac{e^{-x}(\sqrt{5}(\alpha+2)\sin(\sqrt{5}x) + 10\cos(\sqrt{5}x))}{5}$$

✓ Solution by Mathematica

Time used: 0.021 (sec). Leaf size: 42

```
DSolve[{y''[x]+2*y'[x]+6*y[x]==0,{y[0]==2,y'[0]==\[Alpha]}},y[x],x,IncludeSingularSolutions
```

$$y(x) \rightarrow \frac{1}{5}e^{-x} \left(\sqrt{5}(\alpha + 2) \sin(\sqrt{5}x) + 10 \cos(\sqrt{5}x) \right)$$

8.20 problem 26

8.20.1 Existence and uniqueness analysis	2190
8.20.2 Solving as second order linear constant coeff ode	2191
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Internal problem ID [642]

Internal file name [OUTPUT/642_Sunday_June_05_2022_01_46_18_AM_29576795/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 26.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' + 2ay' + (a^2 + 1)y = 0$$

With initial conditions

$$[y(0) = 1, y'(0) = 0]$$

8.20.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 2a$$

$$q(x) = a^2 + 1$$

$$F = 0$$

Hence the ode is

$$y'' + 2ay' + (a^2 + 1)y = 0$$

The domain of $p(x) = 2a$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = a^2 + 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

8.20.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 2a, C = a^2 + 1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 2a\lambda e^{\lambda x} + (a^2 + 1)e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$a^2 + 2a\lambda + \lambda^2 + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 2a, C = a^2 + 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-2a}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{2a^2 - (4)(1)(a^2 + 1)} \\ &= -a \pm i \end{aligned}$$

Hence

$$\lambda_1 = -a + i$$

$$\lambda_2 = -a - i$$

Which simplifies to

$$\lambda_1 = -a + i$$

$$\lambda_2 = -a - i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -a$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-ax}(c_1 \cos(x) + c_2 \sin(x))$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = e^{-ax}(c_1 \cos(x) + c_2 \sin(x)) \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = 0$ in the above gives

$$1 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -a e^{-ax}(c_1 \cos(x) + c_2 \sin(x)) + e^{-ax}(-c_1 \sin(x) + c_2 \cos(x))$$

substituting $y' = 0$ and $x = 0$ in the above gives

$$0 = -ac_1 + c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 1$$

$$c_2 = a$$

Substituting these values back in above solution results in

$$y = e^{-ax}(\sin(x) a + \cos(x))$$

Summary

The solution(s) found are the following

$$y = e^{-ax}(\sin(x)a + \cos(x)) \quad (1)$$

Verification of solutions

$$y = e^{-ax}(\sin(x)a + \cos(x))$$

Verified OK.

8.20.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 2ay' + (a^2 + 1)y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 2a \quad (3)$$

$$C = a^2 + 1$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \quad (6)$$

Comparing the above to (5) shows that

$$s = -1$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = -z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 402: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{2a}{1} dx} \\ &= z_1 e^{-ax} \\ &= z_1 (e^{-ax}) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-ax} \cos(x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{2a}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-2ax}}{(y_1)^2} dx \\ &= y_1 (\tan(x)) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-ax} \cos(x)) + c_2 (e^{-ax} \cos(x) (\tan(x))) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-ax} \cos(x) + c_2 e^{-ax} \sin(x) \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $x = 0$ in the above gives

$$1 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -c_1 a e^{-ax} \cos(x) - c_1 e^{-ax} \sin(x) - c_2 a e^{-ax} \sin(x) + c_2 e^{-ax} \cos(x)$$

substituting $y' = 0$ and $x = 0$ in the above gives

$$0 = -ac_1 + c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 1$$

$$c_2 = a$$

Substituting these values back in above solution results in

$$y = e^{-ax} \sin(x) a + e^{-ax} \cos(x)$$

Which simplifies to

$$y = e^{-ax} (\sin(x) a + \cos(x))$$

Summary

The solution(s) found are the following

$$y = e^{-ax} (\sin(x) a + \cos(x)) \tag{1}$$

Verification of solutions

$$y = e^{-ax} (\sin(x) a + \cos(x))$$

Verified OK.

8.20.4 Maple step by step solution

Let's solve

$$\left[y'' + 2ay' + (a^2 + 1)y = 0, y(0) = 1, y'|_{\{x=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

y''

- Characteristic polynomial of ODE

$$a^2 + 2ar + r^2 + 1 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-2a) \pm (\sqrt{-4})}{2}$$

- Roots of the characteristic polynomial

$$r = (-a - I, -a + I)$$

- 1st solution of the ODE

$$y_1(x) = e^{-\Re(a)x} \cos(|\Im(a) + 1|x)$$

- 2nd solution of the ODE

$$y_2(x) = e^{-\Re(a)x} \sin(|\Im(a) + 1|x)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-\Re(a)x} \cos(|\Im(a) + 1|x) + c_2 e^{-\Re(a)x} \sin(|\Im(a) + 1|x)$$

- Check validity of solution $y = c_1 e^{-\Re(a)x} \cos(|\Im(a) + 1|x) + c_2 e^{-\Re(a)x} \sin(|\Im(a) + 1|x)$

- Use initial condition $y(0) = 1$

$$1 = c_1$$

- Compute derivative of the solution

$$y' = -c_1 \Re(a) e^{-\Re(a)x} \cos(|\Im(a) + 1|x) - c_1 e^{-\Re(a)x} |\Im(a) + 1| \sin(|\Im(a) + 1|x) - c_2 \Re(a) e^{-\Re(a)x} \sin(|\Im(a) + 1|x) + c_2 e^{-\Re(a)x} |\Im(a) + 1| \cos(|\Im(a) + 1|x)$$

- Use the initial condition $y' \Big|_{\{x=0\}} = 0$

$$0 = -c_1 \Re(a) + c_2 |\Im(a) + 1|$$

- Solve for c_1 and c_2

$$\left\{ c_1 = 1, c_2 = \frac{\Re(a)}{|\Im(a)+1|} \right\}$$

- Substitute constant values into general solution and simplify

$$y = \frac{e^{-\Re(a)x} (\text{signum}(\Im(a)+1) \sin(x(\Im(a)+1)) \Re(a) + \cos(x(\Im(a)+1)) |\Im(a)+1|)}{|\Im(a)+1|}$$

- Solution to the IVP

$$y = \frac{e^{-\Re(a)x} (\text{signum}(\Im(a)+1) \sin(x(\Im(a)+1)) \Re(a) + \cos(x(\Im(a)+1)) |\Im(a)+1|)}{|\Im(a)+1|}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve([diff(y(x),x$2)+ 2*a*diff(y(x),x)+(a^2+1)*y(x) = 0,y(0) = 1, D(y)(0) = 0],y(x), sings
```

$$y(x) = e^{-ax}(a \sin(x) + \cos(x))$$

✓ Solution by Mathematica

Time used: 0.031 (sec). Leaf size: 94

```
DSolve[{y''[x]+2*a*y'[x]+(a^2+1)*y[x]==0,{y[0]==1,y'[0]==0}},y[x],x,IncludeSingularSolutions
```

$$y(x) \rightarrow \frac{e^{-((\sqrt{a^2-a-1}+a)x)} \left(a \left(e^{2\sqrt{a^2-a-1}x} - 1 \right) + \sqrt{a^2-a-1} \left(e^{2\sqrt{a^2-a-1}x} + 1 \right) \right)}{2\sqrt{a^2-a-1}}$$

8.21 problem 35

8.21.1 Solving as second order euler ode ode	2199
8.21.2 Solving as second order change of variable on x method 2 ode .	2201
8.21.3 Solving as second order change of variable on x method 1 ode .	2203
8.21.4 Solving as second order change of variable on y method 2 ode .	2205
8.21.5 Solving using Kovacic algorithm	2208
8.21.6 Maple step by step solution	2213

Internal problem ID [643]

Internal file name [OUTPUT/643_Sunday_June_05_2022_01_46_19_AM_92480596/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 35.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_Emden, _Fowler], [_2nd_order, _linear, `_with_symmetry_[0,F(x)]]`]]
```

$$t^2 y'' + t y' + y = 0$$

8.21.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = r t^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} + t r t^{r-1} + t^r = 0$$

Simplifying gives

$$r(r-1)t^r + r t^r + t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r - 1) + r + 1 = 0$$

Or

$$r^2 + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -i$$

$$r_2 = i$$

The roots are complex conjugate of each others. Let the roots be

$$r_1 = \alpha + i\beta$$

$$r_2 = \alpha - i\beta$$

Where in this case $\alpha = 0$ and $\beta = -1$. Hence the solution becomes

$$\begin{aligned} y &= c_1 t^{r_1} + c_2 t^{r_2} \\ &= c_1 t^{\alpha+i\beta} + c_2 t^{\alpha-i\beta} \\ &= t^\alpha (c_1 t^{i\beta} + c_2 t^{-i\beta}) \\ &= t^\alpha (c_1 e^{\ln(t^{i\beta})} + c_2 e^{\ln(t^{-i\beta})}) \\ &= t^\alpha (c_1 e^{i(\beta \ln t)} + c_2 e^{-i(\beta \ln t)}) \end{aligned}$$

Using the values for $\alpha = 0, \beta = -1$, the above becomes

$$y = t^0 (c_1 e^{-i \ln(t)} + c_2 e^{i \ln(t)})$$

Using Euler relation, the expression $c_1 e^{iA} + c_2 e^{-iA}$ is transformed to $c_1 \cos A + c_1 \sin A$ where the constants are free to change. Applying this to the above result gives

$$y = c_1 \cos (\ln (t)) + c_2 \sin (\ln (t))$$

Summary

The solution(s) found are the following

$$y = c_1 \cos (\ln (t)) + c_2 \sin (\ln (t)) \tag{1}$$

Verification of solutions

$$y = c_1 \cos (\ln (t)) + c_2 \sin (\ln (t))$$

Verified OK.

8.21.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$t^2 y'' + ty' + y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = \frac{1}{t}$$
$$q(t) = \frac{1}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t) \tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t) dt)} dt \\ &= \int e^{-(\int \frac{1}{t} dt)} dt \\ &= \int e^{-\ln(t)} dt \\ &= \int \frac{1}{t} dt \\ &= \ln(t) \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{\frac{1}{t^2}}{\frac{1}{t^2}} \\ &= 1 \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + y(\tau) &= 0 \end{aligned}$$

The above ode is now solved for $y(\tau)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(\tau) + By'(\tau) + Cy(\tau) = 0$$

Where in the above $A = 1, B = 0, C = 1$. Let the solution be $y(\tau) = e^{\lambda\tau}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda\tau} + e^{\lambda\tau} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda\tau}$ gives

$$\lambda^2 + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(1)} \\ &= \pm i \end{aligned}$$

Hence

$$\lambda_1 = +i$$

$$\lambda_2 = -i$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= i \\ \lambda_2 &= -i\end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y(\tau) = e^{\alpha\tau}(c_1 \cos(\beta\tau) + c_2 \sin(\beta\tau))$$

Which becomes

$$y(\tau) = e^0(c_1 \cos(\tau) + c_2 \sin(\tau))$$

Or

$$y(\tau) = c_1 \cos(\tau) + c_2 \sin(\tau)$$

The above solution is now transformed back to y using (6) which results in

$$y = c_1 \cos(\ln(t)) + c_2 \sin(\ln(t))$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(\ln(t)) + c_2 \sin(\ln(t)) \quad (1)$$

Verification of solutions

$$y = c_1 \cos(\ln(t)) + c_2 \sin(\ln(t))$$

Verified OK.

8.21.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$t^2 y'' + ty' + y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{1}{t}$$

$$q(t) = \frac{1}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c}\sqrt{q}$$

$$= \frac{\sqrt{\frac{1}{t^2}}}{c} \quad (6)$$

$$\tau'' = -\frac{1}{c\sqrt{\frac{1}{t^2}}t^3}$$

Substituting the above into (4) results in

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2}$$

$$= \frac{-\frac{1}{c\sqrt{\frac{1}{t^2}}t^3} + \frac{1}{t}\frac{\sqrt{\frac{1}{t^2}}}{c}}{\left(\frac{\sqrt{\frac{1}{t^2}}}{c}\right)^2}$$

$$= 0$$

Therefore ode (3) now becomes

$$y(\tau)'' + p_1y(\tau)' + q_1y(\tau) = 0$$

$$\frac{d^2}{d\tau^2}y(\tau) + c^2y(\tau) = 0 \quad (7)$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = c_1 \cos(c\tau) + c_2 \sin(c\tau)$$

Now from (6)

$$\begin{aligned} \tau &= \int \frac{1}{c} \sqrt{q} dt \\ &= \frac{\int \sqrt{\frac{1}{t^2}} dt}{c} \\ &= \frac{\sqrt{\frac{1}{t^2}} t \ln(t)}{c} \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = c_1 \cos(\ln(t)) + c_2 \sin(\ln(t))$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(\ln(t)) + c_2 \sin(\ln(t)) \quad (1)$$

Verification of solutions

$$y = c_1 \cos(\ln(t)) + c_2 \sin(\ln(t))$$

Verified OK.

8.21.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$t^2 y'' + ty' + y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$\begin{aligned} p(t) &= \frac{1}{t} \\ q(t) &= \frac{1}{t^2} \end{aligned}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} + \frac{n}{t^2} + \frac{1}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = i \quad (6)$$

Substituting this value in (3) gives

$$\begin{aligned} v''(t) + \left(\frac{2i}{t} + \frac{1}{t}\right)v'(t) &= 0 \\ v''(t) + \frac{(1+2i)v'(t)}{t} &= 0 \end{aligned} \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{(1+2i)u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{(-1-2i)u}{t} \end{aligned}$$

Where $f(t) = \frac{-1-2i}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= \frac{-1-2i}{t} dt \\ \int \frac{1}{u} du &= \int \frac{-1-2i}{t} dt \\ \ln(u) &= (-1-2i) \ln(t) + c_1 \\ u &= e^{(-1-2i) \ln(t) + c_1} \\ &= c_1 e^{(-1-2i) \ln(t)}\end{aligned}$$

Which simplifies to

$$u(t) = \frac{c_1 t^{-2i}}{t}$$

Now that $u(t)$ is known, then

$$\begin{aligned}v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= \frac{ic_1 t^{-2i}}{2} + c_2\end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\ &= \left(\frac{ic_1 t^{-2i}}{2} + c_2 \right) t^i \\ &= t^i c_2 + \frac{it^{-i} c_1}{2}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \left(\frac{ic_1 t^{-2i}}{2} + c_2 \right) t^i \tag{1}$$

Verification of solutions

$$y = \left(\frac{ic_1 t^{-2i}}{2} + c_2 \right) t^i$$

Verified OK.

8.21.5 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' + t y' + y = 0 \quad (1)$$

$$A y'' + B y' + C y = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t^2 \\ B &= t \\ C &= 1 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = y e^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = r z(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-5}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -5 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(-\frac{5}{4t^2} \right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 404: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{5}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{5}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} + i \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} - i \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{5}{4t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{5}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} + i \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} - i \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{5}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2} + i$	$\frac{1}{2} - i$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2} + i$	$\frac{1}{2} - i$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = \frac{1}{2} - i$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= \frac{1}{2} - i - \left(\frac{1}{2} - i\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{t - c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= \frac{\frac{1}{2} - i}{t} + (-)(0) \\ &= \frac{\frac{1}{2} - i}{t} \\ &= \frac{\frac{1}{2} - i}{t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r)p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2\left(\frac{\frac{1}{2} - i}{t}\right)(0) + \left(\left(\frac{-\frac{1}{2} + i}{t^2}\right) + \left(\frac{\frac{1}{2} - i}{t}\right)^2 - \left(-\frac{5}{4t^2}\right)\right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \frac{\frac{1}{2} - i}{t} dt} \\ &= t^{\frac{1}{2} - i} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\&= z_1 e^{-\int \frac{1}{2} \frac{t}{t^2} dt} \\&= z_1 e^{-\frac{\ln(t)}{2}} \\&= z_1 \left(\frac{1}{\sqrt{t}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = t^{-i}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{t}{t^2} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{-\ln(t)}}{(y_1)^2} dt \\&= y_1 \left(-\frac{it^{2i}}{2} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 (t^{-i}) + c_2 \left(t^{-i} \left(-\frac{it^{2i}}{2} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = t^{-i} c_1 - \frac{ic_2 t^i}{2} \tag{1}$$

Verification of solutions

$$y = t^{-i} c_1 - \frac{ic_2 t^i}{2}$$

Verified OK.

8.21.6 Maple step by step solution

Let's solve

$$y''t^2 + ty' + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{y'}{t} - \frac{y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{y'}{t} + \frac{y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 + ty' + y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right) t^2 + \frac{d}{ds}y(s) + y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) + y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 + 1 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{0 \pm (\sqrt{-4})}{2}$$
- Roots of the characteristic polynomial

$$r = (-I, I)$$
- 1st solution of the ODE

$$y_1(s) = \cos(s)$$
- 2nd solution of the ODE

$$y_2(s) = \sin(s)$$
- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$
- Substitute in solutions

$$y(s) = c_1 \cos(s) + c_2 \sin(s)$$
- Change variables back using $s = \ln(t)$

$$y = c_1 \cos(\ln(t)) + c_2 \sin(\ln(t))$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 15

```
dsolve(t^2*diff(y(t),t$2)+ t*diff(y(t),t)+y(t) = 0,y(t), singsol=all)
```

$$y(t) = c_1 \sin(\ln(t)) + c_2 \cos(\ln(t))$$

✓ Solution by Mathematica

Time used: 0.017 (sec). Leaf size: 18

```
DSolve[t^2*y''[t]+t*y'[t]+y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow c_1 \cos(\log(t)) + c_2 \sin(\log(t))$$

8.22 problem 36

8.22.1 Solving as second order euler ode ode	2217
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Internal problem ID [644]

Internal file name [OUTPUT/644_Sunday_June_05_2022_01_46_20_AM_35632985/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 36.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "exact linear second order ode", "second_order_integrable_as_is", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_1", "second_order_change_of_variable_on_y_method_2", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _exact , _linear , _homogeneous]]
```

$$t^2y'' + 4ty' + 2y = 0$$

8.22.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} + 4trt^{r-1} + 2t^r = 0$$

Simplifying gives

$$r(r-1)t^r + 4rt^r + 2t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r-1) + 4r + 2 = 0$$

Or

$$r^2 + 3r + 2 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -2$$

$$r_2 = -1$$

Since the roots are real and distinct, then the general solution is

$$y = c_1y_1 + c_2y_2$$

Where $y_1 = t^{r_1}$ and $y_2 = t^{r_2}$. Hence

$$y = \frac{c_1}{t^2} + \frac{c_2}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t^2} + \frac{c_2}{t} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{t^2} + \frac{c_2}{t}$$

Verified OK.

8.22.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t)y' + \frac{(p(t)^2 + p'(t))y}{2} = f(t)$$

Where $p(t) = \frac{4}{t}$. Therefore, there is an integrating factor given by

$$\begin{aligned}M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int \frac{4}{t} dx} \\ &= t^2\end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned}(M(x)y)'' &= 0 \\ (t^2y)'' &= 0\end{aligned}$$

Integrating once gives

$$(t^2y)' = c_1$$

Integrating again gives

$$(t^2y) = c_1t + c_2$$

Hence the solution is

$$y = \frac{c_1t + c_2}{t^2}$$

Or

$$y = \frac{c_1}{t} + \frac{c_2}{t^2}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t} + \frac{c_2}{t^2} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{t} + \frac{c_2}{t^2}$$

Verified OK.

8.22.3 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$t^2 y'' + 4ty' + 2y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = \frac{4}{t}$$
$$q(t) = \frac{2}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t) \tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t) dt)} dt \\ &= \int e^{-(\int \frac{4}{t} dt)} dt \\ &= \int e^{-4 \ln(t)} dt \\ &= \int \frac{1}{t^4} dt \\ &= -\frac{1}{3t^3} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{\frac{2}{t^2}}{\frac{1}{t^8}} \\ &= 2t^6 \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + 2t^6y(\tau) &= 0 \end{aligned}$$

But in terms of τ

$$2t^6 = \frac{2}{9\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{2y(\tau)}{9\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$9\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + 2y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$9\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + 2\tau^r = 0$$

Simplifying gives

$$9r(r-1)\tau^r + 0\tau^r + 2\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$9r(r-1) + 0 + 2 = 0$$

Or

$$9r^2 - 9r + 2 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{3}$$

$$r_2 = \frac{2}{3}$$

Since the roots are real and distinct, then the general solution is

$$y(\tau) = c_1 y_1 + c_2 y_2$$

Where $y_1 = \tau^{r_1}$ and $y_2 = \tau^{r_2}$. Hence

$$y(\tau) = c_1 \tau^{\frac{1}{3}} + c_2 \tau^{\frac{2}{3}}$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{c_1 3^{\frac{2}{3}} \left(-\frac{1}{t^3}\right)^{\frac{1}{3}}}{3} + \frac{c_2 3^{\frac{1}{3}} \left(-\frac{1}{t^3}\right)^{\frac{2}{3}}}{3}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 3^{\frac{2}{3}} \left(-\frac{1}{t^3}\right)^{\frac{1}{3}}}{3} + \frac{c_2 3^{\frac{1}{3}} \left(-\frac{1}{t^3}\right)^{\frac{2}{3}}}{3} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 3^{\frac{2}{3}} \left(-\frac{1}{t^3}\right)^{\frac{1}{3}}}{3} + \frac{c_2 3^{\frac{1}{3}} \left(-\frac{1}{t^3}\right)^{\frac{2}{3}}}{3}$$

Verified OK.

8.22.4 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$t^2 y'' + 4ty' + 2y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{4}{t}$$

$$q(t) = \frac{2}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\begin{aligned} \tau' &= \frac{1}{c}\sqrt{q} \\ &= \frac{\sqrt{2}\sqrt{\frac{1}{t^2}}}{c} \\ \tau'' &= -\frac{\sqrt{2}}{c\sqrt{\frac{1}{t^2}}t^3} \end{aligned} \quad (6)$$

Substituting the above into (4) results in

$$\begin{aligned} p_1(\tau) &= \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \\ &= \frac{-\frac{\sqrt{2}}{c\sqrt{\frac{1}{t^2}}t^3} + \frac{4}{t}\frac{\sqrt{2}\sqrt{\frac{1}{t^2}}}{c}}{\left(\frac{\sqrt{2}\sqrt{\frac{1}{t^2}}}{c}\right)^2} \\ &= \frac{3c\sqrt{2}}{2} \end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned} y(\tau)'' + p_1y(\tau)' + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + \frac{3c\sqrt{2}}{2}\left(\frac{d}{d\tau}y(\tau)\right) + c^2y(\tau) &= 0 \end{aligned} \quad (7)$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{-\frac{3\sqrt{2}c\tau}{4}} \left(c_1 \cosh\left(\frac{\sqrt{2}c\tau}{4}\right) + ic_2 \sinh\left(\frac{\sqrt{2}c\tau}{4}\right) \right)$$

Now from (6)

$$\begin{aligned}\tau &= \int \frac{1}{c} \sqrt{q} dt \\ &= \frac{\int \sqrt{2} \sqrt{\frac{1}{t^2}} dt}{c} \\ &= \frac{\sqrt{2} \sqrt{\frac{1}{t^2}} t \ln(t)}{c}\end{aligned}$$

Substituting the above into the solution obtained gives

$$y = \frac{c_1 \cosh\left(\frac{\ln(t)}{2}\right) + ic_2 \sinh\left(\frac{\ln(t)}{2}\right)}{t^{\frac{3}{2}}}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \cosh\left(\frac{\ln(t)}{2}\right) + ic_2 \sinh\left(\frac{\ln(t)}{2}\right)}{t^{\frac{3}{2}}} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 \cosh\left(\frac{\ln(t)}{2}\right) + ic_2 \sinh\left(\frac{\ln(t)}{2}\right)}{t^{\frac{3}{2}}}$$

Verified OK.

8.22.5 Solving as second order change of variable on y method 1 ode

In normal form the given ode is written as

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$\begin{aligned}p(t) &= \frac{4}{t} \\ q(t) &= \frac{2}{t^2}\end{aligned}$$

Calculating the Liouville ode invariant Q given by

$$\begin{aligned}
 Q &= q - \frac{p'}{2} - \frac{p^2}{4} \\
 &= \frac{2}{t^2} - \frac{\left(\frac{4}{t}\right)'}{2} - \frac{\left(\frac{4}{t}\right)^2}{4} \\
 &= \frac{2}{t^2} - \frac{\left(-\frac{4}{t^2}\right)}{2} - \frac{\left(\frac{16}{t^2}\right)}{4} \\
 &= \frac{2}{t^2} - \left(-\frac{2}{t^2}\right) - \frac{4}{t^2} \\
 &= 0
 \end{aligned}$$

Since the Liouville ode invariant does not depend on the independent variable t then the transformation

$$y = v(t) z(t) \quad (3)$$

is used to change the original ode to a constant coefficients ode in v . In (3) the term $z(t)$ is given by

$$\begin{aligned}
 z(t) &= e^{-\left(\int \frac{p(t)}{2} dt\right)} \\
 &= e^{-\int \frac{4}{t} dt} \\
 &= \frac{1}{t^2}
 \end{aligned} \quad (5)$$

Hence (3) becomes

$$y = \frac{v(t)}{t^2} \quad (4)$$

Applying this change of variable to the original ode results in

$$v''(t) = 0$$

Which is now solved for $v(t)$ Integrating twice gives the solution

$$v(t) = c_1 t + c_2$$

Now that $v(t)$ is known, then

$$\begin{aligned}
 y &= v(t) z(t) \\
 &= (c_1 t + c_2) (z(t))
 \end{aligned} \quad (7)$$

But from (5)

$$z(t) = \frac{1}{t^2}$$

Hence (7) becomes

$$y = \frac{c_1 t + c_2}{t^2}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 t + c_2}{t^2} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 t + c_2}{t^2}$$

Verified OK.

8.22.6 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$t^2 y'' + 4ty' + 2y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{4}{t}$$
$$q(t) = \frac{2}{t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} + \frac{4n}{t^2} + \frac{2}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = -1 \quad (6)$$

Substituting this value in (3) gives

$$\begin{aligned} v''(t) + \frac{2v'(t)}{t} &= 0 \\ v''(t) + \frac{2v'(t)}{t} &= 0 \end{aligned} \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{2u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{2u}{t} \end{aligned}$$

Where $f(t) = -\frac{2}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= -\frac{2}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{2}{t} dt \\ \ln(u) &= -2 \ln(t) + c_1 \\ u &= e^{-2 \ln(t) + c_1} \\ &= \frac{c_1}{t^2} \end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned}v'(t) &= u(t) \\v(t) &= \int u(t) dt + c_2 \\&= -\frac{c_1}{t} + c_2\end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\&= \frac{-\frac{c_1}{t} + c_2}{t} \\&= \frac{c_2 t - c_1}{t^2}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{-\frac{c_1}{t} + c_2}{t} \tag{1}$$

Verification of solutions

$$y = \frac{-\frac{c_1}{t} + c_2}{t}$$

Verified OK.

8.22.7 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t t gives

$$\begin{aligned}\int (t^2 y'' + 4t y' + 2y) dt &= 0 \\y' t^2 + 2yt &= c_1\end{aligned}$$

Which is now solved for y .

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$\begin{aligned}p(t) &= \frac{2}{t} \\q(t) &= \frac{c_1}{t^2}\end{aligned}$$

Hence the ode is

$$y' + \frac{2y}{t} = \frac{c_1}{t^2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{2}{t} dt} \\ &= t^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{c_1}{t^2} \right) \\ \frac{d}{dt}(t^2 y) &= (t^2) \left(\frac{c_1}{t^2} \right) \\ d(t^2 y) &= c_1 dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t^2 y &= \int c_1 dt \\ t^2 y &= c_1 t + c_2\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t^2$ results in

$$y = \frac{c_1}{t} + \frac{c_2}{t^2}$$

which simplifies to

$$y = \frac{c_1 t + c_2}{t^2}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 t + c_2}{t^2} \tag{1}$$

Verification of solutions

$$y = \frac{c_1 t + c_2}{t^2}$$

Verified OK.

8.22.8 Solving as type second_order_integrable_as_is (not using ABC version)

Writing the ode as

$$t^2 y'' + 4ty' + 2y = 0$$

Integrating both sides of the ODE w.r.t t gives

$$\int (t^2 y'' + 4ty' + 2y) dt = 0$$
$$y' t^2 + 2yt = c_1$$

Which is now solved for y .

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{2}{t}$$
$$q(t) = \frac{c_1}{t^2}$$

Hence the ode is

$$y' + \frac{2y}{t} = \frac{c_1}{t^2}$$

The integrating factor μ is

$$\mu = e^{\int \frac{2}{t} dt}$$
$$= t^2$$

The ode becomes

$$\frac{d}{dt}(\mu y) = (\mu) \left(\frac{c_1}{t^2} \right)$$
$$\frac{d}{dt}(t^2 y) = (t^2) \left(\frac{c_1}{t^2} \right)$$
$$d(t^2 y) = c_1 dt$$

Integrating gives

$$t^2 y = \int c_1 dt$$
$$t^2 y = c_1 t + c_2$$

Dividing both sides by the integrating factor $\mu = t^2$ results in

$$y = \frac{c_1}{t} + \frac{c_2}{t^2}$$

which simplifies to

$$y = \frac{c_1 t + c_2}{t^2}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 t + c_2}{t^2} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 t + c_2}{t^2}$$

Verified OK.

8.22.9 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' + 4t y' + 2y = 0 \quad (1)$$

$$A y'' + B y' + C y = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t^2 \\ B &= 4t \\ C &= 2 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = y e^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = r z(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 0 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = 0 \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 406: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{4t}{t^2} dt} \\ &= z_1 e^{-2 \ln(t)} \\ &= z_1 \left(\frac{1}{t^2} \right) \end{aligned}$$

Which simplifies to

$$y_1 = \frac{1}{t^2}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{4t}{t^2} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{-4 \ln(t)}}{(y_1)^2} dt \\ &= y_1(t) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}
 y &= c_1 y_1 + c_2 y_2 \\
 &= c_1 \left(\frac{1}{t^2} \right) + c_2 \left(\frac{1}{t^2}(t) \right)
 \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t^2} + \frac{c_2}{t} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{t^2} + \frac{c_2}{t}$$

Verified OK.

8.22.10 Solving as exact linear second order ode ode

An ode of the form

$$p(t) y'' + q(t) y' + r(t) y = s(t)$$

is exact if

$$p''(t) - q'(t) + r(t) = 0 \tag{1}$$

For the given ode we have

$$\begin{aligned}
 p(x) &= t^2 \\
 q(x) &= 4t \\
 r(x) &= 2 \\
 s(x) &= 0
 \end{aligned}$$

Hence

$$\begin{aligned}
 p''(x) &= 2 \\
 q'(x) &= 4
 \end{aligned}$$

Therefore (1) becomes

$$2 - (4) + (2) = 0$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$(p(t) y' + (q(t) - p'(t)) y)' = s(x)$$

Integrating gives

$$p(t) y' + (q(t) - p'(t)) y = \int s(t) dt$$

Substituting the above values for p, q, r, s gives

$$y't^2 + 2yt = c_1$$

We now have a first order ode to solve which is

$$y't^2 + 2yt = c_1$$

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{2}{t}$$
$$q(t) = \frac{c_1}{t^2}$$

Hence the ode is

$$y' + \frac{2y}{t} = \frac{c_1}{t^2}$$

The integrating factor μ is

$$\mu = e^{\int \frac{2}{t} dt}$$
$$= t^2$$

The ode becomes

$$\frac{d}{dt}(\mu y) = (\mu) \left(\frac{c_1}{t^2} \right)$$
$$\frac{d}{dt}(t^2 y) = (t^2) \left(\frac{c_1}{t^2} \right)$$
$$d(t^2 y) = c_1 dt$$

Integrating gives

$$t^2 y = \int c_1 dt$$
$$t^2 y = c_1 t + c_2$$

Dividing both sides by the integrating factor $\mu = t^2$ results in

$$y = \frac{c_1}{t} + \frac{c_2}{t^2}$$

which simplifies to

$$y = \frac{c_1 t + c_2}{t^2}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 t + c_2}{t^2} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 t + c_2}{t^2}$$

Verified OK.

8.22.11 Maple step by step solution

Let's solve

$$y'' t^2 + 4t y' + 2y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{4y'}{t} - \frac{2y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{4y'}{t} + \frac{2y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y'' t^2 + 4t y' + 2y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds} y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s) \right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s) \right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2} \right) t^2 + 4 \frac{d}{ds}y(s) + 2y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) + 3 \frac{d}{ds}y(s) + 2y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 + 3r + 2 = 0$$

- Factor the characteristic polynomial

$$(r + 2)(r + 1) = 0$$

- Roots of the characteristic polynomial

$$r = (-2, -1)$$

- 1st solution of the ODE

$$y_1(s) = e^{-2s}$$

- 2nd solution of the ODE

$$y_2(s) = e^{-s}$$

- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$

- Substitute in solutions

$$y(s) = c_1 e^{-2s} + c_2 e^{-s}$$

- Change variables back using $s = \ln(t)$

$$y = \frac{c_1}{t^2} + \frac{c_2}{t}$$

- Simplify

$$y = \frac{c_1}{t^2} + \frac{c_2}{t}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
<- LODE of Euler type successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 13

```
dsolve(t^2*diff(y(t),t$2)+ 4*t*diff(y(t),t)+2*y(t) = 0,y(t), singsol=all)
```

$$y(t) = \frac{c_2 t + c_1}{t^2}$$

✓ Solution by Mathematica

Time used: 0.016 (sec). Leaf size: 34

```
DSolve[t^2*y''[t]+4*t*y'[t]+y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t^{-\frac{3}{2}-\frac{\sqrt{5}}{2}} \left(c_2 t^{\sqrt{5}} + c_1 \right)$$

8.23 problem 37

8.23.1 Solving as second order euler ode	2239
8.23.2 Solving as second order change of variable on x method 2 ode .	2240
8.23.3 Solving as second order change of variable on x method 1 ode .	2243
8.23.4 Solving as second order change of variable on y method 2 ode .	2245
8.23.5 Solving using Kovacic algorithm	2247
8.23.6 Maple step by step solution	2253

Internal problem ID [645]

Internal file name [OUTPUT/645_Sunday_June_05_2022_01_46_21_AM_42626299/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 37.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_Emden , _Fowler]]
```

$$t^2y'' + 3ty' + \frac{5y}{4} = 0$$

The ode can be written as

$$4t^2y'' + 12ty' + 5y = 0$$

Which shows it is a Euler ODE.

8.23.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$4t^2(r(r-1))t^{r-2} + 12trt^{r-1} + 5t^r = 0$$

Simplifying gives

$$4r(r-1)t^r + 12rt^r + 5t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$4r(r-1) + 12r + 5 = 0$$

Or

$$4r^2 + 8r + 5 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$\begin{aligned} r_1 &= -1 - \frac{i}{2} \\ r_2 &= -1 + \frac{i}{2} \end{aligned}$$

The roots are complex conjugate of each others. Let the roots be

$$\begin{aligned} r_1 &= \alpha + i\beta \\ r_2 &= \alpha - i\beta \end{aligned}$$

Where in this case $\alpha = -1$ and $\beta = -\frac{1}{2}$. Hence the solution becomes

$$\begin{aligned} y &= c_1 t^{r_1} + c_2 t^{r_2} \\ &= c_1 t^{\alpha+i\beta} + c_2 t^{\alpha-i\beta} \\ &= t^\alpha (c_1 t^{i\beta} + c_2 t^{-i\beta}) \\ &= t^\alpha \left(c_1 e^{\ln(t^{i\beta})} + c_2 e^{\ln(t^{-i\beta})} \right) \\ &= t^\alpha \left(c_1 e^{i(\beta \ln t)} + c_2 e^{-i(\beta \ln t)} \right) \end{aligned}$$

Using the values for $\alpha = -1, \beta = -\frac{1}{2}$, the above becomes

$$y = t^{-1} \left(c_1 e^{-\frac{i \ln(t)}{2}} + c_2 e^{\frac{i \ln(t)}{2}} \right)$$

Using Euler relation, the expression $c_1 e^{iA} + c_2 e^{-iA}$ is transformed to $c_1 \cos A + c_1 \sin A$ where the constants are free to change. Applying this to the above result gives

$$y = \frac{1}{t} \left(c_1 \cos \left(\frac{\ln(t)}{2} \right) + c_2 \sin \left(\frac{\ln(t)}{2} \right) \right)$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \cos \left(\frac{\ln(t)}{2} \right) + c_2 \sin \left(\frac{\ln(t)}{2} \right)}{t} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 \cos \left(\frac{\ln(t)}{2} \right) + c_2 \sin \left(\frac{\ln(t)}{2} \right)}{t}$$

Verified OK.

8.23.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$4t^2 y'' + 12ty' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{3}{t}$$

$$q(t) = \frac{5}{4t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t) \tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t)dt)} dt \\ &= \int e^{-(\int \frac{3}{t} dt)} dt \\ &= \int e^{-3\ln(t)} dt \\ &= \int \frac{1}{t^3} dt \\ &= -\frac{1}{2t^2} \end{aligned} \tag{6}$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{\frac{5}{4t^2}}{\frac{1}{t^6}} \\ &= \frac{5t^4}{4} \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + \frac{5t^4y(\tau)}{4} &= 0 \end{aligned}$$

But in terms of τ

$$\frac{5t^4}{4} = \frac{5}{16\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{5y(\tau)}{16\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$16 \left(\frac{d^2}{d\tau^2}y(\tau) \right) \tau^2 + 5y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$16\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + 5\tau^r = 0$$

Simplifying gives

$$16r(r-1)\tau^r + 0\tau^r + 5\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$16r(r-1) + 0 + 5 = 0$$

Or

$$16r^2 - 16r + 5 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2} - \frac{i}{4}$$

$$r_2 = \frac{1}{2} + \frac{i}{4}$$

The roots are complex conjugate of each others. Let the roots be

$$r_1 = \alpha + i\beta$$

$$r_2 = \alpha - i\beta$$

Where in this case $\alpha = \frac{1}{2}$ and $\beta = -\frac{1}{4}$. Hence the solution becomes

$$y(\tau) = c_1\tau^{r_1} + c_2\tau^{r_2}$$

$$= c_1\tau^{\alpha+i\beta} + c_2\tau^{\alpha-i\beta}$$

$$= \tau^\alpha(c_1\tau^{i\beta} + c_2\tau^{-i\beta})$$

$$= \tau^\alpha(c_1e^{\ln(\tau^{i\beta})} + c_2e^{\ln(\tau^{-i\beta})})$$

$$= \tau^\alpha(c_1e^{i(\beta \ln \tau)} + c_2e^{-i(\beta \ln \tau)})$$

Using the values for $\alpha = \frac{1}{2}, \beta = -\frac{1}{4}$, the above becomes

$$y(\tau) = \tau^{\frac{1}{2}} \left(c_1 e^{-\frac{i \ln(\tau)}{4}} + c_2 e^{\frac{i \ln(\tau)}{4}} \right)$$

Using Euler relation, the expression $c_1 e^{iA} + c_2 e^{-iA}$ is transformed to $c_1 \cos A + c_1 \sin A$ where the constants are free to change. Applying this to the above result gives

$$y(\tau) = \sqrt{\tau} \left(c_1 \cos \left(\frac{\ln(\tau)}{4} \right) + c_2 \sin \left(\frac{\ln(\tau)}{4} \right) \right)$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{\left(c_1 \cos \left(-\frac{\ln(2)}{4} + \frac{\ln\left(-\frac{1}{t^2}\right)}{4} \right) + c_2 \sin \left(-\frac{\ln(2)}{4} + \frac{\ln\left(-\frac{1}{t^2}\right)}{4} \right) \right) \sqrt{2} \sqrt{-\frac{1}{t^2}}}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(c_1 \cos \left(-\frac{\ln(2)}{4} + \frac{\ln\left(-\frac{1}{t^2}\right)}{4} \right) + c_2 \sin \left(-\frac{\ln(2)}{4} + \frac{\ln\left(-\frac{1}{t^2}\right)}{4} \right) \right) \sqrt{2} \sqrt{-\frac{1}{t^2}}}{2} \quad (1)$$

Verification of solutions

$$y = \frac{\left(c_1 \cos \left(-\frac{\ln(2)}{4} + \frac{\ln\left(-\frac{1}{t^2}\right)}{4} \right) + c_2 \sin \left(-\frac{\ln(2)}{4} + \frac{\ln\left(-\frac{1}{t^2}\right)}{4} \right) \right) \sqrt{2} \sqrt{-\frac{1}{t^2}}}{2}$$

Verified OK.

8.23.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$4t^2y'' + 12ty' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{3}{t}$$

$$q(t) = \frac{5}{4t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1 \left(\frac{d}{d\tau}y(\tau) \right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\begin{aligned}\tau' &= \frac{1}{c}\sqrt{q} \\ &= \frac{\sqrt{5}\sqrt{\frac{1}{t^2}}}{2c} \\ \tau'' &= -\frac{\sqrt{5}}{2c\sqrt{\frac{1}{t^2}}t^3}\end{aligned}\tag{6}$$

Substituting the above into (4) results in

$$\begin{aligned}p_1(\tau) &= \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \\ &= \frac{-\frac{\sqrt{5}}{2c\sqrt{\frac{1}{t^2}}t^3} + \frac{3}{t}\frac{\sqrt{5}\sqrt{\frac{1}{t^2}}}{2c}}{\left(\frac{\sqrt{5}\sqrt{\frac{1}{t^2}}}{2c}\right)^2} \\ &= \frac{4c\sqrt{5}}{5}\end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned}y(\tau)'' + p_1y(\tau)' + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + \frac{4c\sqrt{5}}{5}\left(\frac{d}{d\tau}y(\tau)\right) + c^2y(\tau) &= 0\end{aligned}\tag{7}$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{-\frac{2\sqrt{5}c\tau}{5}}\left(c_1\cos\left(\frac{\sqrt{5}c\tau}{5}\right) + c_2\sin\left(\frac{\sqrt{5}c\tau}{5}\right)\right)$$

Now from (6)

$$\begin{aligned}\tau &= \int \frac{1}{c}\sqrt{q} dt \\ &= \frac{\int \frac{\sqrt{5}\sqrt{\frac{1}{t^2}}}{2} dt}{c} \\ &= \frac{\sqrt{5}\sqrt{\frac{1}{t^2}}t\ln(t)}{2c}\end{aligned}$$

Substituting the above into the solution obtained gives

$$y = \frac{c_1 \cos\left(\frac{\ln(t)}{2}\right) + c_2 \sin\left(\frac{\ln(t)}{2}\right)}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \cos\left(\frac{\ln(t)}{2}\right) + c_2 \sin\left(\frac{\ln(t)}{2}\right)}{t} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 \cos\left(\frac{\ln(t)}{2}\right) + c_2 \sin\left(\frac{\ln(t)}{2}\right)}{t}$$

Verified OK.

8.23.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$4t^2 y'' + 12ty' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{3}{t}$$
$$q(t) = \frac{5}{4t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variable is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} + \frac{3n}{t^2} + \frac{5}{4t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = -1 + \frac{i}{2} \quad (6)$$

Substituting this value in (3) gives

$$\begin{aligned} v''(t) + \left(\frac{-2+i}{t} + \frac{3}{t} \right) v'(t) &= 0 \\ v''(t) + \frac{(1+i)v'(t)}{t} &= 0 \end{aligned} \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{(1+i)u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{(-1-i)u}{t} \end{aligned}$$

Where $f(t) = \frac{-1-i}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= \frac{-1-i}{t} dt \\ \int \frac{1}{u} du &= \int \frac{-1-i}{t} dt \\ \ln(u) &= (-1-i) \ln(t) + c_1 \\ u &= e^{(-1-i) \ln(t) + c_1} \\ &= c_1 e^{(-1-i) \ln(t)} \end{aligned}$$

Which simplifies to

$$u(t) = \frac{c_1 t^{-i}}{t}$$

Now that $u(t)$ is known, then

$$\begin{aligned}v'(t) &= u(t) \\v(t) &= \int u(t) dt + c_2 \\&= i t^{-i} c_1 + c_2\end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\&= (i t^{-i} c_1 + c_2) t^{-1+\frac{i}{2}} \\&= t^{-1-\frac{i}{2}} (i c_1 + t^i c_2)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = (i t^{-i} c_1 + c_2) t^{-1+\frac{i}{2}} \quad (1)$$

Verification of solutions

$$y = (i t^{-i} c_1 + c_2) t^{-1+\frac{i}{2}}$$

Verified OK.

8.23.5 Solving using Kovacic algorithm

Writing the ode as

$$4t^2 y'' + 12ty' + 5y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= 4t^2 \\B &= 12t \\C &= 5\end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{2t^2} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 2t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(-\frac{1}{2t^2}\right) z(t) \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 408: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 2 - 0 \\
 &= 2
 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 2t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{2t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition

of r given above. Therefore $b = -\frac{1}{2}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} + \frac{i}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} - \frac{i}{2} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{1}{2t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{1}{2}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} + \frac{i}{2} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} - \frac{i}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{1}{2t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2} + \frac{i}{2}$	$\frac{1}{2} - \frac{i}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2} + \frac{i}{2}$	$\frac{1}{2} - \frac{i}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = \frac{1}{2} - \frac{i}{2}$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= \frac{1}{2} - \frac{i}{2} - \left(\frac{1}{2} - \frac{i}{2} \right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{t - c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= \frac{\frac{1}{2} - \frac{i}{2}}{t} + (-)(0) \\ &= \frac{\frac{1}{2} - \frac{i}{2}}{t} \\ &= \frac{\frac{1}{2} - \frac{i}{2}}{t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r)p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2 \left(\frac{\frac{1}{2} - \frac{i}{2}}{t} \right) (0) + \left(\left(\frac{-\frac{1}{2} + \frac{i}{2}}{t^2} \right) + \left(\frac{\frac{1}{2} - \frac{i}{2}}{t} \right)^2 - \left(-\frac{1}{2t^2} \right) \right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \frac{\frac{1}{2} - \frac{i}{2}}{t} dt} \\ &= t^{\frac{1}{2} - \frac{i}{2}} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\&= z_1 e^{-\int \frac{1}{2} \frac{12t}{4t^2} dt} \\&= z_1 e^{-\frac{3 \ln(t)}{2}} \\&= z_1 \left(\frac{1}{t^{\frac{3}{2}}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = t^{-1-\frac{i}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{12t}{4t^2} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{-3 \ln(t)}}{(y_1)^2} dt \\&= y_1 (-it^i)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 \left(t^{-1-\frac{i}{2}} \right) + c_2 \left(t^{-1-\frac{i}{2}} (-it^i) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 t^{-1-\frac{i}{2}} - ic_2 t^{-1+\frac{i}{2}} \quad (1)$$

Verification of solutions

$$y = c_1 t^{-1-\frac{i}{2}} - ic_2 t^{-1+\frac{i}{2}}$$

Verified OK.

8.23.6 Maple step by step solution

Let's solve

$$4y''t^2 + 12ty' + 5y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{3y'}{t} - \frac{5y}{4t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{3y'}{t} + \frac{5y}{4t^2} = 0$$

- Multiply by denominators of the ODE

$$4y''t^2 + 12ty' + 5y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$4\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right)t^2 + 12\frac{d}{ds}y(s) + 5y(s) = 0$$

- Simplify

$$4\frac{d^2}{ds^2}y(s) + 8\frac{d}{ds}y(s) + 5y(s) = 0$$

- Isolate 2nd derivative

$$\frac{d^2}{ds^2}y(s) = -2\frac{d}{ds}y(s) - \frac{5y(s)}{4}$$

- Group terms with $y(s)$ on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$\frac{d^2}{ds^2}y(s) + 2\frac{d}{ds}y(s) + \frac{5y(s)}{4} = 0$$

- Characteristic polynomial of ODE

$$r^2 + 2r + \frac{5}{4} = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-2) \pm (\sqrt{-1})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-1 - \frac{1}{2}, -1 + \frac{1}{2}\right)$$

- 1st solution of the ODE

$$y_1(s) = e^{-s} \cos\left(\frac{s}{2}\right)$$

- 2nd solution of the ODE

$$y_2(s) = e^{-s} \sin\left(\frac{s}{2}\right)$$

- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$

- Substitute in solutions

$$y(s) = c_1 e^{-s} \cos\left(\frac{s}{2}\right) + c_2 e^{-s} \sin\left(\frac{s}{2}\right)$$

- Change variables back using $s = \ln(t)$

$$y = \frac{c_1 \cos\left(\frac{\ln(t)}{2}\right)}{t} + \frac{c_2 \sin\left(\frac{\ln(t)}{2}\right)}{t}$$

- Simplify

$$y = \frac{c_1 \cos\left(\frac{\ln(t)}{2}\right)}{t} + \frac{c_2 \sin\left(\frac{\ln(t)}{2}\right)}{t}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 23

```
dsolve(t^2*diff(y(t),t$2)+ 3*t*diff(y(t),t)+125/100*y(t) = 0,y(t), singsol=all)
```

$$y(t) = \frac{c_1 \sin\left(\frac{\ln(t)}{2}\right) + c_2 \cos\left(\frac{\ln(t)}{2}\right)}{t}$$

✓ Solution by Mathematica

Time used: 0.025 (sec). Leaf size: 30

```
DSolve[t^2*y'[t]+3*t*y'[t]+125/100*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{c_2 \cos\left(\frac{\log(t)}{2}\right) + c_1 \sin\left(\frac{\log(t)}{2}\right)}{t}$$

8.24 problem 38

8.24.1 Solving as second order euler ode ode	2257
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8.24.3 Solving as second order change of variable on y method 2 ode .	2260
8.24.4 Solving as second order integrable as is ode	2263
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Internal problem ID [646]

Internal file name [OUTPUT/646_Sunday_June_05_2022_01_46_22_AM_79982171/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 38.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "exact linear second order ode", "second_order_integrable_as_is", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_2nd_order , _exact , _linear , _homogeneous]]
```

$$t^2y'' - 4ty' - 6y = 0$$

8.24.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} - 4trt^{r-1} - 6t^r = 0$$

Simplifying gives

$$r(r-1)t^r - 4rt^r - 6t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r-1) - 4r - 6 = 0$$

Or

$$r^2 - 5r - 6 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -1$$

$$r_2 = 6$$

Since the roots are real and distinct, then the general solution is

$$y = c_1y_1 + c_2y_2$$

Where $y_1 = t^{r_1}$ and $y_2 = t^{r_2}$. Hence

$$y = \frac{c_1}{t} + c_2t^6$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t} + c_2t^6 \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{t} + c_2t^6$$

Verified OK.

8.24.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$t^2 y'' - 4ty' - 6y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = -\frac{4}{t}$$
$$q(t) = -\frac{6}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t) \tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t) dt)} dt \\ &= \int e^{-(\int -\frac{4}{t} dt)} dt \\ &= \int e^{4 \ln(t)} dt \\ &= \int t^4 dt \\ &= \frac{t^5}{5} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{-\frac{6}{t^2}}{t^8} \\ &= -\frac{6}{t^{10}} \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) - \frac{6y(\tau)}{t^{10}} &= 0 \end{aligned}$$

But in terms of τ

$$-\frac{6}{t^{10}} = -\frac{6}{25\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) - \frac{6y(\tau)}{25\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$25 \left(\frac{d^2}{d\tau^2}y(\tau) \right) \tau^2 - 6y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$25\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} - 6\tau^r = 0$$

Simplifying gives

$$25r(r-1)\tau^r + 0\tau^r - 6\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$25r(r-1) + 0 - 6 = 0$$

Or

$$25r^2 - 25r - 6 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -\frac{1}{5}$$

$$r_2 = \frac{6}{5}$$

Since the roots are real and distinct, then the general solution is

$$y(\tau) = c_1 y_1 + c_2 y_2$$

Where $y_1 = \tau^{r_1}$ and $y_2 = \tau^{r_2}$. Hence

$$y(\tau) = \frac{c_1}{\tau^{\frac{1}{5}}} + c_2 \tau^{\frac{6}{5}}$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{5^{\frac{1}{5}} \left(c_2 5^{\frac{3}{5}} t^5 (t^5)^{\frac{2}{5}} + 25c_1 \right)}{25 (t^5)^{\frac{1}{5}}}$$

Summary

The solution(s) found are the following

$$y = \frac{5^{\frac{1}{5}} \left(c_2 5^{\frac{3}{5}} t^5 (t^5)^{\frac{2}{5}} + 25c_1 \right)}{25 (t^5)^{\frac{1}{5}}} \quad (1)$$

Verification of solutions

$$y = \frac{5^{\frac{1}{5}} \left(c_2 5^{\frac{3}{5}} t^5 (t^5)^{\frac{2}{5}} + 25c_1 \right)}{25 (t^5)^{\frac{1}{5}}}$$

Verified OK.

8.24.3 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$t^2 y'' - 4ty' - 6y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = -\frac{4}{t}$$
$$q(t) = -\frac{6}{t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} - \frac{4n}{t} - \frac{6}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = 6 \quad (6)$$

Substituting this value in (3) gives

$$v''(t) + \frac{8v'(t)}{t} = 0$$
$$v''(t) + \frac{8v'(t)}{t} = 0 \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{8u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned}u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{8u}{t}\end{aligned}$$

Where $f(t) = -\frac{8}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= -\frac{8}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{8}{t} dt \\ \ln(u) &= -8 \ln(t) + c_1 \\ u &= e^{-8 \ln(t) + c_1} \\ &= \frac{c_1}{t^8}\end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned}v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= -\frac{c_1}{7t^7} + c_2\end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\ &= \left(-\frac{c_1}{7t^7} + c_2\right) t^6 \\ &= \frac{7c_2 t^7 - c_1}{7t}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \left(-\frac{c_1}{7t^7} + c_2\right) t^6 \tag{1}$$

Verification of solutions

$$y = \left(-\frac{c_1}{7t^7} + c_2\right) t^6$$

Verified OK.

8.24.4 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t t gives

$$\int (t^2 y'' - 4ty' - 6y) dt = 0$$
$$y't^2 - 6yt = c_1$$

Which is now solved for y .

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{6}{t}$$
$$q(t) = \frac{c_1}{t^2}$$

Hence the ode is

$$y' - \frac{6y}{t} = \frac{c_1}{t^2}$$

The integrating factor μ is

$$\mu = e^{\int -\frac{6}{t} dt}$$
$$= \frac{1}{t^6}$$

The ode becomes

$$\frac{d}{dt}(\mu y) = (\mu) \left(\frac{c_1}{t^2}\right)$$
$$\frac{d}{dt}\left(\frac{y}{t^6}\right) = \left(\frac{1}{t^6}\right) \left(\frac{c_1}{t^2}\right)$$
$$d\left(\frac{y}{t^6}\right) = \left(\frac{c_1}{t^8}\right) dt$$

Integrating gives

$$\frac{y}{t^6} = \int \frac{c_1}{t^8} dt$$
$$\frac{y}{t^6} = -\frac{c_1}{7t^7} + c_2$$

Dividing both sides by the integrating factor $\mu = \frac{1}{t^6}$ results in

$$y = -\frac{c_1}{7t} + c_2t^6$$

Summary

The solution(s) found are the following

$$y = -\frac{c_1}{7t} + c_2t^6 \quad (1)$$

Verification of solutions

$$y = -\frac{c_1}{7t} + c_2t^6$$

Verified OK.

8.24.5 Solving as type second_order_integrable_as_is (not using ABC version)

Writing the ode as

$$t^2y'' - 4ty' - 6y = 0$$

Integrating both sides of the ODE w.r.t t gives

$$\int (t^2y'' - 4ty' - 6y) dt = 0$$
$$y't^2 - 6yt = c_1$$

Which is now solved for y .

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{6}{t}$$
$$q(t) = \frac{c_1}{t^2}$$

Hence the ode is

$$y' - \frac{6y}{t} = \frac{c_1}{t^2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{6}{t} dt} \\ &= \frac{1}{t^6}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{c_1}{t^2}\right) \\ \frac{d}{dt}\left(\frac{y}{t^6}\right) &= \left(\frac{1}{t^6}\right) \left(\frac{c_1}{t^2}\right) \\ d\left(\frac{y}{t^6}\right) &= \left(\frac{c_1}{t^8}\right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y}{t^6} &= \int \frac{c_1}{t^8} dt \\ \frac{y}{t^6} &= -\frac{c_1}{7t^7} + c_2\end{aligned}$$

Dividing both sides by the integrating factor $\mu = \frac{1}{t^6}$ results in

$$y = -\frac{c_1}{7t} + c_2 t^6$$

Summary

The solution(s) found are the following

$$y = -\frac{c_1}{7t} + c_2 t^6 \tag{1}$$

Verification of solutions

$$y = -\frac{c_1}{7t} + c_2 t^6$$

Verified OK.

8.24.6 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' - 4ty' - 6y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= t^2 \\B &= -4t \\C &= -6\end{aligned}\tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t)\tag{4}$$

Where r is given by

$$\begin{aligned}r &= \frac{s}{t} \\&= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}\end{aligned}\tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{12}{t^2}\tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned}s &= 12 \\t &= t^2\end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(\frac{12}{t^2}\right) z(t)\tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 410: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 2 - 0 \\
 &= 2
 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{12}{t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition

of r given above. Therefore $b = 12$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = 4 \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = -3 \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = \frac{12}{t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = 12$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = 4 \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = -3 \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = \frac{12}{t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	4	-3

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	4	-3

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = -3$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= -3 - (-3) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t-c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{t-c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= -\frac{3}{t} + (-)(0) \\ &= -\frac{3}{t} \\ &= -\frac{3}{t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r)p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2\left(-\frac{3}{t}\right)(0) + \left(\left(\frac{3}{t^2}\right) + \left(-\frac{3}{t}\right)^2 - \left(\frac{12}{t^2}\right)\right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int -\frac{3}{t} dt} \\ &= \frac{1}{t^3} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\&= z_1 e^{-\int \frac{1}{2} \frac{-4t}{t^2} dt} \\&= z_1 e^{2 \ln(t)} \\&= z_1 (t^2)\end{aligned}$$

Which simplifies to

$$y_1 = \frac{1}{t}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-4t}{t^2} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{4 \ln(t)}}{(y_1)^2} dt \\&= y_1 \left(\frac{t^7}{7} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 \left(\frac{1}{t} \right) + c_2 \left(\frac{1}{t} \left(\frac{t^7}{7} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t} + \frac{c_2 t^6}{7} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{t} + \frac{c_2 t^6}{7}$$

Verified OK.

8.24.7 Solving as exact linear second order ode

An ode of the form

$$p(t) y'' + q(t) y' + r(t) y = s(t)$$

is exact if

$$p''(t) - q'(t) + r(t) = 0 \tag{1}$$

For the given ode we have

$$\begin{aligned} p(x) &= t^2 \\ q(x) &= -4t \\ r(x) &= -6 \\ s(x) &= 0 \end{aligned}$$

Hence

$$\begin{aligned} p''(x) &= 2 \\ q'(x) &= -4 \end{aligned}$$

Therefore (1) becomes

$$2 - (-4) + (-6) = 0$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$(p(t) y' + (q(t) - p'(t)) y)' = s(x)$$

Integrating gives

$$p(t) y' + (q(t) - p'(t)) y = \int s(t) dt$$

Substituting the above values for p, q, r, s gives

$$y't^2 - 6yt = c_1$$

We now have a first order ode to solve which is

$$y't^2 - 6yt = c_1$$

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{6}{t}$$
$$q(t) = \frac{c_1}{t^2}$$

Hence the ode is

$$y' - \frac{6y}{t} = \frac{c_1}{t^2}$$

The integrating factor μ is

$$\mu = e^{\int -\frac{6}{t} dt}$$
$$= \frac{1}{t^6}$$

The ode becomes

$$\frac{d}{dt}(\mu y) = (\mu) \left(\frac{c_1}{t^2}\right)$$
$$\frac{d}{dt}\left(\frac{y}{t^6}\right) = \left(\frac{1}{t^6}\right) \left(\frac{c_1}{t^2}\right)$$
$$d\left(\frac{y}{t^6}\right) = \left(\frac{c_1}{t^8}\right) dt$$

Integrating gives

$$\frac{y}{t^6} = \int \frac{c_1}{t^8} dt$$
$$\frac{y}{t^6} = -\frac{c_1}{7t^7} + c_2$$

Dividing both sides by the integrating factor $\mu = \frac{1}{t^6}$ results in

$$y = -\frac{c_1}{7t} + c_2 t^6$$

Summary

The solution(s) found are the following

$$y = -\frac{c_1}{7t} + c_2 t^6 \tag{1}$$

Verification of solutions

$$y = -\frac{c_1}{7t} + c_2 t^6$$

Verified OK.

8.24.8 Maple step by step solution

Let's solve

$$y''t^2 - 4ty' - 6y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{4y'}{t} + \frac{6y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{4y'}{t} - \frac{6y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 - 4ty' - 6y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right) t^2 - 4\frac{d}{ds}y(s) - 6y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) - 5\frac{d}{ds}y(s) - 6y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 - 5r - 6 = 0$$

- Factor the characteristic polynomial
 $(r + 1)(r - 6) = 0$
- Roots of the characteristic polynomial
 $r = (-1, 6)$
- 1st solution of the ODE
 $y_1(s) = e^{-s}$
- 2nd solution of the ODE
 $y_2(s) = e^{6s}$
- General solution of the ODE
 $y(s) = c_1 y_1(s) + c_2 y_2(s)$
- Substitute in solutions
 $y(s) = c_1 e^{-s} + c_2 e^{6s}$
- Change variables back using $s = \ln(t)$
 $y = \frac{c_1}{t} + c_2 t^6$
- Simplify
 $y = \frac{c_1}{t} + c_2 t^6$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve(t^2*diff(y(t),t$2)- 4*t*diff(y(t),t)-6*y(t) = 0,y(t), singsol=all)
```

$$y(t) = \frac{c_1 t^7 + c_2}{t}$$

✓ Solution by Mathematica

Time used: 0.01 (sec). Leaf size: 18

```
DSolve[t^2*y''[t]-4*t*y'[t]-6*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{c_2 t^7 + c_1}{t}$$

8.25 problem 39

8.25.1 Solving as second order euler ode	2277
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Internal problem ID [647]

Internal file name [OUTPUT/647_Sunday_June_05_2022_01_46_22_AM_3994398/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 39.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_1", "second_order_change_of_variable_on_y_method_2", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_Emden, _Fowler], [_2nd_order, _linear, `_with_symmetry_[0,F(x)]`]]
```

$$t^2y'' - 4ty' + 6y = 0$$

8.25.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} - 4trt^{r-1} + 6t^r = 0$$

Simplifying gives

$$r(r-1)t^r - 4rt^r + 6t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r-1) - 4r + 6 = 0$$

Or

$$r^2 - 5r + 6 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = 2$$

$$r_2 = 3$$

Since the roots are real and distinct, then the general solution is

$$y = c_1y_1 + c_2y_2$$

Where $y_1 = t^{r_1}$ and $y_2 = t^{r_2}$. Hence

$$y = c_2t^3 + c_1t^2$$

Summary

The solution(s) found are the following

$$y = c_2t^3 + c_1t^2 \tag{1}$$

Verification of solutions

$$y = c_2t^3 + c_1t^2$$

Verified OK.

8.25.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t)y' + \frac{(p(t)^2 + p'(t))y}{2} = f(t)$$

Where $p(t) = -\frac{4}{t}$. Therefore, there is an integrating factor given by

$$\begin{aligned}M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int -\frac{4}{t} dx} \\ &= \frac{1}{t^2}\end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned}(M(x)y)'' &= 0 \\ \left(\frac{y}{t^2}\right)'' &= 0\end{aligned}$$

Integrating once gives

$$\left(\frac{y}{t^2}\right)' = c_1$$

Integrating again gives

$$\left(\frac{y}{t^2}\right) = c_1 t + c_2$$

Hence the solution is

$$y = \frac{c_1 t + c_2}{\frac{1}{t^2}}$$

Or

$$y = c_1 t^3 + c_2 t^2$$

Summary

The solution(s) found are the following

$$y = c_1 t^3 + c_2 t^2 \tag{1}$$

Verification of solutions

$$y = c_1 t^3 + c_2 t^2$$

Verified OK.

8.25.3 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$t^2 y'' - 4ty' + 6y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = -\frac{4}{t}$$
$$q(t) = \frac{6}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t) \tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t) dt)} dt \\ &= \int e^{-(\int -\frac{4}{t} dt)} dt \\ &= \int e^{4 \ln(t)} dt \\ &= \int t^4 dt \\ &= \frac{t^5}{5} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{6}{t^2} \\ &= \frac{6}{t^{10}} \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + \frac{6y(\tau)}{t^{10}} &= 0 \end{aligned}$$

But in terms of τ

$$\frac{6}{t^{10}} = \frac{6}{25\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{6y(\tau)}{25\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$25 \left(\frac{d^2}{d\tau^2}y(\tau) \right) \tau^2 + 6y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$25\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + 6\tau^r = 0$$

Simplifying gives

$$25r(r-1)\tau^r + 0\tau^r + 6\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$25r(r-1) + 0 + 6 = 0$$

Or

$$25r^2 - 25r + 6 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{2}{5}$$

$$r_2 = \frac{3}{5}$$

Since the roots are real and distinct, then the general solution is

$$y(\tau) = c_1 y_1 + c_2 y_2$$

Where $y_1 = \tau^{r_1}$ and $y_2 = \tau^{r_2}$. Hence

$$y(\tau) = c_1 \tau^{\frac{2}{5}} + c_2 \tau^{\frac{3}{5}}$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{c_1 5^{\frac{3}{5}} (t^5)^{\frac{2}{5}}}{5} + \frac{c_2 5^{\frac{2}{5}} (t^5)^{\frac{3}{5}}}{5}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 5^{\frac{3}{5}} (t^5)^{\frac{2}{5}}}{5} + \frac{c_2 5^{\frac{2}{5}} (t^5)^{\frac{3}{5}}}{5} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 5^{\frac{3}{5}} (t^5)^{\frac{2}{5}}}{5} + \frac{c_2 5^{\frac{2}{5}} (t^5)^{\frac{3}{5}}}{5}$$

Verified OK.

8.25.4 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$t^2 y'' - 4ty' + 6y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = -\frac{4}{t}$$

$$q(t) = \frac{6}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\begin{aligned} \tau' &= \frac{1}{c}\sqrt{q} \\ &= \frac{\sqrt{6}\sqrt{\frac{1}{t^2}}}{c} \\ \tau'' &= -\frac{\sqrt{6}}{c\sqrt{\frac{1}{t^2}}t^3} \end{aligned} \quad (6)$$

Substituting the above into (4) results in

$$\begin{aligned} p_1(\tau) &= \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \\ &= \frac{-\frac{\sqrt{6}}{c\sqrt{\frac{1}{t^2}}t^3} - \frac{4}{t}\frac{\sqrt{6}\sqrt{\frac{1}{t^2}}}{c}}{\left(\frac{\sqrt{6}\sqrt{\frac{1}{t^2}}}{c}\right)^2} \\ &= -\frac{5c\sqrt{6}}{6} \end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned} y(\tau)'' + p_1y(\tau)' + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) - \frac{5c\sqrt{6}}{6}\left(\frac{d}{d\tau}y(\tau)\right) + c^2y(\tau) &= 0 \end{aligned} \quad (7)$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{\frac{5\sqrt{6}c\tau}{12}} \left(c_1 \cosh\left(\frac{\sqrt{6}c\tau}{12}\right) + ic_2 \sinh\left(\frac{\sqrt{6}c\tau}{12}\right) \right)$$

Now from (6)

$$\begin{aligned}\tau &= \int \frac{1}{c} \sqrt{q} dt \\ &= \frac{\int \sqrt{6} \sqrt{\frac{1}{t^2}} dt}{c} \\ &= \frac{\sqrt{6} \sqrt{\frac{1}{t^2}} t \ln(t)}{c}\end{aligned}$$

Substituting the above into the solution obtained gives

$$y = t^{\frac{5}{2}} \left(c_1 \cosh \left(\frac{\ln(t)}{2} \right) + ic_2 \sinh \left(\frac{\ln(t)}{2} \right) \right)$$

Summary

The solution(s) found are the following

$$y = t^{\frac{5}{2}} \left(c_1 \cosh \left(\frac{\ln(t)}{2} \right) + ic_2 \sinh \left(\frac{\ln(t)}{2} \right) \right) \quad (1)$$

Verification of solutions

$$y = t^{\frac{5}{2}} \left(c_1 \cosh \left(\frac{\ln(t)}{2} \right) + ic_2 \sinh \left(\frac{\ln(t)}{2} \right) \right)$$

Verified OK.

8.25.5 Solving as second order change of variable on y method 1 ode

In normal form the given ode is written as

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$\begin{aligned}p(t) &= -\frac{4}{t} \\ q(t) &= \frac{6}{t^2}\end{aligned}$$

Calculating the Liouville ode invariant Q given by

$$\begin{aligned}
 Q &= q - \frac{p'}{2} - \frac{p^2}{4} \\
 &= \frac{6}{t^2} - \frac{\left(-\frac{4}{t}\right)'}{2} - \frac{\left(-\frac{4}{t}\right)^2}{4} \\
 &= \frac{6}{t^2} - \frac{\left(\frac{4}{t^2}\right)}{2} - \frac{\left(\frac{16}{t^2}\right)}{4} \\
 &= \frac{6}{t^2} - \left(\frac{2}{t^2}\right) - \frac{4}{t^2} \\
 &= 0
 \end{aligned}$$

Since the Liouville ode invariant does not depend on the independent variable t then the transformation

$$y = v(t) z(t) \quad (3)$$

is used to change the original ode to a constant coefficients ode in v . In (3) the term $z(t)$ is given by

$$\begin{aligned}
 z(t) &= e^{-\left(\int \frac{p(t)}{2} dt\right)} \\
 &= e^{-\int \frac{-4}{2} dt} \\
 &= t^2
 \end{aligned} \quad (5)$$

Hence (3) becomes

$$y = v(t) t^2 \quad (4)$$

Applying this change of variable to the original ode results in

$$t^4 v''(t) = 0$$

Which is now solved for $v(t)$ Integrating twice gives the solution

$$v(t) = c_1 t + c_2$$

Now that $v(t)$ is known, then

$$\begin{aligned}
 y &= v(t) z(t) \\
 &= (c_1 t + c_2) (z(t))
 \end{aligned} \quad (7)$$

But from (5)

$$z(t) = t^2$$

Hence (7) becomes

$$y = (c_1 t + c_2) t^2$$

Summary

The solution(s) found are the following

$$y = (c_1 t + c_2) t^2 \quad (1)$$

Verification of solutions

$$y = (c_1 t + c_2) t^2$$

Verified OK.

8.25.6 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$t^2 y'' - 4ty' + 6y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = -\frac{4}{t}$$
$$q(t) = \frac{6}{t^2}$$

Applying change of variables on the dependent variable $y = v(t) t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p \right) v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q \right) v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} - \frac{4n}{t^2} + \frac{6}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = 3 \tag{6}$$

Substituting this value in (3) gives

$$\begin{aligned} v''(t) + \frac{2v'(t)}{t} &= 0 \\ v''(t) + \frac{2v'(t)}{t} &= 0 \end{aligned} \tag{7}$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{2u(t)}{t} = 0 \tag{8}$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{2u}{t} \end{aligned}$$

Where $f(t) = -\frac{2}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= -\frac{2}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{2}{t} dt \\ \ln(u) &= -2 \ln(t) + c_1 \\ u &= e^{-2 \ln(t) + c_1} \\ &= \frac{c_1}{t^2} \end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned} v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= -\frac{c_1}{t} + c_2 \end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\ &= \left(-\frac{c_1}{t} + c_2\right) t^3 \\ &= (c_2 t - c_1) t^2\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \left(-\frac{c_1}{t} + c_2\right) t^3 \quad (1)$$

Verification of solutions

$$y = \left(-\frac{c_1}{t} + c_2\right) t^3$$

Verified OK.

8.25.7 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' - 4ty' + 6y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= t^2 \\ B &= -4t \\ C &= 6\end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned}r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}\end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 0 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = 0 \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 412: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-4t}{t^2} dt} \\ &= z_1 e^{2 \ln(t)} \\ &= z_1 (t^2) \end{aligned}$$

Which simplifies to

$$y_1 = t^2$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-4t}{t^2} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{4 \ln(t)}}{(y_1)^2} dt \\ &= y_1(t) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (t^2) + c_2 (t^2(t)) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_2 t^3 + c_1 t^2 \quad (1)$$

Verification of solutions

$$y = c_2 t^3 + c_1 t^2$$

Verified OK.

8.25.8 Maple step by step solution

Let's solve

$$y'' t^2 - 4t y' + 6y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{4y'}{t} - \frac{6y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{4y'}{t} + \frac{6y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y'' t^2 - 4t y' + 6y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds} y(s) \right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds} y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2} y(s) \right) s'(t)^2 + s''(t) \left(\frac{d}{ds} y(s) \right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2} y(s)}{t^2} - \frac{\frac{d}{ds} y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2} \right) t^2 - 4 \frac{d}{ds}y(s) + 6y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) - 5 \frac{d}{ds}y(s) + 6y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 - 5r + 6 = 0$$

- Factor the characteristic polynomial

$$(r - 2)(r - 3) = 0$$

- Roots of the characteristic polynomial

$$r = (2, 3)$$

- 1st solution of the ODE

$$y_1(s) = e^{2s}$$

- 2nd solution of the ODE

$$y_2(s) = e^{3s}$$

- General solution of the ODE

$$y(s) = c_1y_1(s) + c_2y_2(s)$$

- Substitute in solutions

$$y(s) = c_1e^{2s} + c_2e^{3s}$$

- Change variables back using $s = \ln(t)$

$$y = c_2t^3 + c_1t^2$$

- Simplify

$$y = t^2(c_2t + c_1)$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
<- LODE of Euler type successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 13

```
dsolve(t^2*diff(y(t),t$2)-4*t*diff(y(t),t)+6*y(t) = 0,y(t), singsol=all)
```

$$y(t) = t^2(c_2t + c_1)$$

✓ Solution by Mathematica

Time used: 0.01 (sec). Leaf size: 16

```
DSolve[t^2*y'[t]-4*t*y'[t]+6*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t^2(c_2t + c_1)$$

8.26 problem 40

8.26.1 Solving as second order euler ode ode	2293
8.26.2 Solving as second order change of variable on x method 2 ode .	2295
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Internal problem ID [648]

Internal file name [OUTPUT/648_Sunday_June_05_2022_01_46_23_AM_19111792/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 40.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_Emden, _Fowler]]
```

$$t^2 y'' - t y' + 5y = 0$$

8.26.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = r t^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} - t r t^{r-1} + 5t^r = 0$$

Simplifying gives

$$r(r-1)t^r - r t^r + 5t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r - 1) - r + 5 = 0$$

Or

$$r^2 - 2r + 5 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = 1 - 2i$$

$$r_2 = 1 + 2i$$

The roots are complex conjugate of each others. Let the roots be

$$r_1 = \alpha + i\beta$$

$$r_2 = \alpha - i\beta$$

Where in this case $\alpha = 1$ and $\beta = -2$. Hence the solution becomes

$$\begin{aligned} y &= c_1 t^{r_1} + c_2 t^{r_2} \\ &= c_1 t^{\alpha+i\beta} + c_2 t^{\alpha-i\beta} \\ &= t^\alpha (c_1 t^{i\beta} + c_2 t^{-i\beta}) \\ &= t^\alpha (c_1 e^{\ln(t^{i\beta})} + c_2 e^{\ln(t^{-i\beta})}) \\ &= t^\alpha (c_1 e^{i(\beta \ln t)} + c_2 e^{-i(\beta \ln t)}) \end{aligned}$$

Using the values for $\alpha = 1, \beta = -2$, the above becomes

$$y = t^1 (c_1 e^{-2i \ln(t)} + c_2 e^{2i \ln(t)})$$

Using Euler relation, the expression $c_1 e^{iA} + c_2 e^{-iA}$ is transformed to $c_1 \cos A + c_1 \sin A$ where the constants are free to change. Applying this to the above result gives

$$y = t(c_1 \cos(2 \ln(t)) + c_2 \sin(2 \ln(t)))$$

Summary

The solution(s) found are the following

$$y = t(c_1 \cos(2 \ln(t)) + c_2 \sin(2 \ln(t))) \tag{1}$$

Verification of solutions

$$y = t(c_1 \cos(2 \ln(t)) + c_2 \sin(2 \ln(t)))$$

Verified OK.

8.26.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$t^2 y'' - ty' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = -\frac{1}{t}$$
$$q(t) = \frac{5}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t) \tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t) dt)} dt \\ &= \int e^{-(\int -\frac{1}{t} dt)} dt \\ &= \int e^{\ln(t)} dt \\ &= \int t dt \\ &= \frac{t^2}{2} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned}q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{\frac{5}{t^2}}{t^2} \\ &= \frac{5}{t^4}\end{aligned}\tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned}\frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + \frac{5y(\tau)}{t^4} &= 0\end{aligned}$$

But in terms of τ

$$\frac{5}{t^4} = \frac{5}{4\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{5y(\tau)}{4\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$4\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + 5y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$4\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + 5\tau^r = 0$$

Simplifying gives

$$4r(r-1)\tau^r + 0\tau^r + 5\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$4r(r-1) + 0 + 5 = 0$$

Or

$$4r^2 - 4r + 5 = 0\tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2} - i$$

$$r_2 = \frac{1}{2} + i$$

The roots are complex conjugate of each others. Let the roots be

$$r_1 = \alpha + i\beta$$

$$r_2 = \alpha - i\beta$$

Where in this case $\alpha = \frac{1}{2}$ and $\beta = -1$. Hence the solution becomes

$$y(\tau) = c_1\tau^{r_1} + c_2\tau^{r_2}$$

$$= c_1\tau^{\alpha+i\beta} + c_2\tau^{\alpha-i\beta}$$

$$= \tau^\alpha (c_1\tau^{i\beta} + c_2\tau^{-i\beta})$$

$$= \tau^\alpha (c_1e^{\ln(\tau^{i\beta})} + c_2e^{\ln(\tau^{-i\beta})})$$

$$= \tau^\alpha (c_1e^{i(\beta \ln \tau)} + c_2e^{-i(\beta \ln \tau)})$$

Using the values for $\alpha = \frac{1}{2}, \beta = -1$, the above becomes

$$y(\tau) = \tau^{\frac{1}{2}} (c_1e^{-i \ln(\tau)} + c_2e^{i \ln(\tau)})$$

Using Euler relation, the expression $c_1e^{iA} + c_2e^{-iA}$ is transformed to $c_1 \cos A + c_1 \sin A$ where the constants are free to change. Applying this to the above result gives

$$y(\tau) = \sqrt{\tau}(c_1 \cos (\ln (\tau)) + c_2 \sin (\ln (\tau)))$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{\sqrt{2}t(c_1 \cos (-\ln (2) + 2 \ln (t)) + c_2 \sin (-\ln (2) + 2 \ln (t)))}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{2}t(c_1 \cos (-\ln (2) + 2 \ln (t)) + c_2 \sin (-\ln (2) + 2 \ln (t)))}{2} \quad (1)$$

Verification of solutions

$$y = \frac{\sqrt{2}t(c_1 \cos (-\ln (2) + 2 \ln (t)) + c_2 \sin (-\ln (2) + 2 \ln (t)))}{2}$$

Verified OK.

8.26.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$t^2 y'' - t y' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = -\frac{1}{t}$$
$$q(t) = \frac{5}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c} \sqrt{q}$$
$$= \frac{\sqrt{5} \sqrt{\frac{1}{t^2}}}{c} \quad (6)$$
$$\tau'' = -\frac{\sqrt{5}}{c \sqrt{\frac{1}{t^2}} t^3}$$

Substituting the above into (4) results in

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2}$$
$$= \frac{-\frac{\sqrt{5}}{c \sqrt{\frac{1}{t^2}} t^3} - \frac{1}{t} \frac{\sqrt{5} \sqrt{\frac{1}{t^2}}}{c}}{\left(\frac{\sqrt{5} \sqrt{\frac{1}{t^2}}}{c} \right)^2}$$
$$= -\frac{2c\sqrt{5}}{5}$$

Therefore ode (3) now becomes

$$\begin{aligned}
 y(\tau)'' + p_1 y(\tau)' + q_1 y(\tau) &= 0 \\
 \frac{d^2}{d\tau^2} y(\tau) - \frac{2c\sqrt{5}}{5} \left(\frac{d}{d\tau} y(\tau) \right) + c^2 y(\tau) &= 0
 \end{aligned} \tag{7}$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{\frac{\sqrt{5}c\tau}{5}} \left(c_1 \cos \left(\frac{2\sqrt{5}c\tau}{5} \right) + c_2 \sin \left(\frac{2\sqrt{5}c\tau}{5} \right) \right)$$

Now from (6)

$$\begin{aligned}
 \tau &= \int \frac{1}{c} \sqrt{q} dt \\
 &= \frac{\int \sqrt{5} \sqrt{\frac{1}{t^2}} dt}{c} \\
 &= \frac{\sqrt{5} \sqrt{\frac{1}{t^2}} t \ln(t)}{c}
 \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = t(c_1 \cos(2 \ln(t)) + c_2 \sin(2 \ln(t)))$$

Summary

The solution(s) found are the following

$$y = t(c_1 \cos(2 \ln(t)) + c_2 \sin(2 \ln(t))) \tag{1}$$

Verification of solutions

$$y = t(c_1 \cos(2 \ln(t)) + c_2 \sin(2 \ln(t)))$$

Verified OK.

8.26.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$t^2 y'' - t y' + 5y = 0 \tag{1}$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \tag{2}$$

Where

$$p(t) = -\frac{1}{t}$$
$$q(t) = \frac{5}{t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} - \frac{n}{t^2} + \frac{5}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = 1 + 2i \quad (6)$$

Substituting this value in (3) gives

$$v''(t) + \left(\frac{2+4i}{t} - \frac{1}{t}\right)v'(t) = 0$$
$$v''(t) + \frac{(1+4i)v'(t)}{t} = 0 \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{(1+4i)u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned}u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{(-1 - 4i)u}{t}\end{aligned}$$

Where $f(t) = \frac{-1-4i}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= \frac{-1 - 4i}{t} dt \\ \int \frac{1}{u} du &= \int \frac{-1 - 4i}{t} dt \\ \ln(u) &= (-1 - 4i) \ln(t) + c_1 \\ u &= e^{(-1-4i)\ln(t)+c_1} \\ &= c_1 e^{(-1-4i)\ln(t)}\end{aligned}$$

Which simplifies to

$$u(t) = \frac{c_1 t^{-4i}}{t}$$

Now that $u(t)$ is known, then

$$\begin{aligned}v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= \frac{ic_1 t^{-4i}}{4} + c_2\end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\ &= \left(\frac{ic_1 t^{-4i}}{4} + c_2 \right) t^{1+2i} \\ &= t^{1+2i} c_2 + \frac{it^{1-2i} c_1}{4}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \left(\frac{ic_1 t^{-4i}}{4} + c_2 \right) t^{1+2i} \quad (1)$$

Verification of solutions

$$y = \left(\frac{ic_1 t^{-4i}}{4} + c_2 \right) t^{1+2i}$$

Verified OK.

8.26.5 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' - t y' + 5y = 0 \quad (1)$$

$$A y'' + B y' + C y = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = t^2$$

$$B = -t \quad (3)$$

$$C = 5$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = y e^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = r z(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-17}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$s = -17$$

$$t = 4t^2$$

Therefore eq. (4) becomes

$$z''(t) = \left(-\frac{17}{4t^2}\right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 414: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{17}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{17}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} + 2i \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} - 2i \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{17}{4t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{17}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} + 2i \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} - 2i \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{17}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2} + 2i$	$\frac{1}{2} - 2i$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2} + 2i$	$\frac{1}{2} - 2i$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to

determine possible non negative integer d from these using

$$d = \alpha_{\infty}^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_{∞}^{\pm} . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_{\infty}^{-} = \frac{1}{2} - 2i$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= \frac{1}{2} - 2i - \left(\frac{1}{2} - 2i \right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{t - c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= \frac{\frac{1}{2} - 2i}{t} + (-)(0) \\ &= \frac{\frac{1}{2} - 2i}{t} \\ &= \frac{\frac{1}{2} - 2i}{t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \tag{1A}$$

Let

$$p(t) = 1 \tag{2A}$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2 \left(\frac{\frac{1}{2} - 2i}{t} \right) (0) + \left(\left(\frac{-\frac{1}{2} + 2i}{t^2} \right) + \left(\frac{\frac{1}{2} - 2i}{t} \right)^2 - \left(-\frac{17}{4t^2} \right) \right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \frac{\frac{1}{2}-2i}{t} dt} \\ &= t^{\frac{1}{2}-2i} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-t}{t^2} dt} \\ &= z_1 e^{\frac{\ln(t)}{2}} \\ &= z_1 (\sqrt{t}) \end{aligned}$$

Which simplifies to

$$y_1 = t^{1-2i}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-t}{t^2} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{\ln(t)}}{(y_1)^2} dt \\ &= y_1 \left(-\frac{it^{4i}}{4} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (t^{1-2i}) + c_2 \left(t^{1-2i} \left(-\frac{it^{4i}}{4} \right) \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = t^{1-2i} c_1 - \frac{ic_2 t^{1+2i}}{4} \quad (1)$$

Verification of solutions

$$y = t^{1-2i} c_1 - \frac{ic_2 t^{1+2i}}{4}$$

Verified OK.

8.26.6 Maple step by step solution

Let's solve

$$y''t^2 - ty' + 5y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{y'}{t} - \frac{5y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{y'}{t} + \frac{5y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 - ty' + 5y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2} \right) t^2 - \frac{d}{ds}y(s) + 5y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) - 2\frac{d}{ds}y(s) + 5y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 - 2r + 5 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{2 \pm (\sqrt{-16})}{2}$$

- Roots of the characteristic polynomial

$$r = (1 - 2I, 1 + 2I)$$

- 1st solution of the ODE

$$y_1(s) = e^s \cos(2s)$$

- 2nd solution of the ODE

$$y_2(s) = e^s \sin(2s)$$

- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$

- Substitute in solutions

$$y(s) = c_1 e^s \cos(2s) + c_2 e^s \sin(2s)$$

- Change variables back using $s = \ln(t)$

$$y = c_1 t \cos(2 \ln(t)) + c_2 t \sin(2 \ln(t))$$

- Simplify

$$y = t(c_1 \cos(2 \ln(t)) + c_2 \sin(2 \ln(t)))$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
<- LODE of Euler type successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 21

```
dsolve(t^2*diff(y(t),t$2)- t*diff(y(t),t)+5*y(t) = 0,y(t), singsol=all)
```

$$y(t) = t(c_1 \sin(2 \ln(t)) + c_2 \cos(2 \ln(t)))$$

✓ Solution by Mathematica

Time used: 0.023 (sec). Leaf size: 24

```
DSolve[t^2*y''[t]-t*y'[t]+5*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t(c_2 \cos(2 \log(t)) + c_1 \sin(2 \log(t)))$$

8.27 problem 41

8.27.1 Solving as second order euler ode ode	2310
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Internal problem ID [649]

Internal file name [OUTPUT/649_Sunday_June_05_2022_01_46_24_AM_12275308/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 41.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2", "second_order_ode_non_constant_coeff_transformation_on_B"

Maple gives the following as the ode type

```
[[_Emden , _Fowler]]
```

$$t^2y'' + 3ty' - 3y = 0$$

8.27.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} + 3trt^{r-1} - 3t^r = 0$$

Simplifying gives

$$r(r-1)t^r + 3rt^r - 3t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r - 1) + 3r - 3 = 0$$

Or

$$r^2 + 2r - 3 = 0 \quad (1)$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -3$$

$$r_2 = 1$$

Since the roots are real and distinct, then the general solution is

$$y = c_1 y_1 + c_2 y_2$$

Where $y_1 = t^{-3}$ and $y_2 = t^1$. Hence

$$y = \frac{c_1}{t^3} + c_2 t$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t^3} + c_2 t \quad (1)$$

Verification of solutions

$$y = \frac{c_1}{t^3} + c_2 t$$

Verified OK.

8.27.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$t^2 y'' + 3ty' - 3y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{3}{t}$$
$$q(t) = -\frac{3}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t)\tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t)dt)} dt \\ &= \int e^{-(\int \frac{3}{t} dt)} dt \\ &= \int e^{-3\ln(t)} dt \\ &= \int \frac{1}{t^3} dt \\ &= -\frac{1}{2t^2} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{-\frac{3}{t^2}}{\frac{1}{t^6}} \\ &= -3t^4 \end{aligned} \quad (7)$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) - 3t^4y(\tau) &= 0 \end{aligned}$$

But in terms of τ

$$-3t^4 = -\frac{3}{4\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) - \frac{3y(\tau)}{4\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$4\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 - 3y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$4\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} - 3\tau^r = 0$$

Simplifying gives

$$4r(r-1)\tau^r + 0\tau^r - 3\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$4r(r-1) + 0 - 3 = 0$$

Or

$$4r^2 - 4r - 3 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -\frac{1}{2}$$
$$r_2 = \frac{3}{2}$$

Since the roots are real and distinct, then the general solution is

$$y(\tau) = c_1y_1 + c_2y_2$$

Where $y_1 = \tau^{r_1}$ and $y_2 = \tau^{r_2}$. Hence

$$y(\tau) = \frac{c_1}{\sqrt{\tau}} + c_2\tau^{\frac{3}{2}}$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{(4c_1t^4 + c_2)\sqrt{2}}{4t^4\sqrt{-\frac{1}{t^2}}}$$

Summary

The solution(s) found are the following

$$y = \frac{(4c_1t^4 + c_2)\sqrt{2}}{4t^4\sqrt{-\frac{1}{t^2}}} \quad (1)$$

Verification of solutions

$$y = \frac{(4c_1t^4 + c_2)\sqrt{2}}{4t^4\sqrt{-\frac{1}{t^2}}}$$

Verified OK.

8.27.3 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$t^2y'' + 3ty' - 3y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{3}{t}$$
$$q(t) = -\frac{3}{t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} + \frac{3n}{t^2} - \frac{3}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = 1 \tag{6}$$

Substituting this value in (3) gives

$$\begin{aligned} v''(t) + \frac{5v'(t)}{t} &= 0 \\ v''(t) + \frac{5v'(t)}{t} &= 0 \end{aligned} \tag{7}$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{5u(t)}{t} = 0 \tag{8}$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{5u}{t} \end{aligned}$$

Where $f(t) = -\frac{5}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= -\frac{5}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{5}{t} dt \\ \ln(u) &= -5 \ln(t) + c_1 \\ u &= e^{-5 \ln(t) + c_1} \\ &= \frac{c_1}{t^5} \end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned} v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= -\frac{c_1}{4t^4} + c_2 \end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\ &= \left(-\frac{c_1}{4t^4} + c_2\right) t \\ &= \left(-\frac{c_1}{4t^4} + c_2\right) t\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \left(-\frac{c_1}{4t^4} + c_2\right) t \quad (1)$$

Verification of solutions

$$y = \left(-\frac{c_1}{4t^4} + c_2\right) t$$

Verified OK.

8.27.4 Solving as second order ode non constant coeff transformation on B ode

Given an ode of the form

$$Ay'' + By' + Cy = F(t)$$

This method reduces the order ode the ODE by one by applying the transformation

$$y = Bv$$

This results in

$$\begin{aligned}y' &= B'v + v'B \\ y'' &= B''v + B'v' + v''B + v'B' \\ &= v''B + 2v' + B' + B''v\end{aligned}$$

And now the original ode becomes

$$\begin{aligned}A(v''B + 2v'B' + B''v) + B(B'v + v'B) + CBv &= 0 \\ ABv'' + (2AB' + B^2)v' + (AB'' + BB' + CB)v &= 0\end{aligned} \quad (1)$$

If the term $AB'' + BB' + CB$ is zero, then this method works and can be used to solve

$$ABv'' + (2AB' + B^2)v' = 0$$

By Using $u = v'$ which reduces the order of the above ode to one. The new ode is

$$ABu' + (2AB' + B^2) u = 0$$

The above ode is first order ode which is solved for u . Now a new ode $v' = u$ is solved for v as first order ode. Then the final solution is obtain from $y = Bv$.

This method works only if the term $AB'' + BB' + CB$ is zero. The given ODE shows that

$$A = t^2$$

$$B = 3t$$

$$C = -3$$

$$F = 0$$

The above shows that for this ode

$$\begin{aligned} AB'' + BB' + CB &= (t^2) (0) + (3t) (3) + (-3) (3t) \\ &= 0 \end{aligned}$$

Hence the ode in v given in (1) now simplifies to

$$3t^3 v'' + (15t^2) v' = 0$$

Now by applying $v' = u$ the above becomes

$$3t^2(u'(t) t + 5u(t)) = 0$$

Which is now solved for u . In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{5u}{t} \end{aligned}$$

Where $f(t) = -\frac{5}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= -\frac{5}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{5}{t} dt \\ \ln(u) &= -5 \ln(t) + c_1 \\ u &= e^{-5 \ln(t) + c_1} \\ &= \frac{c_1}{t^5} \end{aligned}$$

The ode for v now becomes

$$\begin{aligned}v' &= u \\ &= \frac{c_1}{t^5}\end{aligned}$$

Which is now solved for v . Integrating both sides gives

$$\begin{aligned}v(t) &= \int \frac{c_1}{t^5} dt \\ &= -\frac{c_1}{4t^4} + c_2\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y(t) &= Bv \\ &= (3t) \left(-\frac{c_1}{4t^4} + c_2 \right) \\ &= \frac{3c_2t^4 - \frac{3c_1}{4}}{t^3}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{3c_2t^4 - \frac{3c_1}{4}}{t^3} \quad (1)$$

Verification of solutions

$$y = \frac{3c_2t^4 - \frac{3c_1}{4}}{t^3}$$

Verified OK.

8.27.5 Solving using Kovacic algorithm

Writing the ode as

$$t^2y'' + 3ty' - 3y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= t^2 \\ B &= 3t \\ C &= -3\end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{15}{4t^2} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 15 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(\frac{15}{4t^2} \right) z(t) \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 416: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 2 - 0 \\
 &= 2
 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{15}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition

of r given above. Therefore $b = \frac{15}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{5}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{3}{2} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = \frac{15}{4t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = \frac{15}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{5}{2} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{3}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = \frac{15}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{5}{2}$	$-\frac{3}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{5}{2}$	$-\frac{3}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = -\frac{3}{2}$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= -\frac{3}{2} - \left(-\frac{3}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t-c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{t-c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= -\frac{3}{2t} + (-)(0) \\ &= -\frac{3}{2t} \\ &= -\frac{3}{2t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r)p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2\left(-\frac{3}{2t}\right)(0) + \left(\left(\frac{3}{2t^2}\right) + \left(-\frac{3}{2t}\right)^2 - \left(\frac{15}{4t^2}\right) \right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int -\frac{3}{2t} dt} \\ &= \frac{1}{t^{\frac{3}{2}}} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\&= z_1 e^{-\int \frac{1}{2} \frac{3t}{t^2} dt} \\&= z_1 e^{-\frac{3 \ln(t)}{2}} \\&= z_1 \left(\frac{1}{t^{\frac{3}{2}}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = \frac{1}{t^3}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{3t}{t^2} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{-3 \ln(t)}}{(y_1)^2} dt \\&= y_1 \left(\frac{t^4}{4} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 \left(\frac{1}{t^3} \right) + c_2 \left(\frac{1}{t^3} \left(\frac{t^4}{4} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t^3} + \frac{c_2 t}{4} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{t^3} + \frac{c_2 t}{4}$$

Verified OK.

8.27.6 Maple step by step solution

Let's solve

$$y''t^2 + 3ty' - 3y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{3y'}{t} + \frac{3y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{3y'}{t} - \frac{3y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 + 3ty' - 3y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right) t^2 + 3\frac{d}{ds}y(s) - 3y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) + 2\frac{d}{ds}y(s) - 3y(s) = 0$$
- Characteristic polynomial of ODE

$$r^2 + 2r - 3 = 0$$
- Factor the characteristic polynomial

$$(r + 3)(r - 1) = 0$$
- Roots of the characteristic polynomial

$$r = (-3, 1)$$
- 1st solution of the ODE

$$y_1(s) = e^{-3s}$$
- 2nd solution of the ODE

$$y_2(s) = e^s$$
- General solution of the ODE

$$y(s) = c_1y_1(s) + c_2y_2(s)$$
- Substitute in solutions

$$y(s) = c_1e^{-3s} + c_2e^s$$
- Change variables back using $s = \ln(t)$

$$y = \frac{c_1}{t^3} + c_2t$$
- Simplify

$$y = \frac{c_1}{t^3} + c_2t$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve(t^2*diff(y(t),t$2)+ 3*t*diff(y(t),t)-3*y(t) = 0,y(t), singsol=all)
```

$$y(t) = \frac{c_1 t^4 + c_2}{t^3}$$

✓ Solution by Mathematica

Time used: 0.011 (sec). Leaf size: 16

```
DSolve[t^2*y''[t]+3*t*y'[t]-3*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{c_1}{t^3} + c_2 t$$

8.28 problem 42

8.28.1 Solving as second order euler ode ode	2327
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Internal problem ID [650]

Internal file name [OUTPUT/650_Sunday_June_05_2022_01_46_25_AM_47152662/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 42.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_Emden, _Fowler]]
```

$$t^2 y'' + 7ty' + 10y = 0$$

8.28.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} + 7trt^{r-1} + 10t^r = 0$$

Simplifying gives

$$r(r-1)t^r + 7rt^r + 10t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r - 1) + 7r + 10 = 0$$

Or

$$r^2 + 6r + 10 = 0 \quad (1)$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -3 - i$$

$$r_2 = -3 + i$$

The roots are complex conjugate of each others. Let the roots be

$$r_1 = \alpha + i\beta$$

$$r_2 = \alpha - i\beta$$

Where in this case $\alpha = -3$ and $\beta = -1$. Hence the solution becomes

$$\begin{aligned} y &= c_1 t^{r_1} + c_2 t^{r_2} \\ &= c_1 t^{\alpha+i\beta} + c_2 t^{\alpha-i\beta} \\ &= t^\alpha (c_1 t^{i\beta} + c_2 t^{-i\beta}) \\ &= t^\alpha (c_1 e^{\ln(t^{i\beta})} + c_2 e^{\ln(t^{-i\beta})}) \\ &= t^\alpha (c_1 e^{i(\beta \ln t)} + c_2 e^{-i(\beta \ln t)}) \end{aligned}$$

Using the values for $\alpha = -3, \beta = -1$, the above becomes

$$y = t^{-3} (c_1 e^{-i \ln(t)} + c_2 e^{i \ln(t)})$$

Using Euler relation, the expression $c_1 e^{iA} + c_2 e^{-iA}$ is transformed to $c_1 \cos A + c_1 \sin A$ where the constants are free to change. Applying this to the above result gives

$$y = \frac{1}{t^3} (c_1 \cos(\ln(t)) + c_2 \sin(\ln(t)))$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \cos(\ln(t)) + c_2 \sin(\ln(t))}{t^3} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 \cos(\ln(t)) + c_2 \sin(\ln(t))}{t^3}$$

Verified OK.

8.28.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$t^2 y'' + 7ty' + 10y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{7}{t}$$
$$q(t) = \frac{10}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t)\tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t)dt)} dt \\ &= \int e^{-(\int \frac{7}{t} dt)} dt \\ &= \int e^{-7\ln(t)} dt \\ &= \int \frac{1}{t^7} dt \\ &= -\frac{1}{6t^6} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned}q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{\frac{10}{t^2}}{\frac{1}{t^{14}}} \\ &= 10t^{12}\end{aligned}\tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned}\frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + 10t^{12}y(\tau) &= 0\end{aligned}$$

But in terms of τ

$$10t^{12} = \frac{5}{18\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{5y(\tau)}{18\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$18\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + 5y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$18\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + 5\tau^r = 0$$

Simplifying gives

$$18r(r-1)\tau^r + 0\tau^r + 5\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$18r(r-1) + 0 + 5 = 0$$

Or

$$18r^2 - 18r + 5 = 0\tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2} - \frac{i}{6}$$

$$r_2 = \frac{1}{2} + \frac{i}{6}$$

The roots are complex conjugate of each others. Let the roots be

$$r_1 = \alpha + i\beta$$

$$r_2 = \alpha - i\beta$$

Where in this case $\alpha = \frac{1}{2}$ and $\beta = -\frac{1}{6}$. Hence the solution becomes

$$y(\tau) = c_1\tau^{r_1} + c_2\tau^{r_2}$$

$$= c_1\tau^{\alpha+i\beta} + c_2\tau^{\alpha-i\beta}$$

$$= \tau^\alpha(c_1\tau^{i\beta} + c_2\tau^{-i\beta})$$

$$= \tau^\alpha(c_1e^{\ln(\tau^{i\beta})} + c_2e^{\ln(\tau^{-i\beta})})$$

$$= \tau^\alpha(c_1e^{i(\beta \ln \tau)} + c_2e^{-i(\beta \ln \tau)})$$

Using the values for $\alpha = \frac{1}{2}, \beta = -\frac{1}{6}$, the above becomes

$$y(\tau) = \tau^{\frac{1}{2}} \left(c_1 e^{-\frac{i \ln(\tau)}{6}} + c_2 e^{\frac{i \ln(\tau)}{6}} \right)$$

Using Euler relation, the expression $c_1 e^{iA} + c_2 e^{-iA}$ is transformed to $c_1 \cos A + c_1 \sin A$ where the constants are free to change. Applying this to the above result gives

$$y(\tau) = \sqrt{\tau} \left(c_1 \cos \left(\frac{\ln(\tau)}{6} \right) + c_2 \sin \left(\frac{\ln(\tau)}{6} \right) \right)$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{\left(c_1 \cos \left(-\frac{\ln(2)}{6} - \frac{\ln(3)}{6} + \frac{\ln\left(-\frac{1}{t^6}\right)}{6} \right) + c_2 \sin \left(-\frac{\ln(2)}{6} - \frac{\ln(3)}{6} + \frac{\ln\left(-\frac{1}{t^6}\right)}{6} \right) \right) \sqrt{6} \sqrt{-\frac{1}{t^6}}}{6}$$

Summary

The solution(s) found are the following

$$y = \frac{\left(c_1 \cos \left(-\frac{\ln(2)}{6} - \frac{\ln(3)}{6} + \frac{\ln\left(-\frac{1}{t^6}\right)}{6} \right) + c_2 \sin \left(-\frac{\ln(2)}{6} - \frac{\ln(3)}{6} + \frac{\ln\left(-\frac{1}{t^6}\right)}{6} \right) \right) \sqrt{6} \sqrt{-\frac{1}{t^6}}}{6} \quad (1)$$

Verification of solutions

$$y = \frac{\left(c_1 \cos \left(-\frac{\ln(2)}{6} - \frac{\ln(3)}{6} + \frac{\ln\left(-\frac{1}{t^6}\right)}{6} \right) + c_2 \sin \left(-\frac{\ln(2)}{6} - \frac{\ln(3)}{6} + \frac{\ln\left(-\frac{1}{t^6}\right)}{6} \right) \right) \sqrt{6} \sqrt{-\frac{1}{t^6}}}{6}$$

Verified OK.

8.28.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$t^2 y'' + 7ty' + 10y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{7}{t}$$
$$q(t) = \frac{10}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1 \left(\frac{d}{d\tau}y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c} \sqrt{q}$$
$$= \frac{\sqrt{10} \sqrt{\frac{1}{t^2}}}{c} \quad (6)$$
$$\tau'' = -\frac{\sqrt{10}}{c \sqrt{\frac{1}{t^2}} t^3}$$

Substituting the above into (4) results in

$$\begin{aligned}
 p_1(\tau) &= \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \\
 &= \frac{-\frac{\sqrt{10}}{c\sqrt{\frac{1}{t^2}} t^3} + \frac{7}{t} \frac{\sqrt{10}}{c} \sqrt{\frac{1}{t^2}}}{\left(\frac{\sqrt{10}}{c} \sqrt{\frac{1}{t^2}}\right)^2} \\
 &= \frac{3c\sqrt{10}}{5}
 \end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned}
 y(\tau)'' + p_1 y(\tau)' + q_1 y(\tau) &= 0 \\
 \frac{d^2}{d\tau^2} y(\tau) + \frac{3c\sqrt{10}}{5} \left(\frac{d}{d\tau} y(\tau)\right) + c^2 y(\tau) &= 0 \tag{7}
 \end{aligned}$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{-\frac{3\sqrt{10}c\tau}{10}} \left(c_1 \cos\left(\frac{\sqrt{10}c\tau}{10}\right) + c_2 \sin\left(\frac{\sqrt{10}c\tau}{10}\right) \right)$$

Now from (6)

$$\begin{aligned}
 \tau &= \int \frac{1}{c} \sqrt{q} dt \\
 &= \frac{\int \sqrt{10} \sqrt{\frac{1}{t^2}} dt}{c} \\
 &= \frac{\sqrt{10} \sqrt{\frac{1}{t^2}} t \ln(t)}{c}
 \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = \frac{c_1 \cos(\ln(t)) + c_2 \sin(\ln(t))}{t^3}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \cos(\ln(t)) + c_2 \sin(\ln(t))}{t^3} \tag{1}$$

Verification of solutions

$$y = \frac{c_1 \cos(\ln(t)) + c_2 \sin(\ln(t))}{t^3}$$

Verified OK.

8.28.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$t^2 y'' + 7ty' + 10y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{7}{t}$$
$$q(t) = \frac{10}{t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} + \frac{7n}{t} + \frac{10}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = -3 + i \quad (6)$$

Substituting this value in (3) gives

$$v''(t) + \left(\frac{-6 + 2i}{t} + \frac{7}{t}\right)v'(t) = 0$$
$$v''(t) + \frac{(1 + 2i)v'(t)}{t} = 0 \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{(1 + 2i)u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{(-1 - 2i)u}{t} \end{aligned}$$

Where $f(t) = \frac{-1-2i}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= \frac{-1 - 2i}{t} dt \\ \int \frac{1}{u} du &= \int \frac{-1 - 2i}{t} dt \\ \ln(u) &= (-1 - 2i) \ln(t) + c_1 \\ u &= e^{(-1-2i)\ln(t)+c_1} \\ &= c_1 e^{(-1-2i)\ln(t)} \end{aligned}$$

Which simplifies to

$$u(t) = \frac{c_1 t^{-2i}}{t}$$

Now that $u(t)$ is known, then

$$\begin{aligned} v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= \frac{ic_1 t^{-2i}}{2} + c_2 \end{aligned}$$

Hence

$$\begin{aligned} y &= v(t) t^n \\ &= \left(\frac{ic_1 t^{-2i}}{2} + c_2 \right) t^{-3+i} \\ &= c_2 t^{-3+i} + \frac{ic_1 t^{-3-i}}{2} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \left(\frac{ic_1 t^{-2i}}{2} + c_2 \right) t^{-3+i} \quad (1)$$

Verification of solutions

$$y = \left(\frac{ic_1 t^{-2i}}{2} + c_2 \right) t^{-3+i}$$

Verified OK.

8.28.5 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' + 7ty' + 10y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t^2 \\ B &= 7t \\ C &= 10 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-5}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -5 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(-\frac{5}{4t^2}\right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 418: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole

larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{5}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{5}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} + i \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} - i \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{5}{4t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{5}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} + i \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} - i \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{5}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2} + i$	$\frac{1}{2} - i$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2} + i$	$\frac{1}{2} - i$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_\infty^- = \frac{1}{2} - i$ then

$$\begin{aligned} d &= \alpha_\infty^- - (\alpha_{c_1}^-) \\ &= \frac{1}{2} - i - \left(\frac{1}{2} - i\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_\infty$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^-}{t - c_1} \right) + (-)[\sqrt{r}]_\infty \\ &= \frac{\frac{1}{2} - i}{t} + (-)(0) \\ &= \frac{\frac{1}{2} - i}{t} \\ &= \frac{\frac{1}{2} - i}{t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \tag{1A}$$

Let

$$p(t) = 1 \tag{2A}$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(\frac{\frac{1}{2} - i}{t}\right)(0) + \left(\left(\frac{-\frac{1}{2} + i}{t^2}\right) + \left(\frac{\frac{1}{2} - i}{t}\right)^2 - \left(-\frac{5}{4t^2}\right)\right) = 0$$

$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \frac{\frac{1}{2} - i}{t} dt} \\ &= t^{\frac{1}{2} - i} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{7t}{t^2} dt} \\ &= z_1 e^{-\frac{7 \ln(t)}{2}} \\ &= z_1 \left(\frac{1}{t^{\frac{7}{2}}}\right) \end{aligned}$$

Which simplifies to

$$y_1 = t^{-3-i}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{7t}{t^2} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{-7 \ln(t)}}{(y_1)^2} dt \\ &= y_1 \left(-\frac{it^{2i}}{2}\right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (t^{-3-i}) + c_2 \left(t^{-3-i} \left(-\frac{it^{2i}}{2} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 t^{-3-i} - \frac{ic_2 t^{-3+i}}{2} \quad (1)$$

Verification of solutions

$$y = c_1 t^{-3-i} - \frac{ic_2 t^{-3+i}}{2}$$

Verified OK.

8.28.6 Maple step by step solution

Let's solve

$$y''t^2 + 7ty' + 10y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{7y'}{t} - \frac{10y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{7y'}{t} + \frac{10y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 + 7ty' + 10y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds} y(s) \right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s) \right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s) \right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2} \right) t^2 + 7 \frac{d}{ds}y(s) + 10y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) + 6 \frac{d}{ds}y(s) + 10y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 + 6r + 10 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-6) \pm (\sqrt{-4})}{2}$$

- Roots of the characteristic polynomial

$$r = (-3 - I, -3 + I)$$

- 1st solution of the ODE

$$y_1(s) = e^{-3s} \cos(s)$$

- 2nd solution of the ODE

$$y_2(s) = e^{-3s} \sin(s)$$

- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$

- Substitute in solutions

$$y(s) = c_1 e^{-3s} \cos(s) + c_2 e^{-3s} \sin(s)$$

- Change variables back using $s = \ln(t)$

$$y = \frac{c_1 \cos(\ln(t))}{t^3} + \frac{c_2 \sin(\ln(t))}{t^3}$$

- Simplify

$$y = \frac{c_1 \cos(\ln(t))}{t^3} + \frac{c_2 \sin(\ln(t))}{t^3}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
<- LODE of Euler type successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(t^2*diff(y(t),t$2)+ 7*t*diff(y(t),t)+10*y(t) = 0,y(t), singsol=all)
```

$$y(t) = \frac{c_1 \sin(\ln(t)) + c_2 \cos(\ln(t))}{t^3}$$

✓ Solution by Mathematica

Time used: 0.022 (sec). Leaf size: 22

```
DSolve[t^2*y'[t]+7*t*y'[t]+10*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{c_2 \cos(\log(t)) + c_1 \sin(\log(t))}{t^3}$$

8.29 problem 44

8.29.1 Solving as second order change of variable on x method 2 ode . 2344

8.29.2 Solving as second order change of variable on x method 1 ode . 2347

Internal problem ID [651]

Internal file name [OUTPUT/651_Sunday_June_05_2022_01_46_25_AM_16115400/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 44.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2"

Maple gives the following as the ode type

```
[[_2nd_order, _with_linear_symmetries], [_2nd_order, _linear, `
  _with_symmetry_[0,F(x)]`]]
```

$$y'' + ty' + e^{-t^2}y = 0$$

8.29.1 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$y'' + ty' + e^{-t^2}y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = t$$
$$q(t) = e^{-t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t)\tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t)dt)} dt \\ &= \int e^{-(\int t dt)} dt \\ &= \int e^{-\frac{t^2}{2}} dt \\ &= \int e^{-\frac{t^2}{2}} dt \\ &= \frac{\sqrt{\pi} \sqrt{2} \operatorname{erf}\left(\frac{\sqrt{2}t}{2}\right)}{2} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{e^{-t^2}}{e^{-t^2}} \\ &= 1 \end{aligned} \quad (7)$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + y(\tau) &= 0 \end{aligned}$$

The above ode is now solved for $y(\tau)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(\tau) + By'(\tau) + Cy(\tau) = 0$$

Where in the above $A = 1, B = 0, C = 1$. Let the solution be $y(\tau) = e^{\lambda\tau}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda\tau} + e^{\lambda\tau} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda\tau}$ gives

$$\lambda^2 + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(1)} \\ &= \pm i \end{aligned}$$

Hence

$$\lambda_1 = +i$$

$$\lambda_2 = -i$$

Which simplifies to

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y(\tau) = e^{\alpha\tau} (c_1 \cos(\beta\tau) + c_2 \sin(\beta\tau))$$

Which becomes

$$y(\tau) = e^0 (c_1 \cos(\tau) + c_2 \sin(\tau))$$

Or

$$y(\tau) = c_1 \cos(\tau) + c_2 \sin(\tau)$$

The above solution is now transformed back to y using (6) which results in

$$y = c_1 \cos\left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf}\left(\frac{\sqrt{2}t}{2}\right)}{2}\right) + c_2 \sin\left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf}\left(\frac{\sqrt{2}t}{2}\right)}{2}\right)$$

Summary

The solution(s) found are the following

$$y = c_1 \cos\left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf}\left(\frac{\sqrt{2}t}{2}\right)}{2}\right) + c_2 \sin\left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf}\left(\frac{\sqrt{2}t}{2}\right)}{2}\right) \quad (1)$$

Verification of solutions

$$y = c_1 \cos\left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf}\left(\frac{\sqrt{2}t}{2}\right)}{2}\right) + c_2 \sin\left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf}\left(\frac{\sqrt{2}t}{2}\right)}{2}\right)$$

Verified OK.

8.29.2 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$y'' + ty' + e^{-t^2}y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$\begin{aligned} p(t) &= t \\ q(t) &= e^{-t^2} \end{aligned}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\begin{aligned} \tau' &= \frac{1}{c} \sqrt{q} \\ &= \frac{\sqrt{e^{-t^2}}}{c} \\ \tau'' &= -\frac{\sqrt{e^{-t^2}} t}{c} \end{aligned} \quad (6)$$

Substituting the above into (4) results in

$$\begin{aligned} p_1(\tau) &= \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \\ &= \frac{-\frac{\sqrt{e^{-t^2}} t}{c} + t \frac{\sqrt{e^{-t^2}}}{c}}{\left(\frac{\sqrt{e^{-t^2}}}{c}\right)^2} \\ &= 0 \end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned} y(\tau)'' + p_1 y(\tau)' + q_1 y(\tau) &= 0 \\ \frac{d^2}{d\tau^2} y(\tau) + c^2 y(\tau) &= 0 \end{aligned} \quad (7)$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = c_1 \cos(c\tau) + c_2 \sin(c\tau)$$

Now from (6)

$$\begin{aligned} \tau &= \int \frac{1}{c} \sqrt{q} dt \\ &= \frac{\int \sqrt{e^{-t^2}} dt}{c} \\ &= \frac{\sqrt{e^{-t^2}} e^{\frac{t^2}{2}} \sqrt{2} \sqrt{\pi} \operatorname{erf}\left(\frac{\sqrt{2}t}{2}\right)}{2c} \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = c_1 \cos \left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf} \left(\frac{\sqrt{2}t}{2} \right)}{2} \right) + c_2 \sin \left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf} \left(\frac{\sqrt{2}t}{2} \right)}{2} \right)$$

Summary

The solution(s) found are the following

$$y = c_1 \cos \left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf} \left(\frac{\sqrt{2}t}{2} \right)}{2} \right) + c_2 \sin \left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf} \left(\frac{\sqrt{2}t}{2} \right)}{2} \right) \quad (1)$$

Verification of solutions

$$y = c_1 \cos \left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf} \left(\frac{\sqrt{2}t}{2} \right)}{2} \right) + c_2 \sin \left(\frac{\sqrt{\pi} \sqrt{2} \operatorname{erf} \left(\frac{\sqrt{2}t}{2} \right)}{2} \right)$$

Verified OK.

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a symmetry of the form [xi=0, eta=F(x)]
<- linear_1 successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 55

```
dsolve(diff(y(t),t$2)+ t*diff(y(t),t)+exp(-t^2)*y(t) = 0,y(t), singsol=all)
```

$$y(t) = c_1 \operatorname{csgn} \left(e^{\frac{t^2}{2}} \right) \sin \left(\frac{\sqrt{2} \sqrt{\pi} \operatorname{erf} \left(\frac{t\sqrt{2}}{2} \right)}{2} \right) + c_2 \cos \left(\frac{\sqrt{2} \operatorname{csgn} \left(e^{\frac{t^2}{2}} \right) \sqrt{\pi} \operatorname{erf} \left(\frac{t\sqrt{2}}{2} \right)}{2} \right)$$

✓ Solution by Mathematica

Time used: 0.059 (sec). Leaf size: 102

```
DSolve[y''[t]+t*y'[t]+exp(-t^2)*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{-\frac{1}{4}(\sqrt{4\exp+1}+1)t^2} \left(c_1 \text{HermiteH} \left(-\frac{1}{2} - \frac{1}{2\sqrt{4\exp+1}}, \frac{\sqrt[4]{4\exp+1}t}{\sqrt{2}} \right) + c_2 \text{Hypergeometric1F1} \left(\frac{1}{4} \left(1 + \frac{1}{\sqrt{4\exp+1}} \right), \frac{1}{2}, \frac{1}{2} \sqrt{4\exp+1}t^2 \right) \right)$$

8.30 problem 46

- 8.30.1 Solving as second order change of variable on x method 2 ode . 2351
- 8.30.2 Solving as second order change of variable on x method 1 ode . 2355
- 8.30.3 Solving using Kovacic algorithm 2357
- 8.30.4 Maple step by step solution 2363

Internal problem ID [652]

Internal file name [OUTPUT/652_Sunday_June_05_2022_01_46_26_AM_50356826/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.3 Complex Roots of the Characteristic Equation , page 164

Problem number: 46.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$ty'' + (t^2 - 1)y' + yt^3 = 0$$

8.30.1 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$ty'' + (t^2 - 1)y' + yt^3 = 0 \tag{1}$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \tag{2}$$

Where

$$p(t) = \frac{t^2 - 1}{t}$$

$$q(t) = t^2$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t)\tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t)dt)} dt \\ &= \int e^{-\left(\int \frac{t^2-1}{t} dt\right)} dt \\ &= \int e^{-\frac{t^2}{2} + \ln(t)} dt \\ &= \int t e^{-\frac{t^2}{2}} dt \\ &= -e^{-\frac{t^2}{2}} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{t^2}{t^2 e^{-t^2}} \\ &= e^{t^2} \end{aligned} \quad (7)$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + e^{t^2}y(\tau) &= 0 \end{aligned}$$

But in terms of τ

$$e^{t^2} = \frac{1}{\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{y(\tau)}{\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + \tau^r = 0$$

Simplifying gives

$$r(r-1)\tau^r + 0\tau^r + \tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$r(r-1) + 0 + 1 = 0$$

Or

$$r^2 - r + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2} - \frac{i\sqrt{3}}{2}$$
$$r_2 = \frac{1}{2} + \frac{i\sqrt{3}}{2}$$

The roots are complex conjugate of each others. Let the roots be

$$r_1 = \alpha + i\beta$$
$$r_2 = \alpha - i\beta$$

Where in this case $\alpha = \frac{1}{2}$ and $\beta = -\frac{\sqrt{3}}{2}$. Hence the solution becomes

$$\begin{aligned} y(\tau) &= c_1\tau^{r_1} + c_2\tau^{r_2} \\ &= c_1\tau^{\alpha+i\beta} + c_2\tau^{\alpha-i\beta} \\ &= \tau^\alpha(c_1\tau^{i\beta} + c_2\tau^{-i\beta}) \\ &= \tau^\alpha(c_1e^{\ln(\tau^{i\beta})} + c_2e^{\ln(\tau^{-i\beta})}) \\ &= \tau^\alpha(c_1e^{i(\beta \ln \tau)} + c_2e^{-i(\beta \ln \tau)}) \end{aligned}$$

Using the values for $\alpha = \frac{1}{2}, \beta = -\frac{\sqrt{3}}{2}$, the above becomes

$$y(\tau) = \tau^{\frac{1}{2}} \left(c_1 e^{-\frac{i\sqrt{3} \ln(\tau)}{2}} + c_2 e^{\frac{i\sqrt{3} \ln(\tau)}{2}} \right)$$

Using Euler relation, the expression $c_1e^{iA} + c_2e^{-iA}$ is transformed to $c_1 \cos A + c_1 \sin A$ where the constants are free to change. Applying this to the above result gives

$$y(\tau) = \sqrt{\tau} \left(c_1 \cos \left(\frac{\sqrt{3} \ln(\tau)}{2} \right) + c_2 \sin \left(\frac{\sqrt{3} \ln(\tau)}{2} \right) \right)$$

The above solution is now transformed back to y using (6) which results in

$$y = \sqrt{-e^{-\frac{t^2}{2}}} \left(c_1 \cos \left(\frac{\sqrt{3} \ln \left(-e^{-\frac{t^2}{2}} \right)}{2} \right) + c_2 \sin \left(\frac{\sqrt{3} \ln \left(-e^{-\frac{t^2}{2}} \right)}{2} \right) \right)$$

Summary

The solution(s) found are the following

$$y = \sqrt{-e^{-\frac{t^2}{2}}} \left(c_1 \cos \left(\frac{\sqrt{3} \ln \left(-e^{-\frac{t^2}{2}} \right)}{2} \right) + c_2 \sin \left(\frac{\sqrt{3} \ln \left(-e^{-\frac{t^2}{2}} \right)}{2} \right) \right) \quad (1)$$

Verification of solutions

$$y = \sqrt{-e^{-\frac{t^2}{2}}} \left(c_1 \cos \left(\frac{\sqrt{3} \ln \left(-e^{-\frac{t^2}{2}} \right)}{2} \right) + c_2 \sin \left(\frac{\sqrt{3} \ln \left(-e^{-\frac{t^2}{2}} \right)}{2} \right) \right)$$

Verified OK.

8.30.2 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$ty'' + (t^2 - 1)y' + yt^3 = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{t^2 - 1}{t}$$
$$q(t) = t^2$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\begin{aligned} \tau' &= \frac{1}{c}\sqrt{q} \\ &= \frac{\sqrt{t^2}}{c} \\ \tau'' &= \frac{t}{c\sqrt{t^2}} \end{aligned} \quad (6)$$

Substituting the above into (4) results in

$$\begin{aligned} p_1(\tau) &= \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \\ &= \frac{\frac{t}{c\sqrt{t^2}} + \frac{t^2-1}{t}\frac{\sqrt{t^2}}{c}}{\left(\frac{\sqrt{t^2}}{c}\right)^2} \\ &= c \end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned} y(\tau)'' + p_1 y(\tau)' + q_1 y(\tau) &= 0 \\ \frac{d^2}{d\tau^2} y(\tau) + c \left(\frac{d}{d\tau} y(\tau) \right) + c^2 y(\tau) &= 0 \end{aligned} \quad (7)$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{-\frac{c\tau}{2}} \left(c_1 \cos \left(\frac{c\sqrt{3}\tau}{2} \right) + c_2 \sin \left(\frac{c\sqrt{3}\tau}{2} \right) \right)$$

Now from (6)

$$\begin{aligned} \tau &= \int \frac{1}{c} \sqrt{q} dt \\ &= \frac{\int \sqrt{t^2} dt}{c} \\ &= \frac{t\sqrt{t^2}}{2c} \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = e^{-\frac{t^2}{4}} \left(c_1 \cos \left(\frac{\sqrt{3}t^2}{4} \right) + c_2 \sin \left(\frac{\sqrt{3}t^2}{4} \right) \right)$$

Summary

The solution(s) found are the following

$$y = e^{-\frac{t^2}{4}} \left(c_1 \cos \left(\frac{\sqrt{3}t^2}{4} \right) + c_2 \sin \left(\frac{\sqrt{3}t^2}{4} \right) \right) \quad (1)$$

Verification of solutions

$$y = e^{-\frac{t^2}{4}} \left(c_1 \cos \left(\frac{\sqrt{3}t^2}{4} \right) + c_2 \sin \left(\frac{\sqrt{3}t^2}{4} \right) \right)$$

Verified OK.

8.30.3 Solving using Kovacic algorithm

Writing the ode as

$$ty'' + (t^2 - 1)y' + yt^3 = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t \\ B &= t^2 - 1 \\ C &= t^3 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-3t^4 + 3}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -3t^4 + 3 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(\frac{-3t^4 + 3}{4t^2} \right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 420: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 4 \\ &= -2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is -2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Therefore

$$L = [1, 2]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{3t^2}{4} + \frac{3}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = \frac{3}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{3}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is $O_r(\infty) = -2$ then

$$v = \frac{-O_r(\infty)}{2} = \frac{2}{2} = 1$$

$[\sqrt{r}]_\infty$ is the sum of terms involving t^i for $0 \leq i \leq v$ in the Laurent series for \sqrt{r} at ∞ . Therefore

$$\begin{aligned} [\sqrt{r}]_\infty &= \sum_{i=0}^v a_i t^i \\ &= \sum_{i=0}^1 a_i t^i \end{aligned} \tag{8}$$

Let a be the coefficient of $t^v = t^1$ in the above sum. The Laurent series of \sqrt{r} at ∞ is

$$\sqrt{r} \approx \frac{i\sqrt{3}t}{2} - \frac{i\sqrt{3}}{4t^3} - \frac{i\sqrt{3}}{16t^7} - \frac{i\sqrt{3}}{32t^{11}} - \frac{5i\sqrt{3}}{256t^{15}} - \frac{7i\sqrt{3}}{512t^{19}} - \frac{21i\sqrt{3}}{2048t^{23}} - \frac{33i\sqrt{3}}{4096t^{27}} + \dots \tag{9}$$

Comparing Eq. (9) with Eq. (8) shows that

$$a = \frac{i\sqrt{3}}{2}$$

From Eq. (9) the sum up to $v = 1$ gives

$$\begin{aligned} [\sqrt{r}]_\infty &= \sum_{i=0}^1 a_i t^i \\ &= \frac{i\sqrt{3}t}{2} \end{aligned} \tag{10}$$

Now we need to find b , where b be the coefficient of $t^{v-1} = t^0 = 1$ in r minus the coefficient of same term but in $([\sqrt{r}]_\infty)^2$ where $[\sqrt{r}]_\infty$ was found above in Eq (10). Hence

$$([\sqrt{r}]_\infty)^2 = -\frac{3t^2}{4}$$

This shows that the coefficient of 1 in the above is 0. Now we need to find the coefficient of 1 in r . How this is done depends on if $v = 0$ or not. Since $v = 1$ which is not zero, then starting $r = \frac{s}{t}$, we do long division and write this in the form

$$r = Q + \frac{R}{t}$$

Where Q is the quotient and R is the remainder. Then the coefficient of 1 in r will be the coefficient this term in the quotient. Doing long division gives

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{-3t^4 + 3}{4t^2} \\ &= Q + \frac{R}{4t^2} \\ &= \left(-\frac{3t^2}{4}\right) + \left(\frac{3}{4t^2}\right) \\ &= -\frac{3t^2}{4} + \frac{3}{4t^2} \end{aligned}$$

We see that the coefficient of the term t in the quotient is 0. Now b can be found.

$$\begin{aligned} b &= (0) - (0) \\ &= 0 \end{aligned}$$

Hence

$$\begin{aligned} [\sqrt{r}]_{\infty} &= \frac{i\sqrt{3}t}{2} \\ \alpha_{\infty}^+ &= \frac{1}{2} \left(\frac{b}{a} - v \right) = \frac{1}{2} \left(\frac{0}{\frac{i\sqrt{3}}{2}} - 1 \right) = -\frac{1}{2} \\ \alpha_{\infty}^- &= \frac{1}{2} \left(-\frac{b}{a} - v \right) = \frac{1}{2} \left(-\frac{0}{\frac{i\sqrt{3}}{2}} - 1 \right) = -\frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = \frac{-3t^4 + 3}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{3}{2}$	$-\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
-2	$\frac{i\sqrt{3}t}{2}$	$-\frac{1}{2}$	$-\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_\infty^- = -\frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_\infty^- - (\alpha_{c_1}^-) \\ &= -\frac{1}{2} - \left(-\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_\infty$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^-}{t - c_1} \right) + (-)[\sqrt{r}]_\infty \\ &= -\frac{1}{2t} + (-) \left(\frac{i\sqrt{3}t}{2} \right) \\ &= -\frac{1}{2t} - \frac{i\sqrt{3}t}{2} \\ &= \frac{-i\sqrt{3}t^2 - 1}{2t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$(0) + 2 \left(-\frac{1}{2t} - \frac{i\sqrt{3}t}{2} \right) (0) + \left(\left(\frac{1}{2t^2} - \frac{i\sqrt{3}}{2} \right) + \left(-\frac{1}{2t} - \frac{i\sqrt{3}t}{2} \right)^2 - \left(\frac{-3t^4 + 3}{4t^2} \right) \right) = 0$$

$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \left(-\frac{1}{2t} - \frac{i\sqrt{3}t}{2} \right) dt} \\ &= \frac{e^{-\frac{i\sqrt{3}t^2}{4}}}{\sqrt{t}} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{t^2-1}{t} dt} \\ &= z_1 e^{-\frac{t^2}{4} + \frac{\ln(t)}{2}} \\ &= z_1 \left(\sqrt{t} e^{-\frac{t^2}{4}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{t^2(1+i\sqrt{3})}{4}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{t^2-1}{t} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{-\frac{t^2}{2} + \ln(t)}}{(y_1)^2} dt \\ &= y_1 \left(-\frac{i\sqrt{3} e^{\frac{i\sqrt{3}t^2}{2}}}{3} \right) \end{aligned}$$

Therefore the solution is

$$y = c_1 y_1 + c_2 y_2$$

$$= c_1 \left(e^{-\frac{t^2(1+i\sqrt{3})}{4}} \right) + c_2 \left(e^{-\frac{t^2(1+i\sqrt{3})}{4}} \left(-\frac{i\sqrt{3} e^{\frac{i\sqrt{3}t^2}{2}}}{3} \right) \right)$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{t^2(1+i\sqrt{3})}{4}} - \frac{ic_2\sqrt{3} e^{\frac{t^2(i\sqrt{3}-1)}{4}}}{3} \quad (1)$$

Verification of solutions

$$y = c_1 e^{-\frac{t^2(1+i\sqrt{3})}{4}} - \frac{ic_2\sqrt{3} e^{\frac{t^2(i\sqrt{3}-1)}{4}}}{3}$$

Verified OK.

8.30.4 Maple step by step solution

Let's solve

$$ty'' + (t^2 - 1)y' + yt^3 = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{(t^2-1)y'}{t} - yt^2$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{(t^2-1)y'}{t} + yt^2 = 0$$

- Check to see if $t_0 = 0$ is a regular singular point

- Define functions

$$\left[P_2(t) = \frac{t^2-1}{t}, P_3(t) = t^2 \right]$$

- $t \cdot P_2(t)$ is analytic at $t = 0$

$$(t \cdot P_2(t)) \Big|_{t=0} = -1$$

- $t^2 \cdot P_3(t)$ is analytic at $t = 0$

$$(t^2 \cdot P_3(t)) \Big|_{t=0} = 0$$

- $t = 0$ is a regular singular point

Check to see if $t_0 = 0$ is a regular singular point

$$t_0 = 0$$

- Multiply by denominators

$$ty'' + (t^2 - 1)y' + yt^3 = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k t^{k+r}$$

- Rewrite ODE with series expansions

- Convert $t^3 \cdot y$ to series expansion

$$t^3 \cdot y = \sum_{k=0}^{\infty} a_k t^{k+r+3}$$

- Shift index using $k \rightarrow k - 3$

$$t^3 \cdot y = \sum_{k=3}^{\infty} a_{k-3} t^{k+r}$$

- Convert $t^m \cdot y'$ to series expansion for $m = 0..2$

$$t^m \cdot y' = \sum_{k=0}^{\infty} a_k (k+r) t^{k+r-1+m}$$

- Shift index using $k \rightarrow k + 1 - m$

$$t^m \cdot y' = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k+1-m+r) t^{k+r}$$

- Convert $t \cdot y''$ to series expansion

$$t \cdot y'' = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) t^{k+r-1}$$

- Shift index using $k \rightarrow k + 1$

$$t \cdot y'' = \sum_{k=-1}^{\infty} a_{k+1} (k+1+r)(k+r) t^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r(-2+r) t^{-1+r} + a_1 (1+r)(-1+r) t^r + (a_2 (2+r)r + a_0 r) t^{1+r} + (a_3 (3+r)(1+r) + a_1 (1+r)) t^{2+r} + \dots$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$r(-2 + r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \{0, 2\}$$

- The coefficients of each power of t must be 0

$$[a_1(1 + r)(-1 + r) = 0, a_2(2 + r)r + a_0r = 0, a_3(3 + r)(1 + r) + a_1(1 + r) = 0]$$

- Solve for the dependent coefficient(s)

$$\{a_1 = 0, a_2 = -\frac{a_0}{2+r}, a_3 = 0\}$$

- Each term in the series must be 0, giving the recursion relation

$$a_{k+1}(k + 1 + r)(k + r - 1) + a_{k-1}(k + r - 1) + a_{k-3} = 0$$

- Shift index using $k \rightarrow k + 3$

$$a_{k+4}(k + 4 + r)(k + 2 + r) + a_{k+2}(k + 2 + r) + a_k = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+4} = -\frac{ka_{k+2} + ra_{k+2} + a_k + 2a_{k+2}}{(k+4+r)(k+2+r)}$$

- Recursion relation for $r = 0$

$$a_{k+4} = -\frac{ka_{k+2} + a_k + 2a_{k+2}}{(k+4)(k+2)}$$

- Solution for $r = 0$

$$\left[y = \sum_{k=0}^{\infty} a_k t^k, a_{k+4} = -\frac{ka_{k+2} + a_k + 2a_{k+2}}{(k+4)(k+2)}, a_1 = 0, a_2 = -\frac{a_0}{2}, a_3 = 0 \right]$$

- Recursion relation for $r = 2$

$$a_{k+4} = -\frac{ka_{k+2} + a_k + 4a_{k+2}}{(k+6)(k+4)}$$

- Solution for $r = 2$

$$\left[y = \sum_{k=0}^{\infty} a_k t^{k+2}, a_{k+4} = -\frac{ka_{k+2} + a_k + 4a_{k+2}}{(k+6)(k+4)}, a_1 = 0, a_2 = -\frac{a_0}{4}, a_3 = 0 \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k t^k \right) + \left(\sum_{k=0}^{\infty} b_k t^{k+2} \right), a_{k+4} = -\frac{ka_{k+2} + a_k + 2a_{k+2}}{(k+4)(k+2)}, a_1 = 0, a_2 = -\frac{a_0}{2}, a_3 = 0, b_{k+4} = -\frac{kb_k}{(k+6)(k+4)} \right]$$

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacics algorithm  
  A Liouvillian solution exists  
  Group is reducible or imprimitive  
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 34

```
dsolve(t*difff(y(t),t$2)+ (t^2-1)*difff(y(t),t)+t^3*y(t) = 0,y(t), singsol=all)
```

$$y(t) = e^{-\frac{t^2}{4}} \left(c_1 \cos \left(\frac{t^2 \sqrt{3}}{4} \right) + c_2 \sin \left(\frac{t^2 \sqrt{3}}{4} \right) \right)$$

✓ Solution by Mathematica

Time used: 0.03 (sec). Leaf size: 48

```
DSolve[t*y''[t]+(t^2-1)*y'[t]+t^3*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{-\frac{t^2}{4}} \left(c_2 \cos \left(\frac{\sqrt{3}t^2}{4} \right) + c_1 \sin \left(\frac{\sqrt{3}t^2}{4} \right) \right)$$

9 Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

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9.1 problem 1

9.1.1	Solving as second order linear constant coeff ode	2368
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Internal problem ID [653]

Internal file name [OUTPUT/653_Sunday_June_05_2022_01_46_28_AM_41212301/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 1.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - 2y' + y = 0$$

9.1.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -2, C = 1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - 2\lambda e^{\lambda x} + e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - 2\lambda + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -2, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(-2)^2 - (4)(1)(1)} \\ &= 1 \end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = -1$. Therefore the solution is

$$y = c_1 e^x + c_2 x e^x \tag{1}$$

Summary

The solution(s) found are the following

$$y = c_1 e^x + c_2 x e^x \tag{1}$$

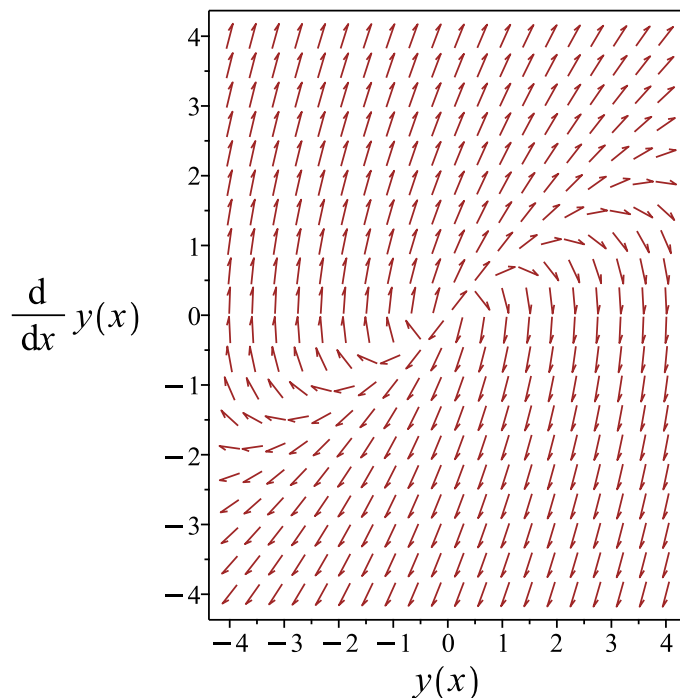


Figure 430: Slope field plot

Verification of solutions

$$y = c_1 e^x + c_2 x e^x$$

Verified OK.

9.1.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(x)y' + \frac{(p(x)^2 + p'(x))y}{2} = f(x)$$

Where $p(x) = -2$. Therefore, there is an integrating factor given by

$$\begin{aligned} M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int -2 dx} \\ &= e^{-x} \end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned} (M(x)y)'' &= 0 \\ (e^{-x}y)'' &= 0 \end{aligned}$$

Integrating once gives

$$(e^{-x}y)' = c_1$$

Integrating again gives

$$(e^{-x}y) = c_1x + c_2$$

Hence the solution is

$$y = \frac{c_1x + c_2}{e^{-x}}$$

Or

$$y = c_1x e^x + c_2e^x$$

Summary

The solution(s) found are the following

$$y = c_1x e^x + c_2e^x \tag{1}$$

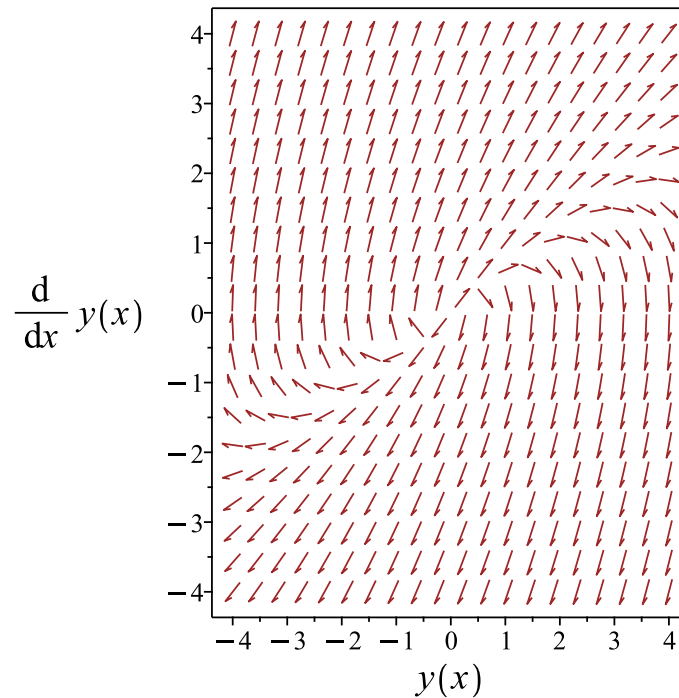


Figure 431: Slope field plot

Verification of solutions

$$y = c_1 x e^x + c_2 e^x$$

Verified OK.

9.1.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 2y' + y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= -2 \\ C &= 1 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 0 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = 0 \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 422: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-2}{1} dx} \\ &= z_1 e^x \\ &= z_1(e^x)\end{aligned}$$

Which simplifies to

$$y_1 = e^x$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{2x}}{(y_1)^2} dx \\ &= y_1(x)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1(e^x) + c_2(e^x(x))\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^x + c_2 x e^x \tag{1}$$

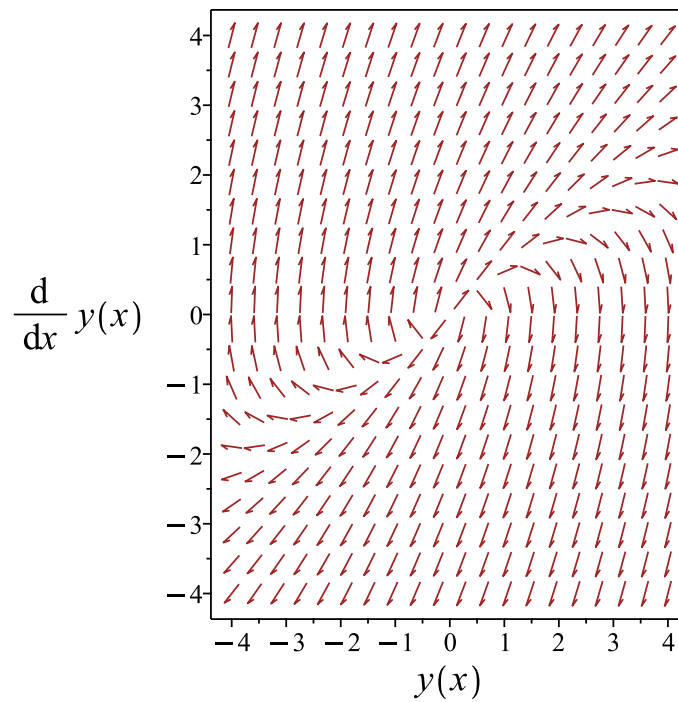


Figure 432: Slope field plot

Verification of solutions

$$y = c_1 e^x + c_2 x e^x$$

Verified OK.

9.1.4 Maple step by step solution

Let's solve

$$y'' - 2y' + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 - 2r + 1 = 0$$

- Factor the characteristic polynomial

$$(r - 1)^2 = 0$$

- Root of the characteristic polynomial

$$r = 1$$

- 1st solution of the ODE

$$y_1(x) = e^x$$

- Repeated root, multiply $y_1(x)$ by x to ensure linear independence

$$y_2(x) = x e^x$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^x + c_2 x e^x$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```


✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 12

```
dsolve(diff(y(x),x$2)-2*diff(y(x),x)+y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^x(c_2x + c_1)$$

✓ Solution by Mathematica

Time used: 0.01 (sec). Leaf size: 16

```
DSolve[y''[x]-2*y'[x]+y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^x(c_2x + c_1)$$

9.2 problem 2

9.2.1	Solving as second order linear constant coeff ode	2377
9.2.2	Solving as linear second order ode solved by an integrating factor ode	2379
9.2.3	Solving using Kovacic algorithm	2380
9.2.4	Maple step by step solution	2384

Internal problem ID [654]

Internal file name [OUTPUT/654_Sunday_June_05_2022_01_46_28_AM_19432812/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 2.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$9y'' + 6y' + y = 0$$

9.2.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 9, B = 6, C = 1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$9\lambda^2 e^{\lambda x} + 6\lambda e^{\lambda x} + e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$9\lambda^2 + 6\lambda + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 9, B = 6, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-6}{(2)(9)} \pm \frac{1}{(2)(9)} \sqrt{(6)^2 - (4)(9)(1)} \\ &= -\frac{1}{3} \end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = \frac{1}{3}$. Therefore the solution is

$$y = c_1 e^{-\frac{x}{3}} + c_2 x e^{-\frac{x}{3}} \quad (1)$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{x}{3}} + c_2 x e^{-\frac{x}{3}} \quad (1)$$

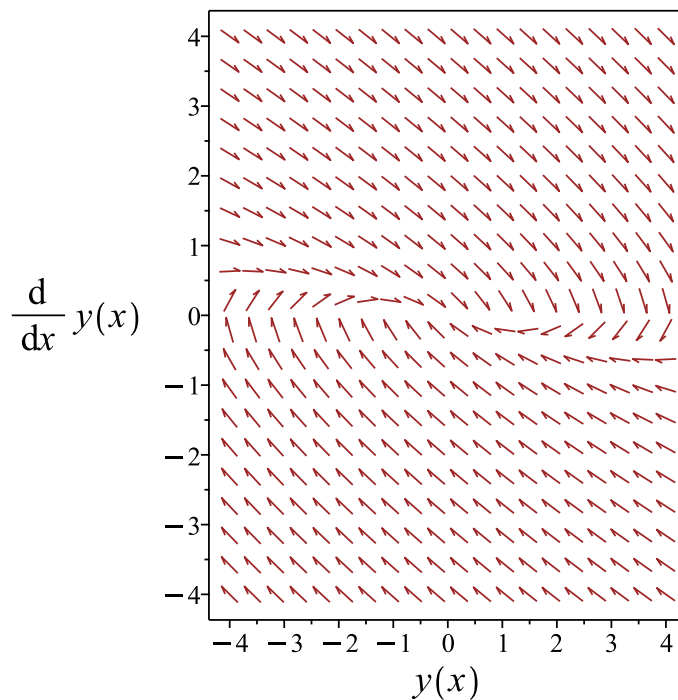


Figure 433: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{x}{3}} + c_2 x e^{-\frac{x}{3}}$$

Verified OK.

9.2.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(x)y' + \frac{(p(x)^2 + p'(x))y}{2} = f(x)$$

Where $p(x) = \frac{2}{3}$. Therefore, there is an integrating factor given by

$$\begin{aligned}M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int \frac{2}{3} dx} \\ &= e^{\frac{x}{3}}\end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned}(M(x)y)'' &= 0 \\ (e^{\frac{x}{3}}y)'' &= 0\end{aligned}$$

Integrating once gives

$$(e^{\frac{x}{3}}y)' = c_1$$

Integrating again gives

$$(e^{\frac{x}{3}}y) = c_1x + c_2$$

Hence the solution is

$$y = \frac{c_1x + c_2}{e^{\frac{x}{3}}}$$

Or

$$y = c_1x e^{-\frac{x}{3}} + c_2e^{-\frac{x}{3}}$$

Summary

The solution(s) found are the following

$$y = c_1x e^{-\frac{x}{3}} + c_2e^{-\frac{x}{3}} \quad (1)$$

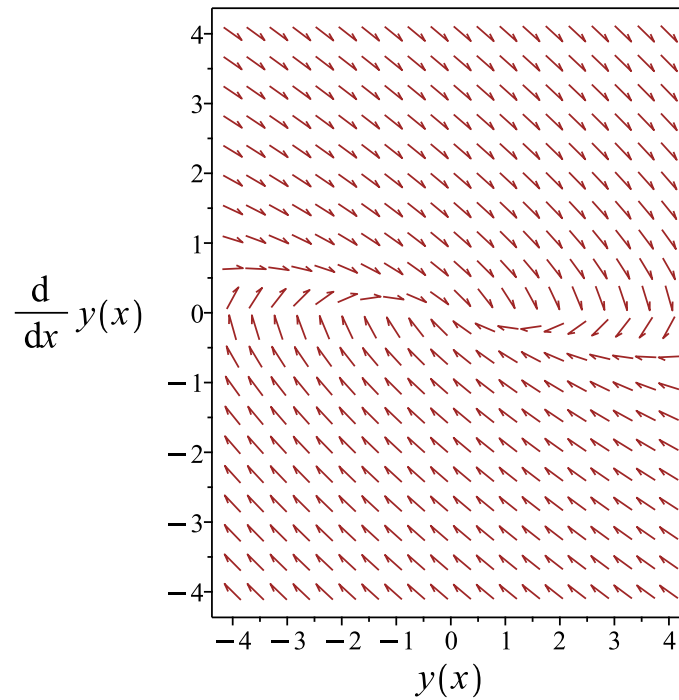


Figure 434: Slope field plot

Verification of solutions

$$y = c_1 x e^{-\frac{x}{3}} + c_2 e^{-\frac{x}{3}}$$

Verified OK.

9.2.3 Solving using Kovacic algorithm

Writing the ode as

$$9y'' + 6y' + y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 9 \tag{3}$$

$$B = 6$$

$$C = 1$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 0 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = 0 \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 424: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{6}{9} dx} \\ &= z_1 e^{-\frac{x}{3}} \\ &= z_1 \left(e^{-\frac{x}{3}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{x}{3}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{6}{9} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-\frac{2x}{3}}}{(y_1)^2} dx \\ &= y_1(x)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-\frac{x}{3}} \right) + c_2 \left(e^{-\frac{x}{3}}(x) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{x}{3}} + c_2 x e^{-\frac{x}{3}} \quad (1)$$

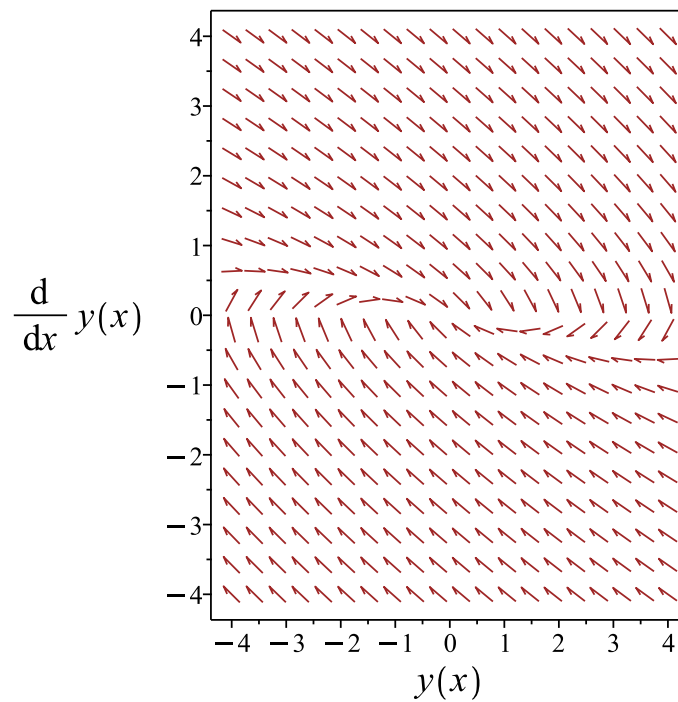


Figure 435: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{x}{3}} + c_2 x e^{-\frac{x}{3}}$$

Verified OK.

9.2.4 Maple step by step solution

Let's solve

$$9y'' + 6y' + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{2y'}{3} - \frac{y}{9}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{2y'}{3} + \frac{y}{9} = 0$$

- Characteristic polynomial of ODE

$$r^2 + \frac{2}{3}r + \frac{1}{9} = 0$$

- Factor the characteristic polynomial

$$\frac{(3r+1)^2}{9} = 0$$

- Root of the characteristic polynomial

$$r = -\frac{1}{3}$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{x}{3}}$$

- Repeated root, multiply $y_1(x)$ by x to ensure linear independence

$$y_2(x) = x e^{-\frac{x}{3}}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-\frac{x}{3}} + c_2 x e^{-\frac{x}{3}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(9*diff(y(x),x$2)+6*diff(y(x),x)+y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-\frac{x}{3}}(c_2x + c_1)$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 20

```
DSolve[9*y''[x]+6*y'[x]+y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x/3}(c_2x + c_1)$$

9.3 problem 3

9.3.1 Solving as second order linear constant coeff ode	2386
9.3.2 Solving using Kovacic algorithm	2388
9.3.3 Maple step by step solution	2392

Internal problem ID [655]

Internal file name [OUTPUT/655_Sunday_June_05_2022_01_46_29_AM_67197608/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 3.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$4y'' - 4y' - 3y = 0$$

9.3.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 4, B = -4, C = -3$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$4\lambda^2 e^{\lambda x} - 4\lambda e^{\lambda x} - 3e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$4\lambda^2 - 4\lambda - 3 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 4, B = -4, C = -3$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{4}{(2)(4)} \pm \frac{1}{(2)(4)} \sqrt{-4^2 - (4)(4)(-3)} \\ &= \frac{1}{2} \pm 1\end{aligned}$$

Hence

$$\lambda_1 = \frac{1}{2} + 1$$

$$\lambda_2 = \frac{1}{2} - 1$$

Which simplifies to

$$\lambda_1 = \frac{3}{2}$$

$$\lambda_2 = -\frac{1}{2}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$y = c_1 e^{(\frac{3}{2})x} + c_2 e^{(-\frac{1}{2})x}$$

Or

$$y = c_1 e^{\frac{3x}{2}} + c_2 e^{-\frac{x}{2}}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{\frac{3x}{2}} + c_2 e^{-\frac{x}{2}} \quad (1)$$

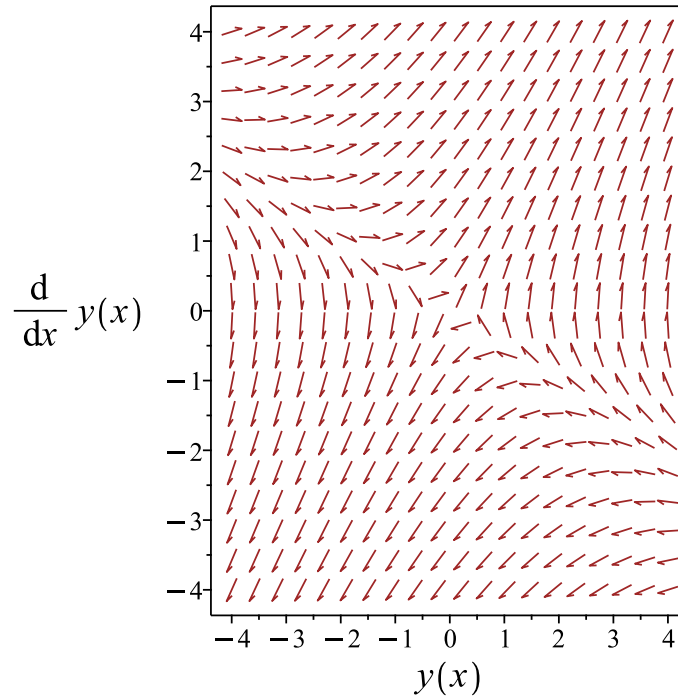


Figure 436: Slope field plot

Verification of solutions

$$y = c_1 e^{\frac{3x}{2}} + c_2 e^{-\frac{x}{2}}$$

Verified OK.

9.3.2 Solving using Kovacic algorithm

Writing the ode as

$$4y'' - 4y' - 3y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 4 \\ B &= -4 \\ C &= -3 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 1$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 426: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-x}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-4}{4} dx} \\ &= z_1 e^{\frac{x}{2}} \\ &= z_1 \left(e^{\frac{x}{2}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{x}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-4}{4} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^x}{(y_1)^2} dx \\ &= y_1 \left(\frac{e^{2x}}{2} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-\frac{x}{2}} \right) + c_2 \left(e^{-\frac{x}{2}} \left(\frac{e^{2x}}{2} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{x}{2}} + \frac{c_2 e^{\frac{3x}{2}}}{2} \quad (1)$$

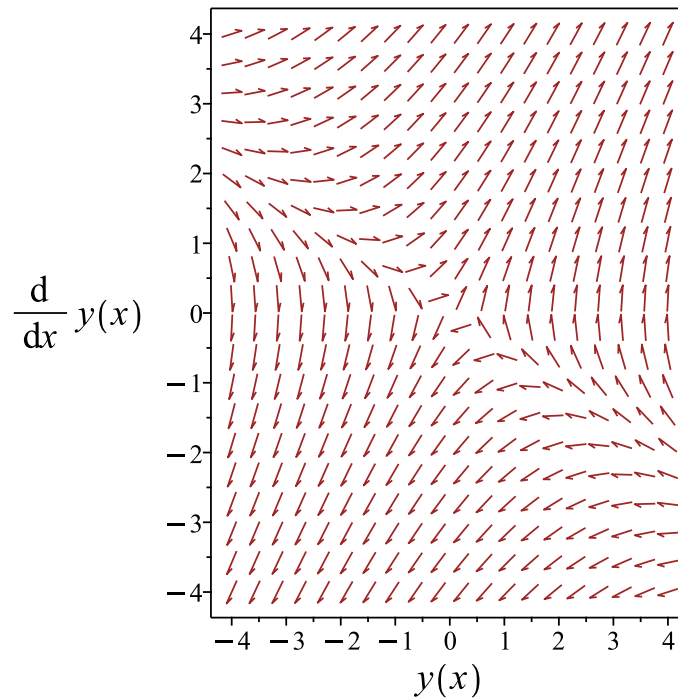


Figure 437: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{x}{2}} + \frac{c_2 e^{\frac{3x}{2}}}{2}$$

Verified OK.

9.3.3 Maple step by step solution

Let's solve

$$4y'' - 4y' - 3y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = y' + \frac{3y}{4}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - y' - \frac{3y}{4} = 0$$

- Characteristic polynomial of ODE

$$r^2 - r - \frac{3}{4} = 0$$

- Factor the characteristic polynomial

$$\frac{(2r+1)(2r-3)}{4} = 0$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{2}, \frac{3}{2}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{x}{2}}$$

- 2nd solution of the ODE

$$y_2(x) = e^{\frac{3x}{2}}$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^{-\frac{x}{2}} + c_2e^{\frac{3x}{2}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(4*diff(y(x),x$2)-4*diff(y(x),x)-3*y(x) = 0,y(x), singsol=all)
```

$$y(x) = c_1 e^{-\frac{x}{2}} + c_2 e^{\frac{3x}{2}}$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 24

```
DSolve[4*y''[x]-4*y'[x]-3*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x/2} (c_2 e^{2x} + c_1)$$

9.4 problem 4

9.4.1	Solving as second order linear constant coeff ode	2394
9.4.2	Solving as linear second order ode solved by an integrating factor ode	2396
9.4.3	Solving using Kovacic algorithm	2397
9.4.4	Maple step by step solution	2401

Internal problem ID [656]

Internal file name [OUTPUT/656_Sunday_June_05_2022_01_46_30_AM_9155155/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 4.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$4y'' + 12y' + 9y = 0$$

9.4.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 4, B = 12, C = 9$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$4\lambda^2 e^{\lambda x} + 12\lambda e^{\lambda x} + 9 e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$4\lambda^2 + 12\lambda + 9 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 4, B = 12, C = 9$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-12}{(2)(4)} \pm \frac{1}{(2)(4)} \sqrt{(12)^2 - (4)(4)(9)} \\ &= -\frac{3}{2} \end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = \frac{3}{2}$. Therefore the solution is

$$y = c_1 e^{-\frac{3x}{2}} + c_2 x e^{-\frac{3x}{2}} \quad (1)$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{3x}{2}} + c_2 x e^{-\frac{3x}{2}} \quad (1)$$

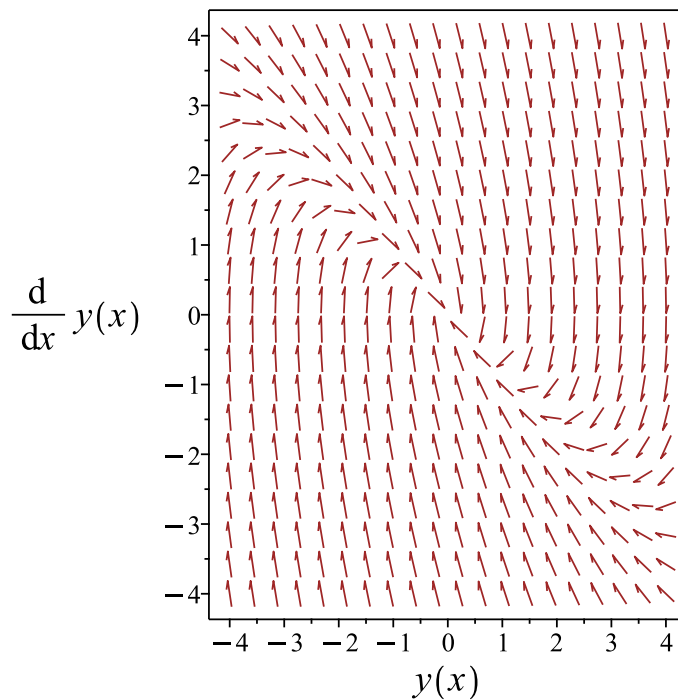


Figure 438: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{3x}{2}} + c_2 x e^{-\frac{3x}{2}}$$

Verified OK.

9.4.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(x)y' + \frac{(p(x)^2 + p'(x))y}{2} = f(x)$$

Where $p(x) = 3$. Therefore, there is an integrating factor given by

$$\begin{aligned} M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int 3 dx} \\ &= e^{\frac{3x}{2}} \end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned} (M(x)y)'' &= 0 \\ \left(e^{\frac{3x}{2}}y\right)'' &= 0 \end{aligned}$$

Integrating once gives

$$\left(e^{\frac{3x}{2}}y\right)' = c_1$$

Integrating again gives

$$\left(e^{\frac{3x}{2}}y\right) = c_1x + c_2$$

Hence the solution is

$$y = \frac{c_1x + c_2}{e^{\frac{3x}{2}}}$$

Or

$$y = c_1x e^{-\frac{3x}{2}} + c_2e^{-\frac{3x}{2}}$$

Summary

The solution(s) found are the following

$$y = c_1x e^{-\frac{3x}{2}} + c_2e^{-\frac{3x}{2}} \quad (1)$$

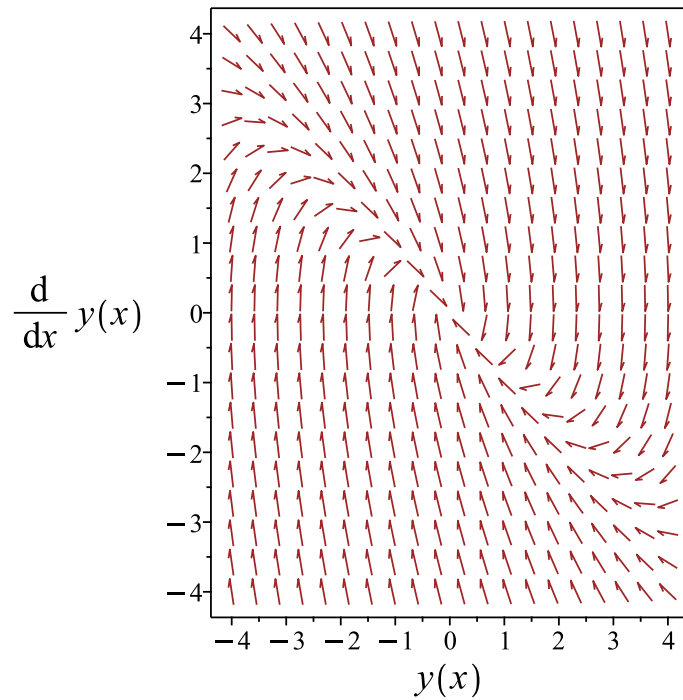


Figure 439: Slope field plot

Verification of solutions

$$y = c_1 x e^{-\frac{3x}{2}} + c_2 e^{-\frac{3x}{2}}$$

Verified OK.

9.4.3 Solving using Kovacic algorithm

Writing the ode as

$$4y'' + 12y' + 9y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 4 \\ B &= 12 \\ C &= 9 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 0$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = 0 \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 428: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{12}{4} dx} \\ &= z_1 e^{-\frac{3x}{2}} \\ &= z_1 \left(e^{-\frac{3x}{2}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{3x}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{12}{4} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-3x}}{(y_1)^2} dx \\ &= y_1(x)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-\frac{3x}{2}} \right) + c_2 \left(e^{-\frac{3x}{2}} (x) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{3x}{2}} + c_2 x e^{-\frac{3x}{2}} \quad (1)$$

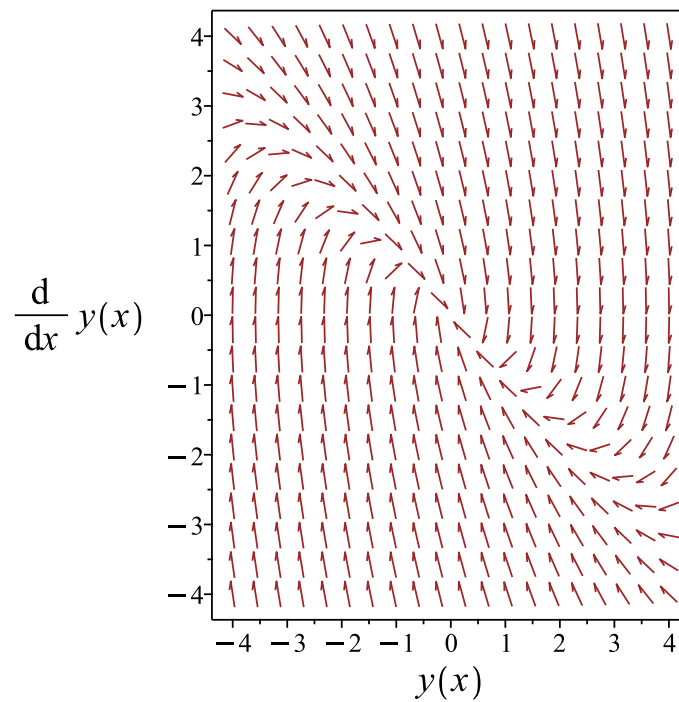


Figure 440: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{3x}{2}} + c_2 x e^{-\frac{3x}{2}}$$

Verified OK.

9.4.4 Maple step by step solution

Let's solve

$$4y'' + 12y' + 9y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -3y' - \frac{9y}{4}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + 3y' + \frac{9y}{4} = 0$$

- Characteristic polynomial of ODE

$$r^2 + 3r + \frac{9}{4} = 0$$

- Factor the characteristic polynomial

$$\frac{(2r+3)^2}{4} = 0$$

- Root of the characteristic polynomial

$$r = -\frac{3}{2}$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{3x}{2}}$$

- Repeated root, multiply $y_1(x)$ by x to ensure linear independence

$$y_2(x) = x e^{-\frac{3x}{2}}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-\frac{3x}{2}} + c_2 x e^{-\frac{3x}{2}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(4*diff(y(x),x$2)+12*diff(y(x),x)+9*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-\frac{3x}{2}}(c_2x + c_1)$$

✓ Solution by Mathematica

Time used: 0.014 (sec). Leaf size: 20

```
DSolve[4*y''[x]+12*y'[x]+9*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-3x/2}(c_2x + c_1)$$

9.5 problem 5

9.5.1 Solving as second order linear constant coeff ode	2403
9.5.2 Solving using Kovacic algorithm	2405
9.5.3 Maple step by step solution	2409

Internal problem ID [657]

Internal file name [OUTPUT/657_Sunday_June_05_2022_01_46_31_AM_43935394/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 5.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - 2y' + 10y = 0$$

9.5.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -2, C = 10$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - 2\lambda e^{\lambda x} + 10e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - 2\lambda + 10 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -2, C = 10$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-2^2 - (4)(1)(10)} \\ &= 1 \pm 3i\end{aligned}$$

Hence

$$\lambda_1 = 1 + 3i$$

$$\lambda_2 = 1 - 3i$$

Which simplifies to

$$\lambda_1 = 1 + 3i$$

$$\lambda_2 = 1 - 3i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 1$ and $\beta = 3$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^x(c_1 \cos(3x) + c_2 \sin(3x))$$

Summary

The solution(s) found are the following

$$y = e^x(c_1 \cos(3x) + c_2 \sin(3x)) \quad (1)$$

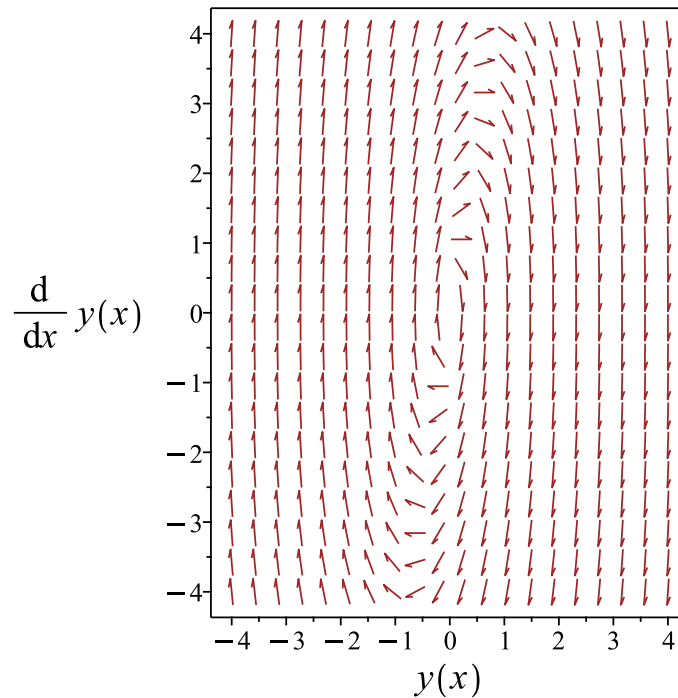


Figure 441: Slope field plot

Verification of solutions

$$y = e^x(c_1 \cos(3x) + c_2 \sin(3x))$$

Verified OK.

9.5.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 2y' + 10y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= -2 \\ C &= 10 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-9}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -9 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -9z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 430: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -9$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(3x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-2}{1} dx} \\ &= z_1 e^x \\ &= z_1 (e^x) \end{aligned}$$

Which simplifies to

$$y_1 = e^x \cos(3x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-2}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{2x}}{(y_1)^2} dx \\ &= y_1 \left(\frac{\tan(3x)}{3} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1(e^x \cos(3x)) + c_2\left(e^x \cos(3x) \left(\frac{\tan(3x)}{3}\right)\right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^x \cos(3x) + \frac{c_2 e^x \sin(3x)}{3} \quad (1)$$

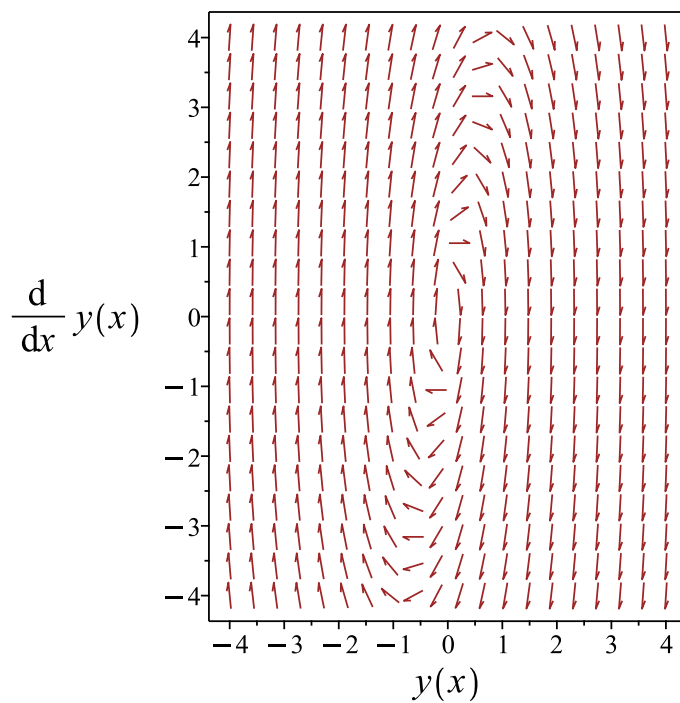


Figure 442: Slope field plot

Verification of solutions

$$y = c_1 e^x \cos(3x) + \frac{c_2 e^x \sin(3x)}{3}$$

Verified OK.

9.5.3 Maple step by step solution

Let's solve

$$y'' - 2y' + 10y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 - 2r + 10 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{2 \pm (\sqrt{-36})}{2}$$

- Roots of the characteristic polynomial

$$r = (1 - 3I, 1 + 3I)$$

- 1st solution of the ODE

$$y_1(x) = e^x \cos(3x)$$

- 2nd solution of the ODE

$$y_2(x) = e^x \sin(3x)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^x \cos(3x) + c_2 e^x \sin(3x)$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

```
dsolve(diff(y(x),x$2)-2*diff(y(x),x)+10*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^x(c_1 \sin(3x) + c_2 \cos(3x))$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 24

```
DSolve[y''[x]-2*y'[x]+10*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^x(c_2 \cos(3x) + c_1 \sin(3x))$$

9.6 problem 6

9.6.1	Solving as second order linear constant coeff ode	2411
9.6.2	Solving as linear second order ode solved by an integrating factor ode	2413
9.6.3	Solving using Kovacic algorithm	2414
9.6.4	Maple step by step solution	2418

Internal problem ID [658]

Internal file name [OUTPUT/658_Sunday_June_05_2022_01_46_32_AM_72033730/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 6.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$y'' - 6y' + 9y = 0$$

9.6.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -6, C = 9$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} - 6\lambda e^{\lambda x} + 9 e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - 6\lambda + 9 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -6, C = 9$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{6}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(-6)^2 - (4)(1)(9)} \\ &= 3 \end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = -3$. Therefore the solution is

$$y = c_1 e^{3x} + c_2 x e^{3x} \quad (1)$$

Summary

The solution(s) found are the following

$$y = c_1 e^{3x} + c_2 x e^{3x} \quad (1)$$

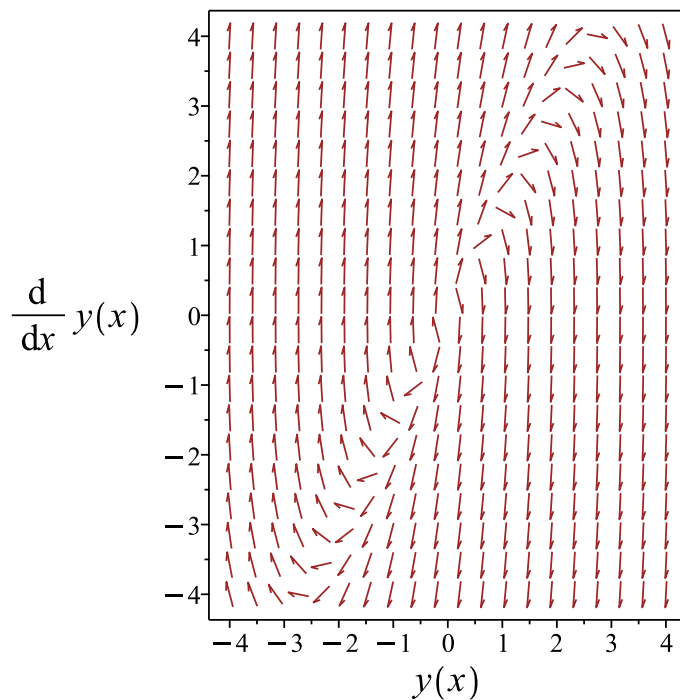


Figure 443: Slope field plot

Verification of solutions

$$y = c_1 e^{3x} + c_2 x e^{3x}$$

Verified OK.

9.6.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(x)y' + \frac{(p(x)^2 + p'(x))y}{2} = f(x)$$

Where $p(x) = -6$. Therefore, there is an integrating factor given by

$$\begin{aligned}M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int -6 dx} \\ &= e^{-3x}\end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned}(M(x)y)'' &= 0 \\ (e^{-3x}y)'' &= 0\end{aligned}$$

Integrating once gives

$$(e^{-3x}y)' = c_1$$

Integrating again gives

$$(e^{-3x}y) = c_1x + c_2$$

Hence the solution is

$$y = \frac{c_1x + c_2}{e^{-3x}}$$

Or

$$y = c_1x e^{3x} + c_2e^{3x}$$

Summary

The solution(s) found are the following

$$y = c_1x e^{3x} + c_2e^{3x} \tag{1}$$

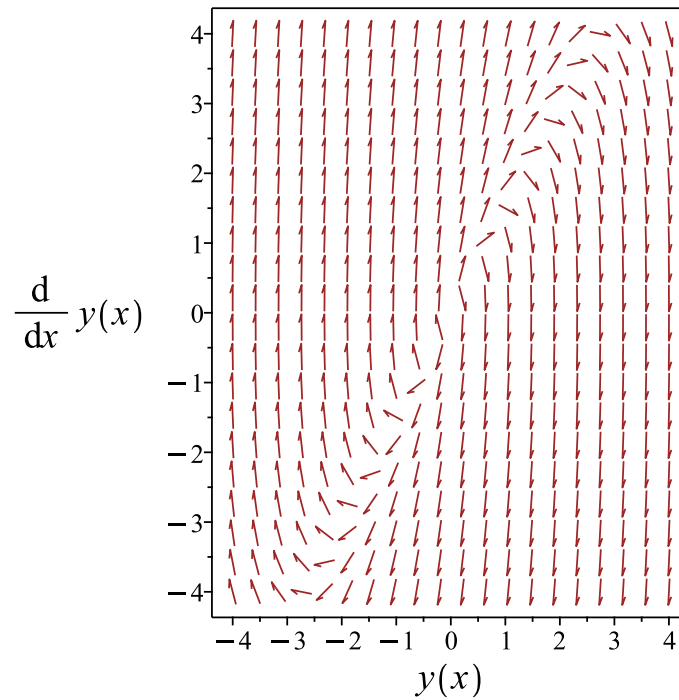


Figure 444: Slope field plot

Verification of solutions

$$y = c_1 x e^{3x} + c_2 e^{3x}$$

Verified OK.

9.6.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 6y' + 9y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = -6 \tag{3}$$

$$C = 9$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 0$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = 0 \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 432: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-6}{1} dx} \\ &= z_1 e^{3x} \\ &= z_1 (e^{3x}) \end{aligned}$$

Which simplifies to

$$y_1 = e^{3x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-6}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{6x}}{(y_1)^2} dx \\ &= y_1(x) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{3x}) + c_2 (e^{3x}(x))\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{3x} + c_2 x e^{3x} \quad (1)$$

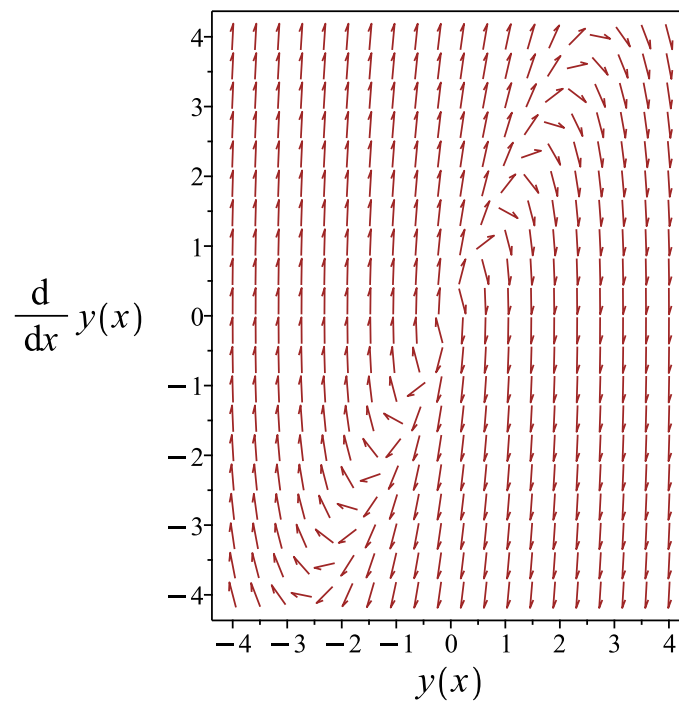


Figure 445: Slope field plot

Verification of solutions

$$y = c_1 e^{3x} + c_2 x e^{3x}$$

Verified OK.

9.6.4 Maple step by step solution

Let's solve

$$y'' - 6y' + 9y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of ODE

$$r^2 - 6r + 9 = 0$$

- Factor the characteristic polynomial

$$(r - 3)^2 = 0$$

- Root of the characteristic polynomial

$$r = 3$$

- 1st solution of the ODE

$$y_1(x) = e^{3x}$$

- Repeated root, multiply $y_1(x)$ by x to ensure linear independence

$$y_2(x) = x e^{3x}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{3x} + c_2 x e^{3x}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(diff(y(x),x$2)-6*diff(y(x),x)+9*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{3x}(c_2x + c_1)$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 18

```
DSolve[y''[x]-6*y'[x]+9*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{3x}(c_2x + c_1)$$

9.7 problem 7

9.7.1 Solving as second order linear constant coeff ode	2420
9.7.2 Solving using Kovacic algorithm	2422
9.7.3 Maple step by step solution	2426

Internal problem ID [659]

Internal file name [OUTPUT/659_Sunday_June_05_2022_01_46_32_AM_89046047/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 7.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$4y'' + 17y' + 4y = 0$$

9.7.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 4, B = 17, C = 4$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$4\lambda^2 e^{\lambda x} + 17\lambda e^{\lambda x} + 4e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$4\lambda^2 + 17\lambda + 4 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 4, B = 17, C = 4$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{-17}{(2)(4)} \pm \frac{1}{(2)(4)} \sqrt{17^2 - (4)(4)(4)} \\ &= -\frac{17}{8} \pm \frac{15}{8}\end{aligned}$$

Hence

$$\begin{aligned}\lambda_1 &= -\frac{17}{8} + \frac{15}{8} \\ \lambda_2 &= -\frac{17}{8} - \frac{15}{8}\end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= -\frac{1}{4} \\ \lambda_2 &= -4\end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned}y &= c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x} \\ y &= c_1 e^{(-\frac{1}{4})x} + c_2 e^{(-4)x}\end{aligned}$$

Or

$$y = c_1 e^{-\frac{x}{4}} + c_2 e^{-4x}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{x}{4}} + c_2 e^{-4x} \quad (1)$$

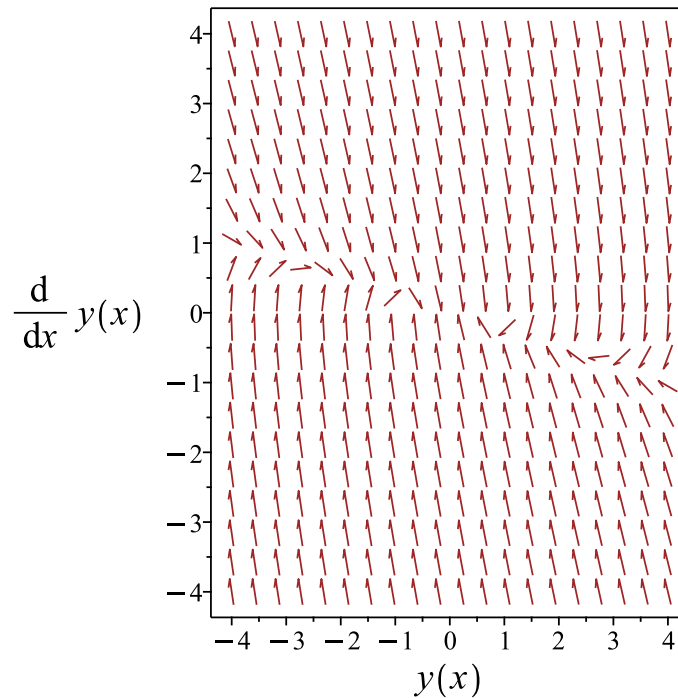


Figure 446: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{x}{4}} + c_2 e^{-4x}$$

Verified OK.

9.7.2 Solving using Kovacic algorithm

Writing the ode as

$$4y'' + 17y' + 4y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 4 \\ B &= 17 \\ C &= 4 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{225}{64} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 225 \\ t &= 64 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{225z(x)}{64} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 434: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{225}{64}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{15x}{8}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{17}{4} dx} \\ &= z_1 e^{-\frac{17x}{8}} \\ &= z_1 \left(e^{-\frac{17x}{8}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-4x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{17}{4} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-\frac{17x}{4}}}{(y_1)^2} dx \\ &= y_1 \left(\frac{4 e^{\frac{15x}{4}}}{15} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-4x}) + c_2 \left(e^{-4x} \left(\frac{4e^{\frac{15x}{4}}}{15} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-4x} + \frac{4c_2 e^{-\frac{x}{4}}}{15} \quad (1)$$

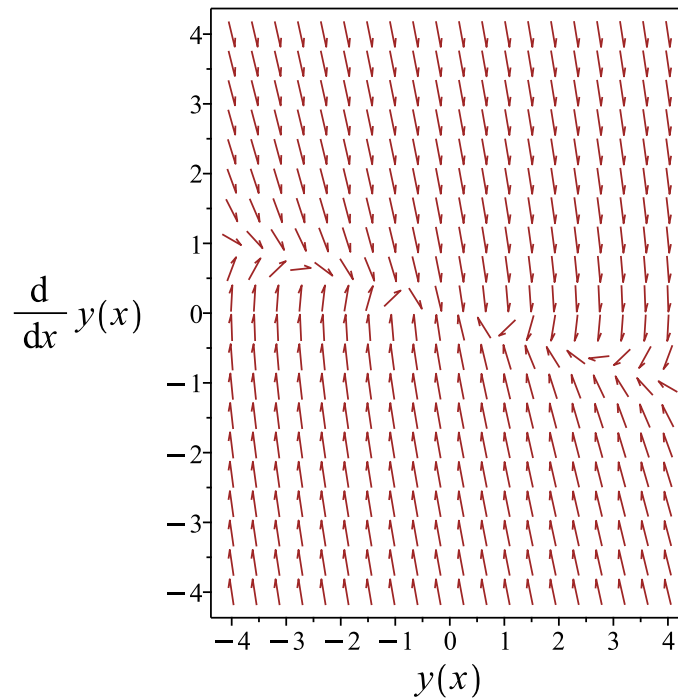


Figure 447: Slope field plot

Verification of solutions

$$y = c_1 e^{-4x} + \frac{4c_2 e^{-\frac{x}{4}}}{15}$$

Verified OK.

9.7.3 Maple step by step solution

Let's solve

$$4y'' + 17y' + 4y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{17y'}{4} - y$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{17y'}{4} + y = 0$$

- Characteristic polynomial of ODE

$$r^2 + \frac{17}{4}r + 1 = 0$$

- Factor the characteristic polynomial

$$\frac{(r+4)(4r+1)}{4} = 0$$

- Roots of the characteristic polynomial

$$r = \left(-4, -\frac{1}{4}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-4x}$$

- 2nd solution of the ODE

$$y_2(x) = e^{-\frac{x}{4}}$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^{-4x} + c_2e^{-\frac{x}{4}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve(4*diff(y(x),x$2)+17*diff(y(x),x)+4*y(x) = 0,y(x), singsol=all)
```

$$y(x) = c_1 e^{-4x} + c_2 e^{-\frac{x}{4}}$$

✓ Solution by Mathematica

Time used: 0.012 (sec). Leaf size: 24

```
DSolve[4*y'[x]+17*y'[x]+4*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-4x} (c_1 e^{15x/4} + c_2)$$

9.8 problem 8

9.8.1	Solving as second order linear constant coeff ode	2428
9.8.2	Solving as linear second order ode solved by an integrating factor ode	2430
9.8.3	Solving using Kovacic algorithm	2431
9.8.4	Maple step by step solution	2435

Internal problem ID [660]

Internal file name [OUTPUT/660_Sunday_June_05_2022_01_46_33_AM_61322730/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 8.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$16y'' + 24y' + 9y = 0$$

9.8.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 16, B = 24, C = 9$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$16\lambda^2 e^{\lambda x} + 24\lambda e^{\lambda x} + 9e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$16\lambda^2 + 24\lambda + 9 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 16, B = 24, C = 9$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-24}{(2)(16)} \pm \frac{1}{(2)(16)} \sqrt{(24)^2 - (4)(16)(9)} \\ &= -\frac{3}{4} \end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = \frac{3}{4}$. Therefore the solution is

$$y = c_1 e^{-\frac{3x}{4}} + c_2 x e^{-\frac{3x}{4}} \quad (1)$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{3x}{4}} + c_2 x e^{-\frac{3x}{4}} \quad (1)$$

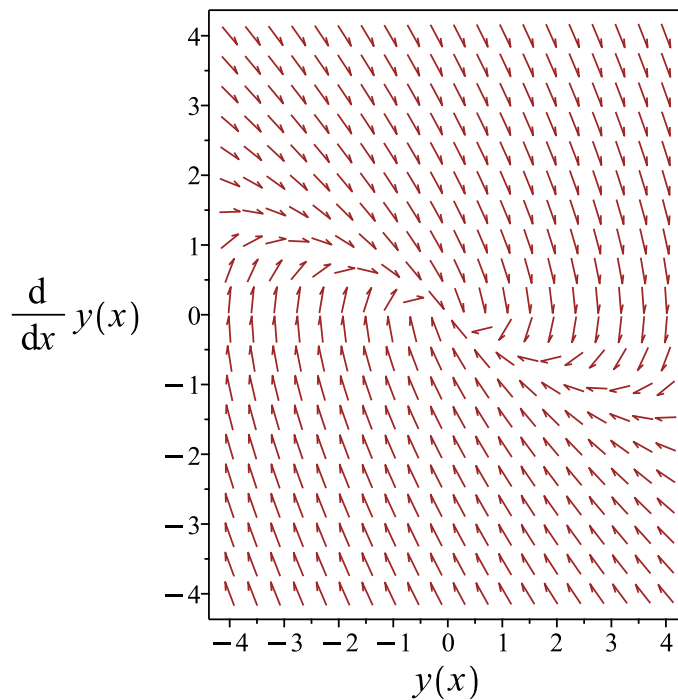


Figure 448: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{3x}{4}} + c_2 x e^{-\frac{3x}{4}}$$

Verified OK.

9.8.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(x)y' + \frac{(p(x)^2 + p'(x))y}{2} = f(x)$$

Where $p(x) = \frac{3}{2}$. Therefore, there is an integrating factor given by

$$\begin{aligned}M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int \frac{3}{2} dx} \\ &= e^{\frac{3x}{4}}\end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned}(M(x)y)'' &= 0 \\ \left(e^{\frac{3x}{4}}y\right)'' &= 0\end{aligned}$$

Integrating once gives

$$\left(e^{\frac{3x}{4}}y\right)' = c_1$$

Integrating again gives

$$\left(e^{\frac{3x}{4}}y\right) = c_1x + c_2$$

Hence the solution is

$$y = \frac{c_1x + c_2}{e^{\frac{3x}{4}}}$$

Or

$$y = c_1x e^{-\frac{3x}{4}} + e^{-\frac{3x}{4}}c_2$$

Summary

The solution(s) found are the following

$$y = c_1x e^{-\frac{3x}{4}} + e^{-\frac{3x}{4}}c_2 \quad (1)$$

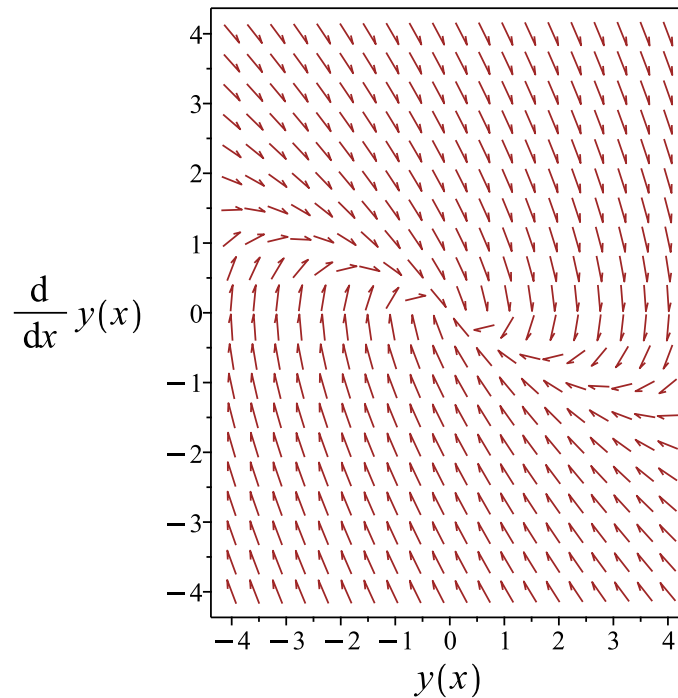


Figure 449: Slope field plot

Verification of solutions

$$y = c_1 x e^{-\frac{3x}{4}} + e^{-\frac{3x}{4}} c_2$$

Verified OK.

9.8.3 Solving using Kovacic algorithm

Writing the ode as

$$16y'' + 24y' + 9y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 16 \\ B &= 24 \\ C &= 9 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} \tag{5}$$

$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 0$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = 0 \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 436: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{24}{16} dx} \\ &= z_1 e^{-\frac{3x}{4}} \\ &= z_1 \left(e^{-\frac{3x}{4}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{3x}{4}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{24}{16} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-\frac{3x}{2}}}{(y_1)^2} dx \\ &= y_1(x) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-\frac{3x}{4}} \right) + c_2 \left(e^{-\frac{3x}{4}} (x) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{3x}{4}} + c_2 x e^{-\frac{3x}{4}} \quad (1)$$

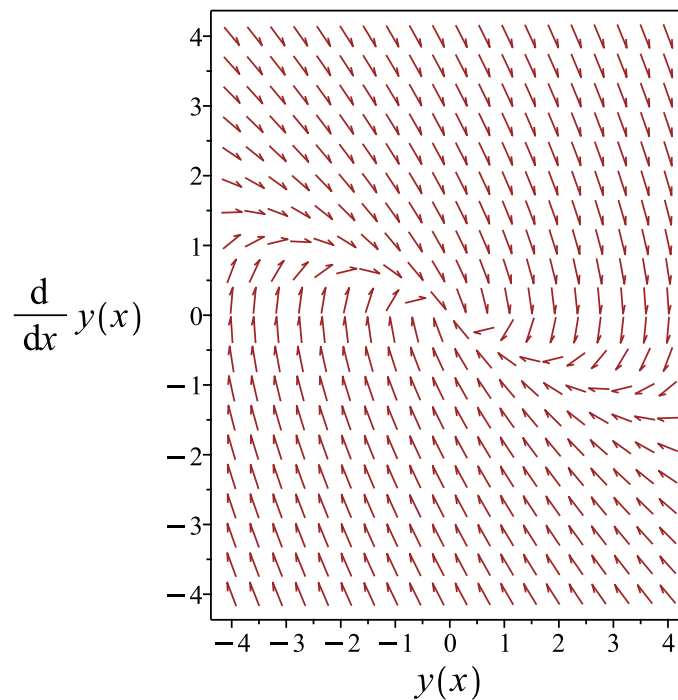


Figure 450: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{3x}{4}} + c_2 x e^{-\frac{3x}{4}}$$

Verified OK.

9.8.4 Maple step by step solution

Let's solve

$$16y'' + 24y' + 9y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{3y'}{2} - \frac{9y}{16}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{3y'}{2} + \frac{9y}{16} = 0$$

- Characteristic polynomial of ODE

$$r^2 + \frac{3}{2}r + \frac{9}{16} = 0$$

- Factor the characteristic polynomial

$$\frac{(4r+3)^2}{16} = 0$$

- Root of the characteristic polynomial

$$r = -\frac{3}{4}$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{3x}{4}}$$

- Repeated root, multiply $y_1(x)$ by x to ensure linear independence

$$y_2(x) = x e^{-\frac{3x}{4}}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-\frac{3x}{4}} + c_2 x e^{-\frac{3x}{4}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(16*diff(y(x),x$2)+24*diff(y(x),x)+9*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-\frac{3x}{4}}(c_2x + c_1)$$

✓ Solution by Mathematica

Time used: 0.014 (sec). Leaf size: 20

```
DSolve[16*y''[x]+24*y'[x]+9*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-3x/4}(c_2x + c_1)$$

9.9 problem 9

9.9.1	Solving as second order linear constant coeff ode	2437
9.9.2	Solving as linear second order ode solved by an integrating factor ode	2439
9.9.3	Solving using Kovacic algorithm	2440
9.9.4	Maple step by step solution	2444

Internal problem ID [661]

Internal file name [OUTPUT/661_Sunday_June_05_2022_01_46_34_AM_88618893/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 9.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$25y'' - 20y' + 4y = 0$$

9.9.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 25$, $B = -20$, $C = 4$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$25\lambda^2 e^{\lambda x} - 20\lambda e^{\lambda x} + 4e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$25\lambda^2 - 20\lambda + 4 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 25, B = -20, C = 4$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{20}{(2)(25)} \pm \frac{1}{(2)(25)} \sqrt{(-20)^2 - (4)(25)(4)} \\ &= \frac{2}{5} \end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = -\frac{2}{5}$. Therefore the solution is

$$y = c_1 e^{\frac{2x}{5}} + c_2 x e^{\frac{2x}{5}} \quad (1)$$

Summary

The solution(s) found are the following

$$y = c_1 e^{\frac{2x}{5}} + c_2 x e^{\frac{2x}{5}} \quad (1)$$

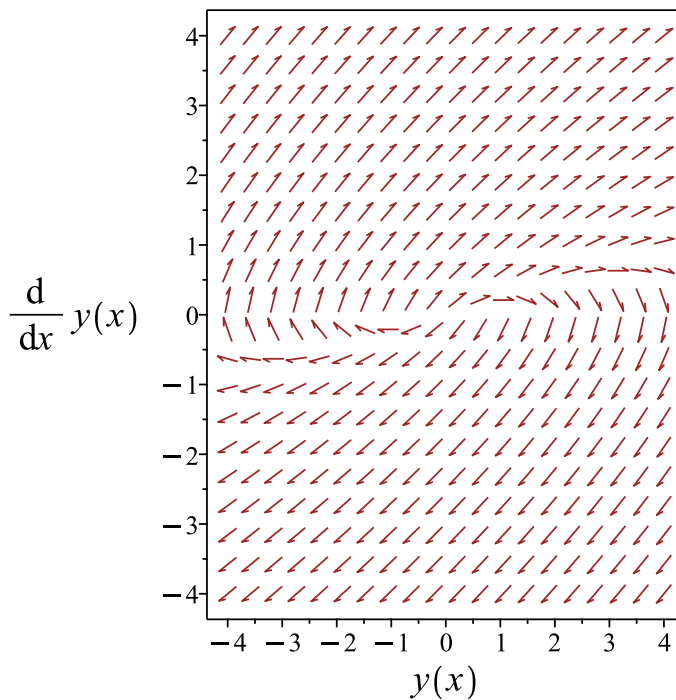


Figure 451: Slope field plot

Verification of solutions

$$y = c_1 e^{\frac{2x}{5}} + c_2 x e^{\frac{2x}{5}}$$

Verified OK.

9.9.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(x)y' + \frac{(p(x)^2 + p'(x))y}{2} = f(x)$$

Where $p(x) = -\frac{4}{5}$. Therefore, there is an integrating factor given by

$$\begin{aligned}M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int -\frac{4}{5} dx} \\ &= e^{-\frac{2x}{5}}\end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned}(M(x)y)'' &= 0 \\ \left(e^{-\frac{2x}{5}}y\right)'' &= 0\end{aligned}$$

Integrating once gives

$$\left(e^{-\frac{2x}{5}}y\right)' = c_1$$

Integrating again gives

$$\left(e^{-\frac{2x}{5}}y\right) = c_1x + c_2$$

Hence the solution is

$$y = \frac{c_1x + c_2}{e^{-\frac{2x}{5}}}$$

Or

$$y = c_1x e^{\frac{2x}{5}} + e^{\frac{2x}{5}} c_2$$

Summary

The solution(s) found are the following

$$y = c_1x e^{\frac{2x}{5}} + e^{\frac{2x}{5}} c_2 \quad (1)$$

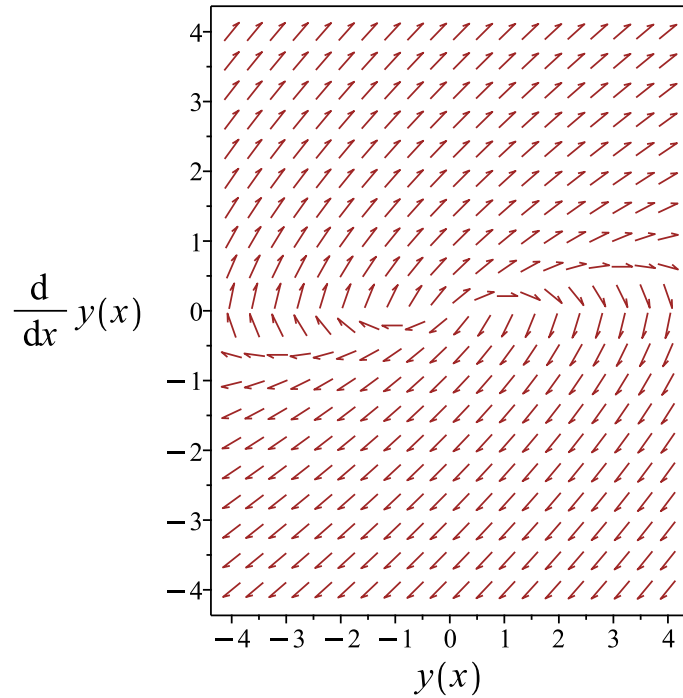


Figure 452: Slope field plot

Verification of solutions

$$y = c_1 x e^{\frac{2x}{5}} + e^{\frac{2x}{5}} c_2$$

Verified OK.

9.9.3 Solving using Kovacic algorithm

Writing the ode as

$$25y'' - 20y' + 4y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 25 \\ B &= -20 \\ C &= 4 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} \tag{5}$$

$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 0$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(x) = 0 \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 438: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-20}{25} dx} \\ &= z_1 e^{\frac{2x}{5}} \\ &= z_1 \left(e^{\frac{2x}{5}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{\frac{2x}{5}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-20}{25} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{\frac{4x}{5}}}{(y_1)^2} dx \\ &= y_1(x) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{\frac{2x}{5}} \right) + c_2 \left(e^{\frac{2x}{5}} (x) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{\frac{2x}{5}} + c_2 x e^{\frac{2x}{5}} \quad (1)$$

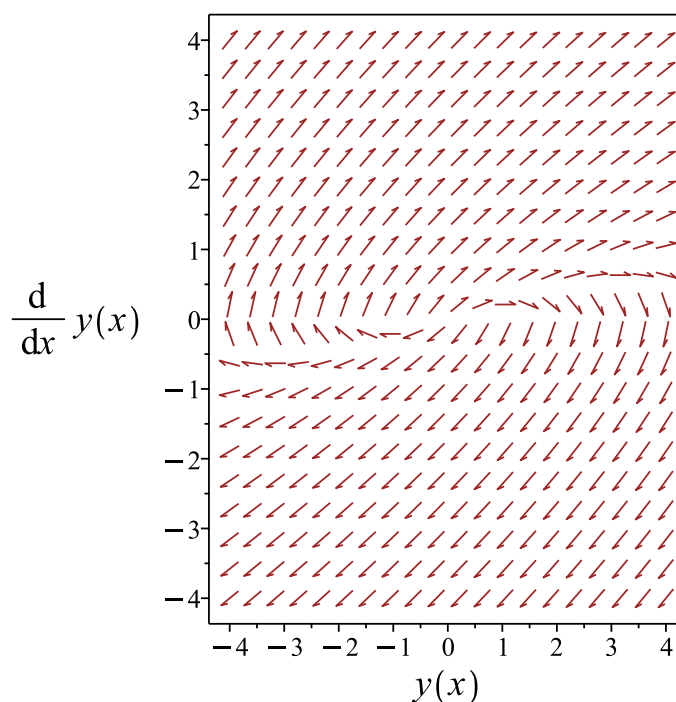


Figure 453: Slope field plot

Verification of solutions

$$y = c_1 e^{\frac{2x}{5}} + c_2 x e^{\frac{2x}{5}}$$

Verified OK.

9.9.4 Maple step by step solution

Let's solve

$$25y'' - 20y' + 4y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{4y'}{5} - \frac{4y}{25}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{4y'}{5} + \frac{4y}{25} = 0$$

- Characteristic polynomial of ODE

$$r^2 - \frac{4}{5}r + \frac{4}{25} = 0$$

- Factor the characteristic polynomial

$$\frac{(5r-2)^2}{25} = 0$$

- Root of the characteristic polynomial

$$r = \frac{2}{5}$$

- 1st solution of the ODE

$$y_1(x) = e^{\frac{2x}{5}}$$

- Repeated root, multiply $y_1(x)$ by x to ensure linear independence

$$y_2(x) = x e^{\frac{2x}{5}}$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{\frac{2x}{5}} + c_2 x e^{\frac{2x}{5}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(25*diff(y(x),x$2)-20*diff(y(x),x)+4*y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{\frac{2x}{5}}(c_2x + c_1)$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 20

```
DSolve[25*y''[x]-20*y'[x]+4*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{2x/5}(c_2x + c_1)$$

9.10 problem 10

9.10.1 Solving as second order linear constant coeff ode	2446
9.10.2 Solving using Kovacic algorithm	2448
9.10.3 Maple step by step solution	2452

Internal problem ID [662]

Internal file name [OUTPUT/662_Sunday_June_05_2022_01_46_34_AM_57523948/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 10.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$2y'' + 2y' + y = 0$$

9.10.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 2, B = 2, C = 1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$2\lambda^2 e^{\lambda x} + 2\lambda e^{\lambda x} + e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$2\lambda^2 + 2\lambda + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 2, B = 2, C = 1$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{-2}{(2)(2)} \pm \frac{1}{(2)(2)} \sqrt{2^2 - (4)(2)(1)} \\ &= -\frac{1}{2} \pm \frac{i}{2}\end{aligned}$$

Hence

$$\begin{aligned}\lambda_1 &= -\frac{1}{2} + \frac{i}{2} \\ \lambda_2 &= -\frac{1}{2} - \frac{i}{2}\end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= -\frac{1}{2} + \frac{i}{2} \\ \lambda_2 &= -\frac{1}{2} - \frac{i}{2}\end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -\frac{1}{2}$ and $\beta = \frac{1}{2}$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x} (c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^{-\frac{x}{2}} \left(c_1 \cos\left(\frac{x}{2}\right) + c_2 \sin\left(\frac{x}{2}\right) \right)$$

Summary

The solution(s) found are the following

$$y = e^{-\frac{x}{2}} \left(c_1 \cos\left(\frac{x}{2}\right) + c_2 \sin\left(\frac{x}{2}\right) \right) \quad (1)$$

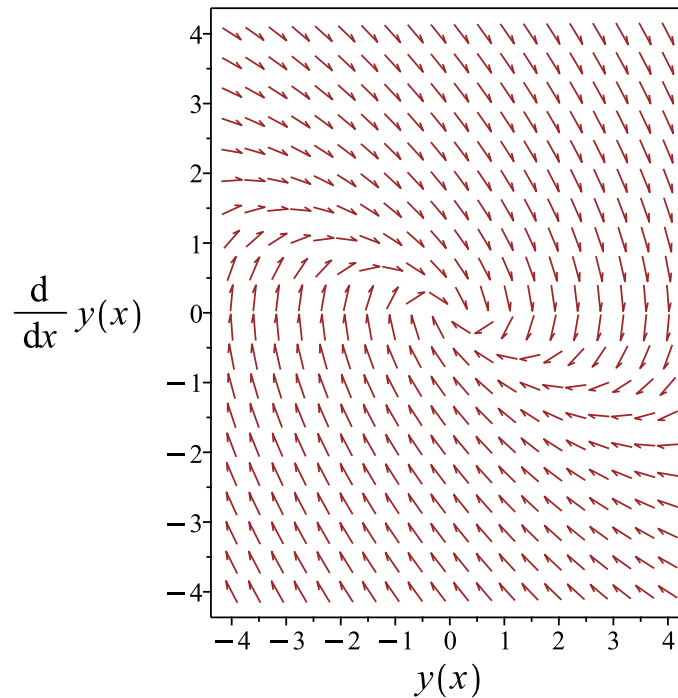


Figure 454: Slope field plot

Verification of solutions

$$y = e^{-\frac{x}{2}} \left(c_1 \cos \left(\frac{x}{2} \right) + c_2 \sin \left(\frac{x}{2} \right) \right)$$

Verified OK.

9.10.2 Solving using Kovacic algorithm

Writing the ode as

$$2y'' + 2y' + y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 2 \\ B &= 2 \\ C &= 1 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -\frac{z(x)}{4} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 440: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -\frac{1}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos\left(\frac{x}{2}\right)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} dx} \\ &= z_1 e^{-\int \frac{1}{2} dx} \\ &= z_1 e^{-\frac{x}{2}} \\ &= z_1 \left(e^{-\frac{x}{2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{x}{2}} \cos\left(\frac{x}{2}\right)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{2}{2} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-x}}{(y_1)^2} dx \\ &= y_1 \left(2 \tan\left(\frac{x}{2}\right) \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-\frac{x}{2}} \cos \left(\frac{x}{2} \right) \right) + c_2 \left(e^{-\frac{x}{2}} \cos \left(\frac{x}{2} \right) \left(2 \tan \left(\frac{x}{2} \right) \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-\frac{x}{2}} \cos \left(\frac{x}{2} \right) + 2c_2 e^{-\frac{x}{2}} \sin \left(\frac{x}{2} \right) \quad (1)$$

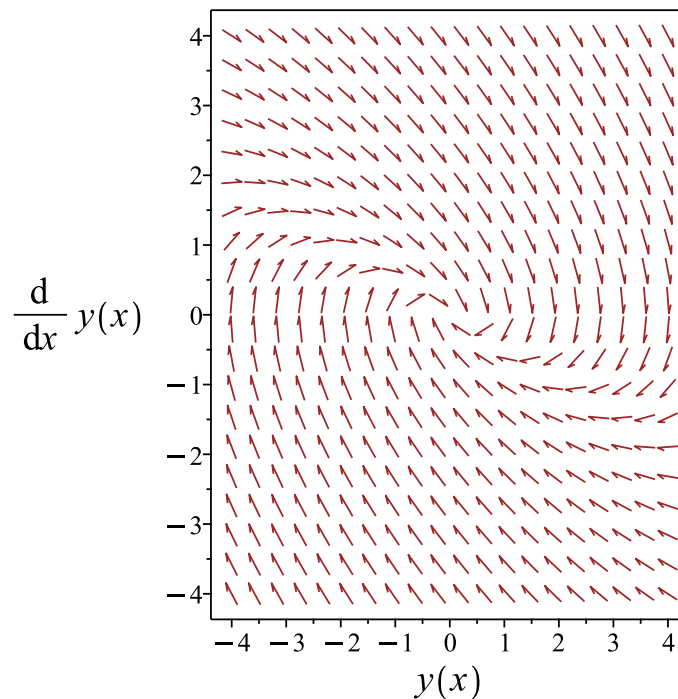


Figure 455: Slope field plot

Verification of solutions

$$y = c_1 e^{-\frac{x}{2}} \cos \left(\frac{x}{2} \right) + 2c_2 e^{-\frac{x}{2}} \sin \left(\frac{x}{2} \right)$$

Verified OK.

9.10.3 Maple step by step solution

Let's solve

$$2y'' + 2y' + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -y' - \frac{y}{2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + y' + \frac{y}{2} = 0$$

- Characteristic polynomial of ODE

$$r^2 + r + \frac{1}{2} = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-1) \pm (\sqrt{-1})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{2} - \frac{i}{2}, -\frac{1}{2} + \frac{i}{2}\right)$$

- 1st solution of the ODE

$$y_1(x) = e^{-\frac{x}{2}} \cos\left(\frac{x}{2}\right)$$

- 2nd solution of the ODE

$$y_2(x) = e^{-\frac{x}{2}} \sin\left(\frac{x}{2}\right)$$

- General solution of the ODE

$$y = c_1 y_1(x) + c_2 y_2(x)$$

- Substitute in solutions

$$y = c_1 e^{-\frac{x}{2}} \cos\left(\frac{x}{2}\right) + c_2 e^{-\frac{x}{2}} \sin\left(\frac{x}{2}\right)$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 22

```
dsolve(2*diff(y(x),x$2)+2*diff(y(x),x)+y(x) = 0,y(x), singsol=all)
```

$$y(x) = e^{-\frac{x}{2}} \left(c_1 \sin\left(\frac{x}{2}\right) + c_2 \cos\left(\frac{x}{2}\right) \right)$$

✓ Solution by Mathematica

Time used: 0.017 (sec). Leaf size: 32

```
DSolve[2*y''[x]+2*y'[x]+y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x/2} \left(c_2 \cos\left(\frac{x}{2}\right) + c_1 \sin\left(\frac{x}{2}\right) \right)$$

9.11 problem 11

9.11.1 Existence and uniqueness analysis	2455
9.11.2 Solving as second order linear constant coeff ode	2455
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9.11.5 Maple step by step solution	2464

Internal problem ID [663]

Internal file name [OUTPUT/663_Sunday_June_05_2022_01_46_35_AM_83642007/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 11.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$9y'' - 12y' + 4y = 0$$

With initial conditions

$$[y(0) = 2, y'(0) = -1]$$

9.11.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(t)y' + q(t)y = F$$

Where here

$$\begin{aligned}p(t) &= -\frac{4}{3} \\q(t) &= \frac{4}{9} \\F &= 0\end{aligned}$$

Hence the ode is

$$y'' - \frac{4y'}{3} + \frac{4y}{9} = 0$$

The domain of $p(t) = -\frac{4}{3}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = \frac{4}{9}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

9.11.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 9, B = -12, C = 4$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$9\lambda^2 e^{\lambda t} - 12\lambda e^{\lambda t} + 4e^{\lambda t} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$9\lambda^2 - 12\lambda + 4 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 9, B = -12, C = 4$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{12}{(2)(9)} \pm \frac{1}{(2)(9)} \sqrt{(-12)^2 - (4)(9)(4)} \\ &= \frac{2}{3}\end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = -\frac{2}{3}$. Therefore the solution is

$$y = c_1 e^{\frac{2t}{3}} + c_2 t e^{\frac{2t}{3}} \quad (1)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\frac{2t}{3}} + c_2 t e^{\frac{2t}{3}} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $t = 0$ in the above gives

$$2 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{2c_1 e^{\frac{2t}{3}}}{3} + c_2 e^{\frac{2t}{3}} + \frac{2c_2 t e^{\frac{2t}{3}}}{3}$$

substituting $y' = -1$ and $t = 0$ in the above gives

$$-1 = \frac{2c_1}{3} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}c_1 &= 2 \\ c_2 &= -\frac{7}{3}\end{aligned}$$

Substituting these values back in above solution results in

$$y = -\frac{7 e^{\frac{2t}{3}} t}{3} + 2 e^{\frac{2t}{3}}$$

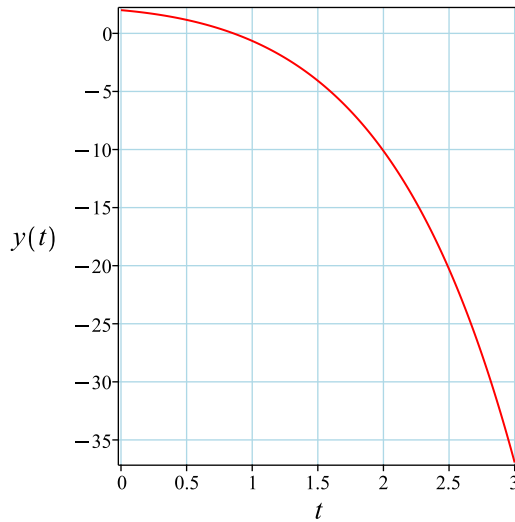
Which simplifies to

$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3} \right)$$

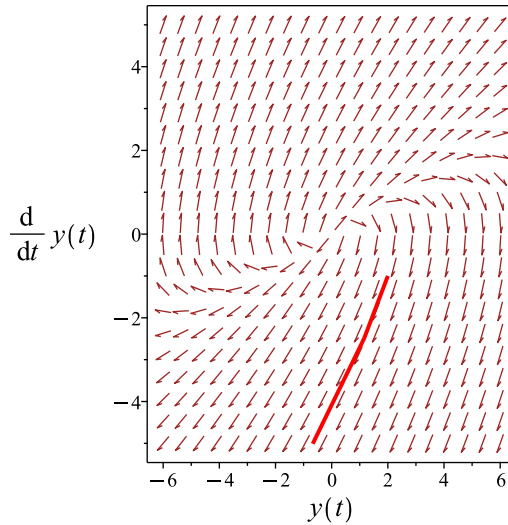
Summary

The solution(s) found are the following

$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3} \right) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3} \right)$$

Verified OK.

9.11.3 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t)y' + \frac{(p(t))^2 + p'(t)}{2}y = f(t)$$

Where $p(t) = -\frac{4}{3}$. Therefore, there is an integrating factor given by

$$\begin{aligned} M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int -\frac{4}{3} dx} \\ &= e^{-\frac{2t}{3}} \end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$(M(x)y)'' = 0$$
$$\left(e^{-\frac{2t}{3}}y\right)'' = 0$$

Integrating once gives

$$\left(e^{-\frac{2t}{3}}y\right)' = c_1$$

Integrating again gives

$$\left(e^{-\frac{2t}{3}}y\right) = c_1t + c_2$$

Hence the solution is

$$y = \frac{c_1t + c_2}{e^{-\frac{2t}{3}}}$$

Or

$$y = c_1t e^{\frac{2t}{3}} + c_2e^{\frac{2t}{3}}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1t e^{\frac{2t}{3}} + c_2e^{\frac{2t}{3}} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $t = 0$ in the above gives

$$2 = c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = c_1e^{\frac{2t}{3}} + \frac{2c_1t e^{\frac{2t}{3}}}{3} + \frac{2c_2e^{\frac{2t}{3}}}{3}$$

substituting $y' = -1$ and $t = 0$ in the above gives

$$-1 = c_1 + \frac{2c_2}{3} \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -\frac{7}{3}$$
$$c_2 = 2$$

Substituting these values back in above solution results in

$$y = -\frac{7e^{\frac{2t}{3}}t}{3} + 2e^{\frac{2t}{3}}$$

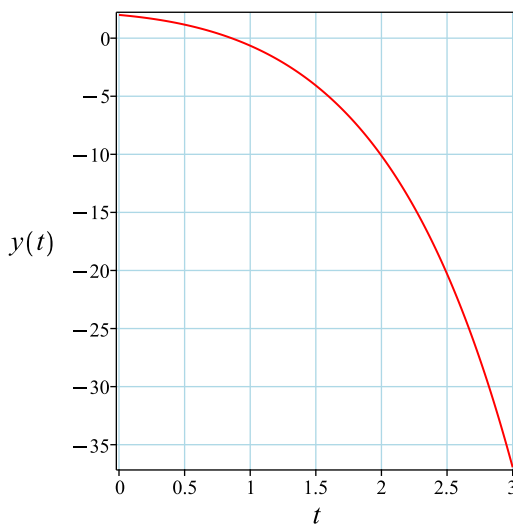
Which simplifies to

$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3} \right)$$

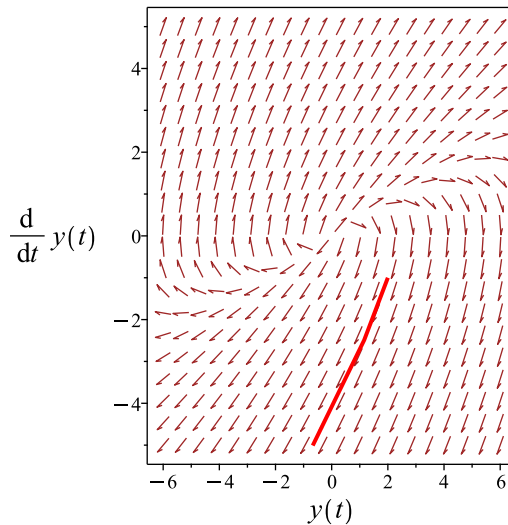
Summary

The solution(s) found are the following

$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3} \right) \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3} \right)$$

Verified OK.

9.11.4 Solving using Kovacic algorithm

Writing the ode as

$$9y'' - 12y' + 4y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 9 \\ B &= -12 \\ C &= 4 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 0 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = 0 \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 442: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-12}{9} dt} \end{aligned}$$

$$\begin{aligned}
&= z_1 e^{\frac{2t}{3}} \\
&= z_1 \left(e^{\frac{2t}{3}} \right)
\end{aligned}$$

Which simplifies to

$$y_1 = e^{\frac{2t}{3}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}
y_2 &= y_1 \int \frac{e^{\int -\frac{-12}{9} dt}}{(y_1)^2} dt \\
&= y_1 \int \frac{e^{\frac{4t}{3}}}{(y_1)^2} dt \\
&= y_1(t)
\end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 \left(e^{\frac{2t}{3}} \right) + c_2 \left(e^{\frac{2t}{3}}(t) \right)
\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\frac{2t}{3}} + c_2 t e^{\frac{2t}{3}} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $t = 0$ in the above gives

$$2 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = \frac{2c_1 e^{\frac{2t}{3}}}{3} + c_2 e^{\frac{2t}{3}} + \frac{2c_2 t e^{\frac{2t}{3}}}{3}$$

substituting $y' = -1$ and $t = 0$ in the above gives

$$-1 = \frac{2c_1}{3} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 2$$
$$c_2 = -\frac{7}{3}$$

Substituting these values back in above solution results in

$$y = -\frac{7e^{\frac{2t}{3}}}{3} + 2e^{\frac{2t}{3}}$$

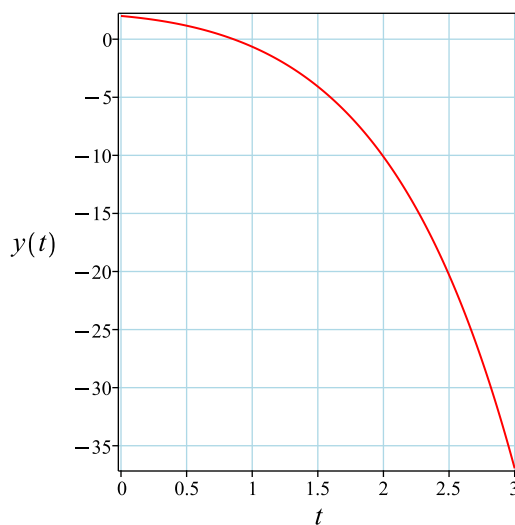
Which simplifies to

$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3} \right)$$

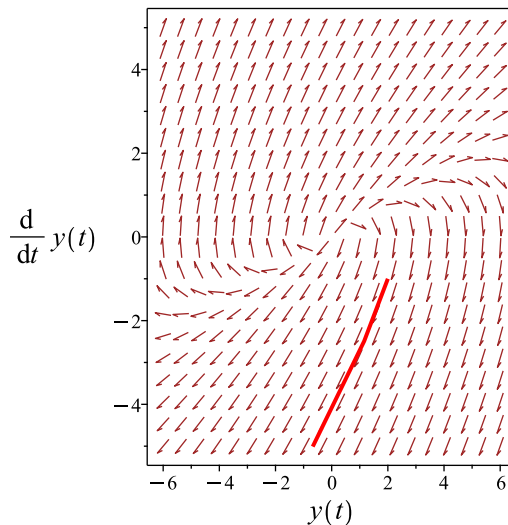
Summary

The solution(s) found are the following

$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3} \right) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3} \right)$$

Verified OK.

9.11.5 Maple step by step solution

Let's solve

$$\left[9y'' - 12y' + 4y = 0, y(0) = 2, y'|_{\{t=0\}} = -1 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{4y'}{3} - \frac{4y}{9}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{4y'}{3} + \frac{4y}{9} = 0$$

- Characteristic polynomial of ODE

$$r^2 - \frac{4}{3}r + \frac{4}{9} = 0$$

- Factor the characteristic polynomial

$$\frac{(3r-2)^2}{9} = 0$$

- Root of the characteristic polynomial

$$r = \frac{2}{3}$$

- 1st solution of the ODE

$$y_1(t) = e^{\frac{2t}{3}}$$

- Repeated root, multiply $y_1(t)$ by t to ensure linear independence

$$y_2(t) = e^{\frac{2t}{3}} t$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t)$$

- Substitute in solutions

$$y = c_1 e^{\frac{2t}{3}} + c_2 t e^{\frac{2t}{3}}$$

- Check validity of solution $y = c_1 e^{\frac{2t}{3}} + c_2 t e^{\frac{2t}{3}}$

- Use initial condition $y(0) = 2$

$$2 = c_1$$
- Compute derivative of the solution
$$y' = \frac{2c_1 e^{\frac{2t}{3}}}{3} + c_2 e^{\frac{2t}{3}} + \frac{2c_2 t e^{\frac{2t}{3}}}{3}$$
- Use the initial condition $y' \Big|_{\{t=0\}} = -1$

$$-1 = \frac{2c_1}{3} + c_2$$
- Solve for c_1 and c_2

$$\{c_1 = 2, c_2 = -\frac{7}{3}\}$$
- Substitute constant values into general solution and simplify
$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3}\right)$$
- Solution to the IVP
$$y = e^{\frac{2t}{3}} \left(2 - \frac{7t}{3}\right)$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 14

```
dsolve([9*dif(y(t),t$2)-12*dif(y(t),t)+4*y(t) = 0,y(0) = 2, D(y)(0) = -1],y(t), singsol=all)
```

$$y(t) = -\frac{e^{\frac{2t}{3}}(-6 + 7t)}{3}$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 15

```
DSolve[{9*y''[t]-12*y'[t]+4*y[t]==0,{y[0]==0,y'[0]==-1}},y[t],t,IncludeSingularSolutions ->
```

$$y(t) \rightarrow -e^{2t/3}t$$

9.12 problem 12

9.12.1 Existence and uniqueness analysis	2468
9.12.2 Solving as second order linear constant coeff ode	2468
9.12.3 Solving as linear second order ode solved by an integrating factor ode	2470
9.12.4 Solving using Kovacic algorithm	2472
9.12.5 Maple step by step solution	2476

Internal problem ID [664]

Internal file name [OUTPUT/664_Sunday_June_05_2022_01_46_36_AM_49512898/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 12.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$y'' - 6y' + 9y = 0$$

With initial conditions

$$[y(0) = 0, y'(0) = 2]$$

9.12.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(t)y' + q(t)y = F$$

Where here

$$p(t) = -6$$

$$q(t) = 9$$

$$F = 0$$

Hence the ode is

$$y'' - 6y' + 9y = 0$$

The domain of $p(t) = -6$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = 9$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

9.12.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = -6, C = 9$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} - 6\lambda e^{\lambda t} + 9e^{\lambda t} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 - 6\lambda + 9 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -6, C = 9$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{6}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(-6)^2 - (4)(1)(9)} \\ &= 3\end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = -3$. Therefore the solution is

$$y = c_1 e^{3t} + c_2 e^{3t} t \quad (1)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{3t} + c_2 e^{3t} t \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $t = 0$ in the above gives

$$0 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$y' = 3c_1 e^{3t} + 3c_2 e^{3t} t + c_2 e^{3t}$$

substituting $y' = 2$ and $t = 0$ in the above gives

$$2 = 3c_1 + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 0$$

$$c_2 = 2$$

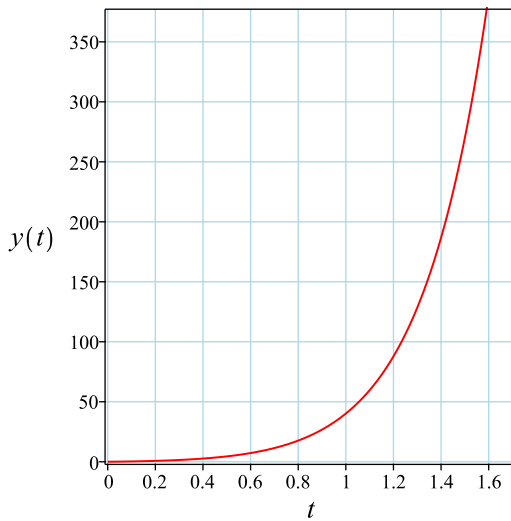
Substituting these values back in above solution results in

$$y = 2 e^{3t} t$$

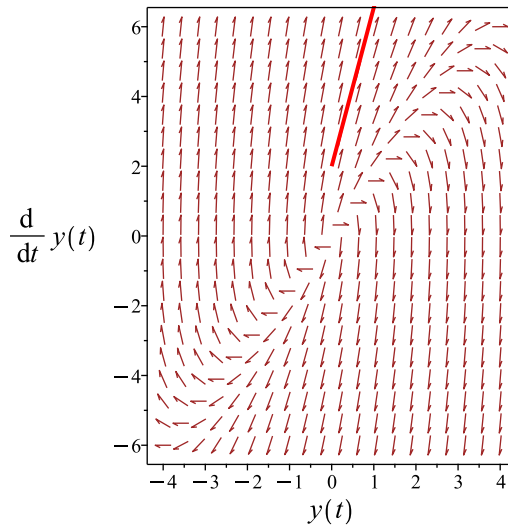
Summary

The solution(s) found are the following

$$y = 2 e^{3t} t \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = 2e^{3t}$$

Verified OK.

9.12.3 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t)y' + \frac{(p(t)^2 + p'(t))y}{2} = f(t)$$

Where $p(t) = -6$. Therefore, there is an integrating factor given by

$$\begin{aligned} M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int -6 dx} \\ &= e^{-3t} \end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned} (M(x)y)'' &= 0 \\ (e^{-3t}y)'' &= 0 \end{aligned}$$

Integrating once gives

$$(e^{-3t}y)' = c_1$$

Integrating again gives

$$(e^{-3t}y) = c_1t + c_2$$

Hence the solution is

$$y = \frac{c_1t + c_2}{e^{-3t}}$$

Or

$$y = c_1te^{3t} + c_2e^{3t}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1te^{3t} + c_2e^{3t} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $t = 0$ in the above gives

$$0 = c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = c_1e^{3t} + 3c_1te^{3t} + 3c_2e^{3t}$$

substituting $y' = 2$ and $t = 0$ in the above gives

$$2 = c_1 + 3c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 2$$

$$c_2 = 0$$

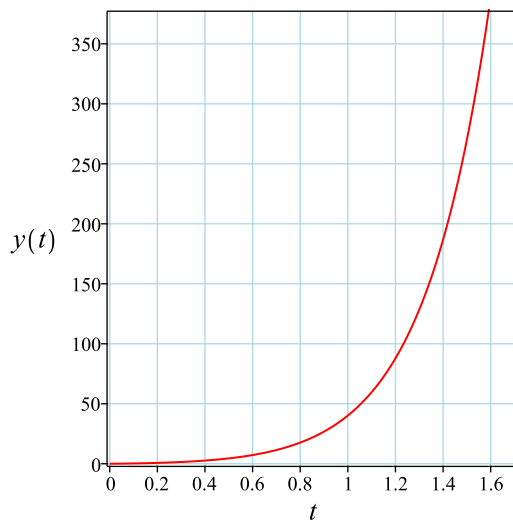
Substituting these values back in above solution results in

$$y = 2e^{3t}t$$

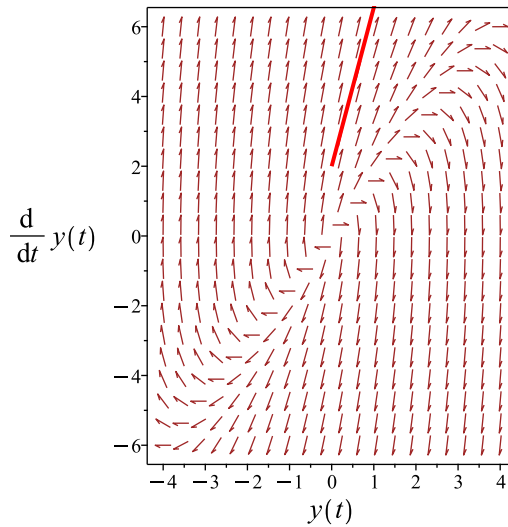
Summary

The solution(s) found are the following

$$y = 2e^{3t}t \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = 2e^{3t}$$

Verified OK.

9.12.4 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 6y' + 9y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = -6 \tag{3}$$

$$C = 9$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} \tag{5}$$

$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 0$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(t) = 0 \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 444: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-6}{1} dt} \\ &= z_1 e^{3t} \\ &= z_1 (e^{3t}) \end{aligned}$$

Which simplifies to

$$y_1 = e^{3t}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-6}{1} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{6t}}{(y_1)^2} dt \\ &= y_1(t) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{3t}) + c_2 (e^{3t} t)\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{3t} + c_2 e^{3t} t \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 0$ and $t = 0$ in the above gives

$$0 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$y' = 3c_1 e^{3t} + 3c_2 e^{3t} t + c_2 e^{3t}$$

substituting $y' = 2$ and $t = 0$ in the above gives

$$2 = 3c_1 + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 0$$

$$c_2 = 2$$

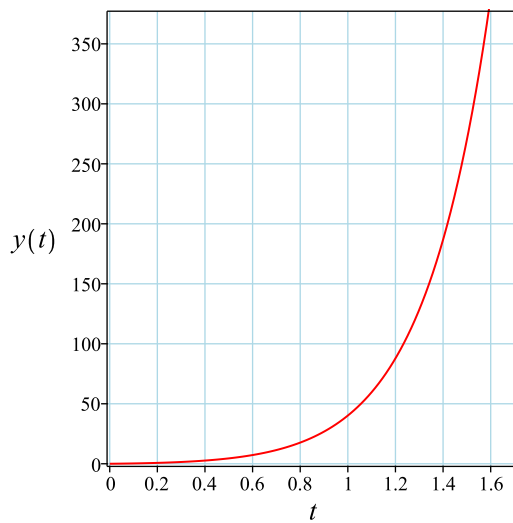
Substituting these values back in above solution results in

$$y = 2 e^{3t} t$$

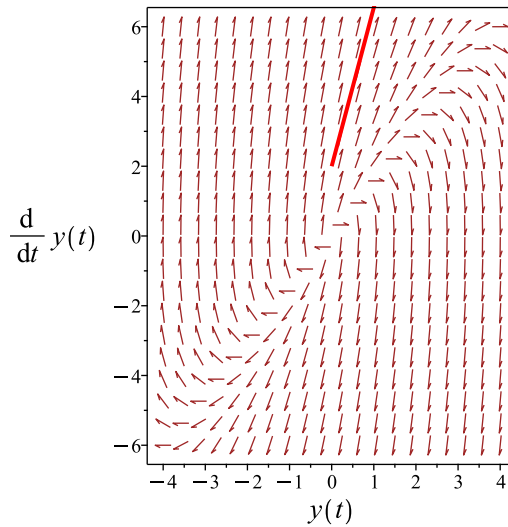
Summary

The solution(s) found are the following

$$y = 2 e^{3t} t \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = 2e^{3t}$$

Verified OK.

9.12.5 Maple step by step solution

Let's solve

$$\left[y'' - 6y' + 9y = 0, y(0) = 0, y' \Big|_{\{t=0\}} = 2 \right]$$

- Highest derivative means the order of the ODE is 2
 y''
- Characteristic polynomial of ODE
 $r^2 - 6r + 9 = 0$
- Factor the characteristic polynomial
 $(r - 3)^2 = 0$
- Root of the characteristic polynomial
 $r = 3$
- 1st solution of the ODE

$$y_1(t) = e^{3t}$$

- Repeated root, multiply $y_1(t)$ by t to ensure linear independence

$$y_2(t) = e^{3t}t$$

- General solution of the ODE

$$y = c_1y_1(t) + c_2y_2(t)$$

- Substitute in solutions

$$y = c_1e^{3t} + c_2e^{3t}t$$

- Check validity of solution $y = c_1e^{3t} + c_2e^{3t}t$

- Use initial condition $y(0) = 0$

$$0 = c_1$$

- Compute derivative of the solution

$$y' = 3c_1e^{3t} + 3c_2e^{3t}t + c_2e^{3t}$$

- Use the initial condition $y' \Big|_{\{t=0\}} = 2$

$$2 = 3c_1 + c_2$$

- Solve for c_1 and c_2

$$\{c_1 = 0, c_2 = 2\}$$

- Substitute constant values into general solution and simplify

$$y = 2e^{3t}t$$

- Solution to the IVP

$$y = 2e^{3t}t$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 11

```
dsolve([diff(y(t),t$2)-6*diff(y(t),t)+9*y(t) = 0,y(0) = 0, D(y)(0) = 2],y(t), singsol=all)
```

$$y(t) = 2e^{3t}$$

✓ Solution by Mathematica

Time used: 0.014 (sec). Leaf size: 13

```
DSolve[{y''[t]-6*y'[t]+9*y[t]==0,{y[0]==0,y'[0]==2}},y[t],t,IncludeSingularSolutions -> True
```

$$y(t) \rightarrow 2e^{3t}$$

9.13 problem 13

9.13.1 Existence and uniqueness analysis	2479
9.13.2 Solving as second order linear constant coeff ode	2480
9.13.3 Solving using Kovacic algorithm	2482
9.13.4 Maple step by step solution	2487

Internal problem ID [665]

Internal file name [OUTPUT/665_Sunday_June_05_2022_01_46_37_AM_75605177/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 13.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$9y'' + 6y' + 82y = 0$$

With initial conditions

$$[y(0) = -1, y'(0) = 2]$$

9.13.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(t)y' + q(t)y = F$$

Where here

$$p(t) = \frac{2}{3}$$
$$q(t) = \frac{82}{9}$$
$$F = 0$$

Hence the ode is

$$y'' + \frac{2y'}{3} + \frac{82y}{9} = 0$$

The domain of $p(t) = \frac{2}{3}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = \frac{82}{9}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

9.13.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 9, B = 6, C = 82$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$9\lambda^2 e^{\lambda t} + 6\lambda e^{\lambda t} + 82 e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$9\lambda^2 + 6\lambda + 82 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 9, B = 6, C = 82$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-6}{(2)(9)} \pm \frac{1}{(2)(9)} \sqrt{6^2 - (4)(9)(82)} \\ &= -\frac{1}{3} \pm 3i \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= -\frac{1}{3} + 3i \\ \lambda_2 &= -\frac{1}{3} - 3i \end{aligned}$$

Which simplifies to

$$\lambda_1 = -\frac{1}{3} + 3i$$

$$\lambda_2 = -\frac{1}{3} - 3i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -\frac{1}{3}$ and $\beta = 3$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha t}(c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$y = e^{-\frac{t}{3}}(c_1 \cos(3t) + c_2 \sin(3t))$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = e^{-\frac{t}{3}}(c_1 \cos(3t) + c_2 \sin(3t)) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = -1$ and $t = 0$ in the above gives

$$-1 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$y' = -\frac{e^{-\frac{t}{3}}(c_1 \cos(3t) + c_2 \sin(3t))}{3} + e^{-\frac{t}{3}}(-3c_1 \sin(3t) + 3c_2 \cos(3t))$$

substituting $y' = 2$ and $t = 0$ in the above gives

$$2 = -\frac{c_1}{3} + 3c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -1$$

$$c_2 = \frac{5}{9}$$

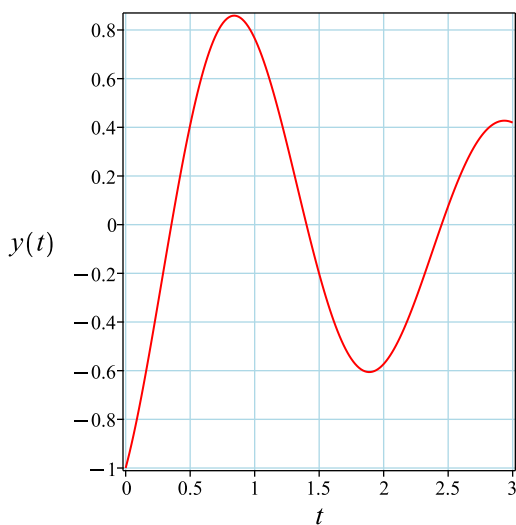
Substituting these values back in above solution results in

$$y = -\frac{e^{-\frac{t}{3}}(9 \cos(3t) - 5 \sin(3t))}{9}$$

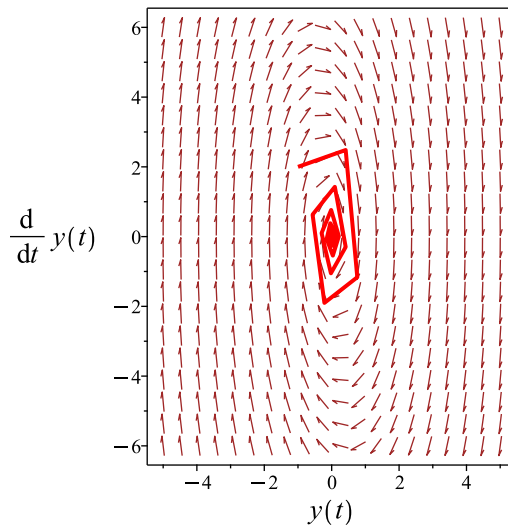
Summary

The solution(s) found are the following

$$y = -\frac{e^{-\frac{t}{3}}(9 \cos(3t) - 5 \sin(3t))}{9} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{e^{-\frac{t}{3}}(9 \cos(3t) - 5 \sin(3t))}{9}$$

Verified OK.

9.13.3 Solving using Kovacic algorithm

Writing the ode as

$$9y'' + 6y' + 82y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= 9 \\B &= 6 \\C &= 82\end{aligned}\tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t)\tag{4}$$

Where r is given by

$$\begin{aligned}r &= \frac{s}{t} \\&= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}\end{aligned}\tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-9}{1}\tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned}s &= -9 \\t &= 1\end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = -9z(t)\tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 446: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -9$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos(3t)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}
 y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\
 &= z_1 e^{-\int \frac{1}{2} \frac{6}{9} dt} \\
 &= z_1 e^{-\frac{t}{3}} \\
 &= z_1 \left(e^{-\frac{t}{3}} \right)
 \end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{t}{3}} \cos(3t)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{6}{9} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{-\frac{2t}{3}}}{(y_1)^2} dt \\ &= y_1 \left(\frac{\tan(3t)}{3} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-\frac{t}{3}} \cos(3t) \right) + c_2 \left(e^{-\frac{t}{3}} \cos(3t) \left(\frac{\tan(3t)}{3} \right) \right) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-\frac{t}{3}} \cos(3t) + \frac{c_2 e^{-\frac{t}{3}} \sin(3t)}{3} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = -1$ and $t = 0$ in the above gives

$$-1 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$y' = -\frac{c_1 e^{-\frac{t}{3}} \cos(3t)}{3} - 3c_1 e^{-\frac{t}{3}} \sin(3t) - \frac{c_2 e^{-\frac{t}{3}} \sin(3t)}{9} + c_2 e^{-\frac{t}{3}} \cos(3t)$$

substituting $y' = 2$ and $t = 0$ in the above gives

$$2 = -\frac{c_1}{3} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -1$$
$$c_2 = \frac{5}{3}$$

Substituting these values back in above solution results in

$$y = -e^{-\frac{t}{3}} \cos(3t) + \frac{5 e^{-\frac{t}{3}} \sin(3t)}{9}$$

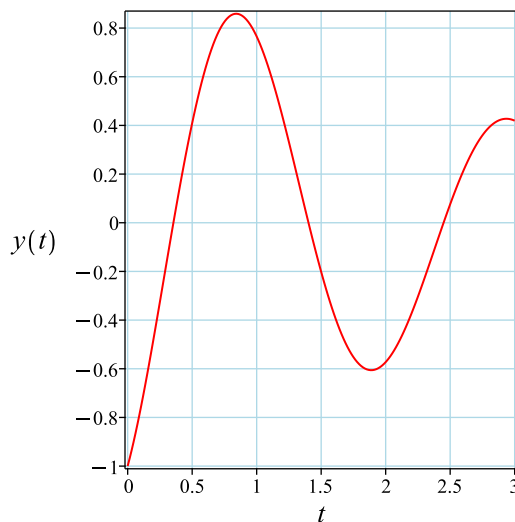
Which simplifies to

$$y = -\frac{e^{-\frac{t}{3}}(9 \cos(3t) - 5 \sin(3t))}{9}$$

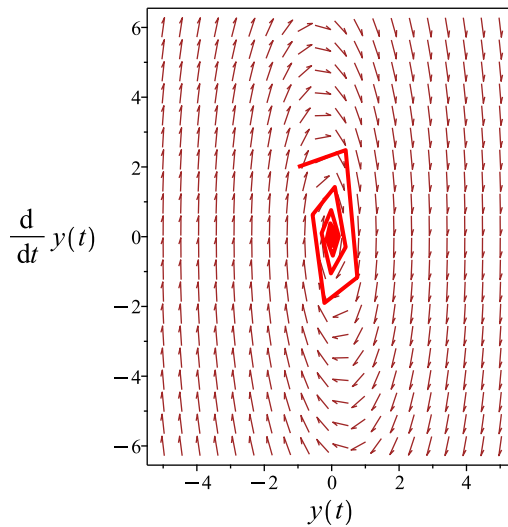
Summary

The solution(s) found are the following

$$y = -\frac{e^{-\frac{t}{3}}(9 \cos(3t) - 5 \sin(3t))}{9} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = -\frac{e^{-\frac{t}{3}}(9 \cos(3t) - 5 \sin(3t))}{9}$$

Verified OK.

9.13.4 Maple step by step solution

Let's solve

$$\left[9y'' + 6y' + 82y = 0, y(0) = -1, y' \Big|_{\{t=0\}} = 2 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{2y'}{3} - \frac{82y}{9}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{2y'}{3} + \frac{82y}{9} = 0$$

- Characteristic polynomial of ODE

$$r^2 + \frac{2}{3}r + \frac{82}{9} = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-\frac{2}{3}) \pm (\sqrt{-36})}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{3} - 3I, -\frac{1}{3} + 3I\right)$$

- 1st solution of the ODE

$$y_1(t) = e^{-\frac{t}{3}} \cos(3t)$$

- 2nd solution of the ODE

$$y_2(t) = e^{-\frac{t}{3}} \sin(3t)$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t)$$

- Substitute in solutions

$$y = c_1 e^{-\frac{t}{3}} \cos(3t) + c_2 e^{-\frac{t}{3}} \sin(3t)$$

- Check validity of solution $y = c_1 e^{-\frac{t}{3}} \cos(3t) + c_2 e^{-\frac{t}{3}} \sin(3t)$

- Use initial condition $y(0) = -1$
 $-1 = c_1$
- Compute derivative of the solution

$$y' = -\frac{c_1 e^{-\frac{t}{3}} \cos(3t)}{3} - 3c_1 e^{-\frac{t}{3}} \sin(3t) - \frac{c_2 e^{-\frac{t}{3}} \sin(3t)}{3} + 3c_2 e^{-\frac{t}{3}} \cos(3t)$$
- Use the initial condition $y' \Big|_{\{t=0\}} = 2$
 $2 = -\frac{c_1}{3} + 3c_2$
- Solve for c_1 and c_2
 $\{c_1 = -1, c_2 = \frac{5}{9}\}$
- Substitute constant values into general solution and simplify

$$y = -\frac{e^{-\frac{t}{3}}(9 \cos(3t) - 5 \sin(3t))}{9}$$
- Solution to the IVP

$$y = -\frac{e^{-\frac{t}{3}}(9 \cos(3t) - 5 \sin(3t))}{9}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 23

```
dsolve([9*dif(y(t),t$2)+6*dif(y(t),t)+82*y(t) = 0,y(0) = -1, D(y)(0) = 2],y(t), singsol=al
```

$$y(t) = \frac{e^{-\frac{t}{3}}(5 \sin(3t) - 9 \cos(3t))}{9}$$

✓ Solution by Mathematica

Time used: 0.019 (sec). Leaf size: 29

```
DSolve[{9*y''[t]+6*y'[t]+82*y[t]==0,{y[0]==-1,y'[0]==2}},y[t],t,IncludeSingularSolutions ->
```

$$y(t) \rightarrow \frac{1}{9}e^{-t/3}(5 \sin(3t) - 9 \cos(3t))$$

9.14 problem 14

9.14.1 Existence and uniqueness analysis	2491
9.14.2 Solving as second order linear constant coeff ode	2491
9.14.3 Solving as linear second order ode solved by an integrating factor ode	2493
9.14.4 Solving using Kovacic algorithm	2495
9.14.5 Maple step by step solution	2499

Internal problem ID [666]

Internal file name [OUTPUT/666_Sunday_June_05_2022_01_46_38_AM_29132401/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 14.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$y'' + 4y' + 4y = 0$$

With initial conditions

$$[y(-1) = 2, y'(-1) = 1]$$

9.14.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 4$$

$$q(x) = 4$$

$$F = 0$$

Hence the ode is

$$y'' + 4y' + 4y = 0$$

The domain of $p(x) = 4$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = -1$ is inside this domain. The domain of $q(x) = 4$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = -1$ is also inside this domain. Hence solution exists and is unique.

9.14.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 4, C = 4$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + 4\lambda e^{\lambda x} + 4e^{\lambda x} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 4\lambda + 4 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 4, C = 4$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{-4}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(4)^2 - (4)(1)(4)} \\ &= -2\end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = 2$. Therefore the solution is

$$y = c_1 e^{-2x} + c_2 e^{-2x} x \quad (1)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-2x} + x e^{-2x} c_2 \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = -1$ in the above gives

$$2 = e^2(c_1 - c_2) \quad (1A)$$

Taking derivative of the solution gives

$$y' = -2c_1 e^{-2x} + c_2 e^{-2x} - 2x e^{-2x} c_2$$

substituting $y' = 1$ and $x = -1$ in the above gives

$$1 = e^2(-2c_1 + 3c_2) \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}c_1 &= 7 e^{-2} \\ c_2 &= 5 e^{-2}\end{aligned}$$

Substituting these values back in above solution results in

$$y = 5x e^{-2x} e^{-2} + 7 e^{-2x} e^{-2}$$

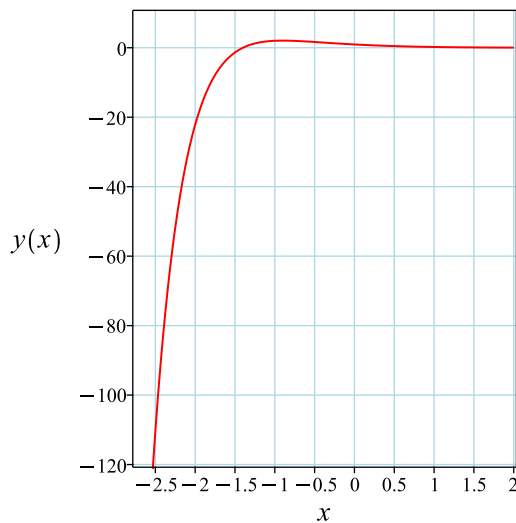
Which simplifies to

$$y = e^{-2x-2}(7 + 5x)$$

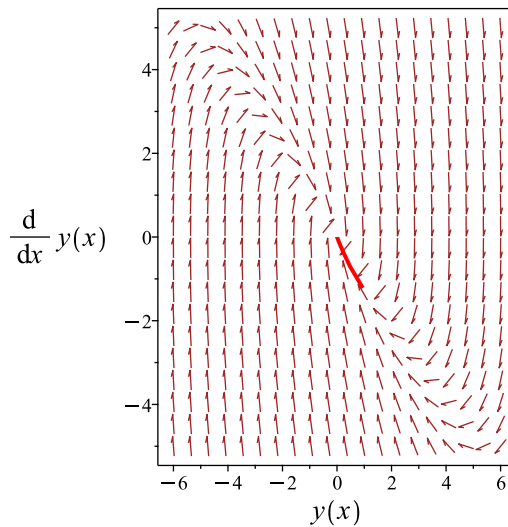
Summary

The solution(s) found are the following

$$y = e^{-2x-2}(7 + 5x) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{-2x-2}(7 + 5x)$$

Verified OK.

9.14.3 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(x)y' + \frac{(p(x)^2 + p'(x))y}{2} = f(x)$$

Where $p(x) = 4$. Therefore, there is an integrating factor given by

$$\begin{aligned} M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int 4 dx} \\ &= e^{2x} \end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned} (M(x)y)'' &= 0 \\ (e^{2x}y)'' &= 0 \end{aligned}$$

Integrating once gives

$$(e^{2x}y)' = c_1$$

Integrating again gives

$$(e^{2x}y) = c_1x + c_2$$

Hence the solution is

$$y = \frac{c_1x + c_2}{e^{2x}}$$

Or

$$y = c_1x e^{-2x} + c_2e^{-2x}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1x e^{-2x} + c_2e^{-2x} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = -1$ in the above gives

$$2 = (-c_1 + c_2) e^2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = c_1e^{-2x} - 2c_1x e^{-2x} - 2c_2e^{-2x}$$

substituting $y' = 1$ and $x = -1$ in the above gives

$$1 = (3c_1 - 2c_2) e^2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 5e^{-2}$$

$$c_2 = 7e^{-2}$$

Substituting these values back in above solution results in

$$y = 5x e^{-2x-2} + 7e^{-2x}e^{-2}$$

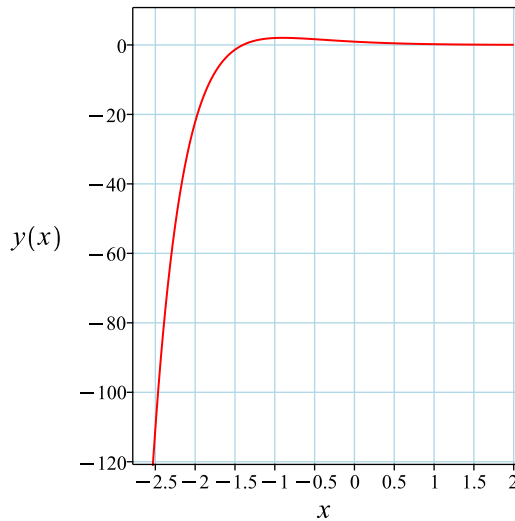
Which simplifies to

$$y = e^{-2x-2}(7 + 5x)$$

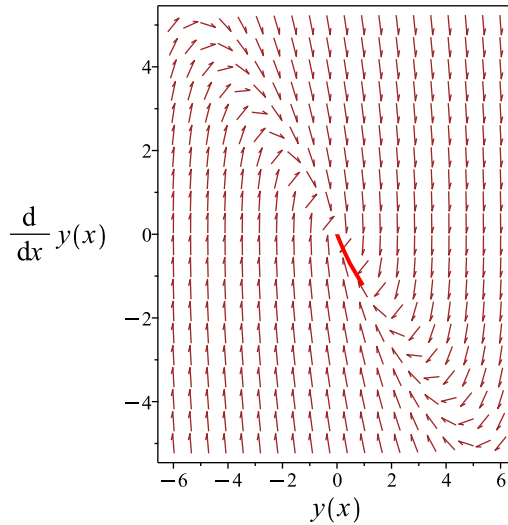
Summary

The solution(s) found are the following

$$y = e^{-2x-2}(7 + 5x) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{-2x-2}(7 + 5x)$$

Verified OK.

9.14.4 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 4y' + 4y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 4 \quad (3)$$

$$C = 4$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 0 \\ t = 1$$

Therefore eq. (4) becomes

$$z''(x) = 0 \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 448: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{4}{1} dx} \\ &= z_1 e^{-2x} \\ &= z_1 (e^{-2x})\end{aligned}$$

Which simplifies to

$$y_1 = e^{-2x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{4}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-4x}}{(y_1)^2} dx \\ &= y_1(x)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-2x}) + c_2 (e^{-2x}(x))\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-2x} + x e^{-2x} c_2 \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $x = -1$ in the above gives

$$2 = e^2(c_1 - c_2) \quad (1A)$$

Taking derivative of the solution gives

$$y' = -2c_1 e^{-2x} + c_2 e^{-2x} - 2x e^{-2x} c_2$$

substituting $y' = 1$ and $x = -1$ in the above gives

$$1 = e^2(-2c_1 + 3c_2) \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}c_1 &= 7 e^{-2} \\ c_2 &= 5 e^{-2}\end{aligned}$$

Substituting these values back in above solution results in

$$y = 5x e^{-2x} e^{-2} + 7 e^{-2x} e^{-2}$$

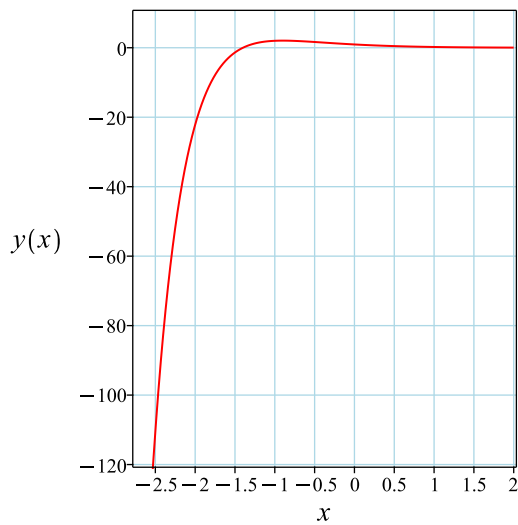
Which simplifies to

$$y = e^{-2x-2}(7 + 5x)$$

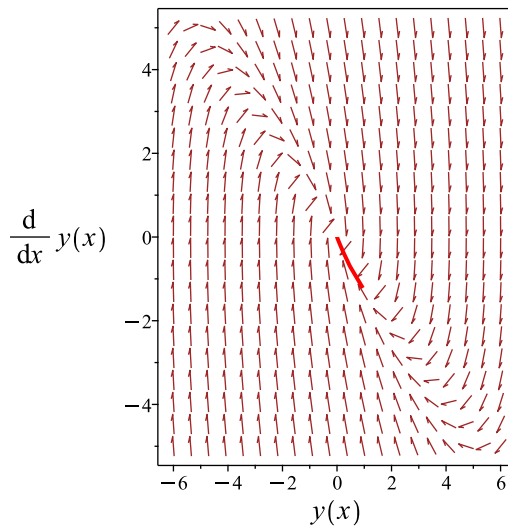
Summary

The solution(s) found are the following

$$y = e^{-2x-2}(7 + 5x) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{-2x-2}(7 + 5x)$$

Verified OK.

9.14.5 Maple step by step solution

Let's solve

$$\left[y'' + 4y' + 4y = 0, y(-1) = 2, y'|_{\{x=-1\}} = 1 \right]$$

- Highest derivative means the order of the ODE is 2
 y''
- Characteristic polynomial of ODE
 $r^2 + 4r + 4 = 0$
- Factor the characteristic polynomial
 $(r + 2)^2 = 0$
- Root of the characteristic polynomial
 $r = -2$
- 1st solution of the ODE

$$y_1(x) = e^{-2x}$$

- Repeated root, multiply $y_1(x)$ by x to ensure linear independence

$$y_2(x) = e^{-2x}x$$

- General solution of the ODE

$$y = c_1y_1(x) + c_2y_2(x)$$

- Substitute in solutions

$$y = c_1e^{-2x} + xe^{-2x}c_2$$

- Check validity of solution $y = c_1e^{-2x} + xe^{-2x}c_2$

- Use initial condition $y(-1) = 2$

$$2 = e^2c_1 - e^2c_2$$

- Compute derivative of the solution

$$y' = -2c_1e^{-2x} + c_2e^{-2x} - 2xe^{-2x}c_2$$

- Use the initial condition $y' \Big|_{\{x=-1\}} = 1$

$$1 = -2e^2c_1 + 3e^2c_2$$

- Solve for c_1 and c_2

$$\left\{ c_1 = \frac{7}{e^2}, c_2 = \frac{5}{e^2} \right\}$$

- Substitute constant values into general solution and simplify

$$y = e^{-2x-2}(7 + 5x)$$

- Solution to the IVP

$$y = e^{-2x-2}(7 + 5x)$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 16

```
dsolve([diff(y(x),x$2)+4*diff(y(x),x)+4*y(x) = 0,y(-1) = 2, D(y)(-1) = 1],y(x), singsol=all)
```

$$y(x) = e^{-2x-2}(5x + 7)$$

✓ Solution by Mathematica

Time used: 0.014 (sec). Leaf size: 18

```
DSolve[{y''[x]+4*y'[x]+4*y[x]==0,{y[-1]==2,y'[-1]==1}},y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-2(x+1)}(5x + 7)$$

9.15 problem 15

9.15.1 Existence and uniqueness analysis	2503
9.15.2 Solving as second order linear constant coeff ode	2503
9.15.3 Solving as linear second order ode solved by an integrating factor ode	2505
9.15.4 Solving using Kovacic algorithm	2508
9.15.5 Maple step by step solution	2512

Internal problem ID [667]

Internal file name [OUTPUT/667_Sunday_June_05_2022_01_46_39_AM_14025215/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 15.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$4y'' + 12y' + 9y = 0$$

With initial conditions

$$[y(0) = 1, y'(0) = -4]$$

9.15.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(t)y' + q(t)y = F$$

Where here

$$p(t) = 3$$

$$q(t) = \frac{9}{4}$$

$$F = 0$$

Hence the ode is

$$y'' + 3y' + \frac{9y}{4} = 0$$

The domain of $p(t) = 3$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = \frac{9}{4}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

9.15.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 4, B = 12, C = 9$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$4\lambda^2 e^{\lambda t} + 12\lambda e^{\lambda t} + 9e^{\lambda t} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$4\lambda^2 + 12\lambda + 9 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 4, B = 12, C = 9$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{-12}{(2)(4)} \pm \frac{1}{(2)(4)} \sqrt{(12)^2 - (4)(4)(9)} \\ &= -\frac{3}{2}\end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = \frac{3}{2}$. Therefore the solution is

$$y = c_1 e^{-\frac{3t}{2}} + c_2 t e^{-\frac{3t}{2}} \quad (1)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-\frac{3t}{2}} + c_2 t e^{-\frac{3t}{2}} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $t = 0$ in the above gives

$$1 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$y' = -\frac{3c_1 e^{-\frac{3t}{2}}}{2} + c_2 e^{-\frac{3t}{2}} - \frac{3c_2 t e^{-\frac{3t}{2}}}{2}$$

substituting $y' = -4$ and $t = 0$ in the above gives

$$-4 = -\frac{3c_1}{2} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}c_1 &= 1 \\ c_2 &= -\frac{5}{2}\end{aligned}$$

Substituting these values back in above solution results in

$$y = -\frac{5 e^{-\frac{3t}{2}} t}{2} + e^{-\frac{3t}{2}}$$

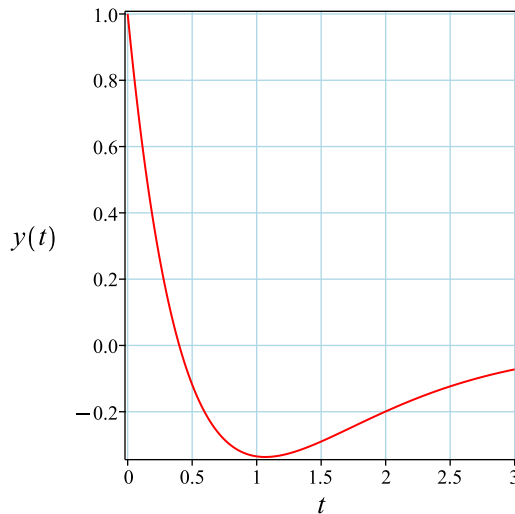
Which simplifies to

$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2}\right)$$

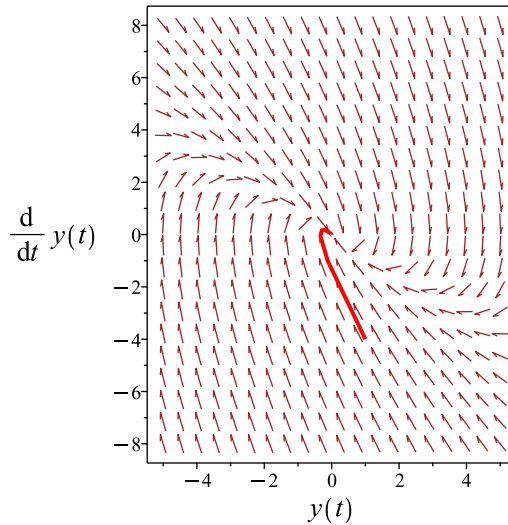
Summary

The solution(s) found are the following

$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2} \right) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2} \right)$$

Verified OK.

9.15.3 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t)y' + \frac{(p(t))^2 + p'(t)}{2}y = f(t)$$

Where $p(t) = 3$. Therefore, there is an integrating factor given by

$$\begin{aligned} M(x) &= e^{\frac{1}{2} \int p \, dx} \\ &= e^{\int 3 \, dx} \\ &= e^{\frac{3t}{2}} \end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$(M(x)y)'' = 0$$
$$\left(e^{\frac{3t}{2}}y\right)'' = 0$$

Integrating once gives

$$\left(e^{\frac{3t}{2}}y\right)' = c_1$$

Integrating again gives

$$\left(e^{\frac{3t}{2}}y\right) = c_1t + c_2$$

Hence the solution is

$$y = \frac{c_1t + c_2}{e^{\frac{3t}{2}}}$$

Or

$$y = c_1t e^{-\frac{3t}{2}} + c_2e^{-\frac{3t}{2}}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1t e^{-\frac{3t}{2}} + c_2e^{-\frac{3t}{2}} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $t = 0$ in the above gives

$$1 = c_2 \tag{1A}$$

Taking derivative of the solution gives

$$y' = c_1e^{-\frac{3t}{2}} - \frac{3c_1t e^{-\frac{3t}{2}}}{2} - \frac{3c_2e^{-\frac{3t}{2}}}{2}$$

substituting $y' = -4$ and $t = 0$ in the above gives

$$-4 = c_1 - \frac{3c_2}{2} \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = -\frac{5}{2}$$
$$c_2 = 1$$

Substituting these values back in above solution results in

$$y = -\frac{5e^{-\frac{3t}{2}}t}{2} + e^{-\frac{3t}{2}}$$

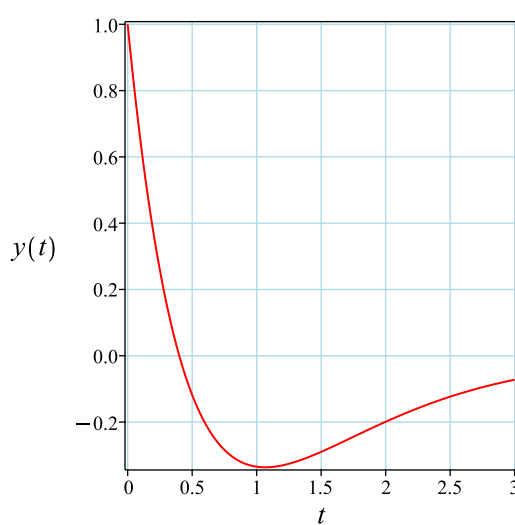
Which simplifies to

$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2}\right)$$

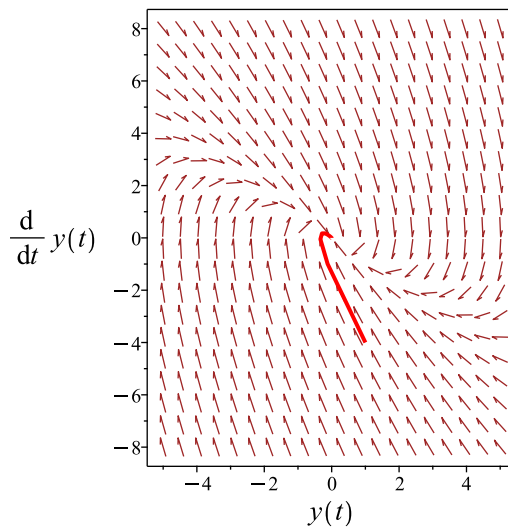
Summary

The solution(s) found are the following

$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2}\right) \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2}\right)$$

Verified OK.

9.15.4 Solving using Kovacic algorithm

Writing the ode as

$$4y'' + 12y' + 9y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 4 \\ B &= 12 \\ C &= 9 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 0 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = 0 \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 450: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{12}{4} dt} \end{aligned}$$

$$\begin{aligned}
&= z_1 e^{-\frac{3t}{2}} \\
&= z_1 \left(e^{-\frac{3t}{2}} \right)
\end{aligned}$$

Which simplifies to

$$y_1 = e^{-\frac{3t}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}
y_2 &= y_1 \int \frac{e^{\int -\frac{12}{4} dt}}{(y_1)^2} dt \\
&= y_1 \int \frac{e^{-3t}}{(y_1)^2} dt \\
&= y_1(t)
\end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 \left(e^{-\frac{3t}{2}} \right) + c_2 \left(e^{-\frac{3t}{2}}(t) \right)
\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{-\frac{3t}{2}} + c_2 t e^{-\frac{3t}{2}} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 1$ and $t = 0$ in the above gives

$$1 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = -\frac{3c_1 e^{-\frac{3t}{2}}}{2} + c_2 e^{-\frac{3t}{2}} - \frac{3c_2 t e^{-\frac{3t}{2}}}{2}$$

substituting $y' = -4$ and $t = 0$ in the above gives

$$-4 = -\frac{3c_1}{2} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 1$$
$$c_2 = -\frac{5}{2}$$

Substituting these values back in above solution results in

$$y = -\frac{5e^{-\frac{3t}{2}}}{2} + e^{-\frac{3t}{2}}$$

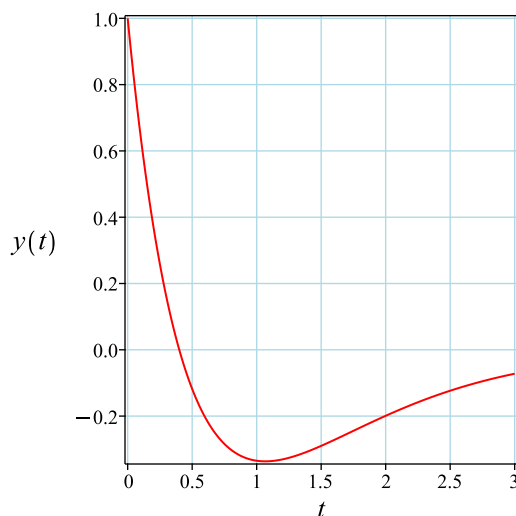
Which simplifies to

$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2}\right)$$

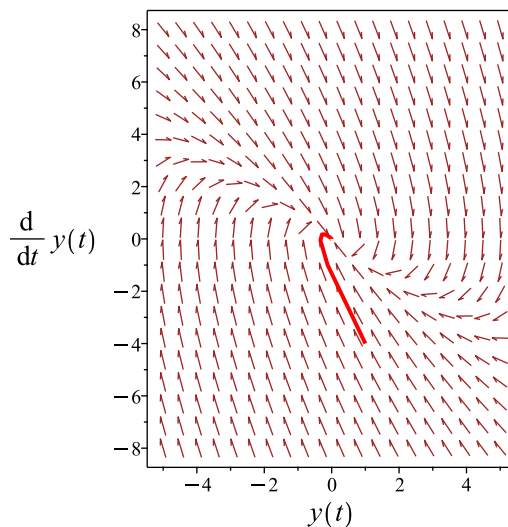
Summary

The solution(s) found are the following

$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2}\right) \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2} \right)$$

Verified OK.

9.15.5 Maple step by step solution

Let's solve

$$\left[4y'' + 12y' + 9y = 0, y(0) = 1, y' \Big|_{\{t=0\}} = -4 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -3y' - \frac{9y}{4}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + 3y' + \frac{9y}{4} = 0$$

- Characteristic polynomial of ODE

$$r^2 + 3r + \frac{9}{4} = 0$$

- Factor the characteristic polynomial

$$\frac{(2r+3)^2}{4} = 0$$

- Root of the characteristic polynomial

$$r = -\frac{3}{2}$$

- 1st solution of the ODE

$$y_1(t) = e^{-\frac{3t}{2}}$$

- Repeated root, multiply $y_1(t)$ by t to ensure linear independence

$$y_2(t) = e^{-\frac{3t}{2}} t$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t)$$

- Substitute in solutions

$$y = c_1 e^{-\frac{3t}{2}} + c_2 t e^{-\frac{3t}{2}}$$

- Check validity of solution $y = c_1 e^{-\frac{3t}{2}} + c_2 t e^{-\frac{3t}{2}}$

- Use initial condition $y(0) = 1$

$$1 = c_1$$
- Compute derivative of the solution
$$y' = -\frac{3c_1 e^{-\frac{3t}{2}}}{2} + c_2 e^{-\frac{3t}{2}} - \frac{3c_2 t e^{-\frac{3t}{2}}}{2}$$
- Use the initial condition $y' \Big|_{\{t=0\}} = -4$

$$-4 = -\frac{3c_1}{2} + c_2$$
- Solve for c_1 and c_2

$$\{c_1 = 1, c_2 = -\frac{5}{2}\}$$
- Substitute constant values into general solution and simplify
$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2}\right)$$
- Solution to the IVP
$$y = e^{-\frac{3t}{2}} \left(1 - \frac{5t}{2}\right)$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 14

```
dsolve([4*dif(y(t),t$2)+12*dif(y(t),t)+9*y(t) = 0,y(0) = 1, D(y)(0) = -4],y(t), singsol=all)
```

$$y(t) = -\frac{e^{-\frac{3t}{2}}(-2 + 5t)}{2}$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 21

```
DSolve[{4*y''[t]+12*y'[t]+9*y[t]==0,{y[0]==1,y'[0]==-4}},y[t],t,IncludeSingularSolutions ->
```

$$y(t) \rightarrow \frac{1}{2}e^{-3t/2}(2 - 5t)$$

9.16 problem 16

9.16.1 Existence and uniqueness analysis	2516
9.16.2 Solving as second order linear constant coeff ode	2516
9.16.3 Solving as linear second order ode solved by an integrating factor ode	2518
9.16.4 Solving using Kovacic algorithm	2520
9.16.5 Maple step by step solution	2523

Internal problem ID [668]

Internal file name [OUTPUT/668_Sunday_June_05_2022_01_46_40_AM_23845844/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 16.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$y'' - y' + \frac{y}{4} = 0$$

With initial conditions

$$[y(0) = 2, y'(0) = b]$$

9.16.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(t)y' + q(t)y = F$$

Where here

$$p(t) = -1$$

$$q(t) = \frac{1}{4}$$

$$F = 0$$

Hence the ode is

$$y'' - y' + \frac{y}{4} = 0$$

The domain of $p(t) = -1$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = \frac{1}{4}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

9.16.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = -1, C = \frac{1}{4}$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} - \lambda e^{\lambda t} + \frac{e^{\lambda t}}{4} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 - \lambda + \frac{1}{4} = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -1, C = \frac{1}{4}$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(-1)^2 - (4)(1) \left(\frac{1}{4}\right)} \\ &= \frac{1}{2}\end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = -\frac{1}{2}$. Therefore the solution is

$$y = c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}} \quad (1)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $t = 0$ in the above gives

$$2 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$y' = \frac{c_1 e^{\frac{t}{2}}}{2} + c_2 e^{\frac{t}{2}} + \frac{c_2 t e^{\frac{t}{2}}}{2}$$

substituting $y' = b$ and $t = 0$ in the above gives

$$b = \frac{c_1}{2} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned}c_1 &= 2 \\ c_2 &= b - 1\end{aligned}$$

Substituting these values back in above solution results in

$$y = (b - 1) t e^{\frac{t}{2}} + 2 e^{\frac{t}{2}}$$

Which simplifies to

$$y = (2 + t(b - 1)) e^{\frac{t}{2}}$$

Summary

The solution(s) found are the following

$$y = (2 + t(b - 1)) e^{\frac{t}{2}} \quad (1)$$

Verification of solutions

$$y = (2 + t(b - 1)) e^{\frac{t}{2}}$$

Verified OK.

9.16.3 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t)y' + \frac{(p(t))^2 + p'(t)}{2}y = f(t)$$

Where $p(t) = -1$. Therefore, there is an integrating factor given by

$$\begin{aligned} M(x) &= e^{\frac{1}{2} \int p \, dx} \\ &= e^{\int -1 \, dx} \\ &= e^{-\frac{t}{2}} \end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned} (M(x)y)'' &= 0 \\ \left(e^{-\frac{t}{2}}y\right)'' &= 0 \end{aligned}$$

Integrating once gives

$$\left(e^{-\frac{t}{2}}y\right)' = c_1$$

Integrating again gives

$$\left(e^{-\frac{t}{2}}y\right) = c_1t + c_2$$

Hence the solution is

$$y = \frac{c_1t + c_2}{e^{-\frac{t}{2}}}$$

Or

$$y = c_1t e^{\frac{t}{2}} + c_2 e^{\frac{t}{2}}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 t e^{\frac{t}{2}} + c_2 e^{\frac{t}{2}} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $t = 0$ in the above gives

$$2 = c_2 \quad (1A)$$

Taking derivative of the solution gives

$$y' = c_1 e^{\frac{t}{2}} + \frac{c_1 t e^{\frac{t}{2}}}{2} + \frac{c_2 e^{\frac{t}{2}}}{2}$$

substituting $y' = b$ and $t = 0$ in the above gives

$$b = c_1 + \frac{c_2}{2} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned} c_1 &= b - 1 \\ c_2 &= 2 \end{aligned}$$

Substituting these values back in above solution results in

$$y = t e^{\frac{t}{2}} b - t e^{\frac{t}{2}} + 2 e^{\frac{t}{2}}$$

Which simplifies to

$$y = (2 + t(b - 1)) e^{\frac{t}{2}}$$

Summary

The solution(s) found are the following

$$y = (2 + t(b - 1)) e^{\frac{t}{2}} \quad (1)$$

Verification of solutions

$$y = (2 + t(b - 1)) e^{\frac{t}{2}}$$

Verified OK.

9.16.4 Solving using Kovacic algorithm

Writing the ode as

$$y'' - y' + \frac{y}{4} = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = -1 \quad (3)$$

$$C = \frac{1}{4}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \quad (5)$$
$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \quad (6)$$

Comparing the above to (5) shows that

$$s = 0$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(t) = 0 \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 452: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-1}{1} dt} \end{aligned}$$

$$\begin{aligned}
&= z_1 e^{\frac{t}{2}} \\
&= z_1 \left(e^{\frac{t}{2}} \right)
\end{aligned}$$

Which simplifies to

$$y_1 = e^{\frac{t}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}
y_2 &= y_1 \int \frac{e^{\int -\frac{-1}{1} dt}}{(y_1)^2} dt \\
&= y_1 \int \frac{e^t}{(y_1)^2} dt \\
&= y_1(t)
\end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1 \left(e^{\frac{t}{2}} \right) + c_2 \left(e^{\frac{t}{2}}(t) \right)
\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$y = c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}} \tag{1}$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $y = 2$ and $t = 0$ in the above gives

$$2 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$y' = \frac{c_1 e^{\frac{t}{2}}}{2} + c_2 e^{\frac{t}{2}} + \frac{c_2 t e^{\frac{t}{2}}}{2}$$

substituting $y' = b$ and $t = 0$ in the above gives

$$b = \frac{c_1}{2} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$\begin{aligned} c_1 &= 2 \\ c_2 &= b - 1 \end{aligned}$$

Substituting these values back in above solution results in

$$y = (b - 1) t e^{\frac{t}{2}} + 2 e^{\frac{t}{2}}$$

Which simplifies to

$$y = (2 + t(b - 1)) e^{\frac{t}{2}}$$

Summary

The solution(s) found are the following

$$y = (2 + t(b - 1)) e^{\frac{t}{2}} \quad (1)$$

Verification of solutions

$$y = (2 + t(b - 1)) e^{\frac{t}{2}}$$

Verified OK.

9.16.5 Maple step by step solution

Let's solve

$$\left[y'' - y' + \frac{y}{4} = 0, y(0) = 2, y' \Big|_{\{t=0\}} = b \right]$$

- Highest derivative means the order of the ODE is 2
 y''
- Characteristic polynomial of ODE
 $r^2 - r + \frac{1}{4} = 0$
- Factor the characteristic polynomial
 $\frac{(2r-1)^2}{4} = 0$
- Root of the characteristic polynomial

$$r = \frac{1}{2}$$

- 1st solution of the ODE

$$y_1(t) = e^{\frac{t}{2}}$$

- Repeated root, multiply $y_1(t)$ by t to ensure linear independence

$$y_2(t) = t e^{\frac{t}{2}}$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t)$$

- Substitute in solutions

$$y = c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}}$$

- Check validity of solution $y = c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}}$

- Use initial condition $y(0) = 2$

$$2 = c_1$$

- Compute derivative of the solution

$$y' = \frac{c_1 e^{\frac{t}{2}}}{2} + c_2 e^{\frac{t}{2}} + \frac{c_2 t e^{\frac{t}{2}}}{2}$$

- Use the initial condition $y' \Big|_{\{t=0\}} = b$

$$b = \frac{c_1}{2} + c_2$$

- Solve for c_1 and c_2

$$\{c_1 = 2, c_2 = b - 1\}$$

- Substitute constant values into general solution and simplify

$$y = (2 + t(b - 1)) e^{\frac{t}{2}}$$

- Solution to the IVP

$$y = (2 + t(b - 1)) e^{\frac{t}{2}}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 16

```
dsolve([diff(y(t),t$2)-diff(y(t),t)+25/100*y(t) = 0,y(0) = 2, D(y)(0) = b],y(t), singsol=all
```

$$y(t) = (2 + t(b - 1)) e^{\frac{t}{2}}$$

✓ Solution by Mathematica

Time used: 0.014 (sec). Leaf size: 20

```
DSolve[{y''[t]-y'[t]+25/100*y[t]==0,{y[0]==2,y'[0]==b}},y[t],t,IncludeSingularSolutions -> T
```

$$y(t) \rightarrow e^{t/2}((b - 1)t + 2)$$

9.17 problem 23

9.17.1 Maple step by step solution 2527

Internal problem ID [669]

Internal file name [OUTPUT/669_Sunday_June_05_2022_01_46_41_AM_12322057/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 23.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "reduction_of_order", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_1", "second_order_change_of_variable_on_y_method_2", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_Emden, _Fowler], [_2nd_order, _linear, `_with_symmetry_[0,F(x)]`]]
```

$$t^2 y'' - 4ty' + 6y = 0$$

Given that one solution of the ode is

$$y_1 = t^2$$

Given one basis solution $y_1(t)$, then the second basis solution is given by

$$y_2(t) = y_1 \left(\int \frac{e^{-\int p dt}}{y_1^2} dt \right)$$

Where $p(x)$ is the coefficient of y' when the ode is written in the normal form

$$y'' + p(t) y' + q(t) y = f(t)$$

Looking at the ode to solve shows that

$$p(t) = -\frac{4}{t}$$

Therefore

$$y_2(t) = t^2 \left(\int \frac{e^{-(\int -\frac{4}{t} dt)}}{t^4} dt \right)$$

$$y_2(t) = t^2 \int \frac{t^4}{t^4} dt$$

$$y_2(t) = t^2 \left(\int 1 dt \right)$$

$$y_2(t) = t^3$$

Hence the solution is

$$\begin{aligned} y &= c_1 y_1(t) + c_2 y_2(t) \\ &= c_2 t^3 + c_1 t^2 \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_2 t^3 + c_1 t^2 \tag{1}$$

Verification of solutions

$$y = c_2 t^3 + c_1 t^2$$

Verified OK.

9.17.1 Maple step by step solution

Let's solve

$$y'' t^2 - 4ty' + 6y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{4y'}{t} - \frac{6y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{4y'}{t} + \frac{6y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 - 4ty' + 6y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right) t^2 - 4\frac{d}{ds}y(s) + 6y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) - 5\frac{d}{ds}y(s) + 6y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 - 5r + 6 = 0$$

- Factor the characteristic polynomial

$$(r - 2)(r - 3) = 0$$

- Roots of the characteristic polynomial

$$r = (2, 3)$$

- 1st solution of the ODE

$$y_1(s) = e^{2s}$$

- 2nd solution of the ODE

$$y_2(s) = e^{3s}$$

- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$

- Substitute in solutions

$$y(s) = c_1 e^{2s} + c_2 e^{3s}$$

- Change variables back using $s = \ln(t)$

$$y = c_2 t^3 + c_1 t^2$$

- Simplify

$$y = t^2(c_2 t + c_1)$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 13

```
dsolve([t^2*diff(y(t),t$2)-4*t*diff(y(t),t)+6*y(t)=0,t^2],singsol=all)
```

$$y(t) = t^2(c_2 t + c_1)$$

✓ Solution by Mathematica

Time used: 0.012 (sec). Leaf size: 16

```
DSolve[t^2*y''[t]-4*t*y'[t]+6*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t^2(c_2 t + c_1)$$

9.18 problem 24

9.18.1 Maple step by step solution 2531

Internal problem ID [670]

Internal file name [OUTPUT/670_Sunday_June_05_2022_01_46_42_AM_6399244/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 24.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "**reduction_of_order**", "**second_order_euler_ode**", "**second_order_change_of_variable_on_x_method_2**", "**second_order_change_of_variable_on_y_method_2**", "**second_order_ode_non_constant_coeff_transformation_on_B**"

Maple gives the following as the ode type

```
[[_Emden , _Fowler]]
```

$$t^2y'' + 2ty' - 2y = 0$$

Given that one solution of the ode is

$$y_1 = t$$

Given one basis solution $y_1(t)$, then the second basis solution is given by

$$y_2(t) = y_1 \left(\int \frac{e^{-\int p dt}}{y_1^2} dt \right)$$

Where $p(x)$ is the coefficient of y' when the ode is written in the normal form

$$y'' + p(t)y' + q(t)y = f(t)$$

Looking at the ode to solve shows that

$$p(t) = \frac{2}{t}$$

Therefore

$$y_2(t) = t \left(\int \frac{e^{-\left(\int \frac{2}{t} dt\right)}}{t^2} dt \right)$$

$$y_2(t) = t \int \frac{1}{t^2} dt$$

$$y_2(t) = t \left(\int \frac{1}{t^4} dt \right)$$

$$y_2(t) = -\frac{1}{3t^2}$$

Hence the solution is

$$\begin{aligned} y &= c_1 y_1(t) + c_2 y_2(t) \\ &= c_1 t - \frac{c_2}{3t^2} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 t - \frac{c_2}{3t^2} \tag{1}$$

Verification of solutions

$$y = c_1 t - \frac{c_2}{3t^2}$$

Verified OK.

9.18.1 Maple step by step solution

Let's solve

$$y''t^2 + 2ty' - 2y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{2y'}{t} + \frac{2y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{2y'}{t} - \frac{2y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 + 2ty' - 2y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right) t^2 + 2\frac{d}{ds}y(s) - 2y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) + \frac{d}{ds}y(s) - 2y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 + r - 2 = 0$$

- Factor the characteristic polynomial

$$(r + 2)(r - 1) = 0$$

- Roots of the characteristic polynomial

$$r = (-2, 1)$$

- 1st solution of the ODE

$$y_1(s) = e^{-2s}$$

- 2nd solution of the ODE

$$y_2(s) = e^s$$

- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$

- Substitute in solutions

$$y(s) = c_1 e^{-2s} + c_2 e^s$$

- Change variables back using $s = \ln(t)$

$$y = \frac{c_1}{t^2} + c_2 t$$

- Simplify

$$y = \frac{c_1}{t^2} + c_2 t$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve([t^2*diff(y(t),t$2)+2*t*diff(y(t),t)-2*y(t)=0,t],singsol=all)
```

$$y(t) = \frac{c_1 t^3 + c_2}{t^2}$$

✓ Solution by Mathematica

Time used: 0.01 (sec). Leaf size: 16

```
DSolve[t^2*y''[t]+2*t*y'[t]-2*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{c_1}{t^2} + c_2 t$$

9.19 problem 25

9.19.1 Maple step by step solution 2535

Internal problem ID [671]

Internal file name [OUTPUT/671_Sunday_June_05_2022_01_46_43_AM_97850935/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 25.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "reduction_of_order", "second_order_euler_ode", "exact linear second order ode", "second_order_integrable_as_is", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_2nd_order , _exact , _linear , _homogeneous]]
```

$$t^2 y'' + 3ty' + y = 0$$

Given that one solution of the ode is

$$y_1 = \frac{1}{t}$$

Given one basis solution $y_1(t)$, then the second basis solution is given by

$$y_2(t) = y_1 \left(\int \frac{e^{-\int p dt}}{y_1^2} dt \right)$$

Where $p(x)$ is the coefficient of y' when the ode is written in the normal form

$$y'' + p(t)y' + q(t)y = f(t)$$

Looking at the ode to solve shows that

$$p(t) = \frac{3}{t}$$

Therefore

$$y_2(t) = \frac{\int e^{-(\int \frac{3}{t} dt)} t^2 dt}{t}$$

$$y_2(t) = \frac{1}{t} \int \frac{1}{t^3}, dt$$

$$y_2(t) = \frac{\int \frac{1}{t} dt}{t}$$

$$y_2(t) = \frac{\ln(t)}{t}$$

Hence the solution is

$$\begin{aligned} y &= c_1 y_1(t) + c_2 y_2(t) \\ &= \frac{c_1}{t} + \frac{c_2 \ln(t)}{t} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t} \quad (1)$$

Verification of solutions

$$y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t}$$

Verified OK.

9.19.1 Maple step by step solution

Let's solve

$$y'' t^2 + 3ty' + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{3y'}{t} - \frac{y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{3y'}{t} + \frac{y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 + 3ty' + y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right) t^2 + 3\frac{d}{ds}y(s) + y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) + 2\frac{d}{ds}y(s) + y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 + 2r + 1 = 0$$

- Factor the characteristic polynomial

$$(r + 1)^2 = 0$$

- Root of the characteristic polynomial

$$r = -1$$

- 1st solution of the ODE

$$y_1(s) = e^{-s}$$

- Repeated root, multiply $y_1(s)$ by s to ensure linear independence

$$y_2(s) = s e^{-s}$$

- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$

- Substitute in solutions

$$y(s) = c_1 e^{-s} + c_2 s e^{-s}$$

- Change variables back using $s = \ln(t)$

$$y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t}$$

- Simplify

$$y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve([t^2*diff(y(t),t$2)+3*t*diff(y(t),t)+y(t)=0,1/t],singsol=all)
```

$$y(t) = \frac{c_2 \ln(t) + c_1}{t}$$

✓ Solution by Mathematica

Time used: 0.016 (sec). Leaf size: 17

```
DSolve[t^2*y''[t]+3*t*y'[t]+y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{c_2 \log(t) + c_1}{t}$$

9.20 problem 26

9.20.1 Maple step by step solution 2539

Internal problem ID [672]

Internal file name [OUTPUT/672_Sunday_June_05_2022_01_46_43_AM_19696878/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 26.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "**reduction_of_order**", "**second_order_change_of_variable_on_y_method_1**", "**second_order_change_of_variable_on_y_method_2**"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$t^2 y'' - t(2+t)y' + (2+t)y = 0$$

Given that one solution of the ode is

$$y_1 = t$$

Given one basis solution $y_1(t)$, then the second basis solution is given by

$$y_2(t) = y_1 \left(\int \frac{e^{-\int p dt}}{y_1^2} dt \right)$$

Where $p(x)$ is the coefficient of y' when the ode is written in the normal form

$$y'' + p(t)y' + q(t)y = f(t)$$

Looking at the ode to solve shows that

$$p(t) = \frac{-t^2 - 2t}{t^2}$$

Therefore

$$y_2(t) = t \left(\int \frac{e^{-\left(\int \frac{-t^2-2t}{t^2} dt\right)}}{t^2} dt \right)$$

$$y_2(t) = t \int \frac{e^{t+2\ln(t)}}{t^2} dt$$

$$y_2(t) = t \left(\int e^t dt \right)$$

$$y_2(t) = t e^t$$

Hence the solution is

$$\begin{aligned} y &= c_1 y_1(t) + c_2 y_2(t) \\ &= c_1 t + c_2 t e^t \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 t + c_2 t e^t \tag{1}$$

Verification of solutions

$$y = c_1 t + c_2 t e^t$$

Verified OK.

9.20.1 Maple step by step solution

Let's solve

$$y'' t^2 + (-t^2 - 2t) y' + (2 + t) y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{(2+t)y}{t^2} + \frac{(2+t)y'}{t}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{(2+t)y'}{t} + \frac{(2+t)y}{t^2} = 0$$

- Check to see if $t_0 = 0$ is a regular singular point

- Define functions

$$\left[P_2(t) = -\frac{2+t}{t}, P_3(t) = \frac{2+t}{t^2} \right]$$

- $t \cdot P_2(t)$ is analytic at $t = 0$

$$(t \cdot P_2(t)) \Big|_{t=0} = -2$$

- $t^2 \cdot P_3(t)$ is analytic at $t = 0$

$$(t^2 \cdot P_3(t)) \Big|_{t=0} = 2$$

- $t = 0$ is a regular singular point

Check to see if $t_0 = 0$ is a regular singular point

$$t_0 = 0$$

- Multiply by denominators

$$y''t^2 - t(2+t)y' + (2+t)y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k t^{k+r}$$

- Rewrite ODE with series expansions

- Convert $t^m \cdot y$ to series expansion for $m = 0..1$

$$t^m \cdot y = \sum_{k=0}^{\infty} a_k t^{k+r+m}$$

- Shift index using $k \rightarrow k - m$

$$t^m \cdot y = \sum_{k=m}^{\infty} a_{k-m} t^{k+r}$$

- Convert $t^m \cdot y'$ to series expansion for $m = 1..2$

$$t^m \cdot y' = \sum_{k=0}^{\infty} a_k (k+r) t^{k+r-1+m}$$

- Shift index using $k \rightarrow k + 1 - m$

$$t^m \cdot y' = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k+1-m+r) t^{k+r}$$

- Convert $t^2 \cdot y''$ to series expansion

$$t^2 \cdot y'' = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) t^{k+r}$$

Rewrite ODE with series expansions

$$a_0(-1+r)(-2+r)t^r + \left(\sum_{k=1}^{\infty} (a_k(k+r-1)(k+r-2) - a_{k-1}(k+r-2))t^{k+r} \right) = 0$$

- a_0 cannot be 0 by assumption, giving the indicial equation
 $(-1+r)(-2+r) = 0$
- Values of r that satisfy the indicial equation
 $r \in \{1, 2\}$
- Each term in the series must be 0, giving the recursion relation
 $(k+r-2)(a_k(k+r-1) - a_{k-1}) = 0$
- Shift index using $k \rightarrow k+1$
 $(k+r-1)(a_{k+1}(k+r) - a_k) = 0$
- Recursion relation that defines series solution to ODE
 $a_{k+1} = \frac{a_k}{k+r}$
- Recursion relation for $r = 1$
 $a_{k+1} = \frac{a_k}{k+1}$
- Solution for $r = 1$
 $\left[y = \sum_{k=0}^{\infty} a_k t^{k+1}, a_{k+1} = \frac{a_k}{k+1} \right]$
- Recursion relation for $r = 2$
 $a_{k+1} = \frac{a_k}{k+2}$
- Solution for $r = 2$
 $\left[y = \sum_{k=0}^{\infty} a_k t^{k+2}, a_{k+1} = \frac{a_k}{k+2} \right]$
- Combine solutions and rename parameters
 $\left[y = \left(\sum_{k=0}^{\infty} a_k t^{k+1} \right) + \left(\sum_{k=0}^{\infty} b_k t^{k+2} \right), a_{k+1} = \frac{a_k}{k+1}, b_{k+1} = \frac{b_k}{k+2} \right]$

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacics algorithm  
  A Liouvillian solution exists  
  Reducible group (found an exponential solution)  
  Reducible group (found another exponential solution)  
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 12

```
dsolve([t^2*diff(y(t),t$2)-t*(t+2)*diff(y(t),t)+(t+2)*y(t)=0,t],singsol=all)
```

$$y(t) = t(c_1 + c_2 e^t)$$

✓ Solution by Mathematica

Time used: 0.025 (sec). Leaf size: 16

```
DSolve[t^2*y'[t]-t*(t+2)*y'[t]+(t+2)*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t(c_2 e^t + c_1)$$

9.21 problem 27

9.21.1 Maple step by step solution 2544

Internal problem ID [673]

Internal file name [OUTPUT/673_Sunday_June_05_2022_01_46_44_AM_69159856/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 27.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "reduction_of_order", "second_order_bessel_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2"

Maple gives the following as the ode type

```
[[_Emden, _Fowler], [_2nd_order, _linear, `_with_symmetry_[0,F(x)]]`]]
```

$$xy'' - y' + 4yx^3 = 0$$

Given that one solution of the ode is

$$y_1 = \sin(x^2)$$

Given one basis solution $y_1(x)$, then the second basis solution is given by

$$y_2(x) = y_1 \left(\int \frac{e^{-\int p dx}}{y_1^2} dx \right)$$

Where $p(x)$ is the coefficient of y' when the ode is written in the normal form

$$y'' + p(x)y' + q(x)y = f(x)$$

Looking at the ode to solve shows that

$$p(x) = -\frac{1}{x}$$

Therefore

$$y_2(x) = \sin(x^2) \left(\int \frac{e^{-(\int -\frac{1}{x} dx)}}{\sin(x^2)^2} dx \right)$$

$$y_2(x) = \sin(x^2) \int \frac{x}{\sin(x^2)^2} dx$$

$$y_2(x) = \sin(x^2) \left(\int \csc(x^2)^2 x dx \right)$$

$$y_2(x) = -\frac{\sin(x^2) \cot(x^2)}{2}$$

Hence the solution is

$$\begin{aligned} y &= c_1 y_1(x) + c_2 y_2(x) \\ &= \sin(x^2) c_1 - \frac{c_2 \sin(x^2) \cot(x^2)}{2} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \sin(x^2) c_1 - \frac{c_2 \sin(x^2) \cot(x^2)}{2} \quad (1)$$

Verification of solutions

$$y = \sin(x^2) c_1 - \frac{c_2 \sin(x^2) \cot(x^2)}{2}$$

Verified OK.

9.21.1 Maple step by step solution

Let's solve

$$y''x - y' + 4yx^3 = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{y'}{x} - 4x^2y$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{y'}{x} + 4x^2y = 0$$

□ Check to see if $x_0 = 0$ is a regular singular point

○ Define functions

$$[P_2(x) = -\frac{1}{x}, P_3(x) = 4x^2]$$

○ $x \cdot P_2(x)$ is analytic at $x = 0$

$$(x \cdot P_2(x)) \Big|_{x=0} = -1$$

○ $x^2 \cdot P_3(x)$ is analytic at $x = 0$

$$(x^2 \cdot P_3(x)) \Big|_{x=0} = 0$$

○ $x = 0$ is a regular singular point

Check to see if $x_0 = 0$ is a regular singular point

$$x_0 = 0$$

• Multiply by denominators

$$y''x - y' + 4yx^3 = 0$$

• Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^{k+r}$$

□ Rewrite ODE with series expansions

○ Convert $x^3 \cdot y$ to series expansion

$$x^3 \cdot y = \sum_{k=0}^{\infty} a_k x^{k+r+3}$$

○ Shift index using $k \rightarrow k - 3$

$$x^3 \cdot y = \sum_{k=3}^{\infty} a_{k-3} x^{k+r}$$

○ Convert y' to series expansion

$$y' = \sum_{k=0}^{\infty} a_k (k+r) x^{k+r-1}$$

○ Shift index using $k \rightarrow k + 1$

$$y' = \sum_{k=-1}^{\infty} a_{k+1} (k+1+r) x^{k+r}$$

○ Convert $x \cdot y''$ to series expansion

$$x \cdot y'' = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) x^{k+r-1}$$

- Shift index using $k \rightarrow k+1$

$$x \cdot y'' = \sum_{k=-1}^{\infty} a_{k+1} (k+1+r)(k+r) x^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r(-2+r) x^{-1+r} + a_1 (1+r)(-1+r) x^r + a_2 (2+r)r x^{1+r} + a_3 (3+r)(1+r) x^{2+r} + \left(\sum_{k=3}^{\infty} a_k \right)$$

- a_0 cannot be 0 by assumption, giving the indicial equation
 $r(-2+r) = 0$
- Values of r that satisfy the indicial equation
 $r \in \{0, 2\}$
- The coefficients of each power of x must be 0
 $[a_1(1+r)(-1+r) = 0, a_2(2+r)r = 0, a_3(3+r)(1+r) = 0]$
- Solve for the dependent coefficient(s)
 $\{a_1 = 0, a_2 = 0, a_3 = 0\}$
- Each term in the series must be 0, giving the recursion relation
 $a_{k+1}(k+1+r)(k+r-1) + 4a_{k-3} = 0$
- Shift index using $k \rightarrow k+3$
 $a_{k+4}(k+4+r)(k+2+r) + 4a_k = 0$
- Recursion relation that defines series solution to ODE
 $a_{k+4} = -\frac{4a_k}{(k+4+r)(k+2+r)}$
- Recursion relation for $r = 0$
 $a_{k+4} = -\frac{4a_k}{(k+4)(k+2)}$
- Solution for $r = 0$
 $\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+4} = -\frac{4a_k}{(k+4)(k+2)}, a_1 = 0, a_2 = 0, a_3 = 0 \right]$
- Recursion relation for $r = 2$
 $a_{k+4} = -\frac{4a_k}{(k+6)(k+4)}$
- Solution for $r = 2$

$$\left[y = \sum_{k=0}^{\infty} a_k x^{k+2}, a_{k+4} = -\frac{4a_k}{(k+6)(k+4)}, a_1 = 0, a_2 = 0, a_3 = 0 \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k x^k \right) + \left(\sum_{k=0}^{\infty} b_k x^{k+2} \right), a_{k+4} = -\frac{4a_k}{(k+4)(k+2)}, a_1 = 0, a_2 = 0, a_3 = 0, b_{k+4} = -\frac{4b_k}{(k+6)(k+4)} \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
<- linear_1 successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve([x*diff(y(x),x$2)-diff(y(x),x)+4*x^3*y(x)=0,sin(x^2)],singsol=all)
```

$$y(x) = c_1 \sin(x^2) + c_2 \cos(x^2)$$

✓ Solution by Mathematica

Time used: 0.019 (sec). Leaf size: 20

```
DSolve[x*y''[x]-y'[x]+4*x^3*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow c_1 \cos(x^2) + c_2 \sin(x^2)$$

9.22 problem 28

9.22.1 Maple step by step solution 2549

Internal problem ID [674]

Internal file name [OUTPUT/674_Sunday_June_05_2022_01_46_45_AM_64926509/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 28.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "**reduction_of_order**", "**second_order_change_of_variable_on_y_method_2**", "**second_order_ode_non_constant_coeff_transformation_on_B**"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(x - 1)y'' - y'x + y = 0$$

Given that one solution of the ode is

$$y_1 = e^x$$

Given one basis solution $y_1(x)$, then the second basis solution is given by

$$y_2(x) = y_1 \left(\int \frac{e^{-(\int p dx)}}{y_1^2} dx \right)$$

Where $p(x)$ is the coefficient of y' when the ode is written in the normal form

$$y'' + p(x)y' + q(x)y = f(x)$$

Looking at the ode to solve shows that

$$p(x) = -\frac{x}{x-1}$$

Therefore

$$y_2(x) = e^x \left(\int e^{-\left(\int -\frac{x}{x-1} dx\right)} e^{-2x} dx \right)$$

$$y_2(x) = e^x \int \frac{e^{x+\ln(x-1)}}{e^{2x}} dx$$

$$y_2(x) = e^x \left(\int (x-1) e^{-x} dx \right)$$

$$y_2(x) = -e^x x e^{-x}$$

Hence the solution is

$$\begin{aligned} y &= c_1 y_1(x) + c_2 y_2(x) \\ &= c_1 e^x - c_2 e^x x e^{-x} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^x - c_2 e^x x e^{-x} \tag{1}$$

Verification of solutions

$$y = c_1 e^x - c_2 e^x x e^{-x}$$

Verified OK.

9.22.1 Maple step by step solution

Let's solve

$$(x-1)y'' - y'x + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{y}{x-1} + \frac{xy'}{x-1}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{xy'}{x-1} + \frac{y}{x-1} = 0$$

- Check to see if $x_0 = 1$ is a regular singular point

- Define functions

$$[P_2(x) = -\frac{x}{x-1}, P_3(x) = \frac{1}{x-1}]$$

- $(x-1) \cdot P_2(x)$ is analytic at $x=1$

$$((x-1) \cdot P_2(x)) \Big|_{x=1} = -1$$
- $(x-1)^2 \cdot P_3(x)$ is analytic at $x=1$

$$((x-1)^2 \cdot P_3(x)) \Big|_{x=1} = 0$$

- $x=1$ is a regular singular point

Check to see if $x_0 = 1$ is a regular singular point

$$x_0 = 1$$

- Multiply by denominators

$$(x-1)y'' - y'x + y = 0$$

- Change variables using $x = u + 1$ so that the regular singular point is at $u = 0$

$$u \left(\frac{d^2}{du^2} y(u) \right) + (-u-1) \left(\frac{d}{du} y(u) \right) + y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k+1-m$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k+1-m+r) u^{k+r}$$

- Convert $u \cdot \left(\frac{d^2}{du^2} y(u) \right)$ to series expansion

$$u \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) u^{k+r-1}$$

- Shift index using $k \rightarrow k+1$

$$u \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=-1}^{\infty} a_{k+1} (k+1+r)(k+r) u^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r(-2+r) u^{-1+r} + \left(\sum_{k=0}^{\infty} (a_{k+1}(k+1+r)(k+r-1) - a_k(k+r-1)) u^{k+r} \right) = 0$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$r(-2+r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \{0, 2\}$$

- Each term in the series must be 0, giving the recursion relation

$$(k+r-1)(a_{k+1}(k+1+r) - a_k) = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+1} = \frac{a_k}{k+1+r}$$

- Recursion relation for $r = 0$

$$a_{k+1} = \frac{a_k}{k+1}$$

- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{k+1} = \frac{a_k}{k+1} \right]$$

- Revert the change of variables $u = x - 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x-1)^k, a_{k+1} = \frac{a_k}{k+1} \right]$$

- Recursion relation for $r = 2$

$$a_{k+1} = \frac{a_k}{k+3}$$

- Solution for $r = 2$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+2}, a_{k+1} = \frac{a_k}{k+3} \right]$$

- Revert the change of variables $u = x - 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x-1)^{k+2}, a_{k+1} = \frac{a_k}{k+3} \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k (x-1)^k \right) + \left(\sum_{k=0}^{\infty} b_k (x-1)^{k+2} \right), a_{k+1} = \frac{a_k}{k+1}, b_{k+1} = \frac{b_k}{k+3} \right]$$

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacics algorithm  
  A Liouvillian solution exists  
  Reducible group (found an exponential solution)  
  Reducible group (found another exponential solution)  
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 12

```
dsolve([(x-1)*diff(y(x),x$2)-x*diff(y(x),x)+y(x)=0,exp(x)],singsol=all)
```

$$y(x) = c_1x + e^x c_2$$

✓ Solution by Mathematica

Time used: 0.031 (sec). Leaf size: 17

```
DSolve[(x-1)*y'[x]-x*y'[x]+y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow c_1 e^x - c_2 x$$

9.23 problem 29

9.23.1 Maple step by step solution 2554

Internal problem ID [675]

Internal file name [OUTPUT/675_Sunday_June_05_2022_01_46_46_AM_68547472/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 29.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "**reduction_of_order**", "**second_order_bessel_ode**"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$x^2 y'' - \left(x - \frac{3}{16}\right) y = 0$$

Given that one solution of the ode is

$$y_1 = x^{\frac{1}{4}} e^{2\sqrt{x}}$$

Given one basis solution $y_1(x)$, then the second basis solution is given by

$$y_2(x) = y_1 \left(\int \frac{e^{-(\int p dx)}}{y_1^2} dx \right)$$

Where $p(x)$ is the coefficient of y' when the ode is written in the normal form

$$y'' + p(x) y' + q(x) y = f(x)$$

Looking at the ode to solve shows that

$$p(x) = 0$$

Therefore

$$y_2(x) = x^{\frac{1}{4}} e^{2\sqrt{x}} \left(\int \frac{e^{-(\int 0 dx)} e^{-4\sqrt{x}}}{\sqrt{x}} dx \right)$$

$$y_2(x) = x^{\frac{1}{4}} e^{2\sqrt{x}} \int \frac{1}{\sqrt{x} e^{4\sqrt{x}}} dx$$

$$y_2(x) = x^{\frac{1}{4}} e^{2\sqrt{x}} \left(\int \frac{e^{-4\sqrt{x}}}{\sqrt{x}} dx \right)$$

$$y_2(x) = -\frac{x^{\frac{1}{4}} e^{2\sqrt{x}} e^{-4\sqrt{x}}}{2}$$

Hence the solution is

$$\begin{aligned} y &= c_1 y_1(x) + c_2 y_2(x) \\ &= x^{\frac{1}{4}} e^{2\sqrt{x}} c_1 - \frac{c_2 x^{\frac{1}{4}} e^{2\sqrt{x}} e^{-4\sqrt{x}}}{2} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = x^{\frac{1}{4}} e^{2\sqrt{x}} c_1 - \frac{c_2 x^{\frac{1}{4}} e^{2\sqrt{x}} e^{-4\sqrt{x}}}{2} \quad (1)$$

Verification of solutions

$$y = x^{\frac{1}{4}} e^{2\sqrt{x}} c_1 - \frac{c_2 x^{\frac{1}{4}} e^{2\sqrt{x}} e^{-4\sqrt{x}}}{2}$$

Verified OK.

9.23.1 Maple step by step solution

Let's solve

$$x^2 y'' + \left(-x + \frac{3}{16}\right) y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{(16x-3)y}{16x^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{(16x-3)y}{16x^2} = 0$$

□ Check to see if $x_0 = 0$ is a regular singular point

○ Define functions

$$[P_2(x) = 0, P_3(x) = -\frac{16x-3}{16x^2}]$$

○ $x \cdot P_2(x)$ is analytic at $x = 0$

$$(x \cdot P_2(x)) \Big|_{x=0} = 0$$

○ $x^2 \cdot P_3(x)$ is analytic at $x = 0$

$$(x^2 \cdot P_3(x)) \Big|_{x=0} = \frac{3}{16}$$

○ $x = 0$ is a regular singular point

Check to see if $x_0 = 0$ is a regular singular point

$$x_0 = 0$$

• Multiply by denominators

$$16x^2y'' + (-16x + 3)y = 0$$

• Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^{k+r}$$

□ Rewrite ODE with series expansions

○ Convert $x^m \cdot y$ to series expansion for $m = 0..1$

$$x^m \cdot y = \sum_{k=0}^{\infty} a_k x^{k+r+m}$$

○ Shift index using $k \rightarrow k - m$

$$x^m \cdot y = \sum_{k=m}^{\infty} a_{k-m} x^{k+r}$$

○ Convert $x^2 \cdot y''$ to series expansion

$$x^2 \cdot y'' = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) x^{k+r}$$

Rewrite ODE with series expansions

$$a_0(-1+4r)(-3+4r)x^r + \left(\sum_{k=1}^{\infty} (a_k(4k+4r-1)(4k+4r-3) - 16a_{k-1}) x^{k+r} \right) = 0$$

• a_0 cannot be 0 by assumption, giving the indicial equation

$$(-1 + 4r)(-3 + 4r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \left\{ \frac{1}{4}, \frac{3}{4} \right\}$$

- Each term in the series must be 0, giving the recursion relation

$$16\left(k + r - \frac{3}{4}\right)\left(k + r - \frac{1}{4}\right)a_k - 16a_{k-1} = 0$$

- Shift index using $k \rightarrow k + 1$

$$16\left(k + \frac{1}{4} + r\right)\left(k + \frac{3}{4} + r\right)a_{k+1} - 16a_k = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+1} = \frac{16a_k}{(4k+1+4r)(4k+3+4r)}$$

- Recursion relation for $r = \frac{1}{4}$

$$a_{k+1} = \frac{16a_k}{(4k+2)(4k+4)}$$

- Solution for $r = \frac{1}{4}$

$$\left[y = \sum_{k=0}^{\infty} a_k x^{k+\frac{1}{4}}, a_{k+1} = \frac{16a_k}{(4k+2)(4k+4)} \right]$$

- Recursion relation for $r = \frac{3}{4}$

$$a_{k+1} = \frac{16a_k}{(4k+4)(4k+6)}$$

- Solution for $r = \frac{3}{4}$

$$\left[y = \sum_{k=0}^{\infty} a_k x^{k+\frac{3}{4}}, a_{k+1} = \frac{16a_k}{(4k+4)(4k+6)} \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k x^{k+\frac{1}{4}} \right) + \left(\sum_{k=0}^{\infty} b_k x^{k+\frac{3}{4}} \right), a_{k+1} = \frac{16a_k}{(4k+2)(4k+4)}, b_{k+1} = \frac{16b_k}{(4k+4)(4k+6)} \right]$$

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacics algorithm  
  A Liouvillian solution exists  
  Group is reducible or imprimitive  
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 25

```
dsolve([x^2*diff(y(x),x$2)-(x-1875/10000)*y(x)=0,x^(1/4)*exp(2*sqrt(x))],singsol=all)
```

$$y(x) = x^{\frac{1}{4}}(c_1 \sinh(2\sqrt{x}) + c_2 \cosh(2\sqrt{x}))$$

✓ Solution by Mathematica

Time used: 0.047 (sec). Leaf size: 41

```
DSolve[x^2*y''[x]-(x-1875/10000)*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2}e^{-2\sqrt{x}}\sqrt[4]{x}\left(2c_1e^{4\sqrt{x}} - c_2\right)$$

9.24 problem 30

9.24.1 Maple step by step solution 2559

Internal problem ID [676]

Internal file name [OUTPUT/676_Sunday_June_05_2022_01_46_46_AM_36615809/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 30.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "**reduction_of_order**", "**second_order_bessel_ode**", "**second_order_change_of_variable_on_y_method_1**"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$x^2 y'' + y' x + \left(x^2 - \frac{1}{4}\right) y = 0$$

Given that one solution of the ode is

$$y_1 = \frac{\sin(x)}{\sqrt{x}}$$

Given one basis solution $y_1(x)$, then the second basis solution is given by

$$y_2(x) = y_1 \left(\int \frac{e^{-\int p dx}}{y_1^2} dx \right)$$

Where $p(x)$ is the coefficient of y' when the ode is written in the normal form

$$y'' + p(x) y' + q(x) y = f(x)$$

Looking at the ode to solve shows that

$$p(x) = \frac{1}{x}$$

Therefore

$$y_2(x) = \frac{\sin(x) \left(\int \frac{e^{-\left(\int \frac{1}{x} dx\right)} x}{\sin(x)^2} dx \right)}{\sqrt{x}}$$

$$y_2(x) = \frac{\sin(x)}{\sqrt{x}} \int \frac{\frac{1}{x}}{\frac{\sin(x)^2}{x}} dx$$

$$y_2(x) = \frac{\sin(x) \left(\int \csc(x)^2 dx \right)}{\sqrt{x}}$$

$$y_2(x) = -\frac{\sin(x) \cot(x)}{\sqrt{x}}$$

Hence the solution is

$$\begin{aligned} y &= c_1 y_1(x) + c_2 y_2(x) \\ &= \frac{\sin(x) c_1}{\sqrt{x}} - \frac{c_2 \sin(x) \cot(x)}{\sqrt{x}} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{\sin(x) c_1}{\sqrt{x}} - \frac{c_2 \sin(x) \cot(x)}{\sqrt{x}} \quad (1)$$

Verification of solutions

$$y = \frac{\sin(x) c_1}{\sqrt{x}} - \frac{c_2 \sin(x) \cot(x)}{\sqrt{x}}$$

Verified OK.

9.24.1 Maple step by step solution

Let's solve

$$x^2 y'' + y' x + \left(x^2 - \frac{1}{4}\right) y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{(4x^2-1)y}{4x^2} - \frac{y'}{x}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{y'}{x} + \frac{(4x^2-1)y}{4x^2} = 0$$

- Check to see if $x_0 = 0$ is a regular singular point

- Define functions

$$\left[P_2(x) = \frac{1}{x}, P_3(x) = \frac{4x^2-1}{4x^2} \right]$$

- $x \cdot P_2(x)$ is analytic at $x = 0$

$$(x \cdot P_2(x)) \Big|_{x=0} = 1$$

- $x^2 \cdot P_3(x)$ is analytic at $x = 0$

$$(x^2 \cdot P_3(x)) \Big|_{x=0} = -\frac{1}{4}$$

- $x = 0$ is a regular singular point

Check to see if $x_0 = 0$ is a regular singular point

$$x_0 = 0$$

- Multiply by denominators

$$4x^2y'' + 4y'x + (4x^2 - 1)y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^{k+r}$$

- Rewrite ODE with series expansions

- Convert $x^m \cdot y$ to series expansion for $m = 0..2$

$$x^m \cdot y = \sum_{k=0}^{\infty} a_k x^{k+r+m}$$

- Shift index using $k \rightarrow k - m$

$$x^m \cdot y = \sum_{k=m}^{\infty} a_{k-m} x^{k+r}$$

- Convert $x \cdot y'$ to series expansion

$$x \cdot y' = \sum_{k=0}^{\infty} a_k (k+r) x^{k+r}$$

- Convert $x^2 \cdot y''$ to series expansion

$$x^2 \cdot y'' = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) x^{k+r}$$

Rewrite ODE with series expansions

$$a_0(1+2r)(-1+2r)x^r + a_1(3+2r)(1+2r)x^{1+r} + \left(\sum_{k=2}^{\infty} (a_k(2k+2r+1)(2k+2r-1) + 4a_{k-2}) \right) x^{k+r}$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$(1+2r)(-1+2r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \left\{ -\frac{1}{2}, \frac{1}{2} \right\}$$

- Each term must be 0

$$a_1(3+2r)(1+2r) = 0$$

- Solve for the dependent coefficient(s)

$$a_1 = 0$$

- Each term in the series must be 0, giving the recursion relation

$$a_k(4k^2 + 8kr + 4r^2 - 1) + 4a_{k-2} = 0$$

- Shift index using $k \rightarrow k+2$

$$a_{k+2}(4(k+2)^2 + 8(k+2)r + 4r^2 - 1) + 4a_k = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+2} = -\frac{4a_k}{4k^2 + 8kr + 4r^2 + 16k + 16r + 15}$$

- Recursion relation for $r = -\frac{1}{2}$

$$a_{k+2} = -\frac{4a_k}{4k^2 + 12k + 8}$$

- Solution for $r = -\frac{1}{2}$

$$\left[y = \sum_{k=0}^{\infty} a_k x^{k-\frac{1}{2}}, a_{k+2} = -\frac{4a_k}{4k^2 + 12k + 8}, a_1 = 0 \right]$$

- Recursion relation for $r = \frac{1}{2}$

$$a_{k+2} = -\frac{4a_k}{4k^2 + 20k + 24}$$

- Solution for $r = \frac{1}{2}$

$$\left[y = \sum_{k=0}^{\infty} a_k x^{k+\frac{1}{2}}, a_{k+2} = -\frac{4a_k}{4k^2 + 20k + 24}, a_1 = 0 \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k x^{k-\frac{1}{2}} \right) + \left(\sum_{k=0}^{\infty} b_k x^{k+\frac{1}{2}} \right), a_{k+2} = -\frac{4a_k}{4k^2 + 12k + 8}, a_1 = 0, b_{k+2} = -\frac{4b_k}{4k^2 + 20k + 24}, b_1 = 0 \right]$$

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacics algorithm  
  A Liouvillian solution exists  
  Group is reducible or imprimitive  
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 17

```
dsolve([x^2*diff(y(x),x$2)+x*diff(y(x),x)+(x^2-25/100)*y(x)=0,x^(-1/2)*sin(x)],singsol=all)
```

$$y(x) = \frac{c_1 \sin(x) + c_2 \cos(x)}{\sqrt{x}}$$

✓ Solution by Mathematica

Time used: 0.042 (sec). Leaf size: 39

```
DSolve[x^2*y''[x]+x*y'[x]+(x^2-25/100)*y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{e^{-ix}(2c_1 - ic_2 e^{2ix})}{2\sqrt{x}}$$

9.25 problem 40

9.25.1 Solving as second order euler ode ode	2563
9.25.2 Solving as second order change of variable on x method 2 ode .	2564
9.25.3 Solving as second order change of variable on x method 1 ode .	2567
9.25.4 Solving as second order change of variable on y method 2 ode .	2569
9.25.5 Solving using Kovacic algorithm	2571
9.25.6 Maple step by step solution	2576

Internal problem ID [677]

Internal file name [OUTPUT/677_Sunday_June_05_2022_01_46_47_AM_3427519/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 40.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_Emden, _Fowler], [_2nd_order, _linear, `_with_symmetry_[0,F(x)]`]]
```

$$t^2 y'' - 3ty' + 4y = 0$$

9.25.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} - 3trt^{r-1} + 4t^r = 0$$

Simplifying gives

$$r(r-1)t^r - 3rt^r + 4t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r - 1) - 3r + 4 = 0$$

Or

$$r^2 - 4r + 4 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = 2$$

$$r_2 = 2$$

Since the roots are equal, then the general solution is

$$y = c_1 y_1 + c_2 y_2$$

Where $y_1 = t^r$ and $y_2 = t^r \ln(t)$. Hence

$$y = c_1 t^2 + c_2 t^2 \ln(t)$$

Summary

The solution(s) found are the following

$$y = c_1 t^2 + c_2 t^2 \ln(t) \tag{1}$$

Verification of solutions

$$y = c_1 t^2 + c_2 t^2 \ln(t)$$

Verified OK.

9.25.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$t^2 y'' - 3ty' + 4y = 0 \tag{1}$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \tag{2}$$

Where

$$p(t) = -\frac{3}{t}$$
$$q(t) = \frac{4}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t)\tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t)dt)} dt \\ &= \int e^{-(\int -\frac{3}{t}dt)} dt \\ &= \int e^{3\ln(t)} dt \\ &= \int t^3 dt \\ &= \frac{t^4}{4} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{4}{t^2} \\ &= \frac{4}{t^8} \end{aligned} \quad (7)$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + \frac{4y(\tau)}{t^8} &= 0 \end{aligned}$$

But in terms of τ

$$\frac{4}{t^8} = \frac{1}{4\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{y(\tau)}{4\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$4\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$4\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + \tau^r = 0$$

Simplifying gives

$$4r(r-1)\tau^r + 0\tau^r + \tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$4r(r-1) + 0 + 1 = 0$$

Or

$$4r^2 - 4r + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2}$$
$$r_2 = \frac{1}{2}$$

Since the roots are equal, then the general solution is

$$y(\tau) = c_1y_1 + c_2y_2$$

Where $y_1 = \tau^r$ and $y_2 = \tau^r \ln(\tau)$. Hence

$$y(\tau) = c_1\sqrt{\tau} + c_2\sqrt{\tau} \ln(\tau)$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{(-2c_2 \ln(2) + c_2 \ln(t^4) + c_1) \sqrt{t^4}}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{(-2c_2 \ln(2) + c_2 \ln(t^4) + c_1) \sqrt{t^4}}{2} \quad (1)$$

Verification of solutions

$$y = \frac{(-2c_2 \ln(2) + c_2 \ln(t^4) + c_1) \sqrt{t^4}}{2}$$

Verified OK.

9.25.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$t^2 y'' - 3ty' + 4y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = -\frac{3}{t}$$
$$q(t) = \frac{4}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c} \sqrt{q}$$
$$= \frac{2\sqrt{\frac{1}{t^2}}}{c} \quad (6)$$
$$\tau'' = -\frac{2}{c\sqrt{\frac{1}{t^2}} t^3}$$

Substituting the above into (4) results in

$$\begin{aligned}
 p_1(\tau) &= \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \\
 &= \frac{-\frac{2}{c\sqrt{\frac{1}{t^2}}t^3} - \frac{3}{t}\frac{2\sqrt{\frac{1}{t^2}}}{c}}{\left(\frac{2\sqrt{\frac{1}{t^2}}}{c}\right)^2} \\
 &= -2c
 \end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned}
 y(\tau)'' + p_1y(\tau)' + q_1y(\tau) &= 0 \\
 \frac{d^2}{d\tau^2}y(\tau) - 2c\left(\frac{d}{d\tau}y(\tau)\right) + c^2y(\tau) &= 0 \tag{7}
 \end{aligned}$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{c\tau}c_1$$

Now from (6)

$$\begin{aligned}
 \tau &= \int \frac{1}{c}\sqrt{q} dt \\
 &= \frac{\int 2\sqrt{\frac{1}{t^2}} dt}{c} \\
 &= \frac{2\sqrt{\frac{1}{t^2}} t \ln(t)}{c}
 \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = c_1t^2$$

Summary

The solution(s) found are the following

$$y = c_1t^2 \tag{1}$$

Verification of solutions

$$y = c_1t^2$$

Verified OK.

9.25.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$t^2 y'' - 3ty' + 4y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = -\frac{3}{t}$$
$$q(t) = \frac{4}{t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} - \frac{3n}{t^2} + \frac{4}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = 2 \quad (6)$$

Substituting this value in (3) gives

$$v''(t) + \frac{v'(t)}{t} = 0$$
$$v''(t) + \frac{v'(t)}{t} = 0 \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{u}{t} \end{aligned}$$

Where $f(t) = -\frac{1}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= -\frac{1}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{1}{t} dt \\ \ln(u) &= -\ln(t) + c_1 \\ u &= e^{-\ln(t)+c_1} \\ &= \frac{c_1}{t} \end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned} v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= c_1 \ln(t) + c_2 \end{aligned}$$

Hence

$$\begin{aligned} y &= v(t) t^n \\ &= (c_1 \ln(t) + c_2) t^2 \\ &= (c_1 \ln(t) + c_2) t^2 \end{aligned}$$

Summary

The solution(s) found are the following

$$y = (c_1 \ln(t) + c_2) t^2 \quad (1)$$

Verification of solutions

$$y = (c_1 \ln(t) + c_2) t^2$$

Verified OK.

9.25.5 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' - 3ty' + 4y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t^2 \\ B &= -3t \\ C &= 4 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(-\frac{1}{4t^2}\right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 462: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{1}{4t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{1}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2}$	$\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = \frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= \frac{1}{2} - \left(\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{t - c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= \frac{1}{2t} + (-)(0) \\ &= \frac{1}{2t} \\ &= \frac{1}{2t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2\left(\frac{1}{2t}\right)(0) + \left(\left(-\frac{1}{2t^2}\right) + \left(\frac{1}{2t}\right)^2 - \left(-\frac{1}{4t^2}\right) \right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= p e^{\int \omega dt} \\ &= e^{\int \frac{1}{2t} dt} \\ &= \sqrt{t} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\&= z_1 e^{-\int \frac{1}{2} \frac{-3t}{t^2} dt} \\&= z_1 e^{\frac{3 \ln(t)}{2}} \\&= z_1 \left(t^{\frac{3}{2}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = t^2$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-3t}{t^2} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{3 \ln(t)}}{(y_1)^2} dt \\&= y_1 (\ln(t))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 (t^2) + c_2 (t^2 (\ln(t)))\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 t^2 + c_2 t^2 \ln(t) \tag{1}$$

Verification of solutions

$$y = c_1 t^2 + c_2 t^2 \ln(t)$$

Verified OK.

9.25.6 Maple step by step solution

Let's solve

$$y''t^2 - 3ty' + 4y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{3y'}{t} - \frac{4y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{3y'}{t} + \frac{4y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 - 3ty' + 4y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right) t^2 - 3\frac{d}{ds}y(s) + 4y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) - 4\frac{d}{ds}y(s) + 4y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 - 4r + 4 = 0$$

- Factor the characteristic polynomial
 $(r - 2)^2 = 0$
- Root of the characteristic polynomial
 $r = 2$
- 1st solution of the ODE
 $y_1(s) = e^{2s}$
- Repeated root, multiply $y_1(s)$ by s to ensure linear independence
 $y_2(s) = s e^{2s}$
- General solution of the ODE
 $y(s) = c_1 y_1(s) + c_2 y_2(s)$
- Substitute in solutions
 $y(s) = c_1 e^{2s} + c_2 s e^{2s}$
- Change variables back using $s = \ln(t)$
 $y = c_1 t^2 + c_2 t^2 \ln(t)$
- Simplify
 $y = t^2(c_2 \ln(t) + c_1)$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(t^2*diff(y(t),t$2)-3*t*diff(y(t),t)+4*y(t)=0,y(t), singsol=all)
```

$$y(t) = t^2(c_2 \ln(t) + c_1)$$

✓ Solution by Mathematica

Time used: 0.016 (sec). Leaf size: 18

```
DSolve[t^2*y''[t]-3*t*y'[t]+4*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t^2(2c_2 \log(t) + c_1)$$

9.26 problem 41

9.26.1 Solving as second order euler ode	2580
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Internal problem ID [678]

Internal file name [OUTPUT/678_Sunday_June_05_2022_01_46_48_AM_2297909/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 41.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_Emden, _Fowler]]
```

$$t^2 y'' + 2ty' + \frac{y}{4} = 0$$

The ode can be written as

$$4t^2 y'' + 8ty' + y = 0$$

Which shows it is a Euler ODE.

9.26.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$4t^2(r(r-1))t^{r-2} + 8trt^{r-1} + t^r = 0$$

Simplifying gives

$$4r(r-1)t^r + 8rt^r + t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$4r(r-1) + 8r + 1 = 0$$

Or

$$4r^2 + 4r + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -\frac{1}{2}$$
$$r_2 = -\frac{1}{2}$$

Since the roots are equal, then the general solution is

$$y = c_1y_1 + c_2y_2$$

Where $y_1 = t^r$ and $y_2 = t^r \ln(t)$. Hence

$$y = \frac{c_1}{\sqrt{t}} + \frac{c_2 \ln(t)}{\sqrt{t}}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{\sqrt{t}} + \frac{c_2 \ln(t)}{\sqrt{t}} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{\sqrt{t}} + \frac{c_2 \ln(t)}{\sqrt{t}}$$

Verified OK.

9.26.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$4t^2y'' + 8ty' + y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{2}{t}$$
$$q(t) = \frac{1}{4t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t)\tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t)dt)} dt \\ &= \int e^{-(\int \frac{2}{t} dt)} dt \\ &= \int e^{-2\ln(t)} dt \\ &= \int \frac{1}{t^2} dt \\ &= -\frac{1}{t} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{\frac{1}{4t^2}}{\frac{1}{t^4}} \\ &= \frac{t^2}{4} \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + \frac{t^2y(\tau)}{4} &= 0 \end{aligned}$$

But in terms of τ

$$\frac{t^2}{4} = \frac{1}{4\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{y(\tau)}{4\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$4\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$4\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + \tau^r = 0$$

Simplifying gives

$$4r(r-1)\tau^r + 0\tau^r + \tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$4r(r-1) + 0 + 1 = 0$$

Or

$$4r^2 - 4r + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2}$$

$$r_2 = \frac{1}{2}$$

Since the roots are equal, then the general solution is

$$y(\tau) = c_1 y_1 + c_2 y_2$$

Where $y_1 = \tau^r$ and $y_2 = \tau^r \ln(\tau)$. Hence

$$y(\tau) = c_1 \sqrt{\tau} + c_2 \sqrt{\tau} \ln(\tau)$$

The above solution is now transformed back to y using (6) which results in

$$y = \left(\ln\left(-\frac{1}{t}\right) c_2 + c_1 \right) \sqrt{-\frac{1}{t}}$$

Summary

The solution(s) found are the following

$$y = \left(\ln\left(-\frac{1}{t}\right) c_2 + c_1 \right) \sqrt{-\frac{1}{t}} \quad (1)$$

Verification of solutions

$$y = \left(\ln\left(-\frac{1}{t}\right) c_2 + c_1 \right) \sqrt{-\frac{1}{t}}$$

Verified OK.

9.26.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$4t^2 y'' + 8ty' + y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{2}{t}$$

$$q(t) = \frac{1}{4t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\begin{aligned} \tau' &= \frac{1}{c}\sqrt{q} \\ &= \frac{\sqrt{\frac{1}{t^2}}}{2c} \\ \tau'' &= -\frac{1}{2c\sqrt{\frac{1}{t^2}}t^3} \end{aligned} \quad (6)$$

Substituting the above into (4) results in

$$\begin{aligned} p_1(\tau) &= \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \\ &= \frac{-\frac{1}{2c\sqrt{\frac{1}{t^2}}t^3} + \frac{2}{t}\frac{\sqrt{\frac{1}{t^2}}}{2c}}{\left(\frac{\sqrt{\frac{1}{t^2}}}{2c}\right)^2} \\ &= 2c \end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned} y(\tau)'' + p_1y(\tau)' + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + 2c\left(\frac{d}{d\tau}y(\tau)\right) + c^2y(\tau) &= 0 \end{aligned} \quad (7)$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{-c\tau}c_1$$

Now from (6)

$$\begin{aligned}\tau &= \int \frac{1}{c} \sqrt{q} dt \\ &= \frac{\int \frac{\sqrt{\frac{1}{t^2}} dt}{c}}{c} \\ &= \frac{\sqrt{\frac{1}{t^2}} t \ln(t)}{2c}\end{aligned}$$

Substituting the above into the solution obtained gives

$$y = \frac{c_1}{\sqrt{t}}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{\sqrt{t}} \quad (1)$$

Verification of solutions

$$y = \frac{c_1}{\sqrt{t}}$$

Verified OK.

9.26.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$4t^2 y'' + 8ty' + y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$\begin{aligned}p(t) &= \frac{2}{t} \\ q(t) &= \frac{1}{4t^2}\end{aligned}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} + \frac{2n}{t^2} + \frac{1}{4t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = -\frac{1}{2} \quad (6)$$

Substituting this value in (3) gives

$$\begin{aligned} v''(t) + \frac{v'(t)}{t} &= 0 \\ v''(t) + \frac{v'(t)}{t} &= 0 \end{aligned} \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{u}{t} \end{aligned}$$

Where $f(t) = -\frac{1}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= -\frac{1}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{1}{t} dt \\ \ln(u) &= -\ln(t) + c_1 \\ u &= e^{-\ln(t)+c_1} \\ &= \frac{c_1}{t} \end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned}v'(t) &= u(t) \\v(t) &= \int u(t) dt + c_2 \\&= c_1 \ln(t) + c_2\end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\&= \frac{c_1 \ln(t) + c_2}{\sqrt{t}} \\&= \frac{c_1 \ln(t) + c_2}{\sqrt{t}}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \ln(t) + c_2}{\sqrt{t}} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 \ln(t) + c_2}{\sqrt{t}}$$

Verified OK.

9.26.5 Solving using Kovacic algorithm

Writing the ode as

$$4t^2 y'' + 8ty' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= 4t^2 \\B &= 8t \\C &= 1\end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(-\frac{1}{4t^2}\right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 464: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{1}{4t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{1}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2}$	$\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_\infty^- = \frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_\infty^- - (\alpha_{c_1}^-) \\ &= \frac{1}{2} - \left(\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t-c} \right) + s(\infty)[\sqrt{r}]_\infty$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^-}{t-c_1} \right) + (-)[\sqrt{r}]_\infty \\ &= \frac{1}{2t} + (-)(0) \\ &= \frac{1}{2t} \\ &= \frac{1}{2t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(\frac{1}{2t}\right)(0) + \left(\left(-\frac{1}{2t^2}\right) + \left(\frac{1}{2t}\right)^2 - \left(-\frac{1}{4t^2}\right)\right) = 0$$

$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \frac{1}{2t} dt} \\ &= \sqrt{t} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{8t}{4t^2} dt} \\ &= z_1 e^{-\ln(t)} \\ &= z_1 \left(\frac{1}{t}\right) \end{aligned}$$

Which simplifies to

$$y_1 = \frac{1}{\sqrt{t}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{8t}{4t^2} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{-2\ln(t)}}{(y_1)^2} dt \\&= y_1(\ln(t))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 \left(\frac{1}{\sqrt{t}} \right) + c_2 \left(\frac{1}{\sqrt{t}} (\ln(t)) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{\sqrt{t}} + \frac{c_2 \ln(t)}{\sqrt{t}} \quad (1)$$

Verification of solutions

$$y = \frac{c_1}{\sqrt{t}} + \frac{c_2 \ln(t)}{\sqrt{t}}$$

Verified OK.

9.26.6 Maple step by step solution

Let's solve

$$4y''t^2 + 8ty' + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{2y'}{t} - \frac{y}{4t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{2y'}{t} + \frac{y}{4t^2} = 0$$

- Multiply by denominators of the ODE

$$4y''t^2 + 8ty' + y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$4\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right)t^2 + 8\frac{d}{ds}y(s) + y(s) = 0$$

- Simplify

$$4\frac{d^2}{ds^2}y(s) + 4\frac{d}{ds}y(s) + y(s) = 0$$

- Isolate 2nd derivative

$$\frac{d^2}{ds^2}y(s) = -\frac{d}{ds}y(s) - \frac{y(s)}{4}$$

- Group terms with $y(s)$ on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$\frac{d^2}{ds^2}y(s) + \frac{d}{ds}y(s) + \frac{y(s)}{4} = 0$$

- Characteristic polynomial of ODE

$$r^2 + r + \frac{1}{4} = 0$$

- Factor the characteristic polynomial

$$\frac{(2r+1)^2}{4} = 0$$

- Root of the characteristic polynomial

$$r = -\frac{1}{2}$$

- 1st solution of the ODE

$$y_1(s) = e^{-\frac{s}{2}}$$

- Repeated root, multiply $y_1(s)$ by s to ensure linear independence

$$y_2(s) = s e^{-\frac{s}{2}}$$

- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$

- Substitute in solutions

$$y(s) = c_1 e^{-\frac{s}{2}} + c_2 s e^{-\frac{s}{2}}$$

- Change variables back using $s = \ln(t)$

$$y = \frac{c_1}{\sqrt{t}} + \frac{c_2 \ln(t)}{\sqrt{t}}$$

- Simplify

$$y = \frac{c_1}{\sqrt{t}} + \frac{c_2 \ln(t)}{\sqrt{t}}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(t^2*diff(y(t),t$2)+2*t*diff(y(t),t)+25/100*y(t)=0,y(t), singsol=all)
```

$$y(t) = \frac{c_2 \ln(t) + c_1}{\sqrt{t}}$$

✓ Solution by Mathematica

Time used: 0.019 (sec). Leaf size: 24

```
DSolve[t^2*y''[t]+2*t*y'[t]+25/100*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{c_2 \log(t) + 2c_1}{2\sqrt{t}}$$

9.27 problem 42

9.27.1 Solving as second order euler ode	2596
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Internal problem ID [679]

Internal file name [OUTPUT/679_Sunday_June_05_2022_01_46_49_AM_62569225/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 42.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2", "second_order_ode_non_constant_coeff_transformation_on_B"

Maple gives the following as the ode type

```
[[_Emden, _Fowler]]
```

$$2t^2y'' - 5ty' + 5y = 0$$

9.27.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$2t^2(r(r-1))t^{r-2} - 5trt^{r-1} + 5t^r = 0$$

Simplifying gives

$$2r(r-1)t^r - 5rt^r + 5t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$2r(r-1) - 5r + 5 = 0$$

Or

$$2r^2 - 7r + 5 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$\begin{aligned} r_1 &= 1 \\ r_2 &= \frac{5}{2} \end{aligned}$$

Since the roots are real and distinct, then the general solution is

$$y = c_1y_1 + c_2y_2$$

Where $y_1 = t^{r_1}$ and $y_2 = t^{r_2}$. Hence

$$y = c_1t + c_2t^{\frac{5}{2}}$$

Summary

The solution(s) found are the following

$$y = c_1t + c_2t^{\frac{5}{2}} \tag{1}$$

Verification of solutions

$$y = c_1t + c_2t^{\frac{5}{2}}$$

Verified OK.

9.27.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$2t^2 y'' - 5ty' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = -\frac{5}{2t}$$
$$q(t) = \frac{5}{2t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t)\tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t)dt)} dt \\ &= \int e^{-(\int -\frac{5}{2t} dt)} dt \\ &= \int e^{\frac{5 \ln(t)}{2}} dt \\ &= \int t^{\frac{5}{2}} dt \\ &= \frac{2t^{\frac{7}{2}}}{7} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned}q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{\frac{5}{2t^2}}{t^5} \\ &= \frac{5}{2t^7}\end{aligned}\tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned}\frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + \frac{5y(\tau)}{2t^7} &= 0\end{aligned}$$

But in terms of τ

$$\frac{5}{2t^7} = \frac{10}{49\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{10y(\tau)}{49\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$49\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + 10y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$49\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + 10\tau^r = 0$$

Simplifying gives

$$49r(r-1)\tau^r + 0\tau^r + 10\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$49r(r-1) + 0 + 10 = 0$$

Or

$$49r^2 - 49r + 10 = 0\tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{2}{7}$$

$$r_2 = \frac{5}{7}$$

Since the roots are real and distinct, then the general solution is

$$y(\tau) = c_1 y_1 + c_2 y_2$$

Where $y_1 = \tau^{r_1}$ and $y_2 = \tau^{r_2}$. Hence

$$y(\tau) = c_1 \tau^{\frac{2}{7}} + c_2 \tau^{\frac{5}{7}}$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{c_1 2^{\frac{2}{7}} 7^{\frac{5}{7}} \left(t^{\frac{7}{2}}\right)^{\frac{2}{7}}}{7} + \frac{c_2 2^{\frac{5}{7}} 7^{\frac{2}{7}} \left(t^{\frac{7}{2}}\right)^{\frac{5}{7}}}{7}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 2^{\frac{2}{7}} 7^{\frac{5}{7}} \left(t^{\frac{7}{2}}\right)^{\frac{2}{7}}}{7} + \frac{c_2 2^{\frac{5}{7}} 7^{\frac{2}{7}} \left(t^{\frac{7}{2}}\right)^{\frac{5}{7}}}{7} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 2^{\frac{2}{7}} 7^{\frac{5}{7}} \left(t^{\frac{7}{2}}\right)^{\frac{2}{7}}}{7} + \frac{c_2 2^{\frac{5}{7}} 7^{\frac{2}{7}} \left(t^{\frac{7}{2}}\right)^{\frac{5}{7}}}{7}$$

Verified OK.

9.27.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$2t^2 y'' - 5ty' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = -\frac{5}{2t}$$

$$q(t) = \frac{5}{2t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c}\sqrt{q}$$

$$= \frac{\sqrt{10}\sqrt{\frac{1}{t^2}}}{2c} \quad (6)$$

$$\tau'' = -\frac{\sqrt{10}}{2c\sqrt{\frac{1}{t^2}}t^3}$$

Substituting the above into (4) results in

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2}$$

$$= \frac{-\frac{\sqrt{10}}{2c\sqrt{\frac{1}{t^2}}t^3} - \frac{5}{2t}\frac{\sqrt{10}\sqrt{\frac{1}{t^2}}}{2c}}{\left(\frac{\sqrt{10}\sqrt{\frac{1}{t^2}}}{2c}\right)^2}$$

$$= -\frac{7c\sqrt{10}}{10}$$

Therefore ode (3) now becomes

$$y(\tau)'' + p_1y(\tau)' + q_1y(\tau) = 0$$

$$\frac{d^2}{d\tau^2}y(\tau) - \frac{7c\sqrt{10}}{10}\left(\frac{d}{d\tau}y(\tau)\right) + c^2y(\tau) = 0 \quad (7)$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{\frac{7\sqrt{10}c\tau}{20}} \left(c_1 \cosh \left(\frac{3\sqrt{10}c\tau}{20} \right) + ic_2 \sinh \left(\frac{3\sqrt{10}c\tau}{20} \right) \right)$$

Now from (6)

$$\begin{aligned} \tau &= \int \frac{1}{c} \sqrt{q} dt \\ &= \frac{\int \frac{\sqrt{10} \sqrt{\frac{1}{t^2}} dt}{2}}{c} \\ &= \frac{\sqrt{10} \sqrt{\frac{1}{t^2}} t \ln(t)}{2c} \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = t^{\frac{7}{4}} \left(c_1 \cosh \left(\frac{3 \ln(t)}{4} \right) + ic_2 \sinh \left(\frac{3 \ln(t)}{4} \right) \right)$$

Summary

The solution(s) found are the following

$$y = t^{\frac{7}{4}} \left(c_1 \cosh \left(\frac{3 \ln(t)}{4} \right) + ic_2 \sinh \left(\frac{3 \ln(t)}{4} \right) \right) \quad (1)$$

Verification of solutions

$$y = t^{\frac{7}{4}} \left(c_1 \cosh \left(\frac{3 \ln(t)}{4} \right) + ic_2 \sinh \left(\frac{3 \ln(t)}{4} \right) \right)$$

Verified OK.

9.27.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$2t^2 y'' - 5ty' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = -\frac{5}{2t}$$
$$q(t) = \frac{5}{2t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} - \frac{5n}{2t^2} + \frac{5}{2t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = \frac{5}{2} \quad (6)$$

Substituting this value in (3) gives

$$v''(t) + \frac{5v'(t)}{2t} = 0$$
$$v''(t) + \frac{5v'(t)}{2t} = 0 \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{5u(t)}{2t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned}u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{5u}{2t}\end{aligned}$$

Where $f(t) = -\frac{5}{2t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= -\frac{5}{2t} dt \\ \int \frac{1}{u} du &= \int -\frac{5}{2t} dt \\ \ln(u) &= -\frac{5 \ln(t)}{2} + c_1 \\ u &= e^{-\frac{5 \ln(t)}{2} + c_1} \\ &= \frac{c_1}{t^{\frac{5}{2}}}\end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned}v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= -\frac{2c_1}{3t^{\frac{3}{2}}} + c_2\end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\ &= \left(-\frac{2c_1}{3t^{\frac{3}{2}}} + c_2\right) t^{\frac{5}{2}} \\ &= c_2 t^{\frac{5}{2}} - \frac{2c_1 t}{3}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \left(-\frac{2c_1}{3t^{\frac{3}{2}}} + c_2\right) t^{\frac{5}{2}} \quad (1)$$

Verification of solutions

$$y = \left(-\frac{2c_1}{3t^{\frac{3}{2}}} + c_2\right) t^{\frac{5}{2}}$$

Verified OK.

9.27.5 Solving as second order ode non constant coeff transformation on B ode

Given an ode of the form

$$Ay'' + By' + Cy = F(t)$$

This method reduces the order ode the ODE by one by applying the transformation

$$y = Bv$$

This results in

$$\begin{aligned}y' &= B'v + v'B \\y'' &= B''v + B'v' + v''B + v'B' \\&= v''B + 2v' + B' + B''v\end{aligned}$$

And now the original ode becomes

$$\begin{aligned}A(v''B + 2v'B' + B''v) + B(B'v + v'B) + CBv &= 0 \\ABv'' + (2AB' + B^2)v' + (AB'' + BB' + CB)v &= 0\end{aligned}\tag{1}$$

If the term $AB'' + BB' + CB$ is zero, then this method works and can be used to solve

$$ABv'' + (2AB' + B^2)v' = 0$$

By Using $u = v'$ which reduces the order of the above ode to one. The new ode is

$$ABu' + (2AB' + B^2)u = 0$$

The above ode is first order ode which is solved for u . Now a new ode $v' = u$ is solved for v as first order ode. Then the final solution is obtain from $y = Bv$.

This method works only if the term $AB'' + BB' + CB$ is zero. The given ODE shows that

$$\begin{aligned}A &= 2t^2 \\B &= -5t \\C &= 5 \\F &= 0\end{aligned}$$

The above shows that for this ode

$$\begin{aligned}AB'' + BB' + CB &= (2t^2)(0) + (-5t)(-5) + (5)(-5t) \\&= 0\end{aligned}$$

Hence the ode in v given in (1) now simplifies to

$$-10t^3v'' + (5t^2)v' = 0$$

Now by applying $v' = u$ the above becomes

$$-10t^3u'(t) + 5t^2u(t) = 0$$

Which is now solved for u . In canonical form the ODE is

$$\begin{aligned}u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{u}{2t}\end{aligned}$$

Where $f(t) = \frac{1}{2t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= \frac{1}{2t} dt \\ \int \frac{1}{u} du &= \int \frac{1}{2t} dt \\ \ln(u) &= \frac{\ln(t)}{2} + c_1 \\ u &= e^{\frac{\ln(t)}{2} + c_1} \\ &= c_1\sqrt{t}\end{aligned}$$

The ode for v now becomes

$$\begin{aligned}v' &= u \\ &= c_1\sqrt{t}\end{aligned}$$

Which is now solved for v . Integrating both sides gives

$$\begin{aligned}v(t) &= \int c_1\sqrt{t} dt \\ &= \frac{2t^{\frac{3}{2}}c_1}{3} + c_2\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y(t) &= Bv \\ &= (-5t) \left(\frac{2t^{\frac{3}{2}}c_1}{3} + c_2 \right) \\ &= -\frac{5t(2t^{\frac{3}{2}}c_1 + 3c_2)}{3}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = -\frac{5t\left(2t^{\frac{3}{2}}c_1 + 3c_2\right)}{3} \quad (1)$$

Verification of solutions

$$y = -\frac{5t\left(2t^{\frac{3}{2}}c_1 + 3c_2\right)}{3}$$

Verified OK.

9.27.6 Solving using Kovacic algorithm

Writing the ode as

$$2t^2y'' - 5ty' + 5y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 2t^2 \\ B &= -5t \\ C &= 5 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{5}{16t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 5 \\ t &= 16t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(\frac{5}{16t^2} \right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 466: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 16t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole

larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{5}{16t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = \frac{5}{16}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{5}{4} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{1}{4} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = \frac{5}{16t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = \frac{5}{16}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{5}{4} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{1}{4} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = \frac{5}{16t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{5}{4}$	$-\frac{1}{4}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{5}{4}$	$-\frac{1}{4}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_\infty^- = -\frac{1}{4}$ then

$$\begin{aligned} d &= \alpha_\infty^- - (\alpha_{c_1}^-) \\ &= -\frac{1}{4} - \left(-\frac{1}{4}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c) [\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty) [\sqrt{r}]_\infty$$

The above gives

$$\begin{aligned} \omega &= \left((-) [\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^-}{t - c_1} \right) + (-) [\sqrt{r}]_\infty \\ &= -\frac{1}{4t} + (-) (0) \\ &= -\frac{1}{4t} \\ &= -\frac{1}{4t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \tag{1A}$$

Let

$$p(t) = 1 \tag{2A}$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(-\frac{1}{4t}\right)(0) + \left(\left(\frac{1}{4t^2}\right) + \left(-\frac{1}{4t}\right)^2 - \left(\frac{5}{16t^2}\right)\right) = 0$$

$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$z_1(t) = pe^{\int \omega dt}$$

$$= e^{\int -\frac{1}{4t} dt}$$

$$= \frac{1}{t^{\frac{1}{4}}}$$

The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt}$$

$$= z_1 e^{-\int \frac{1}{2} \frac{5t}{2t^2} dt}$$

$$= z_1 e^{\frac{5 \ln(t)}{4}}$$

$$= z_1 \left(t^{\frac{5}{4}}\right)$$

Which simplifies to

$$y_1 = t$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$y_2 = y_1 \int \frac{e^{\int -\frac{5t}{2t^2} dt}}{(y_1)^2} dt$$

$$= y_1 \int \frac{e^{\frac{5 \ln(t)}{2}}}{(y_1)^2} dt$$

$$= y_1 \left(\frac{2t^{\frac{3}{2}}}{3}\right)$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1(t) + c_2 \left(t \left(\frac{2t^{\frac{3}{2}}}{3} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 t + \frac{2c_2 t^{\frac{5}{2}}}{3} \quad (1)$$

Verification of solutions

$$y = c_1 t + \frac{2c_2 t^{\frac{5}{2}}}{3}$$

Verified OK.

9.27.7 Maple step by step solution

Let's solve

$$2y''t^2 - 5ty' + 5y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{5y'}{2t} - \frac{5y}{2t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{5y'}{2t} + \frac{5y}{2t^2} = 0$$

- Multiply by denominators of the ODE

$$2y''t^2 - 5ty' + 5y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds} y(s) \right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s) \right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s) \right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$2 \left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2} \right) t^2 - 5 \frac{d}{ds}y(s) + 5y(s) = 0$$

- Simplify

$$2 \frac{d^2}{ds^2}y(s) - 7 \frac{d}{ds}y(s) + 5y(s) = 0$$

- Isolate 2nd derivative

$$\frac{d^2}{ds^2}y(s) = \frac{7 \frac{d}{ds}y(s)}{2} - \frac{5y(s)}{2}$$

- Group terms with $y(s)$ on the lhs of the ODE and the rest on the rhs of the ODE; ODE is lin

$$\frac{d^2}{ds^2}y(s) - \frac{7 \frac{d}{ds}y(s)}{2} + \frac{5y(s)}{2} = 0$$

- Characteristic polynomial of ODE

$$r^2 - \frac{7}{2}r + \frac{5}{2} = 0$$

- Factor the characteristic polynomial

$$\frac{(r-1)(2r-5)}{2} = 0$$

- Roots of the characteristic polynomial

$$r = \left(1, \frac{5}{2} \right)$$

- 1st solution of the ODE

$$y_1(s) = e^s$$

- 2nd solution of the ODE

$$y_2(s) = e^{\frac{5s}{2}}$$

- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$

- Substitute in solutions

$$y(s) = c_1 e^s + c_2 e^{\frac{5s}{2}}$$

- Change variables back using $s = \ln(t)$

$$y = c_1 t + c_2 t^{\frac{5}{2}}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 13

```
dsolve(2*t^2*diff(y(t),t$2)-5*t*diff(y(t),t)+5*y(t)=0,y(t), singsol=all)
```

$$y(t) = c_1 t + c_2 t^{\frac{5}{2}}$$

✓ Solution by Mathematica

Time used: 0.011 (sec). Leaf size: 18

```
DSolve[2*t^2*y'[t]-5*t*y'[t]+5*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t(c_2 t^{3/2} + c_1)$$

9.28 problem 43

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Internal problem ID [680]

Internal file name [OUTPUT/680_Sunday_June_05_2022_01_46_49_AM_855305/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 43.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "exact linear second order ode", "second_order_integrable_as_is", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_2nd_order, _exact, _linear, _homogeneous]]
```

$$t^2y'' + 3ty' + y = 0$$

9.28.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} + 3trt^{r-1} + t^r = 0$$

Simplifying gives

$$r(r-1)t^r + 3rt^r + t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r-1) + 3r + 1 = 0$$

Or

$$r^2 + 2r + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -1$$

$$r_2 = -1$$

Since the roots are equal, then the general solution is

$$y = c_1y_1 + c_2y_2$$

Where $y_1 = t^r$ and $y_2 = t^r \ln(t)$. Hence

$$y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t}$$

Verified OK.

9.28.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$t^2 y'' + 3ty' + y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = \frac{3}{t}$$
$$q(t) = \frac{1}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t) \tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t) dt)} dt \\ &= \int e^{-(\int \frac{3}{t} dt)} dt \\ &= \int e^{-3 \ln(t)} dt \\ &= \int \frac{1}{t^3} dt \\ &= -\frac{1}{2t^2} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned}q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{\frac{1}{t^2}}{\frac{1}{t^6}} \\ &= t^4\end{aligned}\tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned}\frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + t^4y(\tau) &= 0\end{aligned}$$

But in terms of τ

$$t^4 = \frac{1}{4\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{y(\tau)}{4\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$4\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$4\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + \tau^r = 0$$

Simplifying gives

$$4r(r-1)\tau^r + 0\tau^r + \tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$4r(r-1) + 0 + 1 = 0$$

Or

$$4r^2 - 4r + 1 = 0\tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2}$$

$$r_2 = \frac{1}{2}$$

Since the roots are equal, then the general solution is

$$y(\tau) = c_1 y_1 + c_2 y_2$$

Where $y_1 = \tau^r$ and $y_2 = \tau^r \ln(\tau)$. Hence

$$y(\tau) = c_1 \sqrt{\tau} + c_2 \sqrt{\tau} \ln(\tau)$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{\sqrt{2} \sqrt{-\frac{1}{t^2}} (c_1 - c_2 \ln(2) + c_2 \ln(-\frac{1}{t^2}))}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{2} \sqrt{-\frac{1}{t^2}} (c_1 - c_2 \ln(2) + c_2 \ln(-\frac{1}{t^2}))}{2} \quad (1)$$

Verification of solutions

$$y = \frac{\sqrt{2} \sqrt{-\frac{1}{t^2}} (c_1 - c_2 \ln(2) + c_2 \ln(-\frac{1}{t^2}))}{2}$$

Verified OK.

9.28.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$t^2 y'' + 3t y' + y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = \frac{3}{t}$$

$$q(t) = \frac{1}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c}\sqrt{q}$$

$$= \frac{\sqrt{\frac{1}{t^2}}}{c} \quad (6)$$

$$\tau'' = -\frac{1}{c\sqrt{\frac{1}{t^2}}t^3}$$

Substituting the above into (4) results in

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2}$$

$$= \frac{-\frac{1}{c\sqrt{\frac{1}{t^2}}t^3} + \frac{3}{t}\frac{\sqrt{\frac{1}{t^2}}}{c}}{\left(\frac{\sqrt{\frac{1}{t^2}}}{c}\right)^2}$$

$$= 2c$$

Therefore ode (3) now becomes

$$y(\tau)'' + p_1y(\tau)' + q_1y(\tau) = 0$$

$$\frac{d^2}{d\tau^2}y(\tau) + 2c\left(\frac{d}{d\tau}y(\tau)\right) + c^2y(\tau) = 0 \quad (7)$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{-c\tau} c_1$$

Now from (6)

$$\begin{aligned} \tau &= \int \frac{1}{c} \sqrt{q} dt \\ &= \frac{\int \sqrt{\frac{1}{t^2}} dt}{c} \\ &= \frac{\sqrt{\frac{1}{t^2}} t \ln(t)}{c} \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = \frac{c_1}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{t}$$

Verified OK.

9.28.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$t^2 y'' + 3ty' + y = 0 \tag{1}$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \tag{2}$$

Where

$$\begin{aligned} p(t) &= \frac{3}{t} \\ q(t) &= \frac{1}{t^2} \end{aligned}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} + \frac{3n}{t^2} + \frac{1}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = -1 \quad (6)$$

Substituting this value in (3) gives

$$\begin{aligned} v''(t) + \frac{v'(t)}{t} &= 0 \\ v''(t) + \frac{v'(t)}{t} &= 0 \end{aligned} \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{u}{t} \end{aligned}$$

Where $f(t) = -\frac{1}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= -\frac{1}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{1}{t} dt \\ \ln(u) &= -\ln(t) + c_1 \\ u &= e^{-\ln(t)+c_1} \\ &= \frac{c_1}{t}\end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned}v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= c_1 \ln(t) + c_2\end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\ &= \frac{c_1 \ln(t) + c_2}{t} \\ &= \frac{c_1 \ln(t) + c_2}{t}\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \ln(t) + c_2}{t} \tag{1}$$

Verification of solutions

$$y = \frac{c_1 \ln(t) + c_2}{t}$$

Verified OK.

9.28.5 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t t gives

$$\int (t^2 y'' + 3ty' + y) dt = 0$$
$$y't^2 + yt = c_1$$

Which is now solved for y .

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{1}{t}$$
$$q(t) = \frac{c_1}{t^2}$$

Hence the ode is

$$y' + \frac{y}{t} = \frac{c_1}{t^2}$$

The integrating factor μ is

$$\mu = e^{\int \frac{1}{t} dt}$$
$$= t$$

The ode becomes

$$\frac{d}{dt}(\mu y) = (\mu) \left(\frac{c_1}{t^2} \right)$$
$$\frac{d}{dt}(ty) = (t) \left(\frac{c_1}{t^2} \right)$$
$$d(ty) = \left(\frac{c_1}{t} \right) dt$$

Integrating gives

$$ty = \int \frac{c_1}{t} dt$$
$$ty = c_1 \ln(t) + c_2$$

Dividing both sides by the integrating factor $\mu = t$ results in

$$y = \frac{c_1 \ln(t)}{t} + \frac{c_2}{t}$$

which simplifies to

$$y = \frac{c_1 \ln(t) + c_2}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \ln(t) + c_2}{t} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 \ln(t) + c_2}{t}$$

Verified OK.

9.28.6 Solving as type second_order_integrable_as_is (not using ABC version)

Writing the ode as

$$t^2 y'' + 3ty' + y = 0$$

Integrating both sides of the ODE w.r.t t gives

$$\int (t^2 y'' + 3ty' + y) dt = 0$$
$$y't^2 + yt = c_1$$

Which is now solved for y .

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{1}{t}$$
$$q(t) = \frac{c_1}{t^2}$$

Hence the ode is

$$y' + \frac{y}{t} = \frac{c_1}{t^2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{1}{t} dt} \\ &= t\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{c_1}{t^2}\right) \\ \frac{d}{dt}(ty) &= (t) \left(\frac{c_1}{t^2}\right) \\ d(ty) &= \left(\frac{c_1}{t}\right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}ty &= \int \frac{c_1}{t} dt \\ ty &= c_1 \ln(t) + c_2\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t$ results in

$$y = \frac{c_1 \ln(t)}{t} + \frac{c_2}{t}$$

which simplifies to

$$y = \frac{c_1 \ln(t) + c_2}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \ln(t) + c_2}{t} \tag{1}$$

Verification of solutions

$$y = \frac{c_1 \ln(t) + c_2}{t}$$

Verified OK.

9.28.7 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' + 3ty' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t^2 \\ B &= 3t \\ C &= 1 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(-\frac{1}{4t^2}\right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 468: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{1}{4t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{1}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2}$	$\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = \frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= \frac{1}{2} - \left(\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{t - c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= \frac{1}{2t} + (-)(0) \\ &= \frac{1}{2t} \\ &= \frac{1}{2t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2\left(\frac{1}{2t}\right)(0) + \left(\left(-\frac{1}{2t^2}\right) + \left(\frac{1}{2t}\right)^2 - \left(-\frac{1}{4t^2}\right) \right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= p e^{\int \omega dt} \\ &= e^{\int \frac{1}{2t} dt} \\ &= \sqrt{t} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\&= z_1 e^{-\int \frac{1}{2} \frac{3t}{t^2} dt} \\&= z_1 e^{-\frac{3 \ln(t)}{2}} \\&= z_1 \left(\frac{1}{t^{\frac{3}{2}}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = \frac{1}{t}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{3t}{t^2} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{-3 \ln(t)}}{(y_1)^2} dt \\&= y_1 (\ln(t))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 \left(\frac{1}{t} \right) + c_2 \left(\frac{1}{t} (\ln(t)) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t} \tag{1}$$

Verification of solutions

$$y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t}$$

Verified OK.

9.28.8 Solving as exact linear second order ode

An ode of the form

$$p(t) y'' + q(t) y' + r(t) y = s(t)$$

is exact if

$$p''(t) - q'(t) + r(t) = 0 \tag{1}$$

For the given ode we have

$$p(x) = t^2$$

$$q(x) = 3t$$

$$r(x) = 1$$

$$s(x) = 0$$

Hence

$$p''(x) = 2$$

$$q'(x) = 3$$

Therefore (1) becomes

$$2 - (3) + (1) = 0$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$(p(t) y' + (q(t) - p'(t)) y)' = s(x)$$

Integrating gives

$$p(t) y' + (q(t) - p'(t)) y = \int s(t) dt$$

Substituting the above values for p, q, r, s gives

$$y't^2 + yt = c_1$$

We now have a first order ode to solve which is

$$y't^2 + yt = c_1$$

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{1}{t}$$
$$q(t) = \frac{c_1}{t^2}$$

Hence the ode is

$$y' + \frac{y}{t} = \frac{c_1}{t^2}$$

The integrating factor μ is

$$\mu = e^{\int \frac{1}{t} dt}$$
$$= t$$

The ode becomes

$$\frac{d}{dt}(\mu y) = (\mu) \left(\frac{c_1}{t^2} \right)$$
$$\frac{d}{dt}(ty) = (t) \left(\frac{c_1}{t^2} \right)$$
$$d(ty) = \left(\frac{c_1}{t} \right) dt$$

Integrating gives

$$ty = \int \frac{c_1}{t} dt$$
$$ty = c_1 \ln(t) + c_2$$

Dividing both sides by the integrating factor $\mu = t$ results in

$$y = \frac{c_1 \ln(t)}{t} + \frac{c_2}{t}$$

which simplifies to

$$y = \frac{c_1 \ln(t) + c_2}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \ln(t) + c_2}{t} \tag{1}$$

Verification of solutions

$$y = \frac{c_1 \ln(t) + c_2}{t}$$

Verified OK.

9.28.9 Maple step by step solution

Let's solve

$$y''t^2 + 3ty' + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{3y'}{t} - \frac{y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{3y'}{t} + \frac{y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 + 3ty' + y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right) t^2 + 3\frac{d}{ds}y(s) + y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) + 2\frac{d}{ds}y(s) + y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 + 2r + 1 = 0$$

- Factor the characteristic polynomial
 $(r + 1)^2 = 0$
- Root of the characteristic polynomial
 $r = -1$
- 1st solution of the ODE
 $y_1(s) = e^{-s}$
- Repeated root, multiply $y_1(s)$ by s to ensure linear independence
 $y_2(s) = s e^{-s}$
- General solution of the ODE
 $y(s) = c_1 y_1(s) + c_2 y_2(s)$
- Substitute in solutions
 $y(s) = c_1 e^{-s} + c_2 s e^{-s}$
- Change variables back using $s = \ln(t)$
 $y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t}$
- Simplify
 $y = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t}$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 14

```
dsolve(t^2*diff(y(t),t$2)+3*t*diff(y(t),t)+y(t)=0,y(t), singsol=all)
```

$$y(t) = \frac{c_2 \ln(t) + c_1}{t}$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 17

```
DSolve[t^2*y''[t]+3*t*y'[t]+y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{c_2 \log(t) + c_1}{t}$$

9.29 problem 44

9.29.1 Solving as second order euler ode ode	2636
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Internal problem ID [681]

Internal file name [OUTPUT/681_Sunday_June_05_2022_01_46_50_AM_23754413/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 44.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_Emden, _Fowler]]
```

$$4t^2y'' - 8ty' + 9y = 0$$

9.29.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$4t^2(r(r-1))t^{r-2} - 8trt^{r-1} + 9t^r = 0$$

Simplifying gives

$$4r(r-1)t^r - 8rt^r + 9t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$4r(r - 1) - 8r + 9 = 0$$

Or

$$4r^2 - 12r + 9 = 0 \quad (1)$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{3}{2}$$
$$r_2 = \frac{3}{2}$$

Since the roots are equal, then the general solution is

$$y = c_1 y_1 + c_2 y_2$$

Where $y_1 = t^r$ and $y_2 = t^r \ln(t)$. Hence

$$y = t^{\frac{3}{2}} c_1 + c_2 t^{\frac{3}{2}} \ln(t)$$

Summary

The solution(s) found are the following

$$y = t^{\frac{3}{2}} c_1 + c_2 t^{\frac{3}{2}} \ln(t) \quad (1)$$

Verification of solutions

$$y = t^{\frac{3}{2}} c_1 + c_2 t^{\frac{3}{2}} \ln(t)$$

Verified OK.

9.29.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$4t^2 y'' - 8ty' + 9y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = -\frac{2}{t}$$
$$q(t) = \frac{9}{4t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t)\tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t)dt)} dt \\ &= \int e^{-(\int -\frac{2}{t}dt)} dt \\ &= \int e^{2\ln(t)} dt \\ &= \int t^2 dt \\ &= \frac{t^3}{3} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{9}{4t^2} \\ &= \frac{9}{4t^6} \end{aligned} \quad (7)$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + \frac{9y(\tau)}{4t^6} &= 0 \end{aligned}$$

But in terms of τ

$$\frac{9}{4t^6} = \frac{1}{4\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{y(\tau)}{4\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$4\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$4\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + \tau^r = 0$$

Simplifying gives

$$4r(r-1)\tau^r + 0\tau^r + \tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$4r(r-1) + 0 + 1 = 0$$

Or

$$4r^2 - 4r + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2}$$
$$r_2 = \frac{1}{2}$$

Since the roots are equal, then the general solution is

$$y(\tau) = c_1y_1 + c_2y_2$$

Where $y_1 = \tau^r$ and $y_2 = \tau^r \ln(\tau)$. Hence

$$y(\tau) = c_1\sqrt{\tau} + c_2\sqrt{\tau} \ln(\tau)$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{\sqrt{3}\sqrt{t^3}(c_1 + c_2 \ln(t^3)) - c_2 \ln(3)}{3}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{3} \sqrt{t^3} (c_1 + c_2 \ln(t^3) - c_2 \ln(3))}{3} \quad (1)$$

Verification of solutions

$$y = \frac{\sqrt{3} \sqrt{t^3} (c_1 + c_2 \ln(t^3) - c_2 \ln(3))}{3}$$

Verified OK.

9.29.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$4t^2 y'' - 8ty' + 9y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = -\frac{2}{t}$$
$$q(t) = \frac{9}{4t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c} \sqrt{q}$$
$$= \frac{3\sqrt{\frac{1}{t^2}}}{2c} \quad (6)$$
$$\tau'' = -\frac{3}{2c\sqrt{\frac{1}{t^2}} t^3}$$

Substituting the above into (4) results in

$$\begin{aligned}
 p_1(\tau) &= \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \\
 &= \frac{-\frac{3}{2c\sqrt{\frac{1}{t^2}} t^3} - \frac{2}{t} \frac{3\sqrt{\frac{1}{t^2}}}{2c}}{\left(\frac{3\sqrt{\frac{1}{t^2}}}{2c}\right)^2} \\
 &= -2c
 \end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned}
 y(\tau)'' + p_1 y(\tau)' + q_1 y(\tau) &= 0 \\
 \frac{d^2}{d\tau^2} y(\tau) - 2c \left(\frac{d}{d\tau} y(\tau) \right) + c^2 y(\tau) &= 0 \tag{7}
 \end{aligned}$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{c\tau} c_1$$

Now from (6)

$$\begin{aligned}
 \tau &= \int \frac{1}{c} \sqrt{q} dt \\
 &= \frac{\int \frac{3\sqrt{\frac{1}{t^2}}}{2} dt}{c} \\
 &= \frac{3\sqrt{\frac{1}{t^2}} t \ln(t)}{2c}
 \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = t^{\frac{3}{2}} c_1$$

Summary

The solution(s) found are the following

$$y = t^{\frac{3}{2}} c_1 \tag{1}$$

Verification of solutions

$$y = t^{\frac{3}{2}} c_1$$

Verified OK.

9.29.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$4t^2 y'' - 8ty' + 9y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = -\frac{2}{t}$$
$$q(t) = \frac{9}{4t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variable is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} - \frac{2n}{t} + \frac{9}{4t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = \frac{3}{2} \quad (6)$$

Substituting this value in (3) gives

$$v''(t) + \frac{v'(t)}{t} = 0$$
$$v''(t) + \frac{v'(t)}{t} = 0 \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{u}{t} \end{aligned}$$

Where $f(t) = -\frac{1}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= -\frac{1}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{1}{t} dt \\ \ln(u) &= -\ln(t) + c_1 \\ u &= e^{-\ln(t)+c_1} \\ &= \frac{c_1}{t} \end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned} v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= c_1 \ln(t) + c_2 \end{aligned}$$

Hence

$$\begin{aligned} y &= v(t) t^n \\ &= (c_1 \ln(t) + c_2) t^{\frac{3}{2}} \\ &= (c_1 \ln(t) + c_2) t^{\frac{3}{2}} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = (c_1 \ln(t) + c_2) t^{\frac{3}{2}} \quad (1)$$

Verification of solutions

$$y = (c_1 \ln(t) + c_2) t^{\frac{3}{2}}$$

Verified OK.

9.29.5 Solving using Kovacic algorithm

Writing the ode as

$$4t^2y'' - 8ty' + 9y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 4t^2 \\ B &= -8t \\ C &= 9 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(-\frac{1}{4t^2}\right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 470: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{1}{4t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{1}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2}$	$\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = \frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= \frac{1}{2} - \left(\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{t - c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= \frac{1}{2t} + (-)(0) \\ &= \frac{1}{2t} \\ &= \frac{1}{2t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2\left(\frac{1}{2t}\right)(0) + \left(\left(-\frac{1}{2t^2}\right) + \left(\frac{1}{2t}\right)^2 - \left(-\frac{1}{4t^2}\right) \right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= p e^{\int \omega dt} \\ &= e^{\int \frac{1}{2t} dt} \\ &= \sqrt{t} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\&= z_1 e^{-\int \frac{1}{2} \frac{-8t}{4t^2} dt} \\&= z_1 e^{\ln(t)} \\&= z_1(t)\end{aligned}$$

Which simplifies to

$$y_1 = t^{\frac{3}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-8t}{4t^2} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{2\ln(t)}}{(y_1)^2} dt \\&= y_1(\ln(t))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 \left(t^{\frac{3}{2}}\right) + c_2 \left(t^{\frac{3}{2}}(\ln(t))\right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = t^{\frac{3}{2}} c_1 + c_2 t^{\frac{3}{2}} \ln(t) \tag{1}$$

Verification of solutions

$$y = t^{\frac{3}{2}} c_1 + c_2 t^{\frac{3}{2}} \ln(t)$$

Verified OK.

9.29.6 Maple step by step solution

Let's solve

$$4y''t^2 - 8ty' + 9y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{2y'}{t} - \frac{9y}{4t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{2y'}{t} + \frac{9y}{4t^2} = 0$$

- Multiply by denominators of the ODE

$$4y''t^2 - 8ty' + 9y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds}y(s)\right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s)\right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s)\right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$4\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}\right)t^2 - 8\frac{d}{ds}y(s) + 9y(s) = 0$$

- Simplify

$$4\frac{d^2}{ds^2}y(s) - 12\frac{d}{ds}y(s) + 9y(s) = 0$$

- Isolate 2nd derivative

$$\frac{d^2}{ds^2}y(s) = 3\frac{d}{ds}y(s) - \frac{9y(s)}{4}$$

- Group terms with $y(s)$ on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$\frac{d^2}{ds^2}y(s) - 3\frac{d}{ds}y(s) + \frac{9y(s)}{4} = 0$$
- Characteristic polynomial of ODE

$$r^2 - 3r + \frac{9}{4} = 0$$
- Factor the characteristic polynomial

$$\frac{(2r-3)^2}{4} = 0$$
- Root of the characteristic polynomial

$$r = \frac{3}{2}$$
- 1st solution of the ODE

$$y_1(s) = e^{\frac{3s}{2}}$$
- Repeated root, multiply $y_1(s)$ by s to ensure linear independence

$$y_2(s) = s e^{\frac{3s}{2}}$$
- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$
- Substitute in solutions

$$y(s) = c_1 e^{\frac{3s}{2}} + c_2 s e^{\frac{3s}{2}}$$
- Change variables back using $s = \ln(t)$

$$y = t^{\frac{3}{2}} c_1 + c_2 t^{\frac{3}{2}} \ln(t)$$
- Simplify

$$y = (c_2 \ln(t) + c_1) t^{\frac{3}{2}}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
<- LODE of Euler type successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(4*t^2*diff(y(t),t$2)-8*t*diff(y(t),t)+9*y(t)=0,y(t), singsol=all)
```

$$y(t) = (c_2 \ln(t) + c_1) t^{\frac{3}{2}}$$

✓ Solution by Mathematica

Time used: 0.016 (sec). Leaf size: 25

```
DSolve[4*t^2*y''[t]-8*t*y'[t]+9*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{2} t^{3/2} (3c_2 \log(t) + 2c_1)$$

9.30 problem 45

9.30.1 Solving as second order euler ode ode	2652
9.30.2 Solving as second order change of variable on x method 2 ode .	2654
9.30.3 Solving as second order change of variable on x method 1 ode .	2657
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Internal problem ID [682]

Internal file name [OUTPUT/682_Sunday_June_05_2022_01_46_51_AM_41135738/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.4 Repeated roots, reduction of order , page 172

Problem number: 45.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_Emden, _Fowler]]
```

$$t^2 y'' + 5ty' + 13y = 0$$

9.30.1 Solving as second order euler ode ode

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} + 5trt^{r-1} + 13t^r = 0$$

Simplifying gives

$$r(r-1)t^r + 5rt^r + 13t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r - 1) + 5r + 13 = 0$$

Or

$$r^2 + 4r + 13 = 0 \quad (1)$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -2 - 3i$$

$$r_2 = -2 + 3i$$

The roots are complex conjugate of each others. Let the roots be

$$r_1 = \alpha + i\beta$$

$$r_2 = \alpha - i\beta$$

Where in this case $\alpha = -2$ and $\beta = -3$. Hence the solution becomes

$$\begin{aligned} y &= c_1 t^{r_1} + c_2 t^{r_2} \\ &= c_1 t^{\alpha+i\beta} + c_2 t^{\alpha-i\beta} \\ &= t^\alpha (c_1 t^{i\beta} + c_2 t^{-i\beta}) \\ &= t^\alpha (c_1 e^{\ln(t^{i\beta})} + c_2 e^{\ln(t^{-i\beta})}) \\ &= t^\alpha (c_1 e^{i(\beta \ln t)} + c_2 e^{-i(\beta \ln t)}) \end{aligned}$$

Using the values for $\alpha = -2, \beta = -3$, the above becomes

$$y = t^{-2} (c_1 e^{-3i \ln(t)} + c_2 e^{3i \ln(t)})$$

Using Euler relation, the expression $c_1 e^{iA} + c_2 e^{-iA}$ is transformed to $c_1 \cos A + c_1 \sin A$ where the constants are free to change. Applying this to the above result gives

$$y = \frac{1}{t^2} (c_1 \cos(3 \ln(t)) + c_2 \sin(3 \ln(t)))$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \cos(3 \ln(t)) + c_2 \sin(3 \ln(t))}{t^2} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 \cos(3 \ln(t)) + c_2 \sin(3 \ln(t))}{t^2}$$

Verified OK.

9.30.2 Solving as second order change of variable on x method 2 ode

In normal form the ode

$$t^2 y'' + 5ty' + 13y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = \frac{5}{t}$$
$$q(t) = \frac{13}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2} y(\tau) + p_1 \left(\frac{d}{d\tau} y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t) \tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t) dt)} dt \\ &= \int e^{-(\int \frac{5}{t} dt)} dt \\ &= \int e^{-5 \ln(t)} dt \\ &= \int \frac{1}{t^5} dt \\ &= -\frac{1}{4t^4} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{\frac{13}{t^2}}{\frac{1}{t^{10}}} \\ &= 13t^8 \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + 13t^8y(\tau) &= 0 \end{aligned}$$

But in terms of τ

$$13t^8 = \frac{13}{16\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{13y(\tau)}{16\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$16\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + 13y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$16\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + 13\tau^r = 0$$

Simplifying gives

$$16r(r-1)\tau^r + 0\tau^r + 13\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$16r(r-1) + 0 + 13 = 0$$

Or

$$16r^2 - 16r + 13 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2} - \frac{3i}{4}$$

$$r_2 = \frac{1}{2} + \frac{3i}{4}$$

The roots are complex conjugate of each others. Let the roots be

$$r_1 = \alpha + i\beta$$

$$r_2 = \alpha - i\beta$$

Where in this case $\alpha = \frac{1}{2}$ and $\beta = -\frac{3}{4}$. Hence the solution becomes

$$y(\tau) = c_1\tau^{r_1} + c_2\tau^{r_2}$$

$$= c_1\tau^{\alpha+i\beta} + c_2\tau^{\alpha-i\beta}$$

$$= \tau^\alpha (c_1\tau^{i\beta} + c_2\tau^{-i\beta})$$

$$= \tau^\alpha (c_1e^{\ln(\tau^{i\beta})} + c_2e^{\ln(\tau^{-i\beta})})$$

$$= \tau^\alpha (c_1e^{i(\beta \ln \tau)} + c_2e^{-i(\beta \ln \tau)})$$

Using the values for $\alpha = \frac{1}{2}$, $\beta = -\frac{3}{4}$, the above becomes

$$y(\tau) = \tau^{\frac{1}{2}} \left(c_1 e^{-\frac{3i \ln(\tau)}{4}} + c_2 e^{\frac{3i \ln(\tau)}{4}} \right)$$

Using Euler relation, the expression $c_1e^{iA} + c_2e^{-iA}$ is transformed to $c_1 \cos A + c_1 \sin A$ where the constants are free to change. Applying this to the above result gives

$$y(\tau) = \sqrt{\tau} \left(c_1 \cos \left(\frac{3 \ln(\tau)}{4} \right) + c_2 \sin \left(\frac{3 \ln(\tau)}{4} \right) \right)$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{\sqrt{-\frac{1}{t^4}} \left(c_1 \cos \left(-\frac{3 \ln(2)}{2} + \frac{3 \ln\left(-\frac{1}{t^4}\right)}{4} \right) + c_2 \sin \left(-\frac{3 \ln(2)}{2} + \frac{3 \ln\left(-\frac{1}{t^4}\right)}{4} \right) \right)}{2}$$

Summary

The solution(s) found are the following

$$y = \frac{\sqrt{-\frac{1}{t^4}} \left(c_1 \cos \left(-\frac{3 \ln(2)}{2} + \frac{3 \ln\left(-\frac{1}{t^4}\right)}{4} \right) + c_2 \sin \left(-\frac{3 \ln(2)}{2} + \frac{3 \ln\left(-\frac{1}{t^4}\right)}{4} \right) \right)}{2} \quad (1)$$

Verification of solutions

$$y = \frac{\sqrt{-\frac{1}{t^4}} \left(c_1 \cos \left(-\frac{3 \ln(2)}{2} + \frac{3 \ln\left(-\frac{1}{t^4}\right)}{4} \right) + c_2 \sin \left(-\frac{3 \ln(2)}{2} + \frac{3 \ln\left(-\frac{1}{t^4}\right)}{4} \right) \right)}{2}$$

Verified OK.

9.30.3 Solving as second order change of variable on x method 1 ode

In normal form the ode

$$t^2 y'' + 5ty' + 13y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{5}{t}$$
$$q(t) = \frac{13}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1 \left(\frac{d}{d\tau}y(\tau) \right) + q_1 y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c} \sqrt{q}$$
$$= \frac{\sqrt{13} \sqrt{\frac{1}{t^2}}}{c} \quad (6)$$
$$\tau'' = -\frac{\sqrt{13}}{c \sqrt{\frac{1}{t^2}} t^3}$$

Substituting the above into (4) results in

$$\begin{aligned}
 p_1(\tau) &= \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \\
 &= \frac{-\frac{\sqrt{13}}{c\sqrt{\frac{1}{t^2}} t^3} + \frac{5}{t} \frac{\sqrt{13}\sqrt{\frac{1}{t^2}}}{c}}{\left(\frac{\sqrt{13}\sqrt{\frac{1}{t^2}}}{c}\right)^2} \\
 &= \frac{4c\sqrt{13}}{13}
 \end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned}
 y(\tau)'' + p_1 y(\tau)' + q_1 y(\tau) &= 0 \\
 \frac{d^2}{d\tau^2} y(\tau) + \frac{4c\sqrt{13}}{13} \left(\frac{d}{d\tau} y(\tau)\right) + c^2 y(\tau) &= 0 \tag{7}
 \end{aligned}$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{-\frac{2\sqrt{13}c\tau}{13}} \left(c_1 \cos\left(\frac{3\sqrt{13}c\tau}{13}\right) + c_2 \sin\left(\frac{3\sqrt{13}c\tau}{13}\right) \right)$$

Now from (6)

$$\begin{aligned}
 \tau &= \int \frac{1}{c} \sqrt{q} dt \\
 &= \frac{\int \sqrt{13} \sqrt{\frac{1}{t^2}} dt}{c} \\
 &= \frac{\sqrt{13} \sqrt{\frac{1}{t^2}} t \ln(t)}{c}
 \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = \frac{c_1 \cos(3 \ln(t)) + c_2 \sin(3 \ln(t))}{t^2}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \cos(3 \ln(t)) + c_2 \sin(3 \ln(t))}{t^2} \tag{1}$$

Verification of solutions

$$y = \frac{c_1 \cos(3 \ln(t)) + c_2 \sin(3 \ln(t))}{t^2}$$

Verified OK.

9.30.4 Solving as second order change of variable on y method 2 ode

In normal form the ode

$$t^2 y'' + 5ty' + 13y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{5}{t}$$
$$q(t) = \frac{13}{t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} + \frac{5n}{t} + \frac{13}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = -2 + 3i \quad (6)$$

Substituting this value in (3) gives

$$v''(t) + \left(\frac{-4 + 6i}{t} + \frac{5}{t}\right)v'(t) = 0$$
$$v''(t) + \frac{(1 + 6i)v'(t)}{t} = 0 \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{(1 + 6i)u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{(-1 - 6i)u}{t} \end{aligned}$$

Where $f(t) = \frac{-1-6i}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= \frac{-1 - 6i}{t} dt \\ \int \frac{1}{u} du &= \int \frac{-1 - 6i}{t} dt \\ \ln(u) &= (-1 - 6i) \ln(t) + c_1 \\ u &= e^{(-1-6i)\ln(t)+c_1} \\ &= c_1 e^{(-1-6i)\ln(t)} \end{aligned}$$

Which simplifies to

$$u(t) = \frac{c_1 t^{-6i}}{t}$$

Now that $u(t)$ is known, then

$$\begin{aligned} v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= \frac{ic_1 t^{-6i}}{6} + c_2 \end{aligned}$$

Hence

$$\begin{aligned} y &= v(t) t^n \\ &= \left(\frac{ic_1 t^{-6i}}{6} + c_2 \right) t^{-2+3i} \\ &= c_2 t^{-2+3i} + \frac{ic_1 t^{-2-3i}}{6} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \left(\frac{ic_1 t^{-6i}}{6} + c_2 \right) t^{-2+3i} \quad (1)$$

Verification of solutions

$$y = \left(\frac{ic_1 t^{-6i}}{6} + c_2 \right) t^{-2+3i}$$

Verified OK.

9.30.5 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' + 5ty' + 13y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t^2 \\ B &= 5t \\ C &= 13 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-37}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -37 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(-\frac{37}{4t^2} \right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 472: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole

larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{37}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{37}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} + 3i \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} - 3i \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{37}{4t^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{37}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} + 3i \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} - 3i \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{37}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2} + 3i$	$\frac{1}{2} - 3i$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2} + 3i$	$\frac{1}{2} - 3i$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_\infty^- = \frac{1}{2} - 3i$ then

$$\begin{aligned} d &= \alpha_\infty^- - (\alpha_{c_1}^-) \\ &= \frac{1}{2} - 3i - \left(\frac{1}{2} - 3i \right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_\infty$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^-}{t - c_1} \right) + (-)[\sqrt{r}]_\infty \\ &= \frac{\frac{1}{2} - 3i}{t} + (-)(0) \\ &= \frac{\frac{1}{2} - 3i}{t} \\ &= \frac{\frac{1}{2} - 3i}{t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \tag{1A}$$

Let

$$p(t) = 1 \tag{2A}$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(\frac{\frac{1}{2} - 3i}{t}\right) (0) + \left(\left(\frac{-\frac{1}{2} + 3i}{t^2}\right) + \left(\frac{\frac{1}{2} - 3i}{t}\right)^2 - \left(-\frac{37}{4t^2}\right)\right) = 0$$

$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \frac{\frac{1}{2} - 3i}{t} dt} \\ &= t^{\frac{1}{2} - 3i} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{5t}{t^2} dt} \\ &= z_1 e^{-\frac{5 \ln(t)}{2}} \\ &= z_1 \left(\frac{1}{t^{\frac{5}{2}}}\right) \end{aligned}$$

Which simplifies to

$$y_1 = t^{-2-3i}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{5t}{t^2} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{-5 \ln(t)}}{(y_1)^2} dt \\ &= y_1 \left(-\frac{it^{6i}}{6}\right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (t^{-2-3i}) + c_2 \left(t^{-2-3i} \left(-\frac{it^{6i}}{6} \right) \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 t^{-2-3i} - \frac{ic_2 t^{-2+3i}}{6} \quad (1)$$

Verification of solutions

$$y = c_1 t^{-2-3i} - \frac{ic_2 t^{-2+3i}}{6}$$

Verified OK.

9.30.6 Maple step by step solution

Let's solve

$$y''t^2 + 5ty' + 13y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{5y'}{t} - \frac{13y}{t^2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{5y'}{t} + \frac{13y}{t^2} = 0$$

- Multiply by denominators of the ODE

$$y''t^2 + 5ty' + 13y = 0$$

- Make a change of variables

$$s = \ln(t)$$

- Substitute the change of variables back into the ODE

- Calculate the 1st derivative of y with respect to t , using the chain rule

$$y' = \left(\frac{d}{ds} y(s) \right) s'(t)$$

- Compute derivative

$$y' = \frac{\frac{d}{ds}y(s)}{t}$$

- Calculate the 2nd derivative of y with respect to t , using the chain rule

$$y'' = \left(\frac{d^2}{ds^2}y(s) \right) s'(t)^2 + s''(t) \left(\frac{d}{ds}y(s) \right)$$

- Compute derivative

$$y'' = \frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2}$$

Substitute the change of variables back into the ODE

$$\left(\frac{\frac{d^2}{ds^2}y(s)}{t^2} - \frac{\frac{d}{ds}y(s)}{t^2} \right) t^2 + 5 \frac{d}{ds}y(s) + 13y(s) = 0$$

- Simplify

$$\frac{d^2}{ds^2}y(s) + 4 \frac{d}{ds}y(s) + 13y(s) = 0$$

- Characteristic polynomial of ODE

$$r^2 + 4r + 13 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{(-4) \pm (\sqrt{-36})}{2}$$

- Roots of the characteristic polynomial

$$r = (-2 - 3I, -2 + 3I)$$

- 1st solution of the ODE

$$y_1(s) = e^{-2s} \cos(3s)$$

- 2nd solution of the ODE

$$y_2(s) = e^{-2s} \sin(3s)$$

- General solution of the ODE

$$y(s) = c_1 y_1(s) + c_2 y_2(s)$$

- Substitute in solutions

$$y(s) = c_1 e^{-2s} \cos(3s) + c_2 e^{-2s} \sin(3s)$$

- Change variables back using $s = \ln(t)$

$$y = \frac{c_1 \cos(3 \ln(t))}{t^2} + \frac{c_2 \sin(3 \ln(t))}{t^2}$$

- Simplify

$$y = \frac{c_1 \cos(3 \ln(t))}{t^2} + \frac{c_2 \sin(3 \ln(t))}{t^2}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
<- LODE of Euler type successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 23

```
dsolve(t^2*diff(y(t),t$2)+5*t*diff(y(t),t)+13*y(t)=0,y(t), singsol=all)
```

$$y(t) = \frac{c_1 \sin(3 \ln(t)) + c_2 \cos(3 \ln(t))}{t^2}$$

✓ Solution by Mathematica

Time used: 0.024 (sec). Leaf size: 26

```
DSolve[t^2*y''[t]+5*t*y'[t]+13*y[t]==0,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{c_2 \cos(3 \log(t)) + c_1 \sin(3 \log(t))}{t^2}$$

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10.1 problem 1

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Internal file name [OUTPUT/683_Sunday_June_05_2022_01_46_52_AM_10580708/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters. page 190

Problem number: 1.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$y'' - 5y' + 6y = 2e^t$$

10.1.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = -5, C = 6, f(t) = 2e^t$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' - 5y' + 6y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = -5, C = 6$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} - 5\lambda e^{\lambda t} + 6e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 - 5\lambda + 6 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -5, C = 6$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{5}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-5^2 - (4)(1)(6)} \\ &= \frac{5}{2} \pm \frac{1}{2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= \frac{5}{2} + \frac{1}{2} \\ \lambda_2 &= \frac{5}{2} - \frac{1}{2} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= 3 \\ \lambda_2 &= 2 \end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned} y &= c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t} \\ y &= c_1 e^{(3)t} + c_2 e^{(2)t} \end{aligned}$$

Or

$$y = c_1 e^{3t} + c_2 e^{2t}$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 e^{3t} + c_2 e^{2t}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$2e^t$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$\{e^t\}$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\{e^{2t}, e^{3t}\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$y_p = A_1 e^t$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$2A_1 e^t = 2e^t$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = 1]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = e^t$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 e^{3t} + c_2 e^{2t}) + (e^t) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{3t} + c_2 e^{2t} + e^t \quad (1)$$

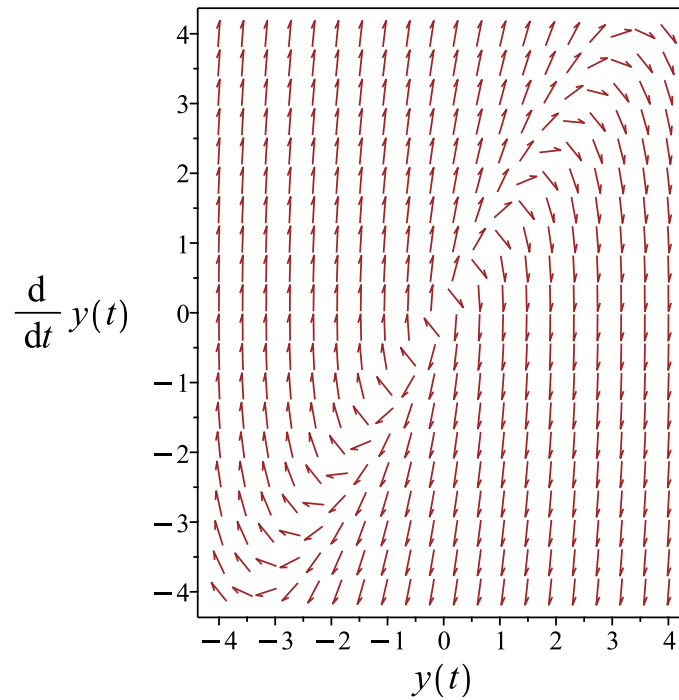


Figure 470: Slope field plot

Verification of solutions

$$y = c_1 e^{3t} + c_2 e^{2t} + e^t$$

Verified OK.

10.1.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 5y' + 6y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = -5 \tag{3}$$

$$C = 6$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{4} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \frac{z(t)}{4} \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 474: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{4}$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = e^{-\frac{t}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-5}{1} dt} \\ &= z_1 e^{\frac{5t}{2}} \\ &= z_1 \left(e^{\frac{5t}{2}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = e^{2t}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-5}{1} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{5t}}{(y_1)^2} dt \\ &= y_1 (e^t)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{2t}) + c_2 (e^{2t}(e^t))\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' - 5y' + 6y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^{2t} + c_2 e^{3t}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$2e^t$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$\{e^t\}$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\{e^{2t}, e^{3t}\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$y_p = A_1 e^t$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$2A_1 e^t = 2e^t$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = 1]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = e^t$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 e^{2t} + c_2 e^{3t}) + (e^t) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{2t} + c_2 e^{3t} + e^t \quad (1)$$

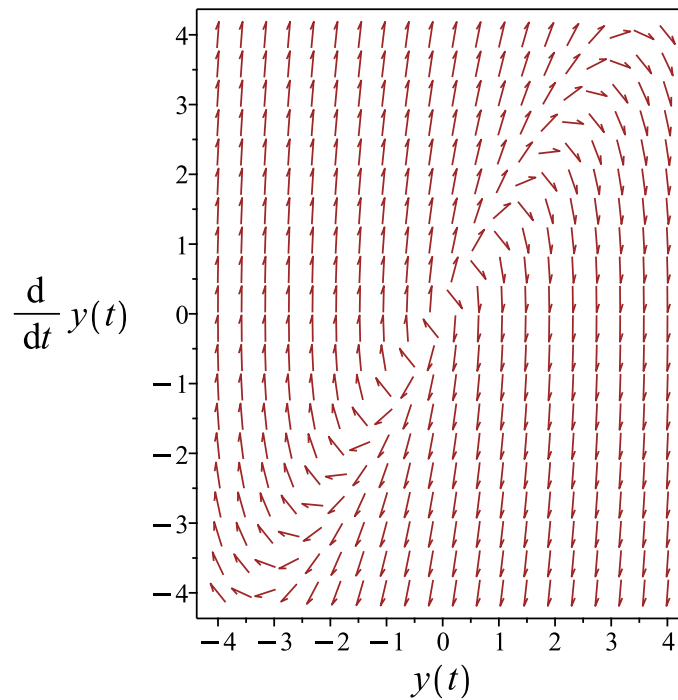


Figure 471: Slope field plot

Verification of solutions

$$y = c_1 e^{2t} + c_2 e^{3t} + e^t$$

Verified OK.

10.1.3 Maple step by step solution

Let's solve

$$y'' - 5y' + 6y = 2e^t$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of homogeneous ODE

$$r^2 - 5r + 6 = 0$$

- Factor the characteristic polynomial

$$(r - 2)(r - 3) = 0$$

- Roots of the characteristic polynomial

$$r = (2, 3)$$

- 1st solution of the homogeneous ODE

$$y_1(t) = e^{2t}$$

- 2nd solution of the homogeneous ODE

$$y_2(t) = e^{3t}$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 e^{2t} + c_2 e^{3t} + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t), y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t), y_2(t))} dt \right), f(t) = 2e^t \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} e^{2t} & e^{3t} \\ 2e^{2t} & 3e^{3t} \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = e^{5t}$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = -2e^{2t} \left(\int e^{-t} dt \right) + 2e^{3t} \left(\int e^{-2t} dt \right)$$

- Compute integrals

$$y_p(t) = e^t$$

- Substitute particular solution into general solution to ODE

$$y = c_1 e^{2t} + c_2 e^{3t} + e^t$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(diff(y(t),t$2)-5*diff(y(t),t)+6*y(t) = 2*exp(t),y(t), singsol=all)
```

$$y(t) = c_2 e^{2t} + c_1 e^{3t} + e^t$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 25

```
DSolve[y''[t]-5*y'[t]+6*y[t] == 2*Exp[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^t (c_1 e^t + c_2 e^{2t} + 1)$$

10.2 problem 2

10.2.1 Solving as second order linear constant coeff ode	2680
10.2.2 Solving using Kovacic algorithm	2683
10.2.3 Maple step by step solution	2688

Internal problem ID [684]

Internal file name [OUTPUT/684_Sunday_June_05_2022_01_46_53_AM_70729508/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 2.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$y'' - y' - 2y = 2e^{-t}$$

10.2.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = -1, C = -2, f(t) = 2e^{-t}$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$.
 y_h is the solution to

$$y'' - y' - 2y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = -1, C = -2$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} - \lambda e^{\lambda t} - 2 e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 - \lambda - 2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -1, C = -2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{1}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-1^2 - (4)(1)(-2)} \\ &= \frac{1}{2} \pm \frac{3}{2} \end{aligned}$$

Hence

$$\lambda_1 = \frac{1}{2} + \frac{3}{2}$$

$$\lambda_2 = \frac{1}{2} - \frac{3}{2}$$

Which simplifies to

$$\lambda_1 = 2$$

$$\lambda_2 = -1$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t}$$

$$y = c_1 e^{(2)t} + c_2 e^{(-1)t}$$

Or

$$y = c_1 e^{2t} + c_2 e^{-t}$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 e^{2t} + c_2 e^{-t}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$2e^{-t}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[e^{-t}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\{e^{-t}, e^{2t}\}$$

Since e^{-t} is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$[te^{-t}]$$

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.

$$y_p = A_1 t e^{-t}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-3A_1 e^{-t} = 2e^{-t}$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = -\frac{2}{3} \right]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = -\frac{2t e^{-t}}{3}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 e^{2t} + c_2 e^{-t}) + \left(-\frac{2t e^{-t}}{3} \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{2t} + c_2 e^{-t} - \frac{2t e^{-t}}{3} \quad (1)$$

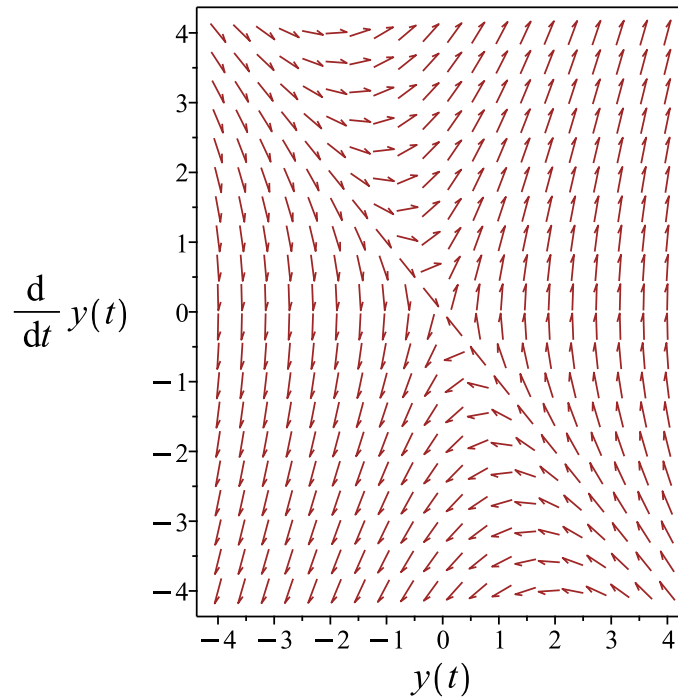


Figure 472: Slope field plot

Verification of solutions

$$y = c_1 e^{2t} + c_2 e^{-t} - \frac{2t e^{-t}}{3}$$

Verified OK.

10.2.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' - y' - 2y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= 1 \\B &= -1 \\C &= -2\end{aligned}\tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t)\tag{4}$$

Where r is given by

$$\begin{aligned}r &= \frac{s}{t} \\&= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}\end{aligned}\tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{9}{4}\tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned}s &= 9 \\t &= 4\end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \frac{9z(t)}{4}\tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 476: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{9}{4}$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = e^{-\frac{3t}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}
 y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\
 &= z_1 e^{-\int \frac{1}{2} \frac{-1}{1} dt} \\
 &= z_1 e^{\frac{t}{2}} \\
 &= z_1 \left(e^{\frac{t}{2}} \right)
 \end{aligned}$$

Which simplifies to

$$y_1 = e^{-t}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{1}{1} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^t}{(y_1)^2} dt \\ &= y_1 \left(\frac{e^{3t}}{3} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-t}) + c_2 \left(e^{-t} \left(\frac{e^{3t}}{3} \right) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$.
 y_h is the solution to

$$y'' - y' - 2y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^{-t} + \frac{c_2 e^{2t}}{3}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$2e^{-t}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[e^{-t}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ \frac{e^{2t}}{3}, e^{-t} \right\}$$

Since e^{-t} is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$[te^{-t}]$$

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.

$$y_p = A_1 t e^{-t}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-3A_1 e^{-t} = 2e^{-t}$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = -\frac{2}{3} \right]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = -\frac{2t e^{-t}}{3}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(c_1 e^{-t} + \frac{c_2 e^{2t}}{3} \right) + \left(-\frac{2t e^{-t}}{3} \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{-t} + \frac{c_2 e^{2t}}{3} - \frac{2t e^{-t}}{3} \quad (1)$$

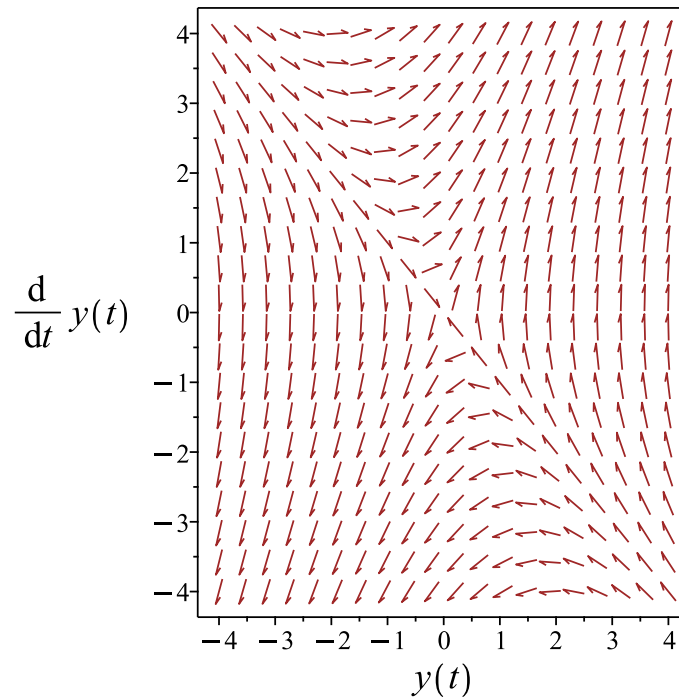


Figure 473: Slope field plot

Verification of solutions

$$y = c_1 e^{-t} + \frac{c_2 e^{2t}}{3} - \frac{2t e^{-t}}{3}$$

Verified OK.

10.2.3 Maple step by step solution

Let's solve

$$y'' - y' - 2y = 2e^{-t}$$

- Highest derivative means the order of the ODE is 2
 y''
- Characteristic polynomial of homogeneous ODE

$$r^2 - r - 2 = 0$$

- Factor the characteristic polynomial

$$(r + 1)(r - 2) = 0$$

- Roots of the characteristic polynomial

$$r = (-1, 2)$$

- 1st solution of the homogeneous ODE

$$y_1(t) = e^{-t}$$

- 2nd solution of the homogeneous ODE

$$y_2(t) = e^{2t}$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 e^{-t} + c_2 e^{2t} + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t), y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t), y_2(t))} dt \right), f(t) = 2e^{-t} \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} e^{-t} & e^{2t} \\ -e^{-t} & 2e^{2t} \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = 3e^t$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = -\frac{2e^{-t}(\int 1 dt)}{3} + \frac{2e^{2t}(\int e^{-3t} dt)}{3}$$

- Compute integrals

$$y_p(t) = -\frac{2(1+3t)e^{-t}}{9}$$

- Substitute particular solution into general solution to ODE

$$y = c_1 e^{-t} + c_2 e^{2t} - \frac{2(1+3t)e^{-t}}{9}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]  
trying a double symmetry of the form [xi=0, eta=F(x)]  
-> Try solving first the homogeneous part of the ODE  
    checking if the LODE has constant coefficients  
    <- constant coefficients successful  
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 24

```
dsolve(diff(y(t),t$2)-diff(y(t),t)-2*y(t) = 2*exp(-t),y(t), singsol=all)
```

$$y(t) = \frac{(-2t + 3c_1)e^{-t}}{3} + c_2e^{2t}$$

✓ Solution by Mathematica

Time used: 0.035 (sec). Leaf size: 32

```
DSolve[y''[t]-y'[t]-2*y[t] == 2*Exp[-t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{9}e^{-t}(-6t + 9c_2e^{3t} - 2 + 9c_1)$$

10.3 problem 3

10.3.1 Solving as second order linear constant coeff ode	2691
10.3.2 Solving as linear second order ode solved by an integrating factor ode	2694
10.3.3 Solving using Kovacic algorithm	2696
10.3.4 Maple step by step solution	2701

Internal problem ID [685]

Internal file name [OUTPUT/685_Sunday_June_05_2022_01_46_54_AM_20733644/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters. page 190

Problem number: 3.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$y'' + 2y' + y = 3e^{-t}$$

10.3.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = 2, C = 1, f(t) = 3e^{-t}$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + 2y' + y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = 2, C = 1$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + 2\lambda e^{\lambda t} + e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + 2\lambda + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 2, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(2)^2 - (4)(1)(1)} \\ &= -1 \end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = 1$. Therefore the solution is

$$y = c_1 e^{-t} + c_2 t e^{-t} \quad (1)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 e^{-t} + c_2 t e^{-t}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$3e^{-t}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{e^{-t}\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\{t e^{-t}, e^{-t}\}$$

Since e^{-t} is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$[\{te^{-t}\}]$$

Since te^{-t} is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$[\{t^2e^{-t}\}]$$

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.

$$y_p = A_1t^2e^{-t}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$2A_1e^{-t} = 3e^{-t}$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = \frac{3}{2}\right]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = \frac{3t^2e^{-t}}{2}$$

Therefore the general solution is

$$\begin{aligned}y &= y_h + y_p \\ &= (c_1e^{-t} + c_2e^{-t}t) + \left(\frac{3t^2e^{-t}}{2}\right)\end{aligned}$$

Which simplifies to

$$y = (c_2t + c_1)e^{-t} + \frac{3t^2e^{-t}}{2}$$

Summary

The solution(s) found are the following

$$y = (c_2 t + c_1) e^{-t} + \frac{3t^2 e^{-t}}{2} \quad (1)$$

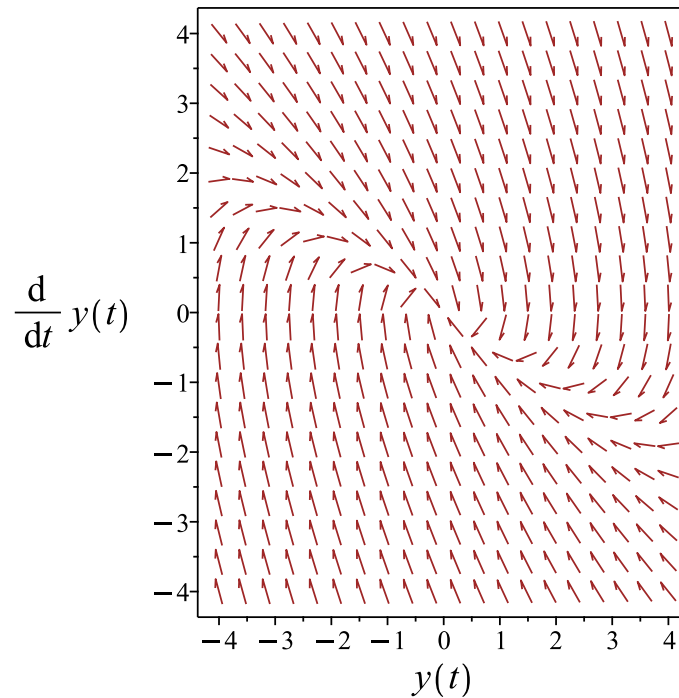


Figure 474: Slope field plot

Verification of solutions

$$y = (c_2 t + c_1) e^{-t} + \frac{3t^2 e^{-t}}{2}$$

Verified OK.

10.3.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t) y' + \frac{(p(t))^2 + p'(t)}{2} y = f(t)$$

Where $p(t) = 2$. Therefore, there is an integrating factor given by

$$\begin{aligned}M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int 2 dx} \\ &= e^t\end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned}(M(x)y)'' &= 3e^{-t}e^t \\ (ye^t)'' &= 3e^{-t}e^t\end{aligned}$$

Integrating once gives

$$(ye^t)' = 3t + c_1$$

Integrating again gives

$$(ye^t) = \frac{t(3t + 2c_1)}{2} + c_2$$

Hence the solution is

$$y = \frac{\frac{t(3t+2c_1)}{2} + c_2}{e^t}$$

Or

$$y = c_1 t e^{-t} + \frac{3t^2 e^{-t}}{2} + c_2 e^{-t}$$

Summary

The solution(s) found are the following

$$y = c_1 t e^{-t} + \frac{3t^2 e^{-t}}{2} + c_2 e^{-t} \quad (1)$$

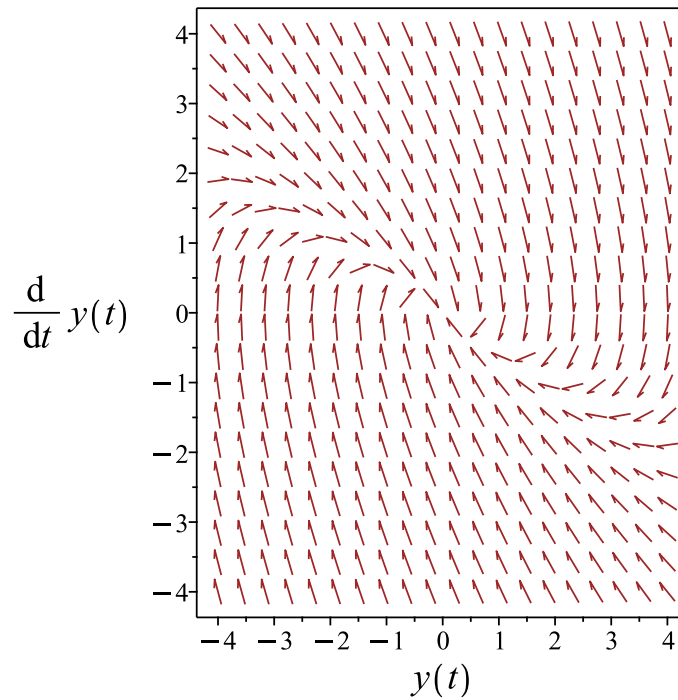


Figure 475: Slope field plot

Verification of solutions

$$y = c_1 t e^{-t} + \frac{3t^2 e^{-t}}{2} + c_2 e^{-t}$$

Verified OK.

10.3.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 2y' + y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 2 \\ C &= 1 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 0$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(t) = 0 \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 478: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{2}{1} dt} \\ &= z_1 e^{-t} \\ &= z_1 (e^{-t})\end{aligned}$$

Which simplifies to

$$y_1 = e^{-t}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{2}{1} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{-2t}}{(y_1)^2} dt \\ &= y_1(t)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-t}) + c_2 (e^{-t} t)\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + 2y' + y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^{-t} + c_2 e^{-t} t$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$3e^{-t}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[{\{e^{-t}\}}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\{t e^{-t}, e^{-t}\}$$

Since e^{-t} is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$[{\{t e^{-t}\}}]$$

Since $t e^{-t}$ is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$[{\{t^2 e^{-t}\}}]$$

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.

$$y_p = A_1 t^2 e^{-t}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$2A_1 e^{-t} = 3 e^{-t}$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = \frac{3}{2} \right]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = \frac{3t^2 e^{-t}}{2}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 e^{-t} + c_2 e^{-t}t) + \left(\frac{3t^2 e^{-t}}{2} \right) \end{aligned}$$

Which simplifies to

$$y = (c_2 t + c_1) e^{-t} + \frac{3t^2 e^{-t}}{2}$$

Summary

The solution(s) found are the following

$$y = (c_2 t + c_1) e^{-t} + \frac{3t^2 e^{-t}}{2} \quad (1)$$

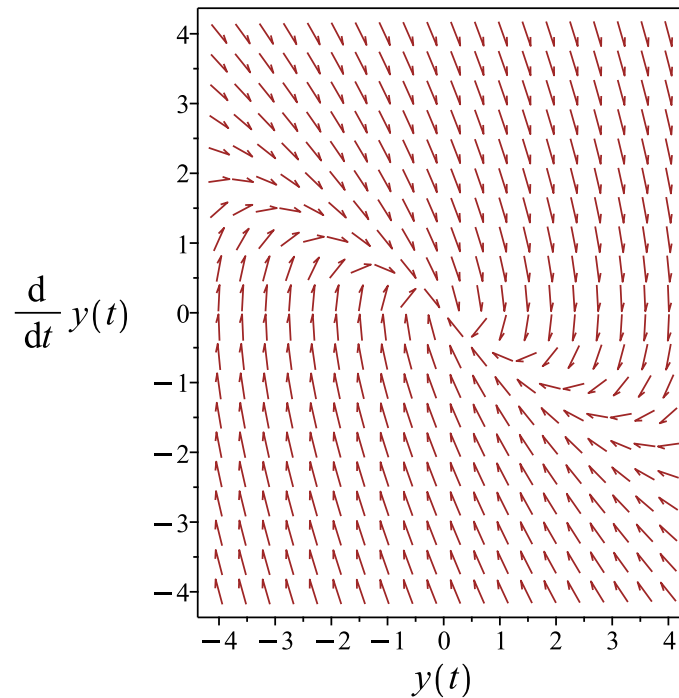


Figure 476: Slope field plot

Verification of solutions

$$y = (c_2 t + c_1) e^{-t} + \frac{3t^2 e^{-t}}{2}$$

Verified OK.

10.3.4 Maple step by step solution

Let's solve

$$y'' + 2y' + y = 3e^{-t}$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of homogeneous ODE

$$r^2 + 2r + 1 = 0$$

- Factor the characteristic polynomial

$$(r + 1)^2 = 0$$

- Root of the characteristic polynomial

$$r = -1$$

- 1st solution of the homogeneous ODE

$$y_1(t) = e^{-t}$$

- Repeated root, multiply $y_1(t)$ by t to ensure linear independence

$$y_2(t) = t e^{-t}$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 e^{-t} + c_2 e^{-t} t + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t), y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t), y_2(t))} dt \right), f(t) = 3e^{-t} \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} e^{-t} & t e^{-t} \\ -e^{-t} & -t e^{-t} + e^{-t} \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = e^{-2t}$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = -3e^{-t} \left(\int t dt - \left(\int 1 dt \right) t \right)$$

- Compute integrals

$$y_p(t) = \frac{3t^2 e^{-t}}{2}$$

- Substitute particular solution into general solution to ODE

$$y = c_2 e^{-t} t + c_1 e^{-t} + \frac{3t^2 e^{-t}}{2}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]  
trying a double symmetry of the form [xi=0, eta=F(x)]  
-> Try solving first the homogeneous part of the ODE  
    checking if the LODE has constant coefficients  
    <- constant coefficients successful  
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(diff(y(t),t$2)+2*diff(y(t),t)+y(t) = 3*exp(-t),y(t), singsol=all)
```

$$y(t) = e^{-t} \left(c_2 + c_1 t + \frac{3}{2} t^2 \right)$$

✓ Solution by Mathematica

Time used: 0.024 (sec). Leaf size: 29

```
DSolve[y''[t]+2*y'[t]+y[t] == 3*Exp[-t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{2} e^{-t} (3t^2 + 2c_2 t + 2c_1)$$

10.4 problem 4

10.4.1 Solving as second order linear constant coeff ode	2704
10.4.2 Solving as linear second order ode solved by an integrating factor ode	2707
10.4.3 Solving using Kovacic algorithm	2709
10.4.4 Maple step by step solution	2714

Internal problem ID [686]

Internal file name [OUTPUT/686_Sunday_June_05_2022_01_46_55_AM_1679039/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters. page 190

Problem number: 4.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$4y'' - 4y' + y = 16e^{\frac{t}{2}}$$

10.4.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 4, B = -4, C = 1, f(t) = 16e^{\frac{t}{2}}$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$4y'' - 4y' + y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 4, B = -4, C = 1$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$4\lambda^2 e^{\lambda t} - 4\lambda e^{\lambda t} + e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$4\lambda^2 - 4\lambda + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 4, B = -4, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{4}{(2)(4)} \pm \frac{1}{(2)(4)} \sqrt{(-4)^2 - (4)(4)(1)} \\ &= \frac{1}{2} \end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = -\frac{1}{2}$. Therefore the solution is

$$y = c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}} \quad (1)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$16 e^{\frac{t}{2}}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$\left[\left\{ e^{\frac{t}{2}} \right\} \right]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ t e^{\frac{t}{2}}, e^{\frac{t}{2}} \right\}$$

Since $e^{\frac{t}{2}}$ is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$\left[\left\{ t e^{\frac{t}{2}} \right\} \right]$$

Since $t e^{\frac{t}{2}}$ is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$\left[\left\{ t^2 e^{\frac{t}{2}} \right\} \right]$$

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.

$$y_p = A_1 t^2 e^{\frac{t}{2}}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$8A_1 e^{\frac{t}{2}} = 16 e^{\frac{t}{2}}$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = 2]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = 2t^2 e^{\frac{t}{2}}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}} \right) + \left(2t^2 e^{\frac{t}{2}} \right) \end{aligned}$$

Which simplifies to

$$y = e^{\frac{t}{2}}(c_2 t + c_1) + 2t^2 e^{\frac{t}{2}}$$

Summary

The solution(s) found are the following

$$y = e^{\frac{t}{2}}(c_2 t + c_1) + 2t^2 e^{\frac{t}{2}} \quad (1)$$

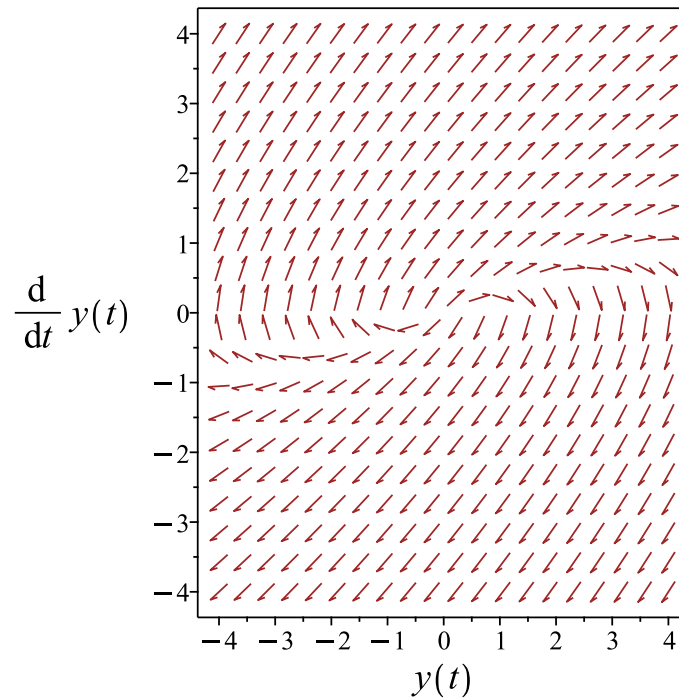


Figure 477: Slope field plot

Verification of solutions

$$y = e^{\frac{t}{2}}(c_2 t + c_1) + 2t^2 e^{\frac{t}{2}}$$

Verified OK.

10.4.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t)y' + \frac{(p(t))^2 + p'(t)}{2}y = f(t)$$

Where $p(t) = -1$. Therefore, there is an integrating factor given by

$$\begin{aligned} M(x) &= e^{\frac{1}{2} \int p \, dx} \\ &= e^{\int -1 \, dx} \\ &= e^{-\frac{t}{2}} \end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$(M(x)y)'' = 4e^{-\frac{t}{2}}e^{\frac{t}{2}}$$

$$\left(e^{-\frac{t}{2}}y\right)'' = 4e^{-\frac{t}{2}}e^{\frac{t}{2}}$$

Integrating once gives

$$\left(e^{-\frac{t}{2}}y\right)' = 4t + c_1$$

Integrating again gives

$$\left(e^{-\frac{t}{2}}y\right) = t(c_1 + 2t) + c_2$$

Hence the solution is

$$y = \frac{t(c_1 + 2t) + c_2}{e^{-\frac{t}{2}}}$$

Or

$$y = c_1 t e^{\frac{t}{2}} + 2t^2 e^{\frac{t}{2}} + c_2 e^{\frac{t}{2}}$$

Summary

The solution(s) found are the following

$$y = c_1 t e^{\frac{t}{2}} + 2t^2 e^{\frac{t}{2}} + c_2 e^{\frac{t}{2}} \quad (1)$$

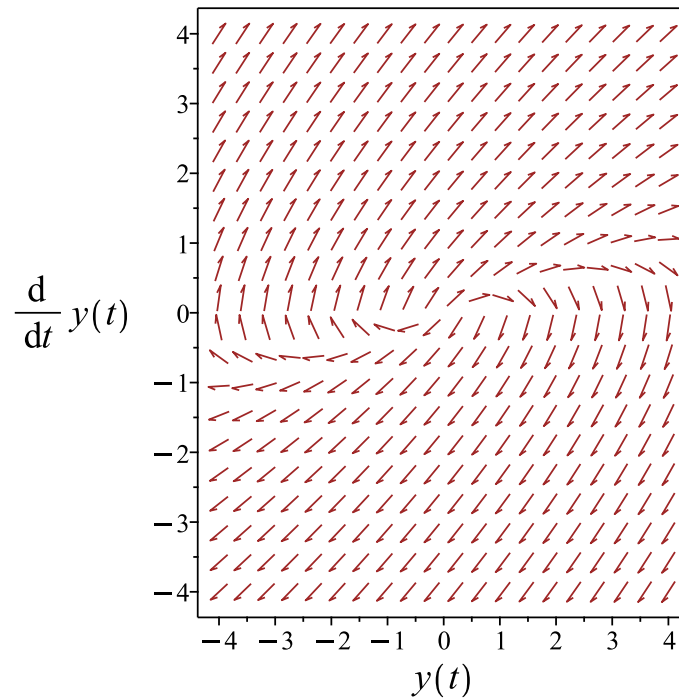


Figure 478: Slope field plot

Verification of solutions

$$y = c_1 t e^{\frac{t}{2}} + 2t^2 e^{\frac{t}{2}} + c_2 e^{\frac{t}{2}}$$

Verified OK.

10.4.3 Solving using Kovacic algorithm

Writing the ode as

$$4y'' - 4y' + y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 4 \\ B &= -4 \\ C &= 1 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 0 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = 0 \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 480: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-4}{4} dt} \\ &= z_1 e^{\frac{t}{2}} \\ &= z_1 \left(e^{\frac{t}{2}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = e^{\frac{t}{2}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-4}{4} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^t}{(y_1)^2} dt \\ &= y_1(t)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{\frac{t}{2}} \right) + c_2 \left(e^{\frac{t}{2}}(t) \right)\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$4y'' - 4y' + y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$16 e^{\frac{t}{2}}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$\left[\left\{ e^{\frac{t}{2}} \right\} \right]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ t e^{\frac{t}{2}}, e^{\frac{t}{2}} \right\}$$

Since $e^{\frac{t}{2}}$ is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$\left[\left\{ t e^{\frac{t}{2}} \right\} \right]$$

Since $t e^{\frac{t}{2}}$ is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$\left[\left\{ t^2 e^{\frac{t}{2}} \right\} \right]$$

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.

$$y_p = A_1 t^2 e^{\frac{t}{2}}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$8A_1 e^{\frac{t}{2}} = 16 e^{\frac{t}{2}}$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = 2]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = 2t^2 e^{\frac{t}{2}}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}} \right) + \left(2t^2 e^{\frac{t}{2}} \right) \end{aligned}$$

Which simplifies to

$$y = e^{\frac{t}{2}}(c_2 t + c_1) + 2t^2 e^{\frac{t}{2}}$$

Summary

The solution(s) found are the following

$$y = e^{\frac{t}{2}}(c_2 t + c_1) + 2t^2 e^{\frac{t}{2}} \quad (1)$$

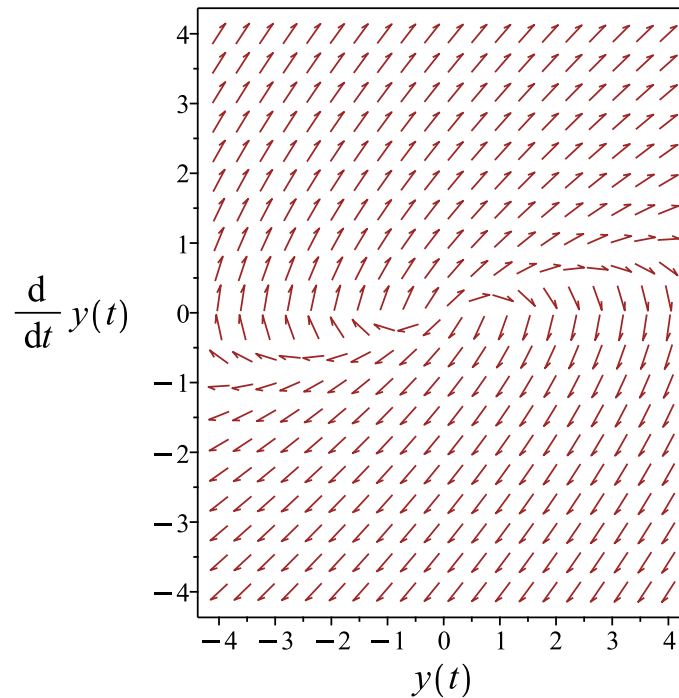


Figure 479: Slope field plot

Verification of solutions

$$y = e^{\frac{t}{2}}(c_2 t + c_1) + 2t^2 e^{\frac{t}{2}}$$

Verified OK.

10.4.4 Maple step by step solution

Let's solve

$$4y'' - 4y' + y = 16e^{\frac{t}{2}}$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = y' - \frac{y}{4} + 4e^{\frac{t}{2}}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - y' + \frac{y}{4} = 4e^{\frac{t}{2}}$$

- Characteristic polynomial of homogeneous ODE

$$r^2 - r + \frac{1}{4} = 0$$

- Factor the characteristic polynomial

$$\frac{(2r-1)^2}{4} = 0$$

- Root of the characteristic polynomial

$$r = \frac{1}{2}$$

- 1st solution of the homogeneous ODE

$$y_1(t) = e^{\frac{t}{2}}$$

- Repeated root, multiply $y_1(t)$ by t to ensure linear independence

$$y_2(t) = t e^{\frac{t}{2}}$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 e^{\frac{t}{2}} + c_2 t e^{\frac{t}{2}} + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t),y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t),y_2(t))} dt \right), f(t) = 4e^{\frac{t}{2}} \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} e^{\frac{t}{2}} & t e^{\frac{t}{2}} \\ \frac{e^{\frac{t}{2}}}{2} & e^{\frac{t}{2}} + \frac{t e^{\frac{t}{2}}}{2} \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = e^t$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = -4e^{\frac{t}{2}} \left(\int t dt - \left(\int 1 dt \right) t \right)$$

- Compute integrals

$$y_p(t) = 2t^2 e^{\frac{t}{2}}$$

- Substitute particular solution into general solution to ODE

$$y = c_2 t e^{\frac{t}{2}} + 2t^2 e^{\frac{t}{2}} + c_1 e^{\frac{t}{2}}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]  
trying a double symmetry of the form [xi=0, eta=F(x)]  
-> Try solving first the homogeneous part of the ODE  
    checking if the LODE has constant coefficients  
    <- constant coefficients successful  
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(4*diff(y(t),t$2)-4*diff(y(t),t)+y(t) = 16*exp(t/2),y(t), singsol=all)
```

$$y(t) = e^{\frac{t}{2}}(c_1 t + 2t^2 + c_2)$$

✓ Solution by Mathematica

Time used: 0.021 (sec). Leaf size: 25

```
DSolve[4*y''[t]-4*y'[t]+y[t]== 16*Exp[t/2],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{t/2}(2t^2 + c_2 t + c_1)$$

10.5 problem 5

10.5.1 Solving as second order linear constant coeff ode	2717
10.5.2 Solving using Kovacic algorithm	2722
10.5.3 Maple step by step solution	2727

Internal problem ID [687]

Internal file name [OUTPUT/687_Sunday_June_05_2022_01_46_56_AM_81877378/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters. page 190

Problem number: 5.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$y'' + y = \tan(t)$$

10.5.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = 0, C = 1, f(t) = \tan(t)$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = 0, C = 1$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(1)} \\ &= \pm i \end{aligned}$$

Hence

$$\lambda_1 = +i$$

$$\lambda_2 = -i$$

Which simplifies to

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha t} (c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$y = e^0 (c_1 \cos(t) + c_2 \sin(t))$$

Or

$$y = c_1 \cos(t) + c_2 \sin(t)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 \cos(t) + c_2 \sin(t)$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \cos(t)$$

$$y_2 = \sin(t)$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \tag{3}$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(t) & \sin(t) \\ \frac{d}{dt}(\cos(t)) & \frac{d}{dt}(\sin(t)) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(t) & \sin(t) \\ -\sin(t) & \cos(t) \end{vmatrix}$$

Therefore

$$W = (\cos(t))(\cos(t)) - (\sin(t))(-\sin(t))$$

Which simplifies to

$$W = \cos(t)^2 + \sin(t)^2$$

Which simplifies to

$$W = 1$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\sin(t) \tan(t)}{1} dt$$

Which simplifies to

$$u_1 = - \int \sin(t) \tan(t) dt$$

Hence

$$u_1 = \sin(t) - \ln(\sec(t) + \tan(t))$$

And Eq. (3) becomes

$$u_2 = \int \frac{\cos(t) \tan(t)}{1} dt$$

Which simplifies to

$$u_2 = \int \sin(t) dt$$

Hence

$$u_2 = -\cos(t)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = (\sin(t) - \ln(\sec(t) + \tan(t))) \cos(t) - \sin(t) \cos(t)$$

Which simplifies to

$$y_p(t) = -\cos(t) \ln(\sec(t) + \tan(t))$$

Therefore the general solution is

$$\begin{aligned}y &= y_h + y_p \\ &= (c_1 \cos(t) + c_2 \sin(t)) + (-\cos(t) \ln(\sec(t) + \tan(t)))\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(t) + c_2 \sin(t) - \cos(t) \ln(\sec(t) + \tan(t)) \quad (1)$$

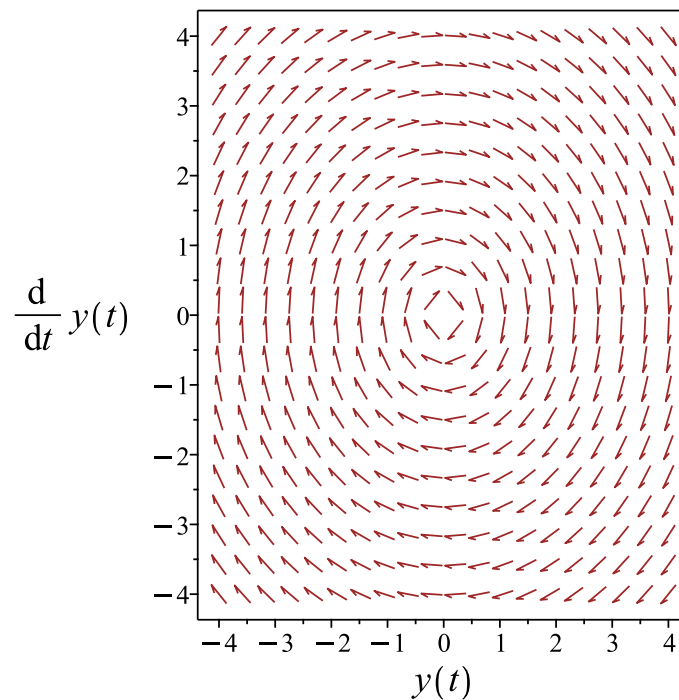


Figure 480: Slope field plot

Verification of solutions

$$y = c_1 \cos(t) + c_2 \sin(t) - \cos(t) \ln(\sec(t) + \tan(t))$$

Verified OK.

10.5.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 0 \quad (3)$$

$$C = 1$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \quad (6)$$

Comparing the above to (5) shows that

$$s = -1$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(t) = -z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 482: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos(t)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}y_1 &= z_1 \\ &= \cos(t)\end{aligned}$$

Which simplifies to

$$y_1 = \cos(t)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Since $B = 0$ then the above becomes

$$\begin{aligned}y_2 &= y_1 \int \frac{1}{y_1^2} dt \\ &= \cos(t) \int \frac{1}{\cos(t)^2} dt \\ &= \cos(t) (\tan(t))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (\cos(t)) + c_2 (\cos(t) (\tan(t)))\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 \cos(t) + c_2 \sin(t)$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \cos(t)$$

$$y_2 = \sin(t)$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(t) & \sin(t) \\ \frac{d}{dt}(\cos(t)) & \frac{d}{dt}(\sin(t)) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(t) & \sin(t) \\ -\sin(t) & \cos(t) \end{vmatrix}$$

Therefore

$$W = (\cos(t))(\cos(t)) - (\sin(t))(-\sin(t))$$

Which simplifies to

$$W = \cos(t)^2 + \sin(t)^2$$

Which simplifies to

$$W = 1$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\sin(t) \tan(t)}{1} dt$$

Which simplifies to

$$u_1 = - \int \sin(t) \tan(t) dt$$

Hence

$$u_1 = \sin(t) - \ln(\sec(t) + \tan(t))$$

And Eq. (3) becomes

$$u_2 = \int \frac{\cos(t) \tan(t)}{1} dt$$

Which simplifies to

$$u_2 = \int \sin(t) dt$$

Hence

$$u_2 = -\cos(t)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = (\sin(t) - \ln(\sec(t) + \tan(t))) \cos(t) - \sin(t) \cos(t)$$

Which simplifies to

$$y_p(t) = -\cos(t) \ln(\sec(t) + \tan(t))$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 \cos(t) + c_2 \sin(t)) + (-\cos(t) \ln(\sec(t) + \tan(t))) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(t) + c_2 \sin(t) - \cos(t) \ln(\sec(t) + \tan(t)) \quad (1)$$

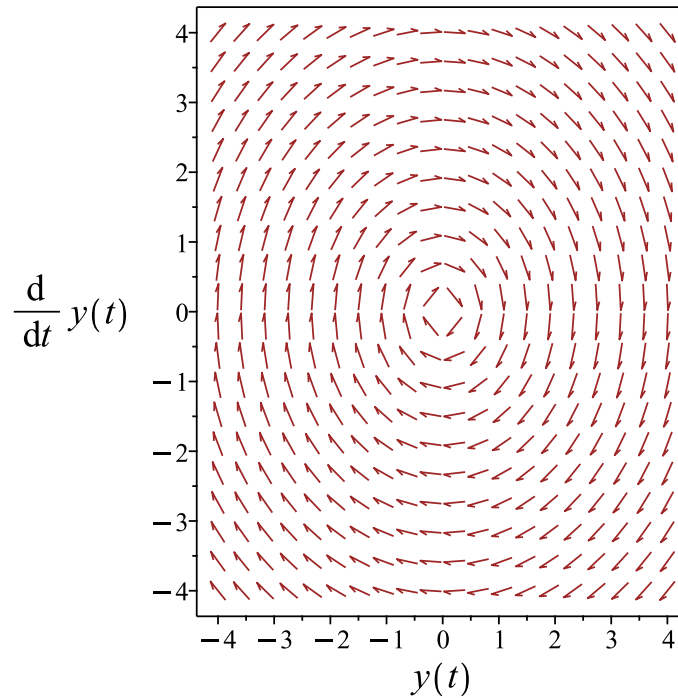


Figure 481: Slope field plot

Verification of solutions

$$y = c_1 \cos(t) + c_2 \sin(t) - \cos(t) \ln(\sec(t) + \tan(t))$$

Verified OK.

10.5.3 Maple step by step solution

Let's solve

$$y'' + y = \tan(t)$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of homogeneous ODE

$$r^2 + 1 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{0 \pm (\sqrt{-4})}{2}$$
- Roots of the characteristic polynomial

$$r = (-I, I)$$
- 1st solution of the homogeneous ODE

$$y_1(t) = \cos(t)$$
- 2nd solution of the homogeneous ODE

$$y_2(t) = \sin(t)$$
- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$
- Substitute in solutions of the homogeneous ODE

$$y = c_1 \cos(t) + c_2 \sin(t) + y_p(t)$$
- Find a particular solution $y_p(t)$ of the ODE
 - Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t), y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t), y_2(t))} dt \right), f(t) = \tan(t) \right]$$
 - Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} \cos(t) & \sin(t) \\ -\sin(t) & \cos(t) \end{bmatrix}$$
 - Compute Wronskian

$$W(y_1(t), y_2(t)) = 1$$
 - Substitute functions into equation for $y_p(t)$

$$y_p(t) = -\cos(t) \left(\int \sin(t) \tan(t) dt \right) + \sin(t) \left(\int \cos(t) dt \right)$$
 - Compute integrals

$$y_p(t) = -\cos(t) \ln(\sec(t) + \tan(t))$$
- Substitute particular solution into general solution to ODE

$$y = c_1 \cos(t) + c_2 \sin(t) - \cos(t) \ln(\sec(t) + \tan(t))$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]  
trying a double symmetry of the form [xi=0, eta=F(x)]  
-> Try solving first the homogeneous part of the ODE  
    checking if the LODE has constant coefficients  
    <- constant coefficients successful  
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 23

```
dsolve(diff(y(t),t$2)+y(t) = tan(t),y(t), singsol=all)
```

$$y(t) = c_2 \sin(t) + \cos(t) c_1 - \cos(t) \ln(\sec(t) + \tan(t))$$

✓ Solution by Mathematica

Time used: 0.04 (sec). Leaf size: 23

```
DSolve[y''[t]+y[t] == Tan[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \cos(t)(-\operatorname{arctanh}(\sin(t))) + c_1 \cos(t) + c_2 \sin(t)$$

10.6 problem 6

10.6.1 Solving as second order linear constant coeff ode	2730
10.6.2 Solving using Kovacic algorithm	2735
10.6.3 Maple step by step solution	2741

Internal problem ID [688]

Internal file name [OUTPUT/688_Sunday_June_05_2022_01_46_58_AM_10201183/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 6.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$y'' + 9y = 9 \sec(3t)^2$$

10.6.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = 0, C = 9, f(t) = 9 \sec(3t)^2$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$.
 y_h is the solution to

$$y'' + 9y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = 0, C = 9$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + 9 e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + 9 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 9$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(9)} \\ &= \pm 3i \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= +3i \\ \lambda_2 &= -3i \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= 3i \\ \lambda_2 &= -3i \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 3$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha t} (c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$y = e^0 (c_1 \cos(3t) + c_2 \sin(3t))$$

Or

$$y = c_1 \cos(3t) + c_2 \sin(3t)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 \cos(3t) + c_2 \sin(3t)$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \cos(3t)$$

$$y_2 = \sin(3t)$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(3t) & \sin(3t) \\ \frac{d}{dt}(\cos(3t)) & \frac{d}{dt}(\sin(3t)) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(3t) & \sin(3t) \\ -3 \sin(3t) & 3 \cos(3t) \end{vmatrix}$$

Therefore

$$W = (\cos(3t))(3 \cos(3t)) - (\sin(3t))(-3 \sin(3t))$$

Which simplifies to

$$W = 3 \cos(3t)^2 + 3 \sin(3t)^2$$

Which simplifies to

$$W = 3$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{9 \sin(3t) \sec(3t)^2}{3} dt$$

Which simplifies to

$$u_1 = - \int 3 \sec(3t) \tan(3t) dt$$

Hence

$$u_1 = - \sec(3t)$$

And Eq. (3) becomes

$$u_2 = \int \frac{9 \cos(3t) \sec(3t)^2}{3} dt$$

Which simplifies to

$$u_2 = \int 3 \sec(3t) dt$$

Hence

$$u_2 = \ln(\sec(3t) + \tan(3t))$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = - \sec(3t) \cos(3t) + \ln(\sec(3t) + \tan(3t)) \sin(3t)$$

Which simplifies to

$$y_p(t) = -1 + \ln(\sec(3t) + \tan(3t)) \sin(3t)$$

Therefore the general solution is

$$\begin{aligned}y &= y_h + y_p \\ &= (c_1 \cos(3t) + c_2 \sin(3t)) + (-1 + \ln(\sec(3t) + \tan(3t)) \sin(3t))\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(3t) + c_2 \sin(3t) - 1 + \ln(\sec(3t) + \tan(3t)) \sin(3t) \quad (1)$$

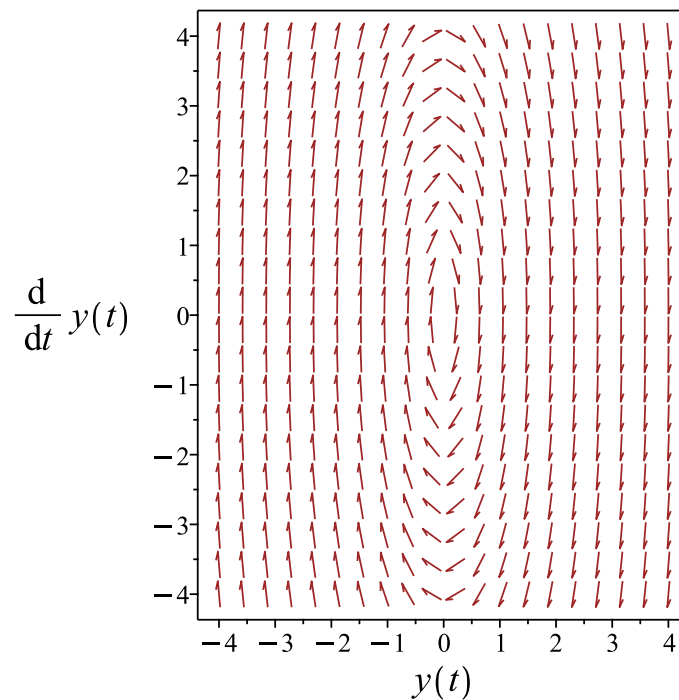


Figure 482: Slope field plot

Verification of solutions

$$y = c_1 \cos(3t) + c_2 \sin(3t) - 1 + \ln(\sec(3t) + \tan(3t)) \sin(3t)$$

Verified OK.

10.6.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 9y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 0 \quad (3)$$

$$C = 9$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-9}{1} \quad (6)$$

Comparing the above to (5) shows that

$$s = -9$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(t) = -9z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 484: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -9$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos(3t)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}y_1 &= z_1 \\ &= \cos(3t)\end{aligned}$$

Which simplifies to

$$y_1 = \cos(3t)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Since $B = 0$ then the above becomes

$$\begin{aligned}y_2 &= y_1 \int \frac{1}{y_1^2} dt \\ &= \cos(3t) \int \frac{1}{\cos(3t)^2} dt \\ &= \cos(3t) \left(\frac{\tan(3t)}{3} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1(\cos(3t)) + c_2 \left(\cos(3t) \left(\frac{\tan(3t)}{3} \right) \right)\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$.
 y_h is the solution to

$$y'' + 9y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 \cos(3t) + \frac{c_2 \sin(3t)}{3}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \cos(3t)$$

$$y_2 = \frac{\sin(3t)}{3}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(3t) & \frac{\sin(3t)}{3} \\ \frac{d}{dt}(\cos(3t)) & \frac{d}{dt}\left(\frac{\sin(3t)}{3}\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(3t) & \frac{\sin(3t)}{3} \\ -3 \sin(3t) & \cos(3t) \end{vmatrix}$$

Therefore

$$W = (\cos(3t))(\cos(3t)) - \left(\frac{\sin(3t)}{3}\right)(-3\sin(3t))$$

Which simplifies to

$$W = \cos(3t)^2 + \sin(3t)^2$$

Which simplifies to

$$W = 1$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{3 \sin(3t) \sec(3t)^2}{1} dt$$

Which simplifies to

$$u_1 = - \int 3 \sec(3t) \tan(3t) dt$$

Hence

$$u_1 = - \sec(3t)$$

And Eq. (3) becomes

$$u_2 = \int \frac{9 \cos(3t) \sec(3t)^2}{1} dt$$

Which simplifies to

$$u_2 = \int 9 \sec(3t) dt$$

Hence

$$u_2 = 3 \ln(\sec(3t) + \tan(3t))$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = - \sec(3t) \cos(3t) + \ln(\sec(3t) + \tan(3t)) \sin(3t)$$

Which simplifies to

$$y_p(t) = -1 + \ln(\sec(3t) + \tan(3t)) \sin(3t)$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(c_1 \cos(3t) + \frac{c_2 \sin(3t)}{3} \right) + (-1 + \ln(\sec(3t) + \tan(3t)) \sin(3t)) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(3t) + \frac{c_2 \sin(3t)}{3} - 1 + \ln(\sec(3t) + \tan(3t)) \sin(3t) \quad (1)$$

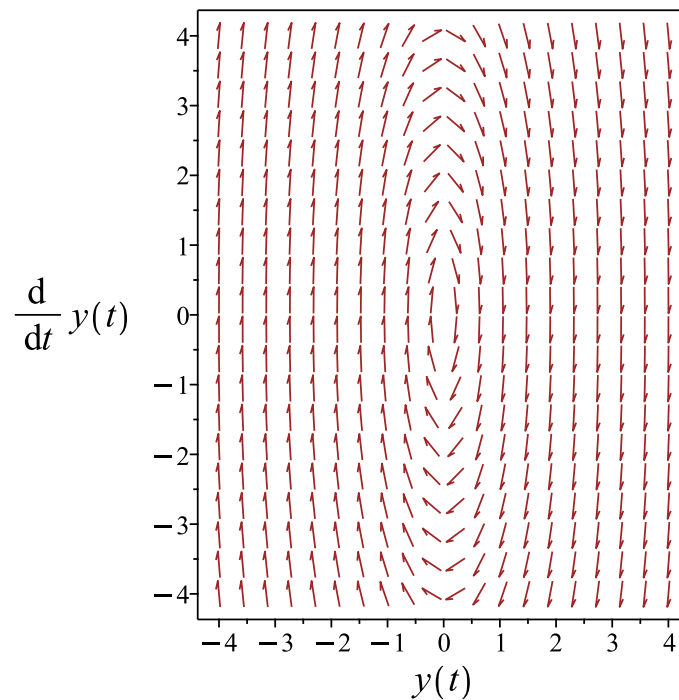


Figure 483: Slope field plot

Verification of solutions

$$y = c_1 \cos(3t) + \frac{c_2 \sin(3t)}{3} - 1 + \ln(\sec(3t) + \tan(3t)) \sin(3t)$$

Verified OK.

10.6.3 Maple step by step solution

Let's solve

$$y'' + 9y = 9 \sec(3t)^2$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of homogeneous ODE

$$r^2 + 9 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{0 \pm (\sqrt{-36})}{2}$$

- Roots of the characteristic polynomial

$$r = (-3I, 3I)$$

- 1st solution of the homogeneous ODE

$$y_1(t) = \cos(3t)$$

- 2nd solution of the homogeneous ODE

$$y_2(t) = \sin(3t)$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 \cos(3t) + c_2 \sin(3t) + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t), y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t), y_2(t))} dt \right), f(t) = 9 \sec(3t)^2 \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} \cos(3t) & \sin(3t) \\ -3 \sin(3t) & 3 \cos(3t) \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = 3$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = -3 \cos(3t) \left(\int \sec(3t) \tan(3t) dt \right) + 3 \sin(3t) \left(\int \sec(3t) dt \right)$$

- Compute integrals

$$y_p(t) = -1 + \ln(\sec(3t) + \tan(3t)) \sin(3t)$$

- Substitute particular solution into general solution to ODE

$$y = c_1 \cos(3t) + c_2 \sin(3t) - 1 + \ln(\sec(3t) + \tan(3t)) \sin(3t)$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 33

```
dsolve(diff(y(t),t$2)+9*y(t) = 9*sec(3*t)^2,y(t), singsol=all)
```

$$y(t) = c_2 \sin(3t) + c_1 \cos(3t) + \ln(\sec(3t) + \tan(3t)) \sin(3t) - 1$$

✓ Solution by Mathematica

Time used: 0.117 (sec). Leaf size: 31

```
DSolve[y''[t]+9*y[t] == 9*Sec[3*t]^2,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow c_1 \cos(3t) + \sin(3t) \coth^{-1}(\sin(3t)) + c_2 \sin(3t) - 1$$

10.7 problem 7

10.7.1 Solving as second order linear constant coeff ode	2743
10.7.2 Solving as linear second order ode solved by an integrating factor ode	2747
10.7.3 Solving using Kovacic algorithm	2749
10.7.4 Maple step by step solution	2755

Internal problem ID [689]

Internal file name [OUTPUT/689_Sunday_June_05_2022_01_46_59_AM_21493213/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters. page 190

Problem number: 7.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$y'' + 4y' + 4y = \frac{e^{-2t}}{t^2}$$

10.7.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = 4, C = 4, f(t) = \frac{e^{-2t}}{t^2}$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + 4y' + 4y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = 4, C = 4$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + 4\lambda e^{\lambda t} + 4e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + 4\lambda + 4 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 4, C = 4$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-4}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(4)^2 - (4)(1)(4)} \\ &= -2 \end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = 2$. Therefore the solution is

$$y = c_1 e^{-2t} + c_2 t e^{-2t} \quad (1)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 e^{-2t} + c_2 t e^{-2t}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned} y_1 &= e^{-2t} \\ y_2 &= t e^{-2t} \end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} e^{-2t} & t e^{-2t} \\ \frac{d}{dt}(e^{-2t}) & \frac{d}{dt}(t e^{-2t}) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} e^{-2t} & t e^{-2t} \\ -2e^{-2t} & e^{-2t} - 2t e^{-2t} \end{vmatrix}$$

Therefore

$$W = (e^{-2t})(e^{-2t} - 2t e^{-2t}) - (t e^{-2t})(-2 e^{-2t})$$

Which simplifies to

$$W = e^{-4t}$$

Which simplifies to

$$W = e^{-4t}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{e^{-4t}}{t}}{e^{-4t}} dt$$

Which simplifies to

$$u_1 = - \int \frac{1}{t} dt$$

Hence

$$u_1 = - \ln(t)$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{e^{-4t}}{t^2}}{e^{-4t}} dt$$

Which simplifies to

$$u_2 = \int \frac{1}{t^2} dt$$

Hence

$$u_2 = -\frac{1}{t}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -\ln(t) e^{-2t} - e^{-2t}$$

Which simplifies to

$$y_p(t) = e^{-2t}(-1 - \ln(t))$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 e^{-2t} + c_2 t e^{-2t}) + (e^{-2t}(-1 - \ln(t))) \end{aligned}$$

Which simplifies to

$$y = e^{-2t}(c_2 t + c_1) + e^{-2t}(-1 - \ln(t))$$

Summary

The solution(s) found are the following

$$y = e^{-2t}(c_2 t + c_1) + e^{-2t}(-1 - \ln(t)) \quad (1)$$

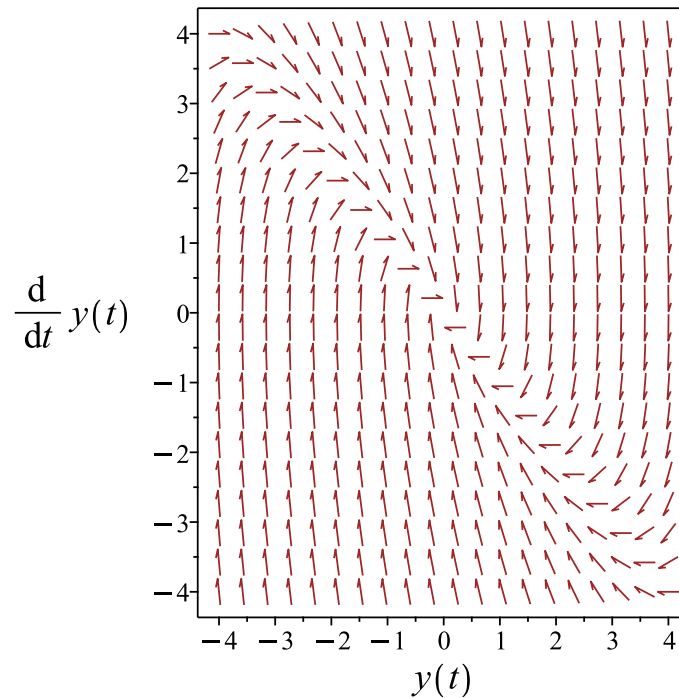


Figure 484: Slope field plot

Verification of solutions

$$y = e^{-2t}(c_2 t + c_1) + e^{-2t}(-1 - \ln(t))$$

Verified OK.

10.7.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t)y' + \frac{(p(t))^2 + p'(t)}{2}y = f(t)$$

Where $p(t) = 4$. Therefore, there is an integrating factor given by

$$\begin{aligned} M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int 4 dx} \\ &= e^{2t} \end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$(M(x)y)'' = \frac{e^{2t}e^{-2t}}{t^2}$$
$$(y e^{2t})'' = \frac{e^{2t}e^{-2t}}{t^2}$$

Integrating once gives

$$(y e^{2t})' = -\frac{1}{t} + c_1$$

Integrating again gives

$$(y e^{2t}) = c_1 t - \ln(t) + c_2$$

Hence the solution is

$$y = \frac{c_1 t - \ln(t) + c_2}{e^{2t}}$$

Or

$$y = c_1 t e^{-2t} + e^{-2t} c_2 - \ln(t) e^{-2t}$$

Summary

The solution(s) found are the following

$$y = c_1 t e^{-2t} + e^{-2t} c_2 - \ln(t) e^{-2t} \quad (1)$$

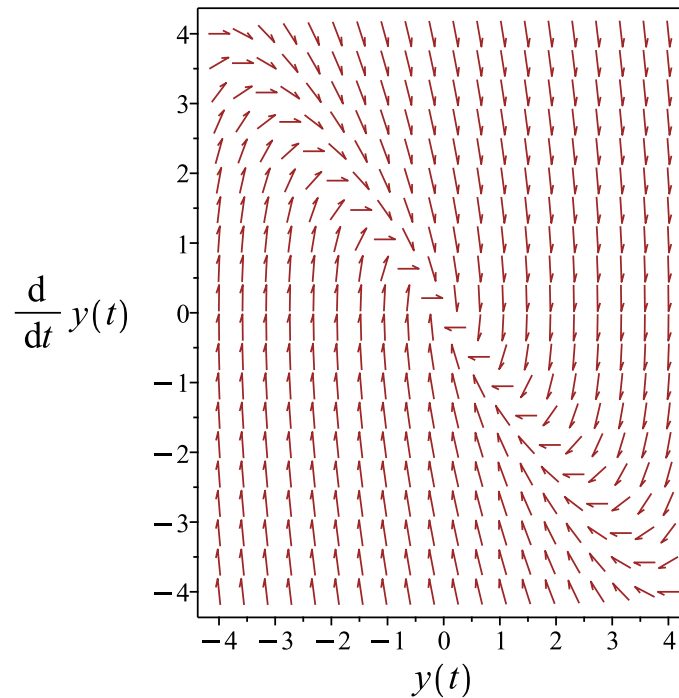


Figure 485: Slope field plot

Verification of solutions

$$y = c_1 t e^{-2t} + e^{-2t} c_2 - \ln(t) e^{-2t}$$

Verified OK.

10.7.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 4y' + 4y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 4 \\ C &= 4 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$r = \frac{s}{t} = \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \tag{6}$$

Comparing the above to (5) shows that

$$s = 0 \\ t = 1$$

Therefore eq. (4) becomes

$$z''(t) = 0 \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 486: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{4}{1} dt} \\ &= z_1 e^{-2t} \\ &= z_1 (e^{-2t}) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-2t}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{4}{1} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{-4t}}{(y_1)^2} dt \\ &= y_1(t) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{-2t}) + c_2 (e^{-2t}(t))\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + 4y' + 4y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^{-2t} + c_2 t e^{-2t}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned}y_1 &= e^{-2t} \\ y_2 &= t e^{-2t}\end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \tag{3}$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} e^{-2t} & t e^{-2t} \\ \frac{d}{dt}(e^{-2t}) & \frac{d}{dt}(t e^{-2t}) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} e^{-2t} & t e^{-2t} \\ -2 e^{-2t} & e^{-2t} - 2t e^{-2t} \end{vmatrix}$$

Therefore

$$W = (e^{-2t})(e^{-2t} - 2t e^{-2t}) - (t e^{-2t})(-2 e^{-2t})$$

Which simplifies to

$$W = e^{-4t}$$

Which simplifies to

$$W = e^{-4t}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{e^{-4t}}{t}}{e^{-4t}} dt$$

Which simplifies to

$$u_1 = - \int \frac{1}{t} dt$$

Hence

$$u_1 = - \ln(t)$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{e^{-4t}}{t^2}}{e^{-4t}} dt$$

Which simplifies to

$$u_2 = \int \frac{1}{t^2} dt$$

Hence

$$u_2 = -\frac{1}{t}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -\ln(t) e^{-2t} - e^{-2t}$$

Which simplifies to

$$y_p(t) = e^{-2t}(-1 - \ln(t))$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 e^{-2t} + c_2 t e^{-2t}) + (e^{-2t}(-1 - \ln(t))) \end{aligned}$$

Which simplifies to

$$y = e^{-2t}(c_2 t + c_1) + e^{-2t}(-1 - \ln(t))$$

Summary

The solution(s) found are the following

$$y = e^{-2t}(c_2 t + c_1) + e^{-2t}(-1 - \ln(t)) \quad (1)$$

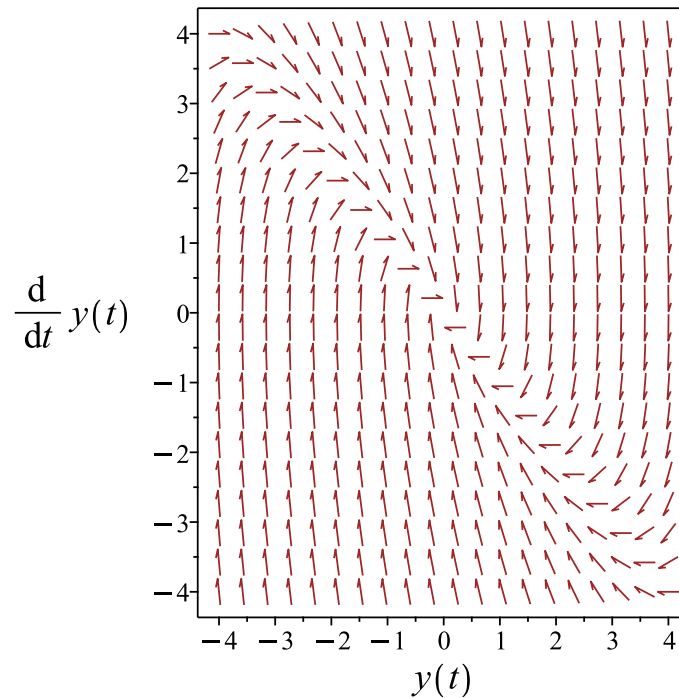


Figure 486: Slope field plot

Verification of solutions

$$y = e^{-2t}(c_2 t + c_1) + e^{-2t}(-1 - \ln(t))$$

Verified OK.

10.7.4 Maple step by step solution

Let's solve

$$y'' + 4y' + 4y = \frac{e^{-2t}}{t^2}$$

- Highest derivative means the order of the ODE is 2
- y''
- Characteristic polynomial of homogeneous ODE
 $r^2 + 4r + 4 = 0$
- Factor the characteristic polynomial
 $(r + 2)^2 = 0$
- Root of the characteristic polynomial

$$r = -2$$

- 1st solution of the homogeneous ODE

$$y_1(t) = e^{-2t}$$

- Repeated root, multiply $y_1(t)$ by t to ensure linear independence

$$y_2(t) = t e^{-2t}$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 e^{-2t} + c_2 t e^{-2t} + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t), y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t), y_2(t))} dt \right), f(t) = \frac{e^{-2t}}{t^2} \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} e^{-2t} & t e^{-2t} \\ -2e^{-2t} & e^{-2t} - 2t e^{-2t} \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = e^{-4t}$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = e^{-2t} \left(- \left(\int \frac{1}{t} dt \right) + \left(\int \frac{1}{t^2} dt \right) t \right)$$

- Compute integrals

$$y_p(t) = -e^{-2t}(1 + \ln(t))$$

- Substitute particular solution into general solution to ODE

$$y = c_1 e^{-2t} + c_2 t e^{-2t} - e^{-2t}(1 + \ln(t))$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]  
trying a double symmetry of the form [xi=0, eta=F(x)]  
-> Try solving first the homogeneous part of the ODE  
    checking if the LODE has constant coefficients  
    <- constant coefficients successful  
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(diff(y(t),t$2)+4*diff(y(t),t)+4*y(t) = t^(-2)*exp(-2*t),y(t), singsol=all)
```

$$y(t) = e^{-2t}(-1 + c_1 t - \ln(t) + c_2)$$

✓ Solution by Mathematica

Time used: 0.028 (sec). Leaf size: 23

```
DSolve[y''[t]+4*y'[t]+4*y[t] == t^(-2)*Exp[-2*t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{-2t}(-\log(t) + c_2 t - 1 + c_1)$$

10.8 problem 8

10.8.1 Solving as second order linear constant coeff ode	2758
10.8.2 Solving using Kovacic algorithm	2763
10.8.3 Maple step by step solution	2769

Internal problem ID [690]

Internal file name [OUTPUT/690_Sunday_June_05_2022_01_47_00_AM_95123719/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters. page 190

Problem number: 8.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$y'' + 4y = 3 \csc(2t)$$

10.8.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = 0, C = 4, f(t) = 3 \csc(2t)$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + 4y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = 0, C = 4$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + 4 e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + 4 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 4$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(4)} \\ &= \pm 2i \end{aligned}$$

Hence

$$\lambda_1 = +2i$$

$$\lambda_2 = -2i$$

Which simplifies to

$$\lambda_1 = 2i$$

$$\lambda_2 = -2i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 2$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha t} (c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$y = e^0 (c_1 \cos(2t) + c_2 \sin(2t))$$

Or

$$y = c_1 \cos(2t) + c_2 \sin(2t)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 \cos(2t) + c_2 \sin(2t)$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \cos(2t)$$

$$y_2 = \sin(2t)$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(2t) & \sin(2t) \\ \frac{d}{dt}(\cos(2t)) & \frac{d}{dt}(\sin(2t)) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(2t) & \sin(2t) \\ -2 \sin(2t) & 2 \cos(2t) \end{vmatrix}$$

Therefore

$$W = (\cos(2t))(2 \cos(2t)) - (\sin(2t))(-2 \sin(2t))$$

Which simplifies to

$$W = 2 \cos (2t)^2 + 2 \sin (2t)^2$$

Which simplifies to

$$W = 2$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{3 \sin (2t) \csc (2t)}{2} dt$$

Which simplifies to

$$u_1 = - \int \frac{3}{2} dt$$

Hence

$$u_1 = -\frac{3t}{2}$$

And Eq. (3) becomes

$$u_2 = \int \frac{3 \cos (2t) \csc (2t)}{2} dt$$

Which simplifies to

$$u_2 = \int \frac{3 \cot (2t)}{2} dt$$

Hence

$$u_2 = -\frac{3 \ln (\cot (2t)^2 + 1)}{8}$$

Which simplifies to

$$u_1 = -\frac{3t}{2}$$
$$u_2 = -\frac{3 \ln (\csc (2t)^2)}{8}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -\frac{3 \cos(2t)t}{2} - \frac{3 \ln(\csc(2t)^2) \sin(2t)}{8}$$

Therefore the general solution is

$$y = y_h + y_p$$

$$= (c_1 \cos(2t) + c_2 \sin(2t)) + \left(-\frac{3 \cos(2t)t}{2} - \frac{3 \ln(\csc(2t)^2) \sin(2t)}{8} \right)$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(2t) + c_2 \sin(2t) - \frac{3 \cos(2t)t}{2} - \frac{3 \ln(\csc(2t)^2) \sin(2t)}{8} \quad (1)$$

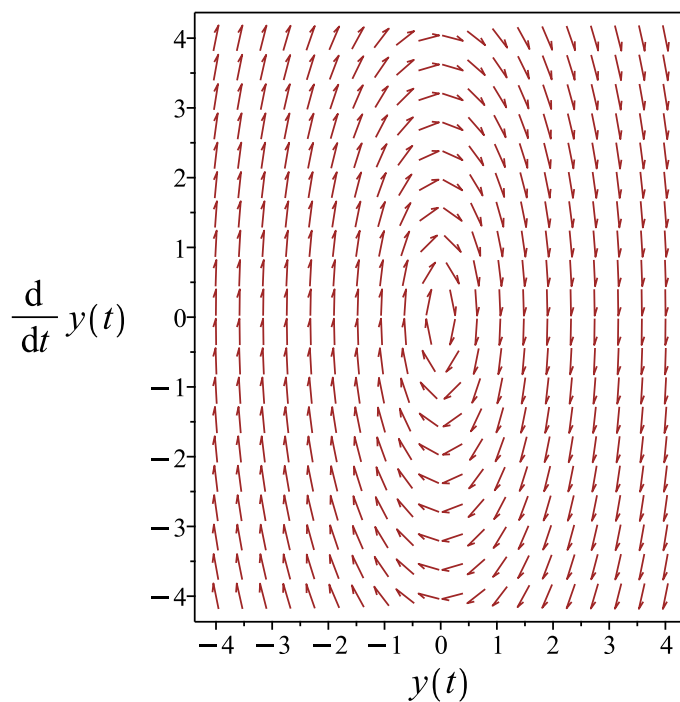


Figure 487: Slope field plot

Verification of solutions

$$y = c_1 \cos(2t) + c_2 \sin(2t) - \frac{3 \cos(2t)t}{2} - \frac{3 \ln(\csc(2t)^2) \sin(2t)}{8}$$

Verified OK.

10.8.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 4y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 0 \quad (3)$$

$$C = 4$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-4}{1} \quad (6)$$

Comparing the above to (5) shows that

$$s = -4$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(t) = -4z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 488: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -4$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos(2t)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}y_1 &= z_1 \\ &= \cos(2t)\end{aligned}$$

Which simplifies to

$$y_1 = \cos(2t)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Since $B = 0$ then the above becomes

$$\begin{aligned}y_2 &= y_1 \int \frac{1}{y_1^2} dt \\ &= \cos(2t) \int \frac{1}{\cos(2t)^2} dt \\ &= \cos(2t) \left(\frac{\tan(2t)}{2} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1(\cos(2t)) + c_2 \left(\cos(2t) \left(\frac{\tan(2t)}{2} \right) \right)\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$.
 y_h is the solution to

$$y'' + 4y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 \cos(2t) + \frac{c_2 \sin(2t)}{2}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \cos(2t)$$

$$y_2 = \frac{\sin(2t)}{2}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(2t) & \frac{\sin(2t)}{2} \\ \frac{d}{dt}(\cos(2t)) & \frac{d}{dt}\left(\frac{\sin(2t)}{2}\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(2t) & \frac{\sin(2t)}{2} \\ -2 \sin(2t) & \cos(2t) \end{vmatrix}$$

Therefore

$$W = (\cos(2t))(\cos(2t)) - \left(\frac{\sin(2t)}{2}\right)(-2\sin(2t))$$

Which simplifies to

$$W = \cos(2t)^2 + \sin(2t)^2$$

Which simplifies to

$$W = 1$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{3 \sin(2t) \csc(2t)}{1} dt$$

Which simplifies to

$$u_1 = - \int \frac{3}{2} dt$$

Hence

$$u_1 = -\frac{3t}{2}$$

And Eq. (3) becomes

$$u_2 = \int \frac{3 \cos(2t) \csc(2t)}{1} dt$$

Which simplifies to

$$u_2 = \int 3 \cot(2t) dt$$

Hence

$$u_2 = -\frac{3 \ln(\cot(2t)^2 + 1)}{4}$$

Which simplifies to

$$u_1 = -\frac{3t}{2}$$
$$u_2 = -\frac{3 \ln(\csc(2t)^2)}{4}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -\frac{3 \cos(2t) t}{2} - \frac{3 \ln(\csc(2t)^2) \sin(2t)}{8}$$

Therefore the general solution is

$$y = y_h + y_p$$

$$= \left(c_1 \cos(2t) + \frac{c_2 \sin(2t)}{2} \right) + \left(-\frac{3 \cos(2t) t}{2} - \frac{3 \ln(\csc(2t)^2) \sin(2t)}{8} \right)$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(2t) + \frac{c_2 \sin(2t)}{2} - \frac{3 \cos(2t) t}{2} - \frac{3 \ln(\csc(2t)^2) \sin(2t)}{8} \quad (1)$$

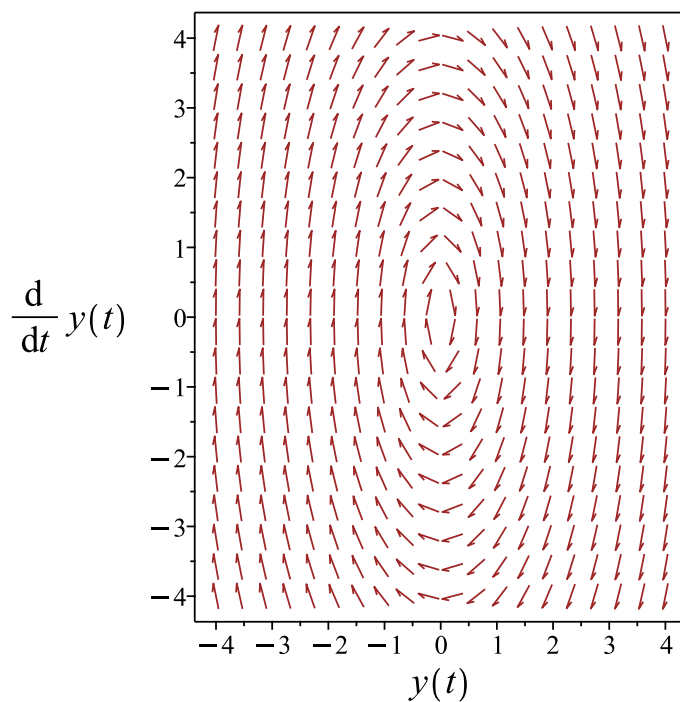


Figure 488: Slope field plot

Verification of solutions

$$y = c_1 \cos(2t) + \frac{c_2 \sin(2t)}{2} - \frac{3 \cos(2t) t}{2} - \frac{3 \ln(\csc(2t)^2) \sin(2t)}{8}$$

Verified OK.

10.8.3 Maple step by step solution

Let's solve

$$y'' + 4y = 3 \csc(2t)$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of homogeneous ODE

$$r^2 + 4 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{0 \pm (\sqrt{-16})}{2}$$

- Roots of the characteristic polynomial

$$r = (-2i, 2i)$$

- 1st solution of the homogeneous ODE

$$y_1(t) = \cos(2t)$$

- 2nd solution of the homogeneous ODE

$$y_2(t) = \sin(2t)$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 \cos(2t) + c_2 \sin(2t) + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t),y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t),y_2(t))} dt \right), f(t) = 3 \csc(2t) \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} \cos(2t) & \sin(2t) \\ -2 \sin(2t) & 2 \cos(2t) \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = 2$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = -\frac{3 \cos(2t) (\int 1 dt)}{2} + \frac{3 \sin(2t) (\int \cot(2t) dt)}{2}$$

- o Compute integrals

$$y_p(t) = -\frac{3 \cos(2t)t}{2} - \frac{3 \ln(\csc(2t)^2) \sin(2t)}{8}$$

- Substitute particular solution into general solution to ODE

$$y = c_1 \cos(2t) + c_2 \sin(2t) - \frac{3 \cos(2t)t}{2} - \frac{3 \ln(\csc(2t)^2) \sin(2t)}{8}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 35

```
dsolve(diff(y(t),t$2)+4*y(t) = 3*csc(2*t),y(t), singsol=all)
```

$$y(t) = -\frac{3 \ln(\csc(2t)) \sin(2t)}{4} + \frac{(-6t + 4c_1) \cos(2t)}{4} + c_2 \sin(2t)$$

✓ Solution by Mathematica

Time used: 0.036 (sec). Leaf size: 39

```
DSolve[y''[t]+4*y[t] ==3*Csc[2*t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \left(-\frac{3t}{2} + c_1\right) \cos(2t) + \frac{1}{4} \sin(2t)(3 \log(\sin(2t)) + 4c_2)$$

10.9 problem 9

10.9.1 Solving as second order linear constant coeff ode	2771
10.9.2 Solving using Kovacic algorithm	2776
10.9.3 Maple step by step solution	2782

Internal problem ID [691]

Internal file name [OUTPUT/691_Sunday_June_05_2022_01_47_02_AM_48256359/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters. page 190

Problem number: 9.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$y'' + y = 2 \sec\left(\frac{t}{2}\right)$$

10.9.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = 0, C = 1, f(t) = 2 \sec\left(\frac{t}{2}\right)$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = 0, C = 1$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(1)} \\ &= \pm i \end{aligned}$$

Hence

$$\lambda_1 = +i$$

$$\lambda_2 = -i$$

Which simplifies to

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha t} (c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$y = e^0 (c_1 \cos(t) + c_2 \sin(t))$$

Or

$$y = c_1 \cos(t) + c_2 \sin(t)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 \cos(t) + c_2 \sin(t)$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \cos(t)$$

$$y_2 = \sin(t)$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(t) & \sin(t) \\ \frac{d}{dt}(\cos(t)) & \frac{d}{dt}(\sin(t)) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(t) & \sin(t) \\ -\sin(t) & \cos(t) \end{vmatrix}$$

Therefore

$$W = (\cos(t))(\cos(t)) - (\sin(t))(-\sin(t))$$

Which simplifies to

$$W = \cos(t)^2 + \sin(t)^2$$

Which simplifies to

$$W = 1$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{2 \sin(t) \sec\left(\frac{t}{2}\right)}{1} dt$$

Which simplifies to

$$u_1 = - \int 4 \sin\left(\frac{t}{2}\right) dt$$

Hence

$$u_1 = 8 \cos\left(\frac{t}{2}\right)$$

And Eq. (3) becomes

$$u_2 = \int \frac{2 \cos(t) \sec\left(\frac{t}{2}\right)}{1} dt$$

Which simplifies to

$$u_2 = \int 2 \cos(t) \sec\left(\frac{t}{2}\right) dt$$

Hence

$$u_2 = -4 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) + 8 \sin\left(\frac{t}{2}\right)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = 8 \cos\left(\frac{t}{2}\right) \cos(t) + \left(-4 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) + 8 \sin\left(\frac{t}{2}\right)\right) \sin(t)$$

Which simplifies to

$$y_p(t) = \cos\left(\frac{t}{2}\right) \left(-8 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) \sin\left(\frac{t}{2}\right) + 8\right)$$

Therefore the general solution is

$$y = y_h + y_p \\ = (c_1 \cos(t) + c_2 \sin(t)) + \left(\cos\left(\frac{t}{2}\right) \left(-8 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) \sin\left(\frac{t}{2}\right) + 8\right)\right)$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(t) + c_2 \sin(t) + \cos\left(\frac{t}{2}\right) \left(-8 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) \sin\left(\frac{t}{2}\right) + 8\right) \quad (1)$$

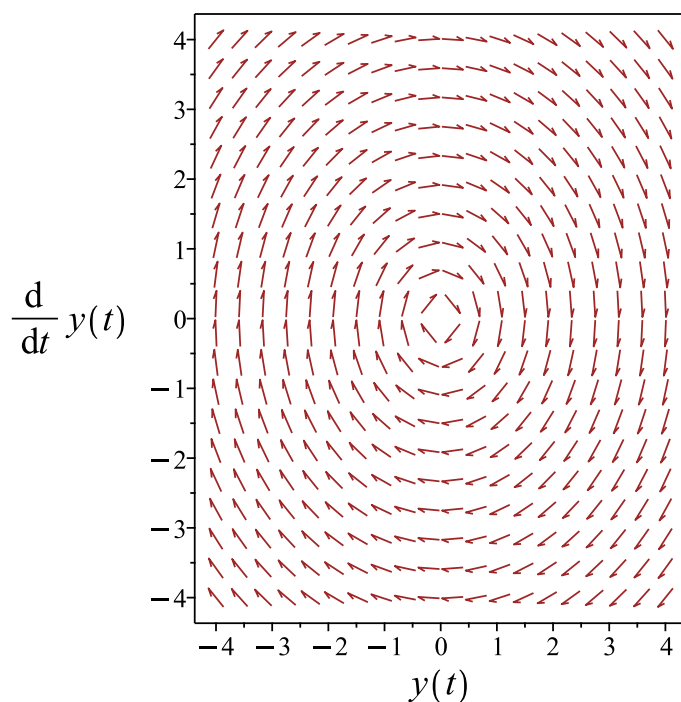


Figure 489: Slope field plot

Verification of solutions

$$y = c_1 \cos(t) + c_2 \sin(t) + \cos\left(\frac{t}{2}\right) \left(-8 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) \sin\left(\frac{t}{2}\right) + 8\right)$$

Verified OK.

10.9.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 0 \quad (3)$$

$$C = 1$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \quad (6)$$

Comparing the above to (5) shows that

$$s = -1$$

$$t = 1$$

Therefore eq. (4) becomes

$$z''(t) = -z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 490: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos(t)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}y_1 &= z_1 \\ &= \cos(t)\end{aligned}$$

Which simplifies to

$$y_1 = \cos(t)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Since $B = 0$ then the above becomes

$$\begin{aligned}y_2 &= y_1 \int \frac{1}{y_1^2} dt \\ &= \cos(t) \int \frac{1}{\cos(t)^2} dt \\ &= \cos(t) (\tan(t))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (\cos(t)) + c_2 (\cos(t) (\tan(t)))\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 \cos(t) + c_2 \sin(t)$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \cos(t)$$

$$y_2 = \sin(t)$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(t) & \sin(t) \\ \frac{d}{dt}(\cos(t)) & \frac{d}{dt}(\sin(t)) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(t) & \sin(t) \\ -\sin(t) & \cos(t) \end{vmatrix}$$

Therefore

$$W = (\cos(t))(\cos(t)) - (\sin(t))(-\sin(t))$$

Which simplifies to

$$W = \cos(t)^2 + \sin(t)^2$$

Which simplifies to

$$W = 1$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{2 \sin(t) \sec\left(\frac{t}{2}\right)}{1} dt$$

Which simplifies to

$$u_1 = - \int 4 \sin\left(\frac{t}{2}\right) dt$$

Hence

$$u_1 = 8 \cos\left(\frac{t}{2}\right)$$

And Eq. (3) becomes

$$u_2 = \int \frac{2 \cos(t) \sec\left(\frac{t}{2}\right)}{1} dt$$

Which simplifies to

$$u_2 = \int 2 \cos(t) \sec\left(\frac{t}{2}\right) dt$$

Hence

$$u_2 = -4 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) + 8 \sin\left(\frac{t}{2}\right)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = 8 \cos\left(\frac{t}{2}\right) \cos(t) + \left(-4 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) + 8 \sin\left(\frac{t}{2}\right)\right) \sin(t)$$

Which simplifies to

$$y_p(t) = \cos\left(\frac{t}{2}\right) \left(-8 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) \sin\left(\frac{t}{2}\right) + 8\right)$$

Therefore the general solution is

$$y = y_h + y_p$$
$$= (c_1 \cos(t) + c_2 \sin(t)) + \left(\cos\left(\frac{t}{2}\right) \left(-8 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) \sin\left(\frac{t}{2}\right) + 8 \right) \right)$$

Summary

The solution(s) found are the following

$$y = c_1 \cos(t) + c_2 \sin(t) + \cos\left(\frac{t}{2}\right) \left(-8 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) \sin\left(\frac{t}{2}\right) + 8 \right) (1)$$

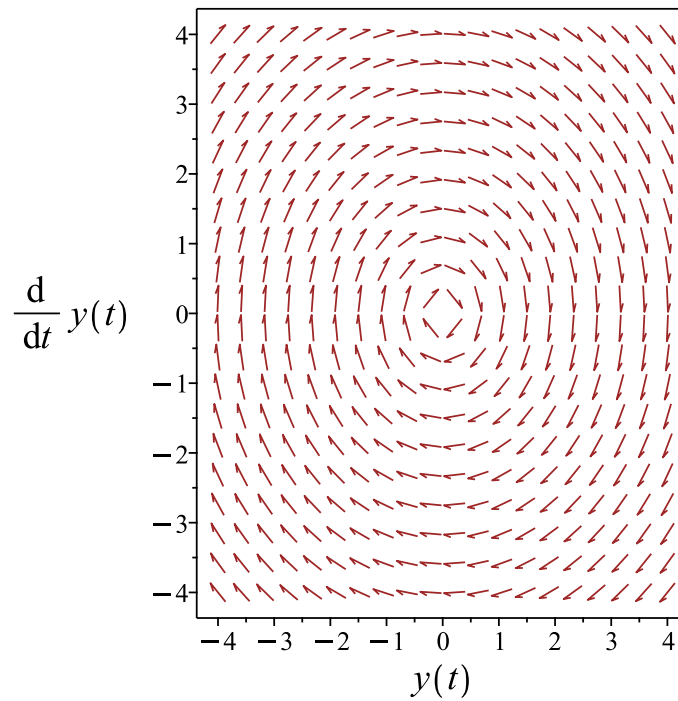


Figure 490: Slope field plot

Verification of solutions

$$y = c_1 \cos(t) + c_2 \sin(t) + \cos\left(\frac{t}{2}\right) \left(-8 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) \sin\left(\frac{t}{2}\right) + 8 \right)$$

Verified OK.

10.9.3 Maple step by step solution

Let's solve

$$y'' + y = 2 \sec\left(\frac{t}{2}\right)$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of homogeneous ODE

$$r^2 + 1 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{0 \pm (\sqrt{-4})}{2}$$

- Roots of the characteristic polynomial

$$r = (-I, I)$$

- 1st solution of the homogeneous ODE

$$y_1(t) = \cos(t)$$

- 2nd solution of the homogeneous ODE

$$y_2(t) = \sin(t)$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 \cos(t) + c_2 \sin(t) + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t), y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t), y_2(t))} dt \right), f(t) = 2 \sec\left(\frac{t}{2}\right) \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} \cos(t) & \sin(t) \\ -\sin(t) & \cos(t) \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = 1$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = -4 \cos(t) \left(\int \sin\left(\frac{t}{2}\right) dt \right) + 2 \sin(t) \left(\int \cos(t) \sec\left(\frac{t}{2}\right) dt \right)$$

- Compute integrals

$$y_p(t) = \cos\left(\frac{t}{2}\right) \left(-8 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) \sin\left(\frac{t}{2}\right) + 8 \right)$$

- Substitute particular solution into general solution to ODE

$$y = c_1 \cos(t) + c_2 \sin(t) + \cos\left(\frac{t}{2}\right) \left(-8 \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) \sin\left(\frac{t}{2}\right) + 8 \right)$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 33

```
dsolve(diff(y(t),t$2)+y(t) = 2*sec(t/2),y(t), singsol=all)
```

$$y(t) = -4 \sin(t) \ln\left(\sec\left(\frac{t}{2}\right) + \tan\left(\frac{t}{2}\right)\right) + c_2 \sin(t) + \cos(t) c_1 + 8 \cos\left(\frac{t}{2}\right)$$

✓ Solution by Mathematica

Time used: 0.114 (sec). Leaf size: 35

```
DSolve[y''[t]+y[t]== 2*Sec[t/2],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -4 \sin(t) \operatorname{arctanh}\left(\sin\left(\frac{t}{2}\right)\right) + 8 \cos\left(\frac{t}{2}\right) + c_1 \cos(t) + c_2 \sin(t)$$

10.10 problem 10

10.10.1 Solving as second order linear constant coeff ode	2784
10.10.2 Solving as linear second order ode solved by an integrating factor ode	2788
10.10.3 Solving using Kovacic algorithm	2789
10.10.4 Maple step by step solution	2794

Internal problem ID [692]

Internal file name [OUTPUT/692_Sunday_June_05_2022_01_47_03_AM_2706801/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 10.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "linear_second_order_ode_solved_by_an_integrating_factor"

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$y'' - 2y' + y = \frac{e^t}{t^2 + 1}$$

10.10.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = -2, C = 1, f(t) = \frac{e^t}{t^2+1}$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$.

y_h is the solution to

$$y'' - 2y' + y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = -2, C = 1$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} - 2\lambda e^{\lambda t} + e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 - 2\lambda + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -2, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{2}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{(-2)^2 - (4)(1)(1)} \\ &= 1 \end{aligned}$$

Hence this is the case of a double root $\lambda_{1,2} = -1$. Therefore the solution is

$$y = c_1 e^t + c_2 t e^t \quad (1)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 e^t + c_2 t e^t$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned} y_1 &= e^t \\ y_2 &= t e^t \end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} e^t & t e^t \\ \frac{d}{dt}(e^t) & \frac{d}{dt}(t e^t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} e^t & t e^t \\ e^t & t e^t + e^t \end{vmatrix}$$

Therefore

$$W = (e^t) (t e^t + e^t) - (t e^t) (e^t)$$

Which simplifies to

$$W = e^{2t}$$

Which simplifies to

$$W = e^{2t}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{t e^{2t}}{e^{2t}} dt$$

Which simplifies to

$$u_1 = - \int \frac{t}{t^2 + 1} dt$$

Hence

$$u_1 = - \frac{\ln(t^2 + 1)}{2}$$

And Eq. (3) becomes

$$u_2 = \int \frac{e^{2t}}{t^2+1} dt$$

Which simplifies to

$$u_2 = \int \frac{1}{t^2+1} dt$$

Hence

$$u_2 = \arctan(t)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -\frac{\ln(t^2+1)e^t}{2} + \arctan(t)te^t$$

Which simplifies to

$$y_p(t) = e^t \left(-\frac{\ln(t^2+1)}{2} + t \arctan(t) \right)$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1e^t + c_2te^t) + \left(e^t \left(-\frac{\ln(t^2+1)}{2} + t \arctan(t) \right) \right) \end{aligned}$$

Which simplifies to

$$y = e^t(c_2t + c_1) + e^t \left(-\frac{\ln(t^2+1)}{2} + t \arctan(t) \right)$$

Summary

The solution(s) found are the following

$$y = e^t(c_2t + c_1) + e^t \left(-\frac{\ln(t^2+1)}{2} + t \arctan(t) \right) \quad (1)$$

Verification of solutions

$$y = e^t(c_2t + c_1) + e^t \left(-\frac{\ln(t^2+1)}{2} + t \arctan(t) \right)$$

Verified OK.

10.10.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t)y' + \frac{(p(t)^2 + p'(t))y}{2} = f(t)$$

Where $p(t) = -2$. Therefore, there is an integrating factor given by

$$\begin{aligned}M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int -2 dx} \\ &= e^{-t}\end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned}(M(x)y)'' &= \frac{e^{-t}e^t}{t^2 + 1} \\ (e^{-t}y)'' &= \frac{e^{-t}e^t}{t^2 + 1}\end{aligned}$$

Integrating once gives

$$(e^{-t}y)' = \arctan(t) + c_1$$

Integrating again gives

$$(e^{-t}y) = c_1t + t \arctan(t) - \frac{\ln(t^2 + 1)}{2} + c_2$$

Hence the solution is

$$y = \frac{c_1t + t \arctan(t) - \frac{\ln(t^2+1)}{2} + c_2}{e^{-t}}$$

Or

$$y = c_1t e^t + t \arctan(t) e^t + c_2 e^t - \frac{e^t \ln(t^2 + 1)}{2}$$

Summary

The solution(s) found are the following

$$y = c_1t e^t + t \arctan(t) e^t + c_2 e^t - \frac{e^t \ln(t^2 + 1)}{2} \quad (1)$$

Verification of solutions

$$y = c_1 t e^t + t \arctan(t) e^t + c_2 e^t - \frac{e^t \ln(t^2 + 1)}{2}$$

Verified OK.

10.10.3 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 2y' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= -2 \\ C &= 1 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 0 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = 0 \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 492: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-2}{1} dt} \\ &= z_1 e^t \\ &= z_1 (e^t) \end{aligned}$$

Which simplifies to

$$y_1 = e^t$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-2}{1} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{2t}}{(y_1)^2} dt \\ &= y_1(t) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^t) + c_2 (e^t(t)) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' - 2y' + y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^t + c_2 t e^t$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = e^t$$

$$y_2 = t e^t$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \tag{3}$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} e^t & t e^t \\ \frac{d}{dt}(e^t) & \frac{d}{dt}(t e^t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} e^t & t e^t \\ e^t & t e^t + e^t \end{vmatrix}$$

Therefore

$$W = (e^t) (t e^t + e^t) - (t e^t) (e^t)$$

Which simplifies to

$$W = e^{2t}$$

Which simplifies to

$$W = e^{2t}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{t e^{2t}}{t^2 + 1} dt$$

Which simplifies to

$$u_1 = - \int \frac{t}{t^2 + 1} dt$$

Hence

$$u_1 = - \frac{\ln(t^2 + 1)}{2}$$

And Eq. (3) becomes

$$u_2 = \int \frac{e^{2t}}{t^2 + 1} dt$$

Which simplifies to

$$u_2 = \int \frac{1}{t^2 + 1} dt$$

Hence

$$u_2 = \arctan(t)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -\frac{e^t \ln(t^2 + 1)}{2} + t \arctan(t) e^t$$

Which simplifies to

$$y_p(t) = e^t \left(-\frac{\ln(t^2 + 1)}{2} + t \arctan(t) \right)$$

Therefore the general solution is

$$\begin{aligned}y &= y_h + y_p \\ &= (c_1 e^t + c_2 t e^t) + \left(e^t \left(-\frac{\ln(t^2 + 1)}{2} + t \arctan(t) \right) \right)\end{aligned}$$

Which simplifies to

$$y = e^t(c_2 t + c_1) + e^t \left(-\frac{\ln(t^2 + 1)}{2} + t \arctan(t) \right)$$

Summary

The solution(s) found are the following

$$y = e^t(c_2 t + c_1) + e^t \left(-\frac{\ln(t^2 + 1)}{2} + t \arctan(t) \right) \quad (1)$$

Verification of solutions

$$y = e^t(c_2 t + c_1) + e^t \left(-\frac{\ln(t^2 + 1)}{2} + t \arctan(t) \right)$$

Verified OK.

10.10.4 Maple step by step solution

Let's solve

$$y'' - 2y' + y = \frac{e^t}{t^2 + 1}$$

- Highest derivative means the order of the ODE is 2
 y''
- Characteristic polynomial of homogeneous ODE
 $r^2 - 2r + 1 = 0$
- Factor the characteristic polynomial
 $(r - 1)^2 = 0$
- Root of the characteristic polynomial
 $r = 1$
- 1st solution of the homogeneous ODE

$$y_1(t) = e^t$$

- Repeated root, multiply $y_1(t)$ by t to ensure linear independence

$$y_2(t) = t e^t$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 e^t + c_2 t e^t + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t),y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t),y_2(t))} dt \right), f(t) = \frac{e^t}{t^2+1} \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} e^t & t e^t \\ e^t & t e^t + e^t \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = e^{2t}$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = e^t \left(- \left(\int \frac{t}{t^2+1} dt \right) + \left(\int \frac{1}{t^2+1} dt \right) t \right)$$

- Compute integrals

$$y_p(t) = - \frac{e^t (-2t \arctan(t) + \ln(t^2+1))}{2}$$

- Substitute particular solution into general solution to ODE

$$y = c_1 e^t + c_2 t e^t - \frac{e^t (-2t \arctan(t) + \ln(t^2+1))}{2}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]  
trying a double symmetry of the form [xi=0, eta=F(x)]  
-> Try solving first the homogeneous part of the ODE  
    checking if the LODE has constant coefficients  
    <- constant coefficients successful  
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 24

```
dsolve(diff(y(t),t$2)-2*diff(y(t),t)+y(t) = exp(t)/(1+t^2),y(t), singsol=all)
```

$$y(t) = e^t \left(c_2 + c_1 t - \frac{\ln(t^2 + 1)}{2} + t \arctan(t) \right)$$

✓ Solution by Mathematica

Time used: 0.029 (sec). Leaf size: 35

```
DSolve[y''[t]-2*y'[t]+y[t] == Exp[t]/(1+t^2),y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{2} e^t (2t \arctan(t) - \log(t^2 + 1) + 2(c_2 t + c_1))$$

10.11 problem 11

10.11.1 Solving as second order linear constant coeff ode	2797
10.11.2 Solving using Kovacic algorithm	2801
10.11.3 Maple step by step solution	2806

Internal problem ID [693]

Internal file name [OUTPUT/693_Sunday_June_05_2022_01_47_04_AM_3995225/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters. page 190

Problem number: 11.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$y'' - 5y' + 6y = g(t)$$

10.11.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = -5, C = 6, f(t) = g(t)$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' - 5y' + 6y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = -5, C = 6$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} - 5\lambda e^{\lambda t} + 6e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 - 5\lambda + 6 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -5, C = 6$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{5}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-5^2 - (4)(1)(6)} \\ &= \frac{5}{2} \pm \frac{1}{2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= \frac{5}{2} + \frac{1}{2} \\ \lambda_2 &= \frac{5}{2} - \frac{1}{2} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= 3 \\ \lambda_2 &= 2 \end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned} y &= c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t} \\ y &= c_1 e^{(3)t} + c_2 e^{(2)t} \end{aligned}$$

Or

$$y = c_1 e^{3t} + c_2 e^{2t}$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 e^{3t} + c_2 e^{2t}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = e^{3t}$$

$$y_2 = e^{2t}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} e^{3t} & e^{2t} \\ \frac{d}{dt}(e^{3t}) & \frac{d}{dt}(e^{2t}) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} e^{3t} & e^{2t} \\ 3e^{3t} & 2e^{2t} \end{vmatrix}$$

Therefore

$$W = (e^{3t})(2e^{2t}) - (e^{2t})(3e^{3t})$$

Which simplifies to

$$W = -e^{3t}e^{2t}$$

Which simplifies to

$$W = -e^{5t}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{e^{2t} g(t)}{-e^{5t}} dt$$

Which simplifies to

$$u_1 = - \int -g(t) e^{-3t} dt$$

Hence

$$u_1 = - \left(\int_0^t -g(\alpha) e^{-3\alpha} d\alpha \right)$$

And Eq. (3) becomes

$$u_2 = \int \frac{e^{3t} g(t)}{-e^{5t}} dt$$

Which simplifies to

$$u_2 = \int -g(t) e^{-2t} dt$$

Hence

$$u_2 = \int_0^t -g(\alpha) e^{-2\alpha} d\alpha$$

Which simplifies to

$$u_1 = \int_0^t g(\alpha) e^{-3\alpha} d\alpha$$
$$u_2 = - \left(\int_0^t g(\alpha) e^{-2\alpha} d\alpha \right)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \left(\int_0^t g(\alpha) e^{-3\alpha} d\alpha \right) e^{3t} - \left(\int_0^t g(\alpha) e^{-2\alpha} d\alpha \right) e^{2t}$$

Therefore the general solution is

$$\begin{aligned}y &= y_h + y_p \\ &= (c_1 e^{3t} + c_2 e^{2t}) + \left(\left(\int_0^t g(\alpha) e^{-3\alpha} d\alpha \right) e^{3t} - \left(\int_0^t g(\alpha) e^{-2\alpha} d\alpha \right) e^{2t} \right)\end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{3t} + c_2 e^{2t} + \left(\int_0^t g(\alpha) e^{-3\alpha} d\alpha \right) e^{3t} - \left(\int_0^t g(\alpha) e^{-2\alpha} d\alpha \right) e^{2t} \quad (1)$$

Verification of solutions

$$y = c_1 e^{3t} + c_2 e^{2t} + \left(\int_0^t g(\alpha) e^{-3\alpha} d\alpha \right) e^{3t} - \left(\int_0^t g(\alpha) e^{-2\alpha} d\alpha \right) e^{2t}$$

Verified OK.

10.11.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' - 5y' + 6y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= 1 \\ B &= -5 \\ C &= 6\end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned}r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}\end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{4} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \frac{z(t)}{4} \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 494: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{4}$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = e^{-\frac{t}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-5}{1} dt} \\ &= z_1 e^{\frac{5t}{2}} \\ &= z_1 \left(e^{\frac{5t}{2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{2t}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-5}{1} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{5t}}{(y_1)^2} dt \\ &= y_1 (e^t) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{2t}) + c_2 (e^{2t} (e^t)) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' - 5y' + 6y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^{2t} + c_2 e^{3t}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = e^{2t}$$

$$y_2 = e^{3t}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \tag{3}$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} e^{2t} & e^{3t} \\ \frac{d}{dt}(e^{2t}) & \frac{d}{dt}(e^{3t}) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} e^{2t} & e^{3t} \\ 2e^{2t} & 3e^{3t} \end{vmatrix}$$

Therefore

$$W = (e^{2t})(3e^{3t}) - (e^{3t})(2e^{2t})$$

Which simplifies to

$$W = e^{3t}e^{2t}$$

Which simplifies to

$$W = e^{5t}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{e^{3t}g(t)}{e^{5t}} dt$$

Which simplifies to

$$u_1 = - \int g(t) e^{-2t} dt$$

Hence

$$u_1 = - \left(\int_0^t g(\alpha) e^{-2\alpha} d\alpha \right)$$

And Eq. (3) becomes

$$u_2 = \int \frac{e^{2t}g(t)}{e^{5t}} dt$$

Which simplifies to

$$u_2 = \int g(t) e^{-3t} dt$$

Hence

$$u_2 = \int_0^t g(\alpha) e^{-3\alpha} d\alpha$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -\left(\int_0^t g(\alpha) e^{-2\alpha} d\alpha\right) e^{2t} + \left(\int_0^t g(\alpha) e^{-3\alpha} d\alpha\right) e^{3t}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 e^{2t} + c_2 e^{3t}) + \left(-\left(\int_0^t g(\alpha) e^{-2\alpha} d\alpha\right) e^{2t} + \left(\int_0^t g(\alpha) e^{-3\alpha} d\alpha\right) e^{3t}\right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1 e^{2t} + c_2 e^{3t} - \left(\int_0^t g(\alpha) e^{-2\alpha} d\alpha\right) e^{2t} + \left(\int_0^t g(\alpha) e^{-3\alpha} d\alpha\right) e^{3t} \quad (1)$$

Verification of solutions

$$y = c_1 e^{2t} + c_2 e^{3t} - \left(\int_0^t g(\alpha) e^{-2\alpha} d\alpha\right) e^{2t} + \left(\int_0^t g(\alpha) e^{-3\alpha} d\alpha\right) e^{3t}$$

Verified OK.

10.11.3 Maple step by step solution

Let's solve

$$y'' - 5y' + 6y = g(t)$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of homogeneous ODE

$$r^2 - 5r + 6 = 0$$

- Factor the characteristic polynomial

$$(r - 2)(r - 3) = 0$$

- Roots of the characteristic polynomial

$$r = (2, 3)$$

- 1st solution of the homogeneous ODE

$$y_1(t) = e^{2t}$$

- 2nd solution of the homogeneous ODE

$$y_2(t) = e^{3t}$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 e^{2t} + c_2 e^{3t} + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t), y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t), y_2(t))} dt \right), f(t) = g(t) \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} e^{2t} & e^{3t} \\ 2e^{2t} & 3e^{3t} \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = e^{5t}$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = -e^{2t} \left(\int g(t) e^{-2t} dt \right) + e^{3t} \left(\int g(t) e^{-3t} dt \right)$$

- Compute integrals

$$y_p(t) = -e^{2t} \left(\int g(t) e^{-2t} dt \right) + e^{3t} \left(\int g(t) e^{-3t} dt \right)$$

- Substitute particular solution into general solution to ODE

$$y = c_1 e^{2t} + c_2 e^{3t} - e^{2t} \left(\int g(t) e^{-2t} dt \right) + e^{3t} \left(\int g(t) e^{-3t} dt \right)$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]  
trying a double symmetry of the form [xi=0, eta=F(x)]  
-> Try solving first the homogeneous part of the ODE  
    checking if the LODE has constant coefficients  
    <- constant coefficients successful  
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 46

```
dsolve(diff(y(t),t$2)-5*diff(y(t),t)+6*y(t) = g(t),y(t), singsol=all)
```

$$y(t) = c_2 e^{2t} + c_1 e^{3t} - \left(\int g(t) e^{-2t} dt \right) e^{2t} + \left(\int g(t) e^{-3t} dt \right) e^{3t}$$

✓ Solution by Mathematica

Time used: 0.064 (sec). Leaf size: 59

```
DSolve[y''[t]-5*y'[t]+6*y[t] == g[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{2t} \left(\int_1^t -e^{-2K[1]} g(K[1]) dK[1] + e^t \int_1^t e^{-3K[2]} g(K[2]) dK[2] + c_2 e^t + c_1 \right)$$

10.12 problem 12

10.12.1 Solving as second order linear constant coeff ode	2809
10.12.2 Solving using Kovacic algorithm	2813
10.12.3 Maple step by step solution	2819

Internal problem ID [694]

Internal file name [OUTPUT/694_Sunday_June_05_2022_01_47_06_AM_78756365/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 12.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff"

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$y'' + 4y = g(t)$$

10.12.1 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = 0, C = 4, f(t) = g(t)$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$.
 y_h is the solution to

$$y'' + 4y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = 0, C = 4$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + 4 e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + 4 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 4$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(4)} \\ &= \pm 2i \end{aligned}$$

Hence

$$\lambda_1 = +2i$$

$$\lambda_2 = -2i$$

Which simplifies to

$$\lambda_1 = 2i$$

$$\lambda_2 = -2i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 2$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha t} (c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$y = e^0 (c_1 \cos(2t) + c_2 \sin(2t))$$

Or

$$y = c_1 \cos(2t) + c_2 \sin(2t)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 \cos(2t) + c_2 \sin(2t)$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \cos(2t)$$

$$y_2 = \sin(2t)$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(2t) & \sin(2t) \\ \frac{d}{dt}(\cos(2t)) & \frac{d}{dt}(\sin(2t)) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(2t) & \sin(2t) \\ -2 \sin(2t) & 2 \cos(2t) \end{vmatrix}$$

Therefore

$$W = (\cos(2t))(2 \cos(2t)) - (\sin(2t))(-2 \sin(2t))$$

Which simplifies to

$$W = 2 \cos (2t)^2 + 2 \sin (2t)^2$$

Which simplifies to

$$W = 2$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\sin (2t) g(t)}{2} dt$$

Which simplifies to

$$u_1 = - \int \frac{\sin (2t) g(t)}{2} dt$$

Hence

$$u_1 = - \left(\int_0^t \frac{\sin (2\alpha) g(\alpha)}{2} d\alpha \right)$$

And Eq. (3) becomes

$$u_2 = \int \frac{\cos (2t) g(t)}{2} dt$$

Which simplifies to

$$u_2 = \int \frac{\cos (2t) g(t)}{2} dt$$

Hence

$$u_2 = \int_0^t \frac{\cos (2\alpha) g(\alpha)}{2} d\alpha$$

Which simplifies to

$$u_1 = - \frac{\left(\int_0^t \sin (2\alpha) g(\alpha) d\alpha \right)}{2}$$
$$u_2 = \frac{\left(\int_0^t \cos (2\alpha) g(\alpha) d\alpha \right)}{2}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -\frac{\left(\int_0^t \sin(2\alpha) g(\alpha) d\alpha\right) \cos(2t)}{2} + \frac{\left(\int_0^t \cos(2\alpha) g(\alpha) d\alpha\right) \sin(2t)}{2}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 \cos(2t) + c_2 \sin(2t)) \\ &\quad + \left(-\frac{\left(\int_0^t \sin(2\alpha) g(\alpha) d\alpha\right) \cos(2t)}{2} + \frac{\left(\int_0^t \cos(2\alpha) g(\alpha) d\alpha\right) \sin(2t)}{2} \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$\begin{aligned} y &= c_1 \cos(2t) + c_2 \sin(2t) - \frac{\left(\int_0^t \sin(2\alpha) g(\alpha) d\alpha\right) \cos(2t)}{2} \\ &\quad + \frac{\left(\int_0^t \cos(2\alpha) g(\alpha) d\alpha\right) \sin(2t)}{2} \end{aligned} \tag{1}$$

Verification of solutions

$$\begin{aligned} y &= c_1 \cos(2t) + c_2 \sin(2t) - \frac{\left(\int_0^t \sin(2\alpha) g(\alpha) d\alpha\right) \cos(2t)}{2} \\ &\quad + \frac{\left(\int_0^t \cos(2\alpha) g(\alpha) d\alpha\right) \sin(2t)}{2} \end{aligned}$$

Verified OK.

10.12.2 Solving using Kovacic algorithm

Writing the ode as

$$y'' + 4y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 0 \\ C &= 4 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-4}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -4 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = -4z(t) \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 496: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -4$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos(2t)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}
 y_1 &= z_1 \\
 &= \cos(2t)
 \end{aligned}$$

Which simplifies to

$$y_1 = \cos(2t)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Since $B = 0$ then the above becomes

$$\begin{aligned} y_2 &= y_1 \int \frac{1}{y_1^2} dt \\ &= \cos(2t) \int \frac{1}{\cos(2t)^2} dt \\ &= \cos(2t) \left(\frac{\tan(2t)}{2} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1(\cos(2t)) + c_2 \left(\cos(2t) \left(\frac{\tan(2t)}{2} \right) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + 4y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 \cos(2t) + \frac{c_2 \sin(2t)}{2}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of

parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \cos(2t)$$

$$y_2 = \frac{\sin(2t)}{2}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(2t) & \frac{\sin(2t)}{2} \\ \frac{d}{dt}(\cos(2t)) & \frac{d}{dt}\left(\frac{\sin(2t)}{2}\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(2t) & \frac{\sin(2t)}{2} \\ -2 \sin(2t) & \cos(2t) \end{vmatrix}$$

Therefore

$$W = (\cos(2t))(\cos(2t)) - \left(\frac{\sin(2t)}{2}\right)(-2 \sin(2t))$$

Which simplifies to

$$W = \cos(2t)^2 + \sin(2t)^2$$

Which simplifies to

$$W = 1$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{\sin(2t)g(t)}{2}}{1} dt$$

Which simplifies to

$$u_1 = - \int \frac{\sin(2t)g(t)}{2} dt$$

Hence

$$u_1 = - \left(\int_0^t \frac{\sin(2\alpha)g(\alpha)}{2} d\alpha \right)$$

And Eq. (3) becomes

$$u_2 = \int \frac{\cos(2t)g(t)}{1} dt$$

Which simplifies to

$$u_2 = \int \cos(2t)g(t) dt$$

Hence

$$u_2 = \int_0^t \cos(2\alpha)g(\alpha) d\alpha$$

Which simplifies to

$$u_1 = - \frac{\left(\int_0^t \sin(2\alpha)g(\alpha) d\alpha \right)}{2}$$

$$u_2 = \int_0^t \cos(2\alpha)g(\alpha) d\alpha$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = - \frac{\left(\int_0^t \sin(2\alpha)g(\alpha) d\alpha \right) \cos(2t)}{2} + \frac{\left(\int_0^t \cos(2\alpha)g(\alpha) d\alpha \right) \sin(2t)}{2}$$

Therefore the general solution is

$$\begin{aligned}
 y &= y_h + y_p \\
 &= \left(c_1 \cos(2t) + \frac{c_2 \sin(2t)}{2} \right) \\
 &\quad + \left(-\frac{\left(\int_0^t \sin(2\alpha) g(\alpha) d\alpha \right) \cos(2t)}{2} + \frac{\left(\int_0^t \cos(2\alpha) g(\alpha) d\alpha \right) \sin(2t)}{2} \right)
 \end{aligned}$$

Summary

The solution(s) found are the following

$$\begin{aligned}
 y &= c_1 \cos(2t) + \frac{c_2 \sin(2t)}{2} - \frac{\left(\int_0^t \sin(2\alpha) g(\alpha) d\alpha \right) \cos(2t)}{2} \\
 &\quad + \frac{\left(\int_0^t \cos(2\alpha) g(\alpha) d\alpha \right) \sin(2t)}{2}
 \end{aligned} \tag{1}$$

Verification of solutions

$$\begin{aligned}
 y &= c_1 \cos(2t) + \frac{c_2 \sin(2t)}{2} - \frac{\left(\int_0^t \sin(2\alpha) g(\alpha) d\alpha \right) \cos(2t)}{2} \\
 &\quad + \frac{\left(\int_0^t \cos(2\alpha) g(\alpha) d\alpha \right) \sin(2t)}{2}
 \end{aligned}$$

Verified OK.

10.12.3 Maple step by step solution

Let's solve

$$y'' + 4y = g(t)$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Characteristic polynomial of homogeneous ODE

$$r^2 + 4 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{0 \pm (\sqrt{-16})}{2}$$

- Roots of the characteristic polynomial

$$r = (-2I, 2I)$$

- 1st solution of the homogeneous ODE

$$y_1(t) = \cos(2t)$$

- 2nd solution of the homogeneous ODE

$$y_2(t) = \sin(2t)$$

- General solution of the ODE

$$y = c_1 y_1(t) + c_2 y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = c_1 \cos(2t) + c_2 \sin(2t) + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t), y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t), y_2(t))} dt \right), f(t) = g(t) \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} \cos(2t) & \sin(2t) \\ -2\sin(2t) & 2\cos(2t) \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = 2$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = -\frac{\cos(2t) \left(\int \sin(2t)g(t)dt \right)}{2} + \frac{\sin(2t) \left(\int \cos(2t)g(t)dt \right)}{2}$$

- Compute integrals

$$y_p(t) = -\frac{\cos(2t) \left(\int \sin(2t)g(t)dt \right)}{2} + \frac{\sin(2t) \left(\int \cos(2t)g(t)dt \right)}{2}$$

- Substitute particular solution into general solution to ODE

$$y = c_1 \cos(2t) + c_2 \sin(2t) - \frac{\cos(2t) \left(\int \sin(2t)g(t)dt \right)}{2} + \frac{\sin(2t) \left(\int \cos(2t)g(t)dt \right)}{2}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]  
trying a double symmetry of the form [xi=0, eta=F(x)]  
-> Try solving first the homogeneous part of the ODE  
    checking if the LODE has constant coefficients  
    <- constant coefficients successful  
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 47

```
dsolve(diff(y(t),t$2)+4*y(t) = g(t),y(t), singsol=all)
```

$$y(t) = c_2 \sin(2t) + c_1 \cos(2t) + \frac{(\int \cos(2t) g(t) dt) \sin(2t)}{2} - \frac{(\int \sin(2t) g(t) dt) \cos(2t)}{2}$$

✓ Solution by Mathematica

Time used: 0.093 (sec). Leaf size: 67

```
DSolve[y''[t]+4*y[t] == g[t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \cos(2t) \int_1^t -\cos(K[1])g(K[1]) \sin(K[1])dK[1] \\ + \sin(2t) \int_1^t \frac{1}{2} \cos(2K[2])g(K[2])dK[2] + c_1 \cos(2t) + c_2 \sin(2t)$$

10.13 problem 13

10.13.1 Solving as second order euler ode ode	2822
10.13.2 Solving as second order integrable as is ode	2826
10.13.3 Solving as type second_order_integrable_as_is (not using ABC version)	2827
10.13.4 Solving using Kovacic algorithm	2829
10.13.5 Solving as exact linear second order ode ode	2836

Internal problem ID [695]

Internal file name [OUTPUT/695_Sunday_June_05_2022_01_47_07_AM_67149286/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 13.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "exact linear second order ode", "second_order_integrable_as_is"

Maple gives the following as the ode type

```
[[_2nd_order, _exact, _linear, _nonhomogeneous]]
```

$$t^2 y'' - 2y = 3t^2 - 1$$

10.13.1 Solving as second order euler ode ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = t^2$, $B = 0$, $C = -2$, $f(t) = 3t^2 - 1$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. Solving for y_h from

$$t^2 y'' - 2y = 0$$

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} + 0rt^{r-1} - 2t^r = 0$$

Simplifying gives

$$r(r-1)t^r + 0t^r - 2t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r-1) + 0 - 2 = 0$$

Or

$$r^2 - r - 2 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -1$$

$$r_2 = 2$$

Since the roots are real and distinct, then the general solution is

$$y = c_1y_1 + c_2y_2$$

Where $y_1 = t^{-1}$ and $y_2 = t^2$. Hence

$$y = \frac{c_1}{t} + c_2t^2$$

Next, we find the particular solution to the ODE

$$t^2y'' - 2y = 3t^2 - 1$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1y_1 + u_2y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \frac{1}{t}$$

$$y_2 = t^2$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{1}{t} & t^2 \\ \frac{d}{dt}\left(\frac{1}{t}\right) & \frac{d}{dt}(t^2) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{1}{t} & t^2 \\ -\frac{1}{t^2} & 2t \end{vmatrix}$$

Therefore

$$W = \left(\frac{1}{t}\right)(2t) - (t^2)\left(-\frac{1}{t^2}\right)$$

Which simplifies to

$$W = 3$$

Which simplifies to

$$W = 3$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{t^2(3t^2 - 1)}{3t^2} dt$$

Which simplifies to

$$u_1 = - \int \left(t^2 - \frac{1}{3}\right) dt$$

Hence

$$u_1 = -\frac{1}{3}t^3 + \frac{1}{3}t$$

And Eq. (3) becomes

$$u_2 = \int \frac{3t^2-1}{3t^2} dt$$

Which simplifies to

$$u_2 = \int \frac{3t^2 - 1}{3t^3} dt$$

Hence

$$u_2 = \ln(t) + \frac{1}{6t^2}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{-\frac{1}{3}t^3 + \frac{1}{3}t}{t} + \left(\ln(t) + \frac{1}{6t^2} \right) t^2$$

Which simplifies to

$$y_p(t) = t^2 \ln(t) - \frac{t^2}{3} + \frac{1}{2}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= t^2 \ln(t) - \frac{t^2}{3} + \frac{1}{2} + \frac{c_1}{t} + c_2 t^2 \end{aligned}$$

Summary

The solution(s) found are the following

$$y = t^2 \ln(t) - \frac{t^2}{3} + \frac{1}{2} + \frac{c_1}{t} + c_2 t^2 \quad (1)$$

Verification of solutions

$$y = t^2 \ln(t) - \frac{t^2}{3} + \frac{1}{2} + \frac{c_1}{t} + c_2 t^2$$

Verified OK.

10.13.2 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t t gives

$$\int (t^2 y'' - 2y) dt = \int (3t^2 - 1) dt$$
$$y' t^2 - 2yt = t^3 - t + c_1$$

Which is now solved for y .

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{2}{t}$$
$$q(t) = \frac{t^3 + c_1 - t}{t^2}$$

Hence the ode is

$$y' - \frac{2y}{t} = \frac{t^3 + c_1 - t}{t^2}$$

The integrating factor μ is

$$\mu = e^{\int -\frac{2}{t} dt}$$
$$= \frac{1}{t^2}$$

The ode becomes

$$\frac{d}{dt}(\mu y) = (\mu) \left(\frac{t^3 + c_1 - t}{t^2} \right)$$
$$\frac{d}{dt} \left(\frac{y}{t^2} \right) = \left(\frac{1}{t^2} \right) \left(\frac{t^3 + c_1 - t}{t^2} \right)$$
$$d \left(\frac{y}{t^2} \right) = \left(\frac{t^3 + c_1 - t}{t^4} \right) dt$$

Integrating gives

$$\frac{y}{t^2} = \int \frac{t^3 + c_1 - t}{t^4} dt$$
$$\frac{y}{t^2} = \ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} + c_2$$

Dividing both sides by the integrating factor $\mu = \frac{1}{t^2}$ results in

$$y = t^2 \left(\ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} \right) + c_2 t^2$$

Summary

The solution(s) found are the following

$$y = t^2 \left(\ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} \right) + c_2 t^2 \quad (1)$$

Verification of solutions

$$y = t^2 \left(\ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} \right) + c_2 t^2$$

Verified OK.

10.13.3 Solving as type second_order_integrable_as_is (not using ABC version)

Writing the ode as

$$t^2 y'' - 2y = 3t^2 - 1$$

Integrating both sides of the ODE w.r.t t gives

$$\int (t^2 y'' - 2y) dt = \int (3t^2 - 1) dt$$
$$y' t^2 - 2y t = t^3 - t + c_1$$

Which is now solved for y .

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = -\frac{2}{t}$$
$$q(t) = \frac{t^3 + c_1 - t}{t^2}$$

Hence the ode is

$$y' - \frac{2y}{t} = \frac{t^3 + c_1 - t}{t^2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{2}{t} dt} \\ &= \frac{1}{t^2}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{t^3 + c_1 - t}{t^2} \right) \\ \frac{d}{dt} \left(\frac{y}{t^2} \right) &= \left(\frac{1}{t^2} \right) \left(\frac{t^3 + c_1 - t}{t^2} \right) \\ d \left(\frac{y}{t^2} \right) &= \left(\frac{t^3 + c_1 - t}{t^4} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y}{t^2} &= \int \frac{t^3 + c_1 - t}{t^4} dt \\ \frac{y}{t^2} &= \ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} + c_2\end{aligned}$$

Dividing both sides by the integrating factor $\mu = \frac{1}{t^2}$ results in

$$y = t^2 \left(\ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} \right) + c_2 t^2$$

Summary

The solution(s) found are the following

$$y = t^2 \left(\ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} \right) + c_2 t^2 \quad (1)$$

Verification of solutions

$$y = t^2 \left(\ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} \right) + c_2 t^2$$

Verified OK.

10.13.4 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' - 2y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t^2 \\ B &= 0 \\ C &= -2 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{2}{t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 2 \\ t &= t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(\frac{2}{t^2}\right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 498: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{2}{t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = 2$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = 2 \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = -1 \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = \frac{2}{t^2}$$

Since the $\text{gcd}(s, t) = 1$. This gives $b = 2$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = 2 \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = -1 \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = \frac{2}{t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	2	-1

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	2	-1

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = -1$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= -1 - (-1) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{t - c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= -\frac{1}{t} + (-)(0) \\ &= -\frac{1}{t} \\ &= -\frac{1}{t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r)p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2\left(-\frac{1}{t}\right)(0) + \left(\left(\frac{1}{t^2}\right) + \left(-\frac{1}{t}\right)^2 - \left(\frac{2}{t^2}\right)\right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int -\frac{1}{t} dt} \\ &= \frac{1}{t} \end{aligned}$$

The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned} y_1 &= z_1 \\ &= \frac{1}{t} \end{aligned}$$

Which simplifies to

$$y_1 = \frac{1}{t}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Since $B = 0$ then the above becomes

$$\begin{aligned} y_2 &= y_1 \int \frac{1}{y_1^2} dt \\ &= \frac{1}{t} \int \frac{1}{\frac{1}{t^2}} dt \\ &= \frac{1}{t} \left(\frac{t^3}{3} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(\frac{1}{t} \right) + c_2 \left(\frac{1}{t} \left(\frac{t^3}{3} \right) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$t^2y'' - 2y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = \frac{c_1}{t} + \frac{c_2t^2}{3}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1y_1 + u_2y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \frac{1}{t}$$

$$y_2 = \frac{t^2}{3}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2f(t)}{aW(t)} \tag{2}$$

$$u_2 = \int \frac{y_1f(t)}{aW(t)} \tag{3}$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{1}{t} & \frac{t^2}{3} \\ \frac{d}{dt}\left(\frac{1}{t}\right) & \frac{d}{dt}\left(\frac{t^2}{3}\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{1}{t} & \frac{t^2}{3} \\ -\frac{1}{t^2} & \frac{2t}{3} \end{vmatrix}$$

Therefore

$$W = \left(\frac{1}{t}\right) \left(\frac{2t}{3}\right) - \left(\frac{t^2}{3}\right) \left(-\frac{1}{t^2}\right)$$

Which simplifies to

$$W = 1$$

Which simplifies to

$$W = 1$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{t^2(3t^2-1)}{t^2} dt$$

Which simplifies to

$$u_1 = - \int \left(t^2 - \frac{1}{3}\right) dt$$

Hence

$$u_1 = -\frac{1}{3}t^3 + \frac{1}{3}t$$

And Eq. (3) becomes

$$u_2 = \int \frac{3t^2-1}{t^2} dt$$

Which simplifies to

$$u_2 = \int \frac{3t^2 - 1}{t^3} dt$$

Hence

$$u_2 = 3 \ln(t) + \frac{1}{2t^2}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{-\frac{1}{3}t^3 + \frac{1}{3}t}{t} + \frac{(3 \ln(t) + \frac{1}{2t^2}) t^2}{3}$$

Which simplifies to

$$y_p(t) = t^2 \ln(t) - \frac{t^2}{3} + \frac{1}{2}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{c_1}{t} + \frac{c_2 t^2}{3} \right) + \left(t^2 \ln(t) - \frac{t^2}{3} + \frac{1}{2} \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t} + \frac{c_2 t^2}{3} + t^2 \ln(t) - \frac{t^2}{3} + \frac{1}{2} \quad (1)$$

Verification of solutions

$$y = \frac{c_1}{t} + \frac{c_2 t^2}{3} + t^2 \ln(t) - \frac{t^2}{3} + \frac{1}{2}$$

Verified OK.

10.13.5 Solving as exact linear second order ode ode

An ode of the form

$$p(t) y'' + q(t) y' + r(t) y = s(t)$$

is exact if

$$p''(t) - q'(t) + r(t) = 0 \quad (1)$$

For the given ode we have

$$\begin{aligned} p(x) &= t^2 \\ q(x) &= 0 \\ r(x) &= -2 \\ s(x) &= 3t^2 - 1 \end{aligned}$$

Hence

$$\begin{aligned}p''(x) &= 2 \\q'(x) &= 0\end{aligned}$$

Therefore (1) becomes

$$2 - (0) + (-2) = 0$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$(p(t)y' + (q(t) - p'(t))y)' = s(x)$$

Integrating gives

$$p(t)y' + (q(t) - p'(t))y = \int s(t) dt$$

Substituting the above values for p, q, r, s gives

$$y't^2 - 2yt = \int 3t^2 - 1 dt$$

We now have a first order ode to solve which is

$$y't^2 - 2yt = t^3 + c_1 - t$$

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$\begin{aligned}p(t) &= -\frac{2}{t} \\q(t) &= \frac{t^3 + c_1 - t}{t^2}\end{aligned}$$

Hence the ode is

$$y' - \frac{2y}{t} = \frac{t^3 + c_1 - t}{t^2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int -\frac{2}{t} dt} \\&= \frac{1}{t^2}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{t^3 + c_1 - t}{t^2} \right) \\ \frac{d}{dt} \left(\frac{y}{t^2} \right) &= \left(\frac{1}{t^2} \right) \left(\frac{t^3 + c_1 - t}{t^2} \right) \\ d \left(\frac{y}{t^2} \right) &= \left(\frac{t^3 + c_1 - t}{t^4} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y}{t^2} &= \int \frac{t^3 + c_1 - t}{t^4} dt \\ \frac{y}{t^2} &= \ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} + c_2\end{aligned}$$

Dividing both sides by the integrating factor $\mu = \frac{1}{t^2}$ results in

$$y = t^2 \left(\ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} \right) + c_2 t^2$$

Summary

The solution(s) found are the following

$$y = t^2 \left(\ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} \right) + c_2 t^2 \quad (1)$$

Verification of solutions

$$y = t^2 \left(\ln(t) - \frac{c_1}{3t^3} + \frac{1}{2t^2} \right) + c_2 t^2$$

Verified OK.

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
<- high order exact linear fully integrable successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 22

```
dsolve(t^2*diff(y(t),t)-2*y(t) = 3*t^2-1,y(t), singsol=all)
```

$$y(t) = t^2 c_2 + \frac{1}{2} + t^2 \ln(t) + \frac{c_1}{t}$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 31

```
DSolve[t^2*y'[t]-2*y[t] == 3*t^2-1,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t^2 \log(t) + \left(-\frac{1}{3} + c_2\right) t^2 + \frac{c_1}{t} + \frac{1}{2}$$

10.14 problem 14

- 10.14.1 Solving as second order change of variable on y method 1 ode . 2840
- 10.14.2 Solving as second order change of variable on y method 2 ode . 2847
- 10.14.3 Solving using Kovacic algorithm 2851

Internal problem ID [696]

Internal file name [OUTPUT/696_Sunday_June_05_2022_01_47_08_AM_14861358/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 14.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_change_of_variable_on_y_method_1", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$t^2 y'' - t(2+t)y' + (2+t)y = 2t^3$$

10.14.1 Solving as second order change of variable on y method 1 ode

This is second order non-homogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$.
 y_h is the solution to

$$t^2 y'' + (-t^2 - 2t)y' + (2+t)y = 0$$

In normal form the given ode is written as

$$y'' + p(t)y' + q(t)y = 0 \tag{2}$$

Where

$$p(t) = \frac{-t^2 - 2t}{t^2}$$

$$q(t) = \frac{2 + t}{t^2}$$

Calculating the Liouville ode invariant Q given by

$$\begin{aligned} Q &= q - \frac{p'}{2} - \frac{p^2}{4} \\ &= \frac{2 + t}{t^2} - \frac{\left(\frac{-t^2 - 2t}{t^2}\right)'}{2} - \frac{\left(\frac{-t^2 - 2t}{t^2}\right)^2}{4} \\ &= \frac{2 + t}{t^2} - \frac{\left(\frac{-2t - 2}{t^2} - \frac{2(-t^2 - 2t)}{t^3}\right)}{2} - \frac{\left(\frac{(-t^2 - 2t)^2}{t^4}\right)}{4} \\ &= \frac{2 + t}{t^2} - \left(\frac{-2t - 2}{2t^2} - \frac{-t^2 - 2t}{t^3}\right) - \frac{(-t^2 - 2t)^2}{4t^4} \\ &= -\frac{1}{4} \end{aligned}$$

Since the Liouville ode invariant does not depend on the independent variable t then the transformation

$$y = v(t) z(t) \tag{3}$$

is used to change the original ode to a constant coefficients ode in v . In (3) the term $z(t)$ is given by

$$\begin{aligned} z(t) &= e^{-\left(\int \frac{p(t)}{2} dt\right)} \\ &= e^{-\int \frac{-t^2 - 2t}{2} dt} \\ &= t e^{\frac{t}{2}} \end{aligned} \tag{5}$$

Hence (3) becomes

$$y = v(t) t e^{\frac{t}{2}} \tag{4}$$

Applying this change of variable to the original ode results in

$$-e^{\frac{t}{2}}(-4v''(t) + v(t)) = 8$$

Which is now solved for $v(t)$ Simplifying the ode gives

$$v''(t) - \frac{v(t)}{4} = 2 e^{-\frac{t}{2}}$$

This is second order non-homogeneous ODE. In standard form the ODE is

$$Av''(t) + Bv'(t) + Cv(t) = f(t)$$

Where $A = 1, B = 0, C = -\frac{1}{4}, f(t) = 2e^{-\frac{t}{2}}$. Let the solution be

$$v(t) = v_h + v_p$$

Where v_h is the solution to the homogeneous ODE $Av''(t) + Bv'(t) + Cv(t) = 0$, and v_p is a particular solution to the non-homogeneous ODE $Av''(t) + Bv'(t) + Cv(t) = f(t)$. v_h is the solution to

$$v''(t) - \frac{v(t)}{4} = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Av''(t) + Bv'(t) + Cv(t) = 0$$

Where in the above $A = 1, B = 0, C = -\frac{1}{4}$. Let the solution be $v(t) = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} - \frac{e^{\lambda t}}{4} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 - \frac{1}{4} = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = -\frac{1}{4}$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1) \left(-\frac{1}{4}\right)} \\ &= \pm \frac{1}{2} \end{aligned}$$

Hence

$$\lambda_1 = +\frac{1}{2}$$

$$\lambda_2 = -\frac{1}{2}$$

Which simplifies to

$$\lambda_1 = \frac{1}{2}$$
$$\lambda_2 = -\frac{1}{2}$$

Since roots are real and distinct, then the solution is

$$v(t) = c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t}$$
$$v(t) = c_1 e^{(\frac{1}{2})t} + c_2 e^{(-\frac{1}{2})t}$$

Or

$$v(t) = c_1 e^{\frac{t}{2}} + c_2 e^{-\frac{t}{2}}$$

Therefore the homogeneous solution v_h is

$$v_h = c_1 e^{\frac{t}{2}} + c_2 e^{-\frac{t}{2}}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$2e^{-\frac{t}{2}}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$\left[\left\{ e^{-\frac{t}{2}} \right\} \right]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ e^{-\frac{t}{2}}, e^{\frac{t}{2}} \right\}$$

Since $e^{-\frac{t}{2}}$ is duplicated in the UC_set, then this basis is multiplied by extra t . The UC_set becomes

$$\left[\left\{ t e^{-\frac{t}{2}} \right\} \right]$$

Since there was duplication between the basis functions in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis function in the above updated UC_set.

$$v_p = A_1 t e^{-\frac{t}{2}}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution v_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-A_1 e^{-\frac{t}{2}} = 2 e^{-\frac{t}{2}}$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = -2]$$

Substituting the above back in the above trial solution v_p , gives the particular solution

$$v_p = -2t e^{-\frac{t}{2}}$$

Therefore the general solution is

$$\begin{aligned} v &= v_h + v_p \\ &= \left(c_1 e^{\frac{t}{2}} + c_2 e^{-\frac{t}{2}} \right) + \left(-2t e^{-\frac{t}{2}} \right) \end{aligned}$$

Now that $v(t)$ is known, then

$$\begin{aligned} y &= v(t) z(t) \\ &= \left(c_1 e^{\frac{t}{2}} + c_2 e^{-\frac{t}{2}} - 2t e^{-\frac{t}{2}} \right) (z(t)) \end{aligned} \tag{7}$$

But from (5)

$$z(t) = t e^{\frac{t}{2}}$$

Hence (7) becomes

$$y = \left(c_1 e^{\frac{t}{2}} + c_2 e^{-\frac{t}{2}} - 2t e^{-\frac{t}{2}} \right) t e^{\frac{t}{2}}$$

Therefore the homogeneous solution y_h is

$$y_h = \left(c_1 e^{\frac{t}{2}} + c_2 e^{-\frac{t}{2}} - 2t e^{-\frac{t}{2}} \right) t e^{\frac{t}{2}}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of

parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = t e^t$$

$$y_2 = e^{-\frac{t}{2}} t e^{\frac{t}{2}}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{a W(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{a W(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t e^t & e^{-\frac{t}{2}} t e^{\frac{t}{2}} \\ \frac{d}{dt}(t e^t) & \frac{d}{dt}(e^{-\frac{t}{2}} t e^{\frac{t}{2}}) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t e^t & e^{-\frac{t}{2}} t e^{\frac{t}{2}} \\ t e^t + e^t & e^{-\frac{t}{2}} e^{\frac{t}{2}} \end{vmatrix}$$

Therefore

$$W = (t e^t) \left(e^{-\frac{t}{2}} e^{\frac{t}{2}} \right) - \left(e^{-\frac{t}{2}} t e^{\frac{t}{2}} \right) (t e^t + e^t)$$

Which simplifies to

$$W = -e^{-\frac{t}{2}} e^{\frac{3t}{2}} t^2$$

Which simplifies to

$$W = -t^2 e^t$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{2e^{-\frac{t}{2}} t^4 e^{\frac{t}{2}}}{-t^4 e^t} dt$$

Which simplifies to

$$u_1 = - \int -2e^{-t} dt$$

Hence

$$u_1 = -2e^{-t}$$

And Eq. (3) becomes

$$u_2 = \int \frac{2t^4 e^t}{-t^4 e^t} dt$$

Which simplifies to

$$u_2 = \int (-2) dt$$

Hence

$$u_2 = -2t$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -2e^{-t} e^t t - 2t^2 e^{-\frac{t}{2}} e^{\frac{t}{2}}$$

Which simplifies to

$$y_p(t) = -2t^2 - 2t$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left((c_1 e^{\frac{t}{2}} + c_2 e^{-\frac{t}{2}} - 2t e^{-\frac{t}{2}}) t e^{\frac{t}{2}} \right) + (-2t^2 - 2t) \end{aligned}$$

Which simplifies to

$$y = t(c_1 e^t + c_2 - 2t) - 2t^2 - 2t$$

Summary

The solution(s) found are the following

$$y = t(c_1 e^t + c_2 - 2t) - 2t^2 - 2t \quad (1)$$

Verification of solutions

$$y = t(c_1 e^t + c_2 - 2t) - 2t^2 - 2t$$

Verified OK.

10.14.2 Solving as second order change of variable on y method 2 ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = t^2$, $B = -t^2 - 2t$, $C = 2 + t$, $f(t) = 2t^3$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. Solving for y_h from

$$t^2 y'' + (-t^2 - 2t) y' + (2 + t) y = 0$$

In normal form the ode

$$t^2 y'' + (-t^2 - 2t) y' + (2 + t) y = 0 \quad (1)$$

Becomes

$$y'' + p(t) y' + q(t) y = 0 \quad (2)$$

Where

$$p(t) = \frac{-t-2}{t}$$
$$q(t) = \frac{2+t}{t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right) v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right) v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} + \frac{n(-t-2)}{t^2} + \frac{2+t}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = 1 \quad (6)$$

Substituting this value in (3) gives

$$\begin{aligned} v''(t) + \left(\frac{2}{t} + \frac{-t-2}{t} \right) v'(t) &= 0 \\ v''(t) - v'(t) &= 0 \end{aligned} \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) - u(t) = 0 \quad (8)$$

The above is now solved for $u(t)$. Integrating both sides gives

$$\begin{aligned} \int \frac{1}{u} du &= t + c_1 \\ \ln(u) &= t + c_1 \\ u &= e^{t+c_1} \\ u &= c_1 e^t \end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned} v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= c_1 e^t + c_2 \end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\ &= (c_1 e^t + c_2) t \\ &= (c_1 e^t + c_2) t\end{aligned}$$

Now the particular solution to this ODE is found

$$t^2 y'' + (-t^2 - 2t) y' + (2 + t) y = 2t^3$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned}y_1 &= t \\ y_2 &= t e^t\end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t & t e^t \\ \frac{d}{dt}(t) & \frac{d}{dt}(t e^t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t & t e^t \\ 1 & t e^t + e^t \end{vmatrix}$$

Therefore

$$W = (t)(t e^t + e^t) - (t e^t) \quad (1)$$

Which simplifies to

$$W = t^2 e^t$$

Which simplifies to

$$W = t^2 e^t$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{2t^4 e^t}{t^4 e^t} dt$$

Which simplifies to

$$u_1 = - \int 2 dt$$

Hence

$$u_1 = -2t$$

And Eq. (3) becomes

$$u_2 = \int \frac{2t^4}{t^4 e^t} dt$$

Which simplifies to

$$u_2 = \int 2 e^{-t} dt$$

Hence

$$u_2 = -2 e^{-t}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -2t^2 - 2 e^{-t} e^t t$$

Which simplifies to

$$y_p(t) = -2t^2 - 2t$$

Therefore the general solution is

$$\begin{aligned}y &= y_h + y_p \\&= ((c_1 e^t + c_2) t) + (-2t^2 - 2t) \\&= -2t^2 - 2t + (c_1 e^t + c_2) t\end{aligned}$$

Which simplifies to

$$y = t(c_1 e^t + c_2 - 2t - 2)$$

Summary

The solution(s) found are the following

$$y = t(c_1 e^t + c_2 - 2t - 2) \tag{1}$$

Verification of solutions

$$y = t(c_1 e^t + c_2 - 2t - 2)$$

Verified OK.

10.14.3 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' + (-t^2 - 2t) y' + (2 + t) y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= t^2 \\B &= -t^2 - 2t \\C &= 2 + t\end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{4} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \frac{z(t)}{4} \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 499: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{4}$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = e^{-\frac{t}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-t^2-2t}{t^2} dt} \\ &= z_1 e^{\frac{t}{2} + \ln(t)} \\ &= z_1 \left(t e^{\frac{t}{2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = t$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-t^2-2t}{t^2} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{t+2\ln(t)}}{(y_1)^2} dt \\ &= y_1 (e^t) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1(t) + c_2(t(e^t))\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$t^2 y'' + (-t^2 - 2t) y' + (2 + t) y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 t + c_2 t e^t$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = t$$

$$y_2 = t e^t$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \tag{3}$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t & t e^t \\ \frac{d}{dt}(t) & \frac{d}{dt}(t e^t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t & t e^t \\ 1 & t e^t + e^t \end{vmatrix}$$

Therefore

$$W = (t)(t e^t + e^t) - (t e^t) \quad (1)$$

Which simplifies to

$$W = t^2 e^t$$

Which simplifies to

$$W = t^2 e^t$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{2t^4 e^t}{t^4 e^t} dt$$

Which simplifies to

$$u_1 = - \int 2 dt$$

Hence

$$u_1 = -2t$$

And Eq. (3) becomes

$$u_2 = \int \frac{2t^4}{t^4 e^t} dt$$

Which simplifies to

$$u_2 = \int 2 e^{-t} dt$$

Hence

$$u_2 = -2e^{-t}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -2t^2 - 2e^{-t}e^t t$$

Which simplifies to

$$y_p(t) = -2t^2 - 2t$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 t + c_2 t e^t) + (-2t^2 - 2t) \end{aligned}$$

Which simplifies to

$$y = t(c_1 + c_2 e^t) - 2t^2 - 2t$$

Summary

The solution(s) found are the following

$$y = t(c_1 + c_2 e^t) - 2t^2 - 2t \quad (1)$$

Verification of solutions

$$y = t(c_1 + c_2 e^t) - 2t^2 - 2t$$

Verified OK.

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    checking if the LODE is of Euler type
    trying a symmetry of the form [xi=0, eta=F(x)]
    checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Reducible group (found an exponential solution)
    Reducible group (found another exponential solution)
<- Kovacics algorithm successful
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.281 (sec). Leaf size: 15

```
dsolve(t^2*diff(y(t),t$2)-t*(t+2)*diff(y(t),t)+(t+2)*y(t) = 2*t^3,y(t), singsol=all)
```

$$y(t) = t(e^t c_1 + c_2 - 2t)$$

✓ Solution by Mathematica

Time used: 0.025 (sec). Leaf size: 20

```
DSolve[t^2*y''[t]-t*(t+2)*y'[t]+(t+2)*y[t] == 2*t^3,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t(-2t + c_2 e^t - 2 + c_1)$$

10.15 problem 15

10.15.1 Solving as second order ode non constant coeff transformation on B ode	2858
10.15.2 Solving using Kovacic algorithm	2863

Internal problem ID [697]

Internal file name [OUTPUT/697_Sunday_June_05_2022_01_47_10_AM_66662417/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 15.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_ode_non_constant_coeff_transformation_on_B"

Maple gives the following as the ode type

```
[[_2nd_order, _with_linear_symmetries]]
```

$$ty'' - (t + 1)y' + y = e^{2t}t^2$$

10.15.1 Solving as second order ode non constant coeff transformation on B ode

Given an ode of the form

$$Ay'' + By' + Cy = F(t)$$

This method reduces the order ode the ODE by one by applying the transformation

$$y = Bv$$

This results in

$$\begin{aligned}y' &= B'v + v'B \\y'' &= B''v + B'v' + v''B + v'B' \\ &= v''B + 2v' + B' + B''v\end{aligned}$$

And now the original ode becomes

$$\begin{aligned} A(v''B + 2v'B' + B''v) + B(B'v + v'B) + CBv &= 0 \\ ABv'' + (2AB' + B^2)v' + (AB'' + BB' + CB)v &= 0 \end{aligned} \quad (1)$$

If the term $AB'' + BB' + CB$ is zero, then this method works and can be used to solve

$$ABv'' + (2AB' + B^2)v' = 0$$

By Using $u = v'$ which reduces the order of the above ode to one. The new ode is

$$ABu' + (2AB' + B^2)u = 0$$

The above ode is first order ode which is solved for u . Now a new ode $v' = u$ is solved for v as first order ode. Then the final solution is obtain from $y = Bv$.

This method works only if the term $AB'' + BB' + CB$ is zero. The given ODE shows that

$$\begin{aligned} A &= t \\ B &= -t - 1 \\ C &= 1 \\ F &= e^{2t}t^2 \end{aligned}$$

The above shows that for this ode

$$\begin{aligned} AB'' + BB' + CB &= (t)(0) + (-t - 1)(-1) + (1)(-t - 1) \\ &= 0 \end{aligned}$$

Hence the ode in v given in (1) now simplifies to

$$-t(t + 1)v'' + (t^2 + 1)v' = 0$$

Now by applying $v' = u$ the above becomes

$$-t(t + 1)u'(t) + (t^2 + 1)u(t) = 0$$

Which is now solved for u . In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{(t^2 + 1)u}{t(t + 1)} \end{aligned}$$

Where $f(t) = \frac{t^2+1}{t(t+1)}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= \frac{t^2 + 1}{t(t+1)} dt \\ \int \frac{1}{u} du &= \int \frac{t^2 + 1}{t(t+1)} dt \\ \ln(u) &= t + \ln(t) - 2 \ln(t+1) + c_1 \\ u &= e^{t+\ln(t)-2\ln(t+1)+c_1} \\ &= c_1 e^{t+\ln(t)-2\ln(t+1)}\end{aligned}$$

Which simplifies to

$$u(t) = \frac{c_1 e^{tt}}{(t+1)^2}$$

The ode for v now becomes

$$\begin{aligned}v' &= u \\ &= \frac{c_1 e^{tt}}{(t+1)^2}\end{aligned}$$

Which is now solved for v . Integrating both sides gives

$$\begin{aligned}v(t) &= \int \frac{c_1 e^{tt}}{(t+1)^2} dt \\ &= \frac{c_1 e^t}{t+1} + c_2\end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}y_h(t) &= Bv \\ &= (-t-1) \left(\frac{c_1 e^t}{t+1} + c_2 \right) \\ &= -c_1 e^t - c_2(t+1)\end{aligned}$$

And now the particular solution $y_p(t)$ will be found. The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = -t - 1$$

$$y_2 = e^t$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} -t - 1 & e^t \\ \frac{d}{dt}(-t - 1) & \frac{d}{dt}(e^t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} -t - 1 & e^t \\ -1 & e^t \end{vmatrix}$$

Therefore

$$W = (-t - 1)(e^t) - (e^t)(-1)$$

Which simplifies to

$$W = -te^t$$

Which simplifies to

$$W = -te^t$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{e^t e^{2t} t^2}{-t^2 e^t} dt$$

Which simplifies to

$$u_1 = - \int -e^{2t} dt$$

Hence

$$u_1 = \frac{e^{2t}}{2}$$

And Eq. (3) becomes

$$u_2 = \int \frac{(-t-1)e^{2t}t^2}{-t^2e^t} dt$$

Which simplifies to

$$u_2 = \int e^t(t+1) dt$$

Hence

$$u_2 = t e^t$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{(-t-1)e^{2t}}{2} + e^{2t}t$$

Which simplifies to

$$y_p(t) = \frac{e^{2t}(-1+t)}{2}$$

Hence the complete solution is

$$\begin{aligned} y(t) &= y_h + y_p \\ &= (-c_1e^t - c_2(t+1)) + \left(\frac{e^{2t}(-1+t)}{2} \right) \\ &= -c_1e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = -c_1e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2} \quad (1)$$

Verification of solutions

$$y = -c_1 e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2}$$

Verified OK.

10.15.2 Solving using Kovacic algorithm

Writing the ode as

$$ty'' + (-t-1)y' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t \\ B &= -t-1 \\ C &= 1 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{t^2 - 2t + 3}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= t^2 - 2t + 3 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(\frac{t^2 - 2t + 3}{4t^2} \right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 500: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 2 \\ &= 0 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Therefore

$$L = [1, 2]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{1}{4} - \frac{1}{2t} + \frac{3}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = \frac{3}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{3}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is $O_r(\infty) = 0$ then

$$v = \frac{-O_r(\infty)}{2} = \frac{0}{2} = 0$$

$[\sqrt{r}]_\infty$ is the sum of terms involving t^i for $0 \leq i \leq v$ in the Laurent series for \sqrt{r} at ∞ . Therefore

$$\begin{aligned} [\sqrt{r}]_\infty &= \sum_{i=0}^v a_i t^i \\ &= \sum_{i=0}^0 a_i t^i \end{aligned} \tag{8}$$

Let a be the coefficient of $t^v = t^0$ in the above sum. The Laurent series of \sqrt{r} at ∞ is

$$\sqrt{r} \approx \frac{1}{2} - \frac{1}{2t} + \frac{1}{2t^2} + \frac{1}{2t^3} + \frac{1}{4t^4} - \frac{1}{4t^5} - \frac{3}{4t^6} - \frac{3}{4t^7} + \dots \tag{9}$$

Comparing Eq. (9) with Eq. (8) shows that

$$a = \frac{1}{2}$$

From Eq. (9) the sum up to $v = 0$ gives

$$\begin{aligned} [\sqrt{r}]_\infty &= \sum_{i=0}^0 a_i t^i \\ &= \frac{1}{2} \end{aligned} \tag{10}$$

Now we need to find b , where b be the coefficient of $t^{v-1} = t^{-1} = \frac{1}{t}$ in r minus the coefficient of same term but in $([\sqrt{r}]_\infty)^2$ where $[\sqrt{r}]_\infty$ was found above in Eq (10). Hence

$$([\sqrt{r}]_\infty)^2 = \frac{1}{4}$$

This shows that the coefficient of $\frac{1}{t}$ in the above is 0. Now we need to find the coefficient of $\frac{1}{t}$ in r . How this is done depends on if $v = 0$ or not. Since $v = 0$ then starting from $r = \frac{s}{t}$ and doing long division in the form

$$r = Q + \frac{R}{t}$$

Where Q is the quotient and R is the remainder. Then the coefficient of $\frac{1}{t}$ in r will be the coefficient in R of the term in t of degree of t minus one, divided by the leading coefficient in t . Doing long division gives

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{t^2 - 2t + 3}{4t^2} \\ &= Q + \frac{R}{4t^2} \\ &= \left(\frac{1}{4}\right) + \left(\frac{-2t + 3}{4t^2}\right) \\ &= \frac{1}{4} + \frac{-2t + 3}{4t^2} \end{aligned}$$

Since the degree of t is 2, then we see that the coefficient of the term t in the remainder R is -2 . Dividing this by leading coefficient in t which is 4 gives $-\frac{1}{2}$. Now b can be found.

$$\begin{aligned} b &= \left(-\frac{1}{2}\right) - (0) \\ &= -\frac{1}{2} \end{aligned}$$

Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= \frac{1}{2} \\ \alpha_\infty^+ &= \frac{1}{2} \left(\frac{b}{a} - v\right) = \frac{1}{2} \left(\frac{-\frac{1}{2}}{\frac{1}{2}} - 0\right) = -\frac{1}{2} \\ \alpha_\infty^- &= \frac{1}{2} \left(-\frac{b}{a} - v\right) = \frac{1}{2} \left(-\frac{-\frac{1}{2}}{\frac{1}{2}} - 0\right) = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = \frac{t^2 - 2t + 3}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{3}{2}$	$-\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
0	$\frac{1}{2}$	$-\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_\infty^+ = -\frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_\infty^+ - (\alpha_{c_1}^-) \\ &= -\frac{1}{2} - \left(-\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c) [\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty) [\sqrt{r}]_\infty$$

Substituting the above values in the above results in

$$\begin{aligned} \omega &= \left((-) [\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^-}{t - c_1} \right) + (+) [\sqrt{r}]_\infty \\ &= -\frac{1}{2t} + \left(\frac{1}{2} \right) \\ &= \frac{1}{2} - \frac{1}{2t} \\ &= \frac{-1 + t}{2t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(\frac{1}{2} - \frac{1}{2t}\right)(0) + \left(\left(\frac{1}{2t^2}\right) + \left(\frac{1}{2} - \frac{1}{2t}\right)^2 - \left(\frac{t^2 - 2t + 3}{4t^2}\right)\right) = 0$$

$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \left(\frac{1}{2} - \frac{1}{2t}\right) dt} \\ &= \frac{e^{\frac{t}{2}}}{\sqrt{t}} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-t-1}{t} dt} \\ &= z_1 e^{\frac{t}{2} + \frac{\ln(t)}{2}} \\ &= z_1 \left(\sqrt{t} e^{\frac{t}{2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^t$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-t-1}{t} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{t+\ln(t)}}{(y_1)^2} dt \\&= y_1(-e^{-t}(t+1))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1(e^t) + c_2(e^t(-e^{-t}(t+1)))\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$ty'' + (-t-1)y' + y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^t + (-t-1)c_2$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned}y_1 &= e^t \\y_2 &= -t-1\end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} e^t & -t - 1 \\ \frac{d}{dt}(e^t) & \frac{d}{dt}(-t - 1) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} e^t & -t - 1 \\ e^t & -1 \end{vmatrix}$$

Therefore

$$W = (e^t)(-1) - (-t - 1)(e^t)$$

Which simplifies to

$$W = t e^t$$

Which simplifies to

$$W = t e^t$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{(-t - 1) e^{2t} t^2}{t^2 e^t} dt$$

Which simplifies to

$$u_1 = - \int -e^t(t + 1) dt$$

Hence

$$u_1 = t e^t$$

And Eq. (3) becomes

$$u_2 = \int \frac{e^t e^{2t} t^2}{t^2 e^t} dt$$

Which simplifies to

$$u_2 = \int e^{2t} dt$$

Hence

$$u_2 = \frac{e^{2t}}{2}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{(-t-1)e^{2t}}{2} + e^{2t}t$$

Which simplifies to

$$y_p(t) = \frac{e^{2t}(-1+t)}{2}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 e^t + (-t-1)c_2) + \left(\frac{e^{2t}(-1+t)}{2} \right) \end{aligned}$$

Which simplifies to

$$y = c_1 e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2}$$

Summary

The solution(s) found are the following

$$y = c_1 e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2} \tag{1}$$

Verification of solutions

$$y = c_1 e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2}$$

Verified OK.

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    checking if the LODE is of Euler type
    trying a symmetry of the form [xi=0, eta=F(x)]
    checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Reducible group (found an exponential solution)
    Reducible group (found another exponential solution)
<- Kovacics algorithm successful
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 23

```
dsolve(t*difff(y(t),t$2)-(1+t)*difff(y(t),t)+y(t) = t^2*exp(2*t),y(t), singsol=all)
```

$$y(t) = (t + 1) c_2 + e^t c_1 + \frac{(t - 1) e^{2t}}{2}$$

✓ Solution by Mathematica

Time used: 0.052 (sec). Leaf size: 31

```
DSolve[t*y'[t]-(1+t)*y'[t]+y[t] == t^2*Exp[2*t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{2} e^{2t} (t - 1) + c_1 e^t - c_2 (t + 1)$$

10.16 problem 16

- 10.16.1 Solving as second order change of variable on y method 2 ode . 2873
- 10.16.2 Solving as second order ode non constant coeff transformation
on B ode 2878
- 10.16.3 Solving using Kovacic algorithm 2883

Internal problem ID [698]

Internal file name [OUTPUT/698_Sunday_June_05_2022_01_47_12_AM_63605988/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 16.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_change_of_variable_on_y_method_2", "second_order_ode_non_constant_coeff_transformation_on_B"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(1 - t)y'' + ty' - y = 2(-1 + t)^2 e^{-t}$$

10.16.1 Solving as second order change of variable on y method 2 ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1 - t$, $B = t$, $C = -1$, $f(t) = 2(-1 + t)^2 e^{-t}$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. Solving for y_h from

$$(1 - t)y'' + ty' - y = 0$$

In normal form the ode

$$(1 - t)y'' + ty' - y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = -\frac{t}{-1+t}$$
$$q(t) = \frac{1}{-1+t}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} - \frac{n}{-1+t} + \frac{1}{-1+t} = 0 \quad (5)$$

Solving (5) for n gives

$$n = 1 \quad (6)$$

Substituting this value in (3) gives

$$v''(t) + \left(\frac{2}{t} - \frac{t}{-1+t}\right)v'(t) = 0$$
$$v''(t) + \left(\frac{2}{t} - \frac{t}{-1+t}\right)v'(t) = 0 \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \left(\frac{2}{t} - \frac{t}{-1+t} \right) u(t) = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{u(t^2 - 2t + 2)}{t(-1+t)} \end{aligned}$$

Where $f(t) = \frac{t^2 - 2t + 2}{t(-1+t)}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= \frac{t^2 - 2t + 2}{t(-1+t)} dt \\ \int \frac{1}{u} du &= \int \frac{t^2 - 2t + 2}{t(-1+t)} dt \\ \ln(u) &= t - 2 \ln(t) + \ln(-1+t) + c_1 \\ u &= e^{t-2 \ln(t)+\ln(-1+t)+c_1} \\ &= c_1 e^{t-2 \ln(t)+\ln(-1+t)} \end{aligned}$$

Which simplifies to

$$u(t) = c_1 \left(-\frac{e^t}{t^2} + \frac{e^t}{t} \right)$$

Now that $u(t)$ is known, then

$$\begin{aligned} v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= \frac{e^t c_1}{t} + c_2 \end{aligned}$$

Hence

$$\begin{aligned} y &= v(t) t^n \\ &= \left(\frac{e^t c_1}{t} + c_2 \right) t \\ &= c_1 e^t + c_2 t \end{aligned}$$

Now the particular solution to this ODE is found

$$(1 - t) y'' + ty' - y = 2(-1 + t)^2 e^{-t}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = t$$

$$y_2 = e^t$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t & e^t \\ \frac{d}{dt}(t) & \frac{d}{dt}(e^t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t & e^t \\ 1 & e^t \end{vmatrix}$$

Therefore

$$W = (t)(e^t) - (e^t)(1)$$

Which simplifies to

$$W = t e^t - e^t$$

Which simplifies to

$$W = e^t(-1 + t)$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{2 e^t(-1 + t)^2 e^{-t}}{(1 - t) e^t (-1 + t)} dt$$

Which simplifies to

$$u_1 = - \int -2 e^{-t} dt$$

Hence

$$u_1 = -2 e^{-t}$$

And Eq. (3) becomes

$$u_2 = \int \frac{2t(-1 + t)^2 e^{-t}}{(1 - t) e^t (-1 + t)} dt$$

Which simplifies to

$$u_2 = \int -2t e^{-2t} dt$$

Hence

$$u_2 = \frac{(1 + 2t) e^{-2t}}{2}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -2t e^{-t} + \frac{(1 + 2t) e^{-2t} e^t}{2}$$

Which simplifies to

$$y_p(t) = e^{-t} \left(-t + \frac{1}{2} \right)$$

Therefore the general solution is

$$\begin{aligned}y &= y_h + y_p \\&= \left(\left(\frac{e^t c_1}{t} + c_2 \right) t \right) + \left(e^{-t} \left(-t + \frac{1}{2} \right) \right) \\&= e^{-t} \left(-t + \frac{1}{2} \right) + \left(\frac{e^t c_1}{t} + c_2 \right) t\end{aligned}$$

Which simplifies to

$$y = e^{-t} \left(-t + \frac{1}{2} \right) + \left(\frac{e^t c_1}{t} + c_2 \right) t$$

Summary

The solution(s) found are the following

$$y = e^{-t} \left(-t + \frac{1}{2} \right) + \left(\frac{e^t c_1}{t} + c_2 \right) t \quad (1)$$

Verification of solutions

$$y = e^{-t} \left(-t + \frac{1}{2} \right) + \left(\frac{e^t c_1}{t} + c_2 \right) t$$

Verified OK.

10.16.2 Solving as second order ode non constant coeff transformation on B ode

Given an ode of the form

$$Ay'' + By' + Cy = F(t)$$

This method reduces the order ode the ODE by one by applying the transformation

$$y = Bv$$

This results in

$$\begin{aligned}y' &= B'v + v'B \\y'' &= B''v + B'v' + v''B + v'B' \\&= v''B + 2v' + B' + B''v\end{aligned}$$

And now the original ode becomes

$$\begin{aligned} A(v''B + 2v'B' + B''v) + B(B'v + v'B) + CBv &= 0 \\ ABv'' + (2AB' + B^2)v' + (AB'' + BB' + CB)v &= 0 \end{aligned} \quad (1)$$

If the term $AB'' + BB' + CB$ is zero, then this method works and can be used to solve

$$ABv'' + (2AB' + B^2)v' = 0$$

By Using $u = v'$ which reduces the order of the above ode to one. The new ode is

$$ABu' + (2AB' + B^2)u = 0$$

The above ode is first order ode which is solved for u . Now a new ode $v' = u$ is solved for v as first order ode. Then the final solution is obtain from $y = Bv$.

This method works only if the term $AB'' + BB' + CB$ is zero. The given ODE shows that

$$\begin{aligned} A &= 1 - t \\ B &= t \\ C &= -1 \\ F &= 2(-1 + t)^2 e^{-t} \end{aligned}$$

The above shows that for this ode

$$\begin{aligned} AB'' + BB' + CB &= (1 - t)(0) + (t)(1) + (-1)(t) \\ &= 0 \end{aligned}$$

Hence the ode in v given in (1) now simplifies to

$$-t(-1 + t)v'' + (t^2 - 2t + 2)v' = 0$$

Now by applying $v' = u$ the above becomes

$$(-t^2 + t)u'(t) + (t^2 - 2t + 2)u(t) = 0$$

Which is now solved for u . In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{u(t^2 - 2t + 2)}{t(-1 + t)} \end{aligned}$$

Where $f(t) = \frac{t^2-2t+2}{t(-1+t)}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= \frac{t^2 - 2t + 2}{t(-1+t)} dt \\ \int \frac{1}{u} du &= \int \frac{t^2 - 2t + 2}{t(-1+t)} dt \\ \ln(u) &= t - 2 \ln(t) + \ln(-1+t) + c_1 \\ u &= e^{t-2\ln(t)+\ln(-1+t)+c_1} \\ &= c_1 e^{t-2\ln(t)+\ln(-1+t)}\end{aligned}$$

Which simplifies to

$$u(t) = c_1 \left(-\frac{e^t}{t^2} + \frac{e^t}{t} \right)$$

The ode for v now becomes

$$\begin{aligned}v' &= u \\ &= c_1 \left(-\frac{e^t}{t^2} + \frac{e^t}{t} \right)\end{aligned}$$

Which is now solved for v . Integrating both sides gives

$$\begin{aligned}v(t) &= \int \frac{c_1 e^t (-1+t)}{t^2} dt \\ &= \frac{e^t c_1}{t} + c_2\end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}y_h(t) &= Bv \\ &= (t) \left(\frac{e^t c_1}{t} + c_2 \right) \\ &= c_1 e^t + c_2 t\end{aligned}$$

And now the particular solution $y_p(t)$ will be found. The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = t$$

$$y_2 = e^t$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t & e^t \\ \frac{d}{dt}(t) & \frac{d}{dt}(e^t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t & e^t \\ 1 & e^t \end{vmatrix}$$

Therefore

$$W = (t)(e^t) - (e^t)(1)$$

Which simplifies to

$$W = te^t - e^t$$

Which simplifies to

$$W = e^t(-1 + t)$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{2e^t(-1+t)^2 e^{-t}}{(1-t)e^t(-1+t)} dt$$

Which simplifies to

$$u_1 = - \int -2 e^{-t} dt$$

Hence

$$u_1 = -2 e^{-t}$$

And Eq. (3) becomes

$$u_2 = \int \frac{2t(-1+t)^2 e^{-t}}{(1-t) e^t (-1+t)} dt$$

Which simplifies to

$$u_2 = \int -2t e^{-2t} dt$$

Hence

$$u_2 = \frac{(1+2t) e^{-2t}}{2}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -2t e^{-t} + \frac{(1+2t) e^{-2t} e^t}{2}$$

Which simplifies to

$$y_p(t) = e^{-t} \left(-t + \frac{1}{2} \right)$$

Hence the complete solution is

$$\begin{aligned} y(t) &= y_h + y_p \\ &= (c_1 e^t + c_2 t) + \left(e^{-t} \left(-t + \frac{1}{2} \right) \right) \\ &= -t e^{-t} + \frac{e^{-t}}{2} + c_1 e^t + c_2 t \end{aligned}$$

Summary

The solution(s) found are the following

$$y = -t e^{-t} + \frac{e^{-t}}{2} + c_1 e^t + c_2 t \quad (1)$$

Verification of solutions

$$y = -te^{-t} + \frac{e^{-t}}{2} + c_1e^t + c_2t$$

Verified OK.

10.16.3 Solving using Kovacic algorithm

Writing the ode as

$$(1-t)y'' + ty' - y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1-t \\ B &= t \\ C &= -1 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{t^2 - 4t + 6}{4(-1+t)^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= t^2 - 4t + 6 \\ t &= 4(-1+t)^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(\frac{t^2 - 4t + 6}{4(-1 + t)^2} \right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 501: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 2 \\ &= 0 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4(-1 + t)^2$. There is a pole at $t = 1$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Therefore

$$L = [1, 2]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{1}{4} - \frac{1}{2(-1+t)} + \frac{3}{4(-1+t)^2}$$

For the pole at $t = 1$ let b be the coefficient of $\frac{1}{(-1+t)^2}$ in the partial fractions decomposition of r given above. Therefore $b = \frac{3}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1+4b} = \frac{3}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1+4b} = -\frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is $O_r(\infty) = 0$ then

$$v = \frac{-O_r(\infty)}{2} = \frac{0}{2} = 0$$

$[\sqrt{r}]_\infty$ is the sum of terms involving t^i for $0 \leq i \leq v$ in the Laurent series for \sqrt{r} at ∞ . Therefore

$$\begin{aligned} [\sqrt{r}]_\infty &= \sum_{i=0}^v a_i t^i \\ &= \sum_{i=0}^0 a_i t^i \end{aligned} \tag{8}$$

Let a be the coefficient of $t^v = t^0$ in the above sum. The Laurent series of \sqrt{r} at ∞ is

$$\sqrt{r} \approx \frac{1}{2} - \frac{1}{2t} + \frac{1}{t^3} + \frac{11}{4t^4} + \frac{21}{4t^5} + \frac{15}{2t^6} + \frac{6}{t^7} - \frac{117}{16t^8} + \dots \tag{9}$$

Comparing Eq. (9) with Eq. (8) shows that

$$a = \frac{1}{2}$$

From Eq. (9) the sum up to $v = 0$ gives

$$\begin{aligned} [\sqrt{r}]_\infty &= \sum_{i=0}^0 a_i t^i \\ &= \frac{1}{2} \end{aligned} \tag{10}$$

Now we need to find b , where b be the coefficient of $t^{v-1} = t^{-1} = \frac{1}{t}$ in r minus the coefficient of same term but in $([\sqrt{r}]_\infty)^2$ where $[\sqrt{r}]_\infty$ was found above in Eq (10). Hence

$$([\sqrt{r}]_\infty)^2 = \frac{1}{4}$$

This shows that the coefficient of $\frac{1}{t}$ in the above is 0. Now we need to find the coefficient of $\frac{1}{t}$ in r . How this is done depends on if $v = 0$ or not. Since $v = 0$ then starting from $r = \frac{s}{t}$ and doing long division in the form

$$r = Q + \frac{R}{t}$$

Where Q is the quotient and R is the remainder. Then the coefficient of $\frac{1}{t}$ in r will be the coefficient in R of the term in t of degree of t minus one, divided by the leading coefficient in t . Doing long division gives

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{t^2 - 4t + 6}{4t^2 - 8t + 4} \\ &= Q + \frac{R}{4t^2 - 8t + 4} \\ &= \left(\frac{1}{4}\right) + \left(\frac{-2t + 5}{4t^2 - 8t + 4}\right) \\ &= \frac{1}{4} + \frac{-2t + 5}{4t^2 - 8t + 4} \end{aligned}$$

Since the degree of t is 2, then we see that the coefficient of the term t in the remainder R is -2 . Dividing this by leading coefficient in t which is 4 gives $-\frac{1}{2}$. Now b can be found.

$$\begin{aligned} b &= \left(-\frac{1}{2}\right) - (0) \\ &= -\frac{1}{2} \end{aligned}$$

Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= \frac{1}{2} \\ \alpha_\infty^+ &= \frac{1}{2} \left(\frac{b}{a} - v\right) = \frac{1}{2} \left(\frac{-\frac{1}{2}}{\frac{1}{2}} - 0\right) = -\frac{1}{2} \\ \alpha_\infty^- &= \frac{1}{2} \left(-\frac{b}{a} - v\right) = \frac{1}{2} \left(-\frac{-\frac{1}{2}}{\frac{1}{2}} - 0\right) = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = \frac{t^2 - 4t + 6}{4(-1 + t)^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
1	2	0	$\frac{3}{2}$	$-\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
0	$\frac{1}{2}$	$-\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_\infty^+ = -\frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_\infty^+ - (\alpha_{c_1}^-) \\ &= -\frac{1}{2} - \left(-\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c) [\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty) [\sqrt{r}]_\infty$$

Substituting the above values in the above results in

$$\begin{aligned} \omega &= \left((-) [\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^-}{t - c_1} \right) + (+) [\sqrt{r}]_\infty \\ &= -\frac{1}{2(-1 + t)} + \left(\frac{1}{2} \right) \\ &= -\frac{1}{2(-1 + t)} + \frac{1}{2} \\ &= \frac{t - 2}{2t - 2} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(-\frac{1}{2(-1+t)} + \frac{1}{2}\right)(0) + \left(\left(\frac{1}{2(-1+t)}\right)^2 + \left(-\frac{1}{2(-1+t)} + \frac{1}{2}\right)^2 - \left(\frac{t^2 - 4t + 6}{4(-1+t)^2}\right)\right) = 0$$

$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \left(-\frac{1}{2(-1+t)} + \frac{1}{2}\right) dt} \\ &= \frac{e^{\frac{t}{2}}}{\sqrt{-1+t}} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{t}{1-t} dt} \\ &= z_1 e^{\frac{t}{2} + \frac{\ln(-1+t)}{2}} \\ &= z_1 \left(\sqrt{-1+t} e^{\frac{t}{2}}\right) \end{aligned}$$

Which simplifies to

$$y_1 = e^t$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{t}{1-t} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{t+\ln(-1+t)}}{(y_1)^2} dt \\&= y_1(-t e^{-t})\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1(e^t) + c_2(e^t(-t e^{-t}))\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$(1-t)y'' + ty' - y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^t - c_2 t$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned}y_1 &= e^t \\y_2 &= -t\end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} e^t & -t \\ \frac{d}{dt}(e^t) & \frac{d}{dt}(-t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} e^t & -t \\ e^t & -1 \end{vmatrix}$$

Therefore

$$W = (e^t)(-1) - (-t)(e^t)$$

Which simplifies to

$$W = t e^t - e^t$$

Which simplifies to

$$W = e^t(-1 + t)$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{-2t(-1+t)^2 e^{-t}}{(1-t)e^t(-1+t)} dt$$

Which simplifies to

$$u_1 = - \int 2t e^{-2t} dt$$

Hence

$$u_1 = \frac{(1 + 2t)e^{-2t}}{2}$$

And Eq. (3) becomes

$$u_2 = \int \frac{2e^t(-1+t)^2 e^{-t}}{(1-t)e^t(-1+t)} dt$$

Which simplifies to

$$u_2 = \int -2e^{-t} dt$$

Hence

$$u_2 = 2e^{-t}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -2te^{-t} + \frac{(1 + 2t)e^{-2t}e^t}{2}$$

Which simplifies to

$$y_p(t) = e^{-t} \left(-t + \frac{1}{2} \right)$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1e^t - c_2t) + \left(e^{-t} \left(-t + \frac{1}{2} \right) \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1e^t - c_2t + e^{-t} \left(-t + \frac{1}{2} \right) \tag{1}$$

Verification of solutions

$$y = c_1e^t - c_2t + e^{-t} \left(-t + \frac{1}{2} \right)$$

Verified OK.

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    checking if the LODE is of Euler type
    trying a symmetry of the form [xi=0, eta=F(x)]
    checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Reducible group (found an exponential solution)
    Reducible group (found another exponential solution)
<- Kovacics algorithm successful
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 25

```
dsolve((1-t)*diff(y(t),t$2)+t*diff(y(t),t)-y(t) = 2*(t-1)^2*exp(-t),y(t), singsol=all)
```

$$y(t) = c_2 t + e^t c_1 - t e^{-t} + \frac{e^{-t}}{2}$$

✓ Solution by Mathematica

Time used: 0.07 (sec). Leaf size: 30

```
DSolve[(1-t)*y'[t]+t*y'[t]-y[t] == 2*(t-1)^2*Exp[-t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow e^{-t} \left(\frac{1}{2} - t \right) + c_1 e^t - c_2 t$$

10.17 problem 17

10.17.1 Solving as second order euler ode	2893
10.17.2 Solving as second order change of variable on x method 2	2897
10.17.3 Solving as second order change of variable on x method 1	2902
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10.17.5 Solving using Kovacic algorithm	2912

Internal problem ID [699]

Internal file name [OUTPUT/699_Sunday_June_05_2022_01_47_14_AM_49289387/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters. page 190

Problem number: 17.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$x^2y'' - 3y'x + 4y = \ln(x)x^2$$

10.17.1 Solving as second order euler ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = f(x)$$

Where $A = x^2$, $B = -3x$, $C = 4$, $f(x) = \ln(x)x^2$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$.

Solving for y_h from

$$x^2 y'' - 3y'x + 4y = 0$$

This is Euler second order ODE. Let the solution be $y = x^r$, then $y' = rx^{r-1}$ and $y'' = r(r-1)x^{r-2}$. Substituting these back into the given ODE gives

$$x^2(r(r-1))x^{r-2} - 3rxr^{r-1} + 4x^r = 0$$

Simplifying gives

$$r(r-1)x^r - 3rx^r + 4x^r = 0$$

Since $x^r \neq 0$ then dividing throughout by x^r gives

$$r(r-1) - 3r + 4 = 0$$

Or

$$r^2 - 4r + 4 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = 2$$

$$r_2 = 2$$

Since the roots are equal, then the general solution is

$$y = c_1 y_1 + c_2 y_2$$

Where $y_1 = x^r$ and $y_2 = x^r \ln(x)$. Hence

$$y = c_1 x^2 + c_2 x^2 \ln(x)$$

Next, we find the particular solution to the ODE

$$x^2 y'' - 3y'x + 4y = \ln(x) x^2$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = x^2$$

$$y_2 = \ln(x) x^2$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \quad (3)$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} x^2 & \ln(x) x^2 \\ \frac{d}{dx}(x^2) & \frac{d}{dx}(\ln(x) x^2) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} x^2 & \ln(x) x^2 \\ 2x & x + 2x \ln(x) \end{vmatrix}$$

Therefore

$$W = (x^2)(x + 2x \ln(x)) - (\ln(x) x^2)(2x)$$

Which simplifies to

$$W = x^3$$

Which simplifies to

$$W = x^3$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\ln(x)^2 x^4}{x^5} dx$$

Which simplifies to

$$u_1 = - \int \frac{\ln(x)^2}{x} dx$$

Hence

$$u_1 = -\frac{\ln(x)^3}{3}$$

And Eq. (3) becomes

$$u_2 = \int \frac{x^4 \ln(x)}{x^5} dx$$

Which simplifies to

$$u_2 = \int \frac{\ln(x)}{x} dx$$

Hence

$$u_2 = \frac{\ln(x)^2}{2}$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = \frac{\ln(x)^3 x^2}{6}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= x^2 \left(\frac{\ln(x)^3}{6} + c_1 + c_2 \ln(x) \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = x^2 \left(\frac{\ln(x)^3}{6} + c_1 + c_2 \ln(x) \right) \quad (1)$$

Verification of solutions

$$y = x^2 \left(\frac{\ln(x)^3}{6} + c_1 + c_2 \ln(x) \right)$$

Verified OK.

10.17.2 Solving as second order change of variable on x method 2 ode

This is second order non-homogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$. y_h is the solution to

$$x^2y'' - 3y'x + 4y = 0$$

In normal form the ode

$$x^2y'' - 3y'x + 4y = 0 \tag{1}$$

Becomes

$$y'' + p(x)y' + q(x)y = 0 \tag{2}$$

Where

$$p(x) = -\frac{3}{x}$$
$$q(x) = \frac{4}{x^2}$$

Applying change of variables $\tau = g(x)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \tag{3}$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(x) + p(x)\tau'(x)}{\tau'(x)^2} \tag{4}$$

$$q_1(\tau) = \frac{q(x)}{\tau'(x)^2} \tag{5}$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(x) + p(x)\tau'(x) = 0$$

This ode is solved resulting in

$$\begin{aligned}
 \tau &= \int e^{-(\int p(x)dx)} dx \\
 &= \int e^{-(\int -\frac{3}{x} dx)} dx \\
 &= \int e^{3\ln(x)} dx \\
 &= \int x^3 dx \\
 &= \frac{x^4}{4}
 \end{aligned} \tag{6}$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned}
 q_1(\tau) &= \frac{q(x)}{\tau'(x)^2} \\
 &= \frac{\frac{4}{x^2}}{x^6} \\
 &= \frac{4}{x^8}
 \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned}
 \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\
 \frac{d^2}{d\tau^2}y(\tau) + \frac{4y(\tau)}{x^8} &= 0
 \end{aligned}$$

But in terms of τ

$$\frac{4}{x^8} = \frac{1}{4\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{y(\tau)}{4\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$4 \left(\frac{d^2}{d\tau^2}y(\tau) \right) \tau^2 + y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$4\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + \tau^r = 0$$

Simplifying gives

$$4r(r - 1)\tau^r + 0\tau^r + \tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$4r(r - 1) + 0 + 1 = 0$$

Or

$$4r^2 - 4r + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{2}$$
$$r_2 = \frac{1}{2}$$

Since the roots are equal, then the general solution is

$$y(\tau) = c_1y_1 + c_2y_2$$

Where $y_1 = \tau^r$ and $y_2 = \tau^r \ln(\tau)$. Hence

$$y(\tau) = c_1\sqrt{\tau} + c_2\sqrt{\tau} \ln(\tau)$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{(c_2 \ln(x^4) - 2c_2 \ln(2) + c_1) \sqrt{x^4}}{2}$$

Therefore the homogeneous solution y_h is

$$y_h = \frac{(c_2 \ln(x^4) - 2c_2 \ln(2) + c_1) \sqrt{x^4}}{2}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1y_1 + u_2y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \sqrt{x^4}$$

$$y_2 = -\ln(2) \sqrt{x^4} + \frac{\ln(x^4) \sqrt{x^4}}{2}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \quad (3)$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \sqrt{x^4} & -\ln(2) \sqrt{x^4} + \frac{\ln(x^4) \sqrt{x^4}}{2} \\ \frac{d}{dx}(\sqrt{x^4}) & \frac{d}{dx} \left(-\ln(2) \sqrt{x^4} + \frac{\ln(x^4) \sqrt{x^4}}{2} \right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \sqrt{x^4} & -\ln(2) \sqrt{x^4} + \frac{\ln(x^4) \sqrt{x^4}}{2} \\ \frac{2x^3}{\sqrt{x^4}} & -\frac{2\ln(2)x^3}{\sqrt{x^4}} + \frac{2\sqrt{x^4}}{x} + \frac{\ln(x^4)x^3}{\sqrt{x^4}} \end{vmatrix}$$

Therefore

$$W = \left(\sqrt{x^4} \right) \left(-\frac{2\ln(2)x^3}{\sqrt{x^4}} + \frac{2\sqrt{x^4}}{x} + \frac{\ln(x^4)x^3}{\sqrt{x^4}} \right) - \left(-\ln(2) \sqrt{x^4} + \frac{\ln(x^4) \sqrt{x^4}}{2} \right) \left(\frac{2x^3}{\sqrt{x^4}} \right)$$

Which simplifies to

$$W = 2x^3$$

Which simplifies to

$$W = 2x^3$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\left(-\ln(2) \sqrt{x^4} + \frac{\ln(x^4)\sqrt{x^4}}{2}\right) \ln(x) x^2}{2x^5} dx$$

Which simplifies for $0 < x$ to

$$u_1 = - \int - \frac{(\ln(2) - 2 \ln(x)) \ln(x)}{2x} dx$$

Hence

$$u_1 = - \frac{\ln(x)^3}{3} + \frac{\ln(2) \ln(x)^2}{4}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\sqrt{x^4} \ln(x) x^2}{2x^5} dx$$

Which simplifies for $0 < x$ to

$$u_2 = \int \frac{\ln(x)}{2x} dx$$

Hence

$$u_2 = \frac{\ln(x)^2}{4}$$

Which simplifies to

$$u_1 = - \frac{\ln(x)^2 (4 \ln(x) - 3 \ln(2))}{12}$$

$$u_2 = \frac{\ln(x)^2}{4}$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = - \frac{\ln(x)^2 (4 \ln(x) - 3 \ln(2)) \sqrt{x^4}}{12} + \frac{\ln(x)^2 \left(-\ln(2) \sqrt{x^4} + \frac{\ln(x^4)\sqrt{x^4}}{2}\right)}{4}$$

Which simplifies to

$$y_p(x) = \frac{\ln(x)^3 x^2}{6}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{(c_2 \ln(x^4) - 2c_2 \ln(2) + c_1) \sqrt{x^4}}{2} \right) + \left(\frac{\ln(x)^3 x^2}{6} \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{(c_2 \ln(x^4) - 2c_2 \ln(2) + c_1) \sqrt{x^4}}{2} + \frac{\ln(x)^3 x^2}{6} \quad (1)$$

Verification of solutions

$$y = \frac{(c_2 \ln(x^4) - 2c_2 \ln(2) + c_1) \sqrt{x^4}}{2} + \frac{\ln(x)^3 x^2}{6}$$

Verified OK. $\{0 < x\}$

10.17.3 Solving as second order change of variable on x method 1 ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = f(x)$$

Where $A = x^2$, $B = -3x$, $C = 4$, $f(x) = \ln(x) x^2$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$. Solving for y_h from

$$x^2 y'' - 3y'x + 4y = 0$$

In normal form the ode

$$x^2 y'' - 3y'x + 4y = 0 \quad (1)$$

Becomes

$$y'' + p(x) y' + q(x) y = 0 \quad (2)$$

Where

$$p(x) = -\frac{3}{x}$$

$$q(x) = \frac{4}{x^2}$$

Applying change of variables $\tau = g(x)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(x) + p(x)\tau'(x)}{\tau'(x)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(x)}{\tau'(x)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c}\sqrt{q}$$

$$= \frac{2\sqrt{\frac{1}{x^2}}}{c} \quad (6)$$

$$\tau'' = -\frac{2}{c\sqrt{\frac{1}{x^2}}x^3}$$

Substituting the above into (4) results in

$$p_1(\tau) = \frac{\tau''(x) + p(x)\tau'(x)}{\tau'(x)^2}$$

$$= \frac{-\frac{2}{c\sqrt{\frac{1}{x^2}}x^3} - \frac{3}{x}\frac{2\sqrt{\frac{1}{x^2}}}{c}}{\left(\frac{2\sqrt{\frac{1}{x^2}}}{c}\right)^2}$$

$$= -2c$$

Therefore ode (3) now becomes

$$y(\tau)'' + p_1y(\tau)' + q_1y(\tau) = 0$$

$$\frac{d^2}{d\tau^2}y(\tau) - 2c\left(\frac{d}{d\tau}y(\tau)\right) + c^2y(\tau) = 0 \quad (7)$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{c\tau} c_1$$

Now from (6)

$$\begin{aligned} \tau &= \int \frac{1}{c} \sqrt{q} dx \\ &= \frac{\int 2\sqrt{\frac{1}{x^2}} dx}{c} \\ &= \frac{2\sqrt{\frac{1}{x^2}} x \ln(x)}{c} \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = c_1 x^2$$

Now the particular solution to this ODE is found

$$x^2 y'' - 3y'x + 4y = \ln(x) x^2$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned} y_1 &= \sqrt{x^4} \\ y_2 &= -\ln(2) \sqrt{x^4} + \frac{\ln(x^4) \sqrt{x^4}}{2} \end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \tag{3}$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \sqrt{x^4} & -\ln(2)\sqrt{x^4} + \frac{\ln(x^4)\sqrt{x^4}}{2} \\ \frac{d}{dx}(\sqrt{x^4}) & \frac{d}{dx}\left(-\ln(2)\sqrt{x^4} + \frac{\ln(x^4)\sqrt{x^4}}{2}\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \sqrt{x^4} & -\ln(2)\sqrt{x^4} + \frac{\ln(x^4)\sqrt{x^4}}{2} \\ \frac{2x^3}{\sqrt{x^4}} & -\frac{2\ln(2)x^3}{\sqrt{x^4}} + \frac{2\sqrt{x^4}}{x} + \frac{\ln(x^4)x^3}{\sqrt{x^4}} \end{vmatrix}$$

Therefore

$$W = (\sqrt{x^4}) \left(-\frac{2\ln(2)x^3}{\sqrt{x^4}} + \frac{2\sqrt{x^4}}{x} + \frac{\ln(x^4)x^3}{\sqrt{x^4}} \right) - \left(-\ln(2)\sqrt{x^4} + \frac{\ln(x^4)\sqrt{x^4}}{2} \right) \left(\frac{2x^3}{\sqrt{x^4}} \right)$$

Which simplifies to

$$W = 2x^3$$

Which simplifies to

$$W = 2x^3$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\left(-\ln(2)\sqrt{x^4} + \frac{\ln(x^4)\sqrt{x^4}}{2} \right) \ln(x) x^2}{2x^5} dx$$

Which simplifies for $0 < x$ to

$$u_1 = - \int -\frac{(\ln(2) - 2\ln(x)) \ln(x)}{2x} dx$$

Hence

$$u_1 = -\frac{\ln(x)^3}{3} + \frac{\ln(2)\ln(x)^2}{4}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\sqrt{x^4} \ln(x) x^2}{2x^5} dx$$

Which simplifies for $0 < x$ to

$$u_2 = \int \frac{\ln(x)}{2x} dx$$

Hence

$$u_2 = \frac{\ln(x)^2}{4}$$

Which simplifies to

$$u_1 = -\frac{\ln(x)^2 (4 \ln(x) - 3 \ln(2))}{12}$$
$$u_2 = \frac{\ln(x)^2}{4}$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = -\frac{\ln(x)^2 (4 \ln(x) - 3 \ln(2)) \sqrt{x^4}}{12} + \frac{\ln(x)^2 \left(-\ln(2) \sqrt{x^4} + \frac{\ln(x^4) \sqrt{x^4}}{2} \right)}{4}$$

Which simplifies to

$$y_p(x) = \frac{\ln(x)^3 x^2}{6}$$

Therefore the general solution is

$$y = y_h + y_p$$
$$= (c_1 x^2) + \left(\frac{\ln(x)^3 x^2}{6} \right)$$
$$= \frac{\ln(x)^3 x^2}{6} + c_1 x^2$$

Which simplifies to

$$y = x^2 \left(\frac{\ln(x)^3}{6} + c_1 \right)$$

Summary

The solution(s) found are the following

$$y = x^2 \left(\frac{\ln(x)^3}{6} + c_1 \right) \quad (1)$$

Verification of solutions

$$y = x^2 \left(\frac{\ln(x)^3}{6} + c_1 \right)$$

Verified OK. $\{0 < x\}$

10.17.4 Solving as second order change of variable on y method 2 ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = f(x)$$

Where $A = x^2, B = -3x, C = 4, f(x) = \ln(x) x^2$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$. Solving for y_h from

$$x^2 y'' - 3y'x + 4y = 0$$

In normal form the ode

$$x^2 y'' - 3y'x + 4y = 0 \quad (1)$$

Becomes

$$y'' + p(x) y' + q(x) y = 0 \quad (2)$$

Where

$$p(x) = -\frac{3}{x}$$
$$q(x) = \frac{4}{x^2}$$

Applying change of variables on the dependent variable $y = v(x)x^n$ to (2) gives the following ode where the dependent variables is $v(x)$ and not y .

$$v''(x) + \left(\frac{2n}{x} + p\right)v'(x) + \left(\frac{n(n-1)}{x^2} + \frac{np}{x} + q\right)v(x) = 0 \quad (3)$$

Let the coefficient of $v(x)$ above be zero. Hence

$$\frac{n(n-1)}{x^2} + \frac{np}{x} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(x)$ and $q(x)$ into (4) gives

$$\frac{n(n-1)}{x^2} - \frac{3n}{x^2} + \frac{4}{x^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = 2 \quad (6)$$

Substituting this value in (3) gives

$$\begin{aligned} v''(x) + \frac{v'(x)}{x} &= 0 \\ v''(x) + \frac{v'(x)}{x} &= 0 \end{aligned} \quad (7)$$

Using the substitution

$$u(x) = v'(x)$$

Then (7) becomes

$$u'(x) + \frac{u(x)}{x} = 0 \quad (8)$$

The above is now solved for $u(x)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{u}{x} \end{aligned}$$

Where $f(x) = -\frac{1}{x}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= -\frac{1}{x} dx \\ \int \frac{1}{u} du &= \int -\frac{1}{x} dx \\ \ln(u) &= -\ln(x) + c_1 \\ u &= e^{-\ln(x)+c_1} \\ &= \frac{c_1}{x}\end{aligned}$$

Now that $u(x)$ is known, then

$$\begin{aligned}v'(x) &= u(x) \\ v(x) &= \int u(x) dx + c_2 \\ &= c_1 \ln(x) + c_2\end{aligned}$$

Hence

$$\begin{aligned}y &= v(x) x^n \\ &= (c_1 \ln(x) + c_2) x^2 \\ &= (c_1 \ln(x) + c_2) x^2\end{aligned}$$

Now the particular solution to this ODE is found

$$x^2 y'' - 3y'x + 4y = \ln(x) x^2$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned}y_1 &= x^2 \\ y_2 &= \ln(x) x^2\end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \quad (3)$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} x^2 & \ln(x) x^2 \\ \frac{d}{dx}(x^2) & \frac{d}{dx}(\ln(x) x^2) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} x^2 & \ln(x) x^2 \\ 2x & x + 2x \ln(x) \end{vmatrix}$$

Therefore

$$W = (x^2)(x + 2x \ln(x)) - (\ln(x) x^2)(2x)$$

Which simplifies to

$$W = x^3$$

Which simplifies to

$$W = x^3$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\ln(x)^2 x^4}{x^5} dx$$

Which simplifies to

$$u_1 = - \int \frac{\ln(x)^2}{x} dx$$

Hence

$$u_1 = - \frac{\ln(x)^3}{3}$$

And Eq. (3) becomes

$$u_2 = \int \frac{x^4 \ln(x)}{x^5} dx$$

Which simplifies to

$$u_2 = \int \frac{\ln(x)}{x} dx$$

Hence

$$u_2 = \frac{\ln(x)^2}{2}$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = \frac{\ln(x)^3 x^2}{6}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= ((c_1 \ln(x) + c_2) x^2) + \left(\frac{\ln(x)^3 x^2}{6} \right) \\ &= \frac{\ln(x)^3 x^2}{6} + (c_1 \ln(x) + c_2) x^2 \end{aligned}$$

Which simplifies to

$$y = x^2 \left(\frac{\ln(x)^3}{6} + c_1 \ln(x) + c_2 \right)$$

Summary

The solution(s) found are the following

$$y = x^2 \left(\frac{\ln(x)^3}{6} + c_1 \ln(x) + c_2 \right) \quad (1)$$

Verification of solutions

$$y = x^2 \left(\frac{\ln(x)^3}{6} + c_1 \ln(x) + c_2 \right)$$

Verified OK. $\{0 < x\}$

10.17.5 Solving using Kovacic algorithm

Writing the ode as

$$x^2 y'' - 3y'x + 4y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= x^2 \\ B &= -3x \\ C &= 4 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4x^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 4x^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \left(-\frac{1}{4x^2}\right) z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 502: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4x^2$. There is a pole at $x = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{4x^2}$$

For the pole at $x = 0$ let b be the coefficient of $\frac{1}{x^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{x^2}$ in the Laurent series expansion of r at ∞ , which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{1}{4x^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{1}{4x^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2}$	$\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = \frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= \frac{1}{2} - \left(\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{x - c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{x - c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= \frac{1}{2x} + (-)(0) \\ &= \frac{1}{2x} \\ &= \frac{1}{2x} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(x)$ of degree $d = 0$ to solve the ode. The polynomial $p(x)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r)p = 0 \quad (1A)$$

Let

$$p(x) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2\left(\frac{1}{2x}\right)(0) + \left(\left(-\frac{1}{2x^2}\right) + \left(\frac{1}{2x}\right)^2 - \left(-\frac{1}{4x^2}\right) \right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(x) &= pe^{\int \omega dx} \\ &= e^{\int \frac{1}{2x} dx} \\ &= \sqrt{x} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\&= z_1 e^{-\int \frac{1}{2} \frac{-3x}{x^2} dx} \\&= z_1 e^{\frac{3 \ln(x)}{2}} \\&= z_1 \left(x^{\frac{3}{2}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = x^2$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-3x}{x^2} dx}}{(y_1)^2} dx \\&= y_1 \int \frac{e^{3 \ln(x)}}{(y_1)^2} dx \\&= y_1 (\ln(x))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 (x^2) + c_2 (x^2 (\ln(x)))\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$.
 y_h is the solution to

$$x^2 y'' - 3y'x + 4y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 x^2 + c_2 x^2 \ln(x)$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = x^2$$

$$y_2 = \ln(x) x^2$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \quad (3)$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} x^2 & \ln(x) x^2 \\ \frac{d}{dx}(x^2) & \frac{d}{dx}(\ln(x) x^2) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} x^2 & \ln(x) x^2 \\ 2x & x + 2x \ln(x) \end{vmatrix}$$

Therefore

$$W = (x^2)(x + 2x \ln(x)) - (\ln(x) x^2)(2x)$$

Which simplifies to

$$W = x^3$$

Which simplifies to

$$W = x^3$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\ln(x)^2 x^4}{x^5} dx$$

Which simplifies to

$$u_1 = - \int \frac{\ln(x)^2}{x} dx$$

Hence

$$u_1 = - \frac{\ln(x)^3}{3}$$

And Eq. (3) becomes

$$u_2 = \int \frac{x^4 \ln(x)}{x^5} dx$$

Which simplifies to

$$u_2 = \int \frac{\ln(x)}{x} dx$$

Hence

$$u_2 = \frac{\ln(x)^2}{2}$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = \frac{\ln(x)^3 x^2}{6}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 x^2 + c_2 x^2 \ln(x)) + \left(\frac{\ln(x)^3 x^2}{6} \right) \end{aligned}$$

Which simplifies to

$$y = x^2(c_1 + c_2 \ln(x)) + \frac{\ln(x)^3 x^2}{6}$$

Summary

The solution(s) found are the following

$$y = x^2(c_1 + c_2 \ln(x)) + \frac{\ln(x)^3 x^2}{6} \quad (1)$$

Verification of solutions

$$y = x^2(c_1 + c_2 \ln(x)) + \frac{\ln(x)^3 x^2}{6}$$

Verified OK. $\{0 < x\}$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]  
trying a double symmetry of the form [xi=0, eta=F(x)]  
-> Try solving first the homogeneous part of the ODE  
    checking if the LODE has constant coefficients  
    checking if the LODE is of Euler type  
    <- LODE of Euler type successful  
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

```
dsolve(x^2*diff(y(x),x$2)-3*x*diff(y(x),x)+4*y(x) = x^2*ln(x),y(x), singsol=all)
```

$$y(x) = x^2 \left(c_2 + \ln(x) c_1 + \frac{\ln(x)^3}{6} \right)$$

✓ Solution by Mathematica

Time used: 0.027 (sec). Leaf size: 27

```
DSolve[x^2*y''[x]-3*x*y'[x]+4*y[x] == x^2*Log[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{6}x^2(\log^3(x) + 12c_2 \log(x) + 6c_1)$$

10.18 problem 20

- 10.18.1 Solving as second order change of variable on y method 1 ode . 2921
- 10.18.2 Solving as second order bessel ode ode 2930
- 10.18.3 Solving using Kovacic algorithm 2933

Internal problem ID [700]

Internal file name [OUTPUT/700_Sunday_June_05_2022_01_47_15_AM_75865028/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 20.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_bessel_ode",
"second_order_change_of_variable_on_y_method_1"

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$x^2y'' + y'x + \left(x^2 - \frac{1}{4}\right)y = g(x)$$

10.18.1 Solving as second order change of variable on y method 1 ode

This is second order non-homogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$.
 y_h is the solution to

$$x^2y'' + y'x + \left(x^2 - \frac{1}{4}\right)y = 0$$

In normal form the given ode is written as

$$y'' + p(x)y' + q(x)y = 0 \tag{2}$$

Where

$$p(x) = \frac{1}{x}$$

$$q(x) = \frac{x^2 - \frac{1}{4}}{x^2}$$

Calculating the Liouville ode invariant Q given by

$$\begin{aligned} Q &= q - \frac{p'}{2} - \frac{p^2}{4} \\ &= \frac{x^2 - \frac{1}{4}}{x^2} - \frac{\left(\frac{1}{x}\right)'}{2} - \frac{\left(\frac{1}{x}\right)^2}{4} \\ &= \frac{x^2 - \frac{1}{4}}{x^2} - \frac{\left(-\frac{1}{x^2}\right)}{2} - \frac{\left(\frac{1}{x^2}\right)}{4} \\ &= \frac{x^2 - \frac{1}{4}}{x^2} - \left(-\frac{1}{2x^2}\right) - \frac{1}{4x^2} \\ &= 1 \end{aligned}$$

Since the Liouville ode invariant does not depend on the independent variable x then the transformation

$$y = v(x) z(x) \tag{3}$$

is used to change the original ode to a constant coefficients ode in v . In (3) the term $z(x)$ is given by

$$\begin{aligned} z(x) &= e^{-\left(\int \frac{p(x)}{2} dx\right)} \\ &= e^{-\int \frac{1}{2x} dx} \\ &= \frac{1}{\sqrt{x}} \end{aligned} \tag{5}$$

Hence (3) becomes

$$y = \frac{v(x)}{\sqrt{x}} \tag{4}$$

Applying this change of variable to the original ode results in

$$x^{\frac{3}{2}}(v''(x) + v(x)) = g(x)$$

Which is now solved for $v(x)$ Simplifying the ode gives

$$v''(x) + v(x) = \frac{g(x)}{x^{\frac{3}{2}}}$$

This is second order non-homogeneous ODE. In standard form the ODE is

$$Av''(x) + Bv'(x) + Cv(x) = f(x)$$

Where $A = 1, B = 0, C = 1, f(x) = \frac{g(x)}{x^{\frac{3}{2}}}$. Let the solution be

$$v(x) = v_h + v_p$$

Where v_h is the solution to the homogeneous ODE $Av''(x) + Bv'(x) + Cv(x) = 0$, and v_p is a particular solution to the non-homogeneous ODE $Av''(x) + Bv'(x) + Cv(x) = f(x)$. v_h is the solution to

$$v''(x) + v(x) = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Av''(x) + Bv'(x) + Cv(x) = 0$$

Where in the above $A = 1, B = 0, C = 1$. Let the solution be $v(x) = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda x} + e^{\lambda x} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(1)} \\ &= \pm i \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= +i \\ \lambda_2 &= -i \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= i \\ \lambda_2 &= -i \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$v(x) = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$v(x) = e^0(c_1 \cos(x) + c_2 \sin(x))$$

Or

$$v(x) = c_1 \cos(x) + c_2 \sin(x)$$

Therefore the homogeneous solution v_h is

$$v_h = c_1 \cos(x) + c_2 \sin(x)$$

The particular solution v_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$v_p(x) = u_1 v_1 + u_2 v_2 \tag{1}$$

Where u_1, u_2 to be determined, and v_1, v_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$v_1 = \cos(x)$$

$$v_2 = \sin(x)$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{v_2 f(x)}{aW(x)} \tag{2}$$

$$u_2 = \int \frac{v_1 f(x)}{aW(x)} \tag{3}$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of v'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} v_1 & v_2 \\ v_1' & v_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \cos(x) & \sin(x) \\ \frac{d}{dx}(\cos(x)) & \frac{d}{dx}(\sin(x)) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \cos(x) & \sin(x) \\ -\sin(x) & \cos(x) \end{vmatrix}$$

Therefore

$$W = (\cos(x))(\cos(x)) - (\sin(x))(-\sin(x))$$

Which simplifies to

$$W = \cos(x)^2 + \sin(x)^2$$

Which simplifies to

$$W = 1$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{\sin(x)g(x)}{x^{\frac{3}{2}}}}{1} dx$$

Which simplifies to

$$u_1 = - \int \frac{\sin(x)g(x)}{x^{\frac{3}{2}}} dx$$

Hence

$$u_1 = - \left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right)$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{\cos(x)g(x)}{x^{\frac{3}{2}}}}{1} dx$$

Which simplifies to

$$u_2 = \int \frac{\cos(x) g(x)}{x^{\frac{3}{2}}} dx$$

Hence

$$u_2 = \int_0^x \frac{\cos(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha$$

Therefore the particular solution, from equation (1) is

$$v_p(x) = - \left(\int_0^x \frac{\sin(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \sin(x)$$

Therefore the general solution is

$$\begin{aligned} v &= v_h + v_p \\ &= (c_1 \cos(x) + c_2 \sin(x)) \\ &\quad + \left(- \left(\int_0^x \frac{\sin(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \sin(x) \right) \end{aligned}$$

Now that $v(x)$ is known, then

$$\begin{aligned} y &= v(x) z(x) \\ &= \left(c_1 \cos(x) + c_2 \sin(x) - \left(\int_0^x \frac{\sin(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \sin(x) \right) (z(x)) \end{aligned} \tag{7}$$

But from (5)

$$z(x) = \frac{1}{\sqrt{x}}$$

Hence (7) becomes

$$y = \frac{c_1 \cos(x) + c_2 \sin(x) - \left(\int_0^x \frac{\sin(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \sin(x)}{\sqrt{x}}$$

Therefore the homogeneous solution y_h is

$$y_h = \frac{c_1 \cos(x) + c_2 \sin(x) - \left(\int_0^x \frac{\sin(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha) g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \sin(x)}{\sqrt{x}}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \frac{\cos(x)}{\sqrt{x}}$$

$$y_2 = \frac{\sin(x)}{\sqrt{x}}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \quad (3)$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{\cos(x)}{\sqrt{x}} & \frac{\sin(x)}{\sqrt{x}} \\ \frac{d}{dx} \left(\frac{\cos(x)}{\sqrt{x}} \right) & \frac{d}{dx} \left(\frac{\sin(x)}{\sqrt{x}} \right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{\cos(x)}{\sqrt{x}} & \frac{\sin(x)}{\sqrt{x}} \\ -\frac{\sin(x)}{\sqrt{x}} - \frac{\cos(x)}{2x^{\frac{3}{2}}} & \frac{\cos(x)}{\sqrt{x}} - \frac{\sin(x)}{2x^{\frac{3}{2}}} \end{vmatrix}$$

Therefore

$$W = \left(\frac{\cos(x)}{\sqrt{x}} \right) \left(\frac{\cos(x)}{\sqrt{x}} - \frac{\sin(x)}{2x^{\frac{3}{2}}} \right) - \left(\frac{\sin(x)}{\sqrt{x}} \right) \left(-\frac{\sin(x)}{\sqrt{x}} - \frac{\cos(x)}{2x^{\frac{3}{2}}} \right)$$

Which simplifies to

$$W = \frac{\cos(x)^2 + \sin(x)^2}{x}$$

Which simplifies to

$$W = \frac{1}{x}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{\sin(x)g(x)}{\sqrt{x}}}{x} dx$$

Which simplifies to

$$u_1 = - \int \frac{\sin(x)g(x)}{x^{\frac{3}{2}}} dx$$

Hence

$$u_1 = - \left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right)$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{\cos(x)g(x)}{\sqrt{x}}}{x} dx$$

Which simplifies to

$$u_2 = \int \frac{\cos(x)g(x)}{x^{\frac{3}{2}}} dx$$

Hence

$$u_2 = \int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = - \frac{\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \cos(x)}{\sqrt{x}} + \frac{\left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \sin(x)}{\sqrt{x}}$$

Which simplifies to

$$y_p(x) = \frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{c_1 \cos(x) + c_2 \sin(x) - \left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}} \right) \\ &\quad + \left(\frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}} \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$\begin{aligned} y &= \frac{c_1 \cos(x) + c_2 \sin(x) - \left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}} \\ &\quad + \frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}} \end{aligned} \quad (1)$$

Verification of solutions

$$\begin{aligned} y &= \frac{c_1 \cos(x) + c_2 \sin(x) - \left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}} \\ &\quad + \frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}} \end{aligned}$$

Verified OK.

10.18.2 Solving as second order Bessel ODE

Writing the ODE as

$$x^2 y'' + y'x + \left(x^2 - \frac{1}{4}\right) y = g(x) \quad (1)$$

Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE and y_p is a particular solution to the non-homogeneous ODE. Bessel ODE has the form

$$x^2 y'' + y'x + (-n^2 + x^2) y = 0 \quad (2)$$

The generalized form of Bessel ODE is given by Bowman (1958) as the following

$$x^2 y'' + (1 - 2\alpha) xy' + (\beta^2 \gamma^2 x^{2\gamma} - n^2 \gamma^2 + \alpha^2) y = 0 \quad (3)$$

With the standard solution

$$y = x^\alpha (c_1 \text{BesselJ}(n, \beta x^\gamma) + c_2 \text{BesselY}(n, \beta x^\gamma)) \quad (4)$$

Comparing (3) to (1) and solving for α, β, n, γ gives

$$\alpha = 0$$

$$\beta = 1$$

$$n = -\frac{1}{2}$$

$$\gamma = 1$$

Substituting all the above into (4) gives the solution as

$$y = \frac{c_1 \sqrt{2} \cos(x)}{\sqrt{\pi} \sqrt{x}} + \frac{c_2 \sqrt{2} \sin(x)}{\sqrt{\pi} \sqrt{x}}$$

Therefore the homogeneous solution y_h is

$$y_h = \frac{c_1 \sqrt{2} \cos(x)}{\sqrt{\pi} \sqrt{x}} + \frac{c_2 \sqrt{2} \sin(x)}{\sqrt{\pi} \sqrt{x}}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of

parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \frac{\cos(x)}{\sqrt{x}}$$

$$y_2 = \frac{\sin(x)}{\sqrt{x}}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \quad (3)$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{\cos(x)}{\sqrt{x}} & \frac{\sin(x)}{\sqrt{x}} \\ \frac{d}{dx} \left(\frac{\cos(x)}{\sqrt{x}} \right) & \frac{d}{dx} \left(\frac{\sin(x)}{\sqrt{x}} \right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{\cos(x)}{\sqrt{x}} & \frac{\sin(x)}{\sqrt{x}} \\ -\frac{\sin(x)}{\sqrt{x}} - \frac{\cos(x)}{2x^{\frac{3}{2}}} & \frac{\cos(x)}{\sqrt{x}} - \frac{\sin(x)}{2x^{\frac{3}{2}}} \end{vmatrix}$$

Therefore

$$W = \left(\frac{\cos(x)}{\sqrt{x}} \right) \left(\frac{\cos(x)}{\sqrt{x}} - \frac{\sin(x)}{2x^{\frac{3}{2}}} \right) - \left(\frac{\sin(x)}{\sqrt{x}} \right) \left(-\frac{\sin(x)}{\sqrt{x}} - \frac{\cos(x)}{2x^{\frac{3}{2}}} \right)$$

Which simplifies to

$$W = \frac{\cos(x)^2 + \sin(x)^2}{x}$$

Which simplifies to

$$W = \frac{1}{x}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{\sin(x)g(x)}{\sqrt{x}}}{x} dx$$

Which simplifies to

$$u_1 = - \int \frac{\sin(x)g(x)}{x^{\frac{3}{2}}} dx$$

Hence

$$u_1 = - \left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right)$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{\cos(x)g(x)}{\sqrt{x}}}{x} dx$$

Which simplifies to

$$u_2 = \int \frac{\cos(x)g(x)}{x^{\frac{3}{2}}} dx$$

Hence

$$u_2 = \int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = - \frac{\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \cos(x)}{\sqrt{x}} + \frac{\left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \sin(x)}{\sqrt{x}}$$

Which simplifies to

$$y_p(x) = \frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{c_1\sqrt{2} \cos(x)}{\sqrt{\pi} \sqrt{x}} + \frac{c_2\sqrt{2} \sin(x)}{\sqrt{\pi} \sqrt{x}} \right) \\ &\quad + \left(\frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}} \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$\begin{aligned} y &= \frac{c_1\sqrt{2} \cos(x)}{\sqrt{\pi} \sqrt{x}} + \frac{c_2\sqrt{2} \sin(x)}{\sqrt{\pi} \sqrt{x}} \\ &\quad + \frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}} \end{aligned} \quad (1)$$

Verification of solutions

$$y = \frac{c_1\sqrt{2} \cos(x)}{\sqrt{\pi} \sqrt{x}} + \frac{c_2\sqrt{2} \sin(x)}{\sqrt{\pi} \sqrt{x}} + \frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}}$$

Verified OK.

10.18.3 Solving using Kovacic algorithm

Writing the ode as

$$x^2 y'' + y'x + \left(x^2 - \frac{1}{4}\right) y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= x^2 \\ B &= x \\ C &= x^2 - \frac{1}{4} \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -z(x) \tag{7}$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 503: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned}
 y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\
 &= z_1 e^{-\int \frac{1}{2} \frac{x}{x^2} dx} \\
 &= z_1 e^{-\frac{\ln(x)}{2}} \\
 &= z_1 \left(\frac{1}{\sqrt{x}} \right)
 \end{aligned}$$

Which simplifies to

$$y_1 = \frac{\cos(x)}{\sqrt{x}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{x}{x^2} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-\ln(x)}}{(y_1)^2} dx \\ &= y_1(\tan(x)) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(\frac{\cos(x)}{\sqrt{x}} \right) + c_2 \left(\frac{\cos(x)}{\sqrt{x}} (\tan(x)) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$. y_h is the solution to

$$x^2 y'' + y'x + \left(x^2 - \frac{1}{4}\right) y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = \frac{c_1 \cos(x)}{\sqrt{x}} + \frac{c_2 \sin(x)}{\sqrt{x}}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of

parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \frac{\cos(x)}{\sqrt{x}}$$

$$y_2 = \frac{\sin(x)}{\sqrt{x}}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \quad (3)$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{\cos(x)}{\sqrt{x}} & \frac{\sin(x)}{\sqrt{x}} \\ \frac{d}{dx} \left(\frac{\cos(x)}{\sqrt{x}} \right) & \frac{d}{dx} \left(\frac{\sin(x)}{\sqrt{x}} \right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{\cos(x)}{\sqrt{x}} & \frac{\sin(x)}{\sqrt{x}} \\ -\frac{\sin(x)}{\sqrt{x}} - \frac{\cos(x)}{2x^{\frac{3}{2}}} & \frac{\cos(x)}{\sqrt{x}} - \frac{\sin(x)}{2x^{\frac{3}{2}}} \end{vmatrix}$$

Therefore

$$W = \left(\frac{\cos(x)}{\sqrt{x}} \right) \left(\frac{\cos(x)}{\sqrt{x}} - \frac{\sin(x)}{2x^{\frac{3}{2}}} \right) - \left(\frac{\sin(x)}{\sqrt{x}} \right) \left(-\frac{\sin(x)}{\sqrt{x}} - \frac{\cos(x)}{2x^{\frac{3}{2}}} \right)$$

Which simplifies to

$$W = \frac{\cos(x)^2 + \sin(x)^2}{x}$$

Which simplifies to

$$W = \frac{1}{x}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{\sin(x)g(x)}{\sqrt{x}}}{x} dx$$

Which simplifies to

$$u_1 = - \int \frac{\sin(x)g(x)}{x^{\frac{3}{2}}} dx$$

Hence

$$u_1 = - \left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right)$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{\cos(x)g(x)}{\sqrt{x}}}{x} dx$$

Which simplifies to

$$u_2 = \int \frac{\cos(x)g(x)}{x^{\frac{3}{2}}} dx$$

Hence

$$u_2 = \int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = - \frac{\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \cos(x)}{\sqrt{x}} + \frac{\left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha \right) \sin(x)}{\sqrt{x}}$$

Which simplifies to

$$y_p(x) = \frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}}$$

Therefore the general solution is

$$y = y_h + y_p$$
$$= \left(\frac{c_1 \cos(x)}{\sqrt{x}} + \frac{c_2 \sin(x)}{\sqrt{x}}\right) + \left(\frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}}\right)$$

Which simplifies to

$$y = \frac{c_1 \cos(x) + c_2 \sin(x)}{\sqrt{x}} + \frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 \cos(x) + c_2 \sin(x)}{\sqrt{x}} + \frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 \cos(x) + c_2 \sin(x)}{\sqrt{x}} + \frac{-\left(\int_0^x \frac{\sin(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \cos(x) + \left(\int_0^x \frac{\cos(\alpha)g(\alpha)}{\alpha^{\frac{3}{2}}} d\alpha\right) \sin(x)}{\sqrt{x}}$$

Verified OK.

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    checking if the LODE is of Euler type
    trying a symmetry of the form [xi=0, eta=F(x)]
    checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Group is reducible or imprimitive
<- Kovacics algorithm successful
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 44

```
dsolve(x^2*diff(y(x),x$2)+x*diff(y(x),x)+(x^2-25/100)*y(x) = g(x),y(x), singsol=all)
```

$$y(x) = \frac{\sin(x) c_2 + \cos(x) c_1 + \left(\int \frac{\cos(x)g(x)}{x^{\frac{3}{2}}} dx \right) \sin(x) - \left(\int \frac{\sin(x)g(x)}{x^{\frac{3}{2}}} dx \right) \cos(x)}{\sqrt{x}}$$

✓ Solution by Mathematica

Time used: 0.203 (sec). Leaf size: 107

```
DSolve[x^2*y''[x]+x*y'[x]+(x^2-25/100)*y[x] == g[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{e^{-ix} \left(2 \int_1^x \frac{ie^{iK[1]}g(K[1])}{2K[1]^{3/2}} dK[1] - ie^{2ix} \int_1^x \frac{e^{-iK[2]}g(K[2])}{K[2]^{3/2}} dK[2] - ic_2e^{2ix} + 2c_1 \right)}{2\sqrt{x}}$$

10.19 problem 29

10.19.1 Solving as second order euler ode ode	2942
10.19.2 Solving as linear second order ode solved by an integrating factor ode	2945
10.19.3 Solving as second order change of variable on x method 2 ode .	2946
10.19.4 Solving as second order change of variable on x method 1 ode .	2951
10.19.5 Solving as second order change of variable on y method 2 ode .	2955
10.19.6 Solving as second order ode non constant coeff transformation on B ode	2960
10.19.7 Solving using Kovacic algorithm	2964

Internal problem ID [701]

Internal file name [OUTPUT/701_Sunday_June_05_2022_01_47_17_AM_20232498/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 29.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2", "linear_second_order_ode_solved_by_an_integrating_factor", "second_order_ode_non_constant_coeff_transformation_on_B"

Maple gives the following as the ode type

```
[[_2nd_order, _with_linear_symmetries]]
```

$$t^2y'' - 2ty' + 2y = 4t^2$$

10.19.1 Solving as second order euler ode ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = t^2$, $B = -2t$, $C = 2$, $f(t) = 4t^2$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. Solving for y_h from

$$t^2y'' - 2ty' + 2y = 0$$

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} - 2trt^{r-1} + 2t^r = 0$$

Simplifying gives

$$r(r-1)t^r - 2rt^r + 2t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r-1) - 2r + 2 = 0$$

Or

$$r^2 - 3r + 2 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = 1$$

$$r_2 = 2$$

Since the roots are real and distinct, then the general solution is

$$y = c_1y_1 + c_2y_2$$

Where $y_1 = t^{r_1}$ and $y_2 = t^{r_2}$. Hence

$$y = c_2t^2 + c_1t$$

Next, we find the particular solution to the ODE

$$t^2y'' - 2ty' + 2y = 4t^2$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = t$$

$$y_2 = t^2$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t & t^2 \\ \frac{d}{dt}(t) & \frac{d}{dt}(t^2) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t & t^2 \\ 1 & 2t \end{vmatrix}$$

Therefore

$$W = (t)(2t) - (t^2) \quad (1)$$

Which simplifies to

$$W = t^2$$

Which simplifies to

$$W = t^2$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{4t^4}{t^4} dt$$

Which simplifies to

$$u_1 = - \int 4dt$$

Hence

$$u_1 = -4t$$

And Eq. (3) becomes

$$u_2 = \int \frac{4t^3}{t^4} dt$$

Which simplifies to

$$u_2 = \int \frac{4}{t} dt$$

Hence

$$u_2 = 4 \ln(t)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -4t^2 + 4t^2 \ln(t)$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= -4t^2 + 4t^2 \ln(t) + c_2t^2 + c_1t \end{aligned}$$

Summary

The solution(s) found are the following

$$y = -4t^2 + 4t^2 \ln(t) + c_2t^2 + c_1t \tag{1}$$

Verification of solutions

$$y = -4t^2 + 4t^2 \ln(t) + c_2t^2 + c_1t$$

Verified OK.

10.19.2 Solving as linear second order ode solved by an integrating factor ode

The ode satisfies this form

$$y'' + p(t)y' + \frac{(p(t))^2 + p'(t)}{2}y = f(t)$$

Where $p(t) = -\frac{2}{t}$. Therefore, there is an integrating factor given by

$$\begin{aligned}M(x) &= e^{\frac{1}{2} \int p dx} \\ &= e^{\int -\frac{2}{t} dx} \\ &= \frac{1}{t}\end{aligned}$$

Multiplying both sides of the ODE by the integrating factor $M(x)$ makes the left side of the ODE a complete differential

$$\begin{aligned}(M(x)y)'' &= \frac{4}{t} \\ \left(\frac{y}{t}\right)'' &= \frac{4}{t}\end{aligned}$$

Integrating once gives

$$\left(\frac{y}{t}\right)' = 4 \ln(t) + c_1$$

Integrating again gives

$$\left(\frac{y}{t}\right) = t(4 \ln(t) + c_1 - 4) + c_2$$

Hence the solution is

$$y = \frac{t(4 \ln(t) + c_1 - 4) + c_2}{\frac{1}{t}}$$

Or

$$y = c_1 t^2 + 4t^2 \ln(t) + c_2 t - 4t^2$$

Summary

The solution(s) found are the following

$$y = c_1 t^2 + 4t^2 \ln(t) + c_2 t - 4t^2 \quad (1)$$

Verification of solutions

$$y = c_1 t^2 + 4t^2 \ln(t) + c_2 t - 4t^2$$

Verified OK.

10.19.3 Solving as second order change of variable on x method 2 ode

This is second order non-homogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$t^2y'' - 2ty' + 2y = 0$$

In normal form the ode

$$t^2y'' - 2ty' + 2y = 0 \tag{1}$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \tag{2}$$

Where

$$p(t) = -\frac{2}{t}$$
$$q(t) = \frac{2}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \tag{3}$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \tag{4}$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \tag{5}$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t)\tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned}
 \tau &= \int e^{-(\int p(t)dt)} dt \\
 &= \int e^{-(\int -\frac{2}{t} dt)} dt \\
 &= \int e^{2\ln(t)} dt \\
 &= \int t^2 dt \\
 &= \frac{t^3}{3}
 \end{aligned} \tag{6}$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned}
 q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\
 &= \frac{\frac{2}{t^2}}{t^4} \\
 &= \frac{2}{t^6}
 \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned}
 \frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\
 \frac{d^2}{d\tau^2}y(\tau) + \frac{2y(\tau)}{t^6} &= 0
 \end{aligned}$$

But in terms of τ

$$\frac{2}{t^6} = \frac{2}{9\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{2y(\tau)}{9\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$9\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + 2y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$9\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + 2\tau^r = 0$$

Simplifying gives

$$9r(r - 1)\tau^r + 0\tau^r + 2\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$9r(r - 1) + 0 + 2 = 0$$

Or

$$9r^2 - 9r + 2 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{3}$$
$$r_2 = \frac{2}{3}$$

Since the roots are real and distinct, then the general solution is

$$y(\tau) = c_1y_1 + c_2y_2$$

Where $y_1 = \tau^{r_1}$ and $y_2 = \tau^{r_2}$. Hence

$$y(\tau) = c_1\tau^{\frac{1}{3}} + c_2\tau^{\frac{2}{3}}$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{c_1 3^{\frac{2}{3}} (t^3)^{\frac{1}{3}}}{3} + \frac{c_2 3^{\frac{1}{3}} (t^3)^{\frac{2}{3}}}{3}$$

Therefore the homogeneous solution y_h is

$$y_h = \frac{c_1 3^{\frac{2}{3}} (t^3)^{\frac{1}{3}}}{3} + \frac{c_2 3^{\frac{1}{3}} (t^3)^{\frac{2}{3}}}{3}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1y_1 + u_2y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = (t^3)^{\frac{1}{3}}$$

$$y_2 = (t^3)^{\frac{2}{3}}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} (t^3)^{\frac{1}{3}} & (t^3)^{\frac{2}{3}} \\ \frac{d}{dt} \left((t^3)^{\frac{1}{3}} \right) & \frac{d}{dt} \left((t^3)^{\frac{2}{3}} \right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} (t^3)^{\frac{1}{3}} & (t^3)^{\frac{2}{3}} \\ \frac{t^2}{(t^3)^{\frac{2}{3}}} & \frac{2t^2}{(t^3)^{\frac{1}{3}}} \end{vmatrix}$$

Therefore

$$W = \left((t^3)^{\frac{1}{3}} \right) \left(\frac{2t^2}{(t^3)^{\frac{1}{3}}} \right) - \left((t^3)^{\frac{2}{3}} \right) \left(\frac{t^2}{(t^3)^{\frac{2}{3}}} \right)$$

Which simplifies to

$$W = t^2$$

Which simplifies to

$$W = t^2$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{4(t^3)^{\frac{2}{3}} t^2}{t^4} dt$$

Which simplifies to

$$u_1 = - \int \frac{4(t^3)^{\frac{2}{3}}}{t^2} dt$$

Hence

$$u_1 = - \frac{4(t^3)^{\frac{2}{3}}}{t}$$

And Eq. (3) becomes

$$u_2 = \int \frac{4(t^3)^{\frac{1}{3}} t^2}{t^4} dt$$

Which simplifies to

$$u_2 = \int \frac{4(t^3)^{\frac{1}{3}}}{t^2} dt$$

Hence

$$u_2 = \frac{4(t^3)^{\frac{1}{3}} \ln(t)}{t}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -4t^2 + 4t^2 \ln(t)$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{c_1 3^{\frac{2}{3}} (t^3)^{\frac{1}{3}}}{3} + \frac{c_2 3^{\frac{1}{3}} (t^3)^{\frac{2}{3}}}{3} \right) + (-4t^2 + 4t^2 \ln(t)) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 3^{\frac{2}{3}} (t^3)^{\frac{1}{3}}}{3} + \frac{c_2 3^{\frac{1}{3}} (t^3)^{\frac{2}{3}}}{3} - 4t^2 + 4t^2 \ln(t) \quad (1)$$

Verification of solutions

$$y = \frac{c_1 3^{\frac{2}{3}} (t^3)^{\frac{1}{3}}}{3} + \frac{c_2 3^{\frac{1}{3}} (t^3)^{\frac{2}{3}}}{3} - 4t^2 + 4t^2 \ln(t)$$

Verified OK.

10.19.4 Solving as second order change of variable on x method 1 ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = t^2$, $B = -2t$, $C = 2$, $f(t) = 4t^2$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. Solving for y_h from

$$t^2 y'' - 2ty' + 2y = 0$$

In normal form the ode

$$t^2 y'' - 2ty' + 2y = 0 \tag{1}$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \tag{2}$$

Where

$$p(t) = -\frac{2}{t}$$
$$q(t) = \frac{2}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \tag{3}$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \tag{4}$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \tag{5}$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\begin{aligned}\tau' &= \frac{1}{c}\sqrt{q} \\ &= \frac{\sqrt{2}\sqrt{\frac{1}{t^2}}}{c} \\ \tau'' &= -\frac{\sqrt{2}}{c\sqrt{\frac{1}{t^2}}t^3}\end{aligned}\tag{6}$$

Substituting the above into (4) results in

$$\begin{aligned}p_1(\tau) &= \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \\ &= \frac{-\frac{\sqrt{2}}{c\sqrt{\frac{1}{t^2}}t^3} - \frac{2}{t}\frac{\sqrt{2}\sqrt{\frac{1}{t^2}}}{c}}{\left(\frac{\sqrt{2}\sqrt{\frac{1}{t^2}}}{c}\right)^2} \\ &= -\frac{3c\sqrt{2}}{2}\end{aligned}$$

Therefore ode (3) now becomes

$$\begin{aligned}y(\tau)'' + p_1y(\tau)' + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) - \frac{3c\sqrt{2}}{2}\left(\frac{d}{d\tau}y(\tau)\right) + c^2y(\tau) &= 0\end{aligned}\tag{7}$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{\frac{3\sqrt{2}c\tau}{4}}\left(c_1 \cosh\left(\frac{\sqrt{2}c\tau}{4}\right) + ic_2 \sinh\left(\frac{\sqrt{2}c\tau}{4}\right)\right)$$

Now from (6)

$$\begin{aligned}\tau &= \int \frac{1}{c}\sqrt{q} dt \\ &= \frac{\int \sqrt{2}\sqrt{\frac{1}{t^2}} dt}{c} \\ &= \frac{\sqrt{2}\sqrt{\frac{1}{t^2}}t \ln(t)}{c}\end{aligned}$$

Substituting the above into the solution obtained gives

$$y = t^{\frac{3}{2}} \left(c_1 \cosh \left(\frac{\ln(t)}{2} \right) + ic_2 \sinh \left(\frac{\ln(t)}{2} \right) \right)$$

Now the particular solution to this ODE is found

$$t^2 y'' - 2ty' + 2y = 4t^2$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = (t^3)^{\frac{1}{3}}$$

$$y_2 = (t^3)^{\frac{2}{3}}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \tag{3}$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} (t^3)^{\frac{1}{3}} & (t^3)^{\frac{2}{3}} \\ \frac{d}{dt} \left((t^3)^{\frac{1}{3}} \right) & \frac{d}{dt} \left((t^3)^{\frac{2}{3}} \right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} (t^3)^{\frac{1}{3}} & (t^3)^{\frac{2}{3}} \\ \frac{t^2}{(t^3)^{\frac{2}{3}}} & \frac{2t^2}{(t^3)^{\frac{1}{3}}} \end{vmatrix}$$

Therefore

$$W = \left((t^3)^{\frac{1}{3}} \right) \left(\frac{2t^2}{(t^3)^{\frac{1}{3}}} \right) - \left((t^3)^{\frac{2}{3}} \right) \left(\frac{t^2}{(t^3)^{\frac{2}{3}}} \right)$$

Which simplifies to

$$W = t^2$$

Which simplifies to

$$W = t^2$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{4(t^3)^{\frac{2}{3}} t^2}{t^4} dt$$

Which simplifies to

$$u_1 = - \int \frac{4(t^3)^{\frac{2}{3}}}{t^2} dt$$

Hence

$$u_1 = - \frac{4(t^3)^{\frac{2}{3}}}{t}$$

And Eq. (3) becomes

$$u_2 = \int \frac{4(t^3)^{\frac{1}{3}} t^2}{t^4} dt$$

Which simplifies to

$$u_2 = \int \frac{4(t^3)^{\frac{1}{3}}}{t^2} dt$$

Hence

$$u_2 = \frac{4(t^3)^{\frac{1}{3}} \ln(t)}{t}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -4t^2 + 4t^2 \ln(t)$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(t^{\frac{3}{2}} \left(c_1 \cosh \left(\frac{\ln(t)}{2} \right) + i c_2 \sinh \left(\frac{\ln(t)}{2} \right) \right) \right) + (-4t^2 + 4t^2 \ln(t)) \\ &= -4t^2 + 4t^2 \ln(t) + t^{\frac{3}{2}} \left(c_1 \cosh \left(\frac{\ln(t)}{2} \right) + i c_2 \sinh \left(\frac{\ln(t)}{2} \right) \right) \end{aligned}$$

Which simplifies to

$$y = i \sinh \left(\frac{\ln(t)}{2} \right) t^{\frac{3}{2}} c_2 + \cosh \left(\frac{\ln(t)}{2} \right) t^{\frac{3}{2}} c_1 + 4t^2 \ln(t) - 4t^2$$

Summary

The solution(s) found are the following

$$y = i \sinh \left(\frac{\ln(t)}{2} \right) t^{\frac{3}{2}} c_2 + \cosh \left(\frac{\ln(t)}{2} \right) t^{\frac{3}{2}} c_1 + 4t^2 \ln(t) - 4t^2 \quad (1)$$

Verification of solutions

$$y = i \sinh \left(\frac{\ln(t)}{2} \right) t^{\frac{3}{2}} c_2 + \cosh \left(\frac{\ln(t)}{2} \right) t^{\frac{3}{2}} c_1 + 4t^2 \ln(t) - 4t^2$$

Verified OK.

10.19.5 Solving as second order change of variable on y method 2 ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = t^2$, $B = -2t$, $C = 2$, $f(t) = 4t^2$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. Solving for y_h from

$$t^2y'' - 2ty' + 2y = 0$$

In normal form the ode

$$t^2y'' - 2ty' + 2y = 0 \tag{1}$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \tag{2}$$

Where

$$p(t) = -\frac{2}{t}$$

$$q(t) = \frac{2}{t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variable is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \tag{3}$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \tag{4}$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} - \frac{2n}{t} + \frac{2}{t^2} = 0 \tag{5}$$

Solving (5) for n gives

$$n = 2 \tag{6}$$

Substituting this value in (3) gives

$$v''(t) + \frac{2v'(t)}{t} = 0$$

$$v''(t) + \frac{2v'(t)}{t} = 0 \tag{7}$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{2u(t)}{t} = 0 \tag{8}$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{2u}{t} \end{aligned}$$

Where $f(t) = -\frac{2}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= -\frac{2}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{2}{t} dt \\ \ln(u) &= -2 \ln(t) + c_1 \\ u &= e^{-2 \ln(t) + c_1} \\ &= \frac{c_1}{t^2} \end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned} v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= -\frac{c_1}{t} + c_2 \end{aligned}$$

Hence

$$\begin{aligned} y &= v(t) t^n \\ &= \left(-\frac{c_1}{t} + c_2\right) t^2 \\ &= (c_2 t - c_1) t \end{aligned}$$

Now the particular solution to this ODE is found

$$t^2y'' - 2ty' + 2y = 4t^2$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1y_1 + u_2y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = t$$

$$y_2 = t^2$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t & t^2 \\ \frac{d}{dt}(t) & \frac{d}{dt}(t^2) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t & t^2 \\ 1 & 2t \end{vmatrix}$$

Therefore

$$W = (t)(2t) - (t^2)(1)$$

Which simplifies to

$$W = t^2$$

Which simplifies to

$$W = t^2$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{4t^4}{t^4} dt$$

Which simplifies to

$$u_1 = - \int 4dt$$

Hence

$$u_1 = -4t$$

And Eq. (3) becomes

$$u_2 = \int \frac{4t^3}{t^4} dt$$

Which simplifies to

$$u_2 = \int \frac{4}{t} dt$$

Hence

$$u_2 = 4 \ln(t)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -4t^2 + 4t^2 \ln(t)$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\left(-\frac{c_1}{t} + c_2 \right) t^2 \right) + (-4t^2 + 4t^2 \ln(t)) \\ &= -4t^2 + 4t^2 \ln(t) + \left(-\frac{c_1}{t} + c_2 \right) t^2 \end{aligned}$$

Which simplifies to

$$y = t(4 \ln(t) t + c_2 t - c_1 - 4t)$$

Summary

The solution(s) found are the following

$$y = t(4 \ln(t) t + c_2 t - c_1 - 4t) \quad (1)$$

Verification of solutions

$$y = t(4 \ln(t) t + c_2 t - c_1 - 4t)$$

Verified OK.

10.19.6 Solving as second order ode non constant coeff transformation on B ode

Given an ode of the form

$$Ay'' + By' + Cy = F(t)$$

This method reduces the order ode the ODE by one by applying the transformation

$$y = Bv$$

This results in

$$\begin{aligned} y' &= B'v + v'B \\ y'' &= B''v + B'v' + v''B + v'B' \\ &= v''B + 2v' + B' + B''v \end{aligned}$$

And now the original ode becomes

$$\begin{aligned} A(v''B + 2v'B' + B''v) + B(B'v + v'B) + CBv &= 0 \\ ABv'' + (2AB' + B^2)v' + (AB'' + BB' + CB)v &= 0 \end{aligned} \quad (1)$$

If the term $AB'' + BB' + CB$ is zero, then this method works and can be used to solve

$$ABv'' + (2AB' + B^2)v' = 0$$

By Using $u = v'$ which reduces the order of the above ode to one. The new ode is

$$ABu' + (2AB' + B^2)u = 0$$

The above ode is first order ode which is solved for u . Now a new ode $v' = u$ is solved for v as first order ode. Then the final solution is obtain from $y = Bv$.

This method works only if the term $AB'' + BB' + CB$ is zero. The given ODE shows that

$$\begin{aligned}A &= t^2 \\B &= -2t \\C &= 2 \\F &= 4t^2\end{aligned}$$

The above shows that for this ode

$$\begin{aligned}AB'' + BB' + CB &= (t^2)(0) + (-2t)(-2) + (2)(-2t) \\&= 0\end{aligned}$$

Hence the ode in v given in (1) now simplifies to

$$-2t^3v'' + (0)v' = 0$$

Now by applying $v' = u$ the above becomes

$$-2t^3u'(t) = 0$$

Which is now solved for u . Integrating both sides gives

$$\begin{aligned}u(t) &= \int 0 \, dt \\&= c_1\end{aligned}$$

The ode for v now becomes

$$\begin{aligned}v' &= u \\&= c_1\end{aligned}$$

Which is now solved for v . Integrating both sides gives

$$\begin{aligned}v(t) &= \int c_1 \, dt \\&= c_1t + c_2\end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}y_h(t) &= Bv \\&= (-2t)(c_1t + c_2) \\&= -2t(c_1t + c_2)\end{aligned}$$

And now the particular solution $y_p(t)$ will be found. The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = t$$

$$y_2 = t^2$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t & t^2 \\ \frac{d}{dt}(t) & \frac{d}{dt}(t^2) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t & t^2 \\ 1 & 2t \end{vmatrix}$$

Therefore

$$W = (t)(2t) - (t^2) \quad (1)$$

Which simplifies to

$$W = t^2$$

Which simplifies to

$$W = t^2$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{4t^4}{t^4} dt$$

Which simplifies to

$$u_1 = - \int 4dt$$

Hence

$$u_1 = -4t$$

And Eq. (3) becomes

$$u_2 = \int \frac{4t^3}{t^4} dt$$

Which simplifies to

$$u_2 = \int \frac{4}{t} dt$$

Hence

$$u_2 = 4 \ln(t)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -4t^2 + 4t^2 \ln(t)$$

Hence the complete solution is

$$\begin{aligned} y(t) &= y_h + y_p \\ &= (-2t(c_1t + c_2)) + (-4t^2 + 4t^2 \ln(t)) \\ &= -2(-2 \ln(t)t + (c_1 + 2)t + c_2)t \end{aligned}$$

Summary

The solution(s) found are the following

$$y = -2(-2 \ln(t)t + (c_1 + 2)t + c_2)t \quad (1)$$

Verification of solutions

$$y = -2(-2 \ln(t)t + (c_1 + 2)t + c_2)t$$

Verified OK.

10.19.7 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' - 2ty' + 2y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t^2 \\ B &= -2t \\ C &= 2 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{0}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 0 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = 0 \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 504: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - -\infty \\ &= \infty \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is *infinity* then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 0$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = 1$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-2t}{t^2} dt} \end{aligned}$$

$$\begin{aligned}
&= z_1 e^{\ln(t)} \\
&= z_1(t)
\end{aligned}$$

Which simplifies to

$$y_1 = t$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}
y_2 &= y_1 \int \frac{e^{\int -\frac{-2t}{t^2} dt}}{(y_1)^2} dt \\
&= y_1 \int \frac{e^{2\ln(t)}}{(y_1)^2} dt \\
&= y_1(t)
\end{aligned}$$

Therefore the solution is

$$\begin{aligned}
y &= c_1 y_1 + c_2 y_2 \\
&= c_1(t) + c_2(t(t))
\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$t^2 y'' - 2ty' + 2y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_2 t^2 + c_1 t$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = t$$

$$y_2 = t^2$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t & t^2 \\ \frac{d}{dt}(t) & \frac{d}{dt}(t^2) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t & t^2 \\ 1 & 2t \end{vmatrix}$$

Therefore

$$W = (t)(2t) - (t^2) \quad (1)$$

Which simplifies to

$$W = t^2$$

Which simplifies to

$$W = t^2$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{4t^4}{t^4} dt$$

Which simplifies to

$$u_1 = - \int 4dt$$

Hence

$$u_1 = -4t$$

And Eq. (3) becomes

$$u_2 = \int \frac{4t^3}{t^4} dt$$

Which simplifies to

$$u_2 = \int \frac{4}{t} dt$$

Hence

$$u_2 = 4 \ln(t)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -4t^2 + 4t^2 \ln(t)$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_2 t^2 + c_1 t) + (-4t^2 + 4t^2 \ln(t)) \end{aligned}$$

Which simplifies to

$$y = t(c_2 t + c_1) - 4t^2 + 4t^2 \ln(t)$$

Summary

The solution(s) found are the following

$$y = t(c_2t + c_1) - 4t^2 + 4t^2 \ln(t) \quad (1)$$

Verification of solutions

$$y = t(c_2t + c_1) - 4t^2 + 4t^2 \ln(t)$$

Verified OK.

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    checking if the LODE is of Euler type
    <- LODE of Euler type successful
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 18

```
dsolve(t^2*diff(y(t),t$2)-2*t*diff(y(t),t)+2*y(t) = 4*t^2,y(t), singsol=all)
```

$$y(t) = t(4t \ln(t) + (c_1 - 4)t + c_2)$$

✓ Solution by Mathematica

Time used: 0.013 (sec). Leaf size: 21

```
DSolve[t^2*y''[t]-2*t*y'[t]+2*y[t] ==4*t^2,y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow t(4t \log(t) + (-4 + c_2)t + c_1)$$

10.20 problem 30

10.20.1 Solving as second order euler ode ode	2971
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10.20.3 Solving as second order change of variable on x method 1 ode .	2979
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Internal problem ID [702]

Internal file name [OUTPUT/702_Sunday_June_05_2022_01_47_18_AM_49530951/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 30.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_euler_ode", "exact linear second order ode", "second_order_integrable_as_is", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second_order_change_of_variable_on_y_method_2"

Maple gives the following as the ode type

```
[[_2nd_order, _exact, _linear, _nonhomogeneous]]
```

$$t^2y'' + 7ty' + 5y = t$$

10.20.1 Solving as second order euler ode ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = t^2, B = 7t, C = 5, f(t) = t$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. Solving for y_h from

$$t^2y'' + 7ty' + 5y = 0$$

This is Euler second order ODE. Let the solution be $y = t^r$, then $y' = rt^{r-1}$ and $y'' = r(r-1)t^{r-2}$. Substituting these back into the given ODE gives

$$t^2(r(r-1))t^{r-2} + 7trt^{r-1} + 5t^r = 0$$

Simplifying gives

$$r(r-1)t^r + 7rt^r + 5t^r = 0$$

Since $t^r \neq 0$ then dividing throughout by t^r gives

$$r(r-1) + 7r + 5 = 0$$

Or

$$r^2 + 6r + 5 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -5$$

$$r_2 = -1$$

Since the roots are real and distinct, then the general solution is

$$y = c_1y_1 + c_2y_2$$

Where $y_1 = t^{r_1}$ and $y_2 = t^{r_2}$. Hence

$$y = \frac{c_1}{t^5} + \frac{c_2}{t}$$

Next, we find the particular solution to the ODE

$$t^2y'' + 7ty' + 5y = t$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \frac{1}{t^5}$$

$$y_2 = \frac{1}{t}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{1}{t^5} & \frac{1}{t} \\ \frac{d}{dt} \left(\frac{1}{t^5} \right) & \frac{d}{dt} \left(\frac{1}{t} \right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{1}{t^5} & \frac{1}{t} \\ -\frac{5}{t^6} & -\frac{1}{t^2} \end{vmatrix}$$

Therefore

$$W = \left(\frac{1}{t^5} \right) \left(-\frac{1}{t^2} \right) - \left(\frac{1}{t} \right) \left(-\frac{5}{t^6} \right)$$

Which simplifies to

$$W = \frac{4}{t^7}$$

Which simplifies to

$$W = \frac{4}{t^7}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{1}{\frac{4}{t^5}} dt$$

Which simplifies to

$$u_1 = - \int \frac{t^5}{4} dt$$

Hence

$$u_1 = -\frac{t^6}{24}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{1}{t^4}}{\frac{4}{t^5}} dt$$

Which simplifies to

$$u_2 = \int \frac{t}{4} dt$$

Hence

$$u_2 = \frac{t^2}{8}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{t}{12}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \frac{t}{12} + \frac{c_1}{t^5} + \frac{c_2}{t} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{t}{12} + \frac{c_1}{t^5} + \frac{c_2}{t} \quad (1)$$

Verification of solutions

$$y = \frac{t}{12} + \frac{c_1}{t^5} + \frac{c_2}{t}$$

Verified OK.

10.20.2 Solving as second order change of variable on x method 2 ode

This is second order non-homogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$.
 y_h is the solution to

$$t^2y'' + 7ty' + 5y = 0$$

In normal form the ode

$$t^2y'' + 7ty' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{7}{t}$$
$$q(t) = \frac{5}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t) \tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(t) + p(t) \tau'(t) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-(\int p(t) dt)} dt \\ &= \int e^{-(\int \frac{7}{t} dt)} dt \\ &= \int e^{-7 \ln(t)} dt \\ &= \int \frac{1}{t^7} dt \\ &= -\frac{1}{6t^6} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(t)}{\tau'(t)^2} \\ &= \frac{\frac{5}{t^2}}{\frac{1}{t^{14}}} \\ &= 5t^{12} \end{aligned} \quad (7)$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2} y(\tau) + q_1 y(\tau) &= 0 \\ \frac{d^2}{d\tau^2} y(\tau) + 5t^{12} y(\tau) &= 0 \end{aligned}$$

But in terms of τ

$$5t^{12} = \frac{5}{36\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{5y(\tau)}{36\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. The ode can be written as

$$36\left(\frac{d^2}{d\tau^2}y(\tau)\right)\tau^2 + 5y(\tau) = 0$$

Which shows it is a Euler ODE. This is Euler second order ODE. Let the solution be $y(\tau) = \tau^r$, then $y' = r\tau^{r-1}$ and $y'' = r(r-1)\tau^{r-2}$. Substituting these back into the given ODE gives

$$36\tau^2(r(r-1))\tau^{r-2} + 0r\tau^{r-1} + 5\tau^r = 0$$

Simplifying gives

$$36r(r-1)\tau^r + 0\tau^r + 5\tau^r = 0$$

Since $\tau^r \neq 0$ then dividing throughout by τ^r gives

$$36r(r-1) + 0 + 5 = 0$$

Or

$$36r^2 - 36r + 5 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = \frac{1}{6}$$

$$r_2 = \frac{5}{6}$$

Since the roots are real and distinct, then the general solution is

$$y(\tau) = c_1y_1 + c_2y_2$$

Where $y_1 = \tau^{r_1}$ and $y_2 = \tau^{r_2}$. Hence

$$y(\tau) = c_1\tau^{\frac{1}{6}} + c_2\tau^{\frac{5}{6}}$$

The above solution is now transformed back to y using (6) which results in

$$y = \frac{c_1 6^{\frac{5}{6}} \left(-\frac{1}{t^6}\right)^{\frac{1}{6}}}{6} + \frac{c_2 6^{\frac{1}{6}} \left(-\frac{1}{t^6}\right)^{\frac{5}{6}}}{6}$$

Therefore the homogeneous solution y_h is

$$y_h = \frac{c_1 6^{\frac{5}{6}} \left(-\frac{1}{t^6}\right)^{\frac{1}{6}}}{6} + \frac{c_2 6^{\frac{1}{6}} \left(-\frac{1}{t^6}\right)^{\frac{5}{6}}}{6}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \left(-\frac{1}{t^6}\right)^{\frac{1}{6}}$$

$$y_2 = \left(-\frac{1}{t^6}\right)^{\frac{5}{6}}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \left(-\frac{1}{t^6}\right)^{\frac{1}{6}} & \left(-\frac{1}{t^6}\right)^{\frac{5}{6}} \\ \frac{d}{dt} \left(\left(-\frac{1}{t^6}\right)^{\frac{1}{6}}\right) & \frac{d}{dt} \left(\left(-\frac{1}{t^6}\right)^{\frac{5}{6}}\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \left(-\frac{1}{t^6}\right)^{\frac{1}{6}} & \left(-\frac{1}{t^6}\right)^{\frac{5}{6}} \\ \frac{1}{\left(-\frac{1}{t^6}\right)^{\frac{5}{6}} t^7} & \frac{5}{\left(-\frac{1}{t^6}\right)^{\frac{1}{6}} t^7} \end{vmatrix}$$

Therefore

$$W = \left(\left(-\frac{1}{t^6} \right)^{\frac{1}{6}} \right) \left(\frac{5}{\left(-\frac{1}{t^6} \right)^{\frac{1}{6}} t^7} \right) - \left(\left(-\frac{1}{t^6} \right)^{\frac{5}{6}} \right) \left(\frac{1}{\left(-\frac{1}{t^6} \right)^{\frac{5}{6}} t^7} \right)$$

Which simplifies to

$$W = \frac{4}{t^7}$$

Which simplifies to

$$W = \frac{4}{t^7}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\left(-\frac{1}{t^6} \right)^{\frac{5}{6}} t}{\frac{4}{t^5}} dt$$

Which simplifies to

$$u_1 = - \int \frac{\left(-\frac{1}{t^6} \right)^{\frac{5}{6}} t^6}{4} dt$$

Hence

$$u_1 = - \frac{\left(-\frac{1}{t^6} \right)^{\frac{5}{6}} t^7}{8}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\left(-\frac{1}{t^6} \right)^{\frac{1}{6}} t}{\frac{4}{t^5}} dt$$

Which simplifies to

$$u_2 = \int \frac{\left(-\frac{1}{t^6} \right)^{\frac{1}{6}} t^6}{4} dt$$

Hence

$$u_2 = \frac{\left(-\frac{1}{t^6} \right)^{\frac{1}{6}} t^7}{24}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{t}{12}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{c_1 6^{\frac{5}{6}} \left(-\frac{1}{t^6}\right)^{\frac{1}{6}}}{6} + \frac{c_2 6^{\frac{1}{6}} \left(-\frac{1}{t^6}\right)^{\frac{5}{6}}}{6} \right) + \left(\frac{t}{12} \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1 6^{\frac{5}{6}} \left(-\frac{1}{t^6}\right)^{\frac{1}{6}}}{6} + \frac{c_2 6^{\frac{1}{6}} \left(-\frac{1}{t^6}\right)^{\frac{5}{6}}}{6} + \frac{t}{12} \quad (1)$$

Verification of solutions

$$y = \frac{c_1 6^{\frac{5}{6}} \left(-\frac{1}{t^6}\right)^{\frac{1}{6}}}{6} + \frac{c_2 6^{\frac{1}{6}} \left(-\frac{1}{t^6}\right)^{\frac{5}{6}}}{6} + \frac{t}{12}$$

Verified OK.

10.20.3 Solving as second order change of variable on x method 1 ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = t^2$, $B = 7t$, $C = 5$, $f(t) = t$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. Solving for y_h from

$$t^2 y'' + 7ty' + 5y = 0$$

In normal form the ode

$$t^2 y'' + 7ty' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{7}{t}$$
$$q(t) = \frac{5}{t^2}$$

Applying change of variables $\tau = g(t)$ to (2) results

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(t)}{\tau'(t)^2} \quad (5)$$

Let $q_1 = c^2$ where c is some constant. Therefore from (5)

$$\tau' = \frac{1}{c}\sqrt{q}$$
$$= \frac{\sqrt{5}\sqrt{\frac{1}{t^2}}}{c} \quad (6)$$
$$\tau'' = -\frac{\sqrt{5}}{c\sqrt{\frac{1}{t^2}}t^3}$$

Substituting the above into (4) results in

$$p_1(\tau) = \frac{\tau''(t) + p(t)\tau'(t)}{\tau'(t)^2}$$
$$= \frac{-\frac{\sqrt{5}}{c\sqrt{\frac{1}{t^2}}t^3} + \frac{7}{t}\frac{\sqrt{5}\sqrt{\frac{1}{t^2}}}{c}}{\left(\frac{\sqrt{5}\sqrt{\frac{1}{t^2}}}{c}\right)^2}$$
$$= \frac{6c\sqrt{5}}{5}$$

Therefore ode (3) now becomes

$$\begin{aligned} y(\tau)'' + p_1 y(\tau)' + q_1 y(\tau) &= 0 \\ \frac{d^2}{d\tau^2} y(\tau) + \frac{6c\sqrt{5}}{5} \left(\frac{d}{d\tau} y(\tau) \right) + c^2 y(\tau) &= 0 \end{aligned} \quad (7)$$

The above ode is now solved for $y(\tau)$. Since the ode is now constant coefficients, it can be easily solved to give

$$y(\tau) = e^{-\frac{3\sqrt{5}c\tau}{5}} \left(c_1 \cosh \left(\frac{2\sqrt{5}c\tau}{5} \right) + ic_2 \sinh \left(\frac{2\sqrt{5}c\tau}{5} \right) \right)$$

Now from (6)

$$\begin{aligned} \tau &= \int \frac{1}{c} \sqrt{q} dt \\ &= \frac{\int \sqrt{5} \sqrt{\frac{1}{t^2}} dt}{c} \\ &= \frac{\sqrt{5} \sqrt{\frac{1}{t^2}} t \ln(t)}{c} \end{aligned}$$

Substituting the above into the solution obtained gives

$$y = \frac{c_1 \cosh(2 \ln(t)) + ic_2 \sinh(2 \ln(t))}{t^3}$$

Now the particular solution to this ODE is found

$$t^2 y'' + 7ty' + 5y = t$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned} y_1 &= \left(-\frac{1}{t^6} \right)^{\frac{1}{6}} \\ y_2 &= \left(-\frac{1}{t^6} \right)^{\frac{5}{6}} \end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \left(-\frac{1}{t^6}\right)^{\frac{1}{6}} & \left(-\frac{1}{t^6}\right)^{\frac{5}{6}} \\ \frac{d}{dt} \left(\left(-\frac{1}{t^6}\right)^{\frac{1}{6}}\right) & \frac{d}{dt} \left(\left(-\frac{1}{t^6}\right)^{\frac{5}{6}}\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \left(-\frac{1}{t^6}\right)^{\frac{1}{6}} & \left(-\frac{1}{t^6}\right)^{\frac{5}{6}} \\ \frac{1}{\left(-\frac{1}{t^6}\right)^{\frac{5}{6}} t^7} & \frac{5}{\left(-\frac{1}{t^6}\right)^{\frac{1}{6}} t^7} \end{vmatrix}$$

Therefore

$$W = \left(\left(-\frac{1}{t^6}\right)^{\frac{1}{6}} \right) \left(\frac{5}{\left(-\frac{1}{t^6}\right)^{\frac{1}{6}} t^7} \right) - \left(\left(-\frac{1}{t^6}\right)^{\frac{5}{6}} \right) \left(\frac{1}{\left(-\frac{1}{t^6}\right)^{\frac{5}{6}} t^7} \right)$$

Which simplifies to

$$W = \frac{4}{t^7}$$

Which simplifies to

$$W = \frac{4}{t^7}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\left(-\frac{1}{t^6}\right)^{\frac{5}{6}} t}{\frac{4}{t^5}} dt$$

Which simplifies to

$$u_1 = - \int \frac{\left(-\frac{1}{t^6}\right)^{\frac{5}{6}} t^6}{4} dt$$

Hence

$$u_1 = - \frac{\left(-\frac{1}{t^6}\right)^{\frac{5}{6}} t^7}{8}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\left(-\frac{1}{t^6}\right)^{\frac{1}{6}} t}{\frac{4}{t^5}} dt$$

Which simplifies to

$$u_2 = \int \frac{\left(-\frac{1}{t^6}\right)^{\frac{1}{6}} t^6}{4} dt$$

Hence

$$u_2 = \frac{\left(-\frac{1}{t^6}\right)^{\frac{1}{6}} t^7}{24}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{t}{12}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{c_1 \cosh(2 \ln(t)) + ic_2 \sinh(2 \ln(t))}{t^3} \right) + \left(\frac{t}{12} \right) \\ &= \frac{t}{12} + \frac{c_1 \cosh(2 \ln(t)) + ic_2 \sinh(2 \ln(t))}{t^3} \end{aligned}$$

Which simplifies to

$$y = \frac{t^6 + (6ic_2 + 6c_1)t^4 - 6ic_2 + 6c_1}{12t^5}$$

Summary

The solution(s) found are the following

$$y = \frac{t^6 + (6ic_2 + 6c_1)t^4 - 6ic_2 + 6c_1}{12t^5} \quad (1)$$

Verification of solutions

$$y = \frac{t^6 + (6ic_2 + 6c_1)t^4 - 6ic_2 + 6c_1}{12t^5}$$

Verified OK.

10.20.4 Solving as second order change of variable on y method 2 ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = t^2$, $B = 7t$, $C = 5$, $f(t) = t$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. Solving for y_h from

$$t^2y'' + 7ty' + 5y = 0$$

In normal form the ode

$$t^2y'' + 7ty' + 5y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = \frac{7}{t}$$
$$q(t) = \frac{5}{t^2}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} + \frac{7n}{t^2} + \frac{5}{t^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = -1 \quad (6)$$

Substituting this value in (3) gives

$$\begin{aligned} v''(t) + \frac{5v'(t)}{t} &= 0 \\ v''(t) + \frac{5v'(t)}{t} &= 0 \end{aligned} \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \frac{5u(t)}{t} = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= -\frac{5u}{t} \end{aligned}$$

Where $f(t) = -\frac{5}{t}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= -\frac{5}{t} dt \\ \int \frac{1}{u} du &= \int -\frac{5}{t} dt \\ \ln(u) &= -5 \ln(t) + c_1 \\ u &= e^{-5 \ln(t) + c_1} \\ &= \frac{c_1}{t^5} \end{aligned}$$

Now that $u(t)$ is known, then

$$\begin{aligned}v'(t) &= u(t) \\v(t) &= \int u(t) dt + c_2 \\&= -\frac{c_1}{4t^4} + c_2\end{aligned}$$

Hence

$$\begin{aligned}y &= v(t) t^n \\&= \frac{-\frac{c_1}{4t^4} + c_2}{t} \\&= \frac{4c_2t^4 - c_1}{4t^5}\end{aligned}$$

Now the particular solution to this ODE is found

$$t^2y'' + 7ty' + 5y = t$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1y_1 + u_2y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned}y_1 &= \frac{1}{t^5} \\y_2 &= \frac{1}{t}\end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \tag{3}$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{1}{t^5} & \frac{1}{t} \\ \frac{d}{dt}\left(\frac{1}{t^5}\right) & \frac{d}{dt}\left(\frac{1}{t}\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{1}{t^5} & \frac{1}{t} \\ -\frac{5}{t^6} & -\frac{1}{t^2} \end{vmatrix}$$

Therefore

$$W = \left(\frac{1}{t^5}\right)\left(-\frac{1}{t^2}\right) - \left(\frac{1}{t}\right)\left(-\frac{5}{t^6}\right)$$

Which simplifies to

$$W = \frac{4}{t^7}$$

Which simplifies to

$$W = \frac{4}{t^7}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{1}{\frac{4}{t^5}} dt$$

Which simplifies to

$$u_1 = - \int \frac{t^5}{4} dt$$

Hence

$$u_1 = -\frac{t^6}{24}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{1}{t^4}}{\frac{4}{t^5}} dt$$

Which simplifies to

$$u_2 = \int \frac{t}{4} dt$$

Hence

$$u_2 = \frac{t^2}{8}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{t}{12}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{-\frac{c_1}{4t^4} + c_2}{t} \right) + \left(\frac{t}{12} \right) \\ &= \frac{t}{12} + \frac{-\frac{c_1}{4t^4} + c_2}{t} \end{aligned}$$

Which simplifies to

$$y = \frac{t}{12} + \frac{-\frac{c_1}{4t^4} + c_2}{t}$$

Summary

The solution(s) found are the following

$$y = \frac{t}{12} + \frac{-\frac{c_1}{4t^4} + c_2}{t} \tag{1}$$

Verification of solutions

$$y = \frac{t}{12} + \frac{-\frac{c_1}{4t^4} + c_2}{t}$$

Verified OK.

10.20.5 Solving as second order integrable as is ode

Integrating both sides of the ODE w.r.t t gives

$$\int (t^2 y'' + 7ty' + 5y) dt = \int t dt$$
$$y't^2 + 5yt = \frac{t^2}{2} + c_1$$

Which is now solved for y .

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{5}{t}$$
$$q(t) = \frac{t^2 + 2c_1}{2t^2}$$

Hence the ode is

$$y' + \frac{5y}{t} = \frac{t^2 + 2c_1}{2t^2}$$

The integrating factor μ is

$$\mu = e^{\int \frac{5}{t} dt}$$
$$= t^5$$

The ode becomes

$$\frac{d}{dt}(\mu y) = (\mu) \left(\frac{t^2 + 2c_1}{2t^2} \right)$$
$$\frac{d}{dt}(t^5 y) = (t^5) \left(\frac{t^2 + 2c_1}{2t^2} \right)$$
$$d(t^5 y) = \left(\frac{(t^2 + 2c_1)t^3}{2} \right) dt$$

Integrating gives

$$t^5 y = \int \frac{(t^2 + 2c_1)t^3}{2} dt$$
$$t^5 y = \frac{1}{12}t^6 + \frac{1}{4}c_1 t^4 + c_2$$

Dividing both sides by the integrating factor $\mu = t^5$ results in

$$y = \frac{\frac{1}{12}t^6 + \frac{1}{4}c_1t^4}{t^5} + \frac{c_2}{t^5}$$

which simplifies to

$$y = \frac{t^6 + 3c_1t^4 + 12c_2}{12t^5}$$

Summary

The solution(s) found are the following

$$y = \frac{t^6 + 3c_1t^4 + 12c_2}{12t^5} \quad (1)$$

Verification of solutions

$$y = \frac{t^6 + 3c_1t^4 + 12c_2}{12t^5}$$

Verified OK.

10.20.6 Solving as type second_order_integrable_as_is (not using ABC version)

Writing the ode as

$$t^2y'' + 7ty' + 5y = t$$

Integrating both sides of the ODE w.r.t t gives

$$\int (t^2y'' + 7ty' + 5y) dt = \int t dt$$
$$y't^2 + 5yt = \frac{t^2}{2} + c_1$$

Which is now solved for y .

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{5}{t}$$
$$q(t) = \frac{t^2 + 2c_1}{2t^2}$$

Hence the ode is

$$y' + \frac{5y}{t} = \frac{t^2 + 2c_1}{2t^2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int \frac{5}{t} dt} \\ &= t^5\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu y) &= (\mu) \left(\frac{t^2 + 2c_1}{2t^2} \right) \\ \frac{d}{dt}(t^5 y) &= (t^5) \left(\frac{t^2 + 2c_1}{2t^2} \right) \\ d(t^5 y) &= \left(\frac{(t^2 + 2c_1)t^3}{2} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}t^5 y &= \int \frac{(t^2 + 2c_1)t^3}{2} dt \\ t^5 y &= \frac{1}{12}t^6 + \frac{1}{4}c_1 t^4 + c_2\end{aligned}$$

Dividing both sides by the integrating factor $\mu = t^5$ results in

$$y = \frac{\frac{1}{12}t^6 + \frac{1}{4}c_1 t^4}{t^5} + \frac{c_2}{t^5}$$

which simplifies to

$$y = \frac{t^6 + 3c_1 t^4 + 12c_2}{12t^5}$$

Summary

The solution(s) found are the following

$$y = \frac{t^6 + 3c_1 t^4 + 12c_2}{12t^5} \tag{1}$$

Verification of solutions

$$y = \frac{t^6 + 3c_1 t^4 + 12c_2}{12t^5}$$

Verified OK.

10.20.7 Solving using Kovacic algorithm

Writing the ode as

$$t^2 y'' + 7ty' + 5y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t^2 \\ B &= 7t \\ C &= 5 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{15}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 15 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(\frac{15}{4t^2} \right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 505: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{15}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = \frac{15}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{5}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{3}{2} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{t^2}$ in the Laurent series expansion of r at ∞ . which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = \frac{15}{4t^2}$$

Since the $\text{gcd}(s, t) = 1$. This gives $b = \frac{15}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{5}{2} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{3}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = \frac{15}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{5}{2}$	$-\frac{3}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{5}{2}$	$-\frac{3}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = -\frac{3}{2}$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= -\frac{3}{2} - \left(-\frac{3}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t-c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{t-c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= -\frac{3}{2t} + (-)(0) \\ &= -\frac{3}{2t} \\ &= -\frac{3}{2t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r)p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2\left(-\frac{3}{2t}\right)(0) + \left(\left(\frac{3}{2t^2}\right) + \left(-\frac{3}{2t}\right)^2 - \left(\frac{15}{4t^2}\right) \right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int -\frac{3}{2t} dt} \\ &= \frac{1}{t^{\frac{3}{2}}} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned}y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\&= z_1 e^{-\int \frac{1}{2} \frac{7t}{t^2} dt} \\&= z_1 e^{-\frac{7 \ln(t)}{2}} \\&= z_1 \left(\frac{1}{t^{\frac{7}{2}}} \right)\end{aligned}$$

Which simplifies to

$$y_1 = \frac{1}{t^{\frac{7}{2}}}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{7t}{t^2} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{-7 \ln(t)}}{(y_1)^2} dt \\&= y_1 \left(\frac{t^4}{4} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1 \left(\frac{1}{t^{\frac{7}{2}}} \right) + c_2 \left(\frac{1}{t^{\frac{7}{2}}} \left(\frac{t^4}{4} \right) \right)\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$t^2y'' + 7ty' + 5y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = \frac{c_1}{t^5} + \frac{c_2}{4t}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1y_1 + u_2y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \frac{1}{t^5}$$

$$y_2 = \frac{1}{4t}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2f(t)}{aW(t)} \tag{2}$$

$$u_2 = \int \frac{y_1f(t)}{aW(t)} \tag{3}$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{1}{t^5} & \frac{1}{4t} \\ \frac{d}{dt}\left(\frac{1}{t^5}\right) & \frac{d}{dt}\left(\frac{1}{4t}\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{1}{t^5} & \frac{1}{4t} \\ -\frac{5}{t^6} & -\frac{1}{4t^2} \end{vmatrix}$$

Therefore

$$W = \left(\frac{1}{t^5}\right) \left(-\frac{1}{4t^2}\right) - \left(\frac{1}{4t}\right) \left(-\frac{5}{t^6}\right)$$

Which simplifies to

$$W = \frac{1}{t^7}$$

Which simplifies to

$$W = \frac{1}{t^7}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{1}{4}}{\frac{1}{t^5}} dt$$

Which simplifies to

$$u_1 = - \int \frac{t^5}{4} dt$$

Hence

$$u_1 = -\frac{t^6}{24}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{1}{t^4}}{\frac{1}{t^5}} dt$$

Which simplifies to

$$u_2 = \int t dt$$

Hence

$$u_2 = \frac{t^2}{2}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{t}{12}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{c_1}{t^5} + \frac{c_2}{4t} \right) + \left(\frac{t}{12} \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \frac{c_1}{t^5} + \frac{c_2}{4t} + \frac{t}{12} \quad (1)$$

Verification of solutions

$$y = \frac{c_1}{t^5} + \frac{c_2}{4t} + \frac{t}{12}$$

Verified OK.

10.20.8 Solving as exact linear second order ode ode

An ode of the form

$$p(t) y'' + q(t) y' + r(t) y = s(t)$$

is exact if

$$p''(t) - q'(t) + r(t) = 0 \quad (1)$$

For the given ode we have

$$\begin{aligned} p(x) &= t^2 \\ q(x) &= 7t \\ r(x) &= 5 \\ s(x) &= t \end{aligned}$$

Hence

$$\begin{aligned} p''(x) &= 2 \\ q'(x) &= 7 \end{aligned}$$

Therefore (1) becomes

$$2 - (7) + (5) = 0$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$(p(t)y' + (q(t) - p'(t))y)' = s(x)$$

Integrating gives

$$p(t)y' + (q(t) - p'(t))y = \int s(t) dt$$

Substituting the above values for p, q, r, s gives

$$y't^2 + 5yt = \int t dt$$

We now have a first order ode to solve which is

$$y't^2 + 5yt = \frac{t^2}{2} + c_1$$

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(t)y = q(t)$$

Where here

$$p(t) = \frac{5}{t}$$
$$q(t) = \frac{t^2 + 2c_1}{2t^2}$$

Hence the ode is

$$y' + \frac{5y}{t} = \frac{t^2 + 2c_1}{2t^2}$$

The integrating factor μ is

$$\mu = e^{\int \frac{5}{t} dt}$$
$$= t^5$$

The ode becomes

$$\frac{d}{dt}(\mu y) = (\mu) \left(\frac{t^2 + 2c_1}{2t^2} \right)$$
$$\frac{d}{dt}(t^5 y) = (t^5) \left(\frac{t^2 + 2c_1}{2t^2} \right)$$
$$d(t^5 y) = \left(\frac{(t^2 + 2c_1)t^3}{2} \right) dt$$

Integrating gives

$$t^5 y = \int \frac{(t^2 + 2c_1)t^3}{2} dt$$
$$t^5 y = \frac{1}{12}t^6 + \frac{1}{4}c_1 t^4 + c_2$$

Dividing both sides by the integrating factor $\mu = t^5$ results in

$$y = \frac{\frac{1}{12}t^6 + \frac{1}{4}c_1 t^4}{t^5} + \frac{c_2}{t^5}$$

which simplifies to

$$y = \frac{t^6 + 3c_1 t^4 + 12c_2}{12t^5}$$

Summary

The solution(s) found are the following

$$y = \frac{t^6 + 3c_1 t^4 + 12c_2}{12t^5} \quad (1)$$

Verification of solutions

$$y = \frac{t^6 + 3c_1 t^4 + 12c_2}{12t^5}$$

Verified OK.

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
<- high order exact linear fully integrable successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 27

```
dsolve(t^2*dif(y(t),t$2)+7*t*dif(y(t),t)+5*y(t) = t,y(t), singsol=all)
```

$$y(t) = \frac{t^6 + 3c_1 t^4 - 4c_1^3 + 12c_2}{12t^5}$$

✓ Solution by Mathematica

Time used: 0.015 (sec). Leaf size: 23

```
DSolve[t^2*y''[t]+7*t*y'[t]+5*y[t]==t,y[t],t,IncludeSingularSolutions->True]
```

$$y(t) \rightarrow \frac{c_1}{t^5} + \frac{t}{12} + \frac{c_2}{t}$$

10.21 problem 31

10.21.1 Solving as second order ode non constant coeff transformation on B ode	3003
10.21.2 Solving using Kovacic algorithm	3008

Internal problem ID [703]

Internal file name [OUTPUT/703_Sunday_June_05_2022_01_47_19_AM_67928080/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 31.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_ode_non_constant_coeff_transformation_on_B"

Maple gives the following as the ode type

```
[[_2nd_order, _with_linear_symmetries]]
```

$$ty'' - (t + 1)y' + y = e^{2t}t^2$$

10.21.1 Solving as second order ode non constant coeff transformation on B ode

Given an ode of the form

$$Ay'' + By' + Cy = F(t)$$

This method reduces the order ode the ODE by one by applying the transformation

$$y = Bv$$

This results in

$$\begin{aligned}y' &= B'v + v'B \\y'' &= B''v + B'v' + v''B + v'B' \\ &= v''B + 2v' + B' + B''v\end{aligned}$$

And now the original ode becomes

$$\begin{aligned} A(v''B + 2v'B' + B''v) + B(B'v + v'B) + CBv &= 0 \\ ABv'' + (2AB' + B^2)v' + (AB'' + BB' + CB)v &= 0 \end{aligned} \quad (1)$$

If the term $AB'' + BB' + CB$ is zero, then this method works and can be used to solve

$$ABv'' + (2AB' + B^2)v' = 0$$

By Using $u = v'$ which reduces the order of the above ode to one. The new ode is

$$ABu' + (2AB' + B^2)u = 0$$

The above ode is first order ode which is solved for u . Now a new ode $v' = u$ is solved for v as first order ode. Then the final solution is obtain from $y = Bv$.

This method works only if the term $AB'' + BB' + CB$ is zero. The given ODE shows that

$$\begin{aligned} A &= t \\ B &= -t - 1 \\ C &= 1 \\ F &= e^{2t}t^2 \end{aligned}$$

The above shows that for this ode

$$\begin{aligned} AB'' + BB' + CB &= (t)(0) + (-t - 1)(-1) + (1)(-t - 1) \\ &= 0 \end{aligned}$$

Hence the ode in v given in (1) now simplifies to

$$-t(t + 1)v'' + (t^2 + 1)v' = 0$$

Now by applying $v' = u$ the above becomes

$$-t(t + 1)u'(t) + (t^2 + 1)u(t) = 0$$

Which is now solved for u . In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{(t^2 + 1)u}{t(t + 1)} \end{aligned}$$

Where $f(t) = \frac{t^2+1}{t(t+1)}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= \frac{t^2 + 1}{t(t+1)} dt \\ \int \frac{1}{u} du &= \int \frac{t^2 + 1}{t(t+1)} dt \\ \ln(u) &= t + \ln(t) - 2 \ln(t+1) + c_1 \\ u &= e^{t+\ln(t)-2\ln(t+1)+c_1} \\ &= c_1 e^{t+\ln(t)-2\ln(t+1)}\end{aligned}$$

Which simplifies to

$$u(t) = \frac{c_1 e^{tt}}{(t+1)^2}$$

The ode for v now becomes

$$\begin{aligned}v' &= u \\ &= \frac{c_1 e^{tt}}{(t+1)^2}\end{aligned}$$

Which is now solved for v . Integrating both sides gives

$$\begin{aligned}v(t) &= \int \frac{c_1 e^{tt}}{(t+1)^2} dt \\ &= \frac{c_1 e^t}{t+1} + c_2\end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}y_h(t) &= Bv \\ &= (-t-1) \left(\frac{c_1 e^t}{t+1} + c_2 \right) \\ &= -c_1 e^t - c_2(t+1)\end{aligned}$$

And now the particular solution $y_p(t)$ will be found. The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = -t - 1$$

$$y_2 = e^t$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} -t - 1 & e^t \\ \frac{d}{dt}(-t - 1) & \frac{d}{dt}(e^t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} -t - 1 & e^t \\ -1 & e^t \end{vmatrix}$$

Therefore

$$W = (-t - 1)(e^t) - (e^t)(-1)$$

Which simplifies to

$$W = -te^t$$

Which simplifies to

$$W = -te^t$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{e^t e^{2t} t^2}{-t^2 e^t} dt$$

Which simplifies to

$$u_1 = - \int -e^{2t} dt$$

Hence

$$u_1 = \frac{e^{2t}}{2}$$

And Eq. (3) becomes

$$u_2 = \int \frac{(-t-1)e^{2t}t^2}{-t^2e^t} dt$$

Which simplifies to

$$u_2 = \int e^t(t+1) dt$$

Hence

$$u_2 = t e^t$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{(-t-1)e^{2t}}{2} + e^{2t}t$$

Which simplifies to

$$y_p(t) = \frac{e^{2t}(-1+t)}{2}$$

Hence the complete solution is

$$\begin{aligned} y(t) &= y_h + y_p \\ &= (-c_1e^t - c_2(t+1)) + \left(\frac{e^{2t}(-1+t)}{2} \right) \\ &= -c_1e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = -c_1e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2} \quad (1)$$

Verification of solutions

$$y = -c_1 e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2}$$

Verified OK.

10.21.2 Solving using Kovacic algorithm

Writing the ode as

$$ty'' + (-t-1)y' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= t \\ B &= -t-1 \\ C &= 1 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{t^2 - 2t + 3}{4t^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= t^2 - 2t + 3 \\ t &= 4t^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(\frac{t^2 - 2t + 3}{4t^2} \right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 506: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 2 \\ &= 0 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4t^2$. There is a pole at $t = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Therefore

$$L = [1, 2]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{1}{4} - \frac{1}{2t} + \frac{3}{4t^2}$$

For the pole at $t = 0$ let b be the coefficient of $\frac{1}{t^2}$ in the partial fractions decomposition of r given above. Therefore $b = \frac{3}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{3}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is $O_r(\infty) = 0$ then

$$v = \frac{-O_r(\infty)}{2} = \frac{0}{2} = 0$$

$[\sqrt{r}]_\infty$ is the sum of terms involving t^i for $0 \leq i \leq v$ in the Laurent series for \sqrt{r} at ∞ . Therefore

$$\begin{aligned} [\sqrt{r}]_\infty &= \sum_{i=0}^v a_i t^i \\ &= \sum_{i=0}^0 a_i t^i \end{aligned} \tag{8}$$

Let a be the coefficient of $t^v = t^0$ in the above sum. The Laurent series of \sqrt{r} at ∞ is

$$\sqrt{r} \approx \frac{1}{2} - \frac{1}{2t} + \frac{1}{2t^2} + \frac{1}{2t^3} + \frac{1}{4t^4} - \frac{1}{4t^5} - \frac{3}{4t^6} - \frac{3}{4t^7} + \dots \tag{9}$$

Comparing Eq. (9) with Eq. (8) shows that

$$a = \frac{1}{2}$$

From Eq. (9) the sum up to $v = 0$ gives

$$\begin{aligned} [\sqrt{r}]_\infty &= \sum_{i=0}^0 a_i t^i \\ &= \frac{1}{2} \end{aligned} \tag{10}$$

Now we need to find b , where b be the coefficient of $t^{v-1} = t^{-1} = \frac{1}{t}$ in r minus the coefficient of same term but in $([\sqrt{r}]_\infty)^2$ where $[\sqrt{r}]_\infty$ was found above in Eq (10). Hence

$$([\sqrt{r}]_\infty)^2 = \frac{1}{4}$$

This shows that the coefficient of $\frac{1}{t}$ in the above is 0. Now we need to find the coefficient of $\frac{1}{t}$ in r . How this is done depends on if $v = 0$ or not. Since $v = 0$ then starting from $r = \frac{s}{t}$ and doing long division in the form

$$r = Q + \frac{R}{t}$$

Where Q is the quotient and R is the remainder. Then the coefficient of $\frac{1}{t}$ in r will be the coefficient in R of the term in t of degree of t minus one, divided by the leading coefficient in t . Doing long division gives

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{t^2 - 2t + 3}{4t^2} \\ &= Q + \frac{R}{4t^2} \\ &= \left(\frac{1}{4}\right) + \left(\frac{-2t + 3}{4t^2}\right) \\ &= \frac{1}{4} + \frac{-2t + 3}{4t^2} \end{aligned}$$

Since the degree of t is 2, then we see that the coefficient of the term t in the remainder R is -2 . Dividing this by leading coefficient in t which is 4 gives $-\frac{1}{2}$. Now b can be found.

$$\begin{aligned} b &= \left(-\frac{1}{2}\right) - (0) \\ &= -\frac{1}{2} \end{aligned}$$

Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= \frac{1}{2} \\ \alpha_\infty^+ &= \frac{1}{2} \left(\frac{b}{a} - v\right) = \frac{1}{2} \left(\frac{-\frac{1}{2}}{\frac{1}{2}} - 0\right) = -\frac{1}{2} \\ \alpha_\infty^- &= \frac{1}{2} \left(-\frac{b}{a} - v\right) = \frac{1}{2} \left(-\frac{-\frac{1}{2}}{\frac{1}{2}} - 0\right) = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = \frac{t^2 - 2t + 3}{4t^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{3}{2}$	$-\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
0	$\frac{1}{2}$	$-\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_\infty^+ = -\frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_\infty^+ - (\alpha_{c_1}^-) \\ &= -\frac{1}{2} - \left(-\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty)[\sqrt{r}]_\infty$$

Substituting the above values in the above results in

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^-}{t - c_1} \right) + (+)[\sqrt{r}]_\infty \\ &= -\frac{1}{2t} + \left(\frac{1}{2}\right) \\ &= \frac{1}{2} - \frac{1}{2t} \\ &= \frac{-1 + t}{2t} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(\frac{1}{2} - \frac{1}{2t}\right)(0) + \left(\left(\frac{1}{2t^2}\right) + \left(\frac{1}{2} - \frac{1}{2t}\right)^2 - \left(\frac{t^2 - 2t + 3}{4t^2}\right)\right) = 0$$

$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \left(\frac{1}{2} - \frac{1}{2t}\right) dt} \\ &= \frac{e^{\frac{t}{2}}}{\sqrt{t}} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-t-1}{t} dt} \\ &= z_1 e^{\frac{t}{2} + \frac{\ln(t)}{2}} \\ &= z_1 \left(\sqrt{t} e^{\frac{t}{2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^t$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned}y_2 &= y_1 \int \frac{e^{\int -\frac{-t-1}{t} dt}}{(y_1)^2} dt \\&= y_1 \int \frac{e^{t+\ln(t)}}{(y_1)^2} dt \\&= y_1(-e^{-t}(t+1))\end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\&= c_1(e^t) + c_2(e^t(-e^{-t}(t+1)))\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$ty'' + (-t-1)y' + y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^t + (-t-1)c_2$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned}y_1 &= e^t \\y_2 &= -t-1\end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} e^t & -t - 1 \\ \frac{d}{dt}(e^t) & \frac{d}{dt}(-t - 1) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} e^t & -t - 1 \\ e^t & -1 \end{vmatrix}$$

Therefore

$$W = (e^t)(-1) - (-t - 1)(e^t)$$

Which simplifies to

$$W = t e^t$$

Which simplifies to

$$W = t e^t$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{(-t - 1) e^{2t} t^2}{t^2 e^t} dt$$

Which simplifies to

$$u_1 = - \int -e^t(t + 1) dt$$

Hence

$$u_1 = t e^t$$

And Eq. (3) becomes

$$u_2 = \int \frac{e^t e^{2t} t^2}{t^2 e^t} dt$$

Which simplifies to

$$u_2 = \int e^{2t} dt$$

Hence

$$u_2 = \frac{e^{2t}}{2}$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = \frac{(-t-1)e^{2t}}{2} + e^{2t}t$$

Which simplifies to

$$y_p(t) = \frac{e^{2t}(-1+t)}{2}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 e^t + (-t-1)c_2) + \left(\frac{e^{2t}(-1+t)}{2} \right) \end{aligned}$$

Which simplifies to

$$y = c_1 e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2}$$

Summary

The solution(s) found are the following

$$y = c_1 e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2} \tag{1}$$

Verification of solutions

$$y = c_1 e^t - c_2(t+1) + \frac{e^{2t}(-1+t)}{2}$$

Verified OK.

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    checking if the LODE is of Euler type
    trying a symmetry of the form [xi=0, eta=F(x)]
    checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Reducible group (found an exponential solution)
    Reducible group (found another exponential solution)
<- Kovacics algorithm successful
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 23

```
dsolve(t*diff(y(t),t$2)-(1+t)*diff(y(t),t)+y(t) = t^2*exp(2*t),y(t), singsol=all)
```

$$y(t) = (t + 1) c_2 + e^t c_1 + \frac{(t - 1) e^{2t}}{2}$$

✓ Solution by Mathematica

Time used: 0.033 (sec). Leaf size: 31

```
DSolve[t*y''[t]-(1+t)*y'[t]+y[t] ==t^2*Exp[2*t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow \frac{1}{2} e^{2t} (t - 1) + c_1 e^t - c_2 (t + 1)$$

10.22 problem 32

- 10.22.1 Solving as second order change of variable on y method 2 ode . 3018
- 10.22.2 Solving as second order ode non constant coeff transformation
on B ode 3023
- 10.22.3 Solving using Kovacic algorithm 3028

Internal problem ID [704]

Internal file name [OUTPUT/704_Sunday_June_05_2022_01_47_21_AM_33934857/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, section 3.6, Variation of Parameters.
page 190

Problem number: 32.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_change_of_variable_on_y_method_2", "second_order_ode_non_constant_coeff_transformation_on_B"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(1 - t)y'' + ty' - y = 2(-1 + t)e^{-t}$$

10.22.1 Solving as second order change of variable on y method 2 ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1 - t$, $B = t$, $C = -1$, $f(t) = 2(-1 + t)e^{-t}$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. Solving for y_h from

$$(1 - t)y'' + ty' - y = 0$$

In normal form the ode

$$(1 - t)y'' + ty' - y = 0 \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = 0 \quad (2)$$

Where

$$p(t) = -\frac{t}{-1+t}$$
$$q(t) = \frac{1}{-1+t}$$

Applying change of variables on the dependent variable $y = v(t)t^n$ to (2) gives the following ode where the dependent variables is $v(t)$ and not y .

$$v''(t) + \left(\frac{2n}{t} + p\right)v'(t) + \left(\frac{n(n-1)}{t^2} + \frac{np}{t} + q\right)v(t) = 0 \quad (3)$$

Let the coefficient of $v(t)$ above be zero. Hence

$$\frac{n(n-1)}{t^2} + \frac{np}{t} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(t)$ and $q(t)$ into (4) gives

$$\frac{n(n-1)}{t^2} - \frac{n}{-1+t} + \frac{1}{-1+t} = 0 \quad (5)$$

Solving (5) for n gives

$$n = 1 \quad (6)$$

Substituting this value in (3) gives

$$v''(t) + \left(\frac{2}{t} - \frac{t}{-1+t}\right)v'(t) = 0$$
$$v''(t) + \left(\frac{2}{t} - \frac{t}{-1+t}\right)v'(t) = 0 \quad (7)$$

Using the substitution

$$u(t) = v'(t)$$

Then (7) becomes

$$u'(t) + \left(\frac{2}{t} - \frac{t}{-1+t} \right) u(t) = 0 \quad (8)$$

The above is now solved for $u(t)$. In canonical form the ODE is

$$\begin{aligned} u' &= F(t, u) \\ &= f(t)g(u) \\ &= \frac{u(t^2 - 2t + 2)}{t(-1+t)} \end{aligned}$$

Where $f(t) = \frac{t^2 - 2t + 2}{t(-1+t)}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned} \frac{1}{u} du &= \frac{t^2 - 2t + 2}{t(-1+t)} dt \\ \int \frac{1}{u} du &= \int \frac{t^2 - 2t + 2}{t(-1+t)} dt \\ \ln(u) &= t - 2 \ln(t) + \ln(-1+t) + c_1 \\ u &= e^{t-2\ln(t)+\ln(-1+t)+c_1} \\ &= c_1 e^{t-2\ln(t)+\ln(-1+t)} \end{aligned}$$

Which simplifies to

$$u(t) = c_1 \left(-\frac{e^t}{t^2} + \frac{e^t}{t} \right)$$

Now that $u(t)$ is known, then

$$\begin{aligned} v'(t) &= u(t) \\ v(t) &= \int u(t) dt + c_2 \\ &= \frac{e^t c_1}{t} + c_2 \end{aligned}$$

Hence

$$\begin{aligned} y &= v(t) t^n \\ &= \left(\frac{e^t c_1}{t} + c_2 \right) t \\ &= c_1 e^t + c_2 t \end{aligned}$$

Now the particular solution to this ODE is found

$$(1 - t)y'' + ty' - y = 2(-1 + t)e^{-t}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1y_1 + u_2y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = t$$

$$y_2 = e^t$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t & e^t \\ \frac{d}{dt}(t) & \frac{d}{dt}(e^t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t & e^t \\ 1 & e^t \end{vmatrix}$$

Therefore

$$W = (t)(e^t) - (e^t)(1)$$

Which simplifies to

$$W = t e^t - e^t$$

Which simplifies to

$$W = e^t(-1 + t)$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{2 e^t(-1 + t) e^{-t}}{(1 - t) e^t (-1 + t)} dt$$

Which simplifies to

$$u_1 = - \int -\frac{2 e^{-t}}{-1 + t} dt$$

Hence

$$u_1 = -2 e^{-1} \expIntegral_1(-1 + t)$$

And Eq. (3) becomes

$$u_2 = \int \frac{2 e^{-t} t(-1 + t)}{(1 - t) e^t (-1 + t)} dt$$

Which simplifies to

$$u_2 = \int -\frac{2t e^{-2t}}{-1 + t} dt$$

Hence

$$u_2 = e^{-2t} + 2 e^{-2} \expIntegral_1(2t - 2)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -2 e^{-1} \expIntegral_1(-1 + t) t + (e^{-2t} + 2 e^{-2} \expIntegral_1(2t - 2)) e^t$$

Which simplifies to

$$y_p(t) = -2 e^{-1} \expIntegral_1(-1 + t) t + 2 \expIntegral_1(2t - 2) e^{t-2} + e^{-t}$$

Therefore the general solution is

$$\begin{aligned}
 y &= y_h + y_p \\
 &= \left(\left(\frac{e^t c_1}{t} + c_2 \right) t \right) + (-2 e^{-1} \exp \text{Integral}_1 (-1+t) t + 2 \exp \text{Integral}_1 (2t-2) e^{t-2} + e^{-t}) \\
 &= -2 e^{-1} \exp \text{Integral}_1 (-1+t) t + 2 \exp \text{Integral}_1 (2t-2) e^{t-2} + e^{-t} + \left(\frac{e^t c_1}{t} + c_2 \right) t
 \end{aligned}$$

Which simplifies to

$$y = -2 e^{-1} \exp \text{Integral}_1 (-1+t) t + 2 \exp \text{Integral}_1 (2t-2) e^{t-2} + c_1 e^t + c_2 t + e^{-t}$$

Summary

The solution(s) found are the following

$$y = -2 e^{-1} \exp \text{Integral}_1 (-1+t) t + 2 \exp \text{Integral}_1 (2t-2) e^{t-2} + c_1 e^t + c_2 t + e^{-t} (1)$$

Verification of solutions

$$y = -2 e^{-1} \exp \text{Integral}_1 (-1+t) t + 2 \exp \text{Integral}_1 (2t-2) e^{t-2} + c_1 e^t + c_2 t + e^{-t}$$

Verified OK.

10.22.2 Solving as second order ode non constant coeff transformation on B ode

Given an ode of the form

$$Ay'' + By' + Cy = F(t)$$

This method reduces the order ode the ODE by one by applying the transformation

$$y = Bv$$

This results in

$$\begin{aligned}
 y' &= B'v + v'B \\
 y'' &= B''v + B'v' + v''B + v'B' \\
 &= v''B + 2v' + B' + B''v
 \end{aligned}$$

And now the original ode becomes

$$\begin{aligned}
 A(v''B + 2v'B' + B''v) + B(B'v + v'B) + CBv &= 0 \\
 ABv'' + (2AB' + B^2)v' + (AB'' + BB' + CB)v &= 0
 \end{aligned} \tag{1}$$

If the term $AB'' + BB' + CB$ is zero, then this method works and can be used to solve

$$ABv'' + (2AB' + B^2)v' = 0$$

By Using $u = v'$ which reduces the order of the above ode to one. The new ode is

$$ABu' + (2AB' + B^2)u = 0$$

The above ode is first order ode which is solved for u . Now a new ode $v' = u$ is solved for v as first order ode. Then the final solution is obtain from $y = Bv$.

This method works only if the term $AB'' + BB' + CB$ is zero. The given ODE shows that

$$\begin{aligned}A &= 1 - t \\B &= t \\C &= -1 \\F &= 2(-1 + t)e^{-t}\end{aligned}$$

The above shows that for this ode

$$\begin{aligned}AB'' + BB' + CB &= (1 - t)(0) + (t)(1) + (-1)(t) \\&= 0\end{aligned}$$

Hence the ode in v given in (1) now simplifies to

$$-t(-1 + t)v'' + (t^2 - 2t + 2)v' = 0$$

Now by applying $v' = u$ the above becomes

$$(-t^2 + t)u'(t) + (t^2 - 2t + 2)u(t) = 0$$

Which is now solved for u . In canonical form the ODE is

$$\begin{aligned}u' &= F(t, u) \\&= f(t)g(u) \\&= \frac{u(t^2 - 2t + 2)}{t(-1 + t)}\end{aligned}$$

Where $f(t) = \frac{t^2-2t+2}{t(-1+t)}$ and $g(u) = u$. Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= \frac{t^2 - 2t + 2}{t(-1+t)} dt \\ \int \frac{1}{u} du &= \int \frac{t^2 - 2t + 2}{t(-1+t)} dt \\ \ln(u) &= t - 2 \ln(t) + \ln(-1+t) + c_1 \\ u &= e^{t-2\ln(t)+\ln(-1+t)+c_1} \\ &= c_1 e^{t-2\ln(t)+\ln(-1+t)}\end{aligned}$$

Which simplifies to

$$u(t) = c_1 \left(-\frac{e^t}{t^2} + \frac{e^t}{t} \right)$$

The ode for v now becomes

$$\begin{aligned}v' &= u \\ &= c_1 \left(-\frac{e^t}{t^2} + \frac{e^t}{t} \right)\end{aligned}$$

Which is now solved for v . Integrating both sides gives

$$\begin{aligned}v(t) &= \int \frac{c_1 e^t (-1+t)}{t^2} dt \\ &= \frac{e^t c_1}{t} + c_2\end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}y_h(t) &= Bv \\ &= (t) \left(\frac{e^t c_1}{t} + c_2 \right) \\ &= c_1 e^t + c_2 t\end{aligned}$$

And now the particular solution $y_p(t)$ will be found. The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = t$$

$$y_2 = e^t$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \quad (3)$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} t & e^t \\ \frac{d}{dt}(t) & \frac{d}{dt}(e^t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} t & e^t \\ 1 & e^t \end{vmatrix}$$

Therefore

$$W = (t) (e^t) - (e^t) (1)$$

Which simplifies to

$$W = t e^t - e^t$$

Which simplifies to

$$W = e^t(-1 + t)$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{2 e^t(-1 + t) e^{-t}}{(1 - t) e^t (-1 + t)} dt$$

Which simplifies to

$$u_1 = - \int -\frac{2e^{-t}}{-1+t} dt$$

Hence

$$u_1 = -2e^{-1} \exp\text{Integral}_1(-1+t)$$

And Eq. (3) becomes

$$u_2 = \int \frac{2e^{-t}t(-1+t)}{(1-t)e^t(-1+t)} dt$$

Which simplifies to

$$u_2 = \int -\frac{2te^{-2t}}{-1+t} dt$$

Hence

$$u_2 = e^{-2t} + 2e^{-2} \exp\text{Integral}_1(2t-2)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -2e^{-1} \exp\text{Integral}_1(-1+t)t + (e^{-2t} + 2e^{-2} \exp\text{Integral}_1(2t-2))e^t$$

Which simplifies to

$$y_p(t) = -2e^{-1} \exp\text{Integral}_1(-1+t)t + 2 \exp\text{Integral}_1(2t-2)e^{t-2} + e^{-t}$$

Hence the complete solution is

$$\begin{aligned} y(t) &= y_h + y_p \\ &= (c_1e^t + c_2t) + (-2e^{-1} \exp\text{Integral}_1(-1+t)t + 2 \exp\text{Integral}_1(2t-2)e^{t-2} + e^{-t}) \\ &= -2e^{-1} \exp\text{Integral}_1(-1+t)t + 2 \exp\text{Integral}_1(2t-2)e^{t-2} + c_1e^t + c_2t + e^{-t} \end{aligned}$$

Summary

The solution(s) found are the following

$$y = -2e^{-1} \exp\text{Integral}_1(-1+t)t + 2 \exp\text{Integral}_1(2t-2)e^{t-2} + c_1e^t + c_2t + e^{-t}(1)$$

Verification of solutions

$$y = -2e^{-1} \exp\text{Integral}_1(-1+t)t + 2 \exp\text{Integral}_1(2t-2)e^{t-2} + c_1e^t + c_2t + e^{-t}$$

Verified OK.

10.22.3 Solving using Kovacic algorithm

Writing the ode as

$$(1 - t)y'' + ty' - y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 - t \\ B &= t \\ C &= -1 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{t^2 - 4t + 6}{4(-1 + t)^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= t^2 - 4t + 6 \\ t &= 4(-1 + t)^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = \left(\frac{t^2 - 4t + 6}{4(-1 + t)^2} \right) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 507: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 2 \\ &= 0 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4(-1 + t)^2$. There is a pole at $t = 1$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Therefore

$$L = [1, 2]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = \frac{1}{4} - \frac{1}{2(-1+t)} + \frac{3}{4(-1+t)^2}$$

For the pole at $t = 1$ let b be the coefficient of $\frac{1}{(-1+t)^2}$ in the partial fractions decomposition of r given above. Therefore $b = \frac{3}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{3}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = -\frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is $O_r(\infty) = 0$ then

$$v = \frac{-O_r(\infty)}{2} = \frac{0}{2} = 0$$

$[\sqrt{r}]_\infty$ is the sum of terms involving t^i for $0 \leq i \leq v$ in the Laurent series for \sqrt{r} at ∞ . Therefore

$$\begin{aligned} [\sqrt{r}]_\infty &= \sum_{i=0}^v a_i t^i \\ &= \sum_{i=0}^0 a_i t^i \end{aligned} \tag{8}$$

Let a be the coefficient of $t^v = t^0$ in the above sum. The Laurent series of \sqrt{r} at ∞ is

$$\sqrt{r} \approx \frac{1}{2} - \frac{1}{2t} + \frac{1}{t^3} + \frac{11}{4t^4} + \frac{21}{4t^5} + \frac{15}{2t^6} + \frac{6}{t^7} - \frac{117}{16t^8} + \dots \tag{9}$$

Comparing Eq. (9) with Eq. (8) shows that

$$a = \frac{1}{2}$$

From Eq. (9) the sum up to $v = 0$ gives

$$\begin{aligned} [\sqrt{r}]_\infty &= \sum_{i=0}^0 a_i t^i \\ &= \frac{1}{2} \end{aligned} \tag{10}$$

Now we need to find b , where b be the coefficient of $t^{v-1} = t^{-1} = \frac{1}{t}$ in r minus the coefficient of same term but in $([\sqrt{r}]_\infty)^2$ where $[\sqrt{r}]_\infty$ was found above in Eq (10). Hence

$$([\sqrt{r}]_\infty)^2 = \frac{1}{4}$$

This shows that the coefficient of $\frac{1}{t}$ in the above is 0. Now we need to find the coefficient of $\frac{1}{t}$ in r . How this is done depends on if $v = 0$ or not. Since $v = 0$ then starting from $r = \frac{s}{t}$ and doing long division in the form

$$r = Q + \frac{R}{t}$$

Where Q is the quotient and R is the remainder. Then the coefficient of $\frac{1}{t}$ in r will be the coefficient in R of the term in t of degree of t minus one, divided by the leading coefficient in t . Doing long division gives

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{t^2 - 4t + 6}{4t^2 - 8t + 4} \\ &= Q + \frac{R}{4t^2 - 8t + 4} \\ &= \left(\frac{1}{4}\right) + \left(\frac{-2t + 5}{4t^2 - 8t + 4}\right) \\ &= \frac{1}{4} + \frac{-2t + 5}{4t^2 - 8t + 4} \end{aligned}$$

Since the degree of t is 2, then we see that the coefficient of the term t in the remainder R is -2 . Dividing this by leading coefficient in t which is 4 gives $-\frac{1}{2}$. Now b can be found.

$$\begin{aligned} b &= \left(-\frac{1}{2}\right) - (0) \\ &= -\frac{1}{2} \end{aligned}$$

Hence

$$\begin{aligned} [\sqrt{r}]_{\infty} &= \frac{1}{2} \\ \alpha_{\infty}^+ &= \frac{1}{2} \left(\frac{b}{a} - v\right) = \frac{1}{2} \left(\frac{-\frac{1}{2}}{\frac{1}{2}} - 0\right) = -\frac{1}{2} \\ \alpha_{\infty}^- &= \frac{1}{2} \left(-\frac{b}{a} - v\right) = \frac{1}{2} \left(-\frac{-\frac{1}{2}}{\frac{1}{2}} - 0\right) = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = \frac{t^2 - 4t + 6}{4(-1 + t)^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
1	2	0	$\frac{3}{2}$	$-\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
0	$\frac{1}{2}$	$-\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_\infty^+ = -\frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_\infty^+ - (\alpha_{c_1}^-) \\ &= -\frac{1}{2} - \left(-\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c) [\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{t - c} \right) + s(\infty) [\sqrt{r}]_\infty$$

Substituting the above values in the above results in

$$\begin{aligned} \omega &= \left((-) [\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^-}{t - c_1} \right) + (+) [\sqrt{r}]_\infty \\ &= -\frac{1}{2(-1+t)} + \left(\frac{1}{2}\right) \\ &= -\frac{1}{2(-1+t)} + \frac{1}{2} \\ &= \frac{t-2}{2t-2} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(t)$ of degree $d = 0$ to solve the ode. The polynomial $p(t)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \quad (1A)$$

Let

$$p(t) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(-\frac{1}{2(-1+t)} + \frac{1}{2}\right) (0) + \left(\left(\frac{1}{2(-1+t)}\right)^2 + \left(-\frac{1}{2(-1+t)} + \frac{1}{2}\right)^2 - \left(\frac{t^2 - 4t + 6}{4(-1+t)^2}\right)\right) = 0$$
$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(t) &= pe^{\int \omega dt} \\ &= e^{\int \left(-\frac{1}{2(-1+t)} + \frac{1}{2}\right) dt} \\ &= \frac{e^{\frac{t}{2}}}{\sqrt{-1+t}} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{t}{1-t} dt} \\ &= z_1 e^{\frac{t}{2} + \frac{\ln(-1+t)}{2}} \\ &= z_1 \left(\sqrt{-1+t} e^{\frac{t}{2}}\right) \end{aligned}$$

Which simplifies to

$$y_1 = e^t$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{t}{1-t} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{t + \ln(-1+t)}}{(y_1)^2} dt \\ &= y_1 (-t e^{-t}) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^t) + c_2 (e^t (-t e^{-t}))\end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$(1 - t)y'' + ty' - y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^t - c_2 t$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on t as well. Let

$$y_p(t) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned}y_1 &= e^t \\ y_2 &= -t\end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(t)}{aW(t)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(t)}{aW(t)} \tag{3}$$

Where $W(t)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} e^t & -t \\ \frac{d}{dt}(e^t) & \frac{d}{dt}(-t) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} e^t & -t \\ e^t & -1 \end{vmatrix}$$

Therefore

$$W = (e^t)(-1) - (-t)(e^t)$$

Which simplifies to

$$W = t e^t - e^t$$

Which simplifies to

$$W = e^t(-1 + t)$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{-2 e^{-t} t(-1 + t)}{(1 - t) e^t (-1 + t)} dt$$

Which simplifies to

$$u_1 = - \int \frac{2t e^{-2t}}{-1 + t} dt$$

Hence

$$u_1 = e^{-2t} + 2 e^{-2} \expIntegral_1(2t - 2)$$

And Eq. (3) becomes

$$u_2 = \int \frac{2 e^t(-1 + t) e^{-t}}{(1 - t) e^t (-1 + t)} dt$$

Which simplifies to

$$u_2 = \int -\frac{2e^{-t}}{-1+t} dt$$

Hence

$$u_2 = 2e^{-1} \expIntegral_1(-1+t)$$

Therefore the particular solution, from equation (1) is

$$y_p(t) = -2e^{-1} \expIntegral_1(-1+t)t + (e^{-2t} + 2e^{-2} \expIntegral_1(2t-2))e^t$$

Which simplifies to

$$y_p(t) = -2e^{-1} \expIntegral_1(-1+t)t + 2 \expIntegral_1(2t-2)e^{t-2} + e^{-t}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1e^t - c_2t) + (-2e^{-1} \expIntegral_1(-1+t)t + 2 \expIntegral_1(2t-2)e^{t-2} + e^{-t}) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = c_1e^t - c_2t - 2e^{-1} \expIntegral_1(-1+t)t + 2 \expIntegral_1(2t-2)e^{t-2} + e^{-t} \quad (1)$$

Verification of solutions

$$y = c_1e^t - c_2t - 2e^{-1} \expIntegral_1(-1+t)t + 2 \expIntegral_1(2t-2)e^{t-2} + e^{-t}$$

Verified OK.

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    checking if the LODE is of Euler type
    trying a symmetry of the form [xi=0, eta=F(x)]
    checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Reducible group (found an exponential solution)
    Reducible group (found another exponential solution)
<- Kovacics algorithm successful
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 39

```
dsolve((1-t)*diff(y(t),t$2)+t*diff(y(t),t)-y(t) = 2*(t-1)*exp(-t),y(t), singsol=all)
```

$$y(t) = -2e^{-1} \operatorname{ExpIntegral}_1(t-1)t + 2 \operatorname{ExpIntegral}_1(2t-2)e^{t-2} + e^t c_1 + c_2 t + e^{-t}$$

✓ Solution by Mathematica

Time used: 0.187 (sec). Leaf size: 47

```
DSolve[(1-t)*y'[t]+t*y'[t]-y[t] ==2*(t-1)*Exp[-t],y[t],t,IncludeSingularSolutions -> True]
```

$$y(t) \rightarrow -2e^{t-2} \operatorname{ExpIntegralEi}(2-2t) + \frac{2t \operatorname{ExpIntegralEi}(1-t)}{e} + e^{-t} + c_1 e^t - c_2 t$$

**11 Chapter 3, Second order linear equations, 3.7
Mechanical and Electrical Vibrations. page 203**

11.1 problem 28 3039
11.2 problem 29 3049

11.1 problem 28

11.1.1 Solving as second order linear constant coeff ode	3039
11.1.2 Solving as second order ode can be made integrable ode	3041
11.1.3 Solving using Kovacic algorithm	3043
11.1.4 Maple step by step solution	3047

Internal problem ID [705]

Internal file name [OUTPUT/705_Sunday_June_05_2022_01_47_23_AM_80288768/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.7 Mechanical and Electrical Vibrations. page 203

Problem number: 28.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "kovacic", "second_order_linear_constant_coeff", "second_order_ode_can_be_made_integrable"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$u'' + 2u = 0$$

11.1.1 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Au''(t) + Bu'(t) + Cu(t) = 0$$

Where in the above $A = 1, B = 0, C = 2$. Let the solution be $u = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + 2e^{\lambda t} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + 2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(2)} \\ &= \pm i\sqrt{2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= +i\sqrt{2} \\ \lambda_2 &= -i\sqrt{2} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= i\sqrt{2} \\ \lambda_2 &= -i\sqrt{2} \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = \sqrt{2}$. Therefore the final solution, when using Euler relation, can be written as

$$u = e^{\alpha t} (c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$u = e^0 (c_1 \cos(\sqrt{2}t) + c_2 \sin(\sqrt{2}t))$$

Or

$$u = c_1 \cos(\sqrt{2}t) + c_2 \sin(\sqrt{2}t)$$

Summary

The solution(s) found are the following

$$u = c_1 \cos(\sqrt{2}t) + c_2 \sin(\sqrt{2}t) \tag{1}$$

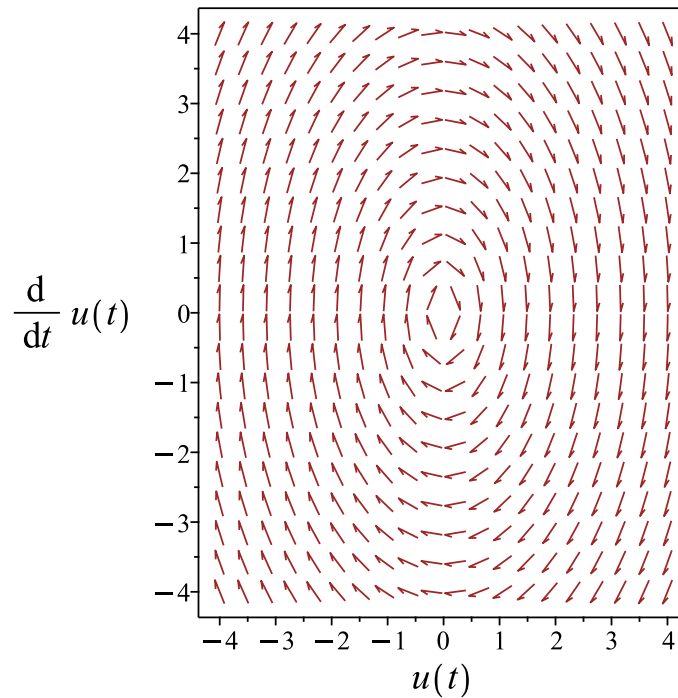


Figure 491: Slope field plot

Verification of solutions

$$u = c_1 \cos(\sqrt{2}t) + c_2 \sin(\sqrt{2}t)$$

Verified OK.

11.1.2 Solving as second order ode can be made integrable ode

Multiplying the ode by u' gives

$$u'u'' + 2u'u = 0$$

Integrating the above w.r.t t gives

$$\int (u'u'' + 2u'u) dt = 0$$

$$\frac{u'^2}{2} + u^2 = c_2$$

Which is now solved for u . Solving the given ode for u' results in 2 differential equations to solve. Each one of these will generate a solution. The equations generated are

$$u' = \sqrt{-2u^2 + 2c_1} \tag{1}$$

$$u' = -\sqrt{-2u^2 + 2c_1} \tag{2}$$

Now each one of the above ODE is solved.

Solving equation (1)

Integrating both sides gives

$$\int \frac{1}{\sqrt{-2u^2 + 2c_1}} du = \int dt$$
$$\frac{\sqrt{2} \arctan\left(\frac{\sqrt{2}u}{\sqrt{-2u^2+2c_1}}\right)}{2} = t + c_2$$

Solving equation (2)

Integrating both sides gives

$$\int -\frac{1}{\sqrt{-2u^2 + 2c_1}} du = \int dt$$
$$-\frac{\sqrt{2} \arctan\left(\frac{\sqrt{2}u}{\sqrt{-2u^2+2c_1}}\right)}{2} = t + c_3$$

Summary

The solution(s) found are the following

$$\frac{\sqrt{2} \arctan\left(\frac{\sqrt{2}u}{\sqrt{-2u^2+2c_1}}\right)}{2} = t + c_2 \quad (1)$$

$$-\frac{\sqrt{2} \arctan\left(\frac{\sqrt{2}u}{\sqrt{-2u^2+2c_1}}\right)}{2} = t + c_3 \quad (2)$$

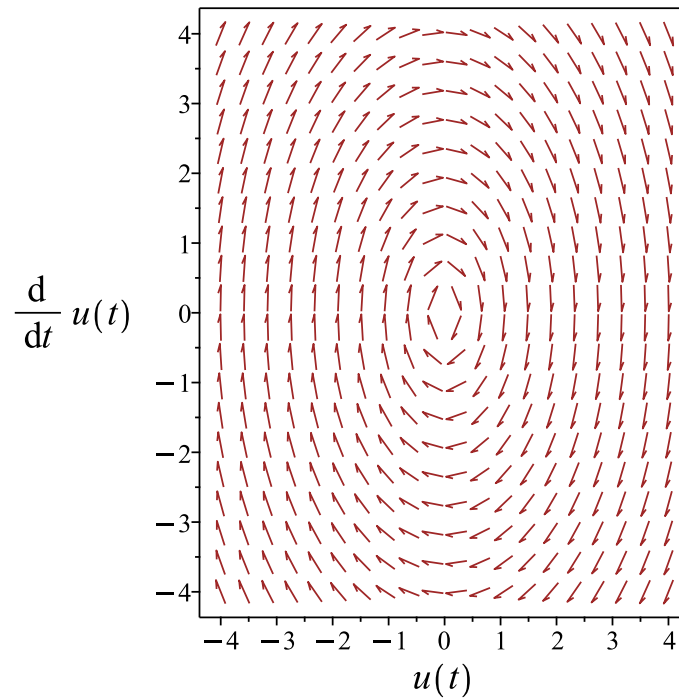


Figure 492: Slope field plot

Verification of solutions

$$\frac{\sqrt{2} \arctan\left(\frac{\sqrt{2}u}{\sqrt{-2u^2+2c_1}}\right)}{2} = t + c_2$$

Verified OK.

$$-\frac{\sqrt{2} \arctan\left(\frac{\sqrt{2}u}{\sqrt{-2u^2+2c_1}}\right)}{2} = t + c_3$$

Verified OK.

11.1.3 Solving using Kovacic algorithm

Writing the ode as

$$u'' + 2u = 0 \tag{1}$$

$$Au'' + Bu' + Cu = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= 1 \\B &= 0 \\C &= 2\end{aligned}\tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ue^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t)\tag{4}$$

Where r is given by

$$\begin{aligned}r &= \frac{s}{t} \\&= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}\end{aligned}\tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-2}{1}\tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned}s &= -2 \\t &= 1\end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = -2z(t)\tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then u is found using the inverse transformation

$$u = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 508: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -2$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos(\sqrt{2}t)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in u is found from

$$u_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned}
 u_1 &= z_1 \\
 &= \cos(\sqrt{2}t)
 \end{aligned}$$

Which simplifies to

$$u_1 = \cos(\sqrt{2}t)$$

The second solution u_2 to the original ode is found using reduction of order

$$u_2 = u_1 \int \frac{e^{\int -\frac{B}{A} dt}}{u_1^2} dt$$

Since $B = 0$ then the above becomes

$$\begin{aligned} u_2 &= u_1 \int \frac{1}{u_1^2} dt \\ &= \cos(\sqrt{2}t) \int \frac{1}{\cos^2(\sqrt{2}t)} dt \\ &= \cos(\sqrt{2}t) \left(\frac{\sqrt{2} \tan(\sqrt{2}t)}{2} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} u &= c_1 u_1 + c_2 u_2 \\ &= c_1 \left(\cos(\sqrt{2}t) \right) + c_2 \left(\cos(\sqrt{2}t) \left(\frac{\sqrt{2} \tan(\sqrt{2}t)}{2} \right) \right) \end{aligned}$$

Summary

The solution(s) found are the following

$$u = c_1 \cos(\sqrt{2}t) + \frac{c_2 \sqrt{2} \sin(\sqrt{2}t)}{2} \tag{1}$$

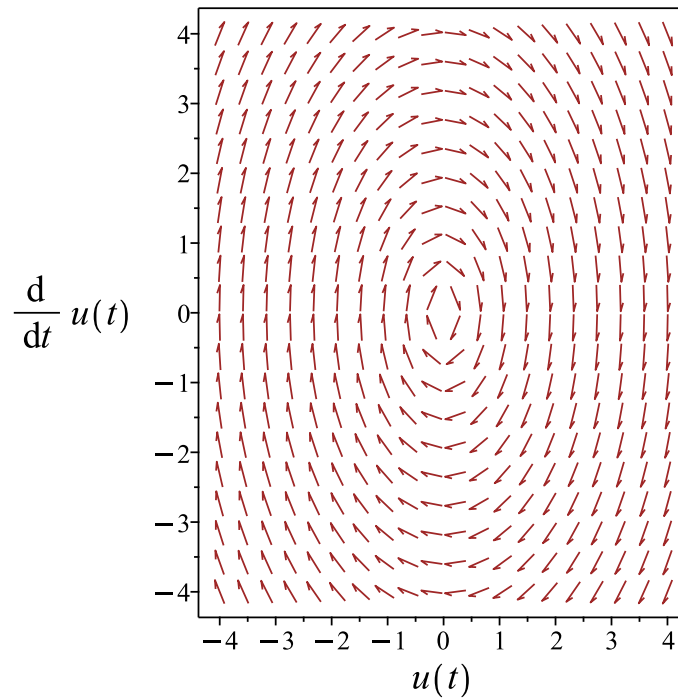


Figure 493: Slope field plot

Verification of solutions

$$u = c_1 \cos(\sqrt{2}t) + \frac{c_2 \sqrt{2} \sin(\sqrt{2}t)}{2}$$

Verified OK.

11.1.4 Maple step by step solution

Let's solve

$$u'' + 2u = 0$$

- Highest derivative means the order of the ODE is 2
- u''
- Characteristic polynomial of ODE
- $r^2 + 2 = 0$
- Use quadratic formula to solve for r
- $r = \frac{0 \pm \sqrt{-8}}{2}$
- Roots of the characteristic polynomial

$$r = (-I\sqrt{2}, I\sqrt{2})$$

- 1st solution of the ODE

$$u_1(t) = \cos(\sqrt{2}t)$$

- 2nd solution of the ODE

$$u_2(t) = \sin(\sqrt{2}t)$$

- General solution of the ODE

$$u = c_1 u_1(t) + c_2 u_2(t)$$

- Substitute in solutions

$$u = c_1 \cos(\sqrt{2}t) + c_2 \sin(\sqrt{2}t)$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 21

```
dsolve(diff(u(t),t$2)+2*u(t) = 0,u(t), singsol=all)
```

$$u(t) = c_1 \sin(t\sqrt{2}) + c_2 \cos(t\sqrt{2})$$

✓ Solution by Mathematica

Time used: 0.017 (sec). Leaf size: 28

```
DSolve[u''[t]+2*u[t] ==0,u[t],t,IncludeSingularSolutions -> True]
```

$$u(t) \rightarrow c_1 \cos(\sqrt{2}t) + c_2 \sin(\sqrt{2}t)$$

11.2 problem 29

11.2.1 Existence and uniqueness analysis	3049
11.2.2 Solving as second order linear constant coeff ode	3050
11.2.3 Solving using Kovacic algorithm	3053
11.2.4 Maple step by step solution	3057

Internal problem ID [706]

Internal file name [OUTPUT/706_Sunday_June_05_2022_01_47_24_AM_3139852/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.7 Mechanical and Electrical Vibrations. page 203

Problem number: 29.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _missing_x]]
```

$$u'' + \frac{u'}{4} + 2u = 0$$

With initial conditions

$$[u(0) = 0, u'(0) = 2]$$

11.2.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$u'' + p(t)u' + q(t)u = F$$

Where here

$$p(t) = \frac{1}{4}$$

$$q(t) = 2$$

$$F = 0$$

Hence the ode is

$$u'' + \frac{u'}{4} + 2u = 0$$

The domain of $p(t) = \frac{1}{4}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = 2$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

11.2.2 Solving as second order linear constant coeff ode

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Au''(t) + Bu'(t) + Cu(t) = 0$$

Where in the above $A = 1, B = \frac{1}{4}, C = 2$. Let the solution be $u = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + \frac{\lambda e^{\lambda t}}{4} + 2e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + \frac{1}{4}\lambda + 2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = \frac{1}{4}, C = 2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-\frac{1}{4}}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{\frac{1^2}{4} - (4)(1)(2)} \\ &= -\frac{1}{8} \pm \frac{i\sqrt{127}}{8} \end{aligned}$$

Hence

$$\lambda_1 = -\frac{1}{8} + \frac{i\sqrt{127}}{8}$$
$$\lambda_2 = -\frac{1}{8} - \frac{i\sqrt{127}}{8}$$

Which simplifies to

$$\lambda_1 = -\frac{1}{8} + \frac{i\sqrt{127}}{8}$$
$$\lambda_2 = -\frac{1}{8} - \frac{i\sqrt{127}}{8}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -\frac{1}{8}$ and $\beta = \frac{\sqrt{127}}{8}$. Therefore the final solution, when using Euler relation, can be written as

$$u = e^{\alpha t} (c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$u = e^{-\frac{t}{8}} \left(c_1 \cos \left(\frac{\sqrt{127} t}{8} \right) + c_2 \sin \left(\frac{\sqrt{127} t}{8} \right) \right)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = e^{-\frac{t}{8}} \left(c_1 \cos \left(\frac{\sqrt{127} t}{8} \right) + c_2 \sin \left(\frac{\sqrt{127} t}{8} \right) \right) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 0$ and $t = 0$ in the above gives

$$0 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$u' = -\frac{e^{-\frac{t}{8}} \left(c_1 \cos \left(\frac{\sqrt{127} t}{8} \right) + c_2 \sin \left(\frac{\sqrt{127} t}{8} \right) \right)}{8} + e^{-\frac{t}{8}} \left(-\frac{c_1 \sqrt{127} \sin \left(\frac{\sqrt{127} t}{8} \right)}{8} + \frac{c_2 \sqrt{127} \cos \left(\frac{\sqrt{127} t}{8} \right)}{8} \right)$$

substituting $u' = 2$ and $t = 0$ in the above gives

$$2 = -\frac{c_1}{8} + \frac{\sqrt{127} c_2}{8} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 0$$
$$c_2 = \frac{16\sqrt{127}}{127}$$

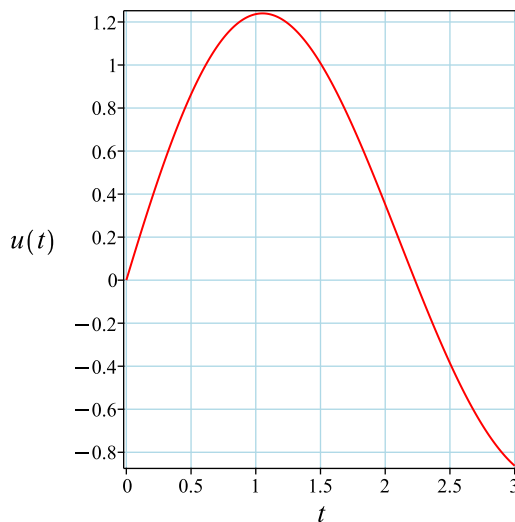
Substituting these values back in above solution results in

$$u = \frac{16\sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127} t}{8}\right)}{127}$$

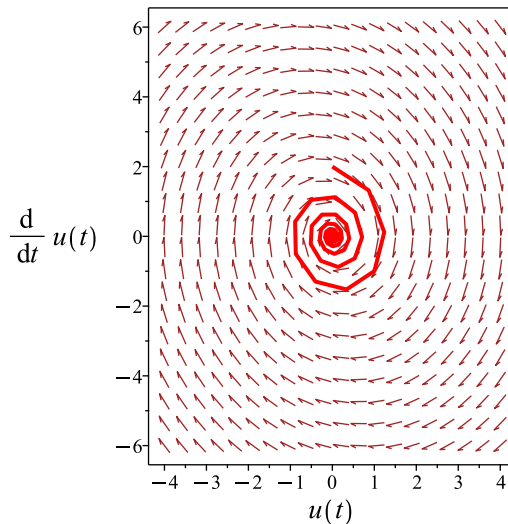
Summary

The solution(s) found are the following

$$u = \frac{16\sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127} t}{8}\right)}{127} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = \frac{16\sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127} t}{8}\right)}{127}$$

Verified OK.

11.2.3 Solving using Kovacic algorithm

Writing the ode as

$$u'' + \frac{u'}{4} + 2u = 0 \quad (1)$$

$$Au'' + Bu' + Cu = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= \frac{1}{4} \\ C &= 2 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ue^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-127}{64} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -127 \\ t &= 64 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = -\frac{127z(t)}{64} \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then u is found using the inverse transformation

$$u = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 510: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -\frac{127}{64}$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos\left(\frac{\sqrt{127}t}{8}\right)$$

Using the above, the solution for the original ode can now be found. The first solution

to the original ode in u is found from

$$\begin{aligned}u_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\&= z_1 e^{-\int \frac{1}{2} \frac{1}{1} dt} \\&= z_1 e^{-\frac{t}{2}} \\&= z_1 \left(e^{-\frac{t}{2}} \right)\end{aligned}$$

Which simplifies to

$$u_1 = e^{-\frac{t}{2}} \cos \left(\frac{\sqrt{127} t}{2} \right)$$

The second solution u_2 to the original ode is found using reduction of order

$$u_2 = u_1 \int \frac{e^{\int -\frac{B}{A} dt}}{u_1^2} dt$$

Substituting gives

$$\begin{aligned}u_2 &= u_1 \int \frac{e^{\int -\frac{1}{2} dt}}{(u_1)^2} dt \\&= u_1 \int \frac{e^{-\frac{t}{2}}}{(u_1)^2} dt \\&= u_1 \left(\frac{8\sqrt{127} \tan \left(\frac{\sqrt{127} t}{2} \right)}{127} \right)\end{aligned}$$

Therefore the solution is

$$\begin{aligned}u &= c_1 u_1 + c_2 u_2 \\&= c_1 \left(e^{-\frac{t}{2}} \cos \left(\frac{\sqrt{127} t}{2} \right) \right) + c_2 \left(e^{-\frac{t}{2}} \cos \left(\frac{\sqrt{127} t}{2} \right) \left(\frac{8\sqrt{127} \tan \left(\frac{\sqrt{127} t}{2} \right)}{127} \right) \right)\end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = c_1 e^{-\frac{t}{8}} \cos\left(\frac{\sqrt{127}t}{8}\right) + \frac{8c_2 \sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127}t}{8}\right)}{127} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 0$ and $t = 0$ in the above gives

$$0 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$u' = -\frac{c_1 e^{-\frac{t}{8}} \cos\left(\frac{\sqrt{127}t}{8}\right)}{8} - \frac{c_1 e^{-\frac{t}{8}} \sqrt{127} \sin\left(\frac{\sqrt{127}t}{8}\right)}{8} - \frac{c_2 \sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127}t}{8}\right)}{127} + c_2 e^{-\frac{t}{8}} \cos\left(\frac{\sqrt{127}t}{8}\right)$$

substituting $u' = 2$ and $t = 0$ in the above gives

$$2 = -\frac{c_1}{8} + c_2 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 0$$

$$c_2 = 2$$

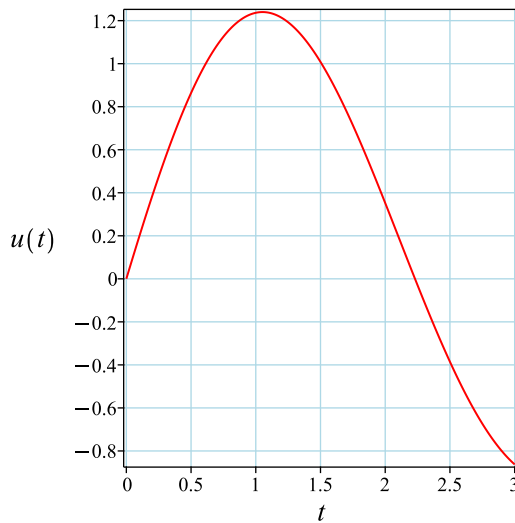
Substituting these values back in above solution results in

$$u = \frac{16\sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127}t}{8}\right)}{127}$$

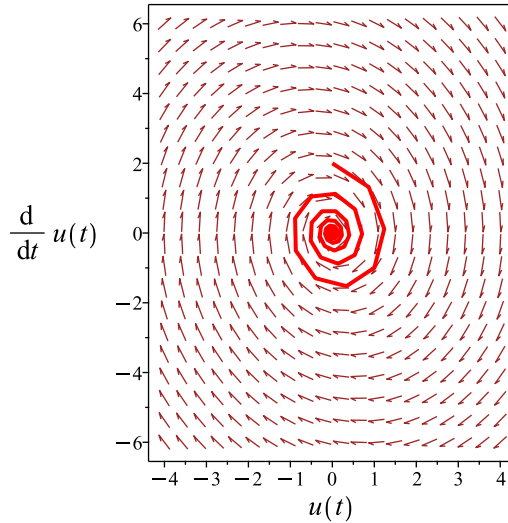
Summary

The solution(s) found are the following

$$u = \frac{16\sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127}t}{8}\right)}{127} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = \frac{16\sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127}t}{8}\right)}{127}$$

Verified OK.

11.2.4 Maple step by step solution

Let's solve

$$\left[u'' + \frac{u'}{4} + 2u = 0, u(0) = 0, u' \Big|_{\{t=0\}} = 2 \right]$$

- Highest derivative means the order of the ODE is 2

u''

- Characteristic polynomial of ODE

$$r^2 + \frac{1}{4}r + 2 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{\left(-\frac{1}{4}\right) \pm \left(\sqrt{-\frac{127}{16}}\right)}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{8} - \frac{\text{I}\sqrt{127}}{8}, -\frac{1}{8} + \frac{\text{I}\sqrt{127}}{8}\right)$$

- 1st solution of the ODE

$$u_1(t) = e^{-\frac{t}{8}} \cos\left(\frac{\sqrt{127}t}{8}\right)$$

- 2nd solution of the ODE

$$u_2(t) = e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127}t}{8}\right)$$

- General solution of the ODE

$$u = c_1 u_1(t) + c_2 u_2(t)$$

- Substitute in solutions

$$u = c_1 e^{-\frac{t}{8}} \cos\left(\frac{\sqrt{127}t}{8}\right) + c_2 \sin\left(\frac{\sqrt{127}t}{8}\right) e^{-\frac{t}{8}}$$

- Check validity of solution $u = c_1 e^{-\frac{t}{8}} \cos\left(\frac{\sqrt{127}t}{8}\right) + c_2 \sin\left(\frac{\sqrt{127}t}{8}\right) e^{-\frac{t}{8}}$

- Use initial condition $u(0) = 0$

$$0 = c_1$$

- Compute derivative of the solution

$$u' = -\frac{c_1 e^{-\frac{t}{8}} \cos\left(\frac{\sqrt{127}t}{8}\right)}{8} - \frac{c_1 e^{-\frac{t}{8}} \sqrt{127} \sin\left(\frac{\sqrt{127}t}{8}\right)}{8} + \frac{c_2 \sqrt{127} \cos\left(\frac{\sqrt{127}t}{8}\right) e^{-\frac{t}{8}}}{8} - \frac{c_2 \sin\left(\frac{\sqrt{127}t}{8}\right) e^{-\frac{t}{8}}}{8}$$

- Use the initial condition $u' \Big|_{\{t=0\}} = 2$

$$2 = -\frac{c_1}{8} + \frac{\sqrt{127} c_2}{8}$$

- Solve for c_1 and c_2

$$\left\{ c_1 = 0, c_2 = \frac{16\sqrt{127}}{127} \right\}$$

- Substitute constant values into general solution and simplify

$$u = \frac{16\sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127}t}{8}\right)}{127}$$

- Solution to the IVP

$$u = \frac{16\sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127}t}{8}\right)}{127}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
<- constant coefficients successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 20

```
dsolve([diff(u(t),t$2)+1/4*diff(u(t),t)+2*u(t) = 0,u(0) = 0, D(u)(0) = 2],u(t), singsol=all)
```

$$u(t) = \frac{16\sqrt{127} e^{-\frac{t}{8}} \sin\left(\frac{\sqrt{127}t}{8}\right)}{127}$$

✓ Solution by Mathematica

Time used: 0.025 (sec). Leaf size: 30

```
DSolve[{u''[t]+1/4*u'[t]+2*u[t] ==0,{u[0]==0,u'[0]==2}},u[t],t,IncludeSingularSolutions -> T
```

$$u(t) \rightarrow \frac{16e^{-t/8} \sin\left(\frac{\sqrt{127}t}{8}\right)}{\sqrt{127}}$$

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12.1 problem 21	3061
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12.1 problem 21

12.1.1 Existence and uniqueness analysis	3061
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Internal problem ID [707]

Internal file name [OUTPUT/707_Sunday_June_05_2022_01_47_25_AM_44410193/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.7 Forced Vibrations. page 217

Problem number: 21.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$u'' + \frac{u'}{8} + 4u = 3 \cos\left(\frac{t}{4}\right)$$

With initial conditions

$$[u(0) = 2, u'(0) = 0]$$

12.1.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$u'' + p(t)u' + q(t)u = F$$

Where here

$$p(t) = \frac{1}{8}$$

$$q(t) = 4$$

$$F = 3 \cos\left(\frac{t}{4}\right)$$

Hence the ode is

$$u'' + \frac{u'}{8} + 4u = 3 \cos\left(\frac{t}{4}\right)$$

The domain of $p(t) = \frac{1}{8}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = 4$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. The domain of $F = 3 \cos\left(\frac{t}{4}\right)$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

12.1.2 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Au''(t) + Bu'(t) + Cu(t) = f(t)$$

Where $A = 1, B = \frac{1}{8}, C = 4, f(t) = 3 \cos\left(\frac{t}{4}\right)$. Let the solution be

$$u = u_h + u_p$$

Where u_h is the solution to the homogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = 0$, and u_p is a particular solution to the non-homogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = f(t)$. u_h is the solution to

$$u'' + \frac{u'}{8} + 4u = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Au''(t) + Bu'(t) + Cu(t) = 0$$

Where in the above $A = 1, B = \frac{1}{8}, C = 4$. Let the solution be $u = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + \frac{\lambda e^{\lambda t}}{8} + 4 e^{\lambda t} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + \frac{1}{8}\lambda + 4 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = \frac{1}{8}, C = 4$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-\frac{1}{8}}{(2)(1)} \pm \frac{1}{(2)(1)}\sqrt{\frac{1^2}{8} - (4)(1)(4)} \\ &= -\frac{1}{16} \pm \frac{i\sqrt{1023}}{16} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= -\frac{1}{16} + \frac{i\sqrt{1023}}{16} \\ \lambda_2 &= -\frac{1}{16} - \frac{i\sqrt{1023}}{16} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= -\frac{1}{16} + \frac{i\sqrt{1023}}{16} \\ \lambda_2 &= -\frac{1}{16} - \frac{i\sqrt{1023}}{16} \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -\frac{1}{16}$ and $\beta = \frac{\sqrt{1023}}{16}$. Therefore the final solution, when using Euler relation, can be written as

$$u = e^{\alpha t}(c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$u = e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023}t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023}t}{16} \right) \right)$$

Therefore the homogeneous solution u_h is

$$u_h = e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023} t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023} t}{16} \right) \right)$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$3 \cos \left(\frac{t}{4} \right)$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$\left[\left\{ \cos \left(\frac{t}{4} \right), \sin \left(\frac{t}{4} \right) \right\} \right]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ e^{-\frac{t}{16}} \cos \left(\frac{\sqrt{1023} t}{16} \right), e^{-\frac{t}{16}} \sin \left(\frac{\sqrt{1023} t}{16} \right) \right\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$u_p = A_1 \cos \left(\frac{t}{4} \right) + A_2 \sin \left(\frac{t}{4} \right)$$

The unknowns $\{A_1, A_2\}$ are found by substituting the above trial solution u_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$\frac{63A_1 \cos \left(\frac{t}{4} \right)}{16} + \frac{63A_2 \sin \left(\frac{t}{4} \right)}{16} - \frac{A_1 \sin \left(\frac{t}{4} \right)}{32} + \frac{A_2 \cos \left(\frac{t}{4} \right)}{32} = 3 \cos \left(\frac{t}{4} \right)$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = \frac{12096}{15877}, A_2 = \frac{96}{15877} \right]$$

Substituting the above back in the above trial solution u_p , gives the particular solution

$$u_p = \frac{12096 \cos \left(\frac{t}{4} \right)}{15877} + \frac{96 \sin \left(\frac{t}{4} \right)}{15877}$$

Therefore the general solution is

$$u = u_h + u_p$$

$$= \left(e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023}t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023}t}{16} \right) \right) \right) + \left(\frac{12096 \cos \left(\frac{t}{4} \right)}{15877} + \frac{96 \sin \left(\frac{t}{4} \right)}{15877} \right)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023}t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023}t}{16} \right) \right) + \frac{12096 \cos \left(\frac{t}{4} \right)}{15877} + \frac{96 \sin \left(\frac{t}{4} \right)}{15877} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 2$ and $t = 0$ in the above gives

$$2 = c_1 + \frac{12096}{15877} \quad (1A)$$

Taking derivative of the solution gives

$$u' = -\frac{e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023}t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023}t}{16} \right) \right)}{16} + e^{-\frac{t}{16}} \left(-\frac{c_1 \sqrt{1023} \sin \left(\frac{\sqrt{1023}t}{16} \right)}{16} + \frac{c_2 \sqrt{1023} \cos \left(\frac{\sqrt{1023}t}{16} \right)}{16} \right)$$

substituting $u' = 0$ and $t = 0$ in the above gives

$$0 = -\frac{c_1}{16} + \frac{\sqrt{1023} c_2}{16} + \frac{24}{15877} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{19658}{15877}$$

$$c_2 = \frac{19274\sqrt{1023}}{16242171}$$

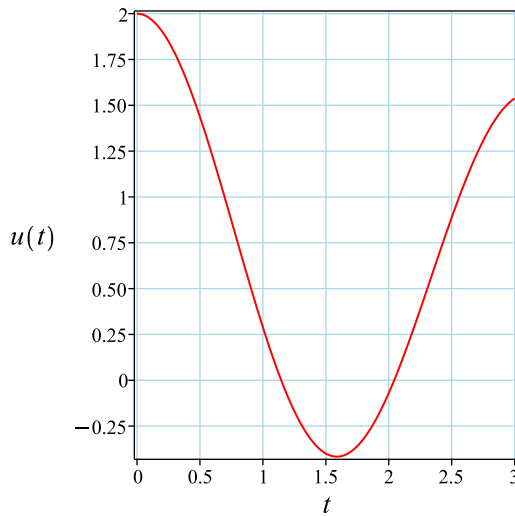
Substituting these values back in above solution results in

$$u = \frac{19658 e^{-\frac{t}{16}} \cos \left(\frac{\sqrt{1023}t}{16} \right)}{15877} + \frac{19274 e^{-\frac{t}{16}} \sin \left(\frac{\sqrt{1023}t}{16} \right) \sqrt{1023}}{16242171} + \frac{12096 \cos \left(\frac{t}{4} \right)}{15877} + \frac{96 \sin \left(\frac{t}{4} \right)}{15877}$$

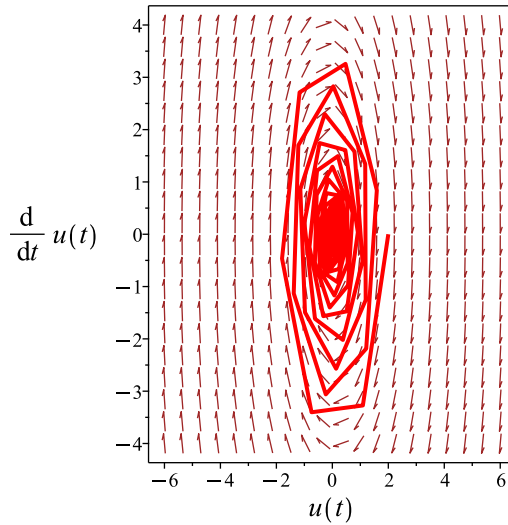
Summary

The solution(s) found are the following

$$u = \frac{19658 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{15877} + \frac{19274 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{16242171} + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = \frac{19658 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{15877} + \frac{19274 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{16242171} + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877}$$

Verified OK.

12.1.3 Solving using Kovacic algorithm

Writing the ode as

$$u'' + \frac{u'}{8} + 4u = 0 \quad (1)$$

$$Au'' + Bu' + Cu = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= \frac{1}{8} \\ C &= 4 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ue^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1023}{256} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1023 \\ t &= 256 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = -\frac{1023z(t)}{256} \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then u is found using the inverse transformation

$$u = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 512: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -\frac{1023}{256}$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

Using the above, the solution for the original ode can now be found. The first solution

to the original ode in u is found from

$$\begin{aligned} u_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{8}{1} dt} \\ &= z_1 e^{-\frac{t}{16}} \\ &= z_1 \left(e^{-\frac{t}{16}} \right) \end{aligned}$$

Which simplifies to

$$u_1 = e^{-\frac{t}{16}} \cos \left(\frac{\sqrt{1023} t}{16} \right)$$

The second solution u_2 to the original ode is found using reduction of order

$$u_2 = u_1 \int \frac{e^{\int -\frac{B}{A} dt}}{u_1^2} dt$$

Substituting gives

$$\begin{aligned} u_2 &= u_1 \int \frac{e^{\int -\frac{1}{8} dt}}{(u_1)^2} dt \\ &= u_1 \int \frac{e^{-\frac{t}{8}}}{(u_1)^2} dt \\ &= u_1 \left(\frac{16\sqrt{1023} \tan \left(\frac{\sqrt{1023} t}{16} \right)}{1023} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} u &= c_1 u_1 + c_2 u_2 \\ &= c_1 \left(e^{-\frac{t}{16}} \cos \left(\frac{\sqrt{1023} t}{16} \right) \right) + c_2 \left(e^{-\frac{t}{16}} \cos \left(\frac{\sqrt{1023} t}{16} \right) \left(\frac{16\sqrt{1023} \tan \left(\frac{\sqrt{1023} t}{16} \right)}{1023} \right) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$u = u_h + u_p$$

Where u_h is the solution to the homogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = 0$, and u_p is a particular solution to the nonhomogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = f(t)$. u_h is the solution to

$$u'' + \frac{u'}{8} + 4u = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$u_h = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + \frac{16c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$3 \cos\left(\frac{t}{4}\right)$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$\left[\left\{ \cos\left(\frac{t}{4}\right), \sin\left(\frac{t}{4}\right) \right\} \right]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right), \frac{16 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} \right\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$u_p = A_1 \cos\left(\frac{t}{4}\right) + A_2 \sin\left(\frac{t}{4}\right)$$

The unknowns $\{A_1, A_2\}$ are found by substituting the above trial solution u_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$\frac{63A_1 \cos\left(\frac{t}{4}\right)}{16} + \frac{63A_2 \sin\left(\frac{t}{4}\right)}{16} - \frac{A_1 \sin\left(\frac{t}{4}\right)}{32} + \frac{A_2 \cos\left(\frac{t}{4}\right)}{32} = 3 \cos\left(\frac{t}{4}\right)$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = \frac{12096}{15877}, A_2 = \frac{96}{15877} \right]$$

Substituting the above back in the above trial solution u_p , gives the particular solution

$$u_p = \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877}$$

Therefore the general solution is

$$\begin{aligned} u &= u_h + u_p \\ &= \left(c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + \frac{16c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} \right) \\ &\quad + \left(\frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877} \right) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + \frac{16c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 2$ and $t = 0$ in the above gives

$$2 = c_1 + \frac{12096}{15877} \quad (1A)$$

Taking derivative of the solution gives

$$u' = -\frac{c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{c_1 e^{-\frac{t}{16}} \sqrt{1023} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + c_2 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

substituting $u' = 0$ and $t = 0$ in the above gives

$$0 = \frac{24}{15877} + c_2 - \frac{c_1}{16} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{19658}{15877}$$

$$c_2 = \frac{9637}{127016}$$

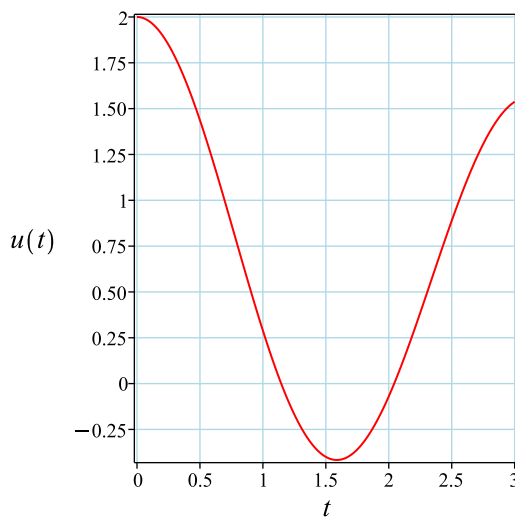
Substituting these values back in above solution results in

$$u = \frac{19658 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{15877} + \frac{19274 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{16242171} + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877}$$

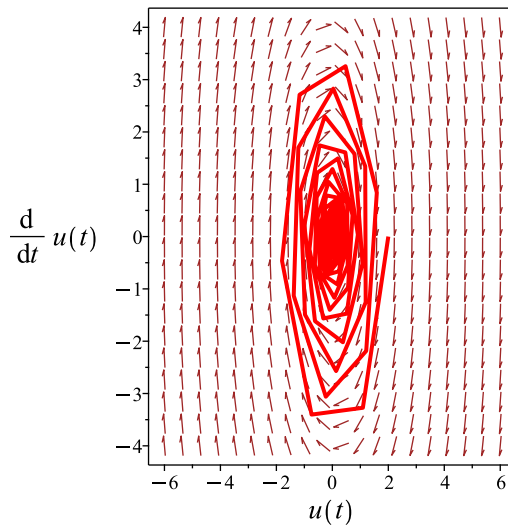
Summary

The solution(s) found are the following

$$u = \frac{19658 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{15877} + \frac{19274 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{16242171} + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = \frac{19658 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{15877} + \frac{19274 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{16242171} + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877}$$

Verified OK.

12.1.4 Maple step by step solution

Let's solve

$$\left[u'' + \frac{u'}{8} + 4u = 3 \cos\left(\frac{t}{4}\right), u(0) = 2, u' \Big|_{\{t=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

u''

- Characteristic polynomial of homogeneous ODE

$$r^2 + \frac{1}{8}r + 4 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{\left(-\frac{1}{8}\right) \pm \left(\sqrt{-\frac{1023}{64}}\right)}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{16} - \frac{i\sqrt{1023}}{16}, -\frac{1}{16} + \frac{i\sqrt{1023}}{16}\right)$$

- 1st solution of the homogeneous ODE

$$u_1(t) = e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

- 2nd solution of the homogeneous ODE

$$u_2(t) = e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right)$$

- General solution of the ODE

$$u = c_1 u_1(t) + c_2 u_2(t) + u_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right) + u_p(t)$$

- Find a particular solution $u_p(t)$ of the ODE

- Use variation of parameters to find u_p here $f(t)$ is the forcing function

$$\left[u_p(t) = -u_1(t) \left(\int \frac{u_2(t)f(t)}{W(u_1(t), u_2(t))} dt \right) + u_2(t) \left(\int \frac{u_1(t)f(t)}{W(u_1(t), u_2(t))} dt \right), f(t) = 3 \cos\left(\frac{t}{4}\right) \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(u_1(t), u_2(t)) = \begin{vmatrix} e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) & e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \\ -\frac{e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{16} & -\frac{e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} + \frac{e^{-\frac{t}{16}} \sqrt{1023} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} \end{vmatrix}$$

- Compute Wronskian

$$W(u_1(t), u_2(t)) = \frac{\sqrt{1023} e^{-\frac{t}{16}}}{16}$$

- Substitute functions into equation for $u_p(t)$

$$u_p(t) = -\frac{16 e^{-\frac{t}{16}} \sqrt{1023} \left(\cos\left(\frac{\sqrt{1023}t}{16}\right) \left(\int \cos\left(\frac{t}{4}\right) e^{\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) dt \right) - \sin\left(\frac{\sqrt{1023}t}{16}\right) \left(\int \cos\left(\frac{t}{4}\right) e^{\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) dt \right) \right)}{341}$$

- Compute integrals

$$u_p(t) = \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877}$$

- Substitute particular solution into general solution to ODE

$$u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right) + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877}$$

- Check validity of solution $u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right) + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877}$

- Use initial condition $u(0) = 2$

$$2 = c_1 + \frac{12096}{15877}$$

- Compute derivative of the solution

$$u' = -\frac{c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{c_1 e^{-\frac{t}{16}} \sqrt{1023} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} + \frac{e^{-\frac{t}{16}} c_2 \sqrt{1023} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{3024}{15877}$$

- Use the initial condition $u' \Big|_{\{t=0\}} = 0$

$$0 = -\frac{c_1}{16} + \frac{\sqrt{1023} c_2}{16} + \frac{24}{15877}$$

- Solve for c_1 and c_2

$$\left\{ c_1 = \frac{19658}{15877}, c_2 = \frac{19274\sqrt{1023}}{16242171} \right\}$$

- Substitute constant values into general solution and simplify

$$u = \frac{19658 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{15877} + \frac{19274 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{16242171} + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877}$$

- Solution to the IVP

$$u = \frac{19658 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{15877} + \frac{19274 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{16242171} + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877}$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]  
trying a double symmetry of the form [xi=0, eta=F(x)]  
-> Try solving first the homogeneous part of the ODE  
    checking if the LODE has constant coefficients  
    <- constant coefficients successful  
<- solving first the homogeneous part of the ODE successful`
```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 46

```
dsolve([diff(u(t),t$2)+125/1000*diff(u(t),t)+4*u(t) = 3*cos(t/4),u(0) = 2, D(u)(0) = 0],u(t))
```

$$u(t) = \frac{19274 e^{-\frac{t}{16}} \sqrt{1023} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16242171} + \frac{19658 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{15877} + \frac{96 \sin\left(\frac{t}{4}\right)}{15877} + \frac{12096 \cos\left(\frac{t}{4}\right)}{15877}$$

✓ Solution by Mathematica

Time used: 0.039 (sec). Leaf size: 71

```
DSolve[{u'[t]+125/1000*u'[t]+4*u[t] ==3*Cos[t/4],{u[0]==0,u'[0]==0}},u[t],t,IncludeSingular
```

$$u(t) \rightarrow \frac{32\left(1023 \sin\left(\frac{t}{4}\right) - 130\sqrt{1023}e^{-t/16} \sin\left(\frac{\sqrt{1023}t}{16}\right) + 128898 \cos\left(\frac{t}{4}\right) - 128898e^{-t/16} \cos\left(\frac{\sqrt{1023}t}{16}\right)\right)}{5414057}$$

12.2 problem 22

12.2.1 Existence and uniqueness analysis	3076
12.2.2 Solving as second order linear constant coeff ode	3077
12.2.3 Solving using Kovacic algorithm	3081
12.2.4 Maple step by step solution	3087

Internal problem ID [708]

Internal file name [OUTPUT/708_Sunday_June_05_2022_01_47_27_AM_67339897/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.7 Forced Vibrations. page 217

Problem number: 22.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$u'' + \frac{u'}{8} + 4u = 3 \cos(2t)$$

With initial conditions

$$[u(0) = 2, u'(0) = 0]$$

12.2.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$u'' + p(t)u' + q(t)u = F$$

Where here

$$p(t) = \frac{1}{8}$$

$$q(t) = 4$$

$$F = 3 \cos(2t)$$

Hence the ode is

$$u'' + \frac{u'}{8} + 4u = 3 \cos(2t)$$

The domain of $p(t) = \frac{1}{8}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = 4$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. The domain of $F = 3 \cos(2t)$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

12.2.2 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Au''(t) + Bu'(t) + Cu(t) = f(t)$$

Where $A = 1, B = \frac{1}{8}, C = 4, f(t) = 3 \cos(2t)$. Let the solution be

$$u = u_h + u_p$$

Where u_h is the solution to the homogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = 0$, and u_p is a particular solution to the non-homogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = f(t)$. u_h is the solution to

$$u'' + \frac{u'}{8} + 4u = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Au''(t) + Bu'(t) + Cu(t) = 0$$

Where in the above $A = 1, B = \frac{1}{8}, C = 4$. Let the solution be $u = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + \frac{\lambda e^{\lambda t}}{8} + 4 e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + \frac{1}{8}\lambda + 4 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = \frac{1}{8}, C = 4$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-\frac{1}{8}}{(2)(1)} \pm \frac{1}{(2)(1)}\sqrt{\frac{1^2}{8} - (4)(1)(4)} \\ &= -\frac{1}{16} \pm \frac{i\sqrt{1023}}{16} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= -\frac{1}{16} + \frac{i\sqrt{1023}}{16} \\ \lambda_2 &= -\frac{1}{16} - \frac{i\sqrt{1023}}{16} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= -\frac{1}{16} + \frac{i\sqrt{1023}}{16} \\ \lambda_2 &= -\frac{1}{16} - \frac{i\sqrt{1023}}{16} \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -\frac{1}{16}$ and $\beta = \frac{\sqrt{1023}}{16}$. Therefore the final solution, when using Euler relation, can be written as

$$u = e^{\alpha t}(c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$u = e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023}t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023}t}{16} \right) \right)$$

Therefore the homogeneous solution u_h is

$$u_h = e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023} t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023} t}{16} \right) \right)$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$3 \cos (2t)$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{\cos (2t), \sin (2t)\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ e^{-\frac{t}{16}} \cos \left(\frac{\sqrt{1023} t}{16} \right), e^{-\frac{t}{16}} \sin \left(\frac{\sqrt{1023} t}{16} \right) \right\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$u_p = A_1 \cos (2t) + A_2 \sin (2t)$$

The unknowns $\{A_1, A_2\}$ are found by substituting the above trial solution u_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-\frac{A_1 \sin (2t)}{4} + \frac{A_2 \cos (2t)}{4} = 3 \cos (2t)$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = 0, A_2 = 12]$$

Substituting the above back in the above trial solution u_p , gives the particular solution

$$u_p = 12 \sin (2t)$$

Therefore the general solution is

$$\begin{aligned} u &= u_h + u_p \\ &= \left(e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023} t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023} t}{16} \right) \right) \right) + (12 \sin (2t)) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023} t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023} t}{16} \right) \right) + 12 \sin (2t) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 2$ and $t = 0$ in the above gives

$$2 = c_1 \quad (1A)$$

Taking derivative of the solution gives

$$u' = -\frac{e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023} t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023} t}{16} \right) \right)}{16} + e^{-\frac{t}{16}} \left(-\frac{c_1 \sqrt{1023} \sin \left(\frac{\sqrt{1023} t}{16} \right)}{16} + \frac{c_2 \sqrt{1023} \cos \left(\frac{\sqrt{1023} t}{16} \right)}{16} \right)$$

substituting $u' = 0$ and $t = 0$ in the above gives

$$0 = -\frac{c_1}{16} + \frac{\sqrt{1023} c_2}{16} + 24 \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 2$$

$$c_2 = -\frac{382\sqrt{1023}}{1023}$$

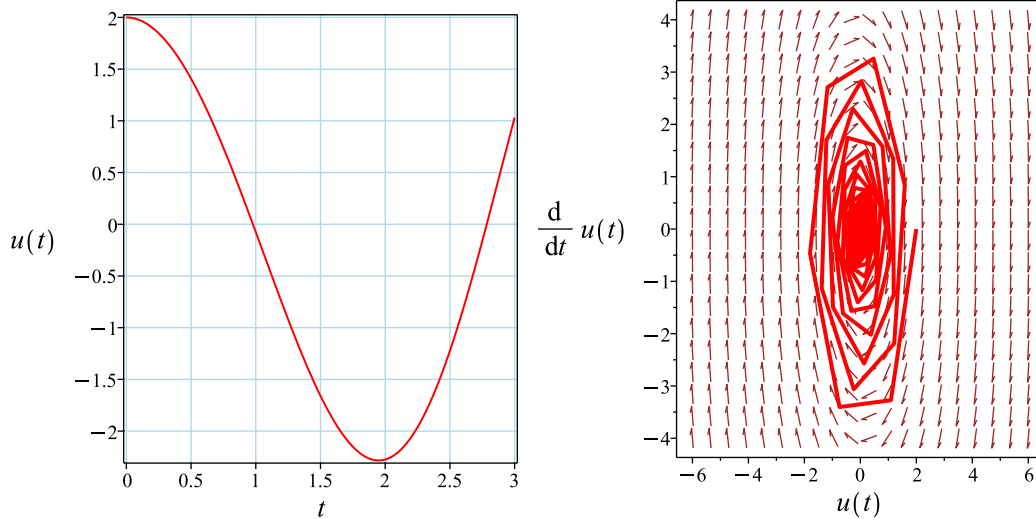
Substituting these values back in above solution results in

$$u = -\frac{382 e^{-\frac{t}{16}} \sin \left(\frac{\sqrt{1023} t}{16} \right) \sqrt{1023}}{1023} + 2 e^{-\frac{t}{16}} \cos \left(\frac{\sqrt{1023} t}{16} \right) + 12 \sin (2t)$$

Summary

The solution(s) found are the following

$$u = -\frac{382 e^{-\frac{t}{16}} \sin \left(\frac{\sqrt{1023} t}{16} \right) \sqrt{1023}}{1023} + 2 e^{-\frac{t}{16}} \cos \left(\frac{\sqrt{1023} t}{16} \right) + 12 \sin (2t) \quad (1)$$



(a) Solution plot

(b) Slope field plot

Verification of solutions

$$u = -\frac{382 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + 2 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + 12 \sin(2t)$$

Verified OK.

12.2.3 Solving using Kovacic algorithm

Writing the ode as

$$u'' + \frac{u'}{8} + 4u = 0 \tag{1}$$

$$Au'' + Bu' + Cu = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= \frac{1}{8} \\ C &= 4 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ue^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1023}{256} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1023 \\ t &= 256 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = -\frac{1023z(t)}{256} \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then u is found using the inverse transformation

$$u = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 514: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -\frac{1023}{256}$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in u is found from

$$\begin{aligned} u_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{8}{1} dt} \\ &= z_1 e^{-\frac{t}{16}} \\ &= z_1 \left(e^{-\frac{t}{16}} \right) \end{aligned}$$

Which simplifies to

$$u_1 = e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

The second solution u_2 to the original ode is found using reduction of order

$$u_2 = u_1 \int \frac{e^{\int -\frac{B}{A} dt}}{u_1^2} dt$$

Substituting gives

$$\begin{aligned}
 u_2 &= u_1 \int \frac{e^{\int -\frac{1}{8} dt}}{(u_1)^2} dt \\
 &= u_1 \int \frac{e^{-\frac{t}{8}}}{(u_1)^2} dt \\
 &= u_1 \left(\frac{16\sqrt{1023} \tan\left(\frac{\sqrt{1023}t}{16}\right)}{1023} \right)
 \end{aligned}$$

Therefore the solution is

$$\begin{aligned}
 u &= c_1 u_1 + c_2 u_2 \\
 &= c_1 \left(e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) \right) + c_2 \left(e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) \left(\frac{16\sqrt{1023} \tan\left(\frac{\sqrt{1023}t}{16}\right)}{1023} \right) \right)
 \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$u = u_h + u_p$$

Where u_h is the solution to the homogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = 0$, and u_p is a particular solution to the nonhomogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = f(t)$. u_h is the solution to

$$u'' + \frac{u'}{8} + 4u = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$u_h = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + \frac{16c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$3 \cos(2t)$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{\cos(2t), \sin(2t)\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right), \frac{16 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} \right\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$u_p = A_1 \cos(2t) + A_2 \sin(2t)$$

The unknowns $\{A_1, A_2\}$ are found by substituting the above trial solution u_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-\frac{A_1 \sin(2t)}{4} + \frac{A_2 \cos(2t)}{4} = 3 \cos(2t)$$

Solving for the unknowns by comparing coefficients results in

$$[A_1 = 0, A_2 = 12]$$

Substituting the above back in the above trial solution u_p , gives the particular solution

$$u_p = 12 \sin(2t)$$

Therefore the general solution is

$$\begin{aligned} u &= u_h + u_p \\ &= \left(c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + \frac{16c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} \right) + (12 \sin(2t)) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + \frac{16c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + 12 \sin(2t) \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 2$ and $t = 0$ in the above gives

$$2 = c_1 \tag{1A}$$

Taking derivative of the solution gives

$$u' = -\frac{c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{c_1 e^{-\frac{t}{16}} \sqrt{1023} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + c_2 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

substituting $u' = 0$ and $t = 0$ in the above gives

$$0 = 24 - \frac{c_1}{16} + c_2 \tag{2A}$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = 2$$

$$c_2 = -\frac{191}{8}$$

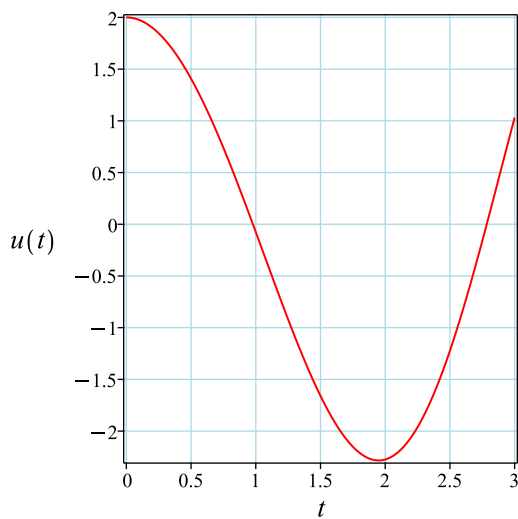
Substituting these values back in above solution results in

$$u = -\frac{382 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + 2 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + 12 \sin(2t)$$

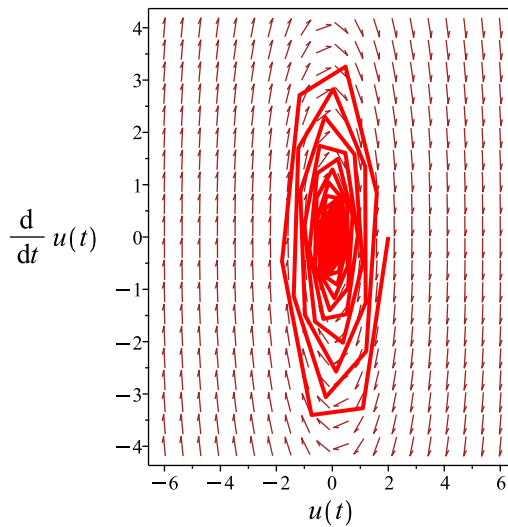
Summary

The solution(s) found are the following

$$u = -\frac{382 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + 2 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + 12 \sin(2t) \tag{1}$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = -\frac{382 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + 2 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + 12 \sin(2t)$$

Verified OK.

12.2.4 Maple step by step solution

Let's solve

$$\left[u'' + \frac{u'}{8} + 4u = 3 \cos(2t), u(0) = 2, u' \Big|_{\{t=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

u''

- Characteristic polynomial of homogeneous ODE

$$r^2 + \frac{1}{8}r + 4 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{\left(-\frac{1}{8}\right) \pm \left(\sqrt{-\frac{1023}{64}}\right)}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{16} - \frac{\sqrt{1023}}{16}, -\frac{1}{16} + \frac{\sqrt{1023}}{16} \right)$$

- 1st solution of the homogeneous ODE

$$u_1(t) = e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

- 2nd solution of the homogeneous ODE

$$u_2(t) = e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right)$$

- General solution of the ODE

$$u = c_1 u_1(t) + c_2 u_2(t) + u_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right) + u_p(t)$$

- Find a particular solution $u_p(t)$ of the ODE

- Use variation of parameters to find u_p here $f(t)$ is the forcing function

$$\left[u_p(t) = -u_1(t) \left(\int \frac{u_2(t)f(t)}{W(u_1(t),u_2(t))} dt \right) + u_2(t) \left(\int \frac{u_1(t)f(t)}{W(u_1(t),u_2(t))} dt \right), f(t) = 3 \cos(2t) \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(u_1(t), u_2(t)) = \begin{vmatrix} e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) & e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \\ -\frac{e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{16} & -\frac{e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} + \frac{e^{-\frac{t}{16}} \sqrt{1023} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} \end{vmatrix}$$

- Compute Wronskian

$$W(u_1(t), u_2(t)) = \frac{\sqrt{1023} e^{-\frac{t}{8}}}{16}$$

- Substitute functions into equation for $u_p(t)$

$$u_p(t) = -\frac{16 e^{-\frac{t}{16}} \sqrt{1023} \left(\cos\left(\frac{\sqrt{1023}t}{16}\right) \left(\int \cos(2t) e^{\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) dt \right) - \sin\left(\frac{\sqrt{1023}t}{16}\right) \left(\int \cos(2t) e^{\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) dt \right) \right)}{341}$$

- Compute integrals

$$u_p(t) = 12 \sin(2t)$$

- Substitute particular solution into general solution to ODE

$$u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right) + 12 \sin(2t)$$

- Check validity of solution $u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right) + 12 \sin(2t)$

- Use initial condition $u(0) = 2$

$$2 = c_1$$

- Compute derivative of the solution

$$u' = -\frac{c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{c_1 e^{-\frac{t}{16}} \sqrt{1023} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} + \frac{e^{-\frac{t}{16}} c_2 \sqrt{1023} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} + 24 c_2$$

- Use the initial condition $u'|_{\{t=0\}} = 0$

$$0 = -\frac{c_1}{16} + \frac{\sqrt{1023}c_2}{16} + 24$$

- Solve for c_1 and c_2

$$\left\{ c_1 = 2, c_2 = -\frac{382\sqrt{1023}}{1023} \right\}$$

- Substitute constant values into general solution and simplify

$$u = -\frac{382 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + 2 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + 12 \sin(2t)$$

- Solution to the IVP

$$u = -\frac{382 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + 2 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + 12 \sin(2t)$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`

```

✓ Solution by Maple

Time used: 0.032 (sec). Leaf size: 40

```
dsolve([diff(u(t),t$2)+125/1000*diff(u(t),t)+4*u(t) = 3*cos(2*t),u(0) = 2, D(u)(0) = 0],u(t))
```

$$u(t) = -\frac{382 e^{-\frac{t}{16}} \sqrt{1023} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{1023} + 2 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + 12 \sin(2t)$$

✓ Solution by Mathematica

Time used: 0.03 (sec). Leaf size: 39

```
DSolve[{u''[t]+125/1000*u'[t]+4*u[t] ==3*Cos[2*t],{u[0]==0,u'[0]==0}},u[t],t,IncludeSingular
```

$$u(t) \rightarrow 12 \sin(2t) - 128 \sqrt{\frac{3}{341}} e^{-t/16} \sin\left(\frac{\sqrt{1023}t}{16}\right)$$

12.3 problem 23

12.3.1 Existence and uniqueness analysis	3091
12.3.2 Solving as second order linear constant coeff ode	3092
12.3.3 Solving using Kovacic algorithm	3096
12.3.4 Maple step by step solution	3102

Internal problem ID [709]

Internal file name [OUTPUT/709_Sunday_June_05_2022_01_47_29_AM_56108713/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.7 Forced Vibrations. page 217

Problem number: 23.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"kovacic", "second_order_linear_constant_coeff"**

Maple gives the following as the ode type

```
[[_2nd_order , _linear , _nonhomogeneous]]
```

$$u'' + \frac{u'}{8} + 4u = 3 \cos(6t)$$

With initial conditions

$$[u(0) = 2, u'(0) = 0]$$

12.3.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$u'' + p(t)u' + q(t)u = F$$

Where here

$$p(t) = \frac{1}{8}$$

$$q(t) = 4$$

$$F = 3 \cos(6t)$$

Hence the ode is

$$u'' + \frac{u'}{8} + 4u = 3 \cos(6t)$$

The domain of $p(t) = \frac{1}{8}$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is inside this domain. The domain of $q(t) = 4$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. The domain of $F = 3 \cos(6t)$ is

$$\{-\infty < t < \infty\}$$

And the point $t_0 = 0$ is also inside this domain. Hence solution exists and is unique.

12.3.2 Solving as second order linear constant coeff ode

This is second order non-homogeneous ODE. In standard form the ODE is

$$Au''(t) + Bu'(t) + Cu(t) = f(t)$$

Where $A = 1, B = \frac{1}{8}, C = 4, f(t) = 3 \cos(6t)$. Let the solution be

$$u = u_h + u_p$$

Where u_h is the solution to the homogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = 0$, and u_p is a particular solution to the non-homogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = f(t)$. u_h is the solution to

$$u'' + \frac{u'}{8} + 4u = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Au''(t) + Bu'(t) + Cu(t) = 0$$

Where in the above $A = 1, B = \frac{1}{8}, C = 4$. Let the solution be $u = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda t} + \frac{\lambda e^{\lambda t}}{8} + 4 e^{\lambda t} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + \frac{1}{8}\lambda + 4 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A}\sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = \frac{1}{8}, C = 4$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-\frac{1}{8}}{(2)(1)} \pm \frac{1}{(2)(1)}\sqrt{\frac{1^2}{8} - (4)(1)(4)} \\ &= -\frac{1}{16} \pm \frac{i\sqrt{1023}}{16} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= -\frac{1}{16} + \frac{i\sqrt{1023}}{16} \\ \lambda_2 &= -\frac{1}{16} - \frac{i\sqrt{1023}}{16} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= -\frac{1}{16} + \frac{i\sqrt{1023}}{16} \\ \lambda_2 &= -\frac{1}{16} - \frac{i\sqrt{1023}}{16} \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -\frac{1}{16}$ and $\beta = \frac{\sqrt{1023}}{16}$. Therefore the final solution, when using Euler relation, can be written as

$$u = e^{\alpha t}(c_1 \cos(\beta t) + c_2 \sin(\beta t))$$

Which becomes

$$u = e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023}t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023}t}{16} \right) \right)$$

Therefore the homogeneous solution u_h is

$$u_h = e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023} t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023} t}{16} \right) \right)$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$3 \cos(6t)$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{\cos(6t), \sin(6t)\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ e^{-\frac{t}{16}} \cos \left(\frac{\sqrt{1023} t}{16} \right), e^{-\frac{t}{16}} \sin \left(\frac{\sqrt{1023} t}{16} \right) \right\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$u_p = A_1 \cos(6t) + A_2 \sin(6t)$$

The unknowns $\{A_1, A_2\}$ are found by substituting the above trial solution u_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-32A_1 \cos(6t) - 32A_2 \sin(6t) - \frac{3A_1 \sin(6t)}{4} + \frac{3A_2 \cos(6t)}{4} = 3 \cos(6t)$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = -\frac{1536}{16393}, A_2 = \frac{36}{16393} \right]$$

Substituting the above back in the above trial solution u_p , gives the particular solution

$$u_p = -\frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$$

Therefore the general solution is

$$u = u_h + u_p$$

$$= \left(e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023} t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023} t}{16} \right) \right) \right) + \left(-\frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393} \right)$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023} t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023} t}{16} \right) \right) - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 2$ and $t = 0$ in the above gives

$$2 = c_1 - \frac{1536}{16393} \quad (1A)$$

Taking derivative of the solution gives

$$u' = -\frac{e^{-\frac{t}{16}} \left(c_1 \cos \left(\frac{\sqrt{1023} t}{16} \right) + c_2 \sin \left(\frac{\sqrt{1023} t}{16} \right) \right)}{16} + e^{-\frac{t}{16}} \left(-\frac{c_1 \sqrt{1023} \sin \left(\frac{\sqrt{1023} t}{16} \right)}{16} + \frac{c_2 \sqrt{1023} \cos \left(\frac{\sqrt{1023} t}{16} \right)}{16} \right)$$

substituting $u' = 0$ and $t = 0$ in the above gives

$$0 = -\frac{c_1}{16} + \frac{\sqrt{1023} c_2}{16} + \frac{216}{16393} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{34322}{16393}$$

$$c_2 = \frac{2806\sqrt{1023}}{1524549}$$

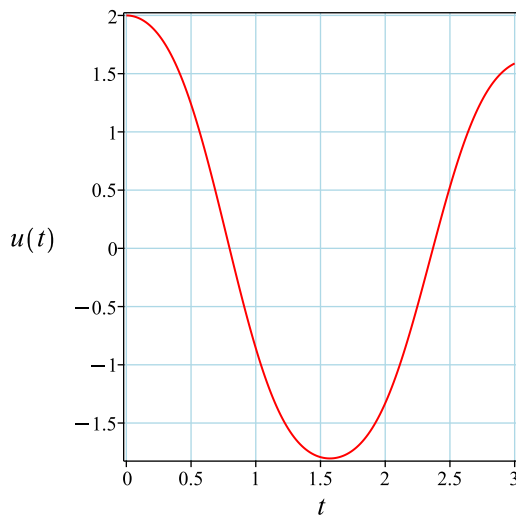
Substituting these values back in above solution results in

$$u = \frac{34322 e^{-\frac{t}{16}} \cos \left(\frac{\sqrt{1023} t}{16} \right)}{16393} + \frac{2806 e^{-\frac{t}{16}} \sin \left(\frac{\sqrt{1023} t}{16} \right) \sqrt{1023}}{1524549} - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$$

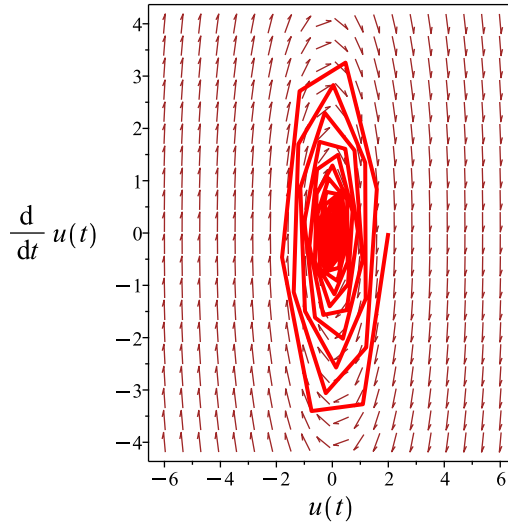
Summary

The solution(s) found are the following

$$u = \frac{34322 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16393} + \frac{2806 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1524549} - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = \frac{34322 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16393} + \frac{2806 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1524549} - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$$

Verified OK.

12.3.3 Solving using Kovacic algorithm

Writing the ode as

$$u'' + \frac{u'}{8} + 4u = 0 \quad (1)$$

$$Au'' + Bu' + Cu = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= 1 \\B &= \frac{1}{8} \\C &= 4\end{aligned}\tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ue^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t)\tag{4}$$

Where r is given by

$$\begin{aligned}r &= \frac{s}{t} \\&= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}\end{aligned}\tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1023}{256}\tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned}s &= -1023 \\t &= 256\end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = -\frac{1023z(t)}{256}\tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then u is found using the inverse transformation

$$u = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 516: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}
 \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\
 &= 0 - 0 \\
 &= 0
 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -\frac{1023}{256}$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in u is found from

$$\begin{aligned}
 u_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\
 &= z_1 e^{-\int \frac{1}{2} \frac{1}{1} dt} \\
 &= z_1 e^{-\frac{t}{16}} \\
 &= z_1 \left(e^{-\frac{t}{16}} \right)
 \end{aligned}$$

Which simplifies to

$$u_1 = e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

The second solution u_2 to the original ode is found using reduction of order

$$u_2 = u_1 \int \frac{e^{\int -\frac{B}{A} dt}}{u_1^2} dt$$

Substituting gives

$$\begin{aligned} u_2 &= u_1 \int \frac{e^{\int -\frac{1}{8} dt}}{(u_1)^2} dt \\ &= u_1 \int \frac{e^{-\frac{t}{8}}}{(u_1)^2} dt \\ &= u_1 \left(\frac{16\sqrt{1023} \tan\left(\frac{\sqrt{1023}t}{16}\right)}{1023} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} u &= c_1 u_1 + c_2 u_2 \\ &= c_1 \left(e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) \right) + c_2 \left(e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) \left(\frac{16\sqrt{1023} \tan\left(\frac{\sqrt{1023}t}{16}\right)}{1023} \right) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$u = u_h + u_p$$

Where u_h is the solution to the homogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = 0$, and u_p is a particular solution to the nonhomogeneous ODE $Au''(t) + Bu'(t) + Cu(t) = f(t)$. u_h is the solution to

$$u'' + \frac{u'}{8} + 4u = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$u_h = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + \frac{16c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$3 \cos(6t)$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{\cos(6t), \sin(6t)\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right), \frac{16 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} \right\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$u_p = A_1 \cos(6t) + A_2 \sin(6t)$$

The unknowns $\{A_1, A_2\}$ are found by substituting the above trial solution u_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-32A_1 \cos(6t) - 32A_2 \sin(6t) - \frac{3A_1 \sin(6t)}{4} + \frac{3A_2 \cos(6t)}{4} = 3 \cos(6t)$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = -\frac{1536}{16393}, A_2 = \frac{36}{16393} \right]$$

Substituting the above back in the above trial solution u_p , gives the particular solution

$$u_p = -\frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$$

Therefore the general solution is

$$\begin{aligned} u &= u_h + u_p \\ &= \left(c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + \frac{16c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} \right) \\ &\quad + \left(-\frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393} \right) \end{aligned}$$

Initial conditions are used to solve for the constants of integration.

Looking at the above solution

$$u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + \frac{16c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393} \quad (1)$$

Initial conditions are now substituted in the above solution. This will generate the required equations to solve for the integration constants. substituting $u = 2$ and $t = 0$ in the above gives

$$2 = c_1 - \frac{1536}{16393} \quad (1A)$$

Taking derivative of the solution gives

$$u' = -\frac{c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{c_1 e^{-\frac{t}{16}} \sqrt{1023} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{c_2 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1023} + c_2 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

substituting $u' = 0$ and $t = 0$ in the above gives

$$0 = \frac{216}{16393} + c_2 - \frac{c_1}{16} \quad (2A)$$

Equations {1A,2A} are now solved for $\{c_1, c_2\}$. Solving for the constants gives

$$c_1 = \frac{34322}{16393}$$

$$c_2 = \frac{15433}{131144}$$

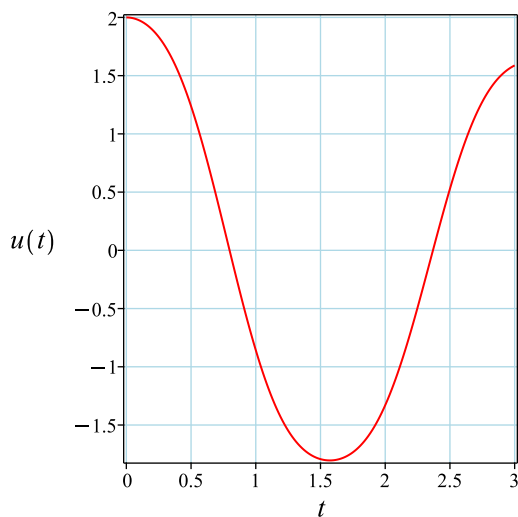
Substituting these values back in above solution results in

$$u = \frac{34322 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16393} + \frac{2806 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1524549} - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$$

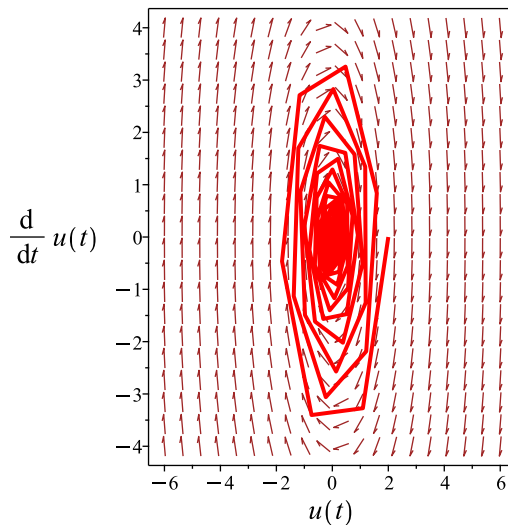
Summary

The solution(s) found are the following

$$u = \frac{34322 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16393} + \frac{2806 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1524549} - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393} \quad (1)$$



(a) Solution plot



(b) Slope field plot

Verification of solutions

$$u = \frac{34322 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16393} + \frac{2806 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1524549} - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$$

Verified OK.

12.3.4 Maple step by step solution

Let's solve

$$\left[u'' + \frac{u'}{8} + 4u = 3 \cos(6t), u(0) = 2, u'|_{\{t=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

u''

- Characteristic polynomial of homogeneous ODE

$$r^2 + \frac{1}{8}r + 4 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{\left(-\frac{1}{8}\right) \pm \left(\sqrt{-\frac{1023}{64}}\right)}{2}$$

- Roots of the characteristic polynomial

$$r = \left(-\frac{1}{16} - \frac{\text{I}\sqrt{1023}}{16}, -\frac{1}{16} + \frac{\text{I}\sqrt{1023}}{16}\right)$$

- 1st solution of the homogeneous ODE

$$u_1(t) = e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)$$

- 2nd solution of the homogeneous ODE

$$u_2(t) = e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right)$$

- General solution of the ODE

$$u = c_1 u_1(t) + c_2 u_2(t) + u_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right) + u_p(t)$$

- Find a particular solution $u_p(t)$ of the ODE

- Use variation of parameters to find u_p here $f(t)$ is the forcing function

$$\left[u_p(t) = -u_1(t) \left(\int \frac{u_2(t)f(t)}{W(u_1(t), u_2(t))} dt \right) + u_2(t) \left(\int \frac{u_1(t)f(t)}{W(u_1(t), u_2(t))} dt \right), f(t) = 3 \cos(6t) \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(u_1(t), u_2(t)) = \begin{vmatrix} e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) & e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \\ -\frac{e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{16} & -\frac{e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} + \frac{e^{-\frac{t}{16}} \sqrt{1023} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} \end{vmatrix}$$

- Compute Wronskian

$$W(u_1(t), u_2(t)) = \frac{\sqrt{1023} e^{-\frac{t}{8}}}{16}$$

- Substitute functions into equation for $u_p(t)$

$$u_p(t) = -\frac{16 e^{-\frac{t}{16}} \sqrt{1023} \left(\cos\left(\frac{\sqrt{1023}t}{16}\right) \left(\int \cos(6t) e^{\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) dt \right) - \sin\left(\frac{\sqrt{1023}t}{16}\right) \left(\int \cos(6t) e^{\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) dt \right) \right)}{341}$$

- Compute integrals

$$u_p(t) = -\frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$$

- Substitute particular solution into general solution to ODE

$$u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right) - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$$

- Check validity of solution $u = c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right) + e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right) - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$

- Use initial condition $u(0) = 2$

$$2 = c_1 - \frac{1536}{16393}$$

- Compute derivative of the solution

$$u' = -\frac{c_1 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{c_1 e^{-\frac{t}{16}} \sqrt{1023} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} - \frac{e^{-\frac{t}{16}} c_2 \sin\left(\frac{\sqrt{1023}t}{16}\right)}{16} + \frac{e^{-\frac{t}{16}} c_2 \sqrt{1023} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16} + \frac{9216}{16}$$

- Use the initial condition $u'|_{\{t=0\}} = 0$

$$0 = -\frac{c_1}{16} + \frac{\sqrt{1023}c_2}{16} + \frac{216}{16393}$$

- Solve for c_1 and c_2

$$\left\{ c_1 = \frac{34322}{16393}, c_2 = \frac{2806\sqrt{1023}}{1524549} \right\}$$

- Substitute constant values into general solution and simplify

$$u = \frac{34322 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16393} + \frac{2806 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1524549} - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$$

- Solution to the IVP

$$u = \frac{34322 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16393} + \frac{2806 e^{-\frac{t}{16}} \sin\left(\frac{\sqrt{1023}t}{16}\right) \sqrt{1023}}{1524549} - \frac{1536 \cos(6t)}{16393} + \frac{36 \sin(6t)}{16393}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`

```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 46

```
dsolve([diff(u(t),t$2)+125/1000*diff(u(t),t)+4*u(t) = 3*cos(6*t),u(0) = 2, D(u)(0) = 0],u(t))
```

$$u(t) = \frac{2806 e^{-\frac{t}{16}} \sqrt{1023} \sin\left(\frac{\sqrt{1023}t}{16}\right)}{1524549} + \frac{34322 e^{-\frac{t}{16}} \cos\left(\frac{\sqrt{1023}t}{16}\right)}{16393} + \frac{36 \sin(6t)}{16393} - \frac{1536 \cos(6t)}{16393}$$

✓ Solution by Mathematica

Time used: 0.032 (sec). Leaf size: 74

```
DSolve[{u'[t]+125/1000*u'[t]+4*u[t] ==3*Cos[6*t],{u[0]==0,u'[0]==0}},u[t],t,IncludeSingular
```

$u(t) \rightarrow$

$$\frac{4e^{-t/16} \left(-3069e^{t/16} \sin(6t) + 160\sqrt{1023} \sin\left(\frac{\sqrt{1023}t}{16}\right) + 130944e^{t/16} \cos(6t) - 130944 \cos\left(\frac{\sqrt{1023}t}{16}\right) \right)}{5590013}$$

12.4 problem 24

Internal problem ID [710]

Internal file name [OUTPUT/710_Sunday_June_05_2022_01_47_31_AM_72098390/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 3, Second order linear equations, 3.7 Forced Vibrations. page 217

Problem number: 24.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "**unknown**"

Maple gives the following as the ode type

[NONE]

Unable to solve or complete the solution.

$$u'' + u' + \frac{u^3}{5} = \cos(t)$$

With initial conditions

$$[u(0) = 2, u'(0) = 0]$$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying 2nd order Liouville
trying 2nd order WeierstrassP
trying 2nd order JacobiSN
differential order: 2; trying a linearization to 3rd order
trying 2nd order ODE linearizable_by_differentiation
trying 2nd order, 2 integrating factors of the form mu(x,y)
trying differential order: 2; missing variables
-> trying 2nd order, dynamical_symmetries, fully reducible to Abel through one integrating f
    --- trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods for dynam
    -> trying 2nd order, dynamical_symmetries, fully reducible to Abel through one integratin
trying 2nd order, integrating factors of the form mu(x,y)/(y)^n, only the singular cases
trying symmetries linear in x and y(x)
trying differential order: 2; exact nonlinear
trying 2nd order, integrating factor of the form mu(y)
trying 2nd order, integrating factor of the form mu(x,y)
trying 2nd order, integrating factor of the form mu(x,y)/(y)^n, only the general case
trying 2nd order, integrating factor of the form mu(y,y)
trying differential order: 2; mu polynomial in y
trying 2nd order, integrating factor of the form mu(x,y)
differential order: 2; looking for linear symmetries
-> trying 2nd order, the S-function method
    -> trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods for the S-
    -> trying 2nd order, the S-function method
    -> trying 2nd order, No Point Symmetries Class V
        --- trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods for dy
        -> trying 2nd order, No Point Symmetries Class V
    -> trying 2nd order, No Point Symmetries Class V
        --- trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods for dy
        -> trying 2nd order, No Point Symmetries Class V
    -> trying 2nd order, No Point Symmetries Class V
        --- trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods for dy
        -> trying 2nd order, No Point Symmetries Class V
trying 2nd order, integrating factor of the form mu(x,y)/(y)^n, only the general case
-> trying 2nd order, dynamical_symmetries, only a reduction of order through one integrating
    --- trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods for dynam
    -> trying 2nd order, dynamical_symmetries, only a reduction of order through one integrat
--- Trying Lie symmetry methods, 2nd order ---
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = 5
`, `-> Computing symmetries using: way = formal`
```

X Solution by Maple

```
dsolve([diff(u(t),t$2)+diff(u(t),t)+1/5*u(t)^3 = cos(t),u(0) = 2, D(u)(0) = 0],u(t), singsol
```

No solution found

X Solution by Mathematica

Time used: 0.0 (sec). Leaf size: 0

```
DSolve[{u''[t]+u'[t]+1/5*u[t]^3 ==3*Cos[t],{u[0]==0,u'[0]==0}},u[t],t,IncludeSingularSolutio
```

Not solved

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13.1 problem 1

13.1.1 Maple step by step solution 3117

Internal problem ID [711]

Internal file name [OUTPUT/711_Sunday_June_05_2022_01_47_34_AM_92887022/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 1.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "**second_order_linear_constant_coeff**", "**second_order_ode_can_be_made_integrable**", "**second order series method. Ordinary point**", "**second order series method. Taylor series method**"

Maple gives the following as the ode type

```
[[_2nd_order, _missing_x]]
```

$$y'' - y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (744)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (745)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= y \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\
 &= y' \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\
 &= y \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\
 &= y' \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\
 &= y
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= y(0) \\
 F_1 &= y'(0) \\
 F_2 &= y(0) \\
 F_3 &= y'(0) \\
 F_4 &= y(0)
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{24}x^4 + \frac{1}{720}x^6\right) y(0) + \left(x + \frac{1}{6}x^3 + \frac{1}{120}x^5\right) y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard

power series Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} a_n x^n \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (3)$$

For $0 \leq n$, the recurrence equation is

$$(n+2) a_{n+2} (n+1) - a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = \frac{a_n}{(n+2)(n+1)} \quad (5)$$

For $n = 0$ the recurrence equation gives

$$2a_2 - a_0 = 0$$

Which after substituting the earlier terms found becomes

$$a_2 = \frac{a_0}{2}$$

For $n = 1$ the recurrence equation gives

$$6a_3 - a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = \frac{a_1}{6}$$

For $n = 2$ the recurrence equation gives

$$12a_4 - a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{24}$$

For $n = 3$ the recurrence equation gives

$$20a_5 - a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_1}{120}$$

For $n = 4$ the recurrence equation gives

$$30a_6 - a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{a_0}{720}$$

For $n = 5$ the recurrence equation gives

$$42a_7 - a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{a_1}{5040}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x + \frac{1}{2} a_0 x^2 + \frac{1}{6} a_1 x^3 + \frac{1}{24} a_0 x^4 + \frac{1}{120} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{24}x^4\right) a_0 + \left(x + \frac{1}{6}x^3 + \frac{1}{120}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{24}x^4\right) c_1 + \left(x + \frac{1}{6}x^3 + \frac{1}{120}x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{24}x^4 + \frac{1}{720}x^6\right) y(0) + \left(x + \frac{1}{6}x^3 + \frac{1}{120}x^5\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{24}x^4\right) c_1 + \left(x + \frac{1}{6}x^3 + \frac{1}{120}x^5\right) c_2 + O(x^6) \quad (2)$$

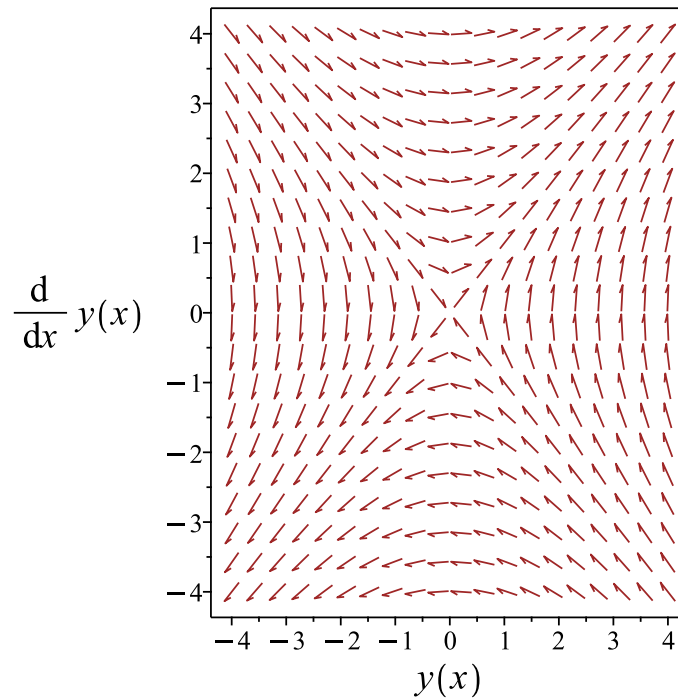


Figure 502: Slope field plot

Verification of solutions

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{24}x^4 + \frac{1}{720}x^6\right) y(0) + \left(x + \frac{1}{6}x^3 + \frac{1}{120}x^5\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{24}x^4\right) c_1 + \left(x + \frac{1}{6}x^3 + \frac{1}{120}x^5\right) c_2 + O(x^6)$$

Verified OK.

13.1.1 Maple step by step solution

Let's solve

$$y'' = y$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - y = 0$$

- Characteristic polynomial of ODE
 $r^2 - 1 = 0$
- Factor the characteristic polynomial
 $(r - 1)(r + 1) = 0$
- Roots of the characteristic polynomial
 $r = (-1, 1)$
- 1st solution of the ODE
 $y_1(x) = e^{-x}$
- 2nd solution of the ODE
 $y_2(x) = e^x$
- General solution of the ODE
 $y = c_1 y_1(x) + c_2 y_2(x)$
- Substitute in solutions
 $y = c_1 e^{-x} + c_2 e^x$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 34

```

Order:=6;
dsolve(diff(y(x),x$2)-y(x)=0,y(x),type='series',x=0);

```

$$y(x) = \left(1 + \frac{1}{2}x^2 + \frac{1}{24}x^4\right) y(0) + \left(x + \frac{1}{6}x^3 + \frac{1}{120}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 42

```
AsymptoticDSolveValue[y''[x]-y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(\frac{x^5}{120} + \frac{x^3}{6} + x \right) + c_1 \left(\frac{x^4}{24} + \frac{x^2}{2} + 1 \right)$$

13.2 problem 2

13.2.1 Maple step by step solution 3127

Internal problem ID [712]

Internal file name [OUTPUT/712_Sunday_June_05_2022_01_47_35_AM_89001034/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 2.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact linear second order ode", "second_order_integrable_as_is", "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order, _exact, _linear, _homogeneous]]
```

$$y'' - y'x - y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (747)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (748)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= y'x + y \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= y'x^2 + yx + 2y' \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= (x^3 + 5x)y' + y(x^2 + 3) \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= (x^4 + 9x^2 + 8)y' + xy(x^2 + 7) \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= (x^5 + 14x^3 + 33x)y' + y(x^4 + 12x^2 + 15)
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= y(0) \\
 F_1 &= 2y'(0) \\
 F_2 &= 3y(0) \\
 F_3 &= 8y'(0) \\
 F_4 &= 15y(0)
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4 + \frac{1}{48}x^6\right)y(0) + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right)y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard

power series Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \sum_{n=1}^{\infty} (-n x^n a_n) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \sum_{n=1}^{\infty} (-n x^n a_n) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 - a_0 = 0$$

$$a_2 = \frac{a_0}{2}$$

For $1 \leq n$, the recurrence equation is

$$(n + 2) a_{n+2}(n + 1) - n a_n - a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = \frac{a_n}{n + 2} \quad (5)$$

For $n = 1$ the recurrence equation gives

$$6a_3 - 2a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = \frac{a_1}{3}$$

For $n = 2$ the recurrence equation gives

$$12a_4 - 3a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{8}$$

For $n = 3$ the recurrence equation gives

$$20a_5 - 4a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_1}{15}$$

For $n = 4$ the recurrence equation gives

$$30a_6 - 5a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{a_0}{48}$$

For $n = 5$ the recurrence equation gives

$$42a_7 - 6a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{a_1}{105}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x + \frac{1}{2} a_0 x^2 + \frac{1}{3} a_1 x^3 + \frac{1}{8} a_0 x^4 + \frac{1}{15} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) a_0 + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) c_1 + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4 + \frac{1}{48}x^6\right) y(0) + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) c_1 + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4 + \frac{1}{48}x^6\right) y(0) + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) c_1 + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) c_2 + O(x^6)$$

Verified OK.

13.2.1 Maple step by step solution

Let's solve

$$y'' = y'x + y$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - y'x - y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite DE with series expansions

- Convert $x \cdot y'$ to series expansion

$$x \cdot y' = \sum_{k=0}^{\infty} a_k k x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k-2 \rightarrow k+2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2} (k+2)(k+1) x^k$$

Rewrite DE with series expansions

$$\sum_{k=0}^{\infty} (a_{k+2}(k+2)(k+1) - a_k(k+1)) x^k = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(k+1)(a_{k+2}(k+2) - a_k) = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = \frac{a_k}{k+2} \right]$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
<- linear_1 successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 34

```
Order:=6;  
dsolve(diff(y(x),x$2)-x*diff(y(x),x)-y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) y(0) + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 42

```
AsymptoticDSolveValue[y''[x]-x*y'[x]-y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(\frac{x^5}{15} + \frac{x^3}{3} + x \right) + c_1 \left(\frac{x^4}{8} + \frac{x^2}{2} + 1 \right)$$

13.3 problem 4

13.3.1 Maple step by step solution 3136

Internal problem ID [713]

Internal file name [OUTPUT/713_Sunday_June_05_2022_01_47_36_AM_6480672/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 4.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second_order_bessel_ode", "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_Emden, _Fowler]]
```

$$y'' + k^2 x^2 y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0) y'(x_0) + \frac{(x - x_0)^2}{2} y''(x_0) + \frac{(x - x_0)^3}{3!} y'''(x_0) + \cdots \\ &= y_0 + x y'_0 + \frac{x^2}{2} f|_{x_0, y_0, y'_0} + \frac{x^3}{3!} f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + x y'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (750)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (751)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -k^2 x^2 y \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\
 &= -k^2 x(2y + y'x) \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\
 &= k^2(x^4 k^2 y - 4y'x - 2y) \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\
 &= ((x^4 k^2 - 6)y' + 8x^3 k^2 y) k^2 \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\
 &= -(-12y'x + y(x^4 k^2 - 30)) k^4 x^2
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= 0 \\
 F_1 &= 0 \\
 F_2 &= -2k^2 y(0) \\
 F_3 &= -6y'(0) k^2 \\
 F_4 &= 0
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 - \frac{x^4 k^2}{12}\right) y(0) + \left(x - \frac{1}{20} k^2 x^5\right) y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard

power series Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = -k^2 x^2 \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=0}^{\infty} x^{n+2} k^2 a_n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

$$\sum_{n=0}^{\infty} x^{n+2} k^2 a_n = \sum_{n=2}^{\infty} a_{n-2} k^2 x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \left(\sum_{n=2}^{\infty} a_{n-2} k^2 x^n \right) = 0 \quad (3)$$

For $2 \leq n$, the recurrence equation is

$$(n+2) a_{n+2} (n+1) + a_{n-2} k^2 = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_{n-2}k^2}{(n+2)(n+1)} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$a_0k^2 + 12a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = -\frac{a_0k^2}{12}$$

For $n = 3$ the recurrence equation gives

$$a_1k^2 + 20a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = -\frac{a_1k^2}{20}$$

For $n = 4$ the recurrence equation gives

$$a_2k^2 + 30a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = 0$$

For $n = 5$ the recurrence equation gives

$$a_3k^2 + 42a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = 0$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - \frac{1}{12} a_0 k^2 x^4 - \frac{1}{20} a_1 k^2 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - \frac{x^4 k^2}{12}\right) a_0 + \left(x - \frac{1}{20} k^2 x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - \frac{x^4 k^2}{12}\right) c_1 + \left(x - \frac{1}{20} k^2 x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{x^4 k^2}{12}\right) y(0) + \left(x - \frac{1}{20} k^2 x^5\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 - \frac{x^4 k^2}{12}\right) c_1 + \left(x - \frac{1}{20} k^2 x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 - \frac{x^4 k^2}{12}\right) y(0) + \left(x - \frac{1}{20} k^2 x^5\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 - \frac{x^4 k^2}{12}\right) c_1 + \left(x - \frac{1}{20} k^2 x^5\right) c_2 + O(x^6)$$

Verified OK.

13.3.1 Maple step by step solution

Let's solve

$$y'' = -k^2 x^2 y$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + k^2 x^2 y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite ODE with series expansions

- Convert $x^2 \cdot y$ to series expansion

$$x^2 \cdot y = \sum_{k=0}^{\infty} a_k x^{k+2}$$

- Shift index using $k \rightarrow k - 2$

$$x^2 \cdot y = \sum_{k=2}^{\infty} a_{k-2} x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k \rightarrow k + 2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2} (k+2)(k+1) x^k$$

Rewrite ODE with series expansions

$$6a_3 x + 2a_2 + \left(\sum_{k=2}^{\infty} (a_{k+2} (k+2)(k+1) + k^2 a_{k-2}) x^k \right) = 0$$

- The coefficients of each power of x must be 0

$$[2a_2 = 0, 6a_3 = 0]$$

- Solve for the dependent coefficient(s)

$$\{a_2 = 0, a_3 = 0\}$$

- Each term in the series must be 0, giving the recursion relation

$$(k^2 + 3k + 2) a_{k+2} + k^2 a_{k-2} = 0$$

- Shift index using $k \rightarrow k + 2$
 $((k + 2)^2 + 3k + 8) a_{k+4} + k^2 a_k = 0$
- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+4} = -\frac{k^2 a_k}{k^2 + 7k + 12}, a_2 = 0, a_3 = 0 \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
    -> Bessel
    <- Bessel successful
<- special function solution successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 30

```

Order:=6;
dsolve(diff(y(x),x$2)+k^2*x^2*y(x)=0,y(x),type='series',x=0);

```

$$y(x) = \left(1 - \frac{k^2 x^4}{12}\right) y(0) + \left(x - \frac{1}{20} k^2 x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 34

```

AsymptoticDSolveValue[y''[x]+k^2*x^2*y[x]==0,y[x],{x,0,5}]

```

$$y(x) \rightarrow c_2 \left(x - \frac{k^2 x^5}{20}\right) + c_1 \left(1 - \frac{k^2 x^4}{12}\right)$$

13.4 problem 5

Internal problem ID [714]

Internal file name [OUTPUT/714_Sunday_June_05_2022_01_47_38_AM_69187182/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 5.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_Emden, _Fowler]]
```

$$(1 - x)y'' + y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \dots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (753)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (754)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= \frac{y}{x-1} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\
 &= \frac{(x-1)y' - y}{(x-1)^2} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\
 &= \frac{(-2x+2)y' + (x+1)y}{(x-1)^3} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\
 &= \frac{(x^2+4x-5)y' + (-4x-2)y}{(x-1)^4} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\
 &= \frac{(-6x^2-12x+18)y' + y(x^2+16x+7)}{(x-1)^5}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -y(0) \\
 F_1 &= -y'(0) - y(0) \\
 F_2 &= -y(0) - 2y'(0) \\
 F_3 &= -2y(0) - 5y'(0) \\
 F_4 &= -7y(0) - 18y'(0)
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 - \frac{1}{2}x^2 - \frac{1}{6}x^3 - \frac{1}{24}x^4 - \frac{1}{60}x^5 - \frac{7}{720}x^6\right) y(0) + \left(x - \frac{1}{6}x^3 - \frac{1}{12}x^4 - \frac{1}{24}x^5 - \frac{1}{40}x^6\right) y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series The ode is normalized to be

$$(1 - x)y'' + y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(1 - x) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\sum_{n=2}^{\infty} (-n x^{n-1} a_n (n-1)) + \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} (-n x^{n-1} a_n (n-1)) = \sum_{n=1}^{\infty} (-(n+1) a_{n+1} n x^n)$$

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\sum_{n=1}^{\infty} (-(n+1) a_{n+1} n x^n) + \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 + a_0 = 0$$

$$a_2 = -\frac{a_0}{2}$$

For $1 \leq n$, the recurrence equation is

$$-(n+1) a_{n+1} n + (n+2) a_{n+2} (n+1) + a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} a_{n+2} &= \frac{n^2 a_{n+1} + n a_{n+1} - a_n}{(n+2)(n+1)} \\ (5) \quad &= -\frac{a_n}{(n+2)(n+1)} + \frac{(n^2 + n) a_{n+1}}{(n+2)(n+1)} \end{aligned}$$

For $n = 1$ the recurrence equation gives

$$-2a_2 + 6a_3 + a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = -\frac{a_0}{6} - \frac{a_1}{6}$$

For $n = 2$ the recurrence equation gives

$$-6a_3 + 12a_4 + a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = -\frac{a_0}{24} - \frac{a_1}{12}$$

For $n = 3$ the recurrence equation gives

$$-12a_4 + 20a_5 + a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = -\frac{a_0}{60} - \frac{a_1}{24}$$

For $n = 4$ the recurrence equation gives

$$-20a_5 + 30a_6 + a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{7a_0}{720} - \frac{a_1}{40}$$

For $n = 5$ the recurrence equation gives

$$-30a_6 + 42a_7 + a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{11a_0}{1680} - \frac{17a_1}{1008}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - \frac{a_0 x^2}{2} + \left(-\frac{a_0}{6} - \frac{a_1}{6}\right) x^3 + \left(-\frac{a_0}{24} - \frac{a_1}{12}\right) x^4 + \left(-\frac{a_0}{60} - \frac{a_1}{24}\right) x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - \frac{1}{2}x^2 - \frac{1}{6}x^3 - \frac{1}{24}x^4 - \frac{1}{60}x^5\right) a_0 + \left(x - \frac{1}{6}x^3 - \frac{1}{12}x^4 - \frac{1}{24}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - \frac{1}{2}x^2 - \frac{1}{6}x^3 - \frac{1}{24}x^4 - \frac{1}{60}x^5\right) c_1 + \left(x - \frac{1}{6}x^3 - \frac{1}{12}x^4 - \frac{1}{24}x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{1}{2}x^2 - \frac{1}{6}x^3 - \frac{1}{24}x^4 - \frac{1}{60}x^5 - \frac{7}{720}x^6\right) y(0) + \left(x - \frac{1}{6}x^3 - \frac{1}{12}x^4 - \frac{1}{24}x^5 - \frac{1}{40}x^6\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 - \frac{1}{2}x^2 - \frac{1}{6}x^3 - \frac{1}{24}x^4 - \frac{1}{60}x^5\right) c_1 + \left(x - \frac{1}{6}x^3 - \frac{1}{12}x^4 - \frac{1}{24}x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 - \frac{1}{2}x^2 - \frac{1}{6}x^3 - \frac{1}{24}x^4 - \frac{1}{60}x^5 - \frac{7}{720}x^6\right) y(0) + \left(x - \frac{1}{6}x^3 - \frac{1}{12}x^4 - \frac{1}{24}x^5 - \frac{1}{40}x^6\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 - \frac{1}{2}x^2 - \frac{1}{6}x^3 - \frac{1}{24}x^4 - \frac{1}{60}x^5\right) c_1 + \left(x - \frac{1}{6}x^3 - \frac{1}{12}x^4 - \frac{1}{24}x^5\right) c_2 + O(x^6)$$

Verified OK.

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
  -> Bessel
  -> elliptic
  -> Legendre
  -> Whittaker
      -> hyper3: Equivalence to 1F1 under a power @ Moebius
  -> hypergeometric
      -> heuristic approach
          <- heuristic approach successful
          <- hypergeometric successful
<- special function solution successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 49

```
Order:=6;
dsolve((1-x)*diff(y(x),x$2)+y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 - \frac{1}{2}x^2 - \frac{1}{6}x^3 - \frac{1}{24}x^4 - \frac{1}{60}x^5\right) y(0) + \left(x - \frac{1}{6}x^3 - \frac{1}{12}x^4 - \frac{1}{24}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 63

```
AsymptoticDSolveValue[(1-x)*y''[x]+y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(-\frac{x^5}{24} - \frac{x^4}{12} - \frac{x^3}{6} + x \right) + c_1 \left(-\frac{x^5}{60} - \frac{x^4}{24} - \frac{x^3}{6} - \frac{x^2}{2} + 1 \right)$$

13.5 problem 6

Internal problem ID [715]

Internal file name [OUTPUT/715_Sunday_June_05_2022_01_47_39_AM_1741303/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 6.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(x^2 + 2)y'' - y'x + 4y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (756)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (757)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= \frac{y'x - 4y}{x^2 + 2} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\
 &= \frac{-4y'x^2 + 4yx - 6y'}{(x^2 + 2)^2} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\
 &= \frac{8x^3y' + 4x^2y + 10y'x + 32y}{(x^2 + 2)^3} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\
 &= \frac{(-12x^4 + 48x^2 + 84)y' + (-48x^3 - 216x)y}{(x^2 + 2)^4} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\
 &= \frac{(-12x^5 - 648x^3 - 828x)y' + (288x^4 + 1032x^2 - 768)y}{(x^2 + 2)^5}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -2y(0) \\
 F_1 &= -\frac{3y'(0)}{2} \\
 F_2 &= 4y(0) \\
 F_3 &= \frac{21y'(0)}{4} \\
 F_4 &= -24y(0)
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 - x^2 + \frac{1}{6}x^4 - \frac{1}{30}x^6\right) y(0) + \left(x - \frac{1}{4}x^3 + \frac{7}{160}x^5\right) y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. The ode is normalized to be

$$(x^2 + 2)y'' - y'x + 4y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(x^2 + 2) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) - \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + 4 \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} x^n a_n n(n-1) \right) + \left(\sum_{n=2}^{\infty} 2n(n-1) a_n x^{n-2} \right) + \sum_{n=1}^{\infty} (-n a_n x^n) + \left(\sum_{n=0}^{\infty} 4a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} 2n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} 2(n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} x^n a_n n(n-1) \right) + \left(\sum_{n=0}^{\infty} 2(n+2) a_{n+2} (n+1) x^n \right) \\ & + \sum_{n=1}^{\infty} (-n a_n x^n) + \left(\sum_{n=0}^{\infty} 4a_n x^n \right) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$4a_2 + 4a_0 = 0$$

$$a_2 = -a_0$$

$n = 1$ gives

$$12a_3 + 3a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{a_1}{4}$$

For $2 \leq n$, the recurrence equation is

$$n a_n (n-1) + 2(n+2) a_{n+2} (n+1) - n a_n + 4a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n (n^2 - 2n + 4)}{2(n+2)(n+1)} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$4a_2 + 24a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{6}$$

For $n = 3$ the recurrence equation gives

$$7a_3 + 40a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{7a_1}{160}$$

For $n = 4$ the recurrence equation gives

$$12a_4 + 60a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{a_0}{30}$$

For $n = 5$ the recurrence equation gives

$$19a_5 + 84a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{19a_1}{1920}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - a_0 x^2 - \frac{1}{4} a_1 x^3 + \frac{1}{6} a_0 x^4 + \frac{7}{160} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - x^2 + \frac{1}{6}x^4\right) a_0 + \left(x - \frac{1}{4}x^3 + \frac{7}{160}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - x^2 + \frac{1}{6}x^4\right) c_1 + \left(x - \frac{1}{4}x^3 + \frac{7}{160}x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - x^2 + \frac{1}{6}x^4 - \frac{1}{30}x^6\right) y(0) + \left(x - \frac{1}{4}x^3 + \frac{7}{160}x^5\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 - x^2 + \frac{1}{6}x^4\right) c_1 + \left(x - \frac{1}{4}x^3 + \frac{7}{160}x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 - x^2 + \frac{1}{6}x^4 - \frac{1}{30}x^6\right) y(0) + \left(x - \frac{1}{4}x^3 + \frac{7}{160}x^5\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 - x^2 + \frac{1}{6}x^4\right) c_1 + \left(x - \frac{1}{4}x^3 + \frac{7}{160}x^5\right) c_2 + O(x^6)$$

Verified OK.

Maple trace **Kovacic algorithm successful**

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Group is reducible or imprimitive
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 34

```
Order:=6;  
dsolve((2+x^2)*diff(y(x),x$2)-x*diff(y(x),x)+4*y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 - x^2 + \frac{1}{6}x^4\right) y(0) + \left(x - \frac{1}{4}x^3 + \frac{7}{160}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 40

```
AsymptoticDSolveValue[(2+x^2)*y'[x]-x*y'[x]+4*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(\frac{7x^5}{160} - \frac{x^3}{4} + x \right) + c_1 \left(\frac{x^4}{6} - x^2 + 1 \right)$$

13.6 problem 7

13.6.1 Maple step by step solution 3163

Internal problem ID [716]

Internal file name [OUTPUT/716_Sunday_June_05_2022_01_47_41_AM_17236857/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 7.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$y'' + y'x + 2y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0) y'(x_0) + \frac{(x - x_0)^2}{2} y''(x_0) + \frac{(x - x_0)^3}{3!} y'''(x_0) + \cdots \\ &= y_0 + x y'_0 + \frac{x^2}{2} f|_{x_0, y_0, y'_0} + \frac{x^3}{3!} f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + x y'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (759)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (760)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -y'x - 2y \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= y'x^2 + 2yx - 3y' \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= -x^3y' - 2x^2y + 7y'x + 8y \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= (x^4 - 12x^2 + 15)y' + 2(x^3 - 9x)y \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= (-x^5 + 18x^3 - 57x)y' - 2y(x^4 - 15x^2 + 24)
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -2y(0) \\
 F_1 &= -3y'(0) \\
 F_2 &= 8y(0) \\
 F_3 &= 15y'(0) \\
 F_4 &= -48y(0)
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 - x^2 + \frac{1}{3}x^4 - \frac{1}{15}x^6\right)y(0) + \left(x - \frac{1}{2}x^3 + \frac{1}{8}x^5\right)y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard

power series Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = - \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x - 2 \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n x^n a_n \right) + \left(\sum_{n=0}^{\infty} 2 a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \left(\sum_{n=1}^{\infty} n x^n a_n \right) + \left(\sum_{n=0}^{\infty} 2 a_n x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 + 2a_0 = 0$$

$$a_2 = -a_0$$

For $1 \leq n$, the recurrence equation is

$$(n + 2) a_{n+2}(n + 1) + na_n + 2a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n}{n + 1} \quad (5)$$

For $n = 1$ the recurrence equation gives

$$6a_3 + 3a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = -\frac{a_1}{2}$$

For $n = 2$ the recurrence equation gives

$$12a_4 + 4a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{3}$$

For $n = 3$ the recurrence equation gives

$$20a_5 + 5a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_1}{8}$$

For $n = 4$ the recurrence equation gives

$$30a_6 + 6a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{a_0}{15}$$

For $n = 5$ the recurrence equation gives

$$42a_7 + 7a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{a_1}{48}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - a_0 x^2 - \frac{1}{2} a_1 x^3 + \frac{1}{3} a_0 x^4 + \frac{1}{8} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - x^2 + \frac{1}{3}x^4\right) a_0 + \left(x - \frac{1}{2}x^3 + \frac{1}{8}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - x^2 + \frac{1}{3}x^4\right) c_1 + \left(x - \frac{1}{2}x^3 + \frac{1}{8}x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - x^2 + \frac{1}{3}x^4 - \frac{1}{15}x^6\right) y(0) + \left(x - \frac{1}{2}x^3 + \frac{1}{8}x^5\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 - x^2 + \frac{1}{3}x^4\right) c_1 + \left(x - \frac{1}{2}x^3 + \frac{1}{8}x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 - x^2 + \frac{1}{3}x^4 - \frac{1}{15}x^6\right) y(0) + \left(x - \frac{1}{2}x^3 + \frac{1}{8}x^5\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 - x^2 + \frac{1}{3}x^4\right) c_1 + \left(x - \frac{1}{2}x^3 + \frac{1}{8}x^5\right) c_2 + O(x^6)$$

Verified OK.

13.6.1 Maple step by step solution

Let's solve

$$y'' = -y'x - 2y$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + y'x + 2y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite DE with series expansions

- Convert $x \cdot y'$ to series expansion

$$x \cdot y' = \sum_{k=0}^{\infty} a_k k x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k- > k+2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2}(k+2)(k+1) x^k$$

Rewrite DE with series expansions

$$\sum_{k=0}^{\infty} (a_{k+2}(k+2)(k+1) + a_k(k+2)) x^k = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(k+2)(ka_{k+2} + a_k + a_{k+2}) = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = -\frac{a_k}{k+1} \right]$$

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacics algorithm  
  A Liouvillian solution exists  
  Reducible group (found an exponential solution)  
  Group is reducible, not completely reducible  
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 34

```
Order:=6;  
dsolve(diff(y(x),x$2)+x*diff(y(x),x)+2*y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 - x^2 + \frac{1}{3}x^4\right) y(0) + \left(x - \frac{1}{2}x^3 + \frac{1}{8}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 40

```
AsymptoticDSolveValue[y''[x]+x*y'[x]+2*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(\frac{x^5}{8} - \frac{x^3}{2} + x \right) + c_1 \left(\frac{x^4}{3} - x^2 + 1 \right)$$

13.7 problem 9

Internal problem ID [717]

Internal file name [OUTPUT/717_Sunday_June_05_2022_01_47_42_AM_81903652/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 9.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(x^2 + 1)y'' - 4y'x + 6y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (762)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (763)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= \frac{4y'x - 6y}{x^2 + 1} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= \frac{6y'x^2 - 12yx - 2y'}{(x^2 + 1)^2} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= 0 \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= 0 \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= 0
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -6y(0) \\
 F_1 &= -2y'(0) \\
 F_2 &= 0 \\
 F_3 &= 0 \\
 F_4 &= 0
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = (-3x^2 + 1)y(0) + \left(-\frac{1}{3}x^3 + x\right)y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. The ode is normalized to be

$$(x^2 + 1) y'' - 4y'x + 6y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(x^2 + 1) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) - 4 \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + 6 \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} x^n a_n n(n-1) \right) + \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \sum_{n=1}^{\infty} (-4n a_n x^n) + \left(\sum_{n=0}^{\infty} 6a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=2}^{\infty} x^n a_n n(n-1) \right) + \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \sum_{n=1}^{\infty} (-4n a_n x^n) + \left(\sum_{n=0}^{\infty} 6a_n x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 + 6a_0 = 0$$

$$a_2 = -3a_0$$

$n = 1$ gives

$$6a_3 + 2a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{a_1}{3}$$

For $2 \leq n$, the recurrence equation is

$$na_n(n-1) + (n+2)a_{n+2}(n+1) - 4na_n + 6a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n(n^2 - 5n + 6)}{(n+2)(n+1)} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$12a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = 0$$

For $n = 3$ the recurrence equation gives

$$20a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = 0$$

For $n = 4$ the recurrence equation gives

$$2a_4 + 30a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = 0$$

For $n = 5$ the recurrence equation gives

$$6a_5 + 42a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = 0$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - 3a_0 x^2 - \frac{1}{3} a_1 x^3 + \dots$$

Collecting terms, the solution becomes

$$y = (-3x^2 + 1) a_0 + \left(-\frac{1}{3} x^3 + x \right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = (-3x^2 + 1) c_1 + \left(-\frac{1}{3} x^3 + x \right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = (-3x^2 + 1) y(0) + \left(-\frac{1}{3} x^3 + x \right) y'(0) + O(x^6) \quad (1)$$

$$y = (-3x^2 + 1) c_1 + \left(-\frac{1}{3} x^3 + x \right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = (-3x^2 + 1) y(0) + \left(-\frac{1}{3}x^3 + x\right) y'(0) + O(x^6)$$

Verified OK.

$$y = (-3x^2 + 1) c_1 + \left(-\frac{1}{3}x^3 + x\right) c_2 + O(x^6)$$

Verified OK.

Maple trace **Kovacic algorithm successful**

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Reducible group (found an exponential solution)
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 25

```
Order:=6;
dsolve((1+x^2)*diff(y(x),x$2)-4*x*diff(y(x),x)+6*y(x)=0,y(x),type='series',x=0);
```

$$y(x) = y(0) + D(y)(0)x - 3y(0)x^2 - \frac{D(y)(0)x^3}{3}$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 26

```
AsymptoticDSolveValue[(1+x^2)*y'[x]-4*x*y'[x]+6*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(x - \frac{x^3}{3}\right) + c_1(1 - 3x^2)$$

13.8 problem 10

Internal problem ID [718]

Internal file name [OUTPUT/718_Sunday_June_05_2022_01_47_43_AM_81935409/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 10.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact linear second order ode", "second_order_integrable_as_is", "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _exact , _linear , _homogeneous]]
```

$$(-x^2 + 4)y'' + 2y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (765)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (766)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= \frac{2y}{x^2 - 4} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\
 &= \frac{2y'x^2 - 4yx - 8y'}{(x^2 - 4)^2} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\
 &= -\frac{8(y'x^2 - 2yx - 4y')x}{(x^2 - 4)^3} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\
 &= \frac{40(x^2 + \frac{4}{5})(-2yx + (x^2 - 4)y')}{(x^2 - 4)^4} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\
 &= -\frac{240(x^2 + \frac{12}{5})(-2yx + (x^2 - 4)y')x}{(x^2 - 4)^5}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -\frac{y(0)}{2} \\
 F_1 &= -\frac{y'(0)}{2} \\
 F_2 &= 0 \\
 F_3 &= -\frac{y'(0)}{2} \\
 F_4 &= 0
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 - \frac{x^2}{4}\right) y(0) + \left(x - \frac{1}{12}x^3 - \frac{1}{240}x^5\right) y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series The ode is normalized to be

$$(-x^2 + 4) y'' + 2y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(-x^2 + 4) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + 2 \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\sum_{n=2}^{\infty} (-x^n a_n n(n-1)) + \left(\sum_{n=2}^{\infty} 4n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=0}^{\infty} 2a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} 4n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} 4(n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\sum_{n=2}^{\infty} (-x^n a_n n(n-1)) + \left(\sum_{n=0}^{\infty} 4(n+2) a_{n+2} (n+1) x^n \right) + \left(\sum_{n=0}^{\infty} 2a_n x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$8a_2 + 2a_0 = 0$$

$$a_2 = -\frac{a_0}{4}$$

$n = 1$ gives

$$24a_3 + 2a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{a_1}{12}$$

For $2 \leq n$, the recurrence equation is

$$-na_n(n-1) + 4(n+2)a_{n+2}(n+1) + 2a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = \frac{(n-2)a_n}{4n+8} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$48a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = 0$$

For $n = 3$ the recurrence equation gives

$$-4a_3 + 80a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = -\frac{a_1}{240}$$

For $n = 4$ the recurrence equation gives

$$-10a_4 + 120a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = 0$$

For $n = 5$ the recurrence equation gives

$$-18a_5 + 168a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{a_1}{2240}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - \frac{1}{4} a_0 x^2 - \frac{1}{12} a_1 x^3 - \frac{1}{240} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - \frac{x^2}{4}\right) a_0 + \left(x - \frac{1}{12} x^3 - \frac{1}{240} x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - \frac{x^2}{4}\right) c_1 + \left(x - \frac{1}{12} x^3 - \frac{1}{240} x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{x^2}{4}\right) y(0) + \left(x - \frac{1}{12} x^3 - \frac{1}{240} x^5\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 - \frac{x^2}{4}\right) c_1 + \left(x - \frac{1}{12} x^3 - \frac{1}{240} x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 - \frac{x^2}{4}\right) y(0) + \left(x - \frac{1}{12}x^3 - \frac{1}{240}x^5\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 - \frac{x^2}{4}\right) c_1 + \left(x - \frac{1}{12}x^3 - \frac{1}{240}x^5\right) c_2 + O(x^6)$$

Verified OK.

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
<- linear_1 successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 29

```
Order:=6;  
dsolve((4-x^2)*diff(y(x),x$2)+2*y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 - \frac{x^2}{4}\right) y(0) + \left(x - \frac{1}{12}x^3 - \frac{1}{240}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 35

```
AsymptoticDSolveValue[(4-x^2)*y'[x]+2*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_1 \left(1 - \frac{x^2}{4}\right) + c_2 \left(-\frac{x^5}{240} - \frac{x^3}{12} + x\right)$$

13.9 problem 11

Internal problem ID [719]

Internal file name [OUTPUT/719_Sunday_June_05_2022_01_47_44_AM_50694559/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 11.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact linear second order ode", "second_order_integrable_as_is", "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _exact , _linear , _homogeneous]]
```

$$(-x^2 + 3)y'' - 3y'x - y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (768)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (769)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -\frac{3y'x + y}{x^2 - 3} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= \frac{11y'x^2 + 5yx + 12y'}{(x^2 - 3)^2} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= \frac{-50x^3y' - 26x^2y - 165y'x - 27y}{(x^2 - 3)^3} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= \frac{(274x^4 + 1821x^2 + 576)y' + (154x^3 + 483x)y}{(x^2 - 3)^4} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= \frac{(-1764x^5 - 19656x^3 - 18711x)y' + (-1044x^4 - 6588x^2 - 2025)y}{(x^2 - 3)^5}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= \frac{y(0)}{3} \\
 F_1 &= \frac{4y'(0)}{3} \\
 F_2 &= y(0) \\
 F_3 &= \frac{64y'(0)}{9} \\
 F_4 &= \frac{25y(0)}{3}
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 + \frac{1}{6}x^2 + \frac{1}{24}x^4 + \frac{5}{432}x^6\right)y(0) + \left(x + \frac{2}{9}x^3 + \frac{8}{135}x^5\right)y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. The ode is normalized to be

$$(-x^2 + 3)y'' - 3y'x - y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(-x^2 + 3) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) - 3 \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x - \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\sum_{n=2}^{\infty} (-x^n a_n n(n-1)) + \left(\sum_{n=2}^{\infty} 3n(n-1) a_n x^{n-2} \right) + \sum_{n=1}^{\infty} (-3n a_n x^n) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} 3n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} 3(n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} \sum_{n=2}^{\infty} (-x^n a_n n(n-1)) + \left(\sum_{n=0}^{\infty} 3(n+2) a_{n+2} (n+1) x^n \right) \\ + \sum_{n=1}^{\infty} (-3n a_n x^n) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$6a_2 - a_0 = 0$$

$$a_2 = \frac{a_0}{6}$$

$n = 1$ gives

$$18a_3 - 4a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = \frac{2a_1}{9}$$

For $2 \leq n$, the recurrence equation is

$$-na_n(n-1) + 3(n+2)a_{n+2}(n+1) - 3na_n - a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = \frac{(n+1)a_n}{3n+6} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$-9a_2 + 36a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{24}$$

For $n = 3$ the recurrence equation gives

$$-16a_3 + 60a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{8a_1}{135}$$

For $n = 4$ the recurrence equation gives

$$-25a_4 + 90a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{5a_0}{432}$$

For $n = 5$ the recurrence equation gives

$$-36a_5 + 126a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{16a_1}{945}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x + \frac{1}{6} a_0 x^2 + \frac{2}{9} a_1 x^3 + \frac{1}{24} a_0 x^4 + \frac{8}{135} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + \frac{1}{6}x^2 + \frac{1}{24}x^4\right) a_0 + \left(x + \frac{2}{9}x^3 + \frac{8}{135}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 + \frac{1}{6}x^2 + \frac{1}{24}x^4\right) c_1 + \left(x + \frac{2}{9}x^3 + \frac{8}{135}x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 + \frac{1}{6}x^2 + \frac{1}{24}x^4 + \frac{5}{432}x^6\right) y(0) + \left(x + \frac{2}{9}x^3 + \frac{8}{135}x^5\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 + \frac{1}{6}x^2 + \frac{1}{24}x^4\right) c_1 + \left(x + \frac{2}{9}x^3 + \frac{8}{135}x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 + \frac{1}{6}x^2 + \frac{1}{24}x^4 + \frac{5}{432}x^6\right) y(0) + \left(x + \frac{2}{9}x^3 + \frac{8}{135}x^5\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 + \frac{1}{6}x^2 + \frac{1}{24}x^4\right) c_1 + \left(x + \frac{2}{9}x^3 + \frac{8}{135}x^5\right) c_2 + O(x^6)$$

Verified OK.

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
<- linear_1 successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 34

```
Order:=6;  
dsolve((3-x^2)*diff(y(x),x$2)-3*x*diff(y(x),x)-y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 + \frac{1}{6}x^2 + \frac{1}{24}x^4\right) y(0) + \left(x + \frac{2}{9}x^3 + \frac{8}{135}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 70

```
AsymptoticDSolveValue[(3-x^2)*y''[x]-3*y'[x]-y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_1 \left(\frac{13x^5}{1080} + \frac{x^4}{36} + \frac{x^3}{18} + \frac{x^2}{6} + 1 \right) + c_2 \left(\frac{49x^5}{1080} + \frac{7x^4}{72} + \frac{2x^3}{9} + \frac{x^2}{2} + x \right)$$

13.10 problem 12

13.10.1 Maple step by step solution 3197

Internal problem ID [720]

Internal file name [OUTPUT/720_Sunday_June_05_2022_01_47_45_AM_90164067/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 12.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second order series method. Ordinary point", "second_order_change_of_variable_on_y_method_2", "second order series method. Taylor series method", "second_order_ode_non_constant_coeff_transformation_on_B"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(1 - x)y'' + y'x - y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using

Taylor series gives

$$\begin{aligned}
 y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\
 &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \dots \\
 &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0}
 \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \tag{1}$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \tag{771}$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \tag{772}$$

$$\begin{aligned}
 \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\
 &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \tag{2}
 \end{aligned}$$

$$\begin{aligned}
 \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\
 &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \tag{3}
 \end{aligned}$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= \frac{-y + y'x}{x - 1} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= \frac{-y + y'x}{x - 1} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= \frac{-y + y'x}{x - 1} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= \frac{-y + y'x}{x - 1} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= \frac{-y + y'x}{x - 1}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$F_0 = y(0)$$

$$F_1 = y(0)$$

$$F_2 = y(0)$$

$$F_3 = y(0)$$

$$F_4 = y(0)$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5 + \frac{1}{720}x^6\right)y(0) + y'(0)x + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. The ode is normalized to be

$$(1 - x)y'' + y'x - y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(1 - x) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x - \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\sum_{n=2}^{\infty} (-n x^{n-1} a_n (n-1)) + \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} (-n x^{n-1} a_n (n-1)) = \sum_{n=1}^{\infty} (-(n+1) a_{n+1} n x^n)$$

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} \sum_{n=1}^{\infty} (-(n+1) a_{n+1} n x^n) + \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) \\ + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$2a_2 - a_0 = 0$$

$$a_2 = \frac{a_0}{2}$$

For $1 \leq n$, the recurrence equation is

$$-(n+1) a_{n+1} n + (n+2) a_{n+2} (n+1) + n a_n - a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} a_{n+2} &= \frac{n^2 a_{n+1} - n a_n + n a_{n+1} + a_n}{(n+2)(n+1)} \\ (5) \quad &= \frac{(-n+1) a_n}{(n+2)(n+1)} + \frac{(n^2+n) a_{n+1}}{(n+2)(n+1)} \end{aligned}$$

For $n = 1$ the recurrence equation gives

$$-2a_2 + 6a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = \frac{a_0}{6}$$

For $n = 2$ the recurrence equation gives

$$-6a_3 + 12a_4 + a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{24}$$

For $n = 3$ the recurrence equation gives

$$-12a_4 + 20a_5 + 2a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_0}{120}$$

For $n = 4$ the recurrence equation gives

$$-20a_5 + 30a_6 + 3a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{a_0}{720}$$

For $n = 5$ the recurrence equation gives

$$-30a_6 + 42a_7 + 4a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{a_0}{5040}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x + \frac{1}{2} a_0 x^2 + \frac{1}{6} a_0 x^3 + \frac{1}{24} a_0 x^4 + \frac{1}{120} a_0 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) a_0 + a_1x + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) c_1 + c_2x + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5 + \frac{1}{720}x^6\right) y(0) + y'(0)x + O(x^6) \quad (1)$$

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) c_1 + c_2x + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5 + \frac{1}{720}x^6\right) y(0) + y'(0)x + O(x^6)$$

Verified OK.

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) c_1 + c_2x + O(x^6)$$

Verified OK.

13.10.1 Maple step by step solution

Let's solve

$$(1 - x)y'' + y'x - y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{y}{x-1} + \frac{xy'}{x-1}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{xy'}{x-1} + \frac{y}{x-1} = 0$$

□ Check to see if $x_0 = 1$ is a regular singular point

○ Define functions

$$\left[P_2(x) = -\frac{x}{x-1}, P_3(x) = \frac{1}{x-1} \right]$$

○ $(x-1) \cdot P_2(x)$ is analytic at $x = 1$

$$\left. ((x-1) \cdot P_2(x)) \right|_{x=1} = -1$$

○ $(x-1)^2 \cdot P_3(x)$ is analytic at $x = 1$

$$\left. ((x-1)^2 \cdot P_3(x)) \right|_{x=1} = 0$$

○ $x = 1$ is a regular singular point

Check to see if $x_0 = 1$ is a regular singular point

$$x_0 = 1$$

• Multiply by denominators

$$(x-1)y'' - y'x + y = 0$$

• Change variables using $x = u + 1$ so that the regular singular point is at $u = 0$

$$u \left(\frac{d^2}{du^2} y(u) \right) + (-u-1) \left(\frac{d}{du} y(u) \right) + y(u) = 0$$

• Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

□ Rewrite ODE with series expansions

○ Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) u^{k+r-1+m}$$

○ Shift index using $k \rightarrow k+1-m$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k+1-m+r) u^{k+r}$$

○ Convert $u \cdot \left(\frac{d^2}{du^2} y(u) \right)$ to series expansion

$$u \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) u^{k+r-1}$$

○ Shift index using $k \rightarrow k+1$

$$u \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=-1}^{\infty} a_{k+1} (k+1+r) (k+r) u^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r (-2+r) u^{-1+r} + \left(\sum_{k=0}^{\infty} (a_{k+1} (k+1+r) (k+r-1) - a_k (k+r-1)) u^{k+r} \right) = 0$$

- a_0 cannot be 0 by assumption, giving the indicial equation
- Values of r that satisfy the indicial equation
- Each term in the series must be 0, giving the recursion relation

$$r(-2+r) = 0$$

$$r \in \{0, 2\}$$

$$(k+r-1)(a_{k+1}(k+1+r) - a_k) = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+1} = \frac{a_k}{k+1+r}$$

- Recursion relation for $r = 0$

$$a_{k+1} = \frac{a_k}{k+1}$$

- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{k+1} = \frac{a_k}{k+1} \right]$$

- Revert the change of variables $u = x - 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x-1)^k, a_{k+1} = \frac{a_k}{k+1} \right]$$

- Recursion relation for $r = 2$

$$a_{k+1} = \frac{a_k}{k+3}$$

- Solution for $r = 2$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+2}, a_{k+1} = \frac{a_k}{k+3} \right]$$

- Revert the change of variables $u = x - 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x-1)^{k+2}, a_{k+1} = \frac{a_k}{k+3} \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k (x-1)^k \right) + \left(\sum_{k=0}^{\infty} b_k (x-1)^{k+2} \right), a_{k+1} = \frac{a_k}{k+1}, b_{k+1} = \frac{b_k}{k+3} \right]$$

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacics algorithm  
  A Liouvillian solution exists  
  Reducible group (found an exponential solution)  
  Reducible group (found another exponential solution)  
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 34

```
Order:=6;  
dsolve((1-x)*diff(y(x),x$2)+x*diff(y(x),x)-y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) y(0) + D(y)(0)x + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 41

```
AsymptoticDSolveValue[(1-x)*y''[x]+x*y'[x]-y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_1 \left(\frac{x^5}{120} + \frac{x^4}{24} + \frac{x^3}{6} + \frac{x^2}{2} + 1 \right) + c_2 x$$

13.11 problem 13

13.11.1 Maple step by step solution 3208

Internal problem ID [721]

Internal file name [OUTPUT/721_Sunday_June_05_2022_01_47_46_AM_4048532/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 13.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$2y'' + y'x + 3y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (774)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (775)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -\frac{y'x}{2} - \frac{3y}{2} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= \frac{y'x^2}{4} + \frac{3yx}{4} - 2y' \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= \frac{(-x^3 + 18x)y'}{8} + \frac{(-3x^2 + 30)y}{8} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= \frac{(x^4 - 30x^2 + 96)y'}{16} + \frac{3xy(x^2 - 22)}{16} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= \frac{(-x^5 + 44x^3 - 348x)y'}{32} - \frac{3y(x^4 - 36x^2 + 140)}{32}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -\frac{3y(0)}{2} \\
 F_1 &= -2y'(0) \\
 F_2 &= \frac{15y(0)}{4} \\
 F_3 &= 6y'(0) \\
 F_4 &= -\frac{105y(0)}{8}
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 - \frac{3}{4}x^2 + \frac{5}{32}x^4 - \frac{7}{384}x^6\right) y(0) + \left(x - \frac{1}{3}x^3 + \frac{1}{20}x^5\right) y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = -\frac{\left(\sum_{n=1}^{\infty} n a_n x^{n-1}\right) x}{2} - \frac{3\left(\sum_{n=0}^{\infty} a_n x^n\right)}{2} \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} 2n(n-1) a_n x^{n-2}\right) + \left(\sum_{n=1}^{\infty} n x^n a_n\right) + \left(\sum_{n=0}^{\infty} 3a_n x^n\right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} 2n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} 2(n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} 2(n+2) a_{n+2} (n+1) x^n\right) + \left(\sum_{n=1}^{\infty} n x^n a_n\right) + \left(\sum_{n=0}^{\infty} 3a_n x^n\right) = 0 \quad (3)$$

$n = 0$ gives

$$4a_2 + 3a_0 = 0$$

$$a_2 = -\frac{3a_0}{4}$$

For $1 \leq n$, the recurrence equation is

$$2(n+2)a_{n+2}(n+1) + na_n + 3a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n(n+3)}{2(n+2)(n+1)} \quad (5)$$

For $n = 1$ the recurrence equation gives

$$12a_3 + 4a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = -\frac{a_1}{3}$$

For $n = 2$ the recurrence equation gives

$$24a_4 + 5a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{5a_0}{32}$$

For $n = 3$ the recurrence equation gives

$$40a_5 + 6a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_1}{20}$$

For $n = 4$ the recurrence equation gives

$$60a_6 + 7a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{7a_0}{384}$$

For $n = 5$ the recurrence equation gives

$$84a_7 + 8a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{a_1}{210}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - \frac{3}{4} a_0 x^2 - \frac{1}{3} a_1 x^3 + \frac{5}{32} a_0 x^4 + \frac{1}{20} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - \frac{3}{4}x^2 + \frac{5}{32}x^4\right) a_0 + \left(x - \frac{1}{3}x^3 + \frac{1}{20}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - \frac{3}{4}x^2 + \frac{5}{32}x^4\right) c_1 + \left(x - \frac{1}{3}x^3 + \frac{1}{20}x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{3}{4}x^2 + \frac{5}{32}x^4 - \frac{7}{384}x^6\right) y(0) + \left(x - \frac{1}{3}x^3 + \frac{1}{20}x^5\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 - \frac{3}{4}x^2 + \frac{5}{32}x^4\right) c_1 + \left(x - \frac{1}{3}x^3 + \frac{1}{20}x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 - \frac{3}{4}x^2 + \frac{5}{32}x^4 - \frac{7}{384}x^6\right) y(0) + \left(x - \frac{1}{3}x^3 + \frac{1}{20}x^5\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 - \frac{3}{4}x^2 + \frac{5}{32}x^4\right) c_1 + \left(x - \frac{1}{3}x^3 + \frac{1}{20}x^5\right) c_2 + O(x^6)$$

Verified OK.

13.11.1 Maple step by step solution

Let's solve

$$y'' = -\frac{y'x}{2} - \frac{3y}{2}$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{y'x}{2} + \frac{3y}{2} = 0$$

- Multiply by denominators

$$2y'' + y'x + 3y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite DE with series expansions

- Convert $x \cdot y'$ to series expansion

$$x \cdot y' = \sum_{k=0}^{\infty} a_k k x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k- \rightarrow k+2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2}(k+2)(k+1) x^k$$

Rewrite DE with series expansions

$$\sum_{k=0}^{\infty} (2a_{k+2}(k+2)(k+1) + a_k(k+3)) x^k = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(2k^2 + 6k + 4) a_{k+2} + a_k(k + 3) = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = -\frac{a_k(k+3)}{2(k^2+3k+2)} \right]$$

Maple trace Kovacic algorithm successful

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
  A Liouvillian solution exists
  Reducible group (found an exponential solution)
  Group is reducible, not completely reducible
  Solution has integrals. Trying a special function solution free of integrals...
-> Trying a solution in terms of special functions:
  -> Bessel
  -> elliptic
  -> Legendre
  <- Kummer successful
<- special function solution successful
  -> Trying to convert hypergeometric functions to elementary form...
  <- elementary form could result into a too large expression - returning special functi
<- Kovacics algorithm successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 34

```
Order:=6;  
dsolve(2*diff(y(x),x$2)+x*diff(y(x),x)+3*y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 - \frac{3}{4}x^2 + \frac{5}{32}x^4\right) y(0) + \left(x - \frac{1}{3}x^3 + \frac{1}{20}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 42

```
AsymptoticDSolveValue[2*y''[x]+x*y'[x]+3*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(\frac{x^5}{20} - \frac{x^3}{3} + x \right) + c_1 \left(\frac{5x^4}{32} - \frac{3x^2}{4} + 1 \right)$$

13.12 problem 15

13.12.1 Existence and uniqueness analysis	3211
13.12.2 Maple step by step solution	3219

Internal problem ID [722]

Internal file name [OUTPUT/722_Sunday_June_05_2022_01_47_48_AM_32860525/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 15.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact linear second order ode", "second_order_integrable_as_is", "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _exact , _linear , _homogeneous]]
```

$$y'' - y'x - y = 0$$

With initial conditions

$$[y(0) = 2, y'(0) = 1]$$

With the expansion point for the power series method at $x = 0$.

13.12.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = -x$$

$$q(x) = -1$$

$$F = 0$$

Hence the ode is

$$y'' - y'x - y = 0$$

The domain of $p(x) = -x$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = -1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \dots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (777)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (778)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= y'x + y \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= y'x^2 + yx + 2y' \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= (x^3 + 5x)y' + y(x^2 + 3) \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= (x^4 + 9x^2 + 8)y' + xy(x^2 + 7) \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= (x^5 + 14x^3 + 33x)y' + y(x^4 + 12x^2 + 15)
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = 2$ and $y'(0) = 1$ gives

$$\begin{aligned}
 F_0 &= 2 \\
 F_1 &= 2 \\
 F_2 &= 6 \\
 F_3 &= 8 \\
 F_4 &= 30
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = x^2 + x + 2 + \frac{x^3}{3} + \frac{x^4}{4} + \frac{x^5}{15} + \frac{x^6}{24} + O(x^6)$$

$$y = x^2 + x + 2 + \frac{x^3}{3} + \frac{x^4}{4} + \frac{x^5}{15} + \frac{x^6}{24} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \sum_{n=1}^{\infty} (-n x^n a_n) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \sum_{n=1}^{\infty} (-n x^n a_n) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 - a_0 = 0$$

$$a_2 = \frac{a_0}{2}$$

For $1 \leq n$, the recurrence equation is

$$(n + 2) a_{n+2}(n + 1) - n a_n - a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = \frac{a_n}{n + 2} \quad (5)$$

For $n = 1$ the recurrence equation gives

$$6a_3 - 2a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = \frac{a_1}{3}$$

For $n = 2$ the recurrence equation gives

$$12a_4 - 3a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{8}$$

For $n = 3$ the recurrence equation gives

$$20a_5 - 4a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_1}{15}$$

For $n = 4$ the recurrence equation gives

$$30a_6 - 5a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{a_0}{48}$$

For $n = 5$ the recurrence equation gives

$$42a_7 - 6a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{a_1}{105}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x + \frac{1}{2} a_0 x^2 + \frac{1}{3} a_1 x^3 + \frac{1}{8} a_0 x^4 + \frac{1}{15} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) a_0 + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) c_1 + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) c_2 + O(x^6)$$

$$y = 2 + x^2 + \frac{x^4}{4} + x + \frac{x^3}{3} + \frac{x^5}{15} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = x^2 + x + 2 + \frac{x^3}{3} + \frac{x^4}{4} + \frac{x^5}{15} + \frac{x^6}{24} + O(x^6) \quad (1)$$

$$y = 2 + x^2 + \frac{x^4}{4} + x + \frac{x^3}{3} + \frac{x^5}{15} + O(x^6) \quad (2)$$

Verification of solutions

$$y = x^2 + x + 2 + \frac{x^3}{3} + \frac{x^4}{4} + \frac{x^5}{15} + \frac{x^6}{24} + O(x^6)$$

Verified OK.

$$y = 2 + x^2 + \frac{x^4}{4} + x + \frac{x^3}{3} + \frac{x^5}{15} + O(x^6)$$

Verified OK.

13.12.2 Maple step by step solution

Let's solve

$$\left[y'' = y'x + y, y(0) = 2, y'|_{\{x=0\}} = 1 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - y'x - y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite DE with series expansions

- Convert $x \cdot y'$ to series expansion

$$x \cdot y' = \sum_{k=0}^{\infty} a_k k x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k- > k+2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2} (k+2)(k+1) x^k$$

Rewrite DE with series expansions

$$\sum_{k=0}^{\infty} (a_{k+2}(k+2)(k+1) - a_k(k+1)) x^k = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(k + 1)(a_{k+2}(k + 2) - a_k) = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = \frac{a_k}{k+2} \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
<- linear_1 successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

```

Order:=6;
dsolve([diff(y(x),x$2)-x*diff(y(x),x)-y(x)=0,y(0) = 2, D(y)(0) = 1],y(x),type='series',x=0);

```

$$y(x) = 2 + x + x^2 + \frac{1}{3}x^3 + \frac{1}{4}x^4 + \frac{1}{15}x^5 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 30

```

AsymptoticDSolveValue[{y'[x]-x*y'[x]-y[x]==0,{y[0]==2,y'[0]==1}},y[x],{x,0,5}]

```

$$y(x) \rightarrow \frac{x^5}{15} + \frac{x^4}{4} + \frac{x^3}{3} + x^2 + x + 2$$

13.13 problem 16

13.13.1 Existence and uniqueness analysis 3221

Internal problem ID [723]

Internal file name [OUTPUT/723_Sunday_June_05_2022_01_47_50_AM_35648494/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 16.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(x^2 + 2)y'' - y'x + 4y = 0$$

With initial conditions

$$[y(0) = -1, y'(0) = 3]$$

With the expansion point for the power series method at $x = 0$.

13.13.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = -\frac{x}{x^2 + 2}$$

$$q(x) = \frac{4}{x^2 + 2}$$

$$F = 0$$

Hence the ode is

$$y'' - \frac{xy'}{x^2 + 2} + \frac{4y}{x^2 + 2} = 0$$

The domain of $p(x) = -\frac{x}{x^2+2}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = \frac{4}{x^2+2}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (780)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (781)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= \frac{y'x - 4y}{x^2 + 2} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\
 &= \frac{-4y'x^2 + 4yx - 6y'}{(x^2 + 2)^2} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\
 &= \frac{8x^3y' + 4x^2y + 10y'x + 32y}{(x^2 + 2)^3} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\
 &= \frac{(-12x^4 + 48x^2 + 84)y' + (-48x^3 - 216x)y}{(x^2 + 2)^4} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\
 &= \frac{(-12x^5 - 648x^3 - 828x)y' + (288x^4 + 1032x^2 - 768)y}{(x^2 + 2)^5}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = -1$ and $y'(0) = 3$ gives

$$\begin{aligned}
 F_0 &= 2 \\
 F_1 &= -\frac{9}{2} \\
 F_2 &= -4 \\
 F_3 &= \frac{63}{4} \\
 F_4 &= 24
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = x^2 + 3x - 1 - \frac{3x^3}{4} - \frac{x^4}{6} + \frac{21x^5}{160} + \frac{x^6}{30} + O(x^6)$$

$$y = x^2 + 3x - 1 - \frac{3x^3}{4} - \frac{x^4}{6} + \frac{21x^5}{160} + \frac{x^6}{30} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series The ode is normalized to be

$$(x^2 + 2)y'' - y'x + 4y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(x^2 + 2) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) - \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + 4 \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} x^n a_n n(n-1) \right) + \left(\sum_{n=2}^{\infty} 2n(n-1) a_n x^{n-2} \right) + \sum_{n=1}^{\infty} (-n a_n x^n) + \left(\sum_{n=0}^{\infty} 4a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} 2n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} 2(n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} x^n a_n n(n-1) \right) + \left(\sum_{n=0}^{\infty} 2(n+2) a_{n+2} (n+1) x^n \right) \\ & + \sum_{n=1}^{\infty} (-n a_n x^n) + \left(\sum_{n=0}^{\infty} 4a_n x^n \right) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$4a_2 + 4a_0 = 0$$

$$a_2 = -a_0$$

$n = 1$ gives

$$12a_3 + 3a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{a_1}{4}$$

For $2 \leq n$, the recurrence equation is

$$n a_n (n-1) + 2(n+2) a_{n+2} (n+1) - n a_n + 4a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n (n^2 - 2n + 4)}{2(n+2)(n+1)} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$4a_2 + 24a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{6}$$

For $n = 3$ the recurrence equation gives

$$7a_3 + 40a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{7a_1}{160}$$

For $n = 4$ the recurrence equation gives

$$12a_4 + 60a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{a_0}{30}$$

For $n = 5$ the recurrence equation gives

$$19a_5 + 84a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{19a_1}{1920}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - a_0 x^2 - \frac{1}{4} a_1 x^3 + \frac{1}{6} a_0 x^4 + \frac{7}{160} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - x^2 + \frac{1}{6}x^4\right) a_0 + \left(x - \frac{1}{4}x^3 + \frac{7}{160}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - x^2 + \frac{1}{6}x^4\right) c_1 + \left(x - \frac{1}{4}x^3 + \frac{7}{160}x^5\right) c_2 + O(x^6)$$

$$y = -1 + x^2 - \frac{x^4}{6} + 3x - \frac{3x^3}{4} + \frac{21x^5}{160} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = x^2 + 3x - 1 - \frac{3x^3}{4} - \frac{x^4}{6} + \frac{21x^5}{160} + \frac{x^6}{30} + O(x^6) \quad (1)$$

$$y = -1 + x^2 - \frac{x^4}{6} + 3x - \frac{3x^3}{4} + \frac{21x^5}{160} + O(x^6) \quad (2)$$

Verification of solutions

$$y = x^2 + 3x - 1 - \frac{3x^3}{4} - \frac{x^4}{6} + \frac{21x^5}{160} + \frac{x^6}{30} + O(x^6)$$

Verified OK.

$$y = -1 + x^2 - \frac{x^4}{6} + 3x - \frac{3x^3}{4} + \frac{21x^5}{160} + O(x^6)$$

Verified OK.

Maple trace **Kovacic algorithm successful**

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Group is reducible or imprimitive
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

Order:=6;

```
dsolve([(2+x^2)*diff(y(x),x)-x*diff(y(x),x)+4*y(x)=0,y(0) = -1, D(y)(0) = 3],y(x),type='se
```

$$y(x) = -1 + 3x + x^2 - \frac{3}{4}x^3 - \frac{1}{6}x^4 + \frac{21}{160}x^5 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 32

```
AsymptoticDSolveValue[{(2+x^2)*y'[x]-x*y'[x]+4*y[x]==0,{y[0]==-1,y'[0]==3}},y[x],{x,0,5}]
```

$$y(x) \rightarrow \frac{21x^5}{160} - \frac{x^4}{6} - \frac{3x^3}{4} + x^2 + 3x - 1$$

13.14 problem 17

13.14.1 Existence and uniqueness analysis 3231

13.14.2 Maple step by step solution 3239

Internal problem ID [724]

Internal file name [OUTPUT/724_Sunday_June_05_2022_01_47_52_AM_6286473/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 17.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$y'' + y'x + 2y = 0$$

With initial conditions

$$[y(0) = 4, y'(0) = -1]$$

With the expansion point for the power series method at $x = 0$.

13.14.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = x$$

$$q(x) = 2$$

$$F = 0$$

Hence the ode is

$$y'' + y'x + 2y = 0$$

The domain of $p(x) = x$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = 2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \dots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (783)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (784)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -y'x - 2y \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= y'x^2 + 2yx - 3y' \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= -x^3y' - 2x^2y + 7y'x + 8y \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= (x^4 - 12x^2 + 15)y' + 2(x^3 - 9x)y \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= (-x^5 + 18x^3 - 57x)y' - 2y(x^4 - 15x^2 + 24)
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = 4$ and $y'(0) = -1$ gives

$$\begin{aligned}
 F_0 &= -8 \\
 F_1 &= 3 \\
 F_2 &= 32 \\
 F_3 &= -15 \\
 F_4 &= -192
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = -4x^2 - x + 4 + \frac{x^3}{2} + \frac{4x^4}{3} - \frac{x^5}{8} - \frac{4x^6}{15} + O(x^6)$$

$$y = -4x^2 - x + 4 + \frac{x^3}{2} + \frac{4x^4}{3} - \frac{x^5}{8} - \frac{4x^6}{15} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = - \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x - 2 \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n x^n a_n \right) + \left(\sum_{n=0}^{\infty} 2 a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \left(\sum_{n=1}^{\infty} n x^n a_n \right) + \left(\sum_{n=0}^{\infty} 2 a_n x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 + 2a_0 = 0$$

$$a_2 = -a_0$$

For $1 \leq n$, the recurrence equation is

$$(n + 2) a_{n+2}(n + 1) + n a_n + 2a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n}{n + 1} \quad (5)$$

For $n = 1$ the recurrence equation gives

$$6a_3 + 3a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = -\frac{a_1}{2}$$

For $n = 2$ the recurrence equation gives

$$12a_4 + 4a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{3}$$

For $n = 3$ the recurrence equation gives

$$20a_5 + 5a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_1}{8}$$

For $n = 4$ the recurrence equation gives

$$30a_6 + 6a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{a_0}{15}$$

For $n = 5$ the recurrence equation gives

$$42a_7 + 7a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{a_1}{48}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - a_0 x^2 - \frac{1}{2} a_1 x^3 + \frac{1}{3} a_0 x^4 + \frac{1}{8} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - x^2 + \frac{1}{3} x^4\right) a_0 + \left(x - \frac{1}{2} x^3 + \frac{1}{8} x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - x^2 + \frac{1}{3} x^4\right) c_1 + \left(x - \frac{1}{2} x^3 + \frac{1}{8} x^5\right) c_2 + O(x^6)$$

$$y = 4 - 4x^2 + \frac{4x^4}{3} - x + \frac{x^3}{2} - \frac{x^5}{8} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = -4x^2 - x + 4 + \frac{x^3}{2} + \frac{4x^4}{3} - \frac{x^5}{8} - \frac{4x^6}{15} + O(x^6) \quad (1)$$

$$y = 4 - 4x^2 + \frac{4x^4}{3} - x + \frac{x^3}{2} - \frac{x^5}{8} + O(x^6) \quad (2)$$

Verification of solutions

$$y = -4x^2 - x + 4 + \frac{x^3}{2} + \frac{4x^4}{3} - \frac{x^5}{8} - \frac{4x^6}{15} + O(x^6)$$

Verified OK.

$$y = 4 - 4x^2 + \frac{4x^4}{3} - x + \frac{x^3}{2} - \frac{x^5}{8} + O(x^6)$$

Verified OK.

13.14.2 Maple step by step solution

Let's solve

$$\left[y'' = -y'x - 2y, y(0) = 4, y'|_{\{x=0\}} = -1 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + y'x + 2y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite DE with series expansions

- Convert $x \cdot y'$ to series expansion

$$x \cdot y' = \sum_{k=0}^{\infty} a_k k x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k- > k+2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2}(k+2)(k+1) x^k$$

Rewrite DE with series expansions

$$\sum_{k=0}^{\infty} (a_{k+2}(k+2)(k+1) + a_k(k+2)) x^k = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(k+2)(ka_{k+2} + a_k + a_{k+2}) = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = -\frac{a_k}{k+1} \right]$$

Maple trace **Kovacic algorithm successful**

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Reducible group (found an exponential solution)
    Group is reducible, not completely reducible
<- Kovacics algorithm successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 20

```
Order:=6;
```

```
dsolve([diff(y(x),x$2)+x*diff(y(x),x)+2*y(x)=0,y(0) = 4, D(y)(0) = -1],y(x),type='series',x=
```

$$y(x) = 4 - x - 4x^2 + \frac{1}{2}x^3 + \frac{4}{3}x^4 - \frac{1}{8}x^5 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 34

```
AsymptoticDSolveValue[{y'[x]+x*y'[x]+2*y[x]==0,{y[0]==4,y'[0]==-1}},y[x],{x,0,5}]
```

$$y(x) \rightarrow -\frac{x^5}{8} + \frac{4x^4}{3} + \frac{x^3}{2} - 4x^2 - x + 4$$

13.15 problem 18

13.15.1 Existence and uniqueness analysis	3241
13.15.2 Maple step by step solution	3249

Internal problem ID [725]

Internal file name [OUTPUT/725_Sunday_June_05_2022_01_47_54_AM_53815552/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 18.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second order series method. Ordinary point", "second_order_change_of_variable_on_y_method_2", "second order series method. Taylor series method", "second_order_ode_non_constant_coeff_transformation_on_B"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(1 - x)y'' + y'x - y = 0$$

With initial conditions

$$[y(0) = -3, y'(0) = 2]$$

With the expansion point for the power series method at $x = 0$.

13.15.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = \frac{x}{1-x}$$
$$q(x) = -\frac{1}{1-x}$$
$$F = 0$$

Hence the ode is

$$y'' + \frac{xy'}{1-x} - \frac{y}{1-x} = 0$$

The domain of $p(x) = \frac{x}{1-x}$ is

$$\{x < 1 \vee 1 < x\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = -\frac{1}{1-x}$ is

$$\{x < 1 \vee 1 < x\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (786)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (787)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= \frac{-y + y'x}{x - 1} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= \frac{-y + y'x}{x - 1} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= \frac{-y + y'x}{x - 1} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= \frac{-y + y'x}{x - 1} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= \frac{-y + y'x}{x - 1}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = -3$ and $y'(0) = 2$ gives

$$F_0 = -3$$

$$F_1 = -3$$

$$F_2 = -3$$

$$F_3 = -3$$

$$F_4 = -3$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = -3 + 2x - \frac{3x^2}{2} - \frac{x^3}{2} - \frac{x^4}{8} - \frac{x^5}{40} - \frac{x^6}{240} + O(x^6)$$

$$y = -3 + 2x - \frac{3x^2}{2} - \frac{x^3}{2} - \frac{x^4}{8} - \frac{x^5}{40} - \frac{x^6}{240} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. The ode is normalized to be

$$(1 - x)y'' + y'x - y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(1 - x) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x - \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\sum_{n=2}^{\infty} (-n x^{n-1} a_n (n-1)) + \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} (-n x^{n-1} a_n (n-1)) = \sum_{n=1}^{\infty} (-(n+1) a_{n+1} n x^n)$$

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} \sum_{n=1}^{\infty} (-(n+1) a_{n+1} n x^n) + \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) \\ + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$2a_2 - a_0 = 0$$

$$a_2 = \frac{a_0}{2}$$

For $1 \leq n$, the recurrence equation is

$$-(n+1) a_{n+1} n + (n+2) a_{n+2} (n+1) + n a_n - a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} a_{n+2} &= \frac{n^2 a_{n+1} - n a_n + n a_{n+1} + a_n}{(n+2)(n+1)} \\ (5) \quad &= \frac{(-n+1) a_n}{(n+2)(n+1)} + \frac{(n^2+n) a_{n+1}}{(n+2)(n+1)} \end{aligned}$$

For $n = 1$ the recurrence equation gives

$$-2a_2 + 6a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = \frac{a_0}{6}$$

For $n = 2$ the recurrence equation gives

$$-6a_3 + 12a_4 + a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{24}$$

For $n = 3$ the recurrence equation gives

$$-12a_4 + 20a_5 + 2a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_0}{120}$$

For $n = 4$ the recurrence equation gives

$$-20a_5 + 30a_6 + 3a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{a_0}{720}$$

For $n = 5$ the recurrence equation gives

$$-30a_6 + 42a_7 + 4a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{a_0}{5040}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x + \frac{1}{2} a_0 x^2 + \frac{1}{6} a_0 x^3 + \frac{1}{24} a_0 x^4 + \frac{1}{120} a_0 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) a_0 + a_1x + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) c_1 + c_2x + O(x^6)$$

$$y = -3 - \frac{3x^2}{2} - \frac{x^3}{2} - \frac{x^4}{8} - \frac{x^5}{40} + 2x + O(x^6)$$

Summary

The solution(s) found are the following

$$y = -3 + 2x - \frac{3x^2}{2} - \frac{x^3}{2} - \frac{x^4}{8} - \frac{x^5}{40} - \frac{x^6}{240} + O(x^6) \quad (1)$$

$$y = -3 - \frac{3x^2}{2} - \frac{x^3}{2} - \frac{x^4}{8} - \frac{x^5}{40} + 2x + O(x^6) \quad (2)$$

Verification of solutions

$$y = -3 + 2x - \frac{3x^2}{2} - \frac{x^3}{2} - \frac{x^4}{8} - \frac{x^5}{40} - \frac{x^6}{240} + O(x^6)$$

Verified OK.

$$y = -3 - \frac{3x^2}{2} - \frac{x^3}{2} - \frac{x^4}{8} - \frac{x^5}{40} + 2x + O(x^6)$$

Verified OK.

13.15.2 Maple step by step solution

Let's solve

$$\left[(1-x)y'' + y'x - y = 0, y(0) = -3, y'|_{\{x=0\}} = 2 \right]$$

- Highest derivative means the order of the ODE is 2
 y''
- Isolate 2nd derivative

$$y'' = -\frac{y}{x-1} + \frac{xy'}{x-1}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{xy'}{x-1} + \frac{y}{x-1} = 0$$

- Check to see if $x_0 = 1$ is a regular singular point

- Define functions

$$[P_2(x) = -\frac{x}{x-1}, P_3(x) = \frac{1}{x-1}]$$

- $(x-1) \cdot P_2(x)$ is analytic at $x = 1$

$$((x-1) \cdot P_2(x)) \Big|_{x=1} = -1$$

- $(x-1)^2 \cdot P_3(x)$ is analytic at $x = 1$

$$((x-1)^2 \cdot P_3(x)) \Big|_{x=1} = 0$$

- $x = 1$ is a regular singular point

Check to see if $x_0 = 1$ is a regular singular point

$$x_0 = 1$$

- Multiply by denominators

$$(x-1)y'' - y'x + y = 0$$

- Change variables using $x = u + 1$ so that the regular singular point is at $u = 0$

$$u \left(\frac{d^2}{du^2} y(u) \right) + (-u-1) \left(\frac{d}{du} y(u) \right) + y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k+1-m$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k+1-m+r) u^{k+r}$$

- Convert $u \cdot \left(\frac{d^2}{du^2} y(u) \right)$ to series expansion

$$u \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) u^{k+r-1}$$

- Shift index using $k \rightarrow k+1$

$$u \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=-1}^{\infty} a_{k+1} (k+1+r)(k+r) u^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r(-2+r) u^{-1+r} + \left(\sum_{k=0}^{\infty} (a_{k+1}(k+1+r)(k+r-1) - a_k(k+r-1)) u^{k+r} \right) = 0$$

- a_0 cannot be 0 by assumption, giving the indicial equation
 $r(-2+r) = 0$
- Values of r that satisfy the indicial equation
 $r \in \{0, 2\}$
- Each term in the series must be 0, giving the recursion relation
 $(k+r-1)(a_{k+1}(k+1+r) - a_k) = 0$
- Recursion relation that defines series solution to ODE

$$a_{k+1} = \frac{a_k}{k+1+r}$$

- Recursion relation for $r = 0$

$$a_{k+1} = \frac{a_k}{k+1}$$

- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{k+1} = \frac{a_k}{k+1} \right]$$

- Revert the change of variables $u = x - 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x-1)^k, a_{k+1} = \frac{a_k}{k+1} \right]$$

- Recursion relation for $r = 2$

$$a_{k+1} = \frac{a_k}{k+3}$$

- Solution for $r = 2$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+2}, a_{k+1} = \frac{a_k}{k+3} \right]$$

- Revert the change of variables $u = x - 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x-1)^{k+2}, a_{k+1} = \frac{a_k}{k+3} \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k (x-1)^k \right) + \left(\sum_{k=0}^{\infty} b_k (x-1)^{k+2} \right), a_{k+1} = \frac{a_k}{k+1}, b_{k+1} = \frac{b_k}{k+3} \right]$$

Maple trace **Kovacic algorithm successful**

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Reducible group (found an exponential solution)
    Reducible group (found another exponential solution)
<- Kovacics algorithm successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

```

Order:=6;
dsolve([(1-x)*diff(y(x),x$2)+x*diff(y(x),x)-y(x)=0,y(0) = -3, D(y)(0) = 2],y(x),type='series

```

$$y(x) = -3 + 2x - \frac{3}{2}x^2 - \frac{1}{2}x^3 - \frac{1}{8}x^4 - \frac{1}{40}x^5 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 36

```

AsymptoticDSolveValue[{{(1-x)*y''[x]+x*y'[x]-y[x]==0,{y[0]==-3,y'[0]==2}},y[x]},{x,0,5}]

```

$$y(x) \rightarrow -\frac{x^5}{40} - \frac{x^4}{8} - \frac{x^3}{2} - \frac{3x^2}{2} + 2x - 3$$

13.16 problem 21

13.16.1 Maple step by step solution 3260

Internal problem ID [726]

Internal file name [OUTPUT/726_Sunday_June_05_2022_01_47_56_AM_30999466/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 21.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$y'' - 2y'x + \lambda y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0) y'(x_0) + \frac{(x - x_0)^2}{2} y''(x_0) + \frac{(x - x_0)^3}{3!} y'''(x_0) + \cdots \\ &= y_0 + x y'_0 + \frac{x^2}{2} f|_{x_0, y_0, y'_0} + \frac{x^3}{3!} f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + x y'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (789)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (790)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = 2y'x - \lambda y$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dx} \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\ &= (4x^2 - \lambda + 2) y' - 2y\lambda x \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dx} \\ &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\ &= (8x^3 - 4\lambda x + 12x) y' - 4\lambda \left(x^2 - \frac{\lambda}{4} + 1 \right) y \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dx} \\ &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\ &= (\lambda^2 + (-12x^2 - 8)\lambda + 16x^4 + 48x^2 + 12) y' - 8 \left(x^2 - \frac{\lambda}{2} + \frac{5}{2} \right) x\lambda y \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dx} \\ &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\ &= (6x\lambda^2 + (-32x^3 - 60x)\lambda + 32x^5 + 160x^3 + 120x) y' - 16 \left(\frac{\lambda^2}{16} + \left(-\frac{3x^2}{4} - \frac{3}{4} \right) \lambda + x^4 + \frac{9x^2}{2} + 2 \right) x\lambda y \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$F_0 = -y(0)\lambda$$

$$F_1 = -y'(0)\lambda + 2y'(0)$$

$$F_2 = y(0)\lambda^2 - 4y(0)\lambda$$

$$F_3 = y'(0)\lambda^2 - 8y'(0)\lambda + 12y'(0)$$

$$F_4 = -y(0)\lambda^3 + 12y(0)\lambda^2 - 32y(0)\lambda$$

Substituting all the above in (7) and simplifying gives the solution as

$$\begin{aligned} y &= \left(1 - \frac{1}{2}\lambda x^2 + \frac{1}{24}x^4\lambda^2 - \frac{1}{6}x^4\lambda - \frac{1}{720}x^6\lambda^3 + \frac{1}{60}x^6\lambda^2 - \frac{2}{45}x^6\lambda \right) y(0) \\ &\quad + \left(x - \frac{1}{6}x^3\lambda + \frac{1}{3}x^3 + \frac{1}{120}x^5\lambda^2 - \frac{1}{15}x^5\lambda + \frac{1}{10}x^5 \right) y'(0) + O(x^6) \end{aligned}$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = 2 \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x - \lambda \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \sum_{n=1}^{\infty} (-2n x^n a_n) + \left(\sum_{n=0}^{\infty} \lambda a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \sum_{n=1}^{\infty} (-2n x^n a_n) + \left(\sum_{n=0}^{\infty} \lambda a_n x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$\lambda a_0 + 2a_2 = 0$$

$$a_2 = -\frac{\lambda a_0}{2}$$

For $1 \leq n$, the recurrence equation is

$$(n+2)a_{n+2}(n+1) - 2na_n + \lambda a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n(\lambda - 2n)}{(n+2)(n+1)} \quad (5)$$

For $n = 1$ the recurrence equation gives

$$\lambda a_1 - 2a_1 + 6a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = -\frac{1}{6}\lambda a_1 + \frac{1}{3}a_1$$

For $n = 2$ the recurrence equation gives

$$\lambda a_2 - 4a_2 + 12a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{1}{24}\lambda^2 a_0 - \frac{1}{6}\lambda a_0$$

For $n = 3$ the recurrence equation gives

$$\lambda a_3 - 6a_3 + 20a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{1}{120}\lambda^2 a_1 - \frac{1}{15}\lambda a_1 + \frac{1}{10}a_1$$

For $n = 4$ the recurrence equation gives

$$\lambda a_4 - 8a_4 + 30a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{1}{720}\lambda^3 a_0 + \frac{1}{60}\lambda^2 a_0 - \frac{2}{45}\lambda a_0$$

For $n = 5$ the recurrence equation gives

$$\lambda a_5 - 10a_5 + 42a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{1}{5040}\lambda^3 a_1 + \frac{1}{280}\lambda^2 a_1 - \frac{23}{1260}\lambda a_1 + \frac{1}{42}a_1$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$\begin{aligned} y &= a_0 + a_1 x - \frac{\lambda a_0 x^2}{2} + \left(-\frac{1}{6}\lambda a_1 + \frac{1}{3}a_1\right) x^3 \\ &\quad + \left(\frac{1}{24}\lambda^2 a_0 - \frac{1}{6}\lambda a_0\right) x^4 + \left(\frac{1}{120}\lambda^2 a_1 - \frac{1}{15}\lambda a_1 + \frac{1}{10}a_1\right) x^5 + \dots \end{aligned}$$

Collecting terms, the solution becomes

$$\begin{aligned} y &= \left(1 - \frac{\lambda x^2}{2} + \left(\frac{1}{24}\lambda^2 - \frac{1}{6}\lambda\right) x^4\right) a_0 \\ &\quad + \left(x + \left(-\frac{\lambda}{6} + \frac{1}{3}\right) x^3 + \left(\frac{1}{120}\lambda^2 - \frac{1}{15}\lambda + \frac{1}{10}\right) x^5\right) a_1 + O(x^6) \end{aligned} \tag{3}$$

At $x = 0$ the solution above becomes

$$\begin{aligned} y &= \left(1 - \frac{\lambda x^2}{2} + \left(\frac{1}{24}\lambda^2 - \frac{1}{6}\lambda\right) x^4\right) c_1 \\ &\quad + \left(x + \left(-\frac{\lambda}{6} + \frac{1}{3}\right) x^3 + \left(\frac{1}{120}\lambda^2 - \frac{1}{15}\lambda + \frac{1}{10}\right) x^5\right) c_2 + O(x^6) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{1}{2}\lambda x^2 + \frac{1}{24}x^4\lambda^2 - \frac{1}{6}x^4\lambda - \frac{1}{720}x^6\lambda^3 + \frac{1}{60}x^6\lambda^2 - \frac{2}{45}x^6\lambda\right) y(0) \\ + \left(x - \frac{1}{6}x^3\lambda + \frac{1}{3}x^3 + \frac{1}{120}x^5\lambda^2 - \frac{1}{15}x^5\lambda + \frac{1}{10}x^5\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 - \frac{\lambda x^2}{2} + \left(\frac{1}{24}\lambda^2 - \frac{1}{6}\lambda\right) x^4\right) c_1 \\ + \left(x + \left(-\frac{\lambda}{6} + \frac{1}{3}\right) x^3 + \left(\frac{1}{120}\lambda^2 - \frac{1}{15}\lambda + \frac{1}{10}\right) x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 - \frac{1}{2}\lambda x^2 + \frac{1}{24}x^4\lambda^2 - \frac{1}{6}x^4\lambda - \frac{1}{720}x^6\lambda^3 + \frac{1}{60}x^6\lambda^2 - \frac{2}{45}x^6\lambda\right) y(0) \\ + \left(x - \frac{1}{6}x^3\lambda + \frac{1}{3}x^3 + \frac{1}{120}x^5\lambda^2 - \frac{1}{15}x^5\lambda + \frac{1}{10}x^5\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 - \frac{\lambda x^2}{2} + \left(\frac{1}{24}\lambda^2 - \frac{1}{6}\lambda\right) x^4\right) c_1 \\ + \left(x + \left(-\frac{\lambda}{6} + \frac{1}{3}\right) x^3 + \left(\frac{1}{120}\lambda^2 - \frac{1}{15}\lambda + \frac{1}{10}\right) x^5\right) c_2 + O(x^6)$$

Verified OK.

13.16.1 Maple step by step solution

Let's solve

$$y'' = 2y'x - \lambda y$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - 2y'x + \lambda y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite DE with series expansions

- Convert $x \cdot y'$ to series expansion

$$x \cdot y' = \sum_{k=0}^{\infty} a_k k x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k \rightarrow k+2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2}(k+2)(k+1) x^k$$

Rewrite DE with series expansions

$$\sum_{k=0}^{\infty} (a_{k+2}(k+2)(k+1) - a_k(2k-\lambda)) x^k = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(k^2 + 3k + 2) a_{k+2} - 2\left(k - \frac{\lambda}{2}\right) a_k = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = \frac{(2k-\lambda)a_k}{k^2+3k+2} \right]$$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
  -> Bessel
  -> elliptic
  -> Legendre
  -> Kummer
      -> hyper3: Equivalence to 1F1 under a power @ Moebius
      <- hyper3 successful: received ODE is equivalent to the 1F1 ODE
      <- Kummer successful
<- special function solution successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 63

```
Order:=6;
dsolve(diff(y(x),x$2)-2*x*diff(y(x),x)+lambda*y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 - \frac{\lambda x^2}{2} + \frac{\lambda(\lambda - 4)x^4}{24}\right) y(0) + \left(x - \frac{(\lambda - 2)x^3}{6} + \frac{(\lambda - 2)(-6 + \lambda)x^5}{120}\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 80

```
AsymptoticDSolveValue[y''[x]-2*x*y'[x]+\[Lambda]*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(\frac{\lambda^2 x^5}{120} - \frac{\lambda x^5}{15} + \frac{x^5}{10} - \frac{\lambda x^3}{6} + \frac{x^3}{3} + x \right) + c_1 \left(\frac{\lambda^2 x^4}{24} - \frac{\lambda x^4}{6} - \frac{\lambda x^2}{2} + 1 \right)$$

13.17 problem 23

13.17.1 Existence and uniqueness analysis	3263
13.17.2 Maple step by step solution	3271

Internal problem ID [727]

Internal file name [OUTPUT/727_Sunday_June_05_2022_01_47_57_AM_94347532/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 23.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact linear second order ode", "second_order_integrable_as_is", "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _exact , _linear , _homogeneous]]
```

$$y'' - y'x - y = 0$$

With initial conditions

$$[y(0) = 1, y'(0) = 0]$$

With the expansion point for the power series method at $x = 0$.

13.17.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = -x$$

$$q(x) = -1$$

$$F = 0$$

Hence the ode is

$$y'' - y'x - y = 0$$

The domain of $p(x) = -x$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = -1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \dots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (792)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (793)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= y'x + y \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= y'x^2 + yx + 2y' \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= (x^3 + 5x)y' + y(x^2 + 3) \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= (x^4 + 9x^2 + 8)y' + xy(x^2 + 7) \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= (x^5 + 14x^3 + 33x)y' + y(x^4 + 12x^2 + 15)
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = 1$ and $y'(0) = 0$ gives

$$\begin{aligned}
 F_0 &= 1 \\
 F_1 &= 0 \\
 F_2 &= 3 \\
 F_3 &= 0 \\
 F_4 &= 15
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = 1 + \frac{x^2}{2} + \frac{x^4}{8} + \frac{x^6}{48} + O(x^6)$$

$$y = 1 + \frac{x^2}{2} + \frac{x^4}{8} + \frac{x^6}{48} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \sum_{n=1}^{\infty} (-n x^n a_n) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \sum_{n=1}^{\infty} (-n x^n a_n) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 - a_0 = 0$$

$$a_2 = \frac{a_0}{2}$$

For $1 \leq n$, the recurrence equation is

$$(n + 2) a_{n+2}(n + 1) - n a_n - a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = \frac{a_n}{n + 2} \quad (5)$$

For $n = 1$ the recurrence equation gives

$$6a_3 - 2a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = \frac{a_1}{3}$$

For $n = 2$ the recurrence equation gives

$$12a_4 - 3a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{8}$$

For $n = 3$ the recurrence equation gives

$$20a_5 - 4a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_1}{15}$$

For $n = 4$ the recurrence equation gives

$$30a_6 - 5a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{a_0}{48}$$

For $n = 5$ the recurrence equation gives

$$42a_7 - 6a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{a_1}{105}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x + \frac{1}{2} a_0 x^2 + \frac{1}{3} a_1 x^3 + \frac{1}{8} a_0 x^4 + \frac{1}{15} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) a_0 + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) c_1 + \left(x + \frac{1}{3}x^3 + \frac{1}{15}x^5\right) c_2 + O(x^6)$$

$$y = 1 + \frac{x^2}{2} + \frac{x^4}{8} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = 1 + \frac{x^2}{2} + \frac{x^4}{8} + \frac{x^6}{48} + O(x^6) \quad (1)$$

$$y = 1 + \frac{x^2}{2} + \frac{x^4}{8} + O(x^6) \quad (2)$$

Verification of solutions

$$y = 1 + \frac{x^2}{2} + \frac{x^4}{8} + \frac{x^6}{48} + O(x^6)$$

Verified OK.

$$y = 1 + \frac{x^2}{2} + \frac{x^4}{8} + O(x^6)$$

Verified OK.

13.17.2 Maple step by step solution

Let's solve

$$\left[y'' = y'x + y, y(0) = 1, y'|_{\{x=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - y'x - y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite DE with series expansions

- Convert $x \cdot y'$ to series expansion

$$x \cdot y' = \sum_{k=0}^{\infty} a_k k x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k- > k+2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2} (k+2)(k+1) x^k$$

Rewrite DE with series expansions

$$\sum_{k=0}^{\infty} (a_{k+2}(k+2)(k+1) - a_k(k+1)) x^k = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(k + 1)(a_{k+2}(k + 2) - a_k) = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = \frac{a_k}{k+2} \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
<- linear_1 successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 14

```

Order:=6;
dsolve([diff(y(x),x$2)-x*diff(y(x),x)-y(x)=0,y(0) = 1, D(y)(0) = 0],y(x),type='series',x=0);

```

$$y(x) = 1 + \frac{1}{2}x^2 + \frac{1}{8}x^4 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 19

```

AsymptoticDSolveValue[{y'[x]-x*y'[x]-y[x]==0,{y[0]==1,y'[0]==0}},y[x],{x,0,5}]

```

$$y(x) \rightarrow \frac{x^4}{8} + \frac{x^2}{2} + 1$$

13.18 problem 24

13.18.1 Existence and uniqueness analysis 3273

Internal problem ID [728]

Internal file name [OUTPUT/728_Sunday_June_05_2022_01_47_59_AM_5624398/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 24.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(x^2 + 2)y'' - y'x + 4y = 0$$

With initial conditions

$$[y(0) = 1, y'(0) = 0]$$

With the expansion point for the power series method at $x = 0$.

13.18.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = -\frac{x}{x^2 + 2}$$

$$q(x) = \frac{4}{x^2 + 2}$$

$$F = 0$$

Hence the ode is

$$y'' - \frac{xy'}{x^2 + 2} + \frac{4y}{x^2 + 2} = 0$$

The domain of $p(x) = -\frac{x}{x^2+2}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = \frac{4}{x^2+2}$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (795)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (796)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= \frac{y'x - 4y}{x^2 + 2} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\
 &= \frac{-4y'x^2 + 4yx - 6y'}{(x^2 + 2)^2} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\
 &= \frac{8x^3y' + 4x^2y + 10y'x + 32y}{(x^2 + 2)^3} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\
 &= \frac{(-12x^4 + 48x^2 + 84)y' + (-48x^3 - 216x)y}{(x^2 + 2)^4} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\
 &= \frac{(-12x^5 - 648x^3 - 828x)y' + (288x^4 + 1032x^2 - 768)y}{(x^2 + 2)^5}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = 1$ and $y'(0) = 0$ gives

$$\begin{aligned}
 F_0 &= -2 \\
 F_1 &= 0 \\
 F_2 &= 4 \\
 F_3 &= 0 \\
 F_4 &= -24
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = 1 - x^2 + \frac{x^4}{6} - \frac{x^6}{30} + O(x^6)$$

$$y = 1 - x^2 + \frac{x^4}{6} - \frac{x^6}{30} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series The ode is normalized to be

$$(x^2 + 2) y'' - y'x + 4y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(x^2 + 2) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) - \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + 4 \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} x^n a_n n(n-1) \right) + \left(\sum_{n=2}^{\infty} 2n(n-1) a_n x^{n-2} \right) + \sum_{n=1}^{\infty} (-n a_n x^n) + \left(\sum_{n=0}^{\infty} 4a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} 2n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} 2(n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=2}^{\infty} x^n a_n n(n-1) \right) + \left(\sum_{n=0}^{\infty} 2(n+2) a_{n+2} (n+1) x^n \right) + \sum_{n=1}^{\infty} (-n a_n x^n) + \left(\sum_{n=0}^{\infty} 4a_n x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$4a_2 + 4a_0 = 0$$

$$a_2 = -a_0$$

$n = 1$ gives

$$12a_3 + 3a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{a_1}{4}$$

For $2 \leq n$, the recurrence equation is

$$na_n(n-1) + 2(n+2)a_{n+2}(n+1) - na_n + 4a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n(n^2 - 2n + 4)}{2(n+2)(n+1)} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$4a_2 + 24a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{6}$$

For $n = 3$ the recurrence equation gives

$$7a_3 + 40a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{7a_1}{160}$$

For $n = 4$ the recurrence equation gives

$$12a_4 + 60a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{a_0}{30}$$

For $n = 5$ the recurrence equation gives

$$19a_5 + 84a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{19a_1}{1920}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - a_0 x^2 - \frac{1}{4} a_1 x^3 + \frac{1}{6} a_0 x^4 + \frac{7}{160} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - x^2 + \frac{1}{6} x^4\right) a_0 + \left(x - \frac{1}{4} x^3 + \frac{7}{160} x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - x^2 + \frac{1}{6} x^4\right) c_1 + \left(x - \frac{1}{4} x^3 + \frac{7}{160} x^5\right) c_2 + O(x^6)$$

$$y = 1 - x^2 + \frac{x^4}{6} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = 1 - x^2 + \frac{x^4}{6} - \frac{x^6}{30} + O(x^6) \quad (1)$$

$$y = 1 - x^2 + \frac{x^4}{6} + O(x^6) \quad (2)$$

Verification of solutions

$$y = 1 - x^2 + \frac{x^4}{6} - \frac{x^6}{30} + O(x^6)$$

Verified OK.

$$y = 1 - x^2 + \frac{x^4}{6} + O(x^6)$$

Verified OK.

Maple trace **Kovacic algorithm successful**

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacics algorithm  
  A Liouvillian solution exists  
  Group is reducible or imprimitive  
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
Order:=6;  
dsolve([(2+x^2)*diff(y(x),x$2)-x*diff(y(x),x)+4*y(x)=0,y(0) = 1, D(y)(0) = 0],y(x),type='ser
```

$$y(x) = 1 - x^2 + \frac{1}{6}x^4 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 17

```
AsymptoticDSolveValue[{(2+x^2)*y''[x]-x*y'[x]+4*y[x]==0,{y[0]==1,y'[0]==0}},y[x],{x,0,5}]
```

$$y(x) \rightarrow \frac{x^4}{6} - x^2 + 1$$

13.19 problem 25

13.19.1 Existence and uniqueness analysis 3283

13.19.2 Maple step by step solution 3291

Internal problem ID [729]

Internal file name [OUTPUT/729_Sunday_June_05_2022_01_48_01_AM_21971426/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 25.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$y'' + y'x + 2y = 0$$

With initial conditions

$$[y(0) = 0, y'(0) = 1]$$

With the expansion point for the power series method at $x = 0$.

13.19.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = x$$

$$q(x) = 2$$

$$F = 0$$

Hence the ode is

$$y'' + y'x + 2y = 0$$

The domain of $p(x) = x$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = 2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \dots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (798)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (799)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -y'x - 2y \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= y'x^2 + 2yx - 3y' \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= -x^3y' - 2x^2y + 7y'x + 8y \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= (x^4 - 12x^2 + 15)y' + 2(x^3 - 9x)y \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= (-x^5 + 18x^3 - 57x)y' - 2y(x^4 - 15x^2 + 24)
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = 0$ and $y'(0) = 1$ gives

$$\begin{aligned}
 F_0 &= 0 \\
 F_1 &= -3 \\
 F_2 &= 0 \\
 F_3 &= 15 \\
 F_4 &= 0
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = x - \frac{x^3}{2} + \frac{x^5}{8} + O(x^6)$$

$$y = x - \frac{x^3}{2} + \frac{x^5}{8} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = - \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x - 2 \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n x^n a_n \right) + \left(\sum_{n=0}^{\infty} 2 a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \left(\sum_{n=1}^{\infty} n x^n a_n \right) + \left(\sum_{n=0}^{\infty} 2 a_n x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 + 2a_0 = 0$$

$$a_2 = -a_0$$

For $1 \leq n$, the recurrence equation is

$$(n + 2) a_{n+2}(n + 1) + na_n + 2a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n}{n + 1} \quad (5)$$

For $n = 1$ the recurrence equation gives

$$6a_3 + 3a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = -\frac{a_1}{2}$$

For $n = 2$ the recurrence equation gives

$$12a_4 + 4a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{3}$$

For $n = 3$ the recurrence equation gives

$$20a_5 + 5a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_1}{8}$$

For $n = 4$ the recurrence equation gives

$$30a_6 + 6a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{a_0}{15}$$

For $n = 5$ the recurrence equation gives

$$42a_7 + 7a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{a_1}{48}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - a_0 x^2 - \frac{1}{2} a_1 x^3 + \frac{1}{3} a_0 x^4 + \frac{1}{8} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - x^2 + \frac{1}{3}x^4\right) a_0 + \left(x - \frac{1}{2}x^3 + \frac{1}{8}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - x^2 + \frac{1}{3}x^4\right) c_1 + \left(x - \frac{1}{2}x^3 + \frac{1}{8}x^5\right) c_2 + O(x^6)$$

$$y = x - \frac{x^3}{2} + \frac{x^5}{8} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = x - \frac{x^3}{2} + \frac{x^5}{8} + O(x^6) \quad (1)$$

$$y = x - \frac{x^3}{2} + \frac{x^5}{8} + O(x^6) \quad (2)$$

Verification of solutions

$$y = x - \frac{x^3}{2} + \frac{x^5}{8} + O(x^6)$$

Verified OK.

$$y = x - \frac{x^3}{2} + \frac{x^5}{8} + O(x^6)$$

Verified OK.

13.19.2 Maple step by step solution

Let's solve

$$\left[y'' = -y'x - 2y, y(0) = 0, y'|_{\{x=0\}} = 1 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + y'x + 2y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite DE with series expansions

- Convert $x \cdot y'$ to series expansion

$$x \cdot y' = \sum_{k=0}^{\infty} a_k k x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k- > k+2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2} (k+2)(k+1) x^k$$

Rewrite DE with series expansions

$$\sum_{k=0}^{\infty} (a_{k+2}(k+2)(k+1) + a_k(k+2)) x^k = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(k+2)(ka_{k+2} + a_k + a_{k+2}) = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = -\frac{a_k}{k+1} \right]$$

Maple trace **Kovacic algorithm successful**

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
    A Liouvillian solution exists
    Reducible group (found an exponential solution)
    Group is reducible, not completely reducible
<- Kovacics algorithm successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
Order:=6;
```

```
dsolve([diff(y(x),x$2)+x*diff(y(x),x)+2*y(x)=0,y(0) = 0, D(y)(0) = 1],y(x),type='series',x=0
```

$$y(x) = x - \frac{1}{2}x^3 + \frac{1}{8}x^5 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 19

```
AsymptoticDSolveValue[{y'[x]+x*y'[x]+2*y[x]==0,{y[0]==0,y'[0]==1}},y[x],{x,0,5}]
```

$$y(x) \rightarrow \frac{x^5}{8} - \frac{x^3}{2} + x$$

13.20 problem 26

13.20.1 Existence and uniqueness analysis	3293
13.20.2 Maple step by step solution	3301

Internal problem ID [730]

Internal file name [OUTPUT/730_Sunday_June_05_2022_01_48_03_AM_91788621/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 26.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(-x^2 + 4)y'' + y'x + 2y = 0$$

With initial conditions

$$[y(0) = 0, y'(0) = 1]$$

With the expansion point for the power series method at $x = 0$.

13.20.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = \frac{x}{-x^2 + 4}$$

$$q(x) = \frac{2}{-x^2 + 4}$$

$$F = 0$$

Hence the ode is

$$y'' + \frac{xy'}{-x^2 + 4} + \frac{2y}{-x^2 + 4} = 0$$

The domain of $p(x) = \frac{x}{-x^2+4}$ is

$$\{-\infty \leq x < -2, -2 < x < 2, 2 < x \leq \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = \frac{2}{-x^2+4}$ is

$$\{-\infty \leq x < -2, -2 < x < 2, 2 < x \leq \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (801)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (802)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= \frac{2y + y'x}{x^2 - 4} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= \frac{2y'x^2 - 2yx - 12y'}{(x^2 - 4)^2} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= \frac{-4x^3y' + 10x^2y + 28y'x - 16y}{(x^2 - 4)^3} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= \frac{(18x^4 - 120x^2 - 48)y' + (-48x^3 + 72x)y}{(x^2 - 4)^4} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= \frac{(-102x^5 + 576x^3 + 1008x)y' + (276x^4 - 168x^2 - 384)y}{(x^2 - 4)^5}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = 0$ and $y'(0) = 1$ gives

$$\begin{aligned}
 F_0 &= 0 \\
 F_1 &= -\frac{3}{4} \\
 F_2 &= 0 \\
 F_3 &= -\frac{3}{16} \\
 F_4 &= 0
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = x - \frac{x^3}{8} - \frac{x^5}{640} + O(x^6)$$

$$y = x - \frac{x^3}{8} - \frac{x^5}{640} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series The ode is normalized to be

$$(-x^2 + 4) y'' + y'x + 2y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(-x^2 + 4) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + 2 \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\sum_{n=2}^{\infty} (-x^n a_n n(n-1)) + \left(\sum_{n=2}^{\infty} 4n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \left(\sum_{n=0}^{\infty} 2a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} 4n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} 4(n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} \sum_{n=2}^{\infty} (-x^n a_n n(n-1)) + \left(\sum_{n=0}^{\infty} 4(n+2) a_{n+2} (n+1) x^n \right) \\ + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \left(\sum_{n=0}^{\infty} 2 a_n x^n \right) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$8a_2 + 2a_0 = 0$$

$$a_2 = -\frac{a_0}{4}$$

$n = 1$ gives

$$24a_3 + 3a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{a_1}{8}$$

For $2 \leq n$, the recurrence equation is

$$-na_n(n-1) + 4(n+2)a_{n+2}(n+1) + na_n + 2a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = \frac{a_n(n^2 - 2n - 2)}{4(n+2)(n+1)} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$2a_2 + 48a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{96}$$

For $n = 3$ the recurrence equation gives

$$-a_3 + 80a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = -\frac{a_1}{640}$$

For $n = 4$ the recurrence equation gives

$$-6a_4 + 120a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{a_0}{1920}$$

For $n = 5$ the recurrence equation gives

$$-13a_5 + 168a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{13a_1}{107520}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - \frac{1}{4} a_0 x^2 - \frac{1}{8} a_1 x^3 + \frac{1}{96} a_0 x^4 - \frac{1}{640} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - \frac{1}{4}x^2 + \frac{1}{96}x^4\right) a_0 + \left(x - \frac{1}{8}x^3 - \frac{1}{640}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - \frac{1}{4}x^2 + \frac{1}{96}x^4\right) c_1 + \left(x - \frac{1}{8}x^3 - \frac{1}{640}x^5\right) c_2 + O(x^6)$$

$$y = x - \frac{x^3}{8} - \frac{x^5}{640} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = x - \frac{x^3}{8} - \frac{x^5}{640} + O(x^6) \quad (1)$$

$$y = x - \frac{x^3}{8} - \frac{x^5}{640} + O(x^6) \quad (2)$$

Verification of solutions

$$y = x - \frac{x^3}{8} - \frac{x^5}{640} + O(x^6)$$

Verified OK.

$$y = x - \frac{x^3}{8} - \frac{x^5}{640} + O(x^6)$$

Verified OK.

13.20.2 Maple step by step solution

Let's solve

$$\left[(-x^2 + 4)y'' + y'x + 2y = 0, y(0) = 0, y'|_{\{x=0\}} = 1 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{xy'}{x^2-4} + \frac{2y}{x^2-4}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{xy'}{x^2-4} - \frac{2y}{x^2-4} = 0$$

- Check to see if x_0 is a regular singular point

- Define functions

$$\left[P_2(x) = -\frac{x}{x^2-4}, P_3(x) = -\frac{2}{x^2-4} \right]$$

- $(2+x) \cdot P_2(x)$ is analytic at $x = -2$

$$\left. ((2+x) \cdot P_2(x)) \right|_{x=-2} = -\frac{1}{2}$$

- $(2+x)^2 \cdot P_3(x)$ is analytic at $x = -2$

$$\left. ((2+x)^2 \cdot P_3(x)) \right|_{x=-2} = 0$$

- $x = -2$ is a regular singular point

Check to see if x_0 is a regular singular point

$$x_0 = -2$$

- Multiply by denominators

$$y''(x^2 - 4) - y'x - 2y = 0$$

- Change variables using $x = u - 2$ so that the regular singular point is at $u = 0$

$$(u^2 - 4u) \left(\frac{d^2}{du^2} y(u) \right) + (-u + 2) \left(\frac{d}{du} y(u) \right) - 2y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k+1-m$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k+1-m+r) u^{k+r}$$

- Convert $u^m \cdot \left(\frac{d^2}{du^2} y(u) \right)$ to series expansion for $m = 1..2$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) u^{k+r-2+m}$$

- Shift index using $k \rightarrow k+2-m$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=-2+m}^{\infty} a_{k+2-m} (k+2-m+r)(k+1-m+r) u^{k+r}$$

Rewrite ODE with series expansions

$$-2a_0r(-3+2r)u^{-1+r} + \left(\sum_{k=0}^{\infty} (-2a_{k+1}(k+1+r)(2k-1+2r) + a_k(k^2+2kr+r^2-2k-2r) \right)$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$-2r(-3+2r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \left\{0, \frac{3}{2}\right\}$$

- Each term in the series must be 0, giving the recursion relation

$$-4\left(k - \frac{1}{2} + r\right)(k+1+r)a_{k+1} + (k^2 + (2r-2)k + r^2 - 2r - 2)a_k = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+1} = \frac{(k^2+2kr+r^2-2k-2r-2)a_k}{2(2k-1+2r)(k+1+r)}$$

- Recursion relation for $r = 0$

$$a_{k+1} = \frac{(k^2-2k-2)a_k}{2(2k-1)(k+1)}$$

- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{k+1} = \frac{(k^2-2k-2)a_k}{2(2k-1)(k+1)} \right]$$

- Revert the change of variables $u = 2 + x$

$$\left[y = \sum_{k=0}^{\infty} a_k (2+x)^k, a_{k+1} = \frac{(k^2-2k-2)a_k}{2(2k-1)(k+1)} \right]$$

- Recursion relation for $r = \frac{3}{2}$

$$a_{k+1} = \frac{(k^2+k-\frac{11}{4})a_k}{2(2k+2)(k+\frac{5}{2})}$$

- Solution for $r = \frac{3}{2}$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+\frac{3}{2}}, a_{k+1} = \frac{(k^2+k-\frac{11}{4})a_k}{2(2k+2)(k+\frac{5}{2})} \right]$$

- Revert the change of variables $u = 2 + x$

$$\left[y = \sum_{k=0}^{\infty} a_k (2+x)^{k+\frac{3}{2}}, a_{k+1} = \frac{(k^2+k-\frac{11}{4})a_k}{2(2k+2)(k+\frac{5}{2})} \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k (2+x)^k \right) + \left(\sum_{k=0}^{\infty} b_k (2+x)^{k+\frac{3}{2}} \right), a_{k+1} = \frac{(k^2-2k-2)a_k}{2(2k-1)(k+1)}, b_{k+1} = \frac{(k^2+k-\frac{11}{4})b_k}{2(2k+2)(k+\frac{5}{2})} \right]$$

Maple trace Kovacic algorithm successful

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacics algorithm
  A Liouvillian solution exists
  Group is reducible or imprimitive
  Solution has integrals. Trying a special function solution free of integrals...
-> Trying a solution in terms of special functions:
  -> Bessel
  -> elliptic
  -> Legendre
  <- Legendre successful
<- special function solution successful
  -> Trying to convert hypergeometric functions to elementary form...
  <- elementary form could result into a too large expression - returning special functi
<- Kovacics algorithm successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
Order:=6;
dsolve([(4-x^2)*diff(y(x),x$2)+x*diff(y(x),x)+2*y(x)=0,y(0) = 0, D(y)(0) = 1],y(x),type='ser
```

$$y(x) = x - \frac{1}{8}x^3 - \frac{1}{640}x^5 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 19

```
AsymptoticDSolveValue[{{(4-x^2)*y'[x]+x*y'[x]+2*y[x]==0,{y[0]==0,y'[0]==1}},y[x],{x,0,5}]
```

$$y(x) \rightarrow -\frac{x^5}{640} - \frac{x^3}{8} + x$$

13.21 problem 27

13.21.1 Existence and uniqueness analysis	3305
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Internal problem ID [731]

Internal file name [OUTPUT/731_Sunday_June_05_2022_01_48_05_AM_12870666/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 27.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "**second_order_bessel_ode**", "**second order series method. Ordinary point**", "**second order series method. Taylor series method**"

Maple gives the following as the ode type

```
[[_Emden, _Fowler]]
```

$$y'' + x^2y = 0$$

With initial conditions

$$[y(0) = 1, y'(0) = 0]$$

With the expansion point for the power series method at $x = 0$.

13.21.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = 0$$

$$q(x) = x^2$$

$$F = 0$$

Hence the ode is

$$y'' + x^2y = 0$$

The domain of $p(x) = 0$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = x^2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \dots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (804)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (805)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -x^2y \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= -(2y + y'x)x \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= yx^4 - 4y'x - 2y \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= y'x^4 + 8yx^3 - 6y' \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= 12x^3y' - x^2y(x^4 - 30)
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = 1$ and $y'(0) = 0$ gives

$$\begin{aligned}
 F_0 &= 0 \\
 F_1 &= 0 \\
 F_2 &= -2 \\
 F_3 &= 0 \\
 F_4 &= 0
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = 1 - \frac{x^4}{12} + O(x^6)$$

$$y = 1 - \frac{x^4}{12} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = -x^2 \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=0}^{\infty} x^{n+2} a_n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

$$\sum_{n=0}^{\infty} x^{n+2} a_n = \sum_{n=2}^{\infty} a_{n-2} x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \left(\sum_{n=2}^{\infty} a_{n-2} x^n \right) = 0 \quad (3)$$

For $2 \leq n$, the recurrence equation is

$$(n+2) a_{n+2} (n+1) + a_{n-2} = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_{n-2}}{(n+2)(n+1)} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$12a_4 + a_0 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = -\frac{a_0}{12}$$

For $n = 3$ the recurrence equation gives

$$20a_5 + a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = -\frac{a_1}{20}$$

For $n = 4$ the recurrence equation gives

$$30a_6 + a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = 0$$

For $n = 5$ the recurrence equation gives

$$42a_7 + a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = 0$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1x - \frac{1}{12}a_0x^4 - \frac{1}{20}a_1x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - \frac{x^4}{12}\right)a_0 + \left(x - \frac{1}{20}x^5\right)a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - \frac{x^4}{12}\right)c_1 + \left(x - \frac{1}{20}x^5\right)c_2 + O(x^6)$$

$$y = 1 - \frac{x^4}{12} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = 1 - \frac{x^4}{12} + O(x^6) \quad (1)$$

$$y = 1 - \frac{x^4}{12} + O(x^6) \quad (2)$$

Verification of solutions

$$y = 1 - \frac{x^4}{12} + O(x^6)$$

Verified OK.

$$y = 1 - \frac{x^4}{12} + O(x^6)$$

Verified OK.

13.21.2 Maple step by step solution

Let's solve

$$\left[y'' = -x^2 y, y(0) = 1, y'|_{\{x=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + x^2 y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite ODE with series expansions

- Convert $x^2 \cdot y$ to series expansion

$$x^2 \cdot y = \sum_{k=0}^{\infty} a_k x^{k+2}$$

- Shift index using $k- > k - 2$

$$x^2 \cdot y = \sum_{k=2}^{\infty} a_{k-2} x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k- > k + 2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2} (k+2)(k+1) x^k$$

Rewrite ODE with series expansions

$$6a_3 x + 2a_2 + \left(\sum_{k=2}^{\infty} (a_{k+2} (k+2)(k+1) + a_{k-2}) x^k \right) = 0$$

- The coefficients of each power of x must be 0

$$[2a_2 = 0, 6a_3 = 0]$$

- Solve for the dependent coefficient(s)

$$\{a_2 = 0, a_3 = 0\}$$

- Each term in the series must be 0, giving the recursion relation

$$(k^2 + 3k + 2) a_{k+2} + a_{k-2} = 0$$

- Shift index using $k \rightarrow k + 2$

$$((k + 2)^2 + 3k + 8) a_{k+4} + a_k = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+4} = -\frac{a_k}{k^2 + 7k + 12}, a_2 = 0, a_3 = 0 \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
    -> Bessel
    <- Bessel successful
<- special function solution successful`

```

✓ Solution by Maple

Time used: 0.282 (sec). Leaf size: 12

```

Order:=6;
dsolve([diff(y(x),x$2)+x^2*y(x)=0,y(0) = 1, D(y)(0) = 0],y(x),type='series',x=0);

```

$$y(x) = 1 - \frac{1}{12}x^4 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 12

```
AsymptoticDSolveValue[{y'[x]+x^2*y[x]==0,{y[0]==1,y'[0]==0}},y[x],{x,0,5}]
```

$$y(x) \rightarrow 1 - \frac{x^4}{12}$$

13.22 problem 28

13.22.1 Existence and uniqueness analysis	3316
13.22.2 Maple step by step solution	3324

Internal problem ID [732]

Internal file name [OUTPUT/732_Sunday_June_05_2022_01_48_08_AM_85496591/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.2, Series Solutions Near an Ordinary Point, Part I. page 263

Problem number: 28.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order, _with_linear_symmetries], [_2nd_order, _linear, `
  _with_symmetry_[0,F(x)]`]]
```

$$(1 - x)y'' + y'x - 2y = 0$$

With initial conditions

$$[y(0) = 0, y'(0) = 1]$$

With the expansion point for the power series method at $x = 0$.

13.22.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = \frac{x}{1-x}$$
$$q(x) = -\frac{2}{1-x}$$
$$F = 0$$

Hence the ode is

$$y'' + \frac{xy'}{1-x} - \frac{2y}{1-x} = 0$$

The domain of $p(x) = \frac{x}{1-x}$ is

$$\{x < 1 \vee 1 < x\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = -\frac{2}{1-x}$ is

$$\{x < 1 \vee 1 < x\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (807)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (808)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= \frac{y'x - 2y}{x - 1} \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\
 &= \frac{(x - 1)y' - 2y}{x - 1} \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\
 &= \frac{((x - 1)y' - 2y)(-2 + x)}{(x - 1)^2} \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\
 &= \frac{(x^2 - 4x + 5)((x - 1)y' - 2y)}{(x - 1)^3} \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\
 &= \frac{(x^3 - 6x^2 + 15x - 16)((x - 1)y' - 2y)}{(x - 1)^4}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = 0$ and $y'(0) = 1$ gives

$$F_0 = 0$$

$$F_1 = 1$$

$$F_2 = 2$$

$$F_3 = 5$$

$$F_4 = 16$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = x + \frac{x^3}{6} + \frac{x^4}{12} + \frac{x^5}{24} + \frac{x^6}{45} + O(x^6)$$

$$y = x + \frac{x^3}{6} + \frac{x^4}{12} + \frac{x^5}{24} + \frac{x^6}{45} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. The ode is normalized to be

$$(1 - x)y'' + y'x - 2y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(1 - x) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x - 2 \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\sum_{n=2}^{\infty} (-n x^{n-1} a_n (n-1)) + \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \sum_{n=0}^{\infty} (-2 a_n x^n) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} (-n x^{n-1} a_n (n-1)) = \sum_{n=1}^{\infty} (-(n+1) a_{n+1} n x^n)$$

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} \sum_{n=1}^{\infty} (-(n+1) a_{n+1} n x^n) + \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) \\ + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \sum_{n=0}^{\infty} (-2a_n x^n) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$2a_2 - 2a_0 = 0$$

$$a_2 = a_0$$

For $1 \leq n$, the recurrence equation is

$$-(n+1) a_{n+1} n + (n+2) a_{n+2} (n+1) + n a_n - 2a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} a_{n+2} &= \frac{n^2 a_{n+1} - n a_n + n a_{n+1} + 2a_n}{(n+2)(n+1)} \\ (5) \quad &= \frac{(-n+2) a_n}{(n+2)(n+1)} + \frac{(n^2+n) a_{n+1}}{(n+2)(n+1)} \end{aligned}$$

For $n = 1$ the recurrence equation gives

$$-2a_2 + 6a_3 - a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = \frac{a_0}{3} + \frac{a_1}{6}$$

For $n = 2$ the recurrence equation gives

$$-6a_3 + 12a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{6} + \frac{a_1}{12}$$

For $n = 3$ the recurrence equation gives

$$-12a_4 + 20a_5 + a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_0}{12} + \frac{a_1}{24}$$

For $n = 4$ the recurrence equation gives

$$-20a_5 + 30a_6 + 2a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{2a_0}{45} + \frac{a_1}{45}$$

For $n = 5$ the recurrence equation gives

$$-30a_6 + 42a_7 + 3a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{13a_0}{504} + \frac{13a_1}{1008}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x + a_0 x^2 + \left(\frac{a_0}{3} + \frac{a_1}{6}\right) x^3 + \left(\frac{a_0}{6} + \frac{a_1}{12}\right) x^4 + \left(\frac{a_0}{12} + \frac{a_1}{24}\right) x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + x^2 + \frac{1}{3}x^3 + \frac{1}{6}x^4 + \frac{1}{12}x^5\right) a_0 + \left(x + \frac{1}{6}x^3 + \frac{1}{12}x^4 + \frac{1}{24}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 + x^2 + \frac{1}{3}x^3 + \frac{1}{6}x^4 + \frac{1}{12}x^5\right) c_1 + \left(x + \frac{1}{6}x^3 + \frac{1}{12}x^4 + \frac{1}{24}x^5\right) c_2 + O(x^6)$$

$$y = x + \frac{x^3}{6} + \frac{x^4}{12} + \frac{x^5}{24} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = x + \frac{x^3}{6} + \frac{x^4}{12} + \frac{x^5}{24} + \frac{x^6}{45} + O(x^6) \quad (1)$$

$$y = x + \frac{x^3}{6} + \frac{x^4}{12} + \frac{x^5}{24} + O(x^6) \quad (2)$$

Verification of solutions

$$y = x + \frac{x^3}{6} + \frac{x^4}{12} + \frac{x^5}{24} + \frac{x^6}{45} + O(x^6)$$

Verified OK.

$$y = x + \frac{x^3}{6} + \frac{x^4}{12} + \frac{x^5}{24} + O(x^6)$$

Verified OK.

13.22.2 Maple step by step solution

Let's solve

$$\left[(1-x)y'' + y'x - 2y = 0, y(0) = 0, y'|_{\{x=0\}} = 1 \right]$$

- Highest derivative means the order of the ODE is 2
 y''
- Isolate 2nd derivative

$$y'' = -\frac{2y}{x-1} + \frac{xy'}{x-1}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' - \frac{xy'}{x-1} + \frac{2y}{x-1} = 0$$

- Check to see if $x_0 = 1$ is a regular singular point

- Define functions

$$[P_2(x) = -\frac{x}{x-1}, P_3(x) = \frac{2}{x-1}]$$

- $(x-1) \cdot P_2(x)$ is analytic at $x = 1$

$$((x-1) \cdot P_2(x)) \Big|_{x=1} = -1$$

- $(x-1)^2 \cdot P_3(x)$ is analytic at $x = 1$

$$((x-1)^2 \cdot P_3(x)) \Big|_{x=1} = 0$$

- $x = 1$ is a regular singular point

Check to see if $x_0 = 1$ is a regular singular point

$$x_0 = 1$$

- Multiply by denominators

$$(x-1)y'' - y'x + 2y = 0$$

- Change variables using $x = u + 1$ so that the regular singular point is at $u = 0$

$$u \left(\frac{d^2}{du^2} y(u) \right) + (-u-1) \left(\frac{d}{du} y(u) \right) + 2y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k+1-m$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k+1-m+r) u^{k+r}$$

- Convert $u \cdot \left(\frac{d^2}{du^2} y(u) \right)$ to series expansion

$$u \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) u^{k+r-1}$$

- Shift index using $k \rightarrow k+1$

$$u \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=-1}^{\infty} a_{k+1} (k+1+r)(k+r) u^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r(-2+r) u^{-1+r} + \left(\sum_{k=0}^{\infty} (a_{k+1} (k+1+r)(k+r-1) - a_k (k+r-2)) u^{k+r} \right) = 0$$

- a_0 cannot be 0 by assumption, giving the indicial equation
 $r(-2+r) = 0$
- Values of r that satisfy the indicial equation
 $r \in \{0, 2\}$
- Each term in the series must be 0, giving the recursion relation
 $a_{k+1} (k+1+r)(k+r-1) - a_k (k+r-2) = 0$
- Recursion relation that defines series solution to ODE

$$a_{k+1} = \frac{a_k (k+r-2)}{(k+1+r)(k+r-1)}$$

- Recursion relation for $r = 0$; series terminates at $k = 2$

$$a_{k+1} = \frac{a_k (k-2)}{(k+1)(k-1)}$$

- Series not valid for $r = 0$, division by 0 in the recursion relation at $k = 1$

$$a_{k+1} = \frac{a_k (k-2)}{(k+1)(k-1)}$$

- Recursion relation for $r = 2$

$$a_{k+1} = \frac{a_k k}{(k+3)(k+1)}$$

- Solution for $r = 2$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+2}, a_{k+1} = \frac{a_k k}{(k+3)(k+1)} \right]$$

- Revert the change of variables $u = x - 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x-1)^{k+2}, a_{k+1} = \frac{a_k k}{(k+3)(k+1)} \right]$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
<- linear_1 successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 16

```
Order:=6;  
dsolve([(1-x)*diff(y(x),x$2)+x*diff(y(x),x)-2*y(x)=0,y(0) = 0, D(y)(0) = 1],y(x),type='series')
```

$$y(x) = x + \frac{1}{6}x^3 + \frac{1}{12}x^4 + \frac{1}{24}x^5 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 26

```
AsymptoticDSolveValue[{(1-x)*y''[x]+x*y'[x]-2*y[x]==0,{y[0]==0,y'[0]==1}},y[x],{x,0,5}]
```

$$y(x) \rightarrow \frac{x^5}{24} + \frac{x^4}{12} + \frac{x^3}{6} + x$$

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14.1 problem 1

- 14.1.1 Existence and uniqueness analysis 3329
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Internal problem ID [733]

Internal file name [OUTPUT/733_Sunday_June_05_2022_01_48_10_AM_75056564/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 1.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact linear second order ode", "second_order_integrable_as_is", "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _exact , _linear , _homogeneous]]
```

$$y'' + y'x + y = 0$$

With initial conditions

$$[y(0) = 1, y'(0) = 0]$$

With the expansion point for the power series method at $x = 0$.

14.1.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = x$$

$$q(x) = 1$$

$$F = 0$$

Hence the ode is

$$y'' + y'x + y = 0$$

The domain of $p(x) = x$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = 1$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \dots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (810)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (811)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -y'x - y \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= y'x^2 + yx - 2y' \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= -x^3y' - x^2y + 5y'x + 3y \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= (x^4 - 9x^2 + 8)y' + xy(x^2 - 7) \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= (-x^5 + 14x^3 - 33x)y' - y(x^4 - 12x^2 + 15)
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = 1$ and $y'(0) = 0$ gives

$$\begin{aligned}
 F_0 &= -1 \\
 F_1 &= 0 \\
 F_2 &= 3 \\
 F_3 &= 0 \\
 F_4 &= -15
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = -\frac{x^2}{2} + 1 + \frac{x^4}{8} - \frac{x^6}{48} + O(x^6)$$

$$y = -\frac{x^2}{2} + 1 + \frac{x^4}{8} - \frac{x^6}{48} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = - \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x - \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n x^n a_n \right) + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \left(\sum_{n=1}^{\infty} n x^n a_n \right) + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 + a_0 = 0$$

$$a_2 = -\frac{a_0}{2}$$

For $1 \leq n$, the recurrence equation is

$$(n+2)a_{n+2}(n+1) + na_n + a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n}{n+2} \quad (5)$$

For $n = 1$ the recurrence equation gives

$$6a_3 + 2a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = -\frac{a_1}{3}$$

For $n = 2$ the recurrence equation gives

$$12a_4 + 3a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{8}$$

For $n = 3$ the recurrence equation gives

$$20a_5 + 4a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_1}{15}$$

For $n = 4$ the recurrence equation gives

$$30a_6 + 5a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{a_0}{48}$$

For $n = 5$ the recurrence equation gives

$$42a_7 + 6a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{a_1}{105}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - \frac{1}{2} a_0 x^2 - \frac{1}{3} a_1 x^3 + \frac{1}{8} a_0 x^4 + \frac{1}{15} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(-\frac{1}{2} x^2 + 1 + \frac{1}{8} x^4 \right) a_0 + \left(x - \frac{1}{3} x^3 + \frac{1}{15} x^5 \right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(-\frac{1}{2} x^2 + 1 + \frac{1}{8} x^4 \right) c_1 + \left(x - \frac{1}{3} x^3 + \frac{1}{15} x^5 \right) c_2 + O(x^6)$$

$$y = -\frac{x^2}{2} + 1 + \frac{x^4}{8} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = -\frac{x^2}{2} + 1 + \frac{x^4}{8} - \frac{x^6}{48} + O(x^6) \quad (1)$$

$$y = -\frac{x^2}{2} + 1 + \frac{x^4}{8} + O(x^6) \quad (2)$$

Verification of solutions

$$y = -\frac{x^2}{2} + 1 + \frac{x^4}{8} - \frac{x^6}{48} + O(x^6)$$

Verified OK.

$$y = -\frac{x^2}{2} + 1 + \frac{x^4}{8} + O(x^6)$$

Verified OK.

14.1.2 Maple step by step solution

Let's solve

$$\left[y'' = -y'x - y, y(0) = 1, y'|_{\{x=0\}} = 0 \right]$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + y'x + y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite DE with series expansions

- Convert $x \cdot y'$ to series expansion

$$x \cdot y' = \sum_{k=0}^{\infty} a_k k x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k- > k+2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2} (k+2)(k+1) x^k$$

Rewrite DE with series expansions

$$\sum_{k=0}^{\infty} (a_{k+2}(k+2)(k+1) + a_k(k+1)) x^k = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(k+1)(a_{k+2}(k+2) + a_k) = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = -\frac{a_k}{k+2} \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
<- linear_1 successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```

Order:=6;
dsolve([diff(y(x),x$2)+x*diff(y(x),x)+y(x)=0,y(0) = 1, D(y)(0) = 0],y(x),type='series',x=0);

```

$$y(x) = 1 - \frac{1}{2}x^2 + \frac{1}{8}x^4 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 19

```

AsymptoticDSolveValue[{y'[x]+x*y'[x]+y[x]==0,{y[0]==1,y'[0]==0}},y[x],{x,0,5}]

```

$$y(x) \rightarrow \frac{x^4}{8} - \frac{x^2}{2} + 1$$

14.2 problem 2

14.2.1 Existence and uniqueness analysis 3339

Internal problem ID [734]

Internal file name [OUTPUT/734_Sunday_June_05_2022_01_48_12_AM_11361222/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 2.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "exact linear second order ode", "second_order_integrable_as_is", "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _exact , _linear , _homogeneous]]
```

$$y'' + \sin(x)y' + \cos(x)y = 0$$

With initial conditions

$$[y(0) = 0, y'(0) = 1]$$

With the expansion point for the power series method at $x = 0$.

14.2.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = \sin(x)$$

$$q(x) = \cos(x)$$

$$F = 0$$

Hence the ode is

$$y'' + \sin(x) y' + \cos(x) y = 0$$

The domain of $p(x) = \sin(x)$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = \cos(x)$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0) y'(x_0) + \frac{(x - x_0)^2}{2} y''(x_0) + \frac{(x - x_0)^3}{3!} y'''(x_0) + \dots \\ &= y_0 + x y'_0 + \frac{x^2}{2} f|_{x_0, y_0, y'_0} + \frac{x^3}{3!} f'|_{x_0, y_0, y'_0} + \dots \\ &= y_0 + x y'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (813)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (814)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = -\sin(x) y' - \cos(x) y$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dx} \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\ &= (\sin(x)^2 - 2\cos(x)) y' + y \sin(x) (\cos(x) + 1) \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dx} \\ &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\ &= (\cos(x)^2 + 5\cos(x) + 2) \sin(x) y' + (\cos(x)^3 + 4\cos(x)^2 - 1) y \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dx} \\ &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\ &= (\cos(x)^4 + 9\cos(x)^3 + 15\cos(x)^2 - 5\cos(x) - 8) y' - \cos(x) \sin(x) y (\cos(x)^2 + 8\cos(x) + 10) \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dx} \\ &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\ &= -\sin(x) (\cos(x)^4 + 14\cos(x)^3 + 50\cos(x)^2 + 35\cos(x) - 13) y' - y (\cos(x)^5 + 13\cos(x)^4 + 39\cos(x)^3 + 25\cos(x)^2 - 13\cos(x) - 8) \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = 0$ and $y'(0) = 1$ gives

$$F_0 = 0$$

$$F_1 = -2$$

$$F_2 = 0$$

$$F_3 = 12$$

$$F_4 = 0$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = -\frac{x^3}{3} + x + \frac{x^5}{10} + O(x^6)$$

$$y = -\frac{x^3}{3} + x + \frac{x^5}{10} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = -\sin(x) \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) - \cos(x) \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Expanding $\sin(x)$ as Taylor series around $x = 0$ and keeping only the first 6 terms gives

$$\begin{aligned} \sin(x) &= x - \frac{1}{6}x^3 + \frac{1}{120}x^5 - \frac{1}{5040}x^7 + \dots \\ &= x - \frac{1}{6}x^3 + \frac{1}{120}x^5 - \frac{1}{5040}x^7 \end{aligned}$$

Expanding $\cos(x)$ as Taylor series around $x = 0$ and keeping only the first 6 terms gives

$$\begin{aligned} \cos(x) &= -\frac{1}{2}x^2 + 1 + \frac{1}{24}x^4 - \frac{1}{720}x^6 + \dots \\ &= -\frac{1}{2}x^2 + 1 + \frac{1}{24}x^4 - \frac{1}{720}x^6 \end{aligned}$$

Hence the ODE in Eq (1) becomes

$$\begin{aligned} &\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(x - \frac{1}{6}x^3 + \frac{1}{120}x^5 - \frac{1}{5040}x^7 \right) \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) \\ &+ \left(-\frac{1}{2}x^2 + 1 + \frac{1}{24}x^4 - \frac{1}{720}x^6 \right) \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \end{aligned}$$

Expanding the second term in (1) gives

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + x \cdot \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) - \frac{x^3}{6} \\ & \cdot \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) + \frac{x^5}{120} \cdot \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) - \frac{x^7}{5040} \\ & \cdot \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) + \left(-\frac{1}{2}x^2 + 1 + \frac{1}{24}x^4 - \frac{1}{720}x^6 \right) \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \end{aligned}$$

Expanding the third term in (1) gives

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + x \cdot \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) - \frac{x^3}{6} \cdot \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) \\ & + \frac{x^5}{120} \cdot \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) - \frac{x^7}{5040} \cdot \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) + -\frac{x^2}{2} \cdot \left(\sum_{n=0}^{\infty} a_n x^n \right) \\ & + 1 \cdot \left(\sum_{n=0}^{\infty} a_n x^n \right) + \frac{x^4}{24} \cdot \left(\sum_{n=0}^{\infty} a_n x^n \right) - \frac{x^6}{720} \cdot \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \end{aligned}$$

Which simplifies to

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \sum_{n=1}^{\infty} \left(-\frac{n x^{n+2} a_n}{6} \right) \\ & + \left(\sum_{n=1}^{\infty} \frac{n x^{n+4} a_n}{120} \right) + \sum_{n=1}^{\infty} \left(-\frac{n x^{n+6} a_n}{5040} \right) + \sum_{n=0}^{\infty} \left(-\frac{x^{n+2} a_n}{2} \right) \\ & + \left(\sum_{n=0}^{\infty} a_n x^n \right) + \left(\sum_{n=0}^{\infty} \frac{x^{n+4} a_n}{24} \right) + \sum_{n=0}^{\infty} \left(-\frac{x^{n+6} a_n}{720} \right) = 0 \end{aligned} \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\begin{aligned} \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} &= \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \\ \sum_{n=1}^{\infty} \left(-\frac{n x^{n+2} a_n}{6} \right) &= \sum_{n=3}^{\infty} \left(-\frac{(n-2) a_{n-2} x^n}{6} \right) \end{aligned}$$

$$\begin{aligned} \sum_{n=1}^{\infty} \frac{n x^{n+4} a_n}{120} &= \sum_{n=5}^{\infty} \frac{(n-4) a_{n-4} x^n}{120} \\ \sum_{n=1}^{\infty} \left(-\frac{n x^{n+6} a_n}{5040} \right) &= \sum_{n=7}^{\infty} \left(-\frac{(n-6) a_{n-6} x^n}{5040} \right) \\ \sum_{n=0}^{\infty} \left(-\frac{x^{n+2} a_n}{2} \right) &= \sum_{n=2}^{\infty} \left(-\frac{a_{n-2} x^n}{2} \right) \\ \sum_{n=0}^{\infty} \frac{x^{n+4} a_n}{24} &= \sum_{n=4}^{\infty} \frac{a_{n-4} x^n}{24} \\ \sum_{n=0}^{\infty} \left(-\frac{x^{n+6} a_n}{720} \right) &= \sum_{n=6}^{\infty} \left(-\frac{a_{n-6} x^n}{720} \right) \end{aligned}$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} &\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \sum_{n=3}^{\infty} \left(-\frac{(n-2) a_{n-2} x^n}{6} \right) \\ &+ \left(\sum_{n=5}^{\infty} \frac{(n-4) a_{n-4} x^n}{120} \right) + \sum_{n=7}^{\infty} \left(-\frac{(n-6) a_{n-6} x^n}{5040} \right) + \sum_{n=2}^{\infty} \left(-\frac{a_{n-2} x^n}{2} \right) \quad (3) \\ &+ \left(\sum_{n=0}^{\infty} a_n x^n \right) + \left(\sum_{n=4}^{\infty} \frac{a_{n-4} x^n}{24} \right) + \sum_{n=6}^{\infty} \left(-\frac{a_{n-6} x^n}{720} \right) = 0 \end{aligned}$$

$n = 0$ gives

$$2a_2 + a_0 = 0$$

$$a_2 = -\frac{a_0}{2}$$

$n = 1$ gives

$$6a_3 + 2a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{a_1}{3}$$

$n = 2$ gives

$$12a_4 + 3a_2 - \frac{a_0}{2} = 0$$

Which after substituting earlier equations, simplifies to

$$a_4 = \frac{a_0}{6}$$

$n = 3$ gives

$$20a_5 + 4a_3 - \frac{2a_1}{3} = 0$$

Which after substituting earlier equations, simplifies to

$$a_5 = \frac{a_1}{10}$$

$n = 4$ gives

$$30a_6 + 5a_4 - \frac{5a_2}{6} + \frac{a_0}{24} = 0$$

Which after substituting earlier equations, simplifies to

$$a_6 = -\frac{31a_0}{720}$$

$n = 5$ gives

$$42a_7 + 6a_5 - a_3 + \frac{a_1}{20} = 0$$

Which after substituting earlier equations, simplifies to

$$a_7 = -\frac{59a_1}{2520}$$

For $7 \leq n$, the recurrence equation is

$$\begin{aligned} (n+2)a_{n+2}(n+1) + na_n - \frac{(n-2)a_{n-2}}{6} + \frac{(n-4)a_{n-4}}{120} \\ - \frac{(n-6)a_{n-6}}{5040} - \frac{a_{n-2}}{2} + a_n + \frac{a_{n-4}}{24} - \frac{a_{n-6}}{720} = 0 \end{aligned} \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned}
 a_{n+2} &= -\frac{5040a_n - a_{n-6} + 42a_{n-4} - 840a_{n-2}}{5040(n+2)} \\
 (5) \qquad &= -\frac{a_n}{n+2} + \frac{a_{n-6}}{5040n+10080} - \frac{a_{n-4}}{120(n+2)} + \frac{a_{n-2}}{6n+12}
 \end{aligned}$$

And so on. Therefore the solution is

$$\begin{aligned}
 y &= \sum_{n=0}^{\infty} a_n x^n \\
 &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots
 \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - \frac{1}{2} a_0 x^2 - \frac{1}{3} a_1 x^3 + \frac{1}{6} a_0 x^4 + \frac{1}{10} a_1 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - \frac{1}{2}x^2 + \frac{1}{6}x^4\right) a_0 + \left(-\frac{1}{3}x^3 + x + \frac{1}{10}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - \frac{1}{2}x^2 + \frac{1}{6}x^4\right) c_1 + \left(-\frac{1}{3}x^3 + x + \frac{1}{10}x^5\right) c_2 + O(x^6)$$

$$y = -\frac{x^3}{3} + x + \frac{x^5}{10} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = -\frac{x^3}{3} + x + \frac{x^5}{10} + O(x^6) \quad (1)$$

$$y = -\frac{x^3}{3} + x + \frac{x^5}{10} + O(x^6) \quad (2)$$

Verification of solutions

$$y = -\frac{x^3}{3} + x + \frac{x^5}{10} + O(x^6)$$

Verified OK.

$$y = -\frac{x^3}{3} + x + \frac{x^5}{10} + O(x^6)$$

Verified OK.

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a symmetry of the form [xi=0, eta=F(x)]  
  One independent solution has integrals. Trying a hypergeometric solution free of integral  
  -> hyper3: Equivalence to 2F1, 1F1 or 0F1 under a power @ Moebius  
No hypergeometric solution was found.  
<- linear_1 successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
Order:=6;  
dsolve([diff(y(x),x$2)+sin(x)*diff(y(x),x)+cos(x)*y(x)=0,y(0) = 0, D(y)(0) = 1],y(x),type='s'
```

$$y(x) = x - \frac{1}{3}x^3 + \frac{1}{10}x^5 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 19

```
AsymptoticDSolveValue[{y'[x]+Sin[x]*y'[x]+Cos[x]*y[x]==0,{y[0]==0,y'[0]==1}},y[x],{x,0,5}]
```

$$y(x) \rightarrow \frac{x^5}{10} - \frac{x^3}{3} + x$$

14.3 problem 3

14.3.1 Existence and uniqueness analysis 3350

Internal problem ID [735]

Internal file name [OUTPUT/735_Sunday_June_05_2022_01_48_14_AM_17099555/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 3.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$x^2 y'' + (x + 1) y' + 3 \ln(x) y = 0$$

With initial conditions

$$[y(1) = 2, y'(1) = 0]$$

With the expansion point for the power series method at $x = 1$.

14.3.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = \frac{x + 1}{x^2}$$
$$q(x) = \frac{3 \ln(x)}{x^2}$$
$$F = 0$$

Hence the ode is

$$y'' + \frac{(x+1)y'}{x^2} + \frac{3\ln(x)y}{x^2} = 0$$

The domain of $p(x) = \frac{x+1}{x^2}$ is

$$\{x < 0 \vee 0 < x\}$$

And the point $x_0 = 1$ is inside this domain. The domain of $q(x) = \frac{3\ln(x)}{x^2}$ is

$$\{0 < x\}$$

And the point $x_0 = 1$ is also inside this domain. Hence solution exists and is unique.

The ode does not have its expansion point at $x = 0$, therefore to simplify the computation of power series expansion, change of variable is made on the independent variable to shift the initial conditions and the expansion point back to zero. The new ode is then solved more easily since the expansion point is now at zero. The solution converted back to the original independent variable. Let

$$t = x - 1$$

The ode is converted to be in terms of the new independent variable t . This results in

$$(t+1)^2 \left(\frac{d^2}{dt^2} y(t) \right) + (2+t) \left(\frac{d}{dt} y(t) \right) + 3\ln(t+1)y(t) = 0$$

With its expansion point and initial conditions now at $t = 0$. With initial conditions now becoming

$$\begin{aligned} y(0) &= 2 \\ y'(0) &= 0 \end{aligned}$$

The transformed ODE is now solved. Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the

case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned}
 y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\
 &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\
 &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0}
 \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \tag{1}$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \tag{816}$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \tag{817}$$

$$\begin{aligned}
 \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\
 &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \tag{2}
 \end{aligned}$$

$$\begin{aligned}
 \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\
 &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \tag{3}
 \end{aligned}$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = -\frac{t\left(\frac{d}{dt}y(t)\right) + 3\ln(t+1)y(t) + 2\frac{d}{dt}y(t)}{(t+1)^2}$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dt} \\ &= \frac{\partial F_0}{\partial t} + \frac{\partial F_0}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_0}{\partial \frac{d}{dt}y(t)} F_0 \\ &= \frac{(-3(t+1)^2\left(\frac{d}{dt}y(t)\right) + (9t+12)y(t)\ln(t+1) + (2t^2+8t+7)\left(\frac{d}{dt}y(t)\right) - 3(t+1)y(t))}{(t+1)^4} \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dt} \\ &= \frac{\partial F_1}{\partial t} + \frac{\partial F_1}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_1}{\partial \frac{d}{dt}y(t)} F_1 \\ &= \frac{9(t+1)^2 y(t) \ln(t+1)^2 + (18(t+1)^2\left(t+\frac{4}{3}\right)\left(\frac{d}{dt}y(t)\right) + (-33t^2-90t-60)y(t)\ln(t+1) + (-12t^2-12t-6)y(t))}{(t+1)^6} \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dt} \\ &= \frac{\partial F_2}{\partial t} + \frac{\partial F_2}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_2}{\partial \frac{d}{dt}y(t)} F_2 \\ &= \frac{9(t+1)^2\left((t+1)^2\left(\frac{d}{dt}y(t)\right) + (-10t-12)y(t)\right)\ln(t+1)^2 + ((-105t^4-492t^3-855t^2-654t-186)y(t))}{(t+1)^8} \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dt} \\ &= \frac{\partial F_3}{\partial t} + \frac{\partial F_3}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_3}{\partial \frac{d}{dt}y(t)} F_3 \\ &= \frac{-27y(t)(t+1)^4\ln(t+1)^3 - 135(t+1)^2\left(\left(t+\frac{6}{5}\right)(t+1)^2\left(\frac{d}{dt}y(t)\right) - \frac{17\left(t^2+\frac{206}{85}t+\frac{124}{85}\right)y(t)}{3}\right)\ln(t+1)^2 + \dots}{(t+1)^{10}} \end{aligned}$$

And so on. Evaluating all the above at initial conditions $t = 0$ and $y(0) = 2$ and $y'(0) = 0$ gives

$$F_0 = 0$$

$$F_1 = -6$$

$$F_2 = 42$$

$$F_3 = -294$$

$$F_4 = 2376$$

Substituting all the above in (7) and simplifying gives the solution as

$$y(t) = -t^3 + 2 + \frac{7t^4}{4} - \frac{49t^5}{20} + \frac{33t^6}{10} + O(t^6)$$

$$y(t) = -t^3 + 2 + \frac{7t^4}{4} - \frac{49t^5}{20} + \frac{33t^6}{10} + O(t^6)$$

Since the expansion point $t = 0$ is an ordinary, we can also solve this using standard power series The ode is normalized to be

$$(t^2 + 2t + 1) \left(\frac{d^2}{dt^2} y(t) \right) + (2 + t) \left(\frac{d}{dt} y(t) \right) + 3 \ln(t + 1) y(t) = 0$$

Let the solution be represented as power series of the form

$$y(t) = \sum_{n=0}^{\infty} a_n t^n$$

Then

$$\begin{aligned} \frac{d}{dt} y(t) &= \sum_{n=1}^{\infty} n a_n t^{n-1} \\ \frac{d^2}{dt^2} y(t) &= \sum_{n=2}^{\infty} n(n-1) a_n t^{n-2} \end{aligned}$$

Substituting the above back into the ode gives

$$(t^2 + 2t + 1) \left(\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2} \right) + (2 + t) \left(\sum_{n=1}^{\infty} n a_n t^{n-1} \right) + 3 \ln(t + 1) \left(\sum_{n=0}^{\infty} a_n t^n \right) = 0 \quad (1)$$

Expanding $3 \ln(t + 1)$ as Taylor series around $t = 0$ and keeping only the first 6 terms gives

$$\begin{aligned} 3 \ln(t + 1) &= 3t - \frac{3}{2}t^2 + t^3 - \frac{3}{4}t^4 + \frac{3}{5}t^5 - \frac{1}{2}t^6 + \dots \\ &= 3t - \frac{3}{2}t^2 + t^3 - \frac{3}{4}t^4 + \frac{3}{5}t^5 - \frac{1}{2}t^6 \end{aligned}$$

Hence the ODE in Eq (1) becomes

$$\begin{aligned} & (t^2 + 2t + 1) \left(\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2} \right) + (2+t) \left(\sum_{n=1}^{\infty} n a_n t^{n-1} \right) \\ & + \left(3t - \frac{3}{2}t^2 + t^3 - \frac{3}{4}t^4 + \frac{3}{5}t^5 - \frac{1}{2}t^6 \right) \left(\sum_{n=0}^{\infty} a_n t^n \right) = 0 \end{aligned}$$

Expanding the third term in (1) gives

$$\begin{aligned} & (t^2 + 2t + 1) \left(\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2} \right) + (2+t) \left(\sum_{n=1}^{\infty} n a_n t^{n-1} \right) \\ & + 3t \cdot \left(\sum_{n=0}^{\infty} a_n t^n \right) - \frac{3t^2}{2} \cdot \left(\sum_{n=0}^{\infty} a_n t^n \right) + t^3 \cdot \left(\sum_{n=0}^{\infty} a_n t^n \right) - \frac{3t^4}{4} \\ & \cdot \left(\sum_{n=0}^{\infty} a_n t^n \right) + \frac{3t^5}{5} \cdot \left(\sum_{n=0}^{\infty} a_n t^n \right) - \frac{t^6}{2} \cdot \left(\sum_{n=0}^{\infty} a_n t^n \right) = 0 \end{aligned}$$

Which simplifies to

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} t^n a_n n(n-1) \right) + \left(\sum_{n=2}^{\infty} 2n t^{n-1} a_n (n-1) \right) \\ & + \left(\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2} \right) + \left(\sum_{n=1}^{\infty} 2n a_n t^{n-1} \right) + \left(\sum_{n=1}^{\infty} n a_n t^n \right) \\ & + \left(\sum_{n=0}^{\infty} 3t^{1+n} a_n \right) + \sum_{n=0}^{\infty} \left(-\frac{3t^{n+2} a_n}{2} \right) + \left(\sum_{n=0}^{\infty} t^{n+3} a_n \right) \\ & + \sum_{n=0}^{\infty} \left(-\frac{3t^{n+4} a_n}{4} \right) + \left(\sum_{n=0}^{\infty} \frac{3t^{n+5} a_n}{5} \right) + \sum_{n=0}^{\infty} \left(-\frac{t^{n+6} a_n}{2} \right) = 0 \end{aligned} \tag{2}$$

The next step is to make all powers of t be n in each summation term. Going over each summation term above with power of t in it which is not already t^n and adjusting the power and the corresponding index gives

$$\begin{aligned} \sum_{n=2}^{\infty} 2n t^{n-1} a_n (n-1) &= \sum_{n=1}^{\infty} 2(1+n) a_{1+n} n t^n \\ \sum_{n=2}^{\infty} n(n-1) a_n t^{n-2} &= \sum_{n=0}^{\infty} (n+2) a_{n+2} (1+n) t^n \end{aligned}$$

$$\begin{aligned}
\sum_{n=1}^{\infty} 2na_n t^{n-1} &= \sum_{n=0}^{\infty} 2(1+n) a_{1+n} t^n \\
\sum_{n=0}^{\infty} 3t^{1+n} a_n &= \sum_{n=1}^{\infty} 3a_{n-1} t^n \\
\sum_{n=0}^{\infty} \left(-\frac{3t^{n+2} a_n}{2} \right) &= \sum_{n=2}^{\infty} \left(-\frac{3a_{n-2} t^n}{2} \right) \\
\sum_{n=0}^{\infty} t^{n+3} a_n &= \sum_{n=3}^{\infty} a_{n-3} t^n \\
\sum_{n=0}^{\infty} \left(-\frac{3t^{n+4} a_n}{4} \right) &= \sum_{n=4}^{\infty} \left(-\frac{3a_{n-4} t^n}{4} \right) \\
\sum_{n=0}^{\infty} \frac{3t^{n+5} a_n}{5} &= \sum_{n=5}^{\infty} \frac{3a_{n-5} t^n}{5} \\
\sum_{n=0}^{\infty} \left(-\frac{t^{n+6} a_n}{2} \right) &= \sum_{n=6}^{\infty} \left(-\frac{a_{n-6} t^n}{2} \right)
\end{aligned}$$

Substituting all the above in Eq (2) gives the following equation where now all powers of t are the same and equal to n .

$$\begin{aligned}
&\left(\sum_{n=2}^{\infty} t^n a_n n(n-1) \right) + \left(\sum_{n=1}^{\infty} 2(1+n) a_{1+n} n t^n \right) \\
&+ \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (1+n) t^n \right) + \left(\sum_{n=0}^{\infty} 2(1+n) a_{1+n} t^n \right) \\
&+ \left(\sum_{n=1}^{\infty} n a_n t^n \right) + \left(\sum_{n=1}^{\infty} 3a_{n-1} t^n \right) + \sum_{n=2}^{\infty} \left(-\frac{3a_{n-2} t^n}{2} \right) + \left(\sum_{n=3}^{\infty} a_{n-3} t^n \right) \\
&+ \sum_{n=4}^{\infty} \left(-\frac{3a_{n-4} t^n}{4} \right) + \left(\sum_{n=5}^{\infty} \frac{3a_{n-5} t^n}{5} \right) + \sum_{n=6}^{\infty} \left(-\frac{a_{n-6} t^n}{2} \right) = 0
\end{aligned} \tag{3}$$

$n = 0$ gives

$$2a_2 + 2a_1 = 0$$

$$a_2 = -a_1$$

$n = 1$ gives

$$8a_2 + 6a_3 + a_1 + 3a_0 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{a_0}{2} + \frac{7a_1}{6}$$

$n = 2$ gives

$$4a_2 + 18a_3 + 12a_4 + 3a_1 - \frac{3a_0}{2} = 0$$

Which after substituting earlier equations, simplifies to

$$a_4 = \frac{7a_0}{8} - \frac{5a_1}{3}$$

$n = 3$ gives

$$9a_3 + 32a_4 + 20a_5 + 3a_2 - \frac{3a_1}{2} + a_0 = 0$$

Which after substituting earlier equations, simplifies to

$$a_5 = -\frac{49a_0}{40} + \frac{71a_1}{30}$$

$n = 4$ gives

$$16a_4 + 50a_5 + 30a_6 + 3a_3 - \frac{3a_2}{2} + a_1 - \frac{3a_0}{4} = 0$$

Which after substituting earlier equations, simplifies to

$$a_6 = \frac{33a_0}{20} - \frac{293a_1}{90}$$

$n = 5$ gives

$$25a_5 + 72a_6 + 42a_7 + 3a_4 - \frac{3a_3}{2} + a_2 - \frac{3a_1}{4} + \frac{3a_0}{5} = 0$$

Which after substituting earlier equations, simplifies to

$$a_7 = -\frac{1843a_0}{840} + \frac{1378a_1}{315}$$

For $6 \leq n$, the recurrence equation is

$$na_n(n-1) + 2(1+n)a_{1+n}n + (n+2)a_{n+2}(1+n) + 2(1+n)a_{1+n} \quad (4)$$

$$+ na_n + 3a_{n-1} - \frac{3a_{n-2}}{2} + a_{n-3} - \frac{3a_{n-4}}{4} + \frac{3a_{n-5}}{5} - \frac{a_{n-6}}{2} = 0$$

Solving for a_{n+2} , gives

$$a_{n+2} = \frac{20n^2a_n + 40n^2a_{1+n} + 80na_{1+n} + 40a_{1+n} - 10a_{n-6} + 12a_{n-5} - 15a_{n-4} + 20a_{n-3} - 30a_{n-2} + 60a_{n-1}}{20(n+2)(1+n)}$$

$$(5) = -\frac{n^2a_n}{(n+2)(1+n)} - \frac{(40n^2 + 80n + 40)a_{1+n}}{20(n+2)(1+n)} + \frac{a_{n-6}}{2(n+2)(1+n)} - \frac{3a_{n-5}}{5(n+2)(1+n)}$$

$$+ \frac{3a_{n-4}}{4(n+2)(1+n)} - \frac{a_{n-3}}{(n+2)(1+n)} + \frac{3a_{n-2}}{2(n+2)(1+n)} - \frac{3a_{n-1}}{(n+2)(1+n)}$$

And so on. Therefore the solution is

$$y(t) = \sum_{n=0}^{\infty} a_n t^n$$

$$= a_3 t^3 + a_2 t^2 + a_1 t + a_0 + \dots$$

Substituting the values for a_n found above, the solution becomes

$$y(t) = a_0 + a_1 t - a_1 t^2 + \left(-\frac{a_0}{2} + \frac{7a_1}{6}\right) t^3 + \left(\frac{7a_0}{8} - \frac{5a_1}{3}\right) t^4 + \left(-\frac{49a_0}{40} + \frac{71a_1}{30}\right) t^5 + \dots$$

Collecting terms, the solution becomes

$$y(t) = \left(1 - \frac{1}{2}t^3 + \frac{7}{8}t^4 - \frac{49}{40}t^5\right) a_0 + \left(t - t^2 + \frac{7}{6}t^3 - \frac{5}{3}t^4 + \frac{71}{30}t^5\right) a_1 + O(t^6) \quad (3)$$

At $t = 0$ the solution above becomes

$$y(t) = \left(1 - \frac{1}{2}t^3 + \frac{7}{8}t^4 - \frac{49}{40}t^5\right) c_1 + \left(t - t^2 + \frac{7}{6}t^3 - \frac{5}{3}t^4 + \frac{71}{30}t^5\right) c_2 + O(t^6)$$

$$y(t) = -t^3 + 2 + \frac{7t^4}{4} - \frac{49t^5}{20} + O(t^6)$$

Replacing t in the above with the original independent variable x using $t = x - 1$ results in

$$y = -(x - 1)^3 + 2 + \frac{7(x - 1)^4}{4} - \frac{49(x - 1)^5}{20} + \frac{33(x - 1)^6}{10} + O((x - 1)^6)$$

Summary

The solution(s) found are the following

$$y = -(x - 1)^3 + 2 + \frac{7(x - 1)^4}{4} - \frac{49(x - 1)^5}{20} + \frac{33(x - 1)^6}{10} + O((x - 1)^6) \quad (1)$$

Verification of solutions

$$y = -(x - 1)^3 + 2 + \frac{7(x - 1)^4}{4} - \frac{49(x - 1)^5}{20} + \frac{33(x - 1)^6}{10} + O((x - 1)^6)$$

Verified OK.

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Heun: Equivalence to the GHE or one of its 4 confluent cases under a power @ Moebius
-> trying a solution of the form  $r_0(x) * Y + r_1(x) * Y$  where  $Y = \exp(\int(r(x), dx)) * 2F1([a$ 
-> Trying changes of variables to rationalize or make the ODE simpler
    trying a symmetry of the form [xi=0, eta=F(x)]
    checking if the LODE is missing y
    -> Heun: Equivalence to the GHE or one of its 4 confluent cases under a power @ Moebius
    -> trying a solution of the form  $r_0(x) * Y + r_1(x) * Y$  where  $Y = \exp(\int(r(x), dx)) * 2F1$ 
        trying a symmetry of the form [xi=0, eta=F(x)]
        trying 2nd order exact linear
        trying symmetries linear in x and y(x)
        trying to convert to a linear ODE with constant coefficients
<- unable to find a useful change of variables
    trying a symmetry of the form [xi=0, eta=F(x)]
    trying differential order: 2; exact nonlinear
    trying symmetries linear in x and y(x)
    trying to convert to a linear ODE with constant coefficients
    trying 2nd order, integrating factor of the form  $\mu(x,y)$ 
    trying a symmetry of the form [xi=0, eta=F(x)]
    checking if the LODE is missing y
    -> Heun: Equivalence to the GHE or one of its 4 confluent cases under a power @ Moebius
    -> trying a solution of the form  $r_0(x) * Y + r_1(x) * Y$  where  $Y = \exp(\int(r(x), dx)) * 2F1$ 
    -> Trying changes of variables to rationalize or make the ODE simpler
        trying a symmetry of the form [xi=0, eta=F(x)]
        checking if the LODE is missing y
        -> Heun: Equivalence to the GHE or one of its 4 confluent cases under a power @ Moebius
        -> trying a solution of the form  $r_0(x) * Y + r_1(x) * Y$  where  $Y = \exp(\int(r(x), dx)) * 2F1$ 
            trying a symmetry of the form [xi=0, eta=F(x)]
            trying 2nd order exact linear
            trying symmetries linear in x and y(x)
            trying to convert to a linear ODE with constant coefficients
<- unable to find a useful change of variables
    trying a symmetry of the form [xi=0, eta=F(x)]
trying to convert to an ODE of Bessel type
-> trying reduction of order to Riccati
    trying Riccati sub-methods:
        trying Riccati_symmetries
        -> trying a symmetry pattern of the form  $[F(x)*G(y), 0]$ 
        -> trying a symmetry pattern of the form  $[0, F(x)*G(y)]$ 
        -> trying a symmetry pattern of the form  $[F(x), G(x)*y+H(x)]$ 
--- Trying Lie symmetry methods, 2nd order ---
-> Computing symmetries using:  $\text{var} = 2; [0, -]$ 
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 16

Order:=6;

```
dsolve([x^2*diff(y(x),x$2)+(1+x)*diff(y(x),x)+3*ln(x)*y(x)=0,y(1) = 2, D(y)(1) = 0],y(x),typ
```

$$y(x) = 2 - (x - 1)^3 + \frac{7}{4}(x - 1)^4 - \frac{49}{20}(x - 1)^5 + O((x - 1)^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 30

```
AsymptoticDSolveValue[{x^2*y''[x]+(1+x)*y'[x]+3*Log[x]*y[x]==0,{y[1]==2,y'[1]==0}},y[x],{x,1
```

$$y(x) \rightarrow -\frac{49}{20}(x - 1)^5 + \frac{7}{4}(x - 1)^4 - (x - 1)^3 + 2$$

14.4 problem 4

14.4.1 Existence and uniqueness analysis 3363

Internal problem ID [736]

Internal file name [OUTPUT/736_Sunday_June_05_2022_01_48_18_AM_81415202/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 4.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$y'' + y'x^2 + \sin(x)y = 0$$

With initial conditions

$$[y(0) = a_0, y'(0) = a_1]$$

With the expansion point for the power series method at $x = 0$.

14.4.1 Existence and uniqueness analysis

This is a linear ODE. In canonical form it is written as

$$y'' + p(x)y' + q(x)y = F$$

Where here

$$p(x) = x^2$$

$$q(x) = \sin(x)$$

$$F = 0$$

Hence the ode is

$$y'' + y'x^2 + \sin(x)y = 0$$

The domain of $p(x) = x^2$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is inside this domain. The domain of $q(x) = \sin(x)$ is

$$\{-\infty < x < \infty\}$$

And the point $x_0 = 0$ is also inside this domain. Hence solution exists and is unique.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \dots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (819)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (820)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = -y'x^2 - \sin(x)y$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dx} \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\ &= (x^4 - 2x - \sin(x))y' + y(\sin(x)x^2 - \cos(x)) \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dx} \\ &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\ &= (-x^6 + 6x^3 + 2\sin(x)x^2 - 2\cos(x) - 2)y' - y(-\sin(x)^2 + (x^4 - 4x - 1)\sin(x) - x^2\cos(x)) \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dx} \\ &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\ &= (\sin(x)^2 + (-3x^4 + 8x + 3)\sin(x) + x^2(x^6 - 12x^3 + 5\cos(x) + 20))y' + (-2\sin(x)^2x^2 + (x^6 - 1 \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dx} \\ &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\ &= (-3\sin(x)^2x^2 + (4x^6 - 30x^3 - 9x^2 + 6\cos(x) + 14)\sin(x) + (-9x^4 + 24x + 4)\cos(x) - x^{10} + 20 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = a_0$ and $y'(0) = a_1$ gives

$$\begin{aligned} F_0 &= 0 \\ F_1 &= -a_0 \\ F_2 &= -4a_1 \\ F_3 &= a_0 \\ F_4 &= 4a_1 + 16a_0 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = xa_1 + a_0 - \frac{a_0x^3}{6} - \frac{a_1x^4}{6} + \frac{a_0x^5}{120} + \frac{x^6a_0}{45} + \frac{x^6a_1}{180} + O(x^6)$$

$$y = xa_1 + a_0 - \frac{a_0x^3}{6} - \frac{a_1x^4}{6} + \frac{a_0x^5}{120} + \frac{x^6a_0}{45} + \frac{x^6a_1}{180} + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = - \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x^2 - \sin(x) \left(\sum_{n=0}^{\infty} a_n x^n \right) \quad (1)$$

Expanding $\sin(x)$ as Taylor series around $x = 0$ and keeping only the first 6 terms gives

$$\begin{aligned} \sin(x) &= x - \frac{1}{6}x^3 + \frac{1}{120}x^5 - \frac{1}{5040}x^7 + \dots \\ &= x - \frac{1}{6}x^3 + \frac{1}{120}x^5 - \frac{1}{5040}x^7 \end{aligned}$$

Hence the ODE in Eq (1) becomes

$$\begin{aligned} &\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x^2 \\ &+ \left(x - \frac{1}{6}x^3 + \frac{1}{120}x^5 - \frac{1}{5040}x^7 \right) \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \end{aligned}$$

Expanding the third term in (1) gives

$$\begin{aligned} &\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x^2 + x \cdot \left(\sum_{n=0}^{\infty} a_n x^n \right) \\ &- \frac{x^3}{6} \cdot \left(\sum_{n=0}^{\infty} a_n x^n \right) + \frac{x^5}{120} \cdot \left(\sum_{n=0}^{\infty} a_n x^n \right) - \frac{x^7}{5040} \cdot \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \end{aligned}$$

Which simplifies to

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n x^{1+n} a_n \right) + \left(\sum_{n=0}^{\infty} x^{1+n} a_n \right) \\ & + \sum_{n=0}^{\infty} \left(-\frac{x^{n+3} a_n}{6} \right) + \left(\sum_{n=0}^{\infty} \frac{x^{n+5} a_n}{120} \right) + \sum_{n=0}^{\infty} \left(-\frac{x^{n+7} a_n}{5040} \right) = 0 \end{aligned} \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\begin{aligned} \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} &= \sum_{n=0}^{\infty} (n+2) a_{n+2} (1+n) x^n \\ \sum_{n=1}^{\infty} n x^{1+n} a_n &= \sum_{n=2}^{\infty} (n-1) a_{n-1} x^n \\ \sum_{n=0}^{\infty} x^{1+n} a_n &= \sum_{n=1}^{\infty} a_{n-1} x^n \\ \sum_{n=0}^{\infty} \left(-\frac{x^{n+3} a_n}{6} \right) &= \sum_{n=3}^{\infty} \left(-\frac{a_{n-3} x^n}{6} \right) \\ \sum_{n=0}^{\infty} \frac{x^{n+5} a_n}{120} &= \sum_{n=5}^{\infty} \frac{a_{n-5} x^n}{120} \\ \sum_{n=0}^{\infty} \left(-\frac{x^{n+7} a_n}{5040} \right) &= \sum_{n=7}^{\infty} \left(-\frac{a_{n-7} x^n}{5040} \right) \end{aligned}$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} & \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (1+n) x^n \right) + \left(\sum_{n=2}^{\infty} (n-1) a_{n-1} x^n \right) + \left(\sum_{n=1}^{\infty} a_{n-1} x^n \right) \\ & + \sum_{n=3}^{\infty} \left(-\frac{a_{n-3} x^n}{6} \right) + \left(\sum_{n=5}^{\infty} \frac{a_{n-5} x^n}{120} \right) + \sum_{n=7}^{\infty} \left(-\frac{a_{n-7} x^n}{5040} \right) = 0 \end{aligned} \quad (3)$$

$n = 1$ gives

$$6a_3 + a_0 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{a_0}{6}$$

$n = 2$ gives

$$12a_4 + 2a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_4 = -\frac{a_1}{6}$$

$n = 3$ gives

$$20a_5 + 3a_2 - \frac{a_0}{6} = 0$$

Which after substituting earlier equations, simplifies to

$$a_5 = \frac{a_0}{120}$$

$n = 4$ gives

$$30a_6 + 4a_3 - \frac{a_1}{6} = 0$$

Which after substituting earlier equations, simplifies to

$$a_6 = \frac{a_0}{45} + \frac{a_1}{180}$$

$n = 5$ gives

$$42a_7 + 5a_4 - \frac{a_2}{6} + \frac{a_0}{120} = 0$$

Which after substituting earlier equations, simplifies to

$$a_7 = -\frac{a_0}{5040} + \frac{5a_1}{252}$$

For $7 \leq n$, the recurrence equation is

$$(n+2)a_{n+2}(1+n) + (n-1)a_{n-1} + a_{n-1} - \frac{a_{n-3}}{6} + \frac{a_{n-5}}{120} - \frac{a_{n-7}}{5040} = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} a_{n+2} &= -\frac{5040na_{n-1} - a_{n-7} + 42a_{n-5} - 840a_{n-3}}{5040(n+2)(1+n)} \\ (5) \quad &= \frac{a_{n-7}}{5040(n+2)(1+n)} - \frac{a_{n-5}}{120(n+2)(1+n)} + \frac{a_{n-3}}{6(n+2)(1+n)} - \frac{na_{n-1}}{(n+2)(1+n)} \end{aligned}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - \frac{1}{6}a_0 x^3 - \frac{1}{6}a_1 x^4 + \frac{1}{120}a_0 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - \frac{1}{6}x^3 + \frac{1}{120}x^5\right) a_0 + \left(x - \frac{1}{6}x^4\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - \frac{1}{6}x^3 + \frac{1}{120}x^5\right) c_1 + \left(x - \frac{1}{6}x^4\right) c_2 + O(x^6)$$

$$y = a_0 - \frac{a_0 x^3}{6} + \frac{a_0 x^5}{120} + x a_1 - \frac{a_1 x^4}{6} + O(x^6)$$

Summary

The solution(s) found are the following

$$y = x a_1 + a_0 - \frac{a_0 x^3}{6} - \frac{a_1 x^4}{6} + \frac{a_0 x^5}{120} + \frac{x^6 a_0}{45} + \frac{x^6 a_1}{180} + O(x^6) \quad (1)$$

$$y = a_0 - \frac{a_0 x^3}{6} + \frac{a_0 x^5}{120} + x a_1 - \frac{a_1 x^4}{6} + O(x^6) \quad (2)$$

Verification of solutions

$$y = xa_1 + a_0 - \frac{a_0x^3}{6} - \frac{a_1x^4}{6} + \frac{a_0x^5}{120} + \frac{x^6a_0}{45} + \frac{x^6a_1}{180} + O(x^6)$$

Verified OK.

$$y = a_0 - \frac{a_0x^3}{6} + \frac{a_0x^5}{120} + xa_1 - \frac{a_1x^4}{6} + O(x^6)$$

Verified OK.

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Heun: Equivalence to the GHE or one of its 4 confluent cases under a power @ Moebius
-> trying a solution of the form r0(x) * Y + r1(x) * Y where Y = exp(int(r(x), dx)) * 2F1([a
-> Trying changes of variables to rationalize or make the ODE simpler
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Heun: Equivalence to the GHE or one of its 4 confluent cases under a power @ Moebius
-> trying a solution of the form r0(x) * Y + r1(x) * Y where Y = exp(int(r(x), dx)) * 2F1
trying a symmetry of the form [xi=0, eta=F(x)]
trying 2nd order exact linear
trying symmetries linear in x and y(x)
trying to convert to a linear ODE with constant coefficients
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Heun: Equivalence to the GHE or one of its 4 confluent cases under a power @ Moebius
-> trying a solution of the form r0(x) * Y + r1(x) * Y where Y = exp(int(r(x), dx)) * 2F1
trying a symmetry of the form [xi=0, eta=F(x)]
trying 2nd order exact linear
trying symmetries linear in x and y(x)
trying to convert to a linear ODE with constant coefficients
-> trying with_periodic_functions in the coefficients
--- Trying Lie symmetry methods, 2nd order ---
`, `-> Computing symmetries using: way = 5`[0, u]
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 24

Order:=6;

```
dsolve([diff(y(x),x$2)+x^2*diff(y(x),x)+sin(x)*y(x)=0,y(0) = a__0, D(y)(0) = a__1],y(x),type
```

$$y(x) = a_0 + a_1x - \frac{1}{6}a_0x^3 - \frac{1}{6}a_1x^4 + \frac{1}{120}a_0x^5 + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 32

```
AsymptoticDSolveValue[{y'[x]+x^2*y'[x]+Sin[x]*y[x]==0,{y[0]==a0,y'[0]==a1}},y[x],{x,0,5}]
```

$$y(x) \rightarrow \frac{a_0x^5}{120} - \frac{a_0x^3}{6} + a_0 - \frac{a_1x^4}{6} + a_1x$$

14.5 problem 5. case $x_0 = 0$

14.5.1 Maple step by step solution 3381

Internal problem ID [737]

Internal file name [OUTPUT/737_Sunday_June_05_2022_01_48_21_AM_72114671/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 5. case $x_0 = 0$.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second_order_airy", "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$y'' + 4y' + 6yx = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (822)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (823)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -4y' - 6yx \\
 F_1 &= \frac{dF_0}{dx} \\
 &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \\
 &= (-6x + 16)y' + (24x - 6)y \\
 F_2 &= \frac{dF_1}{dx} \\
 &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y}y' + \frac{\partial F_1}{\partial y'}F_1 \\
 &= (48x - 76)y' + 36\left(x^2 - \frac{8}{3}x + \frac{2}{3}\right)y \\
 F_3 &= \frac{dF_2}{dx} \\
 &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y}y' + \frac{\partial F_2}{\partial y'}F_2 \\
 &= (36x^2 - 288x + 376)y' - 288\left(x^2 - \frac{11}{6}x + \frac{1}{3}\right)y \\
 F_4 &= \frac{dF_3}{dx} \\
 &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y}y' + \frac{\partial F_3}{\partial y'}F_3 \\
 &= (-432x^2 + 1752x - 1888)y' - 216\left(x^3 - 8x^2 + \frac{118}{9}x - \frac{22}{9}\right)y
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -4y'(0) \\
 F_1 &= 16y'(0) - 6y(0) \\
 F_2 &= -76y'(0) + 24y(0) \\
 F_3 &= 376y'(0) - 96y(0) \\
 F_4 &= -1888y'(0) + 528y(0)
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$\begin{aligned}
 y &= \left(1 - x^3 + x^4 - \frac{4}{5}x^5 + \frac{11}{15}x^6\right)y(0) \\
 &\quad + \left(x - 2x^2 + \frac{8}{3}x^3 - \frac{19}{6}x^4 + \frac{47}{15}x^5 - \frac{118}{45}x^6\right)y'(0) + O(x^6)
 \end{aligned}$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = -4 \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) - 6 \left(\sum_{n=0}^{\infty} a_n x^n \right) x \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} 4n a_n x^{n-1} \right) + \left(\sum_{n=0}^{\infty} 6x^{1+n} a_n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (1+n) x^n$$

$$\sum_{n=1}^{\infty} 4n a_n x^{n-1} = \sum_{n=0}^{\infty} 4(1+n) a_{1+n} x^n$$

$$\sum_{n=0}^{\infty} 6x^{1+n} a_n = \sum_{n=1}^{\infty} 6a_{n-1} x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (1+n) x^n \right) + \left(\sum_{n=0}^{\infty} 4(1+n) a_{1+n} x^n \right) + \left(\sum_{n=1}^{\infty} 6a_{n-1} x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 + 4a_1 = 0$$

$$a_2 = -2a_1$$

For $1 \leq n$, the recurrence equation is

$$(n + 2) a_{n+2}(1 + n) + 4(1 + n) a_{1+n} + 6a_{n-1} = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} a_{n+2} &= -\frac{2(2na_{1+n} + 2a_{1+n} + 3a_{n-1})}{(n + 2)(1 + n)} \\ (5) \quad &= -\frac{2(2n + 2) a_{1+n}}{(n + 2)(1 + n)} - \frac{6a_{n-1}}{(n + 2)(1 + n)} \end{aligned}$$

For $n = 1$ the recurrence equation gives

$$6a_3 + 8a_2 + 6a_0 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = \frac{8a_1}{3} - a_0$$

For $n = 2$ the recurrence equation gives

$$12a_4 + 12a_3 + 6a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = -\frac{19a_1}{6} + a_0$$

For $n = 3$ the recurrence equation gives

$$20a_5 + 16a_4 + 6a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{47a_1}{15} - \frac{4a_0}{5}$$

For $n = 4$ the recurrence equation gives

$$30a_6 + 20a_5 + 6a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{118a_1}{45} + \frac{11a_0}{15}$$

For $n = 5$ the recurrence equation gives

$$42a_7 + 24a_6 + 6a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{1229a_1}{630} - \frac{59a_0}{105}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1 x - 2a_1 x^2 + \left(\frac{8a_1}{3} - a_0\right) x^3 + \left(-\frac{19a_1}{6} + a_0\right) x^4 + \left(\frac{47a_1}{15} - \frac{4a_0}{5}\right) x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - x^3 + x^4 - \frac{4}{5}x^5\right) a_0 + \left(x - 2x^2 + \frac{8}{3}x^3 - \frac{19}{6}x^4 + \frac{47}{15}x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - x^3 + x^4 - \frac{4}{5}x^5\right) c_1 + \left(x - 2x^2 + \frac{8}{3}x^3 - \frac{19}{6}x^4 + \frac{47}{15}x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - x^3 + x^4 - \frac{4}{5}x^5 + \frac{11}{15}x^6\right) y(0) + \left(x - 2x^2 + \frac{8}{3}x^3 - \frac{19}{6}x^4 + \frac{47}{15}x^5 - \frac{118}{45}x^6\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 - x^3 + x^4 - \frac{4}{5}x^5\right) c_1 + \left(x - 2x^2 + \frac{8}{3}x^3 - \frac{19}{6}x^4 + \frac{47}{15}x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 - x^3 + x^4 - \frac{4}{5}x^5 + \frac{11}{15}x^6\right) y(0) + \left(x - 2x^2 + \frac{8}{3}x^3 - \frac{19}{6}x^4 + \frac{47}{15}x^5 - \frac{118}{45}x^6\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 - x^3 + x^4 - \frac{4}{5}x^5\right) c_1 + \left(x - 2x^2 + \frac{8}{3}x^3 - \frac{19}{6}x^4 + \frac{47}{15}x^5\right) c_2 + O(x^6)$$

Verified OK.

14.5.1 Maple step by step solution

Let's solve

$$y'' = -4y' - 6yx$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + 4y' + 6yx = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite ODE with series expansions

- Convert $x \cdot y$ to series expansion

$$x \cdot y = \sum_{k=0}^{\infty} a_k x^{k+1}$$

- Shift index using $k \rightarrow k - 1$

$$x \cdot y = \sum_{k=1}^{\infty} a_{k-1} x^k$$

- Convert y' to series expansion

$$y' = \sum_{k=1}^{\infty} a_k k x^{k-1}$$

- Shift index using $k \rightarrow k + 1$

$$y' = \sum_{k=0}^{\infty} a_{k+1} (k+1) x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k \rightarrow k + 2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2} (k+2)(k+1) x^k$$

Rewrite ODE with series expansions

$$2a_2 + 4a_1 + \left(\sum_{k=1}^{\infty} (a_{k+2}(k+2)(k+1) + 4a_{k+1}(k+1) + 6a_{k-1}) x^k \right) = 0$$

- Each term must be 0

$$2a_2 + 4a_1 = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(k^2 + 3k + 2) a_{k+2} + 4a_{k+1}k + 6a_{k-1} + 4a_{k+1} = 0$$

- Shift index using $k \rightarrow k + 1$

$$((k+1)^2 + 3k + 5) a_{k+3} + 4a_{k+2}(k+1) + 6a_k + 4a_{k+2} = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+3} = -\frac{2(2ka_{k+2} + 3a_k + 4a_{k+2})}{k^2 + 5k + 6}, 2a_2 + 4a_1 = 0 \right]$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacic's algorithm  
<- No Liouvillian solutions exist  
-> Trying a solution in terms of special functions:  
  -> Bessel  
  <- Bessel successful  
<- special function solution successful`
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 47

```
Order:=6;  
dsolve(diff(y(x),x$2)+4*diff(y(x),x)+6*x*y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 - x^3 + x^4 - \frac{4}{5}x^5\right) y(0) + \left(x - 2x^2 + \frac{8}{3}x^3 - \frac{19}{6}x^4 + \frac{47}{15}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 55

```
AsymptoticDSolveValue[y''[x]+4*y'[x]+6*x*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_1 \left(-\frac{4x^5}{5} + x^4 - x^3 + 1 \right) + c_2 \left(\frac{47x^5}{15} - \frac{19x^4}{6} + \frac{8x^3}{3} - 2x^2 + x \right)$$

14.6 problem 5. case $x_0 = 4$

14.6.1 Maple step by step solution 3392

Internal problem ID [738]

Internal file name [OUTPUT/738_Sunday_June_05_2022_01_48_23_AM_64748739/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 5. case $x_0 = 4$.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$y'' + 4y' + 6yx = 0$$

With the expansion point for the power series method at $x = 4$.

The ode does not have its expansion point at $x = 0$, therefore to simplify the computation of power series expansion, change of variable is made on the independent variable to shift the initial conditions and the expansion point back to zero. The new ode is then solved more easily since the expansion point is now at zero. The solution converted back to the original independent variable. Let

$$t = x - 4$$

The ode is converted to be in terms of the new independent variable t . This results in

$$\frac{d^2}{dt^2}y(t) + 4\frac{d}{dt}y(t) + 6y(t)(t + 4) = 0$$

With its expansion point and initial conditions now at $t = 0$. The transformed ODE is now solved. Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (825)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (826)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -6y(t)t - 24y(t) - 4\frac{d}{dt}y(t) \\
 F_1 &= \frac{dF_0}{dt} \\
 &= \frac{\partial F_0}{\partial t} + \frac{\partial F_0}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_0}{\partial \frac{d}{dt}y(t)} F_0 \\
 &= (-6t - 8) \left(\frac{d}{dt}y(t) \right) + (24t + 90) y(t) \\
 F_2 &= \frac{dF_1}{dt} \\
 &= \frac{\partial F_1}{\partial t} + \frac{\partial F_1}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_1}{\partial \frac{d}{dt}y(t)} F_1 \\
 &= (48t + 116) \left(\frac{d}{dt}y(t) \right) + 36y(t) \left(t^2 + \frac{16}{3}t + 6 \right) \\
 F_3 &= \frac{dF_2}{dt} \\
 &= \frac{\partial F_2}{\partial t} + \frac{\partial F_2}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_2}{\partial \frac{d}{dt}y(t)} F_2 \\
 &= (36t^2 - 200) \left(\frac{d}{dt}y(t) \right) - 288 \left(t^2 + \frac{37}{6}t + 9 \right) y(t) \\
 F_4 &= \frac{dF_3}{dt} \\
 &= \frac{\partial F_3}{\partial t} + \frac{\partial F_3}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_3}{\partial \frac{d}{dt}y(t)} F_3 \\
 &= (-432t^2 - 1704t - 1792) \left(\frac{d}{dt}y(t) \right) - 216 \left(t^3 + 4t^2 - \frac{26}{9}t - 14 \right) y(t)
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $t = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -24y(0) - 4y'(0) \\
 F_1 &= -8y'(0) + 90y(0) \\
 F_2 &= 116y'(0) + 216y(0) \\
 F_3 &= -200y'(0) - 2592y(0) \\
 F_4 &= -1792y'(0) + 3024y(0)
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y(t) = \left(1 - 12t^2 + 15t^3 + 9t^4 - \frac{108}{5}t^5 + \frac{21}{5}t^6\right) y(0) + \left(t - 2t^2 - \frac{4}{3}t^3 + \frac{29}{6}t^4 - \frac{5}{3}t^5 - \frac{112}{45}t^6\right) y'(0) + O(t^6)$$

Since the expansion point $t = 0$ is an ordinary, we can also solve this using standard power series Let the solution be represented as power series of the form

$$y(t) = \sum_{n=0}^{\infty} a_n t^n$$

Then

$$\frac{d}{dt}y(t) = \sum_{n=1}^{\infty} n a_n t^{n-1}$$

$$\frac{d^2}{dt^2}y(t) = \sum_{n=2}^{\infty} n(n-1) a_n t^{n-2}$$

Substituting the above back into the ode gives

$$\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2} = -6 \left(\sum_{n=0}^{\infty} a_n t^n \right) t - 24 \left(\sum_{n=0}^{\infty} a_n t^n \right) - 4 \left(\sum_{n=1}^{\infty} n a_n t^{n-1} \right) \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2} \right) + \left(\sum_{n=1}^{\infty} 4n a_n t^{n-1} \right) + \left(\sum_{n=0}^{\infty} 6t^{1+n} a_n \right) + \left(\sum_{n=0}^{\infty} 24a_n t^n \right) = 0 \quad (2)$$

The next step is to make all powers of t be n in each summation term. Going over each summation term above with power of t in it which is not already t^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (1+n) t^n$$

$$\sum_{n=1}^{\infty} 4n a_n t^{n-1} = \sum_{n=0}^{\infty} 4(1+n) a_{1+n} t^n$$

$$\sum_{n=0}^{\infty} 6t^{1+n} a_n = \sum_{n=1}^{\infty} 6a_{n-1} t^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of t are the same and equal to n .

$$\begin{aligned} & \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (1+n) t^n \right) + \left(\sum_{n=0}^{\infty} 4(1+n) a_{1+n} t^n \right) \\ & + \left(\sum_{n=1}^{\infty} 6a_{n-1} t^n \right) + \left(\sum_{n=0}^{\infty} 24a_n t^n \right) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$2a_2 + 4a_1 + 24a_0 = 0$$

$$a_2 = -12a_0 - 2a_1$$

For $1 \leq n$, the recurrence equation is

$$(n+2) a_{n+2} (1+n) + 4(1+n) a_{1+n} + 6a_{n-1} + 24a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} (5) \quad a_{n+2} &= -\frac{2(2na_{1+n} + 12a_n + 2a_{1+n} + 3a_{n-1})}{(n+2)(1+n)} \\ &= -\frac{24a_n}{(n+2)(1+n)} - \frac{2(2n+2)a_{1+n}}{(n+2)(1+n)} - \frac{6a_{n-1}}{(n+2)(1+n)} \end{aligned}$$

For $n = 1$ the recurrence equation gives

$$6a_3 + 8a_2 + 6a_0 + 24a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = 15a_0 - \frac{4a_1}{3}$$

For $n = 2$ the recurrence equation gives

$$12a_4 + 12a_3 + 6a_1 + 24a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = 9a_0 + \frac{29a_1}{6}$$

For $n = 3$ the recurrence equation gives

$$20a_5 + 16a_4 + 6a_2 + 24a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = -\frac{108a_0}{5} - \frac{5a_1}{3}$$

For $n = 4$ the recurrence equation gives

$$30a_6 + 20a_5 + 6a_3 + 24a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{21a_0}{5} - \frac{112a_1}{45}$$

For $n = 5$ the recurrence equation gives

$$42a_7 + 24a_6 + 6a_4 + 24a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{303a_0}{35} + \frac{1061a_1}{630}$$

And so on. Therefore the solution is

$$\begin{aligned} y(t) &= \sum_{n=0}^{\infty} a_n t^n \\ &= a_3 t^3 + a_2 t^2 + a_1 t + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$\begin{aligned} y(t) &= a_0 + a_1 t + (-12a_0 - 2a_1) t^2 + \left(15a_0 - \frac{4a_1}{3}\right) t^3 \\ &\quad + \left(9a_0 + \frac{29a_1}{6}\right) t^4 + \left(-\frac{108a_0}{5} - \frac{5a_1}{3}\right) t^5 + \dots \end{aligned}$$

Collecting terms, the solution becomes

$$y(t) = \left(1 - 12t^2 + 15t^3 + 9t^4 - \frac{108}{5}t^5\right) a_0 + \left(t - 2t^2 - \frac{4}{3}t^3 + \frac{29}{6}t^4 - \frac{5}{3}t^5\right) a_1 + O(t^6) \quad (3)$$

At $t = 0$ the solution above becomes

$$y(t) = \left(1 - 12t^2 + 15t^3 + 9t^4 - \frac{108}{5}t^5\right) c_1 + \left(t - 2t^2 - \frac{4}{3}t^3 + \frac{29}{6}t^4 - \frac{5}{3}t^5\right) c_2 + O(t^6)$$

Replacing t in the above with the original independent variable x using $t = x - 4$ results in

$$\begin{aligned} y &= \left(1 - 12(x-4)^2 + 15(x-4)^3 + 9(x-4)^4 - \frac{108(x-4)^5}{5} + \frac{21(x-4)^6}{5}\right) y(4) \\ &+ \left(x-4 - 2(x-4)^2 - \frac{4(x-4)^3}{3} + \frac{29(x-4)^4}{6} - \frac{5(x-4)^5}{3} - \frac{112(x-4)^6}{45}\right) y'(4) \\ &+ O((x-4)^6) \end{aligned}$$

Summary

The solution(s) found are the following

$$\begin{aligned} y &= \left(1 - 12(x-4)^2 + 15(x-4)^3 + 9(x-4)^4 - \frac{108(x-4)^5}{5} + \frac{21(x-4)^6}{5}\right) y(4) \\ &+ \left(x-4 - 2(x-4)^2 - \frac{4(x-4)^3}{3} + \frac{29(x-4)^4}{6} - \frac{5(x-4)^5}{3} - \frac{112(x-4)^6}{45}\right) y'(4) \\ &+ O((x-4)^6) \end{aligned} \quad (1)$$

Verification of solutions

$$\begin{aligned} y &= \left(1 - 12(x-4)^2 + 15(x-4)^3 + 9(x-4)^4 - \frac{108(x-4)^5}{5} + \frac{21(x-4)^6}{5}\right) y(4) \\ &+ \left(x-4 - 2(x-4)^2 - \frac{4(x-4)^3}{3} + \frac{29(x-4)^4}{6} - \frac{5(x-4)^5}{3} - \frac{112(x-4)^6}{45}\right) y'(4) \\ &+ O((x-4)^6) \end{aligned}$$

Verified OK.

14.6.1 Maple step by step solution

Let's solve

$$y'' + 4y' + 6yx = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^k$$

- Rewrite ODE with series expansions

- Convert $x \cdot y$ to series expansion

$$x \cdot y = \sum_{k=0}^{\infty} a_k x^{k+1}$$

- Shift index using $k \rightarrow k - 1$

$$x \cdot y = \sum_{k=1}^{\infty} a_{k-1} x^k$$

- Convert y' to series expansion

$$y' = \sum_{k=1}^{\infty} a_k k x^{k-1}$$

- Shift index using $k \rightarrow k + 1$

$$y' = \sum_{k=0}^{\infty} a_{k+1} (k+1) x^k$$

- Convert y'' to series expansion

$$y'' = \sum_{k=2}^{\infty} a_k k(k-1) x^{k-2}$$

- Shift index using $k \rightarrow k + 2$

$$y'' = \sum_{k=0}^{\infty} a_{k+2} (k+2)(k+1) x^k$$

Rewrite ODE with series expansions

$$2a_2 + 4a_1 + \left(\sum_{k=1}^{\infty} (a_{k+2}(k+2)(k+1) + 4a_{k+1}(k+1) + 6a_{k-1}) x^k \right) = 0$$

- Each term must be 0

$$2a_2 + 4a_1 = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(k^2 + 3k + 2) a_{k+2} + 4a_{k+1}k + 6a_{k-1} + 4a_{k+1} = 0$$

- Shift index using $k \rightarrow k + 1$

$$((k + 1)^2 + 3k + 5) a_{k+3} + 4a_{k+2}(k + 1) + 6a_k + 4a_{k+2} = 0$$

- Recursion relation that defines the series solution to the ODE

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+3} = -\frac{2(2ka_{k+2} + 3a_k + 4a_{k+2})}{k^2 + 5k + 6}, 2a_2 + 4a_1 = 0 \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
    -> Bessel
    <- Bessel successful
<- special function solution successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 54

```

Order:=6;
dsolve(diff(y(x),x$2)+4*diff(y(x),x)+6*x*y(x)=0,y(x),type='series',x=4);

```

$$y(x) = \left(1 - 12(x - 4)^2 + 15(x - 4)^3 + 9(x - 4)^4 - \frac{108(x - 4)^5}{5} \right) y(4) + \left(x - 4 - 2(x - 4)^2 - \frac{4(x - 4)^3}{3} + \frac{29(x - 4)^4}{6} - \frac{5(x - 4)^5}{3} \right) D(y)(4) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 79

```
AsymptoticDSolveValue[y''[x]+4*y'[x]+6*x*y[x]==0,y[x],{x,4,5}]
```

$$y(x) \rightarrow c_1 \left(-\frac{108}{5}(x-4)^5 + 9(x-4)^4 + 15(x-4)^3 - 12(x-4)^2 + 1 \right) \\ + c_2 \left(-\frac{5}{3}(x-4)^5 + \frac{29}{6}(x-4)^4 - \frac{4}{3}(x-4)^3 - 2(x-4)^2 + x - 4 \right)$$

14.7 problem 6. case $x_0 = 0$

14.7.1 Maple step by step solution 3403

Internal problem ID [739]

Internal file name [OUTPUT/739_Sunday_June_05_2022_01_48_24_AM_30076032/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 6. case $x_0 = 0$.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(x^2 - 2x - 3)y'' + y'x + 4y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (828)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (829)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = -\frac{y'x + 4y}{x^2 - 2x - 3}$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dx} \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\ &= \frac{(-2x^2 + 8x + 15)y' + (12x - 8)y}{(x^2 - 2x - 3)^2} \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dx} \\ &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\ &= \frac{(18x^3 - 64x^2 - 67x + 60)y' - 28y(x^2 - \frac{6}{7}x + \frac{32}{7})}{(x^2 - 2x - 3)^3} \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dx} \\ &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\ &= \frac{(-100x^4 + 400x^3 + 20x^2 - 120x + 945)y' + 40(x^3 + 2x^2 + \frac{65}{2}x - 27)y}{(x^2 - 2x - 3)^4} \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dx} \\ &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\ &= \frac{(540x^5 - 2400x^4 + 2880x^3 - 6480x^2 - 11085x + 11160)y' + 200y(x^4 - 10x^3 - \frac{461}{10}x^2 + \frac{411}{5}x - \frac{408}{5})}{(x^2 - 2x - 3)^5} \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned} F_0 &= \frac{4y(0)}{3} \\ F_1 &= -\frac{8y(0)}{9} + \frac{5y'(0)}{3} \\ F_2 &= \frac{128y(0)}{27} - \frac{20y'(0)}{9} \\ F_3 &= -\frac{40y(0)}{3} + \frac{35y'(0)}{3} \\ F_4 &= \frac{5440y(0)}{81} - \frac{1240y'(0)}{27} \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 + \frac{2}{3}x^2 - \frac{4}{27}x^3 + \frac{16}{81}x^4 - \frac{1}{9}x^5 + \frac{68}{729}x^6\right) y(0) \\ + \left(x + \frac{5}{18}x^3 - \frac{5}{54}x^4 + \frac{7}{72}x^5 - \frac{31}{486}x^6\right) y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series The ode is normalized to be

$$(x^2 - 2x - 3) y'' + y'x + 4y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1} \\ y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(x^2 - 2x - 3) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + 4 \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} x^n a_n n(n-1) \right) + \sum_{n=2}^{\infty} (-2n x^{n-1} a_n (n-1)) \\ + \sum_{n=2}^{\infty} (-3n(n-1) a_n x^{n-2}) + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \left(\sum_{n=0}^{\infty} 4a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the

power and the corresponding index gives

$$\begin{aligned}\sum_{n=2}^{\infty} (-2n x^{n-1} a_n (n-1)) &= \sum_{n=1}^{\infty} (-2(n+1) a_{n+1} n x^n) \\ \sum_{n=2}^{\infty} (-3n(n-1) a_n x^{n-2}) &= \sum_{n=0}^{\infty} (-3(n+2) a_{n+2} (n+1) x^n)\end{aligned}$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned}\left(\sum_{n=2}^{\infty} x^n a_n n(n-1) \right) + \sum_{n=1}^{\infty} (-2(n+1) a_{n+1} n x^n) \\ + \sum_{n=0}^{\infty} (-3(n+2) a_{n+2} (n+1) x^n) + \left(\sum_{n=1}^{\infty} n a_n x^n \right) + \left(\sum_{n=0}^{\infty} 4a_n x^n \right) = 0\end{aligned}\tag{3}$$

$n = 0$ gives

$$-6a_2 + 4a_0 = 0$$

$$a_2 = \frac{2a_0}{3}$$

$n = 1$ gives

$$-4a_2 - 18a_3 + 5a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{4a_0}{27} + \frac{5a_1}{18}$$

For $2 \leq n$, the recurrence equation is

$$na_n(n-1) - 2(n+1)a_{n+1}n - 3(n+2)a_{n+2}(n+1) + na_n + 4a_n = 0\tag{4}$$

Solving for a_{n+2} , gives

$$\begin{aligned}a_{n+2} &= \frac{n^2 a_n - 2n^2 a_{n+1} - 2n a_{n+1} + 4a_n}{3(n+2)(n+1)} \\ (5) \quad &= \frac{(n^2 + 4) a_n}{3(n+2)(n+1)} + \frac{(-2n^2 - 2n) a_{n+1}}{3(n+2)(n+1)}\end{aligned}$$

For $n = 2$ the recurrence equation gives

$$8a_2 - 12a_3 - 36a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{16a_0}{81} - \frac{5a_1}{54}$$

For $n = 3$ the recurrence equation gives

$$13a_3 - 24a_4 - 60a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = -\frac{a_0}{9} + \frac{7a_1}{72}$$

For $n = 4$ the recurrence equation gives

$$20a_4 - 40a_5 - 90a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{68a_0}{729} - \frac{31a_1}{486}$$

For $n = 5$ the recurrence equation gives

$$29a_5 - 60a_6 - 126a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{2143a_0}{30618} + \frac{4307a_1}{81648}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1x + \frac{2a_0x^2}{3} + \left(-\frac{4a_0}{27} + \frac{5a_1}{18}\right)x^3 + \left(\frac{16a_0}{81} - \frac{5a_1}{54}\right)x^4 + \left(-\frac{a_0}{9} + \frac{7a_1}{72}\right)x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + \frac{2}{3}x^2 - \frac{4}{27}x^3 + \frac{16}{81}x^4 - \frac{1}{9}x^5\right)a_0 + \left(x + \frac{5}{18}x^3 - \frac{5}{54}x^4 + \frac{7}{72}x^5\right)a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 + \frac{2}{3}x^2 - \frac{4}{27}x^3 + \frac{16}{81}x^4 - \frac{1}{9}x^5\right)c_1 + \left(x + \frac{5}{18}x^3 - \frac{5}{54}x^4 + \frac{7}{72}x^5\right)c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$\begin{aligned} y &= \left(1 + \frac{2}{3}x^2 - \frac{4}{27}x^3 + \frac{16}{81}x^4 - \frac{1}{9}x^5 + \frac{68}{729}x^6\right)y(0) \\ &\quad + \left(x + \frac{5}{18}x^3 - \frac{5}{54}x^4 + \frac{7}{72}x^5 - \frac{31}{486}x^6\right)y'(0) + O(x^6) \\ y &= \left(1 + \frac{2}{3}x^2 - \frac{4}{27}x^3 + \frac{16}{81}x^4 - \frac{1}{9}x^5\right)c_1 + \left(x + \frac{5}{18}x^3 - \frac{5}{54}x^4 + \frac{7}{72}x^5\right)c_2 + O(x^6) \end{aligned} \quad (1)$$

Verification of solutions

$$\begin{aligned} y &= \left(1 + \frac{2}{3}x^2 - \frac{4}{27}x^3 + \frac{16}{81}x^4 - \frac{1}{9}x^5 + \frac{68}{729}x^6\right)y(0) \\ &\quad + \left(x + \frac{5}{18}x^3 - \frac{5}{54}x^4 + \frac{7}{72}x^5 - \frac{31}{486}x^6\right)y'(0) + O(x^6) \end{aligned}$$

Verified OK.

$$y = \left(1 + \frac{2}{3}x^2 - \frac{4}{27}x^3 + \frac{16}{81}x^4 - \frac{1}{9}x^5\right)c_1 + \left(x + \frac{5}{18}x^3 - \frac{5}{54}x^4 + \frac{7}{72}x^5\right)c_2 + O(x^6)$$

Verified OK.

14.7.1 Maple step by step solution

Let's solve

$$(x^2 - 2x - 3)y'' + y'x + 4y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{4y}{x^2-2x-3} - \frac{xy'}{x^2-2x-3}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{xy'}{x^2-2x-3} + \frac{4y}{x^2-2x-3} = 0$$

- Check to see if x_0 is a regular singular point

- Define functions

$$\left[P_2(x) = \frac{x}{x^2-2x-3}, P_3(x) = \frac{4}{x^2-2x-3} \right]$$

- $(x+1) \cdot P_2(x)$ is analytic at $x = -1$

$$\left. ((x+1) \cdot P_2(x)) \right|_{x=-1} = \frac{1}{4}$$

- $(x+1)^2 \cdot P_3(x)$ is analytic at $x = -1$

$$\left. ((x+1)^2 \cdot P_3(x)) \right|_{x=-1} = 0$$

- $x = -1$ is a regular singular point

Check to see if x_0 is a regular singular point

$$x_0 = -1$$

- Multiply by denominators

$$(x^2 - 2x - 3)y'' + y'x + 4y = 0$$

- Change variables using $x = u - 1$ so that the regular singular point is at $u = 0$

$$(u^2 - 4u) \left(\frac{d^2}{du^2} y(u) \right) + (u - 1) \left(\frac{d}{du} y(u) \right) + 4y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du}y(u)\right) = \sum_{k=0}^{\infty} a_k(k+r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k+1-m$

$$u^m \cdot \left(\frac{d}{du}y(u)\right) = \sum_{k=-1+m}^{\infty} a_{k+1-m}(k+1-m+r) u^{k+r}$$

- Convert $u^m \cdot \left(\frac{d^2}{du^2}y(u)\right)$ to series expansion for $m = 1.2$

$$u^m \cdot \left(\frac{d^2}{du^2}y(u)\right) = \sum_{k=0}^{\infty} a_k(k+r)(k+r-1) u^{k+r-2+m}$$

- Shift index using $k \rightarrow k+2-m$

$$u^m \cdot \left(\frac{d^2}{du^2}y(u)\right) = \sum_{k=-2+m}^{\infty} a_{k+2-m}(k+2-m+r)(k+1-m+r) u^{k+r}$$

Rewrite ODE with series expansions

$$-a_0r(-3+4r)u^{-1+r} + \left(\sum_{k=0}^{\infty} (-a_{k+1}(k+1+r)(4k+1+4r) + a_k(k^2+2kr+r^2+4))u^{k+r}\right) =$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$-r(-3+4r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \left\{0, \frac{3}{4}\right\}$$

- Each term in the series must be 0, giving the recursion relation

$$-4\left(k + \frac{1}{4} + r\right)(k+1+r)a_{k+1} + a_k(k^2+2kr+r^2+4) = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+1} = \frac{a_k(k^2+2kr+r^2+4)}{(4k+1+4r)(k+1+r)}$$

- Recursion relation for $r = 0$

$$a_{k+1} = \frac{a_k(k^2+4)}{(4k+1)(k+1)}$$

- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{k+1} = \frac{a_k(k^2+4)}{(4k+1)(k+1)} \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^k, a_{k+1} = \frac{a_k(k^2+4)}{(4k+1)(k+1)} \right]$$

- Recursion relation for $r = \frac{3}{4}$

$$a_{k+1} = \frac{a_k(k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k+\frac{7}{4})}$$

- Solution for $r = \frac{3}{4}$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+\frac{3}{4}}, a_{k+1} = \frac{a_k(k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k+\frac{7}{4})} \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^{k+\frac{3}{4}}, a_{k+1} = \frac{a_k(k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k+\frac{7}{4})} \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k (x+1)^k \right) + \left(\sum_{k=0}^{\infty} b_k (x+1)^{k+\frac{3}{4}} \right), a_{k+1} = \frac{a_k(k^2+4)}{(4k+1)(k+1)}, b_{k+1} = \frac{b_k(k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k+\frac{7}{4})} \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
  -> Bessel
  -> elliptic
  -> Legendre
  -> Kummer
      -> hyper3: Equivalence to 1F1 under a power @ Moebius
  -> hypergeometric
      -> heuristic approach
          <- heuristic approach successful
          <- hypergeometric successful
<- special function solution successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 49

```
Order:=6;
```

```
dsolve((x^2-2*x-3)*diff(y(x),x$2)+x*diff(y(x),x)+4*y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 + \frac{2}{3}x^2 - \frac{4}{27}x^3 + \frac{16}{81}x^4 - \frac{1}{9}x^5\right) y(0) \\ + \left(x + \frac{5}{18}x^3 - \frac{5}{54}x^4 + \frac{7}{72}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 63

```
AsymptoticDSolveValue[(x^2-2*x-3)*y'[x]+x*y'[x]+4*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(\frac{7x^5}{72} - \frac{5x^4}{54} + \frac{5x^3}{18} + x \right) + c_1 \left(-\frac{x^5}{9} + \frac{16x^4}{81} - \frac{4x^3}{27} + \frac{2x^2}{3} + 1 \right)$$

14.8 problem 6. case $x_0 = 4$ only

14.8.1 Maple step by step solution 3415

Internal problem ID [740]

Internal file name [OUTPUT/740_Sunday_June_05_2022_01_48_25_AM_34591398/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 6. case $x_0 = 4$ only.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(x^2 - 2x - 3)y'' + y'x + 4y = 0$$

With the expansion point for the power series method at $x = 4$.

The ode does not have its expansion point at $x = 0$, therefore to simplify the computation of power series expansion, change of variable is made on the independent variable to shift the initial conditions and the expansion point back to zero. The new ode is then solved more easily since the expansion point is now at zero. The solution converted back to the original independent variable. Let

$$t = x - 4$$

The ode is converted to be in terms of the new independent variable t . This results in

$$((t + 4)^2 - 2t - 11) \left(\frac{d^2}{dt^2} y(t) \right) + \left(\frac{d}{dt} y(t) \right) (t + 4) + 4y(t) = 0$$

With its expansion point and initial conditions now at $t = 0$. The transformed ODE is now solved. Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (831)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (832)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = -\frac{t\left(\frac{d}{dt}y(t)\right) + 4\frac{d}{dt}y(t) + 4y(t)}{t^2 + 6t + 5}$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dt} \\ &= \frac{\partial F_0}{\partial t} + \frac{\partial F_0}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_0}{\partial \frac{d}{dt}y(t)} F_0 \\ &= \frac{(-2t^2 - 8t + 15)\left(\frac{d}{dt}y(t)\right) + (12t + 40)y(t)}{(t^2 + 6t + 5)^2} \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dt} \\ &= \frac{\partial F_1}{\partial t} + \frac{\partial F_1}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_1}{\partial \frac{d}{dt}y(t)} F_1 \\ &= \frac{(18t^3 + 152t^2 + 285t - 80)\left(\frac{d}{dt}y(t)\right) - 28\left(t^2 + \frac{50}{7}t + \frac{120}{7}\right)y(t)}{(t^2 + 6t + 5)^3} \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dt} \\ &= \frac{\partial F_2}{\partial t} + \frac{\partial F_2}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_2}{\partial \frac{d}{dt}y(t)} F_2 \\ &= \frac{(-100t^4 - 1200t^3 - 4780t^2 - 6360t + 785)\left(\frac{d}{dt}y(t)\right) + 40\left(t^3 + 14t^2 + \frac{193}{2}t + 199\right)y(t)}{(t^2 + 6t + 5)^4} \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dt} \\ &= \frac{\partial F_3}{\partial t} + \frac{\partial F_3}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_3}{\partial \frac{d}{dt}y(t)} F_3 \\ &= \frac{(540t^5 + 8400t^4 + 50880t^3 + 143280t^2 + 152115t - 13980)\left(\frac{d}{dt}y(t)\right) + 200\left(t^4 + 6t^3 - \frac{701}{10}t^2 - \frac{2553}{5}t - \frac{120}{5}\right)y(t)}{(t^2 + 6t + 5)^5} \end{aligned}$$

And so on. Evaluating all the above at initial conditions $t = 0$ and $y(0) = y(0)$ and

$y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -\frac{4y(0)}{5} - \frac{4y'(0)}{5} \\
 F_1 &= \frac{8y(0)}{5} + \frac{3y'(0)}{5} \\
 F_2 &= -\frac{96y(0)}{25} - \frac{16y'(0)}{25} \\
 F_3 &= \frac{1592y(0)}{125} + \frac{157y'(0)}{125} \\
 F_4 &= -\frac{34976y(0)}{625} - \frac{2796y'(0)}{625}
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$\begin{aligned}
 y(t) &= \left(1 - \frac{2}{5}t^2 + \frac{4}{15}t^3 - \frac{4}{25}t^4 + \frac{199}{1875}t^5 - \frac{2186}{28125}t^6\right) y(0) \\
 &\quad + \left(t - \frac{2}{5}t^2 + \frac{1}{10}t^3 - \frac{2}{75}t^4 + \frac{157}{15000}t^5 - \frac{233}{37500}t^6\right) y'(0) + O(t^6)
 \end{aligned}$$

Since the expansion point $t = 0$ is an ordinary, we can also solve this using standard power series. The ode is normalized to be

$$\left(\frac{d^2}{dt^2}y(t)\right) (t^2 + 6t + 5) + \left(\frac{d}{dt}y(t)\right) (t + 4) + 4y(t) = 0$$

Let the solution be represented as power series of the form

$$y(t) = \sum_{n=0}^{\infty} a_n t^n$$

Then

$$\begin{aligned}
 \frac{d}{dt}y(t) &= \sum_{n=1}^{\infty} n a_n t^{n-1} \\
 \frac{d^2}{dt^2}y(t) &= \sum_{n=2}^{\infty} n(n-1) a_n t^{n-2}
 \end{aligned}$$

Substituting the above back into the ode gives

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2}\right) (t^2 + 6t + 5) + \left(\sum_{n=1}^{\infty} n a_n t^{n-1}\right) (t + 4) + 4\left(\sum_{n=0}^{\infty} a_n t^n\right) = 0 \quad (1)$$

Which simplifies to

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} t^n a_n n(n-1) \right) + \left(\sum_{n=2}^{\infty} 6n t^{n-1} a_n (n-1) \right) + \left(\sum_{n=2}^{\infty} 5n(n-1) a_n t^{n-2} \right) \\ & + \left(\sum_{n=1}^{\infty} n a_n t^n \right) + \left(\sum_{n=1}^{\infty} 4n a_n t^{n-1} \right) + \left(\sum_{n=0}^{\infty} 4a_n t^n \right) = 0 \end{aligned} \quad (2)$$

The next step is to make all powers of t be n in each summation term. Going over each summation term above with power of t in it which is not already t^n and adjusting the power and the corresponding index gives

$$\begin{aligned} \sum_{n=2}^{\infty} 6n t^{n-1} a_n (n-1) &= \sum_{n=1}^{\infty} 6(n+1) a_{n+1} n t^n \\ \sum_{n=2}^{\infty} 5n(n-1) a_n t^{n-2} &= \sum_{n=0}^{\infty} 5(n+2) a_{n+2} (n+1) t^n \\ \sum_{n=1}^{\infty} 4n a_n t^{n-1} &= \sum_{n=0}^{\infty} 4(n+1) a_{n+1} t^n \end{aligned}$$

Substituting all the above in Eq (2) gives the following equation where now all powers of t are the same and equal to n .

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} t^n a_n n(n-1) \right) + \left(\sum_{n=1}^{\infty} 6(n+1) a_{n+1} n t^n \right) + \left(\sum_{n=0}^{\infty} 5(n+2) a_{n+2} (n+1) t^n \right) \\ & + \left(\sum_{n=1}^{\infty} n a_n t^n \right) + \left(\sum_{n=0}^{\infty} 4(n+1) a_{n+1} t^n \right) + \left(\sum_{n=0}^{\infty} 4a_n t^n \right) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$10a_2 + 4a_1 + 4a_0 = 0$$

$$a_2 = -\frac{2a_0}{5} - \frac{2a_1}{5}$$

$n = 1$ gives

$$20a_2 + 30a_3 + 5a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = \frac{4a_0}{15} + \frac{a_1}{10}$$

For $2 \leq n$, the recurrence equation is

$$na_n(n-1) + 6(n+1)a_{n+1}n + 5(n+2)a_{n+2}(n+1) + na_n + 4(n+1)a_{n+1} + 4a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} a_{n+2} &= -\frac{n^2a_n + 6n^2a_{n+1} + 10na_{n+1} + 4a_n + 4a_{n+1}}{5(n+2)(n+1)} \\ (5) \quad &= -\frac{(n^2+4)a_n}{5(n+2)(n+1)} - \frac{(6n^2+10n+4)a_{n+1}}{5(n+2)(n+1)} \end{aligned}$$

For $n = 2$ the recurrence equation gives

$$8a_2 + 48a_3 + 60a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = -\frac{4a_0}{25} - \frac{2a_1}{75}$$

For $n = 3$ the recurrence equation gives

$$13a_3 + 88a_4 + 100a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{199a_0}{1875} + \frac{157a_1}{15000}$$

For $n = 4$ the recurrence equation gives

$$20a_4 + 140a_5 + 150a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{2186a_0}{28125} - \frac{233a_1}{37500}$$

For $n = 5$ the recurrence equation gives

$$29a_5 + 204a_6 + 210a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{39931a_0}{656250} + \frac{72299a_1}{15750000}$$

And so on. Therefore the solution is

$$\begin{aligned} y(t) &= \sum_{n=0}^{\infty} a_n t^n \\ &= a_3 t^3 + a_2 t^2 + a_1 t + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$\begin{aligned} y(t) &= a_0 + a_1 t + \left(-\frac{2a_0}{5} - \frac{2a_1}{5} \right) t^2 + \left(\frac{4a_0}{15} + \frac{a_1}{10} \right) t^3 \\ &\quad + \left(-\frac{4a_0}{25} - \frac{2a_1}{75} \right) t^4 + \left(\frac{199a_0}{1875} + \frac{157a_1}{15000} \right) t^5 + \dots \end{aligned}$$

Collecting terms, the solution becomes

$$\begin{aligned} y(t) &= \left(1 - \frac{2}{5}t^2 + \frac{4}{15}t^3 - \frac{4}{25}t^4 + \frac{199}{1875}t^5 \right) a_0 \\ &\quad + \left(t - \frac{2}{5}t^2 + \frac{1}{10}t^3 - \frac{2}{75}t^4 + \frac{157}{15000}t^5 \right) a_1 + O(t^6) \end{aligned} \quad (3)$$

At $t = 0$ the solution above becomes

$$y(t) = \left(1 - \frac{2}{5}t^2 + \frac{4}{15}t^3 - \frac{4}{25}t^4 + \frac{199}{1875}t^5 \right) c_1 + \left(t - \frac{2}{5}t^2 + \frac{1}{10}t^3 - \frac{2}{75}t^4 + \frac{157}{15000}t^5 \right) c_2 + O(t^6)$$

Replacing t in the above with the original independent variable x using $t = x - 4$ results in

$$\begin{aligned} y &= \left(1 - \frac{2(x-4)^2}{5} + \frac{4(x-4)^3}{15} - \frac{4(x-4)^4}{25} + \frac{199(x-4)^5}{1875} - \frac{2186(x-4)^6}{28125} \right) y(4) \\ &\quad + \left(x - 4 - \frac{2(x-4)^2}{5} + \frac{(x-4)^3}{10} - \frac{2(x-4)^4}{75} + \frac{157(x-4)^5}{15000} - \frac{233(x-4)^6}{37500} \right) y'(4) \\ &\quad + O((x-4)^6) \end{aligned}$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{2(x-4)^2}{5} + \frac{4(x-4)^3}{15} - \frac{4(x-4)^4}{25} + \frac{199(x-4)^5}{1875} - \frac{2186(x-4)^6}{28125} \right) y(4) \\ + \left(x-4 - \frac{2(x-4)^2}{5} + \frac{(x-4)^3}{10} - \frac{2(x-4)^4}{75} + \frac{157(x-4)^5}{15000} - \frac{233(x-4)^6}{37500} \right) y'(4) \\ + O((x-4)^6)$$

Verification of solutions

$$y = \left(1 - \frac{2(x-4)^2}{5} + \frac{4(x-4)^3}{15} - \frac{4(x-4)^4}{25} + \frac{199(x-4)^5}{1875} - \frac{2186(x-4)^6}{28125} \right) y(4) \\ + \left(x-4 - \frac{2(x-4)^2}{5} + \frac{(x-4)^3}{10} - \frac{2(x-4)^4}{75} + \frac{157(x-4)^5}{15000} - \frac{233(x-4)^6}{37500} \right) y'(4) \\ + O((x-4)^6)$$

Verified OK.

14.8.1 Maple step by step solution

Let's solve

$$(x^2 - 2x - 3)y'' + y'x + 4y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{4y}{x^2-2x-3} - \frac{xy'}{x^2-2x-3}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{xy'}{x^2-2x-3} + \frac{4y}{x^2-2x-3} = 0$$

- Check to see if x_0 is a regular singular point

- Define functions

$$\left[P_2(x) = \frac{x}{x^2-2x-3}, P_3(x) = \frac{4}{x^2-2x-3} \right]$$

- $(x+1) \cdot P_2(x)$ is analytic at $x = -1$

$$\left. ((x+1) \cdot P_2(x)) \right|_{x=-1} = \frac{1}{4}$$

- $(x + 1)^2 \cdot P_3(x)$ is analytic at $x = -1$

$$\left. ((x + 1)^2 \cdot P_3(x)) \right|_{x=-1} = 0$$

- $x = -1$ is a regular singular point

Check to see if x_0 is a regular singular point

$$x_0 = -1$$

- Multiply by denominators

$$(x^2 - 2x - 3)y'' + y'x + 4y = 0$$

- Change variables using $x = u - 1$ so that the regular singular point is at $u = 0$

$$(u^2 - 4u) \left(\frac{d^2}{du^2} y(u) \right) + (u - 1) \left(\frac{d}{du} y(u) \right) + 4y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=0}^{\infty} a_k (k + r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k + 1 - m$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k + 1 - m + r) u^{k+r}$$

- Convert $u^m \cdot \left(\frac{d^2}{du^2} y(u) \right)$ to series expansion for $m = 1..2$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=0}^{\infty} a_k (k + r) (k + r - 1) u^{k+r-2+m}$$

- Shift index using $k \rightarrow k + 2 - m$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=-2+m}^{\infty} a_{k+2-m} (k + 2 - m + r) (k + 1 - m + r) u^{k+r}$$

Rewrite ODE with series expansions

$$-a_0 r(-3 + 4r) u^{-1+r} + \left(\sum_{k=0}^{\infty} (-a_{k+1} (k + 1 + r) (4k + 1 + 4r) + a_k (k^2 + 2kr + r^2 + 4)) u^{k+r} \right) =$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$-r(-3 + 4r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \left\{0, \frac{3}{4}\right\}$$

- Each term in the series must be 0, giving the recursion relation

$$-4\left(k + \frac{1}{4} + r\right)(k + 1 + r)a_{k+1} + a_k(k^2 + 2kr + r^2 + 4) = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+1} = \frac{a_k(k^2 + 2kr + r^2 + 4)}{(4k+1+4r)(k+1+r)}$$

- Recursion relation for $r = 0$

$$a_{k+1} = \frac{a_k(k^2 + 4)}{(4k+1)(k+1)}$$

- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{k+1} = \frac{a_k(k^2 + 4)}{(4k+1)(k+1)} \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x + 1)^k, a_{k+1} = \frac{a_k(k^2 + 4)}{(4k+1)(k+1)} \right]$$

- Recursion relation for $r = \frac{3}{4}$

$$a_{k+1} = \frac{a_k(k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k + \frac{7}{4})}$$

- Solution for $r = \frac{3}{4}$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k + \frac{3}{4}}, a_{k+1} = \frac{a_k(k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k + \frac{7}{4})} \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x + 1)^{k + \frac{3}{4}}, a_{k+1} = \frac{a_k(k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k + \frac{7}{4})} \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k (x + 1)^k \right) + \left(\sum_{k=0}^{\infty} b_k (x + 1)^{k + \frac{3}{4}} \right), a_{k+1} = \frac{a_k(k^2 + 4)}{(4k+1)(k+1)}, b_{k+1} = \frac{b_k(k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k + \frac{7}{4})} \right]$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacic's algorithm  
<- No Liouvillian solutions exist  
-> Trying a solution in terms of special functions:  
  -> Bessel  
  -> elliptic  
  -> Legendre  
  -> Kummer  
    -> hyper3: Equivalence to 1F1 under a power @ Moebius  
  -> hypergeometric  
    -> heuristic approach  
      <- heuristic approach successful  
      <- hypergeometric successful  
<- special function solution successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 54

```
Order:=6;  
dsolve((x^2-2*x-3)*diff(y(x),x$2)+x*diff(y(x),x)+4*y(x)=0,y(x),type='series',x=4);
```

$$y(x) = \left(1 - \frac{2(x-4)^2}{5} + \frac{4(x-4)^3}{15} - \frac{4(x-4)^4}{25} + \frac{199(x-4)^5}{1875}\right) y(4) \\ + \left(x - 4 - \frac{2(x-4)^2}{5} + \frac{(x-4)^3}{10} - \frac{2(x-4)^4}{75} + \frac{157(x-4)^5}{15000}\right) D(y)(4) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 87

```
AsymptoticDSolveValue[(x^2-2*x-3)*y'[x]+x*y'[x]+4*y[x]==0,y[x],{x,4,5}]
```

$$y(x) \rightarrow c_1 \left(\frac{199(x-4)^5}{1875} - \frac{4}{25}(x-4)^4 + \frac{4}{15}(x-4)^3 - \frac{2}{5}(x-4)^2 + 1 \right) \\ + c_2 \left(\frac{157(x-4)^5}{15000} - \frac{2}{75}(x-4)^4 + \frac{1}{10}(x-4)^3 - \frac{2}{5}(x-4)^2 + x - 4 \right)$$

14.9 problem 6. case $x_0 = -4$

14.9.1 Maple step by step solution 3428

Internal problem ID [741]

Internal file name [OUTPUT/741_Sunday_June_05_2022_01_48_27_AM_34703844/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 6. case $x_0 = -4$.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(x^2 - 2x - 3)y'' + y'x + 4y = 0$$

With the expansion point for the power series method at $x = -4$.

The ode does not have its expansion point at $x = 0$, therefore to simplify the computation of power series expansion, change of variable is made on the independent variable to shift the initial conditions and the expansion point back to zero. The new ode is then solved more easily since the expansion point is now at zero. The solution converted back to the original independent variable. Let

$$t = x + 4$$

The ode is converted to be in terms of the new independent variable t . This results in

$$((t - 4)^2 - 2t + 5) \left(\frac{d^2}{dt^2} y(t) \right) + \left(\frac{d}{dt} y(t) \right) (t - 4) + 4y(t) = 0$$

With its expansion point and initial conditions now at $t = 0$. The transformed ODE is now solved. Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (834)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (835)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = -\frac{t\left(\frac{d}{dt}y(t)\right) - 4\frac{d}{dt}y(t) + 4y(t)}{t^2 - 10t + 21}$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dt} \\ &= \frac{\partial F_0}{\partial t} + \frac{\partial F_0}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_0}{\partial \frac{d}{dt}y(t)} F_0 \\ &= \frac{(-2t^2 + 24t - 49)\left(\frac{d}{dt}y(t)\right) + (12t - 56)y(t)}{(t-3)^2(t-7)^2} \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dt} \\ &= \frac{\partial F_1}{\partial t} + \frac{\partial F_1}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_1}{\partial \frac{d}{dt}y(t)} F_1 \\ &= \frac{(18t^3 - 280t^2 + 1309t - 1848)\left(\frac{d}{dt}y(t)\right) - 28\left(t^2 - \frac{62}{7}t + 24\right)y(t)}{(t-3)^3(t-7)^3} \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dt} \\ &= \frac{\partial F_2}{\partial t} + \frac{\partial F_2}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_2}{\partial \frac{d}{dt}y(t)} F_2 \\ &= \frac{(-100t^4 + 2000t^3 - 14380t^2 + 44520t - 49455)\left(\frac{d}{dt}y(t)\right) + 40y(t)\left(t^3 - 10t^2 + \frac{129}{2}t - 189\right)}{(t-3)^4(t-7)^4} \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dt} \\ &= \frac{\partial F_3}{\partial t} + \frac{\partial F_3}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_3}{\partial \frac{d}{dt}y(t)} F_3 \\ &= \frac{(540t^5 - 13200t^4 + 127680t^3 - 617040t^2 + 1484595t - 1399860)\left(\frac{d}{dt}y(t)\right) + 200\left(t^4 - 26t^3 + \frac{1699}{10}t^2 - \dots\right)}{(t-3)^5(t-7)^5} \end{aligned}$$

And so on. Evaluating all the above at initial conditions $t = 0$ and $y(0) = y(0)$ and

$y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -\frac{4y(0)}{21} + \frac{4y'(0)}{21} \\
 F_1 &= -\frac{8y(0)}{63} - \frac{y'(0)}{9} \\
 F_2 &= -\frac{32y(0)}{441} - \frac{88y'(0)}{441} \\
 F_3 &= -\frac{40y(0)}{1029} - \frac{785y'(0)}{3087} \\
 F_4 &= -\frac{800y(0)}{64827} - \frac{22220y'(0)}{64827}
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$\begin{aligned}
 y(t) &= \left(1 - \frac{2}{21}t^2 - \frac{4}{189}t^3 - \frac{4}{1323}t^4 - \frac{1}{3087}t^5 - \frac{10}{583443}t^6\right) y(0) \\
 &+ \left(t + \frac{2}{21}t^2 - \frac{1}{54}t^3 - \frac{11}{1323}t^4 - \frac{157}{74088}t^5 - \frac{1111}{2333772}t^6\right) y'(0) + O(t^6)
 \end{aligned}$$

Since the expansion point $t = 0$ is an ordinary, we can also solve this using standard power series The ode is normalized to be

$$\left(\frac{d^2}{dt^2}y(t)\right) (t^2 - 10t + 21) + \left(\frac{d}{dt}y(t)\right) (t - 4) + 4y(t) = 0$$

Let the solution be represented as power series of the form

$$y(t) = \sum_{n=0}^{\infty} a_n t^n$$

Then

$$\begin{aligned}
 \frac{d}{dt}y(t) &= \sum_{n=1}^{\infty} n a_n t^{n-1} \\
 \frac{d^2}{dt^2}y(t) &= \sum_{n=2}^{\infty} n(n-1) a_n t^{n-2}
 \end{aligned}$$

Substituting the above back into the ode gives

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2}\right) (t^2 - 10t + 21) + \left(\sum_{n=1}^{\infty} n a_n t^{n-1}\right) (t - 4) + 4\left(\sum_{n=0}^{\infty} a_n t^n\right) = 0 \tag{1}$$

Which simplifies to

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} t^n a_n n(n-1) \right) + \sum_{n=2}^{\infty} (-10n t^{n-1} a_n (n-1)) + \left(\sum_{n=2}^{\infty} 21n(n-1) a_n t^{n-2} \right) \\ & + \left(\sum_{n=1}^{\infty} n a_n t^n \right) + \sum_{n=1}^{\infty} (-4n a_n t^{n-1}) + \left(\sum_{n=0}^{\infty} 4a_n t^n \right) = 0 \end{aligned} \quad (2)$$

The next step is to make all powers of t be n in each summation term. Going over each summation term above with power of t in it which is not already t^n and adjusting the power and the corresponding index gives

$$\begin{aligned} \sum_{n=2}^{\infty} (-10n t^{n-1} a_n (n-1)) &= \sum_{n=1}^{\infty} (-10(n+1) a_{n+1} n t^n) \\ \sum_{n=2}^{\infty} 21n(n-1) a_n t^{n-2} &= \sum_{n=0}^{\infty} 21(n+2) a_{n+2} (n+1) t^n \\ \sum_{n=1}^{\infty} (-4n a_n t^{n-1}) &= \sum_{n=0}^{\infty} (-4(n+1) a_{n+1} t^n) \end{aligned}$$

Substituting all the above in Eq (2) gives the following equation where now all powers of t are the same and equal to n .

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} t^n a_n n(n-1) \right) + \sum_{n=1}^{\infty} (-10(n+1) a_{n+1} n t^n) \\ & + \left(\sum_{n=0}^{\infty} 21(n+2) a_{n+2} (n+1) t^n \right) + \left(\sum_{n=1}^{\infty} n a_n t^n \right) \\ & + \sum_{n=0}^{\infty} (-4(n+1) a_{n+1} t^n) + \left(\sum_{n=0}^{\infty} 4a_n t^n \right) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$42a_2 - 4a_1 + 4a_0 = 0$$

$$a_2 = -\frac{2a_0}{21} + \frac{2a_1}{21}$$

$n = 1$ gives

$$-28a_2 + 126a_3 + 5a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{4a_0}{189} - \frac{a_1}{54}$$

For $2 \leq n$, the recurrence equation is

$$na_n(n-1) - 10(n+1)a_{n+1}n + 21(n+2)a_{n+2}(n+1) + na_n - 4(n+1)a_{n+1} + 4a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} a_{n+2} &= -\frac{n^2a_n - 10n^2a_{n+1} - 14na_{n+1} + 4a_n - 4a_{n+1}}{21(n+2)(n+1)} \\ (5) \quad &= -\frac{(n^2+4)a_n}{21(n+2)(n+1)} - \frac{(-10n^2-14n-4)a_{n+1}}{21(n+2)(n+1)} \end{aligned}$$

For $n = 2$ the recurrence equation gives

$$8a_2 - 72a_3 + 252a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = -\frac{4a_0}{1323} - \frac{11a_1}{1323}$$

For $n = 3$ the recurrence equation gives

$$13a_3 - 136a_4 + 420a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = -\frac{a_0}{3087} - \frac{157a_1}{74088}$$

For $n = 4$ the recurrence equation gives

$$20a_4 - 220a_5 + 630a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{10a_0}{583443} - \frac{1111a_1}{2333772}$$

For $n = 5$ the recurrence equation gives

$$29a_5 - 324a_6 + 882a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{83a_0}{19059138} - \frac{48121a_1}{457419312}$$

And so on. Therefore the solution is

$$\begin{aligned} y(t) &= \sum_{n=0}^{\infty} a_n t^n \\ &= a_3 t^3 + a_2 t^2 + a_1 t + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$\begin{aligned} y(t) &= a_0 + a_1 t + \left(-\frac{2a_0}{21} + \frac{2a_1}{21} \right) t^2 + \left(-\frac{4a_0}{189} - \frac{a_1}{54} \right) t^3 \\ &\quad + \left(-\frac{4a_0}{1323} - \frac{11a_1}{1323} \right) t^4 + \left(-\frac{a_0}{3087} - \frac{157a_1}{74088} \right) t^5 + \dots \end{aligned}$$

Collecting terms, the solution becomes

$$\begin{aligned} y(t) &= \left(1 - \frac{2}{21}t^2 - \frac{4}{189}t^3 - \frac{4}{1323}t^4 - \frac{1}{3087}t^5 \right) a_0 \\ &\quad + \left(t + \frac{2}{21}t^2 - \frac{1}{54}t^3 - \frac{11}{1323}t^4 - \frac{157}{74088}t^5 \right) a_1 + O(t^6) \end{aligned} \tag{3}$$

At $t = 0$ the solution above becomes

$$\begin{aligned} y(t) &= \left(1 - \frac{2}{21}t^2 - \frac{4}{189}t^3 - \frac{4}{1323}t^4 - \frac{1}{3087}t^5 \right) c_1 \\ &\quad + \left(t + \frac{2}{21}t^2 - \frac{1}{54}t^3 - \frac{11}{1323}t^4 - \frac{157}{74088}t^5 \right) c_2 + O(t^6) \end{aligned}$$

Replacing t in the above with the original independent variable x using $t = x + 4$ results in

$$y = \left(1 - \frac{2(x+4)^2}{21} - \frac{4(x+4)^3}{189} - \frac{4(x+4)^4}{1323} - \frac{(x+4)^5}{3087} - \frac{10(x+4)^6}{583443} \right) y(-4) \\ + \left(x+4 + \frac{2(x+4)^2}{21} - \frac{(x+4)^3}{54} - \frac{11(x+4)^4}{1323} - \frac{157(x+4)^5}{74088} - \frac{1111(x+4)^6}{2333772} \right) y'(-4) \\ + O((x+4)^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{2(x+4)^2}{21} - \frac{4(x+4)^3}{189} - \frac{4(x+4)^4}{1323} - \frac{(x+4)^5}{3087} - \frac{10(x+4)^6}{583443} \right) y(-4) \\ + \left(x+4 + \frac{2(x+4)^2}{21} - \frac{(x+4)^3}{54} - \frac{11(x+4)^4}{1323} - \frac{157(x+4)^5}{74088} \right. \\ \left. - \frac{1111(x+4)^6}{2333772} \right) y'(-4) + O((x+4)^6) \quad (1)$$

Verification of solutions

$$y = \left(1 - \frac{2(x+4)^2}{21} - \frac{4(x+4)^3}{189} - \frac{4(x+4)^4}{1323} - \frac{(x+4)^5}{3087} - \frac{10(x+4)^6}{583443} \right) y(-4) \\ + \left(x+4 + \frac{2(x+4)^2}{21} - \frac{(x+4)^3}{54} - \frac{11(x+4)^4}{1323} - \frac{157(x+4)^5}{74088} - \frac{1111(x+4)^6}{2333772} \right) y'(-4) \\ + O((x+4)^6)$$

Verified OK.

14.9.1 Maple step by step solution

Let's solve

$$(x^2 - 2x - 3)y'' + y'x + 4y = 0$$

- Highest derivative means the order of the ODE is 2
- $$y''$$
- Isolate 2nd derivative

$$y'' = -\frac{4y}{x^2-2x-3} - \frac{xy'}{x^2-2x-3}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{xy'}{x^2-2x-3} + \frac{4y}{x^2-2x-3} = 0$$

- Check to see if x_0 is a regular singular point

- Define functions

$$[P_2(x) = \frac{x}{x^2-2x-3}, P_3(x) = \frac{4}{x^2-2x-3}]$$

- $(x+1) \cdot P_2(x)$ is analytic at $x = -1$

$$((x+1) \cdot P_2(x)) \Big|_{x=-1} = \frac{1}{4}$$

- $(x+1)^2 \cdot P_3(x)$ is analytic at $x = -1$

$$((x+1)^2 \cdot P_3(x)) \Big|_{x=-1} = 0$$

- $x = -1$ is a regular singular point

Check to see if x_0 is a regular singular point

$$x_0 = -1$$

- Multiply by denominators

$$(x^2 - 2x - 3)y'' + y'x + 4y = 0$$

- Change variables using $x = u - 1$ so that the regular singular point is at $u = 0$

$$(u^2 - 4u) \left(\frac{d^2}{du^2} y(u) \right) + (u - 1) \left(\frac{d}{du} y(u) \right) + 4y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k + 1 - m$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k+1-m+r) u^{k+r}$$

- Convert $u^m \cdot \left(\frac{d^2}{du^2} y(u) \right)$ to series expansion for $m = 1..2$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) (k+r-1) u^{k+r-2+m}$$

- Shift index using $k \rightarrow k+2-m$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=-2+m}^{\infty} a_{k+2-m} (k+2-m+r) (k+1-m+r) u^{k+r}$$

Rewrite ODE with series expansions

$$-a_0 r(-3+4r) u^{-1+r} + \left(\sum_{k=0}^{\infty} (-a_{k+1} (k+1+r) (4k+1+4r) + a_k (k^2 + 2kr + r^2 + 4)) u^{k+r} \right) =$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$-r(-3+4r) = 0$$
- Values of r that satisfy the indicial equation

$$r \in \left\{ 0, \frac{3}{4} \right\}$$
- Each term in the series must be 0, giving the recursion relation

$$-4\left(k + \frac{1}{4} + r\right) (k+1+r) a_{k+1} + a_k (k^2 + 2kr + r^2 + 4) = 0$$
- Recursion relation that defines series solution to ODE

$$a_{k+1} = \frac{a_k (k^2 + 2kr + r^2 + 4)}{(4k+1+4r)(k+1+r)}$$
- Recursion relation for $r = 0$

$$a_{k+1} = \frac{a_k (k^2 + 4)}{(4k+1)(k+1)}$$
- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{k+1} = \frac{a_k (k^2 + 4)}{(4k+1)(k+1)} \right]$$
- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^k, a_{k+1} = \frac{a_k (k^2 + 4)}{(4k+1)(k+1)} \right]$$
- Recursion relation for $r = \frac{3}{4}$

$$a_{k+1} = \frac{a_k (k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k+\frac{7}{4})}$$
- Solution for $r = \frac{3}{4}$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+\frac{3}{4}}, a_{k+1} = \frac{a_k (k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k+\frac{7}{4})} \right]$$
- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^{k+\frac{3}{4}}, a_{k+1} = \frac{a_k (k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k+\frac{7}{4})} \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k (x+1)^k \right) + \left(\sum_{k=0}^{\infty} b_k (x+1)^{k+\frac{3}{4}} \right), a_{k+1} = \frac{a_k (k^2+4)}{(4k+1)(k+1)}, b_{k+1} = \frac{b_k (k^2 + \frac{3}{2}k + \frac{73}{16})}{(4k+4)(k+\frac{7}{4})} \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
  -> Bessel
  -> elliptic
  -> Legendre
  -> Kummer
      -> hyper3: Equivalence to 1F1 under a power @ Moebius
  -> hypergeometric
      -> heuristic approach
          <- heuristic approach successful
          <- hypergeometric successful
<- special function solution successful`

```


✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 54

Order:=6;

```
dsolve((x^2-2*x-3)*diff(y(x),x$2)+x*diff(y(x),x)+4*y(x)=0,y(x),type='series',x=-4);
```

$$y(x) = \left(1 - \frac{2(x+4)^2}{21} - \frac{4(x+4)^3}{189} - \frac{4(x+4)^4}{1323} - \frac{(x+4)^5}{3087}\right) y(-4) \\ + \left(x+4 + \frac{2(x+4)^2}{21} - \frac{(x+4)^3}{54} - \frac{11(x+4)^4}{1323} - \frac{157(x+4)^5}{74088}\right) D(y)(-4) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 87

```
AsymptoticDSolveValue[(x^2-2*x-3)*y'[x]+x*y'[x]+4*y[x]==0,y[x],{x,-4,5}]
```

$$y(x) \rightarrow c_1 \left(-\frac{(x+4)^5}{3087} - \frac{4(x+4)^4}{1323} - \frac{4}{189}(x+4)^3 - \frac{2}{21}(x+4)^2 + 1 \right) \\ + c_2 \left(-\frac{157(x+4)^5}{74088} - \frac{11(x+4)^4}{1323} - \frac{1}{54}(x+4)^3 + \frac{2}{21}(x+4)^2 + x + 4 \right)$$

14.10 problem 7. case $x_0 = 0$

14.10.1 Maple step by step solution 3440

Internal problem ID [742]

Internal file name [OUTPUT/742_Sunday_June_05_2022_01_48_29_AM_48339428/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 7. case $x_0 = 0$.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second order series method. Ordinary point", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(x^3 + 1)y'' + 4y'x + y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (837)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (838)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = -\frac{4y'x + y}{x^3 + 1}$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dx} \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\ &= \frac{(7x^3 + 16x^2 - 5)y' + 3(x + \frac{4}{3})xy}{(x^3 + 1)^2} \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dx} \\ &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\ &= \frac{(-18x^5 - 88x^4 - 64x^3 + 54x^2 + 56x)y' - 12(x^4 + \frac{9}{4}x^3 + \frac{4}{3}x^2 - \frac{1}{2}x - \frac{3}{4})y}{(x^3 + 1)^3} \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dx} \\ &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\ &= \frac{(60x^7 + 485x^6 + 720x^5 - 218x^4 - 1034x^3 - 432x^2 + 114x + 65)y' + 60(x^6 + 3x^5 + \frac{10}{3}x^4 - \frac{8}{15}x^3 - \frac{1}{3}x^2 + \frac{1}{5}x - \frac{1}{15})y}{(x^3 + 1)^4} \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dx} \\ &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\ &= \frac{(-240x^9 - 2970x^8 - 6780x^7 - 688x^6 + 13052x^5 + 12168x^4 - 424x^3 - 4554x^2 - 1212x + 120)y' - 120(x^8 + 8x^7 + \frac{14}{3}x^6 - \frac{14}{3}x^5 - \frac{14}{3}x^4 + \frac{14}{3}x^3 - \frac{14}{3}x^2 + \frac{14}{3}x - \frac{14}{3})y}{(x^3 + 1)^5} \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$F_0 = -y(0)$$

$$F_1 = -5y'(0)$$

$$F_2 = 9y(0)$$

$$F_3 = 6y(0) + 65y'(0)$$

$$F_4 = -153y(0) + 120y'(0)$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 - \frac{1}{2}x^2 + \frac{3}{8}x^4 + \frac{1}{20}x^5 - \frac{17}{80}x^6\right)y(0) + \left(x - \frac{5}{6}x^3 + \frac{13}{24}x^5 + \frac{1}{6}x^6\right)y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series The ode is normalized to be

$$(x^3 + 1) y'' + 4y'x + y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

$$y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(x^3 + 1) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + 4 \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n x^{1+n} a_n (n-1) \right) + \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \left(\sum_{n=1}^{\infty} 4n a_n x^n \right) + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n x^{1+n} a_n (n-1) = \sum_{n=3}^{\infty} (n-1) a_{n-1} (n-2) x^n$$

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (1+n) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} & \left(\sum_{n=3}^{\infty} (n-1) a_{n-1} (n-2) x^n \right) + \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (1+n) x^n \right) \\ & + \left(\sum_{n=1}^{\infty} 4n a_n x^n \right) + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$2a_2 + a_0 = 0$$

$$a_2 = -\frac{a_0}{2}$$

$n = 1$ gives

$$6a_3 + 5a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{5a_1}{6}$$

$n = 2$ gives

$$12a_4 + 9a_2 = 0$$

Which after substituting earlier equations, simplifies to

$$12a_4 - \frac{9a_0}{2} = 0$$

Or

$$a_4 = \frac{3a_0}{8}$$

For $3 \leq n$, the recurrence equation is

$$(n-1) a_{n-1} (n-2) + (n+2) a_{n+2} (1+n) + 4n a_n + a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} a_{n+2} &= -\frac{n^2 a_{n-1} + 4n a_n - 3n a_{n-1} + a_n + 2a_{n-1}}{(n+2)(1+n)} \\ (5) \qquad &= -\frac{(4n+1)a_n}{(n+2)(1+n)} - \frac{(n^2-3n+2)a_{n-1}}{(n+2)(1+n)} \end{aligned}$$

For $n = 3$ the recurrence equation gives

$$2a_2 + 20a_5 + 13a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_0}{20} + \frac{13a_1}{24}$$

For $n = 4$ the recurrence equation gives

$$6a_3 + 30a_6 + 17a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{a_1}{6} - \frac{17a_0}{80}$$

For $n = 5$ the recurrence equation gives

$$12a_4 + 42a_7 + 21a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{37a_0}{280} - \frac{13a_1}{48}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_1x - \frac{a_0x^2}{2} - \frac{5a_1x^3}{6} + \frac{3a_0x^4}{8} + \left(\frac{a_0}{20} + \frac{13a_1}{24}\right)x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 - \frac{1}{2}x^2 + \frac{3}{8}x^4 + \frac{1}{20}x^5\right)a_0 + \left(x - \frac{5}{6}x^3 + \frac{13}{24}x^5\right)a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - \frac{1}{2}x^2 + \frac{3}{8}x^4 + \frac{1}{20}x^5\right)c_1 + \left(x - \frac{5}{6}x^3 + \frac{13}{24}x^5\right)c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{1}{2}x^2 + \frac{3}{8}x^4 + \frac{1}{20}x^5 - \frac{17}{80}x^6\right)y(0) + \left(x - \frac{5}{6}x^3 + \frac{13}{24}x^5 + \frac{1}{6}x^6\right)y'(0) + O(x^7)$$

$$y = \left(1 - \frac{1}{2}x^2 + \frac{3}{8}x^4 + \frac{1}{20}x^5\right)c_1 + \left(x - \frac{5}{6}x^3 + \frac{13}{24}x^5\right)c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 - \frac{1}{2}x^2 + \frac{3}{8}x^4 + \frac{1}{20}x^5 - \frac{17}{80}x^6\right)y(0) + \left(x - \frac{5}{6}x^3 + \frac{13}{24}x^5 + \frac{1}{6}x^6\right)y'(0) + O(x^7)$$

Verified OK.

$$y = \left(1 - \frac{1}{2}x^2 + \frac{3}{8}x^4 + \frac{1}{20}x^5\right)c_1 + \left(x - \frac{5}{6}x^3 + \frac{13}{24}x^5\right)c_2 + O(x^6)$$

Verified OK.

14.10.1 Maple step by step solution

Let's solve

$$(x^3 + 1)y'' + 4y'x + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{4xy'}{x^3+1} - \frac{y}{x^3+1}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{4xy'}{x^3+1} + \frac{y}{x^3+1} = 0$$

- Check to see if x_0 is a regular singular point

- Define functions

$$[P_2(x) = \frac{4x}{x^3+1}, P_3(x) = \frac{1}{x^3+1}]$$

- $(x+1) \cdot P_2(x)$ is analytic at $x = -1$

$$((x+1) \cdot P_2(x)) \Big|_{x=-1} = -\frac{4}{3}$$

- $(x+1)^2 \cdot P_3(x)$ is analytic at $x = -1$

$$((x+1)^2 \cdot P_3(x)) \Big|_{x=-1} = 0$$

- $x = -1$ is a regular singular point

Check to see if x_0 is a regular singular point

$$x_0 = -1$$

- Multiply by denominators

$$(x^3 + 1)y'' + 4y'x + y = 0$$

- Change variables using $x = u - 1$ so that the regular singular point is at $u = 0$

$$(u^3 - 3u^2 + 3u) \left(\frac{d^2}{du^2} y(u) \right) + (4u - 4) \left(\frac{d}{du} y(u) \right) + y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k+1-m$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k+1-m+r) u^{k+r}$$

- Convert $u^m \cdot \left(\frac{d^2}{du^2} y(u) \right)$ to series expansion for $m = 1..3$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) (k+r-1) u^{k+r-2+m}$$

- Shift index using $k \rightarrow k+2-m$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=-2+m}^{\infty} a_{k+2-m} (k+2-m+r) (k+1-m+r) u^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r (-7+3r) u^{-1+r} + (a_1 (1+r) (-4+3r) - a_0 (3r^2 - 7r - 1)) u^r + \left(\sum_{k=1}^{\infty} (a_{k+1} (k+1+r) (3k - \dots) \right)$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$r(-7+3r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \left\{ 0, \frac{7}{3} \right\}$$

- Each term must be 0

$$a_1 (1+r) (-4+3r) - a_0 (3r^2 - 7r - 1) = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(-3a_k + a_{k-1} + 3a_{k+1}) k^2 + ((-6a_k + 2a_{k-1} + 6a_{k+1}) r + 7a_k - 3a_{k-1} - a_{k+1}) k + (-3a_k + a_{k-1} - \dots)$$

- Shift index using $k \rightarrow k+1$

$$(-3a_{k+1} + a_k + 3a_{k+2}) (k+1)^2 + ((-6a_{k+1} + 2a_k + 6a_{k+2}) r + 7a_{k+1} - 3a_k - a_{k+2}) (k+1) + (-\dots)$$

- Recursion relation that defines series solution to ODE

$$a_{k+2} = - \frac{k^2 a_k - 3k^2 a_{k+1} + 2k r a_k - 6k r a_{k+1} + r^2 a_k - 3r^2 a_{k+1} - k a_k + k a_{k+1} - r a_k + r a_{k+1} + 5a_{k+1}}{3k^2 + 6kr + 3r^2 + 5k + 5r - 2}$$

- Recursion relation for $r = 0$

$$a_{k+2} = - \frac{k^2 a_k - 3k^2 a_{k+1} - k a_k + k a_{k+1} + 5a_{k+1}}{3k^2 + 5k - 2}$$

- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{k+2} = - \frac{k^2 a_k - 3k^2 a_{k+1} - k a_k + k a_{k+1} + 5a_{k+1}}{3k^2 + 5k - 2}, -4a_1 + a_0 = 0 \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^k, a_{k+2} = - \frac{k^2 a_k - 3k^2 a_{k+1} - k a_k + k a_{k+1} + 5a_{k+1}}{3k^2 + 5k - 2}, -4a_1 + a_0 = 0 \right]$$

- Recursion relation for $r = \frac{7}{3}$

$$a_{k+2} = - \frac{k^2 a_k - 3k^2 a_{k+1} + \frac{11}{3} k a_k - 13k a_{k+1} + \frac{28}{9} a_k - 9a_{k+1}}{3k^2 + 19k + 26}$$

- Solution for $r = \frac{7}{3}$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+\frac{7}{3}}, a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} + \frac{11}{3} k a_k - 13k a_{k+1} + \frac{28}{9} a_k - 9a_{k+1}}{3k^2 + 19k + 26}, 10a_1 + a_0 = 0 \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x + 1)^{k+\frac{7}{3}}, a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} + \frac{11}{3} k a_k - 13k a_{k+1} + \frac{28}{9} a_k - 9a_{k+1}}{3k^2 + 19k + 26}, 10a_1 + a_0 = 0 \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k (x + 1)^k \right) + \left(\sum_{k=0}^{\infty} b_k (x + 1)^{k+\frac{7}{3}} \right), a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} - k a_k + k a_{k+1} + 5a_{k+1}}{3k^2 + 5k - 2}, -4a_1 + a_0 = 0 \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
  -> Bessel
  -> elliptic
  -> Legendre
  -> Kummer
    -> hyper3: Equivalence to 1F1 under a power @ Moebius
  -> hypergeometric
    -> heuristic approach
    -> hyper3: Equivalence to 2F1, 1F1 or 0F1 under a power @ Moebius
  -> Mathieu
    -> Equivalence to the rational form of Mathieu ODE under a power @ Moebius
trying a solution in terms of MeijerG functions
-> Heun: Equivalence to the GHE or one of its 4 confluent cases under a power @ Moebius
<- Heun successful: received ODE is equivalent to the HeunG ODE, case a <> 0, e <> 0, g <

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 39

```
Order:=6;  
dsolve((1+x^3)*diff(y(x),x$2)+4*x*diff(y(x),x)+y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 - \frac{1}{2}x^2 + \frac{3}{8}x^4 + \frac{1}{20}x^5\right) y(0) + \left(x - \frac{5}{6}x^3 + \frac{13}{24}x^5\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 49

```
AsymptoticDSolveValue[(1+x^3)*y''[x]+4*x*y'[x]+y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(\frac{13x^5}{24} - \frac{5x^3}{6} + x \right) + c_1 \left(\frac{x^5}{20} + \frac{3x^4}{8} - \frac{x^2}{2} + 1 \right)$$

14.11 problem 7. case $x_0 = 2$

14.11.1 Maple step by step solution 3454

Internal problem ID [743]

Internal file name [OUTPUT/743_Sunday_June_05_2022_01_48_32_AM_81422129/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 7. case $x_0 = 2$.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_2nd_order , _with_linear_symmetries]]
```

$$(x^3 + 1)y'' + 4y'x + y = 0$$

With the expansion point for the power series method at $x = 2$.

The ode does not have its expansion point at $x = 0$, therefore to simplify the computation of power series expansion, change of variable is made on the independent variable to shift the initial conditions and the expansion point back to zero. The new ode is then solved more easily since the expansion point is now at zero. The solution converted back to the original independent variable. Let

$$t = -2 + x$$

The ode is converted to be in terms of the new independent variable t . This results in

$$((2 + t)^3 + 1) \left(\frac{d^2}{dt^2} y(t) \right) + 4(2 + t) \left(\frac{d}{dt} y(t) \right) + y(t) = 0$$

With its expansion point and initial conditions now at $t = 0$. The transformed ODE is now solved. Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (840)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (841)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = -\frac{4t\left(\frac{d}{dt}y(t)\right) + 8\frac{d}{dt}y(t) + y(t)}{t^3 + 6t^2 + 12t + 9}$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dt} \\ &= \frac{\partial F_0}{\partial t} + \frac{\partial F_0}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_0}{\partial \frac{d}{dt}y(t)} F_0 \\ &= \frac{(7t^3 + 58t^2 + 148t + 115)\left(\frac{d}{dt}y(t)\right) + 3\left(t + \frac{10}{3}\right)(2+t)y(t)}{(t+3)^2(t^2+3t+3)^2} \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dt} \\ &= \frac{\partial F_1}{\partial t} + \frac{\partial F_1}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_1}{\partial \frac{d}{dt}y(t)} F_1 \\ &= \frac{(-18t^5 - 268t^4 - 1488t^3 - 3882t^2 - 4752t - 2168)\left(\frac{d}{dt}y(t)\right) - 12\left(t^4 + \frac{41}{4}t^3 + \frac{233}{6}t^2 + \frac{383}{6}t + \frac{451}{12}\right)y(t)}{(t+3)^3(t^2+3t+3)^3} \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dt} \\ &= \frac{\partial F_2}{\partial t} + \frac{\partial F_2}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_2}{\partial \frac{d}{dt}y(t)} F_2 \\ &= \frac{(60t^7 + 1325t^6 + 11580t^5 + 52882t^4 + 137222t^3 + 202452t^2 + 156602t + 48565)\left(\frac{d}{dt}y(t)\right) + 60(t^6 + 1)}{(t+3)^4(t^2+3t+3)^4} \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dt} \\ &= \frac{\partial F_3}{\partial t} + \frac{\partial F_3}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_3}{\partial \frac{d}{dt}y(t)} F_3 \\ &= \frac{(-240t^9 - 7290t^8 - 88860t^7 - 589528t^6 - 2379124t^5 - 6091072t^4 - 9900360t^3 - 9820346t^2 - 5354)}{(t+3)^5(t^2+3t+3)^5} \end{aligned}$$

And so on. Evaluating all the above at initial conditions $t = 0$ and $y(0) = y(0)$ and

$y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -\frac{y(0)}{9} - \frac{8y'(0)}{9} \\
 F_1 &= \frac{20y(0)}{81} + \frac{115y'(0)}{81} \\
 F_2 &= -\frac{451y(0)}{729} - \frac{2168y'(0)}{729} \\
 F_3 &= \frac{11510y(0)}{6561} + \frac{48565y'(0)}{6561} \\
 F_4 &= -\frac{322189y(0)}{59049} - \frac{1206632y'(0)}{59049}
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$\begin{aligned}
 y(t) &= \left(1 - \frac{1}{18}t^2 + \frac{10}{243}t^3 - \frac{451}{17496}t^4 + \frac{1151}{78732}t^5 - \frac{322189}{42515280}t^6\right) y(0) \\
 &+ \left(t - \frac{4}{9}t^2 + \frac{115}{486}t^3 - \frac{271}{2187}t^4 + \frac{9713}{157464}t^5 - \frac{150829}{5314410}t^6\right) y'(0) + O(t^6)
 \end{aligned}$$

Since the expansion point $t = 0$ is an ordinary, we can also solve this using standard power series. The ode is normalized to be

$$\left(\frac{d^2}{dt^2}y(t)\right) (t^3 + 6t^2 + 12t + 9) + (8 + 4t) \left(\frac{d}{dt}y(t)\right) + y(t) = 0$$

Let the solution be represented as power series of the form

$$y(t) = \sum_{n=0}^{\infty} a_n t^n$$

Then

$$\begin{aligned}
 \frac{d}{dt}y(t) &= \sum_{n=1}^{\infty} n a_n t^{n-1} \\
 \frac{d^2}{dt^2}y(t) &= \sum_{n=2}^{\infty} n(n-1) a_n t^{n-2}
 \end{aligned}$$

Substituting the above back into the ode gives

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2}\right) (t^3 + 6t^2 + 12t + 9) + (8 + 4t) \left(\sum_{n=1}^{\infty} n a_n t^{n-1}\right) + \left(\sum_{n=0}^{\infty} a_n t^n\right) = 0 \tag{1}$$

Which simplifies to

$$\begin{aligned} & \left(\sum_{n=2}^{\infty} n t^{1+n} a_n (n-1) \right) + \left(\sum_{n=2}^{\infty} 6 t^n a_n n (n-1) \right) + \left(\sum_{n=2}^{\infty} 12 n t^{n-1} a_n (n-1) \right) \\ & + \left(\sum_{n=2}^{\infty} 9 n (n-1) a_n t^{n-2} \right) + \left(\sum_{n=1}^{\infty} 8 n a_n t^{n-1} \right) + \left(\sum_{n=1}^{\infty} 4 n a_n t^n \right) + \left(\sum_{n=0}^{\infty} a_n t^n \right) \\ & = 0 \end{aligned} \quad (2)$$

The next step is to make all powers of t be n in each summation term. Going over each summation term above with power of t in it which is not already t^n and adjusting the power and the corresponding index gives

$$\begin{aligned} \sum_{n=2}^{\infty} n t^{1+n} a_n (n-1) &= \sum_{n=3}^{\infty} (n-1) a_{n-1} (n-2) t^n \\ \sum_{n=2}^{\infty} 12 n t^{n-1} a_n (n-1) &= \sum_{n=1}^{\infty} 12(1+n) a_{1+n} n t^n \\ \sum_{n=2}^{\infty} 9 n (n-1) a_n t^{n-2} &= \sum_{n=0}^{\infty} 9(n+2) a_{n+2} (1+n) t^n \\ \sum_{n=1}^{\infty} 8 n a_n t^{n-1} &= \sum_{n=0}^{\infty} 8(1+n) a_{1+n} t^n \end{aligned}$$

Substituting all the above in Eq (2) gives the following equation where now all powers of t are the same and equal to n .

$$\begin{aligned} & \left(\sum_{n=3}^{\infty} (n-1) a_{n-1} (n-2) t^n \right) + \left(\sum_{n=2}^{\infty} 6 t^n a_n n (n-1) \right) \\ & + \left(\sum_{n=1}^{\infty} 12(1+n) a_{1+n} n t^n \right) + \left(\sum_{n=0}^{\infty} 9(n+2) a_{n+2} (1+n) t^n \right) \\ & + \left(\sum_{n=0}^{\infty} 8(1+n) a_{1+n} t^n \right) + \left(\sum_{n=1}^{\infty} 4 n a_n t^n \right) + \left(\sum_{n=0}^{\infty} a_n t^n \right) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$18a_2 + 8a_1 + a_0 = 0$$

$$a_2 = -\frac{a_0}{18} - \frac{4a_1}{9}$$

$n = 1$ gives

$$40a_2 + 54a_3 + 5a_1 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = \frac{10a_0}{243} + \frac{115a_1}{486}$$

$n = 2$ gives

$$21a_2 + 96a_3 + 108a_4 = 0$$

Which after substituting earlier equations, simplifies to

$$\frac{451a_0}{162} + \frac{1084a_1}{81} + 108a_4 = 0$$

Or

$$a_4 = -\frac{451a_0}{17496} - \frac{271a_1}{2187}$$

For $3 \leq n$, the recurrence equation is

$$(n-1)a_{n-1}(n-2) + 6na_n(n-1) + 12(1+n)a_{1+n}n + 9(n+2)a_{n+2}(1+n) + 8(1+n)a_{1+n} + 4na_n + a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{6n^2a_n + 12n^2a_{1+n} + n^2a_{n-1} - 2na_n + 20na_{1+n} - 3na_{n-1} + a_n + 8a_{1+n} + 2a_{n-1}}{9(n+2)(1+n)}$$

$$(5) \quad = -\frac{(6n^2 - 2n + 1)a_n}{9(n+2)(1+n)} - \frac{(12n^2 + 20n + 8)a_{1+n}}{9(n+2)(1+n)} - \frac{(n^2 - 3n + 2)a_{n-1}}{9(n+2)(1+n)}$$

For $n = 3$ the recurrence equation gives

$$2a_2 + 49a_3 + 176a_4 + 180a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{1151a_0}{78732} + \frac{9713a_1}{157464}$$

For $n = 4$ the recurrence equation gives

$$6a_3 + 89a_4 + 280a_5 + 270a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{322189a_0}{42515280} - \frac{150829a_1}{5314410}$$

For $n = 5$ the recurrence equation gives

$$12a_4 + 141a_5 + 408a_6 + 378a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{4747261a_0}{1339231320} + \frac{30958471a_1}{2678462640}$$

And so on. Therefore the solution is

$$\begin{aligned} y(t) &= \sum_{n=0}^{\infty} a_n t^n \\ &= a_3 t^3 + a_2 t^2 + a_1 t + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$\begin{aligned} y(t) &= a_0 + a_1 t + \left(-\frac{a_0}{18} - \frac{4a_1}{9} \right) t^2 + \left(\frac{10a_0}{243} + \frac{115a_1}{486} \right) t^3 \\ &\quad + \left(-\frac{451a_0}{17496} - \frac{271a_1}{2187} \right) t^4 + \left(\frac{1151a_0}{78732} + \frac{9713a_1}{157464} \right) t^5 + \dots \end{aligned}$$

Collecting terms, the solution becomes

$$\begin{aligned} y(t) &= \left(1 - \frac{1}{18}t^2 + \frac{10}{243}t^3 - \frac{451}{17496}t^4 + \frac{1151}{78732}t^5 \right) a_0 \\ &\quad + \left(t - \frac{4}{9}t^2 + \frac{115}{486}t^3 - \frac{271}{2187}t^4 + \frac{9713}{157464}t^5 \right) a_1 + O(t^6) \end{aligned} \tag{3}$$

At $t = 0$ the solution above becomes

$$y(t) = \left(1 - \frac{1}{18}t^2 + \frac{10}{243}t^3 - \frac{451}{17496}t^4 + \frac{1151}{78732}t^5\right) c_1 \\ + \left(t - \frac{4}{9}t^2 + \frac{115}{486}t^3 - \frac{271}{2187}t^4 + \frac{9713}{157464}t^5\right) c_2 + O(t^6)$$

Replacing t in the above with the original independent variable x using $t = -2 + x$ results in

$$y = \left(1 - \frac{(-2+x)^2}{18} + \frac{10(-2+x)^3}{243} - \frac{451(-2+x)^4}{17496} + \frac{1151(-2+x)^5}{78732} \right. \\ \left. - \frac{322189(-2+x)^6}{42515280}\right) y(2) + \left(-2+x - \frac{4(-2+x)^2}{9} + \frac{115(-2+x)^3}{486} \right. \\ \left. - \frac{271(-2+x)^4}{2187} + \frac{9713(-2+x)^5}{157464} - \frac{150829(-2+x)^6}{5314410}\right) y'(2) + O((-2+x)^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{(-2+x)^2}{18} + \frac{10(-2+x)^3}{243} - \frac{451(-2+x)^4}{17496} + \frac{1151(-2+x)^5}{78732} \right. \\ \left. - \frac{322189(-2+x)^6}{42515280}\right) y(2) + \left(-2+x - \frac{4(-2+x)^2}{9} + \frac{115(-2+x)^3}{486} \right. \\ \left. - \frac{271(-2+x)^4}{2187} + \frac{9713(-2+x)^5}{157464} - \frac{150829(-2+x)^6}{5314410}\right) y'(2) + O((-2+x)^6)$$

Verification of solutions

$$y = \left(1 - \frac{(-2+x)^2}{18} + \frac{10(-2+x)^3}{243} - \frac{451(-2+x)^4}{17496} + \frac{1151(-2+x)^5}{78732} \right. \\ \left. - \frac{322189(-2+x)^6}{42515280}\right) y(2) + \left(-2+x - \frac{4(-2+x)^2}{9} + \frac{115(-2+x)^3}{486} \right. \\ \left. - \frac{271(-2+x)^4}{2187} + \frac{9713(-2+x)^5}{157464} - \frac{150829(-2+x)^6}{5314410}\right) y'(2) + O((-2+x)^6)$$

Verified OK.

14.11.1 Maple step by step solution

Let's solve

$$(x^3 + 1)y'' + 4y'x + y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{4xy'}{x^3+1} - \frac{y}{x^3+1}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{4xy'}{x^3+1} + \frac{y}{x^3+1} = 0$$

- Check to see if x_0 is a regular singular point

- Define functions

$$[P_2(x) = \frac{4x}{x^3+1}, P_3(x) = \frac{1}{x^3+1}]$$

- $(x + 1) \cdot P_2(x)$ is analytic at $x = -1$

$$((x + 1) \cdot P_2(x)) \Big|_{x=-1} = -\frac{4}{3}$$

- $(x + 1)^2 \cdot P_3(x)$ is analytic at $x = -1$

$$((x + 1)^2 \cdot P_3(x)) \Big|_{x=-1} = 0$$

- $x = -1$ is a regular singular point

Check to see if x_0 is a regular singular point

$$x_0 = -1$$

- Multiply by denominators

$$(x^3 + 1)y'' + 4y'x + y = 0$$

- Change variables using $x = u - 1$ so that the regular singular point is at $u = 0$

$$(u^3 - 3u^2 + 3u) \left(\frac{d^2}{du^2} y(u) \right) + (4u - 4) \left(\frac{d}{du} y(u) \right) + y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k+1-m$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k+1-m+r) u^{k+r}$$

- Convert $u^m \cdot \left(\frac{d^2}{du^2} y(u) \right)$ to series expansion for $m = 1..3$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=0}^{\infty} a_k (k+r) (k+r-1) u^{k+r-2+m}$$

- Shift index using $k \rightarrow k+2-m$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=-2+m}^{\infty} a_{k+2-m} (k+2-m+r) (k+1-m+r) u^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r(-7+3r) u^{-1+r} + (a_1(1+r)(-4+3r) - a_0(3r^2 - 7r - 1)) u^r + \left(\sum_{k=1}^{\infty} (a_{k+1}(k+1+r)(3k - \dots) \right)$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$r(-7+3r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \left\{ 0, \frac{7}{3} \right\}$$

- Each term must be 0

$$a_1(1+r)(-4+3r) - a_0(3r^2 - 7r - 1) = 0$$

- Each term in the series must be 0, giving the recursion relation

$$(-3a_k + a_{k-1} + 3a_{k+1})k^2 + ((-6a_k + 2a_{k-1} + 6a_{k+1})r + 7a_k - 3a_{k-1} - a_{k+1})k + (-3a_k + a_{k-1} - a_{k+1}) = 0$$

- Shift index using $k \rightarrow k+1$

$$(-3a_{k+1} + a_k + 3a_{k+2})(k+1)^2 + ((-6a_{k+1} + 2a_k + 6a_{k+2})r + 7a_{k+1} - 3a_k - a_{k+2})(k+1) + (-3a_{k+1} + a_k - a_{k+2}) = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} + 2k r a_k - 6k r a_{k+1} + r^2 a_k - 3r^2 a_{k+1} - k a_k + k a_{k+1} - r a_k + r a_{k+1} + 5a_{k+1}}{3k^2 + 6kr + 3r^2 + 5k + 5r - 2}$$

- Recursion relation for $r = 0$

$$a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} - k a_k + k a_{k+1} + 5a_{k+1}}{3k^2 + 5k - 2}$$

- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} - k a_k + k a_{k+1} + 5a_{k+1}}{3k^2 + 5k - 2}, -4a_1 + a_0 = 0 \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^k, a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} - k a_k + k a_{k+1} + 5a_{k+1}}{3k^2 + 5k - 2}, -4a_1 + a_0 = 0 \right]$$

- Recursion relation for $r = \frac{7}{3}$

$$a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} + \frac{11}{3} k a_k - 13k a_{k+1} + \frac{28}{9} a_k - 9a_{k+1}}{3k^2 + 19k + 26}$$

- Solution for $r = \frac{7}{3}$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+\frac{7}{3}}, a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} + \frac{11}{3} k a_k - 13k a_{k+1} + \frac{28}{9} a_k - 9a_{k+1}}{3k^2 + 19k + 26}, 10a_1 + a_0 = 0 \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^{k+\frac{7}{3}}, a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} + \frac{11}{3} k a_k - 13k a_{k+1} + \frac{28}{9} a_k - 9a_{k+1}}{3k^2 + 19k + 26}, 10a_1 + a_0 = 0 \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k (x+1)^k \right) + \left(\sum_{k=0}^{\infty} b_k (x+1)^{k+\frac{7}{3}} \right), a_{k+2} = -\frac{k^2 a_k - 3k^2 a_{k+1} - k a_k + k a_{k+1} + 5a_{k+1}}{3k^2 + 5k - 2}, -4a_1 + a_0 \right]$$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
  -> Bessel
  -> elliptic
  -> Legendre
  -> Kummer
    -> hyper3: Equivalence to 1F1 under a power @ Moebius
  -> hypergeometric
    -> heuristic approach
    -> hyper3: Equivalence to 2F1, 1F1 or 0F1 under a power @ Moebius
  -> Mathieu
    -> Equivalence to the rational form of Mathieu ODE under a power @ Moebius
trying a solution in terms of MeijerG functions
-> Heun: Equivalence to the GHE or one of its 4 confluent cases under a power @ Moebius
<- Heun successful: received ODE is equivalent to the HeunG ODE, case a <> 0, e <> 0, g <
```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 54

```
Order:=6;
dsolve((1+x^3)*diff(y(x),x$2)+4*x*diff(y(x),x)+y(x)=0,y(x),type='series',x=2);
```

$$y(x) = \left(1 - \frac{(-2+x)^2}{18} + \frac{10(-2+x)^3}{243} - \frac{451(-2+x)^4}{17496} + \frac{1151(-2+x)^5}{78732} \right) y(2) \\ + \left(-2+x - \frac{4(-2+x)^2}{9} + \frac{115(-2+x)^3}{486} - \frac{271(-2+x)^4}{2187} \right. \\ \left. + \frac{9713(-2+x)^5}{157464} \right) D(y)(2) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 87

```
AsymptoticDSolveValue[(1+x^3)*y''[x]+4*x*y'[x]+y[x]==0,y[x],{x,2,5}]
```

$$y(x) \rightarrow c_1 \left(\frac{1151(x-2)^5}{78732} - \frac{451(x-2)^4}{17496} + \frac{10}{243}(x-2)^3 - \frac{1}{18}(x-2)^2 + 1 \right) \\ + c_2 \left(\frac{9713(x-2)^5}{157464} - \frac{271(x-2)^4}{2187} + \frac{115}{486}(x-2)^3 - \frac{4}{9}(x-2)^2 + x - 2 \right)$$

14.12 problem 8

14.12.1 Maple step by step solution 3466

Internal problem ID [744]

Internal file name [OUTPUT/744_Sunday_June_05_2022_01_48_34_AM_11668954/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 8.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

```
[[_Emden , _Fowler]]
```

$$xy'' + y = 0$$

With the expansion point for the power series method at $x = 1$.

The ode does not have its expansion point at $x = 0$, therefore to simplify the computation of power series expansion, change of variable is made on the independent variable to shift the initial conditions and the expansion point back to zero. The new ode is then solved more easily since the expansion point is now at zero. The solution converted back to the original independent variable. Let

$$t = x - 1$$

The ode is converted to be in terms of the new independent variable t . This results in

$$\left(\frac{d^2}{dt^2}y(t)\right)(t+1) + y(t) = 0$$

With its expansion point and initial conditions now at $t = 0$. The transformed ODE is now solved. Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \cdots \\ &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (843)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (844)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$\begin{aligned}
 F_0 &= -\frac{y(t)}{t+1} \\
 F_1 &= \frac{dF_0}{dt} \\
 &= \frac{\partial F_0}{\partial t} + \frac{\partial F_0}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_0}{\partial \frac{d}{dt}y(t)} F_0 \\
 &= \frac{(-t-1) \left(\frac{d}{dt}y(t)\right) + y(t)}{(t+1)^2} \\
 F_2 &= \frac{dF_1}{dt} \\
 &= \frac{\partial F_1}{\partial t} + \frac{\partial F_1}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_1}{\partial \frac{d}{dt}y(t)} F_1 \\
 &= \frac{(2t+2) \left(\frac{d}{dt}y(t)\right) + y(t)(-1+t)}{(t+1)^3} \\
 F_3 &= \frac{dF_2}{dt} \\
 &= \frac{\partial F_2}{\partial t} + \frac{\partial F_2}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_2}{\partial \frac{d}{dt}y(t)} F_2 \\
 &= \frac{(t^2-4t-5) \left(\frac{d}{dt}y(t)\right) + (-4t+2)y(t)}{(t+1)^4} \\
 F_4 &= \frac{dF_3}{dt} \\
 &= \frac{\partial F_3}{\partial t} + \frac{\partial F_3}{\partial y} \frac{d}{dt}y(t) + \frac{\partial F_3}{\partial \frac{d}{dt}y(t)} F_3 \\
 &= \frac{(-6t^2+12t+18) \left(\frac{d}{dt}y(t)\right) - y(t)(t^2-16t+7)}{(t+1)^5}
 \end{aligned}$$

And so on. Evaluating all the above at initial conditions $t = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned}
 F_0 &= -y(0) \\
 F_1 &= -y'(0) + y(0) \\
 F_2 &= -y(0) + 2y'(0) \\
 F_3 &= 2y(0) - 5y'(0) \\
 F_4 &= -7y(0) + 18y'(0)
 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y(t) = \left(1 - \frac{1}{2}t^2 + \frac{1}{6}t^3 - \frac{1}{24}t^4 + \frac{1}{60}t^5 - \frac{7}{720}t^6\right) y(0) \\ + \left(t - \frac{1}{6}t^3 + \frac{1}{12}t^4 - \frac{1}{24}t^5 + \frac{1}{40}t^6\right) y'(0) + O(t^6)$$

Since the expansion point $t = 0$ is an ordinary, we can also solve this using standard power series. The ode is normalized to be

$$\left(\frac{d^2}{dt^2}y(t)\right)(t+1) + y(t) = 0$$

Let the solution be represented as power series of the form

$$y(t) = \sum_{n=0}^{\infty} a_n t^n$$

Then

$$\frac{d}{dt}y(t) = \sum_{n=1}^{\infty} n a_n t^{n-1} \\ \frac{d^2}{dt^2}y(t) = \sum_{n=2}^{\infty} n(n-1) a_n t^{n-2}$$

Substituting the above back into the ode gives

$$\left(\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2}\right)(t+1) + \left(\sum_{n=0}^{\infty} a_n t^n\right) = 0 \quad (1)$$

Which simplifies to

$$\left(\sum_{n=2}^{\infty} n t^{n-1} a_n (n-1)\right) + \left(\sum_{n=2}^{\infty} n(n-1) a_n t^{n-2}\right) + \left(\sum_{n=0}^{\infty} a_n t^n\right) = 0 \quad (2)$$

The next step is to make all powers of t be n in each summation term. Going over each summation term above with power of t in it which is not already t^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n t^{n-1} a_n (n-1) = \sum_{n=1}^{\infty} (n+1) a_{n+1} n t^n \\ \sum_{n=2}^{\infty} n(n-1) a_n t^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) t^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of t are the same and equal to n .

$$\left(\sum_{n=1}^{\infty} (n+1) a_{n+1} n t^n \right) + \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) t^n \right) + \left(\sum_{n=0}^{\infty} a_n t^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$2a_2 + a_0 = 0$$

$$a_2 = -\frac{a_0}{2}$$

For $1 \leq n$, the recurrence equation is

$$(n+1) a_{n+1} n + (n+2) a_{n+2} (n+1) + a_n = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$\begin{aligned} a_{n+2} &= -\frac{n^2 a_{n+1} + n a_{n+1} + a_n}{(n+2)(n+1)} \\ (5) \quad &= -\frac{a_n}{(n+2)(n+1)} - \frac{(n^2 + n) a_{n+1}}{(n+2)(n+1)} \end{aligned}$$

For $n = 1$ the recurrence equation gives

$$2a_2 + 6a_3 + a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = \frac{a_0}{6} - \frac{a_1}{6}$$

For $n = 2$ the recurrence equation gives

$$6a_3 + 12a_4 + a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = -\frac{a_0}{24} + \frac{a_1}{12}$$

For $n = 3$ the recurrence equation gives

$$12a_4 + 20a_5 + a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_0}{60} - \frac{a_1}{24}$$

For $n = 4$ the recurrence equation gives

$$20a_5 + 30a_6 + a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{7a_0}{720} + \frac{a_1}{40}$$

For $n = 5$ the recurrence equation gives

$$30a_6 + 42a_7 + a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = \frac{11a_0}{1680} - \frac{17a_1}{1008}$$

And so on. Therefore the solution is

$$\begin{aligned} y(t) &= \sum_{n=0}^{\infty} a_n t^n \\ &= a_3 t^3 + a_2 t^2 + a_1 t + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y(t) = a_0 + a_1 t - \frac{a_0 t^2}{2} + \left(\frac{a_0}{6} - \frac{a_1}{6}\right) t^3 + \left(-\frac{a_0}{24} + \frac{a_1}{12}\right) t^4 + \left(\frac{a_0}{60} - \frac{a_1}{24}\right) t^5 + \dots$$

Collecting terms, the solution becomes

$$y(t) = \left(1 - \frac{1}{2}t^2 + \frac{1}{6}t^3 - \frac{1}{24}t^4 + \frac{1}{60}t^5\right) a_0 + \left(t - \frac{1}{6}t^3 + \frac{1}{12}t^4 - \frac{1}{24}t^5\right) a_1 + O(t^6) \quad (3)$$

At $t = 0$ the solution above becomes

$$y(t) = \left(1 - \frac{1}{2}t^2 + \frac{1}{6}t^3 - \frac{1}{24}t^4 + \frac{1}{60}t^5\right) c_1 + \left(t - \frac{1}{6}t^3 + \frac{1}{12}t^4 - \frac{1}{24}t^5\right) c_2 + O(t^6)$$

Replacing t in the above with the original independent variable x using $t = x - 1$ results in

$$y = \left(1 - \frac{(x-1)^2}{2} + \frac{(x-1)^3}{6} - \frac{(x-1)^4}{24} + \frac{(x-1)^5}{60} - \frac{7(x-1)^6}{720}\right) y(1) \\ + \left(x-1 - \frac{(x-1)^3}{6} + \frac{(x-1)^4}{12} - \frac{(x-1)^5}{24} + \frac{(x-1)^6}{40}\right) y'(1) + O((x-1)^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{(x-1)^2}{2} + \frac{(x-1)^3}{6} - \frac{(x-1)^4}{24} + \frac{(x-1)^5}{60} - \frac{7(x-1)^6}{720}\right) y(1) \\ + \left(x-1 - \frac{(x-1)^3}{6} + \frac{(x-1)^4}{12} - \frac{(x-1)^5}{24} + \frac{(x-1)^6}{40}\right) y'(1) + O((x-1)^6) \quad (1)$$

Verification of solutions

$$y = \left(1 - \frac{(x-1)^2}{2} + \frac{(x-1)^3}{6} - \frac{(x-1)^4}{24} + \frac{(x-1)^5}{60} - \frac{7(x-1)^6}{720}\right) y(1) \\ + \left(x-1 - \frac{(x-1)^3}{6} + \frac{(x-1)^4}{12} - \frac{(x-1)^5}{24} + \frac{(x-1)^6}{40}\right) y'(1) + O((x-1)^6)$$

Verified OK.

14.12.1 Maple step by step solution

Let's solve

$$y''x + y = 0$$

- Highest derivative means the order of the ODE is 2
- y''
- Isolate 2nd derivative

$$y'' = -\frac{y}{x}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{y}{x} = 0$$

- Check to see if $x_0 = 0$ is a regular singular point

- Define functions

$$[P_2(x) = 0, P_3(x) = \frac{1}{x}]$$

- $x \cdot P_2(x)$ is analytic at $x = 0$

$$(x \cdot P_2(x)) \Big|_{x=0} = 0$$

- $x^2 \cdot P_3(x)$ is analytic at $x = 0$

$$(x^2 \cdot P_3(x)) \Big|_{x=0} = 0$$

- $x = 0$ is a regular singular point

Check to see if $x_0 = 0$ is a regular singular point

$$x_0 = 0$$

- Multiply by denominators

$$y''x + y = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^{k+r}$$

- Rewrite ODE with series expansions

- Convert $x \cdot y''$ to series expansion

$$x \cdot y'' = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) x^{k+r-1}$$

- Shift index using $k- > k+1$

$$x \cdot y'' = \sum_{k=-1}^{\infty} a_{k+1} (k+1+r)(k+r) x^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r(-1+r) x^{-1+r} + \left(\sum_{k=0}^{\infty} (a_{k+1} (k+1+r)(k+r) + a_k) x^{k+r} \right) = 0$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$r(-1+r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \{0, 1\}$$

- Each term in the series must be 0, giving the recursion relation

$$a_{k+1}(k+1+r)(k+r) + a_k = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+1} = -\frac{a_k}{(k+1+r)(k+r)}$$

- Recursion relation for $r = 0$

$$a_{k+1} = -\frac{a_k}{(k+1)k}$$

- Solution for $r = 0$

$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+1} = -\frac{a_k}{(k+1)k} \right]$$

- Recursion relation for $r = 1$

$$a_{k+1} = -\frac{a_k}{(k+2)(k+1)}$$

- Solution for $r = 1$

$$\left[y = \sum_{k=0}^{\infty} a_k x^{k+1}, a_{k+1} = -\frac{a_k}{(k+2)(k+1)} \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k x^k \right) + \left(\sum_{k=0}^{\infty} b_k x^{k+1} \right), a_{k+1} = -\frac{a_k}{(k+1)k}, b_{k+1} = -\frac{b_k}{(k+2)(k+1)} \right]$$

Maple trace

```
`Methods for second order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
checking if the LODE has constant coefficients  
checking if the LODE is of Euler type  
trying a symmetry of the form [xi=0, eta=F(x)]  
checking if the LODE is missing y  
-> Trying a Liouvillian solution using Kovacic's algorithm  
<- No Liouvillian solutions exist  
-> Trying a solution in terms of special functions:  
  -> Bessel  
  <- Bessel successful  
<- special function solution successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 49

```
Order:=6;  
dsolve(x*diff(y(x),x$2)+y(x)=0,y(x),type='series',x=1);
```

$$y(x) = \left(1 - \frac{(x-1)^2}{2} + \frac{(x-1)^3}{6} - \frac{(x-1)^4}{24} + \frac{(x-1)^5}{60}\right) y(1) \\ + \left(x - 1 - \frac{(x-1)^3}{6} + \frac{(x-1)^4}{12} - \frac{(x-1)^5}{24}\right) D(y)(1) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 78

```
AsymptoticDSolveValue[x*y''[x]+y[x]==0,y[x],{x,1,5}]
```

$$y(x) \rightarrow c_1 \left(\frac{1}{60}(x-1)^5 - \frac{1}{24}(x-1)^4 + \frac{1}{6}(x-1)^3 - \frac{1}{2}(x-1)^2 + 1 \right) \\ + c_2 \left(-\frac{1}{24}(x-1)^5 + \frac{1}{12}(x-1)^4 - \frac{1}{6}(x-1)^3 + x - 1 \right)$$

14.13 problem 10

14.13.1 Maple step by step solution 3478

Internal problem ID [745]

Internal file name [OUTPUT/745_Sunday_June_05_2022_01_48_36_AM_47024459/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 10.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : "second order series method. Ordinary point", "second_order_change_of_variable_on_x_method_1", "second_order_change_of_variable_on_x_method_2", "second order series method. Taylor series method"

Maple gives the following as the ode type

```
[_Gegenbauer , [_2nd_order , _linear , ` _with_symmetry_[0,F(x)] `]]
```

$$(-x^2 + 1)y'' - y'x + \alpha^2 y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using

Taylor series gives

$$\begin{aligned}
 y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\
 &= y_0 + xy'_0 + \frac{x^2}{2}f|_{x_0, y_0, y'_0} + \frac{x^3}{3!}f'|_{x_0, y_0, y'_0} + \dots \\
 &= y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0}
 \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \tag{1}$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \tag{846}$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \tag{847}$$

$$\begin{aligned}
 \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\
 &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \tag{2}
 \end{aligned}$$

$$\begin{aligned}
 \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\
 &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \tag{3}
 \end{aligned}$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \tag{4}$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}y' + \frac{\partial f}{\partial y'}f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y}y' + \frac{\partial F_0}{\partial y'}F_0 \end{aligned} \tag{5}$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \end{aligned}$$

⋮

$$\begin{aligned} F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \tag{6}$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \tag{7}$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = \frac{\alpha^2 y - y' x}{x^2 - 1}$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dx} \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\ &= \frac{((\alpha^2 + 2)x^2 - \alpha^2 + 1)y' - 3y\alpha^2 x}{(x^2 - 1)^2} \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dx} \\ &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\ &= \frac{(-6\alpha^2 x^3 + 6\alpha^2 x - 6x^3 - 9x)y' + ((\alpha^2 + 11)x^2 - \alpha^2 + 4)\alpha^2 y}{(x^2 - 1)^3} \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dx} \\ &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\ &= \frac{((\alpha^4 + 35\alpha^2 + 24)x^4 + (-2\alpha^4 - 25\alpha^2 + 72)x^2 + \alpha^4 - 10\alpha^2 + 9)y' - 10((\alpha^2 + 5)x^2 - \alpha^2 + \frac{11}{2})x\alpha^2 y}{(x^2 - 1)^4} \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dx} \\ &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\ &= \frac{15((\alpha^4 + 15\alpha^2 + 8)x^4 + (-2\alpha^4 - 2\alpha^2 + 40)x^2 + \alpha^4 - 13\alpha^2 + 15)xy' - ((\alpha^4 + 85\alpha^2 + 274)x^4 + (-2\alpha^4 - 6\alpha^2 + 15)\alpha^2 y)}{(x^2 - 1)^6} \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$\begin{aligned} F_0 &= -y(0)\alpha^2 \\ F_1 &= -y'(0)\alpha^2 + y'(0) \\ F_2 &= y(0)\alpha^4 - 4y(0)\alpha^2 \\ F_3 &= y'(0)\alpha^4 - 10y'(0)\alpha^2 + 9y'(0) \\ F_4 &= -y(0)\alpha^6 + 20y(0)\alpha^4 - 64y(0)\alpha^2 \end{aligned}$$

Substituting all the above in (7) and simplifying gives the solution as

$$y = \left(1 - \frac{1}{2}\alpha^2 x^2 + \frac{1}{24}\alpha^4 x^4 - \frac{1}{6}\alpha^2 x^4 - \frac{1}{720}x^6 \alpha^6 + \frac{1}{36}x^6 \alpha^4 - \frac{4}{45}x^6 \alpha^2\right) y(0) \\ + \left(x - \frac{1}{6}\alpha^2 x^3 + \frac{1}{6}x^3 + \frac{1}{120}\alpha^4 x^5 - \frac{1}{12}\alpha^2 x^5 + \frac{3}{40}x^5\right) y'(0) + O(x^6)$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series The ode is normalized to be

$$(-x^2 + 1) y'' - y'x + \alpha^2 y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1} \\ y'' = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$$

Substituting the above back into the ode gives

$$(-x^2 + 1) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) - \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + \alpha^2 \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\sum_{n=2}^{\infty} (-x^n a_n n(n-1)) + \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) + \sum_{n=1}^{\infty} (-n a_n x^n) + \left(\sum_{n=0}^{\infty} \alpha^2 a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} \sum_{n=2}^{\infty} (-x^n a_n n(n-1)) + \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) \\ + \sum_{n=1}^{\infty} (-n a_n x^n) + \left(\sum_{n=0}^{\infty} \alpha^2 a_n x^n \right) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$a_0 \alpha^2 + 2a_2 = 0$$

$$a_2 = -\frac{a_0 \alpha^2}{2}$$

$n = 1$ gives

$$a_1 \alpha^2 - a_1 + 6a_3 = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{1}{6} a_1 \alpha^2 + \frac{1}{6} a_1$$

For $2 \leq n$, the recurrence equation is

$$-n a_n (n-1) + (n+2) a_{n+2} (n+1) - n a_n + a_n \alpha^2 = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n (\alpha^2 - n^2)}{(n+2)(n+1)} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$a_2 \alpha^2 - 4a_2 + 12a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{1}{24} \alpha^4 a_0 - \frac{1}{6} a_0 \alpha^2$$

For $n = 3$ the recurrence equation gives

$$a_3\alpha^2 - 9a_3 + 20a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{1}{120}\alpha^4 a_1 - \frac{1}{12}a_1\alpha^2 + \frac{3}{40}a_1$$

For $n = 4$ the recurrence equation gives

$$a_4\alpha^2 - 16a_4 + 30a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{1}{720}\alpha^6 a_0 + \frac{1}{36}\alpha^4 a_0 - \frac{4}{45}a_0\alpha^2$$

For $n = 5$ the recurrence equation gives

$$a_5\alpha^2 - 25a_5 + 42a_7 = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{1}{5040}\alpha^6 a_1 + \frac{1}{144}\alpha^4 a_1 - \frac{37}{720}a_1\alpha^2 + \frac{5}{112}a_1$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$\begin{aligned} y &= a_0 + a_1 x - \frac{a_0\alpha^2 x^2}{2} + \left(-\frac{1}{6}a_1\alpha^2 + \frac{1}{6}a_1\right) x^3 \\ &\quad + \left(\frac{1}{24}\alpha^4 a_0 - \frac{1}{6}a_0\alpha^2\right) x^4 + \left(\frac{1}{120}\alpha^4 a_1 - \frac{1}{12}a_1\alpha^2 + \frac{3}{40}a_1\right) x^5 + \dots \end{aligned}$$

Collecting terms, the solution becomes

$$y = \left(1 - \frac{\alpha^2 x^2}{2} + \left(\frac{1}{24}\alpha^4 - \frac{1}{6}\alpha^2\right)x^4\right) a_0 + \left(x + \left(-\frac{\alpha^2}{6} + \frac{1}{6}\right)x^3 + \left(\frac{1}{120}\alpha^4 - \frac{1}{12}\alpha^2 + \frac{3}{40}\right)x^5\right) a_1 + O(x^6) \quad (3)$$

At $x = 0$ the solution above becomes

$$y = \left(1 - \frac{\alpha^2 x^2}{2} + \left(\frac{1}{24}\alpha^4 - \frac{1}{6}\alpha^2\right)x^4\right) c_1 + \left(x + \left(-\frac{\alpha^2}{6} + \frac{1}{6}\right)x^3 + \left(\frac{1}{120}\alpha^4 - \frac{1}{12}\alpha^2 + \frac{3}{40}\right)x^5\right) c_2 + O(x^6)$$

Summary

The solution(s) found are the following

$$y = \left(1 - \frac{1}{2}\alpha^2 x^2 + \frac{1}{24}\alpha^4 x^4 - \frac{1}{6}\alpha^2 x^4 - \frac{1}{720}x^6 \alpha^6 + \frac{1}{36}x^6 \alpha^4 - \frac{4}{45}x^6 \alpha^2\right) y(0) + \left(x - \frac{1}{6}\alpha^2 x^3 + \frac{1}{6}x^3 + \frac{1}{120}\alpha^4 x^5 - \frac{1}{12}\alpha^2 x^5 + \frac{3}{40}x^5\right) y'(0) + O(x^6) \quad (1)$$

$$y = \left(1 - \frac{\alpha^2 x^2}{2} + \left(\frac{1}{24}\alpha^4 - \frac{1}{6}\alpha^2\right)x^4\right) c_1 + \left(x + \left(-\frac{\alpha^2}{6} + \frac{1}{6}\right)x^3 + \left(\frac{1}{120}\alpha^4 - \frac{1}{12}\alpha^2 + \frac{3}{40}\right)x^5\right) c_2 + O(x^6) \quad (2)$$

Verification of solutions

$$y = \left(1 - \frac{1}{2}\alpha^2 x^2 + \frac{1}{24}\alpha^4 x^4 - \frac{1}{6}\alpha^2 x^4 - \frac{1}{720}x^6 \alpha^6 + \frac{1}{36}x^6 \alpha^4 - \frac{4}{45}x^6 \alpha^2\right) y(0) + \left(x - \frac{1}{6}\alpha^2 x^3 + \frac{1}{6}x^3 + \frac{1}{120}\alpha^4 x^5 - \frac{1}{12}\alpha^2 x^5 + \frac{3}{40}x^5\right) y'(0) + O(x^6)$$

Verified OK.

$$y = \left(1 - \frac{\alpha^2 x^2}{2} + \left(\frac{1}{24}\alpha^4 - \frac{1}{6}\alpha^2\right)x^4\right) c_1 + \left(x + \left(-\frac{\alpha^2}{6} + \frac{1}{6}\right)x^3 + \left(\frac{1}{120}\alpha^4 - \frac{1}{12}\alpha^2 + \frac{3}{40}\right)x^5\right) c_2 + O(x^6)$$

Verified OK.

14.13.1 Maple step by step solution

Let's solve

$$(-x^2 + 1)y'' - y'x + \alpha^2 y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -\frac{xy'}{x^2-1} + \frac{\alpha^2 y}{x^2-1}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{xy'}{x^2-1} - \frac{\alpha^2 y}{x^2-1} = 0$$

- Check to see if x_0 is a regular singular point

- o Define functions

$$\left[P_2(x) = \frac{x}{x^2-1}, P_3(x) = -\frac{\alpha^2}{x^2-1} \right]$$

- o $(x+1) \cdot P_2(x)$ is analytic at $x = -1$

$$\left. ((x+1) \cdot P_2(x)) \right|_{x=-1} = \frac{1}{2}$$

- o $(x+1)^2 \cdot P_3(x)$ is analytic at $x = -1$

$$\left. ((x+1)^2 \cdot P_3(x)) \right|_{x=-1} = 0$$

- o $x = -1$ is a regular singular point

Check to see if x_0 is a regular singular point

$$x_0 = -1$$

- Multiply by denominators

$$(x^2 - 1)y'' + y'x - \alpha^2 y = 0$$

- Change variables using $x = u - 1$ so that the regular singular point is at $u = 0$

$$(u^2 - 2u) \left(\frac{d^2}{du^2} y(u) \right) + (u - 1) \left(\frac{d}{du} y(u) \right) - \alpha^2 y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du}y(u)\right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du}y(u)\right) = \sum_{k=0}^{\infty} a_k(k+r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k+1-m$

$$u^m \cdot \left(\frac{d}{du}y(u)\right) = \sum_{k=-1+m}^{\infty} a_{k+1-m}(k+1-m+r) u^{k+r}$$

- Convert $u^m \cdot \left(\frac{d^2}{du^2}y(u)\right)$ to series expansion for $m = 1..2$

$$u^m \cdot \left(\frac{d^2}{du^2}y(u)\right) = \sum_{k=0}^{\infty} a_k(k+r)(k+r-1) u^{k+r-2+m}$$

- Shift index using $k \rightarrow k+2-m$

$$u^m \cdot \left(\frac{d^2}{du^2}y(u)\right) = \sum_{k=-2+m}^{\infty} a_{k+2-m}(k+2-m+r)(k+1-m+r) u^{k+r}$$

Rewrite ODE with series expansions

$$-a_0r(-1+2r)u^{-1+r} + \left(\sum_{k=0}^{\infty} (-a_{1+k}(1+k+r)(1+2k+2r) - a_k(\alpha+k+r)(\alpha-k-r)) u^{k+r}\right)$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$-r(-1+2r) = 0$$

- Values of r that satisfy the indicial equation

$$r \in \left\{0, \frac{1}{2}\right\}$$

- Each term in the series must be 0, giving the recursion relation

$$-2\left(\frac{1}{2} + k + r\right)(1+k+r)a_{1+k} + a_k(\alpha+k+r)(k+r-\alpha) = 0$$

- Recursion relation that defines series solution to ODE

$$a_{1+k} = -\frac{a_k(\alpha+k+r)(\alpha-k-r)}{(1+2k+2r)(1+k+r)}$$

- Recursion relation for $r = 0$

$$a_{1+k} = -\frac{a_k(\alpha+k)(\alpha-k)}{(1+2k)(1+k)}$$

- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{1+k} = -\frac{a_k(\alpha+k)(\alpha-k)}{(1+2k)(1+k)} \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^k, a_{1+k} = -\frac{a_k(\alpha+k)(\alpha-k)}{(1+2k)(1+k)} \right]$$

- Recursion relation for $r = \frac{1}{2}$

$$a_{1+k} = -\frac{a_k(\alpha+k+\frac{1}{2})(\alpha-k-\frac{1}{2})}{(2+2k)(\frac{3}{2}+k)}$$

- Solution for $r = \frac{1}{2}$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^{k+\frac{1}{2}}, a_{1+k} = -\frac{a_k(\alpha+k+\frac{1}{2})(\alpha-k-\frac{1}{2})}{(2+2k)(\frac{3}{2}+k)} \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^{k+\frac{1}{2}}, a_{1+k} = -\frac{a_k(\alpha+k+\frac{1}{2})(\alpha-k-\frac{1}{2})}{(2+2k)(\frac{3}{2}+k)} \right]$$

- Combine solutions and rename parameters

$$\left[y = \left(\sum_{k=0}^{\infty} a_k (x+1)^k \right) + \left(\sum_{k=0}^{\infty} b_k (x+1)^{k+\frac{1}{2}} \right), a_{k+1} = -\frac{a_k(\alpha+k)(\alpha-k)}{(1+2k)(k+1)}, b_{k+1} = -\frac{b_k(\alpha+k+\frac{1}{2})(\alpha-k-\frac{1}{2})}{(2+2k)(\frac{3}{2}+k)} \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
<- linear_1 successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 71

```

Order:=6;
dsolve((1-x^2)*diff(y(x),x$2)-x*diff(y(x),x)+alpha^2*y(x)=0,y(x),type='series',x=0);

```

$$y(x) = \left(1 - \frac{\alpha^2 x^2}{2} + \frac{\alpha^2(\alpha^2 - 4)x^4}{24} \right) y(0) + \left(x - \frac{(\alpha^2 - 1)x^3}{6} + \frac{(\alpha^4 - 10\alpha^2 + 9)x^5}{120} \right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 88

```
AsymptoticDSolveValue[(1-x^2)*y''[x]-x*y'[x]+\[Alpha]^2*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(\frac{\alpha^4 x^5}{120} - \frac{\alpha^2 x^5}{12} + \frac{3x^5}{40} - \frac{\alpha^2 x^3}{6} + \frac{x^3}{6} + x \right) + c_1 \left(\frac{\alpha^4 x^4}{24} - \frac{\alpha^2 x^4}{6} - \frac{\alpha^2 x^2}{2} + 1 \right)$$

14.14 problem 16

14.14.1 Solving as series ode	3482
14.14.2 Maple step by step solution	3489

Internal problem ID [746]

Internal file name [OUTPUT/746_Sunday_June_05_2022_01_48_37_AM_93764472/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 16.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"first order ode series method. Ordinary point", "first order ode series method. Taylor series method"**

Maple gives the following as the ode type

`[_quadrature]`

$$y' - y = 0$$

With the expansion point for the power series method at $x = 0$.

14.14.1 Solving as series ode

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving first order ode. Let

$$y' = f(x, y)$$

Where $f(x, y)$ is analytic at expansion point x_0 . We can always shift to $x_0 = 0$ if x_0 is not zero. So from now we assume $x_0 = 0$. Assume also that $y(x_0) = y_0$. Using Taylor series

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0) y'(x_0) + \frac{(x - x_0)^2}{2} y''(x_0) + \frac{(x - x_0)^3}{3!} y'''(x_0) + \dots \\ &= y_0 + x f + \frac{x^2}{2} \frac{df}{dx} \Big|_{x_0, y_0} + \frac{x^3}{3!} \frac{d^2 f}{dx^2} \Big|_{x_0, y_0} + \dots \\ &= y_0 + \sum_{n=0}^{\infty} \frac{x^{n+1}}{(n+1)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \quad (1)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y)$ then the above can be written as

$$F_0 = f(x, y) \quad (4)$$

$$\begin{aligned} F_n &= \frac{d}{dx}(F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) F_0 \end{aligned} \quad (5)$$

For example, for $n = 1$ we see that

$$\begin{aligned} F_1 &= \frac{d}{dx}(F_0) \\ &= \frac{\partial}{\partial x} F_0 + \left(\frac{\partial F_0}{\partial y} \right) F_0 \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \end{aligned}$$

Which is (1). And when $n = 2$

$$\begin{aligned} F_2 &= \frac{d}{dx}(F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) F_0 \\ &= \frac{\partial}{\partial x} \left(\frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \right) + \frac{\partial}{\partial y} \left(\frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \right) f \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) f \end{aligned}$$

Which is (2) and so on. Therefore (4,5) can be used from now on along with

$$y(x) = y_0 + \sum_{n=0}^{\infty} \frac{x^{n+1}}{(n+1)!} F_n|_{x_0, y_0} \quad (6)$$

Hence

$$\begin{aligned} F_0 &= y \\ F_1 &= \frac{dF_0}{dx} \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} F_0 \\ &= y \\ F_2 &= \frac{dF_1}{dx} \\ &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} F_1 \\ &= y \\ F_3 &= \frac{dF_2}{dx} \\ &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} F_2 \\ &= y \\ F_4 &= \frac{dF_3}{dx} \\ &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} F_3 \\ &= y \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x(0) = 0$ and $y(0) = y(0)$ gives

$$\begin{aligned} F_0 &= y(0) \\ F_1 &= y(0) \\ F_2 &= y(0) \\ F_3 &= y(0) \\ F_4 &= y(0) \end{aligned}$$

Substituting all the above in (6) and simplifying gives the solution as

$$y = \left(1 + x + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5 \right) y(0) + O(x^6)$$

Since $x = 0$ is also an ordinary point, then standard power series can also be used. Writing the ODE as

$$\begin{aligned}y' + q(x)y &= p(x) \\y' - y &= 0\end{aligned}$$

Where

$$\begin{aligned}q(x) &= -1 \\p(x) &= 0\end{aligned}$$

Next, the type of the expansion point $x = 0$ is determined. This point can be an ordinary point, a regular singular point (also called removable singularity), or irregular singular point (also called non-removable singularity or essential singularity). When $x = 0$ is an ordinary point, then the standard power series is used. If the point is a regular singular point, Frobenius series is used instead. Irregular singular point requires more advanced methods (asymptotic methods) and is not supported now. Hopefully this will be added in the future. $x = 0$ is called an ordinary point $q(x)$ has a Taylor series expansion around the point $x = 0$. $x = 0$ is called a regular singular point if $q(x)$ is not analytic at $x = 0$ but $xq(x)$ has Taylor series expansion. And finally, $x = 0$ is an irregular singular point if the point is not ordinary and not regular singular. This is the most complicated case. Now the expansion point $x = 0$ is checked to see if it is an ordinary point or not. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

Substituting the above back into the ode gives

$$\left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) - \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=1}^{\infty} n a_n x^{n-1} = \sum_{n=0}^{\infty} (n+1) a_{n+1} x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (n+1) a_{n+1} x^n \right) + \sum_{n=0}^{\infty} (-a_n x^n) = 0 \quad (3)$$

For $0 \leq n$, the recurrence equation is

$$(n+1) a_{n+1} - a_n = 0 \quad (4)$$

Solving for a_{n+1} , gives

$$a_{n+1} = \frac{a_n}{n+1} \quad (5)$$

For $n = 0$ the recurrence equation gives

$$a_1 - a_0 = 0$$

Which after substituting the earlier terms found becomes

$$a_1 = a_0$$

For $n = 1$ the recurrence equation gives

$$2a_2 - a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_2 = \frac{a_0}{2}$$

For $n = 2$ the recurrence equation gives

$$3a_3 - a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = \frac{a_0}{6}$$

For $n = 3$ the recurrence equation gives

$$4a_4 - a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{24}$$

For $n = 4$ the recurrence equation gives

$$5a_5 - a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = \frac{a_0}{120}$$

For $n = 5$ the recurrence equation gives

$$6a_6 - a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{a_0}{720}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + a_0 x + \frac{1}{2} a_0 x^2 + \frac{1}{6} a_0 x^3 + \frac{1}{24} a_0 x^4 + \frac{1}{120} a_0 x^5 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + x + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) a_0 + O(x^6) \quad (3)$$

Summary

The solution(s) found are the following

$$y = \left(1 + x + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) y(0) + O(x^6) \quad (1)$$

$$y = \left(1 + x + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) c_1 + O(x^6) \quad (2)$$

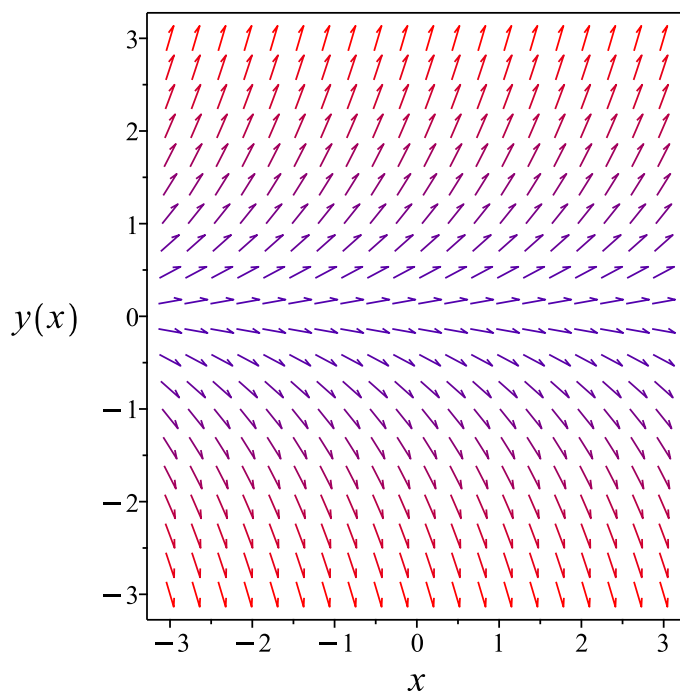


Figure 503: Slope field plot

Verification of solutions

$$y = \left(1 + x + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) y(0) + O(x^6)$$

Verified OK.

$$y = \left(1 + x + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) c_1 + O(x^6)$$

Verified OK.

14.14.2 Maple step by step solution

Let's solve

$$y' - y = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{y} = 1$$

- Integrate both sides with respect to x

$$\int \frac{y'}{y} dx = \int 1 dx + c_1$$

- Evaluate integral

$$\ln(y) = x + c_1$$

- Solve for y

$$y = e^{x+c_1}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 34

```
Order:=6;  
dsolve(diff(y(x),x)-y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 + x + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \frac{1}{24}x^4 + \frac{1}{120}x^5\right) y(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 37

```
AsymptoticDSolveValue[y'[x]-y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_1 \left(\frac{x^5}{120} + \frac{x^4}{24} + \frac{x^3}{6} + \frac{x^2}{2} + x + 1 \right)$$

14.15 problem 17

14.15.1 Solving as series ode	3491
14.15.2 Maple step by step solution	3498

Internal problem ID [747]

Internal file name [OUTPUT/747_Sunday_June_05_2022_01_48_38_AM_46508026/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 17.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"first order ode series method. Ordinary point", "first order ode series method. Taylor series method"**

Maple gives the following as the ode type

`[_separable]`

$$-yx + y' = 0$$

With the expansion point for the power series method at $x = 0$.

14.15.1 Solving as series ode

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving first order ode. Let

$$y' = f(x, y)$$

Where $f(x, y)$ is analytic at expansion point x_0 . We can always shift to $x_0 = 0$ if x_0 is not zero. So from now we assume $x_0 = 0$. Assume also that $y(x_0) = y_0$. Using Taylor series

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\ &= y_0 + xf + \frac{x^2}{2} \left. \frac{df}{dx} \right|_{x_0, y_0} + \frac{x^3}{3!} \left. \frac{d^2f}{dx^2} \right|_{x_0, y_0} + \dots \\ &= y_0 + \sum_{n=0}^{\infty} \frac{x^{n+1}}{(n+1)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \quad (1)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y)$ then the above can be written as

$$F_0 = f(x, y) \quad (4)$$

$$\begin{aligned} F_n &= \frac{d}{dx}(F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) F_0 \end{aligned} \quad (5)$$

For example, for $n = 1$ we see that

$$\begin{aligned} F_1 &= \frac{d}{dx}(F_0) \\ &= \frac{\partial}{\partial x} F_0 + \left(\frac{\partial F_0}{\partial y} \right) F_0 \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \end{aligned}$$

Which is (1). And when $n = 2$

$$\begin{aligned} F_2 &= \frac{d}{dx}(F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) F_0 \\ &= \frac{\partial}{\partial x} \left(\frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \right) + \frac{\partial}{\partial y} \left(\frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \right) f \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) f \end{aligned}$$

Which is (2) and so on. Therefore (4,5) can be used from now on along with

$$y(x) = y_0 + \sum_{n=0}^{\infty} \frac{x^{n+1}}{(n+1)!} F_n|_{x_0, y_0} \quad (6)$$

Hence

$$\begin{aligned} F_0 &= yx \\ F_1 &= \frac{dF_0}{dx} \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} F_0 \\ &= y(x^2 + 1) \\ F_2 &= \frac{dF_1}{dx} \\ &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} F_1 \\ &= yx(x^2 + 3) \\ F_3 &= \frac{dF_2}{dx} \\ &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} F_2 \\ &= y(x^4 + 6x^2 + 3) \\ F_4 &= \frac{dF_3}{dx} \\ &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} F_3 \\ &= yx(x^4 + 10x^2 + 15) \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x(0) = 0$ and $y(0) = y(0)$ gives

$$\begin{aligned} F_0 &= 0 \\ F_1 &= y(0) \\ F_2 &= 0 \\ F_3 &= 3y(0) \\ F_4 &= 0 \end{aligned}$$

Substituting all the above in (6) and simplifying gives the solution as

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) y(0) + O(x^6)$$

Since $x = 0$ is also an ordinary point, then standard power series can also be used. Writing the ODE as

$$\begin{aligned}y' + q(x)y &= p(x) \\ -yx + y' &= 0\end{aligned}$$

Where

$$\begin{aligned}q(x) &= -x \\ p(x) &= 0\end{aligned}$$

Next, the type of the expansion point $x = 0$ is determined. This point can be an ordinary point, a regular singular point (also called removable singularity), or irregular singular point (also called non-removable singularity or essential singularity). When $x = 0$ is an ordinary point, then the standard power series is used. If the point is a regular singular point, Frobenius series is used instead. Irregular singular point requires more advanced methods (asymptotic methods) and is not supported now. Hopefully this will be added in the future. $x = 0$ is called an ordinary point $q(x)$ has a Taylor series expansion around the point $x = 0$. $x = 0$ is called a regular singular point if $q(x)$ is not analytic at $x = 0$ but $xq(x)$ has Taylor series expansion. And finally, $x = 0$ is an irregular singular point if the point is not ordinary and not regular singular. This is the most complicated case. Now the expansion point $x = 0$ is checked to see if it is an ordinary point or not. Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

Substituting the above back into the ode gives

$$-\left(\sum_{n=0}^{\infty} a_n x^n\right) x + \left(\sum_{n=1}^{\infty} n a_n x^{n-1}\right) = 0 \quad (1)$$

Which simplifies to

$$\left(\sum_{n=1}^{\infty} n a_n x^{n-1}\right) + \sum_{n=0}^{\infty} (-x^{1+n} a_n) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=1}^{\infty} n a_n x^{n-1} = \sum_{n=0}^{\infty} (1+n) a_{1+n} x^n$$

$$\sum_{n=0}^{\infty} (-x^{1+n} a_n) = \sum_{n=1}^{\infty} (-a_{n-1} x^n)$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=0}^{\infty} (1+n) a_{1+n} x^n \right) + \sum_{n=1}^{\infty} (-a_{n-1} x^n) = 0 \quad (3)$$

For $1 \leq n$, the recurrence equation is

$$(1+n) a_{1+n} - a_{n-1} = 0 \quad (4)$$

Solving for a_{1+n} , gives

$$a_{1+n} = \frac{a_{n-1}}{1+n} \quad (5)$$

For $n = 1$ the recurrence equation gives

$$2a_2 - a_0 = 0$$

Which after substituting the earlier terms found becomes

$$a_2 = \frac{a_0}{2}$$

For $n = 2$ the recurrence equation gives

$$3a_3 - a_1 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = 0$$

For $n = 3$ the recurrence equation gives

$$4a_4 - a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = \frac{a_0}{8}$$

For $n = 4$ the recurrence equation gives

$$5a_5 - a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = 0$$

For $n = 5$ the recurrence equation gives

$$6a_6 - a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = \frac{a_0}{48}$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0 + \frac{1}{2}a_0 x^2 + \frac{1}{8}a_0 x^4 + \dots$$

Collecting terms, the solution becomes

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) a_0 + O(x^6) \quad (3)$$

Summary

The solution(s) found are the following

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) y(0) + O(x^6) \quad (1)$$

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) c_1 + O(x^6) \quad (2)$$

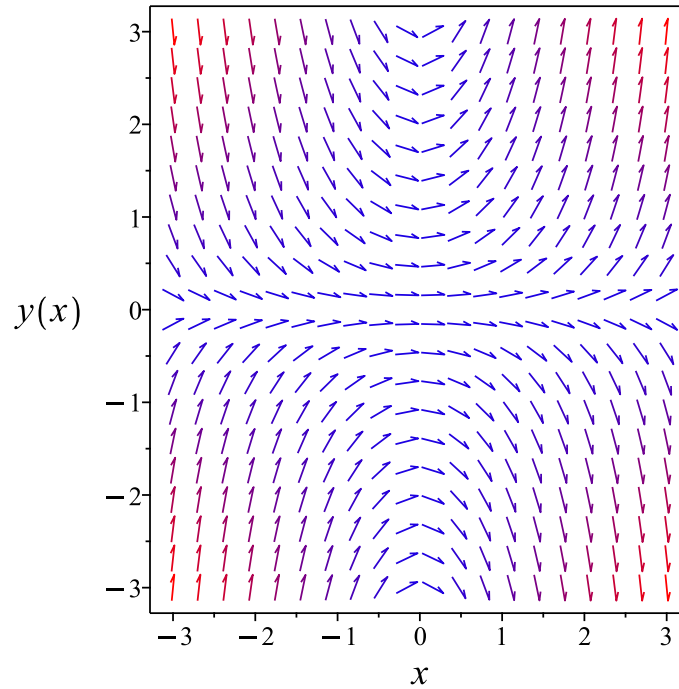


Figure 504: Slope field plot

Verification of solutions

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) y(0) + O(x^6)$$

Verified OK.

$$y = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) c_1 + O(x^6)$$

Verified OK.

14.15.2 Maple step by step solution

Let's solve

$$-yx + y' = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{y} = x$$

- Integrate both sides with respect to x

$$\int \frac{y'}{y} dx = \int x dx + c_1$$

- Evaluate integral

$$\ln(y) = \frac{x^2}{2} + c_1$$

- Solve for y

$$y = e^{\frac{x^2}{2} + c_1}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 21

```
Order:=6;  
dsolve(diff(y(x),x)-x*y(x)=0,y(x),type='series',x=0);
```

$$y(x) = \left(1 + \frac{1}{2}x^2 + \frac{1}{8}x^4\right) y(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 22

```
AsymptoticDSolveValue[y'[x]-x*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_1 \left(\frac{x^4}{8} + \frac{x^2}{2} + 1 \right)$$

14.16 problem 19

14.16.1 Solving as series ode	3500
14.16.2 Maple step by step solution	3507

Internal problem ID [748]

Internal file name [OUTPUT/748_Sunday_June_05_2022_01_48_39_AM_74598292/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 19.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"first order ode series method. Ordinary point", "first order ode series method. Taylor series method"**

Maple gives the following as the ode type

`[_separable]`

$$(1 - x)y' - y = 0$$

With the expansion point for the power series method at $x = 0$.

14.16.1 Solving as series ode

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving first order ode. Let

$$y' = f(x, y)$$

Where $f(x, y)$ is analytic at expansion point x_0 . We can always shift to $x_0 = 0$ if x_0 is not zero. So from now we assume $x_0 = 0$. Assume also that $y(x_0) = y_0$. Using Taylor series

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0)y'(x_0) + \frac{(x - x_0)^2}{2}y''(x_0) + \frac{(x - x_0)^3}{3!}y'''(x_0) + \dots \\ &= y_0 + xf + \frac{x^2}{2} \left. \frac{df}{dx} \right|_{x_0, y_0} + \frac{x^3}{3!} \left. \frac{d^2f}{dx^2} \right|_{x_0, y_0} + \dots \\ &= y_0 + \sum_{n=0}^{\infty} \frac{x^{n+1}}{(n+1)!} \left. \frac{d^n f}{dx^n} \right|_{x_0, y_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \quad (1)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y)$ then the above can be written as

$$F_0 = f(x, y) \quad (4)$$

$$\begin{aligned} F_n &= \frac{d}{dx}(F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) F_0 \end{aligned} \quad (5)$$

For example, for $n = 1$ we see that

$$\begin{aligned} F_1 &= \frac{d}{dx}(F_0) \\ &= \frac{\partial}{\partial x} F_0 + \left(\frac{\partial F_0}{\partial y} \right) F_0 \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \end{aligned}$$

Which is (1). And when $n = 2$

$$\begin{aligned} F_2 &= \frac{d}{dx}(F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) F_0 \\ &= \frac{\partial}{\partial x} \left(\frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \right) + \frac{\partial}{\partial y} \left(\frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} f \right) f \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) f \end{aligned}$$

Which is (2) and so on. Therefore (4,5) can be used from now on along with

$$y(x) = y_0 + \sum_{n=0}^{\infty} \frac{x^{n+1}}{(n+1)!} F_n|_{x_0, y_0} \quad (6)$$

Hence

$$\begin{aligned} F_0 &= -\frac{y}{x-1} \\ F_1 &= \frac{dF_0}{dx} \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} F_0 \\ &= \frac{2y}{(x-1)^2} \\ F_2 &= \frac{dF_1}{dx} \\ &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} F_1 \\ &= -\frac{6y}{(x-1)^3} \\ F_3 &= \frac{dF_2}{dx} \\ &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} F_2 \\ &= \frac{24y}{(x-1)^4} \\ F_4 &= \frac{dF_3}{dx} \\ &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} F_3 \\ &= -\frac{120y}{(x-1)^5} \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x(0) = 0$ and $y(0) = y(0)$ gives

$$\begin{aligned} F_0 &= y(0) \\ F_1 &= 2y(0) \\ F_2 &= 6y(0) \\ F_3 &= 24y(0) \\ F_4 &= 120y(0) \end{aligned}$$

Substituting all the above in (6) and simplifying gives the solution as

$$y = (x^5 + x^4 + x^3 + x^2 + x + 1) y(0) + O(x^6)$$

Since $x = 0$ is also an ordinary point, then standard power series can also be used. Writing the ODE as

$$\begin{aligned} y' + q(x)y &= p(x) \\ y' + \frac{y}{x-1} &= 0 \end{aligned}$$

Where

$$\begin{aligned} q(x) &= \frac{1}{x-1} \\ p(x) &= 0 \end{aligned}$$

Next, the type of the expansion point $x = 0$ is determined. This point can be an ordinary point, a regular singular point (also called removable singularity), or irregular singular point (also called non-removable singularity or essential singularity). When $x = 0$ is an ordinary point, then the standard power series is used. If the point is a regular singular point, Frobenius series is used instead. Irregular singular point requires more advanced methods (asymptotic methods) and is not supported now. Hopefully this will be added in the future. $x = 0$ is called an ordinary point if $q(x)$ has a Taylor series expansion around the point $x = 0$. $x = 0$ is called a regular singular point if $q(x)$ is not analytic at $x = 0$ but $xq(x)$ has Taylor series expansion. And finally, $x = 0$ is an irregular singular point if the point is not ordinary and not regular singular. This is the most complicated case. Now the expansion point $x = 0$ is checked to see if it is an ordinary point or not. Now the ode is normalized by writing it as

$$(x-1)y' + y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$y' = \sum_{n=1}^{\infty} n a_n x^{n-1}$$

Substituting the above back into the ode gives

$$(x - 1) \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\left(\sum_{n=1}^{\infty} n a_n x^n \right) + \sum_{n=1}^{\infty} (-n a_n x^{n-1}) + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=1}^{\infty} (-n a_n x^{n-1}) = \sum_{n=0}^{\infty} (-(n+1) a_{n+1} x^n)$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\left(\sum_{n=1}^{\infty} n a_n x^n \right) + \sum_{n=0}^{\infty} (-(n+1) a_{n+1} x^n) + \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (3)$$

$n = 0$ gives

$$\begin{aligned} -a_1 + a_0 &= 0 \\ a_1 &= a_0 \end{aligned}$$

For $1 \leq n$, the recurrence equation is

$$n a_n - (n+1) a_{n+1} + a_n = 0 \quad (4)$$

Solving for a_{n+1} , gives

$$a_{n+1} = a_n \quad (5)$$

For $n = 1$ the recurrence equation gives

$$2a_1 - 2a_2 = 0$$

Which after substituting the earlier terms found becomes

$$a_2 = a_0$$

For $n = 2$ the recurrence equation gives

$$3a_2 - 3a_3 = 0$$

Which after substituting the earlier terms found becomes

$$a_3 = a_0$$

For $n = 3$ the recurrence equation gives

$$4a_3 - 4a_4 = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = a_0$$

For $n = 4$ the recurrence equation gives

$$5a_4 - 5a_5 = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = a_0$$

For $n = 5$ the recurrence equation gives

$$6a_5 - 6a_6 = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = a_0$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$y = a_0x^5 + a_0x^4 + a_0x^3 + a_0x^2 + a_0x + a_0 + \dots$$

Collecting terms, the solution becomes

$$y = (x^5 + x^4 + x^3 + x^2 + x + 1) a_0 + O(x^6) \quad (3)$$

Summary

The solution(s) found are the following

$$y = (x^5 + x^4 + x^3 + x^2 + x + 1) y(0) + O(x^6) \quad (1)$$

$$y = (x^5 + x^4 + x^3 + x^2 + x + 1) c_1 + O(x^6) \quad (2)$$

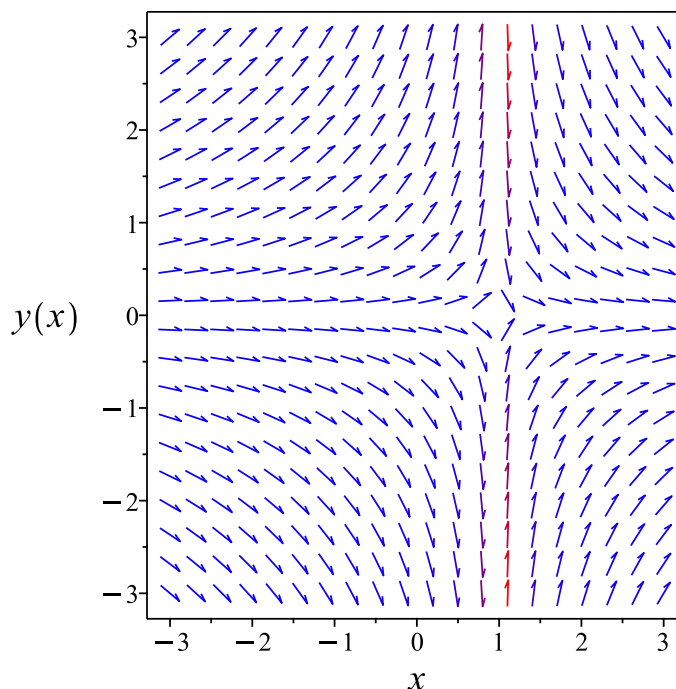


Figure 505: Slope field plot

Verification of solutions

$$y = (x^5 + x^4 + x^3 + x^2 + x + 1) y(0) + O(x^6)$$

Verified OK.

$$y = (x^5 + x^4 + x^3 + x^2 + x + 1) c_1 + O(x^6)$$

Verified OK.

14.16.2 Maple step by step solution

Let's solve

$$(x - 1)y' + y = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Integrate both sides with respect to x

$$\int ((x - 1)y' + y) dx = \int 0 dx + c_1$$

- Evaluate integral

$$y(x - 1) = c_1$$

- Solve for y

$$y = \frac{c_1}{x-1}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 26

```
Order:=6;  
dsolve((1-x)*diff(y(x),x)=y(x),y(x),type='series',x=0);
```

$$y(x) = (x^5 + x^4 + x^3 + x^2 + x + 1)y(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 21

```
AsymptoticDSolveValue[(1-x)*y'[x]==y[x],y[x],{x,0,5}]
```

$$y(x) \rightarrow c_1(x^5 + x^4 + x^3 + x^2 + x + 1)$$

14.17 problem 22

14.17.1 Maple step by step solution 3516

Internal problem ID [749]

Internal file name [OUTPUT/749_Sunday_June_05_2022_01_48_40_AM_29896868/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 5.3, Series Solutions Near an Ordinary Point, Part II. page 269

Problem number: 22.

ODE order: 2.

ODE degree: 1.

The type(s) of ODE detected by this program : **"second order series method. Ordinary point", "second order series method. Taylor series method"**

Maple gives the following as the ode type

[_Gegenbauer]

$$(-x^2 + 1) y'' - 2y'x + \alpha(\alpha + 1) y = 0$$

With the expansion point for the power series method at $x = 0$.

Solving ode using Taylor series method. This gives review on how the Taylor series method works for solving second order ode.

Let

$$y'' = f(x, y, y')$$

Assuming expansion is at $x_0 = 0$ (we can always shift the actual expansion point to 0 by change of variables) and assuming $f(x, y, y')$ is analytic at x_0 which must be the case for an ordinary point. Let initial conditions be $y(x_0) = y_0$ and $y'(x_0) = y'_0$. Using Taylor series gives

$$\begin{aligned} y(x) &= y(x_0) + (x - x_0) y'(x_0) + \frac{(x - x_0)^2}{2} y''(x_0) + \frac{(x - x_0)^3}{3!} y'''(x_0) + \cdots \\ &= y_0 + x y'_0 + \frac{x^2}{2} f|_{x_0, y_0, y'_0} + \frac{x^3}{3!} f'|_{x_0, y_0, y'_0} + \cdots \\ &= y_0 + x y'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} \frac{d^n f}{dx^n} \Big|_{x_0, y_0, y'_0} \end{aligned}$$

But

$$\frac{df}{dx} = \frac{\partial f}{\partial x} \frac{dx}{dx} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} \quad (1)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \quad (852)$$

$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \quad (853)$$

$$\begin{aligned} \frac{d^2 f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{df}{dx} \right) + \frac{\partial}{\partial y} \left(\frac{df}{dx} \right) y' + \frac{\partial}{\partial y'} \left(\frac{df}{dx} \right) f \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{d^3 f}{dx^3} &= \frac{d}{dx} \left(\frac{d^2 f}{dx^2} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{d^2 f}{dx^2} \right) + \left(\frac{\partial}{\partial y} \frac{d^2 f}{dx^2} \right) y' + \frac{\partial}{\partial y'} \left(\frac{d^2 f}{dx^2} \right) f \end{aligned} \quad (3)$$

⋮

And so on. Hence if we name $F_0 = f(x, y, y')$ then the above can be written as

$$F_0 = f(x, y, y') \quad (4)$$

$$\begin{aligned} F_1 &= \frac{df}{dx} \\ &= \frac{dF_0}{dx} \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} y'' \\ &= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} y' + \frac{\partial f}{\partial y'} f \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \end{aligned} \quad (5)$$

$$\begin{aligned} F_2 &= \frac{d}{dx} \left(\frac{d}{dx} f \right) \\ &= \frac{d}{dx} (F_1) \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_1 + \left(\frac{\partial F_1}{\partial y} \right) y' + \left(\frac{\partial F_1}{\partial y'} \right) F_0 \\ &\vdots \\ F_n &= \frac{d}{dx} (F_{n-1}) \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) y'' \\ &= \frac{\partial}{\partial x} F_{n-1} + \left(\frac{\partial F_{n-1}}{\partial y} \right) y' + \left(\frac{\partial F_{n-1}}{\partial y'} \right) F_0 \end{aligned} \quad (6)$$

Therefore (6) can be used from now on along with

$$y(x) = y_0 + xy'_0 + \sum_{n=0}^{\infty} \frac{x^{n+2}}{(n+2)!} F_n|_{x_0, y_0, y'_0} \quad (7)$$

To find $y(x)$ series solution around $x = 0$. Hence

$$F_0 = \frac{\alpha^2 y + \alpha y - 2y'x}{x^2 - 1}$$

$$\begin{aligned} F_1 &= \frac{dF_0}{dx} \\ &= \frac{\partial F_0}{\partial x} + \frac{\partial F_0}{\partial y} y' + \frac{\partial F_0}{\partial y'} F_0 \\ &= \frac{(\alpha^2 x^2 + \alpha x^2 - \alpha^2 + 6x^2 - \alpha + 2) y' - 4y\alpha x(\alpha + 1)}{(x^2 - 1)^2} \end{aligned}$$

$$\begin{aligned} F_2 &= \frac{dF_1}{dx} \\ &= \frac{\partial F_1}{\partial x} + \frac{\partial F_1}{\partial y} y' + \frac{\partial F_1}{\partial y'} F_1 \\ &= \frac{8(x+1)(x-1) \left(x((\alpha^2 + \alpha + 3)x^2 - \alpha^2 - \alpha + 3) y' - \frac{(\alpha+1)((\alpha^2 + \alpha + 18)x^2 - \alpha^2 - \alpha + 6)\alpha y}{8} \right)}{(x^2 - 1)^4} \end{aligned}$$

$$\begin{aligned} F_3 &= \frac{dF_2}{dx} \\ &= \frac{\partial F_2}{\partial x} + \frac{\partial F_2}{\partial y} y' + \frac{\partial F_2}{\partial y'} F_2 \\ &= \frac{(((\alpha^4 + 2\alpha^3 + 59\alpha^2 + 58\alpha + 120)x^4 + (-2\alpha^4 - 4\alpha^3 - 46\alpha^2 - 44\alpha + 240)x^2 + \alpha^4 + 2\alpha^3 - 13\alpha^2 - 13\alpha + 12)) y' + ((\alpha^4 + 2\alpha^3 - 13\alpha^2 - 13\alpha + 12)) y}{(x^2 - 1)^5} \end{aligned}$$

$$\begin{aligned} F_4 &= \frac{dF_3}{dx} \\ &= \frac{\partial F_3}{\partial x} + \frac{\partial F_3}{\partial y} y' + \frac{\partial F_3}{\partial y'} F_3 \\ &= \frac{18 \left(\left(40 + (x^4 - 2x^2 + 1)\alpha^4 + 2(x^4 - 2x^2 + 1)\alpha^3 + \left(-\frac{26}{3}x^2 - 17 + \frac{77}{3}x^4\right)\alpha^2 + 2\left(-9 - \frac{10}{3}x^2 + \frac{37}{3}x\right)\alpha \right) y' + \left((40 + (x^4 - 2x^2 + 1)\alpha^4 + 2(x^4 - 2x^2 + 1)\alpha^3 + \left(-\frac{26}{3}x^2 - 17 + \frac{77}{3}x^4\right)\alpha^2 + 2\left(-9 - \frac{10}{3}x^2 + \frac{37}{3}x\right)\alpha \right) y \right)}{(x^2 - 1)^6} \end{aligned}$$

And so on. Evaluating all the above at initial conditions $x = 0$ and $y(0) = y(0)$ and $y'(0) = y'(0)$ gives

$$F_0 = -y(0) \alpha(\alpha + 1)$$

$$F_1 = -y'(0) \alpha^2 - y'(0) \alpha + 2y'(0)$$

$$F_2 = y(0) \alpha^4 + 2y(0) \alpha^3 - 5y(0) \alpha^2 - 6y(0) \alpha$$

$$F_3 = y'(0) \alpha^4 + 2y'(0) \alpha^3 - 13y'(0) \alpha^2 - 14y'(0) \alpha + 24y'(0)$$

$$F_4 = -y(0) \alpha^6 - 3y(0) \alpha^5 + 23y(0) \alpha^4 + 51y(0) \alpha^3 - 94y(0) \alpha^2 - 120y(0) \alpha$$

Substituting all the above in (7) and simplifying gives the solution as

$$\begin{aligned}
 y = & \left(1 - \frac{1}{2}\alpha^2 x^2 - \frac{1}{2}\alpha x^2 + \frac{1}{24}\alpha^4 x^4 + \frac{1}{12}\alpha^3 x^4 - \frac{5}{24}\alpha^2 x^4 - \frac{1}{4}\alpha x^4 - \frac{1}{720}x^6 \alpha^6 - \frac{1}{240}x^6 \alpha^5 \right. \\
 & \left. + \frac{23}{720}x^6 \alpha^4 + \frac{17}{240}x^6 \alpha^3 - \frac{47}{360}x^6 \alpha^2 - \frac{1}{6}x^6 \alpha \right) y(0) \\
 & + \left(x - \frac{1}{6}\alpha^2 x^3 - \frac{1}{6}\alpha x^3 + \frac{1}{3}x^3 + \frac{1}{120}\alpha^4 x^5 + \frac{1}{60}x^5 \alpha^3 - \frac{13}{120}\alpha^2 x^5 - \frac{7}{60}x^5 \alpha + \frac{1}{5}x^5 \right) y'(0) \\
 & + O(x^6)
 \end{aligned}$$

Since the expansion point $x = 0$ is an ordinary, we can also solve this using standard power series. The ode is normalized to be

$$(-x^2 + 1)y'' - 2y'x + (\alpha^2 + \alpha)y = 0$$

Let the solution be represented as power series of the form

$$y = \sum_{n=0}^{\infty} a_n x^n$$

Then

$$\begin{aligned}
 y' &= \sum_{n=1}^{\infty} n a_n x^{n-1} \\
 y'' &= \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}
 \end{aligned}$$

Substituting the above back into the ode gives

$$(-x^2 + 1) \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) - 2 \left(\sum_{n=1}^{\infty} n a_n x^{n-1} \right) x + (\alpha^2 + \alpha) \left(\sum_{n=0}^{\infty} a_n x^n \right) = 0 \quad (1)$$

Which simplifies to

$$\begin{aligned}
 & \sum_{n=2}^{\infty} (-x^n a_n n(n-1)) + \left(\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \right) \\
 & + \sum_{n=1}^{\infty} (-2n a_n x^n) + \left(\sum_{n=0}^{\infty} (\alpha^2 + \alpha) a_n x^n \right) = 0
 \end{aligned} \quad (2)$$

The next step is to make all powers of x be n in each summation term. Going over each summation term above with power of x in it which is not already x^n and adjusting the power and the corresponding index gives

$$\sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n$$

Substituting all the above in Eq (2) gives the following equation where now all powers of x are the same and equal to n .

$$\begin{aligned} \sum_{n=2}^{\infty} (-x^n a_n n(n-1)) + \left(\sum_{n=0}^{\infty} (n+2) a_{n+2} (n+1) x^n \right) \\ + \sum_{n=1}^{\infty} (-2n a_n x^n) + \left(\sum_{n=0}^{\infty} (\alpha^2 + \alpha) a_n x^n \right) = 0 \end{aligned} \quad (3)$$

$n = 0$ gives

$$2a_2 + a_0\alpha(\alpha + 1) = 0$$

$$a_2 = -\frac{1}{2}a_0\alpha^2 - \frac{1}{2}a_0\alpha$$

$n = 1$ gives

$$6a_3 - 2a_1 + a_1\alpha(\alpha + 1) = 0$$

Which after substituting earlier equations, simplifies to

$$a_3 = -\frac{1}{6}a_1\alpha^2 - \frac{1}{6}a_1\alpha + \frac{1}{3}a_1$$

For $2 \leq n$, the recurrence equation is

$$-na_n(n-1) + (n+2) a_{n+2}(n+1) - 2na_n + a_n\alpha(\alpha + 1) = 0 \quad (4)$$

Solving for a_{n+2} , gives

$$a_{n+2} = -\frac{a_n(\alpha^2 - n^2 + \alpha - n)}{(n+2)(n+1)} \quad (5)$$

For $n = 2$ the recurrence equation gives

$$-6a_2 + 12a_4 + a_2\alpha(\alpha + 1) = 0$$

Which after substituting the earlier terms found becomes

$$a_4 = -\frac{5}{24}a_0\alpha^2 - \frac{1}{4}a_0\alpha + \frac{1}{24}a_0\alpha^4 + \frac{1}{12}a_0\alpha^3$$

For $n = 3$ the recurrence equation gives

$$-12a_3 + 20a_5 + a_3\alpha(\alpha + 1) = 0$$

Which after substituting the earlier terms found becomes

$$a_5 = -\frac{13}{120}a_1\alpha^2 - \frac{7}{60}a_1\alpha + \frac{1}{5}a_1 + \frac{1}{120}a_1\alpha^4 + \frac{1}{60}a_1\alpha^3$$

For $n = 4$ the recurrence equation gives

$$-20a_4 + 30a_6 + a_4\alpha(\alpha + 1) = 0$$

Which after substituting the earlier terms found becomes

$$a_6 = -\frac{47}{360}a_0\alpha^2 - \frac{1}{6}a_0\alpha + \frac{23}{720}a_0\alpha^4 + \frac{17}{240}a_0\alpha^3 - \frac{1}{720}a_0\alpha^6 - \frac{1}{240}a_0\alpha^5$$

For $n = 5$ the recurrence equation gives

$$-30a_5 + 42a_7 + a_5\alpha(\alpha + 1) = 0$$

Which after substituting the earlier terms found becomes

$$a_7 = -\frac{5}{63}a_1\alpha^2 - \frac{37}{420}a_1\alpha + \frac{1}{7}a_1 + \frac{41}{5040}a_1\alpha^4 + \frac{29}{1680}a_1\alpha^3 - \frac{1}{5040}a_1\alpha^6 - \frac{1}{1680}a_1\alpha^5$$

And so on. Therefore the solution is

$$\begin{aligned} y &= \sum_{n=0}^{\infty} a_n x^n \\ &= a_3 x^3 + a_2 x^2 + a_1 x + a_0 + \dots \end{aligned}$$

Substituting the values for a_n found above, the solution becomes

$$\begin{aligned} y = & a_0 + a_1x + \left(-\frac{1}{2}a_0\alpha^2 - \frac{1}{2}a_0\alpha\right)x^2 + \left(-\frac{1}{6}a_1\alpha^2 - \frac{1}{6}a_1\alpha + \frac{1}{3}a_1\right)x^3 \\ & + \left(-\frac{5}{24}a_0\alpha^2 - \frac{1}{4}a_0\alpha + \frac{1}{24}a_0\alpha^4 + \frac{1}{12}a_0\alpha^3\right)x^4 \\ & + \left(-\frac{13}{120}a_1\alpha^2 - \frac{7}{60}a_1\alpha + \frac{1}{5}a_1 + \frac{1}{120}a_1\alpha^4 + \frac{1}{60}a_1\alpha^3\right)x^5 + \dots \end{aligned}$$

Collecting terms, the solution becomes

$$\begin{aligned} y = & \left(1 + \left(-\frac{1}{2}\alpha^2 - \frac{1}{2}\alpha\right)x^2 + \left(-\frac{5}{24}\alpha^2 - \frac{1}{4}\alpha + \frac{1}{24}\alpha^4 + \frac{1}{12}\alpha^3\right)x^4\right)a_0 + \left(x \right. \\ & \left. + \left(-\frac{1}{6}\alpha^2 - \frac{1}{6}\alpha + \frac{1}{3}\right)x^3 + \left(-\frac{13}{120}\alpha^2 - \frac{7}{60}\alpha + \frac{1}{5} + \frac{1}{120}\alpha^4 + \frac{1}{60}\alpha^3\right)x^5\right)a_1 + O(x^6) \end{aligned} \quad (3)$$

At $x = 0$ the solution above becomes

$$\begin{aligned} y = & \left(1 + \left(-\frac{1}{2}\alpha^2 - \frac{1}{2}\alpha\right)x^2 + \left(-\frac{5}{24}\alpha^2 - \frac{1}{4}\alpha + \frac{1}{24}\alpha^4 + \frac{1}{12}\alpha^3\right)x^4\right)c_1 + \left(x \right. \\ & \left. + \left(-\frac{1}{6}\alpha^2 - \frac{1}{6}\alpha + \frac{1}{3}\right)x^3 + \left(-\frac{13}{120}\alpha^2 - \frac{7}{60}\alpha + \frac{1}{5} + \frac{1}{120}\alpha^4 + \frac{1}{60}\alpha^3\right)x^5\right)c_2 + O(x^6) \end{aligned}$$

Summary

The solution(s) found are the following

$$\begin{aligned} y = & \left(1 - \frac{1}{2}\alpha^2x^2 - \frac{1}{2}\alpha x^2 + \frac{1}{24}\alpha^4x^4 + \frac{1}{12}\alpha^3x^4 - \frac{5}{24}\alpha^2x^4 - \frac{1}{4}\alpha x^4 - \frac{1}{720}x^6\alpha^6 - \frac{1}{240}x^6\alpha^5 \right. \\ & \left. + \frac{23}{720}x^6\alpha^4 + \frac{17}{240}x^6\alpha^3 - \frac{47}{360}x^6\alpha^2 - \frac{1}{6}x^6\alpha\right)y(0) + \left(x - \frac{1}{6}\alpha^2x^3 - \frac{1}{6}\alpha x^3 + \frac{1}{3}x^3 \right. \\ & \left. + \frac{1}{120}\alpha^4x^5 + \frac{1}{60}x^5\alpha^3 - \frac{13}{120}\alpha^2x^5 - \frac{7}{60}x^5\alpha + \frac{1}{5}x^5\right)y'(0) + O(x^6) \end{aligned}$$

$$\begin{aligned} y = & \left(1 + \left(-\frac{1}{2}\alpha^2 - \frac{1}{2}\alpha\right)x^2 + \left(-\frac{5}{24}\alpha^2 - \frac{1}{4}\alpha + \frac{1}{24}\alpha^4 + \frac{1}{12}\alpha^3\right)x^4\right)c_1 \\ & + \left(x + \left(-\frac{1}{6}\alpha^2 - \frac{1}{6}\alpha + \frac{1}{3}\right)x^3 + \left(-\frac{13}{120}\alpha^2 - \frac{7}{60}\alpha + \frac{1}{5} + \frac{1}{120}\alpha^4 + \frac{1}{60}\alpha^3\right)x^5\right)c_2 \\ & + O(x^6) \end{aligned}$$

Verification of solutions

$$y = \left(1 - \frac{1}{2}\alpha^2 x^2 - \frac{1}{2}\alpha x^2 + \frac{1}{24}\alpha^4 x^4 + \frac{1}{12}\alpha^3 x^4 - \frac{5}{24}\alpha^2 x^4 - \frac{1}{4}\alpha x^4 - \frac{1}{720}x^6 \alpha^6 - \frac{1}{240}x^6 \alpha^5 \right. \\ \left. + \frac{23}{720}x^6 \alpha^4 + \frac{17}{240}x^6 \alpha^3 - \frac{47}{360}x^6 \alpha^2 - \frac{1}{6}x^6 \alpha \right) y(0) \\ + \left(x - \frac{1}{6}\alpha^2 x^3 - \frac{1}{6}\alpha x^3 + \frac{1}{3}x^3 + \frac{1}{120}\alpha^4 x^5 + \frac{1}{60}x^5 \alpha^3 - \frac{13}{120}\alpha^2 x^5 - \frac{7}{60}x^5 \alpha + \frac{1}{5}x^5 \right) y'(0) \\ + O(x^6)$$

Verified OK.

$$y = \left(1 + \left(-\frac{1}{2}\alpha^2 - \frac{1}{2}\alpha \right) x^2 + \left(-\frac{5}{24}\alpha^2 - \frac{1}{4}\alpha + \frac{1}{24}\alpha^4 + \frac{1}{12}\alpha^3 \right) x^4 \right) c_1 + \left(x \right. \\ \left. + \left(-\frac{1}{6}\alpha^2 - \frac{1}{6}\alpha + \frac{1}{3} \right) x^3 + \left(-\frac{13}{120}\alpha^2 - \frac{7}{60}\alpha + \frac{1}{5} + \frac{1}{120}\alpha^4 + \frac{1}{60}\alpha^3 \right) x^5 \right) c_2 + O(x^6)$$

Verified OK.

14.17.1 Maple step by step solution

Let's solve

$$(-x^2 + 1)y'' - 2y'x + (\alpha^2 + \alpha)y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = \frac{\alpha(\alpha+1)y}{x^2-1} - \frac{2xy'}{x^2-1}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + \frac{2xy'}{x^2-1} - \frac{\alpha(\alpha+1)y}{x^2-1} = 0$$

- Check to see if x_0 is a regular singular point

- Define functions

$$\left[P_2(x) = \frac{2x}{x^2-1}, P_3(x) = -\frac{\alpha(\alpha+1)}{x^2-1} \right]$$

- $(x+1) \cdot P_2(x)$ is analytic at $x = -1$

$$\left((x+1) \cdot P_2(x) \right) \Big|_{x=-1} = 1$$

- $(x+1)^2 \cdot P_3(x)$ is analytic at $x = -1$

$$((x + 1)^2 \cdot P_3(x)) \Big|_{x=-1} = 0$$

- $x = -1$ is a regular singular point

Check to see if x_0 is a regular singular point

$$x_0 = -1$$

- Multiply by denominators

$$(x^2 - 1)y'' + 2y'x - \alpha(\alpha + 1)y = 0$$

- Change variables using $x = u - 1$ so that the regular singular point is at $u = 0$

$$(u^2 - 2u) \left(\frac{d^2}{du^2} y(u) \right) + (2u - 2) \left(\frac{d}{du} y(u) \right) + (-\alpha^2 - \alpha) y(u) = 0$$

- Assume series solution for $y(u)$

$$y(u) = \sum_{k=0}^{\infty} a_k u^{k+r}$$

- Rewrite ODE with series expansions

- Convert $u^m \cdot \left(\frac{d}{du} y(u) \right)$ to series expansion for $m = 0..1$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=0}^{\infty} a_k (k + r) u^{k+r-1+m}$$

- Shift index using $k \rightarrow k + 1 - m$

$$u^m \cdot \left(\frac{d}{du} y(u) \right) = \sum_{k=-1+m}^{\infty} a_{k+1-m} (k + 1 - m + r) u^{k+r}$$

- Convert $u^m \cdot \left(\frac{d^2}{du^2} y(u) \right)$ to series expansion for $m = 1..2$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=0}^{\infty} a_k (k + r) (k + r - 1) u^{k+r-2+m}$$

- Shift index using $k \rightarrow k + 2 - m$

$$u^m \cdot \left(\frac{d^2}{du^2} y(u) \right) = \sum_{k=-2+m}^{\infty} a_{k+2-m} (k + 2 - m + r) (k + 1 - m + r) u^{k+r}$$

Rewrite ODE with series expansions

$$-2a_0 r^2 u^{-1+r} + \left(\sum_{k=0}^{\infty} (-2a_{k+1} (k + 1 + r)^2 - a_k (r + 1 + k + \alpha) (-r - k + \alpha)) u^{k+r} \right) = 0$$

- a_0 cannot be 0 by assumption, giving the indicial equation

$$-2r^2 = 0$$

- Values of r that satisfy the indicial equation

$$r = 0$$

- Each term in the series must be 0, giving the recursion relation

$$-2a_{k+1}(k+1)^2 + a_k(1+k+\alpha)(k-\alpha) = 0$$

- Recursion relation that defines series solution to ODE

$$a_{k+1} = -\frac{a_k(1+k+\alpha)(-k+\alpha)}{2(k+1)^2}$$

- Recursion relation for $r = 0$

$$a_{k+1} = -\frac{a_k(1+k+\alpha)(-k+\alpha)}{2(k+1)^2}$$

- Solution for $r = 0$

$$\left[y(u) = \sum_{k=0}^{\infty} a_k u^k, a_{k+1} = -\frac{a_k(1+k+\alpha)(-k+\alpha)}{2(k+1)^2} \right]$$

- Revert the change of variables $u = x + 1$

$$\left[y = \sum_{k=0}^{\infty} a_k (x+1)^k, a_{k+1} = -\frac{a_k(1+k+\alpha)(-k+\alpha)}{2(k+1)^2} \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
    -> Bessel
    -> elliptic
    -> Legendre
    <- Legendre successful
<- special function solution successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 101

Order:=6;

```
dsolve((1-x^2)*diff(y(x),x$2)-2*x*diff(y(x),x)+alpha*(alpha+1)*y(x)=0,y(x),type='series',x=0
```

$$y(x) = \left(1 - \frac{\alpha(1+\alpha)x^2}{2} + \frac{\alpha(\alpha^3 + 2\alpha^2 - 5\alpha - 6)x^4}{24}\right) y(0) \\ + \left(x - \frac{(\alpha^2 + \alpha - 2)x^3}{6} + \frac{(\alpha^4 + 2\alpha^3 - 13\alpha^2 - 14\alpha + 24)x^5}{120}\right) D(y)(0) + O(x^6)$$

✓ Solution by Mathematica

Time used: 0.001 (sec). Leaf size: 127

```
AsymptoticDSolveValue[(1-x^2)*y'[x]-2*x*y'[x]+\[Alpha]*(\[Alpha]+1)*y[x]==0,y[x],{x,0,5}]
```

$$y(x) \rightarrow c_2 \left(\frac{1}{60}(-\alpha^2 - \alpha)x^5 - \frac{1}{120}(-\alpha^2 - \alpha)(\alpha^2 + \alpha)x^5 - \frac{1}{10}(\alpha^2 + \alpha)x^5 + \frac{x^5}{5} \right. \\ \left. - \frac{1}{6}(\alpha^2 + \alpha)x^3 + \frac{x^3}{3} + x \right) + c_1 \left(\frac{1}{24}(\alpha^2 + \alpha)^2 x^4 - \frac{1}{4}(\alpha^2 + \alpha)x^4 - \frac{1}{2}(\alpha^2 + \alpha)x^2 + 1 \right)$$

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15.1 problem 30

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Internal problem ID [750]

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Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.5, Homogeneous Linear Systems with Constant Coefficients. page 407

Problem number: 30.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -\frac{x_1(t)}{10} + \frac{3x_2(t)}{40} \\x_2'(t) &= \frac{x_1(t)}{10} - \frac{x_2(t)}{5}\end{aligned}$$

With initial conditions

$$[x_1(0) = -17, x_2(0) = -21]$$

15.1.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{1}{10} & \frac{3}{40} \\ \frac{1}{10} & -\frac{1}{5} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{e^{-\frac{t}{4}}}{4} + \frac{3e^{-\frac{t}{20}}}{4} & \frac{3e^{-\frac{t}{20}}}{8} - \frac{3e^{-\frac{t}{4}}}{8} \\ \frac{e^{-\frac{t}{20}}}{2} - \frac{e^{-\frac{t}{4}}}{2} & \frac{3e^{-\frac{t}{4}}}{4} + \frac{e^{-\frac{t}{20}}}{4} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}\vec{x}_h(t) &= e^{At} \vec{x}_0 \\ &= \begin{bmatrix} \frac{e^{-\frac{t}{4}}}{4} + \frac{3e^{-\frac{t}{20}}}{4} & \frac{3e^{-\frac{t}{20}}}{8} - \frac{3e^{-\frac{t}{4}}}{8} \\ \frac{e^{-\frac{t}{20}}}{2} - \frac{e^{-\frac{t}{4}}}{2} & \frac{3e^{-\frac{t}{4}}}{4} + \frac{e^{-\frac{t}{20}}}{4} \end{bmatrix} \begin{bmatrix} -17 \\ -21 \end{bmatrix} \\ &= \begin{bmatrix} \frac{29e^{-\frac{t}{4}}}{8} - \frac{165e^{-\frac{t}{20}}}{8} \\ -\frac{55e^{-\frac{t}{20}}}{4} - \frac{29e^{-\frac{t}{4}}}{4} \end{bmatrix}\end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

15.1.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{1}{10} & \frac{3}{40} \\ \frac{1}{10} & -\frac{1}{5} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} -\frac{1}{10} & \frac{3}{40} \\ \frac{1}{10} & -\frac{1}{5} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -\frac{1}{10} - \lambda & \frac{3}{40} \\ \frac{1}{10} & -\frac{1}{5} - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + \frac{3}{10}\lambda + \frac{1}{80} = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -\frac{1}{4}$$

$$\lambda_2 = -\frac{1}{20}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-\frac{1}{4}$	1	real eigenvalue
$-\frac{1}{20}$	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -\frac{1}{4}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{10} & \frac{3}{40} \\ \frac{1}{10} & -\frac{1}{5} \end{bmatrix} - \left(-\frac{1}{4}\right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} \frac{3}{20} & \frac{3}{40} \\ \frac{1}{10} & \frac{1}{20} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{3}{20} & \frac{3}{40} & 0 \\ \frac{1}{10} & \frac{1}{20} & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{2R_1}{3} \implies \left[\begin{array}{cc|c} \frac{3}{20} & \frac{3}{40} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} \frac{3}{20} & \frac{3}{40} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -\frac{t}{2}\}$

Hence the solution is

$$\begin{bmatrix} -\frac{t}{2} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{t}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{t}{2} \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix} = \begin{bmatrix} -1 \\ 2 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -\frac{1}{20}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{10} & \frac{3}{40} \\ \frac{1}{10} & -\frac{1}{5} \end{bmatrix} - \left(-\frac{1}{20} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -\frac{1}{20} & \frac{3}{40} \\ \frac{1}{10} & -\frac{3}{20} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -\frac{1}{20} & \frac{3}{40} & 0 \\ \frac{1}{10} & -\frac{3}{20} & 0 \end{array} \right]$$

$$R_2 = R_2 + 2R_1 \implies \left[\begin{array}{cc|c} -\frac{1}{20} & \frac{3}{40} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -\frac{1}{20} & \frac{3}{40} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{3t}{2}\}$

Hence the solution is

$$\begin{bmatrix} \frac{3t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{3t}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{3t}{2} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{3}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{3t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{3}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{3t}{2} \\ t \end{bmatrix} = \begin{bmatrix} 3 \\ 2 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-\frac{1}{4}$	1	1	No	$\begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix}$
$-\frac{1}{20}$	1	1	No	$\begin{bmatrix} \frac{3}{2} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue $-\frac{1}{4}$ is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{-\frac{t}{4}} \\ &= \begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix} e^{-\frac{t}{4}}\end{aligned}$$

Since eigenvalue $-\frac{1}{20}$ is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-\frac{t}{20}} \\ &= \begin{bmatrix} \frac{3}{2} \\ 1 \end{bmatrix} e^{-\frac{t}{20}}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} -\frac{e^{-\frac{t}{4}}}{2} \\ e^{-\frac{t}{4}} \end{bmatrix} + c_2 \begin{bmatrix} \frac{3e^{-\frac{t}{20}}}{2} \\ e^{-\frac{t}{20}} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} -\frac{c_1 e^{-\frac{t}{4}}}{2} + \frac{3c_2 e^{-\frac{t}{20}}}{2} \\ c_1 e^{-\frac{t}{4}} + c_2 e^{-\frac{t}{20}} \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x_1(0) = -17 \\ x_2(0) = -21 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} -17 \\ -21 \end{bmatrix} = \begin{bmatrix} -\frac{c_1}{2} + \frac{3c_2}{2} \\ c_1 + c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = -\frac{29}{4} \\ c_2 = -\frac{55}{4} \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{29 e^{-\frac{t}{4}}}{8} - \frac{165 e^{-\frac{t}{20}}}{8} \\ -\frac{55 e^{-\frac{t}{20}}}{4} - \frac{29 e^{-\frac{t}{4}}}{4} \end{bmatrix}$$

The following is the phase plot of the system.

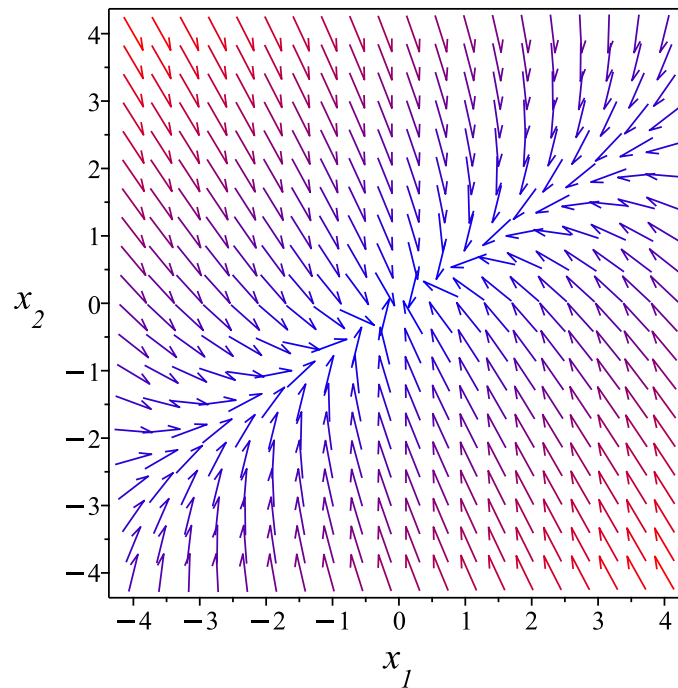
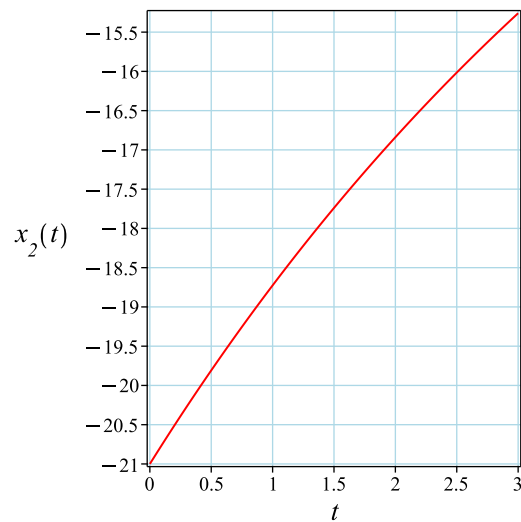
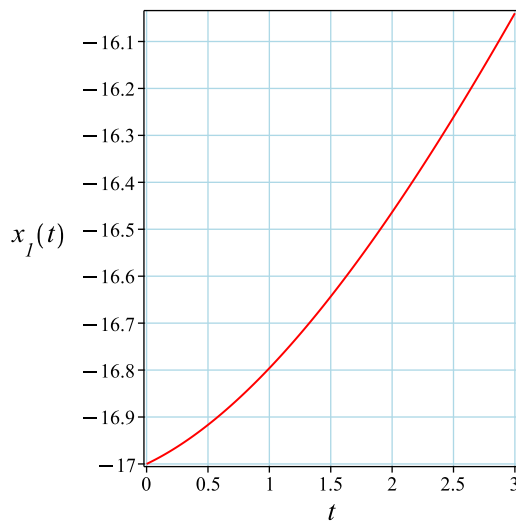


Figure 506: Phase plot

The following are plots of each solution.



✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 34

```
dsolve([diff(x__1(t),t) = -1/10*x__1(t)+3/40*x__2(t), diff(x__2(t),t) = 1/10*x__1(t)-1/5*x__2(t)],{x__1(0)=-17,x__2(0)=-21})
```

$$x_1(t) = -\frac{165 e^{-\frac{t}{20}}}{8} + \frac{29 e^{-\frac{t}{4}}}{8}$$

$$x_2(t) = -\frac{55 e^{-\frac{t}{20}}}{4} - \frac{29 e^{-\frac{t}{4}}}{4}$$

✓ Solution by Mathematica

Time used: 0.007 (sec). Leaf size: 52

```
DSolve[{x1'[t]==-1/10*x1[t]+3/40*x2[t],x2'[t]==1/10*x1[t]-1/5*x2[t]},{x1[0]==-17,x2[0]==-21}]
```

$$x_1(t) \rightarrow \frac{1}{8}e^{-t/4}(29 - 165e^{t/5})$$

$$x_2(t) \rightarrow -\frac{1}{4}e^{-t/4}(55e^{t/5} + 29)$$

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16.1 problem 1

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Internal problem ID [751]

Internal file name [OUTPUT/751_Sunday_June_05_2022_01_48_43_AM_43130140/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$x_1'(t) = 3x_1(t) - 2x_2(t)$$

$$x_2'(t) = 4x_1(t) - x_2(t)$$

16.1.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned} e^{At} &= \begin{bmatrix} e^t \cos(2t) + e^t \sin(2t) & -e^t \sin(2t) \\ 2e^t \sin(2t) & e^t \cos(2t) - e^t \sin(2t) \end{bmatrix} \\ &= \begin{bmatrix} e^t(\cos(2t) + \sin(2t)) & -e^t \sin(2t) \\ 2e^t \sin(2t) & e^t(\cos(2t) - \sin(2t)) \end{bmatrix} \end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^t(\cos(2t) + \sin(2t)) & -e^t \sin(2t) \\ 2e^t \sin(2t) & e^t(\cos(2t) - \sin(2t)) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^t(\cos(2t) + \sin(2t))c_1 - e^t \sin(2t)c_2 \\ 2e^t \sin(2t)c_1 + e^t(\cos(2t) - \sin(2t))c_2 \end{bmatrix} \\
 &= \begin{bmatrix} ((-c_2 + c_1)\sin(2t) + c_1 \cos(2t))e^t \\ e^t(2c_1 - c_2)\sin(2t) + e^t \cos(2t)c_2 \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.1.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 3 - \lambda & -2 \\ 4 & -1 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 2\lambda + 5 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 1 + 2i$$

$$\lambda_2 = 1 - 2i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$1 + 2i$	1	complex eigenvalue
$1 - 2i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = 1 - 2i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} - (1 - 2i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 + 2i & -2 \\ 4 & -2 + 2i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 + 2i & -2 & 0 \\ 4 & -2 + 2i & 0 \end{array} \right]$$

$$R_2 = R_2 + (-1 + i)R_1 \implies \left[\begin{array}{cc|c} 2 + 2i & -2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 + 2i & -2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (\frac{1}{2} - \frac{i}{2}) t\}$

Hence the solution is

$$\begin{bmatrix} (\frac{1}{2} - \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} (\frac{1}{2} - \frac{i}{2}) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (\frac{1}{2} - \frac{i}{2}) t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{2} - \frac{i}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (\frac{1}{2} - \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} \frac{1}{2} - \frac{i}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (\frac{1}{2} - \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} 1 - i \\ 2 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 1 + 2i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} - (1 + 2i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 - 2i & -2 \\ 4 & -2 - 2i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 - 2i & -2 & 0 \\ 4 & -2 - 2i & 0 \end{array} \right]$$

$$R_2 = R_2 + (-1 - i) R_1 \implies \left[\begin{array}{cc|c} 2 - 2i & -2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 - 2i & -2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (\frac{1}{2} + \frac{i}{2}) t\}$

Hence the solution is

$$\begin{bmatrix} (\frac{1}{2} + \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} (\frac{1}{2} + \frac{i}{2}) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (\frac{1}{2} + \frac{i}{2}) t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{2} + \frac{i}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (\frac{1}{2} + \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} \frac{1}{2} + \frac{i}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (\frac{1}{2} + \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} 1 + i \\ 2 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$1 + 2i$	1	1	No	$\begin{bmatrix} \frac{1}{2} + \frac{i}{2} \\ 1 \end{bmatrix}$
$1 - 2i$	1	1	No	$\begin{bmatrix} \frac{1}{2} - \frac{i}{2} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \left(\frac{1}{2} + \frac{i}{2}\right) e^{(1+2i)t} \\ e^{(1+2i)t} \end{bmatrix} + c_2 \begin{bmatrix} \left(\frac{1}{2} - \frac{i}{2}\right) e^{(1-2i)t} \\ e^{(1-2i)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \left(\frac{1}{2} + \frac{i}{2}\right) c_1 e^{(1+2i)t} + \left(\frac{1}{2} - \frac{i}{2}\right) c_2 e^{(1-2i)t} \\ c_1 e^{(1+2i)t} + c_2 e^{(1-2i)t} \end{bmatrix}$$

The following is the phase plot of the system.

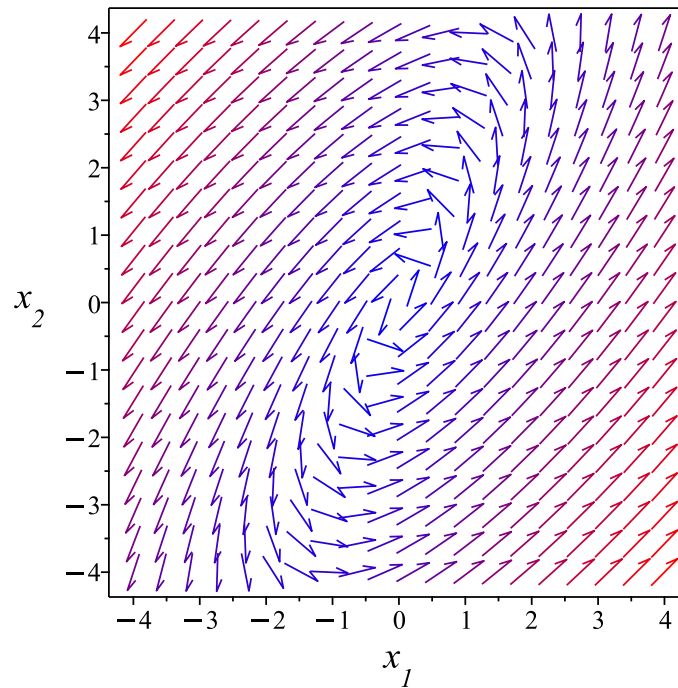


Figure 507: Phase plot

16.1.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 3x_1(t) - 2x_2(t), x_2'(t) = 4x_1(t) - x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[1 - 2I, \begin{bmatrix} \frac{1}{2} - \frac{I}{2} \\ 1 \end{bmatrix} \right], \left[1 + 2I, \begin{bmatrix} \frac{1}{2} + \frac{I}{2} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[1 - 2I, \begin{bmatrix} \frac{1}{2} - \frac{I}{2} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(1-2I)t} \cdot \begin{bmatrix} \frac{1}{2} - \frac{I}{2} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^t \cdot (\cos(2t) - I \sin(2t)) \cdot \begin{bmatrix} \frac{1}{2} - \frac{I}{2} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^t \cdot \begin{bmatrix} \left(\frac{1}{2} - \frac{I}{2}\right) (\cos(2t) - I \sin(2t)) \\ \cos(2t) - I \sin(2t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\underline{x}^{\rightarrow}_1(t) = e^t \cdot \begin{bmatrix} \frac{\cos(2t)}{2} - \frac{\sin(2t)}{2} \\ \cos(2t) \end{bmatrix}, \underline{x}^{\rightarrow}_2(t) = e^t \cdot \begin{bmatrix} -\frac{\sin(2t)}{2} - \frac{\cos(2t)}{2} \\ -\sin(2t) \end{bmatrix} \right]$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1(t) + c_2 \underline{x}^{\rightarrow}_2(t)$$

- Substitute solutions into the general solution

$$\vec{x} = c_1 e^t \cdot \begin{bmatrix} \frac{\cos(2t)}{2} - \frac{\sin(2t)}{2} \\ \cos(2t) \end{bmatrix} + c_2 e^t \cdot \begin{bmatrix} -\frac{\sin(2t)}{2} - \frac{\cos(2t)}{2} \\ -\sin(2t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{e^t((c_1 - c_2)\cos(2t) - \sin(2t)(c_1 + c_2))}{2} \\ e^t(-c_2 \sin(2t) + c_1 \cos(2t)) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{e^t((c_1 - c_2)\cos(2t) - \sin(2t)(c_1 + c_2))}{2}, x_2(t) = e^t(-c_2 \sin(2t) + c_1 \cos(2t)) \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 56

```
dsolve([diff(x__1(t),t)=3*x__1(t)-2*x__2(t),diff(x__2(t),t)=4*x__1(t)-1*x__2(t)],singsol=all
```

$$\begin{aligned} x_1(t) &= e^t(c_1 \sin(2t) + c_2 \cos(2t)) \\ x_2(t) &= -e^t(c_1 \cos(2t) - c_2 \cos(2t) - c_1 \sin(2t) - c_2 \sin(2t)) \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 58

```
DSolve[{x1'[t]==3*x1[t]-2*x2[t],x2'[t]==4*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$\begin{aligned} x1(t) &\rightarrow e^t(c_1 \cos(2t) + (c_1 - c_2) \sin(2t)) \\ x2(t) &\rightarrow e^t(c_2 \cos(2t) + (2c_1 - c_2) \sin(2t)) \end{aligned}$$

16.2 problem 2

- 16.2.1 Solution using Matrix exponential method 3539
- 16.2.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3540
- 16.2.3 Maple step by step solution 3544

Internal problem ID [752]

Internal file name [OUTPUT/752_Sunday_June_05_2022_01_48_44_AM_37520237/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -x_1(t) - 4x_2(t) \\x_2'(t) &= x_1(t) - x_2(t)\end{aligned}$$

16.2.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -1 & -4 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-t} \cos(2t) & -2 e^{-t} \sin(2t) \\ \frac{e^{-t} \sin(2t)}{2} & e^{-t} \cos(2t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{-t} \cos(2t) & -2e^{-t} \sin(2t) \\ \frac{e^{-t} \sin(2t)}{2} & e^{-t} \cos(2t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-t} \cos(2t) c_1 - 2e^{-t} \sin(2t) c_2 \\ \frac{e^{-t} \sin(2t) c_1}{2} + e^{-t} \cos(2t) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-t} (\cos(2t) c_1 - 2 \sin(2t) c_2) \\ \frac{e^{-t} (\sin(2t) c_1 + 2 \cos(2t) c_2)}{2} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.2.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -1 & -4 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} -1 & -4 \\ 1 & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -1 - \lambda & -4 \\ 1 & -1 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 2\lambda + 5 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1 + 2i$$

$$\lambda_2 = -1 - 2i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-1 - 2i$	1	complex eigenvalue
$-1 + 2i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1 - 2i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & -4 \\ 1 & -1 \end{bmatrix} - (-1 - 2i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2i & -4 \\ 1 & 2i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2i & -4 & 0 \\ 1 & 2i & 0 \end{array} \right]$$

$$R_2 = R_2 + \frac{iR_1}{2} \implies \left[\begin{array}{cc|c} 2i & -4 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2i & -4 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -2it\}$

Hence the solution is

$$\begin{bmatrix} -2 I t \\ t \end{bmatrix} = \begin{bmatrix} -2it \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -2 I t \\ t \end{bmatrix} = t \begin{bmatrix} -2i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -2 I t \\ t \end{bmatrix} = \begin{bmatrix} -2i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -1 + 2i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & -4 \\ 1 & -1 \end{bmatrix} - (-1 + 2i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -2i & -4 \\ 1 & -2i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -2i & -4 & 0 \\ 1 & -2i & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{iR_1}{2} \implies \left[\begin{array}{cc|c} -2i & -4 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -2i & -4 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 2it\}$

Hence the solution is

$$\begin{bmatrix} 2 I t \\ t \end{bmatrix} = \begin{bmatrix} 2it \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 2 I t \\ t \end{bmatrix} = t \begin{bmatrix} 2i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 2 I t \\ t \end{bmatrix} = \begin{bmatrix} 2i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-1 + 2i$	1	1	No	$\begin{bmatrix} 2i \\ 1 \end{bmatrix}$
$-1 - 2i$	1	1	No	$\begin{bmatrix} -2i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} 2ie^{(-1+2i)t} \\ e^{(-1+2i)t} \end{bmatrix} + c_2 \begin{bmatrix} -2ie^{(-1-2i)t} \\ e^{(-1-2i)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} -2i(c_2e^{(-1-2i)t} - c_1e^{(-1+2i)t}) \\ c_1e^{(-1+2i)t} + c_2e^{(-1-2i)t} \end{bmatrix}$$

The following is the phase plot of the system.

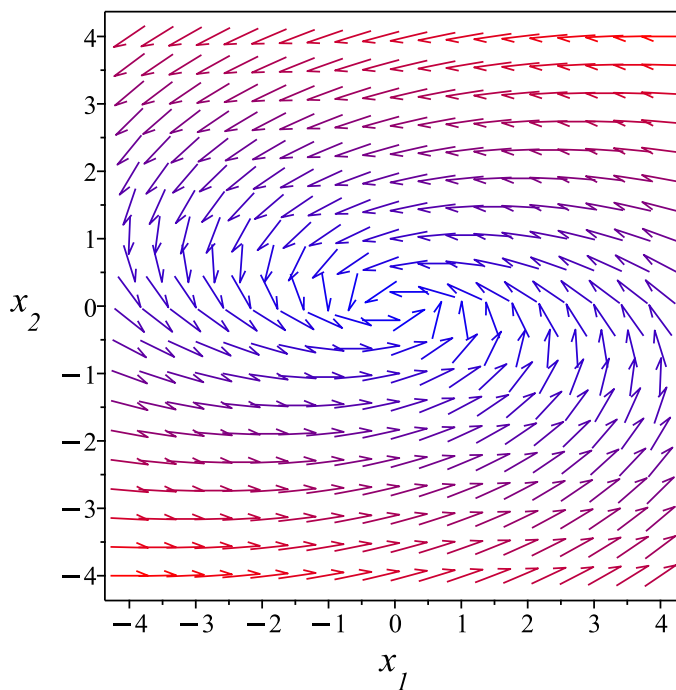


Figure 508: Phase plot

16.2.3 Maple step by step solution

Let's solve

$$[x_1'(t) = -x_1(t) - 4x_2(t), x_2'(t) = x_1(t) - x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -1 & -4 \\ 1 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -1 & -4 \\ 1 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -1 & -4 \\ 1 & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-1 - 2I, \begin{bmatrix} -2I \\ 1 \end{bmatrix} \right], \left[-1 + 2I, \begin{bmatrix} 2I \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-1 - 2I, \begin{bmatrix} -2I \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(-1-2I)t} \cdot \begin{bmatrix} -2I \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{-t} \cdot (\cos(2t) - I \sin(2t)) \cdot \begin{bmatrix} -2I \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{-t} \cdot \begin{bmatrix} -2I(\cos(2t) - I \sin(2t)) \\ \cos(2t) - I \sin(2t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\begin{bmatrix} \underline{x}_1(t) \\ \underline{x}_2(t) \end{bmatrix} = e^{-t} \cdot \begin{bmatrix} -2 \sin(2t) \\ \cos(2t) \end{bmatrix}, \underline{x}_2(t) = e^{-t} \cdot \begin{bmatrix} -2 \cos(2t) \\ -\sin(2t) \end{bmatrix}$$

- General solution to the system of ODEs

$$\underline{x} = c_1 \underline{x}_1(t) + c_2 \underline{x}_2(t)$$

- Substitute solutions into the general solution

$$\underline{x} = c_1 e^{-t} \cdot \begin{bmatrix} -2 \sin(2t) \\ \cos(2t) \end{bmatrix} + c_2 e^{-t} \cdot \begin{bmatrix} -2 \cos(2t) \\ -\sin(2t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} -2 e^{-t}(c_1 \sin(2t) + c_2 \cos(2t)) \\ e^{-t}(-c_2 \sin(2t) + c_1 \cos(2t)) \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = -2 e^{-t}(c_1 \sin(2t) + c_2 \cos(2t)), x_2(t) = e^{-t}(-c_2 \sin(2t) + c_1 \cos(2t))\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 46

```
dsolve([diff(x__1(t),t)=-1*x__1(t)-4*x__2(t),diff(x__2(t),t)=1*x__1(t)-1*x__2(t)],singsol=all)
```

$$\begin{aligned} x_1(t) &= e^{-t}(c_1 \sin(2t) + c_2 \cos(2t)) \\ x_2(t) &= -\frac{e^{-t}(c_1 \cos(2t) - c_2 \sin(2t))}{2} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 55

```
DSolve[{x1'[t]==-1*x1[t]-4*x2[t],x2'[t]==1*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingularSol
```

$$\begin{aligned} x1(t) &\rightarrow e^{-t}(c_1 \cos(2t) - 2c_2 \sin(2t)) \\ x2(t) &\rightarrow \frac{1}{2}e^{-t}(2c_2 \cos(2t) + c_1 \sin(2t)) \end{aligned}$$

16.3 problem 3

16.3.1 Solution using Matrix exponential method	3547
16.3.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3548
16.3.3 Maple step by step solution	3552

Internal problem ID [753]

Internal file name [OUTPUT/753_Sunday_June_05_2022_01_48_46_AM_75742661/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 3.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$\begin{aligned}x_1'(t) &= 2x_1(t) - 5x_2(t) \\x_2'(t) &= x_1(t) - 2x_2(t)\end{aligned}$$

16.3.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} (\cos(t) + 2 \sin(t)) c_1 - 5 \sin(t) c_2 \\ \sin(t) c_1 + (\cos(t) - 2 \sin(t)) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} (2c_1 - 5c_2) \sin(t) + c_1 \cos(t) \\ (c_1 - 2c_2) \sin(t) + c_2 \cos(t) \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.3.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 2 - \lambda & -5 \\ 1 & -2 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-i$	1	complex eigenvalue
i	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - (-i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2+i & -5 \\ 1 & -2+i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2+i & -5 & 0 \\ 1 & -2+i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} + \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2+i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2+i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 - i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = \begin{bmatrix} (2 - i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - (i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 - i & -5 \\ 1 & -2 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 1 & -2 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} - \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 - i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 + i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = \begin{bmatrix} (2 + i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
i	1	1	No	$\begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$
$-i$	1	1	No	$\begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (2+i)e^{it} \\ e^{it} \end{bmatrix} + c_2 \begin{bmatrix} (2-i)e^{-it} \\ e^{-it} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (2+i)c_1e^{it} + (2-i)c_2e^{-it} \\ c_1e^{it} + c_2e^{-it} \end{bmatrix}$$

The following is the phase plot of the system.

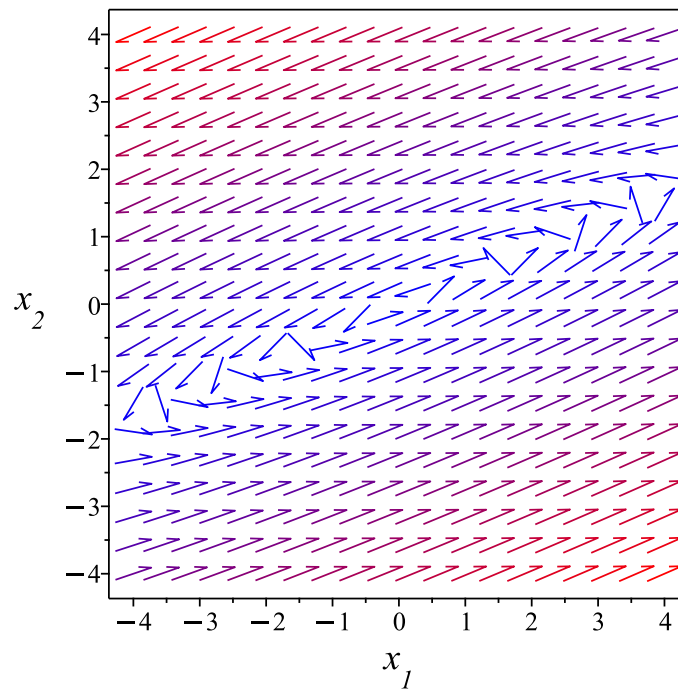


Figure 509: Phase plot

16.3.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 2x_1(t) - 5x_2(t), x_2'(t) = x_1(t) - 2x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right], \left[I, \begin{bmatrix} 2 + I \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{-It} \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$(\cos(t) - I \sin(t)) \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Simplify expression

$$\begin{bmatrix} (2 - I) (\cos (t) - I \sin (t)) \\ \cos (t) - I \sin (t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\vec{x}_{\rightarrow 1}(t) = \begin{bmatrix} 2 \cos (t) - \sin (t) \\ \cos (t) \end{bmatrix}, \vec{x}_{\rightarrow 2}(t) = \begin{bmatrix} -\cos (t) - 2 \sin (t) \\ -\sin (t) \end{bmatrix}$$

- General solution to the system of ODEs

$$\vec{x}_{\rightarrow} = c_1 \vec{x}_{\rightarrow 1}(t) + c_2 \vec{x}_{\rightarrow 2}(t)$$

- Substitute solutions into the general solution

$$\vec{x}_{\rightarrow} = \begin{bmatrix} c_2(-\cos (t) - 2 \sin (t)) + c_1(2 \cos (t) - \sin (t)) \\ -c_2 \sin (t) + c_1 \cos (t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \cos (t) (2c_1 - c_2) - \sin (t) (c_1 + 2c_2) \\ -c_2 \sin (t) + c_1 \cos (t) \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = \cos (t) (2c_1 - c_2) - \sin (t) (c_1 + 2c_2), x_2(t) = -c_2 \sin (t) + c_1 \cos (t)\}$$

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 38

```
dsolve([diff(x__1(t),t)=2*x__1(t)-5*x__2(t),diff(x__2(t),t)=1*x__1(t)-2*x__2(t)],singsol=all
```

$$\begin{aligned} x_1(t) &= c_1 \sin (t) + c_2 \cos (t) \\ x_2(t) &= -\frac{c_1 \cos (t)}{5} + \frac{c_2 \sin (t)}{5} + \frac{2c_1 \sin (t)}{5} + \frac{2c_2 \cos (t)}{5} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 41

```
DSolve[{x1'[t]==2*x1[t]-5*x2[t],x2'[t]==1*x1[t]-2*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$\begin{aligned} x1(t) &\rightarrow c_1(2 \sin (t) + \cos (t)) - 5c_2 \sin (t) \\ x2(t) &\rightarrow c_2 \cos (t) + (c_1 - 2c_2) \sin (t) \end{aligned}$$

16.4 problem 4

- 16.4.1 Solution using Matrix exponential method 3555
- 16.4.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3556
- 16.4.3 Maple step by step solution 3561

Internal problem ID [754]

Internal file name [OUTPUT/754_Sunday_June_05_2022_01_48_47_AM_27946383/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 4.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$\begin{aligned}x_1'(t) &= 2x_1(t) - \frac{5x_2(t)}{2} \\x_2'(t) &= \frac{9x_1(t)}{5} - x_2(t)\end{aligned}$$

16.4.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned}
 e^{At} &= \begin{bmatrix} e^{\frac{t}{2}} \cos\left(\frac{3t}{2}\right) + e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right) & -\frac{5e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{3} \\ \frac{6e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{5} & e^{\frac{t}{2}} \cos\left(\frac{3t}{2}\right) - e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right) \end{bmatrix} \\
 &= \begin{bmatrix} e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) + \sin\left(\frac{3t}{2}\right)\right) & -\frac{5e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{3} \\ \frac{6e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{5} & e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) - \sin\left(\frac{3t}{2}\right)\right) \end{bmatrix}
 \end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) + \sin\left(\frac{3t}{2}\right)\right) & -\frac{5e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{3} \\ \frac{6e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{5} & e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) - \sin\left(\frac{3t}{2}\right)\right) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) + \sin\left(\frac{3t}{2}\right)\right) c_1 - \frac{5e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right) c_2}{3} \\ \frac{6e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right) c_1}{5} + e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) - \sin\left(\frac{3t}{2}\right)\right) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \left(\left(c_1 - \frac{5c_2}{3}\right) \sin\left(\frac{3t}{2}\right) + c_1 \cos\left(\frac{3t}{2}\right)\right) e^{\frac{t}{2}} \\ \frac{e^{\frac{t}{2}} (6c_1 - 5c_2) \sin\left(\frac{3t}{2}\right)}{5} + e^{\frac{t}{2}} \cos\left(\frac{3t}{2}\right) c_2 \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.4.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 2 - \lambda & -\frac{5}{2} \\ \frac{9}{5} & -1 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - \lambda + \frac{5}{2} = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = \frac{1}{2} + \frac{3i}{2}$$

$$\lambda_2 = \frac{1}{2} - \frac{3i}{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$\frac{1}{2} - \frac{3i}{2}$	1	complex eigenvalue
$\frac{1}{2} + \frac{3i}{2}$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = \frac{1}{2} - \frac{3i}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} - \left(\frac{1}{2} - \frac{3i}{2} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} \frac{3}{2} + \frac{3i}{2} & -\frac{5}{2} \\ \frac{9}{5} & -\frac{3}{2} + \frac{3i}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{3}{2} + \frac{3i}{2} & -\frac{5}{2} & 0 \\ \frac{9}{5} & -\frac{3}{2} + \frac{3i}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{3}{5} + \frac{3i}{5}\right) R_1 \implies \left[\begin{array}{cc|c} \frac{3}{2} + \frac{3i}{2} & -\frac{5}{2} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{cc} \frac{3}{2} + \frac{3i}{2} & -\frac{5}{2} \\ 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (\frac{5}{6} - \frac{5i}{6}) t\}$

Hence the solution is

$$\begin{bmatrix} (\frac{5}{6} - \frac{5i}{6}) t \\ t \end{bmatrix} = \begin{bmatrix} (\frac{5}{6} - \frac{5i}{6}) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (\frac{5}{6} - \frac{5i}{6}) t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{5}{6} - \frac{5i}{6} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (\frac{5}{6} - \frac{5i}{6}) t \\ t \end{bmatrix} = \begin{bmatrix} \frac{5}{6} - \frac{5i}{6} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (\frac{5}{6} - \frac{5i}{6}) t \\ t \end{bmatrix} = \begin{bmatrix} 5 - 5i \\ 6 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = \frac{1}{2} + \frac{3i}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\left[\begin{array}{cc} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{array} \right] - \left(\frac{1}{2} + \frac{3i}{2} \right) \left[\begin{array}{cc} 1 & 0 \\ 0 & 1 \end{array} \right] \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\left[\begin{array}{cc} \frac{3}{2} - \frac{3i}{2} & -\frac{5}{2} \\ \frac{9}{5} & -\frac{3}{2} - \frac{3i}{2} \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{3}{2} - \frac{3i}{2} & -\frac{5}{2} & 0 \\ \frac{9}{5} & -\frac{3}{2} - \frac{3i}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{3}{5} - \frac{3i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} \frac{3}{2} - \frac{3i}{2} & -\frac{5}{2} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{cc} \frac{3}{2} - \frac{3i}{2} & -\frac{5}{2} \\ 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (\frac{5}{6} + \frac{5i}{6})t\}$

Hence the solution is

$$\begin{bmatrix} (\frac{5}{6} + \frac{5i}{6})t \\ t \end{bmatrix} = \begin{bmatrix} (\frac{5}{6} + \frac{5i}{6})t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (\frac{5}{6} + \frac{5i}{6})t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{5}{6} + \frac{5i}{6} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (\frac{5}{6} + \frac{5i}{6})t \\ t \end{bmatrix} = \begin{bmatrix} \frac{5}{6} + \frac{5i}{6} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (\frac{5}{6} + \frac{5i}{6})t \\ t \end{bmatrix} = \begin{bmatrix} 5 + 5i \\ 6 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$\frac{1}{2} + \frac{3i}{2}$	1	1	No	$\begin{bmatrix} \frac{5}{6} + \frac{5i}{6} \\ 1 \end{bmatrix}$
$\frac{1}{2} - \frac{3i}{2}$	1	1	No	$\begin{bmatrix} \frac{5}{6} - \frac{5i}{6} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \left(\frac{5}{6} + \frac{5i}{6}\right) e^{\left(\frac{1}{2} + \frac{3i}{2}\right)t} \\ e^{\left(\frac{1}{2} + \frac{3i}{2}\right)t} \end{bmatrix} + c_2 \begin{bmatrix} \left(\frac{5}{6} - \frac{5i}{6}\right) e^{\left(\frac{1}{2} - \frac{3i}{2}\right)t} \\ e^{\left(\frac{1}{2} - \frac{3i}{2}\right)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \left(\frac{5}{6} + \frac{5i}{6}\right) c_1 e^{\left(\frac{1}{2} + \frac{3i}{2}\right)t} + \left(\frac{5}{6} - \frac{5i}{6}\right) c_2 e^{\left(\frac{1}{2} - \frac{3i}{2}\right)t} \\ c_1 e^{\left(\frac{1}{2} + \frac{3i}{2}\right)t} + c_2 e^{\left(\frac{1}{2} - \frac{3i}{2}\right)t} \end{bmatrix}$$

The following is the phase plot of the system.

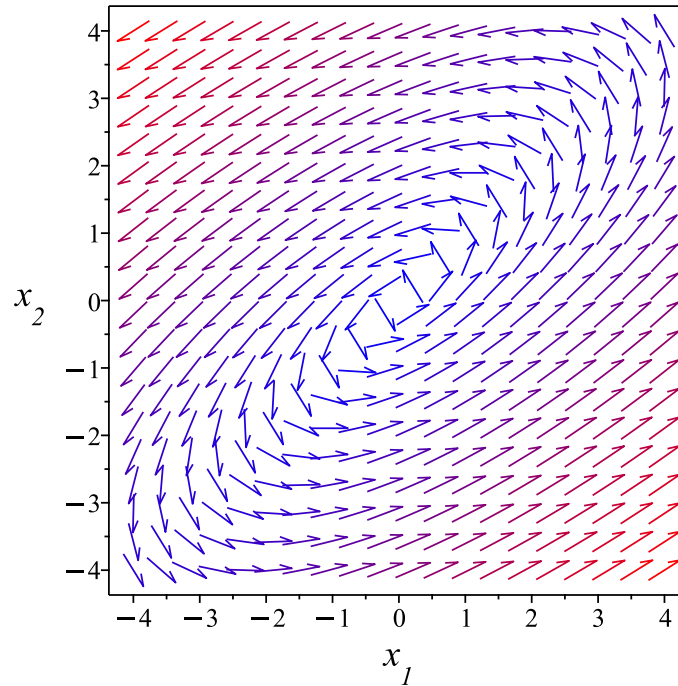


Figure 510: Phase plot

16.4.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = 2x_1(t) - \frac{5x_2(t)}{2}, x_2'(t) = \frac{9x_1(t)}{5} - x_2(t) \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[\frac{1}{2} - \frac{3I}{2}, \begin{bmatrix} \frac{5}{6} - \frac{5I}{6} \\ 1 \end{bmatrix} \right], \left[\frac{1}{2} + \frac{3I}{2}, \begin{bmatrix} \frac{5}{6} + \frac{5I}{6} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[\frac{1}{2} - \frac{3I}{2}, \begin{bmatrix} \frac{5}{6} - \frac{5I}{6} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{\left(\frac{1}{2} - \frac{3I}{2}\right)t} \cdot \begin{bmatrix} \frac{5}{6} - \frac{5I}{6} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{\frac{t}{2}} \cdot \left(\cos\left(\frac{3t}{2}\right) - I \sin\left(\frac{3t}{2}\right) \right) \cdot \begin{bmatrix} \frac{5}{6} - \frac{5I}{6} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{\frac{t}{2}} \cdot \begin{bmatrix} \left(\frac{5}{6} - \frac{5I}{6}\right) \left(\cos\left(\frac{3t}{2}\right) - I \sin\left(\frac{3t}{2}\right)\right) \\ \cos\left(\frac{3t}{2}\right) - I \sin\left(\frac{3t}{2}\right) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\underline{x}^{\rightarrow}_1(t) = e^{\frac{t}{2}} \cdot \begin{bmatrix} \frac{5 \cos(\frac{3t}{2})}{6} - \frac{5 \sin(\frac{3t}{2})}{6} \\ \cos\left(\frac{3t}{2}\right) \end{bmatrix}, \underline{x}^{\rightarrow}_2(t) = e^{\frac{t}{2}} \cdot \begin{bmatrix} -\frac{5 \sin(\frac{3t}{2})}{6} - \frac{5 \cos(\frac{3t}{2})}{6} \\ -\sin\left(\frac{3t}{2}\right) \end{bmatrix} \right]$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1(t) + c_2 \underline{x}^{\rightarrow}_2(t)$$

- Substitute solutions into the general solution

$$\underline{x}^{\rightarrow} = c_1 e^{\frac{t}{2}} \cdot \begin{bmatrix} \frac{5 \cos(\frac{3t}{2})}{6} - \frac{5 \sin(\frac{3t}{2})}{6} \\ \cos(\frac{3t}{2}) \end{bmatrix} + c_2 e^{\frac{t}{2}} \cdot \begin{bmatrix} -\frac{5 \sin(\frac{3t}{2})}{6} - \frac{5 \cos(\frac{3t}{2})}{6} \\ -\sin(\frac{3t}{2}) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{5((c_1 - c_2) \cos(\frac{3t}{2}) - \sin(\frac{3t}{2})(c_1 + c_2)) e^{\frac{t}{2}}}{6} \\ e^{\frac{t}{2}} (c_1 \cos(\frac{3t}{2}) - c_2 \sin(\frac{3t}{2})) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{5((c_1 - c_2) \cos(\frac{3t}{2}) - \sin(\frac{3t}{2})(c_1 + c_2)) e^{\frac{t}{2}}}{6}, x_2(t) = e^{\frac{t}{2}} (c_1 \cos(\frac{3t}{2}) - c_2 \sin(\frac{3t}{2})) \right\}$$

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 58

```
dsolve([diff(x__1(t),t)=2*x__1(t)-5/2*x__2(t),diff(x__2(t),t)=9/5*x__1(t)-1*x__2(t)],singsol
```

$$x_1(t) = e^{\frac{t}{2}} \left(\sin\left(\frac{3t}{2}\right) c_1 + \cos\left(\frac{3t}{2}\right) c_2 \right)$$

$$x_2(t) = \frac{3 e^{\frac{t}{2}} \left(\sin\left(\frac{3t}{2}\right) c_1 + \sin\left(\frac{3t}{2}\right) c_2 - \cos\left(\frac{3t}{2}\right) c_1 + \cos\left(\frac{3t}{2}\right) c_2 \right)}{5}$$

✓ Solution by Mathematica

Time used: 0.006 (sec). Leaf size: 84

```
DSolve[{x1'[t]==2*x1[t]-5/2*x2[t],x2'[t]==9/5*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingular
```

$$x1(t) \rightarrow \frac{1}{3} e^{t/2} \left(3c_1 \cos\left(\frac{3t}{2}\right) + (3c_1 - 5c_2) \sin\left(\frac{3t}{2}\right) \right)$$

$$x2(t) \rightarrow \frac{1}{5} e^{t/2} \left(5c_2 \cos\left(\frac{3t}{2}\right) + (6c_1 - 5c_2) \sin\left(\frac{3t}{2}\right) \right)$$

16.5 problem 5

16.5.1 Solution using Matrix exponential method	3564
16.5.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3565
16.5.3 Maple step by step solution	3570

Internal problem ID [755]

Internal file name [OUTPUT/755_Sunday_June_05_2022_01_48_49_AM_29534035/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 5.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) - x_2(t) \\x_2'(t) &= 5x_1(t) - 3x_2(t)\end{aligned}$$

16.5.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & -1 \\ 5 & -3 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned}e^{At} &= \begin{bmatrix} e^{-t} \cos(t) + 2e^{-t} \sin(t) & -e^{-t} \sin(t) \\ 5e^{-t} \sin(t) & e^{-t} \cos(t) - 2e^{-t} \sin(t) \end{bmatrix} \\ &= \begin{bmatrix} e^{-t}(\cos(t) + 2\sin(t)) & -e^{-t} \sin(t) \\ 5e^{-t} \sin(t) & e^{-t}(\cos(t) - 2\sin(t)) \end{bmatrix}\end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{-t}(\cos(t) + 2\sin(t)) & -e^{-t}\sin(t) \\ 5e^{-t}\sin(t) & e^{-t}(\cos(t) - 2\sin(t)) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-t}(\cos(t) + 2\sin(t))c_1 - e^{-t}\sin(t)c_2 \\ 5e^{-t}\sin(t)c_1 + e^{-t}(\cos(t) - 2\sin(t))c_2 \end{bmatrix} \\
 &= \begin{bmatrix} ((2c_1 - c_2)\sin(t) + c_1\cos(t))e^{-t} \\ ((5c_1 - 2c_2)\sin(t) + c_2\cos(t))e^{-t} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.5.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & -1 \\ 5 & -3 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 1 & -1 \\ 5 & -3 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 1 - \lambda & -1 \\ 5 & -3 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 2\lambda + 2 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1 + i$$

$$\lambda_2 = -1 - i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-1 - i$	1	complex eigenvalue
$-1 + i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1 - i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & -1 \\ 5 & -3 \end{bmatrix} - (-1 - i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 + i & -1 \\ 5 & -2 + i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 + i & -1 & 0 \\ 5 & -2 + i & 0 \end{array} \right]$$

$$R_2 = R_2 + (-2 + i)R_1 \implies \left[\begin{array}{cc|c} 2 + i & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 + i & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (\frac{2}{5} - \frac{i}{5}) t\}$

Hence the solution is

$$\begin{bmatrix} (\frac{2}{5} - \frac{i}{5}) t \\ t \end{bmatrix} = \begin{bmatrix} (\frac{2}{5} - \frac{i}{5}) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (\frac{2}{5} - \frac{i}{5}) t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{2}{5} - \frac{i}{5} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (\frac{2}{5} - \frac{i}{5}) t \\ t \end{bmatrix} = \begin{bmatrix} \frac{2}{5} - \frac{i}{5} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (\frac{2}{5} - \frac{i}{5}) t \\ t \end{bmatrix} = \begin{bmatrix} 2 - i \\ 5 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -1 + i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & -1 \\ 5 & -3 \end{bmatrix} - (-1 + i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 - i & -1 \\ 5 & -2 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 - i & -1 & 0 \\ 5 & -2 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + (-2 - i)R_1 \implies \left[\begin{array}{cc|c} 2 - i & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2-i & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (\frac{2}{5} + \frac{i}{5})t\}$

Hence the solution is

$$\begin{bmatrix} (\frac{2}{5} + \frac{i}{5})t \\ t \end{bmatrix} = \begin{bmatrix} (\frac{2}{5} + \frac{i}{5})t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (\frac{2}{5} + \frac{i}{5})t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{2}{5} + \frac{i}{5} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (\frac{2}{5} + \frac{i}{5})t \\ t \end{bmatrix} = \begin{bmatrix} \frac{2}{5} + \frac{i}{5} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (\frac{2}{5} + \frac{i}{5})t \\ t \end{bmatrix} = \begin{bmatrix} 2+i \\ 5 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-1 + i$	1	1	No	$\begin{bmatrix} \frac{2}{5} + \frac{i}{5} \\ 1 \end{bmatrix}$
$-1 - i$	1	1	No	$\begin{bmatrix} \frac{2}{5} - \frac{i}{5} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (\frac{2}{5} + \frac{i}{5}) e^{(-1+i)t} \\ e^{(-1+i)t} \end{bmatrix} + c_2 \begin{bmatrix} (\frac{2}{5} - \frac{i}{5}) e^{(-1-i)t} \\ e^{(-1-i)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (\frac{2}{5} + \frac{i}{5}) c_1 e^{(-1+i)t} + (\frac{2}{5} - \frac{i}{5}) c_2 e^{(-1-i)t} \\ c_1 e^{(-1+i)t} + c_2 e^{(-1-i)t} \end{bmatrix}$$

The following is the phase plot of the system.

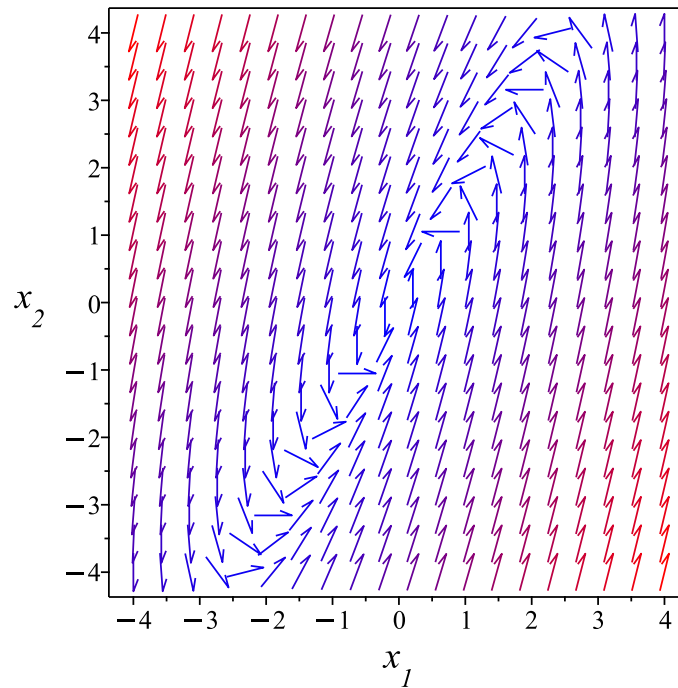


Figure 511: Phase plot

16.5.3 Maple step by step solution

Let's solve

$$[x_1'(t) = x_1(t) - x_2(t), x_2'(t) = 5x_1(t) - 3x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & -1 \\ 5 & -3 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & -1 \\ 5 & -3 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 1 & -1 \\ 5 & -3 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-1 - I, \begin{bmatrix} \frac{2}{5} - \frac{I}{5} \\ 1 \end{bmatrix} \right], \left[-1 + I, \begin{bmatrix} \frac{2}{5} + \frac{I}{5} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-1 - I, \begin{bmatrix} \frac{2}{5} - \frac{I}{5} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(-1-I)t} \cdot \begin{bmatrix} \frac{2}{5} - \frac{I}{5} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{-t} \cdot (\cos(t) - I \sin(t)) \cdot \begin{bmatrix} \frac{2}{5} - \frac{I}{5} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{-t} \cdot \begin{bmatrix} \left(\frac{2}{5} - \frac{I}{5}\right) (\cos(t) - I \sin(t)) \\ \cos(t) - I \sin(t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\underline{x}^{\rightarrow}_1(t) = e^{-t} \cdot \begin{bmatrix} \frac{2 \cos(t)}{5} - \frac{\sin(t)}{5} \\ \cos(t) \end{bmatrix}, \underline{x}^{\rightarrow}_2(t) = e^{-t} \cdot \begin{bmatrix} -\frac{2 \sin(t)}{5} - \frac{\cos(t)}{5} \\ -\sin(t) \end{bmatrix} \right]$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1(t) + c_2 \underline{x}^{\rightarrow}_2(t)$$

- Substitute solutions into the general solution

$$\underline{x}^{\rightarrow} = c_1 e^{-t} \cdot \begin{bmatrix} \frac{2 \cos(t)}{5} - \frac{\sin(t)}{5} \\ \cos(t) \end{bmatrix} + c_2 e^{-t} \cdot \begin{bmatrix} -\frac{2 \sin(t)}{5} - \frac{\cos(t)}{5} \\ -\sin(t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{2 \left((c_1 - \frac{c_2}{2}) \cos(t) - \frac{\sin(t)(c_1 + 2c_2)}{2} \right) e^{-t}}{5} \\ e^{-t} (-c_2 \sin(t) + c_1 \cos(t)) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{2 \left((c_1 - \frac{c_2}{2}) \cos(t) - \frac{\sin(t)(c_1 + 2c_2)}{2} \right) e^{-t}}{5}, x_2(t) = e^{-t} (-c_2 \sin(t) + c_1 \cos(t)) \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 48

```
dsolve([diff(x__1(t),t)=1*x__1(t)-1*x__2(t),diff(x__2(t),t)=5*x__1(t)-3*x__2(t)],singsol=all
```

$$\begin{aligned} x_1(t) &= e^{-t}(c_1 \sin(t) + c_2 \cos(t)) \\ x_2(t) &= -e^{-t}(c_1 \cos(t) - 2c_2 \cos(t) - 2c_1 \sin(t) - c_2 \sin(t)) \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 56

```
DSolve[{x1'[t]==1*x1[t]-1*x2[t],x2'[t]==5*x1[t]-3*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$\begin{aligned} x1(t) &\rightarrow e^{-t}(c_1 \cos(t) + (2c_1 - c_2) \sin(t)) \\ x2(t) &\rightarrow e^{-t}(c_2 \cos(t) + (5c_1 - 2c_2) \sin(t)) \end{aligned}$$

16.6 problem 6

- 16.6.1 Solution using Matrix exponential method 3573
- 16.6.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3574
- 16.6.3 Maple step by step solution 3579

Internal problem ID [756]

Internal file name [OUTPUT/756_Sunday_June_05_2022_01_48_51_AM_70838050/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 6.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) + 2x_2(t) \\x_2'(t) &= -5x_1(t) - x_2(t)\end{aligned}$$

16.6.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & 2 \\ -5 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \cos(3t) + \frac{\sin(3t)}{3} & \frac{2\sin(3t)}{3} \\ -\frac{5\sin(3t)}{3} & \cos(3t) - \frac{\sin(3t)}{3} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} \cos(3t) + \frac{\sin(3t)}{3} & \frac{2\sin(3t)}{3} \\ -\frac{5\sin(3t)}{3} & \cos(3t) - \frac{\sin(3t)}{3} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \left(\cos(3t) + \frac{\sin(3t)}{3}\right) c_1 + \frac{2\sin(3t)c_2}{3} \\ -\frac{5\sin(3t)c_1}{3} + \left(\cos(3t) - \frac{\sin(3t)}{3}\right) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \frac{\sin(3t)(c_1+2c_2)}{3} + c_1 \cos(3t) \\ \frac{(-5c_1-c_2)\sin(3t)}{3} + c_2 \cos(3t) \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.6.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & 2 \\ -5 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 1 & 2 \\ -5 & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 1 - \lambda & 2 \\ -5 & -1 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 9 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 3i$$

$$\lambda_2 = -3i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-3i$	1	complex eigenvalue
$3i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -3i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 2 \\ -5 & -1 \end{bmatrix} - (-3i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 + 3i & 2 \\ -5 & -1 + 3i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 + 3i & 2 & 0 \\ -5 & -1 + 3i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(\frac{1}{2} - \frac{3i}{2} \right) R_1 \implies \left[\begin{array}{cc|c} 1 + 3i & 2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 + 3i & 2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (-\frac{1}{5} + \frac{3i}{5}) t\}$

Hence the solution is

$$\begin{bmatrix} (-\frac{1}{5} + \frac{3i}{5}) t \\ t \end{bmatrix} = \begin{bmatrix} (-\frac{1}{5} + \frac{3i}{5}) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (-\frac{1}{5} + \frac{3i}{5}) t \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{1}{5} + \frac{3i}{5} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (-\frac{1}{5} + \frac{3i}{5}) t \\ t \end{bmatrix} = \begin{bmatrix} -\frac{1}{5} + \frac{3i}{5} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (-\frac{1}{5} + \frac{3i}{5}) t \\ t \end{bmatrix} = \begin{bmatrix} -1 + 3i \\ 5 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 3i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 2 \\ -5 & -1 \end{bmatrix} - (3i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 - 3i & 2 \\ -5 & -1 - 3i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 - 3i & 2 & 0 \\ -5 & -1 - 3i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(\frac{1}{2} + \frac{3i}{2}\right) R_1 \implies \left[\begin{array}{cc|c} 1 - 3i & 2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{cc} 1 - 3i & 2 \\ 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (-\frac{1}{5} - \frac{3i}{5}) t\}$

Hence the solution is

$$\begin{bmatrix} (-\frac{1}{5} - \frac{3i}{5}) t \\ t \end{bmatrix} = \begin{bmatrix} (-\frac{1}{5} - \frac{3i}{5}) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (-\frac{1}{5} - \frac{3i}{5}) t \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{1}{5} - \frac{3i}{5} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (-\frac{1}{5} - \frac{3i}{5}) t \\ t \end{bmatrix} = \begin{bmatrix} -\frac{1}{5} - \frac{3i}{5} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (-\frac{1}{5} - \frac{3i}{5}) t \\ t \end{bmatrix} = \begin{bmatrix} -1 - 3i \\ 5 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$3i$	1	1	No	$\begin{bmatrix} -\frac{1}{5} - \frac{3i}{5} \\ 1 \end{bmatrix}$
$-3i$	1	1	No	$\begin{bmatrix} -\frac{1}{5} + \frac{3i}{5} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \left(-\frac{1}{5} - \frac{3i}{5}\right) e^{3it} \\ e^{3it} \end{bmatrix} + c_2 \begin{bmatrix} \left(-\frac{1}{5} + \frac{3i}{5}\right) e^{-3it} \\ e^{-3it} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \left(-\frac{1}{5} - \frac{3i}{5}\right) c_1 e^{3it} + \left(-\frac{1}{5} + \frac{3i}{5}\right) c_2 e^{-3it} \\ c_1 e^{3it} + c_2 e^{-3it} \end{bmatrix}$$

The following is the phase plot of the system.

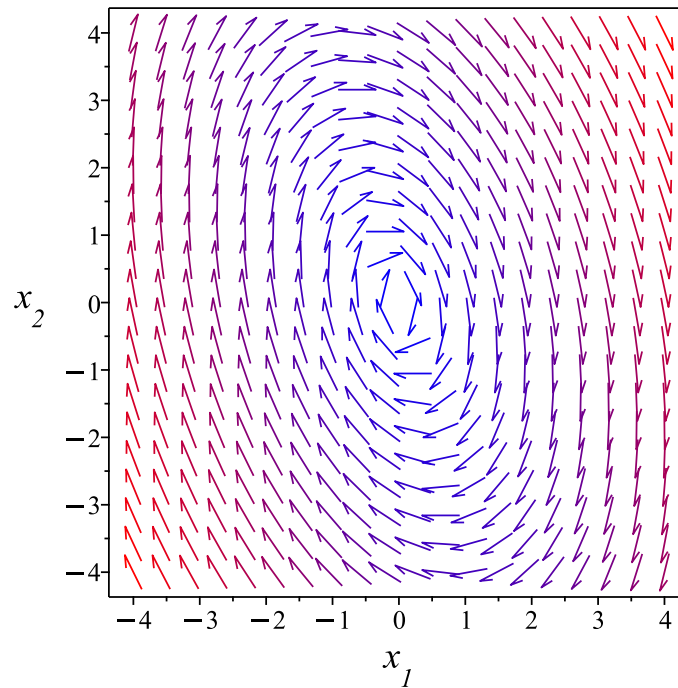


Figure 512: Phase plot

16.6.3 Maple step by step solution

Let's solve

$$[x_1'(t) = x_1(t) + 2x_2(t), x_2'(t) = -5x_1(t) - x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & 2 \\ -5 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & 2 \\ -5 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 1 & 2 \\ -5 & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow}{}'(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-3\mathbf{I}, \begin{bmatrix} -\frac{1}{5} + \frac{3\mathbf{I}}{5} \\ 1 \end{bmatrix} \right], \left[3\mathbf{I}, \begin{bmatrix} -\frac{1}{5} - \frac{3\mathbf{I}}{5} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-3\mathbf{I}, \begin{bmatrix} -\frac{1}{5} + \frac{3\mathbf{I}}{5} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{-3\mathbf{I}t} \cdot \begin{bmatrix} -\frac{1}{5} + \frac{3\mathbf{I}}{5} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$(\cos(3t) - \mathbf{I} \sin(3t)) \cdot \begin{bmatrix} -\frac{1}{5} + \frac{3\mathbf{I}}{5} \\ 1 \end{bmatrix}$$

- Simplify expression

$$\begin{bmatrix} \left(-\frac{1}{5} + \frac{3\mathbf{I}}{5}\right) (\cos(3t) - \mathbf{I} \sin(3t)) \\ \cos(3t) - \mathbf{I} \sin(3t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\underline{x}^{\rightarrow}_1(t) = \begin{bmatrix} -\frac{\cos(3t)}{5} + \frac{3\sin(3t)}{5} \\ \cos(3t) \end{bmatrix}, \underline{x}^{\rightarrow}_2(t) = \begin{bmatrix} \frac{\sin(3t)}{5} + \frac{3\cos(3t)}{5} \\ -\sin(3t) \end{bmatrix}$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1(t) + c_2 \underline{x}^{\rightarrow}_2(t)$$

- Substitute solutions into the general solution

$$\vec{x} = \begin{bmatrix} c_2 \left(\frac{\sin(3t)}{5} + \frac{3 \cos(3t)}{5} \right) + c_1 \left(-\frac{\cos(3t)}{5} + \frac{3 \sin(3t)}{5} \right) \\ -c_2 \sin(3t) + c_1 \cos(3t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(-c_1+3c_2) \cos(3t)}{5} + \frac{3(c_1+\frac{c_2}{3}) \sin(3t)}{5} \\ -c_2 \sin(3t) + c_1 \cos(3t) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{(-c_1+3c_2) \cos(3t)}{5} + \frac{3(c_1+\frac{c_2}{3}) \sin(3t)}{5}, x_2(t) = -c_2 \sin(3t) + c_1 \cos(3t) \right\}$$

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 50

```
dsolve([diff(x__1(t),t)=1*x__1(t)+2*x__2(t),diff(x__2(t),t)=-5*x__1(t)-1*x__2(t)],singsol=all)
```

$$x_1(t) = c_1 \sin(3t) + c_2 \cos(3t)$$

$$x_2(t) = \frac{3c_1 \cos(3t)}{2} - \frac{3c_2 \sin(3t)}{2} - \frac{c_1 \sin(3t)}{2} - \frac{c_2 \cos(3t)}{2}$$

✓ Solution by Mathematica

Time used: 0.004 (sec). Leaf size: 54

```
DSolve[{x1'[t]==1*x1[t]+2*x2[t],x2'[t]==-5*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolutions->True]
```

$$x_1(t) \rightarrow c_1 \cos(3t) + \frac{1}{3}(c_1 + 2c_2) \sin(3t)$$

$$x_2(t) \rightarrow c_2 \cos(3t) - \frac{1}{3}(5c_1 + c_2) \sin(3t)$$

16.7 problem 7

- 16.7.1 Solution using Matrix exponential method 3582
- 16.7.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3583
- 16.7.3 Maple step by step solution 3591

Internal problem ID [757]

Internal file name [OUTPUT/757_Sunday_June_05_2022_01_48_52_AM_27514033/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 7.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) \\x_2'(t) &= 2x_1(t) + x_2(t) - 2x_3(t) \\x_3'(t) &= 3x_1(t) + 2x_2(t) + x_3(t)\end{aligned}$$

16.7.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 2 & 1 & -2 \\ 3 & 2 & 1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned}
 e^{At} &= \begin{bmatrix} e^t & 0 & 0 \\ \frac{3e^t \cos(2t)}{2} + e^t \sin(2t) - \frac{3e^t}{2} & e^t \cos(2t) & -e^t \sin(2t) \\ -e^t \cos(2t) + \frac{3e^t \sin(2t)}{2} + e^t & e^t \sin(2t) & e^t \cos(2t) \end{bmatrix} \\
 &= \begin{bmatrix} e^t & 0 & 0 \\ \frac{e^t(-3+3\cos(2t)+2\sin(2t))}{2} & e^t \cos(2t) & -e^t \sin(2t) \\ -\frac{e^t(-2+2\cos(2t)-3\sin(2t))}{2} & e^t \sin(2t) & e^t \cos(2t) \end{bmatrix}
 \end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^t & 0 & 0 \\ \frac{e^t(-3+3\cos(2t)+2\sin(2t))}{2} & e^t \cos(2t) & -e^t \sin(2t) \\ -\frac{e^t(-2+2\cos(2t)-3\sin(2t))}{2} & e^t \sin(2t) & e^t \cos(2t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ c_3 \end{bmatrix} \\
 &= \begin{bmatrix} e^t c_1 \\ \frac{e^t(-3+3\cos(2t)+2\sin(2t))c_1}{2} + e^t \cos(2t) c_2 - e^t \sin(2t) c_3 \\ -\frac{e^t(-2+2\cos(2t)-3\sin(2t))c_1}{2} + e^t \sin(2t) c_2 + e^t \cos(2t) c_3 \end{bmatrix} \\
 &= \begin{bmatrix} e^t c_1 \\ \frac{3e^t \left((c_1 + \frac{2c_2}{3}) \cos(2t) + \frac{2(c_1 - c_3) \sin(2t)}{3} - c_1 \right)}{2} \\ -\left((c_1 - c_3) \cos(2t) + \left(-\frac{3c_1}{2} - c_2 \right) \sin(2t) - c_1 \right) e^t \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.7.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 2 & 1 & -2 \\ 3 & 2 & 1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 1 & 0 & 0 \\ 2 & 1 & -2 \\ 3 & 2 & 1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 1 - \lambda & 0 & 0 \\ 2 & 1 - \lambda & -2 \\ 3 & 2 & 1 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^3 - 3\lambda^2 + 7\lambda - 5 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 1$$

$$\lambda_2 = 1 + 2i$$

$$\lambda_3 = 1 - 2i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
1	1	real eigenvalue
$1 + 2i$	1	complex eigenvalue
$1 - 2i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = 1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 0 & 0 \\ 2 & 1 & -2 \\ 3 & 2 & 1 \end{bmatrix} - (1) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 & 0 & 0 \\ 2 & 0 & -2 \\ 3 & 2 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} 0 & 0 & 0 & 0 \\ 2 & 0 & -2 & 0 \\ 3 & 2 & 0 & 0 \end{array} \right]$$

Since the current pivot $A(1,1)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 1 and row 2 gives

$$\left[\begin{array}{ccc|c} 2 & 0 & -2 & 0 \\ 0 & 0 & 0 & 0 \\ 3 & 2 & 0 & 0 \end{array} \right]$$

$$R_3 = R_3 - \frac{3R_1}{2} \implies \left[\begin{array}{ccc|c} 2 & 0 & -2 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 2 & 3 & 0 \end{array} \right]$$

Since the current pivot $A(2,2)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 2 and row 3 gives

$$\left[\begin{array}{ccc|c} 2 & 0 & -2 & 0 \\ 0 & 2 & 3 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 & 0 & -2 \\ 0 & 2 & 3 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t, v_2 = -\frac{3t}{2}\}$

Hence the solution is

$$\begin{bmatrix} t \\ -\frac{3t}{2} \\ t \end{bmatrix} = \begin{bmatrix} t \\ -\frac{3t}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ -\frac{3t}{2} \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ -\frac{3}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ -\frac{3t}{2} \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ -\frac{3}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} t \\ -\frac{3t}{2} \\ t \end{bmatrix} = \begin{bmatrix} 2 \\ -3 \\ 2 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 1 - 2i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 0 & 0 \\ 2 & 1 & -2 \\ 3 & 2 & 1 \end{bmatrix} - (1 - 2i) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2i & 0 & 0 \\ 2 & 2i & -2 \\ 3 & 2 & 2i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} 2i & 0 & 0 & 0 \\ 2 & 2i & -2 & 0 \\ 3 & 2 & 2i & 0 \end{array} \right]$$

$$R_2 = iR_1 + R_2 \implies \left[\begin{array}{ccc|c} 2i & 0 & 0 & 0 \\ 0 & 2i & -2 & 0 \\ 3 & 2 & 2i & 0 \end{array} \right]$$

$$R_3 = R_3 + \frac{3iR_1}{2} \implies \left[\begin{array}{ccc|c} 2i & 0 & 0 & 0 \\ 0 & 2i & -2 & 0 \\ 0 & 2 & 2i & 0 \end{array} \right]$$

$$R_3 = iR_2 + R_3 \implies \left[\begin{array}{ccc|c} 2i & 0 & 0 & 0 \\ 0 & 2i & -2 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2i & 0 & 0 \\ 0 & 2i & -2 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 0, v_2 = -it\}$

Hence the solution is

$$\begin{bmatrix} 0 \\ -it \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ -it \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 0 \\ -it \\ t \end{bmatrix} = t \begin{bmatrix} 0 \\ -i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 0 \\ -1t \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ -i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_3 = 1 + 2i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 0 & 0 \\ 2 & 1 & -2 \\ 3 & 2 & 1 \end{bmatrix} - (1 + 2i) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -2i & 0 & 0 \\ 2 & -2i & -2 \\ 3 & 2 & -2i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -2i & 0 & 0 & 0 \\ 2 & -2i & -2 & 0 \\ 3 & 2 & -2i & 0 \end{array} \right]$$

$$R_2 = -iR_1 + R_2 \implies \left[\begin{array}{ccc|c} -2i & 0 & 0 & 0 \\ 0 & -2i & -2 & 0 \\ 3 & 2 & -2i & 0 \end{array} \right]$$

$$R_3 = R_3 - \frac{3iR_1}{2} \implies \left[\begin{array}{ccc|c} -2i & 0 & 0 & 0 \\ 0 & -2i & -2 & 0 \\ 0 & 2 & -2i & 0 \end{array} \right]$$

$$R_3 = -iR_2 + R_3 \implies \left[\begin{array}{ccc|c} -2i & 0 & 0 & 0 \\ 0 & -2i & -2 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -2i & 0 & 0 \\ 0 & -2i & -2 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 0, v_2 = it\}$

Hence the solution is

$$\begin{bmatrix} 0 \\ It \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ it \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 0 \\ It \\ t \end{bmatrix} = t \begin{bmatrix} 0 \\ i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 0 \\ It \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
1	1	1	No	$\begin{bmatrix} 1 \\ -\frac{3}{2} \\ 1 \end{bmatrix}$
$1 + 2i$	1	1	No	$\begin{bmatrix} 0 \\ i \\ 1 \end{bmatrix}$
$1 - 2i$	1	1	No	$\begin{bmatrix} 0 \\ -i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue 1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^t \\ &= \begin{bmatrix} 1 \\ -\frac{3}{2} \\ 1 \end{bmatrix} e^t \end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = c_1 \begin{bmatrix} e^t \\ -\frac{3e^t}{2} \\ e^t \end{bmatrix} + c_2 \begin{bmatrix} 0 \\ ie^{(1+2i)t} \\ e^{(1+2i)t} \end{bmatrix} + c_3 \begin{bmatrix} 0 \\ -ie^{(1-2i)t} \\ e^{(1-2i)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} c_1 e^t \\ -\frac{3c_1 e^t}{2} + i c_2 e^{(1+2i)t} - i c_3 e^{(1-2i)t} \\ c_1 e^t + c_2 e^{(1+2i)t} + c_3 e^{(1-2i)t} \end{bmatrix}$$

16.7.3 Maple step by step solution

Let's solve

$$[x_1'(t) = x_1(t), x_2'(t) = 2x_1(t) + x_2(t) - 2x_3(t), x_3'(t) = 3x_1(t) + 2x_2(t) + x_3(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 1 & 0 & 0 \\ 2 & 1 & -2 \\ 3 & 2 & 1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 1 & 0 & 0 \\ 2 & 1 & -2 \\ 3 & 2 & 1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 1 & 0 & 0 \\ 2 & 1 & -2 \\ 3 & 2 & 1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow}'(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[1, \begin{bmatrix} 1 \\ -\frac{3}{2} \\ 1 \end{bmatrix} \right], \left[1 - 2I, \begin{bmatrix} 0 \\ -I \\ 1 \end{bmatrix} \right], \left[1 + 2I, \begin{bmatrix} 0 \\ I \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[1, \begin{bmatrix} 1 \\ -\frac{3}{2} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^t \cdot \begin{bmatrix} 1 \\ -\frac{3}{2} \\ 1 \end{bmatrix}$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[1 - 2I, \begin{bmatrix} 0 \\ -I \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(1-2I)t} \cdot \begin{bmatrix} 0 \\ -I \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^t \cdot (\cos(2t) - I \sin(2t)) \cdot \begin{bmatrix} 0 \\ -I \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^t \cdot \begin{bmatrix} 0 \\ -I(\cos(2t) - I \sin(2t)) \\ \cos(2t) - I \sin(2t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\vec{x}_2(t) = e^t \cdot \begin{bmatrix} 0 \\ -\sin(2t) \\ \cos(2t) \end{bmatrix}, \vec{x}_3(t) = e^t \cdot \begin{bmatrix} 0 \\ -\cos(2t) \\ -\sin(2t) \end{bmatrix}$$

- General solution to the system of ODEs

$$\vec{x} = c_1 \vec{x}_1 + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

- Substitute solutions into the general solution

$$\vec{x} = c_1 e^t \cdot \begin{bmatrix} 1 \\ -\frac{3}{2} \\ 1 \end{bmatrix} + c_2 e^t \cdot \begin{bmatrix} 0 \\ -\sin(2t) \\ \cos(2t) \end{bmatrix} + c_3 e^t \cdot \begin{bmatrix} 0 \\ -\cos(2t) \\ -\sin(2t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} c_1 e^t \\ -\frac{e^t(2c_3 \cos(2t) + 2c_2 \sin(2t) + 3c_1)}{2} \\ e^t(c_1 + c_2 \cos(2t) - c_3 \sin(2t)) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = c_1 e^t, x_2(t) = -\frac{e^t(2c_3 \cos(2t) + 2c_2 \sin(2t) + 3c_1)}{2}, x_3(t) = e^t(c_1 + c_2 \cos(2t) - c_3 \sin(2t)) \right\}$$

✓ Solution by Maple

Time used: 0.109 (sec). Leaf size: 73

```
dsolve([diff(x__1(t),t)=1*x__1(t)+0*x__2(t)+0*x__3(t),diff(x__2(t),t)=2*x__1(t)+1*x__2(t)-2*
```

$$\begin{aligned} x_1(t) &= c_3 e^t \\ x_2(t) &= \frac{e^t(2c_1 \cos(2t) - 3c_3 \cos(2t) + 2c_2 \sin(2t) - 3c_3)}{2} \\ x_3(t) &= -\frac{e^t(2c_2 \cos(2t) - 2c_1 \sin(2t) + 3c_3 \sin(2t) - 2c_3)}{2} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.007 (sec). Leaf size: 95

```
DSolve[{x1'[t]==1*x1[t]+0*x2[t]+0*x3[t],x2'[t]==2*x1[t]+1*x2[t]-2*x3[t],x3'[t]==3*x1[t]+2*x2[t]}
```

$$x1(t) \rightarrow c_1 e^t$$

$$x2(t) \rightarrow \frac{1}{2} e^t ((3c_1 + 2c_2) \cos(2t) + 2(c_1 - c_3) \sin(2t) - 3c_1)$$

$$x3(t) \rightarrow \frac{1}{2} e^t (-2(c_1 - c_3) \cos(2t) + (3c_1 + 2c_2) \sin(2t) + 2c_1)$$

16.8 problem 8

16.8.1 Solution using Matrix exponential method	3595
16.8.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3596
16.8.3 Maple step by step solution	3604

Internal problem ID [758]

Internal file name [OUTPUT/758_Sunday_June_05_2022_01_48_54_AM_69288655/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 8.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x_1'(t) = -3x_1(t) + 2x_3(t)$$

$$x_2'(t) = x_1(t) - x_2(t)$$

$$x_3'(t) = -2x_1(t) - x_2(t)$$

16.8.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} -3 & 0 & 2 \\ 1 & -1 & 0 \\ -2 & -1 & 0 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{2e^{-2t}}{3} + \frac{e^{-t}\cos(\sqrt{2}t)}{3} - \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} & -\frac{\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} + \frac{2e^{-t}\cos(\sqrt{2}t)}{3} - \frac{2e^{-2t}}{3} & \frac{2e^{-t}\cos(\sqrt{2}t)}{3} + \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} \\ \frac{2e^{-t}\cos(\sqrt{2}t)}{3} + \frac{\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{6} - \frac{2e^{-2t}}{3} & \frac{2e^{-2t}}{3} + \frac{e^{-t}\cos(\sqrt{2}t)}{3} + \frac{\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} & -\frac{2e^{-t}\cos(\sqrt{2}t)}{3} + \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} \\ -\frac{e^{-t}\cos(\sqrt{2}t)}{3} + \frac{e^{-2t}}{3} - \frac{5\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{6} & -\frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} + \frac{e^{-t}\cos(\sqrt{2}t)}{3} - \frac{e^{-2t}}{3} & \frac{4e^{-t}\cos(\sqrt{2}t)}{3} + \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At}\vec{c} \\ &= \begin{bmatrix} \frac{2e^{-2t}}{3} + \frac{e^{-t}\cos(\sqrt{2}t)}{3} - \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} & -\frac{\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} + \frac{2e^{-t}\cos(\sqrt{2}t)}{3} - \frac{2e^{-2t}}{3} & \frac{2e^{-t}\cos(\sqrt{2}t)}{3} + \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} \\ \frac{2e^{-t}\cos(\sqrt{2}t)}{3} + \frac{\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{6} - \frac{2e^{-2t}}{3} & \frac{2e^{-2t}}{3} + \frac{e^{-t}\cos(\sqrt{2}t)}{3} + \frac{\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} & -\frac{2e^{-t}\cos(\sqrt{2}t)}{3} + \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} \\ -\frac{e^{-t}\cos(\sqrt{2}t)}{3} + \frac{e^{-2t}}{3} - \frac{5\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{6} & -\frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} + \frac{e^{-t}\cos(\sqrt{2}t)}{3} - \frac{e^{-2t}}{3} & \frac{4e^{-t}\cos(\sqrt{2}t)}{3} + \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} \end{bmatrix} \\ &= \begin{bmatrix} \left(\frac{2e^{-2t}}{3} + \frac{e^{-t}\cos(\sqrt{2}t)}{3} - \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3}\right) c_1 + \left(-\frac{\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} + \frac{2e^{-t}\cos(\sqrt{2}t)}{3} - \frac{2e^{-2t}}{3}\right) c_2 + \left(\frac{2e^{-t}\cos(\sqrt{2}t)}{3} + \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3}\right) c_3 \\ \left(\frac{2e^{-t}\cos(\sqrt{2}t)}{3} + \frac{\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{6} - \frac{2e^{-2t}}{3}\right) c_1 + \left(\frac{2e^{-2t}}{3} + \frac{e^{-t}\cos(\sqrt{2}t)}{3} + \frac{\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3}\right) c_2 + \left(-\frac{2e^{-t}\cos(\sqrt{2}t)}{3} + \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3}\right) c_3 \\ \left(-\frac{e^{-t}\cos(\sqrt{2}t)}{3} + \frac{e^{-2t}}{3} - \frac{5\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{6}\right) c_1 + \left(-\frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3} + \frac{e^{-t}\cos(\sqrt{2}t)}{3} - \frac{e^{-2t}}{3}\right) c_2 + \left(\frac{4e^{-t}\cos(\sqrt{2}t)}{3} + \frac{2\sqrt{2}e^{-t}\sin(\sqrt{2}t)}{3}\right) c_3 \end{bmatrix} \\ &= \begin{bmatrix} \frac{e^{-t}(c_1+2c_2+2c_3)\cos(\sqrt{2}t)}{3} - \frac{2e^{-t}(c_1+\frac{c_2}{2}-c_3)\sqrt{2}\sin(\sqrt{2}t)}{3} + \frac{2e^{-2t}(c_1-c_2-c_3)}{3} \\ \frac{2e^{-t}(c_1+\frac{c_2}{2}-c_3)\cos(\sqrt{2}t)}{3} + \frac{\sqrt{2}e^{-t}(c_1+2c_2+2c_3)\sin(\sqrt{2}t)}{6} - \frac{2e^{-2t}(c_1-c_2-c_3)}{3} \\ -\frac{e^{-t}(c_1-c_2-4c_3)\cos(\sqrt{2}t)}{3} - \frac{5(c_1+\frac{4c_2}{5}-\frac{2c_3}{5})e^{-t}\sqrt{2}\sin(\sqrt{2}t)}{6} + \frac{e^{-2t}(c_1-c_2-c_3)}{3} \end{bmatrix} \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.8.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} -3 & 0 & 2 \\ 1 & -1 & 0 \\ -2 & -1 & 0 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} -3 & 0 & 2 \\ 1 & -1 & 0 \\ -2 & -1 & 0 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -3 - \lambda & 0 & 2 \\ 1 & -1 - \lambda & 0 \\ -2 & -1 & -\lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^3 + 4\lambda^2 + 7\lambda + 6 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1 + i\sqrt{2}$$

$$\lambda_2 = -1 - i\sqrt{2}$$

$$\lambda_3 = -2$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-1 + i\sqrt{2}$	1	complex eigenvalue
-2	1	real eigenvalue
$-1 - i\sqrt{2}$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -3 & 0 & 2 \\ 1 & -1 & 0 \\ -2 & -1 & 0 \end{bmatrix} - (-2) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 & 0 & 2 \\ 1 & 1 & 0 \\ -2 & -1 & 2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -1 & 0 & 2 & 0 \\ 1 & 1 & 0 & 0 \\ -2 & -1 & 2 & 0 \end{array} \right]$$

$$R_2 = R_2 + R_1 \implies \left[\begin{array}{ccc|c} -1 & 0 & 2 & 0 \\ 0 & 1 & 2 & 0 \\ -2 & -1 & 2 & 0 \end{array} \right]$$

$$R_3 = R_3 - 2R_1 \implies \left[\begin{array}{ccc|c} -1 & 0 & 2 & 0 \\ 0 & 1 & 2 & 0 \\ 0 & -1 & -2 & 0 \end{array} \right]$$

$$R_3 = R_3 + R_2 \implies \left[\begin{array}{ccc|c} -1 & 0 & 2 & 0 \\ 0 & 1 & 2 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 & 0 & 2 \\ 0 & 1 & 2 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 2t, v_2 = -2t\}$

Hence the solution is

$$\begin{bmatrix} 2t \\ -2t \\ t \end{bmatrix} = \begin{bmatrix} 2t \\ -2t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 2t \\ -2t \\ t \end{bmatrix} = t \begin{bmatrix} 2 \\ -2 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 2t \\ -2t \\ t \end{bmatrix} = \begin{bmatrix} 2 \\ -2 \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -1 - i\sqrt{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -3 & 0 & 2 \\ 1 & -1 & 0 \\ -2 & -1 & 0 \end{bmatrix} - (-1 - i\sqrt{2}) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -2 + i\sqrt{2} & 0 & 2 \\ 1 & i\sqrt{2} & 0 \\ -2 & -1 & 1 + i\sqrt{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -2 + i\sqrt{2} & 0 & 2 & 0 \\ 1 & i\sqrt{2} & 0 & 0 \\ -2 & -1 & 1 + i\sqrt{2} & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{R_1}{-2 + i\sqrt{2}} \Rightarrow \left[\begin{array}{ccc|c} -2 + i\sqrt{2} & 0 & 2 & 0 \\ 0 & i\sqrt{2} & -\frac{2}{-2+i\sqrt{2}} & 0 \\ -2 & -1 & 1 + i\sqrt{2} & 0 \end{array} \right]$$

$$R_3 = R_3 + \frac{2R_1}{-2 + i\sqrt{2}} \Rightarrow \left[\begin{array}{ccc|c} -2 + i\sqrt{2} & 0 & 2 & 0 \\ 0 & i\sqrt{2} & -\frac{2}{-2+i\sqrt{2}} & 0 \\ 0 & -1 & -\frac{\sqrt{2}}{2i+\sqrt{2}} & 0 \end{array} \right]$$

$$R_3 = R_3 - \frac{i\sqrt{2}R_2}{2} \Rightarrow \left[\begin{array}{ccc|c} -2 + i\sqrt{2} & 0 & 2 & 0 \\ 0 & i\sqrt{2} & -\frac{2}{-2+i\sqrt{2}} & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{ccc} -2 + i\sqrt{2} & 0 & 2 \\ 0 & i\sqrt{2} & -\frac{2}{-2+i\sqrt{2}} \\ 0 & 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\left\{ v_1 = -\frac{2t}{-2+i\sqrt{2}}, v_2 = -\frac{t}{1+i\sqrt{2}} \right\}$

Hence the solution is

$$\begin{bmatrix} -\frac{2t}{-2+i\sqrt{2}} \\ -\frac{t}{1+i\sqrt{2}} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{2t}{-2+i\sqrt{2}} \\ -\frac{t}{1+i\sqrt{2}} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{2t}{-2+i\sqrt{2}} \\ -\frac{t}{1+i\sqrt{2}} \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{2}{-2+i\sqrt{2}} \\ -\frac{1}{1+i\sqrt{2}} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{2t}{-2+I\sqrt{2}} \\ -\frac{t}{1+I\sqrt{2}} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{2}{-2+i\sqrt{2}} \\ -\frac{1}{1+i\sqrt{2}} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{2t}{-2+I\sqrt{2}} \\ -\frac{t}{1+I\sqrt{2}} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{2}{-2+i\sqrt{2}} \\ -\frac{1}{1+i\sqrt{2}} \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_3 = -1 + i\sqrt{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -3 & 0 & 2 \\ 1 & -1 & 0 \\ -2 & -1 & 0 \end{bmatrix} - (-1 + i\sqrt{2}) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -2 - i\sqrt{2} & 0 & 2 \\ 1 & -i\sqrt{2} & 0 \\ -2 & -1 & 1 - i\sqrt{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -2 - i\sqrt{2} & 0 & 2 & 0 \\ 1 & -i\sqrt{2} & 0 & 0 \\ -2 & -1 & 1 - i\sqrt{2} & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{R_1}{-2 - i\sqrt{2}} \implies \left[\begin{array}{ccc|c} -2 - i\sqrt{2} & 0 & 2 & 0 \\ 0 & -i\sqrt{2} & \frac{2}{2+i\sqrt{2}} & 0 \\ -2 & -1 & 1 - i\sqrt{2} & 0 \end{array} \right]$$

$$R_3 = R_3 + \frac{2R_1}{-2 - i\sqrt{2}} \implies \left[\begin{array}{ccc|c} -2 - i\sqrt{2} & 0 & 2 & 0 \\ 0 & -i\sqrt{2} & \frac{2}{2+i\sqrt{2}} & 0 \\ 0 & -1 & \frac{\sqrt{2}}{2i-\sqrt{2}} & 0 \end{array} \right]$$

$$R_3 = R_3 + \frac{i\sqrt{2}}{2} R_2 \implies \left[\begin{array}{ccc|c} -2 - i\sqrt{2} & 0 & 2 & 0 \\ 0 & -i\sqrt{2} & \frac{2}{2+i\sqrt{2}} & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{ccc} -2 - i\sqrt{2} & 0 & 2 \\ 0 & -i\sqrt{2} & \frac{2}{2+i\sqrt{2}} \\ 0 & 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\left\{ v_1 = \frac{2t}{2+i\sqrt{2}}, v_2 = \frac{t}{-1+i\sqrt{2}} \right\}$

Hence the solution is

$$\begin{bmatrix} \frac{2t}{2+i\sqrt{2}} \\ \frac{t}{-1+i\sqrt{2}} \\ t \end{bmatrix} = \begin{bmatrix} \frac{2t}{2+i\sqrt{2}} \\ \frac{t}{-1+i\sqrt{2}} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{2t}{2+i\sqrt{2}} \\ \frac{t}{-1+i\sqrt{2}} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{2}{2+i\sqrt{2}} \\ \frac{1}{-1+i\sqrt{2}} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{2t}{2+i\sqrt{2}} \\ \frac{t}{-1+i\sqrt{2}} \\ t \end{bmatrix} = \begin{bmatrix} \frac{2}{2+i\sqrt{2}} \\ \frac{1}{-1+i\sqrt{2}} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{2t}{2+i\sqrt{2}} \\ \frac{t}{-1+i\sqrt{2}} \\ t \end{bmatrix} = \begin{bmatrix} \frac{2}{2+i\sqrt{2}} \\ \frac{1}{-1+i\sqrt{2}} \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-1 + i\sqrt{2}$	1	1	No	$\begin{bmatrix} \frac{2}{2+i\sqrt{2}} \\ \frac{i(-2+i\sqrt{2})\sqrt{2}}{6} \\ 1 \end{bmatrix}$
$-1 - i\sqrt{2}$	1	1	No	$\begin{bmatrix} \frac{2}{2-i\sqrt{2}} \\ -\frac{i(-2-i\sqrt{2})\sqrt{2}}{6} \\ 1 \end{bmatrix}$
-2	1	1	No	$\begin{bmatrix} 2 \\ -2 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^{-2t} \\ &= \begin{bmatrix} 2 \\ -2 \\ 1 \end{bmatrix} e^{-2t} \end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{2e^{(-1+i\sqrt{2})t}}{2+i\sqrt{2}} \\ \frac{ie^{(-1+i\sqrt{2})t}(-2+i\sqrt{2})\sqrt{2}}{6} \\ e^{(-1+i\sqrt{2})t} \end{bmatrix} + c_2 \begin{bmatrix} \frac{2e^{(-1-i\sqrt{2})t}}{2-i\sqrt{2}} \\ -\frac{ie^{(-1-i\sqrt{2})t}(-2-i\sqrt{2})\sqrt{2}}{6} \\ e^{(-1-i\sqrt{2})t} \end{bmatrix} + c_3 \begin{bmatrix} 2e^{-2t} \\ -2e^{-2t} \\ e^{-2t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} \frac{(2+i\sqrt{2})c_2e^{-(1+i\sqrt{2})t}}{3} + \frac{c_1(2-i\sqrt{2})e^{(-1+i\sqrt{2})t}}{3} + 2c_3e^{-2t} \\ \frac{(-1+i\sqrt{2})c_2e^{-(1+i\sqrt{2})t}}{3} + \frac{c_1(-1-i\sqrt{2})e^{(-1+i\sqrt{2})t}}{3} - 2c_3e^{-2t} \\ c_1e^{(-1+i\sqrt{2})t} + c_2e^{-(1+i\sqrt{2})t} + c_3e^{-2t} \end{bmatrix}$$

16.8.3 Maple step by step solution

Let's solve

$$[x_1'(t) = -3x_1(t) + 2x_3(t), x_2'(t) = x_1(t) - x_2(t), x_3'(t) = -2x_1(t) - x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} -3 & 0 & 2 \\ 1 & -1 & 0 \\ -2 & -1 & 0 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} -3 & 0 & 2 \\ 1 & -1 & 0 \\ -2 & -1 & 0 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -3 & 0 & 2 \\ 1 & -1 & 0 \\ -2 & -1 & 0 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-2, \begin{bmatrix} 2 \\ -2 \\ 1 \end{bmatrix} \right], \left[-1 - I\sqrt{2}, \begin{bmatrix} \frac{2}{2-I\sqrt{2}} \\ -\frac{1}{6}(-2 - I\sqrt{2})\sqrt{2} \\ 1 \end{bmatrix} \right], \left[-1 + I\sqrt{2}, \begin{bmatrix} \frac{2}{2+I\sqrt{2}} \\ \frac{1}{6}(-2 + I\sqrt{2})\sqrt{2} \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-2, \begin{bmatrix} 2 \\ -2 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{-2t} \cdot \begin{bmatrix} 2 \\ -2 \\ 1 \end{bmatrix}$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-1 - I\sqrt{2}, \begin{bmatrix} \frac{2}{2-I\sqrt{2}} \\ -\frac{1}{6}(-2 - I\sqrt{2})\sqrt{2} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(-1-I\sqrt{2})t} \cdot \begin{bmatrix} \frac{2}{2-I\sqrt{2}} \\ -\frac{1}{6}(-2 - I\sqrt{2})\sqrt{2} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{-t} \cdot (\cos(\sqrt{2}t) - I \sin(\sqrt{2}t)) \cdot \begin{bmatrix} \frac{2}{2-I\sqrt{2}} \\ -\frac{1}{6}(-2 - I\sqrt{2})\sqrt{2} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{-t} \cdot \begin{bmatrix} \frac{2(\cos(\sqrt{2}t) - I \sin(\sqrt{2}t))}{2-I\sqrt{2}} \\ -\frac{1}{6}(\cos(\sqrt{2}t) - I \sin(\sqrt{2}t))(-2 - I\sqrt{2})\sqrt{2} \\ \cos(\sqrt{2}t) - I \sin(\sqrt{2}t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\vec{x}_2(t) = e^{-t} \cdot \begin{bmatrix} \frac{2 \cos(\sqrt{2}t)}{3} + \frac{\sqrt{2} \sin(\sqrt{2}t)}{3} \\ \frac{\sqrt{2}(-\cos(\sqrt{2}t)\sqrt{2} + 2\sin(\sqrt{2}t))}{6} \\ \cos(\sqrt{2}t) \end{bmatrix}, \vec{x}_3(t) = e^{-t} \cdot \begin{bmatrix} \frac{\cos(\sqrt{2}t)\sqrt{2}}{3} - \frac{2 \sin(\sqrt{2}t)}{3} \\ \frac{\sqrt{2}(-2\cos(\sqrt{2}t) - \sqrt{2}\sin(\sqrt{2}t))}{6} \\ -\sin(\sqrt{2}t) \end{bmatrix}$$

- General solution to the system of ODEs

$$\vec{x} = c_1 \vec{x}_1 + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

- Substitute solutions into the general solution

$$\vec{x} = c_1 e^{-2t} \cdot \begin{bmatrix} 2 \\ -2 \\ 1 \end{bmatrix} + c_2 e^{-t} \cdot \begin{bmatrix} \frac{2 \cos(\sqrt{2}t)}{3} + \frac{\sqrt{2} \sin(\sqrt{2}t)}{3} \\ \frac{\sqrt{2}(-\cos(\sqrt{2}t)\sqrt{2} + 2\sin(\sqrt{2}t))}{6} \\ \cos(\sqrt{2}t) \end{bmatrix} + c_3 e^{-t} \cdot \begin{bmatrix} \frac{\cos(\sqrt{2}t)\sqrt{2}}{3} - \frac{2 \sin(\sqrt{2}t)}{3} \\ \frac{\sqrt{2}(-2\cos(\sqrt{2}t) - \sqrt{2}\sin(\sqrt{2}t))}{6} \\ -\sin(\sqrt{2}t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} \frac{2(\frac{c_3\sqrt{2}}{2} + c_2)e^{-t}\cos(\sqrt{2}t)}{3} + \frac{e^{-t}(\sqrt{2}c_2 - 2c_3)\sin(\sqrt{2}t)}{3} + 2c_1e^{-2t} \\ -\frac{e^{-t}(-c_3\sqrt{2} + c_2)\cos(\sqrt{2}t)}{3} + \frac{e^{-t}(\sqrt{2}c_2 + c_3)\sin(\sqrt{2}t)}{3} - 2c_1e^{-2t} \\ c_1e^{-2t} + c_2e^{-t}\cos(\sqrt{2}t) - c_3e^{-t}\sin(\sqrt{2}t) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ \begin{aligned} x_1(t) &= \frac{2(\frac{c_3\sqrt{2}}{2} + c_2)e^{-t}\cos(\sqrt{2}t)}{3} + \frac{e^{-t}(\sqrt{2}c_2 - 2c_3)\sin(\sqrt{2}t)}{3} + 2c_1e^{-2t}, \\ x_2(t) &= -\frac{e^{-t}(-c_3\sqrt{2} + c_2)\cos(\sqrt{2}t)}{3} + \frac{e^{-t}(\sqrt{2}c_2 + c_3)\sin(\sqrt{2}t)}{3} - 2c_1e^{-2t} \end{aligned} \right.$$

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 146

```
dsolve([diff(x__1(t),t)=-3*x__1(t)+0*x__2(t)+2*x__3(t),diff(x__2(t),t)=1*x__1(t)-1*x__2(t)-0
```

$$\begin{aligned}
 x_1(t) &= c_1 e^{-2t} + c_2 e^{-t} \sin(\sqrt{2}t) + c_3 e^{-t} \cos(\sqrt{2}t) \\
 x_2(t) &= -c_1 e^{-2t} - \frac{c_2 e^{-t} \sqrt{2} \cos(\sqrt{2}t)}{2} + \frac{c_3 e^{-t} \sqrt{2} \sin(\sqrt{2}t)}{2} \\
 x_3(t) &= \frac{c_1 e^{-2t}}{2} + c_2 e^{-t} \sin(\sqrt{2}t) + \frac{c_2 e^{-t} \sqrt{2} \cos(\sqrt{2}t)}{2} \\
 &\quad + c_3 e^{-t} \cos(\sqrt{2}t) - \frac{c_3 e^{-t} \sqrt{2} \sin(\sqrt{2}t)}{2}
 \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.038 (sec). Leaf size: 235

```
DSolve[{x1'[t]==-3*x1[t]+0*x2[t]+2*x3[t],x2'[t]==1*x1[t]-1*x2[t]-0*x3[t],x3'[t]==-2*x1[t]-1*
```

$$\begin{aligned}
 x_1(t) &\rightarrow \frac{1}{3} e^{-2t} \left((c_1 + 2(c_2 + c_3)) e^t \cos(\sqrt{2}t) - \sqrt{2} (2c_1 + c_2 - 2c_3) e^t \sin(\sqrt{2}t) \right. \\
 &\quad \left. + 2(c_1 - c_2 - c_3) \right) \\
 x_2(t) &\rightarrow \frac{1}{6} e^{-2t} \left(2(2c_1 + c_2 - 2c_3) e^t \cos(\sqrt{2}t) + \sqrt{2} (c_1 + 2(c_2 + c_3)) e^t \sin(\sqrt{2}t) \right. \\
 &\quad \left. + 4(-c_1 + c_2 + c_3) \right) \\
 x_3(t) &\rightarrow \frac{1}{6} e^{-2t} \left(-2(c_1 - c_2 - 4c_3) e^t \cos(\sqrt{2}t) - \sqrt{2} (5c_1 + 4c_2 - 2c_3) e^t \sin(\sqrt{2}t) \right. \\
 &\quad \left. + 2(c_1 - c_2 - c_3) \right)
 \end{aligned}$$

16.9 problem 9

16.9.1 Solution using Matrix exponential method 3608

16.9.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3609

Internal problem ID [759]

Internal file name [OUTPUT/759_Sunday_June_05_2022_01_48_57_AM_24403091/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 9.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x_1'(t) = x_1(t) - 5x_2(t)$$

$$x_2'(t) = x_1(t) - 3x_2(t)$$

With initial conditions

$$[x_1(0) = 1, x_2(0) = 1]$$

16.9.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned} e^{At} &= \begin{bmatrix} e^{-t} \cos(t) + 2e^{-t} \sin(t) & -5e^{-t} \sin(t) \\ e^{-t} \sin(t) & e^{-t} \cos(t) - 2e^{-t} \sin(t) \end{bmatrix} \\ &= \begin{bmatrix} e^{-t}(\cos(t) + 2\sin(t)) & -5e^{-t} \sin(t) \\ e^{-t} \sin(t) & e^{-t}(\cos(t) - 2\sin(t)) \end{bmatrix} \end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{x}_0 \\
 &= \begin{bmatrix} e^{-t}(\cos(t) + 2\sin(t)) & -5e^{-t}\sin(t) \\ e^{-t}\sin(t) & e^{-t}(\cos(t) - 2\sin(t)) \end{bmatrix} \begin{bmatrix} 1 \\ 1 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-t}(\cos(t) + 2\sin(t)) - 5e^{-t}\sin(t) \\ e^{-t}\sin(t) + e^{-t}(\cos(t) - 2\sin(t)) \end{bmatrix} \\
 &= \begin{bmatrix} (-3\sin(t) + \cos(t))e^{-t} \\ e^{-t}(-\sin(t) + \cos(t)) \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.9.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 1 - \lambda & -5 \\ 1 & -3 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 2\lambda + 2 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1 + i$$

$$\lambda_2 = -1 - i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-1 - i$	1	complex eigenvalue
$-1 + i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1 - i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} - (-1 - i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 + i & -5 \\ 1 & -2 + i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 + i & -5 & 0 \\ 1 & -2 + i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} + \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2 + i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 + i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 - i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = \begin{bmatrix} (2 - i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -1 + i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} - (-1 + i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 - i & -5 \\ 1 & -2 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 1 & -2 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} - \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 - i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 + i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = \begin{bmatrix} (2 + i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-1 + i$	1	1	No	$\begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$
$-1 - i$	1	1	No	$\begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (2+i)e^{(-1+i)t} \\ e^{(-1+i)t} \end{bmatrix} + c_2 \begin{bmatrix} (2-i)e^{(-1-i)t} \\ e^{(-1-i)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (2+i)c_1 e^{(-1+i)t} + (2-i)c_2 e^{(-1-i)t} \\ c_1 e^{(-1+i)t} + c_2 e^{(-1-i)t} \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x_1(0) = 1 \\ x_2(0) = 1 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 1 \\ 1 \end{bmatrix} = \begin{bmatrix} (2+i)c_1 + (2-i)c_2 \\ c_1 + c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = \frac{1}{2} + \frac{i}{2} \\ c_2 = \frac{1}{2} - \frac{i}{2} \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \left(\frac{1}{2} + \frac{3i}{2}\right) e^{(-1+i)t} + \left(\frac{1}{2} - \frac{3i}{2}\right) e^{(-1-i)t} \\ \left(\frac{1}{2} + \frac{i}{2}\right) e^{(-1+i)t} + \left(\frac{1}{2} - \frac{i}{2}\right) e^{(-1-i)t} \end{bmatrix}$$

The following is the phase plot of the system.

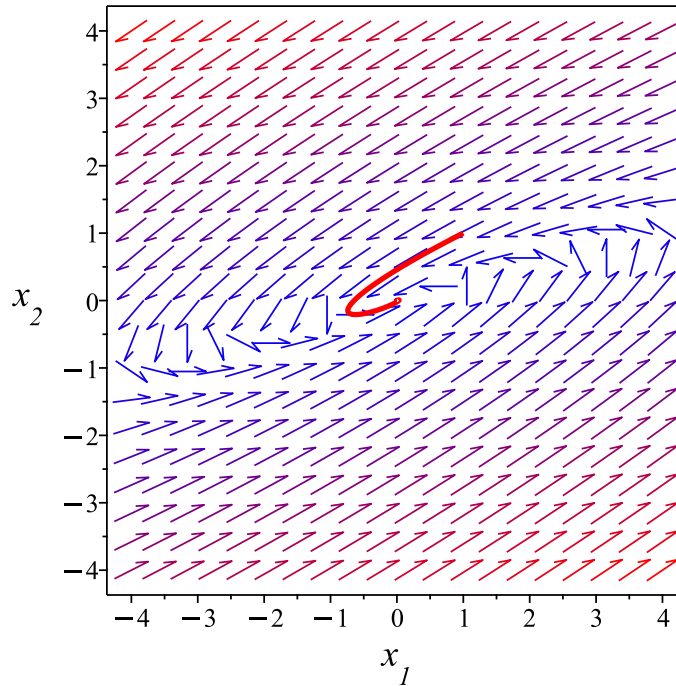


Figure 513: Phase plot

The following are plots of each solution.

✓ Solution by Maple

Time used: 0.125 (sec). Leaf size: 35

```
dsolve([diff(x__1(t),t) = x__1(t)-5*x__2(t), diff(x__2(t),t) = x__1(t)-3*x__2(t), x__1(0) =
```

$$x_1(t) = e^{-t}(-3 \sin(t) + \cos(t))$$

$$x_2(t) = \frac{e^{-t}(5 \cos(t) - 5 \sin(t))}{5}$$

✓ Solution by Mathematica

Time used: 0.006 (sec). Leaf size: 34

```
DSolve[{x1'[t]==1*x1[t]-5*x2[t],x2'[t]==1*x1[t]-3*x2[t]},{x1[0]==1,x2[0]==1},{x1[t],x2[t]},t
```

$$x1(t) \rightarrow e^{-t}(\cos(t) - 3 \sin(t))$$

$$x2(t) \rightarrow e^{-t}(\cos(t) - \sin(t))$$

16.10 problem 10

16.10.1 Solution using Matrix exponential method 3615

16.10.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3616

Internal problem ID [760]

Internal file name [OUTPUT/760_Sunday_June_05_2022_01_48_59_AM_20065990/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 10.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x_1'(t) = -3x_1(t) + 2x_2(t)$$

$$x_2'(t) = -x_1(t) - x_2(t)$$

With initial conditions

$$[x_1(0) = 1, x_2(0) = -2]$$

16.10.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -3 & 2 \\ -1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned} e^{At} &= \begin{bmatrix} e^{-2t} \cos(t) - e^{-2t} \sin(t) & 2e^{-2t} \sin(t) \\ -e^{-2t} \sin(t) & e^{-2t} \cos(t) + e^{-2t} \sin(t) \end{bmatrix} \\ &= \begin{bmatrix} e^{-2t}(-\sin(t) + \cos(t)) & 2e^{-2t} \sin(t) \\ -e^{-2t} \sin(t) & e^{-2t}(\cos(t) + \sin(t)) \end{bmatrix} \end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{x}_0 \\
 &= \begin{bmatrix} e^{-2t}(-\sin(t) + \cos(t)) & 2e^{-2t} \sin(t) \\ -e^{-2t} \sin(t) & e^{-2t}(\cos(t) + \sin(t)) \end{bmatrix} \begin{bmatrix} 1 \\ -2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-2t}(-\sin(t) + \cos(t)) - 4e^{-2t} \sin(t) \\ -e^{-2t} \sin(t) - 2e^{-2t}(\cos(t) + \sin(t)) \end{bmatrix} \\
 &= \begin{bmatrix} e^{-2t}(-5\sin(t) + \cos(t)) \\ e^{-2t}(-3\sin(t) - 2\cos(t)) \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.10.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -3 & 2 \\ -1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -3 & 2 \\ -1 & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -3 - \lambda & 2 \\ -1 & -1 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 4\lambda + 5 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -2 + i$$

$$\lambda_2 = -2 - i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-2 + i$	1	complex eigenvalue
$-2 - i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -2 - i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -3 & 2 \\ -1 & -1 \end{bmatrix} - (-2 - i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 + i & 2 \\ -1 & 1 + i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -1 + i & 2 & 0 \\ -1 & 1 + i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{1}{2} - \frac{i}{2} \right) R_1 \implies \left[\begin{array}{cc|c} -1 + i & 2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 + i & 2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (1 + i)t\}$

Hence the solution is

$$\begin{bmatrix} (1 + i)t \\ t \end{bmatrix} = \begin{bmatrix} (1 + i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (1 + i)t \\ t \end{bmatrix} = t \begin{bmatrix} 1 + i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (1 + i)t \\ t \end{bmatrix} = \begin{bmatrix} 1 + i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -2 + i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -3 & 2 \\ -1 & -1 \end{bmatrix} - (-2 + i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 - i & 2 \\ -1 & 1 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -1 - i & 2 & 0 \\ -1 & 1 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{1}{2} + \frac{i}{2} \right) R_1 \implies \left[\begin{array}{cc|c} -1 - i & 2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 - i & 2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (1 - i)t\}$

Hence the solution is

$$\begin{bmatrix} (1 - i)t \\ t \end{bmatrix} = \begin{bmatrix} (1 - i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (1 - i)t \\ t \end{bmatrix} = t \begin{bmatrix} 1 - i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (1 - i)t \\ t \end{bmatrix} = \begin{bmatrix} 1 - i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-2 + i$	1	1	No	$\begin{bmatrix} 1 - i \\ 1 \end{bmatrix}$
$-2 - i$	1	1	No	$\begin{bmatrix} 1 + i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (1-i)e^{(-2+i)t} \\ e^{(-2+i)t} \end{bmatrix} + c_2 \begin{bmatrix} (1+i)e^{(-2-i)t} \\ e^{(-2-i)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (1-i)c_1 e^{(-2+i)t} + (1+i)c_2 e^{(-2-i)t} \\ c_1 e^{(-2+i)t} + c_2 e^{(-2-i)t} \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x_1(0) = 1 \\ x_2(0) = -2 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 1 \\ -2 \end{bmatrix} = \begin{bmatrix} (1-i)c_1 + (1+i)c_2 \\ c_1 + c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = -1 + \frac{3i}{2} \\ c_2 = -1 - \frac{3i}{2} \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \left(\frac{1}{2} + \frac{5i}{2}\right) e^{(-2+i)t} + \left(\frac{1}{2} - \frac{5i}{2}\right) e^{(-2-i)t} \\ \left(-1 + \frac{3i}{2}\right) e^{(-2+i)t} + \left(-1 - \frac{3i}{2}\right) e^{(-2-i)t} \end{bmatrix}$$

The following is the phase plot of the system.

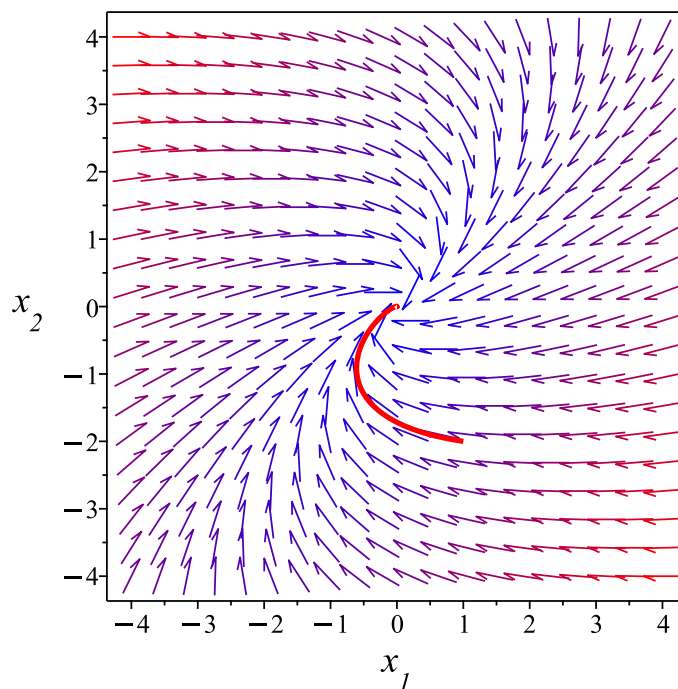


Figure 514: Phase plot

The following are plots of each solution.

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 35

```
dsolve([diff(x__1(t),t) = -3*x__1(t)+2*x__2(t), diff(x__2(t),t) = -x__1(t)-x__2(t), x__1(0)
```

$$x_1(t) = e^{-2t}(-5 \sin(t) + \cos(t))$$

$$x_2(t) = \frac{e^{-2t}(-6 \sin(t) - 4 \cos(t))}{2}$$

✓ Solution by Mathematica

Time used: 0.006 (sec). Leaf size: 27

```
DSolve[{x1'[t]==-3*x1[t]+2*x2[t],x2'[t]==-1*x1[t]-1*x2[t]},{x1[0]==1,x2[0]==1},{x1[t],x2[t]}
```

$$x_1(t) \rightarrow e^{-2t}(\sin(t) + \cos(t))$$

$$x_2(t) \rightarrow e^{-2t} \cos(t)$$

16.11 problem 11

16.11.1 Solution using Matrix exponential method	3622
16.11.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3623
16.11.3 Maple step by step solution	3627

Internal problem ID [761]

Internal file name [OUTPUT/761_Sunday_June_05_2022_01_49_01_AM_57682055/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 11.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$\begin{aligned}x_1'(t) &= \frac{3x_1(t)}{4} - 2x_2(t) \\x_2'(t) &= x_1(t) - \frac{5x_2(t)}{4}\end{aligned}$$

16.11.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} \frac{3}{4} & -2 \\ 1 & -\frac{5}{4} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned}e^{At} &= \begin{bmatrix} e^{-\frac{t}{4}} \cos(t) + e^{-\frac{t}{4}} \sin(t) & -2e^{-\frac{t}{4}} \sin(t) \\ e^{-\frac{t}{4}} \sin(t) & e^{-\frac{t}{4}} \cos(t) - e^{-\frac{t}{4}} \sin(t) \end{bmatrix} \\ &= \begin{bmatrix} e^{-\frac{t}{4}} (\cos(t) + \sin(t)) & -2e^{-\frac{t}{4}} \sin(t) \\ e^{-\frac{t}{4}} \sin(t) & e^{-\frac{t}{4}} (-\sin(t) + \cos(t)) \end{bmatrix}\end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{-\frac{t}{4}}(\cos(t) + \sin(t)) & -2e^{-\frac{t}{4}}\sin(t) \\ e^{-\frac{t}{4}}\sin(t) & e^{-\frac{t}{4}}(-\sin(t) + \cos(t)) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-\frac{t}{4}}(\cos(t) + \sin(t))c_1 - 2e^{-\frac{t}{4}}\sin(t)c_2 \\ e^{-\frac{t}{4}}\sin(t)c_1 + e^{-\frac{t}{4}}(-\sin(t) + \cos(t))c_2 \end{bmatrix} \\
 &= \begin{bmatrix} ((c_1 - 2c_2)\sin(t) + c_1\cos(t))e^{-\frac{t}{4}} \\ ((-c_2 + c_1)\sin(t) + c_2\cos(t))e^{-\frac{t}{4}} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.11.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} \frac{3}{4} & -2 \\ 1 & -\frac{5}{4} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} \frac{3}{4} & -2 \\ 1 & -\frac{5}{4} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} \frac{3}{4} - \lambda & -2 \\ 1 & -\frac{5}{4} - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + \frac{1}{2}\lambda + \frac{17}{16} = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -\frac{1}{4} + i$$

$$\lambda_2 = -\frac{1}{4} - i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-\frac{1}{4} - i$	1	complex eigenvalue
$-\frac{1}{4} + i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -\frac{1}{4} - i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} \frac{3}{4} & -2 \\ 1 & -\frac{5}{4} \end{bmatrix} - \left(-\frac{1}{4} - i \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1+i & -2 \\ 1 & -1+i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1+i & -2 & 0 \\ 1 & -1+i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{1}{2} + \frac{i}{2} \right) R_1 \implies \left[\begin{array}{cc|c} 1+i & -2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1+i & -2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (1 - i)t\}$

Hence the solution is

$$\begin{bmatrix} (1 - i)t \\ t \end{bmatrix} = \begin{bmatrix} (1 - i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (1 - i)t \\ t \end{bmatrix} = t \begin{bmatrix} 1 - i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (1 - i)t \\ t \end{bmatrix} = \begin{bmatrix} 1 - i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -\frac{1}{4} + i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} \frac{3}{4} & -2 \\ 1 & -\frac{5}{4} \end{bmatrix} - \left(-\frac{1}{4} + i \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 - i & -2 \\ 1 & -1 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 - i & -2 & 0 \\ 1 & -1 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{1}{2} - \frac{i}{2} \right) R_1 \implies \left[\begin{array}{cc|c} 1 - i & -2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 - i & -2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (1 + i)t\}$

Hence the solution is

$$\begin{bmatrix} (1 + i)t \\ t \end{bmatrix} = \begin{bmatrix} (1 + i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (1 + i)t \\ t \end{bmatrix} = t \begin{bmatrix} 1 + i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (1 + i)t \\ t \end{bmatrix} = \begin{bmatrix} 1 + i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-\frac{1}{4} + i$	1	1	No	$\begin{bmatrix} 1 + i \\ 1 \end{bmatrix}$
$-\frac{1}{4} - i$	1	1	No	$\begin{bmatrix} 1 - i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (1+i)e^{(-\frac{1}{4}+i)t} \\ e^{(-\frac{1}{4}+i)t} \end{bmatrix} + c_2 \begin{bmatrix} (1-i)e^{(-\frac{1}{4}-i)t} \\ e^{(-\frac{1}{4}-i)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (1+i)c_1e^{(-\frac{1}{4}+i)t} + (1-i)c_2e^{(-\frac{1}{4}-i)t} \\ c_1e^{(-\frac{1}{4}+i)t} + c_2e^{(-\frac{1}{4}-i)t} \end{bmatrix}$$

The following is the phase plot of the system.

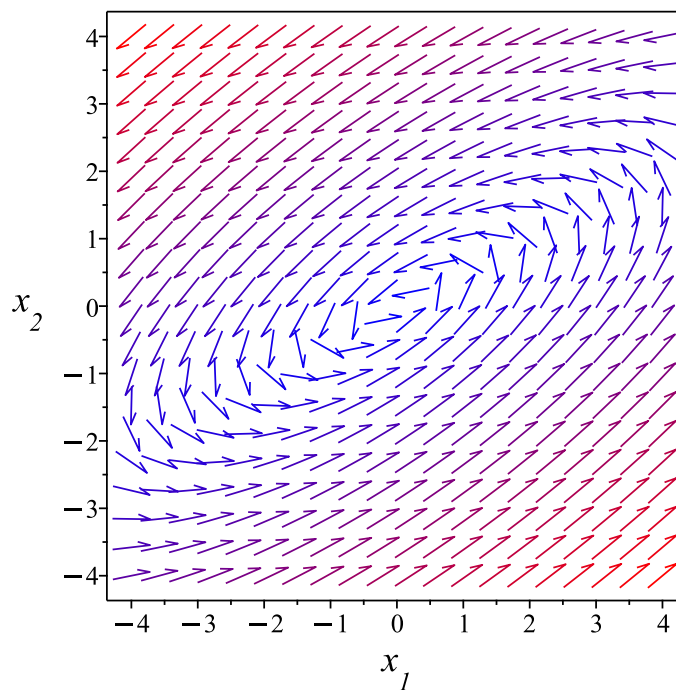


Figure 515: Phase plot

16.11.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = \frac{3x_1(t)}{4} - 2x_2(t), x_2'(t) = x_1(t) - \frac{5x_2(t)}{4} \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} \frac{3}{4} & -2 \\ 1 & -\frac{5}{4} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} \frac{3}{4} & -2 \\ 1 & -\frac{5}{4} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} \frac{3}{4} & -2 \\ 1 & -\frac{5}{4} \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-\frac{1}{4} - I, \begin{bmatrix} 1 - I \\ 1 \end{bmatrix} \right], \left[-\frac{1}{4} + I, \begin{bmatrix} 1 + I \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-\frac{1}{4} - I, \begin{bmatrix} 1 - I \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(-\frac{1}{4}-I)t} \cdot \begin{bmatrix} 1 - I \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{-\frac{t}{4}} \cdot (\cos(t) - I \sin(t)) \cdot \begin{bmatrix} 1 - I \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{-\frac{t}{4}} \cdot \begin{bmatrix} (1 - I) (\cos (t) - I \sin (t)) \\ \cos (t) - I \sin (t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\vec{x}_{\underline{1}}(t) = e^{-\frac{t}{4}} \cdot \begin{bmatrix} -\sin (t) + \cos (t) \\ \cos (t) \end{bmatrix}, \vec{x}_{\underline{2}}(t) = e^{-\frac{t}{4}} \cdot \begin{bmatrix} -\cos (t) - \sin (t) \\ -\sin (t) \end{bmatrix}$$

- General solution to the system of ODEs

$$\vec{x}_{\underline{}} = c_1 \vec{x}_{\underline{1}}(t) + c_2 \vec{x}_{\underline{2}}(t)$$

- Substitute solutions into the general solution

$$\vec{x}_{\underline{}} = c_1 e^{-\frac{t}{4}} \cdot \begin{bmatrix} -\sin (t) + \cos (t) \\ \cos (t) \end{bmatrix} + c_2 e^{-\frac{t}{4}} \cdot \begin{bmatrix} -\cos (t) - \sin (t) \\ -\sin (t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^{-\frac{t}{4}} (\cos (t) (c_1 - c_2) - \sin (t) (c_1 + c_2)) \\ e^{-\frac{t}{4}} (-c_2 \sin (t) + c_1 \cos (t)) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = e^{-\frac{t}{4}} (\cos (t) (c_1 - c_2) - \sin (t) (c_1 + c_2)), x_2(t) = e^{-\frac{t}{4}} (-c_2 \sin (t) + c_1 \cos (t)) \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 46

```
dsolve([diff(x__1(t),t)=3/4*x__1(t)-2*x__2(t),diff(x__2(t),t)=1*x__1(t)-5/4*x__2(t)],singsol
```

$$x_1(t) = e^{-\frac{t}{4}} (c_1 \sin (t) + c_2 \cos (t))$$

$$x_2(t) = \frac{e^{-\frac{t}{4}} (c_1 \sin (t) + c_2 \sin (t) - c_1 \cos (t) + c_2 \cos (t))}{2}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 56

```
DSolve[{x1'[t]==3/4*x1[t]-2*x2[t],x2'[t]==1*x1[t]-5/4*x2[t]},{x1[t],x2[t]},t,IncludeSingular
```

$$x1(t) \rightarrow e^{-t/4} (c_1 \cos(t) + (c_1 - 2c_2) \sin(t))$$

$$x2(t) \rightarrow e^{-t/4} (c_2 \cos(t) + (c_1 - c_2) \sin(t))$$

16.12 problem 12

16.12.1 Solution using Matrix exponential method	3630
16.12.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3631
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Internal problem ID [762]

Internal file name [OUTPUT/762_Sunday_June_05_2022_01_49_02_AM_3546516/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 12.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$\begin{aligned}x_1'(t) &= -\frac{4x_1(t)}{5} + 2x_2(t) \\x_2'(t) &= -x_1(t) + \frac{6x_2(t)}{5}\end{aligned}$$

16.12.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{4}{5} & 2 \\ -1 & \frac{6}{5} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned}e^{At} &= \begin{bmatrix} e^{\frac{t}{5}} \cos(t) - e^{\frac{t}{5}} \sin(t) & 2e^{\frac{t}{5}} \sin(t) \\ -e^{\frac{t}{5}} \sin(t) & e^{\frac{t}{5}} \cos(t) + e^{\frac{t}{5}} \sin(t) \end{bmatrix} \\ &= \begin{bmatrix} e^{\frac{t}{5}}(-\sin(t) + \cos(t)) & 2e^{\frac{t}{5}} \sin(t) \\ -e^{\frac{t}{5}} \sin(t) & e^{\frac{t}{5}}(\cos(t) + \sin(t)) \end{bmatrix}\end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{\frac{t}{5}}(-\sin(t) + \cos(t)) & 2e^{\frac{t}{5}}\sin(t) \\ -e^{\frac{t}{5}}\sin(t) & e^{\frac{t}{5}}(\cos(t) + \sin(t)) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{\frac{t}{5}}(-\sin(t) + \cos(t))c_1 + 2e^{\frac{t}{5}}\sin(t)c_2 \\ -e^{\frac{t}{5}}\sin(t)c_1 + e^{\frac{t}{5}}(\cos(t) + \sin(t))c_2 \end{bmatrix} \\
 &= \begin{bmatrix} ((-c_1 + 2c_2)\sin(t) + c_1\cos(t))e^{\frac{t}{5}} \\ -((-c_2 + c_1)\sin(t) - c_2\cos(t))e^{\frac{t}{5}} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.12.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{4}{5} & 2 \\ -1 & \frac{6}{5} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -\frac{4}{5} & 2 \\ -1 & \frac{6}{5} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -\frac{4}{5} - \lambda & 2 \\ -1 & \frac{6}{5} - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - \frac{2}{5}\lambda + \frac{26}{25} = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = \frac{1}{5} + i$$

$$\lambda_2 = \frac{1}{5} - i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$\frac{1}{5} - i$	1	complex eigenvalue
$\frac{1}{5} + i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = \frac{1}{5} - i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{4}{5} & 2 \\ -1 & \frac{6}{5} \end{bmatrix} - \left(\frac{1}{5} - i \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 + i & 2 \\ -1 & 1 + i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -1 + i & 2 & 0 \\ -1 & 1 + i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{1}{2} - \frac{i}{2} \right) R_1 \implies \left[\begin{array}{cc|c} -1 + i & 2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 + i & 2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (1 + i)t\}$

Hence the solution is

$$\begin{bmatrix} (1 + i)t \\ t \end{bmatrix} = \begin{bmatrix} (1 + i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (1 + i)t \\ t \end{bmatrix} = t \begin{bmatrix} 1 + i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (1 + i)t \\ t \end{bmatrix} = \begin{bmatrix} 1 + i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = \frac{1}{5} + i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{4}{5} & 2 \\ -1 & \frac{6}{5} \end{bmatrix} - \left(\frac{1}{5} + i \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 - i & 2 \\ -1 & 1 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -1 - i & 2 & 0 \\ -1 & 1 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{1}{2} + \frac{i}{2} \right) R_1 \implies \left[\begin{array}{cc|c} -1 - i & 2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 - i & 2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (1 - i)t\}$

Hence the solution is

$$\begin{bmatrix} (1 - i)t \\ t \end{bmatrix} = \begin{bmatrix} (1 - i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (1 - i)t \\ t \end{bmatrix} = t \begin{bmatrix} 1 - i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (1 - i)t \\ t \end{bmatrix} = \begin{bmatrix} 1 - i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$\frac{1}{5} + i$	1	1	No	$\begin{bmatrix} 1 - i \\ 1 \end{bmatrix}$
$\frac{1}{5} - i$	1	1	No	$\begin{bmatrix} 1 + i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (1-i)e^{(\frac{1}{5}+i)t} \\ e^{(\frac{1}{5}+i)t} \end{bmatrix} + c_2 \begin{bmatrix} (1+i)e^{(\frac{1}{5}-i)t} \\ e^{(\frac{1}{5}-i)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (1-i)c_1e^{(\frac{1}{5}+i)t} + (1+i)c_2e^{(\frac{1}{5}-i)t} \\ c_1e^{(\frac{1}{5}+i)t} + c_2e^{(\frac{1}{5}-i)t} \end{bmatrix}$$

The following is the phase plot of the system.

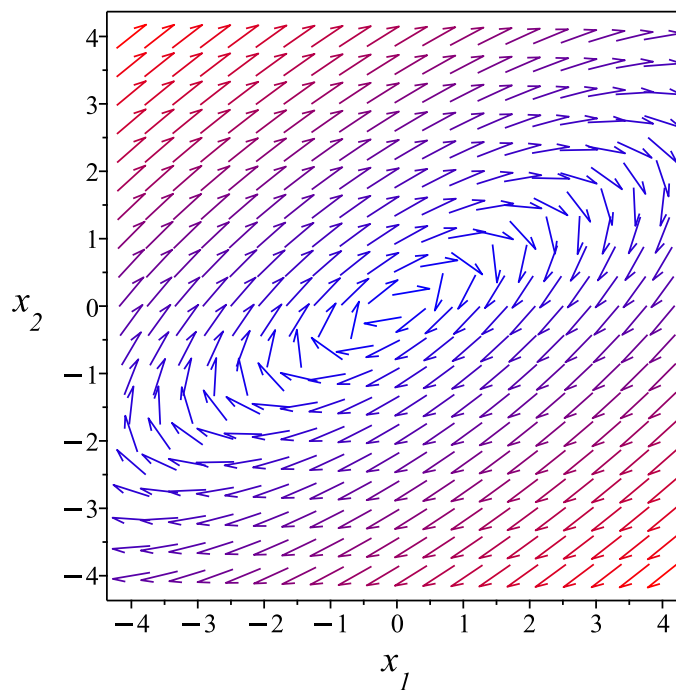


Figure 516: Phase plot

16.12.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = -\frac{4x_1(t)}{5} + 2x_2(t), x_2'(t) = -x_1(t) + \frac{6x_2(t)}{5} \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -\frac{4}{5} & 2 \\ -1 & \frac{6}{5} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -\frac{4}{5} & 2 \\ -1 & \frac{6}{5} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -\frac{4}{5} & 2 \\ -1 & \frac{6}{5} \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[\frac{1}{5} - \text{I}, \begin{bmatrix} 1 + \text{I} \\ 1 \end{bmatrix} \right], \left[\frac{1}{5} + \text{I}, \begin{bmatrix} 1 - \text{I} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[\frac{1}{5} - \text{I}, \begin{bmatrix} 1 + \text{I} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(\frac{1}{5} - \text{I})t} \cdot \begin{bmatrix} 1 + \text{I} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{\frac{t}{5}} \cdot (\cos(t) - \text{I} \sin(t)) \cdot \begin{bmatrix} 1 + \text{I} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{\frac{t}{5}} \cdot \begin{bmatrix} (1 + I) (\cos (t) - I \sin (t)) \\ \cos (t) - I \sin (t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\underline{x}_{\rightarrow 1}(t) = e^{\frac{t}{5}} \cdot \begin{bmatrix} \cos (t) + \sin (t) \\ \cos (t) \end{bmatrix}, \underline{x}_{\rightarrow 2}(t) = e^{\frac{t}{5}} \cdot \begin{bmatrix} -\sin (t) + \cos (t) \\ -\sin (t) \end{bmatrix} \right]$$

- General solution to the system of ODEs

$$\underline{x}_{\rightarrow} = c_1 \underline{x}_{\rightarrow 1}(t) + c_2 \underline{x}_{\rightarrow 2}(t)$$

- Substitute solutions into the general solution

$$\underline{x}_{\rightarrow} = c_1 e^{\frac{t}{5}} \cdot \begin{bmatrix} \cos (t) + \sin (t) \\ \cos (t) \end{bmatrix} + c_2 e^{\frac{t}{5}} \cdot \begin{bmatrix} -\sin (t) + \cos (t) \\ -\sin (t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (\cos (t) (c_1 + c_2) + (c_1 - c_2) \sin (t)) e^{\frac{t}{5}} \\ e^{\frac{t}{5}} (-c_2 \sin (t) + c_1 \cos (t)) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = (\cos (t) (c_1 + c_2) + (c_1 - c_2) \sin (t)) e^{\frac{t}{5}}, x_2(t) = e^{\frac{t}{5}} (-c_2 \sin (t) + c_1 \cos (t)) \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 46

```
dsolve([diff(x__1(t),t)=-4/5*x__1(t)+2*x__2(t),diff(x__2(t),t)=-1*x__1(t)+6/5*x__2(t)],sings
```

$$x_1(t) = e^{\frac{t}{5}} (c_1 \sin (t) + c_2 \cos (t))$$

$$x_2(t) = \frac{e^{\frac{t}{5}} (c_1 \sin (t) - c_2 \sin (t) + c_1 \cos (t) + c_2 \cos (t))}{2}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 56

```
DSolve[{x1'[t]==-4/5*x1[t]+2*x2[t],x2'[t]==-1*x1[t]+6/5*x2[t]},{x1[t],x2[t]},t,IncludeSingul
```

$$x_1(t) \rightarrow e^{t/5} (c_1 \cos(t) - (c_1 - 2c_2) \sin(t))$$

$$x_2(t) \rightarrow e^{t/5} (c_2 (\sin(t) + \cos(t)) - c_1 \sin(t))$$

16.13 problem 23

16.13.1 Solution using Matrix exponential method	3638
16.13.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3639
16.13.3 Maple step by step solution	3646

Internal problem ID [763]

Internal file name [OUTPUT/763_Sunday_June_05_2022_01_49_04_AM_12891500/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 23.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -\frac{x_1(t)}{4} + x_2(t) \\x_2'(t) &= -x_1(t) - \frac{x_2(t)}{4} \\x_3'(t) &= -\frac{x_3(t)}{4}\end{aligned}$$

16.13.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & -\frac{1}{4} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-\frac{t}{4}} \cos(t) & e^{-\frac{t}{4}} \sin(t) & 0 \\ -e^{-\frac{t}{4}} \sin(t) & e^{-\frac{t}{4}} \cos(t) & 0 \\ 0 & 0 & e^{-\frac{t}{4}} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} e^{-\frac{t}{4}} \cos(t) & e^{-\frac{t}{4}} \sin(t) & 0 \\ -e^{-\frac{t}{4}} \sin(t) & e^{-\frac{t}{4}} \cos(t) & 0 \\ 0 & 0 & e^{-\frac{t}{4}} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ c_3 \end{bmatrix} \\ &= \begin{bmatrix} e^{-\frac{t}{4}} \cos(t) c_1 + e^{-\frac{t}{4}} \sin(t) c_2 \\ -e^{-\frac{t}{4}} \sin(t) c_1 + e^{-\frac{t}{4}} \cos(t) c_2 \\ e^{-\frac{t}{4}} c_3 \end{bmatrix} \\ &= \begin{bmatrix} e^{-\frac{t}{4}} (\cos(t) c_1 + \sin(t) c_2) \\ e^{-\frac{t}{4}} (-\sin(t) c_1 + \cos(t) c_2) \\ e^{-\frac{t}{4}} c_3 \end{bmatrix} \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.13.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & -\frac{1}{4} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & -\frac{1}{4} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -\frac{1}{4} - \lambda & 1 & 0 \\ -1 & -\frac{1}{4} - \lambda & 0 \\ 0 & 0 & -\frac{1}{4} - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^3 + \frac{3}{4}\lambda^2 + \frac{19}{16}\lambda + \frac{17}{64} = 0$$

The roots of the above are the eigenvalues.

$$\begin{aligned} \lambda_1 &= -\frac{1}{4} + i \\ \lambda_2 &= -\frac{1}{4} - i \\ \lambda_3 &= -\frac{1}{4} \end{aligned}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-\frac{1}{4} - i$	1	complex eigenvalue
$-\frac{1}{4}$	1	real eigenvalue
$-\frac{1}{4} + i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -\frac{1}{4}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & -\frac{1}{4} \end{bmatrix} - \left(-\frac{1}{4} \right) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Since the current pivot $A(1,1)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 1 and row 2 gives

$$\left[\begin{array}{ccc|c} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 0, v_2 = 0\}$

Hence the solution is

$$\begin{bmatrix} 0 \\ 0 \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 0 \\ 0 \\ t \end{bmatrix} = t \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -\frac{1}{4} - i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & -\frac{1}{4} \end{bmatrix} - \left(-\frac{1}{4} - i \right) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} i & 1 & 0 \\ -1 & i & 0 \\ 0 & 0 & i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} i & 1 & 0 & 0 \\ -1 & i & 0 & 0 \\ 0 & 0 & i & 0 \end{array} \right]$$

$$R_2 = -iR_1 + R_2 \implies \left[\begin{array}{ccc|c} i & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & i & 0 \end{array} \right]$$

Since the current pivot $A(2,3)$ is zero, then the current pivot row is replaced with a

row with a non-zero pivot. Swapping row 2 and row 3 gives

$$\left[\begin{array}{ccc|c} i & 1 & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} i & 1 & 0 \\ 0 & 0 & i \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1, v_3\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = it, v_3 = 0\}$

Hence the solution is

$$\begin{bmatrix} It \\ t \\ 0 \end{bmatrix} = \begin{bmatrix} it \\ t \\ 0 \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} It \\ t \\ 0 \end{bmatrix} = t \begin{bmatrix} i \\ 1 \\ 0 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} It \\ t \\ 0 \end{bmatrix} = \begin{bmatrix} i \\ 1 \\ 0 \end{bmatrix}$$

Considering the eigenvalue $\lambda_3 = -\frac{1}{4} + i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & -\frac{1}{4} \end{bmatrix} - \left(-\frac{1}{4} + i \right) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -i & 1 & 0 \\ -1 & -i & 0 \\ 0 & 0 & -i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -i & 1 & 0 & 0 \\ -1 & -i & 0 & 0 \\ 0 & 0 & -i & 0 \end{array} \right]$$

$$R_2 = iR_1 + R_2 \implies \left[\begin{array}{ccc|c} -i & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & -i & 0 \end{array} \right]$$

Since the current pivot $A(2,3)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 2 and row 3 gives

$$\left[\begin{array}{ccc|c} -i & 1 & 0 & 0 \\ 0 & 0 & -i & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -i & 1 & 0 \\ 0 & 0 & -i \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1, v_3\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -it, v_3 = 0\}$

Hence the solution is

$$\begin{bmatrix} -It \\ t \\ 0 \end{bmatrix} = \begin{bmatrix} -it \\ t \\ 0 \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -It \\ t \\ 0 \end{bmatrix} = t \begin{bmatrix} -i \\ 1 \\ 0 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -It \\ t \\ 0 \end{bmatrix} = \begin{bmatrix} -i \\ 1 \\ 0 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-\frac{1}{4} + i$	1	1	No	$\begin{bmatrix} -i \\ 1 \\ 0 \end{bmatrix}$
$-\frac{1}{4} - i$	1	1	No	$\begin{bmatrix} i \\ 1 \\ 0 \end{bmatrix}$
$-\frac{1}{4}$	1	1	No	$\begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue $-\frac{1}{4}$ is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{-\frac{t}{4}} \\ &= \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} e^{-\frac{t}{4}}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = c_1 \begin{bmatrix} -ie^{(-\frac{1}{4}+i)t} \\ e^{(-\frac{1}{4}+i)t} \\ 0 \end{bmatrix} + c_2 \begin{bmatrix} ie^{(-\frac{1}{4}-i)t} \\ e^{(-\frac{1}{4}-i)t} \\ 0 \end{bmatrix} + c_3 \begin{bmatrix} 0 \\ 0 \\ e^{-\frac{t}{4}} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} -i(c_1 e^{(-\frac{1}{4}+i)t} - c_2 e^{(-\frac{1}{4}-i)t}) \\ c_1 e^{(-\frac{1}{4}+i)t} + c_2 e^{(-\frac{1}{4}-i)t} \\ c_3 e^{-\frac{t}{4}} \end{bmatrix}$$

16.13.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = -\frac{x_1(t)}{4} + x_2(t), x_2'(t) = -x_1(t) - \frac{x_2(t)}{4}, x_3'(t) = -\frac{x_3(t)}{4} \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & -\frac{1}{4} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & -\frac{1}{4} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & -\frac{1}{4} \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-\frac{1}{4}, \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \right], \left[-\frac{1}{4} - \mathbf{I}, \begin{bmatrix} \mathbf{I} \\ 1 \\ 0 \end{bmatrix} \right], \left[-\frac{1}{4} + \mathbf{I}, \begin{bmatrix} -\mathbf{I} \\ 1 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-\frac{1}{4}, \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{-\frac{t}{4}} \cdot \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-\frac{1}{4} - I, \begin{bmatrix} I \\ 1 \\ 0 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(-\frac{1}{4}-I)t} \cdot \begin{bmatrix} I \\ 1 \\ 0 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{-\frac{t}{4}} \cdot (\cos(t) - I \sin(t)) \cdot \begin{bmatrix} I \\ 1 \\ 0 \end{bmatrix}$$

- Simplify expression

$$e^{-\frac{t}{4}} \cdot \begin{bmatrix} I(\cos(t) - I \sin(t)) \\ \cos(t) - I \sin(t) \\ 0 \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\vec{x}_2(t) = e^{-\frac{t}{4}} \cdot \begin{bmatrix} \sin(t) \\ \cos(t) \\ 0 \end{bmatrix}, \vec{x}_3(t) = e^{-\frac{t}{4}} \cdot \begin{bmatrix} \cos(t) \\ -\sin(t) \\ 0 \end{bmatrix} \right]$$

- General solution to the system of ODEs

$$\vec{x} = c_1 \vec{x}_1 + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

- Substitute solutions into the general solution

$$\vec{x} = c_1 e^{-\frac{t}{4}} \cdot \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} + c_2 e^{-\frac{t}{4}} \cdot \begin{bmatrix} \sin(t) \\ \cos(t) \\ 0 \end{bmatrix} + c_3 e^{-\frac{t}{4}} \cdot \begin{bmatrix} \cos(t) \\ -\sin(t) \\ 0 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} e^{-\frac{t}{4}}(c_2 \sin(t) + c_3 \cos(t)) \\ e^{-\frac{t}{4}}(c_2 \cos(t) - c_3 \sin(t)) \\ c_1 e^{-\frac{t}{4}} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = e^{-\frac{t}{4}}(c_2 \sin(t) + c_3 \cos(t)), x_2(t) = e^{-\frac{t}{4}}(c_2 \cos(t) - c_3 \sin(t)), x_3(t) = c_1 e^{-\frac{t}{4}} \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 47

```
dsolve([diff(x__1(t),t)=-1/4*x__1(t)+1*x__2(t)+0*x__3(t),diff(x__2(t),t)=-1*x__1(t)-1/4*x__2
```

$$\begin{aligned} x_1(t) &= e^{-\frac{t}{4}}(c_1 \sin(t) + c_2 \cos(t)) \\ x_2(t) &= -e^{-\frac{t}{4}}(c_2 \sin(t) - c_1 \cos(t)) \\ x_3(t) &= c_3 e^{-\frac{t}{4}} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.023 (sec). Leaf size: 110

```
DSolve[{x1'[t]==-1/4*x1[t]+1*x2[t]+0*x3[t],x2'[t]==-1*x1[t]-1/4*x2[t]+0*x3[t],x3'[t]==0*x1[t
```

$$\begin{aligned} x1(t) &\rightarrow e^{-t/4}(c_1 \cos(t) + c_2 \sin(t)) \\ x2(t) &\rightarrow e^{-t/4}(c_2 \cos(t) - c_1 \sin(t)) \\ x3(t) &\rightarrow c_3 e^{-t/4} \\ x1(t) &\rightarrow e^{-t/4}(c_1 \cos(t) + c_2 \sin(t)) \\ x2(t) &\rightarrow e^{-t/4}(c_2 \cos(t) - c_1 \sin(t)) \\ x3(t) &\rightarrow 0 \end{aligned}$$

16.14 problem 24

- 16.14.1 Solution using Matrix exponential method 3650
- 16.14.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3651
- 16.14.3 Maple step by step solution 3658

Internal problem ID [764]

Internal file name [OUTPUT/764_Sunday_June_05_2022_01_49_05_AM_58372908/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 24.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -\frac{x_1(t)}{4} + x_2(t) \\x_2'(t) &= -x_1(t) - \frac{x_2(t)}{4} \\x_3'(t) &= \frac{x_3(t)}{10}\end{aligned}$$

16.14.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & \frac{1}{10} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-\frac{t}{4}} \cos(t) & e^{-\frac{t}{4}} \sin(t) & 0 \\ -e^{-\frac{t}{4}} \sin(t) & e^{-\frac{t}{4}} \cos(t) & 0 \\ 0 & 0 & e^{\frac{t}{10}} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} e^{-\frac{t}{4}} \cos(t) & e^{-\frac{t}{4}} \sin(t) & 0 \\ -e^{-\frac{t}{4}} \sin(t) & e^{-\frac{t}{4}} \cos(t) & 0 \\ 0 & 0 & e^{\frac{t}{10}} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ c_3 \end{bmatrix} \\ &= \begin{bmatrix} e^{-\frac{t}{4}} \cos(t) c_1 + e^{-\frac{t}{4}} \sin(t) c_2 \\ -e^{-\frac{t}{4}} \sin(t) c_1 + e^{-\frac{t}{4}} \cos(t) c_2 \\ e^{\frac{t}{10}} c_3 \end{bmatrix} \\ &= \begin{bmatrix} e^{-\frac{t}{4}} (\cos(t) c_1 + \sin(t) c_2) \\ e^{-\frac{t}{4}} (-\sin(t) c_1 + \cos(t) c_2) \\ e^{\frac{t}{10}} c_3 \end{bmatrix} \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.14.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & \frac{1}{10} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & \frac{1}{10} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -\frac{1}{4} - \lambda & 1 & 0 \\ -1 & -\frac{1}{4} - \lambda & 0 \\ 0 & 0 & \frac{1}{10} - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^3 + \frac{2}{5}\lambda^2 + \frac{81}{80}\lambda - \frac{17}{160} = 0$$

The roots of the above are the eigenvalues.

$$\begin{aligned}\lambda_1 &= -\frac{1}{4} + i \\ \lambda_2 &= -\frac{1}{4} - i \\ \lambda_3 &= \frac{1}{10}\end{aligned}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-\frac{1}{4} - i$	1	complex eigenvalue
$-\frac{1}{4} + i$	1	complex eigenvalue
$\frac{1}{10}$	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = \frac{1}{10}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & \frac{1}{10} \end{bmatrix} - \left(\frac{1}{10} \right) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -\frac{7}{20} & 1 & 0 \\ -1 & -\frac{7}{20} & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -\frac{7}{20} & 1 & 0 & 0 \\ -1 & -\frac{7}{20} & 0 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{20R_1}{7} \implies \left[\begin{array}{ccc|c} -\frac{7}{20} & 1 & 0 & 0 \\ 0 & -\frac{449}{140} & 0 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -\frac{7}{20} & 1 & 0 \\ 0 & -\frac{449}{140} & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 0, v_2 = 0\}$

Hence the solution is

$$\begin{bmatrix} 0 \\ 0 \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 0 \\ 0 \\ t \end{bmatrix} = t \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -\frac{1}{4} - i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & \frac{1}{10} \end{bmatrix} - \left(-\frac{1}{4} - i \right) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} i & 1 & 0 \\ -1 & i & 0 \\ 0 & 0 & \frac{7}{20} + i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} i & 1 & 0 & 0 \\ -1 & i & 0 & 0 \\ 0 & 0 & \frac{7}{20} + i & 0 \end{array} \right]$$

$$R_2 = -iR_1 + R_2 \implies \left[\begin{array}{ccc|c} i & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & \frac{7}{20} + i & 0 \end{array} \right]$$

Since the current pivot $A(2,3)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 2 and row 3 gives

$$\left[\begin{array}{ccc|c} i & 1 & 0 & 0 \\ 0 & 0 & \frac{7}{20} + i & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{ccc} i & 1 & 0 \\ 0 & 0 & \frac{7}{20} + i \\ 0 & 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1, v_3\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = it, v_3 = 0\}$

Hence the solution is

$$\begin{bmatrix} It \\ t \\ 0 \end{bmatrix} = \begin{bmatrix} it \\ t \\ 0 \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} It \\ t \\ 0 \end{bmatrix} = t \begin{bmatrix} i \\ 1 \\ 0 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} It \\ t \\ 0 \end{bmatrix} = \begin{bmatrix} i \\ 1 \\ 0 \end{bmatrix}$$

Considering the eigenvalue $\lambda_3 = -\frac{1}{4} + i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & \frac{1}{10} \end{bmatrix} - \left(-\frac{1}{4} + i \right) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -i & 1 & 0 \\ -1 & -i & 0 \\ 0 & 0 & \frac{7}{20} - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -i & 1 & 0 & 0 \\ -1 & -i & 0 & 0 \\ 0 & 0 & \frac{7}{20} - i & 0 \end{array} \right]$$

$$R_2 = iR_1 + R_2 \implies \left[\begin{array}{ccc|c} -i & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & \frac{7}{20} - i & 0 \end{array} \right]$$

Since the current pivot $A(2,3)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 2 and row 3 gives

$$\left[\begin{array}{ccc|c} -i & 1 & 0 & 0 \\ 0 & 0 & \frac{7}{20} - i & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -i & 1 & 0 \\ 0 & 0 & \frac{7}{20} - i \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1, v_3\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -it, v_3 = 0\}$

Hence the solution is

$$\begin{bmatrix} -It \\ t \\ 0 \end{bmatrix} = \begin{bmatrix} -it \\ t \\ 0 \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -It \\ t \\ 0 \end{bmatrix} = t \begin{bmatrix} -i \\ 1 \\ 0 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -It \\ t \\ 0 \end{bmatrix} = \begin{bmatrix} -i \\ 1 \\ 0 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-\frac{1}{4} + i$	1	1	No	$\begin{bmatrix} -i \\ 1 \\ 0 \end{bmatrix}$
$-\frac{1}{4} - i$	1	1	No	$\begin{bmatrix} i \\ 1 \\ 0 \end{bmatrix}$
$\frac{1}{10}$	1	1	No	$\begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue $\frac{1}{10}$ is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\frac{t}{10}} \\ &= \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} e^{\frac{t}{10}}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = c_1 \begin{bmatrix} -ie^{(-\frac{1}{4}+i)t} \\ e^{(-\frac{1}{4}+i)t} \\ 0 \end{bmatrix} + c_2 \begin{bmatrix} ie^{(-\frac{1}{4}-i)t} \\ e^{(-\frac{1}{4}-i)t} \\ 0 \end{bmatrix} + c_3 \begin{bmatrix} 0 \\ 0 \\ e^{\frac{t}{10}} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} i(c_2 e^{(-\frac{1}{4}-i)t} - c_1 e^{(-\frac{1}{4}+i)t}) \\ c_1 e^{(-\frac{1}{4}+i)t} + c_2 e^{(-\frac{1}{4}-i)t} \\ c_3 e^{\frac{t}{10}} \end{bmatrix}$$

16.14.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = -\frac{x_1(t)}{4} + x_2(t), x_2'(t) = -x_1(t) - \frac{x_2(t)}{4}, x_3'(t) = \frac{x_3(t)}{10} \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & \frac{1}{10} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & \frac{1}{10} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -\frac{1}{4} & 1 & 0 \\ -1 & -\frac{1}{4} & 0 \\ 0 & 0 & \frac{1}{10} \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[\frac{1}{10}, \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \right], \left[-\frac{1}{4} - \text{I}, \begin{bmatrix} 1 \\ 1 \\ 0 \end{bmatrix} \right], \left[-\frac{1}{4} + \text{I}, \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[\frac{1}{10}, \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{\frac{t}{10}} \cdot \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-\frac{1}{4} - I, \begin{bmatrix} I \\ 1 \\ 0 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(-\frac{1}{4}-I)t} \cdot \begin{bmatrix} I \\ 1 \\ 0 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{-\frac{t}{4}} \cdot (\cos(t) - I \sin(t)) \cdot \begin{bmatrix} I \\ 1 \\ 0 \end{bmatrix}$$

- Simplify expression

$$e^{-\frac{t}{4}} \cdot \begin{bmatrix} I(\cos(t) - I \sin(t)) \\ \cos(t) - I \sin(t) \\ 0 \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\vec{x}_2(t) = e^{-\frac{t}{4}} \cdot \begin{bmatrix} \sin(t) \\ \cos(t) \\ 0 \end{bmatrix}, \vec{x}_3(t) = e^{-\frac{t}{4}} \cdot \begin{bmatrix} \cos(t) \\ -\sin(t) \\ 0 \end{bmatrix} \right]$$

- General solution to the system of ODEs

$$\vec{x} = c_1 \vec{x}_1 + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

- Substitute solutions into the general solution

$$\vec{x} = c_1 e^{\frac{t}{10}} \cdot \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} + c_2 e^{-\frac{t}{4}} \cdot \begin{bmatrix} \sin(t) \\ \cos(t) \\ 0 \end{bmatrix} + c_3 e^{-\frac{t}{4}} \cdot \begin{bmatrix} \cos(t) \\ -\sin(t) \\ 0 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} e^{-\frac{t}{4}}(c_2 \sin(t) + c_3 \cos(t)) \\ e^{-\frac{t}{4}}(c_2 \cos(t) - c_3 \sin(t)) \\ c_1 e^{\frac{t}{10}} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = e^{-\frac{t}{4}}(c_2 \sin(t) + c_3 \cos(t)), x_2(t) = e^{-\frac{t}{4}}(c_2 \cos(t) - c_3 \sin(t)), x_3(t) = c_1 e^{\frac{t}{10}} \right\}$$

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 47

```
dsolve([diff(x__1(t),t)=-1/4*x__1(t)+1*x__2(t)+0*x__3(t),diff(x__2(t),t)=-1*x__1(t)-1/4*x__2
```

$$\begin{aligned} x_1(t) &= e^{-\frac{t}{4}}(c_1 \sin(t) + c_2 \cos(t)) \\ x_2(t) &= -e^{-\frac{t}{4}}(c_2 \sin(t) - c_1 \cos(t)) \\ x_3(t) &= c_3 e^{\frac{t}{10}} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.021 (sec). Leaf size: 110

```
DSolve[{x1'[t]==-1/4*x1[t]+1*x2[t]+0*x3[t],x2'[t]==-1*x1[t]-1/4*x2[t]+0*x3[t],x3'[t]==0*x1[t
```

$$\begin{aligned} x1(t) &\rightarrow e^{-t/4}(c_1 \cos(t) + c_2 \sin(t)) \\ x2(t) &\rightarrow e^{-t/4}(c_2 \cos(t) - c_1 \sin(t)) \\ x3(t) &\rightarrow c_3 e^{t/10} \\ x1(t) &\rightarrow e^{-t/4}(c_1 \cos(t) + c_2 \sin(t)) \\ x2(t) &\rightarrow e^{-t/4}(c_2 \cos(t) - c_1 \sin(t)) \\ x3(t) &\rightarrow 0 \end{aligned}$$

16.15 problem 25

- 16.15.1 Solution using Matrix exponential method 3662
- 16.15.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3663
- 16.15.3 Maple step by step solution 3668

Internal problem ID [765]

Internal file name [OUTPUT/765_Sunday_June_05_2022_01_49_07_AM_11650177/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.6, Complex Eigenvalues. page 417

Problem number: 25.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -\frac{x_1(t)}{2} - \frac{x_2(t)}{8} \\x_2'(t) &= 2x_1(t) - \frac{x_2(t)}{2}\end{aligned}$$

16.15.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) & -\frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right)}{4} \\ 4e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) & e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) & -\frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right)}{4} \\ 4e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) & e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) c_1 - \frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) c_2}{4} \\ 4e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) c_1 + e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) c_1 - \frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) c_2}{4} \\ e^{-\frac{t}{2}} \left(4 \sin\left(\frac{t}{2}\right) c_1 + \cos\left(\frac{t}{2}\right) c_2\right) \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

16.15.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -\frac{1}{2} - \lambda & -\frac{1}{8} \\ 2 & -\frac{1}{2} - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + \lambda + \frac{1}{2} = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -\frac{1}{2} + \frac{i}{2}$$

$$\lambda_2 = -\frac{1}{2} - \frac{i}{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-\frac{1}{2} + \frac{i}{2}$	1	complex eigenvalue
$-\frac{1}{2} - \frac{i}{2}$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -\frac{1}{2} - \frac{i}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} - \left(-\frac{1}{2} - \frac{i}{2} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} \frac{i}{2} & -\frac{1}{8} \\ 2 & \frac{i}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{i}{2} & -\frac{1}{8} & 0 \\ 2 & \frac{i}{2} & 0 \end{array} \right]$$

$$R_2 = 4iR_1 + R_2 \implies \left[\begin{array}{cc|c} \frac{i}{2} & -\frac{1}{8} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} \frac{i}{2} & -\frac{1}{8} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -\frac{it}{4}\}$

Hence the solution is

$$\begin{bmatrix} -\frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} -\frac{it}{4} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{1}{4}t \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{i}{4} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} -\frac{i}{4} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} -i \\ 4 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -\frac{1}{2} + \frac{i}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} - \left(-\frac{1}{2} + \frac{i}{2} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -\frac{i}{2} & -\frac{1}{8} \\ 2 & -\frac{i}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -\frac{i}{2} & -\frac{1}{8} & 0 \\ 2 & -\frac{i}{2} & 0 \end{array} \right]$$

$$R_2 = -4iR_1 + R_2 \implies \left[\begin{array}{cc|c} -\frac{i}{2} & -\frac{1}{8} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -\frac{i}{2} & -\frac{1}{8} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{it}{4}\}$

Hence the solution is

$$\begin{bmatrix} \frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} \frac{it}{4} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{1}{4}t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{i}{4} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} \frac{i}{4} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} i \\ 4 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-\frac{1}{2} + \frac{i}{2}$	1	1	No	$\begin{bmatrix} \frac{i}{4} \\ 1 \end{bmatrix}$
$-\frac{1}{2} - \frac{i}{2}$	1	1	No	$\begin{bmatrix} -\frac{i}{4} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{ie^{(-\frac{1}{2} + \frac{i}{2})t}}{4} \\ e^{(-\frac{1}{2} + \frac{i}{2})t} \end{bmatrix} + c_2 \begin{bmatrix} -\frac{ie^{(-\frac{1}{2} - \frac{i}{2})t}}{4} \\ e^{(-\frac{1}{2} - \frac{i}{2})t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} i(c_2 e^{(-\frac{1}{2} - \frac{i}{2})t} - c_1 e^{(-\frac{1}{2} + \frac{i}{2})t}) \\ c_1 e^{(-\frac{1}{2} + \frac{i}{2})t} + c_2 e^{(-\frac{1}{2} - \frac{i}{2})t} \end{bmatrix}$$

The following is the phase plot of the system.

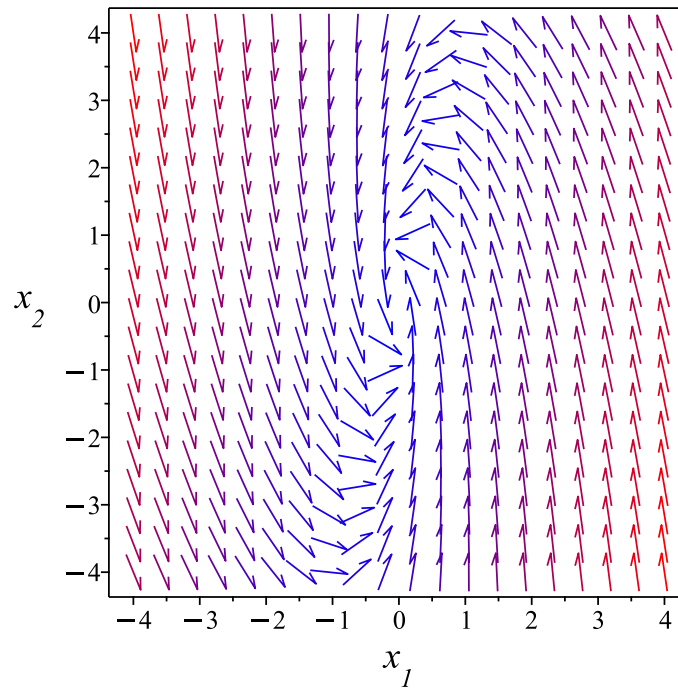


Figure 517: Phase plot

16.15.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = -\frac{x_1(t)}{2} - \frac{x_2(t)}{8}, x_2'(t) = 2x_1(t) - \frac{x_2(t)}{2} \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow}{}'(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-\frac{1}{2} - \frac{1}{2}, \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix} \right], \left[-\frac{1}{2} + \frac{1}{2}, \begin{bmatrix} \frac{1}{4} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-\frac{1}{2} - \frac{1}{2}, \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(-\frac{1}{2}-\frac{1}{2})t} \cdot \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{-\frac{t}{2}} \cdot (\cos(\frac{t}{2}) - I \sin(\frac{t}{2})) \cdot \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{-\frac{t}{2}} \cdot \begin{bmatrix} -\frac{1}{4}(\cos(\frac{t}{2}) - I \sin(\frac{t}{2})) \\ \cos(\frac{t}{2}) - I \sin(\frac{t}{2}) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\underline{x}^{\rightarrow}_1(t) = e^{-\frac{t}{2}} \cdot \begin{bmatrix} -\frac{\sin(\frac{t}{2})}{4} \\ \cos(\frac{t}{2}) \end{bmatrix}, \underline{x}^{\rightarrow}_2(t) = e^{-\frac{t}{2}} \cdot \begin{bmatrix} -\frac{\cos(\frac{t}{2})}{4} \\ -\sin(\frac{t}{2}) \end{bmatrix} \right]$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1(t) + c_2 \underline{x}^{\rightarrow}_2(t)$$

- Substitute solutions into the general solution

$$\underline{x}^{\rightarrow} = c_1 e^{-\frac{t}{2}} \cdot \begin{bmatrix} -\frac{\sin(\frac{t}{2})}{4} \\ \cos(\frac{t}{2}) \end{bmatrix} + c_2 e^{-\frac{t}{2}} \cdot \begin{bmatrix} -\frac{\cos(\frac{t}{2})}{4} \\ -\sin(\frac{t}{2}) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} -\frac{e^{-\frac{t}{2}}(c_1 \sin(\frac{t}{2}) + c_2 \cos(\frac{t}{2}))}{4} \\ e^{-\frac{t}{2}}(c_1 \cos(\frac{t}{2}) - c_2 \sin(\frac{t}{2})) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = -\frac{e^{-\frac{t}{2}}(c_1 \sin(\frac{t}{2}) + c_2 \cos(\frac{t}{2}))}{4}, x_2(t) = e^{-\frac{t}{2}}(c_1 \cos(\frac{t}{2}) - c_2 \sin(\frac{t}{2})) \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 46

```
dsolve([diff(x__1(t),t)=-1/2*x__1(t)-1/8*x__2(t),diff(x__2(t),t)=2*x__1(t)-1/2*x__2(t)],sing
```

$$x_1(t) = e^{-\frac{t}{2}} \left(c_2 \cos\left(\frac{t}{2}\right) + c_1 \sin\left(\frac{t}{2}\right) \right)$$

$$x_2(t) = -4 e^{-\frac{t}{2}} \left(\cos\left(\frac{t}{2}\right) c_1 - \sin\left(\frac{t}{2}\right) c_2 \right)$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 68

```
DSolve[{x1'[t]==-1/2*x1[t]-1/8*x2[t],x2'[t]==2*x1[t]-1/2*x2[t]},{x1[t],x2[t]},t,IncludeSingu
```

$$x1(t) \rightarrow \frac{1}{4} e^{-t/2} \left(4c_1 \cos\left(\frac{t}{2}\right) - c_2 \sin\left(\frac{t}{2}\right) \right)$$

$$x2(t) \rightarrow e^{-t/2} \left(c_2 \cos\left(\frac{t}{2}\right) + 4c_1 \sin\left(\frac{t}{2}\right) \right)$$

17 Chapter 7.8, Repeated Eigenvalues. page 436

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17.1 problem 1

- 17.1.1 Solution using Matrix exponential method 3672
- 17.1.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3673
- 17.1.3 Maple step by step solution 3678

Internal problem ID [766]

Internal file name [OUTPUT/766_Sunday_June_05_2022_01_49_09_AM_14401963/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 3x_1(t) - 4x_2(t) \\x_2'(t) &= x_1(t) - x_2(t)\end{aligned}$$

17.1.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^t(1 + 2t) & -4t e^t \\ t e^t & e^t(1 - 2t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}\vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} e^t(1+2t) & -4te^t \\ te^t & e^t(1-2t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} e^t(1+2t)c_1 - 4te^t c_2 \\ te^t c_1 + e^t(1-2t)c_2 \end{bmatrix} \\ &= \begin{bmatrix} e^t(2tc_1 - 4c_2t + c_1) \\ e^t(tc_1 - 2c_2t + c_2) \end{bmatrix}\end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.1.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 3 - \lambda & -4 \\ 1 & -1 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 2\lambda + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
1	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = 1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} - (1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 & -4 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 & -4 & 0 \\ 1 & -2 & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{R_1}{2} \implies \left[\begin{array}{cc|c} 2 & -4 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 & -4 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 2t\}$

Hence the solution is

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = \begin{bmatrix} 2t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = t \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
1	2	1	Yes	$\begin{bmatrix} 2 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue 1 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

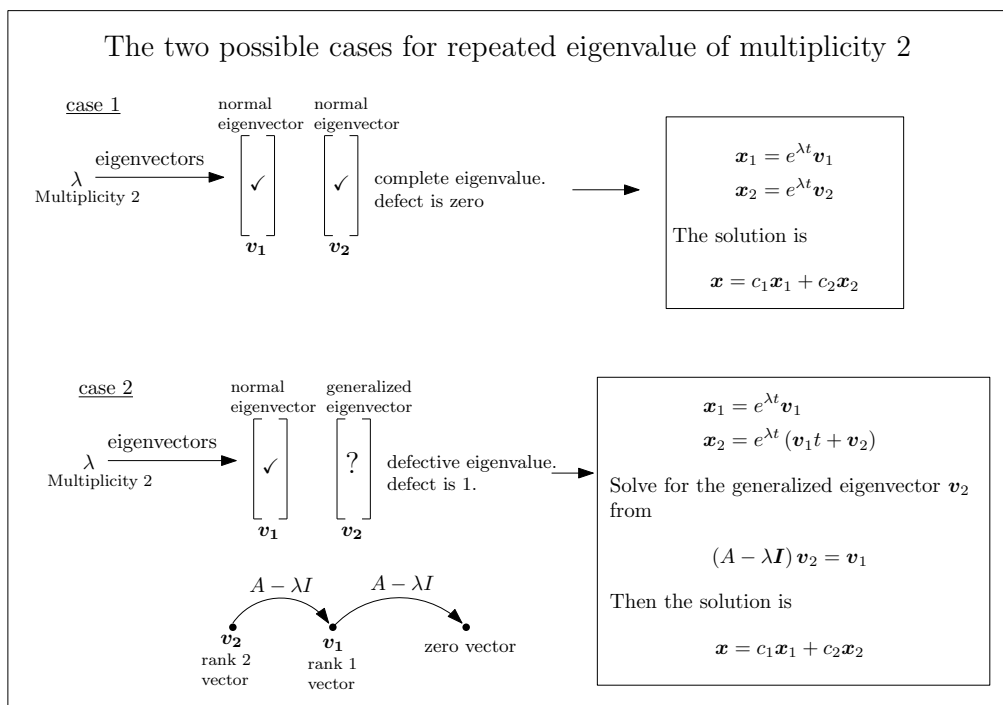


Figure 518: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} - (1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} 2 & -4 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} 3 \\ 1 \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue 1. Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} 2 \\ 1 \end{bmatrix} e^t \\ &= \begin{bmatrix} 2e^t \\ e^t \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} 2 \\ 1 \end{bmatrix} t + \begin{bmatrix} 3 \\ 1 \end{bmatrix} \right) e^t \\ &= \begin{bmatrix} e^t(2t + 3) \\ e^t(t + 1) \end{bmatrix}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} 2e^t \\ e^t \end{bmatrix} + c_2 \begin{bmatrix} e^t(2t + 3) \\ e^t(t + 1) \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} ((2t + 3)c_2 + 2c_1)e^t \\ e^t(c_2 t + c_1 + c_2) \end{bmatrix}$$

The following is the phase plot of the system.

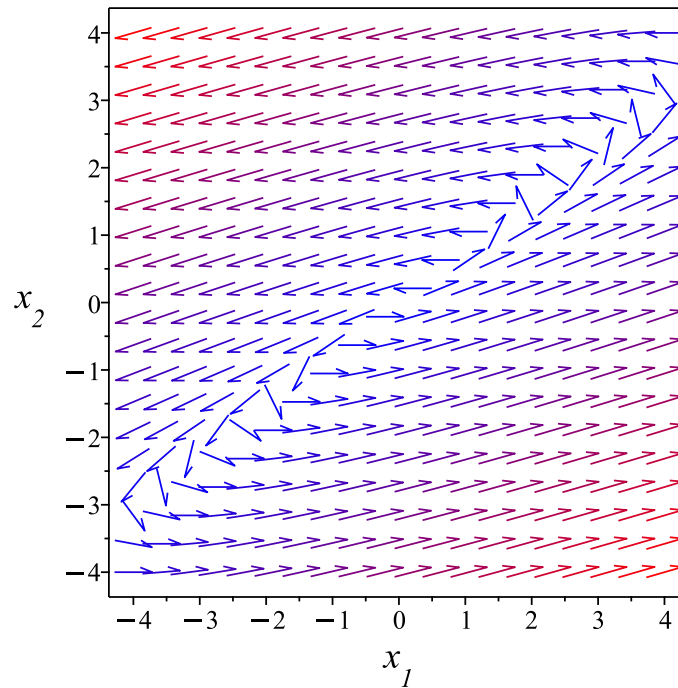


Figure 519: Phase plot

17.1.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 3x_1(t) - 4x_2(t), x_2'(t) = x_1(t) - x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[1, \begin{bmatrix} 2 \\ 1 \end{bmatrix} \right], \left[1, \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair, with eigenvalue of algebraic multiplicity 2

$$\left[1, \begin{bmatrix} 2 \\ 1 \end{bmatrix} \right]$$

- First solution from eigenvalue 1

$$\underline{x}^{\rightarrow}_1(t) = e^t \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

- Form of the 2nd homogeneous solution where \vec{p} is to be solved for, $\lambda = 1$ is the eigenvalue, and

$$\underline{x}^{\rightarrow}_2(t) = e^{\lambda t} (t\vec{v} + \vec{p})$$

- Note that the t multiplying \vec{v} makes this solution linearly independent to the 1st solution obtained
- Substitute $\underline{x}^{\rightarrow}_2(t)$ into the homogeneous system

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = (e^{\lambda t} A) \cdot (t\vec{v} + \vec{p})$$

- Use the fact that \vec{v} is an eigenvector of A

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = e^{\lambda t} (\lambda t\vec{v} + A \cdot \vec{p})$$

- Simplify equation

$$\lambda \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Make use of the identity matrix I

$$(\lambda \cdot I) \cdot \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Condition \vec{p} must meet for $\underline{x}^{\rightarrow}_2(t)$ to be a solution to the homogeneous system

$$(A - \lambda \cdot I) \cdot \vec{p} = \vec{v}$$

- Choose \vec{p} to use in the second solution to the homogeneous system from eigenvalue 1

$$\left(\begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} - 1 \cdot \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \cdot \vec{p} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

- Choice of \vec{p}

$$\vec{p} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

- Second solution from eigenvalue 1

$$\underline{x}_2(t) = e^t \cdot \left(t \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \end{bmatrix} \right)$$

- General solution to the system of ODEs

$$\underline{x} = c_1 \underline{x}_1(t) + c_2 \underline{x}_2(t)$$

- Substitute solutions into the general solution

$$\underline{x} = c_1 e^t \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix} + c_2 e^t \cdot \left(t \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \end{bmatrix} \right)$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^t(2c_2t + 2c_1 + c_2) \\ e^t(c_2t + c_1) \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = e^t(2c_2t + 2c_1 + c_2), x_2(t) = e^t(c_2t + c_1)\}$$

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 31

```
dsolve([diff(x__1(t),t)=3*x__1(t)-4*x__2(t),diff(x__2(t),t)=1*x__1(t)-1*x__2(t)],singsol=all
```

$$x_1(t) = e^t(c_2t + c_1)$$

$$x_2(t) = \frac{e^t(2c_2t + 2c_1 - c_2)}{4}$$

✓ Solution by Mathematica

Time used: 0.002 (sec). Leaf size: 41

```
DSolve[{x1'[t]==3*x1[t]-4*x2[t],x2'[t]==1*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$x1(t) \rightarrow e^t(2c_1t - 4c_2t + c_1)$$

$$x2(t) \rightarrow e^t((c_1 - 2c_2)t + c_2)$$

17.2 problem 2

- 17.2.1 Solution using Matrix exponential method 3682
- 17.2.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3683
- 17.2.3 Maple step by step solution 3688

Internal problem ID [767]

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Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 4x_1(t) - 2x_2(t) \\x_2'(t) &= 8x_1(t) - 4x_2(t)\end{aligned}$$

17.2.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} 4t + 1 & -2t \\ 8t & 1 - 4t \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}\vec{x}_h(t) &= e^{At}\vec{c} \\ &= \begin{bmatrix} 4t+1 & -2t \\ 8t & 1-4t \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} (4t+1)c_1 - 2tc_2 \\ 8tc_1 + (1-4t)c_2 \end{bmatrix} \\ &= \begin{bmatrix} (4c_1 - 2c_2)t + c_1 \\ (8c_1 - 4c_2)t + c_2 \end{bmatrix}\end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.2.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 4-\lambda & -2 \\ 8 & -4-\lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 0$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
0	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = 0$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} - (0) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 4 & -2 & 0 \\ 8 & -4 & 0 \end{array} \right]$$

$$R_2 = R_2 - 2R_1 \implies \left[\begin{array}{cc|c} 4 & -2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 4 & -2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{t}{2}\}$

Hence the solution is

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 2 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
0	2	1	Yes	$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue 0 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

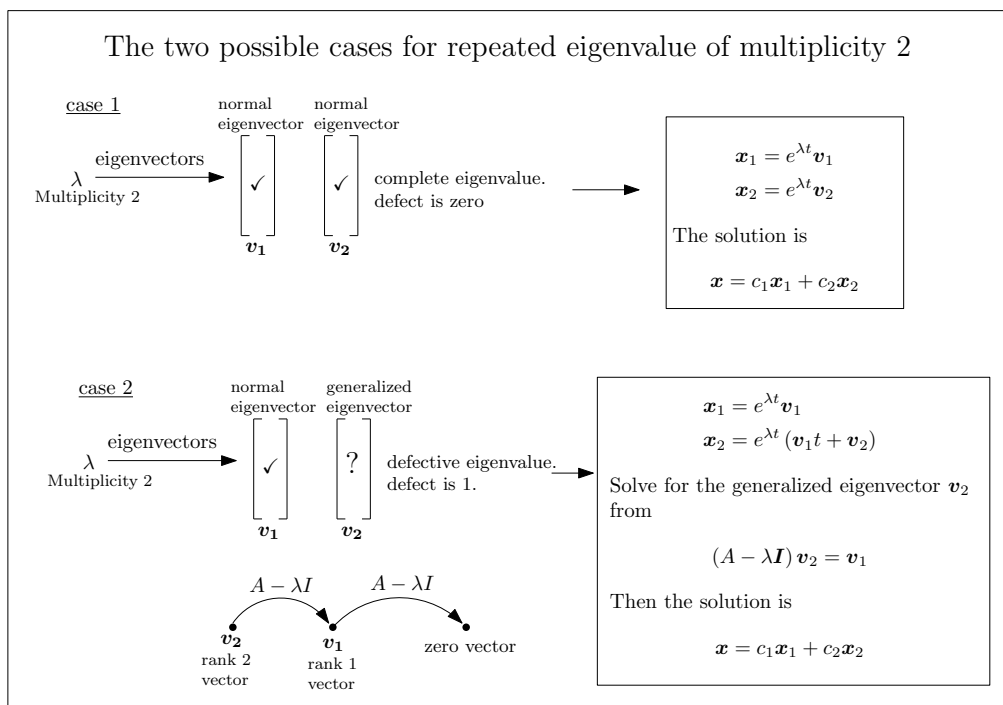


Figure 520: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} - (0) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} 1 \\ \frac{7}{4} \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue 0. Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} e^{0t} \\ &= \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} t + \begin{bmatrix} 1 \\ \frac{7}{4} \end{bmatrix} \right) e^{0t} \\ &= \begin{bmatrix} \frac{t}{2} + 1 \\ t + \frac{7}{4} \end{bmatrix}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} + c_2 \begin{bmatrix} \frac{t}{2} + 1 \\ t + \frac{7}{4} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{1}{2}c_1 + \frac{1}{2}c_2 t + c_2 \\ c_1 + c_2 t + \frac{7}{4}c_2 \end{bmatrix}$$

The following is the phase plot of the system.

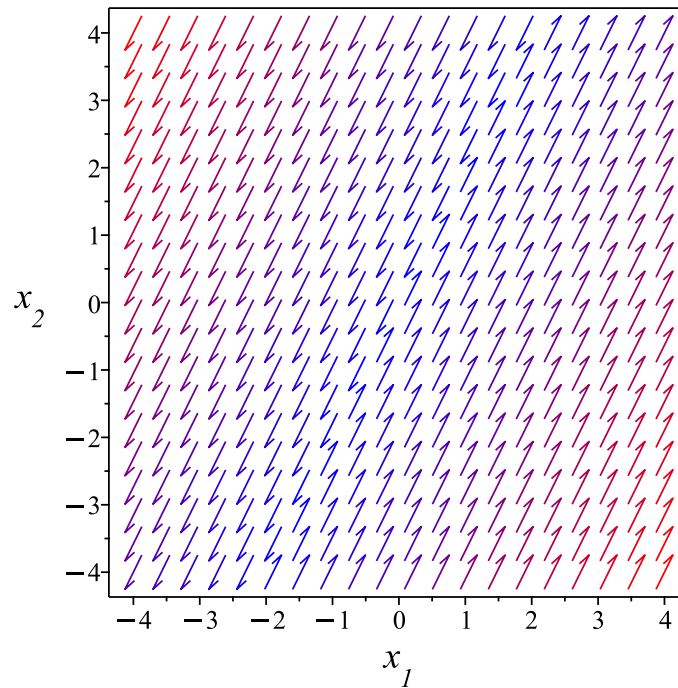


Figure 521: Phase plot

17.2.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 4x_1(t) - 2x_2(t), x_2'(t) = 8x_1(t) - 4x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow}{}'(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[0, \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} \right], \left[0, \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[0, \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}{}_{-1} = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[0, \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}{}_{-2} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}{}_{-1} + c_2 \underline{x}^{\rightarrow}{}_{-2}$$

- Substitute solutions into the general solution

$$\underline{x}^{\rightarrow} = \begin{bmatrix} \frac{c_1}{2} \\ c_1 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{c_1}{2} \\ c_1 \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = \frac{c_1}{2}, x_2(t) = c_1\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 24

```
dsolve([diff(x__1(t),t)=4*x__1(t)-2*x__2(t),diff(x__2(t),t)=8*x__1(t)-4*x__2(t)],singsol=all
```

$$\begin{aligned} x_1(t) &= c_1 t + c_2 \\ x_2(t) &= -\frac{1}{2}c_1 + 2c_1 t + 2c_2 \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.002 (sec). Leaf size: 34

```
DSolve[{x1'[t]==4*x1[t]-2*x2[t],x2'[t]==8*x1[t]-4*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$\begin{aligned} x_1(t) &\rightarrow 4c_1 t - 2c_2 t + c_1 \\ x_2(t) &\rightarrow 8c_1 t - 4c_2 t + c_2 \end{aligned}$$

17.3 problem 3

- 17.3.1 Solution using Matrix exponential method 3691
- 17.3.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3692
- 17.3.3 Maple step by step solution 3697

Internal problem ID [768]

Internal file name [OUTPUT/768_Sunday_June_05_2022_01_49_11_AM_71145901/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 3.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$\begin{aligned}x_1'(t) &= -\frac{3x_1(t)}{2} + x_2(t) \\x_2'(t) &= -\frac{x_1(t)}{4} - \frac{x_2(t)}{2}\end{aligned}$$

17.3.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{3}{2} & 1 \\ -\frac{1}{4} & -\frac{1}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-t}\left(-\frac{t}{2} + 1\right) & t e^{-t} \\ -\frac{t e^{-t}}{4} & e^{-t}\left(\frac{t}{2} + 1\right) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{-t} \left(-\frac{t}{2} + 1\right) & t e^{-t} \\ -\frac{t e^{-t}}{4} & e^{-t} \left(\frac{t}{2} + 1\right) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-t} \left(-\frac{t}{2} + 1\right) c_1 + t e^{-t} c_2 \\ -\frac{t e^{-t} c_1}{4} + e^{-t} \left(\frac{t}{2} + 1\right) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-t} \left(-\frac{1}{2} t c_1 + c_1 + c_2 t\right) \\ -\frac{((-2t-4)c_2 + t c_1) e^{-t}}{4} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.3.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{3}{2} & 1 \\ -\frac{1}{4} & -\frac{1}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} -\frac{3}{2} & 1 \\ -\frac{1}{4} & -\frac{1}{2} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -\frac{3}{2} - \lambda & 1 \\ -\frac{1}{4} & -\frac{1}{2} - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 2\lambda + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{3}{2} & 1 \\ -\frac{1}{4} & -\frac{1}{2} \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -\frac{1}{2} & 1 \\ -\frac{1}{4} & \frac{1}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -\frac{1}{2} & 1 & 0 \\ -\frac{1}{4} & \frac{1}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{R_1}{2} \implies \left[\begin{array}{cc|c} -\frac{1}{2} & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -\frac{1}{2} & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 2t\}$

Hence the solution is

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = \begin{bmatrix} 2t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = t \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-1	2	1	Yes	$\begin{bmatrix} 2 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue -1 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

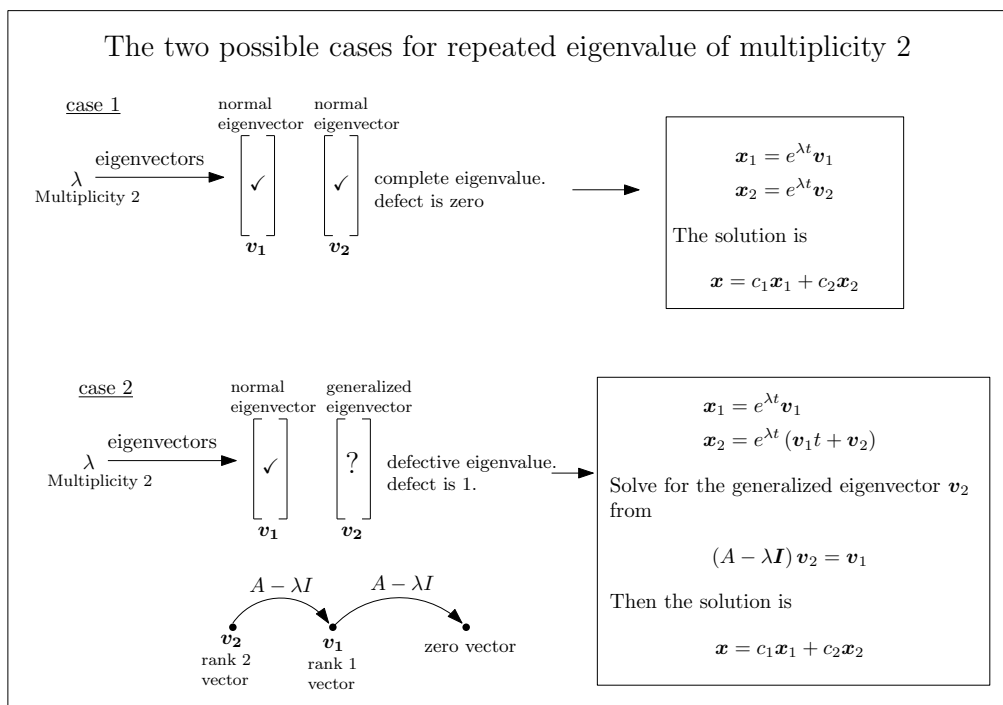


Figure 522: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} -\frac{3}{2} & 1 \\ -\frac{1}{4} & -\frac{1}{2} \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} -\frac{1}{2} & 1 \\ -\frac{1}{4} & \frac{1}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} -2 \\ 1 \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue -1 . Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} 2 \\ 1 \end{bmatrix} e^{-t} \\ &= \begin{bmatrix} 2e^{-t} \\ e^{-t} \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} 2 \\ 1 \end{bmatrix} t + \begin{bmatrix} -2 \\ 1 \end{bmatrix} \right) e^{-t} \\ &= \begin{bmatrix} 2(-1+t)e^{-t} \\ e^{-t}(t+1) \end{bmatrix}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} 2e^{-t} \\ e^{-t} \end{bmatrix} + c_2 \begin{bmatrix} e^{-t}(2t-2) \\ e^{-t}(t+1) \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} 2((-1+t)c_2 + c_1)e^{-t} \\ e^{-t}(c_2 t + c_1 + c_2) \end{bmatrix}$$

The following is the phase plot of the system.

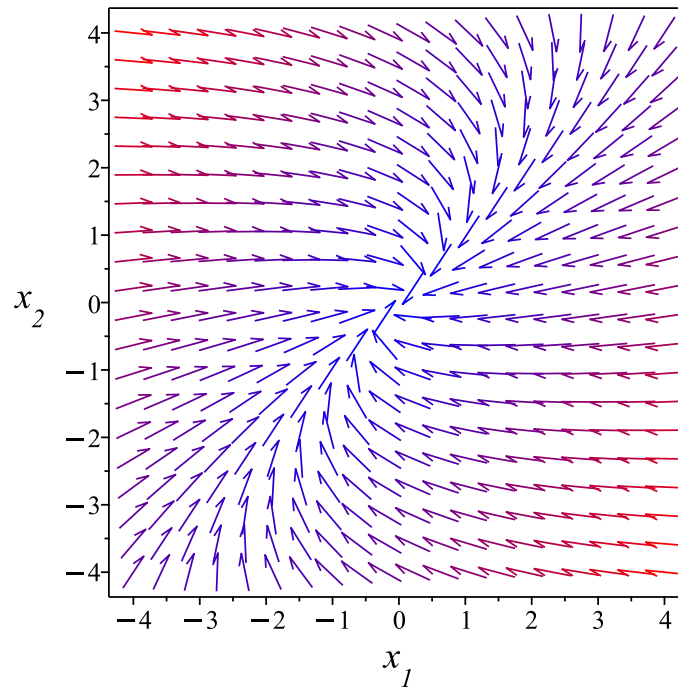


Figure 523: Phase plot

17.3.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = -\frac{3x_1(t)}{2} + x_2(t), x_2'(t) = -\frac{x_1(t)}{4} - \frac{x_2(t)}{2} \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} -\frac{3}{2} & 1 \\ -\frac{1}{4} & -\frac{1}{2} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} -\frac{3}{2} & 1 \\ -\frac{1}{4} & -\frac{1}{2} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -\frac{3}{2} & 1 \\ -\frac{1}{4} & -\frac{1}{2} \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-1, \begin{bmatrix} 2 \\ 1 \end{bmatrix} \right], \left[-1, \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair, with eigenvalue of algebraic multiplicity 2

$$\left[-1, \begin{bmatrix} 2 \\ 1 \end{bmatrix} \right]$$

- First solution from eigenvalue -1

$$\underline{x}^{\rightarrow}_1(t) = e^{-t} \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

- Form of the 2nd homogeneous solution where \vec{p} is to be solved for, $\lambda = -1$ is the eigenvalue, and

$$\underline{x}^{\rightarrow}_2(t) = e^{\lambda t} (t\vec{v} + \vec{p})$$

- Note that the t multiplying \vec{v} makes this solution linearly independent to the 1st solution obtained
- Substitute $\underline{x}^{\rightarrow}_2(t)$ into the homogeneous system

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = (e^{\lambda t} A) \cdot (t\vec{v} + \vec{p})$$

- Use the fact that \vec{v} is an eigenvector of A

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = e^{\lambda t} (\lambda t\vec{v} + A \cdot \vec{p})$$

- Simplify equation

$$\lambda \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Make use of the identity matrix I

$$(\lambda \cdot I) \cdot \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Condition \vec{p} must meet for $\underline{x}^{\rightarrow}_2(t)$ to be a solution to the homogeneous system

$$(A - \lambda \cdot I) \cdot \vec{p} = \vec{v}$$

- Choose \vec{p} to use in the second solution to the homogeneous system from eigenvalue -1

$$\left(\begin{bmatrix} -\frac{3}{2} & 1 \\ -\frac{1}{4} & -\frac{1}{2} \end{bmatrix} - (-1) \cdot \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \cdot \vec{p} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

- Choice of \vec{p}

$$\vec{p} = \begin{bmatrix} -4 \\ 0 \end{bmatrix}$$

- Second solution from eigenvalue -1

$$\underline{x}_2(t) = e^{-t} \cdot \left(t \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix} + \begin{bmatrix} -4 \\ 0 \end{bmatrix} \right)$$

- General solution to the system of ODEs

$$\underline{x} = c_1 \underline{x}_1(t) + c_2 \underline{x}_2(t)$$

- Substitute solutions into the general solution

$$\underline{x} = c_1 e^{-t} \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix} + c_2 e^{-t} \cdot \left(t \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix} + \begin{bmatrix} -4 \\ 0 \end{bmatrix} \right)$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} 2((t-2)c_2 + c_1)e^{-t} \\ (c_2t + c_1)e^{-t} \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = 2((t-2)c_2 + c_1)e^{-t}, x_2(t) = (c_2t + c_1)e^{-t}\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 32

```
dsolve([diff(x__1(t),t)=-3/2*x__1(t)+1*x__2(t),diff(x__2(t),t)=-1/4*x__1(t)-1/2*x__2(t)],sin
```

$$x_1(t) = e^{-t}(c_2t + c_1)$$

$$x_2(t) = \frac{e^{-t}(c_2t + c_1 + 2c_2)}{2}$$

✓ Solution by Mathematica

Time used: 0.002 (sec). Leaf size: 54

```
DSolve[{x1'[t]==-3/2*x1[t]+1*x2[t],x2'[t]==-1/4*x1[t]-1/2*x2[t]},{x1[t],x2[t]},t,IncludeSing
```

$$x1(t) \rightarrow \frac{1}{2}e^{-t}(2c_2t - c_1(t - 2))$$
$$x2(t) \rightarrow \frac{1}{4}e^{-t}(c_1(-t) + 2c_2t + 4c_2)$$

17.4 problem 4

- 17.4.1 Solution using Matrix exponential method 3701
- 17.4.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3702
- 17.4.3 Maple step by step solution 3707

Internal problem ID [769]

Internal file name [OUTPUT/769_Sunday_June_05_2022_01_49_12_AM_14655444/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 4.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -3x_1(t) + \frac{5x_2(t)}{2} \\x_2'(t) &= -\frac{5x_1(t)}{2} + 2x_2(t)\end{aligned}$$

17.4.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -3 & \frac{5}{2} \\ -\frac{5}{2} & 2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-\frac{t}{2}}\left(1 - \frac{5t}{2}\right) & \frac{5te^{-\frac{t}{2}}}{2} \\ -\frac{5te^{-\frac{t}{2}}}{2} & e^{-\frac{t}{2}}\left(1 + \frac{5t}{2}\right) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{-\frac{t}{2}} \left(1 - \frac{5t}{2}\right) & \frac{5t e^{-\frac{t}{2}}}{2} \\ -\frac{5t e^{-\frac{t}{2}}}{2} & e^{-\frac{t}{2}} \left(1 + \frac{5t}{2}\right) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-\frac{t}{2}} \left(1 - \frac{5t}{2}\right) c_1 + \frac{5t e^{-\frac{t}{2}} c_2}{2} \\ -\frac{5t e^{-\frac{t}{2}} c_1}{2} + e^{-\frac{t}{2}} \left(1 + \frac{5t}{2}\right) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \frac{((2-5t)c_1 + 5c_2 t) e^{-\frac{t}{2}}}{2} \\ \frac{((5t+2)c_2 - 5t c_1) e^{-\frac{t}{2}}}{2} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.4.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -3 & \frac{5}{2} \\ -\frac{5}{2} & 2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} -3 & \frac{5}{2} \\ -\frac{5}{2} & 2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -3 - \lambda & \frac{5}{2} \\ -\frac{5}{2} & 2 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + \lambda + \frac{1}{4} = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -\frac{1}{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-\frac{1}{2}$	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -\frac{1}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -3 & \frac{5}{2} \\ -\frac{5}{2} & 2 \end{bmatrix} - \left(-\frac{1}{2}\right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -\frac{5}{2} & \frac{5}{2} \\ -\frac{5}{2} & \frac{5}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -\frac{5}{2} & \frac{5}{2} & 0 \\ -\frac{5}{2} & \frac{5}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} -\frac{5}{2} & \frac{5}{2} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -\frac{5}{2} & \frac{5}{2} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-\frac{1}{2}$	2	1	Yes	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue $-\frac{1}{2}$ is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

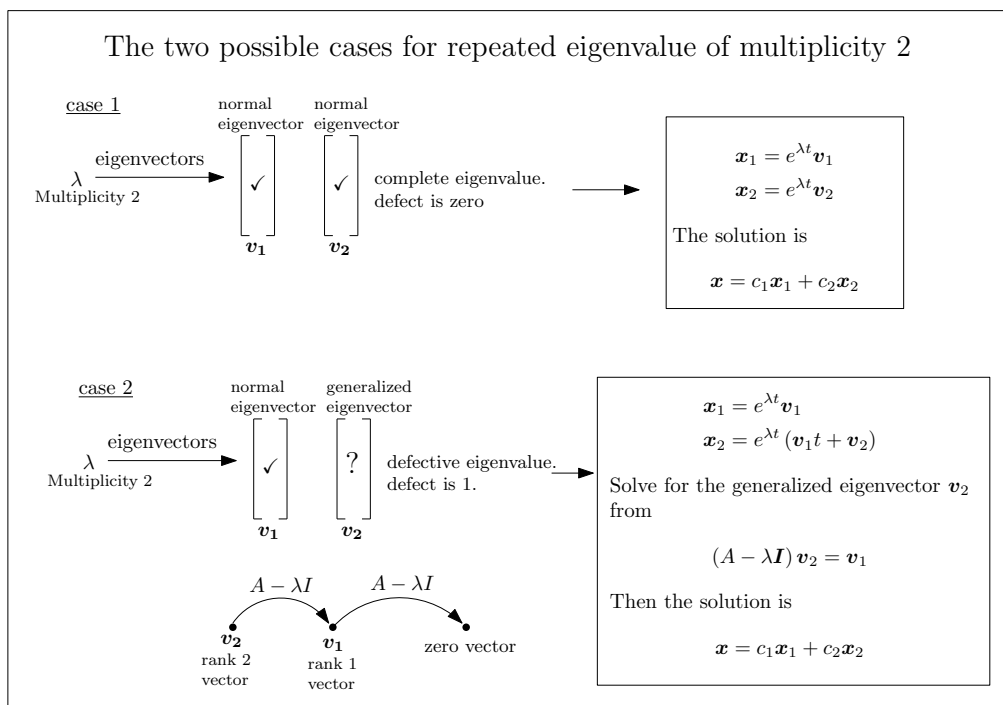


Figure 524: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\left(\begin{bmatrix} -3 & \frac{5}{2} \\ -\frac{5}{2} & 2 \end{bmatrix} - \left(-\frac{1}{2} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right)$$

$$\begin{bmatrix} -\frac{5}{2} & \frac{5}{2} \\ -\frac{5}{2} & \frac{5}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} \frac{3}{5} \\ 1 \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue $-\frac{1}{2}$. Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{-\frac{t}{2}} \\ &= \begin{bmatrix} e^{-\frac{t}{2}} \\ e^{-\frac{t}{2}} \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} 1 \\ 1 \end{bmatrix} t + \begin{bmatrix} \frac{3}{5} \\ 1 \end{bmatrix} \right) e^{-\frac{t}{2}} \\ &= \begin{bmatrix} \frac{e^{-\frac{t}{2}}(5t+3)}{5} \\ e^{-\frac{t}{2}}(t+1) \end{bmatrix}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} e^{-\frac{t}{2}} \\ e^{-\frac{t}{2}} \end{bmatrix} + c_2 \begin{bmatrix} e^{-\frac{t}{2}}(t + \frac{3}{5}) \\ e^{-\frac{t}{2}}(t + 1) \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^{-\frac{t}{2}}(c_1 + c_2 t + \frac{3}{5}c_2) \\ e^{-\frac{t}{2}}(c_2 t + c_1 + c_2) \end{bmatrix}$$

The following is the phase plot of the system.

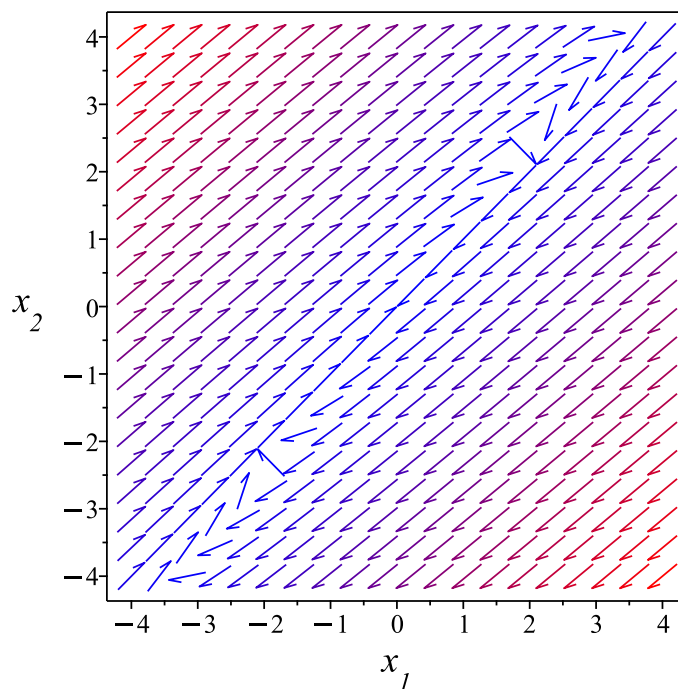


Figure 525: Phase plot

17.4.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = -3x_1(t) + \frac{5x_2(t)}{2}, x_2'(t) = -\frac{5x_1(t)}{2} + 2x_2(t) \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} -3 & \frac{5}{2} \\ -\frac{5}{2} & 2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} -3 & \frac{5}{2} \\ -\frac{5}{2} & 2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -3 & \frac{5}{2} \\ -\frac{5}{2} & 2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-\frac{1}{2}, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right], \left[-\frac{1}{2}, \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair, with eigenvalue of algebraic multiplicity 2

$$\left[-\frac{1}{2}, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right]$$

- First solution from eigenvalue $-\frac{1}{2}$

$$\underline{x}^{\rightarrow}_1(t) = e^{-\frac{t}{2}} \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- Form of the 2nd homogeneous solution where \vec{p} is to be solved for, $\lambda = -\frac{1}{2}$ is the eigenvalue, and

$$\underline{x}^{\rightarrow}_2(t) = e^{\lambda t} (t\vec{v} + \vec{p})$$

- Note that the t multiplying \vec{v} makes this solution linearly independent to the 1st solution obtained

- Substitute $\underline{x}^{\rightarrow}_2(t)$ into the homogeneous system

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = (e^{\lambda t} A) \cdot (t\vec{v} + \vec{p})$$

- Use the fact that \vec{v} is an eigenvector of A

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = e^{\lambda t} (\lambda t\vec{v} + A \cdot \vec{p})$$

- Simplify equation

$$\lambda \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Make use of the identity matrix I

$$(\lambda \cdot I) \cdot \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Condition \vec{p} must meet for $\underline{x}^{\rightarrow}_2(t)$ to be a solution to the homogeneous system

$$(A - \lambda \cdot I) \cdot \vec{p} = \vec{v}$$

- Choose \vec{p} to use in the second solution to the homogeneous system from eigenvalue $-\frac{1}{2}$

$$\left(\begin{bmatrix} -3 & \frac{5}{2} \\ -\frac{5}{2} & 2 \end{bmatrix} - -\frac{1}{2} \cdot \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \cdot \vec{p} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- Choice of \vec{p}

$$\vec{p} = \begin{bmatrix} -\frac{2}{5} \\ 0 \end{bmatrix}$$

- Second solution from eigenvalue $-\frac{1}{2}$

$$\underline{x}_2(t) = e^{-\frac{t}{2}} \cdot \left(t \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix} + \begin{bmatrix} -\frac{2}{5} \\ 0 \end{bmatrix} \right)$$

- General solution to the system of ODEs

$$\underline{x} = c_1 \underline{x}_1(t) + c_2 \underline{x}_2(t)$$

- Substitute solutions into the general solution

$$\underline{x} = c_1 e^{-\frac{t}{2}} \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix} + c_2 e^{-\frac{t}{2}} \cdot \left(t \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix} + \begin{bmatrix} -\frac{2}{5} \\ 0 \end{bmatrix} \right)$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^{-\frac{t}{2}}(c_1 + c_2 t - \frac{2}{5}c_2) \\ e^{-\frac{t}{2}}(c_2 t + c_1) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = e^{-\frac{t}{2}}(c_1 + c_2 t - \frac{2}{5}c_2), x_2(t) = e^{-\frac{t}{2}}(c_2 t + c_1) \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 35

```
dsolve([diff(x__1(t),t)=-3*x__1(t)+5/2*x__2(t),diff(x__2(t),t)=-5/2*x__1(t)+2*x__2(t)],sings
```

$$x_1(t) = e^{-\frac{t}{2}}(c_2 t + c_1)$$

$$x_2(t) = \frac{e^{-\frac{t}{2}}(5c_2 t + 5c_1 + 2c_2)}{5}$$

✓ Solution by Mathematica

Time used: 0.002 (sec). Leaf size: 59

```
DSolve[{x1'[t]==-3*x1[t]+5/2*x2[t],x2'[t]==-5/2*x1[t]+2*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolutions->True]
```

$$x1(t) \rightarrow \frac{1}{2}e^{-t/2}(c_1(2 - 5t) + 5c_2t)$$

$$x2(t) \rightarrow \frac{1}{2}e^{-t/2}(-5c_1t + 5c_2t + 2c_2)$$

17.5 problem 5

- 17.5.1 Solution using Matrix exponential method 3711
- 17.5.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3712
- 17.5.3 Maple step by step solution 3719

Internal problem ID [770]

Internal file name [OUTPUT/770_Sunday_June_05_2022_01_49_14_AM_32078012/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 5.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) + x_2(t) + x_3(t) \\x_2'(t) &= 2x_1(t) + x_2(t) - x_3(t) \\x_3'(t) &= -x_2(t) + x_3(t)\end{aligned}$$

17.5.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 \\ 2 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{e^{-t}}{3} + \frac{2e^{2t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} \\ \frac{(6t+4)e^{2t}}{9} - \frac{4e^{-t}}{9} & \frac{(3t+5)e^{2t}}{9} + \frac{4e^{-t}}{9} & \frac{e^{2t}(3t-4)}{9} + \frac{4e^{-t}}{9} \\ \frac{(-6t+2)e^{2t}}{9} - \frac{2e^{-t}}{9} & \frac{(-3t-2)e^{2t}}{9} + \frac{2e^{-t}}{9} & \frac{(-3t+7)e^{2t}}{9} + \frac{2e^{-t}}{9} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} \frac{e^{-t}}{3} + \frac{2e^{2t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} \\ \frac{(6t+4)e^{2t}}{9} - \frac{4e^{-t}}{9} & \frac{(3t+5)e^{2t}}{9} + \frac{4e^{-t}}{9} & \frac{e^{2t}(3t-4)}{9} + \frac{4e^{-t}}{9} \\ \frac{(-6t+2)e^{2t}}{9} - \frac{2e^{-t}}{9} & \frac{(-3t-2)e^{2t}}{9} + \frac{2e^{-t}}{9} & \frac{(-3t+7)e^{2t}}{9} + \frac{2e^{-t}}{9} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ c_3 \end{bmatrix} \\
 &= \begin{bmatrix} \left(\frac{e^{-t}}{3} + \frac{2e^{2t}}{3}\right) c_1 + \left(\frac{e^{2t}}{3} - \frac{e^{-t}}{3}\right) c_2 + \left(\frac{e^{2t}}{3} - \frac{e^{-t}}{3}\right) c_3 \\ \left(\frac{(6t+4)e^{2t}}{9} - \frac{4e^{-t}}{9}\right) c_1 + \left(\frac{(3t+5)e^{2t}}{9} + \frac{4e^{-t}}{9}\right) c_2 + \left(\frac{e^{2t}(3t-4)}{9} + \frac{4e^{-t}}{9}\right) c_3 \\ \left(\frac{(-6t+2)e^{2t}}{9} - \frac{2e^{-t}}{9}\right) c_1 + \left(\frac{(-3t-2)e^{2t}}{9} + \frac{2e^{-t}}{9}\right) c_2 + \left(\frac{(-3t+7)e^{2t}}{9} + \frac{2e^{-t}}{9}\right) c_3 \end{bmatrix} \\
 &= \begin{bmatrix} \frac{(c_1 - c_2 - c_3)e^{-t}}{3} + \frac{2(c_1 + \frac{c_2}{2} + \frac{c_3}{2})e^{2t}}{3} \\ \frac{((6t+4)c_1 + (3t+5)c_2 + c_3(3t-4))e^{2t}}{9} - \frac{4(c_1 - c_2 - c_3)e^{-t}}{9} \\ \frac{((-6t+2)c_1 + (-3t-2)c_2 + (-3t+7)c_3)e^{2t}}{9} - \frac{2(c_1 - c_2 - c_3)e^{-t}}{9} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.5.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 \\ 2 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 1 & 1 & 1 \\ 2 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 1-\lambda & 1 & 1 \\ 2 & 1-\lambda & -1 \\ 0 & -1 & 1-\lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^3 - 3\lambda^2 + 4 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1$$

$$\lambda_2 = 2$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
2	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 1 & 1 \\ 2 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 & 1 & 1 \\ 2 & 2 & -1 \\ 0 & -1 & 2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} 2 & 1 & 1 & 0 \\ 2 & 2 & -1 & 0 \\ 0 & -1 & 2 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{ccc|c} 2 & 1 & 1 & 0 \\ 0 & 1 & -2 & 0 \\ 0 & -1 & 2 & 0 \end{array} \right]$$

$$R_3 = R_3 + R_2 \implies \left[\begin{array}{ccc|c} 2 & 1 & 1 & 0 \\ 0 & 1 & -2 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 & 1 & 1 \\ 0 & 1 & -2 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -\frac{3t}{2}, v_2 = 2t\}$

Hence the solution is

$$\begin{bmatrix} -\frac{3t}{2} \\ 2t \\ t \end{bmatrix} = \begin{bmatrix} -\frac{3t}{2} \\ 2t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{3t}{2} \\ 2t \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{3}{2} \\ 2 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{3t}{2} \\ 2t \\ t \end{bmatrix} = \begin{bmatrix} -\frac{3}{2} \\ 2 \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{3t}{2} \\ 2t \\ t \end{bmatrix} = \begin{bmatrix} -3 \\ 4 \\ 2 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 1 & 1 \\ 2 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix} - (2) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 & 1 & 1 \\ 2 & -1 & -1 \\ 0 & -1 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -1 & 1 & 1 & 0 \\ 2 & -1 & -1 & 0 \\ 0 & -1 & -1 & 0 \end{array} \right]$$

$$R_2 = R_2 + 2R_1 \implies \left[\begin{array}{ccc|c} -1 & 1 & 1 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & -1 & -1 & 0 \end{array} \right]$$

$$R_3 = R_3 + R_2 \implies \left[\begin{array}{ccc|c} -1 & 1 & 1 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 & 1 & 1 \\ 0 & 1 & 1 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 0, v_2 = -t\}$

Hence the solution is

$$\begin{bmatrix} 0 \\ -t \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ -t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 0 \\ -t \\ t \end{bmatrix} = t \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-1	1	1	No	$\begin{bmatrix} -\frac{3}{2} \\ 2 \\ 1 \end{bmatrix}$
2	2	1	Yes	$\begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{-t} \\ &= \begin{bmatrix} -\frac{3}{2} \\ 2 \\ 1 \end{bmatrix} e^{-t}\end{aligned}$$

eigenvalue 2 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

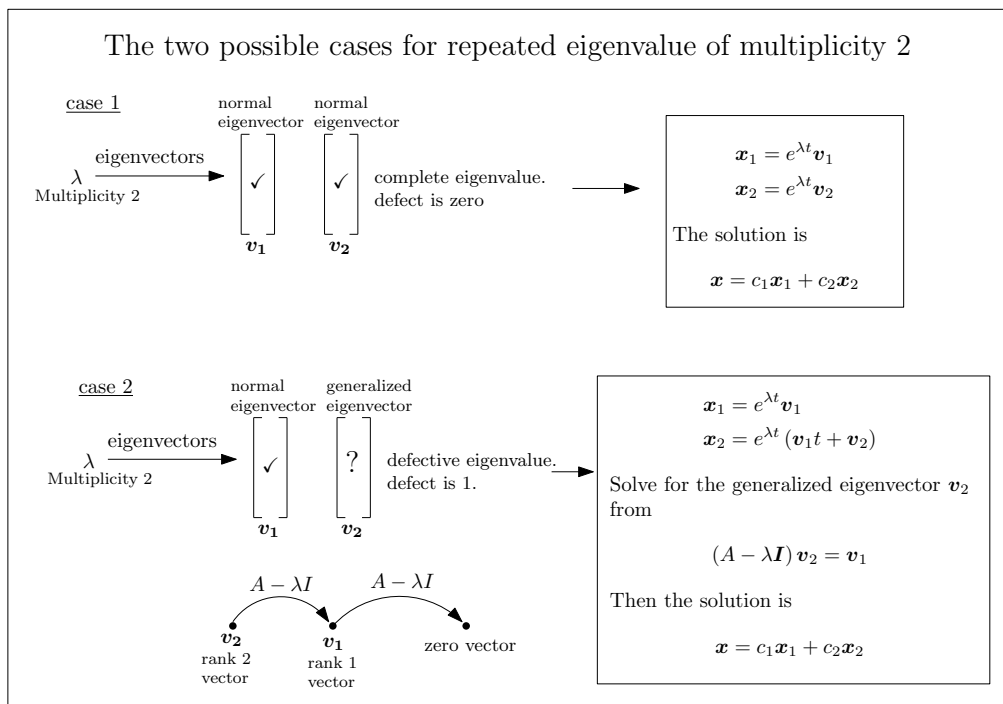


Figure 526: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} 1 & 1 & 1 \\ 2 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix} - (2) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} -1 & 1 & 1 \\ 2 & -1 & -1 \\ 0 & -1 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} -1 \\ -2 \\ 1 \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue 2. Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned} \vec{x}_2(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix} e^{2t} \\ &= \begin{bmatrix} 0 \\ -e^{2t} \\ e^{2t} \end{bmatrix} \end{aligned}$$

And

$$\begin{aligned} \vec{x}_3(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix} t + \begin{bmatrix} -1 \\ -2 \\ 1 \end{bmatrix} \right) e^{2t} \\ &= \begin{bmatrix} -e^{2t} \\ -e^{2t}(2+t) \\ e^{2t}(t+1) \end{bmatrix} \end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = c_1 \begin{bmatrix} -\frac{3e^{-t}}{2} \\ 2e^{-t} \\ e^{-t} \end{bmatrix} + c_2 \begin{bmatrix} 0 \\ -e^{2t} \\ e^{2t} \end{bmatrix} + c_3 \begin{bmatrix} -e^{2t} \\ e^{2t}(-t-2) \\ e^{2t}(t+1) \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} -\frac{3c_1 e^{-t}}{2} - c_3 e^{2t} \\ ((-t-2)c_3 - c_2)e^{2t} + 2c_1 e^{-t} \\ ((t+1)c_3 + c_2)e^{2t} + c_1 e^{-t} \end{bmatrix}$$

17.5.3 Maple step by step solution

Let's solve

$$[x_1'(t) = x_1(t) + x_2(t) + x_3(t), x_2'(t) = 2x_1(t) + x_2(t) - x_3(t), x_3'(t) = -x_2(t) + x_3(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 1 & 1 & 1 \\ 2 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 1 & 1 & 1 \\ 2 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 1 & 1 & 1 \\ 2 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-1, \begin{bmatrix} -\frac{3}{2} \\ 2 \\ 1 \end{bmatrix} \right], \left[2, \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix} \right], \left[2, \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-1, \begin{bmatrix} -\frac{3}{2} \\ 2 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{-t} \cdot \begin{bmatrix} -\frac{3}{2} \\ 2 \\ 1 \end{bmatrix}$$

- Consider eigenpair, with eigenvalue of algebraic multiplicity 2

$$\left[2, \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix} \right]$$

- First solution from eigenvalue 2

$$\underline{x}^{\rightarrow}_2(t) = e^{2t} \cdot \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$$

- Form of the 2nd homogeneous solution where \vec{p} is to be solved for, $\lambda = 2$ is the eigenvalue, and

$$\underline{x}_{\rightarrow 3}(t) = e^{\lambda t} (t\vec{v} + \vec{p})$$

- Note that the t multiplying \vec{v} makes this solution linearly independent to the 1st solution obtained
- Substitute $\underline{x}_{\rightarrow 3}(t)$ into the homogeneous system

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = (e^{\lambda t} A) \cdot (t\vec{v} + \vec{p})$$

- Use the fact that \vec{v} is an eigenvector of A

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = e^{\lambda t} (\lambda t\vec{v} + A \cdot \vec{p})$$

- Simplify equation

$$\lambda \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Make use of the identity matrix I

$$(\lambda \cdot I) \cdot \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Condition \vec{p} must meet for $\underline{x}_{\rightarrow 3}(t)$ to be a solution to the homogeneous system

$$(A - \lambda \cdot I) \cdot \vec{p} = \vec{v}$$

- Choose \vec{p} to use in the second solution to the homogeneous system from eigenvalue 2

$$\left(\begin{bmatrix} 1 & 1 & 1 \\ 2 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix} - 2 \cdot \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \cdot \vec{p} = \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$$

- Choice of \vec{p}

$$\vec{p} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

- Second solution from eigenvalue 2

$$\underline{x}_{\rightarrow 3}(t) = e^{2t} \cdot \left(t \cdot \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \right)$$

- General solution to the system of ODEs

$$\underline{x}_{\rightarrow} = c_1 \underline{x}_{\rightarrow 1} + c_2 \underline{x}_{\rightarrow 2}(t) + c_3 \underline{x}_{\rightarrow 3}(t)$$

- Substitute solutions into the general solution

$$\vec{x} = c_1 e^{-t} \cdot \begin{bmatrix} -\frac{3}{2} \\ 2 \\ 1 \end{bmatrix} + c_2 e^{2t} \cdot \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix} + c_3 e^{2t} \cdot \left(t \cdot \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \right)$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} -\frac{3c_1 e^{-t}}{2} \\ (-c_3 t - c_2) e^{2t} + 2c_1 e^{-t} \\ (c_3 t + c_2) e^{2t} + c_1 e^{-t} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = -\frac{3c_1 e^{-t}}{2}, x_2(t) = (-c_3 t - c_2) e^{2t} + 2c_1 e^{-t}, x_3(t) = (c_3 t + c_2) e^{2t} + c_1 e^{-t} \right\}$$

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 76

```
dsolve([diff(x__1(t),t)=1*x__1(t)+1*x__2(t)+1*x__3(t),diff(x__2(t),t)=2*x__1(t)+1*x__2(t)-1*
```

$$\begin{aligned} x_1(t) &= -\frac{3e^{-t}c_1}{2} - c_3 e^{2t} \\ x_2(t) &= 2e^{-t}c_1 - c_2 e^{2t} - e^{2t}c_3 t - c_3 e^{2t} \\ x_3(t) &= e^{-t}c_1 + c_2 e^{2t} + e^{2t}c_3 t \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.01 (sec). Leaf size: 164

```
DSolve[{x1'[t]==1*x1[t]+1*x2[t]+1*x3[t],x2'[t]==2*x1[t]+1*x2[t]-1*x3[t],x3'[t]==0*x1[t]-1*x2
```

$$\begin{aligned} x_1(t) &\rightarrow \frac{1}{3}e^{-t}(c_1(2e^{3t} + 1) + (c_2 + c_3)(e^{3t} - 1)) \\ x_2(t) &\rightarrow \frac{1}{9}e^{-t}(c_1(e^{3t}(6t + 4) - 4) + c_2(e^{3t}(3t + 5) + 4) + c_3(e^{3t}(3t - 4) + 4)) \\ x_3(t) &\rightarrow \frac{1}{9}e^{-t}(c_1(e^{3t}(2 - 6t) - 2) + c_2(2 - e^{3t}(3t + 2)) - c_3(e^{3t}(3t - 7) - 2)) \end{aligned}$$

17.6 problem 6

- 17.6.1 Solution using Matrix exponential method 3723
- 17.6.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3724
- 17.6.3 Maple step by step solution 3731

Internal problem ID [771]

Internal file name [OUTPUT/771_Sunday_June_05_2022_01_49_16_AM_29590752/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 6.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$x_1'(t) = x_2(t) + x_3(t)$$

$$x_2'(t) = x_1(t) + x_3(t)$$

$$x_3'(t) = x_1(t) + x_2(t)$$

17.6.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{2e^{-t}}{3} + \frac{e^{2t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} \\ \frac{e^{2t}}{3} - \frac{e^{-t}}{3} & \frac{2e^{-t}}{3} + \frac{e^{2t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} \\ \frac{e^{2t}}{3} - \frac{e^{-t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} & \frac{2e^{-t}}{3} + \frac{e^{2t}}{3} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} \frac{2e^{-t}}{3} + \frac{e^{2t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} \\ \frac{e^{2t}}{3} - \frac{e^{-t}}{3} & \frac{2e^{-t}}{3} + \frac{e^{2t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} \\ \frac{e^{2t}}{3} - \frac{e^{-t}}{3} & \frac{e^{2t}}{3} - \frac{e^{-t}}{3} & \frac{2e^{-t}}{3} + \frac{e^{2t}}{3} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ c_3 \end{bmatrix} \\
 &= \begin{bmatrix} \left(\frac{2e^{-t}}{3} + \frac{e^{2t}}{3}\right) c_1 + \left(\frac{e^{2t}}{3} - \frac{e^{-t}}{3}\right) c_2 + \left(\frac{e^{2t}}{3} - \frac{e^{-t}}{3}\right) c_3 \\ \left(\frac{e^{2t}}{3} - \frac{e^{-t}}{3}\right) c_1 + \left(\frac{2e^{-t}}{3} + \frac{e^{2t}}{3}\right) c_2 + \left(\frac{e^{2t}}{3} - \frac{e^{-t}}{3}\right) c_3 \\ \left(\frac{e^{2t}}{3} - \frac{e^{-t}}{3}\right) c_1 + \left(\frac{e^{2t}}{3} - \frac{e^{-t}}{3}\right) c_2 + \left(\frac{2e^{-t}}{3} + \frac{e^{2t}}{3}\right) c_3 \end{bmatrix} \\
 &= \begin{bmatrix} \frac{(2c_1 - c_2 - c_3)e^{-t}}{3} + \frac{e^{2t}(c_1 + c_2 + c_3)}{3} \\ \frac{(-c_1 + 2c_2 - c_3)e^{-t}}{3} + \frac{e^{2t}(c_1 + c_2 + c_3)}{3} \\ \frac{(-c_1 - c_2 + 2c_3)e^{-t}}{3} + \frac{e^{2t}(c_1 + c_2 + c_3)}{3} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.6.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -\lambda & 1 & 1 \\ 1 & -\lambda & 1 \\ 1 & 1 & -\lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^3 - 3\lambda - 2 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 2$$

$$\lambda_2 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
2	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} 1 & 1 & 1 & 0 \\ 1 & 1 & 1 & 0 \\ 1 & 1 & 1 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{ccc|c} 1 & 1 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 0 \end{array} \right]$$

$$R_3 = R_3 - R_1 \implies \left[\begin{array}{ccc|c} 1 & 1 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & 1 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2, v_3\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Let $v_3 = s$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -t - s\}$

Hence the solution is

$$\begin{bmatrix} -t - s \\ t \\ s \end{bmatrix} = \begin{bmatrix} -t - s \\ t \\ s \end{bmatrix}$$

Since there are two free Variable, we have found two eigenvectors associated with this eigenvalue. The above can be written as

$$\begin{aligned} \begin{bmatrix} -t - s \\ t \\ s \end{bmatrix} &= \begin{bmatrix} -t \\ t \\ 0 \end{bmatrix} + \begin{bmatrix} -s \\ 0 \\ s \end{bmatrix} \\ &= t \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} + s \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} \end{aligned}$$

By letting $t = 1$ and $s = 1$ then the above becomes

$$\begin{bmatrix} -t - s \\ t \\ s \end{bmatrix} = \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} + \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix}$$

Hence the two eigenvectors associated with this eigenvalue are

$$\left(\begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix}, \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} \right)$$

Considering the eigenvalue $\lambda_2 = 2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix} - (2) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -2 & 1 & 1 \\ 1 & -2 & 1 \\ 1 & 1 & -2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -2 & 1 & 1 & 0 \\ 1 & -2 & 1 & 0 \\ 1 & 1 & -2 & 0 \end{array} \right]$$

$$R_2 = R_2 + \frac{R_1}{2} \implies \left[\begin{array}{ccc|c} -2 & 1 & 1 & 0 \\ 0 & -\frac{3}{2} & \frac{3}{2} & 0 \\ 1 & 1 & -2 & 0 \end{array} \right]$$

$$R_3 = R_3 + \frac{R_1}{2} \implies \left[\begin{array}{ccc|c} -2 & 1 & 1 & 0 \\ 0 & -\frac{3}{2} & \frac{3}{2} & 0 \\ 0 & \frac{3}{2} & -\frac{3}{2} & 0 \end{array} \right]$$

$$R_3 = R_3 + R_2 \implies \left[\begin{array}{ccc|c} -2 & 1 & 1 & 0 \\ 0 & -\frac{3}{2} & \frac{3}{2} & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -2 & 1 & 1 \\ 0 & -\frac{3}{2} & \frac{3}{2} \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t, v_2 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
2	1	1	No	$\begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$
-1	2	2	No	$\begin{bmatrix} -1 & -1 \\ 0 & 1 \\ 1 & 0 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue 2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^{2t} \\ &= \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} e^{2t} \end{aligned}$$

eigenvalue -1 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

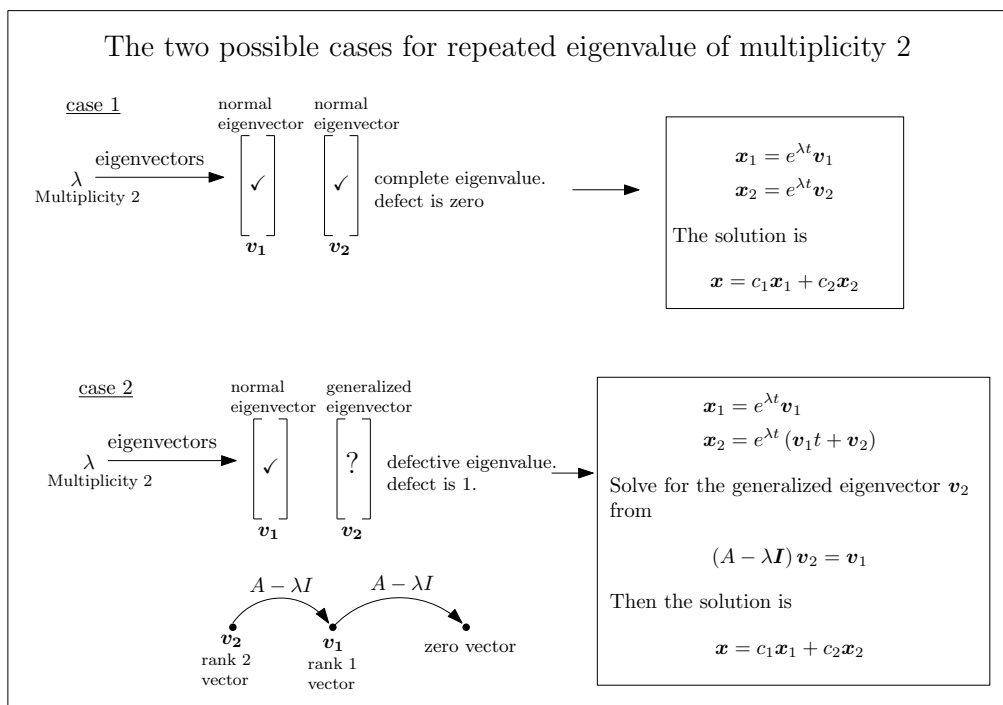


Figure 527: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2 which is the same as its geometric multiplicity 2, then it is complete eigenvalue and this falls into case 1 shown above. Hence the corresponding eigenvector basis are

$$\begin{aligned} \vec{x}_2(t) &= \vec{v}_2 e^{-t} \\ &= \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} e^{-t} \end{aligned}$$

$$\begin{aligned} \vec{x}_3(t) &= \vec{v}_3 e^{-t} \\ &= \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} e^{-t} \end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = c_1 \begin{bmatrix} e^{2t} \\ e^{2t} \\ e^{2t} \end{bmatrix} + c_2 \begin{bmatrix} -e^{-t} \\ 0 \\ e^{-t} \end{bmatrix} + c_3 \begin{bmatrix} -e^{-t} \\ e^{-t} \\ 0 \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} (-c_2 - c_3) e^{-t} + c_1 e^{2t} \\ c_1 e^{2t} + c_3 e^{-t} \\ c_1 e^{2t} + c_2 e^{-t} \end{bmatrix}$$

17.6.3 Maple step by step solution

Let's solve

$$[x'_1(t) = x_2(t) + x_3(t), x'_2(t) = x_1(t) + x_3(t), x'_3(t) = x_1(t) + x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-1, \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} \right], \left[-1, \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} \right], \left[2, \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair, with eigenvalue of algebraic multiplicity 2

$$\left[-1, \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} \right]$$

- First solution from eigenvalue -1

$$\underline{x}^{\rightarrow}_1(t) = e^{-t} \cdot \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix}$$

- Form of the 2nd homogeneous solution where \vec{p} is to be solved for, $\lambda = -1$ is the eigenvalue, and

$$\underline{x}^{\rightarrow}_2(t) = e^{\lambda t} (t\vec{v} + \vec{p})$$

- Note that the t multiplying \vec{v} makes this solution linearly independent to the 1st solution obtained

- Substitute $\underline{x}^{\rightarrow}_2(t)$ into the homogeneous system

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = (e^{\lambda t} A) \cdot (t\vec{v} + \vec{p})$$

- Use the fact that \vec{v} is an eigenvector of A

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = e^{\lambda t} (\lambda t\vec{v} + A \cdot \vec{p})$$

- Simplify equation

$$\lambda \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Make use of the identity matrix I

$$(\lambda \cdot I) \cdot \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Condition \vec{p} must meet for $x_{\underline{2}}(t)$ to be a solution to the homogeneous system

$$(A - \lambda \cdot I) \cdot \vec{p} = \vec{v}$$

- Choose \vec{p} to use in the second solution to the homogeneous system from eigenvalue -1

$$\left(\begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix} - (-1) \cdot \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \cdot \vec{p} = \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix}$$

- Choice of \vec{p}

$$\vec{p} = \begin{bmatrix} -1 \\ 0 \\ 0 \end{bmatrix}$$

- Second solution from eigenvalue -1

$$x_{\underline{2}}(t) = e^{-t} \cdot \left(t \cdot \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} + \begin{bmatrix} -1 \\ 0 \\ 0 \end{bmatrix} \right)$$

- Consider eigenpair

$$\left[2, \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$x_{\underline{3}} = e^{2t} \cdot \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

- General solution to the system of ODEs

$$x_{\underline{}} = c_1 x_{\underline{1}}(t) + c_2 x_{\underline{2}}(t) + c_3 x_{\underline{3}}$$

- Substitute solutions into the general solution

$$\underline{x} \rightarrow = c_1 e^{-t} \cdot \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} + c_2 e^{-t} \cdot \left(t \cdot \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} + \begin{bmatrix} -1 \\ 0 \\ 0 \end{bmatrix} \right) + c_3 e^{2t} \cdot \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} ((-t-1)c_2 - c_1)e^{-t} + c_3 e^{2t} \\ c_3 e^{2t} \\ (c_2 t + c_1)e^{-t} + c_3 e^{2t} \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = ((-t-1)c_2 - c_1)e^{-t} + c_3 e^{2t}, x_2(t) = c_3 e^{2t}, x_3(t) = (c_2 t + c_1)e^{-t} + c_3 e^{2t}\}$$

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 64

```
dsolve([diff(x__1(t),t)=0*x__1(t)+1*x__2(t)+1*x__3(t),diff(x__2(t),t)=1*x__1(t)+0*x__2(t)+1*
```

$$\begin{aligned} x_1(t) &= c_2 e^{-t} + c_3 e^{2t} \\ x_2(t) &= c_2 e^{-t} + c_3 e^{2t} + e^{-t} c_1 \\ x_3(t) &= -2c_2 e^{-t} + c_3 e^{2t} - e^{-t} c_1 \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 124

```
DSolve[{x1'[t]==0*x1[t]+1*x2[t]+1*x3[t],x2'[t]==1*x1[t]+0*x2[t]+1*x3[t],x3'[t]==1*x1[t]+1*x2[t]
```

$$\begin{aligned} x1(t) &\rightarrow \frac{1}{3}e^{-t}(c_1(e^{3t}+2) + (c_2+c_3)(e^{3t}-1)) \\ x2(t) &\rightarrow \frac{1}{3}e^{-t}(c_1(e^{3t}-1) + c_2(e^{3t}+2) + c_3(e^{3t}-1)) \\ x3(t) &\rightarrow \frac{1}{3}e^{-t}(c_1(e^{3t}-1) + c_2(e^{3t}-1) + c_3(e^{3t}+2)) \end{aligned}$$

17.7 problem 7

17.7.1 Solution using Matrix exponential method 3735

17.7.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3736

Internal problem ID [772]

Internal file name [OUTPUT/772_Sunday_June_05_2022_01_49_17_AM_55713821/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 7.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) - 4x_2(t) \\x_2'(t) &= 4x_1(t) - 7x_2(t)\end{aligned}$$

With initial conditions

$$[x_1(0) = 3, x_2(0) = 2]$$

17.7.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-3t}(4t + 1) & -4t e^{-3t} \\ 4t e^{-3t} & e^{-3t}(1 - 4t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{x}_0 \\
 &= \begin{bmatrix} e^{-3t}(4t+1) & -4t e^{-3t} \\ 4t e^{-3t} & e^{-3t}(1-4t) \end{bmatrix} \begin{bmatrix} 3 \\ 2 \end{bmatrix} \\
 &= \begin{bmatrix} 3e^{-3t}(4t+1) - 8t e^{-3t} \\ 12t e^{-3t} + 2e^{-3t}(1-4t) \end{bmatrix} \\
 &= \begin{bmatrix} e^{-3t}(4t+3) \\ (4t+2)e^{-3t} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.7.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 1 - \lambda & -4 \\ 4 & -7 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 6\lambda + 9 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -3$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-3	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -3$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} - (-3) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} 4 & -4 \\ 4 & -4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 4 & -4 & 0 \\ 4 & -4 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} 4 & -4 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 4 & -4 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-3	2	1	Yes	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue -3 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

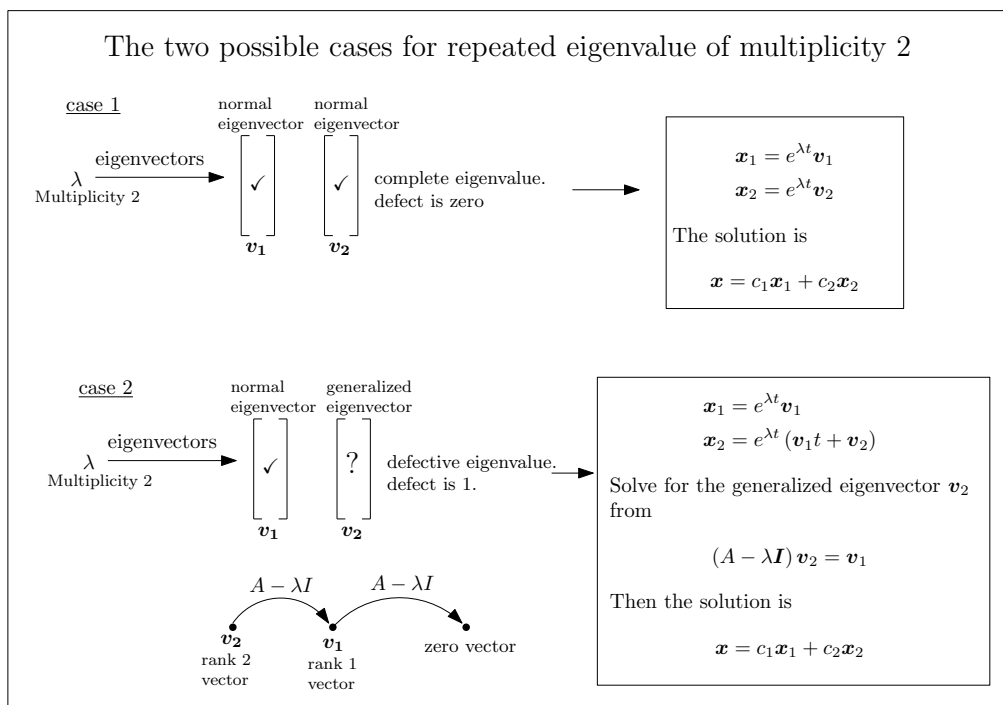


Figure 528: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} - (-3) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} 4 & -4 \\ 4 & -4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} \frac{5}{4} \\ 1 \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue -3 . Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{-3t} \\ &= \begin{bmatrix} e^{-3t} \\ e^{-3t} \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} 1 \\ 1 \end{bmatrix} t + \begin{bmatrix} \frac{5}{4} \\ 1 \end{bmatrix} \right) e^{-3t} \\ &= \begin{bmatrix} \frac{e^{-3t}(4t+5)}{4} \\ e^{-3t}(t+1) \end{bmatrix}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} e^{-3t} \\ e^{-3t} \end{bmatrix} + c_2 \begin{bmatrix} e^{-3t}(t + \frac{5}{4}) \\ e^{-3t}(t + 1) \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^{-3t}(c_1 + c_2 t + \frac{5}{4}c_2) \\ e^{-3t}(c_2 t + c_1 + c_2) \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x_1(0) = 3 \\ x_2(0) = 2 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 3 \\ 2 \end{bmatrix} = \begin{bmatrix} c_1 + \frac{5c_2}{4} \\ c_1 + c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = -2 \\ c_2 = 4 \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^{-3t}(4t + 3) \\ (4t + 2)e^{-3t} \end{bmatrix}$$

The following is the phase plot of the system.

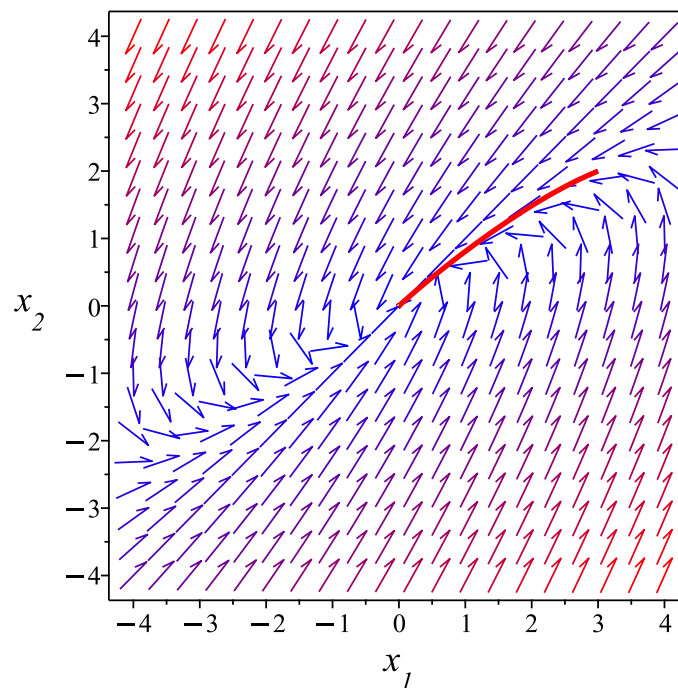
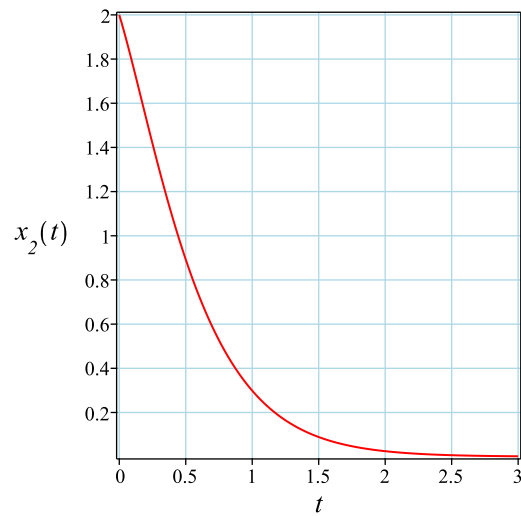
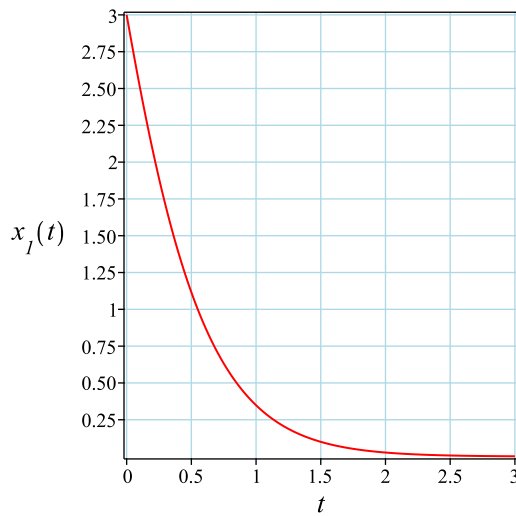


Figure 529: Phase plot

The following are plots of each solution.



✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 29

```
dsolve([diff(x__1(t),t) = x__1(t)-4*x__2(t), diff(x__2(t),t) = 4*x__1(t)-7*x__2(t), x__1(0)
```

$$x_1(t) = e^{-3t}(4t + 3)$$

$$x_2(t) = \frac{e^{-3t}(16t + 8)}{4}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 34

```
DSolve[{x1'[t]==1*x1[t]-4*x2[t],x2'[t]==1*x1[t]-4*x2[t]},{x1[0]==3,x2[0]==2},{x1[t],x2[t]},t
```

$$x1(t) \rightarrow \frac{5e^{-3t}}{3} + \frac{4}{3}$$

$$x2(t) \rightarrow \frac{5e^{-3t}}{3} + \frac{1}{3}$$

17.8 problem 8

17.8.1 Solution using Matrix exponential method 3743

17.8.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3744

Internal problem ID [773]

Internal file name [OUTPUT/773_Sunday_June_05_2022_01_49_18_AM_81518641/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 8.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -\frac{5x_1(t)}{2} + \frac{3x_2(t)}{2} \\x_2'(t) &= -\frac{3x_1(t)}{2} + \frac{x_2(t)}{2}\end{aligned}$$

With initial conditions

$$[x_1(0) = 3, x_2(0) = -1]$$

17.8.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{5}{2} & \frac{3}{2} \\ -\frac{3}{2} & \frac{1}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-t}\left(1 - \frac{3t}{2}\right) & \frac{3te^{-t}}{2} \\ -\frac{3te^{-t}}{2} & e^{-t}\left(1 + \frac{3t}{2}\right) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{x}_0 \\
 &= \begin{bmatrix} e^{-t} \left(1 - \frac{3t}{2}\right) & \frac{3te^{-t}}{2} \\ -\frac{3te^{-t}}{2} & e^{-t} \left(1 + \frac{3t}{2}\right) \end{bmatrix} \begin{bmatrix} 3 \\ -1 \end{bmatrix} \\
 &= \begin{bmatrix} 3e^{-t} \left(1 - \frac{3t}{2}\right) - \frac{3te^{-t}}{2} \\ -\frac{9te^{-t}}{2} - e^{-t} \left(1 + \frac{3t}{2}\right) \end{bmatrix} \\
 &= \begin{bmatrix} (-6t + 3)e^{-t} \\ (-6t - 1)e^{-t} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.8.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{5}{2} & \frac{3}{2} \\ -\frac{3}{2} & \frac{1}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} -\frac{5}{2} & \frac{3}{2} \\ -\frac{3}{2} & \frac{1}{2} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -\frac{5}{2} - \lambda & \frac{3}{2} \\ -\frac{3}{2} & \frac{1}{2} - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 2\lambda + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{5}{2} & \frac{3}{2} \\ -\frac{3}{2} & \frac{1}{2} \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -\frac{3}{2} & \frac{3}{2} \\ -\frac{3}{2} & \frac{3}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -\frac{3}{2} & \frac{3}{2} & 0 \\ -\frac{3}{2} & \frac{3}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} -\frac{3}{2} & \frac{3}{2} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -\frac{3}{2} & \frac{3}{2} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-1	2	1	Yes	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue -1 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

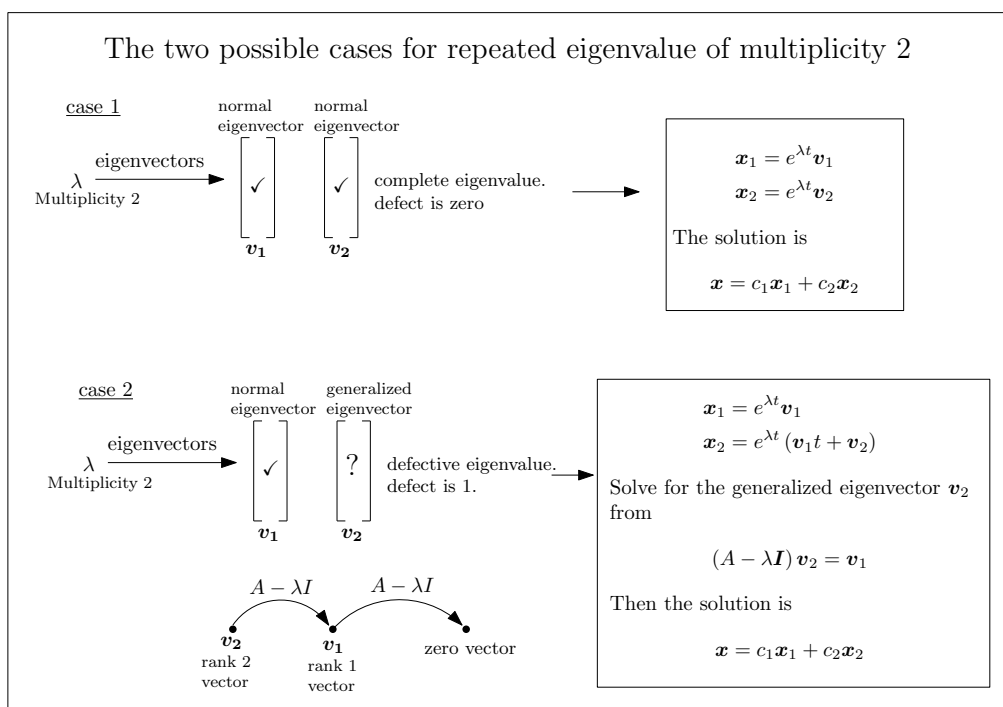


Figure 530: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} -\frac{5}{2} & \frac{3}{2} \\ -\frac{3}{2} & \frac{1}{2} \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} -\frac{3}{2} & \frac{3}{2} \\ -\frac{3}{2} & \frac{3}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue -1 . Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{-t} \\ &= \begin{bmatrix} e^{-t} \\ e^{-t} \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} 1 \\ 1 \end{bmatrix} t + \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} \right) e^{-t} \\ &= \begin{bmatrix} \frac{(1+3t)e^{-t}}{3} \\ e^{-t}(t+1) \end{bmatrix}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} e^{-t} \\ e^{-t} \end{bmatrix} + c_2 \begin{bmatrix} e^{-t}(t + \frac{1}{3}) \\ e^{-t}(t + 1) \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^{-t}(c_1 + c_2 t + \frac{1}{3}c_2) \\ e^{-t}(c_2 t + c_1 + c_2) \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x_1(0) = 3 \\ x_2(0) = -1 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 3 \\ -1 \end{bmatrix} = \begin{bmatrix} c_1 + \frac{c_2}{3} \\ c_1 + c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = 5 \\ c_2 = -6 \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (-6t + 3) e^{-t} \\ (-6t - 1) e^{-t} \end{bmatrix}$$

The following is the phase plot of the system.

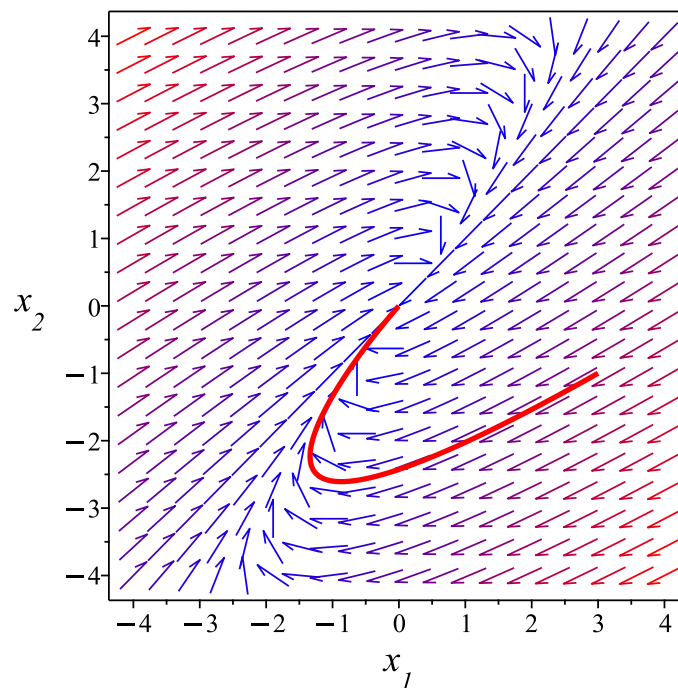
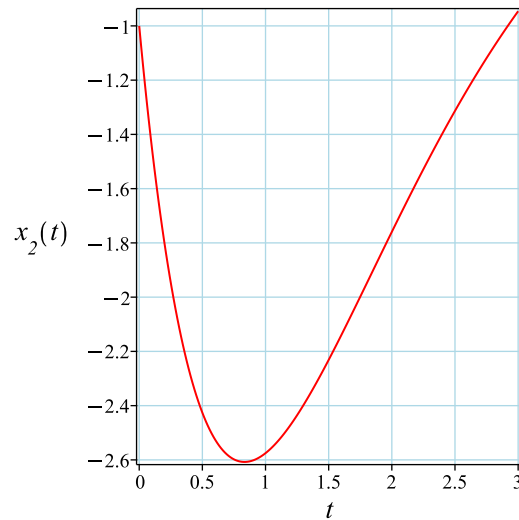
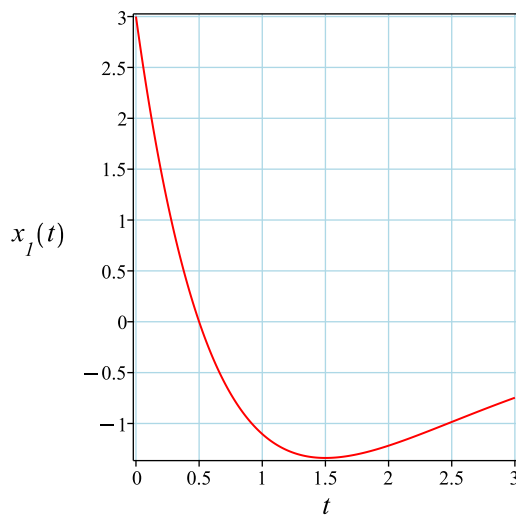


Figure 531: Phase plot

The following are plots of each solution.



✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 29

```
dsolve([diff(x__1(t),t) = -5/2*x__1(t)+3/2*x__2(t), diff(x__2(t),t) = -3/2*x__1(t)+1/2*x__2(t)
```

$$x_1(t) = e^{-t}(-6t + 3)$$

$$x_2(t) = \frac{e^{-t}(-18t - 3)}{3}$$

✓ Solution by Mathematica

Time used: 0.003 (sec). Leaf size: 31

```
DSolve[{x1'[t]==-5/2*x1[t]+3/2*x2[t],x2'[t]==-3/2*x1[t]+1/2*x2[t]},{x1[0]==3,x2[0]==-1},{x1[t]
```

$$x1(t) \rightarrow e^{-t}(3 - 6t)$$

$$x2(t) \rightarrow -e^{-t}(6t + 1)$$

17.9 problem 9

17.9.1 Solution using Matrix exponential method 3751

17.9.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3752

Internal problem ID [774]

Internal file name [OUTPUT/774_Sunday_June_05_2022_01_49_19_AM_6565187/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 9.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 2x_1(t) + \frac{3x_2(t)}{2} \\x_2'(t) &= -\frac{3x_1(t)}{2} - x_2(t)\end{aligned}$$

With initial conditions

$$[x_1(0) = 3, x_2(0) = -2]$$

17.9.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & \frac{3}{2} \\ -\frac{3}{2} & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{\frac{t}{2}} \left(1 + \frac{3t}{2}\right) & \frac{3te^{\frac{t}{2}}}{2} \\ -\frac{3te^{\frac{t}{2}}}{2} & e^{\frac{t}{2}} \left(1 - \frac{3t}{2}\right) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{x}_0 \\
 &= \begin{bmatrix} e^{\frac{t}{2}} \left(1 + \frac{3t}{2}\right) & \frac{3te^{\frac{t}{2}}}{2} \\ -\frac{3te^{\frac{t}{2}}}{2} & e^{\frac{t}{2}} \left(1 - \frac{3t}{2}\right) \end{bmatrix} \begin{bmatrix} 3 \\ -2 \end{bmatrix} \\
 &= \begin{bmatrix} 3e^{\frac{t}{2}} \left(1 + \frac{3t}{2}\right) - 3te^{\frac{t}{2}} \\ -\frac{9te^{\frac{t}{2}}}{2} - 2e^{\frac{t}{2}} \left(1 - \frac{3t}{2}\right) \end{bmatrix} \\
 &= \begin{bmatrix} \frac{3e^{\frac{t}{2}}(2+t)}{2} \\ e^{\frac{t}{2}} \left(-2 - \frac{3t}{2}\right) \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.9.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & \frac{3}{2} \\ -\frac{3}{2} & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 2 & \frac{3}{2} \\ -\frac{3}{2} & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 2 - \lambda & \frac{3}{2} \\ -\frac{3}{2} & -1 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - \lambda + \frac{1}{4} = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = \frac{1}{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$\frac{1}{2}$	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = \frac{1}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & \frac{3}{2} \\ -\frac{3}{2} & -1 \end{bmatrix} - \left(\frac{1}{2}\right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} \frac{3}{2} & \frac{3}{2} \\ -\frac{3}{2} & -\frac{3}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{3}{2} & \frac{3}{2} & 0 \\ -\frac{3}{2} & -\frac{3}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 + R_1 \implies \left[\begin{array}{cc|c} \frac{3}{2} & \frac{3}{2} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} \frac{3}{2} & \frac{3}{2} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -t\}$

Hence the solution is

$$\begin{bmatrix} -t \\ t \end{bmatrix} = \begin{bmatrix} -t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -t \\ t \end{bmatrix} = t \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$\frac{1}{2}$	2	1	Yes	$\begin{bmatrix} -1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue $\frac{1}{2}$ is real and repated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

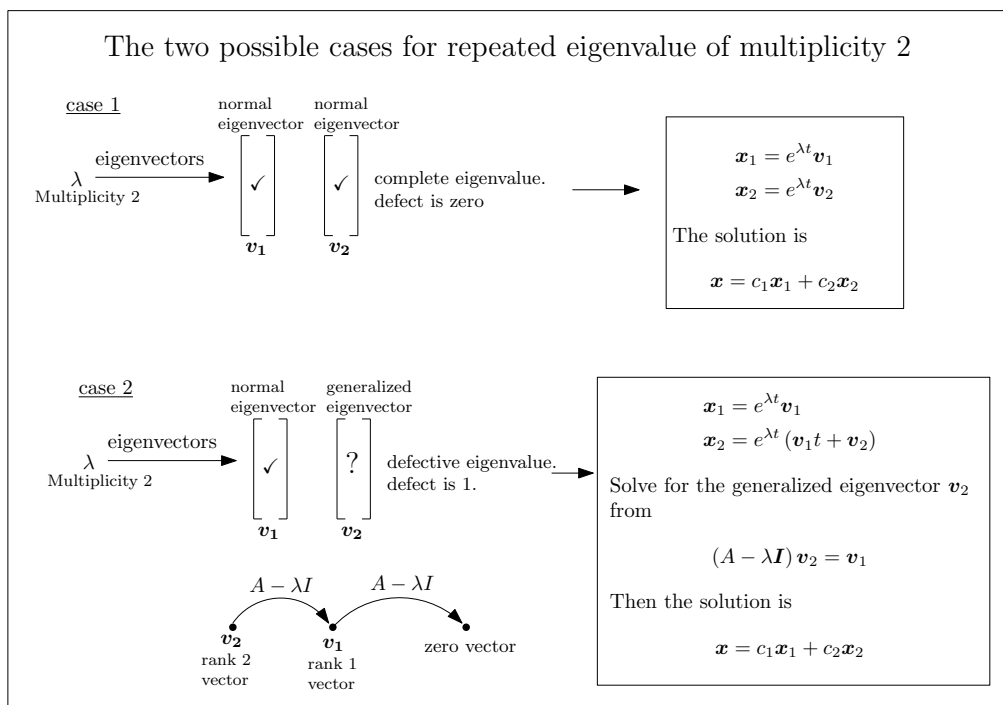


Figure 532: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} 2 & \frac{3}{2} \\ -\frac{3}{2} & -1 \end{bmatrix} - \left(\frac{1}{2} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} \frac{3}{2} & \frac{3}{2} \\ -\frac{3}{2} & -\frac{3}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} -\frac{5}{3} \\ 1 \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue $\frac{1}{2}$. Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} -1 \\ 1 \end{bmatrix} e^{\frac{t}{2}} \\ &= \begin{bmatrix} -e^{\frac{t}{2}} \\ e^{\frac{t}{2}} \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} -1 \\ 1 \end{bmatrix} t + \begin{bmatrix} -\frac{5}{3} \\ 1 \end{bmatrix} \right) e^{\frac{t}{2}} \\ &= \begin{bmatrix} -\frac{e^{\frac{t}{2}}(3t+5)}{3} \\ e^{\frac{t}{2}}(t+1) \end{bmatrix}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} -e^{\frac{t}{2}} \\ e^{\frac{t}{2}} \end{bmatrix} + c_2 \begin{bmatrix} e^{\frac{t}{2}}(-t - \frac{5}{3}) \\ e^{\frac{t}{2}}(t+1) \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^{\frac{t}{2}}(-c_1 - c_2 t - \frac{5}{3}c_2) \\ e^{\frac{t}{2}}(c_2 t + c_1 + c_2) \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x_1(0) = 3 \\ x_2(0) = -2 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 3 \\ -2 \end{bmatrix} = \begin{bmatrix} -c_1 - \frac{5c_2}{3} \\ c_1 + c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = -\frac{1}{2} \\ c_2 = -\frac{3}{2} \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^{\frac{t}{2}} \left(3 + \frac{3t}{2} \right) \\ e^{\frac{t}{2}} \left(-2 - \frac{3t}{2} \right) \end{bmatrix}$$

The following is the phase plot of the system.

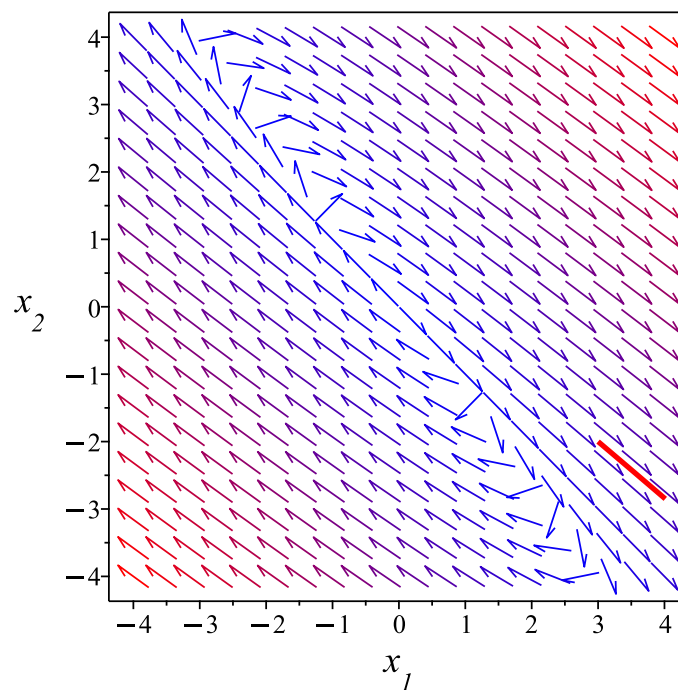
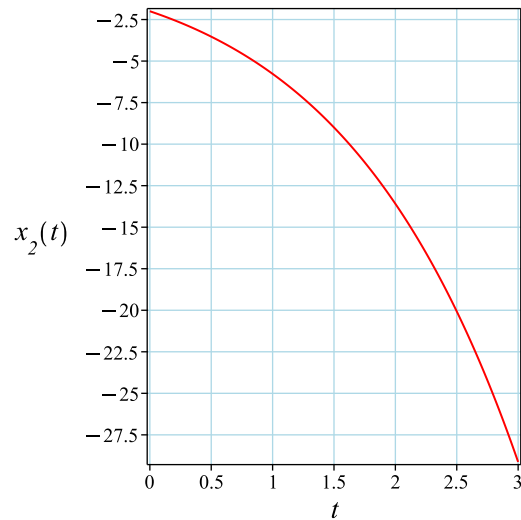
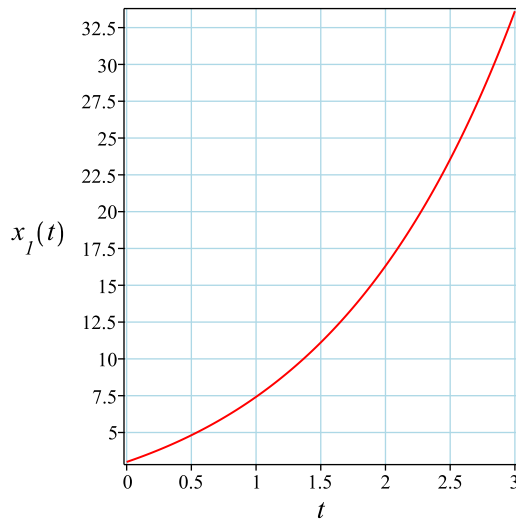


Figure 533: Phase plot

The following are plots of each solution.



✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 29

```
dsolve([diff(x__1(t),t) = 2*x__1(t)+3/2*x__2(t), diff(x__2(t),t) = -3/2*x__1(t)-x__2(t), x__
```

$$x_1(t) = e^{\frac{t}{2}} \left(\frac{3t}{2} + 3 \right)$$

$$x_2(t) = -\frac{e^{\frac{t}{2}} \left(\frac{9t}{2} + 6 \right)}{3}$$

✓ Solution by Mathematica

Time used: 0.003 (sec). Leaf size: 38

```
DSolve[{x1'[t]==2*x1[t]+3/2*x2[t],x2'[t]==-3/2*x1[t]-1*x2[t]},{x1[0]==3,x2[0]==-2},{x1[t],x2
```

$$x_1(t) \rightarrow \frac{3}{2}e^{t/2}(t+2)$$

$$x_2(t) \rightarrow -\frac{1}{2}e^{t/2}(3t+4)$$

17.10 problem 10

17.10.1 Solution using Matrix exponential method 3759

17.10.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3760

Internal problem ID [775]

Internal file name [OUTPUT/775_Sunday_June_05_2022_01_49_21_AM_24868463/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 10.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x_1'(t) = 3x_1(t) + 9x_2(t)$$

$$x_2'(t) = -x_1(t) - 3x_2(t)$$

With initial conditions

$$[x_1(0) = 2, x_2(0) = 4]$$

17.10.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & 9 \\ -1 & -3 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} 1 + 3t & 9t \\ -t & 1 - 3t \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}\vec{x}_h(t) &= e^{At}\vec{x}_0 \\ &= \begin{bmatrix} 1 + 3t & 9t \\ -t & 1 - 3t \end{bmatrix} \begin{bmatrix} 2 \\ 4 \end{bmatrix} \\ &= \begin{bmatrix} 2 + 42t \\ -14t + 4 \end{bmatrix}\end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.10.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & 9 \\ -1 & -3 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 3 & 9 \\ -1 & -3 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 3 - \lambda & 9 \\ -1 & -3 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 0$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
0	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = 0$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 3 & 9 \\ -1 & -3 \end{bmatrix} - (0) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 3 & 9 \\ -1 & -3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 3 & 9 & 0 \\ -1 & -3 & 0 \end{array} \right]$$

$$R_2 = R_2 + \frac{R_1}{3} \implies \left[\begin{array}{cc|c} 3 & 9 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 3 & 9 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -3t\}$

Hence the solution is

$$\begin{bmatrix} -3t \\ t \end{bmatrix} = \begin{bmatrix} -3t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -3t \\ t \end{bmatrix} = t \begin{bmatrix} -3 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -3t \\ t \end{bmatrix} = \begin{bmatrix} -3 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
0	2	1	Yes	$\begin{bmatrix} -3 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue 0 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

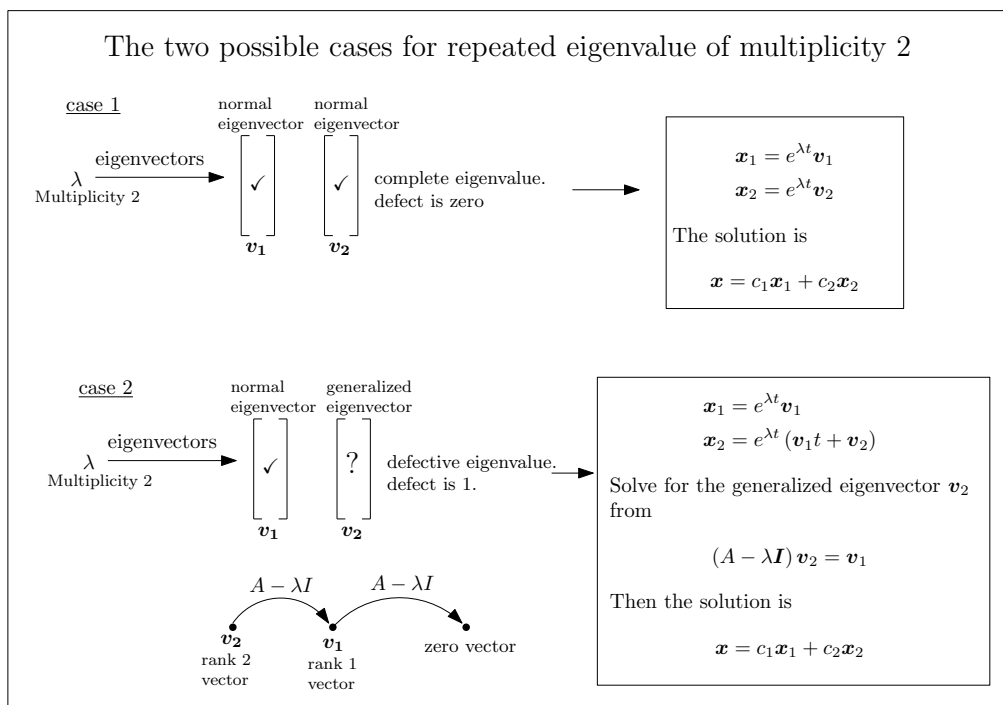


Figure 534: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} 3 & 9 \\ -1 & -3 \end{bmatrix} - (0) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} -3 \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} 3 & 9 \\ -1 & -3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} -3 \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} -4 \\ 1 \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue 0. Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} -3 \\ 1 \end{bmatrix} e^{0t} \\ &= \begin{bmatrix} -3 \\ 1 \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} -3 \\ 1 \end{bmatrix} t + \begin{bmatrix} -4 \\ 1 \end{bmatrix} \right) e^{0t} \\ &= \begin{bmatrix} -4 - 3t \\ t + 1 \end{bmatrix}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} -3 \\ 1 \end{bmatrix} + c_2 \begin{bmatrix} -4 - 3t \\ t + 1 \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} -3c_1 + c_2(-4 - 3t) \\ c_2 t + c_1 + c_2 \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x_1(0) = 2 \\ x_2(0) = 4 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 2 \\ 4 \end{bmatrix} = \begin{bmatrix} -3c_1 - 4c_2 \\ c_1 + c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = 18 \\ c_2 = -14 \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} 2 + 42t \\ -14t + 4 \end{bmatrix}$$

The following is the phase plot of the system.

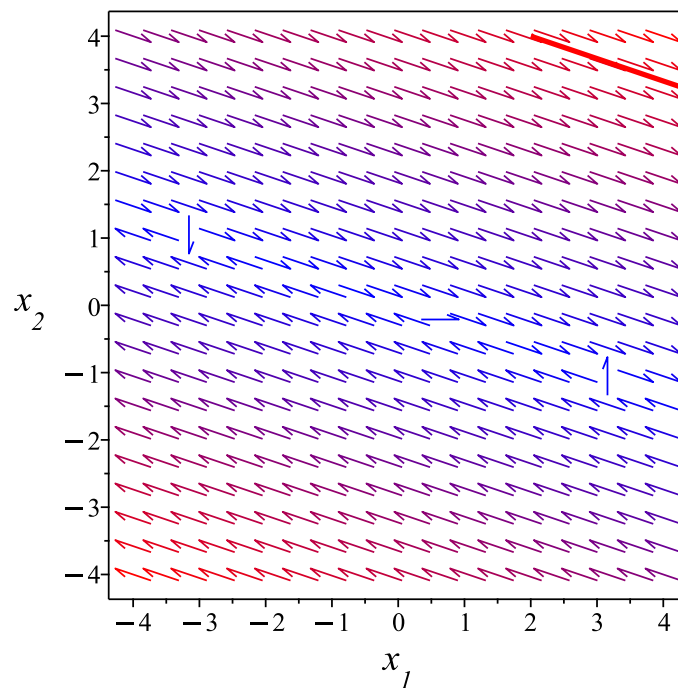
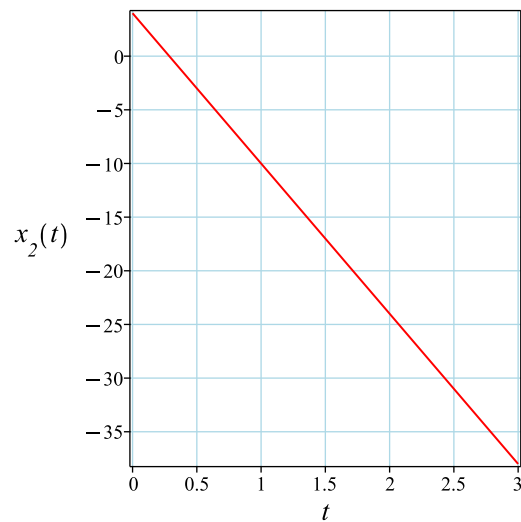
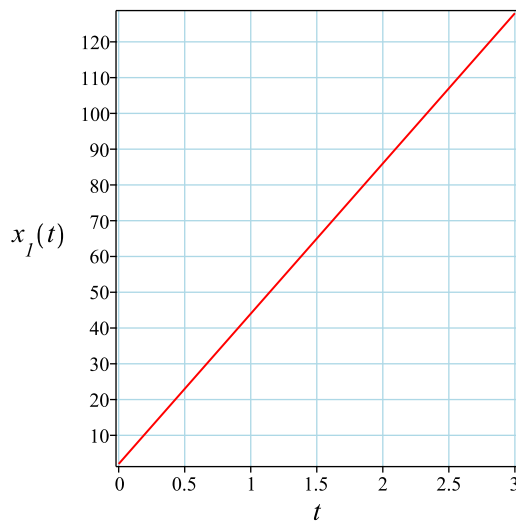


Figure 535: Phase plot

The following are plots of each solution.



✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 18

```
dsolve([diff(x__1(t),t) = 3*x__1(t)+9*x__2(t), diff(x__2(t),t) = -x__1(t)-3*x__2(t), x__1(0)
```

$$\begin{aligned}x_1(t) &= 42t + 2 \\x_2(t) &= 4 - 14t\end{aligned}$$

✓ Solution by Mathematica

Time used: 0.003 (sec). Leaf size: 18

```
DSolve[{x1'[t]==3*x1[t]+9*x2[t],x2'[t]==-1*x1[t]-3*x2[t]},{x1[0]==2,x2[0]==4},{x1[t],x2[t]},
```

$$\begin{aligned}x1(t) &\rightarrow 42t + 2 \\x2(t) &\rightarrow 4 - 14t\end{aligned}$$

17.11 problem 11

17.11.1 Solution using Matrix exponential method 3767

17.11.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3768

Internal problem ID [776]

Internal file name [OUTPUT/776_Sunday_June_05_2022_01_49_22_AM_63754416/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 11.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x_1'(t) = x_1(t)$$

$$x_2'(t) = -4x_1(t) + x_2(t)$$

$$x_3'(t) = 3x_1(t) + 6x_2(t) + 2x_3(t)$$

With initial conditions

$$[x_1(0) = -1, x_2(0) = 2, x_3(0) = -30]$$

17.11.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ -4 & 1 & 0 \\ 3 & 6 & 2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^t & 0 & 0 \\ -4t e^t & e^t & 0 \\ -21 e^{2t} + (24t + 21) e^t & 6 e^{2t} - 6 e^t & e^{2t} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{x}_0 \\ &= \begin{bmatrix} e^t & 0 & 0 \\ -4t e^t & e^t & 0 \\ -21 e^{2t} + (24t + 21) e^t & 6 e^{2t} - 6 e^t & e^{2t} \end{bmatrix} \begin{bmatrix} -1 \\ 2 \\ -30 \end{bmatrix} \\ &= \begin{bmatrix} -e^t \\ 4t e^t + 2 e^t \\ 3 e^{2t} - (24t + 21) e^t - 12 e^t \end{bmatrix} \\ &= \begin{bmatrix} -e^t \\ (4t + 2) e^t \\ 3 e^{2t} + (-24t - 33) e^t \end{bmatrix} \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.11.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ -4 & 1 & 0 \\ 3 & 6 & 2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 1 & 0 & 0 \\ -4 & 1 & 0 \\ 3 & 6 & 2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 1 - \lambda & 0 & 0 \\ -4 & 1 - \lambda & 0 \\ 3 & 6 & 2 - \lambda \end{bmatrix} \right) = 0$$

Since the matrix A is triangular matrix, then the determinant is the product of the elements along the diagonal. Therefore the above becomes

$$(1 - \lambda)(1 - \lambda)(2 - \lambda) = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 2$$

$$\lambda_2 = 1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
1	1	real eigenvalue
2	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = 1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 0 & 0 \\ -4 & 1 & 0 \\ 3 & 6 & 2 \end{bmatrix} - (1) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 & 0 & 0 \\ -4 & 0 & 0 \\ 3 & 6 & 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} 0 & 0 & 0 & 0 \\ -4 & 0 & 0 & 0 \\ 3 & 6 & 1 & 0 \end{array} \right]$$

Since the current pivot $A(1,1)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 1 and row 2 gives

$$\left[\begin{array}{ccc|c} -4 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 3 & 6 & 1 & 0 \end{array} \right]$$

$$R_3 = R_3 + \frac{3R_1}{4} \implies \left[\begin{array}{ccc|c} -4 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 6 & 1 & 0 \end{array} \right]$$

Since the current pivot $A(2,2)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 2 and row 3 gives

$$\left[\begin{array}{ccc|c} -4 & 0 & 0 & 0 \\ 0 & 6 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -4 & 0 & 0 \\ 0 & 6 & 1 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 0, v_2 = -\frac{t}{6}\}$

Hence the solution is

$$\begin{bmatrix} 0 \\ -\frac{t}{6} \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ -\frac{t}{6} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 0 \\ -\frac{t}{6} \\ t \end{bmatrix} = t \begin{bmatrix} 0 \\ -\frac{1}{6} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 0 \\ -\frac{1}{6} \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ -\frac{1}{6} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} 0 \\ -\frac{1}{6} \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ -1 \\ 6 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 0 & 0 \\ -4 & 1 & 0 \\ 3 & 6 & 2 \end{bmatrix} - (2) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 & 0 & 0 \\ -4 & -1 & 0 \\ 3 & 6 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -1 & 0 & 0 & 0 \\ -4 & -1 & 0 & 0 \\ 3 & 6 & 0 & 0 \end{array} \right]$$

$$R_2 = R_2 - 4R_1 \implies \left[\begin{array}{ccc|c} -1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 3 & 6 & 0 & 0 \end{array} \right]$$

$$R_3 = R_3 + 3R_1 \implies \left[\begin{array}{ccc|c} -1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 6 & 0 & 0 \end{array} \right]$$

$$R_3 = R_3 + 6R_2 \implies \left[\begin{array}{ccc|c} -1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 0, v_2 = 0\}$

Hence the solution is

$$\begin{bmatrix} 0 \\ 0 \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 0 \\ 0 \\ t \end{bmatrix} = t \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 0 \\ 0 \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated

with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
2	1	1	No	$\begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$
1	2	1	Yes	$\begin{bmatrix} 0 \\ -\frac{1}{6} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue 2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^{2t} \\ &= \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} e^{2t} \end{aligned}$$

eigenvalue 1 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

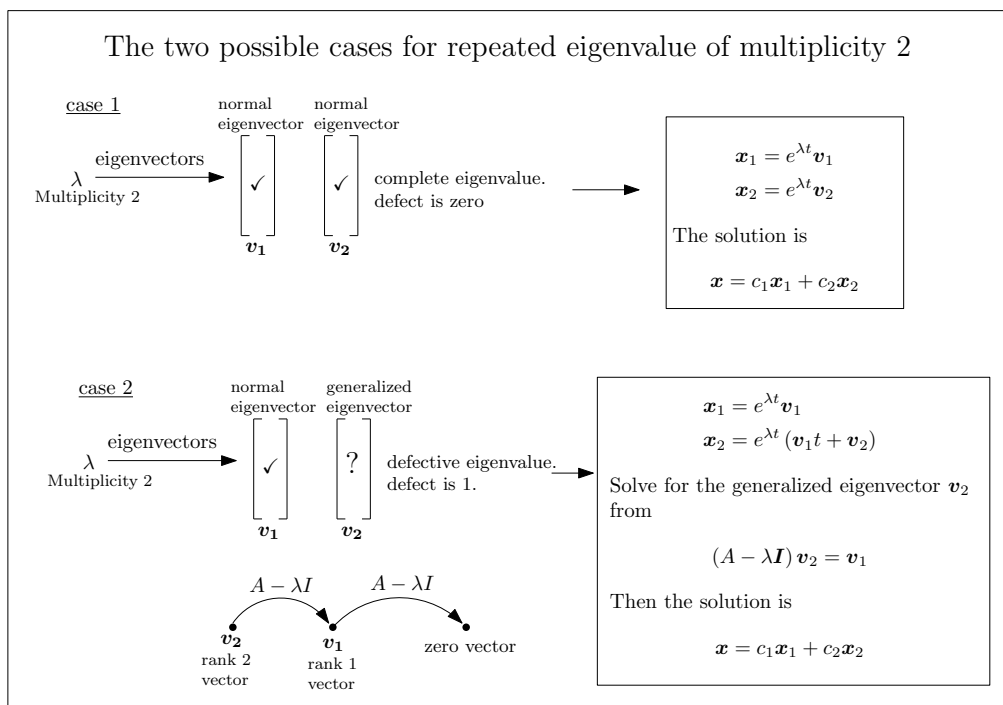


Figure 536: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} 1 & 0 & 0 \\ -4 & 1 & 0 \\ 3 & 6 & 2 \end{bmatrix} - (1) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ -\frac{1}{6} \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} 0 & 0 & 0 \\ -4 & 0 & 0 \\ 3 & 6 & 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ -\frac{1}{6} \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} \frac{1}{24} \\ 1 \\ -\frac{41}{8} \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue 1. Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned} \vec{x}_2(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} 0 \\ -\frac{1}{6} \\ 1 \end{bmatrix} e^t \\ &= \begin{bmatrix} 0 \\ -\frac{e^t}{6} \\ e^t \end{bmatrix} \end{aligned}$$

And

$$\begin{aligned} \vec{x}_3(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} 0 \\ -\frac{1}{6} \\ 1 \end{bmatrix} t + \begin{bmatrix} \frac{1}{24} \\ 1 \\ -\frac{41}{8} \end{bmatrix} \right) e^t \\ &= \begin{bmatrix} \frac{e^t}{24} \\ -\frac{e^t(t-6)}{6} \\ \frac{e^t(8t-41)}{8} \end{bmatrix} \end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = c_1 \begin{bmatrix} 0 \\ 0 \\ e^{2t} \end{bmatrix} + c_2 \begin{bmatrix} 0 \\ -\frac{e^t}{6} \\ e^t \end{bmatrix} + c_3 \begin{bmatrix} \frac{e^t}{24} \\ e^t \left(-\frac{t}{6} + 1\right) \\ e^t \left(t - \frac{41}{8}\right) \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} \frac{c_3 e^t}{24} \\ -\frac{((t-6)c_3 + c_2)e^t}{6} \\ c_1 e^{2t} + \left(t - \frac{41}{8}\right) c_3 + c_2 e^t \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x_1(0) = -1 \\ x_2(0) = 2 \\ x_3(0) = -30 \end{bmatrix} \quad (1)$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} -1 \\ 2 \\ -30 \end{bmatrix} = \begin{bmatrix} \frac{c_3}{24} \\ c_3 - \frac{c_2}{6} \\ c_1 - \frac{41c_3}{8} + c_2 \end{bmatrix}$$

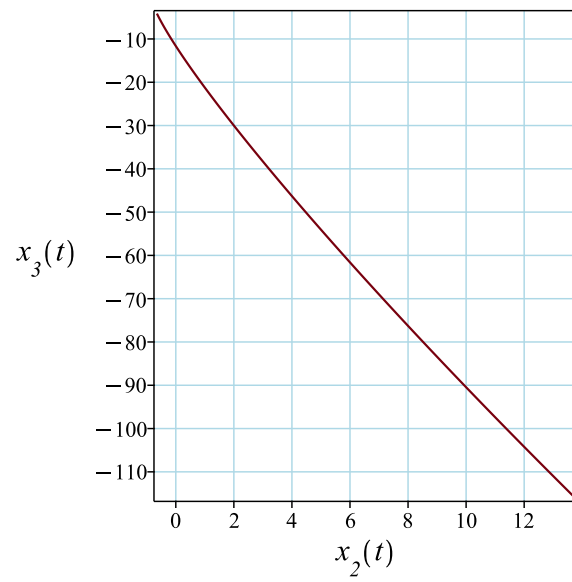
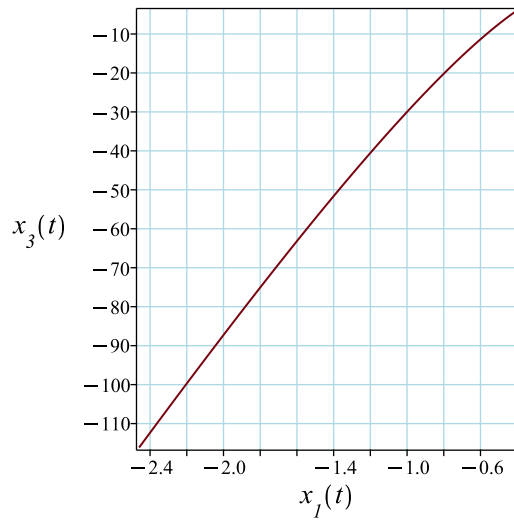
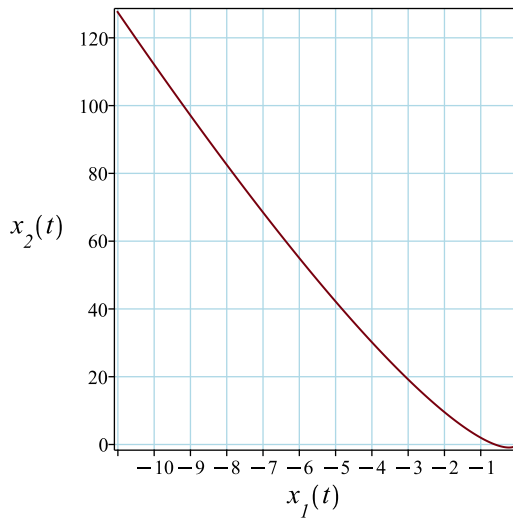
Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = 3 \\ c_2 = -156 \\ c_3 = -24 \end{bmatrix}$$

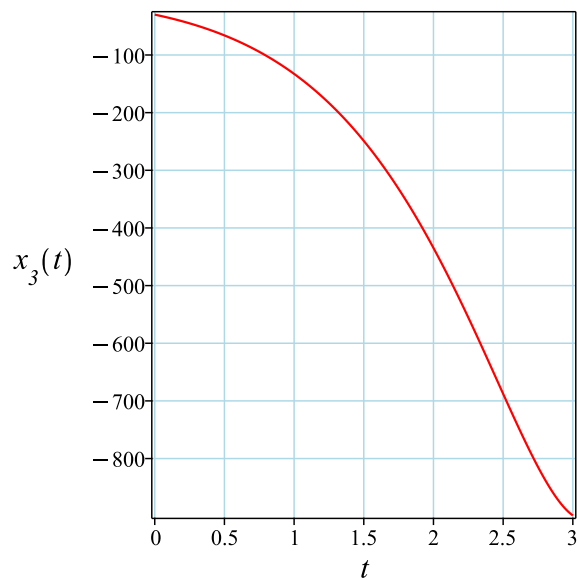
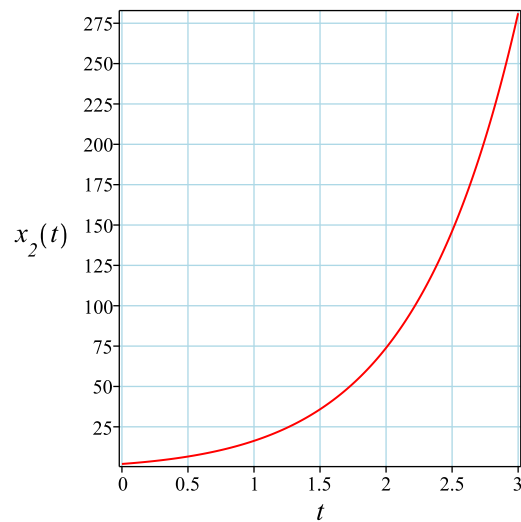
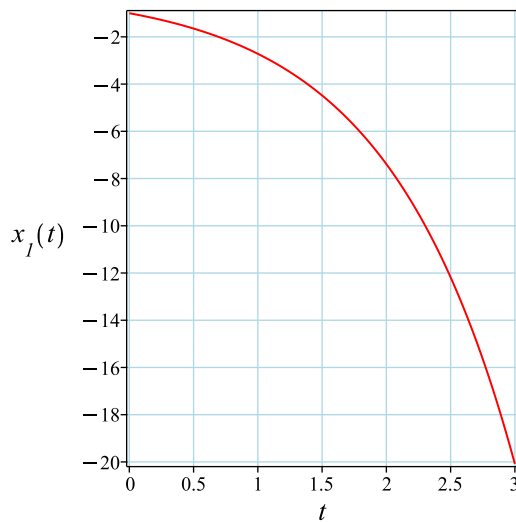
Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} -e^t \\ -\frac{(-24t-12)e^t}{6} \\ 3e^{2t} + (-24t - 33)e^t \end{bmatrix}$$

The following are plots of each solution against another.



The following are plots of each solution.



✓ Solution by Maple

Time used: 0.032 (sec). Leaf size: 39

```
dsolve([diff(x__1(t),t) = x__1(t), diff(x__2(t),t) = -4*x__1(t)+x__2(t), diff(x__3(t),t) = 3
```

$$x_1(t) = -e^t$$

$$x_2(t) = (4t + 2)e^t$$

$$x_3(t) = -24e^t t - 33e^t + 3e^{2t}$$

✓ Solution by Mathematica

Time used: 0.007 (sec). Leaf size: 39

```
DSolve[{x1'[t]==1*x1[t]+0*x2[t]+0*x3[t],x2'[t]==-4*x1[t]+1*x2[t]+0*x3[t],x3'[t]==3*x1[t]+6*x
```

$$x1(t) \rightarrow -e^t$$

$$x2(t) \rightarrow 2e^t(2t + 1)$$

$$x3(t) \rightarrow 3e^t(-8t + e^t - 11)$$

17.12 problem 12

17.12.1 Solution using Matrix exponential method 3780

17.12.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3781

Internal problem ID [777]

Internal file name [OUTPUT/777_Sunday_June_05_2022_01_49_24_AM_20783531/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.8, Repeated Eigenvalues. page 436

Problem number: 12.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"system of linear ODEs"**

Solve

$$x_1'(t) = -\frac{5x_1(t)}{2} + x_2(t) + x_3(t)$$

$$x_2'(t) = x_1(t) - \frac{5x_2(t)}{2} + x_3(t)$$

$$x_3'(t) = x_1(t) + x_2(t) - \frac{5x_3(t)}{2}$$

With initial conditions

$$[x_1(0) = 2, x_2(0) = 3, x_3(0) = -1]$$

17.12.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} -\frac{5}{2} & 1 & 1 \\ 1 & -\frac{5}{2} & 1 \\ 1 & 1 & -\frac{5}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{2e^{-\frac{7t}{2}}}{3} + \frac{e^{-\frac{t}{2}}}{3} & \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} & \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} \\ \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} & \frac{2e^{-\frac{7t}{2}}}{3} + \frac{e^{-\frac{t}{2}}}{3} & \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} \\ \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} & \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} & \frac{2e^{-\frac{7t}{2}}}{3} + \frac{e^{-\frac{t}{2}}}{3} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{x}_0 \\ &= \begin{bmatrix} \frac{2e^{-\frac{7t}{2}}}{3} + \frac{e^{-\frac{t}{2}}}{3} & \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} & \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} \\ \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} & \frac{2e^{-\frac{7t}{2}}}{3} + \frac{e^{-\frac{t}{2}}}{3} & \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} \\ \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} & \frac{e^{-\frac{t}{2}}}{3} - \frac{e^{-\frac{7t}{2}}}{3} & \frac{2e^{-\frac{7t}{2}}}{3} + \frac{e^{-\frac{t}{2}}}{3} \end{bmatrix} \begin{bmatrix} 2 \\ 3 \\ -1 \end{bmatrix} \\ &= \begin{bmatrix} \frac{2e^{-\frac{7t}{2}}}{3} + \frac{4e^{-\frac{t}{2}}}{3} \\ \frac{4e^{-\frac{t}{2}}}{3} + \frac{5e^{-\frac{7t}{2}}}{3} \\ \frac{4e^{-\frac{t}{2}}}{3} - \frac{7e^{-\frac{7t}{2}}}{3} \end{bmatrix} \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

17.12.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \\ x_3'(t) \end{bmatrix} = \begin{bmatrix} -\frac{5}{2} & 1 & 1 \\ 1 & -\frac{5}{2} & 1 \\ 1 & 1 & -\frac{5}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} -\frac{5}{2} & 1 & 1 \\ 1 & -\frac{5}{2} & 1 \\ 1 & 1 & -\frac{5}{2} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -\frac{5}{2} - \lambda & 1 & 1 \\ 1 & -\frac{5}{2} - \lambda & 1 \\ 1 & 1 & -\frac{5}{2} - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^3 + \frac{15}{2}\lambda^2 + \frac{63}{4}\lambda + \frac{49}{8} = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -\frac{1}{2}$$

$$\lambda_2 = -\frac{7}{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-\frac{1}{2}$	1	real eigenvalue
$-\frac{7}{2}$	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -\frac{7}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{5}{2} & 1 & 1 \\ 1 & -\frac{5}{2} & 1 \\ 1 & 1 & -\frac{5}{2} \end{bmatrix} - \left(-\frac{7}{2}\right) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} 1 & 1 & 1 & 0 \\ 1 & 1 & 1 & 0 \\ 1 & 1 & 1 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{ccc|c} 1 & 1 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 0 \end{array} \right]$$

$$R_3 = R_3 - R_1 \implies \left[\begin{array}{ccc|c} 1 & 1 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & 1 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2, v_3\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Let $v_3 = s$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -t - s\}$

Hence the solution is

$$\begin{bmatrix} -t - s \\ t \\ s \end{bmatrix} = \begin{bmatrix} -t - s \\ t \\ s \end{bmatrix}$$

Since there are two free Variable, we have found two eigenvectors associated with this eigenvalue. The above can be written as

$$\begin{aligned} \begin{bmatrix} -t - s \\ t \\ s \end{bmatrix} &= \begin{bmatrix} -t \\ t \\ 0 \end{bmatrix} + \begin{bmatrix} -s \\ 0 \\ s \end{bmatrix} \\ &= t \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} + s \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} \end{aligned}$$

By letting $t = 1$ and $s = 1$ then the above becomes

$$\begin{bmatrix} -t - s \\ t \\ s \end{bmatrix} = \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} + \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix}$$

Hence the two eigenvectors associated with this eigenvalue are

$$\left(\begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix}, \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} \right)$$

Considering the eigenvalue $\lambda_2 = -\frac{1}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{5}{2} & 1 & 1 \\ 1 & -\frac{5}{2} & 1 \\ 1 & 1 & -\frac{5}{2} \end{bmatrix} - \left(-\frac{1}{2} \right) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -2 & 1 & 1 \\ 1 & -2 & 1 \\ 1 & 1 & -2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{ccc|c} -2 & 1 & 1 & 0 \\ 1 & -2 & 1 & 0 \\ 1 & 1 & -2 & 0 \end{array} \right]$$

$$R_2 = R_2 + \frac{R_1}{2} \implies \left[\begin{array}{ccc|c} -2 & 1 & 1 & 0 \\ 0 & -\frac{3}{2} & \frac{3}{2} & 0 \\ 1 & 1 & -2 & 0 \end{array} \right]$$

$$R_3 = R_3 + \frac{R_1}{2} \implies \left[\begin{array}{ccc|c} -2 & 1 & 1 & 0 \\ 0 & -\frac{3}{2} & \frac{3}{2} & 0 \\ 0 & \frac{3}{2} & -\frac{3}{2} & 0 \end{array} \right]$$

$$R_3 = R_3 + R_2 \implies \left[\begin{array}{ccc|c} -2 & 1 & 1 & 0 \\ 0 & -\frac{3}{2} & \frac{3}{2} & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{ccc} -2 & 1 & 1 \\ 0 & -\frac{3}{2} & \frac{3}{2} \\ 0 & 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_3\}$ and the leading variables are $\{v_1, v_2\}$. Let $v_3 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t, v_2 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-\frac{1}{2}$	1	1	No	$\begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$
$-\frac{7}{2}$	2	2	No	$\begin{bmatrix} -1 & -1 \\ 0 & 1 \\ 1 & 0 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue $-\frac{1}{2}$ is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^{-\frac{t}{2}} \\ &= \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} e^{-\frac{t}{2}} \end{aligned}$$

eigenvalue $-\frac{7}{2}$ is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

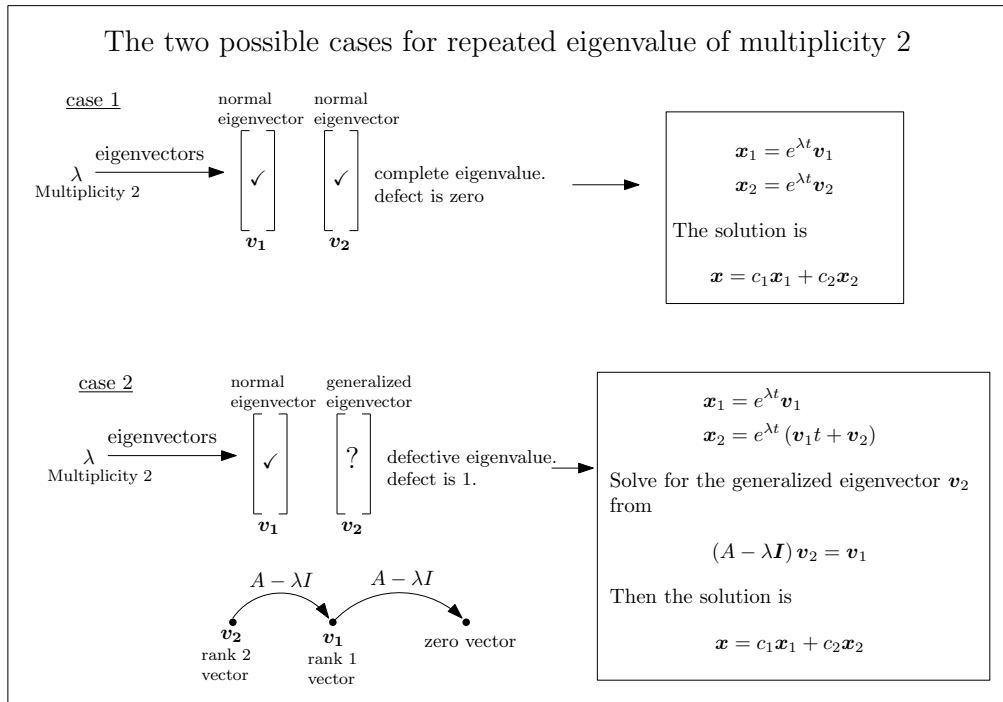


Figure 537: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2 which is the same as its geometric multiplicity 2, then it is complete eigenvalue and this falls into case 1 shown above. Hence the corresponding eigenvector basis are

$$\begin{aligned} \vec{x}_2(t) &= \vec{v}_2 e^{-\frac{7t}{2}} \\ &= \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} e^{-\frac{7t}{2}} \end{aligned}$$

$$\begin{aligned} \vec{x}_3(t) &= \vec{v}_3 e^{-\frac{7t}{2}} \\ &= \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} e^{-\frac{7t}{2}} \end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + c_3 \vec{x}_3(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = c_1 \begin{bmatrix} e^{-\frac{t}{2}} \\ e^{-\frac{t}{2}} \\ e^{-\frac{t}{2}} \end{bmatrix} + c_2 \begin{bmatrix} -e^{-\frac{7t}{2}} \\ 0 \\ e^{-\frac{7t}{2}} \end{bmatrix} + c_3 \begin{bmatrix} -e^{-\frac{7t}{2}} \\ e^{-\frac{7t}{2}} \\ 0 \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} (-c_2 - c_3)e^{-\frac{7t}{2}} + c_1e^{-\frac{t}{2}} \\ c_1e^{-\frac{t}{2}} + c_3e^{-\frac{7t}{2}} \\ c_1e^{-\frac{t}{2}} + c_2e^{-\frac{7t}{2}} \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x_1(0) = 2 \\ x_2(0) = 3 \\ x_3(0) = -1 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 2 \\ 3 \\ -1 \end{bmatrix} = \begin{bmatrix} -c_2 - c_3 + c_1 \\ c_1 + c_3 \\ c_1 + c_2 \end{bmatrix}$$

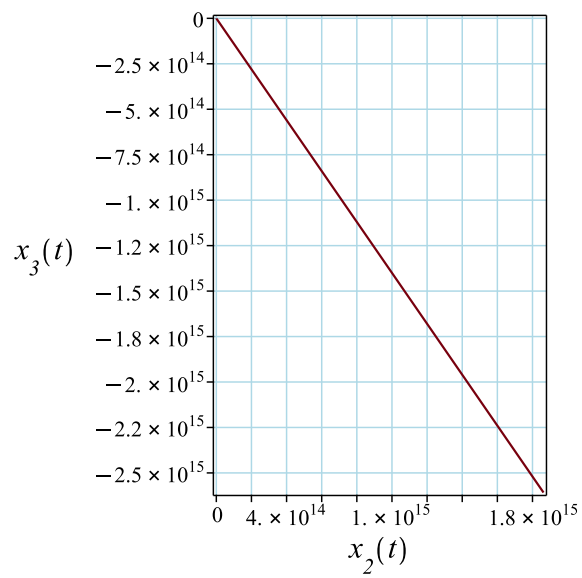
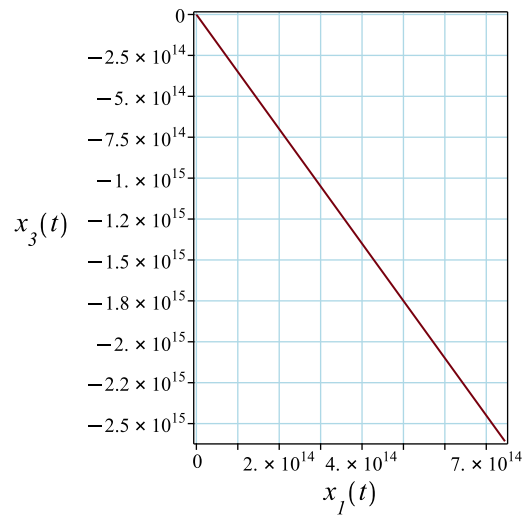
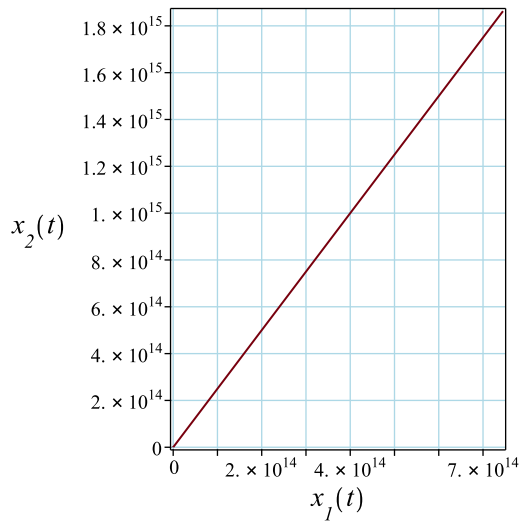
Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = \frac{4}{3} \\ c_2 = -\frac{7}{3} \\ c_3 = \frac{5}{3} \end{bmatrix}$$

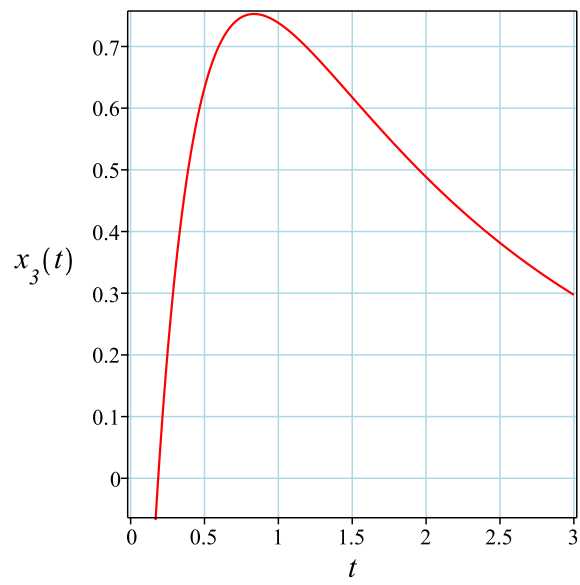
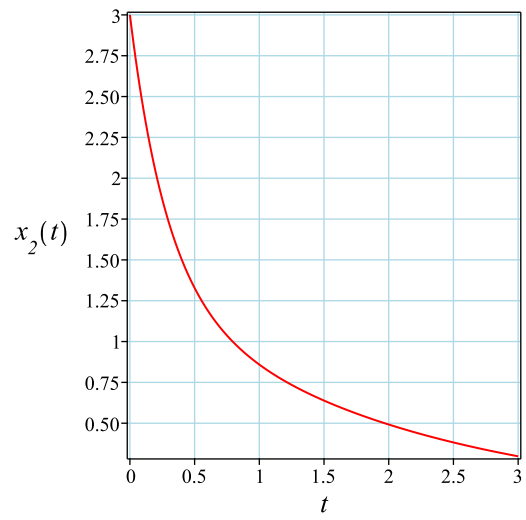
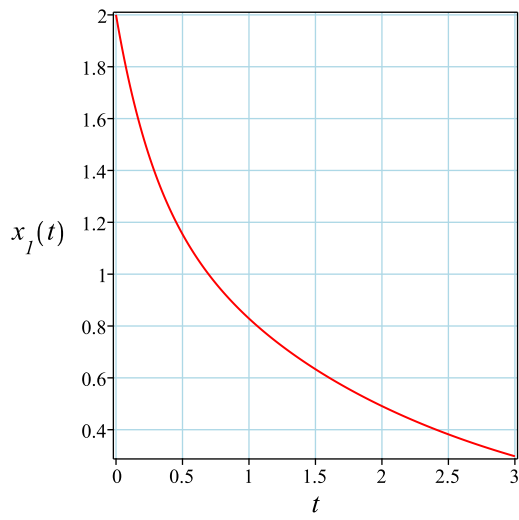
Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} = \begin{bmatrix} \frac{2e^{-\frac{7t}{2}}}{3} + \frac{4e^{-\frac{t}{2}}}{3} \\ \frac{4e^{-\frac{t}{2}}}{3} + \frac{5e^{-\frac{7t}{2}}}{3} \\ \frac{4e^{-\frac{t}{2}}}{3} - \frac{7e^{-\frac{7t}{2}}}{3} \end{bmatrix}$$

The following are plots of each solution against another.



The following are plots of each solution.



✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 50

```
dsolve([diff(x__1(t),t) = -5/2*x__1(t)+x__2(t)+x__3(t), diff(x__2(t),t) = x__1(t)-5/2*x__2(t)
```

$$\begin{aligned}x_1(t) &= \frac{2e^{-\frac{7t}{2}}}{3} + \frac{4e^{-\frac{t}{2}}}{3} \\x_2(t) &= \frac{5e^{-\frac{7t}{2}}}{3} + \frac{4e^{-\frac{t}{2}}}{3} \\x_3(t) &= -\frac{7e^{-\frac{7t}{2}}}{3} + \frac{4e^{-\frac{t}{2}}}{3}\end{aligned}$$

✓ Solution by Mathematica

Time used: 0.008 (sec). Leaf size: 71

```
DSolve[{x1'[t]==-5/2*x1[t]+1*x2[t]+1*x3[t],x2'[t]==1*x1[t]-5/2*x2[t]+1*x3[t],x3'[t]==1*x1[t]
```

$$\begin{aligned}x_1(t) &\rightarrow \frac{2}{3}e^{-7t/2}(2e^{3t} + 1) \\x_2(t) &\rightarrow \frac{1}{3}e^{-7t/2}(4e^{3t} + 5) \\x_3(t) &\rightarrow \frac{1}{3}e^{-7t/2}(4e^{3t} - 7)\end{aligned}$$

18 Chapter 7.9, Nonhomogeneous Linear Systems.

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18.1 problem 1

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Internal problem ID [778]

Internal file name [OUTPUT/778_Sunday_June_05_2022_01_49_25_AM_67064786/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

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Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 2x_1(t) - x_2(t) + e^t \\x_2'(t) &= 3x_1(t) - 2x_2(t) + t\end{aligned}$$

18.1.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} e^t \\ t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} \left(-\frac{e^{-t}}{2} + \frac{3e^t}{2}\right) c_1 + \left(-\frac{e^t}{2} + \frac{e^{-t}}{2}\right) c_2 \\ \left(\frac{3e^t}{2} - \frac{3e^{-t}}{2}\right) c_1 + \left(\frac{3e^{-t}}{2} - \frac{e^t}{2}\right) c_2 \end{bmatrix} \\ &= \begin{bmatrix} \frac{(c_2 - c_1)e^{-t}}{2} + \frac{3(-\frac{c_2}{3} + c_1)e^t}{2} \\ \frac{(-3c_1 + 3c_2)e^{-t}}{2} + \frac{3(-\frac{c_2}{3} + c_1)e^t}{2} \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \frac{3e^{-t}}{2} - \frac{e^t}{2} & \frac{e^t}{2} - \frac{e^{-t}}{2} \\ -\frac{3e^t}{2} + \frac{3e^{-t}}{2} & -\frac{e^{-t}}{2} + \frac{3e^t}{2} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix} \int \begin{bmatrix} \frac{3e^{-t}}{2} - \frac{e^t}{2} & \frac{e^t}{2} - \frac{e^{-t}}{2} \\ -\frac{3e^t}{2} + \frac{3e^{-t}}{2} & -\frac{e^{-t}}{2} + \frac{3e^t}{2} \end{bmatrix} \begin{bmatrix} e^t \\ t \end{bmatrix} dt \\ &= \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix} \begin{bmatrix} \frac{(2t+2)e^{-t}}{4} - \frac{e^{2t}}{4} + \frac{(2t-2)e^t}{4} + \frac{3t}{2} \\ \frac{(2t+2)e^{-t}}{4} - \frac{3e^{2t}}{4} + \frac{(6t-6)e^t}{4} + \frac{3t}{2} \end{bmatrix} \\ &= \begin{bmatrix} \frac{(6t-1)e^t}{4} + t \\ \frac{(6t-3)e^t}{4} + 2t - 1 \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{(-2c_1+2c_2)e^{-t}}{4} + \frac{(6t+6c_1-2c_2-1)e^t}{4} + t \\ \frac{(6c_2-6c_1)e^{-t}}{4} + \frac{(6t+6c_1-2c_2-3)e^t}{4} + 2t - 1 \end{bmatrix}\end{aligned}$$

18.1.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} e^t \\ t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 2 - \lambda & -1 \\ 3 & -2 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 1$$

$$\lambda_2 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
1	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 3 & -1 \\ 3 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 3 & -1 & 0 \\ 3 & -1 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} 3 & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 3 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{t}{3}\}$

Hence the solution is

$$\begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix} = \begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} = \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ 3 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} - (1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & -1 \\ 3 & -3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 & -1 & 0 \\ 3 & -3 & 0 \end{array} \right]$$

$$R_2 = R_2 - 3R_1 \implies \left[\begin{array}{cc|c} 1 & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
1	1	1	No	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$
-1	1	1	No	$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue 1 is real and distinct then the

corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^t \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^t\end{aligned}$$

Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-t} \\ &= \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} e^{-t}\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} e^t \\ e^t \end{bmatrix} + c_2 \begin{bmatrix} \frac{e^{-t}}{3} \\ e^{-t} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} e^t & \frac{e^{-t}}{3} \\ e^t & e^{-t} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} \frac{3e^{-t}}{2} & -\frac{e^{-t}}{2} \\ -\frac{3e^t}{2} & \frac{3e^t}{2} \end{bmatrix}$$

Hence

$$\begin{aligned}
 \vec{x}_p(t) &= \begin{bmatrix} e^t & \frac{e^{-t}}{3} \\ e^t & e^{-t} \end{bmatrix} \int \begin{bmatrix} \frac{3e^{-t}}{2} & -\frac{e^{-t}}{2} \\ -\frac{3e^t}{2} & \frac{3e^t}{2} \end{bmatrix} \begin{bmatrix} e^t \\ t \end{bmatrix} dt \\
 &= \begin{bmatrix} e^t & \frac{e^{-t}}{3} \\ e^t & e^{-t} \end{bmatrix} \int \begin{bmatrix} \frac{3}{2} - \frac{te^{-t}}{2} \\ -\frac{3e^{2t}}{2} + \frac{3te^t}{2} \end{bmatrix} dt \\
 &= \begin{bmatrix} e^t & \frac{e^{-t}}{3} \\ e^t & e^{-t} \end{bmatrix} \begin{bmatrix} \frac{e^{-t}(t+1)}{2} + \frac{3t}{2} \\ -\frac{3e^{2t}}{4} + \frac{(6t-6)e^t}{4} \end{bmatrix} \\
 &= \begin{bmatrix} \frac{(6t-1)e^t}{4} + t \\ \frac{(6t-3)e^t}{4} + 2t - 1 \end{bmatrix}
 \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned}
 \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\
 \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} c_1 e^t \\ c_1 e^t \end{bmatrix} + \begin{bmatrix} \frac{c_2 e^{-t}}{3} \\ c_2 e^{-t} \end{bmatrix} + \begin{bmatrix} \frac{(6t-1)e^t}{4} + t \\ \frac{(6t-3)e^t}{4} + 2t - 1 \end{bmatrix}
 \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{c_2 e^{-t}}{3} + \frac{(-1+6t+4c_1)e^t}{4} + t \\ c_2 e^{-t} + \frac{(6t+4c_1-3)e^t}{4} + 2t - 1 \end{bmatrix}$$

18.1.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 2x_1(t) - x_2(t) + e^t, x_2'(t) = 3x_1(t) - 2x_2(t) + t]$$

- Define vector

$$x_{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$x_{\rightarrow}'(t) = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \cdot x_{\rightarrow}(t) + \begin{bmatrix} e^t \\ t \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} e^t \\ t \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} e^t \\ t \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow}{}'(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-1, \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} \right], \left[1, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-1, \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{-t} \cdot \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[1, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_2 = e^t \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- General solution of the system of ODEs can be written in terms of the particular solution $\vec{x}_p(t)$

$$\vec{x}(t) = c_1 \vec{x}_1 + c_2 \vec{x}_2 + \vec{x}_p(t)$$

□ Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} \frac{e^{-t}}{3} & e^t \\ e^{-t} & e^t \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix

$$\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} \frac{e^{-t}}{3} & e^t \\ e^{-t} & e^t \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} \frac{1}{3} & 1 \\ 1 & 1 \end{bmatrix}}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix}$$

□ Find a particular solution of the system of ODEs using variation of parameters

- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\vec{x}_p(t) = \Phi(t) \cdot \vec{v}(t)$$

- Take the derivative of the particular solution

$$\vec{x}_p'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t)$

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- Plug $\vec{v}(t)$ into the equation for the particular solution

$$\underline{x}_{\rightarrow p}(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- Plug in the fundamental matrix and the forcing function and compute

$$\underline{x}_{\rightarrow p}(t) = \begin{bmatrix} \frac{3e^{-t}}{4} + \frac{(6t-3)e^t}{4} + t \\ \frac{9e^{-t}}{4} + \frac{(-5+6t)e^t}{4} + 2t - 1 \end{bmatrix}$$

- Plug particular solution back into general solution

$$\underline{x}_{\rightarrow}(t) = c_1 \underline{x}_{\rightarrow 1} + c_2 \underline{x}_{\rightarrow 2} + \begin{bmatrix} \frac{3e^{-t}}{4} + \frac{(6t-3)e^t}{4} + t \\ \frac{9e^{-t}}{4} + \frac{(-5+6t)e^t}{4} + 2t - 1 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(9+4c_1)e^{-t}}{12} + \frac{(-3+6t+4c_2)e^t}{4} + t \\ \frac{(9+4c_1)e^{-t}}{4} + \frac{(6t+4c_2-5)e^t}{4} + 2t - 1 \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{(9+4c_1)e^{-t}}{12} + \frac{(-3+6t+4c_2)e^t}{4} + t, x_2(t) = \frac{(9+4c_1)e^{-t}}{4} + \frac{(6t+4c_2-5)e^t}{4} + 2t - 1 \right\}$$

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 54

```
dsolve([diff(x__1(t),t)=2*x__1(t)-1*x__2(t)+exp(t),diff(x__2(t),t)=3*x__1(t)-2*x__2(t)+t],si
```

$$x_1(t) = \frac{e^{-t}c_1}{3} + c_2e^t + \frac{3e^t t}{2} - \frac{e^t}{4} + t$$

$$x_2(t) = c_2e^t + e^{-t}c_1 + \frac{3e^t t}{2} - \frac{3e^t}{4} + 2t - 1$$

✓ Solution by Mathematica

Time used: 0.123 (sec). Leaf size: 97

```
DSolve[{x1'[t]==2*x1[t]-1*x2[t]+Exp[t],x2'[t]==3*x1[t]-2*x2[t]+t},{x1[t],x2[t]},t,IncludeSin
```

$$x1(t) \rightarrow \frac{1}{4}e^{-t}(4e^t t + e^{2t}(6t - 1 + 6c_1 - 2c_2) - 2c_1 + 2c_2)$$

$$x2(t) \rightarrow \frac{1}{4}e^{-t}(e^t(8t - 4) + e^{2t}(6t - 3 + 6c_1 - 2c_2) - 6c_1 + 6c_2)$$

18.2 problem 2

18.2.1 Solution using Matrix exponential method 3805

18.2.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3807

Internal problem ID [779]

Internal file name [OUTPUT/779_Sunday_June_05_2022_01_49_27_AM_63990030/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) + \sqrt{3}x_2(t) + e^t \\x_2'(t) &= \sqrt{3}x_1(t) - x_2(t) + \sqrt{3}e^{-t}\end{aligned}$$

18.2.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & \sqrt{3} \\ \sqrt{3} & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} e^t \\ \sqrt{3}e^{-t} \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{e^{-2t}}{4} + \frac{3e^{2t}}{4} & -\frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} \\ -\frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} & \frac{3e^{-2t}}{4} + \frac{e^{2t}}{4} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \frac{e^{-2t}}{4} + \frac{3e^{2t}}{4} & -\frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} \\ -\frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} & \frac{3e^{-2t}}{4} + \frac{e^{2t}}{4} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} \left(\frac{e^{-2t}}{4} + \frac{3e^{2t}}{4}\right)c_1 - \frac{\sqrt{3}(e^{-2t}-e^{2t})c_2}{4} \\ -\frac{\sqrt{3}(e^{-2t}-e^{2t})c_1}{4} + \left(\frac{3e^{-2t}}{4} + \frac{e^{2t}}{4}\right)c_2 \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \frac{3e^{-2t}}{4} + \frac{e^{2t}}{4} & \frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} \\ \frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} & \frac{e^{-2t}}{4} + \frac{3e^{2t}}{4} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{e^{-2t}}{4} + \frac{3e^{2t}}{4} & -\frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} \\ -\frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} & \frac{3e^{-2t}}{4} + \frac{e^{2t}}{4} \end{bmatrix} \int \begin{bmatrix} \frac{3e^{-2t}}{4} + \frac{e^{2t}}{4} & \frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} \\ \frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} & \frac{e^{-2t}}{4} + \frac{3e^{2t}}{4} \end{bmatrix} \begin{bmatrix} e^t \\ \sqrt{3}e^{-t} \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{e^{-2t}}{4} + \frac{3e^{2t}}{4} & -\frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} \\ -\frac{\sqrt{3}(e^{-2t}-e^{2t})}{4} & \frac{3e^{-2t}}{4} + \frac{e^{2t}}{4} \end{bmatrix} \begin{bmatrix} \frac{(e^{6t}-9e^{4t}-9e^{2t}-3)e^{-3t}}{12} \\ -\frac{\sqrt{3}(e^{6t}-9e^{4t}+3e^{2t}+1)e^{-3t}}{12} \end{bmatrix} \\ &= \begin{bmatrix} -\frac{2e^t}{3} - e^{-t} \\ -\frac{\sqrt{3}(e^t-2e^{-t})}{3} \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{e^{-2t} \left((c_2\sqrt{3}+3c_1)e^{4t} - c_2\sqrt{3} + c_1 - 4e^t - \frac{8e^{3t}}{3} \right)}{4} \\ \frac{\left((c_1\sqrt{3}+c_2)e^{4t} - \frac{4\sqrt{3}e^{3t}}{3} + \left(-c_1 + \frac{8e^t}{3}\right)\sqrt{3} + 3c_2 \right) e^{-2t}}{4} \end{bmatrix}\end{aligned}$$

18.2.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & \sqrt{3} \\ \sqrt{3} & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} e^t \\ \sqrt{3}e^{-t} \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 1 & \sqrt{3} \\ \sqrt{3} & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 1 - \lambda & \sqrt{3} \\ \sqrt{3} & -1 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 4 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -2$$

$$\lambda_2 = 2$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-2	1	real eigenvalue
2	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & \sqrt{3} \\ \sqrt{3} & -1 \end{bmatrix} - (-2) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} 3 & \sqrt{3} \\ \sqrt{3} & 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 3 & \sqrt{3} & 0 \\ \sqrt{3} & 1 & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{\sqrt{3}R_1}{3} \implies \left[\begin{array}{cc|c} 3 & \sqrt{3} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 3 & \sqrt{3} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\left\{ v_1 = -\frac{\sqrt{3}t}{3} \right\}$

Hence the solution is

$$\begin{bmatrix} -\frac{\sqrt{3}t}{3} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{\sqrt{3}t}{3} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{\sqrt{3}t}{3} \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{\sqrt{3}}{3} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{\sqrt{3}}{3} \\ 1 \end{bmatrix} = \begin{bmatrix} -\frac{\sqrt{3}}{3} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{\sqrt{3}}{3} \\ 1 \end{bmatrix} = \begin{bmatrix} -\sqrt{3} \\ 3 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & \sqrt{3} \\ \sqrt{3} & -1 \end{bmatrix} - (2) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 & \sqrt{3} \\ \sqrt{3} & -3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -1 & \sqrt{3} & 0 \\ \sqrt{3} & -3 & 0 \end{array} \right]$$

$$R_2 = R_2 + \sqrt{3}R_1 \implies \left[\begin{array}{cc|c} -1 & \sqrt{3} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 & \sqrt{3} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \sqrt{3}t\}$

Hence the solution is

$$\begin{bmatrix} \sqrt{3}t \\ t \end{bmatrix} = \begin{bmatrix} \sqrt{3}t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \sqrt{3}t \\ t \end{bmatrix} = t \begin{bmatrix} \sqrt{3} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \sqrt{3}t \\ t \end{bmatrix} = \begin{bmatrix} \sqrt{3} \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-2	1	1	No	$\begin{bmatrix} -\frac{\sqrt{3}}{3} \\ 1 \end{bmatrix}$
2	1	1	No	$\begin{bmatrix} \sqrt{3} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{-2t} \\ &= \begin{bmatrix} -\frac{\sqrt{3}}{3} \\ 1 \end{bmatrix} e^{-2t}\end{aligned}$$

Since eigenvalue 2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{2t} \\ &= \begin{bmatrix} \sqrt{3} \\ 1 \end{bmatrix} e^{2t}\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} -\frac{\sqrt{3}e^{-2t}}{3} \\ e^{-2t} \end{bmatrix} + c_2 \begin{bmatrix} \sqrt{3}e^{2t} \\ e^{2t} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} -\frac{\sqrt{3}e^{-2t}}{3} & \sqrt{3}e^{2t} \\ e^{-2t} & e^{2t} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -\frac{\sqrt{3}e^{2t}}{4} & \frac{3e^{2t}}{4} \\ \frac{\sqrt{3}e^{-2t}}{4} & \frac{e^{-2t}}{4} \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} -\frac{\sqrt{3}e^{-2t}}{3} & \sqrt{3}e^{2t} \\ e^{-2t} & e^{2t} \end{bmatrix} \int \begin{bmatrix} -\frac{\sqrt{3}e^{2t}}{4} & \frac{3e^{2t}}{4} \\ \frac{\sqrt{3}e^{-2t}}{4} & \frac{e^{-2t}}{4} \end{bmatrix} \begin{bmatrix} e^t \\ \sqrt{3}e^{-t} \end{bmatrix} dt \\ &= \begin{bmatrix} -\frac{\sqrt{3}e^{-2t}}{3} & \sqrt{3}e^{2t} \\ e^{-2t} & e^{2t} \end{bmatrix} \int \begin{bmatrix} -\frac{\sqrt{3}e^t(e^{2t}-3)}{4} \\ \frac{\sqrt{3}e^{-3t}(e^{2t}+1)}{4} \end{bmatrix} dt \\ &= \begin{bmatrix} -\frac{\sqrt{3}e^{-2t}}{3} & \sqrt{3}e^{2t} \\ e^{-2t} & e^{2t} \end{bmatrix} \begin{bmatrix} \frac{(-e^{3t}+9e^t)\sqrt{3}}{12} \\ -\frac{\sqrt{3}(3e^{2t}+1)e^{-3t}}{12} \end{bmatrix} \\ &= \begin{bmatrix} -\frac{2e^t}{3} - e^{-t} \\ -\frac{\sqrt{3}(e^t-2e^{-t})}{3} \end{bmatrix} \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} -\frac{c_1\sqrt{3}e^{-2t}}{3} \\ c_1e^{-2t} \end{bmatrix} + \begin{bmatrix} c_2\sqrt{3}e^{2t} \\ c_2e^{2t} \end{bmatrix} + \begin{bmatrix} -\frac{2e^t}{3} - e^{-t} \\ -\frac{\sqrt{3}(e^t-2e^{-t})}{3} \end{bmatrix} \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(-3c_2\sqrt{3}e^{4t} + \sqrt{3}c_1 + 3e^t + 2e^{3t})e^{-2t}}{3} \\ \frac{(3e^{4t}c_2 - \sqrt{3}e^{3t} + 2\sqrt{3}e^t + 3c_1)e^{-2t}}{3} \end{bmatrix}$$

✓ Solution by Maple

Time used: 0.141 (sec). Leaf size: 71

```
dsolve([diff(x__1(t),t)=1*x__1(t)+sqrt(3)*x__2(t)+exp(t),diff(x__2(t),t)=sqrt(3)*x__1(t)-1*x__2(t)],t)
```

$$\begin{aligned} x_1(t) &= \sinh(2t)c_2 + \cosh(2t)c_1 - \frac{5\cosh(t)}{3} + \frac{\sinh(t)}{3} \\ x_2(t) &= \frac{\sqrt{3}(\cosh(2t)c_1 - 2\cosh(2t)c_2 - 2\sinh(2t)c_1 + \sinh(2t)c_2 + e^t + 2\sinh(t) - 2\cosh(t))}{3} \end{aligned}$$

✓ Solution by Mathematica

Time used: 2.501 (sec). Leaf size: 313

```
DSolve[{x1'[t]==1*x1[t]+Sqrt[4]*x2[t]+Exp[t],x2'[t]==Sqrt[3]*x1[t]-1*x2[t]+Sqrt[3]*Exp[-t]},
```

$$x1(t) \rightarrow \frac{1}{6} \left(-6e^{-t} - \frac{2(6 + \sqrt{3}) e^t}{1 + 2\sqrt{3}} + \frac{3(\sqrt{1 + 2\sqrt{3}} - 1) c_1 - 6c_2}{\sqrt{1 + 2\sqrt{3}}} e^{-\sqrt{1+2\sqrt{3}}t} \right. \\ \left. + \frac{3\left(\left(1 + \sqrt{1 + 2\sqrt{3}}\right) c_1 + 2c_2\right) e^{\sqrt{1+2\sqrt{3}}t}}{\sqrt{1 + 2\sqrt{3}}} \right)$$
$$x2(t) \rightarrow \frac{1}{4} \left(4e^{-t} - 2e^t + \frac{2\left((6 + \sqrt{3}) c_1 + (1 + 2\sqrt{3}) (\sqrt{1 + 2\sqrt{3}} - 1) c_2\right) e^{\sqrt{1+2\sqrt{3}}t}}{(1 + 2\sqrt{3})^{3/2}} \right. \\ \left. + \frac{\left(2(1 + 2\sqrt{3}) (1 + \sqrt{1 + 2\sqrt{3}}) c_2 - 2(6 + \sqrt{3}) c_1\right) e^{-\sqrt{1+2\sqrt{3}}t}}{(1 + 2\sqrt{3})^{3/2}} \right)$$

18.3 problem 3

- 18.3.1 Solution using Matrix exponential method 3814
- 18.3.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3816
- 18.3.3 Maple step by step solution 3821

Internal problem ID [780]

Internal file name [OUTPUT/780_Sunday_June_05_2022_01_49_29_AM_86783349/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 3.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 2x_1(t) - 5x_2(t) - \cos(t) \\x_2'(t) &= x_1(t) - 2x_2(t) + \sin(t)\end{aligned}$$

18.3.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} -\cos(t) \\ \sin(t) \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} (\cos(t) + 2 \sin(t)) c_1 - 5 \sin(t) c_2 \\ \sin(t) c_1 + (\cos(t) - 2 \sin(t)) c_2 \end{bmatrix} \\ &= \begin{bmatrix} (2c_1 - 5c_2) \sin(t) + c_1 \cos(t) \\ (c_1 - 2c_2) \sin(t) + c_2 \cos(t) \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \cos(t) - 2 \sin(t) & 5 \sin(t) \\ -\sin(t) & \cos(t) + 2 \sin(t) \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \int \begin{bmatrix} \cos(t) - 2 \sin(t) & 5 \sin(t) \\ -\sin(t) & \cos(t) + 2 \sin(t) \end{bmatrix} \begin{bmatrix} -\cos(t) \\ \sin(t) \end{bmatrix} dt \\ &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \begin{bmatrix} -3 \sin(t) \cos(t) + 2t - \cos(t)^2 \\ -\cos(t)^2 - \sin(t) \cos(t) + t + \frac{1}{2} \end{bmatrix} \\ &= \begin{bmatrix} \cos(t) (2t - 1) + \frac{(-2t-5) \sin(t)}{2} \\ -\frac{\cos(t)}{2} - \sin(t) + t \cos(t) \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{(-2t+4c_1-10c_2-5)\sin(t)}{2} + 2\left(t + \frac{c_1}{2} - \frac{1}{2}\right)\cos(t) \\ \frac{(2t+2c_2-1)\cos(t)}{2} + \sin(t)(c_1 - 2c_2 - 1) \end{bmatrix}\end{aligned}$$

18.3.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} -\cos(t) \\ \sin(t) \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 2 - \lambda & -5 \\ 1 & -2 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-i$	1	complex eigenvalue
i	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - (-i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2+i & -5 \\ 1 & -2+i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2+i & -5 & 0 \\ 1 & -2+i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} + \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2+i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2+i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 - i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 - I)t \\ t \end{bmatrix} = \begin{bmatrix} (2 - i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 - I)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 - I)t \\ t \end{bmatrix} = \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - (i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 - i & -5 \\ 1 & -2 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 1 & -2 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} - \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 - i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 + i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 + I) t \\ t \end{bmatrix} = \begin{bmatrix} (2 + i) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 + I) t \\ t \end{bmatrix} = t \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 + I) t \\ t \end{bmatrix} = \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
i	1	1	No	$\begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$
$-i$	1	1	No	$\begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (2 + i) e^{it} \\ e^{it} \end{bmatrix} + c_2 \begin{bmatrix} (2 - i) e^{-it} \\ e^{-it} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -\frac{ie^{-it}}{2} & (\frac{1}{2} + i)e^{-it} \\ \frac{ie^{it}}{2} & (\frac{1}{2} - i)e^{it} \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix} \int \begin{bmatrix} -\frac{ie^{-it}}{2} & (\frac{1}{2} + i)e^{-it} \\ \frac{ie^{it}}{2} & (\frac{1}{2} - i)e^{it} \end{bmatrix} \begin{bmatrix} -\cos(t) \\ \sin(t) \end{bmatrix} dt \\ &= \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix} \int \begin{bmatrix} \frac{((1+2i)\sin(t)+i\cos(t))e^{-it}}{2} \\ -\frac{((-1+2i)\sin(t)+i\cos(t))e^{it}}{2} \end{bmatrix} dt \\ &= \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix} \begin{bmatrix} \frac{((it+i-1)\sin(t)+t\cos(t))e^{-it}}{2} \\ -\frac{((it+i+1)\sin(t)-t\cos(t))e^{it}}{2} \end{bmatrix} \\ &= \begin{bmatrix} (-t-3)\sin(t) + 2t\cos(t) \\ t\cos(t) - \sin(t) \end{bmatrix} \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} (2+i)c_1e^{it} \\ c_1e^{it} \end{bmatrix} + \begin{bmatrix} (2-i)c_2e^{-it} \\ c_2e^{-it} \end{bmatrix} + \begin{bmatrix} (-t-3)\sin(t) + 2t\cos(t) \\ t\cos(t) - \sin(t) \end{bmatrix} \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (2+i)c_1e^{it} + (2-i)c_2e^{-it} + (-t-3)\sin(t) + 2t\cos(t) \\ c_1e^{it} + c_2e^{-it} + t\cos(t) - \sin(t) \end{bmatrix}$$

18.3.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 2x_1(t) - 5x_2(t) - \cos(t), x_2'(t) = x_1(t) - 2x_2(t) + \sin(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} -\cos(t) \\ \sin(t) \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} -\cos(t) \\ \sin(t) \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} -\cos(t) \\ \sin(t) \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right], \left[I, \begin{bmatrix} 2 + I \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{-It} \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$(\cos(t) - I \sin(t)) \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Simplify expression

$$\begin{bmatrix} (2 - I)(\cos(t) - I \sin(t)) \\ \cos(t) - I \sin(t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\begin{array}{l} \underline{x}_{\rightarrow 1}(t) = \begin{bmatrix} 2 \cos(t) - \sin(t) \\ \cos(t) \end{bmatrix}, \underline{x}_{\rightarrow 2}(t) = \begin{bmatrix} -\cos(t) - 2 \sin(t) \\ -\sin(t) \end{bmatrix} \end{array} \right]$$

- General solution of the system of ODEs can be written in terms of the particular solution $\underline{x}_{\rightarrow p}(t)$

$$\underline{x}_{\rightarrow}(t) = c_1 \underline{x}_{\rightarrow 1}(t) + c_2 \underline{x}_{\rightarrow 2}(t) + \underline{x}_{\rightarrow p}(t)$$

□ Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} 2 \cos(t) - \sin(t) & -\cos(t) - 2 \sin(t) \\ \cos(t) & -\sin(t) \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix

$$\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} 2 \cos(t) - \sin(t) & -\cos(t) - 2 \sin(t) \\ \cos(t) & -\sin(t) \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} 2 & -1 \\ 1 & 0 \end{bmatrix}}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix}$$

- Find a particular solution of the system of ODEs using variation of parameters
- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\underline{x}^{\rightarrow}_p(t) = \Phi(t) \cdot \vec{v}(t)$$
 - Take the derivative of the particular solution
$$\underline{x}^{\rightarrow \prime}_p(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$
 - Substitute particular solution and its derivative into the system of ODEs
$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$
 - The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t)$

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$
 - Cancel like terms
$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$
 - Multiply by the inverse of the fundamental matrix
$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$
 - Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$
 - Plug $\vec{v}(t)$ into the equation for the particular solution
$$\underline{x}^{\rightarrow}_p(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$
 - Plug in the fundamental matrix and the forcing function and compute
$$\underline{x}^{\rightarrow}_p(t) = \begin{bmatrix} (-t - 3) \sin(t) + 2t \cos(t) \\ t \cos(t) - \sin(t) \end{bmatrix}$$
 - Plug particular solution back into general solution
$$\underline{x}^{\rightarrow}(t) = c_1 \underline{x}^{\rightarrow}_1(t) + c_2 \underline{x}^{\rightarrow}_2(t) + \begin{bmatrix} (-t - 3) \sin(t) + 2t \cos(t) \\ t \cos(t) - \sin(t) \end{bmatrix}$$
 - Substitute in vector of dependent variables
$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (-c_1 - 2c_2 - t - 3) \sin(t) + 2(t + c_1 - \frac{c_2}{2}) \cos(t) \\ (t + c_1) \cos(t) - \sin(t) (c_2 + 1) \end{bmatrix}$$
 - Solution to the system of ODEs

$$\{x_1(t) = (-c_1 - 2c_2 - t - 3) \sin(t) + 2(t + c_1 - \frac{c_2}{2}) \cos(t), x_2(t) = (t + c_1) \cos(t) - \sin(t) (c_2 + \dots)$$

✓ Solution by Maple

Time used: 0.063 (sec). Leaf size: 60

```
dsolve([diff(x__1(t),t)=2*x__1(t)-5*x__2(t)-cos(t),diff(x__2(t),t)=1*x__1(t)-2*x__2(t)+sin(t)
```

$$x_1(t) = c_2 \sin(t) - \sin(t) t + c_1 \cos(t) + 2 \cos(t) t - \cos(t)$$

$$x_2(t) = \frac{c_1 \sin(t)}{5} + \frac{2c_2 \sin(t)}{5} + \frac{2c_1 \cos(t)}{5} - \frac{c_2 \cos(t)}{5} + \cos(t) t - \cos(t)$$

✓ Solution by Mathematica

Time used: 0.069 (sec). Leaf size: 61

```
DSolve[{x1'[t]==2*x1[t]-5*x2[t]-Cos[t],x2'[t]==1*x1[t]-2*x2[t]+Sin[t]},{x1[t],x2[t]},t,Inclu
```

$$x_1(t) \rightarrow \left(2t - \frac{1}{2} + c_1\right) \cos(t) - (t - 1 - 2c_1 + 5c_2) \sin(t)$$

$$x_2(t) \rightarrow (t - 1 + c_2) \cos(t) + \frac{1}{2}(1 + 2c_1 - 4c_2) \sin(t)$$

18.4 problem 4

18.4.1 Solution using Matrix exponential method	3825
18.4.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3827
18.4.3 Maple step by step solution	3832

Internal problem ID [781]

Internal file name [OUTPUT/781_Sunday_June_05_2022_01_49_32_AM_79929771/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 4.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) + x_2(t) + e^{-2t} \\x_2'(t) &= 4x_1(t) - 2x_2(t) - 2e^t\end{aligned}$$

18.4.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ 4 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} e^{-2t} \\ -2e^t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{(4e^{5t}+1)e^{-3t}}{5} & \frac{(e^{5t}-1)e^{-3t}}{5} \\ \frac{4(e^{5t}-1)e^{-3t}}{5} & \frac{(e^{5t}+4)e^{-3t}}{5} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \frac{(4e^{5t}+1)e^{-3t}}{5} & \frac{(e^{5t}-1)e^{-3t}}{5} \\ \frac{4(e^{5t}-1)e^{-3t}}{5} & \frac{(e^{5t}+4)e^{-3t}}{5} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} \frac{(4e^{5t}+1)e^{-3t}c_1}{5} + \frac{(e^{5t}-1)e^{-3t}c_2}{5} \\ \frac{4(e^{5t}-1)e^{-3t}c_1}{5} + \frac{(e^{5t}+4)e^{-3t}c_2}{5} \end{bmatrix} \\ &= \begin{bmatrix} \frac{e^{-3t}((4c_1+c_2)e^{5t}-c_2+c_1)}{5} \\ \frac{4((c_1+\frac{c_2}{4})e^{5t}+c_2-c_1)e^{-3t}}{5} \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \frac{(e^{5t}+4)e^{-2t}}{5} & -\frac{(e^{5t}-1)e^{-2t}}{5} \\ -\frac{4(e^{5t}-1)e^{-2t}}{5} & \frac{(4e^{5t}+1)e^{-2t}}{5} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{(4e^{5t}+1)e^{-3t}}{5} & \frac{(e^{5t}-1)e^{-3t}}{5} \\ \frac{4(e^{5t}-1)e^{-3t}}{5} & \frac{(e^{5t}+4)e^{-3t}}{5} \end{bmatrix} \int \begin{bmatrix} \frac{(e^{5t}+4)e^{-2t}}{5} & -\frac{(e^{5t}-1)e^{-2t}}{5} \\ -\frac{4(e^{5t}-1)e^{-2t}}{5} & \frac{(4e^{5t}+1)e^{-2t}}{5} \end{bmatrix} \begin{bmatrix} e^{-2t} \\ -2e^t \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{(4e^{5t}+1)e^{-3t}}{5} & \frac{(e^{5t}-1)e^{-3t}}{5} \\ \frac{4(e^{5t}-1)e^{-3t}}{5} & \frac{(e^{5t}+4)e^{-3t}}{5} \end{bmatrix} \begin{bmatrix} \frac{(e^{8t}+2e^{5t}+4e^{3t}-2)e^{-4t}}{10} \\ -\frac{(2e^{8t}+4e^{5t}-2e^{3t}+1)e^{-4t}}{5} \end{bmatrix} \\ &= \begin{bmatrix} \frac{e^t}{2} \\ -e^{-2t} \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{e^{-3t} \left((4c_1 + c_2)e^{5t} + c_1 - c_2 + \frac{5e^{4t}}{2} \right)}{5} \\ \frac{4 \left((c_1 + \frac{c_2}{4})e^{5t} - c_1 + c_2 - \frac{5e^t}{4} \right) e^{-3t}}{5} \end{bmatrix}\end{aligned}$$

18.4.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ 4 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} e^{-2t} \\ -2e^t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 1 & 1 \\ 4 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 1 - \lambda & 1 \\ 4 & -2 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + \lambda - 6 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 2$$

$$\lambda_2 = -3$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-3	1	real eigenvalue
2	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -3$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 1 \\ 4 & -2 \end{bmatrix} - (-3) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 4 & 1 \\ 4 & 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 4 & 1 & 0 \\ 4 & 1 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} 4 & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 4 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -\frac{t}{4}\}$

Hence the solution is

$$\begin{bmatrix} -\frac{t}{4} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{t}{4} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{t}{4} \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix} = \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix} = \begin{bmatrix} -1 \\ 4 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 1 \\ 4 & -2 \end{bmatrix} - (2) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 & 1 \\ 4 & -4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -1 & 1 & 0 \\ 4 & -4 & 0 \end{array} \right]$$

$$R_2 = R_2 + 4R_1 \implies \left[\begin{array}{cc|c} -1 & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
2	1	1	No	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$
-3	1	1	No	$\begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue 2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{2t} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{2t}\end{aligned}$$

Since eigenvalue -3 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-3t} \\ &= \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix} e^{-3t}\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} e^{2t} \\ e^{2t} \end{bmatrix} + c_2 \begin{bmatrix} -\frac{e^{-3t}}{4} \\ e^{-3t} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} e^{2t} & -\frac{e^{-3t}}{4} \\ e^{2t} & e^{-3t} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} \frac{4e^{-2t}}{5} & \frac{e^{-2t}}{5} \\ -\frac{4e^{3t}}{5} & \frac{4e^{3t}}{5} \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} e^{2t} & -\frac{e^{-3t}}{4} \\ e^{2t} & e^{-3t} \end{bmatrix} \int \begin{bmatrix} \frac{4e^{-2t}}{5} & \frac{e^{-2t}}{5} \\ -\frac{4e^{3t}}{5} & \frac{4e^{3t}}{5} \end{bmatrix} \begin{bmatrix} e^{-2t} \\ -2e^t \end{bmatrix} dt \\ &= \begin{bmatrix} e^{2t} & -\frac{e^{-3t}}{4} \\ e^{2t} & e^{-3t} \end{bmatrix} \int \begin{bmatrix} \frac{4e^{-4t}}{5} - \frac{2e^{-t}}{5} \\ -\frac{4e^t}{5} - \frac{8e^{4t}}{5} \end{bmatrix} dt \\ &= \begin{bmatrix} e^{2t} & -\frac{e^{-3t}}{4} \\ e^{2t} & e^{-3t} \end{bmatrix} \begin{bmatrix} -\frac{e^{-4t}}{5} + \frac{2e^{-t}}{5} \\ -\frac{4e^t}{5} - \frac{2e^{4t}}{5} \end{bmatrix} \\ &= \begin{bmatrix} \frac{e^t}{2} \\ -e^{-2t} \end{bmatrix} \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} c_1 e^{2t} \\ c_1 e^{2t} \end{bmatrix} + \begin{bmatrix} -\frac{c_2 e^{-3t}}{4} \\ c_2 e^{-3t} \end{bmatrix} + \begin{bmatrix} \frac{e^t}{2} \\ -e^{-2t} \end{bmatrix} \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(4c_1 e^{5t} + 2e^{4t} - c_2)e^{-3t}}{4} \\ (c_1 e^{5t} - e^t + c_2)e^{-3t} \end{bmatrix}$$

18.4.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = x_1(t) + x_2(t) + \frac{1}{(e^t)^2}, x_2'(t) = 4x_1(t) - 2x_2(t) - 2e^t \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & 1 \\ 4 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} \frac{x_1(t)(e^t)^2 + x_2(t)(e^t)^2 + 1}{(e^t)^2} - x_1(t) - x_2(t) \\ -2e^t \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & 1 \\ 4 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ -2e^t \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} 0 \\ -2e^t \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 1 & 1 \\ 4 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-3, \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix} \right], \left[2, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-3, \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{-3t} \cdot \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[2, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}_{\rightarrow 2} = e^{2t} \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- General solution of the system of ODEs can be written in terms of the particular solution $\underline{x}_{\rightarrow p}$

$$\underline{x}_{\rightarrow}(t) = c_1 \underline{x}_{\rightarrow 1} + c_2 \underline{x}_{\rightarrow 2} + \underline{x}_{\rightarrow p}(t)$$

- Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} -\frac{e^{-3t}}{4} & e^{2t} \\ e^{-3t} & e^{2t} \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix

$$\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} -\frac{e^{-3t}}{4} & e^{2t} \\ e^{-3t} & e^{2t} \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} -\frac{1}{4} & 1 \\ 1 & 1 \end{bmatrix}}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} \frac{(4e^{5t}+1)e^{-3t}}{5} & \frac{(e^{5t}-1)e^{-3t}}{5} \\ \frac{4(e^{5t}-1)e^{-3t}}{5} & \frac{(e^{5t}+4)e^{-3t}}{5} \end{bmatrix}$$

- Find a particular solution of the system of ODEs using variation of parameters

- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\underline{x}_{\rightarrow p}(t) = \Phi(t) \cdot \vec{v}(t)$$

- Take the derivative of the particular solution

$$\underline{x}_{\rightarrow p}'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t)$.

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- Plug $\vec{v}(t)$ into the equation for the particular solution

$$\vec{x}_{\text{part}}(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- Plug in the fundamental matrix and the forcing function and compute

$$\vec{x}_{\text{part}}(t) = \begin{bmatrix} -\frac{(4e^{5t}-5e^{4t}+1)e^{-3t}}{10} \\ -\frac{2(e^{5t}-1)e^{-3t}}{5} \end{bmatrix}$$

- Plug particular solution back into general solution

$$\vec{x}(t) = c_1 \vec{x}_1 + c_2 \vec{x}_2 + \begin{bmatrix} -\frac{(4e^{5t}-5e^{4t}+1)e^{-3t}}{10} \\ -\frac{2(e^{5t}-1)e^{-3t}}{5} \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(20c_2e^{5t}-8e^{5t}+10e^{4t}-5c_1-2)e^{-3t}}{20} \\ \frac{(5c_2e^{5t}-2e^{5t}+5c_1+2)e^{-3t}}{5} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{(20c_2e^{5t}-8e^{5t}+10e^{4t}-5c_1-2)e^{-3t}}{20}, x_2(t) = \frac{(5c_2e^{5t}-2e^{5t}+5c_1+2)e^{-3t}}{5} \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 45

```
dsolve([diff(x__1(t),t)=1*x__1(t)+1*x__2(t)+exp(-2*t),diff(x__2(t),t)=4*x__1(t)-2*x__2(t)-2*exp(-2*t)],{x__1(t),x__2(t)},t)
```

$$\begin{aligned}x_1(t) &= c_2 e^{2t} - \frac{c_1 e^{-3t}}{4} + \frac{e^t}{2} \\x_2(t) &= c_2 e^{2t} + c_1 e^{-3t} - e^{-2t}\end{aligned}$$

✓ Solution by Mathematica

Time used: 0.585 (sec). Leaf size: 84

```
DSolve[{x1'[t]==1*x1[t]+1*x2[t]+Exp[-2*t],x2'[t]==4*x1[t]-2*x2[t]-2*Exp[t]},{x1[t],x2[t]},t]
```

$$\begin{aligned}x_1(t) &\rightarrow \frac{e^t}{2} + \frac{1}{5}(c_1 - c_2)e^{-3t} + \frac{1}{5}(4c_1 + c_2)e^{2t} \\x_2(t) &\rightarrow \frac{1}{5}e^{-3t}(-5e^t + (4c_1 + c_2)e^{5t} - 4c_1 + 4c_2)\end{aligned}$$

18.5 problem 5

- 18.5.1 Solution using Matrix exponential method 3837
- 18.5.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3839
- 18.5.3 Maple step by step solution 3844

Internal problem ID [782]

Internal file name [OUTPUT/782_Sunday_June_05_2022_01_49_34_AM_87329820/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 5.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 4x_1(t) - 2x_2(t) + \frac{1}{t^3} \\x_2'(t) &= 8x_1(t) - 4x_2(t) - \frac{1}{t^2}\end{aligned}$$

18.5.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} \frac{1}{t^3} \\ -\frac{1}{t^2} \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} 4t + 1 & -2t \\ 8t & 1 - 4t \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} 4t + 1 & -2t \\ 8t & 1 - 4t \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} (4t + 1)c_1 - 2tc_2 \\ 8tc_1 + (1 - 4t)c_2 \end{bmatrix} \\ &= \begin{bmatrix} (4c_1 - 2c_2)t + c_1 \\ (8c_1 - 4c_2)t + c_2 \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} 1 - 4t & 2t \\ -8t & 4t + 1 \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} 4t + 1 & -2t \\ 8t & 1 - 4t \end{bmatrix} \int \begin{bmatrix} 1 - 4t & 2t \\ -8t & 4t + 1 \end{bmatrix} \begin{bmatrix} \frac{1}{t^3} \\ -\frac{1}{t^2} \end{bmatrix} dt \\ &= \begin{bmatrix} 4t + 1 & -2t \\ 8t & 1 - 4t \end{bmatrix} \begin{bmatrix} \frac{4}{t} - \frac{1}{2t^2} - 2 \ln(t) \\ \frac{9}{t} - 4 \ln(t) \end{bmatrix} \\ &= \begin{bmatrix} \frac{-2t^2 \ln(t) - 2(t - \frac{1}{2})^2}{t^2} \\ \frac{-4 \ln(t)t - 4t + 5}{t} \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} (4c_1 - 2c_2)t + c_1 + \frac{-2t^2 \ln(t) - 2(t - \frac{1}{2})^2}{t^2} \\ (8c_1 - 4c_2)t + c_2 + \frac{-4 \ln(t)t - 4t + 5}{t} \end{bmatrix}\end{aligned}$$

18.5.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} \frac{1}{t^3} \\ -\frac{1}{t^2} \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 4 - \lambda & -2 \\ 8 & -4 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 0$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
0	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = 0$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} - (0) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 4 & -2 & 0 \\ 8 & -4 & 0 \end{array} \right]$$

$$R_2 = R_2 - 2R_1 \implies \left[\begin{array}{cc|c} 4 & -2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 4 & -2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{t}{2}\}$

Hence the solution is

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{\sqrt{5}} \\ \frac{2}{\sqrt{5}} \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
0	2	1	Yes	$\begin{bmatrix} \frac{1}{\sqrt{5}} \\ \frac{2}{\sqrt{5}} \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue 0 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

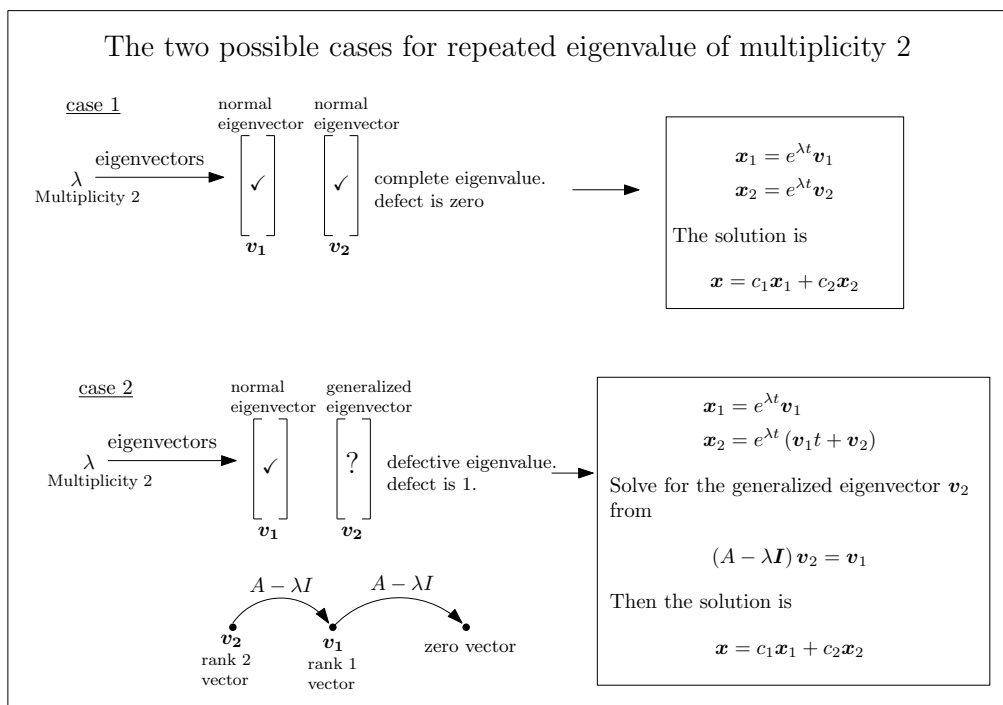


Figure 538: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} - (0) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} 1 \\ \frac{7}{4} \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue 0. Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} 1 \\ &= \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} t + \begin{bmatrix} 1 \\ \frac{7}{4} \end{bmatrix} \right) 1 \\ &= \begin{bmatrix} \frac{t}{2} + 1 \\ t + \frac{7}{4} \end{bmatrix}\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} + c_2 \begin{bmatrix} \frac{t}{2} + 1 \\ t + \frac{7}{4} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} \frac{1}{2} & \frac{t}{2} + 1 \\ 1 & t + \frac{7}{4} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -8t - 14 & 8 + 4t \\ 8 & -4 \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{1}{2} & \frac{t}{2} + 1 \\ 1 & t + \frac{7}{4} \end{bmatrix} \int \begin{bmatrix} -8t - 14 & 8 + 4t \\ 8 & -4 \end{bmatrix} \begin{bmatrix} \frac{1}{t^3} \\ -\frac{1}{t^2} \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{1}{2} & \frac{t}{2} + 1 \\ 1 & t + \frac{7}{4} \end{bmatrix} \int \begin{bmatrix} \frac{-4t^2 - 16t - 14}{t^3} \\ \frac{8 + 4t}{t^3} \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{1}{2} & \frac{t}{2} + 1 \\ 1 & t + \frac{7}{4} \end{bmatrix} \begin{bmatrix} \frac{16}{t} + \frac{7}{t^2} - 4 \ln(t) \\ \frac{-4t - 4}{t^2} \end{bmatrix} \\ &= \begin{bmatrix} \frac{-2t^2 \ln(t) - 2(t - \frac{1}{2})^2}{t^2} \\ \frac{-4 \ln(t)t - 4t + 5}{t} \end{bmatrix} \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} \frac{c_1}{2} \\ c_1 \end{bmatrix} + \begin{bmatrix} c_2 \left(\frac{t}{2} + 1\right) \\ c_2 \left(t + \frac{7}{4}\right) \end{bmatrix} + \begin{bmatrix} \frac{-2t^2 \ln(t) - 2(t - \frac{1}{2})^2}{t^2} \\ \frac{-4 \ln(t)t - 4t + 5}{t} \end{bmatrix} \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{-4t^2 \ln(t) - 1 + c_2 t^3 + (c_1 + 2c_2 - 4)t^2 + 4t}{2t^2} \\ c_1 + c_2 \left(t + \frac{7}{4}\right) + \frac{-4 \ln(t)t - 4t + 5}{t} \end{bmatrix}$$

18.5.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 4x_1(t) - 2x_2(t) + \frac{1}{t^3}, x_2'(t) = 8x_1(t) - 4x_2(t) - \frac{1}{t^2}]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} \frac{4x_1(t)t^3 - 2x_2(t)t^3 + 1}{t^3} - 4x_1(t) + 2x_2(t) \\ \frac{8x_1(t)t^2 - 4x_2(t)t^2 - 1}{t^2} - 8x_1(t) + 4x_2(t) \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 4 & -2 \\ 8 & -4 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[0, \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} \right], \left[0, \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[0, \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[0, \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\vec{x}_2 = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- General solution to the system of ODEs

$$\vec{x} = c_1 \vec{x}_1 + c_2 \vec{x}_2$$

- Substitute solutions into the general solution

$$\vec{x} = \begin{bmatrix} \frac{c_1}{2} \\ c_1 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{c_1}{2} \\ c_1 \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = \frac{c_1}{2}, x_2(t) = c_1\}$$

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 47

```
dsolve([diff(x__1(t),t)=4*x__1(t)-2*x__2(t)+1/(t^3),diff(x__2(t),t)=8*x__1(t)-4*x__2(t)-1/(t
```

$$x_1(t) = -\frac{1}{2t^2} + \frac{2}{t} - 2 \ln(t) + c_1 t + c_2$$

$$x_2(t) = 2c_1 t - 4 \ln(t) - \frac{c_1}{2} + 2c_2 + \frac{5}{t}$$

✓ Solution by Mathematica

Time used: 0.006 (sec). Leaf size: 61

```
DSolve[{x1'[t]==4*x1[t]-2*x2[t]+1/(t^3),x2'[t]==8*x1[t]-4*x2[t]-1/(t^2)},{x1[t],x2[t]},t,Inc
```

$$x1(t) \rightarrow -\frac{1}{2t^2} + \frac{2}{t} - 2 \log(t) + 4c_1 t - 2c_2 t - 2 + c_1$$

$$x2(t) \rightarrow \frac{5}{t} - 4 \log(t) + 8c_1 t - 4c_2 t - 4 + c_2$$

18.6 problem 6

- 18.6.1 Solution using Matrix exponential method 3847
- 18.6.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3849
- 18.6.3 Maple step by step solution 3854

Internal problem ID [783]

Internal file name [OUTPUT/783_Sunday_June_05_2022_01_49_35_AM_71046668/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 6.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -4x_1(t) + 2x_2(t) + \frac{1}{t} \\x_2'(t) &= 2x_1(t) - x_2(t) + \frac{2}{t} + 4\end{aligned}$$

18.6.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -4 & 2 \\ 2 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} \frac{1}{t} \\ \frac{2}{t} + 4 \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{4e^{-5t}}{5} + \frac{1}{5} & \frac{2}{5} - \frac{2e^{-5t}}{5} \\ \frac{2}{5} - \frac{2e^{-5t}}{5} & \frac{e^{-5t}}{5} + \frac{4}{5} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \frac{4e^{-5t}}{5} + \frac{1}{5} & \frac{2}{5} - \frac{2e^{-5t}}{5} \\ \frac{2}{5} - \frac{2e^{-5t}}{5} & \frac{e^{-5t}}{5} + \frac{4}{5} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} \left(\frac{4e^{-5t}}{5} + \frac{1}{5}\right) c_1 + \left(\frac{2}{5} - \frac{2e^{-5t}}{5}\right) c_2 \\ \left(\frac{2}{5} - \frac{2e^{-5t}}{5}\right) c_1 + \left(\frac{e^{-5t}}{5} + \frac{4}{5}\right) c_2 \end{bmatrix} \\ &= \begin{bmatrix} \frac{(4c_1 - 2c_2)e^{-5t}}{5} + \frac{c_1}{5} + \frac{2c_2}{5} \\ \frac{(-2c_1 + c_2)e^{-5t}}{5} + \frac{2c_1}{5} + \frac{4c_2}{5} \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \frac{4e^{5t}}{5} + \frac{1}{5} & -\frac{2e^{5t}}{5} + \frac{2}{5} \\ -\frac{2e^{5t}}{5} + \frac{2}{5} & \frac{e^{5t}}{5} + \frac{4}{5} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{4e^{-5t}}{5} + \frac{1}{5} & \frac{2}{5} - \frac{2e^{-5t}}{5} \\ \frac{2}{5} - \frac{2e^{-5t}}{5} & \frac{e^{-5t}}{5} + \frac{4}{5} \end{bmatrix} \int \begin{bmatrix} \frac{4e^{5t}}{5} + \frac{1}{5} & -\frac{2e^{5t}}{5} + \frac{2}{5} \\ -\frac{2e^{5t}}{5} + \frac{2}{5} & \frac{e^{5t}}{5} + \frac{4}{5} \end{bmatrix} \begin{bmatrix} \frac{1}{t} \\ \frac{2}{t} + 4 \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{4e^{-5t}}{5} + \frac{1}{5} & \frac{2}{5} - \frac{2e^{-5t}}{5} \\ \frac{2}{5} - \frac{2e^{-5t}}{5} & \frac{e^{-5t}}{5} + \frac{4}{5} \end{bmatrix} \begin{bmatrix} \frac{8t}{5} + \ln(5) + \ln(t) - \frac{8e^{5t}}{25} \\ 2 \ln(5) + 2 \ln(t) + \frac{16t}{5} + \frac{4e^{5t}}{25} \end{bmatrix} \\ &= \begin{bmatrix} -\frac{8}{25} + \frac{8t}{5} + \ln(5) + \ln(t) \\ \frac{16t}{5} + 2 \ln(5) + 2 \ln(t) + \frac{4}{25} \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} -\frac{8}{25} + \frac{2(2c_1 - c_2)e^{-5t}}{5} + \frac{8t}{5} + \frac{c_1}{5} + \frac{2c_2}{5} + \ln(5) + \ln(t) \\ \frac{(-2c_1 + c_2)e^{-5t}}{5} + \frac{2c_1}{5} + \frac{4c_2}{5} + \frac{16t}{5} + 2\ln(5) + 2\ln(t) + \frac{4}{25} \end{bmatrix}\end{aligned}$$

18.6.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -4 & 2 \\ 2 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} \frac{1}{t} \\ \frac{2}{t} + 4 \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -4 & 2 \\ 2 & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -4 - \lambda & 2 \\ 2 & -1 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 5\lambda = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -5$$

$$\lambda_2 = 0$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
0	1	real eigenvalue
-5	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -5$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -4 & 2 \\ 2 & -1 \end{bmatrix} - (-5) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & 2 \\ 2 & 4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 & 2 & 0 \\ 2 & 4 & 0 \end{array} \right]$$

$$R_2 = R_2 - 2R_1 \implies \left[\begin{array}{cc|c} 1 & 2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & 2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -2t\}$

Hence the solution is

$$\begin{bmatrix} -2t \\ t \end{bmatrix} = \begin{bmatrix} -2t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -2t \\ t \end{bmatrix} = t \begin{bmatrix} -2 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -2 \\ 1 \end{bmatrix} = \begin{bmatrix} -2 \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 0$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -4 & 2 \\ 2 & -1 \end{bmatrix} - (0) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -4 & 2 \\ 2 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -4 & 2 & 0 \\ 2 & -1 & 0 \end{array} \right]$$

$$R_2 = R_2 + \frac{R_1}{2} \implies \left[\begin{array}{cc|c} -4 & 2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -4 & 2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{t}{2}\}$

Hence the solution is

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ 2 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-5	1	1	No	$\begin{bmatrix} -2 \\ 1 \end{bmatrix}$
0	1	1	No	$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -5 is real and distinct then the

corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{-5t} \\ &= \begin{bmatrix} -2 \\ 1 \end{bmatrix} e^{-5t}\end{aligned}$$

Since eigenvalue 0 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^0 \\ &= \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} e^0\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} -2e^{-5t} \\ e^{-5t} \end{bmatrix} + c_2 \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} -2e^{-5t} & \frac{1}{2} \\ e^{-5t} & 1 \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -\frac{2e^{5t}}{5} & \frac{e^{5t}}{5} \\ \frac{2}{5} & \frac{4}{5} \end{bmatrix}$$

Hence

$$\begin{aligned}
 \vec{x}_p(t) &= \begin{bmatrix} -2e^{-5t} & \frac{1}{2} \\ e^{-5t} & 1 \end{bmatrix} \int \begin{bmatrix} -\frac{2e^{5t}}{5} & \frac{e^{5t}}{5} \\ \frac{2}{5} & \frac{4}{5} \end{bmatrix} \begin{bmatrix} \frac{1}{t} \\ \frac{2}{t} + 4 \end{bmatrix} dt \\
 &= \begin{bmatrix} -2e^{-5t} & \frac{1}{2} \\ e^{-5t} & 1 \end{bmatrix} \int \begin{bmatrix} \frac{4e^{5t}}{5} \\ \frac{2}{t} + \frac{16}{5} \end{bmatrix} dt \\
 &= \begin{bmatrix} -2e^{-5t} & \frac{1}{2} \\ e^{-5t} & 1 \end{bmatrix} \begin{bmatrix} \frac{4e^{5t}}{25} \\ \frac{16t}{5} + 2 \ln(t) \end{bmatrix} \\
 &= \begin{bmatrix} -\frac{8}{25} + \frac{8t}{5} + \ln(t) \\ \frac{4}{25} + \frac{16t}{5} + 2 \ln(t) \end{bmatrix}
 \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned}
 \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\
 \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} -2c_1e^{-5t} \\ c_1e^{-5t} \end{bmatrix} + \begin{bmatrix} \frac{c_2}{2} \\ c_2 \end{bmatrix} + \begin{bmatrix} -\frac{8}{25} + \frac{8t}{5} + \ln(t) \\ \frac{4}{25} + \frac{16t}{5} + 2 \ln(t) \end{bmatrix}
 \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} -2c_1e^{-5t} + \frac{c_2}{2} - \frac{8}{25} + \frac{8t}{5} + \ln(t) \\ c_1e^{-5t} + c_2 + \frac{4}{25} + \frac{16t}{5} + 2 \ln(t) \end{bmatrix}$$

18.6.3 Maple step by step solution

Let's solve

$$[x_1'(t) = -4x_1(t) + 2x_2(t) + \frac{1}{t}, x_2'(t) = 2x_1(t) - x_2(t) + \frac{2}{t} + 4]$$

- Define vector

$$\underline{x} \rightarrow(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x} \rightarrow'(t) = \begin{bmatrix} -4 & 2 \\ 2 & -1 \end{bmatrix} \cdot \underline{x} \rightarrow(t) + \begin{bmatrix} -\frac{4x_1(t)t - 2x_2(t)t - 1}{t} + 4x_1(t) - 2x_2(t) \\ \frac{2x_1(t)t - x_2(t)t + 4t + 2}{t} - 2x_1(t) + x_2(t) \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} -4 & 2 \\ 2 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -4 & 2 \\ 2 & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow}{}'(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-5, \begin{bmatrix} -2 \\ 1 \end{bmatrix} \right], \left[0, \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-5, \begin{bmatrix} -2 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{-5t} \cdot \begin{bmatrix} -2 \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[0, \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_2 = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1 + c_2 \underline{x}^{\rightarrow}_2$$

- Substitute solutions into the general solution

$$\vec{x} = c_1 e^{-5t} \cdot \begin{bmatrix} -2 \\ 1 \end{bmatrix} + \begin{bmatrix} \frac{c_2}{2} \\ c_2 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} -2c_1 e^{-5t} + \frac{c_2}{2} \\ c_1 e^{-5t} + c_2 \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = -2c_1 e^{-5t} + \frac{c_2}{2}, x_2(t) = c_1 e^{-5t} + c_2\}$$

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 45

```
dsolve([diff(x__1(t),t)=-4*x__1(t)+2*x__2(t)+1/t,diff(x__2(t),t)=2*x__1(t)-1*x__2(t)+2/t+4],
```

$$x_1(t) = \ln(-5t) - \frac{c_1 e^{-5t}}{5} + \frac{8t}{5} + c_2$$

$$x_2(t) = \frac{c_1 e^{-5t}}{10} + 2 \ln(-5t) + 2c_2 + \frac{16t}{5} + \frac{4}{5}$$

✓ Solution by Mathematica

Time used: 0.02 (sec). Leaf size: 86

```
DSolve[{x1'[t]==-4*x1[t]+2*x2[t]+1/t,x2'[t]==2*x1[t]-1*x2[t]+2/t+4},{x1[t],x2[t]},t,IncludeS
```

$$x_1(t) \rightarrow \frac{1}{25} (40t + 25 \log(t) + 20c_1 e^{-5t} - 10c_2 e^{-5t} - 8 + 5c_1 + 10c_2)$$

$$x_2(t) \rightarrow \frac{1}{25} (80t + 50 \log(t) - 10c_1 e^{-5t} + 5c_2 e^{-5t} + 4 + 10c_1 + 20c_2)$$

18.7 problem 7

18.7.1 Solution using Matrix exponential method	3857
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18.7.3 Maple step by step solution	3864

Internal problem ID [784]

Internal file name [OUTPUT/784_Sunday_June_05_2022_01_49_37_AM_33602493/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 7.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) + x_2(t) + 2e^t \\x_2'(t) &= 4x_1(t) + x_2(t) - e^t\end{aligned}$$

18.7.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ 4 & 1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 2e^t \\ -e^t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{e^{-t}}{2} + \frac{e^{3t}}{2} & \frac{e^{3t}}{4} - \frac{e^{-t}}{4} \\ e^{3t} - e^{-t} & \frac{e^{-t}}{2} + \frac{e^{3t}}{2} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \frac{e^{-t}}{2} + \frac{e^{3t}}{2} & \frac{e^{3t}}{4} - \frac{e^{-t}}{4} \\ e^{3t} - e^{-t} & \frac{e^{-t}}{2} + \frac{e^{3t}}{2} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} \left(\frac{e^{-t}}{2} + \frac{e^{3t}}{2}\right) c_1 + \left(\frac{e^{3t}}{4} - \frac{e^{-t}}{4}\right) c_2 \\ (e^{3t} - e^{-t}) c_1 + \left(\frac{e^{-t}}{2} + \frac{e^{3t}}{2}\right) c_2 \end{bmatrix} \\ &= \begin{bmatrix} \frac{(2c_1 - c_2)e^{-t}}{4} + \frac{(c_1 + \frac{c_2}{2})e^{3t}}{2} \\ \frac{(-2c_1 + c_2)e^{-t}}{2} + \left(c_1 + \frac{c_2}{2}\right) e^{3t} \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \frac{(e^{4t}+1)e^{-3t}}{2} & -\frac{(e^{4t}-1)e^{-3t}}{4} \\ -(e^{4t}-1)e^{-3t} & \frac{(e^{4t}+1)e^{-3t}}{2} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{e^{-t}}{2} + \frac{e^{3t}}{2} & \frac{e^{3t}}{4} - \frac{e^{-t}}{4} \\ e^{3t} - e^{-t} & \frac{e^{-t}}{2} + \frac{e^{3t}}{2} \end{bmatrix} \int \begin{bmatrix} \frac{(e^{4t}+1)e^{-3t}}{2} & -\frac{(e^{4t}-1)e^{-3t}}{4} \\ -(e^{4t}-1)e^{-3t} & \frac{(e^{4t}+1)e^{-3t}}{2} \end{bmatrix} \begin{bmatrix} 2e^t \\ -e^t \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{e^{-t}}{2} + \frac{e^{3t}}{2} & \frac{e^{3t}}{4} - \frac{e^{-t}}{4} \\ e^{3t} - e^{-t} & \frac{e^{-t}}{2} + \frac{e^{3t}}{2} \end{bmatrix} \begin{bmatrix} \frac{5e^{2t}}{8} - \frac{3e^{-2t}}{8} \\ -\frac{5e^{2t}}{4} - \frac{3e^{-2t}}{4} \end{bmatrix} \\ &= \begin{bmatrix} \frac{e^t}{4} \\ -2e^t \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{(2c_1 - c_2)e^{-t}}{4} + \frac{(2c_1 + c_2)e^{3t}}{4} + \frac{e^t}{4} \\ \frac{(-2c_1 + c_2)e^{-t}}{2} + \frac{(2c_1 + c_2)e^{3t}}{2} - 2e^t \end{bmatrix}\end{aligned}$$

18.7.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ 4 & 1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 2e^t \\ -e^t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 1 & 1 \\ 4 & 1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 1 - \lambda & 1 \\ 4 & 1 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 2\lambda - 3 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 3$$

$$\lambda_2 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
3	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 1 \\ 4 & 1 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 & 1 \\ 4 & 2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 & 1 & 0 \\ 4 & 2 & 0 \end{array} \right]$$

$$R_2 = R_2 - 2R_1 \implies \left[\begin{array}{cc|c} 2 & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -\frac{t}{2}\}$

Hence the solution is

$$\begin{bmatrix} -\frac{t}{2} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{t}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{t}{2} \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix} = \begin{bmatrix} -1 \\ 2 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 3$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 1 \\ 4 & 1 \end{bmatrix} - (3) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -2 & 1 \\ 4 & -2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -2 & 1 & 0 \\ 4 & -2 & 0 \end{array} \right]$$

$$R_2 = R_2 + 2R_1 \implies \left[\begin{array}{cc|c} -2 & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -2 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{t}{2}\}$

Hence the solution is

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ 2 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
3	1	1	No	$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$
-1	1	1	No	$\begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue 3 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{3t} \\ &= \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} e^{3t}\end{aligned}$$

Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-t} \\ &= \begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix} e^{-t}\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{e^{3t}}{2} \\ e^{3t} \end{bmatrix} + c_2 \begin{bmatrix} -\frac{e^{-t}}{2} \\ e^{-t} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} \frac{e^{3t}}{2} & -\frac{e^{-t}}{2} \\ e^{3t} & e^{-t} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} e^{-3t} & \frac{e^{-3t}}{2} \\ -e^t & \frac{e^t}{2} \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{e^{3t}}{2} & -\frac{e^{-t}}{2} \\ e^{3t} & e^{-t} \end{bmatrix} \int \begin{bmatrix} e^{-3t} & \frac{e^{-3t}}{2} \\ -e^t & \frac{e^t}{2} \end{bmatrix} \begin{bmatrix} 2e^t \\ -e^t \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{e^{3t}}{2} & -\frac{e^{-t}}{2} \\ e^{3t} & e^{-t} \end{bmatrix} \int \begin{bmatrix} \frac{3e^{-2t}}{2} \\ -\frac{5e^{2t}}{2} \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{e^{3t}}{2} & -\frac{e^{-t}}{2} \\ e^{3t} & e^{-t} \end{bmatrix} \begin{bmatrix} -\frac{3e^{-2t}}{4} \\ -\frac{5e^{2t}}{4} \end{bmatrix} \\ &= \begin{bmatrix} \frac{e^t}{4} \\ -2e^t \end{bmatrix} \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} \frac{c_1 e^{3t}}{2} \\ c_1 e^{3t} \end{bmatrix} + \begin{bmatrix} -\frac{c_2 e^{-t}}{2} \\ c_2 e^{-t} \end{bmatrix} + \begin{bmatrix} \frac{e^t}{4} \\ -2e^t \end{bmatrix} \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{c_1 e^{3t}}{2} - \frac{c_2 e^{-t}}{2} + \frac{e^t}{4} \\ c_1 e^{3t} + c_2 e^{-t} - 2e^t \end{bmatrix}$$

18.7.3 Maple step by step solution

Let's solve

$$[x_1'(t) = x_1(t) + x_2(t) + 2e^t, x_2'(t) = 4x_1(t) + x_2(t) - e^t]$$

- Define vector

$$\vec{x}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & 1 \\ 4 & 1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 2e^t \\ -e^t \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & 1 \\ 4 & 1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 2e^t \\ -e^t \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} 2e^t \\ -e^t \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 1 & 1 \\ 4 & 1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-1, \begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix} \right], \left[3, \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-1, \begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}_{\rightarrow 1} = e^{-t} \cdot \begin{bmatrix} -\frac{1}{2} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[3, \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}_{\rightarrow 2} = e^{3t} \cdot \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

- General solution of the system of ODEs can be written in terms of the particular solution $\underline{x}_{\rightarrow p}$

$$\underline{x}_{\rightarrow}(t) = c_1 \underline{x}_{\rightarrow 1} + c_2 \underline{x}_{\rightarrow 2} + \underline{x}_{\rightarrow p}(t)$$

- Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} -\frac{e^{-t}}{2} & \frac{e^{3t}}{2} \\ e^{-t} & e^{3t} \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix

$$\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} -\frac{e^{-t}}{2} & \frac{e^{3t}}{2} \\ e^{-t} & e^{3t} \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} -\frac{1}{2} & \frac{1}{2} \\ 1 & 1 \end{bmatrix}}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} \frac{e^{-t}}{2} + \frac{e^{3t}}{2} & \frac{e^{3t}}{4} - \frac{e^{-t}}{4} \\ e^{3t} - e^{-t} & \frac{e^{-t}}{2} + \frac{e^{3t}}{2} \end{bmatrix}$$

- Find a particular solution of the system of ODEs using variation of parameters

- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\underline{x}_{\rightarrow p}(t) = \Phi(t) \cdot \vec{v}(t)$$

- Take the derivative of the particular solution

$$\underline{x}_{\rightarrow p}'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t)$

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- Plug $\vec{v}(t)$ into the equation for the particular solution

$$\underline{x}_p(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- Plug in the fundamental matrix and the forcing function and compute

$$\underline{x}_p(t) = \begin{bmatrix} \frac{e^t}{4} - \frac{5e^{-t}}{8} + \frac{3e^{3t}}{8} \\ \frac{3e^{3t}}{4} + \frac{5e^{-t}}{4} - 2e^t \end{bmatrix}$$

- Plug particular solution back into general solution

$$\underline{x}(t) = c_1 \underline{x}_1 + c_2 \underline{x}_2 + \begin{bmatrix} \frac{e^t}{4} - \frac{5e^{-t}}{8} + \frac{3e^{3t}}{8} \\ \frac{3e^{3t}}{4} + \frac{5e^{-t}}{4} - 2e^t \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(-4c_1-5)e^{-t}}{8} + \frac{(4c_2+3)e^{3t}}{8} + \frac{e^t}{4} \\ \frac{(5+4c_1)e^{-t}}{4} + \frac{(4c_2+3)e^{3t}}{4} - 2e^t \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{(-4c_1-5)e^{-t}}{8} + \frac{(4c_2+3)e^{3t}}{8} + \frac{e^t}{4}, x_2(t) = \frac{(5+4c_1)e^{-t}}{4} + \frac{(4c_2+3)e^{3t}}{4} - 2e^t \right\}$$

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 44

```
dsolve([diff(x__1(t),t)=1*x__1(t)+1*x__2(t)+2*exp(t),diff(x__2(t),t)=4*x__1(t)+1*x__2(t)-exp
```

$$x_1(t) = c_2 e^{3t} + e^{-t} c_1 + \frac{e^t}{4}$$

$$x_2(t) = 2c_2 e^{3t} - 2e^{-t} c_1 - 2e^t$$

✓ Solution by Mathematica

Time used: 0.004 (sec). Leaf size: 80

```
DSolve[{x1'[t]==1*x1[t]+1*x2[t]+2*Exp[t],x2'[t]==4*x1[t]+1*x2[t]-Exp[t]},{x1[t],x2[t]},t,Inc
```

$$x1(t) \rightarrow \frac{1}{4}e^{-t}(e^{2t} + (2c_1 + c_2)e^{4t} + 2c_1 - c_2)$$

$$x2(t) \rightarrow \frac{1}{2}e^{-t}(-4e^{2t} + (2c_1 + c_2)e^{4t} - 2c_1 + c_2)$$

18.8 problem 8

18.8.1 Solution using Matrix exponential method	3869
18.8.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3871
18.8.3 Maple step by step solution	3876

Internal problem ID [785]

Internal file name [OUTPUT/785_Sunday_June_05_2022_01_49_39_AM_83821617/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 8.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 2x_1(t) - x_2(t) + e^t \\x_2'(t) &= 3x_1(t) - 2x_2(t) - e^t\end{aligned}$$

18.8.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} e^t \\ -e^t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} \left(-\frac{e^{-t}}{2} + \frac{3e^t}{2}\right) c_1 + \left(-\frac{e^t}{2} + \frac{e^{-t}}{2}\right) c_2 \\ \left(\frac{3e^t}{2} - \frac{3e^{-t}}{2}\right) c_1 + \left(\frac{3e^{-t}}{2} - \frac{e^t}{2}\right) c_2 \end{bmatrix} \\ &= \begin{bmatrix} \frac{(c_2 - c_1)e^{-t}}{2} + \frac{3(-\frac{c_2}{3} + c_1)e^t}{2} \\ \frac{(-3c_1 + 3c_2)e^{-t}}{2} + \frac{3(-\frac{c_2}{3} + c_1)e^t}{2} \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \frac{3e^{-t}}{2} - \frac{e^t}{2} & \frac{e^t}{2} - \frac{e^{-t}}{2} \\ -\frac{3e^t}{2} + \frac{3e^{-t}}{2} & -\frac{e^{-t}}{2} + \frac{3e^t}{2} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix} \int \begin{bmatrix} \frac{3e^{-t}}{2} - \frac{e^t}{2} & \frac{e^t}{2} - \frac{e^{-t}}{2} \\ -\frac{3e^t}{2} + \frac{3e^{-t}}{2} & -\frac{e^{-t}}{2} + \frac{3e^t}{2} \end{bmatrix} \begin{bmatrix} e^t \\ -e^t \end{bmatrix} dt \\ &= \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix} \begin{bmatrix} 2t - \frac{e^{2t}}{2} \\ 2t - \frac{3e^{2t}}{2} \end{bmatrix} \\ &= \begin{bmatrix} e^t \left(-\frac{1}{2} + 2t\right) \\ e^t \left(-\frac{3}{2} + 2t\right) \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{(c_2 - c_1)e^{-t}}{2} + 2\left(t + \frac{3c_1}{4} - \frac{c_2}{4} - \frac{1}{4}\right)e^t \\ \frac{(-3c_1 + 3c_2)e^{-t}}{2} + 2\left(t + \frac{3c_1}{4} - \frac{c_2}{4} - \frac{3}{4}\right)e^t \end{bmatrix}\end{aligned}$$

18.8.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} e^t \\ -e^t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 2 - \lambda & -1 \\ 3 & -2 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1$$

$$\lambda_2 = 1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
1	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 3 & -1 \\ 3 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 3 & -1 & 0 \\ 3 & -1 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} 3 & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 3 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{t}{3}\}$

Hence the solution is

$$\begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix} = \begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} = \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ 3 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} - (1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & -1 \\ 3 & -3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 & -1 & 0 \\ 3 & -3 & 0 \end{array} \right]$$

$$R_2 = R_2 - 3R_1 \implies \left[\begin{array}{cc|c} 1 & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-1	1	1	No	$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$
1	1	1	No	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -1 is real and distinct then the

corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{-t} \\ &= \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} e^{-t}\end{aligned}$$

Since eigenvalue 1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^t \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^t\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{e^{-t}}{3} \\ e^{-t} \end{bmatrix} + c_2 \begin{bmatrix} e^t \\ e^t \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} \frac{e^{-t}}{3} & e^t \\ e^{-t} & e^t \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -\frac{3e^t}{2} & \frac{3e^t}{2} \\ \frac{3e^{-t}}{2} & -\frac{e^{-t}}{2} \end{bmatrix}$$

Hence

$$\begin{aligned}
 \vec{x}_p(t) &= \begin{bmatrix} \frac{e^{-t}}{3} & e^t \\ e^{-t} & e^t \end{bmatrix} \int \begin{bmatrix} -\frac{3e^t}{2} & \frac{3e^t}{2} \\ \frac{3e^{-t}}{2} & -\frac{e^{-t}}{2} \end{bmatrix} \begin{bmatrix} e^t \\ -e^t \end{bmatrix} dt \\
 &= \begin{bmatrix} \frac{e^{-t}}{3} & e^t \\ e^{-t} & e^t \end{bmatrix} \int \begin{bmatrix} -3e^{2t} \\ 2 \end{bmatrix} dt \\
 &= \begin{bmatrix} \frac{e^{-t}}{3} & e^t \\ e^{-t} & e^t \end{bmatrix} \begin{bmatrix} -\frac{3e^{2t}}{2} \\ 2t \end{bmatrix} \\
 &= \begin{bmatrix} e^t(-\frac{1}{2} + 2t) \\ e^t(-\frac{3}{2} + 2t) \end{bmatrix}
 \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned}
 \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\
 \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} \frac{c_1 e^{-t}}{3} \\ c_1 e^{-t} \end{bmatrix} + \begin{bmatrix} c_2 e^t \\ c_2 e^t \end{bmatrix} + \begin{bmatrix} e^t(-\frac{1}{2} + 2t) \\ e^t(-\frac{3}{2} + 2t) \end{bmatrix}
 \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{c_1 e^{-t}}{3} + 2(t + \frac{c_2}{2} - \frac{1}{4}) e^t \\ c_1 e^{-t} + 2(t + \frac{c_2}{2} - \frac{3}{4}) e^t \end{bmatrix}$$

18.8.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 2x_1(t) - x_2(t) + e^t, x_2'(t) = 3x_1(t) - 2x_2(t) - e^t]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} e^t \\ -e^t \end{bmatrix}$$

- System to solve

$$\vec{x}'(t) = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \cdot \vec{x}(t) + \begin{bmatrix} e^t \\ -e^t \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} e^t \\ -e^t \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\vec{x}'(t) = A \cdot \vec{x}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-1, \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} \right], \left[1, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-1, \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\vec{x}_{-1} = e^{-t} \cdot \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[1, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\vec{x}_{-2} = e^t \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- General solution of the system of ODEs can be written in terms of the particular solution $\vec{x}_p(t)$

$$\vec{x}(t) = c_1 \vec{x}_1 + c_2 \vec{x}_2 + \vec{x}_p(t)$$

□ Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} \frac{e^{-t}}{3} & e^t \\ e^{-t} & e^t \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix

$$\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} \frac{e^{-t}}{3} & e^t \\ e^{-t} & e^t \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} \frac{1}{3} & 1 \\ 1 & 1 \end{bmatrix}}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix}$$

□ Find a particular solution of the system of ODEs using variation of parameters

- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\vec{x}_p(t) = \Phi(t) \cdot \vec{v}(t)$$

- Take the derivative of the particular solution

$$\vec{x}_p'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t)$

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- o Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- o Plug $\vec{v}(t)$ into the equation for the particular solution

$$\vec{x}_{\rightarrow p}(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- o Plug in the fundamental matrix and the forcing function and compute

$$\vec{x}_{\rightarrow p}(t) = \begin{bmatrix} \frac{e^{-t}}{2} - \frac{e^t}{2} + 2t e^t \\ 2t e^t + \frac{3e^{-t}}{2} - \frac{3e^t}{2} \end{bmatrix}$$

- Plug particular solution back into general solution

$$\vec{x}_{\rightarrow}(t) = c_1 \vec{x}_{\rightarrow 1} + c_2 \vec{x}_{\rightarrow 2} + \begin{bmatrix} \frac{e^{-t}}{2} - \frac{e^t}{2} + 2t e^t \\ 2t e^t + \frac{3e^{-t}}{2} - \frac{3e^t}{2} \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(2c_1+3)e^{-t}}{6} + 2\left(t + \frac{c_2}{2} - \frac{1}{4}\right) e^t \\ \frac{(2c_1+3)e^{-t}}{2} + 2\left(t + \frac{c_2}{2} - \frac{3}{4}\right) e^t \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{(2c_1+3)e^{-t}}{6} + 2\left(t + \frac{c_2}{2} - \frac{1}{4}\right) e^t, x_2(t) = \frac{(2c_1+3)e^{-t}}{2} + 2\left(t + \frac{c_2}{2} - \frac{3}{4}\right) e^t \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 45

```
dsolve([diff(x__1(t),t)=2*x__1(t)-1*x__2(t)+exp(t),diff(x__2(t),t)=3*x__1(t)-2*x__2(t)-exp(t)
```

$$\begin{aligned} x_1(t) &= c_2 e^t + e^{-t} c_1 + 2 e^t t \\ x_2(t) &= c_2 e^t + 3 e^{-t} c_1 + 2 e^t t - e^t \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.009 (sec). Leaf size: 80

```
DSolve[{x1'[t]==2*x1[t]-1*x2[t]+Exp[t],x2'[t]==3*x1[t]-2*x2[t]-Exp[t]},{x1[t],x2[t]},t,Inclu
```

$$x1(t) \rightarrow \frac{1}{2}e^{-t}(e^{2t}(4t - 1 + 3c_1 - c_2) - c_1 + c_2)$$

$$x2(t) \rightarrow \frac{1}{2}e^{-t}(e^{2t}(4t - 3 + 3c_1 - c_2) - 3c_1 + 3c_2)$$

18.9 problem 9

- 18.9.1 Solution using Matrix exponential method 3881
- 18.9.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3883
- 18.9.3 Maple step by step solution 3888

Internal problem ID [786]

Internal file name [OUTPUT/786_Sunday_June_05_2022_01_49_41_AM_45101061/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 9.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -\frac{5x_1(t)}{4} + \frac{3x_2(t)}{4} + 2t \\x_2'(t) &= \frac{3x_1(t)}{4} - \frac{5x_2(t)}{4} + e^t\end{aligned}$$

18.9.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{5}{4} & \frac{3}{4} \\ \frac{3}{4} & -\frac{5}{4} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 2t \\ e^t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2} & \frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2} \\ \frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2} & \frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2} & \frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2} \\ \frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2} & \frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} \left(\frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2}\right) c_1 + \left(\frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2}\right) c_2 \\ \left(\frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2}\right) c_1 + \left(\frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2}\right) c_2 \end{bmatrix} \\ &= \begin{bmatrix} \frac{(c_1+c_2)e^{-\frac{t}{2}}}{2} + \frac{(-c_2+c_1)e^{-2t}}{2} \\ \frac{(c_1+c_2)e^{-\frac{t}{2}}}{2} - \frac{(-c_2+c_1)e^{-2t}}{2} \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \frac{\left(e^{\frac{3t}{2}}+1\right)e^{\frac{t}{2}}}{2} & -\frac{\left(e^{\frac{3t}{2}}-1\right)e^{\frac{t}{2}}}{2} \\ -\frac{\left(e^{\frac{3t}{2}}-1\right)e^{\frac{t}{2}}}{2} & \frac{\left(e^{\frac{3t}{2}}+1\right)e^{\frac{t}{2}}}{2} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned}
 \vec{x}_p(t) &= \begin{bmatrix} \frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2} & \frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2} \\ \frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2} & \frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2} \end{bmatrix} \int \begin{bmatrix} \frac{(e^{\frac{3t}{2}}+1)e^{\frac{t}{2}}}{2} & -\frac{(e^{\frac{3t}{2}}-1)e^{\frac{t}{2}}}{2} \\ -\frac{(e^{\frac{3t}{2}}-1)e^{\frac{t}{2}}}{2} & \frac{(e^{\frac{3t}{2}}+1)e^{\frac{t}{2}}}{2} \end{bmatrix} \begin{bmatrix} 2t \\ e^t \end{bmatrix} dt \\
 &= \begin{bmatrix} \frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2} & \frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2} \\ \frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2} & \frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2} \end{bmatrix} \begin{bmatrix} 2(t-2)e^{\frac{t}{2}} + \frac{e^{\frac{3t}{2}}}{3} + \frac{(2t-1)e^{2t}}{4} - \frac{e^{3t}}{6} \\ 2(t-2)e^{\frac{t}{2}} + \frac{e^{\frac{3t}{2}}}{3} + \frac{(1-2t)e^{2t}}{4} + \frac{e^{3t}}{6} \end{bmatrix} \\
 &= \begin{bmatrix} \frac{e^t}{6} + \frac{5t}{2} - \frac{17}{4} \\ \frac{e^t}{2} + \frac{3t}{2} - \frac{15}{4} \end{bmatrix}
 \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}
 \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\
 &= \begin{bmatrix} \frac{e^{-2t}((c_1+c_2)e^{\frac{3t}{2}}+5e^{2t}t+c_1-c_2-\frac{17e^{2t}}{2}+\frac{e^{3t}}{3})}{2} \\ \frac{e^{-2t}((c_1+c_2)e^{\frac{3t}{2}}+3e^{2t}t-c_1+c_2-\frac{15e^{2t}}{2}+e^{3t})}{2} \end{bmatrix}
 \end{aligned}$$

18.9.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{5}{4} & \frac{3}{4} \\ \frac{3}{4} & -\frac{5}{4} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 2t \\ e^t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} -\frac{5}{4} & \frac{3}{4} \\ \frac{3}{4} & -\frac{5}{4} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} -\frac{5}{4} - \lambda & \frac{3}{4} \\ \frac{3}{4} & -\frac{5}{4} - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + \frac{5}{2}\lambda + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -2$$

$$\lambda_2 = -\frac{1}{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-2	1	real eigenvalue
$-\frac{1}{2}$	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{5}{4} & \frac{3}{4} \\ \frac{3}{4} & -\frac{5}{4} \end{bmatrix} - (-2) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} \frac{3}{4} & \frac{3}{4} \\ \frac{3}{4} & \frac{3}{4} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{3}{4} & \frac{3}{4} & 0 \\ \frac{3}{4} & \frac{3}{4} & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} \frac{3}{4} & \frac{3}{4} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} \frac{3}{4} & \frac{3}{4} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -t\}$

Hence the solution is

$$\begin{bmatrix} -t \\ t \end{bmatrix} = \begin{bmatrix} -t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -t \\ t \end{bmatrix} = t \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -t \\ t \end{bmatrix} = \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -\frac{1}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{5}{4} & \frac{3}{4} \\ \frac{3}{4} & -\frac{5}{4} \end{bmatrix} - \left(-\frac{1}{2}\right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -\frac{3}{4} & \frac{3}{4} \\ \frac{3}{4} & -\frac{3}{4} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -\frac{3}{4} & \frac{3}{4} & 0 \\ \frac{3}{4} & -\frac{3}{4} & 0 \end{array} \right]$$

$$R_2 = R_2 + R_1 \implies \left[\begin{array}{cc|c} -\frac{3}{4} & \frac{3}{4} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -\frac{3}{4} & \frac{3}{4} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-2	1	1	No	$\begin{bmatrix} -1 \\ 1 \end{bmatrix}$
$-\frac{1}{2}$	1	1	No	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{-2t} \\ &= \begin{bmatrix} -1 \\ 1 \end{bmatrix} e^{-2t}\end{aligned}$$

Since eigenvalue $-\frac{1}{2}$ is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-\frac{t}{2}} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{-\frac{t}{2}}\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} -e^{-2t} \\ e^{-2t} \end{bmatrix} + c_2 \begin{bmatrix} e^{-\frac{t}{2}} \\ e^{-\frac{t}{2}} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \dots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} -e^{-2t} & e^{-\frac{t}{2}} \\ e^{-2t} & e^{-\frac{t}{2}} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -\frac{e^{2t}}{2} & \frac{e^{2t}}{2} \\ \frac{e^{\frac{t}{2}}}{2} & \frac{e^{\frac{t}{2}}}{2} \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} -e^{-2t} & e^{-\frac{t}{2}} \\ e^{-2t} & e^{-\frac{t}{2}} \end{bmatrix} \int \begin{bmatrix} -\frac{e^{2t}}{2} & \frac{e^{2t}}{2} \\ \frac{e^{\frac{t}{2}}}{2} & \frac{e^{\frac{t}{2}}}{2} \end{bmatrix} \begin{bmatrix} 2t \\ e^t \end{bmatrix} dt \\ &= \begin{bmatrix} -e^{-2t} & e^{-\frac{t}{2}} \\ e^{-2t} & e^{-\frac{t}{2}} \end{bmatrix} \int \begin{bmatrix} -e^{2t}t + \frac{e^{3t}}{2} \\ te^{\frac{t}{2}} + \frac{e^{\frac{3t}{2}}}{2} \end{bmatrix} dt \\ &= \begin{bmatrix} -e^{-2t} & e^{-\frac{t}{2}} \\ e^{-2t} & e^{-\frac{t}{2}} \end{bmatrix} \begin{bmatrix} \frac{(1-2t)e^{2t}}{4} + \frac{e^{3t}}{6} \\ 2(t-2)e^{\frac{t}{2}} + \frac{e^{\frac{3t}{2}}}{3} \end{bmatrix} \\ &= \begin{bmatrix} \frac{e^t}{6} + \frac{5t}{2} - \frac{17}{4} \\ \frac{e^t}{2} + \frac{3t}{2} - \frac{15}{4} \end{bmatrix} \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} -c_1 e^{-2t} \\ c_1 e^{-2t} \end{bmatrix} + \begin{bmatrix} c_2 e^{-\frac{t}{2}} \\ c_2 e^{-\frac{t}{2}} \end{bmatrix} + \begin{bmatrix} \frac{e^t}{6} + \frac{5t}{2} - \frac{17}{4} \\ \frac{e^t}{2} + \frac{3t}{2} - \frac{15}{4} \end{bmatrix} \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(2e^{3t} + 30e^{2t}t - 51e^{2t} + 12c_2e^{\frac{3t}{2}} - 12c_1)e^{-2t}}{12} \\ \frac{(2e^{3t} + 6e^{2t}t - 15e^{2t} + 4c_2e^{\frac{3t}{2}} + 4c_1)e^{-2t}}{4} \end{bmatrix}$$

18.9.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = -\frac{5x_1(t)}{4} + \frac{3x_2(t)}{4} + 2t, x_2'(t) = \frac{3x_1(t)}{4} - \frac{5x_2(t)}{4} + e^t \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -\frac{5}{4} & \frac{3}{4} \\ \frac{3}{4} & -\frac{5}{4} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 2t \\ e^t \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -\frac{5}{4} & \frac{3}{4} \\ \frac{3}{4} & -\frac{5}{4} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 2t \\ e^t \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} 2t \\ e^t \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -\frac{5}{4} & \frac{3}{4} \\ \frac{3}{4} & -\frac{5}{4} \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-2, \begin{bmatrix} -1 \\ 1 \end{bmatrix} \right], \left[-\frac{1}{2}, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-2, \begin{bmatrix} -1 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{-2t} \cdot \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[-\frac{1}{2}, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}_{\rightarrow 2} = e^{-\frac{t}{2}} \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- General solution of the system of ODEs can be written in terms of the particular solution $\underline{x}_{\rightarrow p}$

$$\underline{x}_{\rightarrow}(t) = c_1 \underline{x}_{\rightarrow 1} + c_2 \underline{x}_{\rightarrow 2} + \underline{x}_{\rightarrow p}(t)$$

- Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} -e^{-2t} & e^{-\frac{t}{2}} \\ e^{-2t} & e^{-\frac{t}{2}} \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix

$$\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} -e^{-2t} & e^{-\frac{t}{2}} \\ e^{-2t} & e^{-\frac{t}{2}} \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} -1 & 1 \\ 1 & 1 \end{bmatrix}}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} \frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2} & \frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2} \\ \frac{e^{-\frac{t}{2}}}{2} - \frac{e^{-2t}}{2} & \frac{e^{-2t}}{2} + \frac{e^{-\frac{t}{2}}}{2} \end{bmatrix}$$

- Find a particular solution of the system of ODEs using variation of parameters

- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\underline{x}_{\rightarrow p}(t) = \Phi(t) \cdot \vec{v}(t)$$

- Take the derivative of the particular solution

$$\underline{x}_{\rightarrow p}'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- Plug $\vec{v}(t)$ into the equation for the particular solution

$$\vec{x}_{\text{---}p}(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- Plug in the fundamental matrix and the forcing function and compute

$$\vec{x}_{\text{---}p}(t) = \begin{bmatrix} \frac{(2e^{3t} + 30e^{2t}t - 51e^{2t} + 44e^{\frac{3t}{2}} + 5)e^{-2t}}{12} \\ \frac{(6e^{3t} + 18e^{2t}t - 45e^{2t} + 44e^{\frac{3t}{2}} - 5)e^{-2t}}{12} \end{bmatrix}$$

- Plug particular solution back into general solution

$$\vec{x}_{\text{---}}(t) = c_1 \vec{x}_{\text{---}1} + c_2 \vec{x}_{\text{---}2} + \begin{bmatrix} \frac{(2e^{3t} + 30e^{2t}t - 51e^{2t} + 44e^{\frac{3t}{2}} + 5)e^{-2t}}{12} \\ \frac{(6e^{3t} + 18e^{2t}t - 45e^{2t} + 44e^{\frac{3t}{2}} - 5)e^{-2t}}{12} \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(2e^{3t} + 30e^{2t}t - 51e^{2t} + 12c_2e^{\frac{3t}{2}} + 44e^{\frac{3t}{2}} - 12c_1 + 5)e^{-2t}}{12} \\ \frac{(6e^{3t} + 18e^{2t}t - 45e^{2t} + 12c_2e^{\frac{3t}{2}} + 44e^{\frac{3t}{2}} + 12c_1 - 5)e^{-2t}}{12} \end{bmatrix}$$

- Solution to the system of ODEs

$$\begin{cases} x_1(t) = \frac{(2e^{3t} + 30e^{2t}t - 51e^{2t} + 12c_2e^{\frac{3t}{2}} + 44e^{\frac{3t}{2}} - 12c_1 + 5)e^{-2t}}{12}, & x_2(t) = \frac{(6e^{3t} + 18e^{2t}t - 45e^{2t} + 12c_2e^{\frac{3t}{2}} + 44e^{\frac{3t}{2}} + 12c_1 - 5)e^{-2t}}{12} \end{cases}$$

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 51

```
dsolve([diff(x__1(t),t)=-5/4*x__1(t)+3/4*x__2(t)+2*t,diff(x__2(t),t)=3/4*x__1(t)-5/4*x__2(t)
```

$$\begin{aligned}x_1(t) &= c_2 e^{-2t} + c_1 e^{-\frac{t}{2}} - \frac{17}{4} + \frac{e^t}{6} + \frac{5t}{2} \\x_2(t) &= -c_2 e^{-2t} + c_1 e^{-\frac{t}{2}} + \frac{e^t}{2} - \frac{15}{4} + \frac{3t}{2}\end{aligned}$$

✓ Solution by Mathematica

Time used: 0.349 (sec). Leaf size: 101

```
DSolve[{x1'[t]==-5/4*x1[t]+3/4*x2[t]+2*t,x2'[t]==3/4*x1[t]-5/4*x2[t]+Exp[t]},{x1[t],x2[t]},t
```

$$\begin{aligned}x_1(t) &\rightarrow \frac{1}{12}(30t + 2e^t + 6(c_1 - c_2)e^{-2t} + 6(c_1 + c_2)e^{-t/2} - 51) \\x_2(t) &\rightarrow \frac{1}{4}e^{-2t}(3e^{2t}(2t - 5) + 2e^{3t} + 2(c_1 + c_2)e^{3t/2} - 2c_1 + 2c_2)\end{aligned}$$

18.10 problem 10

18.10.1 Solution using Matrix exponential method 3893

18.10.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3895

Internal problem ID [787]

Internal file name [OUTPUT/787_Sunday_June_05_2022_01_49_43_AM_10721479/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 10.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -3x_1(t) + \sqrt{2}x_2(t) + e^{-t} \\x_2'(t) &= \sqrt{2}x_1(t) - 2x_2(t) - e^{-t}\end{aligned}$$

18.10.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -3 & \sqrt{2} \\ \sqrt{2} & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} e^{-t} \\ -e^{-t} \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{2e^{-4t}}{3} + \frac{e^{-t}}{3} & \frac{\sqrt{2}(e^{-t}-e^{-4t})}{3} \\ \frac{\sqrt{2}(e^{-t}-e^{-4t})}{3} & \frac{e^{-4t}}{3} + \frac{2e^{-t}}{3} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \frac{2e^{-4t}}{3} + \frac{e^{-t}}{3} & \frac{\sqrt{2}(e^{-t}-e^{-4t})}{3} \\ \frac{\sqrt{2}(e^{-t}-e^{-4t})}{3} & \frac{e^{-4t}}{3} + \frac{2e^{-t}}{3} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} \left(\frac{2e^{-4t}}{3} + \frac{e^{-t}}{3}\right) c_1 + \frac{\sqrt{2}(e^{-t}-e^{-4t})c_2}{3} \\ \frac{\sqrt{2}(e^{-t}-e^{-4t})c_1}{3} + \left(\frac{e^{-4t}}{3} + \frac{2e^{-t}}{3}\right) c_2 \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \frac{(2e^{3t}+1)e^t}{3} & -\frac{\sqrt{2}e^t(e^{3t}-1)}{3} \\ -\frac{\sqrt{2}e^t(e^{3t}-1)}{3} & \frac{(e^{3t}+2)e^t}{3} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{2e^{-4t}}{3} + \frac{e^{-t}}{3} & \frac{\sqrt{2}(e^{-t}-e^{-4t})}{3} \\ \frac{\sqrt{2}(e^{-t}-e^{-4t})}{3} & \frac{e^{-4t}}{3} + \frac{2e^{-t}}{3} \end{bmatrix} \int \begin{bmatrix} \frac{(2e^{3t}+1)e^t}{3} & -\frac{\sqrt{2}e^t(e^{3t}-1)}{3} \\ -\frac{\sqrt{2}e^t(e^{3t}-1)}{3} & \frac{(e^{3t}+2)e^t}{3} \end{bmatrix} \begin{bmatrix} e^{-t} \\ -e^{-t} \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{2e^{-4t}}{3} + \frac{e^{-t}}{3} & \frac{\sqrt{2}(e^{-t}-e^{-4t})}{3} \\ \frac{\sqrt{2}(e^{-t}-e^{-4t})}{3} & \frac{e^{-4t}}{3} + \frac{2e^{-t}}{3} \end{bmatrix} \begin{bmatrix} \frac{(2+\sqrt{2})e^{3t}}{9} - \frac{t(\sqrt{2}-1)}{3} \\ \frac{(-1-\sqrt{2})e^{3t}}{9} + \frac{t(-2+\sqrt{2})}{3} \end{bmatrix} \\ &= \begin{bmatrix} \frac{((1-3t)\sqrt{2}+3t+2)e^{-t}}{9} \\ \frac{((-1+3t)\sqrt{2}-6t-1)e^{-t}}{9} \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{((-3t+3c_2+1)\sqrt{2}+3t+3c_1+2)e^{-t}}{9} + \frac{2(-\frac{c_2\sqrt{2}}{2}+c_1)e^{-4t}}{3} \\ \frac{((3t+3c_1-1)\sqrt{2}-6t+6c_2-1)e^{-t}}{9} - \frac{e^{-4t}(c_1\sqrt{2}-c_2)}{3} \end{bmatrix}\end{aligned}$$

18.10.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -3 & \sqrt{2} \\ \sqrt{2} & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} e^{-t} \\ -e^{-t} \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -3 & \sqrt{2} \\ \sqrt{2} & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -3 - \lambda & \sqrt{2} \\ \sqrt{2} & -2 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 5\lambda + 4 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -4$$

$$\lambda_2 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
-4	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -4$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -3 & \sqrt{2} \\ \sqrt{2} & -2 \end{bmatrix} - (-4) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & \sqrt{2} \\ \sqrt{2} & 2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 & \sqrt{2} & 0 \\ \sqrt{2} & 2 & 0 \end{array} \right]$$

$$R_2 = R_2 - \sqrt{2}R_1 \implies \left[\begin{array}{cc|c} 1 & \sqrt{2} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & \sqrt{2} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -\sqrt{2}t\}$

Hence the solution is

$$\begin{bmatrix} -\sqrt{2}t \\ t \end{bmatrix} = \begin{bmatrix} -\sqrt{2}t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\sqrt{2}t \\ t \end{bmatrix} = t \begin{bmatrix} -\sqrt{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\sqrt{2}t \\ t \end{bmatrix} = \begin{bmatrix} -\sqrt{2} \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -3 & \sqrt{2} \\ \sqrt{2} & -2 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -2 & \sqrt{2} \\ \sqrt{2} & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -2 & \sqrt{2} & 0 \\ \sqrt{2} & -1 & 0 \end{array} \right]$$

$$R_2 = R_2 + \frac{\sqrt{2}R_1}{2} \implies \left[\begin{array}{cc|c} -2 & \sqrt{2} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -2 & \sqrt{2} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\left\{v_1 = \frac{\sqrt{2}t}{2}\right\}$

Hence the solution is

$$\begin{bmatrix} \frac{\sqrt{2}t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{\sqrt{2}t}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{\sqrt{2}t}{2} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{\sqrt{2}}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{\sqrt{2}t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{\sqrt{2}}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{\sqrt{2}t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \sqrt{2} \\ 2 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-4	1	1	No	$\begin{bmatrix} -\sqrt{2} \\ 1 \end{bmatrix}$
-1	1	1	No	$\begin{bmatrix} \frac{\sqrt{2}}{2} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -4 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{-4t} \\ &= \begin{bmatrix} -\sqrt{2} \\ 1 \end{bmatrix} e^{-4t}\end{aligned}$$

Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-t} \\ &= \begin{bmatrix} \frac{\sqrt{2}}{2} \\ 1 \end{bmatrix} e^{-t}\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} -\sqrt{2} e^{-4t} \\ e^{-4t} \end{bmatrix} + c_2 \begin{bmatrix} \frac{\sqrt{2} e^{-t}}{2} \\ e^{-t} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} -\sqrt{2} e^{-4t} & \frac{\sqrt{2} e^{-t}}{2} \\ e^{-4t} & e^{-t} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -\frac{\sqrt{2}e^{4t}}{3} & \frac{e^{4t}}{3} \\ \frac{\sqrt{2}e^t}{3} & \frac{2e^t}{3} \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} -\sqrt{2}e^{-4t} & \frac{\sqrt{2}e^{-t}}{2} \\ e^{-4t} & e^{-t} \end{bmatrix} \int \begin{bmatrix} -\frac{\sqrt{2}e^{4t}}{3} & \frac{e^{4t}}{3} \\ \frac{\sqrt{2}e^t}{3} & \frac{2e^t}{3} \end{bmatrix} \begin{bmatrix} e^{-t} \\ -e^{-t} \end{bmatrix} dt \\ &= \begin{bmatrix} -\sqrt{2}e^{-4t} & \frac{\sqrt{2}e^{-t}}{2} \\ e^{-4t} & e^{-t} \end{bmatrix} \int \begin{bmatrix} -\frac{e^{3t}(1+\sqrt{2})}{3} \\ \frac{\sqrt{2}}{3} - \frac{2}{3} \end{bmatrix} dt \\ &= \begin{bmatrix} -\sqrt{2}e^{-4t} & \frac{\sqrt{2}e^{-t}}{2} \\ e^{-4t} & e^{-t} \end{bmatrix} \begin{bmatrix} -\frac{e^{3t}(1+\sqrt{2})}{9} \\ \frac{t(-2+\sqrt{2})}{3} \end{bmatrix} \\ &= \begin{bmatrix} \frac{((1-3t)\sqrt{2}+3t+2)e^{-t}}{9} \\ \frac{((-1+3t)\sqrt{2}-6t-1)e^{-t}}{9} \end{bmatrix} \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} -c_1\sqrt{2}e^{-4t} \\ c_1e^{-4t} \end{bmatrix} + \begin{bmatrix} \frac{c_2\sqrt{2}e^{-t}}{2} \\ c_2e^{-t} \end{bmatrix} + \begin{bmatrix} \frac{((1-3t)\sqrt{2}+3t+2)e^{-t}}{9} \\ \frac{((-1+3t)\sqrt{2}-6t-1)e^{-t}}{9} \end{bmatrix} \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{((-6t+9c_2+2)\sqrt{2}+6t+4)e^{-t}}{18} - c_1\sqrt{2}e^{-4t} \\ \frac{((-1+3t)\sqrt{2}-6t+9c_2-1)e^{-t}}{9} + c_1e^{-4t} \end{bmatrix}$$

✓ Solution by Maple

Time used: 0.063 (sec). Leaf size: 90

```
dsolve([diff(x__1(t),t)=-3*x__1(t)+sqrt(2)*x__2(t)+exp(-t),diff(x__2(t),t)=sqrt(2)*x__1(t)-2
```

$$x_1(t) = c_2 e^{-t} + e^{-4t} c_1 - \frac{t e^{-t} \sqrt{2}}{3} + \frac{t e^{-t}}{3}$$
$$x_2(t) = -\frac{2t e^{-t}}{3} - \frac{e^{-t}}{3} + \sqrt{2} e^{-t} c_2 + \frac{t e^{-t} \sqrt{2}}{3} - \frac{\sqrt{2} e^{-4t} c_1}{2} - \frac{\sqrt{2} e^{-t}}{3}$$

✓ Solution by Mathematica

Time used: 0.035 (sec). Leaf size: 128

```
DSolve[{x1'[t]==-3*x1[t]+Sqrt[2]*x2[t]+Exp[-t],x2'[t]==Sqrt[2]*x1[t]-2*x2[t]-Exp[-t]},{x1[t]
```

$$x_1(t) \rightarrow \frac{1}{9} e^{-4t} \left(e^{3t} \left(-3(\sqrt{2} - 1)t + \sqrt{2} + 2 + 3c_1 + 3\sqrt{2}c_2 \right) + 6c_1 - 3\sqrt{2}c_2 \right)$$
$$x_2(t) \rightarrow \frac{1}{9} e^{-4t} \left(e^{3t} \left(3(\sqrt{2} - 2)t - \sqrt{2} - 1 + 3\sqrt{2}c_1 + 6c_2 \right) - 3\sqrt{2}c_1 + 3c_2 \right)$$

18.11 problem 11

18.11.1 Solution using Matrix exponential method	3902
18.11.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3904
18.11.3 Maple step by step solution	3909

Internal problem ID [788]

Internal file name [OUTPUT/788_Sunday_June_05_2022_01_49_46_AM_47915974/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 11.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 2x_1(t) - 5x_2(t) \\x_2'(t) &= x_1(t) - 2x_2(t) + \cos(t)\end{aligned}$$

18.11.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 0 \\ \cos(t) \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} (\cos(t) + 2 \sin(t)) c_1 - 5 \sin(t) c_2 \\ \sin(t) c_1 + (\cos(t) - 2 \sin(t)) c_2 \end{bmatrix} \\ &= \begin{bmatrix} (2c_1 - 5c_2) \sin(t) + c_1 \cos(t) \\ (c_1 - 2c_2) \sin(t) + c_2 \cos(t) \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \cos(t) - 2 \sin(t) & 5 \sin(t) \\ -\sin(t) & \cos(t) + 2 \sin(t) \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \int \begin{bmatrix} \cos(t) - 2 \sin(t) & 5 \sin(t) \\ -\sin(t) & \cos(t) + 2 \sin(t) \end{bmatrix} \begin{bmatrix} 0 \\ \cos(t) \end{bmatrix} dt \\ &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \begin{bmatrix} \frac{5 \sin(t)^2}{2} \\ \frac{\sin(t) \cos(t)}{2} + \frac{t}{2} - \cos(t)^2 \end{bmatrix} \\ &= \begin{bmatrix} -\frac{5 \sin(t)(t-2)}{2} \\ \frac{(t-2) \cos(t)}{2} + \frac{(-2t+5) \sin(t)}{2} \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{(-5t+4c_1-10c_2+10)\sin(t)}{2} + c_1 \cos(t) \\ \frac{(-2t+2c_1-4c_2+5)\sin(t)}{2} + \frac{\cos(t)(t+2c_2-2)}{2} \end{bmatrix}\end{aligned}$$

18.11.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 0 \\ \cos(t) \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 2 - \lambda & -5 \\ 1 & -2 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-i$	1	complex eigenvalue
i	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - (-i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2+i & -5 \\ 1 & -2+i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2+i & -5 & 0 \\ 1 & -2+i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} + \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2+i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2+i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 - i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 - I)t \\ t \end{bmatrix} = \begin{bmatrix} (2 - i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 - I)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 - I)t \\ t \end{bmatrix} = \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - (i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 - i & -5 \\ 1 & -2 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 1 & -2 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} - \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 - i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 + i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 + I) t \\ t \end{bmatrix} = \begin{bmatrix} (2 + i) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 + I) t \\ t \end{bmatrix} = t \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 + I) t \\ t \end{bmatrix} = \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
i	1	1	No	$\begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$
$-i$	1	1	No	$\begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (2 + i) e^{it} \\ e^{it} \end{bmatrix} + c_2 \begin{bmatrix} (2 - i) e^{-it} \\ e^{-it} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -\frac{ie^{-it}}{2} & (\frac{1}{2} + i)e^{-it} \\ \frac{ie^{it}}{2} & (\frac{1}{2} - i)e^{it} \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix} \int \begin{bmatrix} -\frac{ie^{-it}}{2} & (\frac{1}{2} + i)e^{-it} \\ \frac{ie^{it}}{2} & (\frac{1}{2} - i)e^{it} \end{bmatrix} \begin{bmatrix} 0 \\ \cos(t) \end{bmatrix} dt \\ &= \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix} \int \begin{bmatrix} (\frac{1}{2} + i)e^{-it} \cos(t) \\ (\frac{1}{2} - i)e^{it} \cos(t) \end{bmatrix} dt \\ &= \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix} \begin{bmatrix} (\frac{1}{4} + \frac{i}{2})(it \sin(t) + t \cos(t) + \sin(t))e^{-it} \\ (\frac{1}{4} - \frac{i}{2})t + (-\frac{1}{4} - \frac{i}{8})e^{2it} \end{bmatrix} \\ &= \begin{bmatrix} -\frac{5 \cos(t)}{8} + \frac{5(i-4t) \sin(t)}{8} \\ \frac{\cos(t)(-2-i+4t)}{8} + \frac{(\frac{3}{2}+i-4t) \sin(t)}{4} \end{bmatrix} \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} (2+i)c_1 e^{it} \\ c_1 e^{it} \end{bmatrix} + \begin{bmatrix} (2-i)c_2 e^{-it} \\ c_2 e^{-it} \end{bmatrix} + \begin{bmatrix} -\frac{5 \cos(t)}{8} + \frac{5(i-4t) \sin(t)}{8} \\ \frac{\cos(t)(-2-i+4t)}{8} + \frac{(\frac{3}{2}+i-4t) \sin(t)}{4} \end{bmatrix} \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{((16-8i)c_2-10it-5)e^{-it}}{8} + \frac{5((\frac{8}{5}+\frac{4i}{5})c_1+it)e^{it}}{4} \\ \frac{(8c_1-2-i+4t+8c_2)\cos(t)}{8} + (ic_1 - t + \frac{3}{8} + \frac{1}{4}i - ic_2)\sin(t) \end{bmatrix}$$

18.11.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 2x_1(t) - 5x_2(t), x_2'(t) = x_1(t) - 2x_2(t) + \cos(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ \cos(t) \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ \cos(t) \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} 0 \\ \cos(t) \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right], \left[I, \begin{bmatrix} 2 + I \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{-It} \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$(\cos(t) - I \sin(t)) \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Simplify expression

$$\begin{bmatrix} (2 - I)(\cos(t) - I \sin(t)) \\ \cos(t) - I \sin(t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\begin{bmatrix} \underline{x}_1(t) \\ \underline{x}_2(t) \end{bmatrix} = \begin{bmatrix} 2 \cos(t) - \sin(t) \\ \cos(t) \end{bmatrix}, \begin{bmatrix} \underline{x}_3(t) \\ \underline{x}_4(t) \end{bmatrix} = \begin{bmatrix} -\cos(t) - 2 \sin(t) \\ -\sin(t) \end{bmatrix}$$

- General solution of the system of ODEs can be written in terms of the particular solution $\underline{x}_p(t)$

$$\underline{x}(t) = c_1 \underline{x}_1(t) + c_2 \underline{x}_2(t) + \underline{x}_p(t)$$

□ Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} 2 \cos(t) - \sin(t) & -\cos(t) - 2 \sin(t) \\ \cos(t) & -\sin(t) \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix

$$\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} 2 \cos(t) - \sin(t) & -\cos(t) - 2 \sin(t) \\ \cos(t) & -\sin(t) \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} 2 & -1 \\ 1 & 0 \end{bmatrix}}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix}$$

- Find a particular solution of the system of ODEs using variation of parameters

- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\underline{x}_{\rightarrow p}(t) = \Phi(t) \cdot \vec{v}(t)$$

- Take the derivative of the particular solution

$$\underline{x}_{\rightarrow p}'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t) \cdot \vec{v}(t)$

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- Plug $\vec{v}(t)$ into the equation for the particular solution

$$\underline{x}_{\rightarrow p}(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- Plug in the fundamental matrix and the forcing function and compute

$$\underline{x}_{\rightarrow p}(t) = \begin{bmatrix} -\frac{5t \sin(t)}{2} \\ -t \sin(t) + \frac{\sin(t)}{2} + \frac{t \cos(t)}{2} \end{bmatrix}$$

- Plug particular solution back into general solution

$$\vec{x}(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + \begin{bmatrix} -\frac{5t \sin(t)}{2} \\ -t \sin(t) + \frac{\sin(t)}{2} + \frac{t \cos(t)}{2} \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(-5t-2c_1-4c_2)\sin(t)}{2} + 2\left(c_1 - \frac{c_2}{2}\right) \cos(t) \\ \frac{(-2t-2c_2+1)\sin(t)}{2} + \frac{\cos(t)(2c_1+t)}{2} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{(-5t-2c_1-4c_2)\sin(t)}{2} + 2\left(c_1 - \frac{c_2}{2}\right) \cos(t), x_2(t) = \frac{(-2t-2c_2+1)\sin(t)}{2} + \frac{\cos(t)(2c_1+t)}{2} \right\}$$

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 57

```
dsolve([diff(x__1(t),t)=2*x__1(t)-5*x__2(t)+0,diff(x__2(t),t)=1*x__1(t)-2*x__2(t)+cos(t)],si
```

$$x_1(t) = c_2 \sin(t) + c_1 \cos(t) - \frac{5 \sin(t) t}{2}$$

$$x_2(t) = -\frac{c_2 \cos(t)}{5} + \frac{c_1 \sin(t)}{5} + \frac{\cos(t) t}{2} + \frac{\sin(t)}{2} + \frac{2c_2 \sin(t)}{5} + \frac{2c_1 \cos(t)}{5} - \sin(t) t$$

✓ Solution by Mathematica

Time used: 0.024 (sec). Leaf size: 60

```
DSolve[{x1'[t]==2*x1[t]-5*x2[t]+0,x2'[t]==1*x1[t]-2*x2[t]-Cos[t]},{x1[t],x2[t]},t,IncludeSin
```

$$x1(t) \rightarrow \left(\frac{5}{2} + c_1\right) \cos(t) + \frac{1}{2}(5t + 4c_1 - 10c_2) \sin(t)$$

$$x2(t) \rightarrow \left(-\frac{t}{2} + 1 + c_2\right) \cos(t) + (t + c_1 - 2c_2) \sin(t)$$

18.12 problem 12

18.12.1 Solution using Matrix exponential method	3913
18.12.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3915
18.12.3 Maple step by step solution	3920

Internal problem ID [789]

Internal file name [OUTPUT/789_Sunday_June_05_2022_01_49_48_AM_41360408/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 12.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x_1'(t) = 2x_1(t) - 5x_2(t) + \csc(t)$$

$$x_2'(t) = x_1(t) - 2x_2(t) + \sec(t)$$

18.12.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} \csc(t) \\ \sec(t) \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} (\cos(t) + 2 \sin(t)) c_1 - 5 \sin(t) c_2 \\ \sin(t) c_1 + (\cos(t) - 2 \sin(t)) c_2 \end{bmatrix} \\ &= \begin{bmatrix} (2c_1 - 5c_2) \sin(t) + c_1 \cos(t) \\ (c_1 - 2c_2) \sin(t) + c_2 \cos(t) \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \cos(t) - 2 \sin(t) & 5 \sin(t) \\ -\sin(t) & \cos(t) + 2 \sin(t) \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \int \begin{bmatrix} \cos(t) - 2 \sin(t) & 5 \sin(t) \\ -\sin(t) & \cos(t) + 2 \sin(t) \end{bmatrix} \begin{bmatrix} \csc(t) \\ \sec(t) \end{bmatrix} dt \\ &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \begin{bmatrix} -2t + \ln(\sin(t)) + 5 \ln(\sec(t)) \\ -2 \ln(\cos(t)) \end{bmatrix} \\ &= \begin{bmatrix} (\cos(t) + 2 \sin(t)) (-2t + \ln(\sin(t)) + 5 \ln(\sec(t))) + 10 \sin(t) \ln(\cos(t)) \\ (-2 \cos(t) + 4 \sin(t)) \ln(\cos(t)) - 2 \left(t - \frac{5 \ln(\sec(t))}{2} - \frac{\ln(\sin(t))}{2} \right) \sin(t) \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} (2c_1 - 5c_2) \sin(t) + c_1 \cos(t) + (\cos(t) + 2 \sin(t))(-2t + \ln(\sin(t)) + 5 \ln(\sec(t))) + 10 \sin(t) \\ (-2 \cos(t) + 4 \sin(t)) \ln(\cos(t)) + 5 \sin(t) \ln(\sec(t)) + \sin(t) \ln(\sin(t)) + (-2t + c_1 - 2c_2) \sin(t) \end{bmatrix}\end{aligned}$$

18.12.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} \csc(t) \\ \sec(t) \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A \vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A \vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 2 - \lambda & -5 \\ 1 & -2 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-i$	1	complex eigenvalue
i	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - (-i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2+i & -5 \\ 1 & -2+i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2+i & -5 & 0 \\ 1 & -2+i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} + \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2+i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2+i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 - i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 - I)t \\ t \end{bmatrix} = \begin{bmatrix} (2 - i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 - I)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 - I)t \\ t \end{bmatrix} = \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - (i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 - i & -5 \\ 1 & -2 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 1 & -2 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} - \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 - i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 + i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 + I) t \\ t \end{bmatrix} = \begin{bmatrix} (2 + i) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 + I) t \\ t \end{bmatrix} = t \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 + I) t \\ t \end{bmatrix} = \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
i	1	1	No	$\begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$
$-i$	1	1	No	$\begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (2 + i) e^{it} \\ e^{it} \end{bmatrix} + c_2 \begin{bmatrix} (2 - i) e^{-it} \\ e^{-it} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -\frac{ie^{-it}}{2} & (\frac{1}{2} + i)e^{-it} \\ \frac{ie^{it}}{2} & (\frac{1}{2} - i)e^{it} \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix} \int \begin{bmatrix} -\frac{ie^{-it}}{2} & (\frac{1}{2} + i)e^{-it} \\ \frac{ie^{it}}{2} & (\frac{1}{2} - i)e^{it} \end{bmatrix} \begin{bmatrix} \csc(t) \\ \sec(t) \end{bmatrix} dt \\ &= \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix} \int \begin{bmatrix} -\frac{((-1-2i)\sec(t)+i\csc(t))e^{-it}}{2} \\ \frac{((1-2i)\sec(t)+i\csc(t))e^{it}}{2} \end{bmatrix} dt \\ &= \begin{bmatrix} (2+i)e^{it} & (2-i)e^{-it} \\ e^{it} & e^{-it} \end{bmatrix} \begin{bmatrix} (-1 + \frac{i}{2}) \ln(e^{2it} + 1) - \frac{i \ln(e^{it}-1)}{2} - \frac{i \ln(e^{it}+1)}{2} + 2 \ln(e^{it}) \\ \frac{i \ln(e^{2it}-1)}{2} + (-1 - \frac{i}{2}) \ln(e^{2it} + 1) \end{bmatrix} \\ &= \begin{bmatrix} (\frac{1}{2} + i)e^{-it} \ln(e^{2it} - 1) + \frac{(-5e^{-it}-5e^{it}) \ln(e^{2it}+1)}{2} + 2((\frac{1}{4} - \frac{i}{2}) \ln(e^{it} - 1) + (\frac{1}{4} - \frac{i}{2}) \ln(e^{it} + 1) + \\ \frac{((-2-i)e^{-it}+(-2+i)e^{it}) \ln(e^{2it}+1)}{2} + \frac{ie^{-it} \ln(e^{2it}-1)}{2} - \frac{(i \ln(e^{it}-1)+i \ln(e^{it}+1)-4 \ln(e^{it}))e^{it}}{2} \end{bmatrix} \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} (2+i)c_1e^{it} \\ c_1e^{it} \end{bmatrix} + \begin{bmatrix} (2-i)c_2e^{-it} \\ c_2e^{-it} \end{bmatrix} + \begin{bmatrix} (\frac{1}{2} + i)e^{-it} \ln(e^{2it} - 1) + \frac{(-5e^{-it}-5e^{it}) \ln(e^{2it}+1)}{2} \\ \frac{((-2-i)e^{-it}+(-2+i)e^{it}) \ln(e^{2it}+1)}{2} + \frac{ie^{-it} \ln(e^{2it}-1)}{2} - \frac{(i \ln(e^{it}-1)+i \ln(e^{it}+1)-4 \ln(e^{it}))e^{it}}{2} \end{bmatrix} + \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \left(\frac{1}{2} + i\right) e^{-it} \ln(e^{2it} - 1) + \frac{(-5e^{-it} - 5e^{it}) \ln(e^{2it} + 1)}{2} + \left(\frac{1}{2} - i\right) e^{it} \ln(e^{it} - 1) + \left(\frac{1}{2} - i\right) e^{it} \ln(e^{it} + 1) \\ \frac{((-2-i)e^{-it} + (-2+i)e^{it}) \ln(e^{2it} + 1)}{2} + \frac{ie^{-it} \ln(e^{2it} - 1)}{2} - \frac{ie^{it} \ln(e^{it} - 1)}{2} - \frac{ie^{it} \ln(e^{it} + 1)}{2} \end{bmatrix}$$

18.12.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 2x_1(t) - 5x_2(t) + \csc(t), x_2'(t) = x_1(t) - 2x_2(t) + \sec(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} \csc(t) \\ \sec(t) \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} \csc(t) \\ \sec(t) \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} \csc(t) \\ \sec(t) \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right], \left[I, \begin{bmatrix} 2 + I \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{-It} \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$(\cos(t) - I \sin(t)) \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Simplify expression

$$\begin{bmatrix} (2 - I)(\cos(t) - I \sin(t)) \\ \cos(t) - I \sin(t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\underline{x}_1(t) = \begin{bmatrix} 2 \cos(t) - \sin(t) \\ \cos(t) \end{bmatrix}, \underline{x}_2(t) = \begin{bmatrix} -\cos(t) - 2 \sin(t) \\ -\sin(t) \end{bmatrix} \right]$$

- General solution of the system of ODEs can be written in terms of the particular solution $\underline{x}_p(t)$

$$\underline{x}(t) = c_1 \underline{x}_1(t) + c_2 \underline{x}_2(t) + \underline{x}_p(t)$$

□ Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} 2 \cos(t) - \sin(t) & -\cos(t) - 2 \sin(t) \\ \cos(t) & -\sin(t) \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix

$$\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} 2 \cos(t) - \sin(t) & -\cos(t) - 2 \sin(t) \\ \cos(t) & -\sin(t) \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} 2 & -1 \\ 1 & 0 \end{bmatrix}}$$

- o Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix}$$

- Find a particular solution of the system of ODEs using variation of parameters

- o Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\underline{x}_{\rightarrow p}(t) = \Phi(t) \cdot \vec{v}(t)$$

- o Take the derivative of the particular solution

$$\underline{x}_{\rightarrow p}'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- o Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- o The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t) \cdot \vec{v}(t)$

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- o Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- o Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- o Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- o Plug $\vec{v}(t)$ into the equation for the particular solution

$$\underline{x}_{\rightarrow p}(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- o Plug in the fundamental matrix and the forcing function and compute

$$\underline{x}_{\rightarrow p}(t) = \begin{bmatrix} (\cos(t) + 2 \sin(t)) \left(\int_0^t (-4 \cot(s) - 2 + 5 \sec(s) \csc(s)) ds \right) - 10 \sin(t) \left(\int_0^t \tan(s) ds \right) \\ \left(\int_0^t (-4 \cot(s) - 2 + 5 \sec(s) \csc(s)) ds \right) \sin(t) + 2(\cos(t) - 2 \sin(t)) \left(\int_0^t \tan(s) ds \right) \end{bmatrix}$$

- Plug particular solution back into general solution

$$\vec{x}(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t) + \begin{bmatrix} (\cos(t) + 2 \sin(t)) \left(\int_0^t (-4 \cot(s) - 2 + 5 \sec(s) \csc(s)) ds \right) - 10 \sin(t) \left(\int_0^t \tan(s) ds \right) \\ \left(\int_0^t (-4 \cot(s) - 2 + 5 \sec(s) \csc(s)) ds \right) \sin(t) + 2(\cos(t) - 4 \sin(t)) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (\cos(t) + 2 \sin(t)) \left(\int_0^t (-4 \cot(s) - 2 + 5 \sec(s) \csc(s)) ds \right) - 10 \sin(t) \left(\int_0^t \tan(s) ds \right) \\ \left(\int_0^t (-4 \cot(s) - 2 + 5 \sec(s) \csc(s)) ds \right) \sin(t) + (2 \cos(t) - 4 \sin(t)) \end{bmatrix}$$

- Solution to the system of ODEs

$$\begin{cases} x_1(t) = (\cos(t) + 2 \sin(t)) \left(\int_0^t (-4 \cot(s) - 2 + 5 \sec(s) \csc(s)) ds \right) - 10 \sin(t) \left(\int_0^t \tan(s) ds \right) \\ x_2(t) = \left(\int_0^t (-4 \cot(s) - 2 + 5 \sec(s) \csc(s)) ds \right) \sin(t) + (2 \cos(t) - 4 \sin(t)) \end{cases}$$

✓ Solution by Maple

Time used: 0.235 (sec). Leaf size: 113

```
dsolve([diff(x__1(t),t)=2*x__1(t)-5*x__2(t)+csc(t),diff(x__2(t),t)=1*x__1(t)-2*x__2(t)+sec(t)
```

$$x_1(t) = \ln(\sin(t)) \cos(t) - 5 \cos(t) \ln(\cos(t)) + c_1 \cos(t) - 2 \cos(t) t + 2 \ln(\sin(t)) \sin(t) + c_2 \sin(t) - 4 \sin(t) t + \cos(t)$$

$$x_2(t) = -2 \cos(t) \ln(\cos(t)) + \frac{2c_1 \cos(t)}{5} - \frac{c_2 \cos(t)}{5} + \ln(\sin(t)) \sin(t) - \sin(t) \ln(\cos(t)) + \frac{c_1 \sin(t)}{5} + \frac{2c_2 \sin(t)}{5} - 2 \sin(t) t - \frac{\cos(t)^2}{5 \sin(t)} + \frac{2 \cos(t)}{5} + \frac{\csc(t)}{5}$$

✓ Solution by Mathematica

Time used: 0.018 (sec). Leaf size: 79

```
DSolve[{x1'[t]==2*x1[t]-5*x2[t]+Csc[t],x2'[t]==1*x1[t]-2*x2[t]+Sec[t]},{x1[t],x2[t]},t,Inclu
```

$$x_1(t) \rightarrow \sin(t)(-4t + 2 \log(\sin(t)) + 2c_1 - 5c_2)$$

$$+ \cos(t)(-2t + \log(\sin(t)) - 5 \log(\cos(t)) + c_1)$$

$$x_2(t) \rightarrow \cos(t)(-2 \log(\cos(t)) + c_2) + \sin(t)(-2t + \log(\sin(t)) - \log(\cos(t)) + c_1 - 2c_2)$$

18.13 problem 13

- 18.13.1 Solution using Matrix exponential method 3924
- 18.13.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3926
- 18.13.3 Maple step by step solution 3931

Internal problem ID [790]

Internal file name [OUTPUT/790_Sunday_June_05_2022_01_49_50_AM_26957802/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 13.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -\frac{x_1(t)}{2} - \frac{x_2(t)}{8} + \frac{e^{-\frac{t}{2}}}{2} \\x_2'(t) &= 2x_1(t) - \frac{x_2(t)}{2}\end{aligned}$$

18.13.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} \frac{e^{-\frac{t}{2}}}{2} \\ 0 \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) & -\frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right)}{4} \\ 4e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) & e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) & -\frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right)}{4} \\ 4e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) & e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) c_1 - \frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) c_2}{4} \\ 4e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) c_1 + e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) c_2 \end{bmatrix} \\ &= \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) c_1 - \frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) c_2}{4} \\ e^{-\frac{t}{2}} (4 \sin\left(\frac{t}{2}\right) c_1 + \cos\left(\frac{t}{2}\right) c_2) \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \cos\left(\frac{t}{2}\right) e^{\frac{t}{2}} & \frac{\sin\left(\frac{t}{2}\right) e^{\frac{t}{2}}}{4} \\ -4 \sin\left(\frac{t}{2}\right) e^{\frac{t}{2}} & \cos\left(\frac{t}{2}\right) e^{\frac{t}{2}} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) & -\frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right)}{4} \\ 4e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) & e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) \end{bmatrix} \int \begin{bmatrix} \cos\left(\frac{t}{2}\right) e^{\frac{t}{2}} & \frac{\sin\left(\frac{t}{2}\right) e^{\frac{t}{2}}}{4} \\ -4 \sin\left(\frac{t}{2}\right) e^{\frac{t}{2}} & \cos\left(\frac{t}{2}\right) e^{\frac{t}{2}} \end{bmatrix} \begin{bmatrix} \frac{e^{-\frac{t}{2}}}{2} \\ 0 \end{bmatrix} dt \\ &= \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) & -\frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right)}{4} \\ 4e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) & e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) \end{bmatrix} \begin{bmatrix} \sin\left(\frac{t}{2}\right) \\ 4 \cos\left(\frac{t}{2}\right) \end{bmatrix} \\ &= \begin{bmatrix} 0 \\ 4e^{-\frac{t}{2}} \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} e^{-\frac{t}{2}} \cos\left(\frac{t}{2}\right) c_1 - \frac{e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) c_2}{4} \\ e^{-\frac{t}{2}} \left(4 + 4 \sin\left(\frac{t}{2}\right) c_1 + \cos\left(\frac{t}{2}\right) c_2\right) \end{bmatrix}\end{aligned}$$

18.13.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} \frac{e^{-\frac{t}{2}}}{2} \\ 0 \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -\frac{1}{2} - \lambda & -\frac{1}{8} \\ 2 & -\frac{1}{2} - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + \lambda + \frac{1}{2} = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -\frac{1}{2} + \frac{i}{2}$$

$$\lambda_2 = -\frac{1}{2} - \frac{i}{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-\frac{1}{2} + \frac{i}{2}$	1	complex eigenvalue
$-\frac{1}{2} - \frac{i}{2}$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -\frac{1}{2} - \frac{i}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} - \left(-\frac{1}{2} - \frac{i}{2} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} \frac{i}{2} & -\frac{1}{8} \\ 2 & \frac{i}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{i}{2} & -\frac{1}{8} & 0 \\ 2 & \frac{i}{2} & 0 \end{array} \right]$$

$$R_2 = 4iR_1 + R_2 \implies \left[\begin{array}{cc|c} \frac{i}{2} & -\frac{1}{8} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} \frac{i}{2} & -\frac{1}{8} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -\frac{it}{4}\}$

Hence the solution is

$$\begin{bmatrix} -\frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} -\frac{it}{4} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{1}{4}t \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{i}{4} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} -\frac{i}{4} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} -i \\ 4 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -\frac{1}{2} + \frac{i}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} - \left(-\frac{1}{2} + \frac{i}{2} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -\frac{i}{2} & -\frac{1}{8} \\ 2 & -\frac{i}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -\frac{i}{2} & -\frac{1}{8} & 0 \\ 2 & -\frac{i}{2} & 0 \end{array} \right]$$

$$R_2 = -4iR_1 + R_2 \implies \left[\begin{array}{cc|c} -\frac{i}{2} & -\frac{1}{8} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -\frac{i}{2} & -\frac{1}{8} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{it}{4}\}$

Hence the solution is

$$\begin{bmatrix} \frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} \frac{it}{4} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{1}{4}t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{i}{4} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} \frac{i}{4} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{4}t \\ t \end{bmatrix} = \begin{bmatrix} i \\ 4 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-\frac{1}{2} + \frac{i}{2}$	1	1	No	$\begin{bmatrix} \frac{i}{4} \\ 1 \end{bmatrix}$
$-\frac{1}{2} - \frac{i}{2}$	1	1	No	$\begin{bmatrix} -\frac{i}{4} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{ie^{(-\frac{1}{2} + \frac{i}{2})t}}{4} \\ e^{(-\frac{1}{2} + \frac{i}{2})t} \end{bmatrix} + c_2 \begin{bmatrix} -\frac{ie^{(-\frac{1}{2} - \frac{i}{2})t}}{4} \\ e^{(-\frac{1}{2} - \frac{i}{2})t} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \dots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} \frac{ie^{(-\frac{1}{2} + \frac{i}{2})t}}{4} & -\frac{ie^{(-\frac{1}{2} - \frac{i}{2})t}}{4} \\ e^{(-\frac{1}{2} + \frac{i}{2})t} & e^{(-\frac{1}{2} - \frac{i}{2})t} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -2ie^{(\frac{1}{2} - \frac{i}{2})t} & \frac{e^{(\frac{1}{2} - \frac{i}{2})t}}{2} \\ 2ie^{(\frac{1}{2} + \frac{i}{2})t} & \frac{e^{(\frac{1}{2} + \frac{i}{2})t}}{2} \end{bmatrix}$$

Hence

$$\begin{aligned}
 \vec{x}_p(t) &= \begin{bmatrix} \frac{ie^{(-\frac{1}{2}+\frac{i}{2})t}}{4} & -\frac{ie^{(-\frac{1}{2}-\frac{i}{2})t}}{4} \\ e^{(-\frac{1}{2}+\frac{i}{2})t} & e^{(-\frac{1}{2}-\frac{i}{2})t} \end{bmatrix} \int \begin{bmatrix} -2ie^{(\frac{1}{2}-\frac{i}{2})t} & \frac{e^{(\frac{1}{2}-\frac{i}{2})t}}{2} \\ 2ie^{(\frac{1}{2}+\frac{i}{2})t} & \frac{e^{(\frac{1}{2}+\frac{i}{2})t}}{2} \end{bmatrix} \begin{bmatrix} \frac{e^{-\frac{t}{2}}}{2} \\ 0 \end{bmatrix} dt \\
 &= \begin{bmatrix} \frac{ie^{(-\frac{1}{2}+\frac{i}{2})t}}{4} & -\frac{ie^{(-\frac{1}{2}-\frac{i}{2})t}}{4} \\ e^{(-\frac{1}{2}+\frac{i}{2})t} & e^{(-\frac{1}{2}-\frac{i}{2})t} \end{bmatrix} \int \begin{bmatrix} -ie^{-\frac{it}{2}} \\ ie^{\frac{it}{2}} \end{bmatrix} dt \\
 &= \begin{bmatrix} \frac{ie^{(-\frac{1}{2}+\frac{i}{2})t}}{4} & -\frac{ie^{(-\frac{1}{2}-\frac{i}{2})t}}{4} \\ e^{(-\frac{1}{2}+\frac{i}{2})t} & e^{(-\frac{1}{2}-\frac{i}{2})t} \end{bmatrix} \begin{bmatrix} 2e^{-\frac{it}{2}} \\ 2e^{\frac{it}{2}} \end{bmatrix} \\
 &= \begin{bmatrix} 0 \\ 4e^{-\frac{t}{2}} \end{bmatrix}
 \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned}
 \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\
 \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} \frac{ic_1e^{(-\frac{1}{2}+\frac{i}{2})t}}{4} \\ c_1e^{(-\frac{1}{2}+\frac{i}{2})t} \end{bmatrix} + \begin{bmatrix} -\frac{ic_2e^{(-\frac{1}{2}-\frac{i}{2})t}}{4} \\ c_2e^{(-\frac{1}{2}-\frac{i}{2})t} \end{bmatrix} + \begin{bmatrix} 0 \\ 4e^{-\frac{t}{2}} \end{bmatrix}
 \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{i(c_1e^{(-\frac{1}{2}+\frac{i}{2})t} - c_2e^{(-\frac{1}{2}-\frac{i}{2})t})}{4} \\ c_1e^{(-\frac{1}{2}+\frac{i}{2})t} + c_2e^{(-\frac{1}{2}-\frac{i}{2})t} + 4e^{-\frac{t}{2}} \end{bmatrix}$$

18.13.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = -\frac{x_1(t)}{2} - \frac{x_2(t)}{8} + \frac{e^{-\frac{t}{2}}}{2}, x_2'(t) = 2x_1(t) - \frac{x_2(t)}{2} \right]$$

- Define vector

$$\vec{x}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} \frac{e^{-\frac{t}{2}}}{2} \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} \frac{e^{-\frac{t}{2}}}{2} \\ 0 \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} \frac{e^{-\frac{t}{2}}}{2} \\ 0 \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{8} \\ 2 & -\frac{1}{2} \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-\frac{1}{2} - \frac{1}{2}, \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix} \right], \left[-\frac{1}{2} + \frac{1}{2}, \begin{bmatrix} \frac{1}{4} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-\frac{1}{2} - \frac{1}{2}, \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(-\frac{1}{2}-\frac{1}{2})t} \cdot \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{-\frac{t}{2}} \cdot \left(\cos\left(\frac{t}{2}\right) - I \sin\left(\frac{t}{2}\right) \right) \cdot \begin{bmatrix} -\frac{1}{4} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{-\frac{t}{2}} \cdot \begin{bmatrix} -\frac{1}{4}(\cos(\frac{t}{2}) - I \sin(\frac{t}{2})) \\ \cos(\frac{t}{2}) - I \sin(\frac{t}{2}) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\underline{x}_{\rightarrow 1}(t) = e^{-\frac{t}{2}} \cdot \begin{bmatrix} -\frac{\sin(\frac{t}{2})}{4} \\ \cos(\frac{t}{2}) \end{bmatrix}, \underline{x}_{\rightarrow 2}(t) = e^{-\frac{t}{2}} \cdot \begin{bmatrix} -\frac{\cos(\frac{t}{2})}{4} \\ -\sin(\frac{t}{2}) \end{bmatrix} \right]$$

- General solution of the system of ODEs can be written in terms of the particular solution $\underline{x}_{\rightarrow p}(t)$
 $\underline{x}_{\rightarrow}(t) = c_1 \underline{x}_{\rightarrow 1}(t) + c_2 \underline{x}_{\rightarrow 2}(t) + \underline{x}_{\rightarrow p}(t)$

- Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} -\frac{e^{-\frac{t}{2}} \sin(\frac{t}{2})}{4} & -\frac{e^{-\frac{t}{2}} \cos(\frac{t}{2})}{4} \\ e^{-\frac{t}{2}} \cos(\frac{t}{2}) & -e^{-\frac{t}{2}} \sin(\frac{t}{2}) \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix
 $\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} -\frac{e^{-\frac{t}{2}} \sin(\frac{t}{2})}{4} & -\frac{e^{-\frac{t}{2}} \cos(\frac{t}{2})}{4} \\ e^{-\frac{t}{2}} \cos(\frac{t}{2}) & -e^{-\frac{t}{2}} \sin(\frac{t}{2}) \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} 0 & -\frac{1}{4} \\ 1 & 0 \end{bmatrix}}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} e^{-\frac{t}{2}} \cos(\frac{t}{2}) & -\frac{e^{-\frac{t}{2}} \sin(\frac{t}{2})}{4} \\ 4e^{-\frac{t}{2}} \sin(\frac{t}{2}) & e^{-\frac{t}{2}} \cos(\frac{t}{2}) \end{bmatrix}$$

- Find a particular solution of the system of ODEs using variation of parameters

- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$
 $\underline{x}_{\rightarrow p}(t) = \Phi(t) \cdot \vec{v}(t)$

- Take the derivative of the particular solution

$$\underline{x}_{\rightarrow p}'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- Plug $\vec{v}(t)$ into the equation for the particular solution

$$\underline{x}_{\rightarrow p}(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- Plug in the fundamental matrix and the forcing function and compute

$$\underline{x}_{\rightarrow p}(t) = \begin{bmatrix} e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) \\ -4 e^{-\frac{t}{2}} \left(-1 + \cos\left(\frac{t}{2}\right)\right) \end{bmatrix}$$

- Plug particular solution back into general solution

$$\underline{x}_{\rightarrow}(t) = c_1 \underline{x}_{\rightarrow 1}(t) + c_2 \underline{x}_{\rightarrow 2}(t) + \begin{bmatrix} e^{-\frac{t}{2}} \sin\left(\frac{t}{2}\right) \\ -4 e^{-\frac{t}{2}} \left(-1 + \cos\left(\frac{t}{2}\right)\right) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} -\frac{((-4+c_1)\sin(\frac{t}{2})+c_2\cos(\frac{t}{2}))e^{-\frac{t}{2}}}{4} \\ ((-4+c_1)\cos(\frac{t}{2})-c_2\sin(\frac{t}{2})+4)e^{-\frac{t}{2}} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = -\frac{((-4+c_1)\sin(\frac{t}{2})+c_2\cos(\frac{t}{2}))e^{-\frac{t}{2}}}{4}, x_2(t) = ((-4+c_1)\cos(\frac{t}{2})-c_2\sin(\frac{t}{2})+4)e^{-\frac{t}{2}} \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 47

```
dsolve([diff(x__1(t),t)=-1/2*x__1(t)-1/8*x__2(t)+1/2*exp(-t/2),diff(x__2(t),t)=2*x__1(t)-1/2
```

$$x_1(t) = \frac{e^{-\frac{t}{2}} \left(c_2 \cos\left(\frac{t}{2}\right) - c_1 \sin\left(\frac{t}{2}\right) \right)}{4}$$
$$x_2(t) = e^{-\frac{t}{2}} \left(4 + \cos\left(\frac{t}{2}\right) c_1 + \sin\left(\frac{t}{2}\right) c_2 \right)$$

✓ Solution by Mathematica

Time used: 0.022 (sec). Leaf size: 69

```
DSolve[{x1'[t]==-1/2*x1[t]-1/8*x2[t]+1/2*Exp[-t/2],x2'[t]==2*x1[t]-1/2*x2[t]+0},{x1[t],x2[t]
```

$$x_1(t) \rightarrow \frac{1}{4} e^{-t/2} \left(4c_1 \cos\left(\frac{t}{2}\right) - c_2 \sin\left(\frac{t}{2}\right) \right)$$
$$x_2(t) \rightarrow e^{-t/2} \left(c_2 \cos\left(\frac{t}{2}\right) + 4c_1 \sin\left(\frac{t}{2}\right) + 4 \right)$$

18.14 problem 18

18.14.1 Solution using Matrix exponential method 3936

18.14.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3938

Internal problem ID [791]

Internal file name [OUTPUT/791_Sunday_June_05_2022_01_49_52_AM_56180441/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 7.9, Nonhomogeneous Linear Systems. page 447

Problem number: 18.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -2x_1(t) + x_2(t) + 2e^{-t} \\x_2'(t) &= x_1(t) - 2x_2(t) + 3t\end{aligned}$$

With initial conditions

$$[x_1(0) = \alpha_1, x_2(0) = \alpha_2]$$

18.14.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 2e^{-t} \\ 3t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} & \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} & \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{x}_0 \\ &= \begin{bmatrix} \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} & \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} & \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} \end{bmatrix} \begin{bmatrix} \alpha_1 \\ \alpha_2 \end{bmatrix} \\ &= \begin{bmatrix} \left(\frac{e^{-3t}}{2} + \frac{e^{-t}}{2}\right) \alpha_1 + \left(\frac{e^{-t}}{2} - \frac{e^{-3t}}{2}\right) \alpha_2 \\ \left(\frac{e^{-t}}{2} - \frac{e^{-3t}}{2}\right) \alpha_1 + \left(\frac{e^{-3t}}{2} + \frac{e^{-t}}{2}\right) \alpha_2 \end{bmatrix} \\ &= \begin{bmatrix} \frac{(-\alpha_2 + \alpha_1)e^{-3t}}{2} + \frac{e^{-t}(\alpha_1 + \alpha_2)}{2} \\ \frac{(\alpha_2 - \alpha_1)e^{-3t}}{2} + \frac{e^{-t}(\alpha_1 + \alpha_2)}{2} \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \frac{(e^{2t}+1)e^t}{2} & -\frac{(e^{2t}-1)e^t}{2} \\ -\frac{(e^{2t}-1)e^t}{2} & \frac{(e^{2t}+1)e^t}{2} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} & \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} & \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} \end{bmatrix} \int \begin{bmatrix} \frac{(e^{2t}+1)e^t}{2} & -\frac{(e^{2t}-1)e^t}{2} \\ -\frac{(e^{2t}-1)e^t}{2} & \frac{(e^{2t}+1)e^t}{2} \end{bmatrix} \begin{bmatrix} 2e^{-t} \\ 3t \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} & \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} & \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} \end{bmatrix} \begin{bmatrix} \frac{(1-3t)e^{3t}}{6} + \frac{e^{2t}}{2} + \frac{(9t-9)e^t}{6} + t \\ \frac{(-1+3t)e^{3t}}{6} - \frac{e^{2t}}{2} + \frac{(9t-9)e^t}{6} + t \end{bmatrix} \\ &= \begin{bmatrix} t - \frac{4}{3} + \frac{e^{-t}}{2} + te^{-t} \\ te^{-t} + 2t - \frac{5}{3} - \frac{e^{-t}}{2} \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{(6t+3\alpha_1+3\alpha_2+3)e^{-t}}{6} + \frac{(3\alpha_1-3\alpha_2)e^{-3t}}{6} + t - \frac{4}{3} \\ \frac{(6t+3\alpha_1+3\alpha_2-3)e^{-t}}{6} + \frac{(-3\alpha_1+3\alpha_2)e^{-3t}}{6} + 2t - \frac{5}{3} \end{bmatrix}\end{aligned}$$

18.14.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} 2e^{-t} \\ 3t \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -2 - \lambda & 1 \\ 1 & -2 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 4\lambda + 3 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1$$

$$\lambda_2 = -3$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
-3	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -3$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} - (-3) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 & 1 & 0 \\ 1 & 1 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} 1 & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -t\}$

Hence the solution is

$$\begin{bmatrix} -t \\ t \end{bmatrix} = \begin{bmatrix} -t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -t \\ t \end{bmatrix} = t \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -t \\ t \end{bmatrix} = \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -1 & 1 & 0 \\ 1 & -1 & 0 \end{array} \right]$$

$$R_2 = R_2 + R_1 \implies \left[\begin{array}{cc|c} -1 & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-1	1	1	No	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$
-3	1	1	No	$\begin{bmatrix} -1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^{-t} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{-t} \end{aligned}$$

Since eigenvalue -3 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-3t} \\ &= \begin{bmatrix} -1 \\ 1 \end{bmatrix} e^{-3t}\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} e^{-t} \\ e^{-t} \end{bmatrix} + c_2 \begin{bmatrix} -e^{-3t} \\ e^{-3t} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = [\vec{x}_1 \quad \vec{x}_2 \quad \dots]$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} e^{-t} & -e^{-3t} \\ e^{-t} & e^{-3t} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} \frac{e^t}{2} & \frac{e^t}{2} \\ -\frac{e^{3t}}{2} & \frac{e^{3t}}{2} \end{bmatrix}$$

Hence

$$\begin{aligned}
\vec{x}_p(t) &= \begin{bmatrix} e^{-t} & -e^{-3t} \\ e^{-t} & e^{-3t} \end{bmatrix} \int \begin{bmatrix} \frac{e^t}{2} & \frac{e^t}{2} \\ -\frac{e^{3t}}{2} & \frac{e^{3t}}{2} \end{bmatrix} \begin{bmatrix} 2e^{-t} \\ 3t \end{bmatrix} dt \\
&= \begin{bmatrix} e^{-t} & -e^{-3t} \\ e^{-t} & e^{-3t} \end{bmatrix} \int \begin{bmatrix} 1 + \frac{3te^t}{2} \\ -e^{2t} + \frac{3e^{3t}t}{2} \end{bmatrix} dt \\
&= \begin{bmatrix} e^{-t} & -e^{-3t} \\ e^{-t} & e^{-3t} \end{bmatrix} \begin{bmatrix} \frac{(3t-3)e^t}{2} + t \\ \frac{(-1+3t)e^{3t}}{6} - \frac{e^{2t}}{2} \end{bmatrix} \\
&= \begin{bmatrix} t - \frac{4}{3} + \frac{e^{-t}}{2} + te^{-t} \\ te^{-t} + 2t - \frac{5}{3} - \frac{e^{-t}}{2} \end{bmatrix}
\end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned}
\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\
\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} c_1 e^{-t} \\ c_1 e^{-t} \end{bmatrix} + \begin{bmatrix} -c_2 e^{-3t} \\ c_2 e^{-3t} \end{bmatrix} + \begin{bmatrix} t - \frac{4}{3} + \frac{e^{-t}}{2} + te^{-t} \\ te^{-t} + 2t - \frac{5}{3} - \frac{e^{-t}}{2} \end{bmatrix}
\end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(6t+6c_1+3)e^{-t}}{6} - c_2 e^{-3t} + t - \frac{4}{3} \\ \frac{(6t+6c_1-3)e^{-t}}{6} + c_2 e^{-3t} + 2t - \frac{5}{3} \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x_1(0) = \alpha_1 \\ x_2(0) = \alpha_2 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} \alpha_1 \\ \alpha_2 \end{bmatrix} = \begin{bmatrix} -\frac{5}{6} + c_1 - c_2 \\ -\frac{13}{6} + c_1 + c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = \frac{\alpha_1}{2} + \frac{3}{2} + \frac{\alpha_2}{2} \\ c_2 = \frac{\alpha_2}{2} + \frac{2}{3} - \frac{\alpha_1}{2} \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(6t+3\alpha_1+12+3\alpha_2)e^{-t}}{6} - \left(\frac{\alpha_2}{2} + \frac{2}{3} - \frac{\alpha_1}{2}\right) e^{-3t} + t - \frac{4}{3} \\ \frac{(6t+3\alpha_1+6+3\alpha_2)e^{-t}}{6} + \left(\frac{\alpha_2}{2} + \frac{2}{3} - \frac{\alpha_1}{2}\right) e^{-3t} + 2t - \frac{5}{3} \end{bmatrix}$$

The following are plots of each solution.

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 93

```
dsolve([diff(x__1(t),t) = -2*x__1(t)+x__2(t)+2*exp(-t), diff(x__2(t),t) = x__1(t)-2*x__2(t)+
```

$$\begin{aligned} x_1(t) &= \left(\frac{3}{2} + \frac{\alpha_2}{2} + \frac{\alpha_1}{2}\right) e^{-t} - \left(\frac{2}{3} + \frac{\alpha_2}{2} - \frac{\alpha_1}{2}\right) e^{-3t} + \frac{e^{-t}}{2} + t e^{-t} - \frac{4}{3} + t \\ x_2(t) &= \left(\frac{3}{2} + \frac{\alpha_2}{2} + \frac{\alpha_1}{2}\right) e^{-t} + \left(\frac{2}{3} + \frac{\alpha_2}{2} - \frac{\alpha_1}{2}\right) e^{-3t} + t e^{-t} + 2t - \frac{5}{3} - \frac{e^{-t}}{2} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.061 (sec). Leaf size: 122

```
DSolve[{x1'[t]==-2*x1[t]+1*x2[t]+2*Exp[-t], x2'[t]==1*x1[t]-2*x2[t]+3*t}, {x1[0]==a1, x2[0]==a2
```

$$\begin{aligned} x_1(t) &\rightarrow \frac{1}{6} e^{-3t} (3a_1(e^{2t} + 1) + 3a_2(e^{2t} - 1) + 12e^{2t} - 8e^{3t} + 6e^{2t}t + 6e^{3t}t - 4) \\ x_2(t) &\rightarrow \frac{1}{6} e^{-3t} (3a_1(e^{2t} - 1) + 3a_2(e^{2t} + 1) + 6e^{2t}(t + 1) + 2e^{3t}(6t - 5) + 4) \end{aligned}$$

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19.1 problem 1

- 19.1.1 Solution using Matrix exponential method 3946
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Internal problem ID [792]

Internal file name [OUTPUT/792_Sunday_June_05_2022_01_49_55_AM_22760055/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 3x_1(t) - 2x_2(t) \\x_2'(t) &= 2x_1(t) - 2x_2(t)\end{aligned}$$

19.1.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & -2 \\ 2 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} -\frac{e^{-t}}{3} + \frac{4e^{2t}}{3} & -\frac{2e^{2t}}{3} + \frac{2e^{-t}}{3} \\ \frac{2e^{2t}}{3} - \frac{2e^{-t}}{3} & \frac{4e^{-t}}{3} - \frac{e^{2t}}{3} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} -\frac{e^{-t}}{3} + \frac{4e^{2t}}{3} & -\frac{2e^{2t}}{3} + \frac{2e^{-t}}{3} \\ \frac{2e^{2t}}{3} - \frac{2e^{-t}}{3} & \frac{4e^{-t}}{3} - \frac{e^{2t}}{3} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \left(-\frac{e^{-t}}{3} + \frac{4e^{2t}}{3}\right) c_1 + \left(-\frac{2e^{2t}}{3} + \frac{2e^{-t}}{3}\right) c_2 \\ \left(\frac{2e^{2t}}{3} - \frac{2e^{-t}}{3}\right) c_1 + \left(\frac{4e^{-t}}{3} - \frac{e^{2t}}{3}\right) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \frac{(-c_1 + 2c_2)e^{-t}}{3} + \frac{4(c_1 - \frac{c_2}{2})e^{2t}}{3} \\ \frac{(-2c_1 + 4c_2)e^{-t}}{3} + \frac{2(c_1 - \frac{c_2}{2})e^{2t}}{3} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.1.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & -2 \\ 2 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 3 & -2 \\ 2 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 3 - \lambda & -2 \\ 2 & -2 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - \lambda - 2 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 2$$

$$\lambda_2 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
2	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 3 & -2 \\ 2 & -2 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} 4 & -2 \\ 2 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 4 & -2 & 0 \\ 2 & -1 & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{R_1}{2} \implies \left[\begin{array}{cc|c} 4 & -2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 4 & -2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{t}{2}\}$

Hence the solution is

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{t}{2} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ 2 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 3 & -2 \\ 2 & -2 \end{bmatrix} - (2) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & -2 \\ 2 & -4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 & -2 & 0 \\ 2 & -4 & 0 \end{array} \right]$$

$$R_2 = R_2 - 2R_1 \implies \left[\begin{array}{cc|c} 1 & -2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & -2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 2t\}$

Hence the solution is

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = \begin{bmatrix} 2t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = t \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
2	1	1	No	$\begin{bmatrix} 2 \\ 1 \end{bmatrix}$
-1	1	1	No	$\begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue 2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{2t} \\ &= \begin{bmatrix} 2 \\ 1 \end{bmatrix} e^{2t}\end{aligned}$$

Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-t} \\ &= \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} e^{-t}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} 2 e^{2t} \\ e^{2t} \end{bmatrix} + c_2 \begin{bmatrix} \frac{e^{-t}}{2} \\ e^{-t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} 2c_1 e^{2t} + \frac{c_2 e^{-t}}{2} \\ c_1 e^{2t} + c_2 e^{-t} \end{bmatrix}$$

The following is the phase plot of the system.

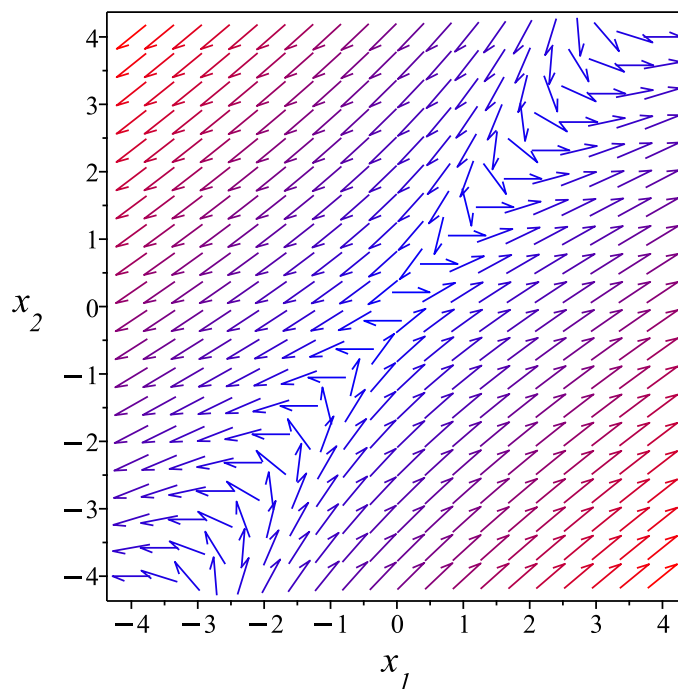


Figure 539: Phase plot

19.1.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 3x_1(t) - 2x_2(t), x_2'(t) = 2x_1(t) - 2x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 3 & -2 \\ 2 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 3 & -2 \\ 2 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 3 & -2 \\ 2 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-1, \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} \right], \left[2, \begin{bmatrix} 2 \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-1, \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{-t} \cdot \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[2, \begin{bmatrix} 2 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_2 = e^{2t} \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1 + c_2 \underline{x}^{\rightarrow}_2$$

- Substitute solutions into the general solution

$$\underline{x}^{\rightarrow} = c_1 e^{-t} \cdot \begin{bmatrix} \frac{1}{2} \\ 1 \end{bmatrix} + c_2 e^{2t} \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{c_1 e^{-t}}{2} + 2c_2 e^{2t} \\ c_1 e^{-t} + c_2 e^{2t} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{c_1 e^{-t}}{2} + 2c_2 e^{2t}, x_2(t) = c_1 e^{-t} + c_2 e^{2t} \right\}$$

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 36

```
dsolve([diff(x__1(t),t)=3*x__1(t)-2*x__2(t),diff(x__2(t),t)=2*x__1(t)-2*x__2(t)],singsol=all
```

$$\begin{aligned} x_1(t) &= e^{-t}c_1 + c_2 e^{2t} \\ x_2(t) &= 2e^{-t}c_1 + \frac{c_2 e^{2t}}{2} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.003 (sec). Leaf size: 73

```
DSolve[{x1'[t]==3*x1[t]-2*x2[t],x2'[t]==2*x1[t]-2*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$\begin{aligned} x_1(t) &\rightarrow \frac{1}{3}e^{-t}(c_1(4e^{3t}-1)-2c_2(e^{3t}-1)) \\ x_2(t) &\rightarrow \frac{1}{3}e^{-t}(2c_1(e^{3t}-1)-c_2(e^{3t}-4)) \end{aligned}$$

19.2 problem 2

19.2.1 Solution using Matrix exponential method	3955
19.2.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3956
19.2.3 Maple step by step solution	3961

Internal problem ID [793]

Internal file name [OUTPUT/793_Sunday_June_05_2022_01_49_56_AM_69247619/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x_1'(t) = 5x_1(t) - x_2(t)$$

$$x_2'(t) = 3x_1(t) + x_2(t)$$

19.2.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 5 & -1 \\ 3 & 1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} -\frac{e^{2t}}{2} + \frac{3e^{4t}}{2} & -\frac{e^{4t}}{2} + \frac{e^{2t}}{2} \\ \frac{3e^{4t}}{2} - \frac{3e^{2t}}{2} & \frac{3e^{2t}}{2} - \frac{e^{4t}}{2} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} -\frac{e^{2t}}{2} + \frac{3e^{4t}}{2} & -\frac{e^{4t}}{2} + \frac{e^{2t}}{2} \\ \frac{3e^{4t}}{2} - \frac{3e^{2t}}{2} & \frac{3e^{2t}}{2} - \frac{e^{4t}}{2} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \left(-\frac{e^{2t}}{2} + \frac{3e^{4t}}{2}\right) c_1 + \left(-\frac{e^{4t}}{2} + \frac{e^{2t}}{2}\right) c_2 \\ \left(\frac{3e^{4t}}{2} - \frac{3e^{2t}}{2}\right) c_1 + \left(\frac{3e^{2t}}{2} - \frac{e^{4t}}{2}\right) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \frac{(c_2 - c_1)e^{2t}}{2} + \frac{3(c_1 - \frac{c_2}{3})e^{4t}}{2} \\ \frac{(-3c_1 + 3c_2)e^{2t}}{2} + \frac{3(c_1 - \frac{c_2}{3})e^{4t}}{2} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.2.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 5 & -1 \\ 3 & 1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 5 & -1 \\ 3 & 1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 5 - \lambda & -1 \\ 3 & 1 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 6\lambda + 8 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 2$$

$$\lambda_2 = 4$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
2	1	real eigenvalue
4	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = 2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 5 & -1 \\ 3 & 1 \end{bmatrix} - (2) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 3 & -1 \\ 3 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 3 & -1 & 0 \\ 3 & -1 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} 3 & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 3 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{t}{3}\}$

Hence the solution is

$$\begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix} = \begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} = \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ 3 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 4$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 5 & -1 \\ 3 & 1 \end{bmatrix} - (4) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & -1 \\ 3 & -3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 & -1 & 0 \\ 3 & -3 & 0 \end{array} \right]$$

$$R_2 = R_2 - 3R_1 \implies \left[\begin{array}{cc|c} 1 & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
2	1	1	No	$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$
4	1	1	No	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue 2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{2t} \\ &= \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} e^{2t}\end{aligned}$$

Since eigenvalue 4 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{4t} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{4t}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{e^{2t}}{3} \\ e^{2t} \end{bmatrix} + c_2 \begin{bmatrix} e^{4t} \\ e^{4t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{c_1 e^{2t}}{3} + c_2 e^{4t} \\ c_1 e^{2t} + c_2 e^{4t} \end{bmatrix}$$

The following is the phase plot of the system.

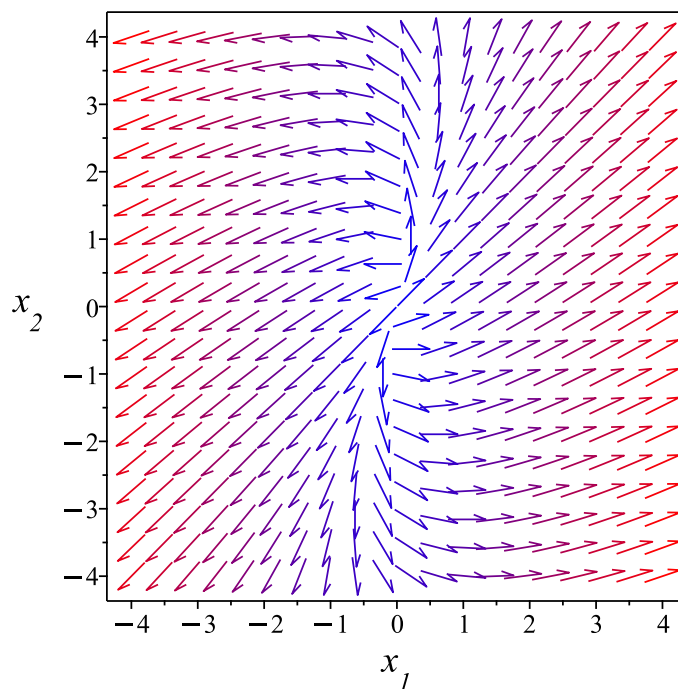


Figure 540: Phase plot

19.2.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 5x_1(t) - x_2(t), x_2'(t) = 3x_1(t) + x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 5 & -1 \\ 3 & 1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 5 & -1 \\ 3 & 1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 5 & -1 \\ 3 & 1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[2, \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} \right], \left[4, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[2, \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{2t} \cdot \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[4, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_2 = e^{4t} \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1 + c_2 \underline{x}^{\rightarrow}_2$$

- Substitute solutions into the general solution

$$\underline{x}^{\rightarrow} = c_1 e^{2t} \cdot \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} + c_2 e^{4t} \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{c_1 e^{2t}}{3} + c_2 e^{4t} \\ c_1 e^{2t} + c_2 e^{4t} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{c_1 e^{2t}}{3} + c_2 e^{4t}, x_2(t) = c_1 e^{2t} + c_2 e^{4t} \right\}$$

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 35

```
dsolve([diff(x__1(t),t)=5*x__1(t)-1*x__2(t),diff(x__2(t),t)=3*x__1(t)+1*x__2(t)],singsol=all
```

$$\begin{aligned} x_1(t) &= c_1 e^{4t} + c_2 e^{2t} \\ x_2(t) &= c_1 e^{4t} + 3c_2 e^{2t} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.003 (sec). Leaf size: 73

```
DSolve[{x1'[t]==5*x1[t]-1*x2[t],x2'[t]==3*x1[t]+1*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$\begin{aligned} x_1(t) &\rightarrow \frac{1}{2} e^{2t} (c_1 (3e^{2t} - 1) - c_2 (e^{2t} - 1)) \\ x_2(t) &\rightarrow \frac{1}{2} e^{2t} (3c_1 (e^{2t} - 1) - c_2 (e^{2t} - 3)) \end{aligned}$$

19.3 problem 3

19.3.1 Solution using Matrix exponential method	3964
19.3.2 Solution using explicit Eigenvalue and Eigenvector method . . .	3965
19.3.3 Maple step by step solution	3970

Internal problem ID [794]

Internal file name [OUTPUT/794_Sunday_June_05_2022_01_49_57_AM_46336794/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 3.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 2x_1(t) - x_2(t) \\x_2'(t) &= 3x_1(t) - 2x_2(t)\end{aligned}$$

19.3.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} -\frac{e^{-t}}{2} + \frac{3e^t}{2} & -\frac{e^t}{2} + \frac{e^{-t}}{2} \\ \frac{3e^t}{2} - \frac{3e^{-t}}{2} & \frac{3e^{-t}}{2} - \frac{e^t}{2} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \left(-\frac{e^{-t}}{2} + \frac{3e^t}{2}\right) c_1 + \left(-\frac{e^t}{2} + \frac{e^{-t}}{2}\right) c_2 \\ \left(\frac{3e^t}{2} - \frac{3e^{-t}}{2}\right) c_1 + \left(\frac{3e^{-t}}{2} - \frac{e^t}{2}\right) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \frac{(c_2 - c_1)e^{-t}}{2} + \frac{3(c_1 - \frac{c_2}{3})e^t}{2} \\ \frac{(-3c_1 + 3c_2)e^{-t}}{2} + \frac{3(c_1 - \frac{c_2}{3})e^t}{2} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.3.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 2 - \lambda & -1 \\ 3 & -2 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1$$

$$\lambda_2 = 1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
1	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 3 & -1 \\ 3 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 3 & -1 & 0 \\ 3 & -1 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} 3 & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 3 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{t}{3}\}$

Hence the solution is

$$\begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix} = \begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{t}{3} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} = \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ 3 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} - (1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & -1 \\ 3 & -3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 & -1 & 0 \\ 3 & -3 & 0 \end{array} \right]$$

$$R_2 = R_2 - 3R_1 \implies \left[\begin{array}{cc|c} 1 & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-1	1	1	No	$\begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$
1	1	1	No	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{-t} \\ &= \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} e^{-t}\end{aligned}$$

Since eigenvalue 1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^t \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^t\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{e^{-t}}{3} \\ e^{-t} \end{bmatrix} + c_2 \begin{bmatrix} e^t \\ e^t \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{c_1 e^{-t}}{3} + c_2 e^t \\ c_1 e^{-t} + c_2 e^t \end{bmatrix}$$

The following is the phase plot of the system.

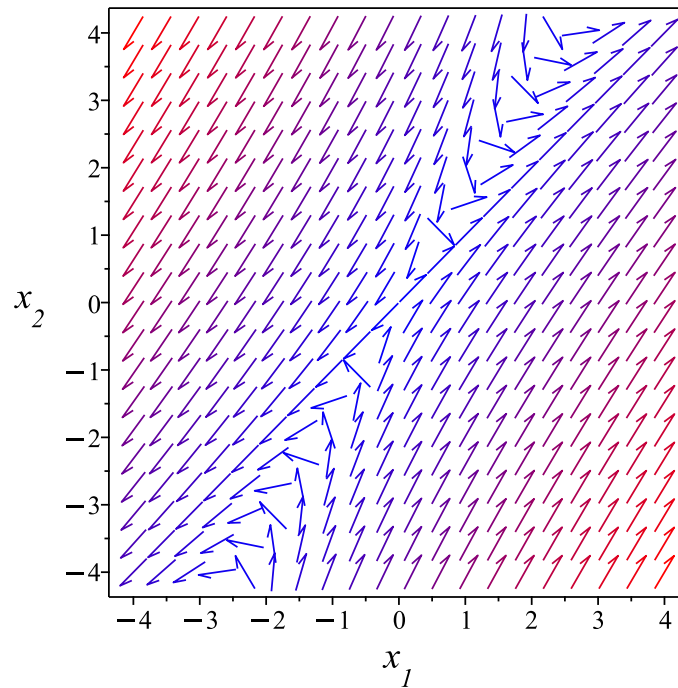


Figure 541: Phase plot

19.3.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 2x_1(t) - x_2(t), x_2'(t) = 3x_1(t) - 2x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 2 & -1 \\ 3 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-1, \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} \right], \left[1, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-1, \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{-t} \cdot \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[1, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_2 = e^t \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1 + c_2 \underline{x}^{\rightarrow}_2$$

- Substitute solutions into the general solution

$$\underline{x}^{\rightarrow} = c_1 e^{-t} \cdot \begin{bmatrix} \frac{1}{3} \\ 1 \end{bmatrix} + c_2 e^t \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{c_1 e^{-t}}{3} + c_2 e^t \\ c_1 e^{-t} + c_2 e^t \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{c_1 e^{-t}}{3} + c_2 e^t, x_2(t) = c_1 e^{-t} + c_2 e^t \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 31

```
dsolve([diff(x__1(t),t)=2*x__1(t)-1*x__2(t),diff(x__2(t),t)=3*x__1(t)-2*x__2(t)],singsol=all
```

$$\begin{aligned} x_1(t) &= c_1 e^t + c_2 e^{-t} \\ x_2(t) &= c_1 e^t + 3c_2 e^{-t} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.003 (sec). Leaf size: 73

```
DSolve[{x1'[t]==2*x1[t]-1*x2[t],x2'[t]==3*x1[t]-2*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$\begin{aligned} x_1(t) &\rightarrow \frac{1}{2} e^{-t} (c_1 (3e^{2t} - 1) - c_2 (e^{2t} - 1)) \\ x_2(t) &\rightarrow \frac{1}{2} e^{-t} (3c_1 (e^{2t} - 1) - c_2 (e^{2t} - 3)) \end{aligned}$$

19.4 problem 4

- 19.4.1 Solution using Matrix exponential method 3973
- 19.4.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3974
- 19.4.3 Maple step by step solution 3979

Internal problem ID [795]

Internal file name [OUTPUT/795_Sunday_June_05_2022_01_49_58_AM_52350690/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 4.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) - 4x_2(t) \\x_2'(t) &= 4x_1(t) - 7x_2(t)\end{aligned}$$

19.4.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-3t}(4t + 1) & -4t e^{-3t} \\ 4t e^{-3t} & e^{-3t}(1 - 4t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{-3t}(4t+1) & -4te^{-3t} \\ 4te^{-3t} & e^{-3t}(1-4t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-3t}(4t+1)c_1 - 4te^{-3t}c_2 \\ 4te^{-3t}c_1 + e^{-3t}(1-4t)c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-3t}(4tc_1 - 4c_2t + c_1) \\ e^{-3t}(4tc_1 - 4c_2t + c_2) \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.4.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 1 - \lambda & -4 \\ 4 & -7 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 6\lambda + 9 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -3$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-3	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -3$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} - (-3) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} 4 & -4 \\ 4 & -4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 4 & -4 & 0 \\ 4 & -4 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} 4 & -4 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 4 & -4 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-3	2	1	Yes	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue -3 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

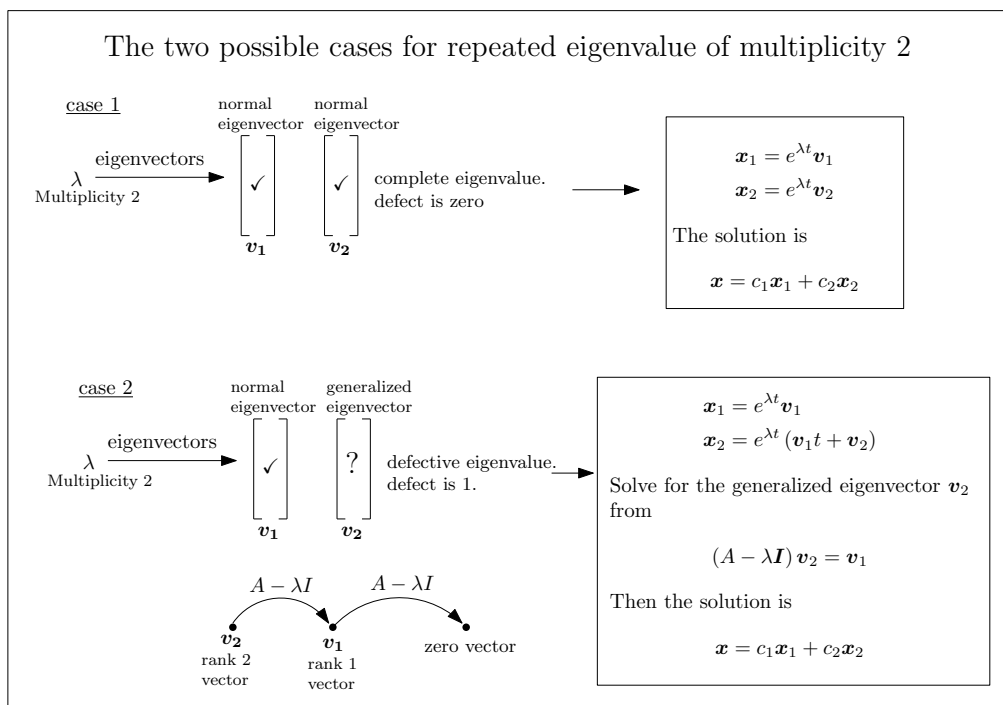


Figure 542: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} - (-3) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} 4 & -4 \\ 4 & -4 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} \frac{5}{4} \\ 1 \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue -3 . Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{-3t} \\ &= \begin{bmatrix} e^{-3t} \\ e^{-3t} \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} 1 \\ 1 \end{bmatrix} t + \begin{bmatrix} \frac{5}{4} \\ 1 \end{bmatrix} \right) e^{-3t} \\ &= \begin{bmatrix} \frac{e^{-3t}(4t+5)}{4} \\ e^{-3t}(t+1) \end{bmatrix}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} e^{-3t} \\ e^{-3t} \end{bmatrix} + c_2 \begin{bmatrix} e^{-3t}(t + \frac{5}{4}) \\ e^{-3t}(t + 1) \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^{-3t}(c_1 + c_2 t + \frac{5}{4}c_2) \\ e^{-3t}(c_2 t + c_1 + c_2) \end{bmatrix}$$

The following is the phase plot of the system.

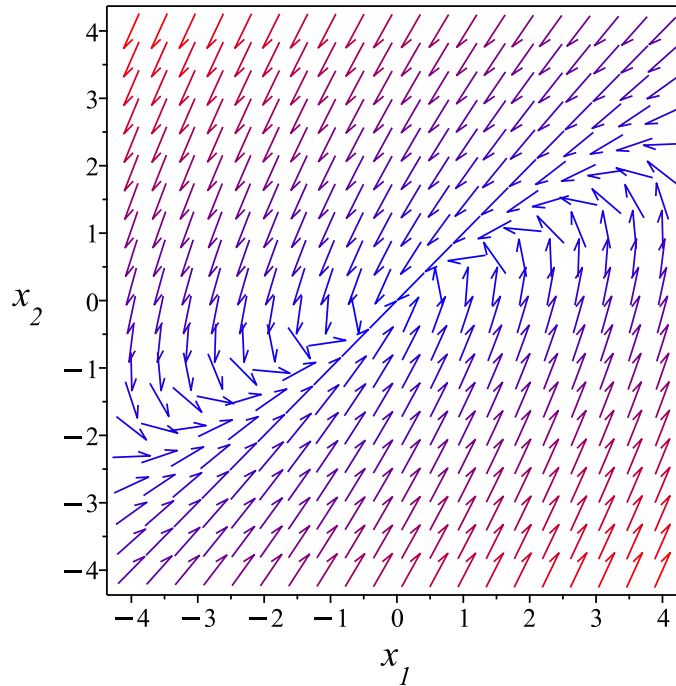


Figure 543: Phase plot

19.4.3 Maple step by step solution

Let's solve

$$[x_1'(t) = x_1(t) - 4x_2(t), x_2'(t) = 4x_1(t) - 7x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-3, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right], \left[-3, \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair, with eigenvalue of algebraic multiplicity 2

$$\left[-3, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right]$$

- First solution from eigenvalue -3

$$\underline{x}^{\rightarrow}_1(t) = e^{-3t} \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- Form of the 2nd homogeneous solution where \vec{p} is to be solved for, $\lambda = -3$ is the eigenvalue, and

$$\underline{x}^{\rightarrow}_2(t) = e^{\lambda t} (t\vec{v} + \vec{p})$$

- Note that the t multiplying \vec{v} makes this solution linearly independent to the 1st solution obtained
- Substitute $\underline{x}^{\rightarrow}_2(t)$ into the homogeneous system

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = (e^{\lambda t} A) \cdot (t\vec{v} + \vec{p})$$

- Use the fact that \vec{v} is an eigenvector of A

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = e^{\lambda t} (\lambda t\vec{v} + A \cdot \vec{p})$$

- Simplify equation

$$\lambda \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Make use of the identity matrix I

$$(\lambda \cdot I) \cdot \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Condition \vec{p} must meet for $\underline{x}^{\rightarrow}_2(t)$ to be a solution to the homogeneous system

$$(A - \lambda \cdot I) \cdot \vec{p} = \vec{v}$$

- Choose \vec{p} to use in the second solution to the homogeneous system from eigenvalue -3

$$\left(\begin{bmatrix} 1 & -4 \\ 4 & -7 \end{bmatrix} - (-3) \cdot \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \cdot \vec{p} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- Choice of \vec{p}

$$\vec{p} = \begin{bmatrix} \frac{1}{4} \\ 0 \end{bmatrix}$$

- Second solution from eigenvalue -3

$$\underline{x}_2(t) = e^{-3t} \cdot \left(t \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix} + \begin{bmatrix} \frac{1}{4} \\ 0 \end{bmatrix} \right)$$

- General solution to the system of ODEs

$$\underline{x} = c_1 \underline{x}_1(t) + c_2 \underline{x}_2(t)$$

- Substitute solutions into the general solution

$$\underline{x} = c_1 e^{-3t} \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix} + e^{-3t} c_2 \cdot \left(t \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix} + \begin{bmatrix} \frac{1}{4} \\ 0 \end{bmatrix} \right)$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^{-3t}(c_1 + c_2 t + \frac{1}{4}c_2) \\ e^{-3t}(c_2 t + c_1) \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = e^{-3t}(c_1 + c_2 t + \frac{1}{4}c_2), x_2(t) = e^{-3t}(c_2 t + c_1)\}$$

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 35

```
dsolve([diff(x__1(t),t)=1*x__1(t)-4*x__2(t),diff(x__2(t),t)=4*x__1(t)-7*x__2(t)],singsol=all
```

$$x_1(t) = e^{-3t}(c_2 t + c_1)$$

$$x_2(t) = \frac{e^{-3t}(4c_2 t + 4c_1 - c_2)}{4}$$

✓ Solution by Mathematica

Time used: 0.002 (sec). Leaf size: 46

```
DSolve[{x1'[t]==1*x1[t]-4*x2[t],x2'[t]==4*x1[t]-7*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$x1(t) \rightarrow e^{-3t}(4c_1t - 4c_2t + c_1)$$

$$x2(t) \rightarrow e^{-3t}(4(c_1 - c_2)t + c_2)$$

19.5 problem 5

- 19.5.1 Solution using Matrix exponential method 3983
- 19.5.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3984
- 19.5.3 Maple step by step solution 3988

Internal problem ID [796]

Internal file name [OUTPUT/796_Sunday_June_05_2022_01_49_59_AM_55876202/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 5.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$x_1'(t) = x_1(t) - 5x_2(t)$$

$$x_2'(t) = x_1(t) - 3x_2(t)$$

19.5.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned} e^{At} &= \begin{bmatrix} e^{-t} \cos(t) + 2e^{-t} \sin(t) & -5e^{-t} \sin(t) \\ e^{-t} \sin(t) & e^{-t} \cos(t) - 2e^{-t} \sin(t) \end{bmatrix} \\ &= \begin{bmatrix} e^{-t}(\cos(t) + 2\sin(t)) & -5e^{-t} \sin(t) \\ e^{-t} \sin(t) & e^{-t}(\cos(t) - 2\sin(t)) \end{bmatrix} \end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{-t}(\cos(t) + 2\sin(t)) & -5e^{-t}\sin(t) \\ e^{-t}\sin(t) & e^{-t}(\cos(t) - 2\sin(t)) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-t}(\cos(t) + 2\sin(t))c_1 - 5e^{-t}\sin(t)c_2 \\ e^{-t}\sin(t)c_1 + e^{-t}(\cos(t) - 2\sin(t))c_2 \end{bmatrix} \\
 &= \begin{bmatrix} ((2c_1 - 5c_2)\sin(t) + c_1\cos(t))e^{-t} \\ ((c_1 - 2c_2)\sin(t) + c_2\cos(t))e^{-t} \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.5.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 1 - \lambda & -5 \\ 1 & -3 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 2\lambda + 2 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1 + i$$

$$\lambda_2 = -1 - i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-1 - i$	1	complex eigenvalue
$-1 + i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1 - i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} - (-1 - i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} 2 + i & -5 \\ 1 & -2 + i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 + i & -5 & 0 \\ 1 & -2 + i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} + \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2 + i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 + i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 - i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = \begin{bmatrix} (2 - i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -1 + i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} - (-1 + i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 - i & -5 \\ 1 & -2 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 1 & -2 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} - \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 - i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 + i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = \begin{bmatrix} (2 + i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-1 + i$	1	1	No	$\begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$
$-1 - i$	1	1	No	$\begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (2+i)e^{(-1+i)t} \\ e^{(-1+i)t} \end{bmatrix} + c_2 \begin{bmatrix} (2-i)e^{(-1-i)t} \\ e^{(-1-i)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (2+i)c_1e^{(-1+i)t} + (2-i)c_2e^{(-1-i)t} \\ c_1e^{(-1+i)t} + c_2e^{(-1-i)t} \end{bmatrix}$$

The following is the phase plot of the system.

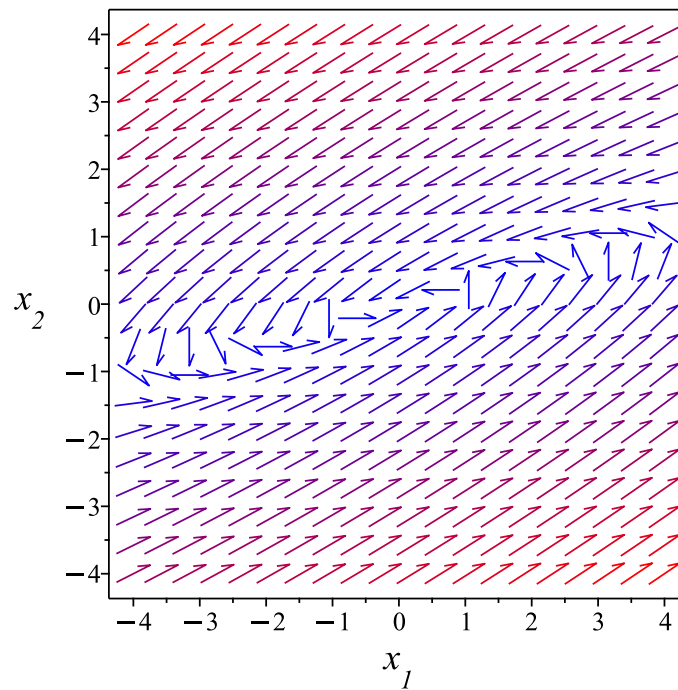


Figure 544: Phase plot

19.5.3 Maple step by step solution

Let's solve

$$[x_1'(t) = x_1(t) - 5x_2(t), x_2'(t) = x_1(t) - 3x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 1 & -5 \\ 1 & -3 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-1 - I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right], \left[-1 + I, \begin{bmatrix} 2 + I \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-1 - I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(-1-I)t} \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{-t} \cdot (\cos(t) - I \sin(t)) \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{-t} \cdot \begin{bmatrix} (2 - I)(\cos(t) - I \sin(t)) \\ \cos(t) - I \sin(t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\vec{x}_{\rightarrow 1}(t) = e^{-t} \cdot \begin{bmatrix} 2 \cos(t) - \sin(t) \\ \cos(t) \end{bmatrix}, \vec{x}_{\rightarrow 2}(t) = e^{-t} \cdot \begin{bmatrix} -\cos(t) - 2 \sin(t) \\ -\sin(t) \end{bmatrix}$$

- General solution to the system of ODEs

$$\vec{x}_{\rightarrow} = c_1 \vec{x}_{\rightarrow 1}(t) + c_2 \vec{x}_{\rightarrow 2}(t)$$

- Substitute solutions into the general solution

$$\vec{x}_{\rightarrow} = c_1 e^{-t} \cdot \begin{bmatrix} 2 \cos(t) - \sin(t) \\ \cos(t) \end{bmatrix} + c_2 e^{-t} \cdot \begin{bmatrix} -\cos(t) - 2 \sin(t) \\ -\sin(t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} 2 \left((c_1 - \frac{c_2}{2}) \cos(t) - \frac{\sin(t)(c_1 + 2c_2)}{2} \right) e^{-t} \\ e^{-t}(-c_2 \sin(t) + c_1 \cos(t)) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = 2 \left((c_1 - \frac{c_2}{2}) \cos(t) - \frac{\sin(t)(c_1 + 2c_2)}{2} \right) e^{-t}, x_2(t) = e^{-t}(-c_2 \sin(t) + c_1 \cos(t)) \right\}$$

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 48

```
dsolve([diff(x__1(t),t)=1*x__1(t)-5*x__2(t),diff(x__2(t),t)=1*x__1(t)-3*x__2(t)],singsol=all
```

$$x_1(t) = e^{-t}(c_1 \sin(t) + c_2 \cos(t))$$

$$x_2(t) = \frac{e^{-t}(-c_1 \cos(t) + c_2 \sin(t) + 2c_1 \sin(t) + 2c_2 \cos(t))}{5}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 54

```
DSolve[{x1'[t]==1*x1[t]-5*x2[t],x2'[t]==1*x1[t]-3*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$x1(t) \rightarrow e^{-t}(c_1 \cos(t) + (2c_1 - 5c_2) \sin(t))$$

$$x2(t) \rightarrow e^{-t}(c_2 \cos(t) + (c_1 - 2c_2) \sin(t))$$

19.6 problem 6

- 19.6.1 Solution using Matrix exponential method 3991
- 19.6.2 Solution using explicit Eigenvalue and Eigenvector method . . . 3992
- 19.6.3 Maple step by step solution 3996

Internal problem ID [797]

Internal file name [OUTPUT/797_Sunday_June_05_2022_01_50_01_AM_59153892/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 6.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 2x_1(t) - 5x_2(t) \\x_2'(t) &= x_1(t) - 2x_2(t)\end{aligned}$$

19.6.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} \cos(t) + 2 \sin(t) & -5 \sin(t) \\ \sin(t) & \cos(t) - 2 \sin(t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} (\cos(t) + 2 \sin(t)) c_1 - 5 \sin(t) c_2 \\ \sin(t) c_1 + (\cos(t) - 2 \sin(t)) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} (2c_1 - 5c_2) \sin(t) + c_1 \cos(t) \\ (c_1 - 2c_2) \sin(t) + c_2 \cos(t) \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.6.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 2 - \lambda & -5 \\ 1 & -2 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-i$	1	complex eigenvalue
i	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - (-i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2+i & -5 \\ 1 & -2+i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2+i & -5 & 0 \\ 1 & -2+i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} + \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2+i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2+i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 - i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = \begin{bmatrix} (2 - i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 - i)t \\ t \end{bmatrix} = \begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} - (i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 - i & -5 \\ 1 & -2 - i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 1 & -2 - i & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{2}{5} - \frac{i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} 2 - i & -5 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 - i & -5 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (2 + i)t\}$

Hence the solution is

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = \begin{bmatrix} (2 + i)t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = t \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (2 + i)t \\ t \end{bmatrix} = \begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
i	1	1	No	$\begin{bmatrix} 2 + i \\ 1 \end{bmatrix}$
$-i$	1	1	No	$\begin{bmatrix} 2 - i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} (2+i)e^{it} \\ e^{it} \end{bmatrix} + c_2 \begin{bmatrix} (2-i)e^{-it} \\ e^{-it} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} (2+i)c_1e^{it} + (2-i)c_2e^{-it} \\ c_1e^{it} + c_2e^{-it} \end{bmatrix}$$

The following is the phase plot of the system.

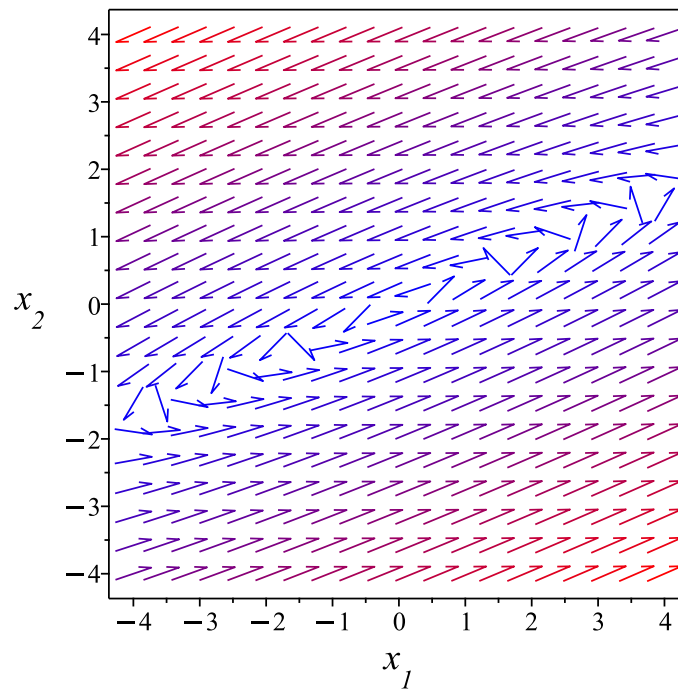


Figure 545: Phase plot

19.6.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 2x_1(t) - 5x_2(t), x_2'(t) = x_1(t) - 2x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 2 & -5 \\ 1 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right], \left[I, \begin{bmatrix} 2 + I \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-I, \begin{bmatrix} 2 - I \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{-It} \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$(\cos(t) - I \sin(t)) \cdot \begin{bmatrix} 2 - I \\ 1 \end{bmatrix}$$

- Simplify expression

$$\begin{bmatrix} (2 - I) (\cos (t) - I \sin (t)) \\ \cos (t) - I \sin (t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\vec{x}_{\underline{1}}(t) = \begin{bmatrix} 2 \cos (t) - \sin (t) \\ \cos (t) \end{bmatrix}, \vec{x}_{\underline{2}}(t) = \begin{bmatrix} -\cos (t) - 2 \sin (t) \\ -\sin (t) \end{bmatrix}$$

- General solution to the system of ODEs

$$\vec{x}_{\underline{}} = c_1 \vec{x}_{\underline{1}}(t) + c_2 \vec{x}_{\underline{2}}(t)$$

- Substitute solutions into the general solution

$$\vec{x}_{\underline{}} = \begin{bmatrix} c_2(-\cos (t) - 2 \sin (t)) + c_1(2 \cos (t) - \sin (t)) \\ -c_2 \sin (t) + c_1 \cos (t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \cos (t) (2c_1 - c_2) - \sin (t) (c_1 + 2c_2) \\ -c_2 \sin (t) + c_1 \cos (t) \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = \cos (t) (2c_1 - c_2) - \sin (t) (c_1 + 2c_2), x_2(t) = -c_2 \sin (t) + c_1 \cos (t)\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 38

```
dsolve([diff(x__1(t),t)=2*x__1(t)-5*x__2(t),diff(x__2(t),t)=1*x__1(t)-2*x__2(t)],singsol=all
```

$$\begin{aligned} x_1(t) &= c_1 \sin (t) + c_2 \cos (t) \\ x_2(t) &= -\frac{c_1 \cos (t)}{5} + \frac{c_2 \sin (t)}{5} + \frac{2c_1 \sin (t)}{5} + \frac{2c_2 \cos (t)}{5} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.004 (sec). Leaf size: 41

```
DSolve[{x1'[t]==2*x1[t]-5*x2[t],x2'[t]==1*x1[t]-2*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$\begin{aligned} x1(t) &\rightarrow c_1(2 \sin (t) + \cos (t)) - 5c_2 \sin (t) \\ x2(t) &\rightarrow c_2 \cos (t) + (c_1 - 2c_2) \sin (t) \end{aligned}$$

19.7 problem 7

19.7.1 Solution using Matrix exponential method	3999
19.7.2 Solution using explicit Eigenvalue and Eigenvector method . . .	4000
19.7.3 Maple step by step solution	4005

Internal problem ID [798]

Internal file name [OUTPUT/798_Sunday_June_05_2022_01_50_03_AM_80643564/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 7.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 3x_1(t) - 2x_2(t) \\x_2'(t) &= 4x_1(t) - x_2(t)\end{aligned}$$

19.7.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned}e^{At} &= \begin{bmatrix} e^t \cos(2t) + e^t \sin(2t) & -e^t \sin(2t) \\ 2e^t \sin(2t) & e^t \cos(2t) - e^t \sin(2t) \end{bmatrix} \\ &= \begin{bmatrix} e^t(\cos(2t) + \sin(2t)) & -e^t \sin(2t) \\ 2e^t \sin(2t) & e^t(\cos(2t) - \sin(2t)) \end{bmatrix}\end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^t(\cos(2t) + \sin(2t)) & -e^t \sin(2t) \\ 2e^t \sin(2t) & e^t(\cos(2t) - \sin(2t)) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^t(\cos(2t) + \sin(2t))c_1 - e^t \sin(2t)c_2 \\ 2e^t \sin(2t)c_1 + e^t(\cos(2t) - \sin(2t))c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^t((-c_2 + c_1)\sin(2t) + c_1 \cos(2t)) \\ e^t(2c_1 - c_2)\sin(2t) + e^t \cos(2t)c_2 \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.7.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} 3 - \lambda & -2 \\ 4 & -1 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 2\lambda + 5 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 1 + 2i$$

$$\lambda_2 = 1 - 2i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$1 + 2i$	1	complex eigenvalue
$1 - 2i$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = 1 - 2i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} - (1 - 2i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 + 2i & -2 \\ 4 & -2 + 2i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 + 2i & -2 & 0 \\ 4 & -2 + 2i & 0 \end{array} \right]$$

$$R_2 = R_2 + (-1 + i)R_1 \implies \left[\begin{array}{cc|c} 2 + 2i & -2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 + 2i & -2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (\frac{1}{2} - \frac{i}{2}) t\}$

Hence the solution is

$$\begin{bmatrix} (\frac{1}{2} - \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} (\frac{1}{2} - \frac{i}{2}) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (\frac{1}{2} - \frac{i}{2}) t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{2} - \frac{i}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (\frac{1}{2} - \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} \frac{1}{2} - \frac{i}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (\frac{1}{2} - \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} 1 - i \\ 2 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 1 + 2i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} - (1 + 2i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 2 - 2i & -2 \\ 4 & -2 - 2i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 - 2i & -2 & 0 \\ 4 & -2 - 2i & 0 \end{array} \right]$$

$$R_2 = R_2 + (-1 - i) R_1 \implies \left[\begin{array}{cc|c} 2 - 2i & -2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 - 2i & -2 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (\frac{1}{2} + \frac{i}{2}) t\}$

Hence the solution is

$$\begin{bmatrix} (\frac{1}{2} + \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} (\frac{1}{2} + \frac{i}{2}) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (\frac{1}{2} + \frac{i}{2}) t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{2} + \frac{i}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (\frac{1}{2} + \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} \frac{1}{2} + \frac{i}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (\frac{1}{2} + \frac{i}{2}) t \\ t \end{bmatrix} = \begin{bmatrix} 1 + i \\ 2 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$1 + 2i$	1	1	No	$\begin{bmatrix} \frac{1}{2} + \frac{i}{2} \\ 1 \end{bmatrix}$
$1 - 2i$	1	1	No	$\begin{bmatrix} \frac{1}{2} - \frac{i}{2} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \left(\frac{1}{2} + \frac{i}{2}\right) e^{(1+2i)t} \\ e^{(1+2i)t} \end{bmatrix} + c_2 \begin{bmatrix} \left(\frac{1}{2} - \frac{i}{2}\right) e^{(1-2i)t} \\ e^{(1-2i)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \left(\frac{1}{2} + \frac{i}{2}\right) c_1 e^{(1+2i)t} + \left(\frac{1}{2} - \frac{i}{2}\right) c_2 e^{(1-2i)t} \\ c_1 e^{(1+2i)t} + c_2 e^{(1-2i)t} \end{bmatrix}$$

The following is the phase plot of the system.

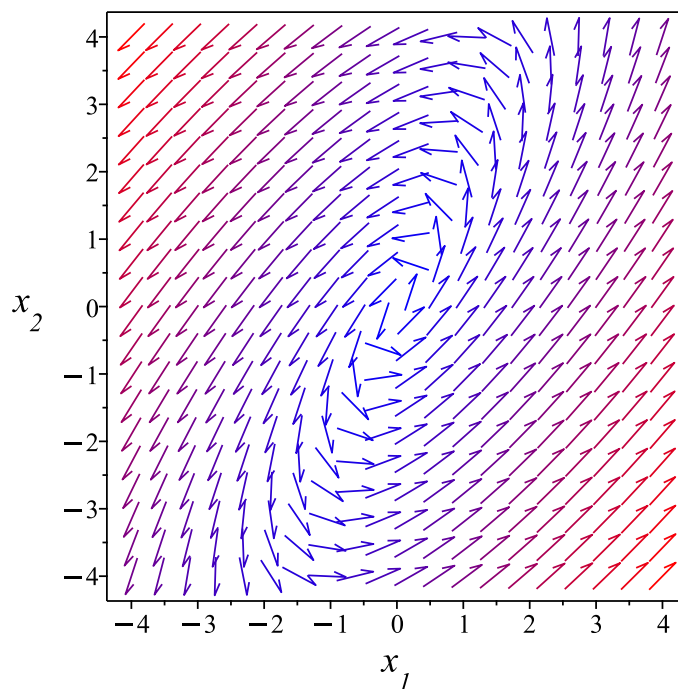


Figure 546: Phase plot

19.7.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 3x_1(t) - 2x_2(t), x_2'(t) = 4x_1(t) - x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 3 & -2 \\ 4 & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[1 - 2I, \begin{bmatrix} \frac{1}{2} - \frac{I}{2} \\ 1 \end{bmatrix} \right], \left[1 + 2I, \begin{bmatrix} \frac{1}{2} + \frac{I}{2} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[1 - 2I, \begin{bmatrix} \frac{1}{2} - \frac{I}{2} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(1-2I)t} \cdot \begin{bmatrix} \frac{1}{2} - \frac{I}{2} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^t \cdot (\cos(2t) - I \sin(2t)) \cdot \begin{bmatrix} \frac{1}{2} - \frac{I}{2} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^t \cdot \begin{bmatrix} \left(\frac{1}{2} - \frac{I}{2}\right) (\cos(2t) - I \sin(2t)) \\ \cos(2t) - I \sin(2t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\underline{x}^{\rightarrow}_1(t) = e^t \cdot \begin{bmatrix} \frac{\cos(2t)}{2} - \frac{\sin(2t)}{2} \\ \cos(2t) \end{bmatrix}, \underline{x}^{\rightarrow}_2(t) = e^t \cdot \begin{bmatrix} -\frac{\sin(2t)}{2} - \frac{\cos(2t)}{2} \\ -\sin(2t) \end{bmatrix} \right]$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1(t) + c_2 \underline{x}^{\rightarrow}_2(t)$$

- Substitute solutions into the general solution

$$\vec{x} = c_1 e^t \cdot \begin{bmatrix} \frac{\cos(2t)}{2} - \frac{\sin(2t)}{2} \\ \cos(2t) \end{bmatrix} + c_2 e^t \cdot \begin{bmatrix} -\frac{\sin(2t)}{2} - \frac{\cos(2t)}{2} \\ -\sin(2t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{e^t((c_1 - c_2)\cos(2t) - \sin(2t)(c_1 + c_2))}{2} \\ e^t(-c_2 \sin(2t) + c_1 \cos(2t)) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{e^t((c_1 - c_2)\cos(2t) - \sin(2t)(c_1 + c_2))}{2}, x_2(t) = e^t(-c_2 \sin(2t) + c_1 \cos(2t)) \right\}$$

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 56

```
dsolve([diff(x__1(t),t)=3*x__1(t)-2*x__2(t),diff(x__2(t),t)=4*x__1(t)-1*x__2(t)],singsol=all
```

$$\begin{aligned} x_1(t) &= e^t(c_1 \sin(2t) + c_2 \cos(2t)) \\ x_2(t) &= -e^t(c_1 \cos(2t) - c_2 \cos(2t) - c_1 \sin(2t) - c_2 \sin(2t)) \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 58

```
DSolve[{x1'[t]==3*x1[t]-2*x2[t],x2'[t]==4*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$\begin{aligned} x1(t) &\rightarrow e^t(c_1 \cos(2t) + (c_1 - c_2) \sin(2t)) \\ x2(t) &\rightarrow e^t(c_2 \cos(2t) + (2c_1 - c_2) \sin(2t)) \end{aligned}$$

19.8 problem 8

19.8.1 Solution using Matrix exponential method	4008
19.8.2 Solution using explicit Eigenvalue and Eigenvector method . . .	4009
19.8.3 Maple step by step solution	4014

Internal problem ID [799]

Internal file name [OUTPUT/799_Sunday_June_05_2022_01_50_04_AM_7663506/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 8.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -x_1(t) - x_2(t) \\x_2'(t) &= -\frac{5x_2(t)}{2}\end{aligned}$$

19.8.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -1 & -1 \\ 0 & -\frac{5}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-t} & -\frac{2e^{-t}}{3} + \frac{2e^{-\frac{5t}{2}}}{3} \\ 0 & e^{-\frac{5t}{2}} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{-t} & -\frac{2e^{-t}}{3} + \frac{2e^{-\frac{5t}{2}}}{3} \\ 0 & e^{-\frac{5t}{2}} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{-t}c_1 + \left(-\frac{2e^{-t}}{3} + \frac{2e^{-\frac{5t}{2}}}{3}\right)c_2 \\ e^{-\frac{5t}{2}}c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \frac{2e^{-\frac{5t}{2}}}{3}c_2 + \left(c_1 - \frac{2c_2}{3}\right)e^{-t} \\ e^{-\frac{5t}{2}}c_2 \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.8.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -1 & -1 \\ 0 & -\frac{5}{2} \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -1 & -1 \\ 0 & -\frac{5}{2} \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -1 - \lambda & -1 \\ 0 & -\frac{5}{2} - \lambda \end{bmatrix}\right) = 0$$

Since the matrix A is triangular matrix, then the determinant is the product of the elements along the diagonal. Therefore the above becomes

$$(-1 - \lambda)\left(-\frac{5}{2} - \lambda\right) = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -\frac{5}{2}$$

$$\lambda_2 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
$-\frac{5}{2}$	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & -1 \\ 0 & -\frac{5}{2} \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 & -1 \\ 0 & -\frac{3}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 0 & -1 & 0 \\ 0 & -\frac{3}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{3R_1}{2} \implies \left[\begin{array}{cc|c} 0 & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 0 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_1\}$ and the leading variables are $\{v_2\}$. Let $v_1 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_2 = 0\}$

Hence the solution is

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = \begin{bmatrix} t \\ 0 \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = t \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -\frac{5}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & -1 \\ 0 & -\frac{5}{2} \end{bmatrix} - \left(-\frac{5}{2}\right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} \frac{3}{2} & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{3}{2} & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} \frac{3}{2} & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = \frac{2t}{3}\}$

Hence the solution is

$$\begin{bmatrix} \frac{2t}{3} \\ t \end{bmatrix} = \begin{bmatrix} \frac{2t}{3} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{2t}{3} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{2}{3} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{2t}{3} \\ t \end{bmatrix} = \begin{bmatrix} \frac{2}{3} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{2t}{3} \\ t \end{bmatrix} = \begin{bmatrix} 2 \\ 3 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-\frac{5}{2}$	1	1	No	$\begin{bmatrix} \frac{2}{3} \\ 1 \end{bmatrix}$
-1	1	1	No	$\begin{bmatrix} 1 \\ 0 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue $-\frac{5}{2}$ is real and distinct then the

corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{-\frac{5t}{2}} \\ &= \begin{bmatrix} \frac{2}{3} \\ 1 \end{bmatrix} e^{-\frac{5t}{2}}\end{aligned}$$

Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-t} \\ &= \begin{bmatrix} 1 \\ 0 \end{bmatrix} e^{-t}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{2e^{-\frac{5t}{2}}}{3} \\ e^{-\frac{5t}{2}} \end{bmatrix} + c_2 \begin{bmatrix} e^{-t} \\ 0 \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{2c_1 e^{-\frac{5t}{2}}}{3} + c_2 e^{-t} \\ c_1 e^{-\frac{5t}{2}} \end{bmatrix}$$

The following is the phase plot of the system.

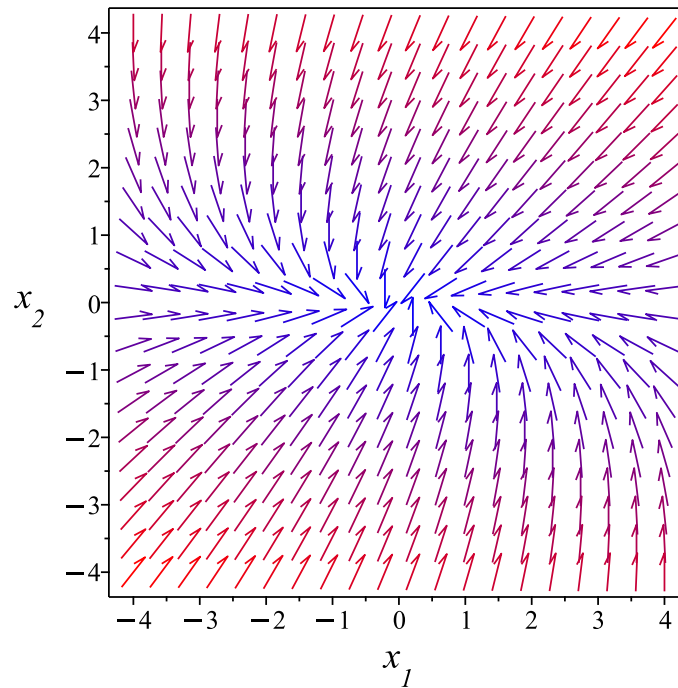


Figure 547: Phase plot

19.8.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = -x_1(t) - x_2(t), x_2'(t) = -\frac{5x_2(t)}{2} \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} -1 & -1 \\ 0 & -\frac{5}{2} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} -1 & -1 \\ 0 & -\frac{5}{2} \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -1 & -1 \\ 0 & -\frac{5}{2} \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[-\frac{5}{2}, \begin{bmatrix} \frac{2}{3} \\ 1 \end{bmatrix} \right], \left[-1, \begin{bmatrix} 1 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-\frac{5}{2}, \begin{bmatrix} \frac{2}{3} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{-\frac{5t}{2}} \cdot \begin{bmatrix} \frac{2}{3} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[-1, \begin{bmatrix} 1 \\ 0 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_2 = e^{-t} \cdot \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1 + c_2 \underline{x}^{\rightarrow}_2$$

- Substitute solutions into the general solution

$$\underline{x}^{\rightarrow} = c_1 e^{-\frac{5t}{2}} \cdot \begin{bmatrix} \frac{2}{3} \\ 1 \end{bmatrix} + c_2 e^{-t} \cdot \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{2c_1 e^{-\frac{5t}{2}}}{3} + c_2 e^{-t} \\ c_1 e^{-\frac{5t}{2}} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{2c_1 e^{-\frac{5t}{2}}}{3} + c_2 e^{-t}, x_2(t) = c_1 e^{-\frac{5t}{2}} \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 28

```
dsolve([diff(x__1(t),t)=-1*x__1(t)-1*x__2(t),diff(x__2(t),t)=0*x__1(t)-25/10*x__2(t)],singsol)
```

$$\begin{aligned} x_1(t) &= \frac{2c_2 e^{-\frac{5t}{2}}}{3} + e^{-t} c_1 \\ x_2(t) &= c_2 e^{-\frac{5t}{2}} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.006 (sec). Leaf size: 47

```
DSolve[{x1'[t]==-1*x1[t]-1*x2[t],x2'[t]==0*x1[t]-25/10*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolutions->True]
```

$$\begin{aligned} x_1(t) &\rightarrow \left(c_1 - \frac{2c_2}{3} \right) e^{-t} + \frac{2}{3} c_2 e^{-5t/2} \\ x_2(t) &\rightarrow c_2 e^{-5t/2} \end{aligned}$$

19.9 problem 9

19.9.1 Solution using Matrix exponential method	4017
19.9.2 Solution using explicit Eigenvalue and Eigenvector method . . .	4018
19.9.3 Maple step by step solution	4023

Internal problem ID [800]

Internal file name [OUTPUT/800_Sunday_June_05_2022_01_50_05_AM_83117480/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 9.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= 3x_1(t) - 4x_2(t) \\x_2'(t) &= x_1(t) - x_2(t)\end{aligned}$$

19.9.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^t(1 + 2t) & -4t e^t \\ t e^t & e^t(1 - 2t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^t(1+2t) & -4te^t \\ te^t & e^t(1-2t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^t(1+2t)c_1 - 4te^tc_2 \\ te^tc_1 + e^t(1-2t)c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^t(2tc_1 - 4c_2t + c_1) \\ e^t(tc_1 - 2c_2t + c_2) \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.9.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 3 - \lambda & -4 \\ 1 & -1 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 2\lambda + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = 1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
1	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = 1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} - (1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} 2 & -4 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 2 & -4 & 0 \\ 1 & -2 & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{R_1}{2} \implies \left[\begin{array}{cc|c} 2 & -4 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 2 & -4 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 2t\}$

Hence the solution is

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = \begin{bmatrix} 2t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = t \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 2t \\ t \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
1	2	1	Yes	$\begin{bmatrix} 2 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue 1 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

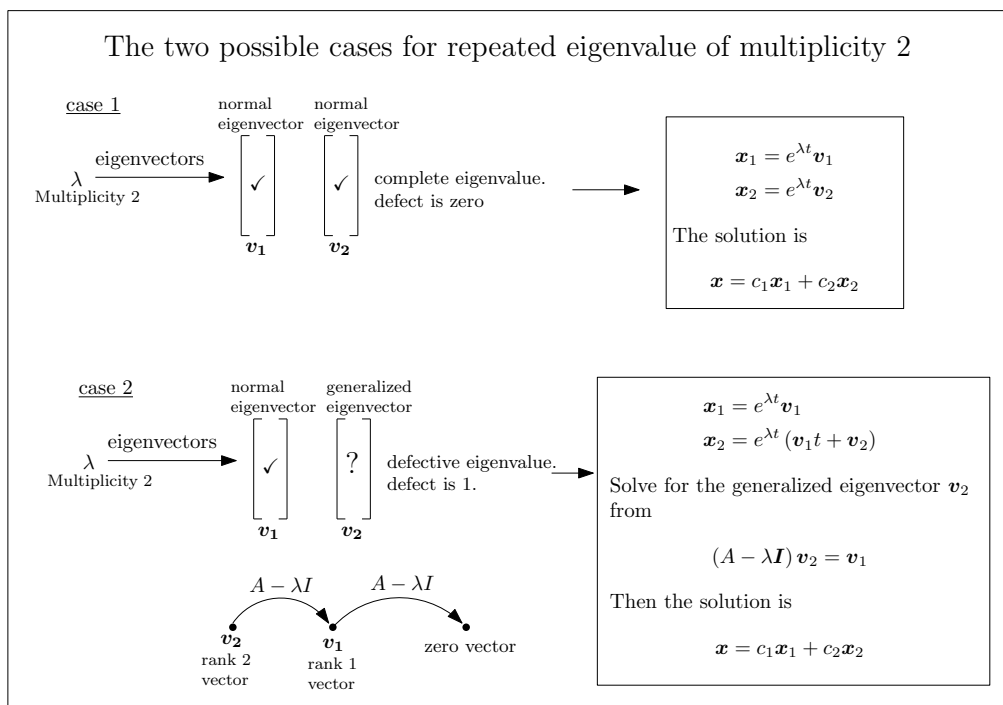


Figure 548: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2, and geometric multiplicity 1, therefore this is defective eigenvalue. The defect is 1. This falls into case 2 shown above. We need to generate the missing additional generalized eigenvector \vec{v}_2 by solving

$$(A - \lambda I) \vec{v}_2 = \vec{v}_1$$

Where \vec{v}_1 is the normal (rank 1) eigenvector found above. Hence we need to solve

$$\left(\begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} - (1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} 2 & -4 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

Solving for \vec{v}_2 gives

$$\vec{v}_2 = \begin{bmatrix} 3 \\ 1 \end{bmatrix}$$

We have found two generalized eigenvectors for eigenvalue 1. Therefore the two basis solution associated with this eigenvalue are

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\lambda t} \\ &= \begin{bmatrix} 2 \\ 1 \end{bmatrix} e^t \\ &= \begin{bmatrix} 2e^t \\ e^t \end{bmatrix}\end{aligned}$$

And

$$\begin{aligned}\vec{x}_2(t) &= (\vec{v}_1 t + \vec{v}_2) e^{\lambda t} \\ &= \left(\begin{bmatrix} 2 \\ 1 \end{bmatrix} t + \begin{bmatrix} 3 \\ 1 \end{bmatrix} \right) e^t \\ &= \begin{bmatrix} e^t(2t + 3) \\ e^t(t + 1) \end{bmatrix}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} 2e^t \\ e^t \end{bmatrix} + c_2 \begin{bmatrix} e^t(2t + 3) \\ e^t(t + 1) \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} ((2t + 3)c_2 + 2c_1)e^t \\ e^t(c_2 t + c_1 + c_2) \end{bmatrix}$$

The following is the phase plot of the system.

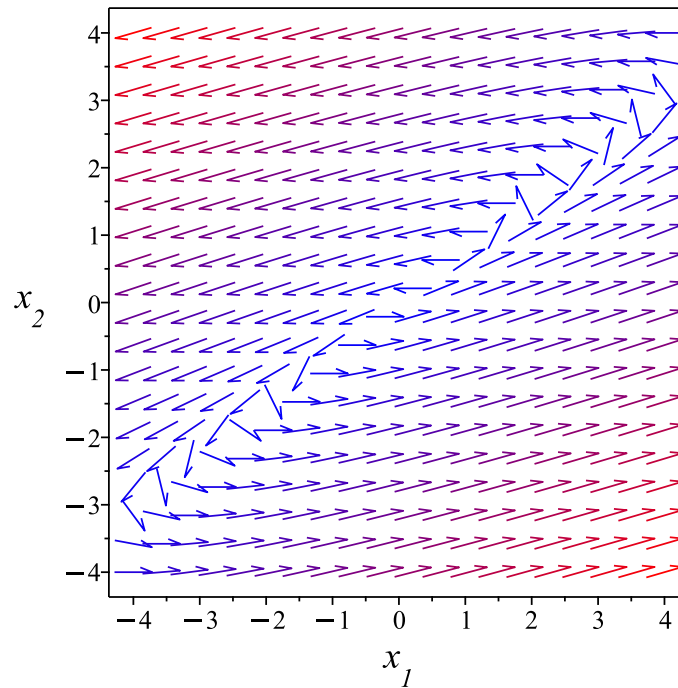


Figure 549: Phase plot

19.9.3 Maple step by step solution

Let's solve

$$[x_1'(t) = 3x_1(t) - 4x_2(t), x_2'(t) = x_1(t) - x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[1, \begin{bmatrix} 2 \\ 1 \end{bmatrix} \right], \left[1, \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair, with eigenvalue of algebraic multiplicity 2

$$\left[1, \begin{bmatrix} 2 \\ 1 \end{bmatrix} \right]$$

- First solution from eigenvalue 1

$$\underline{x}^{\rightarrow}_1(t) = e^t \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

- Form of the 2nd homogeneous solution where \vec{p} is to be solved for, $\lambda = 1$ is the eigenvalue, and

$$\underline{x}^{\rightarrow}_2(t) = e^{\lambda t} (t\vec{v} + \vec{p})$$

- Note that the t multiplying \vec{v} makes this solution linearly independent to the 1st solution obtained
- Substitute $\underline{x}^{\rightarrow}_2(t)$ into the homogeneous system

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = (e^{\lambda t} A) \cdot (t\vec{v} + \vec{p})$$

- Use the fact that \vec{v} is an eigenvector of A

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = e^{\lambda t} (\lambda t\vec{v} + A \cdot \vec{p})$$

- Simplify equation

$$\lambda \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Make use of the identity matrix I

$$(\lambda \cdot I) \cdot \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Condition \vec{p} must meet for $\underline{x}^{\rightarrow}_2(t)$ to be a solution to the homogeneous system

$$(A - \lambda \cdot I) \cdot \vec{p} = \vec{v}$$

- Choose \vec{p} to use in the second solution to the homogeneous system from eigenvalue 1

$$\left(\begin{bmatrix} 3 & -4 \\ 1 & -1 \end{bmatrix} - 1 \cdot \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \cdot \vec{p} = \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$

- Choice of \vec{p}

$$\vec{p} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

- Second solution from eigenvalue 1

$$\underline{x}_2(t) = e^t \cdot \left(t \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \end{bmatrix} \right)$$

- General solution to the system of ODEs

$$\underline{x} = c_1 \underline{x}_1(t) + c_2 \underline{x}_2(t)$$

- Substitute solutions into the general solution

$$\underline{x} = c_1 e^t \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix} + c_2 e^t \cdot \left(t \cdot \begin{bmatrix} 2 \\ 1 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \end{bmatrix} \right)$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} e^t(2c_2t + 2c_1 + c_2) \\ e^t(c_2t + c_1) \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = e^t(2c_2t + 2c_1 + c_2), x_2(t) = e^t(c_2t + c_1)\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 31

```
dsolve([diff(x__1(t),t)=3*x__1(t)-4*x__2(t),diff(x__2(t),t)=1*x__1(t)-1*x__2(t)],singsol=all
```

$$x_1(t) = e^t(c_2t + c_1)$$

$$x_2(t) = \frac{e^t(2c_2t + 2c_1 - c_2)}{4}$$

✓ Solution by Mathematica

Time used: 0.002 (sec). Leaf size: 41

```
DSolve[{x1'[t]==3*x1[t]-4*x2[t],x2'[t]==1*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolu
```

$$x1(t) \rightarrow e^t(2c_1t - 4c_2t + c_1)$$

$$x2(t) \rightarrow e^t((c_1 - 2c_2)t + c_2)$$

19.10 problem 10

19.10.1 Solution using Matrix exponential method	4027
19.10.2 Solution using explicit Eigenvalue and Eigenvector method . . .	4028
19.10.3 Maple step by step solution	4033

Internal problem ID [801]

Internal file name [OUTPUT/801_Sunday_June_05_2022_01_50_06_AM_94844149/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 10.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) + 2x_2(t) \\x_2'(t) &= -5x_1(t)\end{aligned}$$

19.10.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & 2 \\ -5 & 0 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{\frac{t}{2}} \cos\left(\frac{\sqrt{39}t}{2}\right) + \frac{\sqrt{39}e^{\frac{t}{2}} \sin\left(\frac{\sqrt{39}t}{2}\right)}{39} & \frac{4\sqrt{39}e^{\frac{t}{2}} \sin\left(\frac{\sqrt{39}t}{2}\right)}{39} \\ -\frac{10\sqrt{39}e^{\frac{t}{2}} \sin\left(\frac{\sqrt{39}t}{2}\right)}{39} & e^{\frac{t}{2}} \cos\left(\frac{\sqrt{39}t}{2}\right) - \frac{\sqrt{39}e^{\frac{t}{2}} \sin\left(\frac{\sqrt{39}t}{2}\right)}{39} \end{bmatrix}$$

$$= \begin{bmatrix} \frac{e^{\frac{t}{2}} \left(\sqrt{39} \sin\left(\frac{\sqrt{39}t}{2}\right) + 39 \cos\left(\frac{\sqrt{39}t}{2}\right)\right)}{39} & \frac{4\sqrt{39}e^{\frac{t}{2}} \sin\left(\frac{\sqrt{39}t}{2}\right)}{39} \\ -\frac{10\sqrt{39}e^{\frac{t}{2}} \sin\left(\frac{\sqrt{39}t}{2}\right)}{39} & -\frac{e^{\frac{t}{2}} \left(\sqrt{39} \sin\left(\frac{\sqrt{39}t}{2}\right) - 39 \cos\left(\frac{\sqrt{39}t}{2}\right)\right)}{39} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = e^{At} \vec{c}$$

$$= \begin{bmatrix} \frac{e^{\frac{t}{2}} \left(\sqrt{39} \sin\left(\frac{\sqrt{39}t}{2}\right) + 39 \cos\left(\frac{\sqrt{39}t}{2}\right)\right)}{39} & \frac{4\sqrt{39}e^{\frac{t}{2}} \sin\left(\frac{\sqrt{39}t}{2}\right)}{39} \\ -\frac{10\sqrt{39}e^{\frac{t}{2}} \sin\left(\frac{\sqrt{39}t}{2}\right)}{39} & -\frac{e^{\frac{t}{2}} \left(\sqrt{39} \sin\left(\frac{\sqrt{39}t}{2}\right) - 39 \cos\left(\frac{\sqrt{39}t}{2}\right)\right)}{39} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix}$$

$$= \begin{bmatrix} \frac{e^{\frac{t}{2}} \left(\sqrt{39} \sin\left(\frac{\sqrt{39}t}{2}\right) + 39 \cos\left(\frac{\sqrt{39}t}{2}\right)\right) c_1}{39} + \frac{4\sqrt{39}e^{\frac{t}{2}} \sin\left(\frac{\sqrt{39}t}{2}\right) c_2}{39} \\ -\frac{10\sqrt{39}e^{\frac{t}{2}} \sin\left(\frac{\sqrt{39}t}{2}\right) c_1}{39} - \frac{e^{\frac{t}{2}} \left(\sqrt{39} \sin\left(\frac{\sqrt{39}t}{2}\right) - 39 \cos\left(\frac{\sqrt{39}t}{2}\right)\right) c_2}{39} \end{bmatrix}$$

$$= \begin{bmatrix} \frac{\left(\sqrt{39}(c_1 + 4c_2) \sin\left(\frac{\sqrt{39}t}{2}\right) + 39 \cos\left(\frac{\sqrt{39}t}{2}\right) c_1\right) e^{\frac{t}{2}}}{39} \\ -\frac{10e^{\frac{t}{2}} \left(\sqrt{39}(c_1 + \frac{c_2}{10}) \sin\left(\frac{\sqrt{39}t}{2}\right) - \frac{39 \cos\left(\frac{\sqrt{39}t}{2}\right) c_2}{10}\right)}{39} \end{bmatrix}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.10.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & 2 \\ -5 & 0 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 1 & 2 \\ -5 & 0 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 1 - \lambda & 2 \\ -5 & -\lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - \lambda + 10 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = \frac{1}{2} + \frac{i\sqrt{39}}{2}$$

$$\lambda_2 = \frac{1}{2} - \frac{i\sqrt{39}}{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$\frac{1}{2} - \frac{i\sqrt{39}}{2}$	1	complex eigenvalue
$\frac{1}{2} + \frac{i\sqrt{39}}{2}$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = \frac{1}{2} - \frac{i\sqrt{39}}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 2 \\ -5 & 0 \end{bmatrix} - \left(\frac{1}{2} - \frac{i\sqrt{39}}{2} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} \frac{1}{2} + \frac{i\sqrt{39}}{2} & 2 \\ -5 & -\frac{1}{2} + \frac{i\sqrt{39}}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{1}{2} + \frac{i\sqrt{39}}{2} & 2 & 0 \\ -5 & -\frac{1}{2} + \frac{i\sqrt{39}}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 + \frac{5R_1}{\frac{1}{2} + \frac{i\sqrt{39}}{2}} \implies \left[\begin{array}{cc|c} \frac{1}{2} + \frac{i\sqrt{39}}{2} & 2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{cc} \frac{1}{2} + \frac{i\sqrt{39}}{2} & 2 \\ 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\left\{ v_1 = -\frac{4t}{1+i\sqrt{39}} \right\}$

Hence the solution is

$$\begin{bmatrix} -\frac{4t}{1+i\sqrt{39}} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{4t}{1+i\sqrt{39}} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{4t}{1+i\sqrt{39}} \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{4}{1+i\sqrt{39}} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{4}{1+i\sqrt{39}} \\ 1 \end{bmatrix} = \begin{bmatrix} -\frac{4}{1+i\sqrt{39}} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{4}{1+i\sqrt{39}} \\ 1 \end{bmatrix} = \begin{bmatrix} -\frac{4}{1+i\sqrt{39}} \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = \frac{1}{2} + \frac{i\sqrt{39}}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\left[\begin{array}{cc} 1 & 2 \\ -5 & 0 \end{array} \right] - \left(\frac{1}{2} + \frac{i\sqrt{39}}{2} \right) \left[\begin{array}{cc} 1 & 0 \\ 0 & 1 \end{array} \right] \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\left[\begin{array}{cc} \frac{1}{2} - \frac{i\sqrt{39}}{2} & 2 \\ -5 & -\frac{1}{2} - \frac{i\sqrt{39}}{2} \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{1}{2} - \frac{i\sqrt{39}}{2} & 2 & 0 \\ -5 & -\frac{1}{2} - \frac{i\sqrt{39}}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 + \frac{5R_1}{\frac{1}{2} - \frac{i\sqrt{39}}{2}} \implies \left[\begin{array}{cc|c} \frac{1}{2} - \frac{i\sqrt{39}}{2} & 2 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{cc} \frac{1}{2} - \frac{i\sqrt{39}}{2} & 2 \\ 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\left\{ v_1 = \frac{4t}{i\sqrt{39}-1} \right\}$

Hence the solution is

$$\begin{bmatrix} \frac{4t}{i\sqrt{39}-1} \\ t \end{bmatrix} = \begin{bmatrix} \frac{4t}{i\sqrt{39}-1} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{4t}{i\sqrt{39}-1} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{4}{i\sqrt{39}-1} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{4t}{i\sqrt{39}-1} \\ t \end{bmatrix} = \begin{bmatrix} \frac{4}{i\sqrt{39}-1} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{4t}{i\sqrt{39}-1} \\ t \end{bmatrix} = \begin{bmatrix} \frac{4}{i\sqrt{39}-1} \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number

of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$\frac{1}{2} + \frac{i\sqrt{39}}{2}$	1	1	No	$\begin{bmatrix} \frac{2}{-\frac{1}{2} + \frac{i\sqrt{39}}{2}} \\ 1 \end{bmatrix}$
$\frac{1}{2} - \frac{i\sqrt{39}}{2}$	1	1	No	$\begin{bmatrix} \frac{2}{-\frac{1}{2} - \frac{i\sqrt{39}}{2}} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{2e^{\left(\frac{1}{2} + \frac{i\sqrt{39}}{2}\right)t}}{-\frac{1}{2} + \frac{i\sqrt{39}}{2}} \\ e^{\left(\frac{1}{2} + \frac{i\sqrt{39}}{2}\right)t} \end{bmatrix} + c_2 \begin{bmatrix} \frac{2e^{\left(\frac{1}{2} - \frac{i\sqrt{39}}{2}\right)t}}{-\frac{1}{2} - \frac{i\sqrt{39}}{2}} \\ e^{\left(\frac{1}{2} - \frac{i\sqrt{39}}{2}\right)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{i(i-\sqrt{39})c_1 e^{\frac{(1+i\sqrt{39})t}{2}}}{10} + \frac{i(\sqrt{39}+i)c_2 e^{-\frac{(i\sqrt{39}-1)t}{2}}}{10} \\ c_1 e^{\frac{(1+i\sqrt{39})t}{2}} + c_2 e^{-\frac{(i\sqrt{39}-1)t}{2}} \end{bmatrix}$$

The following is the phase plot of the system.

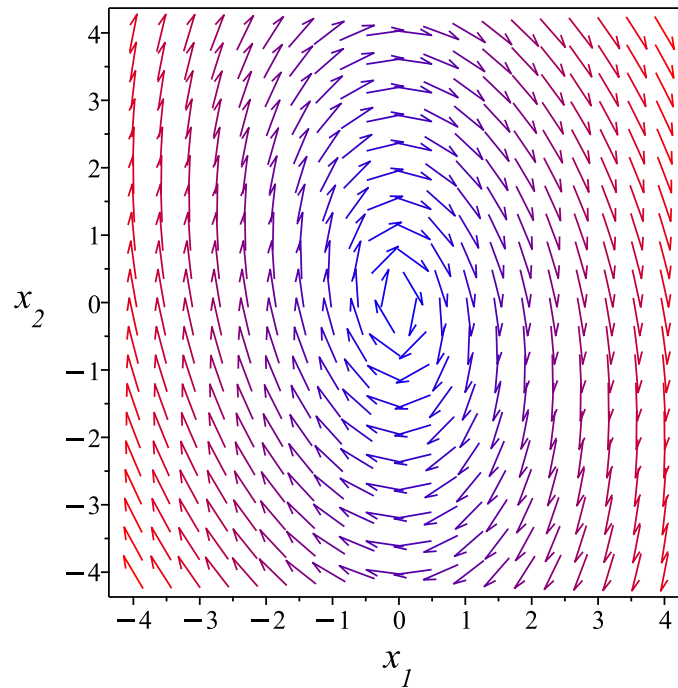


Figure 550: Phase plot

19.10.3 Maple step by step solution

Let's solve

$$[x_1'(t) = x_1(t) + 2x_2(t), x_2'(t) = -5x_1(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 1 & 2 \\ -5 & 0 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 1 & 2 \\ -5 & 0 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 1 & 2 \\ -5 & 0 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}'(t) = A \cdot \underline{x}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[\frac{1}{2} - \frac{I\sqrt{39}}{2}, \begin{bmatrix} \frac{2}{-\frac{1}{2} - \frac{I\sqrt{39}}{2}} \\ 1 \end{bmatrix} \right], \left[\frac{1}{2} + \frac{I\sqrt{39}}{2}, \begin{bmatrix} \frac{2}{-\frac{1}{2} + \frac{I\sqrt{39}}{2}} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[\frac{1}{2} - \frac{I\sqrt{39}}{2}, \begin{bmatrix} \frac{2}{-\frac{1}{2} - \frac{I\sqrt{39}}{2}} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{\left(\frac{1}{2} - \frac{I\sqrt{39}}{2}\right)t} \cdot \begin{bmatrix} \frac{2}{-\frac{1}{2} - \frac{I\sqrt{39}}{2}} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{\frac{t}{2}} \cdot \left(\cos\left(\frac{\sqrt{39}t}{2}\right) - I \sin\left(\frac{\sqrt{39}t}{2}\right) \right) \cdot \begin{bmatrix} \frac{2}{-\frac{1}{2} - \frac{I\sqrt{39}}{2}} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{\frac{t}{2}} \cdot \begin{bmatrix} \frac{2\left(\cos\left(\frac{\sqrt{39}t}{2}\right) - I \sin\left(\frac{\sqrt{39}t}{2}\right)\right)}{-\frac{1}{2} - \frac{I\sqrt{39}}{2}} \\ \cos\left(\frac{\sqrt{39}t}{2}\right) - I \sin\left(\frac{\sqrt{39}t}{2}\right) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\underline{x}_{\rightarrow 1}(t) = e^{\frac{t}{2}} \cdot \begin{bmatrix} -\frac{\cos\left(\frac{\sqrt{39}t}{2}\right)}{10} + \frac{\sqrt{39} \sin\left(\frac{\sqrt{39}t}{2}\right)}{10} \\ \cos\left(\frac{\sqrt{39}t}{2}\right) \end{bmatrix}, \underline{x}_{\rightarrow 2}(t) = e^{\frac{t}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{\sqrt{39}t}{2}\right)\sqrt{39}}{10} + \frac{\sin\left(\frac{\sqrt{39}t}{2}\right)}{10} \\ -\sin\left(\frac{\sqrt{39}t}{2}\right) \end{bmatrix}$$

- General solution to the system of ODEs

$$\underline{x} \rightarrow = c_1 \underline{x}_{\rightarrow 1}(t) + c_2 \underline{x}_{\rightarrow 2}(t)$$

- Substitute solutions into the general solution

$$\vec{x} = c_1 e^{\frac{t}{2}} \cdot \begin{bmatrix} -\frac{\cos\left(\frac{\sqrt{39}t}{2}\right) + \frac{\sqrt{39} \sin\left(\frac{\sqrt{39}t}{2}\right)}{10}}{\cos\left(\frac{\sqrt{39}t}{2}\right)} \end{bmatrix} + c_2 e^{\frac{t}{2}} \cdot \begin{bmatrix} \frac{\cos\left(\frac{\sqrt{39}t}{2}\right)\sqrt{39} + \frac{\sin\left(\frac{\sqrt{39}t}{2}\right)}{10}}{-\sin\left(\frac{\sqrt{39}t}{2}\right)} \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{\left(\left(\sqrt{39}c_2 - c_1\right)\cos\left(\frac{\sqrt{39}t}{2}\right) + \sin\left(\frac{\sqrt{39}t}{2}\right)\left(c_1\sqrt{39} + c_2\right)\right)e^{\frac{t}{2}}}{10}}{e^{\frac{t}{2}}\left(\cos\left(\frac{\sqrt{39}t}{2}\right)c_1 - \sin\left(\frac{\sqrt{39}t}{2}\right)c_2\right)} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{\left(\left(\sqrt{39}c_2 - c_1\right)\cos\left(\frac{\sqrt{39}t}{2}\right) + \sin\left(\frac{\sqrt{39}t}{2}\right)\left(c_1\sqrt{39} + c_2\right)\right)e^{\frac{t}{2}}}{10}, x_2(t) = e^{\frac{t}{2}}\left(\cos\left(\frac{\sqrt{39}t}{2}\right)c_1 - \sin\left(\frac{\sqrt{39}t}{2}\right)c_2\right) \right\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 84

```
dsolve([diff(x__1(t),t)=1*x__1(t)+2*x__2(t),diff(x__2(t),t)=-5*x__1(t)-0*x__2(t)],singsol=all)
```

$$x_1(t) = \frac{e^{\frac{t}{2}}\left(\sin\left(\frac{\sqrt{39}t}{2}\right)\sqrt{39}c_2 - \cos\left(\frac{\sqrt{39}t}{2}\right)\sqrt{39}c_1 - \sin\left(\frac{\sqrt{39}t}{2}\right)c_1 - \cos\left(\frac{\sqrt{39}t}{2}\right)c_2\right)}{10}$$

$$x_2(t) = e^{\frac{t}{2}}\left(\sin\left(\frac{\sqrt{39}t}{2}\right)c_1 + \cos\left(\frac{\sqrt{39}t}{2}\right)c_2\right)$$

✓ Solution by Mathematica

Time used: 0.007 (sec). Leaf size: 54

```
DSolve[{x1'[t]==1*x1[t]+2*x2[t],x2'[t]==-5*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingularSol
```

$$x_1(t) \rightarrow c_1 \cos(3t) + \frac{1}{3}(c_1 + 2c_2) \sin(3t)$$

$$x_2(t) \rightarrow c_2 \cos(3t) - \frac{1}{3}(5c_1 + c_2) \sin(3t)$$

19.11 problem 11

19.11.1 Solution using Matrix exponential method	4036
19.11.2 Solution using explicit Eigenvalue and Eigenvector method . . .	4037
19.11.3 Maple step by step solution	4041

Internal problem ID [802]

Internal file name [OUTPUT/802_Sunday_June_05_2022_01_50_08_AM_89376539/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 11.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x_1'(t) = -x_1(t)$$

$$x_2'(t) = -x_2(t)$$

19.11.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-t} & 0 \\ 0 & e^{-t} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}\vec{x}_h(t) &= e^{At}\vec{c} \\ &= \begin{bmatrix} e^{-t} & 0 \\ 0 & e^{-t} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} e^{-t}c_1 \\ e^{-t}c_2 \end{bmatrix}\end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.11.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} - \lambda\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -1 - \lambda & 0 \\ 0 & -1 - \lambda \end{bmatrix}\right) = 0$$

Since the matrix A is triangular matrix, then the determinant is the product of the elements along the diagonal. Therefore the above becomes

$$(-1 - \lambda)(-1 - \lambda) = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 0 & 0 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_1, v_2\}$ and there are no leading variables. Let $v_1 = t$. Let $v_2 = s$. Since there are two free Variable, we have found two eigenvectors associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ s \end{bmatrix} = \begin{bmatrix} t \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ s \end{bmatrix}$$

$$= t \begin{bmatrix} 1 \\ 0 \end{bmatrix} + s \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

By letting $t = 1$ and $s = 1$ then the above becomes

$$\begin{bmatrix} t \\ s \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

Hence the two eigenvectors associated with this eigenvalue are

$$\left(\begin{bmatrix} 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 1 \end{bmatrix} \right)$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-1	2	2	No	$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. eigenvalue -1 is real and repeated eigenvalue of multiplicity 2. There are two possible cases that can happen. This is illustrated in this diagram

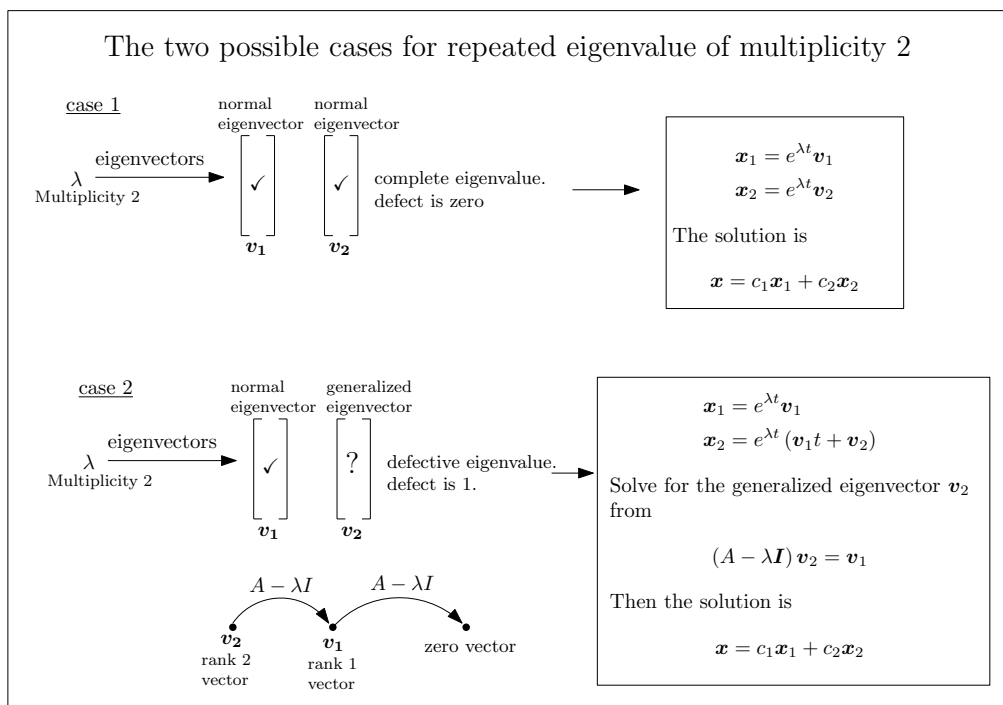


Figure 551: Possible case for repeated λ of multiplicity 2

This eigenvalue has algebraic multiplicity of 2 which is the same as its geometric multiplicity 2, then it is complete eigenvalue and this falls into case 1 shown above. Hence the corresponding eigenvector basis are

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^{-t} \\ &= \begin{bmatrix} 0 \\ 1 \end{bmatrix} e^{-t} \end{aligned}$$

$$\begin{aligned} \vec{x}_2(t) &= \vec{v}_2 e^{-t} \\ &= \begin{bmatrix} 1 \\ 0 \end{bmatrix} e^{-t} \end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} 0 \\ e^{-t} \end{bmatrix} + c_2 \begin{bmatrix} e^{-t} \\ 0 \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} c_2 e^{-t} \\ c_1 e^{-t} \end{bmatrix}$$

The following is the phase plot of the system.

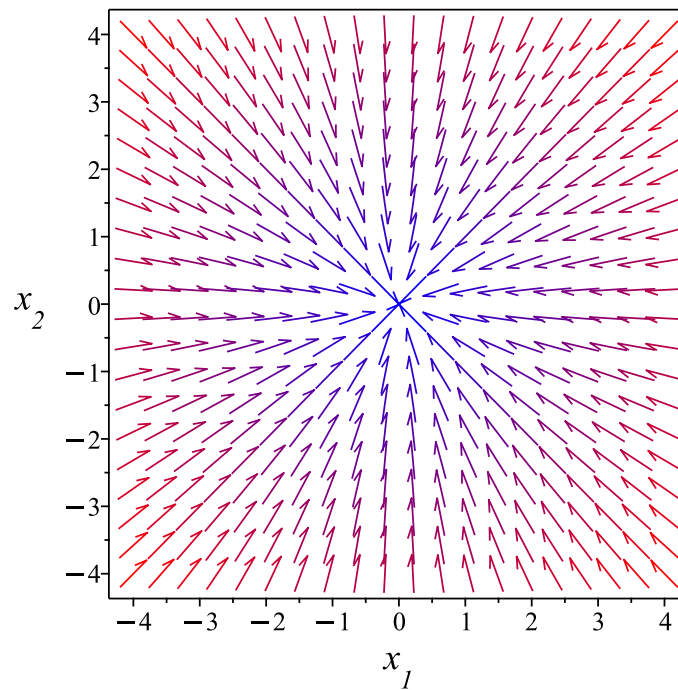


Figure 552: Phase plot

19.11.3 Maple step by step solution

Let's solve

$$[x_1'(t) = -x_1(t), x_2'(t) = -x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-1, \begin{bmatrix} 0 \\ 1 \end{bmatrix} \right], \left[-1, \begin{bmatrix} 1 \\ 0 \end{bmatrix} \right] \right]$$

- Consider eigenpair, with eigenvalue of algebraic multiplicity 2

$$\left[-1, \begin{bmatrix} 0 \\ 1 \end{bmatrix} \right]$$

- First solution from eigenvalue -1

$$\underline{x}^{\rightarrow}_1(t) = e^{-t} \cdot \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

- Form of the 2nd homogeneous solution where \vec{p} is to be solved for, $\lambda = -1$ is the eigenvalue, and

$$\underline{x}^{\rightarrow}_2(t) = e^{\lambda t} (t\vec{v} + \vec{p})$$

- Note that the t multiplying \vec{v} makes this solution linearly independent to the 1st solution obtained

- Substitute $\underline{x}^{\rightarrow}_2(t)$ into the homogeneous system

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = (e^{\lambda t} A) \cdot (t\vec{v} + \vec{p})$$

- Use the fact that \vec{v} is an eigenvector of A

$$\lambda e^{\lambda t} (t\vec{v} + \vec{p}) + e^{\lambda t} \vec{v} = e^{\lambda t} (\lambda t\vec{v} + A \cdot \vec{p})$$

- Simplify equation

$$\lambda \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Make use of the identity matrix I

$$(\lambda \cdot I) \cdot \vec{p} + \vec{v} = A \cdot \vec{p}$$

- Condition \vec{p} must meet for $x_{\underline{\quad}2}(t)$ to be a solution to the homogeneous system

$$(A - \lambda \cdot I) \cdot \vec{p} = \vec{v}$$

- Choose \vec{p} to use in the second solution to the homogeneous system from eigenvalue -1

$$\left(\begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} - (-1) \cdot \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \cdot \vec{p} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

- Choice of \vec{p}

$$\vec{p} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- Second solution from eigenvalue -1

$$x_{\underline{\quad}2}(t) = e^{-t} \cdot \left(t \cdot \begin{bmatrix} 0 \\ 1 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right)$$

- General solution to the system of ODEs

$$x_{\underline{\quad}} = c_1 x_{\underline{\quad}1}(t) + c_2 x_{\underline{\quad}2}(t)$$

- Substitute solutions into the general solution

$$x_{\underline{\quad}} = c_1 e^{-t} \cdot \begin{bmatrix} 0 \\ 1 \end{bmatrix} + c_2 e^{-t} \cdot \left(t \cdot \begin{bmatrix} 0 \\ 1 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right)$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} 0 \\ (c_2 t + c_1) e^{-t} \end{bmatrix}$$

- Solution to the system of ODEs

$$\{x_1(t) = 0, x_2(t) = (c_2 t + c_1) e^{-t}\}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 20

```
dsolve([diff(x__1(t),t)=-1*x__1(t)-0*x__2(t),diff(x__2(t),t)=0*x__1(t)-1*x__2(t)],singsol=all)
```

$$x_1(t) = c_2 e^{-t}$$

$$x_2(t) = e^{-t} c_1$$

✓ Solution by Mathematica

Time used: 0.034 (sec). Leaf size: 65

```
DSolve[{x1'[t]==-1*x1[t]-0*x2[t],x2'[t]==0*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingularSolutions->True]
```

$$x_1(t) \rightarrow c_1 e^{-t}$$

$$x_2(t) \rightarrow c_2 e^{-t}$$

$$x_1(t) \rightarrow c_1 e^{-t}$$

$$x_2(t) \rightarrow 0$$

$$x_1(t) \rightarrow 0$$

$$x_2(t) \rightarrow c_2 e^{-t}$$

$$x_1(t) \rightarrow 0$$

$$x_2(t) \rightarrow 0$$

19.12 problem 12

19.12.1 Solution using Matrix exponential method	4045
19.12.2 Solution using explicit Eigenvalue and Eigenvector method . . .	4046
19.12.3 Maple step by step solution	4051

Internal problem ID [803]

Internal file name [OUTPUT/803_Sunday_June_05_2022_01_50_09_AM_89623807/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 12.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "system of linear ODEs"

Solve

$$\begin{aligned}x_1'(t) &= 2x_1(t) - \frac{5x_2(t)}{2} \\x_2'(t) &= \frac{9x_1(t)}{5} - x_2(t)\end{aligned}$$

19.12.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$\begin{aligned}
 e^{At} &= \begin{bmatrix} e^{\frac{t}{2}} \cos\left(\frac{3t}{2}\right) + e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right) & -\frac{5e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{3} \\ \frac{6e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{5} & e^{\frac{t}{2}} \cos\left(\frac{3t}{2}\right) - e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right) \end{bmatrix} \\
 &= \begin{bmatrix} e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) + \sin\left(\frac{3t}{2}\right)\right) & -\frac{5e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{3} \\ \frac{6e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{5} & e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) - \sin\left(\frac{3t}{2}\right)\right) \end{bmatrix}
 \end{aligned}$$

Therefore the homogeneous solution is

$$\begin{aligned}
 \vec{x}_h(t) &= e^{At} \vec{c} \\
 &= \begin{bmatrix} e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) + \sin\left(\frac{3t}{2}\right)\right) & -\frac{5e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{3} \\ \frac{6e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right)}{5} & e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) - \sin\left(\frac{3t}{2}\right)\right) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\
 &= \begin{bmatrix} e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) + \sin\left(\frac{3t}{2}\right)\right) c_1 - \frac{5e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right) c_2}{3} \\ \frac{6e^{\frac{t}{2}} \sin\left(\frac{3t}{2}\right) c_1}{5} + e^{\frac{t}{2}} \left(\cos\left(\frac{3t}{2}\right) - \sin\left(\frac{3t}{2}\right)\right) c_2 \end{bmatrix} \\
 &= \begin{bmatrix} \left(\left(c_1 - \frac{5c_2}{3}\right) \sin\left(\frac{3t}{2}\right) + c_1 \cos\left(\frac{3t}{2}\right)\right) e^{\frac{t}{2}} \\ \frac{e^{\frac{t}{2}} (6c_1 - 5c_2) \sin\left(\frac{3t}{2}\right)}{5} + e^{\frac{t}{2}} \cos\left(\frac{3t}{2}\right) c_2 \end{bmatrix}
 \end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

19.12.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 2 - \lambda & -\frac{5}{2} \\ \frac{9}{5} & -1 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - \lambda + \frac{5}{2} = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = \frac{1}{2} + \frac{3i}{2}$$

$$\lambda_2 = \frac{1}{2} - \frac{3i}{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$\frac{1}{2} - \frac{3i}{2}$	1	complex eigenvalue
$\frac{1}{2} + \frac{3i}{2}$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = \frac{1}{2} - \frac{3i}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} - \left(\frac{1}{2} - \frac{3i}{2} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} \frac{3}{2} + \frac{3i}{2} & -\frac{5}{2} \\ \frac{9}{5} & -\frac{3}{2} + \frac{3i}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{3}{2} + \frac{3i}{2} & -\frac{5}{2} & 0 \\ \frac{9}{5} & -\frac{3}{2} + \frac{3i}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{3}{5} + \frac{3i}{5}\right) R_1 \implies \left[\begin{array}{cc|c} \frac{3}{2} + \frac{3i}{2} & -\frac{5}{2} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{cc} \frac{3}{2} + \frac{3i}{2} & -\frac{5}{2} \\ 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (\frac{5}{6} - \frac{5i}{6}) t\}$

Hence the solution is

$$\begin{bmatrix} (\frac{5}{6} - \frac{5i}{6}) t \\ t \end{bmatrix} = \begin{bmatrix} (\frac{5}{6} - \frac{5i}{6}) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (\frac{5}{6} - \frac{5i}{6}) t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{5}{6} - \frac{5i}{6} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (\frac{5}{6} - \frac{5i}{6}) t \\ t \end{bmatrix} = \begin{bmatrix} \frac{5}{6} - \frac{5i}{6} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (\frac{5}{6} - \frac{5i}{6}) t \\ t \end{bmatrix} = \begin{bmatrix} 5 - 5i \\ 6 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = \frac{1}{2} + \frac{3i}{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\left(\begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} - \left(\frac{1}{2} + \frac{3i}{2} \right) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right)$$

$$\begin{bmatrix} \frac{3}{2} - \frac{3i}{2} & -\frac{5}{2} \\ \frac{9}{5} & -\frac{3}{2} - \frac{3i}{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} \frac{3}{2} - \frac{3i}{2} & -\frac{5}{2} & 0 \\ \frac{9}{5} & -\frac{3}{2} - \frac{3i}{2} & 0 \end{array} \right]$$

$$R_2 = R_2 + \left(-\frac{3}{5} - \frac{3i}{5} \right) R_1 \implies \left[\begin{array}{cc|c} \frac{3}{2} - \frac{3i}{2} & -\frac{5}{2} & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{cc} \frac{3}{2} - \frac{3i}{2} & -\frac{5}{2} \\ 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = (\frac{5}{6} + \frac{5i}{6}) t\}$

Hence the solution is

$$\begin{bmatrix} (\frac{5}{6} + \frac{5i}{6}) t \\ t \end{bmatrix} = \begin{bmatrix} (\frac{5}{6} + \frac{5i}{6}) t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} (\frac{5}{6} + \frac{5i}{6}) t \\ t \end{bmatrix} = t \begin{bmatrix} \frac{5}{6} + \frac{5i}{6} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} (\frac{5}{6} + \frac{5i}{6}) t \\ t \end{bmatrix} = \begin{bmatrix} \frac{5}{6} + \frac{5i}{6} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} (\frac{5}{6} + \frac{5i}{6}) t \\ t \end{bmatrix} = \begin{bmatrix} 5 + 5i \\ 6 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$\frac{1}{2} + \frac{3i}{2}$	1	1	No	$\begin{bmatrix} \frac{5}{6} + \frac{5i}{6} \\ 1 \end{bmatrix}$
$\frac{1}{2} - \frac{3i}{2}$	1	1	No	$\begin{bmatrix} \frac{5}{6} - \frac{5i}{6} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \left(\frac{5}{6} + \frac{5i}{6}\right) e^{\left(\frac{1}{2} + \frac{3i}{2}\right)t} \\ e^{\left(\frac{1}{2} + \frac{3i}{2}\right)t} \end{bmatrix} + c_2 \begin{bmatrix} \left(\frac{5}{6} - \frac{5i}{6}\right) e^{\left(\frac{1}{2} - \frac{3i}{2}\right)t} \\ e^{\left(\frac{1}{2} - \frac{3i}{2}\right)t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \left(\frac{5}{6} + \frac{5i}{6}\right) c_1 e^{\left(\frac{1}{2} + \frac{3i}{2}\right)t} + \left(\frac{5}{6} - \frac{5i}{6}\right) c_2 e^{\left(\frac{1}{2} - \frac{3i}{2}\right)t} \\ c_1 e^{\left(\frac{1}{2} + \frac{3i}{2}\right)t} + c_2 e^{\left(\frac{1}{2} - \frac{3i}{2}\right)t} \end{bmatrix}$$

The following is the phase plot of the system.

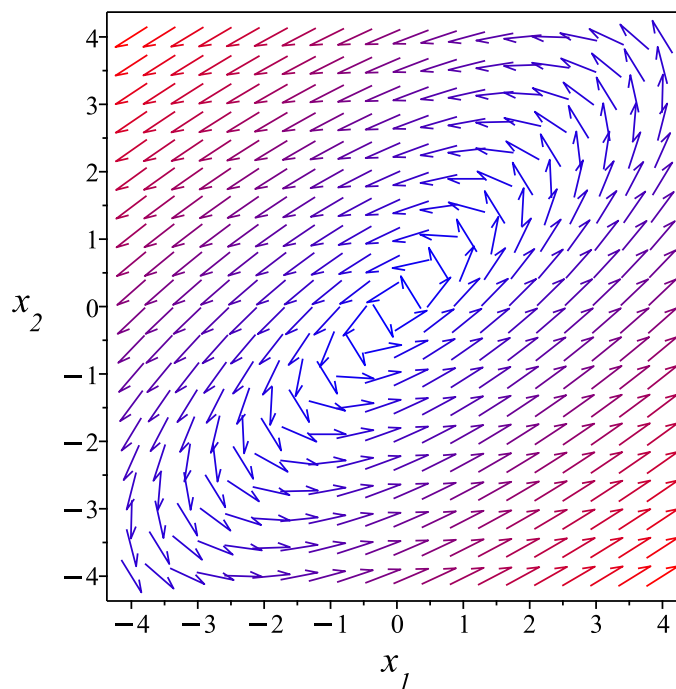


Figure 553: Phase plot

19.12.3 Maple step by step solution

Let's solve

$$\left[x_1'(t) = 2x_1(t) - \frac{5x_2(t)}{2}, x_2'(t) = \frac{9x_1(t)}{5} - x_2(t) \right]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}'(t) = \begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t)$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 2 & -\frac{5}{2} \\ \frac{9}{5} & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t)$$

- To solve the system, find the eigenvalues and eigenvectors of A
- Eigenpairs of A

$$\left[\left[\frac{1}{2} - \frac{3I}{2}, \begin{bmatrix} \frac{5}{6} - \frac{5I}{6} \\ 1 \end{bmatrix} \right], \left[\frac{1}{2} + \frac{3I}{2}, \begin{bmatrix} \frac{5}{6} + \frac{5I}{6} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[\frac{1}{2} - \frac{3I}{2}, \begin{bmatrix} \frac{5}{6} - \frac{5I}{6} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{\left(\frac{1}{2} - \frac{3I}{2}\right)t} \cdot \begin{bmatrix} \frac{5}{6} - \frac{5I}{6} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{\frac{t}{2}} \cdot \left(\cos\left(\frac{3t}{2}\right) - I \sin\left(\frac{3t}{2}\right) \right) \cdot \begin{bmatrix} \frac{5}{6} - \frac{5I}{6} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{\frac{t}{2}} \cdot \begin{bmatrix} \left(\frac{5}{6} - \frac{5I}{6}\right) \left(\cos\left(\frac{3t}{2}\right) - I \sin\left(\frac{3t}{2}\right)\right) \\ \cos\left(\frac{3t}{2}\right) - I \sin\left(\frac{3t}{2}\right) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\underline{x}^{\rightarrow}_1(t) = e^{\frac{t}{2}} \cdot \begin{bmatrix} \frac{5 \cos(\frac{3t}{2})}{6} - \frac{5 \sin(\frac{3t}{2})}{6} \\ \cos\left(\frac{3t}{2}\right) \end{bmatrix}, \underline{x}^{\rightarrow}_2(t) = e^{\frac{t}{2}} \cdot \begin{bmatrix} -\frac{5 \sin(\frac{3t}{2})}{6} - \frac{5 \cos(\frac{3t}{2})}{6} \\ -\sin\left(\frac{3t}{2}\right) \end{bmatrix} \right]$$

- General solution to the system of ODEs

$$\underline{x}^{\rightarrow} = c_1 \underline{x}^{\rightarrow}_1(t) + c_2 \underline{x}^{\rightarrow}_2(t)$$

- Substitute solutions into the general solution

$$\underline{x}^{\rightarrow} = c_1 e^{\frac{t}{2}} \cdot \begin{bmatrix} \frac{5 \cos(\frac{3t}{2})}{6} - \frac{5 \sin(\frac{3t}{2})}{6} \\ \cos(\frac{3t}{2}) \end{bmatrix} + c_2 e^{\frac{t}{2}} \cdot \begin{bmatrix} -\frac{5 \sin(\frac{3t}{2})}{6} - \frac{5 \cos(\frac{3t}{2})}{6} \\ -\sin(\frac{3t}{2}) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{5((c_1 - c_2) \cos(\frac{3t}{2}) - \sin(\frac{3t}{2})(c_1 + c_2)) e^{\frac{t}{2}}}{6} \\ e^{\frac{t}{2}} (c_1 \cos(\frac{3t}{2}) - c_2 \sin(\frac{3t}{2})) \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{5((c_1 - c_2) \cos(\frac{3t}{2}) - \sin(\frac{3t}{2})(c_1 + c_2)) e^{\frac{t}{2}}}{6}, x_2(t) = e^{\frac{t}{2}} (c_1 \cos(\frac{3t}{2}) - c_2 \sin(\frac{3t}{2})) \right\}$$

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 58

```
dsolve([diff(x__1(t),t)=2*x__1(t)-5/2*x__2(t),diff(x__2(t),t)=9/5*x__1(t)-1*x__2(t)],singsol
```

$$x_1(t) = e^{\frac{t}{2}} \left(\sin\left(\frac{3t}{2}\right) c_1 + \cos\left(\frac{3t}{2}\right) c_2 \right)$$

$$x_2(t) = \frac{3 e^{\frac{t}{2}} \left(\sin\left(\frac{3t}{2}\right) c_1 + \sin\left(\frac{3t}{2}\right) c_2 - \cos\left(\frac{3t}{2}\right) c_1 + \cos\left(\frac{3t}{2}\right) c_2 \right)}{5}$$

✓ Solution by Mathematica

Time used: 0.005 (sec). Leaf size: 84

```
DSolve[{x1'[t]==2*x1[t]-5/2*x2[t],x2'[t]==9/5*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingular
```

$$x1(t) \rightarrow \frac{1}{3} e^{t/2} \left(3c_1 \cos\left(\frac{3t}{2}\right) + (3c_1 - 5c_2) \sin\left(\frac{3t}{2}\right) \right)$$

$$x2(t) \rightarrow \frac{1}{5} e^{t/2} \left(5c_2 \cos\left(\frac{3t}{2}\right) + (6c_1 - 5c_2) \sin\left(\frac{3t}{2}\right) \right)$$

19.13 problem 13

19.13.1 Solution using Matrix exponential method	4054
19.13.2 Solution using explicit Eigenvalue and Eigenvector method . . .	4056
19.13.3 Maple step by step solution	4062

Internal problem ID [804]

Internal file name [OUTPUT/804_Sunday_June_05_2022_01_50_11_AM_95295454/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 13.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= x_1(t) + x_2(t) - 2 \\x_2'(t) &= x_1(t) - x_2(t)\end{aligned}$$

19.13.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} -2 \\ 0 \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{(2-\sqrt{2})e^{-\sqrt{2}t}}{4} + \frac{e^{\sqrt{2}t}(2+\sqrt{2})}{4} & \frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} \\ \frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} & \frac{(2+\sqrt{2})e^{-\sqrt{2}t}}{4} - \frac{e^{\sqrt{2}t}(-2+\sqrt{2})}{4} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \frac{(2-\sqrt{2})e^{-\sqrt{2}t}}{4} + \frac{e^{\sqrt{2}t}(2+\sqrt{2})}{4} & \frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} \\ \frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} & \frac{(2+\sqrt{2})e^{-\sqrt{2}t}}{4} - \frac{e^{\sqrt{2}t}(-2+\sqrt{2})}{4} \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} \left(\frac{(2-\sqrt{2})e^{-\sqrt{2}t}}{4} + \frac{e^{\sqrt{2}t}(2+\sqrt{2})}{4} \right) c_1 + \frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}c_2}{4} \\ \frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}c_1}{4} + \left(\frac{(2+\sqrt{2})e^{-\sqrt{2}t}}{4} - \frac{e^{\sqrt{2}t}(-2+\sqrt{2})}{4} \right) c_2 \end{bmatrix} \\ &= \begin{bmatrix} \frac{((-c_1-c_2)\sqrt{2}+2c_1)e^{-\sqrt{2}t}}{4} + \frac{((c_1+c_2)\sqrt{2}+2c_1)e^{\sqrt{2}t}}{4} \\ \frac{(c_2-c_1)\sqrt{2}+2c_2}{4} e^{-\sqrt{2}t} + \frac{e^{\sqrt{2}t}((-c_2+c_1)\sqrt{2}+2c_2)}{4} \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$e^{-At} = (e^{At})^{-1} = \begin{bmatrix} \frac{(2+\sqrt{2})e^{-\sqrt{2}t}}{4} - \frac{e^{\sqrt{2}t}(-2+\sqrt{2})}{4} & -\frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} \\ -\frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} & \frac{(2-\sqrt{2})e^{-\sqrt{2}t}}{4} + \frac{e^{\sqrt{2}t}(2+\sqrt{2})}{4} \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{(2-\sqrt{2})e^{-\sqrt{2}t}}{4} + \frac{e^{\sqrt{2}t}(2+\sqrt{2})}{4} & \frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} \\ \frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} & \frac{(2+\sqrt{2})e^{-\sqrt{2}t}}{4} - \frac{e^{\sqrt{2}t}(-2+\sqrt{2})}{4} \end{bmatrix} \int \begin{bmatrix} \frac{(2+\sqrt{2})e^{-\sqrt{2}t}}{4} - \frac{e^{\sqrt{2}t}(-2+\sqrt{2})}{4} \\ -\frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} \end{bmatrix} \\ &= \begin{bmatrix} \frac{(2-\sqrt{2})e^{-\sqrt{2}t}}{4} + \frac{e^{\sqrt{2}t}(2+\sqrt{2})}{4} & \frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} \\ \frac{(-e^{-\sqrt{2}t}+e^{\sqrt{2}t})\sqrt{2}}{4} & \frac{(2+\sqrt{2})e^{-\sqrt{2}t}}{4} - \frac{e^{\sqrt{2}t}(-2+\sqrt{2})}{4} \end{bmatrix} \begin{bmatrix} \frac{(1+\sqrt{2})e^{-\sqrt{2}t}}{2} - \frac{e^{\sqrt{2}t}(\sqrt{2}-1)}{2} \\ \frac{e^{\sqrt{2}t}}{2} + \frac{e^{-\sqrt{2}t}}{2} \end{bmatrix} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{((-c_1-c_2)\sqrt{2}+2c_1)e^{-\sqrt{2}t}}{4} + \frac{((c_1+c_2)\sqrt{2}+2c_1)e^{\sqrt{2}t}}{4} + 1 \\ \frac{((c_2-c_1)\sqrt{2}+2c_2)e^{-\sqrt{2}t}}{4} + \frac{e^{\sqrt{2}t}((-c_2+c_1)\sqrt{2}+2c_2)}{4} + 1 \end{bmatrix} \end{aligned}$$

19.13.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} -2 \\ 0 \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det \left(\begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) = 0$$

Therefore

$$\det \left(\begin{bmatrix} 1 - \lambda & 1 \\ 1 & -1 - \lambda \end{bmatrix} \right) = 0$$

Which gives the characteristic equation

$$\lambda^2 - 2 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = \sqrt{2}$$
$$\lambda_2 = -\sqrt{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$\sqrt{2}$	1	real eigenvalue
$-\sqrt{2}$	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = \sqrt{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} - (\sqrt{2}) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} 1 - \sqrt{2} & 1 \\ 1 & -1 - \sqrt{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 - \sqrt{2} & 1 & 0 \\ 1 & -1 - \sqrt{2} & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{R_1}{1 - \sqrt{2}} \implies \left[\begin{array}{cc|c} 1 - \sqrt{2} & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{cc} 1 - \sqrt{2} & 1 \\ 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\left\{ v_1 = \frac{t}{\sqrt{2}-1} \right\}$

Hence the solution is

$$\begin{bmatrix} \frac{t}{\sqrt{2}-1} \\ t \end{bmatrix} = \begin{bmatrix} \frac{t}{\sqrt{2}-1} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{t}{\sqrt{2}-1} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{1}{\sqrt{2}-1} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{\sqrt{2}-1} \\ 1 \end{bmatrix} = \begin{bmatrix} \frac{1}{\sqrt{2}-1} \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -\sqrt{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\left(\begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} - (-\sqrt{2}) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \right)$$

$$\begin{bmatrix} 1 + \sqrt{2} & 1 \\ 1 & \sqrt{2} - 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 + \sqrt{2} & 1 & 0 \\ 1 & \sqrt{2} - 1 & 0 \end{array} \right]$$

$$R_2 = R_2 - \frac{R_1}{1 + \sqrt{2}} \implies \left[\begin{array}{cc|c} 1 + \sqrt{2} & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\left[\begin{array}{cc} 1 + \sqrt{2} & 1 \\ 0 & 0 \end{array} \right] \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\left\{ v_1 = -\frac{t}{1+\sqrt{2}} \right\}$

Hence the solution is

$$\begin{bmatrix} -\frac{t}{1+\sqrt{2}} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{t}{1+\sqrt{2}} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{t}{1+\sqrt{2}} \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{1}{1+\sqrt{2}} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{1}{1+\sqrt{2}} \\ 1 \end{bmatrix} = \begin{bmatrix} -\frac{1}{1+\sqrt{2}} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{1}{1+\sqrt{2}} \\ 1 \end{bmatrix} = \begin{bmatrix} -\frac{1}{1+\sqrt{2}} \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$\sqrt{2}$	1	1	No	$\begin{bmatrix} \frac{1}{\sqrt{2}-1} \\ 1 \end{bmatrix}$
$-\sqrt{2}$	1	1	No	$\begin{bmatrix} \frac{1}{-1-\sqrt{2}} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue $\sqrt{2}$ is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_1(t) &= \vec{v}_1 e^{\sqrt{2}t} \\ &= \begin{bmatrix} \frac{1}{\sqrt{2}-1} \\ 1 \end{bmatrix} e^{\sqrt{2}t}\end{aligned}$$

Since eigenvalue $-\sqrt{2}$ is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-\sqrt{2}t} \\ &= \begin{bmatrix} \frac{1}{-1-\sqrt{2}} \\ 1 \end{bmatrix} e^{-\sqrt{2}t}\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{e^{\sqrt{2}t}}{\sqrt{2}-1} \\ e^{\sqrt{2}t} \end{bmatrix} + c_2 \begin{bmatrix} \frac{e^{-\sqrt{2}t}}{-1-\sqrt{2}} \\ e^{-\sqrt{2}t} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} \frac{e^{\sqrt{2}t}}{\sqrt{2}-1} & \frac{e^{-\sqrt{2}t}}{-1-\sqrt{2}} \\ e^{\sqrt{2}t} & e^{-\sqrt{2}t} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} \frac{\sqrt{2}e^{-\sqrt{2}t}}{4} & \frac{\sqrt{2}(\sqrt{2}-1)e^{-\sqrt{2}t}}{4} \\ -\frac{\sqrt{2}e^{\sqrt{2}t}}{4} & \frac{\sqrt{2}e^{\sqrt{2}t}(1+\sqrt{2})}{4} \end{bmatrix}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{e^{\sqrt{2}t}}{\sqrt{2}-1} & \frac{e^{-\sqrt{2}t}}{-1-\sqrt{2}} \\ e^{\sqrt{2}t} & e^{-\sqrt{2}t} \end{bmatrix} \int \begin{bmatrix} \frac{\sqrt{2}e^{-\sqrt{2}t}}{4} & \frac{\sqrt{2}(\sqrt{2}-1)e^{-\sqrt{2}t}}{4} \\ -\frac{\sqrt{2}e^{\sqrt{2}t}}{4} & \frac{\sqrt{2}e^{\sqrt{2}t}(1+\sqrt{2})}{4} \end{bmatrix} \begin{bmatrix} -2 \\ 0 \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{e^{\sqrt{2}t}}{\sqrt{2}-1} & \frac{e^{-\sqrt{2}t}}{-1-\sqrt{2}} \\ e^{\sqrt{2}t} & e^{-\sqrt{2}t} \end{bmatrix} \int \begin{bmatrix} -\frac{\sqrt{2}e^{-\sqrt{2}t}}{2} \\ \frac{\sqrt{2}e^{\sqrt{2}t}}{2} \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{e^{\sqrt{2}t}}{\sqrt{2}-1} & \frac{e^{-\sqrt{2}t}}{-1-\sqrt{2}} \\ e^{\sqrt{2}t} & e^{-\sqrt{2}t} \end{bmatrix} \begin{bmatrix} \frac{e^{-\sqrt{2}t}}{2} \\ \frac{e^{\sqrt{2}t}}{2} \end{bmatrix} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned} \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} \frac{c_1 e^{\sqrt{2}t}}{\sqrt{2}-1} \\ c_1 e^{\sqrt{2}t} \end{bmatrix} + \begin{bmatrix} \frac{c_2 e^{-\sqrt{2}t}}{-1-\sqrt{2}} \\ c_2 e^{-\sqrt{2}t} \end{bmatrix} + \begin{bmatrix} 1 \\ 1 \end{bmatrix} \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} -c_2(\sqrt{2}-1)e^{-\sqrt{2}t} + 1 + c_1(1+\sqrt{2})e^{\sqrt{2}t} \\ c_1 e^{\sqrt{2}t} + c_2 e^{-\sqrt{2}t} + 1 \end{bmatrix}$$

The following is the phase plot of the system.

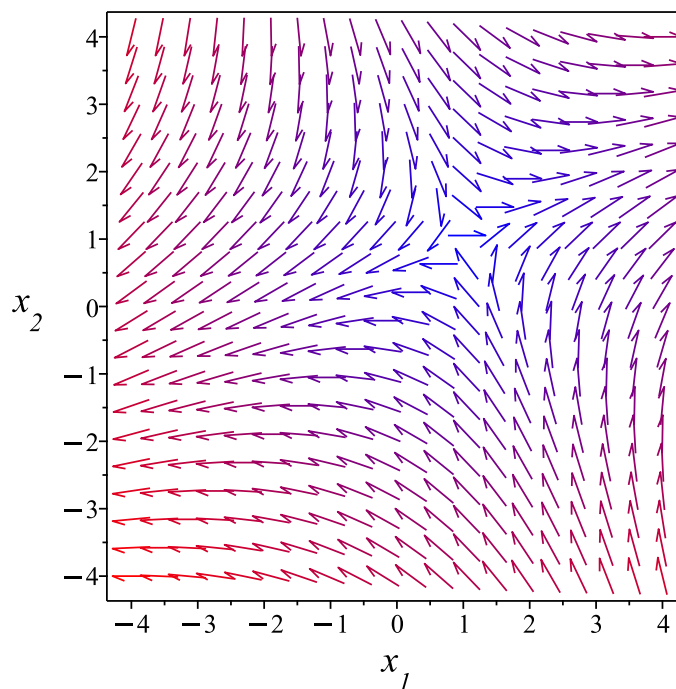


Figure 554: Phase plot

19.13.3 Maple step by step solution

Let's solve

$$[x_1'(t) = x_1(t) + x_2(t) - 2, x_2'(t) = x_1(t) - x_2(t)]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} -2 \\ 0 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} -2 \\ 0 \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} -2 \\ 0 \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow \prime}(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[\sqrt{2}, \begin{bmatrix} \frac{1}{\sqrt{2}-1} \\ 1 \end{bmatrix} \right], \left[-\sqrt{2}, \begin{bmatrix} \frac{1}{-1-\sqrt{2}} \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[\sqrt{2}, \begin{bmatrix} \frac{1}{\sqrt{2}-1} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_1 = e^{\sqrt{2}t} \cdot \begin{bmatrix} \frac{1}{\sqrt{2}-1} \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[-\sqrt{2}, \begin{bmatrix} \frac{1}{-1-\sqrt{2}} \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}_2 = e^{-\sqrt{2}t} \cdot \begin{bmatrix} \frac{1}{-1-\sqrt{2}} \\ 1 \end{bmatrix}$$

- General solution of the system of ODEs can be written in terms of the particular solution $\underline{x}^{\rightarrow}_p$

$$\underline{x}^{\rightarrow}(t) = c_1 \underline{x}^{\rightarrow}_1 + c_2 \underline{x}^{\rightarrow}_2 + \underline{x}^{\rightarrow}_p(t)$$

- Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} \frac{e^{\sqrt{2}t}}{\sqrt{2}-1} & \frac{e^{-\sqrt{2}t}}{-1-\sqrt{2}} \\ e^{\sqrt{2}t} & e^{-\sqrt{2}t} \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix

$$\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} \frac{e^{\sqrt{2}t}}{\sqrt{2}-1} & \frac{e^{-\sqrt{2}t}}{-1-\sqrt{2}} \\ e^{\sqrt{2}t} & e^{-\sqrt{2}t} \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} \frac{1}{\sqrt{2}-1} & \frac{1}{-1-\sqrt{2}} \\ 1 & 1 \end{bmatrix}}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} \frac{((\sqrt{2}-1)e^{-\sqrt{2}t} + e^{\sqrt{2}t}(1+\sqrt{2}))\sqrt{2}}{4} & \frac{(-e^{-\sqrt{2}t} + e^{\sqrt{2}t})\sqrt{2}}{4} \\ \frac{(-e^{-\sqrt{2}t} + e^{\sqrt{2}t})\sqrt{2}}{4} & \frac{(2+\sqrt{2})e^{-\sqrt{2}t}}{4} - \frac{e^{\sqrt{2}t}(-2+\sqrt{2})}{4} \end{bmatrix}$$

- Find a particular solution of the system of ODEs using variation of parameters

- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\underline{x}_p(t) = \Phi(t) \cdot \vec{v}(t)$$

- Take the derivative of the particular solution

$$\underline{x}_p'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t)$

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- Plug $\vec{v}(t)$ into the equation for the particular solution

$$\vec{x}_{\rightarrow p}(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- Plug in the fundamental matrix and the forcing function and compute

$$\vec{x}_{\rightarrow p}(t) = \begin{bmatrix} \frac{(\sqrt{2}-1)e^{-\sqrt{2}t}}{2} + 1 + \frac{(-1-\sqrt{2})e^{\sqrt{2}t}}{2} \\ 1 - \frac{e^{\sqrt{2}t}}{2} - \frac{e^{-\sqrt{2}t}}{2} \end{bmatrix}$$

- Plug particular solution back into general solution

$$\vec{x}_{\rightarrow}(t) = c_1 \vec{x}_{\rightarrow 1} + c_2 \vec{x}_{\rightarrow 2} + \begin{bmatrix} \frac{(\sqrt{2}-1)e^{-\sqrt{2}t}}{2} + 1 + \frac{(-1-\sqrt{2})e^{\sqrt{2}t}}{2} \\ 1 - \frac{e^{\sqrt{2}t}}{2} - \frac{e^{-\sqrt{2}t}}{2} \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{((1-2c_2)\sqrt{2}-1+2c_2)e^{-\sqrt{2}t}}{2} + 1 + \frac{((-1+2c_1)\sqrt{2}+2c_1-1)e^{\sqrt{2}t}}{2} \\ c_1 e^{\sqrt{2}t} + c_2 e^{-\sqrt{2}t} + 1 - \frac{e^{\sqrt{2}t}}{2} - \frac{e^{-\sqrt{2}t}}{2} \end{bmatrix}$$

- Solution to the system of ODEs

$$\begin{cases} x_1(t) = \frac{((1-2c_2)\sqrt{2}-1+2c_2)e^{-\sqrt{2}t}}{2} + 1 + \frac{((-1+2c_1)\sqrt{2}+2c_1-1)e^{\sqrt{2}t}}{2}, x_2(t) = c_1 e^{\sqrt{2}t} + c_2 e^{-\sqrt{2}t} + 1 - \frac{e^{\sqrt{2}t}}{2} \end{cases}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 72

```
dsolve([diff(x__1(t),t)=1*x__1(t)+1*x__2(t)-2,diff(x__2(t),t)=1*x__1(t)-1*x__2(t)],singsol=a
```

$$x_1(t) = e^{\sqrt{2}t}c_2 + e^{-\sqrt{2}t}c_1 + 1$$

$$x_2(t) = \sqrt{2}e^{\sqrt{2}t}c_2 - \sqrt{2}e^{-\sqrt{2}t}c_1 - e^{\sqrt{2}t}c_2 - e^{-\sqrt{2}t}c_1 + 1$$

✓ Solution by Mathematica

Time used: 0.097 (sec). Leaf size: 160

```
DSolve[{x1'[t]==1*x1[t]+1*x2[t]-2,x2'[t]==1*x1[t]-1*x2[t]},{x1[t],x2[t]},t,IncludeSingularSo
```

$$\begin{aligned}x_1(t) &\rightarrow \frac{1}{4}e^{-\sqrt{2}t} \left(4e^{\sqrt{2}t} + \left((2 + \sqrt{2}) c_1 + \sqrt{2}c_2 \right) e^{2\sqrt{2}t} - \left((\sqrt{2} - 2) c_1 \right) - \sqrt{2}c_2 \right) \\x_2(t) &\rightarrow \frac{1}{4}e^{-\sqrt{2}t} \left(4e^{\sqrt{2}t} + \left(\sqrt{2}c_1 - (\sqrt{2} - 2) c_2 \right) e^{2\sqrt{2}t} - \sqrt{2}c_1 + (2 + \sqrt{2}) c_2 \right)\end{aligned}$$

19.14 problem 14

19.14.1 Solution using Matrix exponential method	4067
19.14.2 Solution using explicit Eigenvalue and Eigenvector method . . .	4069
19.14.3 Maple step by step solution	4075

Internal problem ID [805]

Internal file name [OUTPUT/805_Sunday_June_05_2022_01_50_13_AM_61368531/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 14.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -2x_1(t) + x_2(t) - 2 \\x_2'(t) &= x_1(t) - 2x_2(t) + 1\end{aligned}$$

19.14.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} -2 \\ 1 \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} & \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} & \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} & \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} & \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} \end{bmatrix} \begin{bmatrix} C_1 \\ C_2 \end{bmatrix} \\ &= \begin{bmatrix} \left(\frac{e^{-3t}}{2} + \frac{e^{-t}}{2}\right) C_1 + \left(\frac{e^{-t}}{2} - \frac{e^{-3t}}{2}\right) C_2 \\ \left(\frac{e^{-t}}{2} - \frac{e^{-3t}}{2}\right) C_1 + \left(\frac{e^{-3t}}{2} + \frac{e^{-t}}{2}\right) C_2 \end{bmatrix} \\ &= \begin{bmatrix} \frac{(-c_2 + c_1)e^{-3t}}{2} + \frac{e^{-t}(c_1 + c_2)}{2} \\ \frac{(c_2 - c_1)e^{-3t}}{2} + \frac{e^{-t}(c_1 + c_2)}{2} \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \frac{(e^{2t}+1)e^t}{2} & -\frac{(e^{2t}-1)e^t}{2} \\ -\frac{(e^{2t}-1)e^t}{2} & \frac{(e^{2t}+1)e^t}{2} \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} & \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} & \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} \end{bmatrix} \int \begin{bmatrix} \frac{(e^{2t}+1)e^t}{2} & -\frac{(e^{2t}-1)e^t}{2} \\ -\frac{(e^{2t}-1)e^t}{2} & \frac{(e^{2t}+1)e^t}{2} \end{bmatrix} \begin{bmatrix} -2 \\ 1 \end{bmatrix} dt \\ &= \begin{bmatrix} \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} & \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} & \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} \end{bmatrix} \begin{bmatrix} -\frac{e^{3t}}{2} - \frac{e^t}{2} \\ -\frac{e^t}{2} + \frac{e^{3t}}{2} \end{bmatrix} \\ &= \begin{bmatrix} -1 \\ 0 \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} \frac{(-c_2+c_1)e^{-3t}}{2} + \frac{e^{-t}(c_1+c_2)}{2} - 1 \\ \frac{(c_2-c_1)e^{-3t}}{2} + \frac{e^{-t}(c_1+c_2)}{2} \end{bmatrix}\end{aligned}$$

19.14.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} -2 \\ 1 \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -2 - \lambda & 1 \\ 1 & -2 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 4\lambda + 3 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -3$$

$$\lambda_2 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
-3	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -3$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} - (-3) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 & 1 & 0 \\ 1 & 1 & 0 \end{array} \right]$$

$$R_2 = R_2 - R_1 \implies \left[\begin{array}{cc|c} 1 & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -t\}$

Hence the solution is

$$\begin{bmatrix} -t \\ t \end{bmatrix} = \begin{bmatrix} -t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -t \\ t \end{bmatrix} = t \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -t \\ t \end{bmatrix} = \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -1 & 1 & 0 \\ 1 & -1 & 0 \end{array} \right]$$

$$R_2 = R_2 + R_1 \implies \left[\begin{array}{cc|c} -1 & 1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -1 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = t\}$

Hence the solution is

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ t \end{bmatrix} = t \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ t \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-3	1	1	No	$\begin{bmatrix} -1 \\ 1 \end{bmatrix}$
-1	1	1	No	$\begin{bmatrix} 1 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -3 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^{-3t} \\ &= \begin{bmatrix} -1 \\ 1 \end{bmatrix} e^{-3t} \end{aligned}$$

Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-t} \\ &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{-t}\end{aligned}$$

Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} -e^{-3t} \\ e^{-3t} \end{bmatrix} + c_2 \begin{bmatrix} e^{-t} \\ e^{-t} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = \begin{bmatrix} \vec{x}_1 & \vec{x}_2 & \cdots \end{bmatrix}$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} -e^{-3t} & e^{-t} \\ e^{-3t} & e^{-t} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1} \vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -\frac{e^{3t}}{2} & \frac{e^{3t}}{2} \\ \frac{e^t}{2} & \frac{e^t}{2} \end{bmatrix}$$

Hence

$$\begin{aligned}
 \vec{x}_p(t) &= \begin{bmatrix} -e^{-3t} & e^{-t} \\ e^{-3t} & e^{-t} \end{bmatrix} \int \begin{bmatrix} -\frac{e^{3t}}{2} & \frac{e^{3t}}{2} \\ \frac{e^t}{2} & \frac{e^t}{2} \end{bmatrix} \begin{bmatrix} -2 \\ 1 \end{bmatrix} dt \\
 &= \begin{bmatrix} -e^{-3t} & e^{-t} \\ e^{-3t} & e^{-t} \end{bmatrix} \int \begin{bmatrix} \frac{3e^{3t}}{2} \\ -\frac{e^t}{2} \end{bmatrix} dt \\
 &= \begin{bmatrix} -e^{-3t} & e^{-t} \\ e^{-3t} & e^{-t} \end{bmatrix} \begin{bmatrix} \frac{e^{3t}}{2} \\ -\frac{e^t}{2} \end{bmatrix} \\
 &= \begin{bmatrix} -1 \\ 0 \end{bmatrix}
 \end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned}
 \vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\
 \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} -c_1 e^{-3t} \\ c_1 e^{-3t} \end{bmatrix} + \begin{bmatrix} c_2 e^{-t} \\ c_2 e^{-t} \end{bmatrix} + \begin{bmatrix} -1 \\ 0 \end{bmatrix}
 \end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} -c_1 e^{-3t} + c_2 e^{-t} - 1 \\ c_1 e^{-3t} + c_2 e^{-t} \end{bmatrix}$$

The following is the phase plot of the system.

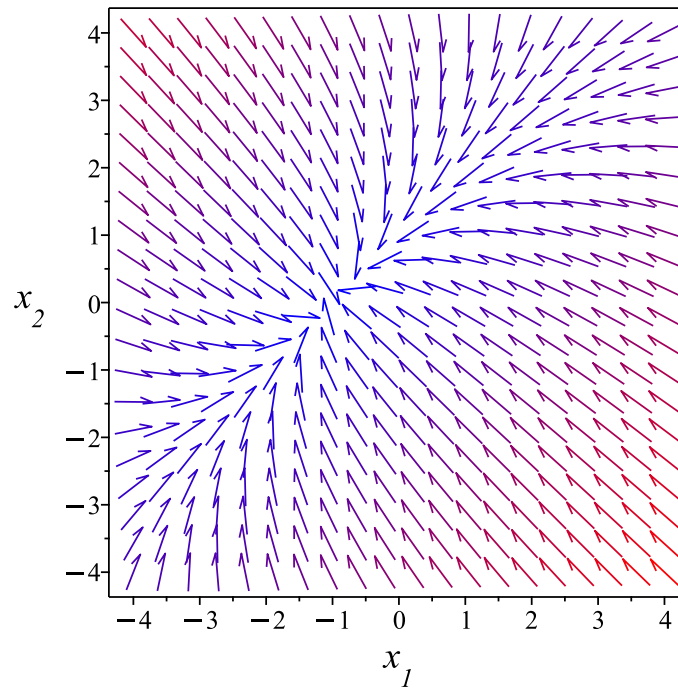


Figure 555: Phase plot

19.14.3 Maple step by step solution

Let's solve

$$[x_1'(t) = -2x_1(t) + x_2(t) - 2, x_2'(t) = x_1(t) - 2x_2(t) + 1]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} -2 \\ 1 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow \prime}(t) = \begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} -2 \\ 1 \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} -2 \\ 1 \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -2 & 1 \\ 1 & -2 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow}{}'(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-3, \begin{bmatrix} -1 \\ 1 \end{bmatrix} \right], \left[-1, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right] \right]$$

- Consider eigenpair

$$\left[-3, \begin{bmatrix} -1 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}{}_{-1} = e^{-3t} \cdot \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

- Consider eigenpair

$$\left[-1, \begin{bmatrix} 1 \\ 1 \end{bmatrix} \right]$$

- Solution to homogeneous system from eigenpair

$$\underline{x}^{\rightarrow}{}_{-2} = e^{-t} \cdot \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- General solution of the system of ODEs can be written in terms of the particular solution $\underline{x}^{\rightarrow}{}_p$

$$\underline{x}^{\rightarrow}(t) = c_1 \underline{x}^{\rightarrow}{}_{-1} + c_2 \underline{x}^{\rightarrow}{}_{-2} + \underline{x}^{\rightarrow}{}_p(t)$$

- Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} -e^{-3t} & e^{-t} \\ e^{-3t} & e^{-t} \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix. $\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} -e^{-3t} & e^{-t} \\ e^{-3t} & e^{-t} \end{bmatrix} \cdot \frac{1}{\begin{bmatrix} -1 & 1 \\ 1 & 1 \end{bmatrix}}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} & \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} & \frac{e^{-3t}}{2} + \frac{e^{-t}}{2} \end{bmatrix}$$

- Find a particular solution of the system of ODEs using variation of parameters

- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\underline{x}_p(t) = \Phi(t) \cdot \vec{v}(t)$$

- Take the derivative of the particular solution

$$\underline{x}_p'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t)$

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- Plug $\vec{v}(t)$ into the equation for the particular solution

$$\vec{x}_{\text{part}}(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- o Plug in the fundamental matrix and the forcing function and compute

$$\vec{x}_{\text{part}}(t) = \begin{bmatrix} \frac{e^{-3t}}{2} - 1 + \frac{e^{-t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \end{bmatrix}$$

- Plug particular solution back into general solution

$$\vec{x}(t) = c_1 \vec{x}_{\text{hom}_1} + c_2 \vec{x}_{\text{hom}_2} + \begin{bmatrix} \frac{e^{-3t}}{2} - 1 + \frac{e^{-t}}{2} \\ \frac{e^{-t}}{2} - \frac{e^{-3t}}{2} \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(-2c_1+1)e^{-3t}}{2} - 1 + \frac{(2c_2+1)e^{-t}}{2} \\ \frac{(-1+2c_1)e^{-3t}}{2} + \frac{(2c_2+1)e^{-t}}{2} \end{bmatrix}$$

- Solution to the system of ODEs

$$\left\{ x_1(t) = \frac{(-2c_1+1)e^{-3t}}{2} - 1 + \frac{(2c_2+1)e^{-t}}{2}, x_2(t) = \frac{(-1+2c_1)e^{-3t}}{2} + \frac{(2c_2+1)e^{-t}}{2} \right\}$$

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 36

```
dsolve([diff(x__1(t),t)=-2*x__1(t)+1*x__2(t)-2,diff(x__2(t),t)=1*x__1(t)-2*x__2(t)+1],singularities)
```

$$\begin{aligned} x_1(t) &= c_2 e^{-t} + c_1 e^{-3t} - 1 \\ x_2(t) &= c_2 e^{-t} - c_1 e^{-3t} \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.022 (sec). Leaf size: 72

```
DSolve[{x1'[t]==-2*x1[t]+1*x2[t]-2,x2'[t]==1*x1[t]-2*x2[t]+1},{x1[t],x2[t]},t,IncludeSingularities->True]
```

$$\begin{aligned} x_1(t) &\rightarrow \frac{1}{2} e^{-3t} (-2e^{3t} + (c_1 + c_2)e^{2t} + c_1 - c_2) \\ x_2(t) &\rightarrow \frac{1}{2} e^{-3t} (c_1(e^{2t} - 1) + c_2(e^{2t} + 1)) \end{aligned}$$

19.15 problem 15

19.15.1 Solution using Matrix exponential method	4079
19.15.2 Solution using explicit Eigenvalue and Eigenvector method . . .	4081
19.15.3 Maple step by step solution	4087

Internal problem ID [806]

Internal file name [OUTPUT/806_Sunday_June_05_2022_01_50_15_AM_35575306/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.1, The Phase Plane: Linear Systems. page 505

Problem number: 15.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$\begin{aligned}x_1'(t) &= -x_1(t) - x_2(t) - 1 \\x_2'(t) &= 2x_1(t) - x_2(t) + 5\end{aligned}$$

19.15.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -1 & -1 \\ 2 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} -1 \\ 5 \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation

of parameters method applied to the fundamental matrix. For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-t} \cos(\sqrt{2}t) & -\frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t)}{2} \\ \sqrt{2}e^{-t} \sin(\sqrt{2}t) & e^{-t} \cos(\sqrt{2}t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned} \vec{x}_h(t) &= e^{At} \vec{c} \\ &= \begin{bmatrix} e^{-t} \cos(\sqrt{2}t) & -\frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t)}{2} \\ \sqrt{2}e^{-t} \sin(\sqrt{2}t) & e^{-t} \cos(\sqrt{2}t) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} \\ &= \begin{bmatrix} e^{-t} \cos(\sqrt{2}t) c_1 - \frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t) c_2}{2} \\ \sqrt{2}e^{-t} \sin(\sqrt{2}t) c_1 + e^{-t} \cos(\sqrt{2}t) c_2 \end{bmatrix} \\ &= \begin{bmatrix} e^{-t} \cos(\sqrt{2}t) c_1 - \frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t) c_2}{2} \\ e^{-t} (\sqrt{2} \sin(\sqrt{2}t) c_1 + \cos(\sqrt{2}t) c_2) \end{bmatrix} \end{aligned}$$

The particular solution given by

$$\vec{x}_p(t) = e^{At} \int e^{-At} \vec{G}(t) dt$$

But

$$\begin{aligned} e^{-At} &= (e^{At})^{-1} \\ &= \begin{bmatrix} \cos(\sqrt{2}t) e^t & \frac{\sin(\sqrt{2}t) \sqrt{2} e^t}{2} \\ -\sin(\sqrt{2}t) \sqrt{2} e^t & \cos(\sqrt{2}t) e^t \end{bmatrix} \end{aligned}$$

Hence

$$\begin{aligned} \vec{x}_p(t) &= \begin{bmatrix} e^{-t} \cos(\sqrt{2}t) & -\frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t)}{2} \\ \sqrt{2}e^{-t} \sin(\sqrt{2}t) & e^{-t} \cos(\sqrt{2}t) \end{bmatrix} \int \begin{bmatrix} \cos(\sqrt{2}t) e^t & \frac{\sin(\sqrt{2}t) \sqrt{2} e^t}{2} \\ -\sin(\sqrt{2}t) \sqrt{2} e^t & \cos(\sqrt{2}t) e^t \end{bmatrix} \begin{bmatrix} -1 \\ 5 \end{bmatrix} dt \\ &= \begin{bmatrix} e^{-t} \cos(\sqrt{2}t) & -\frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t)}{2} \\ \sqrt{2}e^{-t} \sin(\sqrt{2}t) & e^{-t} \cos(\sqrt{2}t) \end{bmatrix} \begin{bmatrix} \frac{e^t (\sqrt{2} \sin(\sqrt{2}t) - 4 \cos(\sqrt{2}t))}{2} \\ e^t (2\sqrt{2} \sin(\sqrt{2}t) + \cos(\sqrt{2}t)) \end{bmatrix} \\ &= \begin{bmatrix} -2 \\ 1 \end{bmatrix} \end{aligned}$$

Hence the complete solution is

$$\begin{aligned}\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\ &= \begin{bmatrix} e^{-t} \cos(\sqrt{2}t) c_1 - \frac{\sqrt{2} e^{-t} \sin(\sqrt{2}t) c_2}{2} - 2 \\ \sqrt{2} e^{-t} \sin(\sqrt{2}t) c_1 + e^{-t} \cos(\sqrt{2}t) c_2 + 1 \end{bmatrix}\end{aligned}$$

19.15.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$$

Or

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -1 & -1 \\ 2 & -1 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} + \begin{bmatrix} -1 \\ 5 \end{bmatrix}$$

Since the system is nonhomogeneous, then the solution is given by

$$\vec{x}(t) = \vec{x}_h(t) + \vec{x}_p(t)$$

Where $\vec{x}_h(t)$ is the homogeneous solution to $\vec{x}'(t) = A\vec{x}(t)$ and $\vec{x}_p(t)$ is a particular solution to $\vec{x}'(t) = A\vec{x}(t) + \vec{G}(t)$. The particular solution will be found using variation of parameters method applied to the fundamental matrix.

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -1 & -1 \\ 2 & -1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -1 - \lambda & -1 \\ 2 & -1 - \lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 2\lambda + 3 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1 + i\sqrt{2}$$

$$\lambda_2 = -1 - i\sqrt{2}$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-1 + i\sqrt{2}$	1	complex eigenvalue
$-1 - i\sqrt{2}$	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1 - i\sqrt{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & -1 \\ 2 & -1 \end{bmatrix} - (-1 - i\sqrt{2}) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} i\sqrt{2} & -1 \\ 2 & i\sqrt{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} i\sqrt{2} & -1 & 0 \\ 2 & i\sqrt{2} & 0 \end{array} \right]$$

$$R_2 = R_2 + i\sqrt{2}R_1 \implies \left[\begin{array}{cc|c} i\sqrt{2} & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} i\sqrt{2} & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\left\{v_1 = -\frac{it\sqrt{2}}{2}\right\}$

Hence the solution is

$$\begin{bmatrix} -\frac{1}{2}t\sqrt{2} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{it\sqrt{2}}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -\frac{1}{2}t\sqrt{2} \\ t \end{bmatrix} = t \begin{bmatrix} -\frac{i\sqrt{2}}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -\frac{1}{2}t\sqrt{2} \\ t \end{bmatrix} = \begin{bmatrix} -\frac{i\sqrt{2}}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} -\frac{1}{2}t\sqrt{2} \\ t \end{bmatrix} = \begin{bmatrix} -i\sqrt{2} \\ 2 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -1 + i\sqrt{2}$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & -1 \\ 2 & -1 \end{bmatrix} - (-1 + i\sqrt{2}) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -i\sqrt{2} & -1 \\ 2 & -i\sqrt{2} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -i\sqrt{2} & -1 & 0 \\ 2 & -i\sqrt{2} & 0 \end{array} \right]$$

$$R_2 = R_2 - i\sqrt{2}R_1 \implies \left[\begin{array}{cc|c} -i\sqrt{2} & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -i\sqrt{2} & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\left\{ v_1 = \frac{it\sqrt{2}}{2} \right\}$

Hence the solution is

$$\begin{bmatrix} \frac{1}{2}t\sqrt{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{it\sqrt{2}}{2} \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} \frac{1}{2}t\sqrt{2} \\ t \end{bmatrix} = t \begin{bmatrix} \frac{i\sqrt{2}}{2} \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} \frac{1}{2}t\sqrt{2} \\ t \end{bmatrix} = \begin{bmatrix} \frac{i\sqrt{2}}{2} \\ 1 \end{bmatrix}$$

Which is normalized to

$$\begin{bmatrix} \frac{1}{2}t\sqrt{2} \\ t \end{bmatrix} = \begin{bmatrix} i\sqrt{2} \\ 2 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
$-1 + i\sqrt{2}$	1	1	No	$\begin{bmatrix} \frac{i\sqrt{2}}{2} \\ 1 \end{bmatrix}$
$-1 - i\sqrt{2}$	1	1	No	$\begin{bmatrix} -\frac{i\sqrt{2}}{2} \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the homogeneous solution is

$$\vec{x}_h(t) = c_1\vec{x}_1(t) + c_2\vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = c_1 \begin{bmatrix} \frac{i\sqrt{2}e^{(-1+i\sqrt{2})t}}{2} \\ e^{(-1+i\sqrt{2})t} \end{bmatrix} + c_2 \begin{bmatrix} -\frac{ie^{(-1-i\sqrt{2})t}\sqrt{2}}{2} \\ e^{(-1-i\sqrt{2})t} \end{bmatrix}$$

Now that we found homogeneous solution above, we need to find a particular solution $\vec{x}_p(t)$. We will use Variation of parameters. The fundamental matrix is

$$\Phi = [\vec{x}_1 \quad \vec{x}_2 \quad \dots]$$

Where \vec{x}_i are the solution basis found above. Therefore the fundamental matrix is

$$\Phi(t) = \begin{bmatrix} \frac{i\sqrt{2}e^{(-1+i\sqrt{2})t}}{2} & -\frac{ie^{(-1-i\sqrt{2})t}\sqrt{2}}{2} \\ e^{(-1+i\sqrt{2})t} & e^{(-1-i\sqrt{2})t} \end{bmatrix}$$

The particular solution is then given by

$$\vec{x}_p(t) = \Phi \int \Phi^{-1}\vec{G}(t) dt$$

But

$$\Phi^{-1} = \begin{bmatrix} -\frac{i\sqrt{2}e^{-(-1+i\sqrt{2})t}}{2} & \frac{e^{-(-1+i\sqrt{2})t}}{2} \\ \frac{i\sqrt{2}e^{(1+i\sqrt{2})t}}{2} & \frac{e^{(1+i\sqrt{2})t}}{2} \end{bmatrix}$$

Hence

$$\begin{aligned}
\vec{x}_p(t) &= \begin{bmatrix} \frac{i\sqrt{2}e^{(-1+i\sqrt{2})t}}{2} & -\frac{ie^{(-1-i\sqrt{2})t}\sqrt{2}}{2} \\ e^{(-1+i\sqrt{2})t} & e^{(-1-i\sqrt{2})t} \end{bmatrix} \int \begin{bmatrix} -\frac{i\sqrt{2}e^{(-1+i\sqrt{2})t}}{2} & \frac{e^{(-1+i\sqrt{2})t}}{2} \\ \frac{i\sqrt{2}e^{(1+i\sqrt{2})t}}{2} & \frac{e^{(1+i\sqrt{2})t}}{2} \end{bmatrix} \begin{bmatrix} -1 \\ 5 \end{bmatrix} dt \\
&= \begin{bmatrix} \frac{i\sqrt{2}e^{(-1+i\sqrt{2})t}}{2} & -\frac{ie^{(-1-i\sqrt{2})t}\sqrt{2}}{2} \\ e^{(-1+i\sqrt{2})t} & e^{(-1-i\sqrt{2})t} \end{bmatrix} \int \begin{bmatrix} \frac{e^{(-1+i\sqrt{2})t}(i\sqrt{2}+5)}{2} \\ -\frac{e^{(1+i\sqrt{2})t}(i\sqrt{2}-5)}{2} \end{bmatrix} dt \\
&= \begin{bmatrix} \frac{i\sqrt{2}e^{(-1+i\sqrt{2})t}}{2} & -\frac{ie^{(-1-i\sqrt{2})t}\sqrt{2}}{2} \\ e^{(-1+i\sqrt{2})t} & e^{(-1-i\sqrt{2})t} \end{bmatrix} \begin{bmatrix} \frac{(1+i\sqrt{2})e^{(-1+i\sqrt{2})t}(i\sqrt{2}+5)}{6} \\ -\frac{e^{(1+i\sqrt{2})t}(i+\sqrt{2})(\sqrt{2}+5i)}{6} \end{bmatrix} \\
&= \begin{bmatrix} -2 \\ 1 \end{bmatrix}
\end{aligned}$$

Now that we found particular solution, the final solution is

$$\begin{aligned}
\vec{x}(t) &= \vec{x}_h(t) + \vec{x}_p(t) \\
\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= \begin{bmatrix} \frac{ic_1\sqrt{2}e^{(-1+i\sqrt{2})t}}{2} \\ c_1e^{(-1+i\sqrt{2})t} \end{bmatrix} + \begin{bmatrix} -\frac{ic_2e^{(-1-i\sqrt{2})t}\sqrt{2}}{2} \\ c_2e^{(-1-i\sqrt{2})t} \end{bmatrix} + \begin{bmatrix} -2 \\ 1 \end{bmatrix}
\end{aligned}$$

Which becomes

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{ic_1\sqrt{2}e^{(-1+i\sqrt{2})t}}{2} - \frac{ic_2e^{(-1-i\sqrt{2})t}\sqrt{2}}{2} - 2 \\ c_1e^{(-1+i\sqrt{2})t} + c_2e^{(-1-i\sqrt{2})t} + 1 \end{bmatrix}$$

The following is the phase plot of the system.

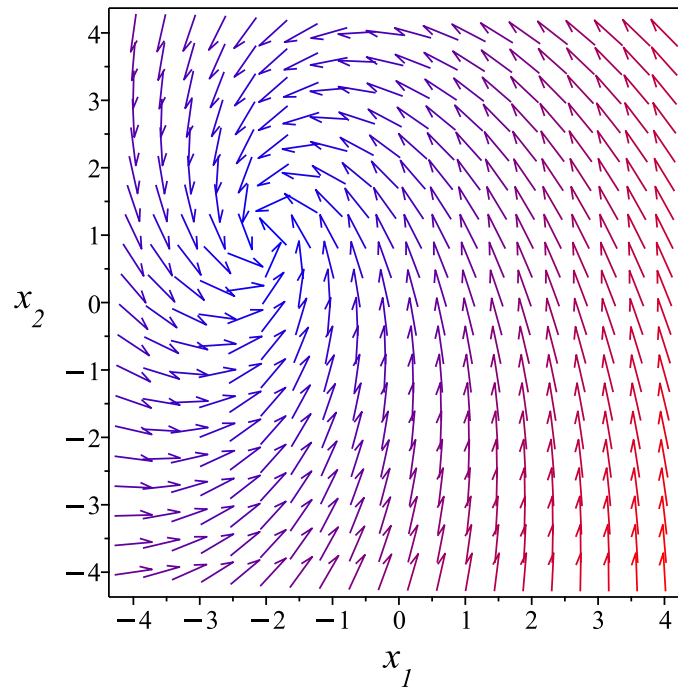


Figure 556: Phase plot

19.15.3 Maple step by step solution

Let's solve

$$[x_1'(t) = -x_1(t) - x_2(t) - 1, x_2'(t) = 2x_1(t) - x_2(t) + 5]$$

- Define vector

$$\underline{x}^{\rightarrow}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$

- Convert system into a vector equation

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} -1 & -1 \\ 2 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} -1 \\ 5 \end{bmatrix}$$

- System to solve

$$\underline{x}^{\rightarrow}{}'(t) = \begin{bmatrix} -1 & -1 \\ 2 & -1 \end{bmatrix} \cdot \underline{x}^{\rightarrow}(t) + \begin{bmatrix} -1 \\ 5 \end{bmatrix}$$

- Define the forcing function

$$\vec{f}(t) = \begin{bmatrix} -1 \\ 5 \end{bmatrix}$$

- Define the coefficient matrix

$$A = \begin{bmatrix} -1 & -1 \\ 2 & -1 \end{bmatrix}$$

- Rewrite the system as

$$\underline{x}^{\rightarrow}{}'(t) = A \cdot \underline{x}^{\rightarrow}(t) + \vec{f}$$

- To solve the system, find the eigenvalues and eigenvectors of A

- Eigenpairs of A

$$\left[\left[-1 - I\sqrt{2}, \begin{bmatrix} -\frac{1}{2}\sqrt{2} \\ 1 \end{bmatrix} \right], \left[-1 + I\sqrt{2}, \begin{bmatrix} \frac{1}{2}\sqrt{2} \\ 1 \end{bmatrix} \right] \right]$$

- Consider complex eigenpair, complex conjugate eigenvalue can be ignored

$$\left[-1 - I\sqrt{2}, \begin{bmatrix} -\frac{1}{2}\sqrt{2} \\ 1 \end{bmatrix} \right]$$

- Solution from eigenpair

$$e^{(-1-I\sqrt{2})t} \cdot \begin{bmatrix} -\frac{1}{2}\sqrt{2} \\ 1 \end{bmatrix}$$

- Use Euler identity to write solution in terms of sin and cos

$$e^{-t} \cdot (\cos(\sqrt{2}t) - I \sin(\sqrt{2}t)) \cdot \begin{bmatrix} -\frac{1}{2}\sqrt{2} \\ 1 \end{bmatrix}$$

- Simplify expression

$$e^{-t} \cdot \begin{bmatrix} -\frac{1}{2}(\cos(\sqrt{2}t) - I \sin(\sqrt{2}t))\sqrt{2} \\ \cos(\sqrt{2}t) - I \sin(\sqrt{2}t) \end{bmatrix}$$

- Both real and imaginary parts are solutions to the homogeneous system

$$\left[\underline{x}^{\rightarrow}_1(t) = e^{-t} \cdot \begin{bmatrix} -\frac{\sqrt{2} \sin(\sqrt{2}t)}{2} \\ \cos(\sqrt{2}t) \end{bmatrix}, \underline{x}^{\rightarrow}_2(t) = e^{-t} \cdot \begin{bmatrix} -\frac{\cos(\sqrt{2}t)\sqrt{2}}{2} \\ -\sin(\sqrt{2}t) \end{bmatrix} \right]$$

- General solution of the system of ODEs can be written in terms of the particular solution $\vec{x}_{\text{part}}(t)$
 $\vec{x}_{\text{part}}(t) = c_1 \vec{x}_{\text{part}1}(t) + c_2 \vec{x}_{\text{part}2}(t) + \vec{x}_{\text{part}p}(t)$

□ Fundamental matrix

- Let $\phi(t)$ be the matrix whose columns are the independent solutions of the homogeneous system

$$\phi(t) = \begin{bmatrix} -\frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t)}{2} & -\frac{\sqrt{2}e^{-t} \cos(\sqrt{2}t)}{2} \\ e^{-t} \cos(\sqrt{2}t) & -e^{-t} \sin(\sqrt{2}t) \end{bmatrix}$$

- The fundamental matrix, $\Phi(t)$ is a normalized version of $\phi(t)$ satisfying $\Phi(0) = I$ where I is the identity matrix

$$\Phi(t) = \phi(t) \cdot \frac{1}{\phi(0)}$$

- Substitute the value of $\phi(t)$ and $\phi(0)$

$$\Phi(t) = \begin{bmatrix} -\frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t)}{2} & -\frac{\sqrt{2}e^{-t} \cos(\sqrt{2}t)}{2} \\ e^{-t} \cos(\sqrt{2}t) & -e^{-t} \sin(\sqrt{2}t) \end{bmatrix} \cdot \begin{bmatrix} 1 & \\ 0 & -\frac{\sqrt{2}}{2} \\ 1 & 0 \end{bmatrix}$$

- Evaluate and simplify to get the fundamental matrix

$$\Phi(t) = \begin{bmatrix} e^{-t} \cos(\sqrt{2}t) & -\frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t)}{2} \\ \sqrt{2}e^{-t} \sin(\sqrt{2}t) & e^{-t} \cos(\sqrt{2}t) \end{bmatrix}$$

□ Find a particular solution of the system of ODEs using variation of parameters

- Let the particular solution be the fundamental matrix multiplied by $\vec{v}(t)$ and solve for $\vec{v}(t)$

$$\vec{x}_{\text{part}p}(t) = \Phi(t) \cdot \vec{v}(t)$$

- Take the derivative of the particular solution

$$\vec{x}_{\text{part}p}'(t) = \Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t)$$

- Substitute particular solution and its derivative into the system of ODEs

$$\Phi'(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- The fundamental matrix has columns that are solutions to the homogeneous system so its derivative is $A \cdot \Phi(t)$

$$A \cdot \Phi(t) \cdot \vec{v}(t) + \Phi(t) \cdot \vec{v}'(t) = A \cdot \Phi(t) \cdot \vec{v}(t) + \vec{f}(t)$$

- Cancel like terms

$$\Phi(t) \cdot \vec{v}'(t) = \vec{f}(t)$$

- Multiply by the inverse of the fundamental matrix

$$\vec{v}'(t) = \frac{1}{\Phi(t)} \cdot \vec{f}(t)$$

- Integrate to solve for $\vec{v}(t)$

$$\vec{v}(t) = \int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds$$

- Plug $\vec{v}(t)$ into the equation for the particular solution

$$\vec{x}_{\text{p}}(t) = \Phi(t) \cdot \left(\int_0^t \frac{1}{\Phi(s)} \cdot \vec{f}(s) ds \right)$$

- Plug in the fundamental matrix and the forcing function and compute

$$\vec{x}_{\text{p}}(t) = \begin{bmatrix} \frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t)}{2} - 2 + 2e^{-t} \cos(\sqrt{2}t) \\ 2\sqrt{2}e^{-t} \sin(\sqrt{2}t) + 1 - e^{-t} \cos(\sqrt{2}t) \end{bmatrix}$$

- Plug particular solution back into general solution

$$\vec{x}(t) = c_1 \vec{x}_{\text{1}}(t) + c_2 \vec{x}_{\text{2}}(t) + \begin{bmatrix} \frac{\sqrt{2}e^{-t} \sin(\sqrt{2}t)}{2} - 2 + 2e^{-t} \cos(\sqrt{2}t) \\ 2\sqrt{2}e^{-t} \sin(\sqrt{2}t) + 1 - e^{-t} \cos(\sqrt{2}t) \end{bmatrix}$$

- Substitute in vector of dependent variables

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{(-\sqrt{2}c_2+4)e^{-t} \cos(\sqrt{2}t)}{2} - 2 - \frac{\sqrt{2}e^{-t}(c_1-1) \sin(\sqrt{2}t)}{2} \\ e^{-t}(c_1-1) \cos(\sqrt{2}t) + 1 - e^{-t}(c_2-2\sqrt{2}) \sin(\sqrt{2}t) \end{bmatrix}$$

- Solution to the system of ODEs

$$\begin{cases} x_1(t) = \frac{(-\sqrt{2}c_2+4)e^{-t} \cos(\sqrt{2}t)}{2} - 2 - \frac{\sqrt{2}e^{-t}(c_1-1) \sin(\sqrt{2}t)}{2}, & x_2(t) = e^{-t}(c_1-1) \cos(\sqrt{2}t) + 1 - e^{-t} \end{cases}$$

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 61

```
dsolve([diff(x__1(t),t)=-1*x__1(t)-1*x__2(t)-1,diff(x__2(t),t)=2*x__1(t)-1*x__2(t)+5],singsol)
```

$$\begin{aligned} x_1(t) &= -2 + e^{-t} \left(\cos(\sqrt{2}t) c_1 + c_2 \sin(\sqrt{2}t) \right) \\ x_2(t) &= 1 - e^{-t} \sqrt{2} \left(c_2 \cos(\sqrt{2}t) - c_1 \sin(\sqrt{2}t) \right) \end{aligned}$$

✓ Solution by Mathematica

Time used: 0.414 (sec). Leaf size: 85

```
DSolve[{x1'[t]==-1*x1[t]-1*x2[t]-1,x2'[t]==2*x1[t]-1*x2[t]+5},{x1[t],x2[t]},t,IncludeSingular
```

$$\begin{aligned}x_1(t) &\rightarrow c_1 e^{-t} \cos(\sqrt{2}t) - \frac{c_2 e^{-t} \sin(\sqrt{2}t)}{\sqrt{2}} - 2 \\x_2(t) &\rightarrow e^{-t} \left(e^t + c_2 \cos(\sqrt{2}t) + \sqrt{2} c_1 \sin(\sqrt{2}t) \right)\end{aligned}$$

20 Chapter 9.2, Autonomous Systems and Stability. page 517

20.1	problem 1	4093
20.2	problem 2 part 1	4101
20.3	problem 2 part 2	4109
20.4	problem 3 part 1	4117
20.5	problem 3 part 2	4124

20.1 problem 1

20.1.1 Solution using Matrix exponential method 4093

20.1.2 Solution using explicit Eigenvalue and Eigenvector method . . . 4094

Internal problem ID [807]

Internal file name [OUTPUT/807_Sunday_June_05_2022_01_50_17_AM_1022039/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.2, Autonomous Systems and Stability. page 517

Problem number: 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x'(t) = -x(t)$$

$$y'(t) = -2y(t)$$

With initial conditions

$$[x(0) = 4, y(0) = 2]$$

20.1.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x'(t) \\ y'(t) \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & -2 \end{bmatrix} \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-t} & 0 \\ 0 & e^{-2t} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}\vec{x}_h(t) &= e^{At}\vec{x}_0 \\ &= \begin{bmatrix} e^{-t} & 0 \\ 0 & e^{-2t} \end{bmatrix} \begin{bmatrix} 4 \\ 2 \end{bmatrix} \\ &= \begin{bmatrix} 4e^{-t} \\ 2e^{-2t} \end{bmatrix}\end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

20.1.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x'(t) \\ y'(t) \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & -2 \end{bmatrix} \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -1 & 0 \\ 0 & -2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -1 - \lambda & 0 \\ 0 & -2 - \lambda \end{bmatrix}\right) = 0$$

Since the matrix A is triangular matrix, then the determinant is the product of the elements along the diagonal. Therefore the above becomes

$$(-1 - \lambda)(-2 - \lambda) = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -2$$

$$\lambda_2 = -1$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
-2	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & 0 \\ 0 & -2 \end{bmatrix} - (-2) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 1 & 0 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 0\}$

Hence the solution is

$$\begin{bmatrix} 0 \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 0 \\ t \end{bmatrix} = t \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & 0 \\ 0 & -2 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 0 & 0 & 0 \\ 0 & -1 & 0 \end{array} \right]$$

Since the current pivot $A(1,2)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 1 and row 2 gives

$$\left[\begin{array}{cc|c} 0 & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 0 & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_1\}$ and the leading variables are $\{v_2\}$. Let $v_1 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_2 = 0\}$

Hence the solution is

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = \begin{bmatrix} t \\ 0 \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = t \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-2	1	1	No	$\begin{bmatrix} 0 \\ 1 \end{bmatrix}$
-1	1	1	No	$\begin{bmatrix} 1 \\ 0 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^{-2t} \\ &= \begin{bmatrix} 0 \\ 1 \end{bmatrix} e^{-2t} \end{aligned}$$

Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{-t} \\ &= \begin{bmatrix} 1 \\ 0 \end{bmatrix} e^{-t}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = c_1 \begin{bmatrix} 0 \\ e^{-2t} \end{bmatrix} + c_2 \begin{bmatrix} e^{-t} \\ 0 \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} c_2 e^{-t} \\ c_1 e^{-2t} \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x(0) = 4 \\ y(0) = 2 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 4 \\ 2 \end{bmatrix} = \begin{bmatrix} c_2 \\ c_1 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = 2 \\ c_2 = 4 \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} 4 e^{-t} \\ 2 e^{-2t} \end{bmatrix}$$

The following is the phase plot of the system.

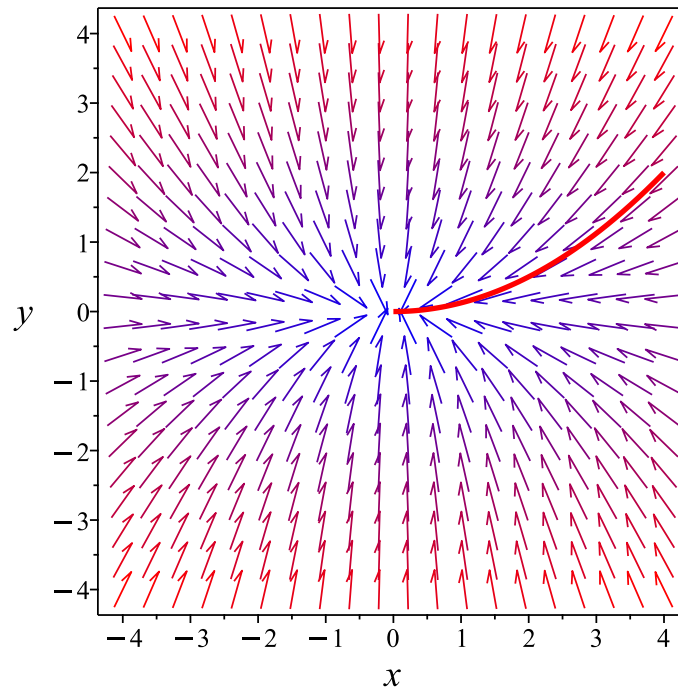
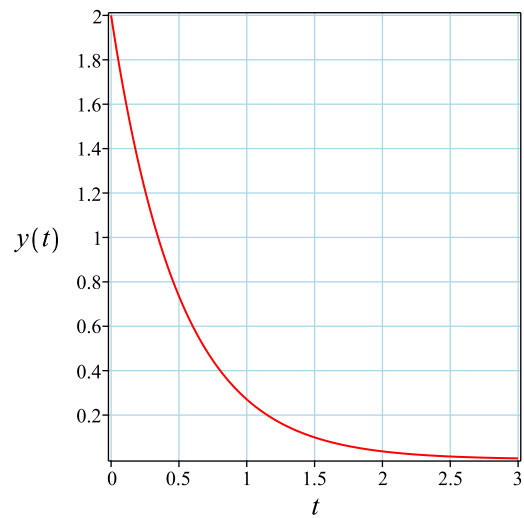
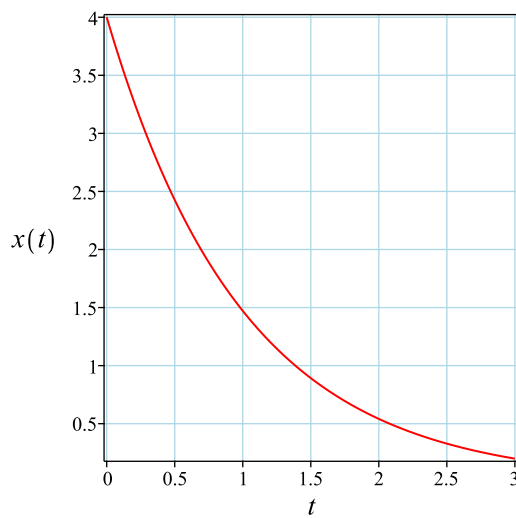


Figure 557: Phase plot

The following are plots of each solution.



✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 20

```
dsolve([diff(x(t),t) = -x(t), diff(y(t),t) = -2*y(t), x(0) = 4, y(0) = 2], singsol=all)
```

$$\begin{aligned}x(t) &= 4e^{-t} \\ y(t) &= 2e^{-2t}\end{aligned}$$

✓ Solution by Mathematica

Time used: 0.038 (sec). Leaf size: 22

```
DSolve[{x'[t]==-1*x[t]+0*y[t],y'[t]==-2*y[t]},{x[0]==4,y[0]==2},{x[t],y[t]},t,IncludeSingularSolutions->True]
```

$$\begin{aligned}x(t) &\rightarrow 4e^{-t} \\ y(t) &\rightarrow 2e^{-2t}\end{aligned}$$

20.2 problem 2 part 1

20.2.1 Solution using Matrix exponential method 4101

20.2.2 Solution using explicit Eigenvalue and Eigenvector method . . . 4102

Internal problem ID [808]

Internal file name [OUTPUT/808_Sunday_June_05_2022_01_50_18_AM_30327055/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.2, Autonomous Systems and Stability. page 517

Problem number: 2 part 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x'(t) = -x(t)$$

$$y'(t) = 2y(t)$$

With initial conditions

$$[x(0) = 4, y(0) = 2]$$

20.2.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x'(t) \\ y'(t) \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & 2 \end{bmatrix} \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-t} & 0 \\ 0 & e^{2t} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}\vec{x}_h(t) &= e^{At}\vec{x}_0 \\ &= \begin{bmatrix} e^{-t} & 0 \\ 0 & e^{2t} \end{bmatrix} \begin{bmatrix} 4 \\ 2 \end{bmatrix} \\ &= \begin{bmatrix} 4e^{-t} \\ 2e^{2t} \end{bmatrix}\end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

20.2.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x'(t) \\ y'(t) \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & 2 \end{bmatrix} \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -1 & 0 \\ 0 & 2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -1 - \lambda & 0 \\ 0 & 2 - \lambda \end{bmatrix}\right) = 0$$

Since the matrix A is triangular matrix, then the determinant is the product of the elements along the diagonal. Therefore the above becomes

$$(-1 - \lambda)(2 - \lambda) = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1$$

$$\lambda_2 = 2$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
2	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & 0 \\ 0 & 2 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 & 0 \\ 0 & 3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 0 & 0 & 0 \\ 0 & 3 & 0 \end{array} \right]$$

Since the current pivot $A(1,2)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 1 and row 2 gives

$$\left[\begin{array}{cc|c} 0 & 3 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 0 & 3 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_1\}$ and the leading variables are $\{v_2\}$. Let $v_1 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_2 = 0\}$

Hence the solution is

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = \begin{bmatrix} t \\ 0 \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = t \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & 0 \\ 0 & 2 \end{bmatrix} - (2) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} -3 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -3 & 0 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -3 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 0\}$

Hence the solution is

$$\begin{bmatrix} 0 \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 0 \\ t \end{bmatrix} = t \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-1	1	1	No	$\begin{bmatrix} 1 \\ 0 \end{bmatrix}$
2	1	1	No	$\begin{bmatrix} 0 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^{-t} \\ &= \begin{bmatrix} 1 \\ 0 \end{bmatrix} e^{-t} \end{aligned}$$

Since eigenvalue 2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{2t} \\ &= \begin{bmatrix} 0 \\ 1 \end{bmatrix} e^{2t}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = c_1 \begin{bmatrix} e^{-t} \\ 0 \end{bmatrix} + c_2 \begin{bmatrix} 0 \\ e^{2t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} c_1 e^{-t} \\ c_2 e^{2t} \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x(0) = 4 \\ y(0) = 2 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 4 \\ 2 \end{bmatrix} = \begin{bmatrix} c_1 \\ c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = 4 \\ c_2 = 2 \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} 4 e^{-t} \\ 2 e^{2t} \end{bmatrix}$$

The following is the phase plot of the system.

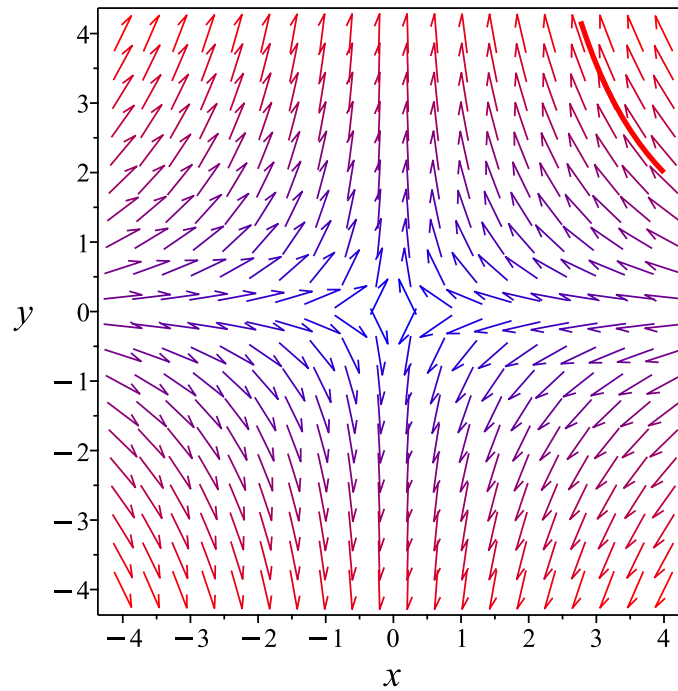
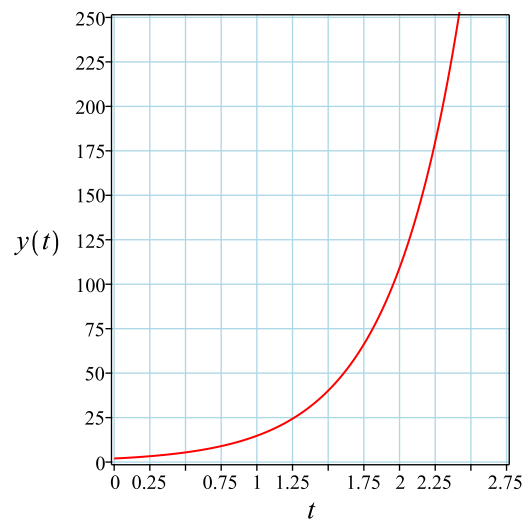
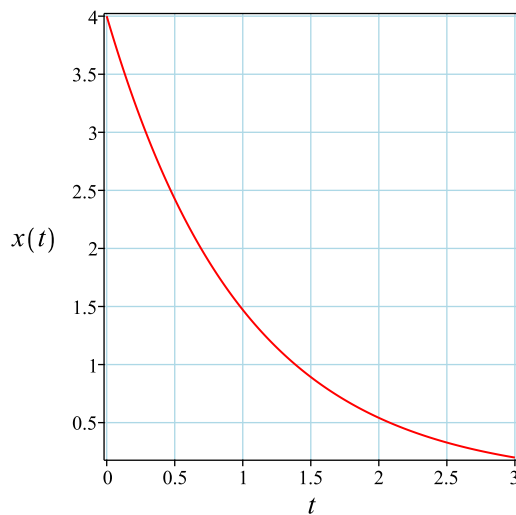


Figure 558: Phase plot

The following are plots of each solution.



✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 20

```
dsolve([diff(x(t),t) = -x(t), diff(y(t),t) = 2*y(t), x(0) = 4, y(0) = 2], singsol=all)
```

$$\begin{aligned}x(t) &= 4e^{-t} \\ y(t) &= 2e^{2t}\end{aligned}$$

✓ Solution by Mathematica

Time used: 0.038 (sec). Leaf size: 22

```
DSolve[{x'[t]==-1*x[t]+0*y[t],y'[t]==0*x[t]+2*y[t]},{x[0]==4,y[0]==2},{x[t],y[t]},t,IncludeS
```

$$\begin{aligned}x(t) &\rightarrow 4e^{-t} \\ y(t) &\rightarrow 2e^{2t}\end{aligned}$$

20.3 problem 2 part 2

20.3.1 Solution using Matrix exponential method 4109

20.3.2 Solution using explicit Eigenvalue and Eigenvector method . . . 4110

Internal problem ID [809]

Internal file name [OUTPUT/809_Sunday_June_05_2022_01_50_20_AM_21084214/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.2, Autonomous Systems and Stability. page 517

Problem number: 2 part 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : "**system of linear ODEs**"

Solve

$$x'(t) = -x(t)$$

$$y'(t) = 2y(t)$$

With initial conditions

$$[x(0) = 4, y(0) = 0]$$

20.3.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x'(t) \\ y'(t) \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & 2 \end{bmatrix} \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} e^{-t} & 0 \\ 0 & e^{2t} \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}\vec{x}_h(t) &= e^{At}\vec{x}_0 \\ &= \begin{bmatrix} e^{-t} & 0 \\ 0 & e^{2t} \end{bmatrix} \begin{bmatrix} 4 \\ 0 \end{bmatrix} \\ &= \begin{bmatrix} 4e^{-t} \\ 0 \end{bmatrix}\end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

20.3.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x'(t) \\ y'(t) \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & 2 \end{bmatrix} \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} -1 & 0 \\ 0 & 2 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -1 - \lambda & 0 \\ 0 & 2 - \lambda \end{bmatrix}\right) = 0$$

Since the matrix A is triangular matrix, then the determinant is the product of the elements along the diagonal. Therefore the above becomes

$$(-1 - \lambda)(2 - \lambda) = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = -1$$

$$\lambda_2 = 2$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
-1	1	real eigenvalue
2	1	real eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -1$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & 0 \\ 0 & 2 \end{bmatrix} - (-1) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 & 0 \\ 0 & 3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} 0 & 0 & 0 \\ 0 & 3 & 0 \end{array} \right]$$

Since the current pivot $A(1,2)$ is zero, then the current pivot row is replaced with a row with a non-zero pivot. Swapping row 1 and row 2 gives

$$\left[\begin{array}{cc|c} 0 & 3 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} 0 & 3 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_1\}$ and the leading variables are $\{v_2\}$. Let $v_1 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_2 = 0\}$

Hence the solution is

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = \begin{bmatrix} t \\ 0 \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = t \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} t \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = 2$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} -1 & 0 \\ 0 & 2 \end{bmatrix} - (2) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} -3 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -3 & 0 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -3 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = 0\}$

Hence the solution is

$$\begin{bmatrix} 0 \\ t \end{bmatrix} = \begin{bmatrix} 0 \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} 0 \\ t \end{bmatrix} = t \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
-1	1	1	No	$\begin{bmatrix} 1 \\ 0 \end{bmatrix}$
2	1	1	No	$\begin{bmatrix} 0 \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Since eigenvalue -1 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned} \vec{x}_1(t) &= \vec{v}_1 e^{-t} \\ &= \begin{bmatrix} 1 \\ 0 \end{bmatrix} e^{-t} \end{aligned}$$

Since eigenvalue 2 is real and distinct then the corresponding eigenvector solution is

$$\begin{aligned}\vec{x}_2(t) &= \vec{v}_2 e^{2t} \\ &= \begin{bmatrix} 0 \\ 1 \end{bmatrix} e^{2t}\end{aligned}$$

Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = c_1 \begin{bmatrix} e^{-t} \\ 0 \end{bmatrix} + c_2 \begin{bmatrix} 0 \\ e^{2t} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} c_1 e^{-t} \\ c_2 e^{2t} \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x(0) = 4 \\ y(0) = 0 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 4 \\ 0 \end{bmatrix} = \begin{bmatrix} c_1 \\ c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = 4 \\ c_2 = 0 \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} 4 e^{-t} \\ 0 \end{bmatrix}$$

The following is the phase plot of the system.

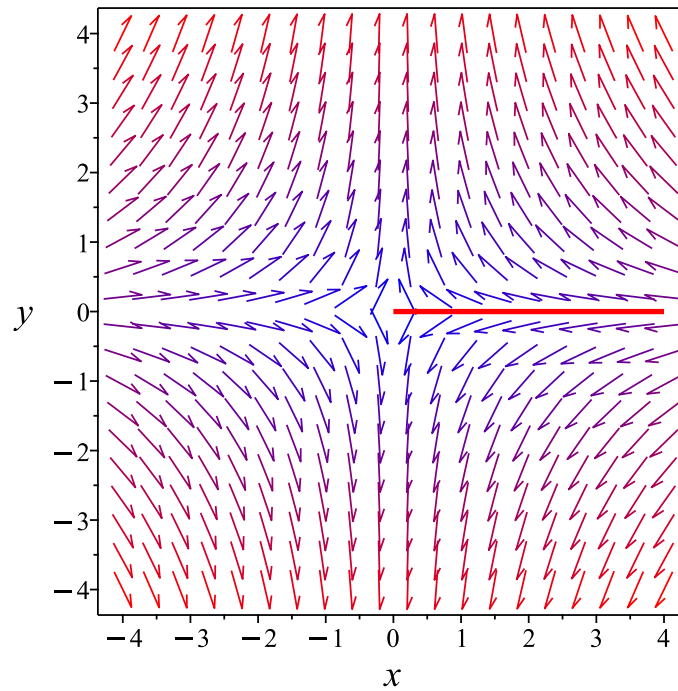
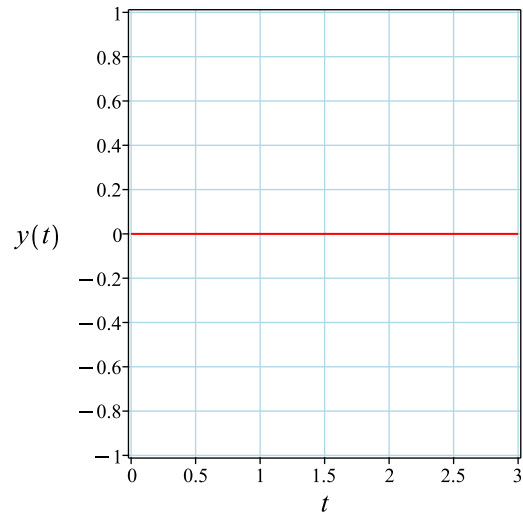
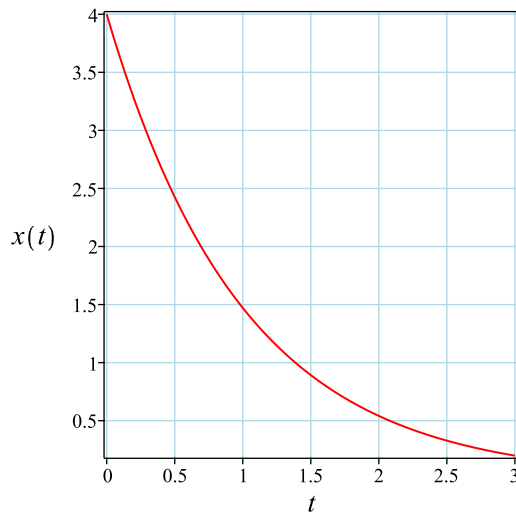


Figure 559: Phase plot

The following are plots of each solution.



✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve([diff(x(t),t) = -x(t), diff(y(t),t) = 2*y(t), x(0) = 4, y(0) = 0], singsol=all)
```

$$\begin{aligned}x(t) &= 4e^{-t} \\ y(t) &= 0\end{aligned}$$

✓ Solution by Mathematica

Time used: 0.042 (sec). Leaf size: 16

```
DSolve[{x'[t]==-1*x[t]+0*y[t],y'[t]==0*x[t]+2*y[t]},{x[0]==4,y[0]==0},{x[t],y[t]},t,IncludeS
```

$$\begin{aligned}x(t) &\rightarrow 4e^{-t} \\ y(t) &\rightarrow 0\end{aligned}$$

20.4 problem 3 part 1

20.4.1 Solution using Matrix exponential method 4117

20.4.2 Solution using explicit Eigenvalue and Eigenvector method . . . 4118

Internal problem ID [810]

Internal file name [OUTPUT/810_Sunday_June_05_2022_01_50_21_AM_58122514/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.2, Autonomous Systems and Stability. page 517

Problem number: 3 part 1.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"system of linear ODEs"**

Solve

$$x'(t) = -y(t)$$

$$y'(t) = x(t)$$

With initial conditions

$$[x(0) = 4, y(0) = 0]$$

20.4.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x'(t) \\ y'(t) \end{bmatrix} = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \cos(t) & -\sin(t) \\ \sin(t) & \cos(t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}\vec{x}_h(t) &= e^{At}\vec{x}_0 \\ &= \begin{bmatrix} \cos(t) & -\sin(t) \\ \sin(t) & \cos(t) \end{bmatrix} \begin{bmatrix} 4 \\ 0 \end{bmatrix} \\ &= \begin{bmatrix} 4\cos(t) \\ 4\sin(t) \end{bmatrix}\end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

20.4.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x'(t) \\ y'(t) \end{bmatrix} = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -\lambda & -1 \\ 1 & -\lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-i$	1	complex eigenvalue
i	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} - (-i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} i & -1 \\ 1 & i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} i & -1 & 0 \\ 1 & i & 0 \end{array} \right]$$

$$R_2 = iR_1 + R_2 \implies \left[\begin{array}{cc|c} i & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} i & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -it\}$

Hence the solution is

$$\begin{bmatrix} -It \\ t \end{bmatrix} = \begin{bmatrix} -it \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -It \\ t \end{bmatrix} = t \begin{bmatrix} -i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -It \\ t \end{bmatrix} = \begin{bmatrix} -i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} - (i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -i & -1 \\ 1 & -i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -i & -1 & 0 \\ 1 & -i & 0 \end{array} \right]$$

$$R_2 = -iR_1 + R_2 \implies \left[\begin{array}{cc|c} -i & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -i & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = it\}$

Hence the solution is

$$\begin{bmatrix} I t \\ t \end{bmatrix} = \begin{bmatrix} i t \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} I t \\ t \end{bmatrix} = t \begin{bmatrix} i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} I t \\ t \end{bmatrix} = \begin{bmatrix} i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
i	1	1	No	$\begin{bmatrix} i \\ 1 \end{bmatrix}$
$-i$	1	1	No	$\begin{bmatrix} -i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = c_1 \begin{bmatrix} i e^{it} \\ e^{it} \end{bmatrix} + c_2 \begin{bmatrix} -i e^{-it} \\ e^{-it} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} -i(c_2e^{-it} - c_1e^{it}) \\ c_1e^{it} + c_2e^{-it} \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x(0) = 4 \\ y(0) = 0 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 4 \\ 0 \end{bmatrix} = \begin{bmatrix} i(c_1 - c_2) \\ c_1 + c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = -2i \\ c_2 = 2i \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} -i(2ie^{-it} + 2ie^{it}) \\ -2ie^{it} + 2ie^{-it} \end{bmatrix}$$

The following is the phase plot of the system.

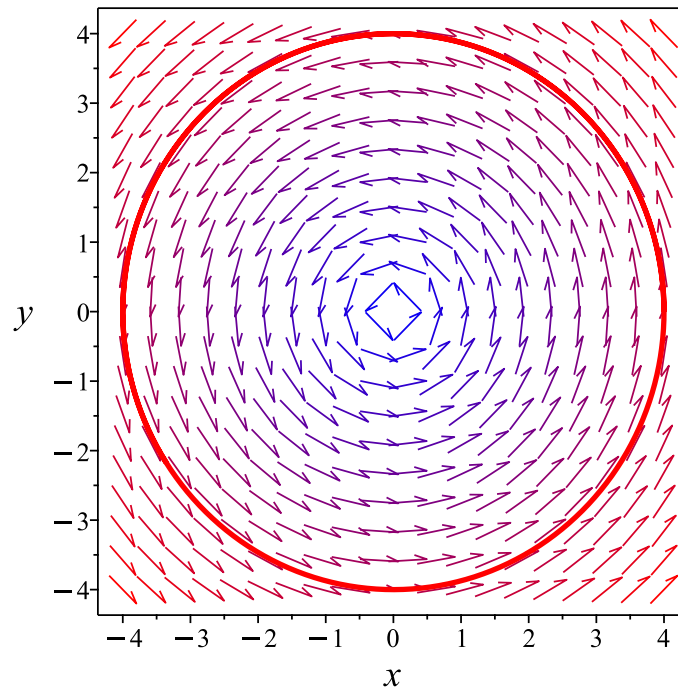


Figure 560: Phase plot

The following are plots of each solution.

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 16

```
dsolve([diff(x(t),t) = -y(t), diff(y(t),t) = x(t), x(0) = 4, y(0) = 0], singsol=all)
```

$$\begin{aligned}x(t) &= 4 \cos(t) \\ y(t) &= 4 \sin(t)\end{aligned}$$

✓ Solution by Mathematica

Time used: 0.002 (sec). Leaf size: 16

```
DSolve[{x'[t]==-0*x[t]-1*y[t],y'[t]==1*x[t]+0*y[t]},{x[0]==4,y[0]==0},{x[t],y[t]},t,IncludeS
```

$$\begin{aligned}x(t) &\rightarrow 4 \cos(t) \\ y(t) &\rightarrow 4 \sin(t)\end{aligned}$$

20.5 problem 3 part 2

20.5.1 Solution using Matrix exponential method 4124

20.5.2 Solution using explicit Eigenvalue and Eigenvector method . . . 4125

Internal problem ID [811]

Internal file name [OUTPUT/811_Sunday_June_05_2022_01_50_22_AM_62168446/index.tex]

Book: Elementary differential equations and boundary value problems, 10th ed., Boyce and DiPrima

Section: Chapter 9.2, Autonomous Systems and Stability. page 517

Problem number: 3 part 2.

ODE order: 1.

ODE degree: 1.

The type(s) of ODE detected by this program : **"system of linear ODEs"**

Solve

$$x'(t) = -y(t)$$

$$y'(t) = x(t)$$

With initial conditions

$$[x(0) = 0, y(0) = 4]$$

20.5.1 Solution using Matrix exponential method

In this method, we will assume we have found the matrix exponential e^{At} already. There are different methods to determine this but will not be shown here. This is a system of linear ODE's given as

$$\vec{x}'(t) = A \vec{x}(t)$$

Or

$$\begin{bmatrix} x'(t) \\ y'(t) \end{bmatrix} = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$$

For the above matrix A , the matrix exponential can be found to be

$$e^{At} = \begin{bmatrix} \cos(t) & -\sin(t) \\ \sin(t) & \cos(t) \end{bmatrix}$$

Therefore the homogeneous solution is

$$\begin{aligned}\vec{x}_h(t) &= e^{At}\vec{x}_0 \\ &= \begin{bmatrix} \cos(t) & -\sin(t) \\ \sin(t) & \cos(t) \end{bmatrix} \begin{bmatrix} 0 \\ 4 \end{bmatrix} \\ &= \begin{bmatrix} -4\sin(t) \\ 4\cos(t) \end{bmatrix}\end{aligned}$$

Since no forcing function is given, then the final solution is $\vec{x}_h(t)$ above.

20.5.2 Solution using explicit Eigenvalue and Eigenvector method

This is a system of linear ODE's given as

$$\vec{x}'(t) = A\vec{x}(t)$$

Or

$$\begin{bmatrix} x'(t) \\ y'(t) \end{bmatrix} = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$$

The first step is find the homogeneous solution. We start by finding the eigenvalues of A . This is done by solving the following equation for the eigenvalues λ

$$\det(A - \lambda I) = 0$$

Expanding gives

$$\det\left(\begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}\right) = 0$$

Therefore

$$\det\left(\begin{bmatrix} -\lambda & -1 \\ 1 & -\lambda \end{bmatrix}\right) = 0$$

Which gives the characteristic equation

$$\lambda^2 + 1 = 0$$

The roots of the above are the eigenvalues.

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

This table summarises the above result

eigenvalue	algebraic multiplicity	type of eigenvalue
$-i$	1	complex eigenvalue
i	1	complex eigenvalue

Now the eigenvector for each eigenvalue are found.

Considering the eigenvalue $\lambda_1 = -i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} - (-i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} i & -1 \\ 1 & i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} i & -1 & 0 \\ 1 & i & 0 \end{array} \right]$$

$$R_2 = iR_1 + R_2 \implies \left[\begin{array}{cc|c} i & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} i & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = -it\}$

Hence the solution is

$$\begin{bmatrix} -It \\ t \end{bmatrix} = \begin{bmatrix} -it \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} -It \\ t \end{bmatrix} = t \begin{bmatrix} -i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} -It \\ t \end{bmatrix} = \begin{bmatrix} -i \\ 1 \end{bmatrix}$$

Considering the eigenvalue $\lambda_2 = i$

We need to solve $A\vec{v} = \lambda\vec{v}$ or $(A - \lambda I)\vec{v} = \vec{0}$ which becomes

$$\left(\begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} - (i) \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -i & -1 \\ 1 & -i \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Now forward elimination is applied to solve for the eigenvector \vec{v} . The augmented matrix is

$$\left[\begin{array}{cc|c} -i & -1 & 0 \\ 1 & -i & 0 \end{array} \right]$$

$$R_2 = -iR_1 + R_2 \implies \left[\begin{array}{cc|c} -i & -1 & 0 \\ 0 & 0 & 0 \end{array} \right]$$

Therefore the system in Echelon form is

$$\begin{bmatrix} -i & -1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

The free variables are $\{v_2\}$ and the leading variables are $\{v_1\}$. Let $v_2 = t$. Now we start back substitution. Solving the above equation for the leading variables in terms of free variables gives equation $\{v_1 = it\}$

Hence the solution is

$$\begin{bmatrix} It \\ t \end{bmatrix} = \begin{bmatrix} it \\ t \end{bmatrix}$$

Since there is one free Variable, we have found one eigenvector associated with this eigenvalue. The above can be written as

$$\begin{bmatrix} It \\ t \end{bmatrix} = t \begin{bmatrix} i \\ 1 \end{bmatrix}$$

Let $t = 1$ the eigenvector becomes

$$\begin{bmatrix} It \\ t \end{bmatrix} = \begin{bmatrix} i \\ 1 \end{bmatrix}$$

The following table gives a summary of this result. It shows for each eigenvalue the algebraic multiplicity m , and its geometric multiplicity k and the eigenvectors associated with the eigenvalue. If $m > k$ then the eigenvalue is defective which means the number of normal linearly independent eigenvectors associated with this eigenvalue (called the geometric multiplicity k) does not equal the algebraic multiplicity m , and we need to determine an additional $m - k$ generalized eigenvectors for this eigenvalue.

eigenvalue	multiplicity		defective?	eigenvectors
	algebraic m	geometric k		
i	1	1	No	$\begin{bmatrix} i \\ 1 \end{bmatrix}$
$-i$	1	1	No	$\begin{bmatrix} -i \\ 1 \end{bmatrix}$

Now that we found the eigenvalues and associated eigenvectors, we will go over each eigenvalue and generate the solution basis. The only problem we need to take care of is if the eigenvalue is defective. Therefore the final solution is

$$\vec{x}_h(t) = c_1 \vec{x}_1(t) + c_2 \vec{x}_2(t)$$

Which is written as

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = c_1 \begin{bmatrix} ie^{it} \\ e^{it} \end{bmatrix} + c_2 \begin{bmatrix} -ie^{-it} \\ e^{-it} \end{bmatrix}$$

Which becomes

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} i(c_1 e^{it} - c_2 e^{-it}) \\ c_1 e^{it} + c_2 e^{-it} \end{bmatrix}$$

Since initial conditions are given, the solution above needs to be updated by solving for the constants of integrations using the given initial conditions

$$\begin{bmatrix} x(0) = 0 \\ y(0) = 4 \end{bmatrix} \tag{1}$$

Substituting initial conditions into the above solution at $t = 0$ gives

$$\begin{bmatrix} 0 \\ 4 \end{bmatrix} = \begin{bmatrix} i(c_1 - c_2) \\ c_1 + c_2 \end{bmatrix}$$

Solving for the constants of integrations gives

$$\begin{bmatrix} c_1 = 2 \\ c_2 = 2 \end{bmatrix}$$

Substituting these constants back in original solution in Eq. (1) gives

$$\begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} i(2e^{it} - 2e^{-it}) \\ 2e^{it} + 2e^{-it} \end{bmatrix}$$

The following is the phase plot of the system.

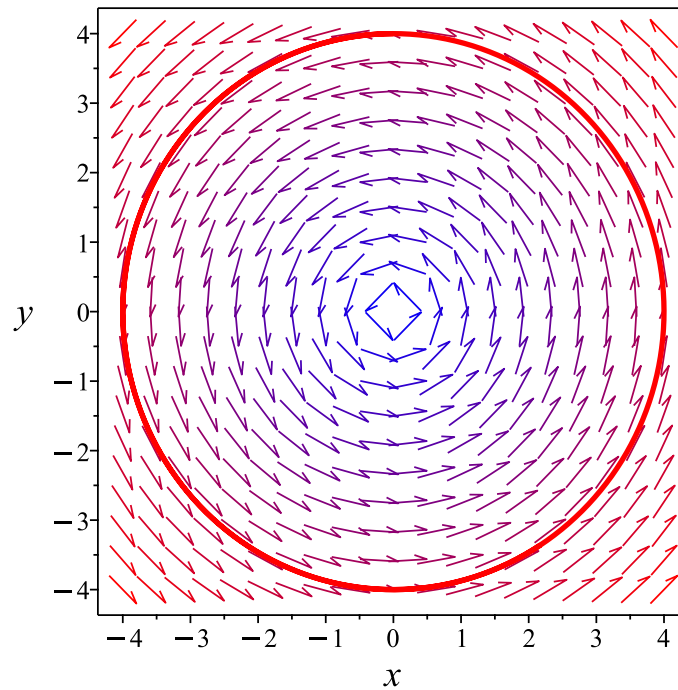


Figure 561: Phase plot

The following are plots of each solution.

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 16

```
dsolve([diff(x(t),t) = -y(t), diff(y(t),t) = x(t), x(0) = 0, y(0) = 4], singsol=all)
```

$$\begin{aligned}x(t) &= -4 \sin(t) \\ y(t) &= 4 \cos(t)\end{aligned}$$

✓ Solution by Mathematica

Time used: 0.002 (sec). Leaf size: 16

```
DSolve[{x'[t]==-0*x[t]-1*y[t],y'[t]==1*x[t]+0*y[t]},{x[0]==0,y[0]==4},{x[t],y[t]},t,IncludeS
```

$$\begin{aligned}x(t) &\rightarrow -4 \sin(t) \\ y(t) &\rightarrow 4 \cos(t)\end{aligned}$$