# HW1, Math 228A 

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## 1 Problem description

$$
\begin{align*}
& \text { Math 228A } \\
& \text { Homework } 1 \\
& \langle u, v\rangle=\int_{0}^{1} u v d x . \\
& \text { (c) It can be shown that the eigenfunctions, } \phi_{j}(x) \text {, form a complete set in } L^{2}[0,1] \text {. This } \\
& \text { means that for any } f \in L^{2}[0,1], f(x)=\sum_{j} \alpha_{j} \phi_{j}(x) \text {. Express the solution to } \\
& u_{x x}=f, u_{x}(0)=u_{x}(1)=0,  \tag{1}\\
& \text { as a series solution of the eigenfunctions. } \\
& \text { (d) Note that equation (1) does not have a solution for all } f \text {. Express the condition for } \\
& \text { existence of a solution in terms of the eigenfunctions of } L \text {. } \\
& \text { 2. Define the functional } F: X \rightarrow \Re \text { by } \\
& F(u)=\int_{0}^{1} \frac{1}{2}\left(u_{x}\right)^{2}+f u d x, \\
& \text { where } X \text { is the space of real valued functions on }[0,1] \text { that have at least one continuous } \\
& \text { derivative and are zero at } x=0 \text { and } x=1 \text {. The Frechet derivative of } F \text { at a point } u \text { is } \\
& \text { defined to be the linear operator } F^{\prime}(u) \text { for which } \\
& F(u+v)=F(u)+F^{\prime}(u) v+R(v), \\
& \text { where } \\
& \lim _{\|v\| \rightarrow 0} \frac{\|R(v)\|}{\|v\|}=0 . \\
& \text { One way to compute the derivative is } \\
& F^{\prime}(u) v=\lim _{\epsilon \rightarrow 0} \frac{F(u+\epsilon v)-F(u)}{\epsilon} . \\
& \text { Note that this looks just like a directional derivative. } \\
& \text { (a) Compute the Frechet derivative of } F \text {. } \\
& \text { (b) } u \in X \text { is a critical point of } F \text { if } F^{\prime}(u) v=0 \text { for all } v \in X \text {. Show that if } u \text { is a solution to } \\
& \text { the Poisson equation } \\
& u_{x x}=f, \quad u(0)=u(1)=0, \\
& \text { then it is a critical point of } F \text {. } \\
& \text { Finite element methods are based on these "weak formulations" of the problem. The Ritz } \\
& \text { method is based on minimizing } F \text { and the Galerkin method is based on finding the critical } \\
& \text { points of } F^{\prime}(u) \text {. }
\end{align*}
$$

Figure 1: problem description

## 2 Problem 1

$L$ is a second order differential operator defined by $L u \equiv u_{x x}$ with boundary conditions on $u$ given as $u_{x}(0)=u_{x}(1)=0$

## 2.1 part a

Let $\phi(x)$ be an eigenfunction of the operator $L$ associated with an eigenvalue $\lambda$. To obtain the eigenfunctions and eigenvalues, we solve an eigenvalue problem $L \phi=\lambda \phi$ where $\lambda$ is scalar. Hence the problem is to solve the differential equation

$$
\begin{equation*}
\phi_{x x}-\lambda \phi=0 \tag{1}
\end{equation*}
$$

with B.C. given as $\phi^{\prime}(0)=\phi^{\prime}(1)=0$. The characteristic equation is

$$
r^{2}-\lambda=0
$$

The roots are $r= \pm \sqrt{\lambda}$, therefore the solution to the eigenvalue problem (1) is

$$
\begin{equation*}
\phi(x)=c_{1} e^{\sqrt{\lambda} x}+c_{2} e^{-\sqrt{\lambda} x} \tag{2}
\end{equation*}
$$

Where $c_{1}, c_{2}$ are constants.

$$
\begin{equation*}
\phi^{\prime}(x)=c_{1} \sqrt{\lambda} e^{\sqrt{\lambda} x}-\sqrt{\lambda} c_{2} e^{-\sqrt{\lambda} x} \tag{3}
\end{equation*}
$$

First we determine the allowed values of the eigenvalues $\lambda$ which satisfies the boundary conditions.

1. Assume $\lambda=0$ The solution (2) becomes $\phi(x)=c_{1}+c_{2}$. Hence the solution is a constant. In other words, when the eigenvalue is zero, the eigenfunction is a constant. Let us now see if this eigenfunction satisfies the B.C. Since $\phi(x)$ is constant, then $\phi^{\prime}(x)=0$, and this does satisfy the B.C. at both $x=0$ and $x=1$. Hence $\lambda=0$ is an eigenvalue with a corresponding eigenfunction being a constant. $\overline{\text { We can take the constant as }} 1$.
2. Assume $\lambda>0$ From the first $B C$ we have, from (3), that $\phi^{\prime}(0)=0=c_{1} \sqrt{\lambda}-\sqrt{\lambda} c_{2}$ or

$$
c_{1}=c_{2}
$$

and from the second BC we have that $\phi^{\prime}(1)=0=c_{1} \sqrt{\lambda} e^{\sqrt{\lambda}}-\sqrt{\lambda} c_{2} e^{-\sqrt{\lambda}}$ or

$$
c_{1} e^{\sqrt{\lambda}}-c_{2} e^{-\sqrt{\lambda}}=0
$$

From the above 2 equations, we find that $e^{\sqrt{\lambda}}=e^{-\sqrt{\lambda}}$ which is not possible for positive $\lambda$. Hence $\lambda$ can not be positive.
3. Assume $\lambda<0$. Let $\lambda=-\beta^{2}$ form some positive $\beta$. Then the solution (2) becomes

$$
\phi(x)=c_{1} e^{i \beta x}+c_{2} e^{-i \beta x}
$$

which can be transformed using the Euler relation to obtain

$$
\begin{align*}
\phi(x) & =c_{1} \cos \beta x+c_{2} \sin \beta x \\
\phi^{\prime}(x) & =-c_{1} \beta \sin \beta x+c_{2} \beta \cos \beta x \tag{4}
\end{align*}
$$

Now consider the BC's. Since $\phi^{\prime}(0)=0$ we obtain $c_{2}=0$ and from $\phi^{\prime}(1)=0$ we obtain $0=c_{1} \beta \sin \beta$ and hence for non trivial solution, i.e. for $c_{1} \neq 0$, we must have that

$$
\sin \beta=0
$$

or

$$
\beta= \pm n \pi
$$

but since $\beta$ is positive, we consider only $\beta_{n}=n \pi$, where $n$ is positive integer $n=1,2,3, \cdots$
Conclusion: The eigenvalues are

$$
\lambda_{n}=-\left(\beta_{n}\right)^{2}=-(n \pi)^{2}=\left\{0,-\pi^{2},-(2 \pi)^{2},-(3 \pi)^{2}, \cdots\right\}
$$

And the corresponding eigenfunctions are $\phi_{n}(x)=\cos \beta_{n} x=\cos n \pi x=\{1, \cos \pi x, \cos 2 \pi x, \cos 3 \pi x, \cdots\}$ where $n=0,1,2, \cdots$

## 2.2 part (b)

Given inner product defined as $\langle u, v\rangle=\int_{0}^{1} u v d x$, then

$$
\begin{aligned}
\left\langle\phi_{n}, \phi_{m}\right\rangle & =\int_{0}^{1}\left(\cos \beta_{n} x\right)\left(\cos \beta_{m} x\right) d x \\
& =\int_{0}^{1}(\cos n \pi x)(\cos m \pi x) d x \\
& = \begin{cases}0 & n \neq m \\
\frac{1}{2} & n=m\end{cases}
\end{aligned}
$$

Also, the first eigenfucntion, $\phi_{0}(x)=1$ is orthogonal to all other eigenfunctions, since $\int_{0}^{1}(\cos n \pi x) d x=\frac{1}{n \pi}[\sin n \pi x]_{0}^{1}=0$ for any integer $n>0$.
Hence all the eigenfunctions are orthogonal to each others in $L^{2}[0,1]$ space.

### 2.3 Part (c)

Given

$$
u_{x x}=f
$$

$u_{x}(0)=u_{x}(1)=0$. This is $L u=f$. We have found the eigenfunctions $\phi(x)$ of $L$ above. These are basis of the function space of $L$ where $f$ resides in. We can express $f$ as a linear combination of the eigenfunctions of the operator $L$, hence we write

$$
f(x)=\sum_{n=0}^{\infty} a_{n} \phi_{n}(x)
$$

where $\phi_{n}(x)$ is the $n^{\text {th }}$ eigenfunction of $L$ and $a_{n}$ is the corresponding coordinate (scalar). Therefore the differential equation above can be written as

$$
\begin{equation*}
L u=f(x)=\sum_{n=0} a_{n} \phi_{n}(x) \tag{1}
\end{equation*}
$$

But since

$$
L \phi_{n}=\lambda_{n} \phi_{n}
$$

Then

$$
L^{-1}=\frac{1}{\lambda_{n}}
$$

Therefore, using (1), the solution is

$$
\begin{equation*}
u(x)=\sum_{n}\left(\frac{a_{n}}{\lambda_{n}}\right) \phi_{n}(x) \tag{2}
\end{equation*}
$$

Now to find $a_{n}$, using $f(x)=\sum_{n} a_{n} \phi_{n}(x)$, we multiply each side by an eigenfunction, say $\phi_{m}(x)$ and integrate

$$
\begin{aligned}
\int_{0}^{1} \phi_{m}(x) f(x) d x & =\int_{0}^{1} \phi_{m}(x) \sum_{n} a_{n} \phi_{n}(x)(x) d x \\
& =\int_{0}^{1} \sum_{n} a_{n} \phi_{m}(x) \phi_{n}(x) d x \\
& =\sum_{n} a_{n} \int_{0}^{1} \phi_{m}(x) \phi_{n}(x) d x
\end{aligned}
$$

The RHS is $1 / 2$ when $n=m$ and zero otherwise, hence the above becomes

$$
\int_{0}^{1} \phi_{n}(x) f(x) d x=\frac{a_{n}}{2}
$$

Or

$$
\begin{equation*}
a_{n}=2 \int_{0}^{1} \cos (n \pi x) f(x) d x \tag{3}
\end{equation*}
$$

Where $a_{n}$ as given by (3).
If we know $f(x)$ we can determine $a_{n}$ and hence the solution is now found.

## 2.4 part (d)

The solution found above

$$
u(x)=\sum_{n}\left(\frac{a_{n}}{\lambda_{n}}\right) \phi_{n}(x)
$$

Is not possible for all $f$. Only an $f$ which has $a_{0}=0$ is possible. This is because $\lambda_{0}=0$, then $a_{0}$ has to be zero to obtain a solution (since $L^{-1}$ does not exist if an eigenvalue is zero).
$a_{0}=0$ implies, by looking at (3) above, that when $n=0$ we have

$$
0=\int_{0}^{1} f(x) d x
$$

So only the functions $f(x)$ which satisfy the above can be a solution to $L u=f$ with the B.C. given.

To review: We found that $\lambda=0$ to be a valid eigenvalue due to the B.C. being Von Neumann boundary conditions. This in resulted in $a_{0}$ having to be zero. This implied that $\int_{0}^{1} f(x) d x=0$.
Having a zero eigenvalue effectively removes one of the space dimensions that $f(x)$ can resides in.

In addition to this restriction, the function $f(x)$ is assumed to meet the Dirichlet conditions for Fourier series expansion, and these are

1. $f(x)$ must have a finite number of extrema in any given interval
2. $f(x)$ must have a finite number of discontinuities in any given interval
3. $f(x)$ must be absolutely integrable over a period.
4. $f(x)$ must be bounded

## 3 Problem 2

### 3.1 Part (a)

Applying the definition given

$$
\begin{equation*}
F^{\prime}(u) v=\lim _{\varepsilon \rightarrow 0} \frac{F(u+\varepsilon v)-F(u)}{\varepsilon} \tag{1}
\end{equation*}
$$

And using $F(u)=\int_{0}^{1} \frac{1}{2}\left(u_{x}\right)^{2}+f u d x$, then (1) becomes

$$
F^{\prime}(u) v=\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon}\left(\int_{0}^{1} \frac{1}{2}\left[(u+\varepsilon v)_{x}\right]^{2}+f(u+\varepsilon v) d x-\int_{0}^{1} \frac{1}{2}\left(u_{x}\right)^{2}+f u d x\right)
$$

Simplify the above, we obtain

$$
F^{\prime}(u) v=\lim _{\varepsilon \rightarrow 0}\left(\int_{0}^{1} \frac{\varepsilon}{2} v_{x}^{2} d x+\int_{0}^{1} u_{x} v_{x} d x+\int_{0}^{1} f v d x\right)
$$

Hence, as $\varepsilon \rightarrow 0$ only the first integral above vanishes (since $v_{x}$ is bounded), and we have

$$
\begin{equation*}
F^{\prime}(u) v=\int_{0}^{1} u_{x} v_{x}+f v d x \tag{1A}
\end{equation*}
$$

### 3.2 Part (b)

The solution to $u_{x x}=f(x)$ with $u(0)=u(1)=0$ was found in class to be

$$
\begin{equation*}
u(x)=\sum_{n}\left(\frac{a_{n}}{\lambda_{n}}\right) \phi_{n}(x) \tag{2}
\end{equation*}
$$

where

$$
\phi_{n}(x)=\sin (n \pi x)
$$

are the eigenfunctions associated with the eigenvalues $\lambda_{n}=-n^{2} \pi^{2}$.
Now we can use this solution in the definition of $F^{\prime}(u) v$ found in (1A) from part (a). Substitute $u(x)$ from (2) into (1A), and also substitute $f=\sum_{n} a_{n} \phi_{n}(x)$ into (1A), we obtain

$$
\begin{equation*}
F^{\prime}(u) v=\int_{0}^{1}\left(\sum_{n}\left(\frac{a_{n}}{\lambda_{n}}\right) \phi_{n}(x)\right)^{\prime} v^{\prime}+\left(\sum_{n} a_{n} \phi_{n}(x)\right) v d x \tag{4}
\end{equation*}
$$

We need to show that the above becomes zero for any $v(x) \in X$.

$$
\begin{align*}
F^{\prime}(u) v & =\int_{0}^{1} \sum_{n} v^{\prime}\left(\frac{a_{n}}{\lambda_{n}}\right) \phi_{n}^{\prime}(x)+\sum_{n} v a_{n} \phi_{n}(x) d x \\
& =\int_{0}^{1} \sum_{n}\left(v^{\prime}\left(\frac{a_{n}}{\lambda_{n}}\right) \phi_{n}^{\prime}(x)+v a_{n} \phi_{n}(x)\right) d x \\
& =\sum_{n} a_{n}\left(\int_{0}^{1} \frac{1}{\lambda_{n}} v^{\prime} \phi_{n}^{\prime}(x)+v \phi_{n}(x) d x\right) \tag{5}
\end{align*}
$$

Now we pay attention to the integral term above. If we can show this is zero, then we are done.

$$
\begin{align*}
I & =\frac{1}{\lambda_{n}} \int_{0}^{1} v^{\prime} \phi_{n}^{\prime}(x)+\int_{0}^{1} v \phi_{n}(x) d x \\
& =I_{1}+I_{2} \tag{6}
\end{align*}
$$

Integrate by parts $I_{1}$

$$
\begin{aligned}
I_{1} & =\frac{1}{\lambda_{n}} \int_{0}^{1} \overbrace{\phi_{n}^{\prime}(x) v^{\prime} d x}^{u d v} \\
& =\frac{1}{\lambda_{n}}\left(\left[\phi_{n}^{\prime}(x) v\right]_{0}^{1}-\int_{0}^{1} v(x) \phi_{n}^{\prime \prime}(x) d x\right) \\
& =\frac{1}{\lambda_{n}}(\overbrace{\left[v(1) \phi_{n}^{\prime}(1)-v(0) \phi_{n}^{\prime}(0)\right]}^{\text {zero due to boundaries on } v(x) \in X}-\int_{0}^{1} v(x) \phi_{n}^{\prime \prime}(x) d x) \\
& =-\frac{1}{\lambda_{n}} \int_{0}^{1} v(x) \phi_{n}^{\prime \prime}(x) d x
\end{aligned}
$$

But since $\phi_{n}(x)=\sin n \pi x$, then $\phi_{n}^{\prime}(x)=n \pi \cos n \pi x$ and $\phi_{n}^{\prime \prime}(x)=-n^{2} \pi^{2} \sin n \pi x=-n^{2} \pi^{2}$ $\phi_{n}(x)$ then

$$
I_{1}=\frac{n^{2} \pi^{2}}{\lambda_{n}} \int_{0}^{1} v(x) \phi_{n}(x) d x
$$

But also $\lambda_{n}=-n^{2} \pi^{2}$ hence the above becomes

$$
I_{1}=-\int_{0}^{1} v(x) \phi_{n}(x) d x
$$

Therefore (6) can be written as

$$
\begin{aligned}
I & =I_{1}+I_{2} \\
& =-\int_{0}^{1} v(x) \phi_{n}(x) d x+\int_{0}^{1} v(x) \phi_{n}(x) d x \\
& =0
\end{aligned}
$$

Therefore, from (5), we see that

$$
F^{\prime}(u) v=0
$$

Hence we showed that if $u$ is solution to $u_{x x}=f$ with $u(0)=u(1)=0$, then $F^{\prime}(u) v=0$.

## 4 Problem (3)

### 4.1 Part (a)

Notations used: let $\tilde{f}$ to mean the approximate discrete solution at a grid point. Let $f$ to mean the exact solution.

Using the method of undetermined coefficients, let the second derivative approximation be

$$
\begin{equation*}
\tilde{f}^{\prime \prime}(x)=a f\left(x-\frac{h}{2}\right)+b f(x)+c f(x+h) \tag{1}
\end{equation*}
$$

Where $a, b, c$ are constants to be found. Now using Taylor expansion, since

$$
f(x+\Delta)=f(x)+\Delta f^{\prime}(x)+\frac{\Delta^{2}}{2!} f^{\prime \prime}(x)+\frac{\Delta^{3}}{3!} f^{\prime \prime \prime}(x)+O\left(h^{4}\right)
$$

Hence apply the above to each of the terms in the RHS of (1) and simplify

$$
\begin{aligned}
f\left(x-\frac{h}{2}\right) & =f(x)-\frac{h}{2} f^{\prime}(x)+\frac{\left(-\frac{h}{2}\right)^{2}}{2!} f^{\prime \prime}(x)+\frac{\left(-\frac{h}{2}\right)^{3}}{3!} f^{\prime \prime}(x)+\frac{\left(-\frac{h}{2}\right)^{4}}{4!} f^{(4)}(x)+O\left(h^{5}\right) \\
f(x) & =f(x) \\
f(x+h) & =f(x)+h f^{\prime}(x)+\frac{h^{2}}{2!} f^{\prime \prime}(x)+\frac{h^{3}}{3!} f^{\prime \prime}(x)+\frac{h^{4}}{4!} f^{(4)}(x)+O\left(h^{5}\right)
\end{aligned}
$$

Substitute the above 3 terms in (1)

$$
\begin{aligned}
\tilde{f}^{\prime \prime}(x) & =a\left(f(x)-\frac{h}{2} f^{\prime}(x)+\frac{h^{2}}{8} f^{\prime \prime}(x)-\frac{h^{3}}{8 \times 6} f^{\prime \prime \prime}(x)+\frac{h^{4}}{16 \times 24} f^{(4)}(x)+O\left(h^{5}\right)\right) \\
& +b f(x) \\
& +c\left(f(x)+h f^{\prime}(x)+\frac{h^{2}}{2!} f^{\prime \prime}(x)+\frac{h^{3}}{6} f^{\prime \prime \prime}(x)+\frac{h^{4}}{24} f^{(4)}(x)+O\left(h^{5}\right)\right)
\end{aligned}
$$

Collect terms

$$
\begin{align*}
\tilde{f}^{\prime \prime}(x) & =(a+b+c) f(x)+f^{\prime}(x) h\left(-\frac{a}{2}+c\right)+f^{\prime \prime}(x) h^{2}\left(\frac{a}{8}+\frac{c}{2}\right)+f^{\prime \prime \prime}(x) h^{3}\left(\frac{-a}{8 \times 6}+\frac{c}{6}\right)  \tag{2}\\
& +f^{(4)} h^{4}\left(\frac{a}{16 \times 24}+\frac{c}{24}\right)+O\left(h^{5}\right)
\end{align*}
$$

Hence for $\tilde{f}^{\prime \prime}(x)$ to best approximate $f^{\prime \prime}(x)$, we need

$$
\begin{aligned}
(a+b+c) & =0 \\
-\frac{a}{2}+c & =0 \\
h^{2}\left(\frac{a}{8}+\frac{c}{2}\right) & =1
\end{aligned}
$$

Solving the above 3 equations we find

$$
\begin{aligned}
a & =\frac{8}{3 h^{2}} \\
b & =-\frac{4}{h^{2}} \\
c & =\frac{4}{3 h^{2}}
\end{aligned}
$$

Hence (1) becomes

$$
\begin{aligned}
\tilde{f}^{\prime \prime}(x) & =a f\left(x-\frac{h}{2}\right)+b f(x)+c f(x+h) \\
& =\frac{8}{3 h^{2}} f\left(x-\frac{h}{2}\right)-\frac{4}{h^{2}} f(x)+\frac{4}{3 h^{2}} f(x+h)
\end{aligned}
$$

To examine the local truncation error, from (2), and using the solution we just found for $a, b, c$ we find

$$
\begin{aligned}
\tilde{f^{\prime \prime}}(x) & =f^{\prime \prime}(x)+f^{\prime \prime \prime}(x) h^{3}\left(\frac{-\left(\frac{8}{3 h^{2}}\right)}{8 \times 6}+\frac{\left(\frac{4}{3 h^{2}}\right)}{6}\right)+f^{(4)} h^{4}\left(\frac{\left(\frac{8}{3 h^{2}}\right)}{16 \times 24}+\frac{\left(\frac{4}{3 h^{2}}\right)}{24}\right)+O\left(h^{5}\right) \\
& =f^{\prime \prime}(x)+f^{\prime \prime \prime}(x) h^{3}\left(\frac{1}{6 h^{2}}\right)+f^{(4)} h^{4}\left(\frac{1}{16 h^{2}}\right)+O\left(h^{5}\right) \\
& =f^{\prime \prime}(x)+f^{\prime \prime \prime}(x)\left(\frac{h}{6}\right)+f^{(4)} h^{2}\left(\frac{1}{16}\right)+O\left(h^{5}\right)
\end{aligned}
$$

We can truncate at either $f^{\prime \prime \prime}(x)$ or $f^{(4)}$. In the first case, we obtain

$$
\tilde{f}^{\prime \prime}(x)=f^{\prime \prime}(x)+O(h)
$$

Where $O(h)=\frac{f^{\prime \prime \prime}(x)}{6} h$, hence $p=1$ in this case, and with the truncation error $\tau=\frac{f^{\prime \prime \prime}\left(x_{j}\right)}{6} h$ at each grid point.
In the second case, we obtain

$$
\tilde{f}^{\prime \prime}(x)=f^{\prime \prime}(x)+\frac{f^{\prime \prime \prime}(x)}{6} h+O\left(h^{2}\right)
$$

Where $O\left(h^{2}\right)=\frac{f^{(4)}}{16} h^{2}$ and $p=2$ in this case, and with the truncation error $\tau=\frac{f^{(4)}\left(x_{j}\right)}{16} h^{2}$ at each grid point. We see that $\tau$ is smaller if we use $p=2$ than $p=1$.

The accuracy then depends on where we decide to truncate. For example, at $p=1$, the error is dominated by $O(h)$, and at $p=2$, it is $O\left(h^{2}\right)$.

## 4.2 part (b) Refinement study

Given $f(x)=\cos (2 \pi x)$, first, let us find the accuracy of this scheme. The finite difference approximation formula found is

$$
\begin{equation*}
\tilde{f}^{\prime \prime}(x)=\frac{8}{3 h^{2}} f\left(x-\frac{h}{2}\right)-\frac{4}{h^{2}} f(x)+\frac{4}{3 h^{2}} f(x+h) \tag{1}
\end{equation*}
$$

And the exact value is

$$
\begin{equation*}
\frac{d^{2}}{d x^{2}} \cos (2 \pi x)=-4 \pi^{2} \cos 2 \pi x \tag{2}
\end{equation*}
$$

To find the local error $\tau$

$$
\tau=\tilde{f}^{\prime \prime}(x)-f^{\prime \prime}(x)
$$

Substitute $f(x)=\cos (2 \pi x)$ in the RHS of (1) to find the approximation of the second derivative and subtract the exact result value of the second derivative from it.

Plug $f(x)=\cos (2 \pi x)$ in RHS of (1) we obtain

$$
\begin{aligned}
\tilde{f}^{\prime \prime}(x) & =\frac{8}{3 h^{2}} \cos \left(2 \pi\left(x-\frac{h}{2}\right)\right)-\frac{4}{h^{2}} \cos (2 \pi x)+\frac{4}{3 h^{2}} \cos (2 \pi(x+h)) \\
& =\frac{8}{3 h^{2}} \cos (2 \pi x-\pi h)-\frac{4}{h^{2}} \cos (2 \pi x)+\frac{4}{3 h^{2}} \cos (2 \pi x+2 \pi h)
\end{aligned}
$$

Hence the local error $\tau$ is

$$
\begin{aligned}
\tau & =\tilde{f}^{\prime \prime}(x)-f^{\prime \prime}(x) \\
& =\left[\frac{8}{3 h^{2}} \cos (2 \pi x-\pi h)-\frac{4}{h^{2}} \cos (2 \pi x)+\frac{4}{3 h^{2}} \cos (2 \pi x+2 \pi h)\right]+4 \pi^{2} \cos 2 \pi x
\end{aligned}
$$

We notice that $\tau$ depends on $h$ and $x$. At $x=1$,

$$
\begin{aligned}
\tau & =\left[\frac{8}{3 h^{2}} \cos (2 \pi-\pi h)-\frac{4}{h^{2}}+\frac{4}{3 h^{2}} \cos (2 \pi+2 \pi h)\right]+4 \pi^{2} \\
& =\frac{4}{3 h^{2}}\left(\cos (2 \pi h)+2 \cos (\pi h)+3 h^{2} \pi^{2}-3\right)
\end{aligned}
$$

In the following we plot local error $\tau$ as a function of $h$ in linear scale and $\log$ scale. Here is the result.


Figure 2: matlab HW1 partb

We notice that the log plot shows the slope $p=2$ and not $p=1$. This is because the $O(h)$ part turned out to be zero at $x=1$ this is because $O(h)=\frac{f^{\prime \prime \prime}(x)}{6} h=\frac{\left(8 \pi^{3} \sin 2 \pi x\right)}{6} h$ and this term is zero at $x=1$, so the dominant error term became $O\left(h^{2}\right)$ which is $\frac{f^{(4)}\left(x_{j}\right)}{16} h^{2}=\frac{16 \pi^{4} \cos 2 \pi x}{16} h^{2}$ or $\pi^{4} h^{2}$ or $O\left(h^{2}\right)$.
This is why we obtained $p=2$ and not $p=1$ at $x=1$.
The following table show the ratio of the local error between each 2 successive halving of the spacing $h$. Each time $h$ is halved, and the ratio of the error (absolute local error) is shown. We see for $x=1$ that the ratio approaches 4 . This indicates that $p=2$.

```
EDU>> nma_HW1_partb()
h error
5.0000E-001 1.8145E+001
2.5000E-001 5.6483E+000 3.2125E+000
    0.0000E+000
1.2500E-001 1.4936E+000 3.7816E+000
6.2500E-002 3.7872E-001 3.9439E+000
3.1250E-002 9.5014E-002 3.9859E+000
1.5625E-002 2.3775E-002 3.9965E+000
7.8125E-003 5.9449E-003 3.9991E+000
3.9063E-003 1.4863E-003 3.9998E+000
```


### 4.3 Part(c)

The refinement study in part (b) showed that the local error became smaller as $h$ become smaller, and the error was $O\left(h^{2}\right)$ since $p=2$ in the log plot.

But this is not a good test as it was done only for one point $x=1$. We need to examine the approximation scheme at other points as well. The reason is the local error at an $x$ location is

$$
\tau=\left[\frac{8}{3 h^{2}} \cos (2 \pi x-\pi h)-\frac{4}{h^{2}} \cos (2 \pi x)+\frac{4}{3 h^{2}} \cos (2 \pi x+2 \pi h)\right]+4 \pi^{2} \cos 2 \pi x
$$

which can be seen to be a function of $x$ and $h$. In (b) we found that at $x=1, \tau=O\left(h^{2}\right)$ and this was because the dominant error term $O(h)$ happened to vanish at $x=1$.
But if we examine $\tau$ at different point, say $x=0.2$, then we will see that $\tau$ is $O(h)$ and $p=1$.

Here is a plot of $\tau$ at $x=0.2$ and at $x=1$. Both showing what happens as $h$ becomes smaller. We see that the at $x=1$ the approximation was more accurate $(p=2)$ but at $x=0.2$ the approximation was less accurate $(p=1)$. What we conclude from this, is that a single test is not sufficient for determine the accuracy for all points. More tests are needed at other points to have more confidence. To verify that at $x=0.2$ we indeed have $p=1$, we generate the error table as shown above, but for $x=0.2$ this time.


Figure 3: matlab HW1 partc

```
EDU>> nma_HW1_partc()
    h
5.0000E-001
2.5000E-001 
1.2500E-001 5.1898E+000
6.2500E-002 2.5508E+000
3.1250E-002 1.2551E+000
1.5625E-002 6.2133E-001 2.0200E+000
7.8125E-003 3.0897E-001 2.0110E+000
3.9063E-003 1.5404E-001 2.0057E+000
```

We see that the ratio becomes 2 this time, not 4 as we half the spacing each time. This mean $p=1$. This means the accuracy of the formula used can depend on the location.

### 4.4 Part (d)

The points that we need to interpolate are $\left[\left[x-\frac{h}{2}, u\left(x-\frac{h}{2}\right)\right],[x, u(x)],[x+h, u(x+h)]\right]$ where $u=\cos (2 \pi x)$

Since we require a quadratic polynomial, then we write

$$
p(x)=a+b x+c x^{2}
$$

Where $p(x)$ is the interpolant. Evaluate the above at each of the 3 points. Choose $x=1$, hence the points are

$$
\left(\begin{array}{c}
1-\frac{h}{2}, u\left(1-\frac{h}{2}\right) \\
1, u(x) \\
1+h, u(1+h)
\end{array}\right)
$$

Evaluate $p(x)$ at each of these points

$$
\begin{aligned}
p\left(1-\frac{h}{2}\right) & =\cos \left(2 \pi\left(1-\frac{h}{2}\right)\right)=a+b\left(1-\frac{h}{2}\right)+c\left(1-\frac{h}{2}\right)^{2} \\
p(1) & =\cos (2 \pi)=a+b+c \\
p(1+h) & =\cos (2 \pi(1+h))=a+b(1+h)+c(1+h)^{2}
\end{aligned}
$$

or

$$
\begin{aligned}
\left(\begin{array}{ccc}
\left(1-\frac{h}{2}\right)^{2} & \left(1-\frac{h}{2}\right) & 1 \\
1 & 1 & 1 \\
(1+h)^{2} & (1+h) & 1
\end{array}\right)\left(\begin{array}{l}
c \\
b \\
a
\end{array}\right) & =\left(\begin{array}{c}
\cos \left(2 \pi\left(1-\frac{h}{2}\right)\right) \\
\cos (2 \pi 1) \\
\cos (2 \pi(1+h))
\end{array}\right) \\
A v & =b
\end{aligned}
$$

Solving the above Vandermonde system, we obtain

$$
\begin{aligned}
a & =\frac{1}{3 h^{2}}(4(1+h) \cos (h \pi)+(h-2)(3+3 h-\cos (2 \pi h))) \\
b & =\frac{-1}{3 h^{2}} 4((h-4) \cos (\pi h)-h-8) \sin ^{2}\left(\frac{\pi h}{2}\right) \\
c & =\frac{2}{3 h^{2}}(2 \cos (\pi h)-3+\cos (2 \pi h))
\end{aligned}
$$

Hence

$$
\begin{align*}
p(x) & =\left[\frac{1}{3 h^{2}}(4(1+h) \cos (h \pi)+(h-2)(3+3 h-\cos (2 \pi h)))\right]  \tag{1}\\
& -\left[\frac{1}{3 h^{2}} 4((h-4) \cos (\pi h)-h-8) \sin ^{2}\left(\frac{\pi h}{2}\right)\right] x \\
& +\left[\frac{2}{3 h^{2}}(2 \cos (\pi h)-3+\cos (2 \pi h))\right] x^{2}
\end{align*}
$$

Recall, that we found, for $u=\cos (2 \pi x)$, the finite difference formula was

$$
\begin{equation*}
\tilde{u}^{\prime \prime}(x)=\left[\frac{8}{3 h^{2}} \cos (2 \pi x-\pi h)-\frac{4}{h^{2}} \cos (2 \pi x)+\frac{4}{3 h^{2}} \cos (2 \pi x+2 \pi h)\right] \tag{2}
\end{equation*}
$$

Take the second derivative of $p(x)$ shown in (1) above

$$
\begin{equation*}
p^{\prime \prime}(x)=\frac{4}{3 h^{2}}(2 \cos (\pi h)+\cos (2 \pi h)-3) \tag{3}
\end{equation*}
$$

But we notice that $\tilde{u}^{\prime \prime}(x)$ evaluated at $x=1$ is

$$
\tilde{u}^{\prime \prime}(1)=\frac{4}{3 h^{2}}(2 \cos (\pi h)+\cos (2 \pi h)-3)
$$

which is the same as $p^{\prime \prime}(x)$.
Therefore, $p^{\prime \prime}(x)$ is the same as as the finite difference approximation evaluated at the central point of the 3 points, used to generate $p$.

In other words, given 3 points

$$
\left(\begin{array}{c}
x_{0}-\frac{h}{2}, u\left(x_{0}\right) \\
x_{0}, u\left(x_{0}\right) \\
x_{0}+h, u\left(x_{0}+h\right)
\end{array}\right)
$$

Where $u(x)$ is some function (here it was $\cos (2 \pi x)$ ), and we generate a quadratic interpolant polynomial $p(x)$ using the above 3 points, then $p^{\prime \prime}(x)$ will given the same value as the finite difference formula evaluated at $x_{0}$.

$$
\left.p^{\prime \prime}(x)\right|_{x=x_{0}}=\left.\tilde{u}(x)\right|_{x=x_{0}}
$$

For this to be valid, $p(x)$ must have been generated with the center point being $x_{0}$. If we pick another center point $x_{1}$, and therefore have the 3 points $x_{1}-h / 2, x_{1}, x_{1}+h$, and then generate a polynomial $q(x)$ as above, then we will find

$$
\left.q^{\prime \prime}(x)\right|_{x=x_{1}}=\left.\tilde{u}(x)\right|_{x=x_{1}}
$$

This is illustrated by the following diagram


Figure 4: prob3 c

## 5 Appendix (Source code)

### 5.1 Matlab

```
%-- by Nasser M. Abbasi, Math 228A, UC Davis, Fall }201
%-- implement part b, problem 3
function nma_HW1_partb()
%-- Generate h values to use, and define tao(h) function
N = 8;
pointAt=1;
data = arrayfun( @(i) [1/(2^i) , local_error(1/(2^i),pointAt)],1:N, ...
    'UniformOutput',false);
data = reshape(cell2mat (data),2,N)';
%-- plot the tao(h) in linear and log scale
set(0,'defaultaxesfontsize',8) ;
set(0,'defaulttextfontsize',8);
subplot(2,1,1);
plot(data(:,1),data(:,2),'-o'); grid on;
title('tao at x=(1), linear scale');
xlabel('spacing h'); ylabel('ABS(error)');
subplot(2,1,2);
loglog(data(:,1),data(:,1),'-0'); grid on;
title('tao at x=(1), log scale');
xlabel('log(h)'); ylabel('log(ABS(error))');
export_fig matlab_HW1_partb.png
%-- now generate the error table, find ratio first
error_ratio = zeros(N,1);
for i=2:N
```

```
    error_ratio(i) = data((i-1),2)/data(i,2);
end
%-- print table
fprintf('h\t\t\t\t error\t\t\t\t ratio\n');
for i=1:N
    fprintf('%6.4E\t\t%6.4E\t\t%6.4E\n',data(i,1),data(i,2),error_ratio(i));
end
end
function tao=local_error(h,x)
tao=8/(3*h^2)*\operatorname{cos}(2*pi*(x-h/2)) -4/h^2*\operatorname{cos}(2*pi*x)+4/(3*h^2)*...
    cos(2*pi*(x+h))+(2*pi)~ 2*cos(2*pi*x);
end
%-- by Nasser M. Abbasi, Math 228A, UC Davis, Fall }201
%-- implement part c, problem 3
function nma_HW1_partc()
%-- Generate h values to use, and define tao(h) function
%-- plot the tao(h) in linear and log scale
set(0,'defaultaxesfontsize',8) ;
set(0,'defaulttextfontsize',8);
%build data, x-axis is spacing h, y-axis is error
N = 8;
pointAt=1.0;
data = arrayfun( @(i) [1/(2^i) , local_error(1/(2^i),pointAt)],1:N, ...
                    'UniformOutput',false);
data = reshape(cell2mat(data),2,N)';
loglog(data(:,1),data(:,1),'-o'); grid off;
title('tao at x=1 and x=0.2, log scale');
xlabel('log(h)'); ylabel('log(ABS(error))');
hold on;
pointAt=0.2;
data = arrayfun( @(i) [1/(2^i) , local_error(1/(2^i),pointAt)],1:N,\ldots
    'UniformOutput',false);
data = reshape(cell2mat(data),2,N)';
loglog(data(:,1), data(:,2),'-s');
legend('p=2,x=1', 'p=1,x=0.2');
export_fig matlab_HW1_partc.png
%-- now generate the error table, find ratio first
error_ratio = zeros(N,1);
for i=2:N
    error_ratio(i) = data((i-1),2)/data(i,2);
end
%-- print table
fprintf('h\t\t\t\t error\t\t\t\t ratio\n');
for i=1:N
    fprintf('%6.4E\t\t%6.4E\t\t%6.4E\n',data(i,1),data(i,2),error_ratio(i));
end
end
function tao=local_error(h,x)
tao=8/(3*h^2)*\operatorname{cos}(2*pi*(x-h/2)) -4/h^2*cos(2*pi*x) +4/(3*h^2)*. . .
    cos(2*pi*(x+h))+(2*pi)^2*\operatorname{cos}(2*\textrm{pi}*\textrm{x});
end
```


### 5.2 Mathematica

HW 1, problem 3, computational part. math 228A UC davis fall 2010
Nasser M. Abbasi
This is the code used to generate the plots and tables used in HW1

## define local error function

```
localError[h_, x_] :=
    \(\operatorname{Module}\left[\left\}, \frac{8}{3 h^{2}} \operatorname{Cos}[2 \pi x-\pi h]-\frac{4}{h^{2}} \operatorname{Cos}[2 \pi x]+\frac{4}{3 h^{2}} \operatorname{Cos}[2 \pi x+2 \pi h]\right]+(2 \pi)^{2} \operatorname{Cos}[2 \pi x] ;\right.\)
```


## define a function to make the plots

```
makePlot[x_, s_, title_, xlabel_, ylabel_, f_] := Module[{data, n = 8},
    data = Table[{1/(2^i), Abs@localError[1/(2^i), x] }, {i, 1, n}];
    f[data, Joined }->\mathrm{ True, AxesOrigin }->{0,0}
    GridLines }->\mathrm{ Automatic, AspectRatio }->\mathrm{ 1, Frame }->\mathrm{ True, PlotRange }->\mathrm{ All,
    FrameLabel }->\mathrm{ {{ylabel, None}, {xlabel, title}}, PlotStyle }->\mathrm{ s, ImageSize }->\mathrm{ Full]
]
```


## make plot for problem 3, part b

title = Style["local error at $x=1$, log scale", 16];
xlabel = Style["h", 16]; ylabel = Style["local error", 16];
p1 = makePlot[1, \{Thick, Dashed\}, title, xlabel, ylabel, ListLogLogPlot];
title = Style["local error at $x=1$, linear scale", 16];
p2 = makePlot[1, \{Thick, Dashed\}, title, xlabel, ylabel, ListPlot];
Framed [Grid[ \{ \{p1, p2\} \}], ImageSize $\rightarrow$ \{600, 300\}]


## Generate error table, problem 3, part b



Generate table for problem 3, part (c)

```
n=14;
x = 0.2;
data = Table[{1/(2^i), Abs@localError[1/(2^i), x] }, {i, 1, n}];
```


t = TableForm[N[data, \$MachinePrecision], TableHeadings $\rightarrow$
\{None, $\{$ " $h$ ", "local error $\tau$ ", "ratio"\}\}, TableSpacing $\rightarrow\{1,6\}$, TableAlignments $\rightarrow$ Left]; Labeled[Framed@ScientificForm[t, \{8, 6\}, NumberFormat $\rightarrow$ (Row[\{\#1, "e", \#3\}] \&),

NumberPadding $\rightarrow$ \{"", "0"\}], Style["local error as function of $h$ at $x=0.2 ", 14]$, Top]
local error as function of $h$ at $x=0.2$

| h | local error $\tau$ | ratio |
| :--- | :--- | :--- |
| $5.000000 \mathrm{e}-1$ | 1.575174 e 1 | 0.000000 e |
| $2.500000 \mathrm{e}-1$ | 1.014949 e 1 | 1.551974 e |
| $1.250000 \mathrm{e}-1$ | 5.189762 e | 1.955675 e |
| $6.250000 \mathrm{e}-2$ | 2.550829 e | 2.034539 e |
| $3.125000 \mathrm{e}-2$ | 1.255100 e | 2.032371 e |
| $1.562500 \mathrm{e}-2$ | $6.213251 \mathrm{e}-1$ | 2.020037 e |
| $7.812500 \mathrm{e}-3$ | $3.089650 \mathrm{e}-1$ | 2.010989 e |
| $3.906250 \mathrm{e}-3$ | $1.540406 \mathrm{e}-1$ | 2.005737 e |
| $1.953125 \mathrm{e}-3$ | $7.690765 \mathrm{e}-2$ | 2.002930 e |
| $9.765625 \mathrm{e}-4$ | $3.842539 \mathrm{e}-2$ | 2.001480 e |
| $4.882812 \mathrm{e}-4$ | $1.920555 \mathrm{e}-2$ | 2.000744 e |
| $2.441406 \mathrm{e}-4$ | $9.600990 \mathrm{e}-3$ | 2.000372 e |
| $1.220703 \mathrm{e}-4$ | $4.800048 \mathrm{e}-3$ | 2.000186 e |
| $6.103516 \mathrm{e}-5$ | $2.399851 \mathrm{e}-3$ | 2.000144 e |

## Generate plot for part (C)

title = Style["local error at different x locations, log scale", 16]; xlabel = Style["h", 16]; ylabel = Style["local error", 16]; p1 = makePlot[1, \{Thick, Dashed\}, title, xlabel, ylabel, ListLogLogPlot]; p2 = makePlot [0.2, \{Thick, Black\}, title, xlabel, ylabel, ListLogLogPlot]; Show [ \{p1, p2\}, ImageSize $\rightarrow$ 500]
local error at different $x$ locations, log scale


