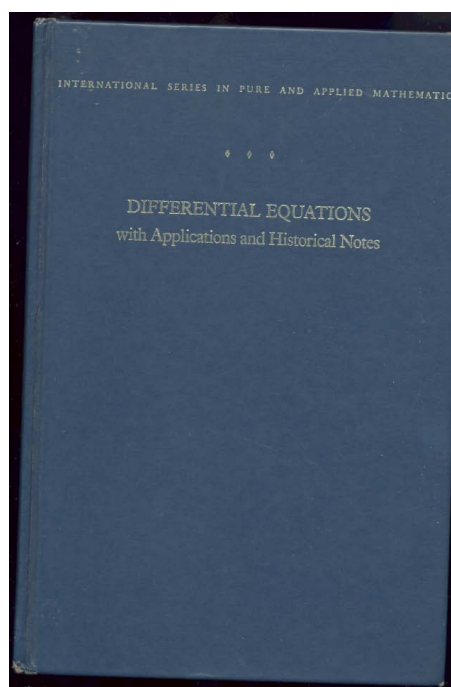


**A Solution Manual For**

**Differential equations with applications  
and historical notes, George F. Simmons,  
1971**



**Nasser M. Abbasi**

May 15, 2024

# Contents

<b>1</b>	<b>Chapter 2, section 7, page 37</b>	<b>2</b>
<b>2</b>	<b>Chapter 2, section 8, page 41</b>	<b>109</b>
<b>3</b>	<b>Chapter 2, section 10, page 47</b>	<b>223</b>
<b>4</b>	<b>Chapter 2, section 11, page 49</b>	<b>314</b>
<b>5</b>	<b>Chapter 2, End of chapter, page 61</b>	<b>392</b>

# 1 Chapter 2, section 7, page 37

1.1	problem 1.a . . . . .	3
1.2	problem 1.b . . . . .	18
1.3	problem 1.c . . . . .	36
1.4	problem 1.d . . . . .	46
1.5	problem 1. . . . .	61
1.6	problem 3.a . . . . .	72
1.7	problem 3.b . . . . .	82
1.8	problem 5.a . . . . .	90
1.9	problem 5.b . . . . .	101

## 1.1 problem 1.a

1.1.1	Solving as homogeneousTypeD2 ode . . . . .	3
1.1.2	Solving as first order ode lie symmetry lookup ode . . . . .	5
1.1.3	Solving as bernoulli ode . . . . .	9
1.1.4	Solving as exact ode . . . . .	13

Internal problem ID [3080]

Internal file name [OUTPUT/2572\_Sunday\_June\_05\_2022\_03\_20\_04\_AM\_84551471/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 7, page 37

**Problem number:** 1.a.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"bernoulli", "homogeneousTypeD2", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _Bernoulli]
```

$$-y^2 + xyy' = -x^2$$

### 1.1.1 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$-u(x)^2 x^2 + x^2 u(x) (u'(x)x + u(x)) = -x^2$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{1}{ux} \end{aligned}$$

Where  $f(x) = -\frac{1}{x}$  and  $g(u) = \frac{1}{u}$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{1}{u}} du &= -\frac{1}{x} dx \\ \int \frac{1}{\frac{1}{u}} du &= \int -\frac{1}{x} dx \\ \frac{u^2}{2} &= -\ln(x) + c_2\end{aligned}$$

The solution is

$$\frac{u(x)^2}{2} + \ln(x) - c_2 = 0$$

Replacing  $u(x)$  in the above solution by  $\frac{y}{x}$  results in the solution for  $y$  in implicit form

$$\begin{aligned}\frac{y^2}{2x^2} + \ln(x) - c_2 &= 0 \\ \frac{y^2}{2x^2} + \ln(x) - c_2 &= 0\end{aligned}$$

### Summary

The solution(s) found are the following

$$\frac{y^2}{2x^2} + \ln(x) - c_2 = 0 \tag{1}$$

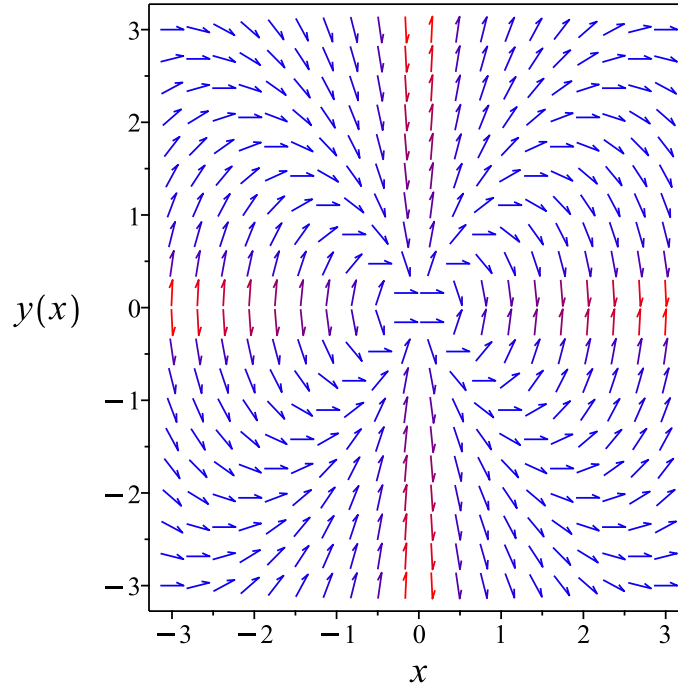


Figure 1: Slope field plot

Verification of solutions

$$\frac{y^2}{2x^2} + \ln(x) - c_2 = 0$$

Verified OK.

### 1.1.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-x^2 + y^2}{xy}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **Bernoulli**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 1: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{x^2}{y}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{x^2}{y}} dy \end{aligned}$$

Which results in

$$S = \frac{y^2}{2x^2}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{-x^2 + y^2}{xy}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -\frac{y^2}{x^3} \\ S_y &= \frac{y}{x^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{x} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{R}$$



The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -\ln(R) + c_1 \quad (4)$$

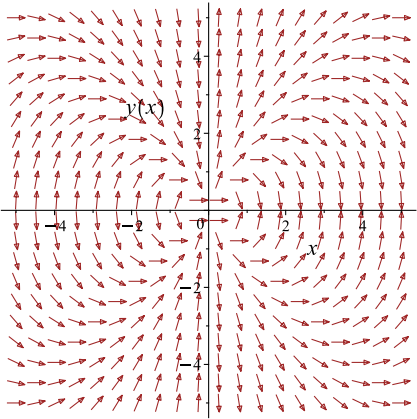
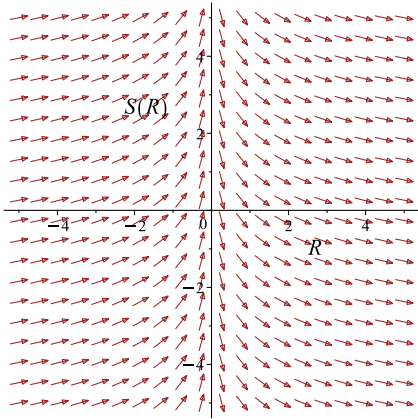
To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{y^2}{2x^2} = -\ln(x) + c_1$$

Which simplifies to

$$\frac{y^2}{2x^2} = -\ln(x) + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{-x^2 + y^2}{xy}$ 	$R = x$ $S = \frac{y^2}{2x^2}$	$\frac{dS}{dR} = -\frac{1}{R}$ 

### Summary

The solution(s) found are the following

$$\frac{y^2}{2x^2} = -\ln(x) + c_1 \quad (1)$$

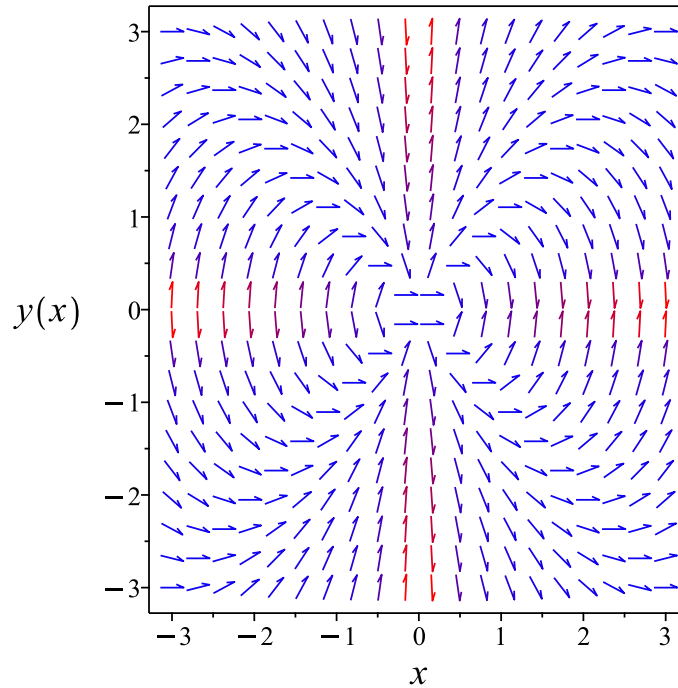


Figure 2: Slope field plot

### Verification of solutions

$$\frac{y^2}{2x^2} = -\ln(x) + c_1$$

Verified OK.

### 1.1.3 Solving as bernoulli ode

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{-x^2 + y^2}{xy} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = \frac{1}{x}y - x\frac{1}{y} \tag{1}$$

The standard Bernoulli ODE has the form

$$y' = f_0(x)y + f_1(x)y^n \tag{2}$$

The first step is to divide the above equation by  $y^n$  which gives

$$\frac{y'}{y^n} = f_0(x)y^{1-n} + f_1(x) \tag{3}$$

The next step is use the substitution  $w = y^{1-n}$  in equation (3) which generates a new ODE in  $w(x)$  which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution  $y(x)$  which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned}f_0(x) &= \frac{1}{x} \\f_1(x) &= -x \\n &= -1\end{aligned}$$

Dividing both sides of ODE (1) by  $y^n = \frac{1}{y}$  gives

$$y'y = \frac{y^2}{x} - x \quad (4)$$

Let

$$\begin{aligned}w &= y^{1-n} \\&= y^2\end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t  $x$  gives

$$w' = 2yy' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned}\frac{w'(x)}{2} &= \frac{w(x)}{x} - x \\w' &= \frac{2w}{x} - 2x\end{aligned} \quad (7)$$

The above now is a linear ODE in  $w(x)$  which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(x) + p(x)w(x) = q(x)$$

Where here

$$\begin{aligned}p(x) &= -\frac{2}{x} \\q(x) &= -2x\end{aligned}$$

Hence the ode is

$$w'(x) - \frac{2w(x)}{x} = -2x$$

The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int -\frac{2}{x} dx} \\ &= \frac{1}{x^2}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu w) &= (\mu)(-2x) \\ \frac{d}{dx}\left(\frac{w}{x^2}\right) &= \left(\frac{1}{x^2}\right)(-2x) \\ d\left(\frac{w}{x^2}\right) &= \left(-\frac{2}{x}\right) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{w}{x^2} &= \int -\frac{2}{x} dx \\ \frac{w}{x^2} &= -2 \ln(x) + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = \frac{1}{x^2}$  results in

$$w(x) = -2 \ln(x) x^2 + c_1 x^2$$

which simplifies to

$$w(x) = x^2(-2 \ln(x) + c_1)$$

Replacing  $w$  in the above by  $y^2$  using equation (5) gives the final solution.

$$y^2 = x^2(-2 \ln(x) + c_1)$$

Solving for  $y$  gives

$$\begin{aligned}y(x) &= \sqrt{-2 \ln(x) + c_1} x \\ y(x) &= -\sqrt{-2 \ln(x) + c_1} x\end{aligned}$$

### Summary

The solution(s) found are the following

$$y = \sqrt{-2 \ln(x) + c_1} x \quad (1)$$

$$y = -\sqrt{-2 \ln(x) + c_1} x \quad (2)$$

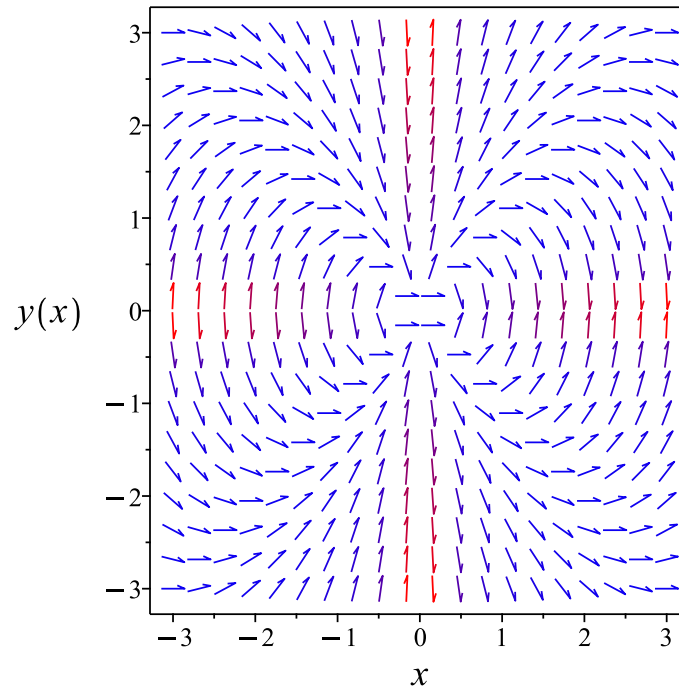


Figure 3: Slope field plot

### Verification of solutions

$$y = \sqrt{-2 \ln(x) + c_1} x$$

Verified OK.

$$y = -\sqrt{-2 \ln(x) + c_1} x$$

Verified OK.

### 1.1.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (xy) dy &= (-x^2 + y^2) dx \\ (x^2 - y^2) dx + (xy) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= x^2 - y^2 \\ N(x, y) &= xy \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(x^2 - y^2) \\ &= -2y\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(xy) \\ &= y\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{xy} ((-2y) - (y)) \\ &= -\frac{3}{x}\end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int -\frac{3}{x} dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-3 \ln(x)} \\ &= \frac{1}{x^3}\end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{x^3}(x^2 - y^2) \\ &= \frac{x^2 - y^2}{x^3}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{x^3}(xy) \\ &= \frac{y}{x^2}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left( \frac{x^2 - y^2}{x^3} \right) + \left( \frac{y}{x^2} \right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{x^2 - y^2}{x^3} dx \\ \phi &= \frac{y^2}{2x^2} + \ln(x) + f(y)\end{aligned} \tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = \frac{y}{x^2} + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{y}{x^2}$ . Therefore equation (4) becomes

$$\frac{y}{x^2} = \frac{y}{x^2} + f'(y) \tag{5}$$



Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = \frac{y^2}{2x^2} + \ln(x) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \frac{y^2}{2x^2} + \ln(x)$$

### Summary

The solution(s) found are the following

$$\frac{y^2}{2x^2} + \ln(x) = c_1 \tag{1}$$

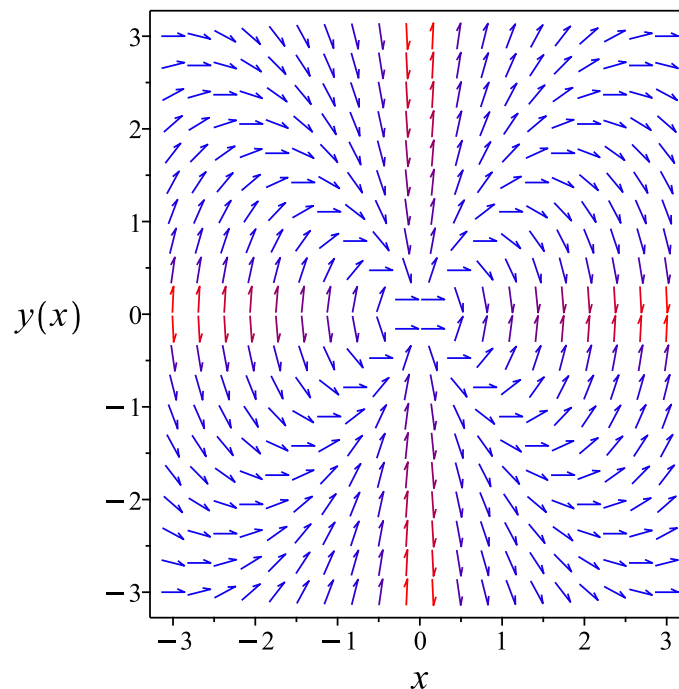


Figure 4: Slope field plot

## Verification of solutions

$$\frac{y^2}{2x^2} + \ln(x) = c_1$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 28

```
dsolve((x^2-y(x)^2)+x*y(x)*diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = \sqrt{-2 \ln(x) + c_1} x$$
$$y(x) = -\sqrt{-2 \ln(x) + c_1} x$$

### ✓ Solution by Mathematica

Time used: 0.186 (sec). Leaf size: 36

```
DSolve[(x^2-y[x]^2)+x*y[x]*y'[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -x\sqrt{-2\log(x) + c_1}$$
$$y(x) \rightarrow x\sqrt{-2\log(x) + c_1}$$

## 1.2 problem 1.b

1.2.1	Solving as homogeneousTypeD2 ode . . . . .	18
1.2.2	Solving as first order ode lie symmetry lookup ode . . . . .	20
1.2.3	Solving as bernoulli ode . . . . .	24
1.2.4	Solving as exact ode . . . . .	28
1.2.5	Solving as riccati ode . . . . .	32

Internal problem ID [3081]

Internal file name [OUTPUT/2573\_Sunday\_June\_05\_2022\_03\_20\_07\_AM\_1694912/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 7, page 37

**Problem number:** 1.b.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "riccati", "bernoulli", "homogeneousTypeD2", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _Bernoulli]
```

$$y'x^2 - 2yx - 2y^2 = 0$$

### 1.2.1 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$(u'(x)x + u(x))x^2 - 2u(x)x^2 - 2u(x)^2x^2 = 0$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{2u^2 + u}{x}\end{aligned}$$

Where  $f(x) = \frac{1}{x}$  and  $g(u) = 2u^2 + u$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{2u^2 + u} du &= \frac{1}{x} dx \\ \int \frac{1}{2u^2 + u} du &= \int \frac{1}{x} dx \\ -\ln(2u + 1) + \ln(u) &= \ln(x) + c_2\end{aligned}$$

Raising both side to exponential gives

$$e^{-\ln(2u+1)+\ln(u)} = e^{\ln(x)+c_2}$$

Which simplifies to

$$\frac{u}{2u + 1} = c_3 x$$

Therefore the solution  $y$  is

$$\begin{aligned}y &= xu \\ &= -\frac{x^2 c_3}{2c_3 x - 1}\end{aligned}$$

### Summary

The solution(s) found are the following

$$y = -\frac{x^2 c_3}{2c_3 x - 1} \tag{1}$$

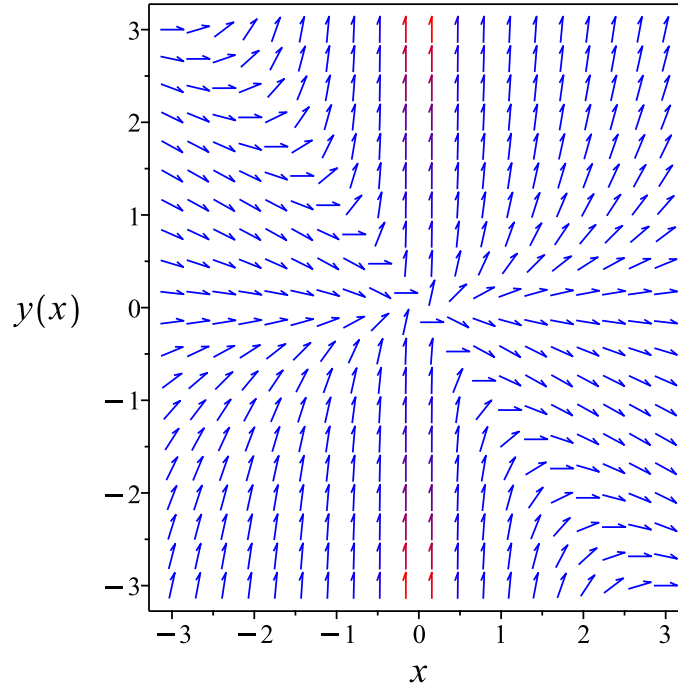


Figure 5: Slope field plot

Verification of solutions

$$y = -\frac{x^2 c_3}{2c_3 x - 1}$$

Verified OK.

### 1.2.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{2y(y+x)}{x^2}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **Bernoulli**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 3: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{y^2}{x^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{y^2}{x^2}} dy \end{aligned}$$

Which results in

$$S = -\frac{x^2}{y}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{2y(y + x)}{x^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -\frac{2x}{y} \\ S_y &= \frac{x^2}{y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 2 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 2$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = 2R + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$-\frac{x^2}{y} = 2x + c_1$$

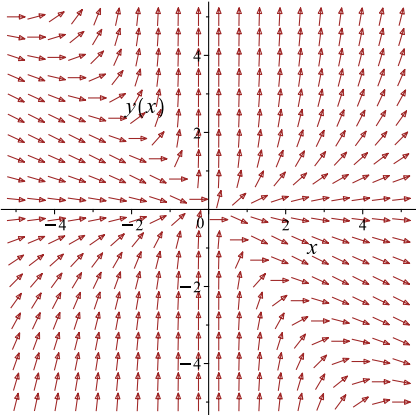
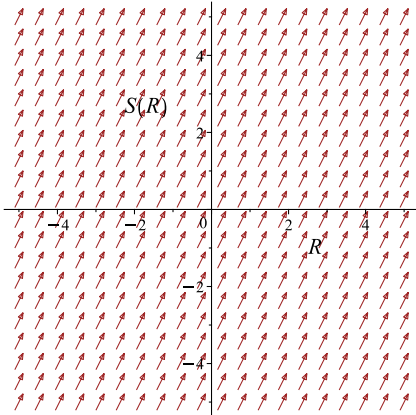
Which simplifies to

$$-\frac{x^2}{y} = 2x + c_1$$

Which gives

$$y = -\frac{x^2}{2x + c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{2y(y+x)}{x^2}$ 	$R = x$ $S = -\frac{x^2}{y}$	$\frac{dS}{dR} = 2$ 



### Summary

The solution(s) found are the following

$$y = -\frac{x^2}{2x + c_1} \quad (1)$$

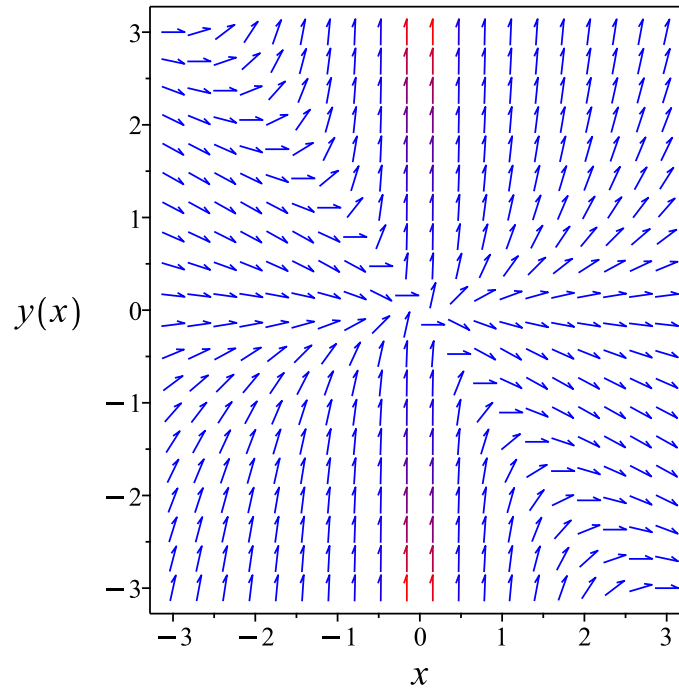


Figure 6: Slope field plot

### Verification of solutions

$$y = -\frac{x^2}{2x + c_1}$$

Verified OK.

### **1.2.3 Solving as bernoulli ode**

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{2y(y+x)}{x^2} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = \frac{2}{x}y + \frac{2}{x^2}y^2 \quad (1)$$

The standard Bernoulli ODE has the form

$$y' = f_0(x)y + f_1(x)y^n \quad (2)$$

The first step is to divide the above equation by  $y^n$  which gives

$$\frac{y'}{y^n} = f_0(x)y^{1-n} + f_1(x) \quad (3)$$

The next step is use the substitution  $w = y^{1-n}$  in equation (3) which generates a new ODE in  $w(x)$  which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution  $y(x)$  which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(x) &= \frac{2}{x} \\ f_1(x) &= \frac{2}{x^2} \\ n &= 2 \end{aligned}$$

Dividing both sides of ODE (1) by  $y^n = y^2$  gives

$$y' \frac{1}{y^2} = \frac{2}{xy} + \frac{2}{x^2} \quad (4)$$

Let

$$\begin{aligned} w &= y^{1-n} \\ &= \frac{1}{y} \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t  $x$  gives

$$w' = -\frac{1}{y^2}y' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} -w'(x) &= \frac{2w(x)}{x} + \frac{2}{x^2} \\ w' &= -\frac{2w}{x} - \frac{2}{x^2} \end{aligned} \quad (7)$$

The above now is a linear ODE in  $w(x)$  which is now solved.

Entering Linear first order ODE solver. In canonical form a linear first order is

$$w'(x) + p(x)w(x) = q(x)$$

Where here

$$p(x) = \frac{2}{x}$$
$$q(x) = -\frac{2}{x^2}$$

Hence the ode is

$$w'(x) + \frac{2w(x)}{x} = -\frac{2}{x^2}$$

The integrating factor  $\mu$  is

$$\mu = e^{\int \frac{2}{x} dx}$$
$$= x^2$$

The ode becomes

$$\frac{d}{dx}(\mu w) = (\mu) \left( -\frac{2}{x^2} \right)$$
$$\frac{d}{dx}(x^2 w) = (x^2) \left( -\frac{2}{x^2} \right)$$
$$d(x^2 w) = -2 dx$$

Integrating gives

$$x^2 w = \int -2 dx$$
$$x^2 w = -2x + c_1$$

Dividing both sides by the integrating factor  $\mu = x^2$  results in

$$w(x) = -\frac{2}{x} + \frac{c_1}{x^2}$$

which simplifies to

$$w(x) = \frac{-2x + c_1}{x^2}$$

Replacing  $w$  in the above by  $\frac{1}{y}$  using equation (5) gives the final solution.

$$\frac{1}{y} = \frac{-2x + c_1}{x^2}$$

Or

$$y = \frac{x^2}{-2x + c_1}$$

Summary

The solution(s) found are the following

$$y = \frac{x^2}{-2x + c_1} \tag{1}$$

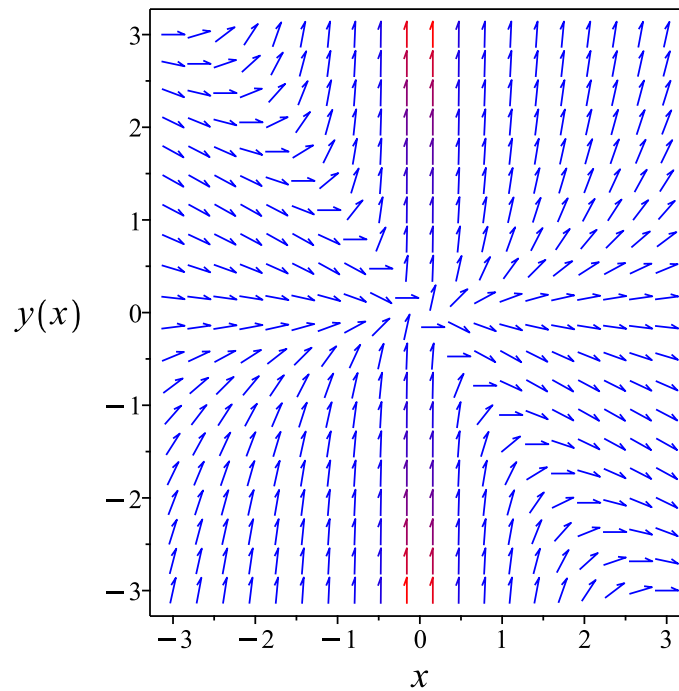


Figure 7: Slope field plot

Verification of solutions

$$y = \frac{x^2}{-2x + c_1}$$

Verified OK.

### 1.2.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (x^2) dy &= (2xy + 2y^2) dx \\ (-2xy - 2y^2) dx + (x^2) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -2xy - 2y^2 \\ N(x, y) &= x^2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-2xy - 2y^2) \\ &= -2x - 4y\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x^2) \\ &= 2x\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{x^2} ((-2x - 4y) - (2x)) \\ &= \frac{-4x - 4y}{x^2}\end{aligned}$$

Since  $A$  depends on  $y$ , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned}B &= \frac{1}{M} \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= -\frac{1}{2y(y+x)} ((2x) - (-2x - 4y)) \\ &= -\frac{2}{y}\end{aligned}$$

Since  $B$  does not depend on  $x$ , it can be used to obtain an integrating factor. Let the integrating factor be  $\mu$ . Then

$$\begin{aligned}\mu &= e^{\int B \, dy} \\ &= e^{\int -\frac{2}{y} \, dy}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-2\ln(y)} \\ &= \frac{1}{y^2}\end{aligned}$$

$M$  and  $N$  are now multiplied by this integrating factor, giving new  $\bar{M}$  and new  $\bar{N}$  which are called  $\bar{M}$  and  $\bar{N}$  so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{y^2}(-2xy - 2y^2) \\ &= \frac{-2y - 2x}{y}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{y^2}(x^2) \\ &= \frac{x^2}{y^2}\end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left( \frac{-2y - 2x}{y} \right) + \left( \frac{x^2}{y^2} \right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{-2y - 2x}{y} dx \\ \phi &= -\frac{x(x + 2y)}{y} + f(y)\end{aligned} \tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\begin{aligned}\frac{\partial\phi}{\partial y} &= -\frac{2x}{y} + \frac{x(x+2y)}{y^2} + f'(y) \\ &= \frac{x^2}{y^2} + f'(y)\end{aligned}\tag{4}$$

But equation (2) says that  $\frac{\partial\phi}{\partial y} = \frac{x^2}{y^2}$ . Therefore equation (4) becomes

$$\frac{x^2}{y^2} = \frac{x^2}{y^2} + f'(y)\tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\frac{x(x+2y)}{y} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\frac{x(x+2y)}{y}$$

The solution becomes

$$y = -\frac{x^2}{2x + c_1}$$

### Summary

The solution(s) found are the following

$$y = -\frac{x^2}{2x + c_1}\tag{1}$$



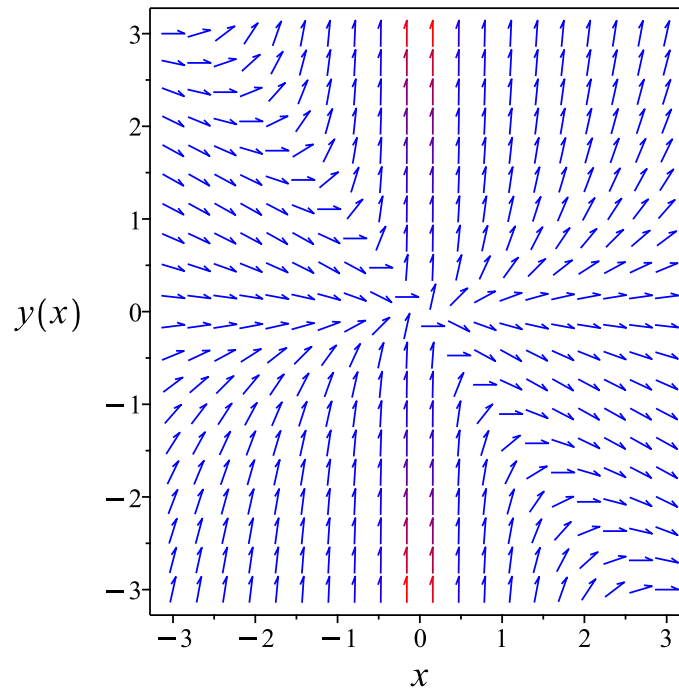


Figure 8: Slope field plot

Verification of solutions

$$y = -\frac{x^2}{2x + c_1}$$

Verified OK.

### 1.2.5 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{2y(y + x)}{x^2} \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = \frac{2y}{x} + \frac{2y^2}{x^2}$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that  $f_0(x) = 0$ ,  $f_1(x) = \frac{2}{x}$  and  $f_2(x) = \frac{2}{x^2}$ . Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{\frac{2u}{x^2}} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for  $u(x)$  which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \tag{2}$$

But

$$\begin{aligned} f_2' &= -\frac{4}{x^3} \\ f_1 f_2 &= \frac{4}{x^3} \\ f_2^2 f_0 &= 0 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$\frac{2u''(x)}{x^2} = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 x + c_2$$

The above shows that

$$u'(x) = c_1$$

Using the above in (1) gives the solution

$$y = -\frac{c_1 x^2}{2(c_1 x + c_2)}$$

Dividing both numerator and denominator by  $c_1$  gives, after renaming the constant  $\frac{c_2}{c_1} = c_3$  the following solution

$$y = -\frac{c_3 x^2}{2c_3 x + 2}$$

### Summary

The solution(s) found are the following

$$y = -\frac{c_3 x^2}{2c_3 x + 2} \quad (1)$$

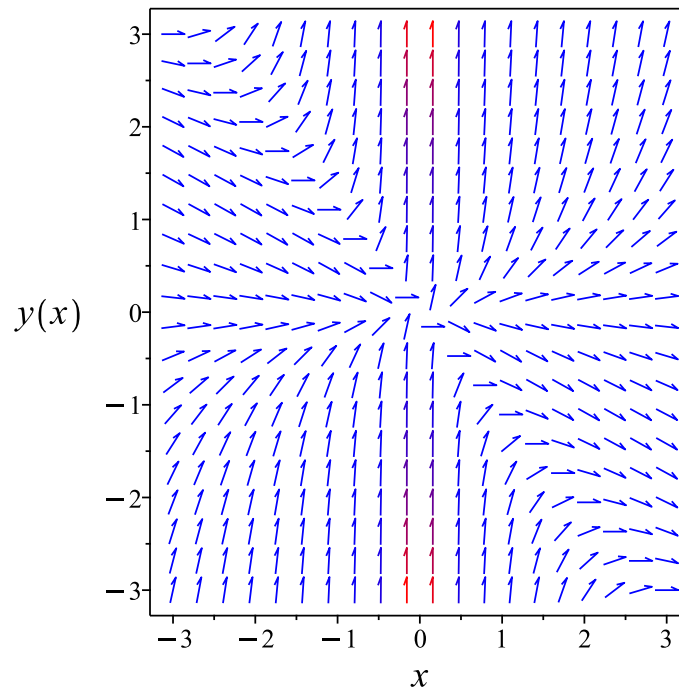


Figure 9: Slope field plot

### Verification of solutions

$$y = -\frac{c_3 x^2}{2c_3 x + 2}$$

Verified OK.

### Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
<- Bernoulli successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve(x^2*diff(y(x),x)-2*x*y(x)-2*y(x)^2=0,y(x), singsol=all)
```

$$y(x) = \frac{x^2}{-2x + c_1}$$

✓ Solution by Mathematica

Time used: 0.143 (sec). Leaf size: 22

```
DSolve[x^2*y'[x]-2*x*y[x]-2*y[x]^2==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{x^2}{-2x + c_1}$$
$$y(x) \rightarrow 0$$

### 1.3 problem 1.c

- 1.3.1 Solving as homogeneousTypeD2 ode . . . . . 36
- 1.3.2 Solving as first order ode lie symmetry calculated ode . . . . . 38

Internal problem ID [3082]

Internal file name [OUTPUT/2574\_Sunday\_June\_05\_2022\_03\_20\_10\_AM\_25744284/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 7, page 37

**Problem number:** 1.c.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD2**", "**first\_order\_ode\_lie\_symmetry\_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _dAlembert]
```

$$y'x^2 - 3(y^2 + x^2) \arctan\left(\frac{y}{x}\right) - yx = 0$$

#### 1.3.1 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$(u'(x)x + u(x))x^2 - 3(u(x)^2x^2 + x^2) \arctan(u(x)) - u(x)x^2 = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{3(u^2 + 1) \arctan(u)}{x} \end{aligned}$$

Where  $f(x) = \frac{3}{x}$  and  $g(u) = \arctan(u)(u^2 + 1)$ . Integrating both sides gives

$$\begin{aligned} \frac{1}{\arctan(u)(u^2 + 1)} du &= \frac{3}{x} dx \\ \int \frac{1}{\arctan(u)(u^2 + 1)} du &= \int \frac{3}{x} dx \\ \ln(\arctan(u)) &= 3 \ln(x) + c_2 \end{aligned}$$

Raising both side to exponential gives

$$\arctan(u) = e^{3\ln(x)+c_2}$$

Which simplifies to

$$\arctan(u) = c_3 x^3$$

Therefore the solution  $y$  is

$$\begin{aligned} y &= xu \\ &= x \tan(c_3 e^{c_2} x^3) \end{aligned}$$

### Summary

The solution(s) found are the following

$$y = x \tan(c_3 e^{c_2} x^3) \tag{1}$$

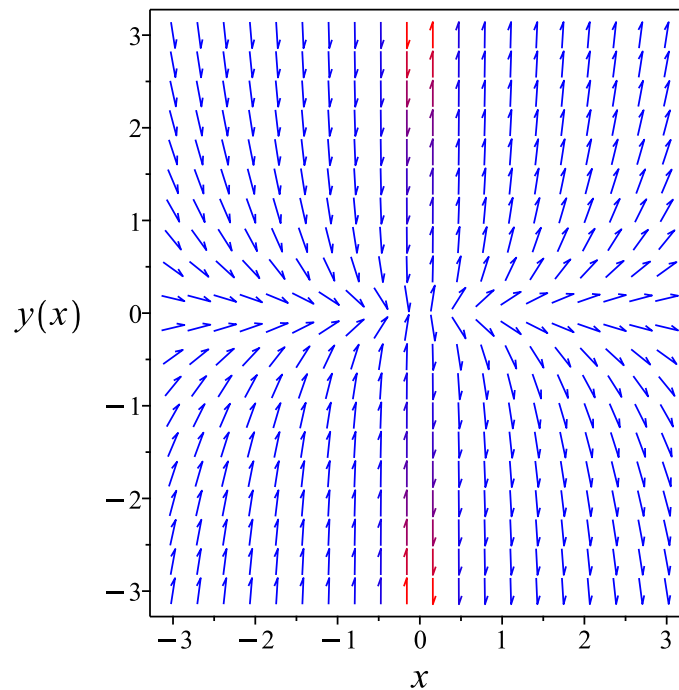


Figure 10: Slope field plot

### Verification of solutions

$$y = x \tan(c_3 e^{c_2} x^3)$$

Verified OK.

### 1.3.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{3 \arctan\left(\frac{y}{x}\right) x^2 + 3 \arctan\left(\frac{y}{x}\right) y^2 + xy}{x^2}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} & b_2 + \frac{(3 \arctan\left(\frac{y}{x}\right) x^2 + 3 \arctan\left(\frac{y}{x}\right) y^2 + xy) (b_3 - a_2)}{x^2} \\ & - \frac{(3 \arctan\left(\frac{y}{x}\right) x^2 + 3 \arctan\left(\frac{y}{x}\right) y^2 + xy)^2 a_3}{x^4} \\ & - \left( \frac{-\frac{3y}{\frac{y^2}{x^2}+1} + 6x \arctan\left(\frac{y}{x}\right) - \frac{3y^3}{x^2\left(\frac{y^2}{x^2}+1\right)} + y}{x^2} \right. \\ & \left. - \frac{2(3 \arctan\left(\frac{y}{x}\right) x^2 + 3 \arctan\left(\frac{y}{x}\right) y^2 + xy)}{x^3} \right) (xa_2 + ya_3 + a_1) \\ & - \frac{\left( \frac{3x}{\frac{y^2}{x^2}+1} + \frac{3y^2}{x\left(\frac{y^2}{x^2}+1\right)} + 6 \arctan\left(\frac{y}{x}\right) y + x \right) (xb_2 + yb_3 + b_1)}{x^2} = 0 \end{aligned} \quad (\text{5E})$$

Putting the above in normal form gives

$$\frac{9 \arctan\left(\frac{y}{x}\right)^2 x^4 a_3 + 18 \arctan\left(\frac{y}{x}\right)^2 x^2 y^2 a_3 + 9 \arctan\left(\frac{y}{x}\right)^2 y^4 a_3 + 3 \arctan\left(\frac{y}{x}\right) x^4 a_2 - 3 \arctan\left(\frac{y}{x}\right) x^4 b_3}{x^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned}
& -9 \arctan\left(\frac{y}{x}\right)^2 x^4 a_3 - 18 \arctan\left(\frac{y}{x}\right)^2 x^2 y^2 a_3 - 9 \arctan\left(\frac{y}{x}\right)^2 y^4 a_3 \\
& - 3 \arctan\left(\frac{y}{x}\right) x^4 a_2 + 3 \arctan\left(\frac{y}{x}\right) x^4 b_3 - 6 \arctan\left(\frac{y}{x}\right) x^3 y a_3 \\
& - 6 \arctan\left(\frac{y}{x}\right) x^3 y b_2 + 3 \arctan\left(\frac{y}{x}\right) x^2 y^2 a_2 - 3 \arctan\left(\frac{y}{x}\right) x^2 y^2 b_3 \\
& - 6 \arctan\left(\frac{y}{x}\right) x^2 y b_1 + 6 \arctan\left(\frac{y}{x}\right) x y^2 a_1 - 3 b_2 x^4 \\
& + 3 x^3 y a_2 - 3 x^3 y b_3 + 3 x^2 y^2 a_3 - 4 x^3 b_1 + 4 x^2 y a_1 = 0
\end{aligned} \tag{6E}$$

Simplifying the above gives

$$\begin{aligned}
& -27 \arctan\left(\frac{y}{x}\right)^2 x^4 y^2 a_3 - 27 \arctan\left(\frac{y}{x}\right)^2 x^2 y^4 a_3 \\
& - 6 \arctan\left(\frac{y}{x}\right) x^5 y a_3 - 6 \arctan\left(\frac{y}{x}\right) x^5 y b_2 - 6 \arctan\left(\frac{y}{x}\right) x^3 y^3 a_3 \\
& - 6 \arctan\left(\frac{y}{x}\right) x^3 y^3 b_2 + 3 \arctan\left(\frac{y}{x}\right) x^2 y^4 a_2 - 3 \arctan\left(\frac{y}{x}\right) x^2 y^4 b_3 \\
& - 6 \arctan\left(\frac{y}{x}\right) x^4 y b_1 + 6 \arctan\left(\frac{y}{x}\right) x^3 y^2 a_1 - 6 \arctan\left(\frac{y}{x}\right) x^2 y^3 b_1 \\
& + 6 \arctan\left(\frac{y}{x}\right) x y^4 a_1 - 9 \arctan\left(\frac{y}{x}\right)^2 x^6 a_3 - 9 \arctan\left(\frac{y}{x}\right)^2 y^6 a_3 \\
& - 3 \arctan\left(\frac{y}{x}\right) x^6 a_2 + 3 \arctan\left(\frac{y}{x}\right) x^6 b_3 - 3 x^6 b_2 - 4 x^5 b_1 \\
& + 3 x^5 y a_2 - 3 x^5 y b_3 + 3 x^4 y^2 a_3 - 3 x^4 y^2 b_2 + 3 x^3 y^3 a_2 \\
& - 3 x^3 y^3 b_3 + 3 x^2 y^4 a_3 + 4 x^4 y a_1 - 4 x^3 y^2 b_1 + 4 x^2 y^3 a_1 = 0
\end{aligned} \tag{6E}$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\left\{ x, y, \arctan\left(\frac{y}{x}\right) \right\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\left\{ x = v_1, y = v_2, \arctan\left(\frac{y}{x}\right) = v_3 \right\}$$



The above PDE (6E) now becomes

$$\begin{aligned}
& -9v_3^2v_1^6a_3 - 27v_3^2v_1^4v_2^2a_3 - 27v_3^2v_1^2v_2^4a_3 - 9v_3^2v_2^6a_3 - 3v_3v_1^6a_2 + 3v_3v_1^2v_2^4a_2 \\
& - 6v_3v_1^5v_2a_3 - 6v_3v_1^3v_2^3a_3 - 6v_3v_1^5v_2b_2 - 6v_3v_1^3v_2^3b_2 + 3v_3v_1^6b_3 \\
& - 3v_3v_1^2v_2^4b_3 + 6v_3v_1^3v_2^2a_1 + 6v_3v_1v_2^4a_1 + 3v_1^5v_2a_2 + 3v_1^3v_2^3a_2 \\
& + 3v_1^4v_2^2a_3 + 3v_1^2v_2^4a_3 - 6v_3v_1^4v_2b_1 - 6v_3v_1^2v_2^3b_1 - 3v_1^6b_2 - 3v_1^4v_2^2b_2 \\
& - 3v_1^5v_2b_3 - 3v_1^3v_2^3b_3 + 4v_1^4v_2a_1 + 4v_1^2v_2^3a_1 - 4v_1^5b_1 - 4v_1^3v_2^2b_1 = 0
\end{aligned} \tag{7E}$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2, v_3\}$$

Equation (7E) now becomes

$$\begin{aligned}
& -9v_3^2v_1^6a_3 + (-3a_2 + 3b_3)v_1^6v_3 - 3v_1^6b_2 + (-6a_3 - 6b_2)v_1^5v_2v_3 \\
& + (3a_2 - 3b_3)v_1^5v_2 - 4v_1^5b_1 - 27v_3^2v_1^4v_2^2a_3 + (3a_3 - 3b_2)v_1^4v_2^2 \\
& - 6v_3v_1^4v_2b_1 + 4v_1^4v_2a_1 + (-6a_3 - 6b_2)v_1^3v_2^3v_3 + (3a_2 - 3b_3)v_1^3v_2^3 \\
& + 6v_3v_1^3v_2^2a_1 - 4v_1^3v_2^2b_1 - 27v_3^2v_1^2v_2^4a_3 + (3a_2 - 3b_3)v_1^2v_2^4v_3 \\
& + 3v_1^2v_2^4a_3 - 6v_3v_1^2v_2^3b_1 + 4v_1^2v_2^3a_1 + 6v_3v_1v_2^4a_1 - 9v_3^2v_2^6a_3 = 0
\end{aligned} \tag{8E}$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}
4a_1 &= 0 \\
6a_1 &= 0 \\
-27a_3 &= 0 \\
-9a_3 &= 0 \\
3a_3 &= 0 \\
-6b_1 &= 0 \\
-4b_1 &= 0 \\
-3b_2 &= 0 \\
-3a_2 + 3b_3 &= 0 \\
3a_2 - 3b_3 &= 0 \\
-6a_3 - 6b_2 &= 0 \\
3a_3 - 3b_2 &= 0
\end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= 0 \\
 a_2 &= b_3 \\
 a_3 &= 0 \\
 b_1 &= 0 \\
 b_2 &= 0 \\
 b_3 &= b_3
 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}
 \xi &= x \\
 \eta &= y
 \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}
 \eta &= \eta - \omega(x, y) \xi \\
 &= y - \left( \frac{3 \arctan\left(\frac{y}{x}\right) x^2 + 3 \arctan\left(\frac{y}{x}\right) y^2 + xy}{x^2} \right) (x) \\
 &= \frac{-3 \arctan\left(\frac{y}{x}\right) x^2 - 3 \arctan\left(\frac{y}{x}\right) y^2}{x} \\
 \xi &= 0
 \end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}
 S &= \int \frac{1}{\eta} dy \\
 &= \int \frac{1}{\frac{-3 \arctan\left(\frac{y}{x}\right) x^2 - 3 \arctan\left(\frac{y}{x}\right) y^2}{x}} dy
 \end{aligned}$$

Which results in

$$S = -\frac{\ln\left(\arctan\left(\frac{y}{x}\right)\right)}{3}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{3 \arctan\left(\frac{y}{x}\right) x^2 + 3 \arctan\left(\frac{y}{x}\right) y^2 + xy}{x^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{y}{3(x^2 + y^2) \arctan\left(\frac{y}{x}\right)} \\ S_y &= -\frac{x}{3(x^2 + y^2) \arctan\left(\frac{y}{x}\right)} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{x} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -\ln(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$-\frac{\ln\left(\arctan\left(\frac{y}{x}\right)\right)}{3} = -\ln(x) + c_1$$

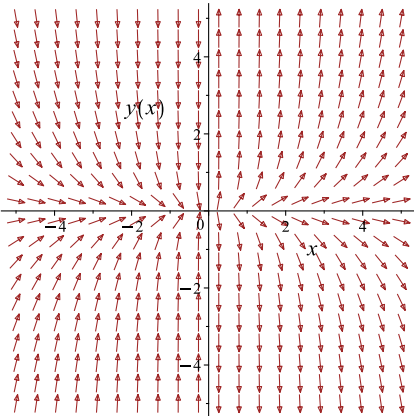
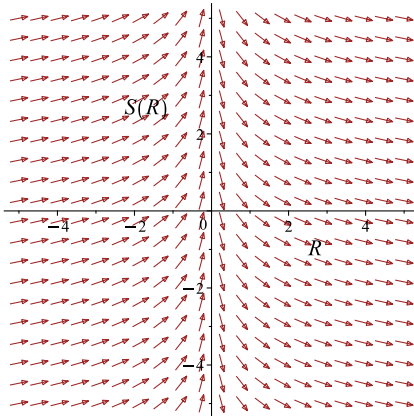
Which simplifies to

$$-\frac{\ln\left(\arctan\left(\frac{y}{x}\right)\right)}{3} = -\ln(x) + c_1$$

Which gives

$$y = \tan\left(e^{-3c_1}x^3\right)x$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{3 \arctan\left(\frac{y}{x}\right)x^2 + 3 \arctan\left(\frac{y}{x}\right)y^2 + xy}{x^2}$ 	$R = x$ $S = -\frac{\ln\left(\arctan\left(\frac{y}{x}\right)\right)}{3}$	$\frac{dS}{dR} = -\frac{1}{R}$ 

### Summary

The solution(s) found are the following

$$y = \tan\left(e^{-3c_1}x^3\right)x \tag{1}$$

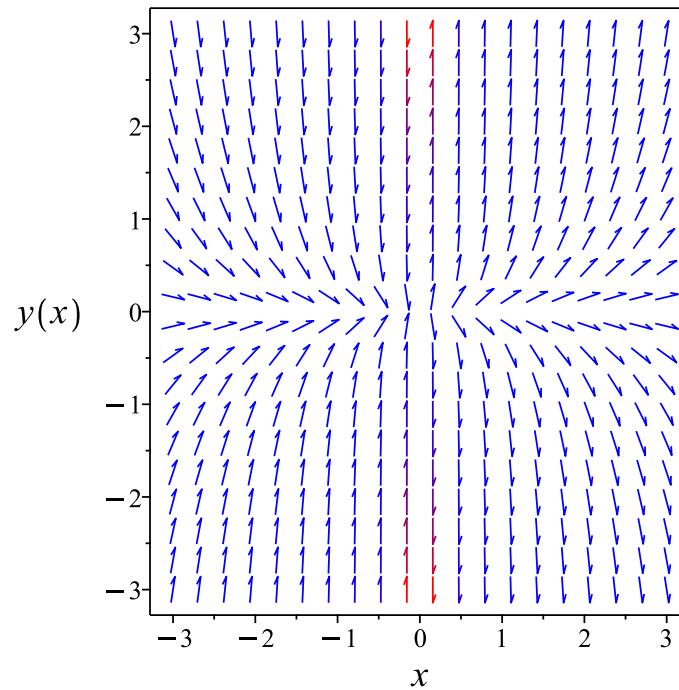


Figure 11: Slope field plot

Verification of solutions

$$y = \tan(e^{-3c_1} x^3) x$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 12

```
dsolve(x^2*diff(y(x),x)=3*(x^2+y(x)^2)*arctan(y(x)/x)+x*y(x),y(x), singsol=all)
```

$$y(x) = \tan(c_1 x^3) x$$

✓ Solution by Mathematica

Time used: 0.179 (sec). Leaf size: 37

```
DSolve[x^2*y'[x]==3*(x^2+y[x]^2)*Arctan[y[x]/x]+x*y[x],y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve} \left[ \int_1^{\frac{y(x)}{x}} \frac{1}{\text{Arctan}(K[1]) (K[1]^2 + 1)} dK[1] = 3 \log(x) + c_1, y(x) \right]$$

## 1.4 problem 1.d

1.4.1	Solving as homogeneousTypeD ode . . . . .	46
1.4.2	Solving as homogeneousTypeD2 ode . . . . .	48
1.4.3	Solving as first order ode lie symmetry lookup ode . . . . .	50
1.4.4	Solving as exact ode . . . . .	55

Internal problem ID [3083]

Internal file name [OUTPUT/2575\_Sunday\_June\_05\_2022\_03\_20\_14\_AM\_70604598/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 7, page 37

**Problem number:** 1.d.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"homogeneousTypeD", "homogeneousTypeD2", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _dAlembert]
```

$$\sin\left(\frac{y}{x}\right)y'x - \sin\left(\frac{y}{x}\right)y = x$$

### 1.4.1 Solving as homogeneousTypeD ode

Writing the ode as

$$y' = \frac{y}{x} + \frac{1}{\sin\left(\frac{y}{x}\right)} \quad (\text{A})$$

The given ode has the form

$$y' = \frac{y}{x} + g(x) f\left(b\frac{y}{x}\right)^{\frac{n}{m}} \quad (1)$$

Where  $b$  is scalar and  $g(x)$  is function of  $x$  and  $n, m$  are integers. The solution is given in Kamke page 20. Using the substitution  $y(x) = u(x)x$  then

$$\frac{dy}{dx} = \frac{du}{dx}x + u$$

Hence the given ode becomes

$$\begin{aligned}\frac{du}{dx}x + u &= u + g(x) f(bu)^{\frac{n}{m}} \\ u' &= \frac{1}{x}g(x) f(bu)^{\frac{n}{m}}\end{aligned}\quad (2)$$

The above ode is always separable. This is easily solved for  $u$  assuming the integration can be resolved, and then the solution to the original ode becomes  $y = ux$ . Comparing the given ode (A) with the form (1) shows that

$$\begin{aligned}g(x) &= 1 \\ b &= 1 \\ f\left(\frac{bx}{y}\right) &= \sin\left(\frac{y}{x}\right)\end{aligned}$$

Substituting the above in (2) results in the  $u(x)$  ode as

$$u'(x) = \frac{1}{x \sin(u(x))}$$

Which is now solved as separable In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{1}{x \sin(u)}\end{aligned}$$

Where  $f(x) = \frac{1}{x}$  and  $g(u) = \frac{1}{\sin(u)}$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{1}{\sin(u)}} du &= \frac{1}{x} dx \\ \int \frac{1}{\frac{1}{\sin(u)}} du &= \int \frac{1}{x} dx \\ -\cos(u) &= \ln(x) + c_1\end{aligned}$$

The solution is

$$-\cos(u(x)) - \ln(x) - c_1 = 0$$

Therefore the solution is found using  $y = ux$ . Hence

$$-\cos\left(\frac{y}{x}\right) - \ln(x) - c_1 = 0$$



### Summary

The solution(s) found are the following

$$-\cos\left(\frac{y}{x}\right) - \ln(x) - c_1 = 0 \quad (1)$$

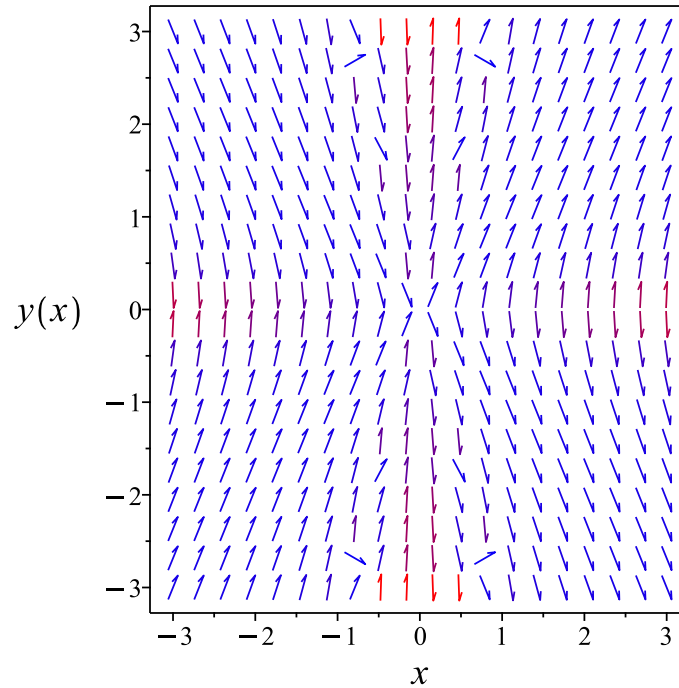


Figure 12: Slope field plot

### Verification of solutions

$$-\cos\left(\frac{y}{x}\right) - \ln(x) - c_1 = 0$$

Verified OK.

#### 1.4.2 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$\sin(u(x)) (u'(x)x + u(x))x - \sin(u(x))u(x)x = x$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{1}{\sin(u)x} \end{aligned}$$

Where  $f(x) = \frac{1}{x}$  and  $g(u) = \frac{1}{\sin(u)}$ . Integrating both sides gives

$$\frac{1}{\frac{1}{\sin(u)}} du = \frac{1}{x} dx$$

$$\int \frac{1}{\frac{1}{\sin(u)}} du = \int \frac{1}{x} dx$$

$$-\cos(u) = \ln(x) + c_2$$

The solution is

$$-\cos(u(x)) - \ln(x) - c_2 = 0$$

Replacing  $u(x)$  in the above solution by  $\frac{y}{x}$  results in the solution for  $y$  in implicit form

$$-\cos\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0$$

$$-\cos\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0$$

### Summary

The solution(s) found are the following

$$-\cos\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0 \tag{1}$$

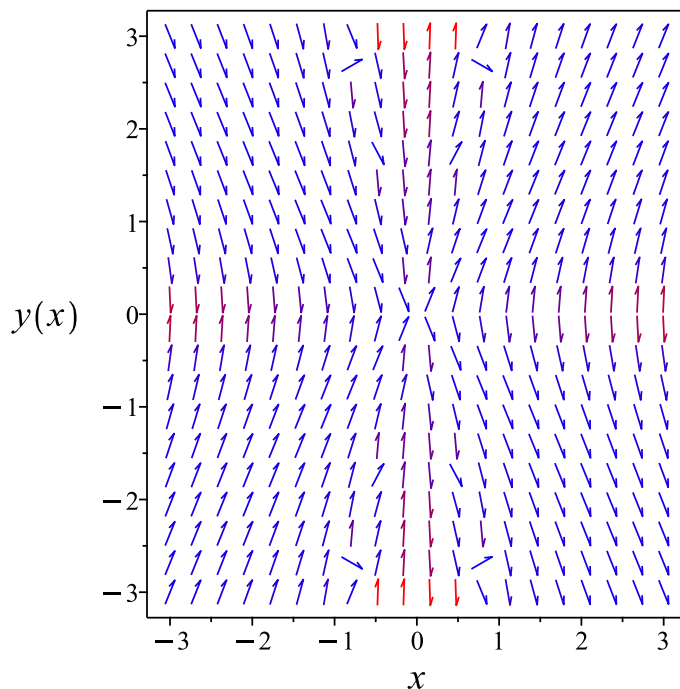


Figure 13: Slope field plot

Verification of solutions

$$-\cos\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0$$

Verified OK.

### 1.4.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{\sin\left(\frac{y}{x}\right) y + x}{\sin\left(\frac{y}{x}\right) x}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **homogeneous Type D**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 5: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= x^2 \\ \eta(x, y) &= xy\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Therefore

$$\begin{aligned}\frac{dy}{dx} &= \frac{\eta}{\xi} \\ &= \frac{xy}{x^2} \\ &= \frac{y}{x}\end{aligned}$$

This is easily solved to give

$$y = c_1 x$$

Where now the coordinate  $R$  is taken as the constant of integration. Hence

$$R = \frac{y}{x}$$

And  $S$  is found from

$$\begin{aligned}dS &= \frac{dx}{\xi} \\ &= \frac{dx}{x^2}\end{aligned}$$

Integrating gives

$$\begin{aligned}S &= \int \frac{dx}{T} \\ &= -\frac{1}{x}\end{aligned}$$

Where the constant of integration is set to zero as we just need one solution. Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{\sin\left(\frac{y}{x}\right)y + x}{\sin\left(\frac{y}{x}\right)x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= -\frac{y}{x^2} \\ R_y &= \frac{1}{x} \\ S_x &= \frac{1}{x^2} \\ S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{\sin\left(\frac{y}{x}\right)}{x} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\sin(R) S(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 e^{\cos(R)} \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$-\frac{1}{x} = c_1 e^{\cos\left(\frac{y}{x}\right)}$$

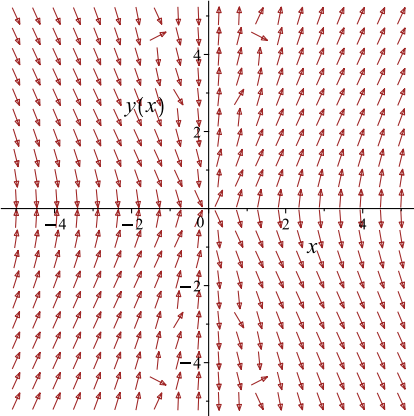
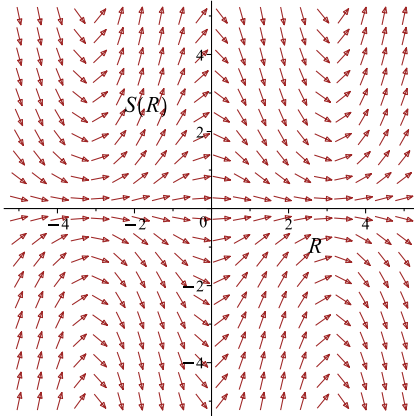
Which simplifies to

$$-\frac{1}{x} = c_1 e^{\cos\left(\frac{y}{x}\right)}$$

Which gives

$$y = \arccos\left(\ln\left(-\frac{1}{c_1 x}\right)\right) x$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{\sin\left(\frac{y}{x}\right)y+x}{\sin\left(\frac{y}{x}\right)x}$ 	$R = \frac{y}{x}$ $S = -\frac{1}{x}$	$\frac{dS}{dR} = -\sin(R) S(R)$ 

### Summary

The solution(s) found are the following

$$y = \arccos \left( \ln \left( -\frac{1}{c_1 x} \right) \right) x \quad (1)$$

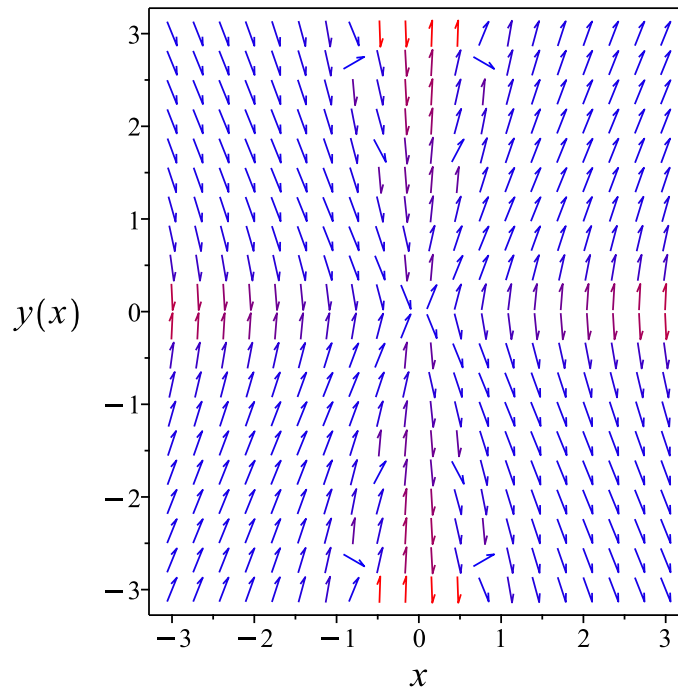


Figure 14: Slope field plot

Verification of solutions

$$y = \arccos \left( \ln \left( -\frac{1}{c_1 x} \right) \right) x$$

Verified OK.

#### 1.4.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$



Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} & \left( \sin \left( \frac{y}{x} \right) x \right) dy = \left( \sin \left( \frac{y}{x} \right) y + x \right) dx \\ \left( -\sin \left( \frac{y}{x} \right) y - x \right) dx + \left( \sin \left( \frac{y}{x} \right) x \right) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -\sin \left( \frac{y}{x} \right) y - x \\ N(x, y) &= \sin \left( \frac{y}{x} \right) x \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left( -\sin \left( \frac{y}{x} \right) y - x \right) \\ &= -\frac{\cos \left( \frac{y}{x} \right) y}{x} - \sin \left( \frac{y}{x} \right)\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left( \sin \left( \frac{y}{x} \right) x \right) \\ &= -\frac{\cos \left( \frac{y}{x} \right) y}{x} + \sin \left( \frac{y}{x} \right)\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{\csc \left( \frac{y}{x} \right)}{x} \left( \left( -\frac{\cos \left( \frac{y}{x} \right) y}{x} - \sin \left( \frac{y}{x} \right) \right) - \left( -\frac{\cos \left( \frac{y}{x} \right) y}{x} + \sin \left( \frac{y}{x} \right) \right) \right) \\ &= -\frac{2}{x}\end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int A \, dx} \\ &= e^{\int -\frac{2}{x} \, dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-2 \ln(x)} \\ &= \frac{1}{x^2}\end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{x^2} \left( -\sin \left( \frac{y}{x} \right) y - x \right) \\ &= \frac{-\sin \left( \frac{y}{x} \right) y - x}{x^2}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{x^2} \left( \sin \left( \frac{y}{x} \right) x \right) \\ &= \frac{\sin \left( \frac{y}{x} \right)}{x}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \overline{M} + \overline{N} \frac{dy}{dx} &= 0 \\ \left( \frac{-\sin\left(\frac{y}{x}\right) y - x}{x^2} \right) + \left( \frac{\sin\left(\frac{y}{x}\right)}{x} \right) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{-\sin\left(\frac{y}{x}\right) y - x}{x^2} dx \\ \phi &= \ln\left(\frac{1}{x}\right) - \cos\left(\frac{y}{x}\right) + f(y) \end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = \frac{\sin\left(\frac{y}{x}\right)}{x} + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{\sin\left(\frac{y}{x}\right)}{x}$ . Therefore equation (4) becomes

$$\frac{\sin\left(\frac{y}{x}\right)}{x} = \frac{\sin\left(\frac{y}{x}\right)}{x} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = \ln\left(\frac{1}{x}\right) - \cos\left(\frac{y}{x}\right) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \ln\left(\frac{1}{x}\right) - \cos\left(\frac{y}{x}\right)$$

### Summary

The solution(s) found are the following

$$\ln\left(\frac{1}{x}\right) - \cos\left(\frac{y}{x}\right) = c_1 \tag{1}$$

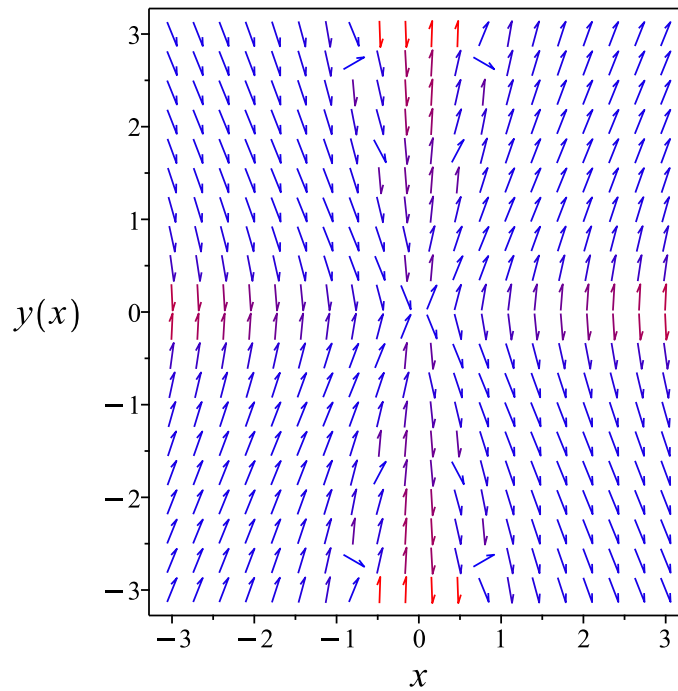


Figure 15: Slope field plot

### Verification of solutions

$$\ln\left(\frac{1}{x}\right) - \cos\left(\frac{y}{x}\right) = c_1$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

### ✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 16

```
dsolve(x*sin(y(x)/x)*diff(y(x),x)=y(x)*sin(y(x)/x)+x,y(x), singsol=all)
```

$$y(x) = \frac{(\pi + 2 \arcsin(\ln(x) + c_1)) x}{2}$$

### ✓ Solution by Mathematica

Time used: 0.435 (sec). Leaf size: 34

```
DSolve[x*Sin[y[x]/x]*y'[x]==y[x]*Sin[y[x]/x]+x,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -x \arccos(-\log(x) - c_1)$$

$$y(x) \rightarrow x \arccos(-\log(x) - c_1)$$

## 1.5 problem 1.

1.5.1	Solving as homogeneousTypeD ode . . . . .	61
1.5.2	Solving as homogeneousTypeD2 ode . . . . .	63
1.5.3	Solving as first order ode lie symmetry lookup ode . . . . .	65

Internal problem ID [3084]

Internal file name [OUTPUT/2576\_Sunday\_June\_05\_2022\_03\_20\_17\_AM\_58252680/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 7, page 37

**Problem number:** 1..

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD**", "**homogeneousTypeD2**", "**first\_order\_ode\_lie\_symmetry\_lookup**"

Maple gives the following as the ode type

```
[[_homogeneous , `class D`]]
```

$$xy' - y - 2e^{-\frac{y}{x}} = 0$$

### 1.5.1 Solving as homogeneousTypeD ode

Writing the ode as

$$y' = \frac{y}{x} + \frac{2e^{-\frac{y}{x}}}{x} \quad (\text{A})$$

The given ode has the form

$$y' = \frac{y}{x} + g(x) f\left(b\frac{y}{x}\right)^{\frac{n}{m}} \quad (1)$$

Where  $b$  is scalar and  $g(x)$  is function of  $x$  and  $n, m$  are integers. The solution is given in Kamke page 20. Using the substitution  $y(x) = u(x)x$  then

$$\frac{dy}{dx} = \frac{du}{dx}x + u$$

Hence the given ode becomes

$$\begin{aligned} \frac{du}{dx}x + u &= u + g(x) f(bu)^{\frac{n}{m}} \\ u' &= \frac{1}{x}g(x) f(bu)^{\frac{n}{m}} \end{aligned} \quad (2)$$

The above ode is always separable. This is easily solved for  $u$  assuming the integration can be resolved, and then the solution to the original ode becomes  $y = ux$ . Comparing the given ode (A) with the form (1) shows that

$$\begin{aligned}g(x) &= \frac{2}{x} \\b &= 1 \\f\left(\frac{bx}{y}\right) &= e^{\frac{y}{x}}\end{aligned}$$

Substituting the above in (2) results in the  $u(x)$  ode as

$$u'(x) = \frac{2e^{-u(x)}}{x^2}$$

Which is now solved as separable In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\&= f(x)g(u) \\&= \frac{2e^{-u}}{x^2}\end{aligned}$$

Where  $f(x) = \frac{2}{x^2}$  and  $g(u) = e^{-u}$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{e^{-u}} du &= \frac{2}{x^2} dx \\ \int \frac{1}{e^{-u}} du &= \int \frac{2}{x^2} dx \\ e^u &= -\frac{2}{x} + c_1\end{aligned}$$

The solution is

$$e^{u(x)} + \frac{2}{x} - c_1 = 0$$

Therefore the solution is found using  $y = ux$ . Hence

$$e^{\frac{y}{x}} + \frac{2}{x} - c_1 = 0$$

### Summary

The solution(s) found are the following

$$e^{\frac{y}{x}} + \frac{2}{x} - c_1 = 0 \tag{1}$$

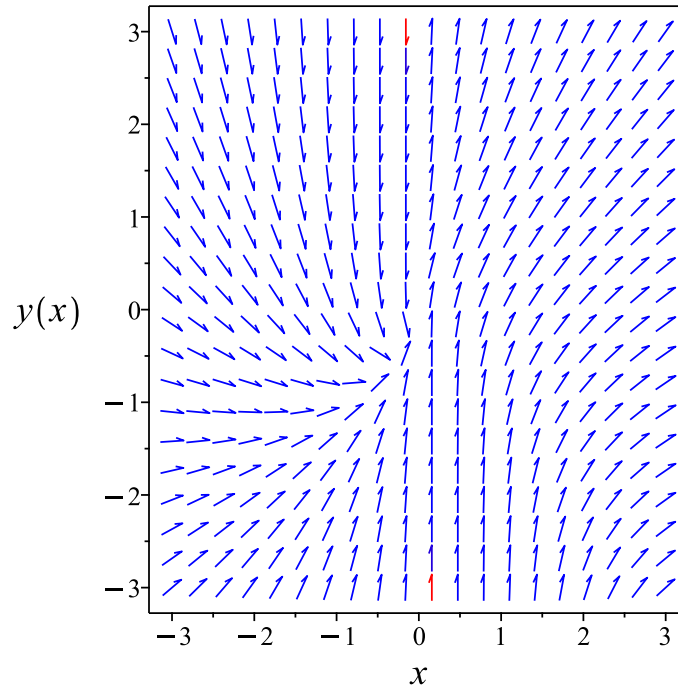


Figure 16: Slope field plot

Verification of solutions

$$e^{\frac{y}{x}} + \frac{2}{x} - c_1 = 0$$

Verified OK.

### 1.5.2 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$x(u'(x)x + u(x)) - u(x)x - 2e^{-u(x)} = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{2e^{-u}}{x^2} \end{aligned}$$

Where  $f(x) = \frac{2}{x^2}$  and  $g(u) = e^{-u}$ . Integrating both sides gives

$$\frac{1}{e^{-u}} du = \frac{2}{x^2} dx$$



$$\int \frac{1}{e^{-u}} du = \int \frac{2}{x^2} dx$$

$$e^u = -\frac{2}{x} + c_2$$

The solution is

$$e^{u(x)} + \frac{2}{x} - c_2 = 0$$

Replacing  $u(x)$  in the above solution by  $\frac{y}{x}$  results in the solution for  $y$  in implicit form

$$e^{\frac{y}{x}} + \frac{2}{x} - c_2 = 0$$

$$\frac{x e^{\frac{y}{x}} - c_2 x + 2}{x} = 0$$

### Summary

The solution(s) found are the following

$$\frac{x e^{\frac{y}{x}} - c_2 x + 2}{x} = 0 \tag{1}$$

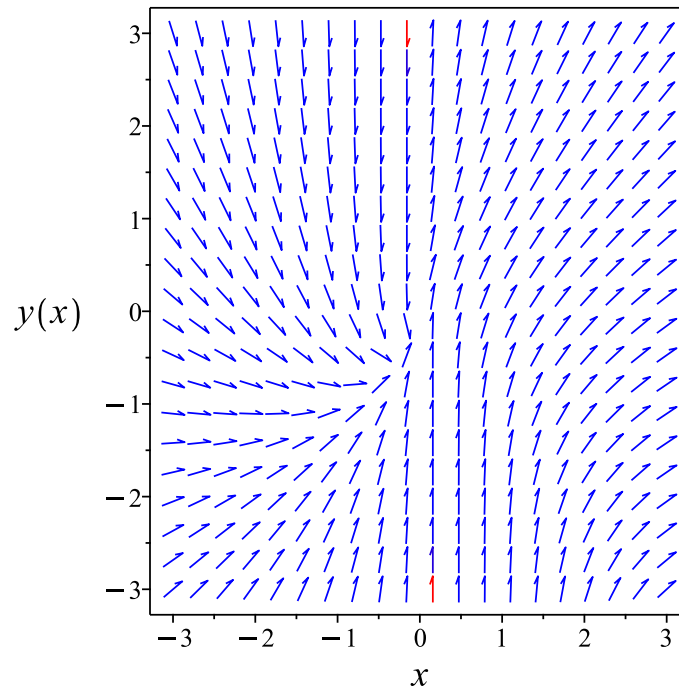


Figure 17: Slope field plot

Verification of solutions

$$\frac{x e^{\frac{y}{x}} - c_2 x + 2}{x} = 0$$

Verified OK.

### 1.5.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{y + 2 e^{-\frac{y}{x}}}{x}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **homogeneous Type D**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 7: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= x^2 \\ \eta(x, y) &= xy\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Therefore

$$\begin{aligned}\frac{dy}{dx} &= \frac{\eta}{\xi} \\ &= \frac{xy}{x^2} \\ &= \frac{y}{x}\end{aligned}$$

This is easily solved to give

$$y = c_1 x$$

Where now the coordinate  $R$  is taken as the constant of integration. Hence

$$R = \frac{y}{x}$$

And  $S$  is found from

$$\begin{aligned}dS &= \frac{dx}{\xi} \\ &= \frac{dx}{x^2}\end{aligned}$$

Integrating gives

$$\begin{aligned}S &= \int \frac{dx}{T} \\ &= -\frac{1}{x}\end{aligned}$$

Where the constant of integration is set to zero as we just need one solution. Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y + 2e^{-\frac{y}{x}}}{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= -\frac{y}{x^2} \\ R_y &= \frac{1}{x} \\ S_x &= \frac{1}{x^2} \\ S_y &= 0\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{e^{\frac{y}{x}}}{2} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{e^R}{2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = \frac{e^R}{2} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$-\frac{1}{x} = \frac{e^{\frac{y}{x}}}{2} + c_1$$

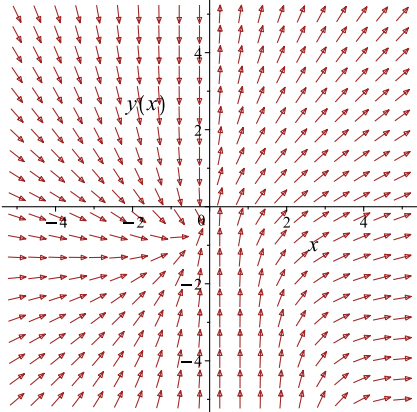
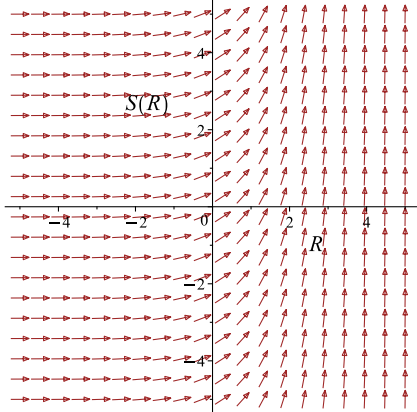
Which simplifies to

$$-\frac{1}{x} = \frac{e^{\frac{y}{x}}}{2} + c_1$$

Which gives

$$y = x \ln \left( -\frac{2(c_1 x + 1)}{x} \right)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{y+2e^{-\frac{y}{x}}}{x}$ 	$R = \frac{y}{x}$ $S = -\frac{1}{x}$	$\frac{dS}{dR} = \frac{e^R}{2}$ 

### Summary

The solution(s) found are the following

$$y = x \ln \left( -\frac{2(c_1 x + 1)}{x} \right) \quad (1)$$

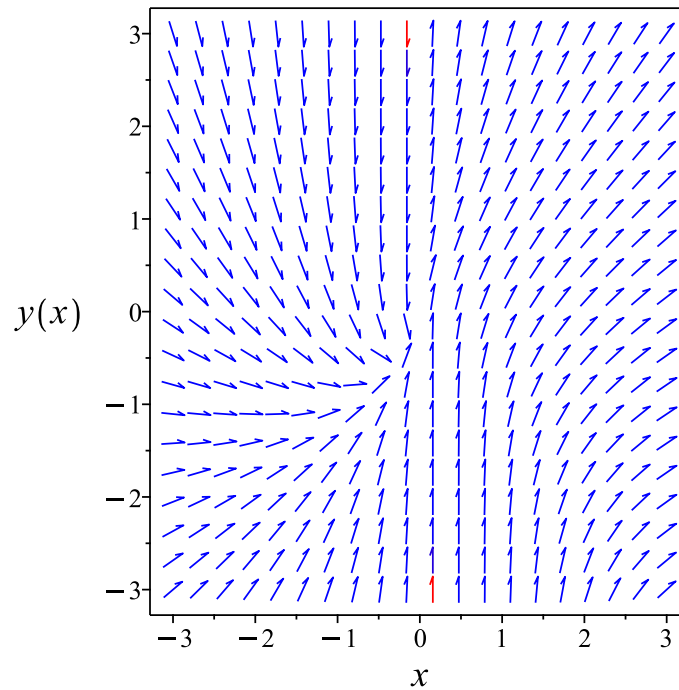


Figure 18: Slope field plot

Verification of solutions

$$y = x \ln \left( -\frac{2(c_1 x + 1)}{x} \right)$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 19

```
dsolve(x*diff(y(x),x)=y(x)+2*exp(- y(x)/x),y(x), singsol=all)
```

$$y(x) = \left( \ln(2) + \ln\left(\frac{c_1 x - 1}{x}\right) \right) x$$

✓ Solution by Mathematica

Time used: 0.618 (sec). Leaf size: 16

```
DSolve[x*y'[x]==y[x]+2*Exp[- y[x]/x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow x \log\left(-\frac{2}{x} + c_1\right)$$



## 1.6 problem 3.a

1.6.1	Solving as homogeneousTypeC ode . . . . .	72
1.6.2	Solving as first order ode lie symmetry lookup ode . . . . .	74
1.6.3	Solving as riccati ode . . . . .	78

Internal problem ID [3085]

Internal file name [OUTPUT/2577\_Sunday\_June\_05\_2022\_03\_20\_19\_AM\_9450860/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 7, page 37

**Problem number:** 3.a.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "riccati", "homogeneousTypeC", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

```
[[_homogeneous, `class C`], _Riccati]
```

$$y' - (y + x)^2 = 0$$

### 1.6.1 Solving as homogeneousTypeC ode

Let

$$z = y + x \tag{1}$$

Then

$$z'(x) = y' + 1$$

Therefore

$$y' = z'(x) - 1$$

Hence the given ode can now be written as

$$z'(x) - 1 = z^2$$

This is separable first order ode. Integrating

$$\int dx = \int \frac{1}{z^2 + 1} dz$$
$$x + c_1 = \arctan(z)$$

Replacing  $z$  back by its value from (1) then the above gives the solution as

$$y = -x + \tan(x + c_1)$$

$$y = -x + \tan(x + c_1)$$

### Summary

The solution(s) found are the following

$$y = -x + \tan(x + c_1) \tag{1}$$

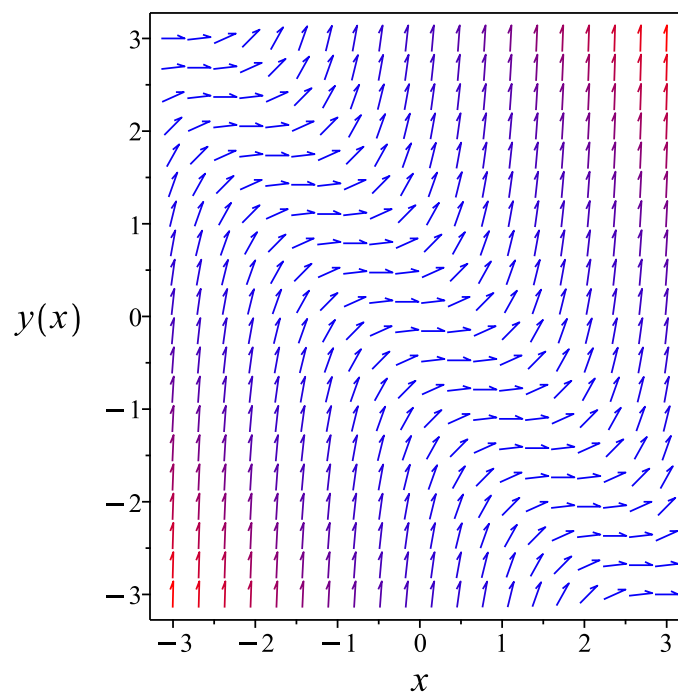


Figure 19: Slope field plot

### Verification of solutions

$$y = -x + \tan(x + c_1)$$

Verified OK.

### 1.6.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = (y + x)^2$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **homogeneous Type C**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 9: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 1 \\ \eta(x, y) &= -1\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Therefore

$$\begin{aligned}\frac{dy}{dx} &= \frac{\eta}{\xi} \\ &= \frac{-1}{1} \\ &= -1\end{aligned}$$

This is easily solved to give

$$y = -x + c_1$$

Where now the coordinate  $R$  is taken as the constant of integration. Hence

$$R = y + x$$

And  $S$  is found from

$$\begin{aligned}dS &= \frac{dx}{\xi} \\ &= \frac{dx}{1}\end{aligned}$$

Integrating gives

$$\begin{aligned}S &= \int \frac{dx}{1} \\ &= x\end{aligned}$$

Where the constant of integration is set to zero as we just need one solution. Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = (y + x)^2$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 1$$

$$S_x = 1$$

$$S_y = 0$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{1 + (y + x)^2} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R^2 + 1}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = \arctan(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$x = \arctan(y + x) + c_1$$

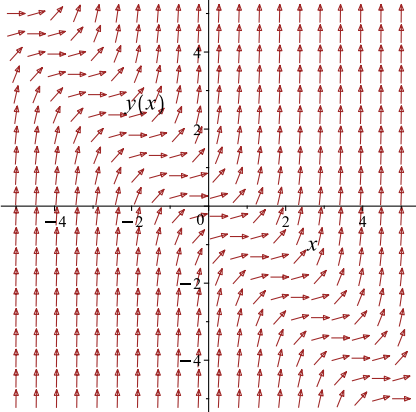
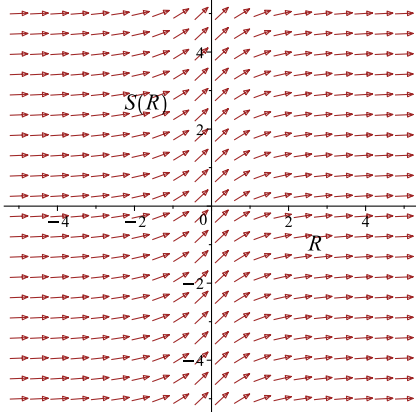
Which simplifies to

$$x = \arctan(y + x) + c_1$$

Which gives

$$y = -x - \tan(-x + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = (y + x)^2$ 	$R = y + x$ $S = x$	$\frac{dS}{dR} = \frac{1}{R^2+1}$ 

### Summary

The solution(s) found are the following

$$y = -x - \tan(-x + c_1) \tag{1}$$

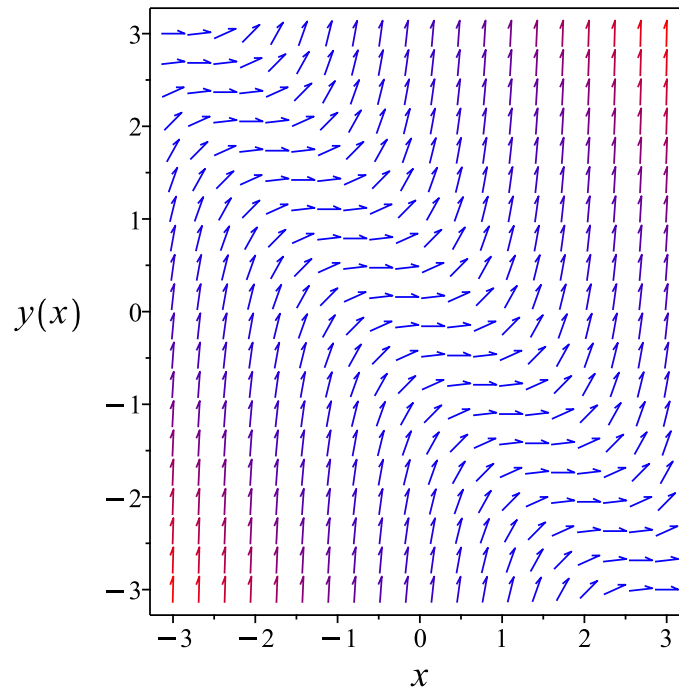


Figure 20: Slope field plot

Verification of solutions

$$y = -x - \tan(-x + c_1)$$

Verified OK.

### 1.6.3 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= (y + x)^2 \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = x^2 + 2xy + y^2$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that  $f_0(x) = x^2$ ,  $f_1(x) = 2x$  and  $f_2(x) = 1$ . Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{u} \end{aligned} \quad (1)$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for  $u(x)$  which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \quad (2)$$

But

$$\begin{aligned} f_2' &= 0 \\ f_1 f_2 &= 2x \\ f_2^2 f_0 &= x^2 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$u''(x) - 2xu'(x) + x^2 u(x) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = e^{\frac{x^2}{2}} (c_1 \cos(x) + c_2 \sin(x))$$

The above shows that

$$u'(x) = e^{\frac{x^2}{2}} ((c_1 x + c_2) \cos(x) + \sin(x) (c_2 x - c_1))$$

Using the above in (1) gives the solution

$$y = -\frac{(c_1 x + c_2) \cos(x) + \sin(x) (c_2 x - c_1)}{c_1 \cos(x) + c_2 \sin(x)}$$

Dividing both numerator and denominator by  $c_1$  gives, after renaming the constant  $\frac{c_2}{c_1} = c_3$  the following solution

$$y = \frac{(-c_3 x - 1) \cos(x) - \sin(x) (-c_3 + x)}{c_3 \cos(x) + \sin(x)}$$



### Summary

The solution(s) found are the following

$$y = \frac{(-c_3x - 1) \cos(x) - \sin(x)(-c_3 + x)}{c_3 \cos(x) + \sin(x)} \quad (1)$$

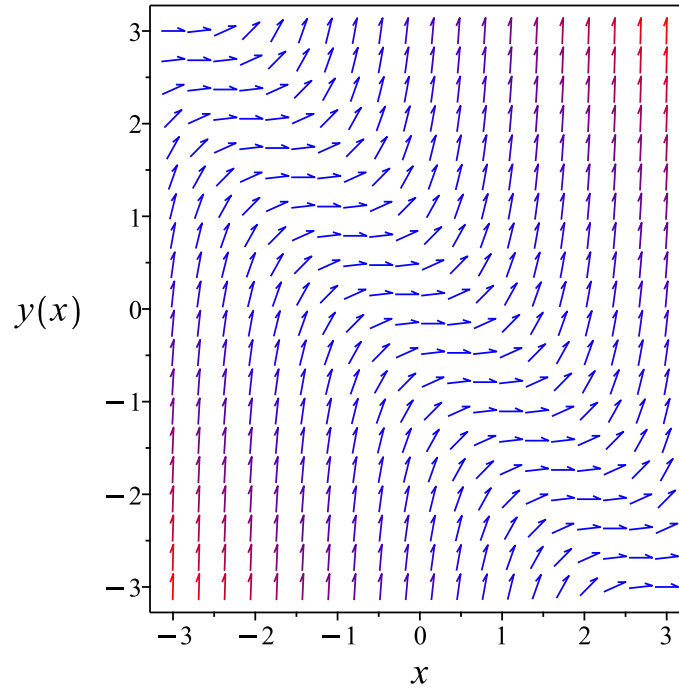


Figure 21: Slope field plot

### Verification of solutions

$$y = \frac{(-c_3x - 1) \cos(x) - \sin(x)(-c_3 + x)}{c_3 \cos(x) + \sin(x)}$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous C  
1st order, trying the canonical coordinates of the invariance group  
<- 1st order, canonical coordinates successful  
<- homogeneous successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 16

```
dsolve(diff(y(x),x)=(x+y(x))^2,y(x), singsol=all)
```

$$y(x) = -x - \tan(c_1 - x)$$

### ✓ Solution by Mathematica

Time used: 0.735 (sec). Leaf size: 14

```
DSolve[y'[x]==(x+y[x])^2,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -x + \tan(x + c_1)$$

## 1.7 problem 3.b

1.7.1 Solving as first order ode lie symmetry calculated ode . . . . . 82

Internal problem ID [3086]

Internal file name [OUTPUT/2578\_Sunday\_June\_05\_2022\_03\_20\_22\_AM\_29566835/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 7, page 37

**Problem number:** 3.b.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**first\_order\_ode\_lie\_symmetry\_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class C`], _dAlembert]
```

$$y' - \sin(x - y + 1)^2 = 0$$

### 1.7.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \sin(x - y + 1)^2$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} & b_2 + \sin(x-y+1)^2(b_3 - a_2) - \sin(x-y+1)^4 a_3 \\ & - 2 \sin(x-y+1) \cos(x-y+1)(xa_2 + ya_3 + a_1) \\ & + 2 \sin(x-y+1) \cos(x-y+1)(xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\begin{aligned} & - \sin(x-y+1)^4 a_3 - 2 \sin(x-y+1) \cos(x-y+1) xa_2 \\ & + 2 \sin(x-y+1) \cos(x-y+1) xb_2 - 2 \sin(x-y+1) \cos(x-y+1) ya_3 \\ & + 2 \sin(x-y+1) \cos(x-y+1) yb_3 - \sin(x-y+1)^2 a_2 + \sin(x-y+1)^2 b_3 \\ & - 2 \sin(x-y+1) \cos(x-y+1) a_1 + 2 \sin(x-y+1) \cos(x-y+1) b_1 + b_2 = 0 \end{aligned}$$

Setting the numerator to zero gives

$$\begin{aligned} & - \sin(x-y+1)^4 a_3 - 2 \sin(x-y+1) \cos(x-y+1) xa_2 \\ & + 2 \sin(x-y+1) \cos(x-y+1) xb_2 - 2 \sin(x-y+1) \cos(x-y+1) ya_3 \\ & + 2 \sin(x-y+1) \cos(x-y+1) yb_3 - \sin(x-y+1)^2 a_2 \\ & + \sin(x-y+1)^2 b_3 - 2 \sin(x-y+1) \cos(x-y+1) a_1 \\ & + 2 \sin(x-y+1) \cos(x-y+1) b_1 + b_2 = 0 \end{aligned} \quad (6E)$$

Simplifying the above gives

$$\begin{aligned} & b_2 - \frac{3a_3}{8} - \frac{a_2}{2} + \frac{b_3}{2} + \frac{a_3 \cos(2x-2y+2)}{2} - \frac{a_3 \cos(4x-4y+4)}{8} \\ & - xa_2 \sin(2x-2y+2) + xb_2 \sin(2x-2y+2) - ya_3 \sin(2x-2y+2) \\ & + yb_3 \sin(2x-2y+2) + \frac{a_2 \cos(2x-2y+2)}{2} - \frac{b_3 \cos(2x-2y+2)}{2} \\ & - a_1 \sin(2x-2y+2) + b_1 \sin(2x-2y+2) = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y, \cos(2x-2y+2), \cos(4x-4y+4), \sin(2x-2y+2)\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2, \cos(2x-2y+2) = v_3, \cos(4x-4y+4) = v_4, \sin(2x-2y+2) = v_5\}$$

The above PDE (6E) now becomes

$$\begin{aligned}
 b_2 - \frac{3}{8}a_3 - \frac{1}{2}a_2 + \frac{1}{2}b_3 + \frac{1}{2}a_3v_3 - \frac{1}{8}a_3v_4 - v_1a_2v_5 + v_1b_2v_5 \\
 - v_2a_3v_5 + v_2b_3v_5 + \frac{1}{2}a_2v_3 - \frac{1}{2}b_3v_3 - a_1v_5 + b_1v_5 = 0
 \end{aligned} \tag{7E}$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2, v_3, v_4, v_5\}$$

Equation (7E) now becomes

$$\begin{aligned}
 b_2 - \frac{3a_3}{8} - \frac{a_2}{2} + \frac{b_3}{2} + \left(\frac{a_3}{2} + \frac{a_2}{2} - \frac{b_3}{2}\right)v_3 - \frac{a_3v_4}{8} \\
 + (-a_1 + b_1)v_5 + (-a_2 + b_2)v_5v_1 + (-a_3 + b_3)v_5v_2 = 0
 \end{aligned} \tag{8E}$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}
 -\frac{a_3}{8} &= 0 \\
 -a_1 + b_1 &= 0 \\
 -a_2 + b_2 &= 0 \\
 -a_3 + b_3 &= 0 \\
 \frac{a_3}{2} + \frac{a_2}{2} - \frac{b_3}{2} &= 0 \\
 b_2 - \frac{3a_3}{8} - \frac{a_2}{2} + \frac{b_3}{2} &= 0
 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= b_1 \\
 a_2 &= 0 \\
 a_3 &= 0 \\
 b_1 &= b_1 \\
 b_2 &= 0 \\
 b_3 &= 0
 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= 1 \\ \eta &= 1\end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= 1 - (\sin(x - y + 1))^2 (1) \\ &= 1 - \sin(x)^2 \cos(y)^2 \cos(1)^2 - 2 \sin(x)^2 \cos(y) \cos(1) \sin(y) \sin(1) + 2 \sin(x) \cos(y) \cos(1)^2 \cos(x) \sin(1) \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{1 - \sin(x)^2 \cos(y)^2 \cos(1)^2 - 2 \sin(x)^2 \cos(y) \cos(1) \sin(y) \sin(1) + 2 \sin(x) \cos(y) \cos(1)^2 \cos(x) \sin(1)} dy\end{aligned}$$

Which results in

$$S = -\tan(x - y + 1)$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y) S_y}{R_x + \omega(x, y) R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \sin(x - y + 1)^2$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 0$$

$$S_x = -\sec(x - y + 1)^2$$

$$S_y = \sec(x - y + 1)^2$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -1 \tag{2A}$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -R + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$-\tan(x - y + 1) = -x + c_1$$

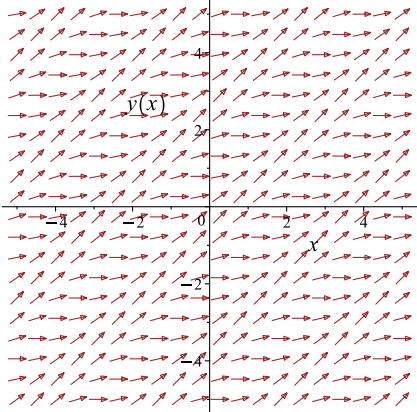
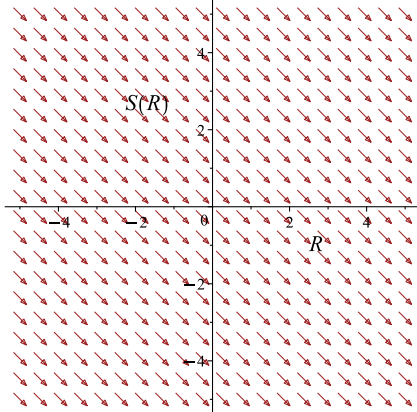
Which simplifies to

$$-\tan(x - y + 1) = -x + c_1$$

Which gives

$$y = x + 1 + \arctan(-x + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \sin(x - y + 1)^2$ 	$R = x$ $S = -\tan(x - y + 1)$	$\frac{dS}{dR} = -1$ 

### Summary

The solution(s) found are the following

$$y = x + 1 + \arctan(-x + c_1) \tag{1}$$



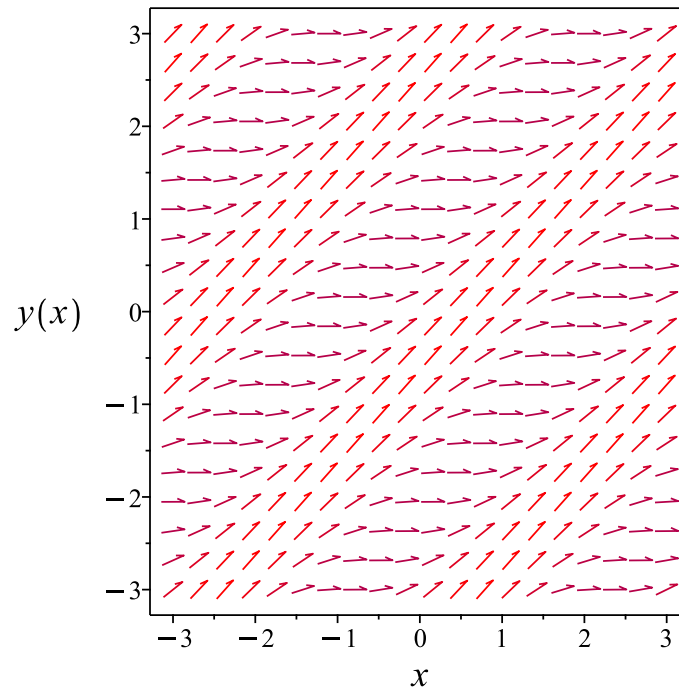


Figure 22: Slope field plot

Verification of solutions

$$y = x + 1 + \arctan(-x + c_1)$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous C
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 13

```
dsolve(diff(y(x),x)=sin(x-y(x)+1)^2,y(x), singsol=all)
```

$$y(x) = x + 1 + \arctan(c_1 - x)$$

✓ Solution by Mathematica

Time used: 0.344 (sec). Leaf size: 33

```
DSolve[y'[x]==Sin[x-y[x]+1]^2,y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve}[2y(x) - 2(\tan(-y(x) + x + 1) - \arctan(\tan(-y(x) + x + 1))) = c_1, y(x)]$$

## 1.8 problem 5.a

- 1.8.1 Solving as homogeneousTypeMapleC ode . . . . . 90
- 1.8.2 Solving as first order ode lie symmetry calculated ode . . . . . 93

Internal problem ID [3087]

Internal file name [OUTPUT/2579\_Sunday\_June\_05\_2022\_03\_20\_34\_AM\_14083881/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 7, page 37

**Problem number:** 5.a.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**homogeneousTypeMapleC**",  
"**first\_order\_ode\_lie\_symmetry\_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class C`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$y' - \frac{x + y + 4}{x - y - 6} = 0$$

### 1.8.1 Solving as homogeneousTypeMapleC ode

Let  $Y = y + y_0$  and  $X = x + x_0$  then the above is transformed to new ode in  $Y(X)$

$$\frac{d}{dX}Y(X) = -\frac{X + x_0 + Y(X) + y_0 + 4}{-X - x_0 + Y(X) + y_0 + 6}$$

Solving for possible values of  $x_0$  and  $y_0$  which makes the above ode a homogeneous ode results in

$$\begin{aligned}x_0 &= 1 \\y_0 &= -5\end{aligned}$$

Using these values now it is possible to easily solve for  $Y(X)$ . The above ode now becomes

$$\frac{d}{dX}Y(X) = -\frac{X + Y(X)}{-X + Y(X)}$$

In canonical form, the ODE is

$$\begin{aligned} Y' &= F(X, Y) \\ &= -\frac{X + Y}{-X + Y} \end{aligned} \quad (1)$$

An ode of the form  $Y' = \frac{M(X, Y)}{N(X, Y)}$  is called homogeneous if the functions  $M(X, Y)$  and  $N(X, Y)$  are both homogeneous functions and of the same order. Recall that a function  $f(X, Y)$  is homogeneous of order  $n$  if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both  $M = X + Y$  and  $N = X - Y$  are both homogeneous and of the same order  $n = 1$ . Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution  $u = \frac{Y}{X}$ , or  $Y = uX$ . Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation  $Y = uX$  to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dX}X + u &= \frac{-u - 1}{u - 1} \\ \frac{du}{dX} &= \frac{\frac{-u(X)-1}{u(X)-1} - u(X)}{X} \end{aligned}$$

Or

$$\frac{d}{dX}u(X) - \frac{\frac{-u(X)-1}{u(X)-1} - u(X)}{X} = 0$$

Or

$$\left(\frac{d}{dX}u(X)\right)Xu(X) - \left(\frac{d}{dX}u(X)\right)X + u(X)^2 + 1 = 0$$

Or

$$X(u(X) - 1) \left(\frac{d}{dX}u(X)\right) + u(X)^2 + 1 = 0$$

Which is now solved as separable in  $u(X)$ . Which is now solved in  $u(X)$ . In canonical form the ODE is

$$\begin{aligned} u' &= F(X, u) \\ &= f(X)g(u) \\ &= -\frac{u^2 + 1}{X(u - 1)} \end{aligned}$$

Where  $f(X) = -\frac{1}{X}$  and  $g(u) = \frac{u^2+1}{u-1}$ . Integrating both sides gives

$$\frac{1}{\frac{u^2+1}{u-1}} du = -\frac{1}{X} dX$$

$$\int \frac{1}{\frac{u^2+1}{u-1}} du = \int -\frac{1}{X} dX$$

$$\frac{\ln(u^2 + 1)}{2} - \arctan(u) = -\ln(X) + c_2$$

The solution is

$$\frac{\ln(u(X)^2 + 1)}{2} - \arctan(u(X)) + \ln(X) - c_2 = 0$$

Now  $u$  in the above solution is replaced back by  $Y$  using  $u = \frac{Y}{X}$  which results in the solution

$$\frac{\ln\left(\frac{Y(X)^2}{X^2} + 1\right)}{2} - \arctan\left(\frac{Y(X)}{X}\right) + \ln(X) - c_2 = 0$$

Using the solution for  $Y(X)$

$$\frac{\ln\left(\frac{Y(X)^2}{X^2} + 1\right)}{2} - \arctan\left(\frac{Y(X)}{X}\right) + \ln(X) - c_2 = 0$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y - 5$$

$$X = x + 1$$

Then the solution in  $y$  becomes

$$\frac{\ln\left(\frac{(y+5)^2}{(x-1)^2} + 1\right)}{2} - \arctan\left(\frac{y+5}{x-1}\right) + \ln(x-1) - c_2 = 0$$

### Summary

The solution(s) found are the following

$$\frac{\ln\left(\frac{(y+5)^2}{(x-1)^2} + 1\right)}{2} - \arctan\left(\frac{y+5}{x-1}\right) + \ln(x-1) - c_2 = 0 \quad (1)$$

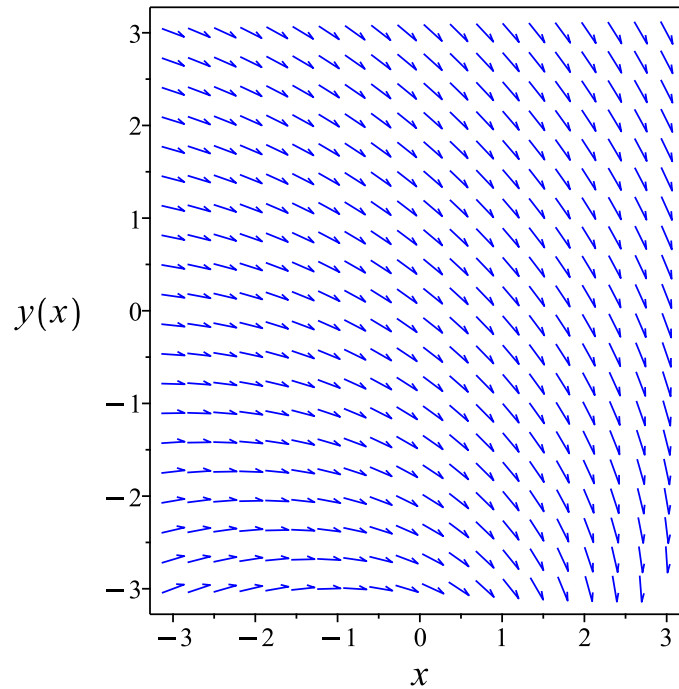


Figure 23: Slope field plot

### Verification of solutions

$$\frac{\ln\left(\frac{(y+5)^2}{(x-1)^2} + 1\right)}{2} - \arctan\left(\frac{y+5}{x-1}\right) + \ln(x-1) - c_2 = 0$$

Verified OK.

### **1.8.2 Solving as first order ode lie symmetry calculated ode**

Writing the ode as

$$y' = -\frac{x+y+4}{-x+y+6}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} b_2 - \frac{(x+y+4)(b_3 - a_2)}{-x+y+6} - \frac{(x+y+4)^2 a_3}{(-x+y+6)^2} \\ - \left( -\frac{1}{-x+y+6} - \frac{x+y+4}{(-x+y+6)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left( -\frac{1}{-x+y+6} + \frac{x+y+4}{(-x+y+6)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (\text{5E})$$

Putting the above in normal form gives

$$\frac{x^2 a_2 + x^2 a_3 + x^2 b_2 - x^2 b_3 - 2xy a_2 + 2xy a_3 + 2xy b_2 + 2xy b_3 - y^2 a_2 - y^2 a_3 - y^2 b_2 + y^2 b_3 - 12xa_2 + 8xa_3 - 12xb_2 + 8xb_3 - 10ya_2 + 10ya_3 - 10yb_2 + 10yb_3 - 10a_1 + 10a_2 - 10a_3 - 10b_1 + 10b_2 - 10b_3}{(x - y + 6)^3} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} -x^2 a_2 - x^2 a_3 - x^2 b_2 + x^2 b_3 + 2xy a_2 - 2xy a_3 - 2xy b_2 - 2xy b_3 + y^2 a_2 \\ + y^2 a_3 + y^2 b_2 - y^2 b_3 + 12xa_2 - 8xa_3 - 2xb_1 - 10xb_2 - 2xb_3 + 2ya_1 + 10ya_2 \\ + 2ya_3 + 12yb_2 - 8yb_3 + 10a_1 + 24a_2 - 16a_3 + 2b_1 + 36b_2 - 24b_3 = 0 \end{aligned} \quad (\text{6E})$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & -a_2v_1^2 + 2a_2v_1v_2 + a_2v_2^2 - a_3v_1^2 - 2a_3v_1v_2 + a_3v_2^2 - b_2v_1^2 - 2b_2v_1v_2 + b_2v_2^2 \\ & + b_3v_1^2 - 2b_3v_1v_2 - b_3v_2^2 + 2a_1v_2 + 12a_2v_1 + 10a_2v_2 - 8a_3v_1 + 2a_3v_2 - 2b_1v_1 \\ & - 10b_2v_1 + 12b_2v_2 - 2b_3v_1 - 8b_3v_2 + 10a_1 + 24a_2 - 16a_3 + 2b_1 + 36b_2 - 24b_3 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & (-a_2 - a_3 - b_2 + b_3)v_1^2 + (2a_2 - 2a_3 - 2b_2 - 2b_3)v_1v_2 \\ & + (12a_2 - 8a_3 - 2b_1 - 10b_2 - 2b_3)v_1 + (a_2 + a_3 + b_2 - b_3)v_2^2 \\ & + (2a_1 + 10a_2 + 2a_3 + 12b_2 - 8b_3)v_2 + 10a_1 + 24a_2 - 16a_3 + 2b_1 + 36b_2 - 24b_3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -a_2 - a_3 - b_2 + b_3 &= 0 \\ a_2 + a_3 + b_2 - b_3 &= 0 \\ 2a_2 - 2a_3 - 2b_2 - 2b_3 &= 0 \\ 2a_1 + 10a_2 + 2a_3 + 12b_2 - 8b_3 &= 0 \\ 12a_2 - 8a_3 - 2b_1 - 10b_2 - 2b_3 &= 0 \\ 10a_1 + 24a_2 - 16a_3 + 2b_1 + 36b_2 - 24b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= -5b_2 - b_3 \\ a_2 &= b_3 \\ a_3 &= -b_2 \\ b_1 &= -b_2 + 5b_3 \\ b_2 &= b_2 \\ b_3 &= b_3 \end{aligned}$$



Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= x - 1 \\ \eta &= y + 5\end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y + 5 - \left( -\frac{x + y + 4}{-x + y + 6} \right) (x - 1) \\ &= \frac{-x^2 - y^2 + 2x - 10y - 26}{x - y - 6} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left( \xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-x^2 - y^2 + 2x - 10y - 26}{x - y - 6}} dy\end{aligned}$$

Which results in

$$S = \frac{\ln(x^2 + y^2 - 2x + 10y + 26)}{2} + \frac{2(1 - x) \arctan\left(\frac{2y+10}{2x-2}\right)}{2x - 2}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y) S_y}{R_x + \omega(x, y) R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{x + y + 4}{-x + y + 6}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{x + y + 4}{x^2 + y^2 - 2x + 10y + 26} \\ S_y &= \frac{-x + y + 6}{x^2 + y^2 - 2x + 10y + 26} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \tag{2A}$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

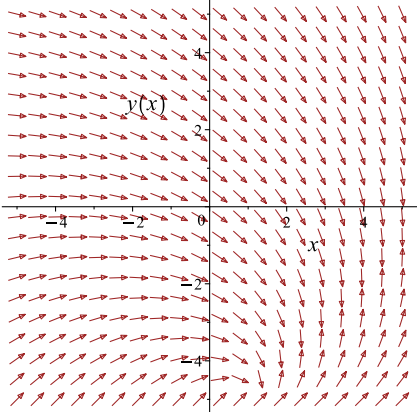
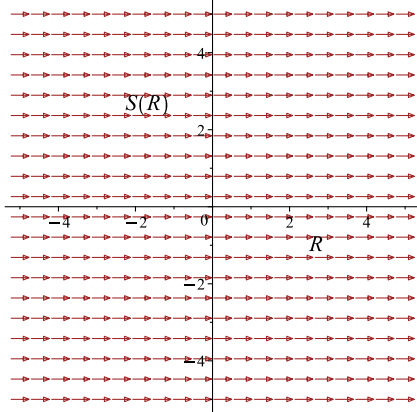
To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{\ln(y^2 + x^2 + 10y - 2x + 26)}{2} - \arctan\left(\frac{y + 5}{x - 1}\right) = c_1$$

Which simplifies to

$$\frac{\ln(y^2 + x^2 + 10y - 2x + 26)}{2} - \arctan\left(\frac{y + 5}{x - 1}\right) = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{x+y+4}{-x+y+6}$ 	$R = x$ $S = \frac{\ln(x^2 + y^2 - 2x + 1)}{2}$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$\frac{\ln(y^2 + x^2 + 10y - 2x + 26)}{2} - \arctan\left(\frac{y + 5}{x - 1}\right) = c_1 \quad (1)$$

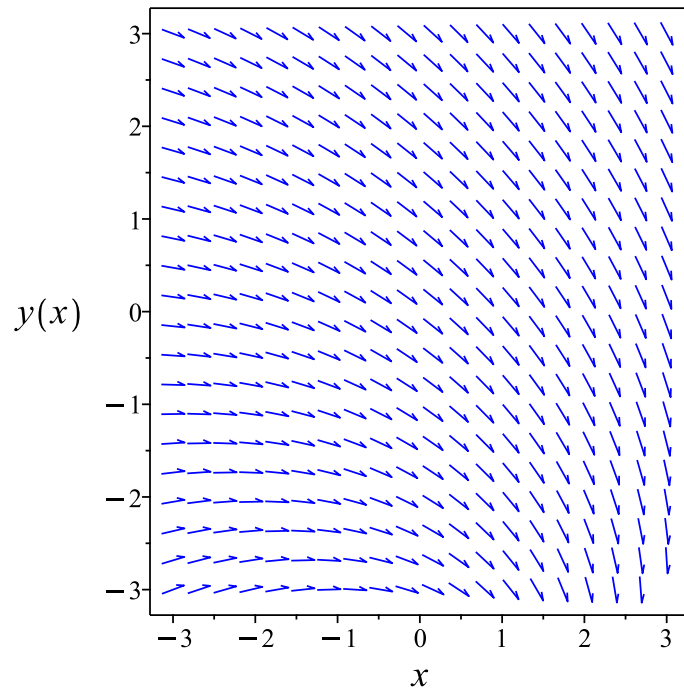


Figure 24: Slope field plot

Verification of solutions

$$\frac{\ln(y^2 + x^2 + 10y - 2x + 26)}{2} - \arctan\left(\frac{y + 5}{x - 1}\right) = c_1$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous C  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful  
<- homogeneous successful`
```

### ✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 31

```
dsolve(diff(y(x),x)=(x+y(x)+4)/(x-y(x)-6),y(x), singsol=all)
```

$$y(x) = -5 - \tan(\text{RootOf}(2\_Z + \ln(\sec(\_Z)^2) + 2\ln(x-1) + 2c_1))(x-1)$$

### ✓ Solution by Mathematica

Time used: 0.057 (sec). Leaf size: 58

```
DSolve[y'[x]==(x+y[x]+4)/(x-y[x]-6),y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve}\left[2 \arctan\left(\frac{y(x) + x + 4}{y(x) - x + 6}\right) + \log\left(\frac{x^2 + y(x)^2 + 10y(x) - 2x + 26}{2(x-1)^2}\right) + 2 \log(x-1) + c_1 = 0, y(x)\right]$$

## 1.9 problem 5.b

1.9.1 Solving as first order ode lie symmetry calculated ode . . . . . 101

Internal problem ID [3088]

Internal file name [OUTPUT/2580\_Sunday\_June\_05\_2022\_03\_20\_38\_AM\_82893986/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 7, page 37

**Problem number:** 5.b.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**first\_order\_ode\_lie\_symmetry\_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class C`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$y' - \frac{x + y + 4}{x + y - 6} = 0$$

### 1.9.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{x + y + 4}{x + y - 6}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} b_2 + \frac{(x+y+4)(b_3-a_2)}{x+y-6} - \frac{(x+y+4)^2 a_3}{(x+y-6)^2} \\ - \left( \frac{1}{x+y-6} - \frac{x+y+4}{(x+y-6)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left( \frac{1}{x+y-6} - \frac{x+y+4}{(x+y-6)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{x^2 a_2 + x^2 a_3 - x^2 b_2 - x^2 b_3 + 2xy a_2 + 2xy a_3 - 2xy b_2 - 2xy b_3 + y^2 a_2 + y^2 a_3 - y^2 b_2 - y^2 b_3 - 12xa_2 + 8xa_3 - 12yb_2 + 8yb_3}{(x+y-6)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} -x^2 a_2 - x^2 a_3 + x^2 b_2 + x^2 b_3 - 2xy a_2 - 2xy a_3 + 2xy b_2 + 2xy b_3 - y^2 a_2 \\ - y^2 a_3 + y^2 b_2 + y^2 b_3 + 12xa_2 - 8xa_3 - 2xb_2 - 2xb_3 + 2ya_2 + 2ya_3 \\ - 12yb_2 + 8yb_3 + 10a_1 + 24a_2 - 16a_3 + 10b_1 + 36b_2 - 24b_3 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} -a_2 v_1^2 - 2a_2 v_1 v_2 - a_2 v_2^2 - a_3 v_1^2 - 2a_3 v_1 v_2 - a_3 v_2^2 + b_2 v_1^2 + 2b_2 v_1 v_2 \\ + b_2 v_2^2 + b_3 v_1^2 + 2b_3 v_1 v_2 + b_3 v_2^2 + 12a_2 v_1 + 2a_2 v_2 - 8a_3 v_1 + 2a_3 v_2 - 2b_2 v_1 \\ - 12b_2 v_2 - 2b_3 v_1 + 8b_3 v_2 + 10a_1 + 24a_2 - 16a_3 + 10b_1 + 36b_2 - 24b_3 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} &(-a_2 - a_3 + b_2 + b_3) v_1^2 + (-2a_2 - 2a_3 + 2b_2 + 2b_3) v_1 v_2 \\ &+ (12a_2 - 8a_3 - 2b_2 - 2b_3) v_1 + (-a_2 - a_3 + b_2 + b_3) v_2^2 \\ &+ (2a_2 + 2a_3 - 12b_2 + 8b_3) v_2 + 10a_1 + 24a_2 - 16a_3 + 10b_1 + 36b_2 - 24b_3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -2a_2 - 2a_3 + 2b_2 + 2b_3 &= 0 \\ -a_2 - a_3 + b_2 + b_3 &= 0 \\ 2a_2 + 2a_3 - 12b_2 + 8b_3 &= 0 \\ 12a_2 - 8a_3 - 2b_2 - 2b_3 &= 0 \\ 10a_1 + 24a_2 - 16a_3 + 10b_1 + 36b_2 - 24b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= -2b_3 - b_1 \\ a_2 &= b_3 \\ a_3 &= b_3 \\ b_1 &= b_1 \\ b_2 &= b_3 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= -1 \\ \eta &= 1 \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= 1 - \left( \frac{x + y + 4}{x + y - 6} \right) (-1) \\ &= \frac{2x + 2y - 2}{x + y - 6} \\ \xi &= 0 \end{aligned}$$



The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{2x+2y-2}{x+y-6}} dy \end{aligned}$$

Which results in

$$S = \frac{y}{2} - \frac{5 \ln(x-1+y)}{2}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x + y + 4}{x + y - 6}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -\frac{5}{2x + 2y - 2} \\ S_y &= \frac{x + y - 6}{2x + 2y - 2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{2} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = \frac{R}{2} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{y}{2} - \frac{5 \ln(y + x - 1)}{2} = \frac{x}{2} + c_1$$

Which simplifies to

$$\frac{y}{2} - \frac{5 \ln(y + x - 1)}{2} = \frac{x}{2} + c_1$$

Which gives

$$y = -5 \text{LambertW} \left( -\frac{e^{-\frac{2x}{5} + \frac{1}{5} - \frac{2c_1}{5}}}{5} \right) - x + 1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{x+y+4}{x+y-6}$	$R = x$ $S = \frac{y}{2} - \frac{5 \ln(x-1+y)}{2}$	$\frac{dS}{dR} = \frac{1}{2}$

Summary

The solution(s) found are the following

$$y = -5 \text{LambertW} \left( -\frac{e^{-\frac{2x}{5} + \frac{1}{5} - \frac{2c_1}{5}}}{5} \right) - x + 1 \tag{1}$$

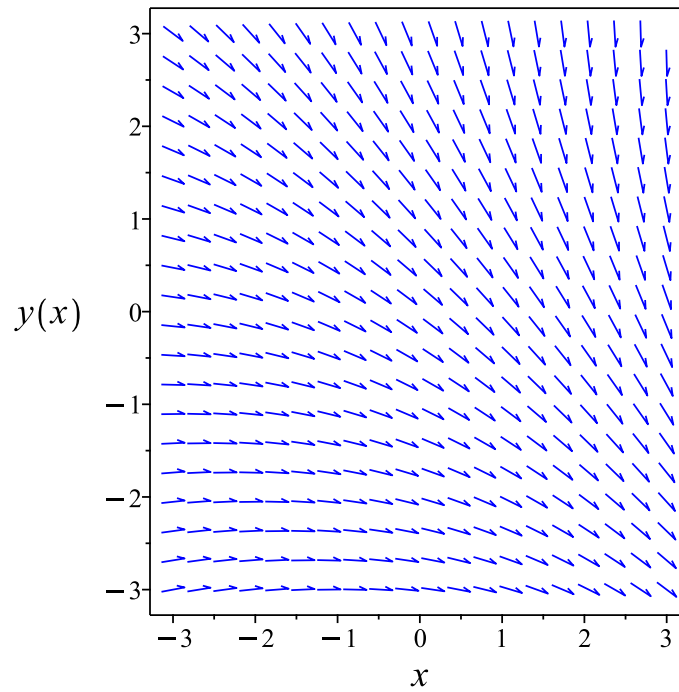


Figure 25: Slope field plot

Verification of solutions

$$y = -5 \operatorname{LambertW} \left( -\frac{e^{-\frac{2x}{5} + \frac{1}{5} - \frac{2c_1}{5}}}{5} \right) - x + 1$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous C  
1st order, trying the canonical coordinates of the invariance group  
<- 1st order, canonical coordinates successful  
<- homogeneous successful`
```

### ✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 21

```
dsolve(diff(y(x),x)=(x+y(x)+4)/(x+y(x)-6),y(x), singsol=all)
```

$$y(x) = -x - 5 \operatorname{LambertW}\left(-\frac{c_1 e^{-\frac{2x}{5} + \frac{1}{5}}}{5}\right) + 1$$

### ✓ Solution by Mathematica

Time used: 4.019 (sec). Leaf size: 35

```
DSolve[y'[x]==(x+y[x]+4)/(x+y[x]-6),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -5W\left(-e^{-\frac{2x}{5}-1+c_1}\right) - x + 1$$
$$y(x) \rightarrow 1 - x$$

## **2 Chapter 2, section 8, page 41**

2.1	problem 1	110
2.2	problem 2	122
2.3	problem 3	129
2.4	problem 4	138
2.5	problem 5	141
2.6	problem 6	157
2.7	problem 7	169
2.8	problem 8	175
2.9	problem 9	189
2.10	problem 10	207
2.11	problem 11	214

## 2.1 problem 1

2.1.1	Solving as first order ode lie symmetry calculated ode . . . . .	110
2.1.2	Solving as exact ode . . . . .	115
2.1.3	Maple step by step solution . . . . .	119

Internal problem ID [3089]

Internal file name [OUTPUT/2581\_Sunday\_June\_05\_2022\_03\_20\_40\_AM\_2961885/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 1.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"exact", "first\_order\_ode\_lie\_symmetry\_calculated"**

Maple gives the following as the ode type

```
[[_homogeneous, `class G`], _exact, _rational, [_Abel, `2nd
  type`, `class B`]]
```

$$\left(x + \frac{2}{y}\right) y' + y = 0$$

### 2.1.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{y^2}{xy + 2}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$b_2 - \frac{y^2(b_3 - a_2)}{xy + 2} - \frac{y^4 a_3}{(xy + 2)^2} - \frac{y^3(xa_2 + ya_3 + a_1)}{(xy + 2)^2} - \left( -\frac{2y}{xy + 2} + \frac{y^2 x}{(xy + 2)^2} \right) (xb_2 + yb_3 + b_1) = 0 \quad (5E)$$

Putting the above in normal form gives

$$\frac{2x^2 y^2 b_2 - 2y^4 a_3 + x y^2 b_1 - y^3 a_1 + 8xyb_2 + 2y^2 a_2 + 2y^2 b_3 + 4yb_1 + 4b_2}{(xy + 2)^2} = 0$$

Setting the numerator to zero gives

$$2x^2 y^2 b_2 - 2y^4 a_3 + x y^2 b_1 - y^3 a_1 + 8xyb_2 + 2y^2 a_2 + 2y^2 b_3 + 4yb_1 + 4b_2 = 0 \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$-2a_3 v_2^4 + 2b_2 v_1^2 v_2^2 - a_1 v_2^3 + b_1 v_1 v_2^2 + 2a_2 v_2^2 + 8b_2 v_1 v_2 + 2b_3 v_2^2 + 4b_1 v_2 + 4b_2 = 0 \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$2b_2 v_1^2 v_2^2 + b_1 v_1 v_2^2 + 8b_2 v_1 v_2 - 2a_3 v_2^4 - a_1 v_2^3 + (2a_2 + 2b_3) v_2^2 + 4b_1 v_2 + 4b_2 = 0 \quad (8E)$$



Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}b_1 &= 0 \\-a_1 &= 0 \\-2a_3 &= 0 \\4b_1 &= 0 \\2b_2 &= 0 \\4b_2 &= 0 \\8b_2 &= 0 \\2a_2 + 2b_3 &= 0\end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}a_1 &= 0 \\a_2 &= -b_3 \\a_3 &= 0 \\b_1 &= 0 \\b_2 &= 0 \\b_3 &= b_3\end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= -x \\ \eta &= y\end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left( -\frac{y^2}{xy + 2} \right) (-x) \\ &= \frac{2y}{xy + 2} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{2y}{xy+2}} dy \end{aligned}$$

Which results in

$$S = \frac{xy}{2} + \ln(y)$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y^2}{xy+2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{y}{2} \\ S_y &= \frac{x}{2} + \frac{1}{y} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{yx}{2} + \ln(y) = c_1$$

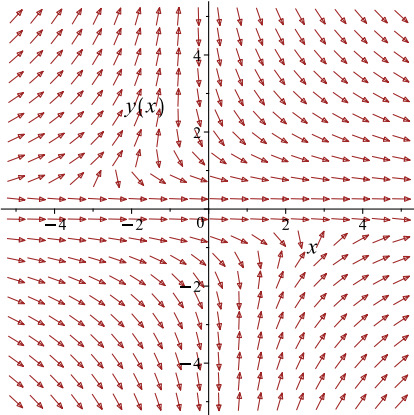
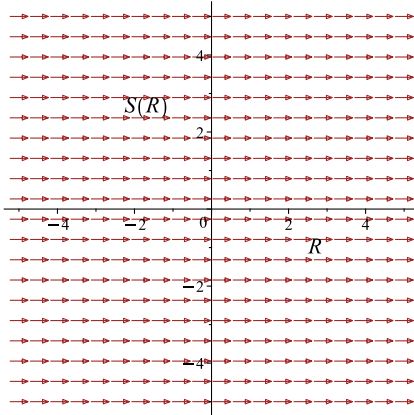
Which simplifies to

$$\frac{yx}{2} + \ln(y) = c_1$$

Which gives

$$y = e^{-\text{LambertW}\left(\frac{e^{c_1}x}{2}\right) + c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{y^2}{xy+2}$ 	$R = x$ $S = \frac{xy}{2} + \ln(y)$	$\frac{dS}{dR} = 0$ 

## Summary

The solution(s) found are the following

$$y = e^{-\text{LambertW}\left(\frac{e^{c_1}x}{2}\right) + c_1} \quad (1)$$

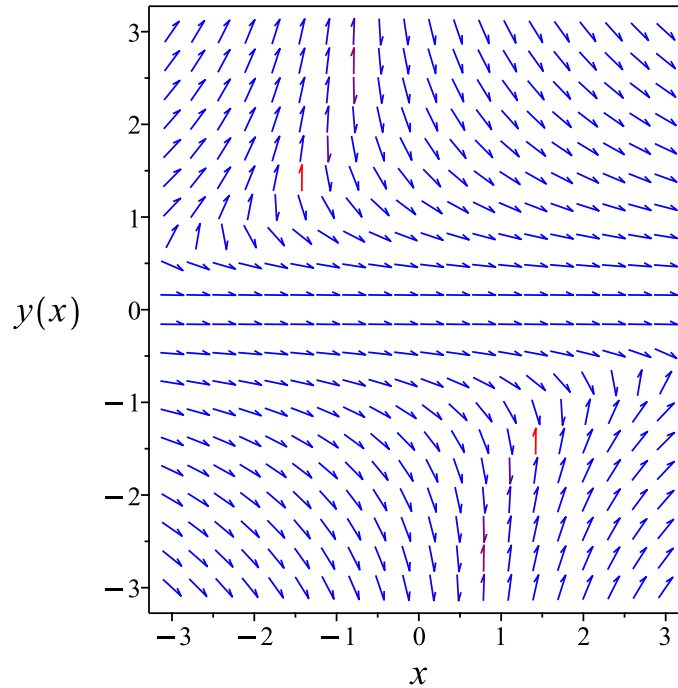


Figure 26: Slope field plot

## Verification of solutions

$$y = e^{-\text{LambertW}\left(\frac{e^{c_1}x}{2}\right) + c_1}$$

Verified OK.

### 2.1.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} \left(x + \frac{2}{y}\right) dy &= (-y) dx \\ (y) dx + \left(x + \frac{2}{y}\right) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= y \\ N(x, y) &= x + \frac{2}{y} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y) \\ &= 1 \end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left( x + \frac{2}{y} \right) \\ &= 1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int y dx \\ \phi &= xy + f(y)\end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = x + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = x + \frac{2}{y}$ . Therefore equation (4) becomes

$$x + \frac{2}{y} = x + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = \frac{2}{y}$$

Integrating the above w.r.t  $y$  gives

$$\begin{aligned}\int f'(y) dy &= \int \left( \frac{2}{y} \right) dy \\ f(y) &= 2 \ln(y) + c_1\end{aligned}$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = xy + 2 \ln(y) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = xy + 2 \ln(y)$$

The solution becomes

$$y = e^{-\text{LambertW}\left(\frac{x e^{\frac{c_1}{2}}}{2}\right) + \frac{c_1}{2}}$$

### Summary

The solution(s) found are the following

$$y = e^{-\text{LambertW}\left(\frac{x e^{\frac{c_1}{2}}}{2}\right) + \frac{c_1}{2}} \tag{1}$$

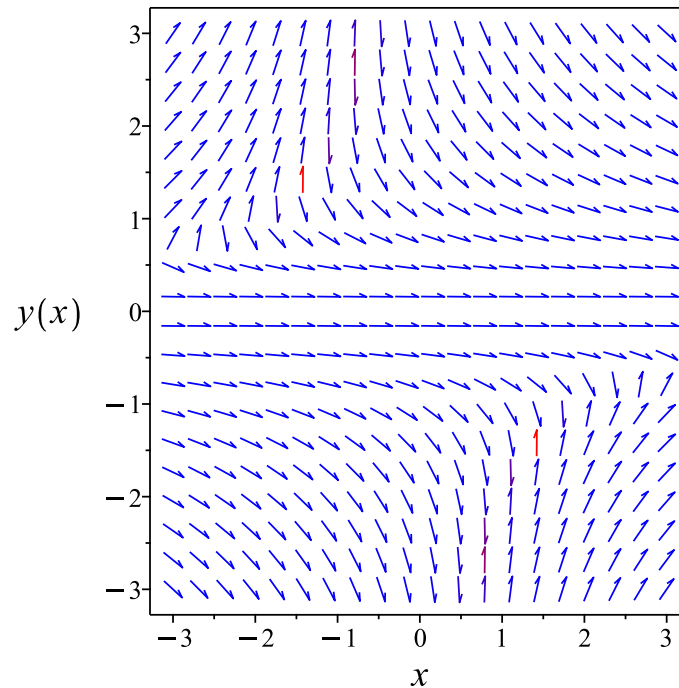


Figure 27: Slope field plot

## Verification of solutions

$$y = e^{-\text{LambertW}\left(\frac{x e^{\frac{c_1}{2}}}{2}\right) + \frac{c_1}{2}}$$

Verified OK.

### 2.1.3 Maple step by step solution

Let's solve

$$\left(x + \frac{2}{y}\right) y' + y = 0$$

- Highest derivative means the order of the ODE is 1  
 $y'$
- Check if ODE is exact
  - ODE is exact if the lhs is the total derivative of a  $C^2$  function  
 $F'(x, y) = 0$
  - Compute derivative of lhs  
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$
  - Evaluate derivatives  
 $1 = 1$
  - Condition met, ODE is exact
- Exact ODE implies solution will be of this form  
 $\left[F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y)\right]$
- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$   
 $F(x, y) = \int y dx + f_1(y)$
- Evaluate integral  
 $F(x, y) = xy + f_1(y)$
- Take derivative of  $F(x, y)$  with respect to  $y$   
 $N(x, y) = \frac{\partial}{\partial y} F(x, y)$
- Compute derivative  
 $x + \frac{2}{y} = x + \frac{d}{dy} f_1(y)$



- Isolate for  $\frac{d}{dy}f_1(y)$   

$$\frac{d}{dy}f_1(y) = \frac{2}{y}$$
- Solve for  $f_1(y)$   

$$f_1(y) = 2 \ln(y)$$
- Substitute  $f_1(y)$  into equation for  $F(x, y)$   

$$F(x, y) = xy + 2 \ln(y)$$
- Substitute  $F(x, y)$  into the solution of the ODE  

$$xy + 2 \ln(y) = c_1$$
- Solve for  $y$   

$$y = e^{-\text{LambertW}\left(\frac{x e^{\frac{c_1}{2}}}{2}\right) + \frac{c_1}{2}}$$

### Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
<- 1st order linear successful
<- inverse linear successful`

```

### ✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 17

```
dsolve((x+2/y(x))*diff(y(x),x)+y(x)=0,y(x), singsol=all)
```

$$y(x) = \frac{2 \text{LambertW}\left(\frac{x e^{\frac{c_1}{2}}}{2}\right)}{x}$$

✓ Solution by Mathematica

Time used: 10.621 (sec). Leaf size: 58

```
DSolve[(x+2/y[x])*y'[x]+y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{2W\left(-\frac{1}{2}\sqrt{e^{c_1}x^2}\right)}{x}$$

$$y(x) \rightarrow \frac{2W\left(\frac{1}{2}\sqrt{e^{c_1}x^2}\right)}{x}$$

$$y(x) \rightarrow 0$$

## 2.2 problem 2

2.2.1 Solving as exact ode . . . . . 122

Internal problem ID [3090]

Internal file name [OUTPUT/2582\_Sunday\_June\_05\_2022\_03\_20\_43\_AM\_35899133/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 2.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**exactWithIntegrationFactor**"

Maple gives the following as the ode type

[`y=\_G(x,y)']

$$\sin(x) \tan(y) + \cos(x) \sec(y)^2 y' = -1$$

### 2.2.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (\cos(x) \sec(y)^2) dy &= (-\sin(x) \tan(y) - 1) dx \\ (\sin(x) \tan(y) + 1) dx + (\cos(x) \sec(y)^2) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= \sin(x) \tan(y) + 1 \\ N(x, y) &= \cos(x) \sec(y)^2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (\sin(x) \tan(y) + 1) \\ &= \sin(x) \sec(y)^2 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (\cos(x) \sec(y)^2) \\ &= -\sin(x) \sec(y)^2 \end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \sec(x) \cos(y)^2 \left( (\sin(x) (1 + \tan(y)^2)) - (-\sin(x) \sec(y)^2) \right) \\ &= 2 \tan(x) \end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int 2 \tan(x) dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-2 \ln(\cos(x))} \\ &= \sec(x)^2\end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \sec(x)^2 (\sin(x) \tan(y) + 1) \\ &= (\sin(x) \tan(y) + 1) \sec(x)^2\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \sec(x)^2 (\cos(x) \sec(y)^2) \\ &= \sec(x) \sec(y)^2\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ ((\sin(x) \tan(y) + 1) \sec(x)^2) + (\sec(x) \sec(y)^2) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int (\sin(x) \tan(y) + 1) \sec(x)^2 dx \\ \phi &= \sec(x) \tan(y) + \tan(x) + f(y)\end{aligned} \tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\begin{aligned}\frac{\partial \phi}{\partial y} &= \sec(x) (1 + \tan(y)^2) + f'(y) \\ &= \sec(x) \sec(y)^2 + f'(y)\end{aligned}\tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \sec(x) \sec(y)^2$ . Therefore equation (4) becomes

$$\sec(x) \sec(y)^2 = \sec(x) \sec(y)^2 + f'(y)\tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = \sec(x) \tan(y) + \tan(x) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \sec(x) \tan(y) + \tan(x)$$

### Summary

The solution(s) found are the following

$$\sec(x) \tan(y) + \tan(x) = c_1\tag{1}$$

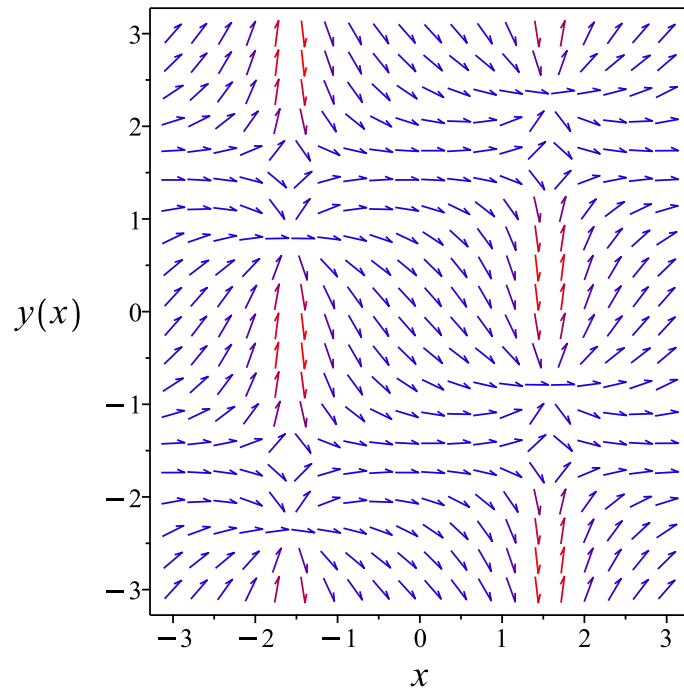


Figure 28: Slope field plot

Verification of solutions

$$\sec(x) \tan(y) + \tan(x) = c_1$$

Verified OK.


## Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
Looking for potential symmetries
trying inverse_Riccati
trying an equivalence to an Abel ODE
differential order: 1; trying a linearization to 2nd order
--- trying a change of variables {x -> y(x), y(x) -> x}
differential order: 1; trying a linearization to 2nd order
trying 1st order ODE linearizable_by_differentiation
--- Trying Lie symmetry methods, 1st order ---
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = 4
`, `-> Computing symmetries using: way = 5
trying symmetry patterns for 1st order ODEs
-> trying a symmetry pattern of the form [F(x)*G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)*G(y)]
-> trying symmetry patterns of the forms [F(x),G(y)] and [G(y),F(x)]
`, `-> Computing symmetries using: way = HINT
-> Calling odsolve with the ODE`, diff(y(x), x)+2*y(x)/sin(2*x), y(x)`
  Methods for first order ODEs:
    --- Trying classification methods ---
    trying a quadrature
    trying 1st order linear
    <- 1st order linear successful
  -> Calling odsolve with the ODE`, diff(y(x), x)+y(x)*sin(x)/cos(x), y(x)`
    Methods for first order ODEs:
      --- Trying classification methods ---
      trying a quadrature
      trying 1st order linear
      <- 1st order linear successful
  `, `-> Computing symmetries using: way = HINT
-> trying a symmetry pattern of the form [F(x),G(x)]
-> trying a symmetry pattern of the form [F(y),G(y)]
-> trying a symmetry pattern of the form [F(x)+G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)+G(y)]
-> trying a symmetry pattern of the form [F(x), G(x)+H(x)]
```

\*\*\* Sublevel

\*\*\* Sublev



 Solution by Maple

```
dsolve((sin(x)*tan(y(x))+1)+(cos(x)*sec(y(x))^2)*diff(y(x),x)=0,y(x), singsol=all)
```

No solution found

 Solution by Mathematica

Time used: 2.318 (sec). Leaf size: 54

```
DSolve[(Sin[x]*Tan[y[x]]+1)+(Cos[x]*Sec[y[x]]^2)*y'[x]==0,y[x],x,IncludeSingularSolutions ->
```

$$y(x) \rightarrow -\arctan(\sin(x) + c_1 \cos(x))$$

$$y(x) \rightarrow -\frac{1}{2}\pi\sqrt{\cos^2(x)}\sec(x)$$

$$y(x) \rightarrow \frac{1}{2}\pi\sqrt{\cos^2(x)}\sec(x)$$

## 2.3 problem 3

2.3.1 Solving as differentialType ode . . . . .	129
2.3.2 Solving as exact ode . . . . .	131
2.3.3 Maple step by step solution . . . . .	134

Internal problem ID [3091]

Internal file name [OUTPUT/2583\_Sunday\_June\_05\_2022\_03\_20\_59\_AM\_88936391/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 3.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exact", "differentialType"

Maple gives the following as the ode type

`[_exact, _rational]`

$$y + (x + y^3) y' = x^3$$

### 2.3.1 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-y + x^3}{x + y^3} \tag{1}$$

Which becomes

$$(y^3) dy = (-x) dy + (x^3 - y) dx \tag{2}$$

But the RHS is complete differential because

$$(-x) dy + (x^3 - y) dx = d\left(\frac{1}{4}x^4 - xy\right)$$

Hence (2) becomes

$$(y^3) dy = d\left(\frac{1}{4}x^4 - xy\right)$$

Integrating both sides gives gives the solution as

$$\frac{y^4}{4} = \frac{x^4}{4} - yx + c_1$$

### Summary

The solution(s) found are the following

$$\frac{y^4}{4} = \frac{x^4}{4} - yx + c_1 \quad (1)$$

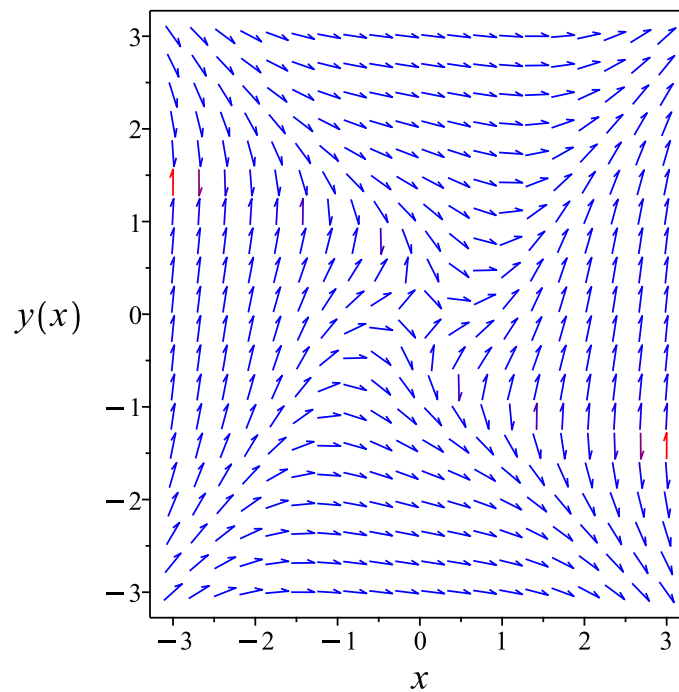


Figure 29: Slope field plot

### Verification of solutions

$$\frac{y^4}{4} = \frac{x^4}{4} - yx + c_1$$

Verified OK.

### 2.3.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (y^3 + x) dy &= (x^3 - y) dx \\ (-x^3 + y) dx + (y^3 + x) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x^3 + y \\ N(x, y) &= y^3 + x \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x^3 + y) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(y^3 + x) \\ &= 1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x^3 + y dx \\ \phi &= -\frac{1}{4}x^4 + xy + f(y)\end{aligned} \tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = x + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = y^3 + x$ . Therefore equation (4) becomes

$$y^3 + x = x + f'(y) \tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = y^3$$

Integrating the above w.r.t  $y$  gives

$$\int f'(y) dy = \int (y^3) dy$$

$$f(y) = \frac{y^4}{4} + c_1$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\frac{1}{4}x^4 + xy + \frac{1}{4}y^4 + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\frac{1}{4}x^4 + xy + \frac{1}{4}y^4$$

### Summary

The solution(s) found are the following

$$\frac{y^4}{4} - \frac{x^4}{4} + yx = c_1 \tag{1}$$

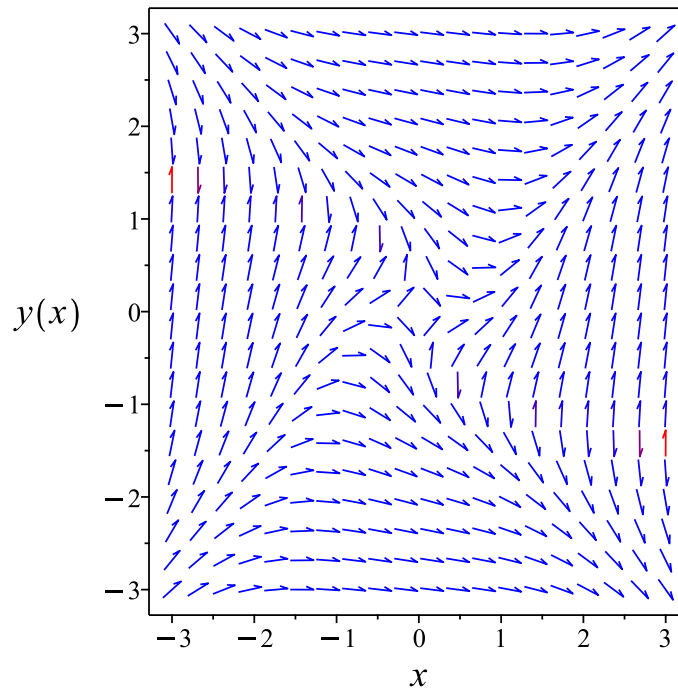


Figure 30: Slope field plot

### Verification of solutions

$$\frac{y^4}{4} - \frac{x^4}{4} + yx = c_1$$

Verified OK.

### 2.3.3 Maple step by step solution

Let's solve

$$y + (x + y^3) y' = x^3$$

- Highest derivative means the order of the ODE is 1  
 $y'$
- Check if ODE is exact
  - ODE is exact if the lhs is the total derivative of a  $C^2$  function  
 $F'(x, y) = 0$
  - Compute derivative of lhs  
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$

- Evaluate derivatives  
 $1 = 1$
- Condition met, ODE is exact
- Exact ODE implies solution will be of this form  

$$\left[ F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$
- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$   

$$F(x, y) = \int (-x^3 + y) dx + f_1(y)$$
- Evaluate integral  

$$F(x, y) = -\frac{x^4}{4} + xy + f_1(y)$$
- Take derivative of  $F(x, y)$  with respect to  $y$   

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$
- Compute derivative  

$$y^3 + x = x + \frac{d}{dy} f_1(y)$$
- Isolate for  $\frac{d}{dy} f_1(y)$   

$$\frac{d}{dy} f_1(y) = y^3$$
- Solve for  $f_1(y)$   

$$f_1(y) = \frac{y^4}{4}$$
- Substitute  $f_1(y)$  into equation for  $F(x, y)$   

$$F(x, y) = -\frac{1}{4}x^4 + xy + \frac{1}{4}y^4$$
- Substitute  $F(x, y)$  into the solution of the ODE  

$$-\frac{1}{4}x^4 + xy + \frac{1}{4}y^4 = c_1$$
- Solve for  $y$   

$$y = \text{RootOf}(\_Z^4 - x^4 + 4\_Zx - 4c_1)$$



## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 20

```
dsolve((y(x)-x^3)+(x+y(x)^3)*diff(y(x),x)=0,y(x), singsol=all)
```

$$-\frac{x^4}{4} + xy(x) + \frac{y(x)^4}{4} + c_1 = 0$$

✓ Solution by Mathematica

Time used: 60.173 (sec). Leaf size: 1210

`DSolve[(y[x]-x^3)+(x+y[x]^3)*y'[x]==0,y[x],x,IncludeSingularSolutions -> True]`

$$y(x) \rightarrow \sqrt{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}} - \frac{\sqrt[3]{3}(x^4+4c_1)}{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}}} + \sqrt{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}}}$$

$$y(x) \rightarrow \sqrt{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}} - \frac{6\sqrt{2}x}{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}}} - \sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}}}$$

$$y(x) \rightarrow \sqrt{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}} - \frac{\sqrt[3]{3}(x^4+4c_1)}{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}}} - \sqrt{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}}}$$

$$y(x) \rightarrow \sqrt{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}} - \frac{\sqrt[3]{3}(x^4+4c_1)}{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}}} + \sqrt{\sqrt[3]{9x^2 + \sqrt{3}\sqrt{27x^4 + (x^4 + 4c_1)^3}}}$$

## 2.4 problem 4

Internal problem ID [3092]

Internal file name [OUTPUT/2584\_Sunday\_June\_05\_2022\_03\_21\_02\_AM\_23832912/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 4.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**unknown**"

Maple gives the following as the ode type

```
[_rational, [_Abel, `2nd type`, `class B`]]
```

Unable to solve or complete the solution.

$$2y^2 - (4 - 2y + 4yx)y' = 4x - 5$$

Unable to determine ODE type.

## Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
trying Abel
Looking for potential symmetries
Looking for potential symmetries
Looking for potential symmetries
trying inverse_Riccati
trying an equivalence to an Abel ODE
differential order: 1; trying a linearization to 2nd order
--- trying a change of variables {x -> y(x), y(x) -> x}
differential order: 1; trying a linearization to 2nd order
trying 1st order ODE linearizable_by_differentiation
--- Trying Lie symmetry methods, 1st order ---
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = 4
`, `-> Computing symmetries using: way = 2
trying symmetry patterns for 1st order ODEs
-> trying a symmetry pattern of the form [F(x)*G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)*G(y)]
-> trying symmetry patterns of the forms [F(x),G(y)] and [G(y),F(x)]
-> trying a symmetry pattern of the form [F(x),G(x)]
-> trying a symmetry pattern of the form [F(y),G(y)]
-> trying a symmetry pattern of the form [F(x)+G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)+G(y)]
-> trying a symmetry pattern of the form [F(x),G(x)*y+H(x)]
-> trying a symmetry pattern of conformal type`
```

**X** Solution by Maple

```
dsolve((2*y(x)^2-4*x+5)=(4-2*y(x)+4*x*y(x))*diff(y(x),x),y(x), singsol=all)
```

No solution found

**X** Solution by Mathematica

Time used: 0.0 (sec). Leaf size: 0

```
DSolve[(2*y[x]^2-4*x+5)==(4-2*y[x]+4*x*y[x])*y'[x],y[x],x,IncludeSingularSolutions -> True]
```

Not solved

## 2.5 problem 5

2.5.1	Solving as separable ode . . . . .	141
2.5.2	Solving as linear ode . . . . .	143
2.5.3	Solving as homogeneousTypeD2 ode . . . . .	144
2.5.4	Solving as differentialType ode . . . . .	146
2.5.5	Solving as first order ode lie symmetry lookup ode . . . . .	147
2.5.6	Solving as exact ode . . . . .	151
2.5.7	Maple step by step solution . . . . .	155

Internal problem ID [3093]

Internal file name [OUTPUT/2585\_Sunday\_June\_05\_2022\_03\_21\_08\_AM\_57565474/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 5.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exact", "linear", "separable", "differentialType", "homogeneousTypeD2", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_separable]

$$y + y \cos(yx) + (x + x \cos(yx)) y' = 0$$

### 2.5.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\frac{y}{x} \end{aligned}$$

Where  $f(x) = -\frac{1}{x}$  and  $g(y) = y$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{y} dy &= -\frac{1}{x} dx \\ \int \frac{1}{y} dy &= \int -\frac{1}{x} dx \\ \ln(y) &= -\ln(x) + c_1 \\ y &= e^{-\ln(x)+c_1} \\ &= \frac{c_1}{x}\end{aligned}$$

### Summary

The solution(s) found are the following

$$y = \frac{c_1}{x} \tag{1}$$

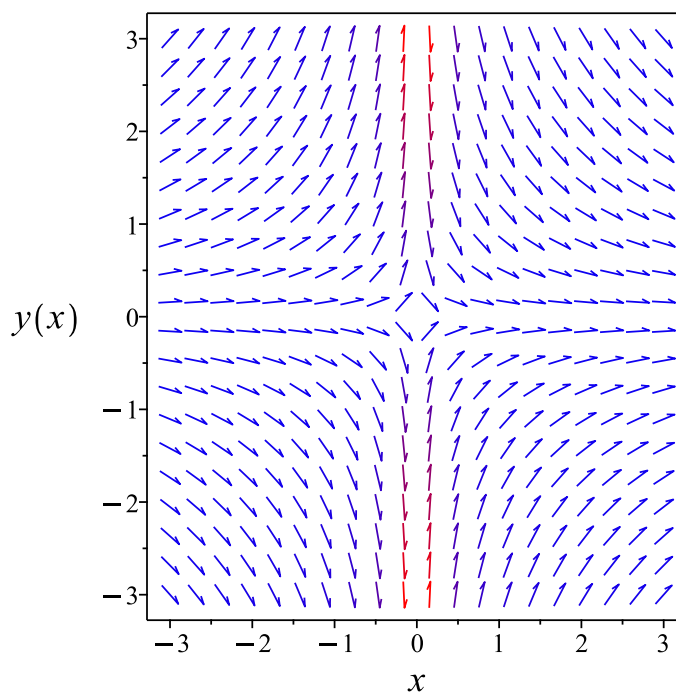


Figure 31: Slope field plot

### Verification of solutions

$$y = \frac{c_1}{x}$$

Verified OK.

### 2.5.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{1}{x}$$
$$q(x) = 0$$

Hence the ode is

$$y' + \frac{y}{x} = 0$$

The integrating factor  $\mu$  is

$$\mu = e^{\int \frac{1}{x} dx}$$
$$= x$$

The ode becomes

$$\frac{d}{dx} \mu y = 0$$
$$\frac{d}{dx} (xy) = 0$$

Integrating gives

$$xy = c_1$$

Dividing both sides by the integrating factor  $\mu = x$  results in

$$y = \frac{c_1}{x}$$

#### Summary

The solution(s) found are the following

$$y = \frac{c_1}{x} \tag{1}$$



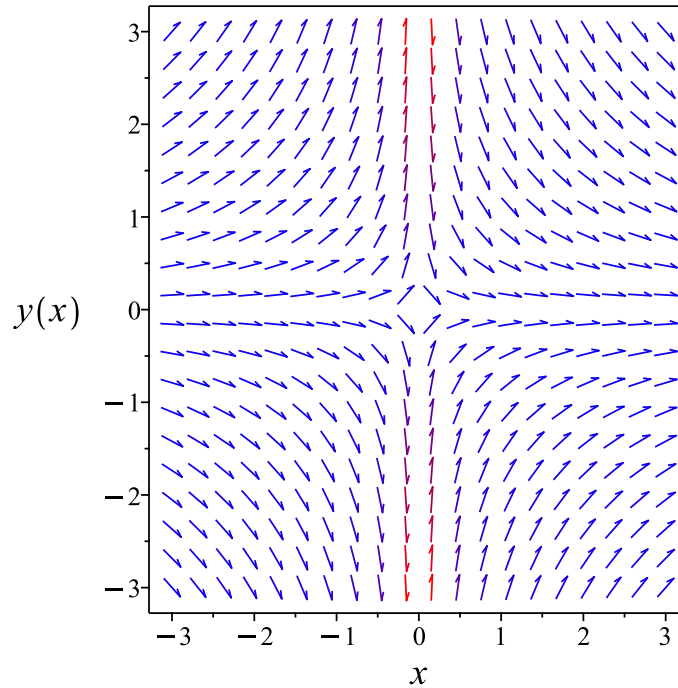


Figure 32: Slope field plot

Verification of solutions

$$y = \frac{c_1}{x}$$

Verified OK.

### 2.5.3 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$u(x)x + u(x)x \cos(u(x)x^2) + (x + x \cos(u(x)x^2))(u'(x)x + u(x)) = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{2u}{x} \end{aligned}$$

Where  $f(x) = -\frac{2}{x}$  and  $g(u) = u$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{u} du &= -\frac{2}{x} dx \\ \int \frac{1}{u} du &= \int -\frac{2}{x} dx \\ \ln(u) &= -2 \ln(x) + c_2 \\ u &= e^{-2 \ln(x) + c_2} \\ &= \frac{c_2}{x^2}\end{aligned}$$

Therefore the solution  $y$  is

$$\begin{aligned}y &= ux \\ &= \frac{c_2}{x}\end{aligned}$$

### Summary

The solution(s) found are the following

$$y = \frac{c_2}{x} \tag{1}$$

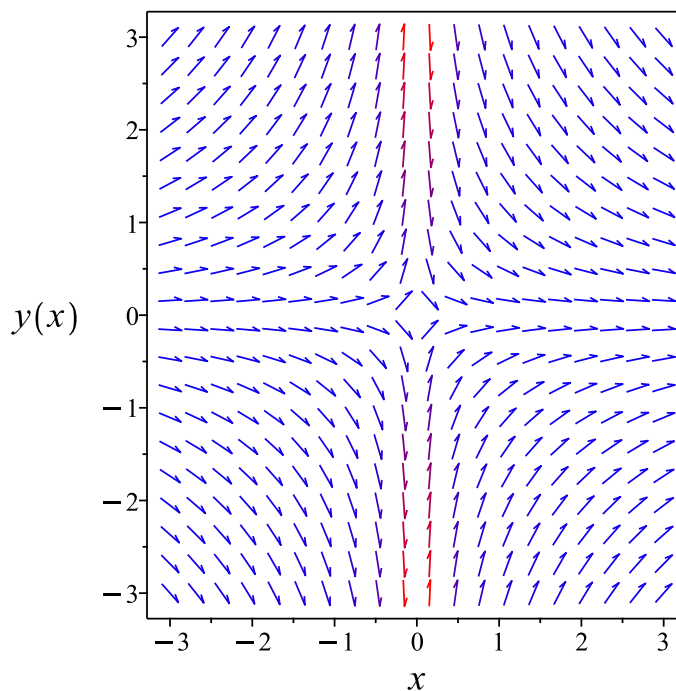


Figure 33: Slope field plot

### Verification of solutions

$$y = \frac{c_2}{x}$$

Verified OK.

### **2.5.4 Solving as differentialType ode**

Writing the ode as

$$y' = \frac{-y - y \cos(yx)}{x + x \cos(yx)} \quad (1)$$

Which becomes

$$0 = (-x) dy + (-y) dx \quad (2)$$

But the RHS is complete differential because

$$(-x) dy + (-y) dx = d(-xy)$$

Hence (2) becomes

$$0 = d(-xy)$$

Integrating both sides gives gives these solutions

$$y = \frac{c_1}{x} + c_1$$

### Summary

The solution(s) found are the following

$$y = \frac{c_1}{x} + c_1 \quad (1)$$

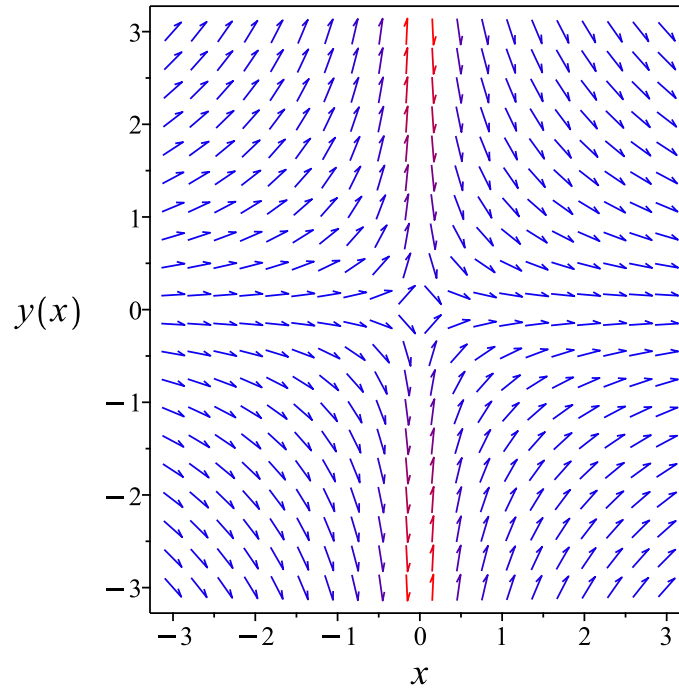


Figure 34: Slope field plot

Verification of solutions

$$y = \frac{c_1}{x} + c_2$$

Verified OK.

### 2.5.5 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{y}{x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 13: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{x}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{x}} dy \end{aligned}$$

Which results in

$$S = xy$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y}{x}$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 0$$

$$S_x = y$$

$$S_y = x$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$yx = c_1$$

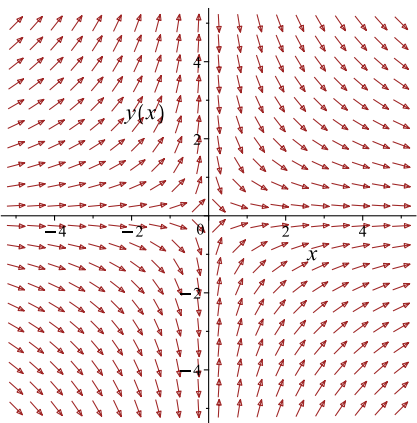
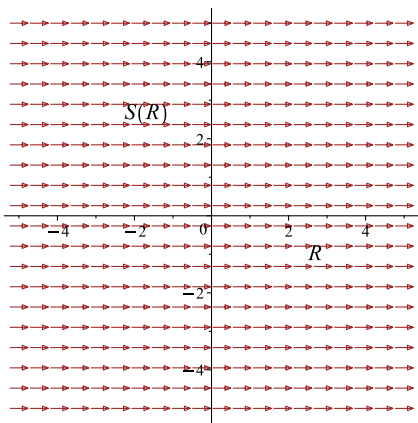
Which simplifies to

$$y = \frac{c_1}{x}$$

Which gives

$$y = \frac{c_1}{x}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{y}{x}$ 	$R = x$ $S = xy$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$y = \frac{c_1}{x} \tag{1}$$

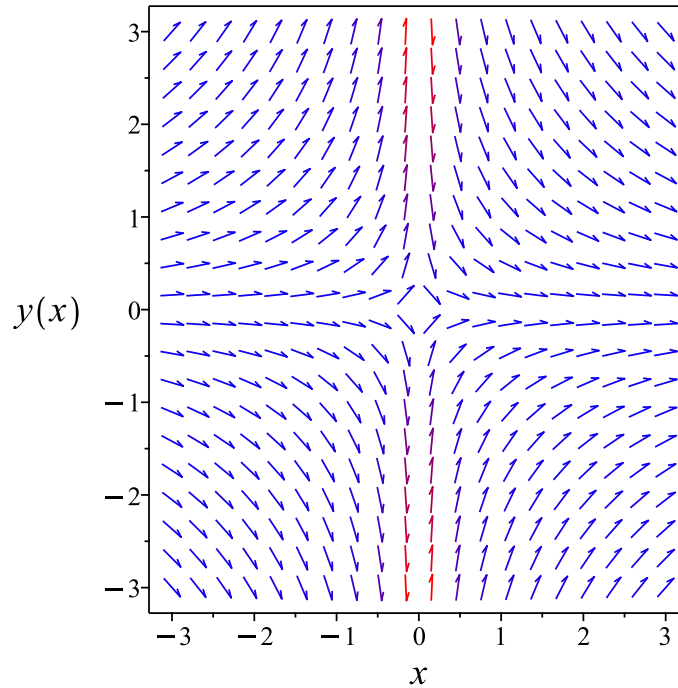


Figure 35: Slope field plot

Verification of solutions

$$y = \frac{c_1}{x}$$

Verified OK.

### 2.5.6 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$



Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(-\frac{1}{y}\right) dy &= \left(\frac{1}{x}\right) dx \\ \left(-\frac{1}{x}\right) dx + \left(-\frac{1}{y}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{1}{x} \\ N(x, y) &= -\frac{1}{y}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{1}{x}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left( -\frac{1}{y} \right) \\ &= 0\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{1}{x} dx \\ \phi &= -\ln(x) + f(y)\end{aligned} \tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = -\frac{1}{y}$ . Therefore equation (4) becomes

$$-\frac{1}{y} = 0 + f'(y) \tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = -\frac{1}{y}$$

Integrating the above w.r.t  $y$  gives

$$\begin{aligned}\int f'(y) dy &= \int \left( -\frac{1}{y} \right) dy \\ f(y) &= -\ln(y) + c_1\end{aligned}$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\ln(x) - \ln(y) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\ln(x) - \ln(y)$$

The solution becomes

$$y = \frac{e^{-c_1}}{x}$$

### Summary

The solution(s) found are the following

$$y = \frac{e^{-c_1}}{x} \tag{1}$$

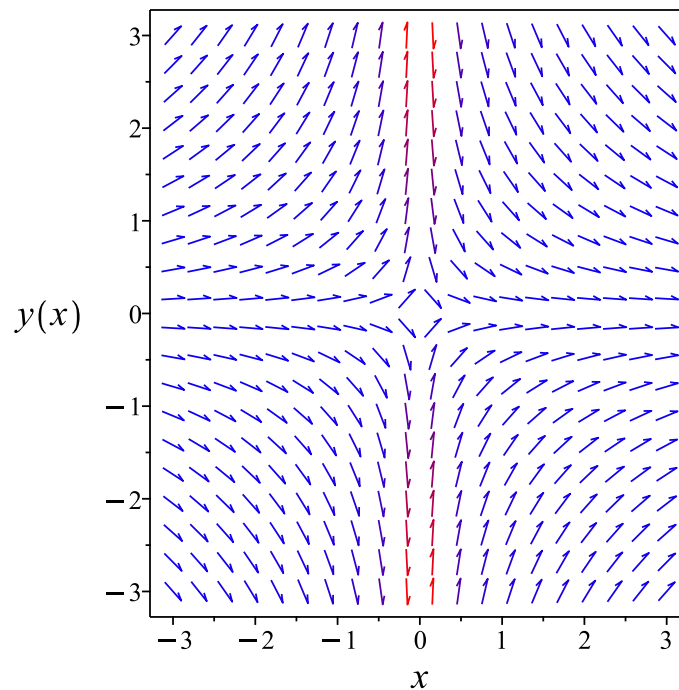


Figure 36: Slope field plot

## Verification of solutions

$$y = \frac{e^{-c_1}}{x}$$

Verified OK.

### 2.5.7 Maple step by step solution

Let's solve

$$y + y \cos(yx) + (x + x \cos(yx)) y' = 0$$

- Highest derivative means the order of the ODE is 1

$y'$

- Integrate both sides with respect to  $x$

$$\int (y + y \cos(yx) + (x + x \cos(yx)) y') dx = \int 0 dx + c_1$$

- Evaluate integral

$$yx + \sin(yx) = c_1$$

## Maple trace

```
`Classification methods on request
Methods to be used are: [exact]
-----
* Tackling ODE using method: exact
--- Trying classification methods ---
trying exact
<- exact successful`
```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 17

```
dsolve((y(x)+y(x)*cos(x*y(x)))+(x+x*cos(x*y(x)))*diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = \frac{\pi}{x}$$
$$y(x) = \frac{c_1}{x}$$

✓ Solution by Mathematica

Time used: 0.033 (sec). Leaf size: 49

```
DSolve[(y[x]+y[x]*Cos[x*y[x]])+(x+x*Cos[x*y[x]])*y'[x]==0,y[x],x,IncludeSingularSolutions ->
```

$$y(x) \rightarrow -\frac{\pi}{x}$$

$$y(x) \rightarrow \frac{\pi}{x}$$

$$y(x) \rightarrow \frac{C_1}{x}$$

$$y(x) \rightarrow -\frac{\pi}{x}$$

$$y(x) \rightarrow \frac{\pi}{x}$$

## 2.6 problem 6

2.6.1	Solving as separable ode . . . . .	157
2.6.2	Solving as first order ode lie symmetry lookup ode . . . . .	159
2.6.3	Solving as exact ode . . . . .	163
2.6.4	Maple step by step solution . . . . .	167

Internal problem ID [3094]

Internal file name [OUTPUT/2586\_Sunday\_June\_05\_2022\_03\_21\_10\_AM\_67179578/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 6.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exact", "separable", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_separable]

$$\cos(x) \cos(y)^2 + 2 \sin(x) \sin(y) \cos(y) y' = 0$$

### 2.6.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= -\frac{\cos(x) \cot(y)}{2 \sin(x)} \end{aligned}$$

Where  $f(x) = -\frac{\cos(x)}{2 \sin(x)}$  and  $g(y) = \cot(y)$ . Integrating both sides gives

$$\begin{aligned} \frac{1}{\cot(y)} dy &= -\frac{\cos(x)}{2 \sin(x)} dx \\ \int \frac{1}{\cot(y)} dy &= \int -\frac{\cos(x)}{2 \sin(x)} dx \\ -\ln(\cos(y)) &= -\frac{\ln(\sin(x))}{2} + c_1 \end{aligned}$$

Raising both side to exponential gives

$$\frac{1}{\cos(y)} = e^{-\frac{\ln(\sin(x))}{2} + c_1}$$

Which simplifies to

$$\sec(y) = \frac{c_2}{\sqrt{\sin(x)}}$$

Which simplifies to

$$y = \operatorname{arcsec}\left(\frac{c_2 e^{c_1}}{\sqrt{\sin(x)}}\right)$$

### Summary

The solution(s) found are the following

$$y = \operatorname{arcsec}\left(\frac{c_2 e^{c_1}}{\sqrt{\sin(x)}}\right) \quad (1)$$

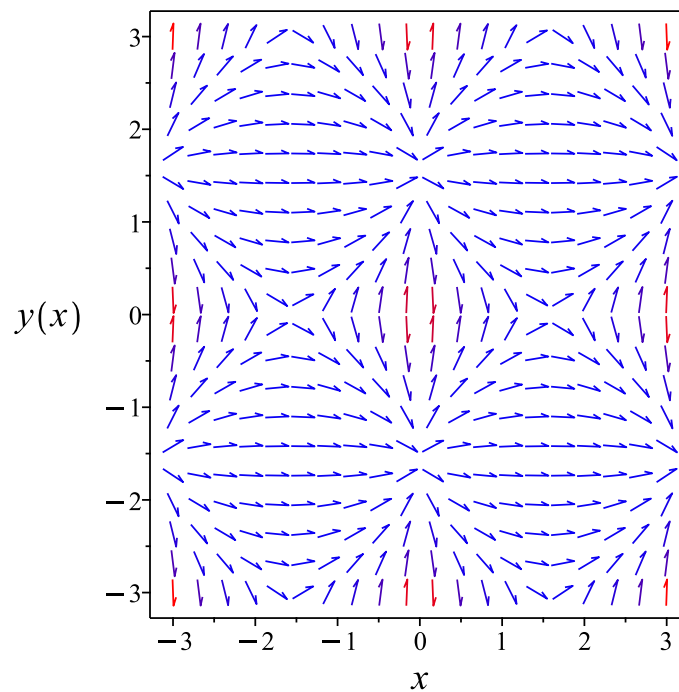


Figure 37: Slope field plot

Verification of solutions

$$y = \operatorname{arcsec} \left( \frac{c_2 e^{c_1}}{\sqrt{\sin(x)}} \right)$$

Verified OK.

### 2.6.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{\cos(x) \cos(y)}{2 \sin(x) \sin(y)}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$



Table 16: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= -\frac{2 \sin(x)}{\cos(x)} \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\eta = 0$  then in this special case

$$R = y$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{-\frac{2 \sin(x)}{\cos(x)}} dx \end{aligned}$$

Which results in

$$S = -\frac{\ln(\sin(x))}{2}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{\cos(x) \cos(y)}{2 \sin(x) \sin(y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 0 \\ R_y &= 1 \\ S_x &= -\frac{\cot(x)}{2} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \tan(y) \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \tan(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -\ln(\cos(R)) + c_1 \quad (4)$$

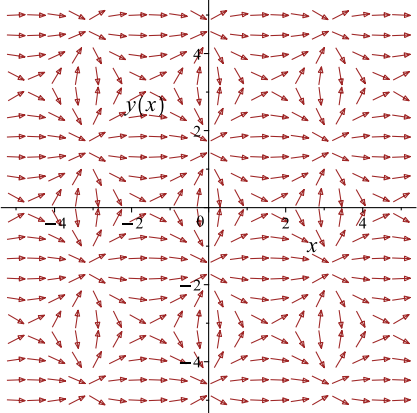
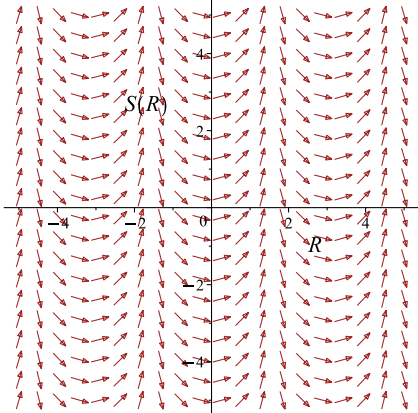
To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$-\frac{\ln(\sin(x))}{2} = -\ln(\cos(y)) + c_1$$

Which simplifies to

$$-\frac{\ln(\sin(x))}{2} = -\ln(\cos(y)) + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{\cos(x)\cos(y)}{2\sin(x)\sin(y)}$ 	$R = y$ $S = -\frac{\ln(\sin(x))}{2}$	$\frac{dS}{dR} = \tan(R)$ 

### Summary

The solution(s) found are the following

$$-\frac{\ln(\sin(x))}{2} = -\ln(\cos(y)) + c_1 \quad (1)$$

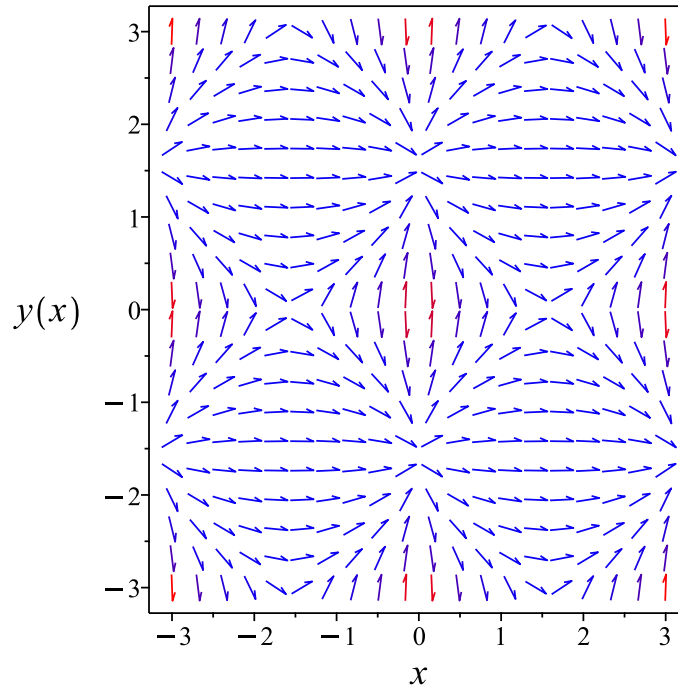


Figure 38: Slope field plot

Verification of solutions

$$-\frac{\ln(\sin(x))}{2} = -\ln(\cos(y)) + c_1$$

Verified OK.

### 2.6.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}\left(-\frac{2 \sin (y)}{\cos (y)}\right) dy &= \left(\frac{\cos (x)}{\sin (x)}\right) dx \\ \left(-\frac{\cos (x)}{\sin (x)}\right) dx + \left(-\frac{2 \sin (y)}{\cos (y)}\right) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{\cos (x)}{\sin (x)} \\ N(x, y) &= -\frac{2 \sin (y)}{\cos (y)}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{\cos (x)}{\sin (x)}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left( -\frac{2 \sin(y)}{\cos(y)} \right) \\ &= 0\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{\cos(x)}{\sin(x)} dx \\ \phi &= -\ln(\sin(x)) + f(y)\end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = -\frac{2 \sin(y)}{\cos(y)}$ . Therefore equation (4) becomes

$$-\frac{2 \sin(y)}{\cos(y)} = 0 + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$\begin{aligned}f'(y) &= -\frac{2 \sin(y)}{\cos(y)} \\ &= -2 \tan(y)\end{aligned}$$

Integrating the above w.r.t  $y$  results in

$$\int f'(y) dy = \int (-2 \tan(y)) dy$$

$$f(y) = 2 \ln(\cos(y)) + c_1$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\ln(\sin(x)) + 2 \ln(\cos(y)) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\ln(\sin(x)) + 2 \ln(\cos(y))$$

### Summary

The solution(s) found are the following

$$-\ln(\sin(x)) + 2 \ln(\cos(y)) = c_1 \tag{1}$$

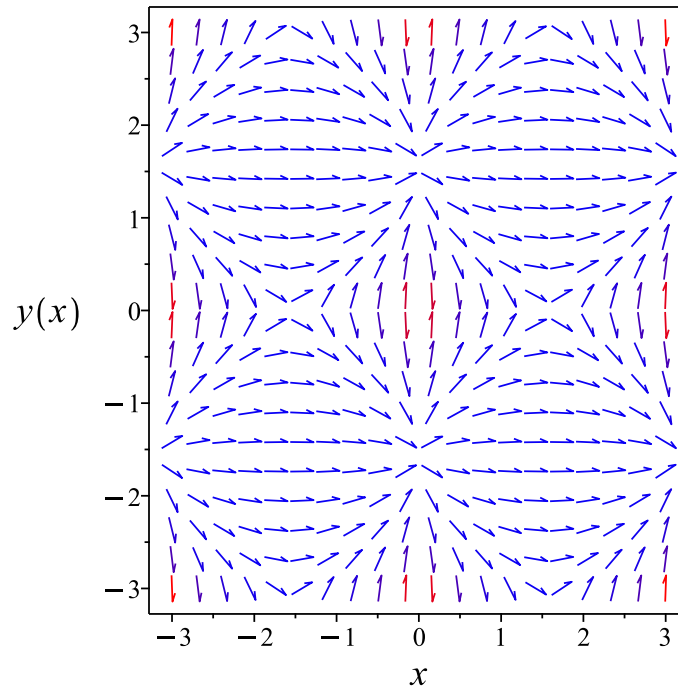


Figure 39: Slope field plot

### Verification of solutions

$$-\ln(\sin(x)) + 2\ln(\cos(y)) = c_1$$

Verified OK.

### 2.6.4 Maple step by step solution

Let's solve

$$\cos(x) \cos(y)^2 + 2 \sin(x) \sin(y) \cos(y) y' = 0$$

- Highest derivative means the order of the ODE is 1

$y'$

- Separate variables

$$\frac{y' \sin(y)}{\cos(y)} = -\frac{\cos(x)}{2 \sin(x)}$$

- Integrate both sides with respect to  $x$

$$\int \frac{y' \sin(y)}{\cos(y)} dx = \int -\frac{\cos(x)}{2 \sin(x)} dx + c_1$$

- Evaluate integral

$$-\ln(\cos(y)) = -\frac{\ln(\sin(x))}{2} + c_1$$

- Solve for  $y$

$$\left\{ y = \pi - \arccos\left(\frac{\sqrt{e^{2c_1} \sin(x)}}{e^{2c_1}}\right), y = \arccos\left(\frac{\sqrt{e^{2c_1} \sin(x)}}{e^{2c_1}}\right) \right\}$$

### Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
<- separable successful`
```



✓ Solution by Maple

Time used: 0.25 (sec). Leaf size: 31

```
dsolve(cos(x)*cos(y(x))^2+(2*sin(x)*sin(y(x))*cos(y(x)))*diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = \frac{\pi}{2}$$

$$y(x) = \arccos\left(\sqrt{c_1 \sin(x)}\right)$$

$$y(x) = \frac{\pi}{2} + \arcsin\left(\sqrt{c_1 \sin(x)}\right)$$

✓ Solution by Mathematica

Time used: 5.526 (sec). Leaf size: 73

```
DSolve[Cos[x]*Cos[y[x]]^2+(2*SIN[x]*Sin[y[x]]*Cos[y[x]])*y'[x]==0,y[x],x,IncludeSingularSolu
```

$$y(x) \rightarrow -\frac{\pi}{2}$$

$$y(x) \rightarrow \frac{\pi}{2}$$

$$y(x) \rightarrow -\arccos\left(-\frac{1}{4}c_1\sqrt{\sin(x)}\right)$$

$$y(x) \rightarrow \arccos\left(-\frac{1}{4}c_1\sqrt{\sin(x)}\right)$$

$$y(x) \rightarrow -\frac{\pi}{2}$$

$$y(x) \rightarrow \frac{\pi}{2}$$

## 2.7 problem 7

2.7.1 Solving as exact ode . . . . .	169
2.7.2 Maple step by step solution . . . . .	172

Internal problem ID [3095]

Internal file name [OUTPUT/2587\_Sunday\_June\_05\_2022\_03\_21\_13\_AM\_3875844/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 7.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"exact"**

Maple gives the following as the ode type

[\_exact]

$$(\sin(x) \sin(y) - x e^y) y' - e^y - \cos(x) \cos(y) = 0$$

### 2.7.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (\sin(x) \sin(y) - x e^y) dy &= (e^y + \cos(x) \cos(y)) dx \\ (-e^y - \cos(x) \cos(y)) dx + (\sin(x) \sin(y) - x e^y) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -e^y - \cos(x) \cos(y) \\ N(x, y) &= \sin(x) \sin(y) - x e^y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (-e^y - \cos(x) \cos(y)) \\ &= \cos(x) \sin(y) - e^y \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (\sin(x) \sin(y) - x e^y) \\ &= \cos(x) \sin(y) - e^y \end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -e^y - \cos(x) \cos(y) dx \\ \phi &= -\sin(x) \cos(y) - x e^y + f(y)\end{aligned}\quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = \sin(x) \sin(y) - x e^y + f'(y)\quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \sin(x) \sin(y) - x e^y$ . Therefore equation (4) becomes

$$\sin(x) \sin(y) - x e^y = \sin(x) \sin(y) - x e^y + f'(y)\quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\sin(x) \cos(y) - x e^y + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\sin(x) \cos(y) - x e^y$$

### Summary

The solution(s) found are the following

$$-\sin(x) \cos(y) - x e^y = c_1\quad (1)$$

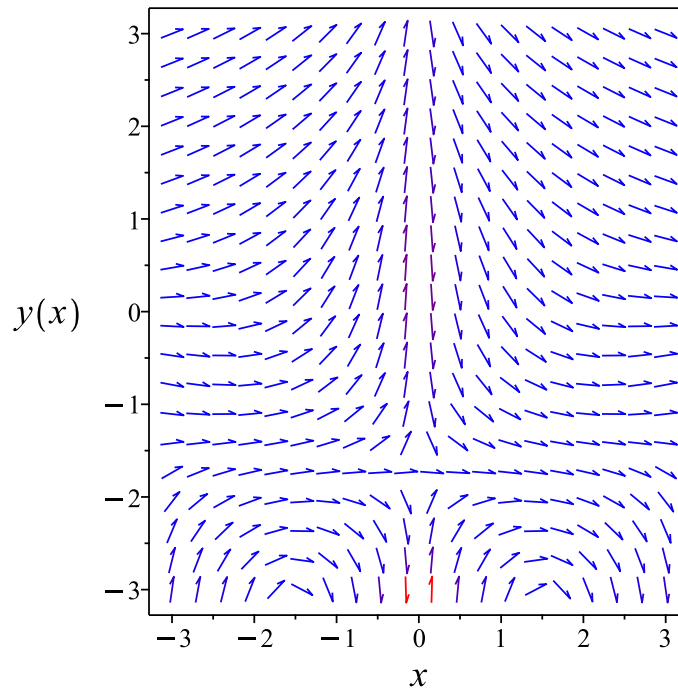


Figure 40: Slope field plot

### Verification of solutions

$$-\sin(x) \cos(y) - x e^y = c_1$$

Verified OK.

### 2.7.2 Maple step by step solution

Let's solve

$$(\sin(x) \sin(y) - x e^y) y' - e^y - \cos(x) \cos(y) = 0$$

- Highest derivative means the order of the ODE is 1  
 $y'$
- Check if ODE is exact
  - ODE is exact if the lhs is the total derivative of a  $C^2$  function  
 $F'(x, y) = 0$
  - Compute derivative of lhs  
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$

- Evaluate derivatives  
 $\cos(x) \sin(y) - e^y = \cos(x) \sin(y) - e^y$
- Condition met, ODE is exact
- Exact ODE implies solution will be of this form  

$$\left[ F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$
- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$   

$$F(x, y) = \int (-e^y - \cos(x) \cos(y)) dx + f_1(y)$$
- Evaluate integral  

$$F(x, y) = -\sin(x) \cos(y) - x e^y + f_1(y)$$
- Take derivative of  $F(x, y)$  with respect to  $y$   

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$
- Compute derivative  

$$\sin(x) \sin(y) - x e^y = \sin(x) \sin(y) - x e^y + \frac{d}{dy} f_1(y)$$
- Isolate for  $\frac{d}{dy} f_1(y)$   

$$\frac{d}{dy} f_1(y) = 0$$
- Solve for  $f_1(y)$   

$$f_1(y) = 0$$
- Substitute  $f_1(y)$  into equation for  $F(x, y)$   

$$F(x, y) = -\sin(x) \cos(y) - x e^y$$
- Substitute  $F(x, y)$  into the solution of the ODE  

$$-\sin(x) \cos(y) - x e^y = c_1$$
- Solve for  $y$   

$$y = \text{RootOf}\left(-Z - \ln\left(-\frac{\sin(x) \cos(-Z) + c_1}{x}\right)\right)$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

### ✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 16

```
dsolve((sin(x)*sin(y(x))-x*exp(y(x)))*diff(y(x),x)=exp(y(x))+cos(x)*cos(y(x))),y(x), singsol=
```

$$c_1 + \sin(x) \cos(y(x)) + e^{y(x)}x = 0$$

### ✓ Solution by Mathematica

Time used: 0.626 (sec). Leaf size: 21

```
DSolve[(Sin[x]*Sin[y[x]]-x*Exp[y[x]])*y'[x]==Exp[y[x]]+Cos[x]*Cos[y[x]],y[x],x,IncludeSingul
```

$$\text{Solve}[2(xe^{y(x)} + \sin(x) \cos(y(x))) = c_1, y(x)]$$

## 2.8 problem 8

2.8.1	Solving as separable ode . . . . .	175
2.8.2	Solving as linear ode . . . . .	177
2.8.3	Solving as homogeneousTypeD2 ode . . . . .	178
2.8.4	Solving as first order ode lie symmetry lookup ode . . . . .	179
2.8.5	Solving as exact ode . . . . .	183
2.8.6	Maple step by step solution . . . . .	187

Internal problem ID [3096]

Internal file name [OUTPUT/2588\_Sunday\_June\_05\_2022\_03\_21\_22\_AM\_91288277/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 8.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_separable]

$$-\frac{\sin\left(\frac{x}{y}\right)}{y} + \frac{x \sin\left(\frac{x}{y}\right) y'}{y^2} = 0$$

### 2.8.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{y}{x} \end{aligned}$$



Where  $f(x) = \frac{1}{x}$  and  $g(y) = y$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{y} dy &= \frac{1}{x} dx \\ \int \frac{1}{y} dy &= \int \frac{1}{x} dx \\ \ln(y) &= \ln(x) + c_1 \\ y &= e^{\ln(x)+c_1} \\ &= c_1 x\end{aligned}$$

### Summary

The solution(s) found are the following

$$y = c_1 x \tag{1}$$

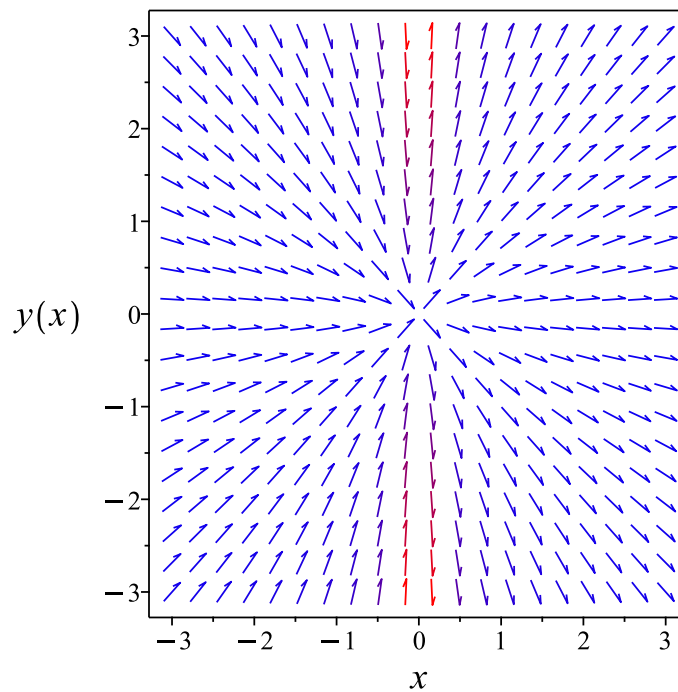


Figure 41: Slope field plot

### Verification of solutions

$$y = c_1 x$$

Verified OK.

### 2.8.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{1}{x}$$
$$q(x) = 0$$

Hence the ode is

$$y' - \frac{y}{x} = 0$$

The integrating factor  $\mu$  is

$$\mu = e^{\int -\frac{1}{x} dx}$$
$$= \frac{1}{x}$$

The ode becomes

$$\frac{d}{dx} \mu y = 0$$
$$\frac{d}{dx} \left( \frac{y}{x} \right) = 0$$

Integrating gives

$$\frac{y}{x} = c_1$$

Dividing both sides by the integrating factor  $\mu = \frac{1}{x}$  results in

$$y = c_1 x$$

#### Summary

The solution(s) found are the following

$$y = c_1 x \tag{1}$$

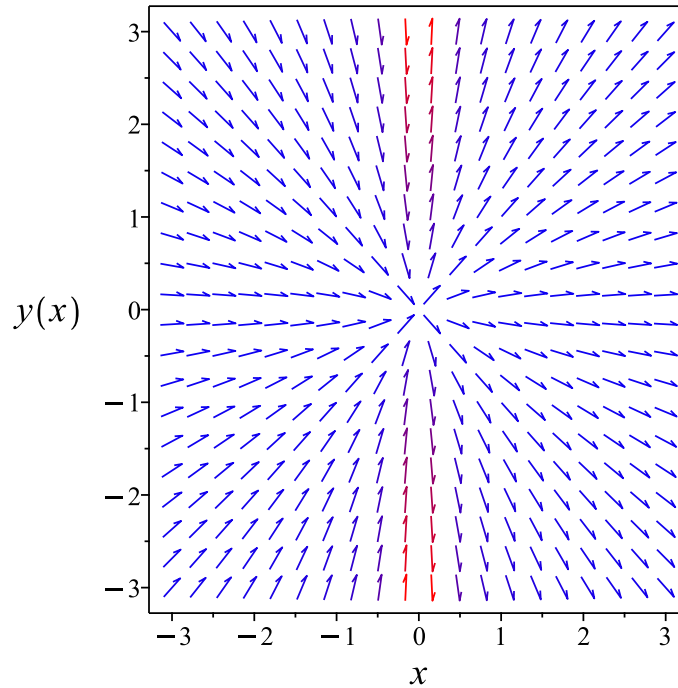


Figure 42: Slope field plot

Verification of solutions

$$y = c_1 x$$

Verified OK.

### 2.8.3 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$-\frac{\sin\left(\frac{1}{u(x)}\right)}{u(x)x} + \frac{\sin\left(\frac{1}{u(x)}\right)(u'(x)x + u(x))}{xu(x)^2} = 0$$

Integrating both sides gives

$$\begin{aligned} u(x) &= \int 0 \, dx \\ &= c_2 \end{aligned}$$

Therefore the solution  $y$  is

$$\begin{aligned} y &= xu \\ &= c_2 x \end{aligned}$$

### Summary

The solution(s) found are the following

$$y = c_2x \quad (1)$$

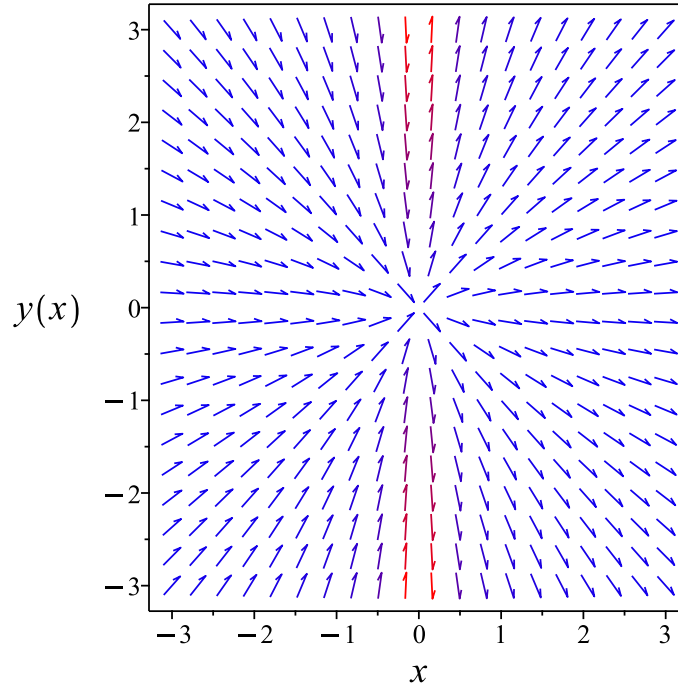


Figure 43: Slope field plot

### Verification of solutions

$$y = c_2x$$

Verified OK.

### **2.8.4 Solving as first order ode lie symmetry lookup ode**

Writing the ode as

$$y' = \frac{y}{x}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 20: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= x\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{x} dy \end{aligned}$$

Which results in

$$S = \frac{y}{x}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y}{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -\frac{y}{x^2} \\ S_y &= \frac{1}{x} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{y}{x} = c_1$$

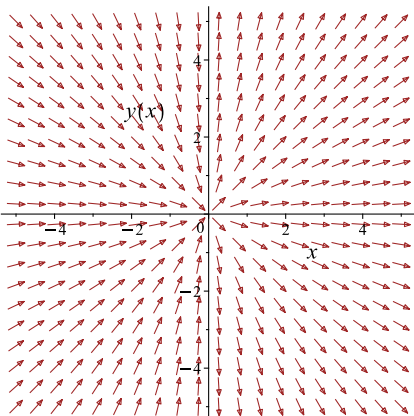
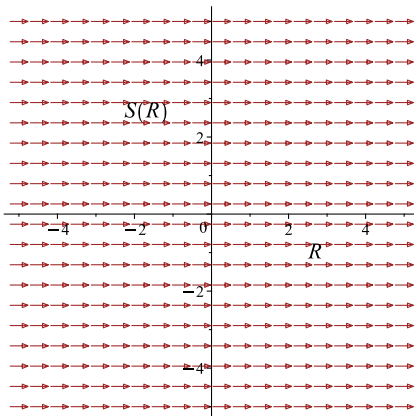
Which simplifies to

$$\frac{y}{x} = c_1$$

Which gives

$$y = c_1 x$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
<div style="text-align: center;"> <math>\frac{dy}{dx} = \frac{y}{x}</math>  </div>	$R = x$ $S = \frac{y}{x}$	<div style="text-align: center;"> <math>\frac{dS}{dR} = 0</math>  </div>

### Summary

The solution(s) found are the following

$$y = c_1 x \tag{1}$$

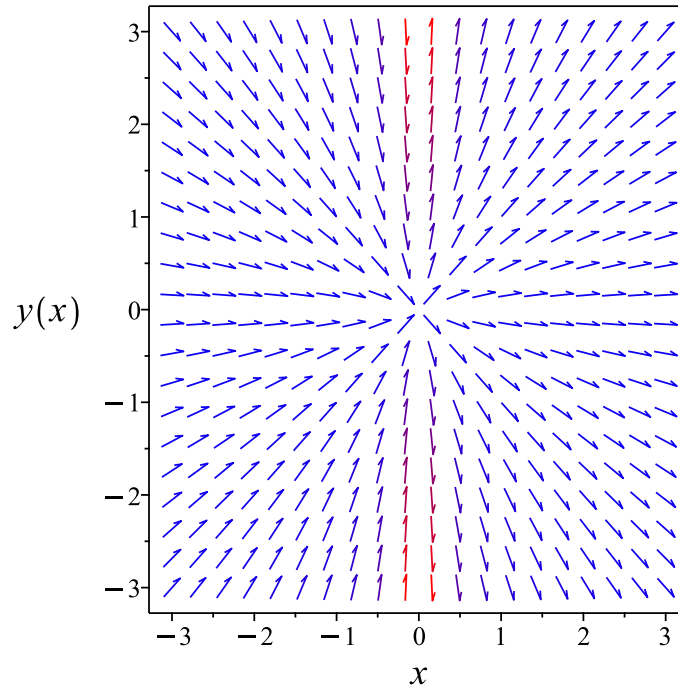


Figure 44: Slope field plot

Verification of solutions

$$y = c_1 x$$

Verified OK.

**2.8.5 Solving as exact ode**

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$



Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}\left(\frac{1}{y}\right) dy &= \left(\frac{1}{x}\right) dx \\ \left(-\frac{1}{x}\right) dx + \left(\frac{1}{y}\right) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{1}{x} \\ N(x, y) &= \frac{1}{y}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{1}{x}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left( \frac{1}{y} \right) \\ &= 0\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{1}{x} dx \\ \phi &= -\ln(x) + f(y)\end{aligned} \tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{1}{y}$ . Therefore equation (4) becomes

$$\frac{1}{y} = 0 + f'(y) \tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = \frac{1}{y}$$

Integrating the above w.r.t  $y$  gives

$$\begin{aligned}\int f'(y) dy &= \int \left( \frac{1}{y} \right) dy \\ f(y) &= \ln(y) + c_1\end{aligned}$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\ln(x) + \ln(y) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\ln(x) + \ln(y)$$

The solution becomes

$$y = e^{c_1} x$$

### Summary

The solution(s) found are the following

$$y = e^{c_1} x \tag{1}$$

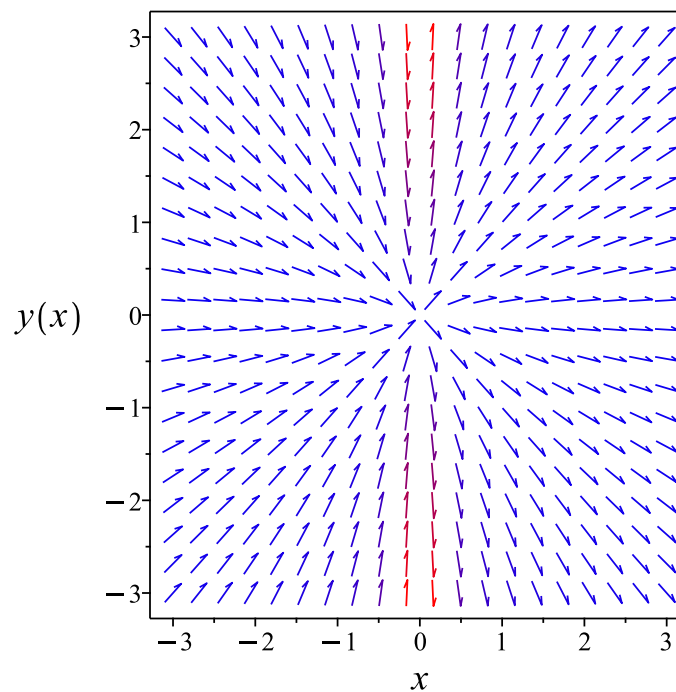


Figure 45: Slope field plot

### Verification of solutions

$$y = e^{c_1} x$$

Verified OK.

## 2.8.6 Maple step by step solution

Let's solve

$$-\frac{\sin\left(\frac{x}{y}\right)}{y} + \frac{x \sin\left(\frac{x}{y}\right)y'}{y^2} = 0$$

- Highest derivative means the order of the ODE is 1

$y'$

- Integrate both sides with respect to  $x$

$$\int \left( -\frac{\sin\left(\frac{x}{y}\right)}{y} + \frac{x \sin\left(\frac{x}{y}\right)y'}{y^2} \right) dx = \int 0 dx + c_1$$

- Evaluate integral

$$\cos\left(\frac{x}{y}\right) = c_1$$

- Solve for  $y$

$$y = \frac{x}{\arccos(c_1)}$$

Maple trace

```
`Classification methods on request
Methods to be used are: [exact]
-----
* Tackling ODE using method: exact
--- Trying classification methods ---
trying exact
<- exact successful`
```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 13

```
dsolve(-1/y(x)*sin(x/y(x))+x/y(x)^2*sin(x/y(x))*diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = \frac{x}{\pi - c_1}$$

✓ Solution by Mathematica

Time used: 0.033 (sec). Leaf size: 19

```
DSolve[-1/y[x]*Sin[x/y[x]]+x/y[x]^2*Sin[x/y[x]]*y'[x]==0,y[x],x,IncludeSingularSolutions ->
```

$$y(x) \rightarrow c_1 x$$

$$y(x) \rightarrow \text{ComplexInfinity}$$

$$y(x) \rightarrow \text{ComplexInfinity}$$

## 2.9 problem 9

2.9.1	Solving as separable ode . . . . .	189
2.9.2	Solving as linear ode . . . . .	191
2.9.3	Solving as homogeneousTypeD2 ode . . . . .	193
2.9.4	Solving as homogeneousTypeMapleC ode . . . . .	194
2.9.5	Solving as first order ode lie symmetry lookup ode . . . . .	197
2.9.6	Solving as exact ode . . . . .	201
2.9.7	Maple step by step solution . . . . .	205

Internal problem ID [3097]

Internal file name [OUTPUT/2589\_Sunday\_June\_05\_2022\_03\_21\_24\_AM\_19231703/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 9.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exact", "linear", "separable", "homogeneousTypeD2", "homogeneousTypeMapleC", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_separable]

$$y + (1 - x)y' = -1$$

### 2.9.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{y + 1}{x - 1}\end{aligned}$$

Where  $f(x) = \frac{1}{x-1}$  and  $g(y) = y + 1$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{y+1} dy &= \frac{1}{x-1} dx \\ \int \frac{1}{y+1} dy &= \int \frac{1}{x-1} dx \\ \ln(y+1) &= \ln(x-1) + c_1\end{aligned}$$

Raising both side to exponential gives

$$y + 1 = e^{\ln(x-1)+c_1}$$

Which simplifies to

$$y + 1 = c_2(x - 1)$$

### Summary

The solution(s) found are the following

$$y = c_2 e^{\ln(x-1)+c_1} - 1 \tag{1}$$

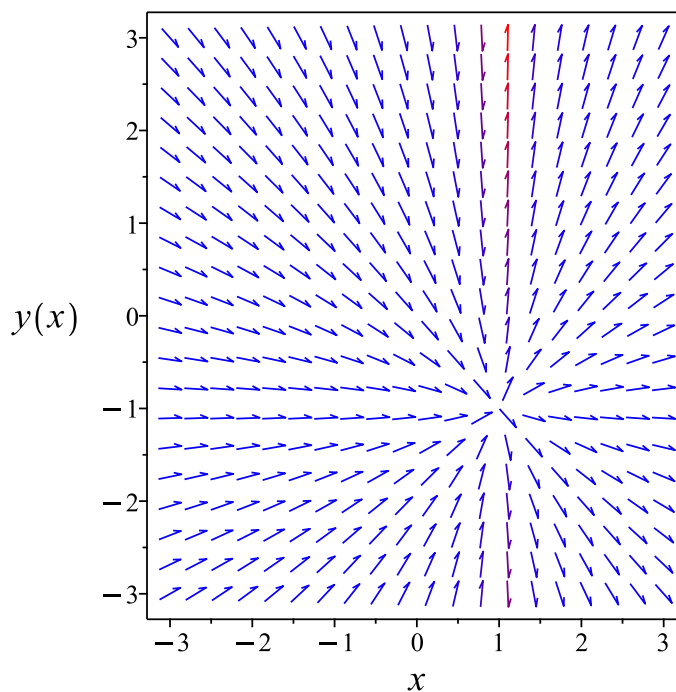


Figure 46: Slope field plot

### Verification of solutions

$$y = c_2 e^{\ln(x-1)+c_1} - 1$$

Verified OK.

### 2.9.2 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{1}{x-1}$$
$$q(x) = \frac{1}{x-1}$$

Hence the ode is

$$y' - \frac{y}{x-1} = \frac{1}{x-1}$$

The integrating factor  $\mu$  is

$$\mu = e^{\int -\frac{1}{x-1} dx}$$
$$= \frac{1}{x-1}$$

The ode becomes

$$\frac{d}{dx}(\mu y) = (\mu) \left( \frac{1}{x-1} \right)$$
$$\frac{d}{dx} \left( \frac{y}{x-1} \right) = \left( \frac{1}{x-1} \right) \left( \frac{1}{x-1} \right)$$
$$d \left( \frac{y}{x-1} \right) = \frac{1}{(x-1)^2} dx$$

Integrating gives

$$\frac{y}{x-1} = \int \frac{1}{(x-1)^2} dx$$
$$\frac{y}{x-1} = -\frac{1}{x-1} + c_1$$



Dividing both sides by the integrating factor  $\mu = \frac{1}{x-1}$  results in

$$y = -1 + c_1(x - 1)$$

### Summary

The solution(s) found are the following

$$y = -1 + c_1(x - 1) \tag{1}$$

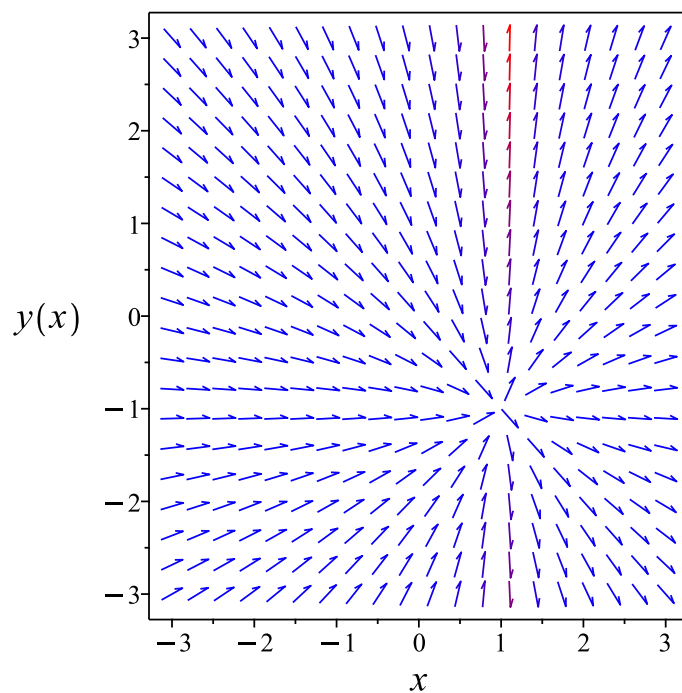


Figure 47: Slope field plot

### Verification of solutions

$$y = -1 + c_1(x - 1)$$

Verified OK.

### 2.9.3 Solving as homogeneous Type D2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$u(x)x + (1-x)(u'(x)x + u(x)) = -1$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{u+1}{x(x-1)}\end{aligned}$$

Where  $f(x) = \frac{1}{x(x-1)}$  and  $g(u) = u+1$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{u+1} du &= \frac{1}{x(x-1)} dx \\ \int \frac{1}{u+1} du &= \int \frac{1}{x(x-1)} dx \\ \ln(u+1) &= \ln(x-1) - \ln(x) + c_2\end{aligned}$$

Raising both side to exponential gives

$$u+1 = e^{\ln(x-1) - \ln(x) + c_2}$$

Which simplifies to

$$u+1 = c_3 e^{\ln(x-1) - \ln(x)}$$

Which simplifies to

$$u(x) = c_3 \left( e^{c_2} - \frac{e^{c_2}}{x} \right) - 1$$

Therefore the solution  $y$  is

$$\begin{aligned}y &= ux \\ &= x \left( c_3 \left( e^{c_2} - \frac{e^{c_2}}{x} \right) - 1 \right)\end{aligned}$$

#### Summary

The solution(s) found are the following

$$y = x \left( c_3 \left( e^{c_2} - \frac{e^{c_2}}{x} \right) - 1 \right) \quad (1)$$

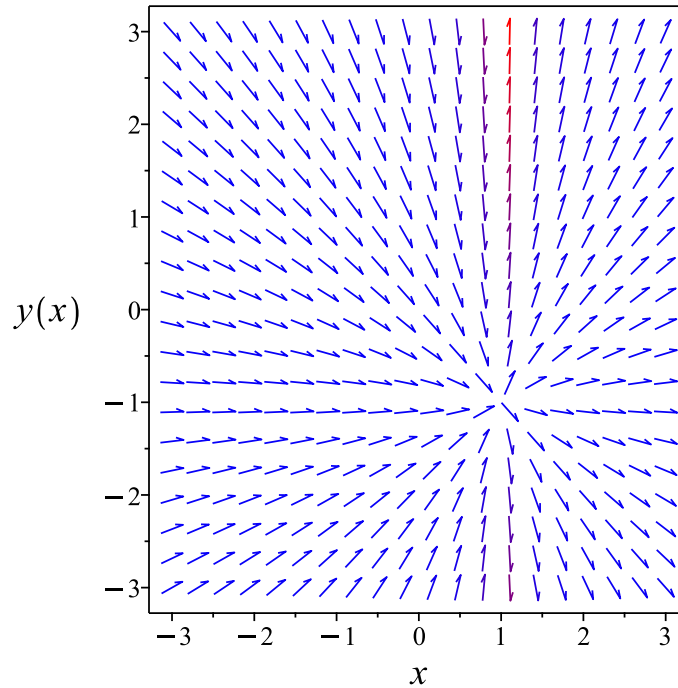


Figure 48: Slope field plot

Verification of solutions

$$y = x \left( c_3 \left( e^{c_2} - \frac{e^{c_2}}{x} \right) - 1 \right)$$

Verified OK.

**2.9.4 Solving as homogeneousTypeMapleC ode**

Let  $Y = y + y_0$  and  $X = x + x_0$  then the above is transformed to new ode in  $Y(X)$

$$\frac{d}{dX}Y(X) = \frac{Y(X) + y_0 + 1}{X + x_0 - 1}$$

Solving for possible values of  $x_0$  and  $y_0$  which makes the above ode a homogeneous ode results in

$$\begin{aligned} x_0 &= 1 \\ y_0 &= -1 \end{aligned}$$

Using these values now it is possible to easily solve for  $Y(X)$ . The above ode now becomes

$$\frac{d}{dX}Y(X) = \frac{Y(X)}{X}$$

In canonical form, the ODE is

$$\begin{aligned} Y' &= F(X, Y) \\ &= \frac{Y}{X} \end{aligned} \tag{1}$$

An ode of the form  $Y' = \frac{M(X, Y)}{N(X, Y)}$  is called homogeneous if the functions  $M(X, Y)$  and  $N(X, Y)$  are both homogeneous functions and of the same order. Recall that a function  $f(X, Y)$  is homogeneous of order  $n$  if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both  $M = Y$  and  $N = X$  are both homogeneous and of the same order  $n = 1$ . Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution  $u = \frac{Y}{X}$ , or  $Y = uX$ . Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation  $Y = uX$  to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dX}X + u &= u \\ \frac{du}{dX} &= 0 \end{aligned}$$

Or

$$\frac{d}{dX}u(X) = 0$$

Which is now solved as separable in  $u(X)$ . Which is now solved in  $u(X)$ . Integrating both sides gives

$$\begin{aligned} u(X) &= \int 0 \, dX \\ &= c_2 \end{aligned}$$

Now  $u$  in the above solution is replaced back by  $Y$  using  $u = \frac{Y}{X}$  which results in the solution

$$Y(X) = Xc_2$$

Using the solution for  $Y(X)$

$$Y(X) = Xc_2$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y - 1$$

$$X = x + 1$$

Then the solution in  $y$  becomes

$$y + 1 = c_2(x - 1)$$

### Summary

The solution(s) found are the following

$$y + 1 = c_2(x - 1) \tag{1}$$

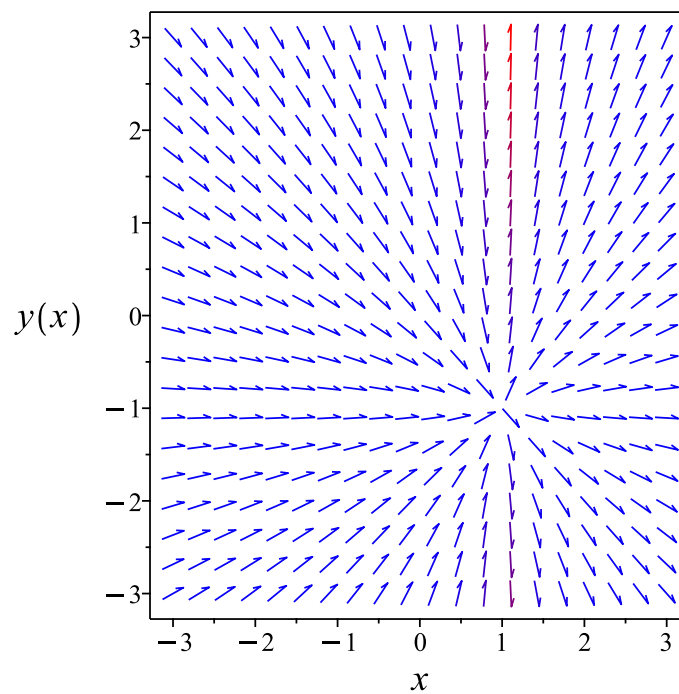


Figure 49: Slope field plot

### Verification of solutions

$$y + 1 = c_2(x - 1)$$

Verified OK.

### 2.9.5 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{y+1}{x-1}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 23: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= x - 1\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{x-1} dy\end{aligned}$$

Which results in

$$S = \frac{y}{x-1}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y+1}{x-1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= -\frac{y}{(x-1)^2} \\S_y &= \frac{1}{x-1}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{(x-1)^2} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{(R-1)^2}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -\frac{1}{R-1} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{y}{x-1} = -\frac{1}{x-1} + c_1$$

Which simplifies to

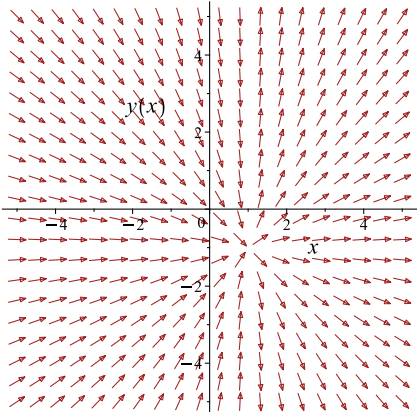
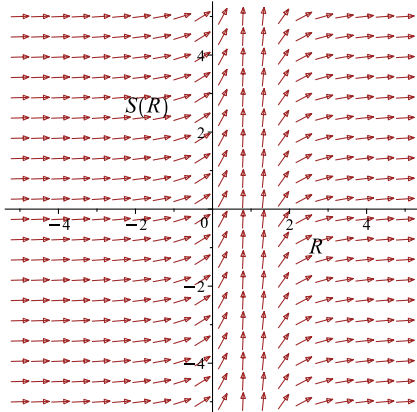
$$\frac{-c_1x + c_1 + y + 1}{x-1} = 0$$

Which gives

$$y = c_1x - c_1 - 1$$



The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{y+1}{x-1}$ 	$R = x$ $S = \frac{y}{x-1}$	$\frac{dS}{dR} = \frac{1}{(R-1)^2}$ 

Summary

The solution(s) found are the following

$$y = c_1 x - c_1 - 1 \tag{1}$$

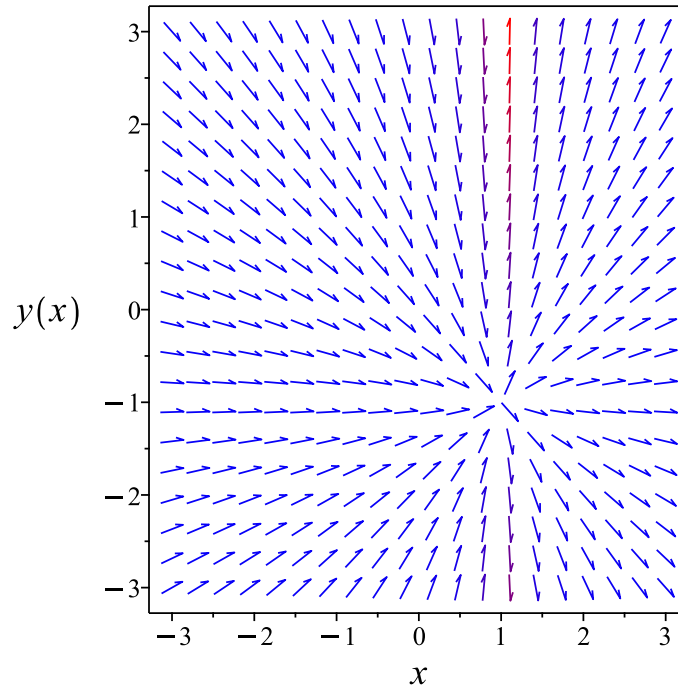


Figure 50: Slope field plot

Verification of solutions

$$y = c_1x - c_1 - 1$$

Verified OK.

### 2.9.6 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}\left(\frac{1}{y+1}\right) dy &= \left(\frac{1}{x-1}\right) dx \\ \left(-\frac{1}{x-1}\right) dx + \left(\frac{1}{y+1}\right) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\frac{1}{x-1} \\ N(x, y) &= \frac{1}{y+1}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{1}{x-1}\right) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left( \frac{1}{y+1} \right) \\ &= 0\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{1}{x-1} dx \\ \phi &= -\ln(x-1) + f(y)\end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{1}{y+1}$ . Therefore equation (4) becomes

$$\frac{1}{y+1} = 0 + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = \frac{1}{y+1}$$

Integrating the above w.r.t  $y$  gives

$$\begin{aligned}\int f'(y) dy &= \int \left( \frac{1}{y+1} \right) dy \\ f(y) &= \ln(y+1) + c_1\end{aligned}$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\ln(x-1) + \ln(y+1) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\ln(x-1) + \ln(y+1)$$

The solution becomes

$$y = e^{c_1}x - e^{c_1} - 1$$

### Summary

The solution(s) found are the following

$$y = e^{c_1}x - e^{c_1} - 1 \tag{1}$$

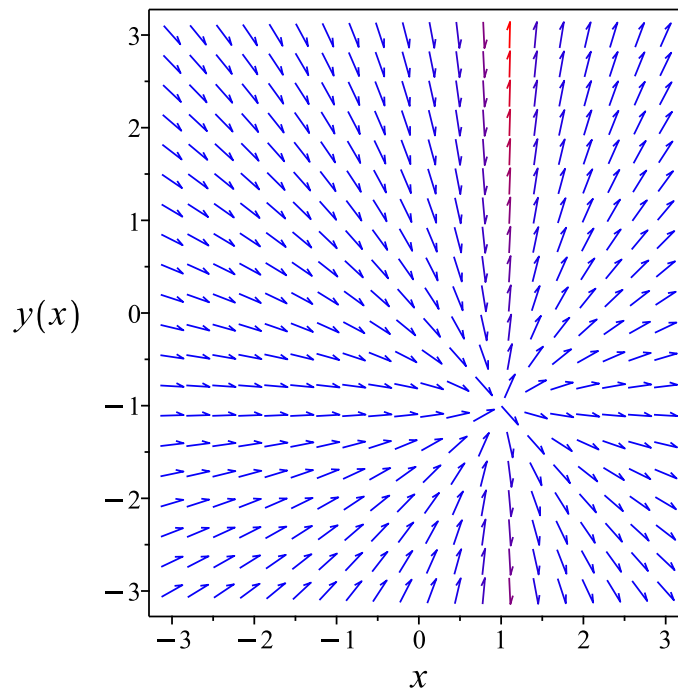


Figure 51: Slope field plot

### Verification of solutions

$$y = e^{c_1}x - e^{c_1} - 1$$

Verified OK.

## 2.9.7 Maple step by step solution

Let's solve

$$y + (1 - x)y' = -1$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Separate variables

$$\frac{y'}{-1-y} = \frac{1}{1-x}$$

- Integrate both sides with respect to  $x$

$$\int \frac{y'}{-1-y} dx = \int \frac{1}{1-x} dx + c_1$$

- Evaluate integral

$$-\ln(-1 - y) = -\ln(1 - x) + c_1$$

- Solve for  $y$

$$y = -\frac{e^{c_1-x+1}}{e^{c_1}}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 11

```
dsolve((1+y(x))+(1-x)*diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = -1 + c_1(x - 1)$$

✓ Solution by Mathematica

Time used: 0.026 (sec). Leaf size: 18

```
DSolve[(1+y[x])+(1-x)*y'[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -1 + c_1(x - 1)$$

$$y(x) \rightarrow -1$$

## 2.10 problem 10

2.10.1 Solving as exact ode . . . . .	207
2.10.2 Maple step by step solution . . . . .	210

Internal problem ID [3098]

Internal file name [OUTPUT/2590\_Sunday\_June\_05\_2022\_03\_21\_26\_AM\_11054784/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 10.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exact"

Maple gives the following as the ode type

```
[_exact , [_1st_order , ` _with_symmetry_[F(x),G(x)*y+H(x)] `]]
```

$$2xy^3 + \cos(x)y + (3y^2x^2 + \sin(x))y' = 0$$

### 2.10.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$



But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (3y^2x^2 + \sin(x)) dy &= (-2xy^3 - \cos(x)y) dx \\ (2xy^3 + \cos(x)y) dx + (3y^2x^2 + \sin(x)) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= 2xy^3 + \cos(x)y \\ N(x, y) &= 3y^2x^2 + \sin(x) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (2xy^3 + \cos(x)y) \\ &= 6xy^2 + \cos(x) \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (3y^2x^2 + \sin(x)) \\ &= 6xy^2 + \cos(x) \end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 2x y^3 + \cos(x) y dx \\ \phi &= y(y^2 x^2 + \sin(x)) + f(y)\end{aligned}\quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = 3y^2 x^2 + \sin(x) + f'(y)\quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = 3y^2 x^2 + \sin(x)$ . Therefore equation (4) becomes

$$3y^2 x^2 + \sin(x) = 3y^2 x^2 + \sin(x) + f'(y)\quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = y(y^2 x^2 + \sin(x)) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = y(y^2 x^2 + \sin(x))$$

### Summary

The solution(s) found are the following

$$y(y^2 x^2 + \sin(x)) = c_1\quad (1)$$

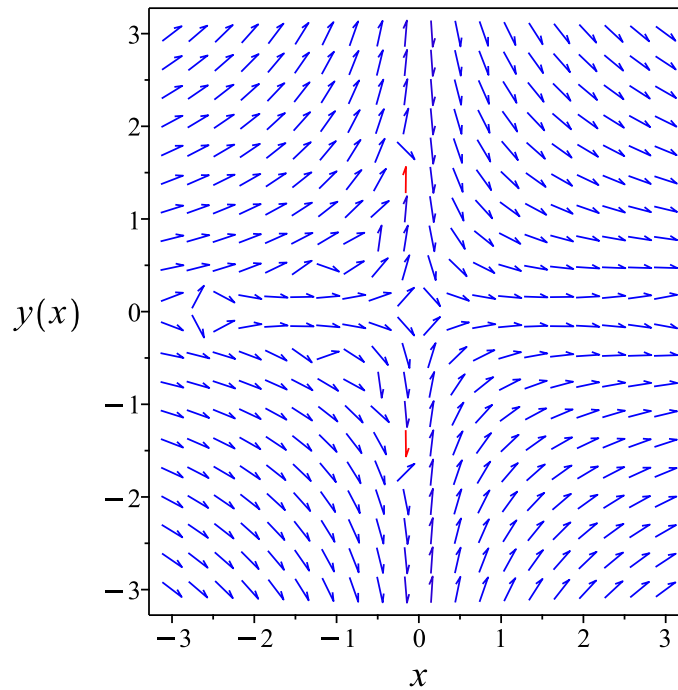


Figure 52: Slope field plot

Verification of solutions

$$y(y^2 x^2 + \sin(x)) = c_1$$

Verified OK.

### 2.10.2 Maple step by step solution

Let's solve

$$2xy^3 + \cos(x)y + (3y^2x^2 + \sin(x))y' = 0$$

- Highest derivative means the order of the ODE is 1  
 $y'$
- Check if ODE is exact
  - ODE is exact if the lhs is the total derivative of a  $C^2$  function  
 $F'(x, y) = 0$
  - Compute derivative of lhs  
 $F'(x, y) + \left(\frac{\partial}{\partial y}F(x, y)\right)y' = 0$

- Evaluate derivatives

$$6x y^2 + \cos(x) = 6x y^2 + \cos(x)$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$\left[ F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$

- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$

$$F(x, y) = \int (2x y^3 + \cos(x) y) dx + f_1(y)$$

- Evaluate integral

$$F(x, y) = \sin(x) y + y^3 x^2 + f_1(y)$$

- Take derivative of  $F(x, y)$  with respect to  $y$

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$

- Compute derivative

$$3y^2 x^2 + \sin(x) = \sin(x) + 3y^2 x^2 + \frac{d}{dy} f_1(y)$$

- Isolate for  $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = 0$$

- Solve for  $f_1(y)$

$$f_1(y) = 0$$

- Substitute  $f_1(y)$  into equation for  $F(x, y)$

$$F(x, y) = \sin(x) y + y^3 x^2$$

- Substitute  $F(x, y)$  into the solution of the ODE

$$\sin(x) y + y^3 x^2 = c_1$$

- Solve for  $y$

$$\left\{ y = \frac{\left(12\sqrt{3} \sqrt{27c_1^2 x^2 + 4 \sin(x)^3} + 108c_1 x\right)^{\frac{1}{3}}}{6x} - \frac{2 \sin(x)}{x \left(12\sqrt{3} \sqrt{27c_1^2 x^2 + 4 \sin(x)^3} + 108c_1 x\right)^{\frac{1}{3}}}, y = -\frac{\left(12\sqrt{3} \sqrt{27c_1^2 x^2 + 4 \sin(x)^3}\right)^{\frac{1}{3}}}{12x}$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 302

```
dsolve((2*x*y(x)^3+y(x)*cos(x))+(3*x^2*y(x)^2+sin(x))*diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = \frac{\left(12\sqrt{3}\sqrt{27c_1^2x^2 + 4\sin(x)^3} - 108c_1x\right)^{\frac{2}{3}} - 12\sin(x)}{6x\left(12\sqrt{3}\sqrt{27c_1^2x^2 + 4\sin(x)^3} - 108c_1x\right)^{\frac{1}{3}}}$$

$$y(x) = \frac{i\sqrt{3}\left(12\sqrt{3}\sqrt{27c_1^2x^2 + 4\sin(x)^3} - 108c_1x\right)^{\frac{2}{3}} + 12i\sqrt{3}\sin(x) + \left(12\sqrt{3}\sqrt{27c_1^2x^2 + 4\sin(x)^3} - 108c_1x\right)^{\frac{2}{3}}}{12x\left(12\sqrt{3}\sqrt{27c_1^2x^2 + 4\sin(x)^3} - 108c_1x\right)^{\frac{1}{3}}}$$

$$y(x) = \frac{i\sqrt{3}\left(12\sqrt{3}\sqrt{27c_1^2x^2 + 4\sin(x)^3} - 108c_1x\right)^{\frac{2}{3}} + 12i\sqrt{3}\sin(x) - \left(12\sqrt{3}\sqrt{27c_1^2x^2 + 4\sin(x)^3} - 108c_1x\right)^{\frac{2}{3}}}{12x\left(12\sqrt{3}\sqrt{27c_1^2x^2 + 4\sin(x)^3} - 108c_1x\right)^{\frac{1}{3}}}$$

✓ Solution by Mathematica

Time used: 32.512 (sec). Leaf size: 339

`DSolve[(2*x*y[x]^3+y[x]*Cos[x])+(3*x^2*y[x]^2+Sin[x])*y'[x]==0,y[x],x,IncludeSingularSolutio`

$$y(x) \rightarrow \frac{\sqrt[3]{9c_1x^4 + \sqrt{12x^6 \sin^3(x) + 81c_1^2x^8}}}{\sqrt[3]{23^{2/3}x^2}} - \frac{\sqrt[3]{\frac{2}{3}} \sin(x)}{\sqrt[3]{9c_1x^4 + \sqrt{12x^6 \sin^3(x) + 81c_1^2x^8}}}$$

$$y(x) \rightarrow \frac{(1 + i\sqrt{3}) \sin(x)}{2^{2/3} \sqrt[3]{27c_1x^4 + 3\sqrt{12x^6 \sin^3(x) + 81c_1^2x^8}}} - \frac{(1 - i\sqrt{3}) \sqrt[3]{27c_1x^4 + \sqrt{108x^6 \sin^3(x) + 729c_1^2x^8}}}{6\sqrt[3]{2}x^2}$$

$$y(x) \rightarrow \frac{(1 - i\sqrt{3}) \sin(x)}{2^{2/3} \sqrt[3]{27c_1x^4 + 3\sqrt{12x^6 \sin^3(x) + 81c_1^2x^8}}} - \frac{(1 + i\sqrt{3}) \sqrt[3]{27c_1x^4 + \sqrt{108x^6 \sin^3(x) + 729c_1^2x^8}}}{6\sqrt[3]{2}x^2}$$

## 2.11 problem 11

2.11.1 Solving as exact ode . . . . .	214
2.11.2 Solving as riccati ode . . . . .	218
2.11.3 Maple step by step solution . . . . .	220

Internal problem ID [3099]

Internal file name [OUTPUT/2591\_Sunday\_June\_05\_2022\_03\_21\_32\_AM\_30495618/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 8, page 41

**Problem number:** 11.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"exact", "riccati"**

Maple gives the following as the ode type

`[_exact, _rational, _Riccati]`

$$\boxed{-\frac{y}{1-y^2x^2} - \frac{xy'}{1-y^2x^2} = -1}$$

### 2.11.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}\left(-\frac{x}{-y^2x^2 + 1}\right) dy &= \left(-1 + \frac{y}{-y^2x^2 + 1}\right) dx \\ \left(1 - \frac{y}{-y^2x^2 + 1}\right) dx &+ \left(-\frac{x}{-y^2x^2 + 1}\right) dy = 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 1 - \frac{y}{-y^2x^2 + 1} \\ N(x, y) &= -\frac{x}{-y^2x^2 + 1}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(1 - \frac{y}{-y^2x^2 + 1}\right) \\ &= \frac{-y^2x^2 - 1}{(y^2x^2 - 1)^2}\end{aligned}$$



And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left( -\frac{x}{-y^2x^2 + 1} \right) \\ &= \frac{-y^2x^2 - 1}{(y^2x^2 - 1)^2}\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 1 - \frac{y}{-y^2x^2 + 1} dx \\ \phi &= x - \frac{\ln(xy + 1)}{2} + \frac{\ln(xy - 1)}{2} + f(y)\end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\begin{aligned}\frac{\partial \phi}{\partial y} &= -\frac{x}{2(xy + 1)} + \frac{x}{2xy - 2} + f'(y) \\ &= \frac{x}{y^2x^2 - 1} + f'(y)\end{aligned} \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = -\frac{x}{-y^2x^2 + 1}$ . Therefore equation (4) becomes

$$-\frac{x}{-y^2x^2 + 1} = \frac{x}{y^2x^2 - 1} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = x - \frac{\ln(xy + 1)}{2} + \frac{\ln(xy - 1)}{2} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = x - \frac{\ln(xy + 1)}{2} + \frac{\ln(xy - 1)}{2}$$

The solution becomes

$$y = -\frac{e^{-2x+2c_1} + 1}{x(e^{-2x+2c_1} - 1)}$$

### Summary

The solution(s) found are the following

$$y = -\frac{e^{-2x+2c_1} + 1}{x(e^{-2x+2c_1} - 1)} \quad (1)$$

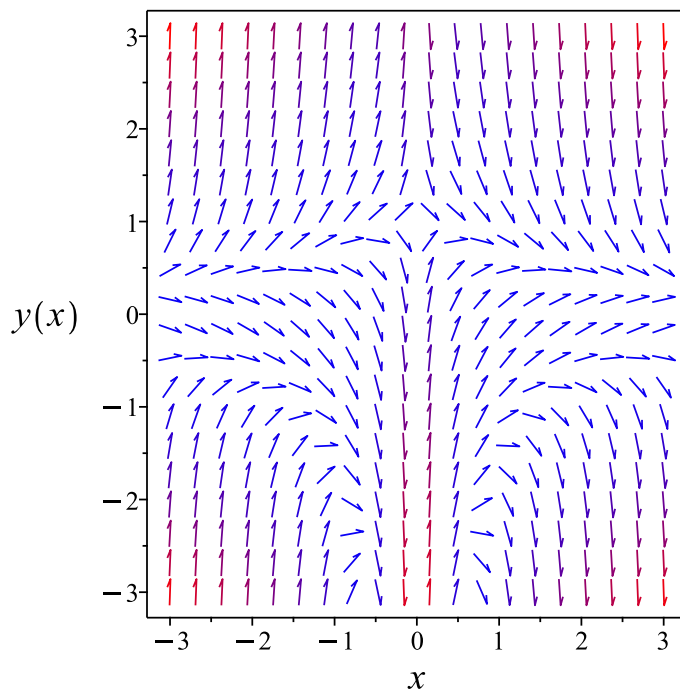


Figure 53: Slope field plot

### Verification of solutions

$$y = -\frac{e^{-2x+2c_1} + 1}{x(e^{-2x+2c_1} - 1)}$$

Verified OK.

### **2.11.2 Solving as riccati ode**

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= -\frac{y^2 x^2 + y - 1}{x}\end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = -x y^2 - \frac{y}{x} + \frac{1}{x}$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that  $f_0(x) = \frac{1}{x}$ ,  $f_1(x) = -\frac{1}{x}$  and  $f_2(x) = -x$ . Let

$$\begin{aligned}y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{-x u}\end{aligned}\tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for  $u(x)$  which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0\tag{2}$$

But

$$\begin{aligned}f_2' &= -1 \\ f_1 f_2 &= 1 \\ f_2^2 f_0 &= x\end{aligned}$$

Substituting the above terms back in equation (2) gives

$$-x u''(x) + x u(x) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 e^{-x} + c_2 e^x$$

The above shows that

$$u'(x) = -c_1 e^{-x} + c_2 e^x$$

Using the above in (1) gives the solution

$$y = \frac{-c_1 e^{-x} + c_2 e^x}{x (c_1 e^{-x} + c_2 e^x)}$$

Dividing both numerator and denominator by  $c_1$  gives, after renaming the constant  $\frac{c_2}{c_1} = c_3$  the following solution

$$y = \frac{e^{2x} - c_3}{x (e^{2x} + c_3)}$$

### Summary

The solution(s) found are the following

$$y = \frac{e^{2x} - c_3}{x (e^{2x} + c_3)} \tag{1}$$

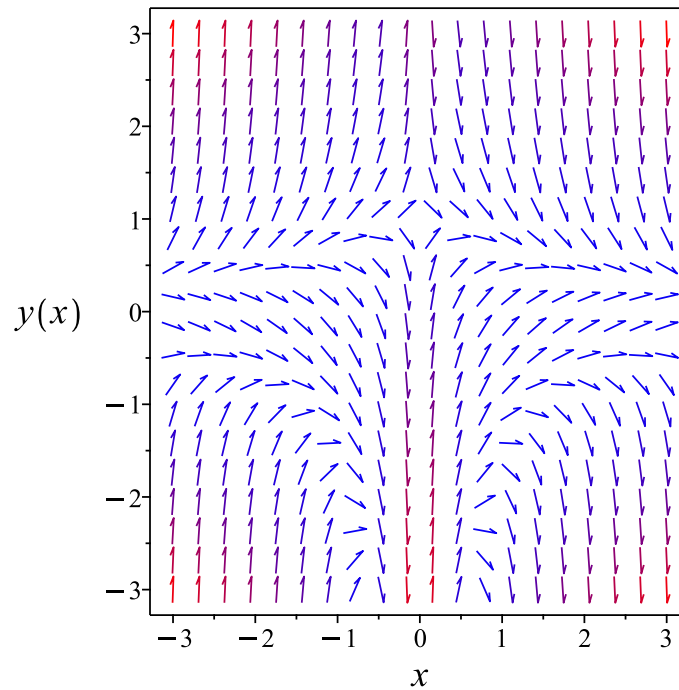


Figure 54: Slope field plot

Verification of solutions

$$y = \frac{e^{2x} - c_3}{x(e^{2x} + c_3)}$$

Verified OK.

### 2.11.3 Maple step by step solution

Let's solve

$$-\frac{y}{1-y^2x^2} - \frac{xy'}{1-y^2x^2} = -1$$

- Highest derivative means the order of the ODE is 1  
 $y'$
- Check if ODE is exact
  - ODE is exact if the lhs is the total derivative of a  $C^2$  function  
 $F'(x, y) = 0$
  - Compute derivative of lhs

$$F'(x, y) + \left( \frac{\partial}{\partial y} F(x, y) \right) y' = 0$$

- Evaluate derivatives

$$-\frac{1}{-y^2x^2+1} - \frac{2y^2x^2}{(-y^2x^2+1)^2} = -\frac{1}{-y^2x^2+1} - \frac{2y^2x^2}{(-y^2x^2+1)^2}$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$\left[ F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$

- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$

$$F(x, y) = \int \left( 1 - \frac{y}{-y^2x^2+1} \right) dx + f_1(y)$$

- Evaluate integral

$$F(x, y) = x - \frac{\ln(xy+1)}{2} + \frac{\ln(xy-1)}{2} + f_1(y)$$

- Take derivative of  $F(x, y)$  with respect to  $y$

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$

- Compute derivative

$$-\frac{x}{-y^2x^2+1} = -\frac{x}{2(xy+1)} + \frac{x}{2(xy-1)} + \frac{d}{dy} f_1(y)$$

- Isolate for  $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = -\frac{x}{-y^2x^2+1} + \frac{x}{2(xy+1)} - \frac{x}{2(xy-1)}$$

- Solve for  $f_1(y)$

$$f_1(y) = 0$$

- Substitute  $f_1(y)$  into equation for  $F(x, y)$

$$F(x, y) = x - \frac{\ln(xy+1)}{2} + \frac{\ln(xy-1)}{2}$$

- Substitute  $F(x, y)$  into the solution of the ODE

$$x - \frac{\ln(xy+1)}{2} + \frac{\ln(xy-1)}{2} = c_1$$

- Solve for  $y$

$$y = -\frac{e^{-2x+2c_1}+1}{x(e^{-2x+2c_1}-1)}$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

### ✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 24

```
dsolve(1=y(x)/(1-x^2*y(x)^2)+x/(1-x^2*y(x)^2)*diff(y(x),x),y(x), singsol=all)
```

$$y(x) = \frac{e^{2x} + c_1}{x(e^{2x} - c_1)}$$

### ✓ Solution by Mathematica

Time used: 0.153 (sec). Leaf size: 18

```
DSolve[1==y[x]/(1-x^2*y[x]^2)+x/(1-x^2*y[x]^2)*y'[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{\tanh(x + ic_1)}{x}$$

### **3 Chapter 2, section 10, page 47**

3.1	problem 2(a)	224
3.2	problem 2(b)	240
3.3	problem 2(c)	246
3.4	problem 4(a)	260
3.5	problem 4(b)	272
3.6	problem 4(c)	284
3.7	problem 4(d)	295
3.8	problem 4(e)	309



### 3.1 problem 2(a)

3.1.1	Solving as homogeneousTypeD2 ode . . . . .	224
3.1.2	Solving as first order ode lie symmetry calculated ode . . . . .	226
3.1.3	Solving as exact ode . . . . .	232

Internal problem ID [3100]

Internal file name [OUTPUT/2592\_Sunday\_June\_05\_2022\_03\_21\_35\_AM\_8164515/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 10, page 47

**Problem number:** 2(a).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"homogeneousTypeD2", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_calculated"**

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _dAlembert]
```

$$(3x^2 - y^2) y' - 2yx = 0$$

#### 3.1.1 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$(3x^2 - u(x)^2 x^2) (u'(x)x + u(x)) - 2u(x)x^2 = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{u^3 - u}{x(u^2 - 3)} \end{aligned}$$

Where  $f(x) = -\frac{1}{x}$  and  $g(u) = \frac{u^3 - u}{u^2 - 3}$ . Integrating both sides gives

$$\frac{1}{\frac{u^3 - u}{u^2 - 3}} du = -\frac{1}{x} dx$$

$$\int \frac{1}{\frac{u^3-u}{u^2-3}} du = \int -\frac{1}{x} dx$$

$$-\ln(u+1) - \ln(u-1) + 3\ln(u) = -\ln(x) + c_2$$

Raising both side to exponential gives

$$e^{-\ln(u+1)-\ln(u-1)+3\ln(u)} = e^{-\ln(x)+c_2}$$

Which simplifies to

$$\frac{u^3}{u^2-1} = \frac{c_3}{x}$$

The solution is

$$\frac{u(x)^3}{u(x)^2-1} = \frac{c_3}{x}$$

Replacing  $u(x)$  in the above solution by  $\frac{y}{x}$  results in the solution for  $y$  in implicit form

$$\frac{y^3}{x^3 \left( \frac{y^2}{x^2} - 1 \right)} = \frac{c_3}{x}$$

$$\frac{y^3}{x(y^2 - x^2)} = \frac{c_3}{x}$$

Which simplifies to

$$-\frac{y^3}{(-y+x)(y+x)} = c_3$$

### Summary

The solution(s) found are the following

$$-\frac{y^3}{(-y+x)(y+x)} = c_3 \quad (1)$$

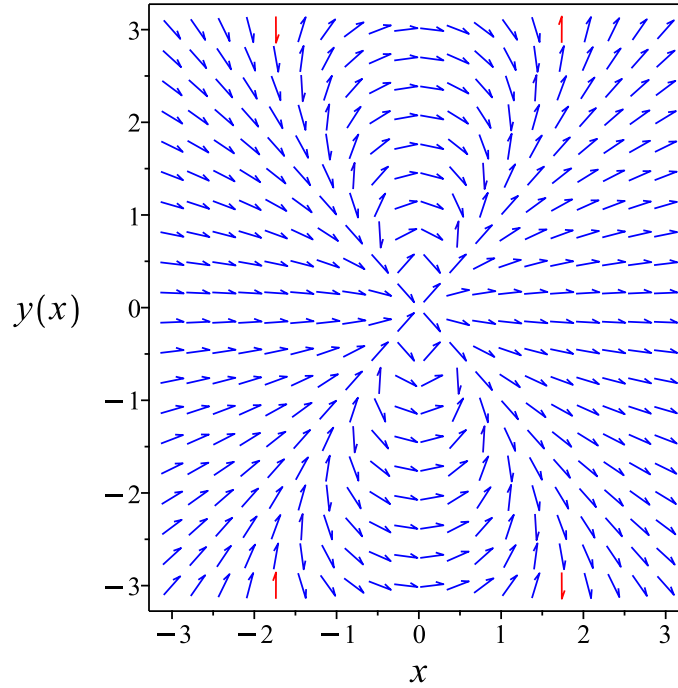


Figure 55: Slope field plot

### Verification of solutions

$$-\frac{y^3}{(-y+x)(y+x)} = c_3$$

Verified OK.

### 3.1.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{2yx}{-3x^2 + y^2}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} b_2 - \frac{2yx(b_3 - a_2)}{-3x^2 + y^2} - \frac{4y^2x^2a_3}{(-3x^2 + y^2)^2} \\ - \left( -\frac{2y}{-3x^2 + y^2} - \frac{12yx^2}{(-3x^2 + y^2)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left( -\frac{2x}{-3x^2 + y^2} + \frac{4y^2x}{(-3x^2 + y^2)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{3x^4b_2 + 2y^2x^2a_3 - 8x^2y^2b_2 + 4xy^3a_2 - 4xy^3b_3 + 2y^4a_3 + y^4b_2 - 6x^3b_1 + 6x^2ya_1 - 2xy^2b_1 + 2y^3a_1}{(3x^2 - y^2)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} 3x^4b_2 + 2y^2x^2a_3 - 8x^2y^2b_2 + 4xy^3a_2 - 4xy^3b_3 + 2y^4a_3 \\ + y^4b_2 - 6x^3b_1 + 6x^2ya_1 - 2xy^2b_1 + 2y^3a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} 4a_2v_1v_2^3 + 2a_3v_1^2v_2^2 + 2a_3v_2^4 + 3b_2v_1^4 - 8b_2v_1^2v_2^2 + b_2v_2^4 \\ - 4b_3v_1v_2^3 + 6a_1v_1^2v_2 + 2a_1v_2^3 - 6b_1v_1^3 - 2b_1v_1v_2^2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} 3b_2v_1^4 - 6b_1v_1^3 + (2a_3 - 8b_2)v_1^2v_2^2 + 6a_1v_1^2v_2 \\ + (4a_2 - 4b_3)v_1v_2^3 - 2b_1v_1v_2^2 + (2a_3 + b_2)v_2^4 + 2a_1v_2^3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 2a_1 &= 0 \\ 6a_1 &= 0 \\ -6b_1 &= 0 \\ -2b_1 &= 0 \\ 3b_2 &= 0 \\ 4a_2 - 4b_3 &= 0 \\ 2a_3 - 8b_2 &= 0 \\ 2a_3 + b_2 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left( -\frac{2yx}{-3x^2 + y^2} \right) (x) \\ &= \frac{x^2y - y^3}{3x^2 - y^2} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left( \xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{x^2y - y^3}{3x^2 - y^2}} dy\end{aligned}$$

Which results in

$$S = -\ln(y + x) - \ln(y - x) + 3 \ln(y)$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{2yx}{-3x^2 + y^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= -\frac{2x}{x^2 - y^2} \\S_y &= -\frac{1}{y+x} + \frac{1}{-y+x} + \frac{3}{y}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \tag{2A}$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

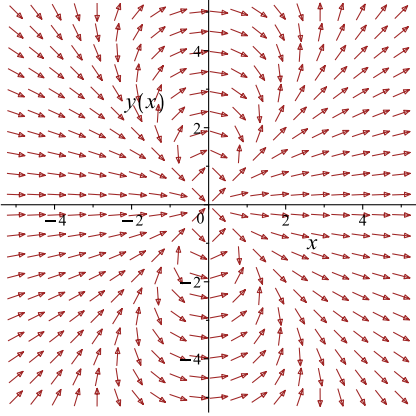
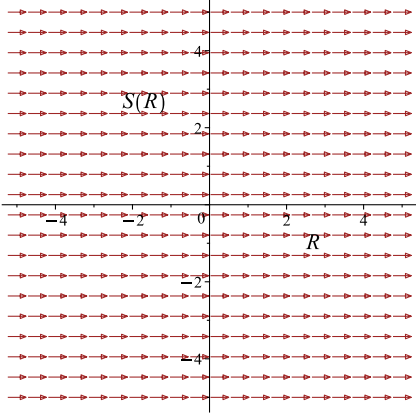
To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$-\ln(y+x) - \ln(y-x) + 3\ln(y) = c_1$$

Which simplifies to

$$-\ln(y+x) - \ln(y-x) + 3\ln(y) = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{2yx}{-3x^2+y^2}$ 	$R = x$ $S = -\ln(y+x) - \ln(y-x)$	$\frac{dS}{dR} = 0$ 

Summary

The solution(s) found are the following

$$-\ln(y+x) - \ln(y-x) + 3\ln(y) = c_1 \tag{1}$$



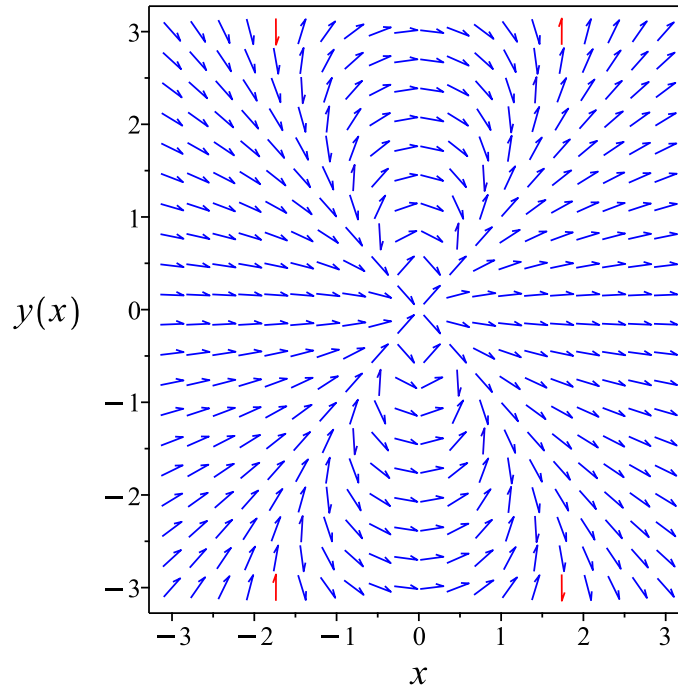


Figure 56: Slope field plot

Verification of solutions

$$-\ln(y+x) - \ln(y-x) + 3\ln(y) = c_1$$

Verified OK.

### 3.1.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(3x^2 - y^2) dy &= (2xy) dx \\ (-2xy) dx + (3x^2 - y^2) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -2xy \\ N(x, y) &= 3x^2 - y^2\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-2xy) \\ &= -2x\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(3x^2 - y^2) \\ &= 6x\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{3x^2 - y^2} ((-2x) - (6x)) \\ &= -\frac{8x}{3x^2 - y^2} \end{aligned}$$

Since  $A$  depends on  $y$ , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= -\frac{1}{2xy} ((6x) - (-2x)) \\ &= -\frac{4}{y} \end{aligned}$$

Since  $B$  does not depend on  $x$ , it can be used to obtain an integrating factor. Let the integrating factor be  $\mu$ . Then

$$\begin{aligned} \mu &= e^{\int B \, dy} \\ &= e^{\int -\frac{4}{y} \, dy} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-4 \ln(y)} \\ &= \frac{1}{y^4} \end{aligned}$$

$M$  and  $N$  are now multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\overline{M}$  and  $\overline{N}$  so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned} \overline{M} &= \mu M \\ &= \frac{1}{y^4} (-2xy) \\ &= -\frac{2x}{y^3} \end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{y^4}(3x^2 - y^2) \\ &= \frac{3x^2 - y^2}{y^4}\end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left(-\frac{2x}{y^3}\right) + \left(\frac{3x^2 - y^2}{y^4}\right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{2x}{y^3} dx \\ \phi &= -\frac{x^2}{y^3} + f(y)\end{aligned} \tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = \frac{3x^2}{y^4} + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{3x^2 - y^2}{y^4}$ . Therefore equation (4) becomes

$$\frac{3x^2 - y^2}{y^4} = \frac{3x^2}{y^4} + f'(y) \tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = -\frac{1}{y^2}$$

Integrating the above w.r.t  $y$  gives

$$\int f'(y) dy = \int \left(-\frac{1}{y^2}\right) dy$$

$$f(y) = \frac{1}{y} + c_1$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\frac{x^2}{y^3} + \frac{1}{y} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\frac{x^2}{y^3} + \frac{1}{y}$$

### Summary

The solution(s) found are the following

$$-\frac{x^2}{y^3} + \frac{1}{y} = c_1 \tag{1}$$

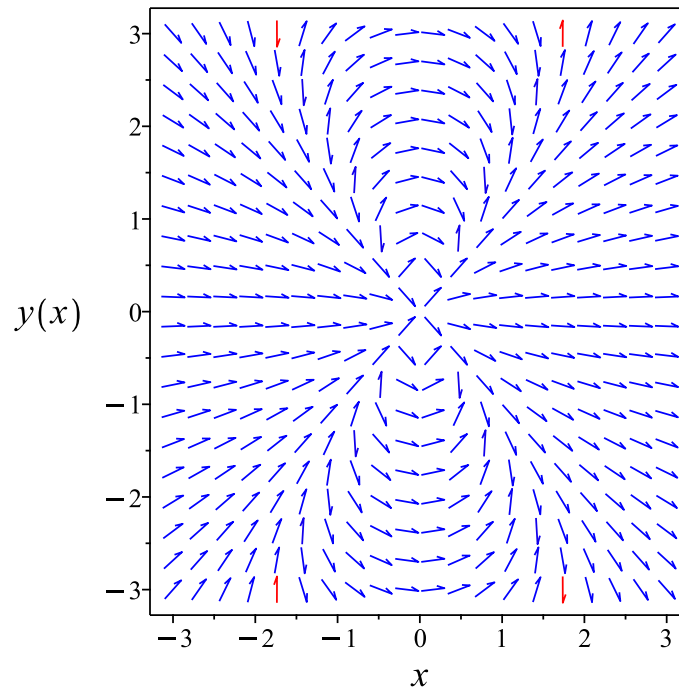


Figure 57: Slope field plot

Verification of solutions

$$-\frac{x^2}{y^3} + \frac{1}{y} = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 317

```
dsolve((3*x^2-y(x)^2)*diff(y(x),x)-2*x*y(x)=0,y(x), singsol=all)
```

$$y(x) = \frac{1 + \frac{\left(12\sqrt{3}x\sqrt{27c_1^2x^2 - 4c_1 - 108c_1^2x^2 + 8}\right)^{\frac{1}{3}}}{2} + \frac{2}{\left(12\sqrt{3}x\sqrt{27c_1^2x^2 - 4c_1 - 108c_1^2x^2 + 8}\right)^{\frac{1}{3}}}}{3c_1}$$

$$y(x) = \frac{(1 + i\sqrt{3}) \left(12\sqrt{3}x\sqrt{27c_1^2x^2 - 4c_1 - 108c_1^2x^2 + 8}\right)^{\frac{2}{3}} - 4i\sqrt{3} - 4 \left(12\sqrt{3}x\sqrt{27c_1^2x^2 - 4c_1 - 108c_1^2x^2 + 8}\right)^{\frac{1}{3}}}{12 \left(12\sqrt{3}x\sqrt{27c_1^2x^2 - 4c_1 - 108c_1^2x^2 + 8}\right)^{\frac{1}{3}} c_1}$$

$$y(x) = \frac{(i\sqrt{3} - 1) \left(12\sqrt{3}x\sqrt{27c_1^2x^2 - 4c_1 - 108c_1^2x^2 + 8}\right)^{\frac{2}{3}} - 4i\sqrt{3} + 4 \left(12\sqrt{3}x\sqrt{27c_1^2x^2 - 4c_1 - 108c_1^2x^2 + 8}\right)^{\frac{1}{3}}}{12 \left(12\sqrt{3}x\sqrt{27c_1^2x^2 - 4c_1 - 108c_1^2x^2 + 8}\right)^{\frac{1}{3}} c_1}$$

✓ Solution by Mathematica

Time used: 60.175 (sec). Leaf size: 458

`DSolve[(3*x^2-y[x]^2)*y'[x]-2*x*y[x]==0,y[x],x,IncludeSingularSolutions -> True]`

$$\begin{aligned}
 y(x) &\rightarrow \frac{1}{3} \left( \frac{\sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}}{\sqrt[3]{2}} \right. \\
 &\quad \left. + \frac{\sqrt[3]{2}e^{2c_1}}{\sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}} - e^{c_1} \right) \\
 y(x) &\rightarrow \frac{i(\sqrt{3} + i) \sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}}{6\sqrt[3]{2}} \\
 &\quad - \frac{i(\sqrt{3} - i) e^{2c_1}}{3 \cdot 2^{2/3} \sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}} - \frac{e^{c_1}}{3} \\
 y(x) &\rightarrow - \frac{i(\sqrt{3} - i) \sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}}{6\sqrt[3]{2}} \\
 &\quad + \frac{i(\sqrt{3} + i) e^{2c_1}}{3 \cdot 2^{2/3} \sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}} - \frac{e^{c_1}}{3}
 \end{aligned}$$



## 3.2 problem 2(b)

3.2.1 Solving as exact ode . . . . . 240

Internal problem ID [3101]

Internal file name [OUTPUT/2593\_Sunday\_June\_05\_2022\_03\_21\_38\_AM\_27792734/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 10, page 47

**Problem number:** 2(b).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**exactWithIntegrationFactor**"

Maple gives the following as the ode type

```
[_rational, [_1st_order, `_with_symmetry_[F(x),G(x)]`], [_Abel,
`2nd type`, `class B`]]
```

$$yx + (x^2 - yx) y' = 1$$

### 3.2.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (x^2 - xy) dy &= (-xy + 1) dx \\ (xy - 1) dx + (x^2 - xy) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= xy - 1 \\ N(x, y) &= x^2 - xy \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(xy - 1) \\ &= x \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x^2 - xy) \\ &= 2x - y \end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{x(-y+x)} ((x) - (2x-y)) \\ &= -\frac{1}{x} \end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int -\frac{1}{x} dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\ln(x)} \\ &= \frac{1}{x}\end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{x}(xy - 1) \\ &= \frac{xy - 1}{x}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{x}(x^2 - xy) \\ &= -y + x\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left( \frac{xy - 1}{x} \right) + (-y + x) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{xy - 1}{x} dx \\ \phi &= xy - \ln(x) + f(y)\end{aligned}\tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = x + f'(y)\tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = -y + x$ . Therefore equation (4) becomes

$$-y + x = x + f'(y)\tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = -y$$

Integrating the above w.r.t  $y$  gives

$$\begin{aligned}\int f'(y) dy &= \int (-y) dy \\ f(y) &= -\frac{y^2}{2} + c_1\end{aligned}$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = xy - \ln(x) - \frac{y^2}{2} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = xy - \ln(x) - \frac{y^2}{2}$$

Summary

The solution(s) found are the following

$$yx - \ln(x) - \frac{y^2}{2} = c_1 \tag{1}$$

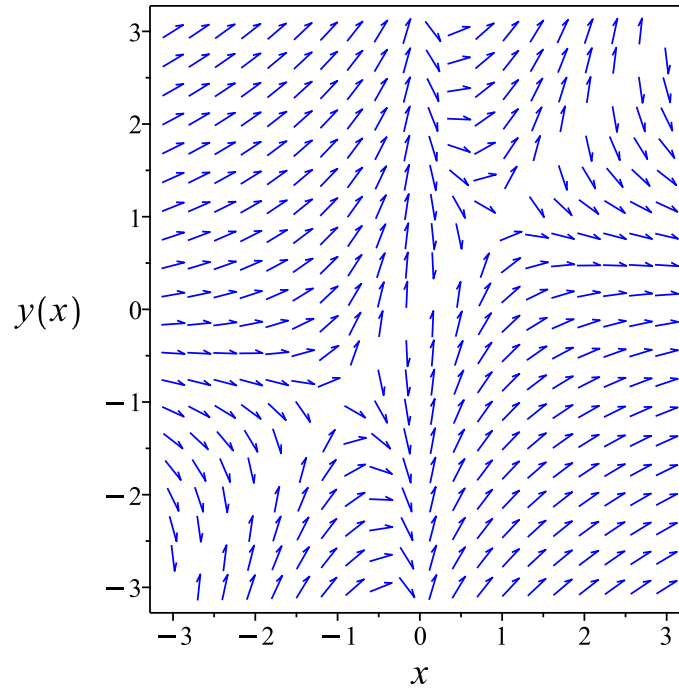


Figure 58: Slope field plot

Verification of solutions

$$yx - \ln(x) - \frac{y^2}{2} = c_1$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

### ✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 39

```
dsolve((x*y(x)-1)+(x^2-x*y(x))*diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = x - \sqrt{x^2 - 2 \ln(x) + 2c_1}$$
$$y(x) = x + \sqrt{x^2 - 2 \ln(x) + 2c_1}$$

### ✓ Solution by Mathematica

Time used: 0.46 (sec). Leaf size: 68

```
DSolve[(x*y[x]-1)+(x^2-x*y[x])*y'[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow x + \sqrt{-\frac{1}{x} \sqrt{-x(x^2 - 2 \log(x) + c_1)}}$$
$$y(x) \rightarrow x + x \left(-\frac{1}{x}\right)^{3/2} \sqrt{-x(x^2 - 2 \log(x) + c_1)}$$

### 3.3 problem 2(c)

3.3.1 Solving as first order ode lie symmetry calculated ode . . . . .	246
3.3.2 Solving as exact ode . . . . .	252

Internal problem ID [3102]

Internal file name [OUTPUT/2594\_Sunday\_June\_05\_2022\_03\_21\_40\_AM\_91788273/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 10, page 47

**Problem number:** 2(c).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exactWithIntegrationFactor",  
"first\_order\_ode\_lie\_symmetry\_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class G`], _rational]
```

$$(x + 3y^4x^3) y' + y = 0$$

#### 3.3.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{y}{x(3x^2y^4 + 1)}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2\xi_y - \omega_x\xi - \omega_y\eta = 0 \tag{A}$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \tag{1E}$$

$$\eta = xb_2 + yb_3 + b_1 \tag{2E}$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} b_2 - \frac{y(b_3 - a_2)}{x(3x^2y^4 + 1)} - \frac{y^2a_3}{x^2(3x^2y^4 + 1)^2} \\ - \left( \frac{y}{x^2(3x^2y^4 + 1)} + \frac{6y^5}{(3x^2y^4 + 1)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left( -\frac{1}{x(3x^2y^4 + 1)} + \frac{12y^4x}{(3x^2y^4 + 1)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{9x^6y^8b_2 - 3x^4y^4b_2 - 6x^3y^5a_2 - 12x^3y^5b_3 - 9x^2y^6a_3 - 9x^3y^4b_1 - 9x^2y^5a_1 + 2b_2x^2 - 2y^2a_3 + xb_1 - ya_1}{x^2(3x^2y^4 + 1)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} 9x^6y^8b_2 - 3x^4y^4b_2 - 6x^3y^5a_2 - 12x^3y^5b_3 - 9x^2y^6a_3 \\ - 9x^3y^4b_1 - 9x^2y^5a_1 + 2b_2x^2 - 2y^2a_3 + xb_1 - ya_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} 9b_2v_1^6v_2^8 - 6a_2v_1^3v_2^5 - 9a_3v_1^2v_2^6 - 3b_2v_1^4v_2^4 - 12b_3v_1^3v_2^5 \\ - 9a_1v_1^2v_2^5 - 9b_1v_1^3v_2^4 - 2a_3v_2^2 + 2b_2v_1^2 - a_1v_2 + b_1v_1 = 0 \end{aligned} \quad (7E)$$



Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$9b_2v_1^6v_2^8 - 3b_2v_1^4v_2^4 + (-6a_2 - 12b_3)v_1^3v_2^5 - 9b_1v_1^3v_2^4 - 9a_3v_1^2v_2^6 - 9a_1v_1^2v_2^5 + 2b_2v_1^2 + b_1v_1 - 2a_3v_2^2 - a_1v_2 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} b_1 &= 0 \\ -9a_1 &= 0 \\ -a_1 &= 0 \\ -9a_3 &= 0 \\ -2a_3 &= 0 \\ -9b_1 &= 0 \\ -3b_2 &= 0 \\ 2b_2 &= 0 \\ 9b_2 &= 0 \\ -6a_2 - 12b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= -2b_3 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= -2x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left( -\frac{y}{x(3x^2y^4 + 1)} \right) (-2x) \\ &= \frac{3x^2y^5 - y}{3x^2y^4 + 1} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left( \xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{3x^2y^5 - y}{3x^2y^4 + 1}} dy\end{aligned}$$

Which results in

$$S = \frac{\ln(3x^2y^4 - 1)}{2} - \ln(y)$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y}{x(3x^2y^4 + 1)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= \frac{3y^4x}{3x^2y^4 - 1} \\S_y &= \frac{3x^2y^4 + 1}{3x^2y^5 - y}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{x} \tag{2A}$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = \ln(R) + c_1 \tag{4}$$

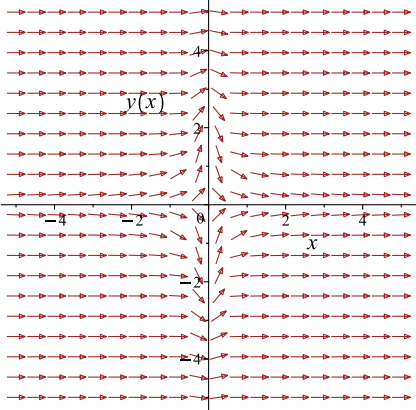
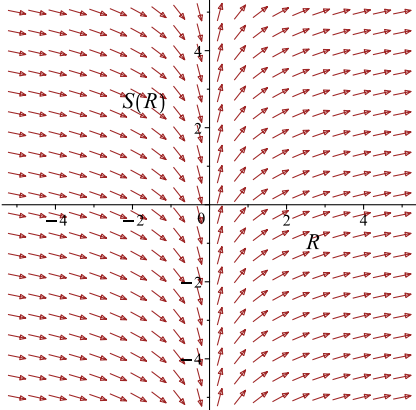
To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{\ln(3x^2y^4 - 1)}{2} - \ln(y) = \ln(x) + c_1$$

Which simplifies to

$$\frac{\ln(3x^2y^4 - 1)}{2} - \ln(y) = \ln(x) + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{y}{x(3x^2y^4+1)}$ 	$R = x$ $S = \frac{\ln(3x^2y^4 - 1)}{2} - \ln(y)$	$\frac{dS}{dR} = \frac{1}{R}$ 

Summary

The solution(s) found are the following

$$\frac{\ln(3x^2y^4 - 1)}{2} - \ln(y) = \ln(x) + c_1 \tag{1}$$

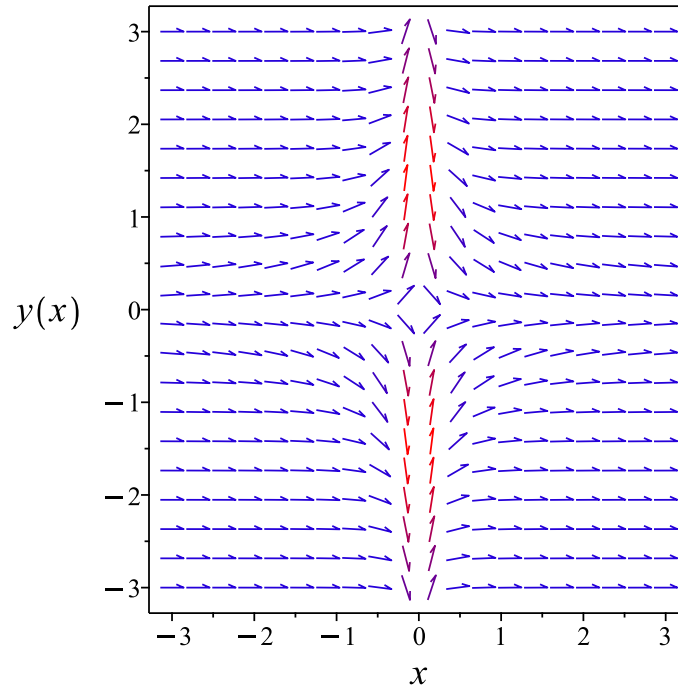


Figure 59: Slope field plot

Verification of solutions

$$\frac{\ln(3x^2y^4 - 1)}{2} - \ln(y) = \ln(x) + c_1$$

Verified OK.

### 3.3.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(3x^3y^4 + x) dy &= (-y) dx \\ (y) dx + (3x^3y^4 + x) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= y \\ N(x, y) &= 3x^3y^4 + x\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(3x^3y^4 + x) \\ &= 9x^2y^4 + 1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{3x^3y^4 + x} ((1) - (9x^2y^4 + 1)) \\ &= -\frac{9xy^4}{3x^2y^4 + 1} \end{aligned}$$

Since  $A$  depends on  $y$ , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= \frac{1}{y} ((9x^2y^4 + 1) - (1)) \\ &= 9y^3x^2 \end{aligned}$$

Since  $B$  depends on  $x$ , it can not be used to obtain an integrating factor. We will now try a third method to find an integrating factor. Let

$$R = \frac{\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y}}{xM - yN}$$

$R$  is now checked to see if it is a function of only  $t = xy$ . Therefore

$$\begin{aligned} R &= \frac{\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y}}{xM - yN} \\ &= \frac{(9x^2y^4 + 1) - (1)}{x(y) - y(3x^3y^4 + x)} \\ &= -\frac{3}{xy} \end{aligned}$$

Replacing all powers of terms  $xy$  by  $t$  gives

$$R = -\frac{3}{t}$$

Since  $R$  depends on  $t$  only, then it can be used to find an integrating factor. Let the integrating factor be  $\mu$  then

$$\begin{aligned} \mu &= e^{\int R dt} \\ &= e^{\int (-\frac{3}{t}) dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-3\ln(t)} \\ &= \frac{1}{t^3}\end{aligned}$$

Now  $t$  is replaced back with  $xy$  giving

$$\mu = \frac{1}{x^3y^3}$$

Multiplying  $M$  and  $N$  by this integrating factor gives new  $\bar{M}$  and new  $\bar{N}$  which are called  $\bar{M}$  and  $\bar{N}$  so not to confuse them with the original  $M$  and  $N$

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{x^3y^3}(y) \\ &= \frac{1}{y^2x^3}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{x^3y^3}(3x^3y^4 + x) \\ &= \frac{3x^2y^4 + 1}{x^2y^3}\end{aligned}$$

A modified ODE is now obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left( \frac{1}{y^2x^3} \right) + \left( \frac{3x^2y^4 + 1}{x^2y^3} \right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$



Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{1}{y^2 x^3} dx \\ \phi &= -\frac{1}{2y^2 x^2} + f(y)\end{aligned}\quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = \frac{1}{y^3 x^2} + f'(y)\quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{3x^2 y^4 + 1}{x^2 y^3}$ . Therefore equation (4) becomes

$$\frac{3x^2 y^4 + 1}{x^2 y^3} = \frac{1}{y^3 x^2} + f'(y)\quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 3y$$

Integrating the above w.r.t  $y$  gives

$$\begin{aligned}\int f'(y) dy &= \int (3y) dy \\ f(y) &= \frac{3y^2}{2} + c_1\end{aligned}$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\frac{1}{2y^2 x^2} + \frac{3y^2}{2} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\frac{1}{2y^2 x^2} + \frac{3y^2}{2}$$

### Summary

The solution(s) found are the following

$$-\frac{1}{2y^2x^2} + \frac{3y^2}{2} = c_1 \quad (1)$$

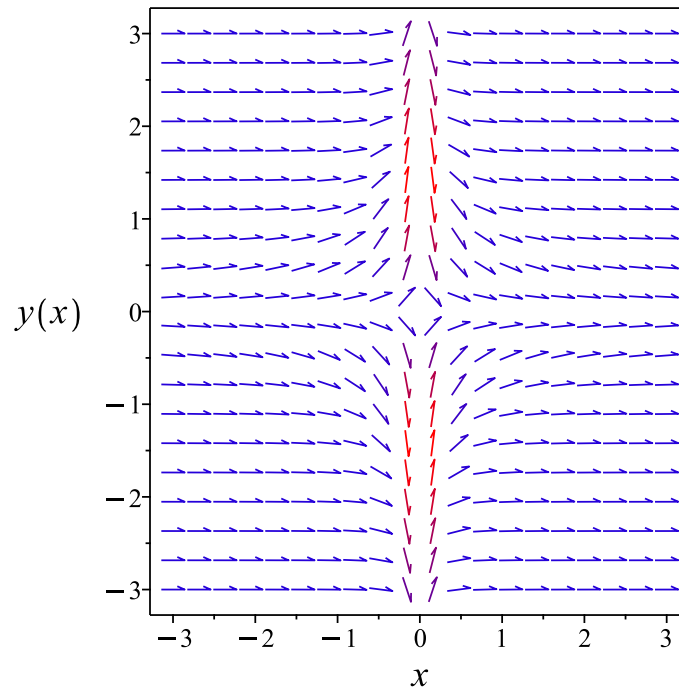


Figure 60: Slope field plot

### Verification of solutions

$$-\frac{1}{2y^2x^2} + \frac{3y^2}{2} = c_1$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous G  
<- homogeneous successful`
```

### ✓ Solution by Maple

Time used: 0.125 (sec). Leaf size: 133

```
dsolve((x+3*x^3*y(x)^4)*diff(y(x),x)+y(x)=0,y(x), singsol=all)
```

$$y(x) = -\frac{\sqrt{6} \sqrt{xc_1 \left( x - \sqrt{12c_1^2 + x^2} \right)}}{6xc_1}$$

$$y(x) = \frac{\sqrt{6} \sqrt{xc_1 \left( x - \sqrt{12c_1^2 + x^2} \right)}}{6xc_1}$$

$$y(x) = -\frac{\sqrt{6} \sqrt{xc_1 \left( x + \sqrt{12c_1^2 + x^2} \right)}}{6xc_1}$$

$$y(x) = \frac{\sqrt{6} \sqrt{xc_1 \left( x + \sqrt{12c_1^2 + x^2} \right)}}{6xc_1}$$

✓ Solution by Mathematica

Time used: 10.044 (sec). Leaf size: 166

```
DSolve[(x+3*x^3*y[x]^4)*y'[x]+y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{\sqrt{c_1 - \frac{\sqrt{x^2(3+c_1^2x^2)}}{x^2}}}{\sqrt{3}}$$

$$y(x) \rightarrow \frac{\sqrt{c_1 - \frac{\sqrt{x^2(3+c_1^2x^2)}}{x^2}}}{\sqrt{3}}$$

$$y(x) \rightarrow -\frac{\sqrt{\frac{\sqrt{x^2(3+c_1^2x^2)}}{x^2} + c_1}}{\sqrt{3}}$$

$$y(x) \rightarrow \frac{\sqrt{\frac{\sqrt{x^2(3+c_1^2x^2)}}{x^2} + c_1}}{\sqrt{3}}$$

$$y(x) \rightarrow 0$$

### 3.4 problem 4(a)

3.4.1	Solving as first order ode lie symmetry calculated ode . . . . .	260
3.4.2	Solving as exact ode . . . . .	265

Internal problem ID [3103]

Internal file name [OUTPUT/2595\_Sunday\_June\_05\_2022\_03\_21\_43\_AM\_85572339/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 10, page 47

**Problem number:** 4(a).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exactWithIntegrationFactor",  
"first\_order\_ode\_lie\_symmetry\_calculated"

Maple gives the following as the ode type

```
[[_1st_order , _with_linear_symmetries], _rational]
```

$(x - 1 - y^2) y' - y = 0$
----------------------------

#### 3.4.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{y}{y^2 - x + 1}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \tag{1E}$$

$$\eta = xb_2 + yb_3 + b_1 \tag{2E}$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$b_2 - \frac{y(b_3 - a_2)}{y^2 - x + 1} - \frac{y^2 a_3}{(y^2 - x + 1)^2} + \frac{y(xa_2 + ya_3 + a_1)}{(y^2 - x + 1)^2} \quad (5E)$$

$$- \left( -\frac{1}{y^2 - x + 1} + \frac{2y^2}{(y^2 - x + 1)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{-y^4 b_2 + 3x y^2 b_2 - y^3 a_2 + 2y^3 b_3 + y^2 b_1 - 2y^2 b_2 + x b_1 + x b_2 - y a_1 - y a_2 - b_1 - b_2}{(-y^2 + x - 1)^2} = 0$$

Setting the numerator to zero gives

$$y^4 b_2 - 3x y^2 b_2 + y^3 a_2 - 2y^3 b_3 - y^2 b_1 + 2y^2 b_2 - x b_1 - x b_2 + y a_1 + y a_2 + b_1 + b_2 = 0 \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$b_2 v_2^4 + a_2 v_2^3 - 3b_2 v_1 v_2^2 - 2b_3 v_2^3 - b_1 v_2^2 + 2b_2 v_2^2 + a_1 v_2 + a_2 v_2 - b_1 v_1 - b_2 v_1 + b_1 + b_2 = 0 \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$-3b_2 v_1 v_2^2 + (-b_1 - b_2) v_1 + b_2 v_2^4 + (a_2 - 2b_3) v_2^3 \quad (8E)$$

$$+ (-b_1 + 2b_2) v_2^2 + (a_1 + a_2) v_2 + b_1 + b_2 = 0$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}
 b_2 &= 0 \\
 -3b_2 &= 0 \\
 a_1 + a_2 &= 0 \\
 a_2 - 2b_3 &= 0 \\
 -b_1 - b_2 &= 0 \\
 -b_1 + 2b_2 &= 0 \\
 b_1 + b_2 &= 0
 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= -2b_3 \\
 a_2 &= 2b_3 \\
 a_3 &= a_3 \\
 b_1 &= 0 \\
 b_2 &= 0 \\
 b_3 &= b_3
 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}
 \xi &= y \\
 \eta &= 0
 \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}
 \eta &= \eta - \omega(x, y) \xi \\
 &= 0 - \left( -\frac{y}{y^2 - x + 1} \right) (y) \\
 &= -\frac{y^2}{-y^2 + x - 1} \\
 \xi &= 0
 \end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{-\frac{y^2}{-y^2+x-1}} dy \end{aligned}$$

Which results in

$$S = y + \frac{x-1}{y}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y}{y^2 - x + 1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{1}{y} \\ S_y &= \frac{y^2 - x + 1}{y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$



The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

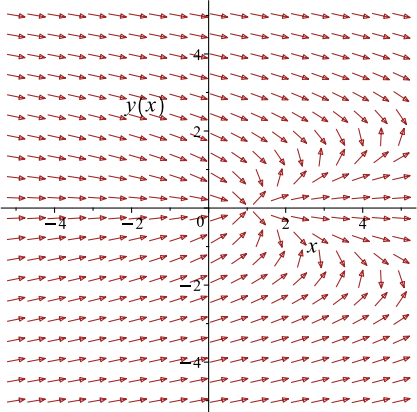
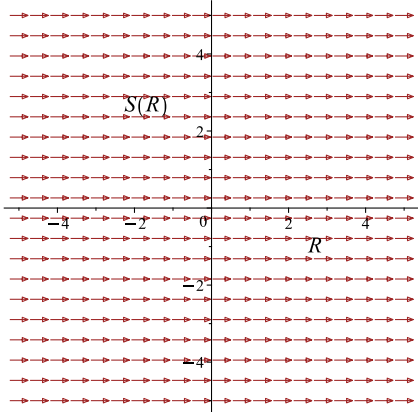
To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{y^2 + x - 1}{y} = c_1$$

Which simplifies to

$$\frac{y^2 + x - 1}{y} = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{y}{y^2 - x + 1}$ 	$R = x$ $S = \frac{y^2 + x - 1}{y}$	$\frac{dS}{dR} = 0$ 

### Summary

The solution(s) found are the following

$$\frac{y^2 + x - 1}{y} = c_1 \tag{1}$$

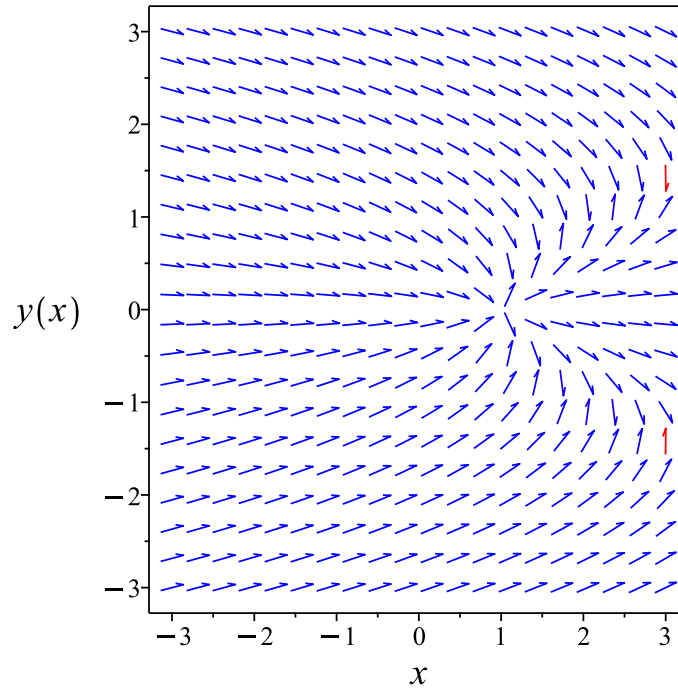


Figure 61: Slope field plot

Verification of solutions

$$\frac{y^2 + x - 1}{y} = c_1$$

Verified OK.

### 3.4.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(-y^2 + x - 1) dy &= (y) dx \\ (-y) dx + (-y^2 + x - 1) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -y \\ N(x, y) &= -y^2 + x - 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-y) \\ &= -1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-y^2 + x - 1) \\ &= 1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{-y^2 + x - 1} ((-1) - (1)) \\ &= -\frac{2}{-y^2 + x - 1} \end{aligned}$$

Since  $A$  depends on  $y$ , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= -\frac{1}{y} ((1) - (-1)) \\ &= -\frac{2}{y} \end{aligned}$$

Since  $B$  does not depend on  $x$ , it can be used to obtain an integrating factor. Let the integrating factor be  $\mu$ . Then

$$\begin{aligned} \mu &= e^{\int B \, dy} \\ &= e^{\int -\frac{2}{y} \, dy} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-2 \ln(y)} \\ &= \frac{1}{y^2} \end{aligned}$$

$M$  and  $N$  are now multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\overline{M}$  and  $\overline{N}$  so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned} \overline{M} &= \mu M \\ &= \frac{1}{y^2} (-y) \\ &= -\frac{1}{y} \end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{y^2}(-y^2 + x - 1) \\ &= \frac{-y^2 + x - 1}{y^2}\end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left(-\frac{1}{y}\right) + \left(\frac{-y^2 + x - 1}{y^2}\right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{1}{y} dx \\ \phi &= -\frac{x}{y} + f(y)\end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = \frac{x}{y^2} + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{-y^2+x-1}{y^2}$ . Therefore equation (4) becomes

$$\frac{-y^2 + x - 1}{y^2} = \frac{x}{y^2} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = -\frac{y^2 + 1}{y^2}$$

Integrating the above w.r.t  $y$  gives

$$\int f'(y) \, dy = \int \left( \frac{-y^2 - 1}{y^2} \right) \, dy$$
$$f(y) = -y + \frac{1}{y} + c_1$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\frac{x}{y} - y + \frac{1}{y} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\frac{x}{y} - y + \frac{1}{y}$$

### Summary

The solution(s) found are the following

$$-\frac{x}{y} - y + \frac{1}{y} = c_1 \tag{1}$$

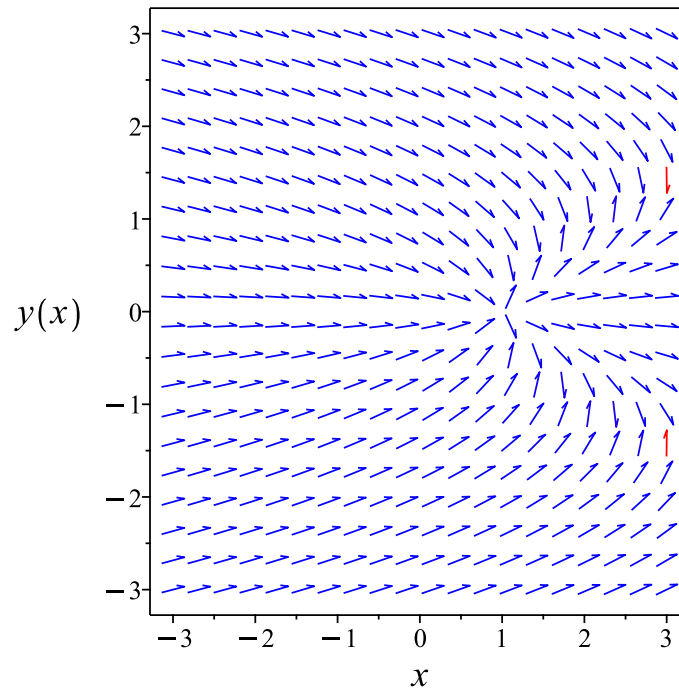


Figure 62: Slope field plot

Verification of solutions

$$-\frac{x}{y} - y + \frac{1}{y} = c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
<- 1st order linear successful
<- inverse linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 39

```
dsolve((x-1-y(x)^2)*diff(y(x),x)-y(x)=0,y(x), singsol=all)
```

$$y(x) = \frac{c_1}{2} - \frac{\sqrt{c_1^2 - 4x + 4}}{2}$$

$$y(x) = \frac{c_1}{2} + \frac{\sqrt{c_1^2 - 4x + 4}}{2}$$

✓ Solution by Mathematica

Time used: 0.304 (sec). Leaf size: 56

```
DSolve[(x-1-y[x]^2)*y'[x]-y[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{2} \left( c_1 - \sqrt{-4x + 4 + c_1^2} \right)$$

$$y(x) \rightarrow \frac{1}{2} \left( \sqrt{-4x + 4 + c_1^2} + c_1 \right)$$

$$y(x) \rightarrow 0$$



### 3.5 problem 4(b)

3.5.1	Solving as separable ode . . . . .	272
3.5.2	Solving as first order ode lie symmetry lookup ode . . . . .	274
3.5.3	Solving as exact ode . . . . .	279
3.5.4	Maple step by step solution . . . . .	282

Internal problem ID [3104]

Internal file name [OUTPUT/2596\_Sunday\_June\_05\_2022\_03\_21\_46\_AM\_66865626/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 10, page 47

**Problem number:** 4(b).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exact", "separable", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_separable]

$$y - (x + xy^3) y' = 0$$

#### 3.5.1 Solving as separable ode

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= f(x)g(y) \\ &= \frac{y}{x(y^3 + 1)} \end{aligned}$$

Where  $f(x) = \frac{1}{x}$  and  $g(y) = \frac{y}{y^3+1}$ . Integrating both sides gives

$$\begin{aligned} \frac{1}{\frac{y}{y^3+1}} dy &= \frac{1}{x} dx \\ \int \frac{1}{\frac{y}{y^3+1}} dy &= \int \frac{1}{x} dx \\ \frac{y^3}{3} + \ln(y) &= \ln(x) + c_1 \end{aligned}$$

Which results in

$$y = \frac{1}{\left(\frac{1}{\text{LambertW}(x^3 e^{3c_1})}\right)^{\frac{1}{3}}}$$

Since  $c_1$  is constant, then exponential powers of this constant are constants also, and these can be simplified to just  $c_1$  in the above solution. Which simplifies to

$$y = \frac{1}{\left(\frac{1}{\text{LambertW}(x^3 e^{3c_1})}\right)^{\frac{1}{3}}}$$

gives

$$y = \frac{1}{\left(\frac{1}{\text{LambertW}(c_1^3 x^3)}\right)^{\frac{1}{3}}}$$

### Summary

The solution(s) found are the following

$$y = \frac{1}{\left(\frac{1}{\text{LambertW}(c_1^3 x^3)}\right)^{\frac{1}{3}}} \tag{1}$$

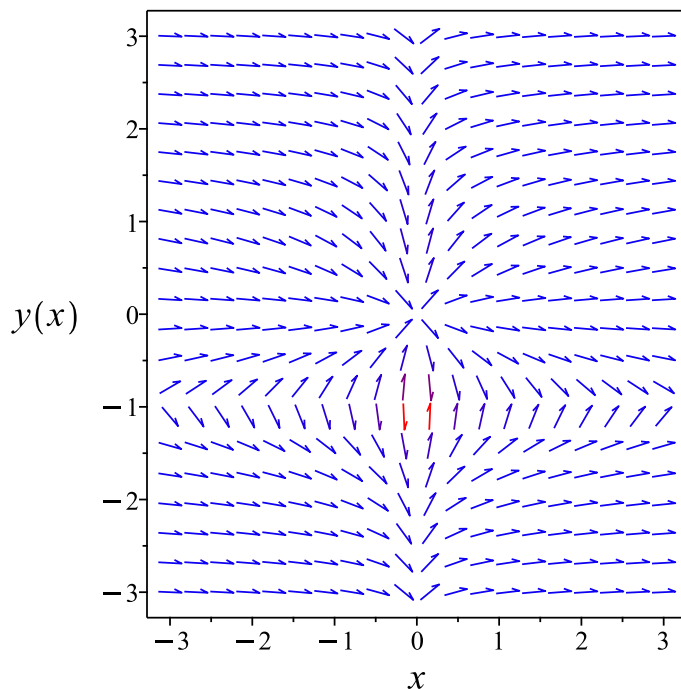


Figure 63: Slope field plot

### Verification of solutions

$$y = \frac{1}{\left(\frac{1}{\text{LambertW}(c_1^3 x^3)}\right)^{\frac{1}{3}}}$$

Verified OK.

### **3.5.2 Solving as first order ode lie symmetry lookup ode**

Writing the ode as

$$y' = \frac{y}{x(y^3 + 1)}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **separable**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 28: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= x \\ \eta(x, y) &= 0\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\eta = 0$  then in this special case

$$R = y$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\xi} dx \\ &= \int \frac{1}{x} dx \end{aligned}$$

Which results in

$$S = \ln(x)$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y}{x(y^3 + 1)}$$

Evaluating all the partial derivatives gives

$$R_x = 0$$

$$R_y = 1$$

$$S_x = \frac{1}{x}$$

$$S_y = 0$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{y^3 + 1}{y} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{R^3 + 1}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = \frac{R^3}{3} + \ln(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\ln(x) = \frac{y^3}{3} + \ln(y) + c_1$$

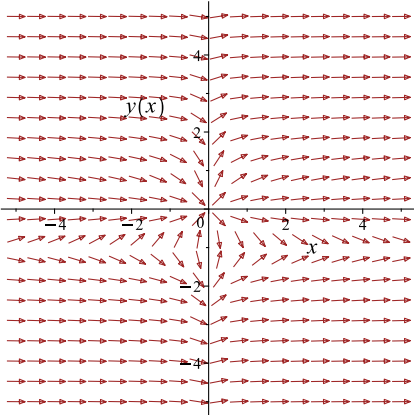
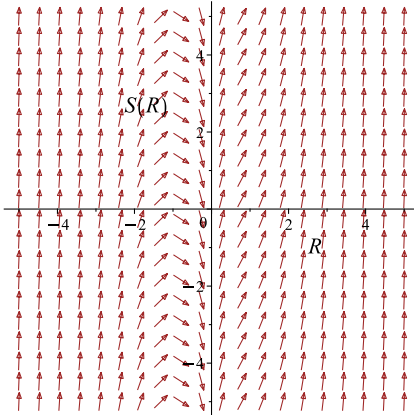
Which simplifies to

$$\ln(x) = \frac{y^3}{3} + \ln(y) + c_1$$

Which gives

$$y = \frac{1}{\left(\frac{1}{\text{LambertW}(e^{-3c_1 x^3})}\right)^{\frac{1}{3}}}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{y}{x(y^3+1)}$ 	$R = y$ $S = \ln(x)$	$\frac{dS}{dR} = \frac{R^3+1}{R}$ 

### Summary

The solution(s) found are the following

$$y = \frac{1}{\left(\frac{1}{\text{LambertW}(e^{-3c_1 x^3})}\right)^{\frac{1}{3}}} \quad (1)$$

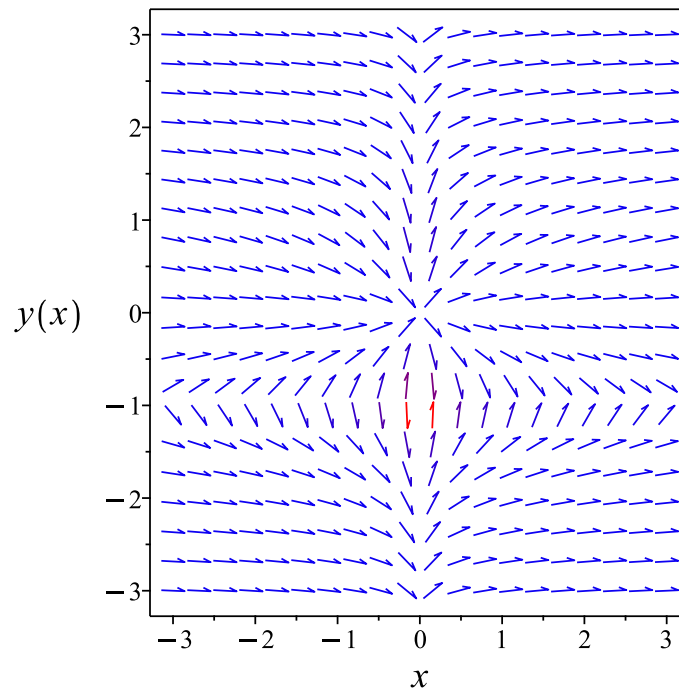


Figure 64: Slope field plot

### Verification of solutions

$$y = \frac{1}{\left(\frac{1}{\text{LambertW}(e^{-3c_1 x^3})}\right)^{\frac{1}{3}}}$$

Verified OK.

### 3.5.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} \left( \frac{y^3 + 1}{y} \right) dy &= \left( \frac{1}{x} \right) dx \\ \left( -\frac{1}{x} \right) dx + \left( \frac{y^3 + 1}{y} \right) dy &= 0 \end{aligned} \quad (\text{2A})$$



Comparing (1A) and (2A) shows that

$$M(x, y) = -\frac{1}{x}$$
$$N(x, y) = \frac{y^3 + 1}{y}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y} \left( -\frac{1}{x} \right)$$
$$= 0$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x} \left( \frac{y^3 + 1}{y} \right)$$
$$= 0$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$
$$\int \frac{\partial \phi}{\partial x} dx = \int -\frac{1}{x} dx$$
$$\phi = -\ln(x) + f(y) \tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{y^3+1}{y}$ . Therefore equation (4) becomes

$$\frac{y^3 + 1}{y} = 0 + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = \frac{y^3 + 1}{y}$$

Integrating the above w.r.t  $y$  gives

$$\int f'(y) dy = \int \left( \frac{y^3 + 1}{y} \right) dy$$

$$f(y) = \frac{y^3}{3} + \ln(y) + c_1$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\ln(x) + \frac{y^3}{3} + \ln(y) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\ln(x) + \frac{y^3}{3} + \ln(y)$$

The solution becomes

$$y = \frac{1}{\left( \frac{1}{\text{LambertW}(x^3 e^{3c_1})} \right)^{\frac{1}{3}}}$$

### Summary

The solution(s) found are the following

$$y = \frac{1}{\left( \frac{1}{\text{LambertW}(x^3 e^{3c_1})} \right)^{\frac{1}{3}}} \quad (1)$$

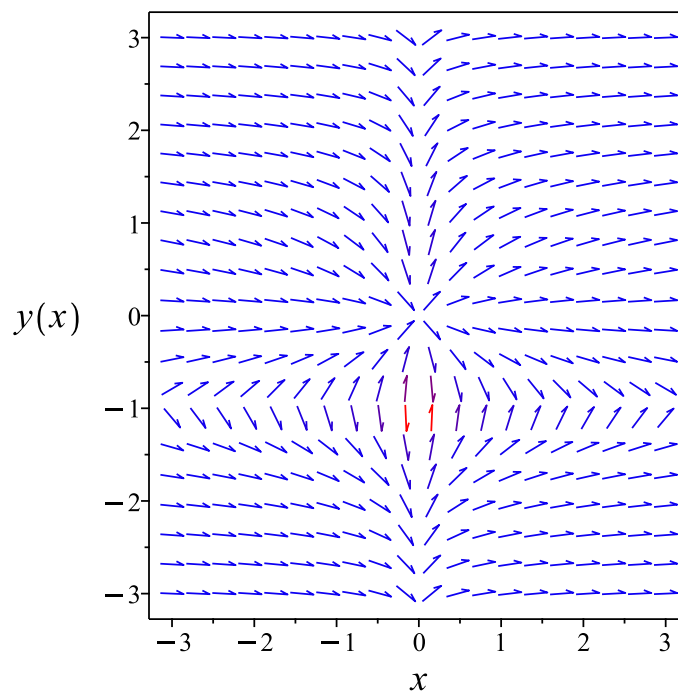


Figure 65: Slope field plot

### Verification of solutions

$$y = \frac{1}{\left(\frac{1}{\text{LambertW}(x^3 e^{3c_1})}\right)^{\frac{1}{3}}}$$

Verified OK.

### 3.5.4 Maple step by step solution

Let's solve

$$y - (x + xy^3) y' = 0$$

- Highest derivative means the order of the ODE is 1

$y'$

- Separate variables

$$\frac{y'(y+1)(y^2-y+1)}{y} = \frac{1}{x}$$

- Integrate both sides with respect to  $x$

$$\int \frac{y'(y+1)(y^2-y+1)}{y} dx = \int \frac{1}{x} dx + c_1$$

- Evaluate integral  

$$\frac{y^3}{3} + \ln(y) = \ln(x) + c_1$$

- Solve for  $y$

$$y = \frac{x}{e^{\frac{\text{LambertW}(x^3 e^{3c_1})}{3} - c_1}}$$

### Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

#### ✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 14

```
dsolve(y(x)-(x+x*y(x)^3)*diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = \frac{1}{\left(\frac{1}{\text{LambertW}(c_1 x^3)}\right)^{\frac{1}{3}}}$$

#### ✓ Solution by Mathematica

Time used: 4.377 (sec). Leaf size: 76

```
DSolve[y[x]-(x+x*y[x]^3)*y'[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$\begin{aligned}
 y(x) &\rightarrow \sqrt[3]{W(e^{3c_1} x^3)} \\
 y(x) &\rightarrow -\sqrt[3]{-1} \sqrt[3]{W(e^{3c_1} x^3)} \\
 y(x) &\rightarrow (-1)^{2/3} \sqrt[3]{W(e^{3c_1} x^3)} \\
 y(x) &\rightarrow 0
 \end{aligned}$$

### 3.6 problem 4(c)

3.6.1	Solving as homogeneousTypeD2 ode . . . . .	284
3.6.2	Solving as exact ode . . . . .	286
3.6.3	Solving as riccati ode . . . . .	291

Internal problem ID [3105]

Internal file name [OUTPUT/2597\_Sunday\_June\_05\_2022\_03\_21\_48\_AM\_92210252/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 10, page 47

**Problem number:** 4(c).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"riccati", "exactByInspection", "homogeneousTypeD2"**

Maple gives the following as the ode type

```
[[_homogeneous, `class D`], _rational, _Riccati]
```

$$xy' - y^2x^3 - y = x^5$$

#### 3.6.1 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$x(u'(x)x + u(x)) - u(x)^2x^5 - u(x)x = x^5$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= x^3(u^2 + 1)\end{aligned}$$

Where  $f(x) = x^3$  and  $g(u) = u^2 + 1$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{u^2 + 1} du &= x^3 dx \\ \int \frac{1}{u^2 + 1} du &= \int x^3 dx \\ \arctan(u) &= \frac{x^4}{4} + c_2\end{aligned}$$

The solution is

$$\arctan(u(x)) - \frac{x^4}{4} - c_2 = 0$$

Replacing  $u(x)$  in the above solution by  $\frac{y}{x}$  results in the solution for  $y$  in implicit form

$$\arctan\left(\frac{y}{x}\right) - \frac{x^4}{4} - c_2 = 0$$

$$\arctan\left(\frac{y}{x}\right) - \frac{x^4}{4} - c_2 = 0$$

### Summary

The solution(s) found are the following

$$\arctan\left(\frac{y}{x}\right) - \frac{x^4}{4} - c_2 = 0 \quad (1)$$

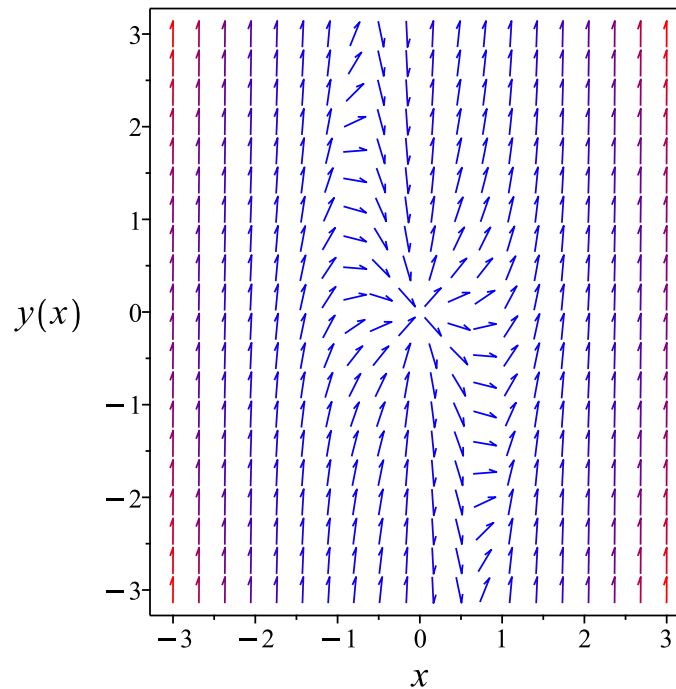


Figure 66: Slope field plot

### Verification of solutions

$$\arctan\left(\frac{y}{x}\right) - \frac{x^4}{4} - c_2 = 0$$

Verified OK.

### 3.6.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (x) dy &= (x^5 + x^3y^2 + y) dx \\ (-x^5 - x^3y^2 - y) dx + (x) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x^5 - x^3y^2 - y \\ N(x, y) &= x \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x^5 - x^3y^2 - y) \\ &= -2x^3y - 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x) \\ &= 1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. By inspection  $\frac{1}{y^2+x^2}$  is an integrating factor. Therefore by multiplying  $M = -y^2x^3 - x^5 - y$  and  $N = x$  by this integrating factor the ode becomes exact. The new  $M, N$  are

$$\begin{aligned}M &= \frac{-y^2x^3 - x^5 - y}{y^2 + x^2} \\ N &= \frac{x}{y^2 + x^2}\end{aligned}$$

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$



But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left( \frac{x}{x^2 + y^2} \right) dy &= \left( -\frac{-x^5 - x^3 y^2 - y}{x^2 + y^2} \right) dx \\ \left( \frac{-x^5 - x^3 y^2 - y}{x^2 + y^2} \right) dx + \left( \frac{x}{x^2 + y^2} \right) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= \frac{-x^5 - x^3 y^2 - y}{x^2 + y^2} \\ N(x, y) &= \frac{x}{x^2 + y^2} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left( \frac{-x^5 - x^3 y^2 - y}{x^2 + y^2} \right) \\ &= \frac{-x^2 + y^2}{(x^2 + y^2)^2} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left( \frac{x}{x^2 + y^2} \right) \\ &= \frac{-x^2 + y^2}{(x^2 + y^2)^2} \end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{-x^5 - x^3 y^2 - y}{x^2 + y^2} dx \\ \phi &= -\frac{x^4}{4} - \arctan\left(\frac{x}{y}\right) + f(y) \end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\begin{aligned} \frac{\partial \phi}{\partial y} &= \frac{x}{y^2 \left(\frac{x^2}{y^2} + 1\right)} + f'(y) \\ &= \frac{x}{x^2 + y^2} + f'(y) \end{aligned} \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{x}{x^2 + y^2}$ . Therefore equation (4) becomes

$$\frac{x}{x^2 + y^2} = \frac{x}{x^2 + y^2} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\frac{x^4}{4} - \arctan\left(\frac{x}{y}\right) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\frac{x^4}{4} - \arctan\left(\frac{x}{y}\right)$$

The solution becomes

$$y = -\frac{x}{\tan\left(\frac{x^4}{4} + c_1\right)}$$

### Summary

The solution(s) found are the following

$$y = -\frac{x}{\tan\left(\frac{x^4}{4} + c_1\right)} \quad (1)$$

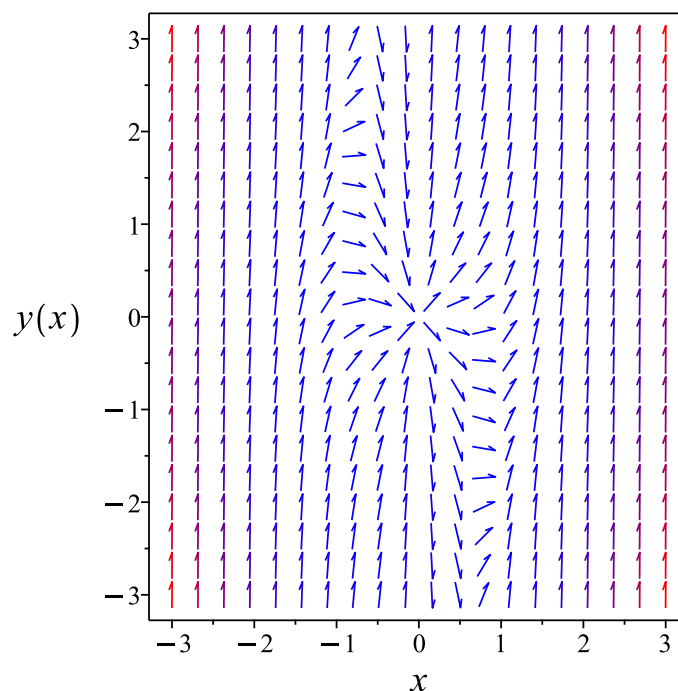


Figure 67: Slope field plot

### Verification of solutions

$$y = -\frac{x}{\tan\left(\frac{x^4}{4} + c_1\right)}$$

Verified OK.

### 3.6.3 Solving as riccati ode

In canonical form the ODE is

$$\begin{aligned}y' &= F(x, y) \\ &= \frac{x^5 + x^3y^2 + y}{x}\end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = x^4 + y^2x^2 + \frac{y}{x}$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that  $f_0(x) = x^4$ ,  $f_1(x) = \frac{1}{x}$  and  $f_2(x) = x^2$ . Let

$$\begin{aligned}y &= \frac{-u'}{f_2u} \\ &= \frac{-u'}{x^2u}\end{aligned}\tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for  $u(x)$  which is

$$f_2u''(x) - (f_2' + f_1f_2)u'(x) + f_2^2f_0u(x) = 0\tag{2}$$

But

$$\begin{aligned}f_2' &= 2x \\ f_1f_2 &= x \\ f_2^2f_0 &= x^8\end{aligned}$$

Substituting the above terms back in equation (2) gives

$$x^2u''(x) - 3xu'(x) + x^8u(x) = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 \sin\left(\frac{x^4}{4}\right) + c_2 \cos\left(\frac{x^4}{4}\right)$$

The above shows that

$$u'(x) = x^3 \left( c_1 \cos \left( \frac{x^4}{4} \right) - c_2 \sin \left( \frac{x^4}{4} \right) \right)$$

Using the above in (1) gives the solution

$$y = - \frac{x \left( c_1 \cos \left( \frac{x^4}{4} \right) - c_2 \sin \left( \frac{x^4}{4} \right) \right)}{c_1 \sin \left( \frac{x^4}{4} \right) + c_2 \cos \left( \frac{x^4}{4} \right)}$$

Dividing both numerator and denominator by  $c_1$  gives, after renaming the constant  $\frac{c_2}{c_1} = c_3$  the following solution

$$y = \frac{\left( -c_3 \cos \left( \frac{x^4}{4} \right) + \sin \left( \frac{x^4}{4} \right) \right) x}{c_3 \sin \left( \frac{x^4}{4} \right) + \cos \left( \frac{x^4}{4} \right)}$$

### Summary

The solution(s) found are the following

$$y = \frac{\left( -c_3 \cos \left( \frac{x^4}{4} \right) + \sin \left( \frac{x^4}{4} \right) \right) x}{c_3 \sin \left( \frac{x^4}{4} \right) + \cos \left( \frac{x^4}{4} \right)} \quad (1)$$

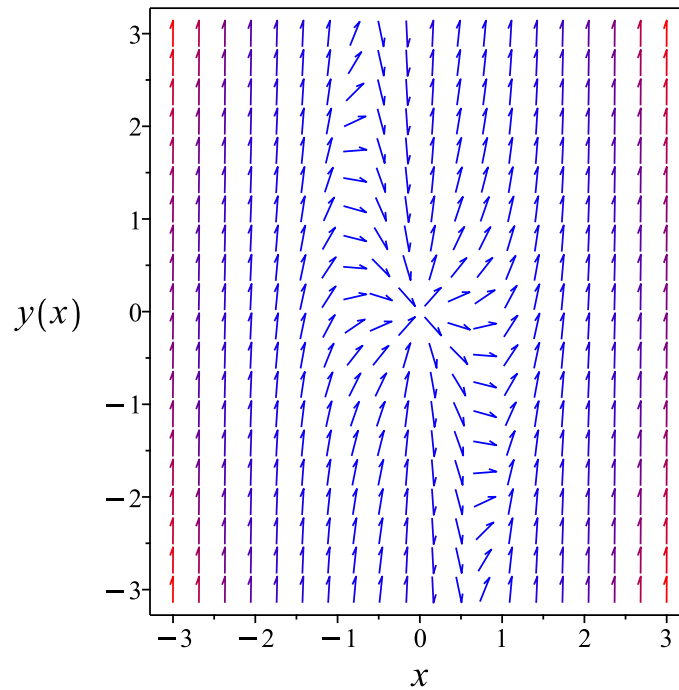


Figure 68: Slope field plot

Verification of solutions

$$y = \frac{\left(-c_3 \cos\left(\frac{x^4}{4}\right) + \sin\left(\frac{x^4}{4}\right)\right) x}{c_3 \sin\left(\frac{x^4}{4}\right) + \cos\left(\frac{x^4}{4}\right)}$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 14

```
dsolve(x*diff(y(x),x)=x^5+x^3*y(x)^2+y(x),y(x), singsol=all)
```

$$y(x) = \tan\left(\frac{x^4}{4} + c_1\right) x$$

✓ Solution by Mathematica

Time used: 0.212 (sec). Leaf size: 18

```
DSolve[x*y'[x]==x^5+x^3*y[x]^2+y[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow x \tan\left(\frac{x^4}{4} + c_1\right)$$

### 3.7 problem 4(d)

3.7.1	Solving as homogeneousTypeD2 ode . . . . .	295
3.7.2	Solving as first order ode lie symmetry calculated ode . . . . .	297
3.7.3	Solving as exact ode . . . . .	302

Internal problem ID [3106]

Internal file name [OUTPUT/2598\_Sunday\_June\_05\_2022\_03\_21\_51\_AM\_23105553/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 10, page 47

**Problem number:** 4(d).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"exactByInspection", "homogeneousTypeD2", "first\_order\_ode\_lie\_symmetry\_calculated"**

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$(y + x)y' - y = -x$$

#### 3.7.1 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$(u(x)x + x)(u'(x)x + u(x)) - u(x)x = -x$$

In canonical form the ODE is

$$\begin{aligned}u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{u^2 + 1}{x(u + 1)}\end{aligned}$$



Where  $f(x) = -\frac{1}{x}$  and  $g(u) = \frac{u^2+1}{u+1}$ . Integrating both sides gives

$$\frac{1}{\frac{u^2+1}{u+1}} du = -\frac{1}{x} dx$$

$$\int \frac{1}{\frac{u^2+1}{u+1}} du = \int -\frac{1}{x} dx$$

$$\frac{\ln(u^2 + 1)}{2} + \arctan(u) = -\ln(x) + c_2$$

The solution is

$$\frac{\ln(u(x)^2 + 1)}{2} + \arctan(u(x)) + \ln(x) - c_2 = 0$$

Replacing  $u(x)$  in the above solution by  $\frac{y}{x}$  results in the solution for  $y$  in implicit form

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} + \arctan\left(\frac{y}{x}\right) + \ln(x) - c_2 = 0$$

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} + \arctan\left(\frac{y}{x}\right) + \ln(x) - c_2 = 0$$

### Summary

The solution(s) found are the following

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} + \arctan\left(\frac{y}{x}\right) + \ln(x) - c_2 = 0 \quad (1)$$

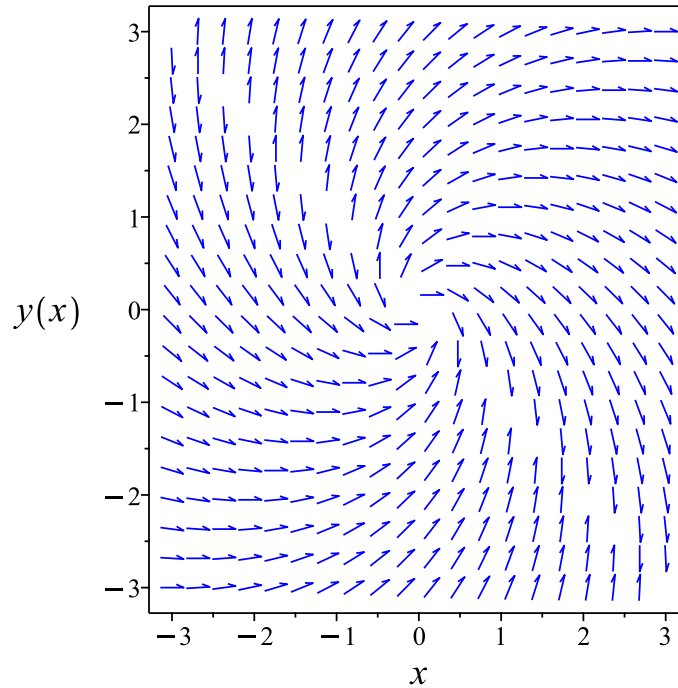


Figure 69: Slope field plot

### Verification of solutions

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} + \arctan\left(\frac{y}{x}\right) + \ln(x) - c_2 = 0$$

Verified OK.

### 3.7.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{y - x}{y + x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$b_2 + \frac{(y-x)(b_3 - a_2)}{y+x} - \frac{(y-x)^2 a_3}{(y+x)^2} - \left( -\frac{1}{y+x} - \frac{y-x}{(y+x)^2} \right) (xa_2 + ya_3 + a_1) \quad (5E)$$

$$- \left( \frac{1}{y+x} - \frac{y-x}{(y+x)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{x^2 a_2 - x^2 a_3 - x^2 b_2 - x^2 b_3 + 2xy a_2 + 2xy a_3 + 2xy b_2 - 2xy b_3 - y^2 a_2 + y^2 a_3 + y^2 b_2 + y^2 b_3 - 2xb_1 + 2ya_1}{(y+x)^2} = 0$$

Setting the numerator to zero gives

$$x^2 a_2 - x^2 a_3 - x^2 b_2 - x^2 b_3 + 2xy a_2 + 2xy a_3 + 2xy b_2 - 2xy b_3 - y^2 a_2 + y^2 a_3 + y^2 b_2 + y^2 b_3 - 2xb_1 + 2ya_1 = 0 \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$a_2 v_1^2 + 2a_2 v_1 v_2 - a_2 v_2^2 - a_3 v_1^2 + 2a_3 v_1 v_2 + a_3 v_2^2 - b_2 v_1^2 + 2b_2 v_1 v_2 + b_2 v_2^2 - b_3 v_1^2 - 2b_3 v_1 v_2 + b_3 v_2^2 + 2a_1 v_2 - 2b_1 v_1 = 0 \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$(a_2 - a_3 - b_2 - b_3) v_1^2 + (2a_2 + 2a_3 + 2b_2 - 2b_3) v_1 v_2 - 2b_1 v_1 + (-a_2 + a_3 + b_2 + b_3) v_2^2 + 2a_1 v_2 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 2a_1 &= 0 \\ -2b_1 &= 0 \\ -a_2 + a_3 + b_2 + b_3 &= 0 \\ a_2 - a_3 - b_2 - b_3 &= 0 \\ 2a_2 + 2a_3 + 2b_2 - 2b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= -b_2 \\ b_1 &= 0 \\ b_2 &= b_2 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= y - \left( \frac{y-x}{y+x} \right) (x) \\ &= \frac{x^2 + y^2}{y+x} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{x^2+y^2}{y+x}} dy \end{aligned}$$

Which results in

$$S = \frac{\ln(x^2 + y^2)}{2} + \arctan\left(\frac{y}{x}\right)$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y - x}{y + x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{-y + x}{x^2 + y^2} \\ S_y &= \frac{y + x}{x^2 + y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

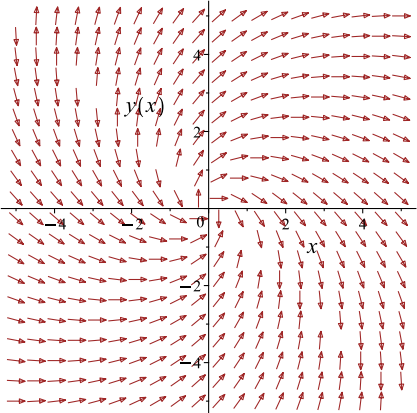
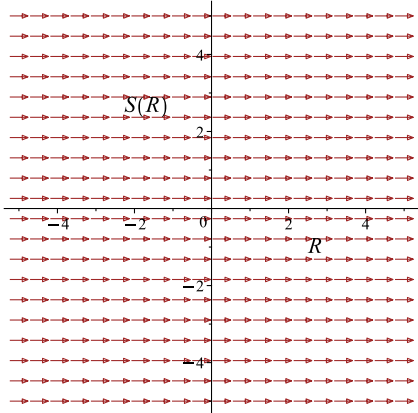
To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{\ln(y^2 + x^2)}{2} + \arctan\left(\frac{y}{x}\right) = c_1$$

Which simplifies to

$$\frac{\ln(y^2 + x^2)}{2} + \arctan\left(\frac{y}{x}\right) = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{y-x}{y+x}$ 	$R = x$ $S = \frac{\ln(x^2 + y^2)}{2} + \arctan\left(\frac{y}{x}\right)$	$\frac{dS}{dR} = 0$ 

### Summary

The solution(s) found are the following

$$\frac{\ln(y^2 + x^2)}{2} + \arctan\left(\frac{y}{x}\right) = c_1 \quad (1)$$

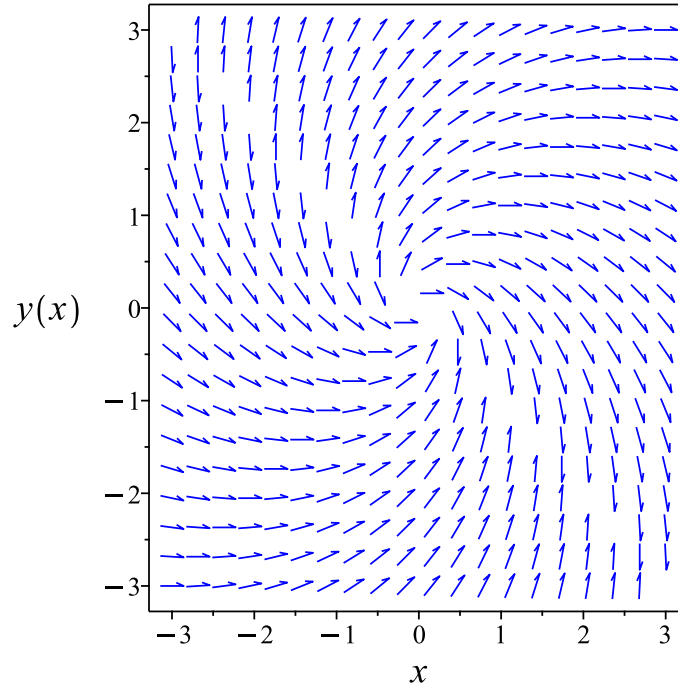


Figure 70: Slope field plot

### Verification of solutions

$$\frac{\ln(y^2 + x^2)}{2} + \arctan\left(\frac{y}{x}\right) = c_1$$

Verified OK.

### **3.7.3 Solving as exact ode**

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (A)$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (y + x) dy &= (y - x) dx \\ (-y + x) dx + (y + x) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -y + x \\ N(x, y) &= y + x \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-y + x) \\ &= -1 \end{aligned}$$



And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(y+x) \\ &= 1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. By inspection  $\frac{1}{y^2+x^2}$  is an integrating factor. Therefore by multiplying  $M = -y+x$  and  $N = y+x$  by this integrating factor the ode becomes exact. The new  $M, N$  are

$$\begin{aligned}M &= \frac{-y+x}{y^2+x^2} \\ N &= \frac{y+x}{y^2+x^2}\end{aligned}$$

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might

or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left(\frac{y+x}{x^2+y^2}\right) dy &= \left(-\frac{-y+x}{x^2+y^2}\right) dx \\ \left(\frac{-y+x}{x^2+y^2}\right) dx + \left(\frac{y+x}{x^2+y^2}\right) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= \frac{-y+x}{x^2+y^2} \\ N(x, y) &= \frac{y+x}{x^2+y^2} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{-y+x}{x^2+y^2}\right) \\ &= \frac{-x^2 - 2xy + y^2}{(x^2+y^2)^2} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{y+x}{x^2+y^2}\right) \\ &= \frac{-x^2 - 2xy + y^2}{(x^2+y^2)^2} \end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{-y+x}{x^2+y^2} dx \\ \phi &= \frac{\ln(x^2+y^2)}{2} - \arctan\left(\frac{x}{y}\right) + f(y)\end{aligned}\quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\begin{aligned}\frac{\partial \phi}{\partial y} &= \frac{y}{x^2+y^2} + \frac{x}{y^2\left(\frac{x^2}{y^2}+1\right)} + f'(y) \\ &= \frac{y+x}{x^2+y^2} + f'(y)\end{aligned}\quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{y+x}{x^2+y^2}$ . Therefore equation (4) becomes

$$\frac{y+x}{x^2+y^2} = \frac{y+x}{x^2+y^2} + f'(y)\quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = \frac{\ln(x^2+y^2)}{2} - \arctan\left(\frac{x}{y}\right) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \frac{\ln(x^2+y^2)}{2} - \arctan\left(\frac{x}{y}\right)$$

### Summary

The solution(s) found are the following

$$\frac{\ln(y^2 + x^2)}{2} - \arctan\left(\frac{x}{y}\right) = c_1 \quad (1)$$

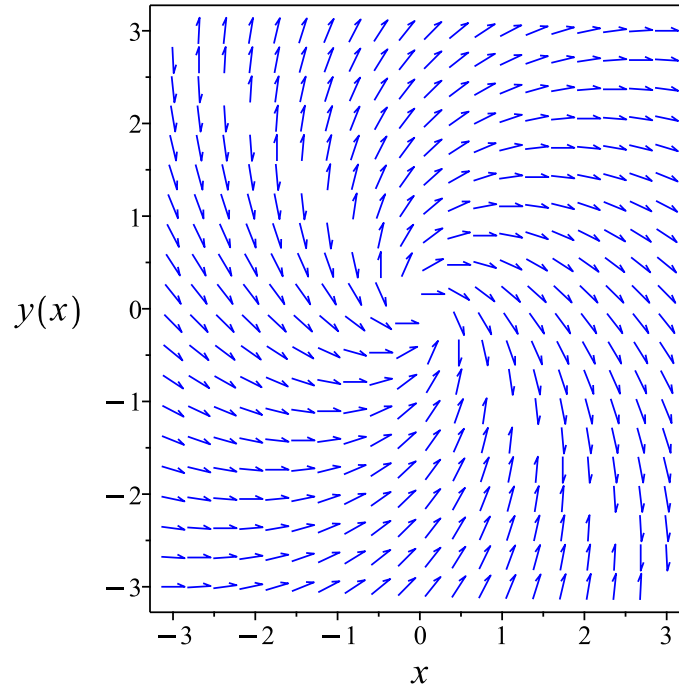


Figure 71: Slope field plot

### Verification of solutions

$$\frac{\ln(y^2 + x^2)}{2} - \arctan\left(\frac{x}{y}\right) = c_1$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 24

```
dsolve((y(x)+x)*diff(y(x),x)=(y(x)-x),y(x), singsol=all)
```

$$y(x) = \tan \left( \text{RootOf} \left( 2\_Z + \ln \left( \sec \left( \_Z \right)^2 \right) + 2 \ln (x) + 2c_1 \right) \right) x$$

### ✓ Solution by Mathematica

Time used: 0.033 (sec). Leaf size: 34

```
DSolve[(y[x]+x)*y'[x]==(y[x]-x),y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve} \left[ \arctan \left( \frac{y(x)}{x} \right) + \frac{1}{2} \log \left( \frac{y(x)^2}{x^2} + 1 \right) = -\log(x) + c_1, y(x) \right]$$

### 3.8 problem 4(e)

3.8.1 Solving as homogeneousTypeD2 ode . . . . .	309
3.8.2 Solving as riccati ode . . . . .	310

Internal problem ID [3107]

Internal file name [OUTPUT/2599\_Sunday\_June\_05\_2022\_03\_21\_53\_AM\_63742245/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, section 10, page 47

**Problem number:** 4(e).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**riccati**", "**homogeneousTypeD2**"

Maple gives the following as the ode type

```
[[_homogeneous, `class D`], _rational, _Riccati]
```

$$xy' - y - 9y^2 = x^2$$

#### 3.8.1 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$x(u'(x)x + u(x)) - u(x)x - 9u(x)^2x^2 = x^2$$

Integrating both sides gives

$$\int \frac{1}{9u^2 + 1} du = x + c_2$$
$$\frac{\arctan(3u)}{3} = x + c_2$$

Solving for  $u$  gives these solutions

$$u_1 = \frac{\tan(3x + 3c_2)}{3}$$

Therefore the solution  $y$  is

$$y = xu$$
$$= \frac{x \tan(3x + 3c_2)}{3}$$

### Summary

The solution(s) found are the following

$$y = \frac{x \tan(3x + 3c_2)}{3} \quad (1)$$

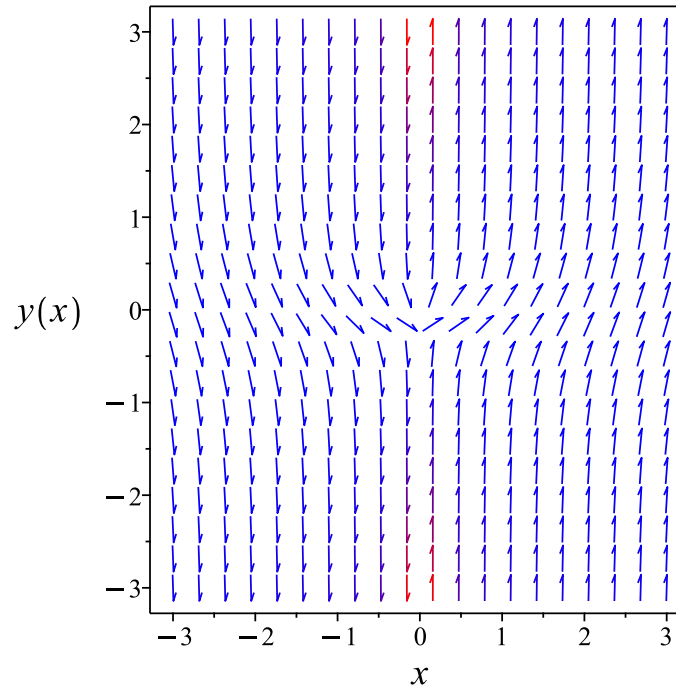


Figure 72: Slope field plot

### Verification of solutions

$$y = \frac{x \tan(3x + 3c_2)}{3}$$

Verified OK.

### **3.8.2 Solving as riccati ode**

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{x^2 + 9y^2 + y}{x} \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = x + \frac{9y^2}{x} + \frac{y}{x}$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that  $f_0(x) = x$ ,  $f_1(x) = \frac{1}{x}$  and  $f_2(x) = \frac{9}{x}$ . Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{\frac{9u}{x}} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for  $u(x)$  which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \tag{2}$$

But

$$\begin{aligned} f_2' &= -\frac{9}{x^2} \\ f_1 f_2 &= \frac{9}{x^2} \\ f_2^2 f_0 &= \frac{81}{x} \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$\frac{9u''(x)}{x} + \frac{81u(x)}{x} = 0$$

Solving the above ODE (this ode solved using Maple, not this program), gives

$$u(x) = c_1 \sin(3x) + c_2 \cos(3x)$$

The above shows that

$$u'(x) = 3c_1 \cos(3x) - 3c_2 \sin(3x)$$

Using the above in (1) gives the solution

$$y = -\frac{(3c_1 \cos(3x) - 3c_2 \sin(3x)) x}{9(c_1 \sin(3x) + c_2 \cos(3x))}$$

Dividing both numerator and denominator by  $c_1$  gives, after renaming the constant  $\frac{c_2}{c_1} = c_3$  the following solution



$$y = \frac{(-c_3 \cos(3x) + \sin(3x)) x}{3c_3 \sin(3x) + 3 \cos(3x)}$$

Summary

The solution(s) found are the following

$$y = \frac{(-c_3 \cos(3x) + \sin(3x)) x}{3c_3 \sin(3x) + 3 \cos(3x)} \tag{1}$$

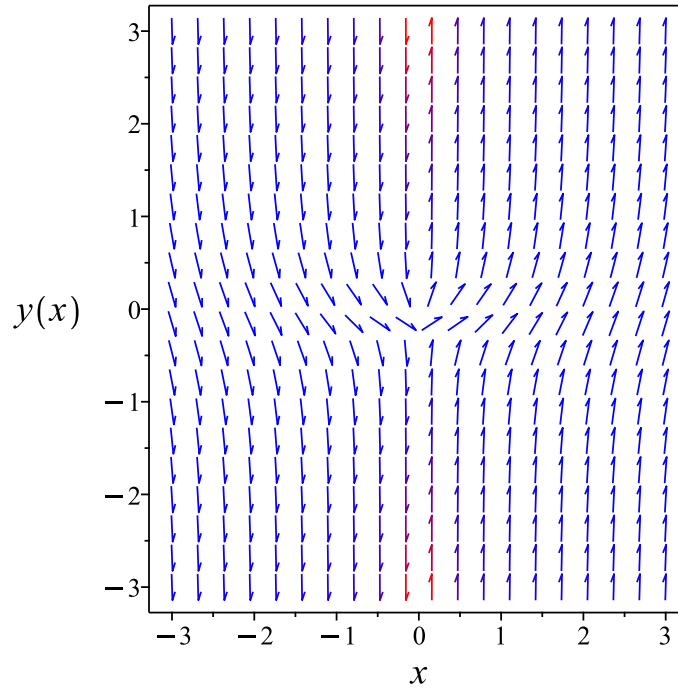


Figure 73: Slope field plot

Verification of solutions

$$y = \frac{(-c_3 \cos(3x) + \sin(3x)) x}{3c_3 \sin(3x) + 3 \cos(3x)}$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 15

```
dsolve(x*diff(y(x),x)=y(x)+x^2+9*y(x)^2,y(x), singsol=all)
```

$$y(x) = \frac{\tan(3x + 3c_1)x}{3}$$

### ✓ Solution by Mathematica

Time used: 0.285 (sec). Leaf size: 17

```
DSolve[x*y'[x]==y[x]+x^2+9*y[x]^2,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{1}{3}x \tan(3(x + c_1))$$

## 4 Chapter 2, section 11, page 49

4.1	problem 2(a)	315
4.2	problem 2(b)	328
4.3	problem 2(c)	341
4.4	problem 2(d)	353
4.5	problem 2(e)	366
4.6	problem 2(f)	379

## 4.1 problem 2(a)

4.1.1	Solving as linear ode . . . . .	315
4.1.2	Solving as first order ode lie symmetry lookup ode . . . . .	317
4.1.3	Solving as exact ode . . . . .	321
4.1.4	Maple step by step solution . . . . .	326

Internal problem ID [3108]

Internal file name [OUTPUT/2600\_Sunday\_June\_05\_2022\_03\_21\_56\_AM\_40981428/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 11, page 49

**Problem number:** 2(a).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

[\_linear]

$$xy' - 3y = x^4$$

### 4.1.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{3}{x}$$
$$q(x) = x^3$$

Hence the ode is

$$y' - \frac{3y}{x} = x^3$$

The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int -\frac{3}{x} dx} \\ &= \frac{1}{x^3}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu)(x^3) \\ \frac{d}{dx}\left(\frac{y}{x^3}\right) &= \left(\frac{1}{x^3}\right)(x^3) \\ d\left(\frac{y}{x^3}\right) &= dx\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y}{x^3} &= \int dx \\ \frac{y}{x^3} &= x + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = \frac{1}{x^3}$  results in

$$y = c_1 x^3 + x^4$$

which simplifies to

$$y = x^3(x + c_1)$$

### Summary

The solution(s) found are the following

$$y = x^3(x + c_1) \tag{1}$$

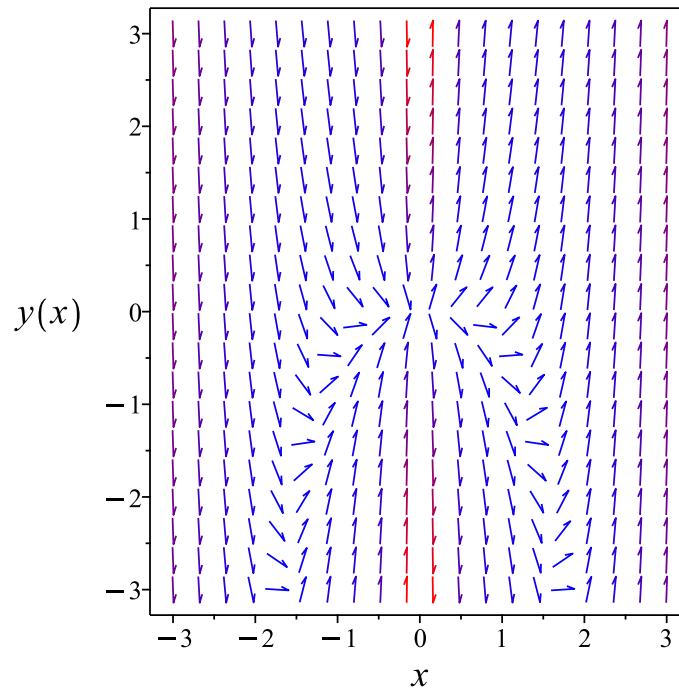


Figure 74: Slope field plot

Verification of solutions

$$y = x^3(x + c_1)$$

Verified OK.

#### 4.1.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{x^4 + 3y}{x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 31: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= x^3\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{x^3} dy \end{aligned}$$

Which results in

$$S = \frac{y}{x^3}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x^4 + 3y}{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -\frac{3y}{x^4} \\ S_y &= \frac{1}{x^3} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 1 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 1$$



The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = R + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{y}{x^3} = x + c_1$$

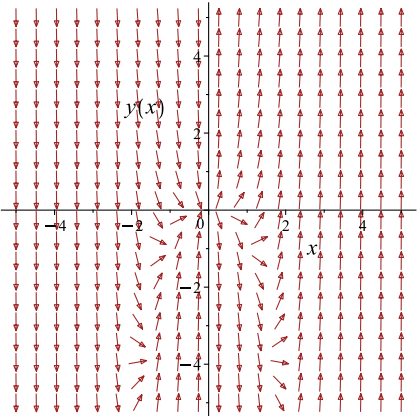
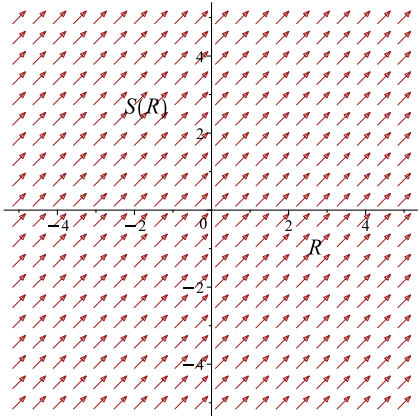
Which simplifies to

$$\frac{y}{x^3} = x + c_1$$

Which gives

$$y = x^3(x + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{x^4 + 3y}{x}$ 	$R = x$ $S = \frac{y}{x^3}$	$\frac{dS}{dR} = 1$ 

### Summary

The solution(s) found are the following

$$y = x^3(x + c_1) \tag{1}$$

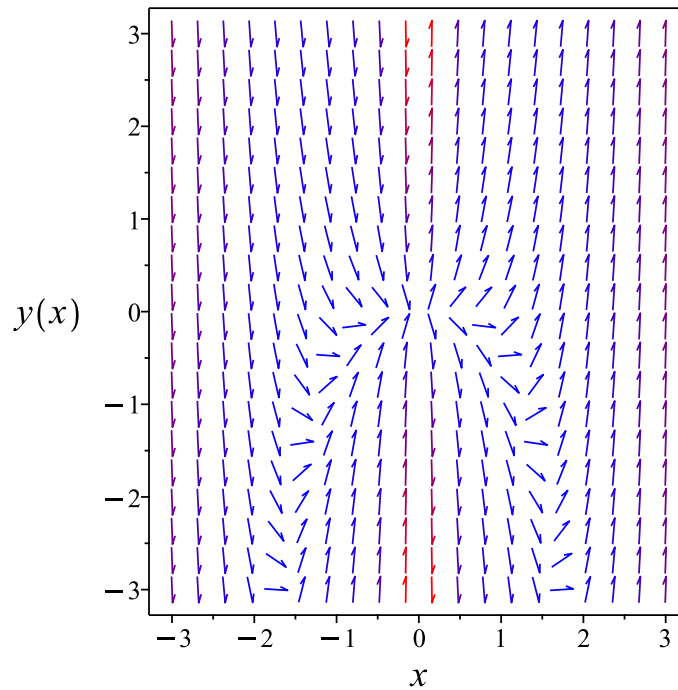


Figure 75: Slope field plot

Verification of solutions

$$y = x^3(x + c_1)$$

Verified OK.

### 4.1.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(x) dy &= (x^4 + 3y) dx \\ (-x^4 - 3y) dx + (x) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -x^4 - 3y \\ N(x, y) &= x\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x^4 - 3y) \\ &= -3\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x) \\ &= 1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{x} ((-3) - (1)) \\ &= -\frac{4}{x} \end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned} \mu &= e^{\int A \, dx} \\ &= e^{\int -\frac{4}{x} \, dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-4 \ln(x)} \\ &= \frac{1}{x^4} \end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{1}{x^4} (-x^4 - 3y) \\ &= \frac{-x^4 - 3y}{x^4} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \frac{1}{x^4} (x) \\ &= \frac{1}{x^3} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left( \frac{-x^4 - 3y}{x^4} \right) + \left( \frac{1}{x^3} \right) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{-x^4 - 3y}{x^4} dx \\ \phi &= -x + \frac{y}{x^3} + f(y) \end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = \frac{1}{x^3} + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{1}{x^3}$ . Therefore equation (4) becomes

$$\frac{1}{x^3} = \frac{1}{x^3} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -x + \frac{y}{x^3} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -x + \frac{y}{x^3}$$

The solution becomes

$$y = x^3(x + c_1)$$

Summary

The solution(s) found are the following

$$y = x^3(x + c_1) \tag{1}$$

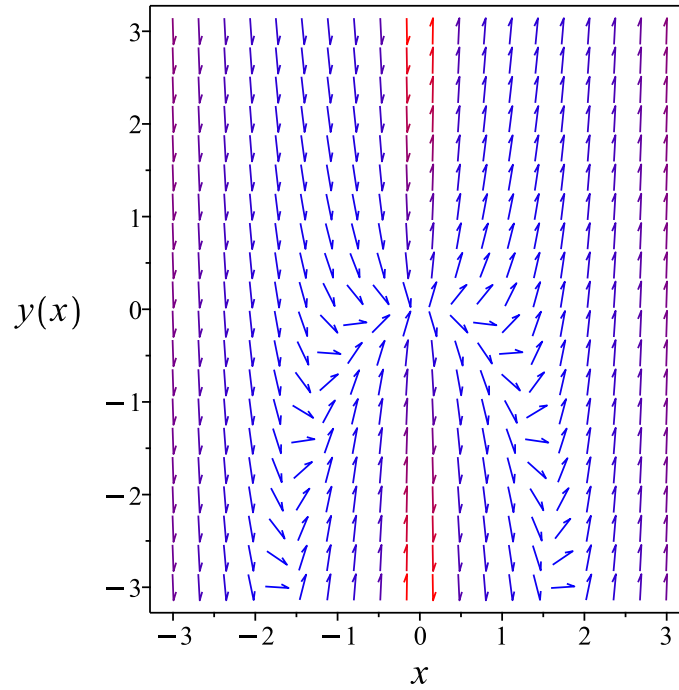


Figure 76: Slope field plot

Verification of solutions

$$y = x^3(x + c_1)$$

Verified OK.

#### 4.1.4 Maple step by step solution

Let's solve

$$xy' - 3y = x^4$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = \frac{3y}{x} + x^3$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' - \frac{3y}{x} = x^3$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x) \left( y' - \frac{3y}{x} \right) = \mu(x) x^3$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$

$$\mu(x) \left( y' - \frac{3y}{x} \right) = \mu'(x)y + \mu(x)y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = -\frac{3\mu(x)}{x}$$

- Solve to find the integrating factor

$$\mu(x) = \frac{1}{x^3}$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x)y) \right) dx = \int \mu(x) x^3 dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int \mu(x) x^3 dx + c_1$$

- Solve for  $y$

$$y = \frac{\int \mu(x) x^3 dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = \frac{1}{x^3}$

$$y = x^3 \left( \int 1 dx + c_1 \right)$$

- Evaluate the integrals on the rhs

$$y = x^3(x + c_1)$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 11

```
dsolve(x*diff(y(x),x)-3*y(x)=x^4,y(x), singsol=all)
```

$$y(x) = (c_1 + x) x^3$$

### ✓ Solution by Mathematica

Time used: 0.026 (sec). Leaf size: 13

```
DSolve[x*y'[x]-3*y[x]==x^4,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow x^3(x + c_1)$$



## 4.2 problem 2(b)

4.2.1	Solving as linear ode . . . . .	328
4.2.2	Solving as first order ode lie symmetry lookup ode . . . . .	330
4.2.3	Solving as exact ode . . . . .	334
4.2.4	Maple step by step solution . . . . .	338

Internal problem ID [3109]

Internal file name [OUTPUT/2601\_Sunday\_June\_05\_2022\_03\_21\_58\_AM\_39847362/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 11, page 49

**Problem number:** 2(b).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$y' + y = \frac{1}{e^{2x} + 1}$$

### 4.2.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = 1$$
$$q(x) = \frac{1}{e^{2x} + 1}$$

Hence the ode is

$$y' + y = \frac{1}{e^{2x} + 1}$$

The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int 1 dx} \\ &= e^x\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) \left( \frac{1}{e^{2x} + 1} \right) \\ \frac{d}{dx}(e^x y) &= (e^x) \left( \frac{1}{e^{2x} + 1} \right) \\ d(e^x y) &= \left( \frac{e^x}{e^{2x} + 1} \right) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}e^x y &= \int \frac{e^x}{e^{2x} + 1} dx \\ e^x y &= \arctan(e^x) + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = e^x$  results in

$$y = \arctan(e^x) e^{-x} + c_1 e^{-x}$$

which simplifies to

$$y = e^{-x}(\arctan(e^x) + c_1)$$

### Summary

The solution(s) found are the following

$$y = e^{-x}(\arctan(e^x) + c_1) \tag{1}$$

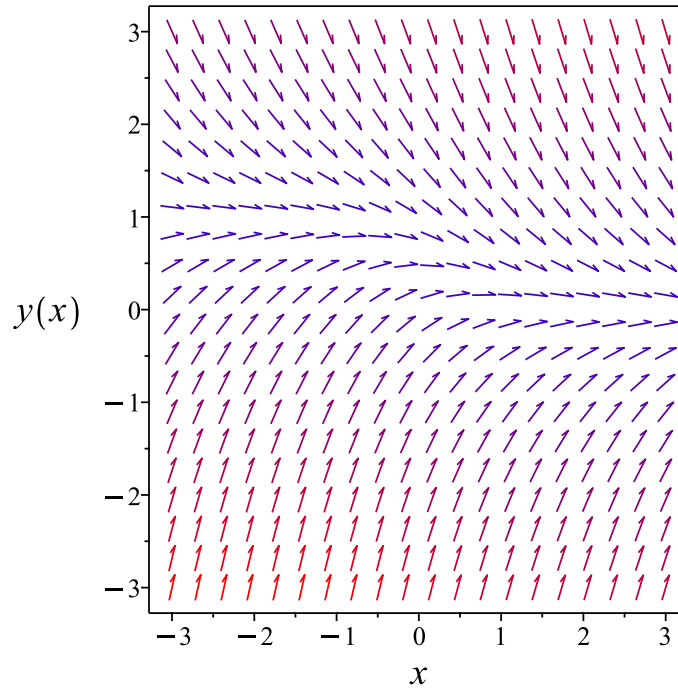


Figure 77: Slope field plot

Verification of solutions

$$y = e^{-x}(\arctan(e^x) + c_1)$$

Verified OK.

#### 4.2.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{e^{2x}y + y - 1}{e^{2x} + 1}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 34: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= e^{-x}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-x}} dy \end{aligned}$$

Which results in

$$S = e^x y$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{e^{2x}y + y - 1}{e^{2x} + 1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= e^x y \\ S_y &= e^x \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{e^x}{e^{2x} + 1} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{e^R}{e^{2R} + 1}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = \arctan(e^R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$e^x y = \arctan(e^x) + c_1$$

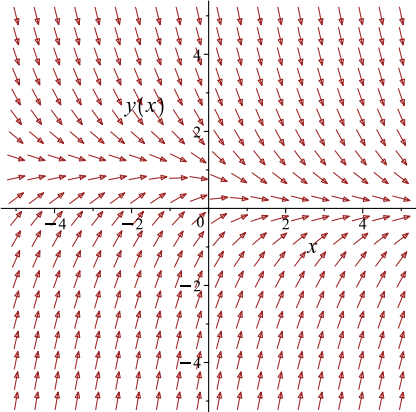
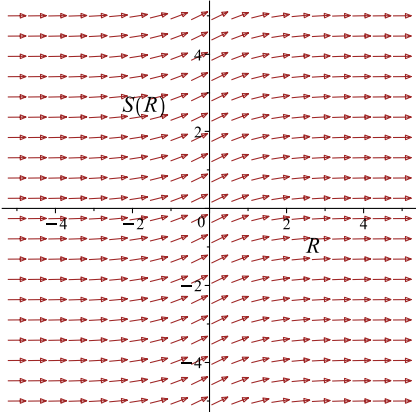
Which simplifies to

$$e^x y = \arctan(e^x) + c_1$$

Which gives

$$y = e^{-x}(\arctan(e^x) + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{e^{2x}y+y-1}{e^{2x}+1}$ 	$R = x$ $S = e^x y$	$\frac{dS}{dR} = \frac{e^R}{e^{2R}+1}$ 

### Summary

The solution(s) found are the following

$$y = e^{-x}(\arctan(e^x) + c_1) \quad (1)$$

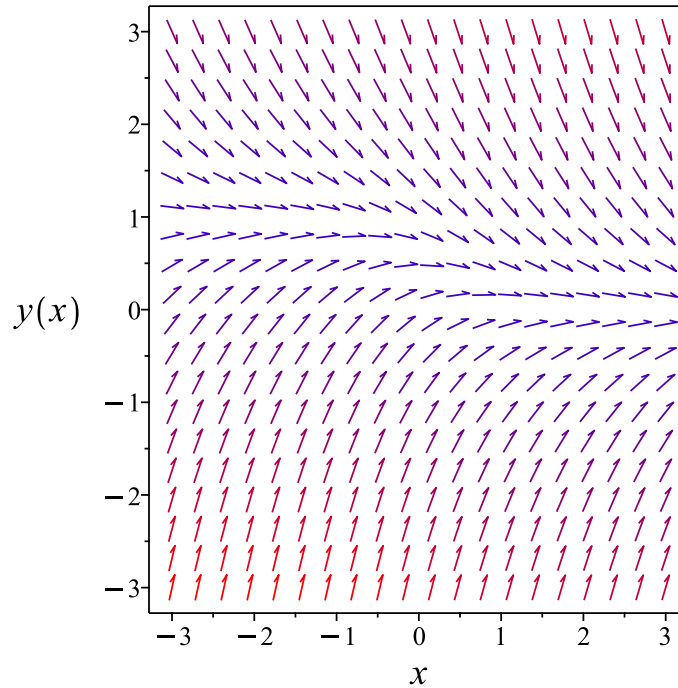


Figure 78: Slope field plot

Verification of solutions

$$y = e^{-x}(\arctan(e^x) + c_1)$$

Verified OK.

### 4.2.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dy &= \left(-y + \frac{1}{e^{2x} + 1}\right) dx \\ \left(y - \frac{1}{e^{2x} + 1}\right) dx + dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= y - \frac{1}{e^{2x} + 1} \\ N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(y - \frac{1}{e^{2x} + 1}\right) \\ &= 1\end{aligned}$$



And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((1) - (0)) \\ &= 1\end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int 1 dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^x \\ &= e^x\end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned}\bar{M} &= \mu M \\ &= e^x \left( y - \frac{1}{e^{2x} + 1} \right) \\ &= \frac{e^x(e^{2x}y + y - 1)}{e^{2x} + 1}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^x(1) \\ &= e^x\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left( \frac{e^x(e^{2x}y + y - 1)}{e^{2x} + 1} \right) + (e^x) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \overline{M} dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int \frac{e^x(e^{2x}y + y - 1)}{e^{2x} + 1} dx$$

$$\phi = e^x y - \arctan(e^x) + f(y) \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = e^x + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = e^x$ . Therefore equation (4) becomes

$$e^x = e^x + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = e^x y - \arctan(e^x) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = e^x y - \arctan(e^x)$$

### Summary

The solution(s) found are the following

$$e^x y - \arctan(e^x) = c_1 \quad (1)$$

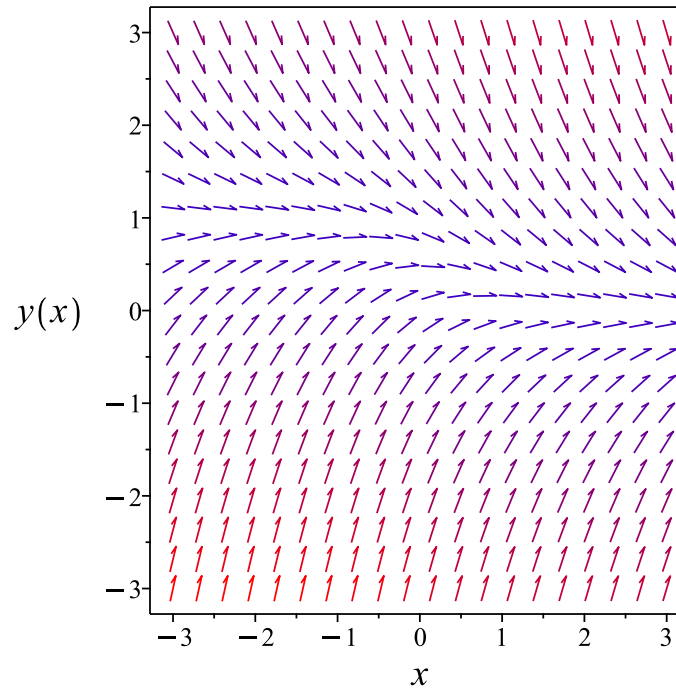


Figure 79: Slope field plot

### Verification of solutions

$$e^x y - \arctan(e^x) = c_1$$

Verified OK.

### 4.2.4 Maple step by step solution

Let's solve

$$y' + y = \frac{1}{e^{2x} + 1}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -y + \frac{1}{e^{2x} + 1}$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + y = \frac{1}{e^{2x+1}}$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x)(y' + y) = \frac{\mu(x)}{e^{2x+1}}$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$

$$\mu(x)(y' + y) = \mu'(x)y + \mu(x)y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = \mu(x)$$

- Solve to find the integrating factor

$$\mu(x) = e^x$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x)y) \right) dx = \int \frac{\mu(x)}{e^{2x+1}} dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int \frac{\mu(x)}{e^{2x+1}} dx + c_1$$

- Solve for  $y$

$$y = \frac{\int \frac{\mu(x)}{e^{2x+1}} dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = e^x$

$$y = \frac{\int \frac{e^x}{e^{2x+1}} dx + c_1}{e^x}$$

- Evaluate the integrals on the rhs

$$y = \frac{\arctan(e^x) + c_1}{e^x}$$

- Simplify

$$y = e^{-x}(\arctan(e^x) + c_1)$$

### Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(diff(y(x),x)+y(x)=1/(1+exp(2*x)),y(x), singsol=all)
```

$$y(x) = (\arctan(e^x) + c_1) e^{-x}$$

✓ Solution by Mathematica

Time used: 0.081 (sec). Leaf size: 18

```
DSolve[y'[x]+y[x]==1/(1+Exp[2*x]),y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x}(\arctan(e^x) + c_1)$$

### 4.3 problem 2(c)

4.3.1	Solving as linear ode . . . . .	341
4.3.2	Solving as first order ode lie symmetry lookup ode . . . . .	343
4.3.3	Solving as exact ode . . . . .	347
4.3.4	Maple step by step solution . . . . .	351

Internal problem ID [3110]

Internal file name [OUTPUT/2602\_Sunday\_June\_05\_2022\_03\_22\_00\_AM\_85419095/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 11, page 49

**Problem number:** 2(c).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"exact", "linear", "first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$2yx + (x^2 + 1)y' = \cot(x)$$

#### 4.3.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{2x}{x^2 + 1}$$
$$q(x) = \frac{\cot(x)}{x^2 + 1}$$

Hence the ode is

$$y' + \frac{2xy}{x^2 + 1} = \frac{\cot(x)}{x^2 + 1}$$

The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int \frac{-2x}{x^2+1} dx} \\ &= x^2 + 1\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) \left( \frac{\cot(x)}{x^2 + 1} \right) \\ \frac{d}{dx}((x^2 + 1) y) &= (x^2 + 1) \left( \frac{\cot(x)}{x^2 + 1} \right) \\ d((x^2 + 1) y) &= \cot(x) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}(x^2 + 1) y &= \int \cot(x) dx \\ (x^2 + 1) y &= \ln(\sin(x)) + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = x^2 + 1$  results in

$$y = \frac{\ln(\sin(x))}{x^2 + 1} + \frac{c_1}{x^2 + 1}$$

which simplifies to

$$y = \frac{\ln(\sin(x)) + c_1}{x^2 + 1}$$

Summary

The solution(s) found are the following

$$y = \frac{\ln(\sin(x)) + c_1}{x^2 + 1} \tag{1}$$

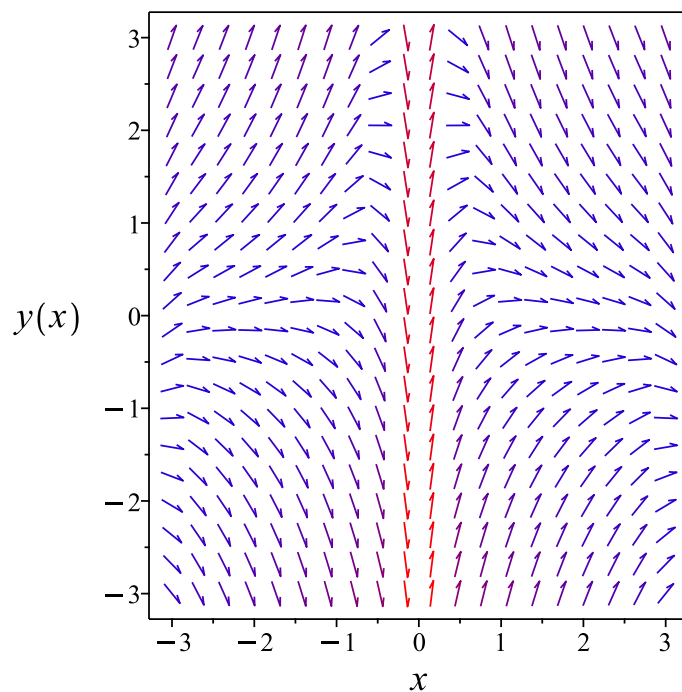


Figure 80: Slope field plot

Verification of solutions

$$y = \frac{\ln(\sin(x)) + c_1}{x^2 + 1}$$

Verified OK.

### 4.3.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-2xy + \cot(x)}{x^2 + 1}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$



Table 37: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{x^2 + 1}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{x^2+1}} dy \end{aligned}$$

Which results in

$$S = (x^2 + 1) y$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{-2xy + \cot(x)}{x^2 + 1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= 2xy \\ S_y &= x^2 + 1 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \cot(x) \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \cot(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = \ln(\sin(R)) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$(x^2 + 1)y = \ln(\sin(x)) + c_1$$

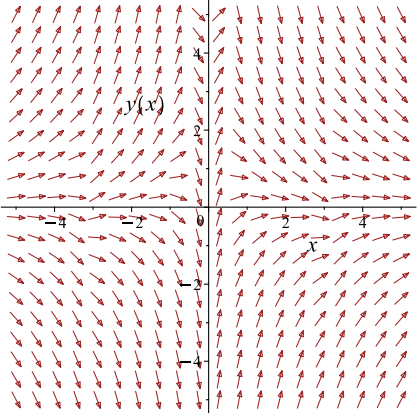
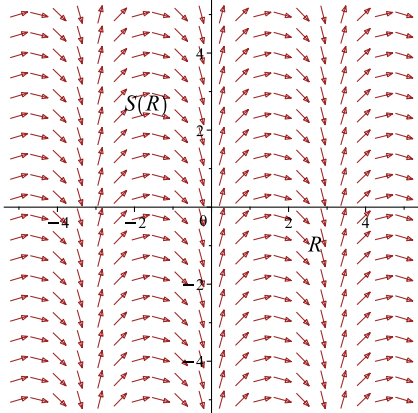
Which simplifies to

$$(x^2 + 1)y = \ln(\sin(x)) + c_1$$

Which gives

$$y = \frac{\ln(\sin(x)) + c_1}{x^2 + 1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{-2xy + \cot(x)}{x^2 + 1}$ 	$R = x$ $S = (x^2 + 1)y$	$\frac{dS}{dR} = \cot(R)$ 

### Summary

The solution(s) found are the following

$$y = \frac{\ln(\sin(x)) + c_1}{x^2 + 1} \quad (1)$$

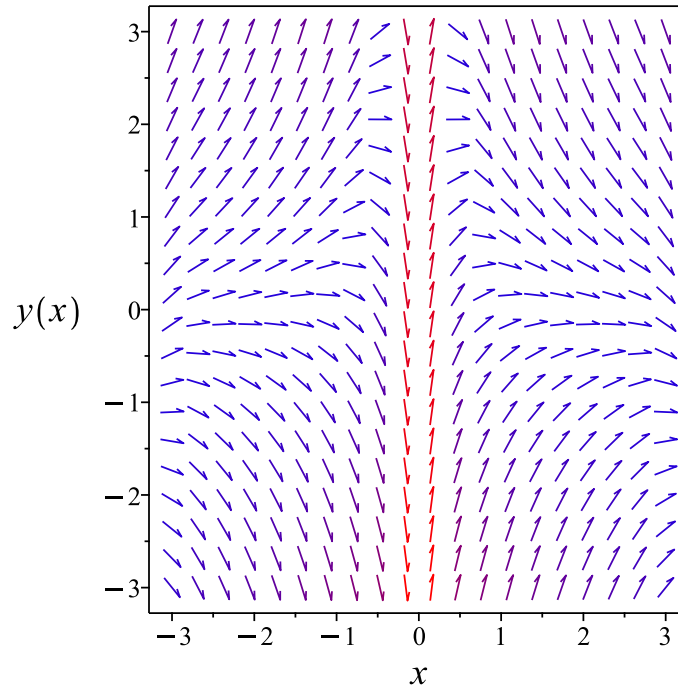


Figure 81: Slope field plot

Verification of solutions

$$y = \frac{\ln(\sin(x)) + c_1}{x^2 + 1}$$

Verified OK.

### 4.3.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(x^2 + 1) dy &= (-2xy + \cot(x)) dx \\ (2xy - \cot(x)) dx + (x^2 + 1) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 2xy - \cot(x) \\ N(x, y) &= x^2 + 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2xy - \cot(x)) \\ &= 2x\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x^2 + 1) \\ &= 2x\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 2xy - \cot(x) dx \\ \phi &= x^2y - \ln(\sin(x)) + f(y) \end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = x^2 + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = x^2 + 1$ . Therefore equation (4) becomes

$$x^2 + 1 = x^2 + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 1$$

Integrating the above w.r.t  $y$  gives

$$\begin{aligned} \int f'(y) dy &= \int (1) dy \\ f(y) &= y + c_1 \end{aligned}$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = x^2y - \ln(\sin(x)) + y + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = x^2y - \ln(\sin(x)) + y$$

The solution becomes

$$y = \frac{\ln(\sin(x)) + c_1}{x^2 + 1}$$

### Summary

The solution(s) found are the following

$$y = \frac{\ln(\sin(x)) + c_1}{x^2 + 1} \tag{1}$$

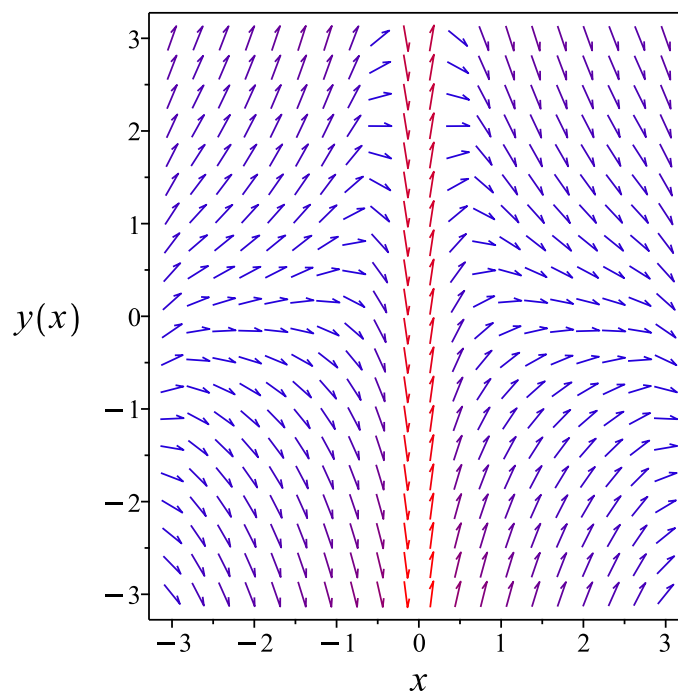


Figure 82: Slope field plot

### Verification of solutions

$$y = \frac{\ln(\sin(x)) + c_1}{x^2 + 1}$$

Verified OK.

#### 4.3.4 Maple step by step solution

Let's solve

$$2yx + (x^2 + 1)y' = \cot(x)$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{2xy}{x^2+1} + \frac{\cot(x)}{x^2+1}$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{2xy}{x^2+1} = \frac{\cot(x)}{x^2+1}$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x) \left( y' + \frac{2xy}{x^2+1} \right) = \frac{\mu(x)\cot(x)}{x^2+1}$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$

$$\mu(x) \left( y' + \frac{2xy}{x^2+1} \right) = \mu'(x)y + \mu(x)y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = \frac{2\mu(x)x}{x^2+1}$$

- Solve to find the integrating factor

$$\mu(x) = x^2 + 1$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x)y) \right) dx = \int \frac{\mu(x)\cot(x)}{x^2+1} dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int \frac{\mu(x)\cot(x)}{x^2+1} dx + c_1$$

- Solve for  $y$

$$y = \frac{\int \frac{\mu(x)\cot(x)}{x^2+1} dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = x^2 + 1$

$$y = \frac{\int \cot(x) dx + c_1}{x^2+1}$$

- Evaluate the integrals on the rhs

$$y = \frac{\ln(\sin(x)) + c_1}{x^2+1}$$



## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve((1+x^2)*diff(y(x),x)+2*x*y(x)=cot(x),y(x), singsol=all)
```

$$y(x) = \frac{\ln(\sin(x)) + c_1}{x^2 + 1}$$

### ✓ Solution by Mathematica

Time used: 0.056 (sec). Leaf size: 19

```
DSolve[(1+x^2)*y'[x]+2*x*y[x]==Cot[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{\log(\sin(x)) + c_1}{x^2 + 1}$$

## 4.4 problem 2(d)

4.4.1	Solving as linear ode . . . . .	353
4.4.2	Solving as first order ode lie symmetry lookup ode . . . . .	355
4.4.3	Solving as exact ode . . . . .	359
4.4.4	Maple step by step solution . . . . .	363

Internal problem ID [3111]

Internal file name [OUTPUT/2603\_Sunday\_June\_05\_2022\_03\_22\_02\_AM\_53367141/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 11, page 49

**Problem number:** 2(d).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

```
[[_linear, `class A`]]
```

$$y' + y = 2x e^{-x} + x^2$$

### 4.4.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = 1$$

$$q(x) = x(2e^{-x} + x)$$

Hence the ode is

$$y' + y = x(2e^{-x} + x)$$

The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int 1 dx} \\ &= e^x\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) (x(2e^{-x} + x)) \\ \frac{d}{dx}(e^x y) &= (e^x) (x(2e^{-x} + x)) \\ d(e^x y) &= (x(xe^x + 2)) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}e^x y &= \int x(xe^x + 2) dx \\ e^x y &= x^2 e^x - 2x e^x + 2e^x + x^2 + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = e^x$  results in

$$y = e^{-x}(x^2 e^x - 2x e^x + 2e^x + x^2) + c_1 e^{-x}$$

which simplifies to

$$y = (x^2 + c_1) e^{-x} + x^2 - 2x + 2$$

### Summary

The solution(s) found are the following

$$y = (x^2 + c_1) e^{-x} + x^2 - 2x + 2 \tag{1}$$

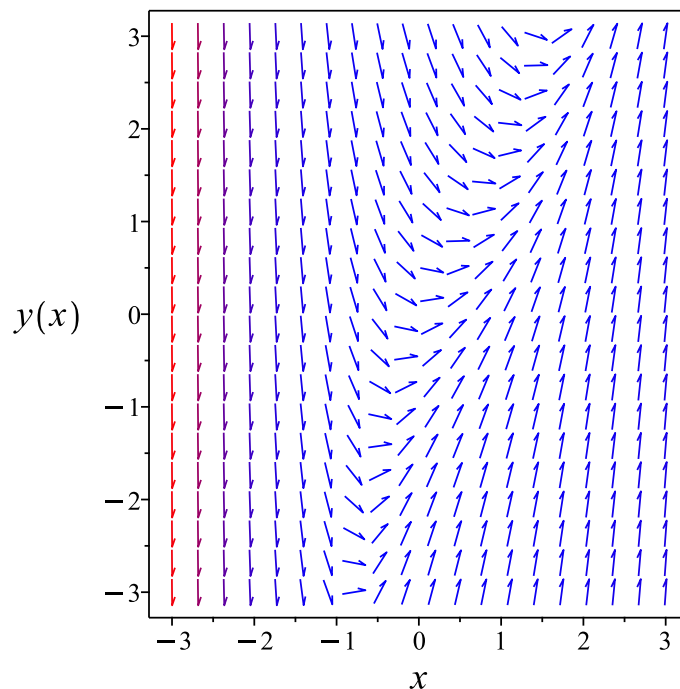


Figure 83: Slope field plot

### Verification of solutions

$$y = (x^2 + c_1) e^{-x} + x^2 - 2x + 2$$

Verified OK.

### **4.4.2 Solving as first order ode lie symmetry lookup ode**

Writing the ode as

$$\begin{aligned}y' &= -y + 2x e^{-x} + x^2 \\y' &= \omega(x, y)\end{aligned}$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 40: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= e^{-x}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-x}} dy \end{aligned}$$

Which results in

$$S = e^x y$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -y + 2x e^{-x} + x^2$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= e^x y \\ S_y &= e^x \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = x(x e^x + 2) \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R(R e^R + 2)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = e^R R^2 - 2R e^R + 2e^R + R^2 + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$e^x y = x^2 e^x - 2x e^x + 2e^x + x^2 + c_1$$

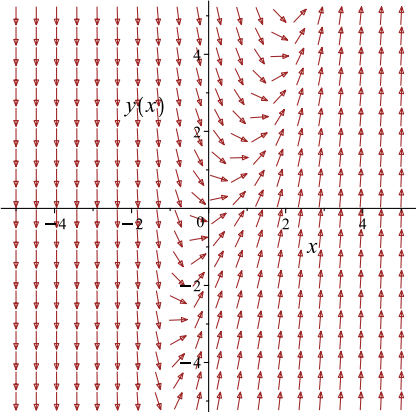
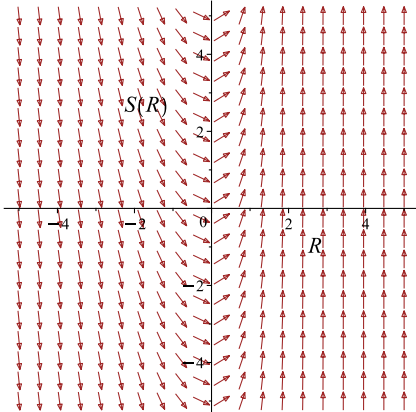
Which simplifies to

$$e^x y = x^2 e^x - 2x e^x + 2e^x + x^2 + c_1$$

Which gives

$$y = (x^2 e^x - 2x e^x + 2e^x + x^2 + c_1) e^{-x}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -y + 2x e^{-x} + x^2$ 	$R = x$ $S = e^x y$	$\frac{dS}{dR} = R(R e^R + 2)$ 

### Summary

The solution(s) found are the following

$$y = (x^2 e^x - 2x e^x + 2e^x + x^2 + c_1) e^{-x} \quad (1)$$

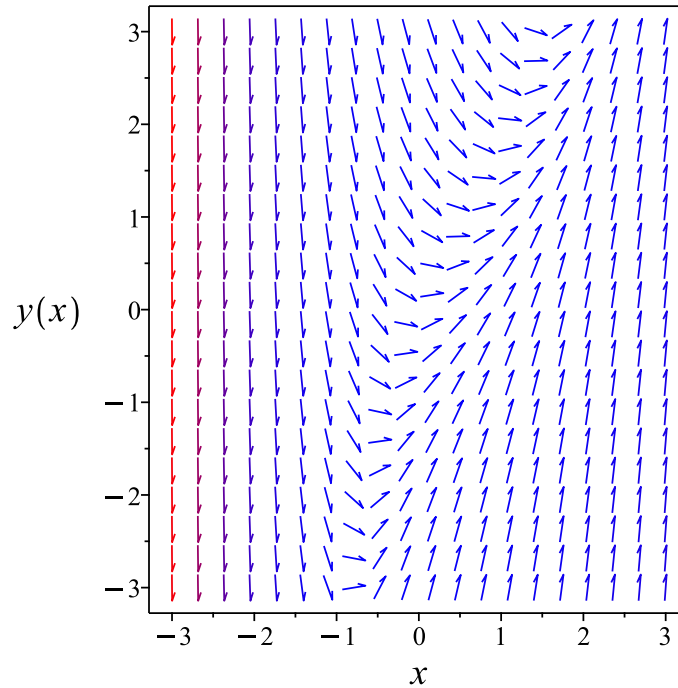


Figure 84: Slope field plot

Verification of solutions

$$y = (x^2 e^x - 2x e^x + 2e^x + x^2 + c_1) e^{-x}$$

Verified OK.

#### 4.4.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$



Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dy &= (-y + 2x e^{-x} + x^2) dx \\ (y - 2x e^{-x} - x^2) dx + dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= y - 2x e^{-x} - x^2 \\ N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (y - 2x e^{-x} - x^2) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (1) \\ &= 0\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((1) - (0)) \\ &= 1 \end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned} \mu &= e^{\int A dx} \\ &= e^{\int 1 dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^x \\ &= e^x \end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^x(y - 2x e^{-x} - x^2) \\ &= (-x^2 + y) e^x - 2x \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^x(1) \\ &= e^x \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ ((-x^2 + y) e^x - 2x) + (e^x) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int (-x^2 + y) e^x - 2x dx \\ \phi &= (-x^2 + 2x + y - 2) e^x - x^2 + f(y)\end{aligned}\quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = e^x + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = e^x$ . Therefore equation (4) becomes

$$e^x = e^x + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = (-x^2 + 2x + y - 2) e^x - x^2 + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = (-x^2 + 2x + y - 2) e^x - x^2$$

The solution becomes

$$y = (x^2 e^x - 2x e^x + 2 e^x + x^2 + c_1) e^{-x}$$

### Summary

The solution(s) found are the following

$$y = (x^2 e^x - 2x e^x + 2 e^x + x^2 + c_1) e^{-x} \quad (1)$$

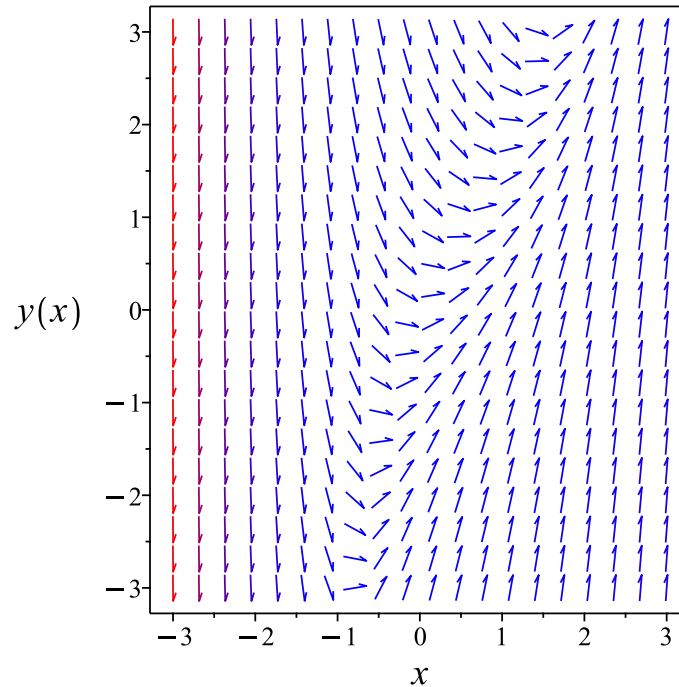


Figure 85: Slope field plot

### Verification of solutions

$$y = (x^2 e^x - 2x e^x + 2 e^x + x^2 + c_1) e^{-x}$$

Verified OK.

#### 4.4.4 Maple step by step solution

Let's solve

$$y' + y = 2x e^{-x} + x^2$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -y + 2x e^{-x} + x^2$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + y = 2x e^{-x} + x^2$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x) (y' + y) = \mu(x) (2x e^{-x} + x^2)$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x) y)$

$$\mu(x) (y' + y) = \mu'(x) y + \mu(x) y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = \mu(x)$$

- Solve to find the integrating factor

$$\mu(x) = e^x$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x) y) \right) dx = \int \mu(x) (2x e^{-x} + x^2) dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x) y = \int \mu(x) (2x e^{-x} + x^2) dx + c_1$$

- Solve for  $y$

$$y = \frac{\int \mu(x) (2x e^{-x} + x^2) dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = e^x$

$$y = \frac{\int e^x (2x e^{-x} + x^2) dx + c_1}{e^x}$$

- Evaluate the integrals on the rhs

$$y = \frac{x^2 e^x - 2x e^x + 2e^x + x^2 + c_1}{e^x}$$

- Simplify

$$y = (x^2 + c_1) e^{-x} + x^2 - 2x + 2$$

### Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 22

```
dsolve(diff(y(x),x)+y(x)=2*x*exp(-x)+x^2,y(x), singsol=all)
```

$$y(x) = (x^2 + c_1) e^{-x} + x^2 - 2x + 2$$

✓ Solution by Mathematica

Time used: 0.1 (sec). Leaf size: 29

```
DSolve[y'[x]+y[x]==2*x*Exp[-x]+x^2,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x}(x^2 + e^x(x^2 - 2x + 2) + c_1)$$

## 4.5 problem 2(e)

4.5.1	Solving as linear ode . . . . .	366
4.5.2	Solving as first order ode lie symmetry lookup ode . . . . .	368
4.5.3	Solving as exact ode . . . . .	372
4.5.4	Maple step by step solution . . . . .	376

Internal problem ID [3112]

Internal file name [OUTPUT/2604\_Sunday\_June\_05\_2022\_03\_22\_07\_AM\_72034493/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 11, page 49

**Problem number:** 2(e).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"linear", "exactWithIntegrationFactor", "first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

[\_linear]

$$y' + y \cot(x) = 2 \csc(x) x$$

### 4.5.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \cot(x)$$

$$q(x) = 2 \csc(x) x$$

Hence the ode is

$$y' + y \cot(x) = 2 \csc(x) x$$

The integrating factor  $\mu$  is

$$\begin{aligned} \mu &= e^{\int \cot(x) dx} \\ &= \sin(x) \end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) (2 \csc(x) x) \\ \frac{d}{dx}(\sin(x) y) &= (\sin(x)) (2 \csc(x) x) \\ d(\sin(x) y) &= (2x) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}\sin(x) y &= \int 2x dx \\ \sin(x) y &= x^2 + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = \sin(x)$  results in

$$y = \csc(x) x^2 + c_1 \csc(x)$$

which simplifies to

$$y = \csc(x) (x^2 + c_1)$$

### Summary

The solution(s) found are the following

$$y = \csc(x) (x^2 + c_1) \tag{1}$$

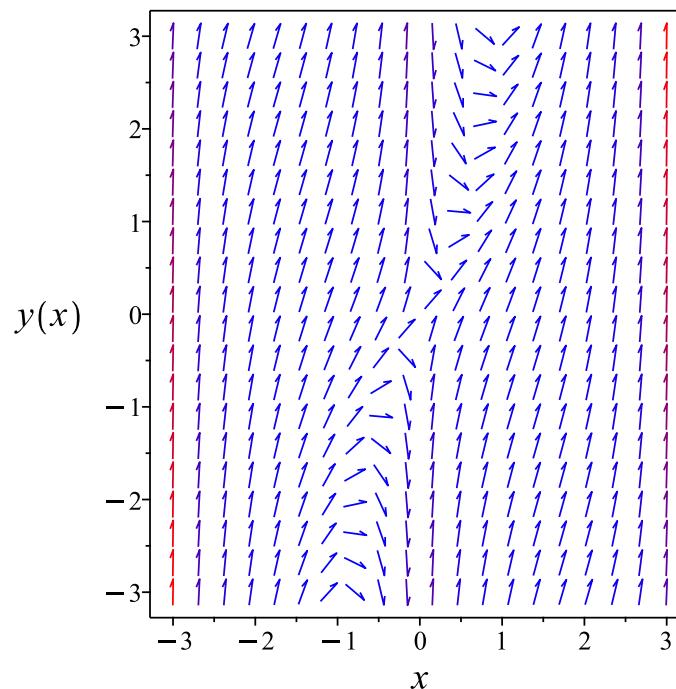


Figure 86: Slope field plot



### Verification of solutions

$$y = \csc(x) (x^2 + c_1)$$

Verified OK.

### **4.5.2 Solving as first order ode lie symmetry lookup ode**

Writing the ode as

$$\begin{aligned}y' &= -y \cot(x) + 2 \csc(x) x \\y' &= \omega(x, y)\end{aligned}$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 43: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{\sin(x)}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right)S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{\sin(x)}} dy \end{aligned}$$

Which results in

$$S = \sin(x) y$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -y \cot(x) + 2 \csc(x) x$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \cos(x) y \\ S_y &= \sin(x) \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 2x \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 2R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = R^2 + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$y \sin(x) = x^2 + c_1$$

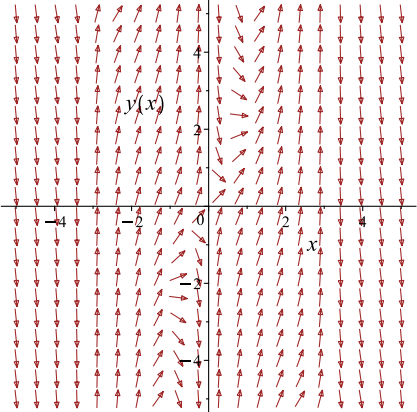
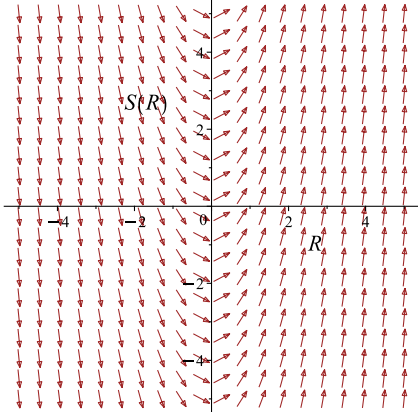
Which simplifies to

$$y \sin(x) = x^2 + c_1$$

Which gives

$$y = \frac{x^2 + c_1}{\sin(x)}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -y \cot(x) + 2 \csc(x) x$ 	$R = x$ $S = \sin(x) y$	$\frac{dS}{dR} = 2R$ 

### Summary

The solution(s) found are the following

$$y = \frac{x^2 + c_1}{\sin(x)} \quad (1)$$

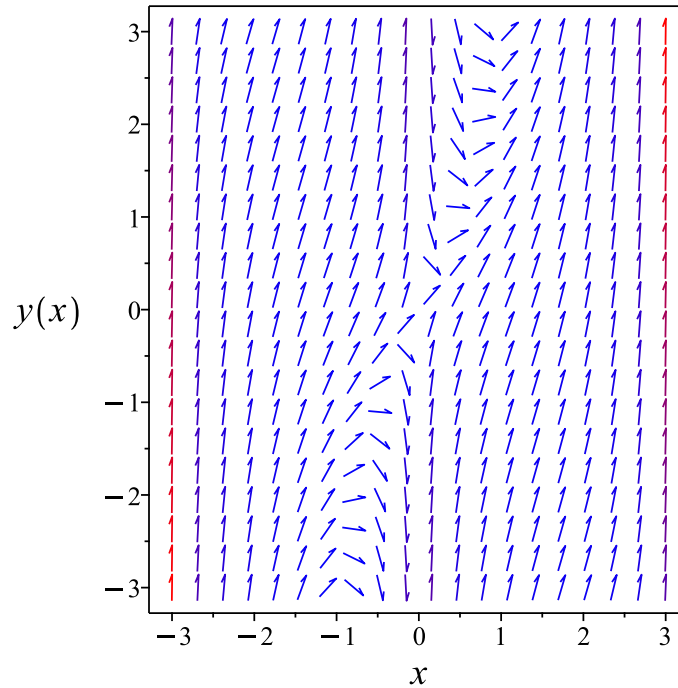


Figure 87: Slope field plot

Verification of solutions

$$y = \frac{x^2 + c_1}{\sin(x)}$$

Verified OK.

### 4.5.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dy &= (-y \cot(x) + 2 \csc(x) x) dx \\ (y \cot(x) - 2 \csc(x) x) dx + dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= y \cot(x) - 2 \csc(x) x \\ N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y \cot(x) - 2 \csc(x) x) \\ &= \cot(x)\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((\cot(x)) - (0)) \\ &= \cot(x) \end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned} \mu &= e^{\int A dx} \\ &= e^{\int \cot(x) dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\ln(\sin(x))} \\ &= \sin(x) \end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \sin(x) (y \cot(x) - 2 \csc(x) x) \\ &= \cos(x) y - 2x \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \sin(x) (1) \\ &= \sin(x) \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (\cos(x) y - 2x) + (\sin(x)) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \overline{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \cos(x) y - 2x dx \\ \phi &= \sin(x) y - x^2 + f(y)\end{aligned}\tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = \sin(x) + f'(y)\tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \sin(x)$ . Therefore equation (4) becomes

$$\sin(x) = \sin(x) + f'(y)\tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = \sin(x) y - x^2 + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \sin(x) y - x^2$$

The solution becomes

$$y = \frac{x^2 + c_1}{\sin(x)}$$



### Summary

The solution(s) found are the following

$$y = \frac{x^2 + c_1}{\sin(x)} \quad (1)$$

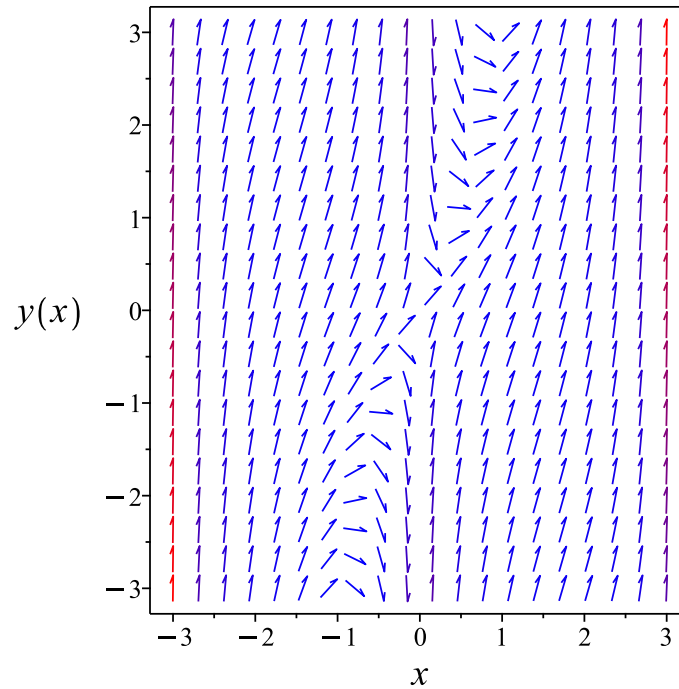


Figure 88: Slope field plot

### Verification of solutions

$$y = \frac{x^2 + c_1}{\sin(x)}$$

Verified OK.

#### 4.5.4 Maple step by step solution

Let's solve

$$y' + y \cot(x) = 2 \csc(x) x$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -y \cot(x) + 2 \csc(x) x$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + y \cot(x) = 2 \csc(x) x$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x) (y' + y \cot(x)) = 2\mu(x) \csc(x) x$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x) y)$

$$\mu(x) (y' + y \cot(x)) = \mu'(x) y + \mu(x) y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = \mu(x) \cot(x)$$

- Solve to find the integrating factor

$$\mu(x) = \sin(x)$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x) y) \right) dx = \int 2\mu(x) \csc(x) x dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x) y = \int 2\mu(x) \csc(x) x dx + c_1$$

- Solve for  $y$

$$y = \frac{\int 2\mu(x) \csc(x) x dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = \sin(x)$

$$y = \frac{\int 2 \csc(x) x \sin(x) dx + c_1}{\sin(x)}$$

- Evaluate the integrals on the rhs

$$y = \frac{x^2 + c_1}{\sin(x)}$$

- Simplify

$$y = \csc(x) (x^2 + c_1)$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 12

```
dsolve(diff(y(x),x)+y(x)*cot(x)=2*x*csc(x),y(x), singsol=all)
```

$$y(x) = \csc(x) (x^2 + c_1)$$

### ✓ Solution by Mathematica

Time used: 0.043 (sec). Leaf size: 14

```
DSolve[y'[x]+y[x]*Cot[x]==2*x*Csc[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow (x^2 + c_1) \csc(x)$$

## 4.6 problem 2(f)

4.6.1	Solving as linear ode . . . . .	379
4.6.2	Solving as first order ode lie symmetry lookup ode . . . . .	381
4.6.3	Solving as exact ode . . . . .	385
4.6.4	Maple step by step solution . . . . .	390

Internal problem ID [3113]

Internal file name [OUTPUT/2605\_Sunday\_June\_05\_2022\_03\_22\_09\_AM\_5462670/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, section 11, page 49

**Problem number:** 2(f).

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

[\_linear]

$$2y - xy' = x^3$$

### 4.6.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{2}{x}$$
$$q(x) = -x^2$$

Hence the ode is

$$y' - \frac{2y}{x} = -x^2$$

The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int -\frac{2}{x} dx} \\ &= \frac{1}{x^2}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu)(-x^2) \\ \frac{d}{dx}\left(\frac{y}{x^2}\right) &= \left(\frac{1}{x^2}\right)(-x^2) \\ d\left(\frac{y}{x^2}\right) &= -1 dx\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y}{x^2} &= \int -1 dx \\ \frac{y}{x^2} &= -x + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = \frac{1}{x^2}$  results in

$$y = c_1 x^2 - x^3$$

which simplifies to

$$y = x^2(-x + c_1)$$

### Summary

The solution(s) found are the following

$$y = x^2(-x + c_1) \tag{1}$$

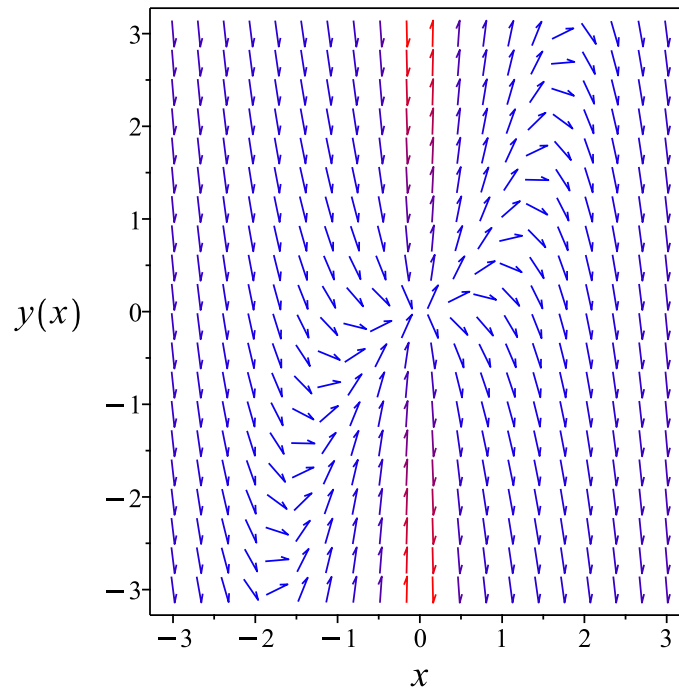


Figure 89: Slope field plot

Verification of solutions

$$y = x^2(-x + c_1)$$

Verified OK.

#### 4.6.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-x^3 + 2y}{x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 46: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= x^2\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{x^2} dy \end{aligned}$$

Which results in

$$S = \frac{y}{x^2}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{-x^3 + 2y}{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -\frac{2y}{x^3} \\ S_y &= \frac{1}{x^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -1 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -1$$



The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -R + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{y}{x^2} = -x + c_1$$

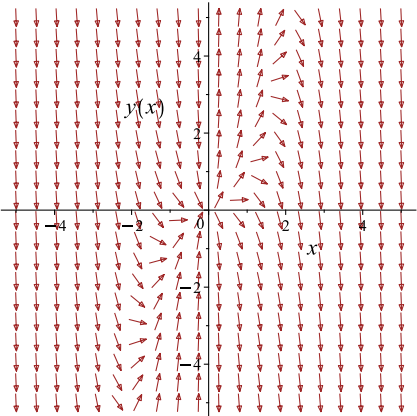
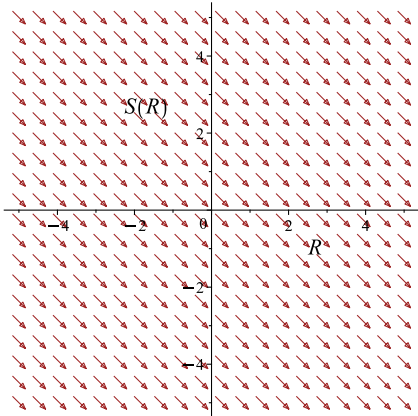
Which simplifies to

$$\frac{y}{x^2} = -x + c_1$$

Which gives

$$y = x^2(-x + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{-x^3 + 2y}{x}$ 	$R = x$ $S = \frac{y}{x^2}$	$\frac{dS}{dR} = -1$ 

### Summary

The solution(s) found are the following

$$y = x^2(-x + c_1) \quad (1)$$

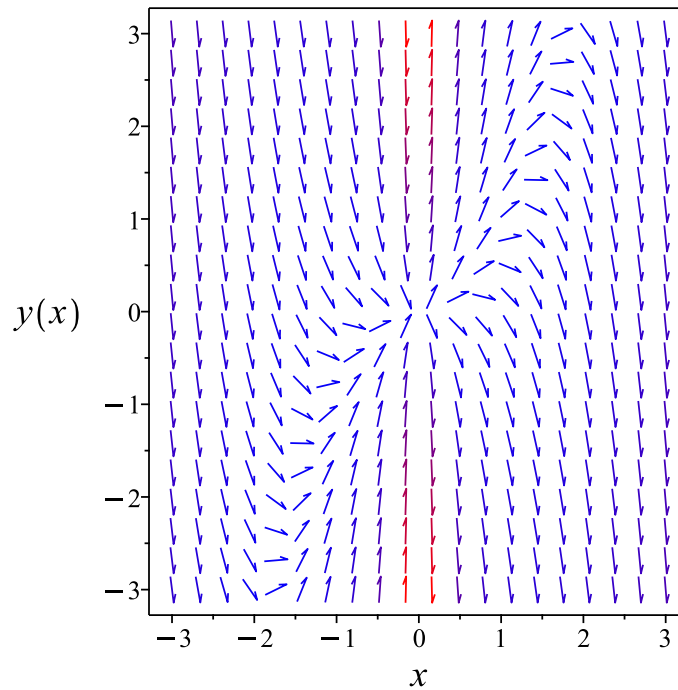


Figure 90: Slope field plot

Verification of solutions

$$y = x^2(-x + c_1)$$

Verified OK.

### 4.6.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(-x) dy &= (x^3 - 2y) dx \\ (-x^3 + 2y) dx + (-x) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -x^3 + 2y \\ N(x, y) &= -x\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x^3 + 2y) \\ &= 2\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x) \\ &= -1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= -\frac{1}{x} ((2) - (-1)) \\ &= -\frac{3}{x} \end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned} \mu &= e^{\int A dx} \\ &= e^{\int -\frac{3}{x} dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-3 \ln(x)} \\ &= \frac{1}{x^3} \end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{1}{x^3} (-x^3 + 2y) \\ &= \frac{-x^3 + 2y}{x^3} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \frac{1}{x^3} (-x) \\ &= -\frac{1}{x^2} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left( \frac{-x^3 + 2y}{x^3} \right) + \left( -\frac{1}{x^2} \right) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{-x^3 + 2y}{x^3} dx \\ \phi &= -x - \frac{y}{x^2} + f(y) \end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = -\frac{1}{x^2} + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = -\frac{1}{x^2}$ . Therefore equation (4) becomes

$$-\frac{1}{x^2} = -\frac{1}{x^2} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -x - \frac{y}{x^2} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -x - \frac{y}{x^2}$$

The solution becomes

$$y = -(x + c_1) x^2$$

Summary

The solution(s) found are the following

$$y = -(x + c_1) x^2 \tag{1}$$

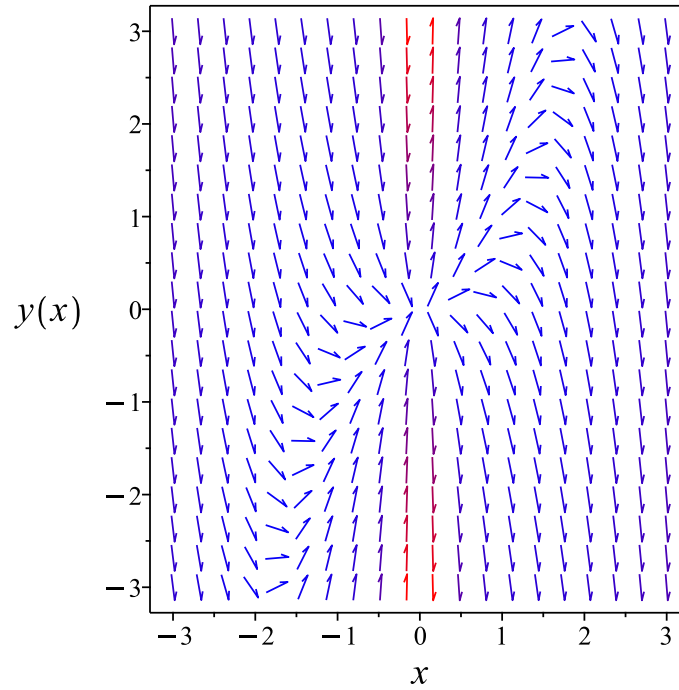


Figure 91: Slope field plot

Verification of solutions

$$y = -(x + c_1) x^2$$

Verified OK.

#### 4.6.4 Maple step by step solution

Let's solve

$$2y - xy' = x^3$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = \frac{2y}{x} - x^2$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' - \frac{2y}{x} = -x^2$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x) \left( y' - \frac{2y}{x} \right) = -\mu(x) x^2$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$

$$\mu(x) \left( y' - \frac{2y}{x} \right) = \mu'(x)y + \mu(x)y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = -\frac{2\mu(x)}{x}$$

- Solve to find the integrating factor

$$\mu(x) = \frac{1}{x^2}$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x)y) \right) dx = \int -\mu(x) x^2 dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int -\mu(x) x^2 dx + c_1$$

- Solve for  $y$

$$y = \frac{\int -\mu(x)x^2 dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = \frac{1}{x^2}$

$$y = x^2 \left( \int (-1) dx + c_1 \right)$$

- Evaluate the integrals on the rhs

$$y = x^2(-x + c_1)$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 13

```
dsolve((2*y(x)-x^3)=x*diff(y(x),x),y(x), singsol=all)
```

$$y(x) = (c_1 - x) x^2$$

### ✓ Solution by Mathematica

Time used: 0.026 (sec). Leaf size: 15

```
DSolve[(2*y[x]-x^3)==x*y'[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow x^2(-x + c_1)$$



## 5 Chapter 2, End of chapter, page 61

5.1	problem 2	393
5.2	problem 3	405
5.3	problem 4	415
5.4	problem 5	423
5.5	problem 6	432
5.6	problem 8	435
5.7	problem 9	447
5.8	problem 10	461
5.9	problem 12	468
5.10	problem 13	482
5.11	problem 14	494
5.12	problem 15	502
5.13	problem 17	514
5.14	problem 18	520
5.15	problem 19	532
5.16	problem 20	545
5.17	problem 21	559
5.18	problem 22	573
5.19	problem 24	579
5.20	problem 25	585

## 5.1 problem 2

- 5.1.1 Solving as first order ode lie symmetry calculated ode . . . . . 393
- 5.1.2 Solving as exact ode . . . . . 398

Internal problem ID [3114]

Internal file name [OUTPUT/2606\_Sunday\_June\_05\_2022\_03\_22\_11\_AM\_87162369/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 2.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exactWithIntegrationFactor",  
"first\_order\_ode\_lie\_symmetry\_calculated"

Maple gives the following as the ode type

```
[[_homogeneous, `class G`], _rational, [_Abel, `2nd type`, `class B`]]
```

$$(-yx + 1)y' - y^2 = 0$$

### 5.1.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{y^2}{xy - 1}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \tag{1E}$$

$$\eta = xb_2 + yb_3 + b_1 \tag{2E}$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$b_2 - \frac{y^2(b_3 - a_2)}{xy - 1} - \frac{y^4 a_3}{(xy - 1)^2} - \frac{y^3(xa_2 + ya_3 + a_1)}{(xy - 1)^2} - \left( -\frac{2y}{xy - 1} + \frac{y^2 x}{(xy - 1)^2} \right) (xb_2 + yb_3 + b_1) = 0 \quad (5E)$$

Putting the above in normal form gives

$$\frac{2x^2 y^2 b_2 - 2y^4 a_3 + x y^2 b_1 - y^3 a_1 - 4xyb_2 - y^2 a_2 - y^2 b_3 - 2yb_1 + b_2}{(xy - 1)^2} = 0$$

Setting the numerator to zero gives

$$2x^2 y^2 b_2 - 2y^4 a_3 + x y^2 b_1 - y^3 a_1 - 4xyb_2 - y^2 a_2 - y^2 b_3 - 2yb_1 + b_2 = 0 \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$-2a_3 v_2^4 + 2b_2 v_1^2 v_2^2 - a_1 v_2^3 + b_1 v_1 v_2^2 - a_2 v_2^2 - 4b_2 v_1 v_2 - b_3 v_2^2 - 2b_1 v_2 + b_2 = 0 \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$2b_2 v_1^2 v_2^2 + b_1 v_1 v_2^2 - 4b_2 v_1 v_2 - 2a_3 v_2^4 - a_1 v_2^3 + (-a_2 - b_3) v_2^2 - 2b_1 v_2 + b_2 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}b_1 &= 0 \\b_2 &= 0 \\-a_1 &= 0 \\-2a_3 &= 0 \\-2b_1 &= 0 \\-4b_2 &= 0 \\2b_2 &= 0 \\-a_2 - b_3 &= 0\end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}a_1 &= 0 \\a_2 &= -b_3 \\a_3 &= 0 \\b_1 &= 0 \\b_2 &= 0 \\b_3 &= b_3\end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= -x \\ \eta &= y\end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left( -\frac{y^2}{xy - 1} \right) (-x) \\ &= -\frac{y}{xy - 1} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{-\frac{y}{xy-1}} dy \end{aligned}$$

Which results in

$$S = -xy + \ln(y)$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y^2}{xy-1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -y \\ S_y &= -x + \frac{1}{y} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$-yx + \ln(y) = c_1$$

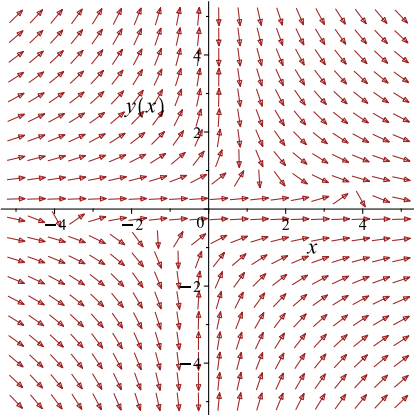
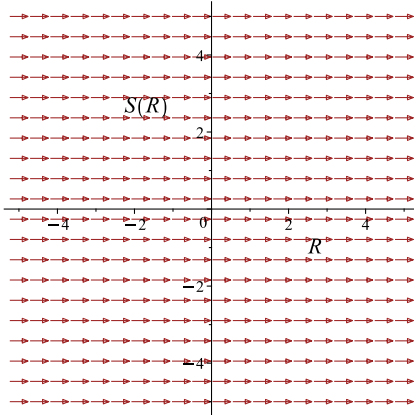
Which simplifies to

$$-yx + \ln(y) = c_1$$

Which gives

$$y = e^{-\text{LambertW}(-e^{c_1}x) + c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{y^2}{xy-1}$ 	$R = x$ $S = -xy + \ln(y)$	$\frac{dS}{dR} = 0$ 

### Summary

The solution(s) found are the following

$$y = e^{-\text{LambertW}(-e^{c_1}x)+c_1} \quad (1)$$

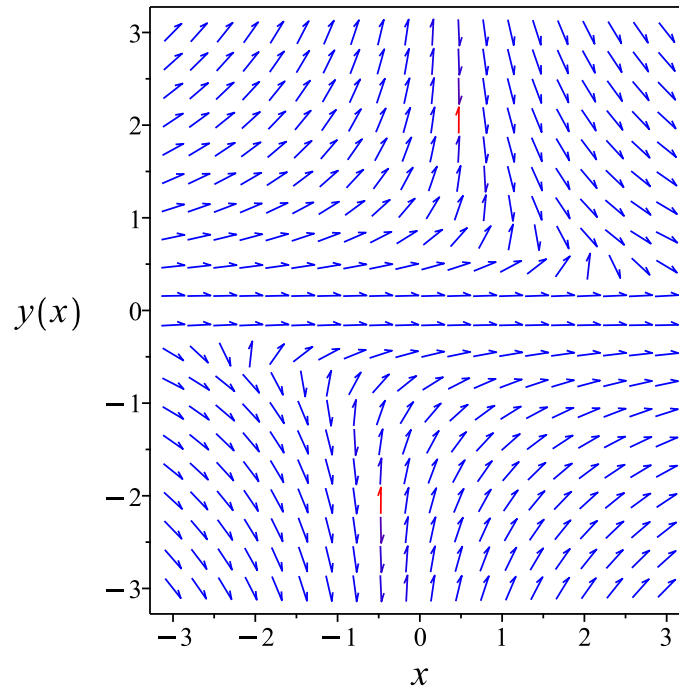


Figure 92: Slope field plot

### Verification of solutions

$$y = e^{-\text{LambertW}(-e^{c_1}x)+c_1}$$

Verified OK.

### 5.1.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (-xy + 1) dy &= (y^2) dx \\ (-y^2) dx + (-xy + 1) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -y^2 \\ N(x, y) &= -xy + 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-y^2) \\ &= -2y \end{aligned}$$



And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-xy + 1) \\ &= -y\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{-xy + 1} ((-2y) - (-y)) \\ &= \frac{y}{xy - 1}\end{aligned}$$

Since  $A$  depends on  $y$ , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned}B &= \frac{1}{M} \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= -\frac{1}{y^2} ((-y) - (-2y)) \\ &= -\frac{1}{y}\end{aligned}$$

Since  $B$  does not depend on  $x$ , it can be used to obtain an integrating factor. Let the integrating factor be  $\mu$ . Then

$$\begin{aligned}\mu &= e^{\int B \, dy} \\ &= e^{\int -\frac{1}{y} \, dy}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\ln(y)} \\ &= \frac{1}{y}\end{aligned}$$

$M$  and  $N$  are now multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{y}(-y^2) \\ &= -y\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{y}(-xy + 1) \\ &= \frac{-xy + 1}{y}\end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (-y) + \left( \frac{-xy + 1}{y} \right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -y dx \\ \phi &= -xy + f(y)\end{aligned} \tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = -x + f'(y) \tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{-xy+1}{y}$ . Therefore equation (4) becomes

$$\frac{-xy + 1}{y} = -x + f'(y) \tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = \frac{1}{y}$$

Integrating the above w.r.t  $y$  gives

$$\int f'(y) dy = \int \left(\frac{1}{y}\right) dy$$
$$f(y) = \ln(y) + c_1$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -xy + \ln(y) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -xy + \ln(y)$$

The solution becomes

$$y = e^{-\text{LambertW}(-e^{c_1}x)+c_1}$$

### Summary

The solution(s) found are the following

$$y = e^{-\text{LambertW}(-e^{c_1}x)+c_1} \tag{1}$$

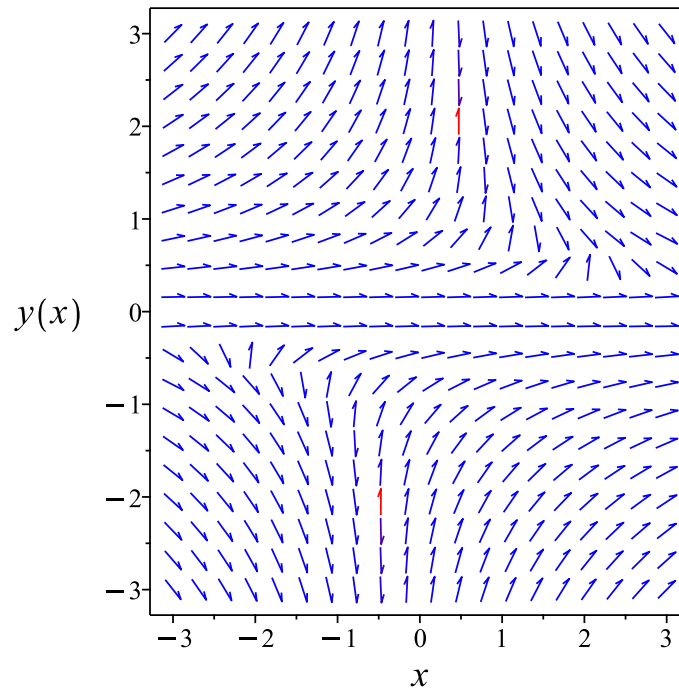


Figure 93: Slope field plot

Verification of solutions

$$y = e^{-\text{LambertW}(-e^{c_1}x)+c_1}$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
<- 1st order linear successful
<- inverse linear successful`

```

✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve((1-x*y(x))*diff(y(x),x)=y(x)^2,y(x), singsol=all)
```

$$y(x) = -\frac{\text{LambertW}(-x e^{-c_1})}{x}$$

✓ Solution by Mathematica

Time used: 2.155 (sec). Leaf size: 25

```
DSolve[(1-x*y[x])*y'[x]==y[x]^2,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{W(-e^{-c_1}x)}{x}$$
$$y(x) \rightarrow 0$$

## 5.2 problem 3

- 5.2.1 Solving as homogeneousTypeMapleC ode . . . . . 405
- 5.2.2 Solving as first order ode lie symmetry calculated ode . . . . . 408

Internal problem ID [3115]

Internal file name [OUTPUT/2607\_Sunday\_June\_05\_2022\_03\_22\_14\_AM\_19060358/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 3.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"homogeneousTypeMapleC",  
"first\_order\_ode\_lie\_symmetry\_calculated"**

Maple gives the following as the ode type

```
[[_homogeneous, `class C`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$3y + (2y - 3x + 5)y' = -1 - 2x$$

### 5.2.1 Solving as homogeneousTypeMapleC ode

Let  $Y = y + y_0$  and  $X = x + x_0$  then the above is transformed to new ode in  $Y(X)$

$$\frac{d}{dX}Y(X) = -\frac{2X + 2x_0 + 3Y(X) + 3y_0 + 1}{2Y(X) + 2y_0 - 3X - 3x_0 + 5}$$

Solving for possible values of  $x_0$  and  $y_0$  which makes the above ode a homogeneous ode results in

$$\begin{aligned}x_0 &= 1 \\y_0 &= -1\end{aligned}$$

Using these values now it is possible to easily solve for  $Y(X)$ . The above ode now becomes

$$\frac{d}{dX}Y(X) = -\frac{2X + 3Y(X)}{2Y(X) - 3X}$$

In canonical form, the ODE is

$$\begin{aligned} Y' &= F(X, Y) \\ &= -\frac{2X + 3Y}{2Y - 3X} \end{aligned} \quad (1)$$

An ode of the form  $Y' = \frac{M(X, Y)}{N(X, Y)}$  is called homogeneous if the functions  $M(X, Y)$  and  $N(X, Y)$  are both homogeneous functions and of the same order. Recall that a function  $f(X, Y)$  is homogeneous of order  $n$  if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both  $M = 2X + 3Y$  and  $N = -2Y + 3X$  are both homogeneous and of the same order  $n = 1$ . Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution  $u = \frac{Y}{X}$ , or  $Y = uX$ . Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation  $Y = uX$  to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dX}X + u &= \frac{-3u - 2}{2u - 3} \\ \frac{du}{dX} &= \frac{\frac{-3u(X)-2}{2u(X)-3} - u(X)}{X} \end{aligned}$$

Or

$$\frac{d}{dX}u(X) - \frac{\frac{-3u(X)-2}{2u(X)-3} - u(X)}{X} = 0$$

Or

$$2\left(\frac{d}{dX}u(X)\right)Xu(X) - 3\left(\frac{d}{dX}u(X)\right)X + 2u(X)^2 + 2 = 0$$

Or

$$2 + X(2u(X) - 3)\left(\frac{d}{dX}u(X)\right) + 2u(X)^2 = 0$$

Which is now solved as separable in  $u(X)$ . Which is now solved in  $u(X)$ . In canonical form the ODE is

$$\begin{aligned} u' &= F(X, u) \\ &= f(X)g(u) \\ &= -\frac{2(u^2 + 1)}{X(2u - 3)} \end{aligned}$$

Where  $f(X) = -\frac{2}{X}$  and  $g(u) = \frac{u^2+1}{2u-3}$ . Integrating both sides gives

$$\frac{1}{\frac{u^2+1}{2u-3}} du = -\frac{2}{X} dX$$

$$\int \frac{1}{\frac{u^2+1}{2u-3}} du = \int -\frac{2}{X} dX$$

$$\ln(u^2 + 1) - 3 \arctan(u) = -2 \ln(X) + c_2$$

The solution is

$$\ln(u(X)^2 + 1) - 3 \arctan(u(X)) + 2 \ln(X) - c_2 = 0$$

Now  $u$  in the above solution is replaced back by  $Y$  using  $u = \frac{Y}{X}$  which results in the solution

$$\ln\left(\frac{Y(X)^2}{X^2} + 1\right) - 3 \arctan\left(\frac{Y(X)}{X}\right) + 2 \ln(X) - c_2 = 0$$

Using the solution for  $Y(X)$

$$\ln\left(\frac{Y(X)^2}{X^2} + 1\right) - 3 \arctan\left(\frac{Y(X)}{X}\right) + 2 \ln(X) - c_2 = 0$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y - 1$$

$$X = x + 1$$

Then the solution in  $y$  becomes

$$\ln\left(\frac{(y+1)^2}{(x-1)^2} + 1\right) - 3 \arctan\left(\frac{y+1}{x-1}\right) + 2 \ln(x-1) - c_2 = 0$$

### Summary

The solution(s) found are the following

$$\ln\left(\frac{(y+1)^2}{(x-1)^2} + 1\right) - 3 \arctan\left(\frac{y+1}{x-1}\right) + 2 \ln(x-1) - c_2 = 0 \quad (1)$$



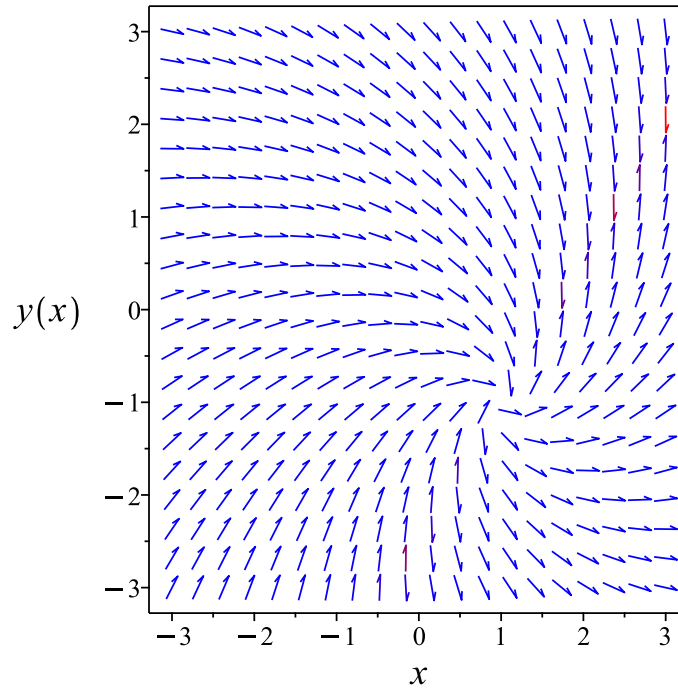


Figure 94: Slope field plot

### Verification of solutions

$$\ln \left( \frac{(y+1)^2}{(x-1)^2} + 1 \right) - 3 \arctan \left( \frac{y+1}{x-1} \right) + 2 \ln(x-1) - c_2 = 0$$

Verified OK.

### 5.2.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{2x + 3y + 1}{2y - 3x + 5}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} b_2 - \frac{(2x + 3y + 1)(b_3 - a_2)}{2y - 3x + 5} - \frac{(2x + 3y + 1)^2 a_3}{(2y - 3x + 5)^2} \\ - \left( -\frac{2}{2y - 3x + 5} - \frac{3(2x + 3y + 1)}{(2y - 3x + 5)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left( -\frac{3}{2y - 3x + 5} + \frac{4x + 6y + 2}{(2y - 3x + 5)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{6x^2a_2 + 4x^2a_3 + 4x^2b_2 - 6x^2b_3 - 8xya_2 + 12xya_3 + 12xyb_2 + 8xyb_3 - 6y^2a_2 - 4y^2a_3 - 4y^2b_2 + 6y^2b_3 - \dots}{\dots} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} -6x^2a_2 - 4x^2a_3 - 4x^2b_2 + 6x^2b_3 + 8xya_2 - 12xya_3 - 12xyb_2 - 8xyb_3 + 6y^2a_2 \\ + 4y^2a_3 + 4y^2b_2 - 6y^2b_3 + 20xa_2 - 4xa_3 - 13xb_1 - 17xb_2 - 7xb_3 + 13ya_1 \\ + 17ya_2 + 7ya_3 + 20yb_2 - 4yb_3 + 13a_1 + 5a_2 - a_3 + 13b_1 + 25b_2 - 5b_3 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} -6a_2v_1^2 + 8a_2v_1v_2 + 6a_2v_2^2 - 4a_3v_1^2 - 12a_3v_1v_2 + 4a_3v_2^2 - 4b_2v_1^2 \\ - 12b_2v_1v_2 + 4b_2v_2^2 + 6b_3v_1^2 - 8b_3v_1v_2 - 6b_3v_2^2 + 13a_1v_2 \\ + 20a_2v_1 + 17a_2v_2 - 4a_3v_1 + 7a_3v_2 - 13b_1v_1 - 17b_2v_1 + 20b_2v_2 \\ - 7b_3v_1 - 4b_3v_2 + 13a_1 + 5a_2 - a_3 + 13b_1 + 25b_2 - 5b_3 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} &(-6a_2 - 4a_3 - 4b_2 + 6b_3)v_1^2 + (8a_2 - 12a_3 - 12b_2 - 8b_3)v_1v_2 \\ &+ (20a_2 - 4a_3 - 13b_1 - 17b_2 - 7b_3)v_1 + (6a_2 + 4a_3 + 4b_2 - 6b_3)v_2^2 \\ &+ (13a_1 + 17a_2 + 7a_3 + 20b_2 - 4b_3)v_2 + 13a_1 + 5a_2 - a_3 + 13b_1 + 25b_2 - 5b_3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -6a_2 - 4a_3 - 4b_2 + 6b_3 &= 0 \\ 6a_2 + 4a_3 + 4b_2 - 6b_3 &= 0 \\ 8a_2 - 12a_3 - 12b_2 - 8b_3 &= 0 \\ 13a_1 + 17a_2 + 7a_3 + 20b_2 - 4b_3 &= 0 \\ 20a_2 - 4a_3 - 13b_1 - 17b_2 - 7b_3 &= 0 \\ 13a_1 + 5a_2 - a_3 + 13b_1 + 25b_2 - 5b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= -b_3 - b_2 \\ a_2 &= b_3 \\ a_3 &= -b_2 \\ b_1 &= b_3 - b_2 \\ b_2 &= b_2 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x - 1 \\ \eta &= y + 1 \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= y + 1 - \left( -\frac{2x + 3y + 1}{2y - 3x + 5} \right) (x - 1) \\ &= \frac{-2x^2 - 2y^2 + 4x - 4y - 4}{-2y + 3x - 5} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-2x^2 - 2y^2 + 4x - 4y - 4}{-2y + 3x - 5}} dy \end{aligned}$$

Which results in

$$S = \frac{\ln(x^2 + y^2 - 2x + 2y + 2)}{2} + \frac{(-3x + 3) \arctan\left(\frac{2+2y}{2x-2}\right)}{2x - 2}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{2x + 3y + 1}{2y - 3x + 5}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{2x + 3y + 1}{2x^2 + 2y^2 - 4x + 4y + 4} \\ S_y &= \frac{2y - 3x + 5}{2x^2 + 2y^2 - 4x + 4y + 4} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \tag{2A}$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

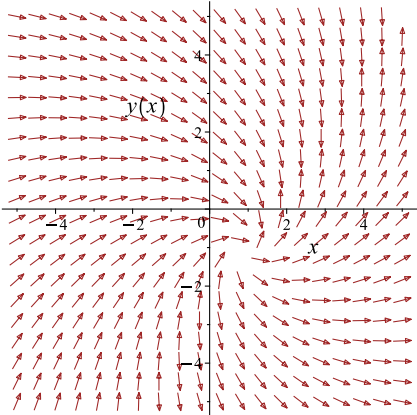
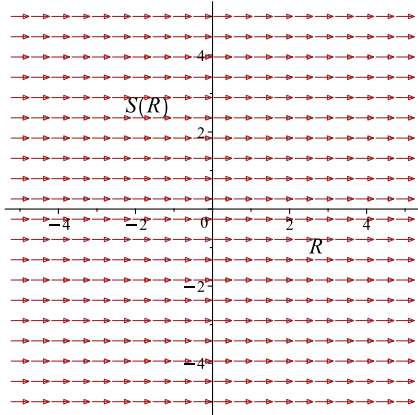
To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{\ln(y^2 + x^2 + 2y - 2x + 2)}{2} - \frac{3 \arctan\left(\frac{y+1}{x-1}\right)}{2} = c_1$$

Which simplifies to

$$\frac{\ln(y^2 + x^2 + 2y - 2x + 2)}{2} - \frac{3 \arctan\left(\frac{y+1}{x-1}\right)}{2} = c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{2x+3y+1}{2y-3x+5}$ 	$R = x$ $S = \frac{\ln(x^2 + y^2 - 2x + 2)}{2}$	$\frac{dS}{dR} = 0$ 

### Summary

The solution(s) found are the following

$$\frac{\ln(y^2 + x^2 + 2y - 2x + 2)}{2} - \frac{3 \arctan\left(\frac{y+1}{x-1}\right)}{2} = c_1 \quad (1)$$

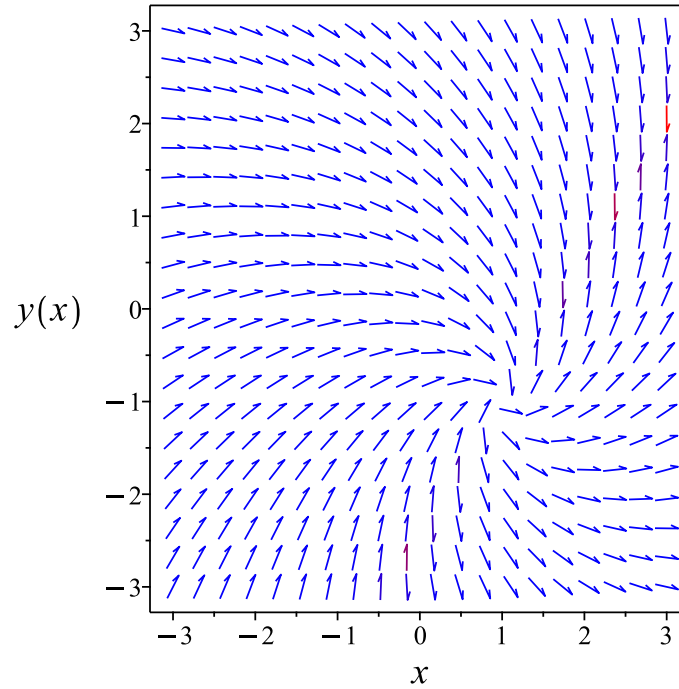


Figure 95: Slope field plot

### Verification of solutions

$$\frac{\ln(y^2 + x^2 + 2y - 2x + 2)}{2} - \frac{3 \arctan\left(\frac{y+1}{x-1}\right)}{2} = c_1$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous C  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful  
<- homogeneous successful`
```

### ✓ Solution by Maple

Time used: 0.046 (sec). Leaf size: 31

```
dsolve((2*x+3*y(x)+1)+(2*y(x)-3*x+5)*diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = -1 - \tan(\text{RootOf}(3\_Z + \ln(\sec(\_Z)^2) + 2 \ln(x - 1) + 2c_1))(x - 1)$$

### ✓ Solution by Mathematica

Time used: 0.065 (sec). Leaf size: 68

```
DSolve[(2*x+3*y[x]+1)+(2*y[x]-3*x+5)*y'[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve} \left[ 54 \arctan \left( \frac{3y(x) + 2x + 1}{2y(x) - 3x + 5} \right) + 18 \log \left( \frac{4(x^2 + y(x)^2 + 2y(x) - 2x + 2)}{13(x - 1)^2} \right) + 36 \log(x - 1) + 13c_1 = 0, y(x) \right]$$

## 5.3 problem 4

5.3.1 Solving as first order ode lie symmetry calculated ode . . . . . 415

Internal problem ID [3116]

Internal file name [OUTPUT/2608\_Sunday\_June\_05\_2022\_03\_22\_18\_AM\_69094642/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 4.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**first\_order\_ode\_lie\_symmetry\_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, _dAlembert]
```

$$xy' - \sqrt{y^2 + x^2} = 0$$

### 5.3.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{\sqrt{x^2 + y^2}}{x}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$



Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$b_2 + \frac{\sqrt{x^2 + y^2} (b_3 - a_2)}{x} - \frac{(x^2 + y^2) a_3}{x^2} \quad (5E)$$

$$- \left( \frac{1}{\sqrt{x^2 + y^2}} - \frac{\sqrt{x^2 + y^2}}{x^2} \right) (x a_2 + y a_3 + a_1) - \frac{y(x b_2 + y b_3 + b_1)}{x \sqrt{x^2 + y^2}} = 0$$

Putting the above in normal form gives

$$\frac{\sqrt{x^2 + y^2} x^2 a_3 - b_2 \sqrt{x^2 + y^2} x^2 + \sqrt{x^2 + y^2} y^2 a_3 + x^3 a_2 - x^3 b_3 + x^2 y b_2 - y^3 a_3 + x y b_1 - y^2 a_1}{\sqrt{x^2 + y^2} x^2} = 0$$

Setting the numerator to zero gives

$$-\sqrt{x^2 + y^2} x^2 a_3 + b_2 \sqrt{x^2 + y^2} x^2 - \sqrt{x^2 + y^2} y^2 a_3 \quad (6E)$$

$$- x^3 a_2 + x^3 b_3 - x^2 y b_2 + y^3 a_3 - x y b_1 + y^2 a_1 = 0$$

Simplifying the above gives

$$x(x^2 + y^2) b_3 + (x^2 + y^2) y a_3 - \sqrt{x^2 + y^2} x^2 a_3 \quad (6E)$$

$$+ b_2 \sqrt{x^2 + y^2} x^2 - \sqrt{x^2 + y^2} y^2 a_3 - x^3 a_2 - x^2 y a_3$$

$$- x^2 y b_2 - x y^2 b_3 + (x^2 + y^2) a_1 - x^2 a_1 - x y b_1 = 0$$

Since the PDE has radicals, simplifying gives

$$-\sqrt{x^2 + y^2} x^2 a_3 + b_2 \sqrt{x^2 + y^2} x^2 - \sqrt{x^2 + y^2} y^2 a_3$$

$$- x^3 a_2 + x^3 b_3 - x^2 y b_2 + y^3 a_3 - x y b_1 + y^2 a_1 = 0$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y, \sqrt{x^2 + y^2}\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2, \sqrt{x^2 + y^2} = v_3\}$$

The above PDE (6E) now becomes

$$-v_1^3 a_2 - v_3 v_1^2 a_3 + v_2^3 a_3 - v_3 v_2^2 a_3 - v_1^2 v_2 b_2 + b_2 v_3 v_1^2 + v_1^3 b_3 + v_2^2 a_1 - v_1 v_2 b_1 = 0 \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2, v_3\}$$

Equation (7E) now becomes

$$(b_3 - a_2) v_1^3 - v_1^2 v_2 b_2 + (-a_3 + b_2) v_1^2 v_3 - v_1 v_2 b_1 + v_2^3 a_3 - v_3 v_2^2 a_3 + v_2^2 a_1 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} a_1 &= 0 \\ a_3 &= 0 \\ -a_3 &= 0 \\ -b_1 &= 0 \\ -b_2 &= 0 \\ -a_3 + b_2 &= 0 \\ b_3 - a_2 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left( \frac{\sqrt{x^2 + y^2}}{x} \right) (x) \\ &= y - \sqrt{x^2 + y^2} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left( \xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{y - \sqrt{x^2 + y^2}} dy\end{aligned}$$

Which results in

$$S = -\frac{y^2}{2x^2} - \frac{y\sqrt{x^2 + y^2}}{2x^2} - \frac{\ln(y + \sqrt{x^2 + y^2})}{2}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{\sqrt{x^2 + y^2}}{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}
 R_x &= 1 \\
 R_y &= 0 \\
 S_x &= -\frac{(-x^2y - 4y^3)\sqrt{x^2 + y^2} + x^4 - 3y^2x^2 - 4y^4}{2\sqrt{x^2 + y^2}(y + \sqrt{x^2 + y^2})x^3} \\
 S_y &= -\frac{y\sqrt{x^2 + y^2} + x^2 + y^2}{\sqrt{x^2 + y^2}x^2}
 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{3(y\sqrt{x^2 + y^2} + x^2 + y^2)}{2x\sqrt{x^2 + y^2}(y + \sqrt{x^2 + y^2})} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{3}{2R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -\frac{3 \ln(R)}{2} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{-\ln(y + \sqrt{y^2 + x^2})x^2 - y(y + \sqrt{y^2 + x^2})}{2x^2} = -\frac{3 \ln(x)}{2} + c_1$$

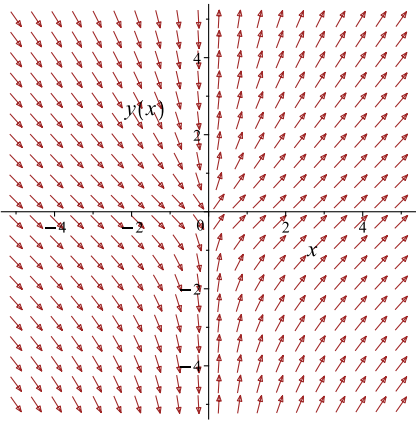
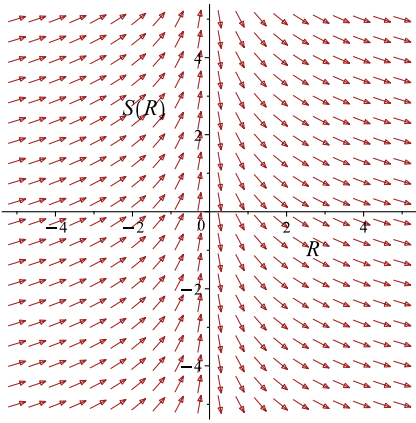
Which simplifies to

$$\frac{-\ln(y + \sqrt{y^2 + x^2})x^2 - y(y + \sqrt{y^2 + x^2})}{2x^2} = -\frac{3 \ln(x)}{2} + c_1$$

Which gives

$$y = \frac{x^2(\text{LambertW}(x^4 e^{-4c_1+1}) - 1) \sqrt{\frac{1}{x^2 \text{LambertW}(x^4 e^{-4c_1+1})}}}{2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{\sqrt{x^2+y^2}}{x}$ 	$R = x$ $S = \frac{-\ln(y + \sqrt{x^2 + y^2})}{2}$	$\frac{dS}{dR} = -\frac{3}{2R}$ 

Summary

The solution(s) found are the following

$$y = \frac{x^2(\text{LambertW}(x^4 e^{-4c_1+1}) - 1) \sqrt{\frac{1}{x^2 \text{LambertW}(x^4 e^{-4c_1+1})}}}{2} \tag{1}$$

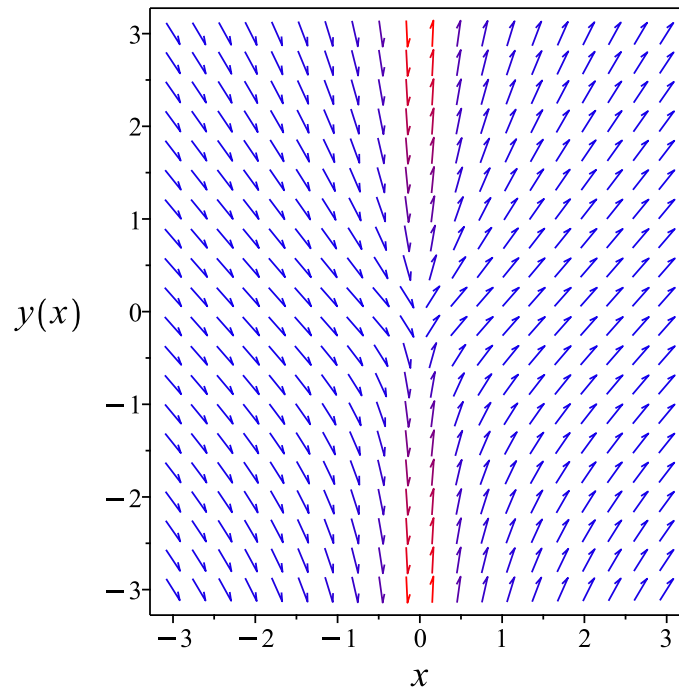


Figure 96: Slope field plot

Verification of solutions

$$y = \frac{x^2(\text{LambertW}(x^4 e^{-4c_1+1}) - 1) \sqrt{\frac{1}{x^2 \text{LambertW}(x^4 e^{-4c_1+1})}}}{2}$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying homogeneous types:
trying homogeneous G
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 51

```
dsolve(x*diff(y(x),x)=sqrt(x^2+y(x)^2),y(x), singsol=all)
```

$$\frac{y(x)^2 + y(x) \sqrt{x^2 + y(x)^2} + \left( \ln \left( y(x) + \sqrt{x^2 + y(x)^2} \right) - c_1 - 3 \ln(x) \right) x^2}{x^2} = 0$$

✓ Solution by Mathematica

Time used: 0.313 (sec). Leaf size: 66

```
DSolve[x*y'[x]==Sqrt[x^2+y[x]^2],y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve} \left[ \frac{1}{2} \left( \frac{y(x) \left( \sqrt{\frac{y(x)^2}{x^2} + 1} + \frac{y(x)}{x} \right)}{x} - \log \left( \sqrt{\frac{y(x)^2}{x^2} + 1} - \frac{y(x)}{x} \right) \right) = \log(x) + c_1, y(x) \right]$$

## 5.4 problem 5

5.4.1 Solving as first order ode lie symmetry calculated ode . . . . . 423

Internal problem ID [3117]

Internal file name [OUTPUT/2609\_Sunday\_June\_05\_2022\_03\_22\_21\_AM\_13020875/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 5.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**first\_order\_ode\_lie\_symmetry\_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class G`], _rational, [_Abel, `2nd type`, `class B`]]
```

$$y^2 - (x^3 - yx) y' = 0$$

### 5.4.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{y^2}{x(-x^2 + y)}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$



Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} b_2 - \frac{y^2(b_3 - a_2)}{x(-x^2 + y)} - \frac{y^4 a_3}{x^2(-x^2 + y)^2} \\ - \left( \frac{y^2}{x^2(-x^2 + y)} - \frac{2y^2}{(-x^2 + y)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left( -\frac{2y}{x(-x^2 + y)} + \frac{y^2}{x(-x^2 + y)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{x^6 b_2 - 4x^4 y b_2 + 2x^3 y^2 a_2 - x^3 y^2 b_3 + 3x^2 y^3 a_3 - 2x^3 y b_1 + 3x^2 y^2 a_1 + 2x^2 y^2 b_2 - 2y^4 a_3 + x y^2 b_1 - y^3 a_1}{x^2(x^2 - y)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} x^6 b_2 - 4x^4 y b_2 + 2x^3 y^2 a_2 - x^3 y^2 b_3 + 3x^2 y^3 a_3 - 2x^3 y b_1 \\ + 3x^2 y^2 a_1 + 2x^2 y^2 b_2 - 2y^4 a_3 + x y^2 b_1 - y^3 a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} b_2 v_1^6 + 2a_2 v_1^3 v_2^2 + 3a_3 v_1^2 v_2^3 - 4b_2 v_1^4 v_2 - b_3 v_1^3 v_2^2 + 3a_1 v_1^2 v_2^2 \\ - 2a_3 v_2^4 - 2b_1 v_1^3 v_2 + 2b_2 v_1^2 v_2^2 - a_1 v_2^3 + b_1 v_1 v_2^2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} b_2 v_1^6 - 4b_2 v_1^4 v_2 + (2a_2 - b_3) v_1^3 v_2^2 - 2b_1 v_1^3 v_2 + 3a_3 v_1^2 v_2^3 \\ + (3a_1 + 2b_2) v_1^2 v_2^2 + b_1 v_1 v_2^2 - 2a_3 v_2^4 - a_1 v_2^3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} b_1 &= 0 \\ b_2 &= 0 \\ -a_1 &= 0 \\ -2a_3 &= 0 \\ 3a_3 &= 0 \\ -2b_1 &= 0 \\ -4b_2 &= 0 \\ 3a_1 + 2b_2 &= 0 \\ 2a_2 - b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= a_2 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= 2a_2 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= 2y \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= 2y - \left( -\frac{y^2}{x(-x^2 + y)} \right) (x) \\ &= \frac{2x^2y - 3y^2}{x^2 - y} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left( \xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{2x^2y - 3y^2}{x^2 - y}} dy\end{aligned}$$

Which results in

$$S = -\frac{\ln(-2x^2 + 3y)}{6} + \frac{\ln(y)}{2}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{y^2}{x(-x^2 + y)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= -\frac{2x}{6x^2 - 9y} \\S_y &= \frac{x^2 - y}{2x^2y - 3y^2}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{3R} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{3R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -\frac{\ln(R)}{3} + c_1 \quad (4)$$

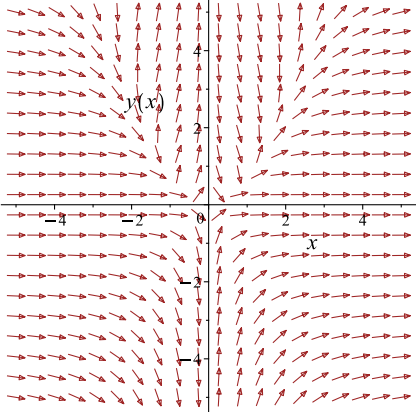
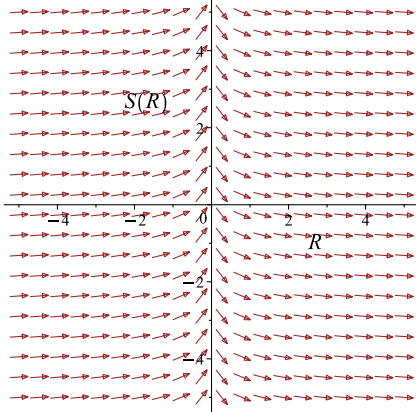
To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$-\frac{\ln(-2x^2 + 3y)}{6} + \frac{\ln(y)}{2} = -\frac{\ln(x)}{3} + c_1$$

Which simplifies to

$$-\frac{\ln(-2x^2 + 3y)}{6} + \frac{\ln(y)}{2} = -\frac{\ln(x)}{3} + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{y^2}{x(-x^2+y)}$ 	$R = x$ $S = -\frac{\ln(-2x^2 + 3y)}{6} + \frac{1}{3}$	$\frac{dS}{dR} = -\frac{1}{3R}$ 

Summary

The solution(s) found are the following

$$-\frac{\ln(-2x^2 + 3y)}{6} + \frac{\ln(y)}{2} = -\frac{\ln(x)}{3} + c_1 \tag{1}$$

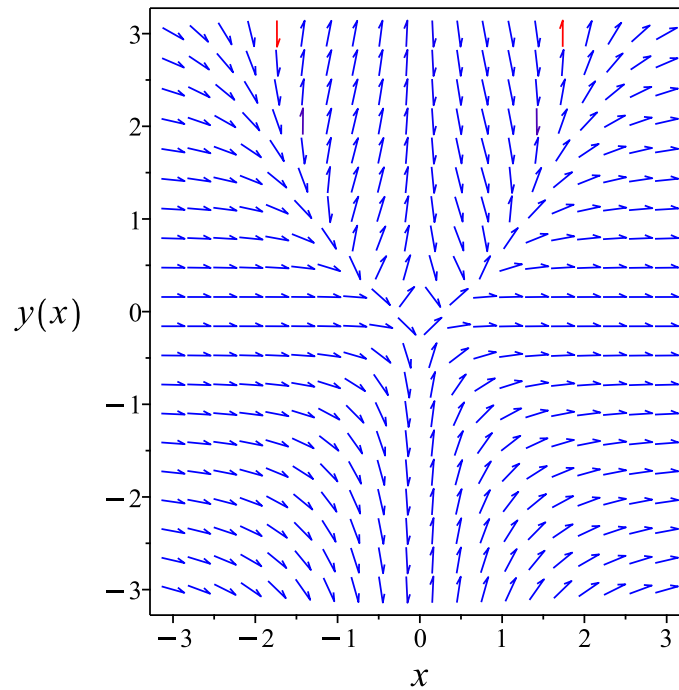


Figure 97: Slope field plot

Verification of solutions

$$-\frac{\ln(-2x^2 + 3y)}{6} + \frac{\ln(y)}{2} = -\frac{\ln(x)}{3} + c_1$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous G
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.5 (sec). Leaf size: 211

```
dsolve(y(x)^2=(x^3-x*y(x))*diff(y(x),x),y(x), singsol=all)
```

$$y(x) = \frac{c_1 \left( \left( -x^3 + \sqrt{x^6 - c_1^3} \right)^{\frac{2}{3}} + c_1 \right)}{x \left( -x^3 + \sqrt{x^6 - c_1^3} \right)^{\frac{1}{3}}}$$

$$y(x) = -\frac{c_1 \left( i\sqrt{3} \left( -x^3 + \sqrt{x^6 - c_1^3} \right)^{\frac{2}{3}} - i\sqrt{3} c_1 + \left( -x^3 + \sqrt{x^6 - c_1^3} \right)^{\frac{2}{3}} + c_1 \right)}{2x \left( -x^3 + \sqrt{x^6 - c_1^3} \right)^{\frac{1}{3}}}$$

$$y(x) = -\frac{c_1 \left( -i\sqrt{3} \left( -x^3 + \sqrt{x^6 - c_1^3} \right)^{\frac{2}{3}} + i\sqrt{3} c_1 + \left( -x^3 + \sqrt{x^6 - c_1^3} \right)^{\frac{2}{3}} + c_1 \right)}{2x \left( -x^3 + \sqrt{x^6 - c_1^3} \right)^{\frac{1}{3}}}$$

✓ Solution by Mathematica

Time used: 60.13 (sec). Leaf size: 820

`DSolve[y[x]^2==(x^3-x*y[x])*y'[x],y[x],x,IncludeSingularSolutions -> True]`

$$y(x) \rightarrow x^2$$

$$\frac{9 \sqrt[3]{x^{12} \left( -\cosh\left(\frac{3c_1}{4}\right) \right) - x^{12} \sinh\left(\frac{3c_1}{4}\right) + 2x^6 \cosh\left(\frac{3c_1}{8}\right) + 2x^6 \sinh\left(\frac{3c_1}{8}\right) + \sqrt{x^6 \left( \cosh\left(\frac{15c_1}{16}\right) - 1 \right)}}{x^6 \cosh\left(\frac{3c_1}{8}\right) + x^6 \sinh\left(\frac{3c_1}{8}\right) - 1}$$

$$y(x) \rightarrow x^2$$

$$\frac{9i(\sqrt{3}+i) \sqrt[3]{x^{12} \left( -\cosh\left(\frac{3c_1}{4}\right) \right) - x^{12} \sinh\left(\frac{3c_1}{4}\right) + 2x^6 \cosh\left(\frac{3c_1}{8}\right) + 2x^6 \sinh\left(\frac{3c_1}{8}\right) + \sqrt{x^6 \left( \cosh\left(\frac{15c_1}{16}\right) - 1 \right)}}{x^6 \cosh\left(\frac{3c_1}{8}\right) + x^6 \sinh\left(\frac{3c_1}{8}\right) - 1}$$

$$y(x) \rightarrow x^2$$

$$\frac{9i(\sqrt{3}-i) \sqrt[3]{x^{12} \left( -\cosh\left(\frac{3c_1}{4}\right) \right) - x^{12} \sinh\left(\frac{3c_1}{4}\right) + 2x^6 \cosh\left(\frac{3c_1}{8}\right) + 2x^6 \sinh\left(\frac{3c_1}{8}\right) + \sqrt{x^6 \left( \cosh\left(\frac{15c_1}{16}\right) - 1 \right)}}{x^6 \cosh\left(\frac{3c_1}{8}\right) + x^6 \sinh\left(\frac{3c_1}{8}\right) - 1}$$



## 5.5 problem 6

Internal problem ID [3118]

Internal file name [OUTPUT/2610\_Sunday\_June\_05\_2022\_03\_22\_25\_AM\_64609755/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 6.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**unknown**"

Maple gives the following as the ode type

`[_rational]`

Unable to solve or complete the solution.

$$y^3 + y - (y^2x^3 - x)y' = -x^2$$

Unable to determine ODE type.

## Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
Looking for potential symmetries
trying inverse_Riccati
trying an equivalence to an Abel ODE
differential order: 1; trying a linearization to 2nd order
--- trying a change of variables {x -> y(x), y(x) -> x}
differential order: 1; trying a linearization to 2nd order
trying 1st order ODE linearizable_by_differentiation
--- Trying Lie symmetry methods, 1st order ---
`, `-> Computing symmetries using: way = 2
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = 4
trying symmetry patterns for 1st order ODEs
-> trying a symmetry pattern of the form [F(x)*G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)*G(y)]
-> trying symmetry patterns of the forms [F(x),G(y)] and [G(y),F(x)]
-> trying a symmetry pattern of the form [F(x),G(x)]
-> trying a symmetry pattern of the form [F(y),G(y)]
-> trying a symmetry pattern of the form [F(x)+G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)+G(y)]
-> trying a symmetry pattern of the form [F(x),G(x)*y+H(x)]
-> trying a symmetry pattern of conformal type`
```

## **X** Solution by Maple

```
dsolve((x^2+y(x)^3+y(x))=(x^3*y(x)^2-x)*diff(y(x),x),y(x), singsol=all)
```

No solution found

**X** Solution by Mathematica

Time used: 0.0 (sec). Leaf size: 0

```
DSolve[(x^2+y[x]^3+y[x])==(x^3*y[x]^2-x)*y'[x],y[x],x,IncludeSingularSolutions -> True]
```

Not solved

## 5.6 problem 8

5.6.1	Solving as linear ode . . . . .	435
5.6.2	Solving as first order ode lie symmetry lookup ode . . . . .	437
5.6.3	Solving as exact ode . . . . .	441
5.6.4	Maple step by step solution . . . . .	445

Internal problem ID [3119]

Internal file name [OUTPUT/2611\_Sunday\_June\_05\_2022\_03\_22\_28\_AM\_67308557/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 8.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"exact", "linear", "first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$xy' + y = \cos(x)x$$

### 5.6.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{1}{x}$$
$$q(x) = \cos(x)$$

Hence the ode is

$$y' + \frac{y}{x} = \cos(x)$$

The integrating factor  $\mu$  is

$$\mu = e^{\int \frac{1}{x} dx}$$
$$= x$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) (\cos(x)) \\ \frac{d}{dx}(xy) &= (x) (\cos(x)) \\ d(xy) &= (\cos(x) x) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}xy &= \int \cos(x) x dx \\ xy &= x \sin(x) + \cos(x) + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = x$  results in

$$y = \frac{x \sin(x) + \cos(x)}{x} + \frac{c_1}{x}$$

which simplifies to

$$y = \frac{x \sin(x) + \cos(x) + c_1}{x}$$

Summary

The solution(s) found are the following

$$y = \frac{x \sin(x) + \cos(x) + c_1}{x} \tag{1}$$

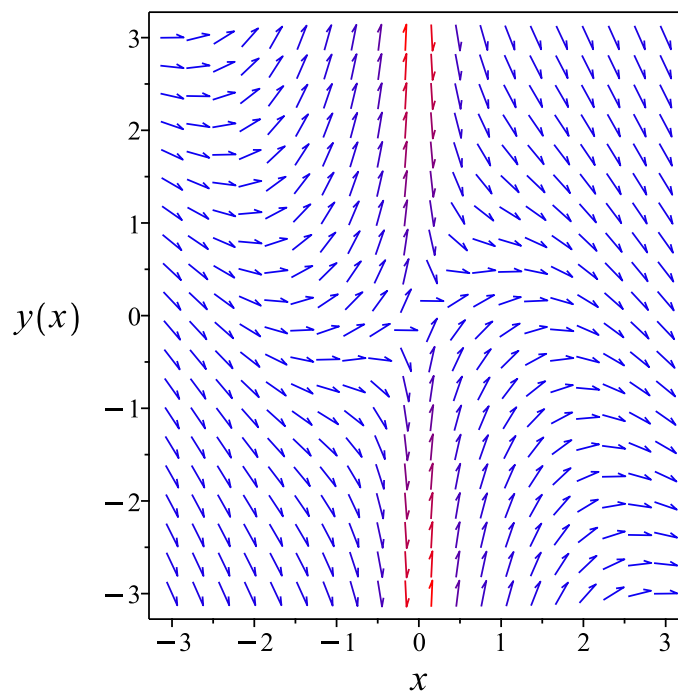


Figure 98: Slope field plot

Verification of solutions

$$y = \frac{x \sin(x) + \cos(x) + c_1}{x}$$

Verified OK.

### 5.6.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-y + \cos(x)x}{x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 49: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{x}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{x}} dy \end{aligned}$$

Which results in

$$S = xy$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{-y + \cos(x)x}{x}$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 0$$

$$S_x = y$$

$$S_y = x$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \cos(x)x \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \cos(R)R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by



integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = \cos(R) + \sin(R) R + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$yx = x \sin(x) + \cos(x) + c_1$$

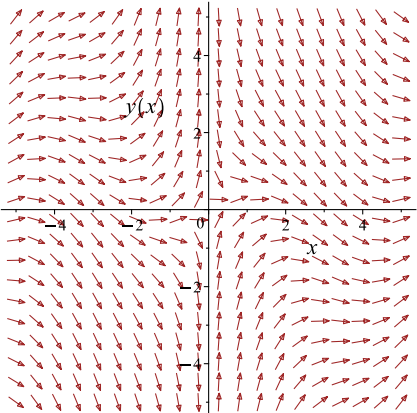
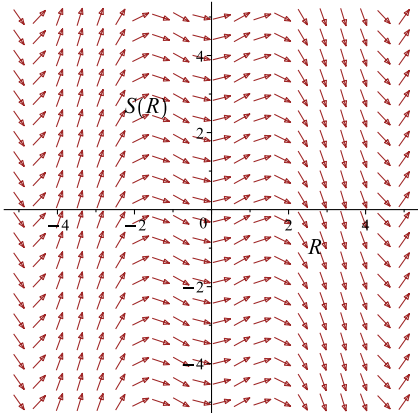
Which simplifies to

$$yx = x \sin(x) + \cos(x) + c_1$$

Which gives

$$y = \frac{x \sin(x) + \cos(x) + c_1}{x}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{-y + \cos(x)x}{x}$ 	$R = x$ $S = xy$	$\frac{dS}{dR} = \cos(R) R$ 

### Summary

The solution(s) found are the following

$$y = \frac{x \sin(x) + \cos(x) + c_1}{x} \quad (1)$$

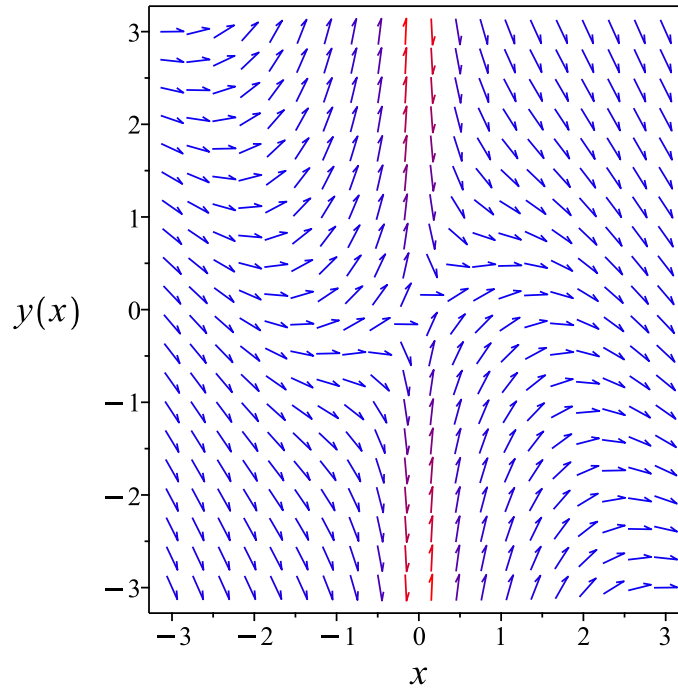


Figure 99: Slope field plot

Verification of solutions

$$y = \frac{x \sin(x) + \cos(x) + c_1}{x}$$

Verified OK.

### 5.6.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(x) dy &= (-y + \cos(x) x) dx \\ (y - \cos(x) x) dx + (x) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= y - \cos(x) x \\ N(x, y) &= x\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y - \cos(x) x) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x) \\ &= 1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int y - \cos(x) x dx \\ \phi &= xy - \cos(x) - x \sin(x) + f(y) \end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = x + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = x$ . Therefore equation (4) becomes

$$x = x + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = xy - \cos(x) - x \sin(x) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = xy - \cos(x) - x \sin(x)$$

The solution becomes

$$y = \frac{x \sin(x) + \cos(x) + c_1}{x}$$

Summary

The solution(s) found are the following

$$y = \frac{x \sin(x) + \cos(x) + c_1}{x} \tag{1}$$

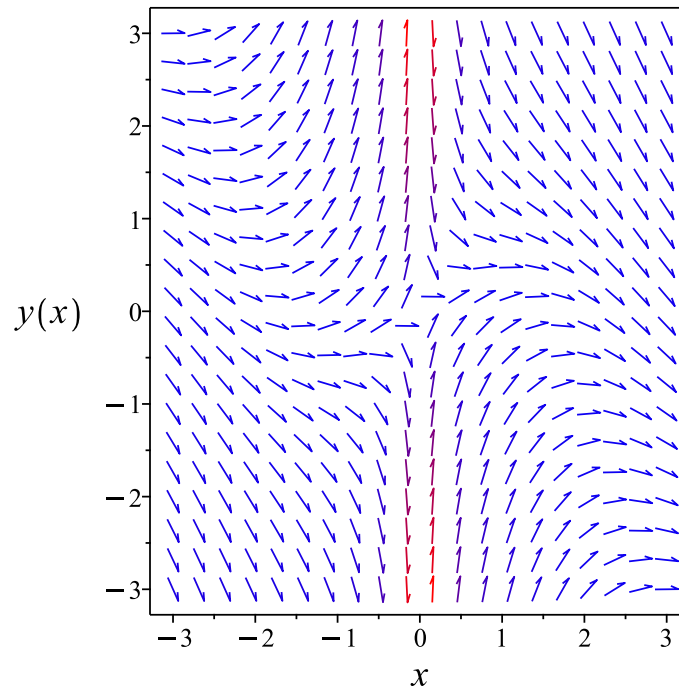


Figure 100: Slope field plot

Verification of solutions

$$y = \frac{x \sin(x) + \cos(x) + c_1}{x}$$

Verified OK.

#### 5.6.4 Maple step by step solution

Let's solve

$$xy' + y = \cos(x) x$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{y}{x} + \cos(x)$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{y}{x} = \cos(x)$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x) \left( y' + \frac{y}{x} \right) = \mu(x) \cos(x)$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x) y)$

$$\mu(x) \left( y' + \frac{y}{x} \right) = \mu'(x) y + \mu(x) y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = \frac{\mu(x)}{x}$$

- Solve to find the integrating factor

$$\mu(x) = x$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x) y) \right) dx = \int \mu(x) \cos(x) dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x) y = \int \mu(x) \cos(x) dx + c_1$$

- Solve for  $y$

$$y = \frac{\int \mu(x) \cos(x) dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = x$

$$y = \frac{\int \cos(x) x dx + c_1}{x}$$

- Evaluate the integrals on the rhs

$$y = \frac{x \sin(x) + \cos(x) + c_1}{x}$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 16

```
dsolve(x*diff(y(x),x)+y(x)=x*cos(x),y(x), singsol=all)
```

$$y(x) = \frac{x \sin(x) + \cos(x) + c_1}{x}$$

### ✓ Solution by Mathematica

Time used: 0.034 (sec). Leaf size: 18

```
DSolve[x*y'[x]+y[x]==x*Cos[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{x \sin(x) + \cos(x) + c_1}{x}$$

## 5.7 problem 9

5.7.1	Solving as homogeneousTypeD2 ode . . . . .	447
5.7.2	Solving as first order ode lie symmetry calculated ode . . . . .	449
5.7.3	Solving as exact ode . . . . .	454

Internal problem ID [3120]

Internal file name [OUTPUT/2612\_Sunday\_June\_05\_2022\_03\_22\_30\_AM\_99694529/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 9.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"exactByInspection", "homogeneousTypeD2", "first\_order\_ode\_lie\_symmetry\_calculated"**

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class B`]]
```

$$(yx - x^2) y' - y^2 = 0$$

### 5.7.1 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$(u(x)x^2 - x^2)(u'(x)x + u(x)) - u(x)^2x^2 = 0$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= \frac{u}{x(u-1)} \end{aligned}$$



Where  $f(x) = \frac{1}{x}$  and  $g(u) = \frac{u}{u-1}$ . Integrating both sides gives

$$\frac{1}{\frac{u}{u-1}} du = \frac{1}{x} dx$$

$$\int \frac{1}{\frac{u}{u-1}} du = \int \frac{1}{x} dx$$

$$u - \ln(u) = \ln(x) + c_2$$

The solution is

$$u(x) - \ln(u(x)) - \ln(x) - c_2 = 0$$

Replacing  $u(x)$  in the above solution by  $\frac{y}{x}$  results in the solution for  $y$  in implicit form

$$\frac{y}{x} - \ln\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0$$

$$\frac{y}{x} - \ln\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0$$

### Summary

The solution(s) found are the following

$$\frac{y}{x} - \ln\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0 \tag{1}$$

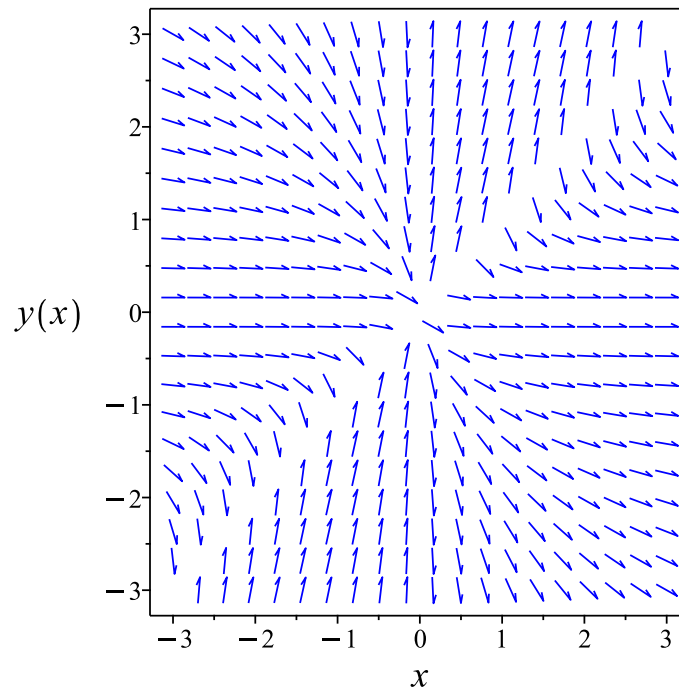


Figure 101: Slope field plot

### Verification of solutions

$$\frac{y}{x} - \ln\left(\frac{y}{x}\right) - \ln(x) - c_2 = 0$$

Verified OK.

### 5.7.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{y^2}{x(y-x)}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$b_2 + \frac{y^2(b_3 - a_2)}{x(y-x)} - \frac{y^4 a_3}{x^2(y-x)^2} - \left( -\frac{y^2}{x^2(y-x)} + \frac{y^2}{x(y-x)^2} \right) (xa_2 + ya_3 + a_1) \quad (\text{5E})$$
$$- \left( \frac{2y}{x(y-x)} - \frac{y^2}{x(y-x)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{x^4 b_2 - x^2 y^2 a_2 + x^2 y^2 b_3 - 2x y^3 a_3 + 2x^2 y b_1 - 2x y^2 a_1 - x y^2 b_1 + y^3 a_1}{x^2 (-y + x)^2} = 0$$

Setting the numerator to zero gives

$$x^4b_2 - x^2y^2a_2 + x^2y^2b_3 - 2xy^3a_3 + 2x^2yb_1 - 2xy^2a_1 - xy^2b_1 + y^3a_1 = 0 \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$-a_2v_1^2v_2^2 - 2a_3v_1v_2^3 + b_2v_1^4 + b_3v_1^2v_2^2 - 2a_1v_1v_2^2 + a_1v_2^3 + 2b_1v_1^2v_2 - b_1v_1v_2^2 = 0 \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$b_2v_1^4 + (b_3 - a_2)v_1^2v_2^2 + 2b_1v_1^2v_2 - 2a_3v_1v_2^3 + (-2a_1 - b_1)v_1v_2^2 + a_1v_2^3 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} a_1 &= 0 \\ b_2 &= 0 \\ -2a_3 &= 0 \\ 2b_1 &= 0 \\ -2a_1 - b_1 &= 0 \\ b_3 - a_2 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= x \\ \eta &= y\end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left( \frac{y^2}{x(y-x)} \right) (x) \\ &= \frac{xy}{-y+x} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left( \xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{xy}{-y+x}} dy\end{aligned}$$

Which results in

$$S = \ln(y) - \frac{y}{x}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y) S_y}{R_x + \omega(x, y) R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y^2}{x(y-x)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{y}{x^2} \\ S_y &= \frac{-y+x}{xy} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \tag{2A}$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{\ln(y)x - y}{x} = c_1$$

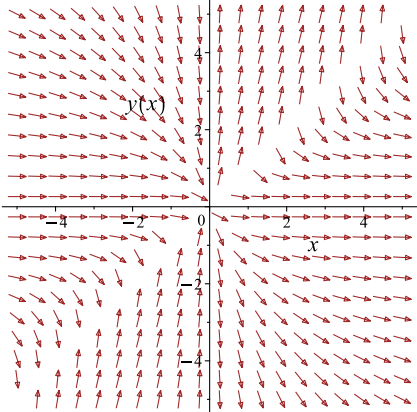
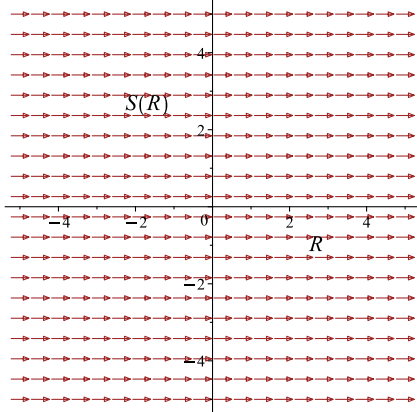
Which simplifies to

$$\frac{\ln(y)x - y}{x} = c_1$$

Which gives

$$y = e^{-\text{LambertW}\left(-\frac{e^{c_1}}{x}\right) + c_1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{y^2}{x(y-x)}$ 	$R = x$ $S = \frac{\ln(y) x - y}{x}$	$\frac{dS}{dR} = 0$ 

### Summary

The solution(s) found are the following

$$y = e^{-\text{LambertW}\left(-\frac{e^{c_1}}{x}\right) + c_1} \quad (1)$$

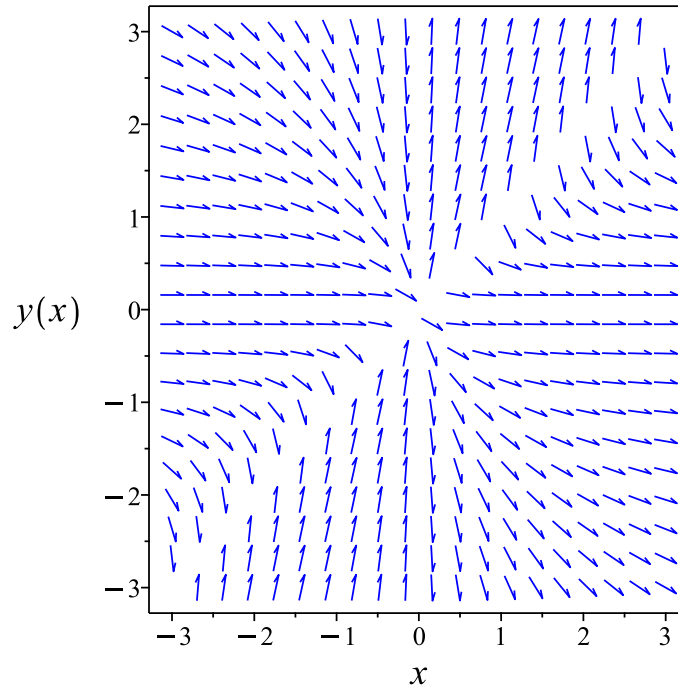


Figure 102: Slope field plot

Verification of solutions

$$y = e^{-\text{LambertW}\left(-\frac{e^{c_1}}{x}\right)+c_1}$$

Verified OK.

**5.7.3 Solving as exact ode**

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(-x^2 + xy) dy &= (y^2) dx \\ (-y^2) dx + (-x^2 + xy) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -y^2 \\ N(x, y) &= -x^2 + xy\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-y^2) \\ &= -2y\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x^2 + xy) \\ &= -2x + y\end{aligned}$$



Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. By inspection  $\frac{1}{x^2y}$  is an integrating factor. Therefore by multiplying  $M = -y^2$  and  $N = yx - x^2$  by this integrating factor the ode becomes exact. The new  $M, N$  are

$$M = -\frac{y}{x^2}$$

$$N = \frac{yx - x^2}{x^2y}$$

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$

$$\frac{\partial \phi}{\partial y} = N$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} \left(\frac{-x^2 + xy}{x^2y}\right) dy &= \left(\frac{y}{x^2}\right) dx \\ \left(-\frac{y}{x^2}\right) dx + \left(\frac{-x^2 + xy}{x^2y}\right) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -\frac{y}{x^2} \\ N(x, y) &= \frac{-x^2 + xy}{x^2y} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{y}{x^2}\right) \\ &= -\frac{1}{x^2} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{-x^2 + xy}{x^2y}\right) \\ &= -\frac{1}{x^2} \end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -\frac{y}{x^2} dx \\ \phi &= \frac{y}{x} + f(y) \end{aligned} \tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = \frac{1}{x} + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \frac{-x^2 + xy}{x^2y}$ . Therefore equation (4) becomes

$$\frac{-x^2 + xy}{x^2y} = \frac{1}{x} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = -\frac{1}{y}$$

Integrating the above w.r.t  $y$  gives

$$\int f'(y) dy = \int \left(-\frac{1}{y}\right) dy$$

$$f(y) = -\ln(y) + c_1$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = \frac{y}{x} - \ln(y) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \frac{y}{x} - \ln(y)$$

The solution becomes

$$y = e^{-\text{LambertW}\left(-\frac{e^{-c_1}}{x}\right) - c_1}$$

### Summary

The solution(s) found are the following

$$y = e^{-\text{LambertW}\left(-\frac{e^{-c_1}}{x}\right) - c_1} \quad (1)$$

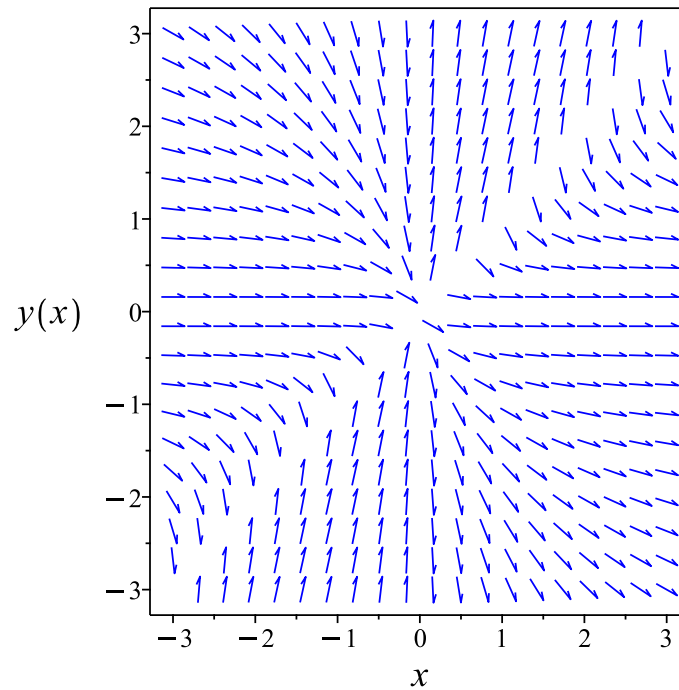


Figure 103: Slope field plot

Verification of solutions

$$y = e^{-\text{LambertW}\left(-\frac{e^{-c_1}}{x}\right) - c_1}$$

Verified OK.

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 17

```
dsolve((x*y(x)-x^2)*diff(y(x),x)=y(x)^2,y(x), singsol=all)
```

$$y(x) = -x \operatorname{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

✓ Solution by Mathematica

Time used: 2.286 (sec). Leaf size: 25

```
DSolve[(x*y[x]-x^2)*y'[x]==y[x]^2,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -xW\left(-\frac{e^{-c_1}}{x}\right)$$
$$y(x) \rightarrow 0$$

## 5.8 problem 10

5.8.1 Solving as exact ode . . . . .	461
5.8.2 Maple step by step solution . . . . .	464

Internal problem ID [3121]

Internal file name [OUTPUT/2613\_Sunday\_June\_05\_2022\_03\_22\_33\_AM\_29737813/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 10.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exact"

Maple gives the following as the ode type

```
[_exact , [_1st_order , ` _with_symmetry_[F(x),G(x)*y+H(x)] `]]
```

$$(e^x - 3y^2x^2) y' + e^xy - 2xy^3 = 0$$

### 5.8.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (e^x - 3y^2x^2) dy &= (-e^xy + 2xy^3) dx \\ (-2xy^3 + e^xy) dx + (e^x - 3y^2x^2) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -2xy^3 + e^xy \\ N(x, y) &= e^x - 3y^2x^2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (-2xy^3 + e^xy) \\ &= e^x - 6xy^2 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (e^x - 3y^2x^2) \\ &= e^x - 6xy^2 \end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -2xy^3 + e^x y dx \\ \phi &= y(-y^2x^2 + e^x) + f(y)\end{aligned}\tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = e^x - 3y^2x^2 + f'(y)\tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = e^x - 3y^2x^2$ . Therefore equation (4) becomes

$$e^x - 3y^2x^2 = e^x - 3y^2x^2 + f'(y)\tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = y(-y^2x^2 + e^x) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = y(-y^2x^2 + e^x)$$

### Summary

The solution(s) found are the following

$$y(-y^2x^2 + e^x) = c_1\tag{1}$$



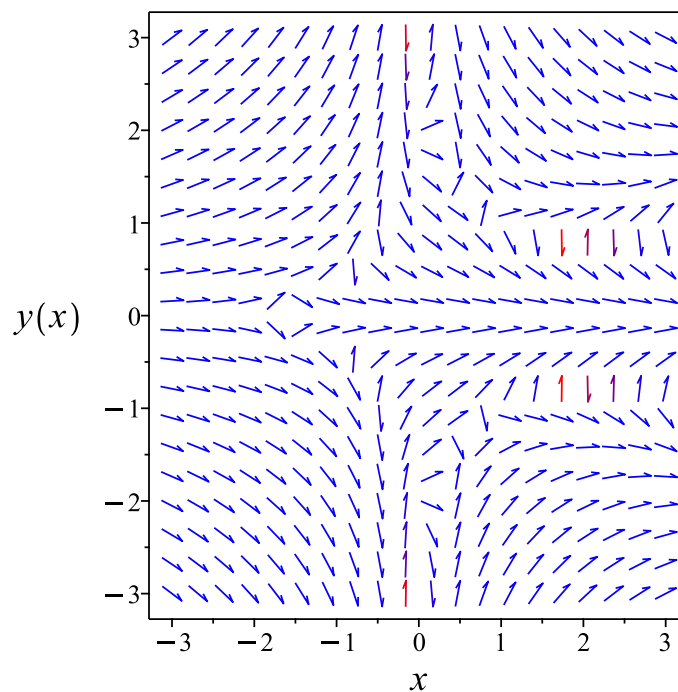


Figure 104: Slope field plot

### Verification of solutions

$$y(-y^2x^2 + e^x) = c_1$$

Verified OK.

### 5.8.2 Maple step by step solution

Let's solve

$$(e^x - 3y^2x^2) y' + e^xy - 2xy^3 = 0$$

- Highest derivative means the order of the ODE is 1  
 $y'$
- Check if ODE is exact
  - ODE is exact if the lhs is the total derivative of a  $C^2$  function  
 $F'(x, y) = 0$
  - Compute derivative of lhs  
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$

- Evaluate derivatives

$$e^x - 6xy^2 = e^x - 6xy^2$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$\left[ F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$

- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$

$$F(x, y) = \int (-2xy^3 + e^xy) dx + f_1(y)$$

- Evaluate integral

$$F(x, y) = e^xy - y^3x^2 + f_1(y)$$

- Take derivative of  $F(x, y)$  with respect to  $y$

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$

- Compute derivative

$$e^x - 3y^2x^2 = e^x - 3y^2x^2 + \frac{d}{dy} f_1(y)$$

- Isolate for  $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = 0$$

- Solve for  $f_1(y)$

$$f_1(y) = 0$$

- Substitute  $f_1(y)$  into equation for  $F(x, y)$

$$F(x, y) = e^xy - y^3x^2$$

- Substitute  $F(x, y)$  into the solution of the ODE

$$e^xy - y^3x^2 = c_1$$

- Solve for  $y$

$$\left\{ y = \frac{\left( -108c_1x + 12\sqrt{81c_1^2x^2 - 12(e^x)^3} \right)^{\frac{1}{3}}}{6x} + \frac{2e^x}{x \left( -108c_1x + 12\sqrt{81c_1^2x^2 - 12(e^x)^3} \right)^{\frac{1}{3}}}, y = -\frac{\left( -108c_1x + 12\sqrt{81c_1^2x^2 - 12(e^x)^3} \right)^{\frac{1}{3}}}{12x} \right.$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

### ✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 278

```
dsolve((exp(x)-3*x^2*y(x)^2)*diff(y(x),x)+y(x)*exp(x)=2*x*y(x)^3,y(x), singsol=all)
```

$$y(x) = \frac{\left(108c_1x + 12\sqrt{81c_1^2x^2 - 12e^{3x}}\right)^{\frac{2}{3}} + 12e^x}{6\left(108c_1x + 12\sqrt{81c_1^2x^2 - 12e^{3x}}\right)^{\frac{1}{3}}x}$$

$$y(x) = \frac{-i\sqrt{3}\left(108c_1x + 12\sqrt{81c_1^2x^2 - 12e^{3x}}\right)^{\frac{2}{3}} + 12ie^x\sqrt{3} - \left(108c_1x + 12\sqrt{81c_1^2x^2 - 12e^{3x}}\right)^{\frac{2}{3}} - 12e^x}{12\left(108c_1x + 12\sqrt{81c_1^2x^2 - 12e^{3x}}\right)^{\frac{1}{3}}x}$$

$$y(x) = \frac{-i\sqrt{3}\left(108c_1x + 12\sqrt{81c_1^2x^2 - 12e^{3x}}\right)^{\frac{2}{3}} + 12ie^x\sqrt{3} + \left(108c_1x + 12\sqrt{81c_1^2x^2 - 12e^{3x}}\right)^{\frac{2}{3}} + 12e^x}{12\left(108c_1x + 12\sqrt{81c_1^2x^2 - 12e^{3x}}\right)^{\frac{1}{3}}x}$$

✓ Solution by Mathematica

Time used: 48.841 (sec). Leaf size: 364

`DSolve[(Exp[x]-3*x^2*y[x]^2)*y'[x]+y[x]*Exp[x]==2*x*y[x]^3,y[x],x,IncludeSingularSolutions`

$$y(x) \rightarrow \frac{2\sqrt[3]{3}e^x x^2 + \sqrt[3]{2}(9c_1 x^4 + \sqrt{-12e^{3x}x^6 + 81c_1^2 x^8})^{2/3}}{6^{2/3}x^2 \sqrt[3]{9c_1 x^4 + \sqrt{-12e^{3x}x^6 + 81c_1^2 x^8}}}$$

$$y(x) \rightarrow \frac{i(\sqrt{3} + i) \sqrt[3]{9c_1 x^4 + \sqrt{-12e^{3x}x^6 + 81c_1^2 x^8}}}{2\sqrt[3]{2}3^{2/3}x^2 (\sqrt{3} + 3i) e^x}$$

$$- \frac{2^{2/3}3^{5/6} \sqrt[3]{9c_1 x^4 + \sqrt{-12e^{3x}x^6 + 81c_1^2 x^8}}}{2\sqrt[3]{2}3^{2/3}x^2 (\sqrt{3} + 3i) e^x}$$

$$y(x) \rightarrow \frac{(-1 - i\sqrt{3}) \sqrt[3]{9c_1 x^4 + \sqrt{-12e^{3x}x^6 + 81c_1^2 x^8}}}{2\sqrt[3]{2}3^{2/3}x^2 (\sqrt{3} - 3i) e^x}$$

$$- \frac{2^{2/3}3^{5/6} \sqrt[3]{9c_1 x^4 + \sqrt{-12e^{3x}x^6 + 81c_1^2 x^8}}}{2\sqrt[3]{2}3^{2/3}x^2 (\sqrt{3} - 3i) e^x}$$

## 5.9 problem 12

5.9.1	Solving as linear ode . . . . .	468
5.9.2	Solving as homogeneousTypeD2 ode . . . . .	470
5.9.3	Solving as first order ode lie symmetry lookup ode . . . . .	471
5.9.4	Solving as exact ode . . . . .	475
5.9.5	Maple step by step solution . . . . .	480

Internal problem ID [3122]

Internal file name [OUTPUT/2614\_Sunday\_June\_05\_2022\_03\_22\_37\_AM\_79544205/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 12.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"linear"**, **"homogeneousTypeD2"**, **"exactWithIntegrationFactor"**, **"first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$-xy' + y = -x^2$$

### 5.9.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = -\frac{1}{x}$$

$$q(x) = x$$

Hence the ode is

$$y' - \frac{y}{x} = x$$

The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int -\frac{1}{x} dx} \\ &= \frac{1}{x}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu)(x) \\ \frac{d}{dx}\left(\frac{y}{x}\right) &= \left(\frac{1}{x}\right)(x) \\ d\left(\frac{y}{x}\right) &= dx\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y}{x} &= \int dx \\ \frac{y}{x} &= x + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = \frac{1}{x}$  results in

$$y = c_1 x + x^2$$

which simplifies to

$$y = x(x + c_1)$$

### Summary

The solution(s) found are the following

$$y = x(x + c_1) \tag{1}$$

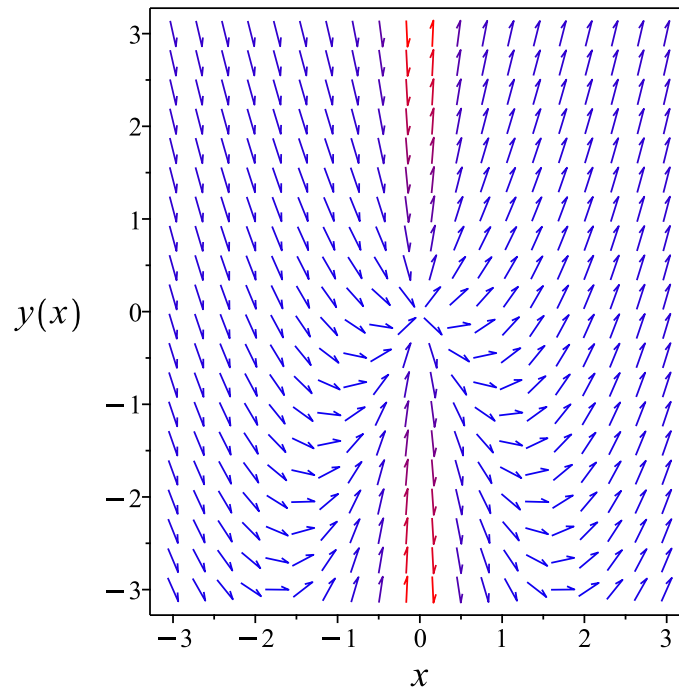


Figure 105: Slope field plot

Verification of solutions

$$y = x(x + c_1)$$

Verified OK.

### 5.9.2 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$-x(u'(x)x + u(x)) + u(x)x = -x^2$$

Integrating both sides gives

$$\begin{aligned} u(x) &= \int 1 \, dx \\ &= x + c_2 \end{aligned}$$

Therefore the solution  $y$  is

$$\begin{aligned} y &= ux \\ &= x(x + c_2) \end{aligned}$$

### Summary

The solution(s) found are the following

$$y = x(x + c_2) \tag{1}$$

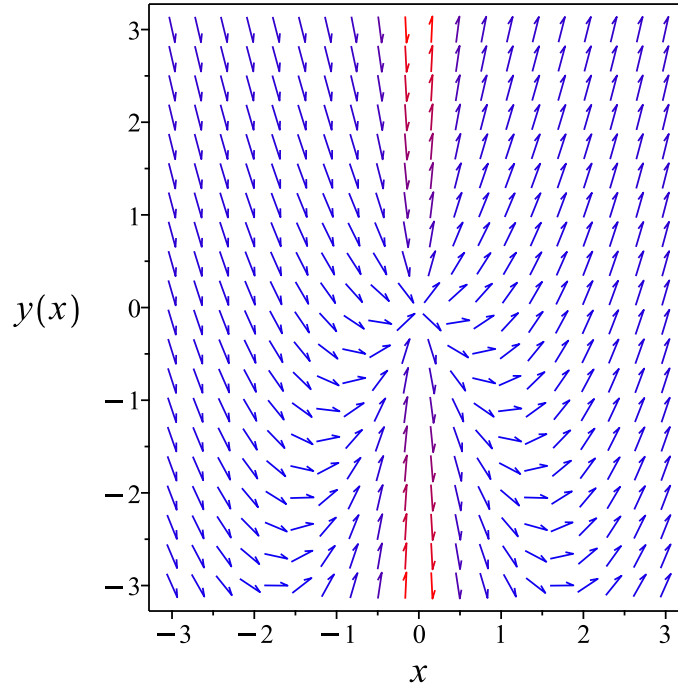


Figure 106: Slope field plot

### Verification of solutions

$$y = x(x + c_2)$$

Verified OK.

### **5.9.3 Solving as first order ode lie symmetry lookup ode**

Writing the ode as

$$y' = \frac{x^2 + y}{x}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \tag{A}$$



The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 53: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= x\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{x} dy \end{aligned}$$

Which results in

$$S = \frac{y}{x}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{x^2 + y}{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -\frac{y}{x^2} \\ S_y &= \frac{1}{x} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 1 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = R + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$\frac{y}{x} = x + c_1$$

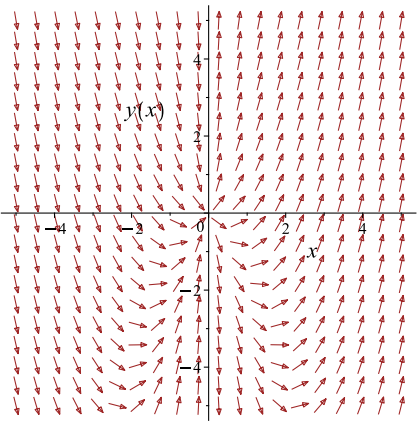
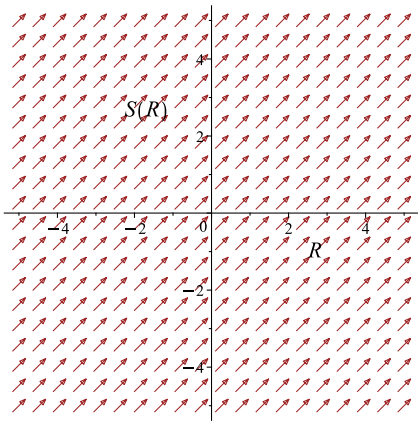
Which simplifies to

$$\frac{y}{x} = x + c_1$$

Which gives

$$y = x(x + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{x^2 + y}{x}$ 	$R = x$ $S = \frac{y}{x}$	$\frac{dS}{dR} = 1$ 

### Summary

The solution(s) found are the following

$$y = x(x + c_1) \tag{1}$$

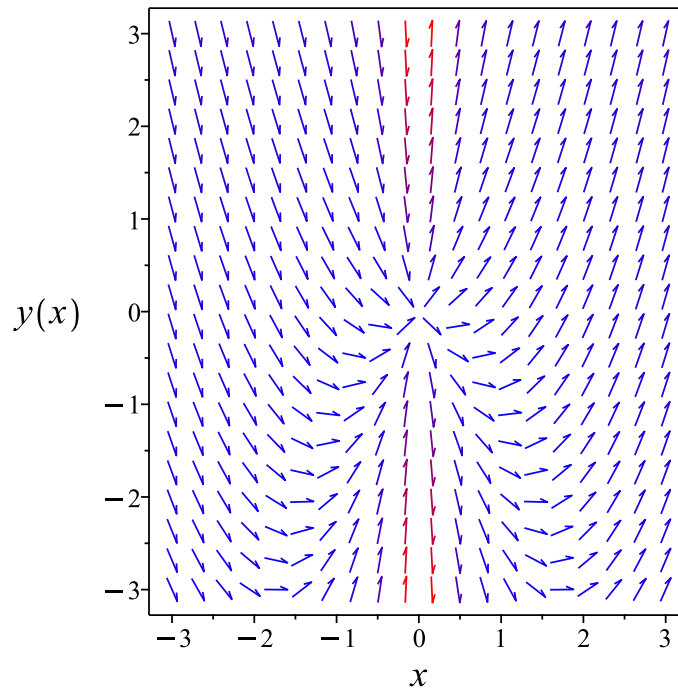


Figure 107: Slope field plot

Verification of solutions

$$y = x(x + c_1)$$

Verified OK.

#### 5.9.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(-x) dy &= (-x^2 - y) dx \\ (x^2 + y) dx + (-x) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= x^2 + y \\ N(x, y) &= -x\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(x^2 + y) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x) \\ &= -1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= -\frac{1}{x} ((1) - (-1)) \\ &= -\frac{2}{x} \end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned} \mu &= e^{\int A dx} \\ &= e^{\int -\frac{2}{x} dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-2 \ln(x)} \\ &= \frac{1}{x^2} \end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{1}{x^2} (x^2 + y) \\ &= \frac{x^2 + y}{x^2} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \frac{1}{x^2} (-x) \\ &= -\frac{1}{x} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left( \frac{x^2 + y}{x^2} \right) + \left( -\frac{1}{x} \right) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \overline{M} dx$$
$$\int \frac{\partial \phi}{\partial x} dx = \int \frac{x^2 + y}{x^2} dx$$
$$\phi = x - \frac{y}{x} + f(y) \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = -\frac{1}{x} + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = -\frac{1}{x}$ . Therefore equation (4) becomes

$$-\frac{1}{x} = -\frac{1}{x} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = x - \frac{y}{x} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = x - \frac{y}{x}$$

The solution becomes

$$y = -(-x + c_1) x$$

Summary

The solution(s) found are the following

$$y = -(-x + c_1) x \tag{1}$$

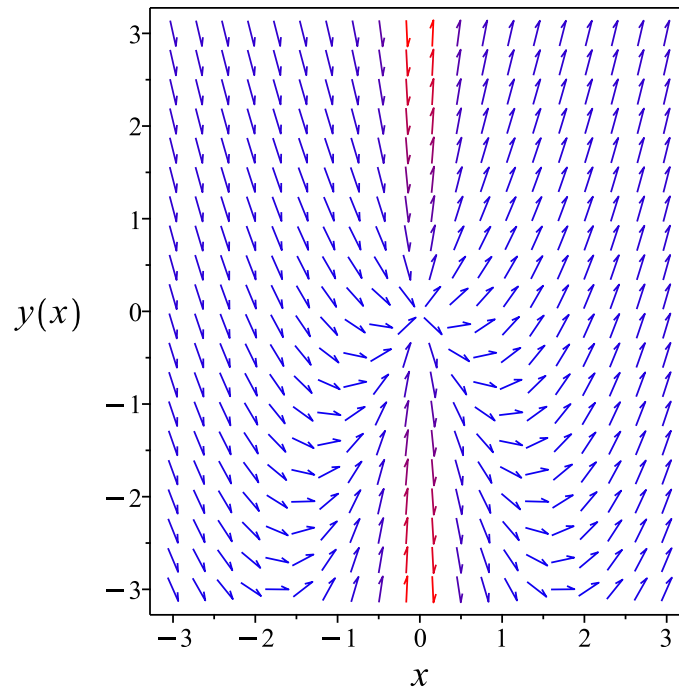


Figure 108: Slope field plot

Verification of solutions

$$y = -(-x + c_1) x$$

Verified OK.



### 5.9.5 Maple step by step solution

Let's solve

$$-xy' + y = -x^2$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = \frac{y}{x} + x$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' - \frac{y}{x} = x$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x) \left( y' - \frac{y}{x} \right) = \mu(x) x$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x) y)$

$$\mu(x) \left( y' - \frac{y}{x} \right) = \mu'(x) y + \mu(x) y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = -\frac{\mu(x)}{x}$$

- Solve to find the integrating factor

$$\mu(x) = \frac{1}{x}$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x) y) \right) dx = \int \mu(x) x dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x) y = \int \mu(x) x dx + c_1$$

- Solve for  $y$

$$y = \frac{\int \mu(x) x dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = \frac{1}{x}$

$$y = x \left( \int 1 dx + c_1 \right)$$

- Evaluate the integrals on the rhs

$$y = x(x + c_1)$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 9

```
dsolve((x^2+y(x))=x*diff(y(x),x),y(x), singsol=all)
```

$$y(x) = (c_1 + x)x$$

### ✓ Solution by Mathematica

Time used: 0.025 (sec). Leaf size: 11

```
DSolve[(x^2+y[x])=x*y'[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow x(x + c_1)$$

## 5.10 problem 13

5.10.1 Solving as linear ode . . . . .	482
5.10.2 Solving as first order ode lie symmetry lookup ode . . . . .	484
5.10.3 Solving as exact ode . . . . .	488
5.10.4 Maple step by step solution . . . . .	492

Internal problem ID [3123]

Internal file name [OUTPUT/2615\_Sunday\_June\_05\_2022\_03\_22\_39\_AM\_13834837/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 13.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exact", "linear", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_linear]

$$xy' + y = x^2 \cos(x)$$

### 5.10.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{1}{x}$$
$$q(x) = \cos(x) x$$

Hence the ode is

$$y' + \frac{y}{x} = \cos(x) x$$

The integrating factor  $\mu$  is

$$\mu = e^{\int \frac{1}{x} dx}$$
$$= x$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) (\cos(x) x) \\ \frac{d}{dx}(xy) &= (x) (\cos(x) x) \\ d(xy) &= (x^2 \cos(x)) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}xy &= \int x^2 \cos(x) dx \\ xy &= x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = x$  results in

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x}{x} + \frac{c_1}{x}$$

which simplifies to

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1}{x}$$

Summary

The solution(s) found are the following

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1}{x} \tag{1}$$

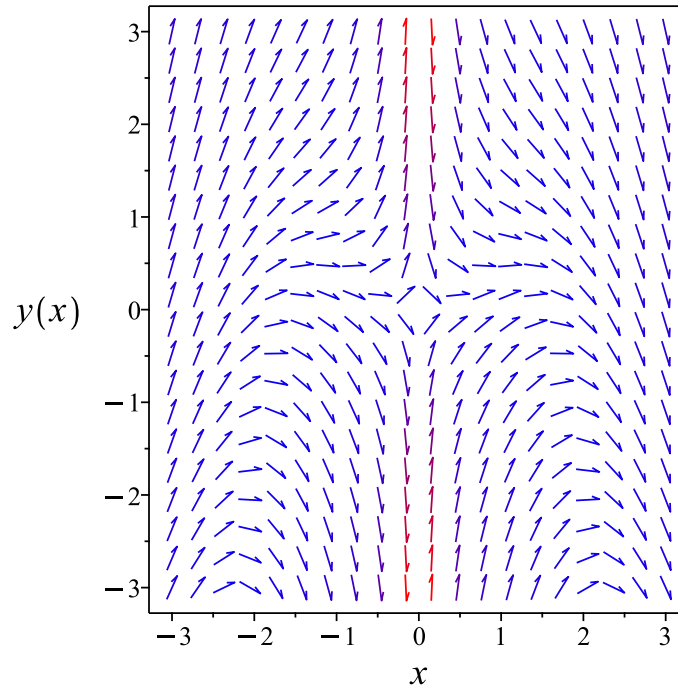


Figure 109: Slope field plot

Verification of solutions

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x)x + c_1}{x}$$

Verified OK.

### 5.10.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = \frac{-y + x^2 \cos(x)}{x}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 56: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{x}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right)S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the

canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{x}} dy \end{aligned}$$

Which results in

$$S = xy$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{-y + x^2 \cos(x)}{x}$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 0$$

$$S_x = y$$

$$S_y = x$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = x^2 \cos(x) \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R^2 \cos(R)$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = R^2 \sin(R) - 2 \sin(R) + 2R \cos(R) + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$yx = x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1$$

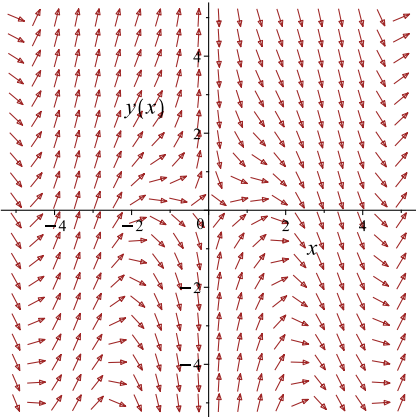
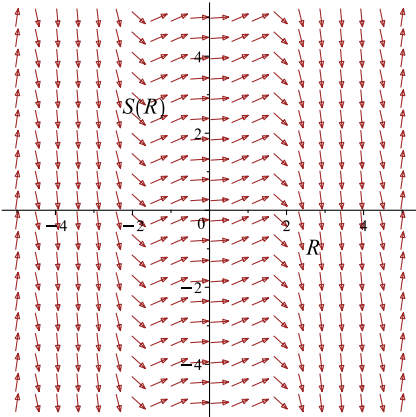
Which simplifies to

$$yx = x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1$$

Which gives

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1}{x}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{-y + x^2 \cos(x)}{x}$ 	$R = x$ $S = xy$	$\frac{dS}{dR} = R^2 \cos(R)$ 

### Summary

The solution(s) found are the following

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1}{x} \quad (1)$$



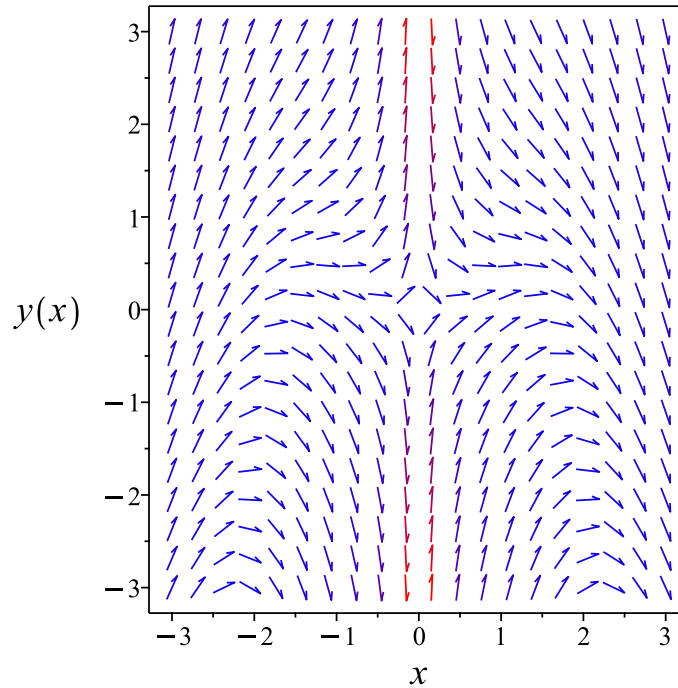


Figure 110: Slope field plot

Verification of solutions

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1}{x}$$

Verified OK.

### 5.10.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(x) dy &= (-y + x^2 \cos(x)) dx \\ (y - x^2 \cos(x)) dx + (x) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= y - x^2 \cos(x) \\ N(x, y) &= x\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y - x^2 \cos(x)) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x) \\ &= 1\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int y - x^2 \cos(x) dx \\ \phi &= (-x^2 + 2) \sin(x) + x(-2 \cos(x) + y) + f(y) \end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = x + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = x$ . Therefore equation (4) becomes

$$x = x + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = (-x^2 + 2) \sin(x) + x(-2 \cos(x) + y) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = (-x^2 + 2) \sin(x) + x(-2 \cos(x) + y)$$

The solution becomes

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1}{x}$$

### Summary

The solution(s) found are the following

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1}{x} \quad (1)$$

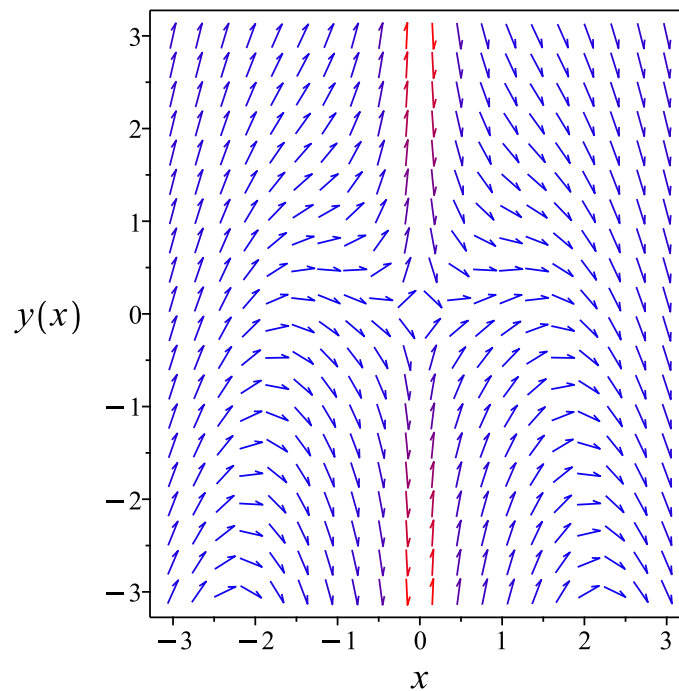


Figure 111: Slope field plot

### Verification of solutions

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1}{x}$$

Verified OK.

### 5.10.4 Maple step by step solution

Let's solve

$$xy' + y = x^2 \cos(x)$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{y}{x} + \cos(x) x$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{y}{x} = \cos(x) x$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x) \left( y' + \frac{y}{x} \right) = \mu(x) \cos(x) x$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x) y)$

$$\mu(x) \left( y' + \frac{y}{x} \right) = \mu'(x) y + \mu(x) y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = \frac{\mu(x)}{x}$$

- Solve to find the integrating factor

$$\mu(x) = x$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x) y) \right) dx = \int \mu(x) \cos(x) x dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x) y = \int \mu(x) \cos(x) x dx + c_1$$

- Solve for  $y$

$$y = \frac{\int \mu(x) \cos(x) x dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = x$

$$y = \frac{\int x^2 \cos(x) dx + c_1}{x}$$

- Evaluate the integrals on the rhs

$$y = \frac{x^2 \sin(x) - 2 \sin(x) + 2 \cos(x) x + c_1}{x}$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 25

```
dsolve(x*diff(y(x),x)+y(x)=x^2*cos(x),y(x), singsol=all)
```

$$y(x) = \frac{\sin(x)x^2 - 2\sin(x) + 2x\cos(x) + c_1}{x}$$

### ✓ Solution by Mathematica

Time used: 0.066 (sec). Leaf size: 25

```
DSolve[x*y'[x]+y[x]==x^2*Cos[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{(x^2 - 2)\sin(x) + 2x\cos(x) + c_1}{x}$$

## 5.11 problem 14

5.11.1 Solving as first order ode lie symmetry calculated ode . . . . . 494

Internal problem ID [3124]

Internal file name [OUTPUT/2616\_Sunday\_June\_05\_2022\_03\_22\_42\_AM\_16397768/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 14.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**first\_order\_ode\_lie\_symmetry\_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class C`], _rational, [_Abel, `2nd type`, `class A`]]
```

$$4y + (3x + 2y + 2)y' = -3 - 6x$$

### 5.11.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{6x + 4y + 3}{3x + 2y + 2}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} b_2 - \frac{(6x+4y+3)(b_3-a_2)}{3x+2y+2} - \frac{(6x+4y+3)^2 a_3}{(3x+2y+2)^2} \\ - \left( -\frac{6}{3x+2y+2} + \frac{18x+12y+9}{(3x+2y+2)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left( -\frac{4}{3x+2y+2} + \frac{12x+8y+6}{(3x+2y+2)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{18x^2a_2 - 36x^2a_3 + 9x^2b_2 - 18x^2b_3 + 24xya_2 - 48xya_3 + 12xyb_2 - 24xyb_3 + 8y^2a_2 - 16y^2a_3 + 4y^2b_2 - 8y^2b_3 + 24xa_2 - 36xa_3 + 14xb_2 - 21xb_3 + 14ya_2 - 21ya_3 + 8yb_2 - 12yb_3 + 3a_1 + 6a_2 - 9a_3 + 2b_1 + 4b_2 - 6b_3}{(3x+2y+2)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} 18x^2a_2 - 36x^2a_3 + 9x^2b_2 - 18x^2b_3 + 24xya_2 - 48xya_3 + 12xyb_2 - 24xyb_3 \\ + 8y^2a_2 - 16y^2a_3 + 4y^2b_2 - 8y^2b_3 + 24xa_2 - 36xa_3 + 14xb_2 - 21xb_3 \\ + 14ya_2 - 21ya_3 + 8yb_2 - 12yb_3 + 3a_1 + 6a_2 - 9a_3 + 2b_1 + 4b_2 - 6b_3 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} 18a_2v_1^2 + 24a_2v_1v_2 + 8a_2v_2^2 - 36a_3v_1^2 - 48a_3v_1v_2 - 16a_3v_2^2 + 9b_2v_1^2 + 12b_2v_1v_2 \\ + 4b_2v_2^2 - 18b_3v_1^2 - 24b_3v_1v_2 - 8b_3v_2^2 + 24a_2v_1 + 14a_2v_2 - 36a_3v_1 - 21a_3v_2 \\ + 14b_2v_1 + 8b_2v_2 - 21b_3v_1 - 12b_3v_2 + 3a_1 + 6a_2 - 9a_3 + 2b_1 + 4b_2 - 6b_3 = 0 \end{aligned} \quad (7E)$$



Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & (18a_2 - 36a_3 + 9b_2 - 18b_3) v_1^2 + (24a_2 - 48a_3 + 12b_2 - 24b_3) v_1 v_2 \\ & + (24a_2 - 36a_3 + 14b_2 - 21b_3) v_1 + (8a_2 - 16a_3 + 4b_2 - 8b_3) v_2^2 \\ & + (14a_2 - 21a_3 + 8b_2 - 12b_3) v_2 + 3a_1 + 6a_2 - 9a_3 + 2b_1 + 4b_2 - 6b_3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 8a_2 - 16a_3 + 4b_2 - 8b_3 &= 0 \\ 14a_2 - 21a_3 + 8b_2 - 12b_3 &= 0 \\ 18a_2 - 36a_3 + 9b_2 - 18b_3 &= 0 \\ 24a_2 - 48a_3 + 12b_2 - 24b_3 &= 0 \\ 24a_2 - 36a_3 + 14b_2 - 21b_3 &= 0 \\ 3a_1 + 6a_2 - 9a_3 + 2b_1 + 4b_2 - 6b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= a_1 \\ a_2 &= \frac{3a_3}{2} \\ a_3 &= a_3 \\ b_1 &= -\frac{3a_1}{2} \\ b_2 &= -3a_3 \\ b_3 &= -2a_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= 1 \\ \eta &= -\frac{3}{2} \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= -\frac{3}{2} - \left( -\frac{6x + 4y + 3}{3x + 2y + 2} \right) \quad (1) \\ &= \frac{3x + 2y}{6x + 4y + 4} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left( \xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{3x+2y}{6x+4y+4}} dy\end{aligned}$$

Which results in

$$S = 2y + 2 \ln(3x + 2y)$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{6x + 4y + 3}{3x + 2y + 2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= \frac{6}{3x + 2y} \\S_y &= 2 + \frac{4}{3x + 2y}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -4 \tag{2A}$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -4$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -4R + c_1 \tag{4}$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$2y + 2 \ln(3x + 2y) = -4x + c_1$$

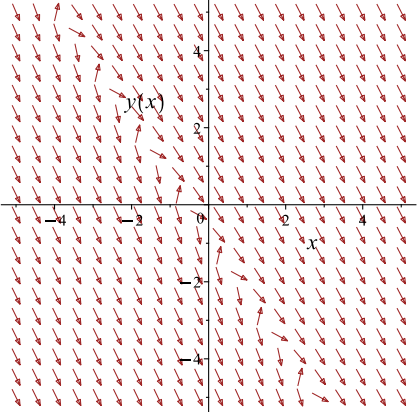
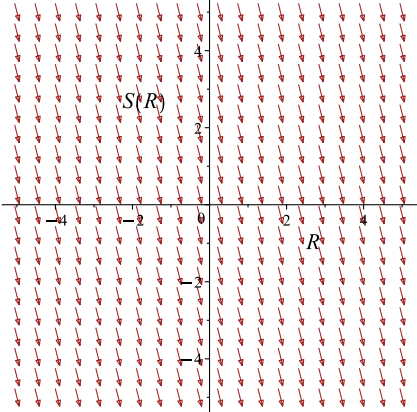
Which simplifies to

$$2y + 2 \ln(3x + 2y) = -4x + c_1$$

Which gives

$$y = -\frac{3x}{2} + \text{LambertW}\left(\frac{e^{-\frac{x}{2} + \frac{c_1}{2}}}{2}\right)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{6x+4y+3}{3x+2y+2}$ 	$R = x$ $S = 2y + 2 \ln(3x + 2y)$	$\frac{dS}{dR} = -4$ 

Summary

The solution(s) found are the following

$$y = -\frac{3x}{2} + \text{LambertW}\left(\frac{e^{-\frac{x}{2} + \frac{c_1}{2}}}{2}\right) \tag{1}$$

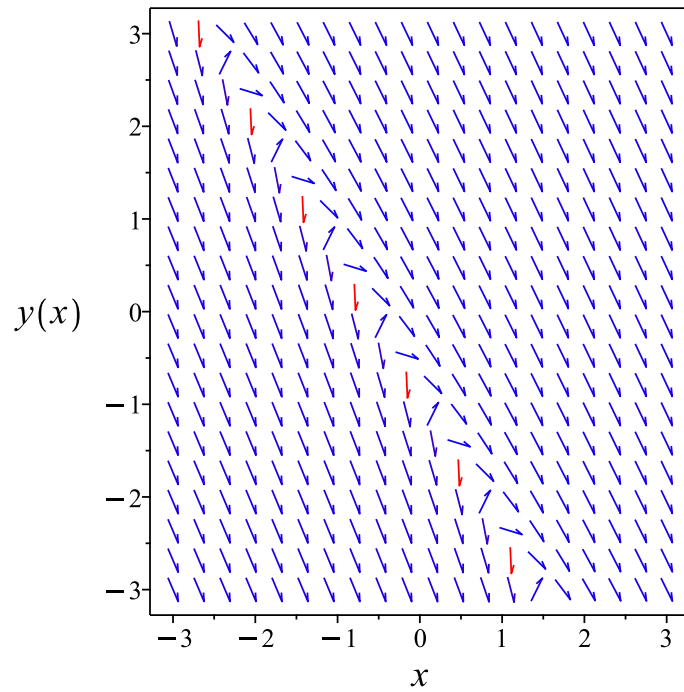


Figure 112: Slope field plot

Verification of solutions

$$y = -\frac{3x}{2} + \text{LambertW}\left(\frac{e^{-\frac{x}{2} + \frac{c_1}{2}}}{2}\right)$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous C  
1st order, trying the canonical coordinates of the invariance group  
<- 1st order, canonical coordinates successful  
<- homogeneous successful`
```

### ✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 15

```
dsolve((6*x+4*y(x)+3)+(3*x+2*y(x)+2)*diff(y(x),x)=0,y(x), singsol=all)
```

$$y(x) = -\frac{3x}{2} + \text{LambertW}(c_1 e^{-\frac{x}{2}})$$

### ✓ Solution by Mathematica

Time used: 4.333 (sec). Leaf size: 34

```
DSolve[(6*x+4*y[x]+3)+(3*x+2*y[x]+2)*y'[x]==0,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -\frac{3x}{2} + W(-e^{-\frac{x}{2}-1+c_1})$$
$$y(x) \rightarrow -\frac{3x}{2}$$

## 5.12 problem 15

5.12.1 Solving as first order ode lie symmetry calculated ode . . . . .	502
5.12.2 Solving as exact ode . . . . .	508
5.12.3 Maple step by step solution . . . . .	512

Internal problem ID [3125]

Internal file name [OUTPUT/2617\_Sunday\_June\_05\_2022\_03\_22\_44\_AM\_15140119/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 15.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"exact", "first\_order\_ode\_lie\_symmetry\_calculated"**

Maple gives the following as the ode type

```
[[_1st_order , _with_linear_symmetries] , _exact]
```

$$\cos(y+x) - x \sin(y+x) - x \sin(y+x) y' = 0$$

### 5.12.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{-x \sin(y+x) + \cos(y+x)}{x \sin(y+x)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \tag{1E}$$

$$\eta = xb_2 + yb_3 + b_1 \tag{2E}$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} & b_2 + \frac{(-x \sin(y+x) + \cos(y+x))(b_3 - a_2)}{x \sin(y+x)} \\ & - \frac{(-x \sin(y+x) + \cos(y+x))^2 a_3}{x^2 \sin(y+x)^2} \\ & - \left( \frac{-2 \sin(y+x) - x \cos(y+x)}{x \sin(y+x)} - \frac{-x \sin(y+x) + \cos(y+x)}{x^2 \sin(y+x)} \right. \\ & \left. - \frac{(-x \sin(y+x) + \cos(y+x)) \cos(y+x)}{x \sin(y+x)^2} \right) (xa_2 + ya_3 + a_1) \\ & - \left( \frac{-x \cos(y+x) - \sin(y+x)}{x \sin(y+x)} \right. \\ & \left. - \frac{(-x \sin(y+x) + \cos(y+x)) \cos(y+x)}{x \sin(y+x)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\begin{aligned} & \cos(y+x)^2 x^2 a_2 + \cos(y+x)^2 x^2 b_2 + \cos(y+x)^2 xy a_3 + \cos(y+x)^2 xy b_3 + 2 \sin(y+x)^2 x^2 a_2 - \sin(y+x)^2 x^2 a_3 \\ & - \sin(y+x)^2 x^2 b_3 + \sin(y+x)^2 xy a_3 + \sin(y+x)^2 xy b_3 \\ & + \cos(y+x)^2 xa_1 + \cos(y+x)^2 xb_1 + 2 \cos(y+x) \sin(y+x) xa_3 \\ & + \cos(y+x) \sin(y+x) xb_3 + \cos(y+x) \sin(y+x) ya_3 + \sin(y+x)^2 xa_1 \\ & + \sin(y+x)^2 xb_1 - \cos(y+x)^2 a_3 + \cos(y+x) \sin(y+x) a_1 = 0 \end{aligned}$$

Setting the numerator to zero gives

$$\begin{aligned} & \cos(y+x)^2 x^2 a_2 + \cos(y+x)^2 x^2 b_2 + \cos(y+x)^2 xy a_3 + \cos(y+x)^2 xy b_3 \\ & + 2 \sin(y+x)^2 x^2 a_2 - \sin(y+x)^2 x^2 a_3 + 2b_2 x^2 \sin(y+x)^2 \\ & - \sin(y+x)^2 x^2 b_3 + \sin(y+x)^2 xy a_3 + \sin(y+x)^2 xy b_3 \\ & + \cos(y+x)^2 xa_1 + \cos(y+x)^2 xb_1 + 2 \cos(y+x) \sin(y+x) xa_3 \\ & + \cos(y+x) \sin(y+x) xb_3 + \cos(y+x) \sin(y+x) ya_3 + \sin(y+x)^2 xa_1 \\ & + \sin(y+x)^2 xb_1 - \cos(y+x)^2 a_3 + \cos(y+x) \sin(y+x) a_1 = 0 \end{aligned} \quad (6E)$$



Simplifying the above gives

$$\begin{aligned}
& -\frac{a_3}{2} + xy a_3 + xy b_3 - \frac{a_3 \cos(2x+2y)}{2} + \frac{a_1 \sin(2x+2y)}{2} + \frac{3x^2 a_2}{2} \\
& - \frac{x^2 a_2 \cos(2x+2y)}{2} + \frac{3x^2 b_2}{2} - \frac{x^2 b_2 \cos(2x+2y)}{2} - \frac{x^2 a_3}{2} \\
& + \frac{x^2 a_3 \cos(2x+2y)}{2} - \frac{x^2 b_3}{2} + \frac{x^2 b_3 \cos(2x+2y)}{2} + x a_1 + x b_1 \\
& + \frac{y a_3 \sin(2x+2y)}{2} + x a_3 \sin(2x+2y) + \frac{x b_3 \sin(2x+2y)}{2} = 0
\end{aligned} \tag{6E}$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y, \cos(2x+2y), \sin(2x+2y)\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2, \cos(2x+2y) = v_3, \sin(2x+2y) = v_4\}$$

The above PDE (6E) now becomes

$$\begin{aligned}
& -\frac{1}{2} a_3 + v_1 v_2 a_3 + v_1 v_2 b_3 - \frac{1}{2} a_3 v_3 + \frac{1}{2} a_1 v_4 + \frac{3}{2} v_1^2 a_2 \\
& - \frac{1}{2} v_1^2 a_2 v_3 + \frac{3}{2} v_1^2 b_2 - \frac{1}{2} v_1^2 b_2 v_3 - \frac{1}{2} v_1^2 a_3 + \frac{1}{2} v_1^2 a_3 v_3 - \frac{1}{2} v_1^2 b_3 \\
& + \frac{1}{2} v_1^2 b_3 v_3 + v_1 a_1 + v_1 b_1 + \frac{1}{2} v_2 a_3 v_4 + v_1 a_3 v_4 + \frac{1}{2} v_1 b_3 v_4 = 0
\end{aligned} \tag{7E}$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2, v_3, v_4\}$$

Equation (7E) now becomes

$$\begin{aligned}
& -\frac{a_3}{2} + (a_1 + b_1) v_1 - \frac{a_3 v_3}{2} + \frac{a_1 v_4}{2} + \left( \frac{3a_2}{2} + \frac{3b_2}{2} - \frac{a_3}{2} - \frac{b_3}{2} \right) v_1^2 \\
& + (a_3 + b_3) v_2 v_1 + \left( -\frac{a_2}{2} - \frac{b_2}{2} + \frac{a_3}{2} + \frac{b_3}{2} \right) v_3 v_1^2 + \frac{v_2 a_3 v_4}{2} + \left( a_3 + \frac{b_3}{2} \right) v_4 v_1 = 0
\end{aligned} \tag{8E}$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}
 \frac{a_1}{2} &= 0 \\
 -\frac{a_3}{2} &= 0 \\
 \frac{a_3}{2} &= 0 \\
 a_1 + b_1 &= 0 \\
 a_3 + b_3 &= 0 \\
 a_3 + \frac{b_3}{2} &= 0 \\
 -\frac{a_2}{2} - \frac{b_2}{2} + \frac{a_3}{2} + \frac{b_3}{2} &= 0 \\
 \frac{3a_2}{2} + \frac{3b_2}{2} - \frac{a_3}{2} - \frac{b_3}{2} &= 0
 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= 0 \\
 a_2 &= -b_2 \\
 a_3 &= 0 \\
 b_1 &= 0 \\
 b_2 &= b_2 \\
 b_3 &= 0
 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}
 \xi &= -x \\
 \eta &= x
 \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}
 \eta &= \eta - \omega(x, y) \xi \\
 &= x - \left( \frac{-x \sin(y+x) + \cos(y+x)}{x \sin(y+x)} \right) (-x) \\
 &= \frac{\cos(x) \cos(y) - \sin(x) \sin(y)}{\cos(x) \sin(y) + \sin(x) \cos(y)} \\
 \xi &= 0
 \end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{\cos(x)\cos(y) - \sin(x)\sin(y)}{\cos(x)\sin(y) + \sin(x)\cos(y)}} dy \end{aligned}$$

Which results in

$$S = -\ln(\cos(x)\cos(y) - \sin(x)\sin(y))$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{-x \sin(y+x) + \cos(y+x)}{x \sin(y+x)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{\cos(x)\sin(y) + \sin(x)\cos(y)}{\cos(x)\cos(y) - \sin(x)\sin(y)} \\ S_y &= \frac{\cos(x)\sin(y) + \sin(x)\cos(y)}{\cos(x)\cos(y) - \sin(x)\sin(y)} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{x} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = \ln(R) + c_1 \tag{4}$$

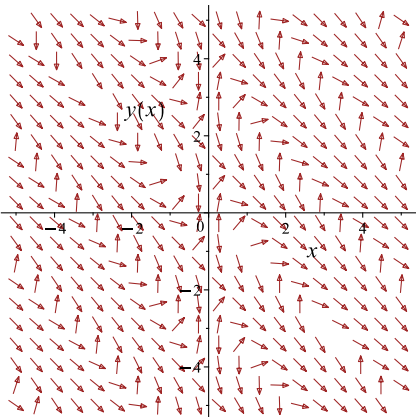
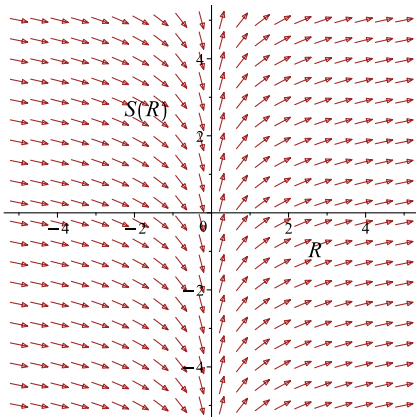
To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$-\ln(\cos(x)\cos(y) - \sin(x)\sin(y)) = \ln(x) + c_1$$

Which simplifies to

$$-\ln(\cos(x)\cos(y) - \sin(x)\sin(y)) = \ln(x) + c_1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{-x \sin(y+x) + \cos(y+x)}{x \sin(y+x)}$ 	$R = x$ $S = -\ln(\cos(x)\cos(y) - \sin(x)\sin(y)) -$	$\frac{dS}{dR} = \frac{1}{R}$ 

### Summary

The solution(s) found are the following

$$-\ln(\cos(x)\cos(y) - \sin(x)\sin(y)) = \ln(x) + c_1 \quad (1)$$

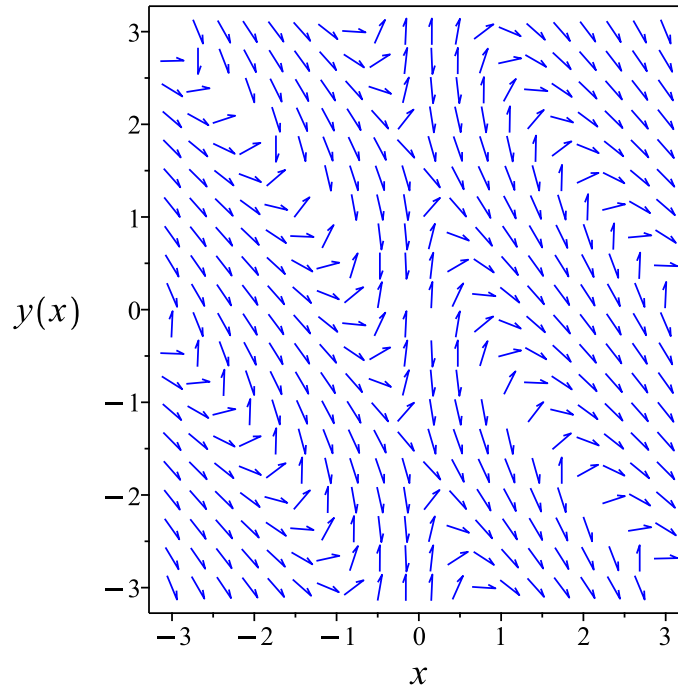


Figure 113: Slope field plot

### Verification of solutions

$$-\ln(\cos(x)\cos(y) - \sin(x)\sin(y)) = \ln(x) + c_1$$

Verified OK.

### **5.12.2 Solving as exact ode**

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (-x \sin(y+x)) dy &= (-\cos(y+x) + x \sin(y+x)) dx \\ (-x \sin(y+x) + \cos(y+x)) dx &+ (-x \sin(y+x)) dy = 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x \sin(y+x) + \cos(y+x) \\ N(x, y) &= -x \sin(y+x) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x \sin(y+x) + \cos(y+x)) \\ &= -x \cos(y+x) - \sin(y+x) \end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x \sin(y+x)) \\ &= -x \cos(y+x) - \sin(y+x)\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x \sin(y+x) + \cos(y+x) dx \\ \phi &= x \cos(y+x) + f(y)\end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = -x \sin(y+x) + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = -x \sin(y+x)$ . Therefore equation (4) becomes

$$-x \sin(y+x) = -x \sin(y+x) + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = x \cos(y+x) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = x \cos(y + x)$$

### Summary

The solution(s) found are the following

$$x \cos(y + x) = c_1 \tag{1}$$

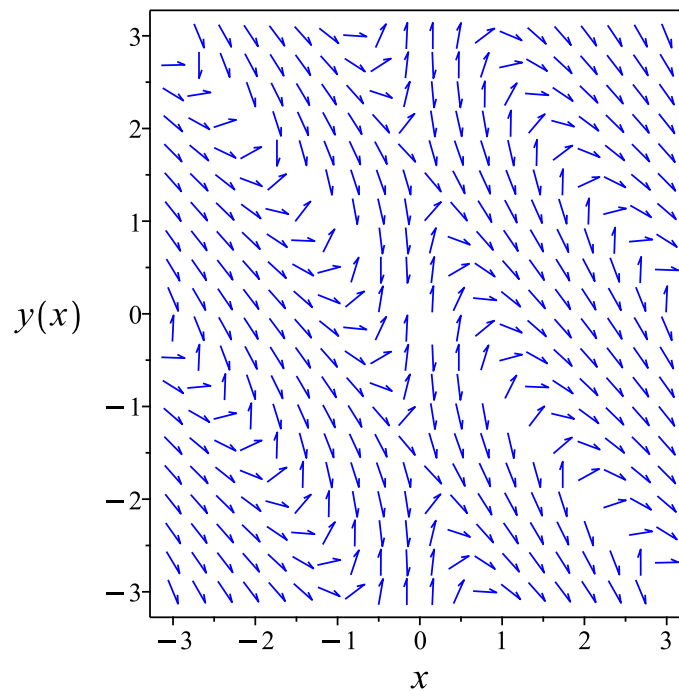


Figure 114: Slope field plot

### Verification of solutions

$$x \cos(y + x) = c_1$$

Verified OK.



### 5.12.3 Maple step by step solution

Let's solve

$$\cos(y+x) - x \sin(y+x) - x \sin(y+x) y' = 0$$

- Highest derivative means the order of the ODE is 1  
 $y'$
- Check if ODE is exact
  - ODE is exact if the lhs is the total derivative of a  $C^2$  function  
 $F'(x, y) = 0$
  - Compute derivative of lhs  
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$
  - Evaluate derivatives  
 $-x \cos(y+x) - \sin(y+x) = -x \cos(y+x) - \sin(y+x)$
  - Condition met, ODE is exact
- Exact ODE implies solution will be of this form  
 $\left[ F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$
- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$   
 $F(x, y) = \int (-x \sin(y+x) + \cos(y+x)) dx + f_1(y)$
- Evaluate integral  
 $F(x, y) = -y \cos(y+x) + (y+x) \cos(y+x) + f_1(y)$
- Take derivative of  $F(x, y)$  with respect to  $y$   
 $N(x, y) = \frac{\partial}{\partial y} F(x, y)$
- Compute derivative  
 $-x \sin(y+x) = y \sin(y+x) - (y+x) \sin(y+x) + \frac{d}{dy} f_1(y)$
- Isolate for  $\frac{d}{dy} f_1(y)$   
 $\frac{d}{dy} f_1(y) = -x \sin(y+x) - y \sin(y+x) + (y+x) \sin(y+x)$
- Solve for  $f_1(y)$   
 $f_1(y) = 0$
- Substitute  $f_1(y)$  into equation for  $F(x, y)$

$$F(x, y) = -y \cos(y + x) + (y + x) \cos(y + x)$$

- Substitute  $F(x, y)$  into the solution of the ODE

$$-y \cos(y + x) + (y + x) \cos(y + x) = c_1$$

- Solve for  $y$

$$y = -x + \arccos\left(\frac{c_1}{x}\right)$$

### Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
differential order: 1; found: 1 linear symmetries. Trying reduction of order
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful`

```

### ✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 14

```
dsolve(cos(x+y(x))-x*sin(x+y(x))=x*sin(x+y(x))*diff(y(x),x),y(x), singsol=all)
```

$$y(x) = -x + \arccos\left(\frac{c_1}{x}\right)$$

### ✓ Solution by Mathematica

Time used: 10.102 (sec). Leaf size: 35

```
DSolve[Cos[x+y[x]]-x*Sin[x+y[x]]==x*Sin[x+y[x]]*y'[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow -x - \arccos\left(-\frac{c_1}{x}\right)$$

$$y(x) \rightarrow -x + \arccos\left(-\frac{c_1}{x}\right)$$

## 5.13 problem 17

5.13.1 Solving as exact ode . . . . .	514
5.13.2 Maple step by step solution . . . . .	517

Internal problem ID [3126]

Internal file name [OUTPUT/2618\_Sunday\_June\_05\_2022\_03\_22\_49\_AM\_71482201/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 17.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"exact"**

Maple gives the following as the ode type

[\_exact]

$$y^2 e^{yx} + (e^{yx} + e^{yx} yx) y' = -\cos(x)$$

### 5.13.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (e^{xy} + y e^{xy} x) dy &= (-y^2 e^{xy} - \cos(x)) dx \\ (y^2 e^{xy} + \cos(x)) dx &+ (e^{xy} + y e^{xy} x) dy = 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= y^2 e^{xy} + \cos(x) \\ N(x, y) &= e^{xy} + y e^{xy} x \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (y^2 e^{xy} + \cos(x)) \\ &= y e^{xy} (xy + 2) \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (e^{xy} + y e^{xy} x) \\ &= y e^{xy} (xy + 2) \end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int y^2 e^{xy} + \cos(x) dx \\ \phi &= y e^{xy} + \sin(x) + f(y)\end{aligned}\quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\begin{aligned}\frac{\partial \phi}{\partial y} &= e^{xy} + y e^{xy} x + f'(y) \\ &= e^{xy}(xy + 1) + f'(y)\end{aligned}\quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = e^{xy} + y e^{xy} x$ . Therefore equation (4) becomes

$$e^{xy} + y e^{xy} x = e^{xy}(xy + 1) + f'(y)\quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = y e^{xy} + \sin(x) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = y e^{xy} + \sin(x)$$

The solution becomes

$$y = \frac{\text{LambertW}(x(c_1 - \sin(x)))}{x}$$

### Summary

The solution(s) found are the following

$$y = \frac{\text{LambertW}(x(c_1 - \sin(x)))}{x} \quad (1)$$

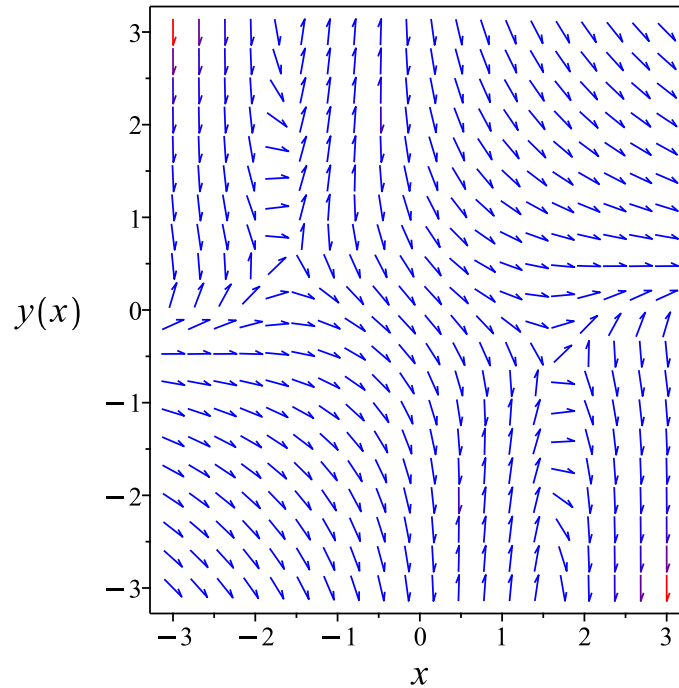


Figure 115: Slope field plot

### Verification of solutions

$$y = \frac{\text{LambertW}(x(c_1 - \sin(x)))}{x}$$

Verified OK.

### 5.13.2 Maple step by step solution

Let's solve

$$y^2 e^{yx} + (e^{yx} + e^{yx} yx) y' = -\cos(x)$$

- Highest derivative means the order of the ODE is 1  
 $y'$
- Check if ODE is exact

- ODE is exact if the lhs is the total derivative of a  $C^2$  function  

$$F'(x, y) = 0$$
- Compute derivative of lhs  

$$F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$$
- Evaluate derivatives  

$$2y e^{xy} + e^{xy} x y^2 = 2y e^{xy} + e^{xy} x y^2$$
- Condition met, ODE is exact
- Exact ODE implies solution will be of this form  

$$\left[ F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$
- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$   

$$F(x, y) = \int (y^2 e^{xy} + \cos(x)) dx + f_1(y)$$
- Evaluate integral  

$$F(x, y) = y e^{xy} + \sin(x) + f_1(y)$$
- Take derivative of  $F(x, y)$  with respect to  $y$   

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$
- Compute derivative  

$$e^{xy} + y e^{xy} x = e^{xy} + y e^{xy} x + \frac{d}{dy} f_1(y)$$
- Isolate for  $\frac{d}{dy} f_1(y)$   

$$\frac{d}{dy} f_1(y) = 0$$
- Solve for  $f_1(y)$   

$$f_1(y) = 0$$
- Substitute  $f_1(y)$  into equation for  $F(x, y)$   

$$F(x, y) = y e^{xy} + \sin(x)$$
- Substitute  $F(x, y)$  into the solution of the ODE  

$$y e^{xy} + \sin(x) = c_1$$
- Solve for  $y$   

$$y = \frac{\text{LambertW}(x(c_1 - \sin(x)))}{x}$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

### ✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 16

```
dsolve((y(x)^2*exp(x*y(x))+cos(x))+(exp(x*y(x))+x*y(x)*exp(x*y(x)))*diff(y(x),x)=0,y(x), sin
```

$$y(x) = \frac{\text{LambertW}(-x(\sin(x) + c_1))}{x}$$

### ✓ Solution by Mathematica

Time used: 60.266 (sec). Leaf size: 19

```
DSolve[(y[x]^2*Exp[x*y[x]]+Cos[x])+(Exp[x*y[x]]+x*y[x]*Exp[x*y[x]])*y'[x]==0,y[x],x,IncludeS
```

$$y(x) \rightarrow \frac{W(x(-\sin(x) + c_1))}{x}$$



## 5.14 problem 18

5.14.1 Solving as first order ode lie symmetry calculated ode . . . . .	520
5.14.2 Solving as exact ode . . . . .	526
5.14.3 Maple step by step solution . . . . .	530

Internal problem ID [3127]

Internal file name [OUTPUT/2619\_Sunday\_June\_05\_2022\_03\_22\_53\_AM\_32827643/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 18.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**exact**", "**first\_order\_ode\_lie\_symmetry\_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class C`], _exact, _dAlembert]
```

$$y' \ln(-y + x) - \ln(-y + x) = 1$$

### 5.14.1 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = \frac{1 + \ln(-y + x)}{\ln(-y + x)}$$
$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \tag{A}$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \tag{1E}$$

$$\eta = xb_2 + yb_3 + b_1 \tag{2E}$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} b_2 + \frac{(1 + \ln(-y + x))(b_3 - a_2)}{\ln(-y + x)} - \frac{(1 + \ln(-y + x))^2 a_3}{\ln(-y + x)^2} \\ - \left( \frac{1}{(-y + x) \ln(-y + x)} - \frac{1 + \ln(-y + x)}{\ln(-y + x)^2 (-y + x)} \right) (xa_2 + ya_3 + a_1) \\ - \left( -\frac{1}{(-y + x) \ln(-y + x)} + \frac{1 + \ln(-y + x)}{\ln(-y + x)^2 (-y + x)} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\begin{aligned} \frac{\ln(-y + x)^2 xa_2 + \ln(-y + x)^2 xa_3 - b_2 \ln(-y + x)^2 x - \ln(-y + x)^2 xb_3 - \ln(-y + x)^2 ya_2 - \ln(-y + x)^2 yb_3 + a_1 - b_1}{\ln(-y + x)^2} \\ = 0 \end{aligned}$$

Setting the numerator to zero gives

$$\begin{aligned} -\ln(-y + x)^2 xa_2 - \ln(-y + x)^2 xa_3 + b_2 \ln(-y + x)^2 x \\ + \ln(-y + x)^2 xb_3 + \ln(-y + x)^2 ya_2 + \ln(-y + x)^2 yb_3 \\ - b_2 \ln(-y + x)^2 y - \ln(-y + x)^2 yb_3 - \ln(-y + x) xa_2 \\ - 2 \ln(-y + x) xa_3 + \ln(-y + x) xb_3 + \ln(-y + x) ya_2 + 2 \ln(-y + x) yb_3 \\ - \ln(-y + x) yb_3 + xa_2 - xa_3 - xb_2 + 2ya_3 - yb_3 + a_1 - b_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y, \ln(-y + x)\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2, \ln(-y + x) = v_3\}$$

The above PDE (6E) now becomes

$$\begin{aligned} -v_3^2 v_1 a_2 + v_3^2 v_2 a_2 - v_3^2 v_1 a_3 + v_3^2 v_2 a_3 + b_2 v_3^2 v_1 - b_2 v_3^2 v_2 + v_3^2 v_1 b_3 \\ - v_3^2 v_2 b_3 - v_3 v_1 a_2 + v_3 v_2 a_2 - 2v_3 v_1 a_3 + 2v_3 v_2 a_3 + v_3 v_1 b_3 \\ - v_3 v_2 b_3 + v_1 a_2 - v_1 a_3 + 2v_2 a_3 - v_1 b_2 - v_2 b_3 + a_1 - b_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2, v_3\}$$

Equation (7E) now becomes

$$\begin{aligned} &(-a_2 - a_3 + b_2 + b_3) v_1 v_3^2 + (-a_2 - 2a_3 + b_3) v_1 v_3 + (a_2 - a_3 - b_2) v_1 \\ &+ (a_2 + a_3 - b_2 - b_3) v_2 v_3^2 + (a_2 + 2a_3 - b_3) v_2 v_3 + (2a_3 - b_3) v_2 + a_1 - b_1 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} a_1 - b_1 &= 0 \\ 2a_3 - b_3 &= 0 \\ -a_2 - 2a_3 + b_3 &= 0 \\ a_2 - a_3 - b_2 &= 0 \\ a_2 + 2a_3 - b_3 &= 0 \\ -a_2 - a_3 + b_2 + b_3 &= 0 \\ a_2 + a_3 - b_2 - b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= b_1 \\ a_2 &= 0 \\ a_3 &= -b_2 \\ b_1 &= b_1 \\ b_2 &= b_2 \\ b_3 &= -2b_2 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= 1 \\ \eta &= 1 \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= 1 - \left( \frac{1 + \ln(-y + x)}{\ln(-y + x)} \right) (1) \\ &= -\frac{1}{\ln(-y + x)} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{-\ln(-y+x)} dy \end{aligned}$$

Which results in

$$S = (-y + x) \ln(-y + x) + y - x$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = \frac{1 + \ln(-y + x)}{\ln(-y + x)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \ln(-y + x) \\ S_y &= -\ln(-y + x) \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -1 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -R + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$(\ln(-y + x) - 1)(-y + x) = -x + c_1$$

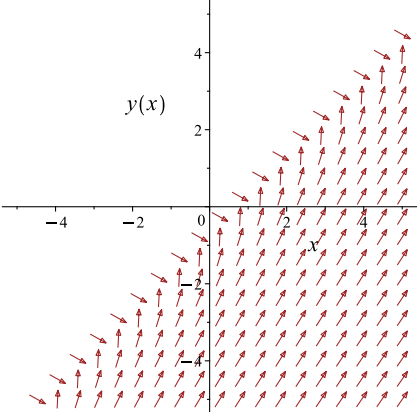
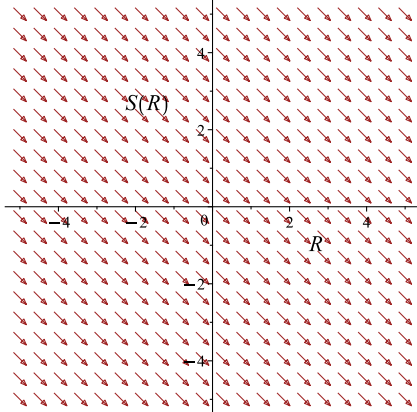
Which simplifies to

$$(-y + x) \ln(-y + x) + y - c_1 = 0$$

Which gives

$$y = -e^{\text{LambertW}((-x+c_1)e^{-1})+1} + x$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = \frac{1 + \ln(-y+x)}{\ln(-y+x)}$ 	$R = x$ $S = (\ln(-y+x) - 1) (-1)$	$\frac{dS}{dR} = -1$ 

Summary

The solution(s) found are the following

$$y = -e^{\text{LambertW}((-x+c_1)e^{-1})+1} + x \tag{1}$$

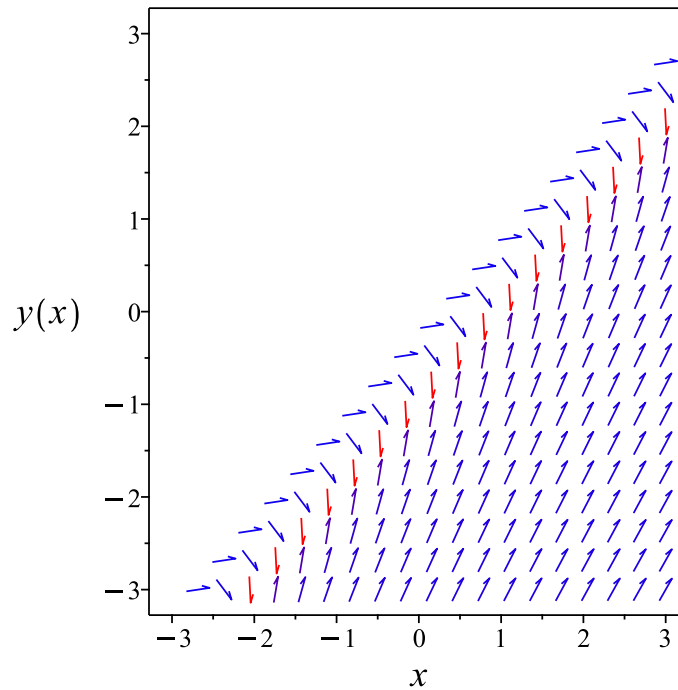


Figure 116: Slope field plot

### Verification of solutions

$$y = -e^{\text{LambertW}((-x+c_1)e^{-1})+1} + x$$

Verified OK.

### 5.14.2 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(\ln(-y+x)) dy &= (1 + \ln(-y+x)) dx \\ (-1 - \ln(-y+x)) dx &+ (\ln(-y+x)) dy = 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -1 - \ln(-y+x) \\ N(x, y) &= \ln(-y+x)\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-1 - \ln(-y+x)) \\ &= \frac{1}{-y+x}\end{aligned}$$



And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(\ln(-y+x)) \\ &= \frac{1}{-y+x}\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -1 - \ln(-y+x) dx \\ \phi &= (y-x)\ln(-y+x) - y + f(y)\end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\begin{aligned}\frac{\partial \phi}{\partial y} &= \ln(-y+x) - \frac{y-x}{-y+x} - 1 + f'(y) \\ &= \ln(-y+x) + f'(y)\end{aligned} \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \ln(-y+x)$ . Therefore equation (4) becomes

$$\ln(-y+x) = \ln(-y+x) + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = (y - x) \ln(-y + x) - y + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = (y - x) \ln(-y + x) - y$$

The solution becomes

$$y = -e^{\text{LambertW}(-(x+c_1)e^{-1})+1} + x$$

### Summary

The solution(s) found are the following

$$y = -e^{\text{LambertW}(-(x+c_1)e^{-1})+1} + x \tag{1}$$

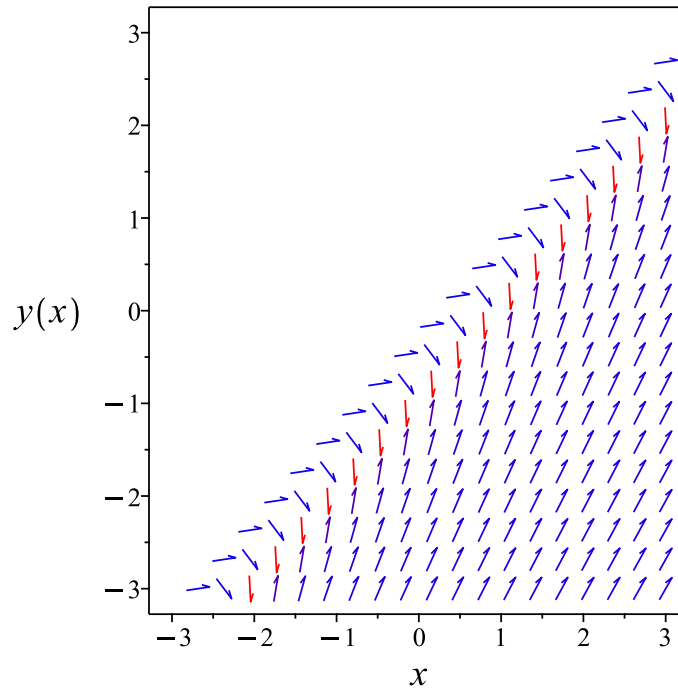


Figure 117: Slope field plot

### Verification of solutions

$$y = -e^{\text{LambertW}(-(x+c_1)e^{-1})+1} + x$$

Verified OK.

### 5.14.3 Maple step by step solution

Let's solve

$$y' \ln(-y + x) - \ln(-y + x) = 1$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Check if ODE is exact

- ODE is exact if the lhs is the total derivative of a  $C^2$  function

$$F'(x, y) = 0$$

- Compute derivative of lhs

$$F'(x, y) + \left( \frac{\partial}{\partial y} F(x, y) \right) y' = 0$$

- Evaluate derivatives

$$\frac{1}{-y+x} = \frac{1}{-y+x}$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$\left[ F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$

- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$

$$F(x, y) = \int (-1 - \ln(-y + x)) dx + f_1(y)$$

- Evaluate integral

$$F(x, y) = -(-y + x) \ln(-y + x) - y + f_1(y)$$

- Take derivative of  $F(x, y)$  with respect to  $y$

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$

- Compute derivative

$$\ln(-y + x) = \ln(-y + x) + \frac{d}{dy} f_1(y)$$

- Isolate for  $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = 0$$

- Solve for  $f_1(y)$

$$f_1(y) = 0$$

- Substitute  $f_1(y)$  into equation for  $F(x, y)$

$$F(x, y) = -(-y + x) \ln(-y + x) - y$$

- Substitute  $F(x, y)$  into the solution of the ODE

$$-(-y + x) \ln(-y + x) - y = c_1$$

- Solve for  $y$

$$y = -e^{\text{LambertW}\left(-\frac{x+c_1}{e}\right)+1} + x$$

### Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous C
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful
<- homogeneous successful`

```

### ✓ Solution by Maple

Time used: 0.031 (sec). Leaf size: 32

```
dsolve(diff(y(x),x)*ln(x-y(x))=1+ln(x-y(x)),y(x), singsol=all)
```

$$y(x) = \frac{x \text{LambertW}((c_1 - x) e^{-1}) - c_1 + x}{\text{LambertW}((c_1 - x) e^{-1})}$$

### ✓ Solution by Mathematica

Time used: 0.127 (sec). Leaf size: 26

```
DSolve[y'[x]*Log[x-y[x]]==1+Log[x-y[x]],y[x],x,IncludeSingularSolutions -> True]
```

$$\text{Solve}[(x - y(x))(-\log(x - y(x))) - y(x) = c_1, y(x)]$$

## 5.15 problem 19

5.15.1 Solving as linear ode . . . . .	532
5.15.2 Solving as first order ode lie symmetry lookup ode . . . . .	534
5.15.3 Solving as exact ode . . . . .	538
5.15.4 Maple step by step solution . . . . .	543

Internal problem ID [3128]

Internal file name [OUTPUT/2620\_Sunday\_June\_05\_2022\_03\_22\_55\_AM\_39789059/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 19.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"linear"**, **"exactWithIntegrationFactor"**, **"first\_order\_ode\_lie\_symmetry\_lookup"**

Maple gives the following as the ode type

`[_linear]`

$$y' + 2yx = e^{-x^2}$$

### 5.15.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = 2x$$

$$q(x) = e^{-x^2}$$

Hence the ode is

$$y' + 2yx = e^{-x^2}$$

The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int 2x dx} \\ &= e^{x^2}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) (e^{-x^2}) \\ \frac{d}{dx}(e^{x^2} y) &= (e^{x^2}) (e^{-x^2}) \\ d(e^{x^2} y) &= dx\end{aligned}$$

Integrating gives

$$\begin{aligned}e^{x^2} y &= \int dx \\ e^{x^2} y &= x + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = e^{x^2}$  results in

$$y = x e^{-x^2} + c_1 e^{-x^2}$$

which simplifies to

$$y = e^{-x^2}(x + c_1)$$

### Summary

The solution(s) found are the following

$$y = e^{-x^2}(x + c_1) \tag{1}$$

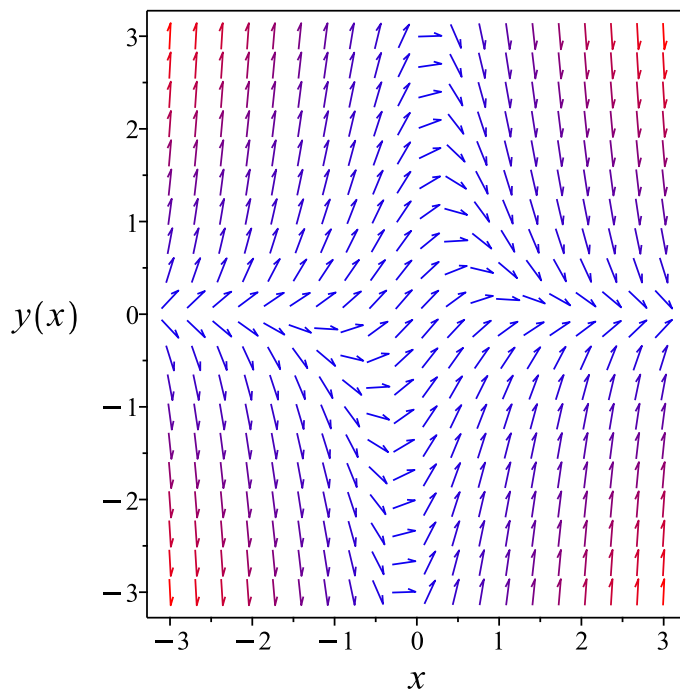


Figure 118: Slope field plot

### Verification of solutions

$$y = e^{-x^2}(x + c_1)$$

Verified OK.

### 5.15.2 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$\begin{aligned}y' &= -2xy + e^{-x^2} \\y' &= \omega(x, y)\end{aligned}$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 62: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= e^{-x^2}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the



canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{e^{-x^2}} dy \end{aligned}$$

Which results in

$$S = e^{x^2} y$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -2xy + e^{-x^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= 2x e^{x^2} y \\ S_y &= e^{x^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 1 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by

integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = R + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$y e^{x^2} = x + c_1$$

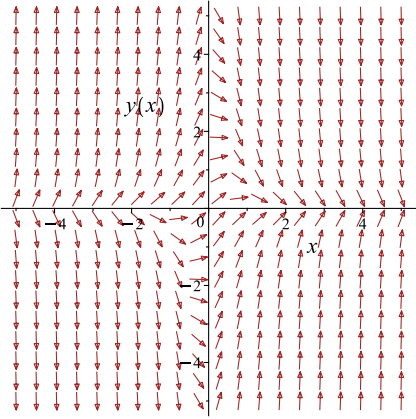
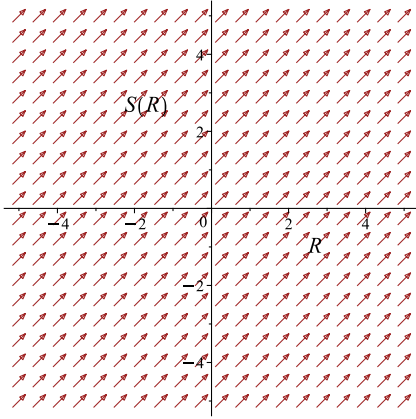
Which simplifies to

$$y e^{x^2} = x + c_1$$

Which gives

$$y = e^{-x^2}(x + c_1)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -2xy + e^{-x^2}$ 	$R = x$ $S = e^{x^2} y$	$\frac{dS}{dR} = 1$ 

### Summary

The solution(s) found are the following

$$y = e^{-x^2}(x + c_1) \quad (1)$$

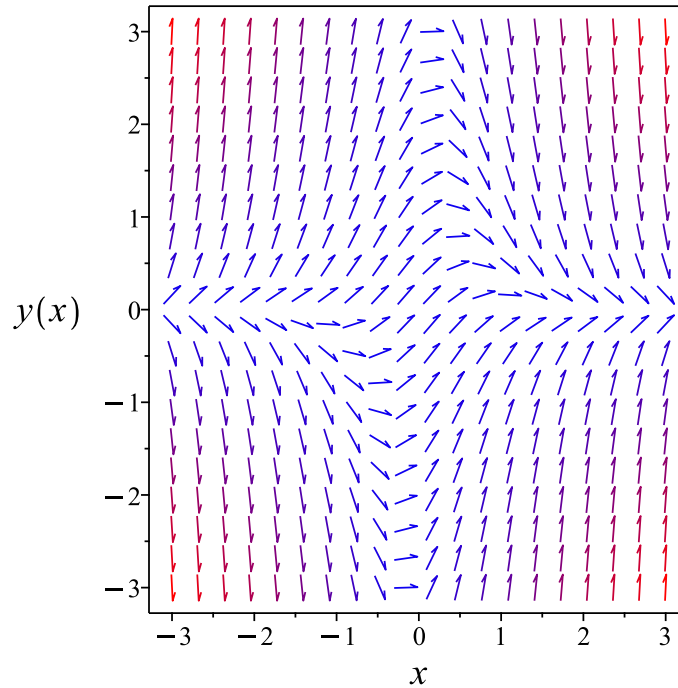


Figure 119: Slope field plot

Verification of solutions

$$y = e^{-x^2}(x + c_1)$$

Verified OK.

### 5.15.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dy &= (-2xy + e^{-x^2}) dx \\ (2xy - e^{-x^2}) dx + dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 2xy - e^{-x^2} \\ N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2xy - e^{-x^2}) \\ &= 2x\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((2x) - (0)) \\ &= 2x \end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned} \mu &= e^{\int A \, dx} \\ &= e^{\int 2x \, dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{x^2} \\ &= e^{x^2} \end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{x^2} (2xy - e^{-x^2}) \\ &= 2x e^{x^2} y - 1 \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{x^2} (1) \\ &= e^{x^2} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (2x e^{x^2} y - 1) + (e^{x^2}) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \overline{M} dx$$
$$\int \frac{\partial \phi}{\partial x} dx = \int 2x e^{x^2} y - 1 dx$$
$$\phi = -x + e^{x^2} y + f(y) \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = e^{x^2} + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = e^{x^2}$ . Therefore equation (4) becomes

$$e^{x^2} = e^{x^2} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -x + e^{x^2} y + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -x + e^{x^2} y$$

The solution becomes

$$y = e^{-x^2}(x + c_1)$$

### Summary

The solution(s) found are the following

$$y = e^{-x^2}(x + c_1) \tag{1}$$

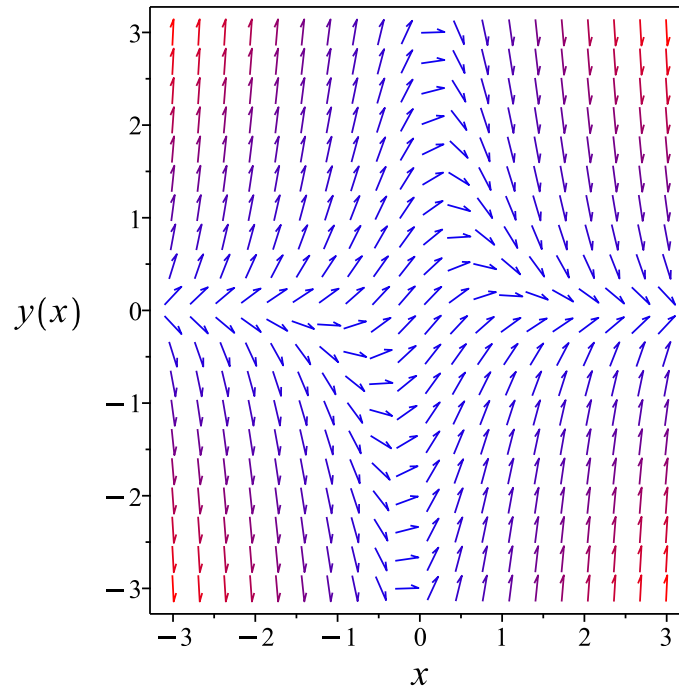


Figure 120: Slope field plot

### Verification of solutions

$$y = e^{-x^2}(x + c_1)$$

Verified OK.

### 5.15.4 Maple step by step solution

Let's solve

$$y' + 2yx = e^{-x^2}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -2yx + e^{-x^2}$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + 2yx = e^{-x^2}$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x)(y' + 2yx) = \mu(x)e^{-x^2}$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$

$$\mu(x)(y' + 2yx) = \mu'(x)y + \mu(x)y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = 2\mu(x)x$$

- Solve to find the integrating factor

$$\mu(x) = e^{x^2}$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x)y) \right) dx = \int \mu(x)e^{-x^2} dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int \mu(x)e^{-x^2} dx + c_1$$

- Solve for  $y$

$$y = \frac{\int \mu(x)e^{-x^2} dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = e^{x^2}$

$$y = \frac{\int e^{-x^2} e^{x^2} dx + c_1}{e^{x^2}}$$

- Evaluate the integrals on the rhs

$$y = \frac{x + c_1}{e^{x^2}}$$

- Simplify



$$y = e^{-x^2}(x + c_1)$$

### Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

#### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 14

```
dsolve(diff(y(x),x)+2*x*y(x)=exp(-x^2),y(x), singsol=all)
```

$$y(x) = (c_1 + x)e^{-x^2}$$

#### ✓ Solution by Mathematica

Time used: 0.055 (sec). Leaf size: 17

```
DSolve[y'[x]+2*x*y[x]==Exp[-x^2],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow e^{-x^2}(x + c_1)$$

## 5.16 problem 20

5.16.1 Solving as homogeneousTypeD2 ode . . . . .	545
5.16.2 Solving as first order ode lie symmetry calculated ode . . . . .	547
5.16.3 Solving as exact ode . . . . .	553

Internal problem ID [3129]

Internal file name [OUTPUT/2621\_Sunday\_June\_05\_2022\_03\_22\_57\_AM\_39946048/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 20.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**homogeneousTypeD2**", "**exactWithIntegrationFactor**", "**first\_order\_ode\_lie\_symmetry\_calculated**"

Maple gives the following as the ode type

```
[[_homogeneous, `class A`], _rational, [_Abel, `2nd type`, `class B`]]
```

$$y^2 - 3yx - (x^2 - yx) y' = 2x^2$$

### 5.16.1 Solving as homogeneousTypeD2 ode

Using the change of variables  $y = u(x)x$  on the above ode results in new ode in  $u(x)$

$$u(x)^2 x^2 - 3u(x)x^2 - (x^2 - u(x)x^2)(u'(x)x + u(x)) = 2x^2$$

In canonical form the ODE is

$$\begin{aligned} u' &= F(x, u) \\ &= f(x)g(u) \\ &= -\frac{2(u^2 - 2u - 1)}{x(u - 1)} \end{aligned}$$

Where  $f(x) = -\frac{2}{x}$  and  $g(u) = \frac{u^2-2u-1}{u-1}$ . Integrating both sides gives

$$\begin{aligned}\frac{1}{\frac{u^2-2u-1}{u-1}} du &= -\frac{2}{x} dx \\ \int \frac{1}{\frac{u^2-2u-1}{u-1}} du &= \int -\frac{2}{x} dx \\ \frac{\ln(u^2 - 2u - 1)}{2} &= -2 \ln(x) + c_2\end{aligned}$$

Raising both side to exponential gives

$$\sqrt{u^2 - 2u - 1} = e^{-2\ln(x)+c_2}$$

Which simplifies to

$$\sqrt{u^2 - 2u - 1} = \frac{c_3}{x^2}$$

Which simplifies to

$$\sqrt{u(x)^2 - 2u(x) - 1} = \frac{c_3 e^{c_2}}{x^2}$$

The solution is

$$\sqrt{u(x)^2 - 2u(x) - 1} = \frac{c_3 e^{c_2}}{x^2}$$

Replacing  $u(x)$  in the above solution by  $\frac{y}{x}$  results in the solution for  $y$  in implicit form

$$\begin{aligned}\sqrt{\frac{y^2}{x^2} - \frac{2y}{x} - 1} &= \frac{c_3 e^{c_2}}{x^2} \\ \sqrt{\frac{y^2 - 2yx - x^2}{x^2}} &= \frac{c_3 e^{c_2}}{x^2}\end{aligned}$$

### Summary

The solution(s) found are the following

$$\sqrt{\frac{y^2 - 2yx - x^2}{x^2}} = \frac{c_3 e^{c_2}}{x^2} \quad (1)$$

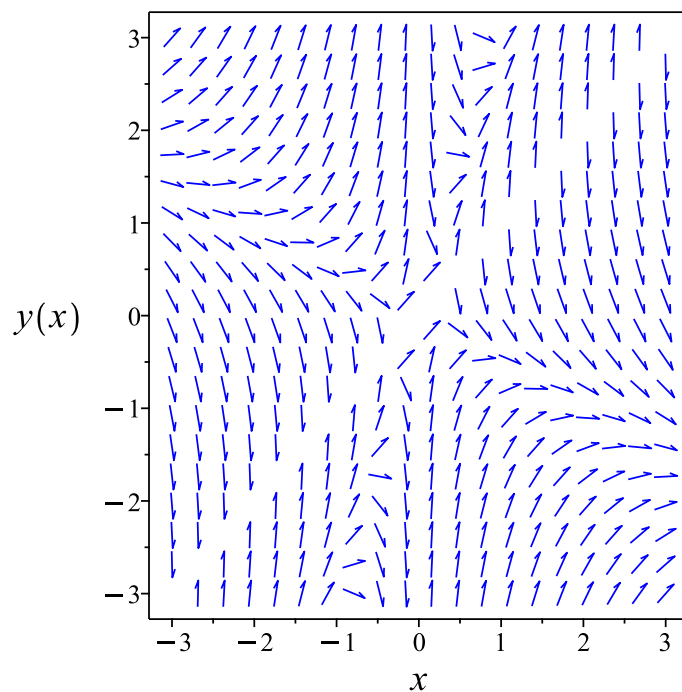


Figure 121: Slope field plot

### Verification of solutions

$$\sqrt{\frac{y^2 - 2yx - x^2}{x^2}} = \frac{c_3 e^{c_2}}{x^2}$$

Verified OK.

### 5.16.2 Solving as first order ode lie symmetry calculated ode

Writing the ode as

$$y' = -\frac{-2x^2 - 3xy + y^2}{x(y - x)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is not in the lookup table. To determine  $\xi, \eta$  then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and  $\omega$  into (A) gives

$$\begin{aligned} b_2 - \frac{(-2x^2 - 3xy + y^2)(b_3 - a_2)}{x(y-x)} - \frac{(-2x^2 - 3xy + y^2)^2 a_3}{x^2(y-x)^2} \\ - \left( -\frac{4x - 3y}{x(y-x)} + \frac{-2x^2 - 3xy + y^2}{x^2(y-x)} - \frac{-2x^2 - 3xy + y^2}{x(y-x)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left( -\frac{3x + 2y}{x(y-x)} + \frac{-2x^2 - 3xy + y^2}{x(y-x)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{2x^4a_2 - 4x^4a_3 + 6x^4b_2 - 2x^4b_3 - 4x^3ya_2 - 12x^3ya_3 - 4x^3yb_2 + 4x^3yb_3 - 2x^2y^2a_2 - 10x^2y^2a_3 + 2x^2y^2b_2 - 2x^2y^2b_3 + 8xy^3a_3 - 2y^4a_3 + 5x^3b_1 - 5x^2ya_1 - 2x^2yb_1 + 2xy^2a_1 + xy^2b_1 - y^3a_1}{x^2(-y+x)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} 2x^4a_2 - 4x^4a_3 + 6x^4b_2 - 2x^4b_3 - 4x^3ya_2 - 12x^3ya_3 - 4x^3yb_2 \\ + 4x^3yb_3 - 2x^2y^2a_2 - 10x^2y^2a_3 + 2x^2y^2b_2 + 2x^2y^2b_3 + 8xy^3a_3 \\ - 2y^4a_3 + 5x^3b_1 - 5x^2ya_1 - 2x^2yb_1 + 2xy^2a_1 + xy^2b_1 - y^3a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with  $\{x, y\}$  in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with  $\{x, y\}$  in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} 2a_2v_1^4 - 4a_2v_1^3v_2 - 2a_2v_1^2v_2^2 - 4a_3v_1^4 - 12a_3v_1^3v_2 - 10a_3v_1^2v_2^2 + 8a_3v_1v_2^3 \\ - 2a_3v_1^4 + 6b_2v_1^4 - 4b_2v_1^3v_2 + 2b_2v_1^2v_2^2 - 2b_3v_1^4 + 4b_3v_1^3v_2 + 2b_3v_1^2v_2^2 \\ - 5a_1v_1^2v_2 + 2a_1v_1v_2^2 - a_1v_2^3 + 5b_1v_1^3 - 2b_1v_1^2v_2 + b_1v_1v_2^2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms  $v_i$  introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & (2a_2 - 4a_3 + 6b_2 - 2b_3)v_1^4 + (-4a_2 - 12a_3 - 4b_2 + 4b_3)v_1^3v_2 + 5b_1v_1^3 \\ & + (-2a_2 - 10a_3 + 2b_2 + 2b_3)v_1^2v_2^2 + (-5a_1 - 2b_1)v_1^2v_2 \\ & + 8a_3v_1v_2^3 + (2a_1 + b_1)v_1v_2^2 - 2a_3v_2^4 - a_1v_2^3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -a_1 &= 0 \\ -2a_3 &= 0 \\ 8a_3 &= 0 \\ 5b_1 &= 0 \\ -5a_1 - 2b_1 &= 0 \\ 2a_1 + b_1 &= 0 \\ -4a_2 - 12a_3 - 4b_2 + 4b_3 &= 0 \\ -2a_2 - 10a_3 + 2b_2 + 2b_3 &= 0 \\ 2a_2 - 4a_3 + 6b_2 - 2b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make  $\xi = 0$  in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left( -\frac{-2x^2 - 3xy + y^2}{x(y-x)} \right) (x) \\ &= \frac{2x^2 + 4xy - 2y^2}{-y+x} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that  $\left( \xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{2x^2 + 4xy - 2y^2}{-y+x}} dy\end{aligned}$$

Which results in

$$S = \frac{\ln(-x^2 - 2xy + y^2)}{4}$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{-2x^2 - 3xy + y^2}{x(y-x)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{y+x}{2x^2+4xy-2y^2} \\ S_y &= \frac{-y+x}{2x^2+4xy-2y^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{2x} \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{2R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = -\frac{\ln(R)}{2} + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

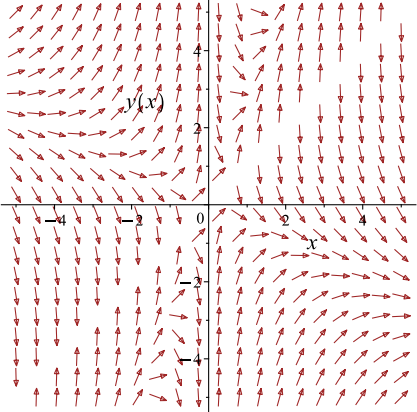
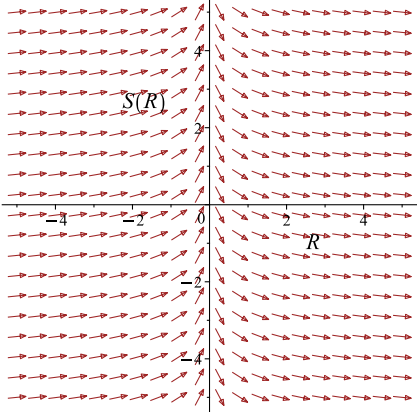
$$\frac{\ln(y^2 - 2yx - x^2)}{4} = -\frac{\ln(x)}{2} + c_1$$

Which simplifies to

$$\frac{\ln(y^2 - 2yx - x^2)}{4} = -\frac{\ln(x)}{2} + c_1$$



The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{-2x^2 - 3xy + y^2}{x(y-x)}$ 	$R = x$ $S = \frac{\ln(-x^2 - 2xy + y^2)}{4}$	$\frac{dS}{dR} = -\frac{1}{2R}$ 

Summary

The solution(s) found are the following

$$\frac{\ln(y^2 - 2yx - x^2)}{4} = -\frac{\ln(x)}{2} + c_1 \tag{1}$$

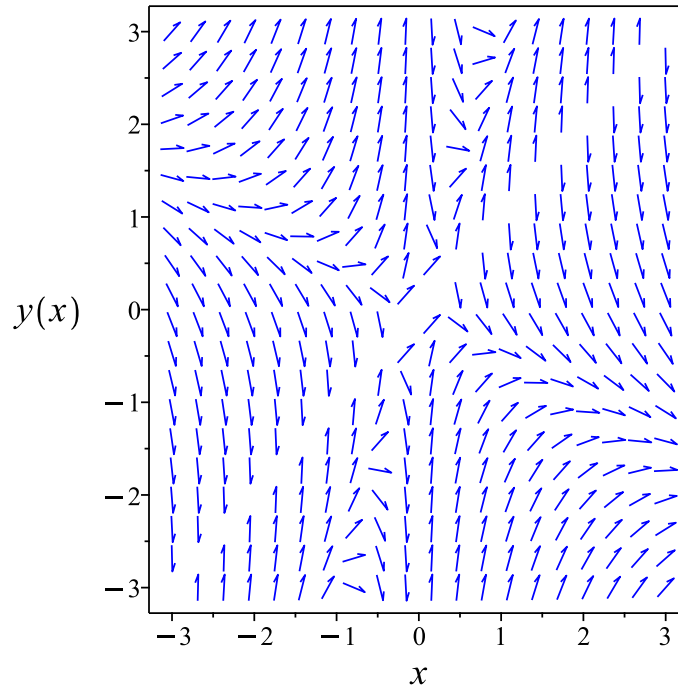


Figure 122: Slope field plot

Verification of solutions

$$\frac{\ln(y^2 - 2yx - x^2)}{4} = -\frac{\ln(x)}{2} + c_1$$

Verified OK.

### 5.16.3 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(-x^2 + xy) dy &= (2x^2 + 3xy - y^2) dx \\ (-2x^2 - 3xy + y^2) dx + (-x^2 + xy) dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -2x^2 - 3xy + y^2 \\ N(x, y) &= -x^2 + xy\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-2x^2 - 3xy + y^2) \\ &= -3x + 2y\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x^2 + xy) \\ &= -2x + y\end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= -\frac{1}{x(-y+x)} ((-3x+2y) - (-2x+y)) \\ &= \frac{1}{x} \end{aligned}$$

Since  $A$  does not depend on  $y$ , then it can be used to find an integrating factor. The integrating factor  $\mu$  is

$$\begin{aligned} \mu &= e^{\int A dx} \\ &= e^{\int \frac{1}{x} dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\ln(x)} \\ &= x \end{aligned}$$

$M$  and  $N$  are multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\bar{M}$  and  $\bar{N}$  for now so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned} \bar{M} &= \mu M \\ &= x(-2x^2 - 3xy + y^2) \\ &= -2x^3 - 3x^2y + xy^2 \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= x(-x^2 + xy) \\ &= -x^2(-y + x) \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (-2x^3 - 3x^2y + xy^2) + (-x^2(-y + x)) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -2x^3 - 3x^2y + xy^2 dx \\ \phi &= -\frac{1}{2}x^4 - x^3y + \frac{1}{2}y^2x^2 + f(y) \end{aligned} \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\begin{aligned} \frac{\partial \phi}{\partial y} &= -x^3 + x^2y + f'(y) \\ &= -x^2(-y + x) + f'(y) \end{aligned} \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = -x^2(-y + x)$ . Therefore equation (4) becomes

$$-x^2(-y + x) = -x^2(-y + x) + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\frac{1}{2}x^4 - x^3y + \frac{1}{2}y^2x^2 + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\frac{1}{2}x^4 - x^3y + \frac{1}{2}y^2x^2$$

### Summary

The solution(s) found are the following

$$-\frac{x^4}{2} - yx^3 + \frac{y^2x^2}{2} = c_1 \quad (1)$$

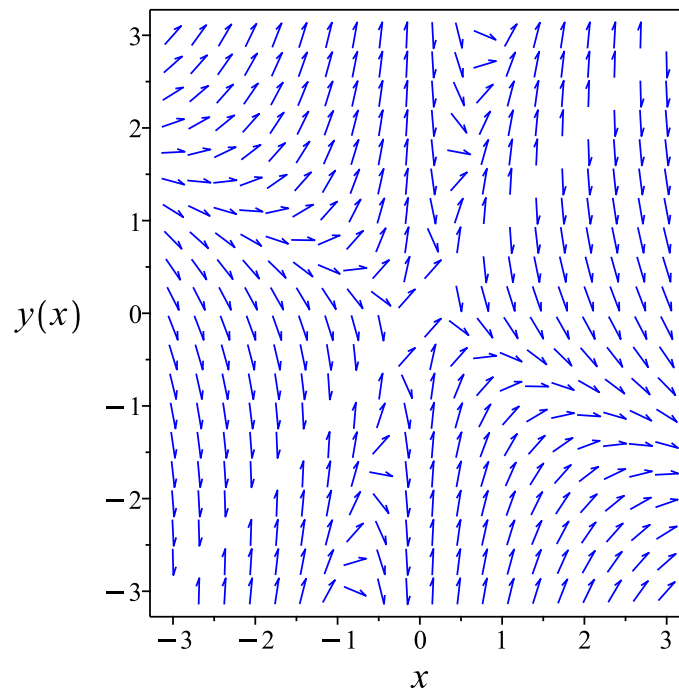


Figure 123: Slope field plot

### Verification of solutions

$$-\frac{x^4}{2} - yx^3 + \frac{y^2x^2}{2} = c_1$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying homogeneous D  
<- homogeneous successful`
```

### ✓ Solution by Maple

Time used: 0.047 (sec). Leaf size: 59

```
dsolve((y(x)^2-3*x*y(x)-2*x^2)=(x^2-x*y(x))*diff(y(x),x),y(x), singsol=all)
```

$$y(x) = \frac{c_1 x^2 - \sqrt{2c_1^2 x^4 + 1}}{c_1 x}$$
$$y(x) = \frac{c_1 x^2 + \sqrt{2c_1^2 x^4 + 1}}{c_1 x}$$

### ✓ Solution by Mathematica

Time used: 0.701 (sec). Leaf size: 99

```
DSolve[(y[x]^2-3*x*y[x]-2*x^2)==(x^2-x*y[x])*y'[x],y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow x - \frac{\sqrt{2x^4 + e^{2c_1}}}{x}$$
$$y(x) \rightarrow x + \frac{\sqrt{2x^4 + e^{2c_1}}}{x}$$
$$y(x) \rightarrow x - \frac{\sqrt{2}\sqrt{x^4}}{x}$$
$$y(x) \rightarrow \frac{\sqrt{2}\sqrt{x^4}}{x} + x$$

## 5.17 problem 21

5.17.1 Solving as linear ode . . . . .	559
5.17.2 Solving as differentialType ode . . . . .	561
5.17.3 Solving as first order ode lie symmetry lookup ode . . . . .	563
5.17.4 Solving as exact ode . . . . .	567
5.17.5 Maple step by step solution . . . . .	571

Internal problem ID [3130]

Internal file name [OUTPUT/2622\_Sunday\_June\_05\_2022\_03\_23\_00\_AM\_414119/index.tex]

**Book:** Differential equations with applications and historial notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 21.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "exact", "linear", "differentialType", "first\_order\_ode\_lie\_symmetry\_lookup"

Maple gives the following as the ode type

[\_linear]

$$2yx + (x^2 + 1)y' = 4x^3$$

### 5.17.1 Solving as linear ode

Entering Linear first order ODE solver. In canonical form a linear first order is

$$y' + p(x)y = q(x)$$

Where here

$$p(x) = \frac{2x}{x^2 + 1}$$
$$q(x) = \frac{4x^3}{x^2 + 1}$$

Hence the ode is

$$y' + \frac{2xy}{x^2 + 1} = \frac{4x^3}{x^2 + 1}$$



The integrating factor  $\mu$  is

$$\begin{aligned}\mu &= e^{\int \frac{2x}{x^2+1} dx} \\ &= x^2 + 1\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= (\mu) \left( \frac{4x^3}{x^2 + 1} \right) \\ \frac{d}{dx}((x^2 + 1) y) &= (x^2 + 1) \left( \frac{4x^3}{x^2 + 1} \right) \\ d((x^2 + 1) y) &= (4x^3) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}(x^2 + 1) y &= \int 4x^3 dx \\ (x^2 + 1) y &= x^4 + c_1\end{aligned}$$

Dividing both sides by the integrating factor  $\mu = x^2 + 1$  results in

$$y = \frac{x^4}{x^2 + 1} + \frac{c_1}{x^2 + 1}$$

which simplifies to

$$y = \frac{x^4 + c_1}{x^2 + 1}$$

### Summary

The solution(s) found are the following

$$y = \frac{x^4 + c_1}{x^2 + 1} \tag{1}$$

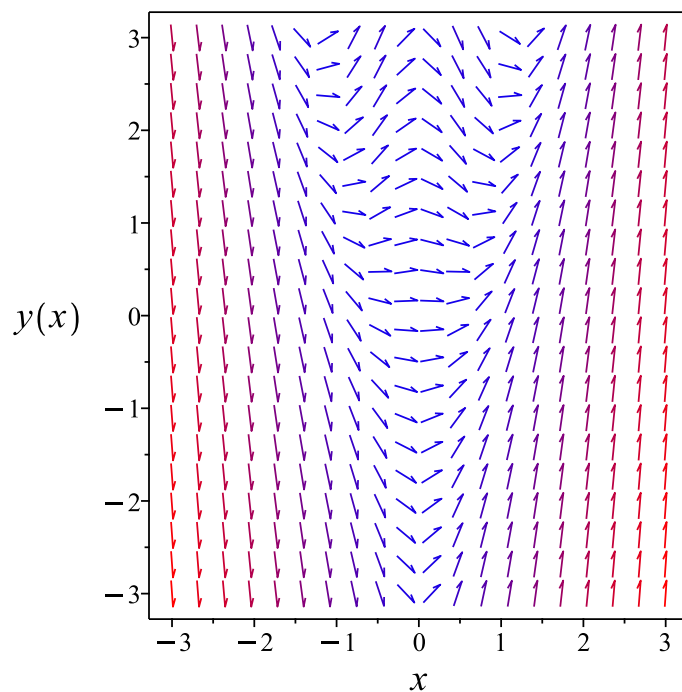


Figure 124: Slope field plot

Verification of solutions

$$y = \frac{x^4 + c_1}{x^2 + 1}$$

Verified OK.

### 5.17.2 Solving as differentialType ode

Writing the ode as

$$y' = \frac{-2yx + 4x^3}{x^2 + 1} \tag{1}$$

Which becomes

$$0 = (-x^2 - 1) dy + (2x(2x^2 - y)) dx \tag{2}$$

But the RHS is complete differential because

$$(-x^2 - 1) dy + (2x(2x^2 - y)) dx = d\left(\frac{(-2x^2 + y)^2}{4} - \frac{y^2}{4} - y\right)$$

Hence (2) becomes

$$0 = d\left(\frac{(-2x^2 + y)^2}{4} - \frac{y^2}{4} - y\right)$$

Integrating both sides gives gives these solutions

$$y = \frac{x^4 + c_1}{x^2 + 1} + c_1$$

### Summary

The solution(s) found are the following

$$y = \frac{x^4 + c_1}{x^2 + 1} + c_1 \tag{1}$$

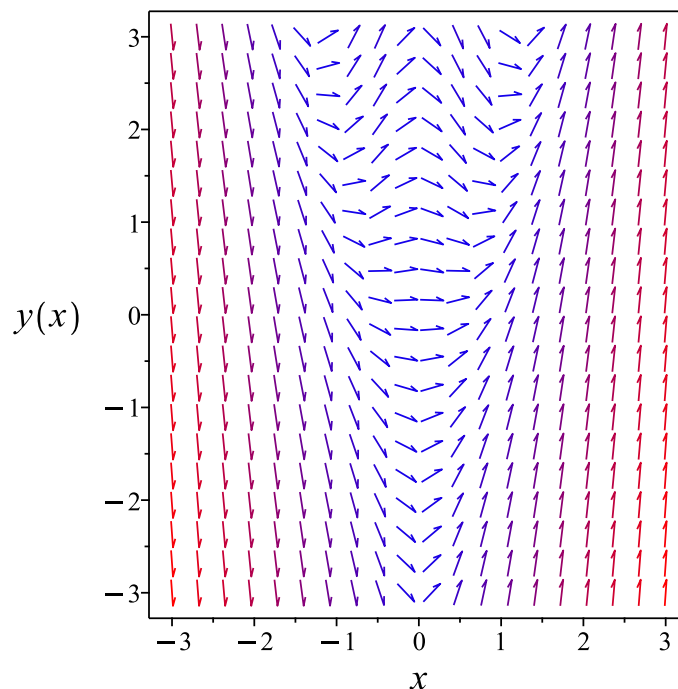


Figure 125: Slope field plot

### Verification of solutions

$$y = \frac{x^4 + c_1}{x^2 + 1} + c_1$$

Verified OK.

### 5.17.3 Solving as first order ode lie symmetry lookup ode

Writing the ode as

$$y' = -\frac{2x(-2x^2 + y)}{x^2 + 1}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

The type of this ode is known. It is of type **linear**. Therefore we do not need to solve the PDE (A), and can just use the lookup table shown below to find  $\xi, \eta$

Table 65: Lie symmetry infinitesimal lookup table for known first order ODE's

ODE class	Form	$\xi$	$\eta$
linear ode	$y' = f(x)y(x) + g(x)$	0	$e^{\int f dx}$
separable ode	$y' = f(x)g(y)$	$\frac{1}{f}$	0
quadrature ode	$y' = f(x)$	0	1
quadrature ode	$y' = g(y)$	1	0
homogeneous ODEs of Class A	$y' = f\left(\frac{y}{x}\right)$	$x$	$y$
homogeneous ODEs of Class C	$y' = (a + bx + cy)^{\frac{n}{m}}$	1	$-\frac{b}{c}$
homogeneous class D	$y' = \frac{y}{x} + g(x)F\left(\frac{y}{x}\right)$	$x^2$	$xy$
First order special form ID 1	$y' = g(x)e^{h(x)+by} + f(x)$	$\frac{e^{-\int bf(x)dx-h(x)}}{g(x)}$	$\frac{f(x)e^{-\int bf(x)dx-h(x)}}{g(x)}$
polynomial type ode	$y' = \frac{a_1x+b_1y+c_1}{a_2x+b_2y+c_2}$	$\frac{a_1b_2x-a_2b_1x-b_1c_2+b_2c_1}{a_1b_2-a_2b_1}$	$\frac{a_1b_2y-a_2b_1y-a_1c_2-a_2c_1}{a_1b_2-a_2b_1}$
Bernoulli ode	$y' = f(x)y + g(x)y^n$	0	$e^{-\int (n-1)f(x)dx}y^n$
Reduced Riccati	$y' = f_1(x)y + f_2(x)y^2$	0	$e^{-\int f_1 dx}$

The above table shows that

$$\begin{aligned}\xi(x, y) &= 0 \\ \eta(x, y) &= \frac{1}{x^2 + 1}\end{aligned}\tag{A1}$$

The next step is to determine the canonical coordinates  $R, S$ . The canonical coordinates map  $(x, y) \rightarrow (R, S)$  where  $(R, S)$  are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS\tag{1}$$

The above comes from the requirements that  $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$ . Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable  $R$  in the canonical coordinates, where  $S(R)$ . Since  $\xi = 0$  then in this special case

$$R = x$$

$S$  is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{1}{x^2+1}} dy\end{aligned}$$

Which results in

$$S = (x^2 + 1) y$$

Now that  $R, S$  are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y}\tag{2}$$

Where in the above  $R_x, R_y, S_x, S_y$  are all partial derivatives and  $\omega(x, y)$  is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{2x(-2x^2 + y)}{x^2 + 1}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= 2xy \\S_y &= x^2 + 1\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 4x^3 \quad (2A)$$

We now need to express the RHS as function of  $R$  only. This is done by solving for  $x, y$  in terms of  $R, S$  from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 4R^3$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates  $R, S$ . Integrating the above gives

$$S(R) = R^4 + c_1 \quad (4)$$

To complete the solution, we just need to transform (4) back to  $x, y$  coordinates. This results in

$$(x^2 + 1)y = x^4 + c_1$$

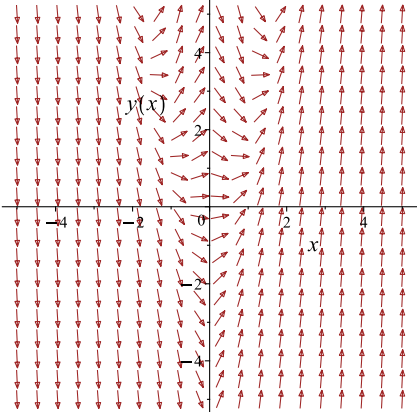
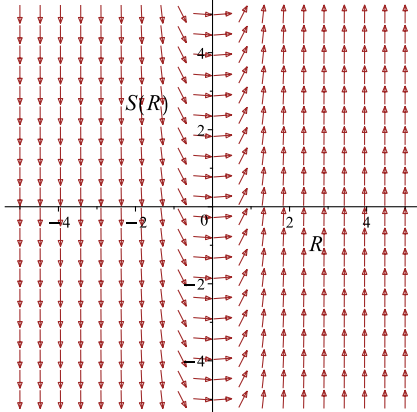
Which simplifies to

$$(x^2 + 1)y = x^4 + c_1$$

Which gives

$$y = \frac{x^4 + c_1}{x^2 + 1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in $x, y$ coordinates	Canonical coordinates transformation	ODE in canonical coordinates $(R, S)$
$\frac{dy}{dx} = -\frac{2x(-2x^2+y)}{x^2+1}$ 	$R = x$ $S = (x^2 + 1) y$	$\frac{dS}{dR} = 4R^3$ 

Summary

The solution(s) found are the following

$$y = \frac{x^4 + c_1}{x^2 + 1} \tag{1}$$

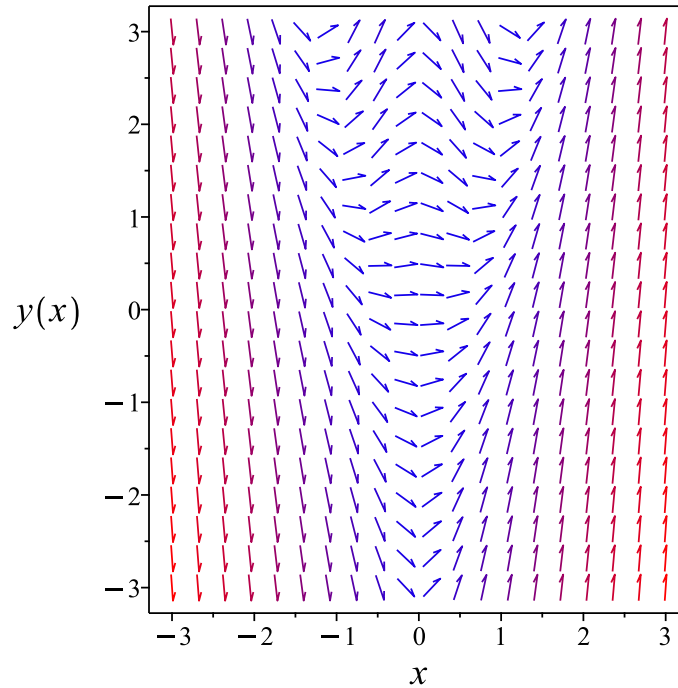


Figure 126: Slope field plot

Verification of solutions

$$y = \frac{x^4 + c_1}{x^2 + 1}$$

Verified OK.

#### 5.17.4 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$



Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(x^2 + 1) dy &= (4x^3 - 2xy) dx \\ (-4x^3 + 2xy) dx + (x^2 + 1) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -4x^3 + 2xy \\ N(x, y) &= x^2 + 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-4x^3 + 2xy) \\ &= 2x\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x^2 + 1) \\ &= 2x\end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact. The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int -4x^3 + 2xy dx$$

$$\phi = -\frac{(2x^2 - y)^2}{4} + f(y) \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = x^2 - \frac{y}{2} + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = x^2 + 1$ . Therefore equation (4) becomes

$$x^2 + 1 = x^2 - \frac{y}{2} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = \frac{y}{2} + 1$$

Integrating the above w.r.t  $y$  gives

$$\int f'(y) dy = \int \left(\frac{y}{2} + 1\right) dy$$

$$f(y) = \frac{1}{4}y^2 + y + c_1$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = -\frac{(2x^2 - y)^2}{4} + \frac{y^2}{4} + y + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = -\frac{(2x^2 - y)^2}{4} + \frac{y^2}{4} + y$$

The solution becomes

$$y = \frac{x^4 + c_1}{x^2 + 1}$$

### Summary

The solution(s) found are the following

$$y = \frac{x^4 + c_1}{x^2 + 1} \tag{1}$$

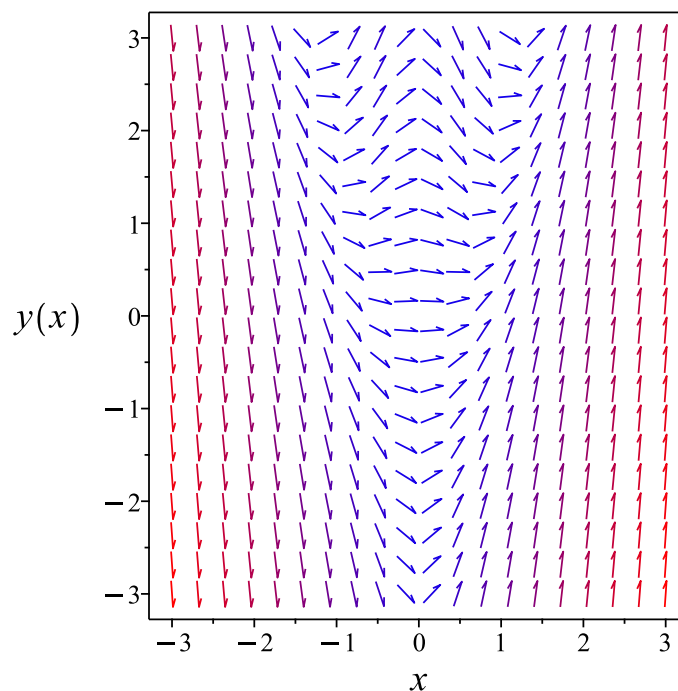


Figure 127: Slope field plot

### Verification of solutions

$$y = \frac{x^4 + c_1}{x^2 + 1}$$

Verified OK.

### 5.17.5 Maple step by step solution

Let's solve

$$2yx + (x^2 + 1)y' = 4x^3$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Isolate the derivative

$$y' = -\frac{2xy}{x^2+1} + \frac{4x^3}{x^2+1}$$

- Group terms with  $y$  on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{2xy}{x^2+1} = \frac{4x^3}{x^2+1}$$

- The ODE is linear; multiply by an integrating factor  $\mu(x)$

$$\mu(x) \left( y' + \frac{2xy}{x^2+1} \right) = \frac{4\mu(x)x^3}{x^2+1}$$

- Assume the lhs of the ODE is the total derivative  $\frac{d}{dx}(\mu(x)y)$

$$\mu(x) \left( y' + \frac{2xy}{x^2+1} \right) = \mu'(x)y + \mu(x)y'$$

- Isolate  $\mu'(x)$

$$\mu'(x) = \frac{2\mu(x)x}{x^2+1}$$

- Solve to find the integrating factor

$$\mu(x) = x^2 + 1$$

- Integrate both sides with respect to  $x$

$$\int \left( \frac{d}{dx}(\mu(x)y) \right) dx = \int \frac{4\mu(x)x^3}{x^2+1} dx + c_1$$

- Evaluate the integral on the lhs

$$\mu(x)y = \int \frac{4\mu(x)x^3}{x^2+1} dx + c_1$$

- Solve for  $y$

$$y = \frac{\int \frac{4\mu(x)x^3}{x^2+1} dx + c_1}{\mu(x)}$$

- Substitute  $\mu(x) = x^2 + 1$

$$y = \frac{\int 4x^3 dx + c_1}{x^2 + 1}$$

- Evaluate the integrals on the rhs

$$y = \frac{x^4 + c_1}{x^2 + 1}$$

### Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

### ✓ Solution by Maple

Time used: 0.0 (sec). Leaf size: 17

```
dsolve((1+x^2)*diff(y(x),x)+2*x*y(x)=4*x^3,y(x), singsol=all)
```

$$y(x) = \frac{x^4 + c_1}{x^2 + 1}$$

### ✓ Solution by Mathematica

Time used: 0.032 (sec). Leaf size: 19

```
DSolve[(1+x^2)*y'[x]+2*x*y[x]==4*x^3,y[x],x,IncludeSingularSolutions -> True]
```

$$y(x) \rightarrow \frac{x^4 + c_1}{x^2 + 1}$$

## 5.18 problem 22

5.18.1 Solving as exact ode . . . . .	573
5.18.2 Maple step by step solution . . . . .	576

Internal problem ID [3131]

Internal file name [OUTPUT/2623\_Sunday\_June\_05\_2022\_03\_23\_03\_AM\_57615570/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 22.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**exact**"

Maple gives the following as the ode type

[\_exact]

$$e^x \sin(y) - y \sin(yx) + (\cos(y) e^x - x \sin(yx)) y' = 0$$

### 5.18.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (\cos(y) e^x - x \sin(xy)) dy &= (-\sin(y) e^x + y \sin(xy)) dx \\ (\sin(y) e^x - y \sin(xy)) dx &+ (\cos(y) e^x - x \sin(xy)) dy = 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= \sin(y) e^x - y \sin(xy) \\ N(x, y) &= \cos(y) e^x - x \sin(xy) \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} (\sin(y) e^x - y \sin(xy)) \\ &= \cos(y) e^x - \sin(xy) - y \cos(xy) x \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (\cos(y) e^x - x \sin(xy)) \\ &= \cos(y) e^x - \sin(xy) - y \cos(xy) x \end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \sin(y) e^x - y \sin(xy) dx \\ \phi &= \sin(y) e^x + \cos(xy) + f(y)\end{aligned}\tag{3}$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = \cos(y) e^x - x \sin(xy) + f'(y)\tag{4}$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = \cos(y) e^x - x \sin(xy)$ . Therefore equation (4) becomes

$$\cos(y) e^x - x \sin(xy) = \cos(y) e^x - x \sin(xy) + f'(y)\tag{5}$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where  $c_1$  is constant of integration. Substituting this result for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = \sin(y) e^x + \cos(xy) + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = \sin(y) e^x + \cos(xy)$$

### Summary

The solution(s) found are the following

$$e^x \sin(y) + \cos(yx) = c_1\tag{1}$$



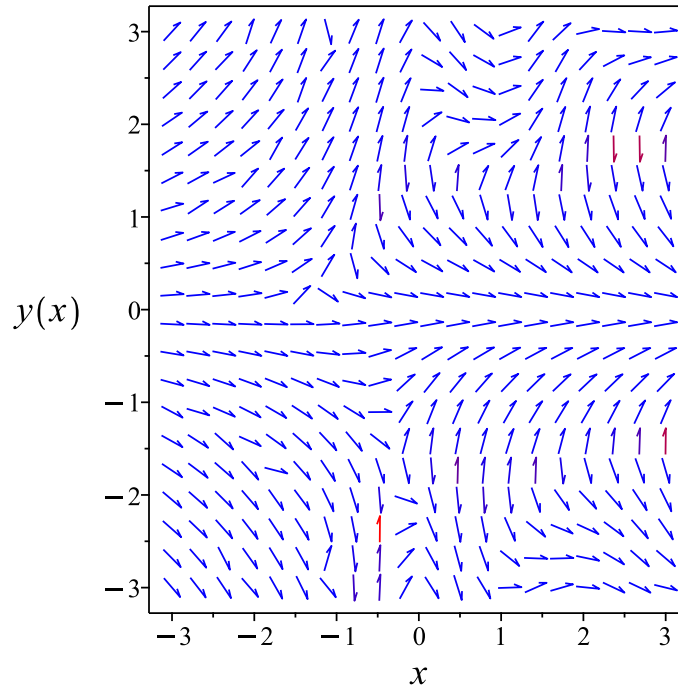


Figure 128: Slope field plot

### Verification of solutions

$$e^x \sin(y) + \cos(yx) = c_1$$

Verified OK.

### 5.18.2 Maple step by step solution

Let's solve

$$e^x \sin(y) - y \sin(yx) + (\cos(y) e^x - x \sin(yx)) y' = 0$$

- Highest derivative means the order of the ODE is 1  
 $y'$
- Check if ODE is exact
  - ODE is exact if the lhs is the total derivative of a  $C^2$  function  
 $F'(x, y) = 0$
  - Compute derivative of lhs  
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$

- Evaluate derivatives
 
$$\cos(y) e^x - \sin(xy) - y \cos(xy) x = \cos(y) e^x - \sin(xy) - y \cos(xy) x$$
- Condition met, ODE is exact
- Exact ODE implies solution will be of this form
 
$$\left[ F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$
- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$ 

$$F(x, y) = \int (\sin(y) e^x - y \sin(xy)) dx + f_1(y)$$
- Evaluate integral
 
$$F(x, y) = \sin(y) e^x + \cos(xy) + f_1(y)$$
- Take derivative of  $F(x, y)$  with respect to  $y$ 

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$
- Compute derivative
 
$$\cos(y) e^x - x \sin(xy) = \cos(y) e^x - x \sin(xy) + \frac{d}{dy} f_1(y)$$
- Isolate for  $\frac{d}{dy} f_1(y)$ 

$$\frac{d}{dy} f_1(y) = 0$$
- Solve for  $f_1(y)$ 

$$f_1(y) = 0$$
- Substitute  $f_1(y)$  into equation for  $F(x, y)$ 

$$F(x, y) = \sin(y) e^x + \cos(xy)$$
- Substitute  $F(x, y)$  into the solution of the ODE
 
$$\sin(y) e^x + \cos(xy) = c_1$$
- Solve for  $y$ 

$$y = \text{RootOf}(\_Zx - \arccos(-\sin(\_Z) e^x + c_1))$$

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

### ✓ Solution by Maple

Time used: 0.265 (sec). Leaf size: 16

```
dsolve((exp(x)*sin(y(x))-y(x)*sin(x*y(x)))+(exp(x)*cos(y(x))-x*sin(x*y(x)))*diff(y(x),x)=0,y(x),x,
```

$$e^x \sin(y(x)) + \cos(xy(x)) + c_1 = 0$$

### ✓ Solution by Mathematica

Time used: 0.58 (sec). Leaf size: 19

```
DSolve[(Exp[x]*Sin[y[x]]-y[x]*Sin[x*y[x]])+(Exp[x]*Cos[y[x]]-x*Ssin[x*y[x]])*y'[x]==0,y[x],x,
```

$$\text{Solve}[e^x \sin(y(x)) + \cos(xy(x)) = c_1, y(x)]$$

## 5.19 problem 24

5.19.1 Solving as exact ode . . . . .	579
5.19.2 Maple step by step solution . . . . .	583

Internal problem ID [3132]

Internal file name [OUTPUT/2624\_Sunday\_June\_05\_2022\_03\_23\_11\_AM\_89760593/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 24.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : **"exact"**

Maple gives the following as the ode type

[\_exact]

$$(x e^y + y - x^2) y' - 2yx + e^y = -x$$

### 5.19.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned} (x e^y + y - x^2) dy &= (2xy - e^y - x) dx \\ (-2xy + e^y + x) dx + (x e^y + y - x^2) dy &= 0 \end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -2xy + e^y + x \\ N(x, y) &= x e^y + y - x^2 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-2xy + e^y + x) \\ &= e^y - 2x \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x e^y + y - x^2) \\ &= e^y - 2x \end{aligned}$$

Since  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ , then the ODE is exact The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t.  $x$  gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -2xy + e^y + x dx \\ \phi &= x e^y - \left(y - \frac{1}{2}\right) x^2 + f(y)\end{aligned}\quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\begin{aligned}\frac{\partial \phi}{\partial y} &= x e^y - x^2 + f'(y) \\ &= x(e^y - x) + f'(y)\end{aligned}\quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = x e^y + y - x^2$ . Therefore equation (4) becomes

$$x e^y + y - x^2 = x(e^y - x) + f'(y)\quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = y$$

Integrating the above w.r.t  $y$  gives

$$\begin{aligned}\int f'(y) dy &= \int (y) dy \\ f(y) &= \frac{y^2}{2} + c_1\end{aligned}$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = x e^y - \left(y - \frac{1}{2}\right) x^2 + \frac{y^2}{2} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = x e^y - \left(y - \frac{1}{2}\right) x^2 + \frac{y^2}{2}$$

### Summary

The solution(s) found are the following

$$x e^y - \left(y - \frac{1}{2}\right) x^2 + \frac{y^2}{2} = c_1 \quad (1)$$

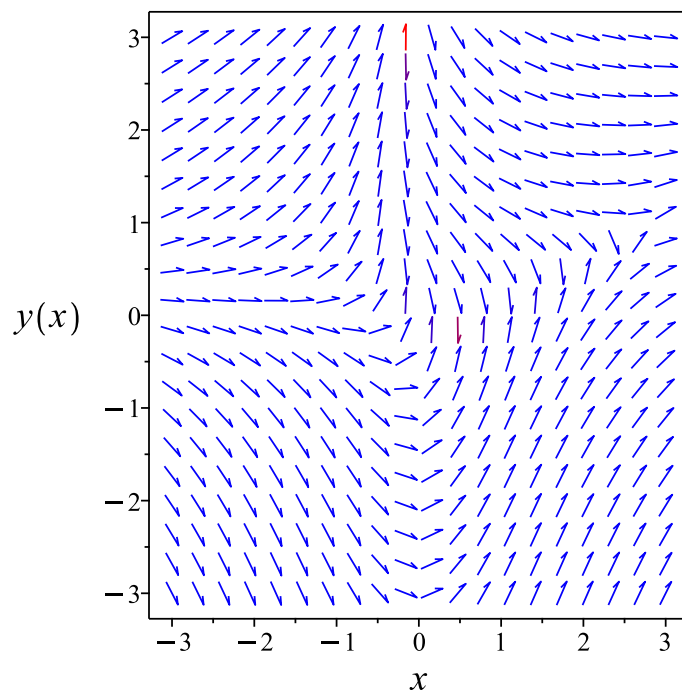


Figure 129: Slope field plot

### Verification of solutions

$$x e^y - \left(y - \frac{1}{2}\right) x^2 + \frac{y^2}{2} = c_1$$

Verified OK.

### 5.19.2 Maple step by step solution

Let's solve

$$(x e^y + y - x^2) y' - 2yx + e^y = -x$$

- Highest derivative means the order of the ODE is 1

$y'$

- Check if ODE is exact

- ODE is exact if the lhs is the total derivative of a  $C^2$  function

$$F'(x, y) = 0$$

- Compute derivative of lhs

$$F'(x, y) + \left( \frac{\partial}{\partial y} F(x, y) \right) y' = 0$$

- Evaluate derivatives

$$e^y - 2x = e^y - 2x$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$\left[ F(x, y) = c_1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$

- Solve for  $F(x, y)$  by integrating  $M(x, y)$  with respect to  $x$

$$F(x, y) = \int (-2xy + e^y + x) dx + f_1(y)$$

- Evaluate integral

$$F(x, y) = -x^2y + x e^y + \frac{x^2}{2} + f_1(y)$$

- Take derivative of  $F(x, y)$  with respect to  $y$

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$

- Compute derivative

$$x e^y + y - x^2 = -x^2 + x e^y + \frac{d}{dy} f_1(y)$$

- Isolate for  $\frac{d}{dy} f_1(y)$

$$\frac{d}{dy} f_1(y) = y$$

- Solve for  $f_1(y)$

$$f_1(y) = \frac{y^2}{2}$$

- Substitute  $f_1(y)$  into equation for  $F(x, y)$



$$F(x, y) = -x^2y + xe^y + \frac{x^2}{2} + \frac{y^2}{2}$$

- Substitute  $F(x, y)$  into the solution of the ODE

$$-x^2y + xe^y + \frac{x^2}{2} + \frac{y^2}{2} = c_1$$

- Solve for  $y$

$$y = \text{RootOf}(2x^2Z - 2e^{-Z}x - Z^2 - x^2 + 2c_1)$$

### Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying Chini
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```

### ✓ Solution by Maple

Time used: 0.015 (sec). Leaf size: 28

```
dsolve((x*exp(y(x))+y(x)-x^2)*diff(y(x),x)=(2*x*y(x) -exp(y(x))-x),y(x), singsol=all)
```

$$-x^2y(x) + e^{y(x)}x + \frac{x^2}{2} + \frac{y(x)^2}{2} + c_1 = 0$$

### ✓ Solution by Mathematica

Time used: 0.315 (sec). Leaf size: 35

```
DSolve[(x*Exp[y[x]]+y[x]-x^2)*y'[x]==(2*x*y[x] -Exp[y[x]]-x),y[x],x,IncludeSingularSolutions
```

$$\text{Solve}\left[x^2(-y(x)) + \frac{x^2}{2} + xe^{y(x)} + \frac{y(x)^2}{2} = c_1, y(x)\right]$$

## 5.20 problem 25

5.20.1 Solving as exact ode . . . . . 585

Internal problem ID [3133]

Internal file name [OUTPUT/2625\_Sunday\_June\_05\_2022\_03\_23\_15\_AM\_43739337/index.tex]

**Book:** Differential equations with applications and historical notes, George F. Simmons, 1971

**Section:** Chapter 2, End of chapter, page 61

**Problem number:** 25.

**ODE order:** 1.

**ODE degree:** 1.

The type(s) of ODE detected by this program : "**exactWithIntegrationFactor**"

Maple gives the following as the ode type

[`y=\_G(x,y)']

$$-(x e^x - e^y y) y' = -e^x (x + 1)$$

### 5.20.1 Solving as exact ode

Entering Exact first order ODE solver. (Form one type)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function  $\phi(x, y) = c$  where  $c$  is constant, that satisfies the ode. Taking derivative of  $\phi$  w.r.t.  $x$  gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine  $\phi(x, y)$  but at least we know now that we can do that since the condition  $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$  is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (-x e^x + y e^y) dy &= (-e^x(x + 1)) dx \\ (e^x(x + 1)) dx + (-x e^x + y e^y) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= e^x(x + 1) \\ N(x, y) &= -x e^x + y e^y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(e^x(x + 1)) \\ &= 0 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x e^x + y e^y) \\ &= e^x(-x - 1) \end{aligned}$$

Since  $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$ , then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left( \frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= -\frac{1}{x e^x - y e^y} ((0) - (-x e^x - e^x)) \\ &= -\frac{e^x(x + 1)}{x e^x - y e^y} \end{aligned}$$

Since  $A$  depends on  $y$ , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= \frac{e^{-x}}{x+1} ((-x e^x - e^x) - (0)) \\ &= -1 \end{aligned}$$

Since  $B$  does not depend on  $x$ , it can be used to obtain an integrating factor. Let the integrating factor be  $\mu$ . Then

$$\begin{aligned} \mu &= e^{\int B \, dy} \\ &= e^{\int -1 \, dy} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-y} \\ &= e^{-y} \end{aligned}$$

$M$  and  $N$  are now multiplied by this integrating factor, giving new  $M$  and new  $N$  which are called  $\overline{M}$  and  $\overline{N}$  so not to confuse them with the original  $M$  and  $N$ .

$$\begin{aligned} \overline{M} &= \mu M \\ &= e^{-y}(e^x(x+1)) \\ &= (x+1)e^{-y+x} \end{aligned}$$

And

$$\begin{aligned} \overline{N} &= \mu N \\ &= e^{-y}(-x e^x + y e^y) \\ &= -x e^{-y+x} + y \end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned} \overline{M} + \overline{N} \frac{dy}{dx} &= 0 \\ ((x+1)e^{-y+x}) + (-x e^{-y+x} + y) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function  $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t.  $x$  gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \bar{M} dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int (x + 1) e^{-y+x} dx$$

$$\phi = x e^{-y+x} + f(y) \quad (3)$$

Where  $f(y)$  is used for the constant of integration since  $\phi$  is a function of both  $x$  and  $y$ . Taking derivative of equation (3) w.r.t  $y$  gives

$$\frac{\partial \phi}{\partial y} = -x e^{-y+x} + f'(y) \quad (4)$$

But equation (2) says that  $\frac{\partial \phi}{\partial y} = -x e^{-y+x} + y$ . Therefore equation (4) becomes

$$-x e^{-y+x} + y = -x e^{-y+x} + f'(y) \quad (5)$$

Solving equation (5) for  $f'(y)$  gives

$$f'(y) = y$$

Integrating the above w.r.t  $y$  gives

$$\int f'(y) dy = \int (y) dy$$

$$f(y) = \frac{y^2}{2} + c_1$$

Where  $c_1$  is constant of integration. Substituting result found above for  $f(y)$  into equation (3) gives  $\phi$

$$\phi = x e^{-y+x} + \frac{y^2}{2} + c_1$$

But since  $\phi$  itself is a constant function, then let  $\phi = c_2$  where  $c_2$  is new constant and combining  $c_1$  and  $c_2$  constants into new constant  $c_1$  gives the solution as

$$c_1 = x e^{-y+x} + \frac{y^2}{2}$$

### Summary

The solution(s) found are the following

$$x e^{-y+x} + \frac{y^2}{2} = c_1 \quad (1)$$

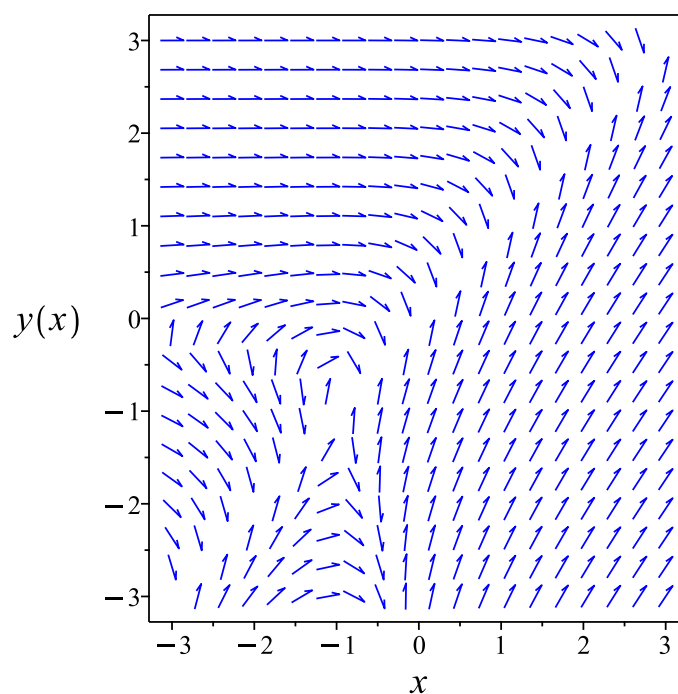


Figure 130: Slope field plot

### Verification of solutions

$$x e^{-y+x} + \frac{y^2}{2} = c_1$$

Verified OK.

## Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
trying Bernoulli  
trying separable  
trying inverse linear  
trying homogeneous types:  
trying Chini  
differential order: 1; looking for linear symmetries  
trying exact  
<- exact successful`
```

### ✓ Solution by Maple

Time used: 0.016 (sec). Leaf size: 20

```
dsolve(exp(x)*(1+x)=(x*exp(x)-y(x)*exp(y(x)))*diff(y(x),x),y(x), singsol=all)
```

$$x e^{x-y(x)} + \frac{y(x)^2}{2} + c_1 = 0$$

### ✓ Solution by Mathematica

Time used: 0.307 (sec). Leaf size: 26

```
DSolve[Exp[x]*(1+x)==(x*Exp[x]-y[x]*Exp[y[x]])*y'[x],y[x],x,IncludeSingularSolutions -> True
```

$$\text{Solve} \left[ -\frac{1}{2}y(x)^2 - x e^{x-y(x)} = c_1, y(x) \right]$$